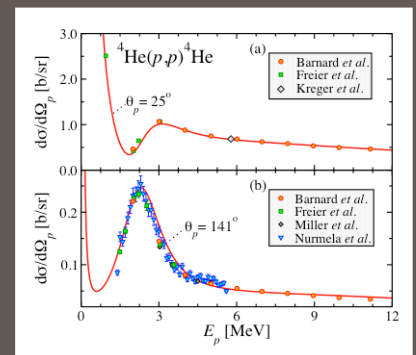
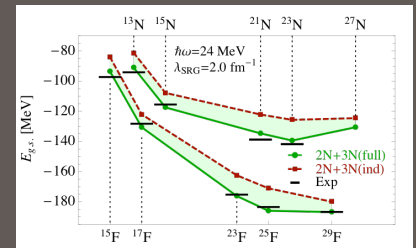


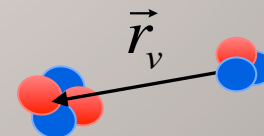
Ab initio calculations in nuclear physics

International Conference on Science and Technology for FAIR in Europe
 Worms, Germany
 October 13-17, 2014

Petr Navratil | TRIUMF



- What is meant by *ab initio* in nuclear physics
- *Ab initio* nuclear structure and reaction approaches
 - Exact few-body calculations ($A=3,4$)
 - Quantum Monte Carlo ($A \leq 12$)
 - Nuclear Lattice EFT ($A=4,8,12,16, 20, 24, 28$)
 - Coupled Cluster Method ($A \leq 132$, magic, semi-magic)
 - In-medium Similarity Renormalization Group ($A \leq 90$, open shells)
 - Self-Consistent Green's Function Method ($A \leq 78$, open shells)
- No-core shell model ($A \leq 26$, hypernuclei)
- Including the continuum with the resonating group method
 - NCSM/RGM
 - NCSM with continuum
- Outlook



What is meant by *ab initio* in nuclear physics?

- **First principles for Nuclear Physics:**

- QCD**

- Non-perturbative at low energies
 - Lattice QCD in the future

- **Degrees of freedom: NUCLEONS**

- Nuclei made of nucleons
 - Interacting by nucleon-nucleon and three-nucleon potentials

- *Ab initio*
 - ✧ All nucleons are active
 - ✧ Exact Pauli principle
 - ✧ Realistic inter-nucleon interactions
 - ✧ Accurate description of NN (and 3N) data
 - ✧ Controllable approximations

Chiral Effective Field Theory

- **First principles for Nuclear Physics:**

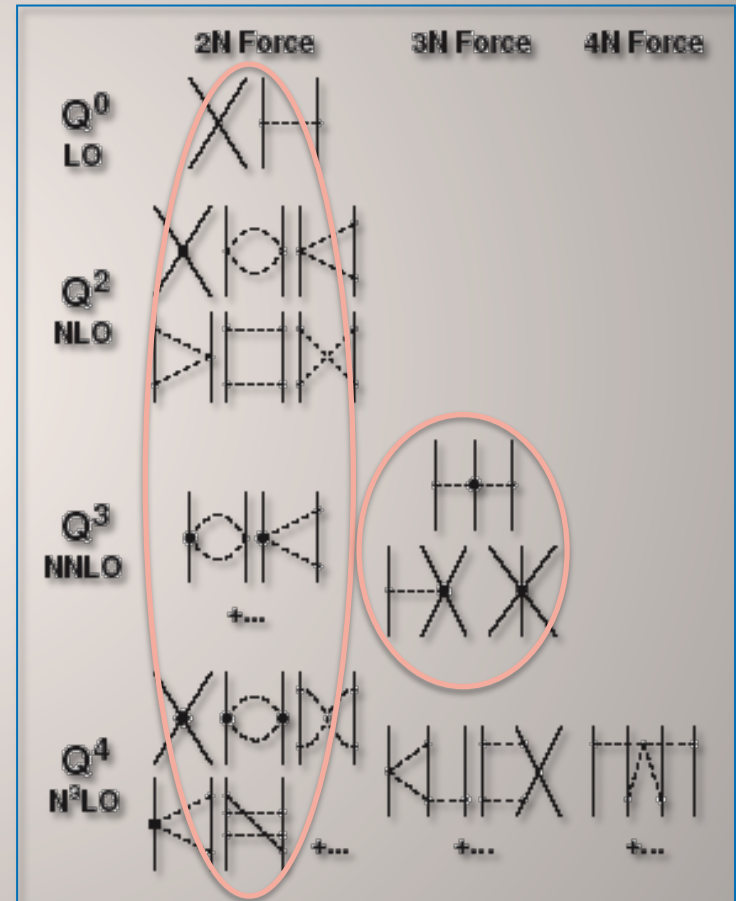
QCD

- Non-perturbative at low energies
- Lattice QCD in the future

- ***For now a good place to start:***

- **Inter-nucleon forces from chiral effective field theory**

- Based on the symmetries of QCD
 - Chiral symmetry of QCD ($m_u \approx m_d \approx 0$), spontaneously broken with pion as the Goldstone boson
 - Degrees of freedom: nucleons + pions
- Systematic low-momentum expansion to a given order (Q/Λ_χ)
- Hierarchy
- Consistency
- Low energy constants (LEC)
 - Fitted to data
 - Can be calculated by lattice QCD



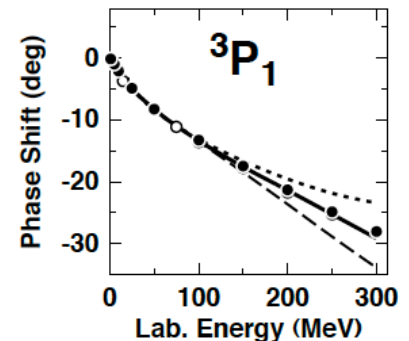
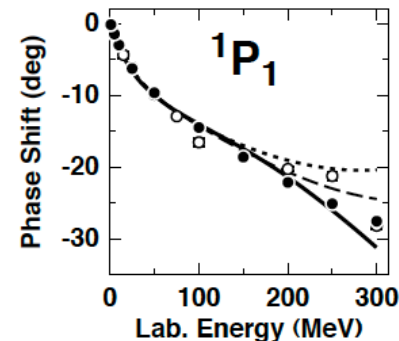
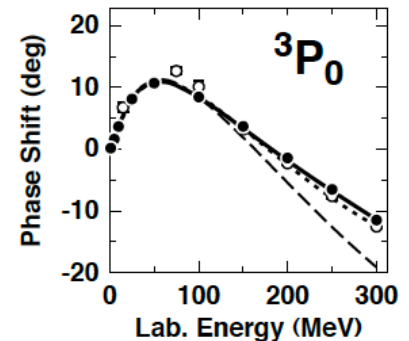
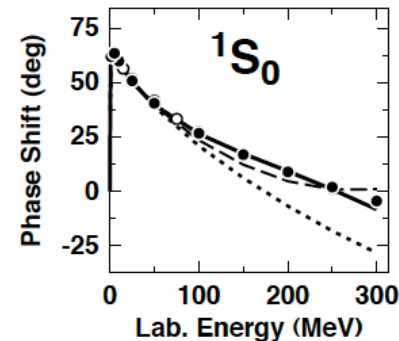
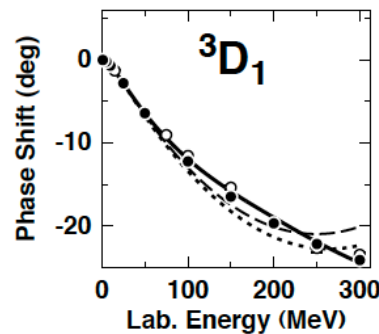
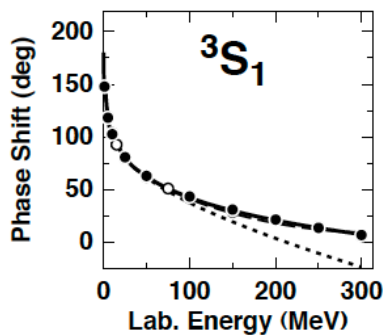
$\Lambda_\chi \sim 1 \text{ GeV}$:
Chiral symmetry breaking scale

The NN interaction from chiral EFT

PHYSICAL REVIEW C **68**, 041001(R) (2003)

Accurate charge-dependent nucleon-nucleon potential at fourth order of chiral perturbation theory

D. R. Entem^{1,2,*} and R. Machleidt^{1,†}



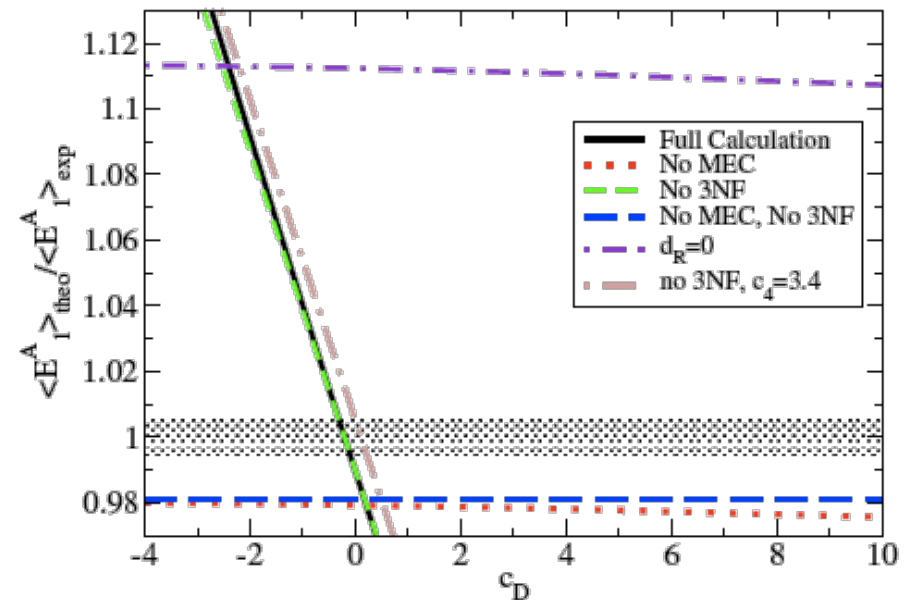
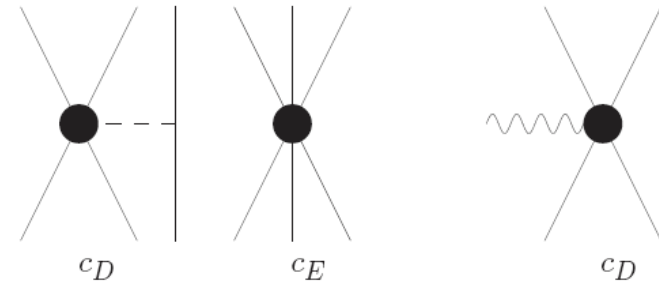
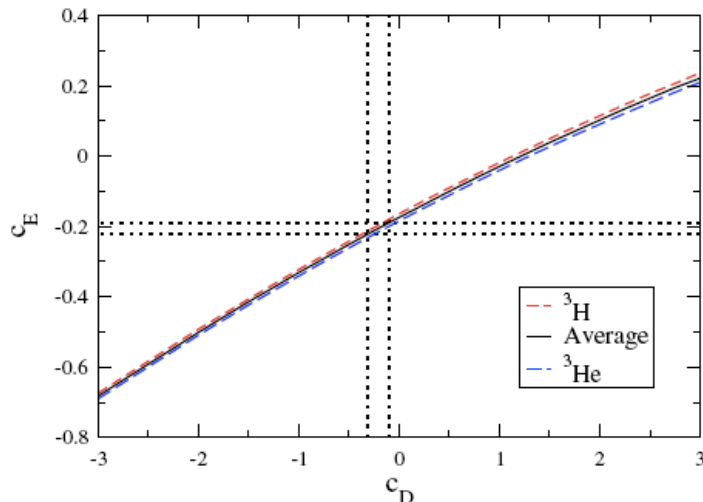
- 24 LECs fitted to the np scattering data and the deuteron properties
 - Including c_i LECs ($i=1-4$) from pion-nucleon Lagrangian

Determination of NNN LECs c_D and c_E from the triton binding energy and the half life

- **Chiral EFT:** c_D also in the two-nucleon contact vertex with an external probe
- Calculate $\langle E_1^A \rangle = |\langle {}^3\text{He} || E_1^A || {}^3\text{H} \rangle|$
 - Leading order GT
 - N²LO: one-pion exchange plus contact

- **A=3 binding energy constraint:**

$$c_D = -0.2 \pm 0.1 \quad c_E = -0.205 \pm 0.015$$



PRL 103, 102502 (2009) PHYSICAL REVIEW LETTERS week ending 4 SEPTEMBER 2009

Three-Nucleon Low-Energy Constants from the Consistency of Interactions and Currents in Chiral Effective Field Theory

Doron Gazit

Institute for Nuclear Theory, University of Washington, Box 351550, Seattle, Washington 98195, USA

Sofia Quaglioni and Petr Navrátil

Lawrence Livermore National Laboratory, P.O. Box 808, L-414, Livermore, California 94551, USA

Exact few-body calculations ($A=3,4$)

