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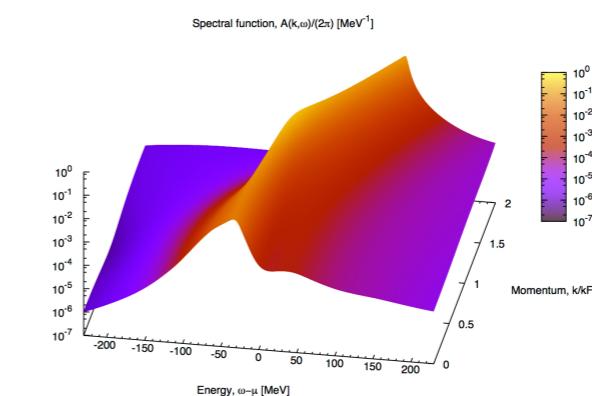
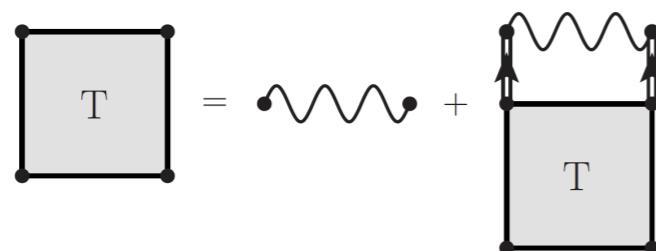
Alexander von Humboldt
Stiftung / Foundation



Self-consistent Green's functions with three-nucleon interactions

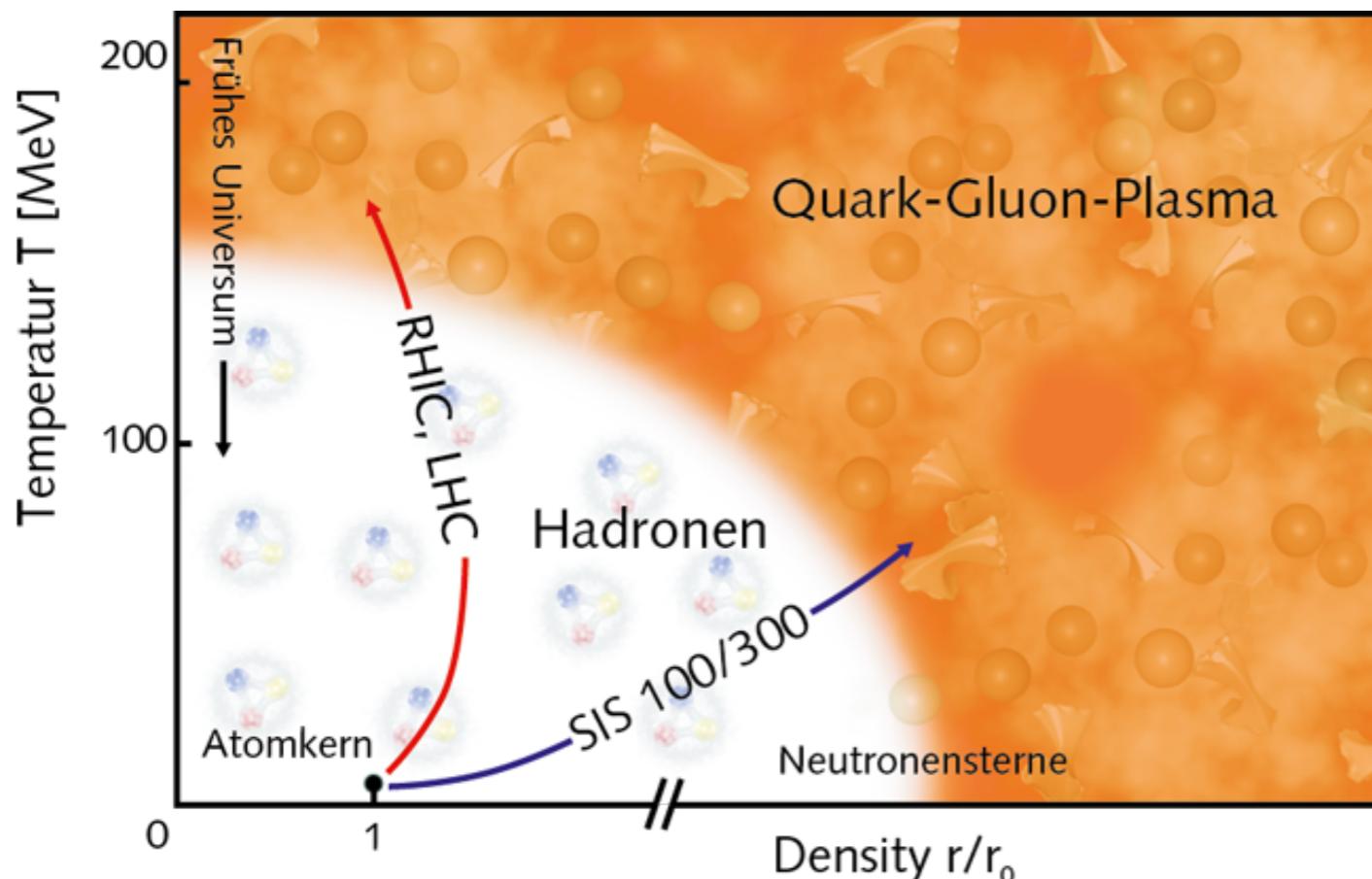
Arianna Carbone - TU Darmstadt
EMMI Workshop 13-16 October 2015

Cold dense nuclear matter: from short-range nuclear correlations to neutron stars



Nuclear matter covers wide ranges of density and temperature

The phase diagram of hadronic matter

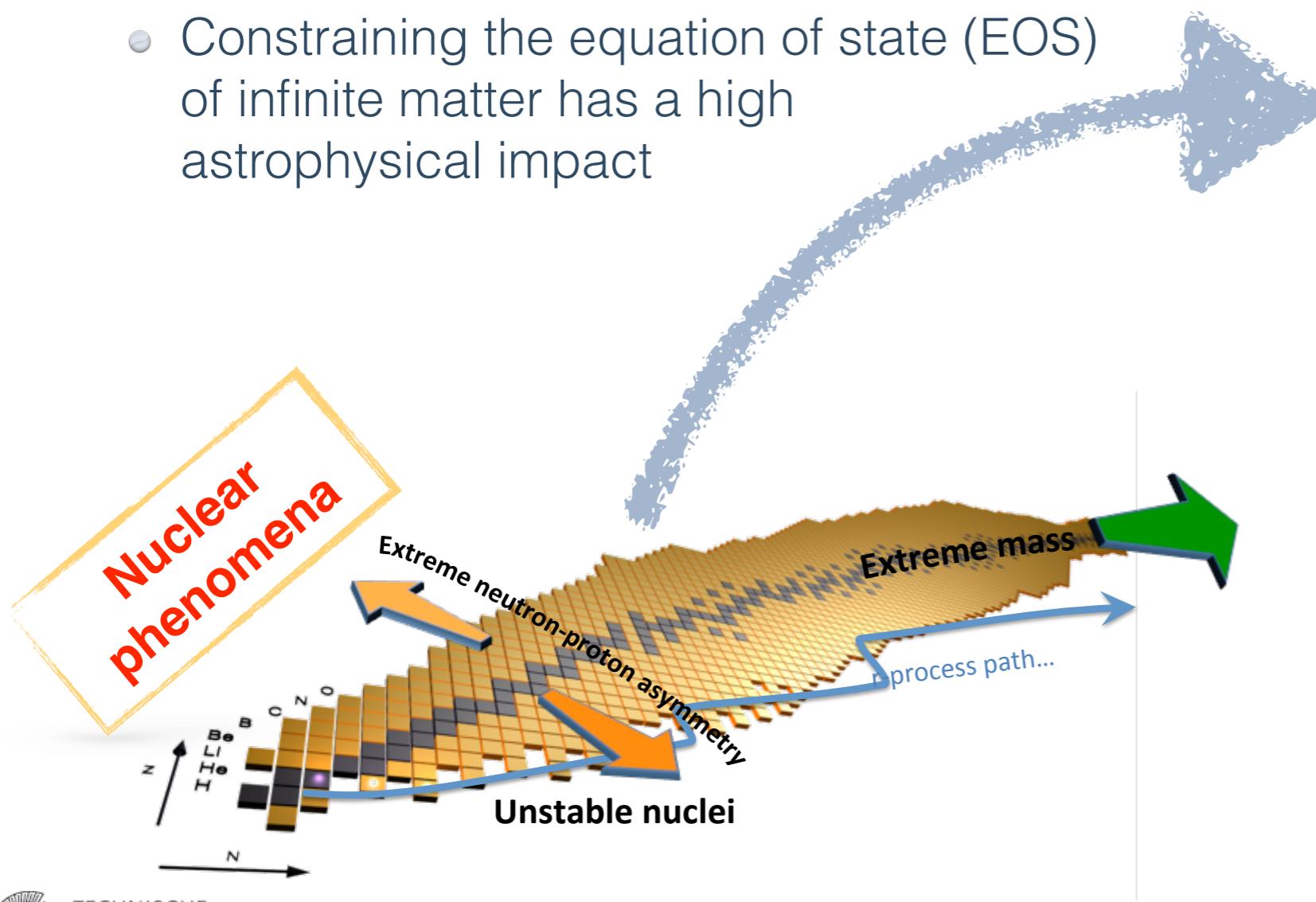


- Matter interacting via the strong force appears in diverse forms
- Experiments try to fill in the pieces of the phase diagram puzzle
- High-energy nucleus-nucleus collisions will probe the neutron-rich region
- What's the contribution from nuclear many-body theory?

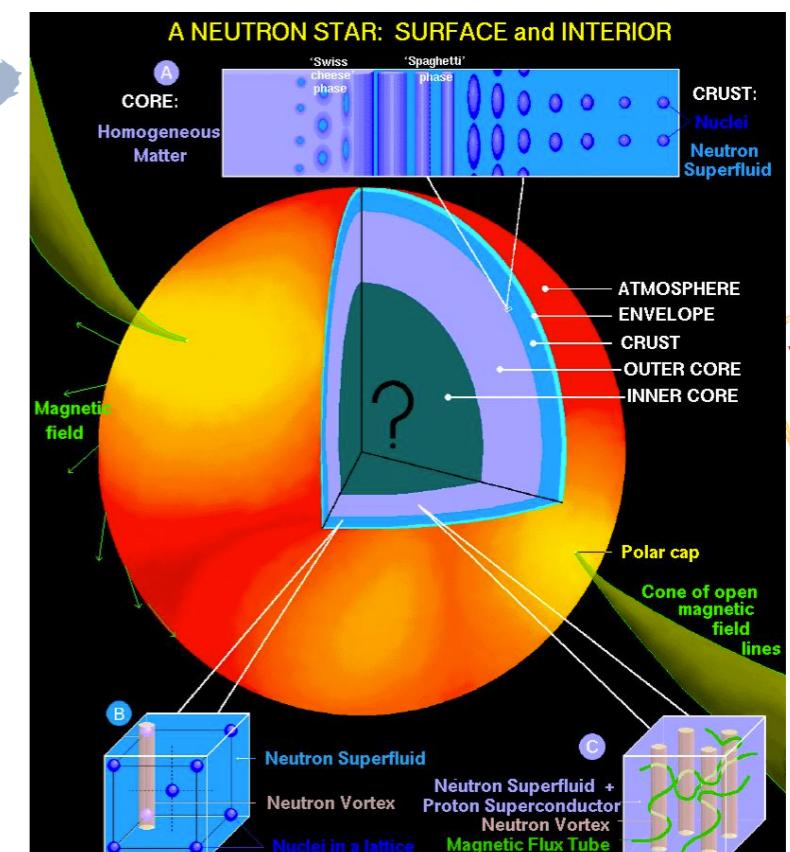
https://www.gsi.de/en/start/fair/forschung_an_fair/kernmateriephysik.htm

The nuclear many-body problem

- Build reliable methods with predictive power
- The study of exotic nuclei is probing the limits of the nuclear landscape
- Constraining the equation of state (EOS) of infinite matter has a high astrophysical impact



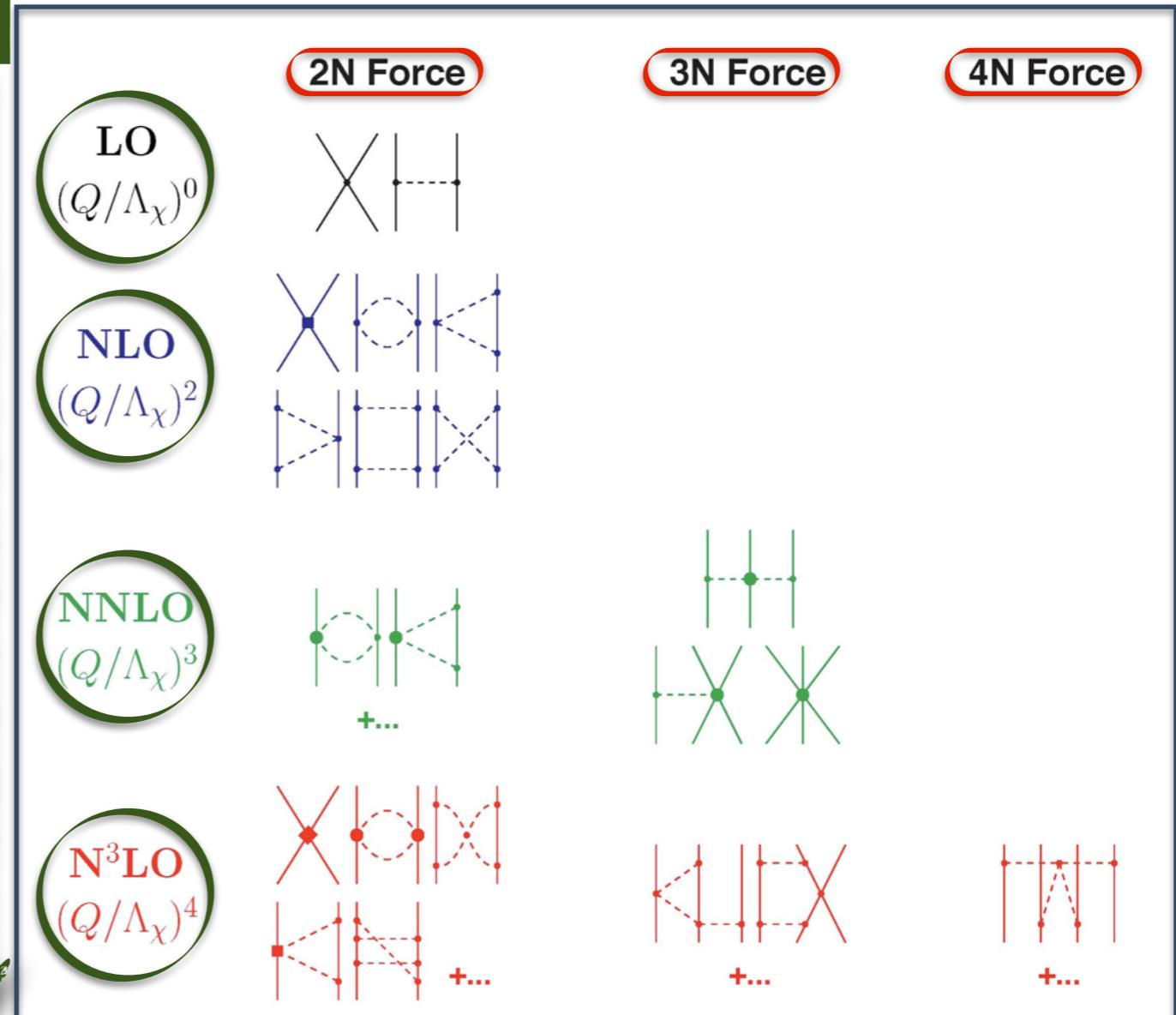
From nuclei to nuclear matter



Why nuclear matter from chiral EFT?

