

Neutron Skins with α -Clusters

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NUSTAR Annual Meeting 2015

Outline

- **Introduction**

Neutron Skins and Neutron Matter Equation of State, Density Dependence of Symmetry Energy, Correlations in Low-Density Nuclear Matter

- **Generalized Relativistic Density Functional**

Details of gRDF Model, Effective Interaction, Neutron Star Matter

- **α -Clusters on the Surface of Nuclei**

Application of gRDF Model, Nuclei with α -Clusters

- **Experimental Test**

Quasi-Free ($p,p\alpha$) Knockout Reactions, Kinematics, Cross Sections

- **Conclusions**

Details:

S. Typel, Phys. Rev. C 89 (2014) 064321

S. Typel, H.H. Wolter, G. Röpke, D. Blaschke, Eur. Phys. J. A 50 (2014) 17

M.D. Voskresenskaya, S. Typel, Nucl. Phys. A 887 (2012) 42

S. Typel, G. Röpke, T. Klähn, D. Blaschke, H.H. Wolter, Phys. Rev. C 81 (2010) 015803

Introduction

Neutron Skins and Neutron Matter Equation of State

- neutron-rich nuclei \Rightarrow development of **neutron skin** with thickness:

$$\Delta r_{np} = S = \sqrt{\langle r_n^2 \rangle} - \sqrt{\langle r_p^2 \rangle}$$

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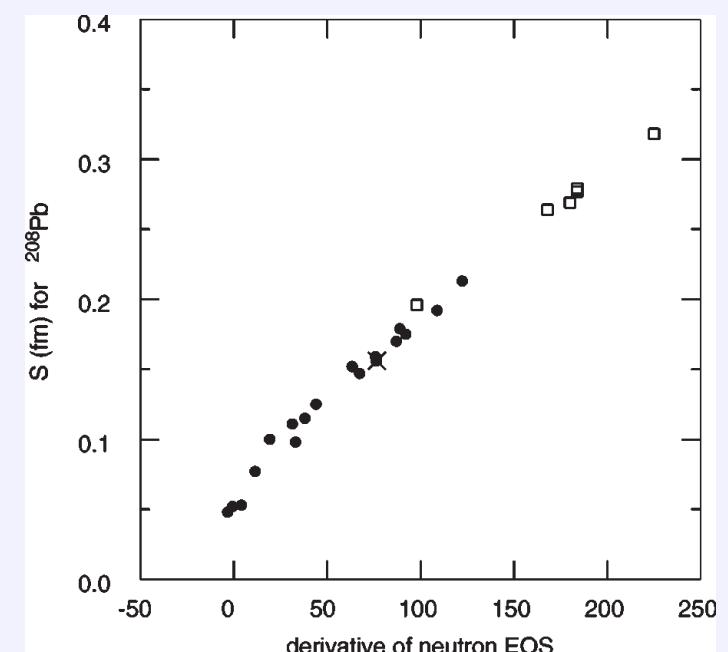
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$$\frac{d(E/A)}{dn} \Big|_{n=n_0} = \frac{p_0}{n_0^2} \quad \text{at density } n_0 = 0.1 \text{ fm}^{-3}$$

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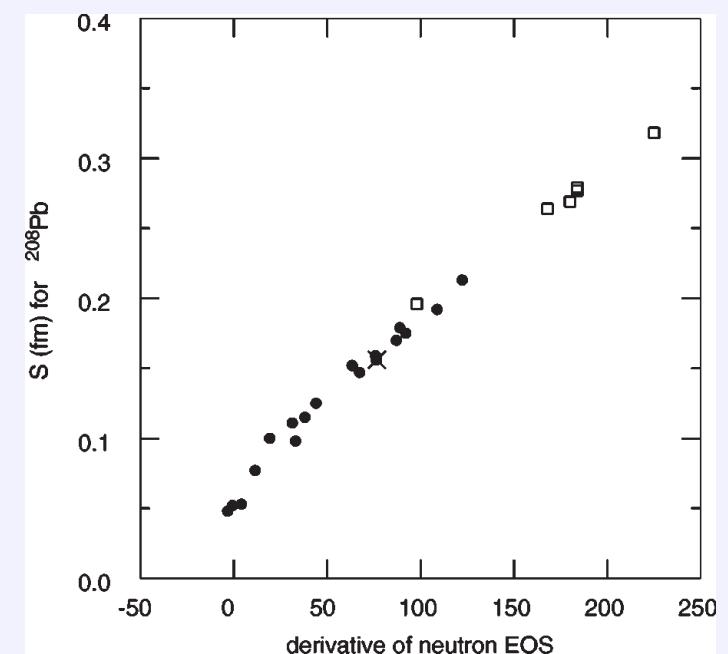
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- extension to relativistic mean-field models
(S. Typel and B. A. Brown, Phys. Rev. C 64 (2001) 027302)



Neutron Skins and Symmetry Energy

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 - symmetry energy at saturation

$$J = E_s(n_{\text{sat}})$$

○ slope coefficient

$$L = 3n \frac{d}{dn} E_s \Big|_{n=n_{\text{sat}}}$$

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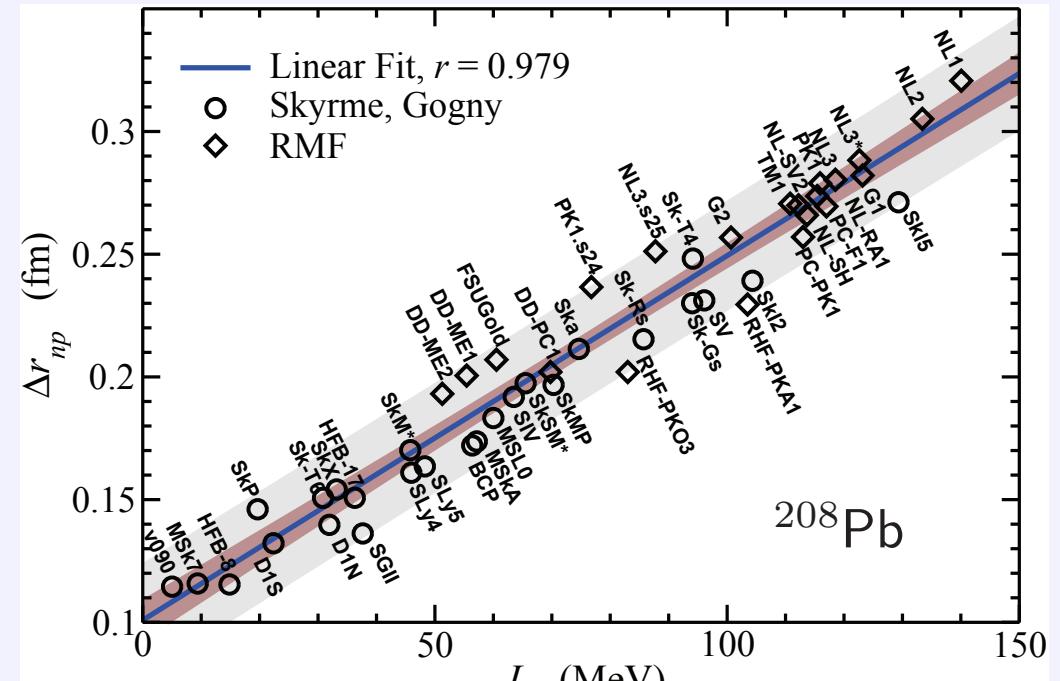
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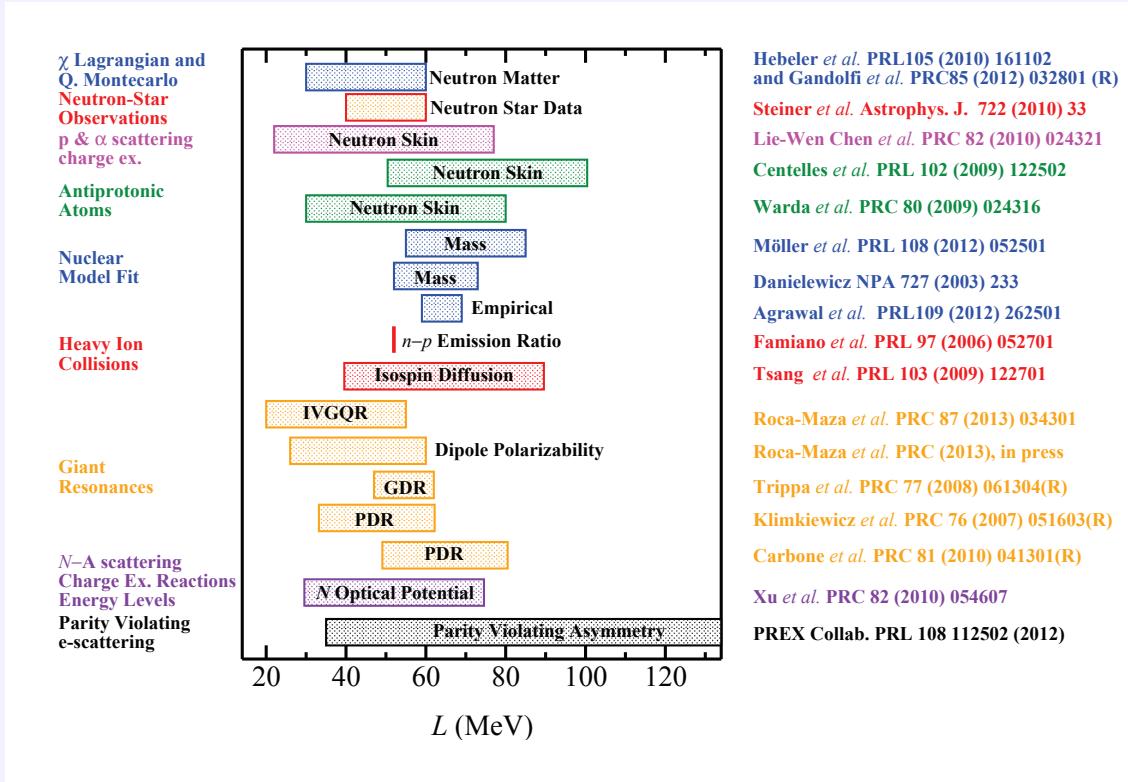
$$L = 3n \frac{d}{dn} E_s \Big|_{n=n_{\text{sat}}}$$

- determine L from experimental measurement of Δr_{np}
 \Rightarrow constraint for neutron matter EoS

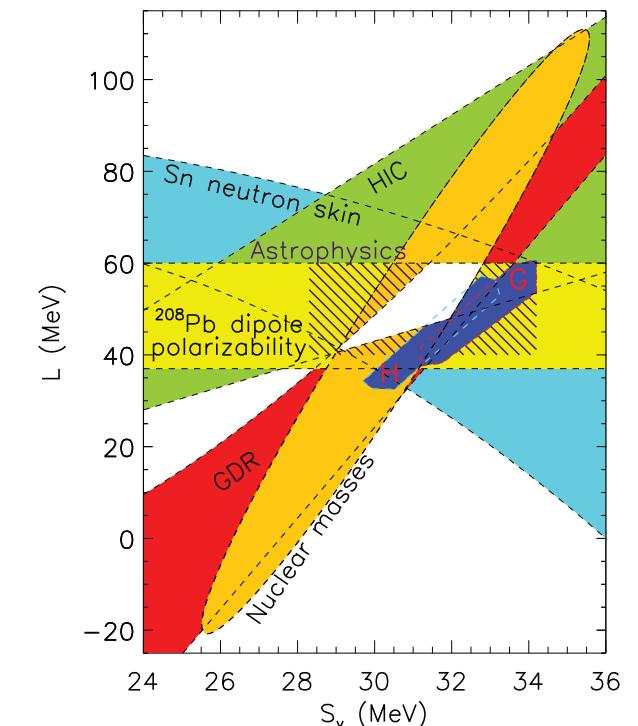


Symmetry Energy Parameters

- many attempts to determine symmetry energy $J = S_0 = S_v$ and slope coefficient L experimentally



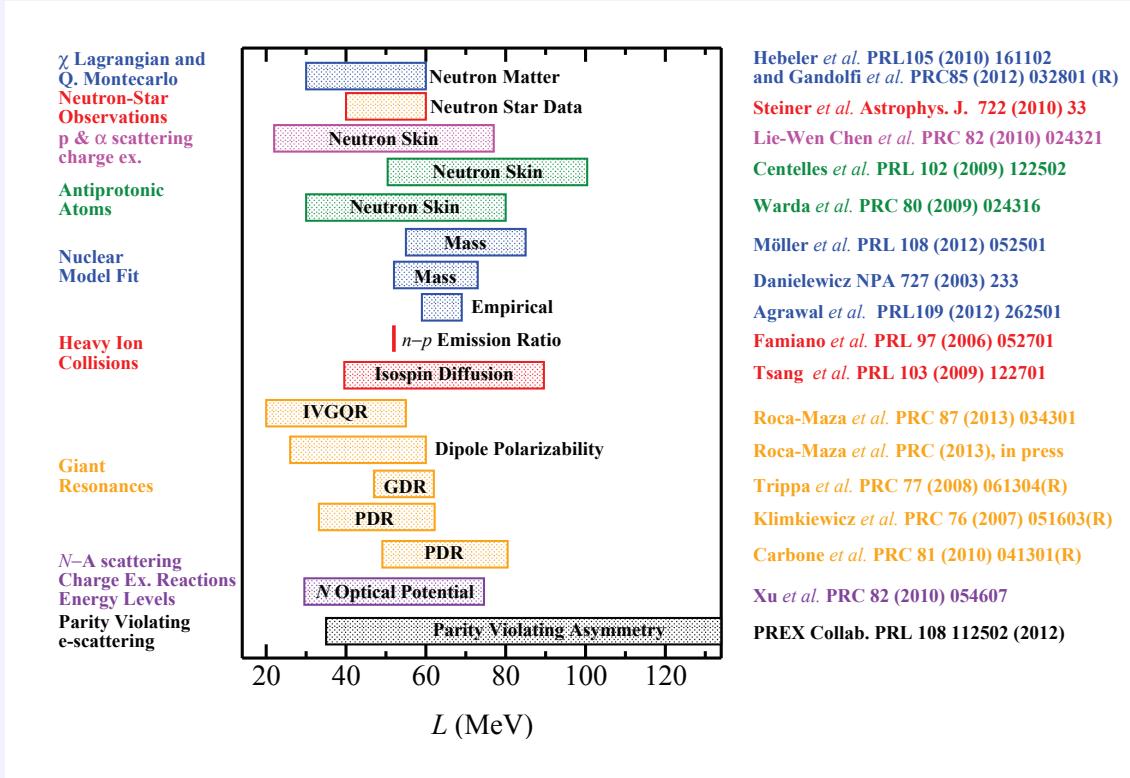
(X. Viñas *et al.*, Eur. Phys. J. A50 (2014) 27)



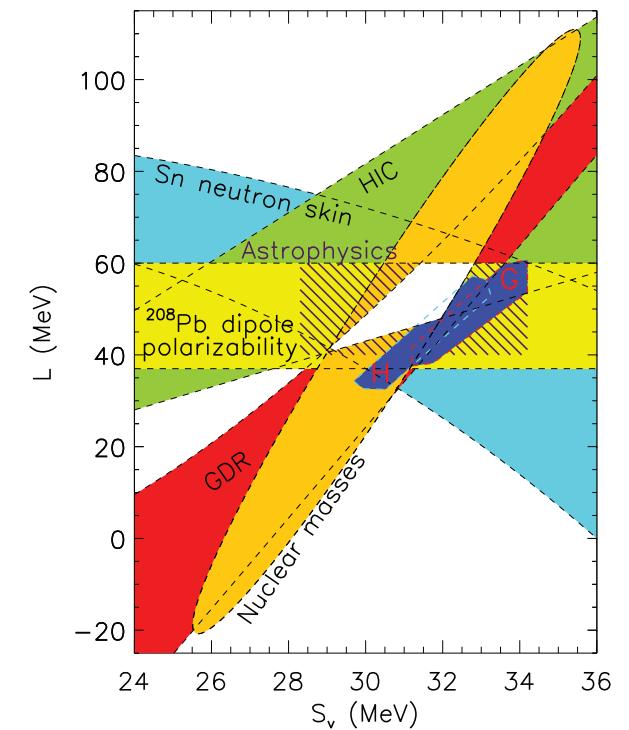
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(J.M. Lattimer, Y. Lim, ApJ. 771 (2013) 51)

- measurement of neutron skin thickness (e.g. PREX), correlation with L from mean-field calculations: effects of clustering correlations on surface of nuclei?

