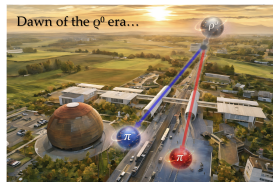


Theoretical Uncertainties and Systematic Effects in Exclusive Vector Meson Production

Kornelija Passek-K.

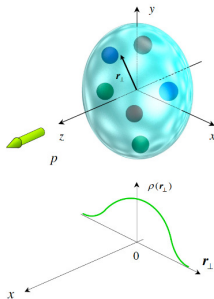
Zavod za teorijsku fiziku, IRB



Impact of Vector Mesons . . .
CERN, June 15th 2026

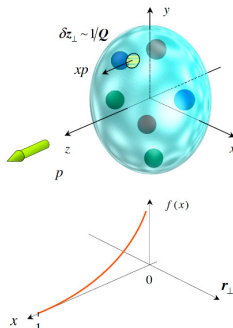
Hard exclusive processes and GPDs

Elastic scattering



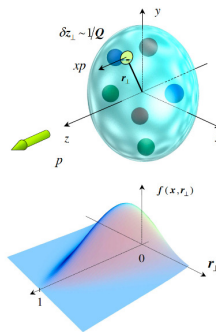
Form factors

Deep inelastic scattering



Parton distributions

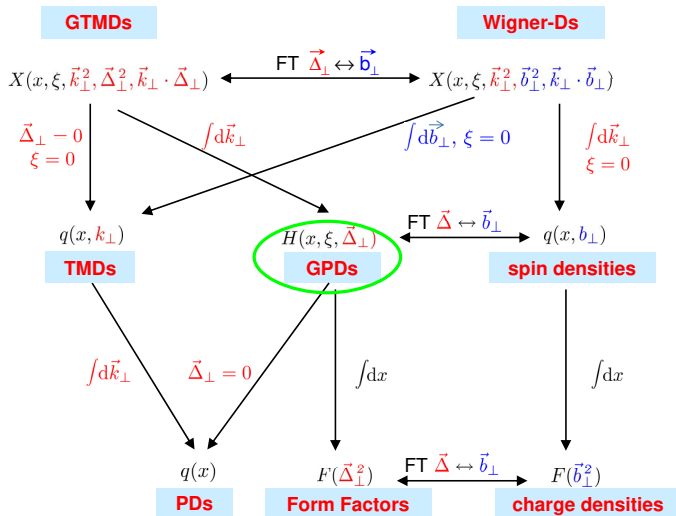
Hard exclusive processes



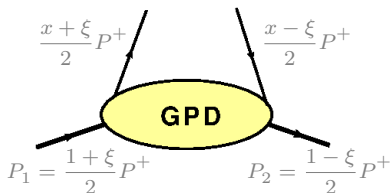
Generalized Parton Distributions (GPDs)

▷ GPDs combine momentum and spatial information

Nucleon distributions



Generalized Parton Distributions



$$P = P_1 + P_2$$

$$\Delta = P_2 - P_1$$

GPDs: $H^a(x, \xi, t; \mu) \quad a \in \{q, g\}$

x average longitudinal momentum fraction

$\xi = -\frac{\Delta^+}{P^+}$ longitudinal momentum transfer (skewness)

$t = \Delta^2$ momentum transfer

[Müller et al. '94, Radyushkin, Ji '96]

Generalized parton distributions: properties

GPDs: $F^a(x, \xi, t), \quad a \in \{q, g\}$

(twist-2) GPDs		Partons		
		helicity conserving (vector)	helicity conserving (axial)	helicity flip (transversity)
Nucleon	non-flip	H	$\tilde{H}$	$\bar{E}_T, \tilde{E}_T$
	spin flip	E	$\tilde{E}$	$H_T, \tilde{H}_T$

▷ $H, \tilde{H}, H_T(x, 0, 0) \rightarrow \text{PDFs}(x)$

▷ $\int_{-1}^1 dx F(x, \xi, t) \rightarrow \text{form factors}(t)$

▷ higher moments $\rightarrow$ angular momentum and EMT

How do we unravell GPDs?

- ▶ theory input:
 - factorization
 - perturbative calculation
 - GPD models
 - meson distribution amplitudes (DAs)
- ▶ deconvolution problem (GPDs are not measured directly)
 - amplitudes = convolutions of GPDs with hard part:

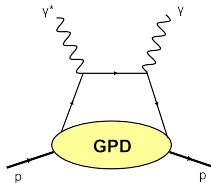
$$\mathcal{A} \sim \int dx \, C(x, \xi, Q^2) H(x, \xi, t)$$

→ observables $\sim |\mathcal{A}|^2$

- ▶ GPDs are universal (same GPDs → different processes)

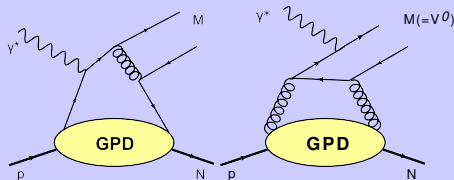
Hard exclusive processes: focus on meson production

