CHIRAL DYNAMICS and NUCLEAR MATTER



Wolfram Weise

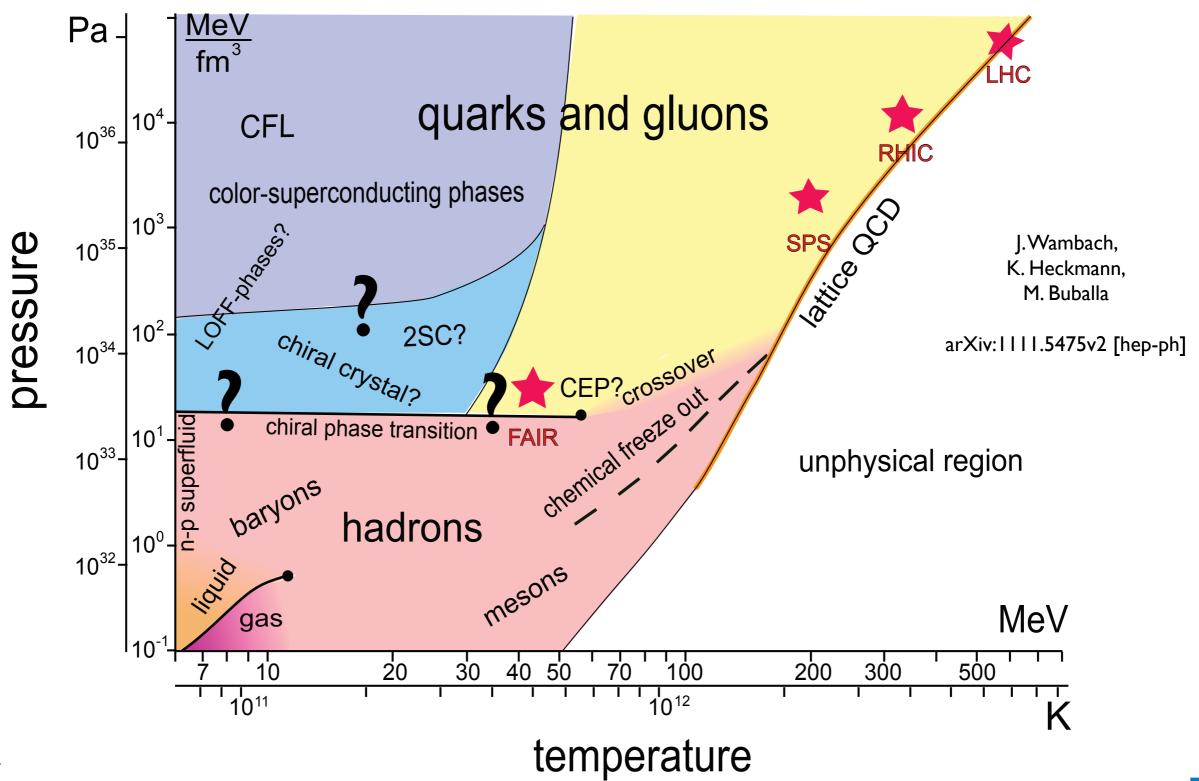
ECT* Trento and Technische Universität München



- Chiral EFT approaches to nuclear many-body systems
- Beyond mean field: fluctuations and Functional Renormalisation Group
- Nuclear matter, neutron matter and neutron stars
- Pion mass in the nuclear medium
- Thermodynamics of the chiral order parameter
- Outlook:
 Chiral SU(3) dynamics and hypernuclear matter

PHASES and STRUCTURES of QCD

- facts and visions -

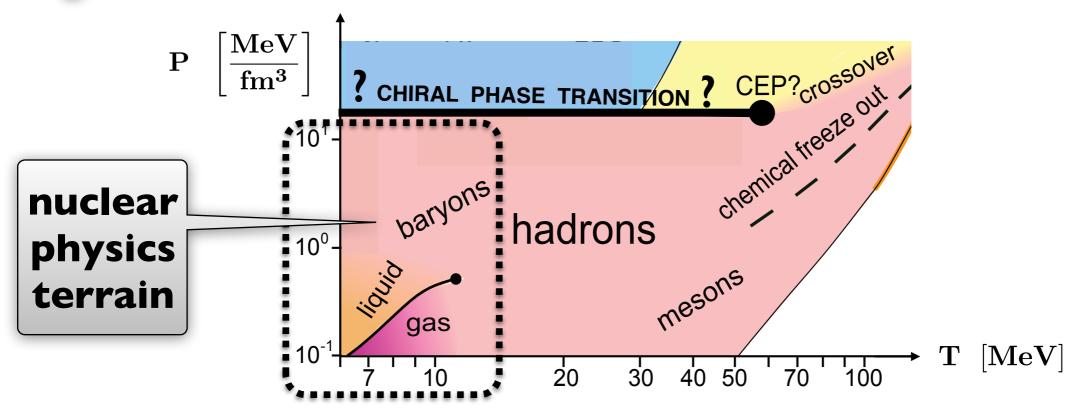




CHIRAL SYMMETRY RESTORATION

from Nambu-Goldstone to Wigner-Weyl Realisation of Chiral Symmetry

PHASE TRANSITION or smooth CROSSOVER?



- Chiral first-order phase transition and critical point?
- ... based on chiral quark models which do not respect nuclear physics constraints
 - Needed: systematic approach to nuclear thermodynamics



PIONS, NUCLEONS and NUCLEI in the context of LOW-ENERGY QCD

- CONFINEMENT of quarks and gluons in hadrons
- Spontaneously broken CHIRAL SYMMETRY
- LOW-ENERGY QCD with light (u,d) quarks:
 Effective Field Theory of (weakly) interacting
 Nambu-Goldstone Bosons (pions)
- Chiral EFT represents QCD at energy/momentum scales

$$\mathbf{Q} << 4\pi\,\mathbf{f}_\pi \sim \,\mathbf{1\,GeV}$$

Strategies at the interface between QCD and nuclear physics :

In-medium Chiral Perturbation Theory based on non-linear sigma model (with inclusion of nucleons)

expansion of free energy density in powers of Fermi momentum

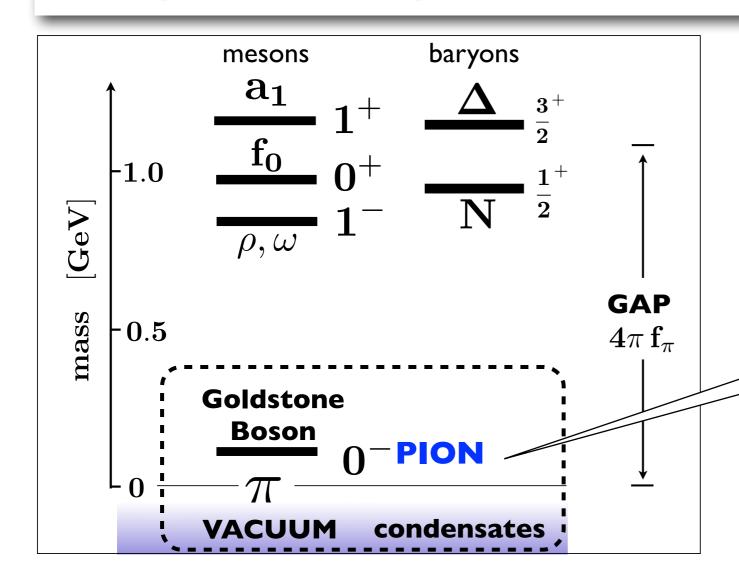
Chiral Nucleon-Meson model based on linear sigma model

non-perturbative Renormalization Group approach





Spontaneously Broken CHIRAL SYMMETRY



Triplet
of
NAMBU GOLDSTONE
BOSONS:

 $\pi^{+}, \ \pi^{0}, \ \pi^{-}$

Characteristic Symmetry Breaking SCALE:

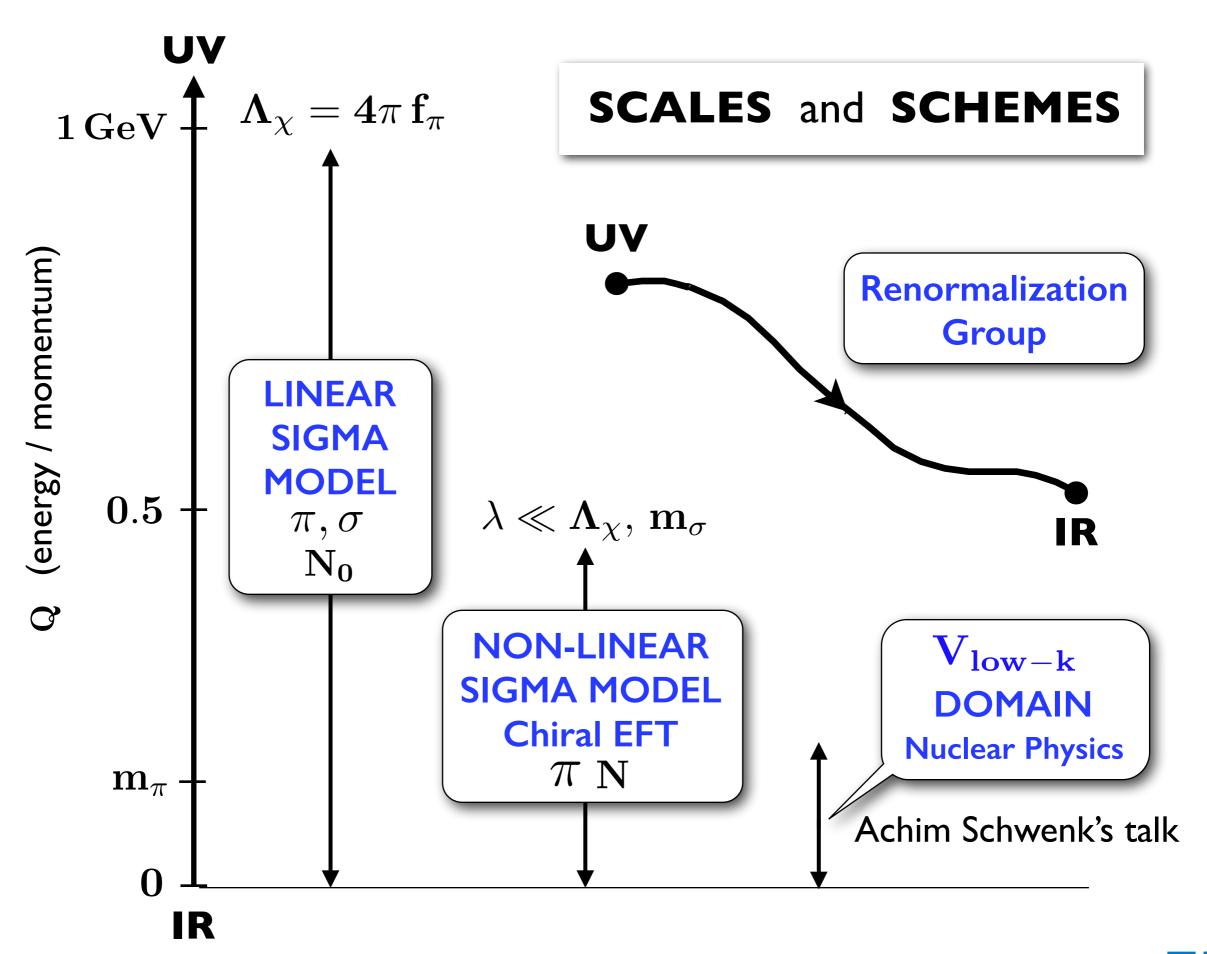
 $\mathbf{4}\pi\,\mathbf{f}_\pi\sim\mathbf{1}\,\,\mathbf{GeV}$

PION DECAY CONSTANT

Axial current
$$\mu$$
 $f_{\pi}=92.2\pm0.2~{
m MeV}$

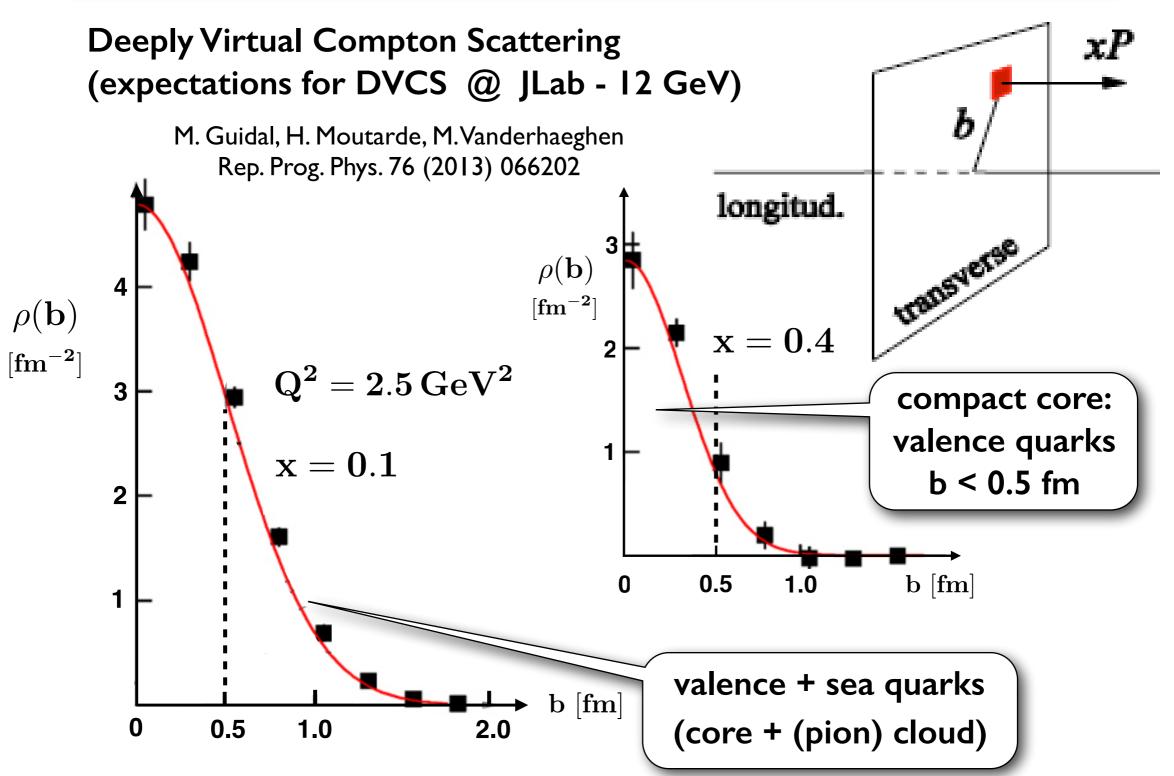








Transverse distributions of quarks in the proton core - plus - cloud structure?

