

# Maximum Mass and Radial modes of hybrid star in presence of Magnetic field



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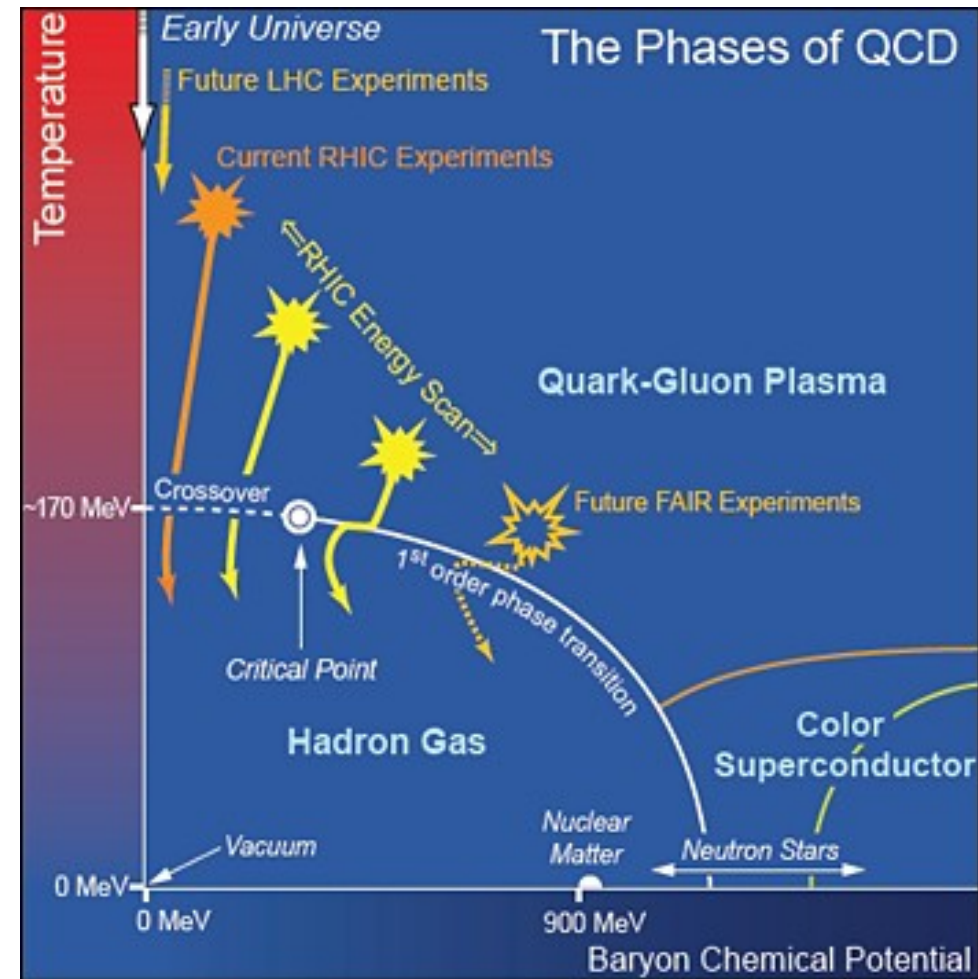
## Talk Outline

- Motivation
- Introduction
- Hadronic EOS
- Quark EOS
- Inclusion of Magnetic field
- Magnetic field in hadronic phase and quark phase
- Mixed phase
- Radial modes
- Summary

# Motivation



- Formation of QGP at very large temperatures or densities.
- QGP plays an important role in the development of the early universe and basic structures of some astrophysical objects like neutron stars.
- A simple phase diagram would show confined quarks in a region of low temperature and chemical potential, surrounded by a phase boundary and QGP outside.
- Ambiguities regarding the physics and the existence of the critical point (CP) on the QCD phase boundary still exist and the mist regarding the conjectured QCD phase boundary has not yet cleared.
- Intensive search for a proper and realistic equations of state (EOS) is still continued for studying the phase diagram existing between quark gluon plasma (QGP) and hadron gas (HG) phases.



**The conjectured phase diagram**

**Reference:-**The Frontiers of Nuclear Science, 2007 NSAC Long Range Plan

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At FAIR energies using baryon meson interaction in relativistic theoretical model based on QCD approach may be proposed. Applying this model we can :-

- \* Study phase transition from Hadron phase to Quark phase
- \* Investigate for the signatures of QGP formation
- \* Verify the observed high precision mass of neutron star and its gross properties.

# Introduction



Compact stars which has only nuclear matter, basically neutron and proton are called **Neutron stars** .

$$M_n \approx 1.5 - 2.1 M_0, R \approx 10 - 15 \text{ km}, \rho \approx 5 - 10 \rho_{\text{nuclear}}$$

**Reference** -[1]P C C Freire, C Bassa, N Wex et al, Mon. Not. R. Astron. Soc. 412, 2763 (2010)  
[2] P Demorest, T Pennucci, S Ransom, M Roberts and J Hessels, Nature 467, 1081 (2010)

Since central density of neutron star exceeds nuclear saturation density so it may contain deconfined quark matter.

