



# Adaptive triggering for scintillation signals

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-The optimum choice of detector depends on the requirements of the experiment.

-HPGe – the highest resolution.

-Scintillator based detectors:

+Good timing properties and efficiencies at reasonable cost

- not great resolution.



The detection of ionizing radiation by the scintillation light (produced in certain materials) is one of the oldest techniques used .

The scintillation process is however one of the most useful methods nowadays available for the detection and spectroscopy of a wide assortment of radiation.



The ideal scintillation material should possess the following properties:

- ❑ It should convert the kinetic energy of charged particles into detectable light with a high scintillation efficiency.
- ❑ The conversion should be linear.
- ❑ The medium should be transparent to the wavelength of its own emission.
- ❑ The decay time of the induced luminescence should be short – Fast signal pulses.
- ❑ The material should be of good optical quality/subject to manufacture in various sizes.
- ❑ Its index of refraction should be near that of glass (in case of a coupling to a PM tube).

## No material simultaneously meets all the criteria

The most widely applied scintillators include the inorganic alkali halide crystals (of which sodium iodide is the favorite) and organic-based liquids and plastics.

Inorganic halide crystals + The best light output and linearity  
- Relatively slow in their response time



Organic scintillators + Generally faster  
- Yield less light



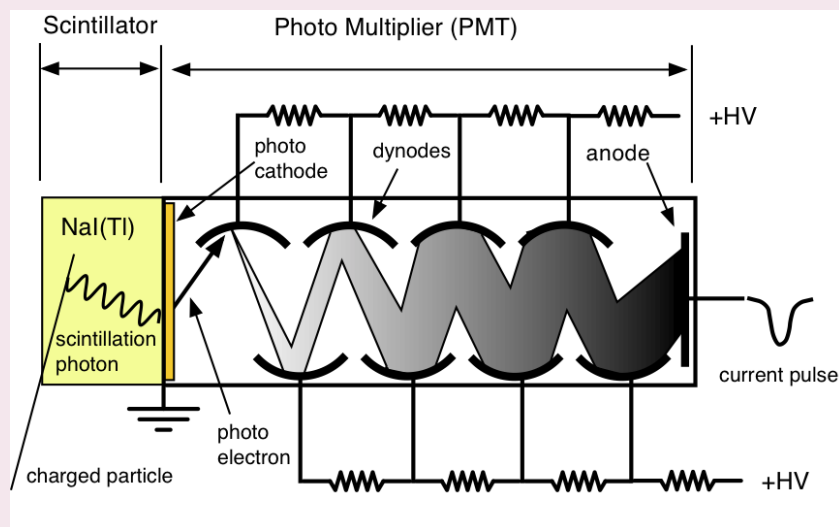
(Plastic scintillators) (+ Can be easily shaped)

The high Z-value of the constituents and high density of inorganic crystals - Gamma ray spectroscopy

Organics - Beta spectroscopy/Fast neutron detection.

Typical pulse from a scintillator – a few hundred (thousand) photons – **Too small to be measured.**

The PM tube convert a weak pulse from the scintillator into a corresponding electrical signal –  $10^7$ - $10^{10}$  electrons (Without adding a large amount of noise to the signal).



## Signal evolution:

1. Energy is absorbed in scintillator.
2. Population of states that emit photons.
3. Decay of radiative states: Governed by Poisson distribution:

$$P(x) = \frac{(pn)^x e^{-pn}}{x!} = \frac{\mu^x e^{-\mu}}{x!}$$

$$\mu = \sum_{x=0}^n xP(x) = pn$$

$$\sigma^2 = \sum_{x=0}^n (x - \mu)^2 P(x) = pn$$

4. Photons absorbed in photocatode by expelling photoelectrons.
5. Finally the photoelectrons are multiplied through the gain structure.

$$I_{anode}(t) = G \cdot Q \cdot N_o \cdot e^{-t/\tau}$$

## Precision of signal magnitude is limited by fluctuations

1. Fluctuations due to the stochastic nature of the pulse creation.
2. Baseline fluctuations due to the electronics – „Electronic noise“

### Main contribution to the fluctuations in scintillator signals - Statistical fluctuations

The output pulses are not all the same shape but rather „noised“ with statistical fluctuations on the pulse tails -Afterpulsing.

### Caution - Some non statistical processes can also give rise to the afterpulsing

- The emission of light from the latter stages of the multiplier structure which comes back to the photocatode.
- Imperfect vacuum in the tube – Traces of residual gas can be ionized.

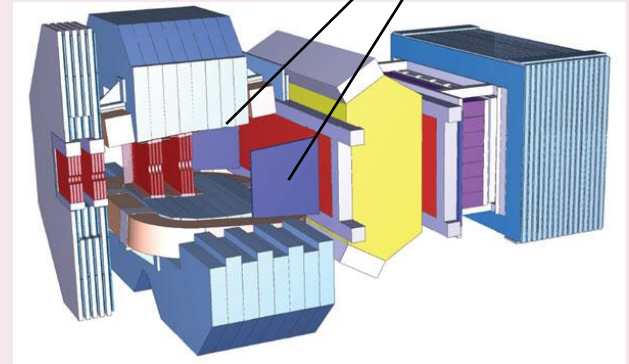
# Scintillators @



TOF measurements with relativistic heavy – ion beams



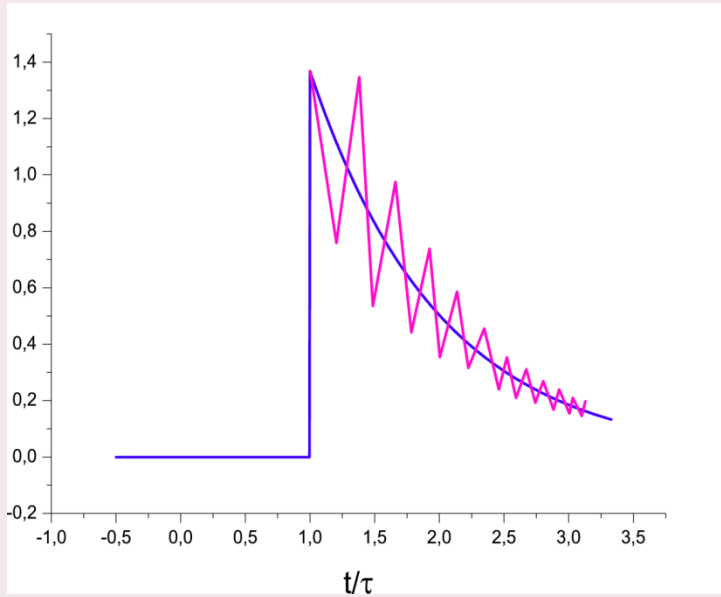
Plastic scintillator Time-Of-Flight system. (PANDA)



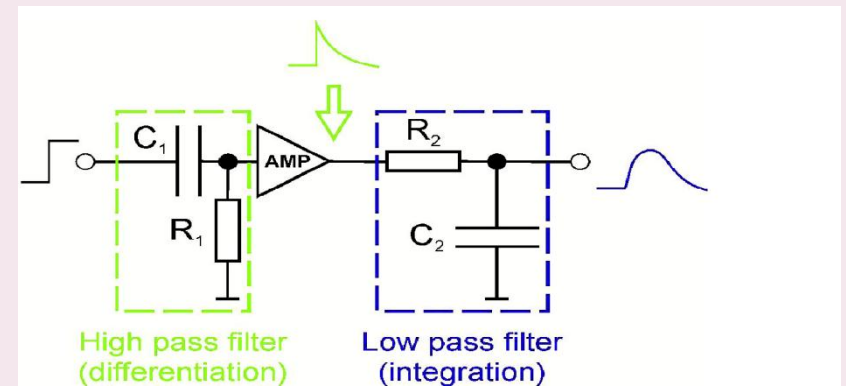
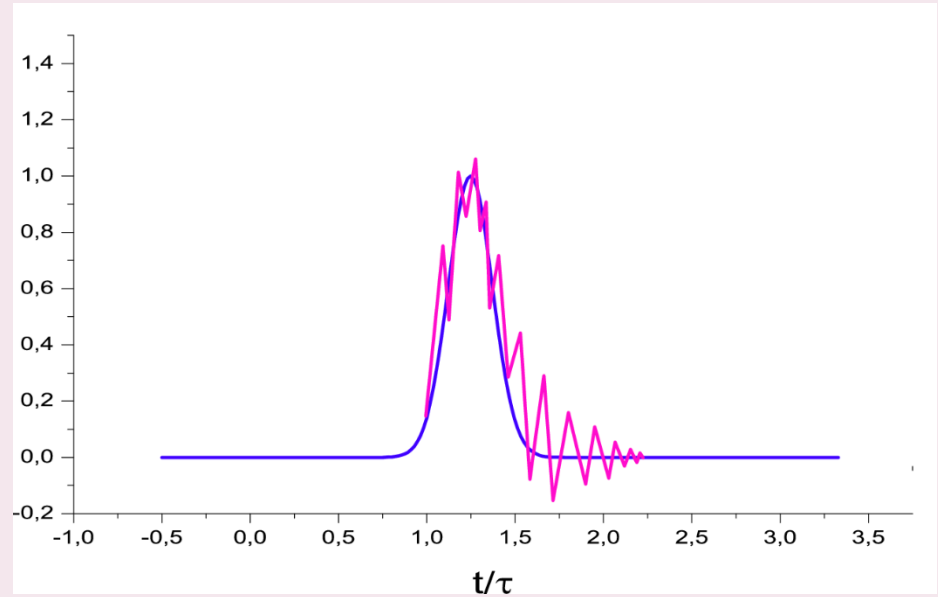


# Afterpulsing+Analog electronics

Area=Energy



Amplitude=Energy



## Moving window deconvolution: (MWD)

The single exponential decay with a starting time  $t=0$  is written as:

$$f(t) = Ae^{-t/\tau}, t \geq 0$$
$$0, t < 0$$

Knowing the value  $f(t_n)$  at time  $t_n$  the initial amplitude can be found as:

$$A = f(t_n) + A - f(t_n) = f(t_n) + A(1 - e^{-\frac{t_n}{\tau}}) = f(t_n) + \frac{\int_0^{t_n} f(t) dt}{\tau}.$$

In digital domain:

$$A[n] = x[n] + \frac{\sum_{k=0}^{n-1} x[k]}{\tau} = x[n] - \left(1 - \frac{1}{\tau}\right)x[n-1] + A[n-1].$$

Differentiate the deconvoluted signal – obtain the MWD equation with a window  $M$ :

$$MWD_M[n] = A[n] - A[n-M] = x[n] - x[n-M] + \frac{1}{\tau} \sum_{k=n-M}^{n-1} x[k].$$