Compton scattering
→ deeply virtual (DVCS)



$$\gamma^* p \rightarrow \gamma p$$

meson production
→ deeply virtual (DVMP)



$$\gamma^* p \rightarrow MN$$

→ meson **distribution amplitude (DA)**: $\phi_M(z)$
 z ... longitudinal momentum fraction

Factorization for hard exclusive processes

$$\text{Amplitude} = \text{Hard-scattering amplitude (parton interactions)} \otimes \text{Hadron wave functions (GPD, DA)}$$

$\Rightarrow$ observables

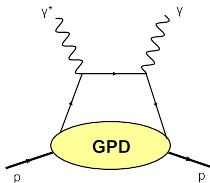
- hard scale (e.g. Q^2 , s , $M_{J/\Psi}^2$, ...) $\rightarrow$ perturbative QCD
 - expansion in $\alpha_s(Q)$
 - power corrections $\sim 1/Q$ (higher twist)

Generic questions

- factorization proven/assumed? broken?
- pQCD calculation:
 - accessible energies high enough ?
 - finite order predictions $\Rightarrow$ theoretical uncertainties ?, dependence on factorization and renormalization scales ?
 - higher-orders in α_S ?
 - higher-twist effects ?
- nonperturbative input:
 - form/modeling of GPDs and DAs

Hard exclusive processes: focus on meson production

Compton scattering
→ deeply virtual (DVCS)



$$\gamma_T^* N \rightarrow \gamma N$$

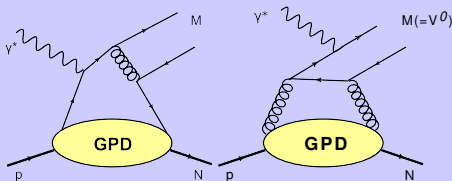
[Collins, Freund '99]

$$H^q, E^q, \tilde{H}^q, \tilde{E}^q$$

NLO:

$$H^G, E^G, \tilde{H}^G, \tilde{E}^G, F_T^G$$

meson production
→ deeply virtual (DVMP)



$$\gamma_L^* N \rightarrow M N' \quad \text{twist-2}$$

[Collins, Frankfurt, Strikman '97]

$$M = V_L: H^{q_i}, E^{q_i}; H^G, E^G$$

$$M = PS: \tilde{H}^{q_i}, \tilde{E}^{q_i}$$

$$\gamma_T^* N \rightarrow M N' \quad \text{twist-3}$$

$$M_{\text{twist-3}} \Rightarrow F_T^q$$

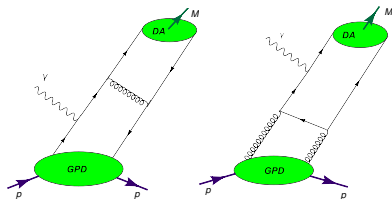
Meson production status

- experimental data from HERA, COMPASS and JLab, expected from EIC ...
- $DV\rho^0P$, $DV\phi P$, $DV\omega P$:
 - data show importance of γ_L^* contributions ($Q^2 < 100 \text{ GeV}^2$)
→ twist-2 predictions to describe σ_L
- in contrast to $DV\pi P$:
 - data show suppression of γ_L^* contributions ($Q^2 < 10 \text{ GeV}^2$)
→ twist-3 predictions describing γ_T^* contributions needed

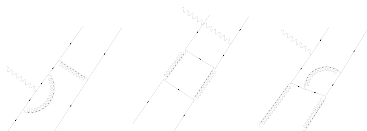
$$\mathcal{A} \sim \text{hard amplitude}(\mu_R, \mu_F) \otimes GPD(\mu_F) \otimes DA(\mu_F)$$

$$\mu_R, \mu_F \rightarrow f(Q)$$

$$\text{LO} \sim \alpha_S(\mu_R)$$



$$\text{NLO} \sim \alpha_S^2(\mu_R)$$



[Belitsky, Müller '01, Ivanov et al '04,
Duplanić, Müller, KPK '17]

- large NLO corrections [Diehl, Kugler '07, Müller, Lautenschläger, KPK, Schäfer '14] ,
(and large $\ln(1/\xi)$ terms for $\xi \ll 1$)
- strong renormalization scale dependence at LO but improved stability at NLO (due to $\alpha_S^2(\mu_R) \log \mu_R$ terms)
- GPD, DA (NLO) evolution to factorization scale μ_F
- few DVMP analysis to NLO

[Lautenschläger, Müller, KPK, Schäfer '14

(unpublished), Čuić, Duplanić, Kumerički, KPK

'23, Guo et al. '24 '25]

From x-space to j-space: CPaW formalism

$$\mathcal{A} \sim \text{hard amplitude} \otimes GPD \otimes DA$$

$$\otimes \equiv \int_{-1}^1 \frac{dx}{2\xi} \quad \text{x-space (momentum fraction)}$$

$$\rightarrow 2 \sum_{j=0}^{\infty} \xi^{-j-1} \quad \text{j-space (conformal moments)}$$

(as Mellin moments in DIS

$$x^n \rightarrow C_n^{3/2}(x), C_n^{5/2}(x))$$

$$\rightarrow \otimes \equiv \frac{1}{2i} \int_{c-i\infty}^{c+i\infty} dj \, \xi^{-j-1} \left[i \pm \left\{ \frac{\tan}{\cot} \right\} \left(\frac{\pi j}{2} \right) \right]$$

