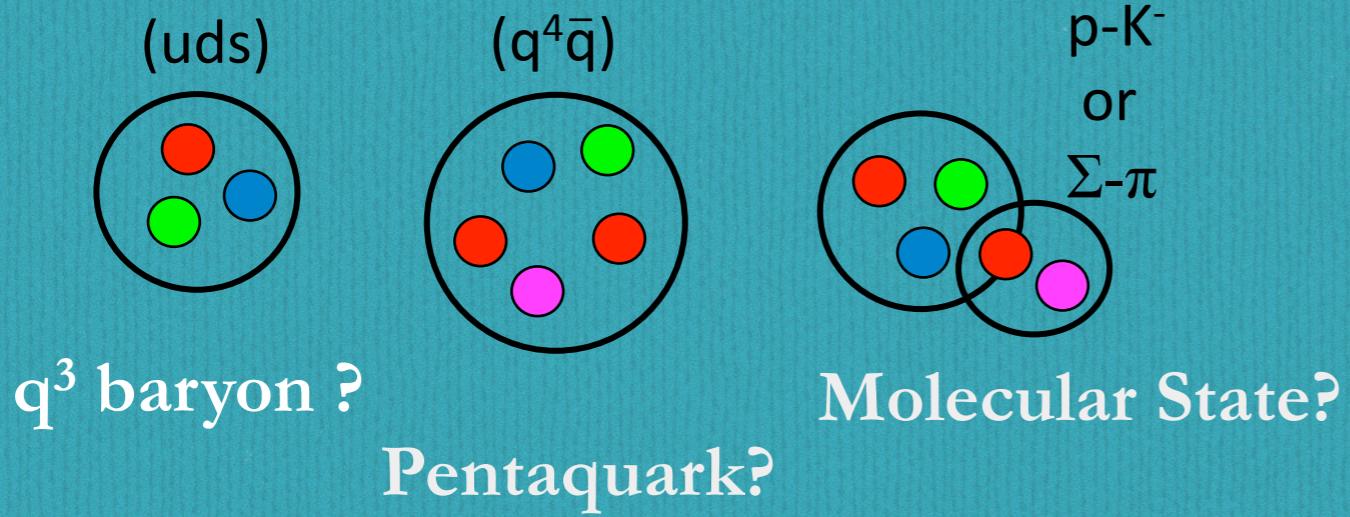


$\Lambda(1405)$ Observations in p+p and K⁻-induced reactions

L. Fabbietti

Technische Universitaet Muenchen
Excellence Cluster “Origin and Structure of the Universe”
<https://www.e12.ph.tum.de/groups/kcluster>

$\Lambda(1405)$



$M \approx 1406 \text{ MeV}/c^2$

$\Gamma \approx 50 \text{ MeV}/c^2$

$I(J^P)=0 \ (1/2^+)$

Decay Channels:

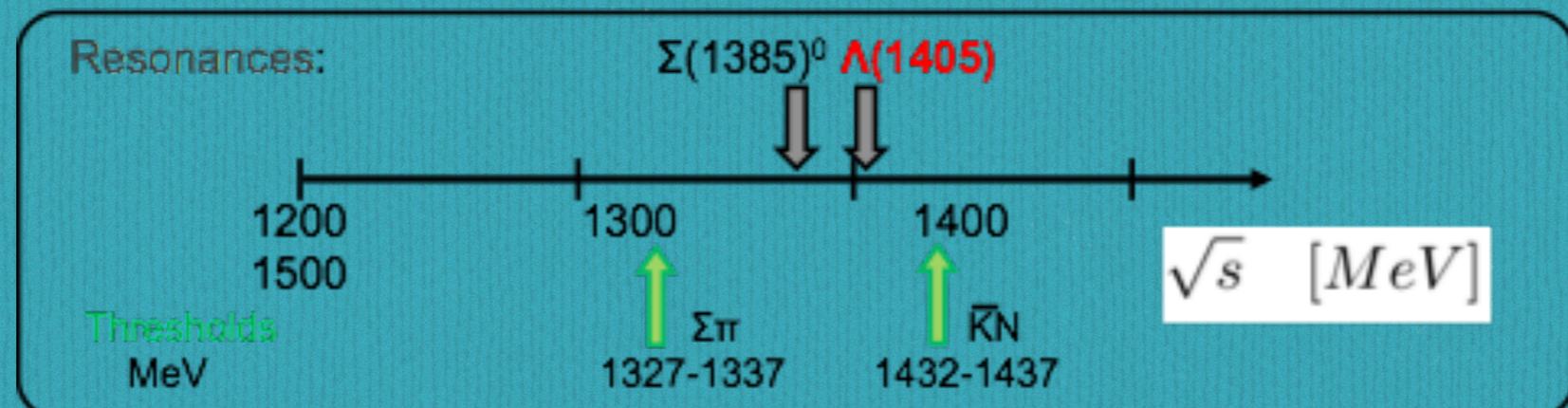
$\Sigma^0\pi^0, \Sigma^+\pi^-, \Sigma^-\pi^+$

q^3 :

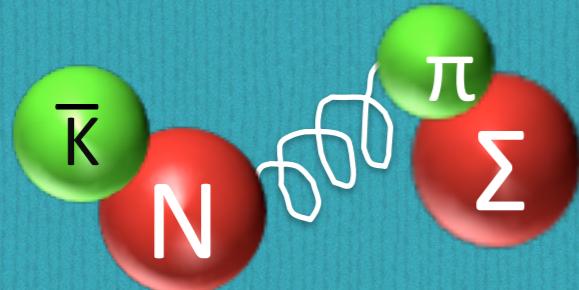
Expected Mass $\sim 1700 \text{ MeV}/c^2$

Pentaquark:

more observed excited baryons



Double Pole Structure



From calculation constrained by scattering data and kaonic atoms:

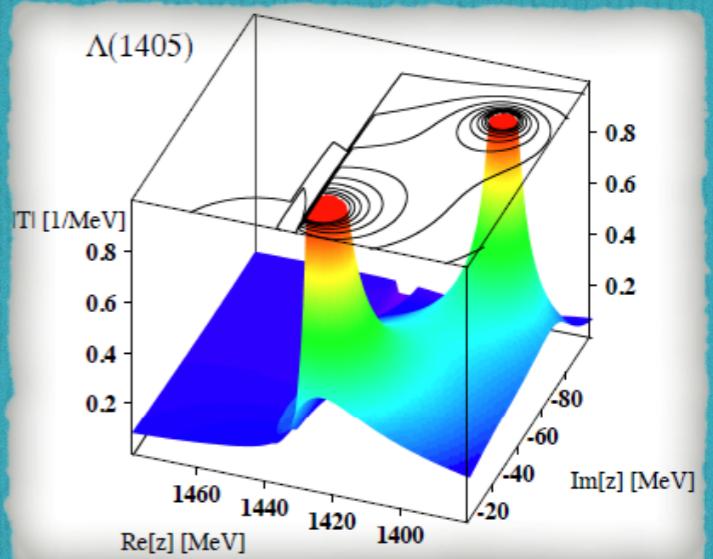
Nucl.Phys. A881 (2012) 98-114

$$\begin{aligned} m_{0,1} = \text{Re}(z_1) &= 1424^{+7}_{-23} \text{ MeV} \\ \Gamma_{0,1} = 2*\text{Im}(z_1) &= 52^{+6}_{-28} \text{ MeV} \\ m_{0,2} = \text{Re}(z_2) &= 1381^{+18}_{-6} \text{ MeV} \\ \Gamma_{0,2} = 2*\text{Im}(z_2) &= 162^{+38}_{-16} \text{ MeV} \end{aligned}$$

Nucl.Phys. A900 (2013) 51 - 64

$$\begin{aligned} m_{0,1} = \text{Re}(z_1) &= 1428^{+2}_{-1} \text{ MeV} \\ \Gamma_{0,1} = 2*\text{Im}(z_1) &= 16^{+4}_{-4} \text{ MeV} \\ m_{0,2} = \text{Re}(z_2) &= 1497^{+11}_{-7} \text{ MeV} \\ \Gamma_{0,2} = 2*\text{Im}(z_2) &= 150^{+18}_{-18} \text{ MeV} \end{aligned}$$

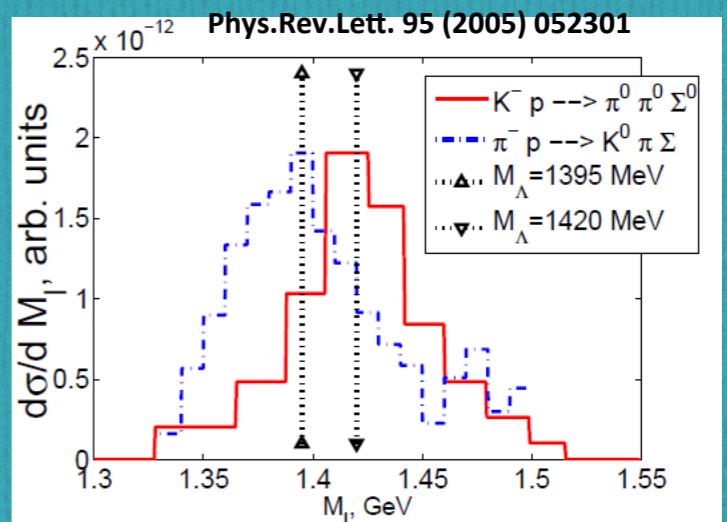
Prog.Part.Nucl.Phys. 67 (2012) 55-98



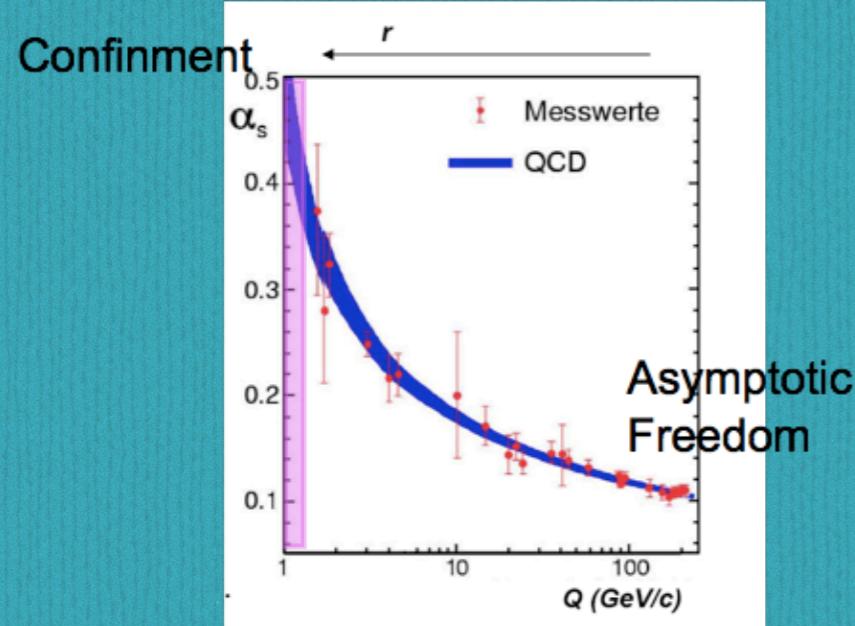
Experimental $\Lambda(1405)$ in $\Sigma\pi$ spectrum:

- * $K^- + p$ reactions excite mainly **1420 pole**
- * $\pi^- + p$ reactions excite mainly **1390 pole**

→ $\Lambda(1405)$ correlated with KN dynamics!



Effective Hadron-Hadron Interaction

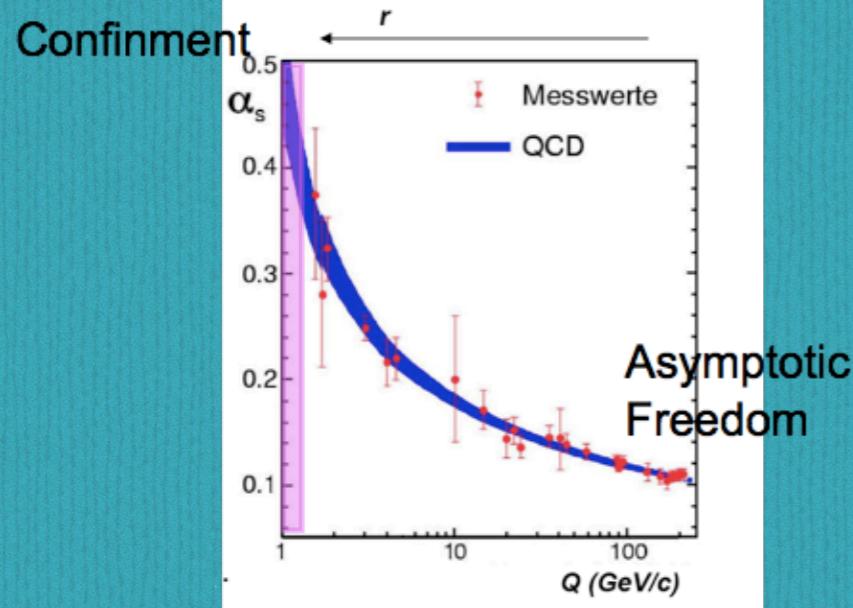


Small Q (~ 1 GeV)
 Large Distances (1 fm)

Effective Field Theory of interacting Bosons

- Test-bed of the strong interaction in few body systems
- Strange quarks are intermediate between “light” and “heavy”
- -> Interplay between spontaneous and explicit chiral symmetry breaking in low energy QCD.
- Testing ground: K-N and \bar{K} -N interactions

Effective Hadron-Hadron Interaction



Small Q (~ 1 GeV)
Large Distances (1 fm)

Effective Field Theory of interacting Bosons

- Test-bed of the strong interaction in few body systems
- Strange quarks are intermediate between “light” and “heavy”
-> Interplay between spontaneous and explicit chiral symmetry breaking in low energy QCD.
- Testing ground: K-N and \bar{K} -N interactions

Effective Hadron-Hadron Interaction

Antikaon

ChPT in SU(3) (exact Theory) does not work for the KN system

-> **Coupled-Channel Ansatz** based on Chiral Dynamics

Kaon

Effective Hadron-Hadron Interaction

Antikaon

ChPT in SU(3) (exact Theory) does not work for the KN system

-> Coupled-Channel Ansatz based on Chiral Dynamics

Ex: $\Lambda(1405)$ dynamically generated like a QUASI-BOUND K-p (I=0) State

R.H. Dalitz et al., Phys. Rev. 153 (1967) 1617

Kaon

Effective Hadron-Hadron Interaction

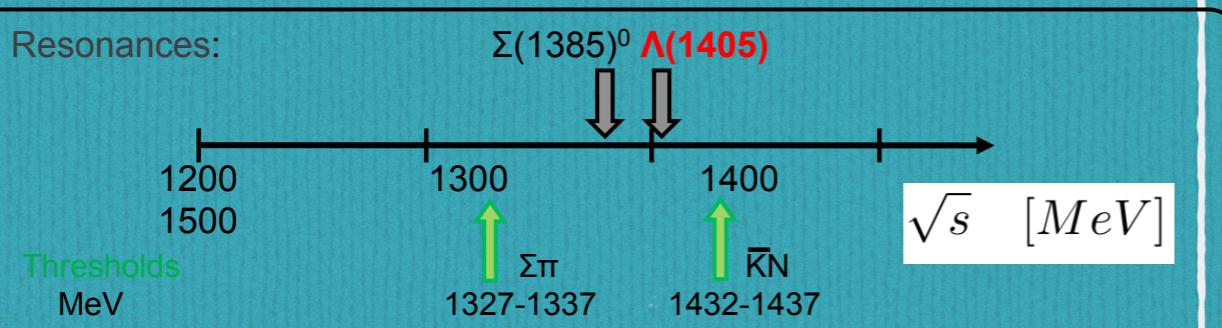
Antikaon

ChPT in SU(3) (exact Theory) does not work for the KN system

-> **Coupled-Channel Ansatz** based on Chiral Dynamics

Ex: $\Lambda(1405)$ dynamically generated like a QUASI-BOUND K-p (I=0) State

R.H. Dalitz et al., Phys. Rev. 153 (1967) 1617



Kaon

Effective Hadron-Hadron Interaction

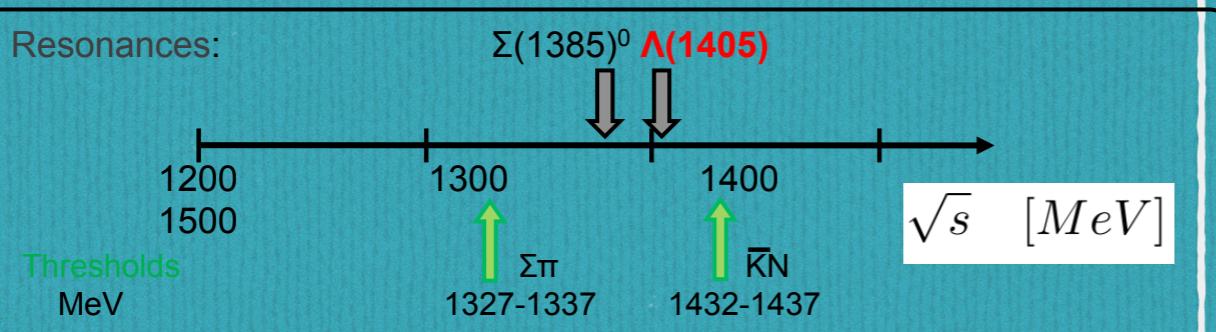
Antikaon

ChPT in SU(3) (exact Theory) does not work for the KN system

-> Coupled-Channel Ansatz based on Chiral Dynamics

Ex: $\Lambda(1405)$ dynamically generated like a QUASI-BOUND K-p (I=0) State

R.H. Dalitz et al., Phys. Rev. 153 (1967) 1617



Kaon



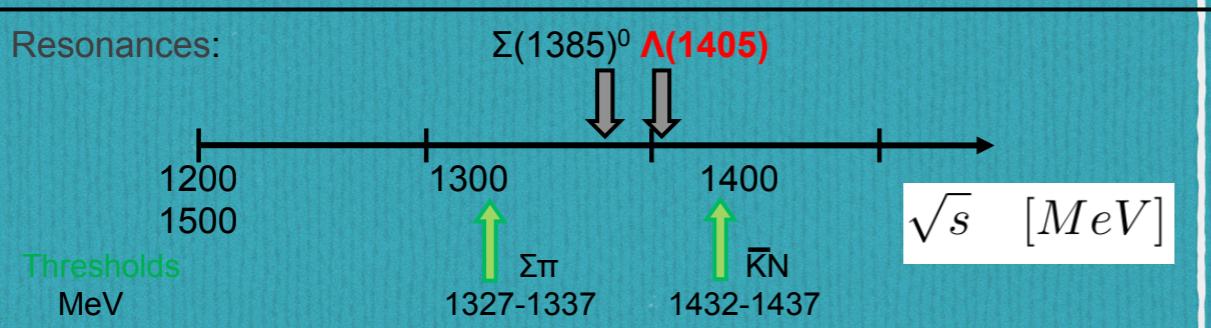
Effective Hadron-Hadron Interaction

Antikaon

ChPT in SU(3) (exact Theory) does not work for the KN system
 -> **Coupled-Channel Ansatz** based on Chiral Dynamics

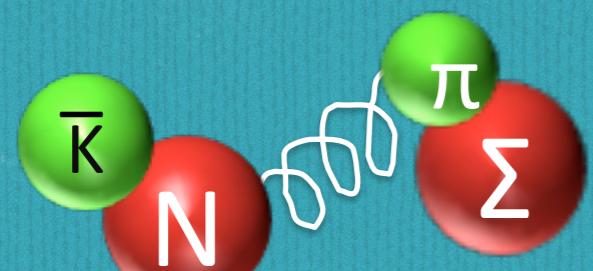
Ex: $\Lambda(1405)$ dynamically generated like a QUASI-BOUND K-p (I=0) State

R.H. Dalitz et al., Phys. Rev. 153 (1967) 1617



Kaon

ChPT is used to describe K-N interaction since Kaons do not get absorbed generating resonances



p+p in the GeV Energy Range

Fixed Target experiments, $E_{\text{kin}} \sim A \text{GeV}$

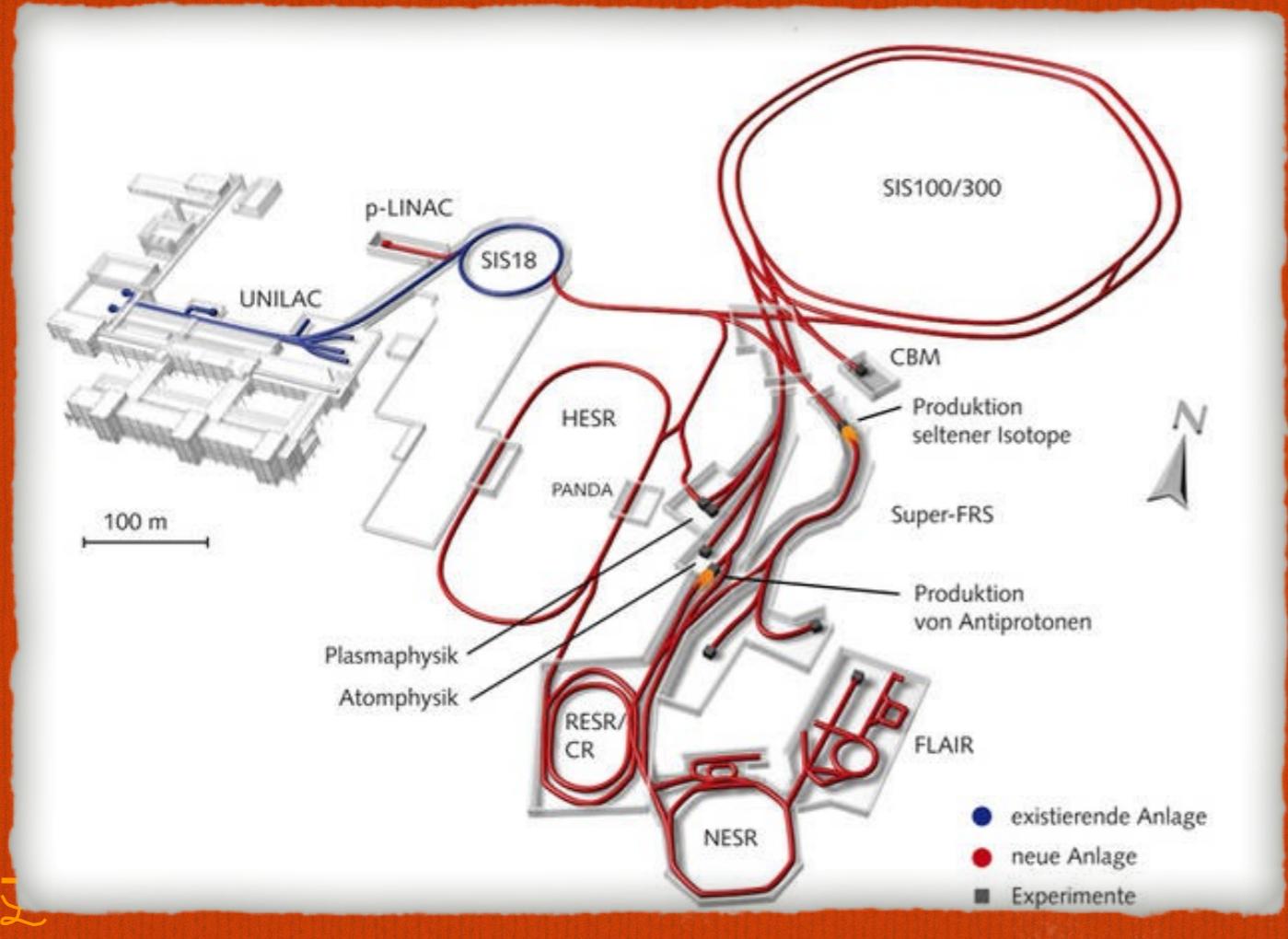
proton-proton



p+p in the GeV Energy Range

Fixed Target experiments, $E_{\text{kin}} \sim A \text{GeV}$

proton-proton



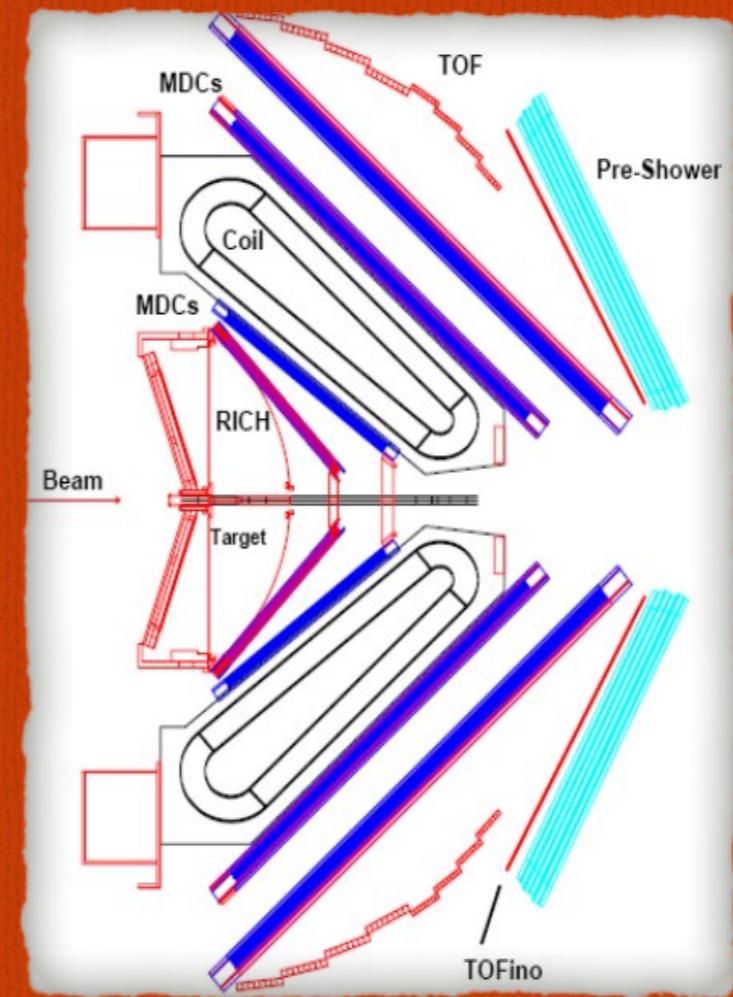
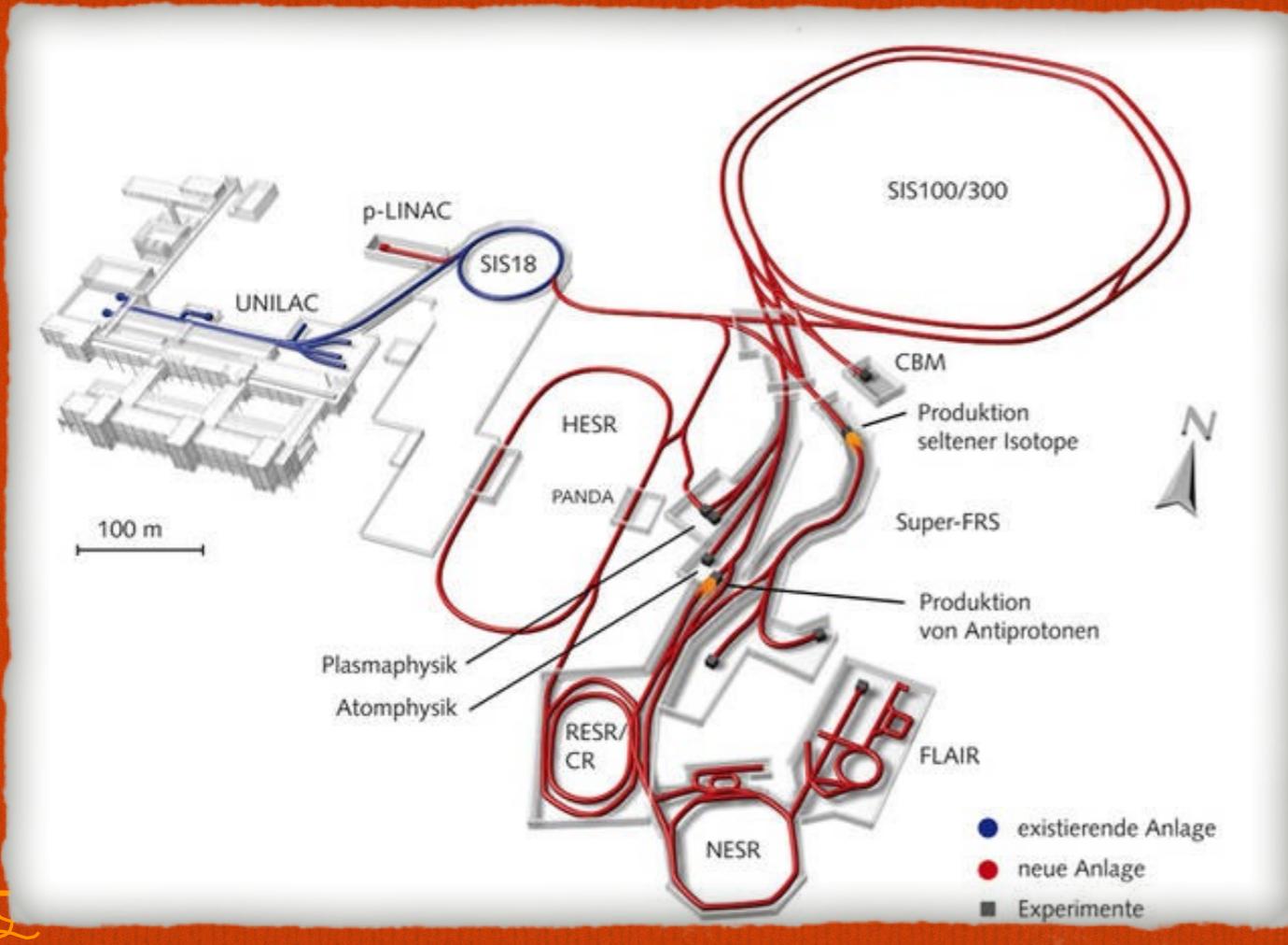
p+p in the GeV Energy Range

