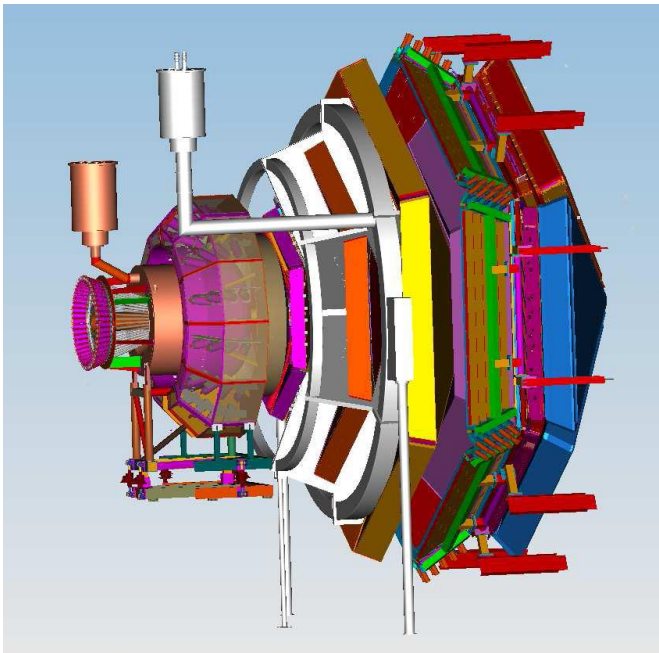
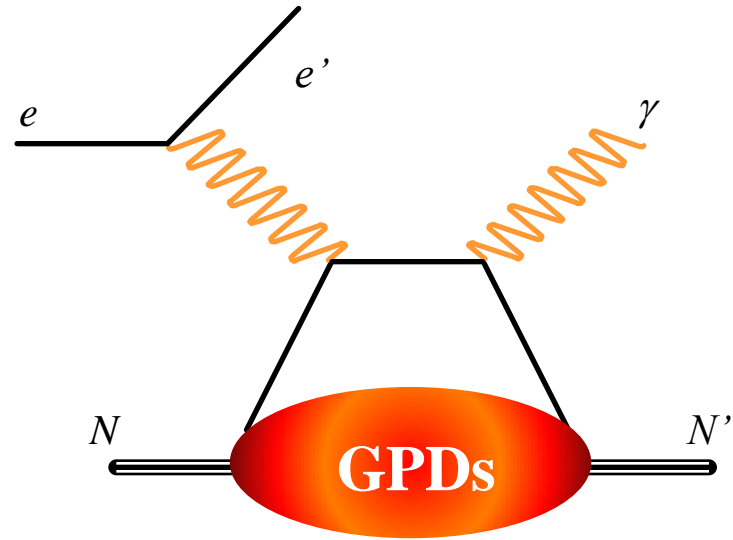


Experimental studies of Generalized Parton Distributions

*Silvia Niccolai, IPN Orsay
for the CLAS Collaboration*

EUNPC2015
4/9/2015, Groningen



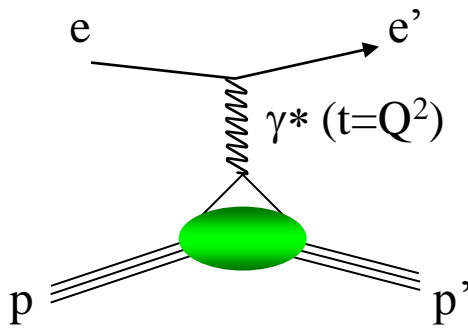
- Interest of GPDs
- GPDs and Deeply Virtual Compton Scattering
 - New DVCS results from Jefferson Lab
 - The JLab 12 GeV upgrade
 - Future JLab experiments on DVCS



Electron-nucleon scattering to study nucleon structure

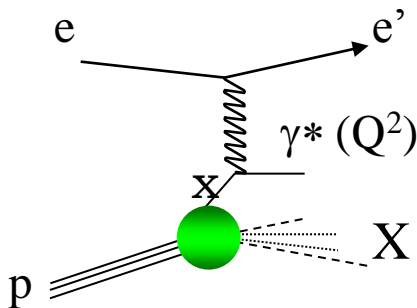
Electrons are **structureless** and interact only **electromagnetically**

➤ 1950: **Elastic scattering** $ep \rightarrow e'p'$ (Hofstadter, Nobel prize 1961)

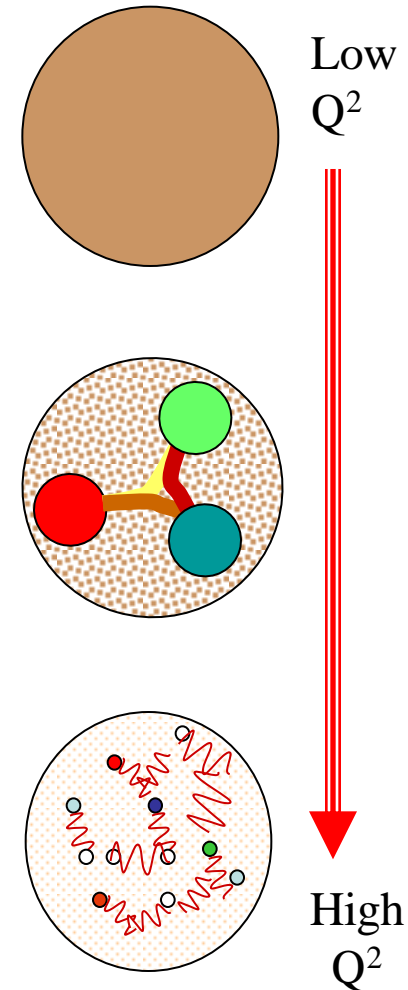


- The proton is **not a point-like object**
- Measurement of charge and current distributions of the proton: **form factors** ($F_1(t), F_2(t)$)

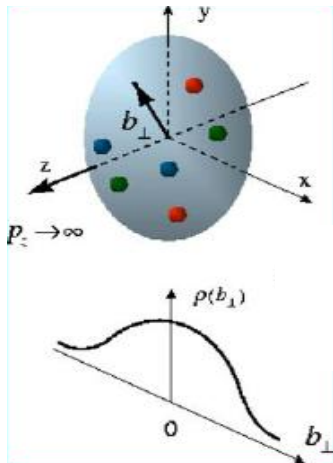
➤ 1967: **Deep inelastic scattering** (DIS) $ep \rightarrow e'X$
(Friedman, Kendall, Taylor, Nobel prize 1990)



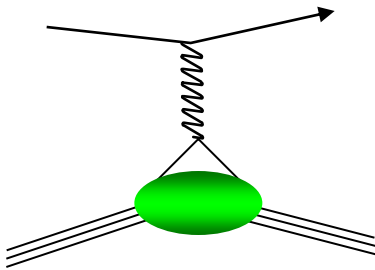
- Discovery of the **quarks** (or “partons”)
- Measurement of the momentum and spin distributions of the partons: $q(x), \Delta q(x)$



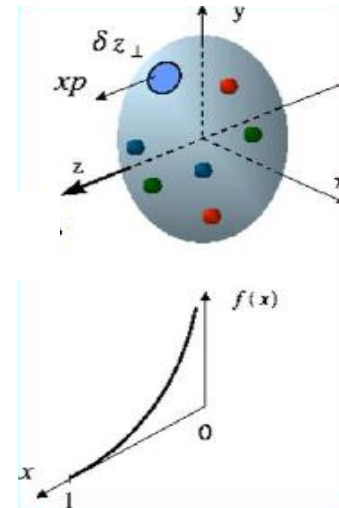
Electron-proton scattering: yesterday



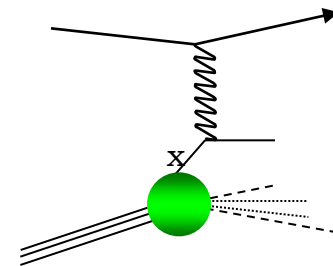
Form factors:
transverse quark
distribution in
coordinate space



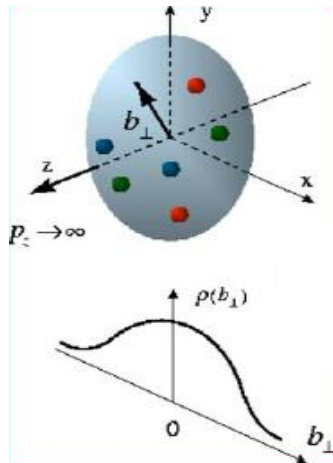
?



Parton distributions:
longitudinal
quark distribution
in momentum space

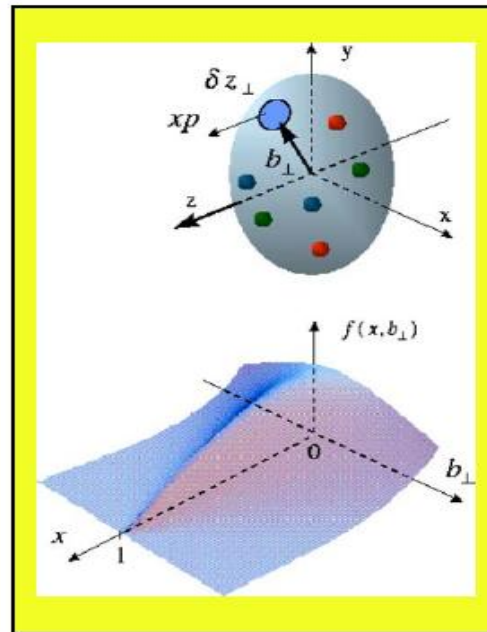


Electron-proton scattering: today



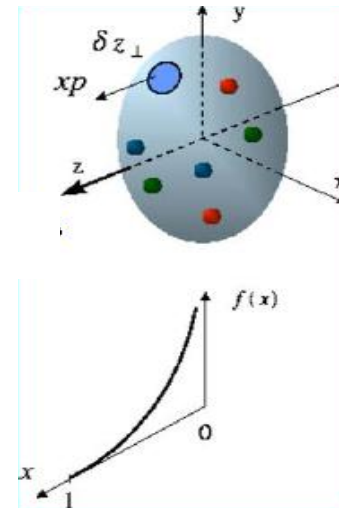
Form factors:
transverse quark
distribution in
coordinate space

GPDs: $H, E, \tilde{H}, \tilde{E}$
Fully correlated quark
distributions in both coordinate
and momentum space

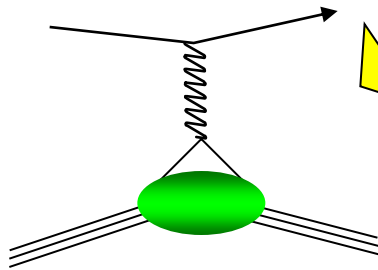


Accessible in
hard exclusive processes

High Q^2 Final state known

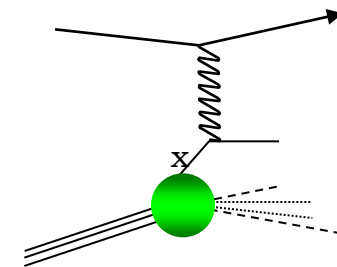


Parton distributions:
longitudinal
quark distribution
in momentum space



$$\int H(x, \xi, t) dx = F_1(t) \quad (\forall \xi)$$

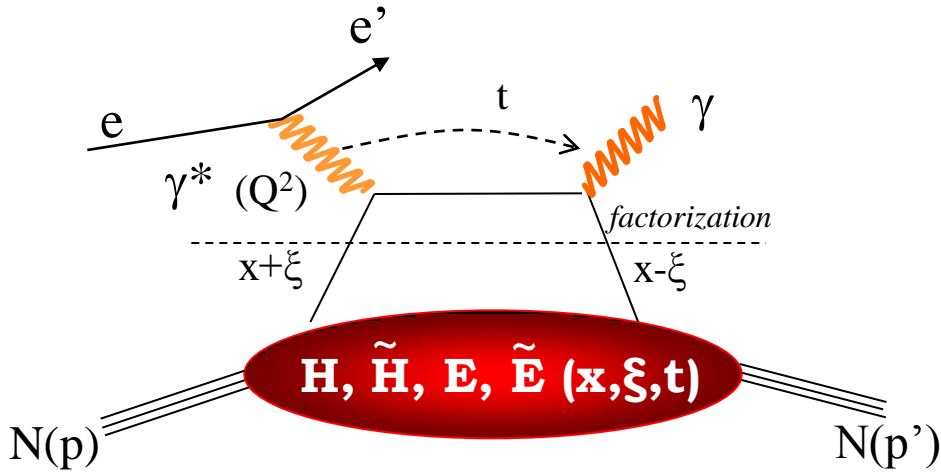
$$\int E(x, \xi, t) dx = F_2(t) \quad (\forall \xi)$$



$$H(x, 0, 0) = q(x),$$

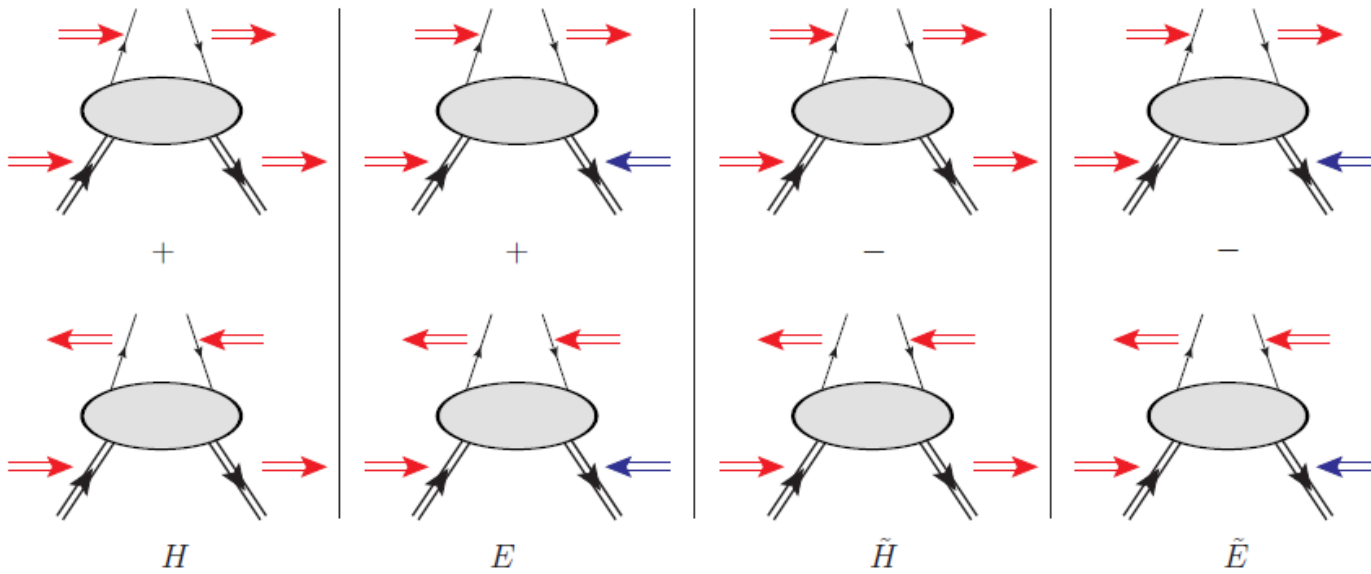
$$\tilde{H}(x, 0, 0) = \Delta q(x)$$

Deeply Virtual Compton Scattering and GPDs



- $Q^2 = -(e-e')^2$
- $x_B = Q^2/2Mv$ $v = E_e - E_{e'}$
- $x+\xi, x-\xi$ longitudinal momentum fractions
- $t = \Delta^2 = (p-p')^2$
- $\xi \cong x_B/(2-x_B)$

« Handbag » factorization valid in the **Bjorken regime**:
high Q^2 , v (fixed x_B), $t \ll Q^2$



GPDs: Fourier transforms of **non-local, non-diagonal QCD operators**

conserve nucleon spin Vector: $H(x, \xi, t)$ Axial-Vector: $\tilde{H}(x, \xi, t)$
flip nucleon spin Tensor: $E(x, \xi, t)$ Pseudoscalar: $\tilde{E}(x, \xi, t)$

At leading order QCD, twist 2, chiral-even (quark helicity is conserved), quark sector
 → 4 GPDs for each quark flavor

Properties and “virtues” of GPDs

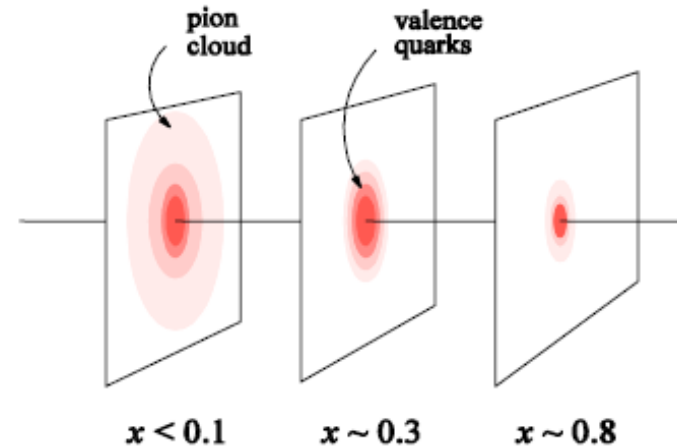
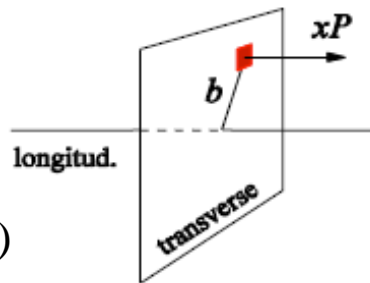
$$\left. \begin{aligned} \int H(x, \xi, t) dx &= F_1(t) \quad \forall \xi \\ \int E(x, \xi, t) dx &= F_2(t) \quad \forall \xi \\ \int \tilde{H}(x, \xi, t) dx &= G_A(t) \quad \forall \xi \\ \int \tilde{E}(x, \xi, t) dx &= G_P(t) \quad \forall \xi \end{aligned} \right\} \text{Link with FFs}$$

$$\left. \begin{aligned} H(x, 0, 0) &= q(x) \\ \tilde{H}(x, 0, 0) &= \Delta q(x) \end{aligned} \right\} \text{Forward limit: PDFs (not for } E, \tilde{E})$$

Nucleon tomography

$$q(x, \mathbf{b}_\perp) = \int_0^\infty \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{i\Delta_\perp \mathbf{b}_\perp} H(x, 0, -\Delta_\perp^2)$$

$$\Delta q(x, \mathbf{b}_\perp) = \int_0^\infty \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{i\Delta_\perp \mathbf{b}_\perp} \tilde{H}(x, 0, -\Delta_\perp^2)$$



M. Burkardt, PRD 62, 71503 (2000)

Quark angular momentum (Ji's sum rule)

$$\frac{1}{2} \int_{-1}^1 x dx (H(x, \xi, t=0) + E(x, \xi, t=0)) = J = \frac{1}{2} \Delta\Sigma + \Delta L$$

X. Ji, Phy.Rev.Lett.78,610(1997)

$$\text{Nucleon spin: } \frac{1}{2} = \underbrace{\frac{1}{2} \Delta\Sigma + \Delta L}_{\mathbf{J}} + \Delta G$$

Intrinsic spin of the quarks $\Delta\Sigma \approx 25\%$

Intrinsic spin on the gluons $\Delta G \approx 0$ (??)