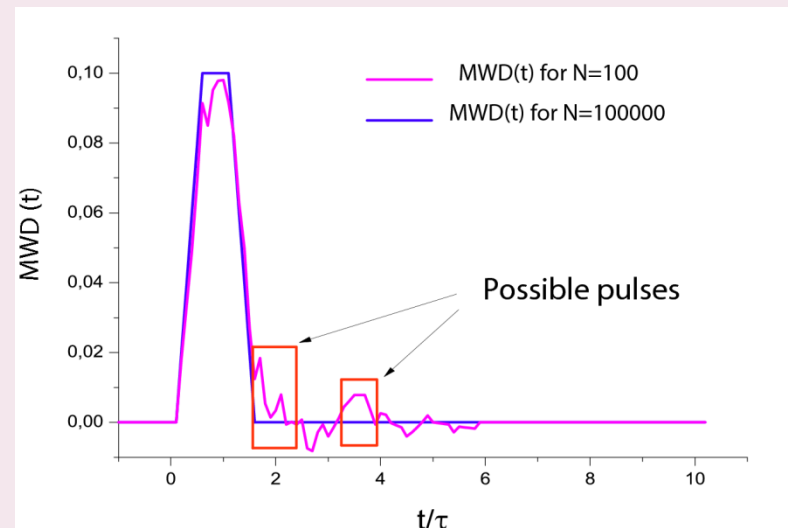
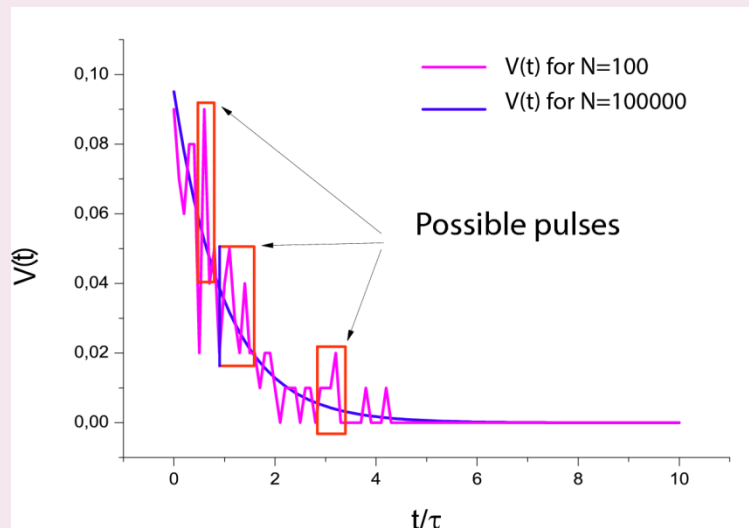
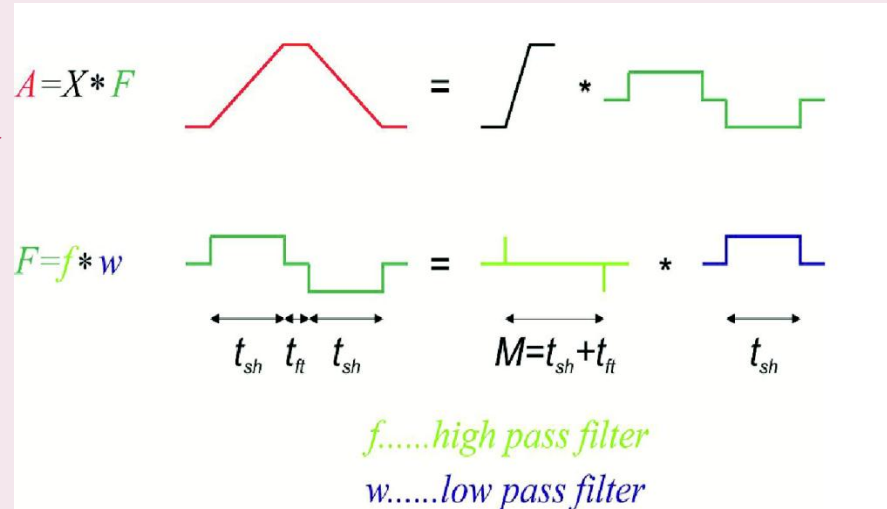
The MWD algorithm converts an exponentially decaying signal into a step signal of length  $M$ .

Signal to noise ratio?

To improve the signal to noise ratio, a low pass filter is applied after the MWD module.

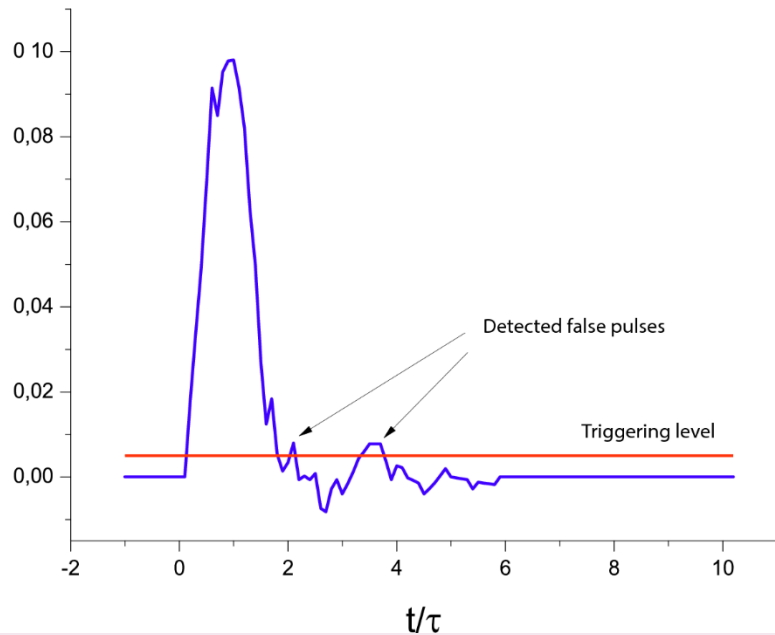
$$f_n = \begin{cases} 1 & n = 1 \\ \frac{1}{\tau} & 1 < n < M \\ \frac{1}{\tau} - 1 & n = M \end{cases} \quad M = t_{sh} + t_{ft}$$

$$w_n = \begin{cases} 1/N_{sh} & 1 \leq n \leq N_{sh} \\ 0 & \text{otherwise} \end{cases} \quad N_{sh} = t_{sh}$$

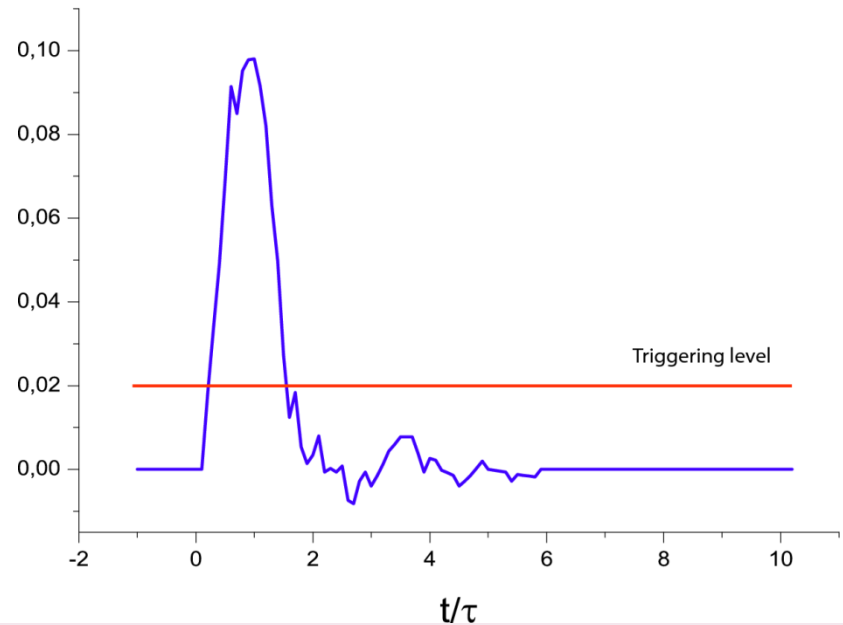


# State of art – Increase trigger level/Deadtime

Low triggering level:

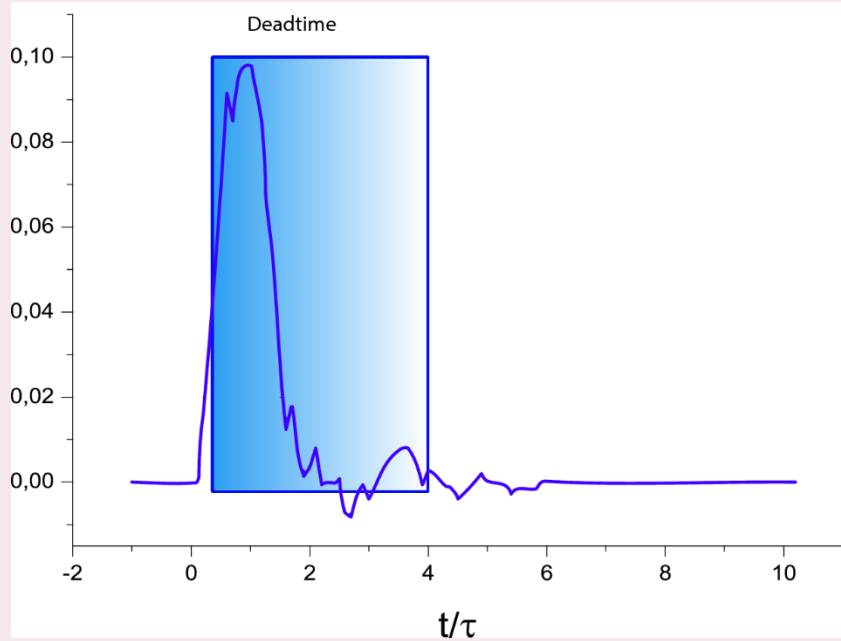


Increase triggering level:



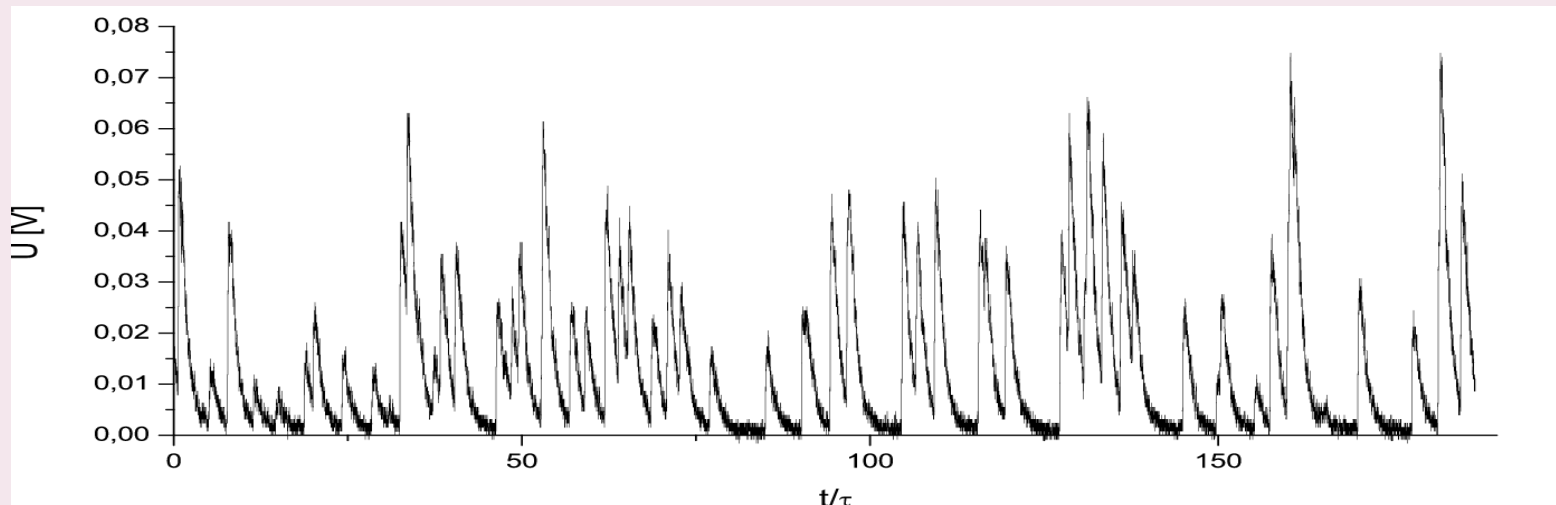
Lowers the dynamic range of the detector.

# Deadtime



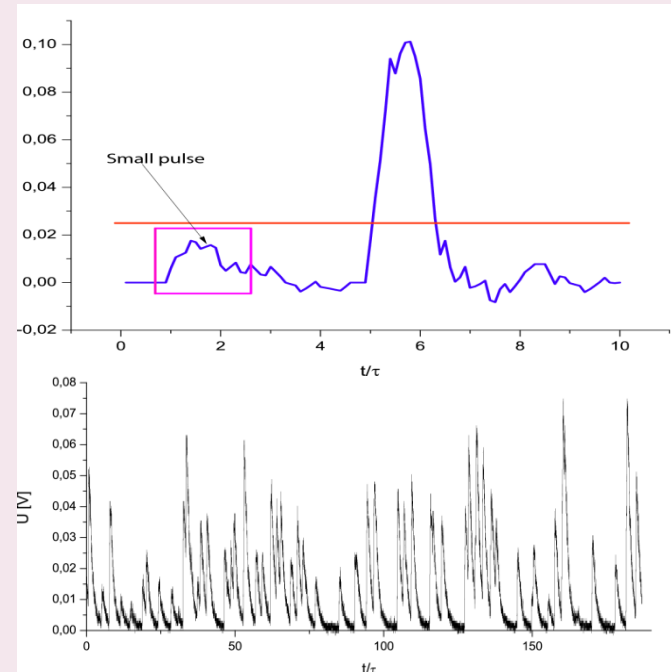
Introduce a **deadtime** after each pulse.