Fixed Target experiments, $E_{\text{kin}} \sim A \text{GeV}$

proton-proton



HADES
High Acceptance Di-Electron Spectrometer
Fixed Target Experiment
SIS18, $E_{\text{kin}} = 1-3 \text{ GeV/nucl}$
Full azimuthal coverage, $18^\circ - 85^\circ$ in polar angle
 $\delta p/p \sim 1-3 \%$

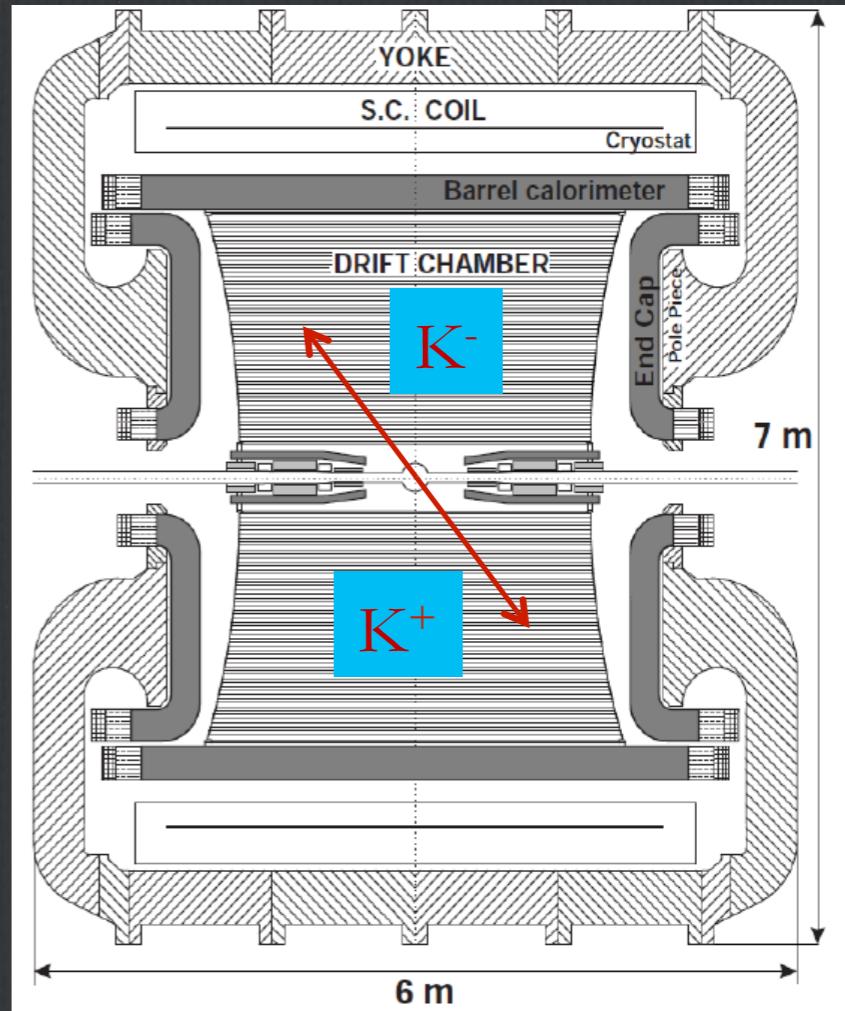


K- ${}^4\text{He}$ reactions at Daφne



- K^- Momentum = 127 MeV/c
- $\sigma p/p \sim 0.4 \text{ MeV}/c$
- 96% geometrical acceptance
- Calorimeter for γs : $\sigma_m \sim 18 \text{ MeV}/c^2$
- Vertex resolution: 1 mm
- Gas: 90% He, 10% C_4H_{10}

KLOE Experiment

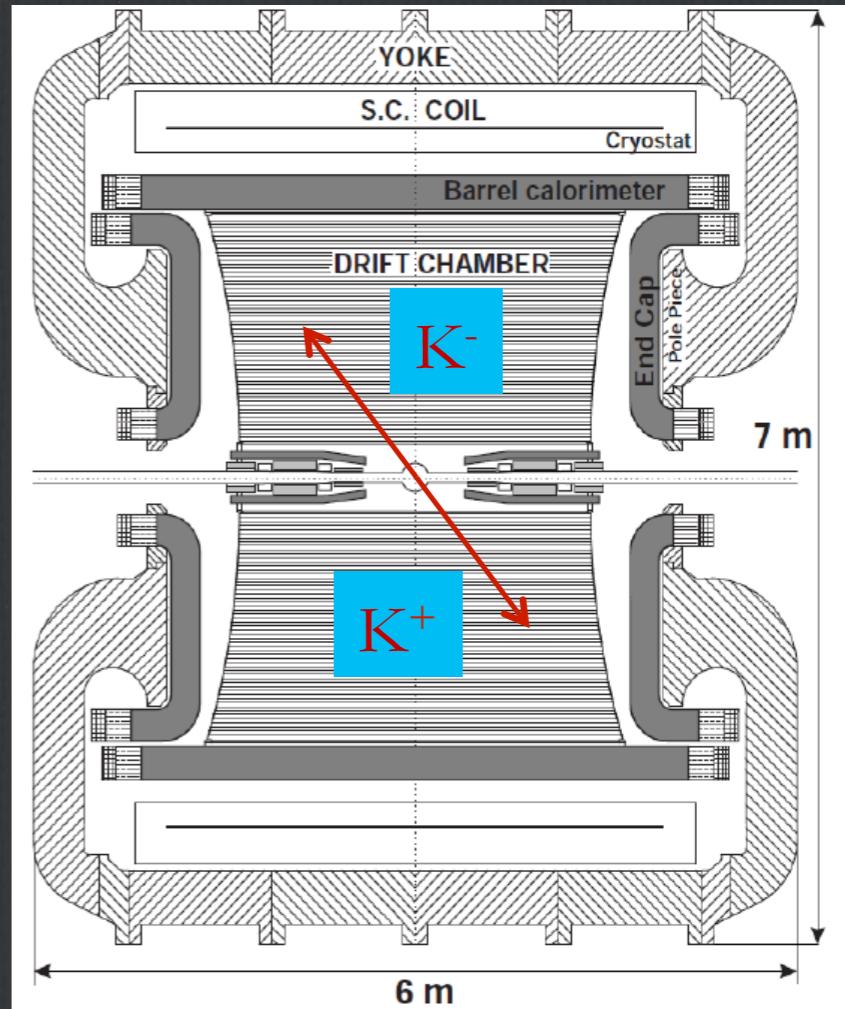


K- ${}^4\text{He}$ reactions at Daφne

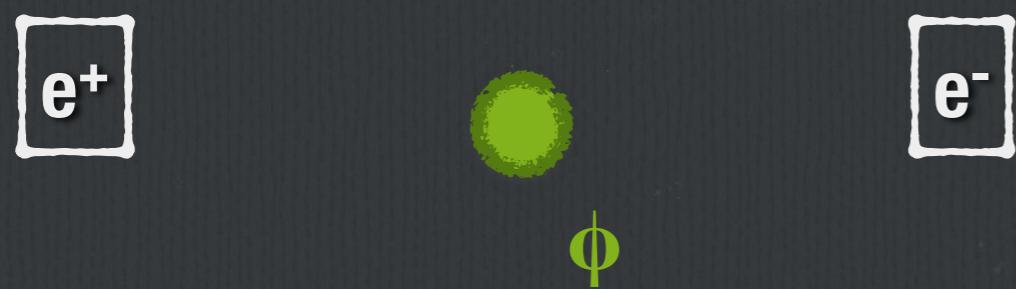


- K- Momentum = 127 MeV/c
- $\sigma p/p \sim 0.4 \text{ MeV}/c$
- 96% geometrical acceptance
- Calorimeter for γs : $\sigma_m \sim 18 \text{ MeV}/c^2$
- Vertex resolution: 1 mm
- Gas: 90% He, 10% C_4H_{10}

KLOE Experiment

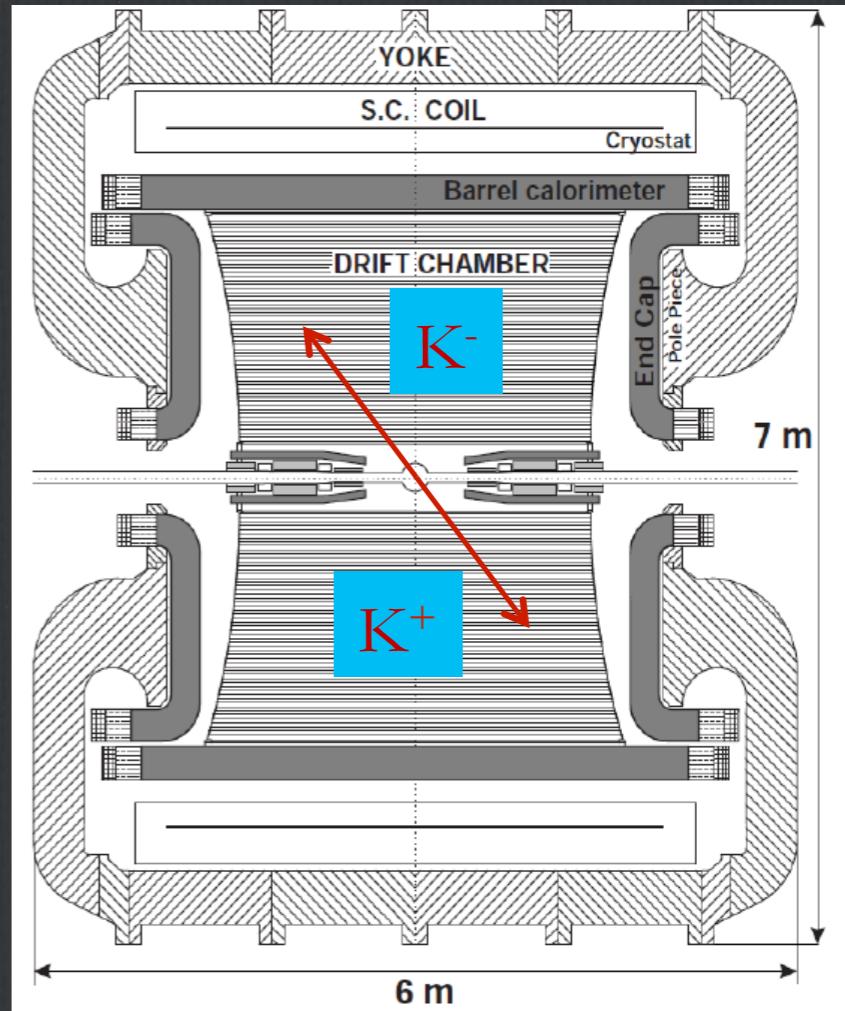


K- ${}^4\text{He}$ reactions at Daφne

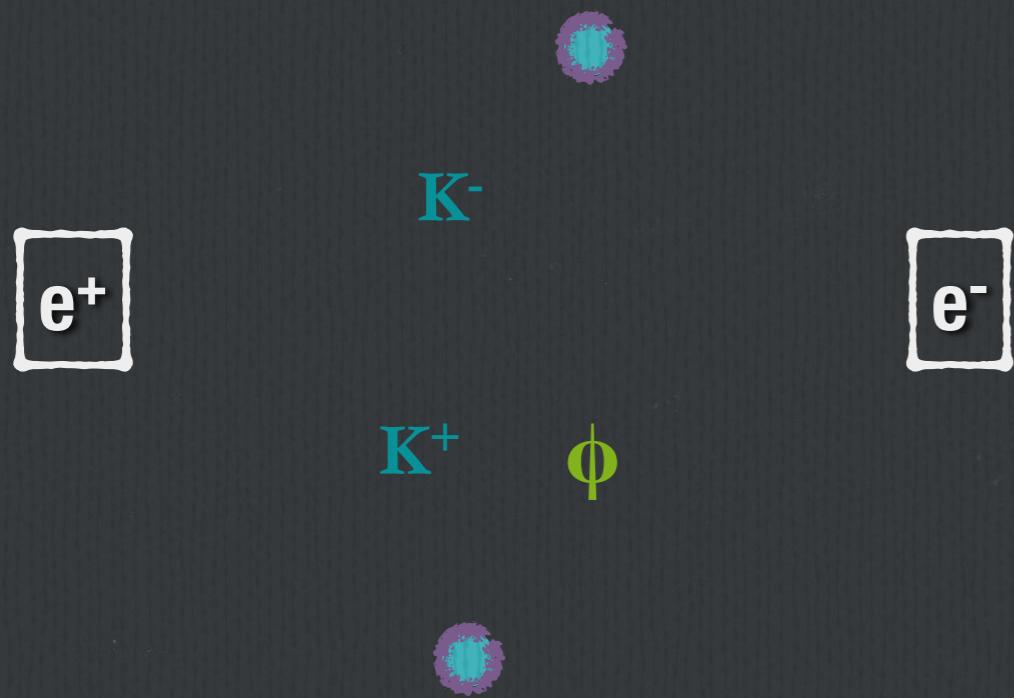


- K- Momentum = 127 MeV/c
- $\sigma p/p \sim 0.4 \text{ MeV}/c$
- 96% geometrical acceptance
- Calorimeter for γs : $\sigma_m \sim 18 \text{ MeV}/c^2$
- Vertex resolution: 1 mm
- Gas: 90% He, 10% C_4H_{10}

KLOE Experiment

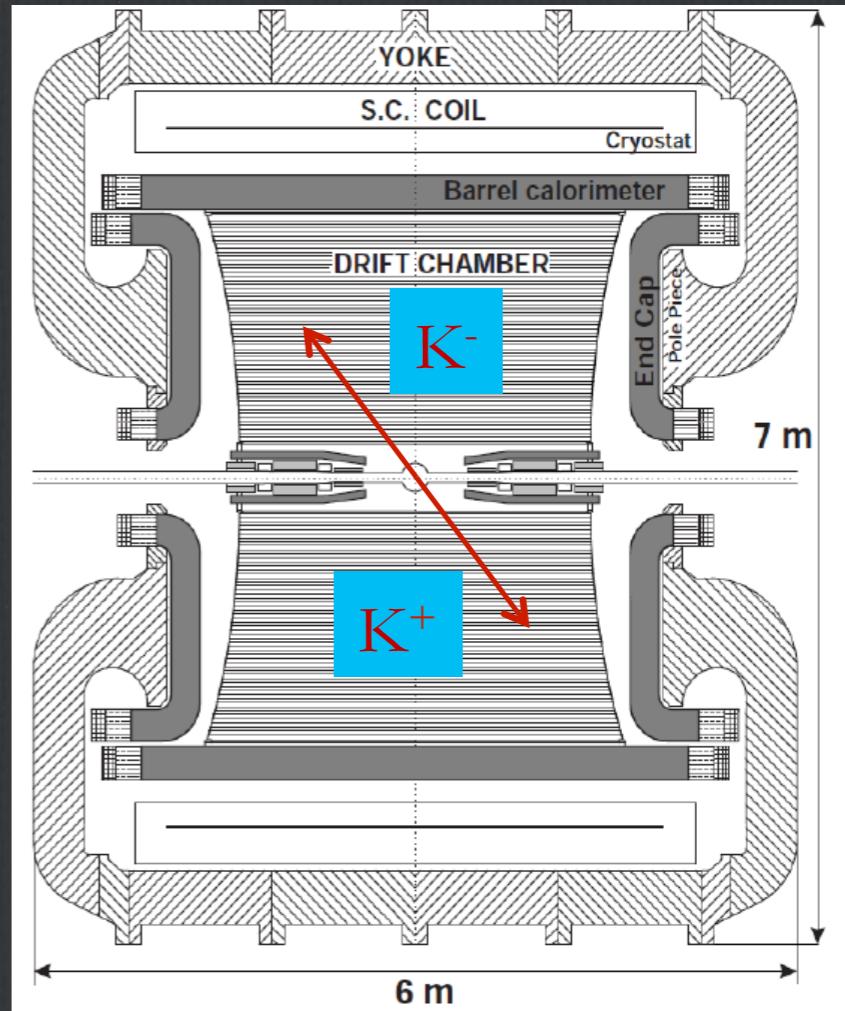


K- ${}^4\text{He}$ reactions at Daφne



- K- Momentum = 127 MeV/c
- $\sigma p/p \sim 0.4 \text{ MeV}/c$
- 96% geometrical acceptance
- Calorimeter for γs : $\sigma_m \sim 18 \text{ MeV}/c^2$
- Vertex resolution: 1 mm
- Gas: 90% He, 10% C_4H_{10}

KLOE Experiment



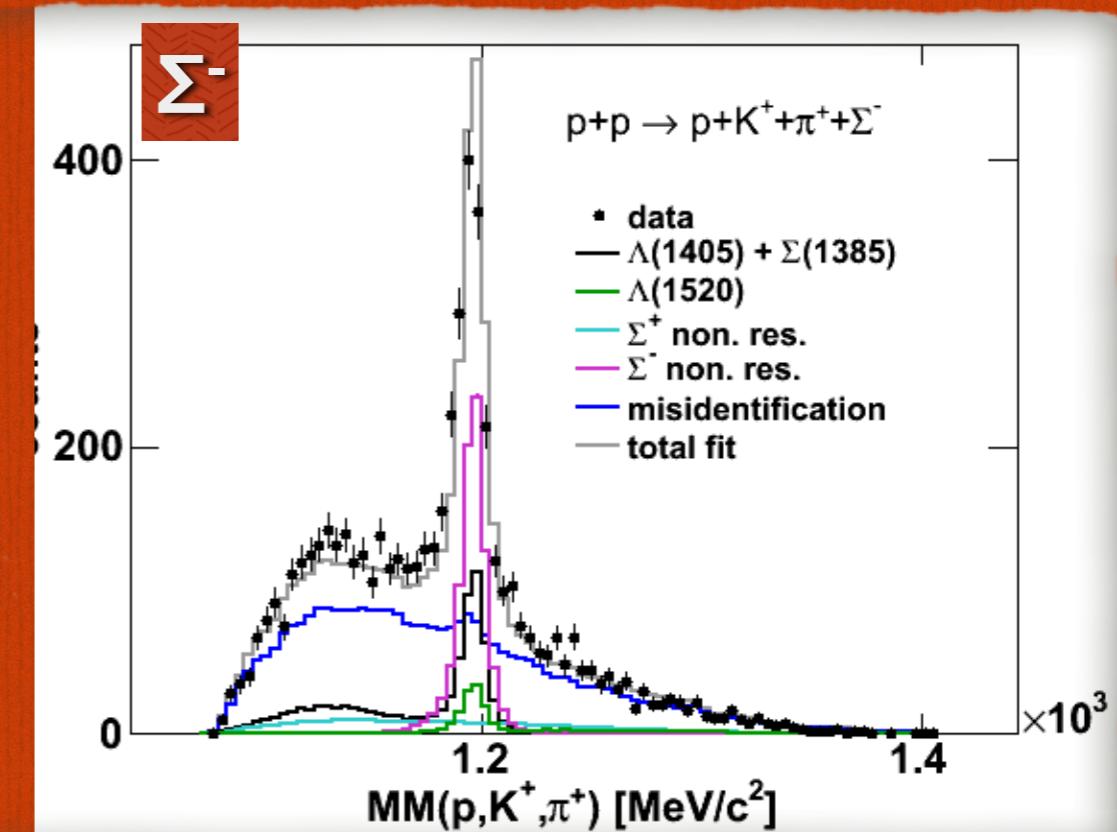
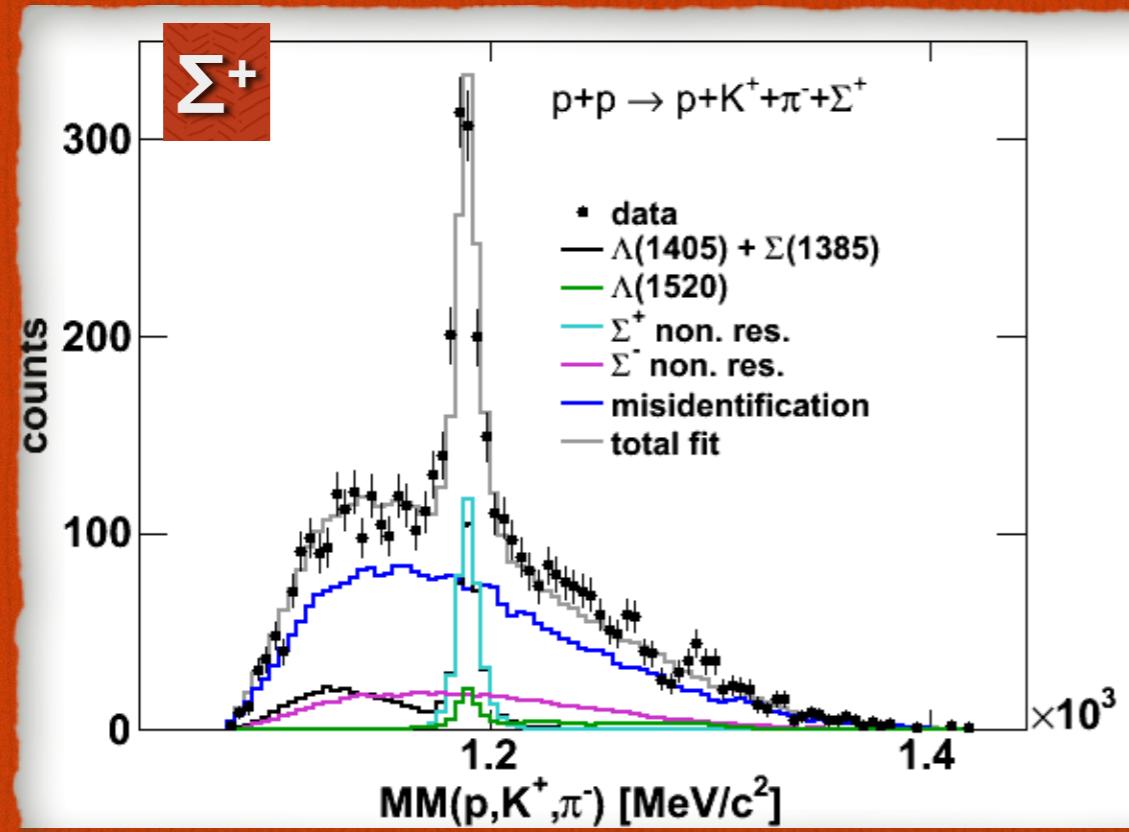
p+p collisions

proton-proton



$\sigma(\Lambda(1405))/\sigma(\Sigma(1385)^0) \approx 2.8$, known from analysis of neutral decay channels

Conservative estimation: $\sigma(\Sigma(1385)^0) = 0.5\sigma(\Sigma(1385)^+)$
 $\rightarrow \sigma(\Lambda(1405))/\sigma(\Sigma(1385)^0) \approx 1.0 \rightarrow$ no influence on pole mass or line shape of $\Lambda(1405)$



