Stylized features of nuclear momentum distributions and the quest for nuclear short-range correlations through the mass table

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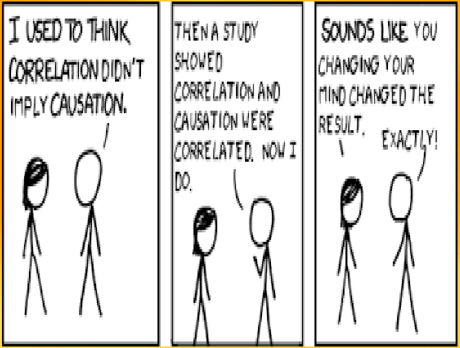
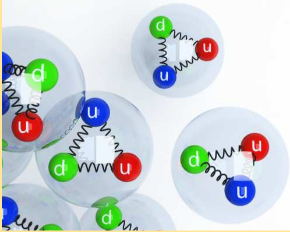
Talking about nuclear correlations



- Whole is different from the sum of the "parts"
- "Parts" can be effective degrees of freedom
- In nuclei: "Parts" are quasi-nucleons moving in a mean-field potential *(scheme dependent)*

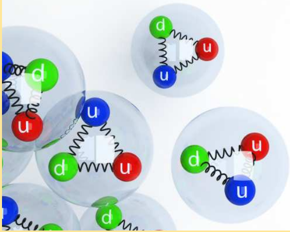
- Momentum correlations: $P^{(2)}\left(\vec{p}_1, \vec{p}_2\right) \neq P^{(1)}\left(\vec{p}_1\right) P^{(1)}\left(\vec{p}_2\right)$
- Spatial correlations: $P^{(2)}\left(\vec{r}_1, \vec{r}_2\right) \neq P^{(1)}\left(\vec{r}_1\right) P^{(1)}\left(\vec{r}_2\right)$
  1. short-range: $P^{(2)}\left(\vec{r}_1, \vec{r}_2\right) \neq 0$ for $\left|\vec{r}_1 - \vec{r}_2\right| \approx R_N$ (nucleon radius)
  2. long-range: $P^{(2)}\left(\vec{r}_1, \vec{r}_2\right) \neq 0$ for $\left|\vec{r}_1 - \vec{r}_2\right| \approx R_A$ (nuclear radius)

Nuclei in all their facets: IPM, SRC, LRC



Independent Particle Model (IPM)

- Nucleons have an identity: $\alpha_i(n_i, l_i, j_i, m_i, t_i)$ and $\psi_{\alpha_i}(\vec{r})$
- Average quantities: $\langle T_p \rangle$, $\langle U_{pot} \rangle$, $\langle \rho \rangle$, . . .



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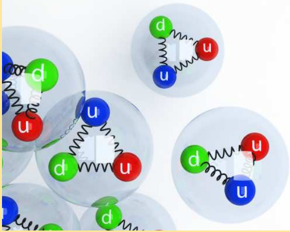

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Long Range Correlations (LRC)

- Nucleons loose their identity
- Spatio-temporal fluctuations: $\Delta T_p, \Delta U_{pot}, \Delta \rho, \ldots$
- “Most” nucleons get involved ($\sim R_A$)
- Energy scale $\Delta E \approx$10 MeV
- Experimentally observed and theoretically understood [giant resonances in $\gamma^{(*)}(A, X)$]

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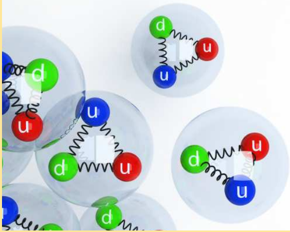

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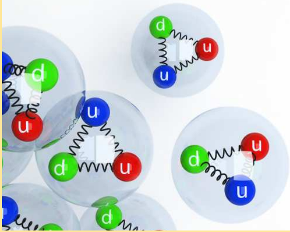

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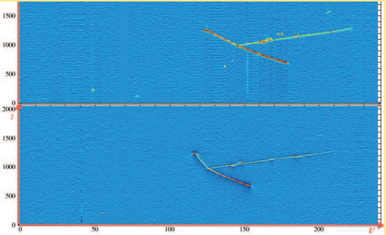
- Average quantities: $\langle T_p \rangle$, $\langle U_{pot} \rangle$, $\langle \rho \rangle$, ...

Short Range Correlations (SRC)

- Nucleons loose their identity
- Spatio-temporal fluctuations: $\Delta T_p, \Delta U_{pot}, \Delta\rho, \ldots$
- "Few" nucleons get involved ($\sim R_N$)
- Energy scale $\Delta E \approx$100 MeV
- Experimentally observed and theoretically understood [2N knockout in $A(e, e'X)$]

Research goals: comprehensive picture of SRC



"hammer events" in $(\nu_\mu, \mu^- pp)$ (arXiv:1405.4261)

- Develop an approximate flexible method for computing nuclear momentum distributions
- Study the mass and isospin dependence of SRC
- Provide a unified framework to establish connections with measurable quantities that are sensitive to SRC
  1. Inclusive $A(e, e')$ at $x_B \gtrsim 1.5$
  2. Magnitude of the EMC effect
  3. Two-nucleon knockout: $A(e, e'pN)$, $A(\nu_\mu, \mu^- pp)$
- Learn about SRC physics (nuclear structure AND reactions) in a unified framework

Nuclear correlation operators (I)

- Shift complexity from wave functions to operators

$$\mid \Psi\rangle = \frac{1}{\sqrt{\mathcal{N}}}\widehat{\mathcal{G}} \mid \Phi\rangle \qquad \text{with,} \qquad \mathcal{N} \equiv \langle \Phi \mid \widehat{\mathcal{G}}^{\dagger}\widehat{\mathcal{G}} \mid \Phi\rangle$$

$\mid \Phi\rangle$ is an IPM single Slater determinant

- Nuclear correlation operator $\widehat{\mathcal{G}}$

$$\widehat{\mathcal{G}} \approx \widehat{\mathcal{S}} \left(\prod_{i<j=1}^{A} \left[1 + \hat{l}(i,j) \right] \right) ,$$

- Major source of correlations: central (Jastrow), tensor and spin-isospin

$$\hat{l}(i,j) = -g_c(r_{ij}) + f_{\sigma\tau}(r_{ij})\vec{\sigma}_i \cdot \vec{\sigma}_j \vec{\tau}_i \cdot \vec{\tau}_j + f_{t\tau}(r_{ij})\widehat{S}_{ij}\vec{\tau}_i \cdot \vec{\tau}_j \ .$$

Nuclear correlation operators (II)

- Expectation values between correlated states Ψ can be turned into expectation values between uncorrelated states Φ

$$\langle \Psi \mid \widehat{\Omega} \mid \Psi \rangle = \frac{1}{\mathcal{N}} \langle \Phi \mid \widehat{\Omega}^{\text{eff}} \mid \Phi \rangle$$

- "Conservation Law of Misery": multi-body operators

$$\widehat{\Omega}^{\text{eff}} = \widehat{\mathcal{G}}^{\dagger} \, \widehat{\Omega} \, \widehat{\mathcal{G}} = \left(\sum_{i<j=1}^{A} \left[1 - \hat{l}(i,j) \right] \right)^{\dagger} \widehat{\Omega} \left(\sum_{k<l=1}^{A} \left[1 - \hat{l}(k,l) \right] \right)$$

