

EMMI RRTF - TOP3 discussion session

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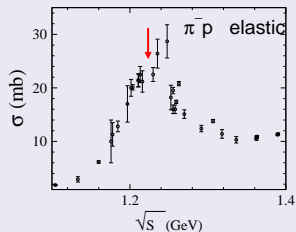


Baryon resonance analysis technique

50's...60's: discovery of p , n , e^+ , e^- , π
what about new particles?

try scattering experiments

$\pi^- N$ elastic by E. Fermi:

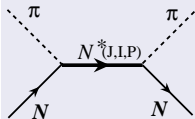


- peak in the elastic cross section
- is it a new particle?
- if yes, which properties ?
- why it appears as a broad peak ?

to identify new a particle
one needs to know production amplitude

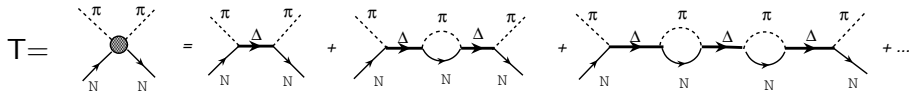
Baryon resonance analysis technique

Resonance production: reaction amplitude



- production vertices are unknown
- if particle is created it should propagate
- scattering amplitude $T \sim (s - m_R^2 + i\epsilon)^{-1}$,

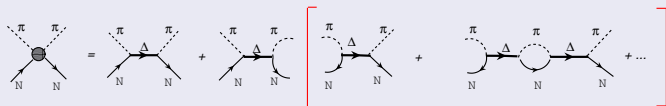
Actual amplitude could be more complicated



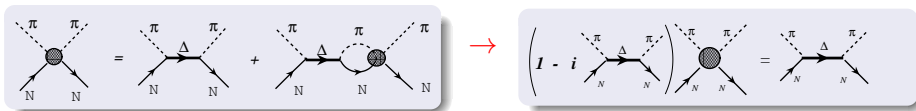
How to solve the problem ?

Baryon resonance analysis technique

$\pi^- N$ scattering amplitude :



infinite sum of all diagrams in the brackets - it is a full amplitude!

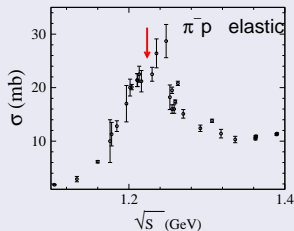


$$\begin{aligned}
 \text{Diagram} &= \frac{\text{Diagram with self-energy loop on } \Delta}{1 - i} = \frac{\frac{\Gamma_{\pi NN^*} \Gamma_{\pi NN^*}}{s - m_R^2}}{1 - i \frac{\Gamma_{\pi NN^*} \Gamma_{\pi NN^*}}{s - m_R^2}} = \frac{\Gamma_{\pi NN^*} \Gamma_{\pi NN^*}}{s - m_R^2 - i \Gamma_{\pi NN^*} \Gamma_{\pi NN^*}}
 \end{aligned}$$

Breit-Wigner form: finite width for non-vanishing $\Gamma_{\pi NN^*}$

Baryon resonance analysis technique

$\pi^- N$ elastic by E. Fermi



Scattering amplitude with Δ

$$T = \frac{\Gamma_{\pi NN^*} \Gamma_{\pi NN^*}}{E^2 - m^2 + i\Sigma(\sqrt{s})}$$

pole on the second Riemann sheet

$$E_{\text{pole}} = \sqrt{m^2 - i\Sigma(E_{\text{pole}})}$$

$$\sigma_{\text{tot}}(\sqrt{s}) \sim \frac{F(s)}{(s - m^2)^2 + \Sigma^2(\sqrt{s})}$$

particle with a short lifetime (resonance) \rightarrow peak in the σ_{tot}

... however: still not enough to identify the peak as a resonance excitation

Baryon resonance analysis: isospin of the $\Delta(1232)$.

How to extract isospin of $\Delta(1232)$ from experimental data?

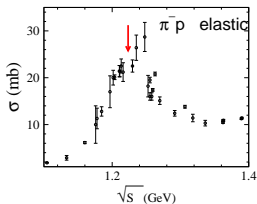
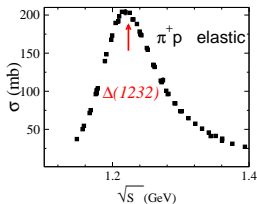
$$|\pi^+ p\rangle = \left| \frac{3}{2} \right\rangle$$
$$|\pi^- p\rangle = \sqrt{\frac{1}{3}} \left| \frac{3}{2} \right\rangle + \sqrt{\frac{2}{3}} \left| \frac{1}{2} \right\rangle$$

$$T(\pi^+ p \rightarrow \pi^+ p) = \langle \frac{3}{2} | T | \frac{3}{2} \rangle = T^{\frac{3}{2}}$$

$$T(\pi^- p \rightarrow \pi^- p) = \frac{1}{3}(T^{\frac{3}{2}} + 2T^{\frac{1}{2}})$$

if $I_{\Delta(1232)} = \frac{3}{2}$ then $T^{\frac{1}{2}} = 0$ and

$$\frac{\sigma_{tot}^{el}(\pi^+ p \rightarrow \pi^+ p)}{\sigma_{tot}^{el}(\pi^- p \rightarrow \pi^- p)} = 9$$



