Partial-Wave Analysis of Pion Scattering Reactions

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- Spectroscopy of Baryons.
- SAID for Baryon Spectroscopy.
- πN Elastic Scattering.
- $-N(1440)1/2^+$.
- πN inElastic Scattering.
 - *−*π⁻p→ηp.
 - $-\pi^{-}p \rightarrow KY.$
 - $-\pi \rightarrow 2\pi$.
- Pion Photoproduction.
- Pion Electroproduction.













PHYSICAL REVIEW

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Angular Distribution of Pions Scattered by Hydrogen*

H. L. ANDERSON, E. FERMI, R. MARTIN, AND D. E. NAGLE Institute for Nuclear Studies, University of Chicago, Chicago, Illinois (Received March 6, 1953)



The results have been interpreted in terms of phase shift analysis on the assumption that the scattering is mainly due to states of isotopic spins $\frac{3}{2}$ and $\frac{1}{2}$ and angular momenta s_i , p_i and p_i .

PHYSICS REPORTS (Review Section of Physics Letters) 96, Nos. 2 & 3 (1983) 71-204. North-Holland Publishing Company

BARYON SPECTROSCOPY

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Received 29 September 1982

In 1952 Fermi and coworkers (Andersen et al. [1952]) discovered the first baryon resonance – the Δ (1238). Since then, hundreds of resonances have been identified and nuclear democracy has given way to fundamental quarks. Baryon spectroscopy is now thirty years old and perhaps approaching a mid-life crisis. For it is inevitable in such a fast-moving field as high energy particle physics, that experiments have moved on beyond the resonance region to higher energies and different priorities. Thus it is probably no exaggeration to say that we now have essentially *all* the experimental data relevant to the low-energy baryon spectrum, that we are *ever* likely to obtain. It is therefore timely to review both the accumulated mass of resonance data, together with the techniques used in its analysis, and also our theoretical framework for understanding the results. The latter is inevitably based on quarks and, by and large, on a

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Spectroscopy of Baryons



``It is clear that we still need much more information about the existence and parameters of many baryon states, especially in the N=2 mass region, before this question of non-minimal SU(6) x O(3) super-multiplet can be settled." **Dick Dalitz**, **1976**.

``The first problem is the notion of a resonance is not well defined. The ideal case is a narrow resonance far away from the thresholds, superimposed on slowly varying background. It can be described by a Breit-Wigner formula and is characterized by a pole in the analytic continuation of the partial wave amplitude into the low half of energy plane.'' **Gerhard Hoehler, 1987**.





"Why N*s are important – The *first* is that nucleons are the stuff of which our world is made. My *second* reason is that they are simplest system in which the quintessentially non-Abelian character of QCD is manifest. The *third* reason is that history has taught us that, while relatively simple, Baryons are sufficiently complex to reveal physics hidden from us in the mesons." **Nathan Isgur, 2000**.

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Baryon Sector at PDG

[J. Beringer et a/[PDG] Phys Rev D 86, 010001 (2012)]

	$\begin{array}{llllllllllllllllllllllllllllllllllll$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} \Xi^{0} & 1/2^{+} & \cdots \\ \Xi^{-} & 1/2^{+} & \cdots \\ \Xi(1530) & \Xi(1620) & \vdots \\ \Xi(1690) & \vdots \\ \Xi(1960) & \Xi(1820) & \vdots \\ \Xi(2250) & \Xi(2250) & \vdots \\ \Xi(2250) & \vdots \\ \Xi(2250) & \vdots \\ \Xi(2250) & \vdots \\ \Xi(2500) & \Xi(2500) & \vdots \\ \Xi(2500) & \Xi(2500) & \vdots \\ \Xi(2500) & \Xi(250$	$\begin{array}{ccccccc} \Lambda_c^+ & 1/2^+ & & & & & & \\ \Lambda_c(2595)^+ & 1/2^- & & & & & \\ \Lambda_c(265)^+ & & & & & & \\ \Lambda_c(2800)^+ & & & & & & & \\ \Lambda_c(2940)^+ & & & & & & \\ \Sigma_c(2455) & 1/2^+ & & & & & \\ \Sigma_c(2520) & & & & & & & \\ \Sigma_c(2520) & & & & & & \\ \Xi_c^+ & & 1/2^+ & & & & \\ \Xi_c^0 & & & & & & \\ \Xi_c^+ & & & & & & \\ \Xi_c^+ & & & & & & \\ \Xi_c(2645) & & & & & \\ \Xi_c(2900) & & & & & \\ \Xi_c(2900) & & & & & \\ \Xi_c(2900) & & & & \\ \Xi_c(3055) & & & & \\ \Xi_c(3055) & & & & \\ \Xi_c(3060) & & & & \\ \Xi_c(3123) & & & \\ \Omega_c^0 & & & & & \\ \Xi_c^0 & & & & & \\ \Omega_c^0 & & & & & \\ \Xi_c^0 & & & & & \\ \Omega_c^0 & & & & & \\ \Xi_c^0 & & & & \\ \Xi_c^0 & & & & \\ \Xi_c^0 & & & & \\ \Omega_c^0 & & & & & \\ \Xi_c^0 & & & & \\ \Xi_c^0 & & & & \\ \Omega_c^0 & & & & \\ \Xi_c^0 & & $	 PDG12 has 112 Baryon Resonances (58 4* & 3* of them). For example for SU(6) x O(3), it would be 434 resonances, if all revealed three 70- and four 56- multiplets were filled in. There are many more states in the QCD inspired models than currently observed. A quick check of the PDG Listings reveals that resonance parameters of many established states are not well determined.
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Prolific source of $N^* \& \Delta^*$ baryons

Measure many channels with different combinations of quantum numbers.