Power counting

- Effective theory of QCD
- Nucleons & pions as degrees of freedom
- Power counting expansion
- Hierarchy of many-body forces
- Enables estimates of theoretical uncertainties



Epelbaum *et al.*, Rev. Mod. Phys. 81, 1773(2009)
Machleidt *et al.*, Phys. Rep. 503, 1 (2011)

Chiral nuclear forces

Power counting

Approach used:

2NF: N2LO or N3LO

Entem & Machleidt, PRC 68, 041001 (2003)

Epelbaum *et al.*, NPA **747**, 362 (2005)

Ekström *et al.*, PRL **110**, 192502 (2013)

Recent improvements:

Epelbaum *et al.*, EPJA **51**, 53 (2015)

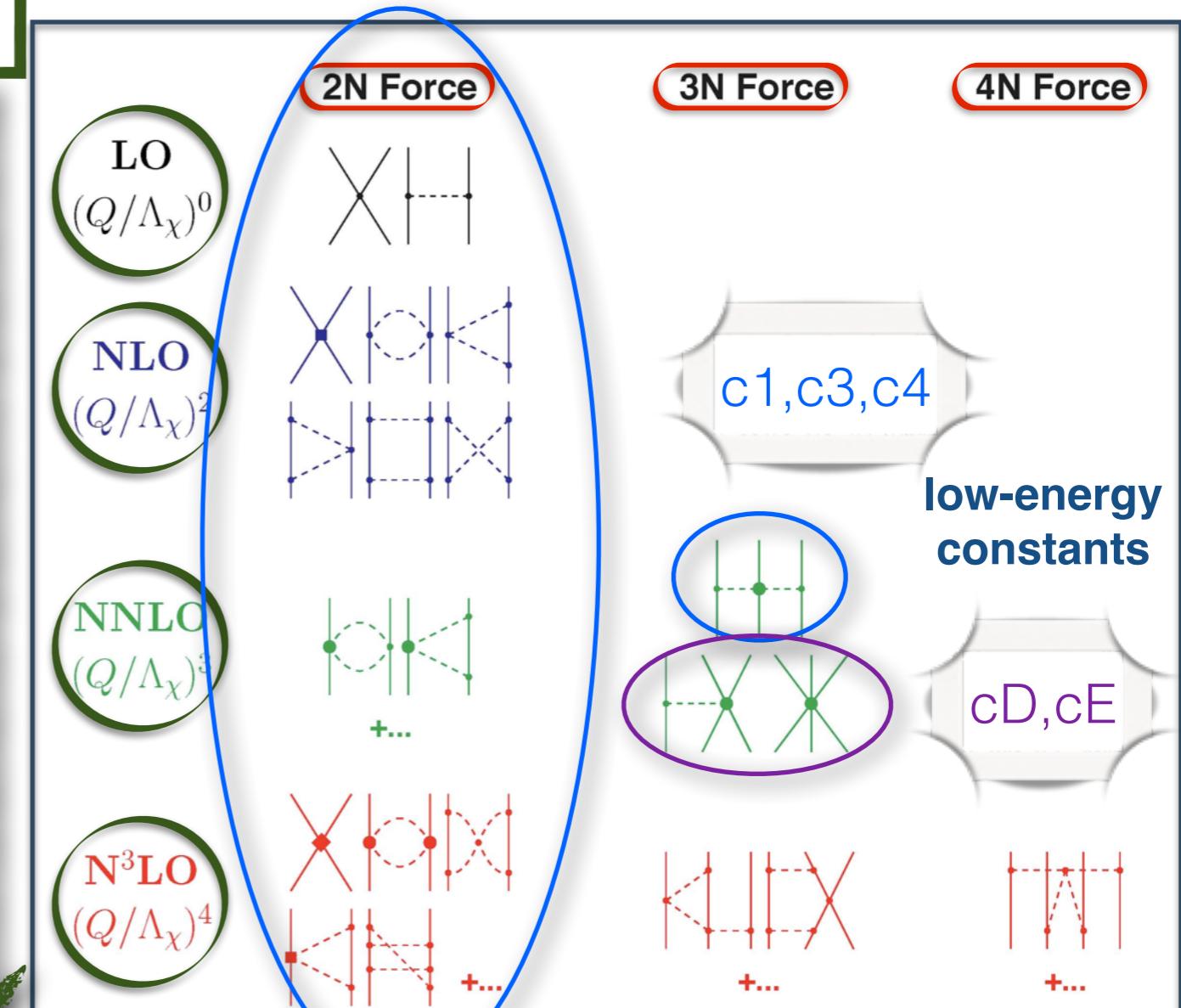
Ekström *et al.*, PRC **91**, 051301(R) (2015)

3NF: N2LO

Van Kolck, PRC **49**, 2932 (1994)

Epelbaum *et al.*, PRC 66, 064001 (2002)

At N2LO two new couplings need to
be fit to few-body data



Epelbaum *et al.*, Rev. Mod. Phys. 81, 1773(2009)
Machleidt *et al.*, Phys. Rep. 503, 1 (2011)

Chiral nuclear forces

Power counting

First calculations with full N3LO:

3NF: N3LO

Ishikawa *et al.*, PRC **76**, 014006 (2007)

Bernard *et al.*, PRC **77**, 064004 (2008)

Bernard *et al.*, PRC **84**, 054001 (2011)

Important recent improvements

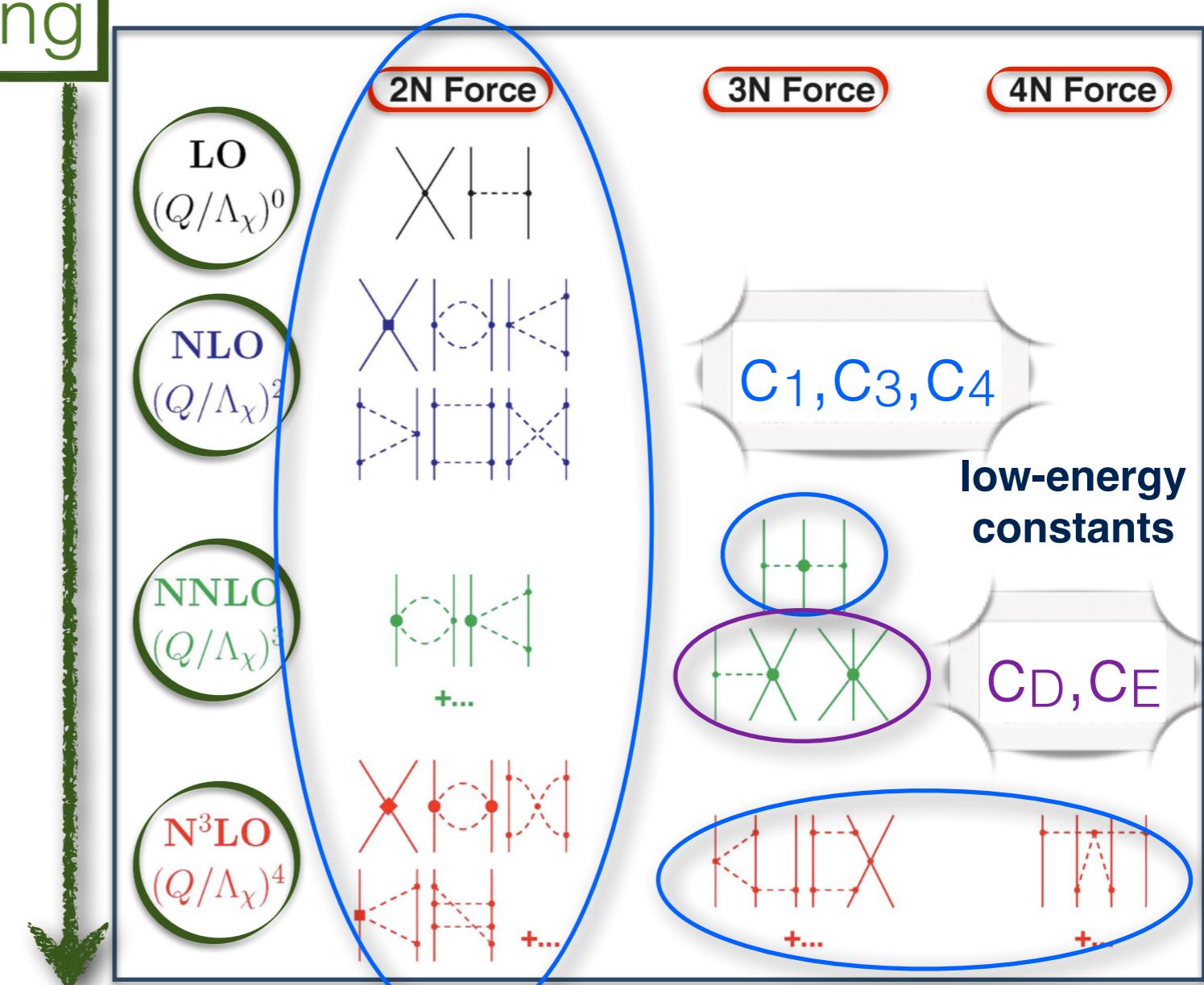
Hebeler *et al.*, PRC **91**, 044001 (2015)

4NF: N3LO

Epelbaum, PLB **639**, 456 (2006)

Epelbaum, EPJA **34**, 197 (2007)

Many-body forces have to be included because they are part of the theory!



Epelbaum *et al.*, Rev. Mod. Phys. 81, 1773(2009)
Machleidt *et al.*, Phys. Rep. 503, 1 (2011)

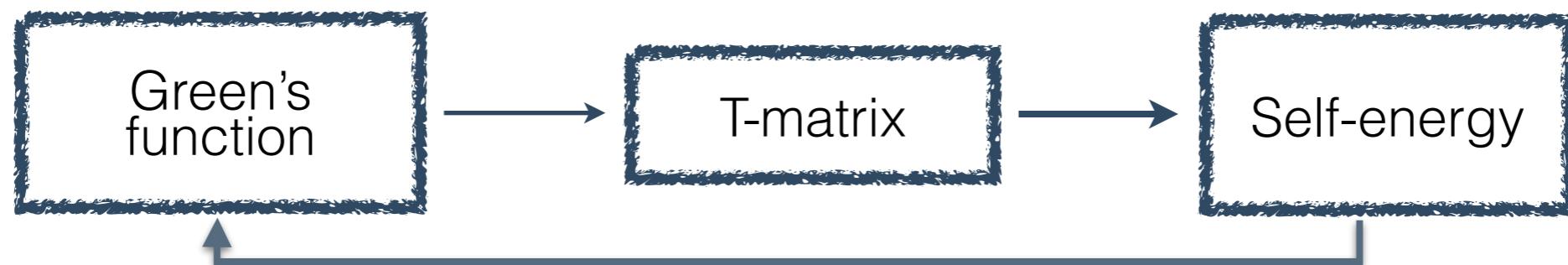
Self-consistent Green's functions

Dickhoff & Barbieri, PPNP **52**, 377 (2004)

- The Green's function as a tool to solve the nuclear many-body problem:

$$G_{\alpha\beta}(\omega) = \sum_m \frac{\langle \Psi_0^N | a_\alpha | \Psi_m^{N+1} \rangle \langle \Psi_m^{N+1} | a_\beta^\dagger | \Psi_0^N \rangle}{\omega - (E_m^{N+1} - E_0^N) + i\eta} + \sum_n \frac{\langle \Psi_0^N | a_\beta^\dagger | \Psi_n^{N-1} \rangle \langle \Psi_n^{N-1} | a_\alpha | \Psi_0^N \rangle}{\omega - (E_0 - E_n^{N-1}) + i\eta}$$

- Self-consistent nonperturbative method:



- Nonrelativistic Hamiltonian:

$$\hat{H} = \sum_\alpha \varepsilon_\alpha^0 a_\alpha^\dagger a_\alpha - \sum_{\alpha\beta} U_{\alpha\beta} a_\alpha^\dagger a_\beta + \frac{1}{4} \sum_{\substack{\alpha\gamma \\ \beta\delta}} V_{\alpha\gamma,\beta\delta} a_\alpha^\dagger a_\gamma^\dagger a_\delta a_\beta + \frac{1}{36} \sum_{\substack{\alpha\gamma\epsilon \\ \beta\delta\eta}} W_{\alpha\gamma\epsilon,\beta\delta\eta} a_\alpha^\dagger a_\gamma^\dagger a_\epsilon^\dagger a_\eta a_\delta a_\beta$$

