

Observational Probes of the Neutron Star Equation of State with Hyperons, Sexaquarks, and Quark Matter

Avraham's talk: deeply bound H-dibaryon 

Davood Rafiei Karkevandi

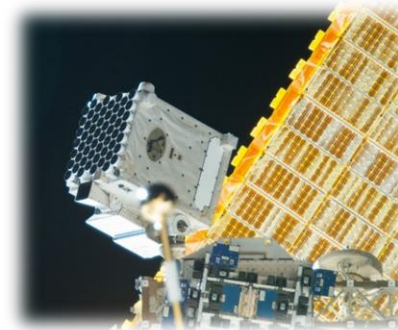
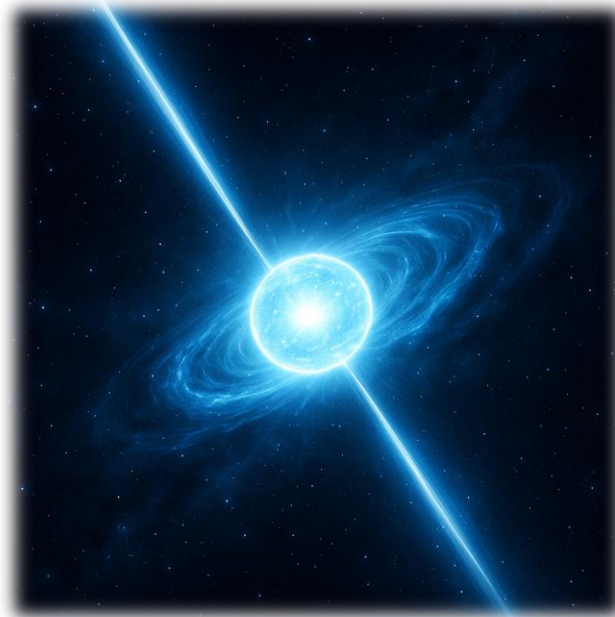
**Institute of Theoretical Physics,
Wroclaw University, Poland**

EMMI Workshop

5th Workshop Anti-Matter, Hyper-Matter and Exotica

Production at the LHC,

Nov 10-14, 2025, Salerno, Italy



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Observational Probes of the Neutron Star Equation of State with Hyperons, Bosonic Dark Matter, and Quark Matter

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arXiv:2402.18686v2

Accepted for publication in
Astronomy & Astrophysics

Probing Strange Dark Matter through f -mode Oscillations of Neutron Stars with Hyperons and Quark Matter

Mahboubeh Shahrbafl^{1,*}, Prashant Thakur^{2,†}, and Davood Rafiei Karkevandi^{1‡}

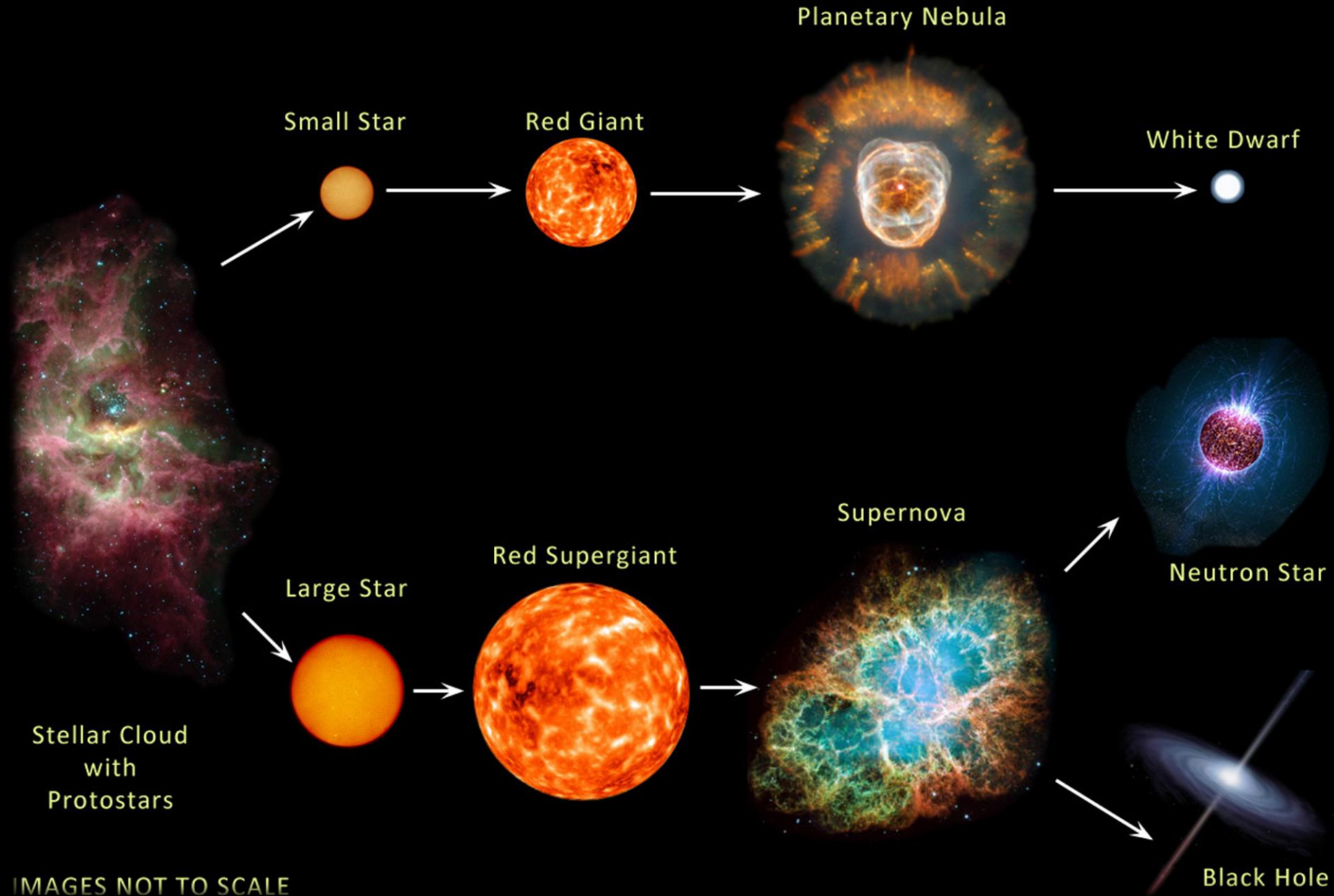
¹*Institute of Theoretical Physics, University of Wrocław, 50-204 Wrocław, Poland and*