 $\pi N \rightarrow \pi N, \pi \pi N, ... \square$ $\gamma N \rightarrow \pi N, \pi \pi N, ... \square$ $\gamma^* N \rightarrow \pi N, \pi \pi N, ... \square$ $p p \rightarrow p p \pi^0, p p \pi \pi, ... \square$ $J/\Psi \rightarrow p \overline{p} \pi^0, p \overline{n} \pi, ... \square$

ер → ерХ



- Most of PDG info comes from these sources and PWA is main of them.
- πN elastic scattering is highly constrained.
- Resonance structure is correlated.
- Two-body final state, fewer amplitudes.



PWA for non-Strange Baryons & SAID Database



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\mathcal{N}^* and \varDelta^* States coupled to $\pi \mathcal{N}$

[SAID: http://gwdac.phys.gwu.edu/]



GWM Program















SAID for $\pi \mathcal{N} \rightarrow \pi \mathcal{N} \ll \pi^{-} p \rightarrow \eta n$

[R. Arndt, W. Briscoe, IS, R. Workman, Phys Rev C 74, 045205 (2006)]

Energy dependent SP06/WI08 and associated SES [W = 1078 - 2460 MeV]= 0 - 2600 MeV $[\pi N, \pi \Delta, \rho N, \eta N]$ • (4) channel Chew-Mandelstam K-matrix parameterization 3 mapping variables: $g^2/4\pi$, $a[\pi-p]$, Eth • PWs = 30 πN {15 [I=1/2] + 15 [I=3/2]} + 4 ηN [] < 9] • Prms =99 [I=1/2] + 89 [I=3/2] 1st generation ('57–'79) Reaction Data Used by CMB79 and KH84 analyses. 27,136 13,354 10k π^{\pm} p each & 1.5k CXS. $\pi^+ \mathbf{p} \rightarrow \pi^+ \mathbf{p}$ 17% data is polarized. [0 – 2600 MeV] → 10 data/MeV 11,978 22,632 **π⁻p→**π⁻p 2nd generation ('80–'06) 3,115 6.068 *π*⁻**p**→*π*⁰n SAID fits: $[550 - 800 \text{ MeV}] \rightarrow 1 \text{ data/MeV}$

13k π^{\pm} p each, **3k** CXS & **0.3k** π^{-} p \rightarrow nn 25% data is polarized.

Meson Factories [LAMPF, TRIUMF, & **PSI**] are the main source of new

measurements.

There is no discrimination against data

• <u>3rd generation</u> (07'+) New data may come from

J-PARC, HADES, EPECUR, etc

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257 650 *π*⁻**p→**η**n** DR constraint 2,775 671 Jotal 31,479 57,157

> **DRs** have been derived from the first principles.



27 σ^{tot} & **37** P data



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above 800 MeV \rightarrow 0.03 data/MeV

Fit to *TN* Elastic Scattering Data



- Some of structures in 35-year old solutions.
- [KH and CMB] are still considered as resonances.
- W > 1.5 GeV is less constrained by data.







New Observables

[R. Arndt, W. Briscoe, IS, R. Workman, Phys Rev C 74, 045205 (2006)]



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Where is a Resonance?

[R. Arndt, W. Briscoe, IS, R. Workman, Phys Rev C 74, 045205 (2006)]



<u>GWDAC</u> Search for \mathcal{N}^* and Δ^*

We are considering a resonance as a **Pole** in the complex plane which is not far away from the physical axis.

- <u>Applied</u> directly to the data via BW + Bckgr
- <u>Assume</u>: $S \rightarrow S_R S_B$ $S_R = 1 + 2iT_R$ $T_R = (\Gamma_e/2) / [W_R - W - i(\Gamma_e/2 + \Gamma_I/2)]$ $\Gamma = \Gamma_e + \Gamma_I \quad \Gamma_e = \rho_e \Gamma R \quad \Gamma_I = \rho_i \Gamma (1 - R)$ $T_B = K_B (1 - iK_B)^{-1} \quad K_B = a + b(W - W_R) + c$
- <u>Map</u> $\chi^2[W_R,\Gamma]$ while searching all other PW parameters Look for significant improvement
- <u>Subjective variables are</u>
 - Energy binning
 - Strength of constraints
 - Which PW to be searched

 <u>Standard PWA</u>
 Tends (by construction) to miss narrow Resonances with Γ < 30 MeV
 Reveals only wide Resonances, but not too wide [Γ < 500 MeV] and possessing not too small BR [BR > 0.04]
 <u>Modified PWA</u>
 Allows to put a resonance by hands with subsequent refitting the data Then the search will allow to see how reliable/tolerable it is



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Complex Energy Plane vs BW fits

I = 1/2





• There are shifts between **Pole** and **BW** prms.