Self-consistent Green's functions

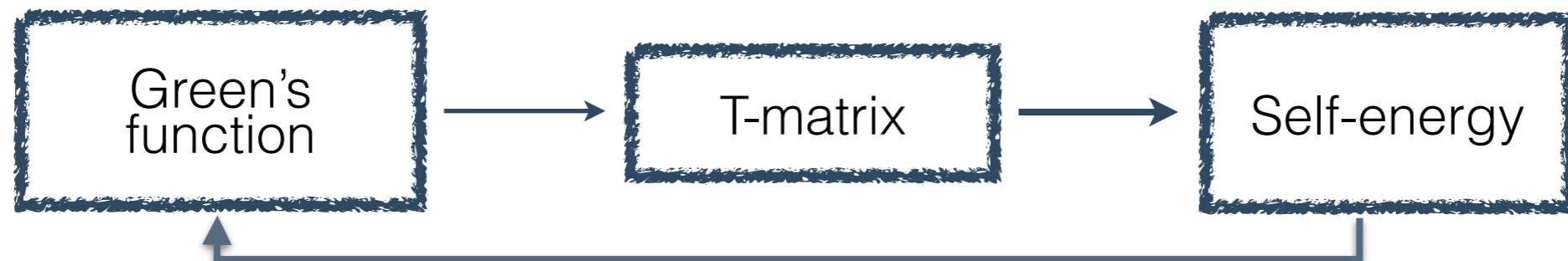
Dickhoff & Barbieri, PPNP **52**, 377 (2004)

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energy with an added particle *energy with a removed particle*

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Self-consistent Green's functions

Dickhoff & Barbieri, PPNP **52**, 377 (2004)

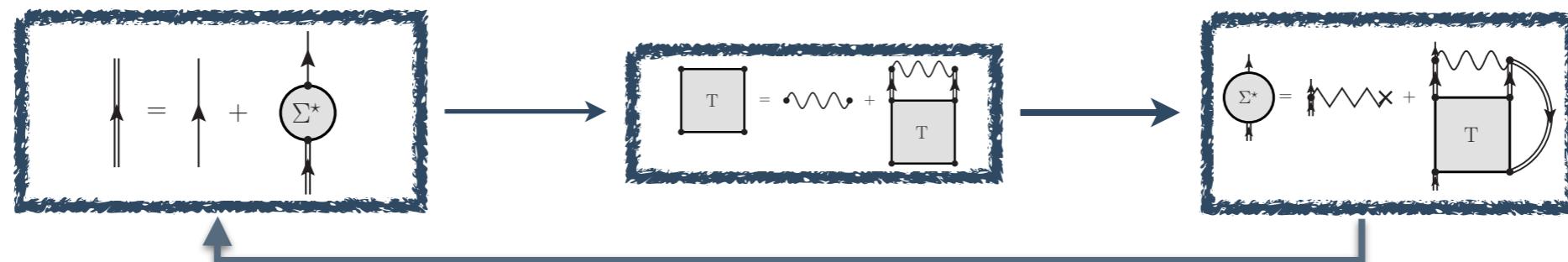
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Self-consistent Green's functions

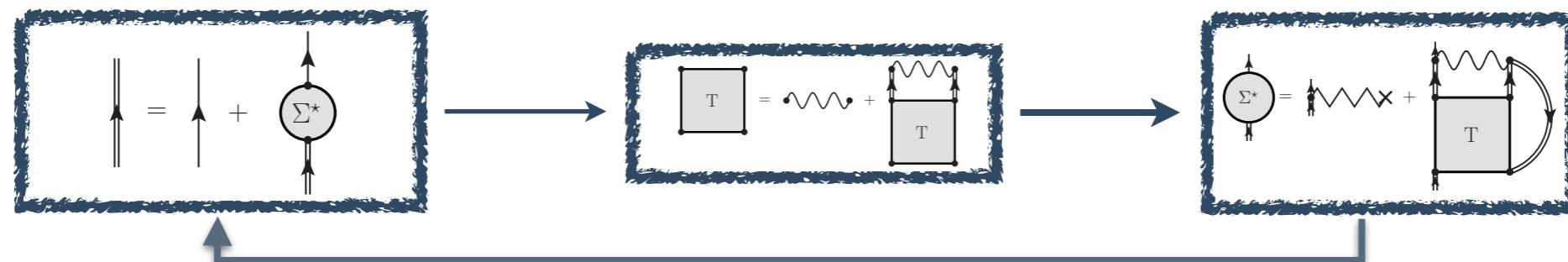
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\hat{H}_0 $\cdots \times \cdots \cdots \cdots \cdots$

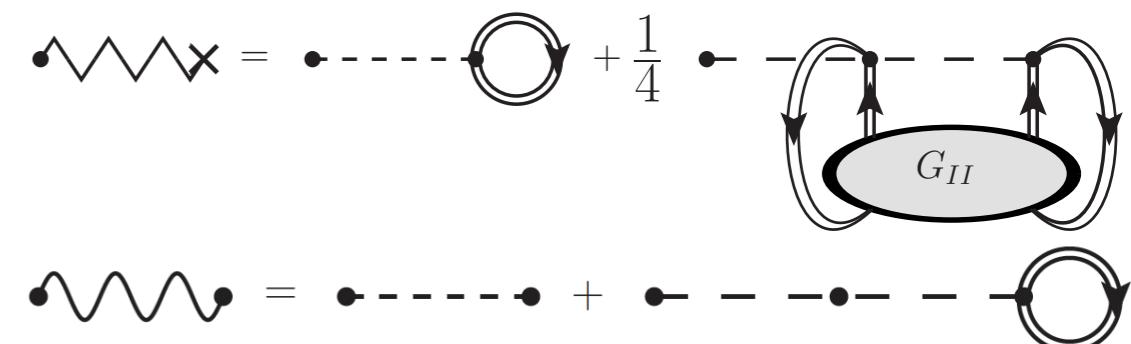
Extend the SCGF approach

Carbone, Cipollone, Barbieri, Rios, Polls, PRC **88**, 054326 (2013)

2B **2B + 3B**

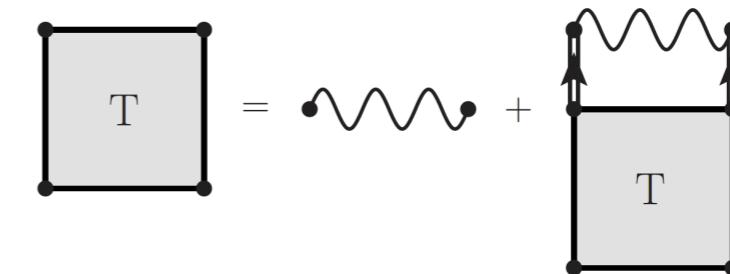
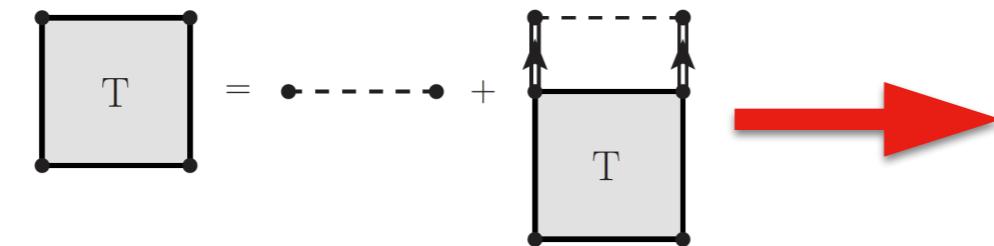
1. define **effective interactions** to include correctly 3B terms:

Interaction



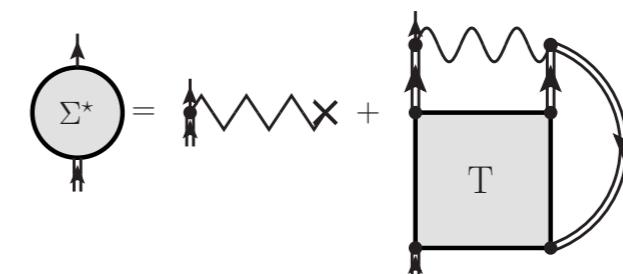
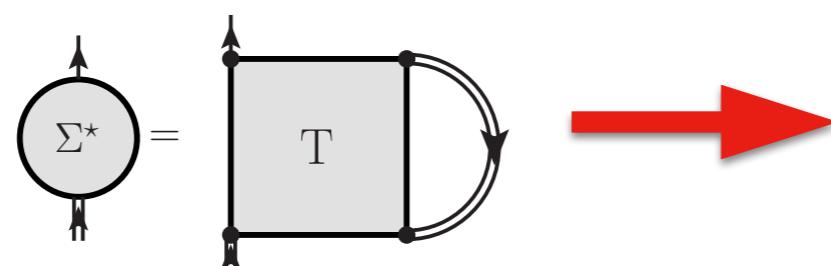
2. calculate T-matrix with effective 2B term:

T-matrix

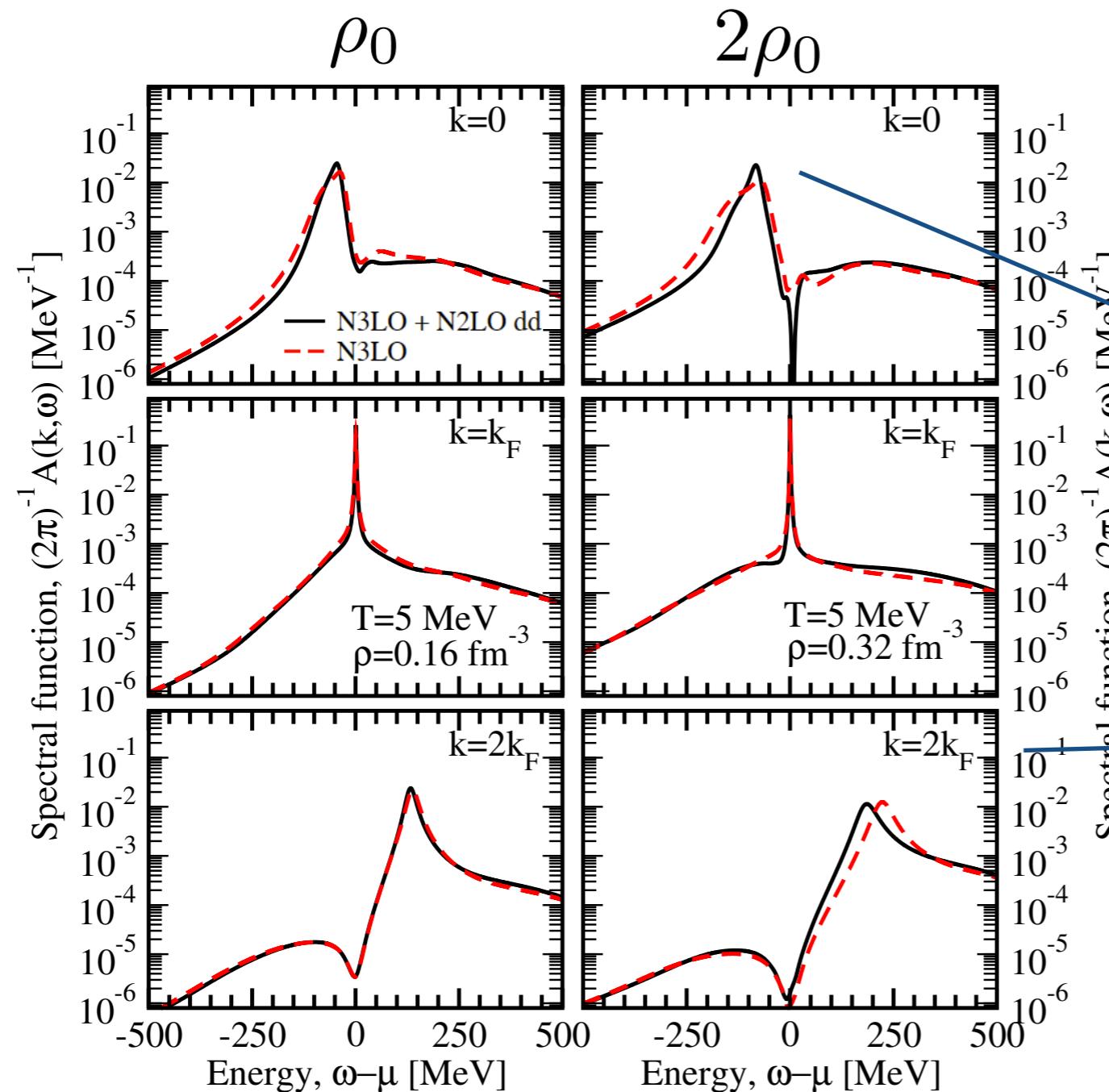


3. calculate self-energy distinguishing the effective terms:

Self-energy



The spectral function



Slight 3BF effect in general....