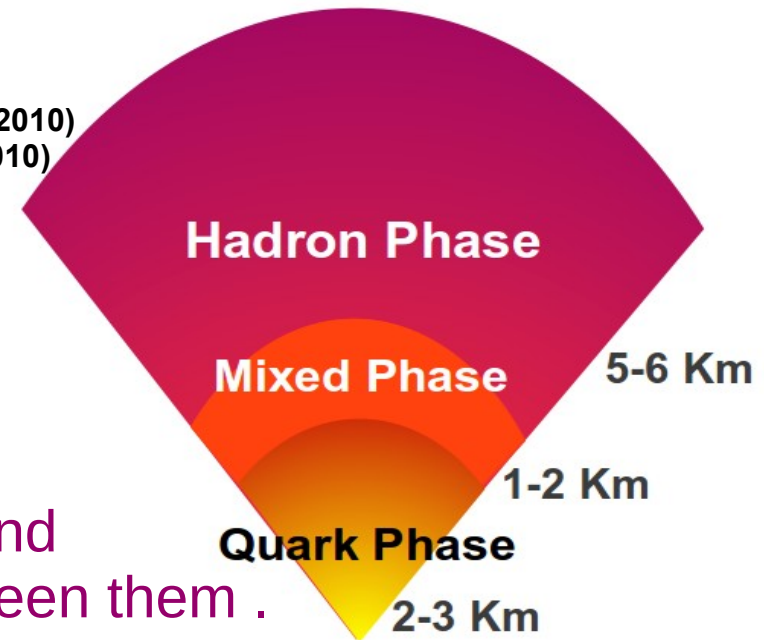
Hybrid stars are NS having interior quark phase and exterior hadron phase and a mixed phase in between them .

**Surface magnetic field =  $10^{14} - 10^{15} G$  , Central magnetic field =  $10^{17} - 10^{18} G$**

**Reference**-D M Palmer, S Barthelmy, N Gehrels et al, Nature 434, 1107 (2005)

The high magnetic field modifies the properties of nuclear and deconfined quark matter of hybrid star

## HYBRID STAR



# Hadronic EOS



The Lagrangian density is :

$$\begin{aligned}\mathcal{L} = & \sum_b \bar{\Psi}_b [ \gamma_\nu (i \partial^\mu - g_{\omega b} \omega^\mu - \frac{1}{2} g_{\rho b} \vec{\tau} \cdot \vec{\rho}^\mu) - (m_b - g_{\rho b} \rho) ] \Psi_b \\ & + \frac{1}{2} (\partial_\mu \sigma \partial^\mu \sigma - m_\sigma^2 \sigma^2) - \frac{1}{4} \omega_{\mu\nu} \omega^{\mu\nu} + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu - \frac{1}{4} \vec{\rho}_{\mu\nu} \cdot \vec{\rho}^{\mu\nu} \\ & + \frac{1}{2} m_\rho^2 \vec{\rho}_\mu \cdot \vec{\rho}^\mu - \frac{1}{3} b m_n (g_\sigma \sigma)^3 - \frac{1}{4} c (g_\sigma \sigma)^4 + \frac{1}{4} d (\omega_\mu \omega^\mu)^2 \\ & + \sum_L \bar{\Psi}_L [ i \gamma_\mu \partial^\mu - m_L ] \Psi_L\end{aligned}$$

\*The leptons L are assumed to be non-interacting,

\*The baryons are coupled to the meson (scalar  $\sigma$  , isoscalar-vector  $\omega$  and isovector-vector  $\rho$  ).

**Reference**-PRAMANA journal of physics, Vol. 82, No. 5, May 2014, pp. 797–807

**Reference**- Nucl. Phys. A 921 (2014) 96-113.

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$$\begin{aligned}\mathcal{L}_{hy} = & \frac{1}{2} (\partial_\nu \sigma^* \partial^\nu \sigma^* - m_{\sigma^*}^2 \sigma^{*2}) - \frac{1}{4} S_{\mu\nu} S^{\mu\nu} \\ & + \frac{1}{2} m_\phi^2 \phi_\mu \phi^\mu - \sum_b g_{\sigma^* b} \bar{\Psi}_b \Psi_b \sigma^* \\ & - \sum_b g_{\phi b} \bar{\Psi}_b \gamma_\mu \Psi_b \phi^\mu\end{aligned}$$

Reference-J. Schaffner, I.N. Mishustin, Phys. Rev. C, Nucl. Phys. 53 (1996) 1416.

Energy density of hadron phase :

$$\begin{aligned}\varepsilon_{hp} = & \frac{1}{2} m_\omega^2 \omega_0^2 + \frac{1}{2} m_\rho^2 \rho_0^2 + \frac{1}{2} m_\sigma^2 \sigma_0^2 + \frac{1}{2} m_{\sigma^*}^2 \sigma^{*2} + \frac{1}{2} m_\phi^2 \phi_0^2 \\ & + \frac{3}{4} d \omega_0^4 + U(\sigma) + \sum_b \varepsilon_b + \sum_l \varepsilon_l\end{aligned}$$

Pressure of hadron phase  $P_{hp} = \sum_i \mu_i n_i - \varepsilon$

$U(\sigma)$  stands for scalar self interactions

# Quark EOS



Taking density dependence MIT bag model

$$\varepsilon_{qp} = \sum_{i=u,d,s} \frac{g_i}{2\pi^2} \int_0^{P_F^i} dp p^2 \sqrt{m_i^2 + p^2} + B_G$$

$$P_{qp} = \sum_{i=u,d,s} \frac{g_i}{6\pi^2} \int_0^{P_F^i} dp \frac{p^4}{\sqrt{m_i^2 + p^2}} - B_G$$

$$B_G(n_b) = B_\infty + (B_g - B_\infty) \exp\left[-\beta \left(\frac{n_b}{n_0}\right)^2\right]$$

# Inclusion of Magnetic field



For the beta equilibrated matter :  $\mu_i = b_i \mu_B + q_i \mu_e$

For charge neutrality condition :  $\rho_c = \sum q_i n_i$

For the magnetic field along Z-direction :  $A^\mu \equiv (0, -yB, 0, 0)$ ,  $\vec{B} = B \hat{k}$

$$\mathbf{B} = \mathbf{B}_s + \mathbf{B}_0 \left\{ 1 - e^{-\alpha(n_b/n_0)^{\gamma}} \right\}$$

