# Multiplicity fluctuations within nonequilibrium chiral fluid dynamics

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based on CH, Nahrgang, Yan, Kobdaj, arXiv: 1407.8277, accepted by JPG

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# Nucleons and Quark-gluon-plasma



# The QCD critical point



# Finding the CP - I

#### 1. From the QCD Lagrangian

- Solve partition function Z on a lattice (sign problem)
- Solve Dyson-Schwinger equations



<sup>(</sup>Fischer, Luecker, Phys. Lett. B 718 (2013) 1036-1043)

## Finding the CP - II

#### 2. From effective models

- Respect chiral symmetry (Sigma model, NJL model, ...)
- Existence/location of CP not universal!



## Finding the CP - III

#### 3. From experiment

Fluctuations sensitive to critical region



(STAR collaboration, Phys. Rev. Lett. 112 (2014) 032302)



(NA49 collaboration, Nucl. Phys. A 830 (2009))

$$\sigma^{2} = \langle \delta N^{2} \rangle \sim \xi^{2}$$
$$S\sigma = \frac{\langle \delta N^{3} \rangle}{\langle \delta N^{2} \rangle} \sim \xi^{2.5}$$
$$s\sigma^{2} = \frac{\langle \delta N^{4} \rangle}{\langle \delta N^{2} \rangle} - 3 \langle \delta N^{2} \rangle \sim \xi^{5}$$

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(Stephanov, Phys. Rev. Lett. 102 (2009))

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#### CP and first-order phase transition



- Net quark number susceptibility at CP (solid) and first-order phase transition (dashed) from NJL model
- Change of universality class,  $\gamma = 2/3$  for CP,  $\gamma = 1/2$  for first-order

$$\chi_{\mu\mu} \sim (\mu - \mu_0)^{-\gamma}$$

(Sasaki, Friman, Redlcih, Phys. Rev. D 77 (2008))

#### Latent heat in heavy-ion collisions?



 Latent heat might influence directed flow v<sub>1</sub>, strength of expansion Measured by STAR collaboration

(STAR collaboration, Phys. Rev. Lett. 112 (2014) 162301)

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Data

Start from effective chiral model with CEP and first-order phase transition 1. Study thermodynamics:

- Calculate susceptibility and compare with NJL
- Calculate kurtosis
- Determine critical indices
- 2. Study fluid dynamics (heavy-ion collisions):
  - Quark and chiral fields  $\rightarrow$  quark fluid and explicitly propagated fields
  - Ensemble fluctuations event-by-event

#### A chiral model with dilatons

Potential and equation of state from

$$\mathcal{L} = \overline{q} \left[ i \left( \gamma^{\mu} \partial_{\mu} - i g_{s} \gamma^{0} A_{0} \right) - g \sigma \right] q + \frac{1}{2} \left( \partial_{\mu} \sigma \right)^{2} + \frac{1}{2} \left( \partial_{\mu} \chi \right)^{2} - U(\sigma) - \mathcal{U}(\chi)$$

(Sasaki, Mishustin, Phys. Rev. C 85 (2012) 025202)



#### Spinodal instabilities at T = 40 MeV



- Phase transition in the presence of spinodal instabilities
- Mechanically instable region in the equation of state

#### Spinodal instabilities at T = 40 MeV



•  $\frac{\chi_{q}}{T^{2}} = \frac{1}{VT^{3}} \langle \delta N_{q}^{2} \rangle$  proportional to quark number fluctuations •  $\kappa = \frac{\langle \delta N_{q}^{4} \rangle}{\langle \delta N_{q}^{2} \rangle} - 3 \langle \delta N_{q}^{2} \rangle$ 

Expect: Enhancement of fluctuations at CP AND 1st order transition

### **Critical indices**



$$\chi_{
m q}\sim(\mu-\mu_0)^{-\gamma}$$
 with  $\gamma=1/2$  for first-order and  $\gamma=2/3$  for CEP

$$\kappa \sim (\mu - \mu_0)^{-\zeta}$$

 $\lambda - \zeta$ 

with  $\zeta = 2$  for first-order and  $\zeta = 2$  for CEP

### A chiral model with dilatons ... dynamically

Ingredients for fully dynamical model:

- Hot medium (quarks)
- Fluctuations (chiral fields) Chiral fluid dynamics

$$-\frac{\delta S_{\rm cl}}{\delta \sigma} - D = \xi , \ \partial_{\mu} T_{\rm q}^{\mu\nu} = S_{\sigma}^{\nu}$$

(Nahrgang, Leupold, Herold, Bleicher, Phys. Rev. C 84 (2011))

#### Potential and equation of state from

$$\mathcal{L} = \overline{q} \left[ i \left( \gamma^{\mu} \partial_{\mu} - i g_{s} \gamma^{0} A_{0} \right) - g \sigma \right] q + \frac{1}{2} \left( \partial_{\mu} \sigma \right)^{2} + \frac{1}{2} \left( \partial_{\mu} \chi \right)^{2} - U(\sigma) - \mathcal{U}(\chi)$$

(Sasaki, Mishustin, Phys. Rev. C 85 (2012) 025202)





#### Event-by-event fluctuations: variance



Fixed volume vs. rapidity (y < 0.5) and  $p_T$  cut (100 MeV/fm<sup>3</sup>  $< p_T < 500$  MeV/fm<sup>3</sup>)

- Different scales due to different volumes
- Fluctuating volume for rapidity and momentum cut

#### Event-by-event fluctuations: kurtosis



Fixed volume vs. rapidity (y < 0.5) and  $p_T$  cut (100 MeV/fm<sup>3</sup>  $< p_T < 500$  MeV/fm<sup>3</sup>)

- Different scales due baryon number conservation
- Ratios of cumulants depend on fraction of measured to total baryons

#### From fluctuations to observables ...

- ... some more things need to be considered in the future
  - Freeze out over hypersurface with constant energy density or temperature
  - Final state interactions
  - Evolution of fluctuations in the hadronic phase

## SU(3) chiral quark-hadron (QH) model

include

- 3 quarks (*u*, *d*, *s*), baryon octet
- scalar mesons  $\sigma$ ,  $\zeta$ , vector meson  $\omega$

$$\mathcal{L} = \sum_{i} \overline{\psi}_{i} \left( i \gamma^{\mu} \partial_{\mu} - \gamma^{0} g_{i\omega} \omega - M_{i} \right) \psi_{i} + \frac{1}{2} \left( \partial_{\mu} \sigma \right)^{2} - U(\sigma, \zeta, \omega) - \mathcal{U}(\ell)$$

with effective masses generated by  $\sigma$  and  $\ell$ 

$$\begin{split} M_q &= g_{q\sigma}\sigma + g_{q\zeta}\zeta + M_{0q} + g_{q\ell}(1-\ell) \\ M_B &= g_{B\sigma}\sigma + g_{B\zeta}\zeta + M_{0B} + g_{B\ell}\ell^2 \end{split}$$

(Dexheimer, Schramm, Phys. Rev. C 81 (2010), 045201)

Improves equation of state, plus:

- Chiral phase transition at larger chemical potentials
- Disentangled from liquid-gas phase transition

## QH model - phase diagram



- Chiral phase transition at larger chemical potentials
- Disentangled from liquid-gas phase transition

(Herold, Limphirat, Kobdaj, Yan, SPC 2014)

### Summary and Conclusions

- Event-by-event fluctuations become enhanced in hydrodynamic phase for CEP and first-order phase transition
- Effects of hadronic phase have to be taken into account for reliable predictions

# THANK YOU