



Quarkonium and Quark Gluon Plasma: from SPS to LHC energies

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INFN Torino

GSI Colloquium, May 26th 2015

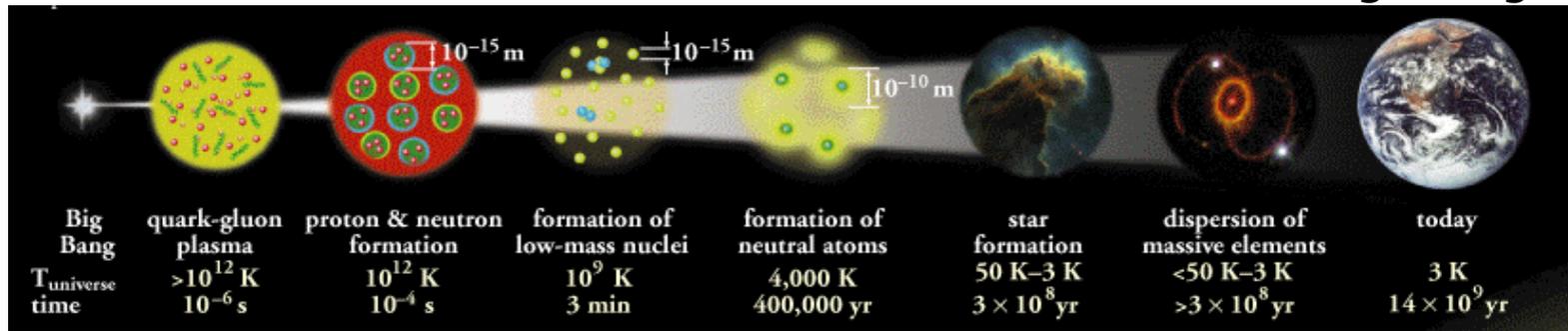


Outlook

- The Quark Gluon Plasma: a very short introduction!
- Why a qq bound state (quarkonium) is one of the most important probes of Quark Gluon Plasma formation?
- Quarkonium studies in heavy ion collisions: from low energy experiments to LHC

HEAVY IONS AND QUARK GLUON PLASMA

- ➔ Quark Gluon Plasma is a state of strongly interacting matter in which quarks and gluons are no more confined into hadrons
- ➔ QGP is formed at high temperatures and/or density → conditions similar to those achieved few micro-seconds after the Big-Bang

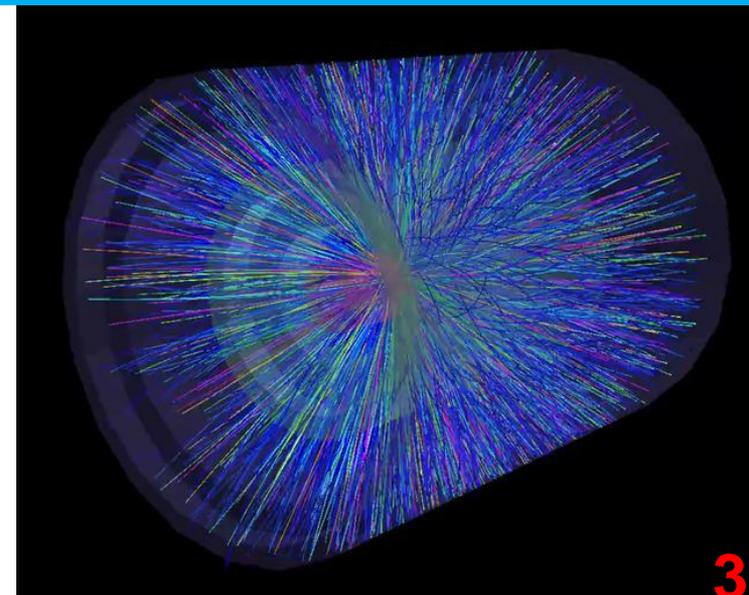


➔ How can QGP be produced in laboratory?

heavy-ion collisions

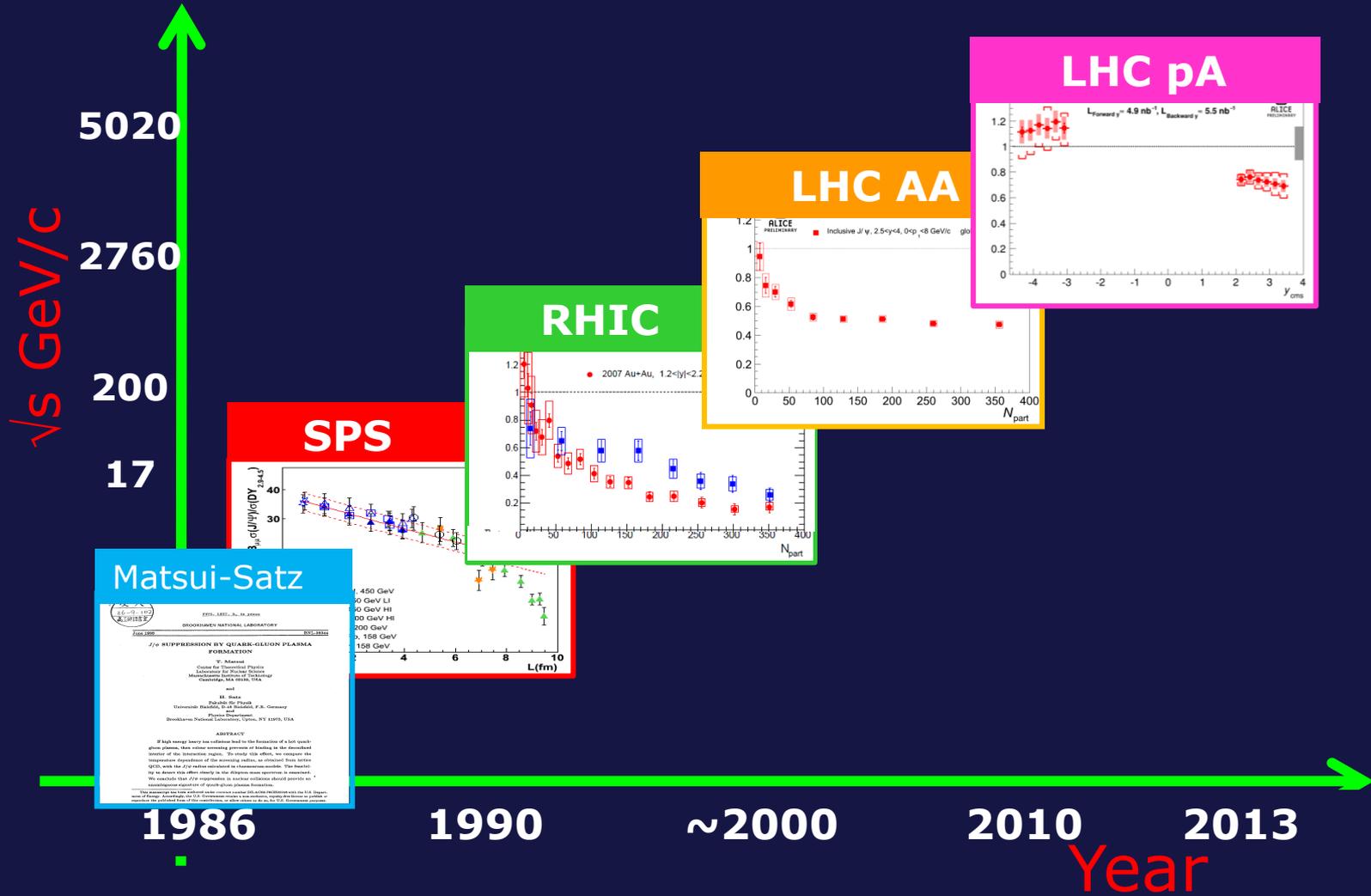
➔ How to understand the properties of the created hot medium?

study specific probes as jets, open heavy flavors, quarkonium...



HISTORY OF HEAVY-ION QUARKONIUM STUDIES

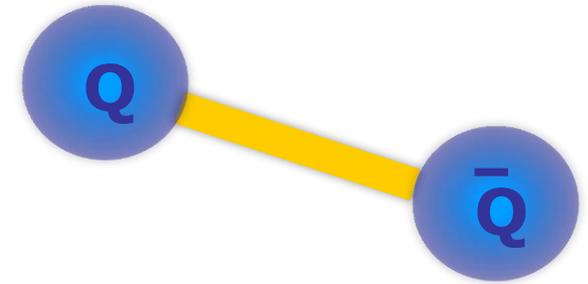
➔ The experimental investigation of quarkonium as one of the most striking signatures for QGP formation is a **30 years long story!**



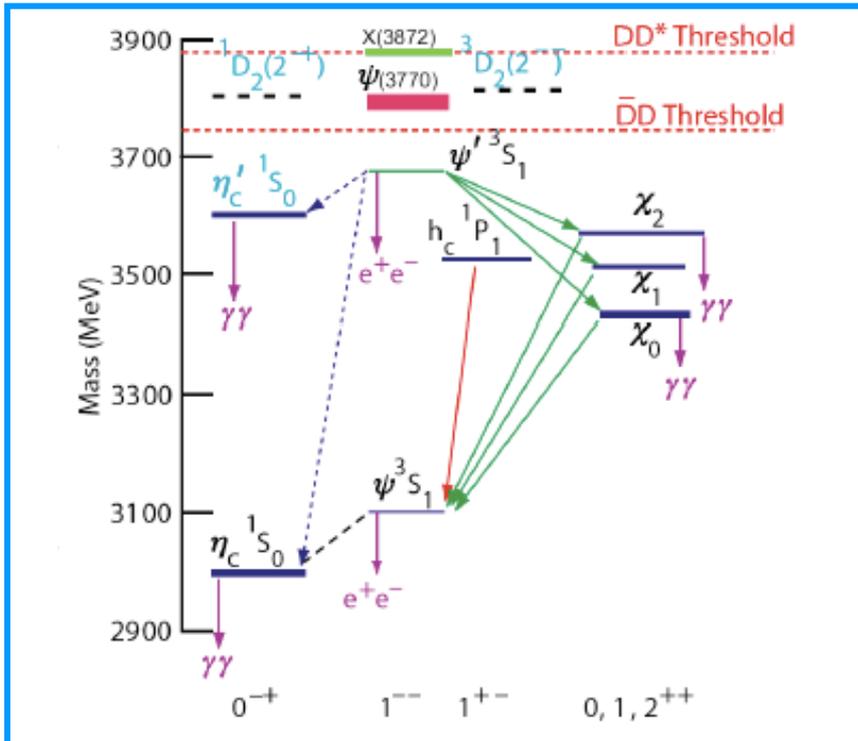
WHAT IS QUARKONIUM?

➔ Quarkonium is a bound state of Q and \bar{Q} with $m_{QQ} < 2m_D(m_B)$

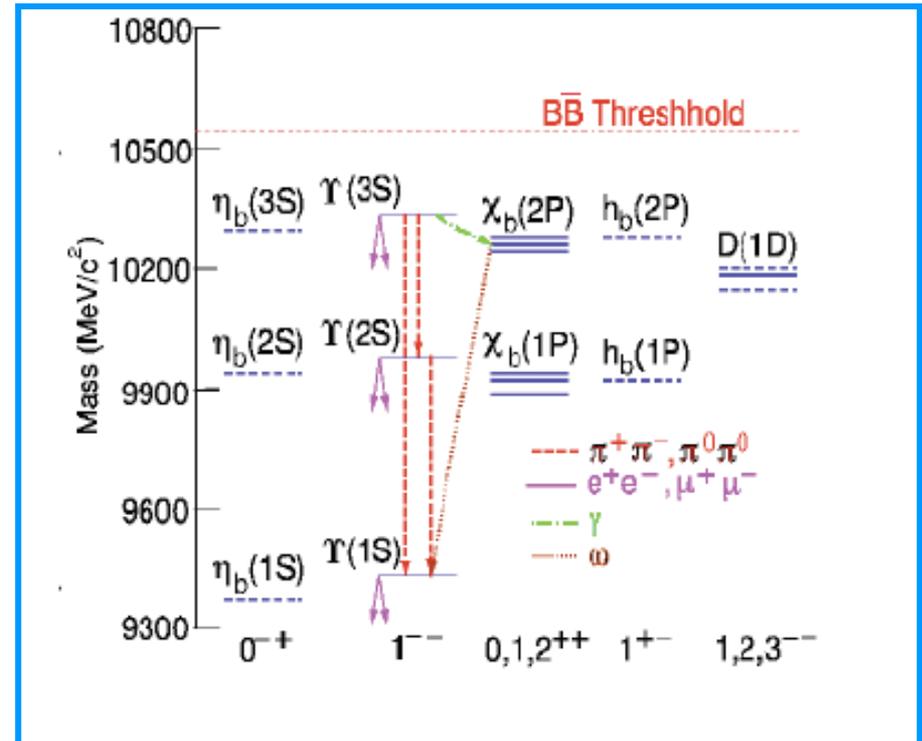
➔ Several quarkonium states exist, characterized by different quantum numbers



Charmonium ($c\bar{c}$) family



Bottomonium ($b\bar{b}$) family



QUARKONIUM AT T=0

→ At T=0, the binding of the Q and \bar{Q} quarks can be expressed using the Cornell potential:

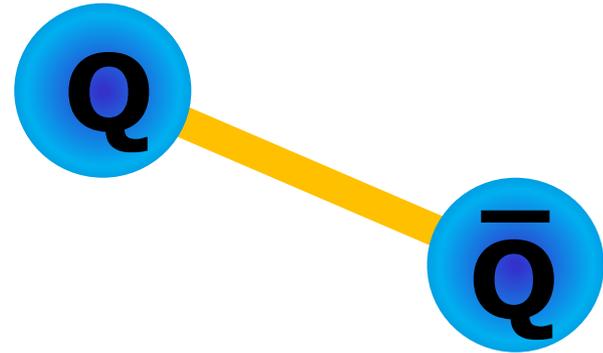
$$V(r) = -\frac{\alpha}{r} + kr$$



coulombian contribution,
induced by a g exchange
between Q and \bar{Q}



confinement
term



QUARKONIUM IN A QGP

➔ What happens to a $q\bar{q}$ pair placed in the QGP?

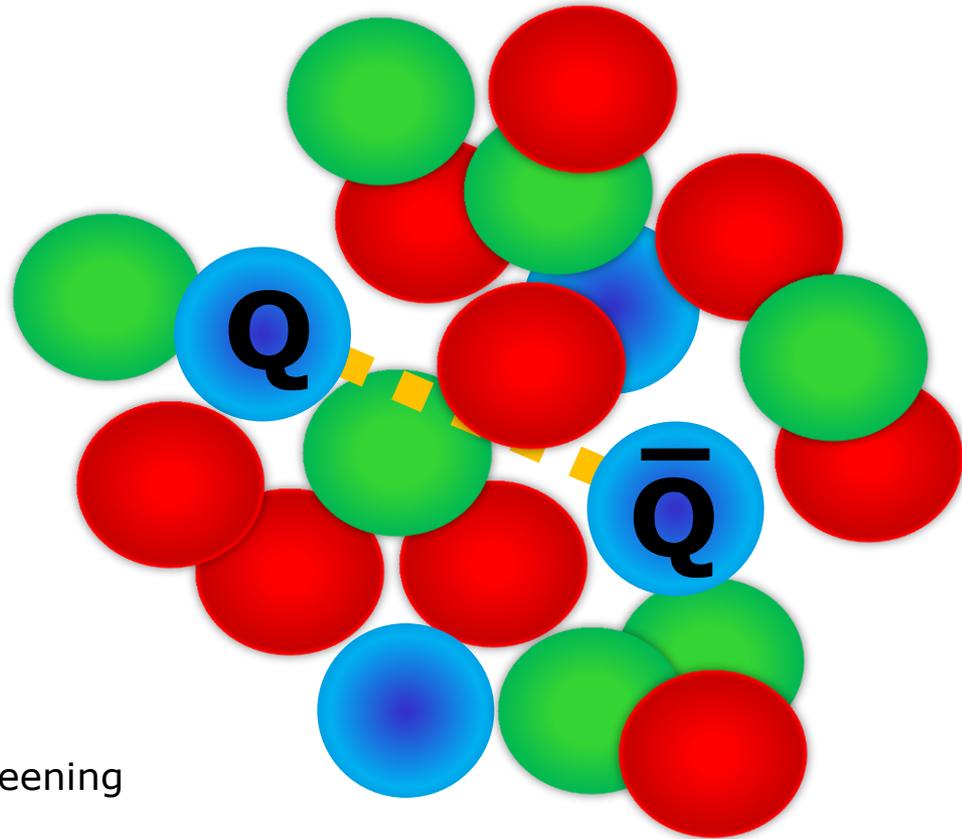
The QGP consists of deconfined colour charges

$$V(r) = -\frac{\alpha}{r} + kr$$



$$V(r) = -\frac{\alpha}{r} e^{-r/\lambda_D}$$

λ_D : screening radius

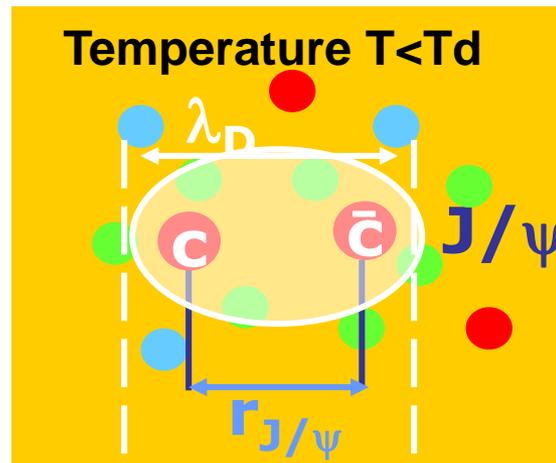
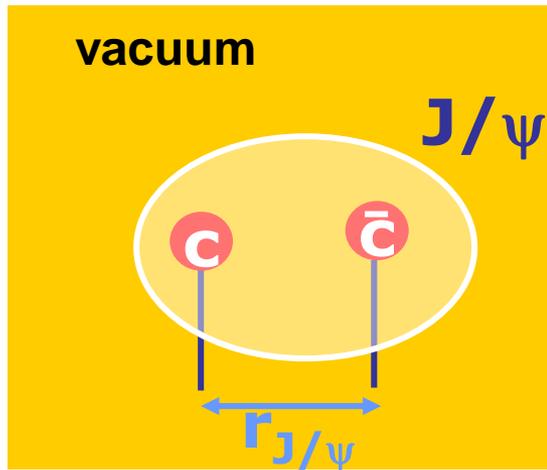


➔ the binding of a $q\bar{q}$ pair is subject to the effects of colour screening:

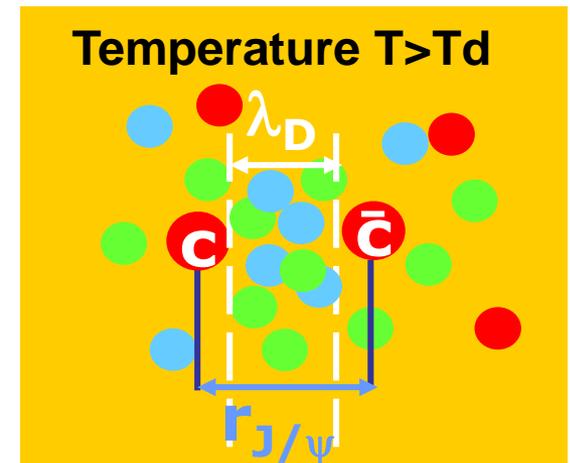
- the “confinement” contribution disappears
- the coulombian term of the potential is screened by the high color density

DEBYE SCREENING

➔ The screening radius $\lambda_D(T)$ (i.e. the maximum distance which allows the formation of a bound $Q\bar{Q}$ pair) decreases with the temperature T



if resonance radius
 $< \lambda_D(T)$
→ resonance can be
formed



if resonance radius
 $> \lambda_D(T)$
→ no resonance can
be formed

At a given T :

QUARKONIUM SUPPRESSION



PHYS. LETT. B, in press

BROOKHAVEN NATIONAL LABORATORY

June 1986

BNL-38344

J/ψ SUPPRESSION BY QUARK-GLUON PLASMA FORMATION

T. Matsui

Center for Theoretical Physics
Laboratory for Nuclear Science
Massachusetts Institute of Technology
Cambridge, MA 02139, USA

and

H. Satz

Fakultät für Physik
Universität Bielefeld, D-48 Bielefeld, F.R. Germany
and
Physics Department
Brookhaven National Laboratory, Upton, NY 11973, USA

ABSTRACT

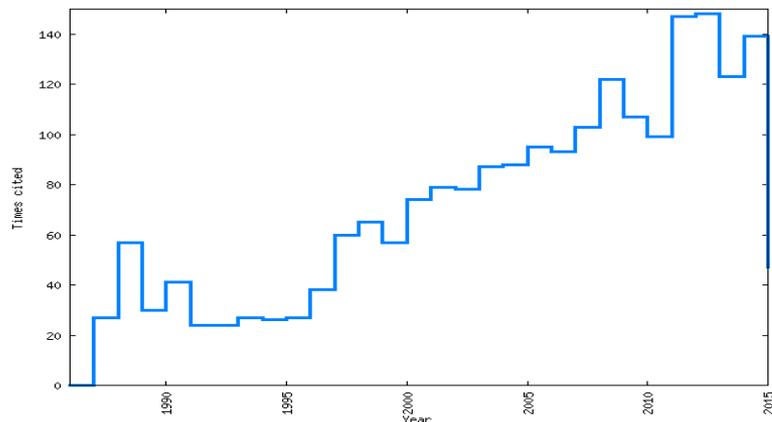
If high energy heavy ion collisions lead to the formation of a hot quark-gluon plasma, then colour screening prevents $c\bar{c}$ binding in the deconfined interior of the interaction region. To study this effect, we compare the temperature dependence of the screening radius, as obtained from lattice QCD, with the J/ψ radius calculated in charmonium models. The feasibility to detect this effect clearly in the dilepton mass spectrum is examined. We conclude that J/ψ suppression in nuclear collisions should provide an unambiguous signature of quark-gluon plasma formation.

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→ This is the idea behind the suggestion (by Matsui and Satz) of the J/ψ as a signature of QGP formation (~30 years ago!)



→ very famous paper cited ~2150 times!



SEQUENTIAL SUPPRESSION

- Strongly bound states have smaller sizes
- Debye screening condition $r_0 > \lambda_D$ will occur at different T



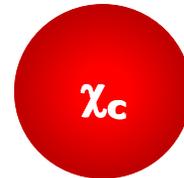
Differences in the binding energies of the quarkonium states lead to a sequential melting of the states with increasing temperature



thermometer of the initial QGP temperature

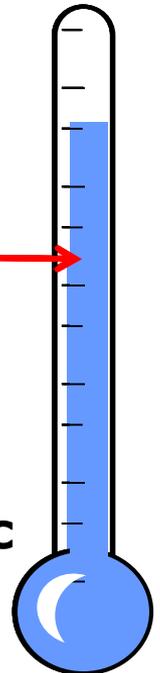
state	J/ψ	χ _c	ψ(2S)
Mass(GeV)	3.10	3.51	3.69
ΔE (GeV)	0.64	0.22	0.05
r ₀ (fm)	0.50	0.72	0.90

state	Υ(1S)	Υ(2S)	Υ(3S)
Mass(GeV)	9.46	10.0	10.36
ΔE (GeV)	1.10	0.54	0.20
r ₀ (fm)	0.28	0.56	0.78



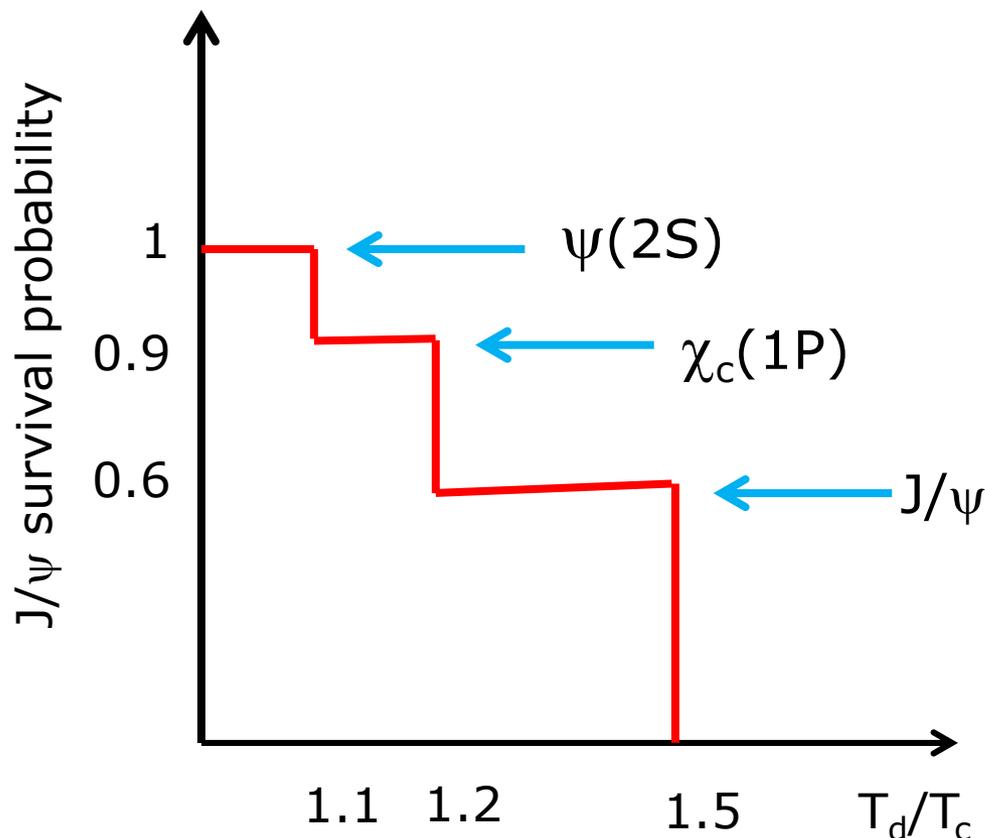
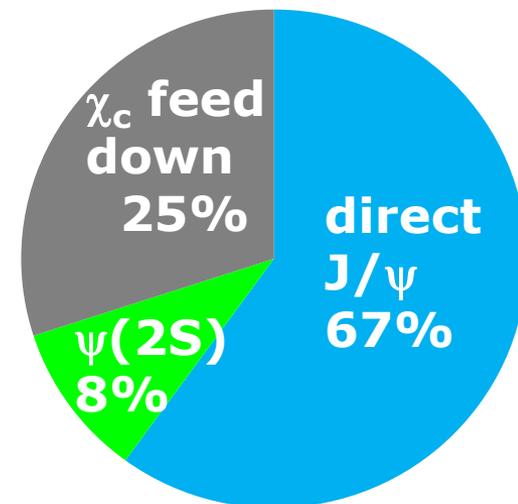
T_c →

