

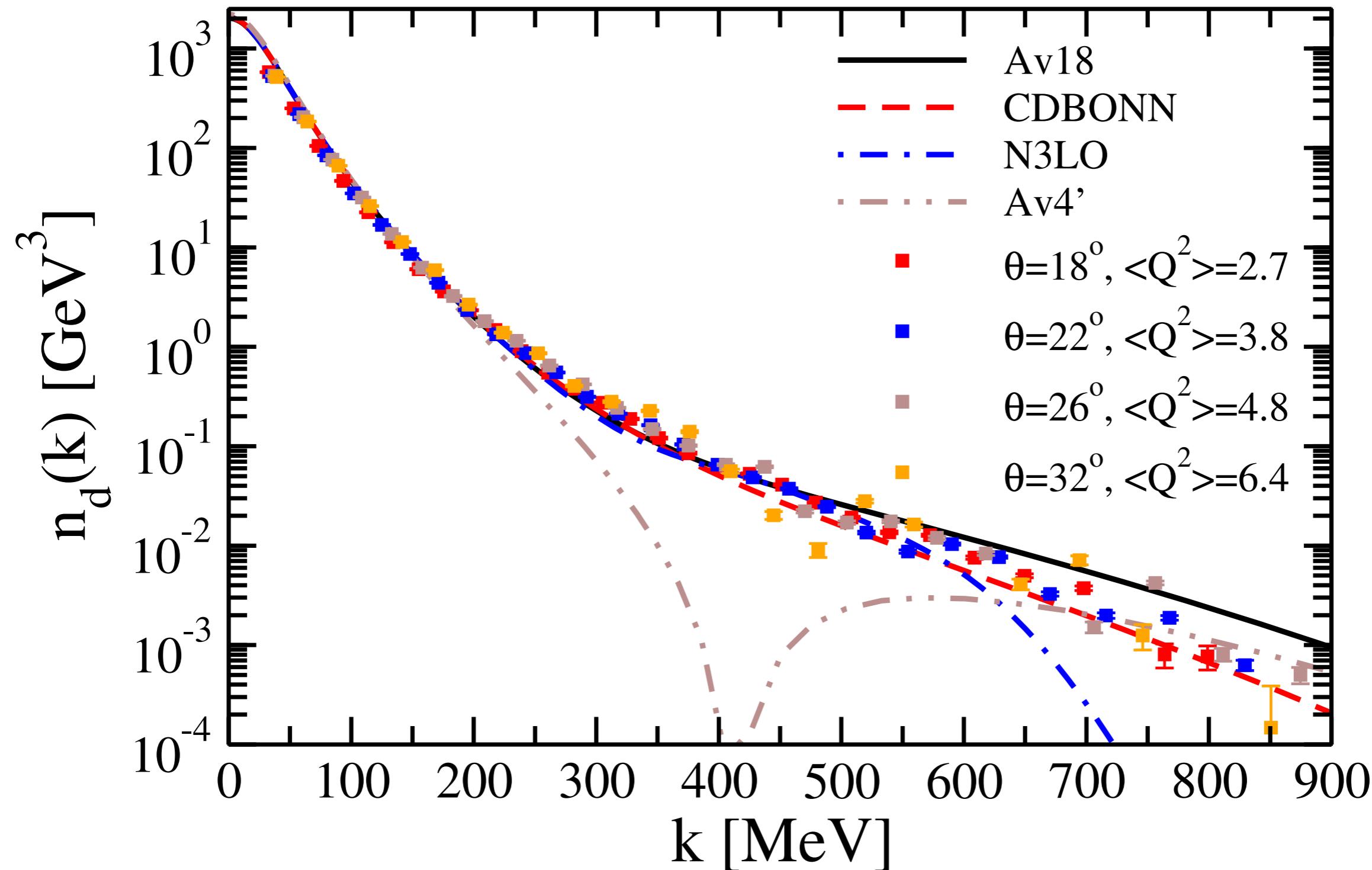
Density & isospin asymmetry dependence of correlations

Arnau Rios Huguet
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Department of Physics
University of Surrey

1. Depletion of the nuclear Fermi sea
2. High-momentum components
3. Implications for symmetry energy
4. Pairing in neutron stars

Why short-range correlations?

Deuteron momentum distribution

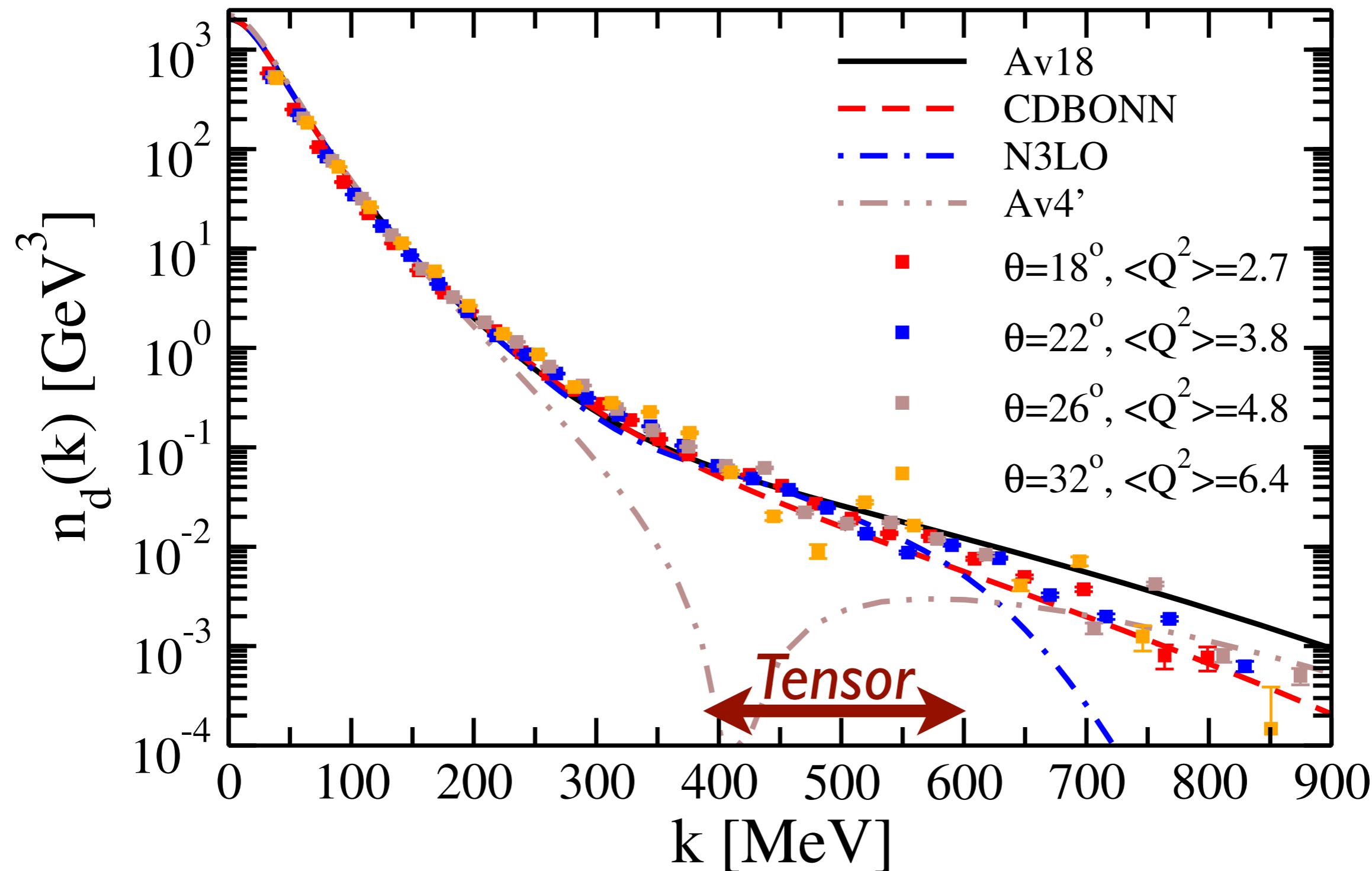


Fomin et al., PRL 108 092502 (2012)

Inclusive quasi-elastic e^- scattering vs NN potential theory

Why short-range correlations?

Deuteron momentum distribution

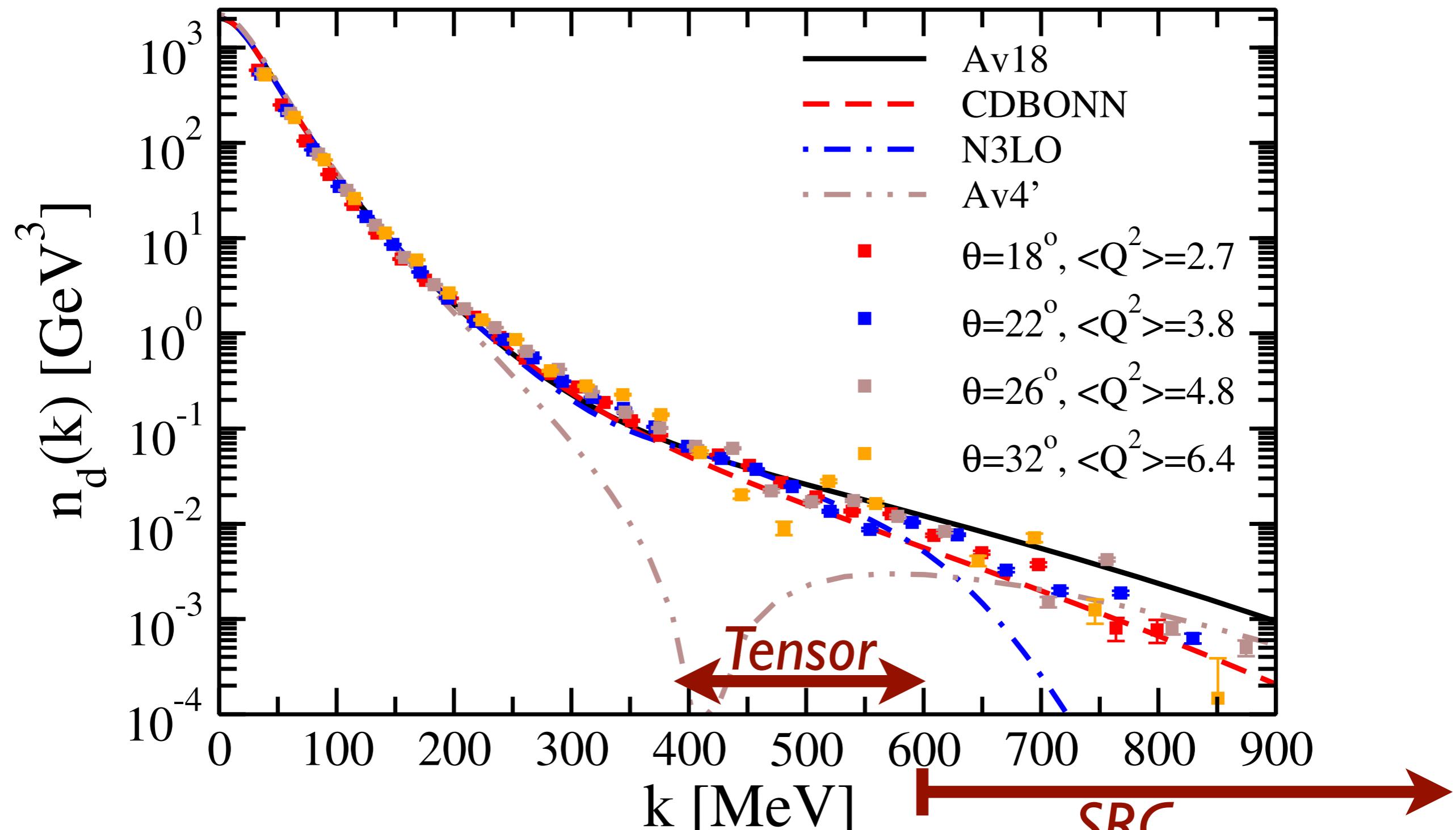


Fomin et al., PRL 108 092502 (2012)

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Fomin et al., PRL 108 092502 (2012)

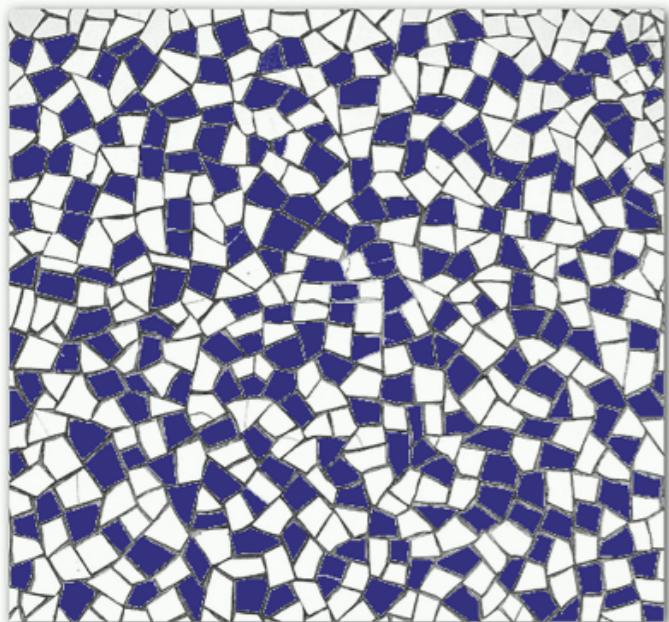
Inclusive quasi-elastic e^- scattering vs NN potential theory

Nuclear vs neutron matter

Nuclear “trencadís”

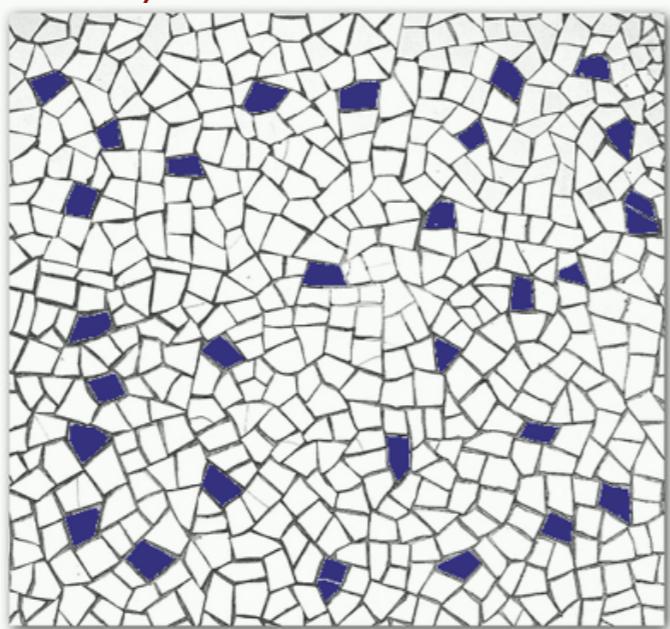
$N=Z$, $\beta=0$, $x_p=0.5$

Symmetric matter



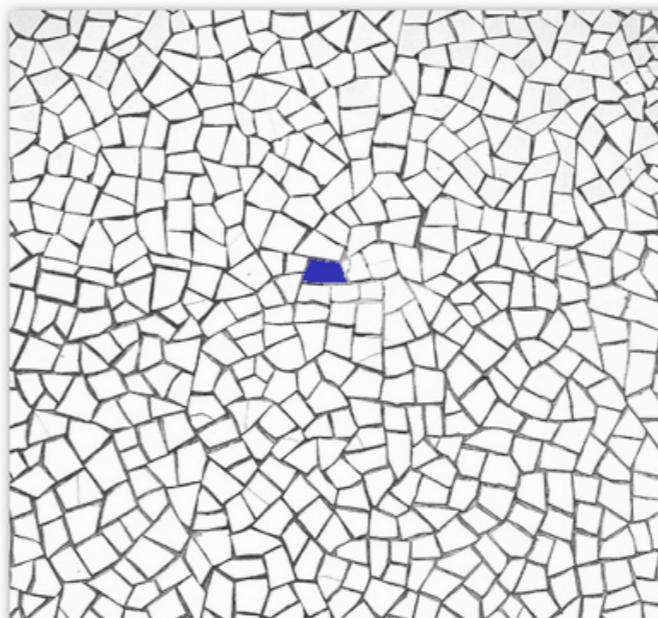
$\beta \neq 0$, $x_p < 0.5$

Asymmetric nuclei



$\beta \approx 1$, $x_p \approx 0$

Polaron

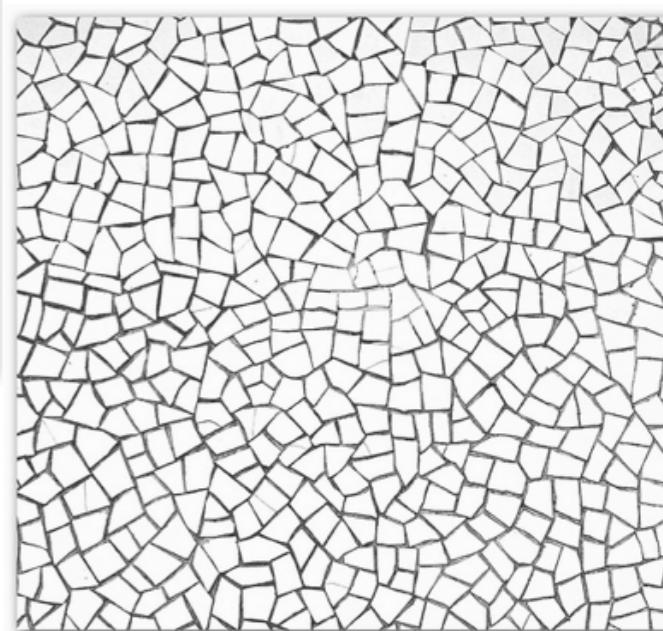


Neutron stars

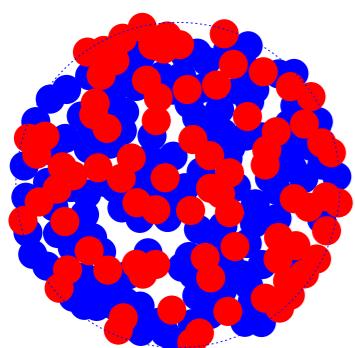


$\beta = 1$, $x_p = 0$

Neutron matter



Nuclei



Lead 208

$$\beta = (126 - 82) / 208 = 0.2$$
$$x_p = 82 / 208 = 0.39$$

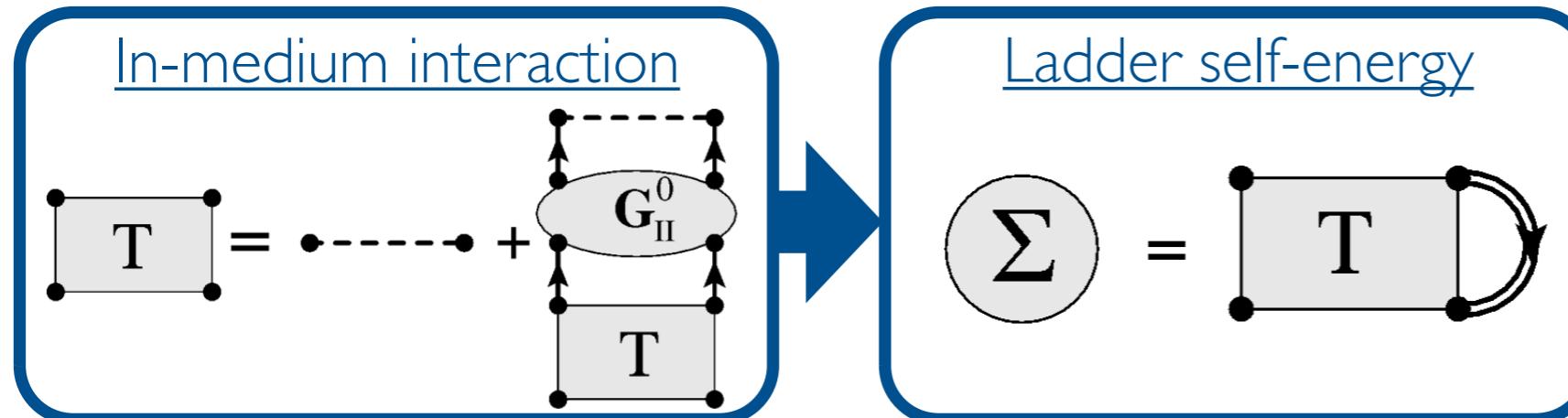
$$\beta = \frac{N - Z}{N + Z}$$

$$x_p = \frac{Z}{N + Z}$$

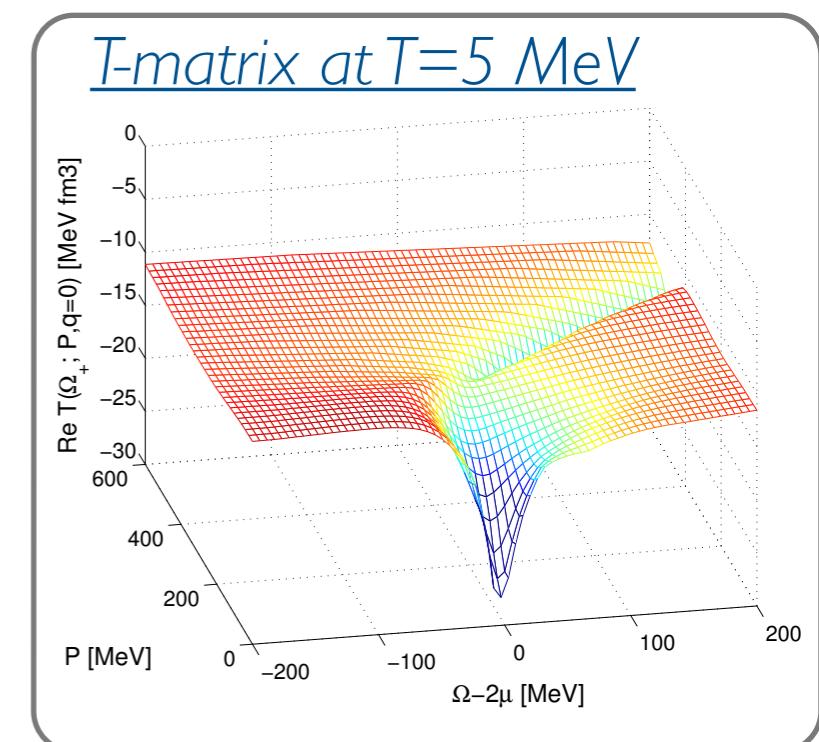
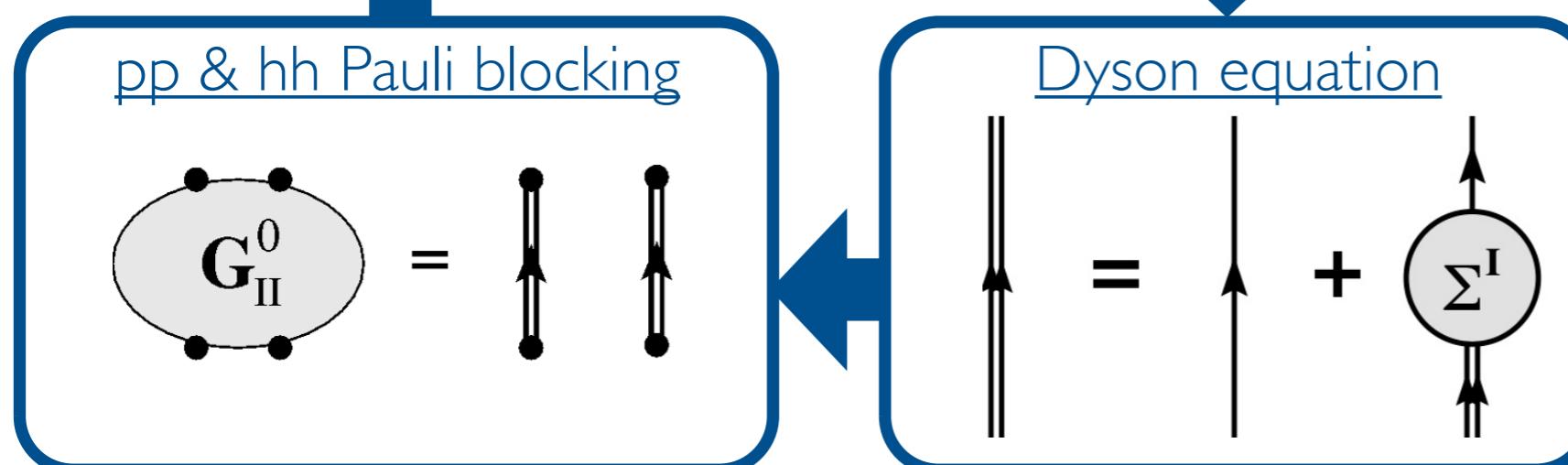
Experimentally unknown!



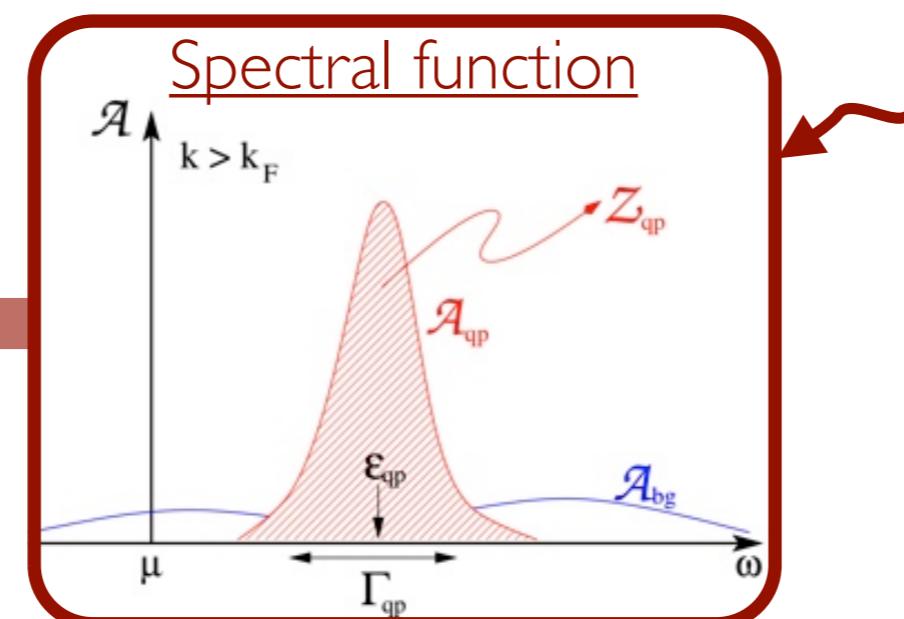
SCGF Ladder approximation



- Self-consistent resummation
- Energy and momentum integral
- @Finite T (Matsubara)

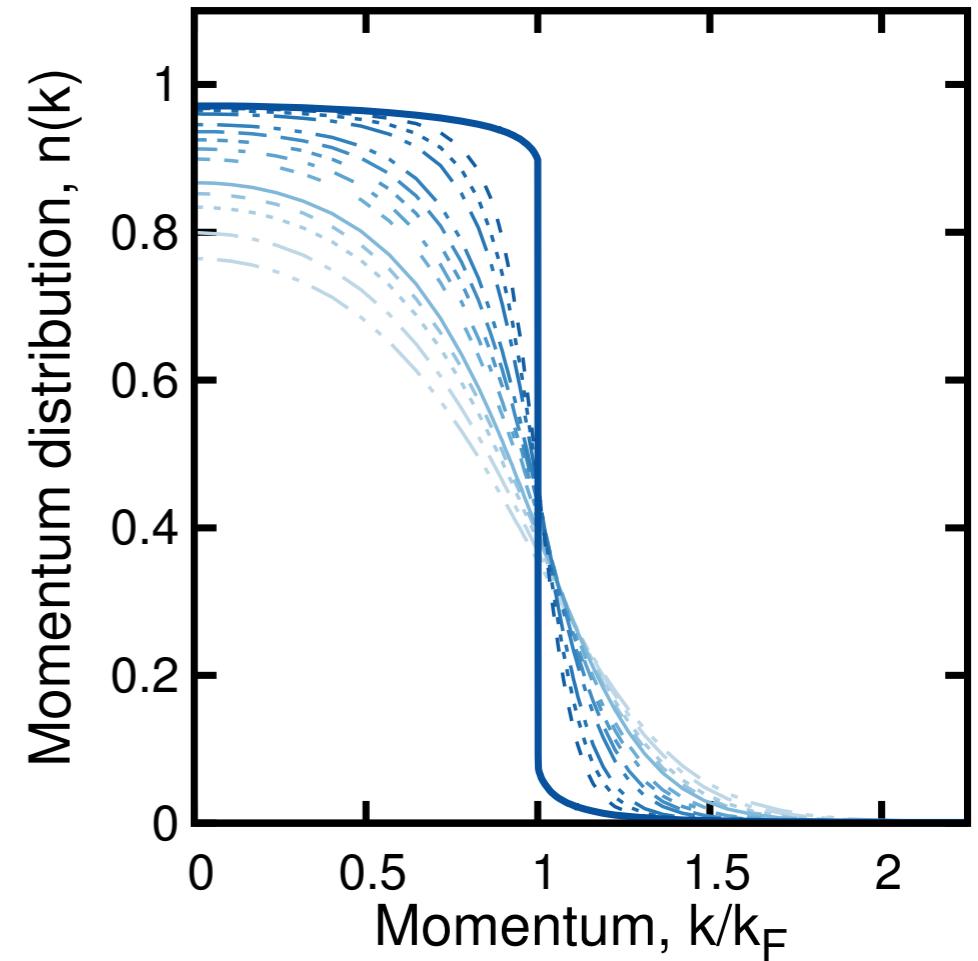
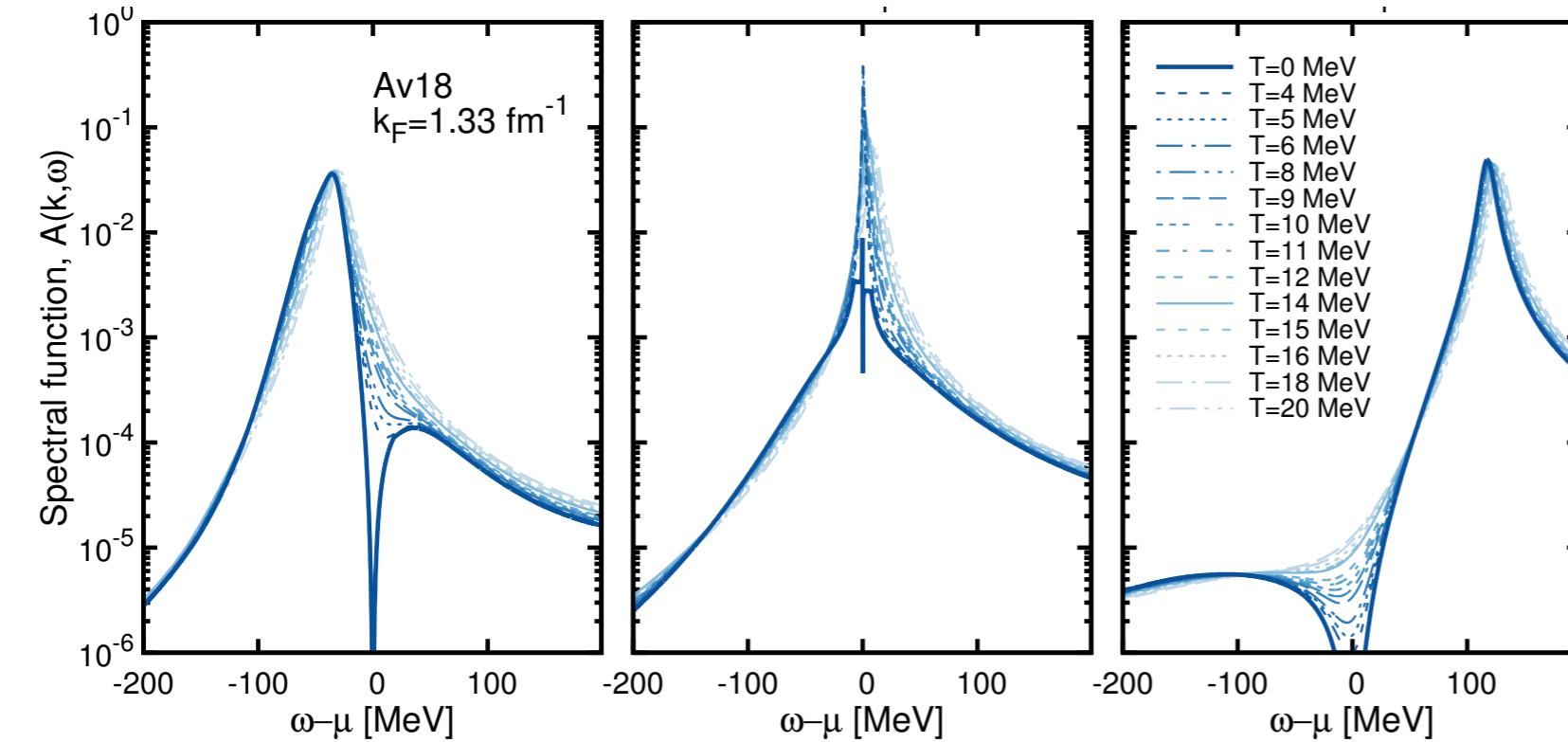
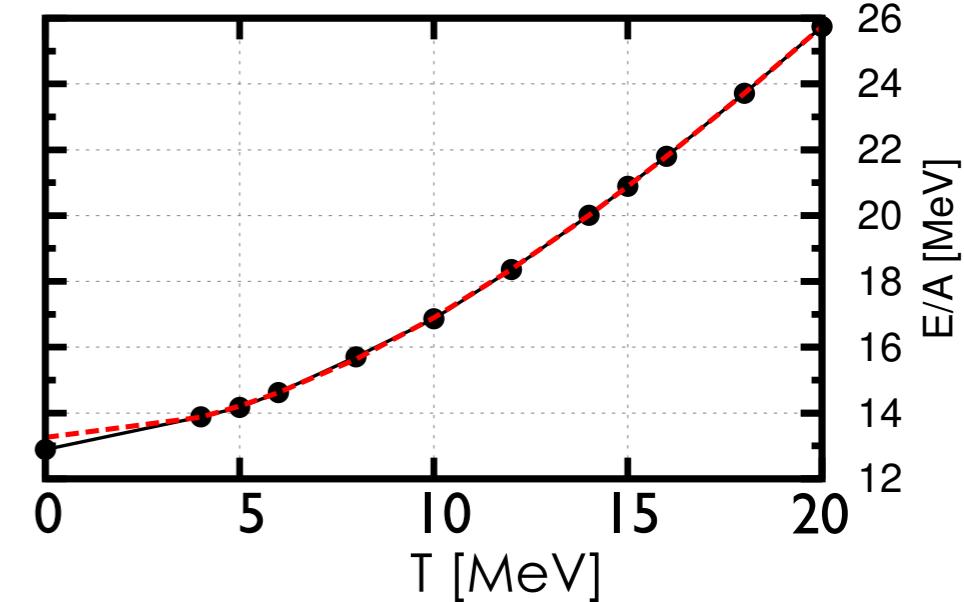
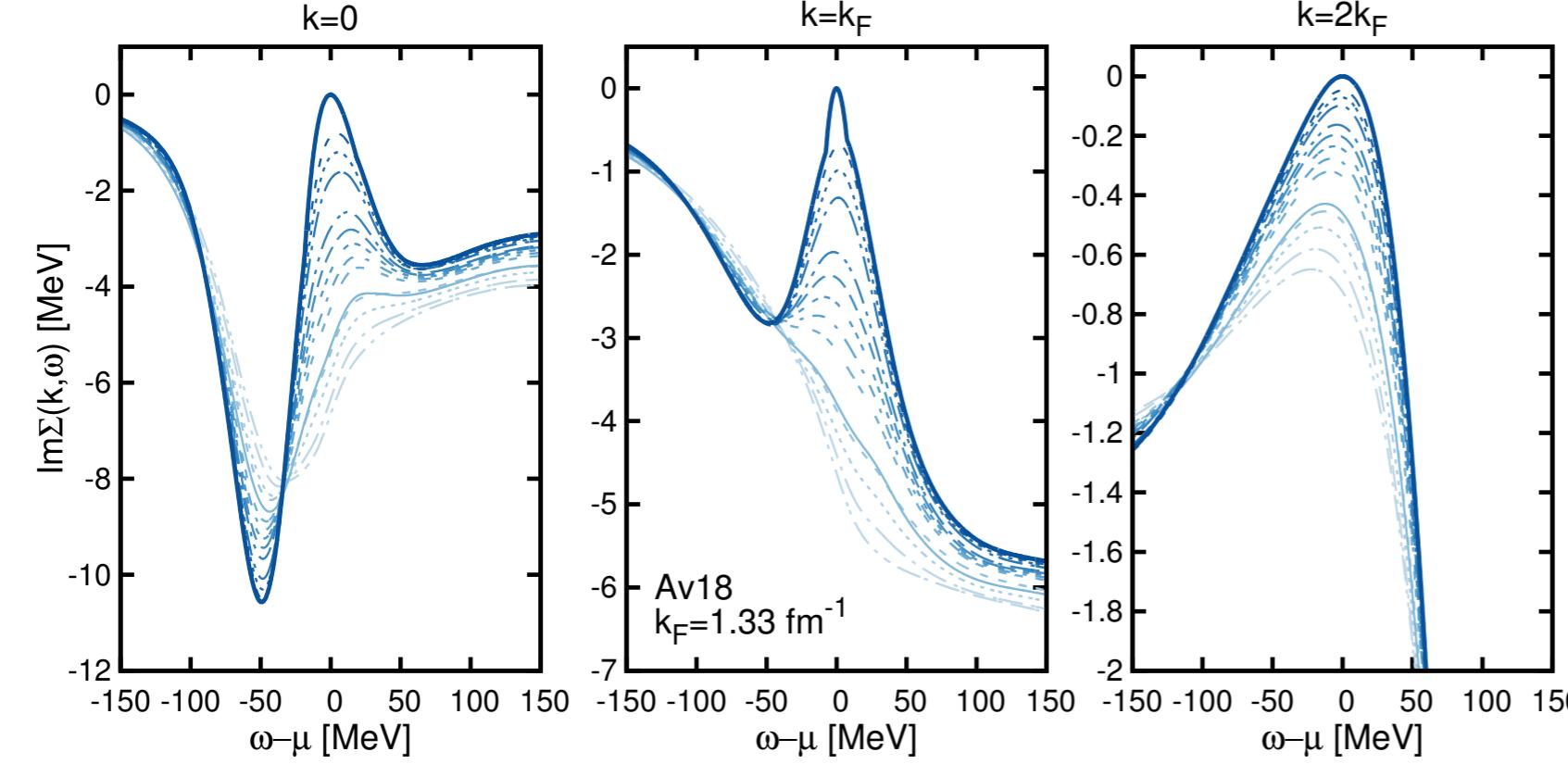