Experiment

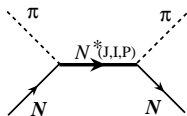
$$\frac{\sigma_{\pi^+ p}}{\sigma_{\pi^- p}} \Big|_{\sqrt{s}=1.23} = \frac{200 \text{ mb}}{22 \text{ mb}} \approx 9$$

$\Delta(1232)$: Isospin $I = \frac{3}{2}$

Isospin decomposition: $T^{\frac{3}{2}} = T(\pi^+ p \rightarrow \pi^+ p)$ and

$$T^{\frac{1}{2}} = \frac{1}{2}[3T(\pi^- p \rightarrow \pi^- p) - T(\pi^+ p \rightarrow \pi^+ p)]$$

Baryon resonance analysis technique: Spin of the $\Delta(1232)$.



$\Delta(1232)$: **Spin:** $\pi N \rightarrow \Delta(1232) \rightarrow \pi N$ -reaction

:
angular distribution is defined by total spin J ;
how to extract ?

$\pi N \rightarrow \pi N$ differential cross section:

$$\frac{d\sigma}{d\cos(\theta)} \sim \sum_{J\lambda\lambda'} |d_{\lambda\lambda'}^J(\theta) T_{\lambda\lambda'}^J(\sqrt{s})|^2$$

There are two independent amplitudes

$$\left\{ T_{\frac{1}{2}, \frac{1}{2}}^J, T_{\frac{1}{2}, -\frac{1}{2}}^J \right\} \leftrightarrow \{ T^{J+}, T^{J-} \}$$

$$T_{\frac{1}{2}, \frac{1}{2}}^J = T^{J+} + T^{J-}$$

$$T_{\frac{1}{2}, -\frac{1}{2}}^J = T^{J+} - T^{J-}$$

-rewrite in terms of T^{J-} and T^{J+} parity conserved amplitudes

$$\frac{d\sigma}{d\cos(\theta)} \sim \left[\left(d_{\frac{1}{2}, \frac{1}{2}}^J(\theta) \right)^2 + \left(d_{\frac{1}{2}, -\frac{1}{2}}^J(\theta) \right)^2 \right] |T^{J\pm}(\sqrt{s})|^2$$

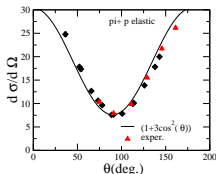
Spin of the $\Delta(1232)$.

assuming $J = \frac{3}{2}$:

$$\frac{d\sigma}{d\cos(\theta)} \sim \left[\left(d_{\frac{1}{2}, \frac{1}{2}}^{\frac{3}{2}}(\theta) \right)^2 + \left(d_{\frac{1}{2}, -\frac{1}{2}}^{\frac{3}{2}}(\theta) \right)^2 \right] T^{\frac{3}{2}\pm}(\sqrt{s})$$

$$\frac{d\sigma}{d\cos(\theta)} \sim \text{const} \times (1 + 3\cos^2(\theta))$$

$$\sqrt{s} = 1.232 \text{ GeV}$$



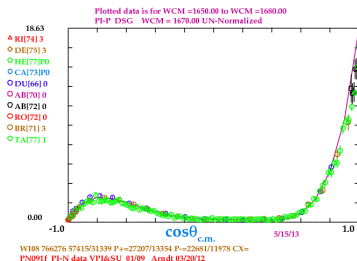
Finally: $\Delta(1232)$ Isospin $I = \frac{3}{2}$, Spin $J = \frac{3}{2}$

Inverse task: obtain contributions of given spin from exp. data

$$T_{\lambda_N \lambda'_N}^J(\sqrt{s}) = \int d\theta d_{\lambda \lambda'}^J(\theta) \frac{d\sigma}{d\cos(\theta)}$$

Higher energies, many states

$\pi^- p$ elastic at higher energies



spectrum could be reach: many states + non-resonant background

- write down scattering amplitudes as sum

$$T(\sqrt{s}, \theta) = \sum_i \left(\frac{\Gamma_i}{s - m_{R_i}^2 + im_i \Gamma_i} + T^{non-pole} \right) d_{\lambda\lambda'}^{J_i}(\theta)$$