  $\widehat{\Omega}^{\text{eff}}$ is an A-body operator
- Truncation procedure for short-distance phenomena:
  K. Wilson's OPE: $\Psi^{\dagger}(\vec{R} - \frac{\vec{r}}{2})\Psi(\vec{R} + \frac{\vec{r}}{2}) = \sum_n c_n(\vec{r}) O_n(\vec{R}) \quad (\mid \vec{r} \mid \approx 0)$
- **Low-order correlation operator approximation (LCA)**
- LCA: N-body operators receive SRC-induced $(N+1)$-body corrections

Including SRC: LCA method for one-body operators

- LCA effective operator corresponding with a one-body operator $\sum_{i=1}^{A} \widehat{\Omega}^{[1]}(i)$ (corrects for SRC)

$$\widehat{\Omega}^{\text{eff}} \approx \widehat{\Omega}^{\text{LCA}} = \sum_{i=1}^{A} \widehat{\Omega}^{[1]}(i) + \sum_{i<j=1}^{A} \left\{ \widehat{\Omega}^{[1],\text{l}}(i,j) + \left[\widehat{\Omega}^{[1],\text{l}}(i,j) \right]^{\dagger} + \widehat{\Omega}^{[1],\text{q}}(i,j) \right\}$$

- Two types of SRC corrections (two-body)
  1. linear in the correlation operator:

$$\widehat{\Omega}^{[1],\text{l}}(i,j) = \left[\Omega^{[1]}(i) + \Omega^{[1]}(j) \right] \hat{l}(i,j)$$

  2. quadratic in the correlation operator:

$$\widehat{\Omega}^{[1],\text{q}}(i,j) = \hat{l}^{\dagger}(i,j) \left[\widehat{\Omega}^{[1]}(i) + \widehat{\Omega}^{[1]}(j) \right] \hat{l}(i,j).$$

Norm $\mathcal{N} \equiv \langle \Phi \mid \widehat{\mathcal{G}}^{\dagger}\widehat{\mathcal{G}} \mid \Phi \rangle$: aggregated SRC effect

- LCA expansion of the norm $\mathcal{N}$

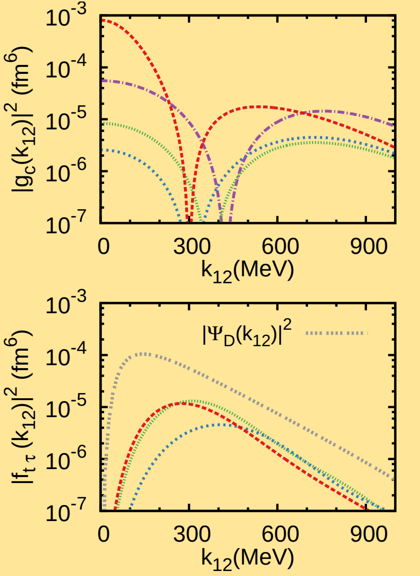
$$\mathcal{N} = 1 + \frac{2}{A} \sum_{\alpha<\beta} {}_{\text{nas}}\langle \alpha\beta \mid \hat{l}^{\dagger}(1,2) + \hat{l}^{\dagger}(1,2)\hat{l}(1,2) + \hat{l}(1,2) \mid \alpha\beta \rangle_{\text{nas}}.$$

  1. $\mid \alpha\beta\rangle_{\text{nas}}$: normalized and anti-symmetrized two-nucleon IPM-state
  2. $\sum_{\alpha<\beta}$ extends over all IPM states $\mid \alpha \rangle \equiv \mid n_\alpha l_\alpha j_\alpha m_{j_\alpha} t_\alpha \rangle$,
- $(\mathcal{N} - 1)$: measure for aggregated effect of SRC in the ground state
- Aggregated quantitative effect of SRC in A relative to ^{2}H

$$R_2(A/^2\text{H}) = \frac{\mathcal{N}(A) - 1}{\mathcal{N}(^2\text{H}) - 1} = \frac{\text{measure for SRC effect in } A}{\text{measure for SRC effect in } ^2\text{H}}.$$

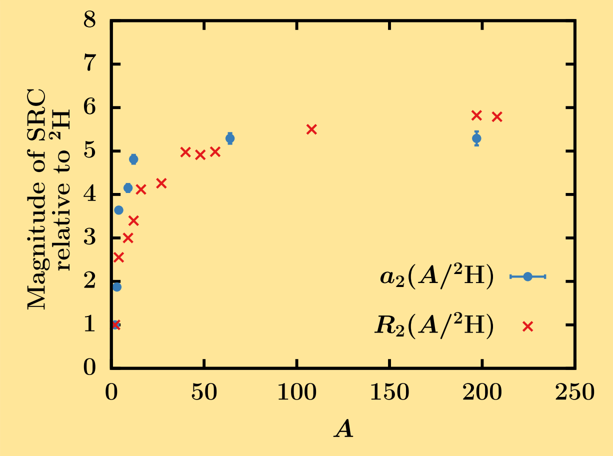
- Input to the calculations for $R_2(A/^2\text{H})$:
  1. HO IPM states with $\hbar\omega = 45A^{-1/3} - 25A^{-2/3}$
  2. A-independent universal correlation functions $[g_c(r), f_{t\tau}(r), f_{\sigma\tau}(r)]$

Central and tensor correlation function



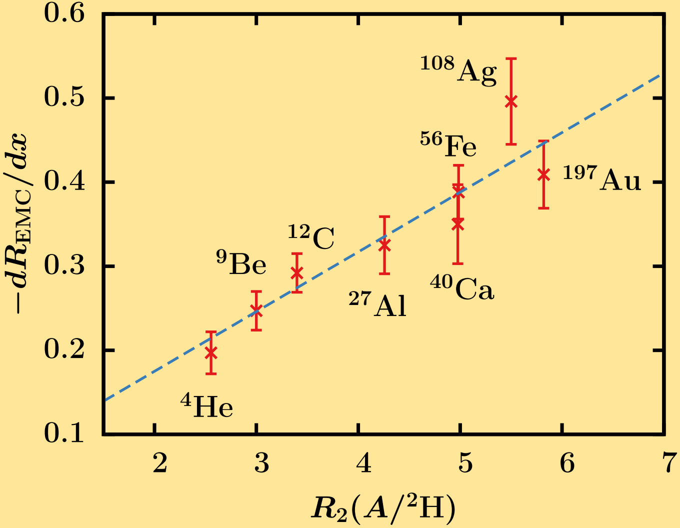
- the $g_C(k_{12})$ looks like the correlation function of a monoatomic classical liquid (reflects finite-size effects)
- the $g_c(k_{12})$ are ill constrained
- $|f_{t\tau}(k_{12})|^2$ is well constrained! (*D*-state deuteron wave function)
- $|f_{t\tau}(k_{12})|^2 \sim |\Psi_D(k_{12})|^2$
- very high relative pair momenta: central correlations
- moderate relative pair momenta: tensor correlations

$a_2(A/^2\text{H})$ from $A(e,e')$ at $x_B \gtrsim 1.5$ and $R_2(A/^2\text{H})$



1. $A \lesssim 40$: strong mass dependence in SRC effect
2. $A > 40$: soft mass dependence
3. SRC effect saturates for A large *(for large A aggregated SRC effect per nucleon is about 5× larger than in ^2H)*

Magnitude of EMC effect versus $R_2(A/^2\mathrm{H})$



LCA can predict magnitude of EMC effect for any $A(N,Z) \geq 4$

Single-nucleon momentum distribution $n^{[1]}(p)$

- Probability to find a nucleon with momentum p

$$
\begin{aligned}
n^{[1]}(p) \quad = \quad & \int \frac{d^2\Omega_p}{(2\pi)^3} \int d^3\vec{r}_1 \; d^3\vec{r}_1{}' \; d^{3(A-1)}\{\vec{r}_{2-A}\} e^{-i\vec{p}\cdot(\vec{r}_1{}'-\vec{r}_1)} \\
& \times \Psi^*(\vec{r}_1, \vec{r}_{2-A}) \Psi(\vec{r}_1{}', \vec{r}_{2-A}).
\end{aligned}
$$