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 - study dependence of results on neutron excess and on isovector part of interaction
 - experimental test of predictions?

Generalized Relativistic Density Functional

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- **grand canonical approach**
 - extension of relativistic mean-field models with density-dependent meson-nucleon couplings
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- **selected model features**

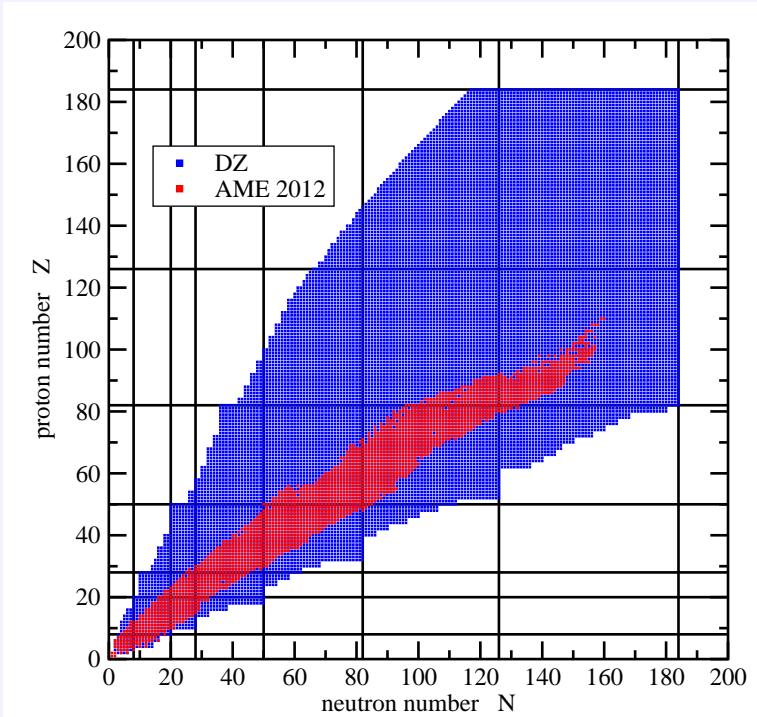
- [extended set of constituents](#): nucleons, light clusters (^2H , ^3H , ^3He , ^4He) and heavy nuclei

- experimental binding energies: AME 2012

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- extension: DZ10 predictions

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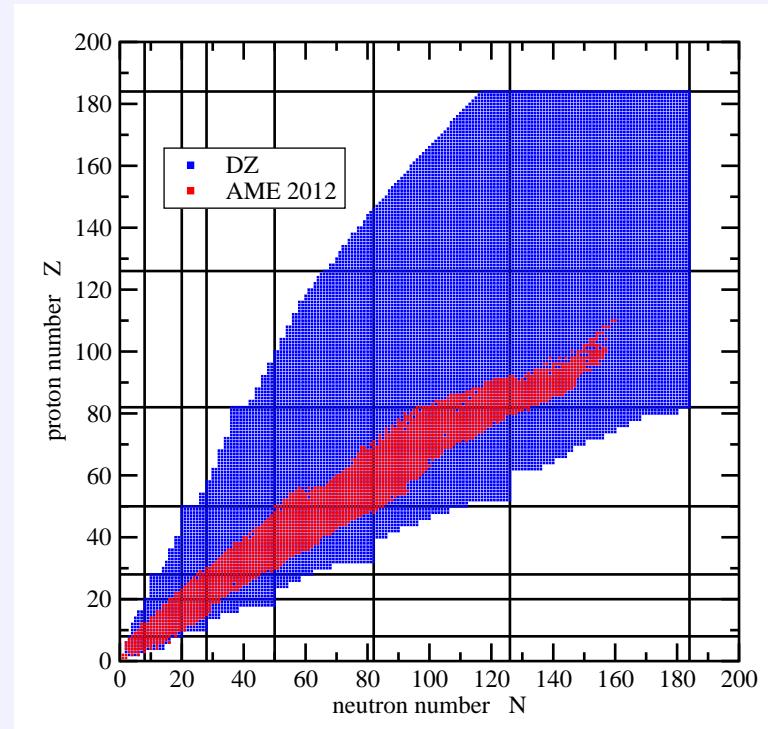
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- **medium modifications** of composite particles (mass shifts, internal excitations)
- scattering correlations considered (essential for **correct low-density limit**)
- **thermodynamically consistent** approach (⇒ "rearrangement" contributions)
- model parameters from fit to properties of finite nuclei



Effective Interaction

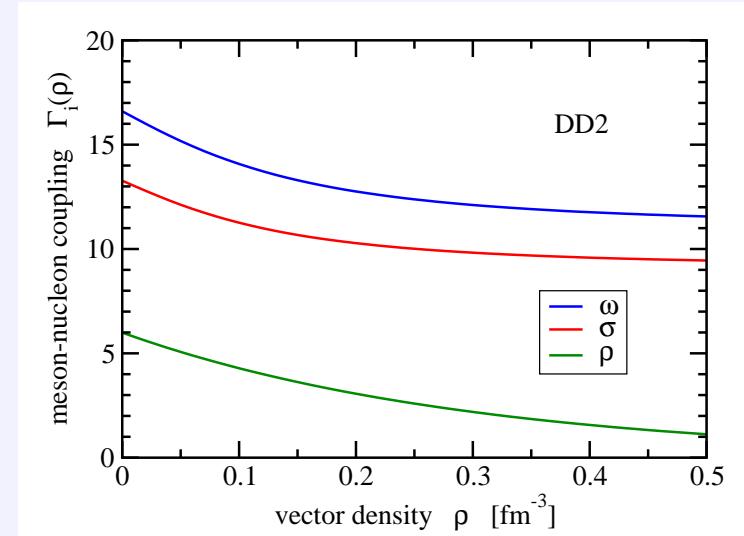
exchange of

- Lorentz scalar mesons $m \in \mathcal{S} = \{\sigma, \delta, \sigma_*, \dots\}$
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- represented by (classical) fields A_m with mass m_m
- coupling to constituents: $\Gamma_{im} = g_{im}\Gamma_m$
 - scaling factors g_{im}
 - density dependent $\Gamma_m = \Gamma_m(\varrho)$
 $\varrho = \sum_i (N_i + Z_i)n_i$ with parametrization DD2
(S. Typel et al., Phys. Rev. C 81 (2010) 015803)



nuclear matter parameters

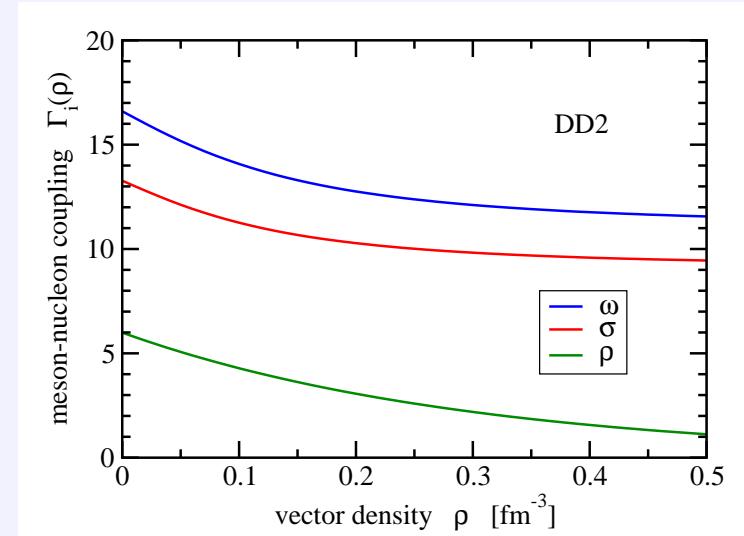
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- scalar potential $S_i = \sum_{m \in \mathcal{S}} \Gamma_{im} A_m - \Delta m_i$
with medium-dependent mass shift $\Delta m_i(T, n_j)$
 - from microscopic calculations
 - mainly action of Pauli principle

⇒ dissolution of clusters at high densities



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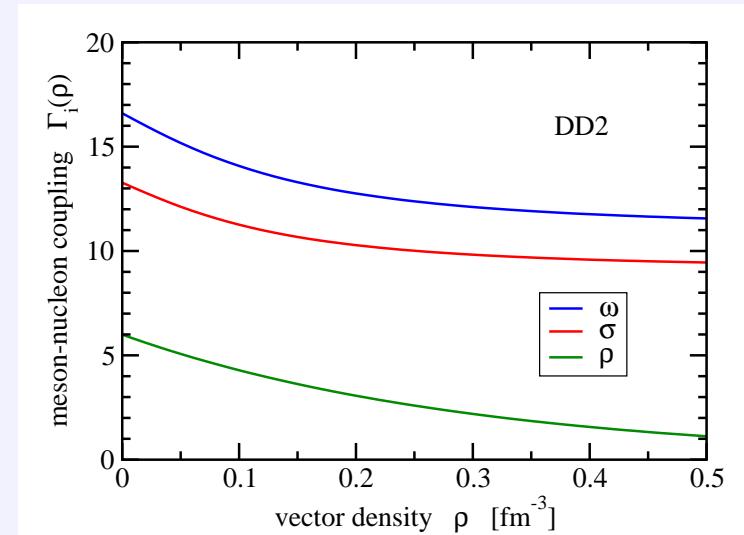
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○ vector potential $V_i = \sum_{m \in \mathcal{V}} \Gamma_{im} A_m + V_i^{(r)}$

with “rearrangement” contribution $V_i^{(r)}$



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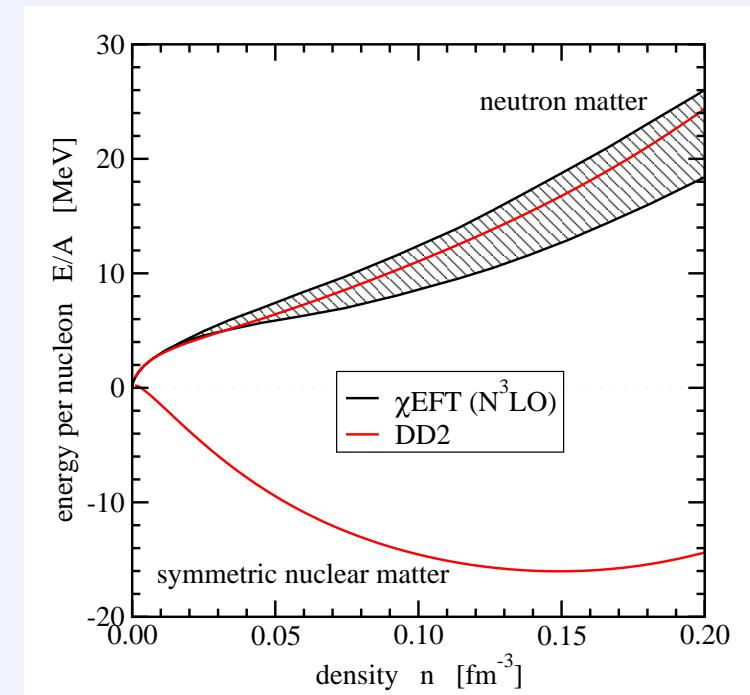
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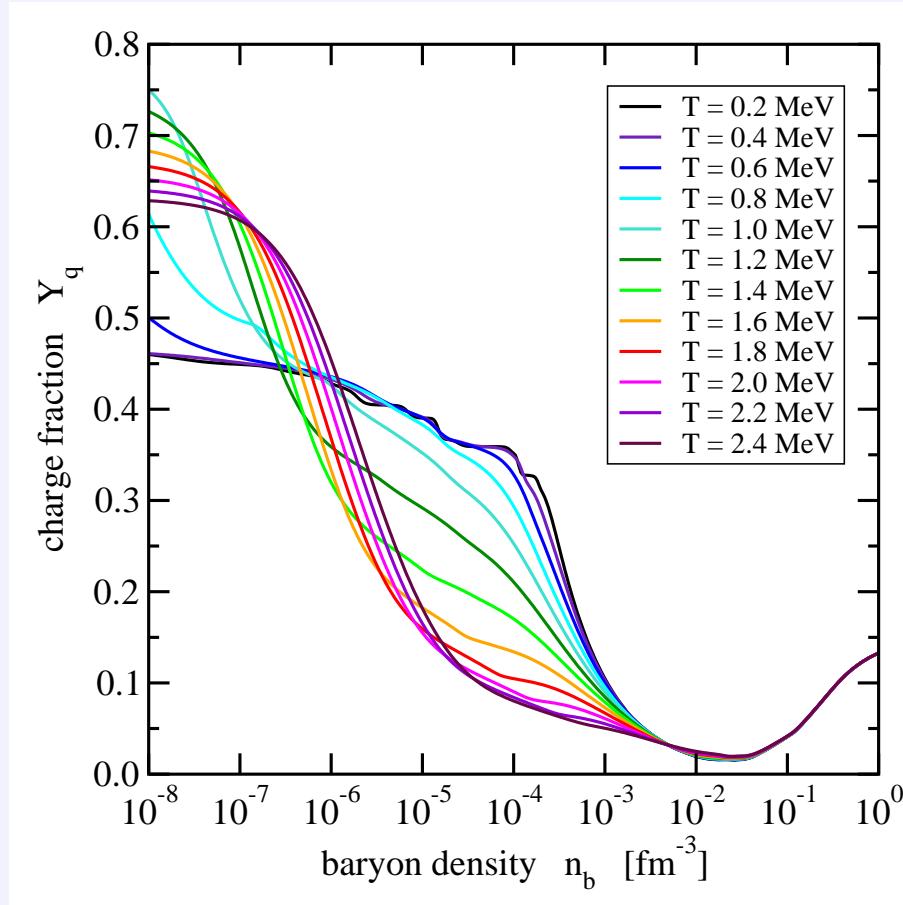
χ EFT(N^3LO):

