HMC on Lefschetz thimbles-- A study of the residual sign problem

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in collaboration with

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based on

arXiv:1309.4371; JHEP10(2013)147

Feb. 20, 2014 @ GSI

Plan

I. Lattice models on Lefschetz thimbles (brief rev.)

- Pahm's result (Morse theory)
- Gradient flow, Critical points, Lefschetz thimbles
- * Residual sign problem: extra phase factor / Tangent spaces

2. An algorithm of HMC on Lefschetz thimbles

- a. how to parametrize/generate field conf. on the thimble
- b. how to formulate/solve the molecular dynamics on the thimble
- c. how to measure observables: reweighting the residual phase?

3. Test in the $\lambda \phi^4 \mu$ model

4. Summary & Discussions

Lattice models with complex-valued actions

- QCD with finite chemical potential
- [$e^{\mu a}$ a la Hasenfratz and Karsch]

Chiral gauge theories

[exact chiral gauge symmetry thanks to Ginsparg-Wilson rel.]

Chiral Yukawa theories

[reflection positivity in spite of Ginsparg-Wilson rel.]

..., etc.

physically well-defined, but the state-of-art Monte Carlo methods do not apply straightforwardly

Approaches to Lattice models with complex-valued actions

highly desirable to have a stochastic method which is based on a sound theoretical basis and applicable to these models

many methods proposed (and many analyses of the sign problem): rewighting; histgram; dual variables/worm algorithm; Tayler expansion in μ ; analytic continuation in μ (imaginary μ), etc.

One possible approach is to complexify the lattice models

$$\phi_x \in \mathbb{R} \longrightarrow z_x \in \mathbb{C}$$
 $U_{x\mu} = e^{iA_{x\mu}^a T^a} \in SU(3) \longrightarrow e^{iZ_{x\mu}^a T^a} \in SL(3, \mathbb{C})$

complexified Langevin dynamics

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Parisi (1983), Klauder (1983), ... (the old and classic approach) I.-O. Stamateschu et al., Phys. Rev. D75 045007 (2007), etc. G. Aarts, PRL 102(2009) 131601 ( \lambda\phi^4_{\,\,\mu} ) D. Sexty, arXiv:1307.7748 (QCD_{\mu})
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Path-Integral contours deformed to Lefschetz thimbles

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F. Pham (1983); E. Witten, arXiv:1001.2933;
M. Cristoforetti, F. Di Renzo, A. Mukherjee, L. Scorzato (AuroraScience Collaboration) Phys. Rev. D 86, 074506 (2012), arXiv:1205.3996 Phys. Rev. D88, 051501(2013), arXiv:1303.7204; 051502, arXiv:1308.0233;
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 F. Pham (1983); E.Witten, arXiv:1001.2933;
 M. Cristoforetti, F. Di Renzo, A. Mukherjee, L. Scorzato (AuroraScience Im(S) = const. !, HMC?! ;-)

Phys. Rev. D 86, 074506 (2012), arXiv:1205.3996
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Lattice models on Lefschetz thimbles

$$x \in \mathcal{C}_{\mathbb{R}} (\subseteq \mathbb{R}^n) \longrightarrow x + iy = z \in \mathbb{C}^n$$

$$S[x] \to S[x + iy] = S[z]$$

$$Z = \int_{\mathcal{C}_{\mathbb{R}}} \mathcal{D}[x] \exp\{-S[x]\} = \int_{\mathcal{C}} \mathcal{D}[z] \exp\{-S[z]\} \qquad (\mathcal{D}[x] = d^n x)$$

the contour of path-integration is selected by using the result of Morse theory [F. Pham (1983)]

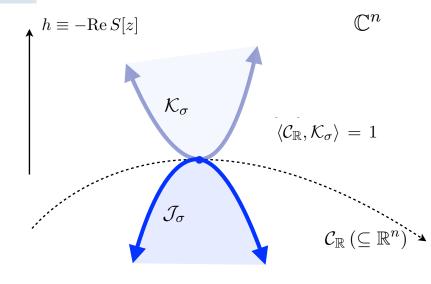
$$\mathcal{C}_{\mathbb{R}} = \sum_{\sigma \in \Sigma} n_{\sigma} \mathcal{J}_{\sigma}, \qquad n_{\sigma} = \langle \mathcal{C}_{\mathbb{R}}, \mathcal{K}_{\sigma} \rangle$$

$$h \equiv -\text{Re}\,S[z]$$

$$\frac{d}{dt}z(t) = \frac{\partial \bar{S}[\bar{z}]}{\partial \bar{z}}, \qquad \frac{d}{dt}\bar{z}(t) = \frac{\partial S[z]}{\partial z}, \qquad t \in \mathbb{R}$$

critical points
$$\mathbf{z}_{\sigma}$$
: $\frac{\partial S[z]}{\partial z}\Big|_{z=z_{\sigma}}=0$

Lefschetz thimble $\mathcal{J}_{\sigma}(\mathcal{K}_{\sigma})$ (n-dim. real mfd.) = the union of all down(up)ward flows which trace back to z_{σ} in the limit t goes to $-\infty$



$$\langle \mathcal{J}_{\sigma}, \mathcal{K}_{\tau} \rangle = \delta_{\sigma\tau}$$
 (intersection numbers)

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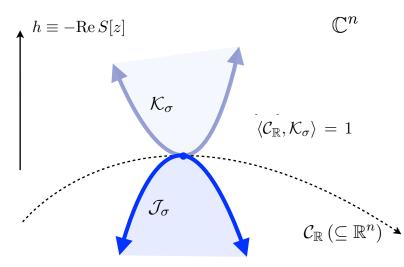
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$$\frac{d}{dt}h = -\frac{1}{2} \left\{ \frac{\partial S[z]}{\partial z} \cdot \frac{d}{dt} z(t) + \frac{\partial \bar{S}[\bar{z}]}{\partial \bar{z}} \cdot \frac{d}{dt} \bar{z}(t) \right\} = -\left| \frac{\partial S[z]}{\partial z} \right|^2 \le 0$$

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$$\frac{d}{dt}\operatorname{Im} S[z] = \frac{1}{2i} \left\{ \frac{\partial S[z]}{\partial z} \cdot \frac{d}{dt} z(t) - \frac{\partial \bar{S}[\bar{z}]}{\partial \bar{z}} \cdot \frac{d}{dt} \bar{z}(t) \right\} = 0$$



Partition function

$$Z = \sum_{\sigma \in \Sigma} n_{\sigma} \exp\{-S[z_{\sigma}]\} Z_{\sigma}, \qquad n_{\sigma} = \langle \mathcal{C}_{\mathbb{R}}, \mathcal{K}_{\sigma} \rangle$$
$$Z_{\sigma} = \int_{\mathcal{J}_{\sigma}} \mathcal{D}[z] \exp\{-\operatorname{Re}(S[z] - S[z_{\sigma}])\}$$

Observables

$$\langle O[z] \rangle = \frac{1}{Z} \sum_{\sigma \in \Sigma} n_{\sigma} \exp\{-S[z_{\sigma}]\} Z_{\sigma} \langle O[z] \rangle_{\mathcal{J}_{\sigma}}$$

$$\langle O[z] \rangle_{\mathcal{J}_{\sigma}} = \frac{1}{Z_{\sigma}} \int_{\mathcal{J}_{\sigma}} \mathcal{D}[z] \exp\{-\operatorname{Re}(S[z] - S[z_{\sigma}])\} O[z]$$

$$\langle \mathcal{C}_{\mathbb{R}}, \mathcal{K}_{\sigma} \rangle = 0$$

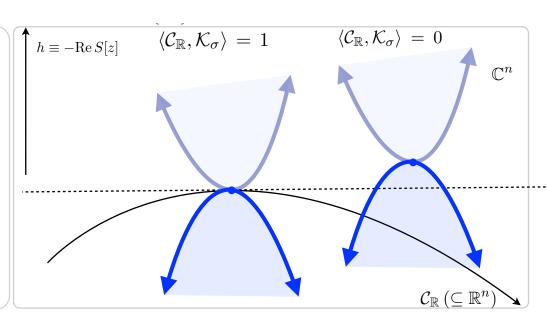
 $\{z_{\sigma}\} \text{ satisfying } -\text{Re}S[z_{\sigma}] > \max\{-\text{Re}S[x]\} (x \in \mathcal{C}_{\mathbb{R}})$

$$\langle \mathcal{C}_{\mathbb{R}}, \mathcal{K}_{\sigma} \rangle = 1$$

 $\{z_{\sigma}\}$ in the original cycle $\mathcal{C}_{\mathbb{R}}$

the relative weights proportional to $\exp(-S[z_{\sigma}])$

$$z_{\text{vac}} \in \mathcal{C}_{\mathbb{R}} - \text{Re} S[z_{\text{vac}}] = \max \{-\text{Re} S[x]\} (x \in \mathcal{C}_{\mathbb{R}})$$



$$\langle O[z] \rangle = \frac{1}{Z} \sum_{\sigma \in \Sigma} n_{\sigma} \exp\{-S[z_{\sigma}]\} Z_{\sigma} \langle O[z] \rangle_{\mathcal{J}_{\sigma}}$$

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It is not straightforward to compute the sum, in general

$$Z_{\sigma} = 1/\sqrt{\det K}$$

$$K_{ij} \equiv \partial_i \partial_j S[z]|_{z=z_\sigma}$$

in the saddle point approximation

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in the saddle point approximation

The functional measure should be specified by **the tangent spaces** of the thimble It may give rise to **an extra phase factor!** >> **residual sign problem**

if $\{U_z^{\alpha}\}$ is an orthonormal basis of the tangent space

$$\delta z = U_z^{\alpha} \delta \xi^{\alpha} \quad |\delta z|^2 = \delta \xi^2$$

$$d^n z |_{\mathcal{J}_{\sigma}} = d^n \delta \xi \det U_z$$

$$e^{i\phi_z} = \det U_z = \frac{\det V_z}{|\det V_z|}$$

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$$Z_{\sigma} = 1/\sqrt{\det K}$$

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in the saddle point approximation

Since Im(S) stays constant, this part may be evaluated by **MC**, but with the residual phase factor **reweighted** The functional measure should be specified by the tangent spaces of the thimble It may give rise to an extra phase factor!