One-body properties
Momentum distribution
Thermodynamics & EoS
Transport



- Ramos, Polls & Dickhoff, NPA **503** 1 (1989)
 Alm et al., PRC **53** 2181 (1996)
 Dewulf et al., PRL **90** 152501 (2003)
 Frick & Muther, PRC **68** 034310 (2003)
 Rios, PhD Thesis, U. Barcelona (2007)
 Soma & Bozek, PRC **78** 054003 (2008)
 Rios & Soma PRL **108** 012501 (2012)

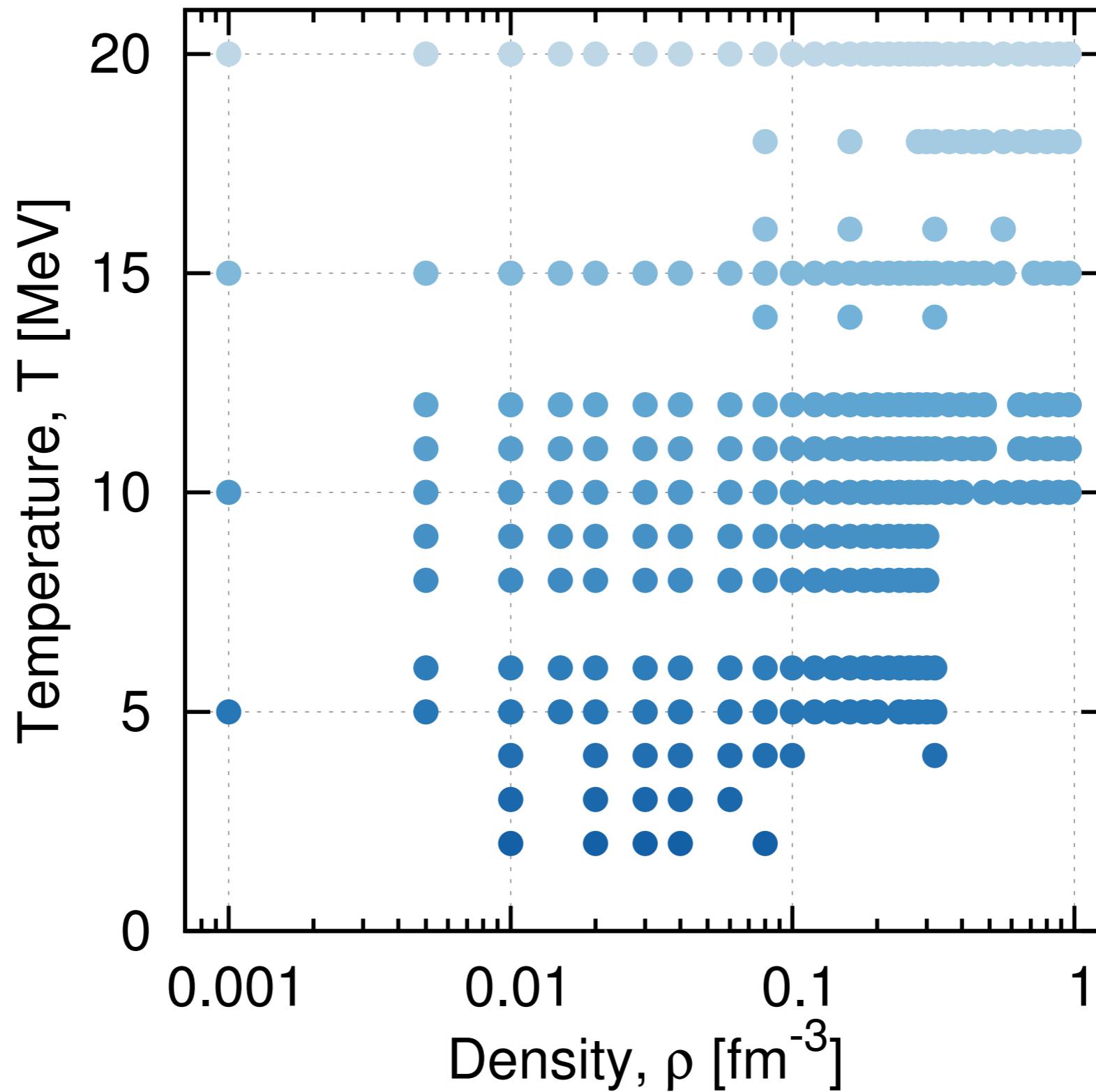
$T=0$ extrapolations



- **Extrapolate** $\Sigma(k, \omega; T)$ (13 Gb worth data)
- **Constraining with macroscopic** properties

Available data

Av18



Self-energy, spectral function & thermodynamics

- **Advantages**

- All kinds of NN interactions ✓
- **3N** interactions too (Arianna's talk) ✓
- Short-range & tensor **correlations** ✓
- Density & **isospin** dependence ✓
- Access to **off-shell** spectral function ✓
- No artificial **separation** of MF vs correlations ✓
- **Thermodynamically** consistent ✓

- **Limitations**

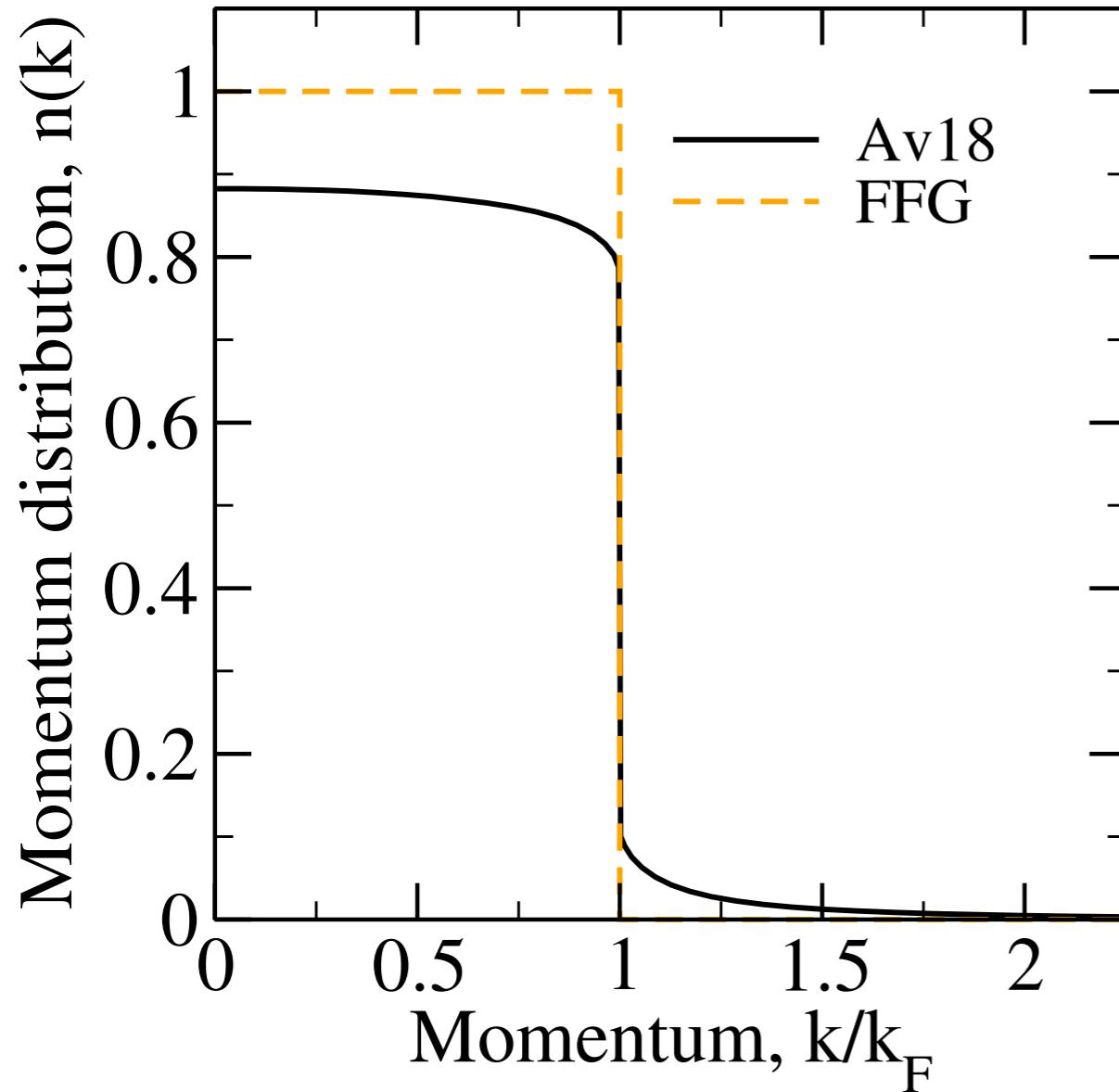
- Non-relativistic ✗
- Missing diagrams? ✗
- No real-space picture (e.g. correlators) ✗
- No nuclear surface ✗

Momentum distribution

Single-particle occupation

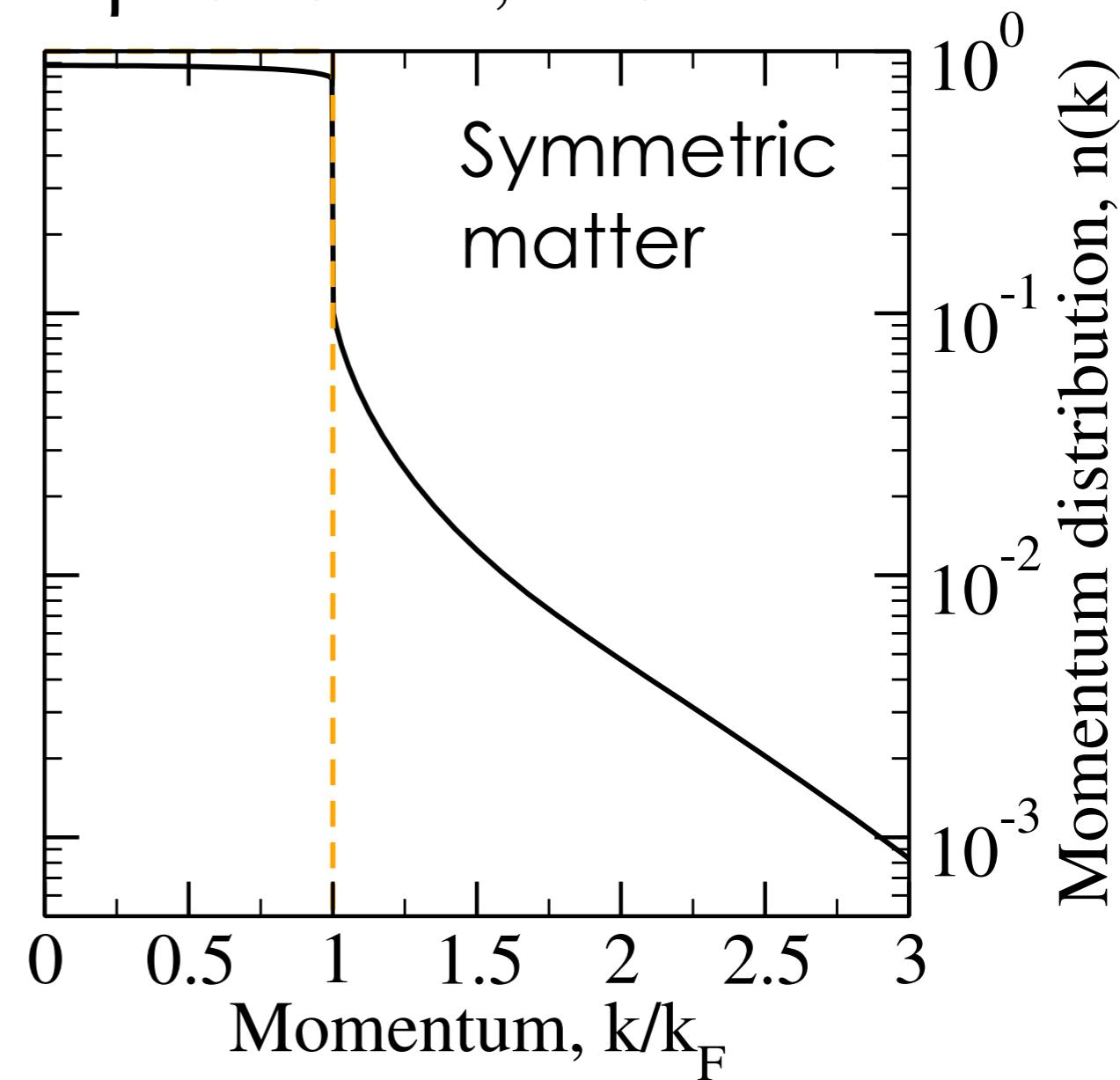
$$n(k) = \langle a_k^\dagger a_k \rangle$$

$\rho=0.16 \text{ fm}^{-3}$, $T=0 \text{ MeV}$



$$\nu \int \frac{d^3k}{(2\pi)^3} n(k) = \rho$$

$\rho=0.16 \text{ fm}^{-3}$, $T=0 \text{ MeV}$



- SNM: 11-13% depletion at low k , population at high k
- Dependence on NN interaction under control
- PNM: 4-5% depletion at low k

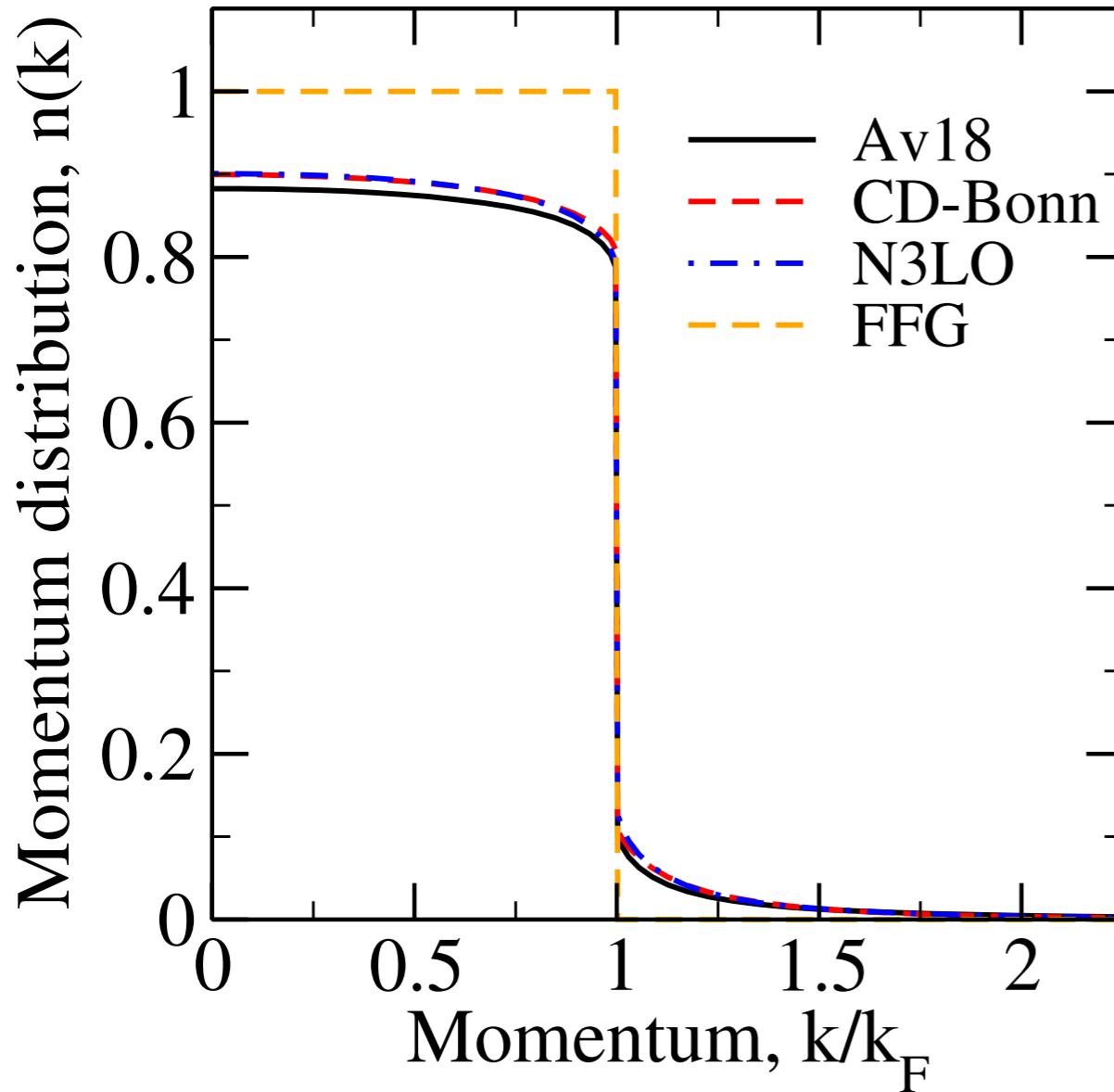


Momentum distribution

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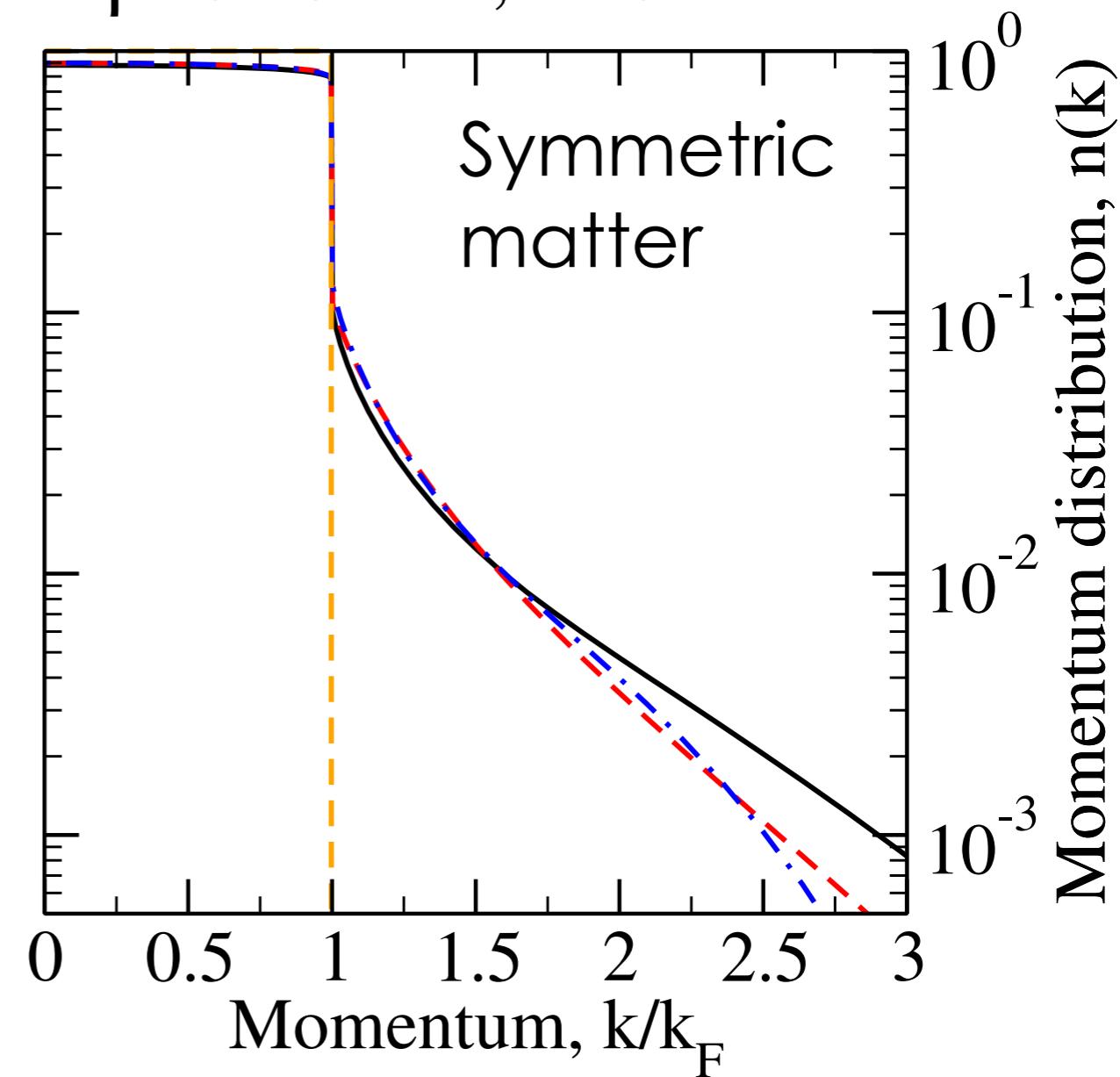
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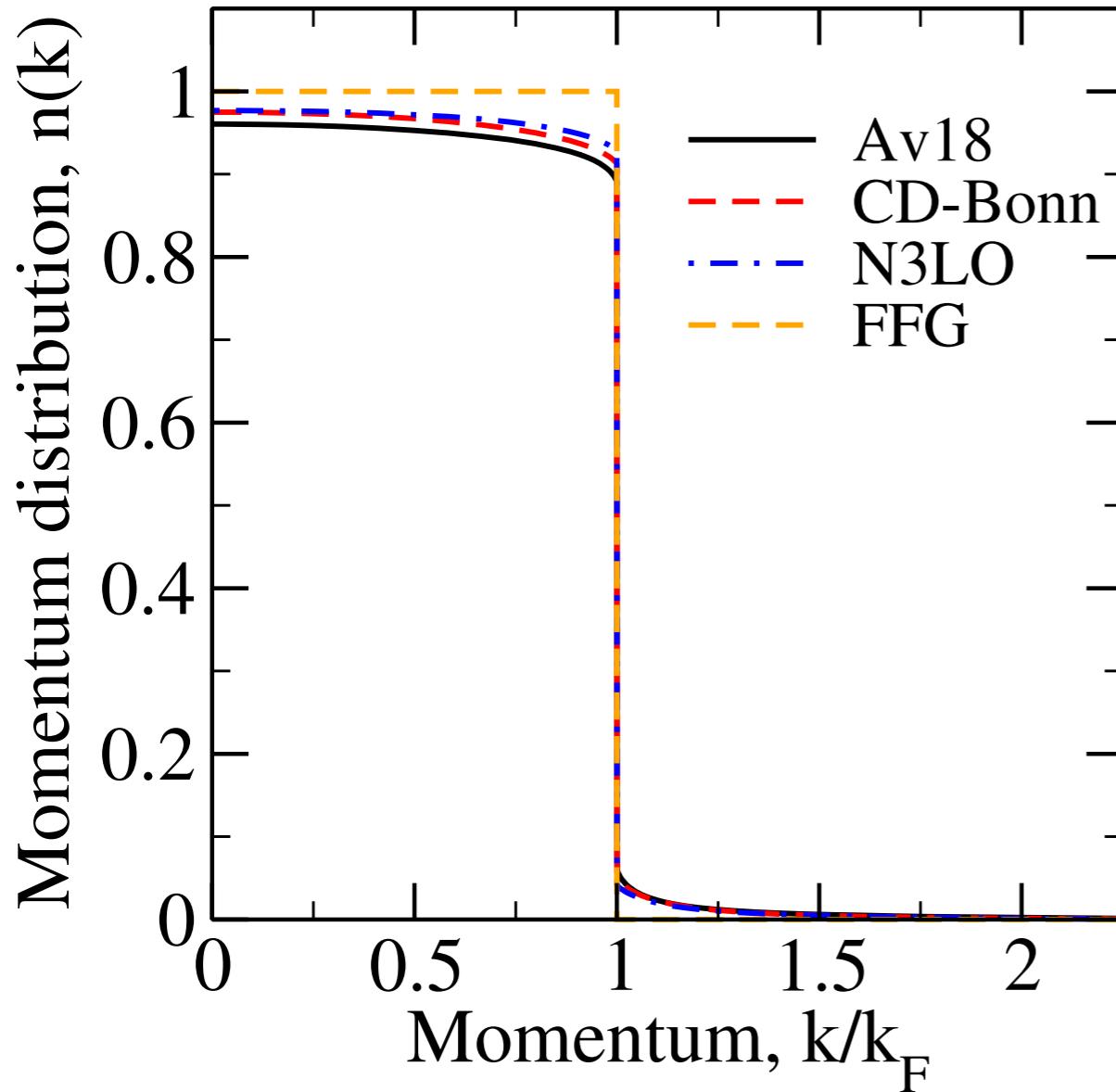
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Momentum distribution

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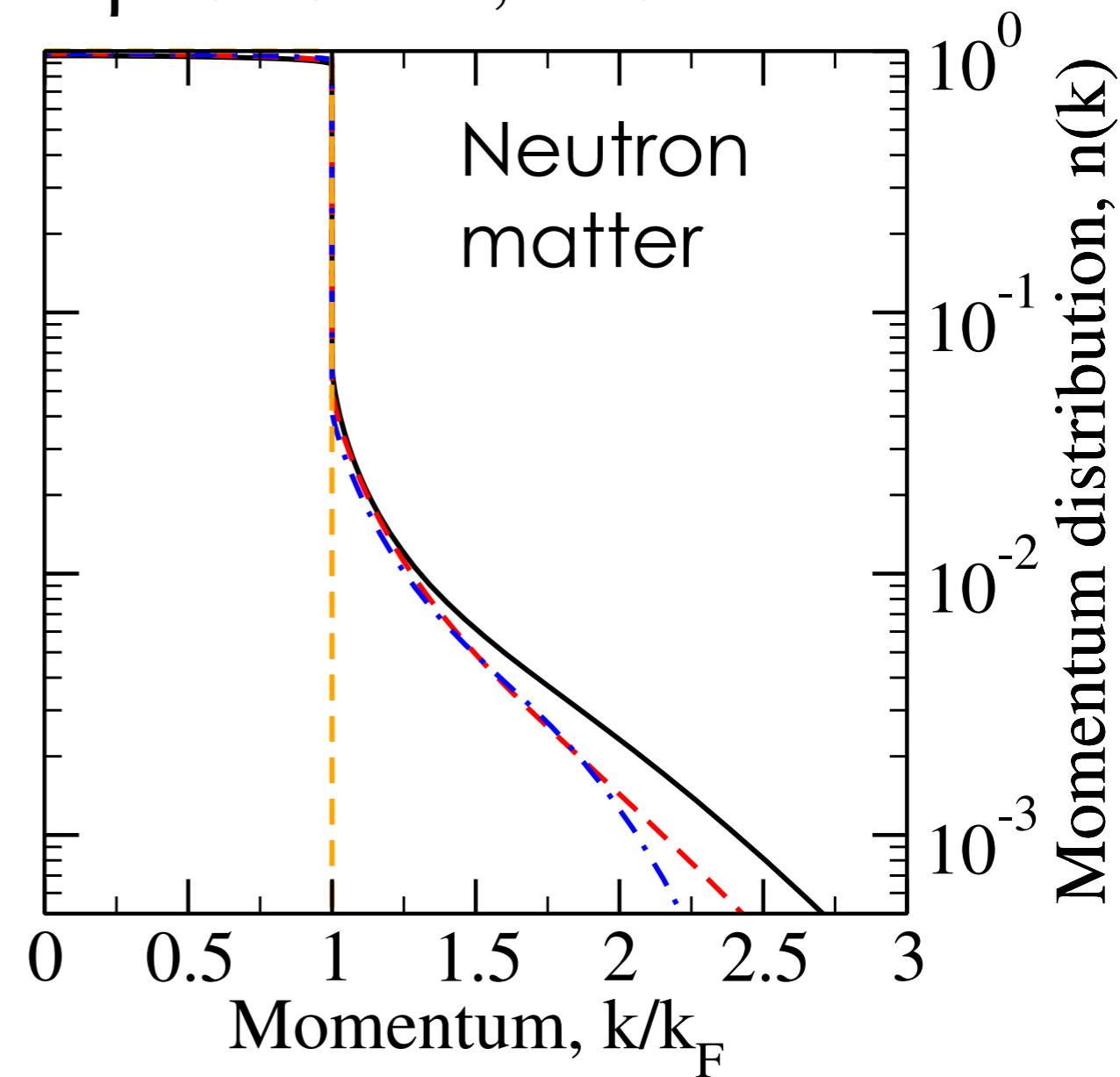
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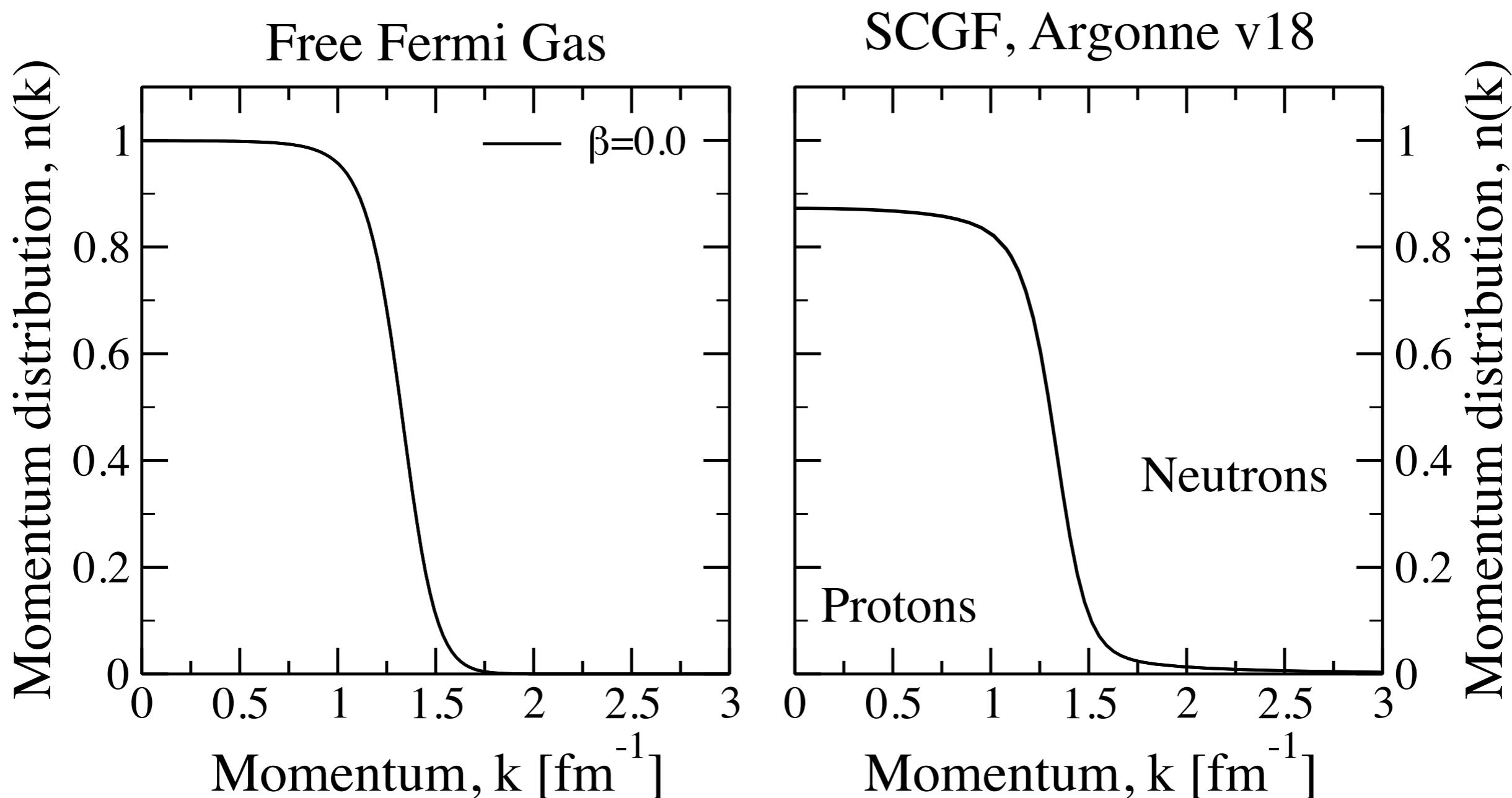
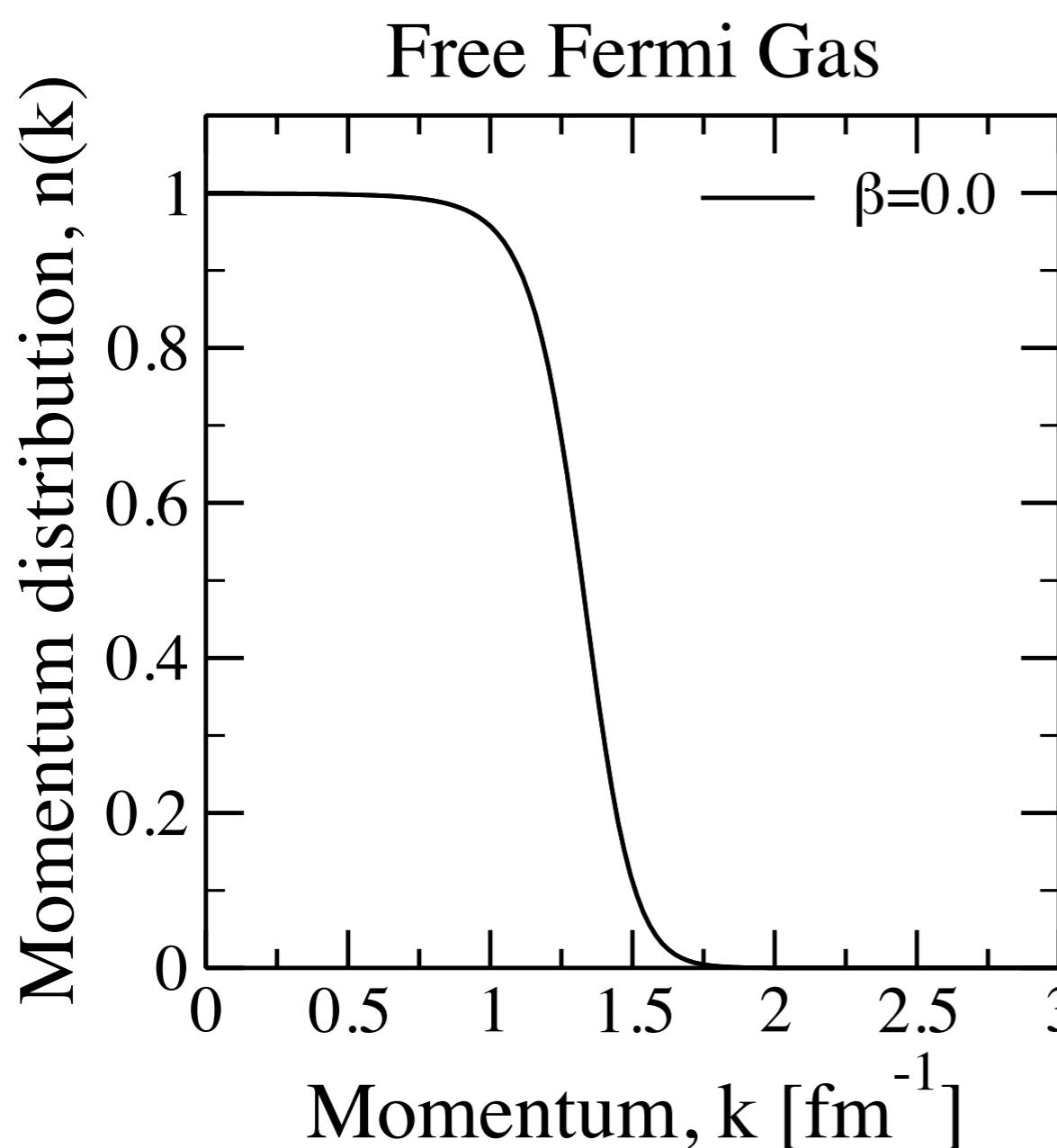