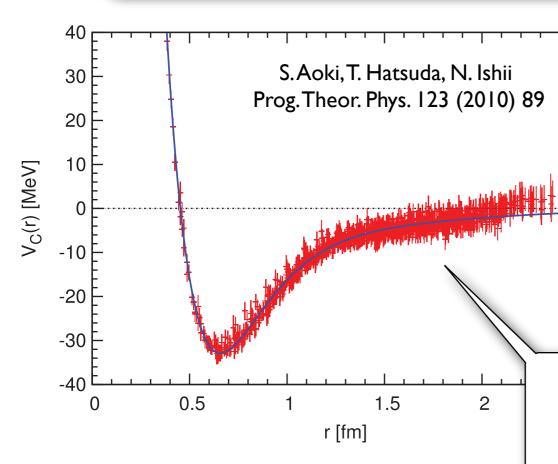




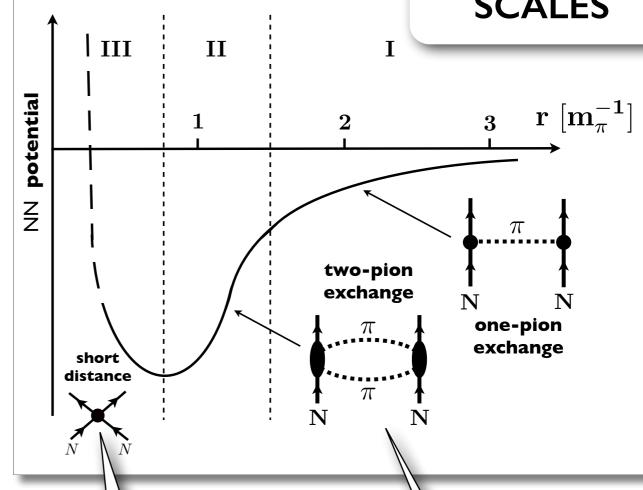


NN Interaction

Chiral Effective
Field Theory
&
Lattice QCD



Hierarchy of SCALES



contact terms

explicit treatment of two-pion exchange

NN Central Potential from Lattice QCD

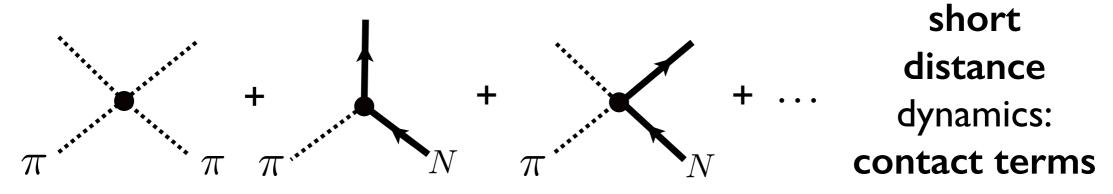


CHIRAL EFFECTIVE FIELD THEORY

- Realization of Low-Energy QCD
 based on Non-Linear Sigma Model plus (heavy) baryons
- Interacting systems of
 PIONS (light / fast) and NUCLEONS (heavy / slow):

$$\mathcal{L}_{eff} = \mathcal{L}_{\pi}(U, \partial U) + \mathcal{L}_{N}(\Psi_{N}, U, ...)$$
$$U(x) = exp[i\tau_{a}\pi_{a}(x)/f_{\pi}]$$

Construction of Effective Lagrangian: Symmetries







NUCLEAR INTERACTIONS from CHIRAL EFFECTIVE FIELD THEORY

Weinberg

Bedaque & van Kolck

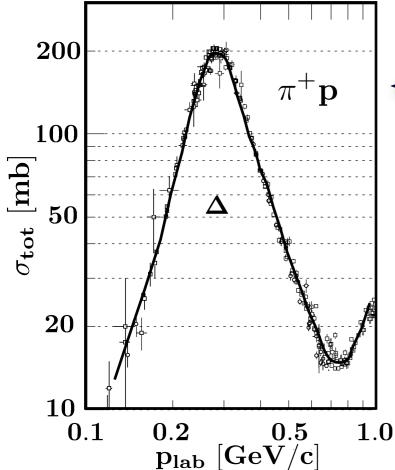
Bernard, Epelbaum, Kaiser, Meißner; ...

		Two-nucleon force	Three-nucleon force	Four-nucleon force
$\mathcal{O}\left(rac{\mathbf{Q^0}}{\mathbf{\Lambda^0}} ight)$	LO	X - 		
$\mathcal{O}\left(rac{\mathbf{Q^2}}{\mathbf{\Lambda^2}} ight)$	NLO	X科科科I	<u></u>	
(/	N ² LO	★ ★	 - - - - - - - 	
$\mathcal{O}\left(rac{\mathbf{Q^4}}{\mathbf{\Lambda^4}} ight)$	N ³ LO		本	





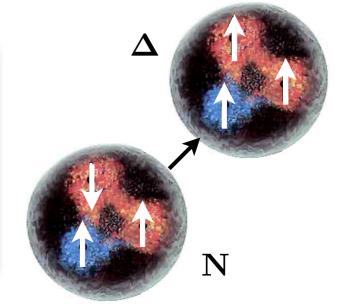
Explicit $\Delta(1230)$ DEGREES of FREEDOM



- Large spin-isospin polarizability of the Nucleon
- Dominance of $\Delta(1230)$ in pion-nucleon scattering

$$eta_{\Delta} = rac{\mathbf{g_A^2}}{\mathbf{f_\pi^2}(\mathbf{M_\Delta} - \mathbf{M_N})} \sim 5\,\mathrm{fm^3}$$
 $\mathbf{M_\Delta} - \mathbf{M_N} \simeq 2\,\,\mathbf{m_\pi} << 4\pi\,\mathbf{f_\pi}$

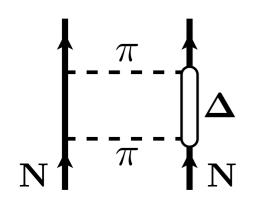
(small scale)



Pionic Van der Waals - type intermediate range central potential

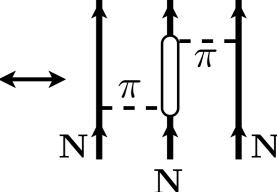
N. Kaiser, S. Gerstendörfer, W.W., NPA637 (1998) 395

N. Kaiser, S. Fritsch, W.W., NPA750 (2005) 259



I. Fujita, H. Miyazawa (1957)

$$egin{aligned} \mathbf{\Delta} & \mathbf{V_c(r)} = -rac{9\,\mathbf{g_A^2}}{32\pi^2\,\mathbf{f_\pi^2}}\,eta_\Delta\,rac{\mathbf{e^{-2m_\pi r}}}{\mathbf{r^6}}\,\mathbf{P(m_\pi r)} \end{aligned} egin{aligned} \longleftrightarrow & \mathbf{P(m_\pi r)} \end{aligned}$$

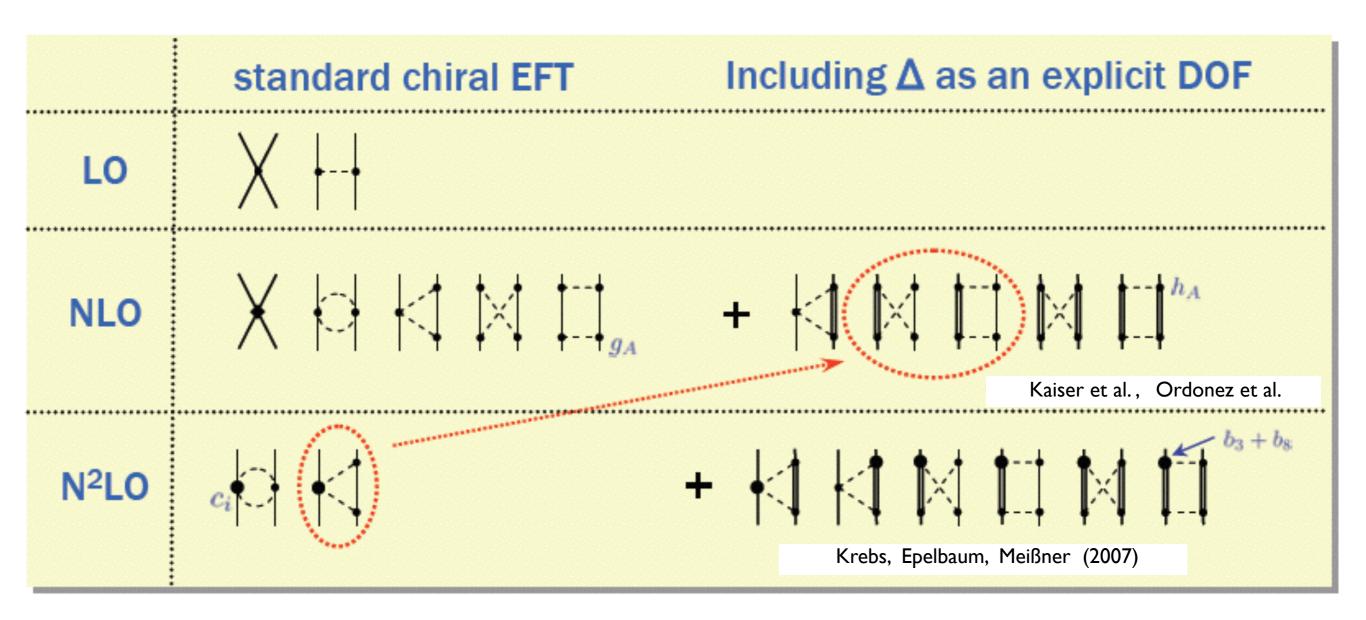


Pieper, Pandharipande, Wiringa, Carlson (2001)

strong 3-body interaction



Explicit $\Delta(1230)$ DEGREES of FREEDOM (contd.)