Orbital angular momentum of the quarks ΔL ?

Accessing GPDs through DVCS

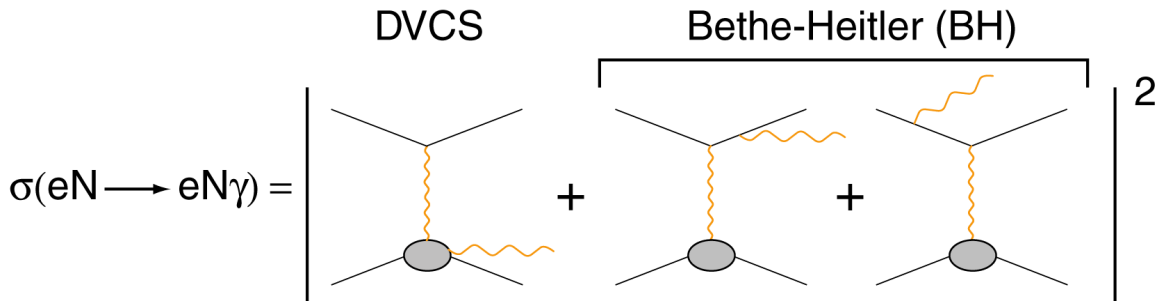
DVCS allows access to 4 complex GPDs-related quantities: **Compton Form Factors (ξ, t)**

$$T^{DVCS} \sim P \int_{-1}^{+1} \frac{GPDs(x, \xi, t)}{x \pm \xi} dx \pm i\pi GPDs(\pm \xi, \xi, t) + \dots$$

Only ξ and t are accessible experimentally

$$Re\mathcal{H}_q = e_q^2 P \int_0^{+1} (H^q(x, \xi, t) - H^q(-x, \xi, t)) \left[\frac{1}{\xi - x} + \frac{1}{\xi + x} \right] dx$$

$$Im\mathcal{H}_q = \pi e_q^2 [H^q(\xi, \xi, t) - H^q(-\xi, \xi, t)]$$

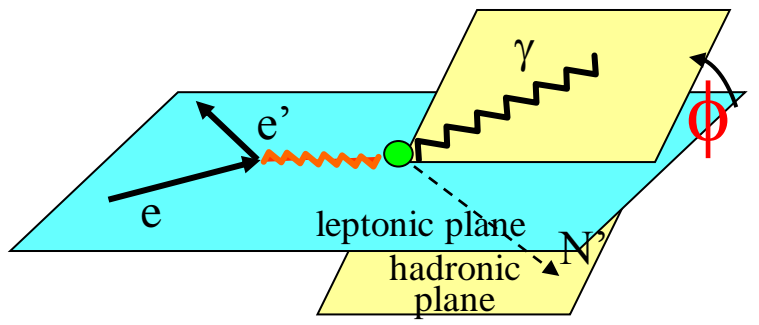


BH is calculable
(electromagnetic FFs)

$$\sigma \sim |T^{DVCS} + T^{BH}|^2 \rightarrow \text{Re}(CFFs) \quad (\text{also DSA})$$

$$\Delta\sigma = \sigma^+ - \sigma^- \propto I(DVCS \cdot BH) \rightarrow \text{Im}(CFFs)$$

$$A = \frac{\Delta\sigma}{2\sigma} \propto \frac{I(DVCS \cdot BH)}{|BH|^2 + |DVCS|^2 + I}$$



Sensitivity to CFFs of DVCS spin observables

$$A_{LU(UL)} = \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-} \propto \frac{s_{1,unp(UL)}^I \sin \phi}{c_{0,unp}^{BH} + c_{0,unp}^I + (c_{1,unp}^{BH} + c_{1,unp}^I) \cos \phi}$$

$$A_{LL} = \frac{\sigma^{++} + \sigma^{+-} - \sigma^{+-} - \sigma^{--}}{\sigma^{++} + \sigma^{+-} + \sigma^{+-} + \sigma^{--}} \propto \frac{c_{0,LP}^{BH} + c_{0,LP}^I + (c_{1,LP}^{BH} + c_{1,LP}^I) \cos \phi}{c_{0,unp}^{BH} + c_{0,unp}^I + (c_{1,unp}^{BH} + c_{1,unp}^I) \cos \phi}$$

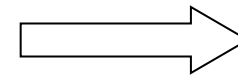
Twist-2
approximation
($-t \ll Q^2$)

$$(\xi = x_B / (2 - x_B) \quad k = -t / 4M^2)$$

Proton Neutron

Polarized beam, unpolarized target:

$$s_{1,unp}^I \sim \sin \phi \operatorname{Im}\{F_1 \mathcal{H} + \xi(F_1 + F_2) \tilde{\mathcal{H}} - kF_2 \mathcal{E}\}$$

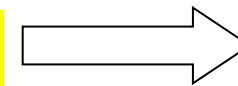


$$\operatorname{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p, \mathcal{E}_p\}$$

$$\operatorname{Im}\{\mathcal{H}_n, \tilde{\mathcal{H}}_n, \mathcal{E}_n\}$$

Unpolarized beam, longitudinal target:

$$s_{1,UL}^I \sim \sin \phi \operatorname{Im}\{F_1 \tilde{\mathcal{H}} + \xi(F_1 + F_2)(\mathcal{H} + x_B / 2 \mathcal{E}) - \xi k F_2 \tilde{\mathcal{E}} + \dots\}$$

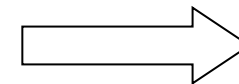


$$\operatorname{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\}$$

$$\operatorname{Im}\{\mathcal{H}_n, \mathcal{E}_n\}$$

Polarized beam, longitudinal target:

$$c_{1,LP}^I \sim (A + B \cos \phi) \operatorname{Re}\{F_1 \tilde{\mathcal{H}} + \xi(F_1 + F_2)(\mathcal{H} + x_B / 2 \mathcal{E}) \dots\}$$

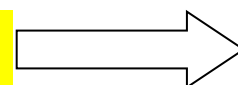


$$\operatorname{Re}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\}$$

$$\operatorname{Re}\{\mathcal{H}_n, \mathcal{E}_n\}$$

Unpolarized beam, transverse target:

$$\Delta \sigma_{UT} \sim \cos \phi \sin(\phi_s - \phi) \operatorname{Im}\{k(F_2 \mathcal{H} - F_1 \mathcal{E}) + \dots\}$$

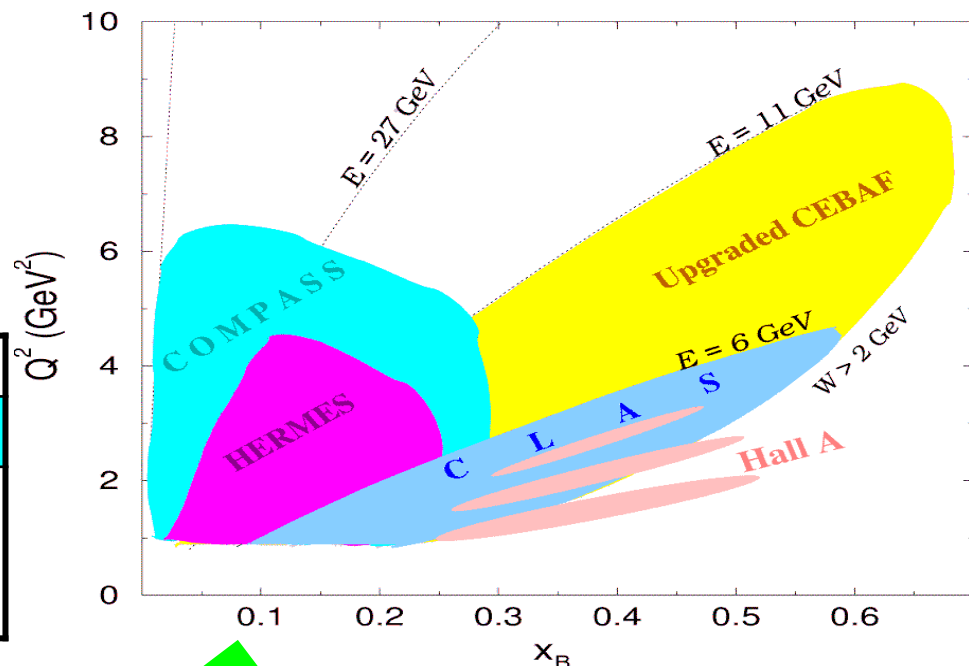


$$\operatorname{Im}\{\mathcal{H}_p, \mathcal{E}_p\}$$

$$\operatorname{Im}\{\mathcal{H}_n\}$$

DVCS experiments worldwide

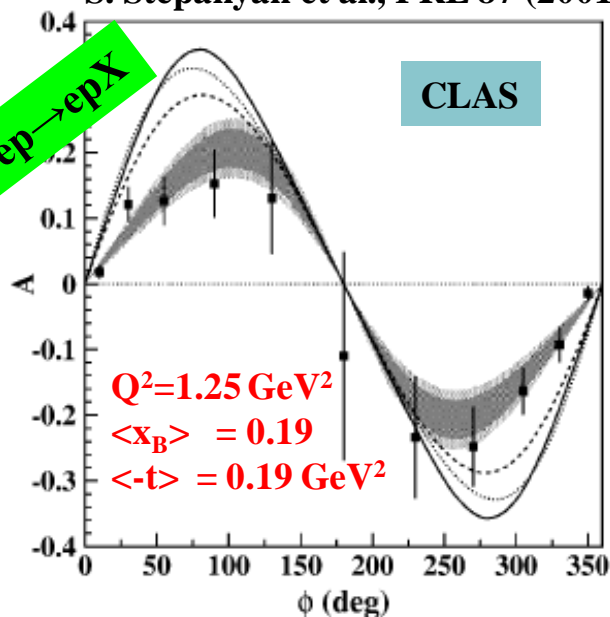
JLAB	
<i>Hall A</i>	<i>CLAS (Hall B)</i>
p,n-DVCS (Bpol.) CS	p-DVCS BSA,ITSA,DSA,CS



DESY	
<i>HERMES</i>	<i>H1/ZEUS</i>
p-DVCS BSA,BCA, tTSA,ITSA,DSA	p-DVCS CS,BCA

CERN
<i>COMPASS</i>
p-DVCS CS,BSA,BCA, tTSA,ITSA,DSA

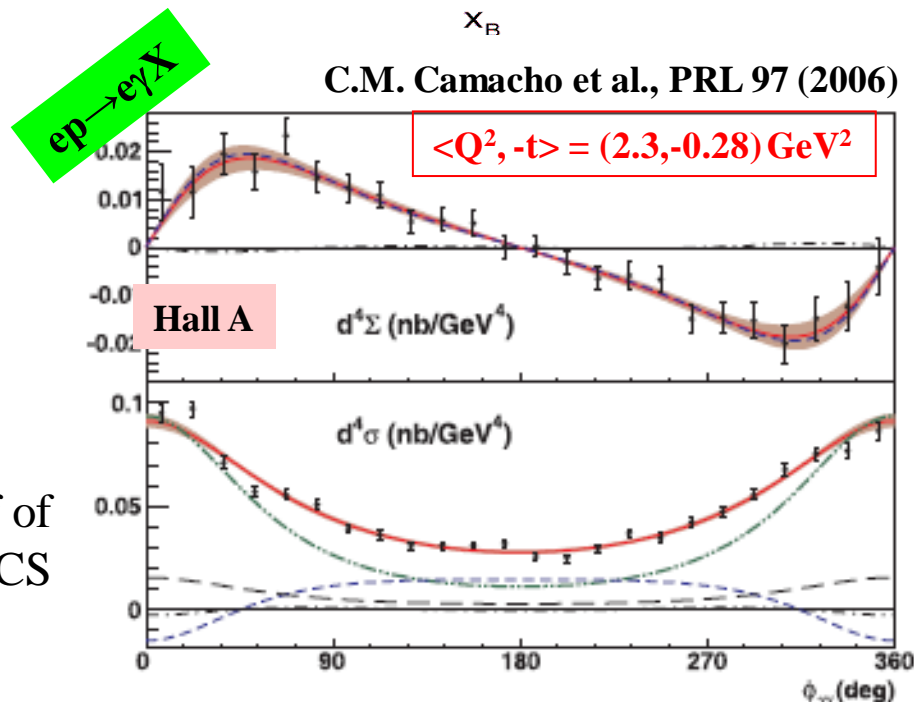
S. Stepanyan et al., PRL 87 (2001)



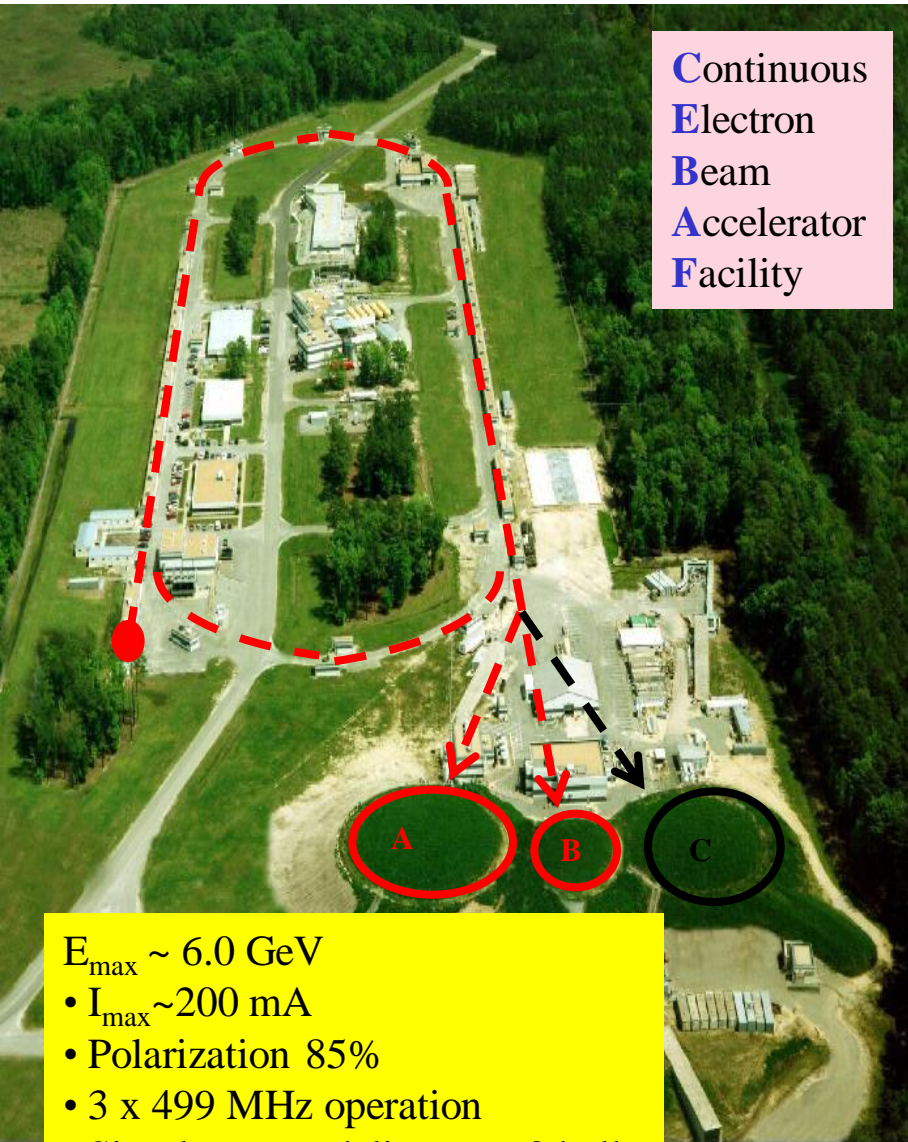
CLAS: first observation of DVCS-BH interference

Hall A: proof of scaling for DVCS

C.M. Camacho et al., PRL 97 (2006)