What about a high count rates ?



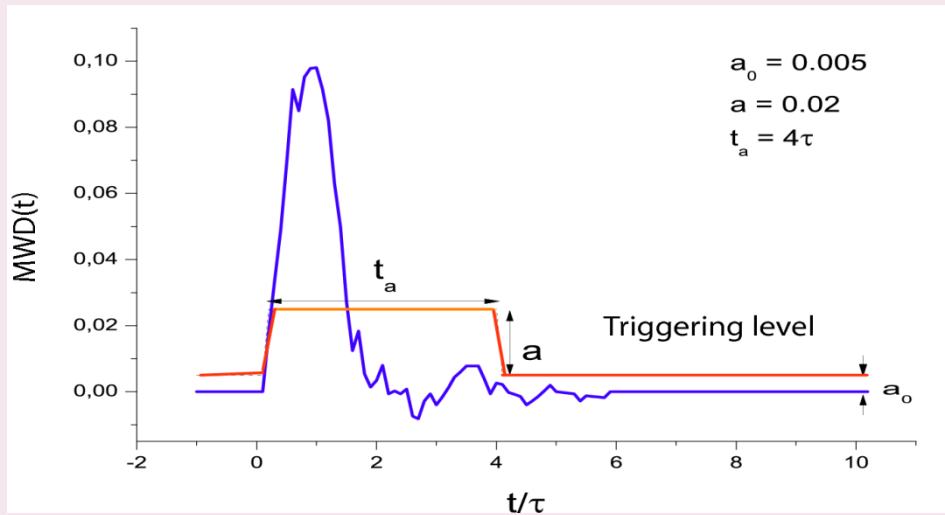
Increasing trigger level-Lowers the dynamic range.

Introducing the deadtime - Makes the detector inefficient at higher count rates.

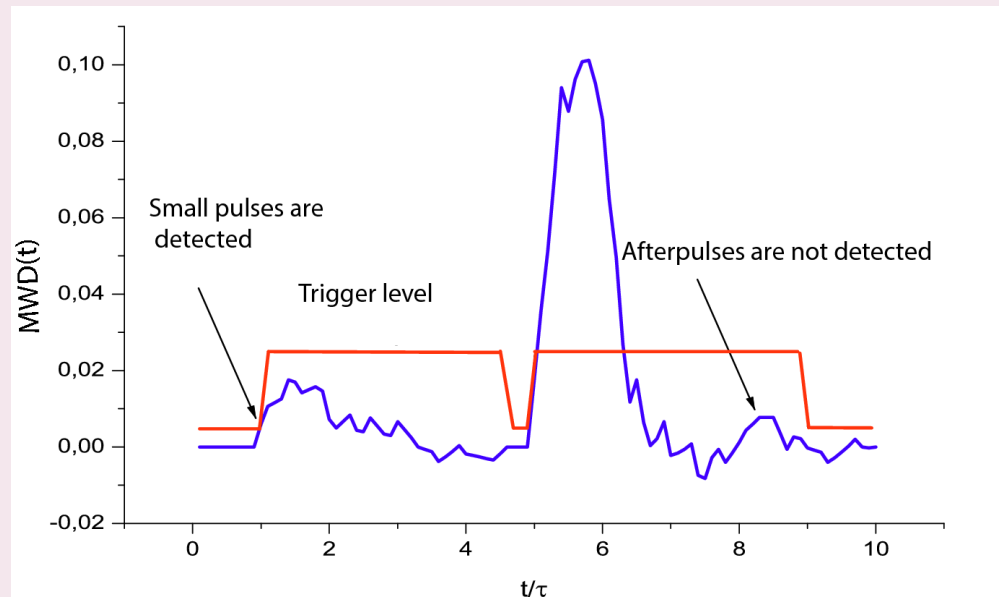


Solution – Raise the trigger by a finite value and for a finite amount of time after each pulse.

# Adaptive triggering for scintillation signals



$$\text{THR}(t) \begin{cases} a_0, t \leq t_0 \\ a_0 + a, t_0 < t < t_0 + t_a \\ a_0, t_0 + t_a \leq t \end{cases}$$



# Simulation

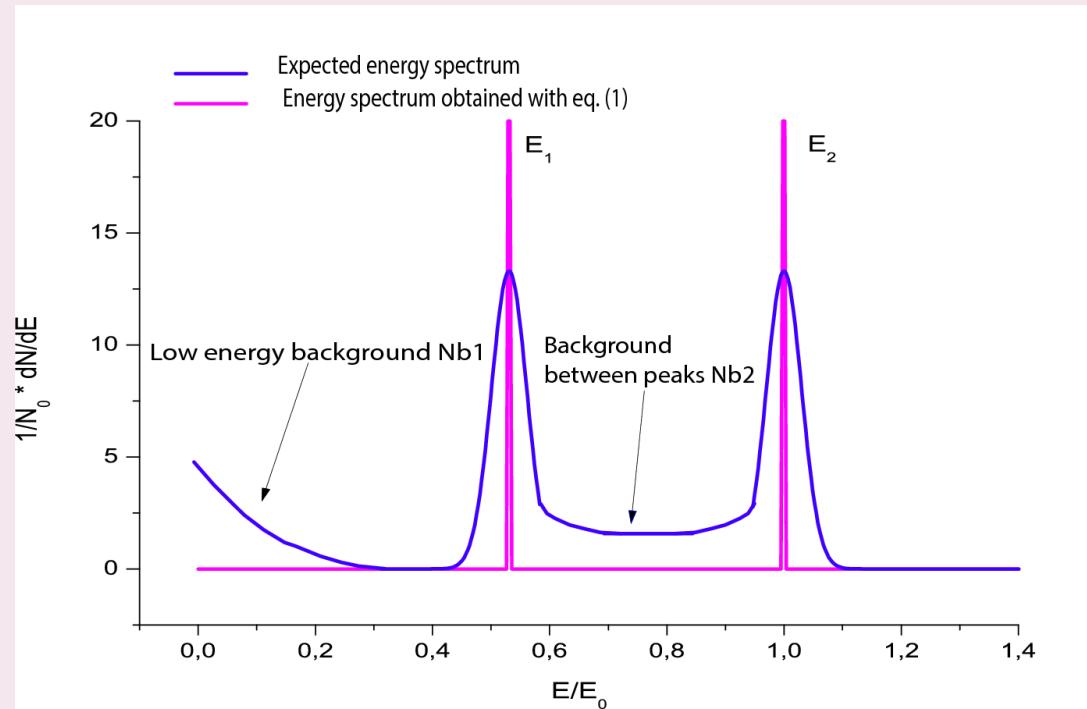
1. Simple energy spectrum given by the equation:

$$\frac{1}{N_0} \frac{dN(E)}{dE} = \frac{1}{2} (c\delta(E - E_1) + (1 - c)\delta(E - E_2)), \quad c = \frac{N_1}{N_0}$$

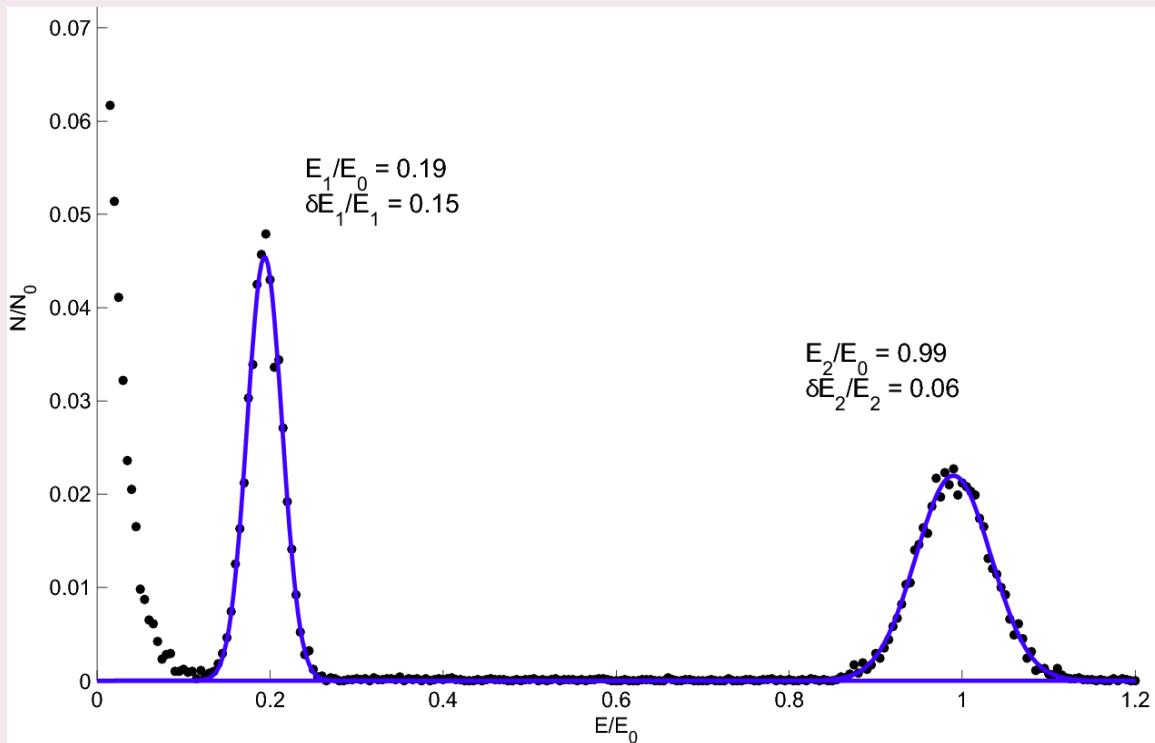
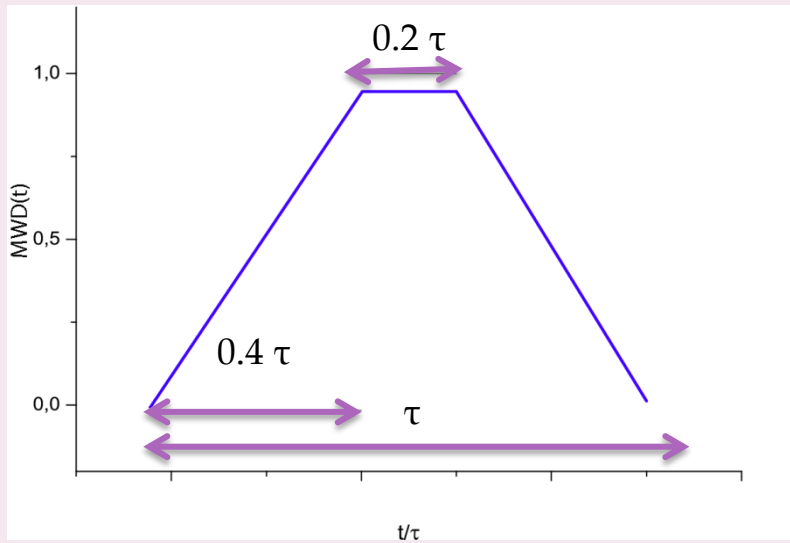
2. Every event is described with the time of event and energy (Energy is randomly sampled from the energy spectrum, time is sampled from uniform time distribution).
3. Every event excites  $n$  states proportional to its energy (Decay times are sampled from exponential distribution).
4. Every photon gets time stamp.
5. Put the time stamps into histogram with a bandwidth  $\omega$ .
6. Histogram values - Amplitude of a current pulse from a photomultiplier.



| Simulation parameters   | Parameter values   |
|---|--|
| Energy spectrum   | $\frac{1}{N_0} \frac{dN(E)}{dE} = \frac{1}{2} (c\delta(E - E_1) + (1 - c)\delta(E - E_2))$ (1) |
| Beam frequency  | $10^5$ Hz  |
| Decay time  | 250 ns   |
| Proportionality constant<br>(Reference energy corresponds to 1000<br>detected photoelectrons) | 1000   |
| Bandwidth   | 200 MHz  |