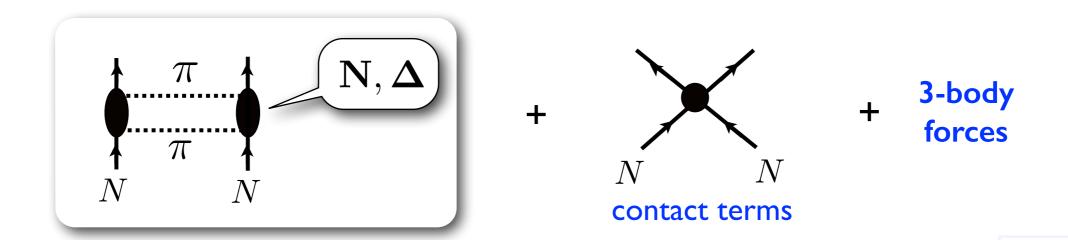


- lacksquare Important physics of $oldsymbol{\Delta}(1230)$ promoted to NLO
 - Improved convergence



Important:

Explicit treatment of two-pion exchange dynamics

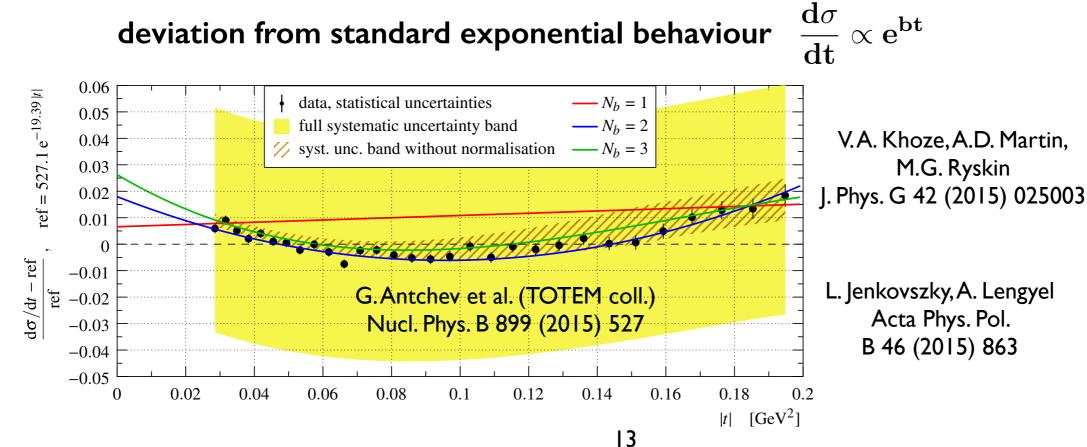


Short digression:

"Discovery" of two-pion exchange at LHC: elastic pp scattering at $\sqrt{s}=8\,\mathrm{TeV}$

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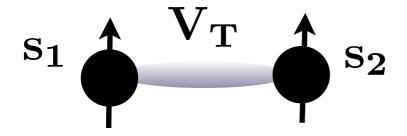
deviation from standard exponential behaviour



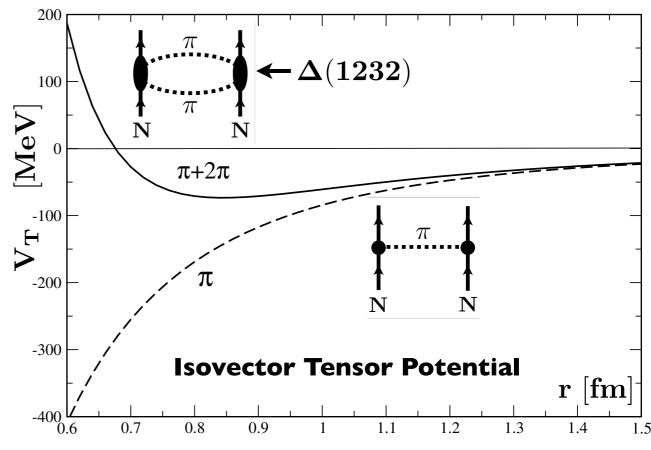
Important pieces of the

CHIRAL NUCLEON-NUCLEON INTERACTION

ISOVECTORTENSOR FORCE

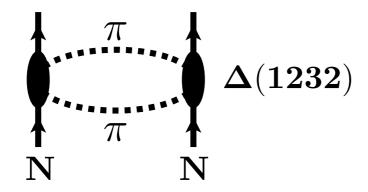


 \bullet note: **no** ρ meson



N. Kaiser, S. Gerstendörfer, W.W.: Nucl. Phys. A 637 (1998) 395

CENTRAL ATTRACTION from TWO-PION EXCHANGE



ullet note: **no** σ boson

Van der WAALS - like force:

$$\mathbf{V_c}(\mathbf{r}) \propto -rac{\exp[-2\mathbf{m_\pi r}]}{\mathbf{r}^6}\mathbf{P}(\mathbf{m_\pi r})$$

... at intermediate and long distance





IN-MEDIUM CHIRAL PERTURBATION THEORY

Small scales:

energy, momentum, $m_\pi,~k_F << 4\pi f_\pi \sim 1\,GeV$

"Medium insertion" in the nucleon propagator:

$$(\gamma_{\mu}\mathbf{p}^{\mu} + \mathbf{M_{N}}) \left[\frac{\mathbf{i}}{\mathbf{p^{2} - M_{N}^{2} + i\varepsilon}} - 2\pi \, \delta(\mathbf{p^{2} - M_{N}^{2}}) \, \theta(\mathbf{p^{0}}) \, \theta(\mathbf{k_{F}} \cdot |\vec{\mathbf{p}}|) \right]$$

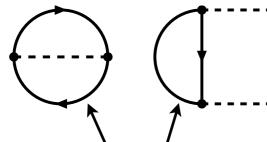
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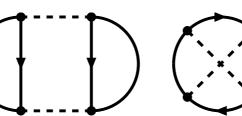
Loop expansion of (In-Medium) Chiral Perturbation Theory

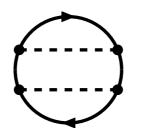
 \longleftrightarrow Expansion of ENERGY DENSITY $\mathcal{E}(\mathbf{k_F})$ in powers of Fermi momentum [modulo functions $\mathbf{f_n}(\mathbf{k_F}/\mathbf{m_\pi})$]



Nuclear thermodynamics: compute free energy density





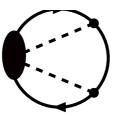


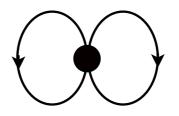
(3-loop order)

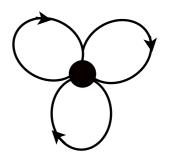
N. Kaiser, S. Fritsch, W.W. (2002-2004)

in-medium nucleon propagators incl. Pauli blocking











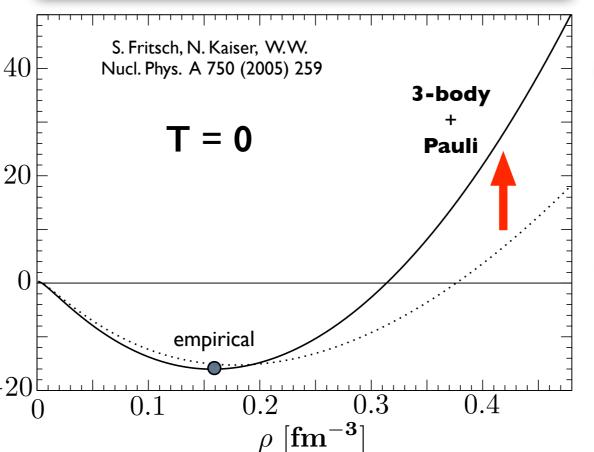
NUCLEAR MATTER

In-medium ChPT

3-loop $(\pi, \mathbf{N}, \boldsymbol{\Delta})$

Input parameters: few contact terms

basically: analytic calculation



- Binding, saturation
- Asymmetry energy
- Nuclear thermodynamics: liquid-gas phase transition
- Realistic (complex, momentum dependent) single-particle potential ... satisfying Hugenholtz van Hove and Luttinger theorems (!)
- Fermi Liquid Theory:
 Quasiparticle interaction and Landau parameters
- Nuclear Energy Density Functional and finite nuclei

J.W. Holt, N. Kaiser, W.W. (2011 - 2013)

C.Wellenhofer, J.W. Holt, N. Kaiser, W.W. Phys. Rev. C 89 (2014) 064009

Recent reviews: J.W. Holt, N. Kaiser, W.W. J.W. Holt, M. Rho, W.W.

 ${f E}/{f A}$ [MeV]

Prog. Part. Nucl. Phys. 73 (2013) 35 arXiv:1411.6681, Phys. Reports (2015)



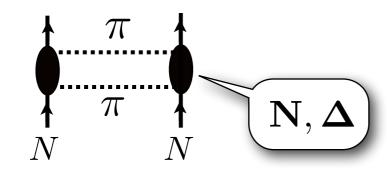


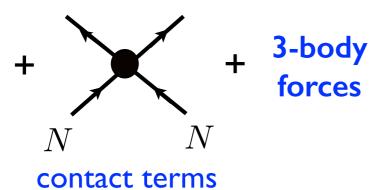
NUCLEAR THERMODYNAMICS

NUCLEAR
CHIRAL (PION) DYNAMICS

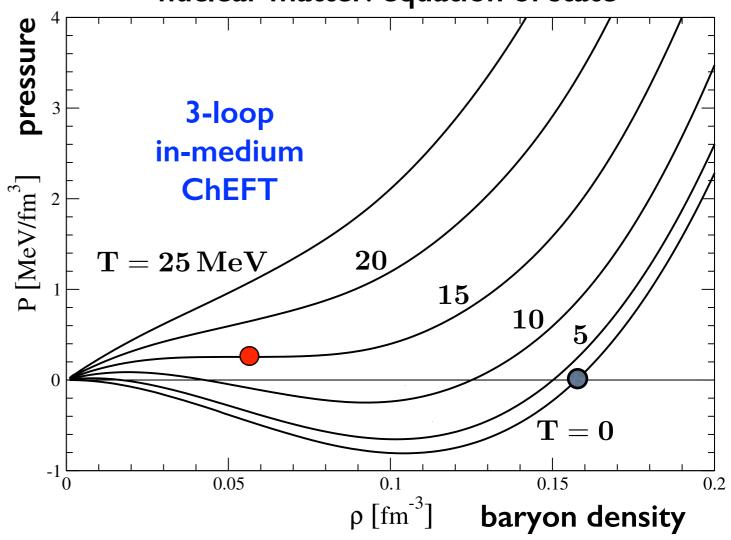
BINDING & SATURATION:

Van der Waals + Pauli





nuclear matter: equation of state

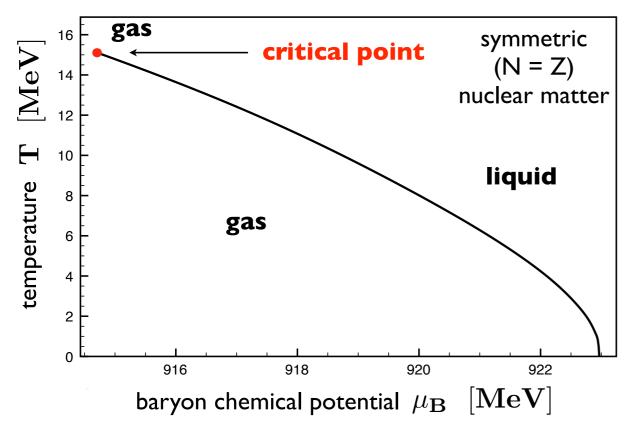


Liquid - Gas Transition at Critical Temperature $T_c = 15 \text{ MeV}$ (empirical: $T_c = 16 - 18 \text{ MeV}$)





PHASE DIAGRAM of NUCLEAR MATTER

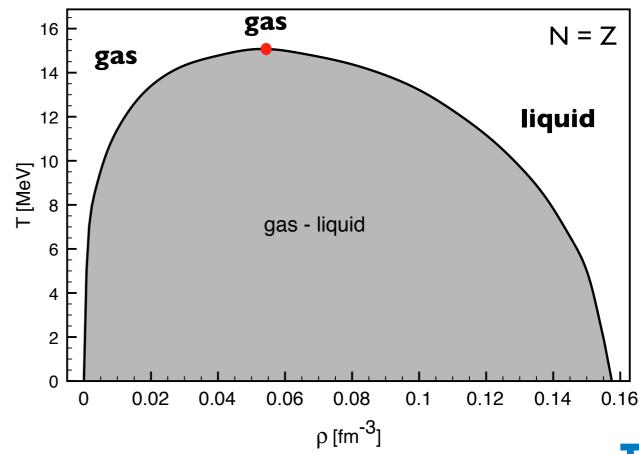


- Pion-nucleon dynamics including delta isobars
- Short-distanceNN contact terms
- Three-body forces

In-medium
 chiral effective field theory
 (3-loop calculation of free energy density)

S. Fritsch, N. Kaiser, W.W. Nucl. Phys. A 750 (2005) 259

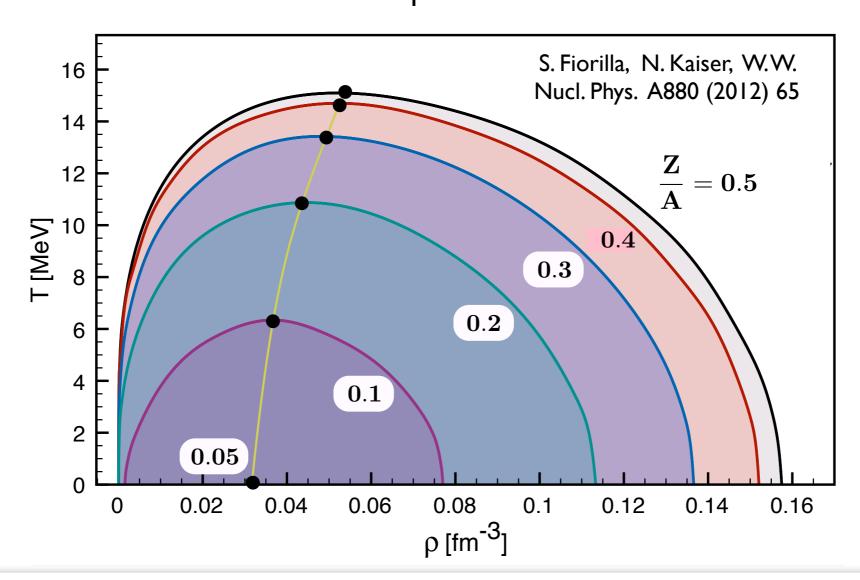
S. Fiorilla, N. Kaiser, W.W. Nucl. Phys. A 880 (2012) 65