I. Tews et al., Phys. Rev. Lett 110 (2013) 032504

T. Krüger et al., Phys. Rev. C 88 (2013) 025802

Neutron Star Matter I

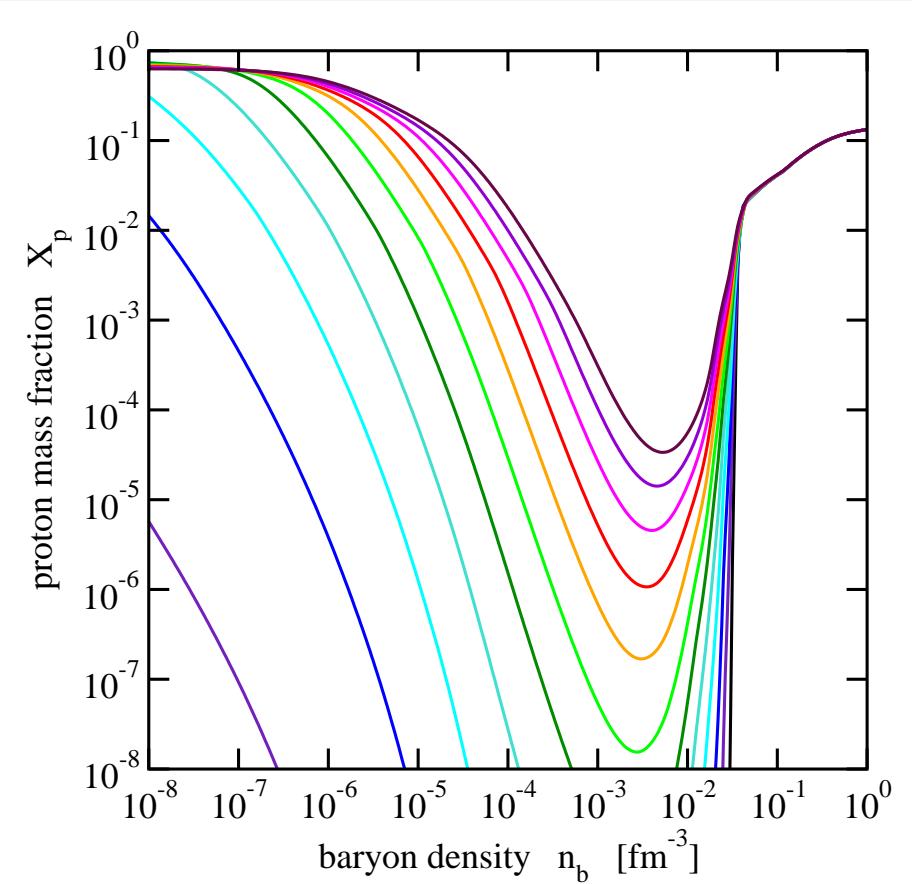
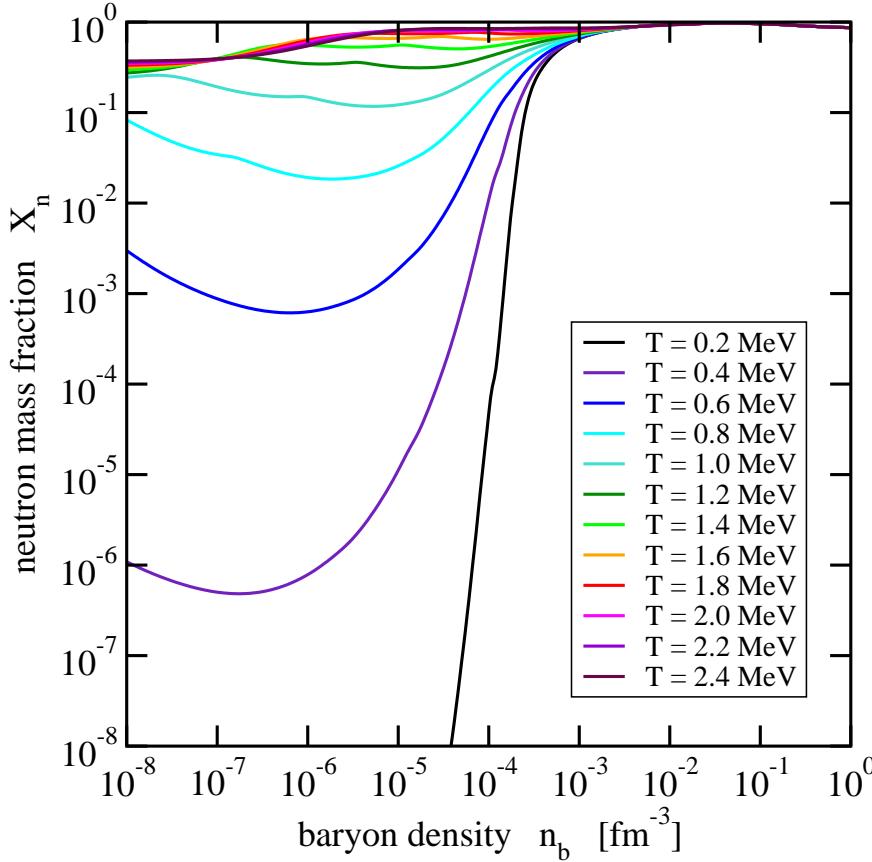
- degrees of freedom: nucleons, light & heavy nuclei, electrons, muons
- charge neutrality $n_p^{\text{total}} = n_e + n_\mu$
- β equilibrium $\mu_n = \mu_p + \mu_e = \mu_p + \mu_\mu$
- charge fraction $Y_q = (n_e + n_\mu)/n_b$



preliminary results:
parametrisation of mass shifts
still under discussion

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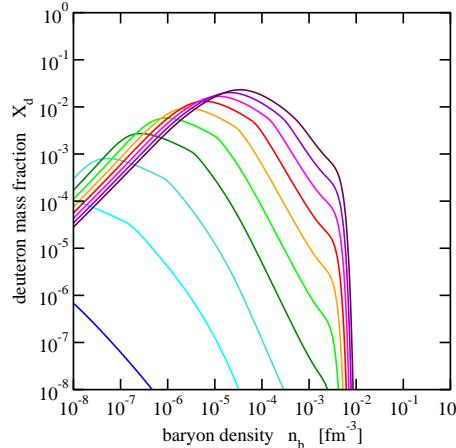
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- nucleon mass fractions $X_n = n_n/n_b$, $X_p = n_p/n_b$



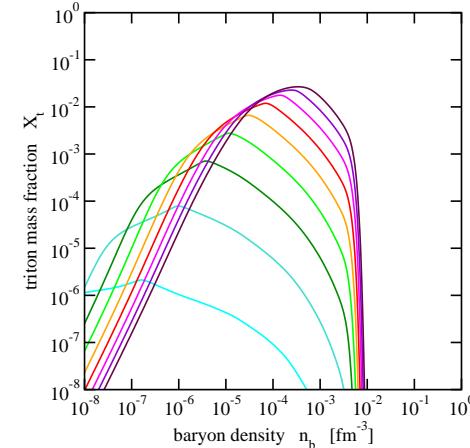
Neutron Star Matter II

- degrees of freedom: nucleons, light & heavy nuclei, electrons, muons
- charge neutrality $n_p^{\text{total}} = n_e + n_\mu$
- β equilibrium $\mu_n = \mu_p + \mu_e = \mu_p + \mu_\mu$
- mass fractions of light nuclei $X_i = A_i n_i / n_b$ ($i = d, t, h, \alpha$)

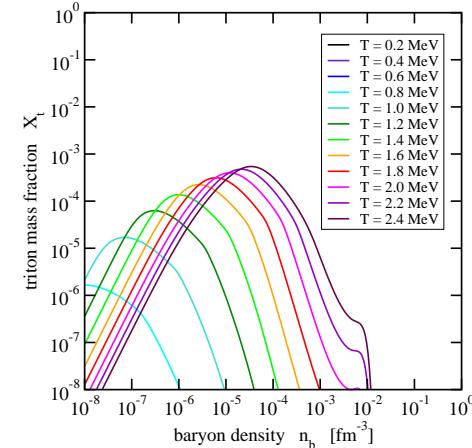
^2H



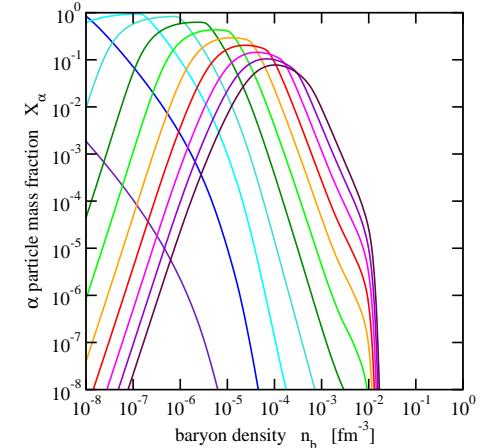
^3H



^3He

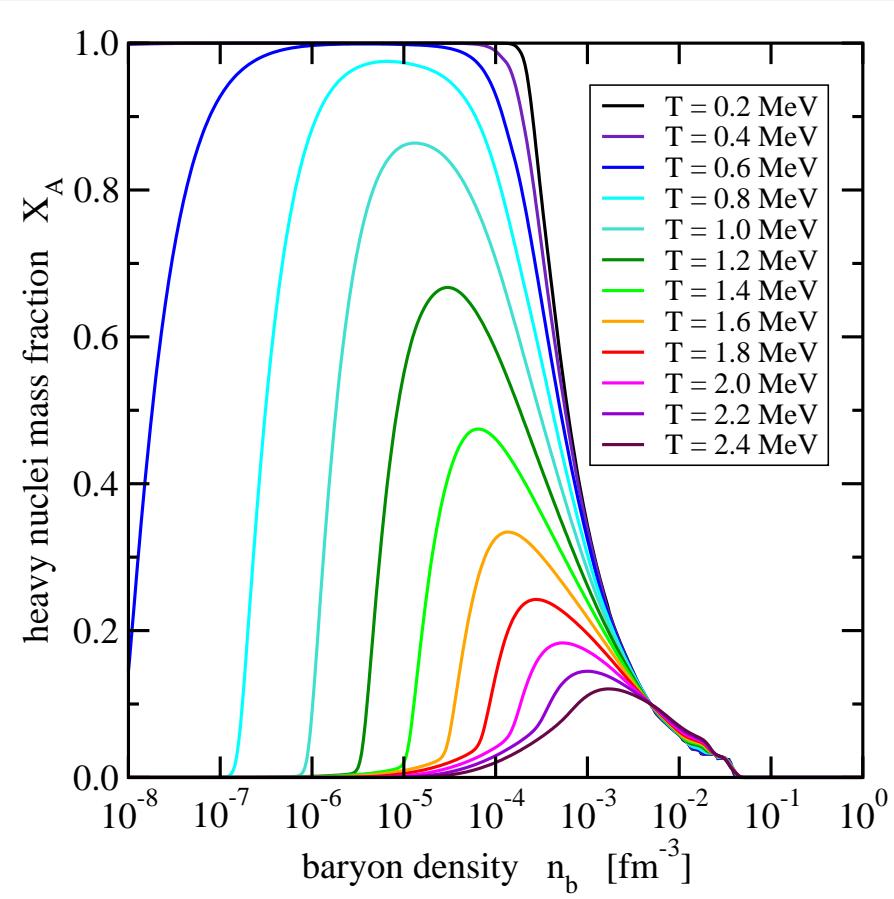


^4He



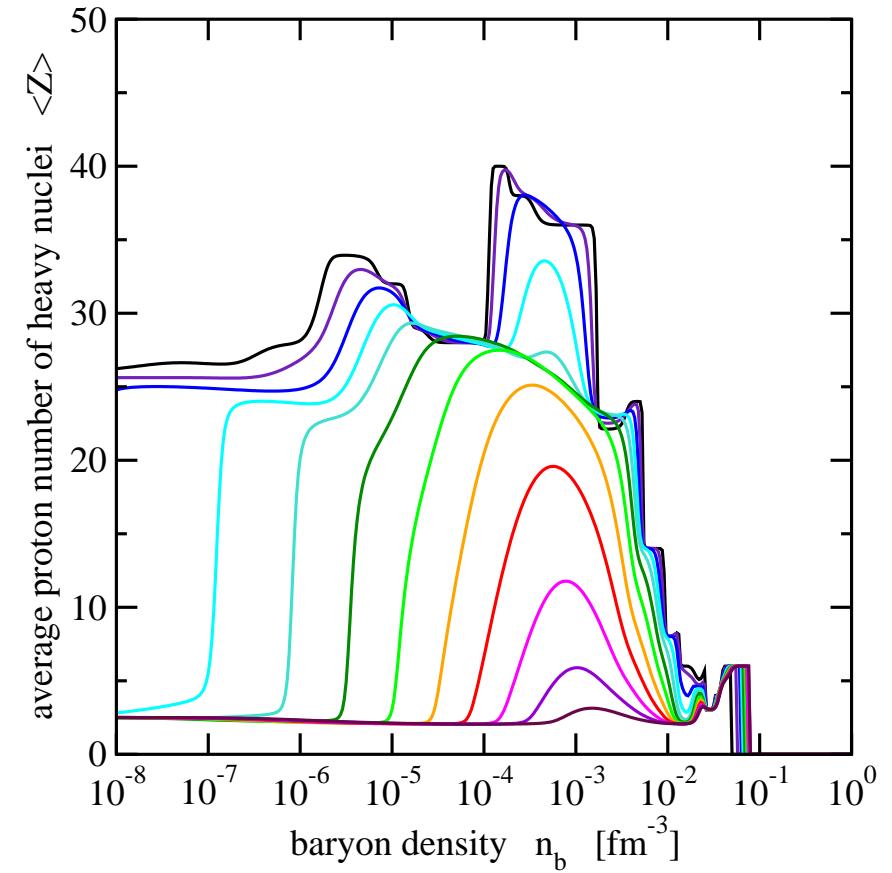
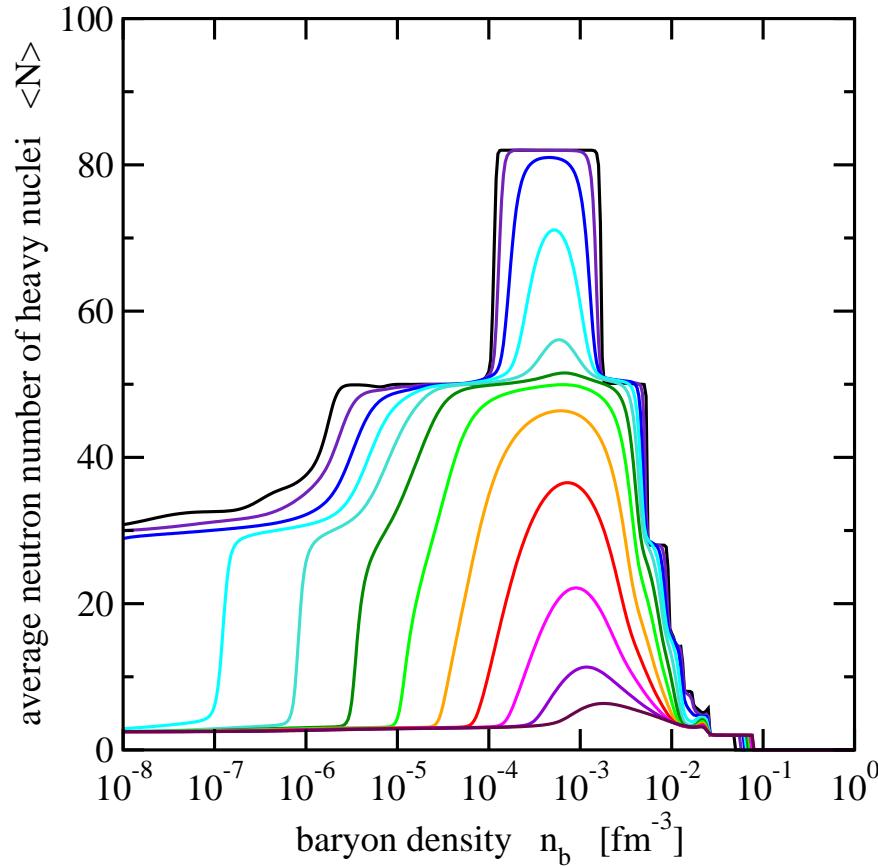
Neutron Star Matter III

- degrees of freedom: nucleons, light & heavy nuclei, electrons, muons
- charge neutrality $n_p^{\text{total}} = n_e + n_\mu$
- β equilibrium $\mu_n = \mu_p + \mu_e = \mu_p + \mu_\mu$
- mass fractions of heavy nuclei $X_i = A_i n_i / n_b$