JLab@6 GeV



$E_{\max} \sim 6.0 \text{ GeV}$

• $I_{\max} \sim 200 \text{ mA}$

• Polarization 85%

• 3 x 499 MHz operation

• Simultaneous delivery to 3 halls

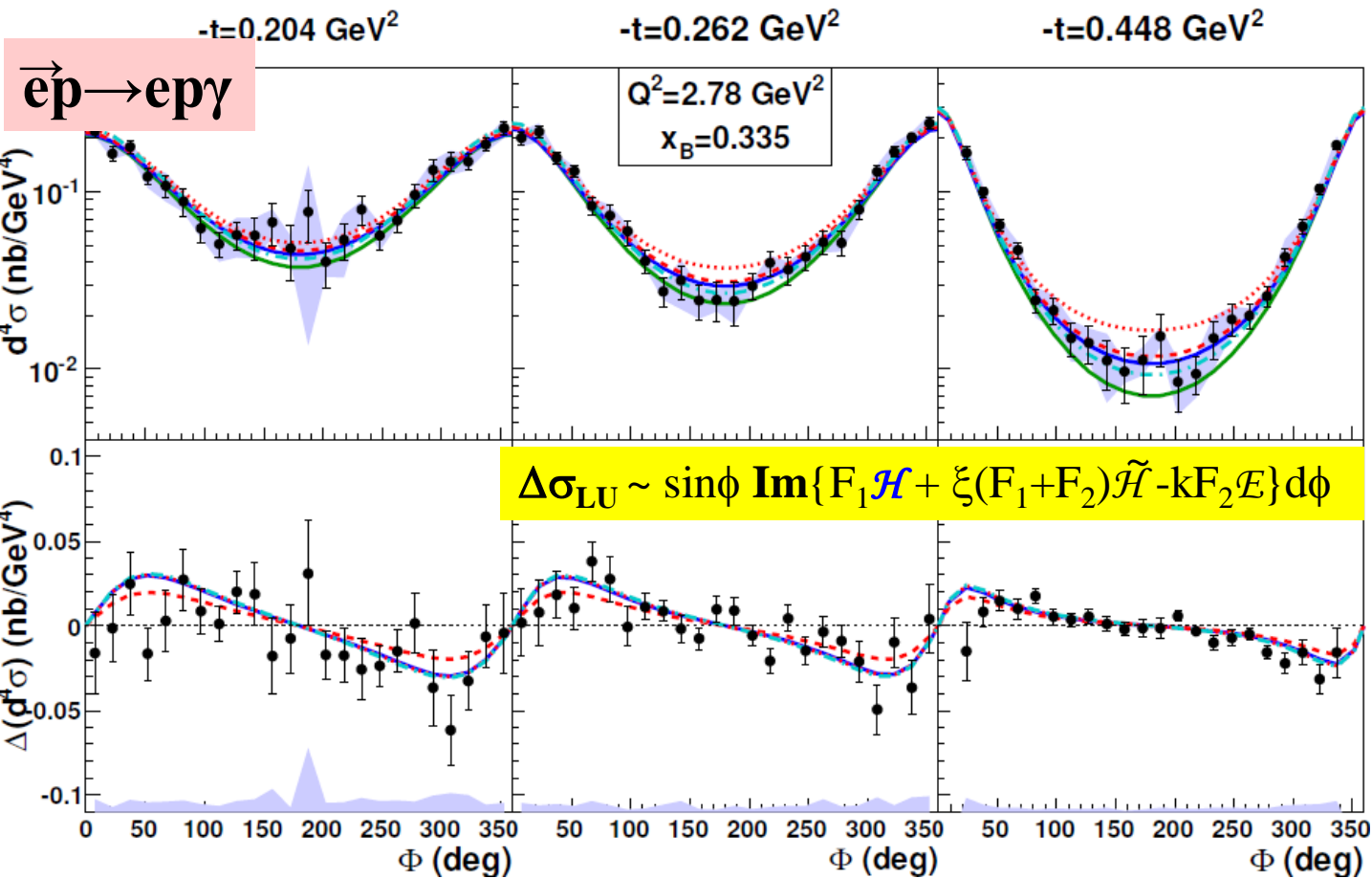
• Shutdown in May 2012



Hall B: CLAS
Large acceptance
Suited for multi-particle
final states
 $L \sim 10^{34}$



CLAS: unpolarized and beam-polarized cross sections



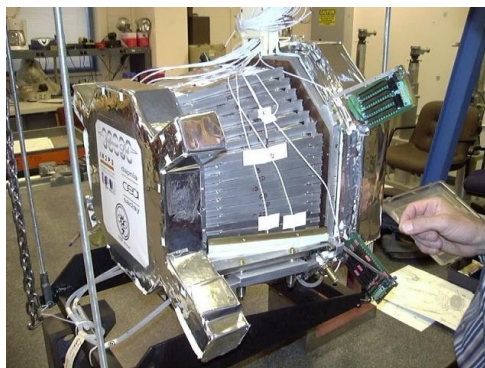
*H.S. Jo et al.,
[arXiv:1504.02009](https://arxiv.org/abs/1504.02009),
submitted to PRL*

- Largest kinematic ever covered
- Two observables extracted

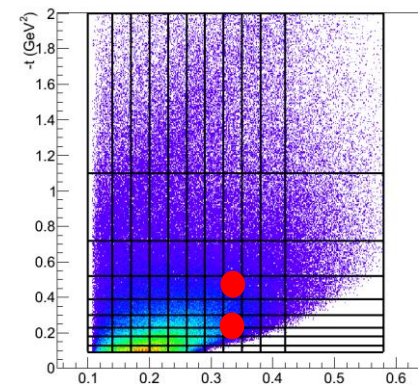
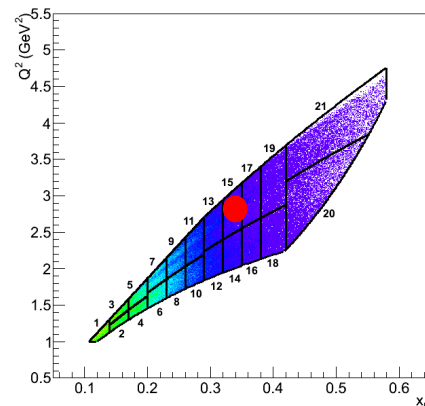
— BH — VGG
..... KM10 --- KM10a

KM10a model fits
Hall A 2006 data using
« anomalous » $\tilde{\mathcal{H}}$

- Data taken in 2005, e1-dvcs
- Beam energy $\sim 5.75 \text{ GeV}$
- Beam polarization $\sim 80\%$
- Target LH_2
- Inner Calorimeter (IC)



21 Q^2 - x_B bins, 9 t bins, 24 ϕ bins



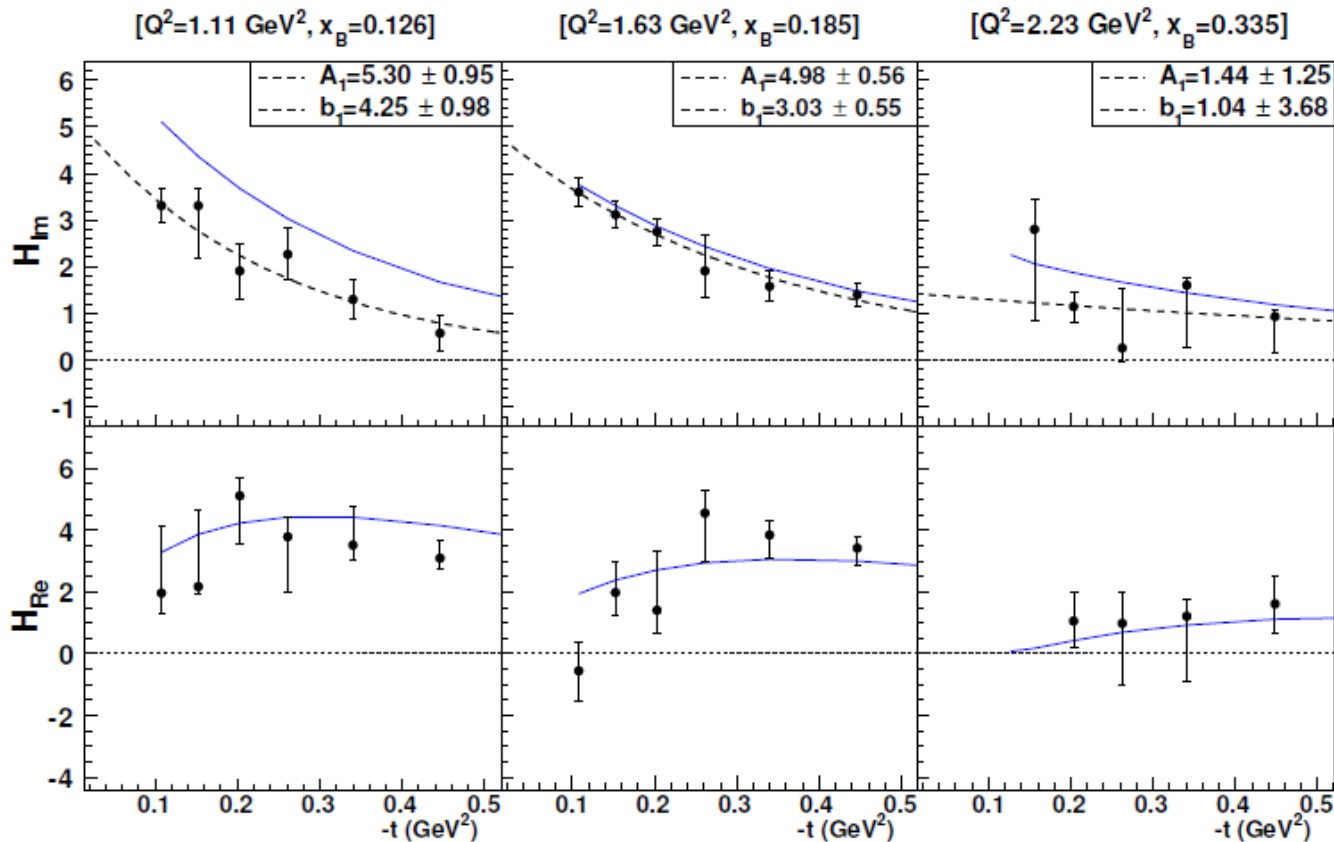
Extraction of Compton Form Factors from DVCS observables

GPDs cannot directly be extracted from DVCS observables, one can access
Compton Form Factors:

$$\begin{array}{l} \text{8 CFF} \left\{ \begin{array}{l} \text{Re}(\mathcal{H}) = P \int_0^1 dx [H(x, \xi, t) - H(-x, \xi, t)] C^+(x, \xi) \\ \text{Re}(\mathcal{E}) = P \int_0^1 dx [E(x, \xi, t) - E(-x, \xi, t)] C^+(x, \xi) \\ \text{Re}(\tilde{\mathcal{H}}) = P \int_0^1 dx [\tilde{H}(x, \xi, t) + \tilde{H}(-x, \xi, t)] C^-(x, \xi) \\ \text{Re}(\tilde{\mathcal{E}}) = P \int_0^1 dx [\tilde{E}(x, \xi, t) + \tilde{E}(-x, \xi, t)] C^-(x, \xi) \\ \text{Im}(\mathcal{H}) = H(\xi, \xi, t) - H(-\xi, \xi, t) \\ \text{Im}(\mathcal{E}) = E(\xi, \xi, t) - E(-\xi, \xi, t) \\ \text{Im}(\tilde{\mathcal{H}}) = \tilde{H}(\xi, \xi, t) - \tilde{H}(-\xi, \xi, t) \\ \text{Im}(\tilde{\mathcal{E}}) = \tilde{E}(\xi, \xi, t) - \tilde{E}(-\xi, \xi, t) \end{array} \right. \\ \text{with } C^\pm(x, \xi) = \frac{1}{x - \xi} \pm \frac{1}{x + \xi} \end{array}$$

M. Guidal: Model-independent fit, at fixed Q^2 , x_B and t of DVCS observables
8 parameters (the CFFs), loosely bound (± 5 x VGG prediction)
M. Guidal, Eur. Phys. J. A 37 (2008) 319 & many other papers...

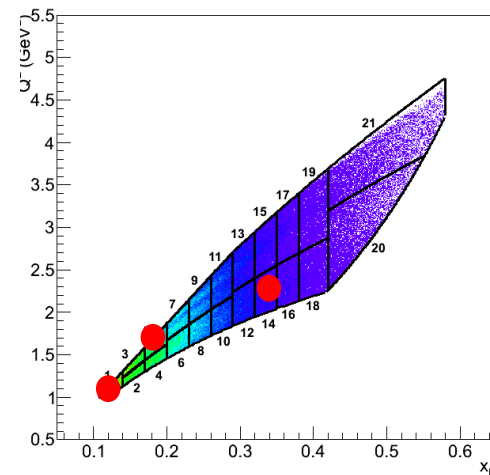
Extraction of CFFs from e1dvc's pol. and unpol. cross sections



*CFF fits by M. Guidal
(H and \tilde{H} only)*

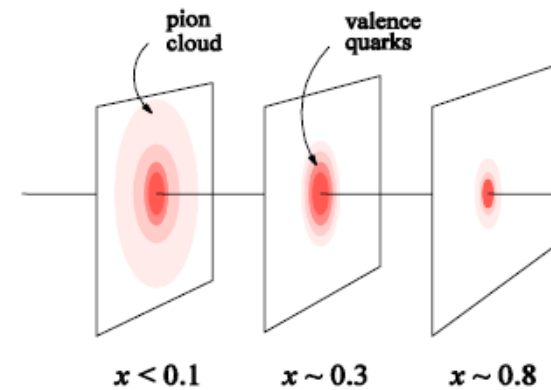
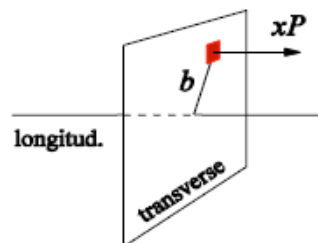
--- Ae^{-bt} fit

— VGG predictions



$$q(x, b_{\perp}) = \int_0^{\infty} \frac{d^2\Delta_{\perp}}{(2\pi)^2} e^{i\Delta_{\perp}b_{\perp}} H(x, 0, -\Delta_{\perp}^2)$$

**$Im(\mathcal{H}_p)$, flatter t slope at high x_B : faster quarks (valence) at the core of the nucleon, slower quarks (sea) at its periphery
→ PROTON TOMOGRAPHY**

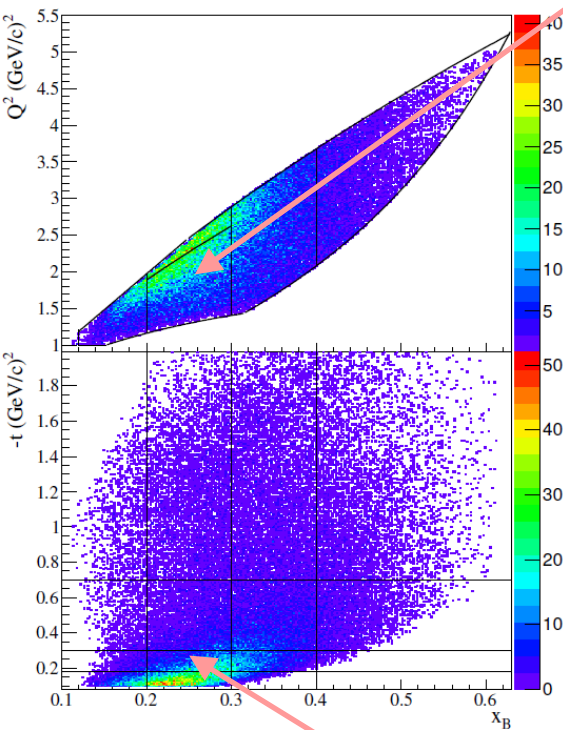


CLAS: DVCS on longitudinally polarized target

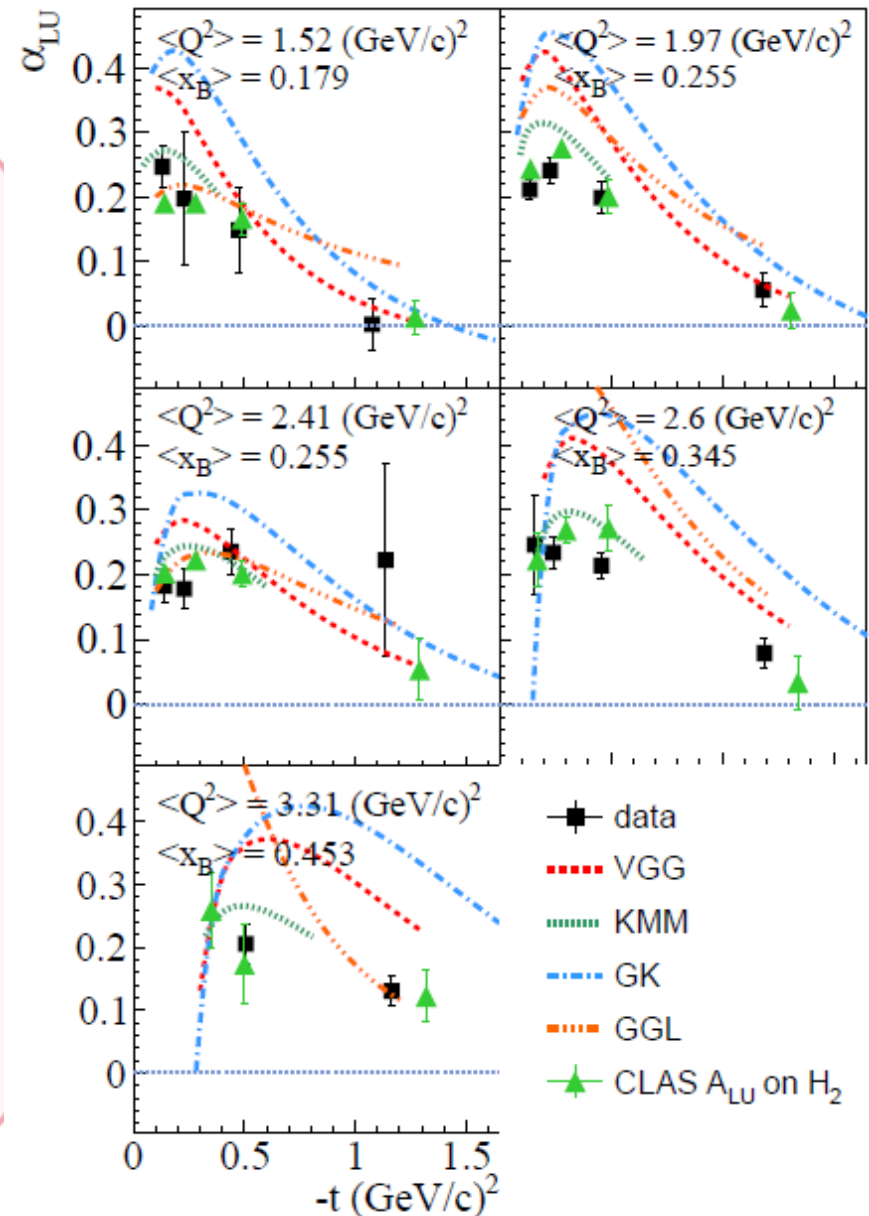
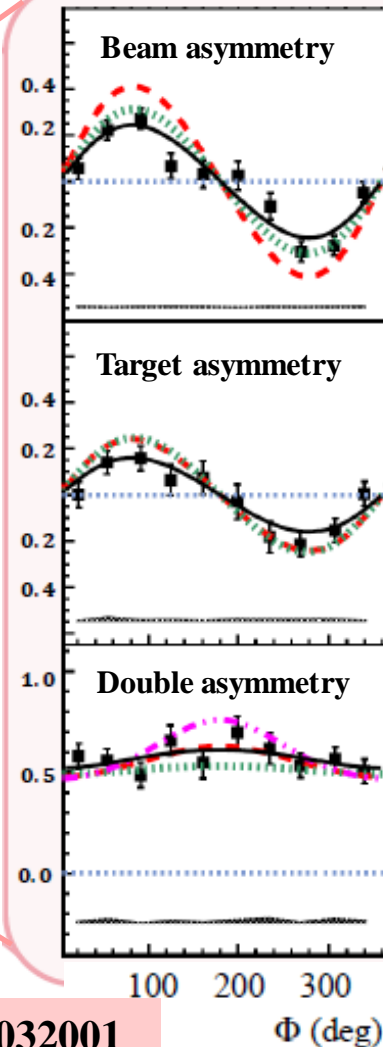
$$\vec{e}p \rightarrow epy$$

$$\text{BSA} \sim \text{Im}\{\mathcal{H}_p\}$$

- Data taken in 2009, eg1-dvcs
- Beam energy ~ 5.9 GeV
- CLAS + IC to detect forward photons
- Target: longitudinally polarized NH_3 ($P \sim 80\%$)
- **3 DVCS observables**