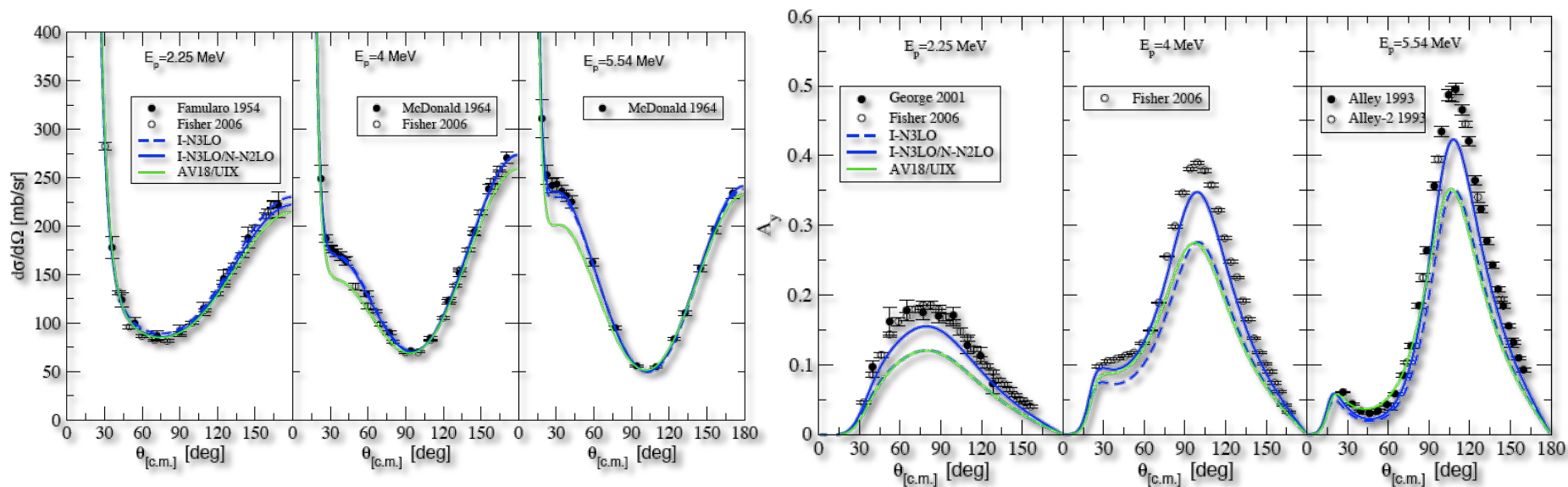
Proton- ^3He elastic scattering with $\chi\text{EFT NN+NNN}$

- **Hyperspherical-harmonics variational calculations**

- M. Viviani, L. Girlanda, A. Kievski, L. E. Marcucci, and S. Rosati, EPJ Web Conf. **3** (2010) 05011; Few Body Syst. **54** (2013) 885

- **A_y puzzle (almost) resolved with the chiral N^3LO NN plus local chiral N^2LO NNN**

- *used with the NCSM and other methods*



Quantum Monte Carlo

Variational Monte Carlo (VMC): construct Ψ_V that

- Are fully **antisymmetric** and **translationally invariant**
- Have **cluster structure** and correct asymptotic form
- Contain non-commuting 2- & 3-body **operator correlations** from v_{ij} & V_{ijk}
- Are orthogonal for multiple J^π states
- Minimize $E_V = \langle \Psi_V | H | \Psi_V \rangle \geq E$ integrating by Metropolis Monte Carlo

These are $\sim 2^A \binom{A}{Z}$ component (270,336 for ^{12}C) spin-isospin vectors in $3A$ dimensions

Green's function Monte Carlo (GFMC): project out the exact eigenfunction

- $\Psi(\tau) = \exp[-(H - E_0)\tau]\Psi_V = \sum_n \exp[-(E_n - E_0)\tau]a_n\Psi_n \Rightarrow \Psi_0$ at large τ
- Propagation done stochastically in small time slices $\Delta\tau$
- Exact $\langle H \rangle$ for local potentials; mixed estimates for other $\langle O \rangle$
- **Constrained-path propagation** controls fermion sign problem for $A \geq 8$
- Multiple excited states for same J^π stay orthogonal

Many tests demonstrate 1–2% accuracy for realistic $\langle H \rangle$

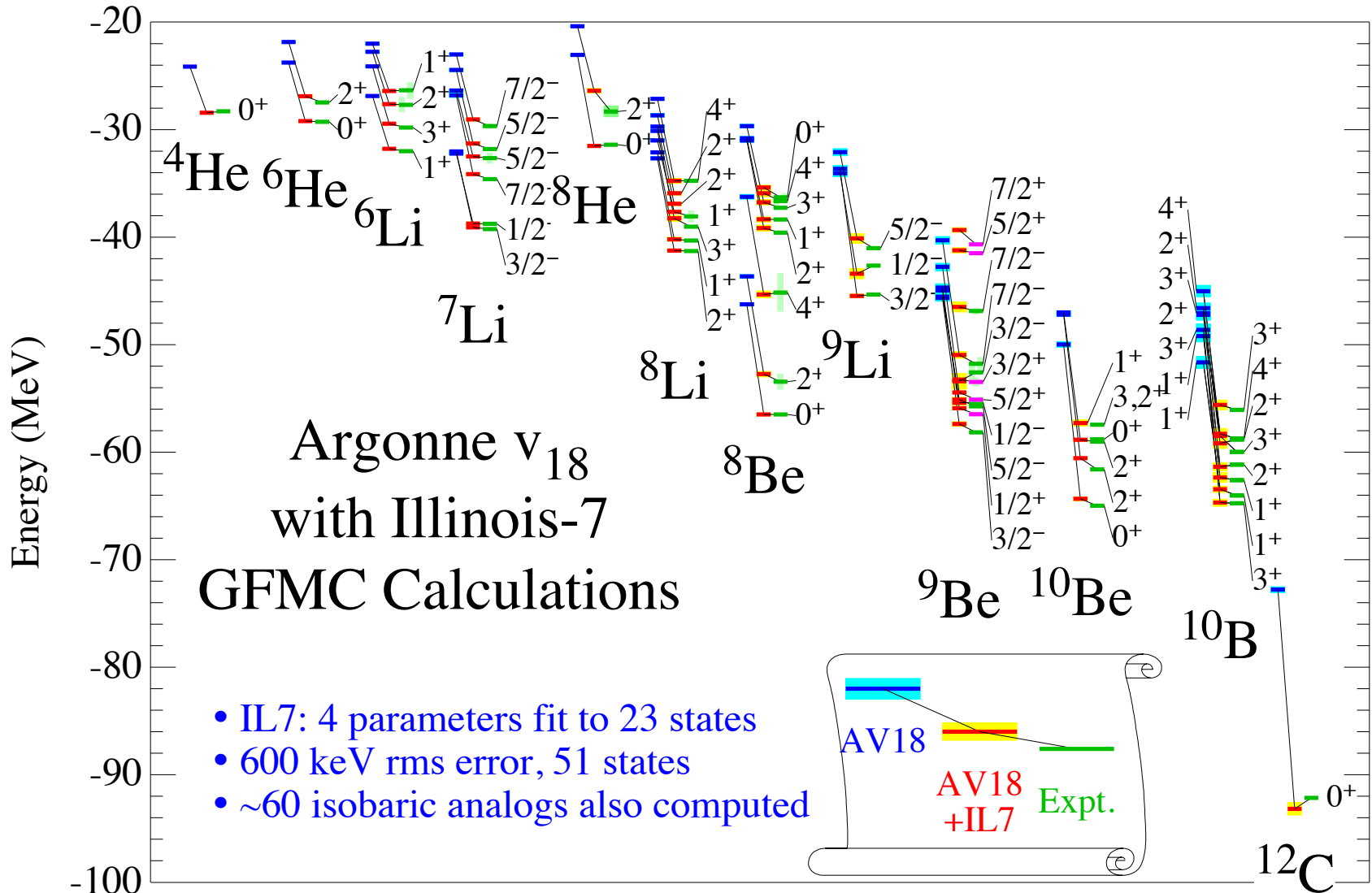
Wiringa, Pieper, Carlson, & Pandharipande, PRC **62**, 014001 (2000)

Pieper, Varga, & Wiringa, PRC **66**, 044310 (2002)

Pieper, Wiringa, & Carlson, PRC **70**, 054325 (2004)

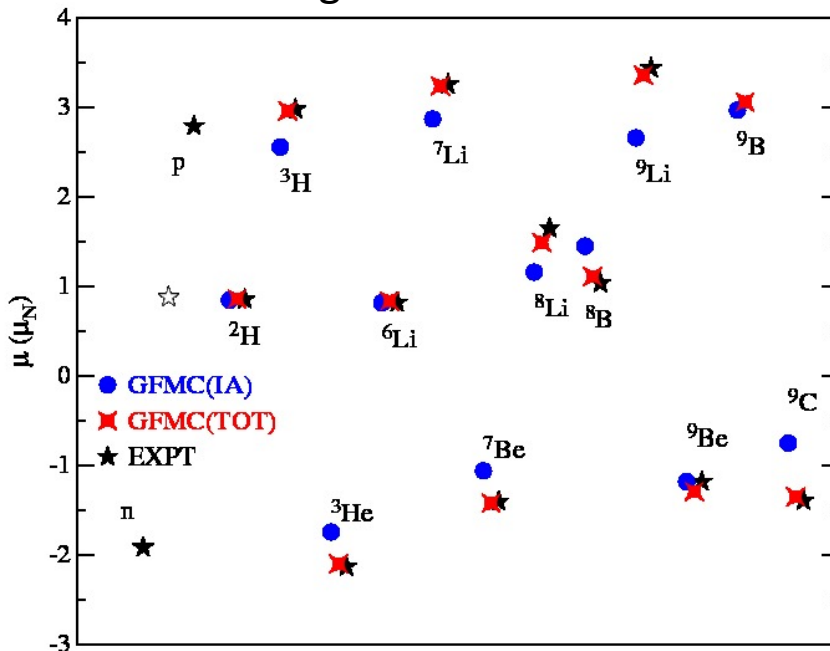
Pieper, NPA **751**, 516c (2005)

Quantum Monte Carlo: Eigenenergies of light nuclei



Quantum Monte Carlo: Magnetic moments and transitions light nuclei

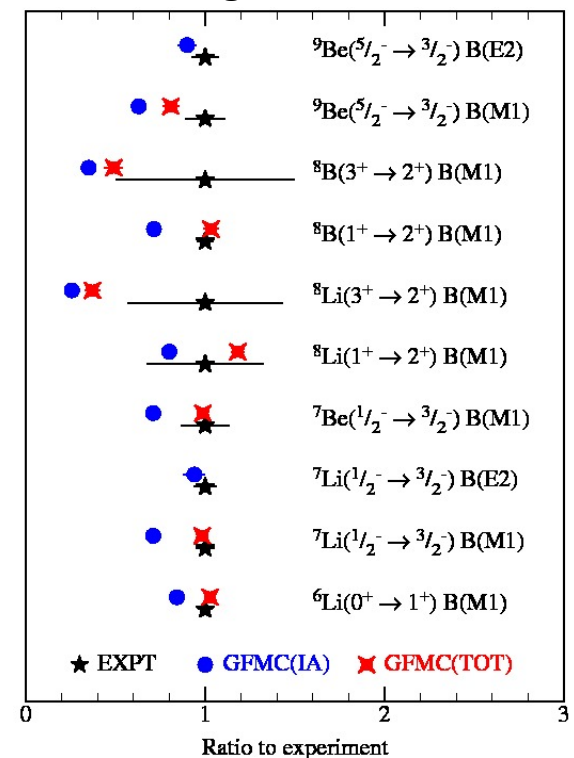
Magnetic Moments



Green's Function Monte Carlo (GFMC) calculations of light nuclei give accurate energies but a **lowest-order theory of one-body currents (blue)** disagrees with experiment (black).

Including **two-nucleon currents based on effective field theory (red)** improves all predictions!