>> residual sign problem

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in the saddle point approximation

Since Im(S) stays constant, this part may be evaluated by **MC**, but with the residual phase factor **reweighted**

a possible approximation : take a single thimble \mathcal{J}_{vac}

$$\langle O[z] \rangle = \langle O[z] \rangle_{\mathcal{J}_{\text{vac}}}$$

(AuroraScience Collaboration)

The functional measure should be specified by **the tangent spaces** of the thimble It may give rise to **an extra phase factor!** >> **residual sign problem**

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Geometric properties of Lefschetz thimbles

a) Tangent spaces of Lefschetz thimbles

basis of tangent vectors $\{V_z^{\alpha}\}(\alpha=1,\cdots,n)$

at a generic point z on \mathcal{J}_{σ}

$$\frac{d}{dt}V_{zi}^{\alpha}(t) = \bar{\partial}_i\bar{\partial}_j\bar{S}[\bar{z}] \ \bar{V}_{zj}^{\alpha}(t) \qquad (\alpha = 1, \dots, n)$$

In the vicinity of critical point z_{σ}

linearized flow equation and its solution:

$$\frac{d}{dt}(z_i(t) - z_{\sigma i}) = \bar{K}_{ij}(\bar{z}_j(t) - \bar{z}_{\sigma j}), \qquad K_{ij} \equiv \partial_i \partial_j S[z]|_{z=z_{\sigma}}$$

$$z_i(t) - z_{\sigma i} = v_i^{\alpha} \exp\left(\kappa^{\alpha}(t - t_0)\right) \xi_0^{\alpha}, \qquad \xi_0^{\alpha} \in \mathbb{R} \ (\alpha = 1, \dots, n)$$

$$\{v^{\alpha}\}(\alpha=1,\cdots,n)$$
 spans the tangent space $T_{z_{\sigma}}$

$$\{V_z\partial + \bar{V}_z\bar{\partial}\}V_z' - \{V_z'\partial + \bar{V}_z'\bar{\partial}\}V_z = 0$$

$$g \equiv \bar{\partial}\bar{S}[\bar{z}]$$

$$\{g\partial + \bar{g}\bar{\partial}\}V_z^{\alpha} - \{V_z^{\alpha}\partial + \bar{V}_z^{\alpha}\bar{\partial}\}g = 0$$

$$v_i^{\alpha} K_{ij} v_j^{\beta} = \kappa^{\alpha} \delta^{\alpha\beta}$$

 $\kappa^{\alpha} \ge 0 \ (\alpha = 1, \dots, n)$
 $v_i^{\alpha} (\alpha = 1, \dots, n)$ are orthonormal

$$\bar{V}_{zi}^{\alpha}V_{zi}^{\beta} - \bar{V}_{zi}^{\beta}V_{zi}^{\alpha} = 0 \qquad (\alpha, \beta = 1, \dots, n)$$

$$V_z^{\alpha} = U_z^{\beta} E^{\beta \alpha}$$
 $\{U_z^{\alpha}\}$ is an orthonormal basis

E is a real upper triangle matrix

$$\frac{d}{dt}\operatorname{Im}\{\bar{V}_{z}^{\alpha}(t)V_{z}^{\beta}(t)\}$$

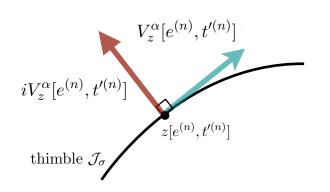
$$=\operatorname{Im}\{V_{z}^{\alpha}\partial^{2}S[z]V_{z}^{\beta}(t)+\bar{V}_{z}^{\alpha}\bar{\partial}^{2}\bar{S}[\bar{z}]\bar{V}_{z}^{\beta}(t)\}=0$$

b) Normal directions of thimbles

the set of normal vectors

$$\{iU_z^{\alpha}\}\ \text{or}\ \{iV_z^{\alpha}\}(\alpha=1,\cdots,n)$$

$$\operatorname{Re}\left\{ (-i)\bar{V}_{zi}^{\alpha}\,V_{zi}^{\beta}\right\} = 0$$



c) Parametrization of points z on thimbles

Asymptotic solutions of Flow equations

$$z(t) \simeq z_{\sigma} + v^{\alpha} \exp(\kappa^{\alpha} t) e^{\alpha}; \qquad e^{\alpha} e^{\alpha} = n$$

 $V_{z}^{\alpha}(t) \simeq v^{\alpha} \exp(\kappa^{\alpha} t),$

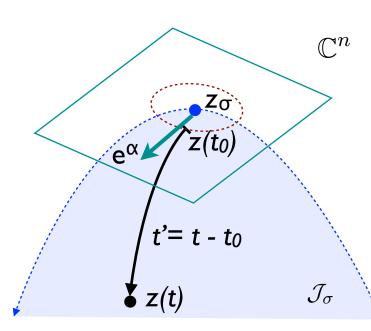
the **direction** of the flow : e^{α} $(\alpha = 1, \dots, n; ||e||^2 = n)$

the **time** of the flow : $t' = t - t_0$

$$z[e, t'] : (e^{\alpha}, t') \to z \in \mathcal{J}_{\sigma}$$

$$z[e, t'] = z(t)|_{t=t'+t_0}$$

$$\delta z[e, t'] = V_z^{\alpha}[e, t'] \left(\delta e^{\alpha} + \kappa^{\alpha} e^{\alpha} \delta t'\right)$$



Algorithm of HMC on Lefschetz thimbles

the saddle-point structures!

a) To generate a thimble

use the parameterization
$$z[e,t']:(e^{\alpha},t')\to z\in\mathcal{J}_{\sigma}$$
 solve the flow eqs. for **both** $z[e,t']$ & $V_z^{\alpha}[e,t']$ by 4th-order RK

- b) To formulate / solve the molecular dynamics introduce a dynamical system constrained to the thimble use 2nd-order constraint-preserving symmetric integrator
- c) To measure observables try to reweight the residual sign factors

$$\langle O[z] \rangle_{\mathcal{J}_{\sigma}} = \frac{\langle \mathrm{e}^{i\phi_z} O[z] \rangle_{\mathcal{J}_{\sigma}}'}{\langle \mathrm{e}^{i\phi_z} \rangle_{\mathcal{J}_{\sigma}}'} \qquad \text{where} \qquad \langle o[z] \rangle_{\mathcal{J}_{\sigma}}' = \frac{1}{N_{\mathrm{conf}}} \sum_{k=1}^{N_{\mathrm{conf}}} o[z^{(k)}]$$

$$\mathrm{e}^{i\phi_z} = \det U_z = \frac{\det V_z}{|\det V_z|}$$

 $\{\langle \mathrm{e}^{i\phi_z}\rangle_{\mathcal{J}_\sigma}'\}(\sigma\in\Sigma)$ should not be vanishingly small

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numerically very demanding!