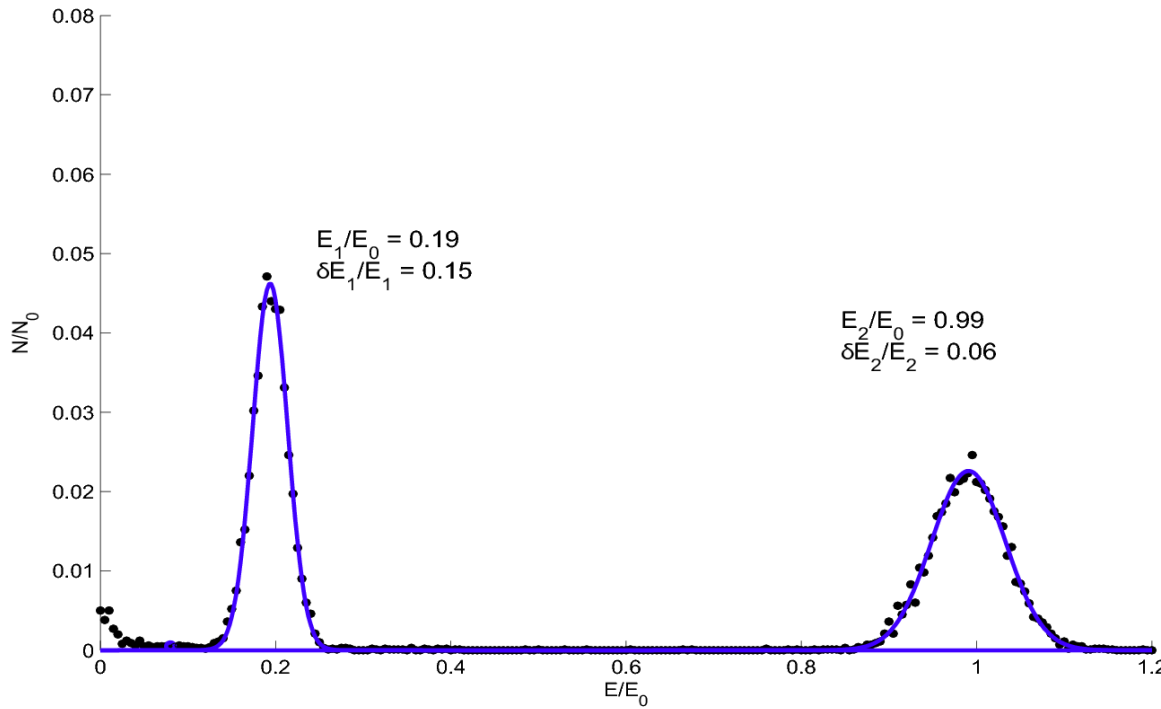
MWD parameters:



$E_1/E_0=0.2$   $E_2/E_0=1$ ,  $c=0.5$

Adaptive triggering OFF  
 THR=0.025  $E_0$

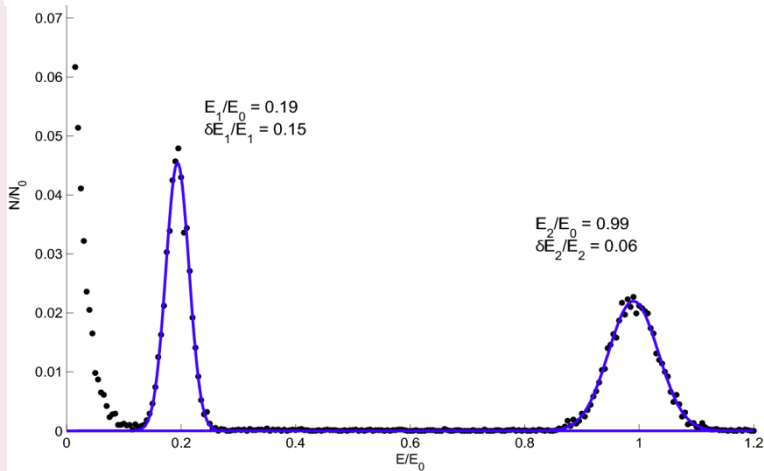
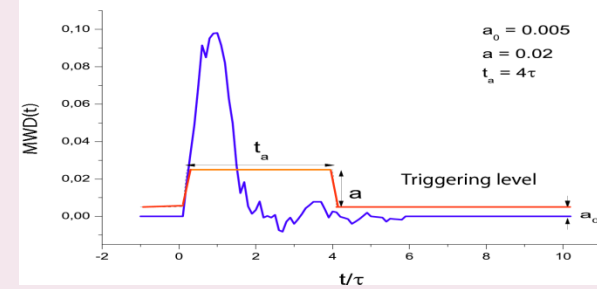
$E_1/E_0=0.2$   $E_2/E_0=1$ ,  $c=0.5$



Adaptive triggering ON

$a=0.1 E_0$

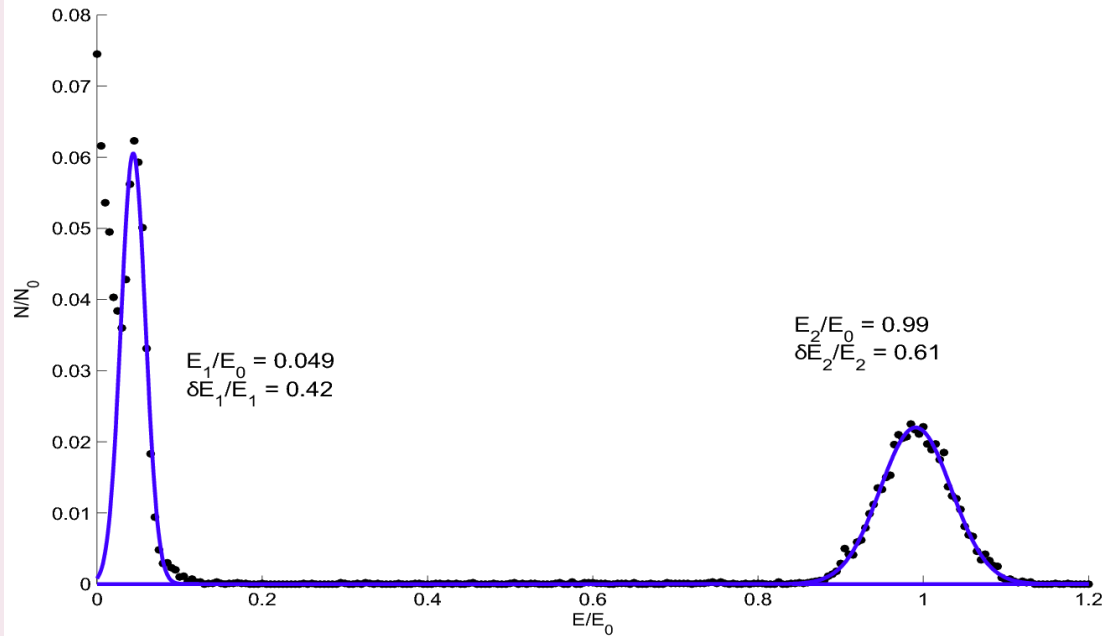
$t_a=4\tau$



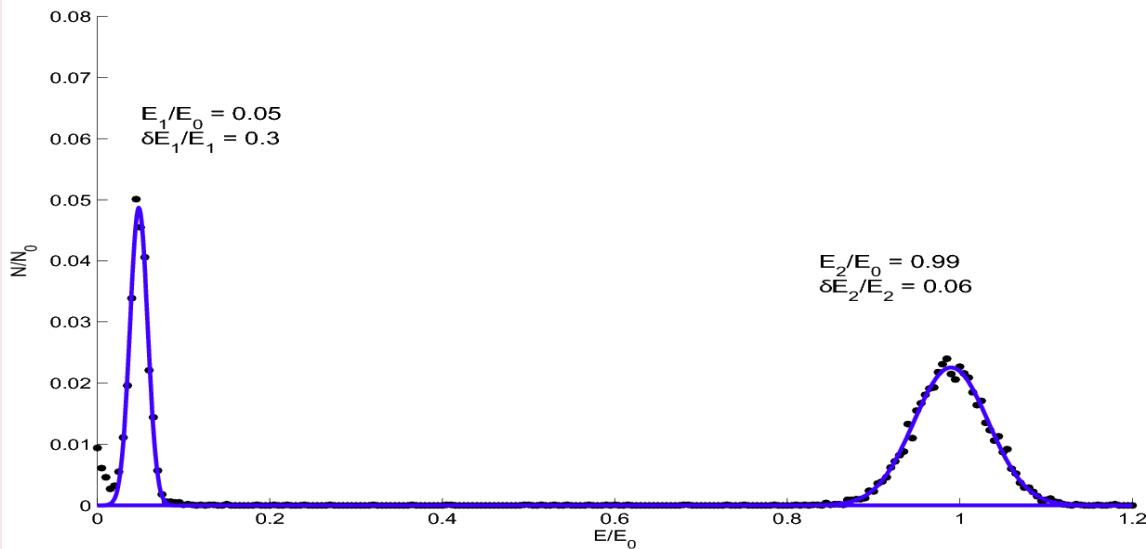
Low energy background – Due to the afterpulses

Background between peaks – Real pulses that follow the afterpulses.

