

BB and $B\bar{B}$ four-quark systems from lattice QCD

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[arXiv:1505.00613](https://arxiv.org/abs/1505.00613)

[arXiv:1510.03441](https://arxiv.org/abs/1510.03441)

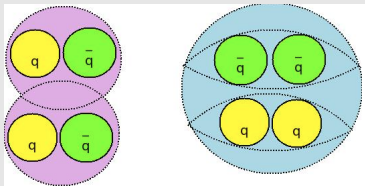


FAIRNESS 2016, Garmisch-Partenkirchen

Study of tetraquark states

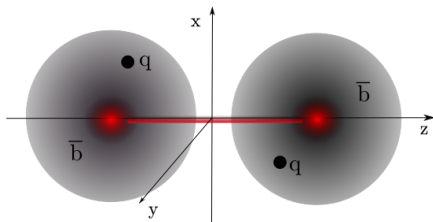
Motivation

- A number of mesons found in particle detectors (LHCb, Belle) are not well understood.
- E.g. charged charmonium and bottomonium states (Z_c^\pm and Z_b^\pm)
- They include bottomonium $b\bar{b}$ or charmonium $c\bar{c}$, but are also charged: must be 4-quark states



Approach

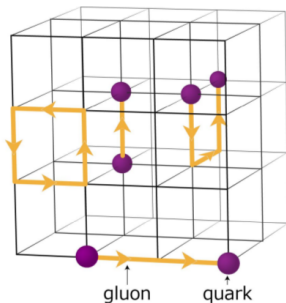
- Computation of 4-quark state very difficult
- If 2 quarks are heavy and 2 quarks are light: Treat degrees of freedom independently in two steps (Born-Oppenheimer approximation)
 - 1 Lattice computation of the potential of two static quarks in the presence of two light quarks



- 2 Solve Schrödinger's equation to check whether potentials are sufficiently attractive to form a bound state.

Lattice QCD

- QCD: Theory to describe quarks and gluons and the forces between them
- No analytical solution for low energy observables
- No perturbative approach on the potential
 \rightsquigarrow numerical solution (Lattice QCD)



[<http://www.jicfus.jp/en/promotion/pr/mj/guido-cossu/>, April 28, 2015]

Hadron spectroscopy I

- A hadron is completely described by its isospin (flavor content), total angular momentum J , parity P and charge conjugation C .
- Application of a suitable operator $\mathcal{O}(t)$ on the vacuum generates field excitations which are similar to the hadron of interest.
- Here: Use two static quarks and two quarks of finite mass
 - Static quarks: no spin, no contribution to total angular momentum and isospin
 - BB : $\bar{Q}\bar{Q}qq$ and $B\bar{B}$: $Q\bar{Q}q\bar{q}$ with $Q = b$ and $q \in \{u, d, s, c\}$
 - e.g. $B\bar{B}$ four-quark operator

$$\mathcal{O}_{B\bar{B}}(t) = \Gamma_{AB}\tilde{\Gamma}_{CD} \underbrace{\bar{Q}_C^a(\vec{x}, t)u_A^a(\vec{x}, t)}_{B \text{ meson at } \vec{x}} \underbrace{\bar{d}_B^b(\vec{y}, t)Q_D^b(\vec{y}, t)}_{\bar{B} \text{ meson at } \vec{y}}$$

- Obtain the correlation function in time for each separation r of the static quarks via

$$C(t, r) = \langle \Omega | \mathcal{O}^\dagger(t) \mathcal{O}(0) | \Omega \rangle$$

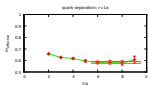
Requires $\mathcal{O}(\text{months})$ of computing time on high performance computers.

Hadron spectroscopy II

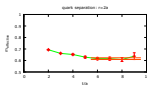
- For large t one finds the potential $V(r)$ of the hadronic state $\mathcal{O}(t)$:

$$\lim_{t \rightarrow \infty} \langle \Omega | \mathcal{O}^\dagger(t) \mathcal{O}(0) | \Omega \rangle \propto \exp(-V(r)t)$$

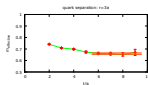
- Get a value of the potential for each quark separation r and obtain the complete potential:



$r=1a$

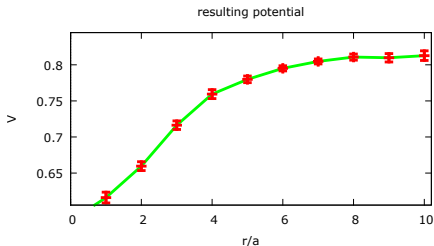


$r=2a$



$r=3a$

and so on...