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 $\{\langle \mathrm{e}^{i\phi_z}\rangle_{\mathcal{J}_\sigma}'\}(\sigma\in\Sigma)$ should not be vanishingly small

A possible sign problem! Need a careful and systematic study!

b) To formulate/solve Molecular Dynamics on the thimble

Constrained dynamical system

Equations of motion:

$$\dot{z}_i = w_i,$$

$$\dot{w}_i = -\bar{\partial}_i \bar{S}[\bar{z}] - iV_{zi}^{\alpha} \lambda^{\alpha} \qquad \lambda^{\alpha} \in \mathbb{R} \ (\alpha = 1, \dots, n)$$

Constraints:

$$z_i = z_i[e, t']$$
 $w_i = V_{zi}^{\alpha}[e, t'] w^{\alpha}, \quad w^{\alpha} \in \mathbb{R}$

A conserved Hamiltonian:

$$H = \frac{1}{2}\bar{w}_i w_i + \frac{1}{2} \left\{ S[z] + \bar{S}[\bar{z}] \right\}$$

$$\dot{H} = \frac{1}{2} \{ \dot{\bar{w}}_i w_i + \bar{w}_i \dot{w}_i \} + \frac{1}{2} \{ \partial_i S[z] \dot{z}_i + \bar{\partial}_i \bar{S}[\bar{z}] \dot{\bar{z}}_i \}$$

$$= \frac{1}{2} \{ (+i\bar{V}_{zi}^{\alpha} \lambda^{\alpha}) w_i + \bar{w}_i (-iV_{zi}^{\alpha} \lambda^{\alpha}) \}$$

$$= \frac{i}{2} \lambda^{\alpha} w^{\beta} \{ \bar{V}_{zi}^{\alpha} V_{zi}^{\beta} - \bar{V}_{zi}^{\beta} V_{zi}^{\alpha} \} = 0.$$

b) To formulate/solve Molecular Dynamics on the thimble

Second-order constraint-preserving symmetric integrator

$$z^{n} = z[e^{(n)}, t'^{(n)}],$$

 $w^{n} = V_{z}^{\alpha}[e^{(n)}, t'^{(n)}] w^{\alpha(n)}, \quad w^{\alpha(n)} \in \mathbb{R}.$

$$w^{n+1/2} = w^n - \frac{1}{2} \Delta \tau \, \bar{\partial} \bar{S}[\bar{z}^n] - \frac{1}{2} \Delta \tau \, iV_z^{\alpha}[e^{(n)}, t'^{(n)}] \, \lambda_{[r]}^{\alpha},$$

$$z^{n+1} = z^n + \Delta \tau \, w^{n+1/2},$$

$$w^{n+1} = w^{n+1/2} - \frac{1}{2} \Delta \tau \, \bar{\partial} \bar{S}[\bar{z}^{n+1}] - \frac{1}{2} \Delta \tau \, iV_z^{\alpha}[e^{(n+1)}, t'^{(n+1)}] \, \lambda_{[v]}^{\alpha}$$

 $\lambda^{\alpha}_{[r]}$ and $\lambda^{\alpha}_{[v]}$ are fixed by

$$z^{n+1} = z[e^{(n+1)}, t'^{(n+1)}],$$

$$w^{n+1} = V_z^{\alpha}[e^{(n+1)}, t'^{(n+1)}] w^{\alpha(n+1)}, \quad w^{\alpha(n+1)} \in \mathbb{R}$$

$$v_z^{\alpha}[e^{(n)}, t'^{(n)}] \lambda_{[r]}^{\alpha}$$

$$v_z^{\alpha}[e^{(n)}, t'^{(n)}]$$

$$v_z^{\alpha}[e^{(n)}, t'^{(n)}]$$

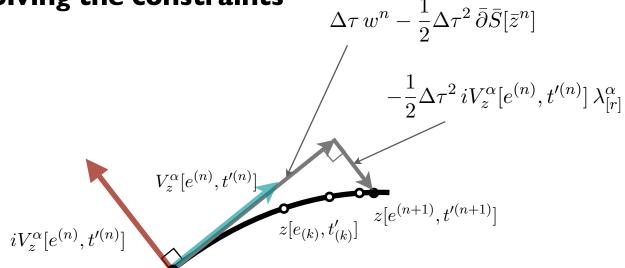
$$v_z^{\alpha}[e^{(n)}, t'^{(n)}]$$

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thimble \mathcal{J}_{σ}

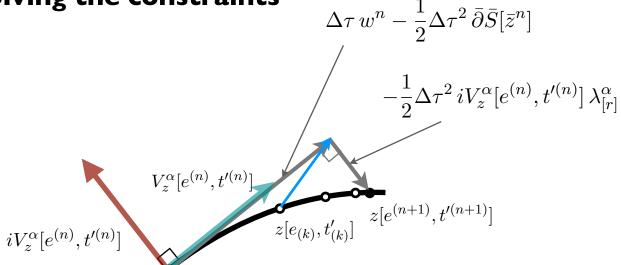
 $z[e^{(n)},t'^{(n)}]$

the sequences
$$(e_{(k)}^{\alpha}, t_{(k)}')$$
 $(k = 0, 1, \dots)$ with $(e_{(0)}^{\alpha}, t_{(0)}') = (e^{\alpha(n)}, t'^{(n)})$

$$\Delta e^{\alpha}_{(k)} = e^{\alpha}_{(k+1)} - e^{\alpha}_{(k)}, \qquad \sum_{\alpha=1} \Delta e^{\alpha}_{(k)} e^{\alpha(n)} = 0,$$

$$\Delta t'_{(k)} = t'_{(k+1)} - t'_{(k)},$$





 $z[e^{(n)},t'^{(n)}]$

the sequences
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 $(k = 0, 1, \dots)$ with $(e_{(0)}^{\alpha}, t_{(0)}') = (e^{\alpha(n)}, t'^{(n)})$

$$\Delta e^{\alpha}_{(k)} = e^{\alpha}_{(k+1)} - e^{\alpha}_{(k)}, \qquad \sum_{\alpha=1} \Delta e^{\alpha}_{(k)} e^{\alpha(n)} = 0,$$

$$\Delta t'_{(k)} = t'_{(k+1)} - t'_{(k)},$$



 $z[e^{(n)}, t'^{(n)}]$

$$\Delta \tau \, w^n - \frac{1}{2} \Delta \tau^2 \, \bar{\partial} \bar{S}[\bar{z}^n]$$

$$-\frac{1}{2} \Delta \tau^2 \, i V_z^{\alpha}[e^{(n)}, t'^{(n)}] \, \lambda_{[r]}^{\alpha}$$

$$i V_z^{\alpha}[e^{(n)}, t'^{(n)}] \qquad z[e^{(n+1)}, t'^{(n+1)}]$$

thimble \mathcal{J}_{σ}

the sequences $(e_{(k)}^{\alpha}, t_{(k)}')$ $(k = 0, 1, \dots)$ with $(e_{(0)}^{\alpha}, t_{(0)}') = (e^{\alpha(n)}, t'^{(n)})$

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$$\Delta z_{(k)}[e^{(n)}, t'^{(n)}] \equiv z_i[e^{(n)}, t'^{(n)}] + \Delta \tau \, w_i^n - \frac{1}{2} \Delta \tau^2 \, \bar{\partial}_i \bar{S}[\bar{z}^n] - z_i[e_{(k)}, t'_{(k)}]$$



$$\Delta \tau \, w^n - \frac{1}{2} \Delta \tau^2 \, \bar{\partial} \bar{S}[\bar{z}^n] \\ - \frac{1}{2} \Delta \tau^2 \, i V_z^\alpha [e^{(n)}, t'^{(n)}] \, \lambda_{[r]}^\alpha \\ V_z^\alpha [e^{(n)}, t'^{(n)}] \\ z[e_{(k)}, t'_{(k)}] \\ z[e^{(n+1)}, t'^{(n+1)}]$$

$$z[e^{(n)},t'^{(n)}]$$
 thimble \mathcal{J}_{σ}

 $iV_z^{\alpha}[e^{(n)},t'^{(n)}]$

the sequences
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 $z[e^{(n)}, t'^{(n)}]$

$$\Delta \tau \, w^n - \frac{1}{2} \Delta \tau^2 \, \bar{\partial} \bar{S}[\bar{z}^n]$$

$$-\frac{1}{2} \Delta \tau^2 \, i V_z^\alpha [e^{(n)}, t'^{(n)}] \, \lambda_{[r]}^\alpha$$

$$V_z^\alpha [e^{(n)}, t'^{(n)}]$$

$$z[e_{(k)}, t'_{(k)}] \, z[e^{(n+1)}, t'^{(n+1)}]$$

$$iV_z^{\alpha}[e^{(n)},t'^{(n)}]$$

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$$\Delta \tau w^{n} - \frac{1}{2} \Delta \tau^{2} \, \bar{\partial} \bar{S}[\bar{z}^{n}]$$

$$-\frac{1}{2} \Delta \tau^{2} \, iV_{z}^{\alpha}[e^{(n)}, t'^{(n)}] \, \lambda_{[r]}^{\alpha}$$

$$V_z^{\alpha}[e^{(n)}, t'^{(n)}]$$

 $z[e^{(n)}, t'^{(n)}]$

 $iV_z^{\alpha}[e^{(n)},t'^{(n)}]$

$$z[e_{(k)}, t'_{(k)}]$$
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$$\Delta \tau \, w^n - \frac{1}{2} \Delta \tau^2 \, \bar{\partial} \bar{S}[\bar{z}^n]$$

 $z[e_{(k)}, t'_{(k)}]$ $z[e^{(n+1)}, t'^{(n+1)}]$

$$-\frac{1}{2} \Delta \tau^2 i V_z^{\alpha} [e^{(n)}, t'^{(n)}] \lambda_{[r]}^{\alpha}$$

$$V_z^{\alpha}[e^{(n)}, t'^{(n)}]$$

 $z[e^{(n)},t'^{(n)}]$

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$$\Delta z_{(k)}[e^{(n)}, t'^{(n)}]_{\perp} = iV_z^{\alpha}[e^{(n)}, t'^{(n)}] \left(\frac{1}{2} \Delta \tau^2 \lambda_{[r](k)}^{\alpha}\right)$$



$$\Delta \tau \, w^n - \frac{1}{2} \Delta \tau^2 \, \bar{\partial} \bar{S}[\bar{z}^n]$$