Electromagnetic Transitions



PHYSICAL REVIEW C **87**, 035503 (2013)



Quantum Monte Carlo calculations of electromagnetic moments and transitions in $A \leq 9$ nuclei with meson-exchange currents derived from chiral effective field theory

S. Pastore,^{1,*} Steven C. Pieper,^{1,†} R. Schiavilla,^{2,3,‡} and R. B. Wiringa^{1,§}

Nuclear Lattice Effective Field Theory Calculations

E. Epelbaum, H. Krebs, T. Lahde, D. Lee, U.-G. Meissner

Discretized version of
chiral EFT for nuclear
dynamics

$$\left[\left(\sum_{i=1}^A \frac{-\vec{\nabla}_i^2}{2m_N} + \mathcal{O}(m_N^{-3}) \right) + \underbrace{V_{2N} + V_{3N} + V_{4N} + \dots}_{\text{derivable within ChPT}} \right] |\Psi\rangle = E|\Psi\rangle$$

Borasoy, E.E., Krebs, Lee, Meißner, Eur. Phys. J. A31 (07) 105,

Eur. Phys. J. A34 (07) 185,

Eur. Phys. J. A35 (08) 343,

Eur. Phys. J. A35 (08) 357,

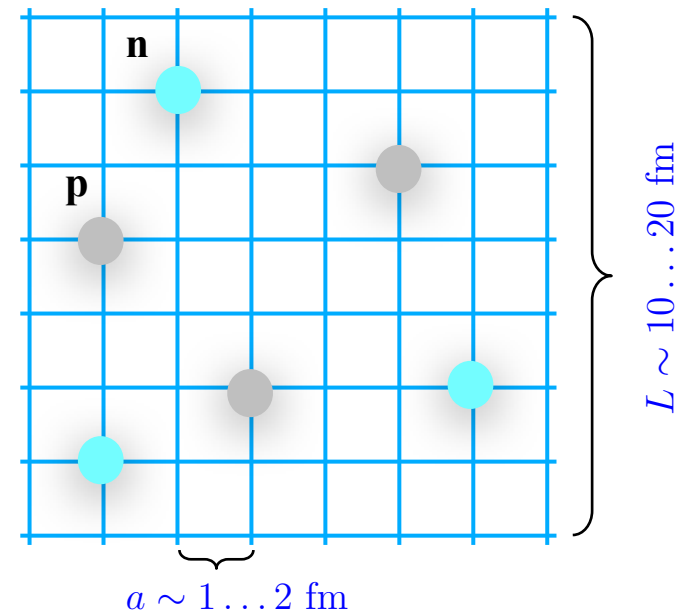
E.E., Krebs, Lee, Meißner, Eur. Phys. J. A40 (09) 199,

Eur. Phys. J. A41 (09) 125,

Phys. Rev. Lett 104 (10) 142501,

Eur. Phys. J. 45 (10) 335,

Phys. Rev. Lett. 106 (11) 192501





Physics Letters B 732 (2014) 110–115

Contents lists available at ScienceDirect

Physics Letters B

www.elsevier.com/locate/physletb

Lattice effective field theory for medium-mass nuclei

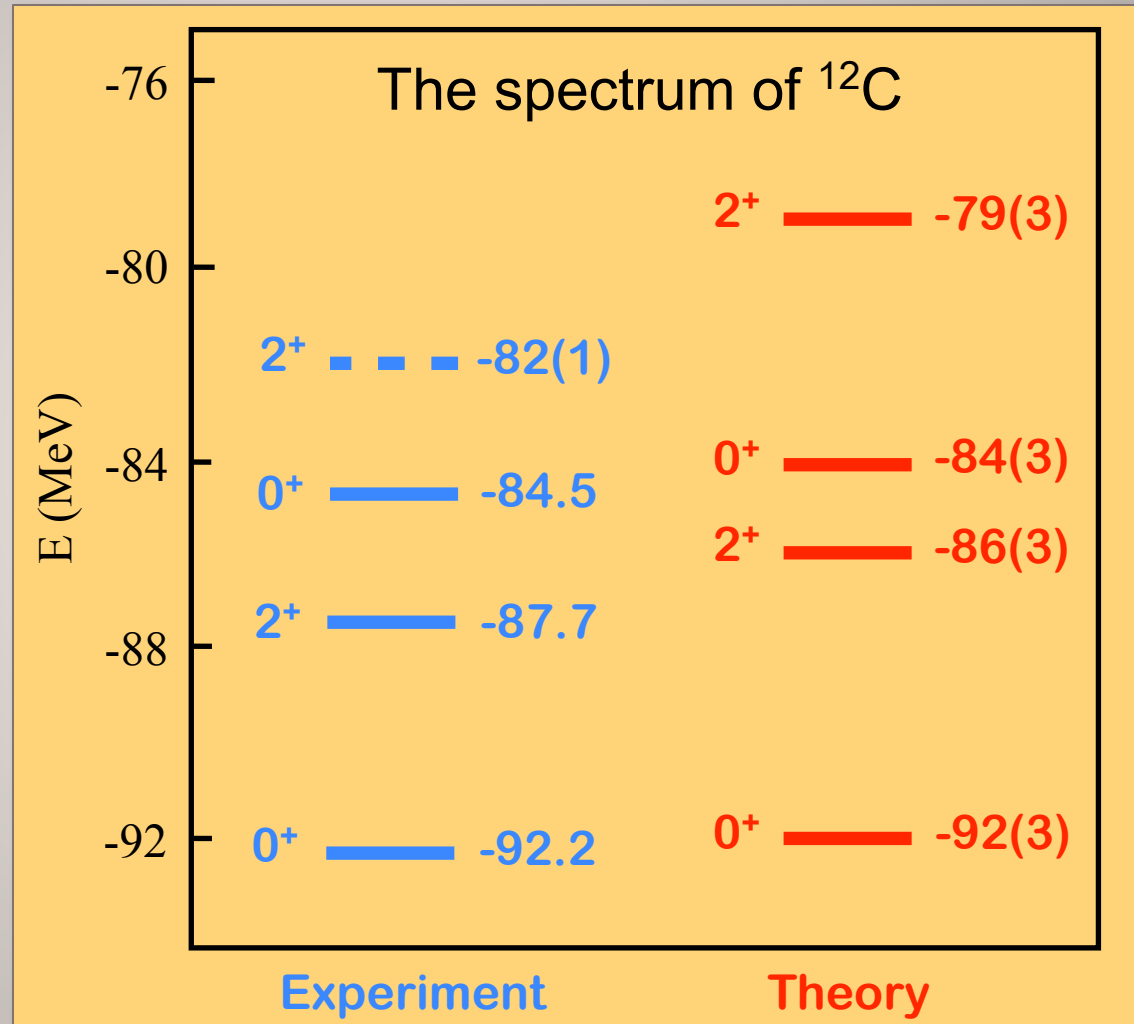
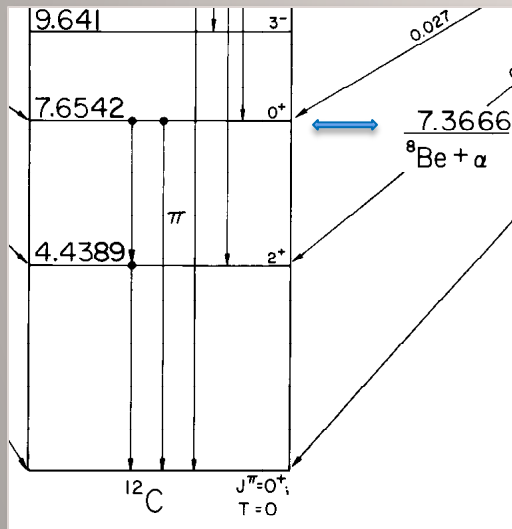
Timo A. Lähde^{a,*}, Evgeny Epelbaum^b, Hermann Krebs^b, Dean Lee^c, Ulf-G. Meißner^{a,d,e}, Gautam Rupak^f

Nuclear Lattice Effective Field Theory Calculations

E. Epelbaum, H. Krebs, T. Lahde, D. Lee, U.-G. Meissner

The Hoyle state

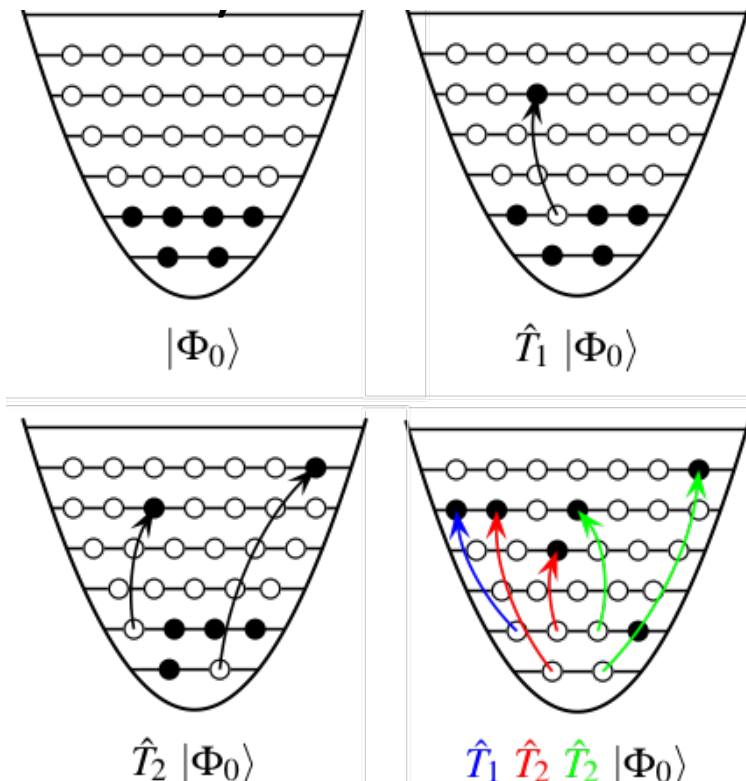
^{12}C production
in the Universe



Coupled-Cluster Method

- **exponential Ansatz** for wave operator $|\Psi\rangle = e^{\hat{T}} |\Phi_0\rangle$

- **CCSD**: truncate \hat{T} at the **2p2h excitation** level, $\hat{T} = \hat{T}_1 + \hat{T}_2$

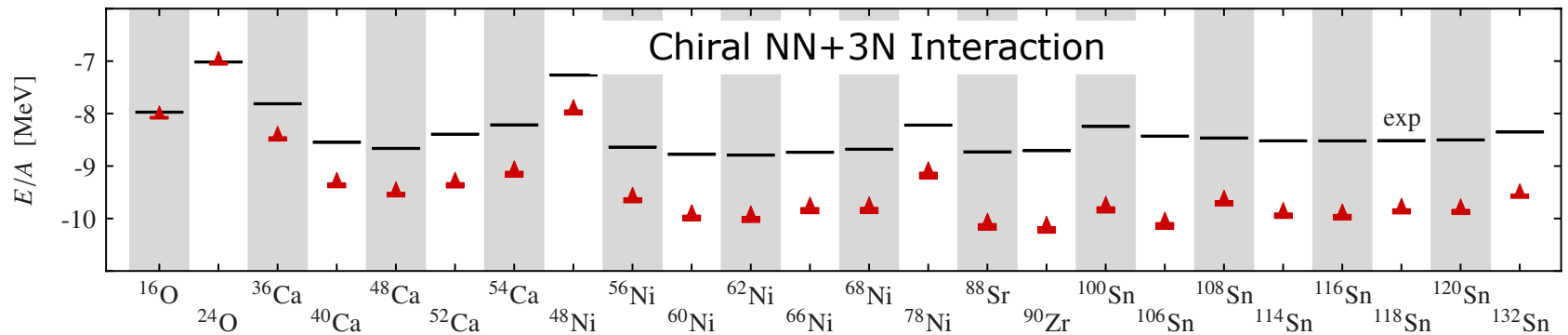


$$\hat{T}_n = \frac{1}{(n!)^2} \sum_{\substack{ijk\dots \\ abc\dots}} t_{ijk\dots} \{ \hat{a}_a^\dagger \hat{a}_b^\dagger \hat{a}_c^\dagger \dots \hat{a}_k \hat{a}_j \hat{a}_i \}$$

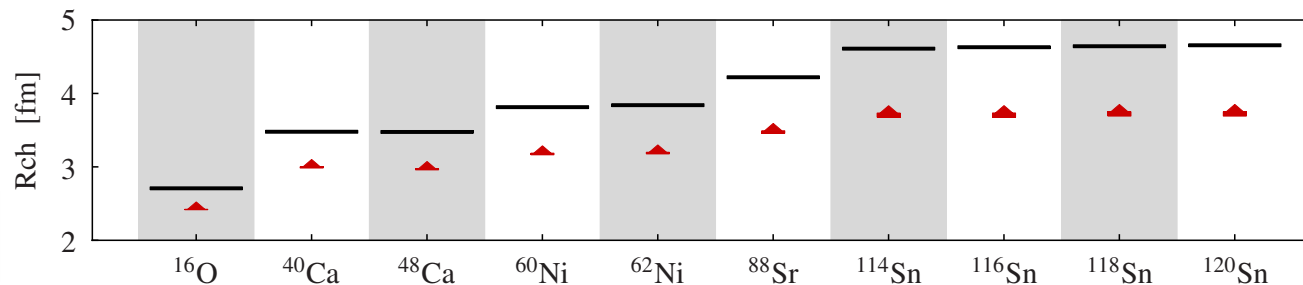
- effects of T_3 clusters **included approximately** in ground-state calculations via Λ CCSD(T) or CR-CC(2,3) method

State-of-the-art: Λ -CCSD(T) with 3N interaction

Coupled-Cluster calculations for heavy nuclei with chiral interactions



- **current** chiral Hamiltonians capable of describing the **experimental trend** of binding energies
- systematic overbinding indicates that there are still **deficiencies**
 - ➔ **consistent 3N** interaction at N^3LO , and **4N** interactions
- **charge radii** are considerably **too small**