- calculate exp. observables
- compares to the data, fix parameters, extract N^* parameters (poles)

TWO MAIN INGREDIENTS:

- Scattering amplitude (theory).
- Experiment

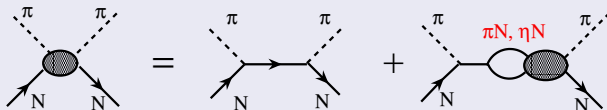
DEFINE THE SCATTERING AMPLITUDE

- parametrizations: $T = (\text{Breit-Wigner} + \text{non-pole terms})$, Chew-Mandelstam formalism, + unitarization, dispersion relations: less model dependence
- dynamical models: calculate T using an effective field theory: Lippman-Schwinger/Bethe-Salpeter equations, coupled-channel, Lagrangian input - coupling constants.

Construct the model for the $\pi N \rightarrow \pi/\eta N$ transitions

take rescattering in the πN and ηN channels into account

The interaction potentials $V_{\pi N \rightarrow \eta N}$ and $V_{\pi N \rightarrow \pi N}$ enter to



Coupled-channel problem for $\pi N \rightarrow \pi N$ scattering:

$$\begin{aligned} T_{\pi N \rightarrow \pi N} &= V_{\pi N \rightarrow \pi N} + \int d^4 p V_{\pi N \rightarrow \pi N} G_{\pi N}(p) T_{\pi N \rightarrow \pi N} \\ &\quad + \int d^4 p V_{\pi N \rightarrow \eta N} G_{\eta N}(p) T_{\eta N \rightarrow \pi N} \\ T_{\eta N \rightarrow \pi N} &= V_{\eta N \rightarrow \pi N} + \int d^4 p V_{\eta N \rightarrow \pi N} G_{\pi N}(p) T_{\pi N \rightarrow \pi N} \\ &\quad + \int d^4 p V_{\eta N \rightarrow \eta N} G_{\eta N}(p) T_{\eta N \rightarrow \pi N} \end{aligned}$$

coupled channel problem:

There are four equations for the $T_{\pi N, \pi N}$ $T_{\pi N, \eta N}$ $T_{\eta N, \pi N}$ and $T_{\eta N, \eta N}$. They can be written in the matrix form :

$$\begin{pmatrix} T_{\pi N, \pi N} & T_{\pi N, \eta N} \\ T_{\eta N, \pi N} & T_{\eta N, \eta N} \end{pmatrix} = \begin{pmatrix} V_{\pi N, \pi N} & V_{\pi N, \eta N} \\ V_{\eta N, \pi N} & V_{\eta N, \eta N} \end{pmatrix} + \int d^4 p \begin{pmatrix} V_{\pi N, \pi N} & V_{\pi N, \eta N} \\ V_{\eta N, \pi N} & V_{\eta N, \eta N} \end{pmatrix} \begin{pmatrix} G_{\pi N} & 0 \\ 0 & G_{\eta N} \end{pmatrix} \begin{pmatrix} T_{\pi N, \pi N} & T_{\pi N, \eta N} \\ T_{\eta N, \pi N} & T_{\eta N, \eta N} \end{pmatrix}$$

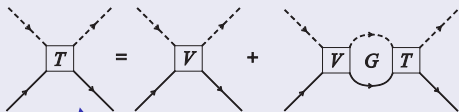
or more compact:

$$[\hat{T}] = [\hat{V}] + i \int \frac{d^4 k}{(2\pi)^4} [\hat{V}] G_{\text{mB}}^{\hat{}} [\hat{T}]$$

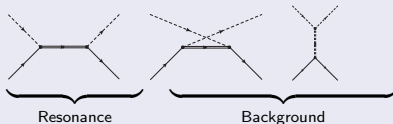
- can be easily generalized for any number of channels
- PWD + K-matrix approx \rightarrow algebraic matrix equations

Bethe-Salpeter in K -matrix: dynamical model: based on eff. L_{mBB}

T-matrix



Interaction terms V_{ij}



multidimensional T-matrix

$$T = \begin{pmatrix} T_{\gamma N, \gamma N} & T_{\gamma N, \pi N} & T_{\gamma N, K\Lambda} & \cdots \\ T_{\pi N, \gamma N} & T_{\pi N, \pi N} & T_{\pi N, K\Lambda} & \cdots \\ T_{K\Lambda, \gamma N} & T_{K\Lambda, \pi N} & T_{K\Lambda, K\Lambda} & \cdots \\ \cdots & \cdots & \cdots & \cdots \end{pmatrix}$$

How many channels?