(Mellin-Barnes integral)

► widely used, intuitive

→ see P. Sznajder talk

(PARTONS software)

► simpler evolution

► efficient phenomenology

CPaW(conformal partial wave): [Müller '06, Müller, Schäfer '06],

[Kumerički, Müller, KPK, Schäfer '07, Müller, Lautenschläger, KPK, Schäfer '14]

(Gepard software) [Kumerički '22]

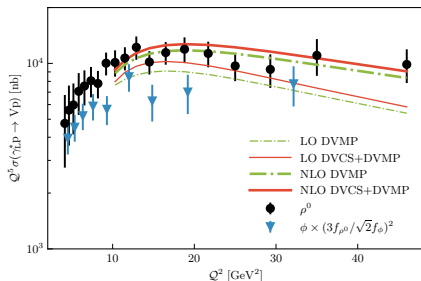
→ existing NLO DVMP phenomenology: conformal representation

- fits to small- x HERA collider data
- CPaW formalism
- small- x kinematics $\Rightarrow$ only dominant GPD H^a , $a \in \{\text{sea}, G\}$,
updated [Kumerički, Müller '10] model, 11 parameters
- LT separation: $R \equiv \frac{\sigma_L^{\rho^0}}{\sigma_T^{\rho^0}} \rightarrow R(W, Q^2)$ fit

NLO analysis: global NLO fits DIS, DVCS, $DV\rho_L^0P$

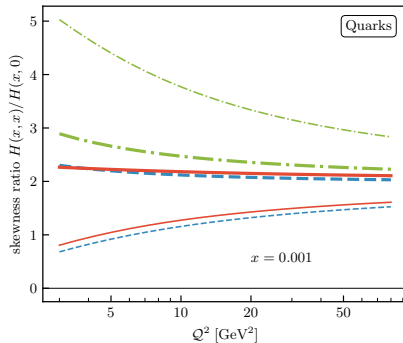
[Čuić, Duplančić, Kumerički, KPK '23]

- ▶ successful description of Q^2 dependence at NLO



- ▶ ! only NLO fit successful

goodness of GPD extraction:



- σ_L asymptotically $\sim 1/Q^6$
- exp: $\sim 1/Q^4$ for fixed x_B ,
 $\sim 1/Q^5$ for fixed W

--- LO DVCS
 --- LO DVMP
 --- LO DVCS+DVMP
 --- NLO DVCS
 --- NLO DVMP
 --- NLO DVCS+DVMP

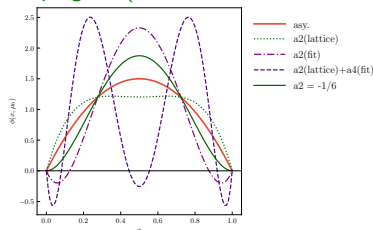
NLO analysis: global NLO fits DIS, DVCS, $DV\rho_L^0P$

- global DIS+DVCS+DVMP fits show importance of NLO
 $\Rightarrow$ universal GPDs

- DVV_LP can only be described at NLO

- ! meson DA additional nontrivial nonperturbative input
 $\Rightarrow$ end-point suppressed ρ_L^0 DA favoured?

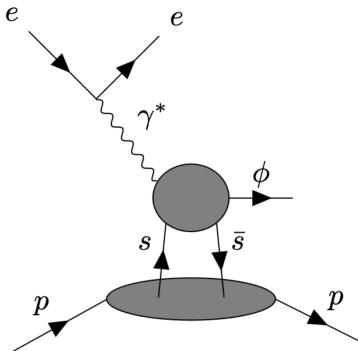
$\rightarrow$ work in progress (Kishore, Kumerički, KPK)



Near threshold kinematics: DV ϕ production proposed

$$\gamma^* p \rightarrow \phi p$$

- near threshold $\Rightarrow$ large ξ ($\xi \rightarrow 1$)
 $\Rightarrow$ dominance of first conformal moment $\rightarrow \sim D(t)$ $\rightarrow$ pressure in p
- $\phi \simeq s\bar{s} \Rightarrow$ **sensitivity to $D_s(t)$**
- complementary to DVCS (u, d) and J/ψ production (gluons) [Guo, Ji, Yuan '24]
- proposed measurement at JLab



[Hatta, Klest, KPK, Schoenleber '25]

$$D(0) = \underbrace{D_g(0)}_{J/\psi} + \underbrace{D_u(0) + D_d(0)}_{\text{DVCS}} + \underbrace{D_s(0)}_{\phi \text{ DVMP}} + \dots$$

Beyond leading-twist collinear DVMP

→ x-space representation

- **Modified perturbative approach (MPA)** [Goloskokov, Kroll '05]
 - meson-side $k_{\perp}$ + Sudakov suppression; nucleon side unchanged
 - realistic Q^2 dependence already at LO
 - broad phenomenology and GPD-model development
 - difficult to extend consistently to NLO
- **Higher twist /transverse photons**
 - [Anikin, Teryaev '02] $\gamma_T^* N \rightarrow V_T N'$: endpoint singularities; collinear factorization breaks; no pheno
 - [Goloskokov, Kroll '13] $\gamma_T^* N \rightarrow V_L N'$: partial twist-3 + MPA; phenomenology for ρ , ω and polarization observables

