

First Spectroscopy of the Hyperfine Interval of Positronium Using Millimeter Waves

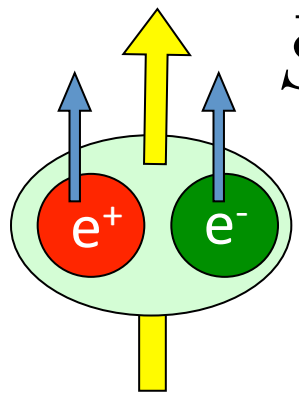
Akira Miyazaki (The University of Tokyo→CERN)

T. Yamazaki, T. Suehara, T. Namba
S. Asai, T. Kobayashi, H. Saito
(The University of Tokyo)

Y. Tatematsu, I. Ogawa, T. Idehara
(Fukui University)

INTRODUCTION

Spin Eigenstates of ground-state Positronium (Ps)



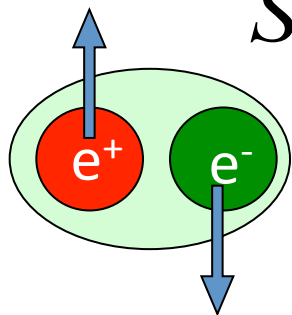
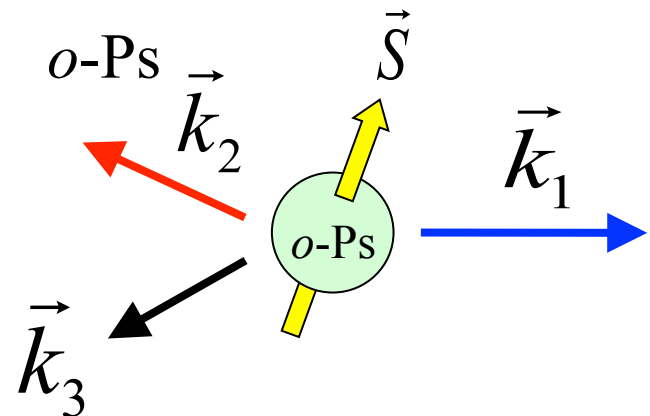
$\vec{S} = 1$ (Triplet)

ortho-positronium (*o-Ps*)

three photon decay

$$o\text{-Ps} \rightarrow 3\gamma \text{ (, } 5\gamma, \dots)$$

lifetime = 142 ns



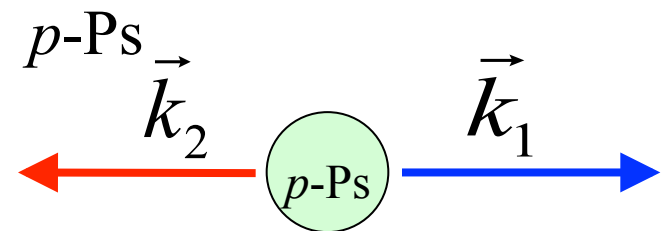
$\vec{S} = 0$ (Singlet)

para-positronium (*p-Ps*)

two photon decay

$$p\text{-Ps} \rightarrow 2\gamma \text{ (, } 4\gamma, \dots)$$

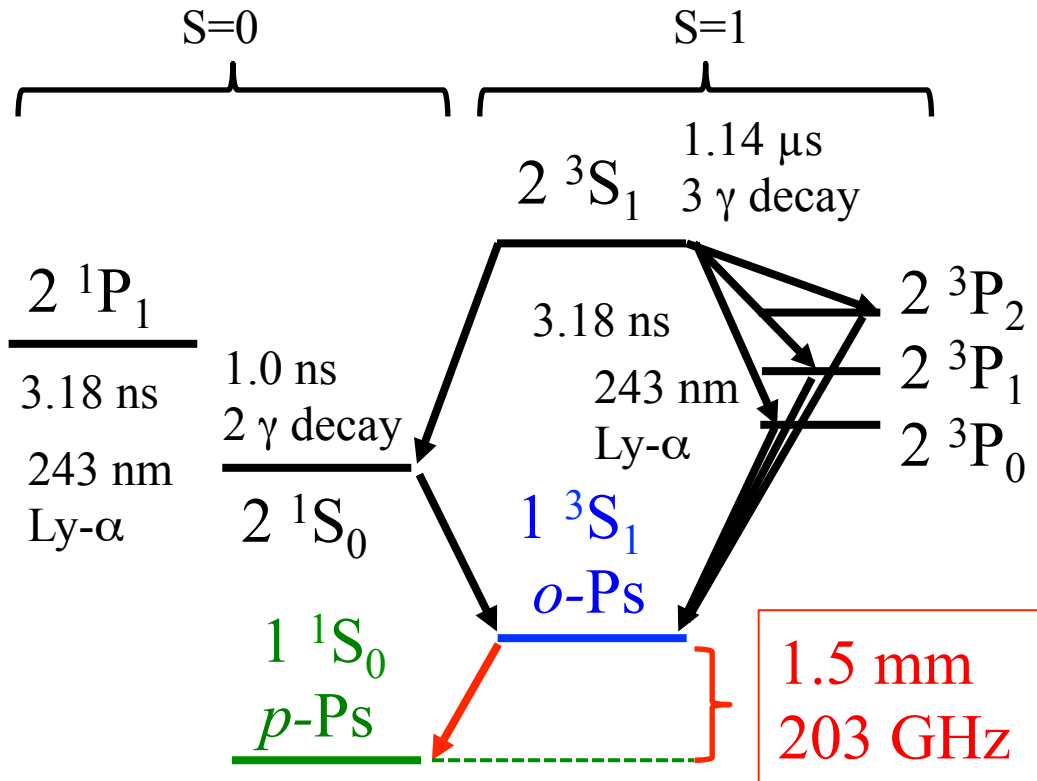
lifetime = 125 ps



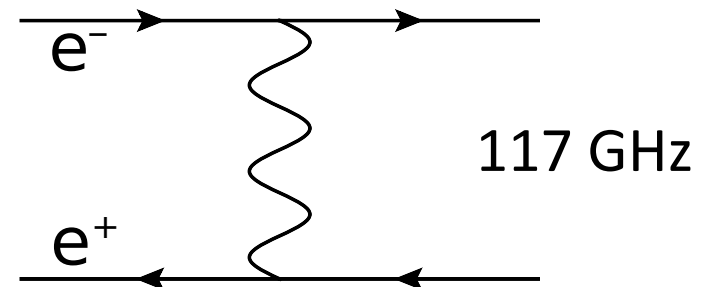
back-to-back 511 keV ³

Hyperfine structure of Ps (Δ_{HFS})

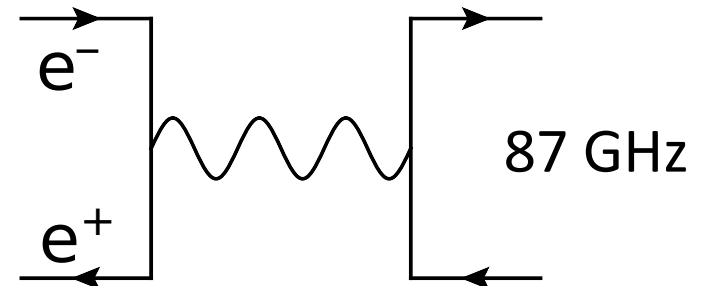
Energy Level Diagram of Ps



1. Spin-spin interaction



2. Vacuum oscillation



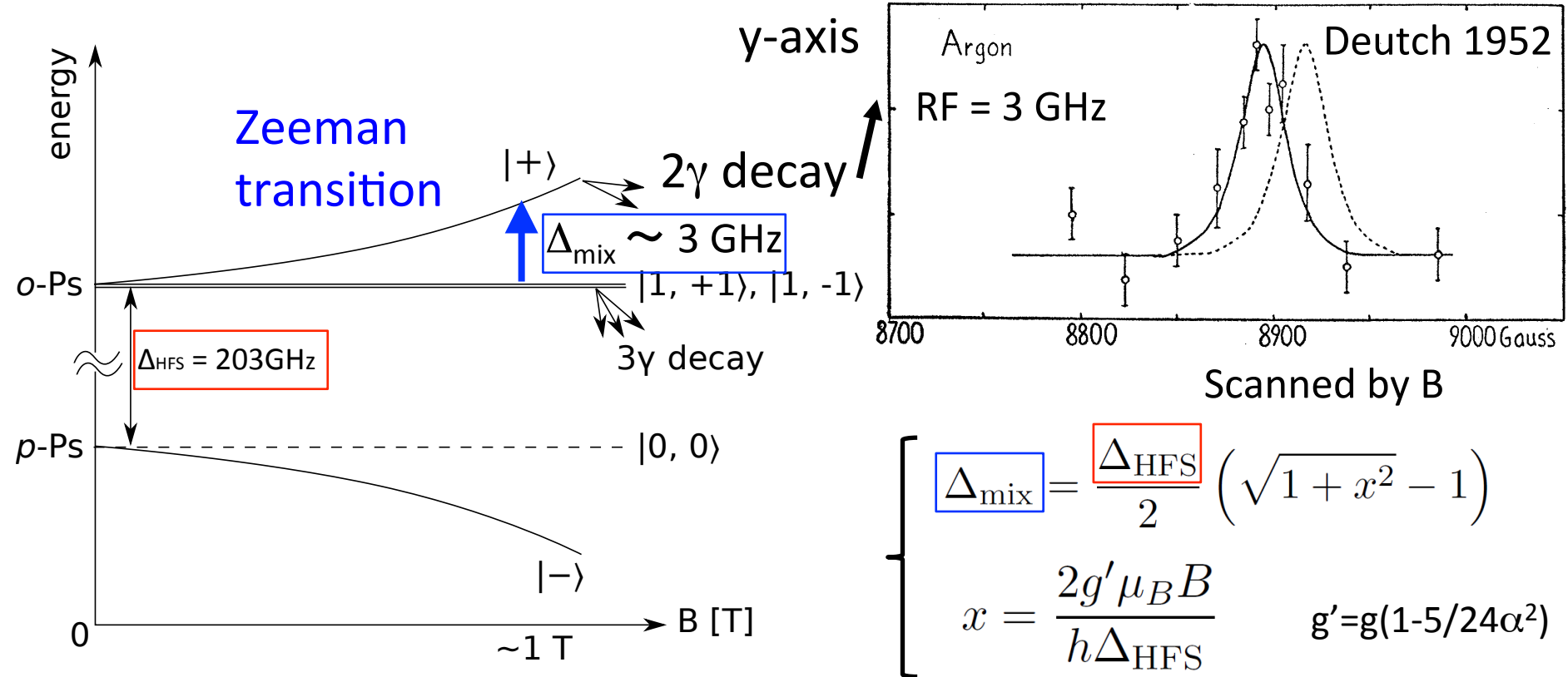
The spin-spin interaction and vacuum oscillation raises o -Ps energy from p -Ps one by $\Delta_{\text{HFS}} = 203 \text{ GHz}$ (\gg hydrogen's HFS 1.4GHz)

No established method to treat millimeter waves

→ How was Δ_{HFS} determined so far?