- 5 Q^2 - x_B bins
- 4 t bins
- 10 ϕ bins

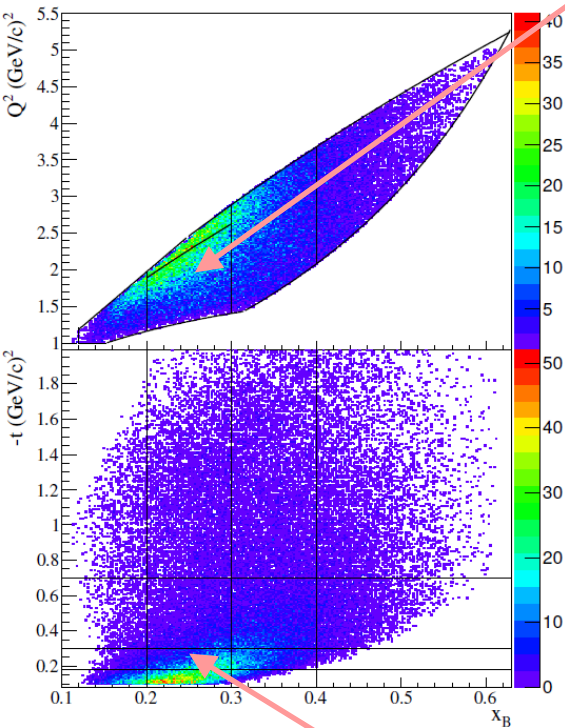


CLAS: DVCS on longitudinally polarized target

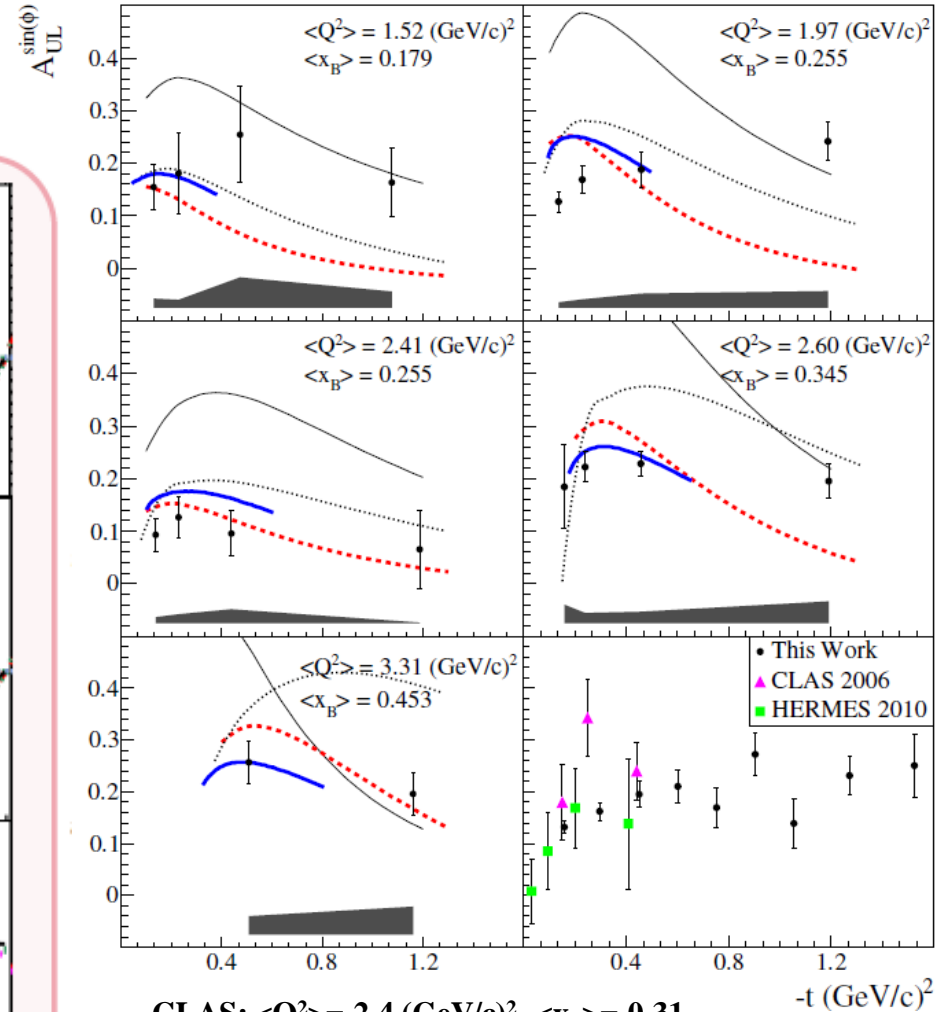
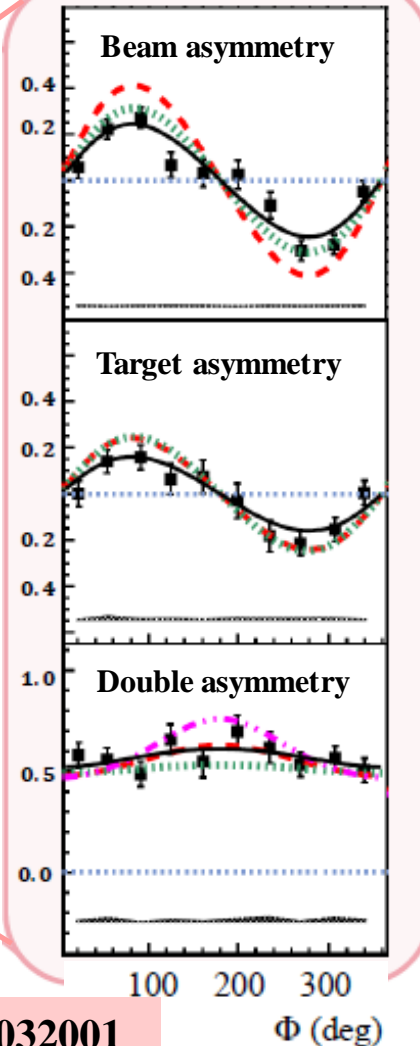
$$\vec{e}\vec{p} \rightarrow e\gamma$$

- Data taken in 2009, eg1-dvcs
- Beam energy ~ 5.9 GeV
- CLAS + IC to detect forward photons
- Target: longitudinally polarized NH_3 ($P \sim 80\%$)
- **3 DVCS observables**

$$\text{TSA} \sim \text{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\}$$



- 5 Q^2 - x_B bins
- 4 t bins
- 10 ϕ bins



CLAS: $\langle Q^2 \rangle = 2.4$ (GeV/c^2), $\langle x_B \rangle = 0.31$

HERMES: $\langle Q^2 \rangle = 2.459$ (GeV/c^2), $\langle x_B \rangle = 0.096$

CLAS2006: $\langle Q^2 \rangle = 1.82$ (GeV/c^2), $\langle x_B \rangle = 0.28$

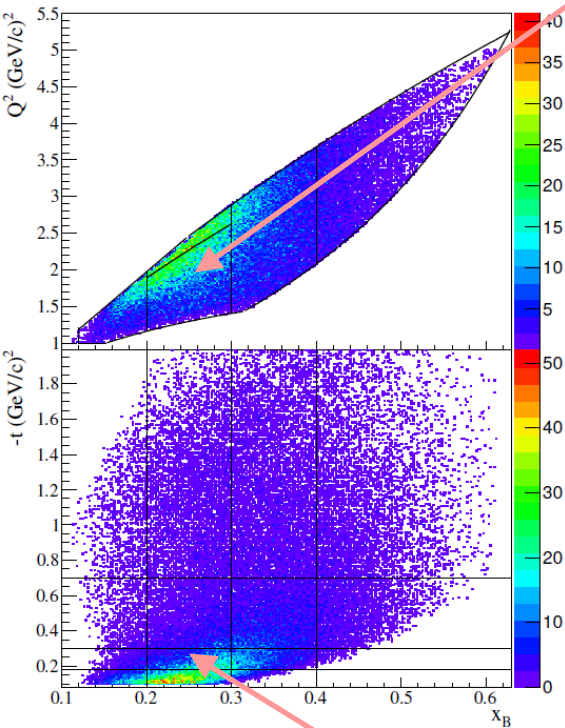
- Improved statistics x10 at low $-t$
- Extended kinematic coverage

CLAS: DVCS on longitudinally polarized target

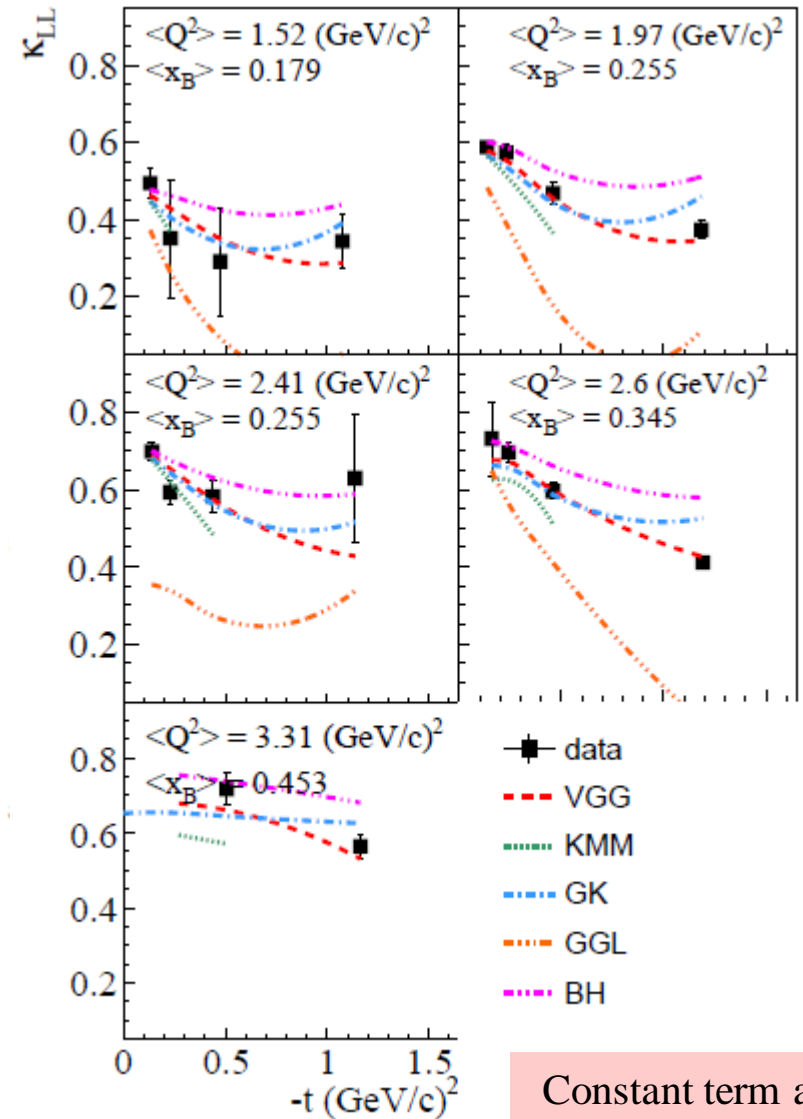
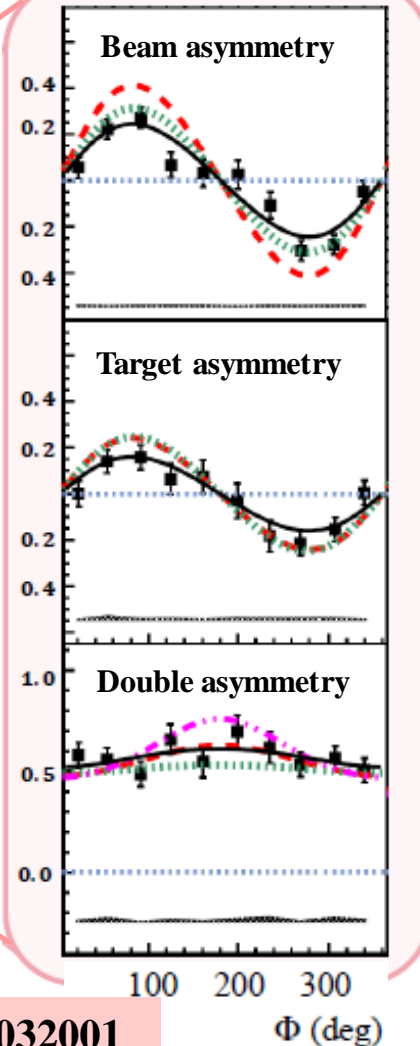
$$\vec{e}p^{\rightarrow} \rightarrow e\gamma$$

- Data taken in 2009, eg1-dvcs
- Beam energy ~ 5.9 GeV
- CLAS + IC to detect forward photons
- Target: longitudinally polarized NH_3 ($P \sim 80\%$)
- **3 DVCS observables**

$$\text{DSA} \sim \text{Re}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\}$$



- 5 Q^2 - x_B bins
- 4 t bins
- 10 ϕ bins



Constant term and $\cos\phi$ term are dominated by BH

Extraction of CFFs from eg1-dvcs TSA, BSA, DSA

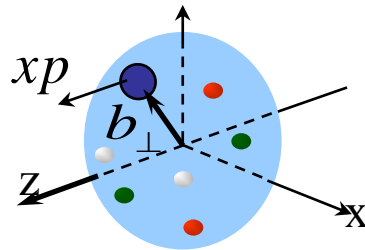
CFFs fitting code by M. Guidal
(7 CFFs included in the fit)

$Im\mathcal{H}$ has steeper t-slope than $Im\tilde{\mathcal{H}}$: is axial charge more “concentrated” than the electromagnetic charge?
→ PROTON TOMOGRAPHY

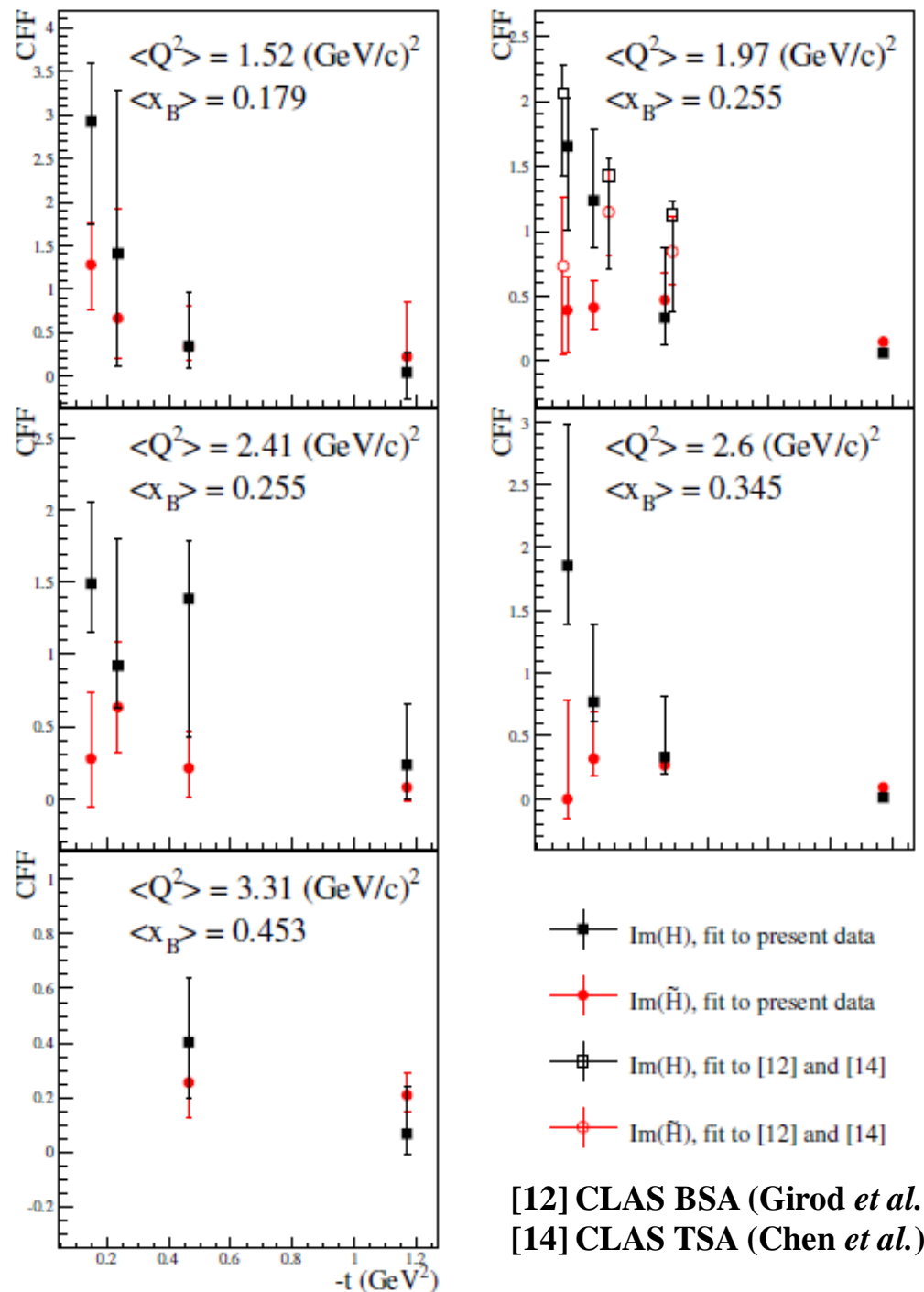
$$\Delta q(x, b_{\perp}) = \int_0^{\infty} \frac{d^2\Delta_{\perp}}{(2\pi)^2} e^{i\Delta_{\perp} b_{\perp}} \tilde{H}(x, 0, -\Delta_{\perp}^2)$$

$$\int H(x, \xi, t) dx = F_1(t)$$

$$\int \tilde{H}(x, \xi, t) dx = G_A(t)$$



S. Pisano et al., PRD 91, 052014 (2015)



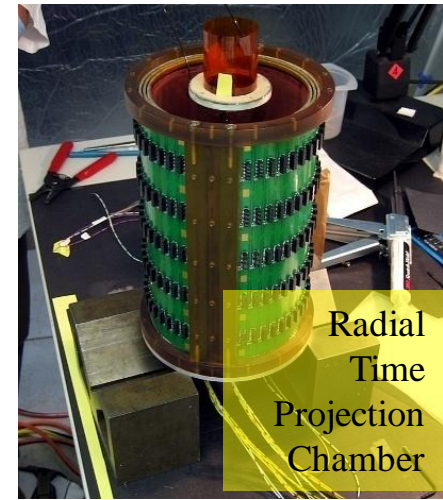
[12] CLAS BSA (Girod *et al.*)
[14] CLAS TSA (Chen *et al.*)