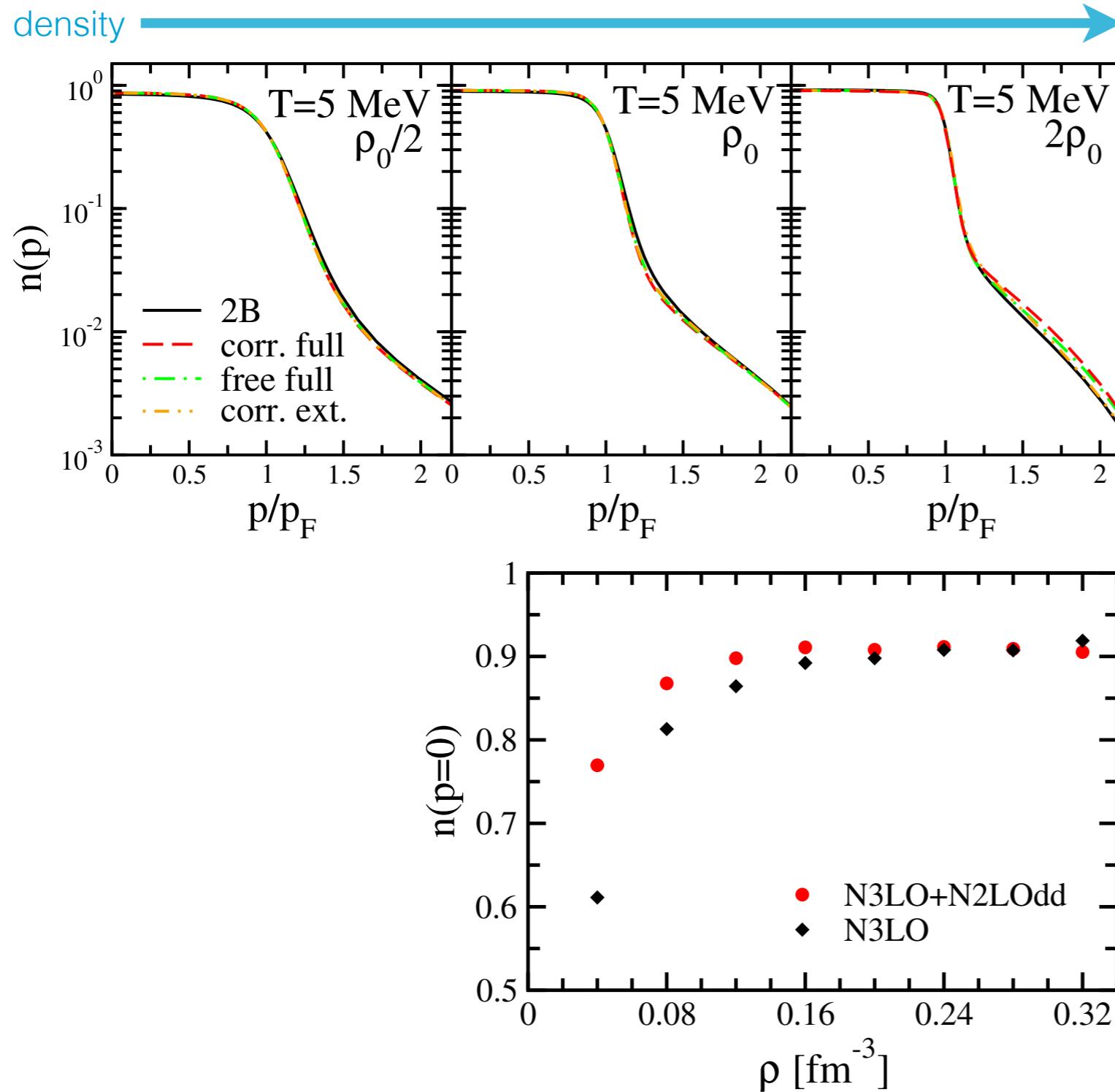
Narrower quasi-particle peak at low momenta

Lower energy for quasi-particle with 3NF because of the rescaling

Carbone, Rios, Polls, PRC 88, 044302 (2013)

Momentum distribution

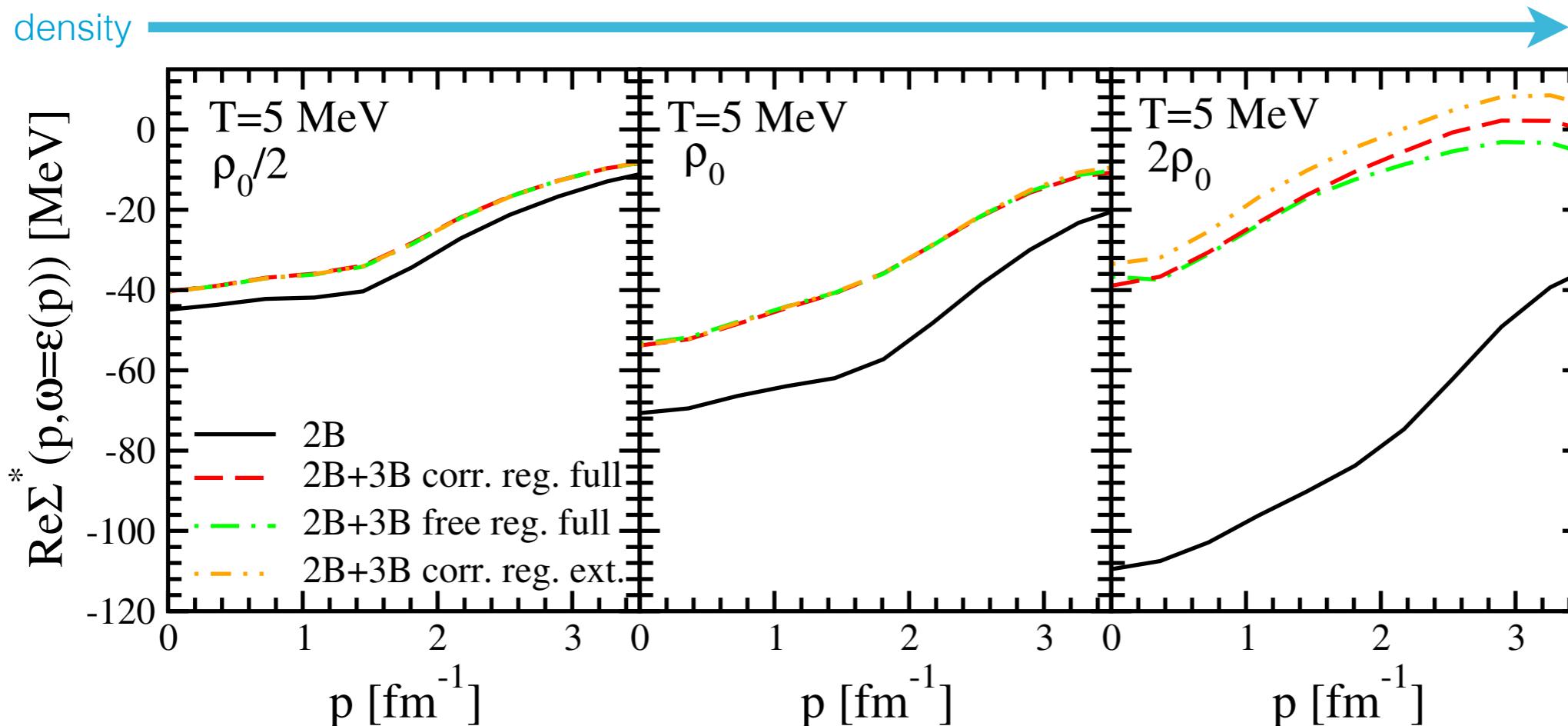
Carbone, Rios, Polls, PRC 90, 054322 (2014)



$$n(p) = \int \frac{d\omega}{2\pi} \mathcal{A}(p, \omega) f(\omega)$$

- small changes due to 3BF
- high-momentum components
- depletion density-dependent
- small averaging differences

Single-particle potential

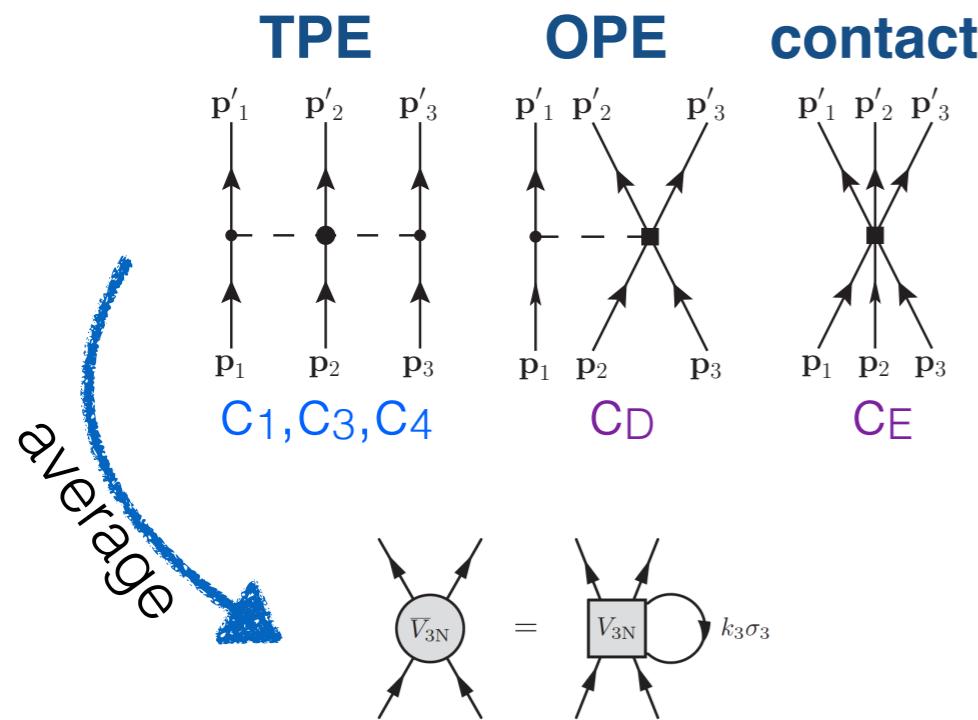


- strong effect of 3-body forces
- repulsion rises with density
- modifications due to averaging procedure visible at high density

$$\varepsilon_{qp}(p) = \frac{p^2}{2m} + \text{Re}\Sigma^\star(p, \varepsilon_{qp}(p))$$

Single-particle spectra

The need for 3-body nuclear forces



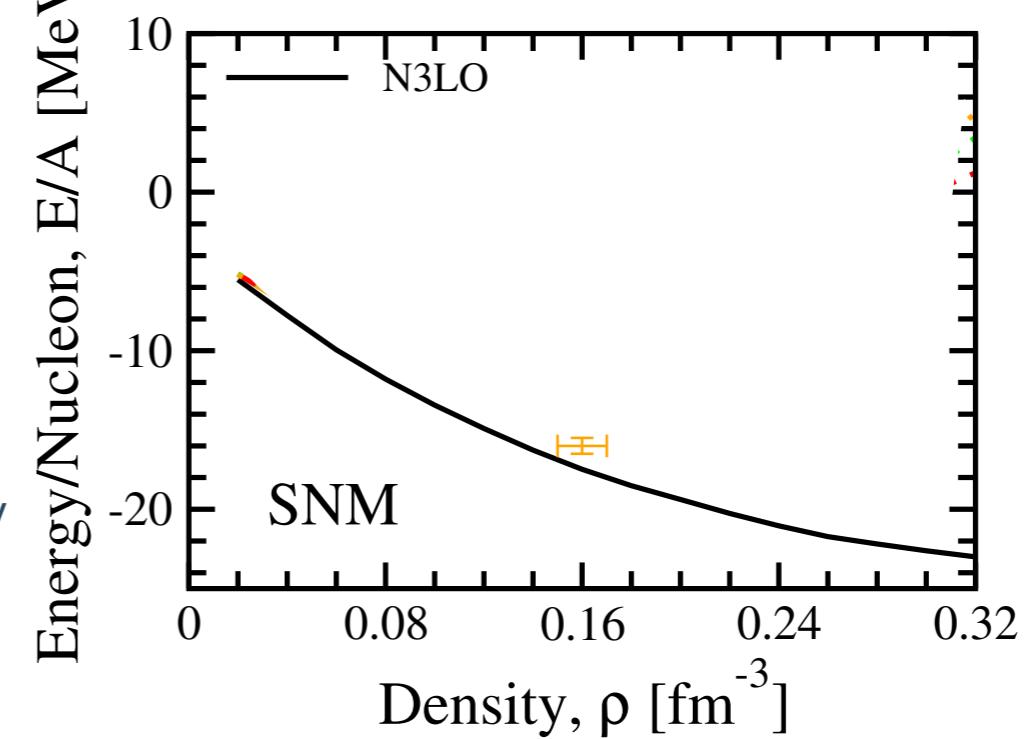
J.W. Holt *et al.*, PRC 81, 024002 (2010)
 Hebeler *et al.*, PRC 82, 014314 (2010)
 Carbone *et al.*, PRC 90, 054322 (2014)

- Repulsion due to 3BF
- Improved prediction of saturation density
- Small averaging dependence
- However saturation energy underbound

The Koltun sumrule

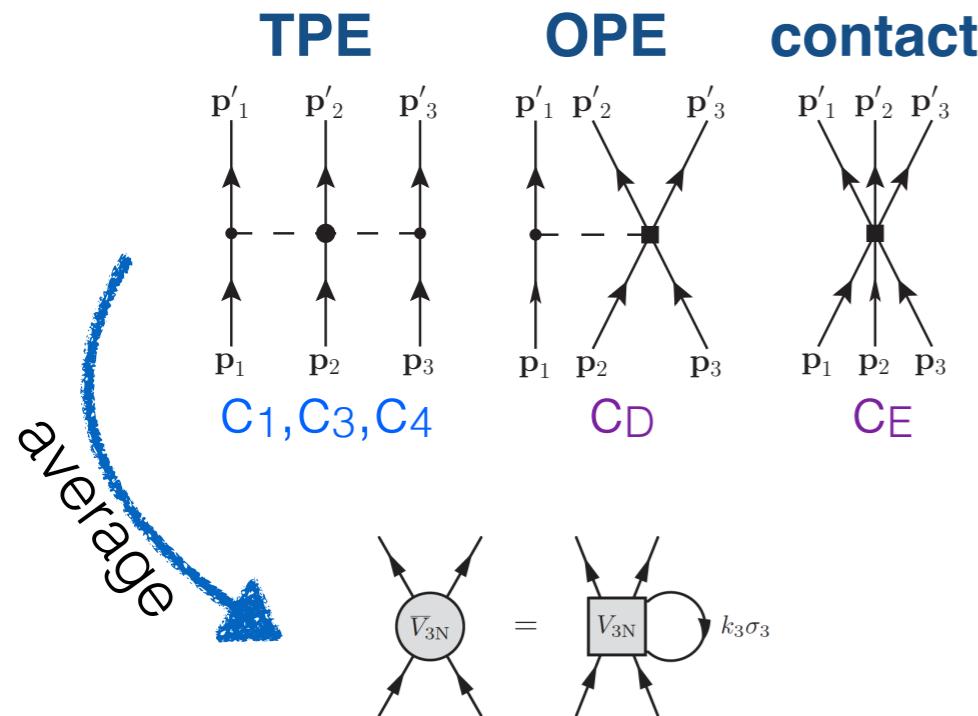
$$\frac{E}{A} = \frac{\nu}{\rho} \int \frac{d^3 p}{(2\pi)^3} \int \frac{d\omega}{2\pi} \frac{1}{2} \left\{ \frac{p^2}{2m} + \omega \right\} \mathcal{A}(p, \omega) f(\omega) - \frac{1}{2} \langle \Psi^N | \hat{W} | \Psi^N \rangle$$