 $M = ReW_p$ $\Gamma = 2*ImW_p$



Partial Waves [L₍₂₁₎₍₂₁₎]

[R. Arndt, W. Briscoe, IS, R. Workman, Phys Rev C 74, 045205 (2006)]



97(1440)1/2+



Courtesy of Donna Arndt, 2010





$\mathcal{N}(1440)^{****} - What is Known$

[J. Beringeret al [PDG] Phys Rev D 86, 010001 (2012)]



Two-faced Janus Roman God of Gates & Doors Dick Arndt: ``This is one of mysterious Resonances"

PDG- PWA-Po	le Ref	Re(MeV)	–2xlm(MeV)
	BnGa12 SAID-SP06 KH93 CMU80	1370 ± 4 1359 1388 1385 1375±30	190 ± 7 164 166 164 180 ± 40	1 st Riemann sheet 2 nd Riemann sheet
PDG- PWA-BV	/ Ref	Mass(MeV)	Width(Me	V) BR

$$f(E) = \frac{k}{\left(E^2 - M^2\right)^2 + M^2\Gamma^2} \begin{bmatrix} \mathsf{M} = & \mathsf{ReW}_\mathsf{p} \\ \Gamma = \mathsf{2*ImW}_\mathsf{p} \end{bmatrix}$$





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Discovery and First Direct Measurements of $\mathcal{N}(1440)$





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Complex Energy Plane for **P**₁₁

[R. Arndt, W. Briscoe, IS, R. Workman, Phys Rev C 74, 045205 (2006)]



Two Pole Observation



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Direct Measurements of $\mathcal{N}(1440)$: Hadronic Probes





Direct Measurements of
$$\mathcal{N}(1440)$$
: EM Probes

Relative contributions of various singularities may be different in different processes



 $\mathcal{N}(1440)$ Puzzle for Jlab CLAS12

 Most of analyses of N(1440) are based on its BW parameterization, which assumes that the Res is related to an isolated Pole

- However, the latest GW PWA for the elastic πN scattering gives evidence that N(1440) corresponds to a more complicated case of several nearby singularities in the amplitude
- Then, the BW description is only an efficient one for N(1440), which could be different in different processes
- Some inelastic data indirectly support this point:

they give the N(1440) BW mass and width essentially different from the PDG BW values



 Since Q²-dependences for contributions of different singularities may be different, the set of several singularities might provide the N(1440) BW mass and width depending on the Q²

• This problem can be studied in future measurements with JLab CLAS12





 $N(2700) 13/2^+$

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Status of Non-strange Resonances

J. Beringer *et al* (PDG) Phys Rev D **86**, 010001 (2012)

• More than half of states have poor evidence.

Most of states need more work to do.

<u>GW SAI</u>	<u>) Co</u>	ion Status as seen I = 1/2						Most			
Particle J^P	overa	Statu II πN	γN	$N\eta$	Νσ	Nω	ΛK	Σŀ	Κ Νρ	$\Delta \pi$	• Whe
$N = 1/2^+$	****										
$N(1440) 1/2^+$	****	****	****		***				*	***	1 = 3/2
$N(1520) 3/2^{-}$	****	****	****	***					***	***	
$N(1535) 1/2^{-}$	****	****	****	****					**	*	Particle J
$N(1650) 1/2^{-}$	****	****	***	***			***	**	**	***	A(1232) 3/
$N(1675) 5/2^{-1}$	****	****	***	*			*		*	***	$\Delta(1202) 0_{1}$
$N(1680) 5/2^+$	****	****	****	*	**				***	***	$\Delta(1600) 0$
N(1685) ??	*										$\Delta(1020)$ 1/ $\Delta(1700)$ 3/
$N(1700) 3/2^{-1}$	***	***	**	*			*	*	*	***	$\Delta(1750) 1$
$N(1710) 1/2^+$	***	***	***	***		**	***	**	*	**	$\Delta(1750) 1$
$N(1720) 3/2^+$	****	****	***	***			**	**	**	*	$\Delta(1900) 1/$
$N(1860) 5/2^+$	**	**							*	*	$\Delta(1905) 5/$
$N(1875) 3/2^{-1}$	***	*	***			**	***	**		***	$\Delta(1910) 1/$
$N(1880) 1/2^+$	**	*	*		**		*				$\Delta(1920) 3/$
$N(1895) 1/2^{-}$	**	*	**	**			**	*			$\Delta(1930) 5/$
$N(1900) 3/2^+$	***	**	***	**		**	***	**	*	**	$\Delta(1940)$ 3
$N(1990) 7/2^+$	**	**	**					*			$\Delta(1950) 7_{0}$
$N(2000) 5/2^+$	**	*	**	**			**	*	**		$\Delta(2000) 5/$
$N(2040) 3/2^+$	*										$\Delta(2150) 1_{0}$
$N(2060) 5/2^{-1}$	**	**	**	*				**			$\Delta(2200) 7/$
$N(2100) 1/2^+$	*										$\Delta(2300) 9/$
$N(2150) 3/2^{-}$	**	**	**	26	*		**			**	$\Delta(2350) 5/$
$N(2190) 7/2^{-1}$	****	****	***	201	V		**		*		$\Delta(2390) 7$
$N(2220) 9/2^+$	****	****		11 *	***						$\Delta(2400) 9$
$N(2250) 9/2^{-}$	****	****		5	***						$\Delta(2420)$ 11
N(2600) 11/2-	\mathbf{A}	***		7	**						Δ(2750) 13

3 *

• Most of QCD models predict more states than observed.

• Where are missing resonances?