Single particle energy eigenvalue in presence of magnetic field

$$E_i = \sqrt{p_i^2 + m_i^2 + 2 \tilde{\nu} |q_i| B}, \text{ where } \tilde{\nu} = 2n + s + 1$$

The energy density of charged particles :

$$\varepsilon_i = \frac{|q_i| B}{4 \pi^2} \sum_{\tilde{\nu}} \left[ E_F^i p_{F, \tilde{\nu}}^i + \tilde{m}_{\tilde{\nu}}^{i2} \ln \left( \frac{E_F^i + p_{F, \tilde{\nu}}^i}{\tilde{m}_{\tilde{\nu}}^i} \right) \right]$$

$$p_{F, \tilde{\nu}}^i{}^2 = E_F^i{}^2 - \tilde{m}_{\tilde{\nu}}^i{}^2$$

$$\tilde{m}_{\tilde{\nu}}^i{}^2 = m_i^2 + 2 \tilde{\nu} |q_i| B$$



# Magnetic field in hadronic phase and quark phase



Total energy density and pressure of hadronic phase :

$$\begin{aligned} \varepsilon_{\text{HP}} = & \frac{1}{2} m_{\omega}^2 \omega_0^2 + \frac{1}{2} m_{\rho}^2 \rho_0^2 + \frac{1}{2} m_{\sigma}^2 \sigma_0^2 + \frac{1}{2} m_{\sigma^*}^2 \sigma^{*2} + \frac{1}{2} m_{\varphi}^2 \varphi_0^2 \\ & + \frac{3}{4} d \omega_0^4 + U(\sigma) + \sum_b \varepsilon_b + \sum_l \varepsilon_l + \frac{B^2}{8\pi^2} \end{aligned}$$

$$P_{\text{HP}} = \sum_i \mu_i n_i - \varepsilon$$

Thermodynamic potential in presence of strong magnetic field at zero temperature:

$$\Omega_i = -\frac{2 g_i |q_i| B}{4\pi^2} \sum_{\tilde{\nu}} \int_{\sqrt{m_i^2 + 2\nu|q_i|B}}^{\mu} dE_i \sqrt{E_i^2 - m_i^2 - 2\tilde{\nu}|q_i|B}$$

Energy density and Pressure of quark phase

$$\varepsilon_{QP} = \sum_i \Omega_i + B_G + \sum_i n_i \mu_i, \quad P_{QP} = -\sum_i \Omega_i - B_G$$