Concluding remarks

Theoretical uncertainties

- **Factorization:** controlled channel $\gamma_L^* N \rightarrow V_L N'$ (twist-2)
- **Perturbative accuracy:** NLO and GPD/DA evolution needed $\rightarrow$ scale dependence, possibly resummation needed
- **Higher-twist:** transverse photons $\gamma_T^* N \rightarrow V N'$ $\rightarrow$ endpoint singularities, factorization breaking
- **Nonperturbative input:** GPD ansatz, vector-meson DA

Systematics effects

- **Measured vs. controlled object:** experiments measure $\sigma_T + \epsilon\sigma_L$; theory controls σ_L
- **Separation assumptions:** $R = \sigma_L/\sigma_T$, SDMEs, SCHC
- **Input to GPD fits:** the extracted σ_L carries the assumptions used in the separation

Concluding remarks

Theoretical uncertainties

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Systematics effects

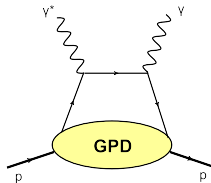
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- **Input to GPD fits:** the extracted σ_L carries the assumptions used in the separation

Thank you.

Appendix

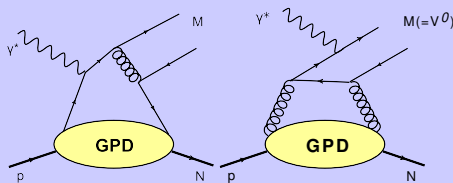
Hard exclusive processes: focus on meson production

Compton scattering
→ deeply virtual
(DVCS)



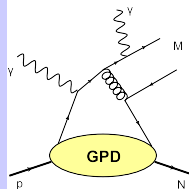
$$\gamma^* p \rightarrow \gamma p$$

meson production
→ deeply virtual (DVMP)



$$\gamma^* p \rightarrow MN$$

(γM)
photoproduction



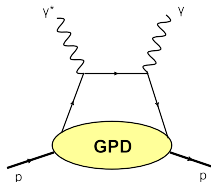
$$\gamma p \rightarrow (\gamma M) N$$

→ meson **distribution amplitude (DA)**: $\phi_M(z)$

$z \dots$ longitudinal momentum fraction

Hard exclusive processes: focus on meson production

Compton scattering
→ deeply virtual
(DVCS)



$$\gamma_T^* N \rightarrow \gamma N$$

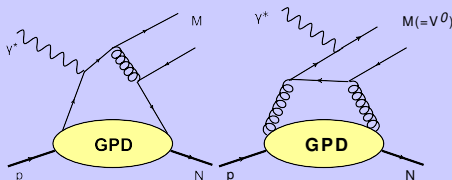
[Collins, Freund '99]

$$H^q, E^q, \tilde{H}^q, \tilde{E}^q$$

NLO:

$$H^G, E^G, \tilde{H}^G, \tilde{E}^G, F_T^G$$

meson production
→ deeply virtual (DVMP)



$$\gamma_L^* N \rightarrow M N' \quad \text{twist-2}$$

[Collins, Frankfurt, Strikman '97]

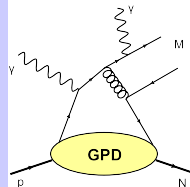
$$M = V_L: H^{q_i}, E^{q_i}; H^G, E^G$$

$$M = PS: \tilde{H}^{q_i}, \tilde{E}^{q_i}$$

$$\gamma_T^* N \rightarrow M N' \quad \text{twist-3}$$

$$M_{\text{twist-3}} \Rightarrow F_T^q$$

(γM)
photopro-
duction



$$\gamma N \rightarrow (\gamma M) N'$$

[Qiu, Yu '22]

(quarks)

$$F^a, F_T^a$$

Meson production status

- experimental data from HERA, JLab and COMPASS (CERN), expected from EIC ...
- $DV\rho^0P$:
 - data show importance of γ_L^* contributions ($Q^2 < 100 \text{ GeV}^2$)
→ twist-2 predictions describe σ_L
 - global DIS+DVCS+DVV_LP fits at NLO [Čuić, Duplančić, Kumerički, P-K. '23]
- $DV\pi P$:
 - data show suppression of γ_L^* contributions ($Q^2 < 10 \text{ GeV}^2$)
→ twist-3 predictions describing γ_T^* contributions needed
 - twist-3 contributions analyzed [Duplančić, Kroll, P-K., Szymanowski '24]

Global NLO fits (DIS+DVCS+DVV_LP)

small- x global fits to HERA collider data (ρ_0)

- only NLO predecessor: [Lautenschlager, Müller, Schäfer '13 unpublished]
- hard scattering amplitude corrected [Duplanić, Müller, P-K. '17]
- new NLO fit [Čuić, Duplanić, Kumerički, P-K. '23]:
improved treatment of experimental data

GPD model: [Kumerički, Müller, P-K., Schäfer '07, Kumerički, Müller '10]

$$\bullet H_j^a(\xi, t) = q_j^a \frac{1 + j - \alpha_0^a}{1 + j - \alpha_0^a - \alpha'^a t} \left(1 - \frac{t}{m_a^2}\right)^{-2} (1 + s_2^a \xi^2 + s_4^a \xi^4)$$