Solution in literature

Measure the Zeeman shifted levels $\rightarrow \Delta_{\text{HFS}}$ was then calculated

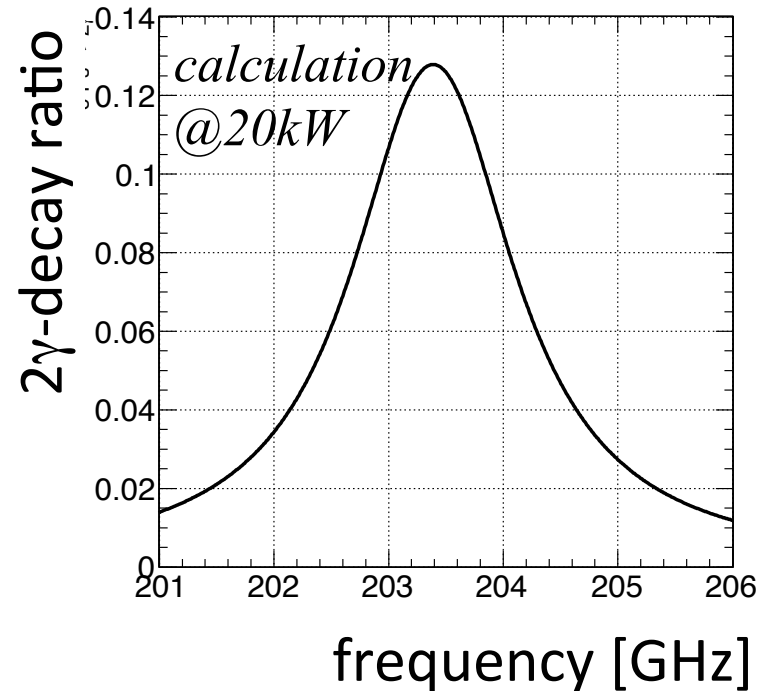
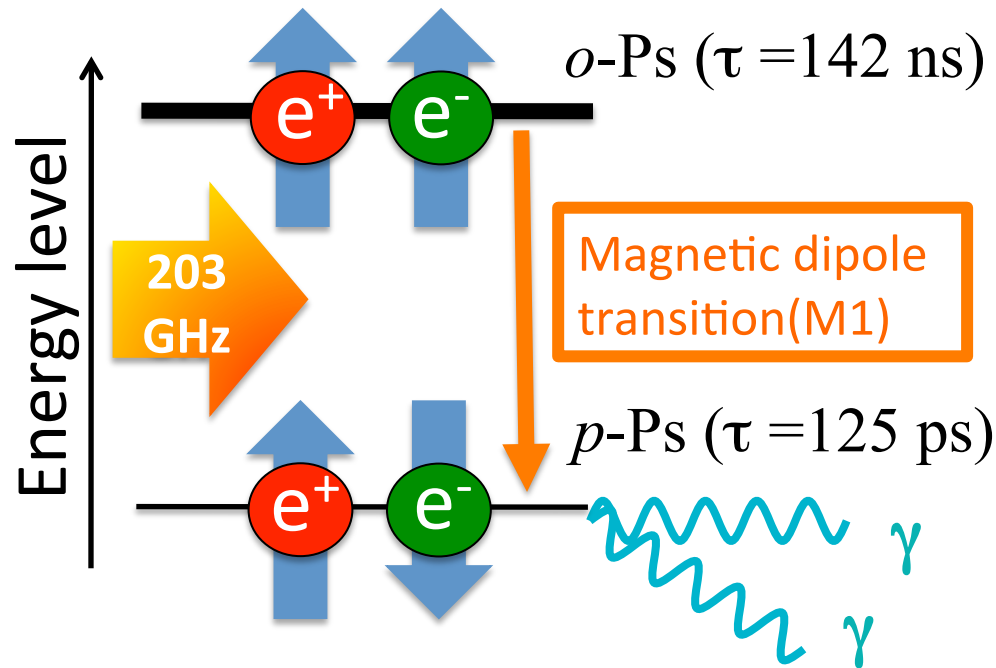


In such a way, Δ_{HFS} is determined with accuracy of about a few ppm.

However, Δ_{HFS} is originally derived in a fundamental way for **free Ps**

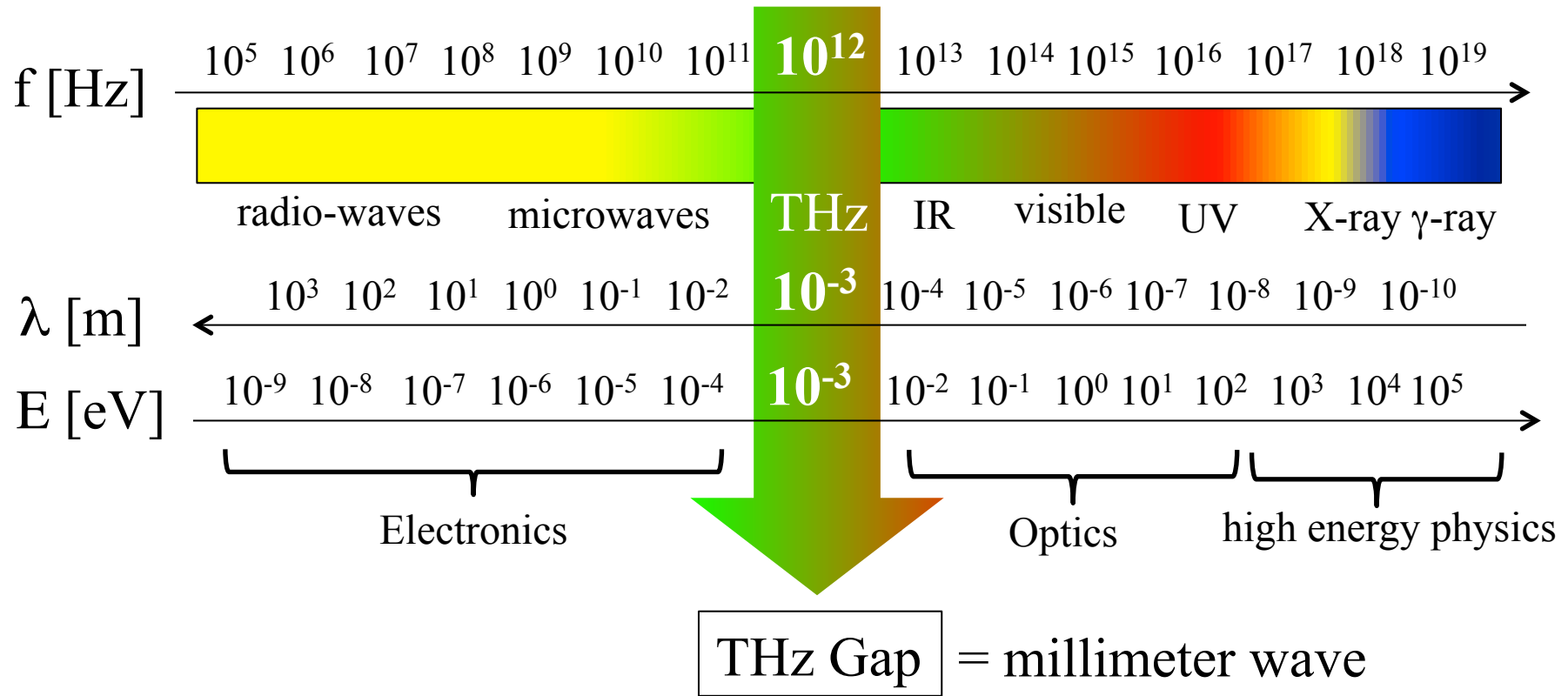
\rightarrow Can we directly measure Δ_{HFS} with leading-edge technology today?

Simple method, but...



- 203GHz radiation induces the transition from *o*-Ps to *p*-Ps (M1)
- The transiting *p*-Ps decays into 2 γ rays (511 keV) in 125 ps
- The 2 γ -decay ratio becomes the Breit-Wigner curve (mean = Δ_{HFS})
- High power of over 10 kW is required because of short lifetime of Ps
- Frequency should be scanned from 201 GHz to 206 GHz
- Are they difficult to achieve today??