T_c < T < T_c



SUPPRESSION PATTERN

- ➔ **Feed-down process:** charmonium “ground state” resonances can be produced through decay of larger mass quarkonia (J/ψ production from B decays neglected)
- ➔ Effect : ~30-40% for J/ψ, ~50% for Υ(1S)



- ➔ Dissociations temperatures are computed with lattice QCD or potential models calculations

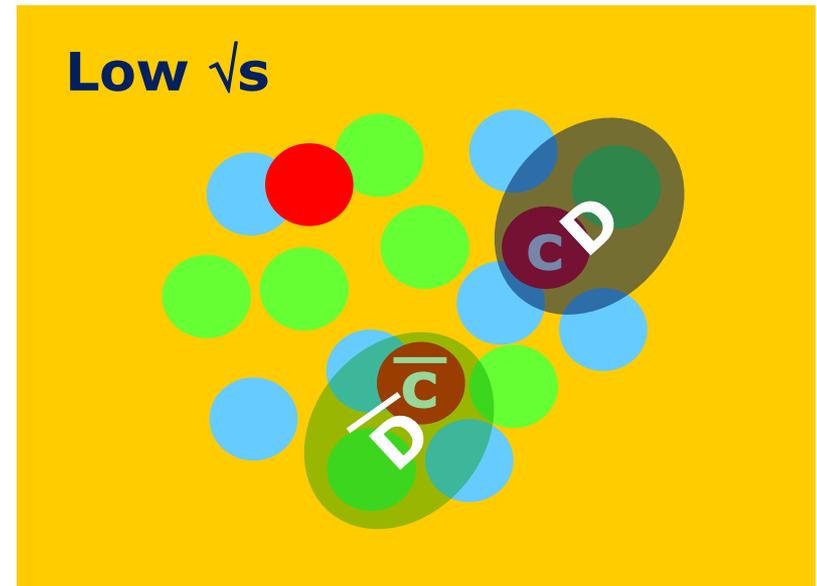
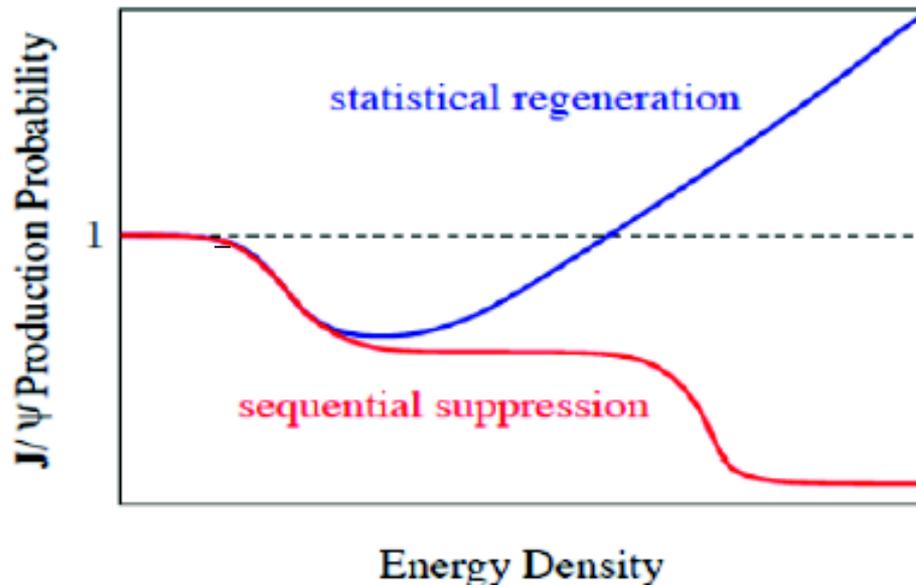
	J/ψ	χ _c (1P)	ψ(2S)
T_d/T_c	~1.5	~1.1	~1.1

H.Satz, arXiv:1310.1209

FROM SUPPRESSION... TO (RE)COMBINATION

➔ Increasing the energy of the collision the cc pair multiplicity increases

In most central AA collisions	SPS 20 GeV	RHIC 200GeV	LHC 2.76TeV
$N_{c\bar{c}}/\text{event}$	~0.2	~10	~75

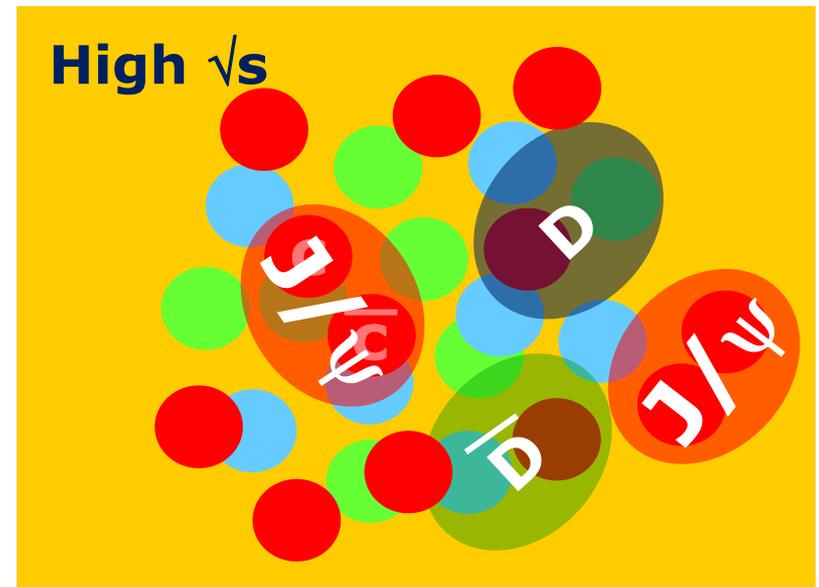
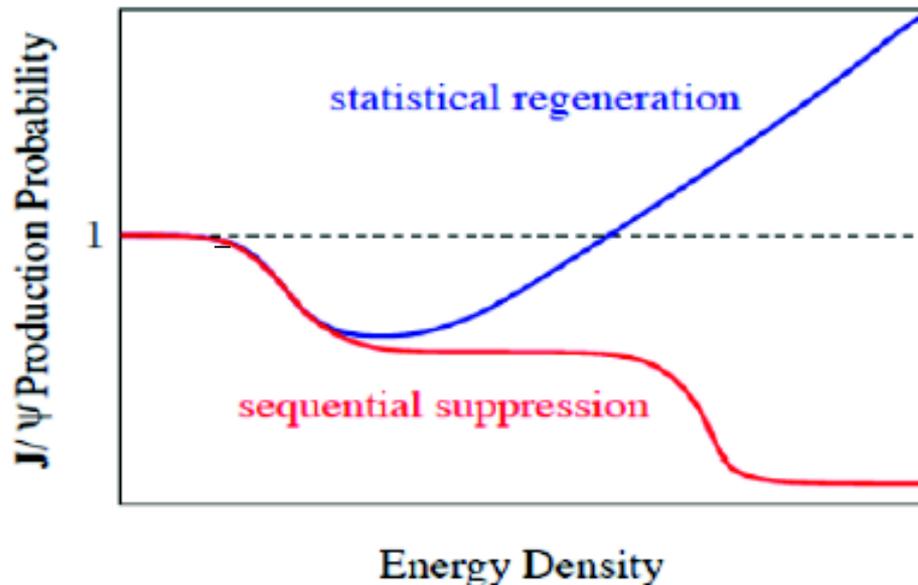


➔ An enhancement via (re)combination of cc pairs producing quarkonia can take place at hadronization or during QGP stage

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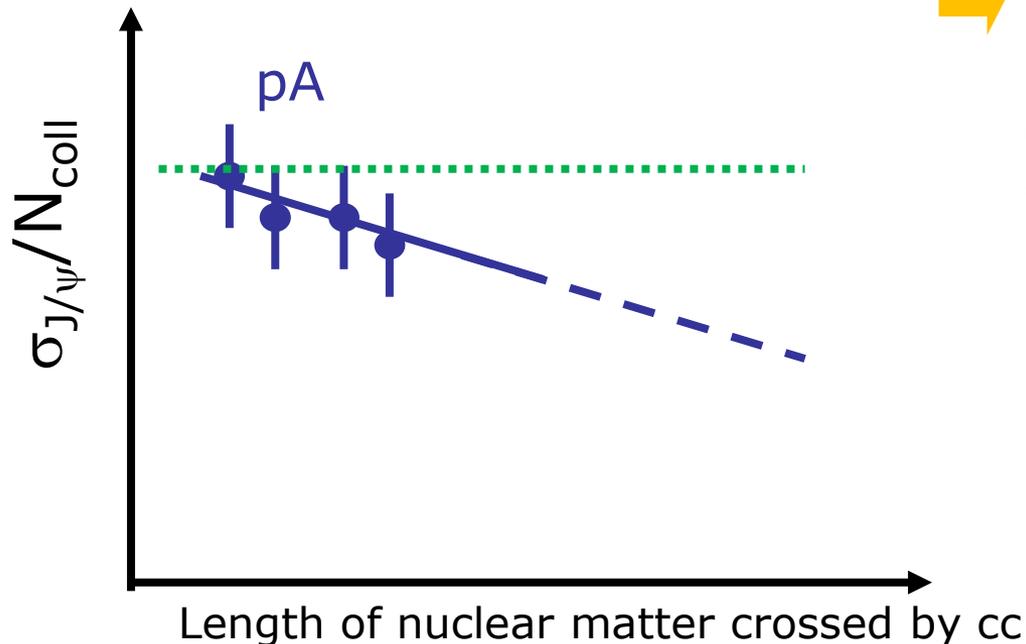
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WHAT ABOUT COLD NUCLEAR MATTER?

➔ What is the fate of a J/ψ placed in a cold nuclear medium?

➔ in principle, no J/ψ suppression

➔ however a reduction of the yield per nucleon-nucleon collisions is observed



➔ Charmonium is modified by cold nuclear matter effects as

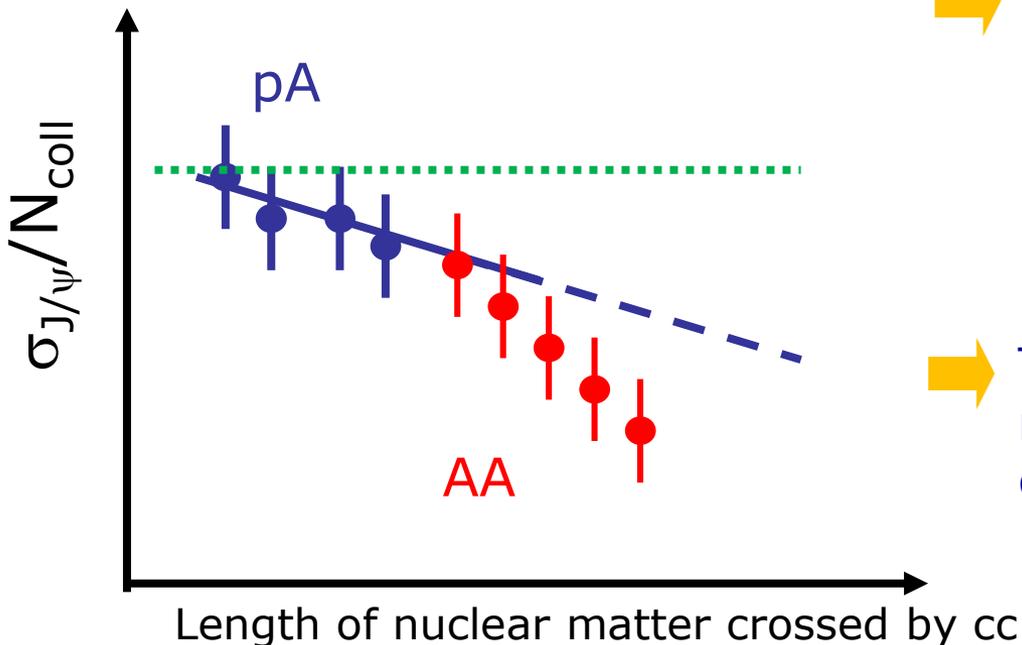
- nuclear parton shadowing
- energy loss
- cc dissociation in the medium...

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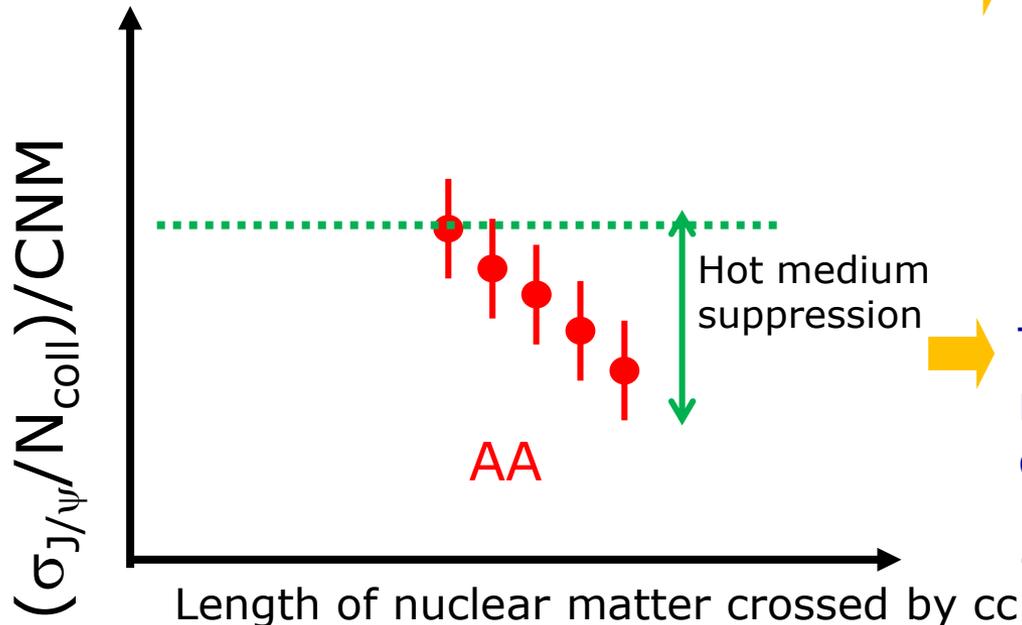
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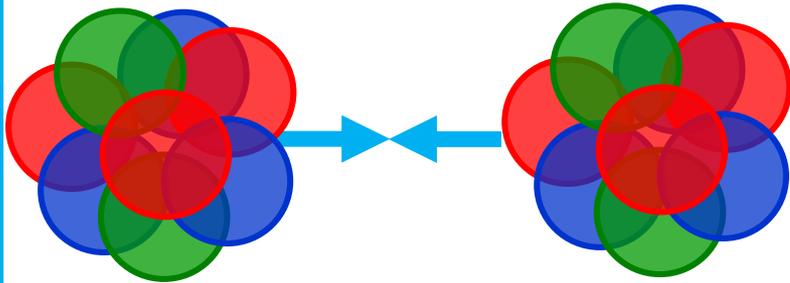
➔ The J/ψ suppression in cold nuclear matter (CNM) can mask genuine QGP effects in AA

➔ These CNM effects need to be calibrated and factorized out

➔ CNM can be studied in pA collisions

HOW CAN WE STUDY QUARKONIUM IN HI?

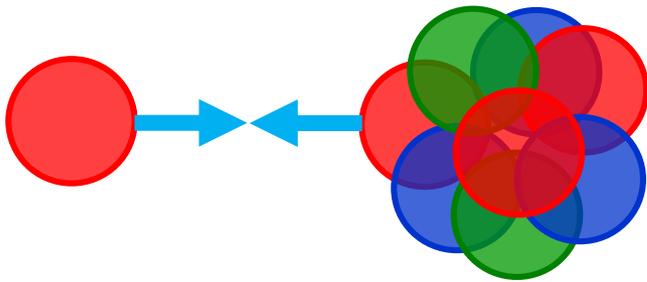
A-A



Quarkonium strongly affected by the hot matter (QGP):
suppression vs regeneration

Studies are done as a function of the collision centrality

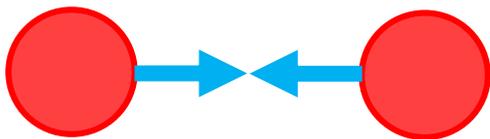
p-A



Quarkonium affected by cold nuclear matter effects (CNM)

Crucial tool to disentangle genuine QGP effect is AA collisions

p-p



Reference process to understand behaviour in pA, AA collisions

Useful to investigate production mechanisms can be investigated (NRQCD, CEM models...)

HOW CAN WE MEASURE MEDIUM EFFECTS?

→ Nuclear modification factor R_{AA} :

compare quarkonium cross sections in AA and pp, scaled by the nucleus mass number A

$$R_{AA}^{J/\psi}(p_T) = \frac{dN_{AA}^{J/\psi} / dp_T}{N_{coll} \times dN_{pp}^{J/\psi} / dp_T}$$

→ If yield scales with the number of binary collisions

$$\rightarrow R_{AA} = 1$$

→ If there are medium effects

$$\rightarrow R_{AA} \neq 1$$

Cold Nuclear Matter effects (CNM):

- Nuclear parton shadowing
- Parton energy loss
- cc in medium dissociation



Hot Medium effects:

- quarkonium suppression
- enhancement due to recombination



(as long as the total charm cross section remains unmodified)

→ knowledge of CNM effects fundamental to disentangle genuine QGP induced suppression in AA

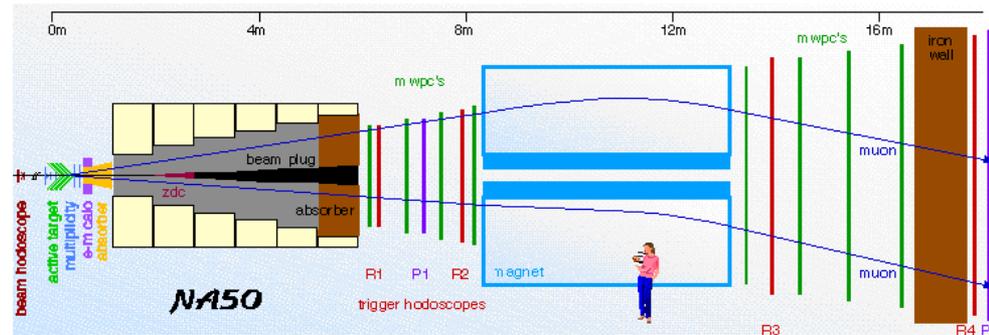
HOW IS QUARKONIUM DETECTED?

➔ Experimentally quarkonia are detected through the decay channels:

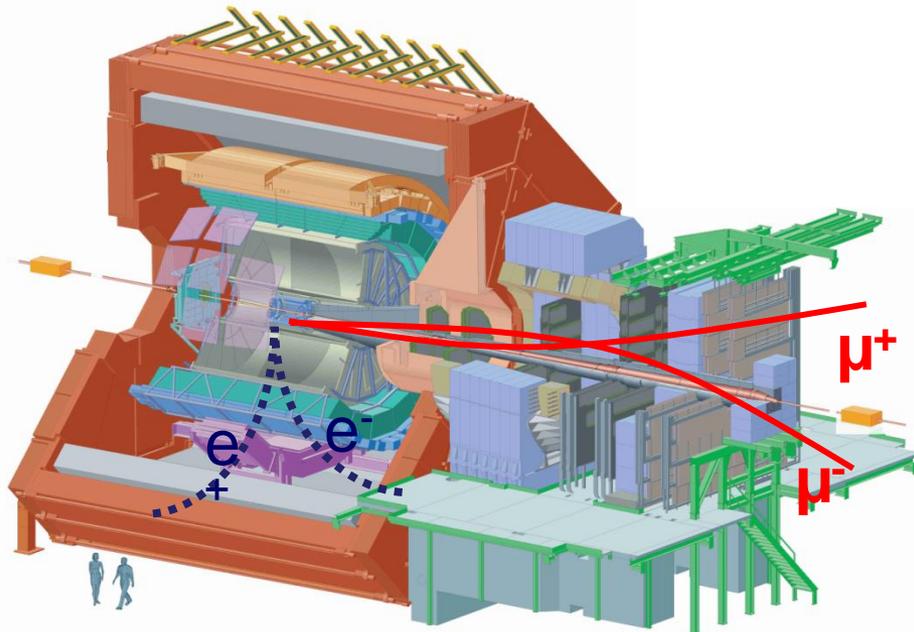
$$J/\psi \rightarrow \mu^+ \mu^- \quad (\text{B.R.} = 5.93\%)$$

$$J/\psi \rightarrow e^+ e^- \quad (\text{B.R.} = 5.94\%)$$

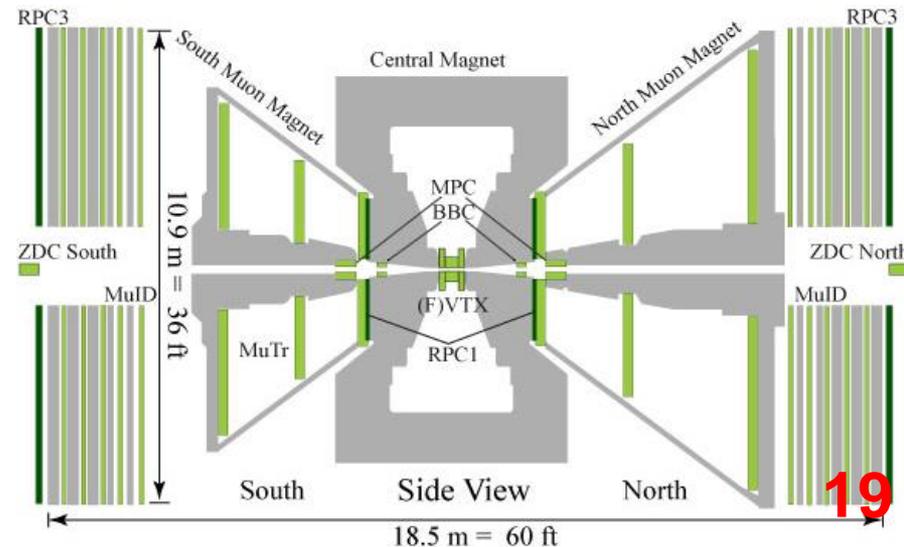
NA50



ALICE



PHENIX



FROM SPS TO LHC

➡ 30 years of data taking: counting rooms...



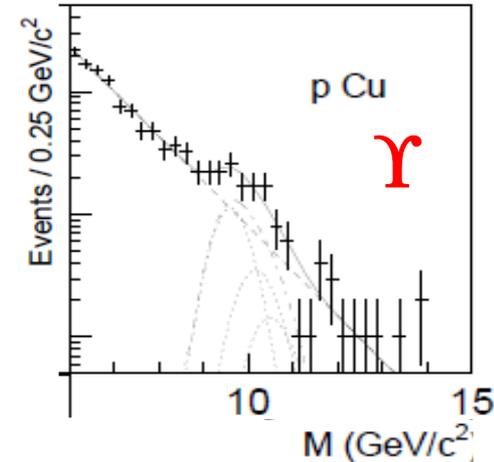
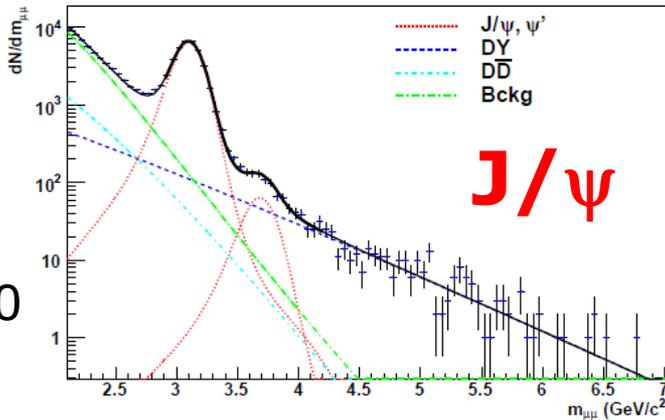
FROM SPS TO LHC

➡ Invariant mass spectra...