DVCS on nuclei: the CLAS eg6 experiment

- Data taken in the fall 2009
- CLAS+IC+RTPC+⁴He target; E~6.065 GeV
- **Coherent** and **incoherent** DVCS: **nuclear GPDs**, **EMC effect**

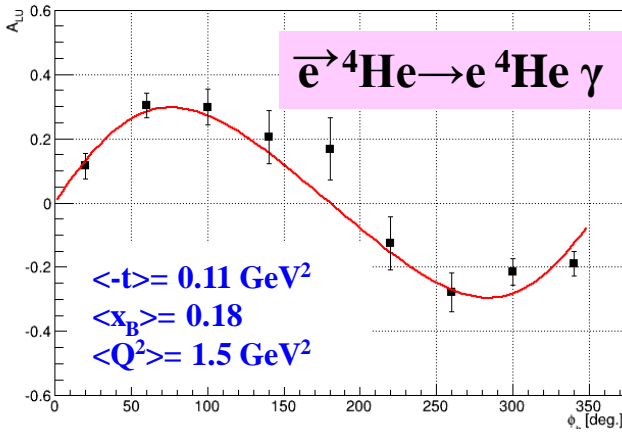
⁴He is a spin-0 nucleus: at twist-2 only one CFF in DVCS BSA

$$A_{LU}^{4\text{He}}(\varphi) = \frac{\alpha_0(\varphi) F_A(t) \Im m[\mathcal{H}_A]}{\alpha_1(\varphi) F_A^2(t) + \alpha_2(\varphi) F_A(t) \Re e[\mathcal{H}_A] + \alpha_3(\varphi) \Re e[\mathcal{H}_A]^2 + \alpha_3(\varphi) \Im m[\mathcal{H}_A]^2}$$

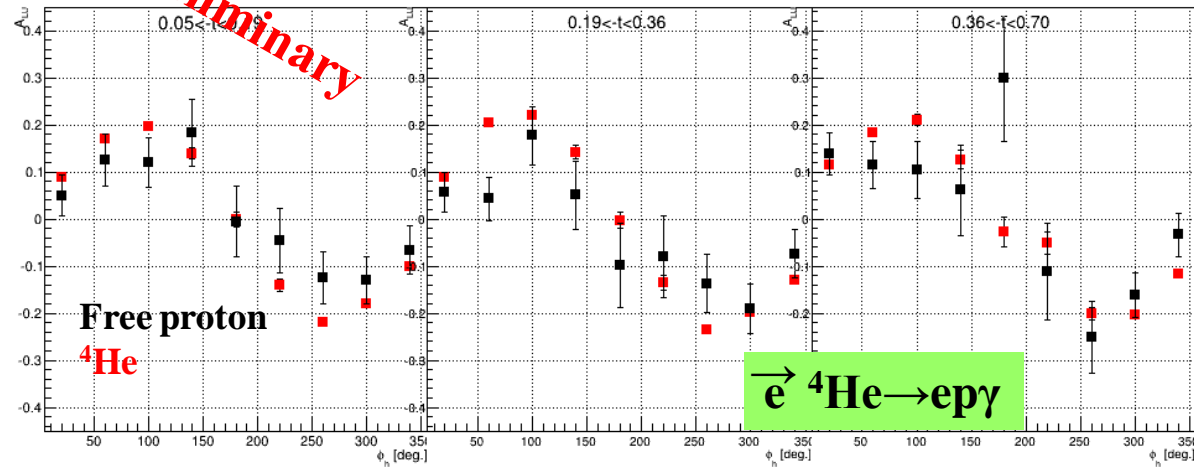
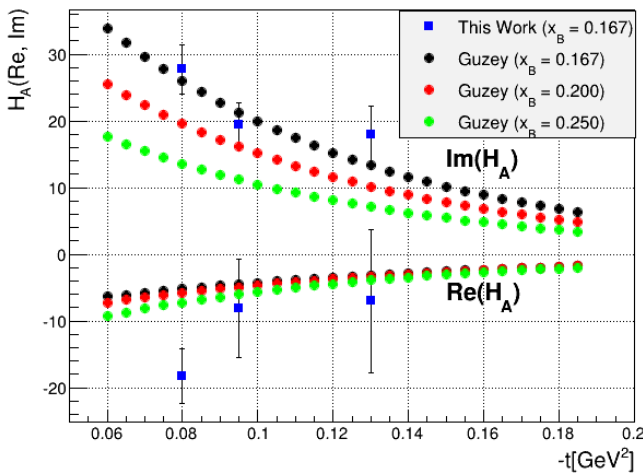


Work by M. Hattawy, IPNO

CLAS preliminary



H_A vs. -t

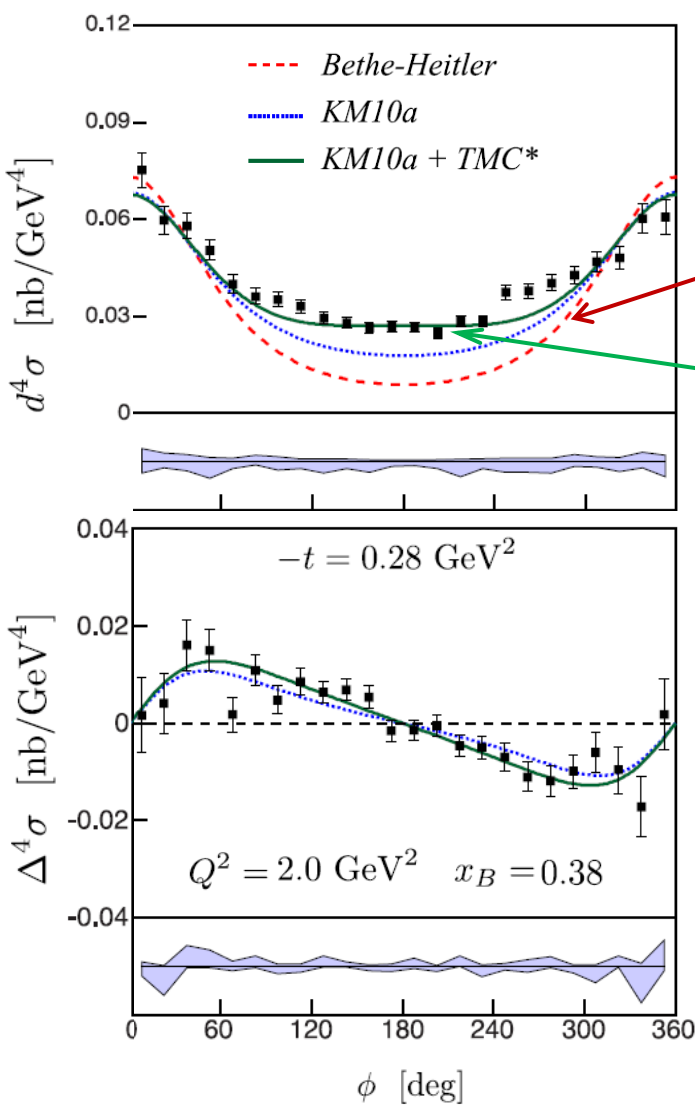


- **Small -t**: asymmetry for bound proton lower than the ⁴He one
- **High -t**: the two asymmetries tend to become compatible

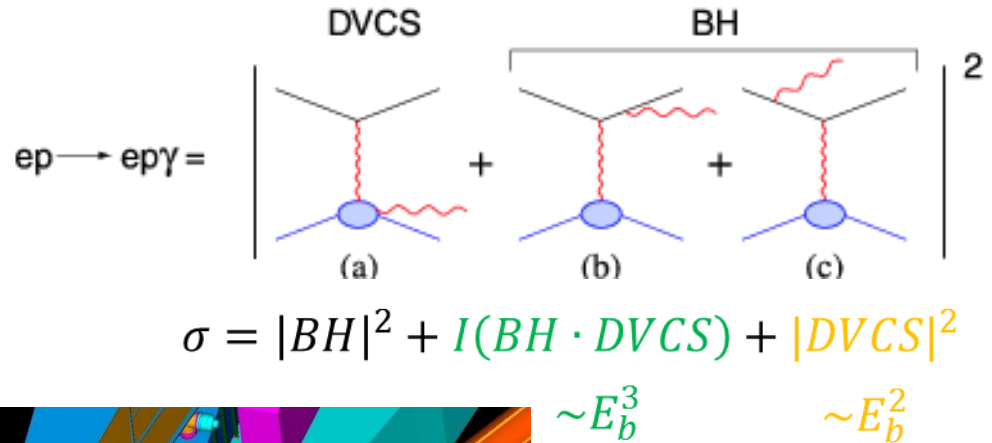
DVCS on the proton in Hall A

Results from E00-110

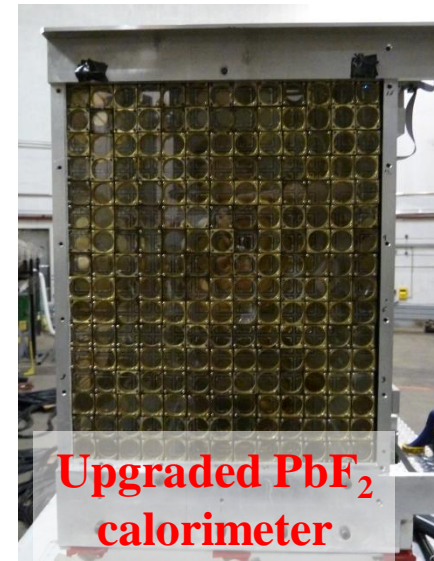
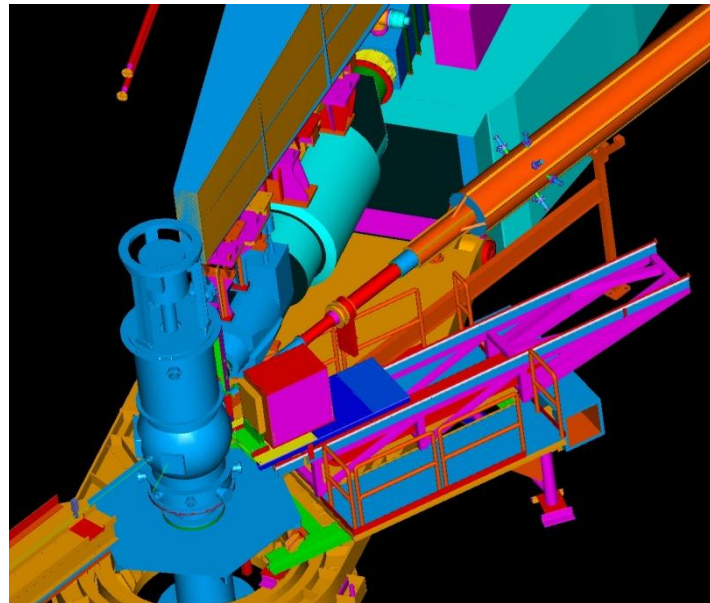
M. Defurne et. al., arXiv 1504.05453 (2015)



- Significant deviation from Bethe-Heitler (BH)
- Both $I(\text{BH} \cdot \text{DVCS})$ and DVCS^2 contribute to the cross section
- Twist-4 corrections (TMC) may be necessary to fully describe data

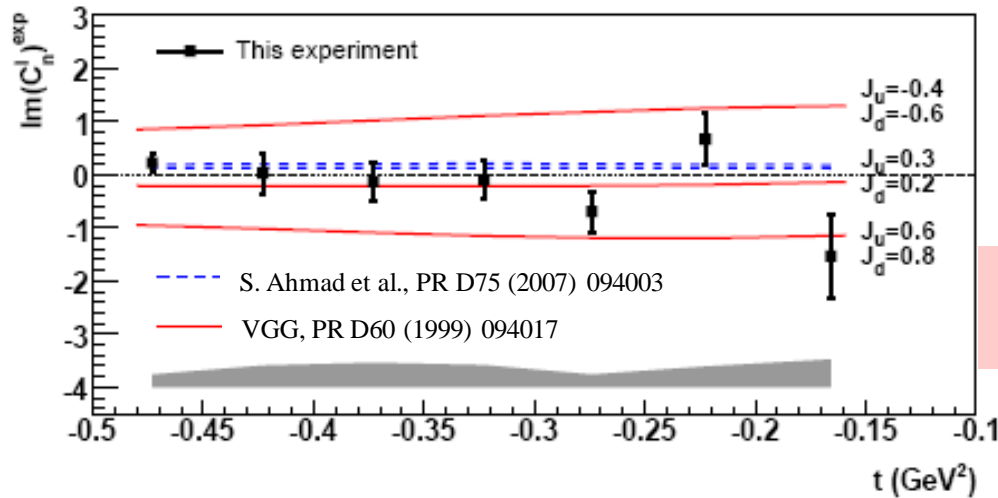


Beam-energy separation
at constant Q^2 , x_B and t :
experiment E07-007
(Analysis ongoing)



DVCS on the neutron in Hall A

M. Mazouz et al., PRL 99 (2007) 242501



Proton and neutron GPDs (and CFFs) are **linear combinations** of quark GPDs

$$\mathcal{H}_p(\xi, t) = \frac{4}{9} \mathcal{H}_u(\xi, t) + \frac{1}{9} \mathcal{H}_d(\xi, t)$$

$$\mathcal{H}_n(\xi, t) = \frac{1}{9} \mathcal{H}_u(\xi, t) + \frac{4}{9} \mathcal{H}_d(\xi, t)$$

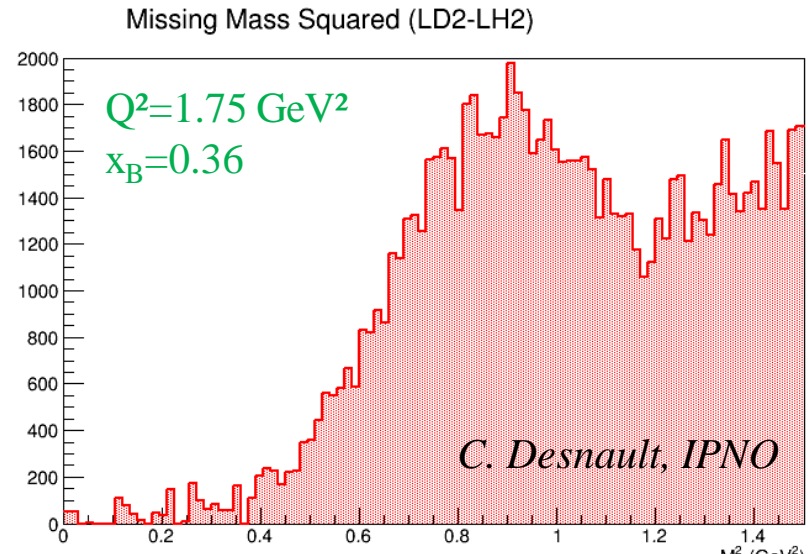
A **combined analysis** of DVCS observables for **proton and neutron targets** is necessary to perform a **quark-flavor separation** of the GPDs

$$\Delta\sigma_{LU} \sim \sin\phi \operatorname{Im}\{F_1\mathcal{H} + \xi(F_1+F_2)\tilde{\mathcal{H}} - kF_2\mathcal{E}\}$$

$$\frac{1}{2} \int_{-1}^1 dx (H(x, \xi, t=0) + E(x, \xi, t=0)) = J$$

• **E03-106: First-time measurement of $\Delta\sigma_{LU}$ for nDVCS, model-dependent extraction of J_u, J_d**

- **E08-025: Beam-energy « Rosenbluth » separation of nDVCS CS using an LD2 target**
- Data taken in fall 2010 concurrently with E07-007, **analysis ongoing**