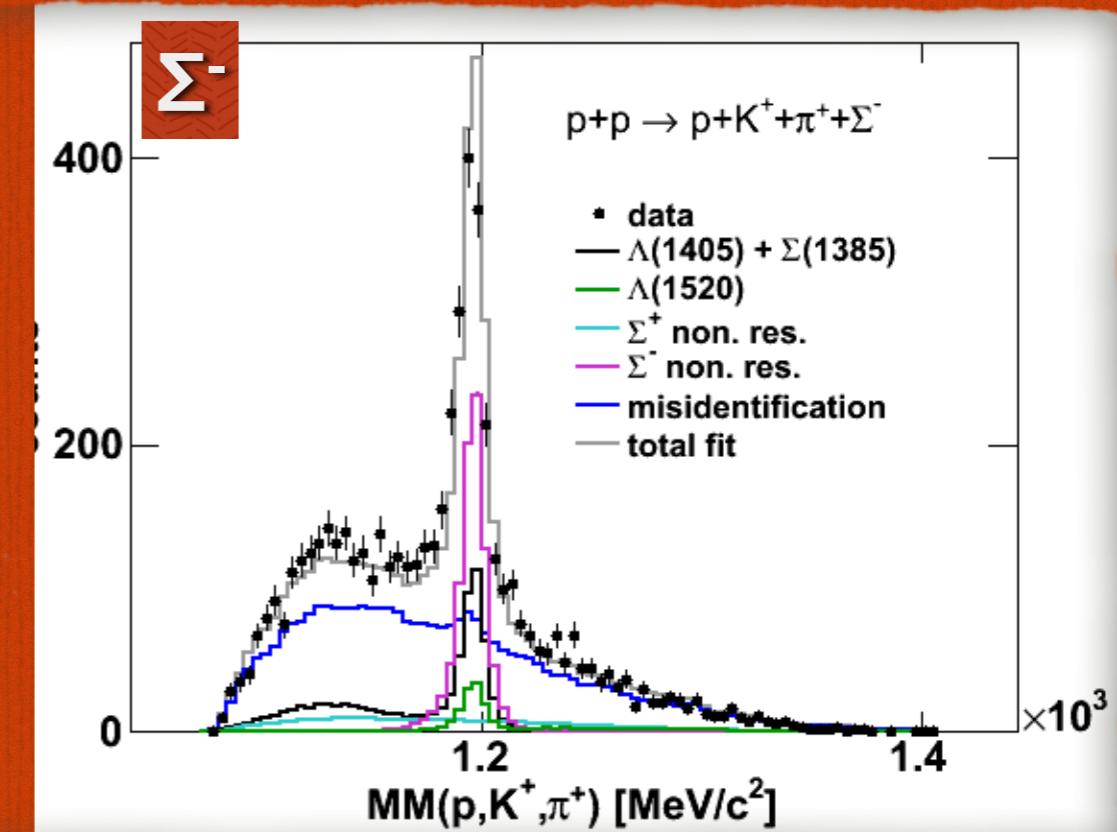
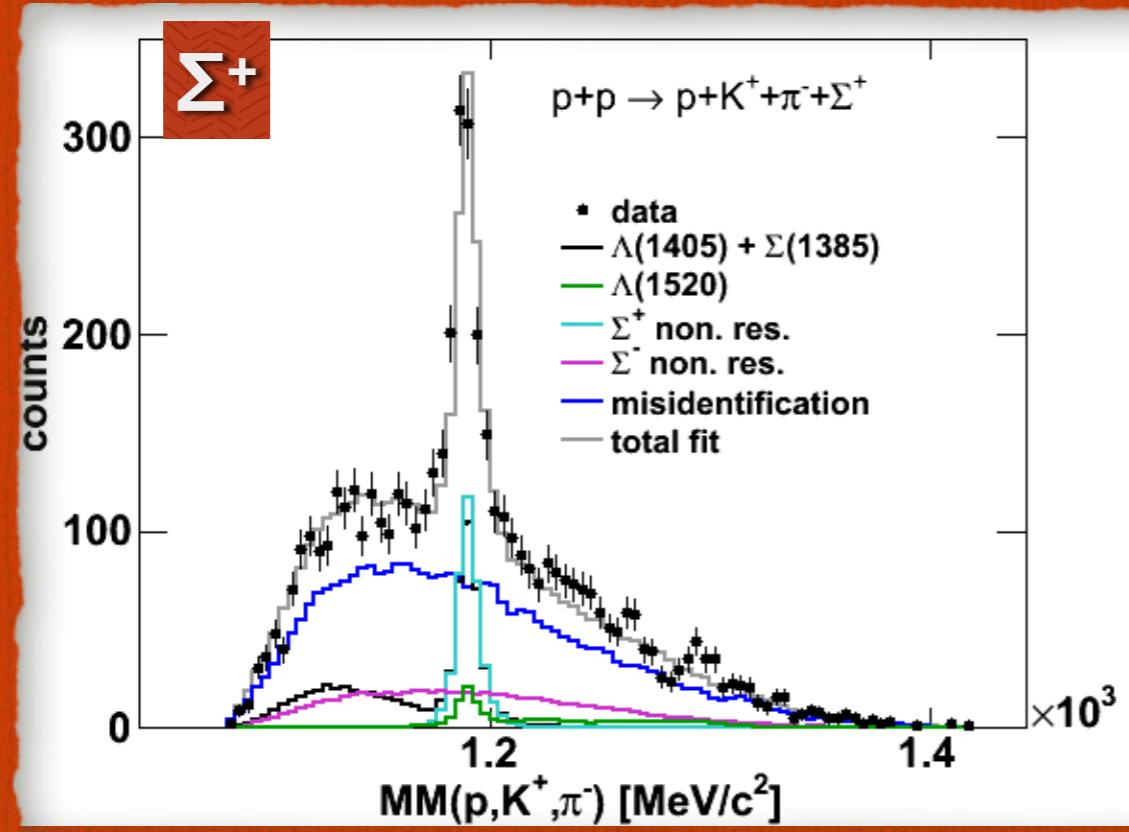
p+p collisions

proton-proton



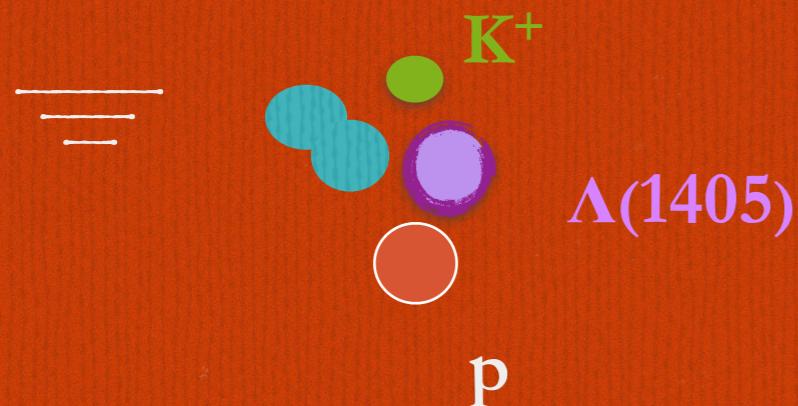
$\sigma(\Lambda(1405))/\sigma(\Sigma(1385)^0) \approx 2.8$, known from analysis of neutral decay channels

Conservative estimation: $\sigma(\Sigma(1385)^0) = 0.5\sigma(\Sigma(1385)^+ \rightarrow \Sigma^-)$
 $\Rightarrow \sigma(\Lambda(1405))/\sigma(\Sigma(1385)^0) \approx 1.0 \rightarrow$ no influence on pole mass or line shape of $\Lambda(1405)$



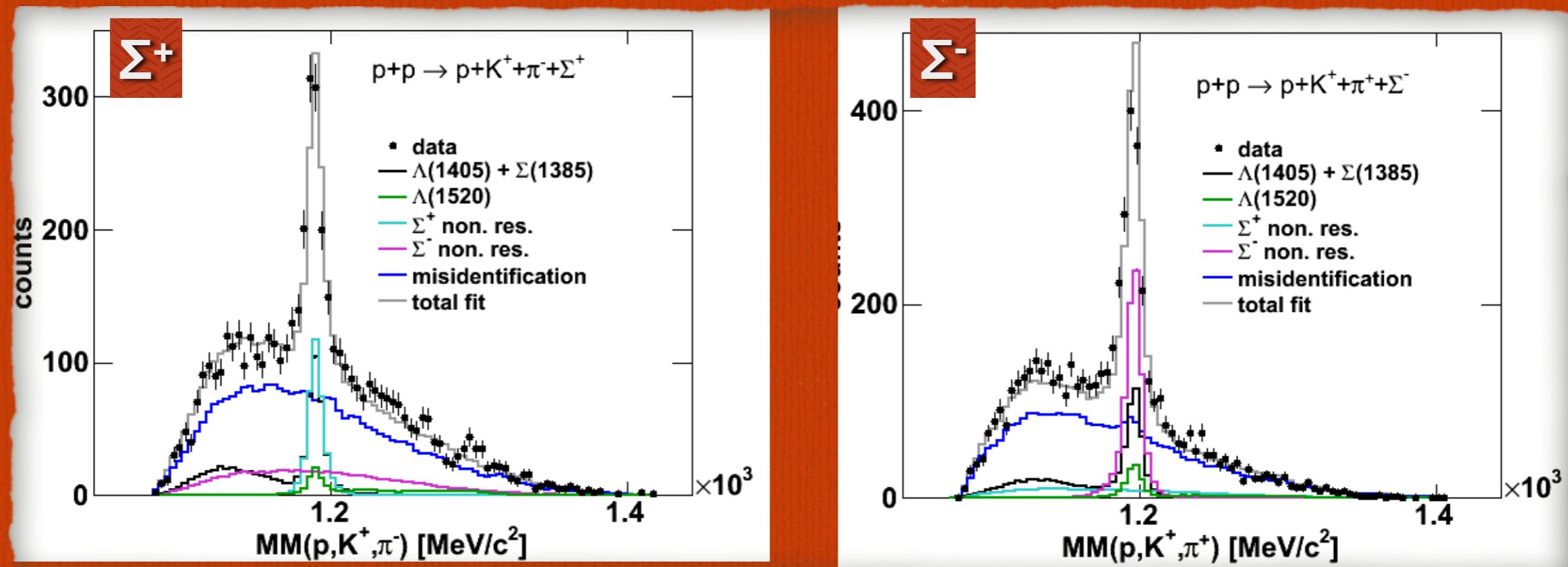
p+p collisions

proton-proton



$\sigma(\Lambda(1405))/\sigma(\Sigma(1385)^0) \approx 2.8$, known from analysis of neutral decay channels

Conservative estimation: $\sigma(\Sigma(1385)^0) = 0.5\sigma(\Sigma(1385)^+ \rightarrow \Sigma^-)$
 $\Rightarrow \sigma(\Lambda(1405))/\sigma(\Sigma(1385)^0) \approx 1.0 \rightarrow$ no influence on pole mass or line shape of $\Lambda(1405)$

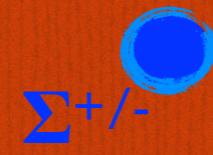


p+p collisions

proton-proton



$\Lambda(1405)$

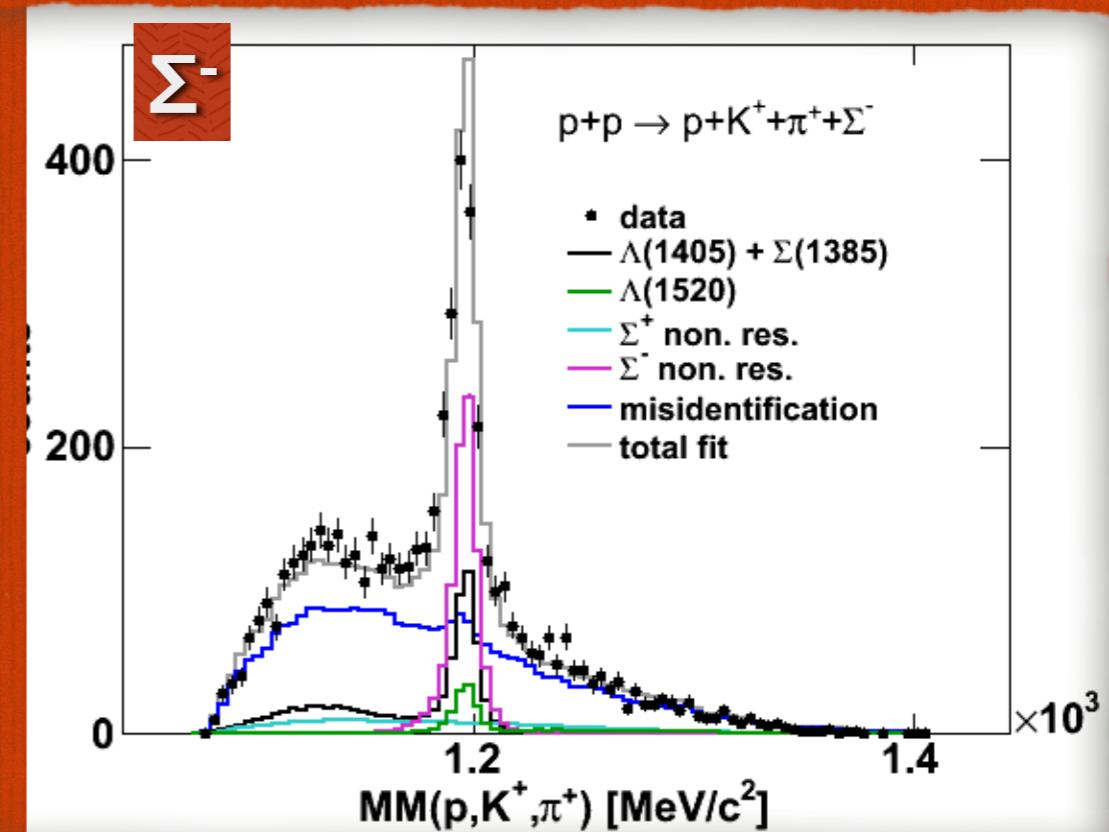
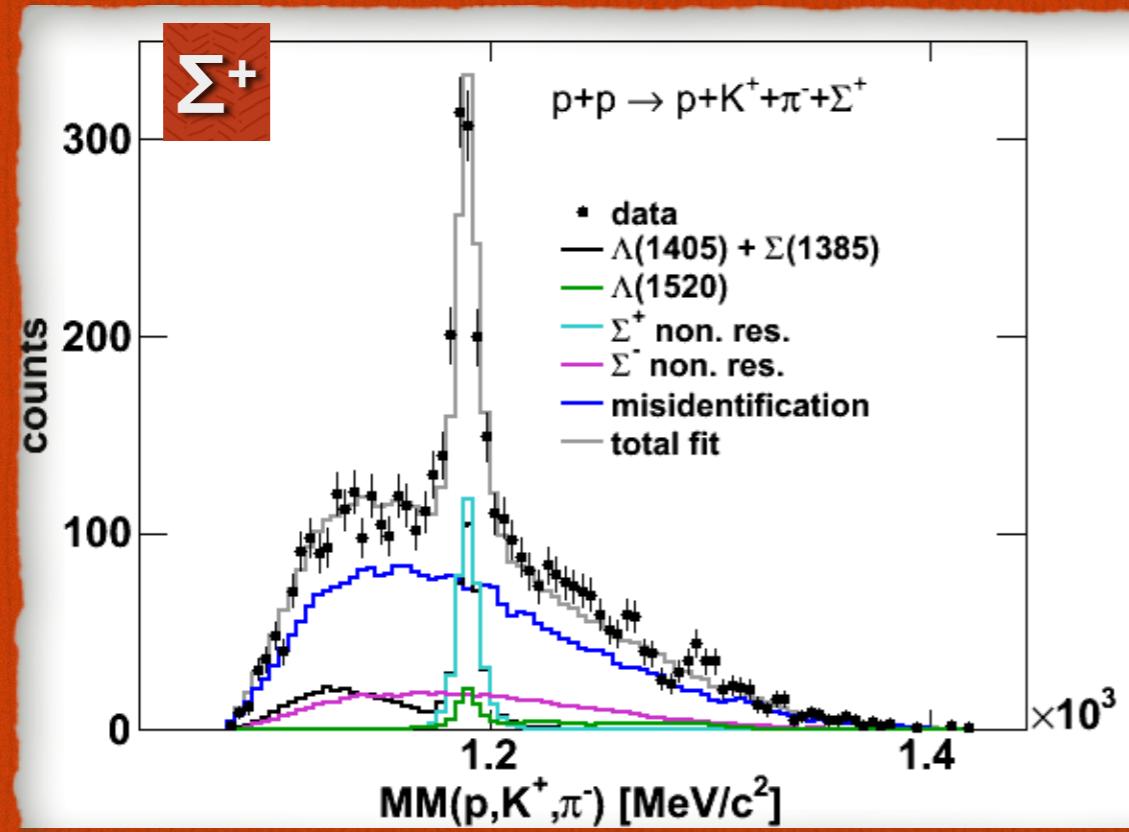


$\pi^{-/+}$



$\sigma(\Lambda(1405))/\sigma(\Sigma(1385)^0) \approx 2.8$, known from analysis of neutral decay channels

Conservative estimation: $\sigma(\Sigma(1385)^0) = 0.5\sigma(\Sigma(1385)^+ \rightarrow \Sigma^-)$
 $\Rightarrow \sigma(\Lambda(1405))/\sigma(\Sigma(1385)^0) \approx 1.0 \rightarrow$ no influence on pole mass or line shape of $\Lambda(1405)$



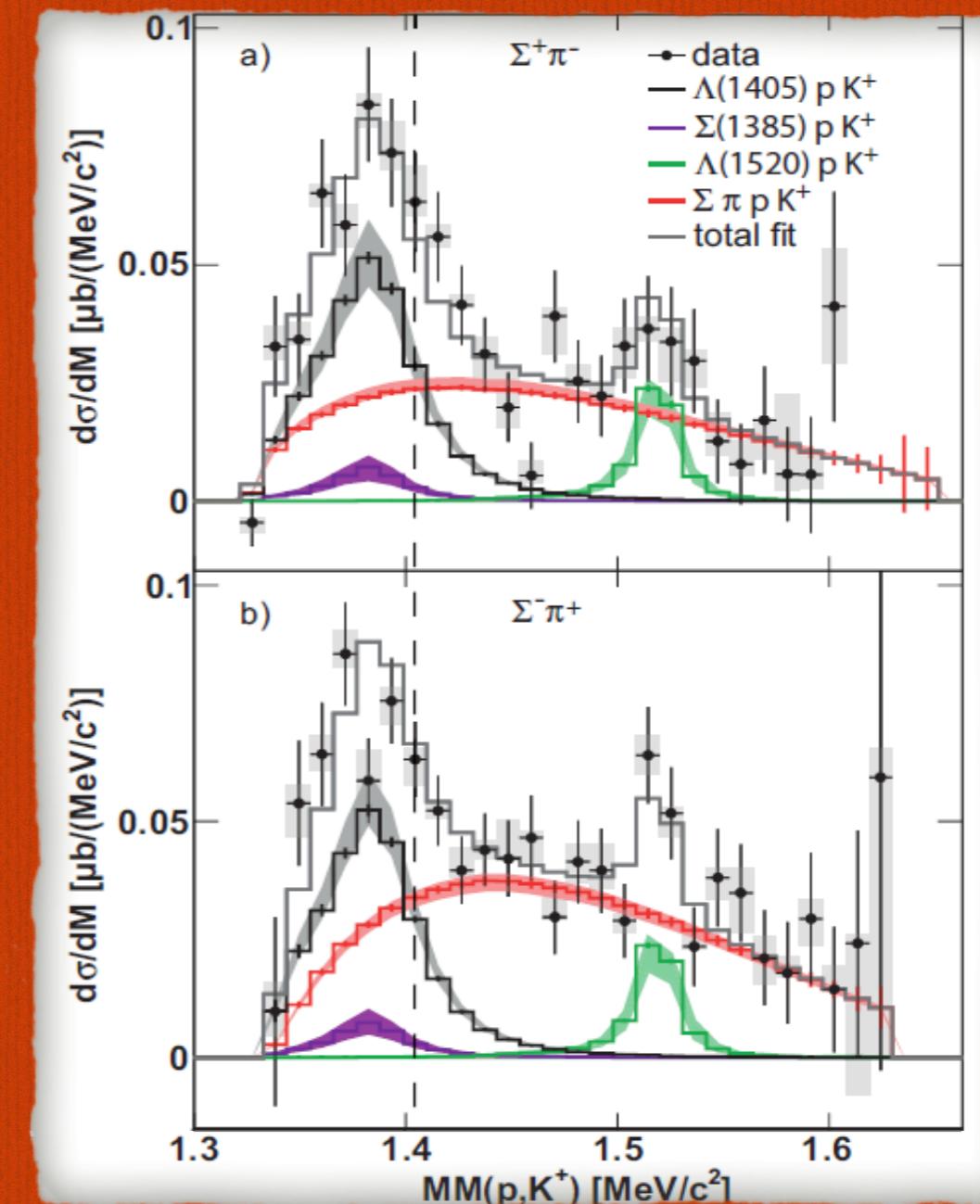
Our Special Resonance

J. Siebenson, E. Epple

Efficiency and Acceptance Corrected data

Channel	Cross section
$p+p \rightarrow \Lambda(1405)+p+K^+$	$9.2 \pm 0.9 \pm 0.7^{+3.3}_{-1.0} \mu b$
$p+p \rightarrow \Sigma(1385)^0+p+K^+$	$5.6 \pm 0.5^{+2.0}_{-1.1} \mu b$
$p+p \rightarrow \Lambda(1520)+p+K^+$	$5.6 \pm 1.1 \pm 0.4^{+1.1}_{-1.6} \mu b$
$p+p \rightarrow \Sigma^- + \Delta^{++}(1232) + K^+$	$7.7 \pm 0.9 \pm 0.5^{+0.3}_{-0.9} \mu b$
$p+p \rightarrow \Sigma^+ + \pi^- + p+K^+$	$5.4 \pm 0.5 \pm 0.4^{+1.0}_{-2.1} \mu b$

Phys. Rev. C 87, 025201 (2013)

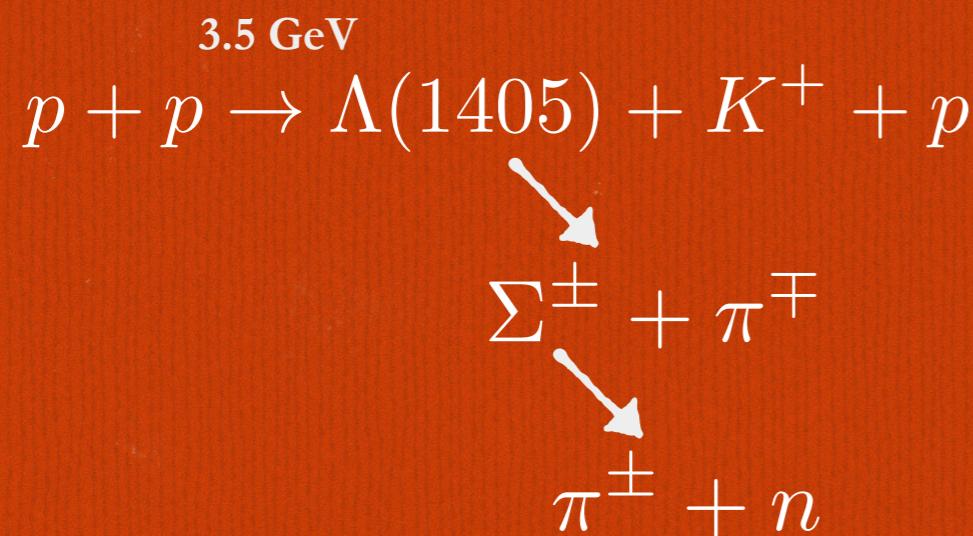


Our Special Resonance

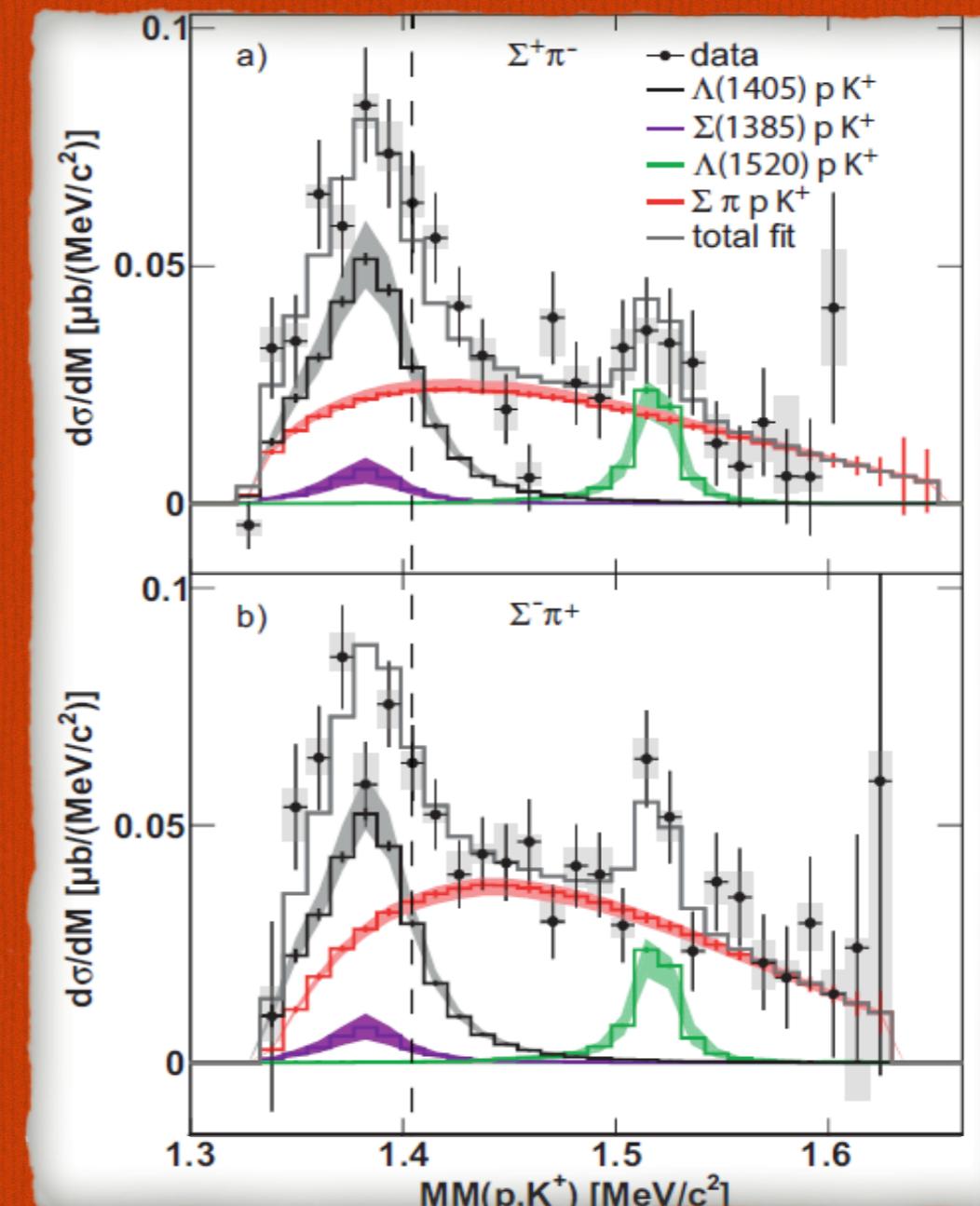
J. Siebenson, E. Epple

Efficiency and Acceptance Corrected data

Channel	Cross section
$p+p \rightarrow \Lambda(1405)+p+K^+$	$9.2 \pm 0.9 \pm 0.7^{+3.3}_{-1.0} \mu b$
$p+p \rightarrow \Sigma(1385)^0+p+K^+$	$5.6 \pm 0.5^{+2.0}_{-1.1} \mu b$
$p+p \rightarrow \Lambda(1520)+p+K^+$	$5.6 \pm 1.1 \pm 0.4^{+1.1}_{-1.6} \mu b$
$p+p \rightarrow \Sigma^- + \Delta^{++}(1232) + K^+$	$7.7 \pm 0.9 \pm 0.5^{+0.3}_{-0.9} \mu b$
$p+p \rightarrow \Sigma^+ + \pi^- + p+K^+$	$5.4 \pm 0.5 \pm 0.4^{+1.0}_{-2.1} \mu b$