$E_1/E_0=0.05$   $E_2/E_0=1$ ,  $c=0.5$

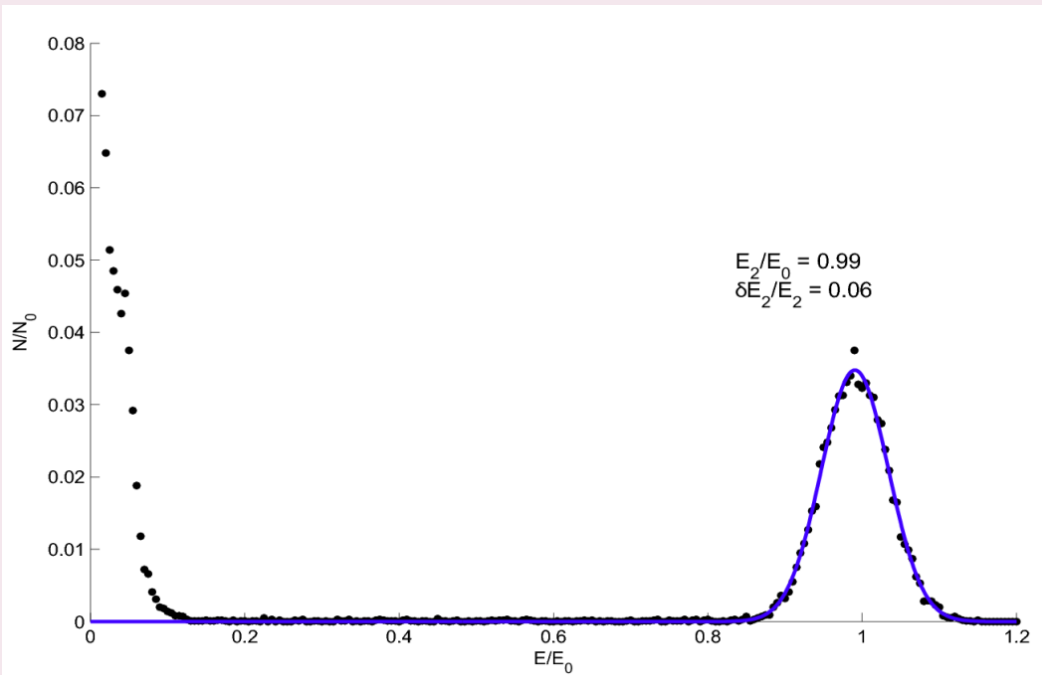


Adaptive triggering OFF  
THR=0.025  $E_0$

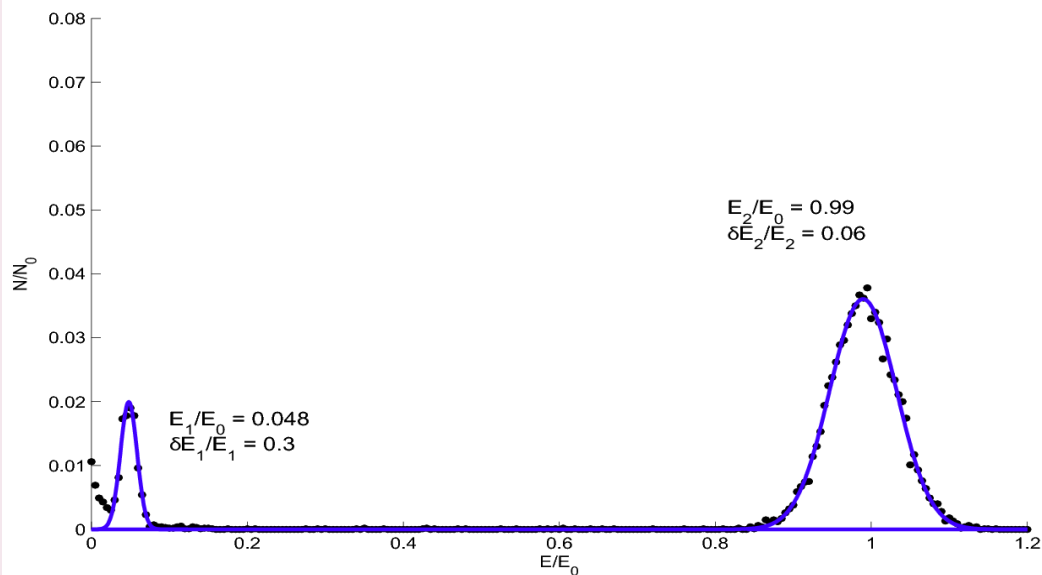


Adaptive triggering ON  
 $a=0.1 E_0$   
 $t_a=4\tau$

$E_1/E_0=0.05, E_2/E_0=1, c=0.2$

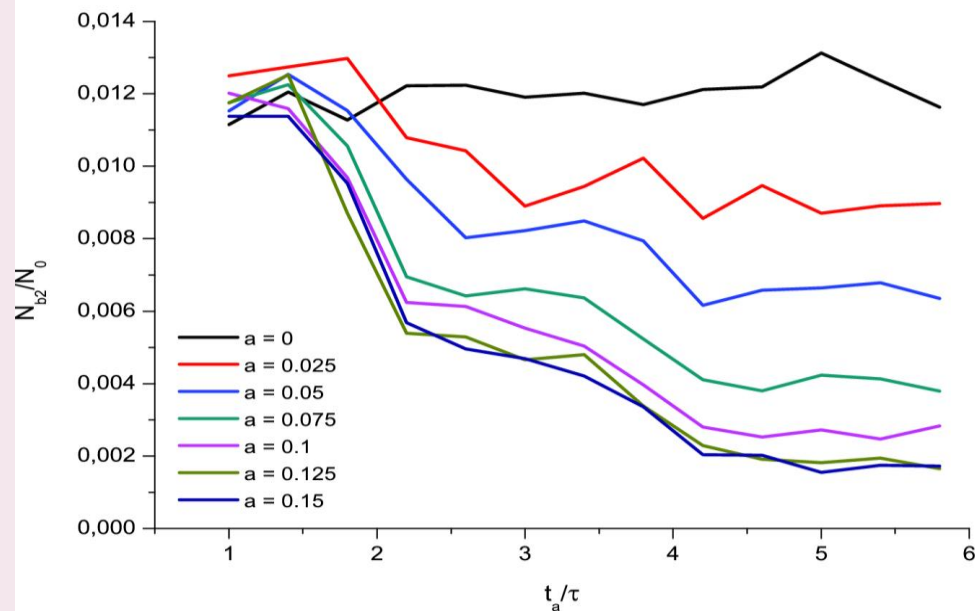
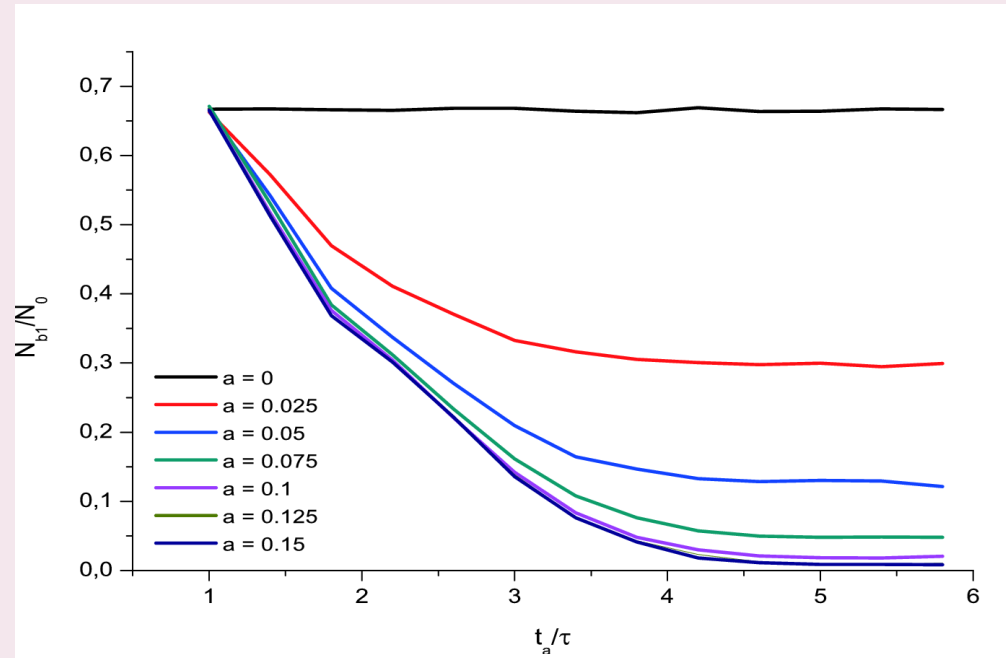


Adaptive triggering OFF  
THR=0.025  $E_0$



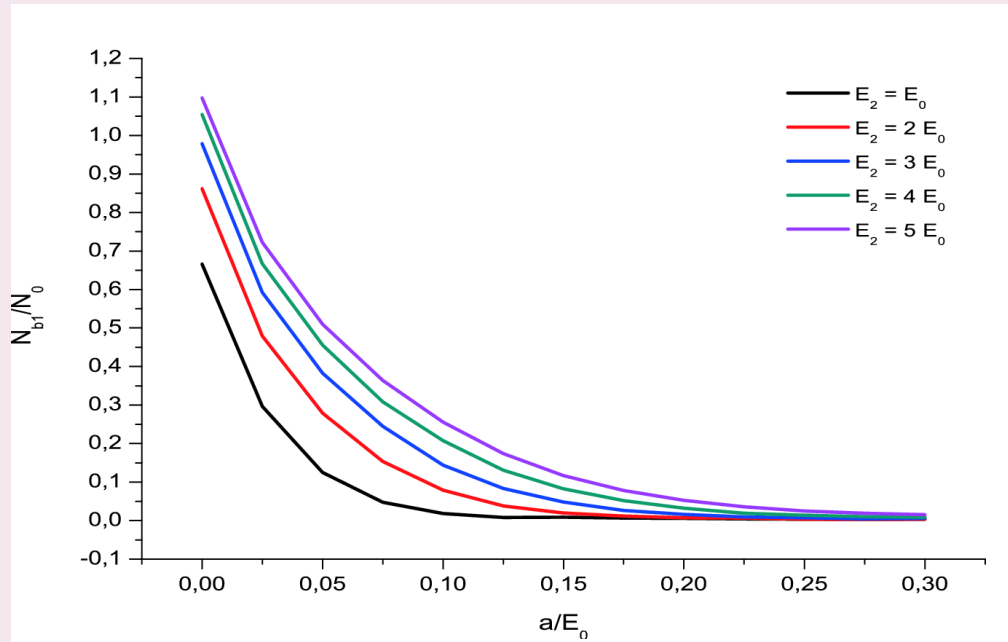
Adaptive triggering ON  
 $a=0.1 E_0$   
 $t_a=4\tau$