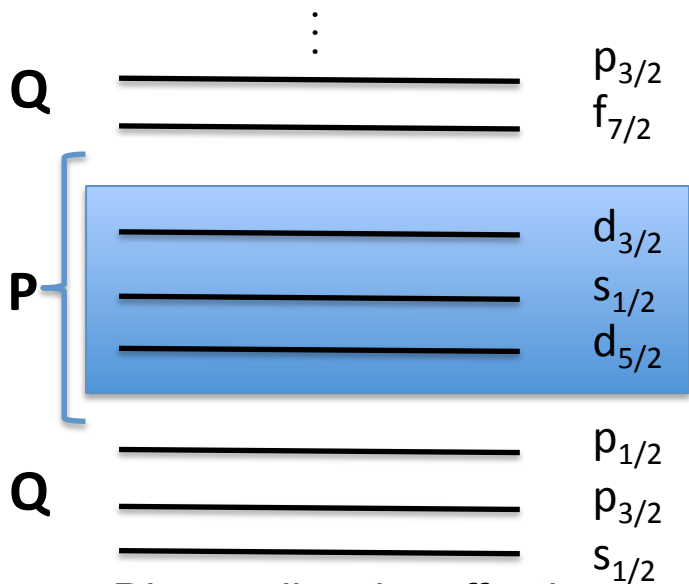


Coupled-cluster effective interactions (CCEI) for the shell model

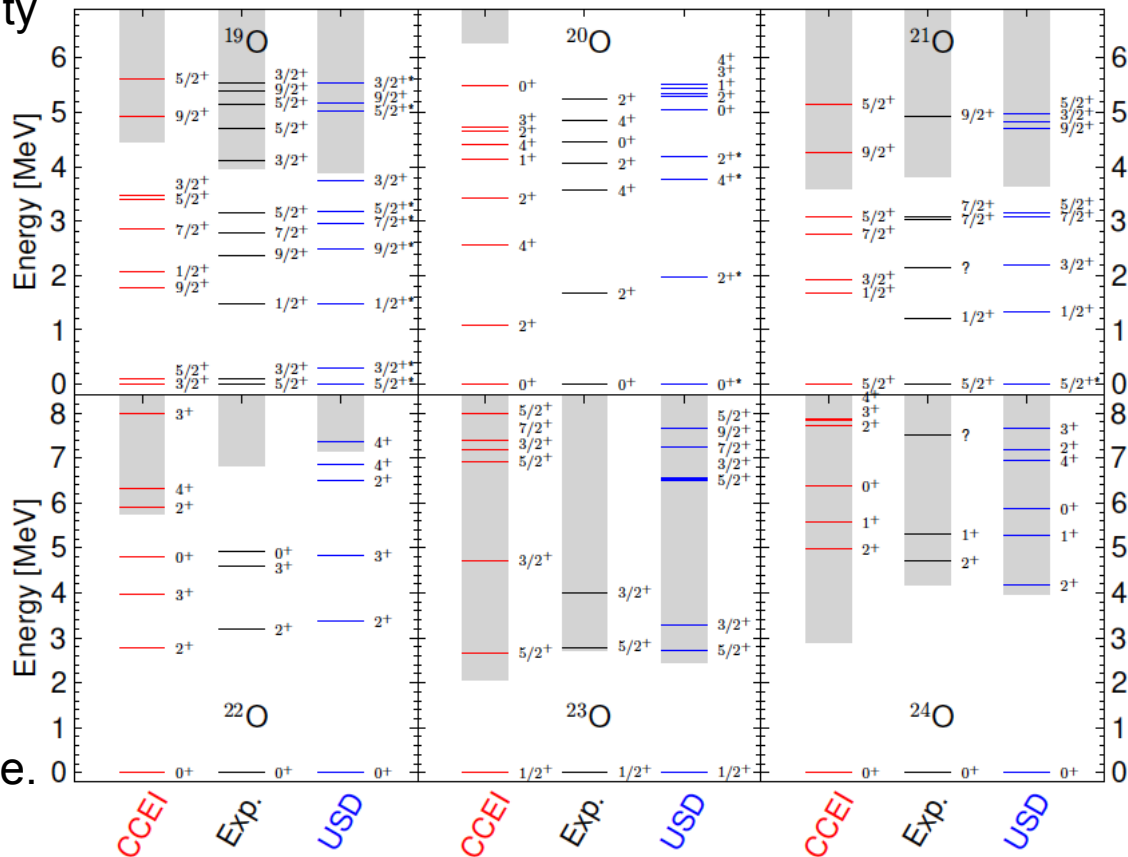
G. R. Jansen, J. Engel, G. Hagen, P. Navratil, A. Signoracci, Phys. Rev. Lett. **113**, 142502 (2014).

- Start from chiral NN(N3LO_{EM}) + 3NF(N2LO) interactions
- Solve for A+1 and A+2 using CC. Project A+1 and A+2 CC wave functions onto the s-d model space using Lee-Suzuki similarity transformation.

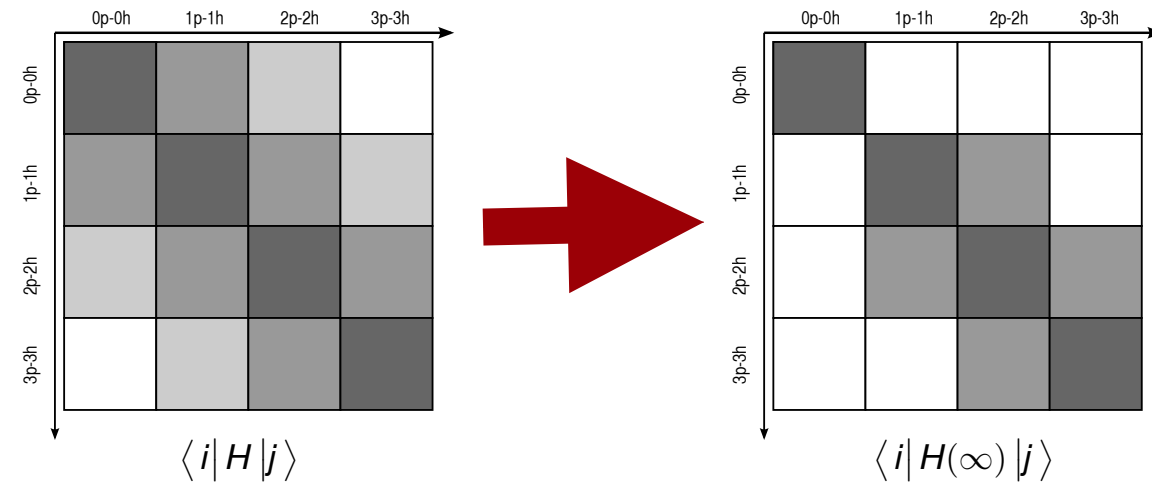
Spectra of oxygen isotopes computed with coupled-cluster effective interaction (CCEI), and compared to experimental data and the phenomenological USD shell model interaction.



- Diagonalize the effective hamiltonian in the valence space.



In-medium SRG approach: Application to Oxygen isotopes



aim: decouple reference state (0p-0h) from excitations

$$\frac{d}{ds} H(s) = [\eta(s), H(s)], \quad \eta(s) = \frac{dU(s)}{ds} U^\dagger(s) = -\eta^\dagger(s)$$

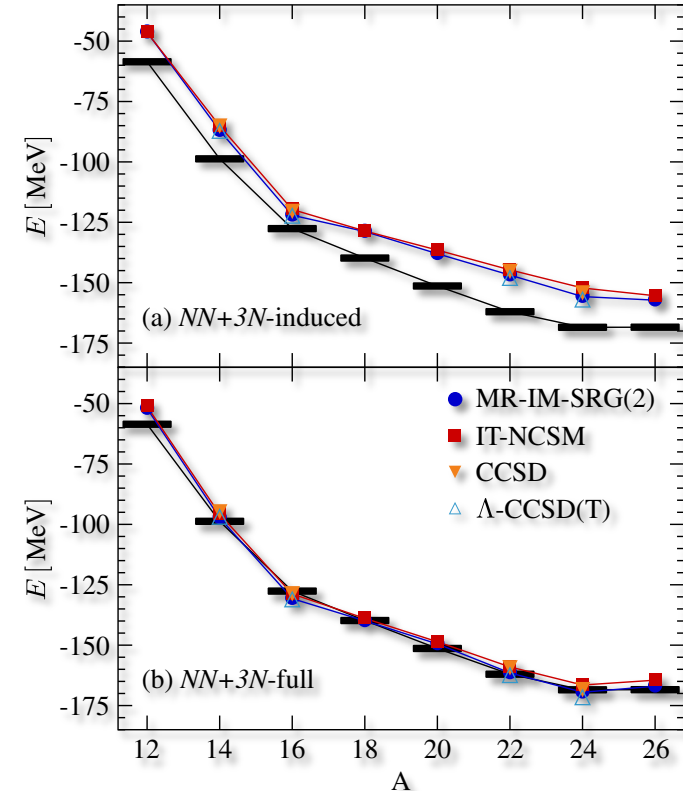
- Wegner

$$\eta^I = [H^d, H^{od}]$$

- White (J. Chem. Phys. 117, 7472)

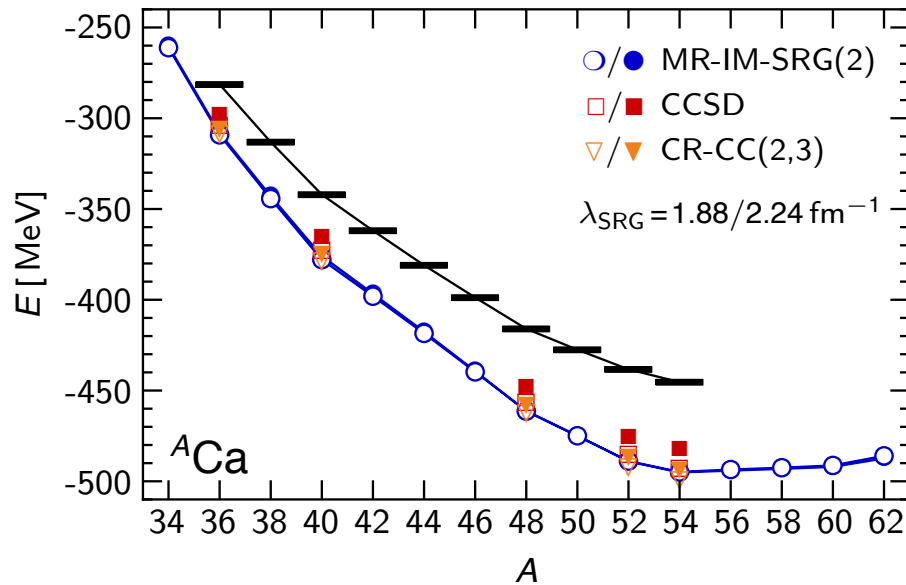
$$\eta^{II} = \sum_{ph} \frac{f_h^p}{E_p - E_h} : A_h^p : + \sum_{pp'hh'} \frac{\Gamma_{hh'}^{pp'}}{E_{pp'} - E_{hh'}} : A_{hh'}^{pp'} : + \text{H.c.}$$

$E_p - E_h, E_{pp'} - E_{hh'}$: approx. 1p1h, 2p2h excitation energies

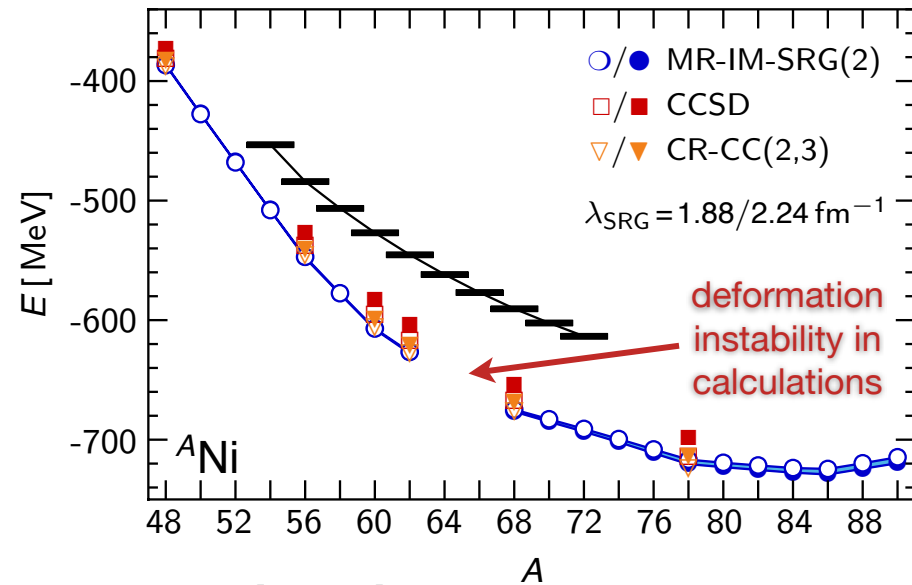


In-medium SRG approach: Application to Ca and Ni isotopes

NN + 3N-full(400)



NN + 3N-full(400)



H. Hergert et al., arXiv:1408.6555 [nucl-th]

- IM-SRG calculations for $A \sim 100$ are routine, tin isotopes in progress
- controlled uncertainties & consistent results for different ab-initio methods
- systematic overbinding due to current chiral Hamiltonians - results for new generation of chiral Hamiltonians soon