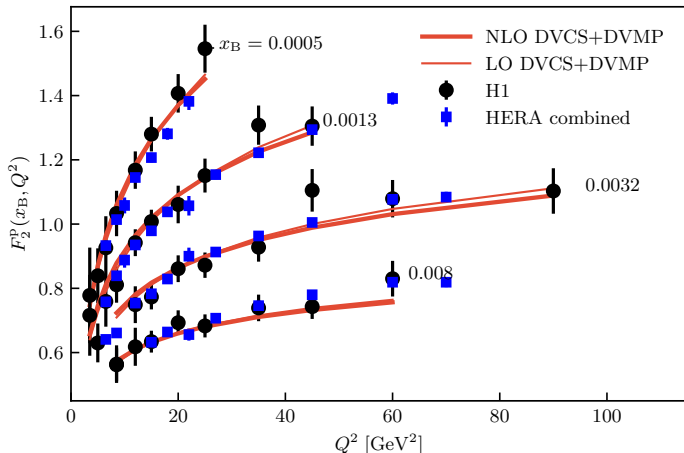
$$q_j^a = N_a \frac{B(1 - \alpha_0^a + j, \beta^a + 1)}{B(2 - \alpha_0^a, \beta^a + 1)}$$

- small- x kinematics $\Rightarrow a \in \{\text{sea}, G\}$, only dominant H GPD

Fit parameters:

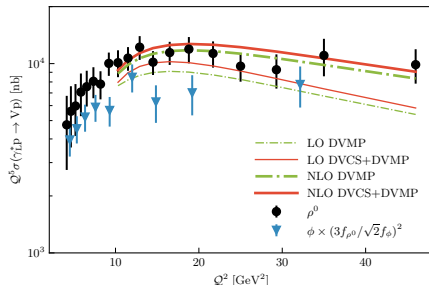
- DIS: $\{N_{\text{sea}}, \alpha_0^{\text{sea}}, \alpha_0^G\}$
- DVCS+DVMP: $\{\alpha'^{\text{sea}}, \alpha'^G, m_{\text{sea}}^2, m_G^2, s_2^{\text{sea}}, s_2^G, s_4^{\text{sea}}, s_4^G\}$

Global NLO fits (DIS+DVCS+DVV_LP)



● may seem trivial, but not all popular models describe DIS

Global NLO fits (DIS+DVCS+DVV_LP)



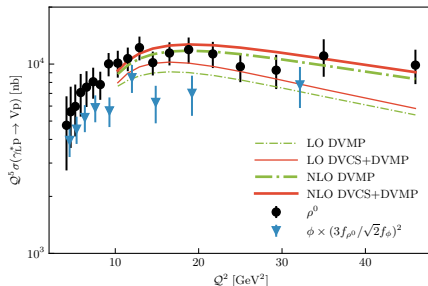
σ_L asymptotically: $\frac{1}{Q^6}$

experimental data for fixed x_B : $\approx \frac{1}{Q^4}$, for fixed W : $\approx \frac{1}{Q^5}$

- successful description of Q^2 dependence

NLO analysis: global NLO fits DIS, DVCS, $DV_{\rho_L^0}P$

[Čuić, Duplanić, Kumerički, KPK '23]



σ_L asymptotically $\frac{1}{Q^6}$ while experimental data $\approx \frac{1}{Q^5}$

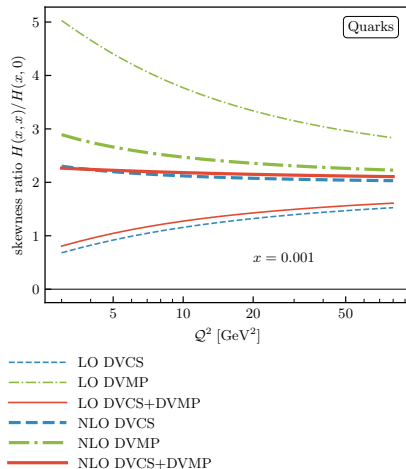
→ ! successful description of Q^2 dependence at NLO

NLO analysis: global NLO fits DIS, DVCS, $DV\rho_L^0P$

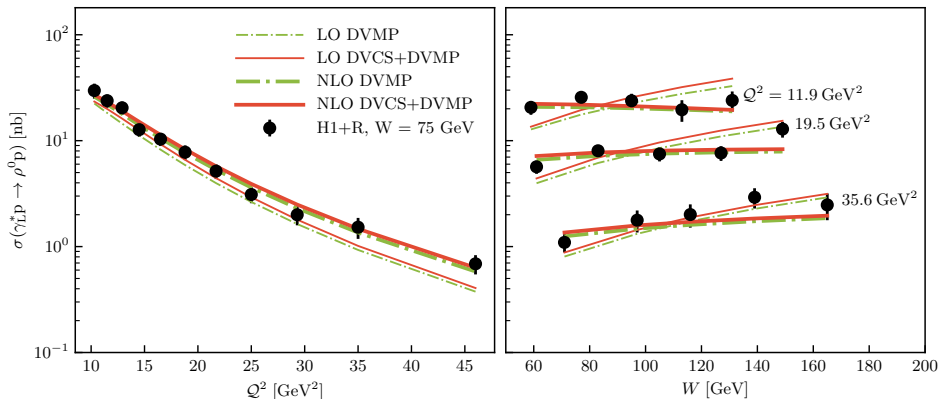
[Čuić, Duplančić, Kumerički, KPK '23]