PHASE DIAGRAM of NUCLEAR MATTER

Trajectory of CRITICAL POINT for asymmetric matter as function of proton fraction Z/A



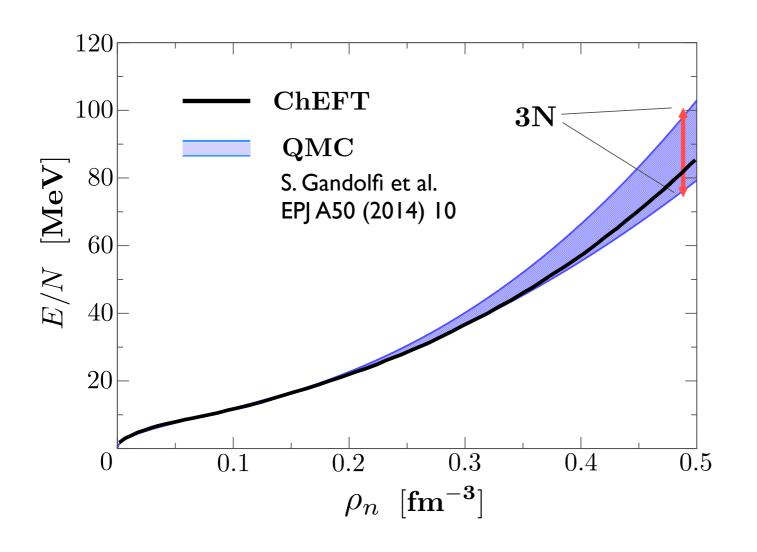
... determined almost completely by isospin dependent (one- and two-) pion exchange dynamics





NEUTRON MATTER

In-medium chiral effective field theory (3-loop) with resummation of short distance contact terms (large nn scattering length, $a_{\rm s}=19~{
m fm})$



- Neutron matter behaves almost (but not quite) like a unitary Fermi gas
- Bertsch parameter

$$\xi = rac{ar{ ext{E}}}{ ext{E}_{ ext{Fermi gas}}} \simeq 0.5$$

J.W. Holt, N. Kaiser, W.W. Phys. Rev. C 87 (2013) 014338

agreement with sophisticated many-body calculations
 (e.g. recent Quantum Monte Carlo computations)

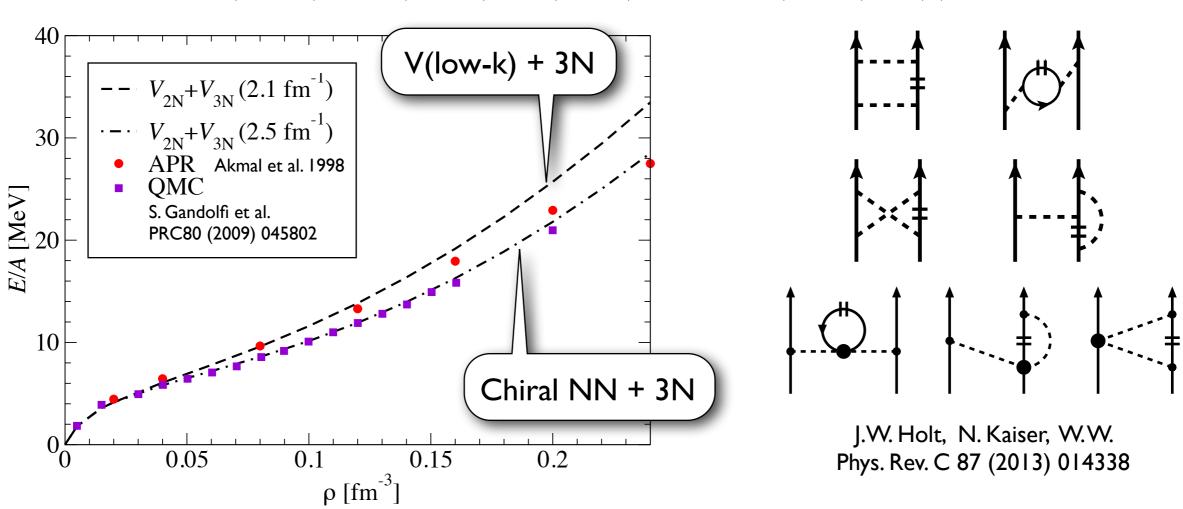


Chiral Fermi Liquid Approach to Neutron Matter

Quasiparticle interaction based on accurate NNLO chiral nucleon-nucleon interaction including three-body forces:

$$\delta \mathcal{E} = \sum_{\vec{p}st} \epsilon_{\vec{p}} \, \delta n_{\vec{p}st} + \frac{1}{2} \sum_{\substack{\vec{p}_1 s_1 t_1 \\ \vec{p}_2 s_2 t_2}} \mathcal{F}(\vec{p}_1 s_1 t_1; \, \vec{p}_2 s_2 t_2) \delta n_{\vec{p}_1 s_1 t_1} \delta n_{\vec{p}_2 s_2 t_2} + \cdots,$$

$$\mathcal{F}(\vec{p}_1, \vec{p}_2) = f(\vec{p}_1, \vec{p}_2) + g(\vec{p}_1, \vec{p}_2)\vec{\sigma}_1 \cdot \vec{\sigma}_2 + h(\vec{p}_1, \vec{p}_2)S_{12}(\hat{q}) + \dots$$



Quantum Monte Carlo calculations with ChEFT Interactions
 A. Gezerlis et al.
 exhibit systematic order-by-order convergence
 Phys. Rev. Lett. 111 (2013) 032501



Chiral Nucleon-Meson Model

(based on Linear Sigma Model)

and

Functional Renormalization Group





Mesons, Nucleons, Nuclear Matter and

Functional Renormalization Group

Chiral nucleon - meson model

$$\psi = (\psi_p, \psi_n)^T$$

$$\mathcal{L} = \bar{\psi}i\gamma_{\mu}\partial^{\mu}\psi + \frac{1}{2}\partial_{\mu}\sigma\,\partial^{\mu}\sigma + \frac{1}{2}\partial_{\mu}\boldsymbol{\pi}\cdot\partial^{\mu}\boldsymbol{\pi}$$

$$-\bar{\psi}\left[g(\sigma + i\gamma_{5}\boldsymbol{\tau}\cdot\boldsymbol{\pi}) + \gamma_{\mu}(g_{\omega}\,\omega^{\mu} + g_{\rho}\boldsymbol{\tau}\cdot\boldsymbol{\rho}^{\mu})\right]\psi$$

$$-\frac{1}{4}F_{\mu\nu}^{(\omega)}F^{(\omega)\mu\nu} - \frac{1}{4}\boldsymbol{F}_{\mu\nu}^{(\rho)}\cdot\boldsymbol{F}^{(\rho)\mu\nu}$$

$$+\frac{1}{2}m_{V}^{2}\left(\omega_{\mu}\omega^{\mu} + \boldsymbol{\rho}_{\mu}\cdot\boldsymbol{\rho}^{\mu}\right) - \mathcal{U}(\sigma,\boldsymbol{\pi})$$

- Effective potential constructed to reproduce standard nuclear thermodynamics around equilibrium
- Mean field calculations
 S. Floerchinger, Ch. Wetterich: Nucl. Phys. A 890-891 (2012) 11
- Mesonic and nucleonic particle-hole fluctuations treated non-perturbatively using FRG

M. Drews, T. Hell, B. Klein, W.W. Phys. Rev. D88 (2013) 096011M. Drews, W.W. Phys. Lett. B738 (2014) 187 Phys. Rev. C91 (2015) 035802

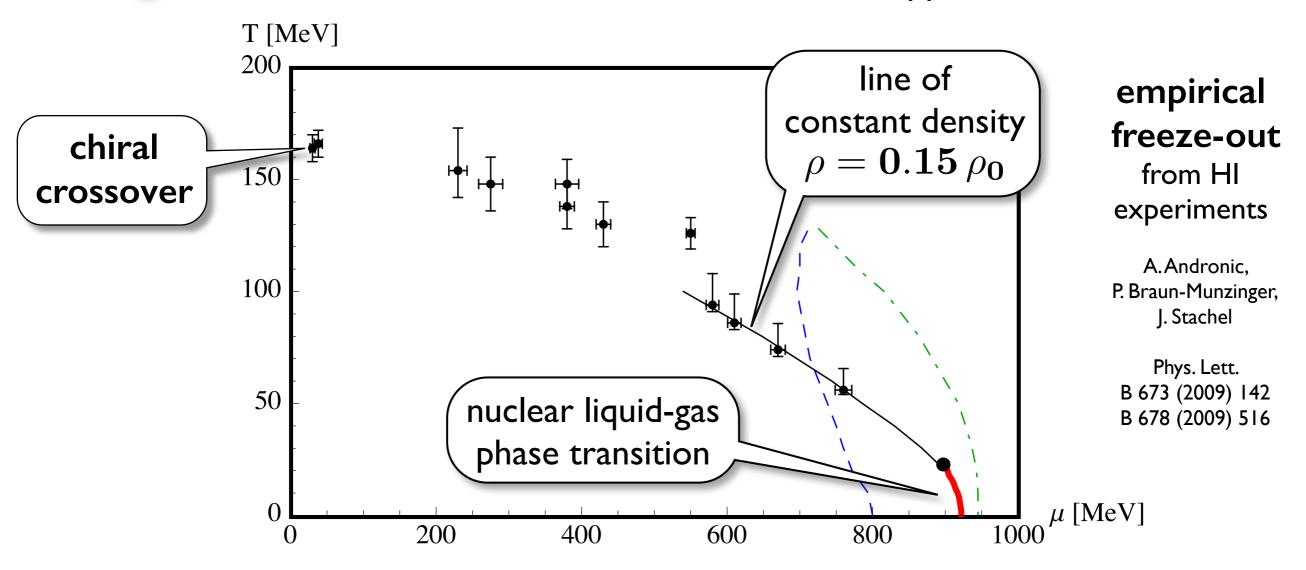




CHEMICAL FREEZE-OUT

S. Floerchinger, Ch. Wetterich: Nucl. Phys. A 890-891 (2012) 11

Chiral nucleon - meson model in mean-field approximation



 Chemical freeze-out in baryonic matter at T < 100 MeV is not associated with (chiral) phase transition or rapid crossover



Fixing the input: some comments

• Potential $\mathcal{U}(\sigma, \boldsymbol{\pi}) = \mathcal{U}_0(\chi) - m_\pi^2 f_\pi(\sigma - f_\pi)$

chiral invariant part parametrized in powers of $\chi = \frac{1}{2}(\sigma^2 + \pi^2)$

explicit chiral symmetry breaking

- Scalar ("sigma") field has mean-field (chiral order parameter) and fluctuating pieces. σ mass: NOT to be identified with " $\sigma(500)$ " pole in I = 0 s-wave pion-pion T matrix Nucleon mass: $m_N^2=2g\,\chi\,\ldots$ in vacuum: $m_N=g\,f_\pi$
- Vector fields encode short-distance NN dynamics, self-consistently determined background mean fields (non-fluctuating) (NOT to be identified with physical ω and ρ mesons)

Effective chemical potentials $\mu_{n,p}^{\rm eff}=\mu_{n,p}-g_\omega\,\omega_0\,\pm g_\rho\,\rho_0^3$

Relevant quantities: $G_{\rho}=\frac{g_{\rho}^2}{m_V^2}$, $G_{\omega}=\frac{g_{\omega}^2}{m_V^2}$ \longleftrightarrow contact terms in ChEFT

• Parameters: 2 coefficients in \mathcal{U}_0 , $m_\sigma \simeq 0.8\, GeV$, $G_\rho \sim G_\omega/4 \simeq 1\, fm^2$ determined by nuclear matter properties and symmetry energy





Renormalization Group strategies Chiral nucleon-meson model beyond mean-field

M. Drews, T. Hell, B. Klein, W.W.

Phys. Rev. D 88 (2013) 096011

Fluctuations: Wetterich's RG flow equations

effective action

full propagator

$$k \frac{\partial \Gamma_k}{\partial k} = \underbrace{ = \frac{1}{2} \operatorname{Tr} \frac{k \frac{\partial R_k}{\partial k}}{\Gamma_k^{(2)} + R_k}}$$

scale regulator:
$$R_k(p^2) = (k^2 - p^2) \theta(k^2 - p^2)$$

Thermodynamics:

nucleons

pions

$$k \,\partial_k \bar{\Gamma}_k(T,\mu) = \left(\bigotimes + \bigotimes \right) \Big|_{T,\mu_p,\mu_n}$$

$$-\left(\bigotimes + \bigotimes \right)$$

C. Wetterich: Phys. Lett. B 301 (1993) 90

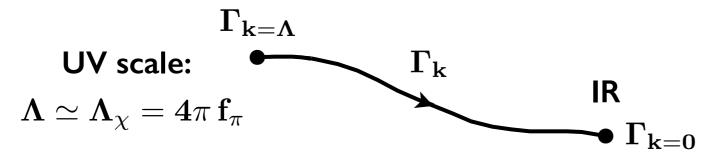
> Non-perturbative treatment of

- multi-pion exchange processes
- nucleon-hole excitations
- multi-nucleon correlations