# Mixed phase



Performing Glendenning construction and taking the Gibb's condition

$$P_{\text{HP}}(\mu_e, \mu_B) = P_{\text{QP}}(\mu_e, \mu_B) = P_{\text{MP}}$$

$$\chi \rho_c^{\text{QP}} + (1 - \chi) \rho_c^{\text{HP}} = 0$$

$\chi = 1$  : Quark Phase

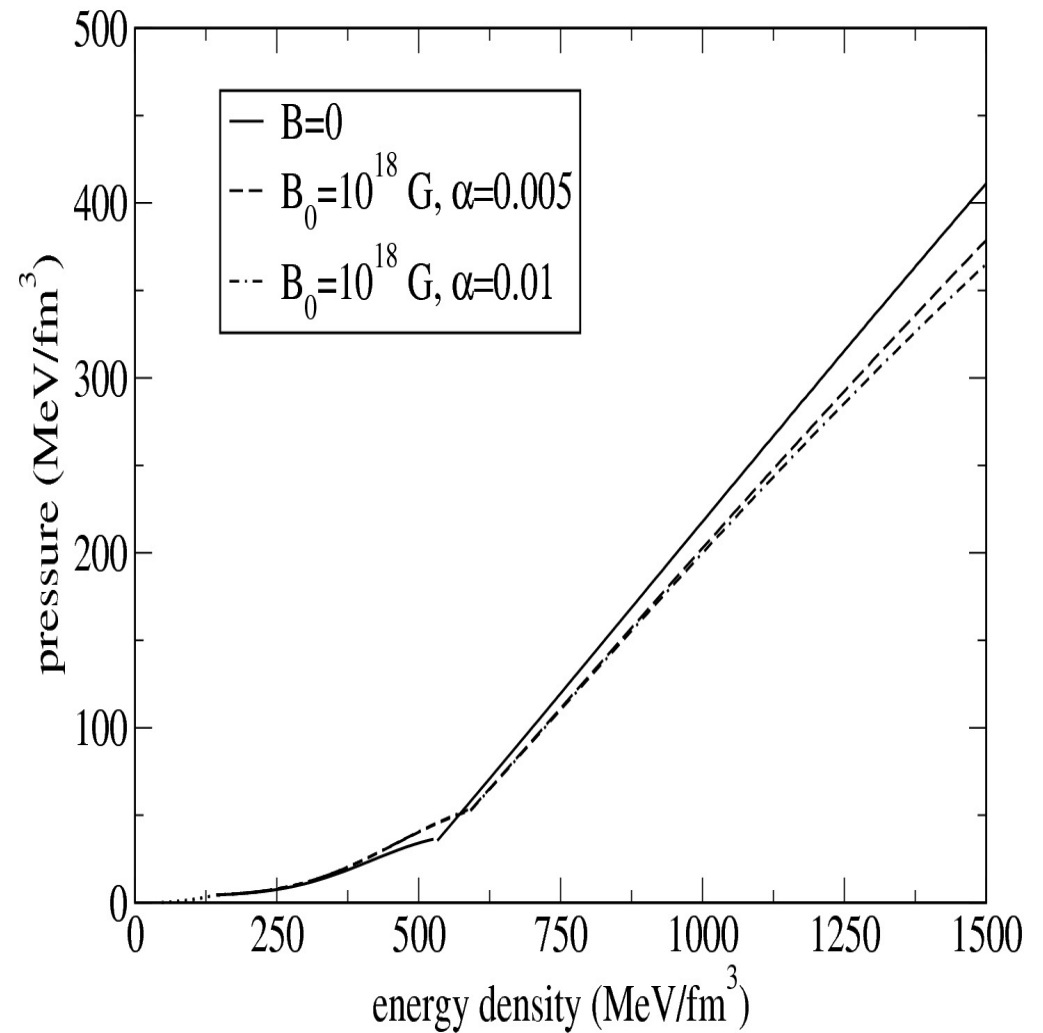
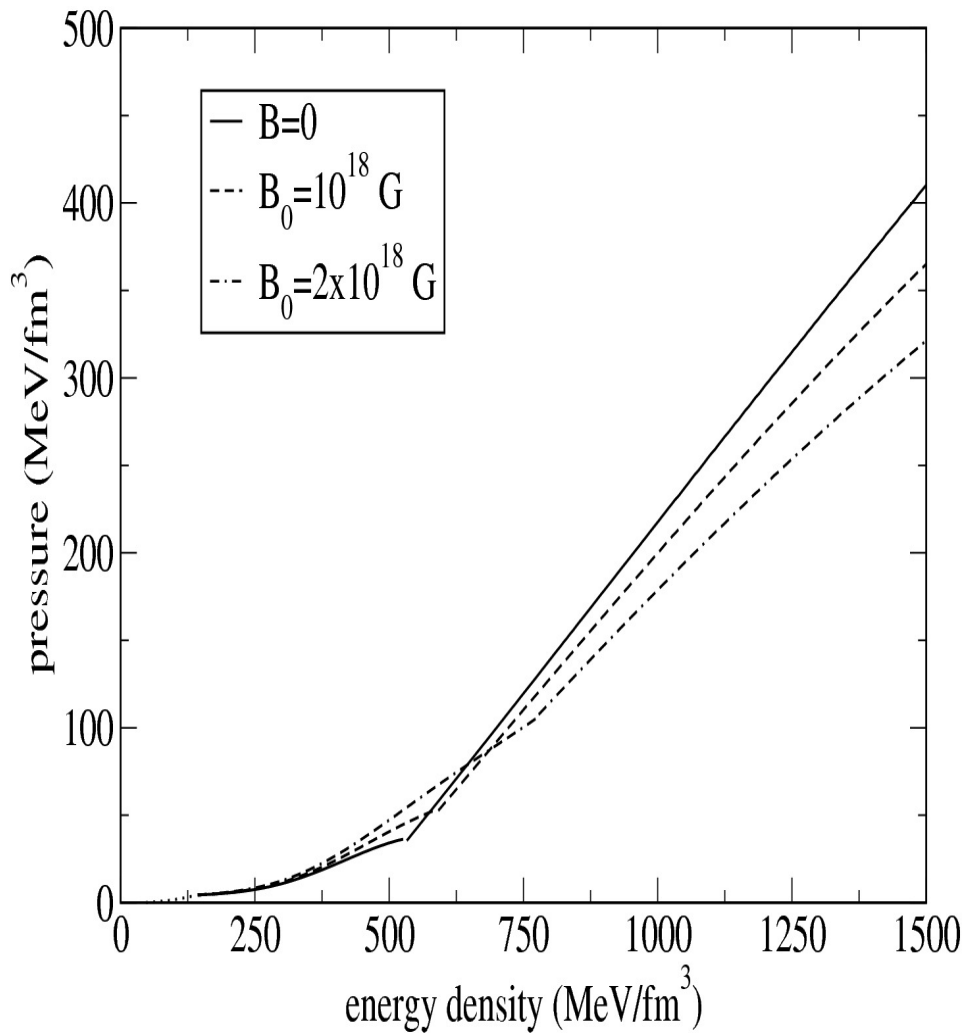
$\chi = 0$  : Hadron Phase

$$\varepsilon_{\text{MP}} = \chi \varepsilon_{\text{QP}} + (1 - \chi) \varepsilon_{\text{HP}}$$