 $z[e_{(k)}, t'_{(k)}]$ $z[e^{(n+1)}, t'^{(n+1)}]$

$$-\frac{1}{2}\Delta\tau^2 iV_z^{\alpha}[e^{(n)}, t'^{(n)}] \lambda_{[r]}^{\alpha}$$

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 $z[e^{(n)}, t'^{(n)}]$

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$$\Delta z_{(k)}[e^{(n)}, t'^{(n)}]_{\perp} = iV_z^{\alpha}[e^{(n)}, t'^{(n)}] \left(\frac{1}{2} \Delta \tau^2 \lambda_{[r](k)}^{\alpha}\right)$$

$$\left\| V_z^{\alpha}[e^{(n)}, t'^{(n)}] \left(\Delta e^{\alpha}_{(k)} + e^{\alpha(n)} \kappa^{\alpha} \Delta t'_{(k)} \right) \right\|^2 \le n \epsilon'^2$$



$$\Delta \tau \, w^n - \frac{1}{2} \Delta \tau^2 \, \bar{\partial} \bar{S}[\bar{z}^n]$$

 $z[e_{(k)}, t'_{(k)}]$ $z[e^{(n+1)}, t'^{(n+1)}]$

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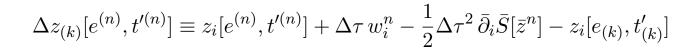
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thimble \mathcal{J}_{σ}

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the sequences
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the constraints to be solved

$$z[e^{(n+1)}, t'^{(n+1)}] - z[e^{(n)}, t'^{(n)}] = \Delta \tau w^n - \frac{1}{2} \Delta \tau^2 \, \bar{\partial} \bar{S}[\bar{z}^n] - \frac{1}{2} \Delta \tau^2 \, iV_z^{\alpha}[e^{(n)}, t'^{(n)}] \, \lambda_{[r]}^{\alpha}$$

the sequences
$$(e_{(k)}^{\alpha}, t'_{(k)})$$
 $(k = 0, 1, \cdots)$ with $(e_{(0)}^{\alpha}, t'_{(0)}) = (e^{\alpha(n)}, t'^{(n)})$
$$\Delta e_{(k)}^{\alpha} = e_{(k+1)}^{\alpha} - e_{(k)}^{\alpha}, \qquad \sum_{\alpha=1} \Delta e_{(k)}^{\alpha} e^{\alpha(n)} = 0,$$

$$\Delta t'_{(k)} = t'_{(k+1)} - t'_{(k)},$$

where

$$\Delta e^{\alpha}{}_{(k)} + e^{\alpha(n)} \kappa^{\alpha} \Delta t'_{(k)} = \text{Re} \left[\{ V_z^{-1} [e^{(n)}, t'^{(n)}] \}_i^{\alpha} \times \left[z_i [e^{(n)}, t'^{(n)}] + \Delta \tau \, w_i^n - \frac{1}{2} \Delta \tau^2 \, \bar{\partial}_i \bar{S}[\bar{z}^n] - z_i [e_{(k)}, t'_{(k)}] \right] \right]$$

$$\frac{1}{2} \Delta \tau^2 \, \lambda_{[r]_{(k)}}^{\alpha} = \text{Im} \left[\{ V_z^{-1} [e^{(n)}, t'^{(n)}] \}_i^{\alpha} \left(z_i [e^{(n)}, t'^{(n)}] - z_i [e_{(k)}, t'_{(k)}] \right) \right]$$

stopping cond.:

$$\left\| V_z^{\alpha}[e^{(n)}, t'^{(n)}] \left(\Delta e^{\alpha}_{(k)} + e^{\alpha(n)} \kappa^{\alpha} \Delta t'_{(k)} \right) \right\|^2 \le n \epsilon'^2$$

$$\frac{1}{2}\Delta\tau \,\lambda_{[v]}^{\alpha} = \operatorname{Im}\left[\left\{V_{z}^{-1}[e^{(n+1)}, t'^{(n+1)}]\right\}_{i}^{\alpha}\left(w_{i}^{n+1/2} - \frac{1}{2}\Delta\tau \,\bar{\partial}_{i}\bar{S}[\bar{z}^{n+1}]\right)\right]$$

a HMC update

A hybrid Monte Carlo update then consists of the following steps for a given trajectory length τ_{traj} and a number of steps n_{step} :

1. Set the initial field configuration z_i :

$$\{e^{\alpha(0)}, t'^{(0)}\} = \{e^{\alpha}, t'\}, \qquad z^0 = z[e, t'].$$

2. Refresh the momenta w_i by generating n pairs of unit gaussian random numbers (ξ_i, η_i) , setting tentatively $w_i = \xi_i + i\eta_i$, and chopping the non-tangential parts:

$$w^{0} = V_{z}^{\alpha} \operatorname{Re}[\{V_{z}^{-1}\}_{j}^{\alpha} (\xi_{j} + i\eta_{j})] = U_{z}^{\alpha} \operatorname{Re}[\{U_{z}^{-1}\}_{j}^{\alpha} (\xi_{j} + i\eta_{j})].$$

- 3. Repeat n_{step} times of the second order symmetric integration the step size $\Delta \tau = \tau_{\text{traj}}/n_{\text{step}}$.
- 4. Accept or reject by $\Delta H = H[w^{n_{\text{step}}}, z^{n_{\text{step}}}] H[w^0, z^0].$

As for the initialization procedure, one may generate unit gaussian random numbers $\eta^{\alpha}(\alpha=1,\cdots,n)$, set

$$e^{\alpha} = \eta^{\alpha} \sqrt{\frac{n}{\sum_{\beta=1}^{n} \eta^{\beta} \eta^{\beta}}}, \qquad t' = -t_0,$$

and then prepare z[e,t'], $\{V_z^{\alpha}[e,t']\},$ and the inverse matrix $V_z^{-1}[e,t'].$

Test in the $\lambda \phi^4 \mu$ model

Complex Langevin simulation

G.Aarts, PRL 102:131601, 2009, arXiv:0810.2089

Dual variables / worm algorithm

C. Gattringer and T. Kolber, NP B869 (2013) 56, arXiv:1206.2954

$$\varphi(x) = (\phi_1(x) + i\phi_2(x))/\sqrt{2}$$

$$\phi_a(x) \in \mathbb{R} \ (a = 1, 2)$$

$$\phi_a(x) \to z_a(x) \in \mathbb{C} \ (a=1,2)$$

$$S[z] = \sum_{x \in \mathbb{L}^4} \left\{ + \frac{1}{2} z_a(x) z_a(x) + \frac{\lambda_0}{4} \left(z_a(x) z_a(x) \right)^2 - K_0 \sum_{k=1}^3 z_a(x) z_a(x + \hat{k}) - K_0 z_a(x) z_b(x + \hat{0}) \left[\delta_{ab} \cosh(\mu) - i \epsilon_{ab} \sinh(\mu) \right] \right\}.$$

where
$$K_0 = \frac{1}{(2D+\kappa)}$$
, $\lambda_0 = K_0^2 \lambda$

$$K=1.0, \lambda=1.0, \mu=0.0\sim1.8$$

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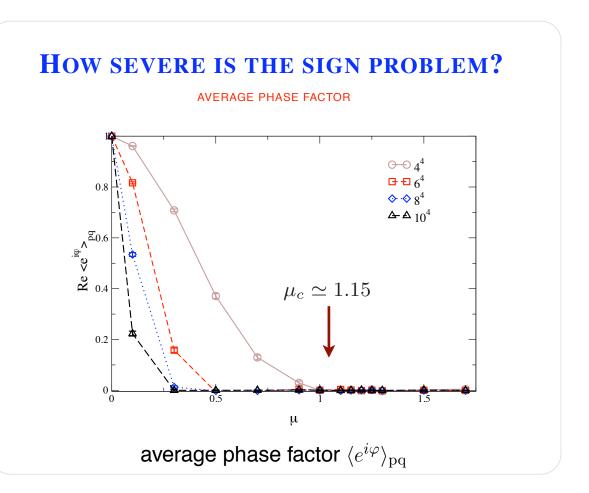
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$$S[z] = \sum_{x \in \mathbb{L}^4} \left\{ + \frac{1}{2} z_a(x) z_a(x) + \frac{\lambda_0}{4} \left(z_a(x) z_a(x) z_a(x) z_b(x) \right) \right\}$$

$$-K_0 z_a(x) z_b(x + \hat{0}) \left[\delta_{ab} \cos \theta \right]$$
where $K_0 = \frac{1}{(2D + \kappa)}$, $\lambda_0 = K_0^2 \lambda$

$$K=1.0, \lambda=1.0, \mu=0.0\sim1.8$$

L=4 (, ... 12)



G. Aarts, PRL 102:131601, 2009, arXiv:0810.2089

Test in the $\lambda \phi^4 \mu$ model (cont'd)

critical points with constant field $z_a(x)=z_a$

$$\left. \frac{\partial S[z]}{\partial z_a(x)} \right|_{z_a(x)=z_a} = (1 - 6K_0 - 2K_0 \cosh(\mu)) z_a + \lambda_0 (z_1^2 + z_2^2) z_a = 0 \quad (a = 1, 2).$$

critical value of
$$\mu$$
 (classical) $\tilde{\mu}_c = \ln \left[\left(\frac{1 - 6K_0}{2K_0} \right) + \sqrt{\left(\frac{1 - 6K_0}{2K_0} \right)^2 - 1} \right]$

