

# $Z_c(3900)/Z_c(3885)$ : what have we really seen?

[M. A., F. K. Guo, C. Hidalgo-Duque, J. Nieves, arXiv:1512.03638. Phys. Lett. B (to appear)]

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# Outline

- 1 Introduction
- 2 Formalism
- 3 Results
- 4 Improvements
- 5 Conclusions

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1 **Introduction**

2 Formalism

3 Results

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5 Conclusions

## Light sector before the heavy...

Take three examples from the light sector:

- ① Scalar–isoscalar sector in 1 – 2 GeV:  $f_0(1370)$ ,  $f_0(1500)$ ,  $f_0(1710)$  (and  $f_0(1790)$ )
- ② Pseudoscalar sector 1 – 2 GeV
- ③ Light scalar sector:  $f_0(500)$ ,  $f_0(980)$ ,  $a_0(980)$ ,  $\kappa$

### Scalar–isoscalar in 1 – 2 GeV

- Three candidates for the two isoscalars in the nonet:  $f_0(1370)$ ,  $f_0(1500)$ ,  $f_0(1710)$
- Again, they can be dynamically generated within the Unitary Chiral Approach ( $PP$  interactions + unitarity)
- One extra member makes it a candidate to be a glueball ( $f_0(1500)$  and  $f_0(1710)$ )
- (See the note of the PDG about the scalar sector)

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- ② Pseudoscalar sector 1 – 2 GeV
- ③ Light scalar sector:  $f_0(500), f_0(980), a_0(980), \kappa$

### Excited Pseudoscalars above 1 GeV

- $I = 1$   $\pi(1300), \pi(1800)$
- $I = 1/2$   $K(1460), K(1630)$
- $I = 0$

$\eta(1295), \eta(1405), \eta(1475)$   
 $\underbrace{\hspace{10em}}_{\eta(1440)?}$

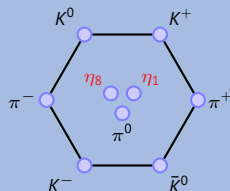
[Masoni *et al.*, JPG, 32('06)]

[Klempt, Zaitsev, PR, 454('07)]

[Wu *et al.*, PRL, 108('11)]

- $I = 0$   $X(1835)?$
- Exotic quantum numbers? Glueballs?

[Chanowitz, PRL, 46, 981('81)][Ishikawa, PRL, 46, 978('81)][Close, PRD, 55, 5749('97)]



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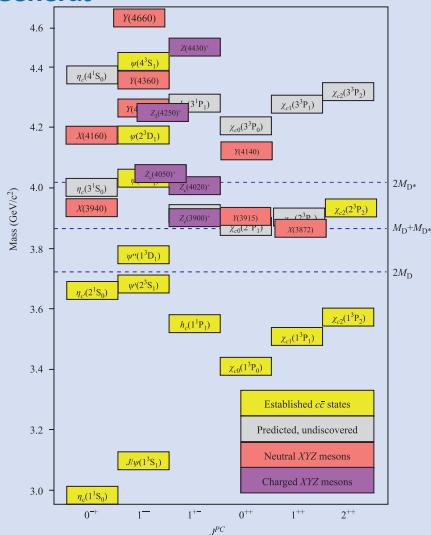
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### Light scalar sector

- A full nonet (no extra members):  $\sigma \equiv f_0(500), f_0(980), a_0(980), \kappa$
- They can be dynamically generated within the Unitary Chiral Approach ( $PP$  interactions + unitarity)
- Still, what are they?
  - MIT Bag Model predicts tetraquark  $0^{++}$  nonet, with masses  $m_\sigma \simeq 600$  MeV
  - For the  $\sigma$  meson,  $\sqrt{\langle r^2 \rangle_\sigma} = 0.44$  fm, compare with  $\sqrt{\langle r^2 \rangle_\pi} = 0.65$  fm.  
[M.A., J.A. Oller, *Phys.Rev. D* 86,034003(2012)]
  - The  $\sigma$  meson is very compact. Tetraquark?