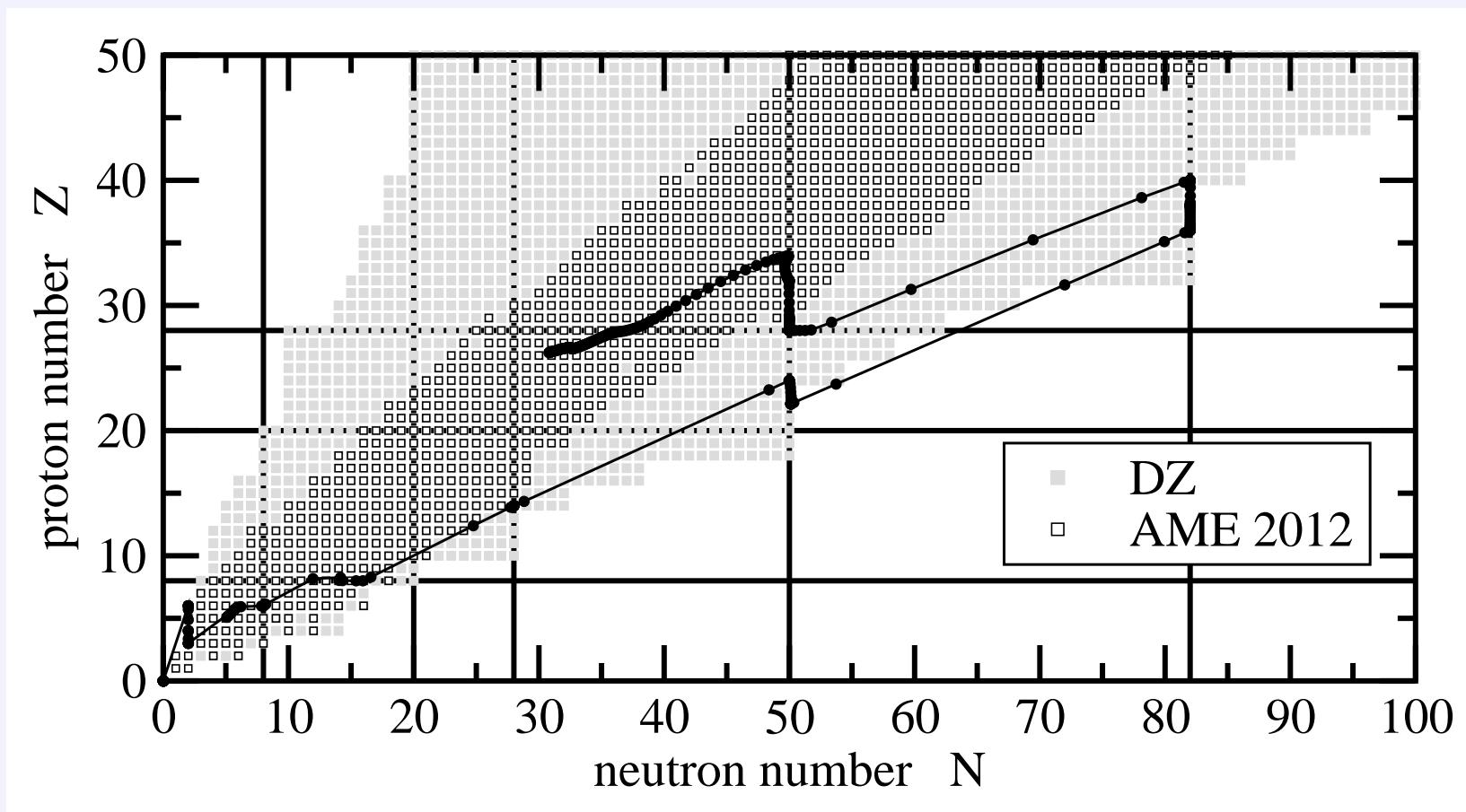
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- **average neutron and proton numbers of heavy nuclei**



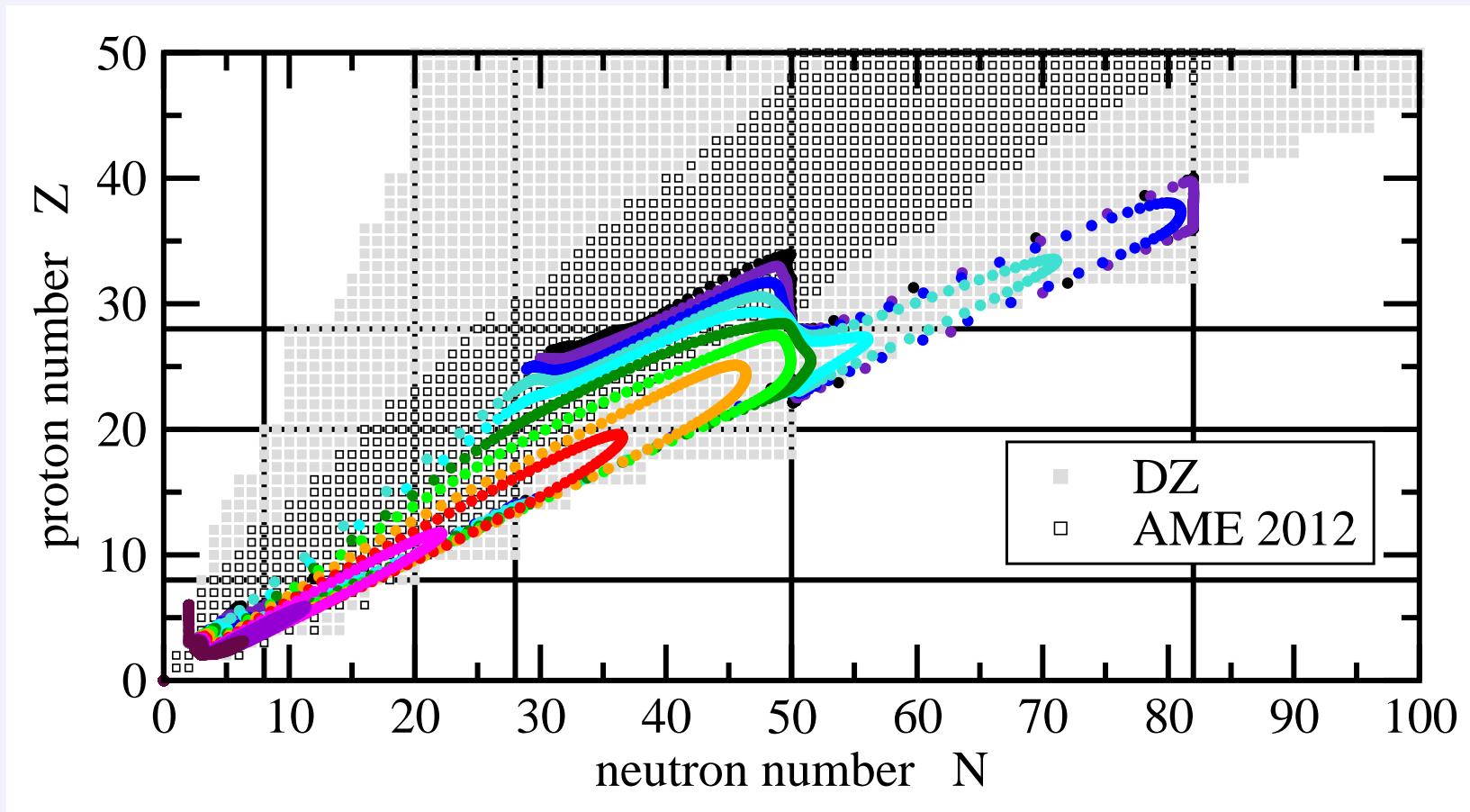
Neutron Star Matter IV

- degrees of freedom: nucleons, light & heavy nuclei, electrons, muons
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Neutron Star Matter IV

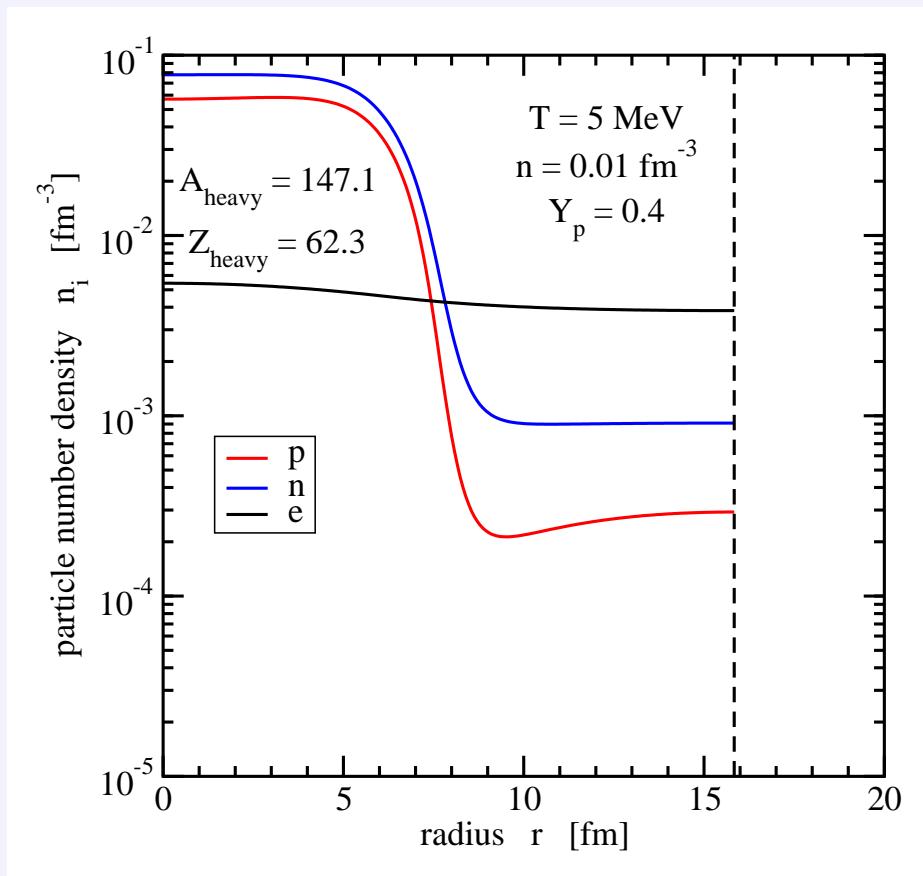
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α -Clusters on the Surface of Nuclei

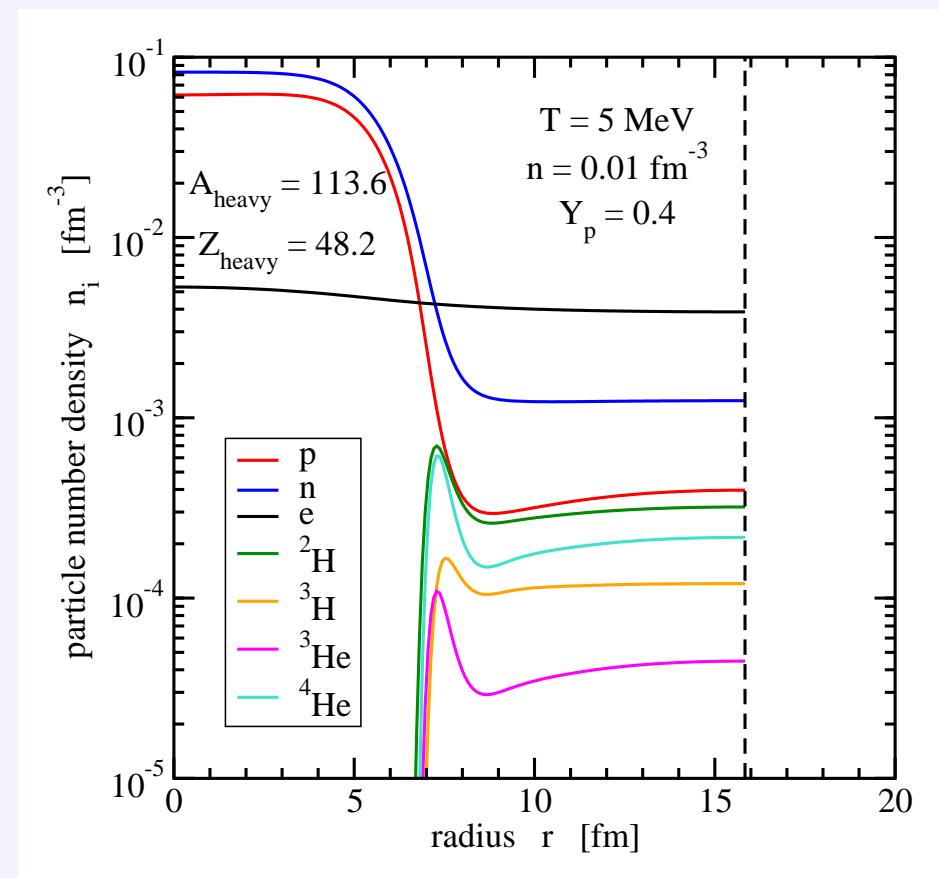
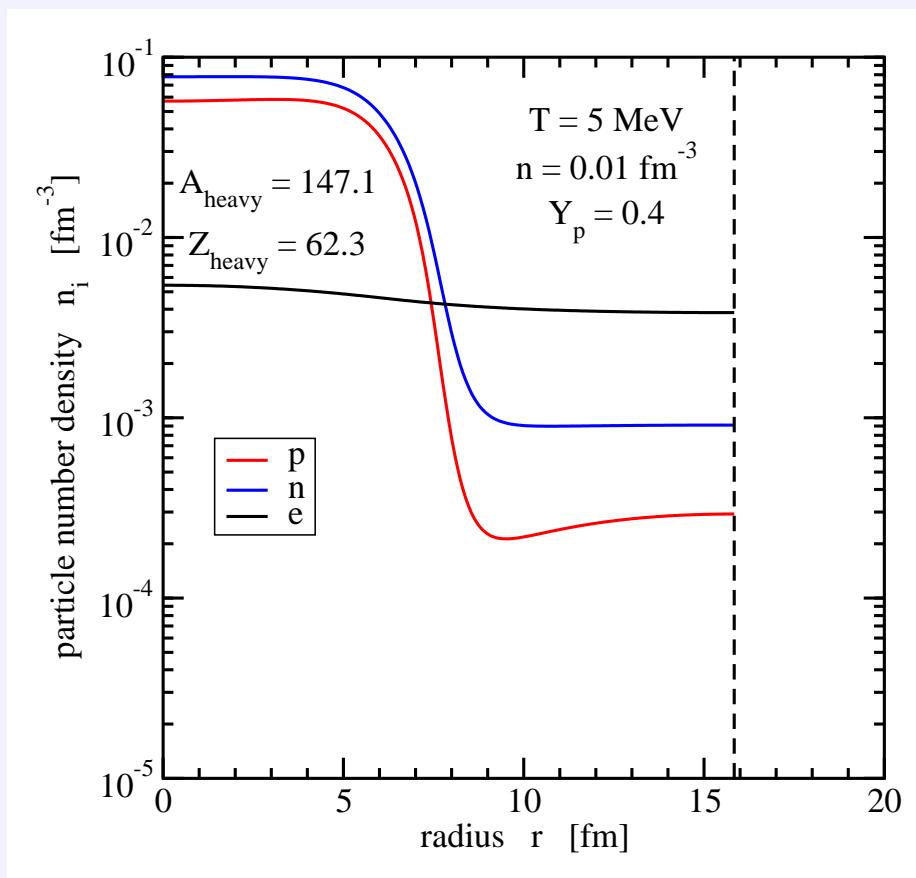
Application of gRDF Model

- finite temperature gRDF calculations in spherical Wigner-Seitz cell, extended Thomas-Fermi approximation without light clusters