NA50/NA60 @ SPS

Fixed target experiments

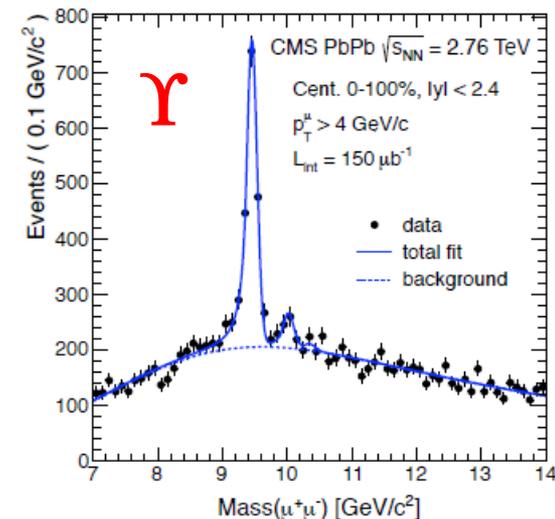
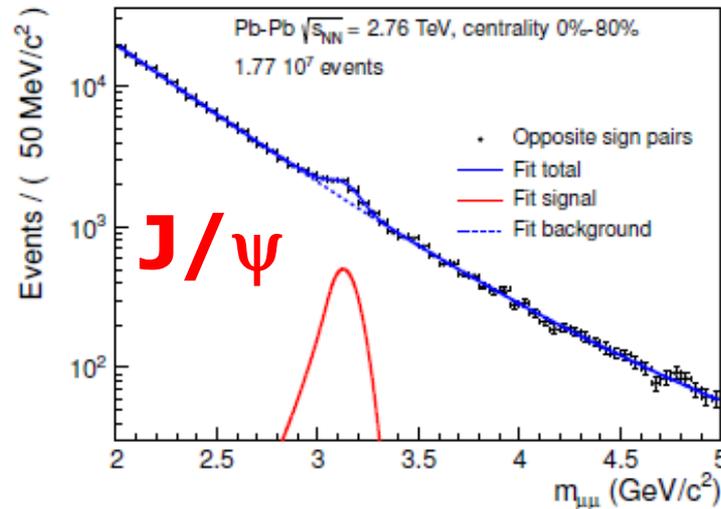
- high J/ψ statistics:
 10^5 in NA50, $4 \cdot 10^4$ in NA60
- $\sim 300 \Upsilon$ in pA (limited resolution)



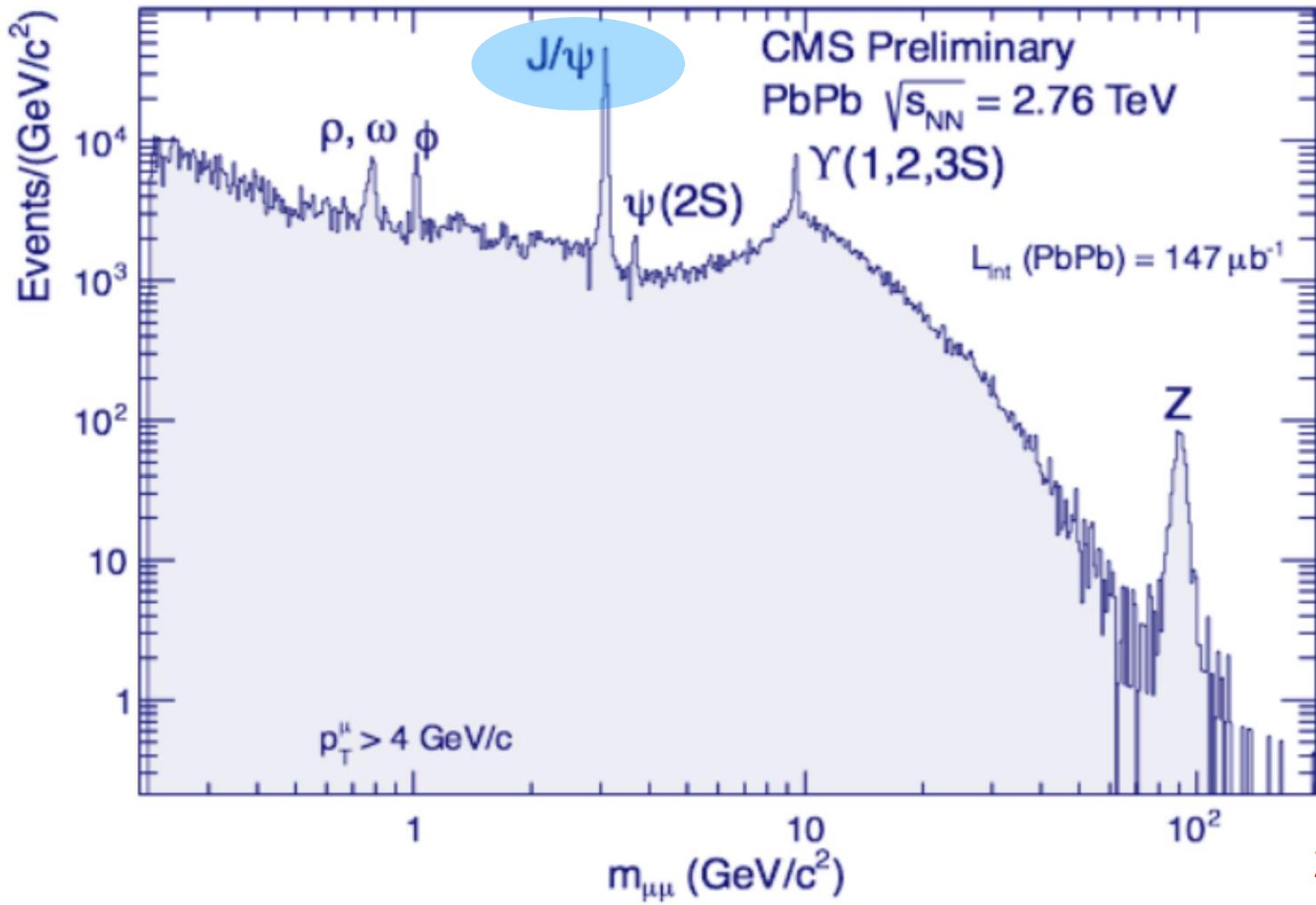
ALICE/CMS @ LHC

Collider experiments

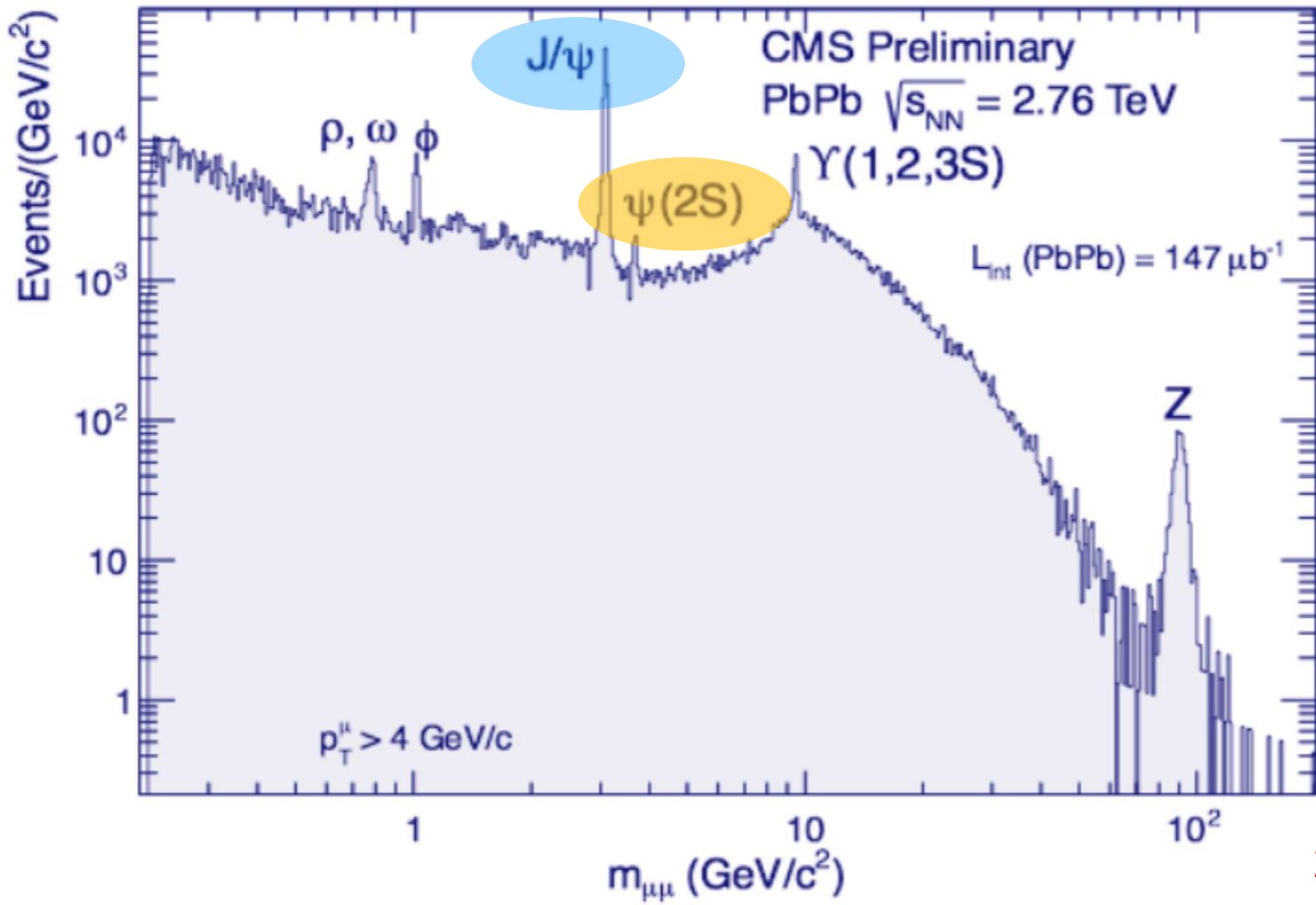
- max J/ψ statistics
in ALICE: ~ 25000
in CMS: ~ 8000
(2000 $\Upsilon(1S)$)



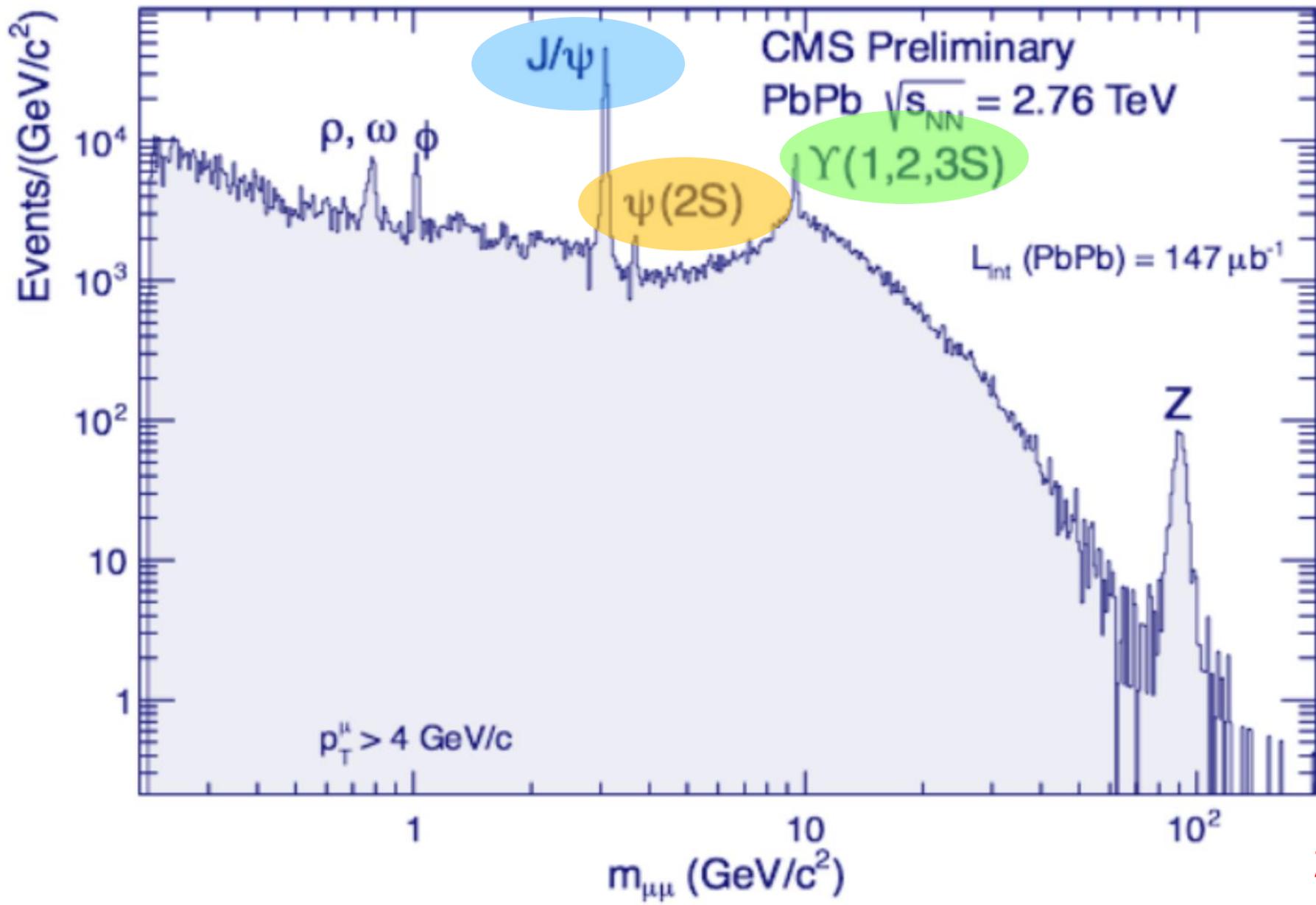
QUARKONIUM



QUARKONIUM



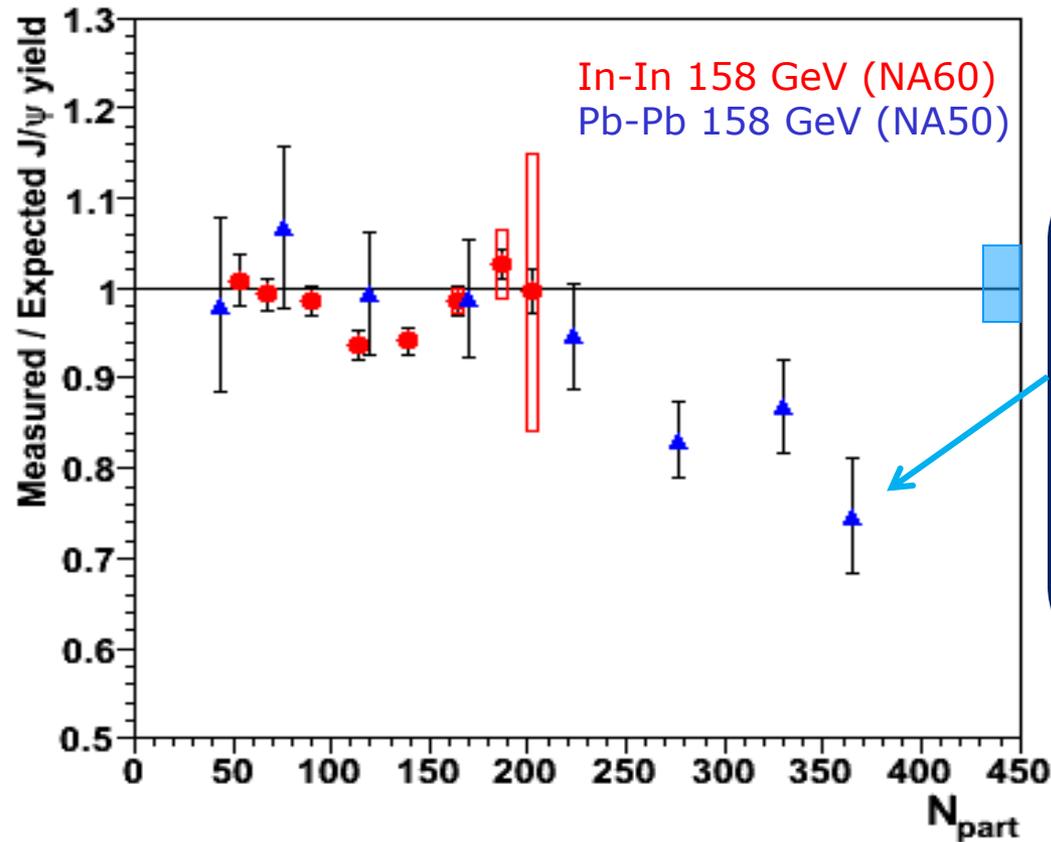
QUARKONIUM



FIRST J/ψ MEASUREMENTS AT LOW ENERGY: SPS

Charmonium production deeply investigated at

SPS (NA38, NA50, NA60) $\sqrt{s_{NN}} = 17$ GeV



SPS:

first evidence of anomalous suppression (i.e. beyond CNM expectations) in Pb-Pb at $\sqrt{s} = 17$ GeV

~30% suppression compatible with $\psi(2S)$ and χ_c decays



FIRST J/ψ MEASUREMENTS AT LOW ENERGY: RHIC

Charmonium production deeply investigated at

RHIC (PHENIX, STAR) $\sqrt{s_{NN}} = 39, 62.4, 200 \text{ GeV}$

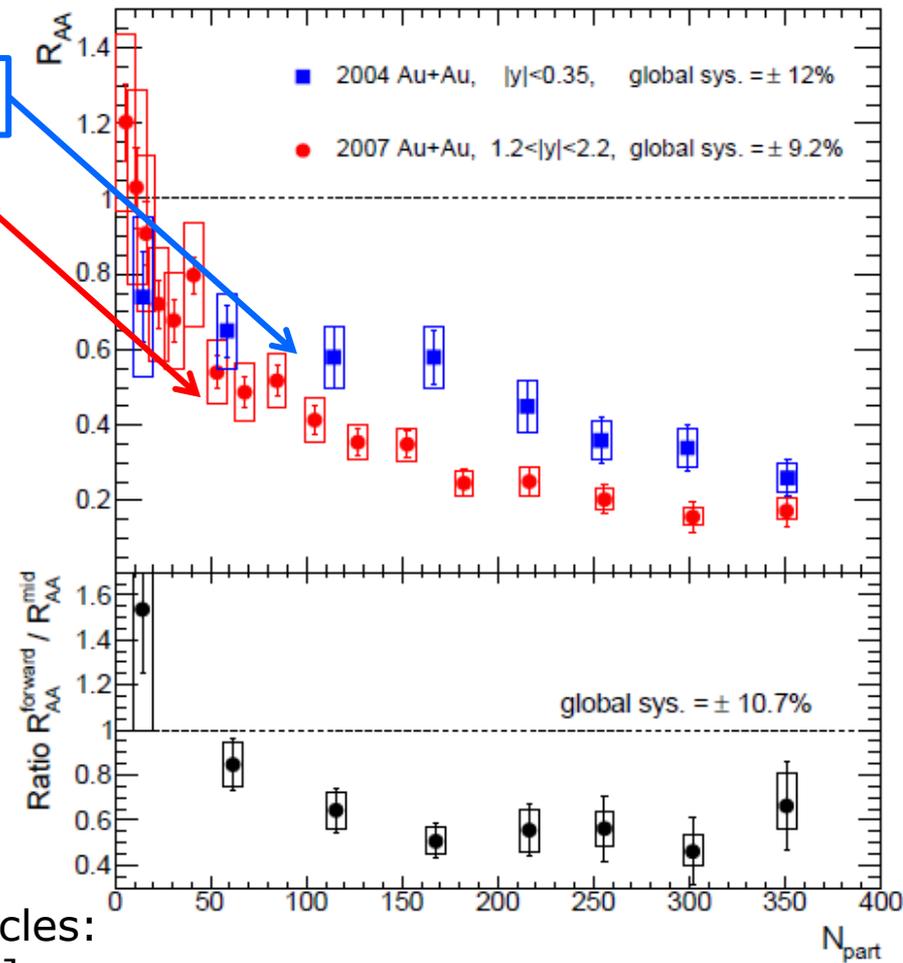
Mid-rapidity

Forward-rapidity

RHIC:

suppression, strongly rapidity dependent, in Au-Au at $\sqrt{s} = 200 \text{ GeV}$

Stronger suppression at forward y
 \rightarrow not expected if suppression increases with energy density (larger at mid- y)



Rapidity $y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right)$

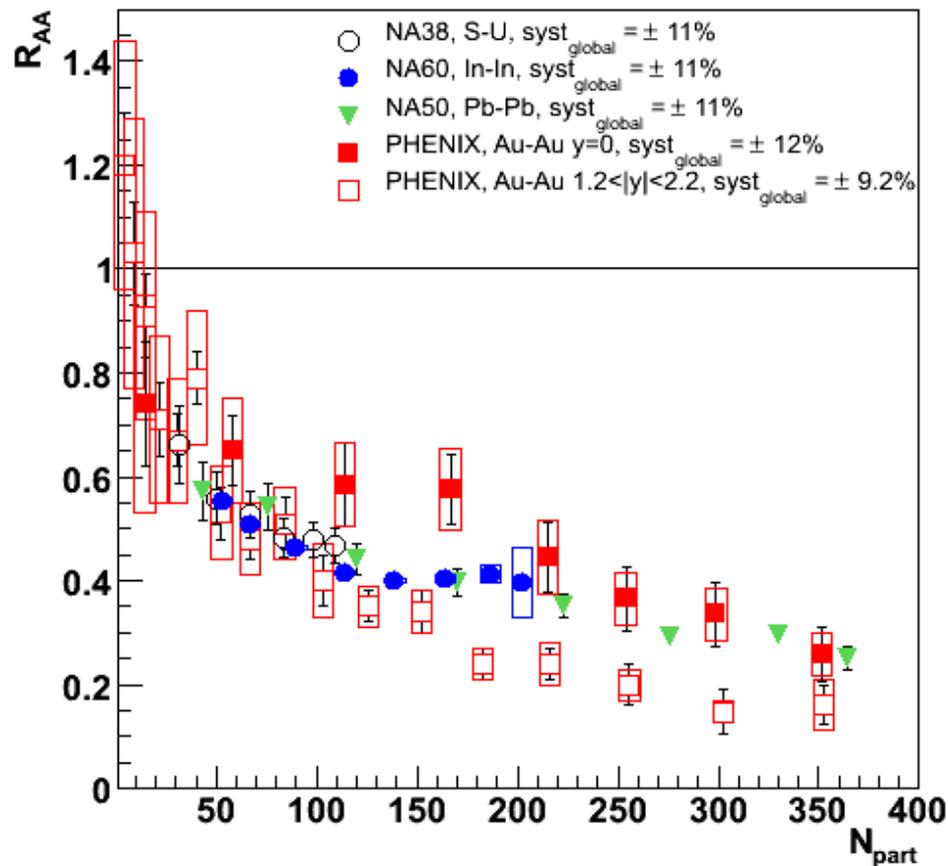
For relativistic particles:
 $y \sim \eta = -\ln[\tan(\theta/2)]$

COMPARISON OF SPS AND RHIC RESULTS

Charmonium production deeply investigated at

SPS (NA50, NA60) $\sqrt{s_{NN}} = 17$ GeV

RHIC (PHENIX, STAR) $\sqrt{s_{NN}} = 39, 62.4, 200$ GeV



➔ Puzzles from SPS and RHIC

- **RHIC:** stronger suppression at forward rapidities
- **SPS vs. RHIC:** similar R_{AA} pattern versus centrality

➔ Hint for (re)combination at RHIC?
No final theoretical explanation

QUARKONIUM AT LHC

➔ Decisive inputs expected from LHC results, having access to:

higher energies

→ stronger quarkonium suppression?

more charm

→ larger (re)combination?

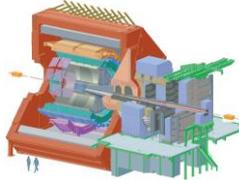
more bottom

→ Υ can be investigated

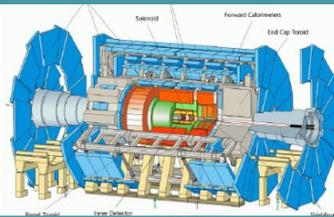


QUARKONIUM AT LHC

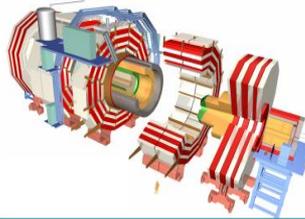
ALICE $J/\psi, \psi(2S) \rightarrow \mu^+\mu^-$
 $\Upsilon \rightarrow \mu^+\mu^-$
 $J/\psi \rightarrow e^+e^-$



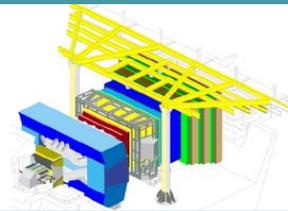
ATLAS $J/\psi \rightarrow \mu^+\mu^-$



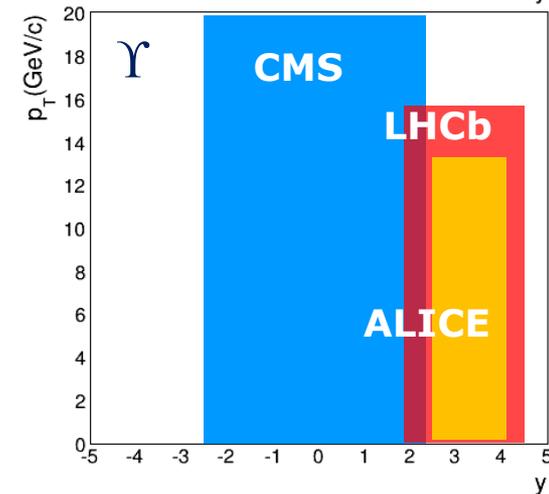
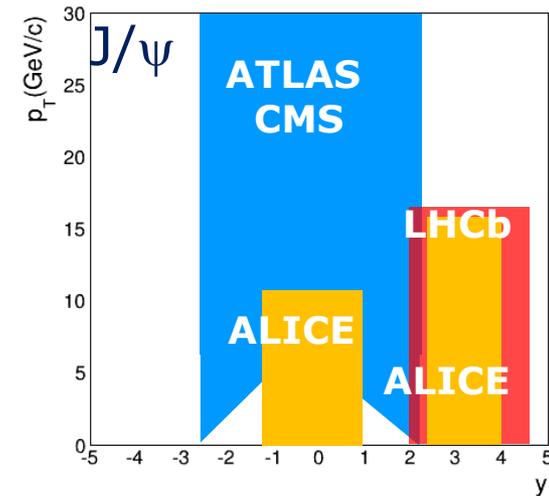
CMS $J/\psi, \psi(2S) \rightarrow \mu^+\mu^-$
 $\Upsilon \rightarrow \mu^+\mu^-$



LHCb $J/\psi, \Upsilon \rightarrow \mu^+\mu^-$
 (no heavy ion physics program)



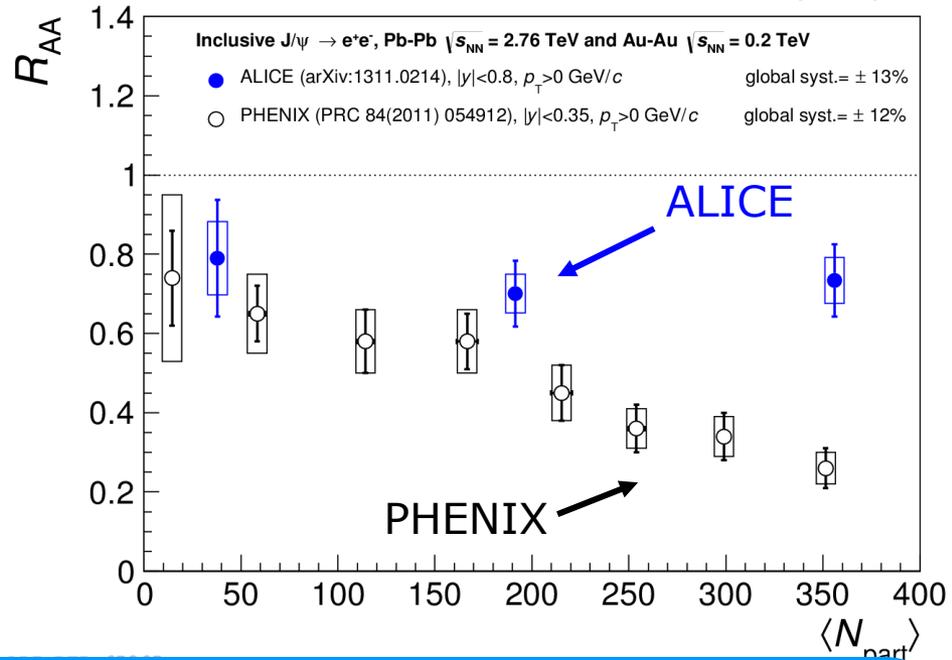
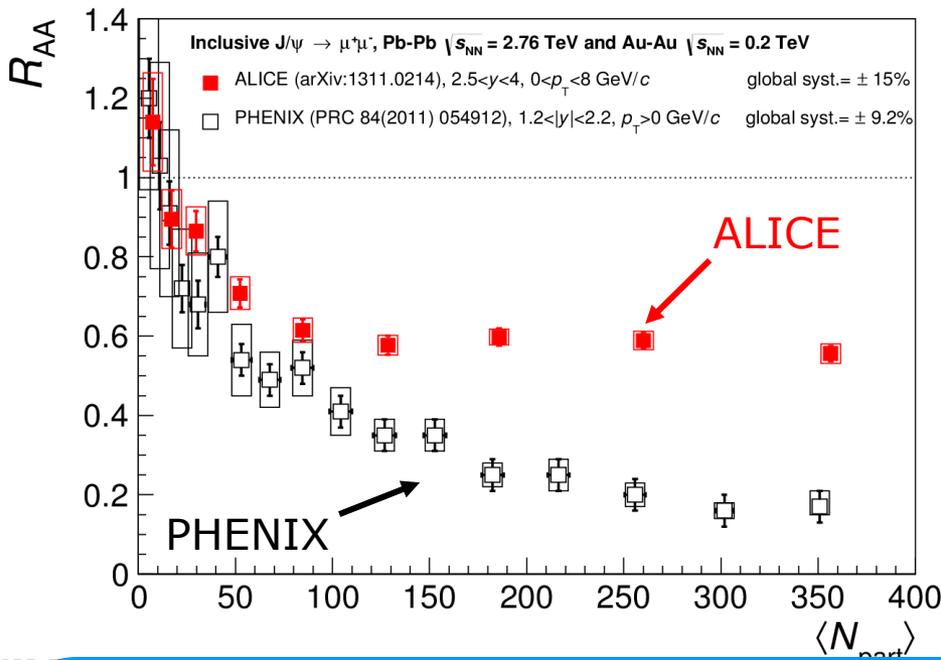
➔ Kinematic coverage of quarkonium measurements:



➔ Complementary quarkonium results from LHC experiments!