Yes! Millimeter waves are still very difficult

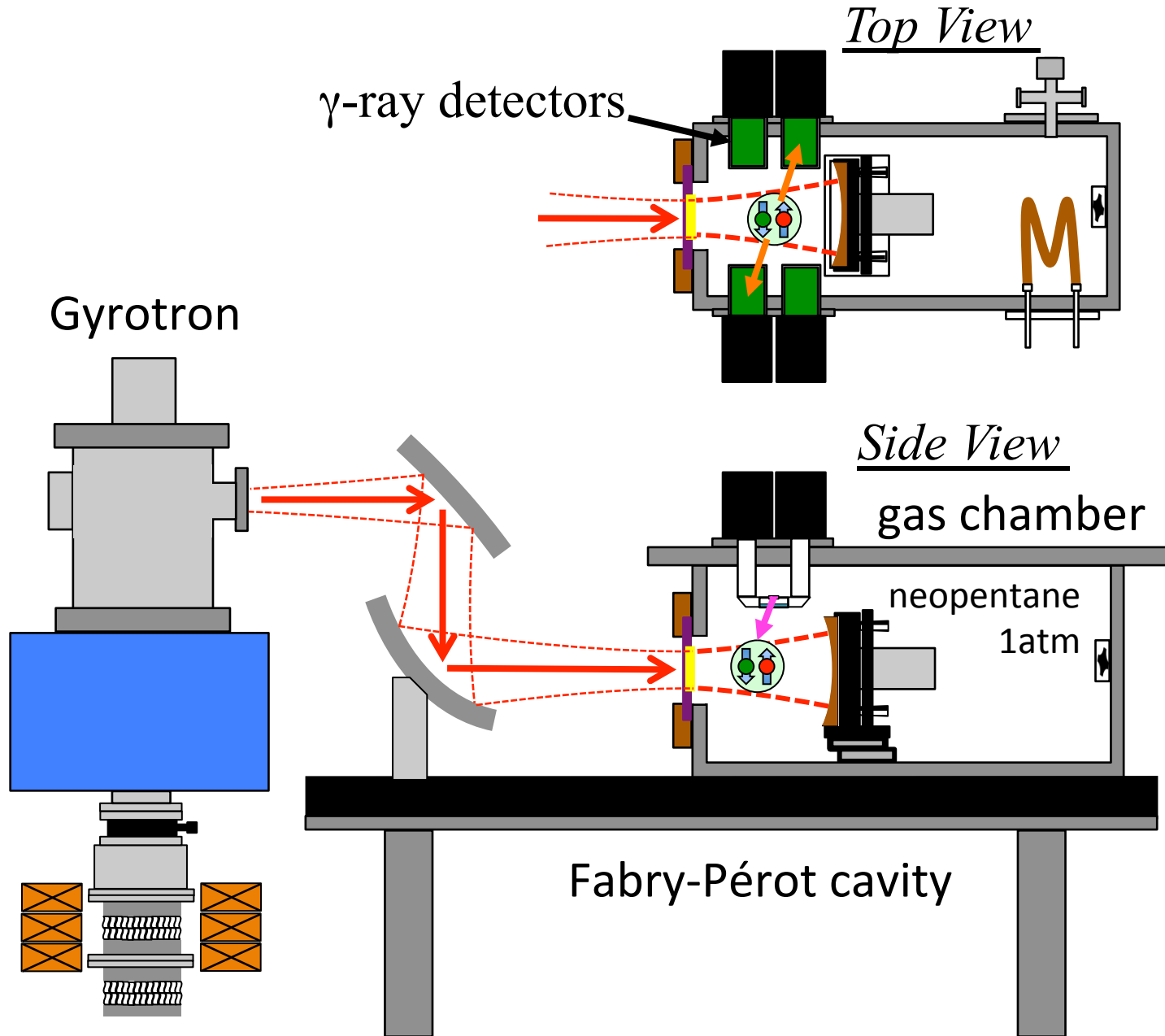


★ **High-power millimeter-wave spectroscopy is quite unique**

Newly developed two devices:

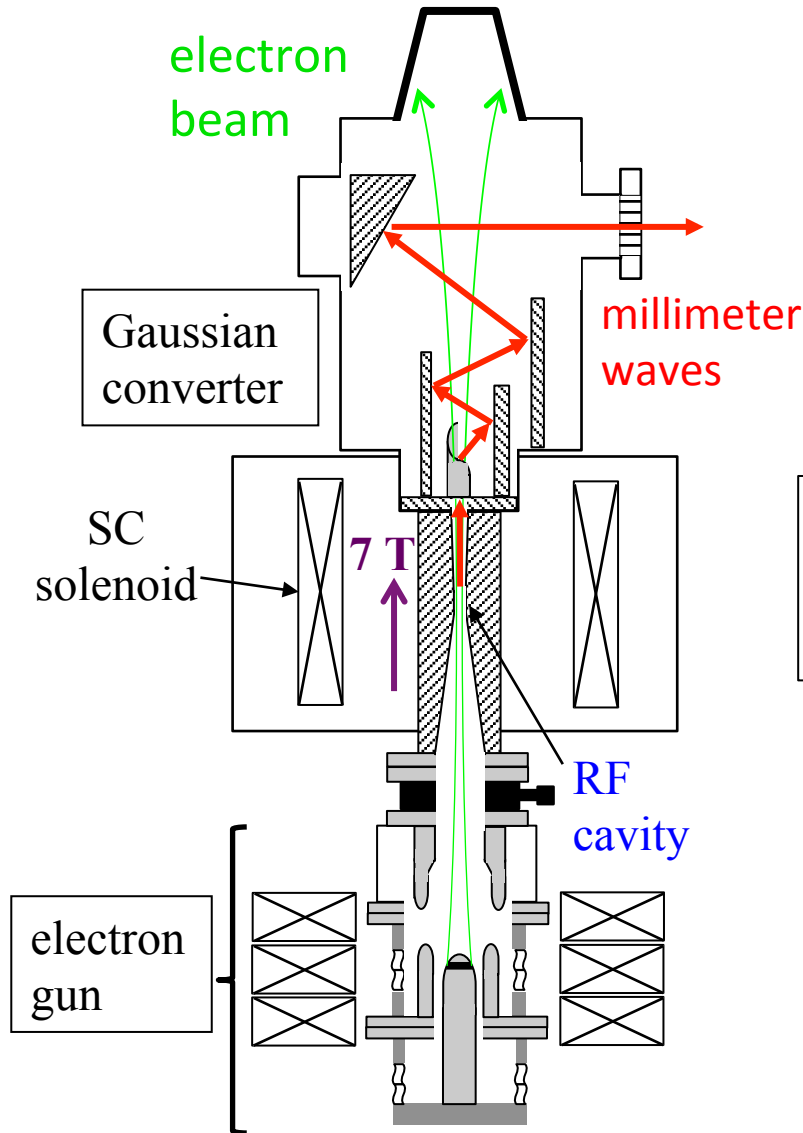
1. A millimeter-wave oscillator of over **100 W** (**Gyrotron**)
2. A resonant cavity accumulating over **20 kW** (**Fabry-Pérot cavity**)

Experimental setup



MILLIMETER-WAVE DEVICES

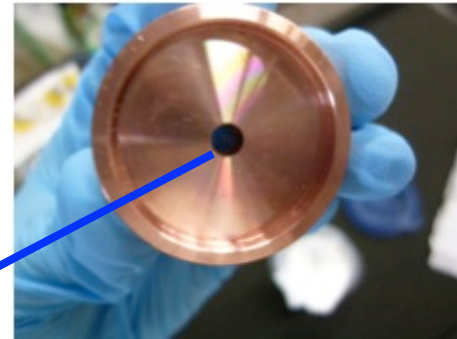
Gyrotron oscillator



A gyrotron is a high-power CW oscillator of millimeter waves.

The **electrons** thermally emitted from the gun move in a cyclotron motion in a magnetic field

$$\omega_c = \frac{eB}{m_e \gamma} \sim 200 \text{ GHz}$$

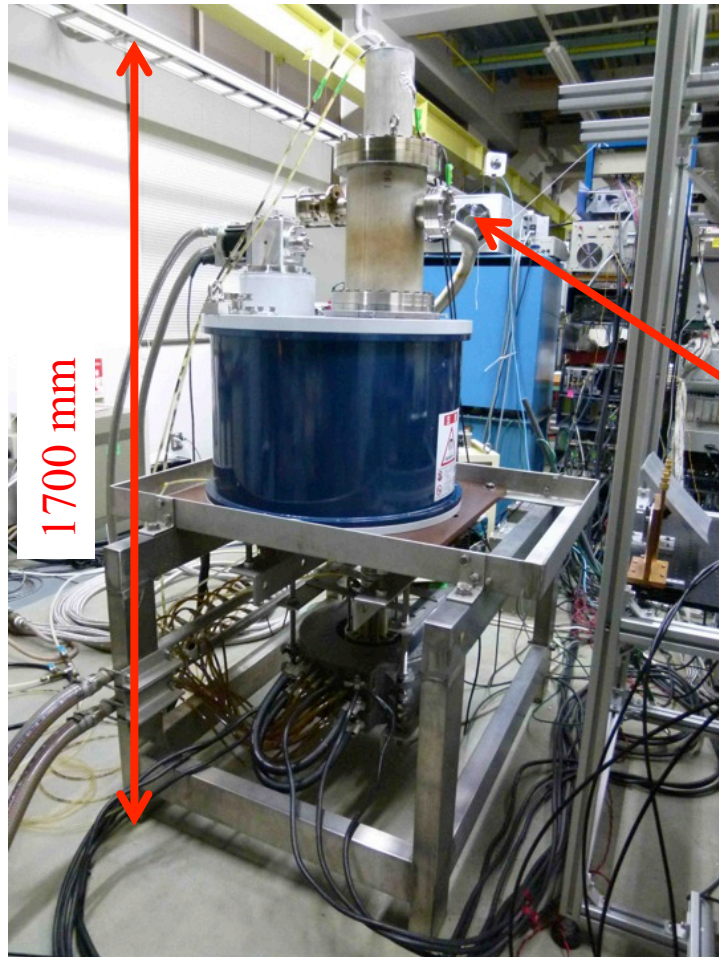


In the **RF cavity** (an open-ended cavity of $\phi 5\text{mm}$), electron beam couples and excites a TE modes ($\sim \omega_c$)