Self-Consistent Green's Function Method: Oxygen, Fluorine, Nitrogen isotopes

Magic and semi-magic nuclei

$$g_{\alpha\beta}(\omega) = \sum_n \frac{\langle \Psi_0^A | c_\alpha | \Psi_n^{A+1} \rangle \langle \Psi_n^{A+1} | c_\beta^\dagger | \Psi_0^A \rangle}{\omega - \varepsilon_n^{A+1} + i\eta} + \sum_k \frac{\langle \Psi_0^A | c_\beta^\dagger | \Psi_k^{A-1} \rangle \langle \Psi_k^{A-1} | c_\alpha | \Psi_0^A \rangle}{\omega - \varepsilon_k^{A-1} - i\eta},$$

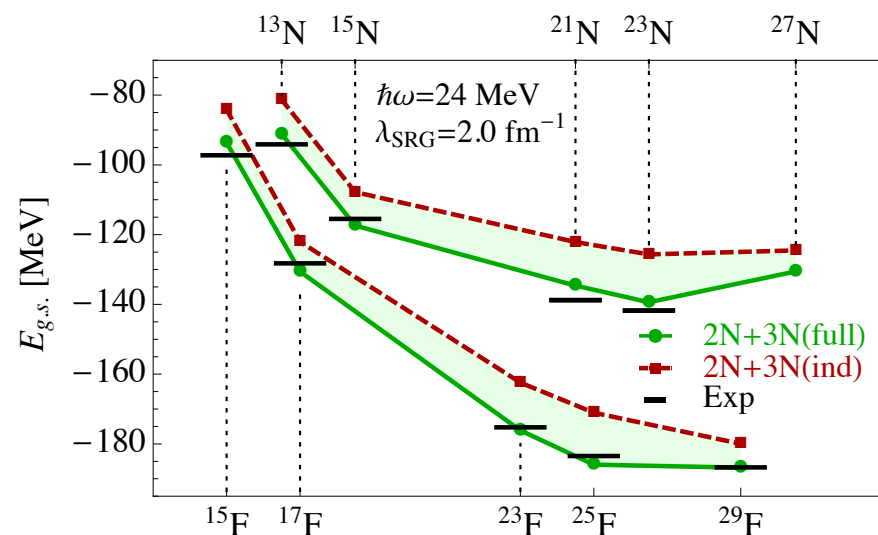
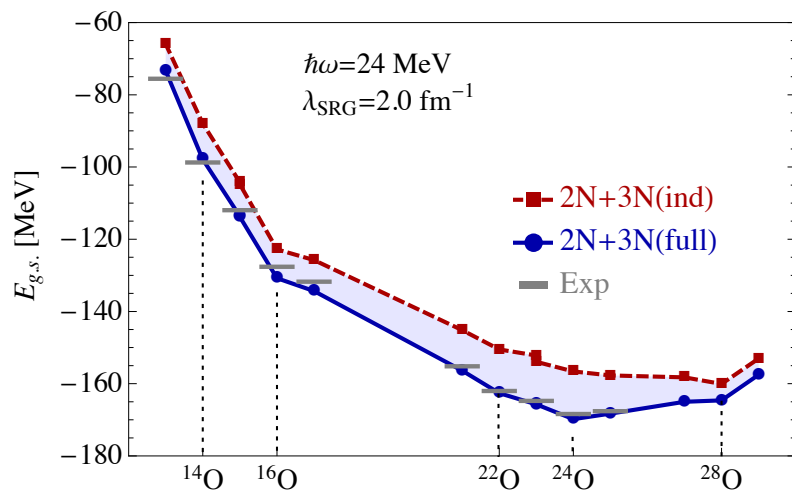
PRL 111, 062501 (2013)

PHYSICAL REVIEW LETTERS

week ending
9 AUGUST 2013

Isotopic Chains Around Oxygen from Evolved Chiral Two- and Three-Nucleon Interactions

A. Cipollone,¹ C. Barbieri,^{1,*} and P. Navrátil²

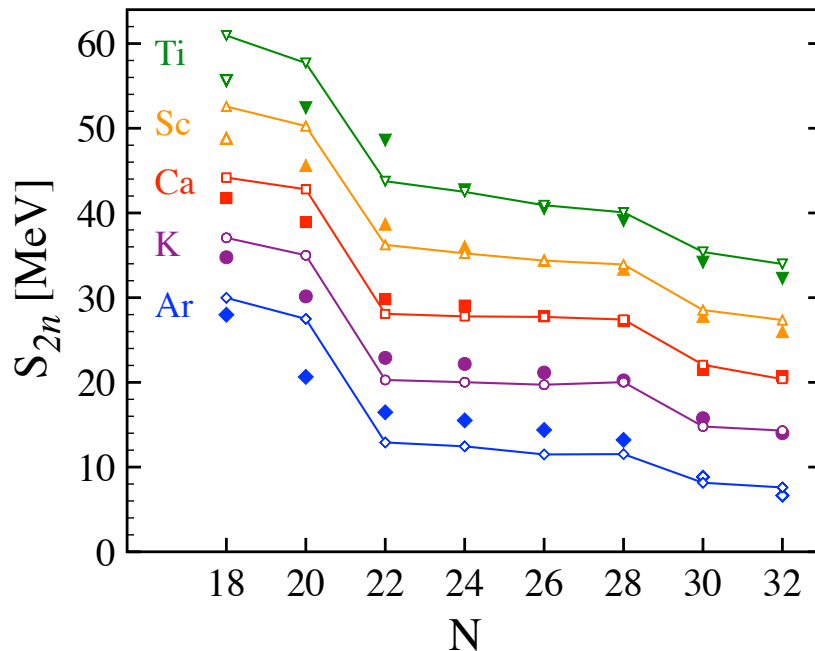


- 3NF crucial for reproducing binding energies and driplines around oxygen
- $d_{3/2}$ raised by genuine 3NF

Green's functions in medium-mass nuclei

Gorkov GF go **beyond standard expansion schemes** and are not limited to doubly closed-shells

- Expansion around a Bogoliubov vacuum
- **From few tens to hundreds** of medium-mass open-shell systems (→ complete chains)



Open shell nuclei

PHYSICAL REVIEW C **89**, 061301(R) (2014)

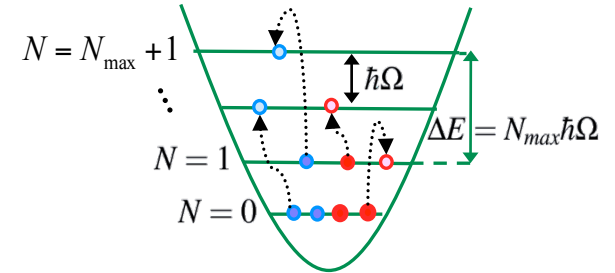
Chiral two- and three-nucleon forces along medium-mass isotope chains

V. Somà,^{1,2,3,*} A. Cipollone,⁴ C. Barbieri,^{4,†} P. Navrátil,⁵ and T. Duguet^{3,6,‡}

- Systematic overbinding of medium-mass nuclei (in **agreement with other ab initio methods**)
- initial (full) 3NF are **necessary** to reproduce relative trends
- Relative energies (S_{2n}) well reproduced

No-core shell model

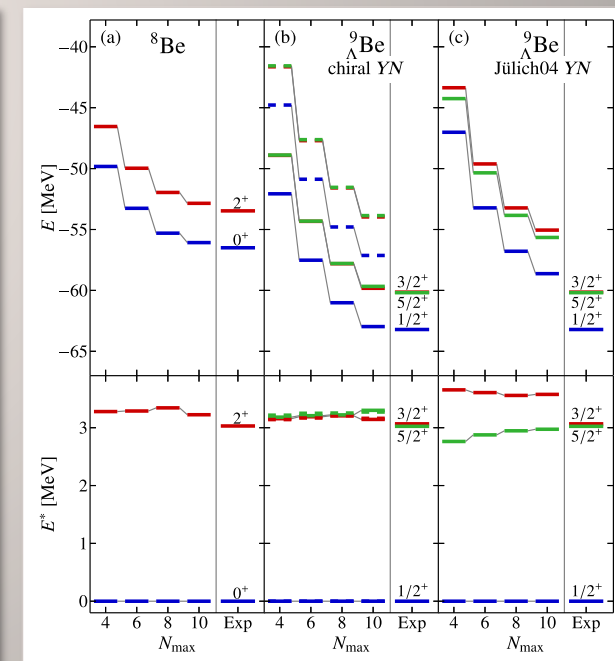
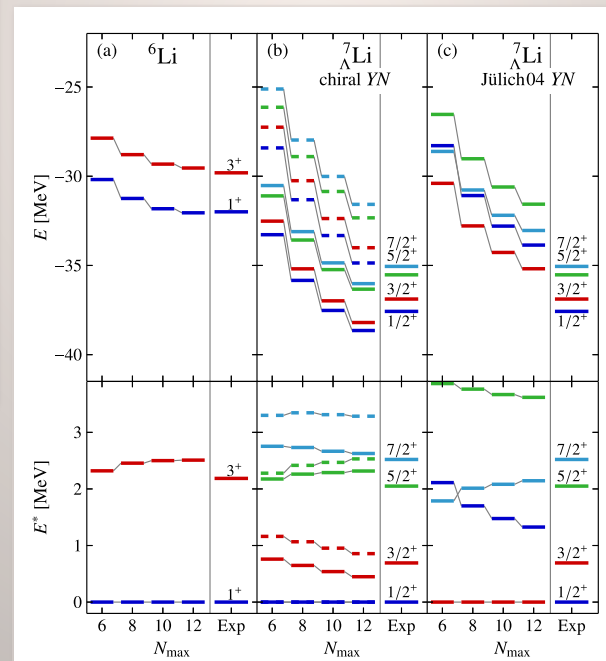
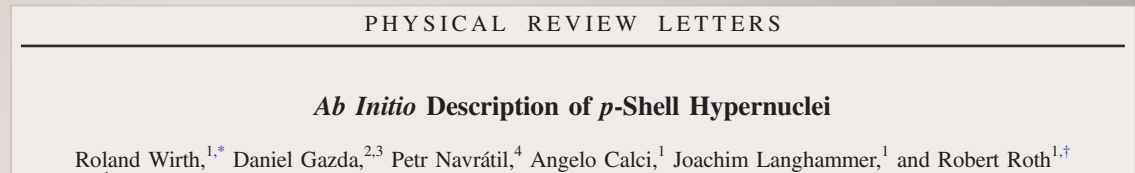
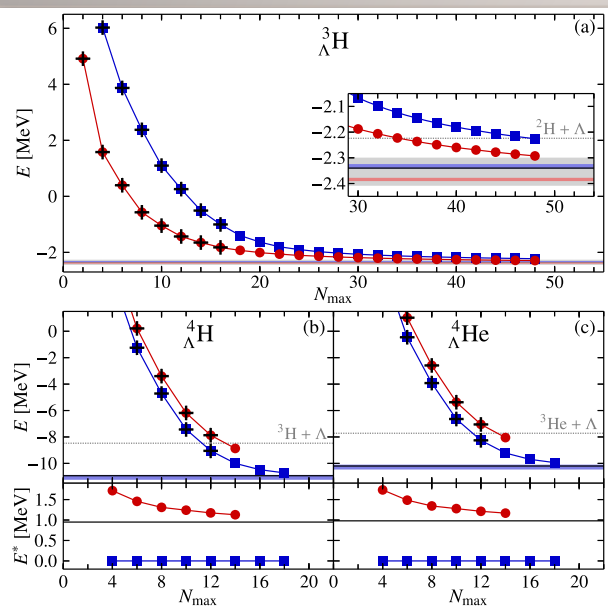
- No-core shell model (NCSM)
 - A -nucleon wave function expansion in the harmonic-oscillator (HO) basis
 - short- and medium range correlations
 - Bound-states, narrow resonances



$$\Psi^A = \sum_{N=0}^{N_{\max}} \sum_i c_{Ni} \Phi_{Ni}^A$$

NCSM calculations for light nuclei and *hypernuclei*

- Flexible approach capable performing exact calculations for few-nucleon systems and accurate calculations for nuclei with $A \leq 24$ & **hypernuclei**



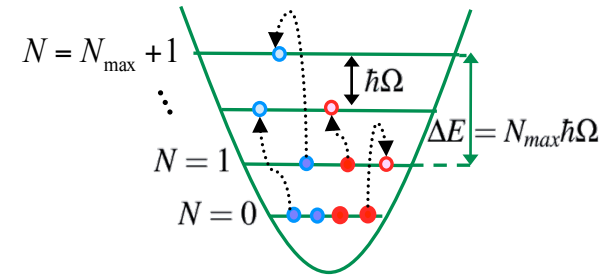
Testing chiral LO NY potentials with Λ - Σ mixing included

...outperform the Julich '04 YN potential

No-core shell model with continuum

- No-core shell model (NCSM)

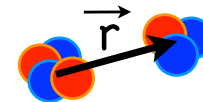
- A -nucleon wave function expansion in the harmonic-oscillator (HO) basis
- short- and medium range correlations
- Bound-states, narrow resonances



$$\Psi^A = \sum_{N=0}^{N_{\max}} \sum_i c_{Ni} \Phi_{Ni}^A$$

- NCSM with Resonating Group Method (NCSM/RGM)