J/ψ R_{AA} VS CENTRALITY: ALICE VS PHENIX

➔ Centrality dependence of the J/ψ inclusive R_{AA} studied by ALICE in both central and forward rapidities down to zero p_T ALICE Coll. PLB 734 (2014) 314



➔ ALICE results:

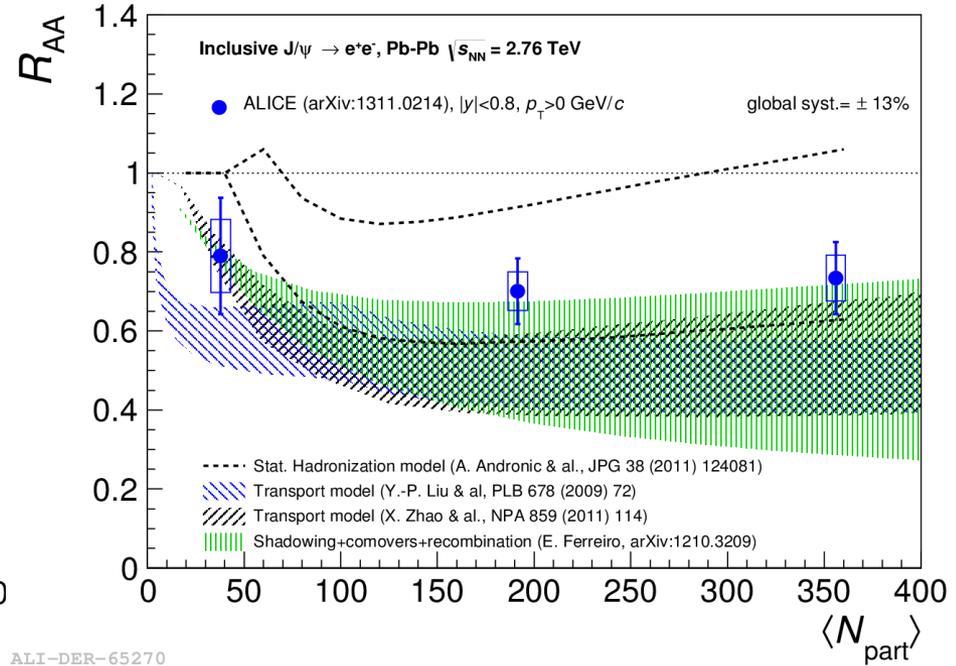
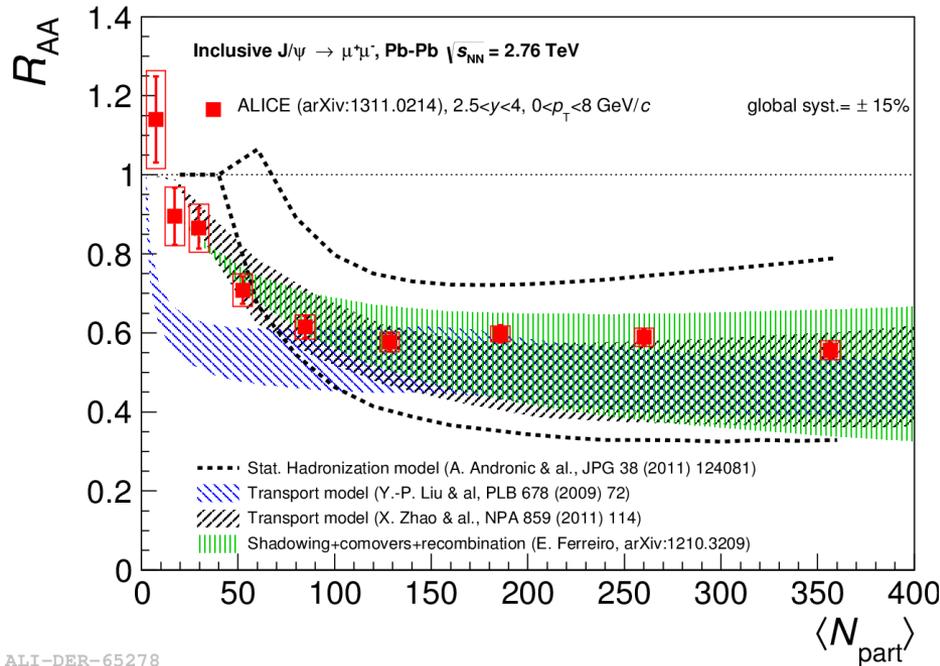
➔ clear J/ψ suppression with almost no centrality dependence for $N_{part} > 100$

➔ Comparison with PHENIX:

➔ ALICE results show weaker centrality dependence and smaller suppression for central events

➔ Behaviour expected in a (re)combination scenario

J/ψ R_{AA} VS CENTRALITY: THEORY COMPARISON



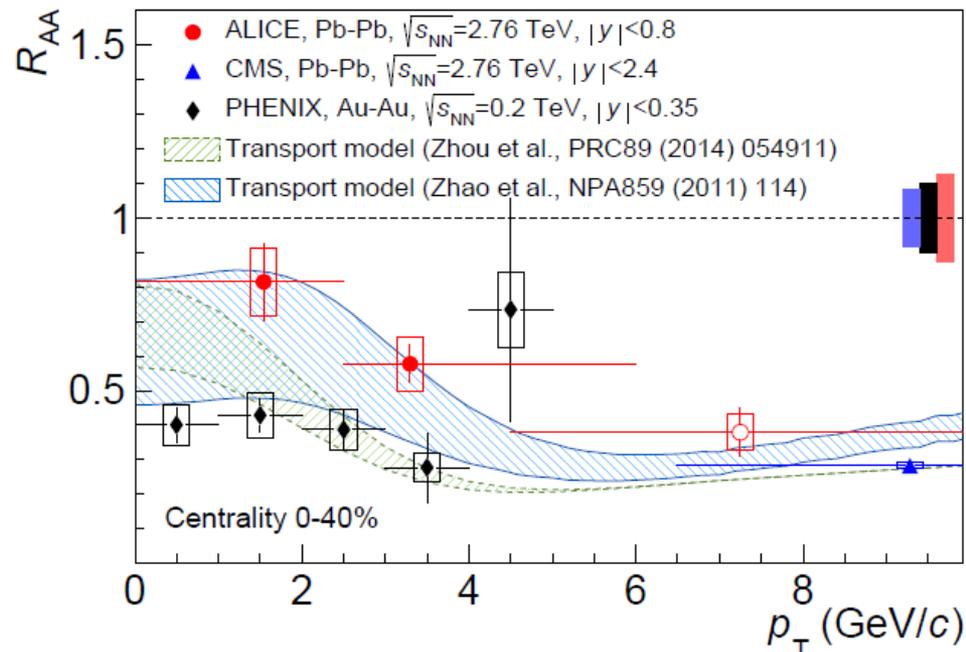
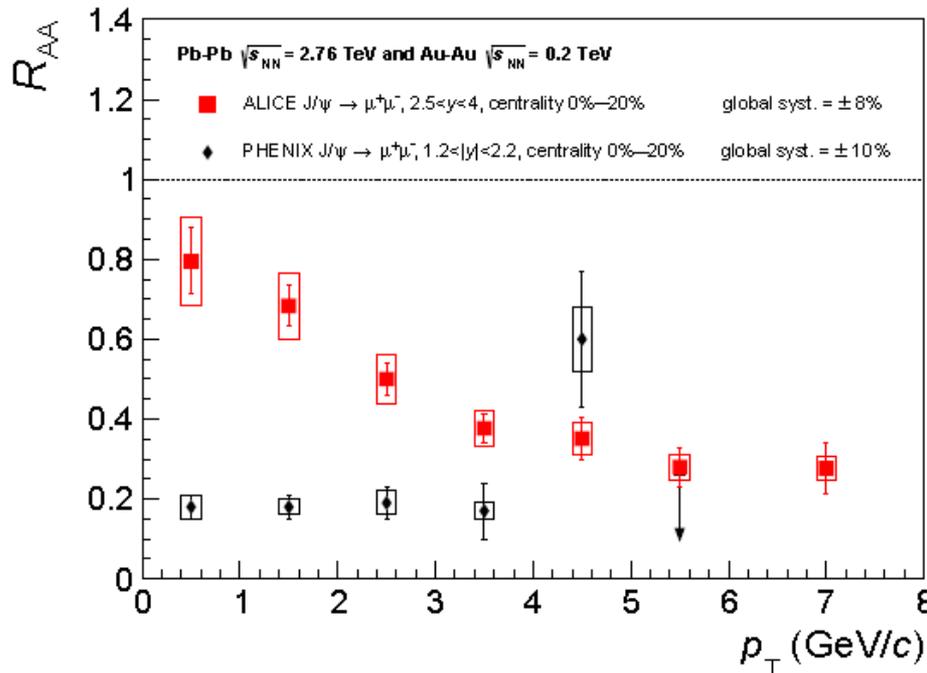
Comparison to theory calculations:

- ➔ Models including a large fraction (> 50% in central collisions) of J/ψ produced from (re)combination or models with all J/ψ produced at hadronization provide a reasonable description of ALICE results
- ➔ Still rather large theory uncertainties: models will benefit from a precise measurement of σ_{cc} and from cold nuclear matter evaluation

ALICE: R_{AA} VS p_T

➔ J/ψ production via (re)combination should be more important at low transverse momentum ➔ p_T region accessible by ALICE

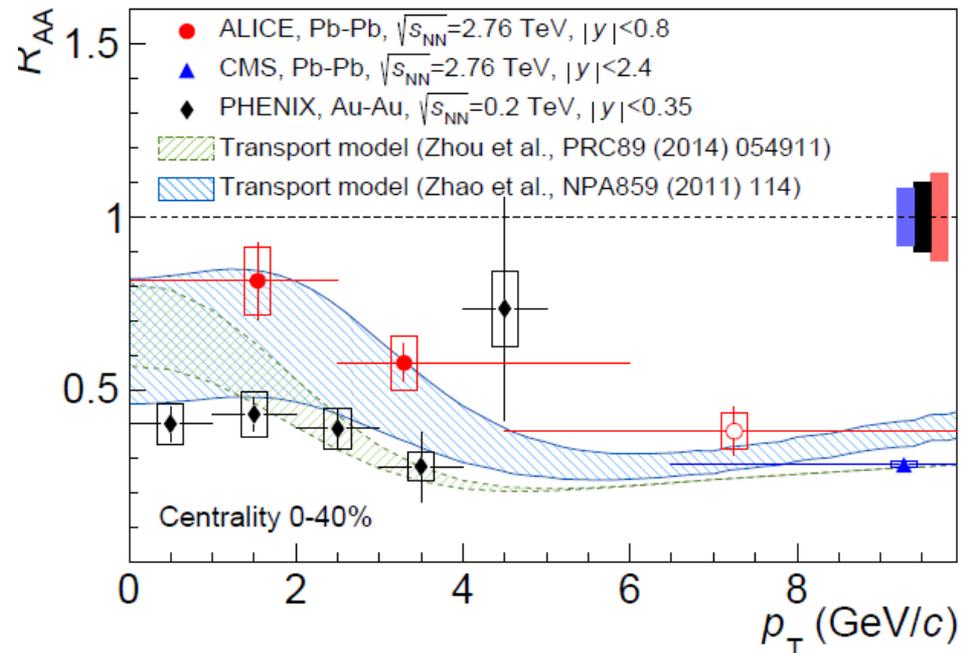
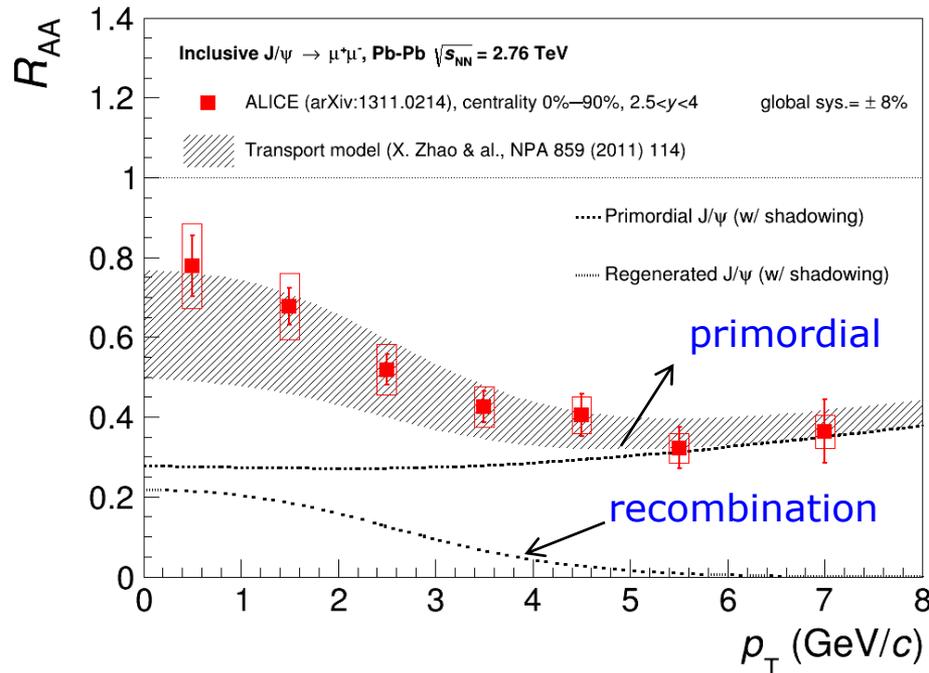
arXiv:1504.07151



- ➔ Different suppression for low and high p_T J/ψ
- ➔ Smaller R_{AA} for high p_T J/ψ
- ➔ Striking difference, at low p_T , between PHENIX and ALICE patterns

ALICE: R_{AA} VS p_T

➔ J/ψ production via (re)combination should be more important at low transverse momentum ➔ p_T region accessible by ALICE



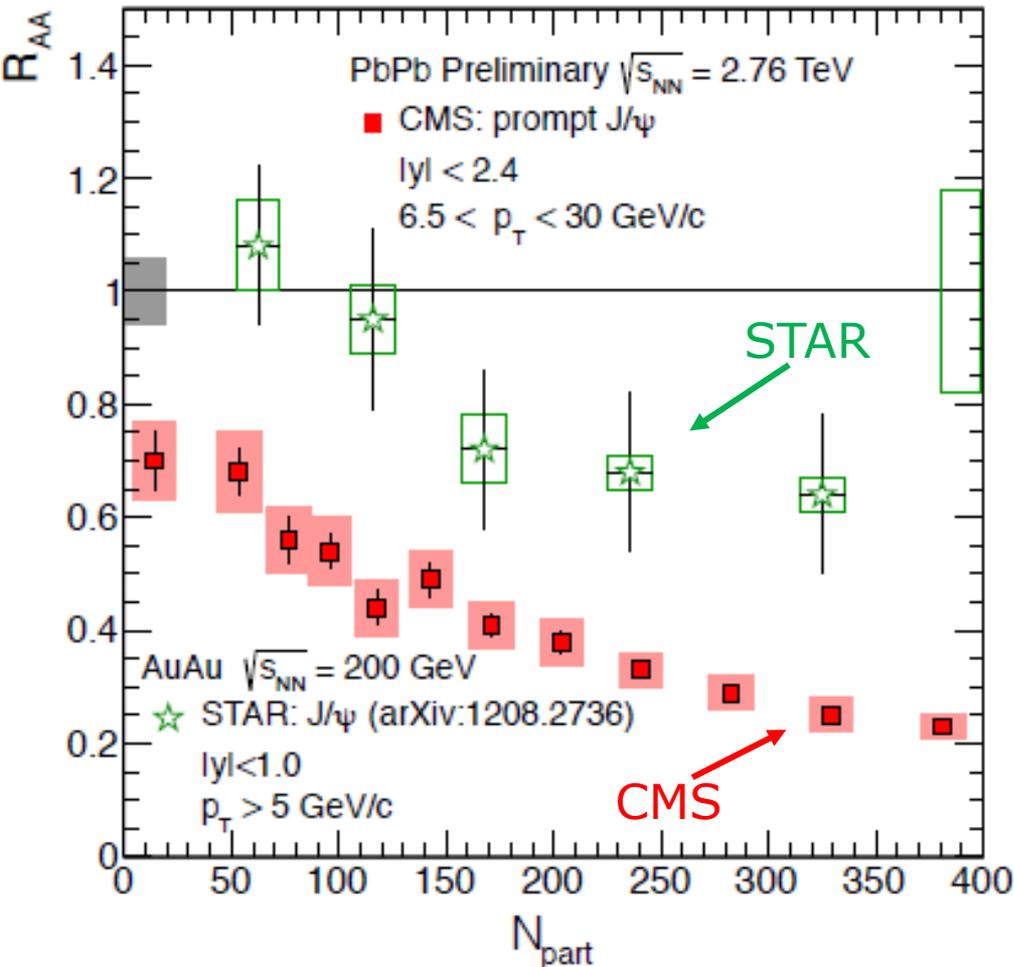
- ➔ Different suppression for low and high p_T J/ψ
- ➔ Smaller R_{AA} for high p_T J/ψ
- ➔ Models: ~50% of low- p_T J/ψ are produced via (re)combination, while at high p_T the contribution is negligible

HIGH p_T J/ψ : CMS & STAR

➔ At LHC high p_T J/ψ have been investigated by CMS

➔ Limits in the CMS low- p_T J/ψ acceptance since muons need to overcome the magnetic field and energy loss in the absorber:

- mid- y : $p_T > 6.5$ GeV/c
- forward y : $p_T > 3$ GeV/c

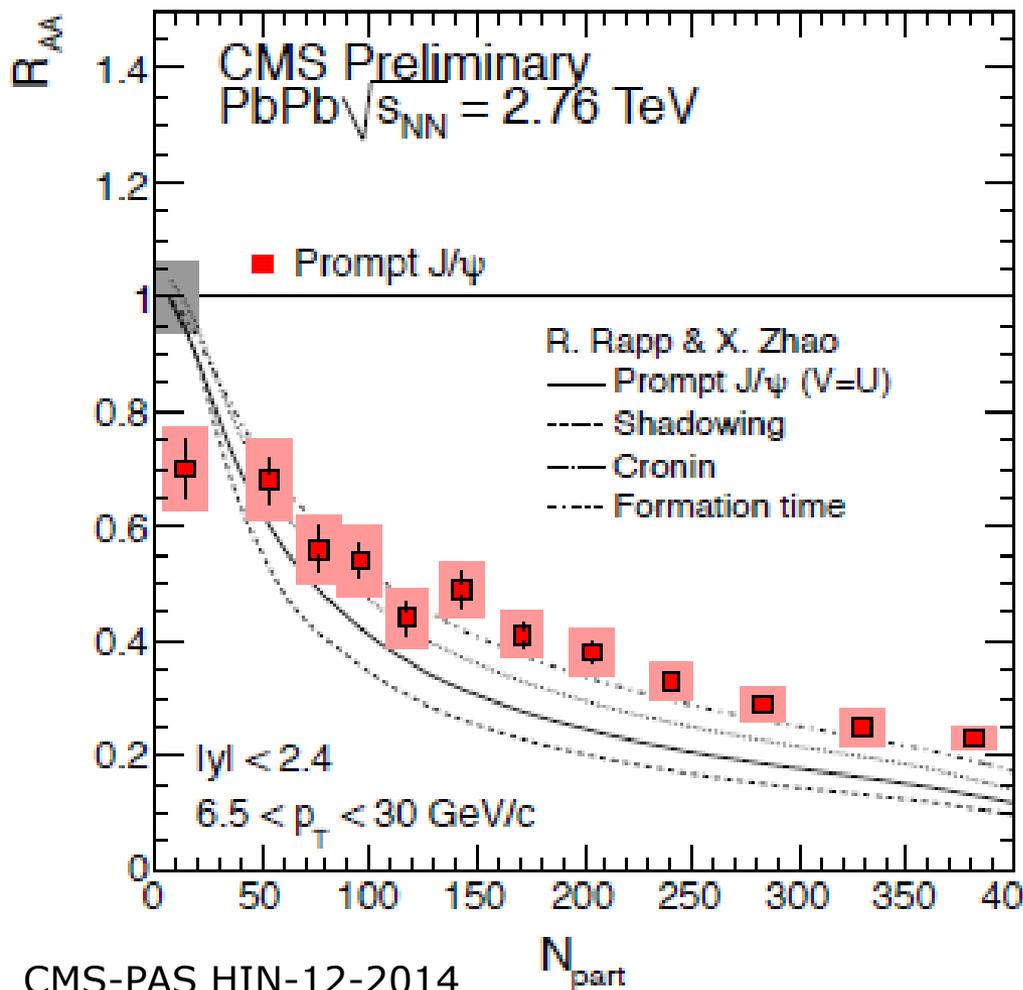


➔ Opposite behavior when compared to ALICE low- p_T results

➔ Suppression is stronger at LHC energy (by a factor ~ 3 compared to RHIC for central events)

HIGH p_T J/ψ : CMS & STAR

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- mid- y : $p_T > 6.5$ GeV/c
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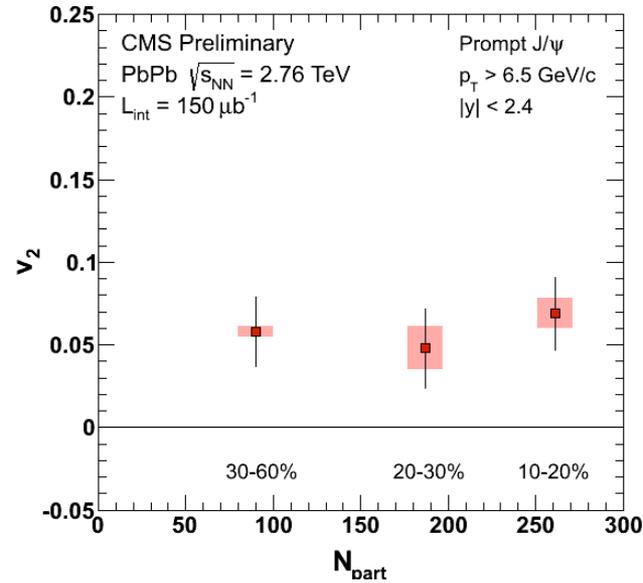
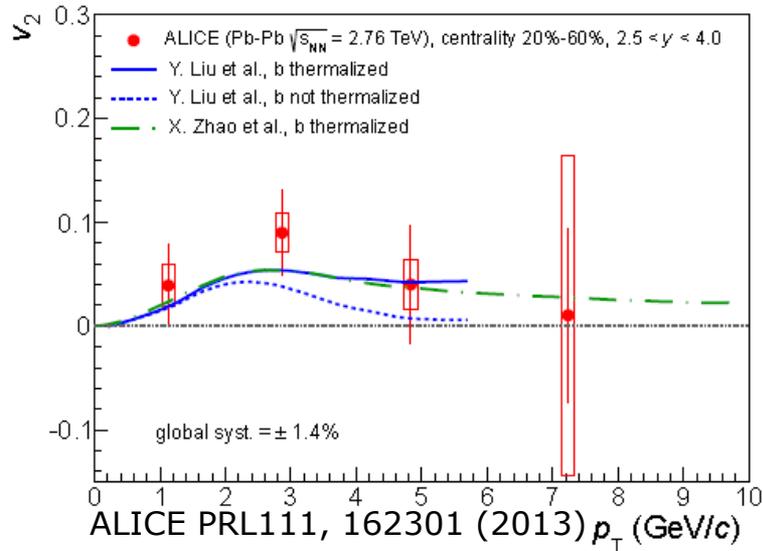
➔ Opposite behavior when compared to ALICE low- p_T results

➔ Suppression is stronger at LHC energy (by a factor ~ 3 compared to RHIC for central events)

➔ negligible (re)generation effects expected at high p_T

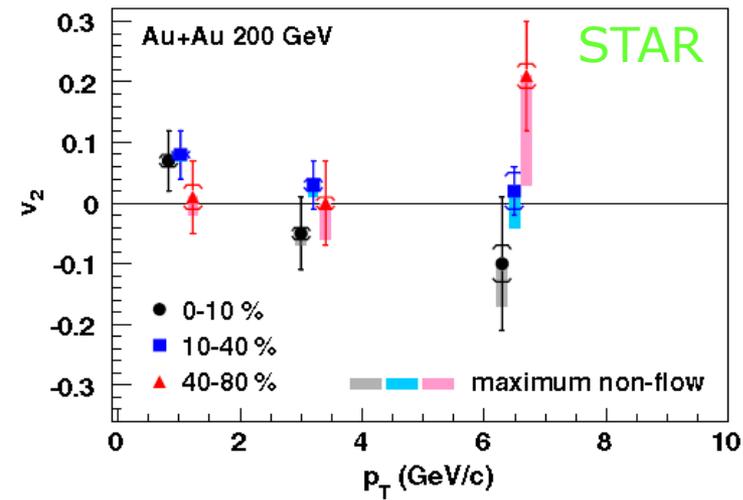
J/ψ FLOW

➔ The contribution of J/ψ from (re)combination should lead to a significant elliptic flow signal at LHC energy



D.Moon, HP2013

➔ Hint for J/ψ flow at LHC, contrary to $v_2 \sim 0$ observed at RHIC!