Flow equations in practice



"full" effective action effective potential $\mathbf{U}_{\mathbf{k}=\mathbf{0}}$

$$k \frac{\partial U_k}{\partial k} (T, \mu_p, \mu_n, \chi, \omega_0, \rho_0^3) = \bigotimes + \bigotimes + \bigotimes + \bigotimes + \underbrace{\frac{\partial U_k}{\partial k}}_{E_{\sigma}} \left\{ \frac{1 + 2n_{\rm B}(E_{\sigma})}{E_{\sigma}} + \frac{3[1 + 2n_{\rm B}(E_{\pi})]}{E_{\pi}} - 4\sum_{i=n,p} \frac{1 - n_{\rm F}(E_{\rm N} - \mu_{i,\text{eff}})}{E_{\rm N}} \right\}$$

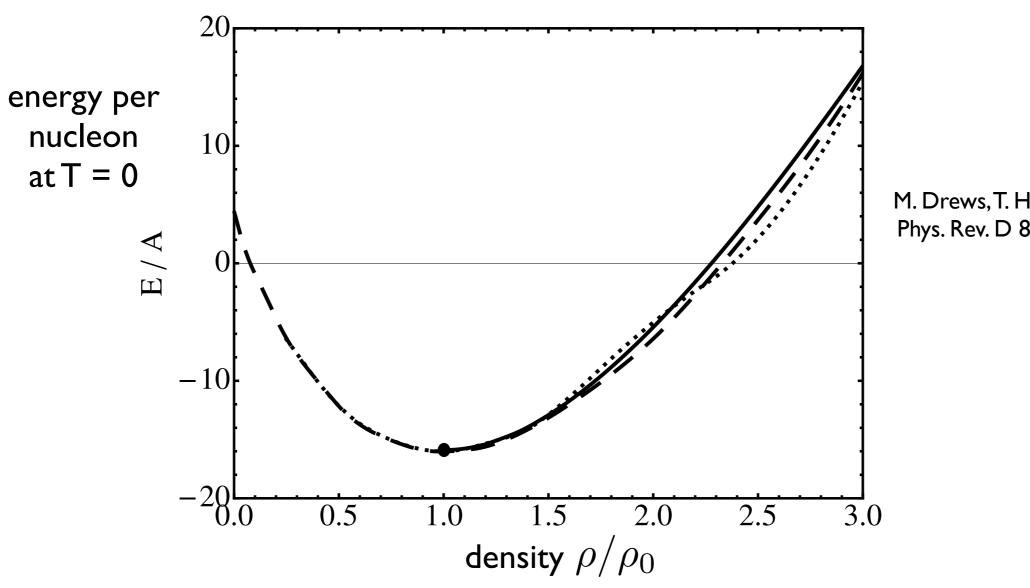
$$\begin{split} E_{\pi}^2 &= k^2 + U_k'(\chi) \,, \ E_{\sigma}^2 = k^2 + U_k'(\chi) + 2\chi \, U_k''(\chi) \,, \qquad n_{\rm B}(E) = \frac{1}{{\rm e}^{E/T} - 1} \,, \\ U_k'(\chi) &= \frac{\partial U_k(\chi)}{\partial \chi} \,, \quad E_{\rm N}^2 = k^2 + 2g^2 \chi \,, \\ \mu_{n,p}^{\rm eff}(k) &= \mu_{n,p} - g_{\omega} \, \omega_0(k) \pm g_{\rho} \, \rho_0^3(k) \,, \end{split} \qquad n_{\rm F}(E) &= \frac{1}{{\rm e}^{E/T} + 1} \,. \end{split}$$

... plus vector field equations, then full system of equations solved on a grid.





Symmetric nuclear matter in the chiral FRG approach



M. Drews, T. Hell, B. Klein, W.W. Phys. Rev. D 88 (2013) 096011

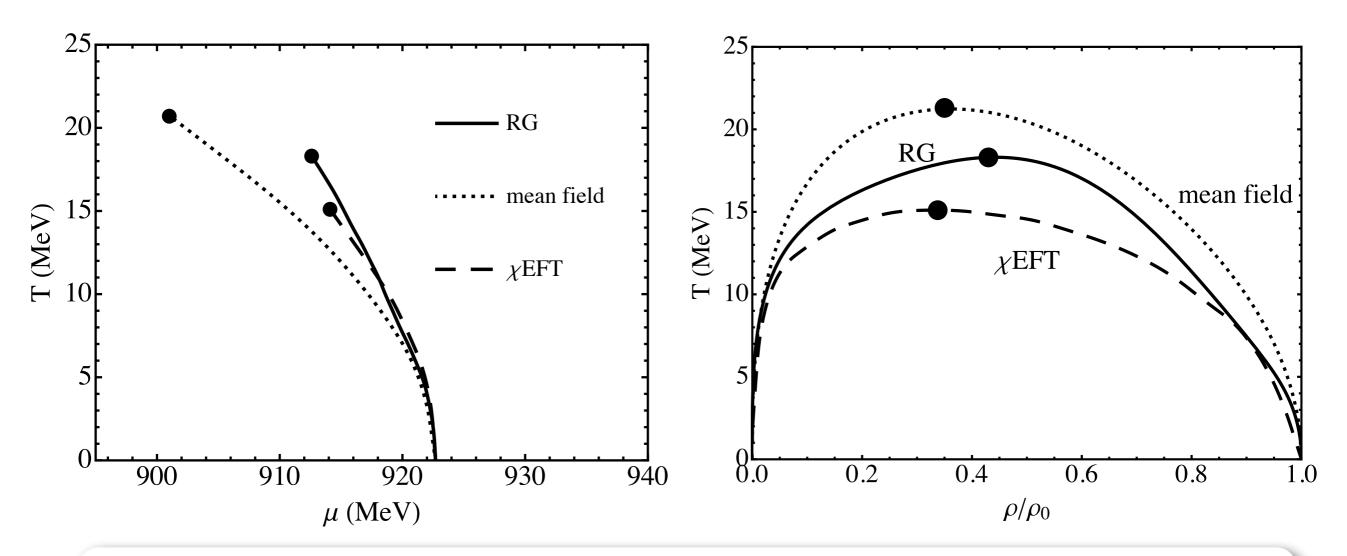
 FRG-Nucleon-Meson-Model (solid curve) in comparison with advanced many-body (variational and QMC) computations



Results: Liquid - Gas Transition

- symmetric nuclear matter -

M. Drews, T. Hell, B. Klein, W.W. Phys. Rev. D 88 (2013) 096011



close correspondence between (perturbative) in-medium ChEFT
 and (non-perturbative) FRG results

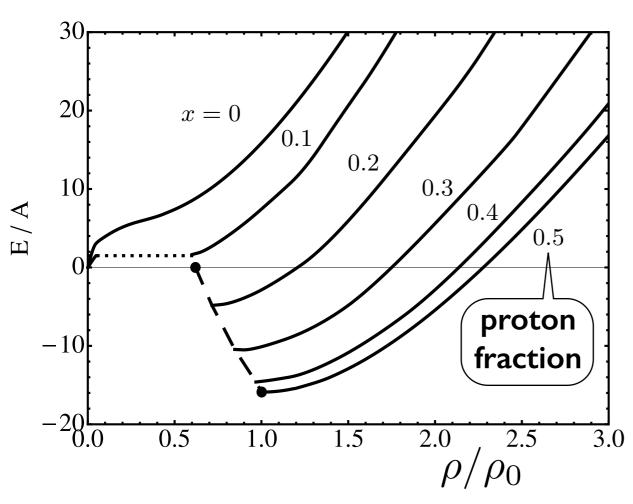


Asymmetric nuclear matter in the chiral FRG approach

M. Drews, W.W.

Phys. Lett. B738 (2014) 187

Phys. Rev. C91 (2015) 035802

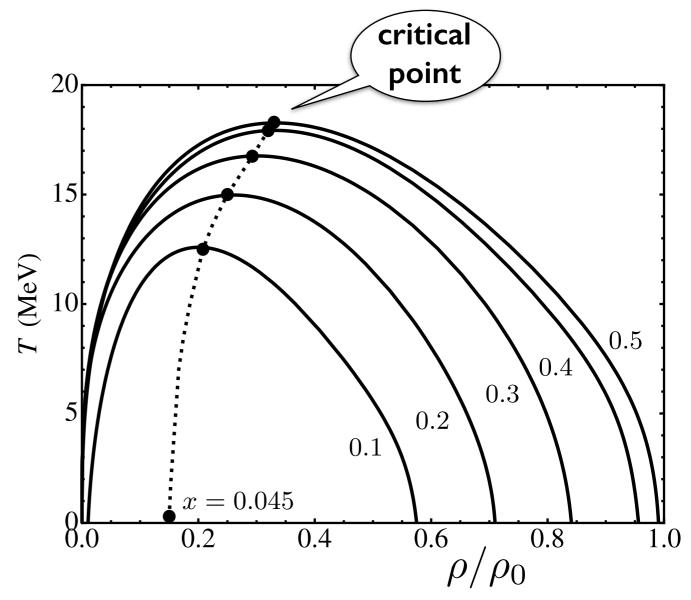


FRG results

 (non-perturbative)
 are remarkably similar to
 (perturbative) in-medium
 Chiral EFT calculations

Liquid-gas phase transition:

evolution of coexistence regions from symmetric to asymmetric nuclear matter

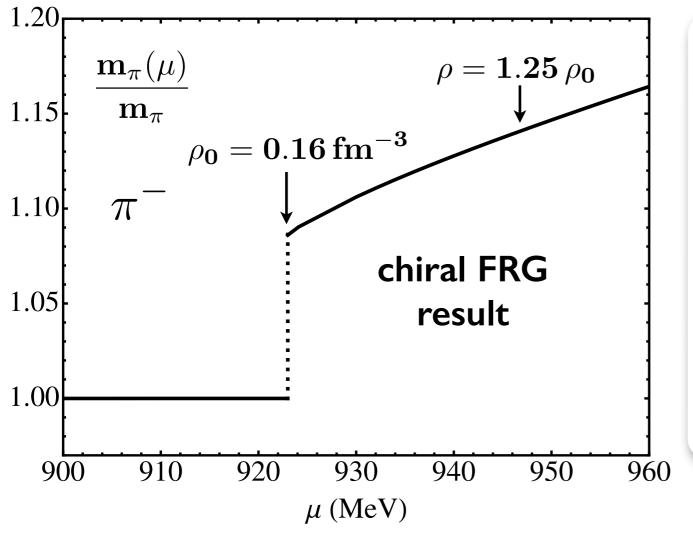


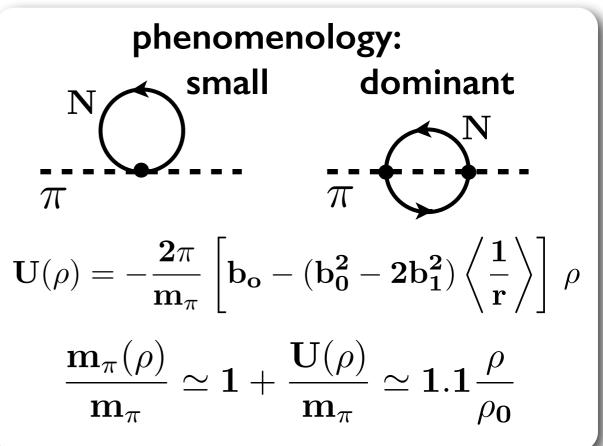




In-medium pion mass

Test case and contact with phenomenology:
 compare with s-wave pion-nuclear optical potential from pionic atoms



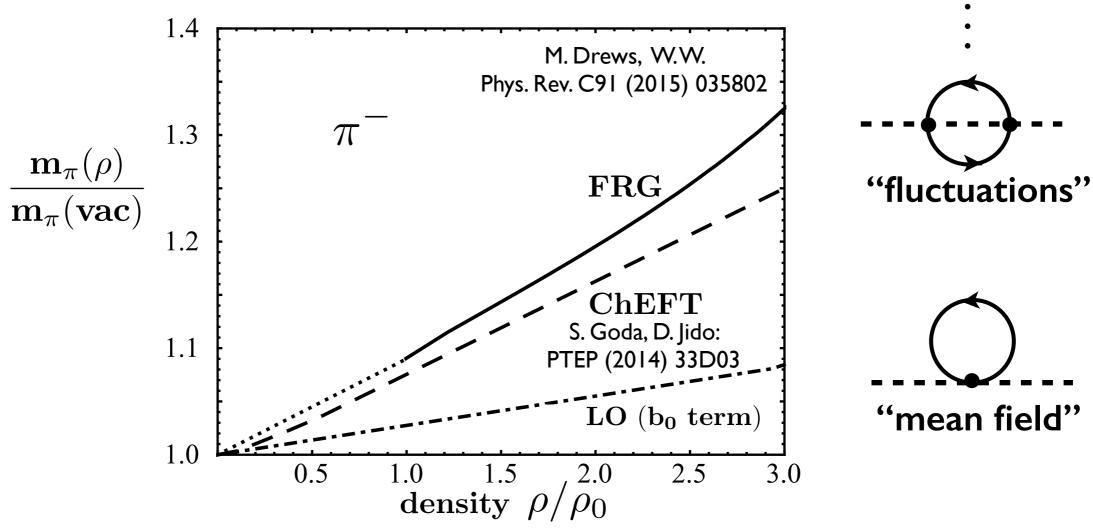


 Good agreement of FRG calculation with empirical in-medium pion mass shift, both in sign and magnitude



In-medium pion mass (contd.)