- cluster expansion
- proper asymptotic behavior
- long-range correlations



$$\Psi^A = \sum_{\nu} \int d\vec{r} \varphi_{\nu}(\vec{r}) \mathcal{A}_{\nu} \Phi_{1\nu}^{(A-a)} \Phi_{2\nu}^{(a)} \delta(\vec{r} - \vec{r}_{A-a,a})$$

- NCSM with continuum (NCSMC)

- unified description of bound and unbound states



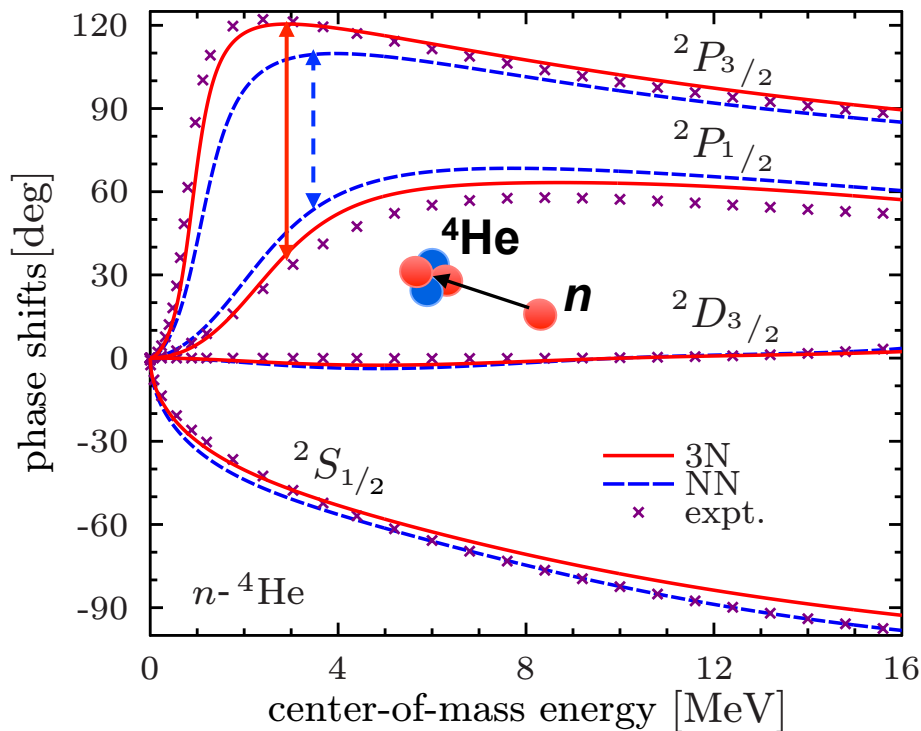
Coupled NCSMC equations

$$\begin{array}{c}
 \begin{array}{c}
 \boxed{E_{\lambda}^{NCSM} \delta_{\lambda\lambda'}} \\
 \downarrow \text{blue} \\
 \left(\begin{array}{cc}
 H_{NCSM} & h \\
 h & H_{RGM}
 \end{array} \right) \begin{pmatrix} \textcircled{C} \\ \textcircled{\gamma} \end{pmatrix}
 \end{array}
 & = E &
 \begin{array}{c}
 \boxed{\delta_{\lambda\lambda'}} \\
 \downarrow \text{blue} \\
 \left(\begin{array}{cc}
 1_{NCSM} & g \\
 g & N_{RGM}
 \end{array} \right) \begin{pmatrix} \textcircled{C} \\ \textcircled{\gamma} \end{pmatrix}
 \end{array}
 \\
 & &
 \begin{array}{c}
 \downarrow \text{green} \\
 \begin{array}{c}
 \boxed{\langle (A) \left| H \hat{A}_v \right| (a) (A-a) \rangle} \\
 \downarrow \text{green} \\
 \begin{array}{c}
 \boxed{\langle (A-a) (a) \left| \hat{A}_v H \hat{A}_v \right| (a) (A-a) \rangle} \\
 \uparrow \text{red}
 \end{array}
 \end{array}
 & &
 \begin{array}{c}
 \downarrow \text{green} \\
 \begin{array}{c}
 \boxed{\langle (A) \left| \hat{A}_v \right| (a) (A-a) \rangle} \\
 \downarrow \text{green} \\
 \begin{array}{c}
 \boxed{\langle (A-a) (a) \left| \hat{A}_v \hat{A}_v \right| (a) (A-a) \rangle} \\
 \uparrow \text{red}
 \end{array}
 \end{array}
 \end{array}
 \end{array}$$

Scattering matrix (and observables) from matching solutions to known asymptotic with microscopic R -matrix on Lagrange mesh

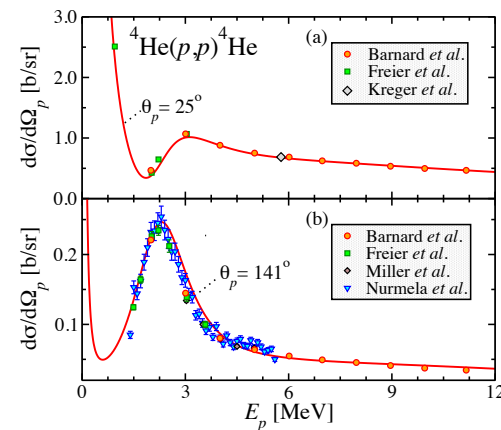
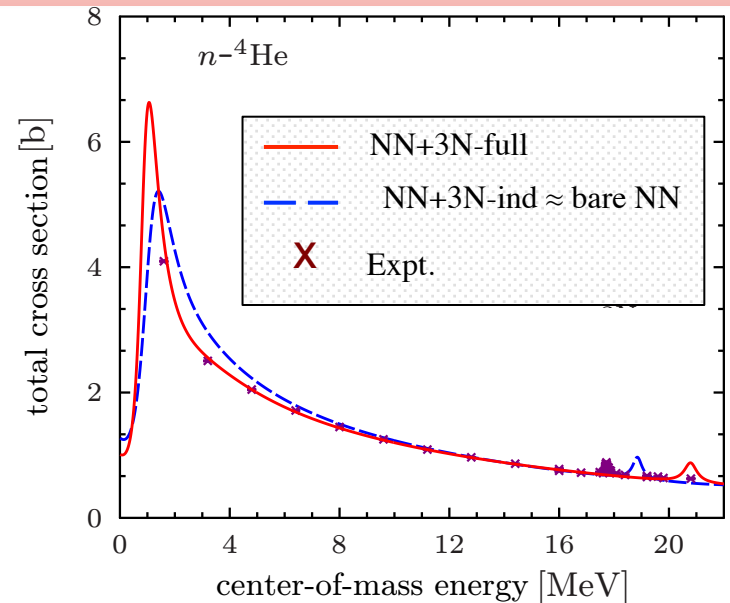
n - ^4He & p - ^4He scattering within NCSMC

n - ^4He scattering phase-shifts for chiral NN and NN+3N potential



3N force enhances $1/2^- \leftrightarrow 3/2^-$ splitting; essential at low energies!

Total n - ^4He cross section with NN and NN+3N potentials



Differential p - ^4He cross section with NN+3N potentials

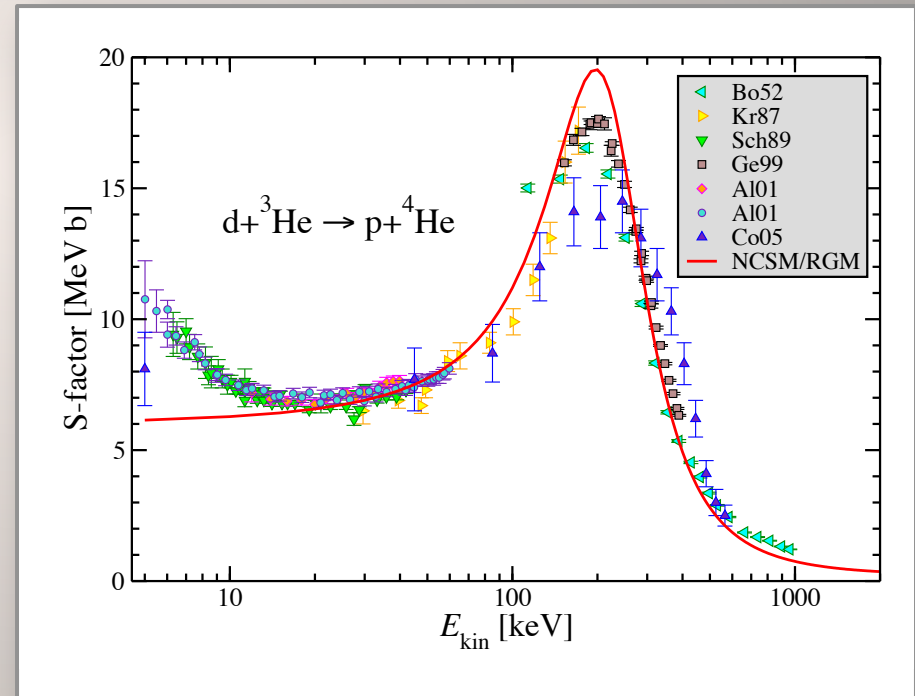
NCSM/RGM calculations of transfer reactions

$$\int dr r^2 \left(\begin{array}{c} \left\langle \begin{array}{c} \mathbf{r}' \\ \alpha \\ n \end{array} \left| \hat{A}_1 (H - E) \hat{A}_1 \right| \begin{array}{c} \mathbf{r} \\ \alpha \\ n \end{array} \right\rangle \\ \left\langle \begin{array}{c} \mathbf{r}' \\ \alpha \\ n \end{array} \left| \hat{A}_1 (H - E) \hat{A}_2 \right| \begin{array}{c} \mathbf{r} \\ \alpha \\ n \end{array} \right\rangle \\ \left\langle \begin{array}{c} \mathbf{r}' \\ d \text{ } ^3\text{H} \end{array} \left| \hat{A}_2 (H - E) \hat{A}_1 \right| \begin{array}{c} \mathbf{r} \\ \alpha \\ n \end{array} \right\rangle \\ \left\langle \begin{array}{c} \mathbf{r}' \\ d \text{ } ^3\text{H} \end{array} \left| \hat{A}_2 (H - E) \hat{A}_2 \right| \begin{array}{c} \mathbf{r} \\ d \text{ } ^3\text{H} \end{array} \right\rangle \end{array} \right) \begin{pmatrix} \frac{g_1(r)}{r} \\ \frac{g_2(r)}{r} \end{pmatrix} = 0$$

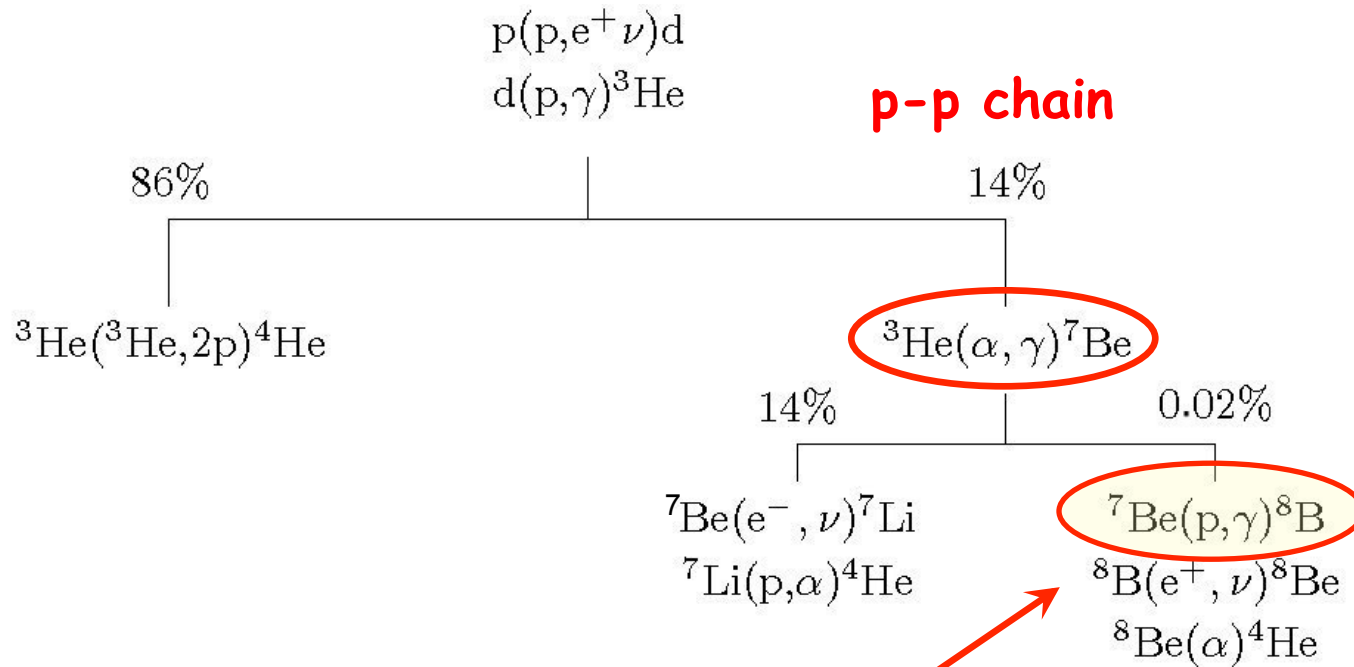
Straightforward to couple different mass partitions in the NCSM/RGM formalism