$E_1/E_0=0.2, c=0.2$

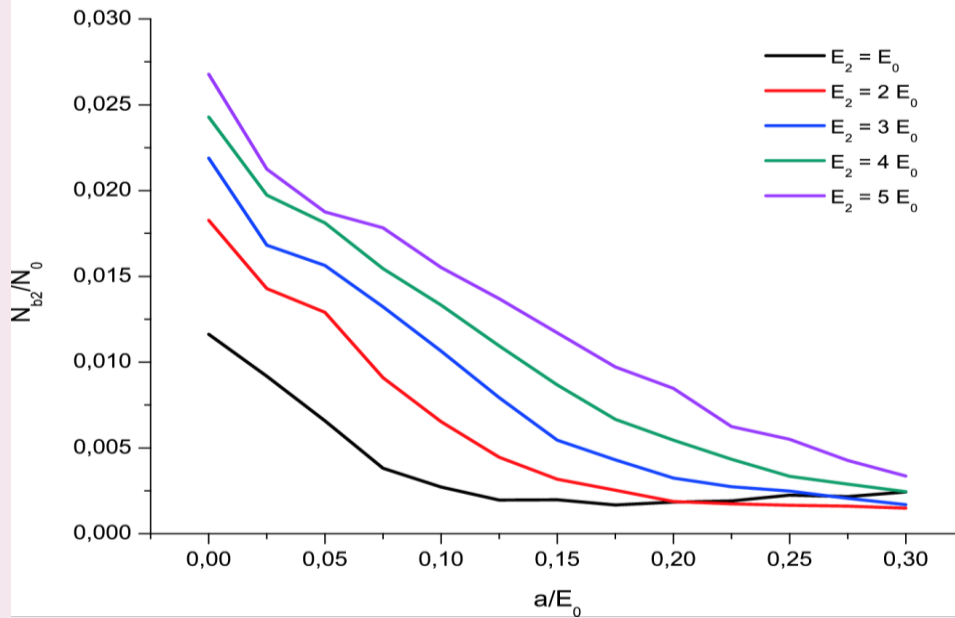


THR=0.025  $E_0$

$$E_1/E_0=0.2, c=0.2$$

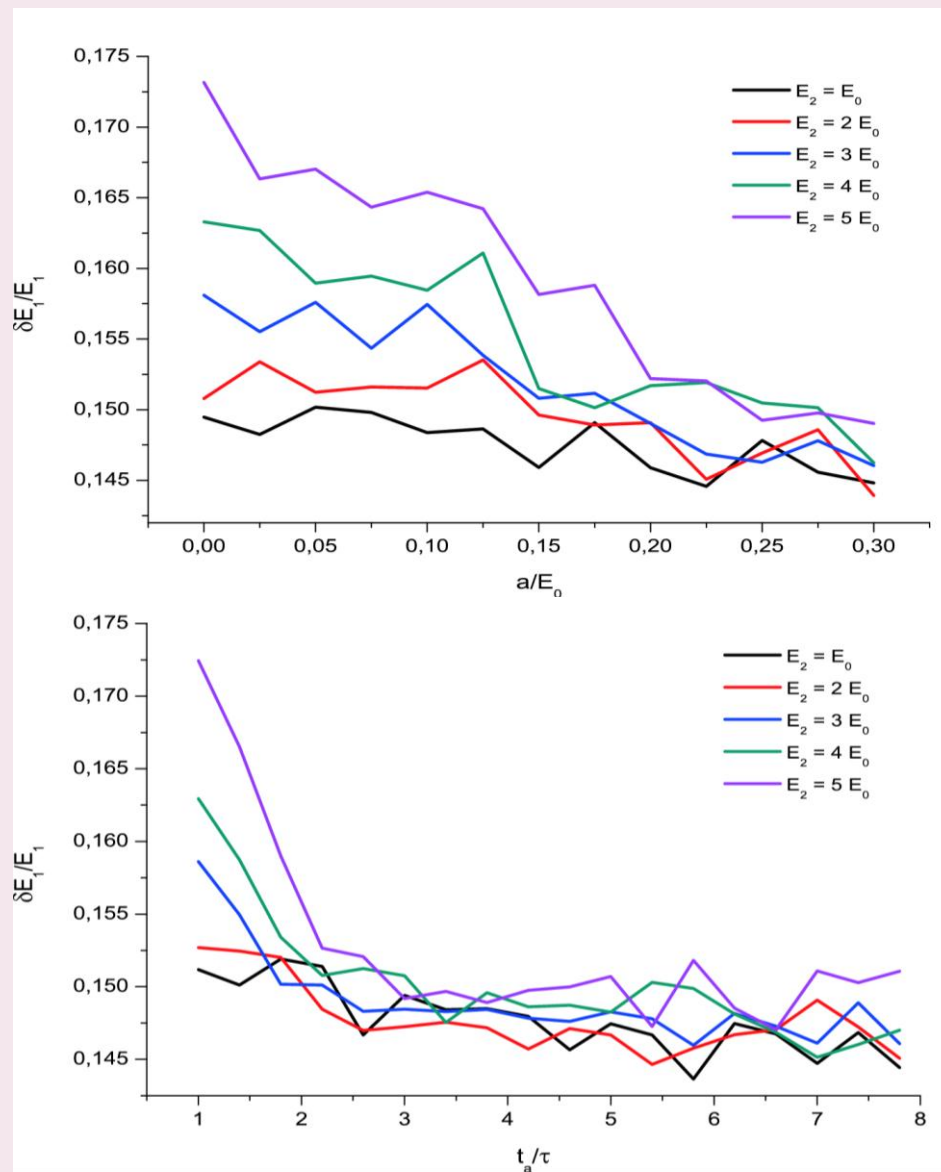


$t=\tau$



$\text{THR}=0.025 E_0$

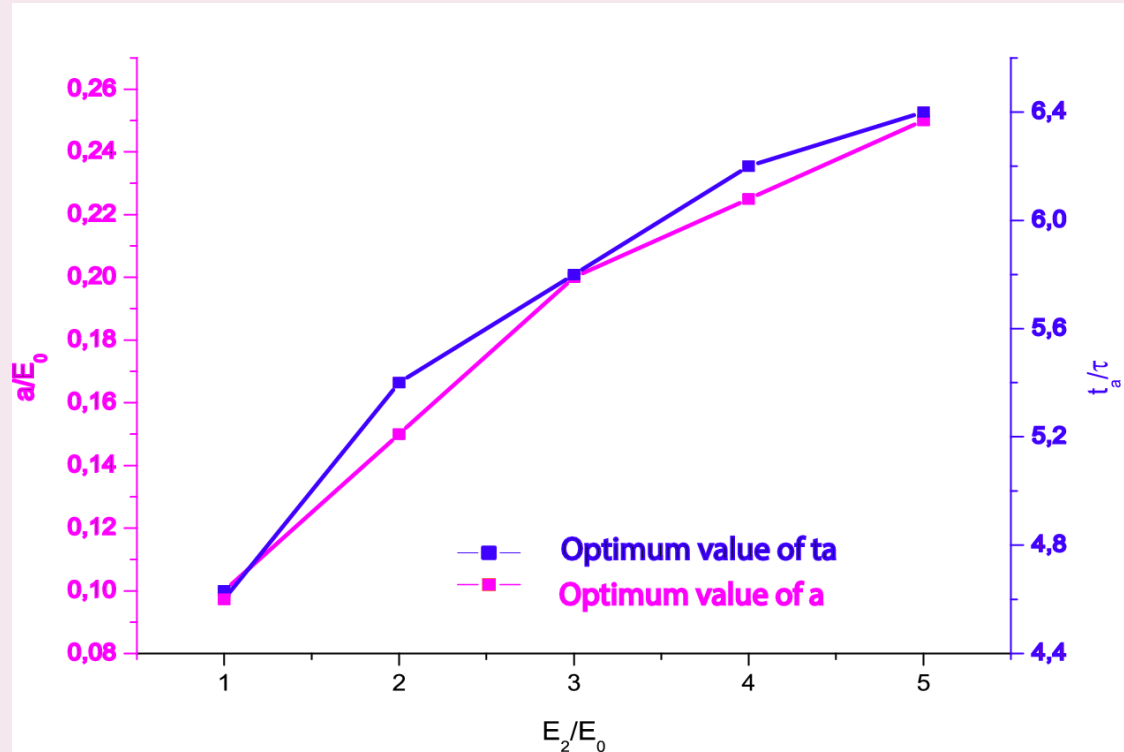
$$E_1/E_0=0.2, c=0.2$$





# The optimal triggering parameters

$E_1/E_0=0.2, c=0.2$



$$\frac{t_{opt}}{\tau} = 3.76 + 0.95 \frac{E}{E_0} - 0.086 \left( \frac{E}{E_0} \right)^2$$

$$\frac{a_{opt}}{E_0} = 0.035 + 0.07 \frac{E}{E_0} - 0.005 \left( \frac{E}{E_0} \right)^2$$

## Test of the optimal triggering parameters – simulated signal

Simulated sintetic signal using an energy spectrum given by eq.:

$$\frac{1}{N_0} h_{ideal}(E) = \frac{1}{N_0} \frac{dN(E)}{d(E)} = \begin{cases} \frac{1}{E_1} & ; E < E_1 \\ 0 & ; E > E_1 \end{cases}$$

The quality of parameters was deduced using:

$$\sigma_{spect} = \left( \frac{1}{n-1} \sum_{i=1}^n (h(E_i) - h_{ideal}(E_i))^2 \right)^{1/2}$$

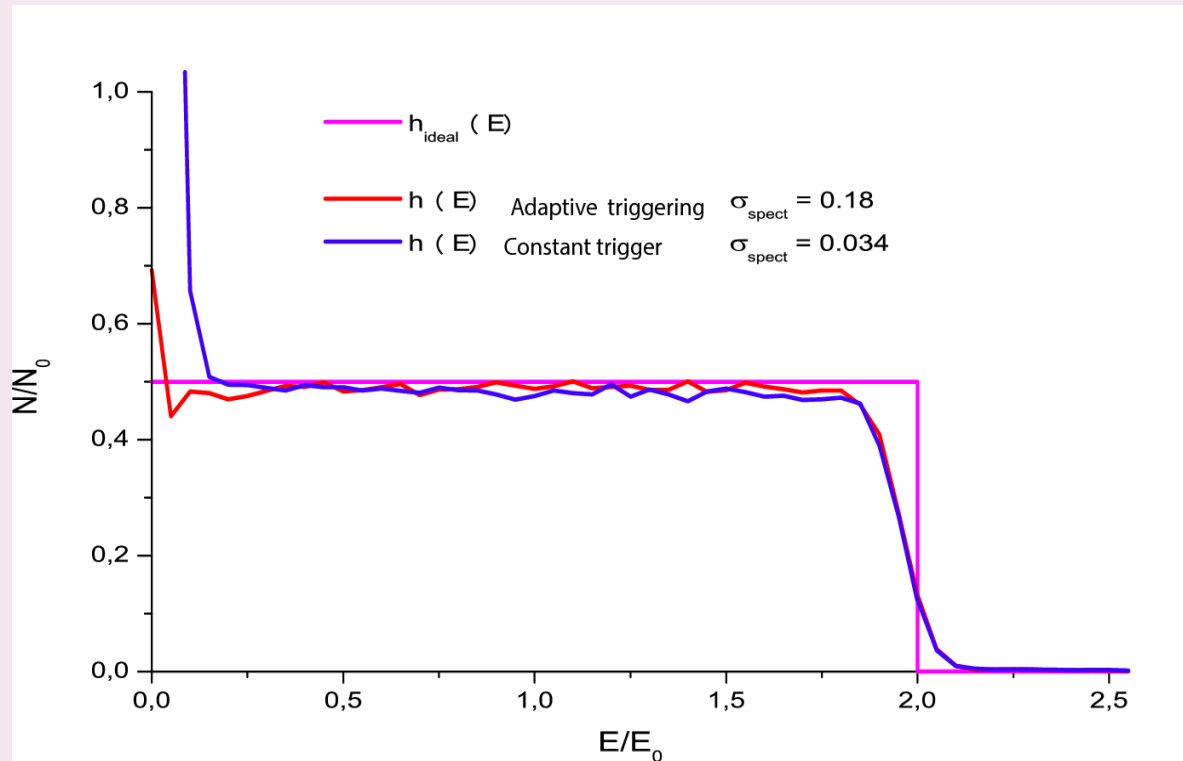
Adaptive triggering OFF:  $\sigma_{spect}$ .

Adaptive triggering ON: With every simulated pulse we have chosen new  $t$  and  $a$ ,  $\sigma_{spect}/30$ .

$$\left( \begin{array}{l} \frac{t_{opt}}{\tau} = 3.76 + 0.95 \frac{E}{E_0} - 0.086 \left( \frac{E}{E_0} \right)^2 \\ \frac{a_{opt}}{E_0} = 0.035 + 0.07 \frac{E}{E_0} - 0.005 \left( \frac{E}{E_0} \right)^2 \end{array} \right)$$

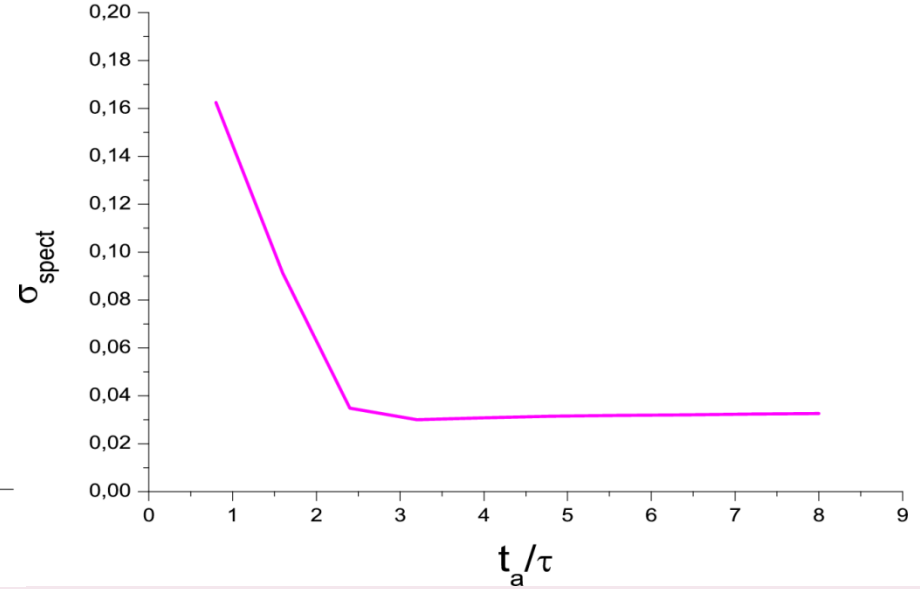
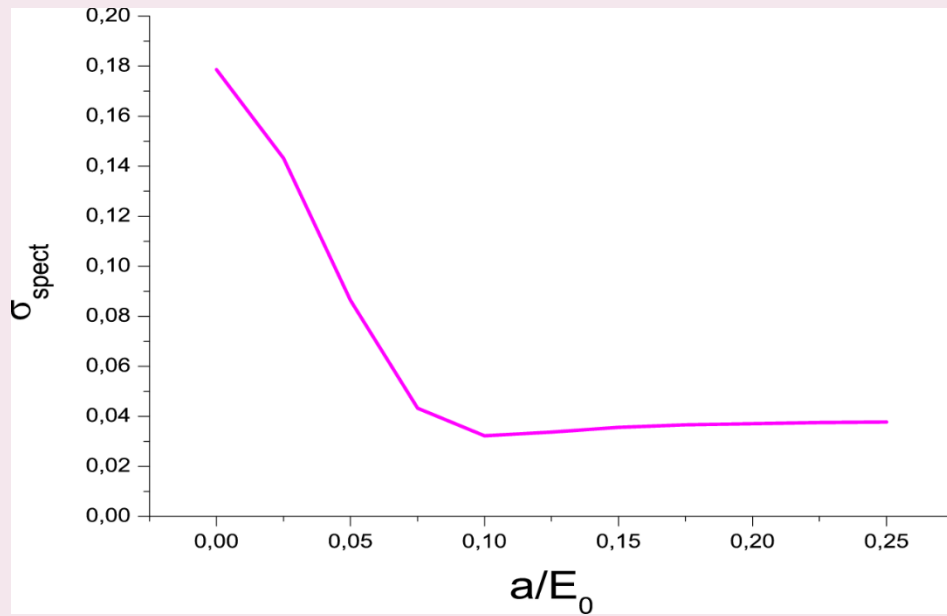
# Simulated signal – With and without adaptive triggering

$$\frac{1}{N_0} h_{ideal}(E) = \frac{1}{N_0} \frac{dN(E)}{d(E)} = \begin{cases} \frac{1}{E} & ; E < 2E_0 \\ 0 & ; E > 2E_0 \end{cases}$$



$$t_a = 4\tau$$

$$\text{THR} = 0.125 E_0$$



Instead of simultaneously adapting  $a$  and  $t_a$ , choose the values of  $a$  and  $t_a$  for the highest energy in the spectrum!