 Non-perturbative FRG result in comparison with in-medium Chiral Perturbation Theory



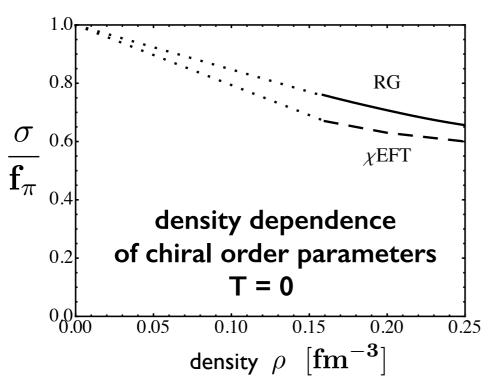
in-medium ChEFT (NLO):

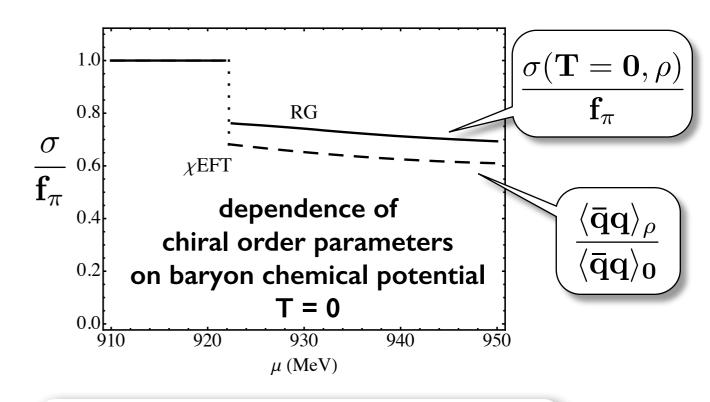
$$m_{\pi}^{*} = m_{\pi} \left\{ 1 + \rho \frac{b^{+}}{2m_{\pi}^{2}} - \frac{g_{A}^{2}k_{F}^{4}}{24\pi^{4}f_{\pi}^{4}}F(\frac{m_{\pi}}{2k_{F}}) + \left[\frac{1}{8} + m_{\pi}^{2} \left(\frac{b^{+}}{2m_{\pi}^{2}} - \frac{g_{A}^{2}}{8m_{N}} \right)^{2} \right] \frac{2k_{F}^{4}}{\pi^{4}f_{\pi}^{4}} \right\}$$

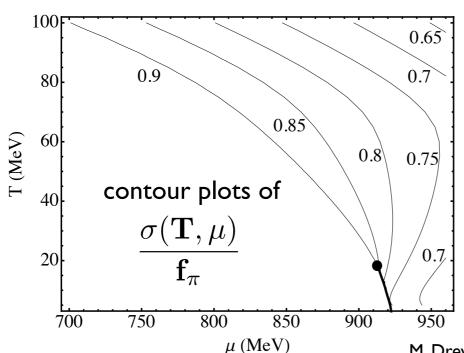


Chiral Order Parameters

Comparison of chiral effective field theory and NM-FRG results



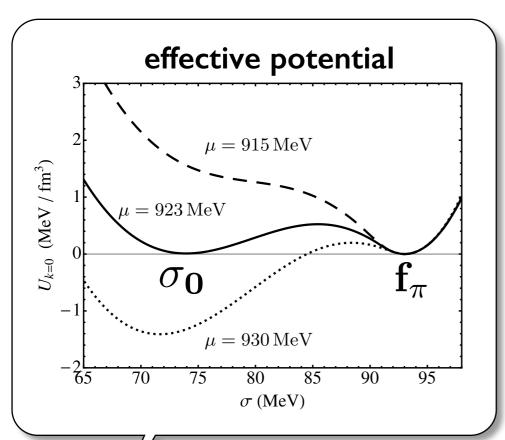




ullet No tendency towards chiral phase transition for baryon chemical potentials $\mu \lesssim 1~{
m GeV}$ and temperatures ${
m T} \lesssim 100~{
m MeV}$





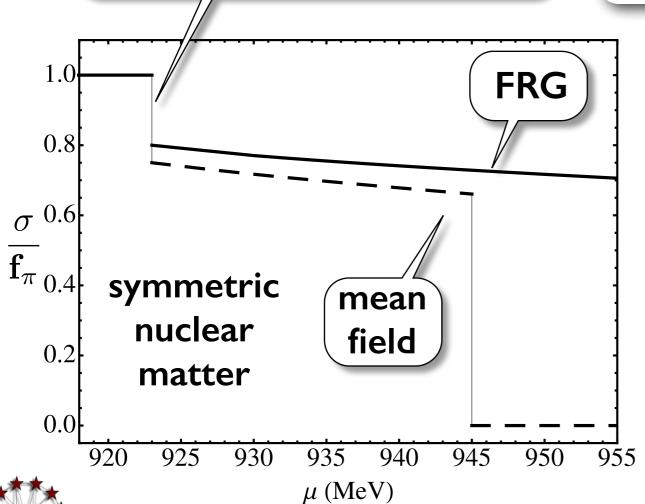


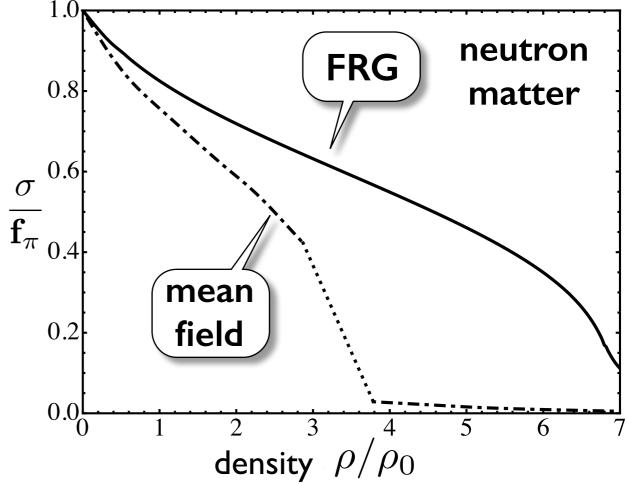
Chiral Order Parameter

M. Drews, W.W. Phys. Rev. C91 (2015) 035802

important role of **fluctuations** beyond mean-field approximation:

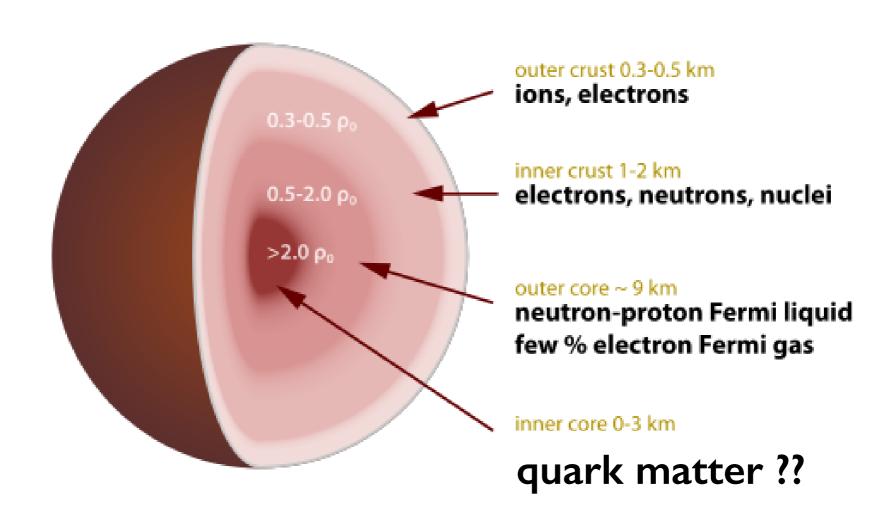
DISAPPEARANCE of first-order chiral phase transition