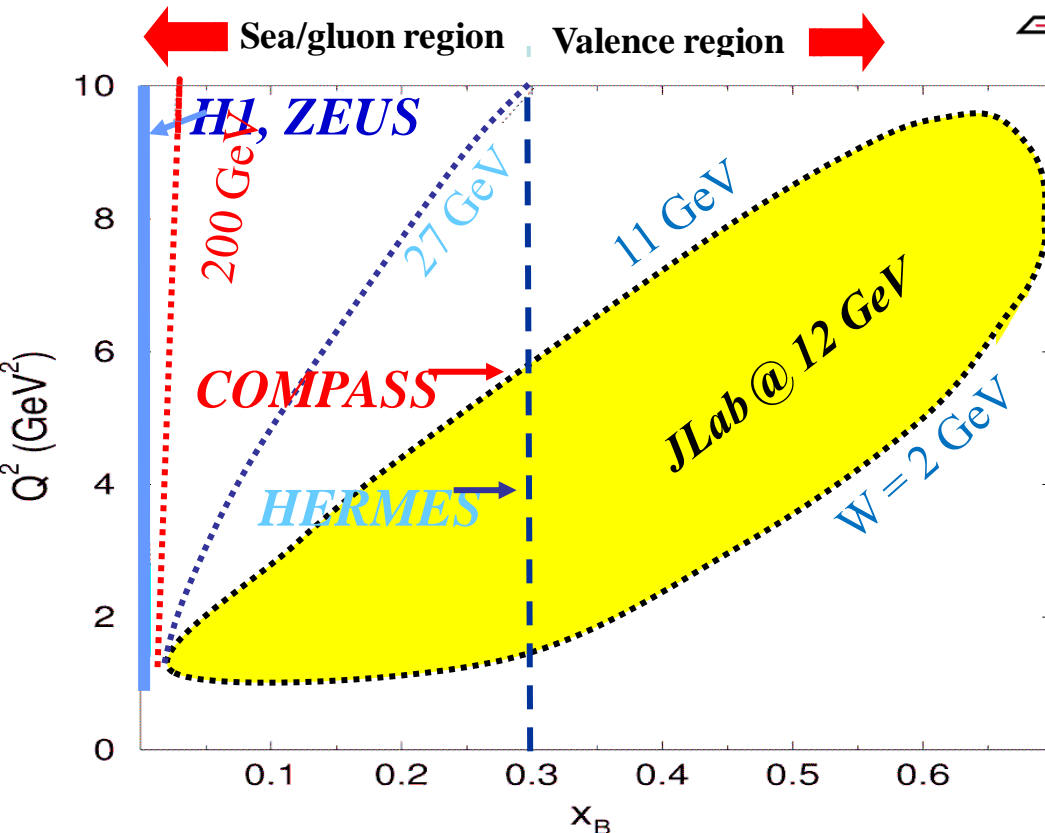
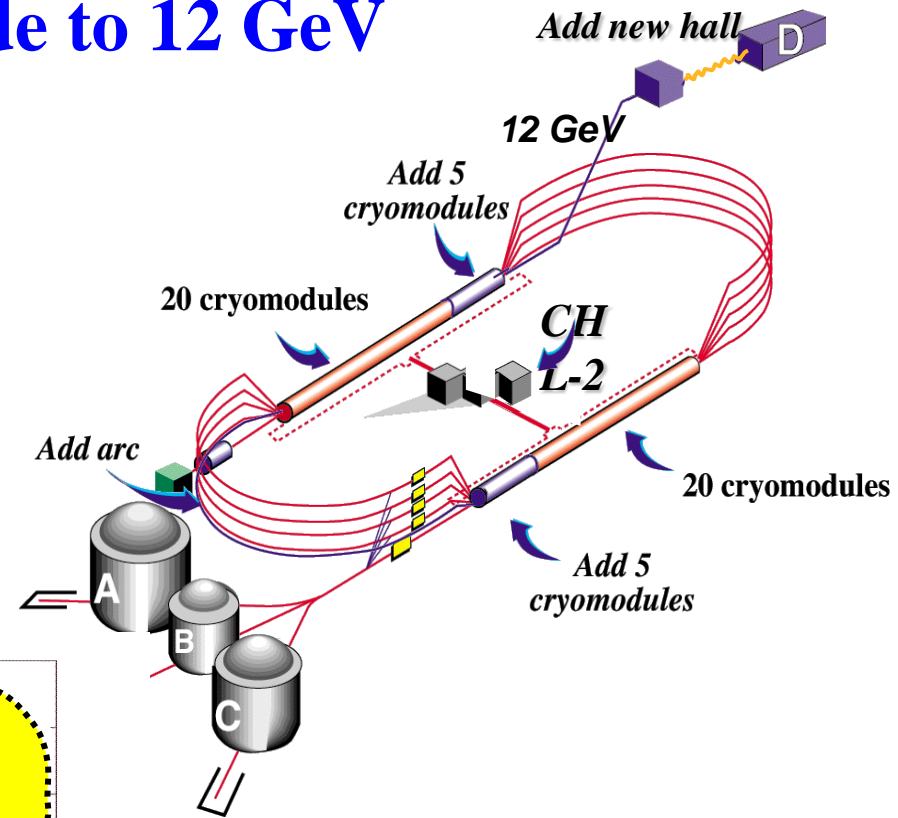


JLab upgrade to 12 GeV

$E = 2.2, 4.4, 6.6, 8.8, 11$ GeV
for the Halls A, B, C

Beam polarization $> 80\%$

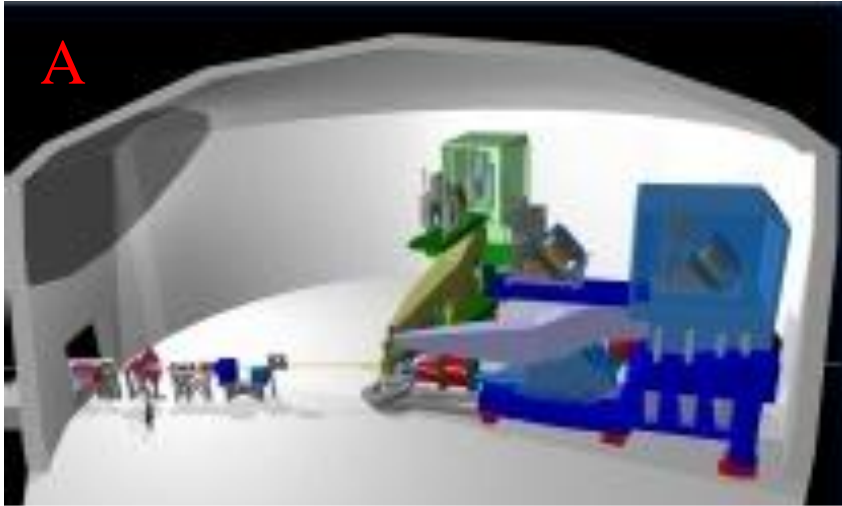
Upgrade completed in 2014



Study of high x_B domain
requires high luminosity

The 12-GeV upgrade is
well matched to studies in
the valence-quark regime

New capabilities in Halls A, B & C



**High Resolution Spectrometer (HRS) pair
and specialized large installation experiments**



CLAS12: large acceptance, high luminosity



**Super High Momentum Spectrometer (SHMS)
at high luminosity and forward angles**

DVCS experiments at 11 GeV have been approved for each of these **three halls**.

Complementary programs:

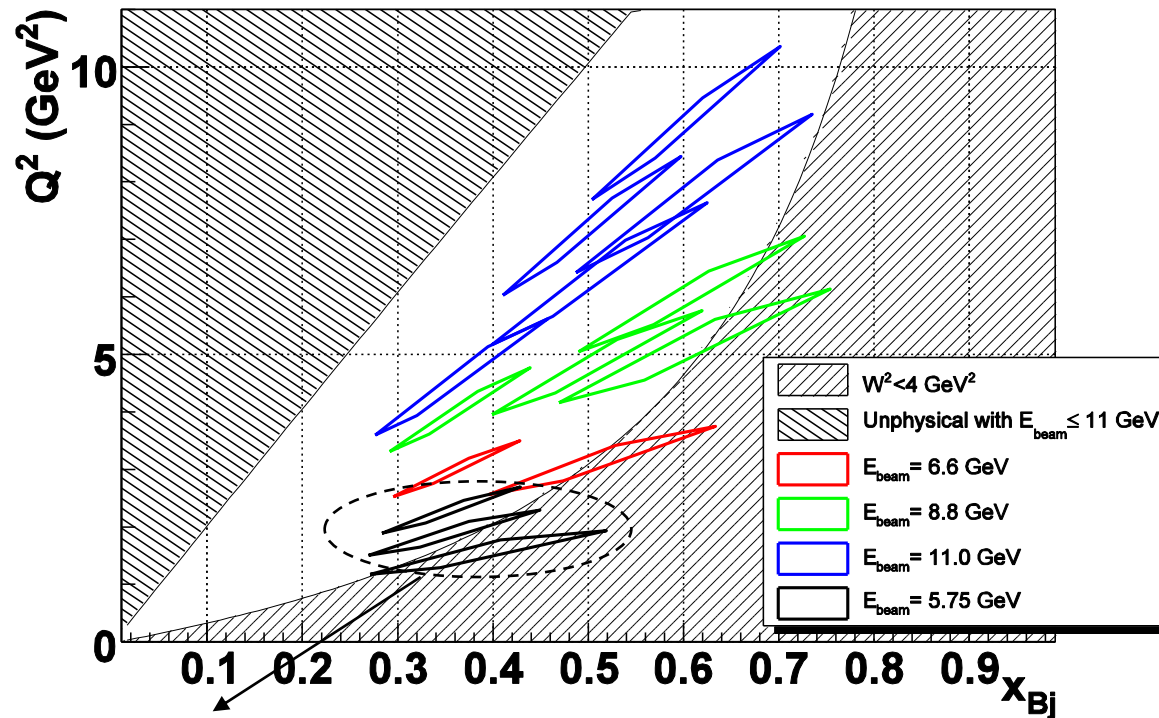
- different kinematic coverage
- different precisions/resolutions
- focus on different observables

E12-06-114: DVCS at 11 GeV in Hall A

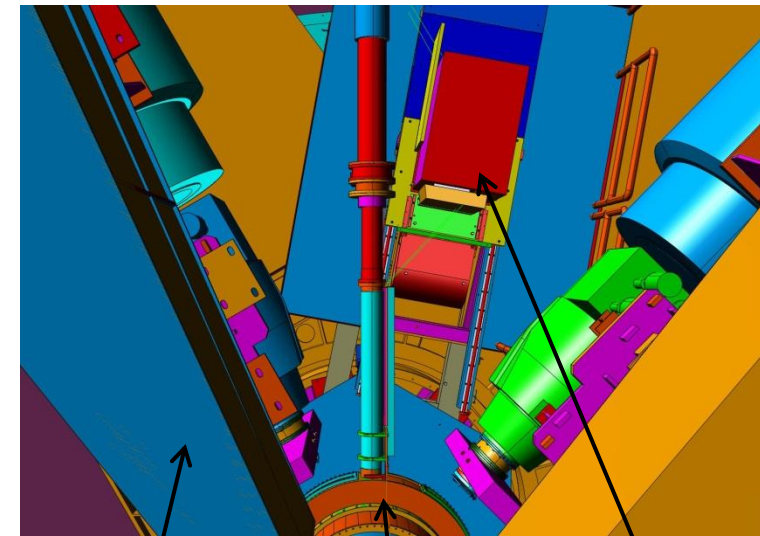
- Absolute cross-section measurements
- Test of scaling: Q^2 dependence of $d\sigma$ at fixed x_{Bj}
- Increased kinematical coverage

JLab12 with 3, 4, 5 pass beam (6.6, 8.8, 11.0 GeV)

DVCS measurements in Hall A/JLab



JLab @ 6 GeV



L-HRS

Scattering chamber

PbF2 calorimeter

1st experiment to run after the 12-GeV upgrade

Started in fall 2014, to be continued in 2015-2016

E12-13-010: DVCS at 11 GeV in Hall C

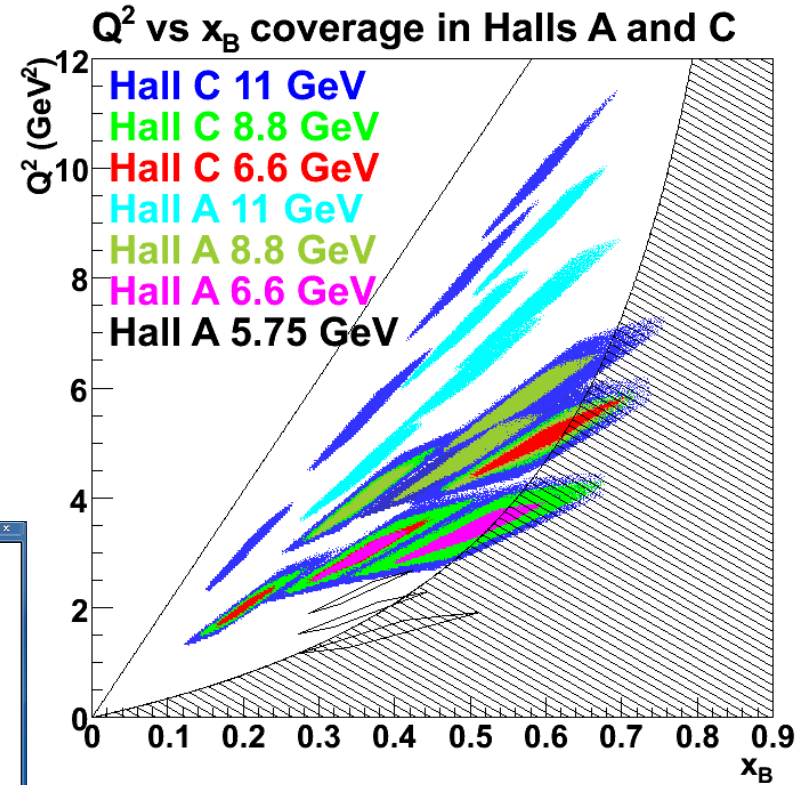
- Energy separation of the DVCS cross section
- Higher Q^2 : measurement of higher twist contributions
- Low- x_B extension (thanks to sweeping magnet)

Sweeping magnet

1116-block PbWO_4 calorimeter

Hall C
HMS

Tentative running:
~ 2019-20



DVCS BSA and TSA with CLAS12 & 11 GeV beam

85 days of beam time

$P_{\text{beam}} = 85\%$

$L = 10^{35} \text{ cm}^{-2}\text{s}^{-1}$

$1 < Q^2 < 10 \text{ GeV}^2$

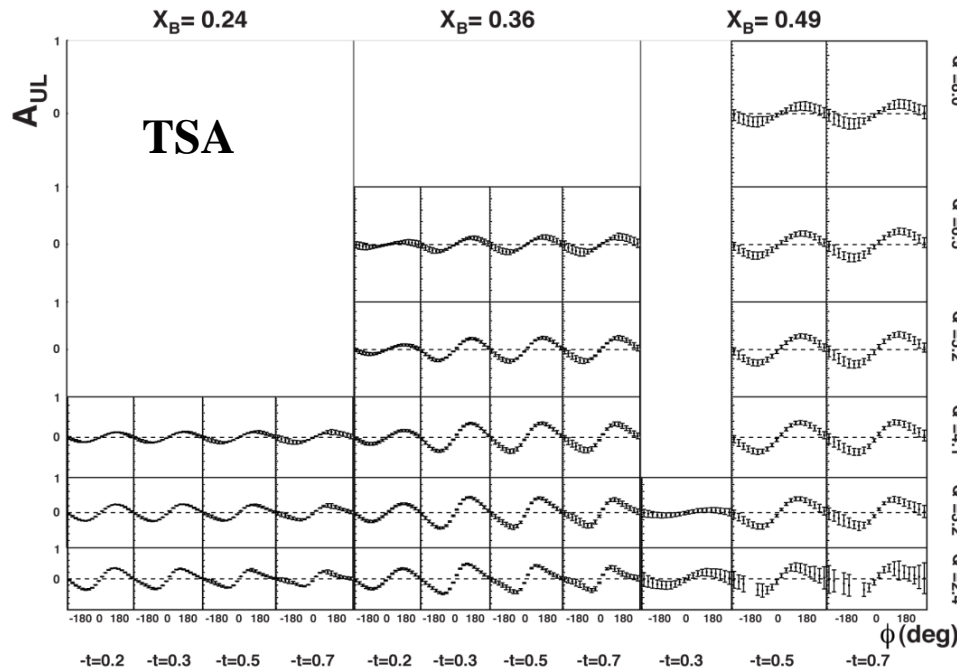
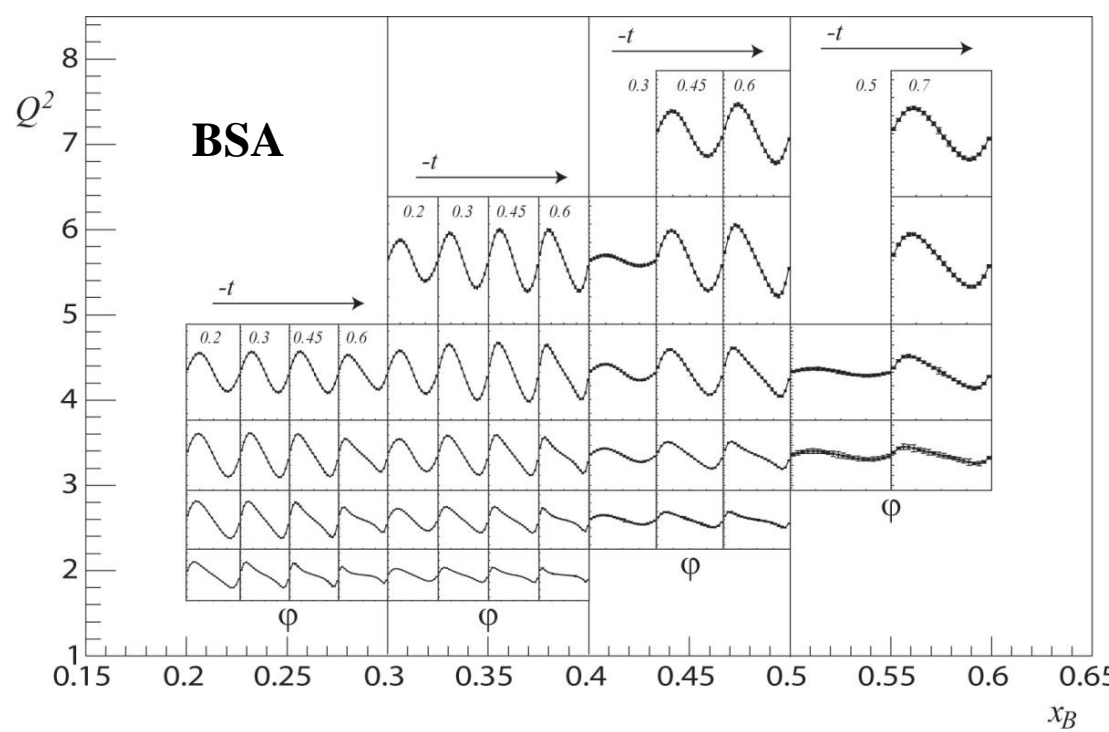
$0.1 < x_B < 0.65$