Self-consistent Green's functions



Carbone *et al.*, PRC 90, 054322 (2014)

The need for 3-body nuclear forces

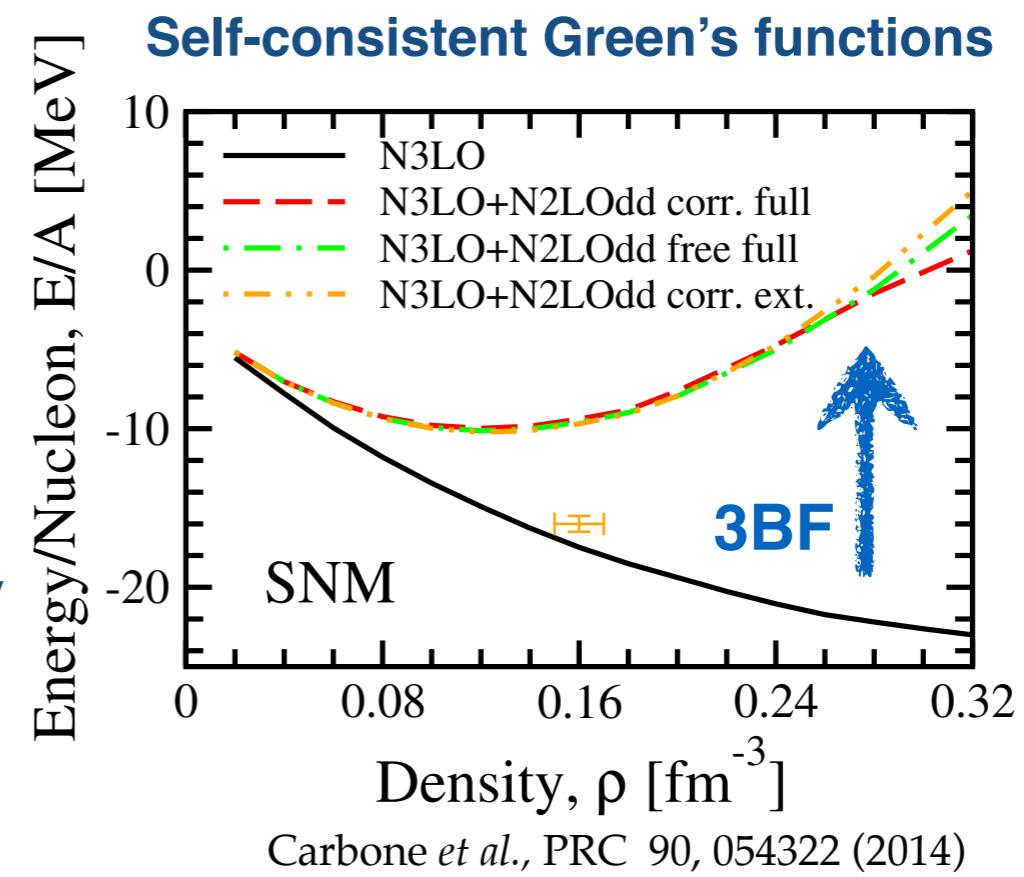


J.W. Holt *et al.*, PRC 81, 024002 (2010)
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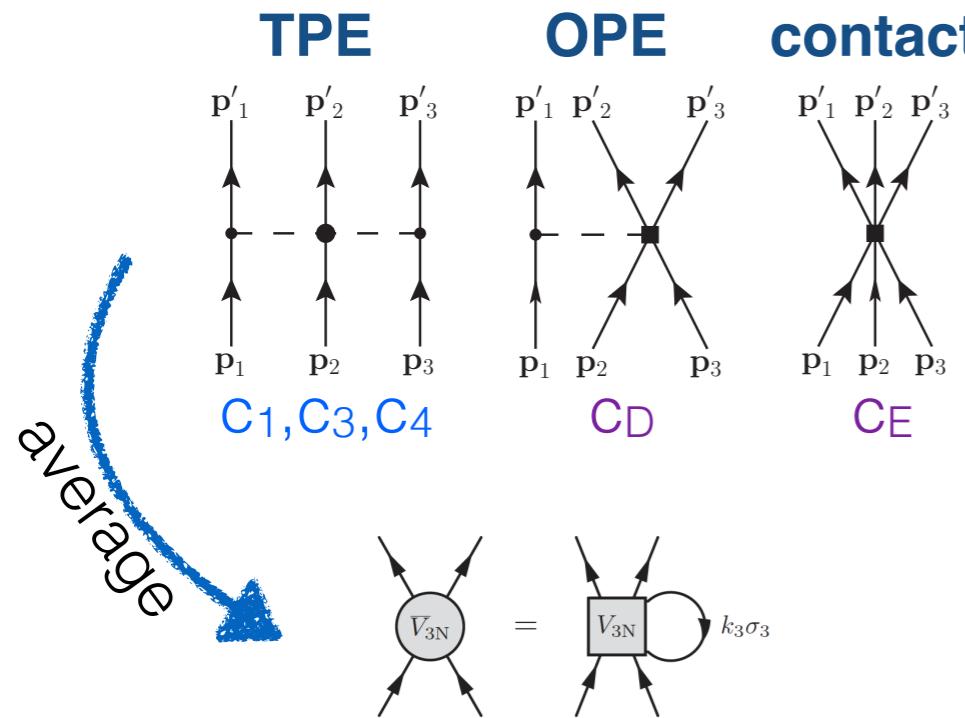
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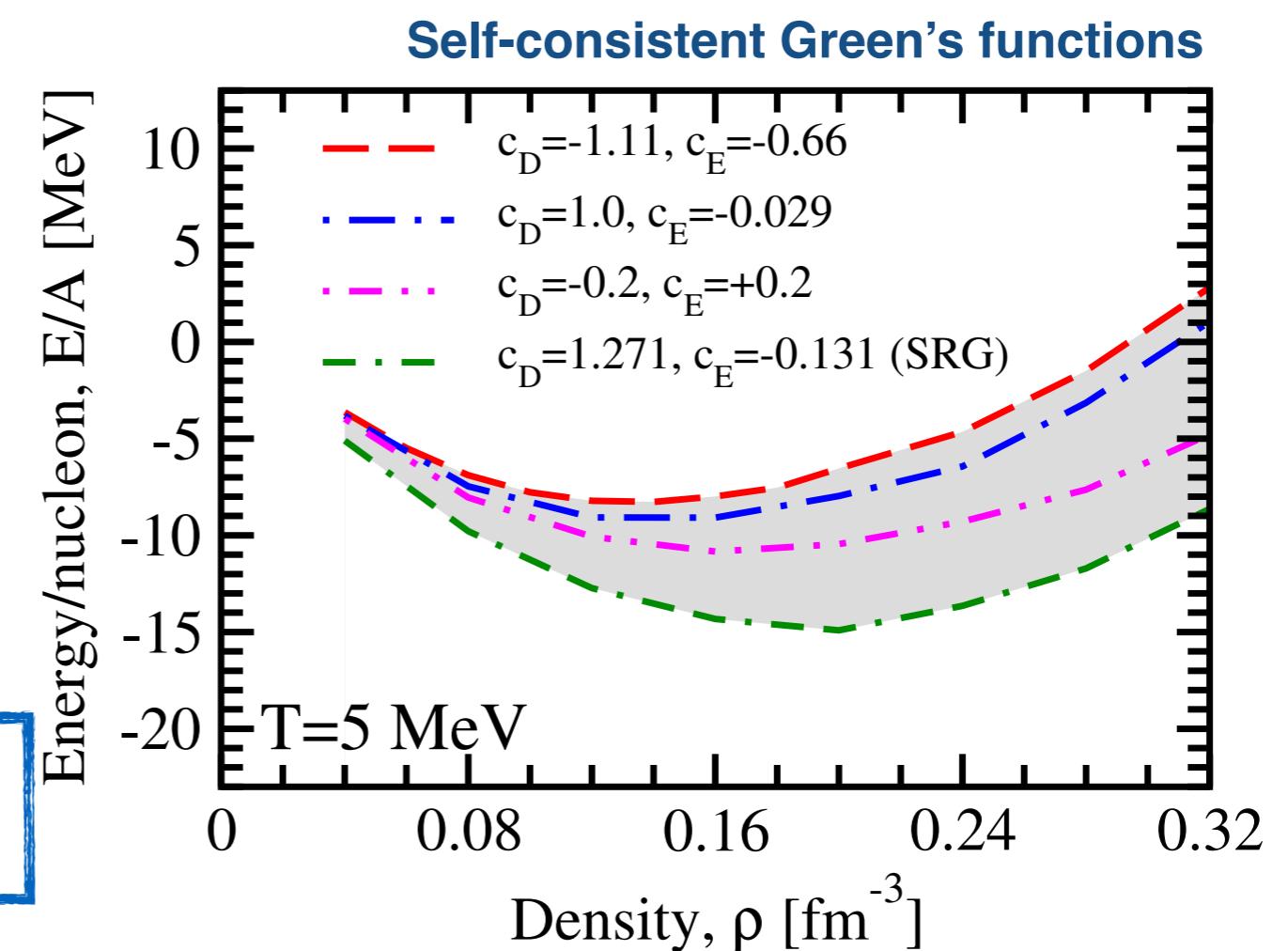
Theoretical uncertainties



J.W. Holt *et al.*, PRC 81, 024002 (2010)
 Hebeler *et al.*, PRC 82, 014314 (2010)
 Carbone *et al.*, PRC 90, 054322 (2014)

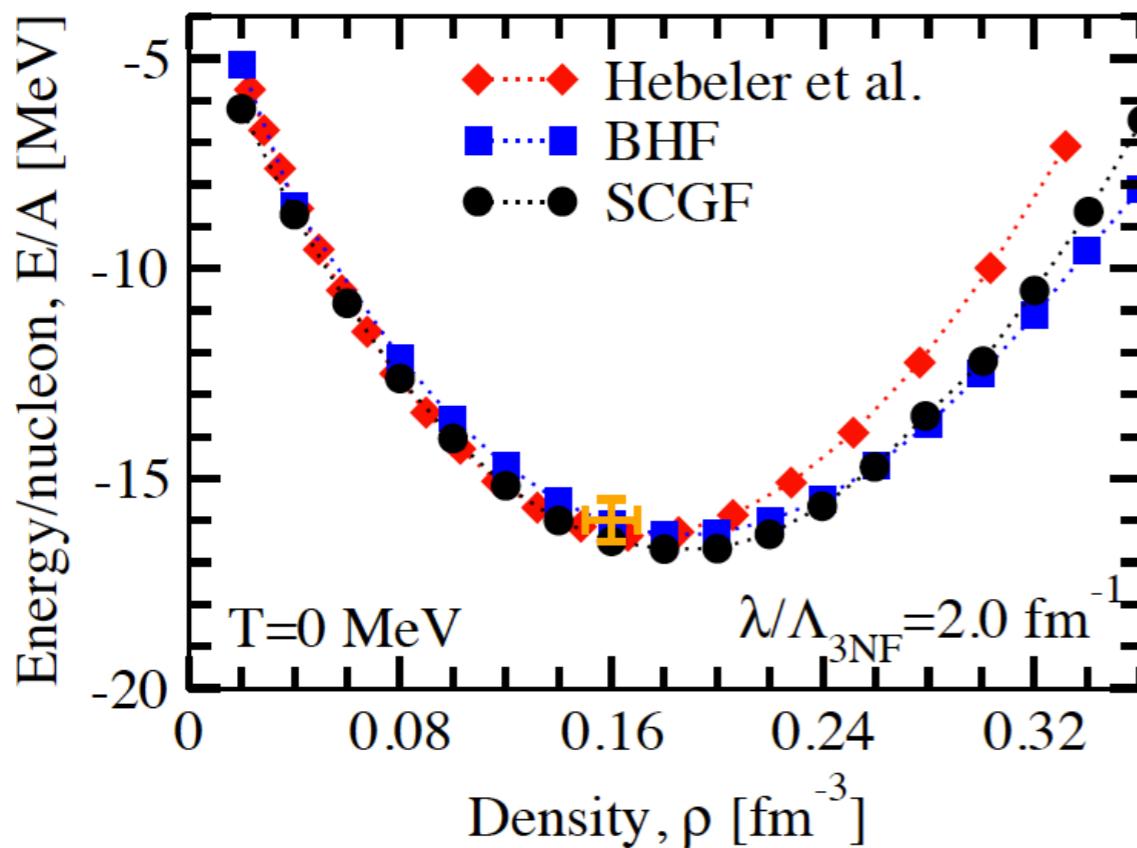
c_D and c_E fit to few-body properties

- Band represents the theoretical uncertainties
- The uncertainty dependence increases with density as the effect of 3NF