- strong constraints on GPDs
- ! only NLO fit successful

measures goodness of GPD extraction:



Global NLO fits (DIS+DVCS+DVV_LP)



$$R \equiv \frac{\sigma_L^{\rho^0}}{\sigma_T^{\rho^0}} \rightarrow R(W, Q^2) = \frac{Q^2}{m_{\rho^0}^2} \left(1 + a \frac{Q^2}{m_{\rho^0}^2} \right)^{-p} \left(1 + b \frac{Q^2}{W} \right) \text{ fit}$$

Global NLO fits (DIS+DVCS+DVV_LP)

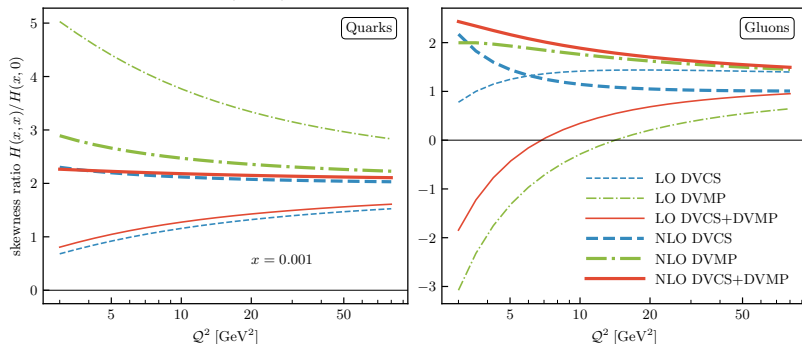
Dataset	Refs.	n_{pts}	LO-			NLO-		
			DVCS	DVMP	DVCS-DVMP	DVCS	DVMP	DVCS-DVMP
DIS	[90]	85	0.6	0.6	0.6	0.8	0.8	0.8
DVCS	[92–95]	27	0.4	$\gg 1$	0.6	0.6	$\gg 1$	0.8
DVMP	[88, 89]	45	$\gg 1$	3.1	3.3	$\gg 1$	1.5	1.8
Total		157	$\gg 1$	$\gg 1$	1.4	3.7	$\gg 1$	1.1

Table 3. Values of χ^2/n_{pts} for each LO or NLO model (columns) for the total DIS + DVCS + DVMP dataset and for subsets corresponding to different processes (rows). (The values denoted by $\gg 1$ are greater than 10.).

- NLO DVCS-DVMP fit describes the data well

Global NLO fits (DIS+DVCS+DVV_LP)

$$\text{Skewness ratio } r = \frac{H(x, x)}{H(x, 0)}$$



- r measures goodness of GPD extraction $\Rightarrow$ NLO fit successful

NLO for DV V_L production

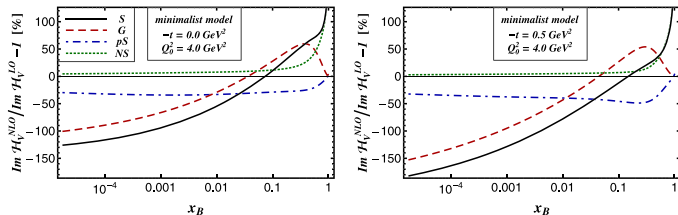


Fig. 6. Relative NLO corrections to the imaginary part of the flavor singlet TFF $\mathcal{F}_V^S$ (solid) broken down to the gluon (dashed), pure singlet quark (dash-dotted) and 'non-singlet' quark (dotted) at $t = 0 \text{ GeV}^2$ (left panel) and $t = -0.5 \text{ GeV}^2$ (right panel) at the initial scale $Q_0^2 = 4 \text{ GeV}^2$.

[Müller, Lautenschlager, P-K., Schäfer '14]

- large NLO corrections for small x_B , i.e., ξ

↑

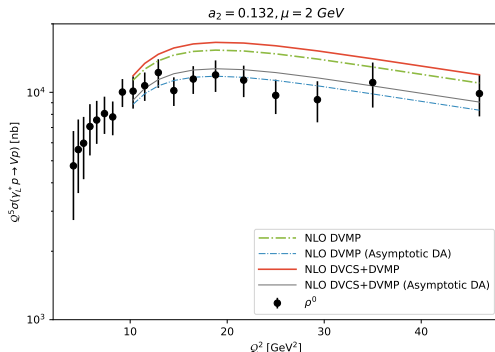
big $\ln(1/\xi)$ terms for $\xi \ll 1$, i.e., $j = 0$ pole, in gluon evolution and gluon coefficient function