Momentum distribution

Asymmetric matter

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$T=5 \text{ MeV}$



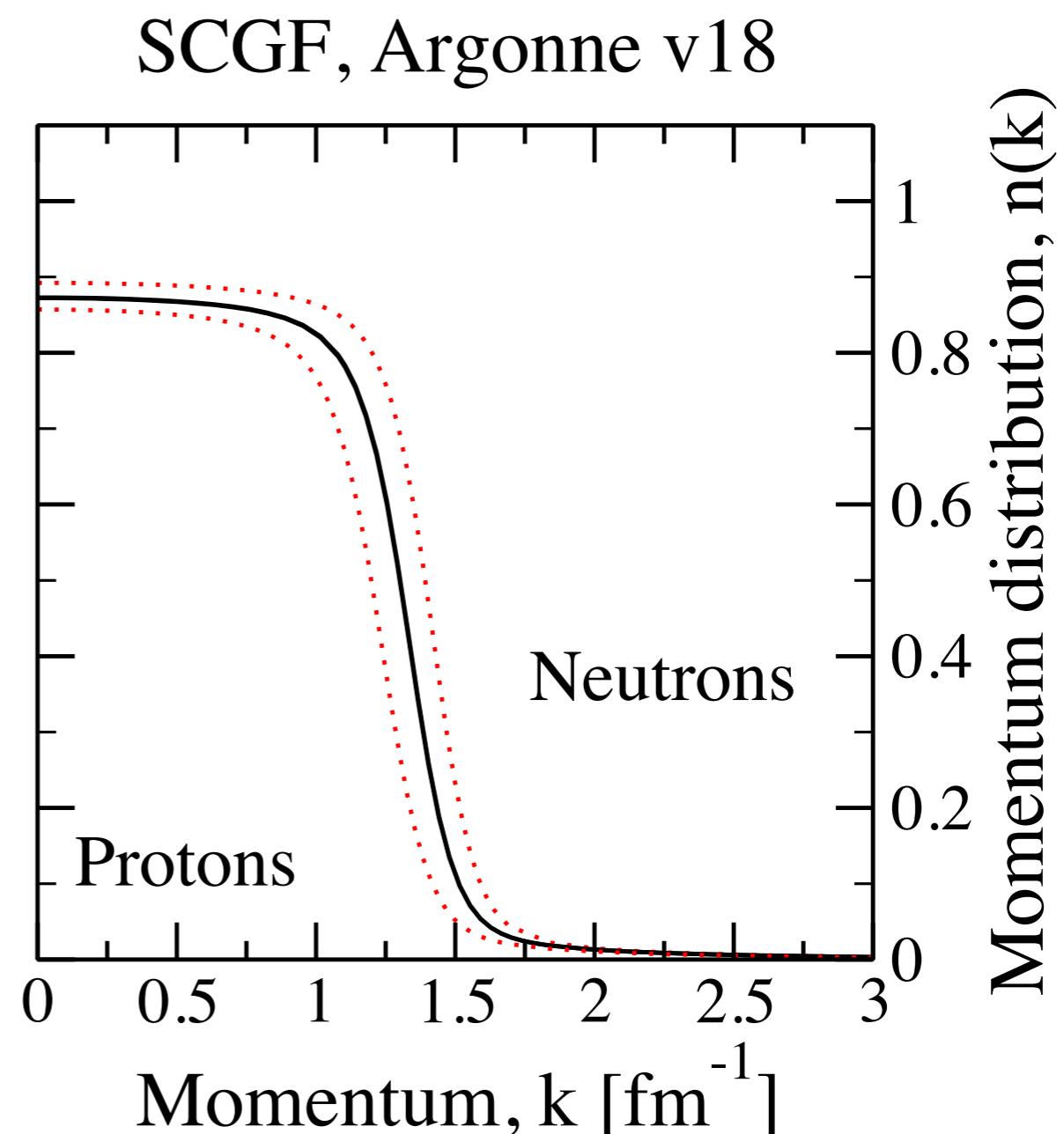
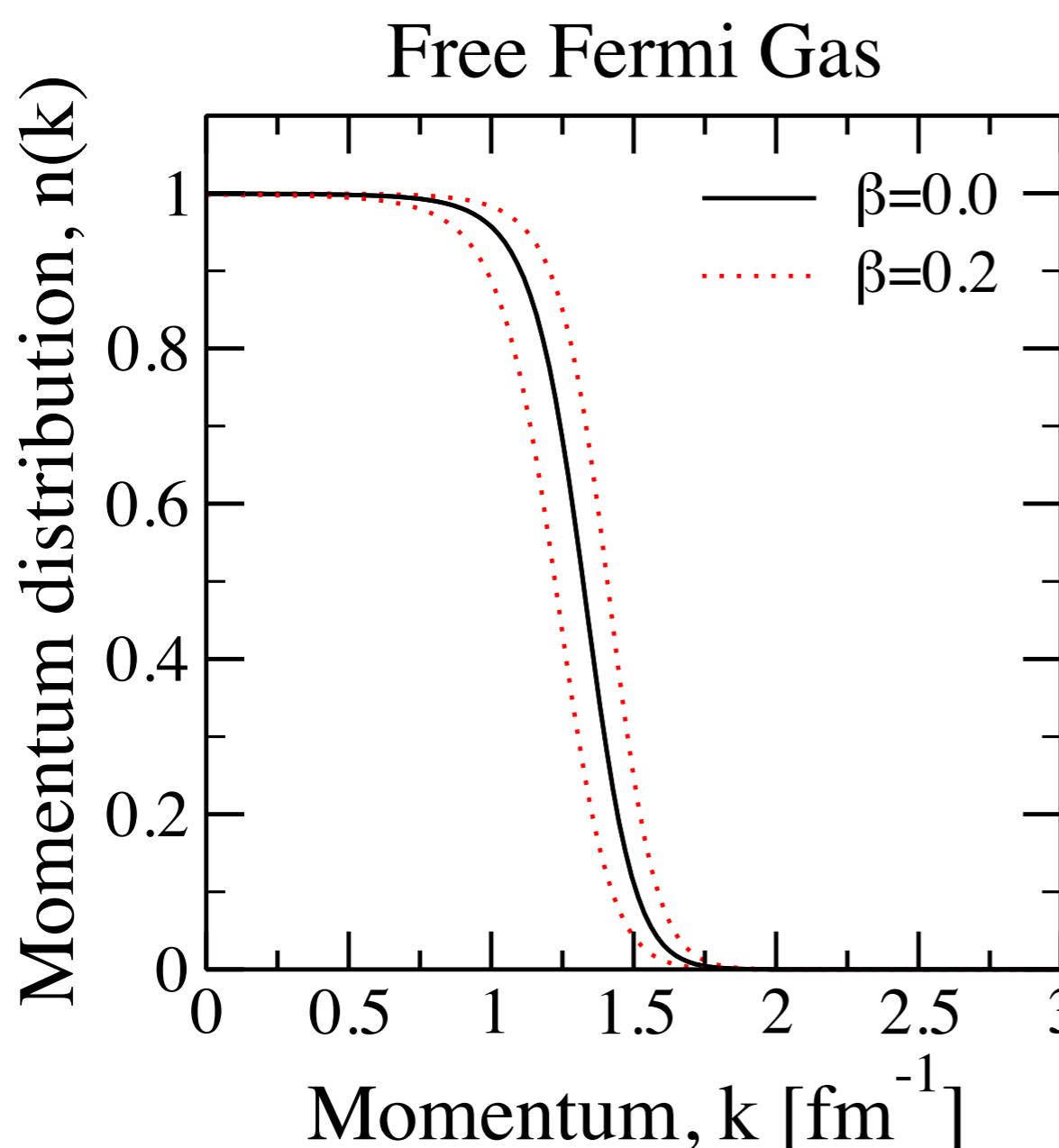
- Correlations affect depletion \Rightarrow non-perturbative effect

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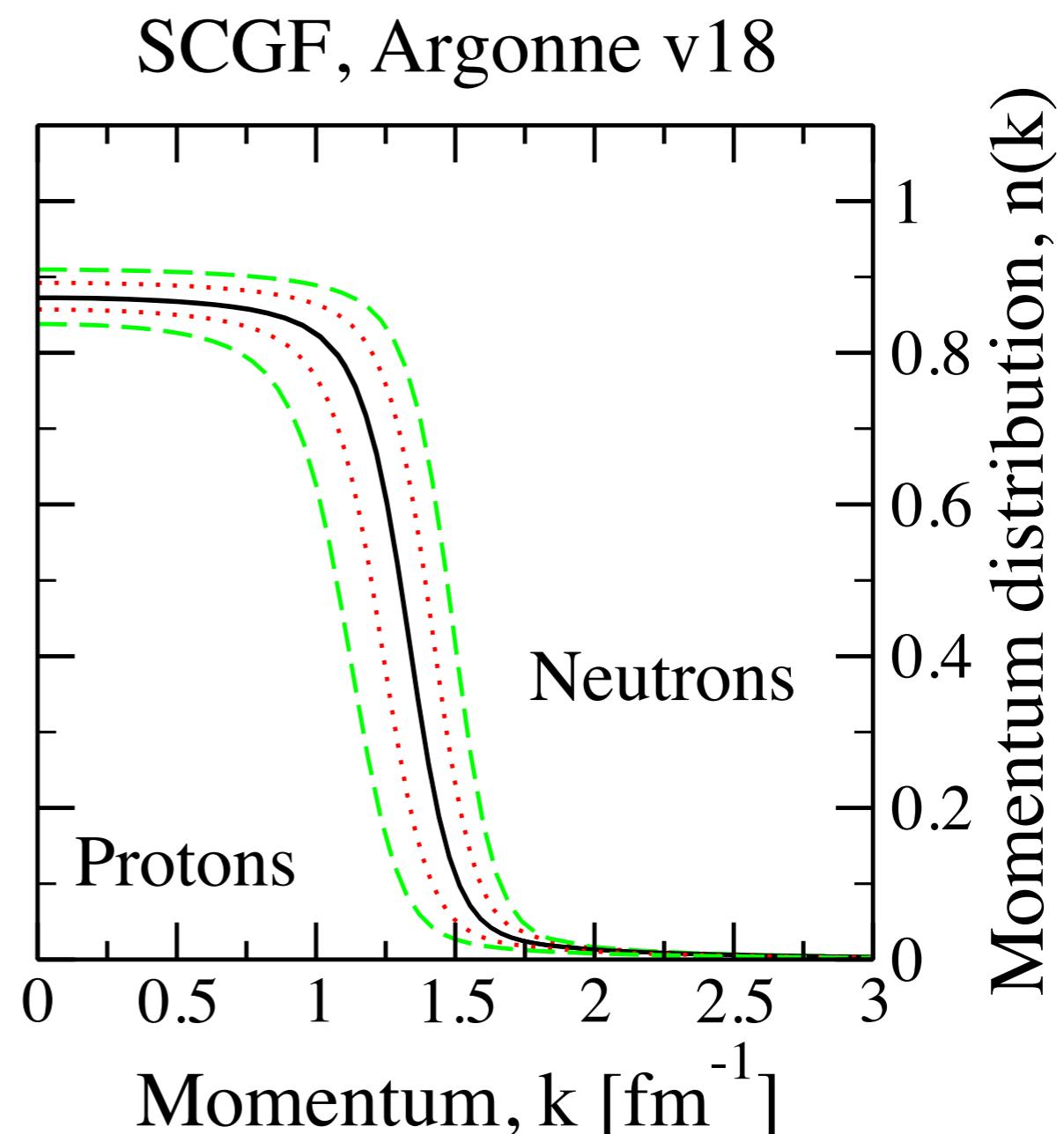
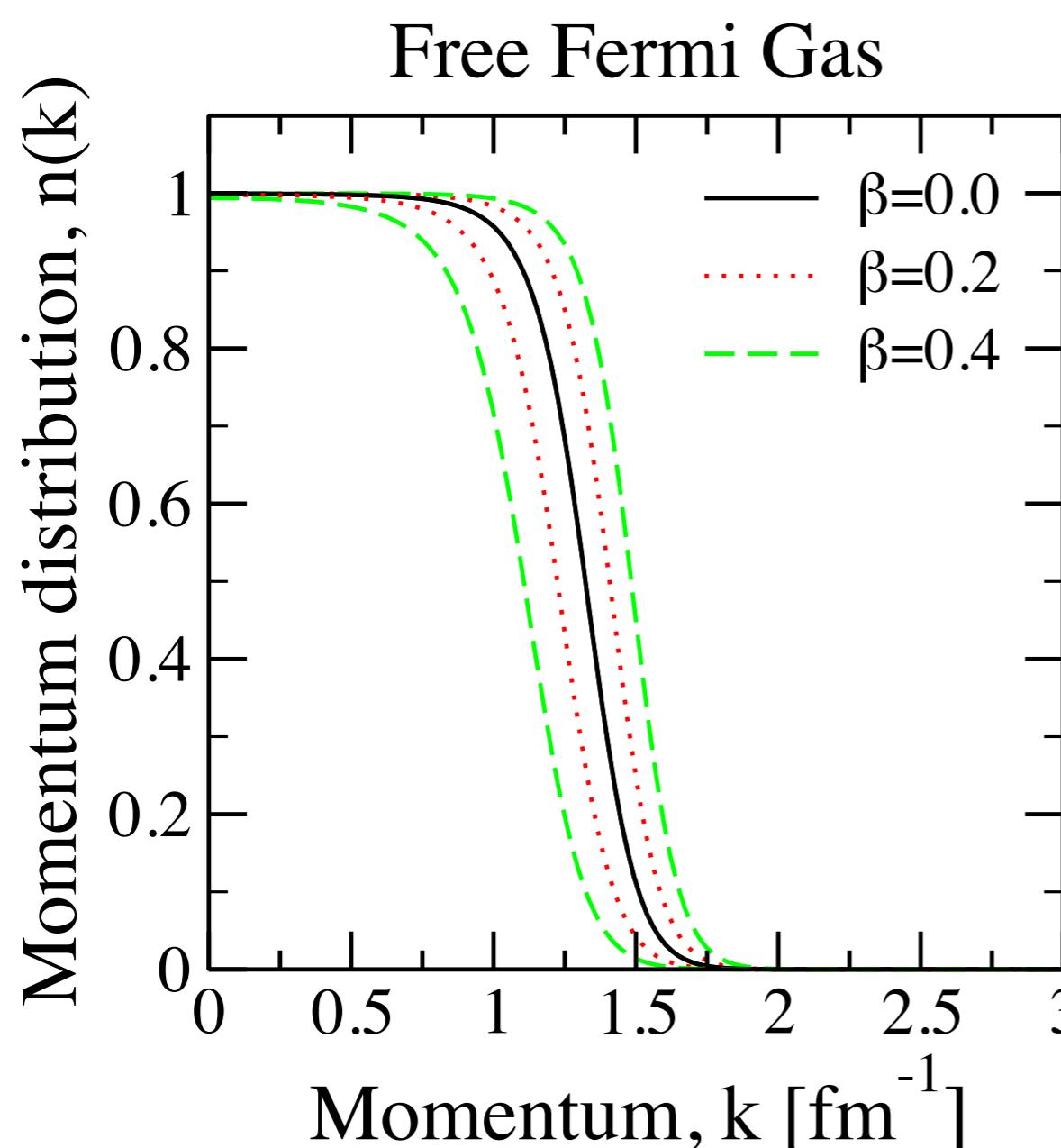
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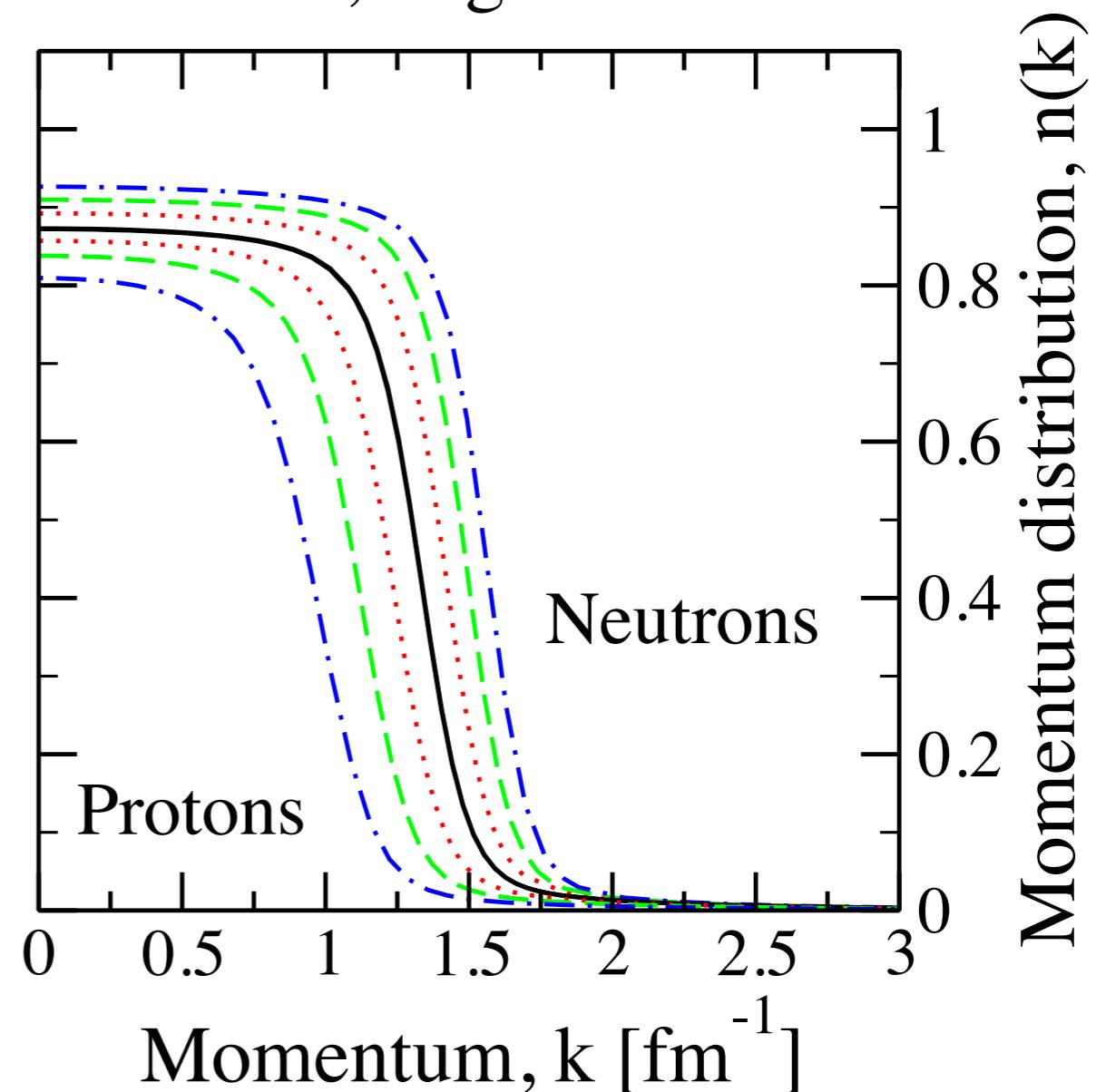
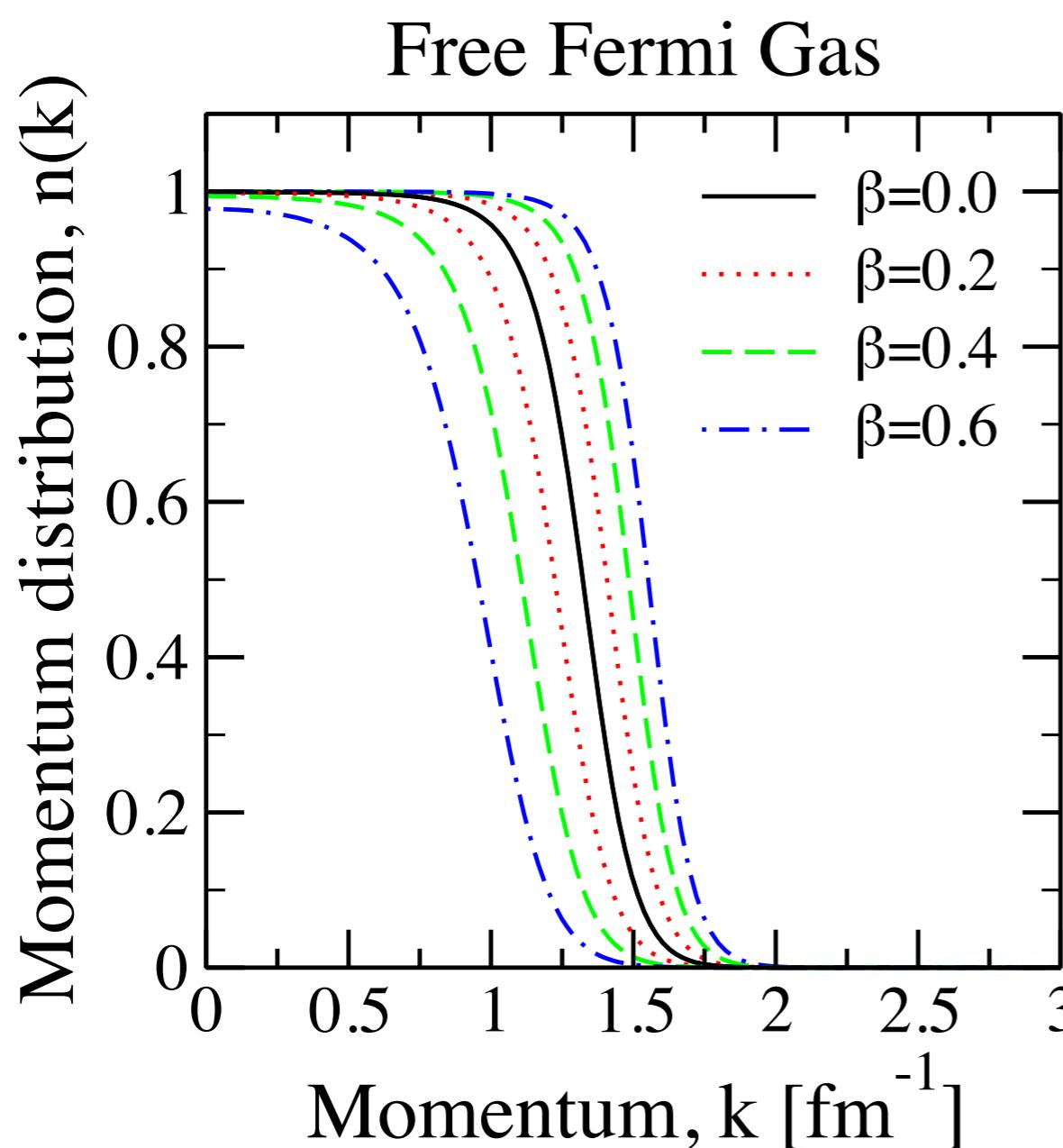
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SCGF, Argonne v18



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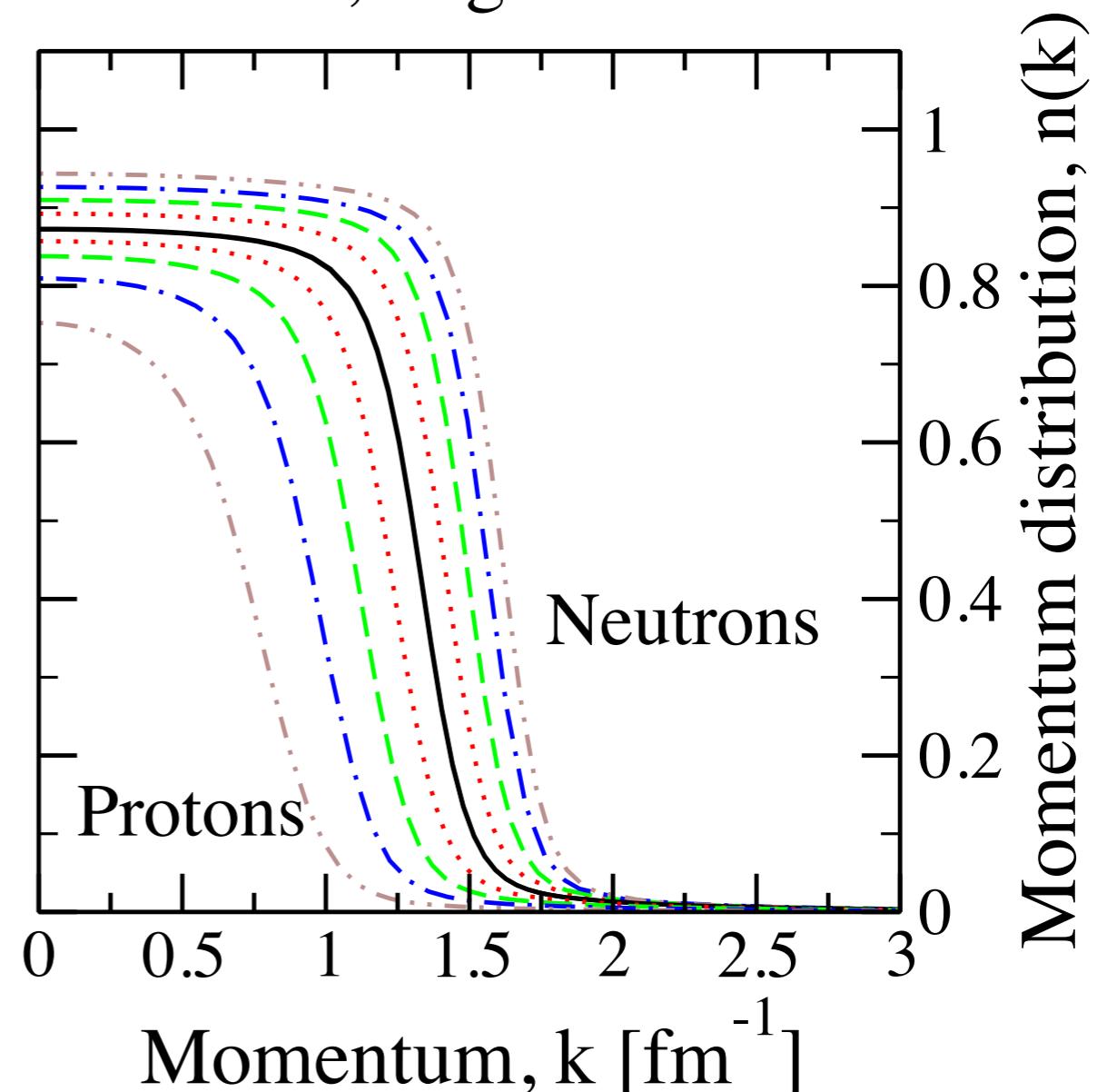
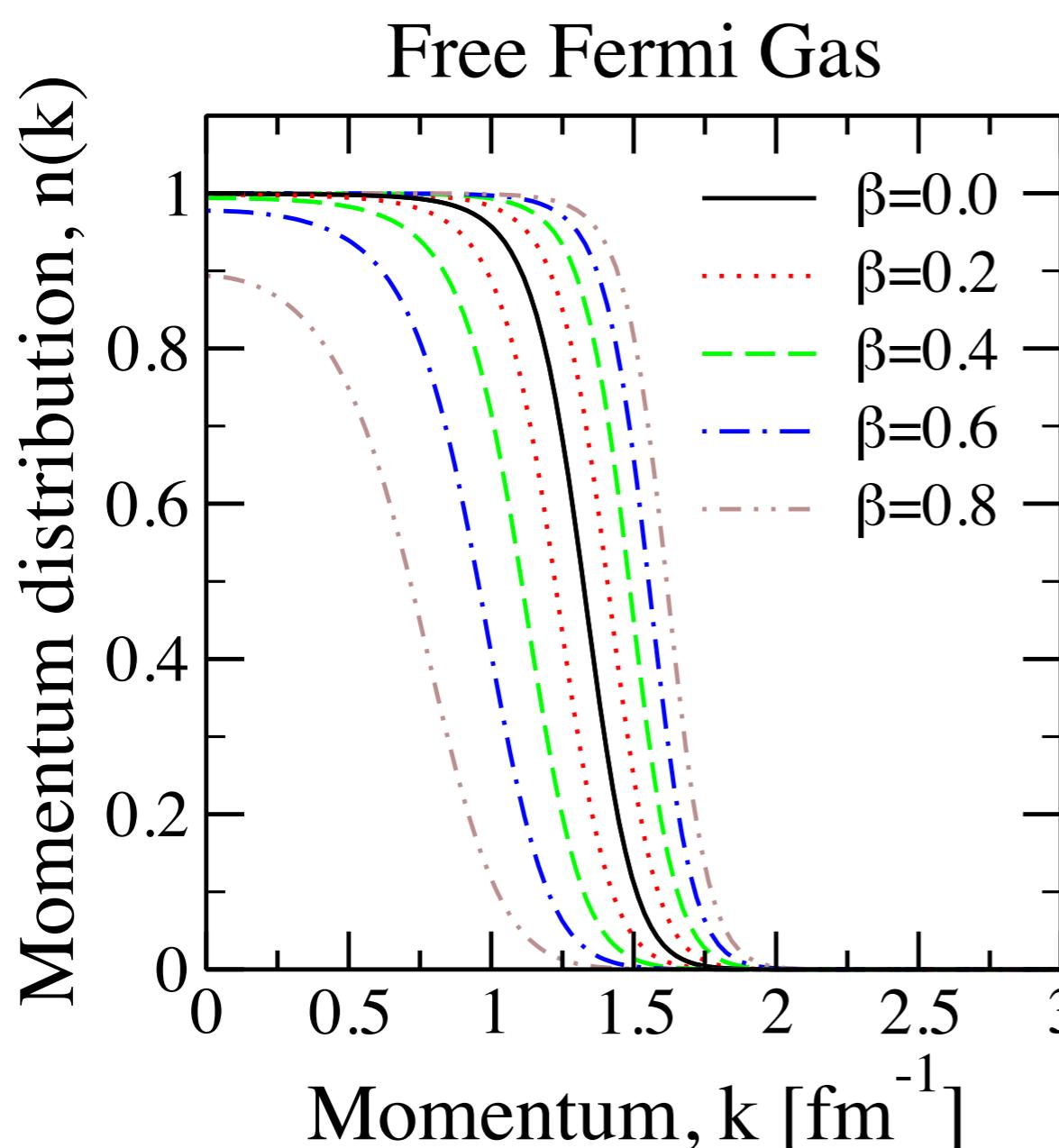
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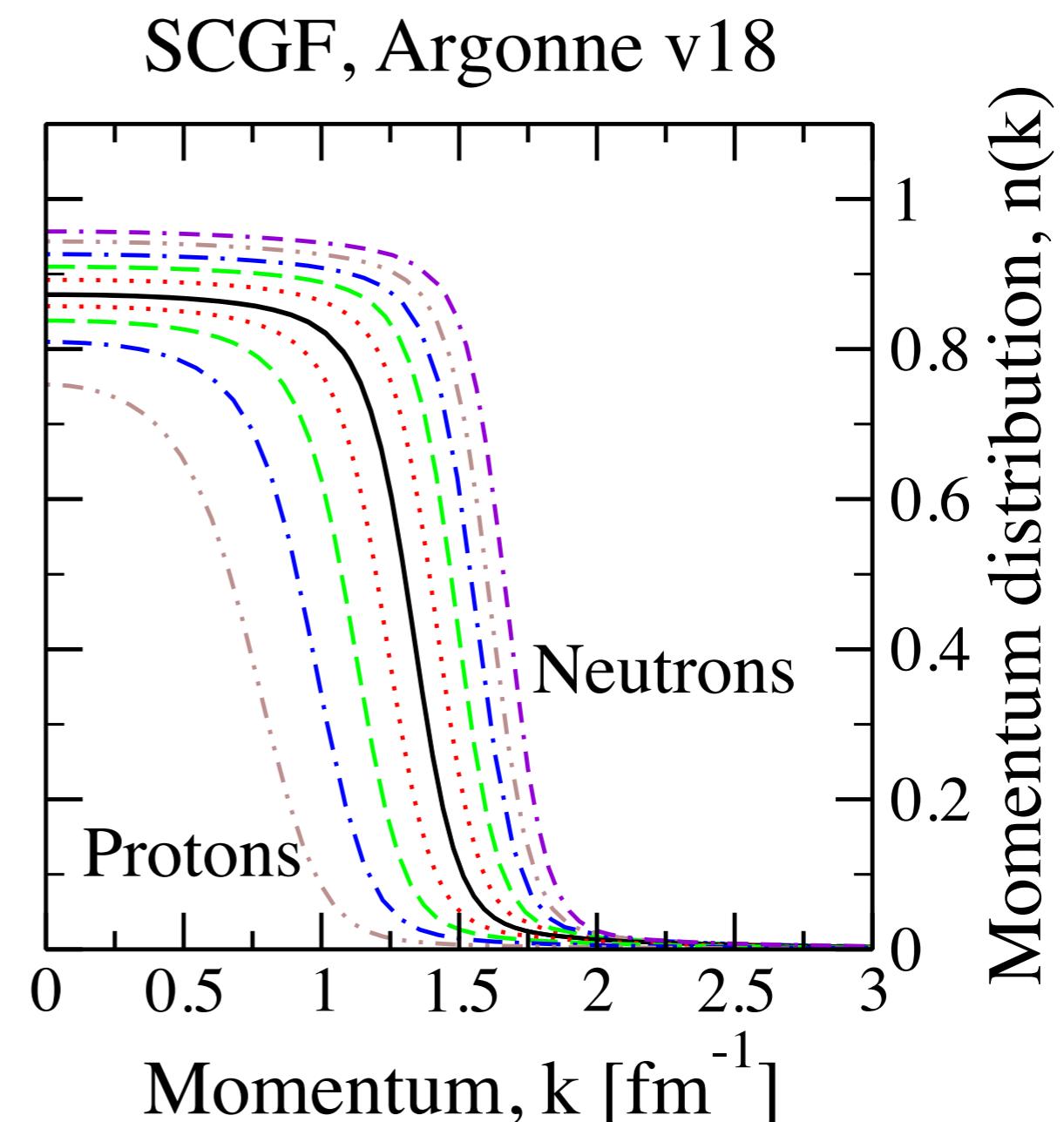
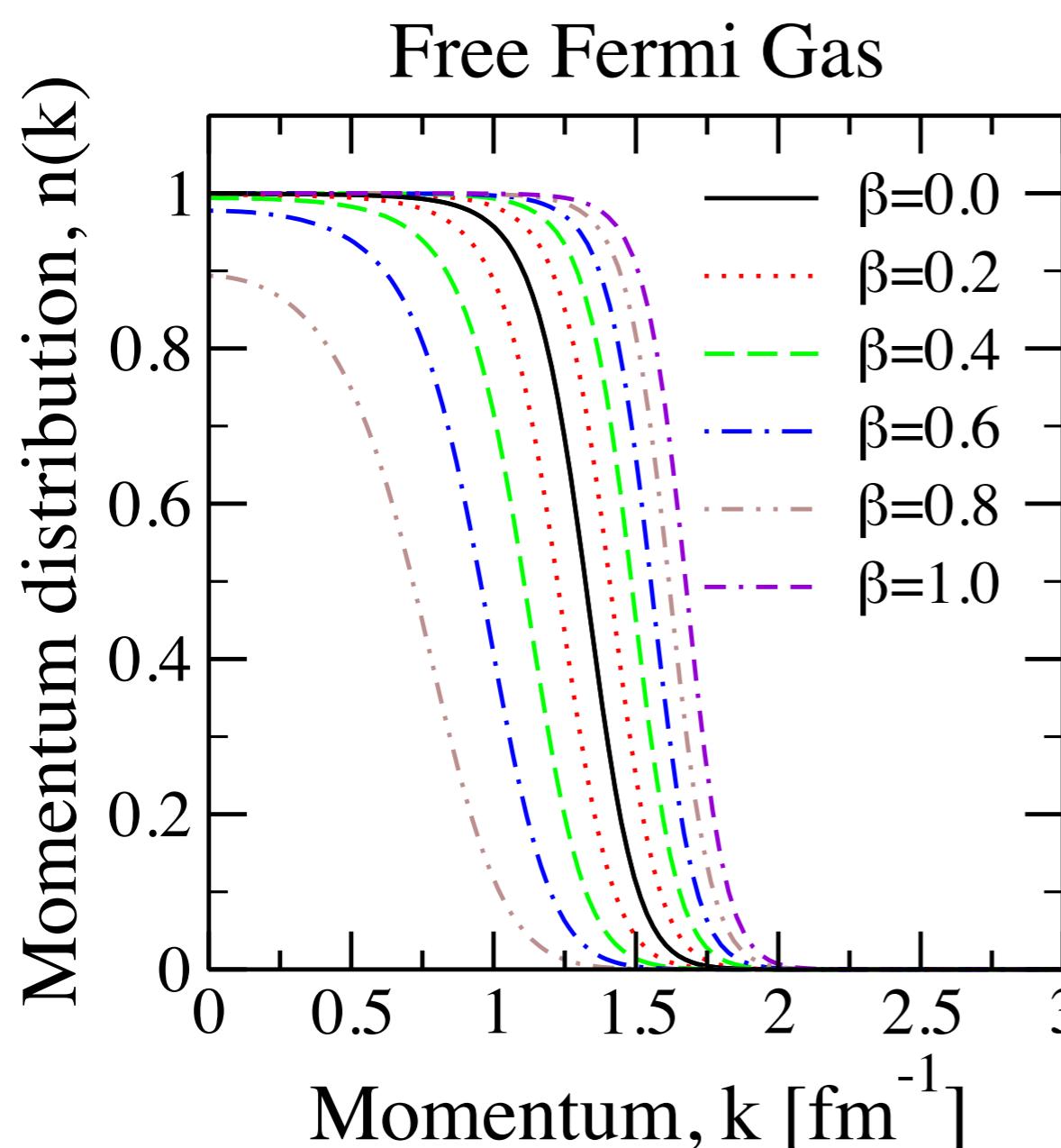
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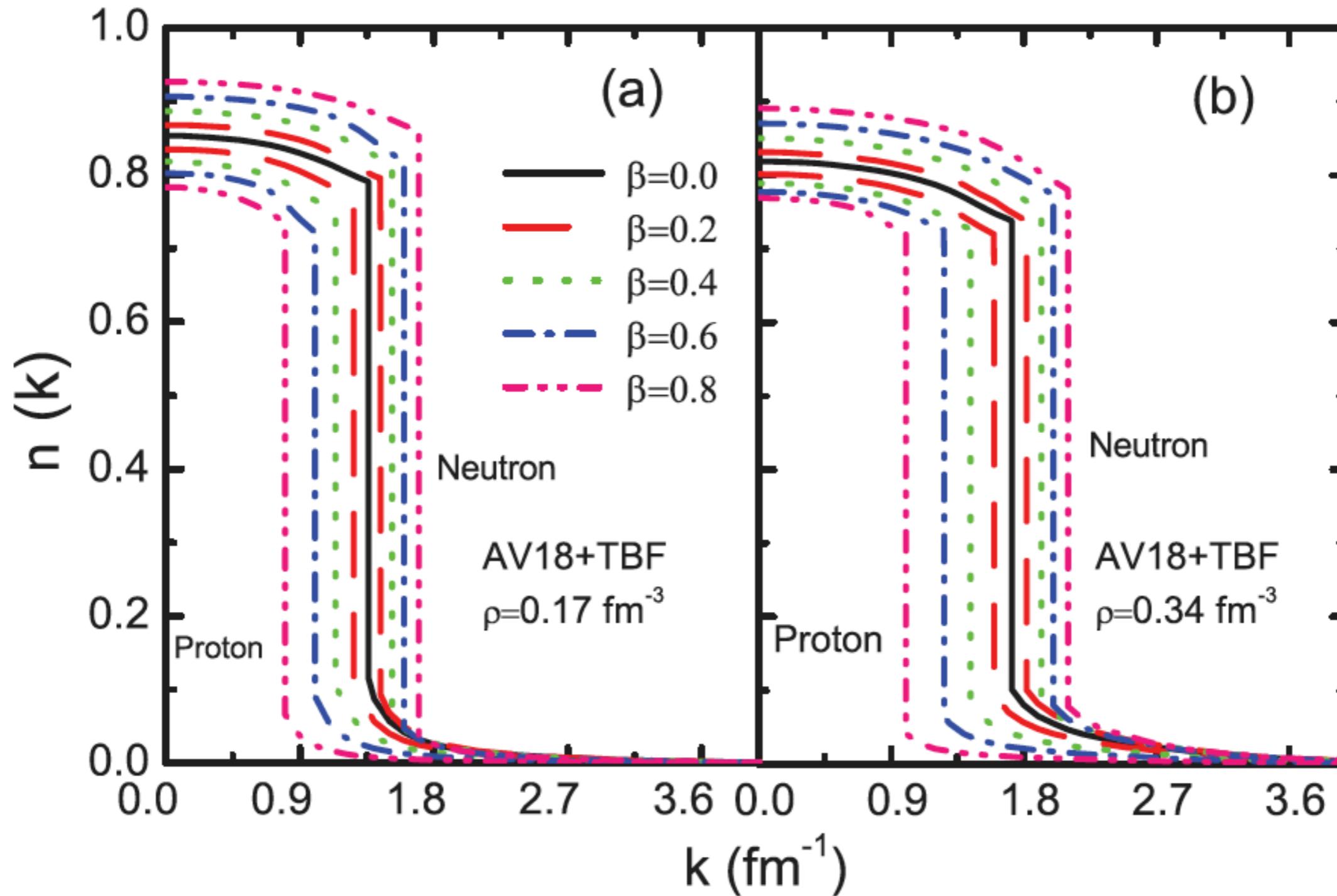