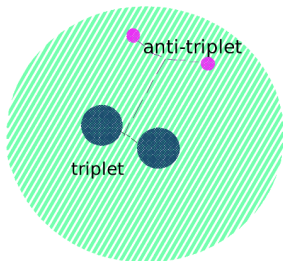


BB systems

The BB system - Expectations

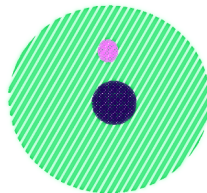
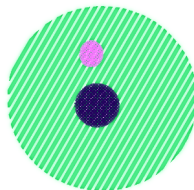
small separations of the static antiquarks:

- interaction due to 1-gluon exchange
- bound state: static $\bar{Q}\bar{Q}$ pair in a color triplet (attractive) \longrightarrow antiquark



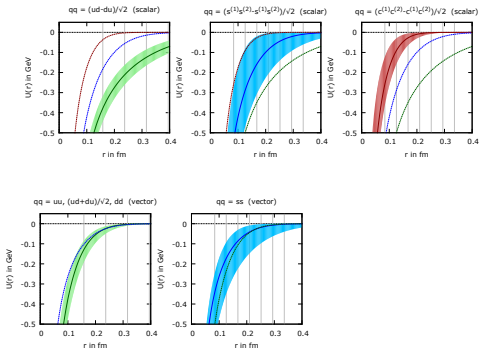
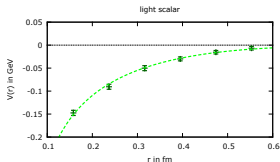
large separations of the static antiquarks:

- screening of the antiquark-antiquark interaction due to light quarks (stronger, the more massive the light quarks)
- basically 2 static-light mesons



Fit an ansatz to the potentials:

$$V(r) = - \underbrace{\frac{\alpha}{r}}_{\text{Coulomb-like}} \underbrace{e^{-\left(\frac{r}{a}\right)^2}}_{\text{colour screening}}$$



Observation: Potentials show more promise for binding the **less massive** the light quarks are! Strongest attraction for **scalar isosinglet** $qq = \frac{ud-du}{\sqrt{2}}$.

Solve Schrödinger's equation

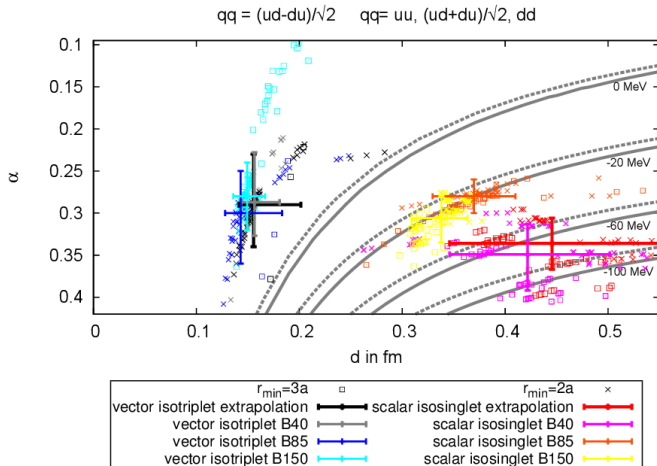
- notice: Born-Oppenheimer approximation is valid for $m_q \ll m_Q$
- solve Schrödinger's equation:

$$\left(-\frac{1}{2\mu} \frac{d^2}{dr^2} + V(r) \right) R(r) = E_B R(r) \quad , \quad \psi(r) = R(r)/r$$

- lowest eigenvalue $E_B < 0$ binding (4-quark bound state), $E_B > 0$ no binding (2 mesons)

Most promising channel for a bound state: scalar $ud\bar{b}\bar{b}$

Extrapolation to the physical quark mass: **Binding increases**



Summary BB

- BB systems are experimentally hard to observe, but theoretically easier to investigate
- Here: BB with light quarks $qq = ud$ with quantum numbers $I(J^P) = 0(1^+)$ (attractive channel)
- We find $E_B = -90_{-36}^{+43} \rightarrow$ **binding** with more than 2σ confidence level
- $qq = ss, cc$: no binding

$B\bar{B}$ systems

The $B\bar{B}$ system - Expectations

- For the experimentally most interesting case of isospin $|I_z| = 1$, $B\bar{B}$ has the same quantum numbers as $Q\bar{Q}$ and π .
- So there are **several states** to be taken into account:
 - four-quark state
 - 2 mesons
 - $Q\bar{Q} + \pi$

Consequently, one wonders...