		*	***	l = 3/2					S	tatus	as se	en in	_			this
		*** **	*** *	Particle J^F	overa	Statu ll πN	γN	$N\eta$	Νσ	Nω	ΛK	ΣK	Νρ	$\Delta \pi$	Ī	e for
*** * *** *** *** **	** * ** ** ** ** **	** ** * ** * * *	* *** *** *** ** ** ** ** ** ** ** **	$\begin{array}{c} \text{Particle } J^F \\ \hline \Delta(1232) \ 3/2 \\ \Delta(1600) \ 3/2 \\ \Delta(1600) \ 3/2 \\ \Delta(1620) \ 1/2 \\ \Delta(1700) \ 3/2 \\ \Delta(1750) \ 1/2 \\ \Delta(1900) \ 1/2 \\ \Delta(1900) \ 1/2 \\ \Delta(1900) \ 3/2 \\ \Delta(1910) \ 1/2 \\ \Delta(1910) \ 1/2 \\ \Delta(1920) \ 3/2 \\ \Delta(1940) \ 3/2 \\ \Delta(1950) \ 7/2 \\ \Delta(2000) \ 5/2 \\ \Delta(2150) \ 1/2 \\ \Delta(2200) \ 7/2 \\ \Delta(2300) \ 9/2 \\ \Delta(2350) \ 5/2 \\ \Delta($	overal 2+ **** 2+ **** 2- **** 2- **** 2+ **** 2+ **** 2+ **** 2+ **** 2+ **** 2+ **** 2+ **** 2+ **** 2+ **** 2+ **** 2+ **** 2+ **** 2+ **** 2+ *** 2- * 2+ *** 2- * 2+ *** 2- * 2+ *** 2- * 2- * 2- * 2- * 2- * 2- * 2- * 2- * 2- * 2- * 2- * 2- * 2- *	11 πN **** *** *** *** *** *** ***	γN **** *** *** *** *** *** *** *** ***	Nη F o	Nσ r b i	d d d	AK n	*** *** (seei ***	Nρ * ** ** * * *	$\Delta \pi$ **** *** *** *** *** $\Delta \eta$) **** **		t GWU analysis (ARNDT 06) finds no evidence fo
**		*		$\begin{array}{c} \Delta(2390) \ 7/2 \\ \Delta(2400) \ 9/2 \\ \Delta(2420) \ 11 \\ \Delta(2750) \ 13 \\ \Delta(2950) \ 15 \end{array}$	2 ⁺ * 2 ⁻ ** /2 ⁺ **** /2 ⁻ /2 ⁺	* ** **** ** **	22 7 ' 3 ' 7 ' 5 '	∆ * **** *** **	<u>D</u>		n n		Stat	tes		The latest





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Status of Non-strange Resonances Twenty Years Ago – Rosenfeld Tables

L. Montanet et al (PDG) Phys Rev D 50, 1173 (1994)

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		0 1			Stat	us as s	een in -		
Particle	$L_{2I\cdot 2J}$	Overall status	$N\pi$	$N\eta$	ΛK	ΣK	$\Delta \pi$	$N\rho$	$N\gamma$
N(939)	P11	****							
N(1440)	P11	****	****	*			***	*	***
N(1520)	D_{13}	****	****	*			****	****	****
N(1535)	S11	****	****	****			*	**	***
N(1650)	S11	****	****	*	***	**	***	**	***
N(1675)	D_{15}	****	****	*	*		** **	*	****
N(1680)	F_{15}	****	****				****	****	****
N(1700)	D_{13}	***	***	*	**	*	**	*	**
N(1710)	P11	***	***	**	**	*	**	*	***
N(1720)	P_{13}	****	****	*	**	*	*	**	**
N(1900)	P_{13}	**	**					*	
N(1990)	F_{17}	**	**	*	*	*			*
N(2000)	F_{15}	**	**	*	*	*	*	**	
N(2080)	D_{13}	**	**	*	*				*
N(2090)	S11	*	*						
N(2100)	P11	*	*						
N(2190)	G_{17}	****	****	*	*	*		*	*
N(2200)	D_{15}	**	**	*	*				
N(2220)	H_{19}	****	****	*		22 N	*		
N(2250)	G_{19}	****	****	*		11 *	<mark>***</mark>		
N(2600)	I111	***	***				e ale ale		
N(2700)	K_{113}	**	**			3 1	• ~ ~		
		\mathbf{A}				6 *	• *		
						2 1			

Particle	$L_{2I\cdot 2j}$	Overall status	$N\pi$	$N\eta$	ΛK	ΣK	$\Delta \pi$	$N\rho$	$N\gamma$
$\Delta(1232)$	P_{33}	****	****	F					** **
$\Delta(1600)$	P_{33}	***	***	0			***	*	**
$\Delta(1620)$	S_{31}	****	****	r			****	****	***
$\Delta(1700)$	D_{33}	****	****	ь		*	***	**	***
$\Delta(1750)$	P_{31}	*	*	i					
$\Delta(1900)$	S_{31}	***	***	c	i	*	*	**	*
$\Delta(1905)$	F_{35}	****	****		d	*	**	**	***
$\Delta(1910)$	P_{31}	****	****		е	*	*	*	*
$\Delta(1920)$	P_{33}	***	***		n	*	**		*
$\Delta(1930)$	D_{35}	***	***			*			**
$\Delta(1940)$	D_{33}	*	*	F					
$\Delta(1950)$	F_{37}	****	****	0		*	****	*	** **
$\Delta(2000)$	F_{35}	**		r				**	
$\Delta(2150)$	S_{31}	*	*	ь					
$\Delta(2200)$	G_{37}	*	*	i					
$\Delta(2300)$	H_{39}	**	**	ć	1				
$\Delta(2350)$	D_{35}	*	*		d				
$\Delta(2390)$	F_{37}	*	*		<				
$\Delta(2400)$	G_{39}	**	**		22	Δ*			
$\Delta(2420)$	H_{311}	****	****		7	****			*
$\Delta(2750)$	I313	**	**						
$\Delta(2950)$	K_{315}	**	**		4	<u> </u>			
		\frown			- 5	**			
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Courtesy of Vitaly Shklyar, Spring 2013