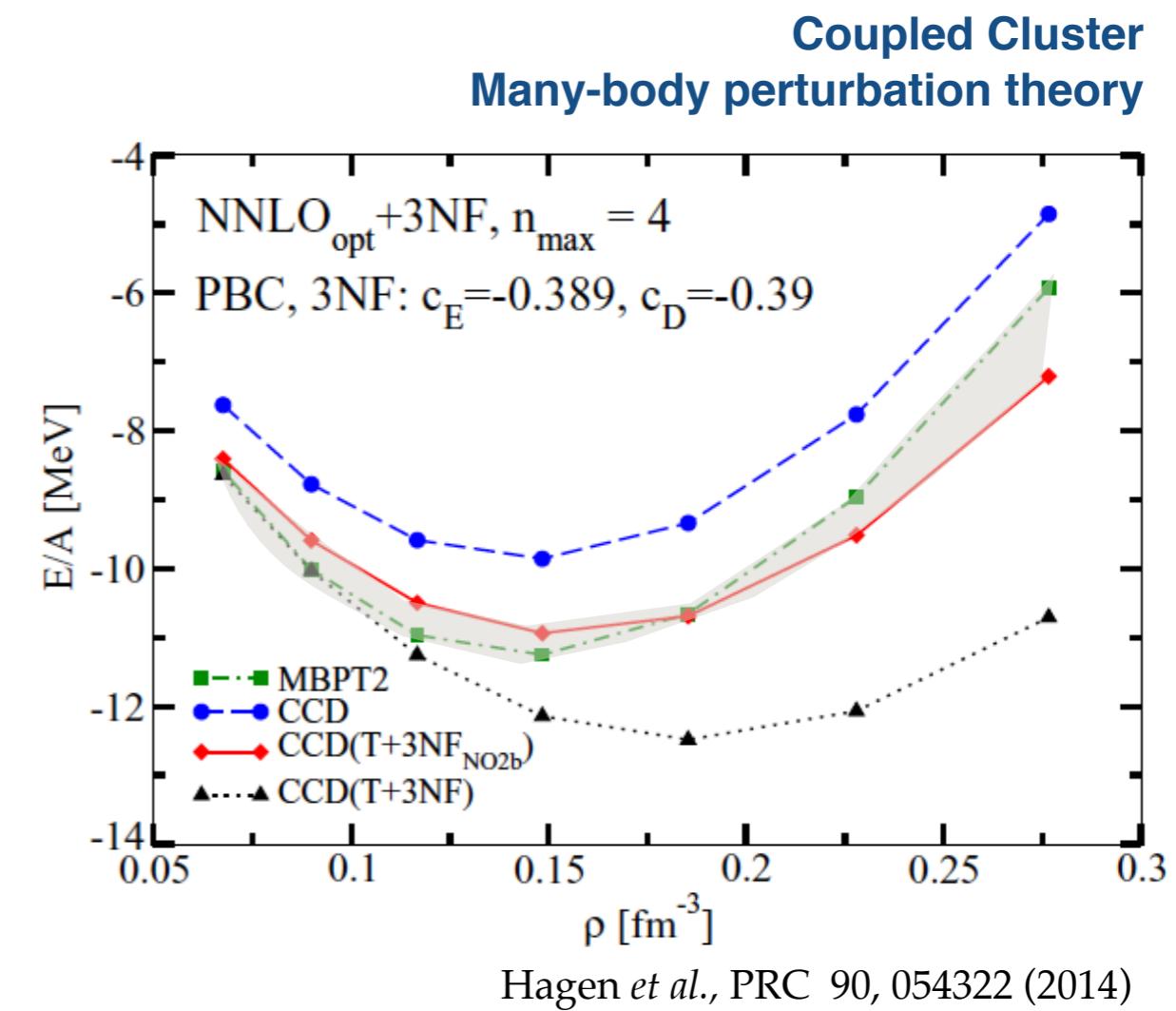
Many-body methods comparison

Many-body perturbation theory
Brueckner-Hartree-Fock
Self-consistent Green's functions



Carbone *et al.*, PRC 88, 044302 (2013)

Hebele *et al.*, PRC 83, 031301(R) (2011)



Further results:

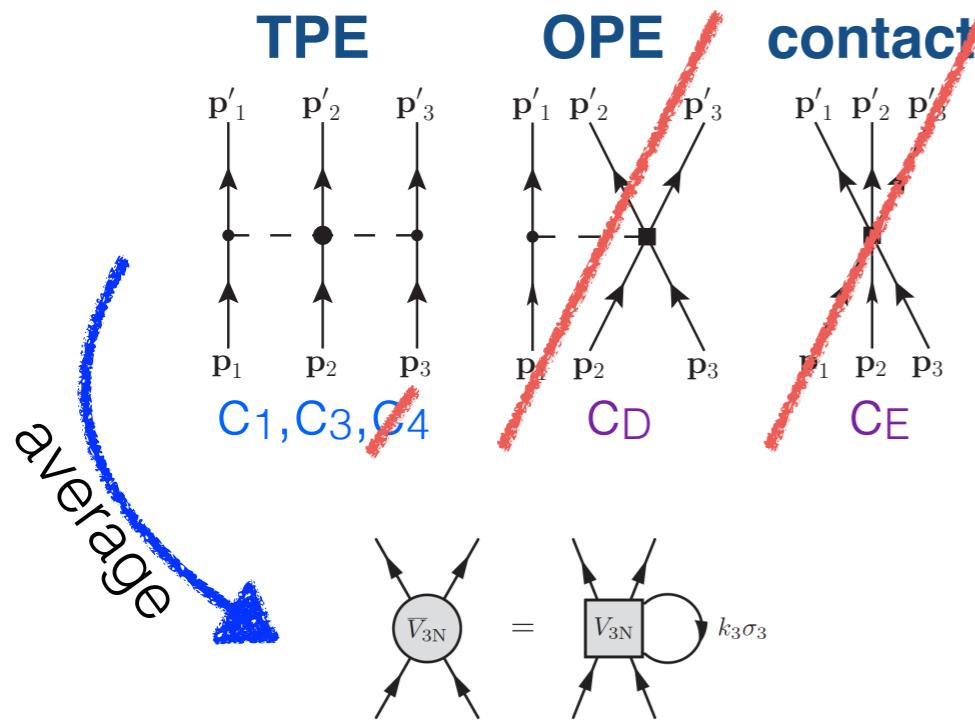
In-medium Chiral PT J.W. Holt *et al.*, PPNP 73, 35 (2013),

Lacour *et al.*, Ann. Phys. 326, 241 (2011)

MBPT Coraggio *et al.*, PRC 89, 044321 (2014)

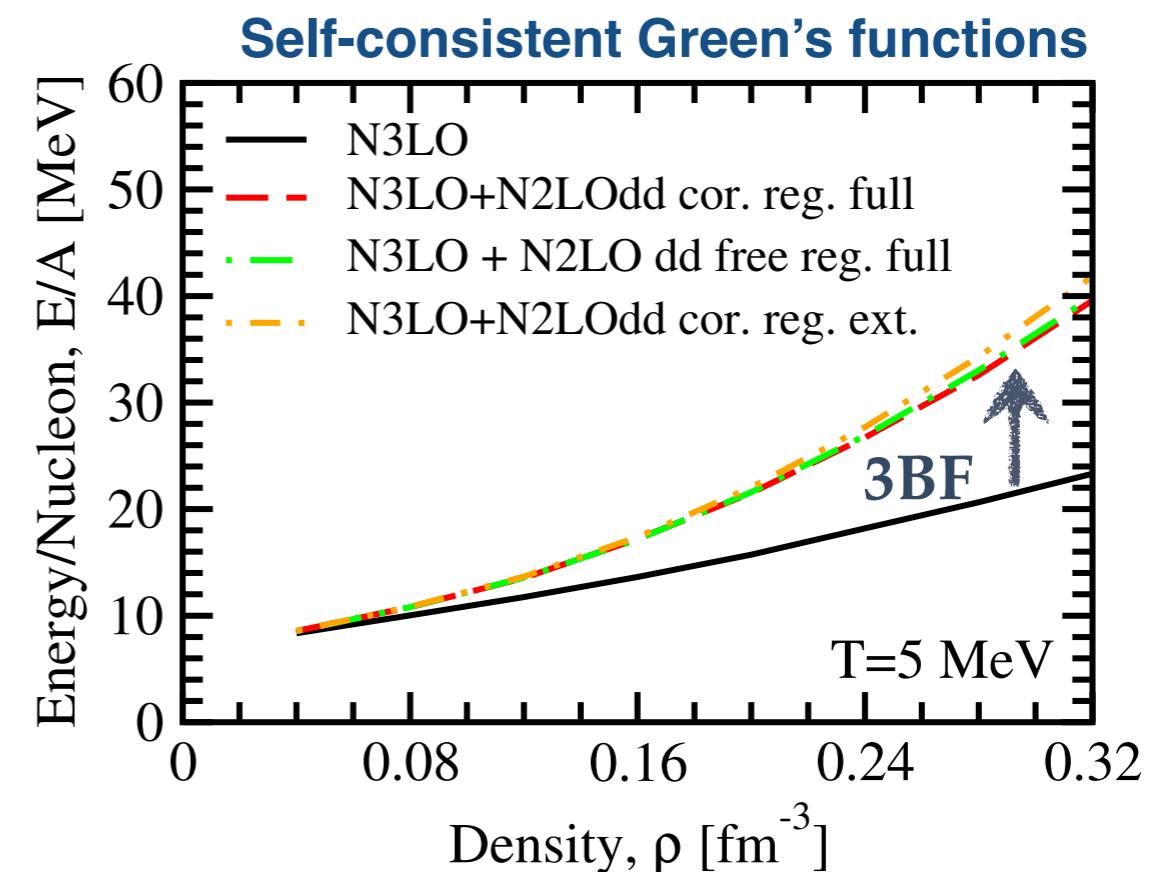
Uncertainty from the many-body calculations is smaller than the one estimated from chiral EFT

How neutron matter energy stiffens



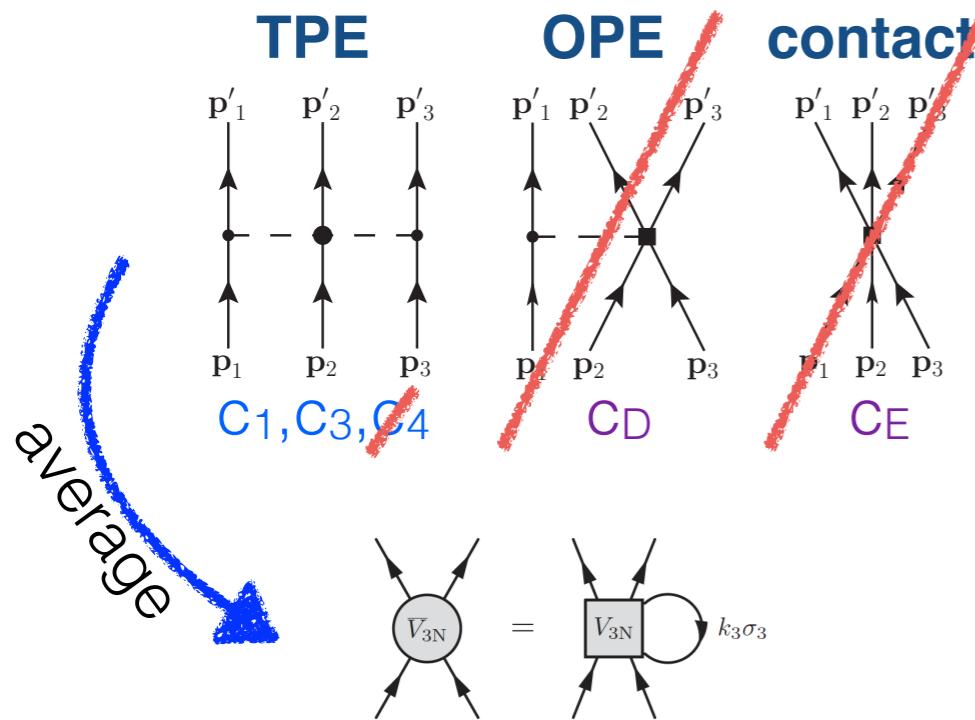
- Global repulsive effect due to 3bf
- Repulsion of 4 MeV at 0.16 fm^{-3} to 15 MeV at 0.32 fm^{-3}
- Small dependence on averaging procedures

- 3NFs fully predicted
- no need to fit to few-body properties



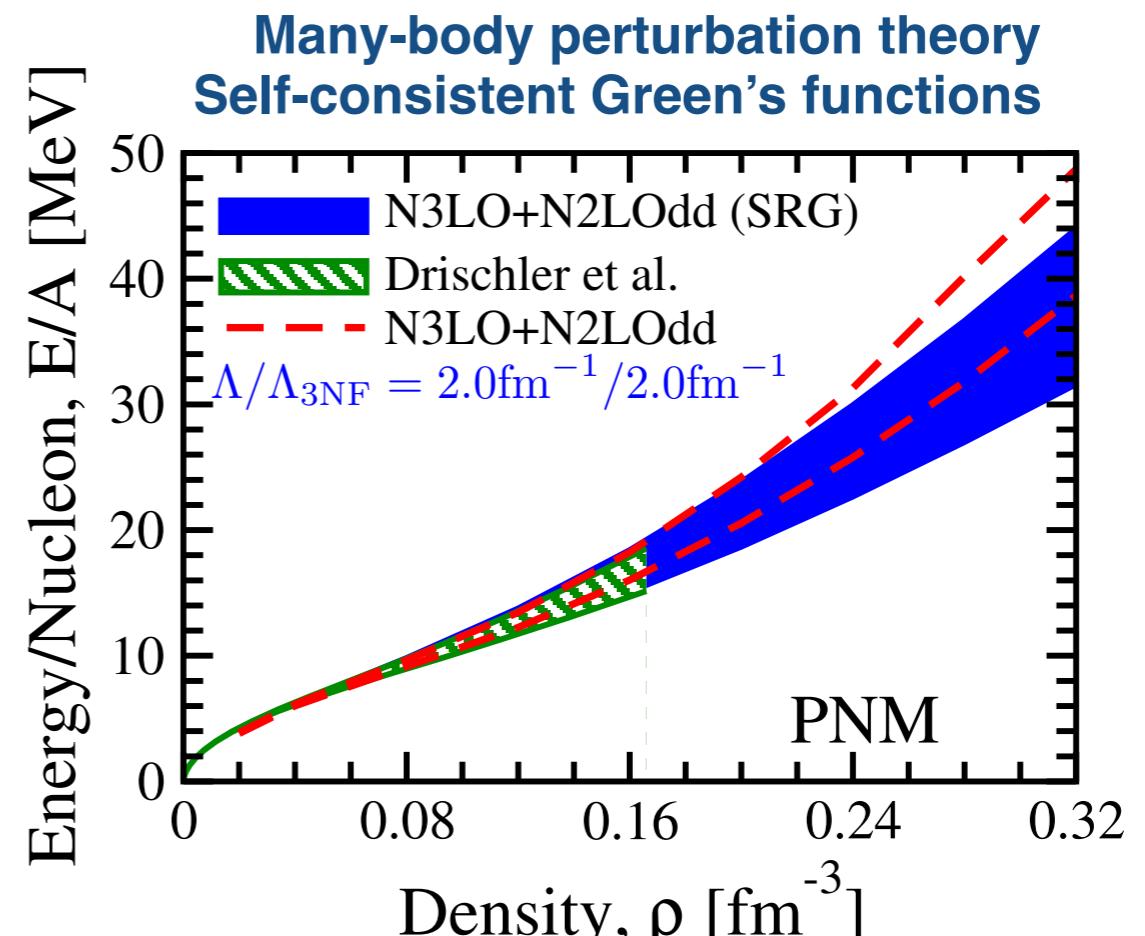
Carbone *et al.*, PRC 90, 054322 (2014)