- Corresponding single-nucleon operator $\hat{n}_p$

$$
\hat{n}_p = \frac{1}{A} \sum_{i=1}^{A} \int \frac{d^2\Omega_p}{(2\pi)^3} e^{-i\vec{p}\cdot(\vec{r}_i{}'-\vec{r}_i)} = \sum_{i=1}^{A} \hat{n}_p^{[1]}(i).
$$

- Effective correlated operator $\hat{n}_p^{\mathrm{LCA}}$
  (SRC-induced corrections to IPM $\hat{n}_p$ are of two-body type)
- Normalization property $\int dp \; p^2 n^{[1]}(p) = 1$ can be preserved by evaluating $\mathcal{N}$ in LCA

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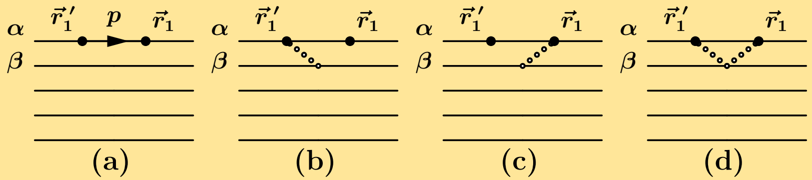
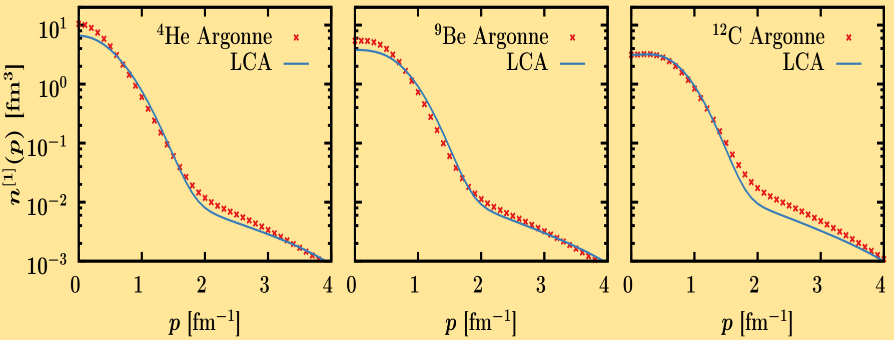
$$\times \Psi^*(\vec{r}_1, \vec{r}_{2-A})\Psi(\vec{r}_1{}', \vec{r}_{2-A}).$$



(a): IPM contribution

(b)-(d): SRC contributions

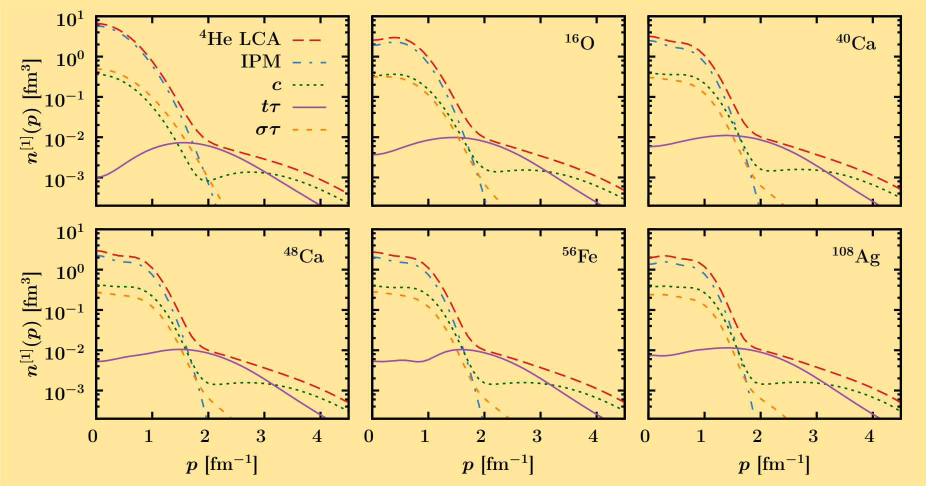
$n^{[1]}(p)$ for light nuclei: LCA (Ghent) vs QMC (Argonne)



QMC: PRC89(2014)024305 LCA: JPG42(2015)055104

1. $p \lesssim p_F = 1.25$ fm^{-1}: $n^{[1]}(p)$ is “Gaussian” (IPM PART)
2. $p \gtrsim p_F$: $n^{[1]}(p)$ has an “exponential” fat tail (CORRELATED PART)
3. fat tail in QMC and LCA are in reasonable agreement

Major source of correlated strength in $n^{[1]}(p)$?



1. $1.5 \lesssim p \lesssim 3$ fm^{-1} is dominated by tensor correlations
2. central correlations substantial at $p \gtrsim 3.5$ fm^{-1}

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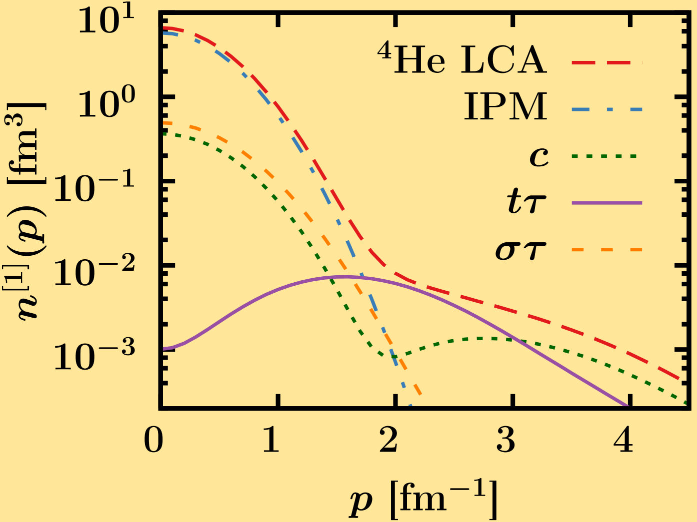

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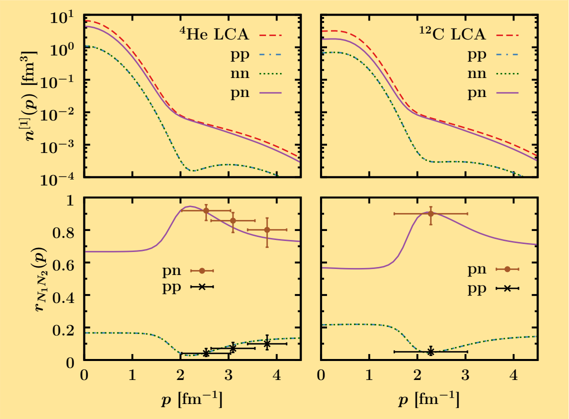
Isospin dependence of correlations: pp, nn and pn

$$n^{[1]}(p) \equiv n^{[1]}_{pp}(p) + n^{[1]}_{nn}(p) + n^{[1]}_{pn}(p) \qquad r_{N_1 N_2}(p) \equiv n^{[1]}_{N_1 N_2}(p)/n^{[1]}(p)$$



The fat tail is dominated by "pn" (momentum dependent)