➔ ALICE: qualitative agreement with transport models including regeneration

➔ CMS: path-length dependence of energy loss?

STAR, PRL 052301(2013)

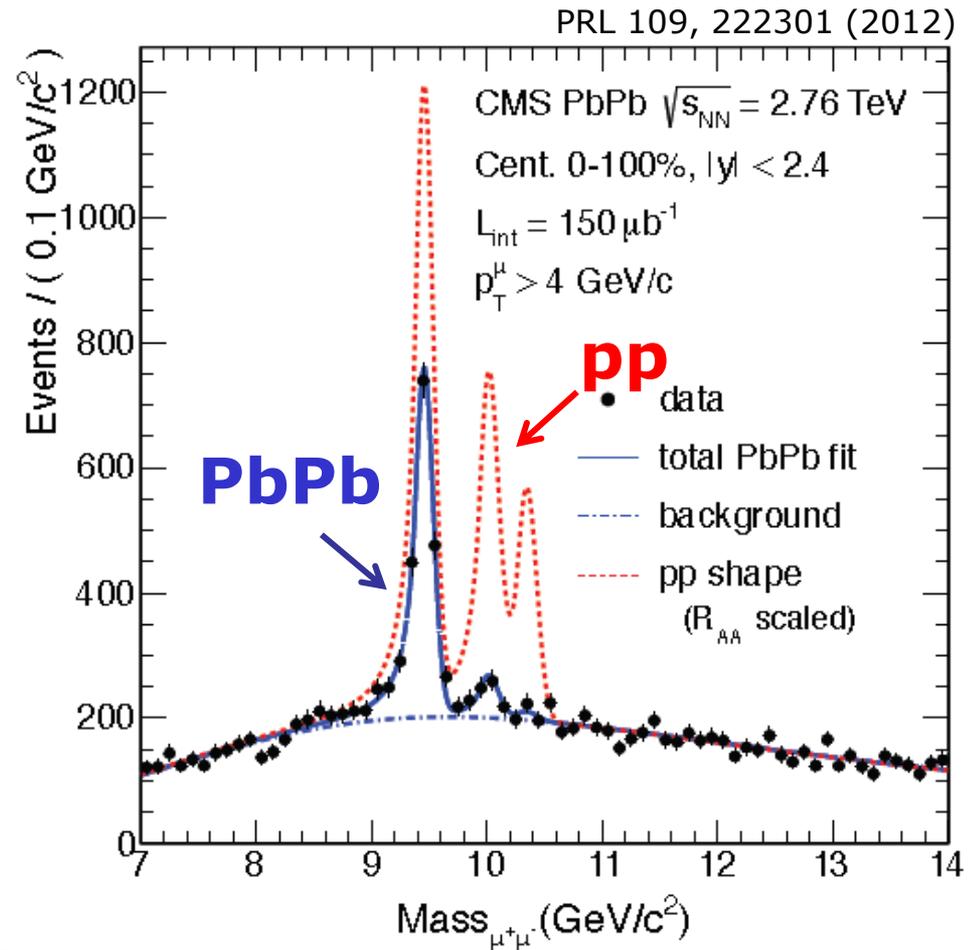
$\Upsilon(1S)$ PRODUCTION IN Pb-Pb COLLISIONS

➔ LHC is the machine for studying bottomonium in AA collisions

Main features of bottomonium production wrt charmonia:

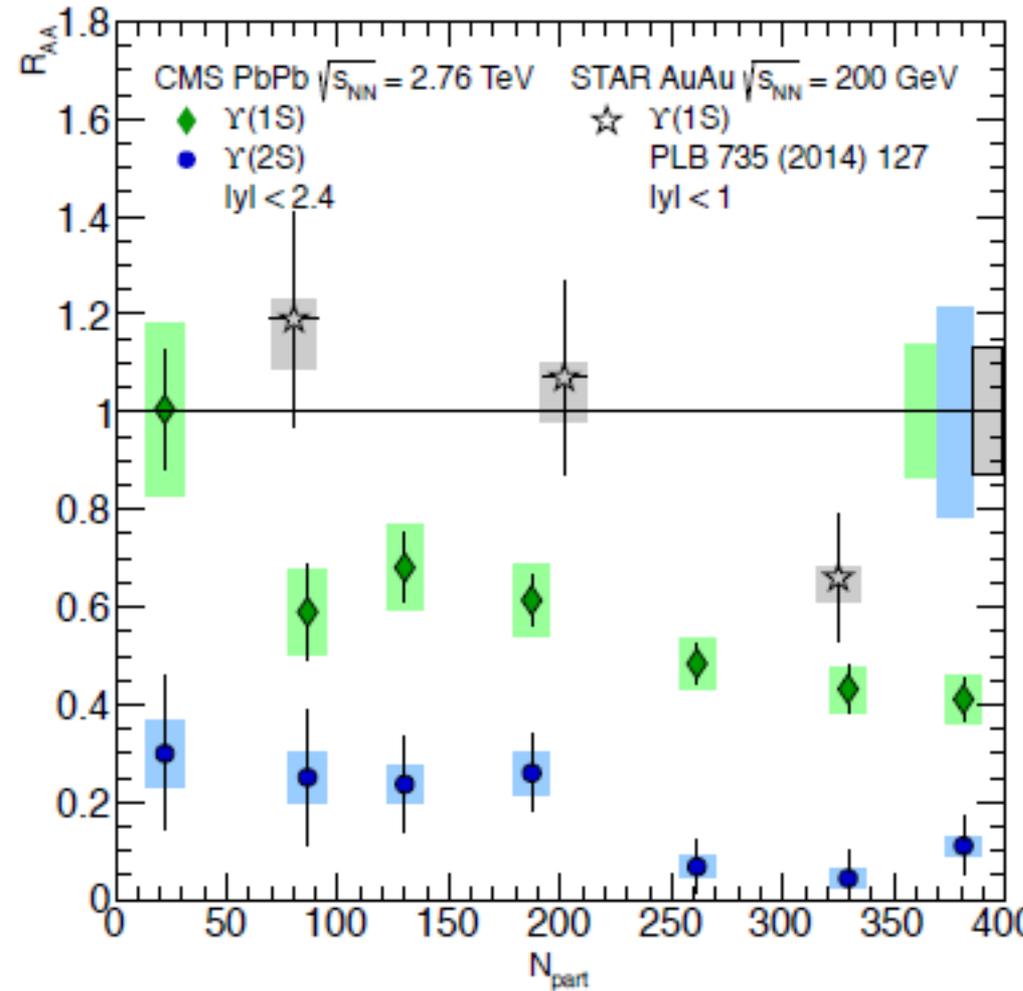
- no B hadron feed-down
- gluon shadowing effect are smaller
- (re)combination expected to be smaller
- theoretical predictions more robust due to the higher mass of b quark

with a drawback...smaller production cross-section



➔ Clear suppression of Υ states in PbPb with respect to pp collisions

$\Upsilon(1S)$ PRODUCTION IN PB-PB COLLISIONS



- ➔ Clear suppression of $\Upsilon(2S)$
- ➔ $\Upsilon(1S)$ suppression compatible with suppression of excited states (50% feed-down)
- ➔ Sequential suppression of the three Υ states according to their binding energy:

$$R_{AA}^{\Upsilon(3S)} < R_{AA}^{\Upsilon(2S)} < R_{AA}^{\Upsilon(1S)}$$

$$R_{AA}(\Upsilon(1S)) = 0.56 \pm 0.08 \text{ (stat)} \pm 0.07 \text{ (syst)}$$

$$R_{AA}(\Upsilon(2S)) = 0.12 \pm 0.04 \text{ (stat)} \pm 0.02 \text{ (syst)}$$

$$R_{AA}(\Upsilon(3S)) < 0.1 \text{ (at 95\% C.L.)}$$

- ➔ Suppression at LHC is stronger than at RHIC

QUARKONIUM IN AA: WHERE ARE WE?

→ ~30 years after first suppression prediction, this is observed with very good accuracy!

→ Two main mechanisms at play:

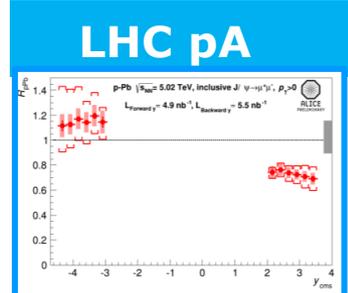
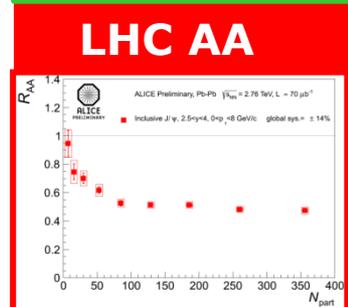
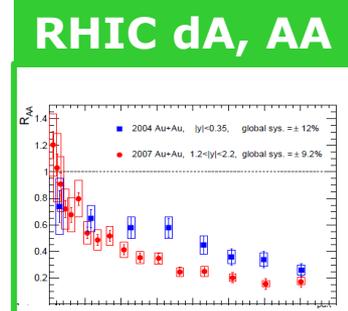
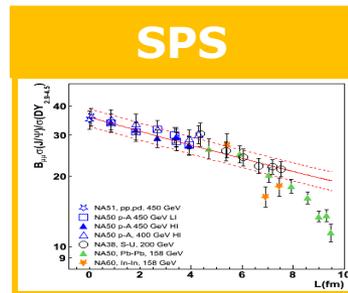
1. Suppression in a deconfined medium
2. (charmonium) re-combination at high \sqrt{s} and low p_T

can qualitatively explain the main features of the results

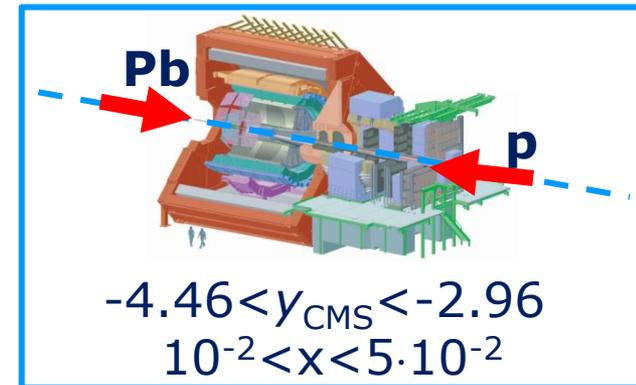
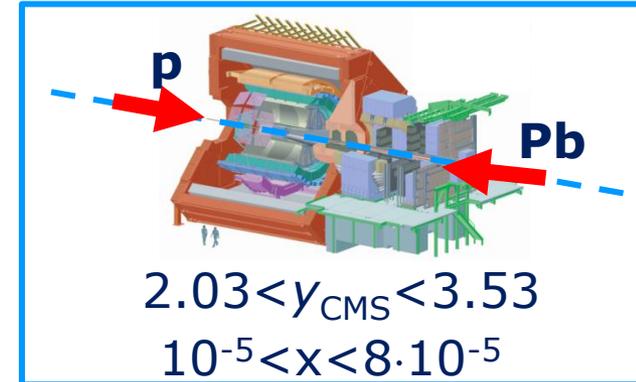
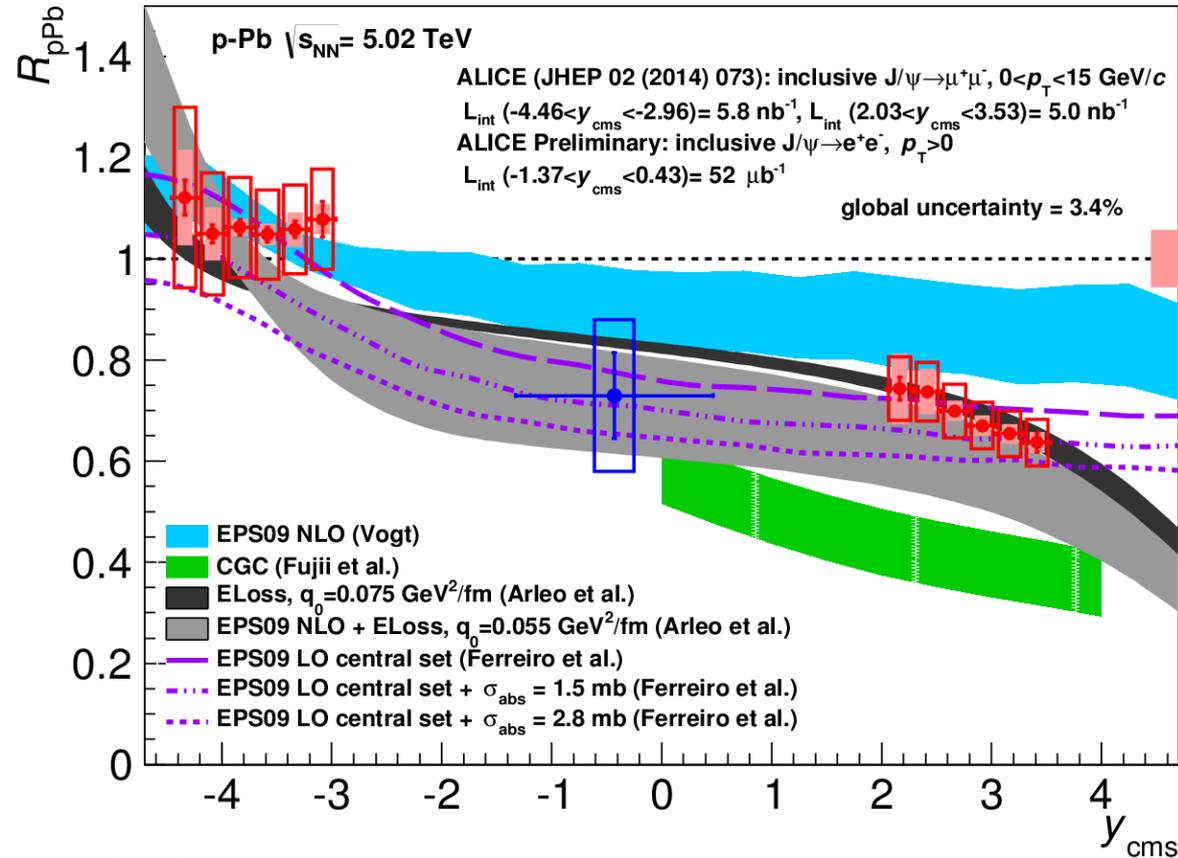
→ To move towards a more quantitative understanding, a precise knowledge of CNM effects is crucial!

→ pA results, where no hot medium should be formed, are needed to:

1. investigate initial/final state CNM effects
2. build a reference for AA collisions



J/ψ IN pA COLLISIONS



 J/ψ production is modified also in pA because of CNM effects: R_{pA} decreases towards forward y

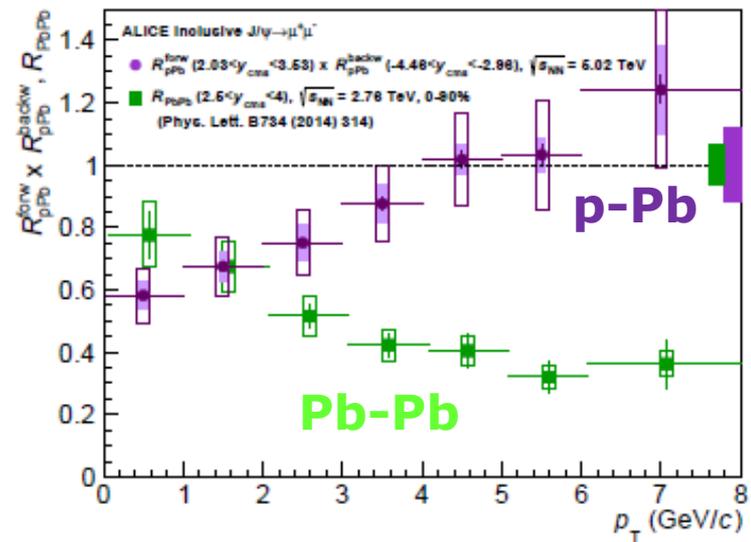
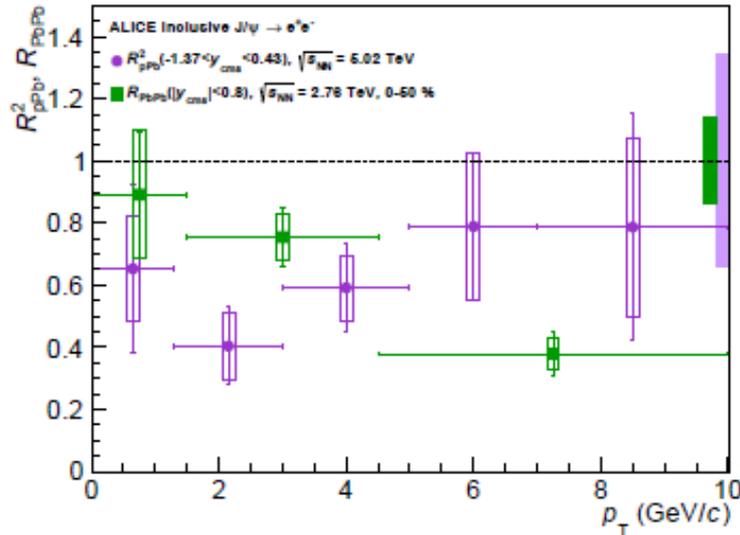
 **Theoretical predictions:** reasonable agreement with

- shadowing calculations and models including coherent parton energy loss
- CGC description seems not to be favoured

CNM EFFECTS FROM p-Pb TO Pb-Pb

➔ Once CNM effects are measured in pA, what can we learn on J/ψ production in PbPb?

- Hypothesis:
- 2→1 kinematics for J/ψ production
 - CNM effects (dominated by shadowing) factorize in p-A
 - CNM obtained as $R_{pA} \times R_{Ap}$ (R_{pA}^2), similar x-coverage as PbPb



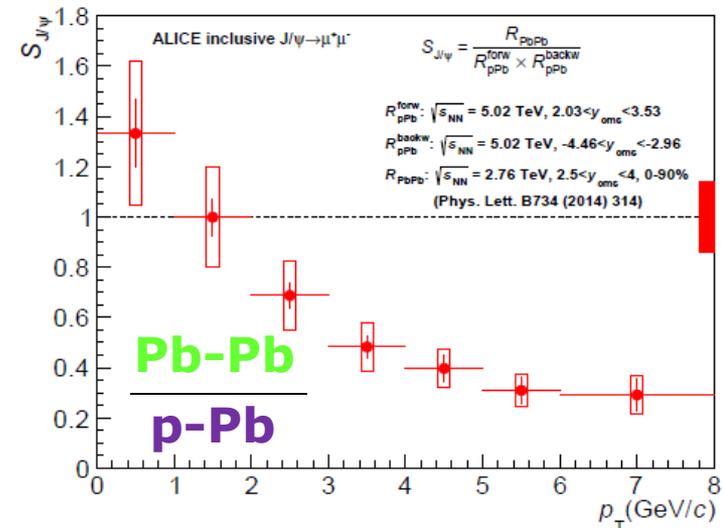
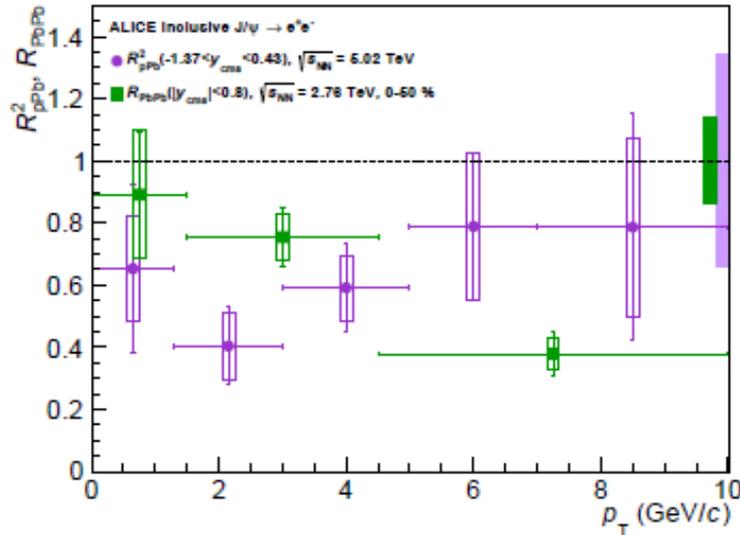
➔ Sizeable p_T dependent suppression still visible
 → CNM effects not enough to explain AA data at high p_T

➔ we get rid of CNM effects, by doing the ratio **AA / pA**

CNM EFFECTS FROM p-Pb TO Pb-Pb

➔ Once CNM effects are measured in pA, what can we learn on J/ψ production in PbPb?

- Hypothesis:
- 2→1 kinematics for J/ψ production
 - CNM effects (dominated by shadowing) factorize in p-A
 - CNM obtained as $R_{pA} \times R_{Ap}$ (R_{pA}^2), similar x-coverage as PbPb

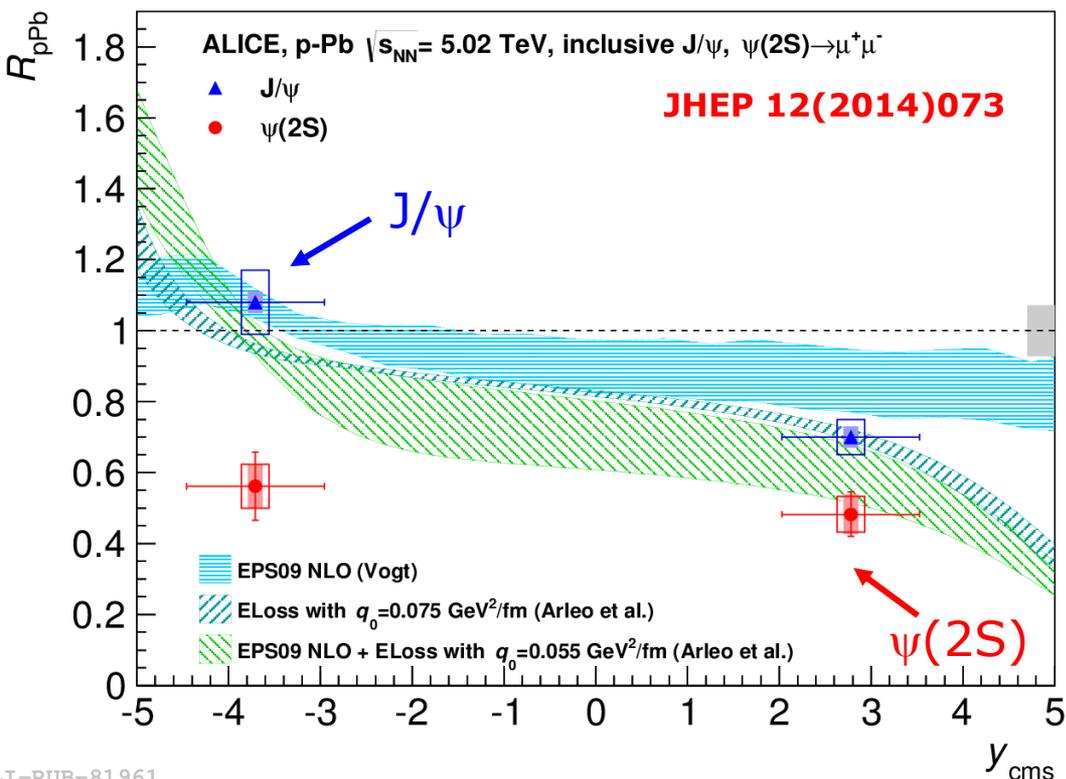


➔ Sizeable p_T dependent suppression still visible
 → CNM effects not enough to explain AA data at high p_T

➔ we get rid of CNM effects, by doing the ratio **AA / pA**

$\psi(2S)/J/\psi$ IN p-Pb

A strong decrease of the $\psi(2S)$ production in p-Pb, relative to J/ψ , is observed with respect to the pp measurement ($2.5 < y_{\text{cms}} < 4$, $\sqrt{s} = 7\text{TeV}$)



Similar effect seen by PHENIX in d-Au at $\sqrt{s_{\text{NN}}} = 200\text{ GeV}$

same initial state CNM effects (shadowing & coherent energy loss) for J/ψ and $\psi(2S)$

theoretical predictions in disagreement with $\psi(2S)$ result

Final state effects related to the (hadronic) medium created in the p-Pb collisions?

AT THE END OF LHC RUN-I...

Large wealth of results at LHC complementing SPS and RHIC measurements!