## General

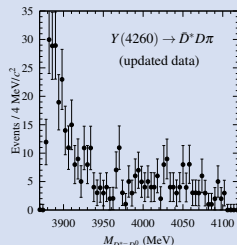
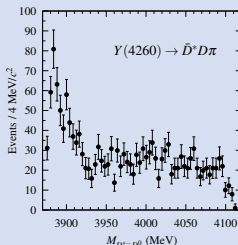
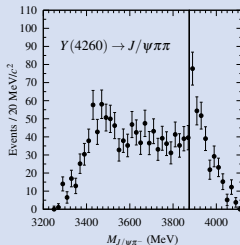


(Taken from: [Olsen, Front. Phys. **10**, 101401])

- Two recent reviews (2015):
  - [Olsen, Front. Phys. **10**, 101401]
  - [Chen *et al.*, arXiv:1601.02092]
- All the  $c\bar{c}$  states predicted by QM below  $D\bar{D}$  threshold have been found
- In 2003,  $X(3872)$  is discovered [Belle Collab., PRL,91,262001]
  - Very close to  $D^0\bar{D}^0$  threshold.
  - Close to (but lower)  $\chi_{c1}(2^3P_1)$ .
- **Lattice QCD:** [Prelovsek, Leskovec, PRL,111,192001] candidate for  $X(3872)$  only if  $c\bar{c} + D\bar{D}^*$  components are considered **together**

## Introduction: experimental information on $Z_c(3885)/Z_c(3900)$

- $Z_c(3900)$  first seen by **BESIII** and **Belle** Collabs. in  $J/\psi\pi^\pm$  invariant mass spectrum in  $e^+e^- \rightarrow Y(4260) \rightarrow J/\psi\pi^+\pi^-$   
[Phys. Rev. Lett. **110**, 252001 (2013), Phys. Rev. Lett. **110**, 252002 (2013)]
- Later on, **CLEO-c** data confirmed  $Z_c(3900)$  in  $e^+e^- \rightarrow \psi(4160) \rightarrow J/\psi\pi^+\pi^-$   
[Phys. Lett. B **727**, 366 (2013)]
- **BESIII** analyses  $e^+e^- \rightarrow Y(4260) \rightarrow \bar{D}^* D\pi$ , and sees  $Z_c(3885)$  in  $\bar{D}^* D$  invariant mass spectrum.  $J^P = 1^+$  favoured. [Phys. Rev. Lett. **112**, 022001 (2014)]
- **BESIII** confirms  $Z_c(3885)$  in  $\bar{D}^* D$  spectrum at different  $e^+e^-$  c.m. energies  
[Phys. Rev. D **92**, 092006 (2015)]
- If they are the same object, **Ratio:**  $\frac{\Gamma(Z_c \rightarrow D\bar{D}^*)}{\Gamma(Z_c \rightarrow J/\psi\pi)} = 6.2 \pm 2.9$





## Introduction: theoretical speculation

- “One of the most interesting resonances”: couples strongly to charmonium ( $\sim \bar{c}c$ ) and yet it has charge ( $\sim \bar{u}d$ ). Minimal quark constituent is **four** [ $\bar{c}c\bar{u}d$ ].
- Many different interpretations have been given:
  - [Olsen, *Front. Phys.* **10**, 101401][Chen *et al.*, arXiv:1601.02092]
    - Tetraquark
    - $\bar{D}^*D$  molecular state
    - Simply a kinematical effect
    - Hadrocharmonium
    - It has also been searched for in lattice QCD



**What is still missing?**

A **joint study of both reactions** in which the  $Z_c$  structure has been seen

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## Coupling $\bar{D}^*D$ and $J/\psi\pi$ channels

**Coupled channel** formalism is needed, because  $Z_c(3900)$ :

- is expected to be dynamically generated in  $\bar{D}^*D$  channel (#2),
- but it is also seen in  $J/\psi\pi$  channel (#1).

$$T = (\mathbb{I} - V \cdot G)^{-1} \cdot V ,$$

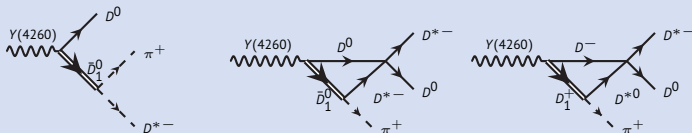
$$V_{ij} = 4\sqrt{m_{i1}m_{i2}}\sqrt{m_{j1}m_{j2}} e^{-q_i^2/\Lambda_i^2} e^{-q_j^2/\Lambda_j^2} C_{ij} ,$$

