

Experimental landscape: vector-meson DVMP data

Exp.	years	V	Q^2 [GeV ²]	ξ	W [GeV]
ZEUS	96–00	$\rho^0, \phi, J/\psi$	2–160	10^{-5} –0.08	32–180
H1	96–00	$\rho^0, \phi, J/\psi$	2.5–60	10^{-5} –0.02	35–180
HERMES	96–07	ρ^0, ω, ϕ	0.5–7	10^{-2} –0.3	3–7
COMPASS	12/16–17	ρ^0, ω, ϕ	1–10	10^{-3} –0.2	5–17
JLab CLAS	01–08	ρ^0, ω, ϕ	1–6	0.06–0.5	2–3
JLab CLAS12	18–	$\phi, \omega, J/\psi$	1–12	0.03–0.7	thr.–4.5

Theory description: vector-meson DVMP

subprocess	accuracy	scheme	reference	comment
$\gamma_L^* \rightarrow V_L$	tw-2, LO/NLO	collinear	[Lautenschläger, Müller, K _{PK} , Schäfer '14 (unpublished)], [Čuić, Duplanić, Kumerički, K _{PK} '23], [Guo et al. '24 '25]	GPD and DA evolution important; <u>small-x pheno</u> (conf. rep.)
	tw-2, LO	MPA	[Goloskokov, Kroll '05]	meson-side $k_\perp$ + Sudakov; hard to extend consistently to NLO
$\gamma_T^* \rightarrow V_L$	partial tw-3	MPA	[Goloskokov, Kroll '13]	transversity GPDs
$\gamma_L^* \rightarrow V_T$	partial tw-3	MPA	[Goloskokov, Kroll '07]	usually small
$\gamma_T^* \rightarrow V_T$	partial tw-3	MPA	[Goloskokov, Kroll '07]	(gluon transversity GPDs for double flip)
	twist-3	collinear	[Anikin, Teryaev '02]	endpoint singularities; <u>no pheno</u>

DVMP kinematics

$$\gamma^*(q) + N(p) \rightarrow V(q_V) + N'(p')$$

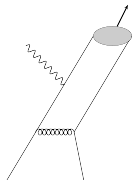
$$\begin{aligned} Q^2 &= -q^2 & t &= (p' - p)^2 = (q - q_V)^2 & W^2 &= (p + q)^2 \\ x_B &= \frac{Q^2}{2p \cdot q} = \frac{Q^2}{W^2 + Q^2 - M_N^2} & \xi_{m_V=0} &= \frac{x_B}{2 - x_B} = \frac{Q^2}{2(W^2 - M_N^2) + Q^2} \\ p_T^2 &= (1 - \xi^2)(t_{\min} - t) & t_{\min} &= -\frac{4\xi^2 M_N^2}{1 - \xi^2} \end{aligned}$$

$$\sigma_U = \sigma_T + \epsilon \sigma_L \quad R = \frac{\sigma_L}{\sigma_T}$$

We need $\sigma_L(Q^2, \xi, t)$ i.e. $\sigma_L(Q^2, W, p_T^2)$

Appendix

Subprocess amplitudes $\mathcal{H}$: projectors

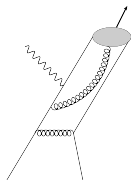


$q\bar{q} \rightarrow \pi$ projector

[Beneke, Feldmann '00]

$$(\tau q' + k_{\perp}) + (\bar{\tau} q' - k_{\perp}) = q'$$

$$\begin{aligned} \mathcal{P}_2^{\pi} \sim & f_{\pi} \left\{ \gamma_5 q' \phi_{\pi}(\tau, \mu_F) \right. \\ & + \mu_{\pi}(\mu_F) \left[\gamma_5 \phi_{\pi p}(\tau, \mu_F) \right. \\ & - \frac{i}{6} \gamma_5 \sigma_{\mu\nu} \frac{q'^{\mu} n^{\nu}}{q' \cdot n} \phi'_{\pi\sigma}(\tau, \mu_F) \\ & \left. \left. + \frac{i}{6} \gamma_5 \sigma_{\mu\nu} q'^{\mu} \phi_{\pi\sigma}(\tau, \mu_F) \frac{\partial}{\partial k_{\perp\nu}} \right] \right\}_{k_{\perp} \rightarrow 0} \end{aligned}$$



$q\bar{q}g \rightarrow \pi$ projector

[Kroll, P-K '18]

$$\tau_a q' + \tau_b q' + \tau_g q' = q'$$

$$\mathcal{P}_3^{\pi} \sim f_{3\pi}(\mu_F) \frac{i}{g} \gamma_5 \sigma_{\mu\nu} q'^{\mu} g_{\perp}^{\nu\rho} \frac{\phi_{3\pi}(\tau_a, \tau_b, \tau_g, \mu_F)}{\tau_g}, \quad f_{3\pi} \sim \mu_{\pi}$$

$$\mu_{\pi} = m_{\pi}^2 / (m_u + m_d) \cong 2 \text{ GeV}$$

DAs and EOMs

$$\tau \phi_{\pi p}(\tau) + \frac{\tau}{6} \phi'_{\pi\sigma}(\tau) - \frac{1}{3} \phi_{\pi\sigma}(\tau) = \phi_{\pi}^{EOM}(\bar{\tau})$$

$$\bar{\tau} \phi_{\pi p}(\tau) - \frac{\bar{\tau}}{6} \phi'_{\pi\sigma}(\tau) - \frac{1}{3} \phi_{\pi\sigma}(\tau) = \phi_{\pi}^{EOM}(\tau)$$

$$\phi_{\pi}^{EOM}(\tau) = 2 \frac{f_{3\pi}}{f_{\pi} \mu_{\pi}} \int_0^{\bar{\tau}} \frac{d\tau_g}{\tau_g} \phi_{3\pi}(\tau, \bar{\tau} - \tau_g, \tau_g)$$

- EOMs and symmetry properties
⇒ the subprocess amplitudes in terms of two twist-3 DAs and 2- and 3-body contributions combined
- combined EOMs → first order differential equation ⇒ from known form of $\phi_{3\pi}$ [Braun, Filyanov '90] one determines $\phi_{\pi p}$ (and $\phi_{\pi\sigma}$)

Note: $q\bar{q}g$ projector and EOMs were derived using light-cone gauge for constituent gluon