Many-body methods comparison



- SCGF: agreement up to 0.20 fm^{-3} with the use of different Hamiltonians
- MBPT vs SCGF: agreement up to saturation density
- Low-density neutron matter perturbative

- Bands from c_1 and c_3 uncertainties

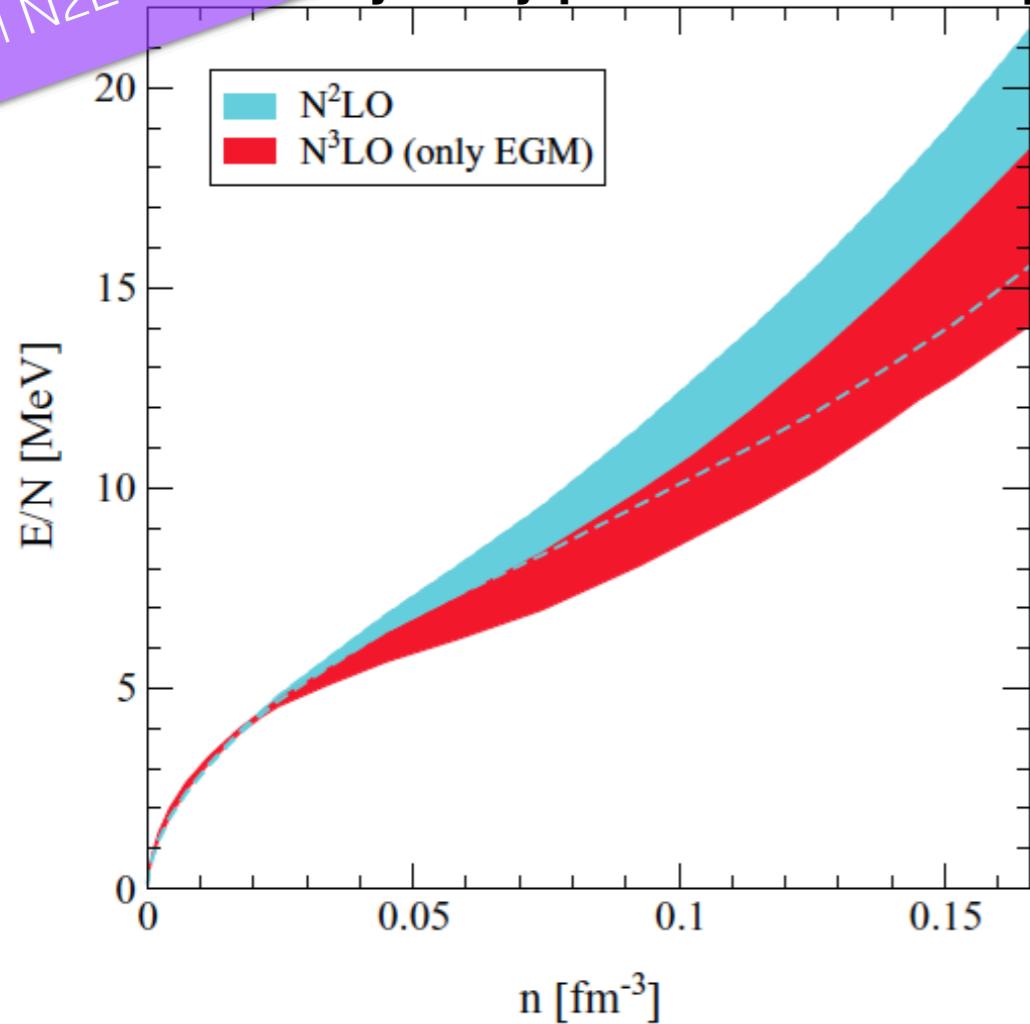


Carbone *et al.*, PRC 90, 054322 (2014)

Low-density neutron matter

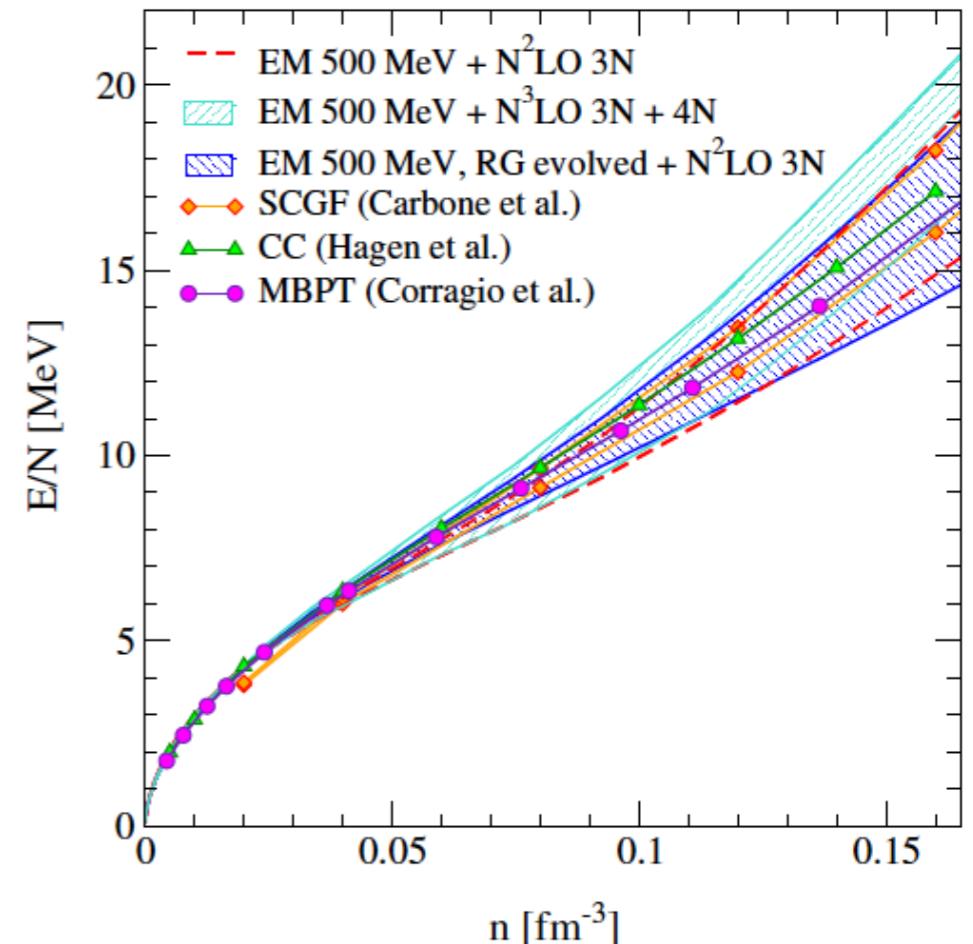
Full N²LO vs N³LO

Many-body perturbation theory



Krüger *et al.*, PRC 88, 025802 (2013)

Remarkable agreement between many-body methods and different Hamiltonians



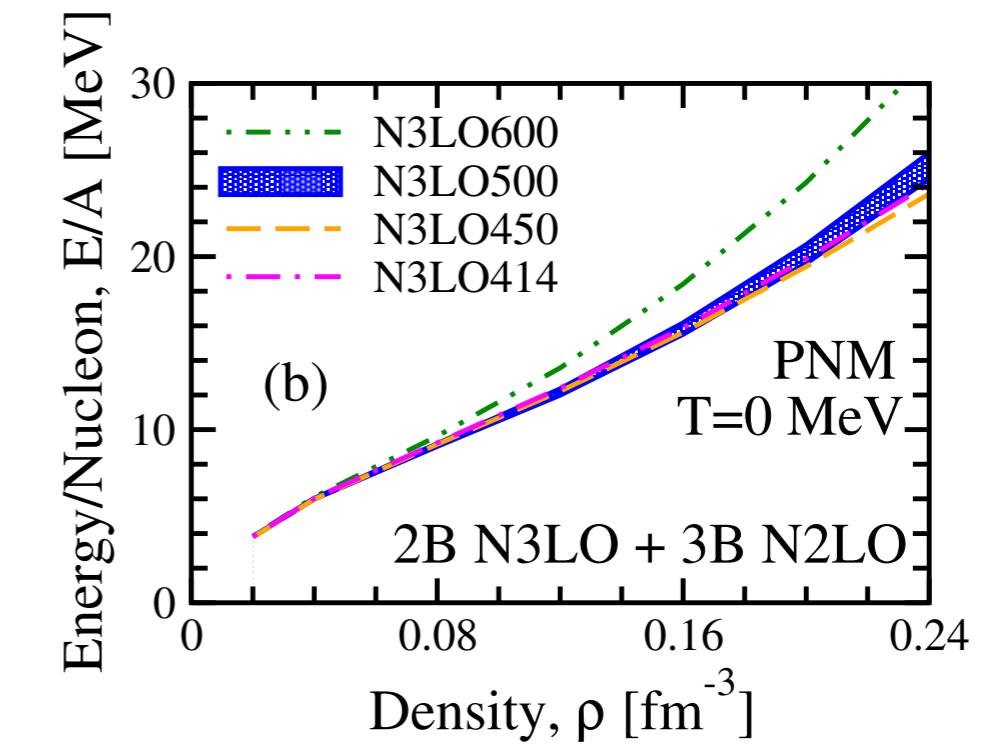
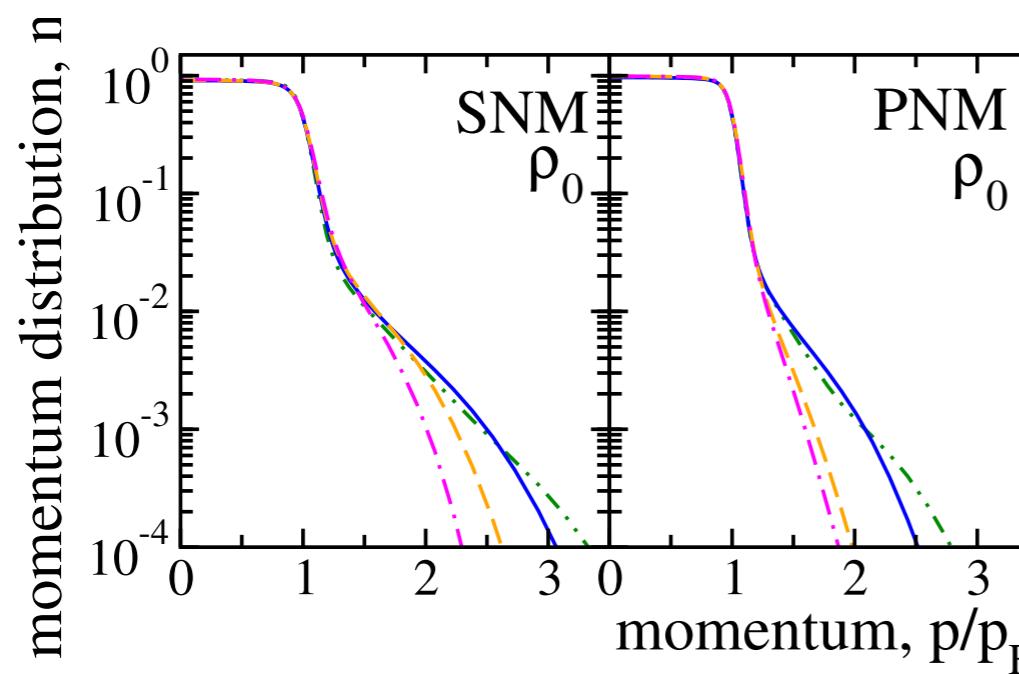
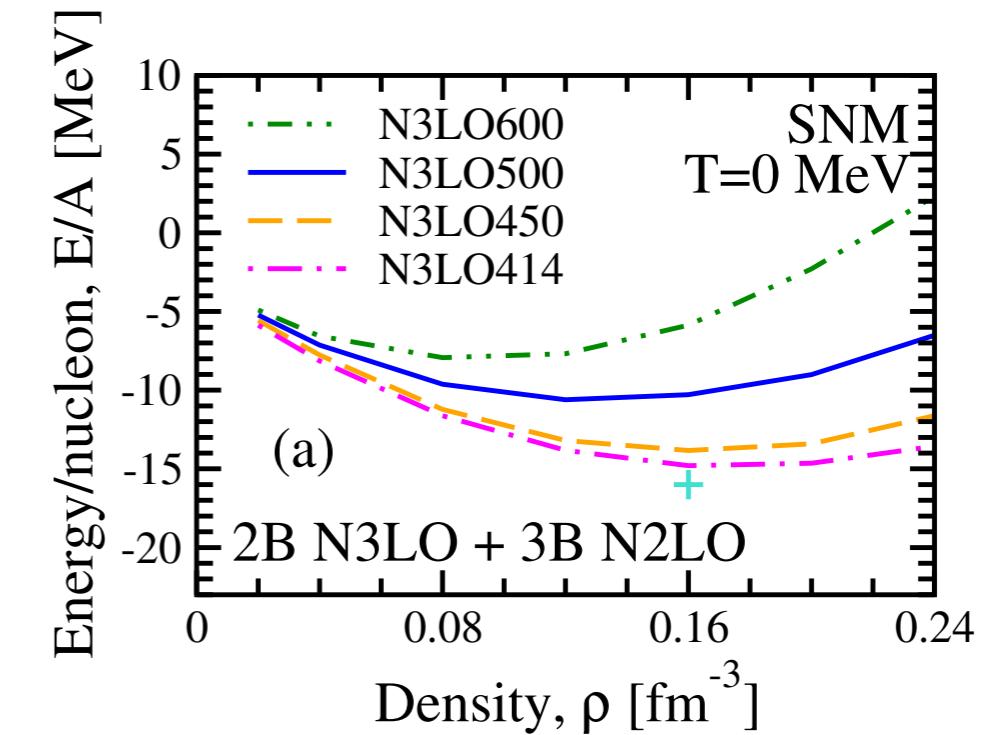
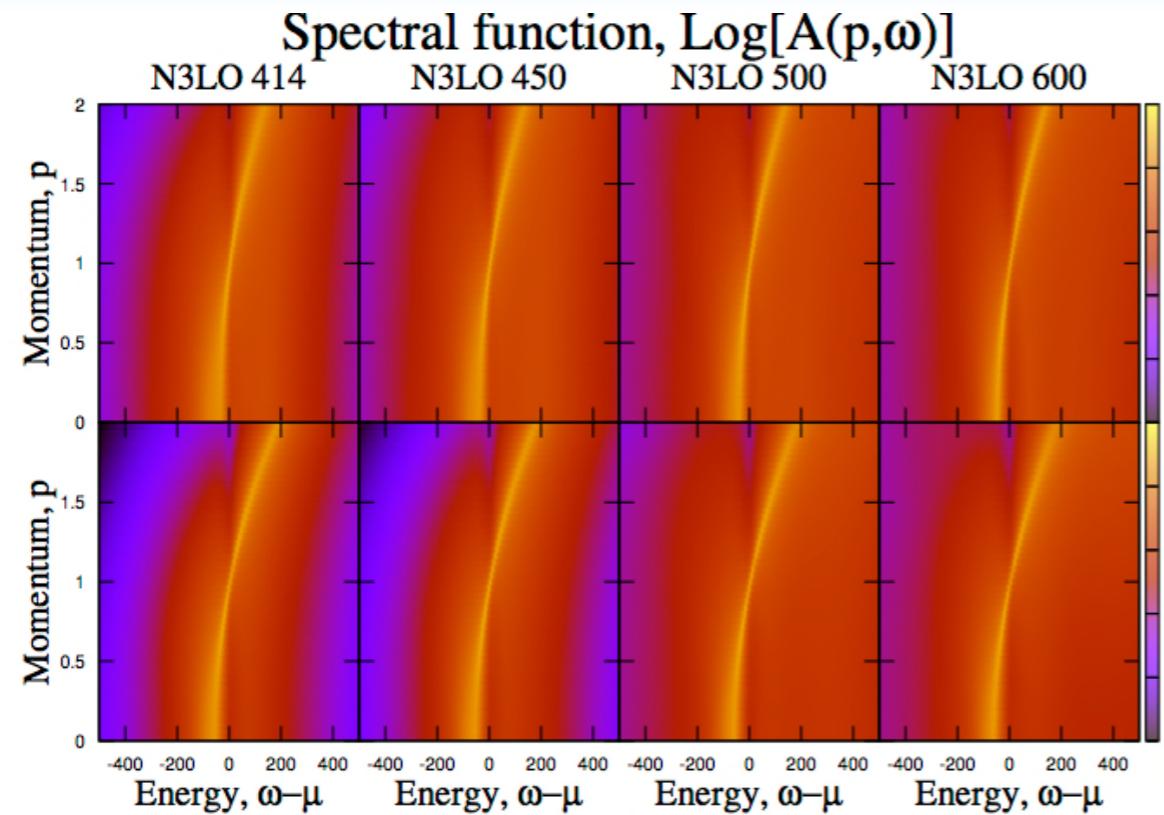
Hebeler *et al.*, Ann. Rev. Nucl. Part. Sci. 65, 457 (2015)