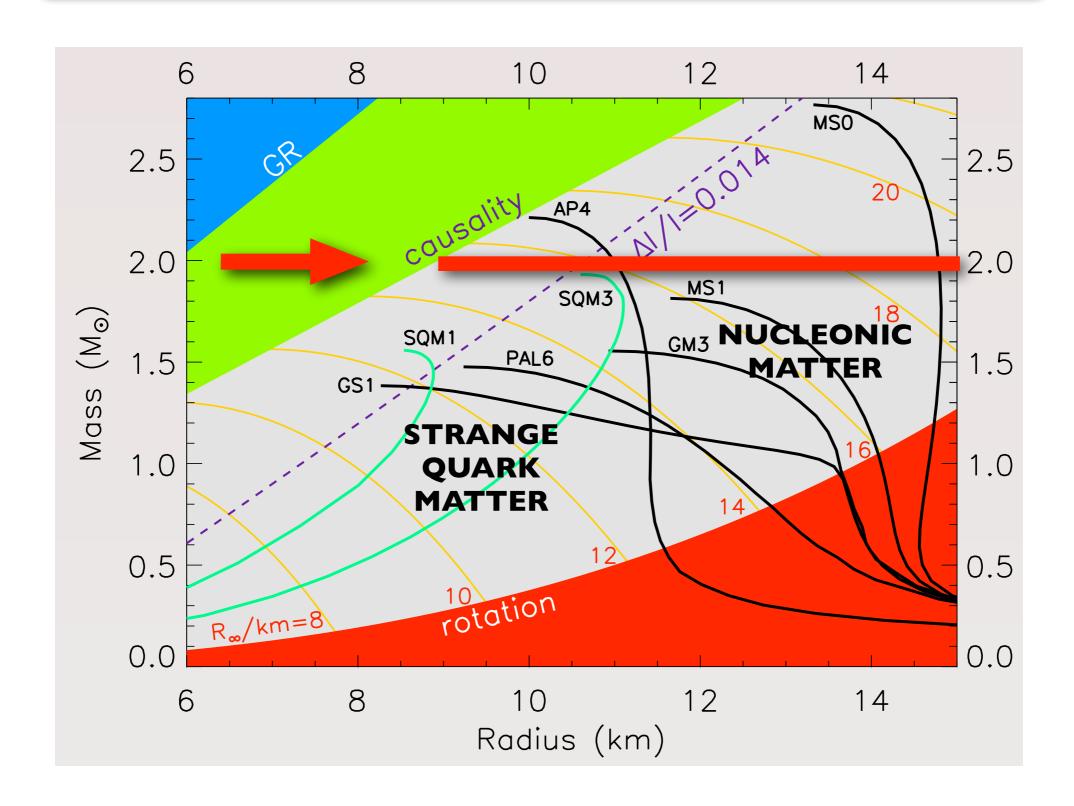


News from NEUTRON STARS





New constraints from 2-solar-mass NEUTRON STARS

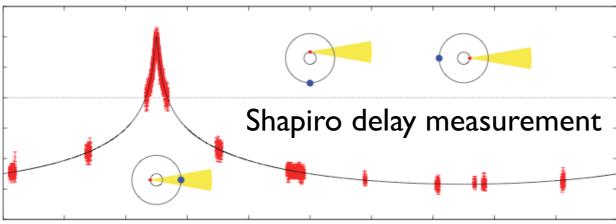


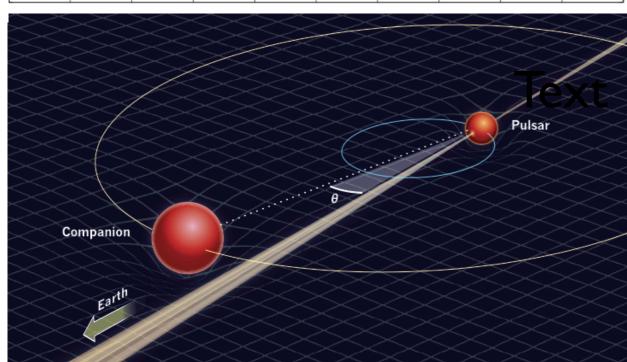




New constraints from NEUTRON STARS

P.B. Demorest et al. Nature 467 (2010) 1081

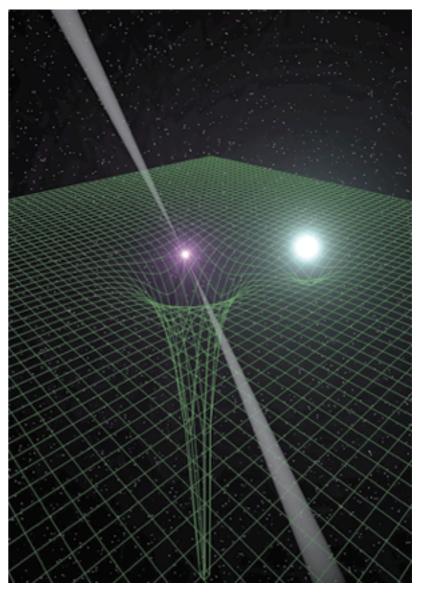




PSR J1614+2230

$$M=1.97\pm0.04~M_{\odot}$$

J. Antoniadis et al. Science 340 (2013) 6131



PSR J0348+0432

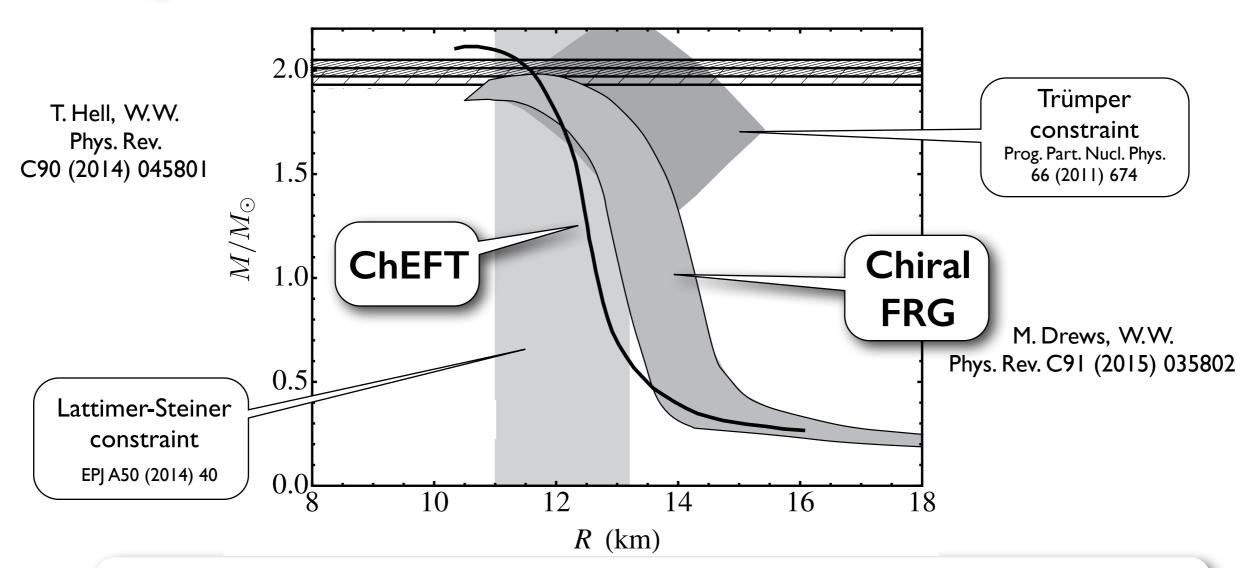
$$\mathbf{M} = 2.01 \pm 0.04 \ \mathrm{M}_{\odot}$$



NEUTRON STAR MATTER

from Chiral Nucleon-Meson Approach and FRG

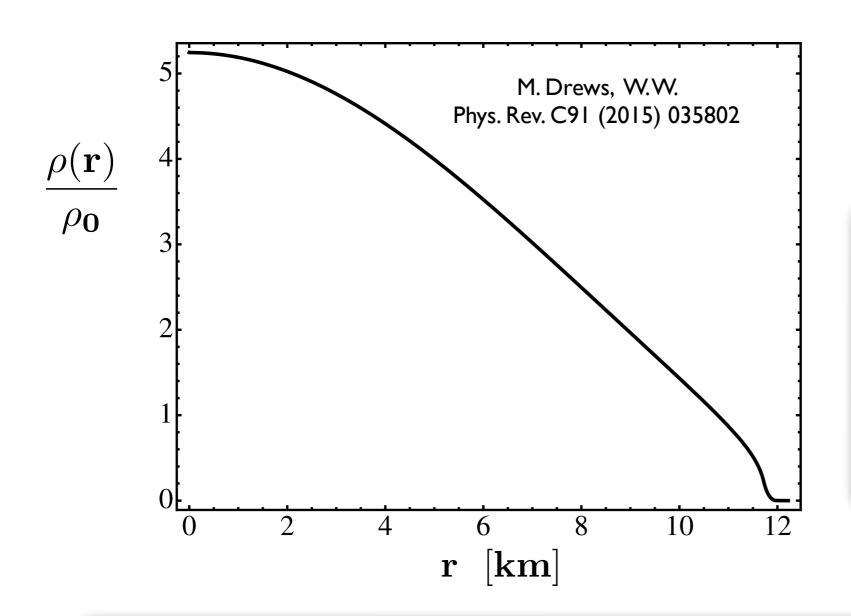
- Neutron matter plus proton admixture (beta equilibrium)
- Symmetry energy range: 30 37 MeV



Chiral many-body dynamics using "conventional" (pion & nucleon) degrees of freedom is consistent with neutron star constraints



Density profile of two-solar-mass neutron star



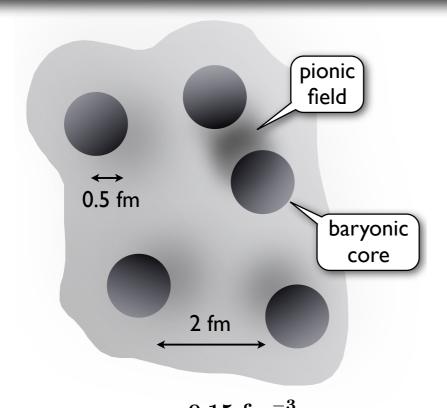
Chiral FRG calculation

maximum n-star core density:

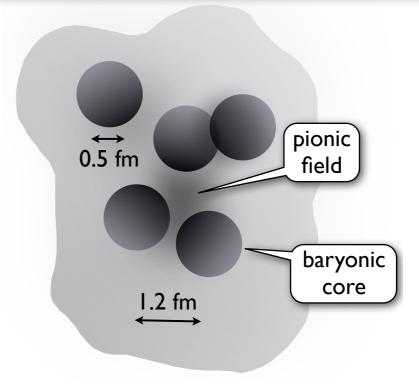
$$ho_{\mathbf{max}} \sim \mathbf{5} \,
ho_{\mathbf{0}}$$
 $(
ho_{\mathbf{0}} = \mathbf{0.16} \; \mathbf{fm^{-3}})$

No ultrahigh densities in the neutron star core

Densities and Scales in Compressed Baryonic Matter



 $ho_{
m B} = 0.15~{
m fm^{-3}}$ normal nuclear matter: dilute

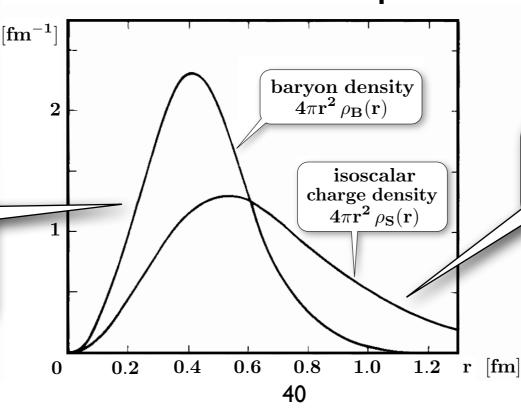


 $ho_{
m B}=0.6~{
m fm}^{-3}$

neutron star core matter: compressed but not superdense

chiral (soliton) modelof the nucleon

compact $extbf{baryonic core} \ \langle extbf{r}^2
angle_{ extbf{B}}^{1/2} \simeq 0.5 \,\, ext{fm}$



N. Kaiser, U.-G. Meißner, W.W. Nucl. Phys. A 466 (1987) 685

mesonic cloud

$$\langle {
m r^2}
angle_{
m E,isoscalar}^{1/2} \simeq 0.8 {
m \ fm}$$

in-medium pion field treated properly in chiral EFT



SUMMARY

- Chiral Effective Field Theory and Functional Renormalization Group:
 Nuclear Chiral Dynamics and Thermodynamics
 from symmetric to asymmetric nuclear matter and neutron (star) matter
 - Fluctuations beyond mean field include important multi-pion exchange mechanisms and low-energy nucleonic particle-hole excitations
 - Ist order phase transition: Fermi liquid interacting Fermi gas
- No indication of first-order chiral phase transition
 - Fluctuations work against early restoration of chiral symmetry
- New constraints from neutron stars for the equation-of-state of dense & cold baryonic matter:
 - Mass radius relation: stiff equation of state required ! No ultrahigh densities $(arrho_{f max}\sim {f 5}\,arrho_{f 0})$
 - Conventional (nucleon-meson, "non-exotic") EoS meets constraints lssue of strangeness: suppression of hyperons in neutron stars?)

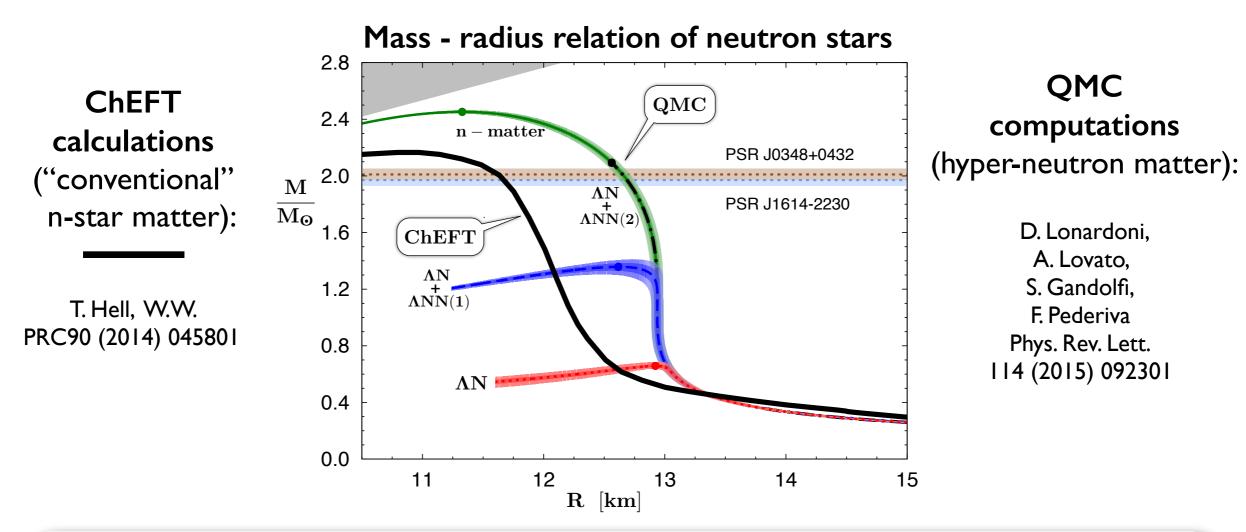


Appendix:

HYPERONS in NEUTRON STARS

NEUTRON STAR MATTER including HYPERONS

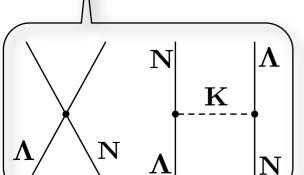
New Quantum Monte Carlo calculations using phenomenological hyperon-nucleon and hyperon-NN three-body interactions constrained by hypernuclei

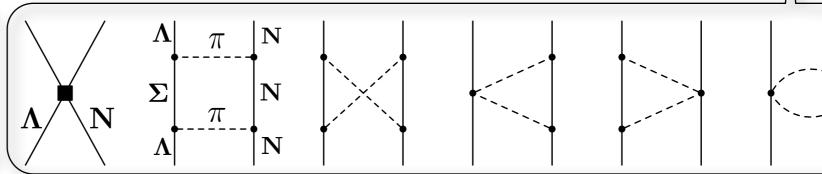


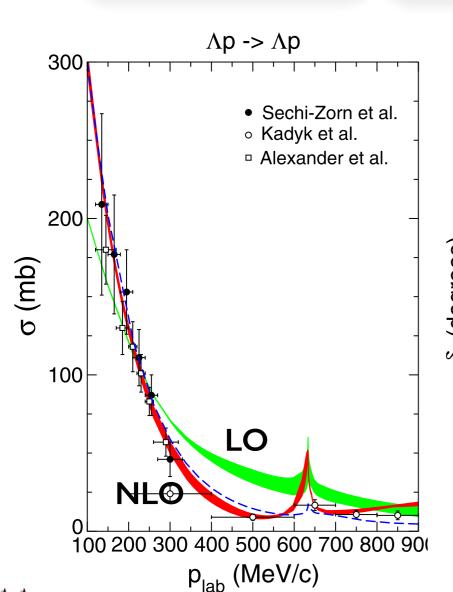
with inclusion of hyperons: EoS too soft to support 2-solar-mass star unless strong short-range repulsion in YN and / or YNN interactions