# EXPERIMENT

$^{137}\text{Cs}$  as a radiation source ( $A=88.4$  GBq).

NaI(Tl) scintillator Photonis XP2040/PC PMT

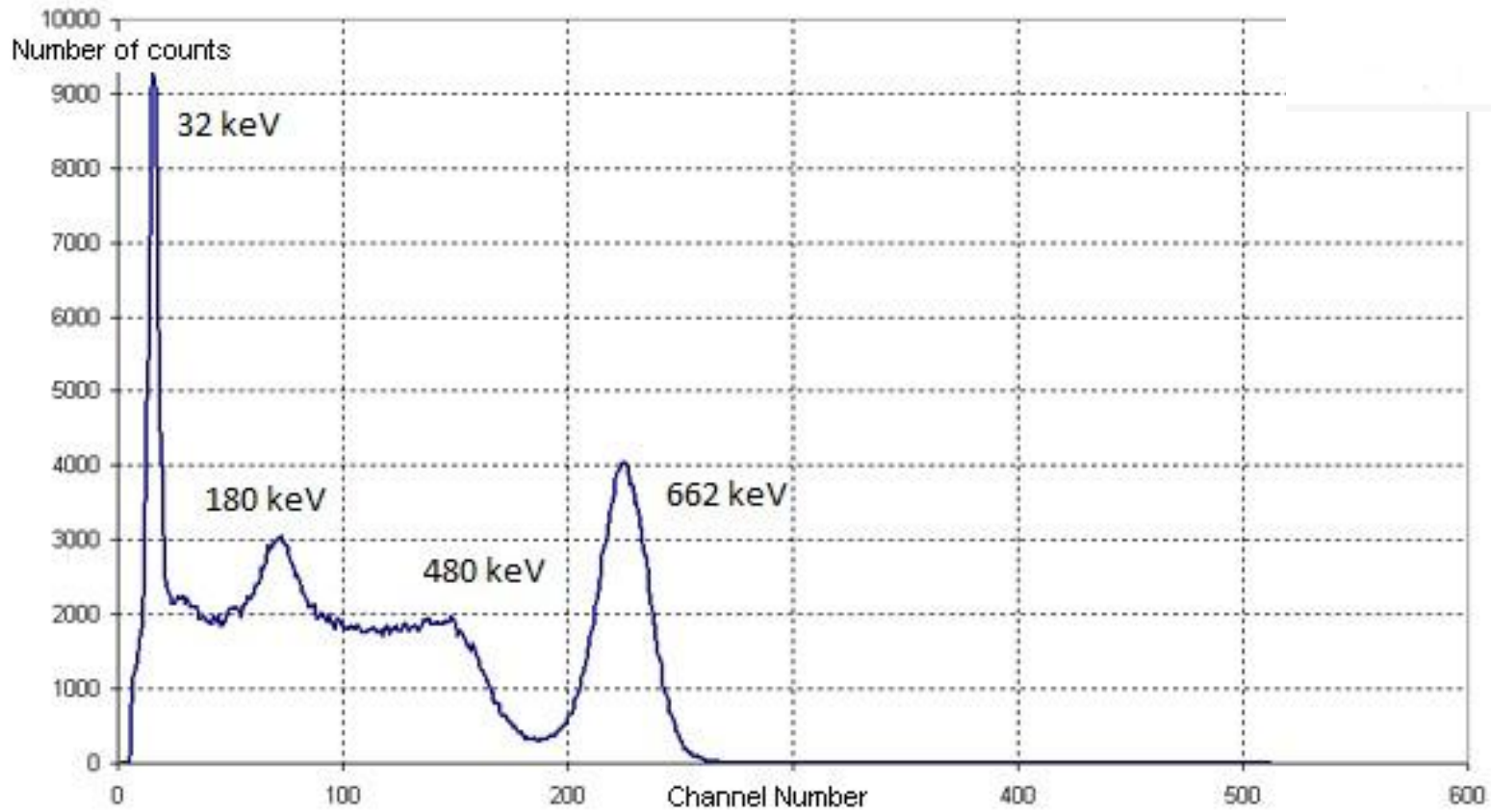
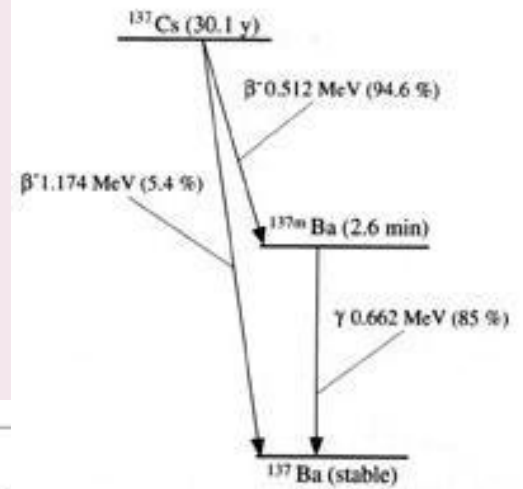
Raw anode output signal was sampled with digital PicoScope oscilloscope.

|   |                  |
|---|------------------|
| Scintillation material                  | NaI(Tl) (2"*2" ) |
| Voltage between catode and anode of PMT | 716 V            |
| Sampling period                         | 200 ps           |
| Resolution                              | 8 bits           |

Two setups:

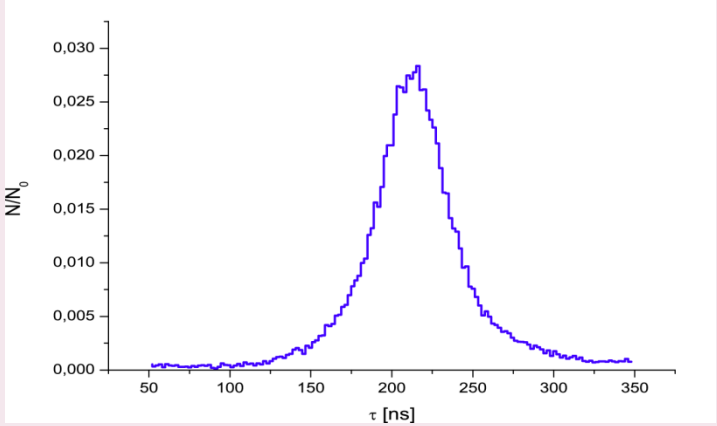
| Distance between the source and detector | Pulse frequency [Hz] |
|--|----------------------|
| $d_1=3.95$ m                             | $7.3 \cdot 10^5$     |
| $d_2=8.67$ m                             | $1.6 \cdot 10^5$     |

# $^{137}\text{Cs}$ Spectrum



# Step 1. Deduce a reference spectrum (At lower pulse frequency)

1. Deduce  $\tau$  – Fit the exponential tail of every raw pulse,  $\tau=214$  ns

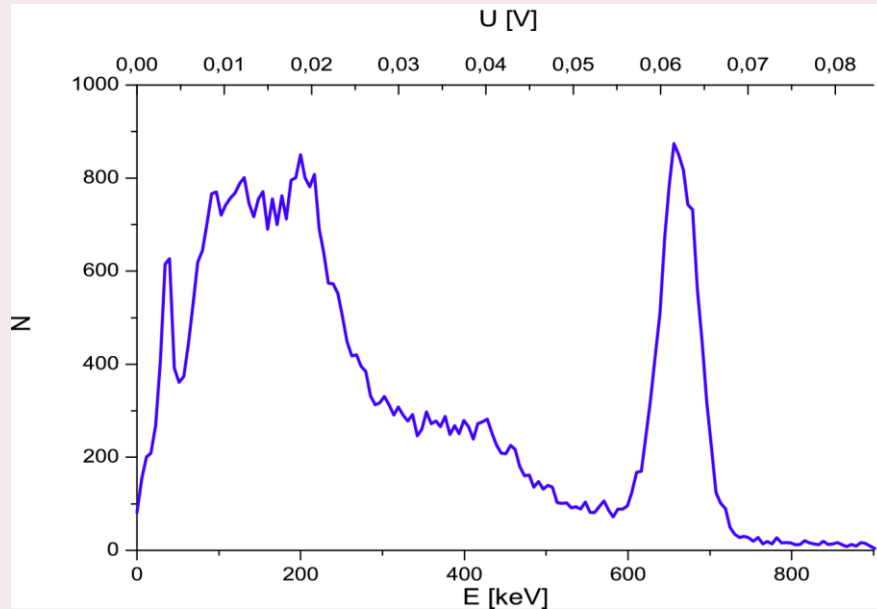


2. Deduce timing and energy

| MWD parameters   | Timestamp* | Energy** |
|------------------|------------|----------|
| Shaping time     | 100 ns     | 300 ns   |
| Flattop          | 0 ns       | 300 ns   |
| MWD pulse length | 200 ns     | 900 ns   |

- \* Only pulses that are more than  $3\tau$  separated are taken into account.
- \*\* THR=0.002 V, a=0.005 V,  $t_a=300$  ns.

## Reference spectrum :

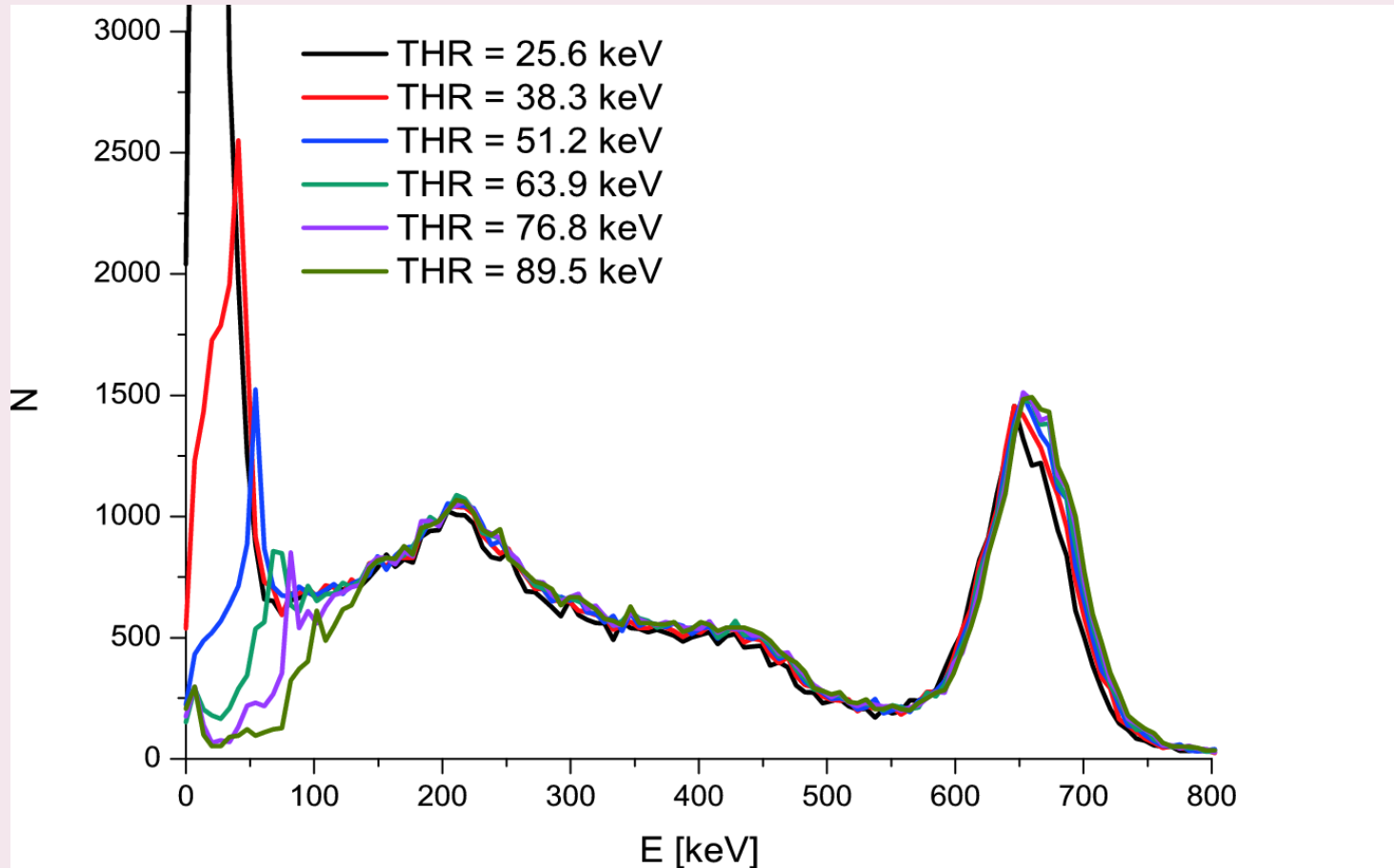


## Spectrum at higher beam frequency:

| MWD parameters   | Values |
|------------------|--------|
| Shaping time     | 50 ns  |
| Flattop time     | 50 ns  |
| MWD pulse length | 150 ns |