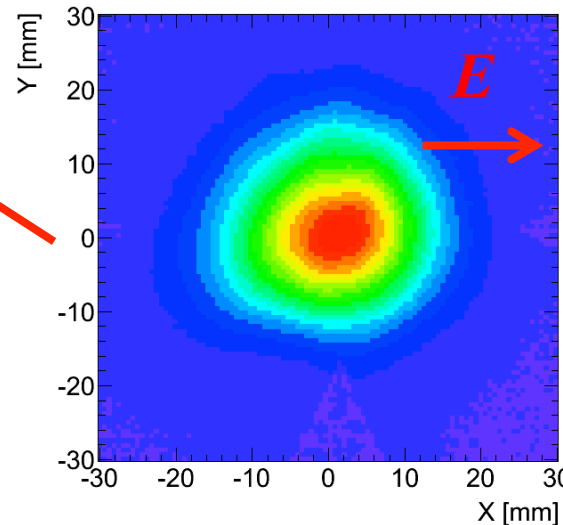
→ cyclotron-maser resonance

Developed gyrotron

collaboration with **Fukui University**



Power = 100 – 600 W

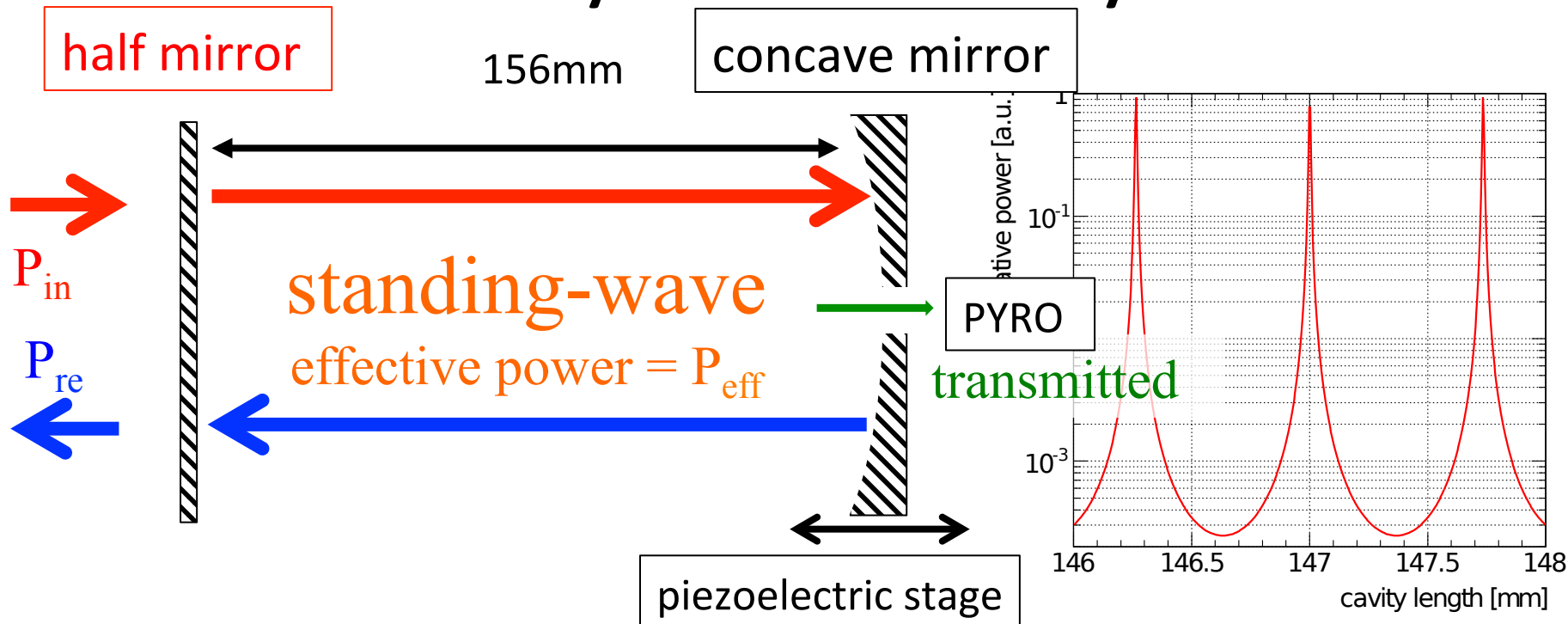


Output radiation is converted from a TE_{52} mode to TEM_{00} using geometrical optics

Output power is limited by the power supply (CW, 18kV, 500 mA)

→ A **resonant cavity** is required.

Fabry-Pérot cavity

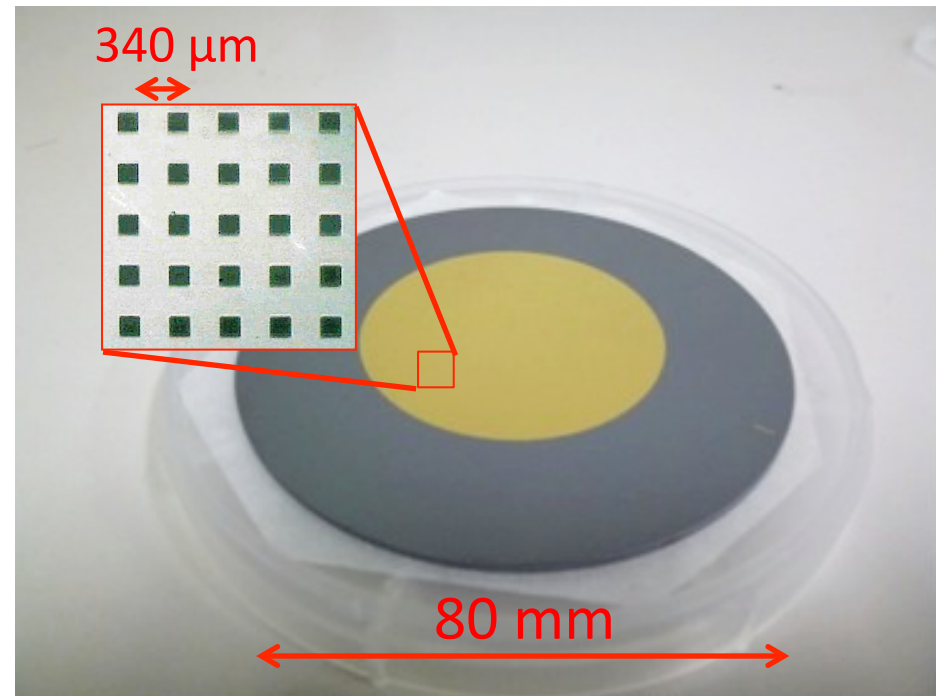
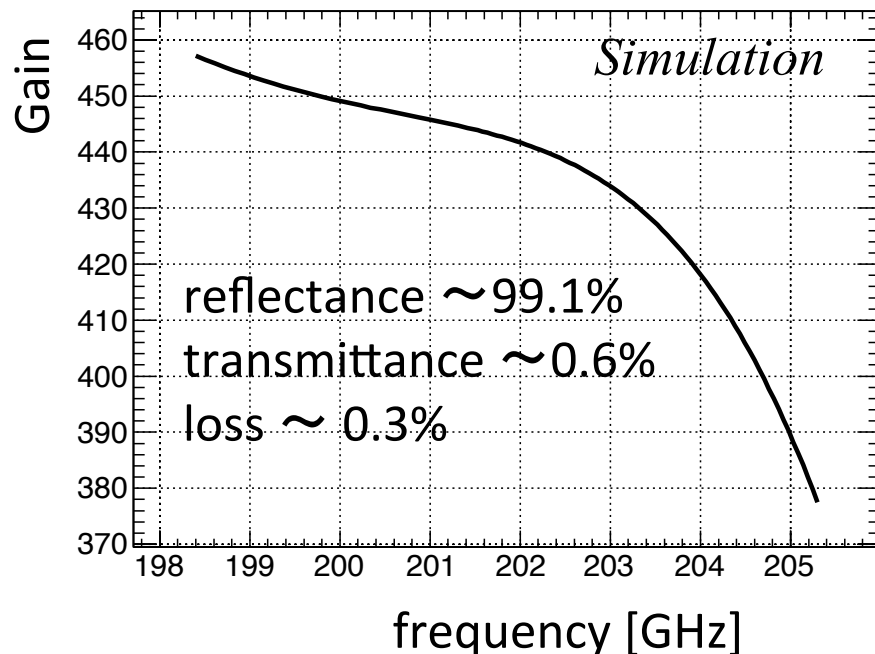
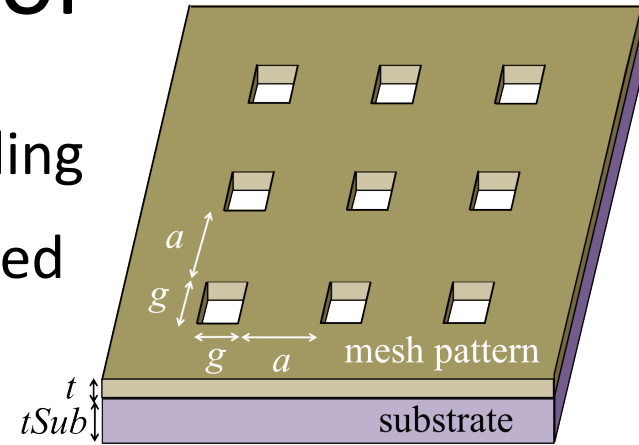


- A Fabry-Pérot cavity resonates when cavity length matches $\lambda/2 \times p$
- The resonance is monitored by a pyroelectric detector (PYRO) through a small hole ($\phi 0.6$ mm) on a copper concave mirror.
- The cavity length is controlled ($< 1\mu\text{m}$) with a piezoelectric stage

★ **High-reflectance and low-loss half mirror are required for high-Q**
 → I designed a **gold mesh mirror** as an efficient half mirror!

Gold mesh mirror

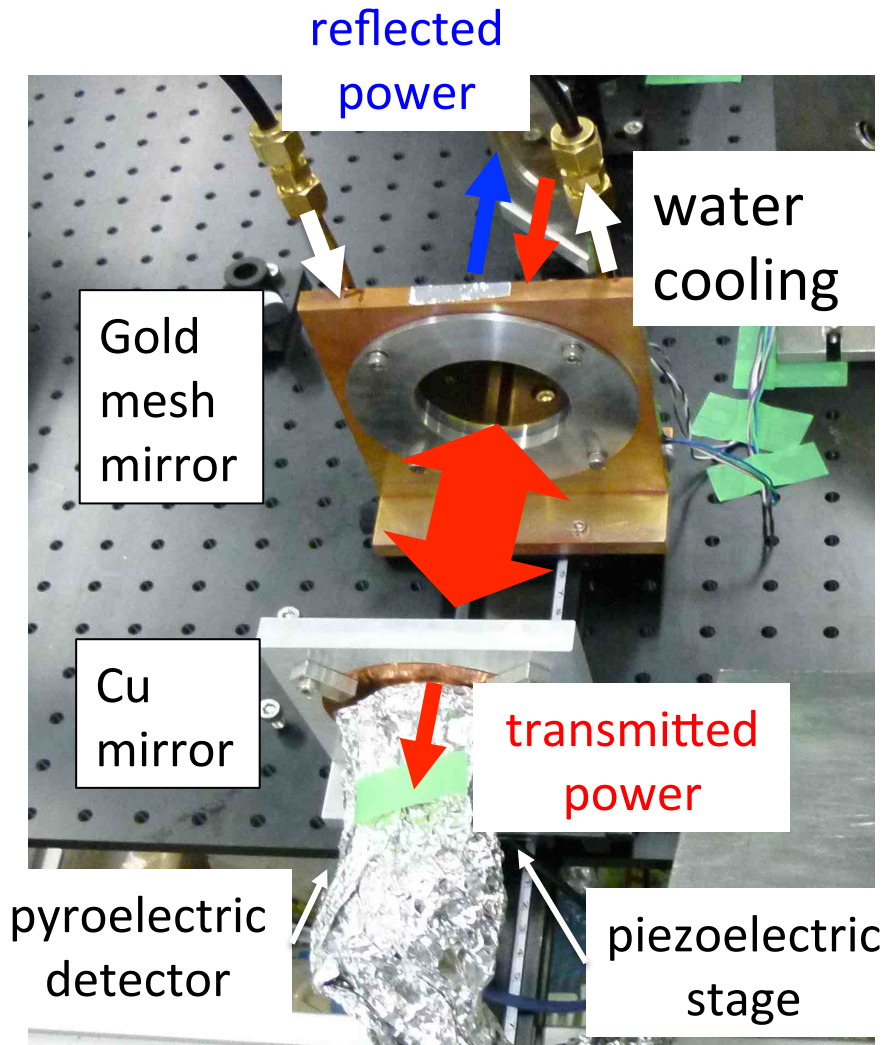
- Thin ($1\mu\text{m}$) gold film is evaporated on a high-resistive silicon base (1.96mm) with water-cooling
- Using CST MW Studio, gain ($P_{\text{eff}}/P_{\text{in}}$) is optimized to have small frequency dependence ($<10\%$) ($a=200\mu\text{m}$, $g=140\mu\text{m}$)