- Is the potential contaminated by $Q\bar{Q} + \pi$?
- How to find a four-quark signal that is precise enough to be distinguishable from $Q\bar{Q} + \pi$?
- How to exclude "inadvertent" and unnecessary expensive computation of the $Q\bar{Q} + \pi$ instead of four-quark state?

Solution:

Build a 2×2 -matrix $C(t)$

- $C_{00}(t) = \langle \mathcal{O}_{B\bar{B}}^\dagger(t) \mathcal{O}_{B\bar{B}}(0) \rangle$
- $C_{01}(t) = \langle \mathcal{O}_{B\bar{B}}^\dagger(t) \mathcal{O}_{Q\bar{Q}+\pi}(0) \rangle$
- $C_{10}(t) = \langle \mathcal{O}_{Q\bar{Q}+\pi}^\dagger(t) \mathcal{O}_{B\bar{B}}(0) \rangle$
- $C_{11}(t) = \langle \mathcal{O}_{Q\bar{Q}+\pi}^\dagger(t) \mathcal{O}_{Q\bar{Q}+\pi}(0) \rangle$

and extract the **first excited state** with the Generalized Eigenvalue Problem (GEP).

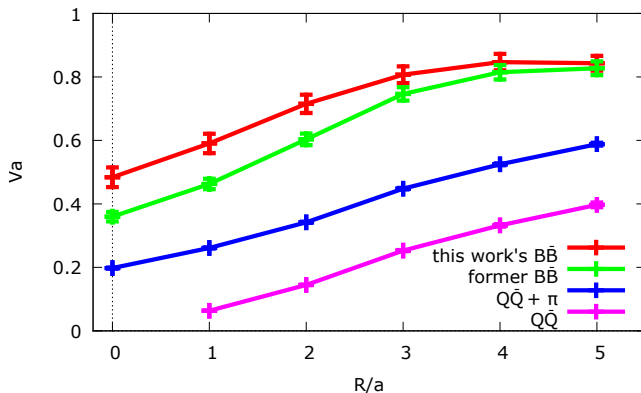
The matrix $C(t)$

$$C(t) = \begin{array}{|c|c|} \hline \begin{array}{|c|c|} \hline \text{Diagram 1} & \text{Diagram 2} \\ \hline \end{array} & \begin{array}{|c|c|} \hline \text{Diagram 3} & \text{Diagram 4} \\ \hline \end{array} \\ \hline \begin{array}{|c|c|} \hline \text{Diagram 5} & \text{Diagram 6} \\ \hline \end{array} & \begin{array}{|c|c|} \hline \text{Diagram 7} & \text{Diagram 8} \\ \hline \end{array} \\ \hline \end{array}$$

The diagram shows a 2x2 grid of lattice diagrams representing different elements of the matrix $C(t)$. Each diagram is enclosed in a square frame. The top-left diagram shows two separate, irregular shapes. The top-right diagram shows two shapes connected by a horizontal line. The bottom-left diagram shows two shapes connected by a vertical line. The bottom-right diagram shows a single irregular shape and a square.

- Derive different elements $C_{ij}(t)$ of the matrix analytically.
- Implement them.
- Identify the symmetries of $C_{ij}(t)$ (time-reversal, parity, charge conjugation, hermiticity and cubic rotations).
- Average according to the symmetries to increase statistics.

Preliminary results I



Potentials obtained

- the $Q\bar{Q}$ potential for comparison
- $Q\bar{Q} + \pi$: ground state
- $B\bar{B}$: includes contributions of $Q\bar{Q} + \pi$
- first excited state of the 2×2 matrix: free of contributions of $Q\bar{Q} + \pi$

Preliminary results II

Binding energy

A very preliminary analysis yields for quantum numbers $I(J^P) = 1(1^+)$:

$$E_B = (-170 \pm 100) \text{ MeV}$$

Summary $B\bar{B}$ system

- $B\bar{B}$ is experimentally more easy to investigate than BB .
- $B\bar{B}$ with light quarks $qq = u\bar{d}$ with quantum numbers $I(J^P) = 1(1^+)$ is a **tetraquark candidate**.
- Work in progress

Summary BB and $B\bar{B}$ systems

- BB systems with light quarks are able to form a **bound state**.
- $B\bar{B}$ systems are experimentally **more easy to access** than BB systems, but theoretically **more challenging**.
- Candidate for a binding $B\bar{B}$ state is currently investigated.