$-\text{t}_{\text{min}} < -\text{t} < 2.5 \text{ GeV}^2$

Statistical error: 1% to 10%

on $\sin\phi$ moments

Systematic uncertainties: $\sim 6\text{-}8\%$



120 days of beam time

$P_{\text{beam}} = 85\%$, $P_{\text{target}} = 80\%$

$L = 2.10^{35} \text{ cm}^{-2}\text{s}^{-1}$

$1 < Q^2 < 10 \text{ GeV}^2$

$0.1 < x_B < 0.65$

$-\text{t}_{\text{min}} < -\text{t} < 2.5 \text{ GeV}^2$

Statistical error: 2% to 15%

on $\sin\phi$ moments

Systematic uncertainties: $\sim 6\text{-}8\%$

DVCS BSA and TSA with CLAS12 & 11 GeV beam

85 days of beam time

$P_{\text{beam}} = 85\%$

$L = 10^{35} \text{ cm}^{-2}\text{s}^{-1}$

$1 < Q^2 < 10 \text{ GeV}^2$

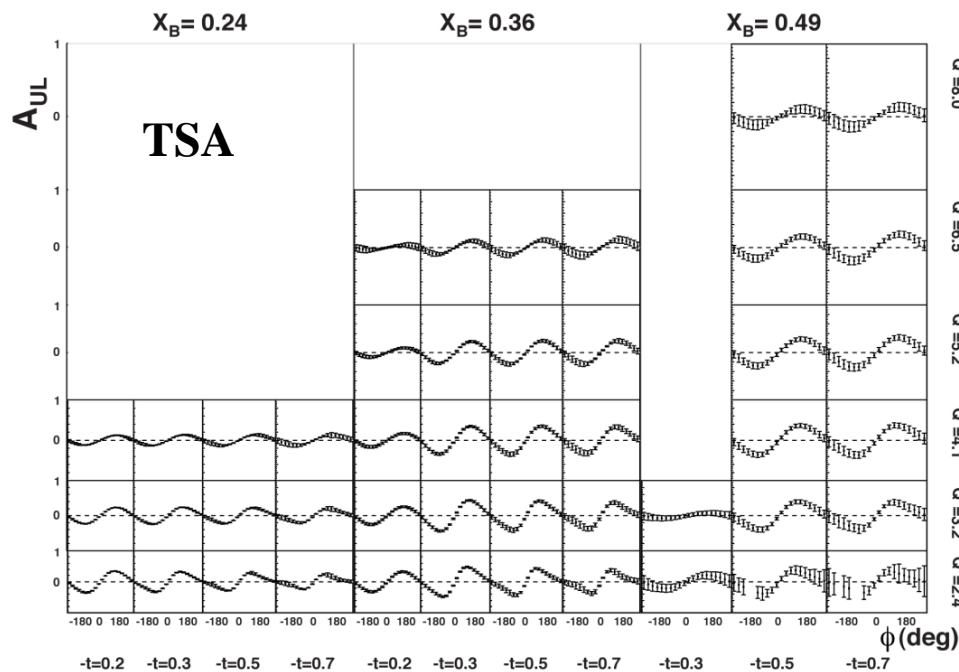
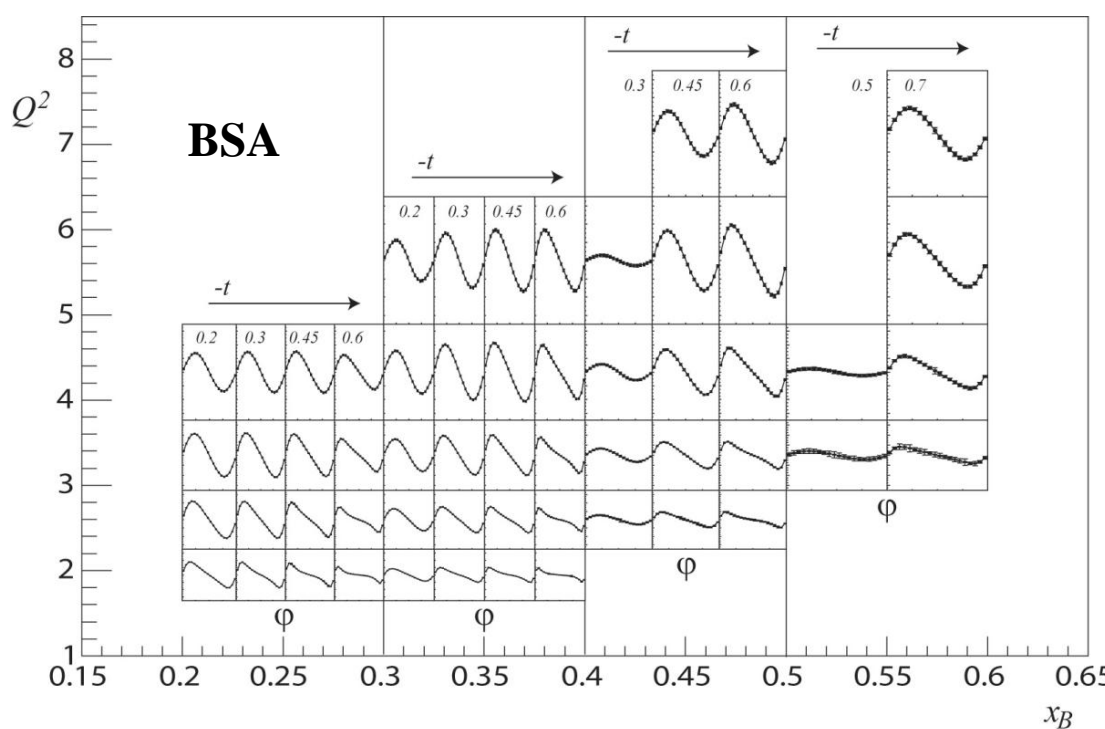
$0.1 < x_B < 0.65$

$-t_{\text{min}} < -t < 2.5 \text{ GeV}^2$

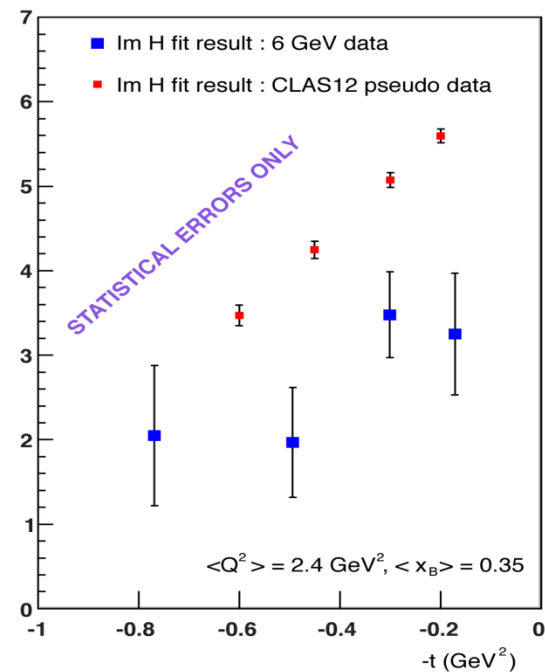
Statistical error: 1% to 10%

on $\sin\phi$ moments

Systematic uncertainties: $\sim 6\text{-}8\%$



Impact of CLAS12 DVCS-BSA data on model-independent fit to extract $\text{Im}(H)$



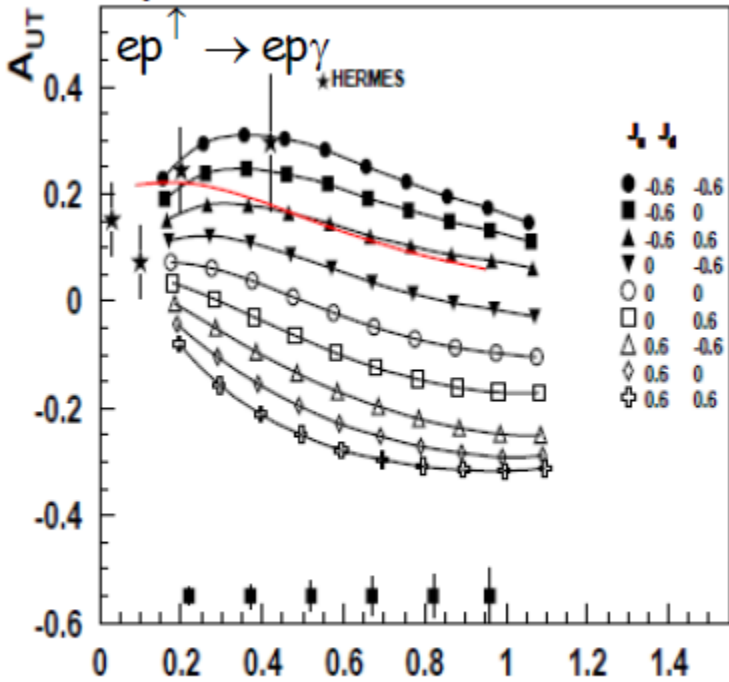
CLAS12: p-DVCS *transverse* target-spin asymmetry

100 days of beam time

Beam pol. = 80% ; target pol. (HDIce) = 60% ; Luminosity = $5 \times 10^{33} \text{ cm}^{-2}\text{s}^{-1}$

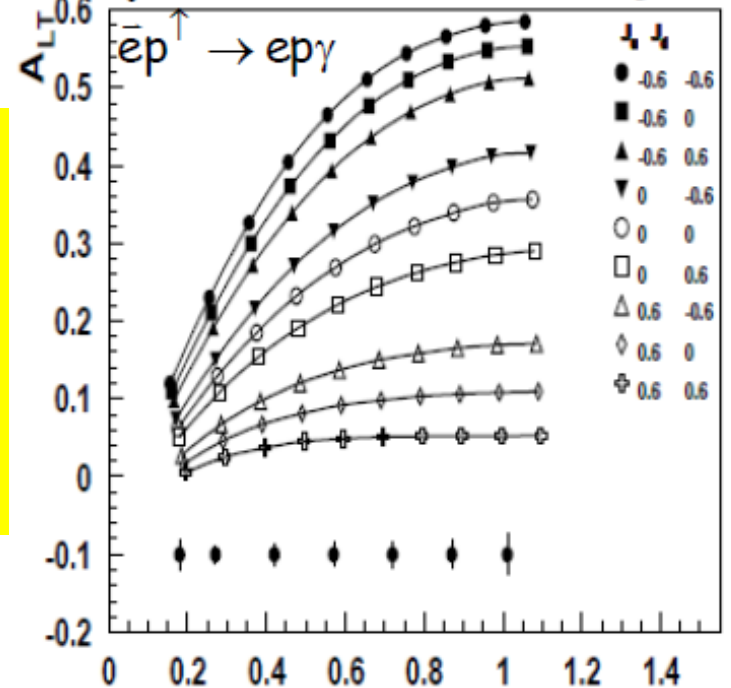
$1 < Q^2 < 10 \text{ GeV}^2$, $0.06 < x_B < 0.66$, $-t_{\text{min}} < -t < 1.5 \text{ GeV}^2$

Projections for $Q^2=2.5 \text{ GeV}^2$, $x_B = 0.2$



Transverse-target spin asymmetry for p-DVCS is **highly sensitive** to the **u-quark** contributions to proton spin.

Projections for $Q^2=2.5 \text{ GeV}^2$, $x_B = 0.2$



$$\Delta\sigma_{\text{UT}} \Rightarrow \text{Im}\{\mathcal{H}_p, \mathcal{E}_p\}$$

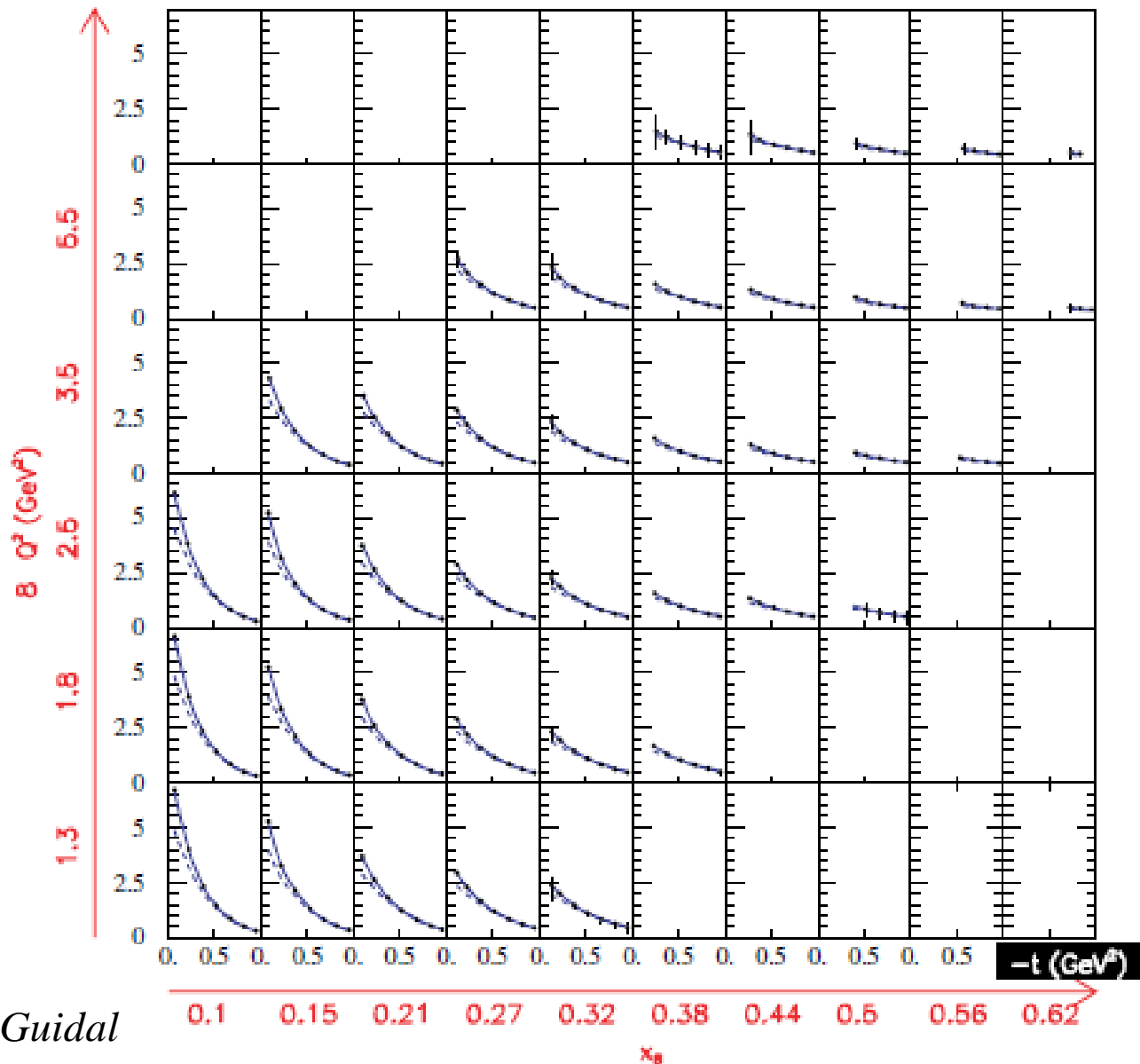
$$\Delta\sigma_{\text{LT}} \Rightarrow \text{Re}\{\mathcal{H}_p, \mathcal{E}_p\}$$

JLab PAC:
high-impact
experiment

Proposal conditionally approved by PAC39

Projections for CLAS12 for H_{Im}

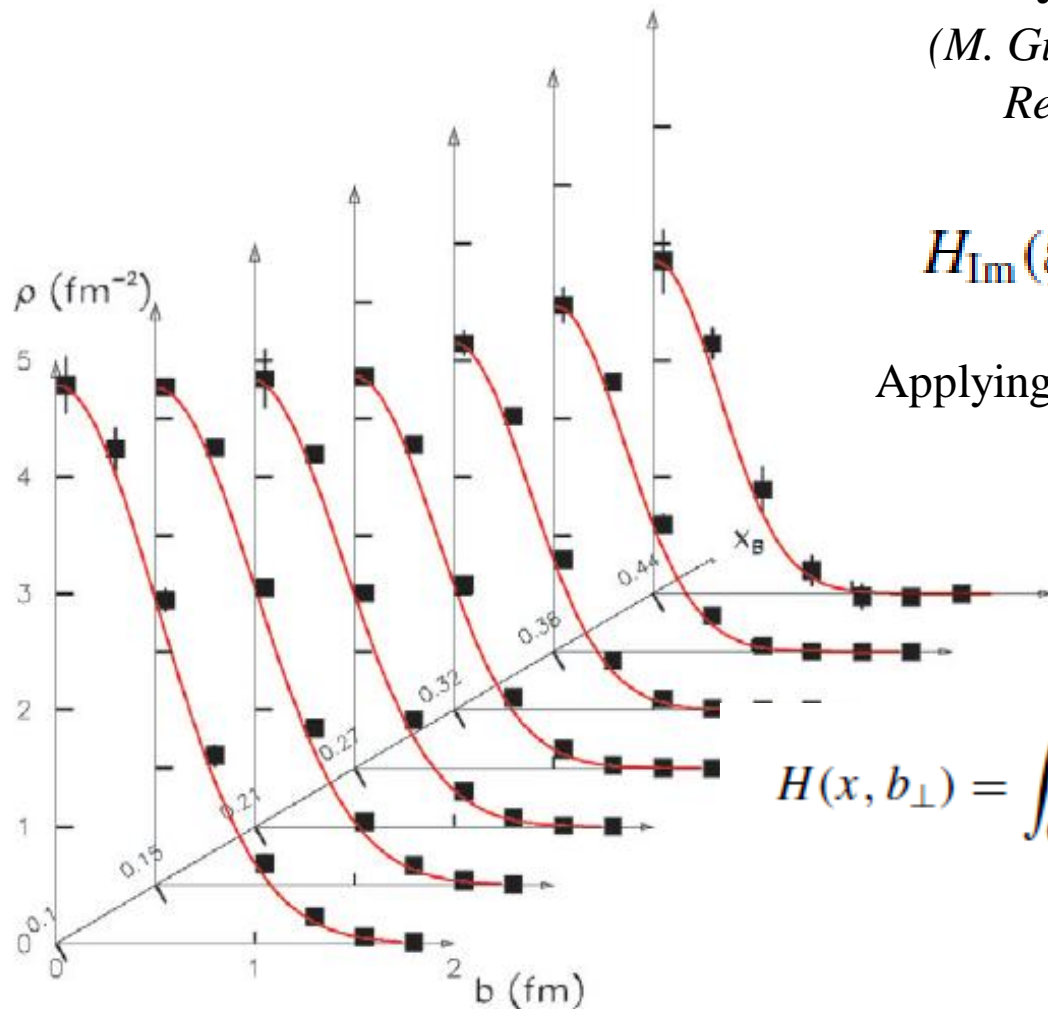
H_{Im} fit of $\sigma_{tot}, A_{LU}, A_{UL}, A_{LL}, A_{Ux}, A_{Uy}, A_{Lx}, A_{Ly}$



From CFFs to spatial densities

How to go from momentum coordinates (t)
to space-time coordinates (b) ?