$$n_{\text{MP}} = \chi n_{\text{QP}} + (1 - \chi) n_{\text{HP}}$$

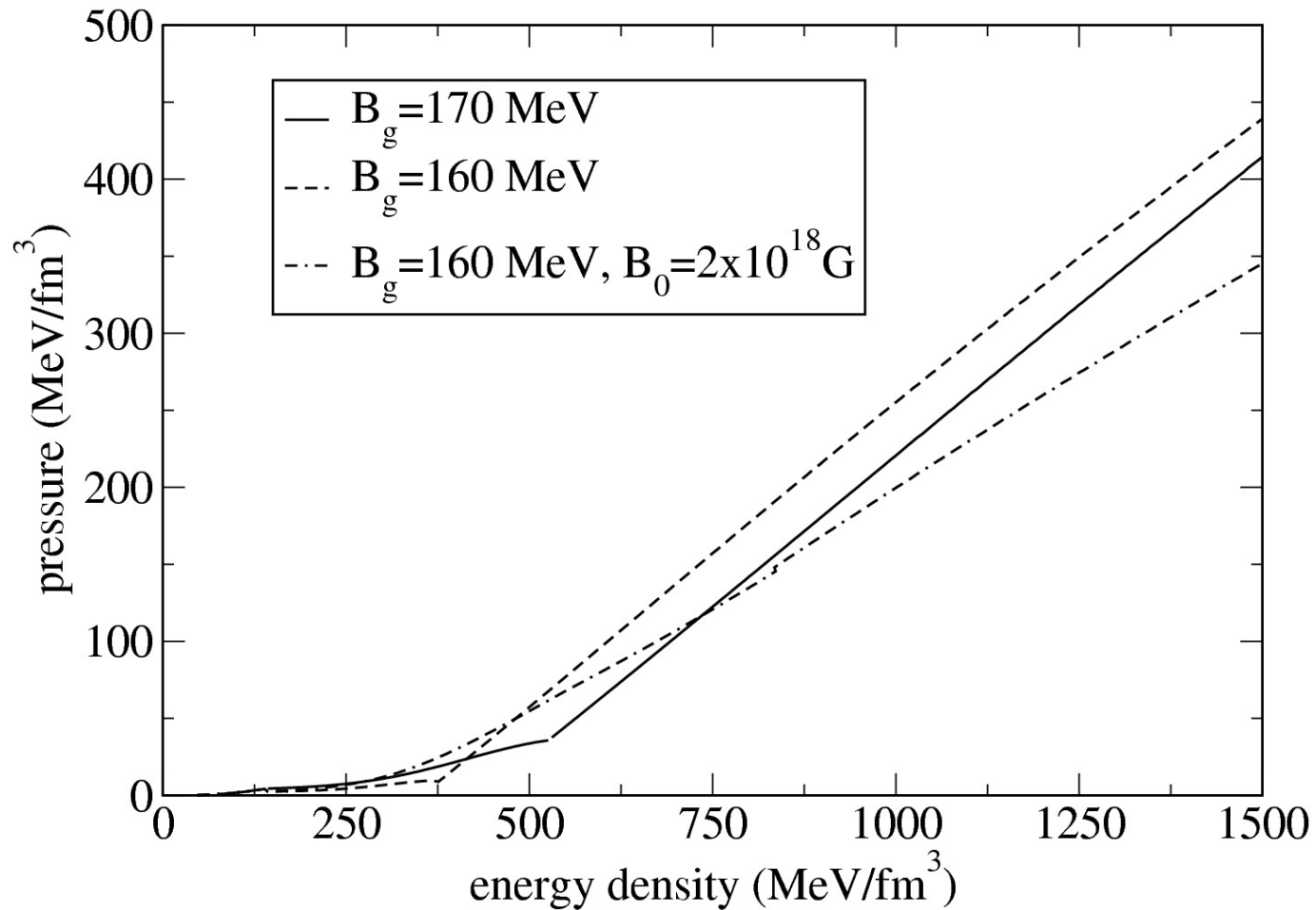
# Results

## Energy density vs Pressure



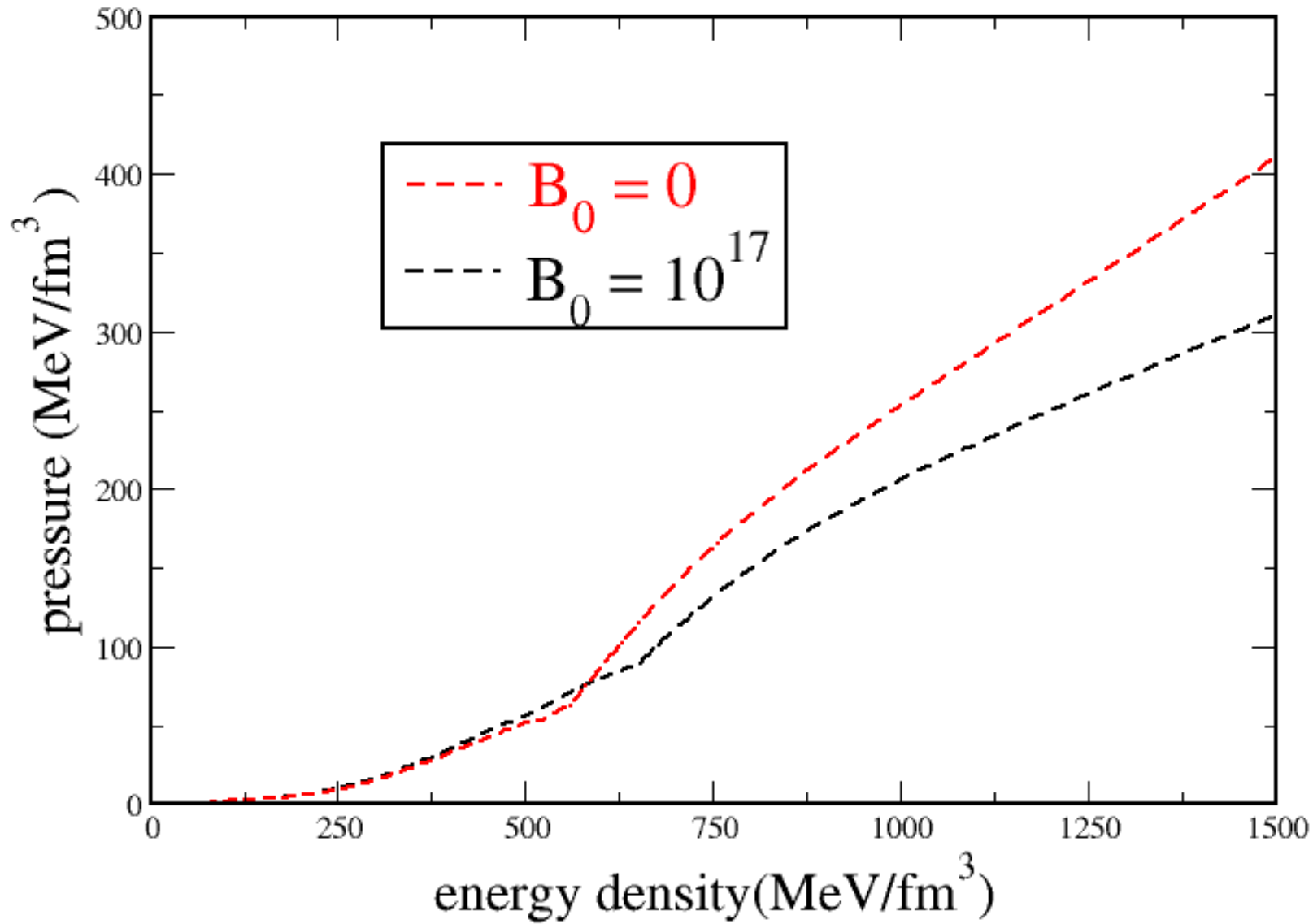
# Results

## Energy density vs Pressure



# Results

## Energy density vs Pressure



# Radial Modes



$$X \frac{d^2 \xi}{dr^2} + Y \frac{d \xi}{dr} + Z \xi = \tilde{\sigma}^2 \xi$$

$\xi(r)$  is the Lagrangian fluid displacement  $c \tilde{\sigma}$  is the characteristic eigenfrequency