- ➔ Very interesting observations, qualitative understanding of the main quarkonium features in A-A:
 - important role of charmonium (re)generation processes at low p_T
 - bottomonium sequential suppression observed

- ➔ In p-A collisions:
 - interplay of shadowing and coherent energy loss can satisfactorily describe the J/ψ results
 - loosely bound $\psi(2S)$ is likely influenced by the hadronic final state

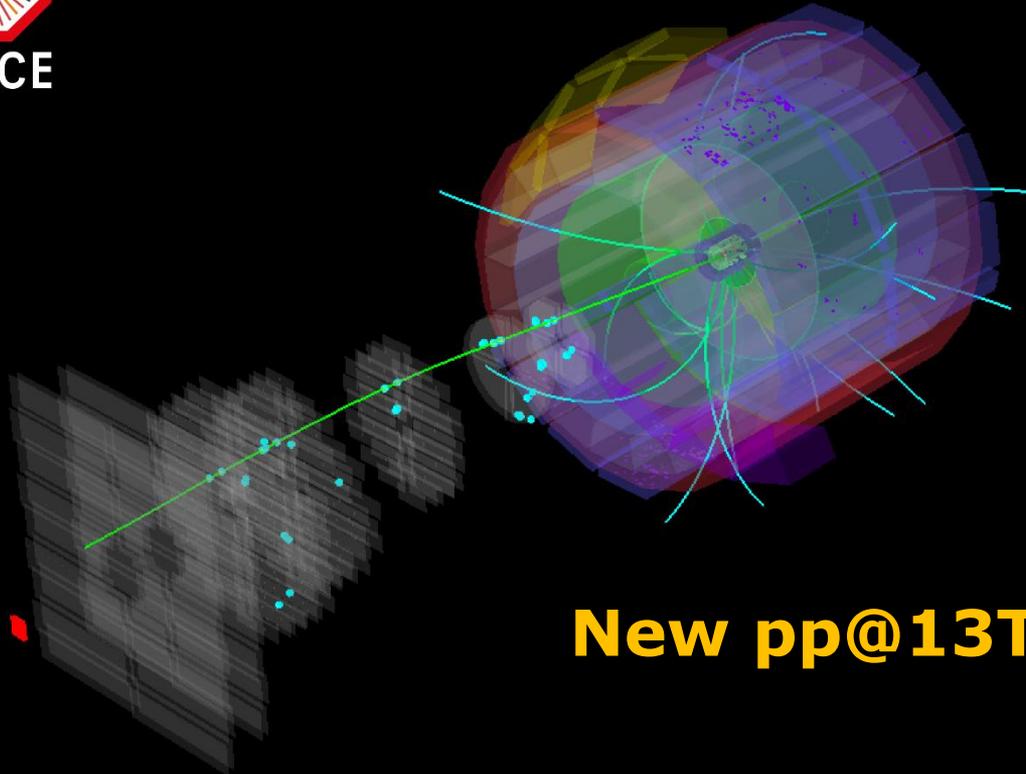
➔ **Results from LHC Run2 eagerly awaited!**

- Energy increase ($\sqrt{s_{NN}}=5\text{TeV}$) will allow for confirmation of the (re) combination role at low p_T
- Statistics increase will allow to sharpen Run-I results

AND NOW. . .LHC RUN-II

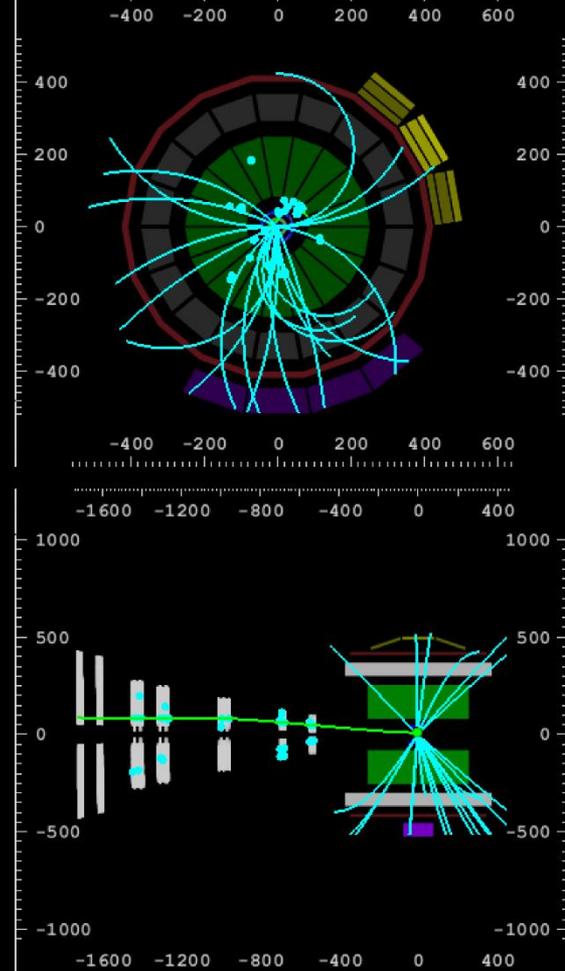


ALICE



New pp@13TeV!!!

Run: 223327
LHC fill: 3746
Timestamp: 2015-05-21 09:30:17 (UTC)

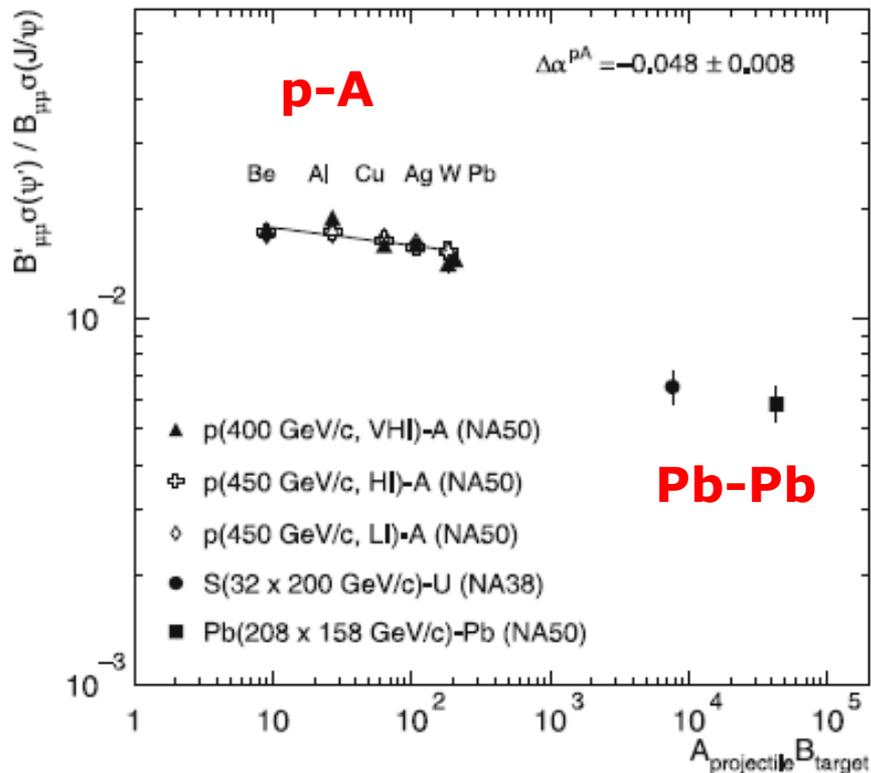


Backup slides

LOW ENERGY RESULTS: $\psi(2S)$ FROM SPS & RHIC

→ SPS (NA50) pA, AA @ $\sqrt{s_{NN}} = 17$ GeV

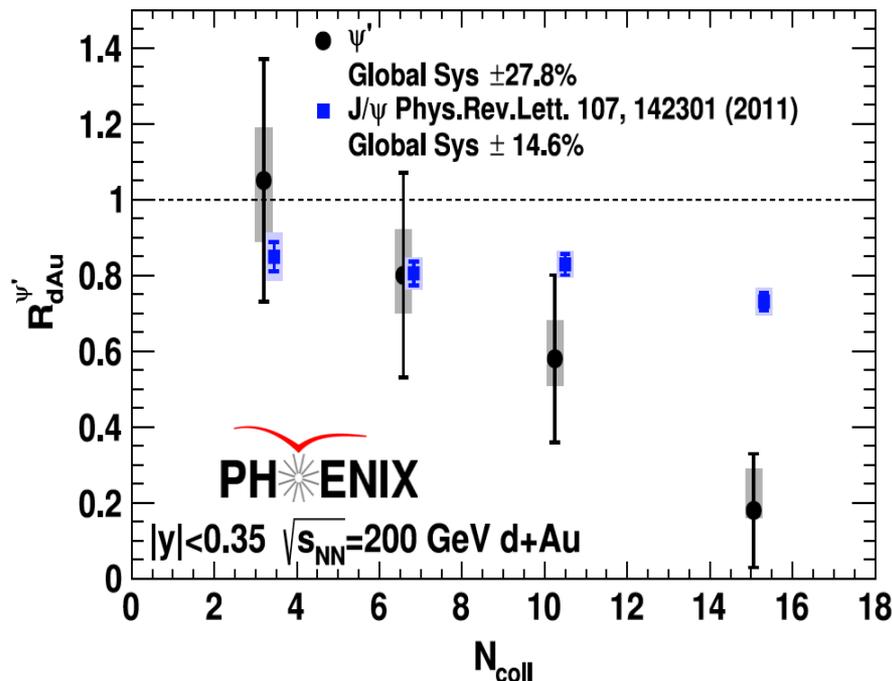
Eur. Phys. J. C 49, 559 (2007)



→ $\psi(2S)$ is more suppressed than J/ψ already in pA collisions and the suppression increases in Pb-Pb

→ RHIC (PHENIX)
d-Au @ $\sqrt{s_{NN}} = 200$ GeV

PRL 111, 202301 (2013)

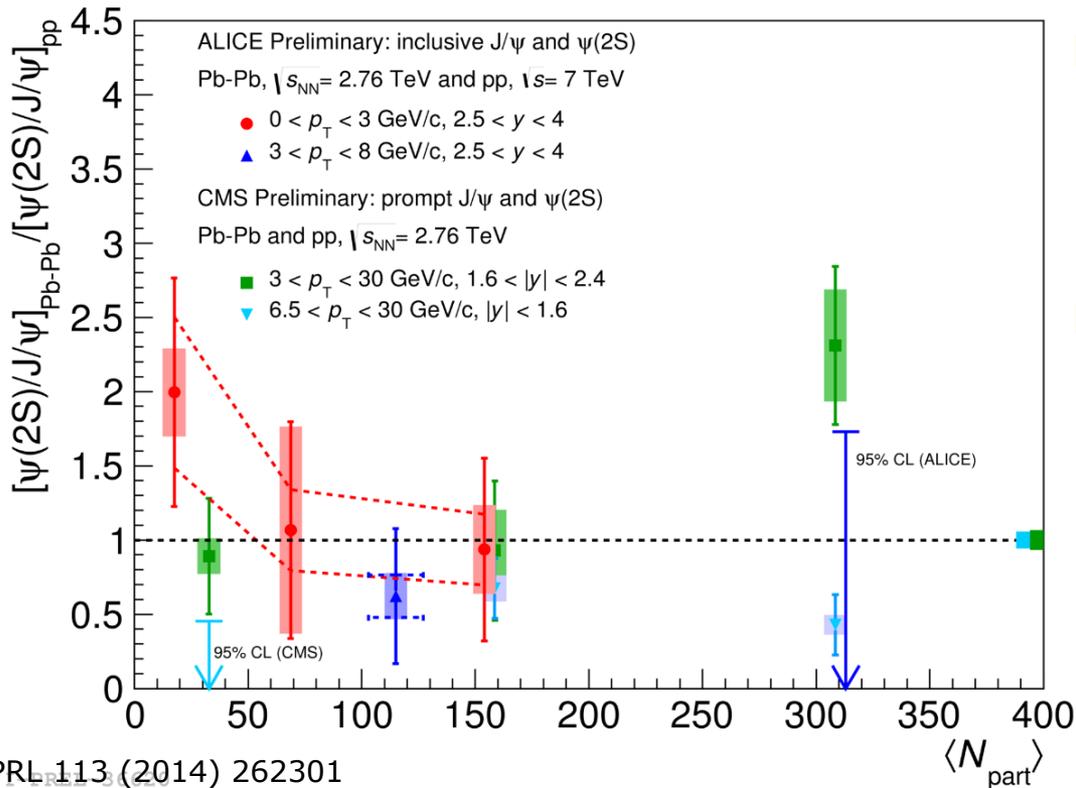


→ unexpected $\psi(2S)$ suppression, stronger than the J/ψ one in d-Au

$\psi(2S)/J/\psi$ IN Pb-Pb @LHC

Being a more weakly bound state than the J/ψ , the $\psi(2S)$ is another interesting probe to investigate charmonium behaviour in the medium

The $\psi(2S)$ yield is compared to the J/ψ one in Pb-Pb and in pp



ALICE: reference pp@ $\sqrt{s}=7$ TeV

low p_T ($0 < p_T < 3$ GeV/c) →
 $\psi(2S)$ more suppressed than J/ψ

CMS: reference pp@ $\sqrt{s}=2.76$ TeV

$p_T > 3$ GeV/c & $1.6 < |y| < 2.4$ →
 $\psi(2S)$ less suppressed than J/ψ

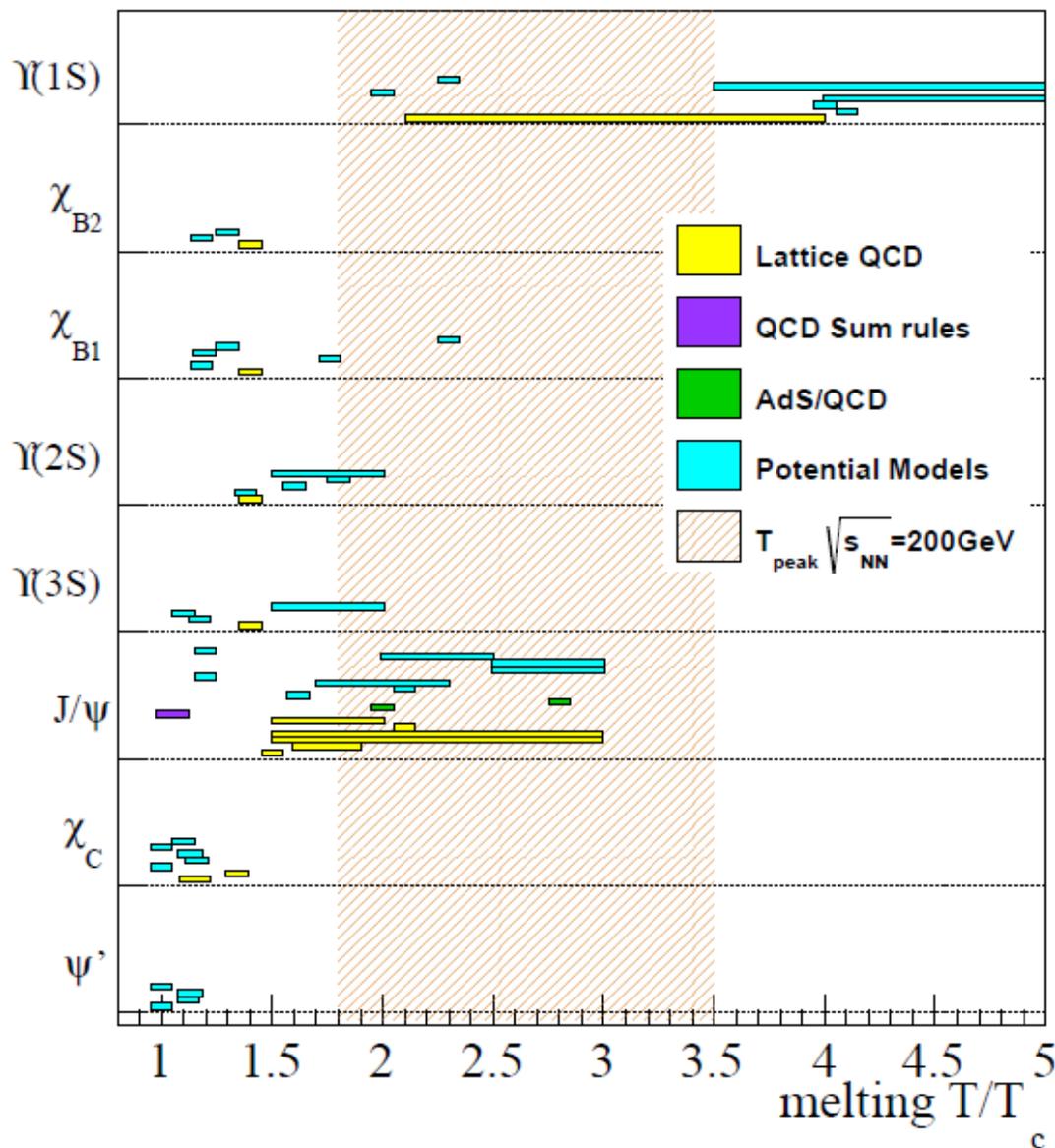
$p_T > 6.5$ GeV/c & $|y| < 1.6$ →
 $\psi(2S)$ more suppressed than J/ψ

PRL 113, (2014) 262301

Improved agreement between ALICE and CMS data (new pp CMS reference)

Large statistics and systematic uncertainties prevent a firm conclusion on the $\psi(2S)$ trend vs centrality

DISSOCIATION TEMPERATURES



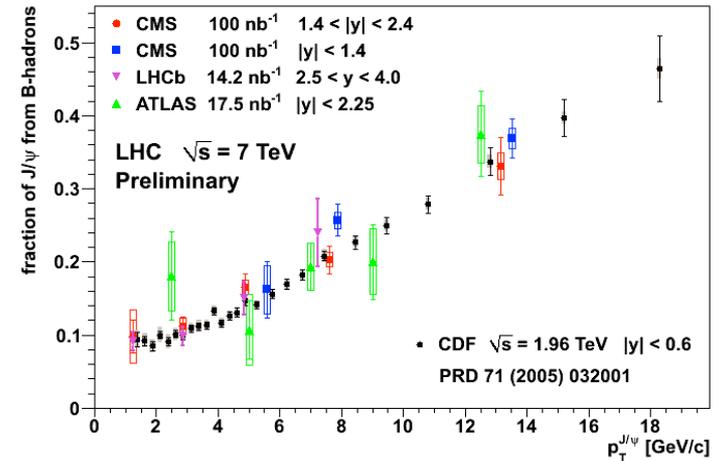
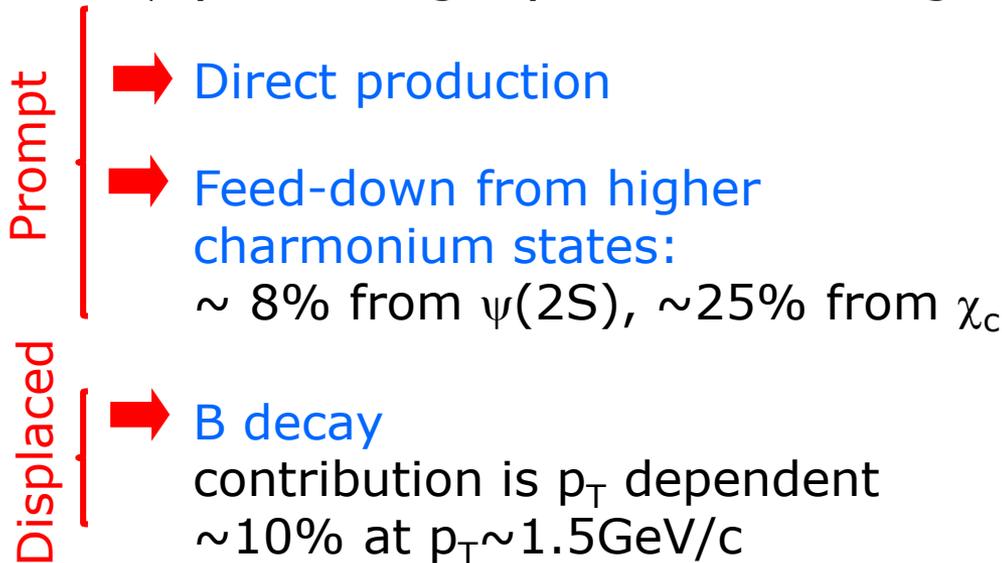
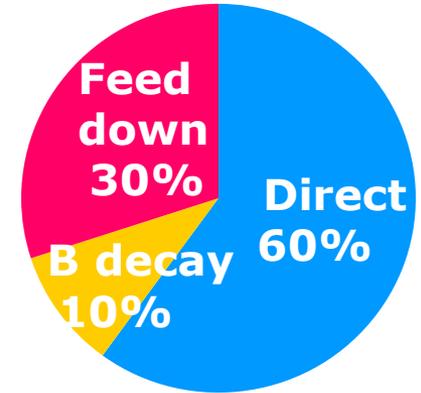
Quarkonium production and decay

J/ψ production

Quarkonium production can proceed:

- directly in the interaction of the initial partons
- via the decay of heavier hadrons (feed-down)

For J/ψ (LHC energies) the contributing mechanisms are:



J/ψ decay

J/ψ can be studied through its decays:

$$J/\psi \rightarrow \mu^+\mu^-$$

$$J/\psi \rightarrow e^+e^-$$

(~6% branching ratio)

Nuclear shadowing

PDF in nuclei are strongly modified with respect to those in a free nucleon

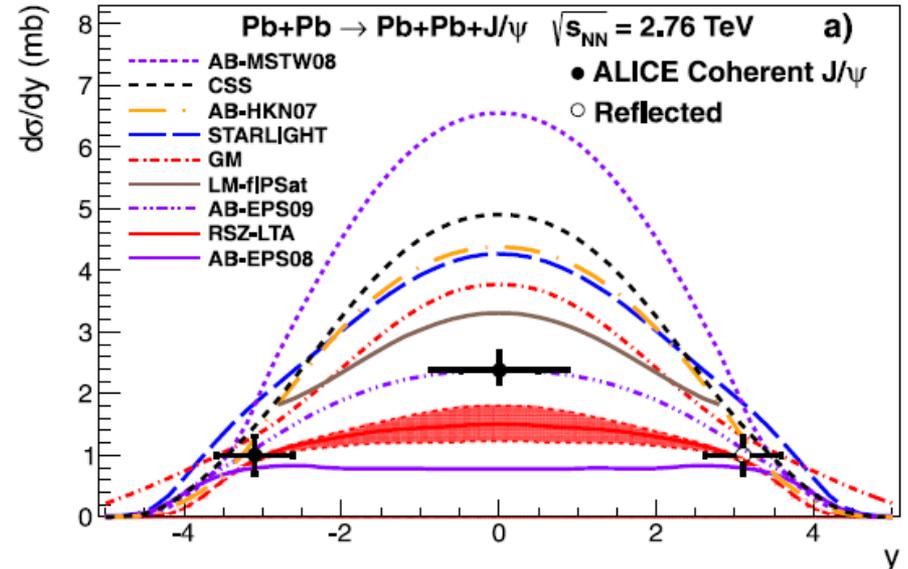
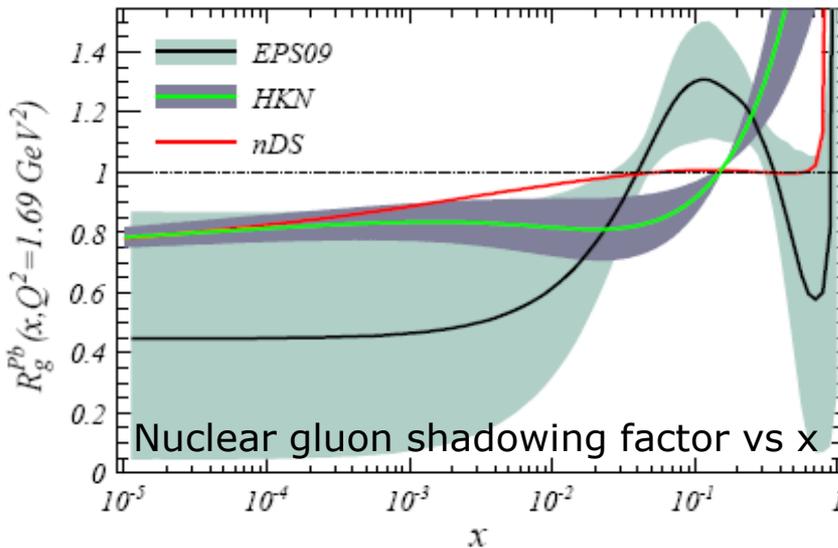
$$f_i^A(x, Q^2) = R_i(A, x, Q^2) \times f_i^p(x, Q^2)$$

nPDF: PDF of proton in a nucleus

free proton PDF

$$R_i(A, x, Q^2) = \frac{f_i^A(x, Q^2)}{f_i^p(x, Q^2)}$$

LHC data cover the low x domain (Bjorken $x \sim 10^{-2} - 10^{-5}$)



J/ψ photoproduction cross section is a powerful tool to constrain gluon shadowing

ALICE: low p_T J/ψ

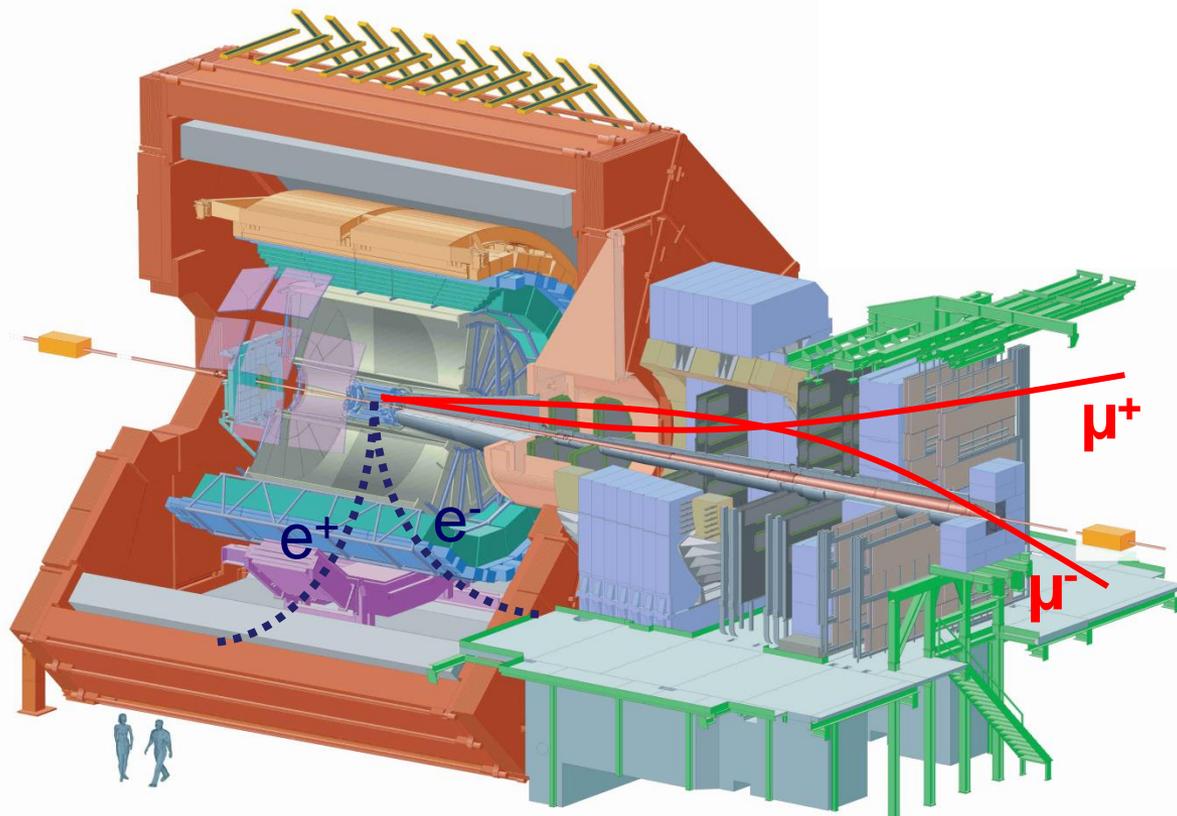
Central Barrel
($|y_{\text{LAB}}| < 0.9$)