(M. Guidal, H. Moutarde, M. Vanderhagen,
Rept.Prog.Phys. 76 (2013) 066202)



$$H_{1m}(\xi, t) \equiv H(\xi, \xi, t) - H(-\xi, \xi, t)$$

Applying a model-dependent “deskewing” factor:

$$\frac{H(\xi, 0, t)}{H(\xi, \xi, t)}$$

$$H(x, b_{\perp}) = \int_0^{\infty} \frac{d\Delta_{\perp}}{2\pi} \Delta_{\perp} J_0(b_{\perp} \Delta_{\perp}) H(x, 0, -\Delta_{\perp}^2)$$

Burkardt (2000)

Projections for CLAS12

E12-11-003: BSA for DVCS on the neutron with CLAS12

$$(H, E)_u(\xi, \xi, t) = \frac{9}{15} [4(H, E)_p(\xi, \xi, t) - (H, E)_n(\xi, \xi, t)]$$

$$(H, E)_d(\xi, \xi, t) = \frac{9}{15} [4(H, E)_n(\xi, \xi, t) - (H, E)_p(\xi, \xi, t)]$$

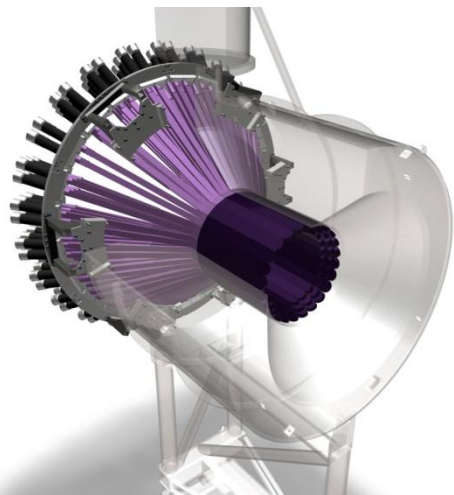
$$\Delta\sigma_{\text{LU}} \sim \sin\phi \operatorname{Im}\{F_1\mathcal{H} + \xi(F_1+F_2)\tilde{\mathcal{H}} - kF_2\mathcal{E}\} d\phi$$

The most sensitive observable to the GPD \mathcal{E}

$ed \rightarrow e(p)\gamma$

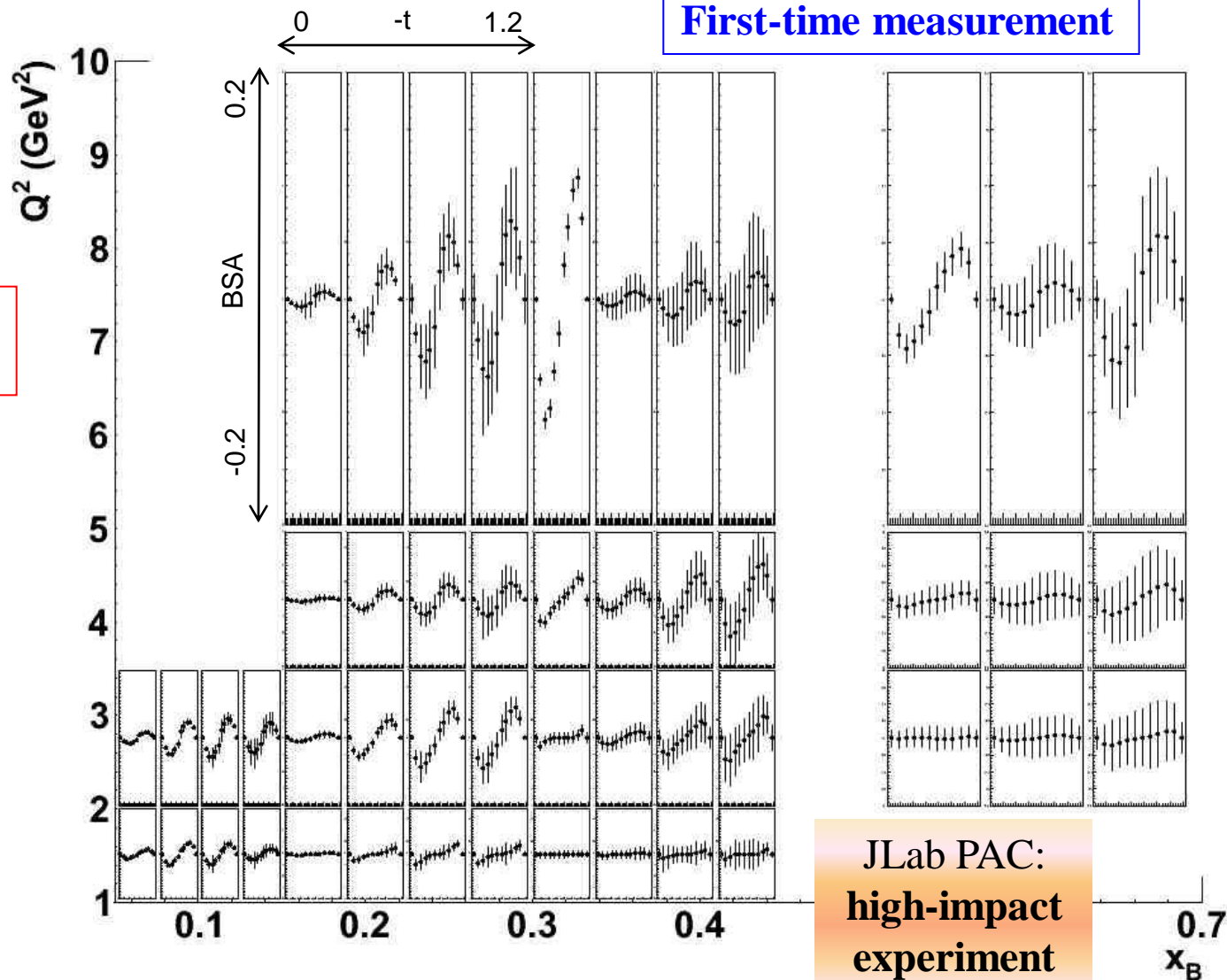
CLAS12 +
Forward Calorimeter +
Neutron Detector

80 days of data taking
 $L = 10^{35} \text{ cm}^{-2}\text{s}^{-1}/\text{nucleon}$



Construction completed
(IPN Orsay)

First-time measurement



JLab PAC:
high-impact
experiment

E12-11-003: BSA for DVCS on the neutron with CLAS12

$$(H, E)_u(\xi, \xi, t) = \frac{9}{15} [4(H, E)_p(\xi, \xi, t) - (H, E)_n(\xi, \xi, t)]$$

$$(H, E)_d(\xi, \xi, t) = \frac{9}{15} [4(H, E)_n(\xi, \xi, t) - (H, E)_p(\xi, \xi, t)]$$

$$\Delta\sigma_{\text{LU}} \sim \sin\phi \operatorname{Im}\{F_1\mathcal{H} + \xi(F_1+F_2)\tilde{\mathcal{H}} - kF_2\mathcal{E}\}d\phi$$

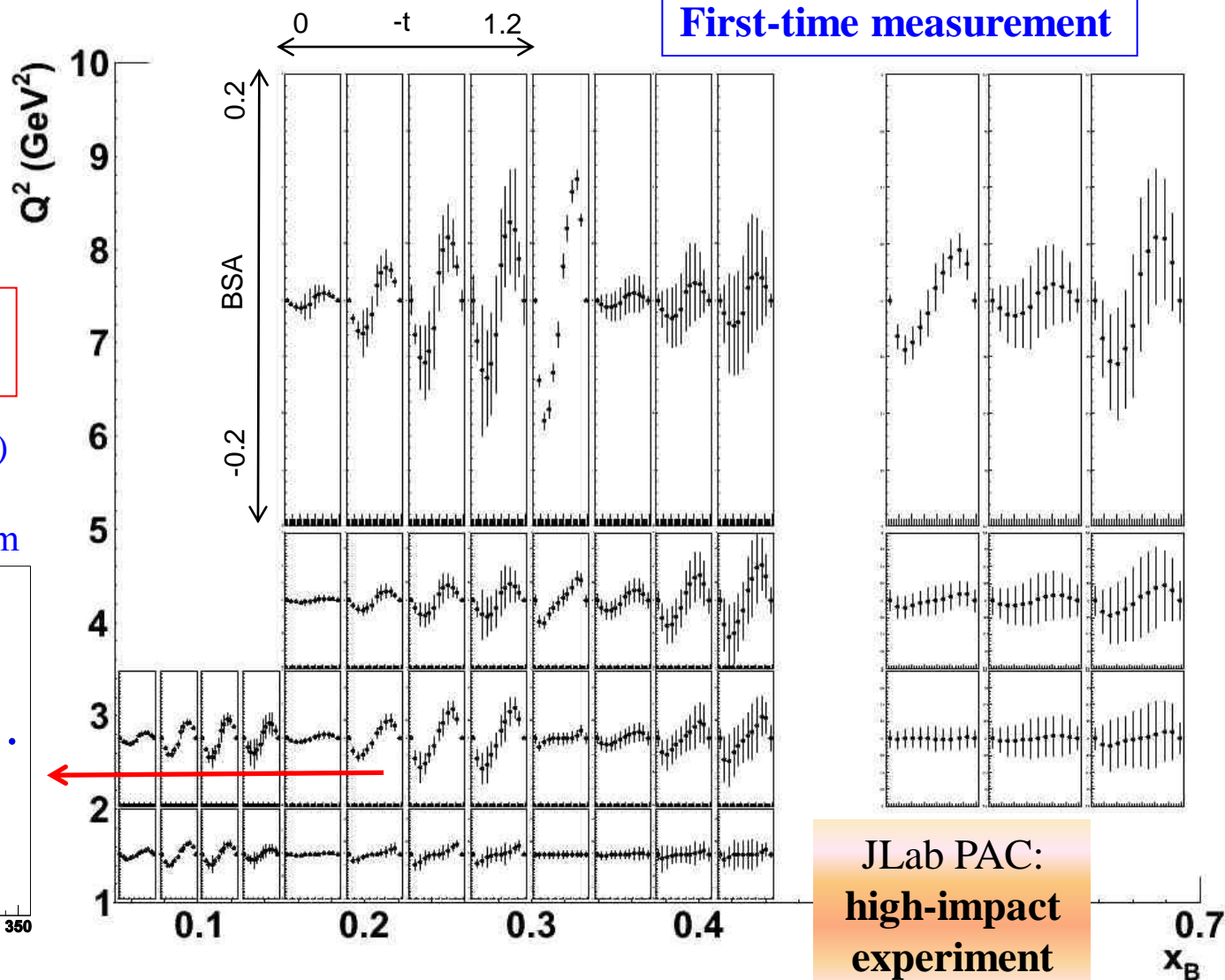
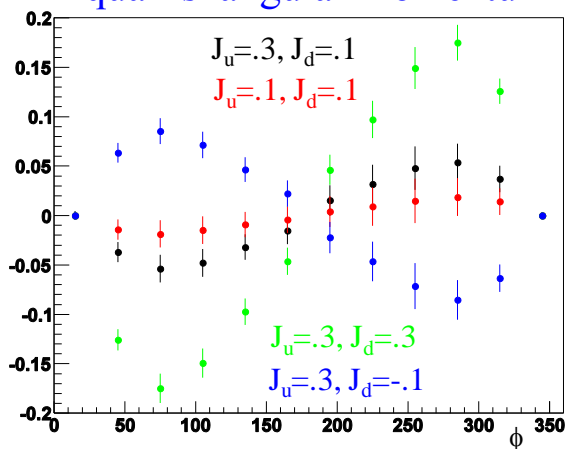
The most sensitive observable to the GPD \mathcal{E}

$ed \rightarrow e(p)\gamma$

CLAS12 +
Forward Calorimeter +
Neutron Detector

80 days of data taking
 $L = 10^{35} \text{ cm}^{-2}\text{s}^{-1}/\text{nucleon}$

Model predictions (VGG)
for different values of
quarks' angular momentum



JLab PAC:
high-impact
experiment

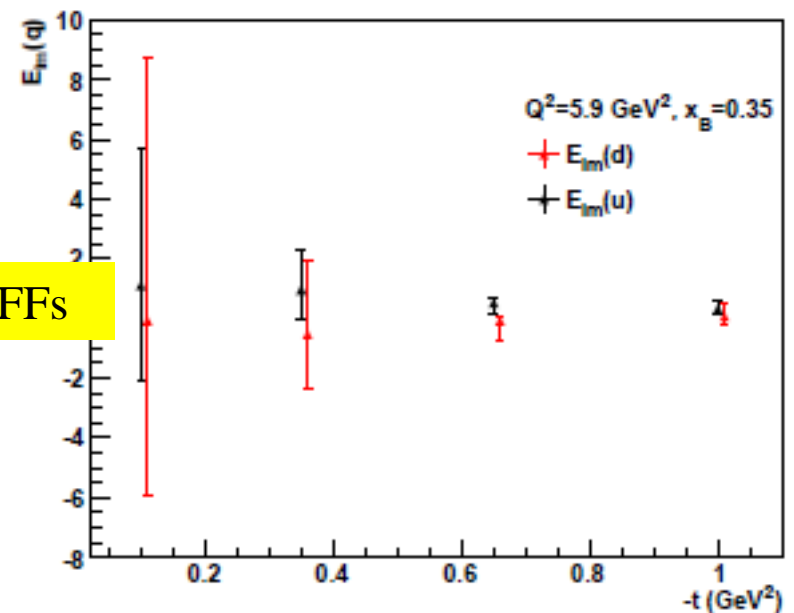
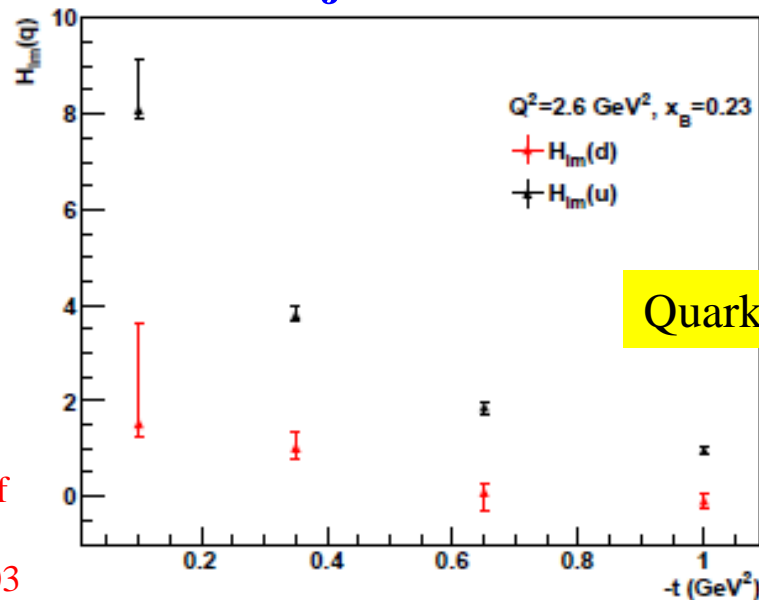
DVCS: global view

Observable (target)	Target	Sensitivity to CFFs	Completed experiments	12-GeV experiments
$\Delta\sigma_{beam}(p)$	Unpolarized hydrogen	$\Im m\mathcal{H}_p$	Hall A, CLAS	Hall A, CLAS12, Hall C
BSA(p)	Unpolarized hydrogen	$\Im m\mathcal{H}_p$	HERMES, CLAS	CLAS12
TSA(p)	Long. pol. NH3	$\Im m\mathcal{H}_p, \Im m\mathcal{H}_p$	HERMES, CLAS	CLAS12
DSA(p)	Long. pol. NH3	$\Re e\mathcal{H}_p, \Re e\mathcal{H}_p$	HERMES, CLAS	CLAS12
\bar{t} TSA(p)	Transv. pol. protons	$\Im m\mathcal{H}_p, \Im m\mathcal{E}_p$	HERMES	CLAS12
$\Delta\sigma_{beam}(n)$	Unpolarized deuterium	$\Im m\mathcal{E}_n$	Hall A	
BSA(n)	Unpolarized deuterium	$\Im m\mathcal{E}_n$		CLAS12
TSA(n)	Long. pol. ND3	$\Im m\mathcal{H}_n$		PR12-15-004
DSA(n)	Long. pol. ND3	$\Re e\mathcal{H}_n$		PR12-15-004

Projections for flavor separation ($\Im m\mathcal{H}, \Im m\mathcal{E}$)

Fit using all the projected pDVCS asymmetries of the CLAS12 program

Fit using the projected nDVCS asymmetries of PR12-15-004 and E12-11-003



Conclusions

- GPDs are a unique tool to explore the **internal dynamics of the nucleon**:
 - **3D** quark/gluon **imaging** of the nucleon
 - **orbital angular** momentum carried by quarks
- Their extraction from experimental data is **very difficult**:
 - there are **4 GPDs for each quark flavor**
 - they depend on **3 variables**, only two (ξ , t) experimentally accessible via DVCS
- ✓ Recently-developed fitting methods allow to **extract CFFs from DVCS observables**. Need to measure several **p-DVCS** and **n-DVCS observables** over a **wide phase space**
- ✓ A wealth of **new results** on various DVCS observables is coming from recent **CLAS and Hall-A experiments** (on the proton, deuterium and ^4He targets)
- ✓ First **tomographic interpretations** of the quarks in the **proton**:
 - ✓ **valence quarks** are concentrated in its **center**, **sea quarks** at its **periphery**
 - ✓ **axial charge** more concentrated than the **electric** one
- The 12-GeV-upgraded JLab will be **the only facility** to perform DVCS experiments **in the valence region**, for Q^2 up to 11 GeV
- DVCS experiments on both **proton** and **deuterium** targets (polarized and unpolarized) are planned for **3 of the 4 Halls at JLab@12 GeV**: **quarks' spatial densities, quark-flavor separation, quarks' orbital angular momentum...**
- **Beyond DVCS: double DVCS (x dependence), TCS, exclusive meson production, ...**