Improving DA description

→ work in progress with Raj Kishore, K. Kumerički

$$\phi^p(u, \mu_F) = 6u(1-u) \left[1 + a_2(\mu_F) C_2^{3/2}(2u-1) + \dots \right]$$

$$a_2(\mu_0) = 0.132, \mu_0 = 2 \text{ GeV} \text{ [Braun et al. '16]}$$



⇒ significant impact of the DA form

Improving DA description

→ work in progress with Raj Kishore, K. Kumerički

- new GPD fit:

Dataset	Refs.	n_{pts}	LO	NLO			
			asy.	asy.	a_2 (lat.)	a_2 (fit)	a_2 (lat.) + a_4 (fit)
DIS	[5]	85	0.6	0.8	0.8	0.8	0.8
DVCS	[6, 7, 8, 9]	27	0.6	0.8	0.9	0.9	0.9
DVMP	[1, 10]	45	3.3	1.8	2.1	0.9	0.7
Total		157	1.4	1.1	1.2	0.9	0.8

- $a_2(\text{fit}) = -0.369$

- $a_2[\text{lattice: Braun et al. '16}] = 0.132, a_4(\text{fit}) = -0.519$

⇒ dominant influence of LO $\sim \sum a_n$

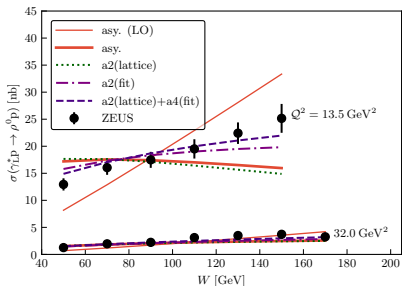
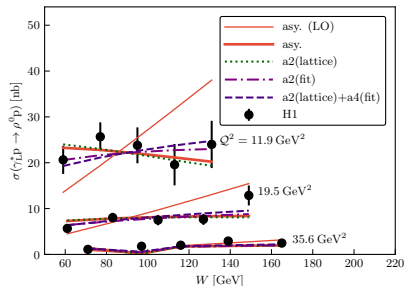
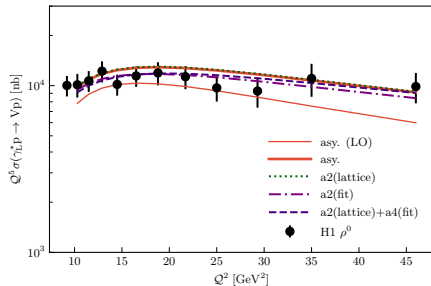
⇒ indication of end-point suppressed ρ_L DA [Liu, Shuryak, Zahed '24]

⇒ GPDs not changed much

(similar skewness ratio, W dependence improved)

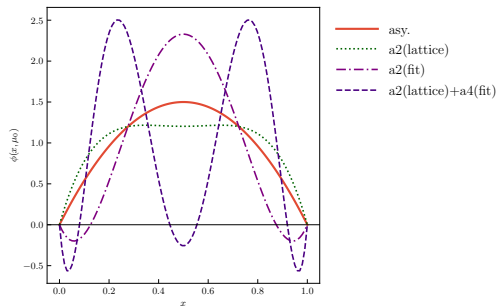
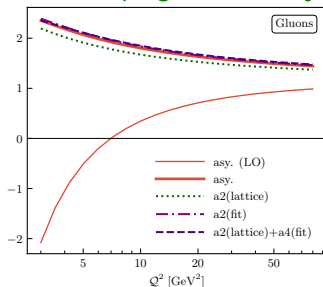
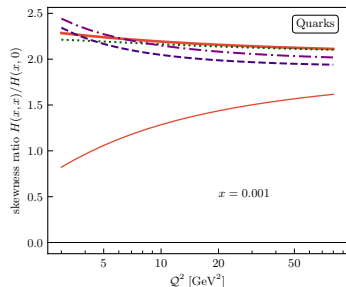
Improving DA description

→ work in progress with Raj Kishore, K. Kumerički



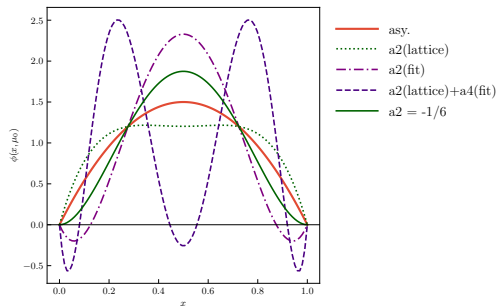
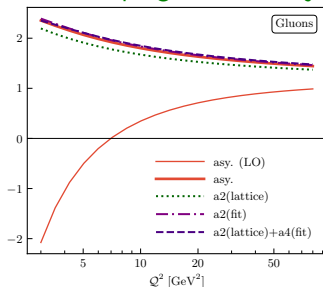
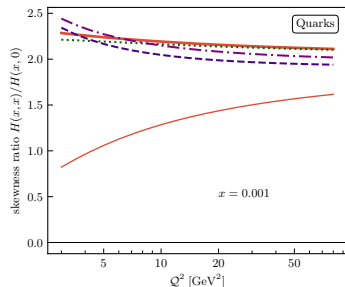
Improving DA description

→ work in progress with Raj Kishore, K. Kumerički



Improving DA description

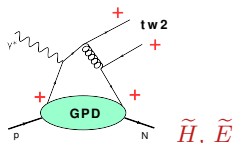
→ work in progress with Raj Kishore, K. Kumerički