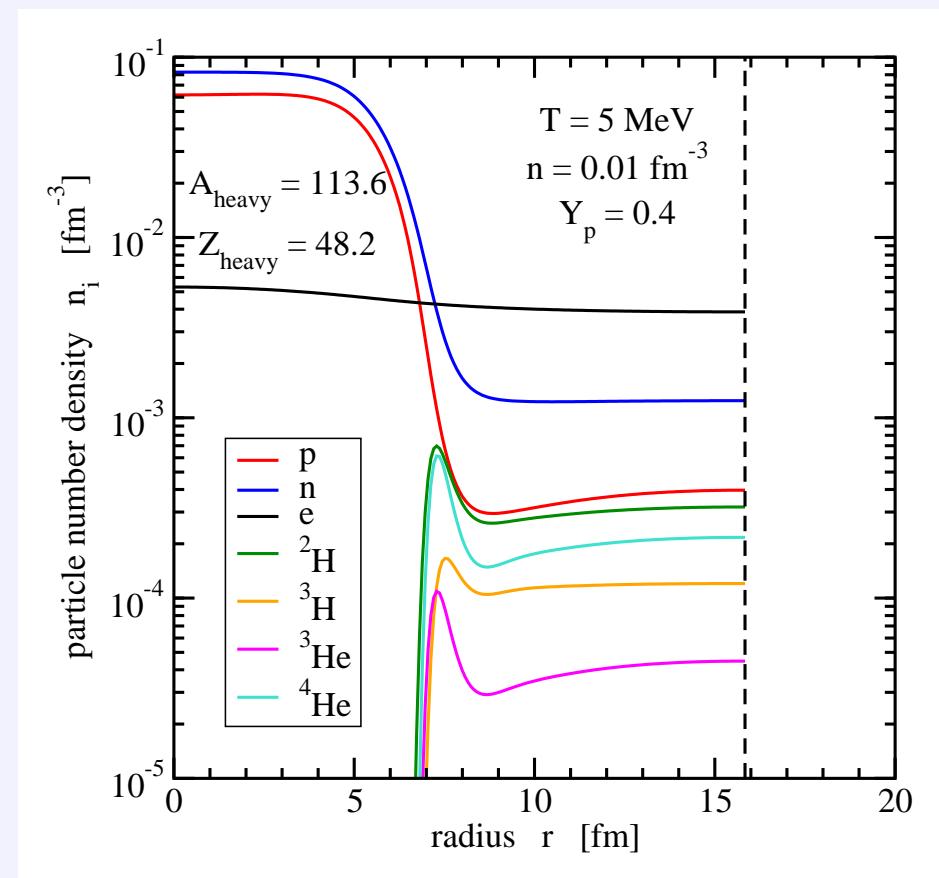
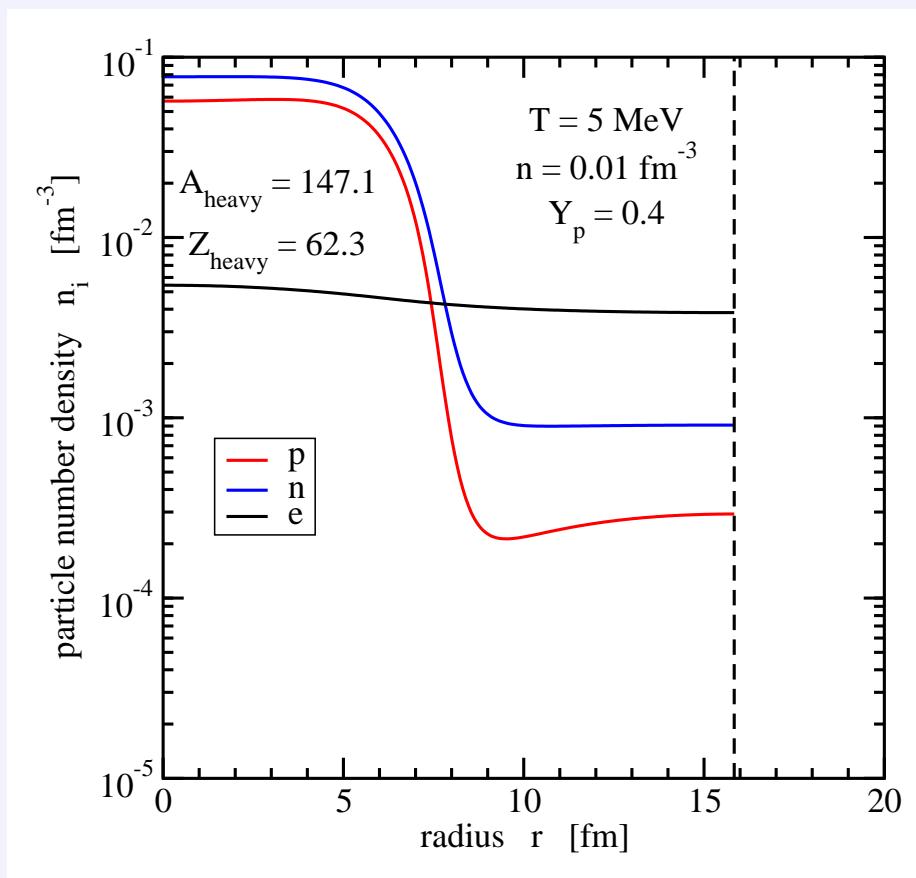
Application of gRDF Model

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 \Rightarrow enhanced cluster probability at surface of heavy nuclei,
effects for **heavy nuclei in vacuum at zero temperature?**



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⇒ modified parametrizations, ^{208}Pb nucleus

parametrization	symmetry energy J [MeV]	slope coefficient L [MeV]	ρ -meson coupling $\Gamma_\rho(n_{\text{ref}})$	ρ -meson parameter a_ρ
DD2 ⁺⁺⁺	35.34	100.00	4.109251	0.063577
DD2 ⁺⁺	34.12	85.00	3.966652	0.193151
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DD2	31.67	55.04	3.626940	0.518903
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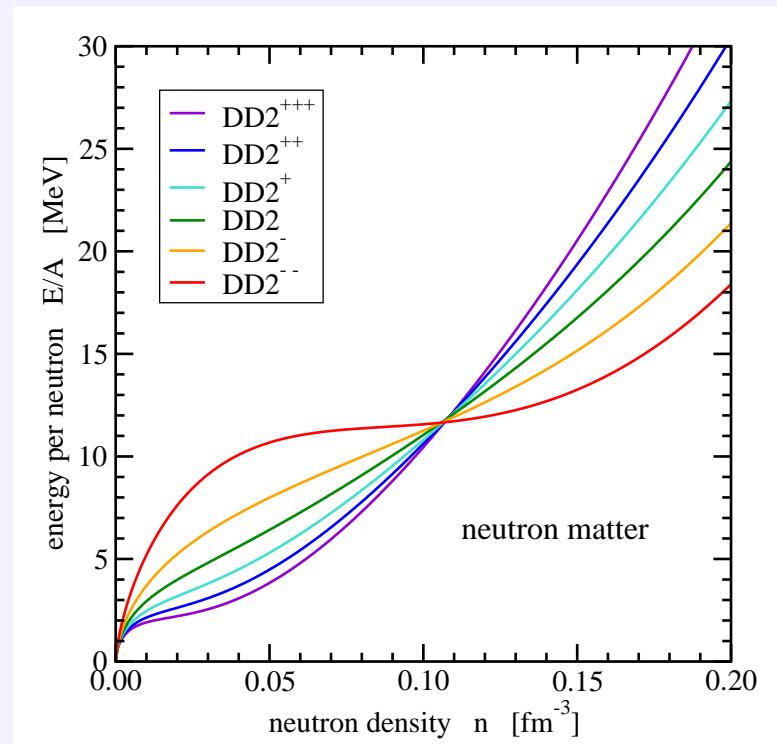
$$\Gamma_\rho(n) = \Gamma_\rho(n_{\text{ref}}) \exp \left[-a_\rho \left(\frac{n}{n_{\text{ref}}} - 1 \right) \right]$$

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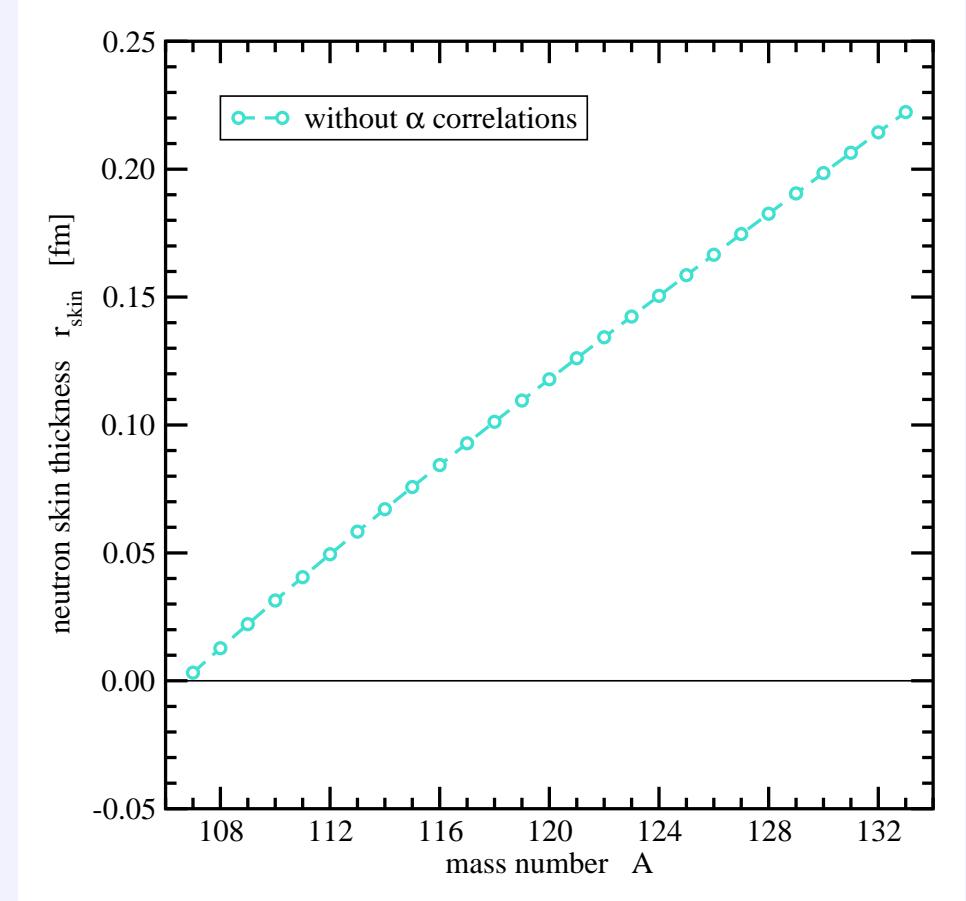
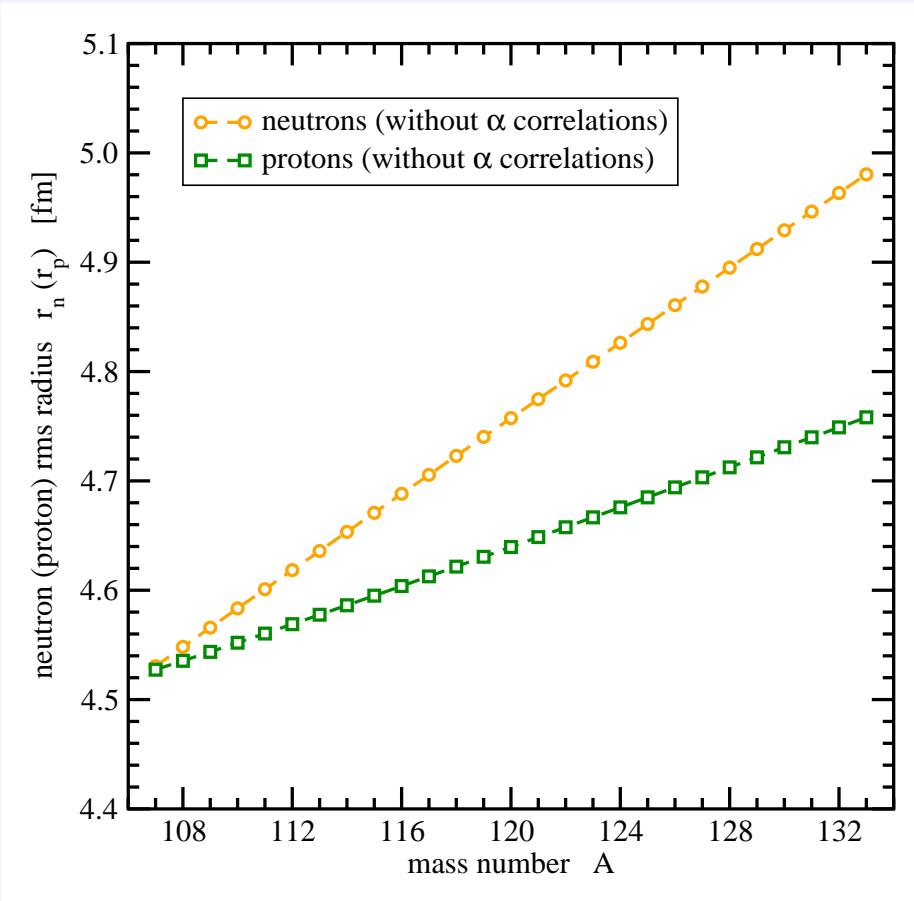
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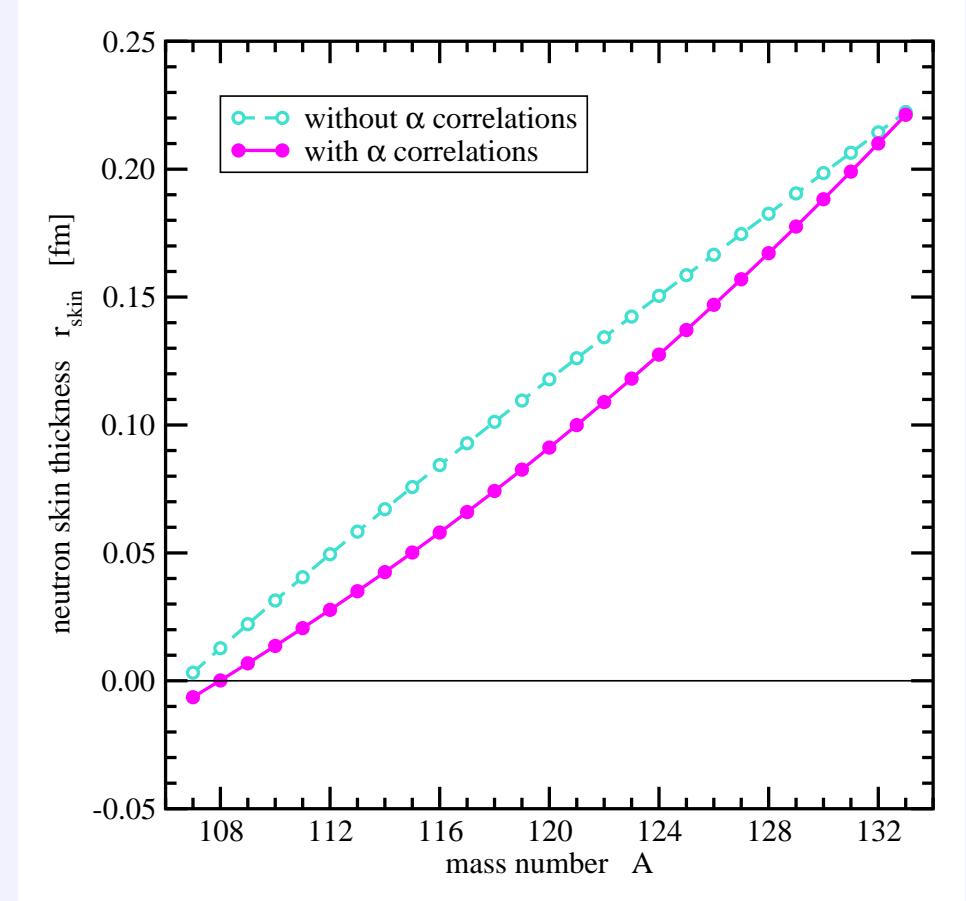
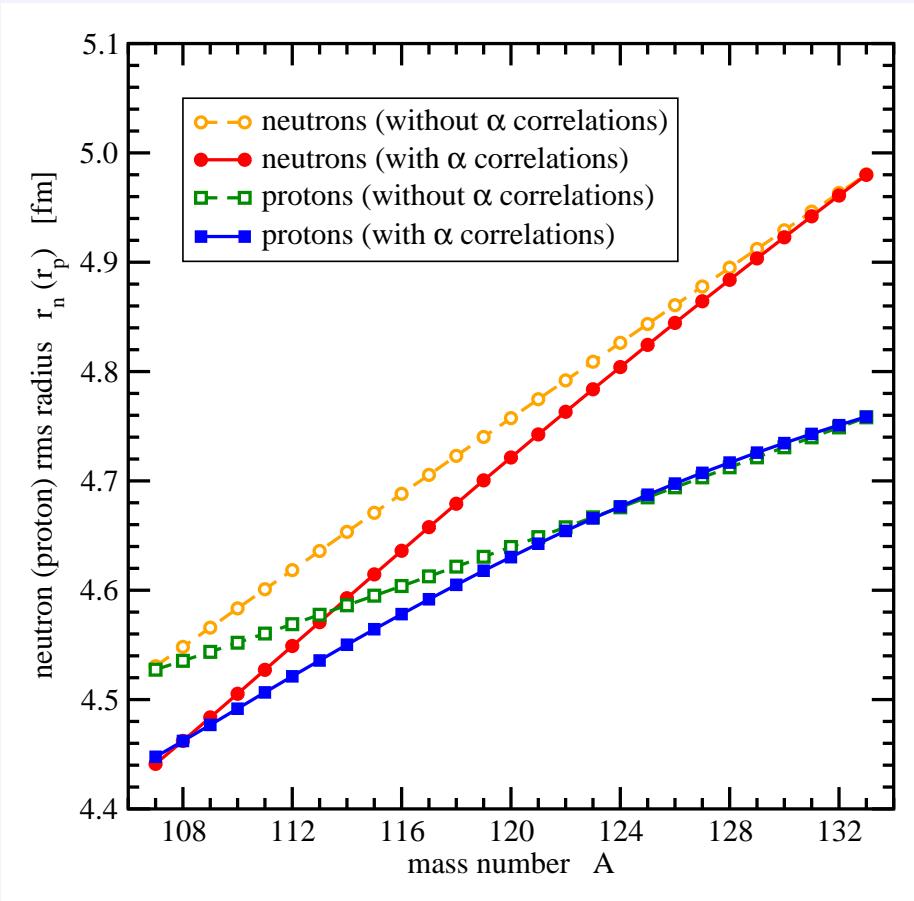
Neutron Skin of Sn Nuclei