Phys. Rev. C 87, 025201 (2013)



Our Special Resonance

J. Siebenson, E. Epple

Efficiency and Acceptance Corrected data

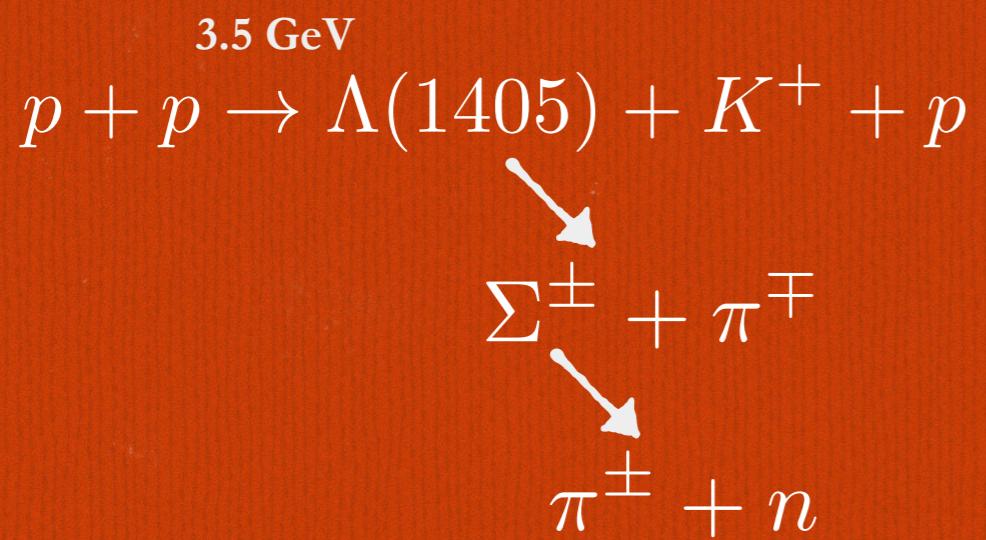
Channel	Cross section
$p+p \rightarrow \Lambda(1405)+p+K^+$	$9.2 \pm 0.9 \pm 0.7^{+3.3}_{-1.0} \mu b$
$p+p \rightarrow \Sigma(1385)^0+p+K^+$	$5.6 \pm 0.5^{+2.0}_{-1.1} \mu b$
$p+p \rightarrow \Lambda(1520)+p+K^+$	$5.6 \pm 1.1 \pm 0.4^{+1.1}_{-1.6} \mu b$
$p+p \rightarrow \Sigma^- + \Delta^{++}(1232) + K^+$	$7.7 \pm 0.9 \pm 0.5^{+0.3}_{-0.9} \mu b$
$p+p \rightarrow \Sigma^+ + \pi^- + p + K^+$	$5.4 \pm 0.5 \pm 0.4^{+1.0}_{-2.1} \mu b$

Our Special Resonance

J. Siebenson, E. Epple

Efficiency and Acceptance Corrected data

Channel	Cross section
$p+p \rightarrow \Lambda(1405) + p + K^+$	$9.2 \pm 0.9 \pm 0.7^{+3.3}_{-1.0} \mu b$
$p+p \rightarrow \Sigma(1385)^0 + p + K^+$	$5.6 \pm 0.5^{+2.0}_{-1.1} \mu b$
$p+p \rightarrow \Lambda(1520) + p + K^+$	$5.6 \pm 1.1 \pm 0.4^{+1.1}_{-1.6} \mu b$
$p+p \rightarrow \Sigma^- + \Delta^{++}(1232) + K^+$	$7.7 \pm 0.9 \pm 0.5^{+0.3}_{-0.9} \mu b$
$p+p \rightarrow \Sigma^+ + \pi^- + p + K^+$	$5.4 \pm 0.5 \pm 0.4^{+1.0}_{-2.1} \mu b$

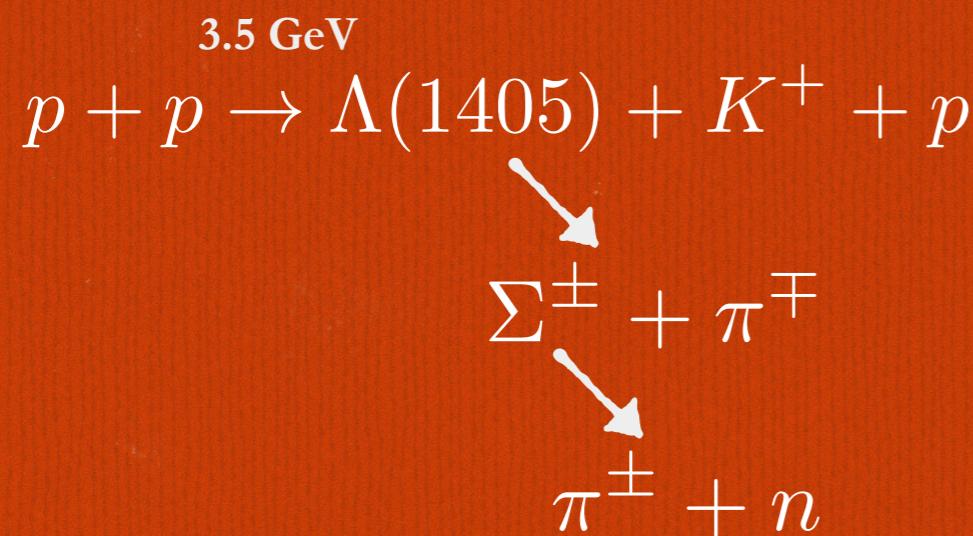


Our Special Resonance

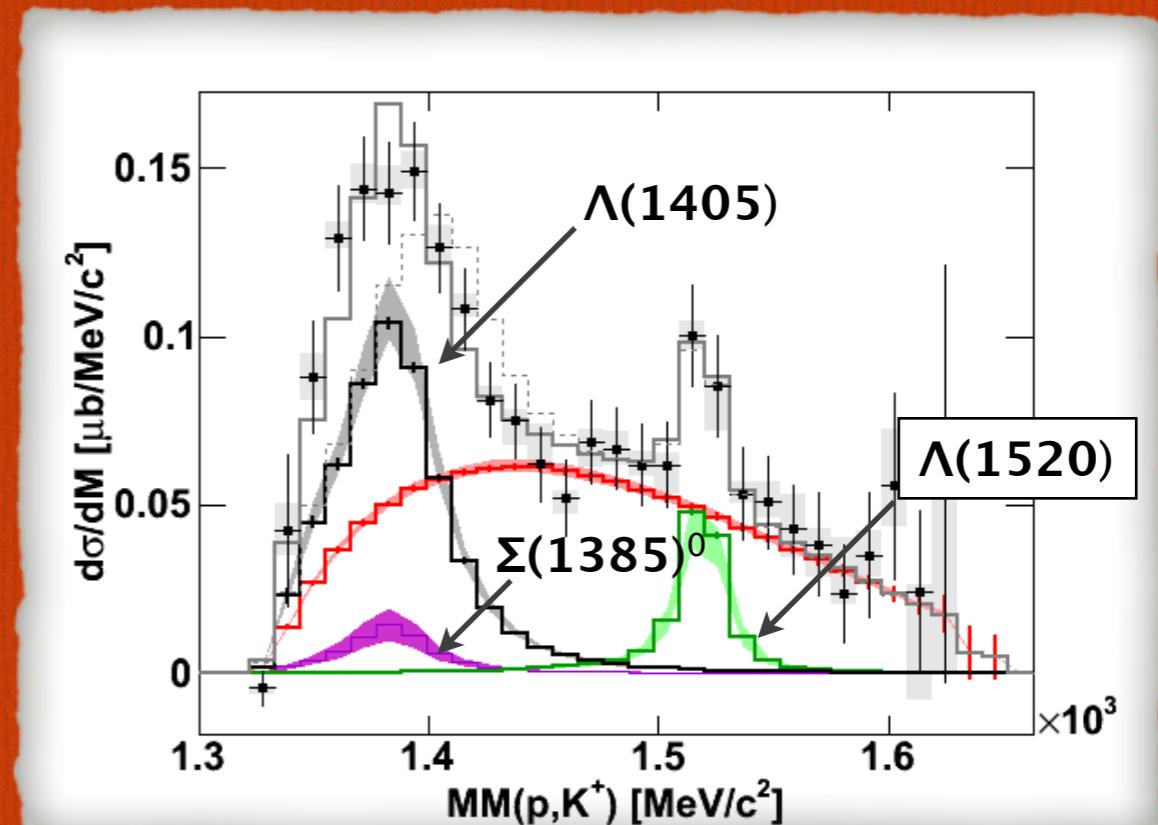
J. Siebenson, E. Epple

Efficiency and Acceptance Corrected data

Channel	Cross section
$p+p \rightarrow \Lambda(1405) + p + K^+$	$9.2 \pm 0.9 \pm 0.7^{+3.3}_{-1.0} \mu b$
$p+p \rightarrow \Sigma(1385)^0 + p + K^+$	$5.6 \pm 0.5^{+2.0}_{-1.1} \mu b$
$p+p \rightarrow \Lambda(1520) + p + K^+$	$5.6 \pm 1.1 \pm 0.4^{+1.1}_{-1.6} \mu b$
$p+p \rightarrow \Sigma^- + \Delta^{++}(1232) + K^+$	$7.7 \pm 0.9 \pm 0.5^{+0.3}_{-0.9} \mu b$
$p+p \rightarrow \Sigma^+ + \pi^- + p + K^+$	$5.4 \pm 0.5 \pm 0.4^{+1.0}_{-2.1} \mu b$



G. Agakishiev et al. [HADES] Phys. Rev. C 87 (2013) 025201.
G. Agakishiev et al. [HADES] Nucl. Phys. A 881 (2012) 178–186.

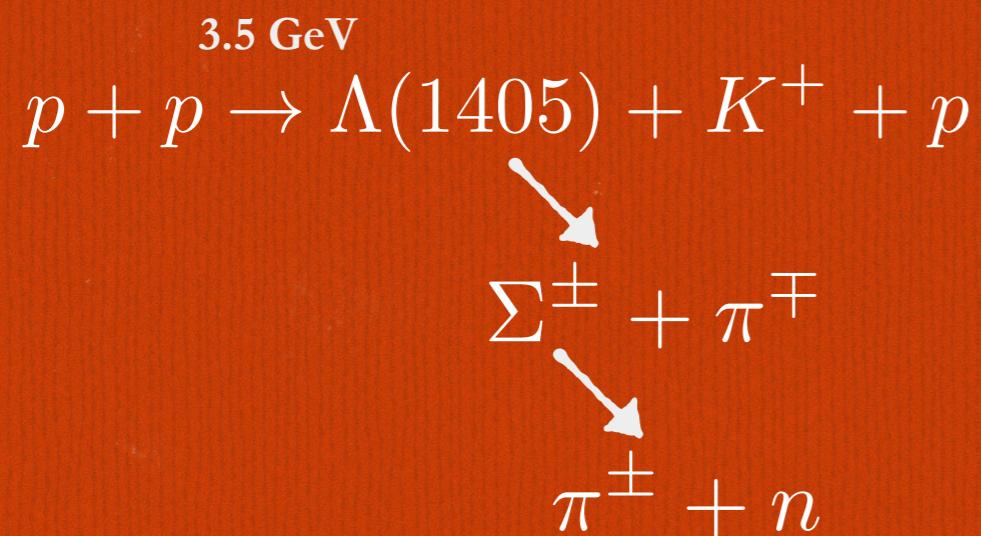


Our Special Resonance

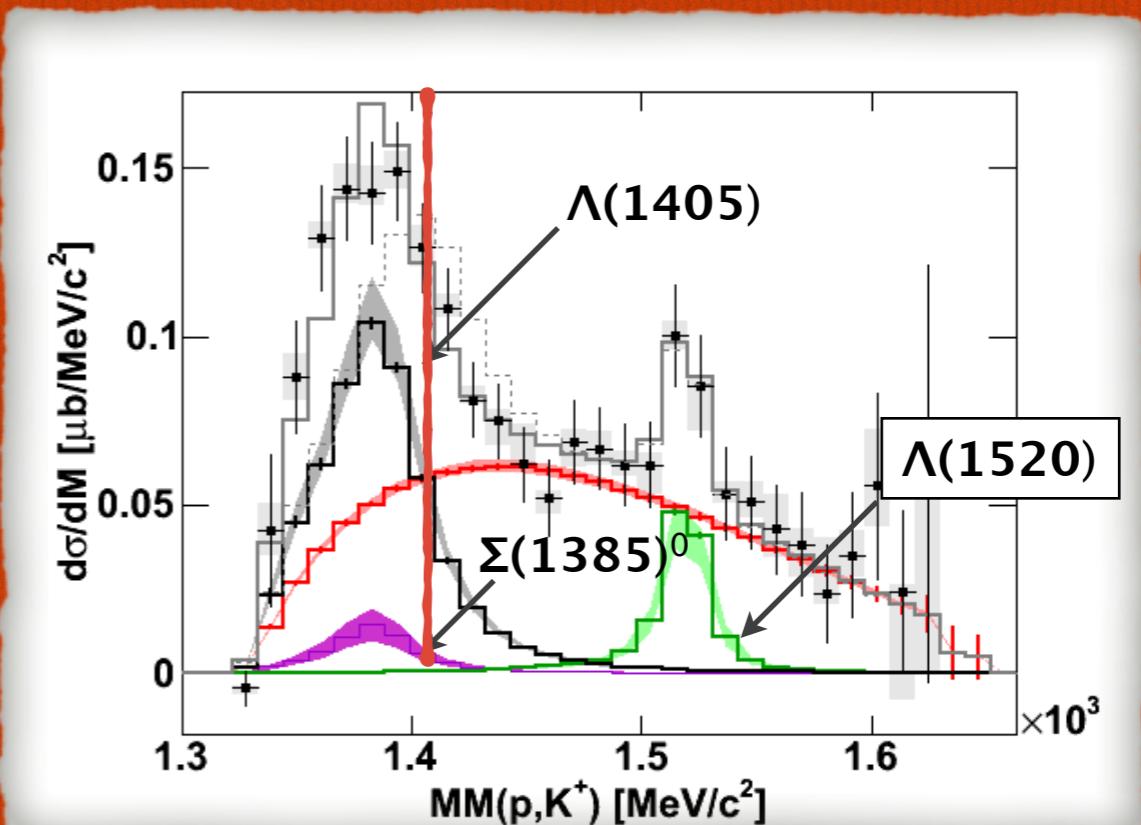
J. Siebenson, E. Epple

Efficiency and Acceptance Corrected data

Channel	Cross section
$p+p \rightarrow \Lambda(1405) + p + K^+$	$9.2 \pm 0.9 \pm 0.7^{+3.3}_{-1.0} \mu b$
$p+p \rightarrow \Sigma(1385)^0 + p + K^+$	$5.6 \pm 0.5^{+2.0}_{-1.1} \mu b$
$p+p \rightarrow \Lambda(1520) + p + K^+$	$5.6 \pm 1.1 \pm 0.4^{+1.1}_{-1.6} \mu b$
$p+p \rightarrow \Sigma^- + \Delta^{++}(1232) + K^+$	$7.7 \pm 0.9 \pm 0.5^{+0.3}_{-0.9} \mu b$
$p+p \rightarrow \Sigma^+ + \pi^- + p + K^+$	$5.4 \pm 0.5 \pm 0.4^{+1.0}_{-2.1} \mu b$



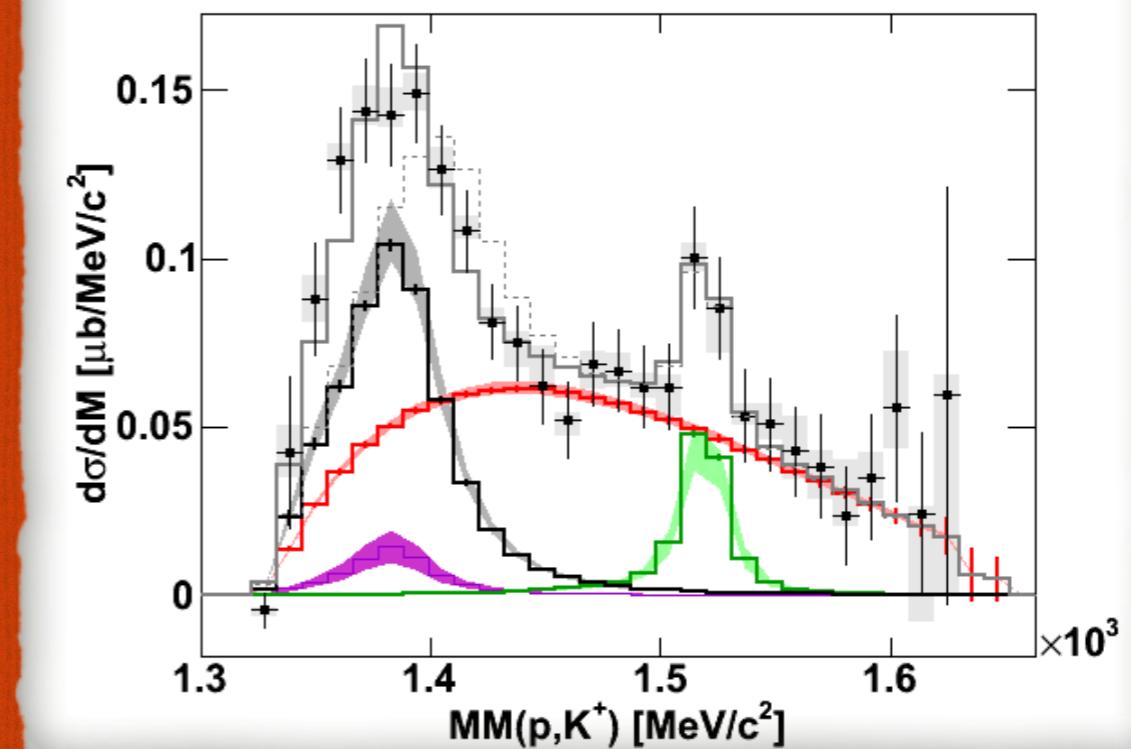
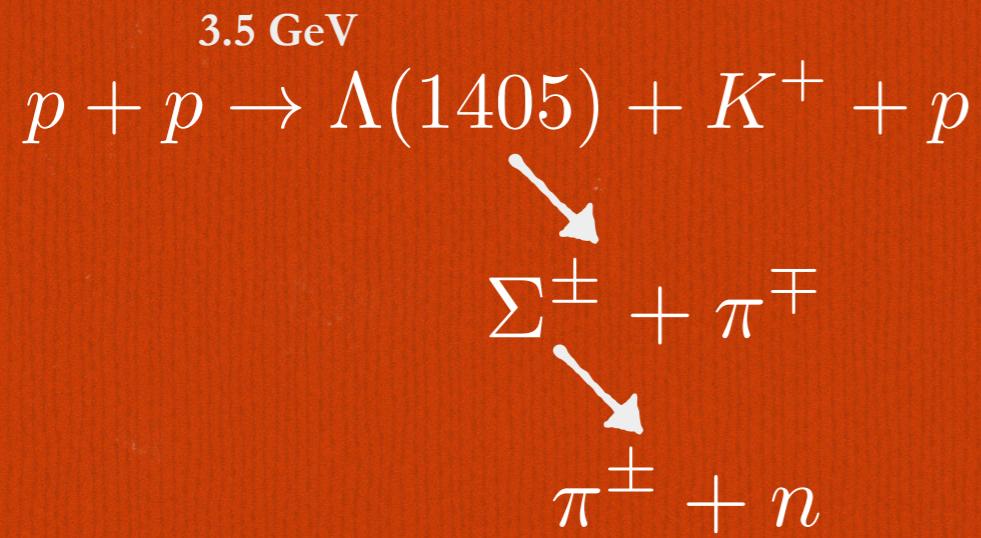
G. Agakishiev et al. [HADES] Phys. Rev. C 87 (2013) 025201.
G. Agakishiev et al. [HADES] Nucl. Phys. A 881 (2012) 178–186.



Our Special Resonance

J. Siebenson, E. Epple

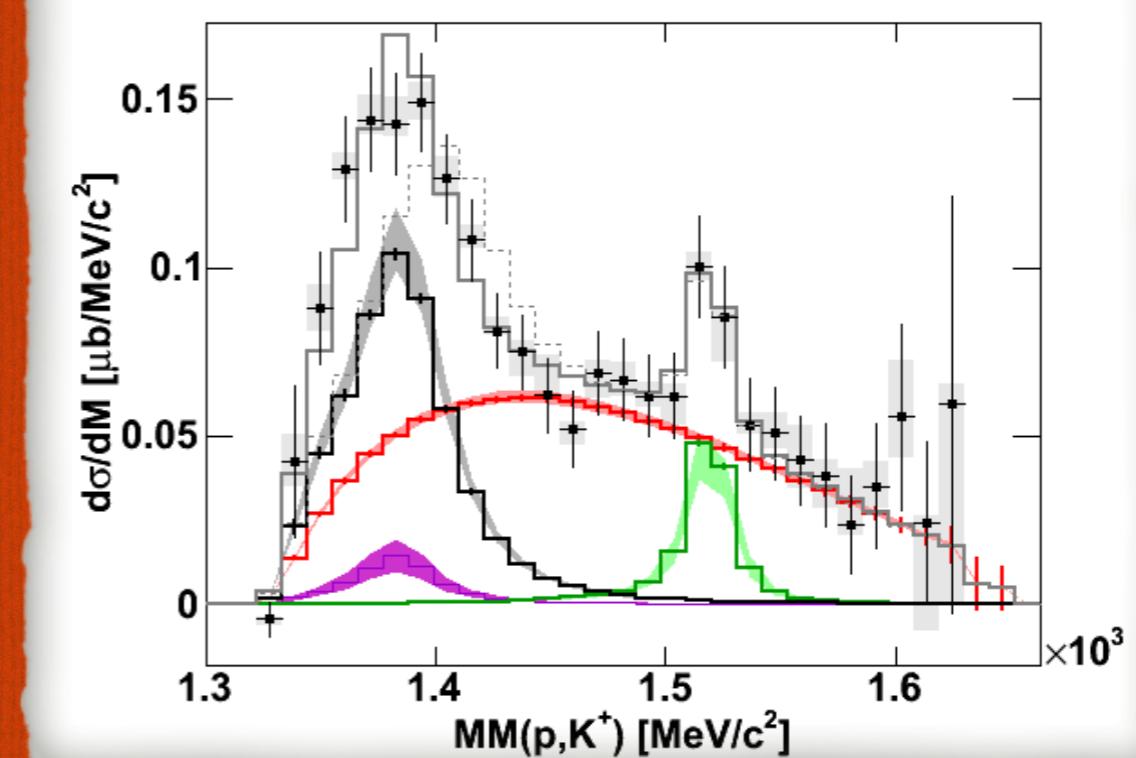
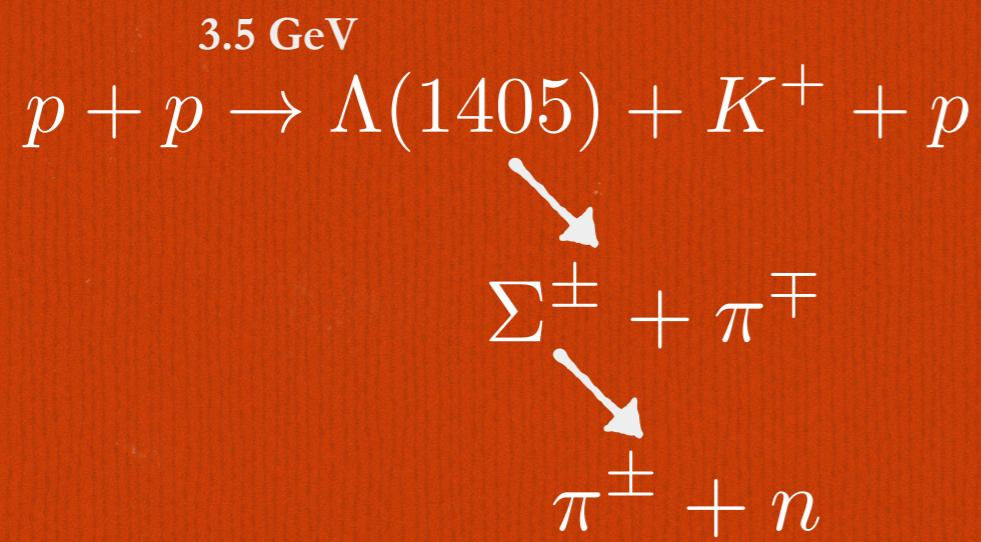
$$\Lambda(1405) : M = 1405 \text{ MeV}/c^2 (?) \Gamma = 50 \text{ MeV}/c^2 (B = 1, S = -1 J^P = 1/2^-)$$