$J/\psi \rightarrow e^+e^-$

Electrons tracked using ITS and TPC
Particle identification: TPC, TOF, TRD

Forward muon arm $J/\psi \rightarrow \mu^+\mu^-$
($2.5 < y_{\text{LAB}} < 4$)

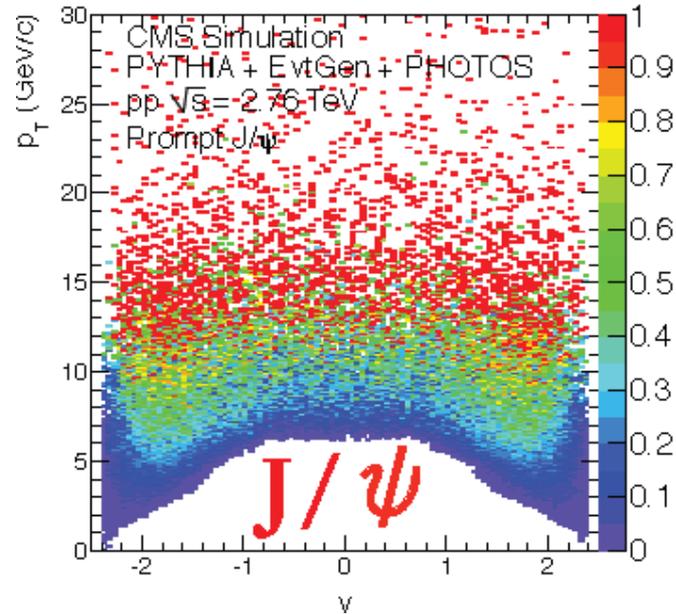
Muons identified and tracked in the
muon spectrometer



➔ Acceptance coverage
in both y regions
down to zero p_T

➔ ALICE results refer
to inclusive J/ψ
production

CMS: high p_T J/ψ

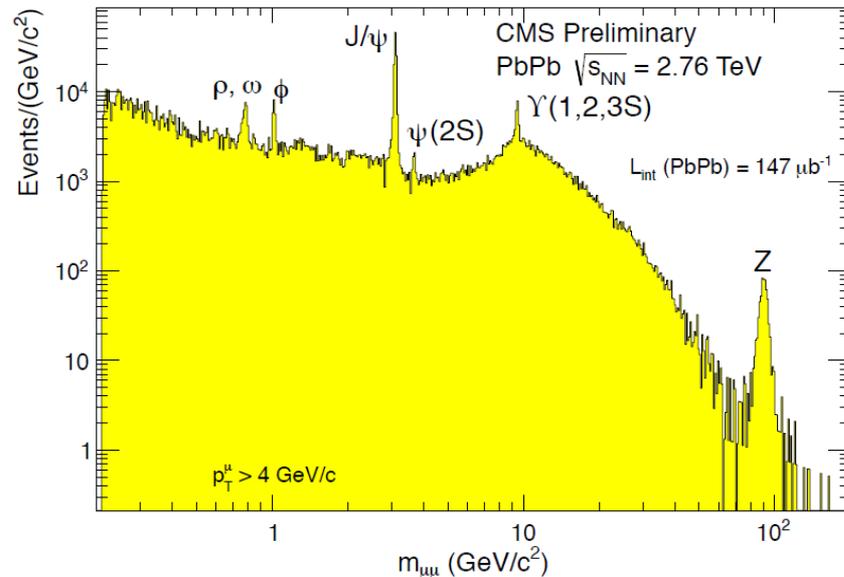
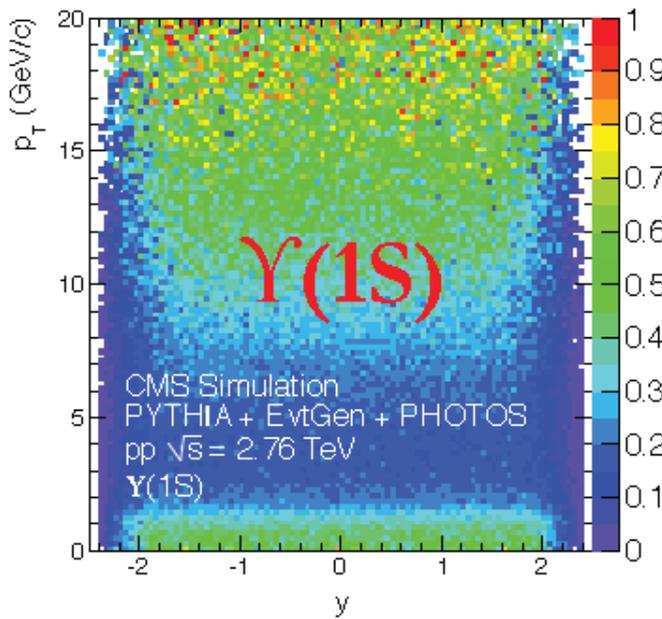


➔ Muons need to overcome the magnetic field and energy loss in the absorber:
 $p_{\min} \sim 3\text{-}5$ GeV/c to reach muon stations

➔ Limits J/ψ acceptance:

- Midrapidity: $p_T > 6.5$ GeV/c
- Forward rapidity: $p_T > 3$ GeV/c

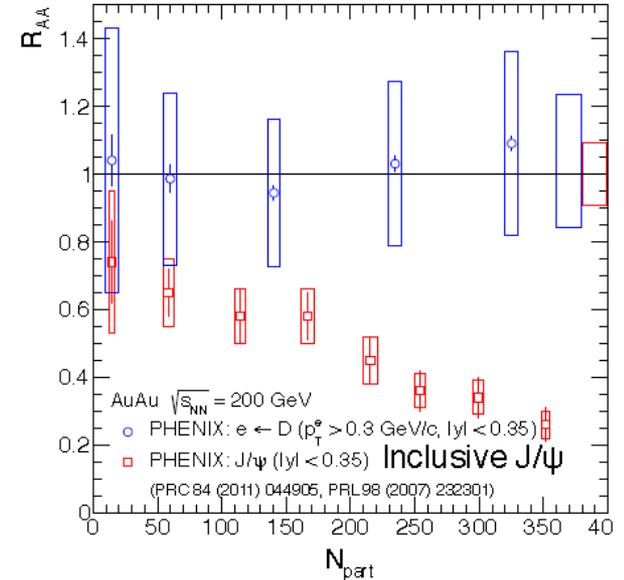
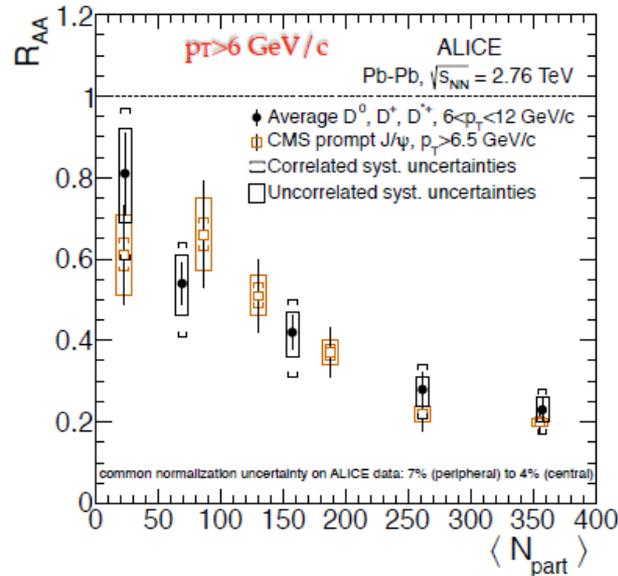
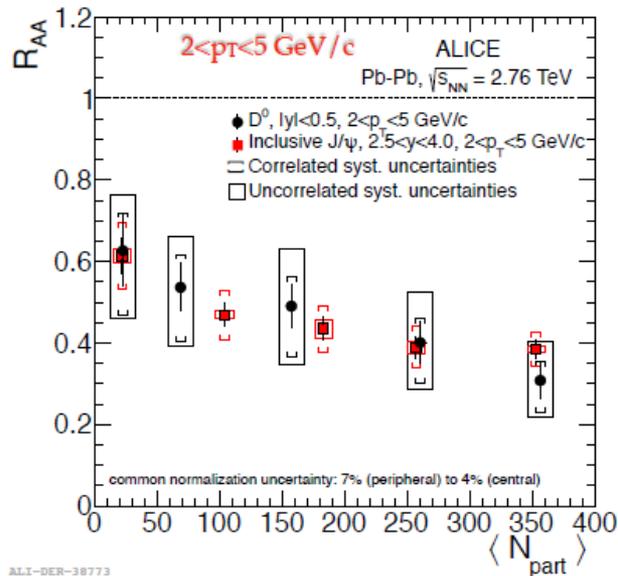
..but not the Υ one ($p_T > 0$ everywhere)



➔ Prompt and B-decay J/ψ are measured

J/ψ vs D in AA collisions

➔ Open charm should be a very good reference to study J/ψ suppression (à la Satz)



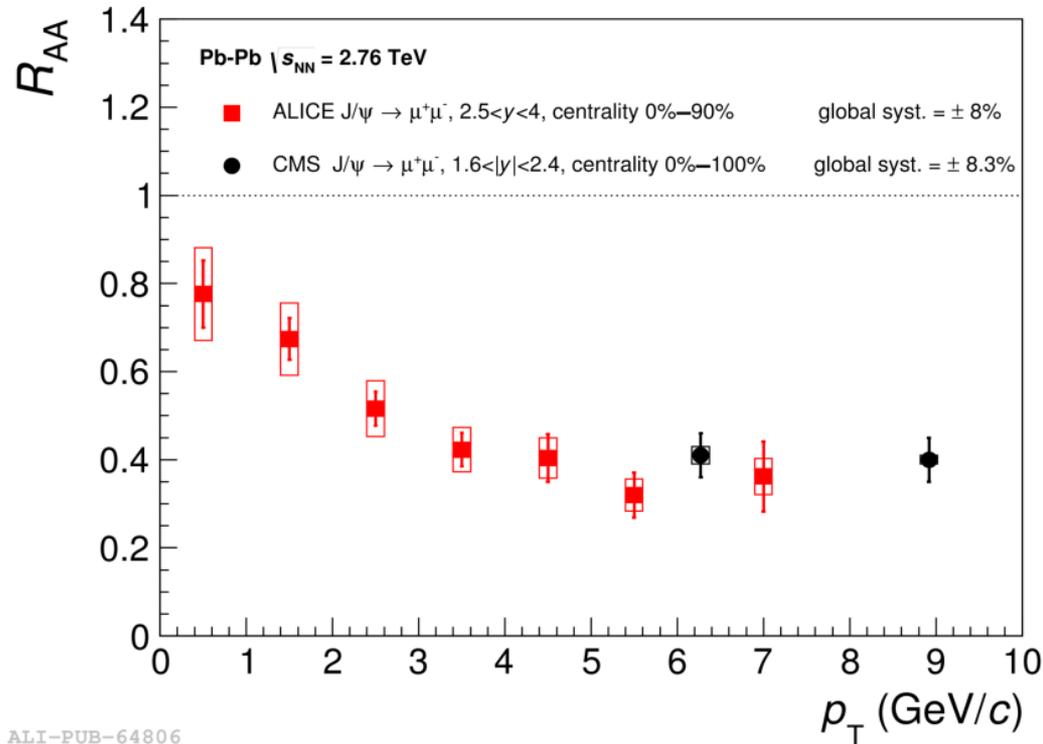
➔ Interesting comparison between ALICE and CMS J/ψ compared to D

Caveat:
complicate to compare J/ψ and D R_{AA} at LHC because of restricted kinematic regions.
Low p_T D not accessible for the moment

➔ Different trend observed at low p_T at RHIC.
At high p_T trend is similar to the LHC one

CMS: high p_T J/ ψ

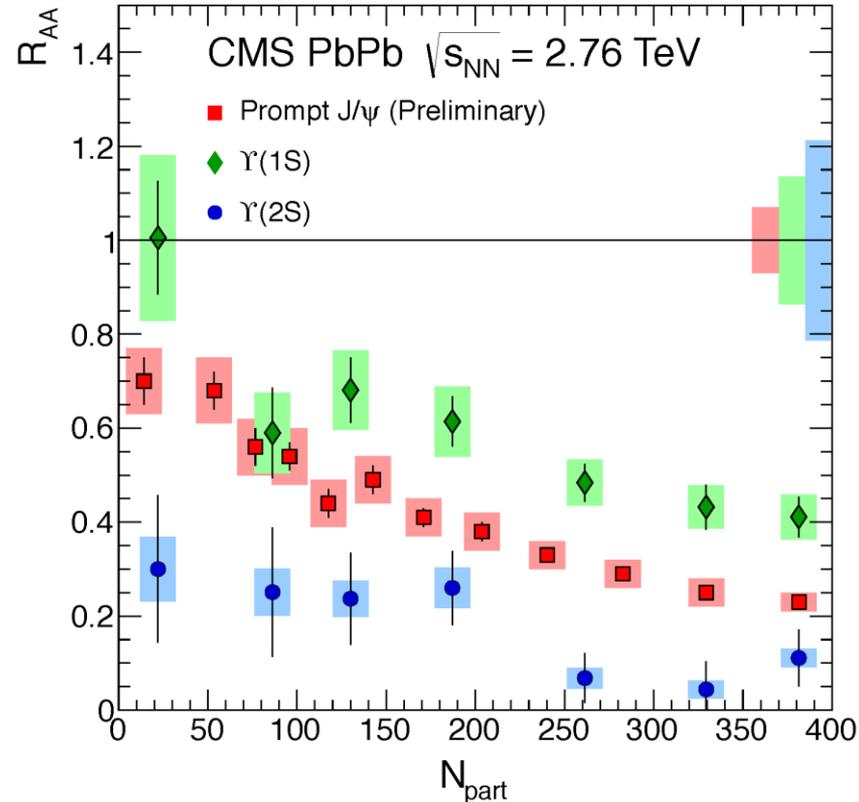
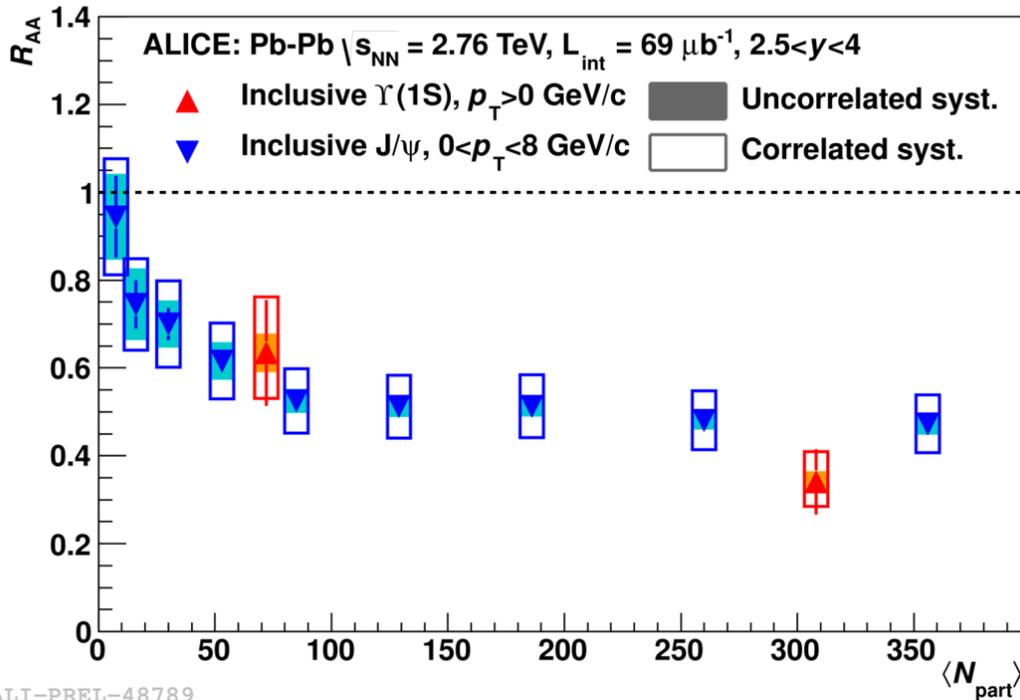
➔ The high p_T region can be investigated by CMS!



ALI-PUB-64806

➔ Good agreement with ALICE (at high p_T) in spite of the different rapidity range

Comparison Υ and J/ψ



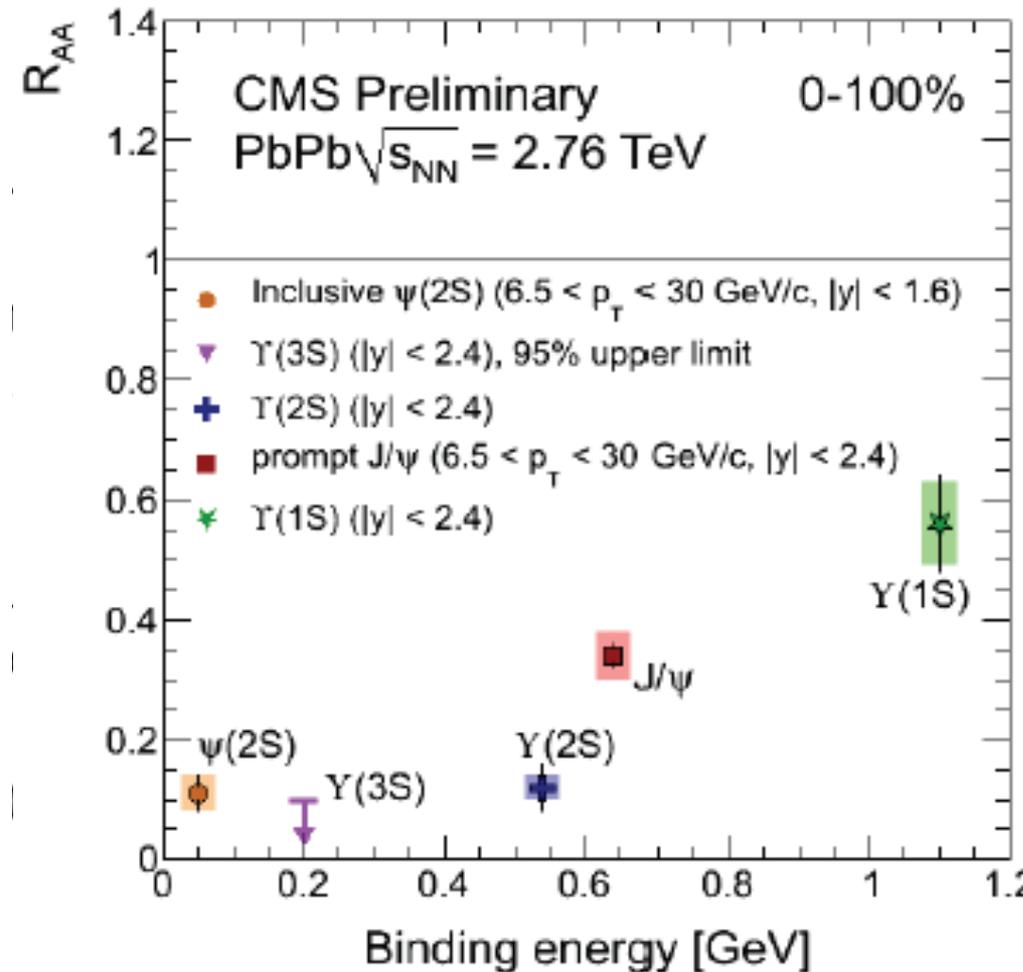
➔ Similar R_{AA} for low p_T inclusive J/ψ and $\Upsilon(1S)$

➔ Sequential suppression observed for prompt J/ψ and $\Upsilon(nS)$ at high p_T

➔ interplay of the competing mechanisms for J/ψ and Υ can be different and dependent on kinematics!

Where are we?

27 years after first suppression prediction, this is finally observed also in the Υ sector with very good accuracy!



Two main mechanisms at play:

1. Suppression in a deconfined medium
2. Re-combination (for charmonium) at high \sqrt{s}

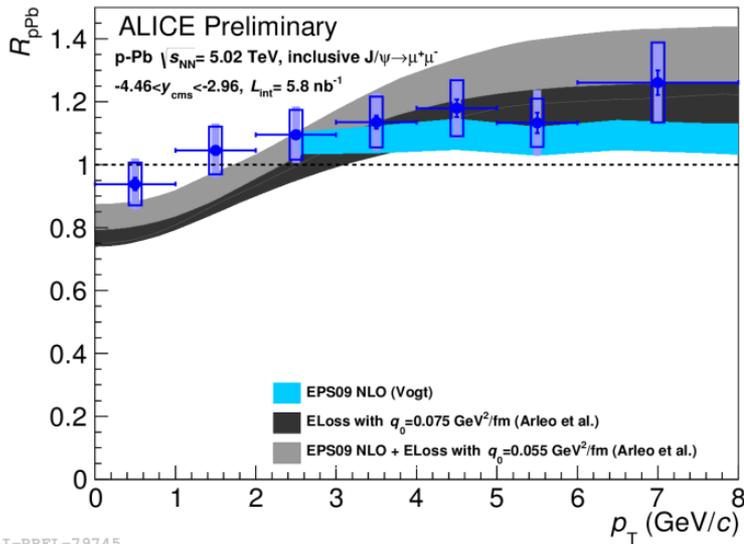
can qualitatively explain the main features of the results

R_{AA} vs binding energy: looser bound states more suppressed than the tighter ones

however hot and cold effects not yet disentangled...need pA results!