- 3NF N³LO subleading terms important
- Uncertainty band narrows
- Inclusion of explicit Δ will improve band

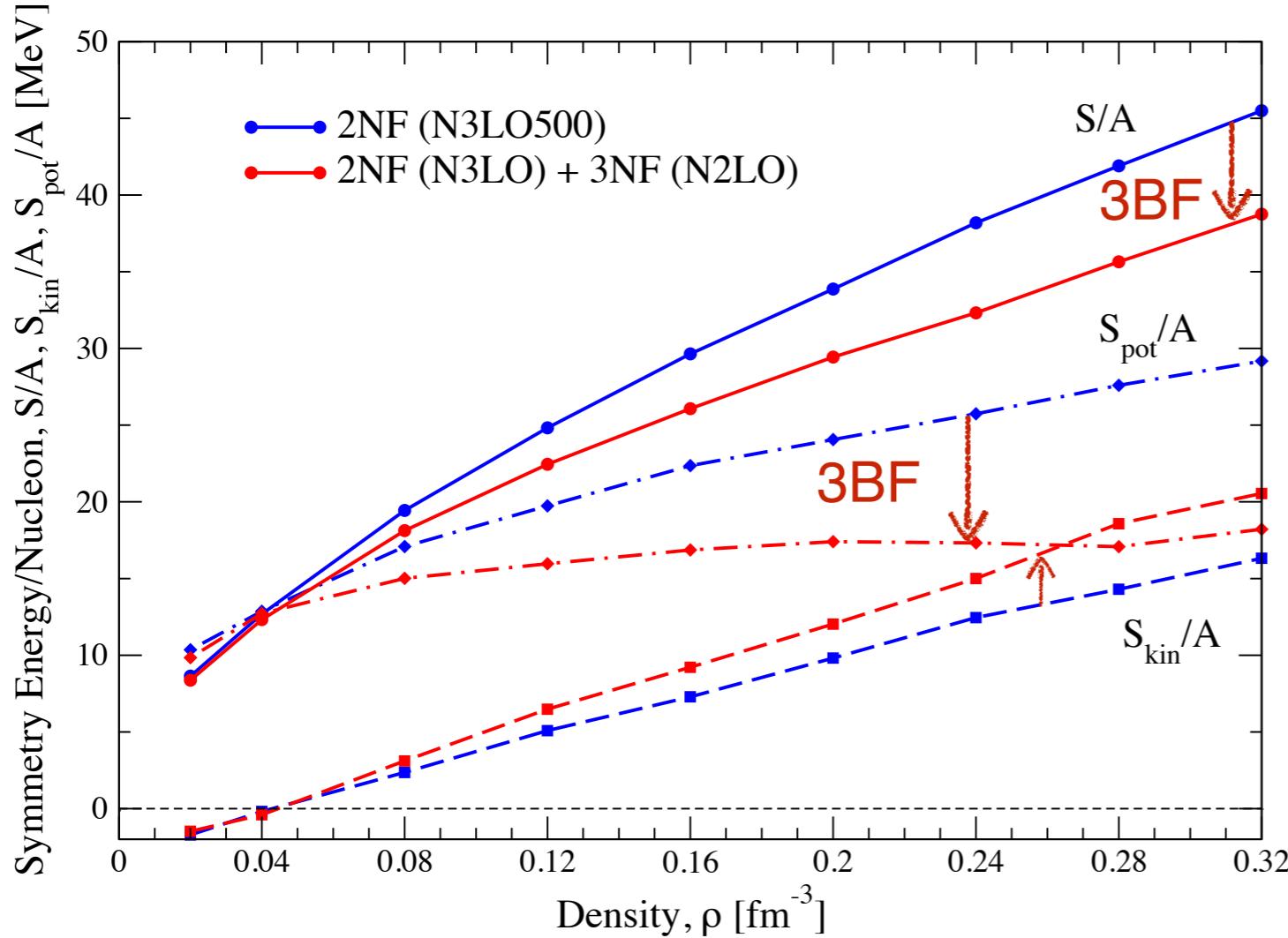
Further results:

AFDMC - Gezerlis et al., PRC 90, 054323 (2014)
 Lattice EFT - Epelbaum et al., EPJA 40, 199 (2009)
 In-medium Chiral PT - J.W. Holt et al., PPNP 73, 35 (2013),
 Lacour et al., Ann. Phys. 326, 241 (2011)

Cutoff dependence



The symmetry energy

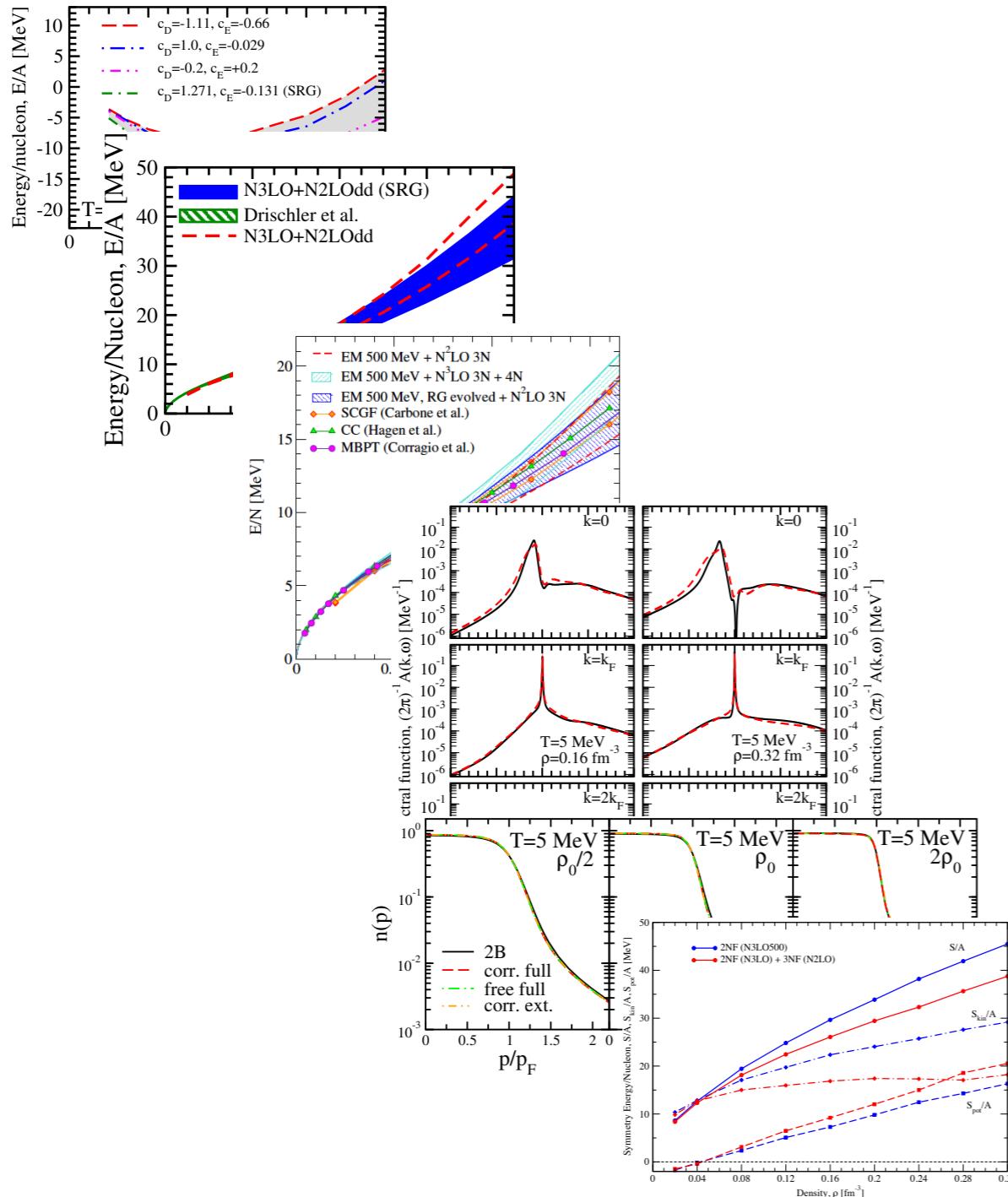


$$\frac{S}{A}(\rho) \sim \frac{E_{pnm}}{A}(\rho) - \frac{E_{snm}}{A}(\rho)$$

- 3NFs decrease symmetry energy
- Main effect from potential term
- The kinetic term stiffens
- 3NFs lowers high momentum components in SNM

$\rho_0 = 0.16 \text{ fm}^{-3}$	S_{tot}	E_{pnm}	E_{snm}	S_{kin}	K_{pnm}	K_{snm}	S_{pot}	P_{pnm}	P_{snm}
2NF [MeV]	29.64	13.64	-16.00	7.29	40.17	32.88	22.36	-26.53	-48.89
2NF+3NF [MeV]	26.07	17.20	-8.87	9.21	41.21	32.00	16.86	-24.00	-40.87

Summary



- A consistent extension of the SCGF method to include 3NFs is accomplished
- Nuclear and neutron matter with theoretical uncertainties can be calculated reliably using *ab initio* methods based on chiral Hamiltonians
- Uncertainties in the many-body methods are under control
- Small overall effect of 3NFs on the momentum distribution
- 3NFs lowers the correlations in SNM

Outlook

- Chiral EFT Hamiltonian: power counting, theoretical uncertainties, limits of chiral EFT, etc.
- Many-body approximation methods: include irreducible 3B terms, improve the effective interactions, include particle-hole diagrams, asymmetric matter, etc.
- Reliable finite temperature results from *ab initio* theory: high astrophysical impact (EOSs at finite T, dynamics of neutron star merger simulations, core-collapse supernovae, etc.)

Collaborators:



A. Polls



C. Barbieri
A. Rios



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UNIVERSITÄT
DARMSTADT

C. Drischler, T. Krüger, I. Tews, P. Klos
K. Hebeler, A. Schwenk

Thank you for your attention!

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