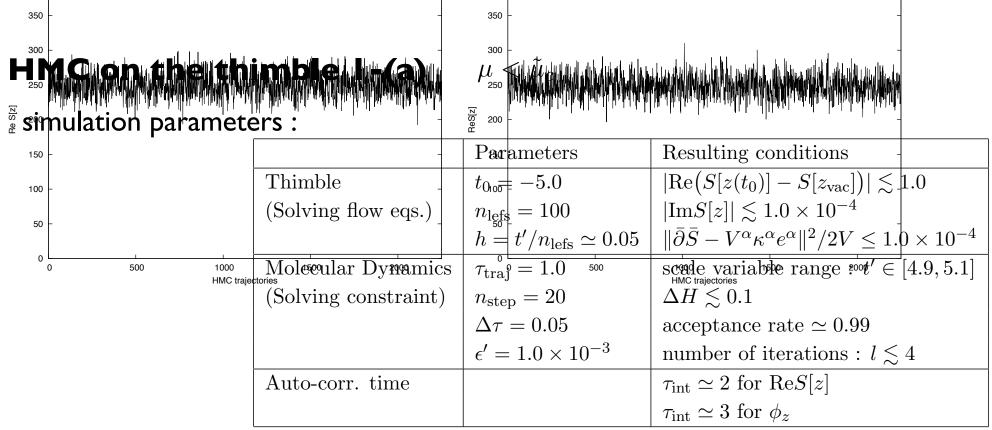
- 1. For $\mu \leq \tilde{\mu}_c$,
 - (a) $z_1 = z_2 = 0$; S[z] = 0,
 - (b) $z_1 = i\phi_0 \cos \theta$, $z_2 = i\phi_0 \sin \theta$; $S[z] = -L^4 \frac{\lambda_0}{4} \phi_0^4$, where $\phi_0 = \sqrt{\frac{+(1-6K_0-2K_0\cosh(\mu))}{\lambda_0}}$.
- 2. For $\mu > \tilde{\mu}_c$,
 - (a) $z_1 = z_2 = 0$; S[z] = 0,
 - (b) $z_1 = \phi_0 \cos \theta$, $z_2 = \phi_0 \sin \theta$; $S[z] = -L^4 \frac{\lambda_0}{4} \phi_0^4$, where $\phi_0 = \sqrt{\frac{-(1 6K_0 2K_0 \cosh(\mu))}{\lambda_0}}$.

 $\mu_c \sim 0.962$ for K=1.0, λ =1.0

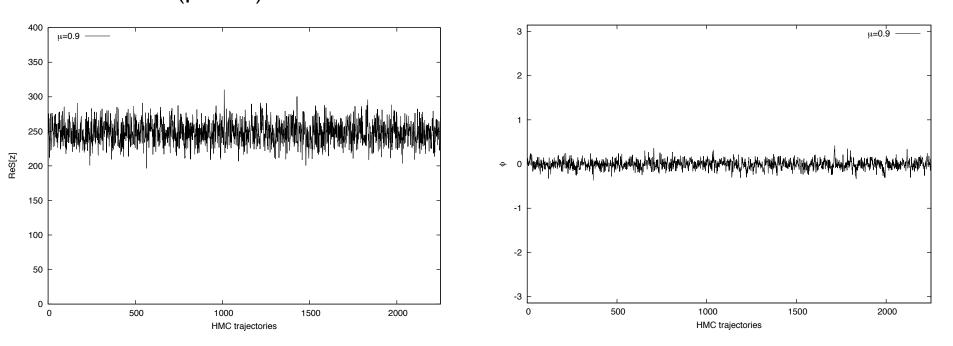
the thimble I-(a)

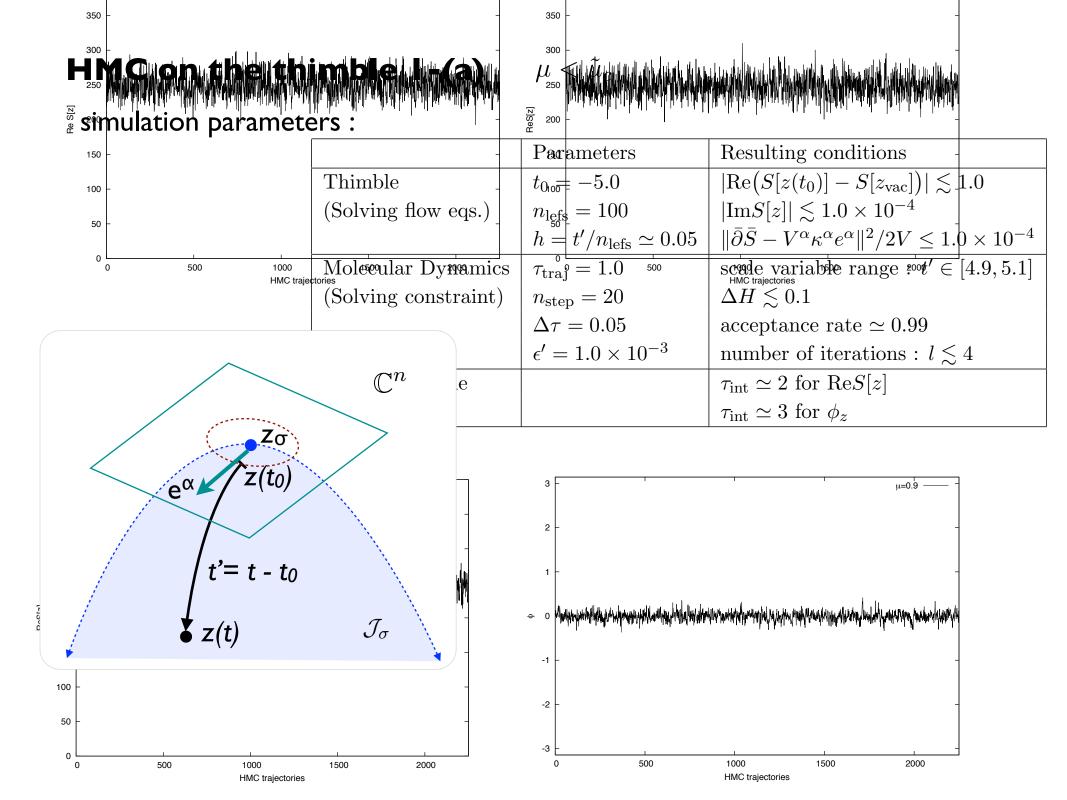
 \leftarrow the thimble 2-(a)

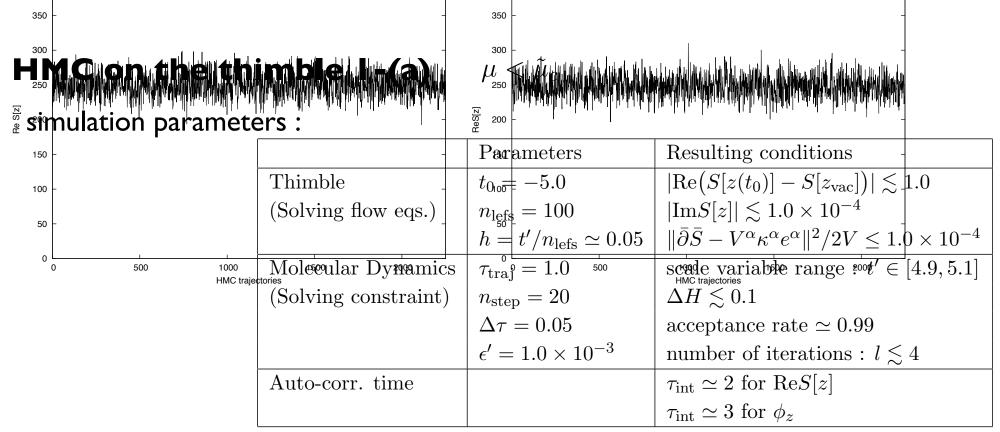
 \leftarrow the thimble 2-(b)



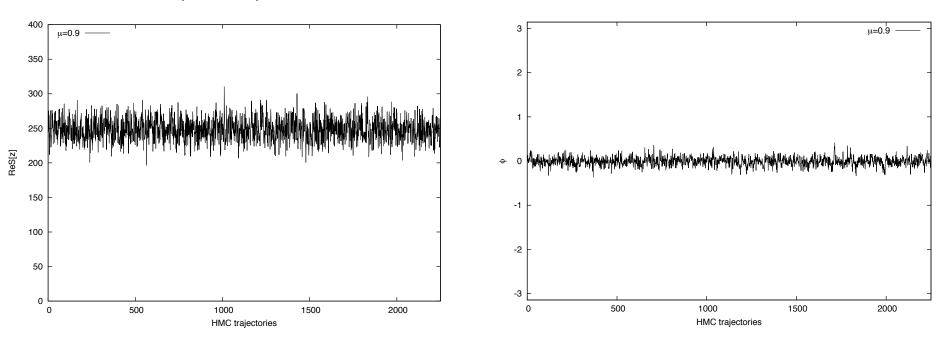
HMC histories ($\mu = 0.9$)