- $G(E)$  are loop functions (Regularized with standard gaussian regulator)
- $J/\psi\pi \rightarrow J/\psi\pi$ : known to be tiny,  $C_{11} = 0$ .
- $\bar{D}^*D \rightarrow J/\psi\pi$ : we make the simplest possible assumption,  $C_{12} \equiv \tilde{C}$  (constant)
- $\bar{D}^*D \rightarrow \bar{D}^*D$ : In a momentum expansion (HQSS), simply a constant,  $C_{22} \equiv C_{1Z}$ .
- **Problem**: no resonance in the complex plane above threshold with only constant potentials (even with coupled channels).
- We introduce some energy dependence,

$$C_{22}(E) = C_{1Z} + b(E - m_D - m_{D^*}) .$$

Amplitudes:  $Y(4260) \rightarrow (J/\psi\pi^-)\pi^+, (D^{*-}D^0)\pi^+$

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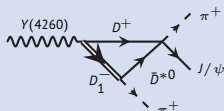
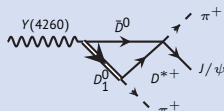
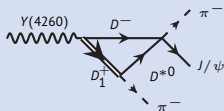
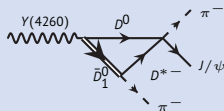
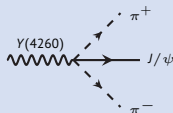


$$|\overline{\mathcal{M}}_2(s, t)|^2 = \left| \frac{1}{t - m_{D_1}^2} + I_3(s) T_{22}(s) \right|^2 q_\pi^4(s) + |\beta (1 + T_{22}(s) G_{22}(s))|^2$$

- $s$  (Mandelstam)  $\bar{D}^* D$  invariant mass squared
- $I_3(s)$ : three meson loop propagator
- $\bar{D}^* D$  rescattering enters through  $T_{22}(s)$
- $q_\pi^2(s) = \lambda(M_Y^2, s, m_\pi^2)/(4M_Y^2)$

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- The decay proceeds mainly through  $[T_{12}(s)]$   
 $Y \rightarrow (\bar{D}^* D) \pi \rightarrow (J/\psi \pi) \pi$
- Some direct production included through  $\alpha$
- $s, t$  (Mandelstam)  $J/\psi \pi^-, J/\psi \pi^+$  invariant mass squared

$$|\overline{\mathcal{M}}_1(s, t)|^2 = |\tau(s)|^2 q_\pi^4(s) + |\tau(t)|^2 q_\pi^4(t) + \frac{3 \cos^2 \theta - 1}{4} \left( \tau(s) \tau(t)^* + \tau(s)^* \tau(t) \right) q_\pi^2(s) q_\pi^2(t),$$

$$\tau(s) = \sqrt{2} I_3(s) T_{12}(s) + \alpha$$

## Events distributions and Experimental data

- Events distributions  $\mathcal{N}_i$ :

$$\mathcal{N}_i(s) = K_i (\mathcal{A}_i(s) + \mathcal{B}_i(s))$$

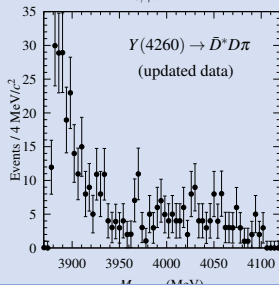
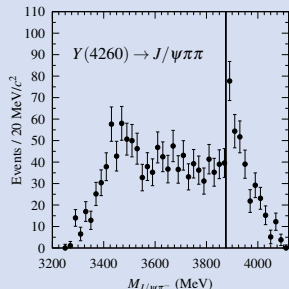
$$\mathcal{A}_i(s) = \int_{t_{i,-}}^{t_{i,+}} dt |\overline{\mathcal{M}}_i(s, t)|^2$$

- $K_i$  (unknown) global normalization constants
- $B_i$  are background functions (parametrized as in the experimental analyses)
- “Branching ratio”:

$$R_{\text{exp}} = \frac{\Gamma(Z_c \rightarrow D\bar{D}^*)}{\Gamma(Z_c \rightarrow J/\psi\pi)} = 6.2 \pm 2.9$$