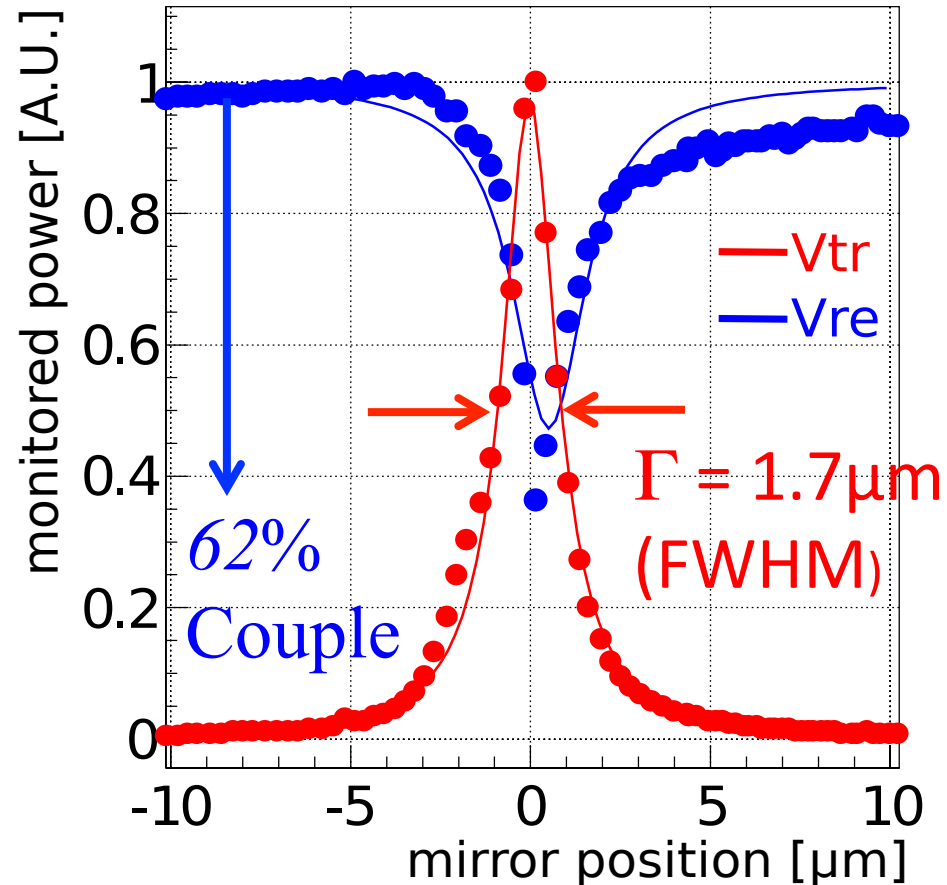
AM, et al., J. Infrared Milli. Terahz Waves 35, 1, 91 (2014)

Designed gain 400 \rightarrow 100W input means 40kW effective power 13

Test of the Fabry-Pérot cavity



Reflected & transmitted power



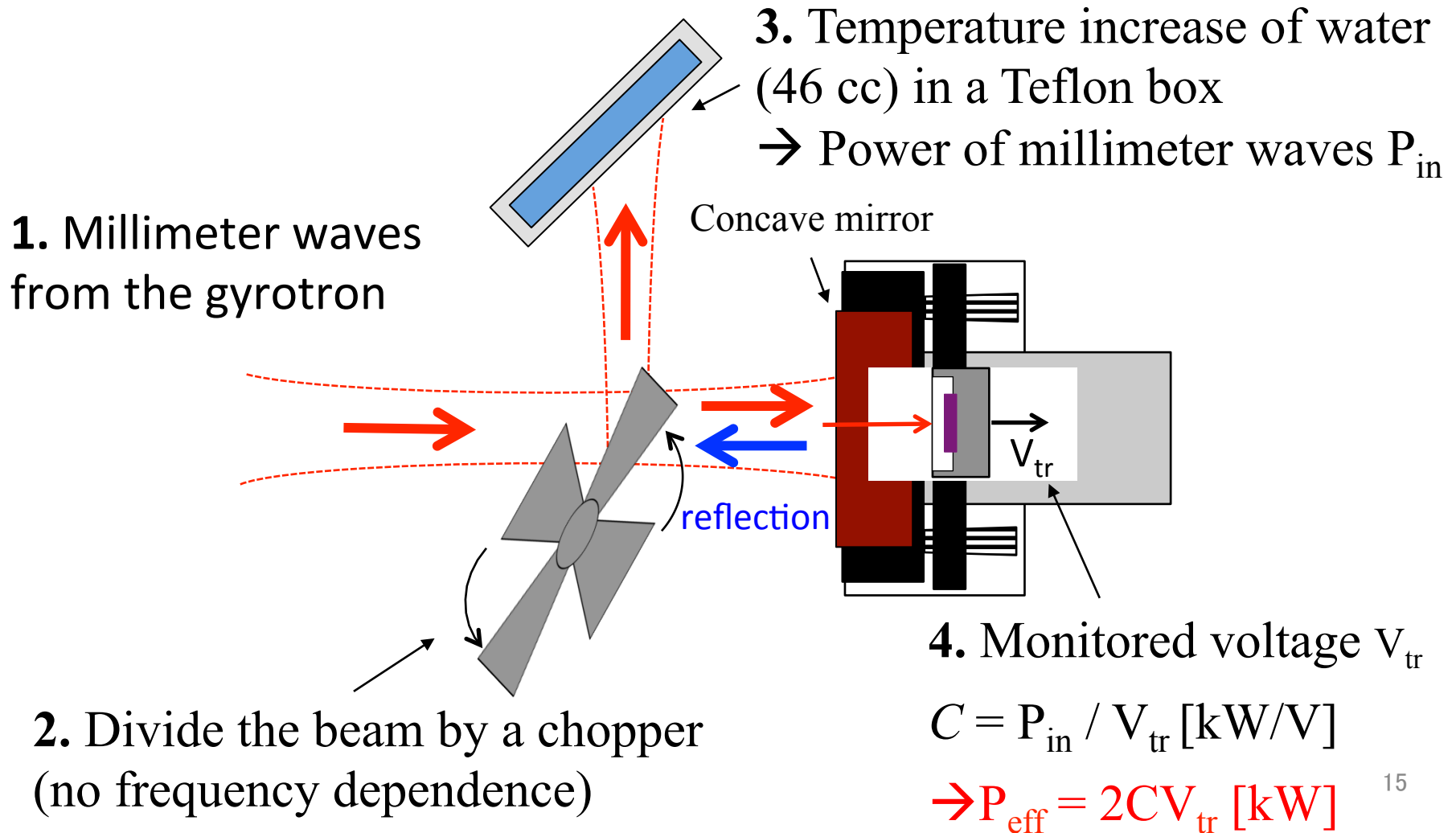
Γ and Coupling \rightarrow Obtained P_{eff} is inaccurate ($\Delta P_{\text{eff}} \gg 50\%$)

\rightarrow A better calibration method ($\Delta P_{\text{eff}} \sim 20\%$) is required.

Estimation of P_{eff}

Transmitted power samples stored *effective* power

→ calibration coefficient C (accuracy $\sim 20\%$)



Experiments were done at eight different frequencies

The frequency is scanned by changing the RF cavity in the gyrotron.

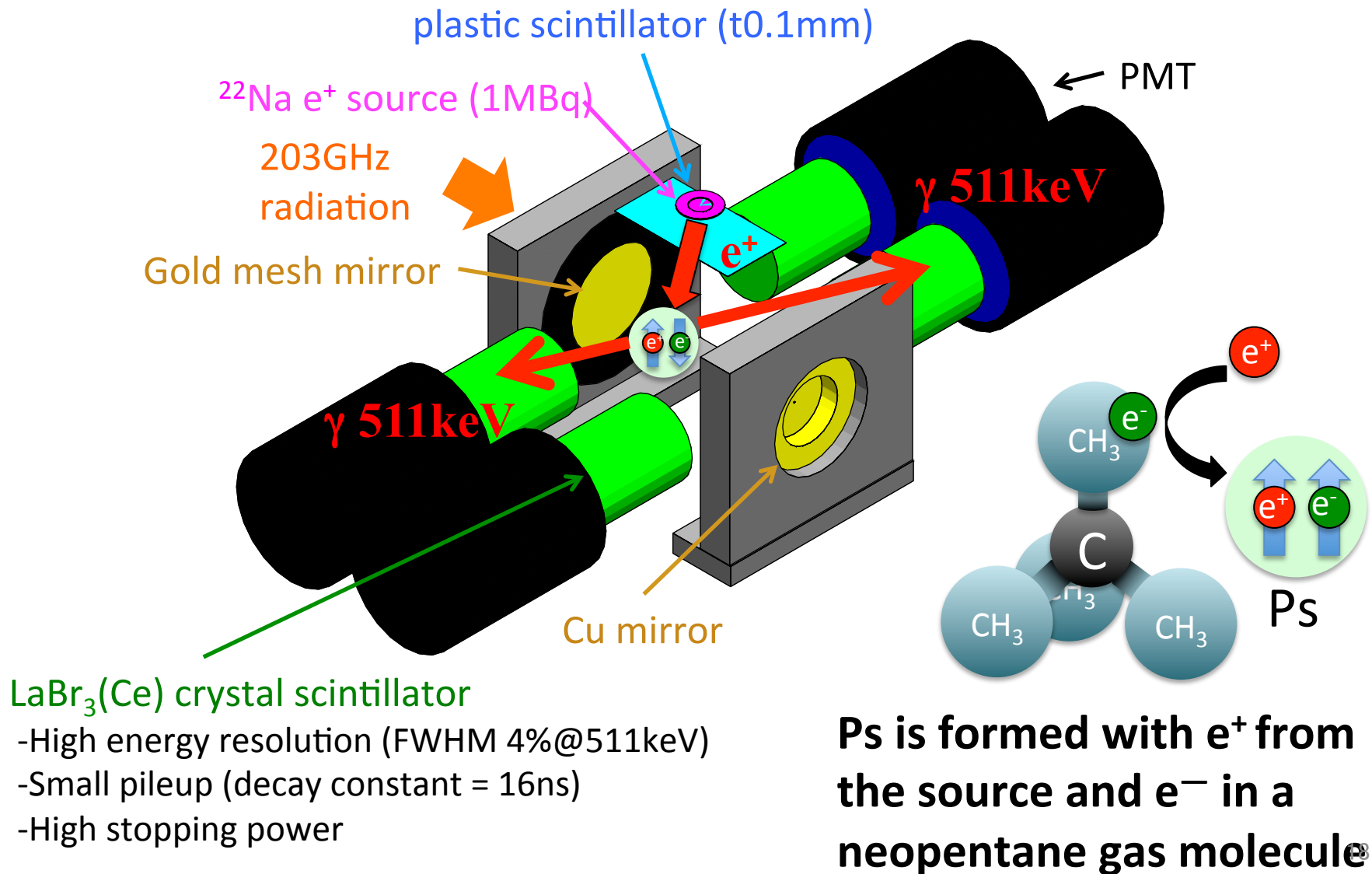
→ Frequency tuning is new in the high-power millimeter-wave range.