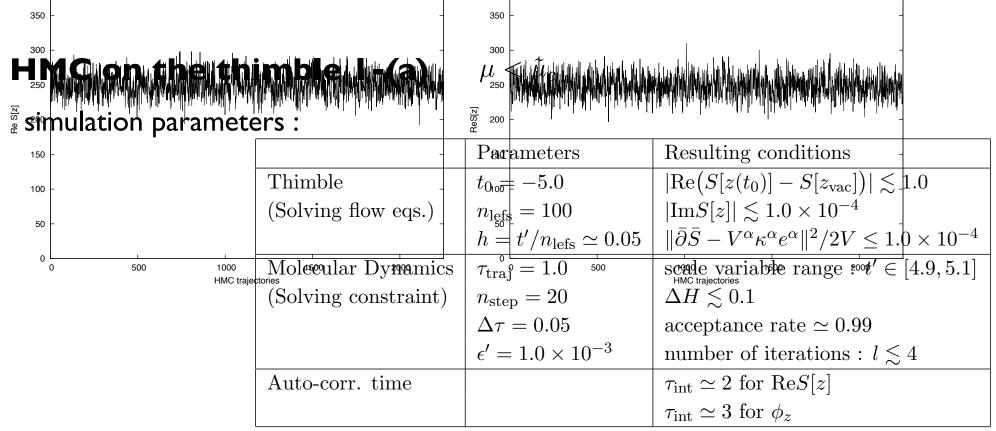




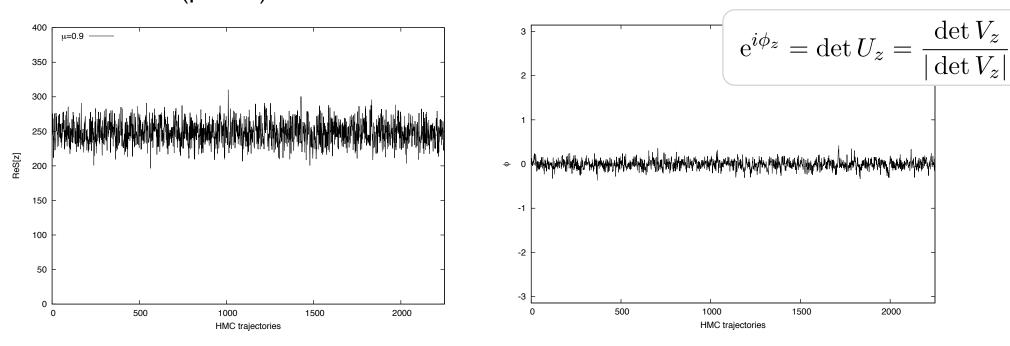


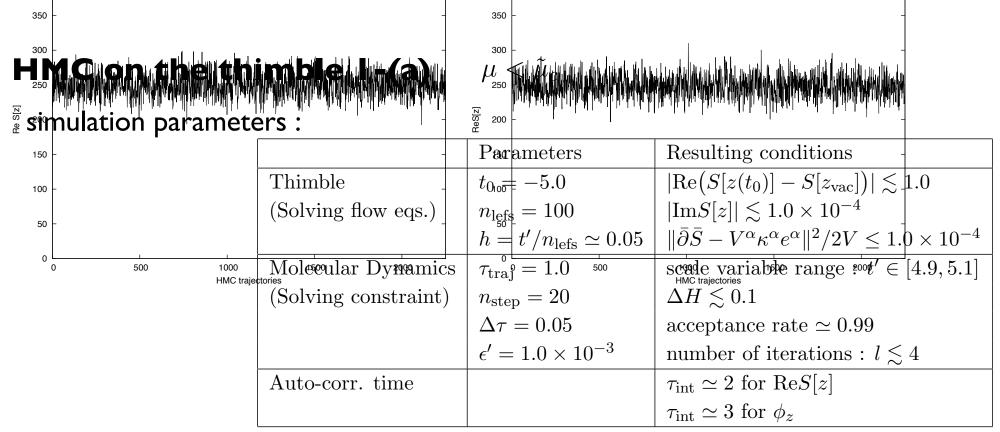
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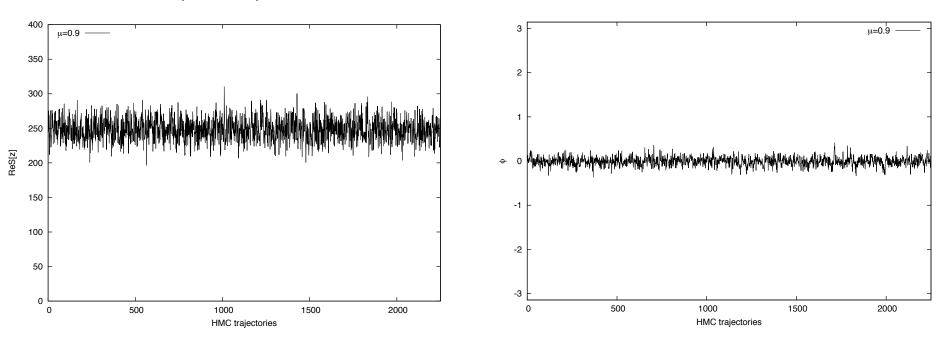


HMC histories ($\mu = 0.9$)





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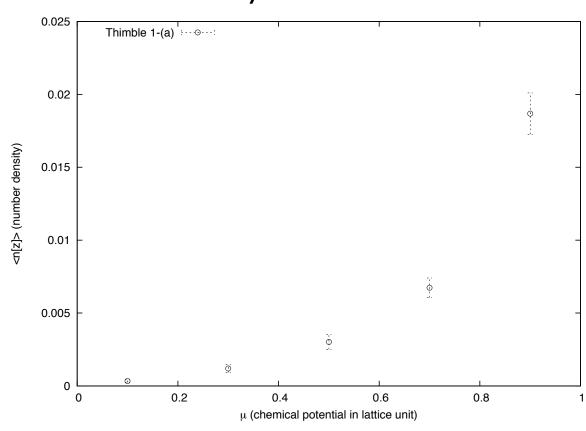


generated 4,250 traj. sampling 300 conf. with the separation of 10

residual phase:

μ	$\langle \mathrm{e}^{i\phi_z} angle_{\mathcal{J}_{\mathrm{vac}}}'$
0.1	$(9.99e-01, -1.15e-03) \pm (5.7e-02, 7.4e-04)$
0.3	$(9.99e-01, -1.03e-03) \pm (5.7e-02, 2.1e-03)$
0.5	$(9.98e-01, -2.68e-03) \pm (5.7e-02, 3.3e-03)$
0.7	$(9.97e-01, 5.24e-04) \pm (5.7e-02, 4.3e-03)$
0.9	$(9.94e-01, -7.40e-03) \pm (5.7e-02, 5.9e-03)$

$$e^{i\phi_z} = \det U_z = \frac{\det V_z}{|\det V_z|}$$



$$n[z] = \frac{1}{L^4} \sum_{x} K_0 z_a(x) z_b(x + \hat{0}) \left[\delta_{ab} \sinh(\mu) - i\epsilon_{ab} \cosh(\mu) \right]$$

HMC on the thimble 2-(b) $\mu > \mu_c$

Critical region of real dimension one : $\theta \in [0, 2\pi]$

$$z_a(x;t) \simeq R_{ab}(\theta) \left\{ \delta_{b1} \phi_0 + \sum_{\beta=1}^{2V-1} v_b(x)^{\beta} \exp(\kappa^{\beta} t) e^{\beta} \right\} \qquad (t \ll 0)$$
$$\delta z_a(x;t) = V_a(x;t)^0 \left(\phi_0 \sqrt{V} \delta \theta \right) + \sum_{\beta=1}^{2V-1} V_b(x;t)^{\beta} (\delta e^{\beta} + \kappa^{\beta} e^{\beta} \delta t)$$

zero mode $\kappa^0 = 0$ $v_a(x)^0 = \delta_{a2}/\sqrt{V}$

Critical fluctuation : lowest mode
$$\frac{\kappa^1=2\lambda_0\phi_0^2}{v_a(x)^1=\delta_{a1}/\sqrt{V}}$$
 gets very light! $(\mu\gtrsim \tilde{\mu}_c)$

$$z_a(x;t) \simeq R_{ab}(\theta) \left\{ \delta_{b1} \frac{\phi_0}{\sqrt{1 - \frac{2}{\sqrt{V}\phi_0}} e^1 \exp(\kappa^1 t)} + \sum_{\beta=2}^{2V-1} v_b(x)^\beta \exp(\kappa^\beta t) e^\beta \right\} \qquad \sum_{\beta=2}^{2V-1} e^\beta e^\beta = 2V-2$$

$$V_{a}(x;t)^{0} \simeq R_{ab}(\theta) v_{b}(x)^{0} \frac{1}{\sqrt{1 - \frac{2}{\sqrt{V}\phi_{0}} e^{1} \exp(\kappa^{1}t)}},$$

$$V_{a}(x;t)^{1} \simeq R_{ab}(\theta) v_{b}(x)^{1} \frac{\exp(\kappa^{1}t)}{\left(1 - \frac{2}{\sqrt{V}\phi_{0}} e^{1} \exp(\kappa^{1}t)\right)^{3/2}},$$

$$V_{a}(x;t)^{\beta} \simeq R_{ab}(\theta) v_{b}(x)^{\beta} \exp(\kappa^{\beta}t) \qquad (\beta = 2, \dots, 2V - 1)$$

the global flow mode $z_a(x;t) = z_a(t)$

$$\frac{d}{dt}z_a(t) = \bar{\partial}_{ax}\bar{S}[\bar{z}]\big|_{z_a(x;t)=z_a(t)}$$

$$= \lambda_0 (\bar{z}_b(t)\bar{z}_b(t) - \phi_0^2)\bar{z}_a(t)$$

HMC on the thimble 2-(b) $\mu > \tilde{\mu}_c$

simulation parameters:

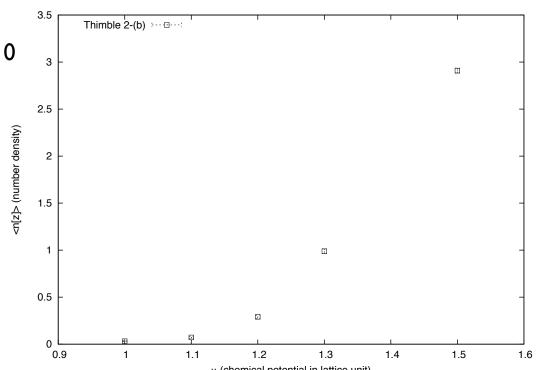
	Parameters	Resulting conditions
Thimble	$t_0 = -3.0$	$ \operatorname{Re}(S[z(t_0)] - S[z_{\operatorname{vac}}]) \lesssim 2.0 \times 10^1$
	$n_{\mathrm{lefs}} = 100$	$ \operatorname{Im}(S[z] - S[z_{\operatorname{vac}}]) \lesssim 5.0 \times 10^{-2}$
	$h = t'/n_{\mathrm{lefs}} \simeq 0.03$	$\ \bar{\partial}\bar{S} - V^{\alpha}\kappa^{\alpha}e^{\alpha}\ ^2/2V \le 3.0 \times 10^{-2}$
MD	$ au_{ m traj} = 0.3$	$t' \in [2.5, 3.5]$
	$n_{\text{step}} = 10, 30 \ (\mu = 1.0, 1.1)$	$\Delta H \lesssim 0.05$
	$\Delta \tau = 0.03, 0.01 \ (\mu = 1.0, 1.1)$	Acceptance rate $\simeq 0.99$
	$\epsilon' = \sqrt{10} \times 10^{-3}$	$l \lesssim 4, 6 \ (\mu = 1.0), 14 \ (\mu = 1.1)$
Auto-corr. time	(for $ReS[z]$)	$\tau_{\rm int} \simeq 10, 14 \; (\mu = 1.0, 1.1)$
	$(\text{for }\phi_z)$	$\tau_{\text{int}} \simeq 15, 14 \ (\mu = 1.0), 28 \ (\mu = 1.1)$

number density:

generated 11,250 traj. sampling 1,000 conf. with the separation of 10

residual phase averages:

μ	$\langle { m e}^{i\phi_z} angle_{{\cal J}_{ m vac}}'$
1.0	$(9.94e-01, -8.77e-03) \pm (3.1e-02, 3.1e-03)$
1.1	$(9.94e-01, -3.21e-03) \pm (3.1e-02, 3.4e-03)$
1.2	$(9.95e-01, -8.25e-04) \pm (3.1e-02, 3.0e-03)$
1.3	$(9.97e-01, -3.08e-03) \pm (3.1e-02, 2.2e-03)$
1.5	$(9.99e-01, -1.06e-03) \pm (3.1e-02, 1.0e-03)$



HMC on the thimble 2-(b) $\mu > \tilde{\mu}_c$

simulation parameters:

	Parameters	Resulting conditions
Thimble	$t_0 = -3.0$	$ \operatorname{Re}(S[z(t_0)] - S[z_{\operatorname{vac}}]) \lesssim 2.0 \times 10^1$
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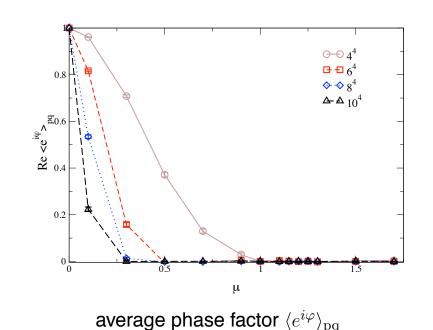
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HOW SEVERE IS THE SIGN PROBLEM?

AVERAGE PHASE FACTOR



HMC on the thimble 2-(b) $\mu > \tilde{\mu}_c$

simulation parameters:

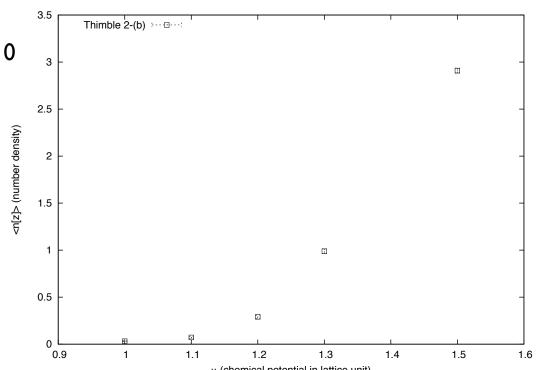
	Parameters	Resulting conditions
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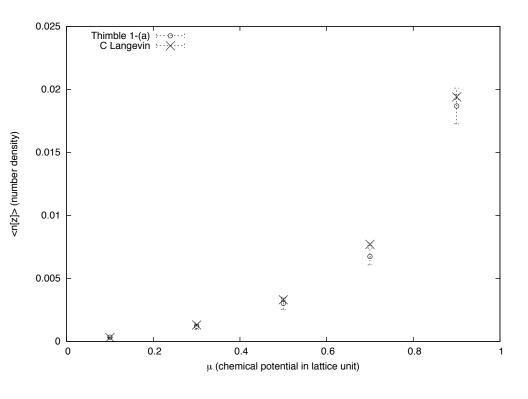


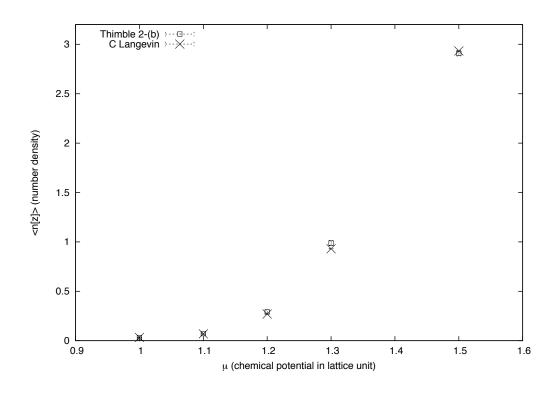
Comparison to Complex Langevin simulations

$$\frac{dz(t)}{dt} = -\frac{\partial S[z]}{\partial z} + \eta(t); \quad \langle \eta(t)\eta(t') \rangle = 2\delta(t - t')$$

$$\langle \mathcal{O} \rangle = \lim_{t \to \infty} \frac{1}{t} \int_0^t dt' \, \mathcal{O}(z(t'))$$

parameters of CL simulations: step size ϵ =5.0 x 10⁻⁵ , 5,000,000 time steps sampling 10,000 configurations with the separation of 500



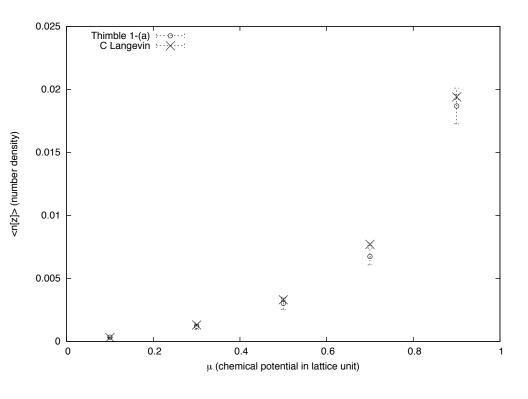


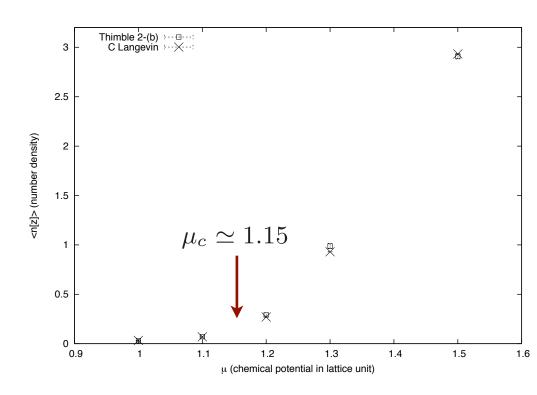
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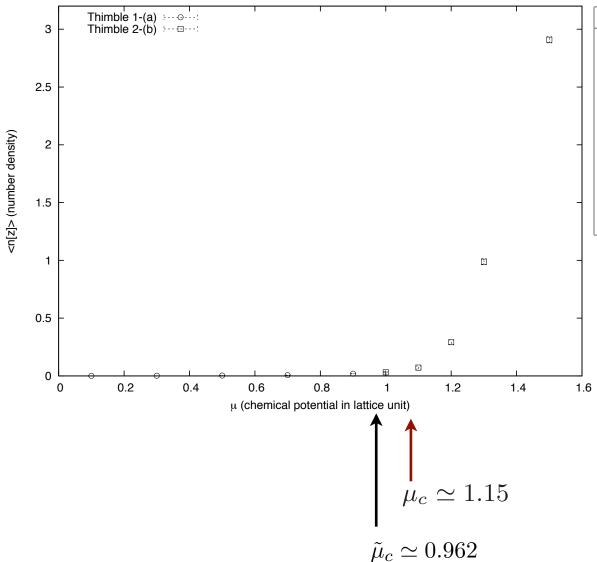
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HMC on the thimbles I-(a) & 2-(b)