Our Special Resonance

J. Siebenson, E. Epple

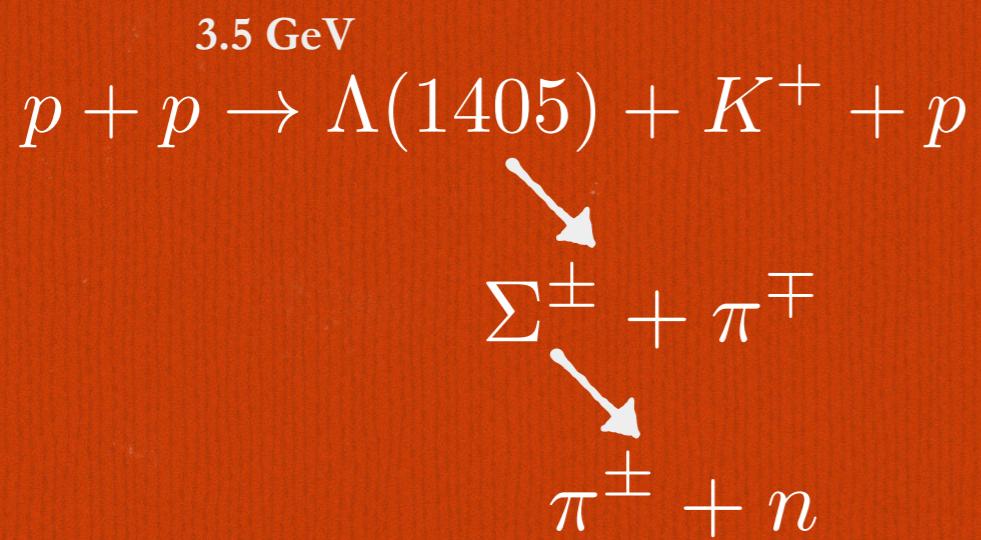
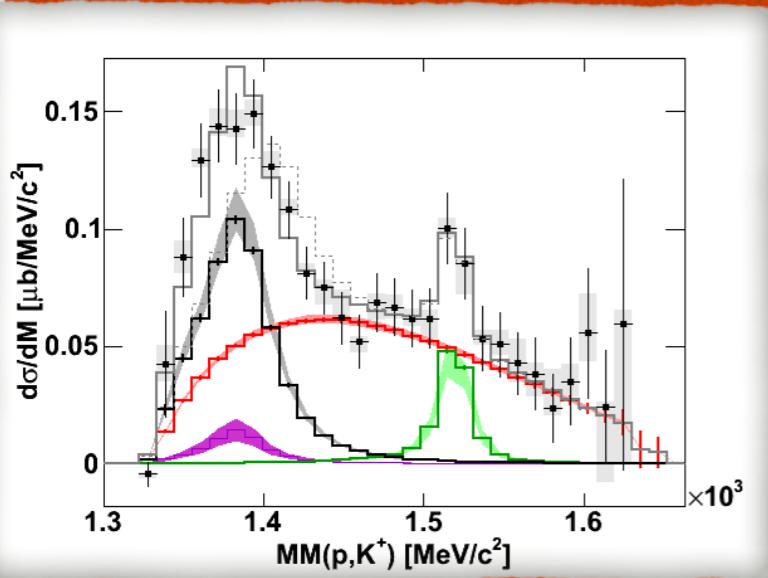
$\Lambda(1405) : M = 1405 \text{ MeV}/c^2 (?) \Gamma = 50 \text{ MeV}/c^2 (B = 1, S = -1 J^P = 1/2^-)$



Our Special Resonance

J. Siebenson, E. Eppe

$\Lambda(1405) : M = 1405 \text{ MeV}/c^2 (?) \Gamma = 50 \text{ MeV}/c^2 (B = 1, S = -1 J^P = 1/2^-)$

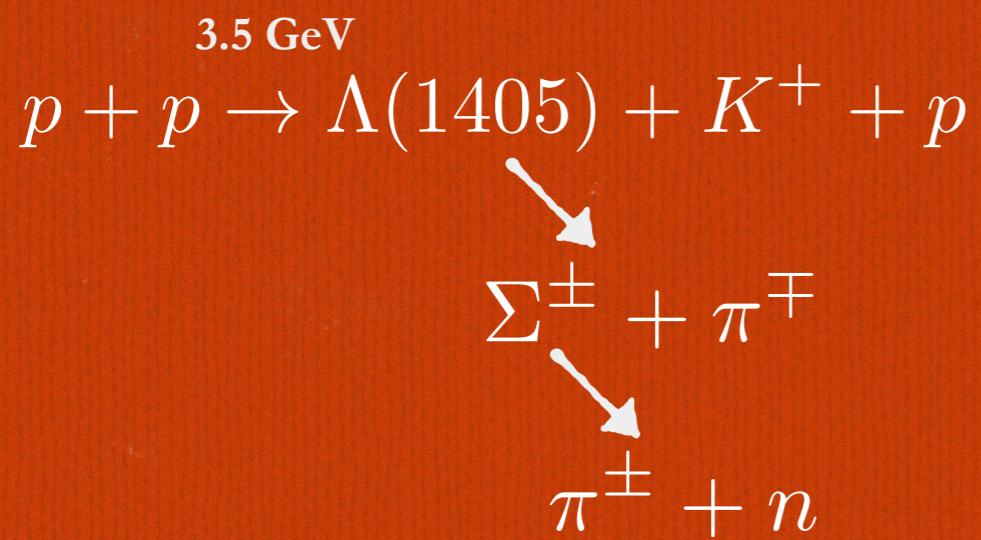
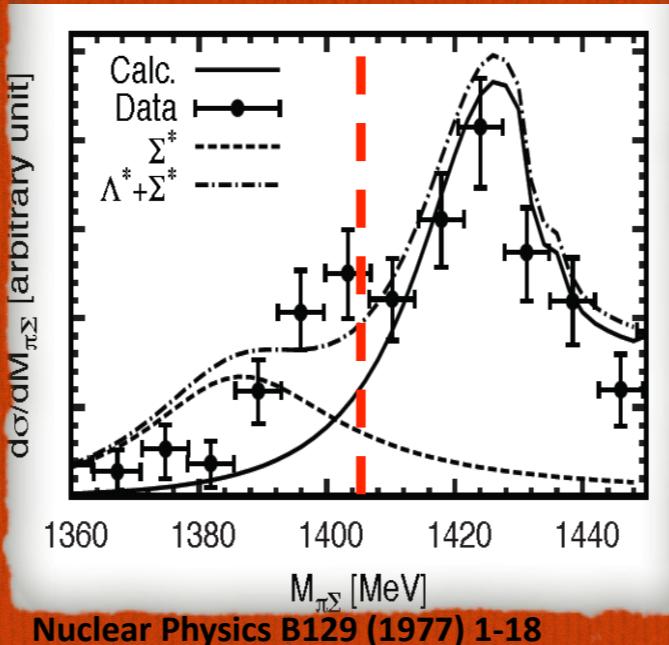
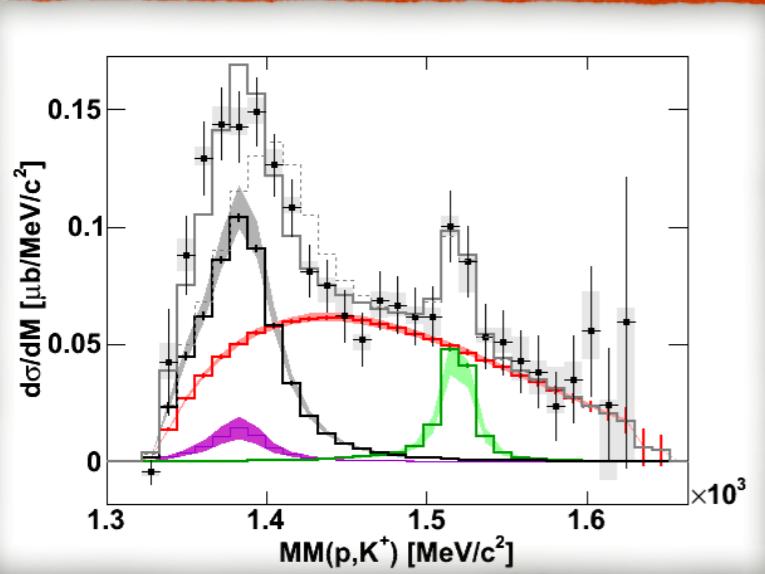


Our Special Resonance

J. Siebenson, E. Epple

$\Lambda(1405) : M = 1405 \text{ MeV}/c^2 (?) \Gamma = 50 \text{ MeV}/c^2 (B = 1, S = -1 J^P = 1/2^-)$

K+d 0.68 – 0.84 GeV/c



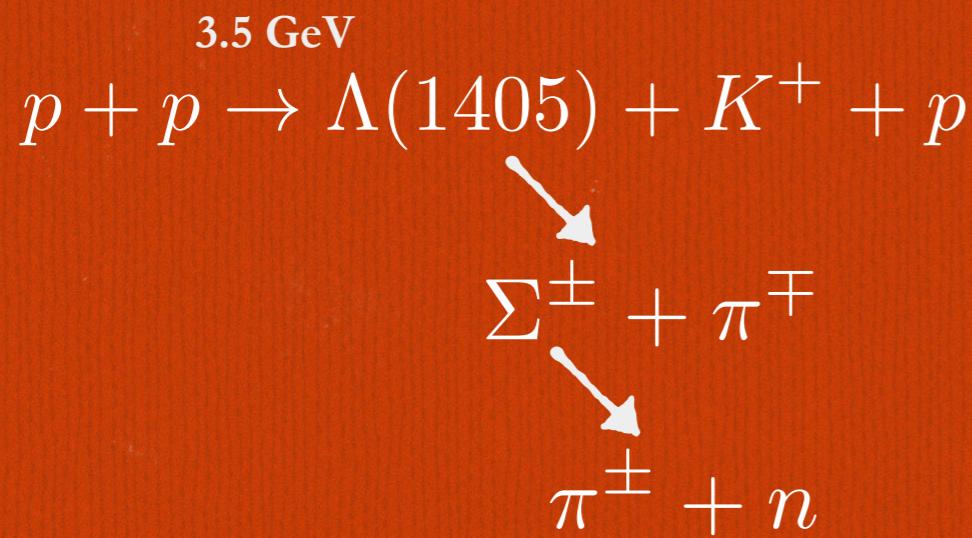
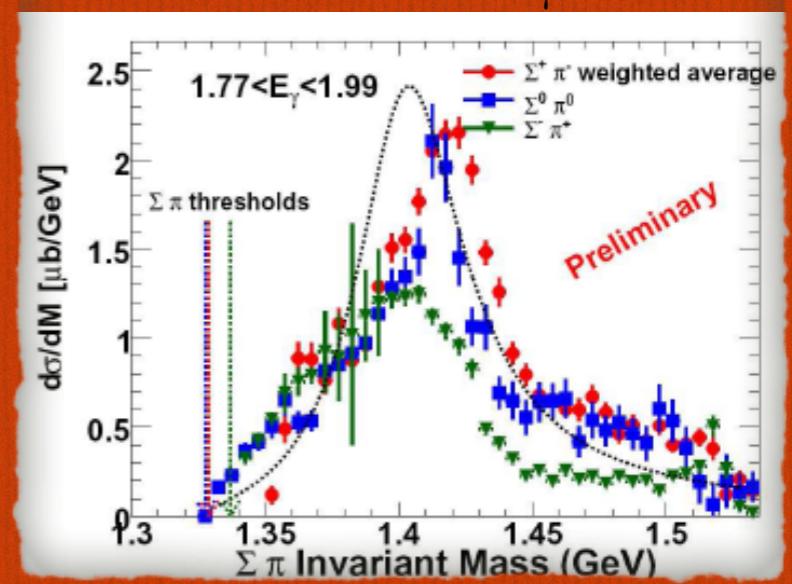
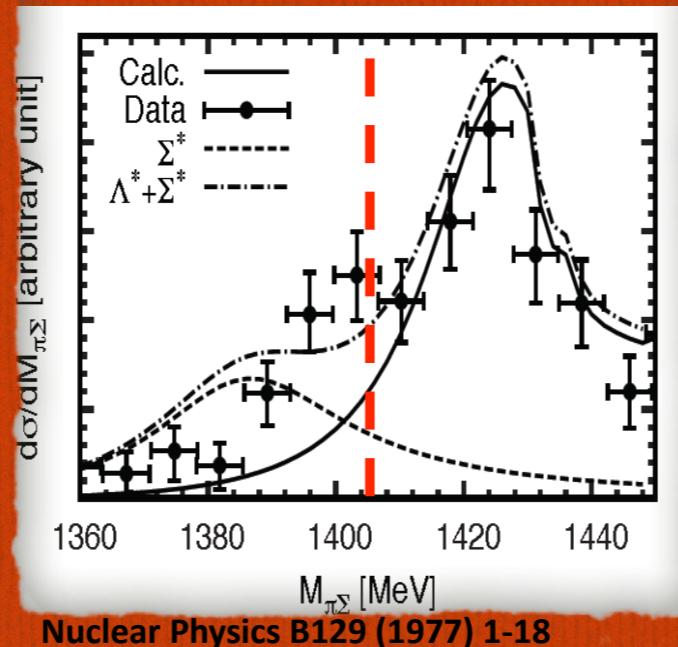
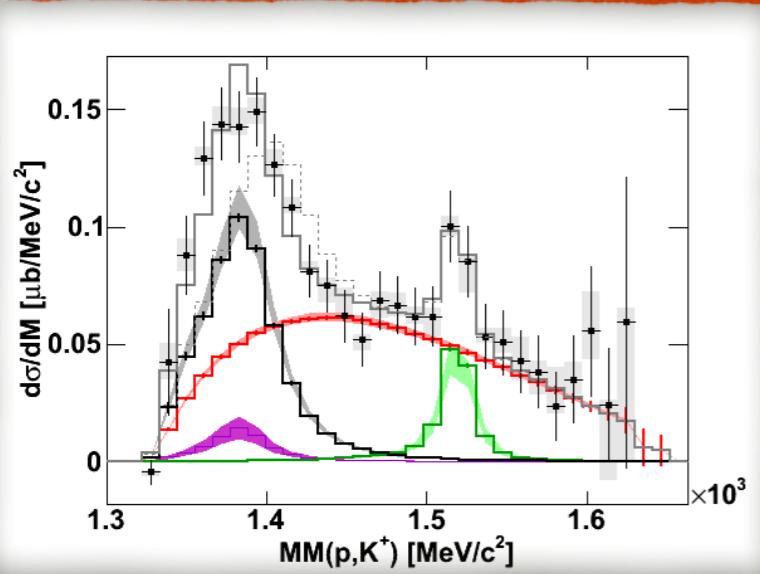
Our Special Resonance

J. Siebenson, E. Epple

$\Lambda(1405) : M = 1405 \text{ MeV}/c^2 (?) \Gamma = 50 \text{ MeV}/c^2 (B = 1, S = -1 J^P = 1/2^-)$

K+d 0.68 – 0.84 GeV/c

$\gamma + p \quad 1.77 < p_\gamma < 1.99$



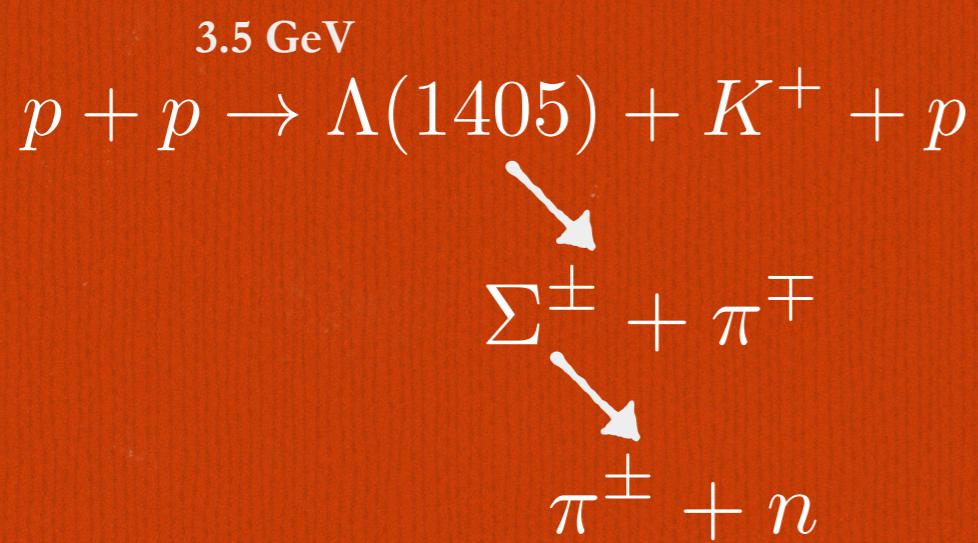
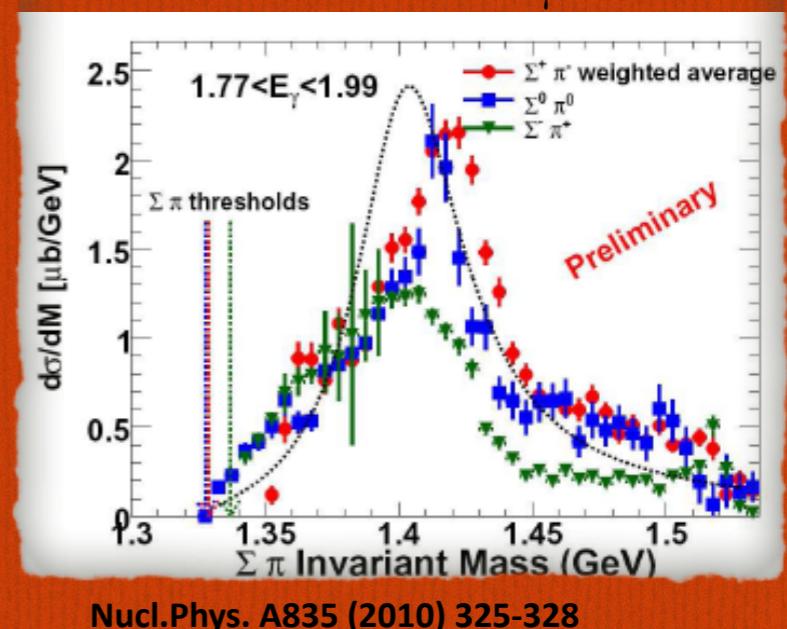
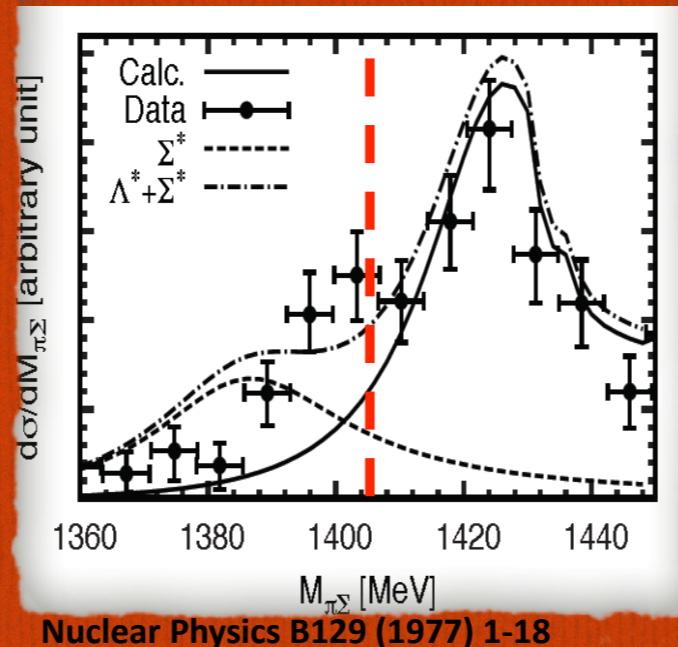
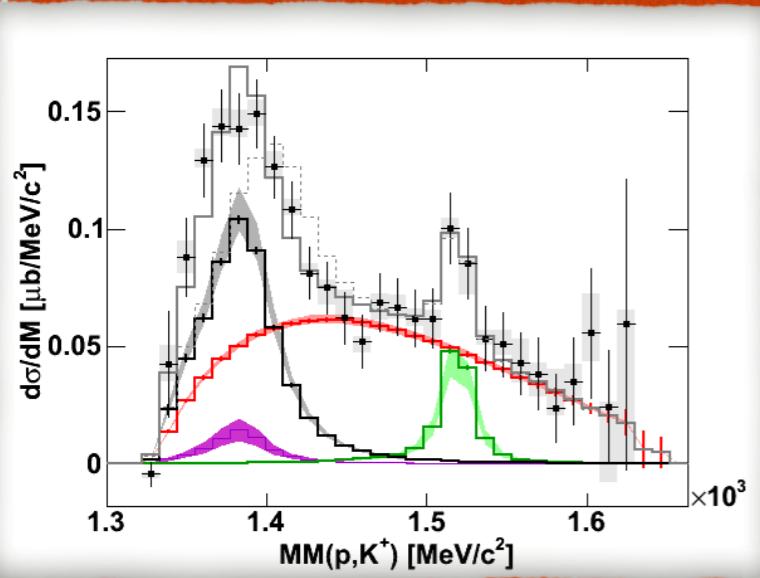
Our Special Resonance

J. Siebenson, E. Epple

$\Lambda(1405) : M = 1405 \text{ MeV}/c^2 (?) \Gamma = 50 \text{ MeV}/c^2 (B = 1, S = -1 J^P = 1/2^-)$

K+d 0.68 – 0.84 GeV/c

$\gamma + p \quad 1.77 < p_\gamma < 1.99$



The spectral shape of the $\Lambda(1405)$ looks different for different initial state reactions.

p+p and π+p reactions look similar
Different coupling of the 2 poles??

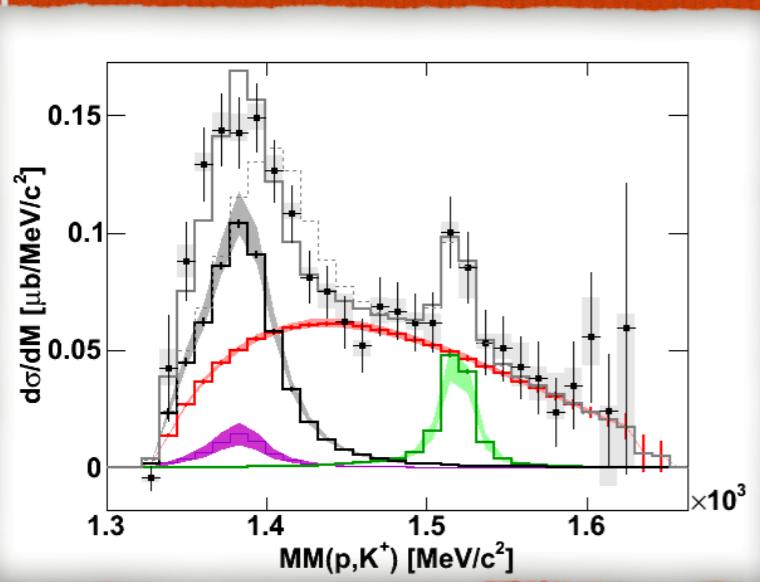
Our Special Resonance

J. Siebenson, E. Epple

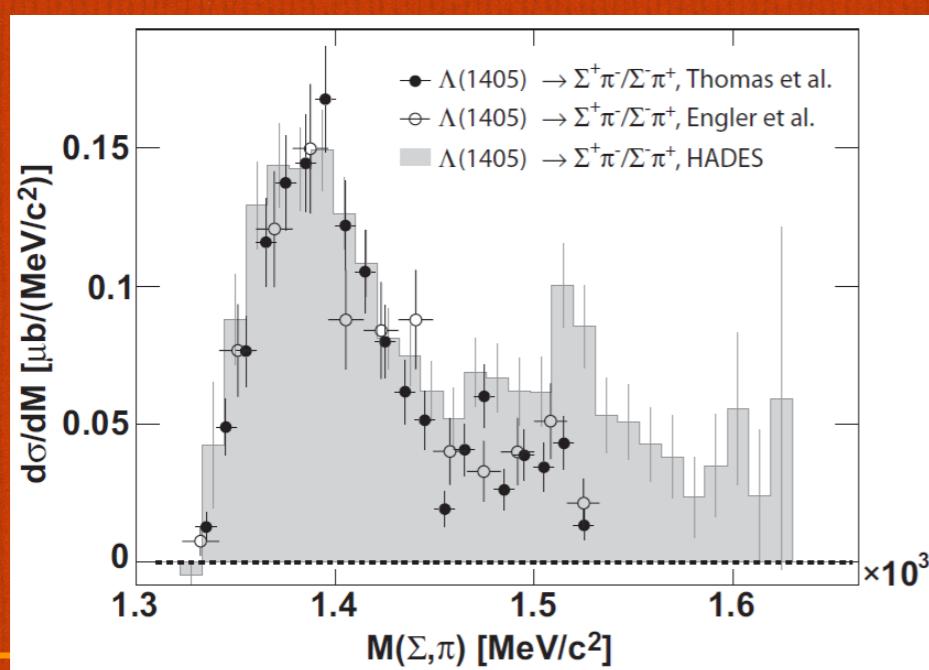
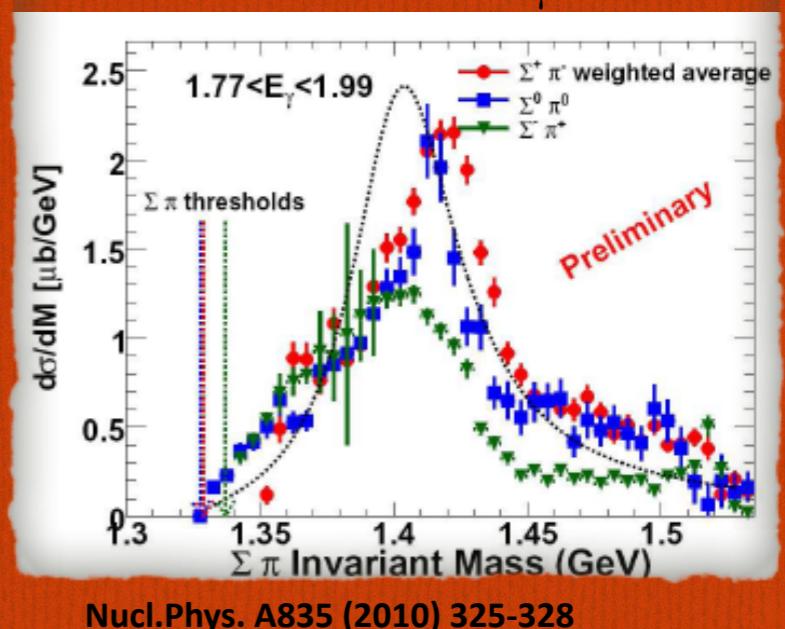
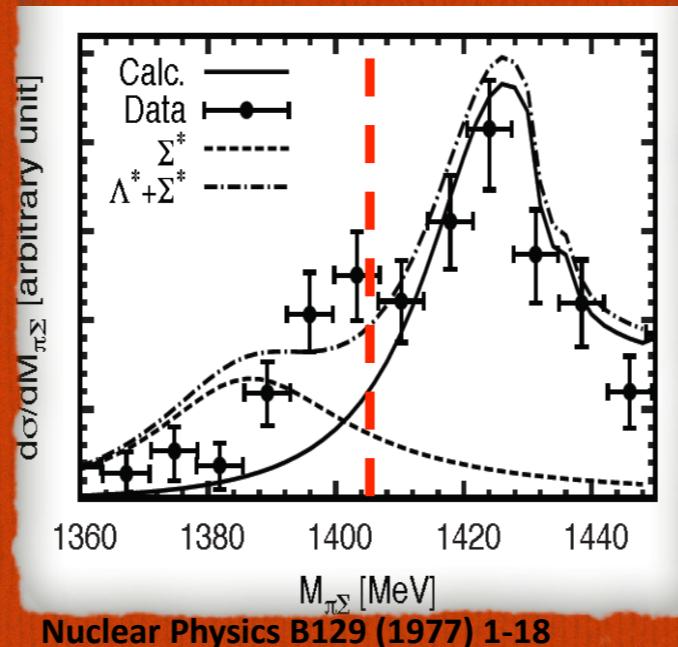
$\Lambda(1405) : M = 1405 \text{ MeV}/c^2 (?) \Gamma = 50 \text{ MeV}/c^2 (B = 1, S = -1 J^P = 1/2^-)$