cavity radius	f	P _{eff}
(*)2.453 mm	180.59 GHz	41 kW
2.481 mm	201.83 GHz	22 kW
2.475 mm	202.64 GHz	23 kW
2.467 mm	203.00 GHz	21 kW
(**)2.467 mm	203.25 GHz	21 kW
2.463 mm	203.51 GHz	41 kW
2.453 mm	204.56 GHz	20 kW
2.443 mm	205.31 GHz	24 kW

(*) different resonant mode TE₄₂ (**) Radius was expanded during the measurement

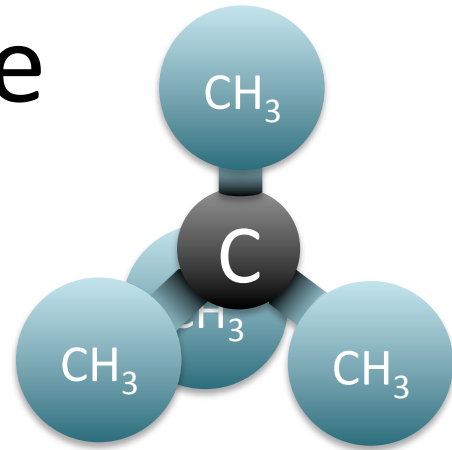
SETUP FOR PS PRODUCTION AND DETECTION

Ps production assembly & γ -ray detector



Gift from Neopentane

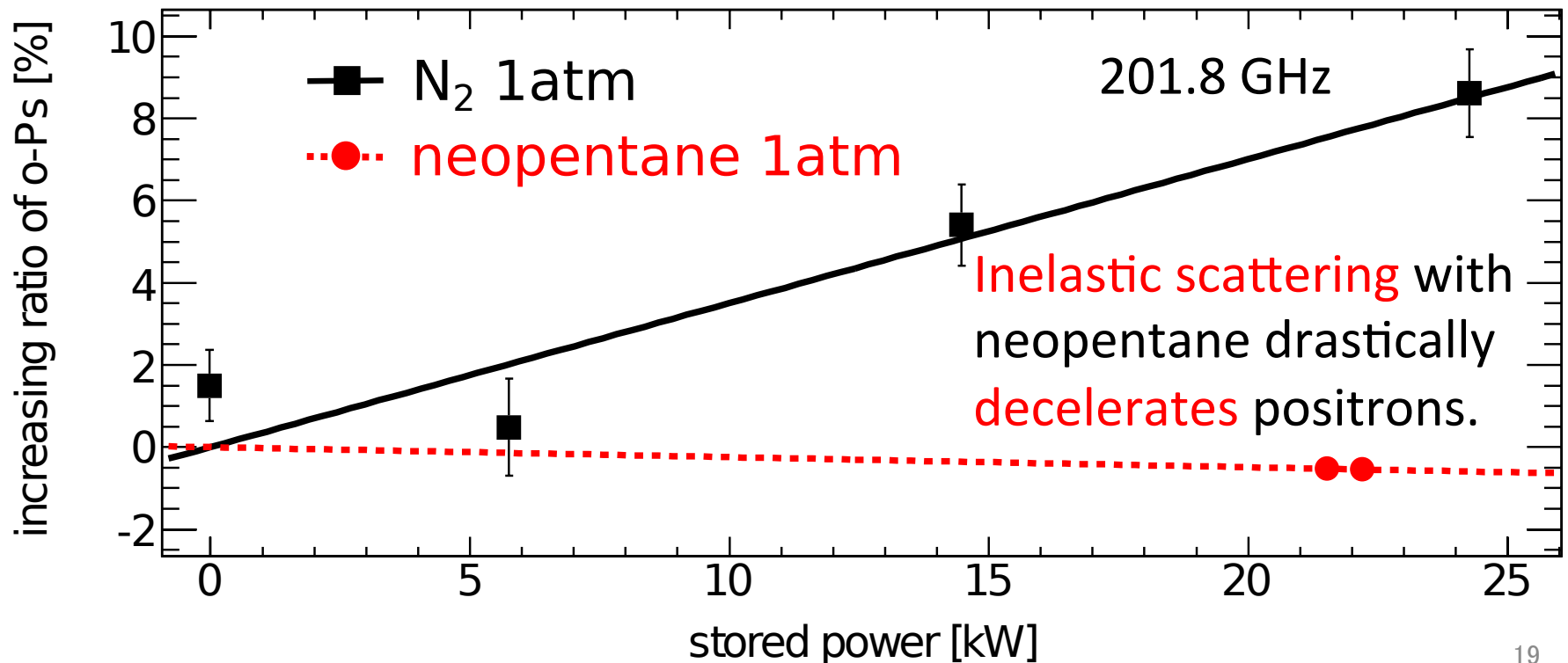
- Large background is found in N₂ gas
- **Ps production ratio** increases with stored power



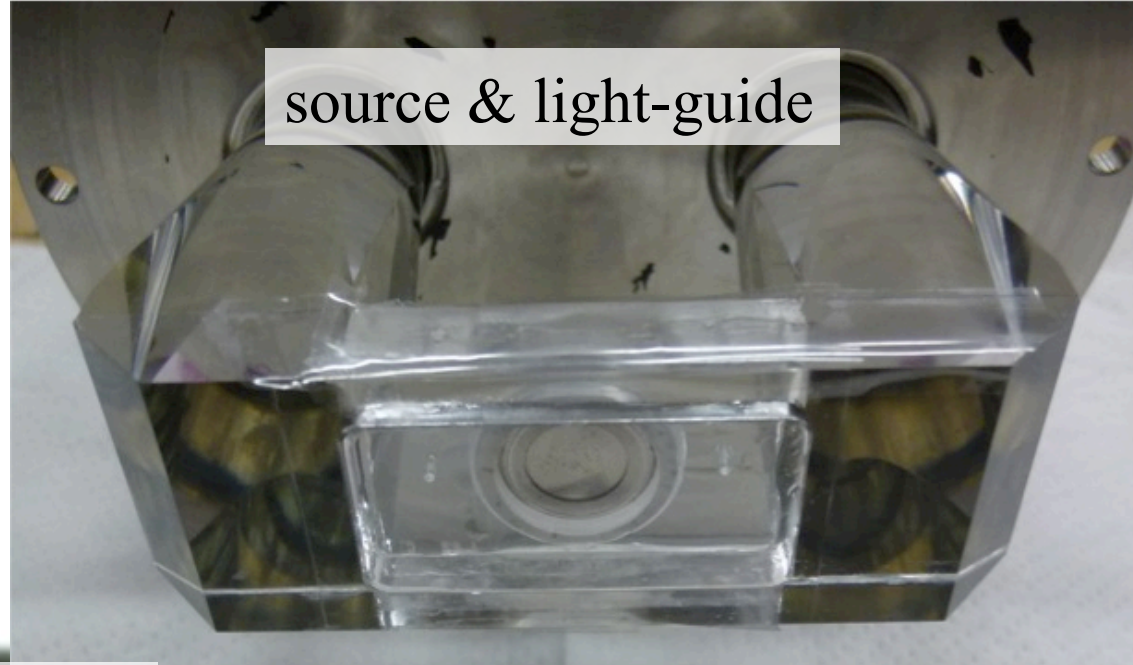
Reason

Strong E-field in the cavity **accelerates** positrons

→ Probability of Ps production increases → Big offset for the signal

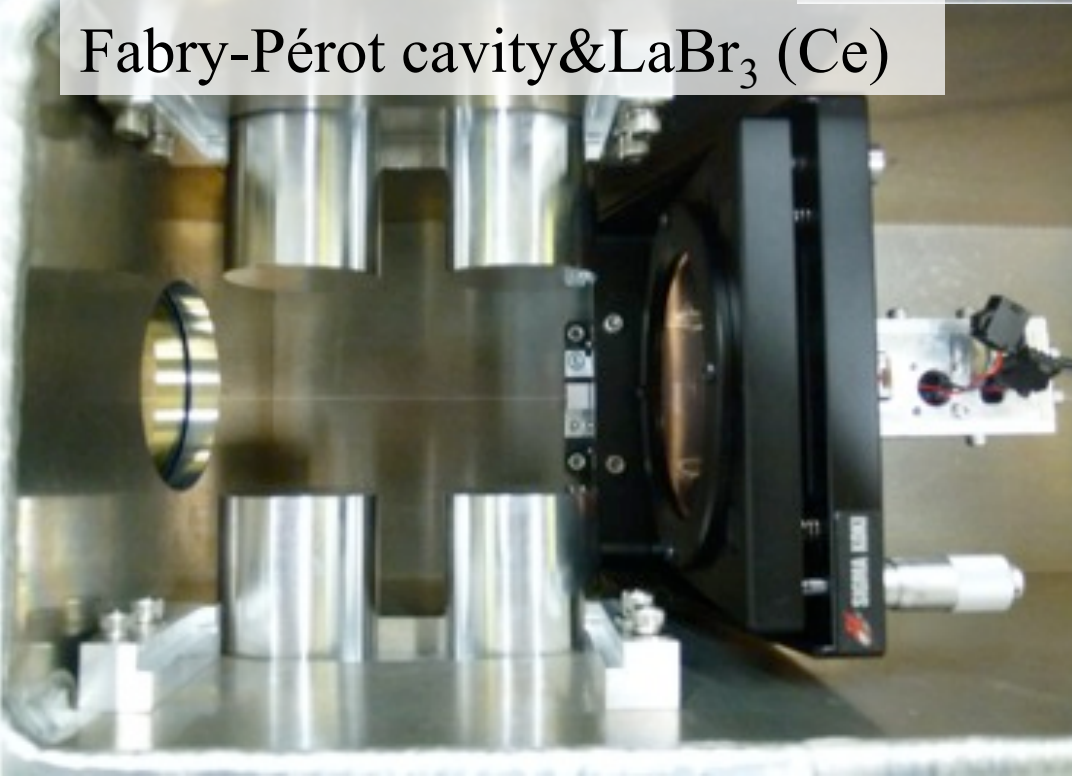


Photograph of the gas chamber



source & light-guide

Fabry-Pérot cavity & LaBr_3 (Ce)