- $r_{N_1 N_2}(p)$: relative contribution of $N_1 N_2$ pairs to $n^{[1]}(p)$ at p
- Naive IPM:
  $r_{pp} = \frac{Z(Z-1)}{A(A-1)}$,
  $r_{nn} = \frac{N(N-1)}{A(A-1)}$,
  $r_{pn} = \frac{2NZ}{A(A-1)}$.
- Data extracted from ^{4}He$(e, e'pp)$/^{4}He$(e, e'pn)$ (PRL 113, 022501) and $\frac{^{12}\text{C}(p,ppn)}{^{12}\text{C}(p,pp)}$ (Science 320, 1476) assuming that $r_{pp} \approx r_{nn}$

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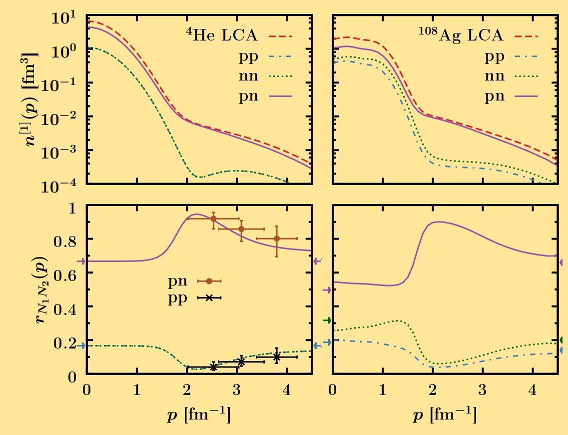

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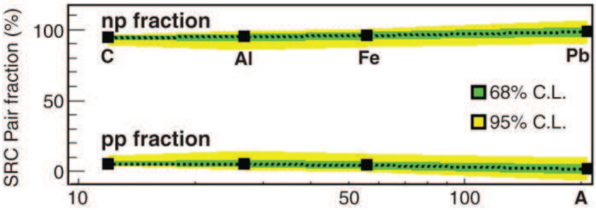
Momentum sharing in imbalanced Fermi systems

O. Hen,[1*] M. Sargsian,[2] L. B. Weinstein,[3] E. Piasetzky,[1] H. Hakobyan,[4,5] D. W. Higinbotham,[6] M



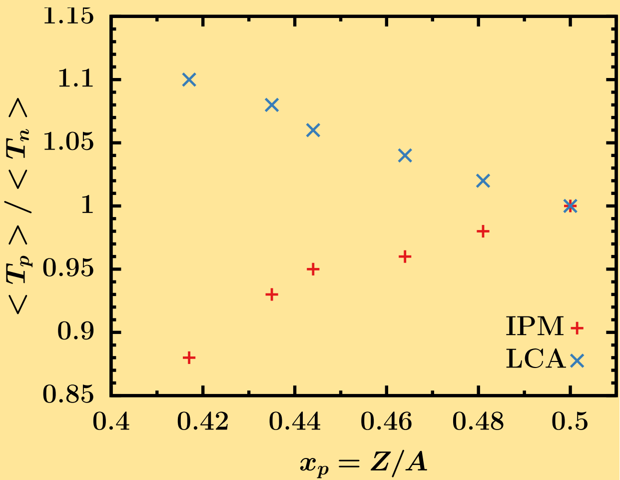
LCA predicts that ≈90% of correlated pairs is “pn”, and ≈5% is”pp” (A independent)

Average kinetic energy per nucleon $\langle T_N \rangle$

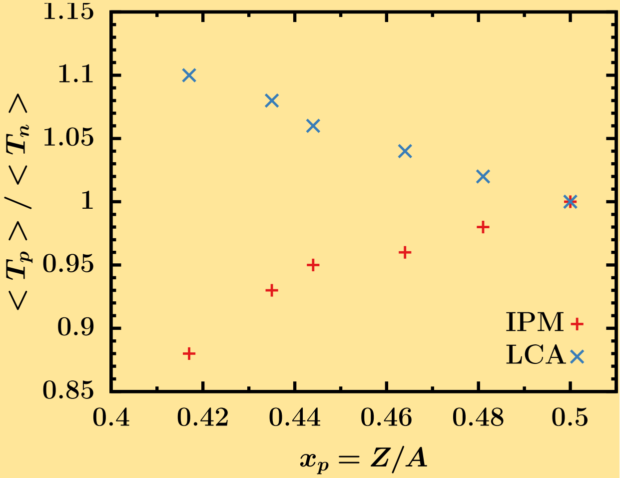
A	$x_p = \frac{Z}{A}$	$\langle T_N \rangle$ (MeV)						$\langle T_p \rangle / \langle T_n \rangle$	
		IPM (p)	IPM (n)	LCA (p)	LCA(n)	Perug	UCOM	IPM	LCA
^{2}H	0.500	14.95	14.93	20.95	20.91			1.00	1.00
^{4}He	0.500	13.80	13.78	25.28	25.23		19.63	1.00	1.00
^{9}Be	0.444	15.81	16.58	28.91	27.33			0.95	1.06
^{12}C	0.500	16.08	16.06	28.96	28.92	32.4	22.38	1.00	1.00
^{16}O	0.500	15.61	15.59	29.48	29.43	30.9	23.81	1.00	1.00
^{27}Al	0.481	16.61	16.92	30.93	30.26		25.12	0.98	1.02
^{40}Ca	0.500	16.44	16.42	31.23	31.18	33.8	27.72	1.00	1.00
^{48}Ca	0.417	15.64	17.84	33.04	30.06		27.05	0.88	1.10
^{56}Fe	0.464	16.71	17.45	32.33	31.13	32.7		0.96	1.04
^{108}Ag	0.435	16.48	17.81	33.55	31.16			0.93	1.08

1. SRC substantially increase $\langle T_N \rangle$ (factor of about 2)
2. after including SRC: minority component has largest $\langle T_N \rangle$

Predictions for $\langle T_p \rangle / \langle T_n \rangle$ ratio



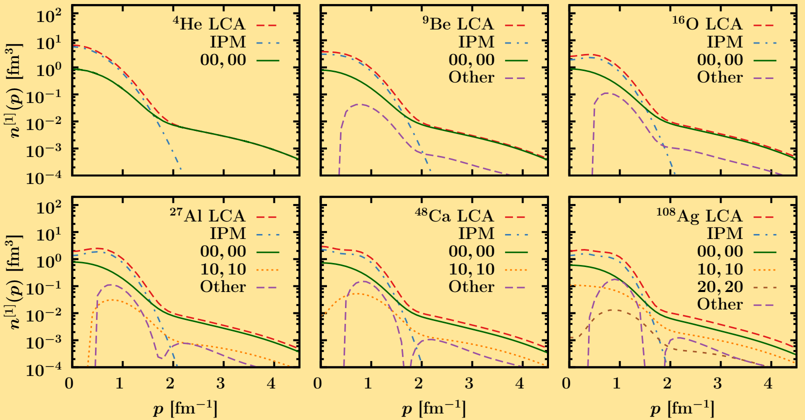
Predictions for $\langle T_p \rangle / \langle T_n \rangle$ ratio



SRC turn the IPM predictions upside down

Quantum numbers of SRC-susceptible IPM pairs?