$\gamma N \rightarrow \gamma N$	$\pi N \rightarrow \pi N$
$\gamma N \rightarrow \pi N$	$\pi N \rightarrow 2\pi N$
$\gamma N \rightarrow \eta N$	$\pi N \rightarrow \eta N$
$\gamma N \rightarrow \omega N$	$\pi N \rightarrow \omega N$
$\gamma N \rightarrow K\Lambda$	$\pi N \rightarrow K\Lambda$
$\gamma N \rightarrow K\Sigma$	$\pi N \rightarrow K\Sigma$

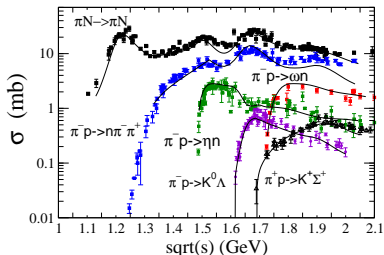
Partial wave version of optical theorem

constraints on partial wave cross sections

$$\text{Im} T_{\pi N \rightarrow \pi N}^{JP} = \frac{k^2}{4\pi} (\sigma_{\pi N \rightarrow \pi N}^{JP} + \sigma_{\pi N \rightarrow 2\pi N}^{JP} + \sigma_{\pi N \rightarrow \eta N}^{JP} + \sigma_{\pi N \rightarrow \omega N}^{JP} + \sigma_{\pi N \rightarrow K\Lambda}^{JP} + \sigma_{\pi N \rightarrow K\Sigma}^{JP} + \dots)$$

all reaction data are linked

→ need for coupled-channel unitary analysis



$$T = \begin{pmatrix} T_{\gamma N, \gamma N} & T_{\gamma N, \pi N} & T_{\gamma N, K\Lambda} & \dots \\ T_{\pi N, \gamma N} & T_{\pi N, \pi N} & T_{\pi N, K\Lambda} & \dots \\ T_{K\Lambda, \gamma N} & T_{K\Lambda, \pi N} & T_{K\Lambda, K\Lambda} & \dots \\ \dots & \dots & \dots & \dots \end{pmatrix}$$

← Giessen Model vs
experimental data

pion beam experiments vs. photoproduction

"Missing states" and photoproduction:

Pion beam experiment

- Most of our knowledge about N^* properties comes from single channel partial wave analysis (PWA) of πN elastic scattering
- Isospin decomposition is straightforward

surprisingly ... but pion experiments were stopped in 70s

why ? main reasons:

- problems in identification of neutral particles in $\pi^- p \rightarrow \eta n$, $K\Lambda$, ωN , ϕN etc scattering: the most of experiments with charged particles in final states

Main argument against pion-beams

- we know everything from hadronic experiments!
- "missing resonances" are weakly coupled to πN : can only be seen in photoproduction !

$N^* \rightarrow \pi N$ decays

PDG 2012 main contribution from the analysis of the πN elastic scattering

N^*	$L_{2I} 2J$	Overall	$\text{Br}(N^* \rightarrow \pi N)$
$N(939)$	P_{11}	★★★★	—
$N(1440)$	P_{11}	★★★★	55...67 %
$N(1520)$	D_{13}	★★★★	55...65 %
$N(1535)$	S_{11}	★★★★	35...55 %
$N(1650)$	S_{11}	★★★★	50...90 %
$N(1675)$	D_{15}	★★★★	35...45 %
$N(1680)$	F_{15}	★★★★	65...70 %
$N(1700)$	D_{13}	★★★	8 ... 17 %
$N(1710)$	P_{11}	★★★	5 ... 20 %
$N(1720)$	P_{13}	★★★★	9...14 %
$N(1870)$	D_{13}	★★★	10 ... 22 %
$N(1900)$	P_{13}	★★★	10 %
$N(2000)$	F_{15}	★★	9 %

- for $\text{Br}(N^* \rightarrow \pi N) < 20\%$ no general agreement between different analyses !

- indication for the existence is smaller for higher masses (more degrees of freedom at higher energies, many open channels)

- resonance/background separation is difficult

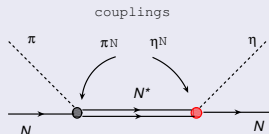
- πN elastic scattering: increase exp. resolution

Another possibility: inelastic reactions

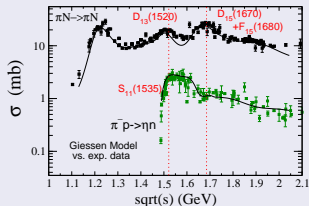
inelastic $\pi N \rightarrow \eta N, \omega N, \rho N \dots$ etc scattering

My argument:

- resonance contribution to e.g. η -production: $\frac{d\sigma}{d\Omega} \sim g_{\pi NN^*}^2 g_{\eta NN^*}^2$
- signals from N^* with small πN coupling **can be visible** provided $g_{\eta NN^*}^2$ is large



- signals from N^* with small πN coupling **less screening** by contributions from N^* with large πN coupling: **no clean signal** from $D_{13}(1520)$, $D_{15}(1680)$, $F_{15}(1680)$ in $\pi N \rightarrow \eta N$



Short summary :

Main argument against pion-beams

- ~~"missing resonances"~~ are weakly coupled to πN : can only be seen in photoproduction **NO!**
- The pion-induced inelastic reactions **provide great possibility** to study N^* spectra !