- Correlations affect depletion \Rightarrow non-perturbative effect
- Neutrons become less correlated
- Protons become more correlated

Theoretical confirmation

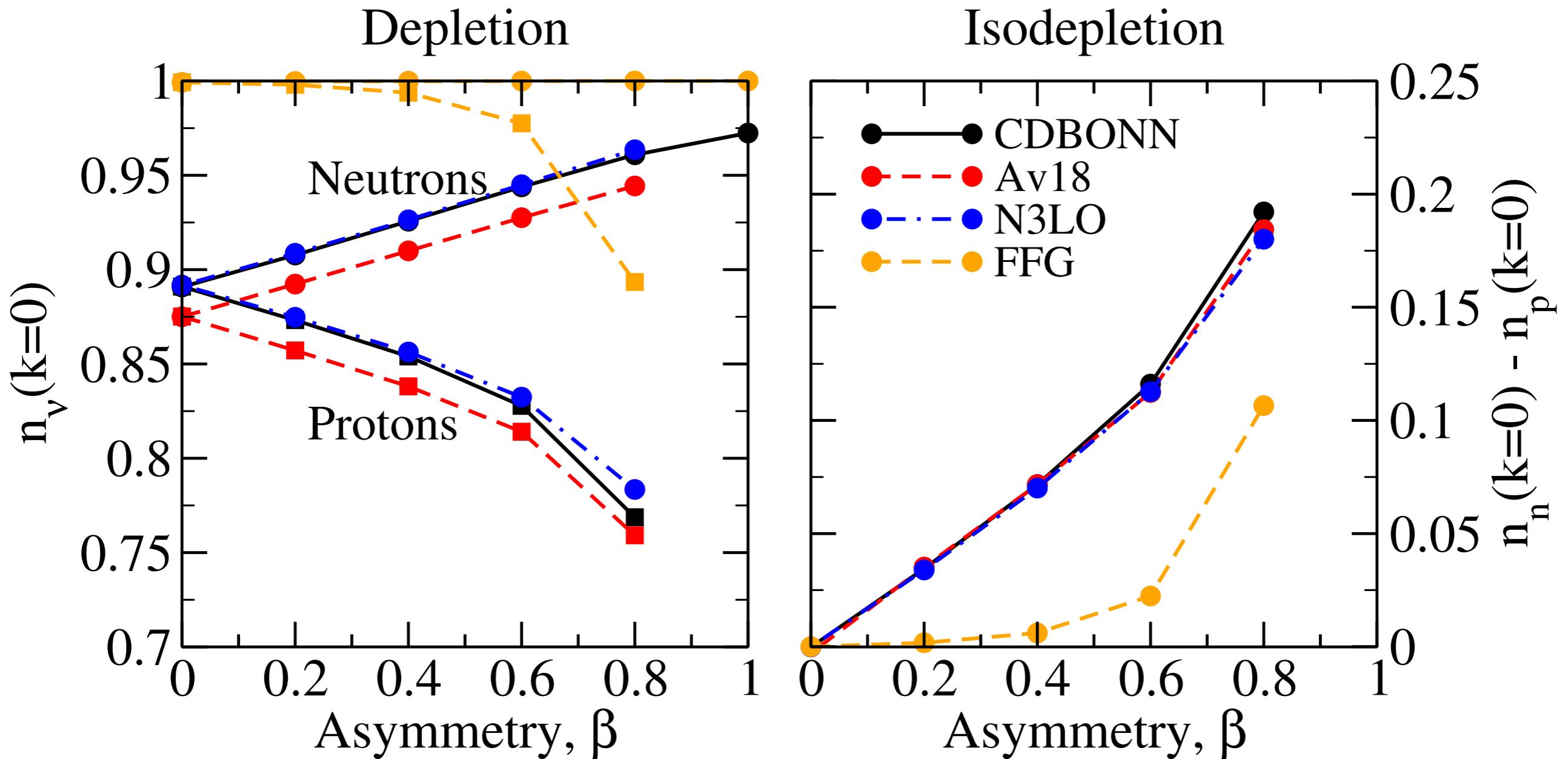
Asymmetric matter momentum distribution

EBHF calculation, $T = 0$



Behaviour at low momentum

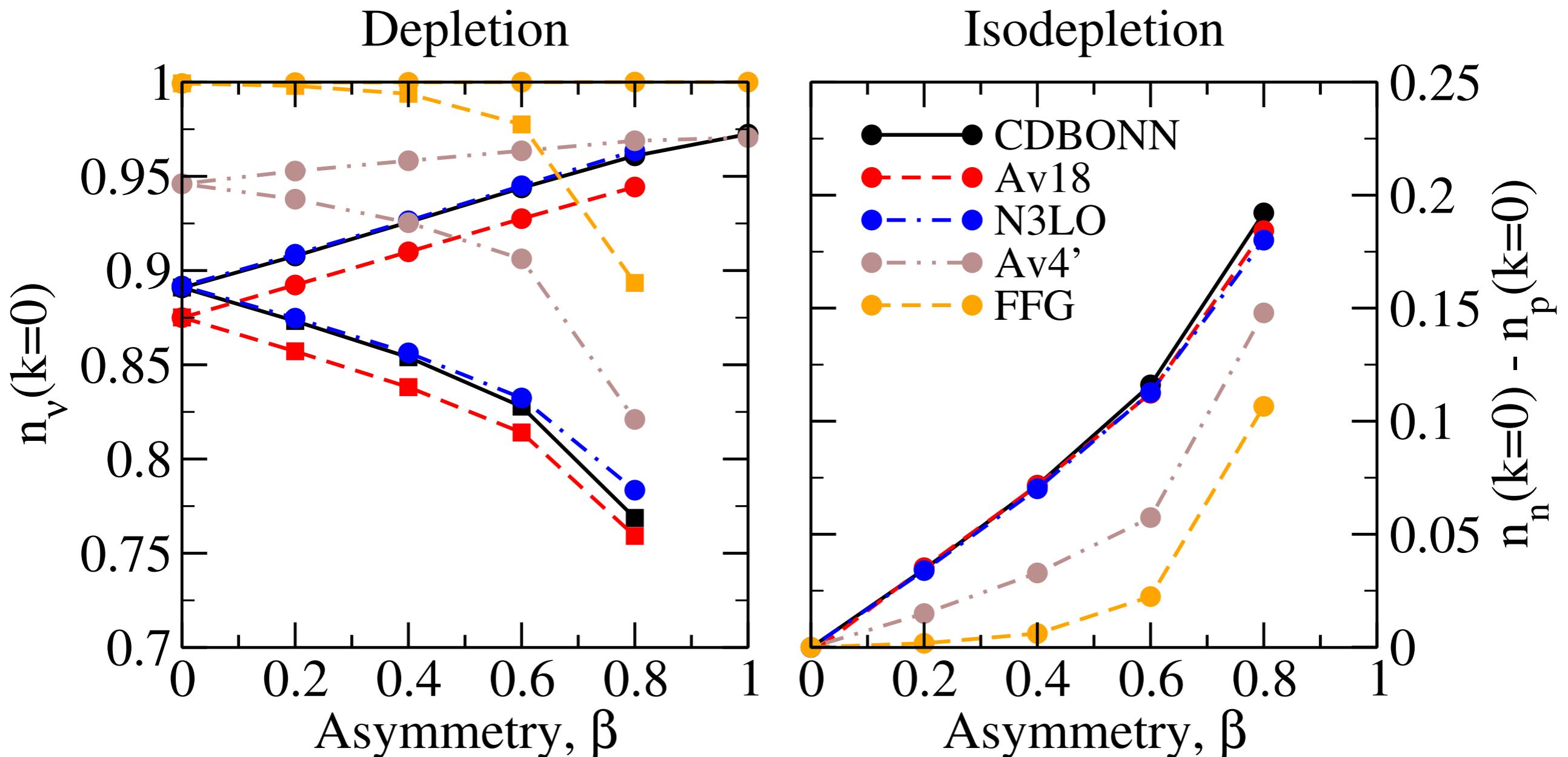
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- Correlations vs asymmetry

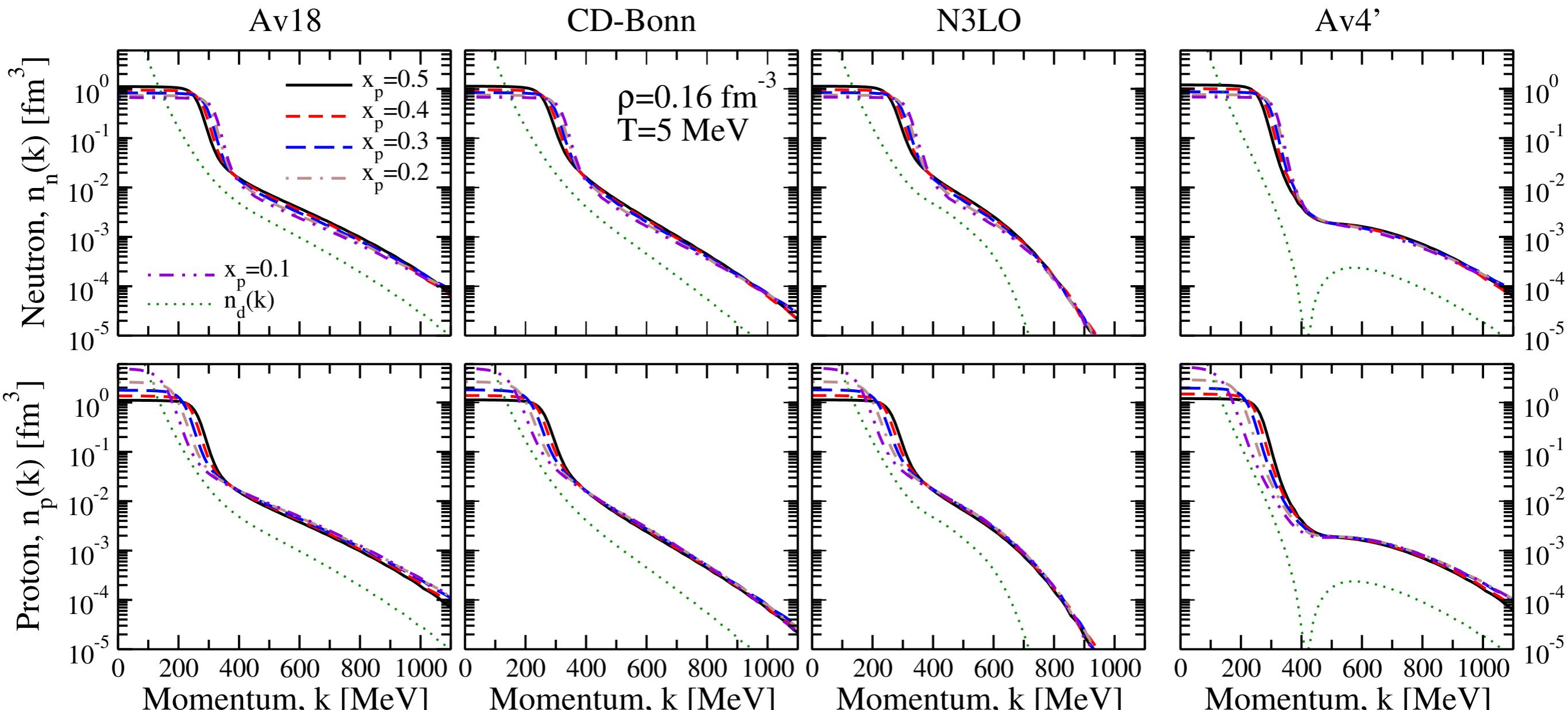
Behaviour at low momentum

$\rho=0.16 \text{ fm}^{-3}$ $T=5 \text{ MeV}$



- Correlations vs asymmetry
- All potentials lie in a narrow iso-depletion band
- Tensor needed!

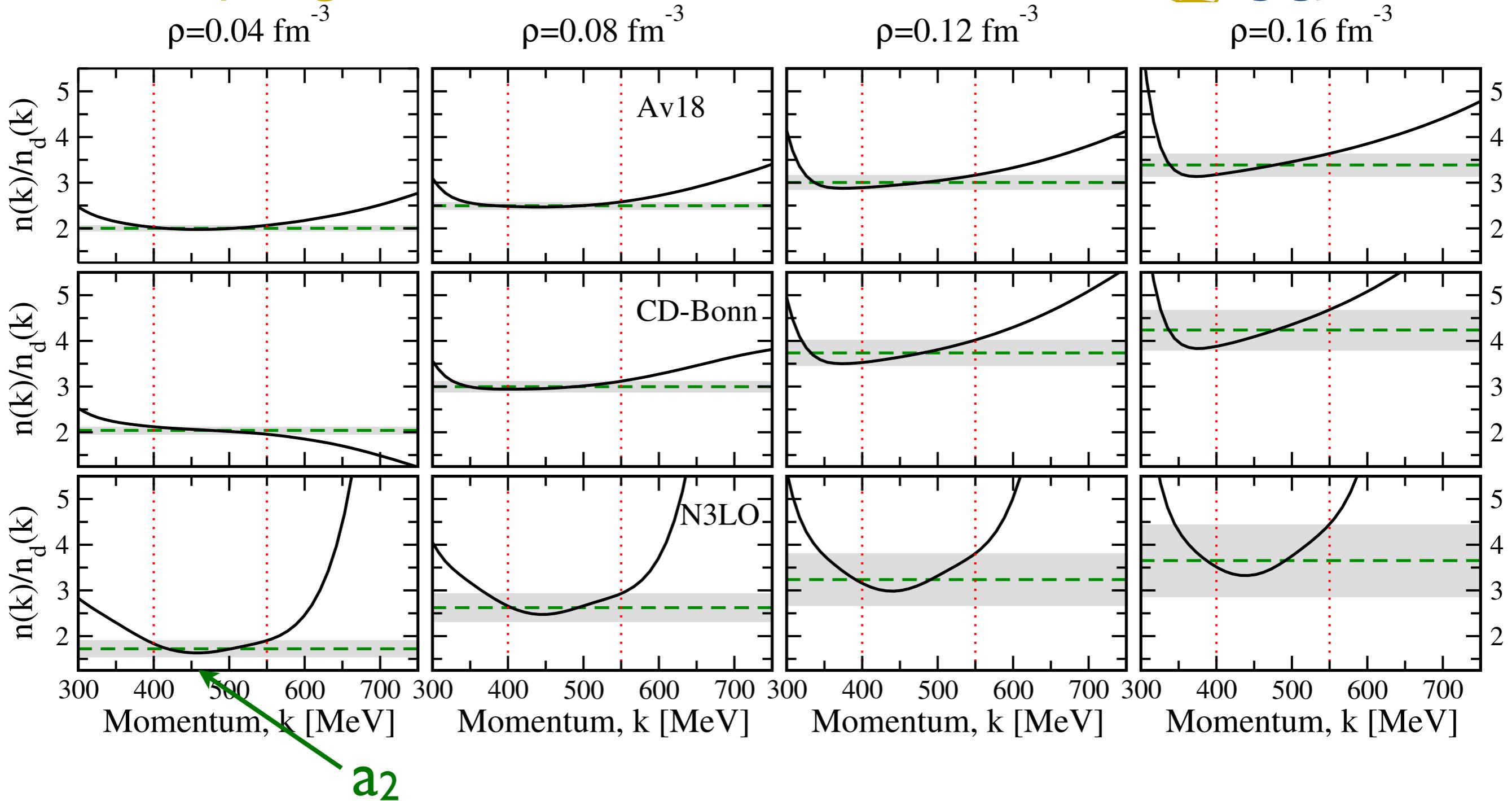
Quantifying tensor-like correlations



$$4\pi \int_0^\infty dk k^2 n(k) = 1$$

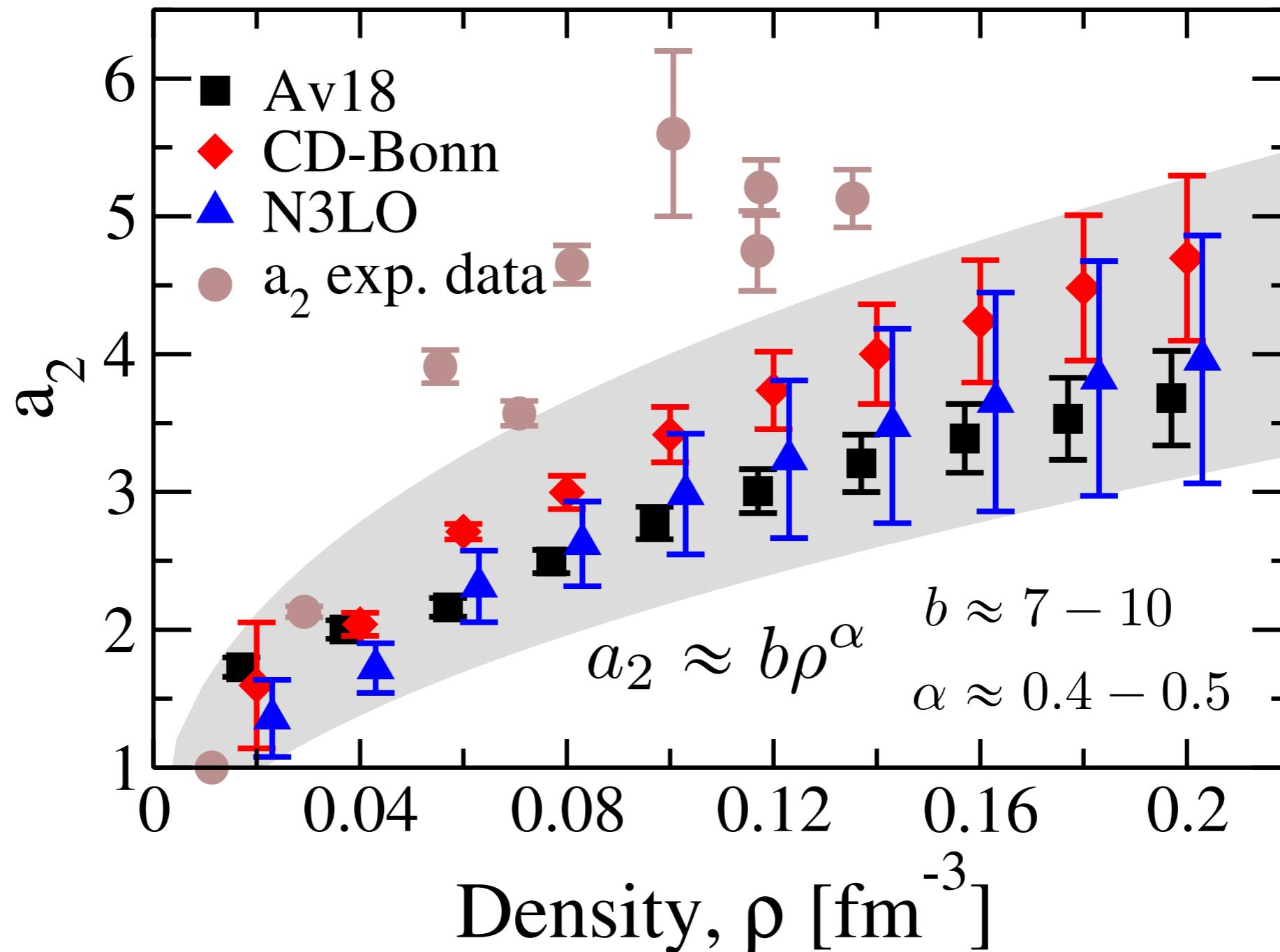
- SNM vs deuteron momentum distribution
- Very similar in tensor-like area
- Compare to empirical estimates Arrington et al. PPNP 67 898 (2012)

Quantifying tensor-like correlations



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Quantifying tensor-like correlations

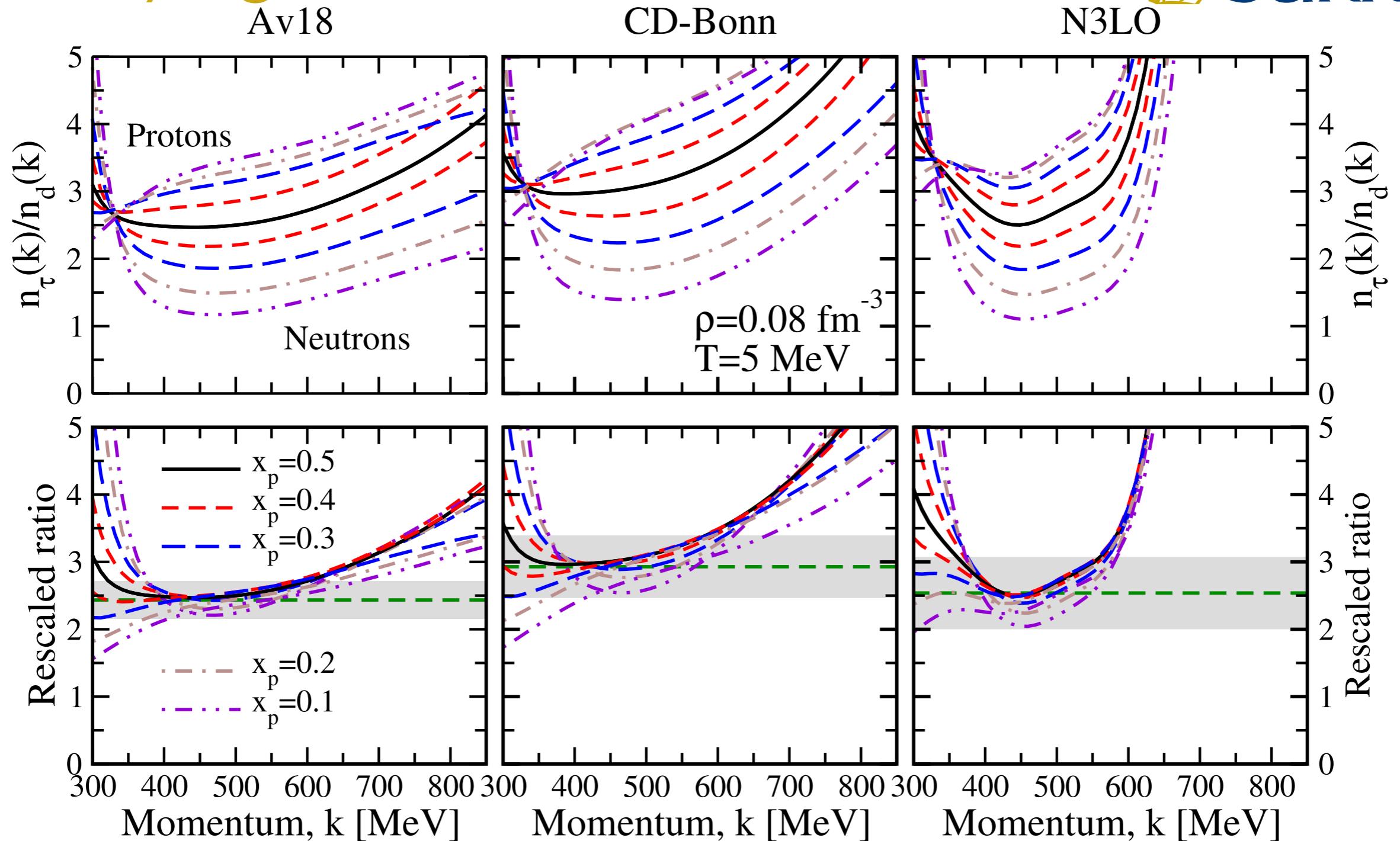


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Quantifying tensor-like correlations II



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SURREY



$$a_2^\tau(\rho, \beta) = a_2^\tau(\rho, 0) [1 \mp \gamma(\rho)\beta]$$

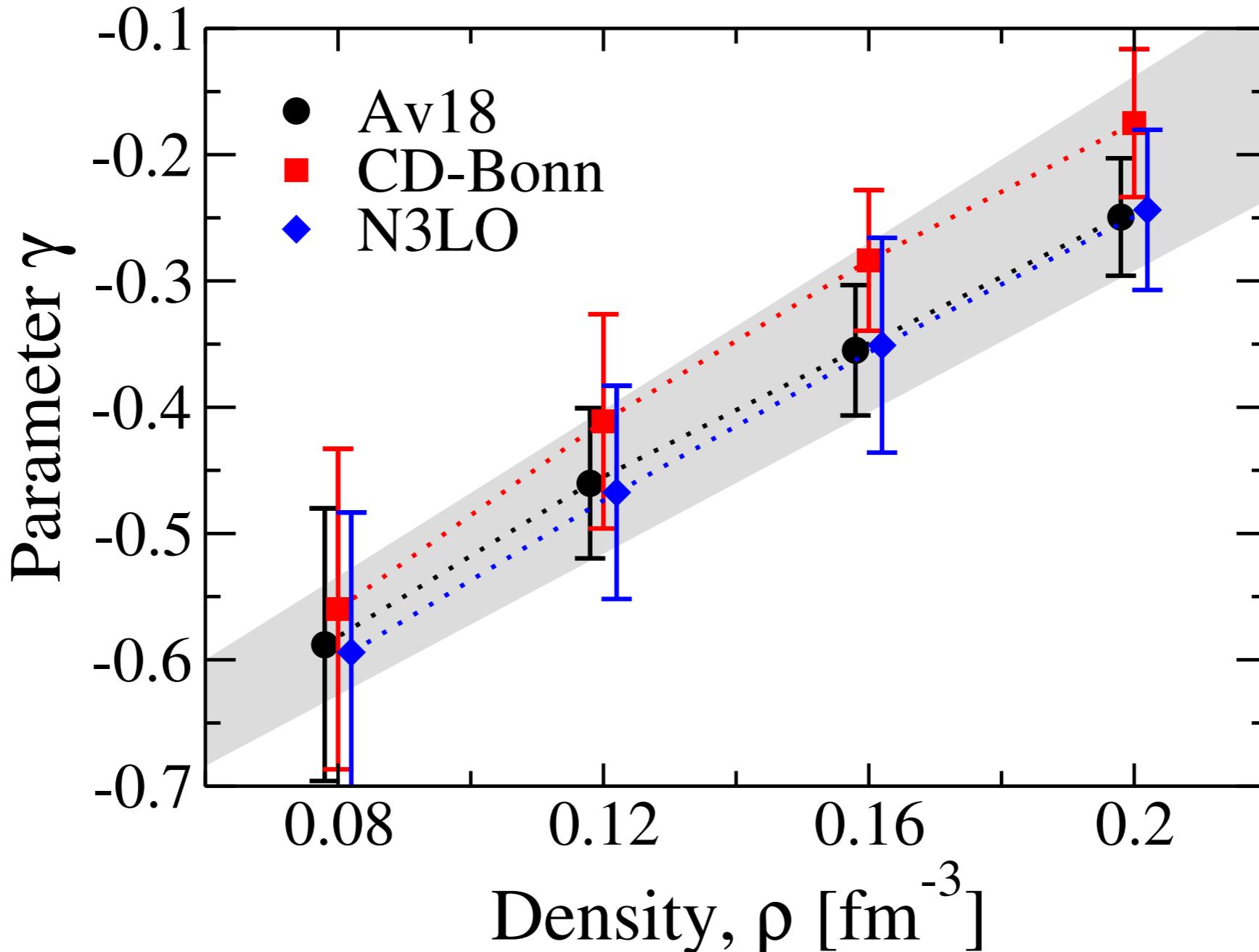
Asymmetry dependence factorizes

Predicted by DFG (Lee-Yang) model $\gamma = \frac{a_{nn}^2 - 2a_{np}^2}{a_{nn}^2 + 2a_{np}^2}$

Quantifying tensor-like correlations II



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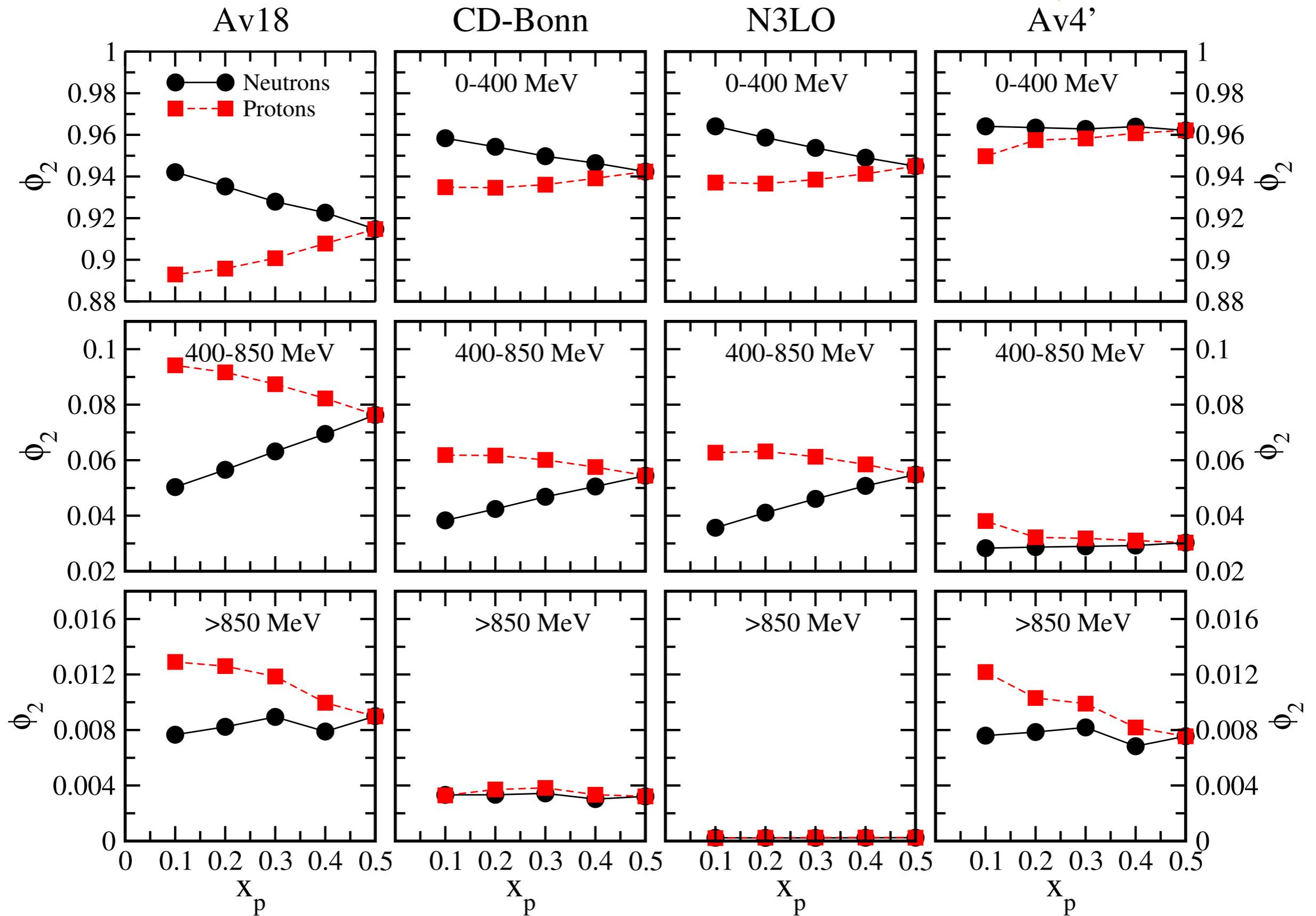


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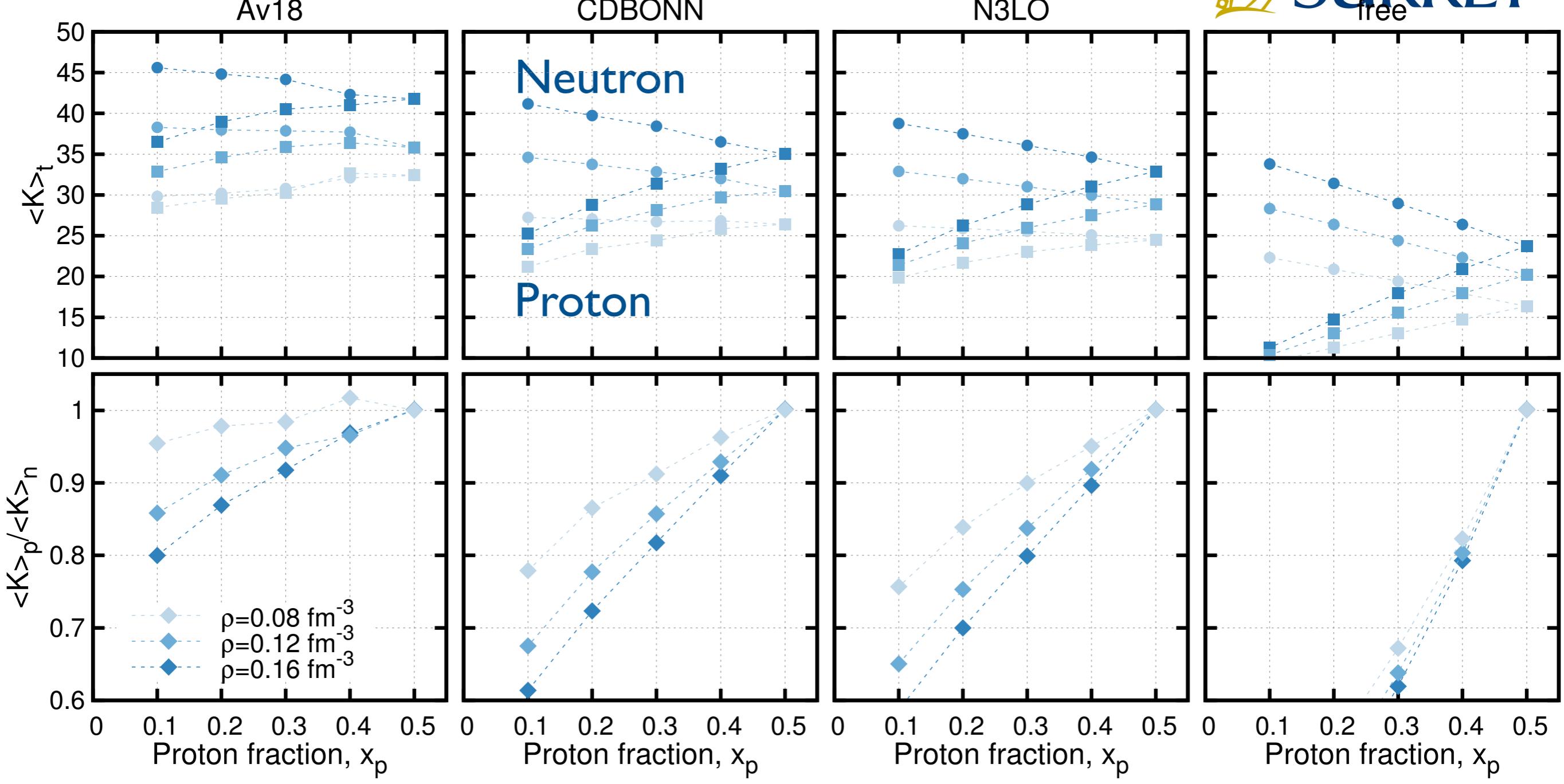
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Minority protons KE?