- neutron and protons rms radii r_n and r_p
- neutron skin thickness $r_{\text{skin}} = r_n - r_p$



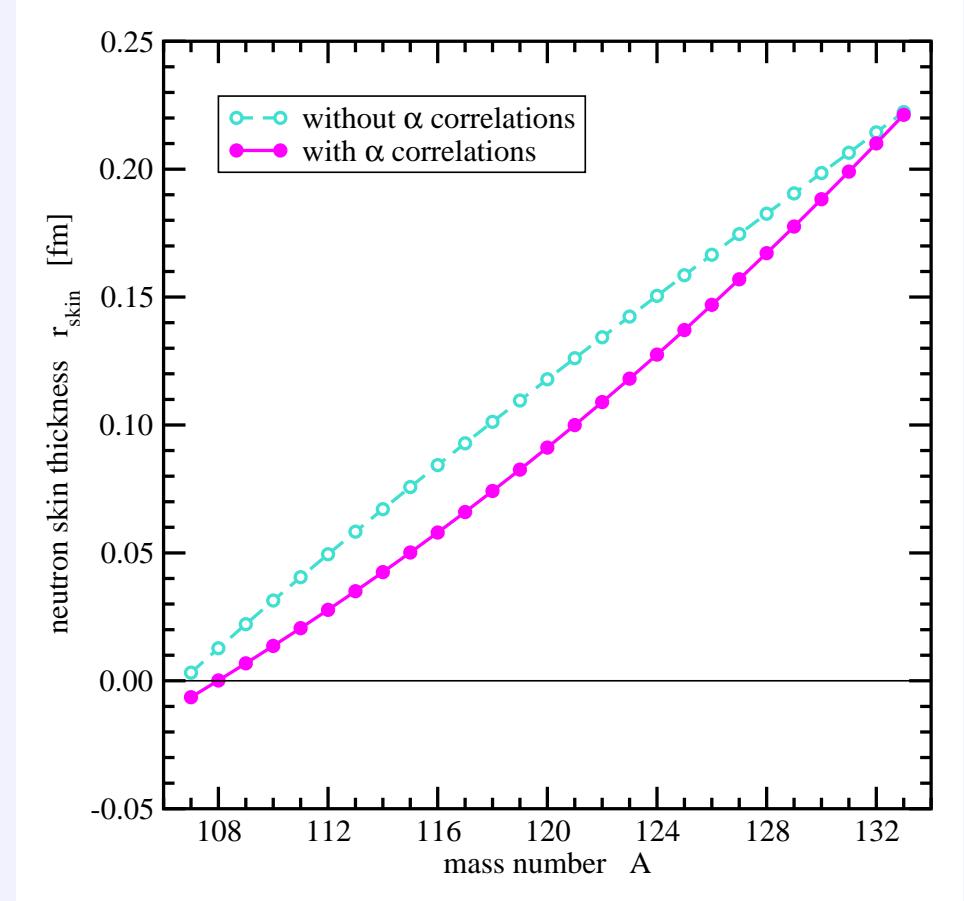
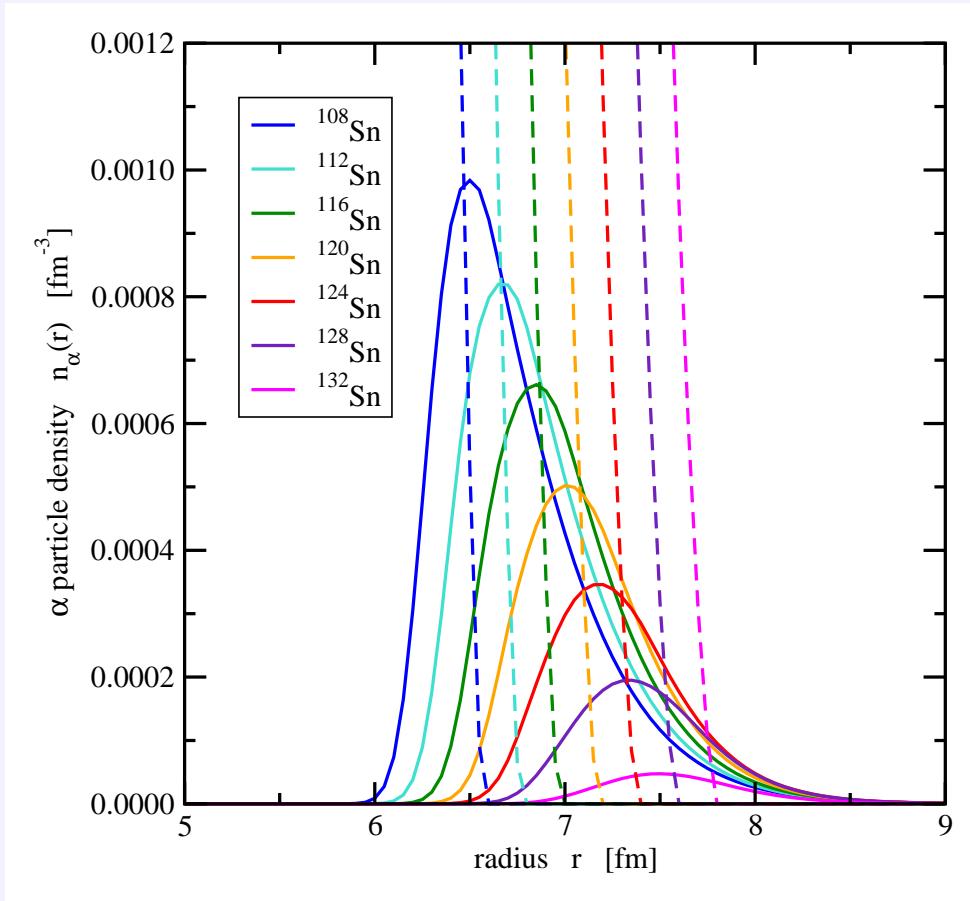
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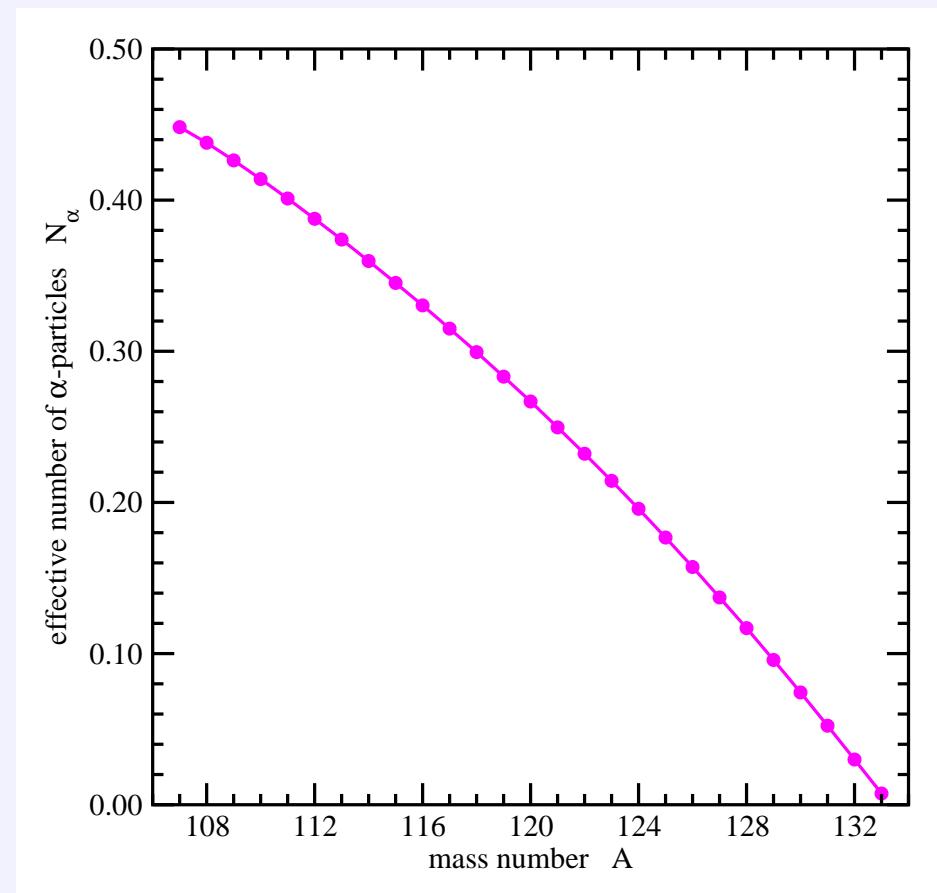
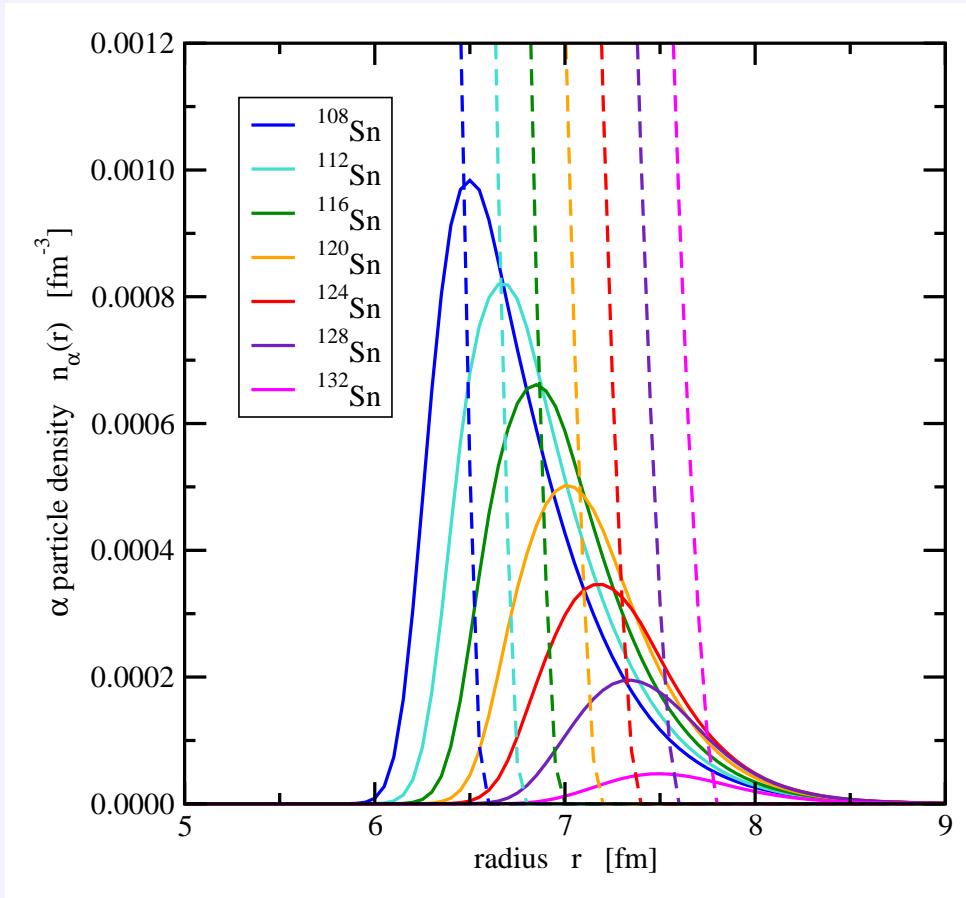
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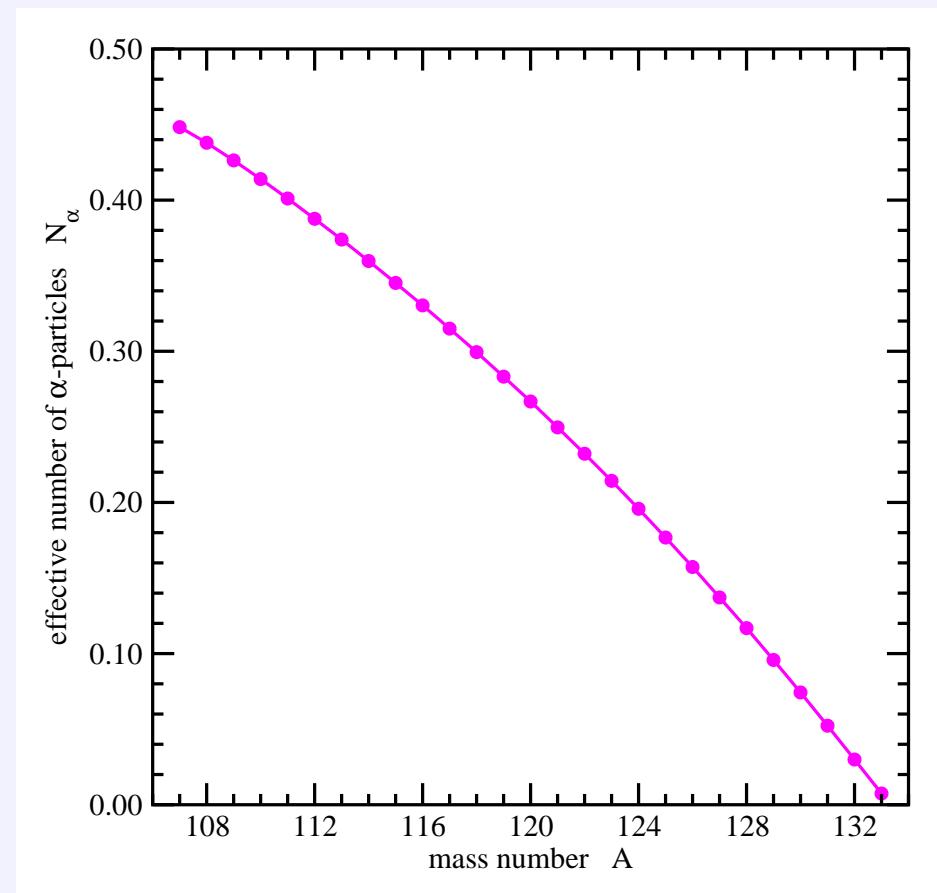
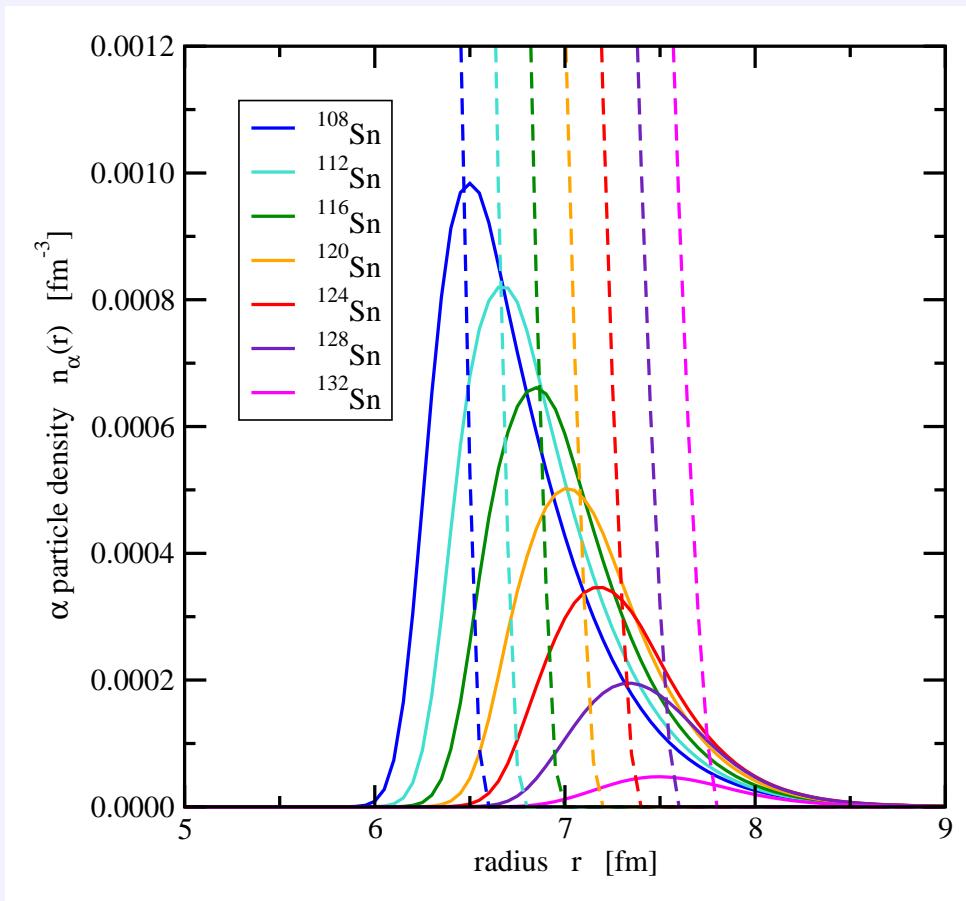
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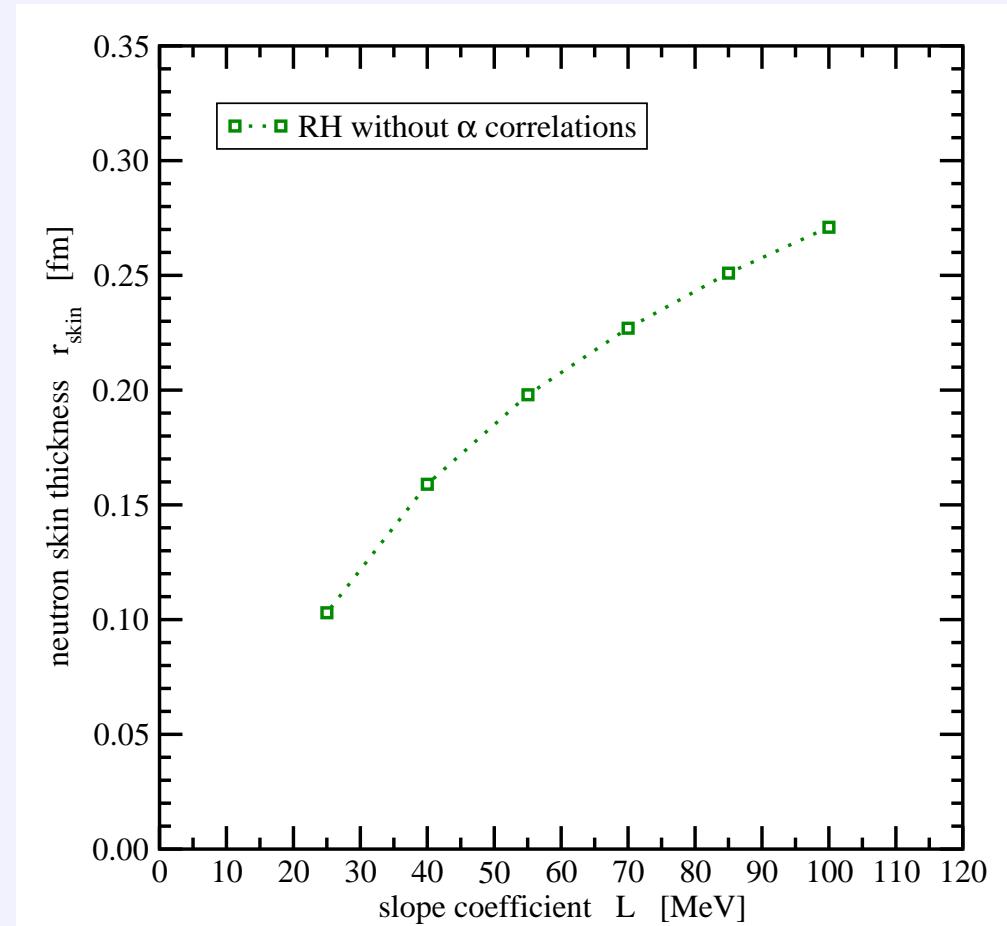
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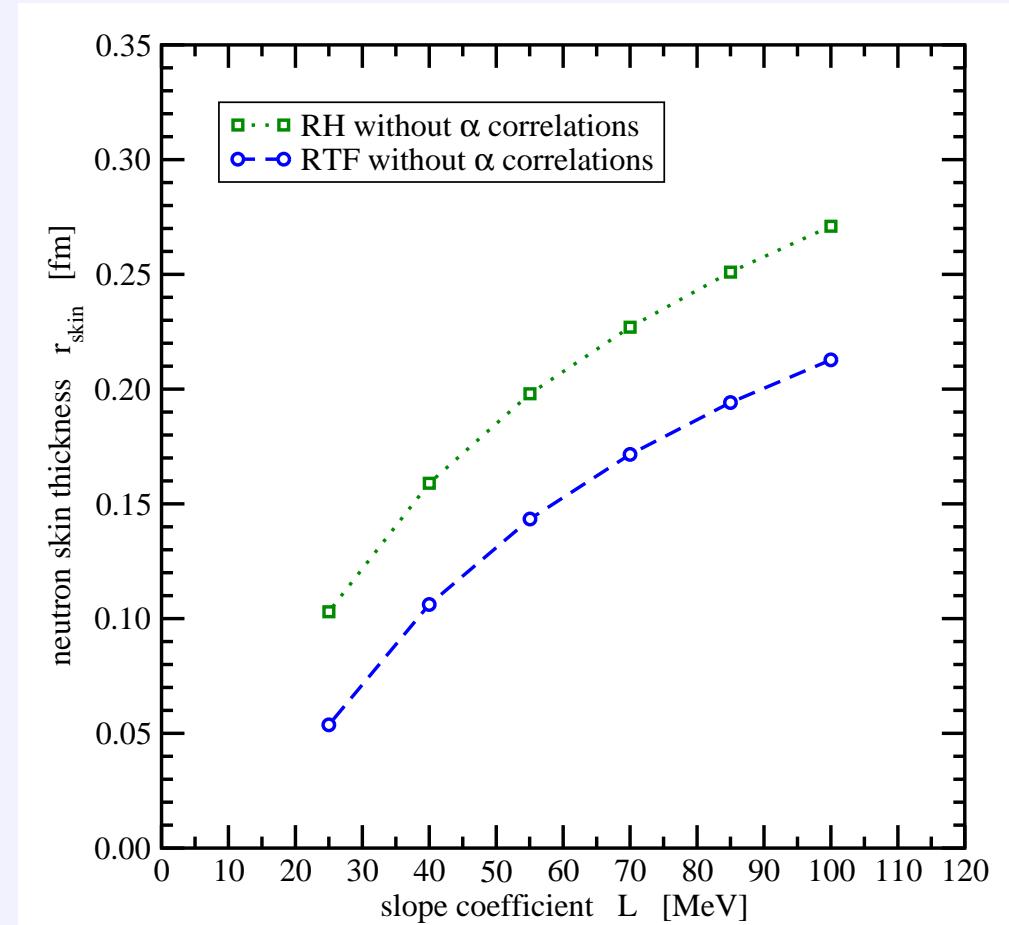
Neutron Skin of ^{208}Pb