But

- May be we know everything from old πN experiments ?

ωN -meson dynamics in nuclear medium

- in-medium modification of omega-mesons in nuclei: (HADES, CB-ELSA/TAPS etc)
- broadening of the omega meson in nuclear medium but no mass shift
- strong absorption in the nuclear medium

but

- large collisional broadening
- what about chiral symmetry restoration ?

Microscopical model is needed

Building block: ωN scattering amplitude

ωN scattering length

- $\bar{a} = -0.026 + i0.28$ fm, Giessen (coupled-channel) NPA780 187
- $\bar{a} = -0.44 + i0.20$ fm, Lutz, et al(coupled-channel, low partial waves) NPA706:431
- $\bar{a} = +1.60 + i0.30$ fm, Kling, Weise (single channel) NPA630:299

Common feature of above analysis:

- constrained by the $\pi N \rightarrow \omega N$ experimental data
- agrees on the value of the imaginary part of the scattering lengths

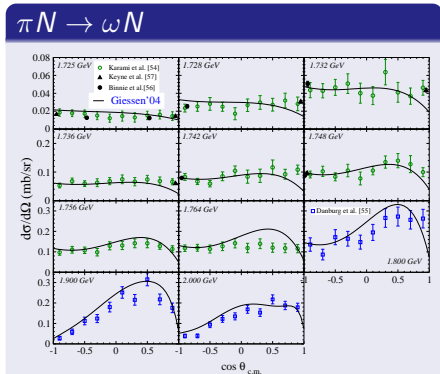
low density theorem: $i0.28$ corresponds to ≈ 60 MeV broadening
but too small to explain the strong absorption of ω in medium

- theory: take in-medium corrections into account
- experiment: is everything clear with old $\pi N \rightarrow \omega N$ data?

Giessen model. Results for the $(\pi, \gamma)N \rightarrow \omega N$ reactions

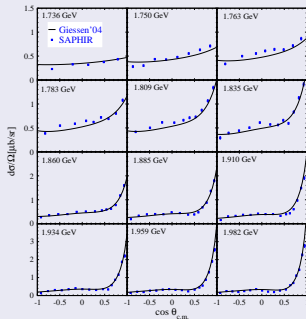
ωN : coupled channel analysis Shklyar et al PRC 71:055206:

Aim: extract resonance coupling to ωN



few measurements, low statistic

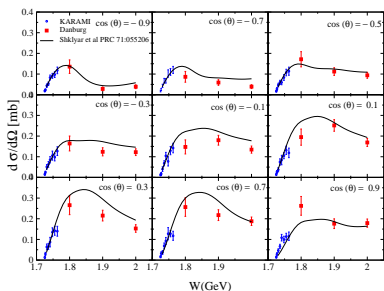
$\gamma N \rightarrow \omega N$



strong t -channel pion exchange
screens other reaction
mechanisms

- $W=1.72$ to 1.76 GeV: H. Karami, et al NPB154 503 (1979) : 80 datapoints threshold region
- $W=1.8$ to 2.1 GeV: J.S. Danburg, PR2, 2564(1970) from $\pi^+ D \rightarrow \pi^+ \pi^- \pi^0 p(p)$: 41 datapoints Fermi-motion, final state interaction!

Shklyar et al,
PRC 71:055206,2005



Difficulties:

- ωN has three helicities: need ω -polarization measurements
- Karami data - close to threshold
- region 1.76...2.0 GeV is almost empty - standard PWA not possible
- no polarization measurements
- Problem: N^* extraction ...

$$(\pi/\gamma)N \rightarrow \omega N$$

Summary of $(\pi/\gamma)N \rightarrow \omega N$ reactions

- $\gamma p \rightarrow \omega p$: strong t -channel background \rightarrow other reaction mechanisms are shadowed: hard to see any resonance contributions
- $\pi N \rightarrow \omega N$: almost NO data in the region 1.76...2.0 GeV - standard PWA not possible
- contributions from many groups: Lutz, Wolf, Friman, Titov, Sibirtsev, Zhao, Shklyar, Mosel, Penner - no general conclusion on N^* contributions

NEED $\pi^- p \rightarrow \omega p$ measurements in order to

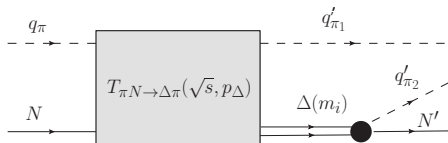
- get information on N^* couplings to ωN - fill white pages in PDG
- construct microscopical model of ω -dynamics in nuclear medium; explain large collisional broadening

$\pi N \rightarrow 2\pi N$ reactions

investigate cascade reactions e.g. $N^* \rightarrow \pi N^* \rightarrow \pi\pi N$ etc. :
multiparticle production

Analysis of $\pi N \rightarrow 2\pi N$: Manley, Arndt, Goradia, Teplitz
PRD**30**,(1984) 904.