The function  $X, Y, Z$  depend on equilibrium profile of pressure 'P' and energy density ' $\epsilon$ ' and metric functions  $\lambda(r)$  and  $\nu(r)$

**TOV equations :** 
$$\frac{dm(r)}{dr} = 4 \pi r^2 \epsilon(r)$$

$$\frac{dP(r)}{dr} = \frac{-Gm(r)\epsilon(r)[1+P(r)/\epsilon(r)][1+4\pi r^3 P(r)/m(r)]}{r^2 (1-2Gm(r)/r)}$$

$$\lambda(r) = \ln(1 - 2Gm/rc^2) \quad \frac{d\nu(r)}{dr} = \frac{2G}{r^2 c^2} \frac{(m + 4\pi r^3 P/c^2)}{1 - 2Gm/rc^2}$$

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$$\xi(r=0)=0$$

$$\delta P(r) = -\xi \frac{dP}{dr} - \Gamma P \frac{e^{\nu/2}}{r^2} \frac{\partial}{\partial r} (r^2 e^{\nu/2} \xi)$$

$$\delta P(r=R)=0$$

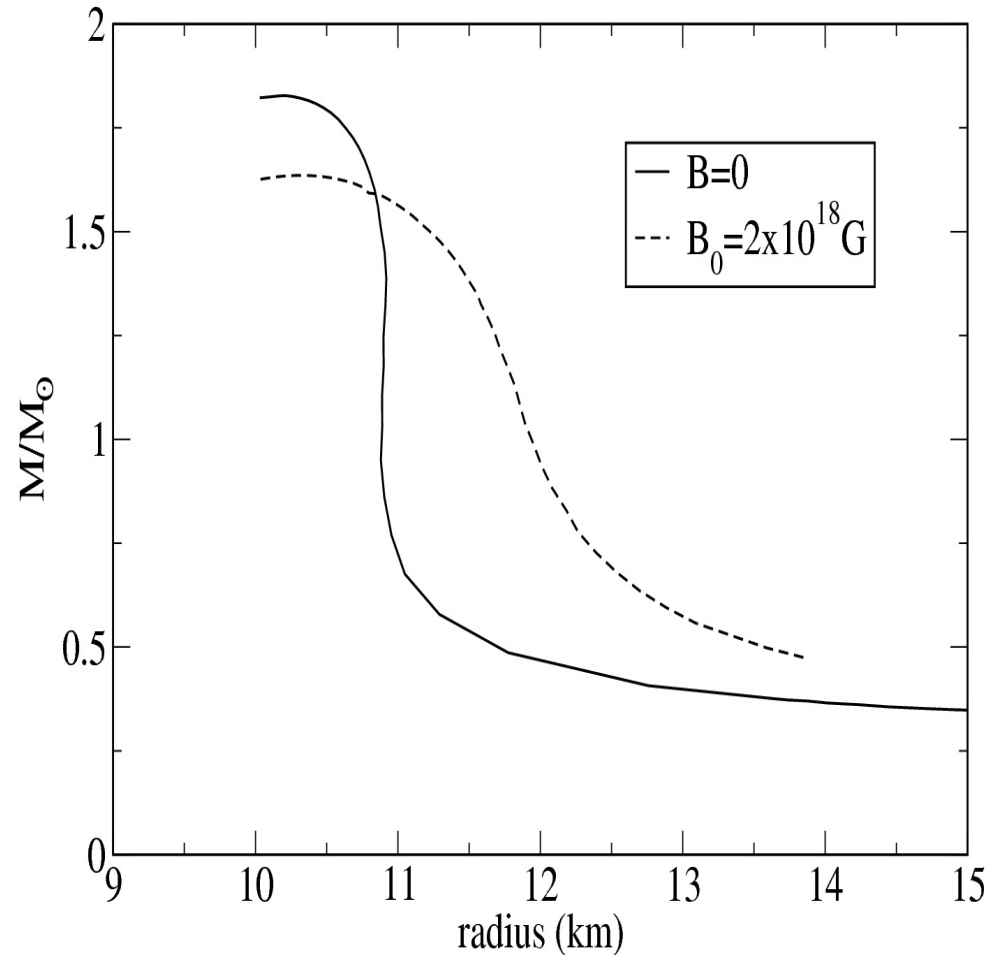
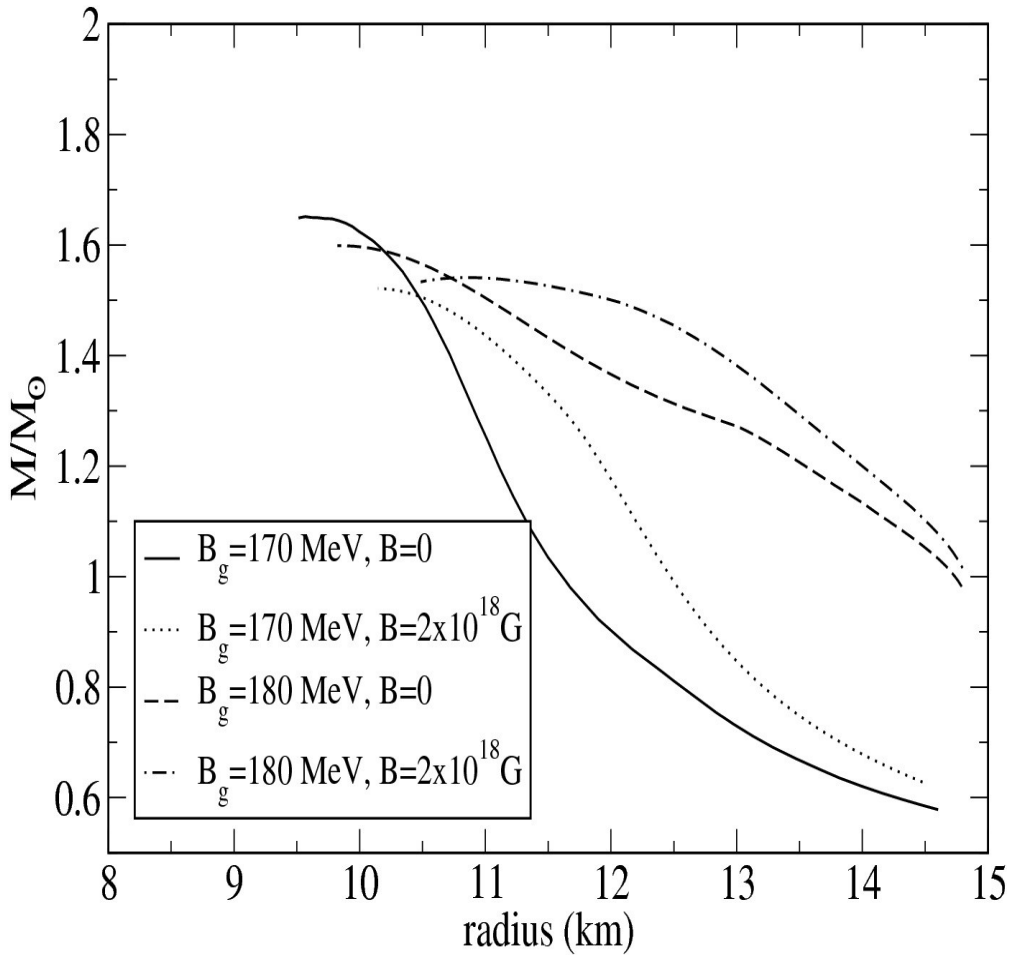
$$\tilde{\sigma}_0^2 < \tilde{\sigma}_1^2 < \tilde{\sigma}_2^2 < \dots < \tilde{\sigma}_n^2$$