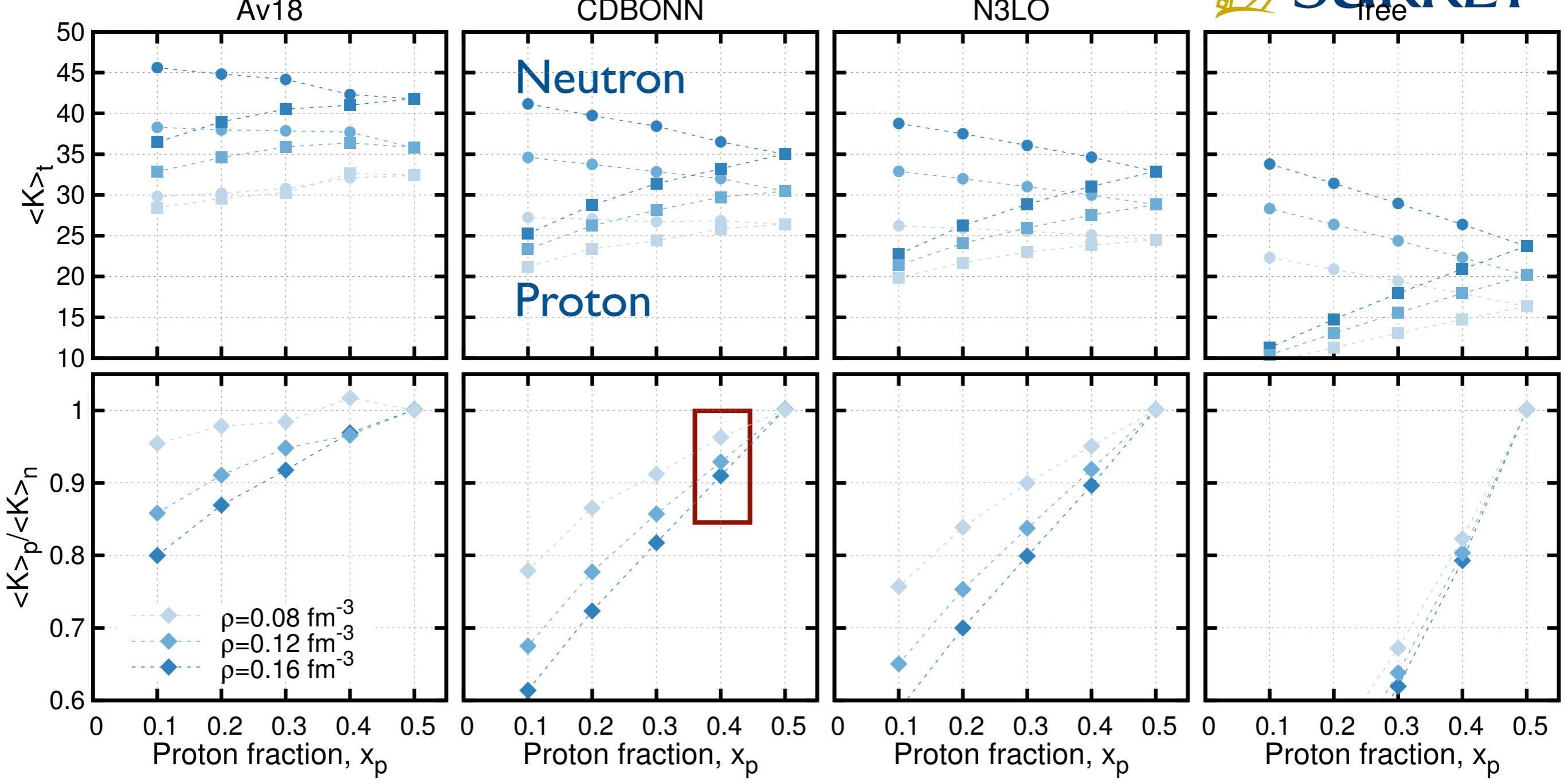
But are protons more energetic?



	x_p	0.5	0.4	0.3	0.2	0.1
Av18	K_n	41.8	42.3	44.2	44.8	45.6
	K_p	41.8	41.0	40.5	39.0	36.5
CD-Bonn	K_n	35.0	36.5	38.4	39.8	41.2
	K_p	35.0	33.2	31.4	28.7	25.3
N3LO	K_n	32.9	34.6	36.1	37.5	39.0
	K_p	32.9	31.1	28.8	26.2	22.7

10% effect in ^{208}Pb

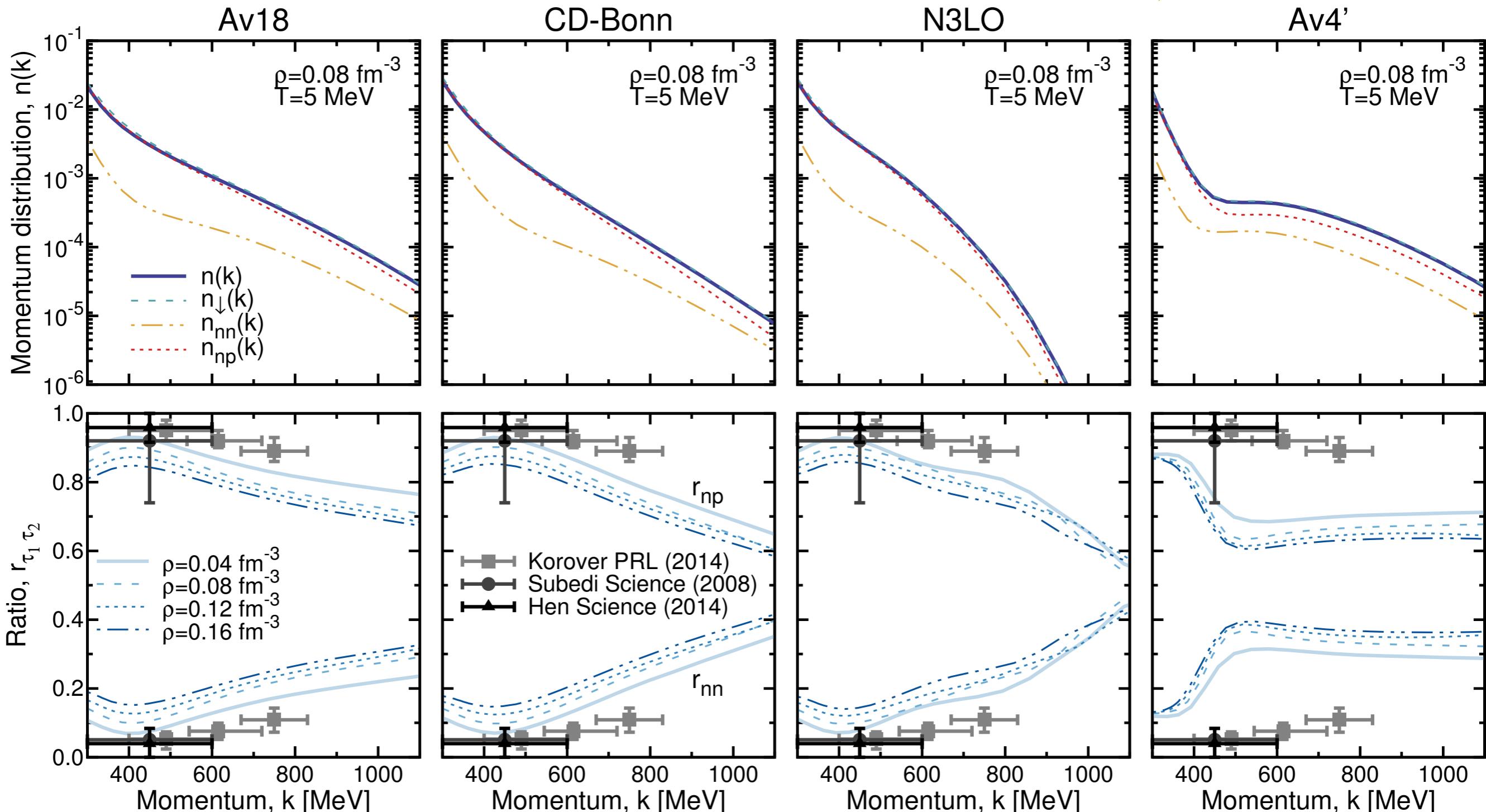
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np vs nn components



$$n_\downarrow(k) \approx \partial_\omega \text{Re}\Sigma_\uparrow(k, \omega)|_{\omega=\varepsilon(k)}$$

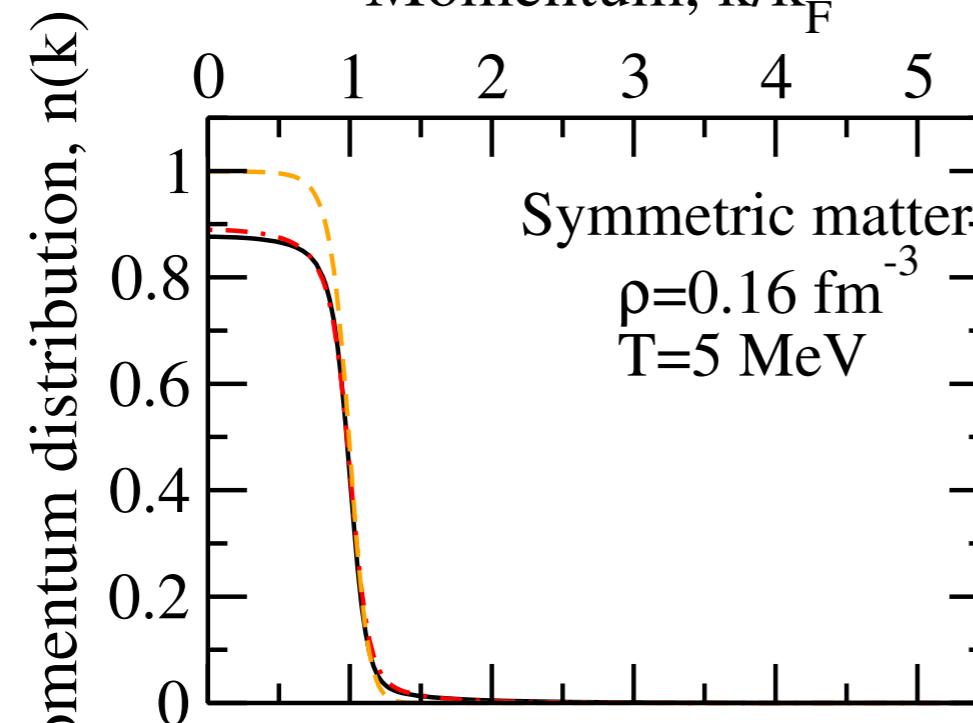
$$\text{Re}\Sigma_\uparrow(k, \omega) = \text{Re}\Sigma_\uparrow^{nn}(k, \omega) + \text{Re}\Sigma_\uparrow^{np}(k, \omega)$$

$$r_{\tau_1 \tau_2} = \frac{n_{\tau_1 \tau_2}(k)}{n(k)}$$

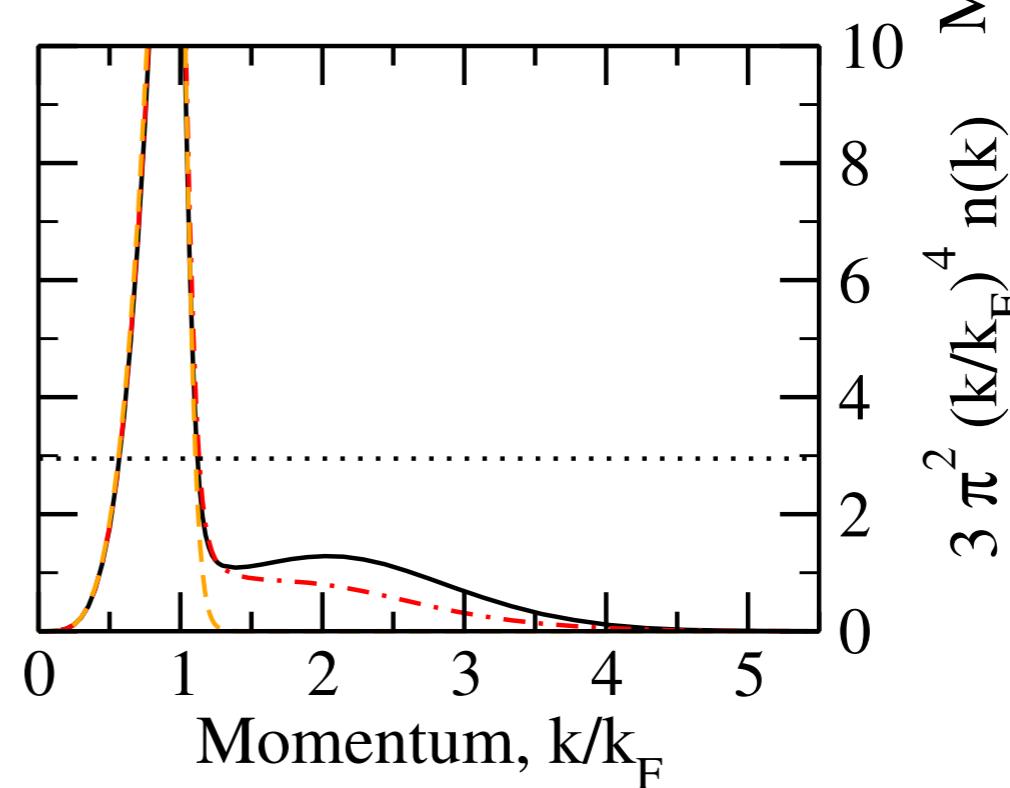
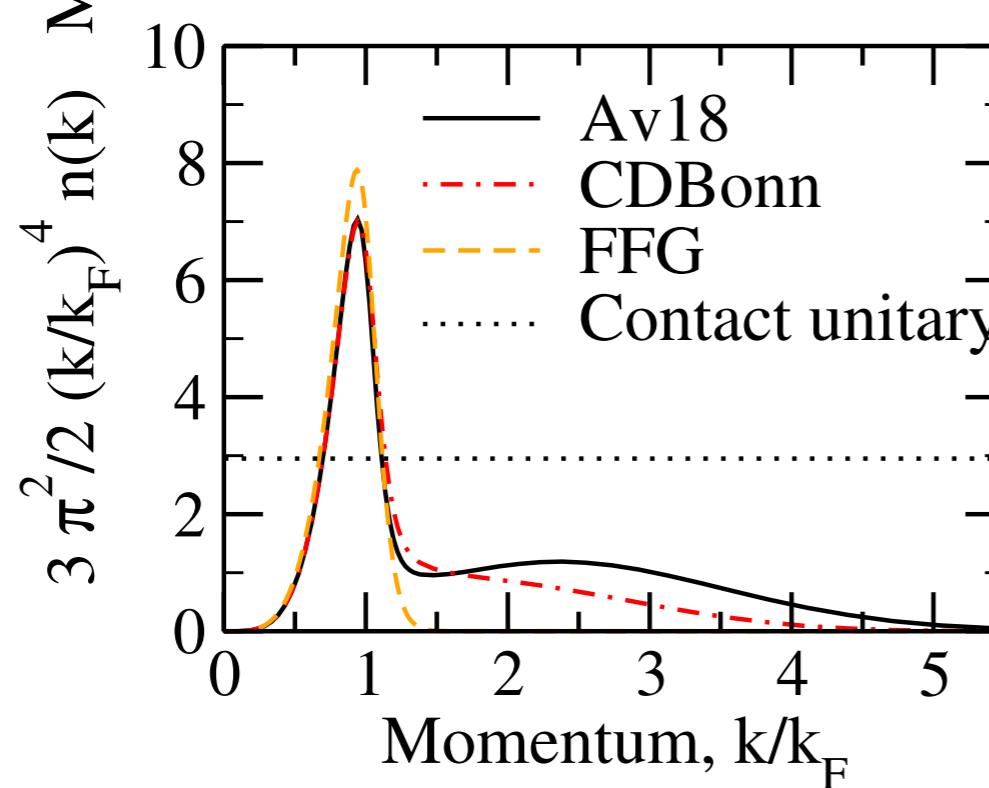
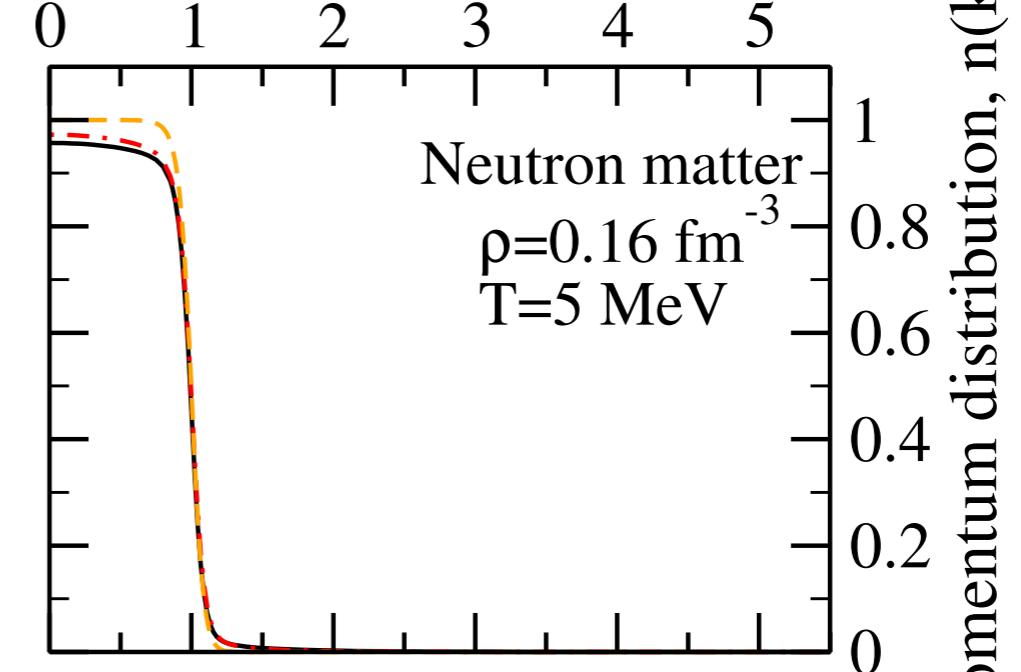
Kinetic symmetry energy I

$$\frac{K}{A} \approx \int \frac{d^3k}{(2\pi)^3} \frac{k^2}{2m} n(k) \approx \int_0^\infty dk k^4 n(k)$$

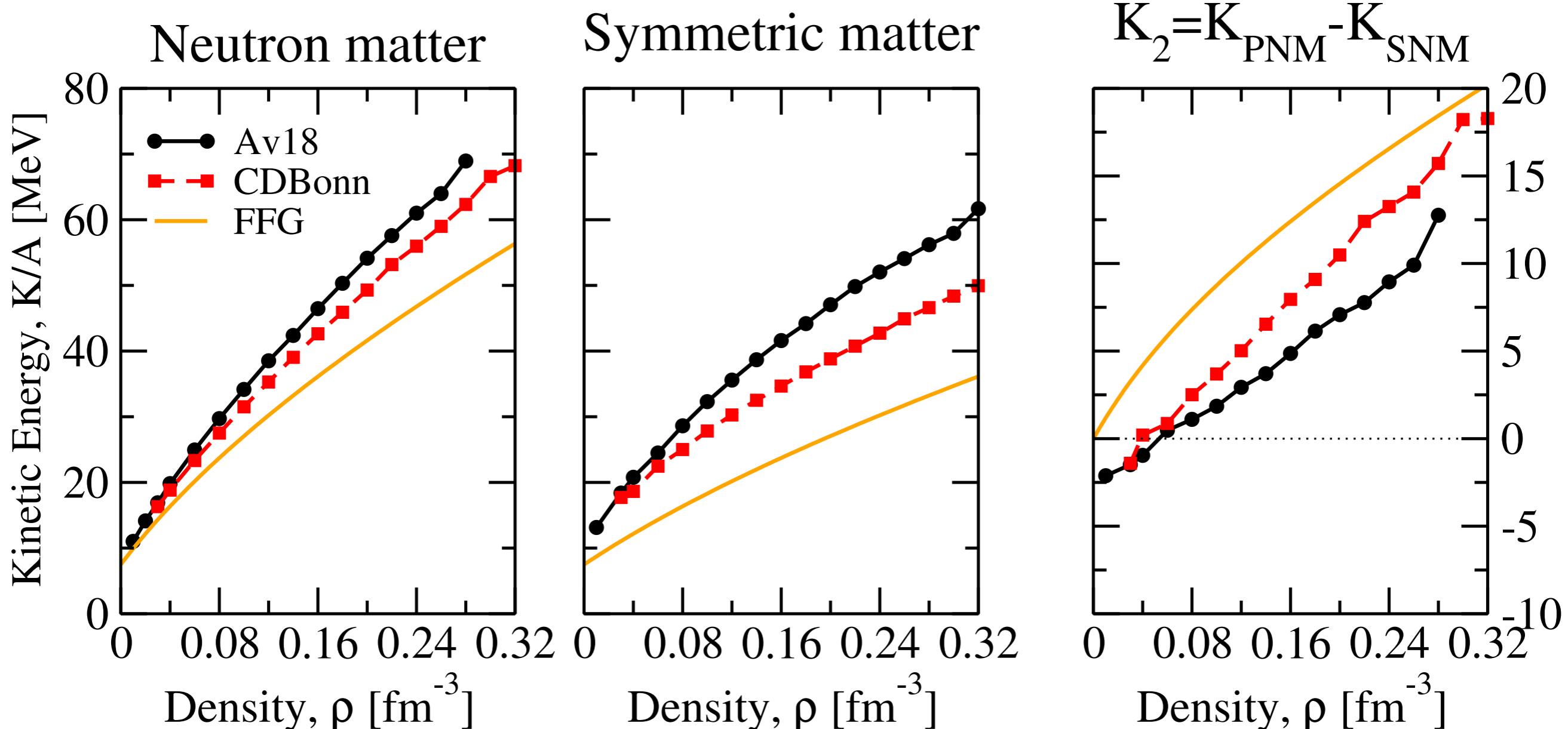
Momentum, k/k_F



Momentum, k/k_F

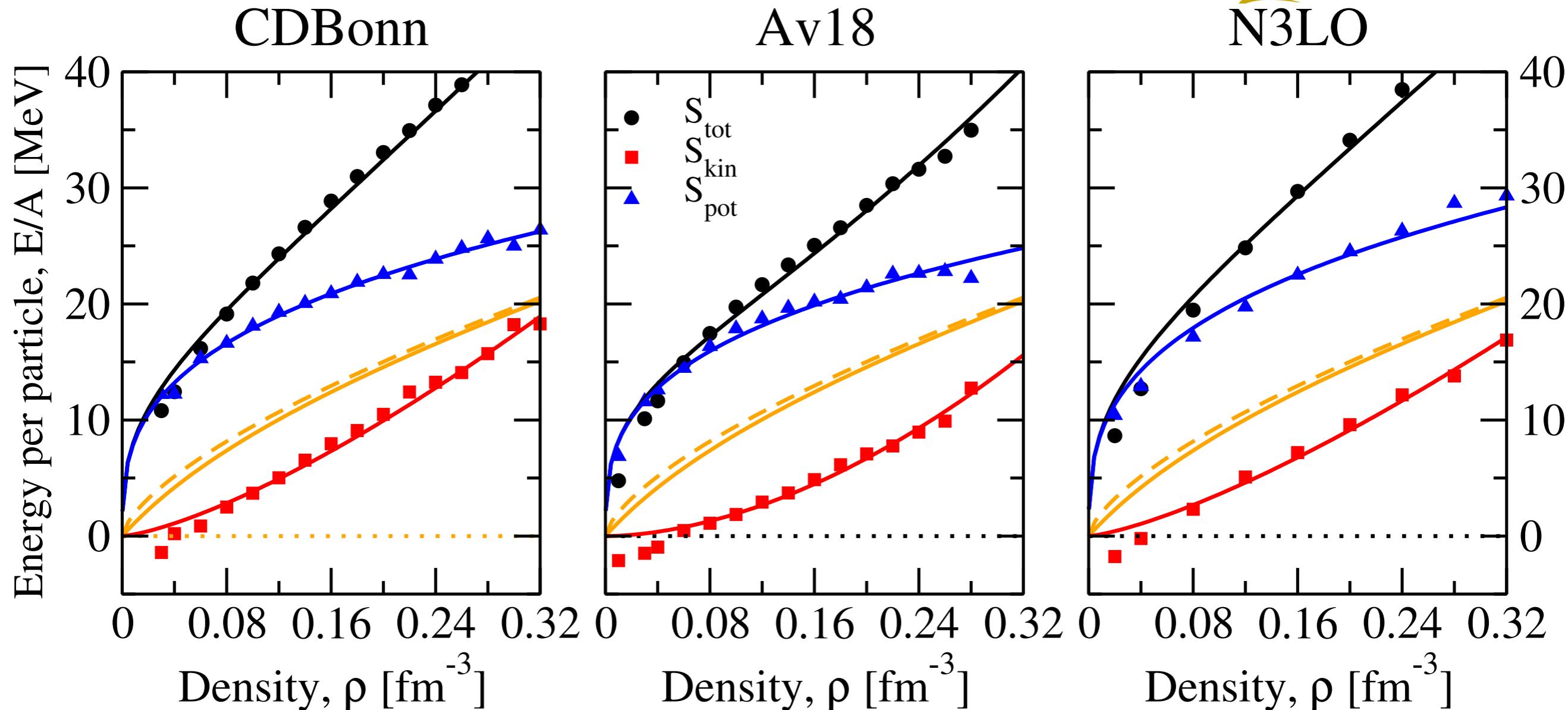


Symmetry energy consequences



- Kinetic component is reduced by correlations
- Implications for observables?

Symmetry energy



$$S(\rho) = S_k \left(\frac{\rho}{\rho_0} \right)^\alpha + S_p \left(\frac{\rho}{\rho_0} \right)^\gamma$$

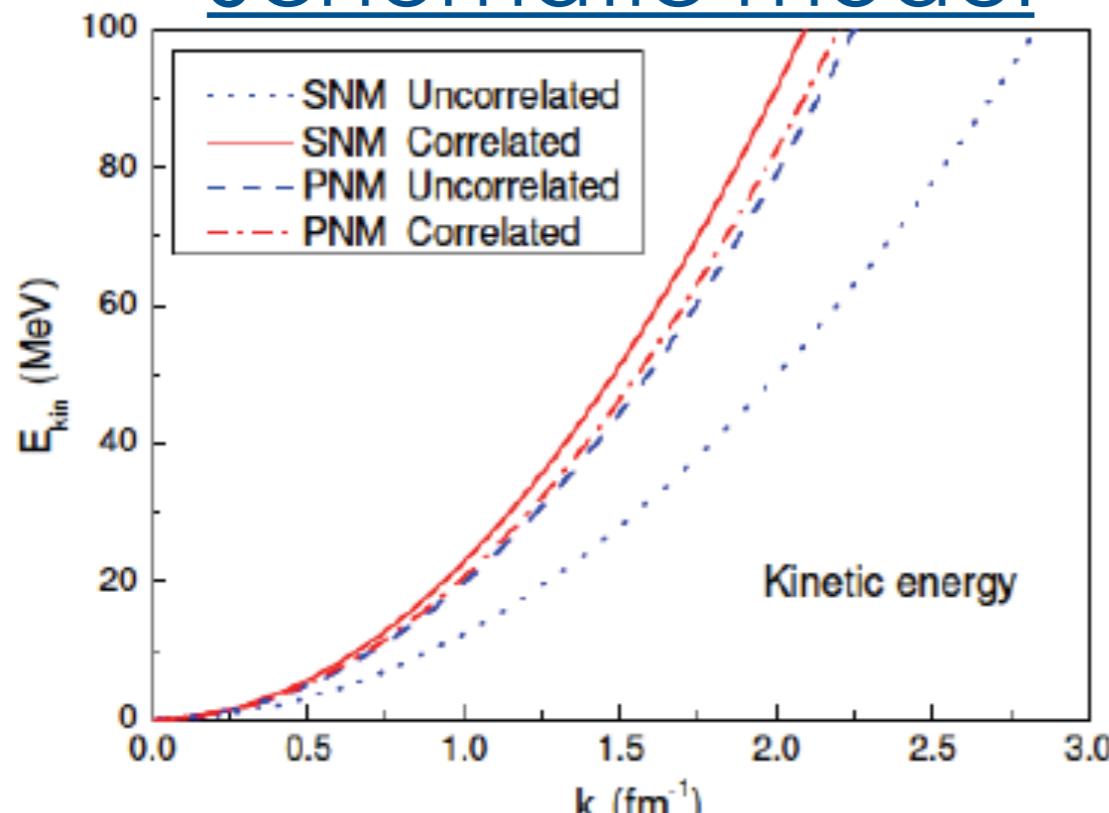
$$S_k = 4.5 - 7.5 \text{ MeV} \quad \alpha = 1.3 - 1.8$$

$$S_p = 20 - 22.5 \text{ MeV} \quad \gamma = 0.33$$

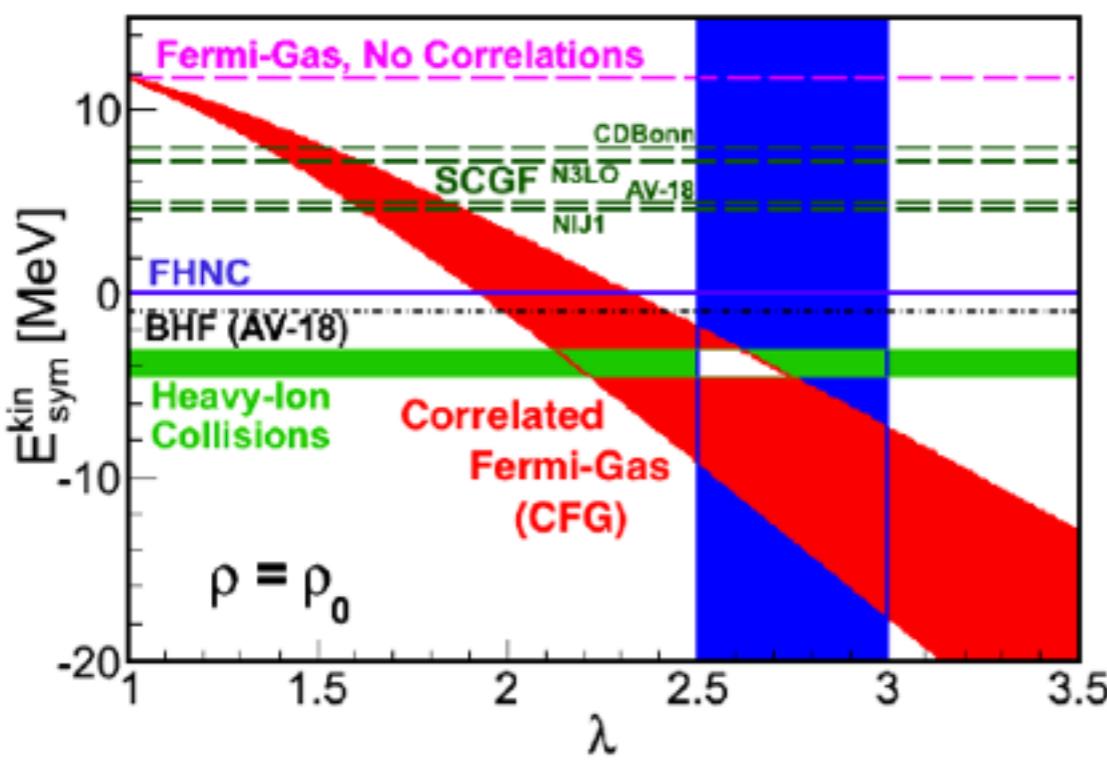
- Arbitrary / non-observable separation of energies
- Implications for observables? See **Or's** and **Bao-An's** talk

Confirmed by other approaches

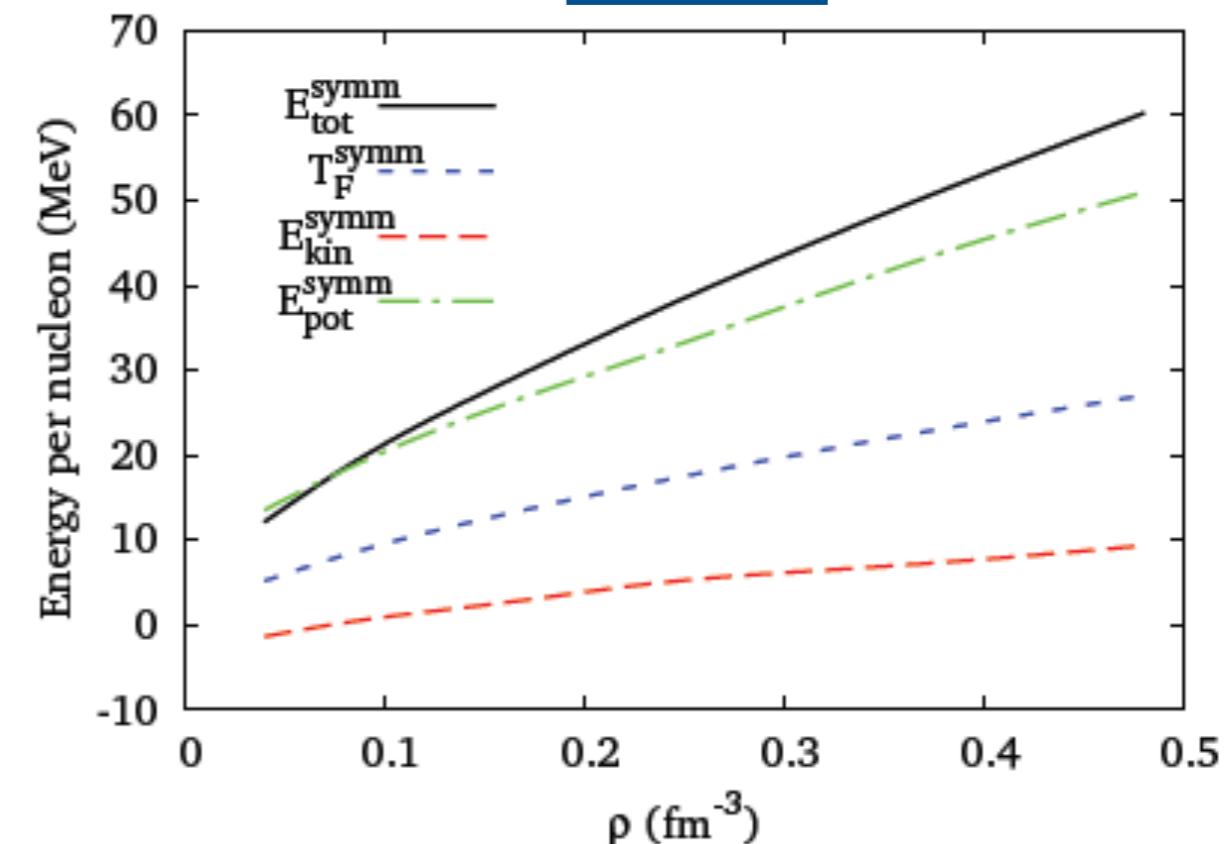
Schematic model



Xu & Li, arxiv:1104.2075



Hen et al., arxiv:1408.0772



A. Lovato, private communication

Av18+3BF, $\rho = 0.16$ fm $^{-3}$

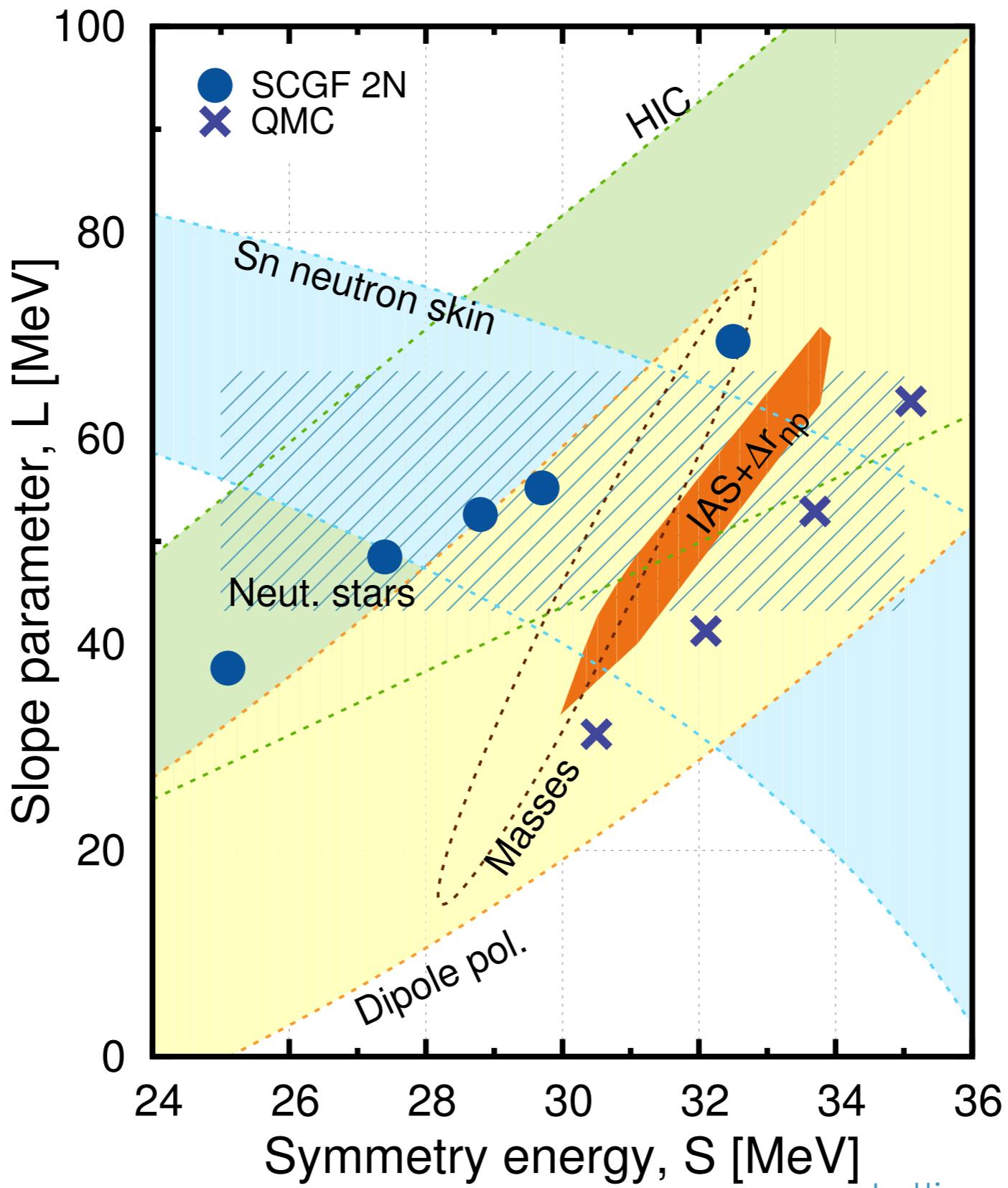
BHF + Hellman-Feynman theorem

	E_{PNM}	E_{SNM}	S_{tot}	L
K/A	53.3	54.3	-1.0	14.9
U/A	-34.2	-69.5	35.3	51.6
Total	19.1	-15.2	34.3	66.5