- Theoretically estimated as the (physical) ratio of areas around  $Z_c$ (3900) mass

$$R_{\text{th}} = \frac{\int ds A_2(s)}{\int ds A_1(s)}$$



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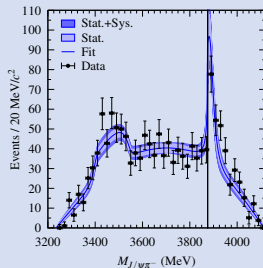
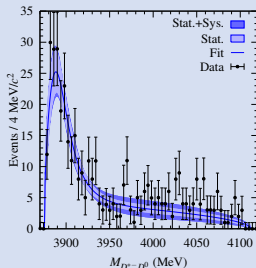
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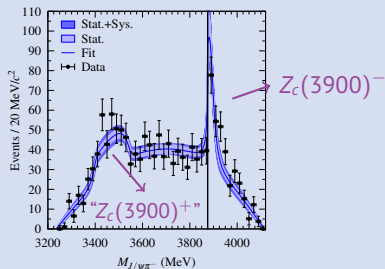
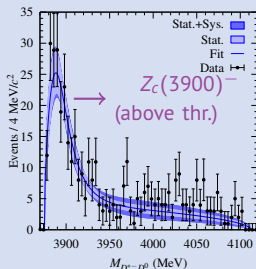
## Results: comparison with experiment(s)



$\Lambda_2$ (GeV)	$C_{1Z}$ (fm <sup>2</sup> )	$b$ (fm <sup>3</sup> )	$\tilde{C}$ (fm <sup>2</sup> )	$\chi^2/\text{dof}$	$R_{\text{th}}$
1.0	$-0.19 \pm 0.08 \pm 0.01$	$-2.0 \pm 0.7 \pm 0.4$	$0.39 \pm 0.10 \pm 0.02$	1.02	$6.0 \pm 3.5 \pm 0.5$
0.5	$+0.01 \pm 0.21 \pm 0.03$	$-7.0 \pm 0.4 \pm 1.4$	$0.64 \pm 0.16 \pm 0.02$	1.09	$6.5 \pm 3.6 \pm 0.2$
1.0	$-0.27 \pm 0.08 \pm 0.07$	0 (fixed)	$0.34 \pm 0.14 \pm 0.01$	1.31	$10.3 \pm 9.0 \pm 1.1$
0.5	$-0.27 \pm 0.16 \pm 0.13$	0 (fixed)	$0.54 \pm 0.16 \pm 0.02$	1.36	$10.9 \pm 9.0 \pm 2.5$

- Four different fits:  $b = \{\text{free}, 0\}$ ,  $\Lambda_2 = \{0.5, 1.0\}$  GeV
- Only the  $T$ -matrix parameters are shown (not shown: normalization, ...)
- All fits have  $\hat{\chi}^2 \simeq 1$  ( $\simeq 1.4$  for  $b = 0$ ), and are within the error band of the best one
- Reproduction of the data is **excellent**

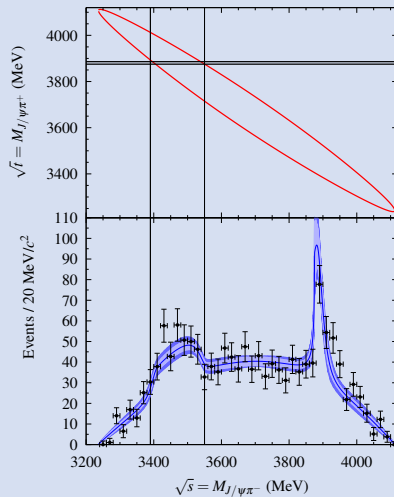
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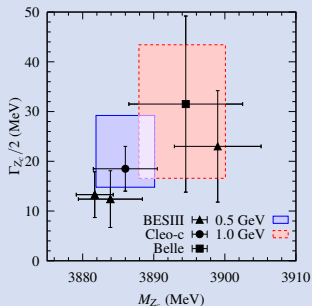
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## Reflection of threshold and $Z_c(3900)$