$n^{[1],\mathrm{corr}}$ stems from correlation operators acting on IPM pairs.
What are relative quantum numbers (nl) of those IPM pairs?



$$\sum_{nl} \sum_{n'l'} n^{[1],\mathrm{corr}}_{nl,n'l'}(p) = n^{[1],\mathrm{corr}}(p)$$

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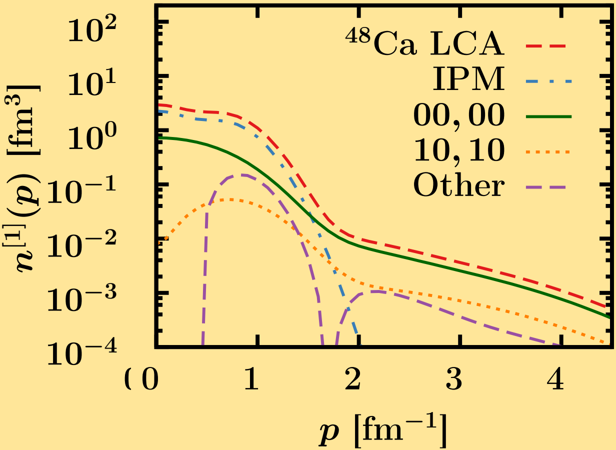
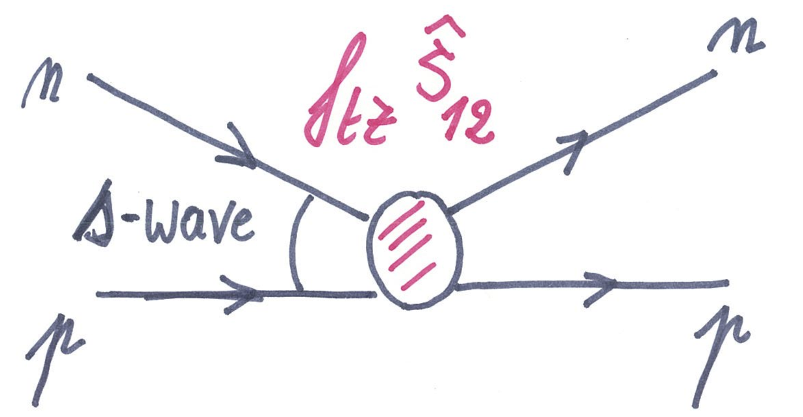


Major source of SRC: correlations acting on ($n = 0$ $l = 0$) IPM pairs

- Physical picture from LCA: for $1.5 \lesssim p \lesssim 3$ fm^{-1} the SRC are mainly due to tensor-induced scattering between IPM pn pairs in a relative *s*-state



Stylized features of nuclear SRC

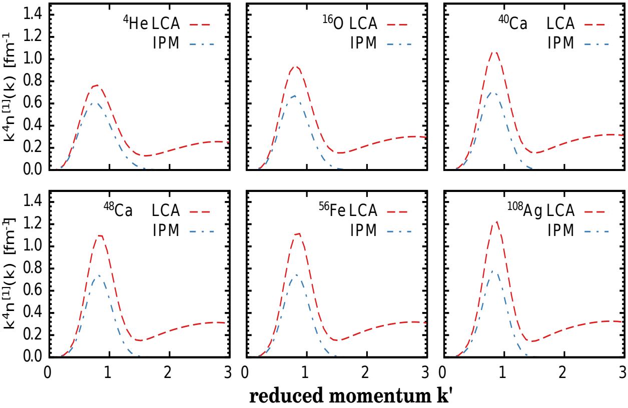
- Physical picture from LCA: for $1.5 \lesssim p \lesssim 3$ fm^{-1} the SRC are mainly due to tensor-induced scattering between IPM pn pairs in a relative *s*-state

- In tensor-dominated momentum range: nuclear Hamiltonian can be captured by the stylized Hamiltonian

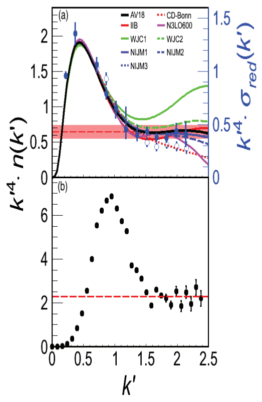
$$\begin{aligned}\widehat{\mathcal{H}} &\approx \sum_{\tau=p,n} \int d^3\vec{r}\psi_\tau^\dagger(\vec{r}) \left[-\frac{\hbar^2}{2m_N} \nabla_{\vec{r}}^2 + U_\tau(\vec{r}) \right] \psi_\tau(\vec{r}) \\ &+ \int d^3\vec{r} d^3\vec{R} \psi_p^\dagger \left(\vec{R} + \frac{\vec{r}}{2} \right) \psi_n^\dagger \left(\vec{R} - \frac{\vec{r}}{2} \right) \psi_n \left(\vec{R} \right) \psi_p \left(\vec{R} \right) \lambda_{t\tau} \left(\vec{r} \right)\end{aligned}$$

- Physics of a two-component and strongly correlated Fermi gas subject to an *s*-wave contact interaction is described by Tan (Ann. of Phys. 322 (2008) 2971)
- Landmark of a contact interaction: $n^{[1]}(p) \sim Cp^{-4}$

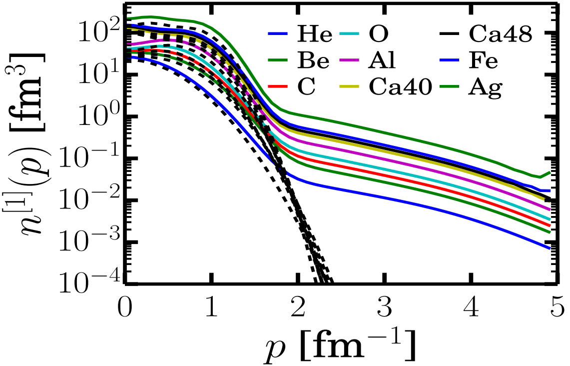
Approximate p^4 scaling of the $n^{[1]}(p)$



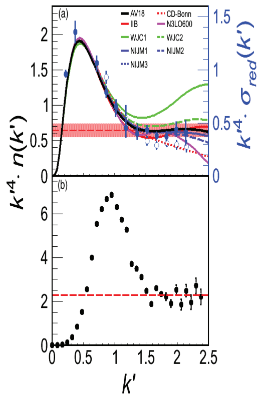
arXiv:1407.8175



- Momentum scale: $k' \equiv \frac{p}{p_F}$
- IPM is approximately Gaussian: stochastic collisions
- Fat tail is the landmark of strong correlations
- ***p*-dependence of tail is universal**



arXiv:1407.8175



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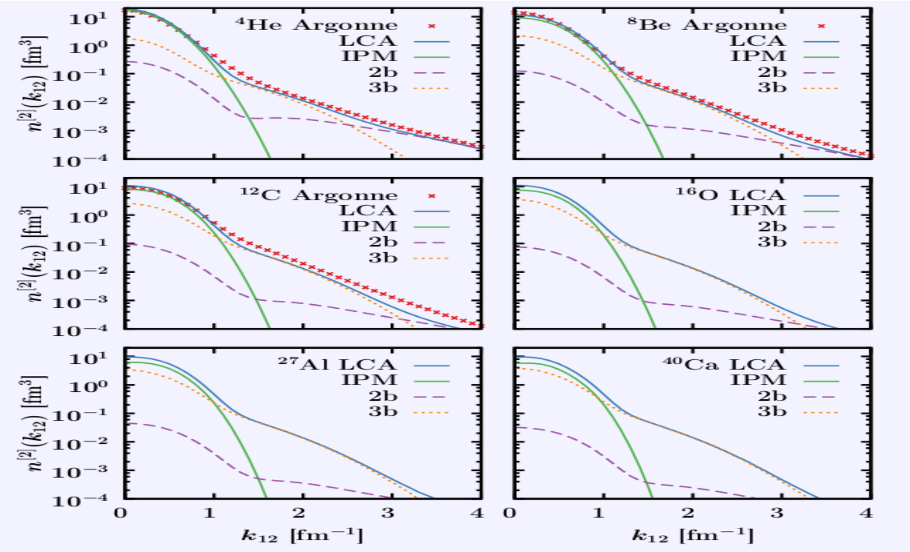
Two-nucleon momentum distribution (TNMD) $n^{[2]}\left(\vec{k}_{12},\vec{P}_{12}\right)$