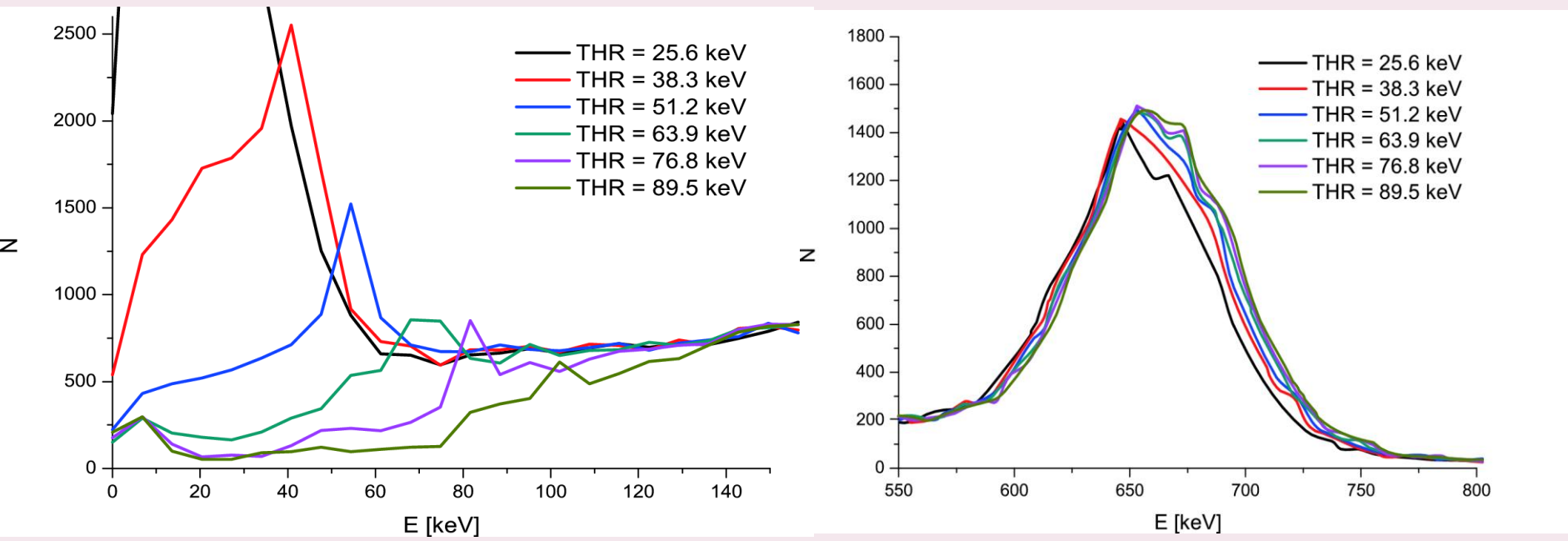


## Results – Adaptive triggering OFF

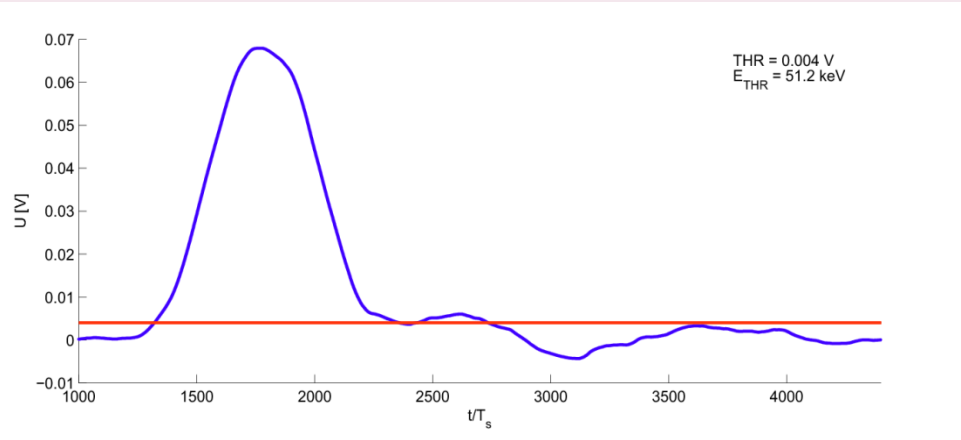


-E=32 keV is not visible.

# Adaptive triggering OFF –Enhanced spectrum

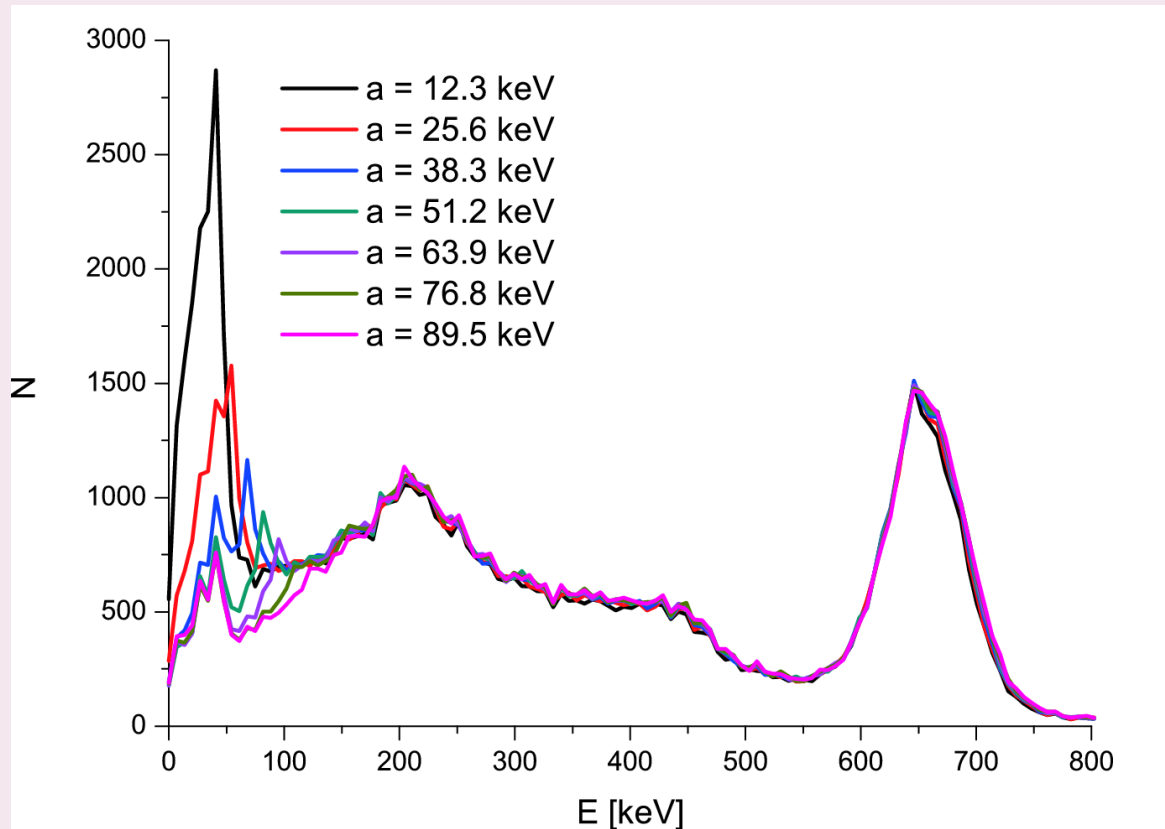


False peak with energy approximately equal to the THR energy occurs.



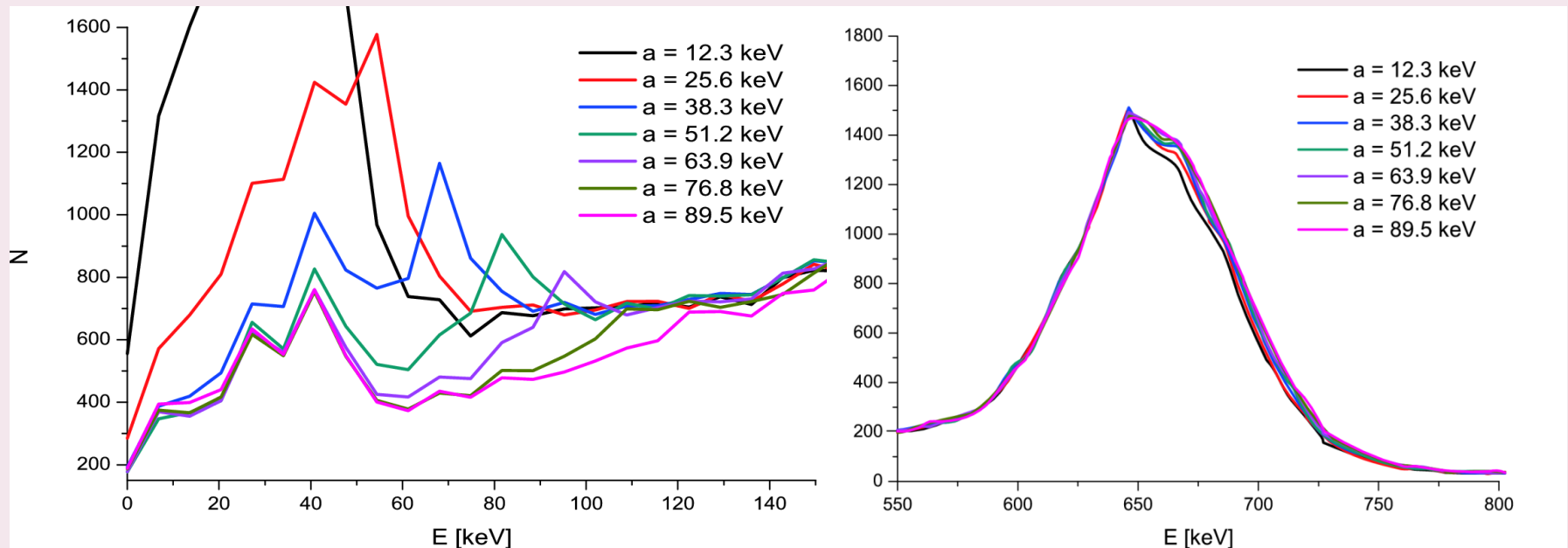
## Results – Adaptive triggering ON

$t_a=920$  ns, THR=26 keV



- Low a - Low energy peak is piled up with the afterpulses.
- High a- Low energy peak is clearly visible.

# Adaptive triggering ON – Enhanced spectrum



- False peak with energy approximately equal to the  $\text{THR}+a$  energy.

The optimal triggering parameters –  $E_{\text{max}}=662$  keV,  $a_{\text{opt}}=76$  keV,  $t_{\text{opt}}=920$  ns

(Seems ok!)

# Applications

- High throughput in-beam experiments
- QA systems in medical radiotherapies
- Active interrogation analysis / homeland security
- Industrial construction imaging

## Conclusions:

- The adaptive triggering is a relatively simple method that significantly extends the useful energy range on the low energy side of the spectrum

(when post-pulse inhibition is not applicable due to high count rates)

- This method is crucial if the desired information is presented there.

Future work:

1. More complicated synthesised signals
2. Dependence of this method on:
  - MWD parameters
  - Initial pulse frequency

# NUCLEAR STRUCTURE AND DYNAMICS III

June 14<sup>th</sup>-19<sup>th</sup>

Portorož-Portorose, Slovenia