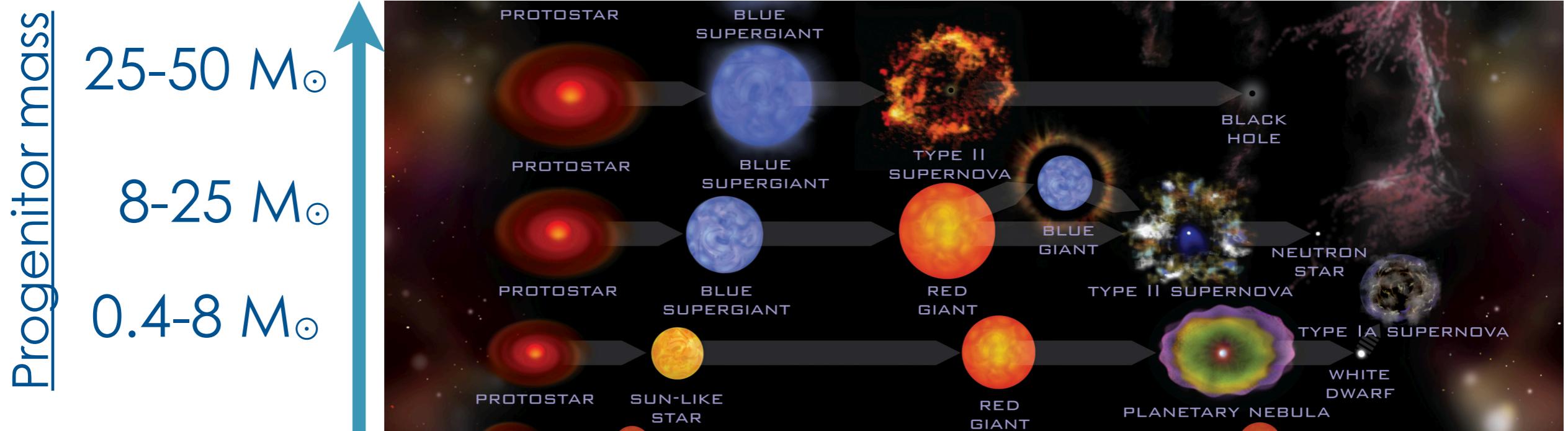
A. Carbone, et al. EPJA **50** 13 (2014)

Comparison to phenomenology

Experimental constrains

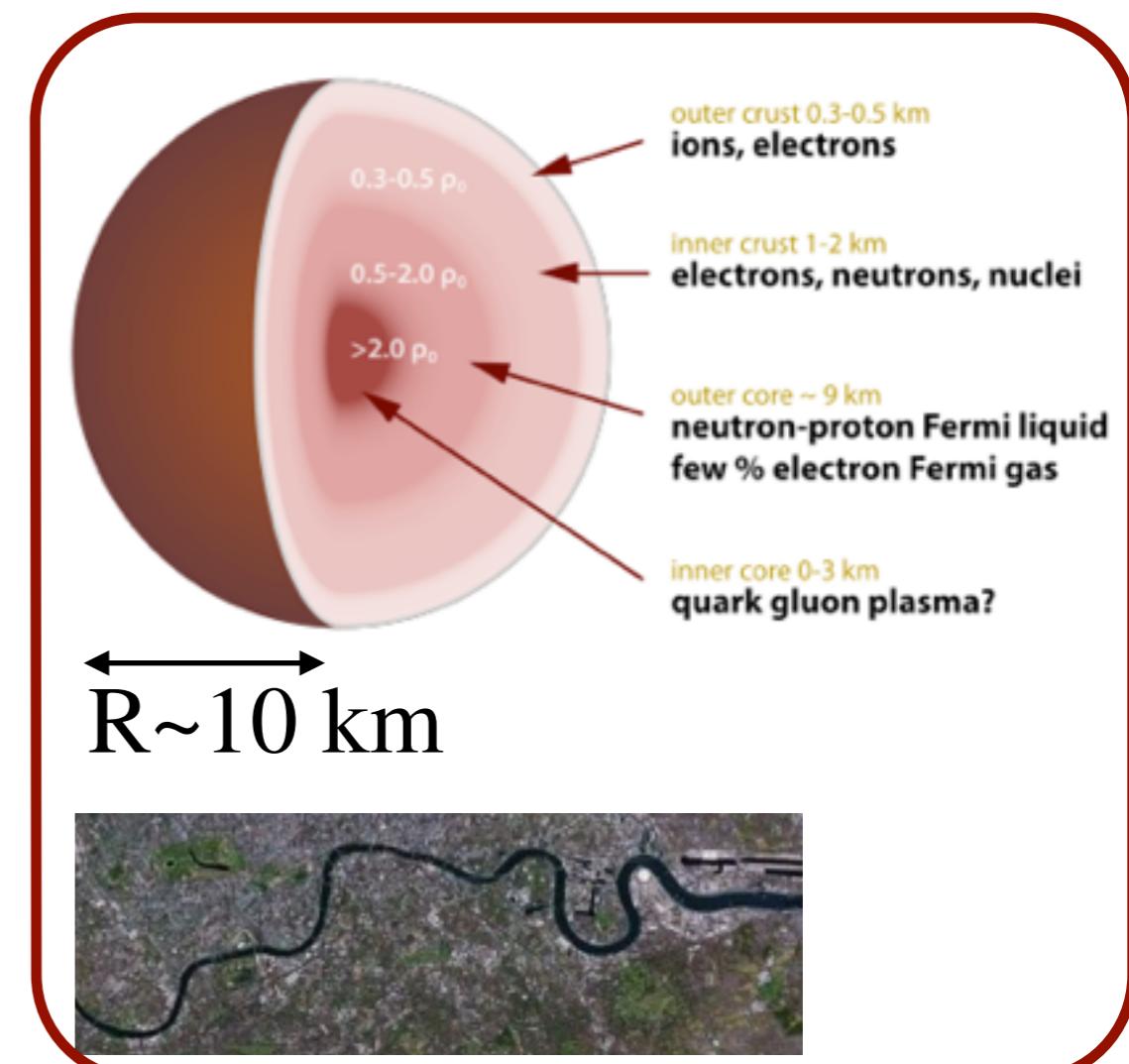
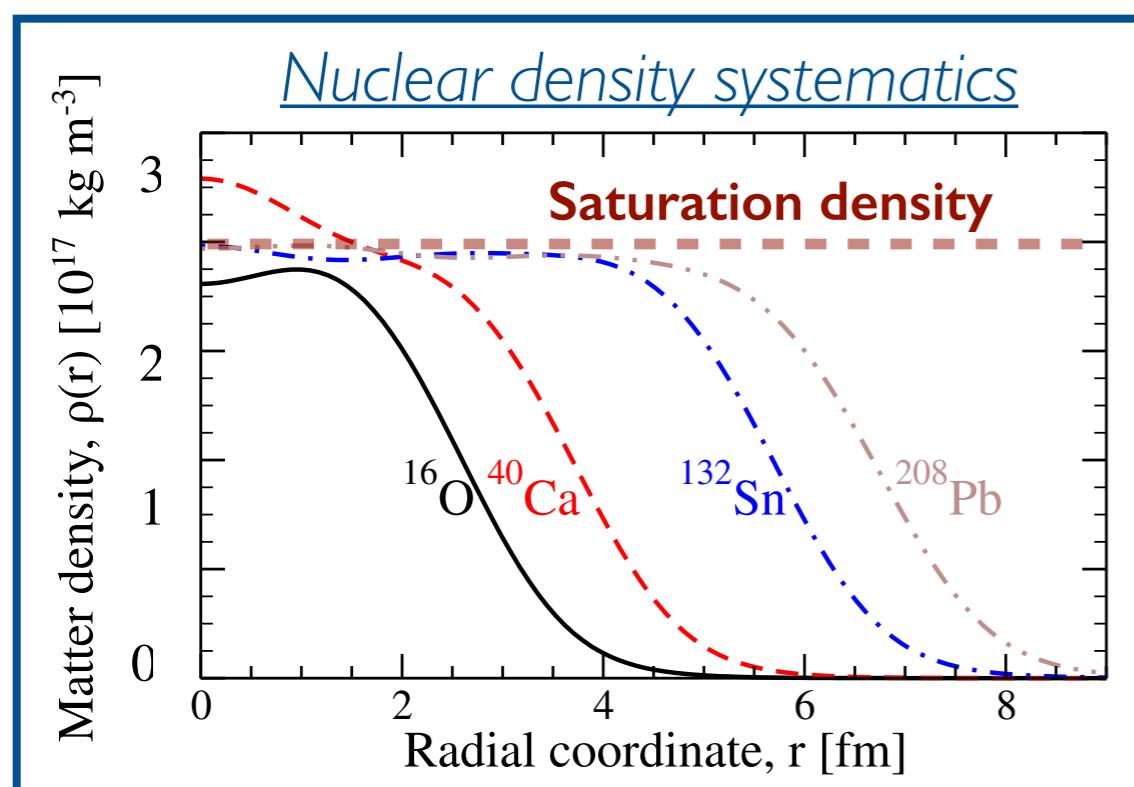


What is a neutron star?

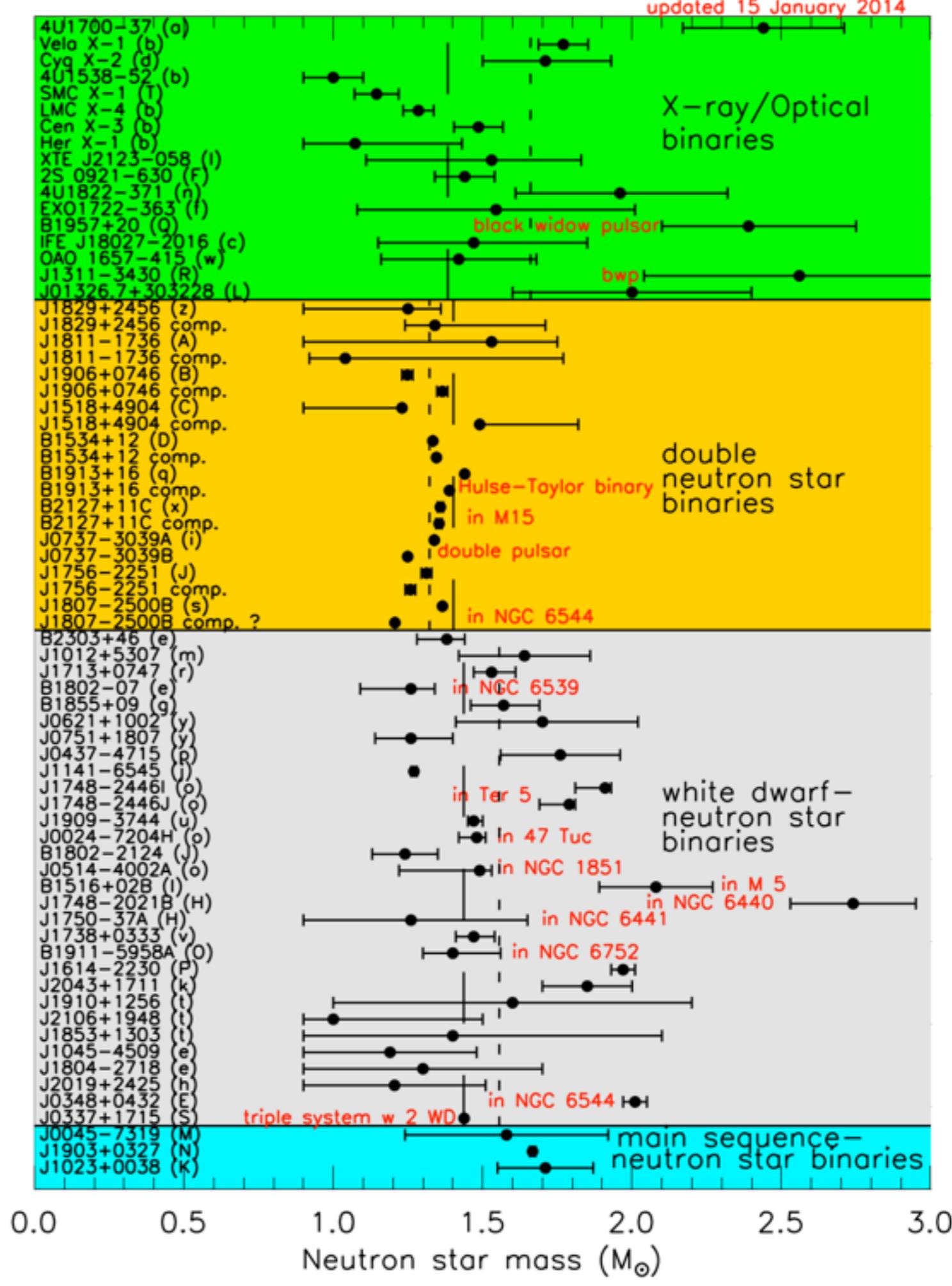
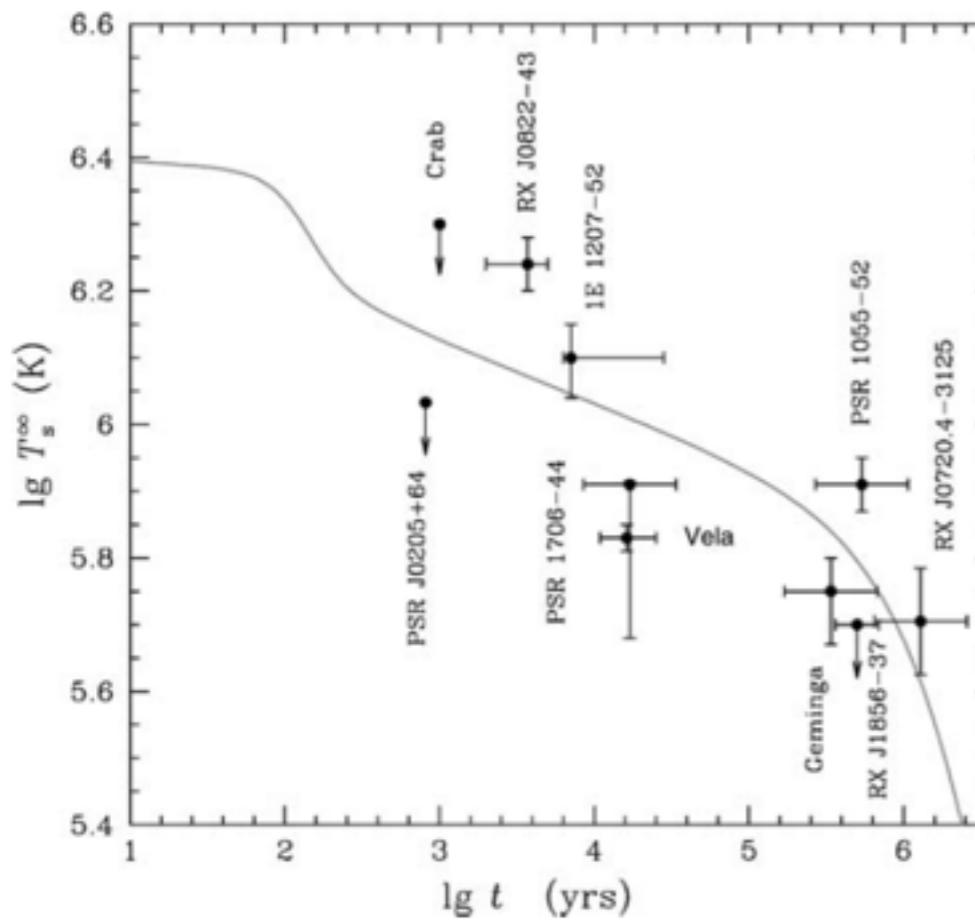
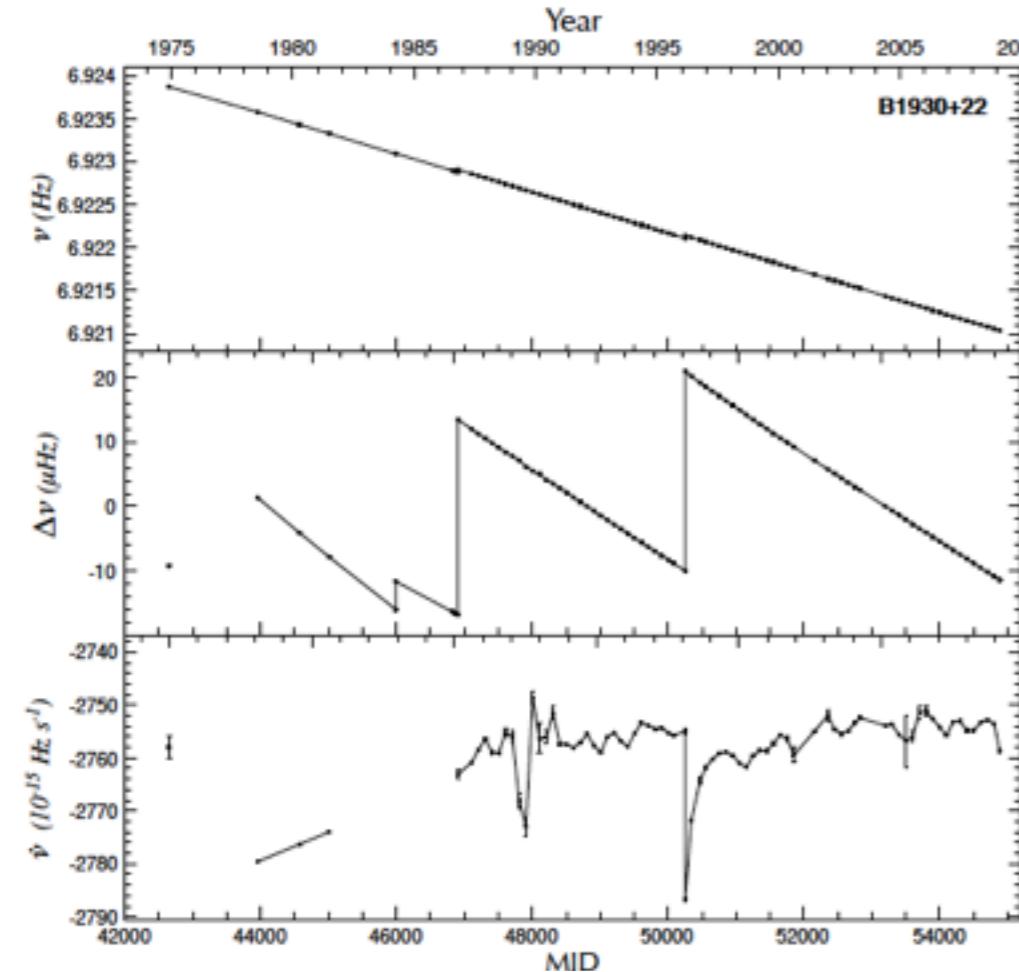


Mass $M \sim 1.5M_{\odot} = 3 \times 10^{30}$ kg

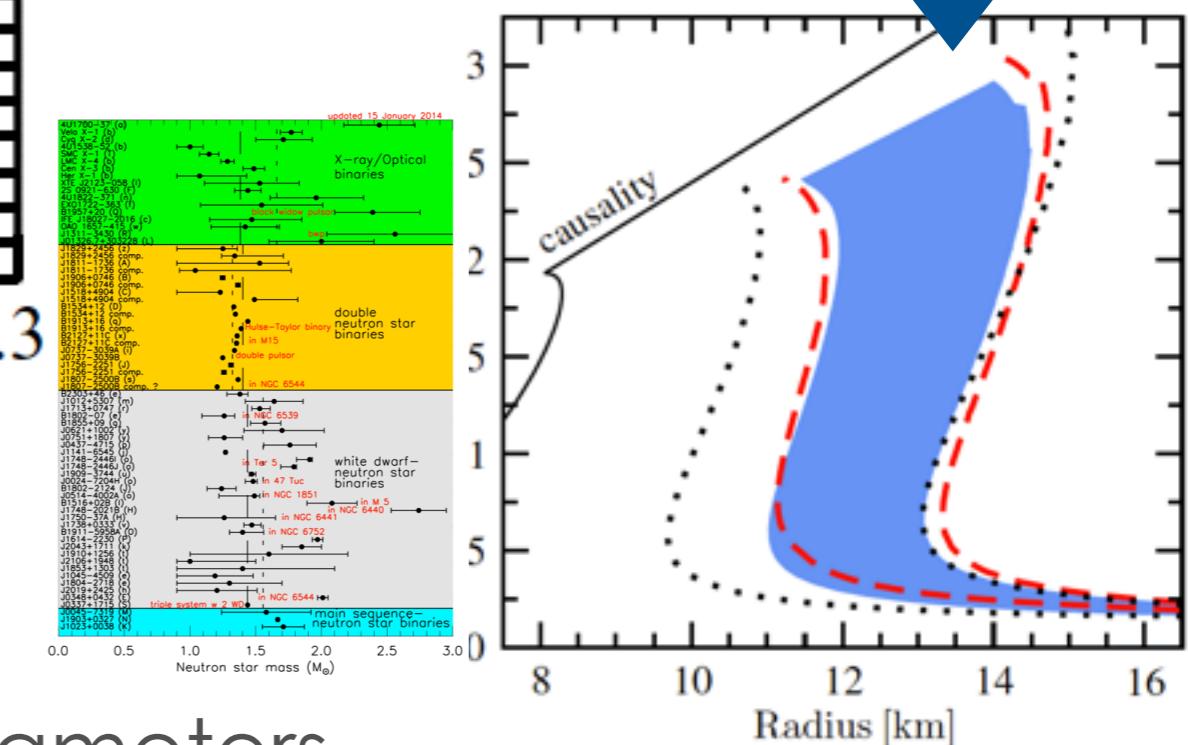
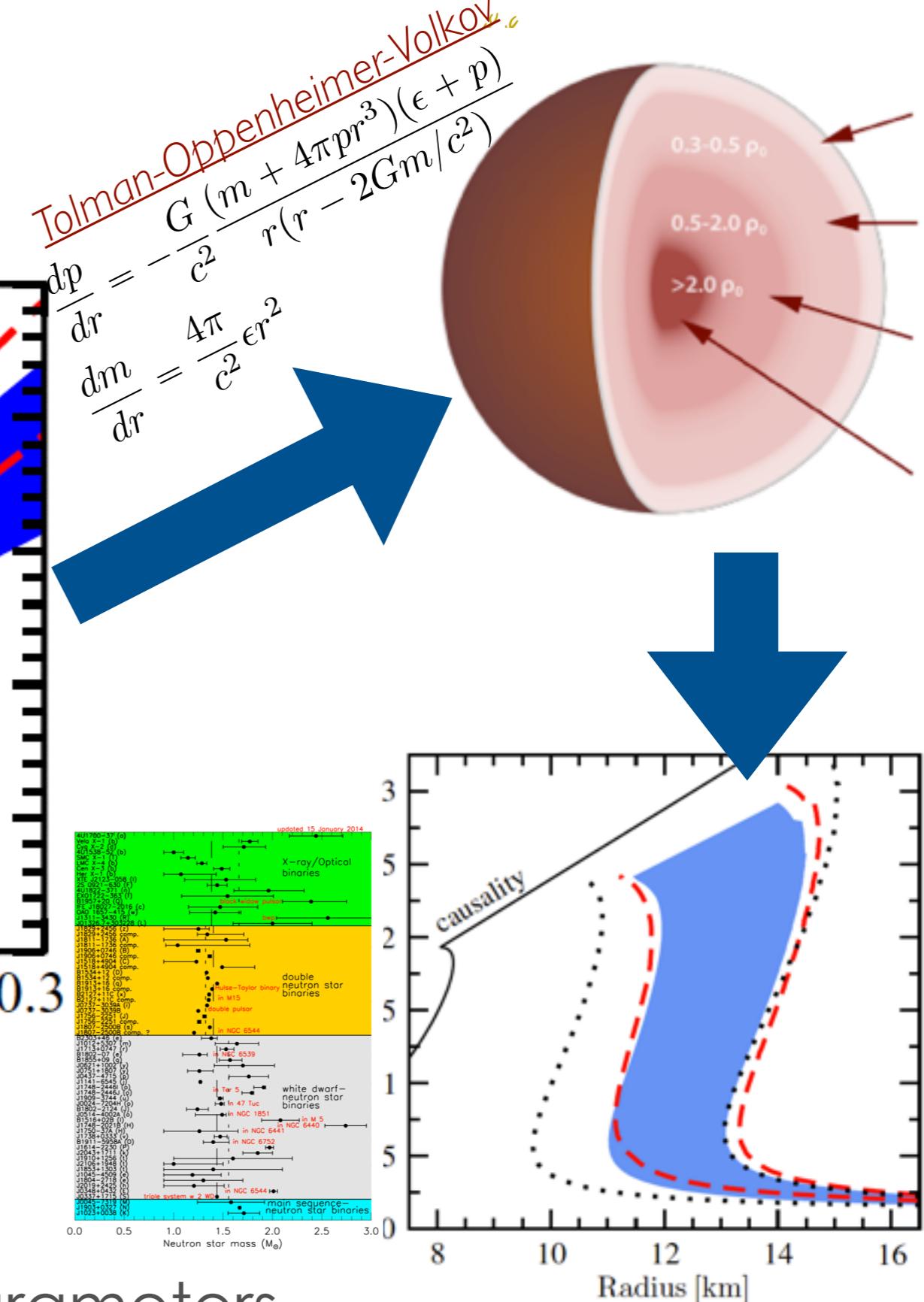
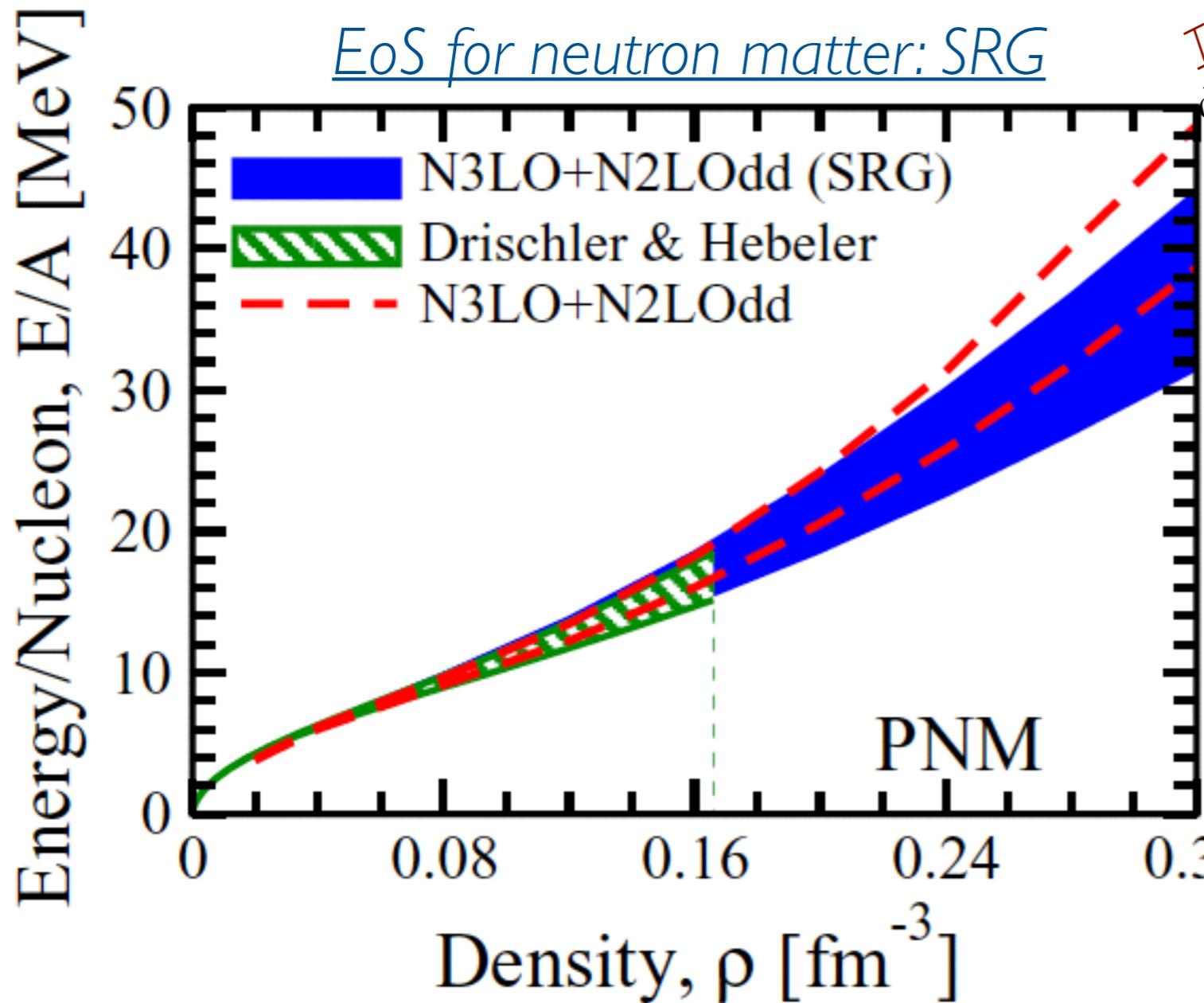
Mass density $\rho = \frac{M}{V} \approx 7.5 \times 10^{17} \text{ kg m}^{-3}$



NS data!



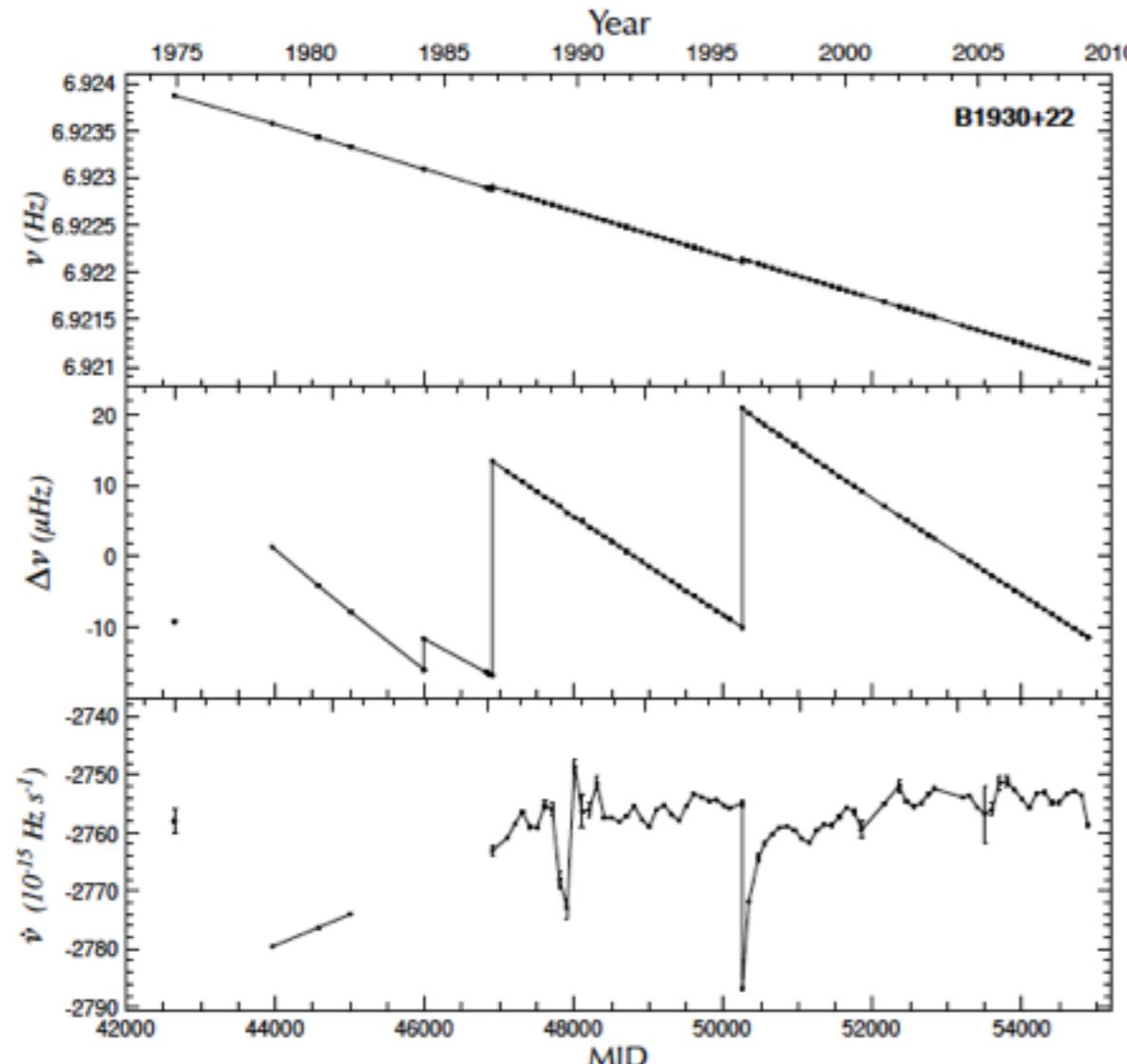
Neutron matter



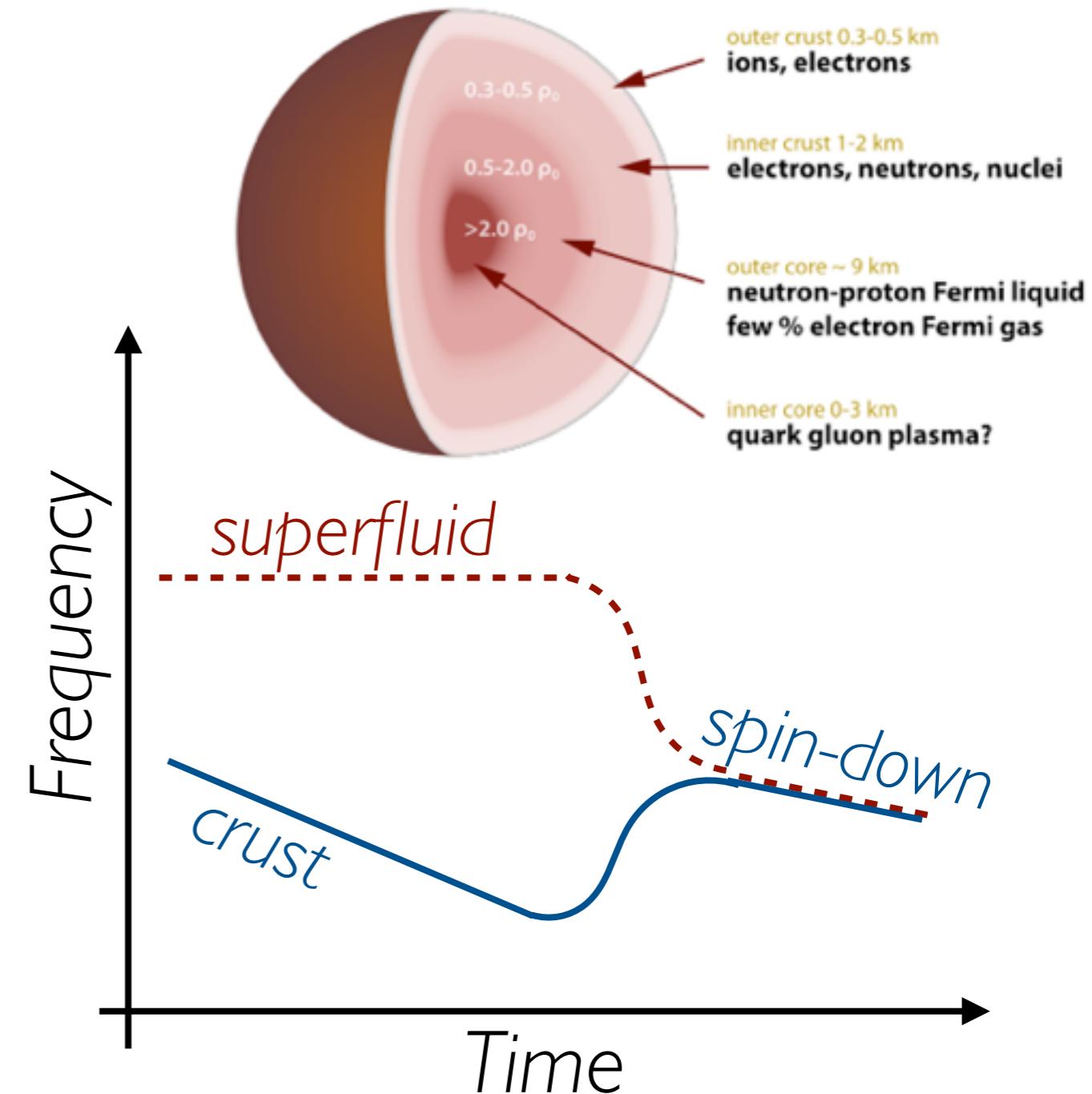
- Error band from fits in ChPT c_1, c_3 parameters
- Finite temperature & higher densities available