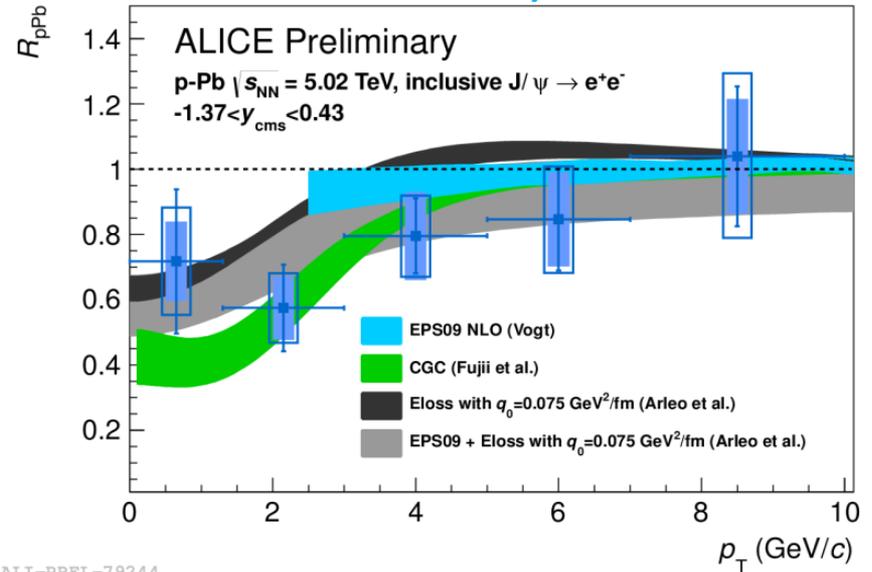
p-Pb: ROLE OF CNM EFFECTS ON J/ψ

backward-y



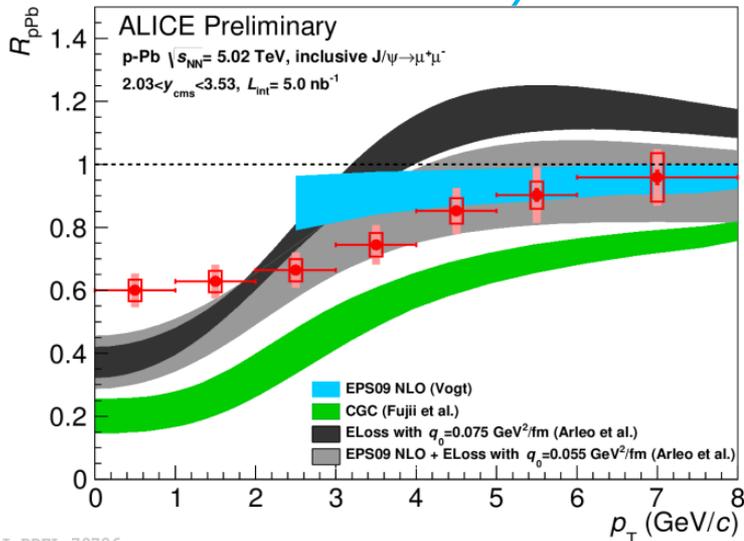
ALI-PREL-79745

mid-y



ALI-PREL-79244

forward-y

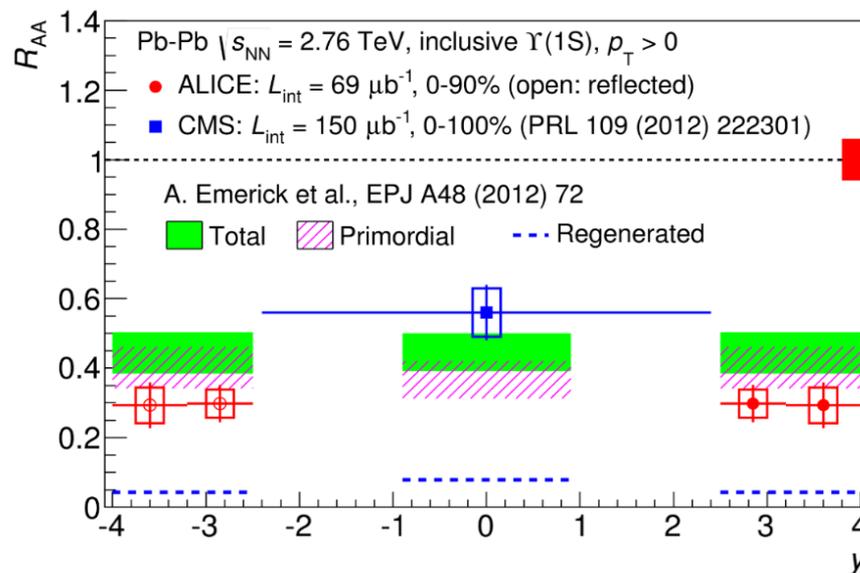
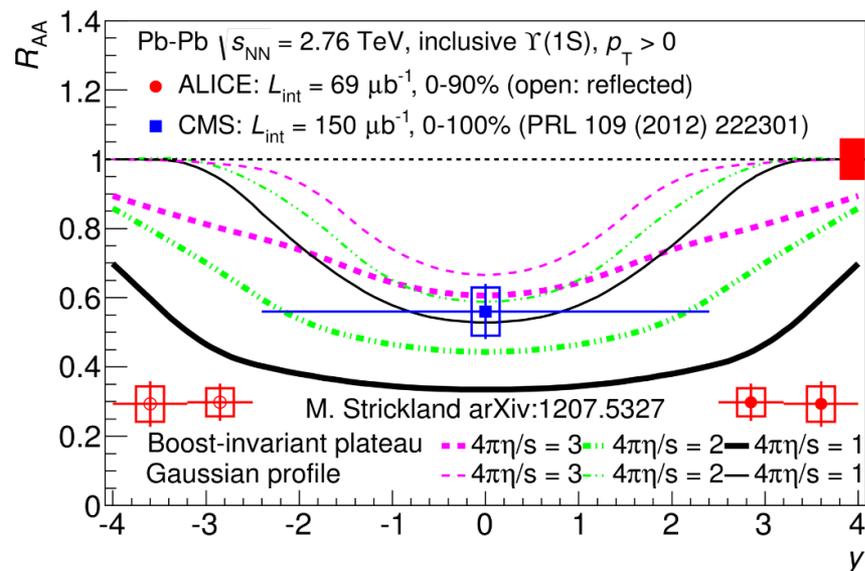


ALI-PREL-79726

➔ R_{pA} p_T dependence:

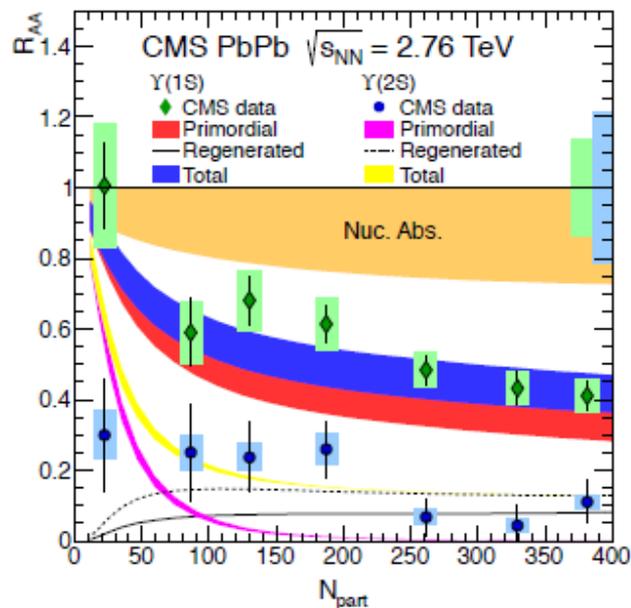
fair agreements with models based on shadowing + energy loss except at forward-y and low p_T

COMPARISON WITH THEORY



ALI-PUB-85796

ALI-PUB-85792

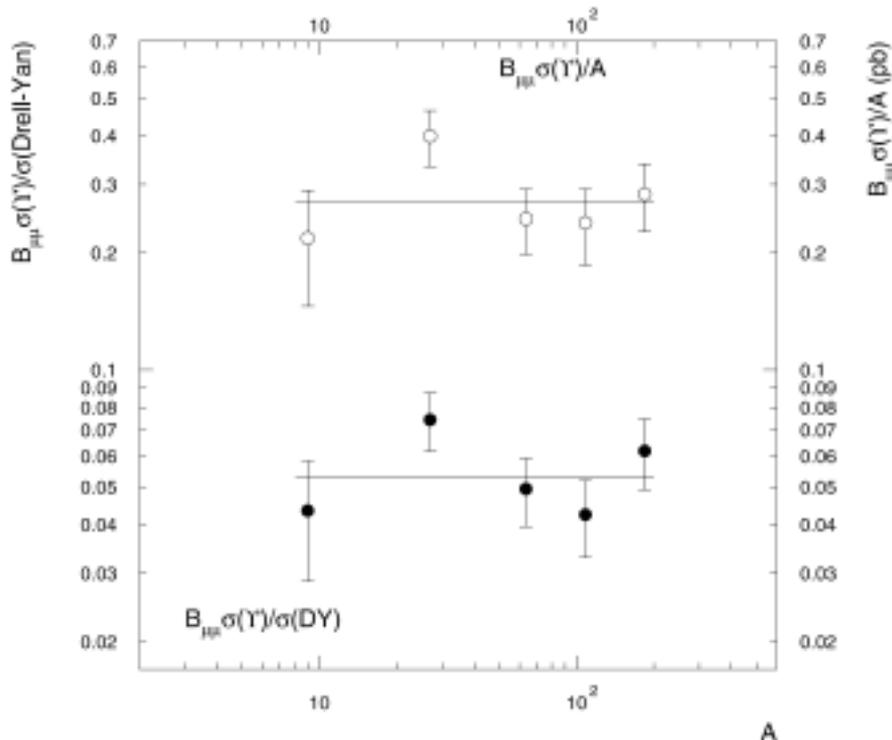


➔ Stronger suppression at forward rapidity (ALICE) than at mid-rapidity (CMS)

➔ Theory still meets difficulties in describing simultaneously the R_{AA} centrality and rapidity dependence (suppression slightly overestimated at forward-y, while better reproduced at mid-y)

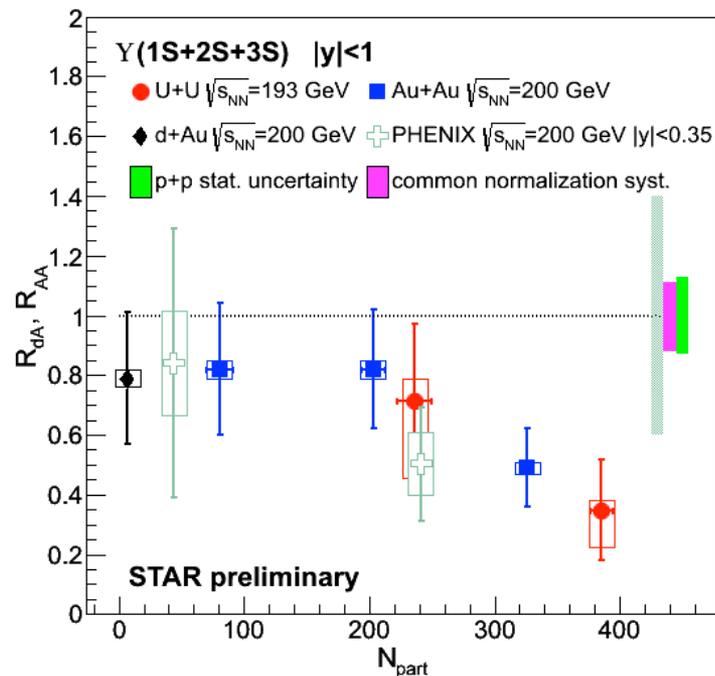
LOW ENERGY RESULTS: Υ FROM SPS & RHIC

→ SPS (NA50) pA, $\sqrt{s_{NN}}=29$ GeV



→ First Υ measurement at SPS energies.
Hint for no strong medium effects on $\Upsilon(1S+2S+3S)$ in pA

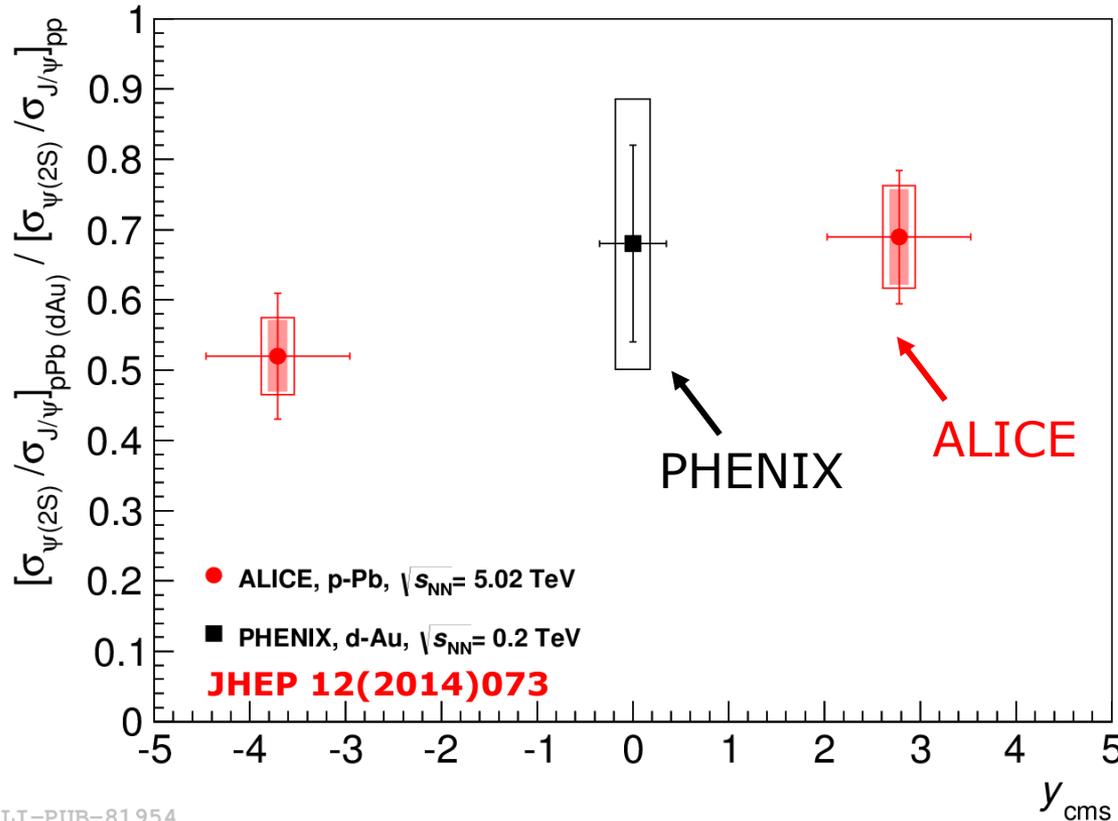
→ RHIC (PHENIX, STAR) dAu, Au-Au $\sqrt{s_{NN}} = 200$ GeV



→ ΥR_{AA} compatible with suppression of excited states but large uncertainties prevents further insights

$\psi(2S)/J/\psi$ IN p-Pb

➔ A strong decrease of the $\psi(2S)$ production in p-Pb, relative to J/ψ , is observed with respect to the pp measurement ($2.5 < y_{\text{cms}} < 4$, $\sqrt{s} = 7\text{TeV}$)



➔ Double ratio allows a direct comparison of the J/ψ and $\psi(2S)$ production yields between experiments

➔ Similar effect seen by PHENIX in d-Au collisions, at mid- y , at $\sqrt{s_{\text{NN}}} = 200\text{ GeV}$

ALI-PUB-81954

$[\psi(2S)/J/\psi]_{\text{pp}}$ variation between ($\sqrt{s} = 7\text{TeV}$, $2.5 < y < 4$) and ($\sqrt{s} = 5.02\text{TeV}$, $2.03 < y < 3.53$ or $-4.46 < y < -2.96$) based on CDF and LHCb data ($\sim 8\%$ included in the systematic uncertainty) **61**

$\psi(2S)$ R_{pA} VS RAPIDITY

➔ $\psi(2S)$ suppression stronger than the J/ψ one, reaching a factor ~ 2 wrt pp

can the stronger $\psi(2S)$ suppression be due to break-up of the fully formed resonance in CNM?

➔ possible if:

formation (τ_f) < crossing time (τ_c)

forward- y :

$\tau_c \sim 10^{-4}$ fm/c

backward- y :

$\tau_c \sim 10^{-1}$ fm/c

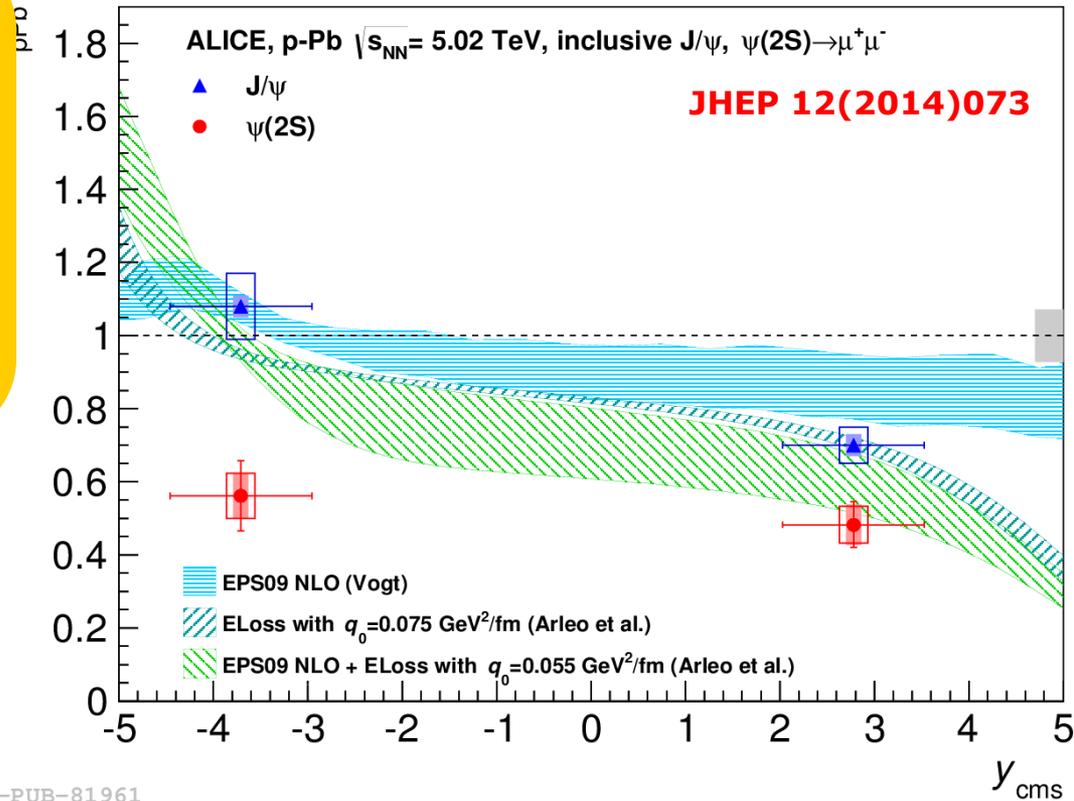
while $\tau_f \sim 0.05-0.15$ fm/c

➔ **forward- y :**

break-up effects excluded

➔ **backward- y :**

$\tau_f \sim \tau_c$, hence break-up in CNM hardly explains the strong J/ψ and $\psi(2S)$ difference



.I-PUB-81961

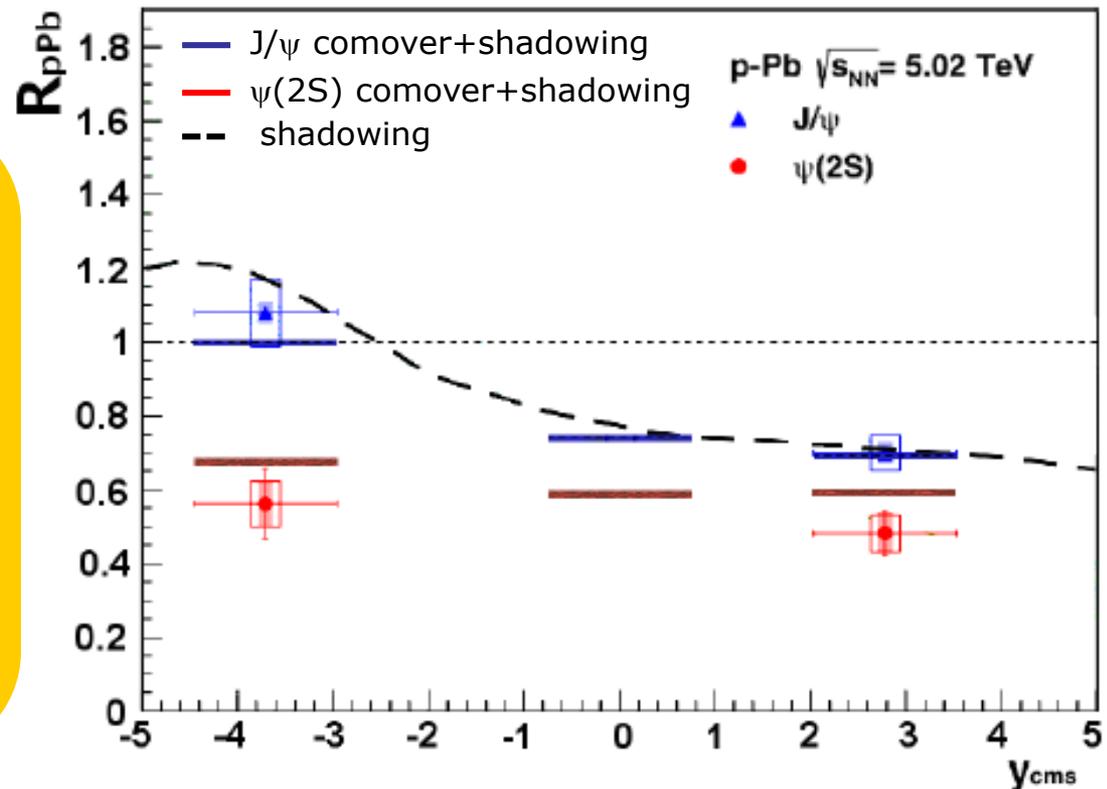
y_{cms}

$\psi(2S)$ R_{pA} VS RAPIDITY: COMOVERS?

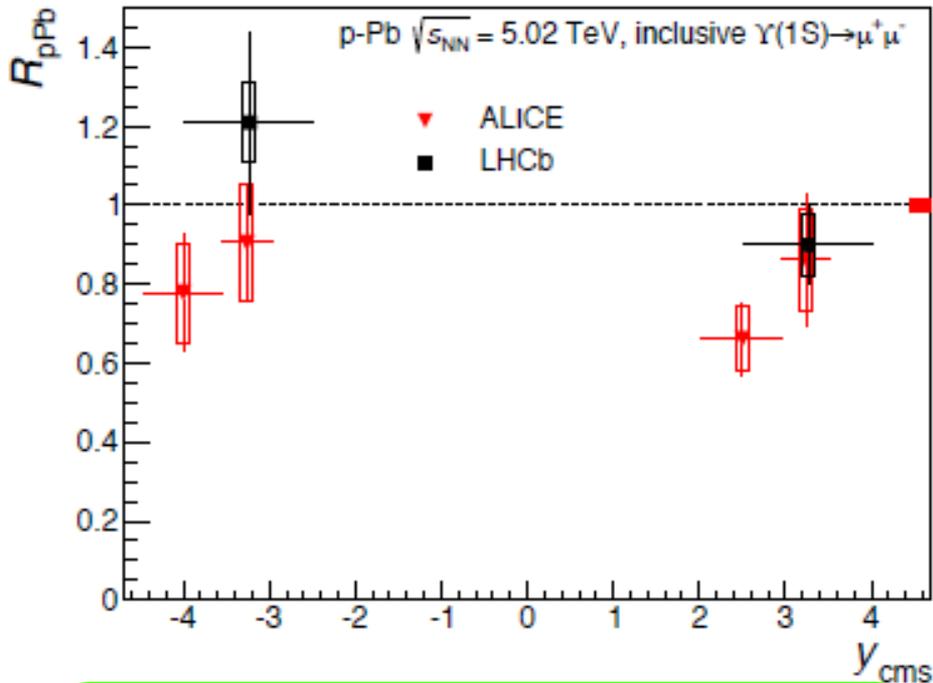
➔ Final state effects related to the (hadronic) medium created in the p-Pb collisions?

➔ Charmonium interaction with comoving particles:

- Comovers dissociation affects more strongly the loosely bound $\psi(2S)$ than the J/ψ
- Comovers density larger at backward rapidity



$\Upsilon(1S)$ PRODUCTION IN p-Pb

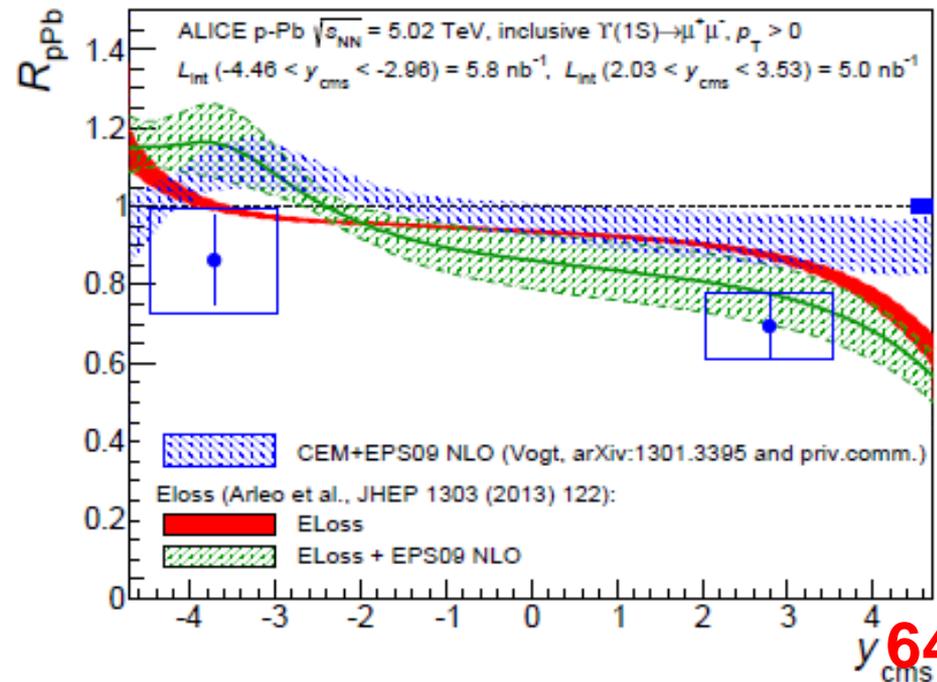


→ $\Upsilon(1S)$ measured at mid-y by CMS and at forward-y by both ALICE and LHCb

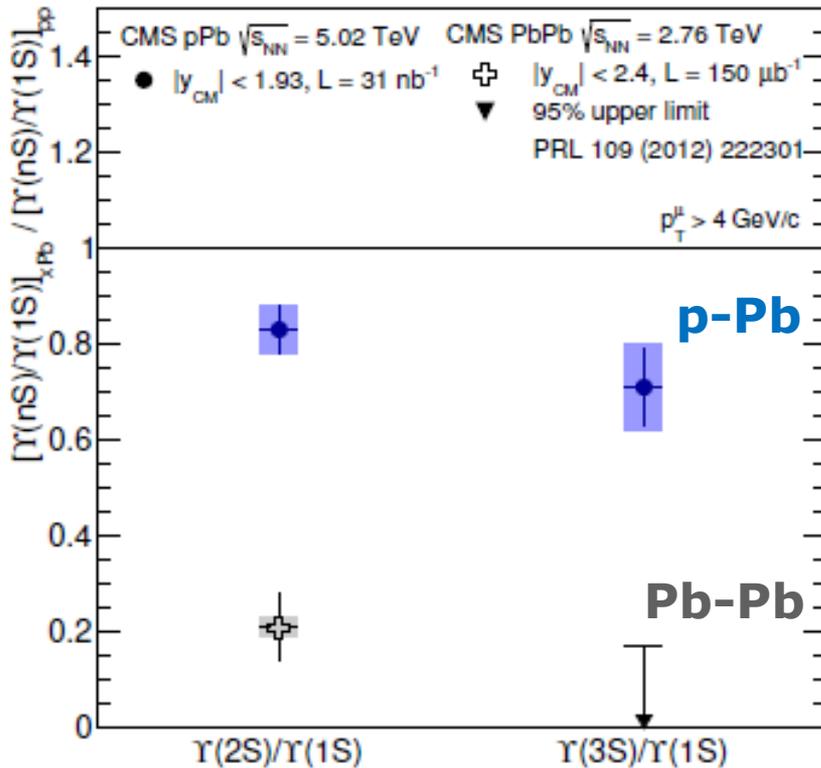
→ Compatible R_{pA} results within uncertainties (but LHCb systematically higher)

→ Hint for stronger suppression at forward-y (similarly to J/ψ)

→ Theoretical calculations based on initial state effects seem not to describe simultaneously forward and backward y



$\Upsilon(nS)/\Upsilon(1S)$ PRODUCTION IN p-Pb



Initial state effects similar for the three Υ states

p-Pb vs pp @mid-y:
different/stronger final states effects in p-Pb affecting the excited states

p-Pb vs PbPb @mid-y :
even stronger suppression of excited states in PbPb

CMS HIN-13-003, JHEP 04 (2014) 103, PRL 109 (2012)

ALICE (and LHCb) observes:

$\Upsilon(2S)/\Upsilon(1S)$ (ALICE)
 $2.03 < y < 3.53$: $0.27 \pm 0.08 \pm 0.04$
 $-4.46 < y < -2.96$: $0.26 \pm 0.09 \pm 0.04$

Compatible with pp results
 0.26 ± 0.08 (ALICE, pp@7TeV)

CMS analyses the double ratio $[\Upsilon(2S)/\Upsilon(1S)]/[\Upsilon(nS)/\Upsilon(1S)]_{pp}$ and finds

$0.83 \pm 0.05 \pm 0.05$