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Coupled Channel Fit for S₁₁ & D₁₃: 7n BRS [R. Arndt, W. Briscoe, IS, R. Workman, A. Gridnev, Phys Rev C 72, 045202 (2005)]





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$\pi \not \rightarrow \eta \pi Puzzle$

[R. Arndt, W. Briscoe, IS, R. Workman, Phys Rev C 74, 045205 (2006)]



- Most of Nimrod data do not satisfy requirements
 [systematics (10% or more), momentum err (up to 50 MeV/c), and so on]
- For that reason, SAID is not able to use them in $\pi^-p \rightarrow \pi^-p$, π^0n , & ηn PWAs



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 $\pi \rho \rightarrow \eta \pi Puzzle$

[R. Arndt, W. Briscoe, IS, R. Workman, Phys Rev C 74, 045205 (2006)] [R. Arndt, W. Briscoe, IS, R. Workman, A. Gridnev, Phys Rev C 72, 045202 (2005)]



 $\pi \mathcal{N} \rightarrow \mathcal{K} \mathcal{Y} \mathcal{P} uzzle$



 $\pi \mathcal{N} \to \mathcal{K} A$



RAL78

0.6

0.2

-0.2

-0.6

-1.5

-1



ANL75

RAL78

RAL78

0.2

-0.2

0.6

Courtesy of Eberhard Klempt, PWA7 Workshop, Sept 2013

0.6

1



RRTF Workshop, Darmstadt, Germany, Oct 2013

-1.5

-1

-0.6



$\pi \mathcal{N} \rightarrow \mathcal{K} \mathcal{A} Amplitudes$



- Points with error bars: SES of Shresta and Manley. Curves: BnGa2011.
- Shresta and Manley start from a model-dependant fit and select the solution of the SES which is closest to the energy-dependent fit.
- They first freeze the S_{11} wave, then S_{11} and P_{11} , to the energy dependent solution.

Courtesy of Eberhard Klempt, PWA7 Workshop, Sept 2013







Our knowledge of $\pi\Delta$, ρN , and other quasi-two-body $\pi\pi N$ channels comes mainly from **Isobar-model** analyses of the $\pi N \rightarrow \pi\pi N$ data.





Previous $\pi / \rightarrow \pi \pi / Measurements$



$\pi \mathbb{N} \to \pi \pi \mathbb{N}$ in Isobar Model

[D.M. Manley, R. Arndt, Y. Goradia, V. Teplitz, Phys Rev D 30, 904 (1984)]



• $\pi N \rightarrow \pi \pi N$ is the dominant inelastic reaction in πN scattering above **1300** MeV, $\sigma_{\text{unel}} \sim \sigma(\pi \pi N)$

• **Drawbacks** – analysis of 3-body final states is complicated (many partial waves are involved).



The **total amplitude** for a given charge channel can be written as a coherent sum over all **isobars** and **partial waves**.

 Many of the **3**- and **4**-star resonances have large decay branching ratios to ππN channels.

• There remains a strong need for detailed new measurements in all charge channels!





$\pi / \rightarrow \pi \pi /$ in Isobar Model at low Energies

[V. Kozhevnikov & S. Sherman, Phys Atom Nucl 71, 1860 (2008)]



ANL-OSAKA (EBAC) Dynamical Coupled-Channels Study of $\pi N \rightarrow \pi \pi N$ Reactions

[H. Kamano, et al Phys Rev C 79, 025206 (2005)]



Resonance $\rightarrow \mathcal{N}\rho$ Branching Ratios

	GiBUU12	UrQMD09	KSU12	KSU92	BnGa12	CLAS12	PDG12	
N(1520)3/2 ⁻	21	15	20.9(7)	21(4)	10(3)	12.7(4.3)	20(5)	D13
N(1720)3/2+	87	73	1.4(5)	87(5)	10(13)	47.5(21.5	77.5(7.5)	P13
∆(1620)1/2⁻	29	5	26(2)	25(6)	12(9)	37(12)	16(9)	S 31
∆ (1905)5/2 ⁺	87	80	<6	86(3)	42(8)		>60	F35

Partial courtesy of Piotr Salabura, Sept 2013

CLAS12: V. Mokeev *et al*, Phys Rev C 86, 035203 (2012); V. Mokeev, PC
BnGa12: A.V, Anisovich *et al*, Eur Phys J A 48, 15 (2012)
GiBUU12: J. Weil *et al*, Eur Phys J A 48, 111 (2012); J. Weil, PC
KSU92: D.M. Manley and E.M. Saleski, Phys Rev D 45, 055203 (1992)
KSU12: M. Shrestha and D.M. Manley, Phys Rev D 86, 055203 (2012)
PDG12: J. Beringer *et al* [RPP] Phys Rev D 86, 010001 (2012)
UrQMD09: K. Schmidt *et al*, Phys Rev C 79, 4002 (2009)





CLAS \mathcal{N}^* candidate at 1720 MeV in $p\pi^+\pi^-$?



Origin of the structure at W~1.7 GeV for the first time Jefferson Lab

observed in the CLAS $M\pi\pi$ data

• P₁₃(1720) state with hadronic decays fit to the CLAS data 2.94 $<\chi^2$ /dp<3.15

	M (GeV)	Γ _{tot} (MeV)	BR(πΔ) (%)	BR(ρp) (%)
P ₁₃ (1720) CLAS	1.728± 0.005	133±19	66±26	16±11
P ₁₃ (1720) PDG	1.70 -1.75	150-400	comp. with 0.	70-85

Courtesy of Victor Mokeev, Sept 2013

 3/2⁺(1725) candidate state, hadronic parameters of others N*'s are within PDG uncertainties 2.78<χ²/dp<2.9



Precise data on $\pi\Delta \& \rho p$ hadronic decays of P₁₃(1720) is critical in order to understand origin of the structure at W~1.7 GeV in N $\pi\pi$ electroproduction cross sections. This information can be obtained in the studies of $\pi N \rightarrow \pi \pi N$ reactions at **J-PARC** and **HADES**...