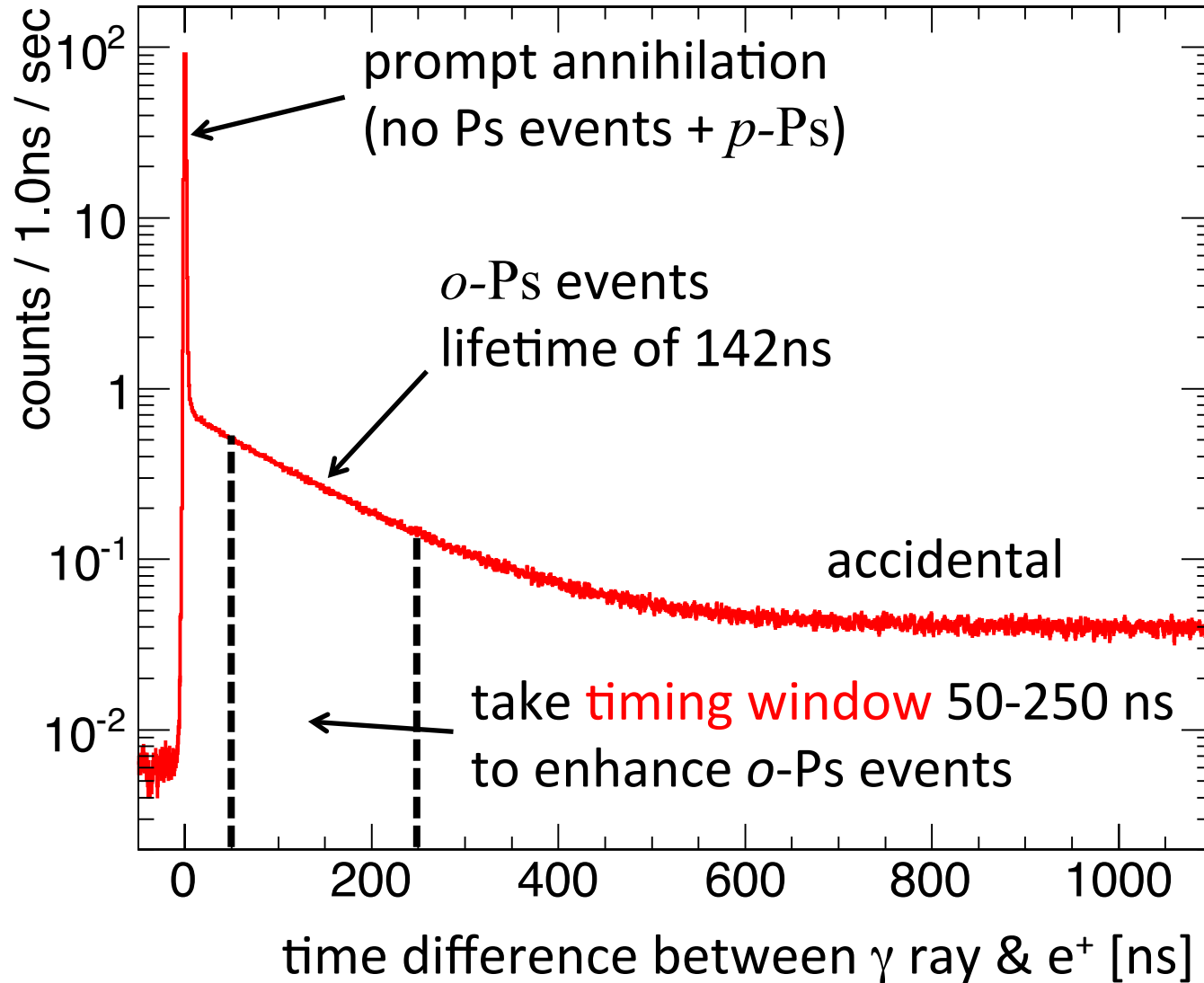
The silicon base for the mesh

- water-cooling
- window of the chamber
- light shielding

DATA ANALYSIS

Timing window

signal: $o\text{-Ps}$ (long lifetime of 142ns) $\rightarrow p\text{-Ps} \rightarrow 2\gamma$ (back-to-back 511keV)

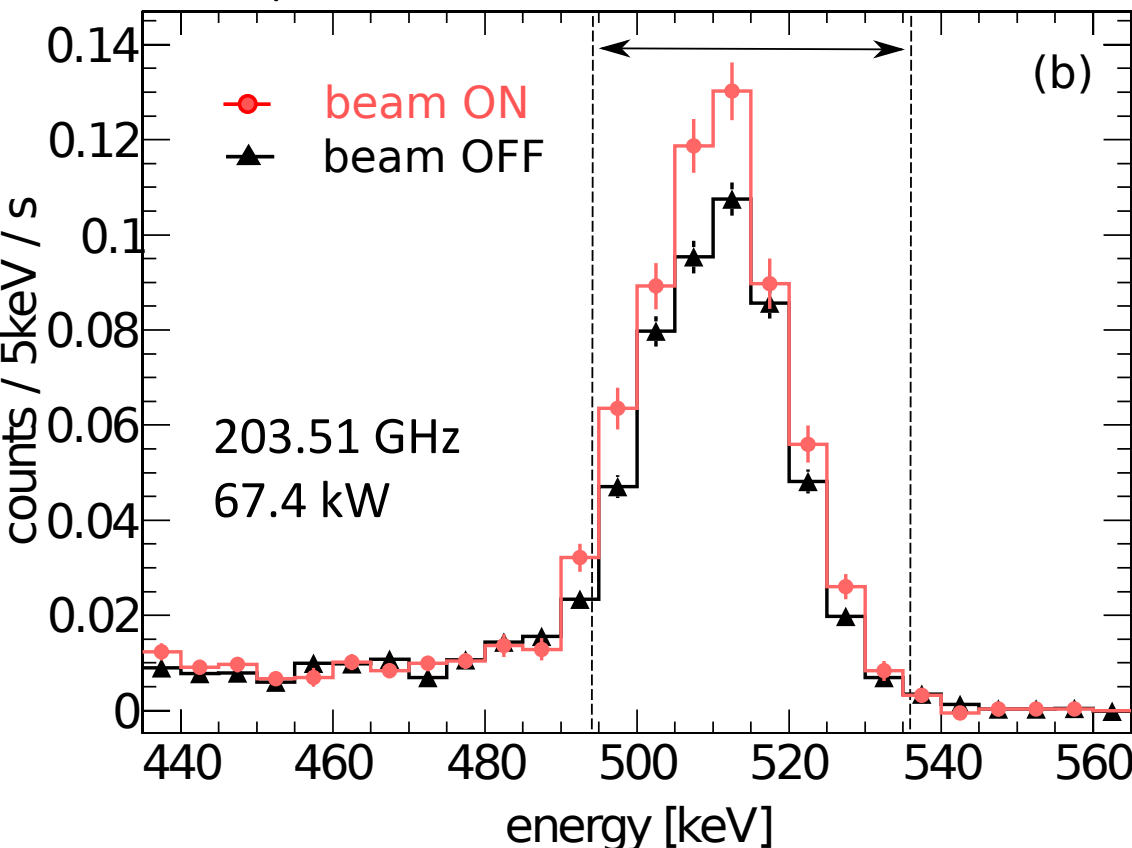


Energy cut & background estimation

signal: o-Ps (long lifetime of 142ns) \rightarrow p-Ps \rightarrow 2 γ (back-to-back 511keV)

Back-to-back and energy selection of 511 keV $+3\sigma/-2\sigma$
 \rightarrow enhance “ \rightarrow 2 γ ” events

γ -ray Energy spectrum



- Remaining background

- \rightarrow Make the gyrotron work in pulsed operation (duty ratio 30%, repetition rate 5Hz)

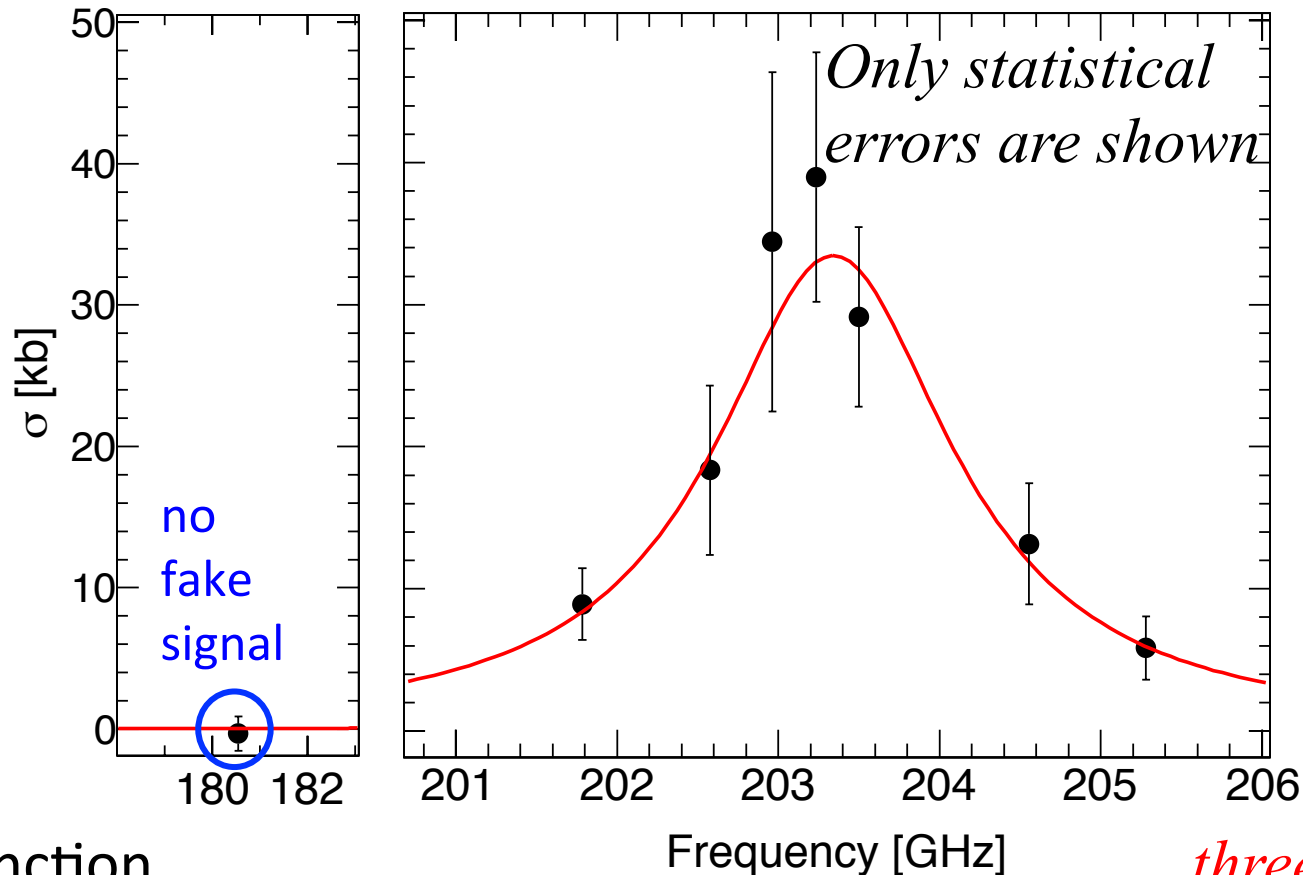
- \rightarrow Use beam-OFF events to estimate background