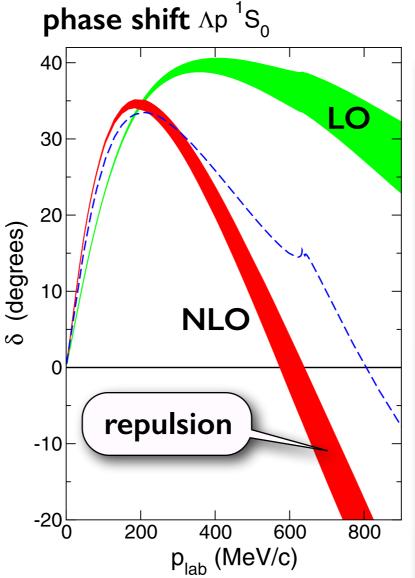


Hyperon - Nucleon Interaction from CHIRAL SU(3) Effective Field Theory









J. Haidenbauer, S. Petschauer, N. Kaiser, U.-G. Meißner, A. Nogga, W.W. Nucl. Phys. A 915 (2013) 24 moderate attraction at low momenta

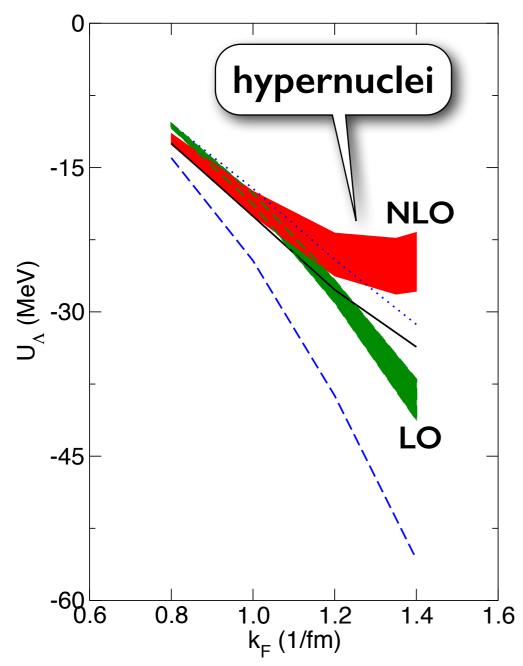
NLO

- relevant for hypernuclei
- strong repulsion
 at higher momenta
 - → relevant for dense baryonic matter

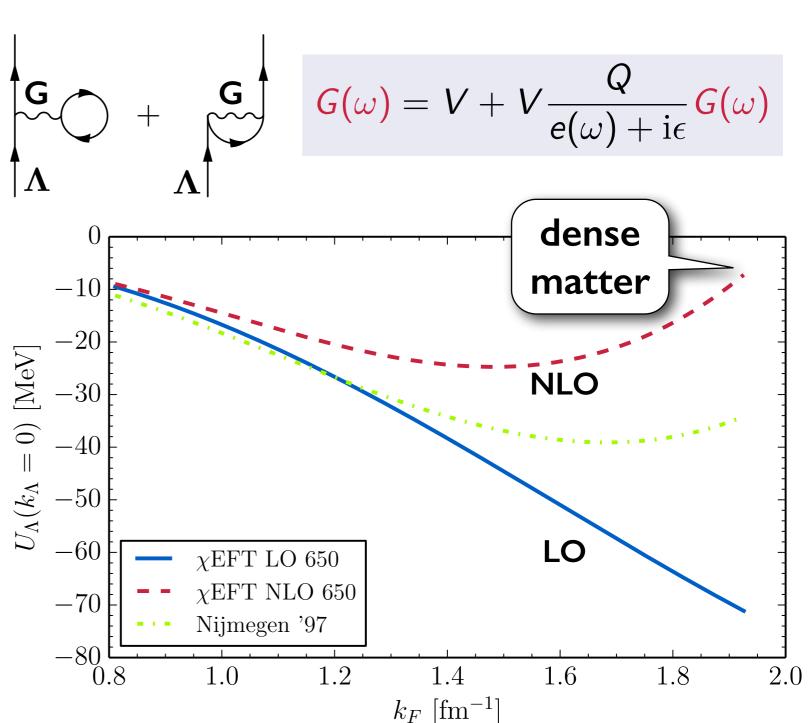


Density dependence of Λ single particle potential

Brueckner calculations using chiral SU(3) interaction



J. Haidenbauer, U.-G. Meißner, Nucl. Phys. A 936 (2015) 29

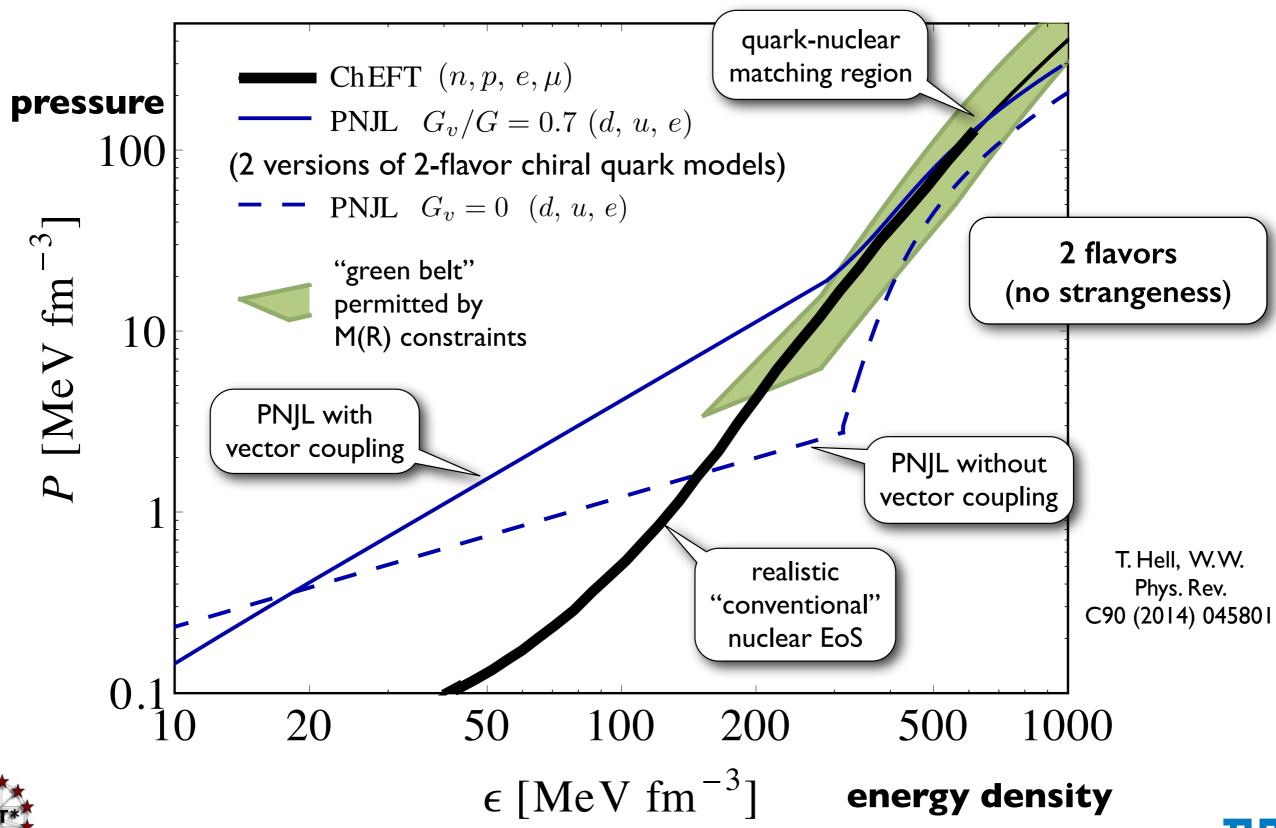


S. Petschauer, J. Haidenbauer, N. Kaiser, U.-G. Meißner, W.W. EPJA (2015); arXiv:1507.08808 [nucl-th]



Supplementary Materials

NEUTRON STAR Equation of State

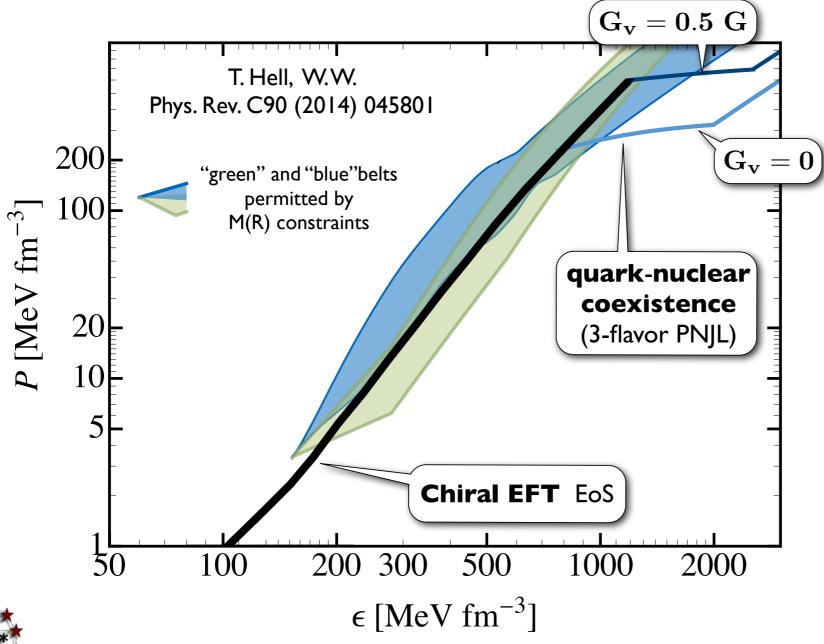




NEUTRON STAR MATTER

Equation of State

- In-medium Chiral Effective Field Theory up to 3 loops (reproducing thermodynamics of normal nuclear matter)
- 3-flavor PNJL model at high densities (incl. strange quarks)



- beta equilibrium $\mathbf{n} \leftrightarrow \mathbf{p} + e, \mu$
- charge conservation
- coexistence region:Gibbs conditions
- o quark-nuclear coexistence occurs (if at all) at baryon densities $ho > {f 5} \,\,
 ho_0$

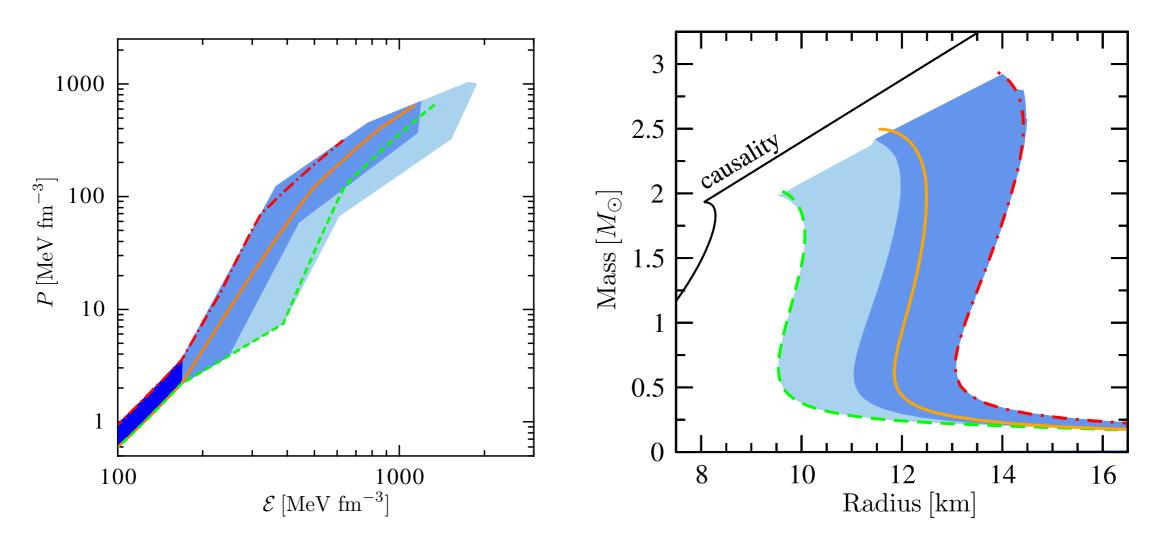




NEUTRON STAR MATTER

Equation of State

 Further independent analysis of new constraints from n-stars (Chiral EFT EoS and polytrope extrapolation)



K. Hebeler, J.M. Lattimer, C.J. Pethick, A. Schwenk

Astrophys. J. 773 (2013) 11