- isobar approximation $\pi N \rightarrow 2\pi N$ via $\sigma N, \rho N, \pi \Delta \rightarrow 2\pi$

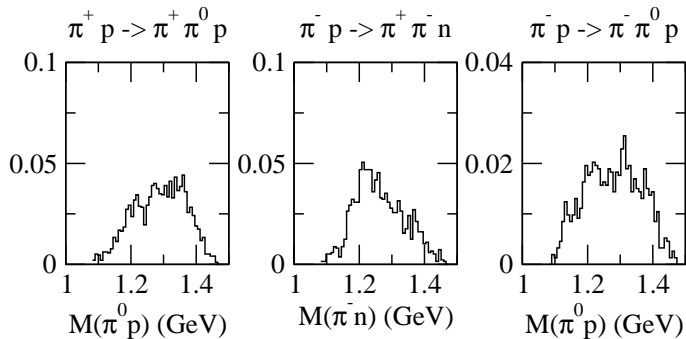


$$T_{\pi N \rightarrow 2\pi N}^{JP} = T_{\pi N \rightarrow \Delta\pi}^{JP}(\sqrt{s}) S_\Delta(p_\Delta, m_\Delta) \Gamma_{\Delta\pi N}(q'_{\pi_2}, N')$$

Potential problems:

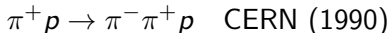
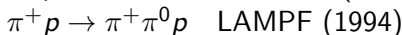
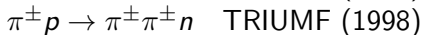
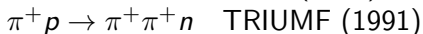
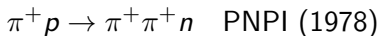
- no three-body unitarity
- no dependence on isobar mass (momentum)
- poor database based on 240000 events from old bubble-chamber experiments $W = 1.2 \dots 2 \text{ GeV}$: ≈ 9000 events per energy/angular (θ, ϕ) bin for $\pi^- p \rightarrow \pi^+ \pi^- n, \pi^- p \rightarrow \pi^0 \pi^- p, \pi^+ p \rightarrow \pi^0 \pi^+ p, \pi^+ p \rightarrow \pi^+ \pi^+ n$
 $\approx 2000 \dots 3000$ events per energy bin for each reaction

$$\pi N \rightarrow 2\pi N$$

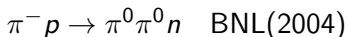


New data came later (most of them are total X-sections)
(I.Strakovsky, GWU) but **not suited for $N^* \rightarrow \rho N$**

- **W=1221 to 1356 MeV**



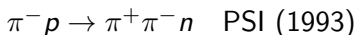
- **W=1213 to 1527 MeV**



- **W=1257 to 1302 MeV**



- **W=1300 to 1302 MeV**



Why pion beam experiment for $2\pi N$ is production important?

Pion-induced reactions:

$$\begin{aligned}\pi^- p &\rightarrow \pi^+ \pi^- n, \pi^- p \rightarrow \pi^0 \pi^- p, \\ \pi^+ p &\rightarrow \pi^0 \pi^+ p, \pi^+ p \rightarrow \pi^+ \pi^+ n \\ \pi^- p &\rightarrow \pi^0 \pi^0 n\end{aligned}$$

- Isospin decomposition :
4 independent isospin amplitudes (in isobar approximation)
- optical theorem $Im T_{\pi N \rightarrow \pi N}^{JP} = \frac{k^2}{4\pi} (\sigma_{\pi N \rightarrow \pi N}^{JP} + \sigma_{\pi N \rightarrow 2\pi N}^{JP} + \dots)$

Photon-induced reactions:

$$\begin{aligned}\gamma p &\rightarrow \pi^+ \pi^- n, \gamma p \rightarrow \pi^0 \pi^- p, \\ \gamma p &\rightarrow \pi^0 \pi^0 p\end{aligned}$$

- No isospin decomposition is possible (separation between $l = \frac{1}{2}$ and $\frac{3}{2}$ states is more difficult)
- difficulties with the gauge invariance
- need input from hadronic reactions