## Results: Spectroscopy



$M_{Z_c}$ (MeV)	$\Gamma_{Z_c}/2$ (MeV)	Ref.	Final state
$3899 \pm 6$	$23 \pm 11$	▲(BESIII)	$J/\psi \pi$
$3895 \pm 8$	$32 \pm 18$	■(Belle)	$J/\psi \pi$
$3886 \pm 5$	$19 \pm 5$	●(CLEO-c)	$J/\psi \pi$
$3884 \pm 5$	$12 \pm 6$	▲(BESIII)	$\bar{D}^* D$
$3882 \pm 3$	$13 \pm 5$	▲(BESIII)	$\bar{D}^* D$
$3894 \pm 6 \pm 1$	$30 \pm 12 \pm 6$	■( $\Lambda = 1.0$ GeV)	both
$3886 \pm 4 \pm 1$	$22 \pm 6 \pm 4$	■( $\Lambda = 0.5$ GeV)	both
$3831 \pm 26^{+7}_{-28}$	virtual state	( $\Lambda = 1.0$ GeV)	both
$3844 \pm 19^{+12}_{-21}$	virtual state	( $\Lambda = 0.5$ GeV)	both

Two different scenarios:

- $(b \neq 0)$   $Z_c$  is a  $\bar{D}^* D$  **resonance** very close to threshold  
(Differences with experiments are related to Breit-Wigner parametrizations)
- $(b = 0)$   $Z_c$  is a **virtual state**

In both scenarios,

- Data are very well reproduced
- A single structure (not two)  $Z_c(3885)/Z_c(3900)$  is needed

## Bound state, resonance, virtual ...

### Well known example: $NN$ scattering and the deuteron

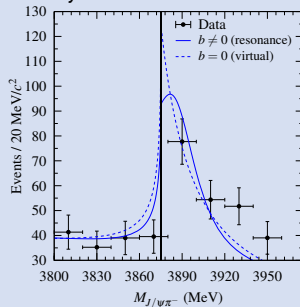
Triplet ( ${}^3S_1 - {}^3D_1$ ):

- $a_t \simeq 5$  fm.
- In this wave there is a **bound state**. The deuteron is a well known, really physical particle.

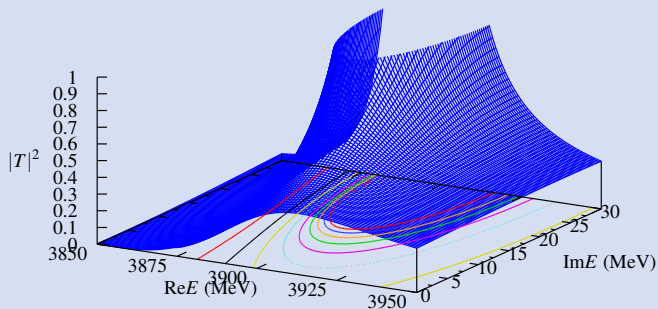
Singlet ( ${}^1S_0$ ):

- $a_s \simeq -24$  fm.
- In this wave there is a **virtual state**.

- A **virtual state** does not correspond to a real particle. (Wavefunction not localized.)
- It produces effects at the threshold similar to those of a bound state or a nearby resonance.

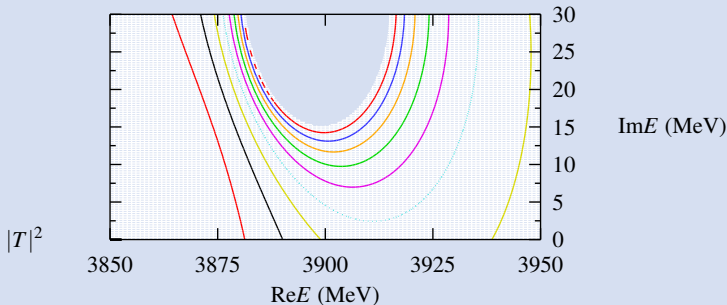


## Complex plane & poles: First scenario (resonance)



- Pole located at  $3894 - i30$  MeV
- Plot: unphysical Riemann sheet connected to the physical one above  $D^*\bar{D}$
- Shift of the pole towards higher energies (interference!)