10/6/2013









Single **Pion** Photoproduction

- An accurate evaluation of the EM couplings $N^*(\Delta^*) \rightarrow \gamma N$ from meson photoproduction data remains a paramount task in hadron physics.
- Only with good data on both the proton and neutron targets, one can hope to disentangle the isoscalar & isovector EM couplings of various N*& Δ* resonances,
 K.M. Watson, Phys Rev 95, 228 (1954); R.L. Walker, Phys Rev 182, 1729 (1969) as well as the isospin properties of non-resonant background amplitudes.



 The lack of the γn→π⁻p & γn→π⁰n data does not allow us to be as confident about the determination of neutron couplings relative to those of the • proton.

• The radiative decay width of neutral baryons may be extracted from $\pi^- \& \pi^0$ photoproduction off the neutron, which involves a bound neutron target and needs the use of model-dependent nuclear (FSI) corrections.

A.B. Migdal, JETP 1, 2 (1955); K.M. Watson, Phys Rev 95, 228 (1954)



Where We Are Now

Some of the N* baryons [N(1675)D₁₅, for instance] have stronger EM coupling to the neutron than to the proton but parameters are very uncertain.

$$N(1675) 5/2^{-1} I(J^P) = \frac{1}{2}(\frac{5}{2})$$
 Status: ***



• **PDG** estimates for the $A_{1/2} \& A_{3/2}$ decay amplitudes of the N(1720)P₁₃ state are consistent with *zero*, while the recent **SAID** determination gives small but non-vanishing values.

$$N(1720) \ 3/2^+ I(J^P) = \frac{1}{2}(\frac{3}{2}^+) \ \text{Status:} \ * * * P_{13} N(1720)3/2^+ \rightarrow p\gamma, \ \text{helicity-1/2 ampl, A1/2. -0.01 to 0.11 +0.095\pm0.002} N(1720)3/2^+ \rightarrow p\gamma, \ \text{helicity-3/2 ampl, A3/3: -0.019\pm0.029} -0.048\pm0.002 O(1720)3/2^+ \rightarrow p\gamma, \ \text{helicity-3/2 ampl, A3/3: -0.019\pm0.029} -0.048\pm0.002 O(1720)3/2^+ \rightarrow p\gamma, \ \text{helicity-3/2 ampl, A3/3: -0.019\pm0.029} -0.048\pm0.002 O(1720)3/2^+ \rightarrow p\gamma, \ \text{helicity-3/2 ampl, A3/3: -0.019\pm0.029} -0.048\pm0.002 O(1720)3/2^+ \rightarrow p\gamma, \ \text{helicity-3/2 ampl, A3/3: -0.019\pm0.029} -0.048\pm0.002 O(1720)3/2^+ \rightarrow p\gamma, \ \text{helicity-3/2 ampl, A3/3: -0.019\pm0.029} -0.048\pm0.002 O(1720)3/2^+ \rightarrow p\gamma, \ \text{helicity-3/2 ampl, A3/3: -0.019\pm0.029} -0.048\pm0.002 O(1720)3/2^+ \rightarrow p\gamma, \ \text{helicity-3/2 ampl, A3/3: -0.019\pm0.029} -0.048\pm0.002 O(1720)3/2^+ \rightarrow p\gamma, \ \text{helicity-3/2 ampl, A3/3: -0.019\pm0.029} -0.048\pm0.002 O(1720)3/2^+ \rightarrow p\gamma, \ \text{helicity-3/2 ampl, A3/3: -0.019\pm0.029} -0.048\pm0.002 O(1720)3/2^+ \rightarrow p\gamma, \ \text{helicity-3/2 ampl, A3/3: -0.019\pm0.029} -0.048\pm0.002 O(1720)3/2^+ \rightarrow p\gamma, \ \text{helicity-3/2 ampl, A3/3: -0.019\pm0.029} -0.048\pm0.002 O(1720)3/2^+ \rightarrow p\gamma$$

• Other unresolved issues relate to the second P_{11} , $N(1710)P_{11}$, that are not seen in the recent $\pi N PWA$, contrary to other PWAs used by the PDG12.

$$N(1710) 1/2^{+} I(J^{P}) = \frac{1}{2}(\frac{1}{2}^{+}) \text{ Status: } * * P_{11} \text{ The latest GWU analysis (ARNDT 06) finds no evidence for this resonance.}$$

$$R. Arndt, Ya. Azimov, M. Polyakov, IS, R. Workman, Phys Rev C 69, 035208 (2004) \longrightarrow N(1680) 1/2^{+}$$

$$N(1680) 1/2^{+} I(J^{P}) = \frac{1}{2}(\frac{1}{2}^{+}) \text{ Status: } * * P_{11} \text{ The latest GWU analysis (ARNDT 06) finds no evidence for this resonance.}$$

$$N(1680) 1/2^{+} I(J^{P}) = \frac{1}{2}(\frac{1}{2}^{+}) \text{ Status: } * * P_{11} \text{ The latest GWU analysis (ARNDT 06) finds no evidence for this resonance.}$$