Hebeler, Lattimer, Pethick, Schwenk
ApJ 773 11 (2013)

Pulsar glitches



Espinoza, Lyne, Stappers & Kramer
MNRAS **414** 1679 (2011)



- **Crystalline** crust + dripped neutron **superfluid**
- Crust **slows down** due to magnetic braking
- **Superfluid** can only spin if **vortices** move out
- If vortices are **pinned** to nuclear lattice, they experience a time lag
- At some critical **pile-up**, a lot of vortices are **released** and crust spins up

Superfluid in the core

RESEARCH ARTICLE

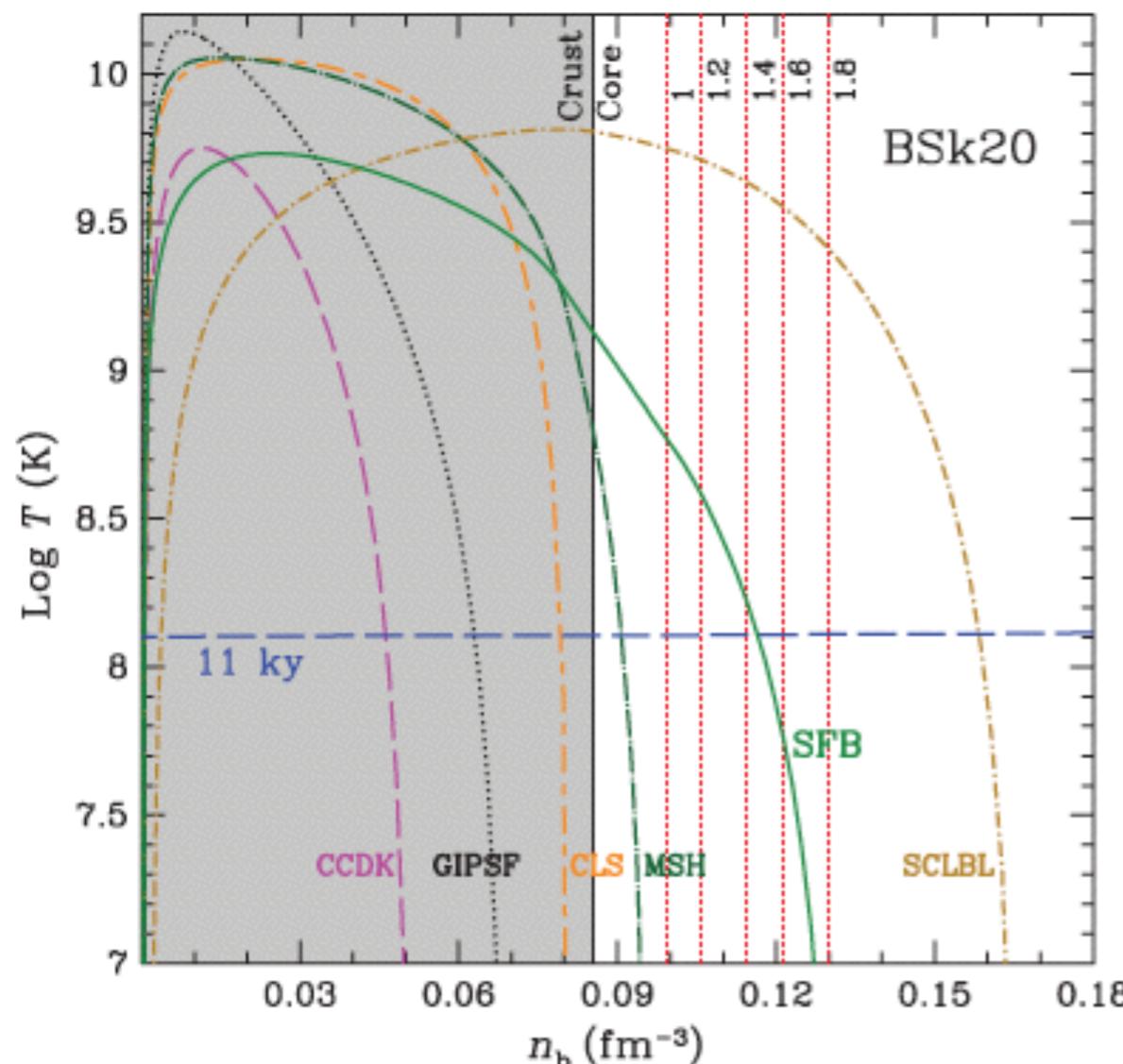


Fig. 2. Temperature dependence of neutron superfluid models as a function of baryon number density for the BSk20 nuclear equation of state. Thick curved lines are the superfluid critical temperature for the (labeled) models from (23). The vertical solid line indicates the separation between the crust (shaded region) and the core. Vertical dotted lines denote the density at which the superfluid moment of inertia (using the SFB superfluid model) is 1.6% of the total stellar moment of inertia for neutron stars of different mass (labeled in units of solar mass). The (nearly horizontal) dashed line is the temperature of a $1.4 M_{\odot}$ neutron star at an age of 11,000 years.

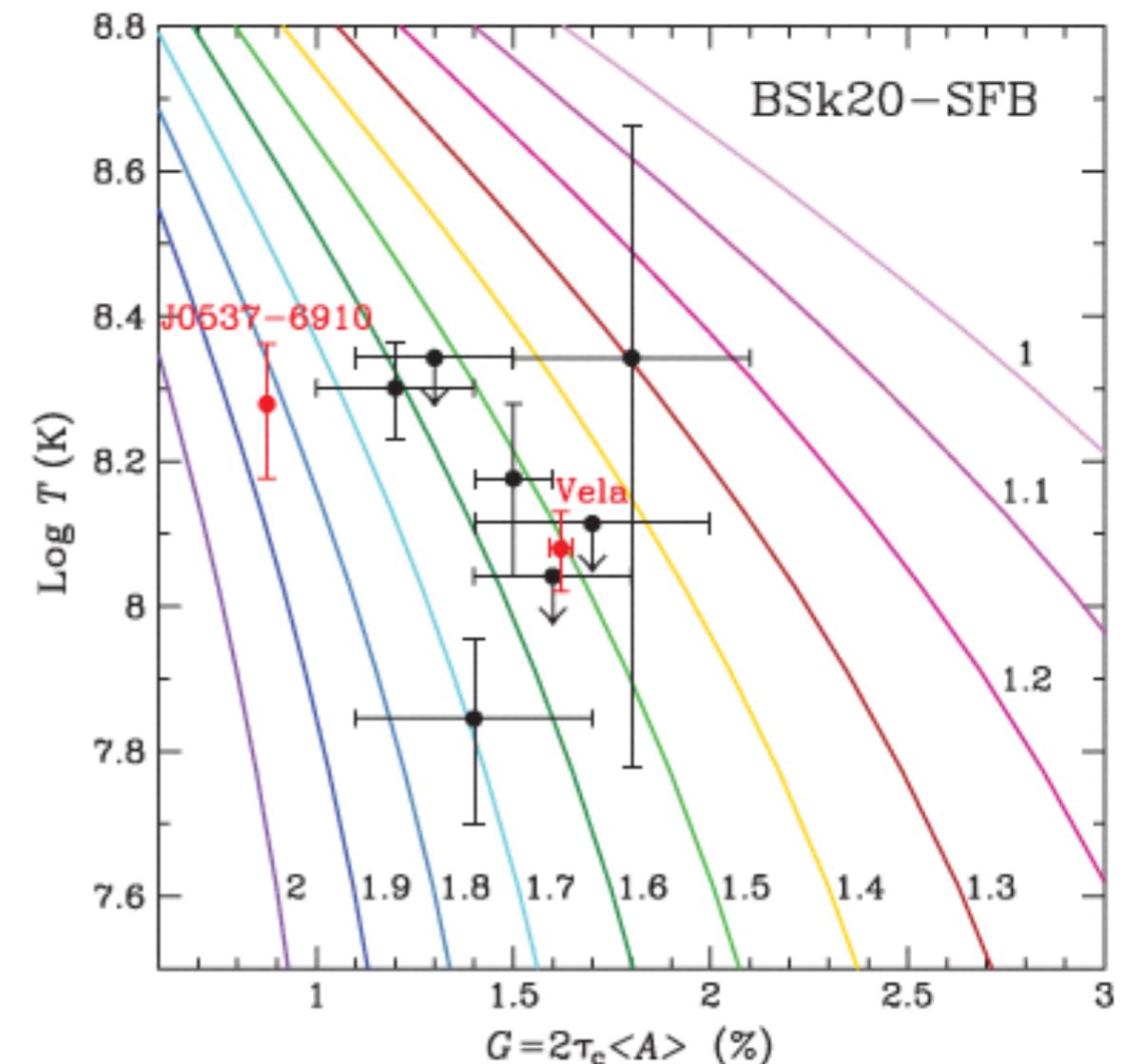
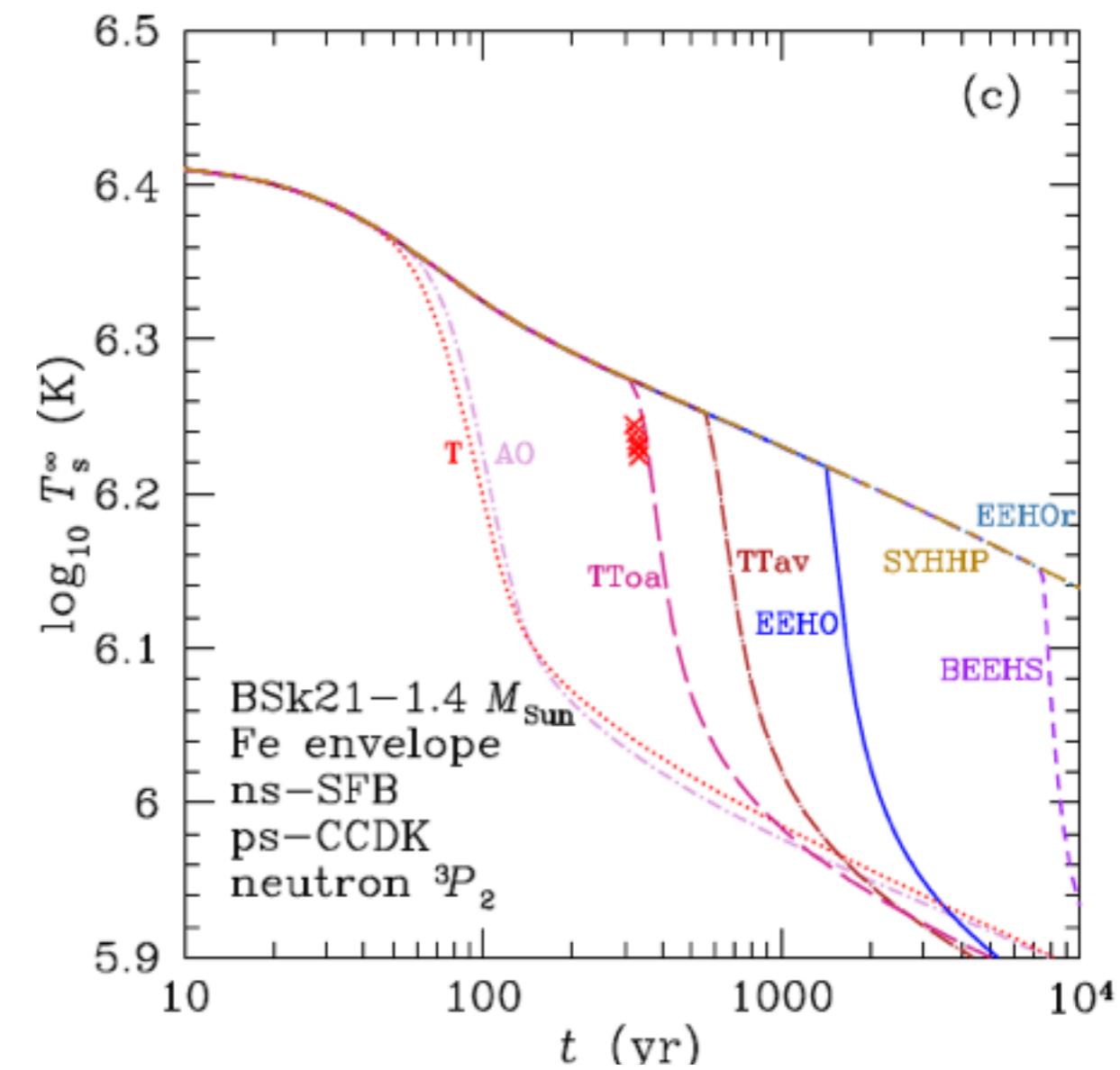
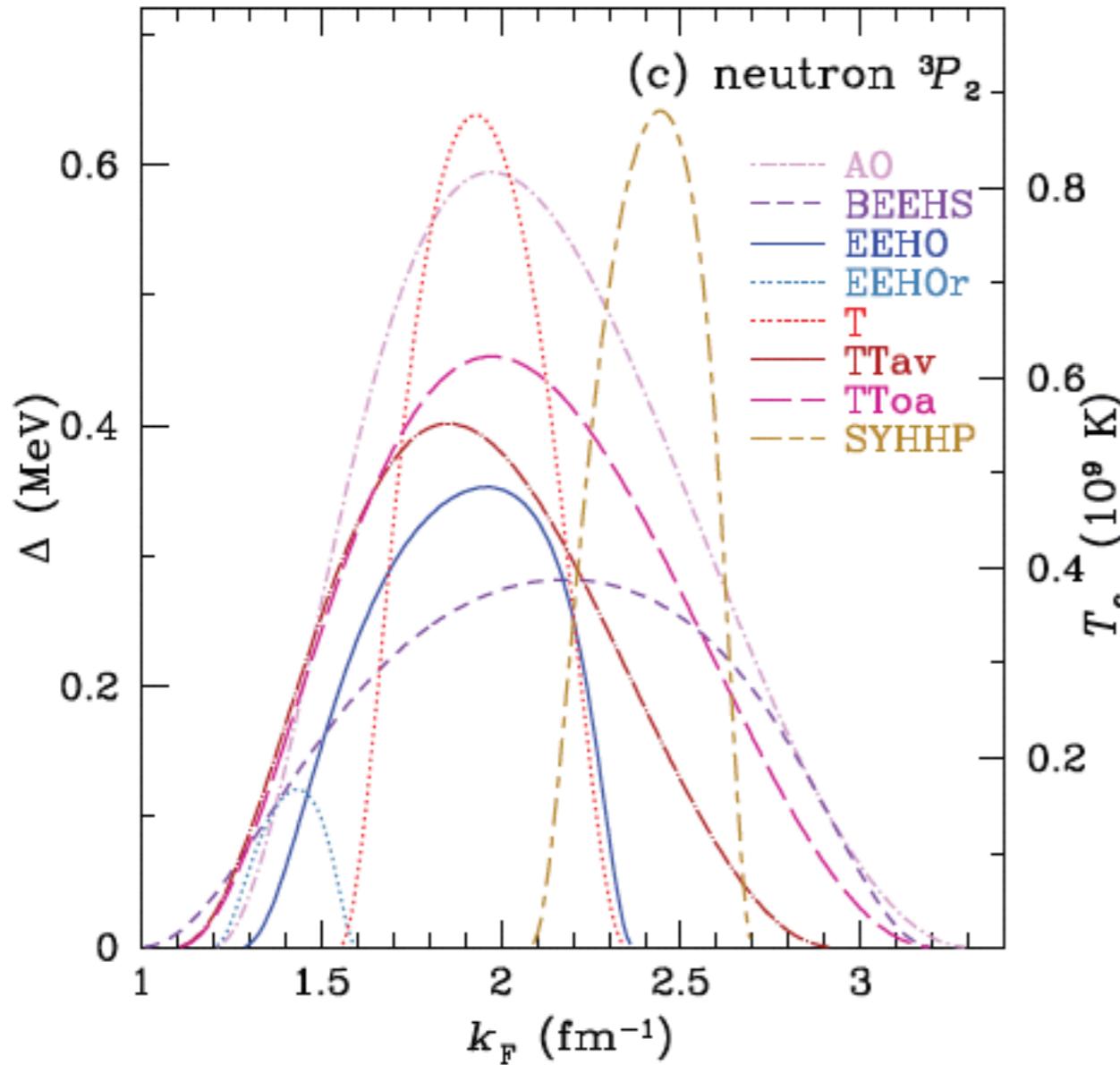


Fig. 3. Neutron star mass from pulsar observables G and interior temperature T . Data points are for pulsars with measured G from glitches and T from an age or surface temperature observation (see Table 1). Lines (labeled by neutron star mass, in units of solar mass) are the theoretical prediction for G and T using the BSk20 nuclear equation of state and SFB neutron superfluid models.

pulsar; dashed line) for a $1.4 M_{\odot}$ neutron star built using the BSk20

Cooling of Cassiopea A



Ingredients

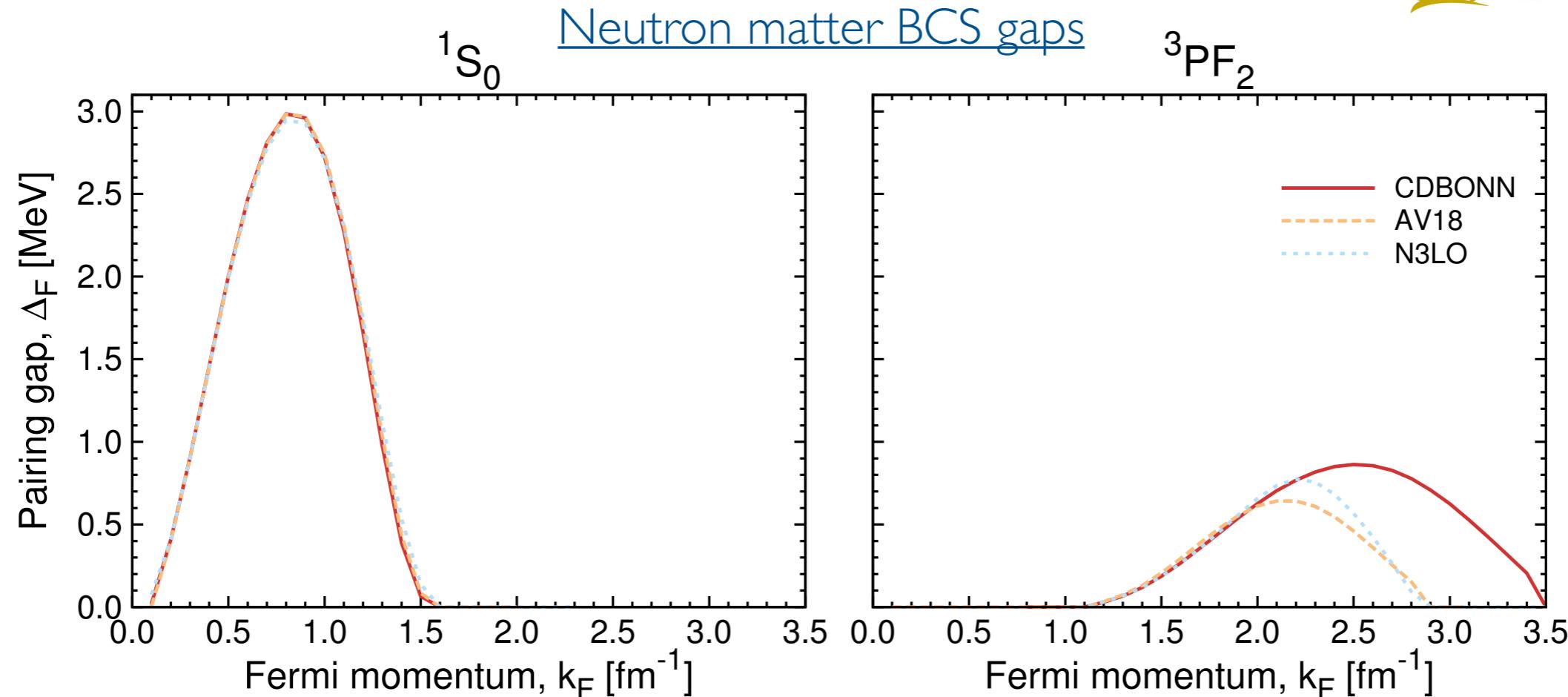
- (a) Mass of pulsar
- (b) EoS (determines radius)
- (c) Internal composition
- (d) **Pairing gaps** (1S_0 & 3P_2 channels)
- (e) Atmosphere composition

Name	Process	Emissivity (erg cm $^{-3}$ s $^{-1}$)
Modified Urca (neutron branch)	$n + n \rightarrow n + p + e^- + \bar{\nu}_e$ $n + p + e^- \rightarrow n + n + \nu_e$	$\sim 2 \times 10^{21} R T_9^8$
Modified Urca (proton branch)	$p + n \rightarrow p + p + e^- + \bar{\nu}_e$ $p + p + e^- \rightarrow p + n + \nu_e$ $n + n \rightarrow n + n + \nu + \bar{\nu}$	$\sim 10^{21} R T_9^8$
Bremsstrahlungs	$n + p \rightarrow n + p + \nu + \bar{\nu}$ $p + p \rightarrow p + p + \nu + \bar{\nu}$	$\sim 10^{19} R T_9^8$
Cooper pair	$n + n \rightarrow [nn] + \nu + \bar{\nu}$ $p + p \rightarrow [pp] + \nu + \bar{\nu}$	$\sim 5 \times 10^{21} R T_9^7$ $\sim 5 \times 10^{19} R T_9^7$
Direct Urca (nucleons)	$n \rightarrow p + e^- + \bar{\nu}_e$ $p + e^- \rightarrow n + \nu_e$	$\sim 10^{27} R T_9^6$

Bardeen-Cooper-Schrieffer pairing



UNIVERSITY OF
SURREY



Dean & Hjorth-Jensen, *Rev. Mod. Phys.* **75** 607 (2003)

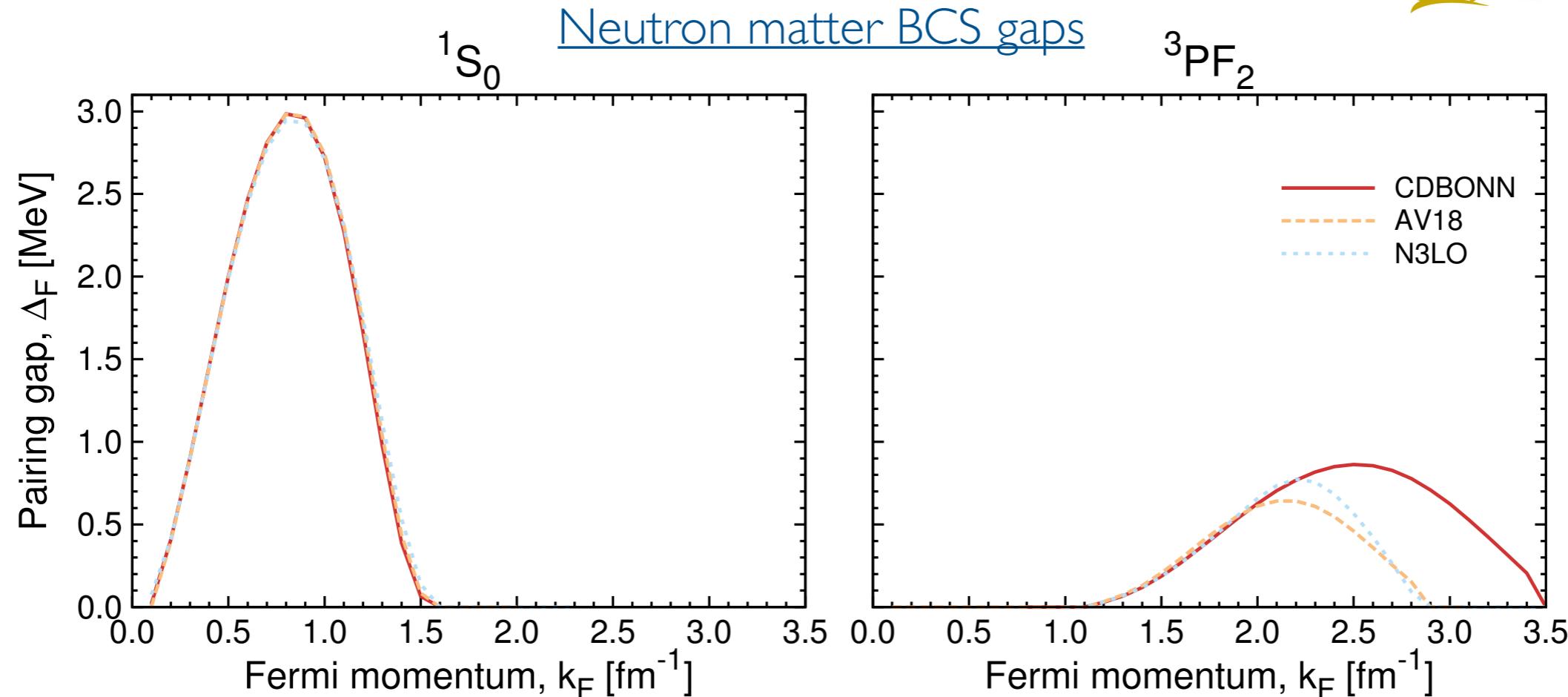
BCS equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} + \chi_k = \varepsilon_k - \mu$$

- Single-particle spectrum choice: $\varepsilon_k = \frac{k^2}{2m} + U(k) - \mu$

- Angular gap dependence: $|\Delta_k|^2 = \sum_L |\Delta_k^L|^2$

Bardeen-Cooper-Schrieffer pairing



Dean & Hjorth-Jensen, *Rev. Mod. Phys.* **75** 607 (2003)

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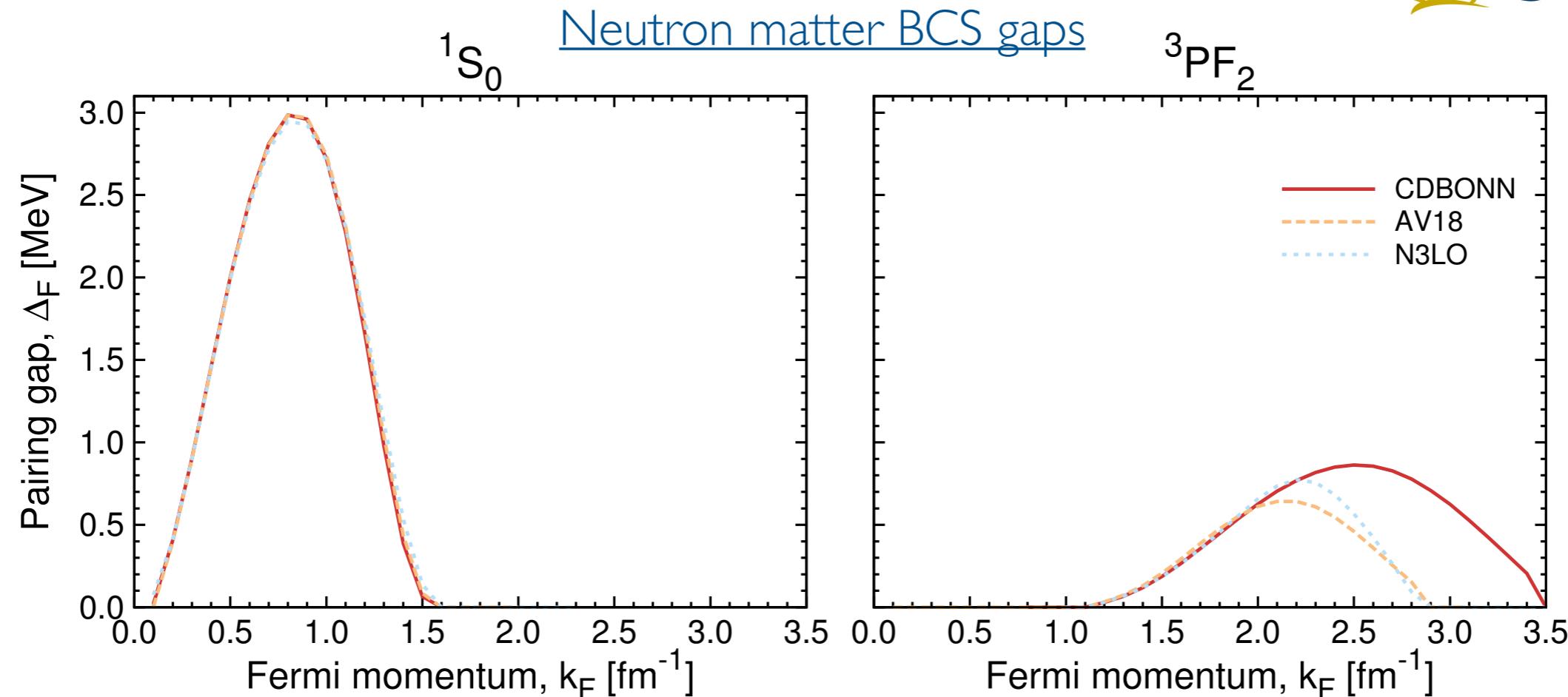
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UNIVERSITY OF
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$$(B) iF(1,2) = \langle T\{\psi(1)\psi(2)\} \rangle = \text{---} = \text{---} + \text{---}$$

Normal state

$$(B') iF^\dagger(1,2) = \langle T\{\psi^\dagger(1)\psi^\dagger(2)\} \rangle = \text{---} = \text{---} + \text{---}$$

$$(C) iG(1,2) = \langle T\{\psi(1)\psi^\dagger(2)\} \rangle = \text{---} = \text{---} + \text{---} + \text{---}$$

$$(D) \text{---} \Sigma \text{---} = \text{---} K \text{---}$$

$$(E) \text{---} \Delta \text{---} = \text{---} K \text{---}$$

Superfluid
 $\Delta(k_F)$

BCS+SRC equation

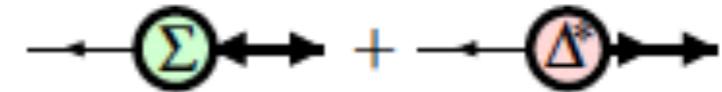
$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} +$$

$$\frac{1}{2\chi_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A_s(k, \omega')$$

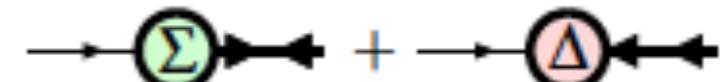
- **BCS** is lowest order in Gorkov Green's function expansion
- T-matrix can be extended to paired systems
- But full self-consistency is still missing



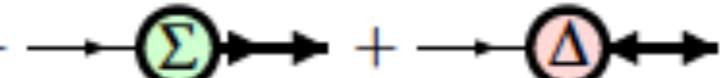
$$(B) iF(1,2) = \langle T\{\psi(1)\psi(2)\} \rangle = \text{---} =$$



$$(B') iF^\dagger(1,2) = \langle T\{\psi^\dagger(1)\psi^\dagger(2)\} \rangle = \text{---} =$$

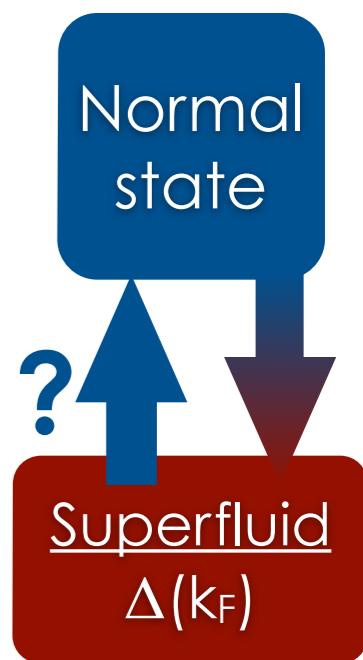


$$(C) iG(1,2) = \langle T\{\psi(1)\psi^\dagger(2)\} \rangle = \text{---} = \text{---} + \text{---} + \text{---}$$



$$(D) \rightarrow \Sigma \rightarrow = \rightarrow K \rightarrow$$

$$(E) \rightarrow \Delta \leftarrow = \rightarrow K \leftarrow$$



BCS+SRC equation

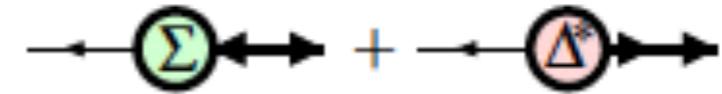
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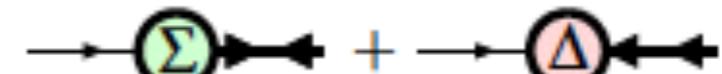
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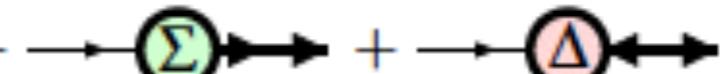
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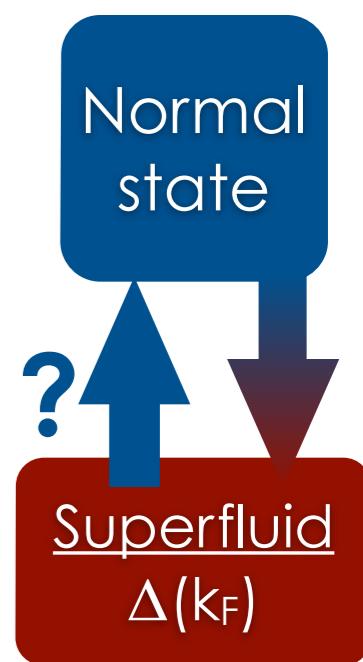


$$(C) iG(1,2) = \langle T\{\psi(1)\psi^\dagger(2)\} \rangle = \text{---} = \text{---} + \text{---} + \text{---}$$



$$(D) \rightarrow \Sigma \rightarrow = \rightarrow K \rightarrow$$

$$(E) \rightarrow \Delta \leftarrow = \rightarrow K \leftarrow$$



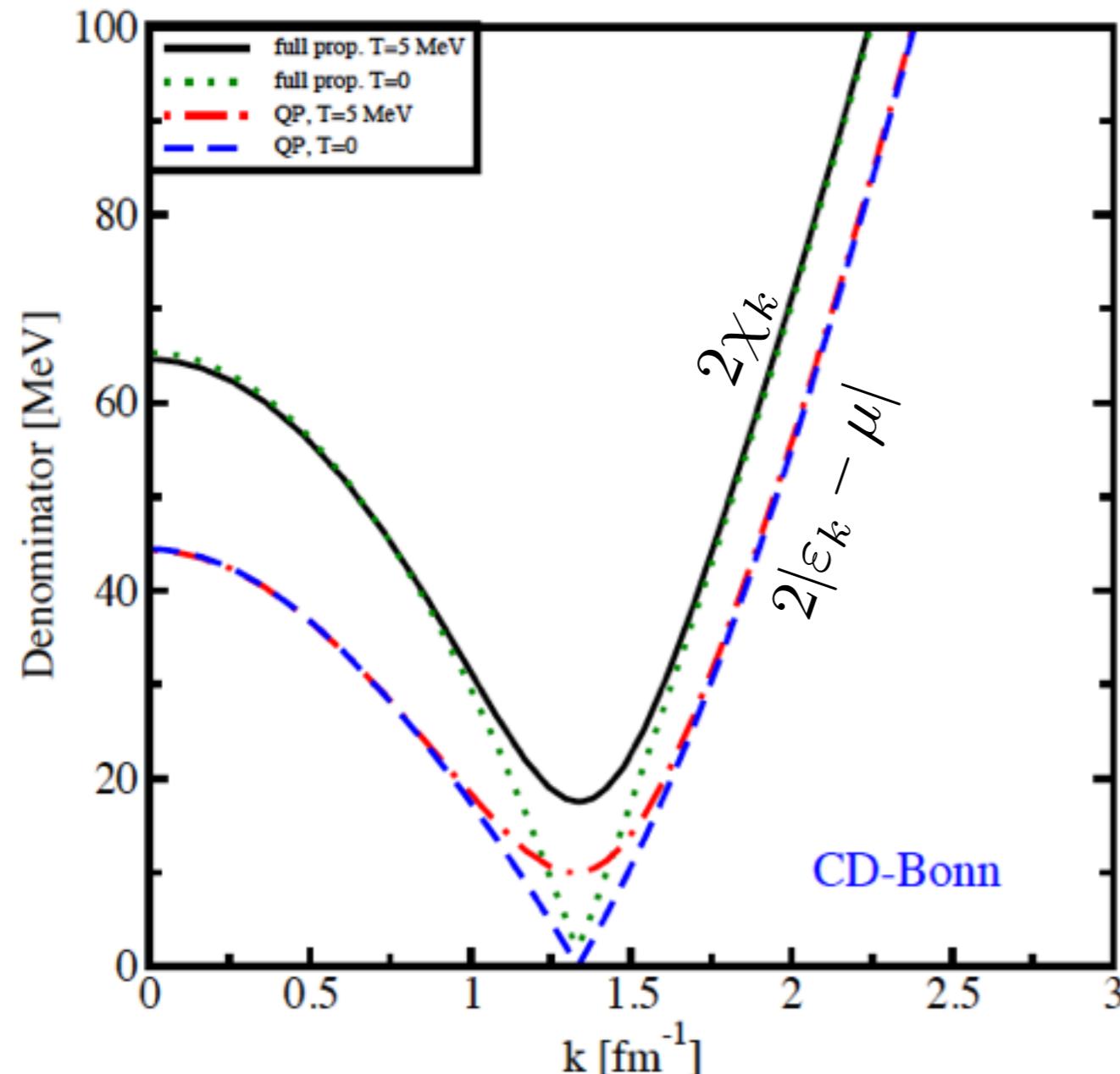
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- **BCS** is lowest order in Gorkov Green's function expansion
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- But full self-consistency is still missing