$N(1520) D_{13}$ state

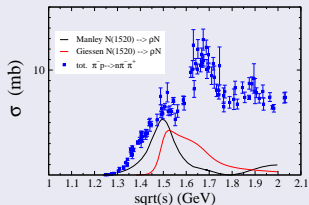
Manley et al: PRD(1984)

$$M_R = 1.52\text{MeV}$$

$$\Gamma_{\text{tot}} = 120\text{MeV}$$

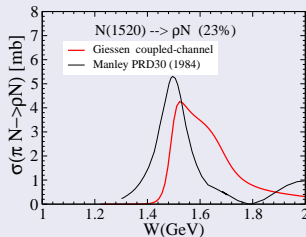
strong $N(1520) \rightarrow 2\pi N$

$$\text{Br}(\rho N) \approx 20\%$$



- need diff. x-sections of $\pi^- p \rightarrow \pi\pi N$

Giessen Model (CC): $\pi N \rightarrow \rho N$



- Giessen : overlapping of spectral functions of $N^*(1520)$ and ρ -meson: non-symmetric
- Giessen: no effect below 1.4 GeV
- Manley: no ρ -spectral function: should be updated

Summary of the $\pi N \rightarrow 2\pi N$ reactions

- important for understanding ρ -meson dynamics and resonance couplings
- could solve many puzzles in non-strange baryon spectroscopy: origin and properties of the $P_{11}(1440)$, $P_{11}(1710)$, $D_{13}(1520)$ etc.

Theory

- analysis of Manley et. al. should be updated!

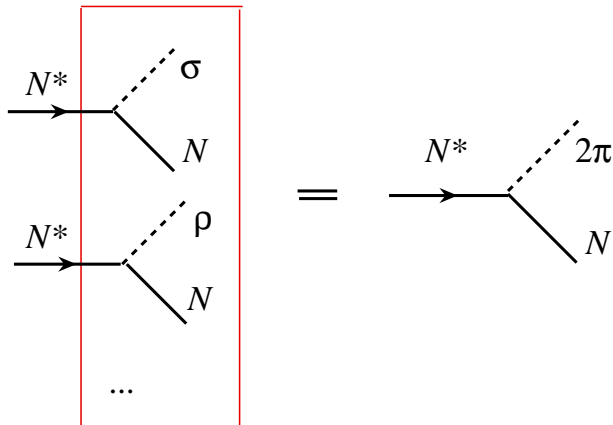
Experiment

- need for new measurements $\pi N \rightarrow 2\pi N$ in region 1.2...2.GeV

Next step: improve description of the $2\pi N$ channel

so far: N^* decay into 'generic' 2π channel

- take $2\pi N$ inelastic flux into account
- $N^* \rightarrow 2\pi N$ couplings constrained by $\sigma_{\pi N \rightarrow 2\pi N}^{JJ}$



Roper resonance $N(1440)$ properties:

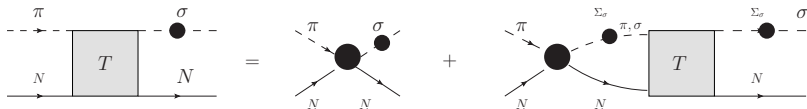
- Manley & Saleski PRD30 904, $Br(\Delta\pi) = 22\%$ $Br(\sigma N) = 9\%$
- Vrana et al PRPL328, $Br(\Delta\pi) = 16\%$ $Br(\sigma N) = 12\%$
- Sarantsev et al PLB659,94, $Br(\Delta\pi) = 17\%$ $Br(\sigma N) = 21\%$

- Julich Model: PRC62: pion exchange is responsible for a large amount of attraction: $P(1440)$ is dynamically generated
- Crystal Ball PRL91(2003): PWA of the $2\pi^0$ -subsystem: σ -meson production via pion exchange is small
- Crystal Ball PRL69(2004): measurement of the $\pi N \rightarrow 2\pi^0 N$ -reaction: no direct evidence for a strong σN subchannel

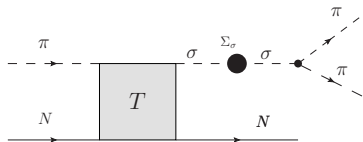
$\pi N \rightarrow 2\pi$ channel in the first resonance energy region

BSE in the isobar approximation:

system of coupled-channel integral equations



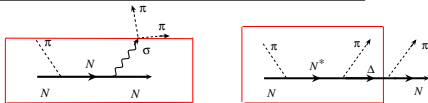
$\pi N \rightarrow 2\pi N$ amplitude from BSE