Complete Experiment in **Pion** Photo Production



Igor Strakovsky 48 🕔

RRTF Workshop, Darmstadt, Germany, Oct 2013

10/6/2013

Direct Amplitude Reconstruction in Pion Photo Production

$\gamma \quad \mathbf{N} \rightarrow \mathbf{N} \ \pi$





SAID for Pion Photoproduction

[W. Chen et al, Phys Rev C 86, 015206 (2012)]



10/6/2013



Recent SAID Progress in PionPR

DU13:	included recent CLAS π^{0} p & π^{+} n Σ	[M. Dugger et al, arXiv:1308.4028[nucl-ex]))	
GB12/GZ12:	included recent CLAS π^-p d σ /d Ω	[W. Chen <i>et al,</i> Phys Rev C 86 , 015206 (2012)]	$\mathbf{>}$	H
CM12:	CM parameterization for $T_{\pi N}$	[R. Workman <i>et al,</i> Phys Rev C 86 , 015202 (2012)]		Γ
SN11/SK11:	included recent GRAAL $\pi^-p \otimes \pi^0n$	Σ	-	
	<i>LEPS</i> π ⁰ p dσ/dΩ	[R. Workman et al, Phys Rev C 85, 025201 (2012)]		
M =	= (Born + A)(1 + $iT_{\pi N}$) + $BT_{\pi N}$ + (C + i	$D)(ImT_{\pi N} - T_{\pi N} ^2)$		
SP09:	included recent CLAS π⁺n d o /dΩ	[M. Dugger <i>et al,</i> Phys Rev C 79 , 065206 (2009)]		
M :	= (Born + α_R)(1 + iT _{πN}) + $\alpha_R T_{\pi N}$ + hig	her terms		1
			_	

Solution	Energy Limit (MeV)	χ^2/N_{Data}	N _{Data}	
DU13	2700	2.23	27,265	_
GB12	2700	2.09	26,179	<hr/>
CM12	2700	2.01	25,814	
SN11	2700	2.08	25,553	
SP09	2700	2.05	24,912	
FA06	3000	2.18	25,524	
SM02	2000	2.01	17,571	/
SM95	2000	2.37	13,415	

- The overall SAID χ² has remained stable against the growing database, which has increased by a factor of 2 since 1995.
- Most of this increase coming from photon-tagging facilities.



Proton Multipoles from DU13 & CM12

[R. Workman et al, Phys Rev C 86, 015202 (2012); M. Dugger et al, arXiv:1308.4028 [nucl-ex]



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<u>CLAS</u> Data Impact for <u>Proton</u> $S = 0 \& I = \frac{1}{2}$ Couplings [M. Dugger et al, arXiv:1308.4028 [nucl-ex]



 $\gamma \pi \rightarrow \pi \rho$ Experiment

- The existing $\gamma n \rightarrow \pi^- p$ database contains mainly differential cross sections (17%) of which are from **polarized** measurements.)
- Many of these are old **bremsstrahlung** measurements with limited **angular** (θ = 40 - 140°) coverage and large **energy** binning (E_{γ} = 100 - 200 MeV.) In several cases, the **systematic** uncertainties have not been given.

- At lower energies (E_{γ} < 700 MeV,) there are data sets for the inverse π^- photoproduction reaction: $\pi^-p \rightarrow \gamma n$. This process is free from complications associated with a deuteron target.
- However, the disadvantage of using $\pi^- p \rightarrow \gamma n$ is the large background because of the **5** to **500** times larger cross section for $\pi^- p \rightarrow \pi^0 n \rightarrow \gamma \gamma n$.







World Neutral and Charged PionPR Data

[SAID: http://gwdac.phys.gwu.edu/]

Future exp activity will fill empty spots specifically for n-target.



FSI and $\gamma d \rightarrow \pi \bar{\rho} p \longrightarrow \gamma \pi \rightarrow \pi \bar{\rho}$ [V. Tarasov, A. Kudryavtsev, W. Briscoe, H. Gao, & IS, Phys Rev C 84, 035203 (2011)]

• FSI plays a critical role in the state-of-the-art analysis of $\gamma n \rightarrow \pi N$ data. • For $\gamma n \rightarrow \pi^- p$ the effect: 5% - 60%. It depends on (E, θ).



$$R = (d\sigma/d\Omega_{\pi p})/(d\sigma^{IA}/d\Omega_{\pi p}) \qquad \Longrightarrow \qquad \frac{d\sigma}{d\Omega}(\gamma n)$$





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MAMI-B for $\gamma n \rightarrow \pi^- p$ around the Δ

[W.J. Briscoe et al, Phys Rev C 86, 065207 (2012)]

• **MAMI-B** data for $\gamma n \rightarrow \pi^- p$ (including **FSI** corrections) and previous hadronic data for $\pi^- p \rightarrow n\gamma$ appear to **agree well**.





Neutron Multipoles from SAID GB12 & SN11

[W. Chen et al, Phys Rev C 86, 015206 (2012); R. Workman et al, Phys Rev C 85, 025201 (2012)]]



CLAS Data Impact for Neutron $S = 0 \ll I = \frac{1}{2}$ Couplings [W. Chen *et al*, Phys Rev C **86**, 015206 (2012)]

• BnGa13 and SAID GB12 used the same (almost) data to fit them while BnGa13 has several new Ad Hoc resonances.





 $N(2700) 13/2^+$

10/6/2013

Status of Non-strange Resonances

J. Beringer *et al* (PDG) Phys Rev D **86**, 010001 (2012)

• More than half of states have poor evidence.

Most of states need more work to do.