Nuclear matter, $\rho=0.16 \text{ fm}^{-3}$



Quasi-particle \leftrightarrow BCS

$$\frac{1}{2\bar{\chi}_k} = \frac{1}{2|\varepsilon_k - \mu|}$$

Full off-shell

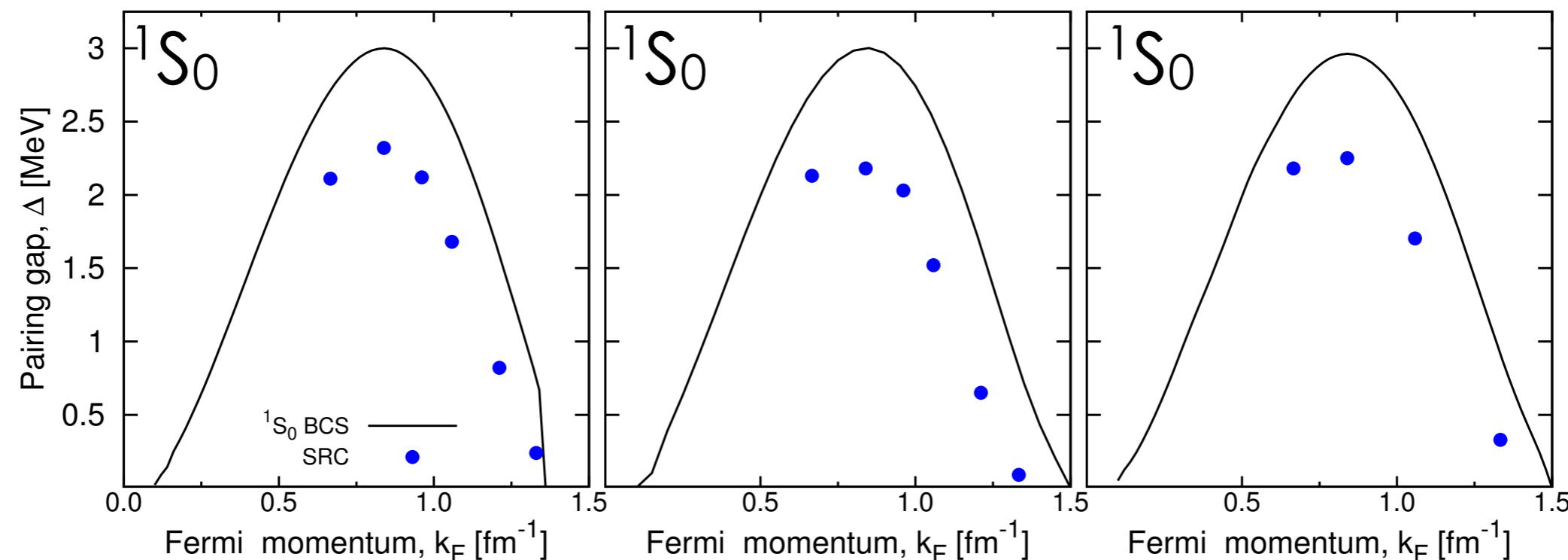
$$\frac{1}{2\bar{\chi}_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A(k, \omega')}$$

- How does **removal** of **strength** affect pairing?
- Necessarily **decreases** gap
- We use **normal** properties throughout

CDBonn

Av18

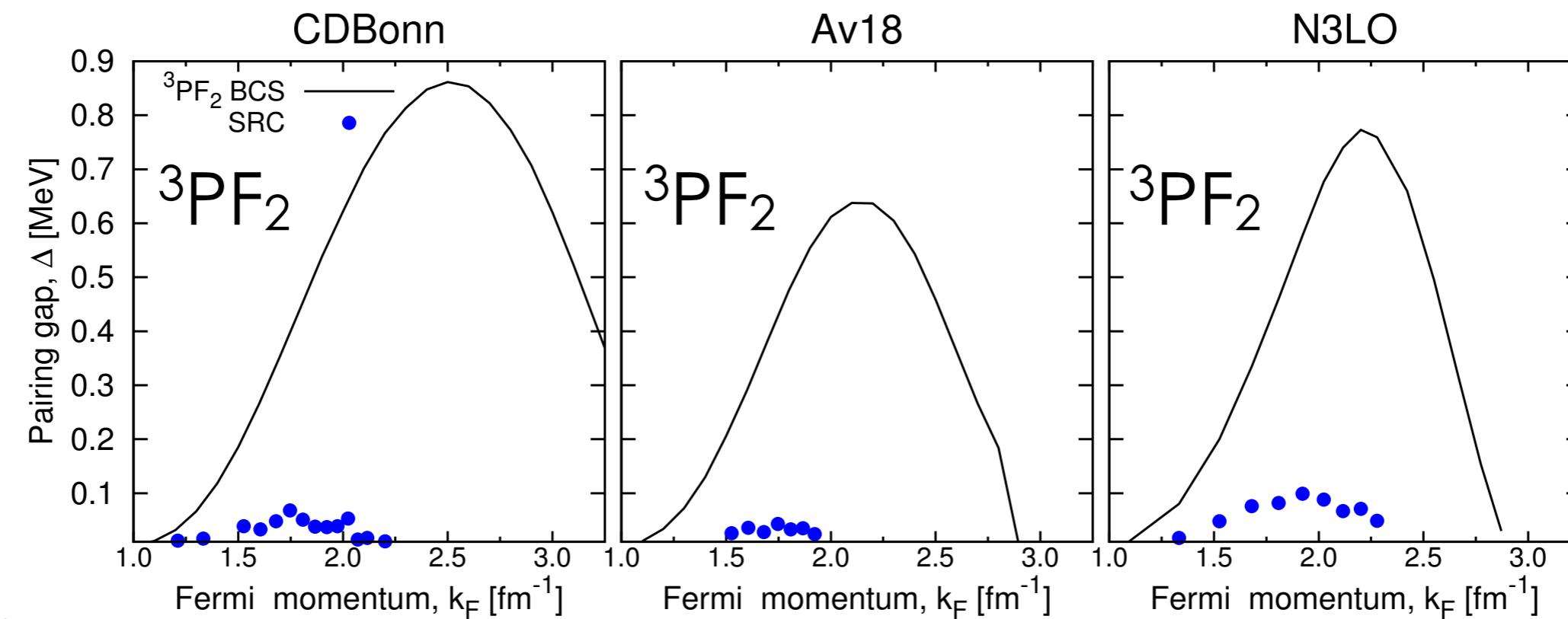
N3LO



CDBonn

Av18

N3LO

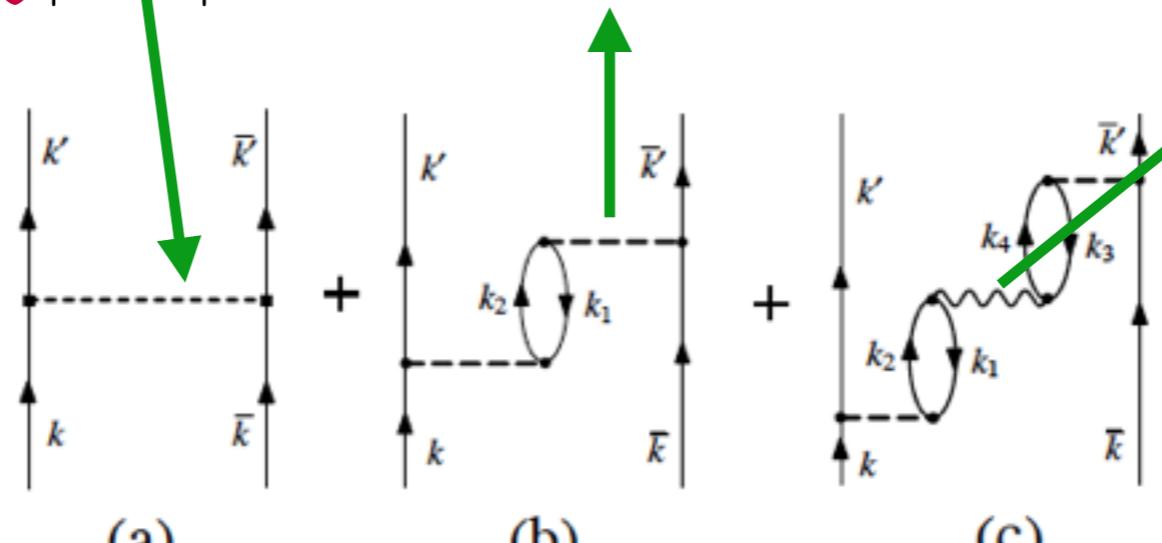


BCS+SRC

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} + \frac{1}{2\chi_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A_s(k, \omega')$$

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} \quad \text{ph recoupled G-matrix}$$

$\mathcal{V}_{\text{pair}} =$



?

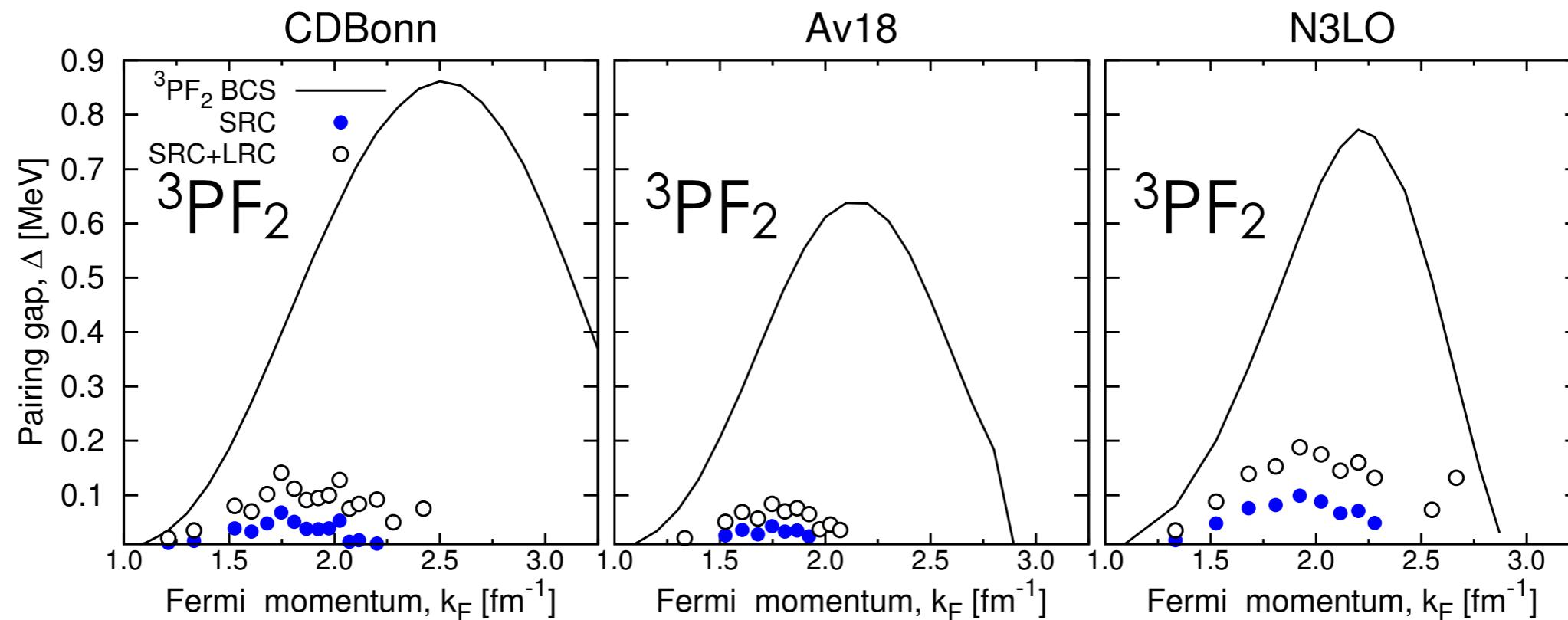
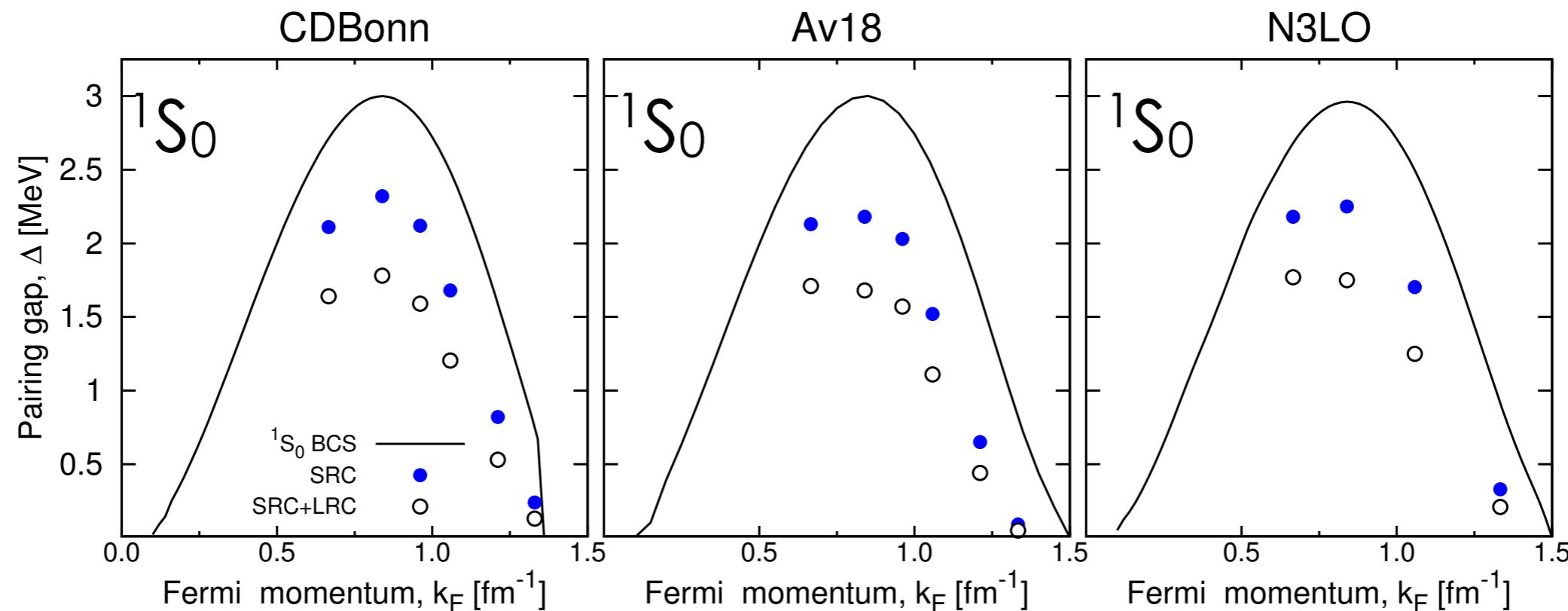
Effective Landau parameters

$$\langle 1\bar{1} | \mathcal{V} | 1\bar{1} \rangle = \frac{1}{4} \sum_{2,2'} \sum_{S,T} (-)^S (2S+1) \langle 12 | G_{ST}^{\text{ph}} | 1'2' \rangle_A \langle 2'\bar{1} | G_{ST}^{\text{ph}} | 2\bar{1}' \rangle_A \Lambda(22')$$

$$\Lambda_{ST}(q) = \frac{\Lambda_{ST}^0(q)}{1 - \Lambda_{ST}^0(q) \times F_{ST}}$$

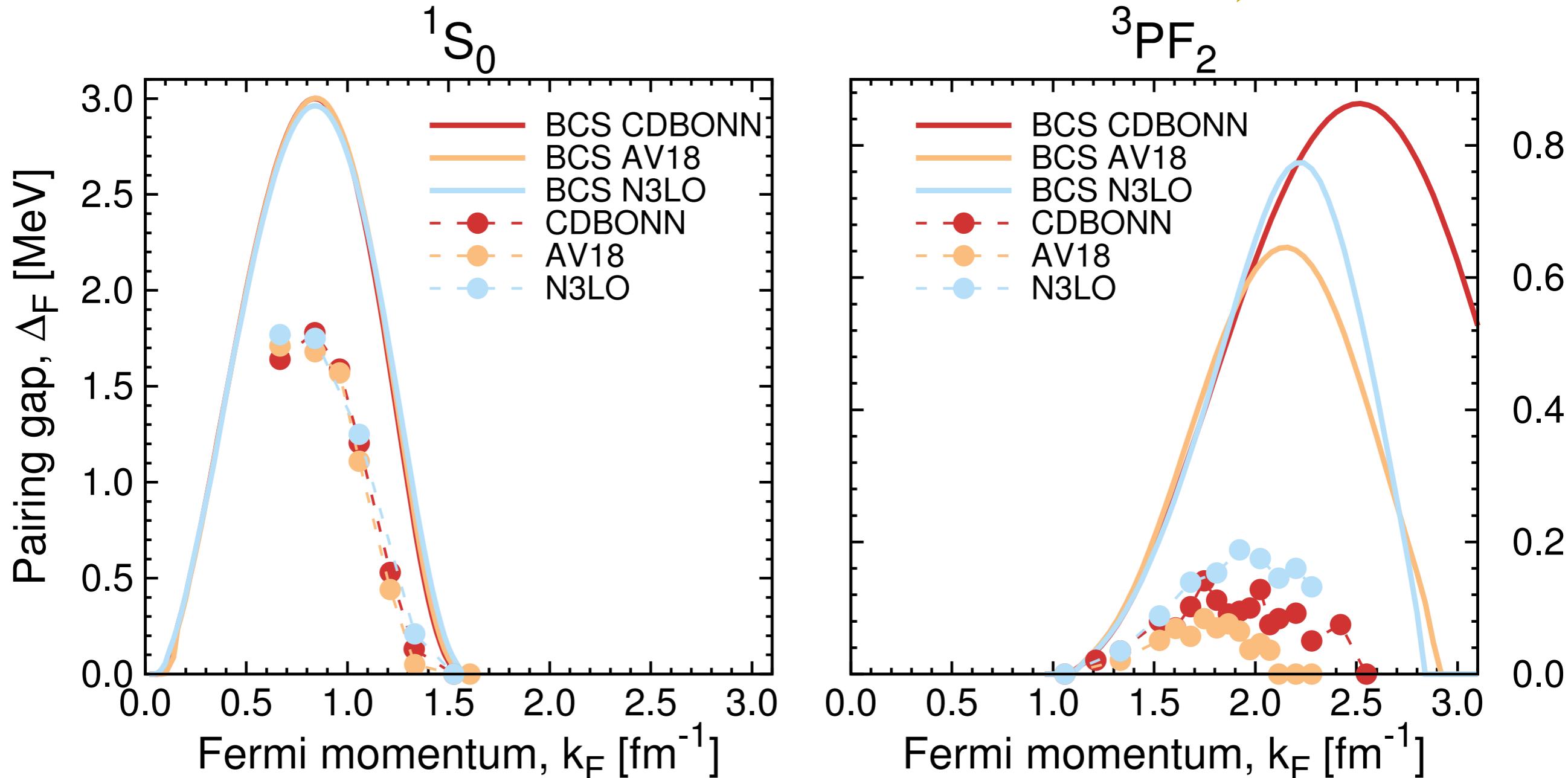
- Bare NN potential only is not the only possible interaction
- Diagram (a): nuclear interaction
- Diagram (b): in-medium interaction, **density** and spin fluctuations
- Diagram (c): included by **Landau parameters**

Beyond BCS 201: results 1S_0 Neutron matter

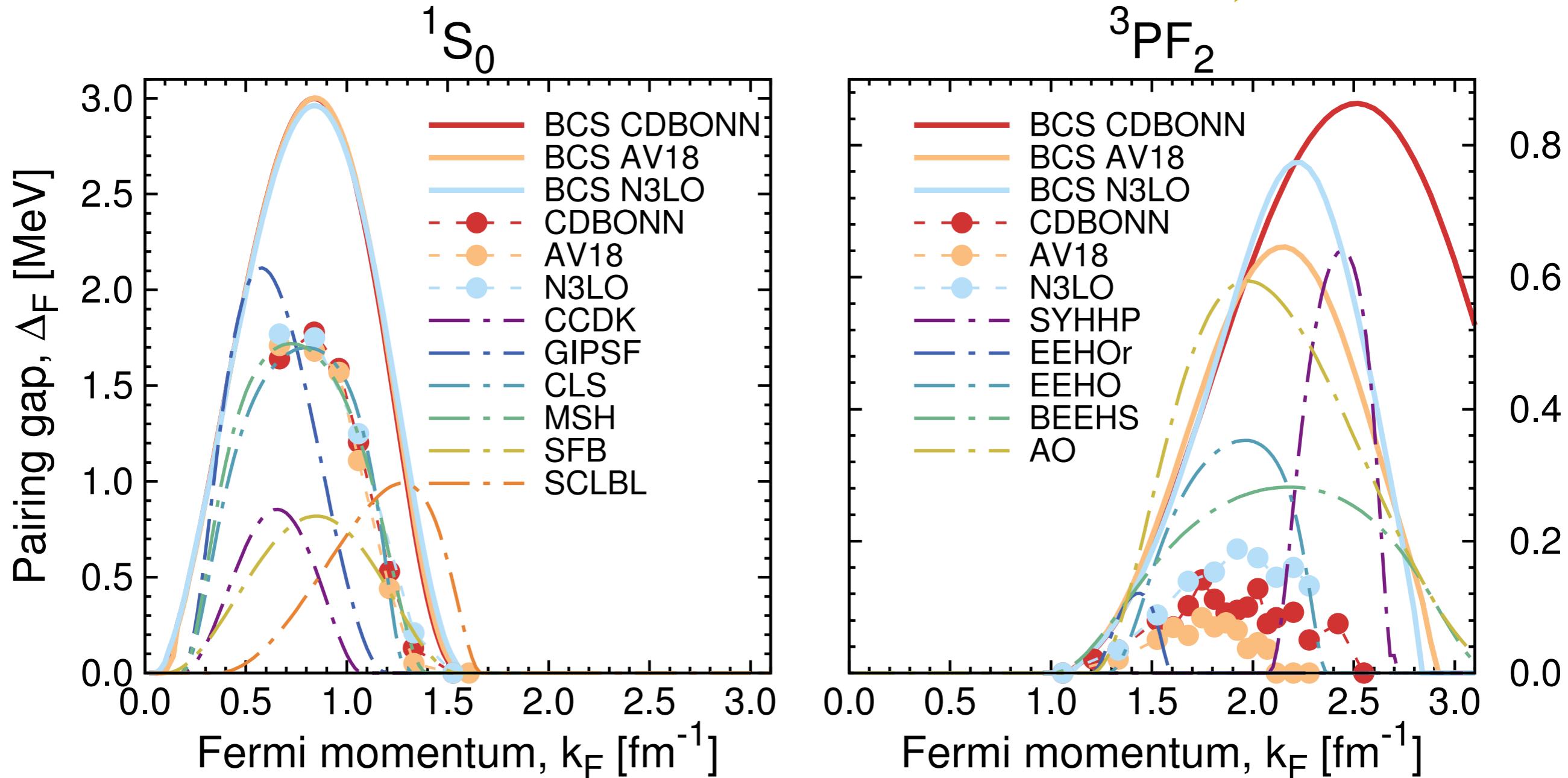


- LRC 1S_0 (3P_2) produces (anti-)screening

Beyond BCS 201: overview



- Effect is **robust**: independent of NN potential
- 3NF effect **not** included in SRC, BCS indicates **small**
- **Singlet** channel under control in **astrophysical** situations



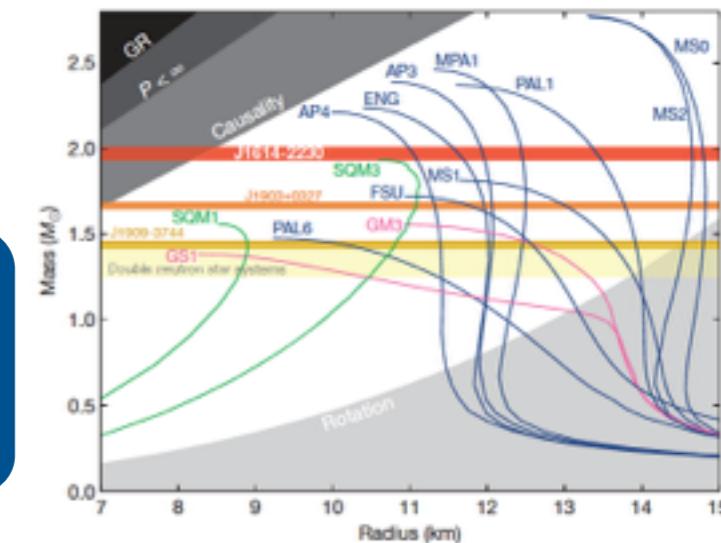
- Effect is **robust**: independent of NN potential
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Neutron star modeling

Input #1
EoS

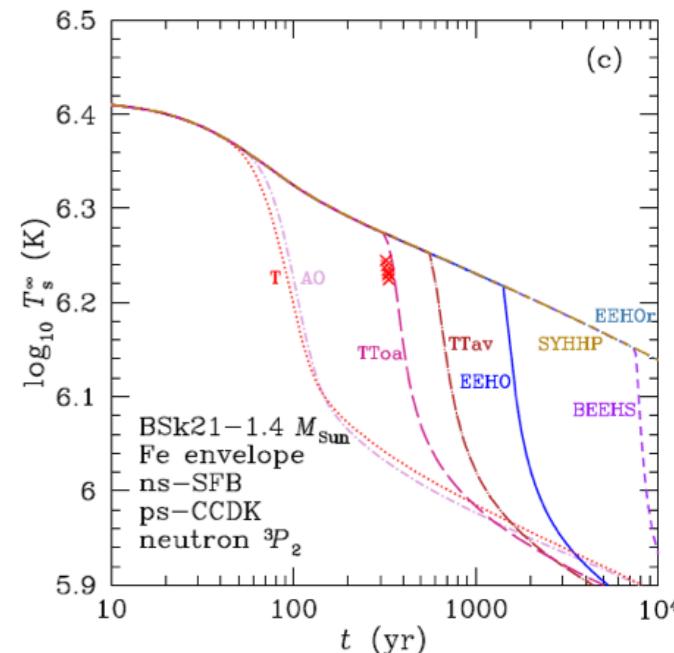
Observable #1
Mass-Radius relation

See Arianna's talk



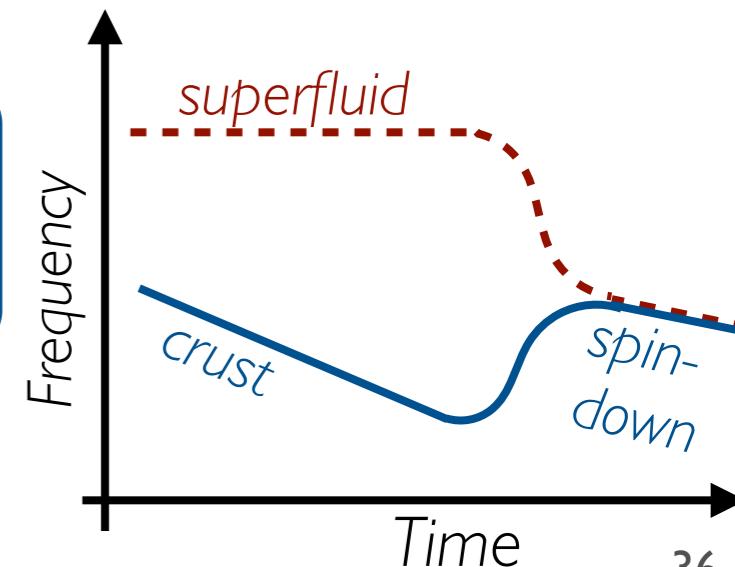
Input #2
Pairing gap

Observable #2
Cooling curve

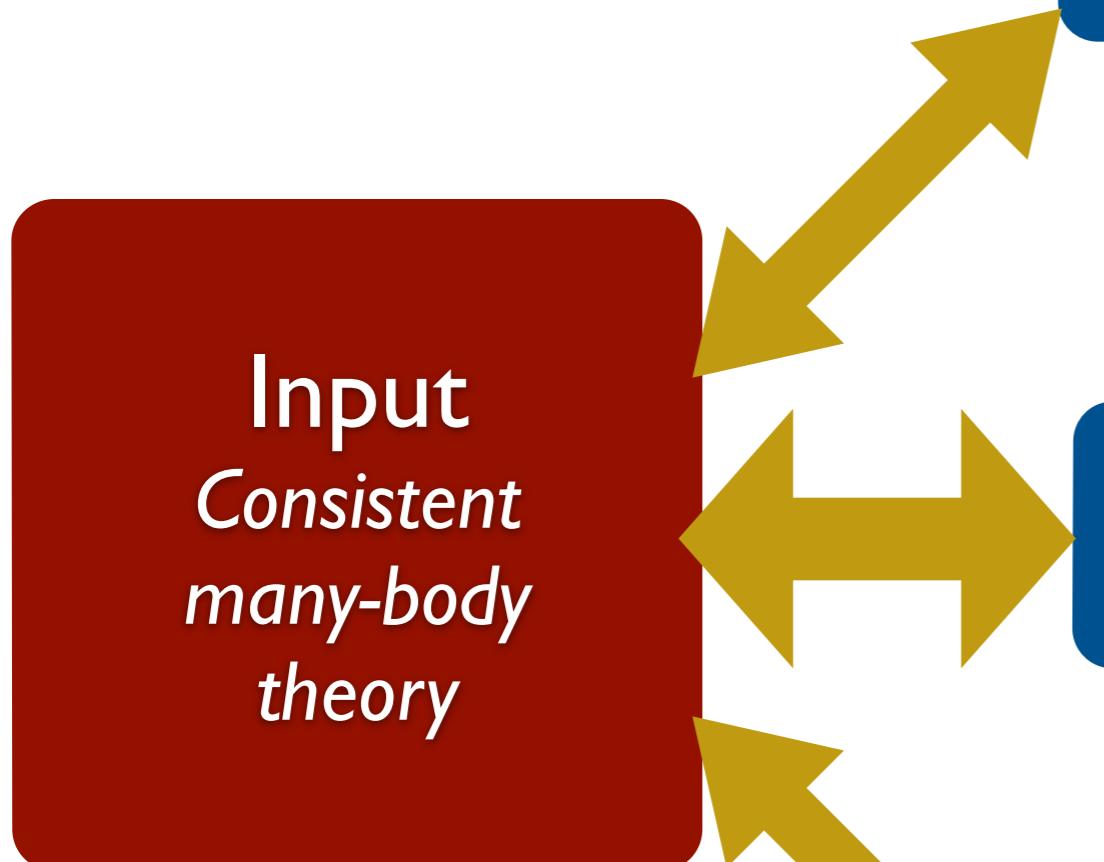


Input #3
Crust-core

Observable #3
Glitching

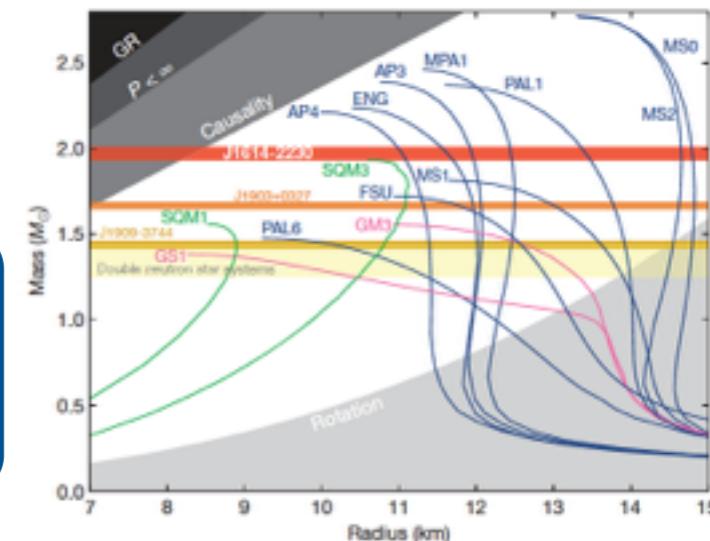


Neutron star modeling

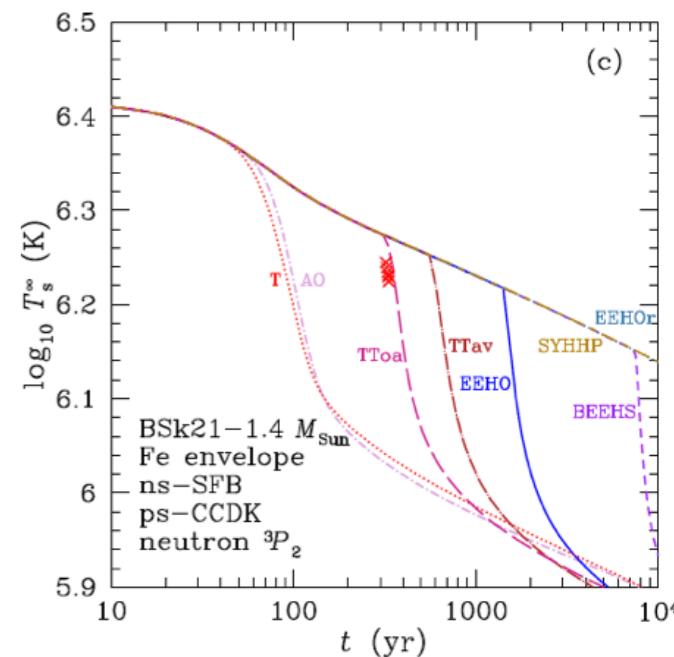


Observable #1
Mass-Radius relation

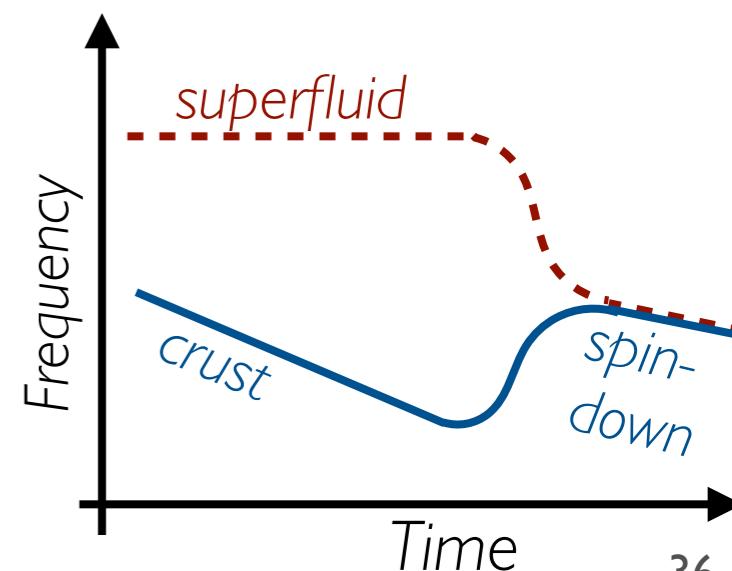
See Arianna's talk



Observable #2
Cooling curve



Observable #3
Glitching



Collaborators



+A. Carbone



TECHNISCHE
UNIVERSITÄT
DARMSTADT



Alexander von Humboldt
Stiftung/Foundation

+D. Ding, W. H. Dickhoff, H. Dussan



+A. Polls



+C. Barbieri



+V. Somà



Conclusions

- Ab initio nuclear **theory** to treat **correlations**
- **Talk** to us if you need **quantitative** answers!
- **Different** NN forces give robust predictions
- Approximations introduced **meaningfully**
- Challenges ahead:
 - **2 body** propagators and momentum distributions
 - **Pairing** in isospin **asymmetric** matter
 - **Consistent** treatment of **cooling, glitch & EoS**