Applications to (d,p) and (d,n) reactions
Example: $^3\text{He}(d,p)^4\text{He}$

Work in progress:
 $^7\text{Li}(d,p)^8\text{Li}$ & $^8\text{Li}(d,p)^9\text{Li}$



Solar p-p chain



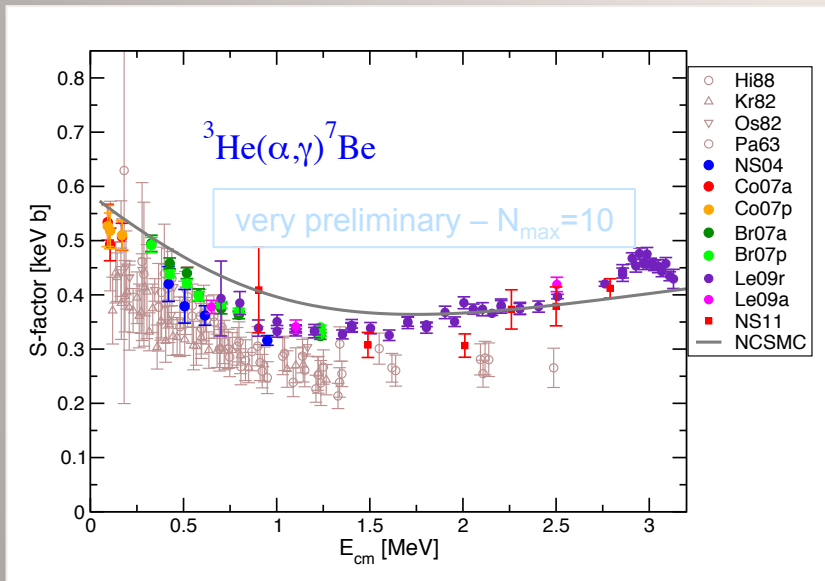
Solar neutrinos
 $E_\nu < 15 \text{ MeV}$

$^3\text{He}(^4\text{He}, \gamma)^7\text{Be}$ & $^7\text{Be}(p, \gamma)^8\text{B}$ radiative capture

- NCSMC & NCSM/RGM calculations
 - Soft NN potential (chiral SRG- $N^3\text{LO}$ with $\Lambda = 2.1 \text{ fm}^{-1}$ & $\Lambda = 1.86 \text{ fm}^{-1}$)

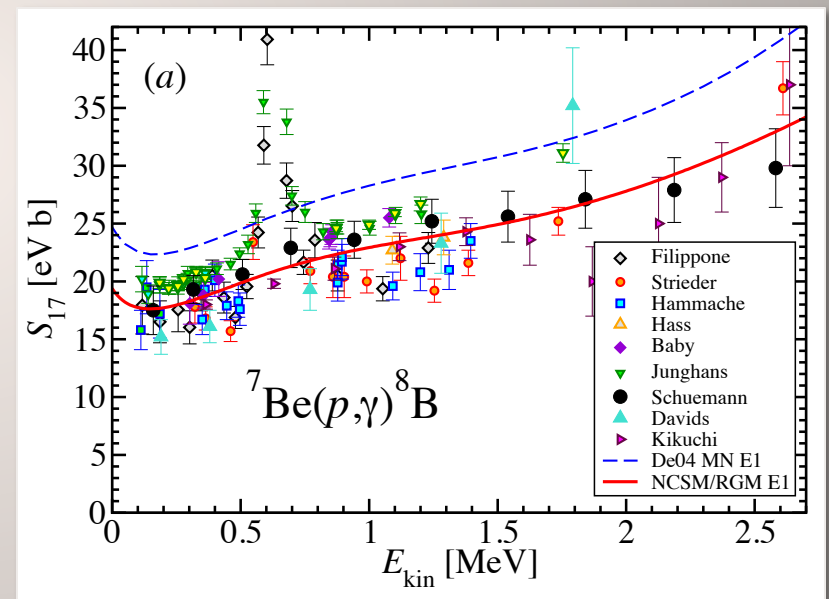
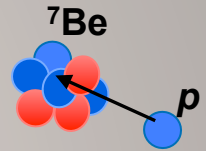
In progress

J. Dohet-Eraly, P.N., S. Quaglioni, W. Horiuchi, G. Hupin



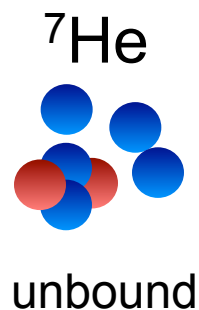
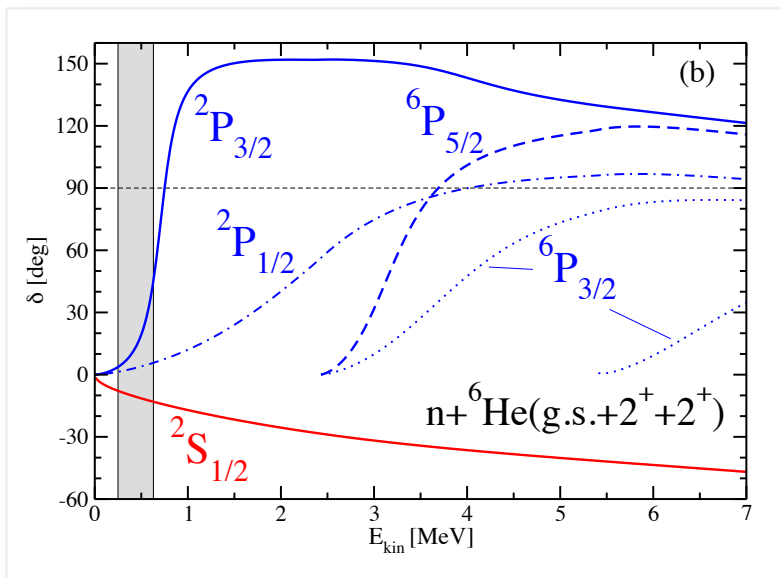
Preliminary: $N_{\text{max}}=10$
 $E_{\text{th}}(^7\text{Be}) = -1.32 \text{ MeV}$ (Expt. -1.59 MeV)

P.N., R. Roth, S. Quaglioni,
 Physics Letters B 704 (2011) 379



$^8\text{B } 2^+$ g.s. **bound** by 136 keV (expt. 137 keV)
 $S(0) \sim 19.4(0.7) \text{ eV b}$
 Current data evaluation: $S(0) = 20.8(2.1) \text{ eV b}$

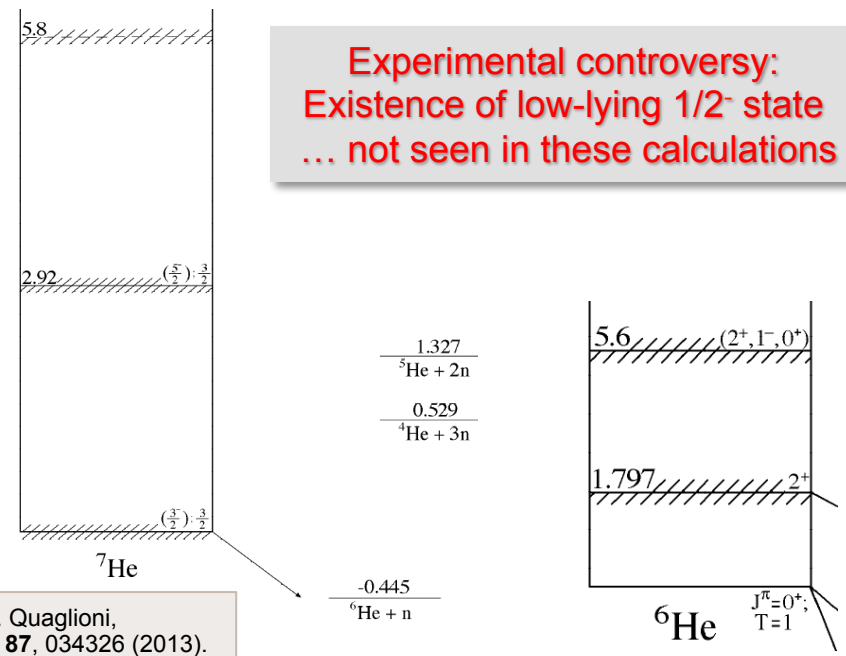
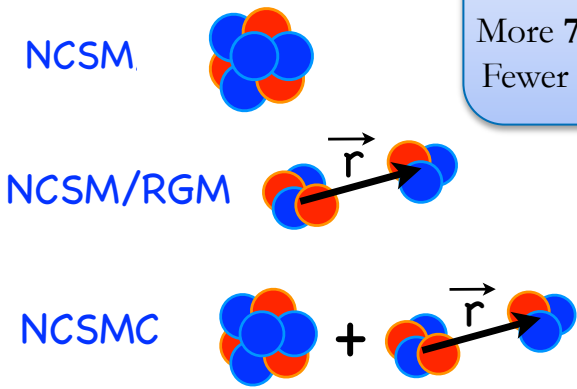
NCSM with continuum: ${}^7\text{He} \leftrightarrow {}^6\text{He}+n$



| J^π | experiment | | | NCSMC | |
|---------|------------|----------|------|-------|----------|
| | E_R | Γ | Ref. | E_R | Γ |
| $3/2^-$ | 0.430(3) | 0.182(5) | [2] | 0.71 | 0.30 |
| $5/2^-$ | 3.35(10) | 1.99(17) | [40] | 3.13 | 1.07 |
| $1/2^-$ | 3.03(10) | 2 | [11] | 2.39 | 2.89 |
| | 3.53 | 10 | [15] | | |
| | 1.0(1) | 0.75(8) | [5] | | |

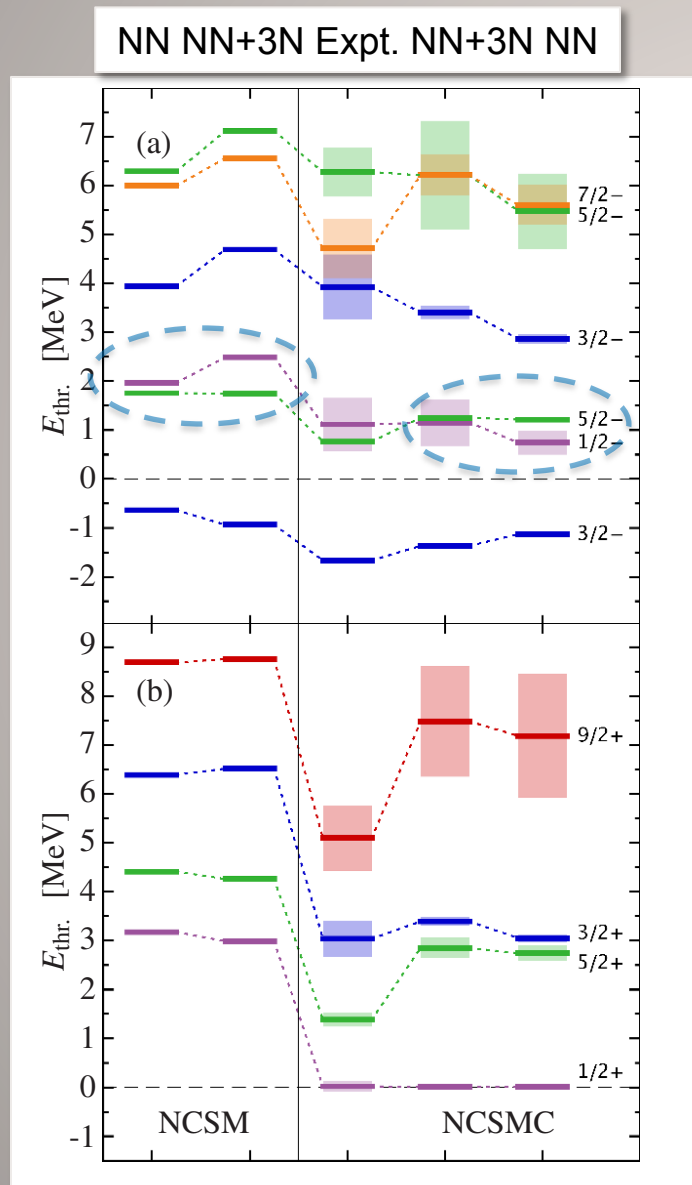
[11] A. H. Wuosmaa *et al.*, Phys. Rev. C **72**, 061301 (2005).