+ $\pi N \rightarrow \rho N, \pi N \rightarrow \pi \Delta$ etc

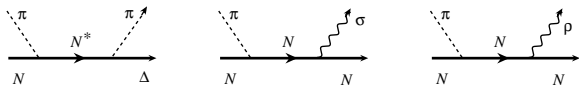
Next step: improve description of the $2\pi N$ channel

$\pi N \rightarrow 2\pi N$ reaction via ρN , $\pi\Delta$ channels

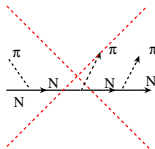


Assumptions

- decays $N^* \rightarrow \rho N$, σN , $\pi\Delta$ drive the $\pi N \rightarrow 2\pi N$ channel



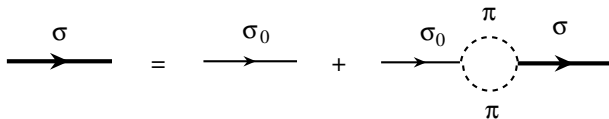
- two-step diagrams are neglected



σ -meson dynamics

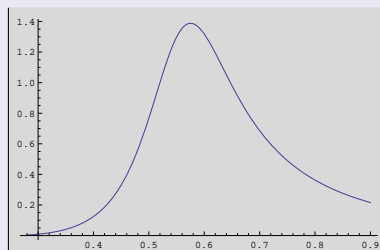
propagator of the σ -meson

$$D(s) = \frac{1}{s - m_\sigma^2 + i\Sigma_\sigma(s)}$$

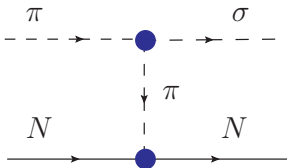


$$A_\sigma(s) = \frac{1}{\pi} \frac{\Sigma_\sigma(s)}{(s - m_\sigma^2)^2 + \Sigma_\sigma(s)^2}$$

spectral function A_σ



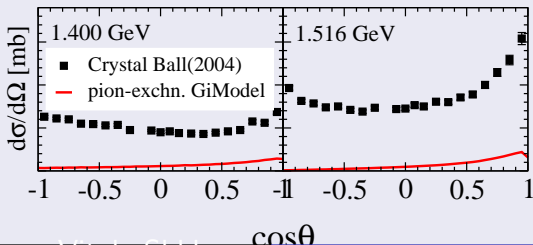
t-channel pion exchange: σN how large?



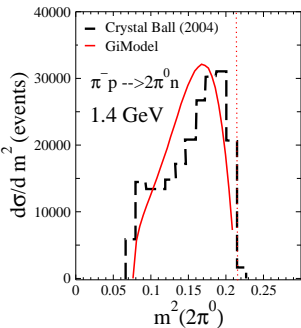
- coupling constants are well fixed
- $g_{\pi NN} = 13$, $g_{\sigma\pi\pi} = 2$ correspond to $m_{\sigma}^0 = 600\text{MeV}$, $\Gamma_{\sigma\pi\pi} = 600\text{MeV}$
- contribution from the t-channel diagram is well fixed

- shed light on the σ -meson dynamics
- background mechanism in $\pi N \rightarrow 2\pi N$ reaction

$\pi^- p \rightarrow \pi^0 \pi^0 n$: t-channel pion exchange: very small !



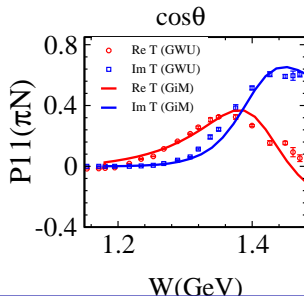
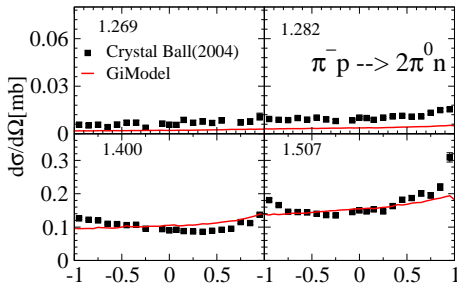
Giessen Model vs. Crystal Ball data



- good description of the $\pi^- p \rightarrow 2\pi^0 n$ data
- three-body unitarity is maintained

$$\text{Im } T_{\pi N}^{11} = \frac{k^2}{4\pi} (\sigma_{\pi N}^{11} + \sigma_{2\pi^0 N}^{11})$$

Roper resonance



GIModel for $\pi^- p \rightarrow \pi^0 \pi^0 n$ reaction

- model space is extended to include σN , $\pi \Delta$, and ρN channels
- t -channel pion exchange in σN channel is very weak - underestimate the data.
- do not rule out dynamical pole; however if it exists the contribution to the production cross section should be small
- calculation with a genuine Roper resonance: nice description of the CB-measurements