Significant signal

(ON – OFF)
 $= 93.3 \pm 14.4$ mHz

Breit-Wigner Resonance

Transition signal (**ON**-OFF) is interpreted as cross-section σ of the resonance using P_{eff} , and Ps position distribution from MC simulations



Fitting function

$$g(\omega) = 3A \frac{\pi c^2}{\hbar^2 \omega_0^2} \cdot \frac{1}{\pi} \frac{\Gamma/2}{(\omega - \omega_0)^2 + (\Gamma/2)^2}$$

$$\omega_0 = 2\pi \Delta_{\text{HFS}}^{\text{Ps}}$$

$$\Gamma = A + \Gamma_{\text{o-Ps}} + \Gamma_{\text{p-Ps}} \sim \Gamma_{\text{p-Ps}}$$

three fitting parameters

Systematic uncertainties

Source	$\Delta_{\text{HFS}}^{\text{Ps}}$	$\Gamma_{\text{p-Ps}}$	A
Power estimation	430 ppm	10.0 %	7.2 %
Stark effect	460 ppm	—	—
Monte Carlo simulation	280 ppm	5.5 %	3.0 %
Total	540 ppm	11.4 %	7.8 %

- Uncertainty of calibration factor C
- The Stark effect from gas molecules shifts Δ_{HFS}
- Monte Carlo simulation for widely spread Ps distribution

Results

value \pm statistical error \pm systematic error

Parameter	This experiment	Theory
$\Delta_{\text{HFS}}^{\text{Ps}}$ [GHz]	$203.39^{+0.15}_{-0.14} \pm 0.11$	203.391 91(22)
$\Gamma_{\text{p-Ps}}$ [ns ⁻¹]	$11.2^{+1.9}_{-2.3} \pm 1.3$	(*) 7.9894 76(13)
A [$\times 10^{-8}$ s ⁻¹]	$3.69 \pm 0.48 \pm 0.29$	(**) 3.37

(*) A. Kniehl, and A. A. Penin, Phys. Rev. Lett. 85, 1210 (2000);
K. Melnikov and A. Yelkhovsky, Phys. Rev. D. 62, 116003 (2000).

(**) P. Wallyn, et al., Astrophys. J. 465, 473 (1996).

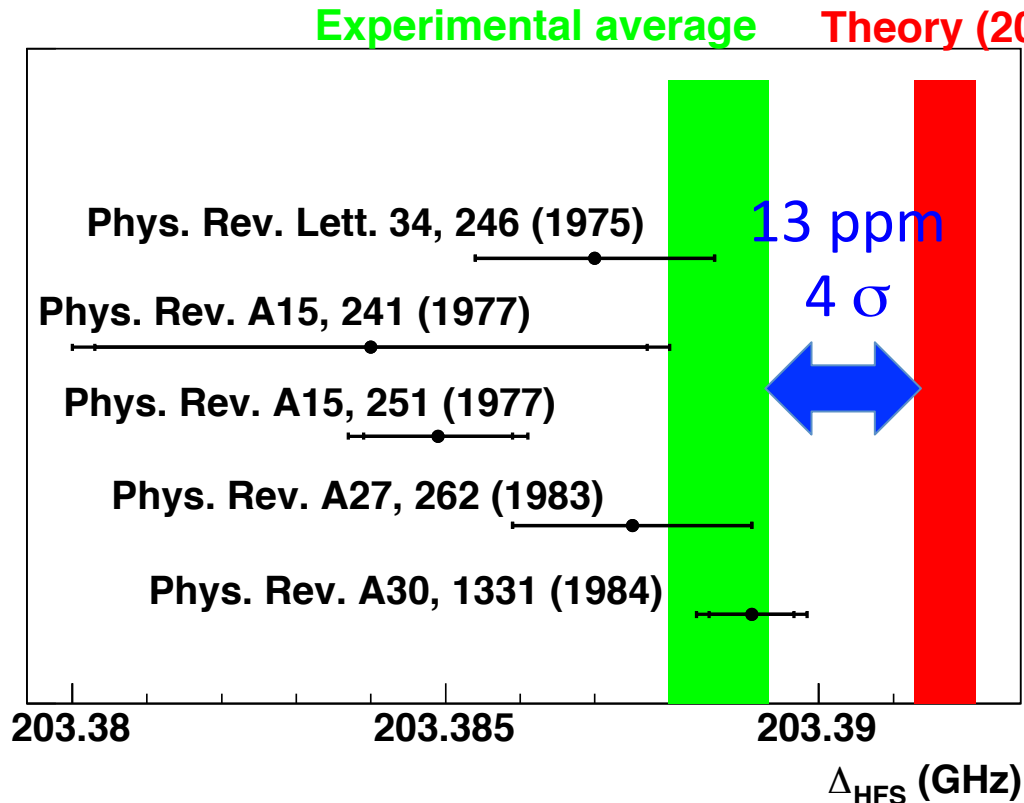
These 2 parameters are firstly determined with a direct way.

★ We firstly demonstrate Δ_{HFS} can be directly determined!

PROBLEM IN HYPERFINE STRUCTURE & FUTURE PROSPECT

Displacement

Higher order QED corrections up to $O(\alpha^3 \ln \alpha^{-1})$ are available in 2000. Results of some $O(\alpha^3)$ diagrams were reported (c.f. prof. Adkins).



4 σ Discrepancy!

Possible reasons

- i) Magnetic field
- ii) Material effect
- iii) (QED)
- iv) Physics BSM

→ Study of systematic uncertainties is required

Theoretical papers:

- B.A. Kniehl et al. Phys. Rev. Lett. 85, 5094 (200)
- K. Melnikov et al., Phys. Rev. Lett. 86, 1498 (2001)
- R.K. Hill, Phys. Rev. Lett. 86, 3280 (2001)
- M. Baker, et al. Phys. Rev. Lett. 112, 120407 (2014)
- G.S. Adkins et al., Phys. Rev. A 89, 052518 (2014)

Some recent studies

Y. Sasaki and AM et al, Phys. Lett. B 697 121 (2011)

- Quantum oscillation between **Zeeman shifted levels**
- 200 ppm accuracy but shown to be improved to 10 ppm level

D.B. Cassidy et al, Phys. Rev. Lett. 109 073401 (2012)

- Saturation absorption spectroscopy between **Zeeman shifted 1S and 2P levels**
- 2% accuracy but promising improvements because of established laser technology

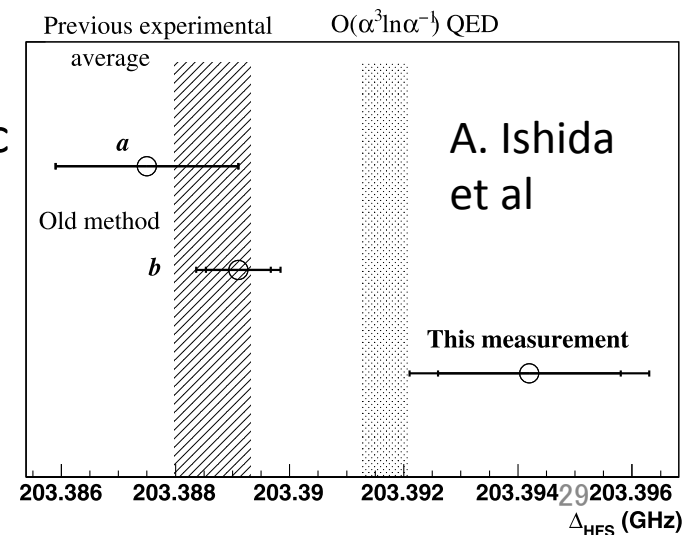
A. Ishida et al, Phys. Lett. B 734 338 (2014)

- Conventional method i.e. transition between **Zeeman shifted 1S levels** by RF

- i) Uniform B-field with a big SC solenoid ($\phi 80\text{cm}$).
- ii) An interpolation to vacuum is studied considering slow thermalization of Ps in gas (Ps is more energetic than thermal temperature for long time).

→ favors QED by 2.7σ

☆ **Direct measurement can be a complementary check of methods using the Zeeman effect.**



Toward precision measurement

The discrepancy $13 \text{ ppm} < \text{current accuracy } 900 \text{ ppm}$

Three major improvements

- Positron beam for better statistics ($\times 10^4$)
 - Well collimated beam can improve fraction of positrons existing in the Fabry-Pérot cavity [$10^8 \text{ Hz} \gg 10^3 \text{ Hz}$ from ^{22}Na source (1MBq)]
 - Positron beam and Ps converter are available in KEK
- Ps is formed in vacuum to eliminate gas effects
 - Efficient Ps converter (a material in which e^+ is converted to Ps extracted into vacuum; such as heated alumina efficiency=50%)
 - No stark effect and non-thermalization effect of Ps
- MW-class gyrotron
 - No Fabry-Pérot cavity \rightarrow Much better power estimation (0.3%) by temperature increase of water-flow for cooling power dump
 - Well controlled power ($< 1\%$)

Estimated future uncertainties

Source	$\Delta_{\text{HFS}}^{\text{Ps}}$
Statistics	5 ppm
Power estimation	3 ppm
Power control	4 ppm
Frequency stability	< 1 ppm
Power stability	4 ppm
Doppler shift	1 ppm
Total	8 ppm

- Uncertainty from **the magnetic field** can be the largest systematic uncertainty for the Zeeman-type measurement in vacuum
- Instead of the magnetic field, systematic uncertainties from **millimeter waves** can be dominant in the direct measurement.
 - They are complementary with each other.

Conclusion

- Though Ps is a good system to study QED, its hyperfine structure has never been directly measured because of difficulties regarding the use of millimeter waves.
- We developed new millimeter-wave devices: gyrotron and Fabry-Pérot cavity; $\Delta_{\text{HFS}}^{\text{Ps}}$ was firstly measured with the direct transition method.
- value: $\Delta_{\text{HFS}}^{\text{Ps}} = 203.39^{+0.15}_{-0.14}(\text{stat.}) \pm 0.11 (\text{syst.}) \text{ GHz}$
- p -Ps decay width was also measured
- value: $\Gamma_{p\text{-Ps}} = 11.2^{+1.9}_{-2.3}(\text{stat.}) \pm 1.3 (\text{syst.}) \text{ ns}^{-1}$
- arXiv:1403.0312
- PhD thesis: http://www.icepp.s.u-tokyo.ac.jp/papers/ps/thesis/doctor/miyazaki_thesis.pdf
(It will be published by Springer Thesis in one year)