- dependence on symmetry energy slope coefficient L
⇒ use parametrizations DD2⁺⁺⁺, ..., DD2⁻⁻
- relativistic Hartree (RH) calculation used in original fit of model parameters
(correlation $r_{\text{skin}} \leftrightarrow L$ not linear because no complete refit of model parameters, only of effective isovector interaction)



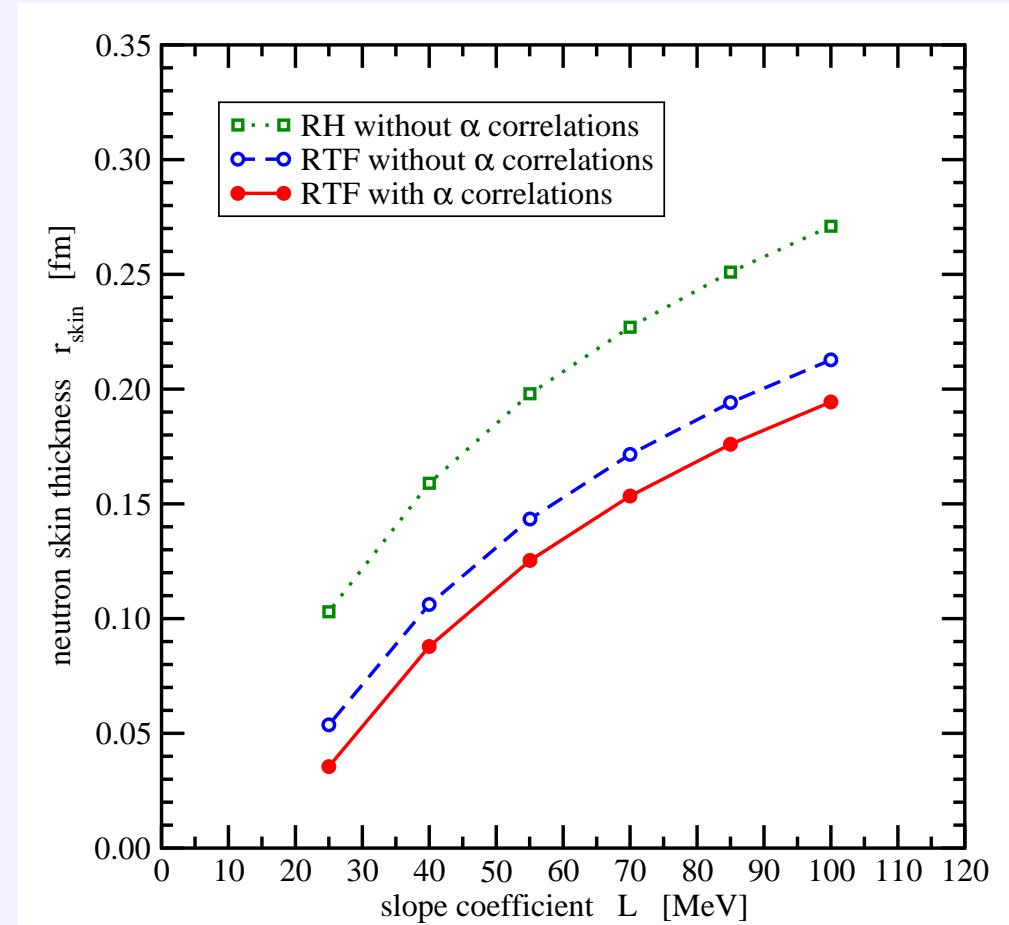
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- with α -particles at surface
⇒ systematic reduction of neutron skin



Experimental Test

Experimental Study of α -Clustering at Nuclear Surface

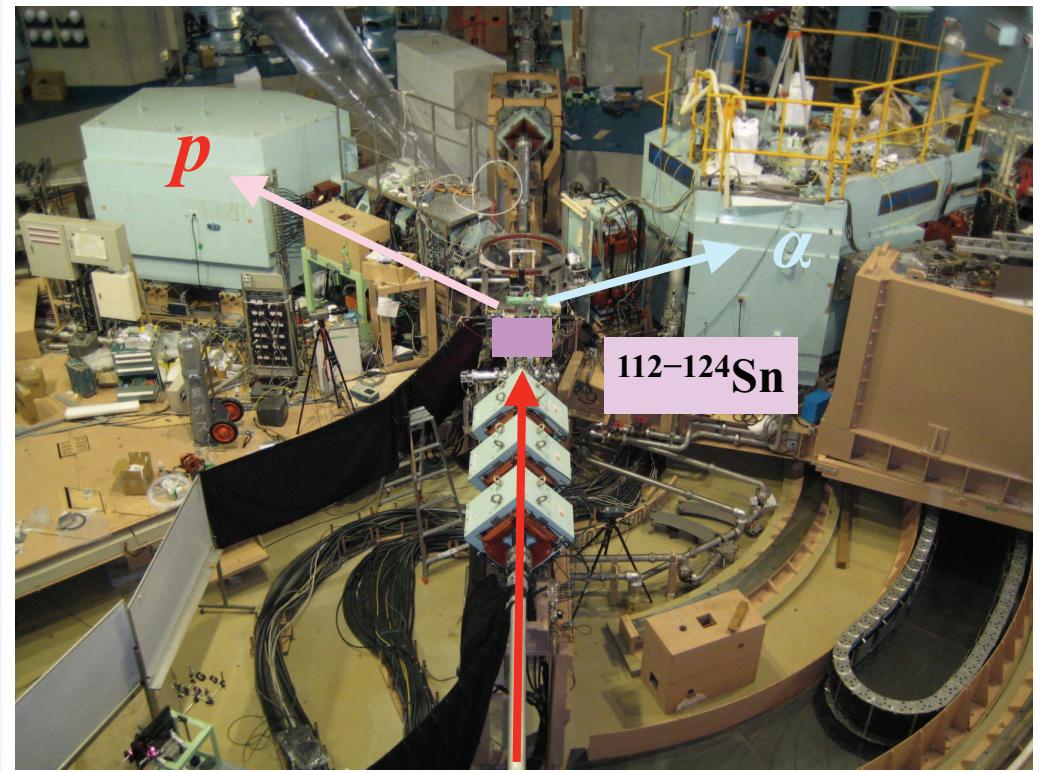
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Experimental Study of α -Clustering at Nuclear Surface