K+d 0.68 – 0.84 GeV/c

$\gamma + p \quad 1.77 < p_\gamma < 1.99$



Nucl.Phys. B56 (1973) 46-51, Phys.Rev.Lett. 15 (1965) 224



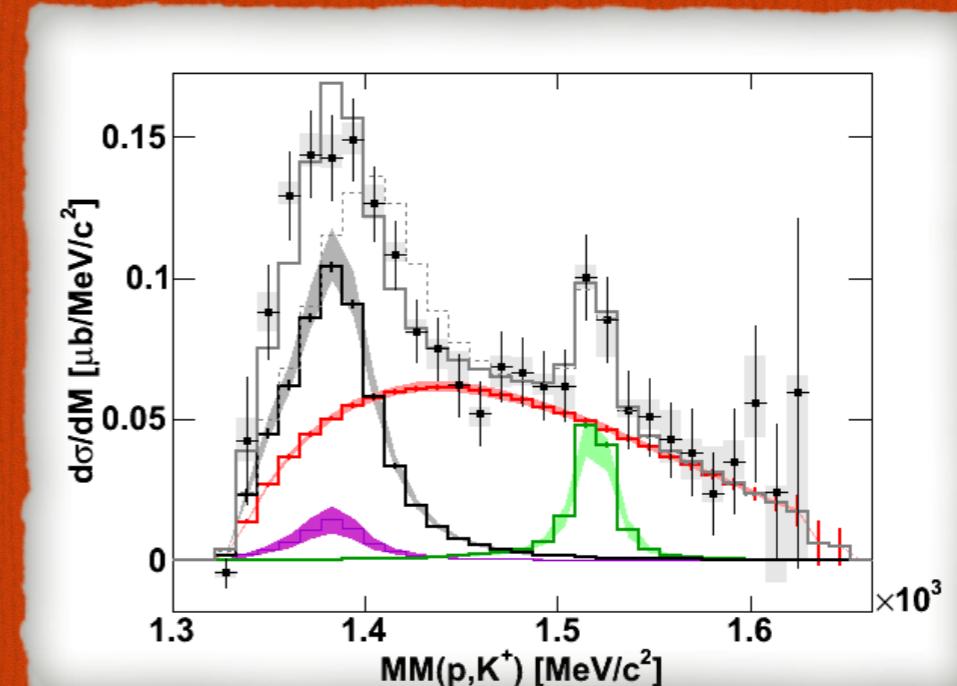
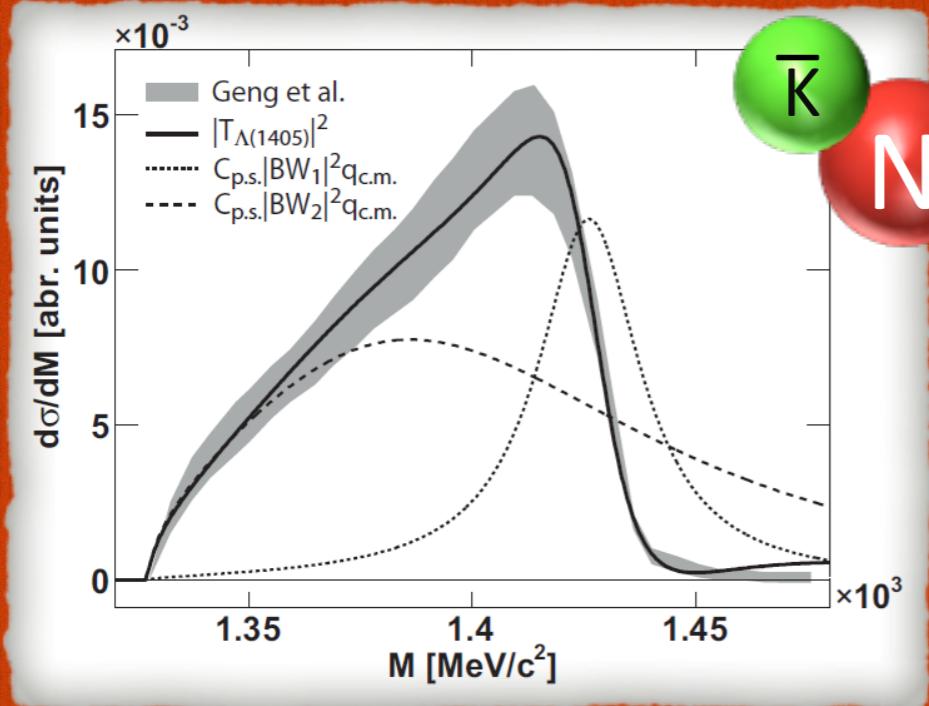
The spectral shape of the $\Lambda(1405)$ looks different for different initial state reactions.
 $p + p$ and $\pi + p$ reactions look similar
 Different coupling of the 2 poles??

Our Special Resonance

J. Siebenson

Chiral Ansatz

Nucl.Phys. A881 (2012) 98-114



Nucl.Phys. A881 (2012) 98-114

$$\begin{aligned} m_{0,1} &= \text{Re}(z_1) = 1424^{+7}_{-23} \text{ MeV} \\ \Gamma_{0,1} &= 2 \cdot \text{Im}(z_1) = 52^{+6}_{-28} \text{ MeV} \\ m_{0,2} &= \text{Re}(z_2) = 1381^{+18}_{-6} \text{ MeV} \\ \Gamma_{0,2} &= 2 \cdot \text{Im}(z_2) = 162^{+38}_{-16} \text{ MeV} \end{aligned}$$

Two poles combined such to reproduce the calculation for p+p

- S.F. for the $\Sigma^0\pi^0$ decay
- lower energy than for HADES

Calculation does not match the HADES data

Our Special Resonance

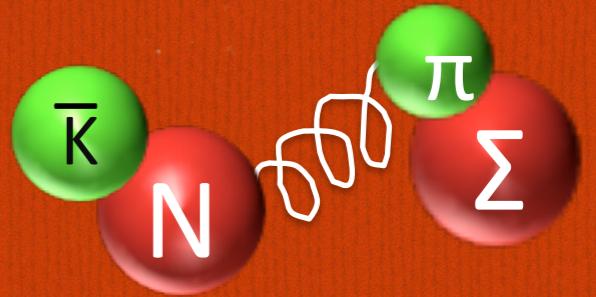
J. Siebenson

$$m_{0,1} = \text{Re}(z_1) = 1424^{+7}_{-23} \text{ MeV}$$

$$\Gamma_{0,1} = 2 * \text{Im}(z_1) = 52^{+6}_{-28} \text{ MeV}$$

$$m_{0,2} = \text{Re}(z_2) = 1381^{+18}_{-6} \text{ MeV}$$

$$\Gamma_{0,2} = 2 * \text{Im}(z_2) = 162^{+38}_{-16} \text{ MeV}$$



Fit to the data leaving the coupling of the two poles free and also width and mass within the error -> Broader $K\pi$ pole and larger amplitude for $\Sigma\pi$

Our Special Resonance

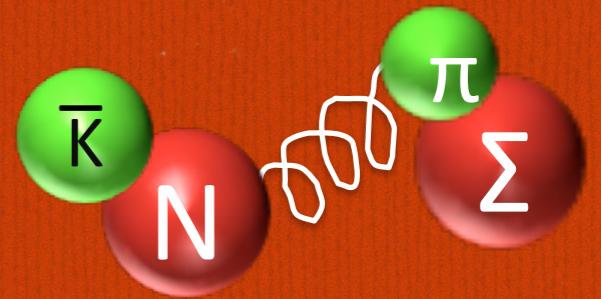
J. Siebenson

$$m_{0,1} = \text{Re}(z_1) = 1424^{+7}_{-23} \text{ MeV}$$

$$\Gamma_{0,1} = 2 * \text{Im}(z_1) = 52^{+6}_{-28} \text{ MeV}$$

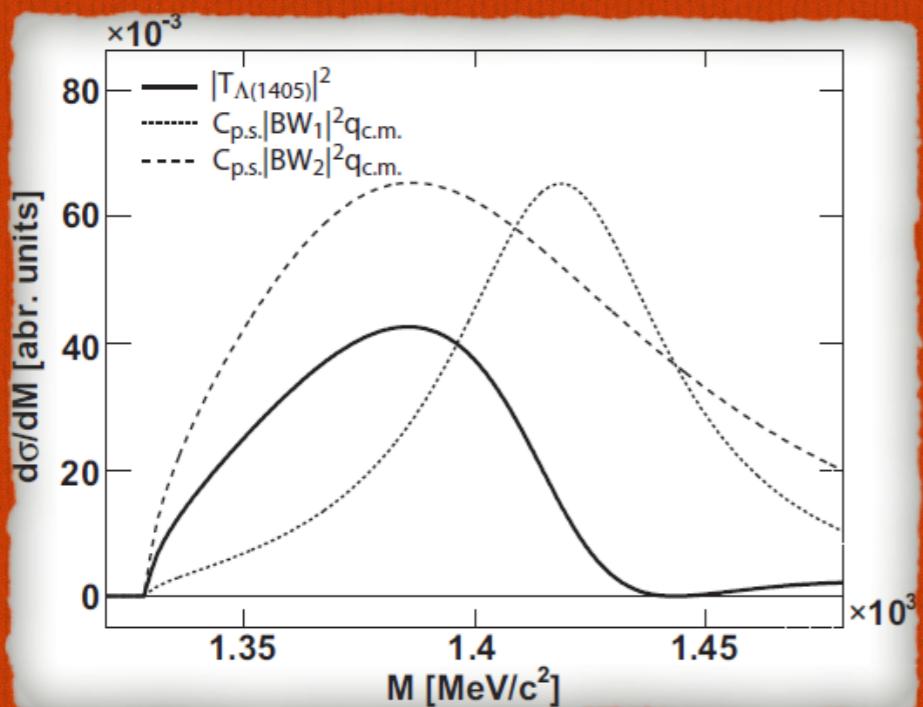
$$m_{0,2} = \text{Re}(z_2) = 1381^{+18}_{-6} \text{ MeV}$$

$$\Gamma_{0,2} = 2 * \text{Im}(z_2) = 162^{+38}_{-16} \text{ MeV}$$



Fit to the data leaving the coupling of the two poles free and also width and mass within the error -> Broader K_p pole and larger amplitude for Σπ

Our Ansatz



Our Special Resonance

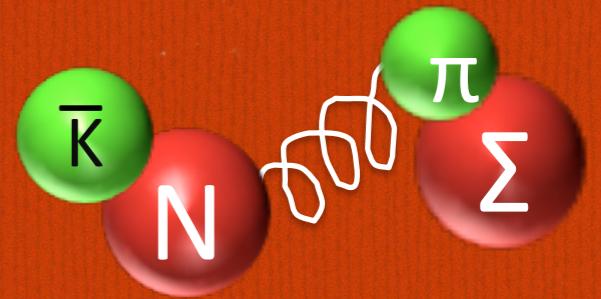
J. Siebenson

$$m_{0,1} = \text{Re}(z_1) = 1424^{+7}_{-23} \text{ MeV}$$

$$\Gamma_{0,1} = 2 * \text{Im}(z_1) = 52^{+6}_{-28} \text{ MeV}$$

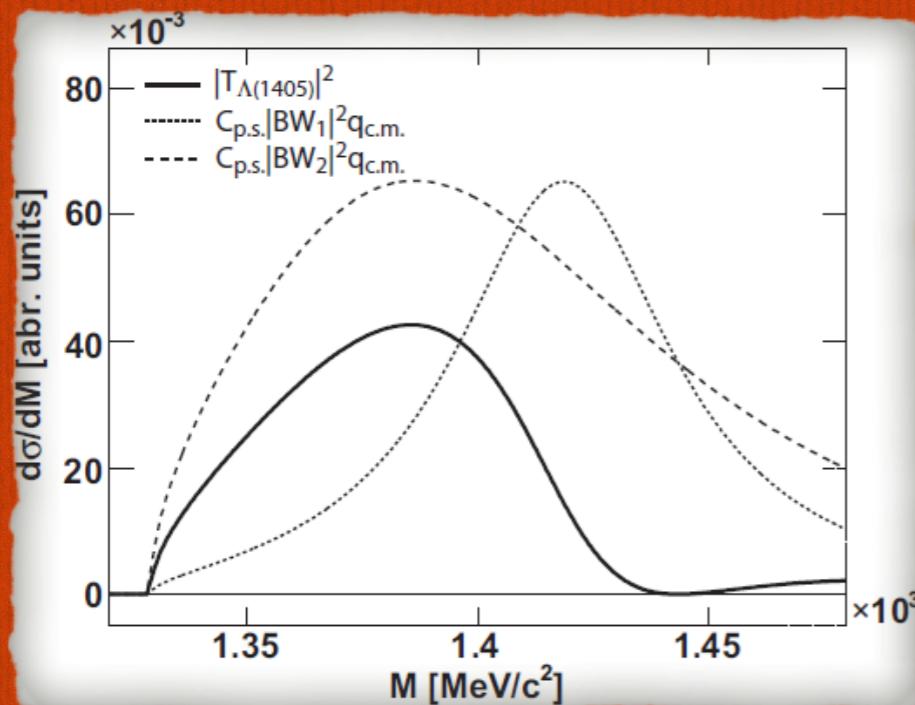
$$m_{0,2} = \text{Re}(z_2) = 1381^{+18}_{-6} \text{ MeV}$$

$$\Gamma_{0,2} = 2 * \text{Im}(z_2) = 162^{+38}_{-16} \text{ MeV}$$

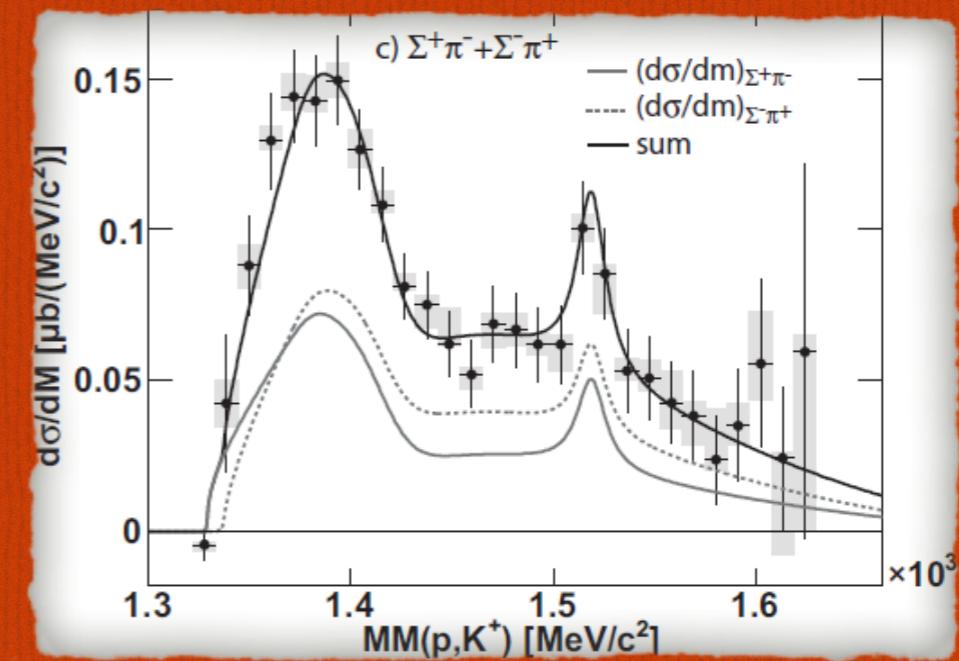


Fit to the data leaving the coupling of the two poles free and also width and mass within the error -> Broader K_p pole and larger amplitude for Σπ

Our Ansatz



J. Siebenson, L. Fabbietti, Phys. Rev. C 88 (2013) 055201.



Our Special Resonance

J. Siebenbon

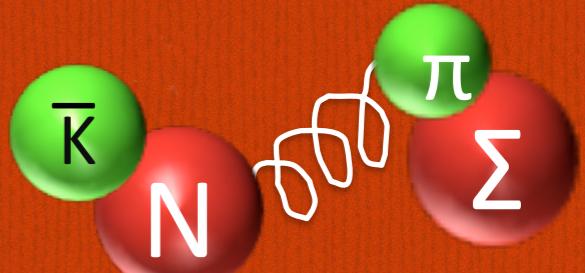
Nucl.Phys. A900 (2013) 51 - 64

$$m_{0,1} = \text{Re}(z_1) = 1428^{+2}_{-1} \text{ MeV}$$

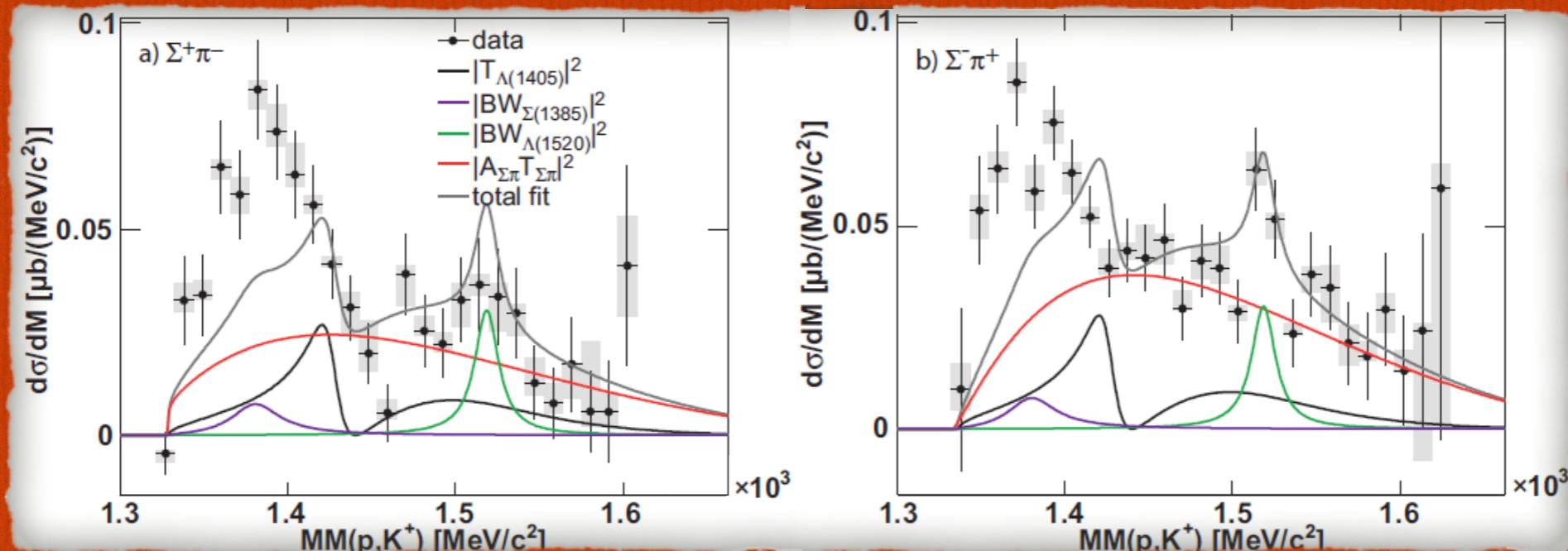
$$\Gamma_{0,1} = 2 * \text{Im}(z_1) = 16^{+4}_{-4} \text{ MeV}$$

$$m_{0,2} = \text{Re}(z_2) = 1497^{+11}_{-7} \text{ MeV}$$

$$\Gamma_{0,2} = 2 * \text{Im}(z_2) = 150^{+18}_{-18} \text{ MeV}$$

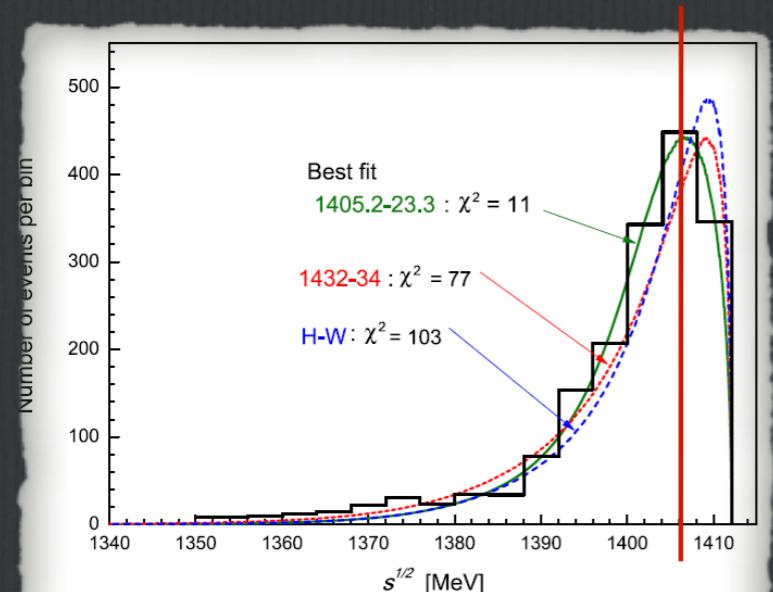


Assuming these two poles the data cannot be reproduced even if advocating interferences with the non resonant part



K⁻ ⁴He

D. Riley et al., Phys. Rev. D 11(1975) 3065
 Esmaili et al., Phys. Lett B686 (2010) 23-28



g. 6. Detailed differences in $M_{\Sigma\pi}$ spectra among the Hyodo-Weise prediction and the present model predictions.

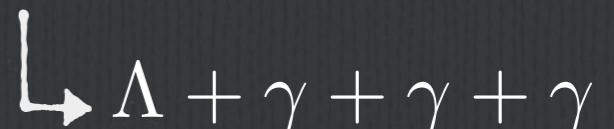
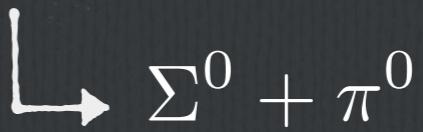
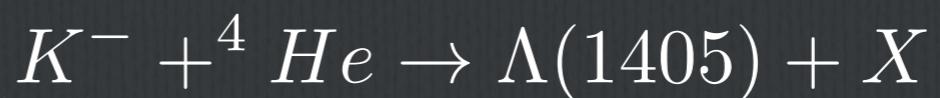
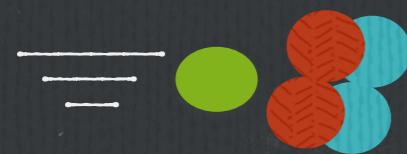
Previous data at rest with $\pi\Sigma$ mass below the kinematical limit

Theoretical interpretation of the $\Lambda(1405)$ as a pure K-p pole

-> Data for in-flight processes have access to the high mass region

$K_{mom} \sim 100 \text{ MeB} \rightarrow M_{\pi\Sigma}$ up to $1425 \text{ MeV}/c^2$

K⁻ ⁴He



3 γ s have to be detected in the calorimeter

K⁻ ⁴He

D. Riley et al., Phys. Rev. D 11(1975) 3065
 Esmaili et al., Phys. Lett B686 (2010) 23-28

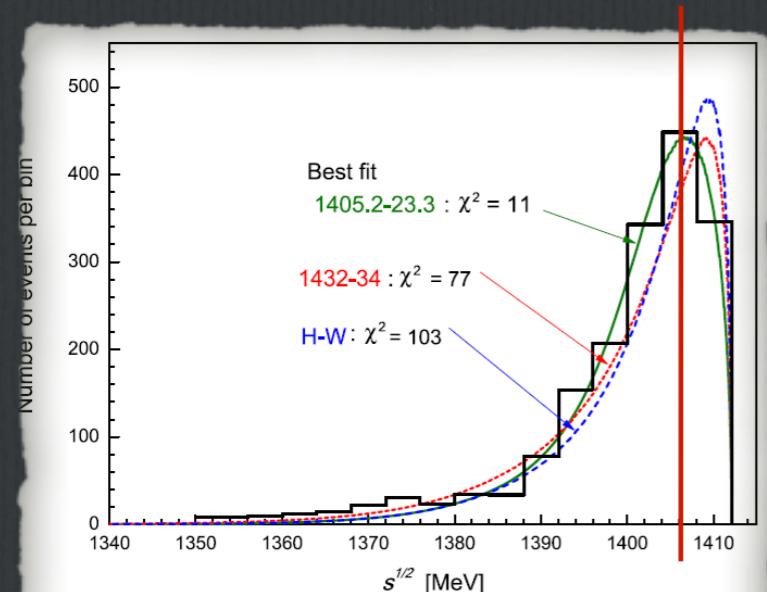


Fig. 6. Detailed differences in $M_{\Sigma\pi}$ spectra among the Hyodo-Weise prediction and the present model predictions.

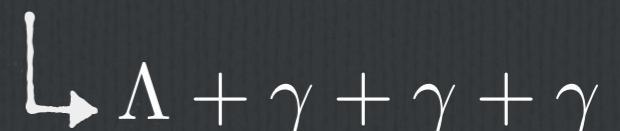
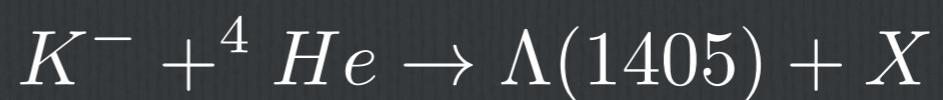
Previous data at rest with $\pi\Sigma$ mass below the kinematical limit

Theoretical interpretation of the $\Lambda(1405)$ as a pure K-p pole

-> Data for in-flight processes have access to the high mass region

$K_{mom} \sim 100 \text{ MeB} \rightarrow M_{\pi\Sigma}$ up to $1425 \text{ MeV}/c^2$

K⁻ ⁴He



3 γ s have to be detected in the calorimeter

K⁻ ${}^4\text{He}$

D. Riley et al., Phys. Rev. D 11(1975) 3065
 Esmaili et al., Phys. Lett B686 (2010) 23-28

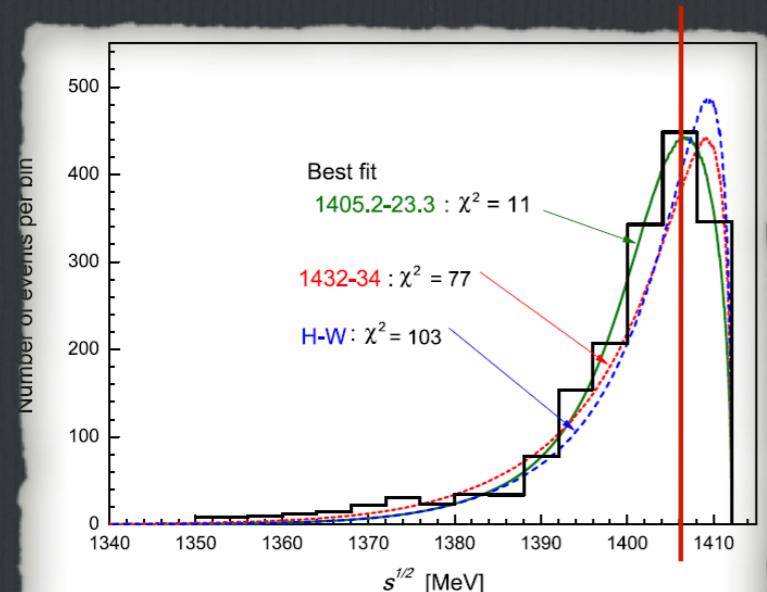


Fig. 6. Detailed differences in $M_{\Sigma\pi}$ spectra among the Hyodo-Weise prediction and the present model predictions.