<u>GW SAI</u>	<u>) Co</u>	ion Status as seen I = 1/2						l /2	 Most 		
Particle J^P	overa	Statu II πN	γN	Νη	Νσ	Nω	ΛK	Σŀ	Κ Νρ	$\Delta \pi$	• Whe
$N = 1/2^+$	****										1 2/2
$N(1440) 1/2^+$	****	****	****		***				*	***	1 = 3/2
$N(1520) 3/2^{-}$	****	****	****	***					***	***	
$N(1535) 1/2^{-}$	****	****	****	****					**	*	Particle J
$N(1650) 1/2^{-}$	****	****	***	***			***	**	**	***	A(1232) 3
$N(1675) 5/2^{-1}$	****	****	***	*			*		*	***	$\Delta(1202) 0$
$N(1680) 5/2^+$	****	****	****	*	**				***	***	$\Delta(1600) 0$
N(1685) ??	*										$\Delta(1020)$ 1/ $\Delta(1700)$ 3
$N(1700) 3/2^{-1}$	***	***	**	*			*	*	*	***	$\Delta(1750) 1$
$N(1710) 1/2^+$	***	***	***	***		**	***	**	*	**	$\Delta(1750) 1$
$N(1720) 3/2^+$	****	****	***	***			**	**	**	*	$\Delta(1900) 1$ $\Delta(1005) 5$
$N(1860) 5/2^+$	**	**							*	*	$\Delta(1905) 5$
$N(1875) 3/2^{-1}$	***	*	***			**	***	**		***	$\Delta(1910) 1$
$N(1880) 1/2^+$	**	*	*		**		*				$\Delta(1920) 3_{0}$
$N(1895) 1/2^{-1}$	**	*	**	**			**	*			$\Delta(1930) 5/$
$N(1900) 3/2^+$	***	**	***	**		**	***	**	*	**	$\Delta(1940)$ 3
$N(1990) 7/2^+$	**	**	**					*			$\Delta(1950)$ 7
$N(2000) 5/2^+$	**	*	**	**			**	*	**		$\Delta(2000) 5_{0}$
$N(2040) 3/2^+$	*										$\Delta(2150) 1_{0}$
$N(2060) 5/2^{-1}$	**	**	**	*				**			$\Delta(2200) 7/$
$N(2100) 1/2^+$	*										$\Delta(2300) 9/$
$N(2150) 3/2^{-}$	**	**	**	26	*		**			**	$\Delta(2350) 5/$
$N(2190) 7/2^{-}$	****	****	***	201	N ala ala ata		**		*		$\Delta(2390) 7_{0}$
$N(2220) 9/2^+$	****	****		11 *	***						$\Delta(2400) 9$
$N(2250) 9/2^{-}$	****	****		5 ³	* * *						$\Delta(2420)$ 1
$N(2600) 11/2^{-1}$	\mathbf{A}	***		7 :	**						$\Delta(2750)$ 13

3 *

• Most of QCD models predict more states than observed.

• Where are missing resonances?

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			*	***	l = 3/2					S	tatus	as se	en in	—			:his
			***	***			Statu	s								1	7
			**	*	Particle J^F	overal	$1 \pi N$	γN	$N\eta$	Nσ	$N\omega$	ΛK	ΣK	Nρ	$\Delta \pi$		p 2
	***	**	**	***	$\Delta(1232) 3/2$	2+ ****	****	****	F							Ť	S NC
	*		*	***	$\Delta(1600) 3/2$	2+ ***	***	***	0					*	***		ide
			***	***	$\Delta(1620) 1/2$	2 ****	****	***		r				***	***		ę
					$\Delta(1700) 3/2$	2 ****	****	****		ь				**	***		2
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	*				$\Delta(1920) 3/2$	2' ***	***	**				n	***		**		Ģ
	**	*			$\Delta(1930) 5/2$	2 ***	***										전
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	**			**	$\Delta(2350) 5/2$	2 *	*			_	d		N				5
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					$\Delta(2400) 9/2$	2- **	**	7 :	****			n		tat	toc		ate
					$\Delta(2420)$ 11	/2+ ****	****		Ja ala ala			_		la	les		<u></u>
					$\Delta(2750)$ 13	/2 **	**	3	ዮ ጥ ጥ								Ě
					A(2050) 15		**	7	**								
<u> </u>					5(2500) 15		**	_ <u> </u>	*							L	-
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GWM Program







Jefferson Lab

Helicity Amplitudes : $S_{11}(1535)$



Constituent Quark Model Predictions

[E. Santopinto and M. Giannini, Phys Rev C 86, 065202 (2012)]

∆(1232)3/2⁺



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N(1520)3/2-





Igor Strakovsky



Timelike $\gamma^* \mathcal{N} \rightarrow \Delta$ form factors and Δ Dalitz decay

[G. Ramalho and M.T. Pena, Phys Rev D 85, 113014 (2012)]



Inverse Pion Electroproduction (IPE)

• IPE is the only process which allows the determination of EM nucleon & pion formfactors in the intervals

 $0 < k^2 < 4 M^2 = 3.53 \text{ GeV}^2$ $0 < k^2 < 4 m_{\pi} = 0.08 \text{ GeV}^2$

which are kinematically unattainable from e^+e^- initial states.



• IPE $\pi^- p \rightarrow e^+ e^- n$ measurements will significantly complement the current electroproduction $\gamma^* N \rightarrow \pi N$ study for the evolution of baryon properties with increasing momentum transfer by investigation of the case for the *time-like virtual photon*.

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Phenomenology for Baryon Resonances