When the simple picture fails: twist-3 effects in DVMP

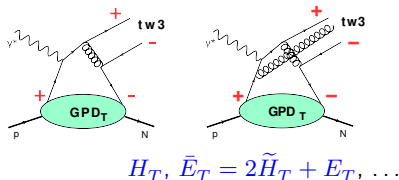
$$\gamma^* p \rightarrow \pi p$$

twist-2 ... non-flip subprocess amplitudes



► twist-2 contribution insufficient

twist-3 ... flip subprocess amplitudes



► twist-3 contributions $\sim \mu_\pi/Q$
($\mu_\pi = 2 \text{ GeV}$) important

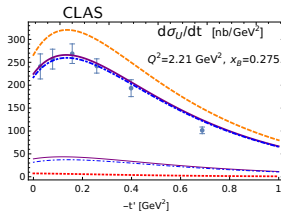
► factorization challenging

DV π P at twist-3

! σ_T (twist-3) dominant

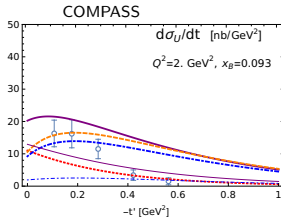
! σ_L (twist-2) negligible

(except for COMPASS
kinematics at $x_B < 0.1$)



red curves: tw2

blue curves: tw3



exp data:

full circles [CLAS '14]

open circles [COMPASS '19]

[Duplanić, Kroll, KPK, Szymanowski '24]

Treatment of end-point singularities: MPA

⇒ Modified perturbative approach (MPA) [Goloskov, Kroll, '10]

- $k_{\perp}$ quark transverse momenta in pion

$$\frac{1}{((x + \xi)\tau - k_T^2/Q^2(2\xi) + i\epsilon)} \frac{1}{(x + \xi + i\epsilon)}$$

- $\phi_{\pi} \rightarrow$ light-cone wave function $\Psi_{\pi} \sim \phi_{\pi} \exp[-a_{\pi}^2 k_{\perp}^2]$

- $\int_0^1 d\tau \rightarrow \int d^2\mathbf{k}_T \int_0^1 d\tau \xrightarrow{\text{FT}} \int d^2\mathbf{b} \int_0^1 d\tau$

- Sudakov form factor $\exp[-S(\tau, \mathbf{b}, Q^2)]$

- ▶ consistently treated 2- and 3-body tw3 contributions, as well as tw2
- ▶ involved multidimensional integrations
- ▶ calculation of NLO corrections would be complicated

Appendix: EMT matrix element

$$\begin{aligned}\langle p', s' | T_a^{\mu\nu} | p, s \rangle &= \bar{u}(p', s') \left[A_a(t) \gamma^{(\mu} P^{\nu)} + B_a(t) \frac{P^{(\mu} i \sigma^{\nu)\rho} \Delta_\rho}{2M} \right. \\ &\quad \left. + D_a(t) \frac{\Delta^\mu \Delta^\nu - g^{\mu\nu} \Delta^2}{4M} + \bar{C}_a(t) M g^{\mu\nu} \right] u(p, s)\end{aligned}$$

$$P = \frac{p + p'}{2}, \quad \Delta = p' - p, \quad t = \Delta^2$$

- $A_a \rightarrow$ momentum / energy
- $A_a + B_a \rightarrow$ angular momentum
- $D_a \rightarrow$ mechanical properties

Appendix: GPD moments and GFFs

$$\int_{-1}^1 dx \, x \, H^a(x, \xi, t) = A_a(t) + \xi^2 D_a(t)$$

$$\int_{-1}^1 dx \, x \, E^a(x, \xi, t) = B_a(t) - \xi^2 D_a(t)$$

$$\int_{-1}^1 dx \, x \, [H^a + E^a] = A_a(t) + B_a(t)$$

$$J_a(t) = \frac{1}{2} [A_a(t) + B_a(t)]$$

$$D(t) = \sum_a D_a(t)$$

DVMP differential cross-sections

$$\frac{d^4\sigma}{dW^2 dQ^2 dt d\varphi} = \frac{\alpha_{em}(W^2 - m_N^2)}{16\pi^2 E_L^2 m_N^2 Q^2 (1 - \varepsilon)} \left(\frac{d\sigma_T}{dt} + \varepsilon \frac{d\sigma_L}{dt} + \varepsilon \cos(2\varphi) \frac{d\sigma_{TT}}{dt} + \sqrt{2\varepsilon(1 + \varepsilon)} \cos\varphi \frac{d\sigma_{LT}}{dt} \right)$$

$$\frac{d\sigma_U}{dt} = \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt}$$

$$\frac{d\sigma_L}{dt} : \tilde{H}, \tilde{E} \quad \frac{d\sigma_T}{dt} : H_T, \bar{E}_T \quad \frac{d\sigma_{TT}}{dt} : \bar{E}_T \quad \frac{d\sigma_{LT}}{dt} : \tilde{E}, H_T$$

$$\left| \frac{d\sigma_{TT}}{dt} \right| \leq \frac{d\sigma_T}{dt}$$