The density-of-states approach for dense matter Monte-Carlo simulations

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with:

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The density-of-states approach for dense matter Monte-Carlo simulations - p. 1/24

Consider a theory with complex action:

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- Let us *drop* it:

$$Z_{\rm mod}(\mu) = \int \mathcal{D}\phi \, \exp\{S_R[\phi,\mu]\}$$

- How big is the error?
 - \Rightarrow apparently small for small chemical potential μ

• Quantifying the problem:

We have: $Z(\mu), Z_{mod}(\mu) > 0$

$$O(\mu) := \frac{Z(\mu)}{Z_{\text{mod}}(\mu)} , \quad Z(\mu) = O(\mu) Z_{\text{mod}}(\mu)$$

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• Density:

$$\rho(\mu) = \Delta \rho(\mu) + \rho_{\rm mod}(\mu) ,$$

 $\rho = \frac{1}{V} \frac{d}{d\mu} \ln Z(\mu):$ full density $\Delta \rho = \frac{1}{V} \frac{d}{d\mu} \ln O(\mu):$ overlap contribution $\rho_{\text{mod}} = \frac{1}{V} \frac{d}{d\mu} \ln Z_{\text{mod}}(\mu):$ easy to calculate

• How important is the overlap $O(\mu)$? [physics grounds] $|\det_{\mu}[U]|^2 = \det_{\mu}[U] \det_{-\mu}[U]$

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Impact -

Pions condense for $\mu \approx 70 \text{ MeV}$ $O(\mu)$ needs to remove pion condensation and provide a density of baryons for $\mu \approx 300 \text{ MeV}$



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, but...

F: free energy, Δf : difference in free energy density

$$O(\mu) = \exp\{-[F(\mu) - F_{mod}(\mu)]\} = \exp\{-\Delta f V\} \ll 1$$

V: volume

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overlap problem !

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Strong coupling expansion

Nuclear Physics [deForcrand, Fromm, 2010] [Philipsen et al, 2011]

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The density-of-states approach (LLR-algorithm)

[Langfeld, Lucini, Rago, PRL 109 (2012) 111601]

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 ⇒ positive measure (even for complex action systems)

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- Pros: △ does not rely on a reformulation
 [aiming at a Universal solution]

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- Cons: \triangle does it work? [need an *exponential* error suppression]

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- Cons: \triangle does it work? [need an *exponential* error suppression]
- LLR-algorithm has all prerequisites...

[Langfeld, Lucini, Rago, PRL 109 (2012) 111601]

⇒ generically solves overlap problems
but is it good enough? [potentially theory dependent]

The density of states - SU(3) versus SU(2)

• Definition:

$$\rho(E) = \int \mathcal{D}\phi \, \delta\Big(E - S[\phi]\Big)$$

The density of states - SU(3) versus SU(2)



The density of states - compact U(1)

study phase transition in U(1):

```
△ weakly first order
[ Arnold, Bunk, Lippert,
Schilling,
Nucl.Proc.Proc.Suppl.
119 (2003) 119]
```

 $\triangle P_{\beta}(E) = \rho(E) \exp\{\beta E\}$

[Langfeld,Lucini,Pellegrini,Rago, in preparation]

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[benchmark testing!]

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Which theory should we choose?

- \Rightarrow needs to be "solvable" by standard techniques
 - [benchmark testing!]

...choose a theory that admits dualisation

worm algorithm!

[Prokof'ev, Svistuno, 2001] [Chandrasekharan 2010] [Gattringer, Evertz, 2011]

The quantum O(2) model:

Results - phase diagram



quantum : scaling limit
of the 2nd order transition
m : mass gap
mu : chemical potential
T : temperature

superfluidity!

in the cold, but dense regime

[KL, PRD D87 (2013) 114504]

Theory - degrees of freedom: $P(x) \in Z_3$

$$S[P] = \tau \sum_{x,\nu} [P_x P_{x+\nu}^* + cc] + \sum_x [\eta P_x + \bar{\eta} P_x^*]$$

 τ : "temperature"

$$\eta = \kappa \exp(\mu)$$

$$\bar{\eta} = \kappa \exp(-\mu)$$

\Rightarrow solvable by a **worm** algorithm

[Delago Mercado, Evertz, Gattringer, PRL 106 (2011) 222001]



 \Rightarrow my choice here!