Hence if fundamental radial mode of a star is stable ( $\tilde{\sigma}_0^2 > 0$ ), then all radial modes are stable

$$\text{The period of oscillation} = \frac{2\pi}{c\tilde{\sigma}}$$

# Results

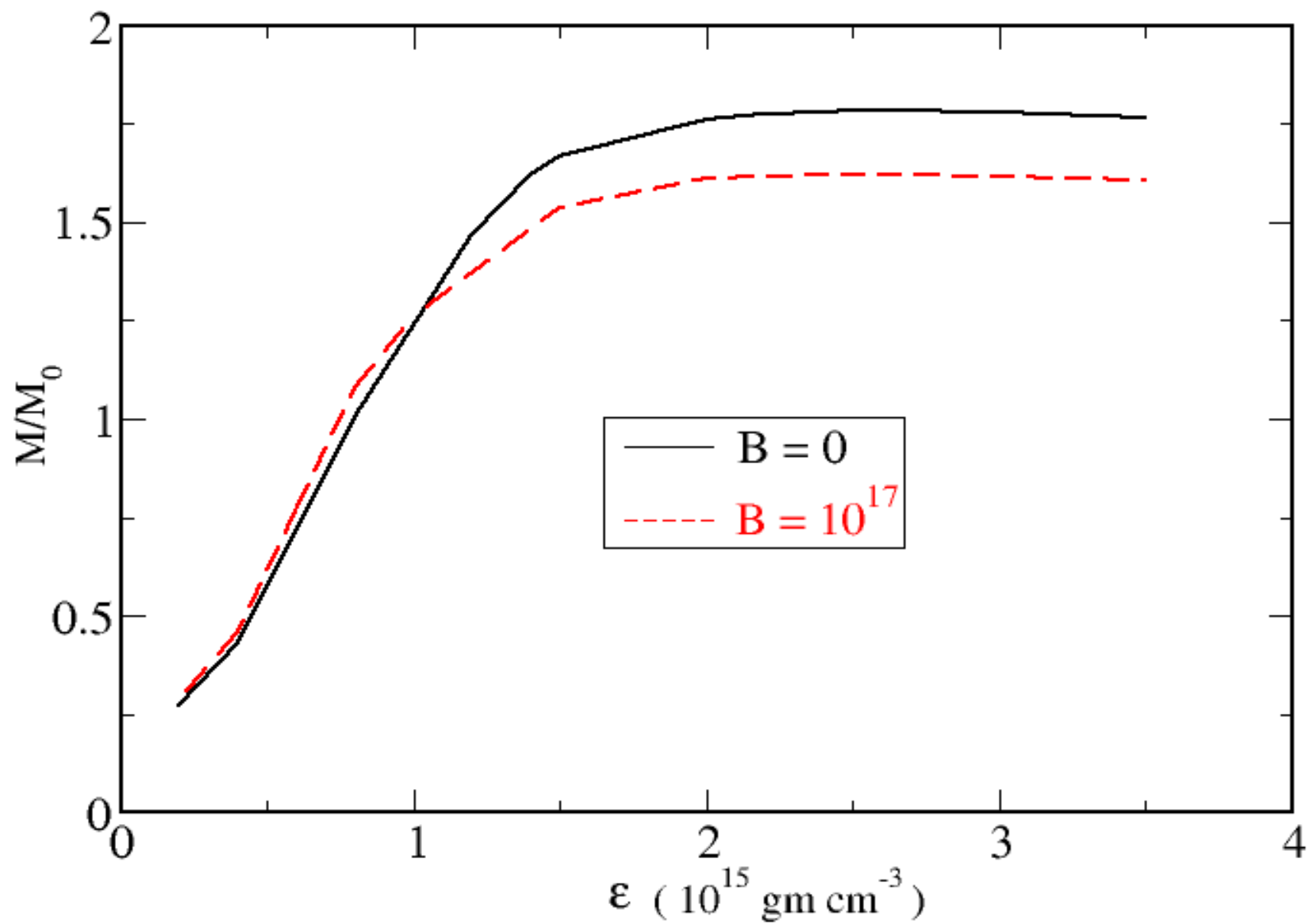
## Gravitational mass vs. Radius





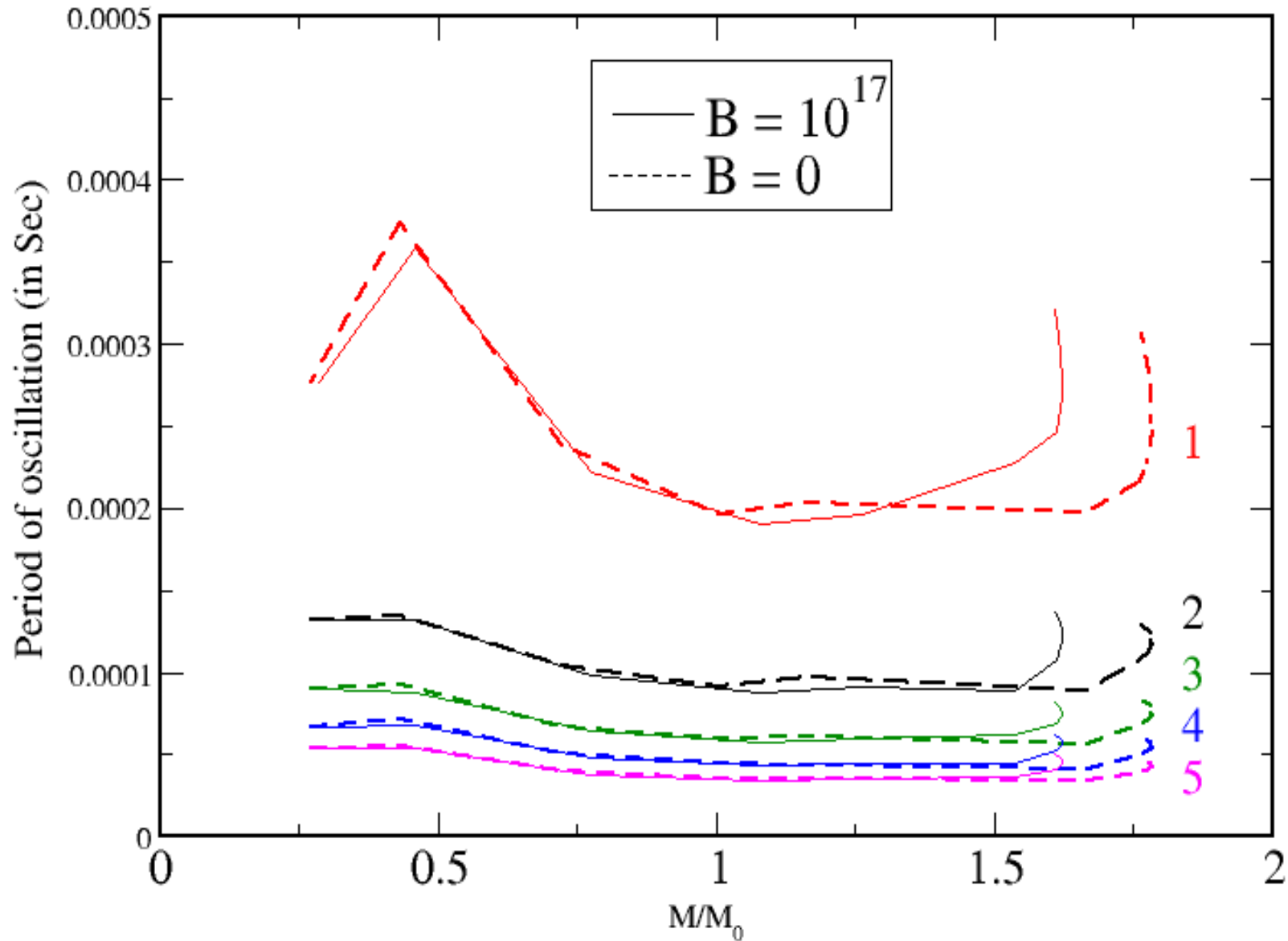
# Results

## Density vs. Gravitational mass



# Results

## Gravitational mass vs. Period of oscillation



# Summary & Remarks



- The effect of magnetic field is much more pronounced in quark matter .
- The period of oscillation shows a kink around the point where mixed phase starts, in primary as well as in the higher modes. Which is the distinct signature of quark matter onset in neutron star.
- The presence of magnetic field increases period of oscillation of fundamental as well as in higher mode at maximum mass but the effect is significant in fundamental mode.
- The presence of magnetic field broadens the mixed phase region .

Collaborators: P. K. Sahu, K. Mohanta

*Thank You*