- Belongs to the class of four-point correlation functions (two tagged nucleons)
- Corresponding two-nucleon operator $\widehat{n}_{k_{12}P_{12}}$
- In LCA: effective correlated operator $\widehat{n}^{LCA}_{k_{12}P_{12}}$ (SRC-induced corrections are two-body ("2b") and three-body ("3b") operators)
- Relative TNMD: distribution of the relative momentum of the tagged pair

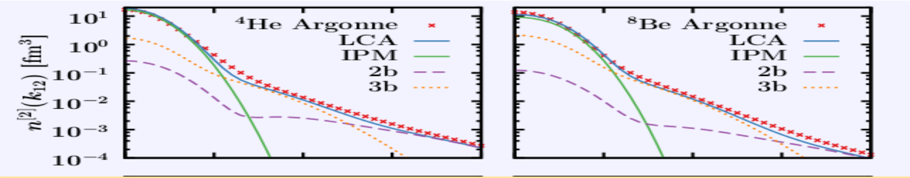
$$n^{[2]}(k_{12}) = \int d^3\vec{P}_{12} d^2\Omega_{k_{12}} n^{[2]}\left(\vec{k}_{12},\vec{P}_{12}\right)$$

- **No direct connection between $n^{[2]}\left(\vec{k}_{12},\vec{P}_{12}\right)$ and SRC dominated two-nucleon knockout cross sections**

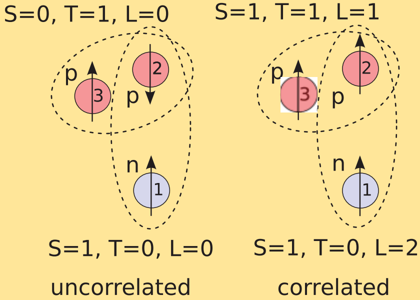
Relative TNMD: tail is dominated by “3-body” effects



Relative TNMD: tail is dominated by “3-body” effects

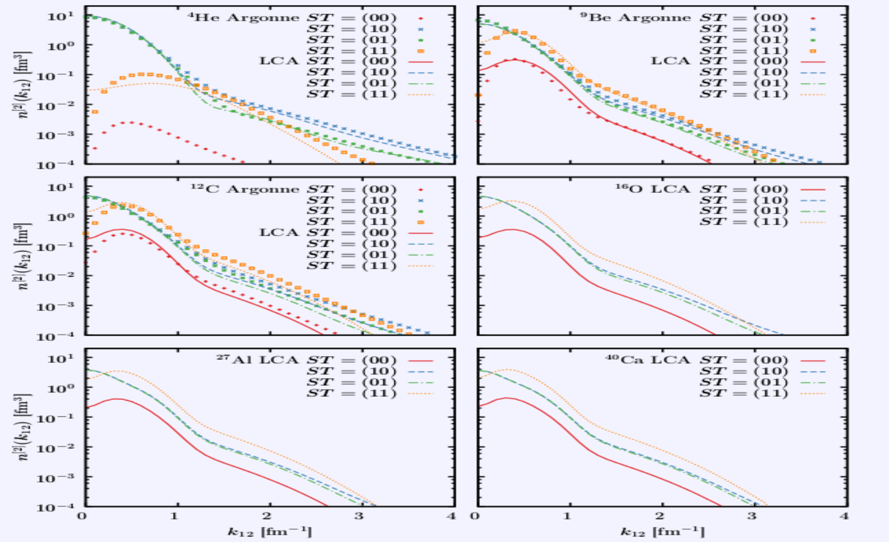


Correlations through the mediation of a third particle:

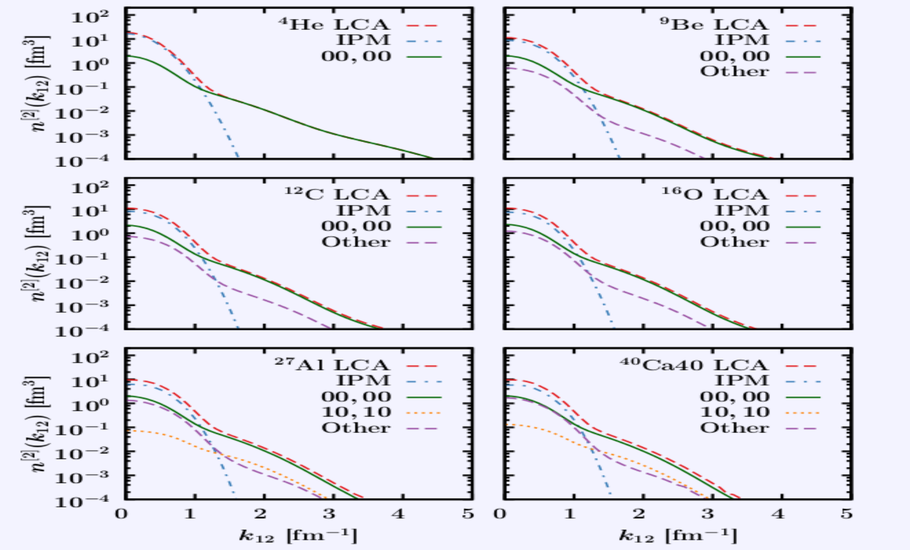


Feldmeier *et al.*, PRC 84 (2011), 054003

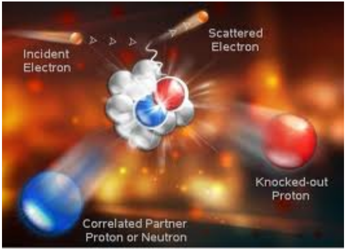
Relative TNMD: quantum numbers of tagged pairs $\neq$ quantum numbers of correlated pair

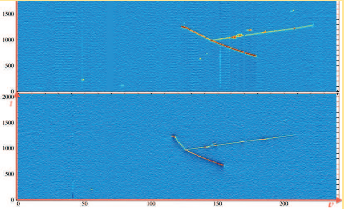


Correlated part of relative TNMD: dominated by *s*-wave scattering!