... thanks to Ydalia and Christof!

• How **strong** is the sign problem?

Study the overlap factor:

$$O(\mu) := \frac{Z(\mu)}{Z_{\text{mod}}(\mu)}$$

- \Rightarrow use the worm algorithm to solve the sign problem
- \Rightarrow Still need to solve an overlap problem

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- Use the **snake algorithm**:

$$\frac{Z(\mu + \Delta \mu)}{Z(\mu)} = \left\langle \exp\{S(\mu + \Delta \mu) - S(\mu)\}\right\rangle_{\mu}.$$
$$Z(5\Delta\mu) = \frac{Z(5\Delta\mu)}{Z(4\Delta\mu)} \frac{Z(4\Delta\mu)}{Z(3\Delta\mu)} \frac{Z(3\Delta\mu)}{Z(2\Delta\mu)} \frac{Z(2\Delta\mu)}{Z(\Delta\mu)} \frac{Z(\Delta\mu)}{Z(0)} Z(0)$$

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• need to caculate $Z(\mu)$ and $Z_{
m mod}(\mu)$

Results



Results



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• A bit of theory:

$$P \in Z_3: P = \exp\{i\frac{2\pi}{3}m\}$$

$$P \in \{1, (-1 + i\sqrt{3})/2, (-1 - i\sqrt{3})/2\}.$$

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• The imaginary part of the action:

$$\sum_{x} \left(\mathrm{e}^{\mu} P(x) + \mathrm{e}^{-\mu} P^{*}(x) \right) = \left(3N_{0} - V \right) \, \cosh(\mu) \, + \, i \, \sqrt{3} \, \sinh(\mu) \, \Delta N$$

- N_0 : number with P = 1
- N_+ : number with $P = (-1 + i\sqrt{3})/2$
- N_{-} : number with $P = (-1 i\sqrt{3})/2$
- $\Delta N = N_+ N_-$
- $V = N_0 + N_+ + N_-$: volume

• Infroduce - generalised density-of-states:

$$R(s) = \int \mathcal{D}P \,\delta\Big(s - \Delta N\Big) \,\exp\{S_R[P,\mu]\}$$
.

symmetries: R(s) = R(-s)

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• Overlap factor:

$$O(\mu) = \int ds \ R(s) \ \cos\left(\sqrt{3} \sinh(\mu) \ s\right) / \int ds \ R(s)$$

[no need to calculate the normalisation of R]

Results: histogram versus density-of-states method



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• The overlap $O(\mu)$:

Can we do rapidly oscillating integrals? $Z(\mu) = \int ds R(s) \cos(\sqrt{3} \sinh(\mu) s)$. $\mu = 2, \ \Delta N = s = 0 \dots 5000$

The overlap $O(\mu)$:

Can we do rapidly oscillating integrals?

 $Z(\mu) = \int ds R(\mathbf{f})$ $\mu = 2, \ \Delta N = s =$

yes!

> [if the integrand is boring enough]



• Method and error analysis:

key observation: $\ln R(s)$ very smooth function

$$R(s) = \prod_{s=1}^{s_0} e^{a(s_0)} \exp\{a(s_0)(s-s_0)\} \text{ for } s \in [s, s_0 + 1]$$

LLR-method $\Rightarrow a(s0)$

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LLR-method $\Rightarrow a(s0)$

• Bootstrap error analysis:

1. Choose a set of $a(s_0)$, $s_0 = 1 \dots s_{max} \approx 5000$

2. Fit a Polynomial: $\ln R(s) = \sum_{i \text{ even}}^{p} c_i s^i$, for p = 2, 4, 6, 8...

3. Get an answer for $O(\mu)$

Repeat steps 1-3 times and produce the boostrap average & standard deviation for $O(\mu)$.

Results for the overlap:



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 standard Solution: worm + snake algorithm

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 ⇒ standard Solution: worm + snake algorithm
- The density-of-states + LLR approach seems to work well
 ⇒ I am carefully optimistic
 - \Rightarrow Next test case: quantum O(2) model