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- experiment at RCNP Cyclotron Facility, Osaka, Japan
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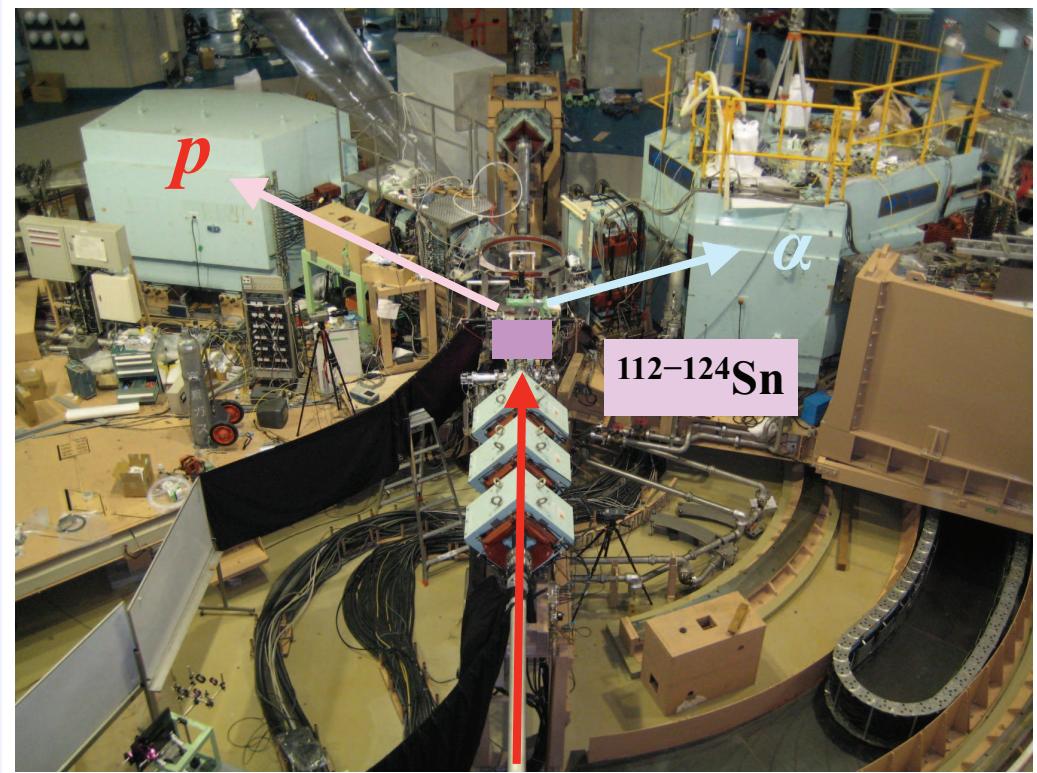
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 - targets: ^{112}Sn , ^{116}Sn , ^{120}Sn , ^{124}Sn
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 - several spectrometer settings
 - experimental signatures:
 - dependence of effective α -particle number (\Rightarrow cross sections) on neutron excess $N - Z$
 - localisation of α -particles on surface of nucleus \Rightarrow broad momentum distribution



Quasi-Free ($p,p\alpha$) Reactions on Sn Nuclei @ 300 MeV

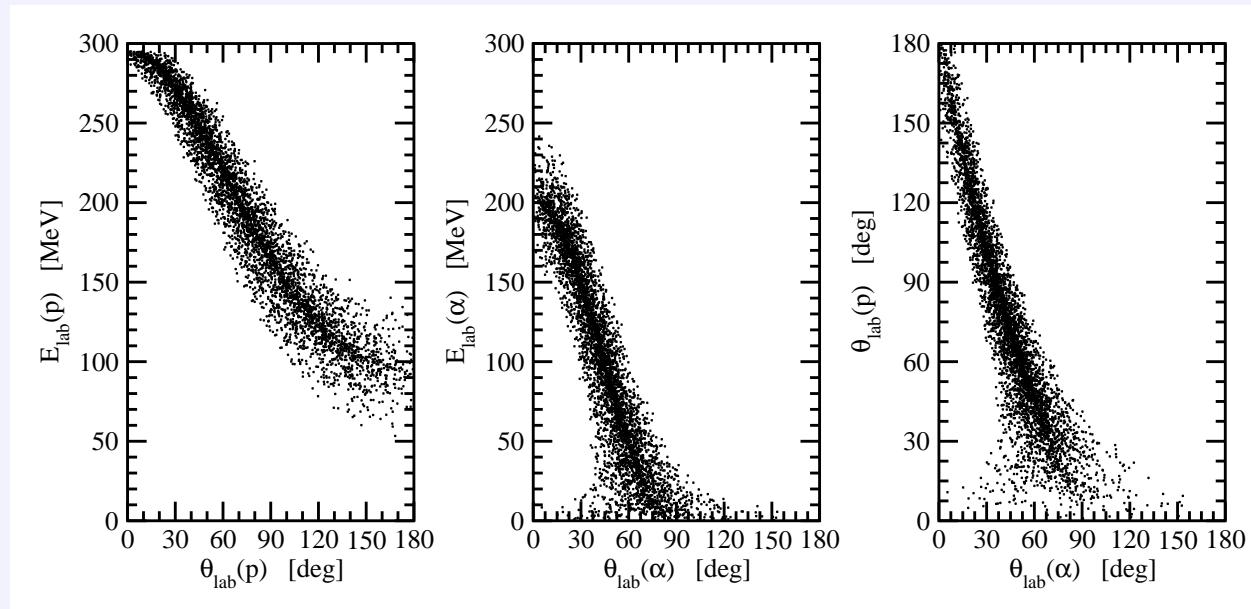
kinematics

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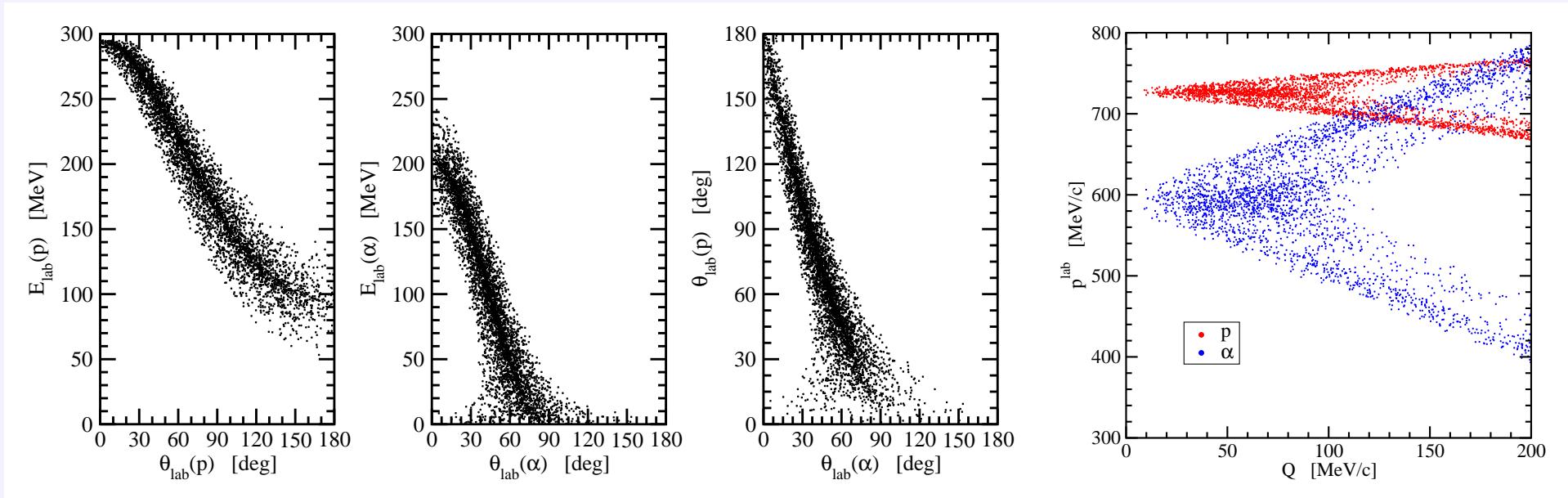
- low momentum transfer to residual Cd nucleus ⇒
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Quasi-Free ($p, p\alpha$) Reactions on Sn Nuclei @ 300 MeV

kinematics

- low momentum transfer to residual Cd nucleus ⇒
 - strong correlation of angles/energies of emitted protons and α -particles
 - select pair of angles, e.g., $\theta_{\text{lab}}(p) = 45^\circ$ and $\theta_{\text{lab}}(\alpha) = 60^\circ$
 - choose spectrometer settings to cover different ranges of intrinsic α -particle momenta Q within acceptance (p, Grand Raiden: 5%, α , LAS: 30%)



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cross sections

- relativistic distorted-wave impulse approximation

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⇒ factorization

$$\frac{d^5\sigma}{dQd\Omega_Qd\Omega'_p} = K \times \frac{d^2\sigma}{d\Omega'_p} \times W_\alpha(\vec{Q}) \times R$$

- kinematic factor K

Quasi-Free ($p,p\alpha$) Reactions on Sn Nuclei @ 300 MeV

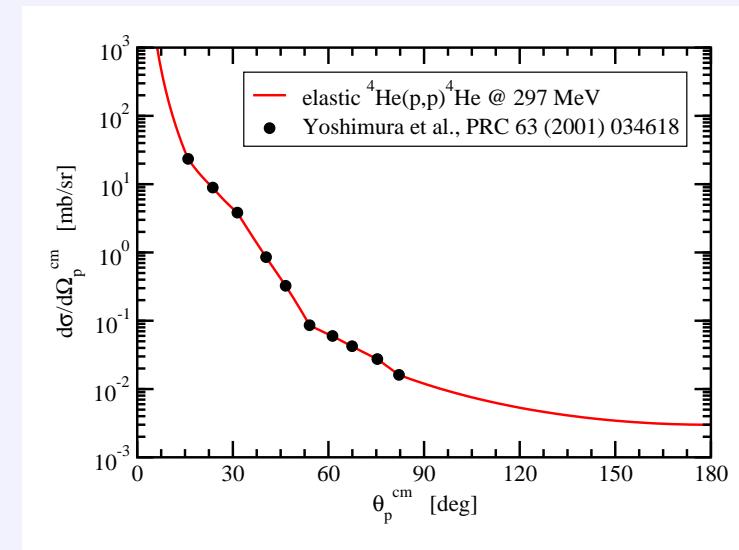
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⇒ use parametrized experimental
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$$\frac{d^2\sigma}{d\Omega'_p}$$



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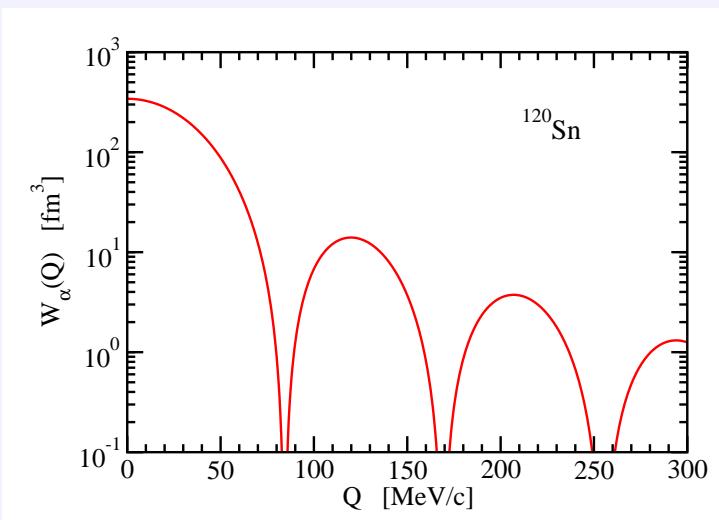
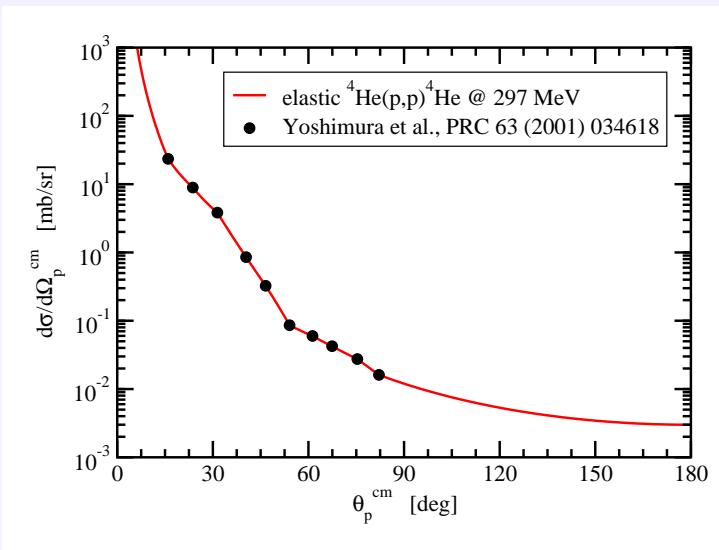
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- α -particle in Sn nucleus
- reduction factor R due to absorption



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 \Rightarrow factorization

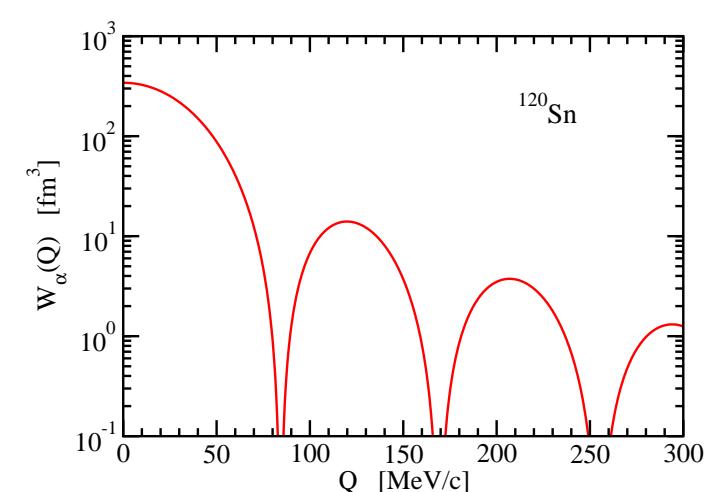
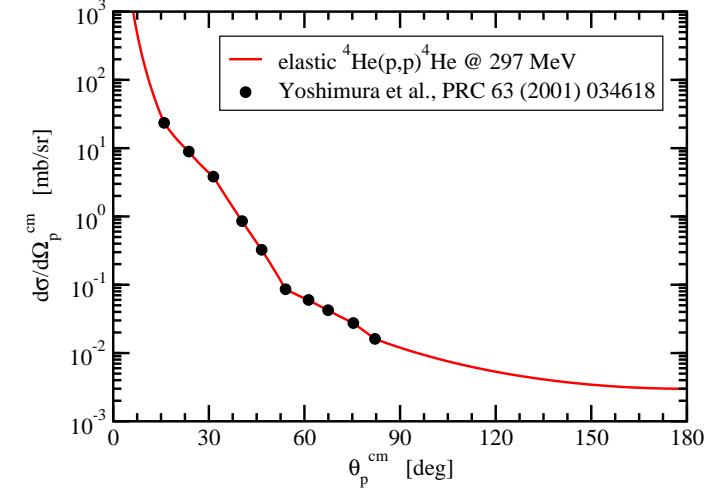
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of α -particle in Sn nucleus

- reduction factor R due to absorption
- Monte Carlo simulation of experiment
 \Rightarrow estimate of count rates



Conclusions

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- many-body correlations essential in low-density nuclear matter
 - formation of clusters/nuclei
 - conditions at surface of heavy nuclei
 - prerequisite for explanation of α -decay
- generalized relativistic density functional (gRDF) for equation of state calculations
 - model with explicit cluster degrees of freedom, quasiparticles with medium-dependent properties
 - effective interaction with density-dependent couplings, well-constrained parameters
- application of gRDF approach to heavy nuclei
 - predicts formation of α -clusters at surface of heavy nuclei
 - ⇒ reduction of neutron skin thickness
 - ⇒ affects correlation with slope coefficient of symmetry energy
 - ⇒ systematic variation of effect with neutron excess of nucleus and with isovector part of effective interaction
- experimental test of predictions
 - quasi-free ($p,p\alpha$) knockout reactions
 - ⇒ experiment with Sn nuclei planned at RCNP, Osaka