Exclusive two-nucleon knockout $A(e, e'NN)$

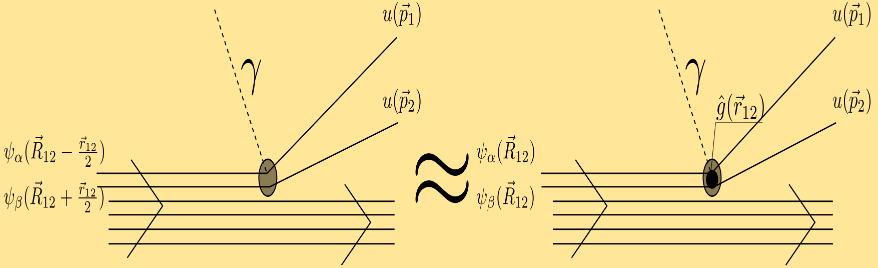




"hammer events" in $(\nu_\mu, \mu^- pp)$ (arXiv:1405.4261)

- The (virtual) photon-nucleon interaction is a one-body operator
- Two-nucleon knockout is the hallmark of SRC (one hits a nucleon and its correlated partner)
  1. $A(e, e'pN)$
  2. $A(\nu_\mu, \mu^- pp)$
  3. $A(p, pNN)$

The fact that SRC-prone IPM NN pairs are mostly in a close-proximity ($n_{12} = 0, l_{12} = 0$) state has important consequences for the EXCLUSIVE $A(e, e'NN)$ cross sections
[PLB 383,1 (1996) and PRC 89, 024603 (2014)]



ZRA: Zero range approximation

Exclusive $A(e, e'NN)$ along the LCA lines

The fact that SRC-prone IPM NN pairs are mostly in a close-proximity ($n_{12} = 0, l_{12} = 0$) state has important consequences for the EXCLUSIVE $A(e, e'NN)$ cross sections
[PLB 383,1 (1996) and PRC 89, 024603 (2014)]

1. **$A(e, e'NN)$ cross section factorizes according to**

$$\frac{d^8\sigma}{d\epsilon' d\Omega_{\epsilon'} d\Omega_1 d\Omega_2 dT_{p_2}}(e, e'NN) = K\sigma_{eNN}(k_+, k_-, q) F^{(D)}_{h_1,h_2}(P)$$

   $F^{(D)}_{h_1,h_2}(P)$: FSI corrected conditional probability to find a dinucleon with c.m. momentum P in a relative ($n_{12} = 0, l_{12} = 0$) state

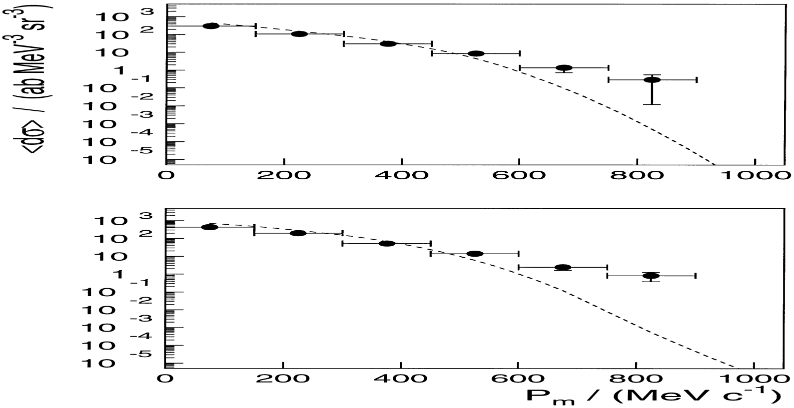
2. **A dependence of the $A(e, e'pp)$ cross sections is soft**
   (much softer than predicted by naive $Z(Z-1)$ counting)

$$\frac{A(e, e'pp)}{^{12}\mathrm{C}(e, e'pp)} \approx \frac{N_{pp}(A)}{N_{pp}(^{12}\mathrm{C})} \times \left(\frac{T_A(e, e'p)}{T_{^{12}\mathrm{C}}(e, e'p)}\right)^{1-2}$$

3. **C.m. width of SRC susceptible pairs is “large” (in p-space)**

Factorization of the $A(e, e'pp)$ cross sections

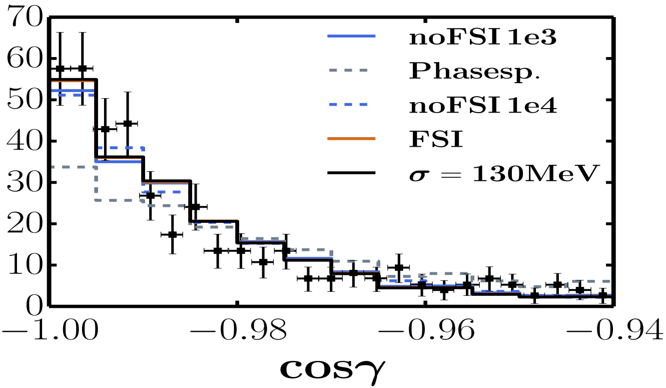
$^{12}C(e, e'pp)$ @ MAMI (Mainz) (Physics Letters B **421** (1998) 71.)



For P $\lesssim$ 0.5 GeV c.m. motion of correlated pairs in ^{12}C is mean-field like $\left(\exp \frac{-P^2}{2\sigma^2_{c.m}}\right)$! Data prove the proposed factorization in terms of $F^{(D)}_{h_1,h_2}(P)$.

$A(e, e'NN)$: Effect of the final-state interactions?

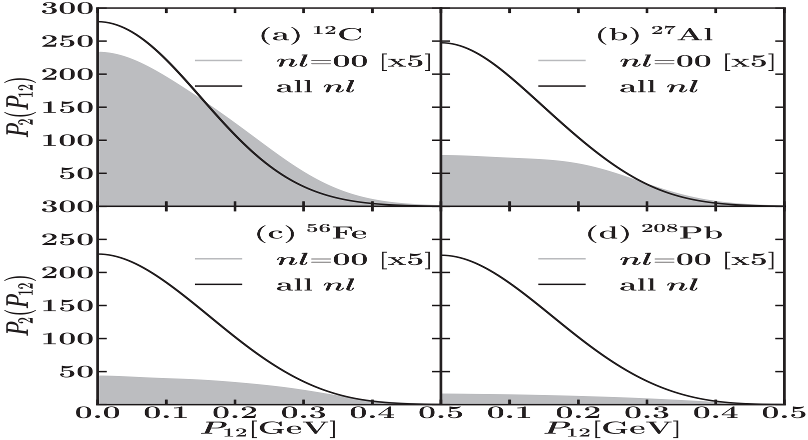
Opening-angle distribution of ^{4}He$(e, e'pp)$



1. **FSI (eikonal model) reduces the cross sections**
2. **FSI marginally affects the angular distributions**
   (FSI preserves factorization properties)

C.m. motion of correlated pp pairs

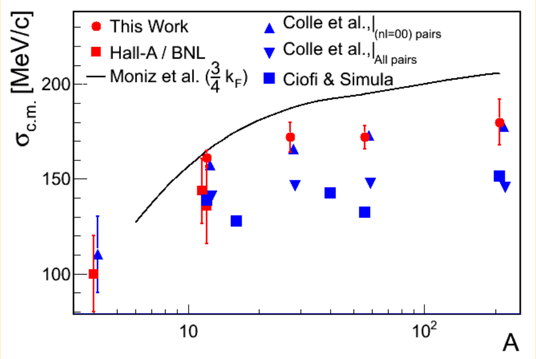
PHYSICAL REVIEW C **89**, 024603 (2014)



Width of c.m. distribution is a lever to discriminate between SRC-prone IPM pairs and the other IPM pairs

C.m. motion of correlated pp pairs

DATA IS PRELIMINARY! (COURTESY OF O. HEN AND E. PIASETZKY)



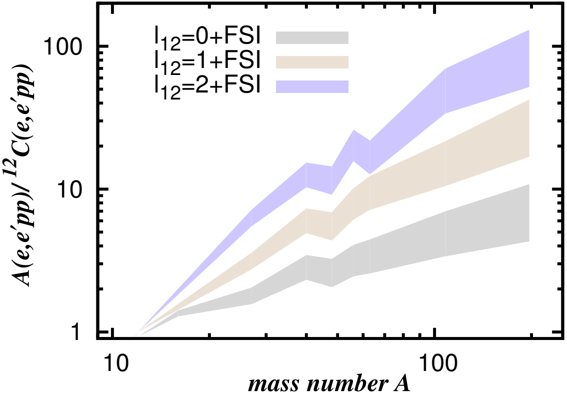
- Analysis of exclusive $A(e, e'pp)$ for ^{12}C, ^{27}Al, ^{56}Fe, ^{208}Pb by Data Mining Collaboration at Jefferson Lab
- Distribution of events against P is fairly Gaussian
- $\sigma_{c.m.}$: Gaussian widths from a fit to measured c.m. distributions