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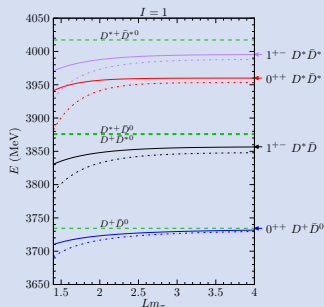
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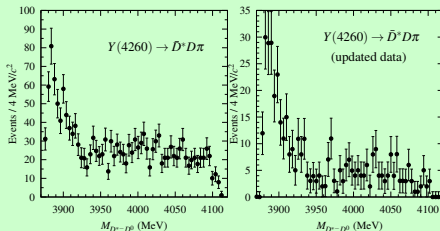
## $Z_c(3900)$ on the lattice

- Two recent works:
  - [Prelovsek *et al*, Phys.Rev.D91,014504(2015)] ( $m_\pi = 266$  MeV)  
“no additional candidate”
  - [HAL QCD, arXiv:1602.03465] ( $m_\pi \geq 410$  MeV)  
Virtual poles with very low masses and deep in the complex plane.
- Results are not conclusive (large pion masses, etc...)
- We can predict energy levels in a finite box. It might be helpful to understand these (and future) lattice simulations



# Experimental improvements for $Z_c(3900)$

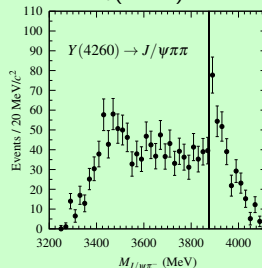
## New data on $\bar{D}^* D$ from BESIII



- Better statistics
- Lower “background”
- double  $D$ -tag

## New data on $J/\psi \pi$

The spectrum of  $J/\psi \pi$  with **narrower bins** is highly desirable in order to better elaborate on the nature of  $Z_c(3900)$



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## Conclusions (this work)

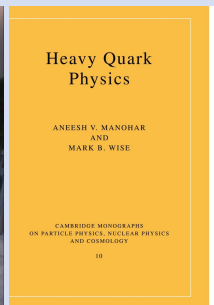
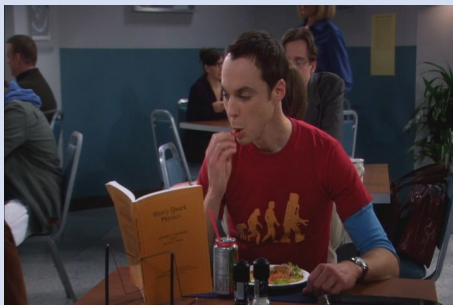
- $Z_c(3900)$  is a most-interesting, exotic, structure. A candidate for “tetraquark”
- We have presented the first simultaneous study of the two decays  $(Y(4260) \rightarrow J/\psi\pi\pi, \bar{D}^*D\pi)$  in which  $Z_c(3900)$  is seen
- Data are well reproduced in all fits ( $\hat{\chi}^2 \simeq 1$ )
- Two different scenarios are found:
  - ① ( $b \neq 0$ )  $Z_c(3900)$  is a  $\bar{D}^*D$  resonance
  - ② ( $b = 0$ )  $Z_c(3900)$  is a virtual state
- In any case, a single structure for  $Z_c(3885)/Z_c(3900)$  is needed.
- Improved data on  $J/\psi\pi$  invariant mass spectrum are necessary

## Conclusions (general)

- Charmonium spectrum, well known below  $D\bar{D}$  threshold.
- Since 2003, the charmonium(-like) spectrum increases continuously ( $\simeq 1$  state/year), but we do not fully understand: there are  $c\bar{c}$ , there are meson-meson **molecules**, there are **tetraquarks**, and many others.
- They must be **mixing**, specially around thresholds.
- Lattice still must go down to physical masses.
- We shall all be studying **Heavy Quark Physics**...

## Conclusions (general)

- Charmonium spectrum, well known below  $D\bar{D}$  threshold.
- Since 2003, the charmonium(-like) spectrum increases continuously ( $\simeq 1$  state/year), but we do not fully understand: there are  $c\bar{c}$ , there are meson-meson **molecules**, there are **tetraquarks**, and many others.
- They must be **mixing**, specially around thresholds.
- Lattice still must go down to physical masses.
- We shall all be studying **Heavy Quark Physics**...



# $Z_c(3900)/Z_c(3885)$ : what have we really seen?

[M. A., F. K. Guo, C. Hidalgo-Duque, J. Nieves, arXiv:1512.03638. Phys. Lett. B (to appear)]







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**Thanks for your attention**



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