NCSMC
with three ${}^6\text{He}$ states
and ten ${}^7\text{He}$ eigenstates
More 7-nucleon correlations
Fewer ${}^6\text{He}$ -core states needed



S. Baroni, P. N., and S. Quaglioni,
PRL **110**, 022505 (2013); PRC **87**, 034326 (2013).

Structure of ${}^9\text{Be}$: bound states and resonances



${}^9\text{Be}$ is a stable nucleus
 ... but all its excited states unbound
 A proper description requires to include
 effects of continuum

Three-nucleon interaction *and* continuum
 improve agreement with experiment for
 negative parity states

Continuum crucial for the description of
 positive-parity states

J. Langhammer, P. N., G. Hupin, S. Quaglioni, A. Calci, R. Roth,
 in preparation

p+¹⁰C scattering: structure of ¹¹N resonances

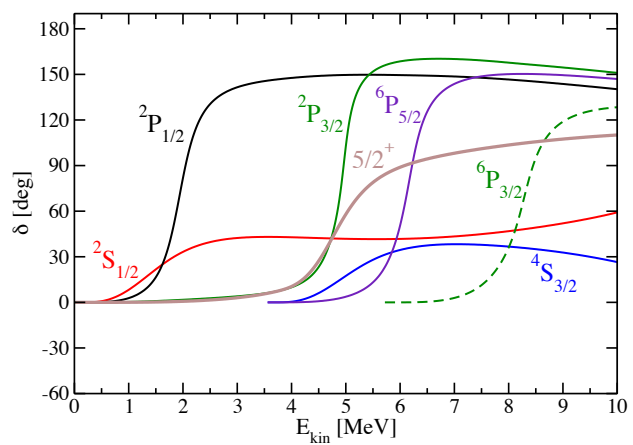
¹¹N from chiral NN+3N within NCSMC

¹¹N Expt. (TUNL evaluation)

– Preliminary

| J^π | T | E_{res} [MeV] | E_x [MeV] | Γ [keV] |
|--------------------|-----|------------------------|-------------|----------------|
| | | | | “4100” |
| ✓ 1/2 ⁺ | 3/2 | 1.35 | 0 | 580 |
| ✓ 1/2 ⁻ | 3/2 | 1.94 | 0.59 | 280 |
| ✓ 3/2 ⁻ | 3/2 | 4.69 | 3.34 | 1790 |
| 5/2 ⁺ | 3/2 | 4.75 | 3.40 | “4760” |
| 3/2 ⁺ | 3/2 | 4.95 | 3.60 | 470 |
| 5/2 ⁻ | 3/2 | 5.95 | 4.60 | 620 |
| 3/2 ⁻ | 3/2 | 7.68 | 6.33 | |

| E_{res} (MeV ± keV) | E_x (MeV ± keV) | $J^\pi; T$ | Γ (keV) |
|------------------------------|-------------------|------------------------------|----------------|
| 1.49 ± 60 | 0 | $\frac{1}{2}^+; \frac{3}{2}$ | 830 ± 30 |
| 2.22 ± 30 | 0.73 ± 70 | $\frac{1}{2}^-$ | 600 ± 100 |
| 3.06 ± 80 | (1.57 ± 80) | | < 100 |
| 3.69 ± 30 | 2.20 ± 70 | $\frac{5}{2}^+$ | 540 ± 40 |
| 4.35 ± 30 | 2.86 ± 70 | $\frac{3}{2}^-$ | 340 ± 40 |
| 5.12 ± 80 | (3.63 ± 100) | $(\frac{5}{2}^-)$ | < 220 |
| 5.91 ± 30 | 4.42 ± 70 | $(\frac{5}{2}^-)$ | |
| 6.57 ± 100 | 5.08 ± 120 | $(\frac{3}{2}^-)$ | 100 ± 60 |



$$\Gamma = \frac{2}{\left. \frac{\partial \delta(E_{\text{kin}})}{\partial E_{\text{kin}}} \right|_{E_{\text{kin}}=E_R}}$$

Negative parity 1/2⁻ and 3/2⁻ resonances in a good agreement with the current evaluation

Positive parity resonances too broad
– N_{max} convergence

p+¹⁰C scattering: structure of ¹¹N resonances

¹¹N from chiral NN+3N within NCSMC

¹¹N Expt. (TUNL evaluation)

– Preliminary

| J^π | T | E_{res} [MeV] | E_x [MeV] | Γ [keV] |
|--------------------|-----|------------------------|-------------|----------------|
| | | | | “4100” |
| ✓ 1/2 ⁻ | 3/2 | 1.94 | 0.59 | 580 |
| ✓ 3/2 ⁻ | 3/2 | 4.69 | 3.34 | 280 |
| 5/2 ⁺ | 3/2 | 4.75 | 3.40 | 1790 |
| 3/2 ⁺ | 3/2 | 4.95 | 3.60 | “4760” |
| 5/2 ⁻ | 3/2 | 5.95 | 4.60 | 470 |
| 3/2 ⁻ | 3/2 | 7.68 | 6.33 | 620 |

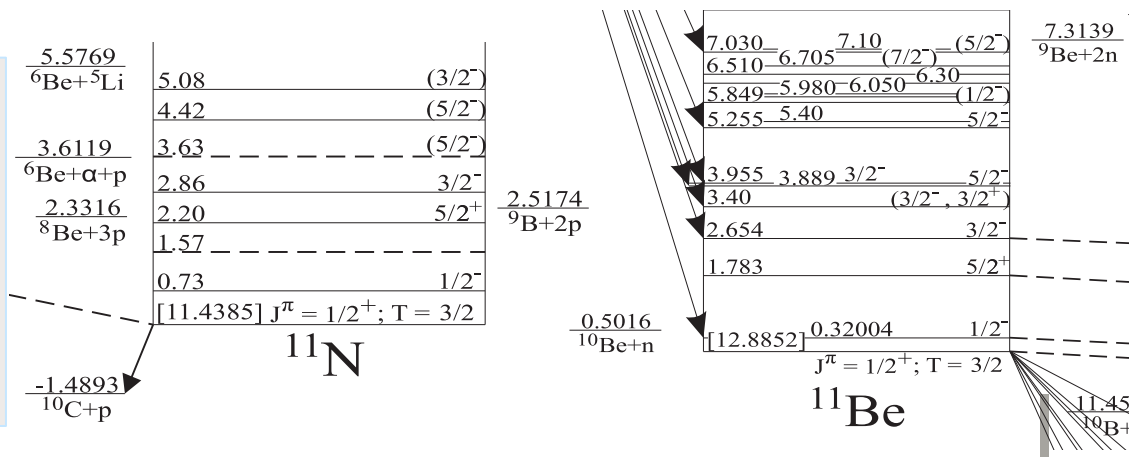
| E_{res} (MeV ± keV) | E_x (MeV ± keV) | $J^\pi; T$ | Γ (keV) |
|------------------------------|-------------------|------------------------------|----------------|
| 1.49 ± 60 | 0 | $\frac{1}{2}^+; \frac{3}{2}$ | 830 ± 30 |
| 2.22 ± 30 | 0.73 ± 70 | $\frac{1}{2}^-$ | 600 ± 100 |
| → 3.06 ± 80 | (1.57 ± 80) | | < 100 |
| 3.69 ± 30 | 2.20 ± 70 | $\frac{5}{2}^+$ | 540 ± 40 |
| 4.35 ± 30 | 2.86 ± 70 | $\frac{3}{2}^-$ | 340 ± 40 |
| → 5.12 ± 80 | (3.63 ± 100) | $(\frac{5}{2}^-)$ | < 220 |
| → 5.91 ± 30 | 4.42 ± 70 | $(\frac{5}{2}^-)$ | |
| 6.57 ± 100 | 5.08 ± 120 | $(\frac{3}{2}^-)$ | 100 ± 60 |

No candidate for 3.06 MeV resonance

We predict only one 5/2⁻ resonance below the 3/2₂⁻

Calculations suggest that either 5.12 MeV or 5.91 MeV resonance might be 3/2⁺ instead

NCSMC resonance predictions more in line with assignments in ¹¹Be



NCSM/RGM for three-body clusters: Structure of ${}^6\text{He}$

${}^4\text{He} + n + n$

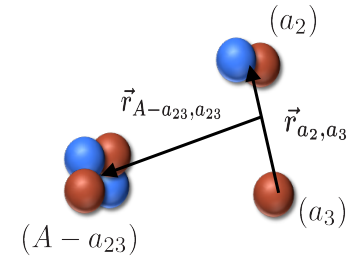
PRL 113, 032503 (2014) PHYSICAL REVIEW LETTERS week ending 18 JULY 2014

${}^4\text{He} + n + n$ Continuum within an *Ab initio* Framework

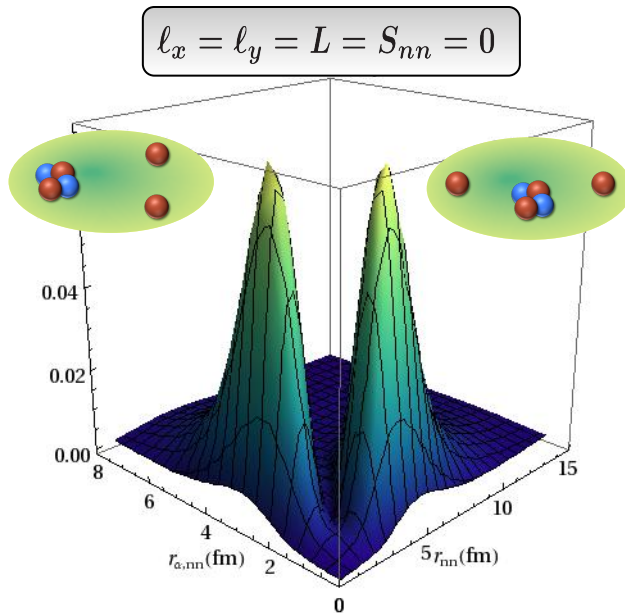
Carolina Romero-Redondo,^{1,*} Sofia Quaglioni,^{2,†} Petr Navrátil,^{1,‡} and Guillaume Hupin^{2,§}

¹TRIUMF, 4004 Wesbrook Mall, Vancouver, British Columbia V6T 2A3, Canada

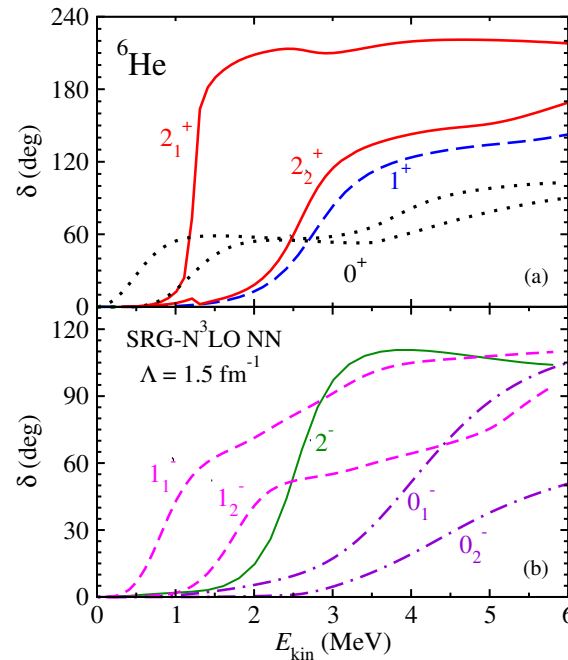
²Lawrence Livermore National Laboratory, P.O. Box 808, L-414, Livermore, California 94551, USA



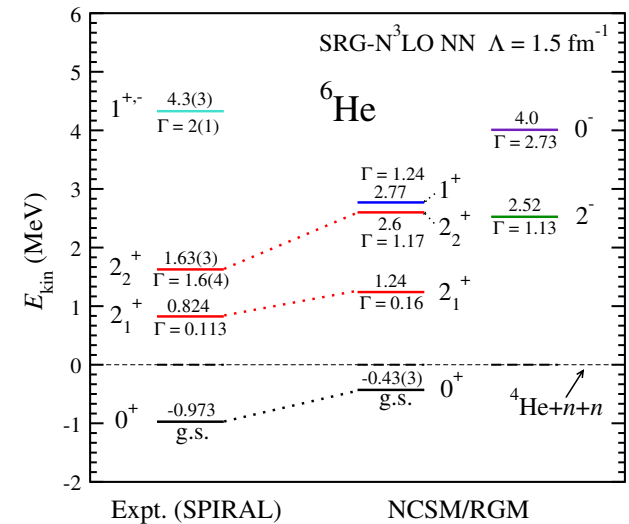
${}^6\text{He}$ bound 0^+ ground state



${}^6\text{He}$ resonances and continuum



Comparison to recent experiment



${}^5\text{H} \approx {}^4\text{He} + n + n$ in progress

Conclusions and Outlook

- *Ab initio* calculations of nuclear structure & reactions is a dynamic field with rapid advances
- Several exact methods applicable to few-nucleon systems ($A=3,4$)
- Significant progress in *ab initio* approaches for p -shell nuclei
- New very successful approaches to medium mass nuclei
- We developed a new unified approach to nuclear bound and unbound states
 - Merging of the NCSM and the NCSM/RGM = **NCSMC**
- **Outlook:**
 - Applications to astrophysics
 - nuclear reactions important for astrophysics (and fusion energy generation)
 - equation of state, symmetry energy
 - Neutrino physics
 - neutrino-nucleus cross sections
 - double beta decay nuclear matrix elements
 - Fundamental symmetries
 - nuclear corrections (CKM unitarity...)
 - Strangeness
 - hypernuclei