Mass dependence of the $A(e, e'pp)$ cross sections

PREDICTION: A dependence of $A(e, e'pp)$ c.s. is soft

(much softer than predicted by naive $Z(Z-1)$ counting)

$$\frac{A(e,e'pp)}{^{12}\mathrm{C}(e,e'pp)} \approx \frac{N_{pp}(A)}{N_{pp}\left(^{12}\mathrm{C}\right)} \times \left(\frac{T_A(e,e'p)}{T_{^{12}\mathrm{C}}(e,e'p)} \right)^{1-2}$$



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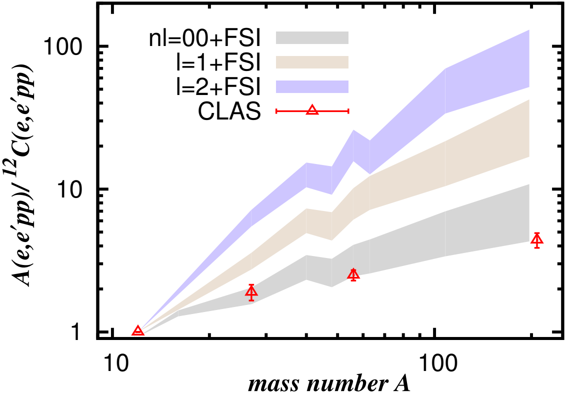
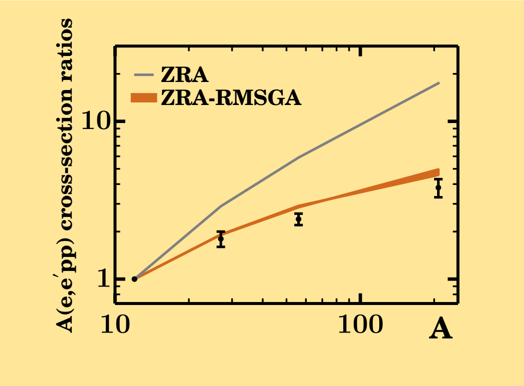
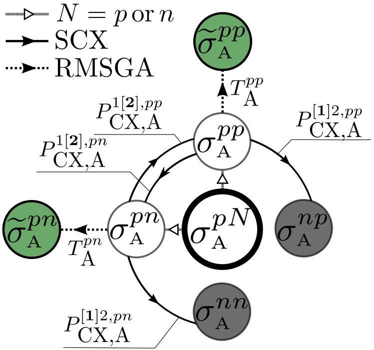


DATA COMPATIBLE WITH ABSORPTION ON $(n_{12} = 0, l_{12} = 0)$ PAIRS

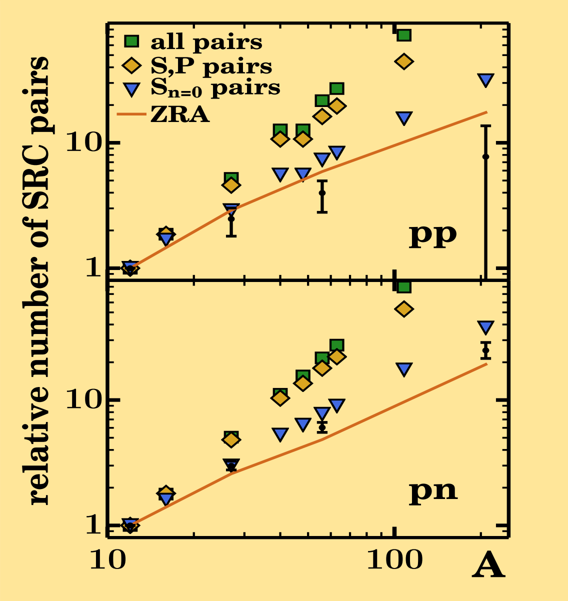
arXiv:1503.06050





1. Effect of final-state interactions in the eikonal approximation
2. Effect of single-charge exchange (SCX) included
3. Pb to C cross-section ratio for $(e, e'pp) \approx 4$

A dependence of number of pp and pn SRC pairs



- Analysis of $A(e, e'pp)$ and $A(e, e'p)$ (A=^{12}C, ^{27}Al, ^{56}Fe, ^{208}Pb) in "SRC" kinematics (Data Mining Collaboration @JLAB)
- FSI corrections applied to the data
- Reaction-model calculations in the large phase space: importance sampling
- Relative number of SRC pp-pairs and pn-pairs

CONCLUSIONS (I)

Stylized features of nuclear SRC: The mass and isospin dependence of the magnitude of the 2N and 3N correlations can be captured by some general principles

- LCA: efficient and realistic way of computing the SRC contributions to nuclear momentum distributions (NMD)
  1. Magnitude of EMC effect and $A(e,e')/D(e,e')$ scaling factor ($x_B \gtrsim 1.5$) can be predicted in LCA
  2. Light nuclei: LCA predictions for tails are in line with those of QMC
  3. LCA predictions for $\langle T_N \rangle$ and radii are "realistic" (consistency checks)
  4. Natural explanation for the universal behavior of the NMD tails
- Number of SRC-prone pairs in a nucleus $A(N,Z)$ is proportional with the number of pairs in a relative $(n_{12}=0, l_{12}=0)$ state

CONCLUSIONS (II)

- Insights from study of SRC contribution to NMD has implications for exclusive $A(e, e'NN)$:
  1. Scaling behavior of cross section ($\sim F(P)$) (CONFIRMED!)
  2. Very soft mass dependence of cross section (CONFIRMED!)
  3. Peculiar c.m. width of the SRC-susceptible pairs (CONFIRMED!)
- Aggregated effect of SRC: A independent correlation operators acting on close-proximity pairs in a nodeless relative S state
- Generally applicable techniques for quantifying SRC: two-body effects in neutrino reactions, role of SRC in exotic forms of hadronic matter, . . .
- SRC induced spatio-temporal fluctuations are measurable and significant (scales are set)


THANK YOU!

Selected publications

- J. Ryckebusch, M. Vanhalst, W. Cosyn
  "Stylized features of single-nucleon momentum distributions"
  arXiv:1405.3814 and Journal of Physics G **42** (2015) 055104.
- C. Colle, O. Hen, W. Cosyn, I. Korover, E. Piasetzky, J. Ryckebusch, L.B. Weinstein
  "Extracting the Mass Dependence and Quantum Numbers of Short-Range Correlated Pairs from $A(e,e'p)$ and $A(e,e'pp)$ Scattering"
  arXiv:1503.06050 and Physical Review C **92** (2015), 024604.
- C. Colle, W. Cosyn, J. Ryckebusch, M. Vanhalst
  "Factorization of electroinduced two-nucleon knockout reactions"
  arXiv:1311.1980 and Physical Review C **89** (2014), 024603.
- Maarten Vanhalst, Jan Ryckebusch, Wim Cosyn
  "Quantifying short-range correlations in nuclei"
  arXiv:1206.5151 and Physical Review C **86** (2012), 044619.
- Maarten Vanhalst, Wim Cosyn, Jan Ryckebusch
  "Counting the amount of correlated pairs in a nucleus"
  arXiv:1105.1038 and Physical Review C **84** (2011), 031302(R).