Previous data at rest with $\pi\Sigma$ mass below the kinematical limit

Theoretical interpretation of the $\Lambda(1405)$ as a pure K-p pole

-> Data for in-flight processes have access to the high mass region

$K_{\text{mom}} \sim 100 \text{ MeB} \rightarrow M_{\pi\Sigma}$ up to $1425 \text{ MeV}/c^2$

K⁻ ${}^4\text{He}$



$\Lambda(1405)$

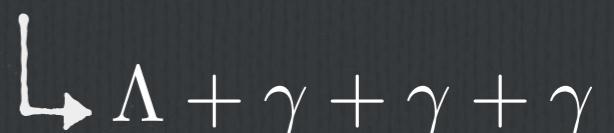
Σ^0

π^0

X



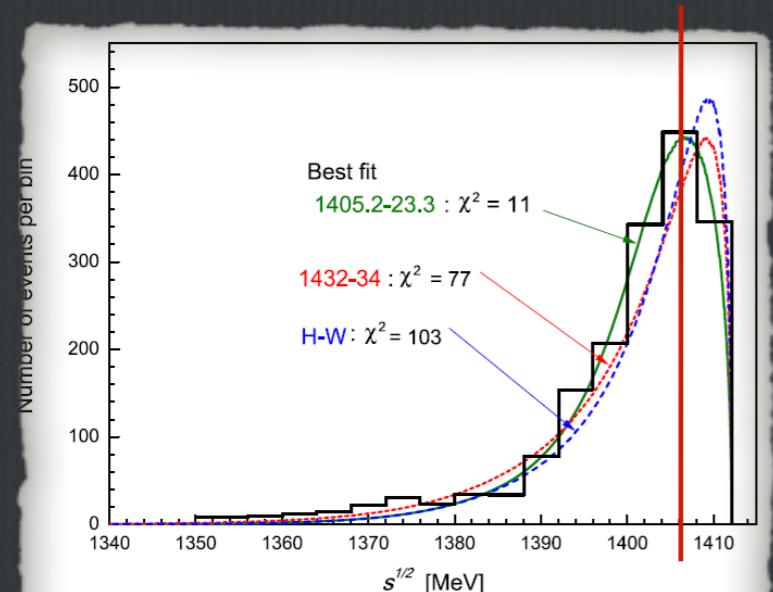
$K^- + {}^4\text{He} \rightarrow \Lambda(1405) + X$



3 γ s have to be detected in the calorimeter

K⁻ ⁴He

D. Riley et al., Phys. Rev. D 11(1975) 3065
 Esmaili et al., Phys. Lett B686 (2010) 23-28



g. 6. Detailed differences in $M_{\Sigma\pi}$ spectra among the Hyodo-Weise prediction and the present model predictions.

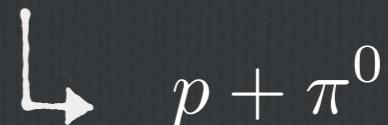
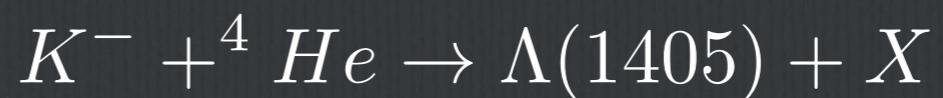
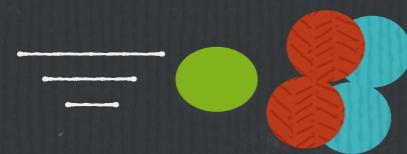
Previous data at rest with $\pi\Sigma$ mass below the kinematical limit

Theoretical interpretation of the $\Lambda(1405)$ as a pure K-p pole

-> Data for in-flight processes have access to the high mass region

$K_{mom} \sim 100 \text{ MeB} \rightarrow M_{\pi\Sigma}$ up to $1425 \text{ MeV}/c^2$

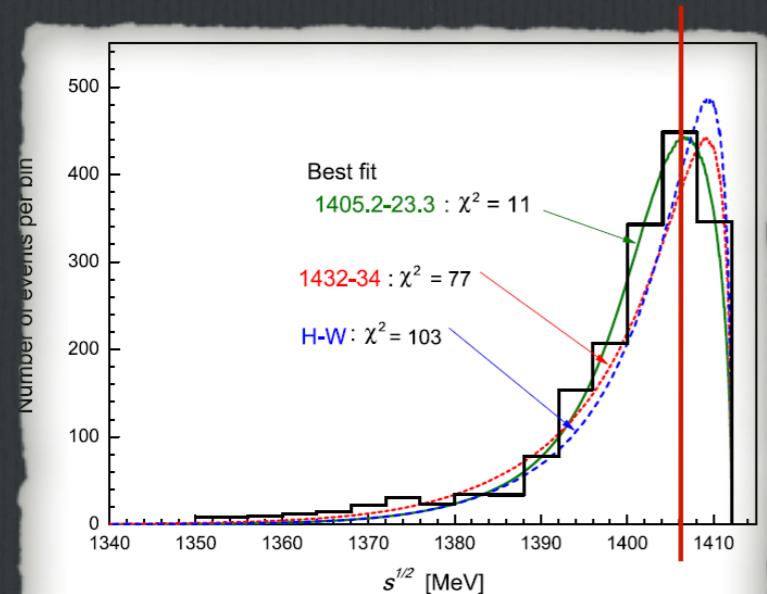
K⁻ ⁴He



2 γ s have to be detected in the calorimeter

K⁻ ⁴He

D. Riley et al., Phys. Rev. D 11(1975) 3065
 Esmaili et al., Phys. Lett B686 (2010) 23-28



g. 6. Detailed differences in $M_{\Sigma\pi}$ spectra among the Hyodo-Weise prediction and the present model predictions.

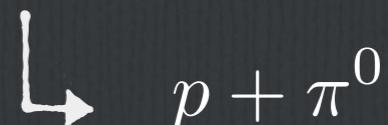
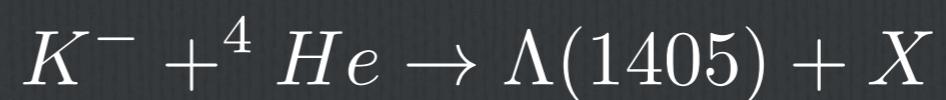
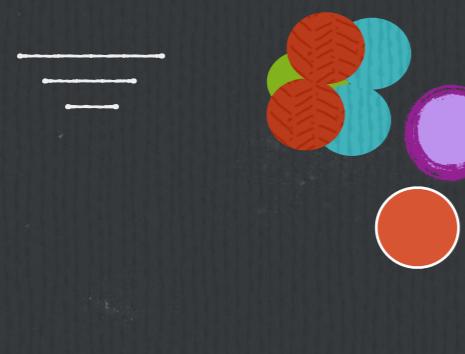
Previous data at rest with $\pi\Sigma$ mass below the kinematical limit

Theoretical interpretation of the $\Lambda(1405)$ as a pure K-p pole

-> Data for in-flight processes have access to the high mass region

$K_{mom} \sim 100 \text{ MeB} \rightarrow M_{\pi\Sigma}$ up to $1425 \text{ MeV}/c^2$

K⁻ ⁴He



2 γ s have to be detected in the calorimeter

K⁻ ${}^4\text{He}$

D. Riley et al., Phys. Rev. D 11(1975) 3065
 Esmaili et al., Phys. Lett B686 (2010) 23-28

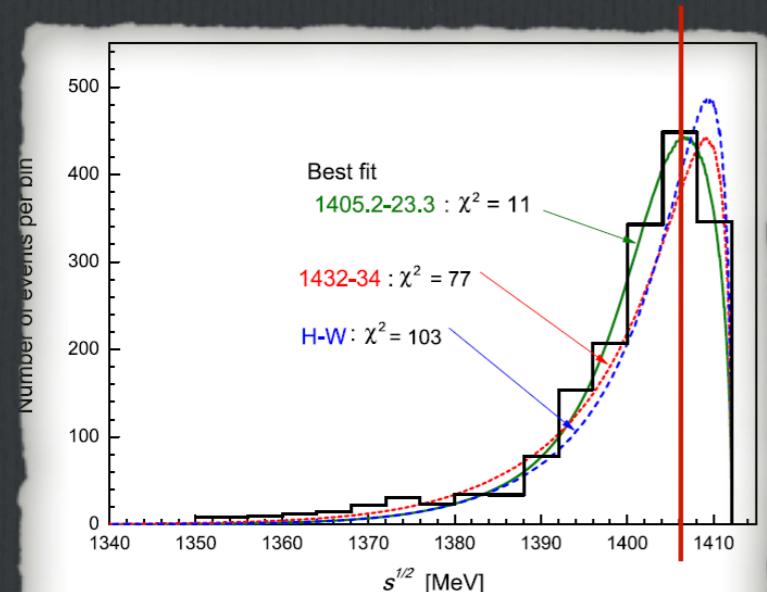


Fig. 6. Detailed differences in $M_{\Sigma\pi}$ spectra among the Hyodo-Weise prediction and the present model predictions.

Previous data at rest with $\pi\Sigma$ mass below the kinematical limit

Theoretical interpretation of the $\Lambda(1405)$ as a pure K-p pole

-> Data for in-flight processes have access to the high mass region

$K_{\text{mom}} \sim 100 \text{ MeB} \rightarrow M_{\pi\Sigma}$ up to $1425 \text{ MeV}/c^2$

K⁻ ${}^4\text{He}$



$\Lambda(1405)$

Σ^+

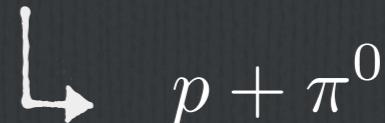
π^-

π^0

X



π^0



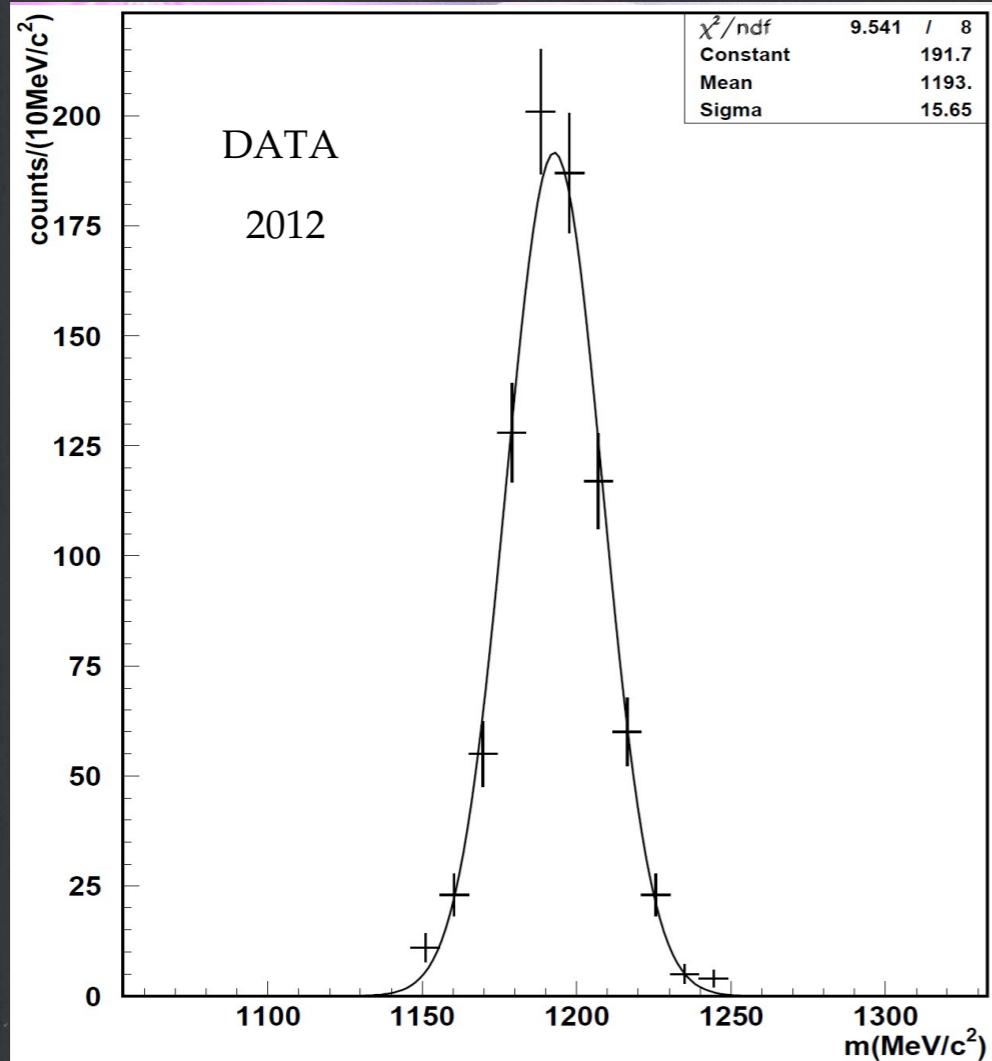
2 γ s have to be detected in the calorimeter

Particle Identification in KLOE

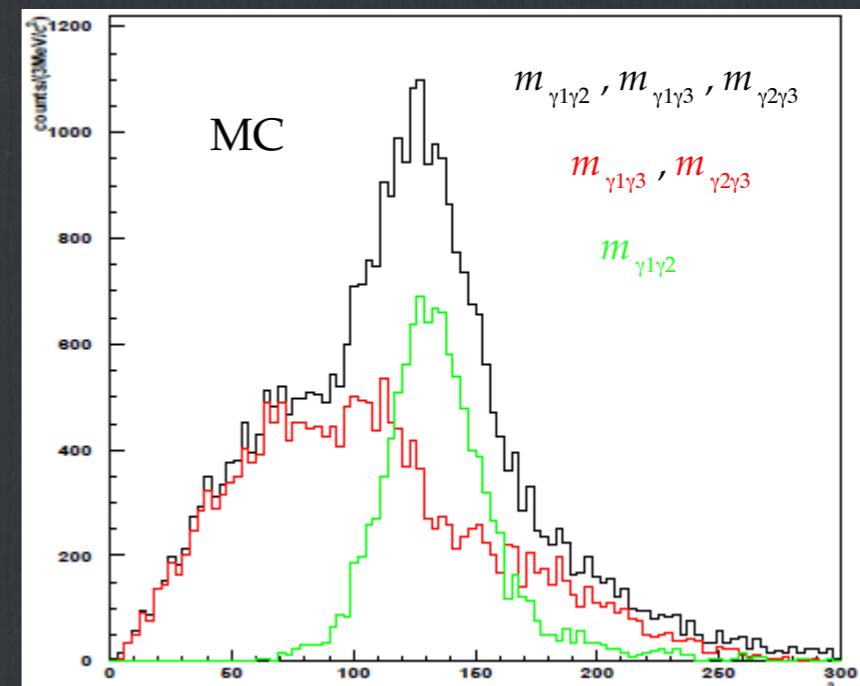
C. Curceanu, L. Fabbietti, K. Piscicchia, A. Scordo, I. Tucakovic, O. Vasquez-Doce + KLOE

After the Λ selection..

Σ^0



π^0



$m_{\Lambda\gamma}$

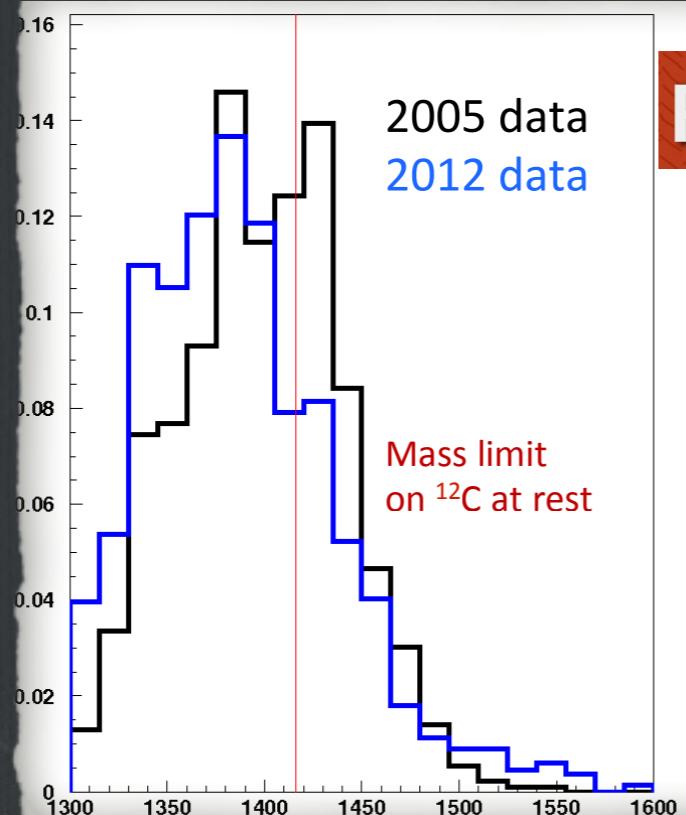
80% efficiency for $\Sigma^0\pi^0$

$m_{\gamma\gamma}$

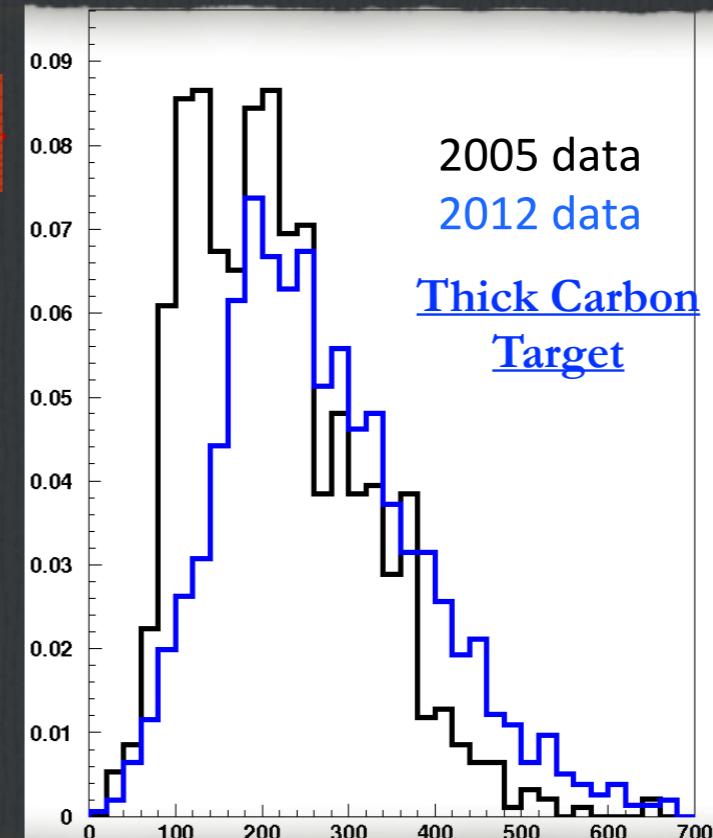
$\Lambda(1405) \rightarrow \pi^0 \Sigma^0$

At rest Data
 In-flight + at Rest Data
 In-flight: larger masses are accessible

a.u.



Preliminary



$$m_{\pi^0 \Sigma^0} (\text{MeV}/c^2)$$

Mass resolution $\sigma_m = 32 \text{ MeV}/c^2$



$p_{\pi^0 \Sigma^0} (\text{MeV}/c)$
 Momentum resolution $\sigma_p = 20 \text{ MeV}/c$

$\Lambda(1405) \rightarrow \pi^0 \Sigma^0$

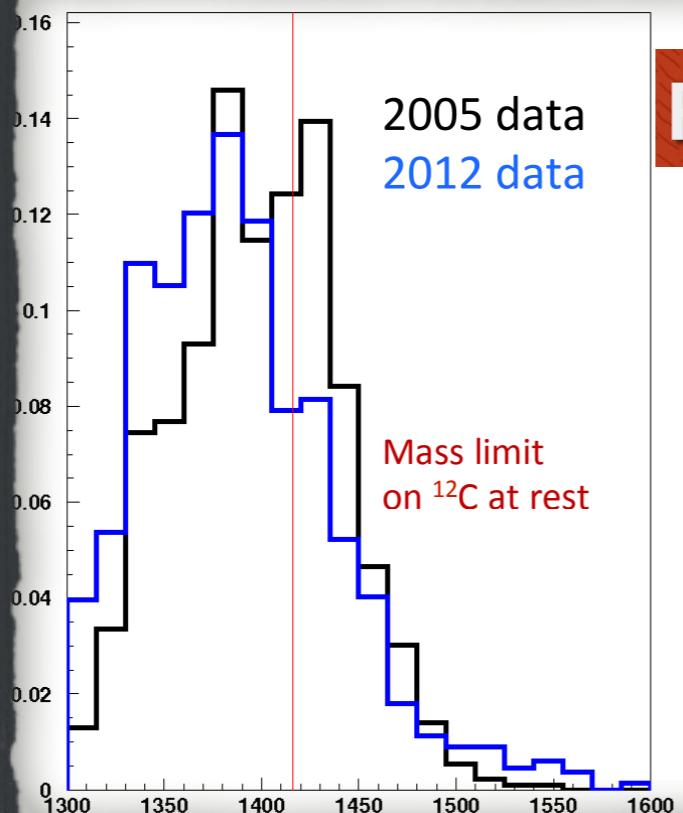
At rest Data

In-flight + at Rest Data

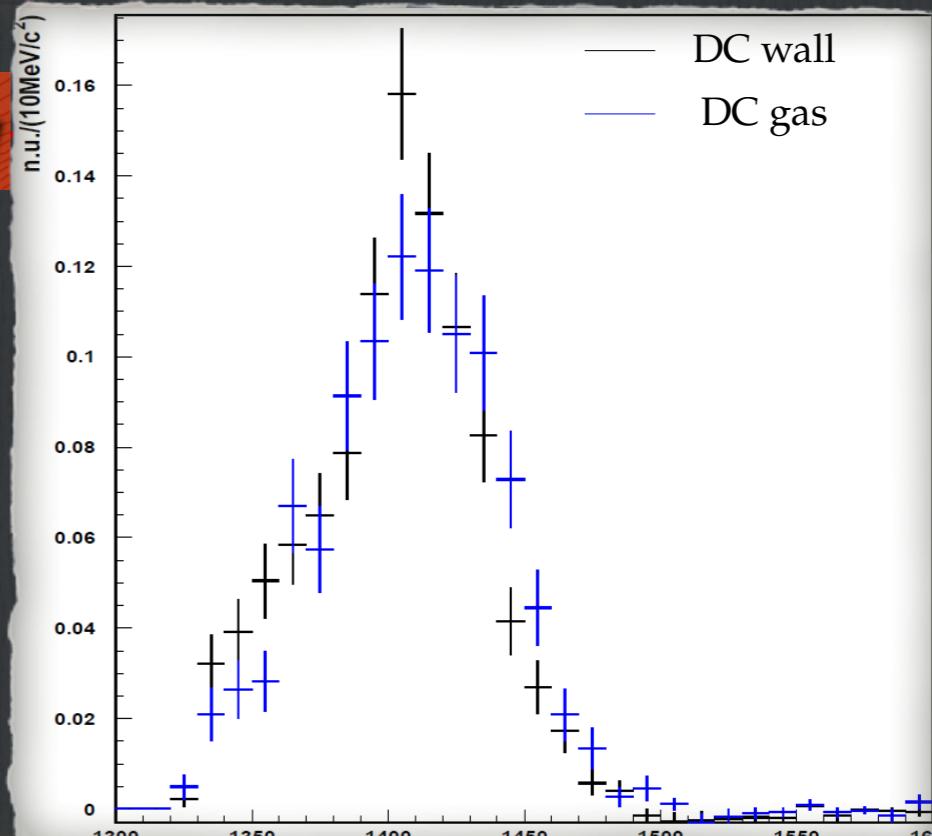
In-flight: larger masses are accessible

Nominal mass for Σ^0 and π^0

a.u.