μ	Re $\langle n[z] \rangle_{\mathcal{J}_{\text{vac}}}$ (jk. error)	Re $\langle e^{i\phi_z} n[z] \rangle'_{\mathcal{J}_{vac}}$	Re $\langle n[z] \rangle'_{\mathcal{J}_{\text{vac}}}$
0.1	3.34e-04 (9.2e-05)	3.35e-04	2.15e-04
0.3	1.20e-03 (2.7e-04)	1.19e-03	8.56e-04
0.5	3.02e-03 (5.0e-04)	3.01e-03	2.44e-03
0.7	6.74e-03 (6.7e-04)	6.71e-03	5.91e-03
0.9	1.89e-02 (1.4e-03)	1.85e-02	1.73e-02
1.0	3.14e-02 (4.3e-03)	3.12e-02	3.00e-02
1.1	7.17e-02 (1.3e-02)	7.12e-02	7.01e-02
1.2	2.92e-01 (1.8e-02)	2.90e-01	2.90e-01
1.3	9.88e-01 (2.6e-02)	9.85e-01	9.87e-01
1.5	2.91e-00 (2.7e-02)	2.90e-00	2.90e-00

- We have formulated a HMC algorithm which is applicable to lattice models defined on Lefschetz thimbles
- We have tested the algorithm in the $\lambda\phi^4\,_{\mu}$ model on the lattice V=4^4
 - the thimbles associated with the classical vacua
 - the residual phase factors reweighted successfully
 - known results of the number density reproduced (cf. CL, dual v.)
 - Need the careful study of the systematic errors
 - setup of the asymptotic regions
 - contributions of other thimbles, ex. thimble 2-(a), ...
 - Need the study of the residual sign problem on larger lattices
- Numerical cost per traj.: literally, scales as $O(V^3 \times n_{step})$ solving flow eqs. (all tangent vectors): $O(V^2 \times n_{Lefs})$ computing V^{-1} , det V (residual sign factors): $O(V^3)$
- Dynamical fermions:
 possible applications to QCD μ cf. D. Sexty, arXiv:1307.7748

Test in the $\lambda \phi^4 \mu$ model (cont'd)

critical points with constant field $z_a(x)=z_a$

$$\left. \frac{\partial S[z]}{\partial z_a(x)} \right|_{z_a(x)=z_a} = (1 - 6K_0 - 2K_0 \cosh(\mu)) z_a + \lambda_0 (z_1^2 + z_2^2) z_a = 0 \quad (a = 1, 2)$$

critical value of
$$\mu$$
 (classical) $\tilde{\mu}_c = \ln \left[\left(\frac{1 - 6K_0}{2K_0} \right) + \sqrt{\left(\frac{1 - 6K_0}{2K_0} \right)^2 - 1} \right]$

- 1. For $\mu \leq \tilde{\mu}_c$,
 - (a) $z_1 = z_2 = 0$; S[z] = 0,
 - (b) $z_1 = i\phi_0 \cos \theta$, $z_2 = i\phi_0 \sin \theta$; $S[z] = -L^4 \frac{\lambda_0}{4} \phi_0^4$, where $\phi_0 = \sqrt{\frac{+(1-6K_0-2K_0 \cosh(\mu))}{\lambda_0}}$.
- 2. For $\mu > \tilde{\mu}_c$,
 - (a) $z_1 = z_2 = 0$; S[z] = 0,
 - (b) $z_1 = \phi_0 \cos \theta$, $z_2 = \phi_0 \sin \theta$; $S[z] = -L^4 \frac{\lambda_0}{4} \phi_0^4$, where $\phi_0 = \sqrt{\frac{-(1 6K_0 2K_0 \cosh(\mu))}{\lambda_0}}$.

 $\mu_c \sim 0.962$ for K=1.0, λ =1.0

 \leftarrow the thimble I-(a)

 \leftarrow the thimble 2-(a)

 \leftarrow the thimble 2-(b)

- We have formulated a HMC algorithm which is applicable to lattice models defined on Lefschetz thimbles
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 possible applications to QCD μ cf. D. Sexty, arXiv:1307.7748

- We have formulated a HMC algorithm which is applicable to lattice models defined on Lefschetz thimbles
- We have tested the algorithm in the $\lambda \phi^4$ μ model for V=4⁴
 - the thimbles associated with the classical vacua
 - the residual phase factors reweighted successfully
 - known results of the number density reproduced (cf. CL, dual v.)
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 - contributions of other thimbles, ex. thimble 2-(a), ...
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- Numerical cost per traj.: but, actually $O(V \times n_{Lefs} \times n_{step})$ solving flow eqs. (all tangent vectors): $O(X^2 \times n_{Lefs})$ computing X^1 , det V (residual sign factors): $O(V^3)$
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 possible applications to QCD μ cf. D. Sexty, arXiv:1307.7748

- We have formulated a HMC algorithm which is applicable to lattice models defined on Lefschetz thimbles
- We have tested the algorithm in the $\lambda \phi^4$ μ model for V=4 4
 - the thimbles associated with the classical vacua
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 - contributions of other thimbles, ex. thimble 2-(a), ...
 - Need the study of the residual sign problem on larger lattices
- Numerical cost per traj.: but, actually $O(V \times n_{Lefs} \times n_{step})$ solving flow eqs. (all tangent vectors): $O(X^2 \times n_{Lefs}) \times CG \times V^2(?)$ computing X^1 , det V (residual sign factors): $O(V^3)$
- Dynamical fermions: psuedo fermions can be implemented possible applications to QCD µ cf. D. Sexty, arXiv:1307.7748



$$\Delta \tau \, w^n - \frac{1}{2} \Delta \tau^2 \, \bar{\partial} \bar{S}[\bar{z}^n]$$

$$-\frac{1}{2}\Delta\tau^2 iV_z^{\alpha}[e^{(n)}, t'^{(n)}] \lambda_{[r]}^{\alpha}$$

$$V_z^{\alpha}[e^{(n)}, t'^{(n)}]$$

 $z[e^{(n)}, t'^{(n)}]$

 $iV_z^{\alpha}[e^{(n)},t'^{(n)}]$

thimble \mathcal{J}_{σ}

$$z[e_{(k)}, t'_{(k)}]$$
 $z[e^{(n+1)}, t'^{(n+1)}]$

the sequences
$$(e_{(k)}^{\alpha}, t_{(k)}')$$
 $(k = 0, 1, \dots)$ with $(e_{(0)}^{\alpha}, t_{(0)}') = (e^{\alpha(n)}, t'^{(n)})$

$$\Delta e^{\alpha}_{(k)} = e^{\alpha}_{(k+1)} - e^{\alpha}_{(k)}, \qquad \sum_{\alpha=1} \Delta e^{\alpha}_{(k)} e^{\alpha(n)} = 0.$$

$$\Delta t'_{(k)} = t'_{(k+1)} - t'_{(k)},$$

$$\Delta z_{(k)}[e^{(n)}, t'^{(n)}] \equiv z_i[e^{(n)}, t'^{(n)}] + \Delta \tau \, w_i^n - \frac{1}{2} \Delta \tau^2 \, \bar{\partial}_i \bar{S}[\bar{z}^n] - z_i[e_{(k)}, t'_{(k)}]$$

$$\Delta z_{(k)}[e^{(n)}, t'^{(n)}]_{\parallel} = V_z^{\alpha}[e^{(n)}, t'^{(n)}] \left(\Delta e^{\alpha}_{(k)} + e^{\alpha(n)} \kappa^{\alpha} \Delta t'_{(k)}\right)$$

$$\Delta z_{(k)}[e^{(n)}, t'^{(n)}]_{\perp} = iV_z^{\alpha}[e^{(n)}, t'^{(n)}] \left(\frac{1}{2} \Delta \tau^2 \lambda_{[r](k)}^{\alpha}\right)$$

$$\left\| V_z^{\alpha}[e^{(n)}, t'^{(n)}] \left(\Delta e^{\alpha}_{(k)} + e^{\alpha(n)} \kappa^{\alpha} \Delta t'_{(k)} \right) \right\|^2 \le n \epsilon'^2$$