Preliminary



$$m_{\pi^0 \Sigma^0} (\text{MeV}/c^2)$$

Mass resolution $\sigma_m = 32 \text{ MeV}/c^2$



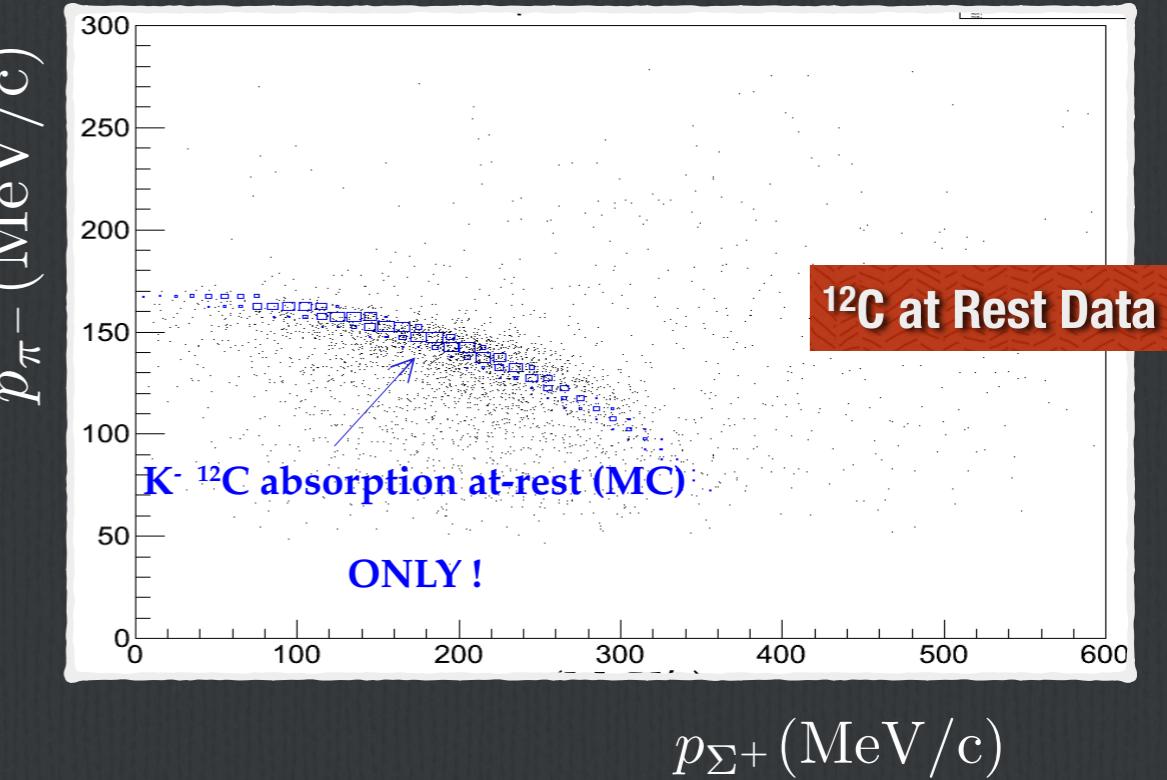
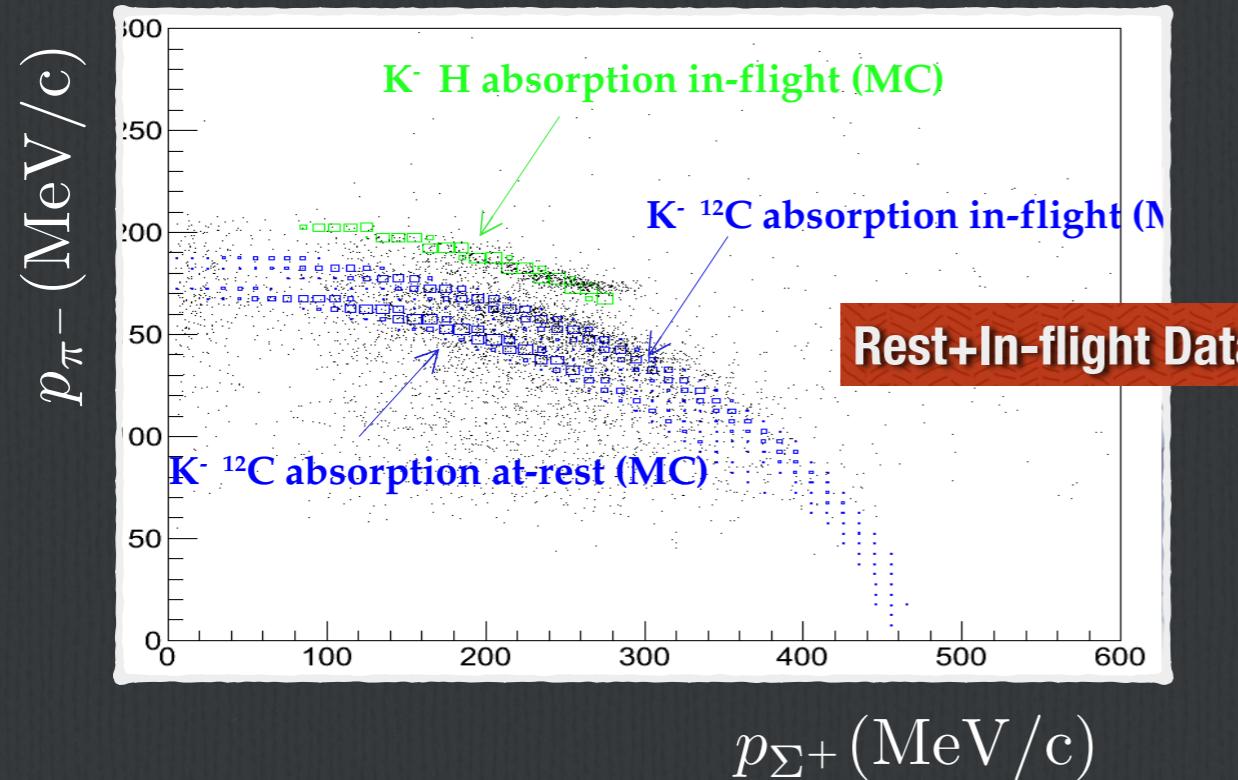
Momentum resolution $\sigma_p = 20 \text{ MeV}/c$

Fit Components

- Resonant component $K^- + C/{}^4He$ at-rest/in-flight (M, Γ) = (1390- 1430, 5-52 MeV/c²)
- Non resonant $\Sigma^0\pi^0$ $K^- + H$ production in-flight
- Non resonant $\Sigma^0\pi^0$ $K^- + C/{}^4H$ production at-rest/in-flight
- $\Lambda\pi^0$
- misidentification background

$\Sigma^+ \pi^-$ Selection: Hydrogen Contamination

$$\Sigma^+ \pi^- \rightarrow (p\pi^0)\pi^-$$

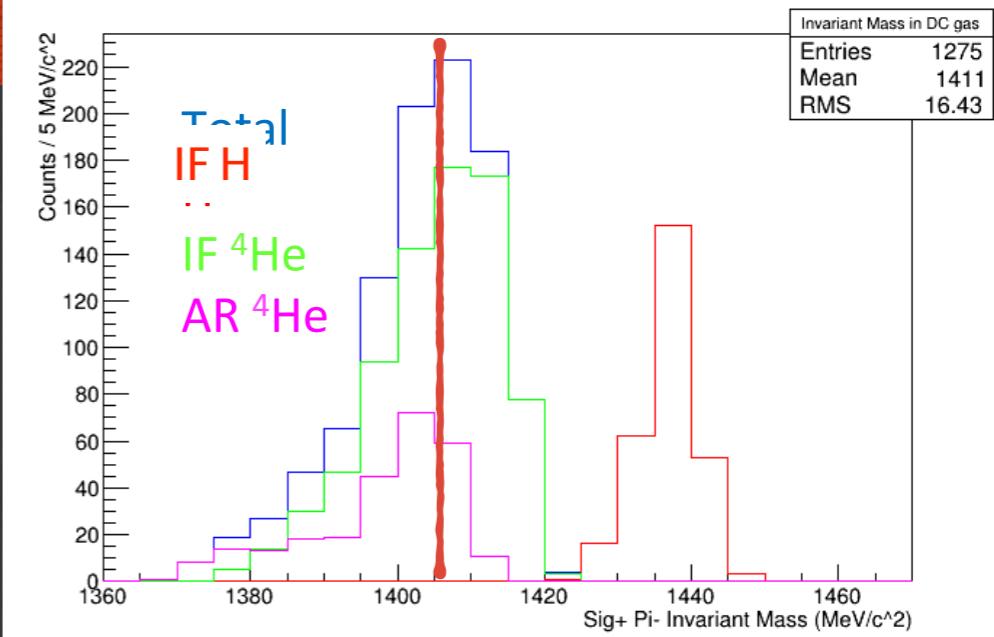
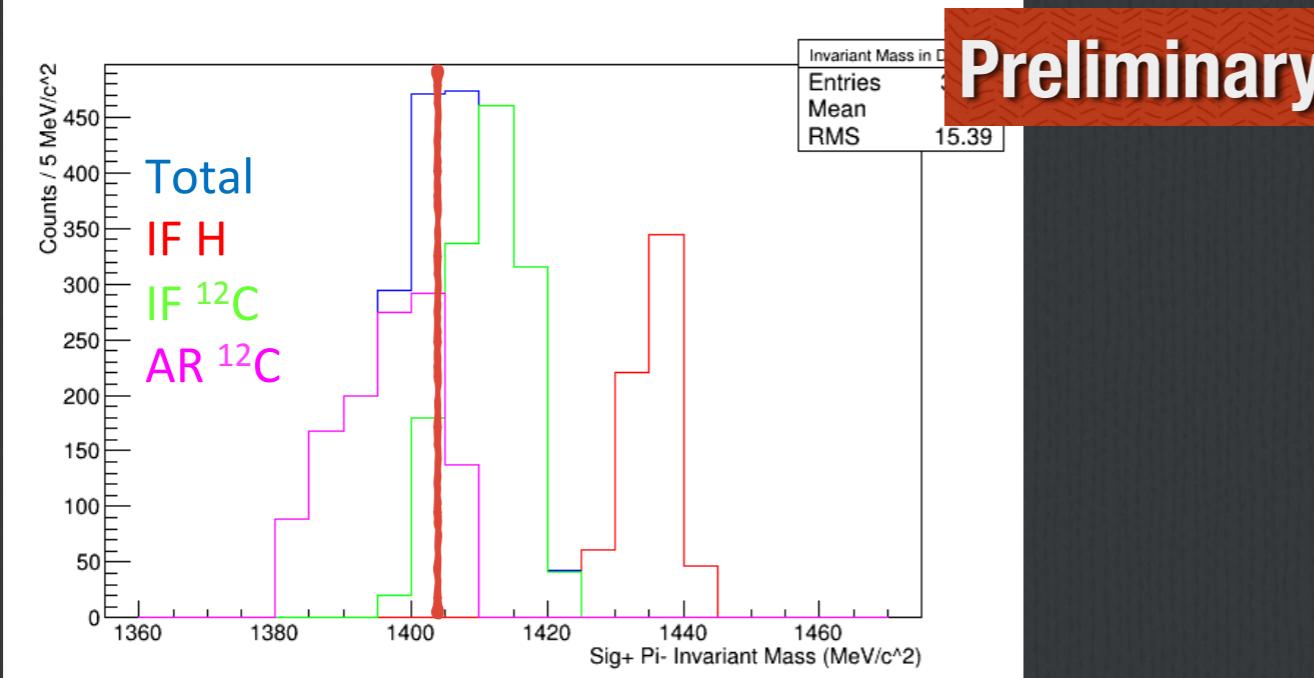


Momentum resolution $\pi^- < 1$ MeV

This channel offers a bench mark for the $K^- + H$ contribution

Shifted-Resonance or Quasi-free?

Preliminary



$$\text{IF/AR} = 1.16 \pm 0.05$$

$$\text{IF/AR} = 2.9 \pm 0.3$$

The in-flight $\Sigma^0\pi^0$ invariant mass distribution reaches out to 1420 MeV/c²
 Is that the $\Lambda(1405)$ or a quasi-free process?

Conclusions and Summary

$\Lambda(1405)$ measured in p+p at 3.5 GeV with HADES in the charged decay channel. Shift to lower masses

Theoretical calculations unable to explain this data
Interferences? Different coupling of the two poles?

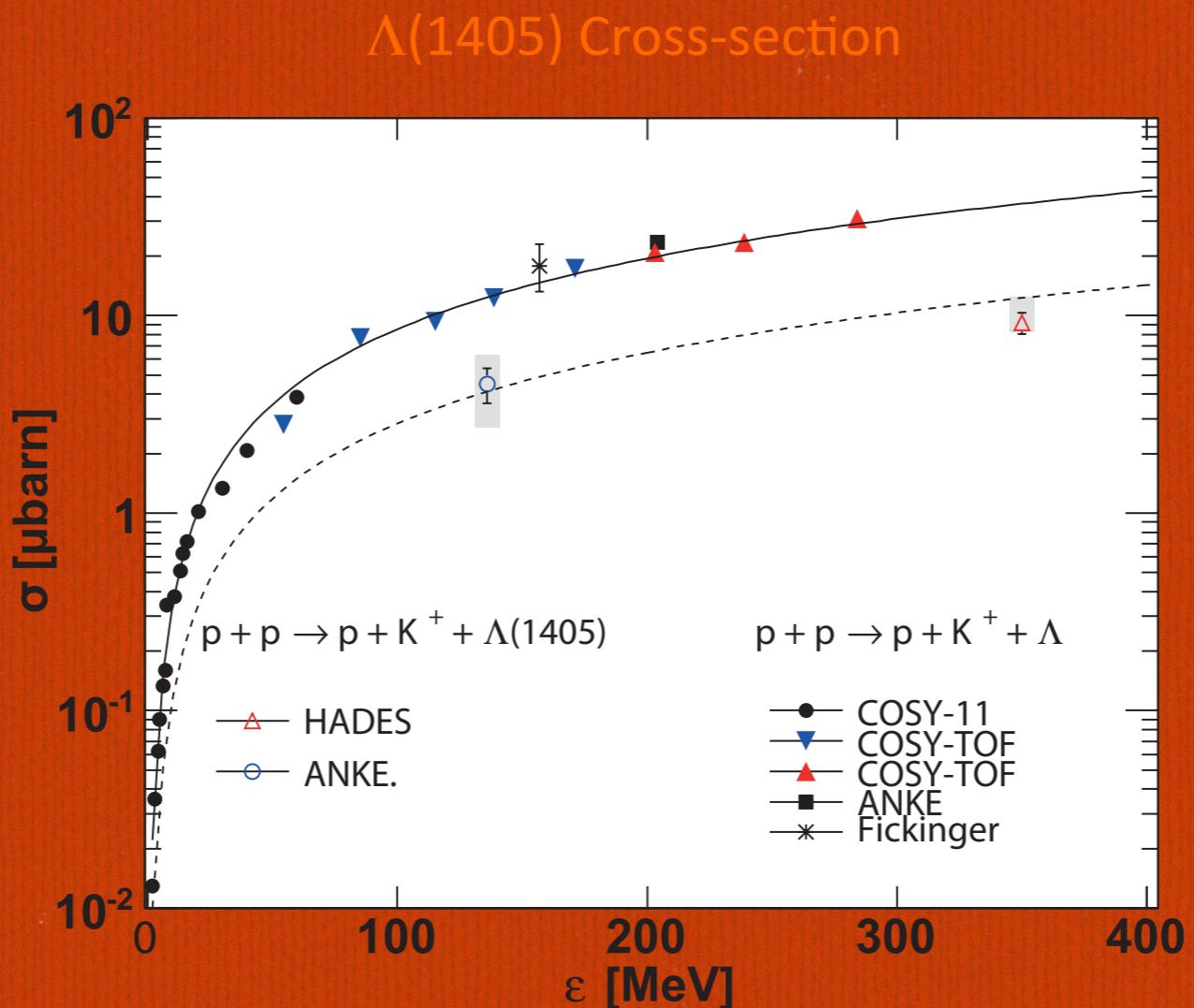
$\Lambda(1405)$ measured in K- + 4He/C with KLOE in the neutral and 1 charged channel

at-rest and in-flight component

Different contributions not disentangled yet
In-flight component shifted?? or Quasi-free?

Calculations are needed

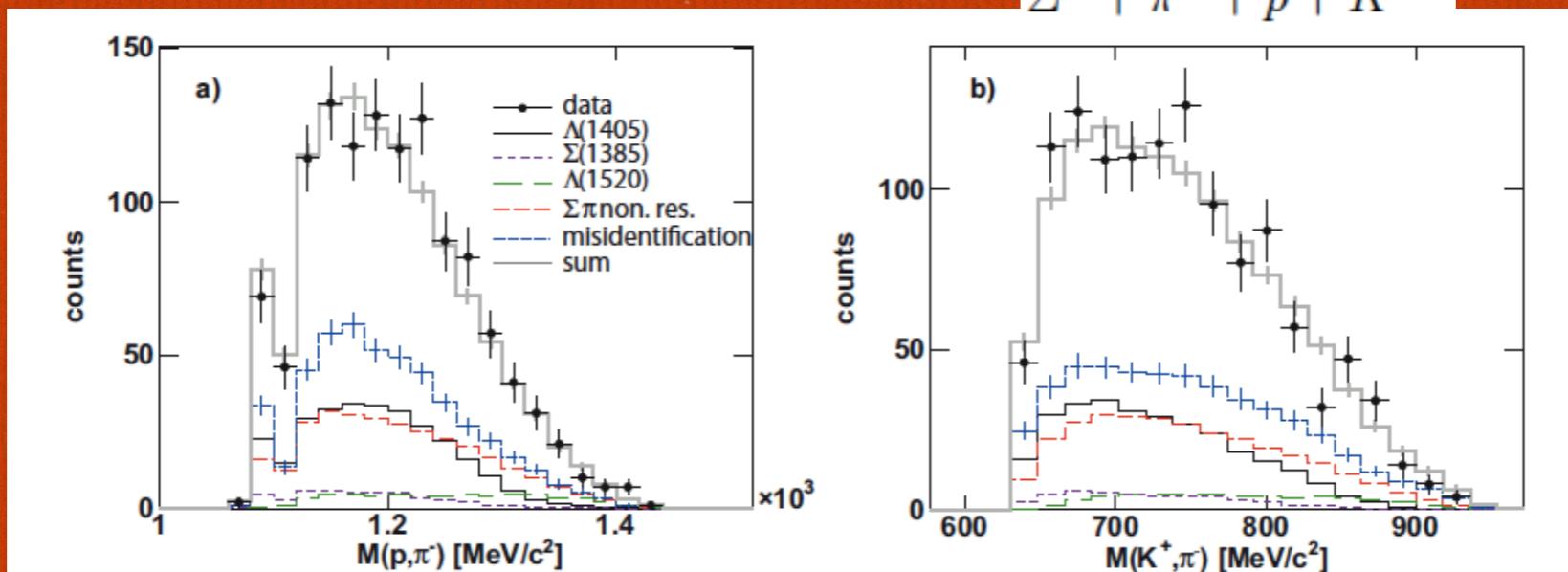
Total Cross-section



Non resonant Contribution

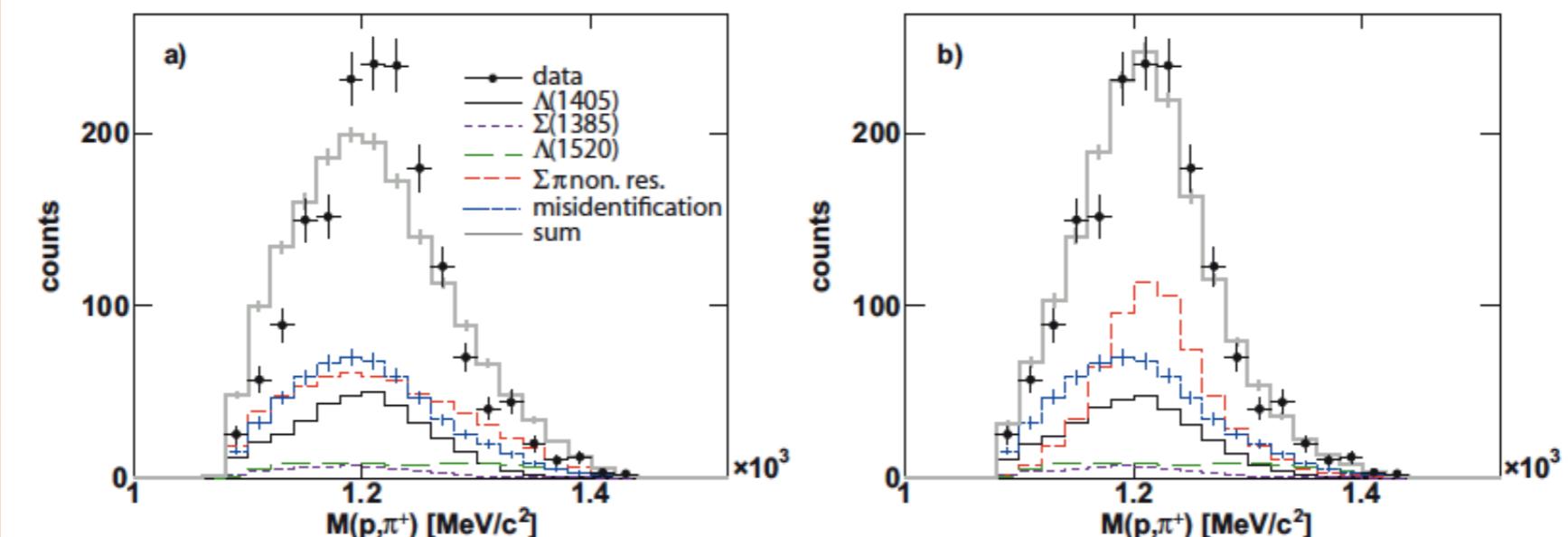
Δ^{++}

$\Sigma^+ + \pi^- + p + K^+$



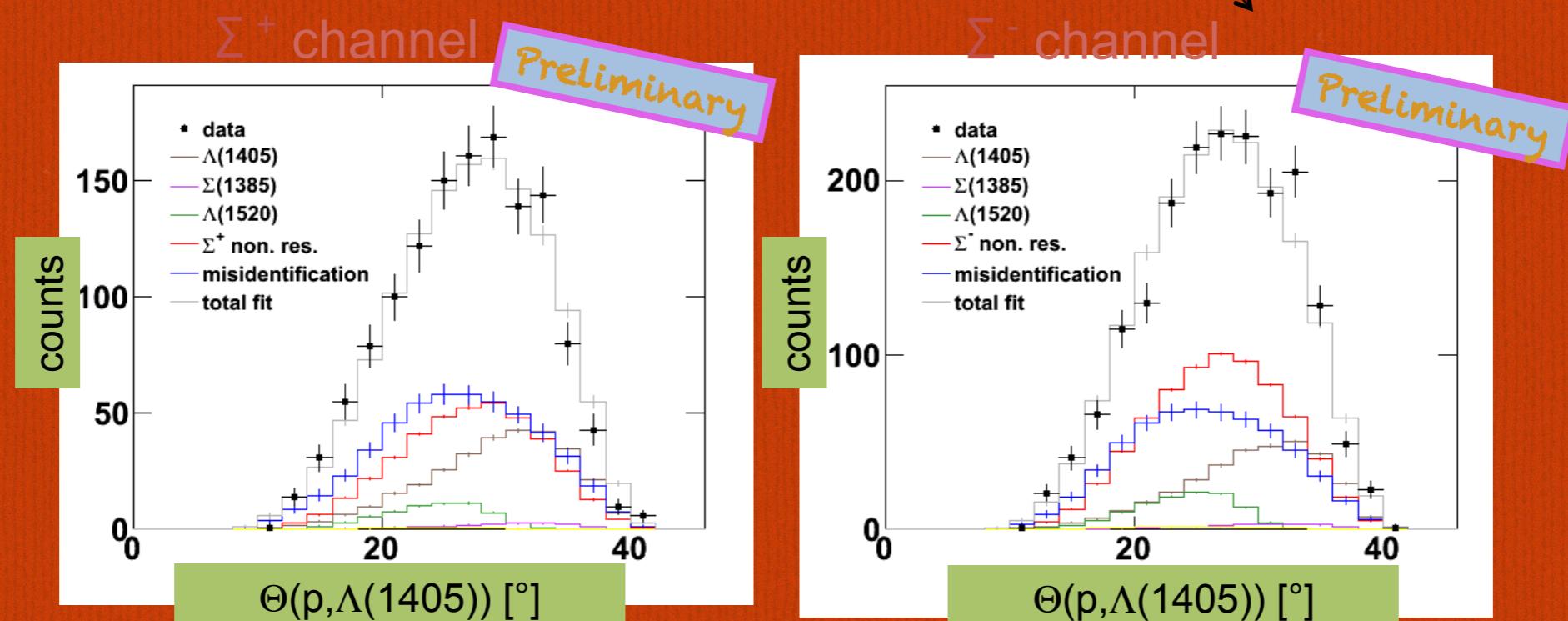
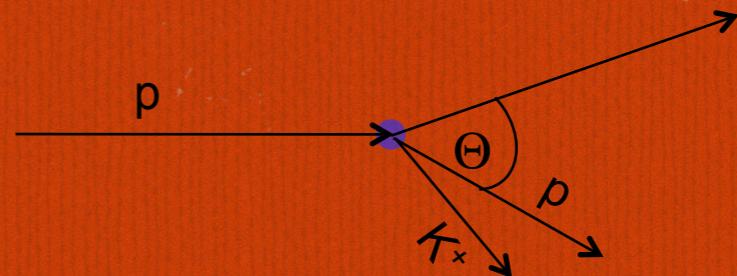
$\Sigma^- + \pi^+ + p + K^+$

$\Sigma^- + K^+ + \Delta^{++}(1232) \rightarrow \Sigma^- + K^+ + (p\pi^+)$



$\Lambda(1405)$ -p angle

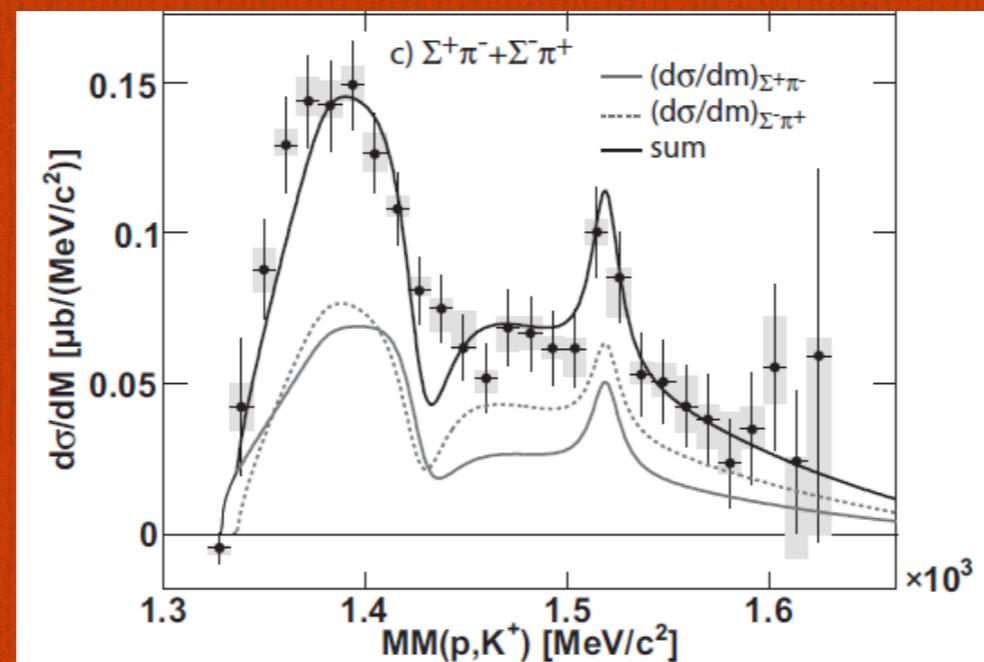
Opening angle in the laboratory system:



No significant deviation from phase space

Interference with non resonant

2. Double pole nature of $\Lambda(1405)$ and include interference effects with non-resonant channels



$A_{\Lambda(1405)}$	$A_{\Sigma^+\pi^-}$	$A_{\Sigma^-\pi^+}$	α	β
1.05	0.94	1.05	68°	111°

Interference with non resonant

- * The $p+p \rightarrow \Sigma^- + \Delta^{++}(1232) + K^+$ reaction is a very probable candidate for the non-resonant part of the $\Sigma^- + \pi^+ + p + K^+$ spectrum and has different quantum numbers than the $\Lambda(1405)$ contribution. This does not lead to interferences
- * Why should both charged decay channels interfere in the same way?
- * Why we dont see any shift in $p+p \rightarrow \Sigma(1385)^+ + K^+ + n$ or for the $\Lambda(1520)$
- * $p+p(@3.5\text{GeV}) \rightarrow \pi^+ + p + p/\pi^- + \pi^+ + p + p$ Incoherent analysis of 14 N* with angular distribution reproduce perfectly the data and is consistent with dilepton yield. This looks like small room for interferences..