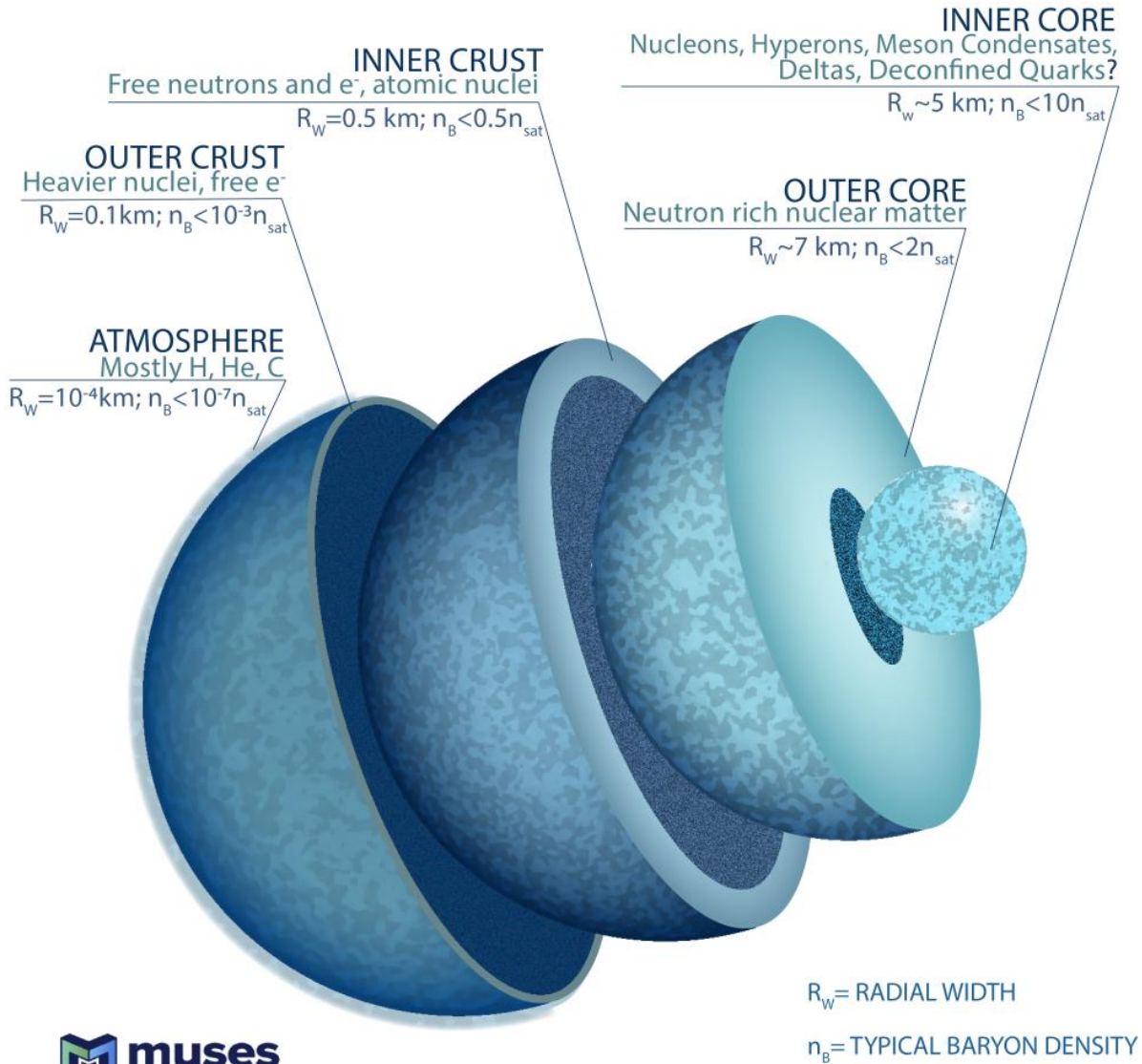
²*Department of Physics, Yonsei University, Seoul, 03722, South Korea*

arXiv:2510.08115

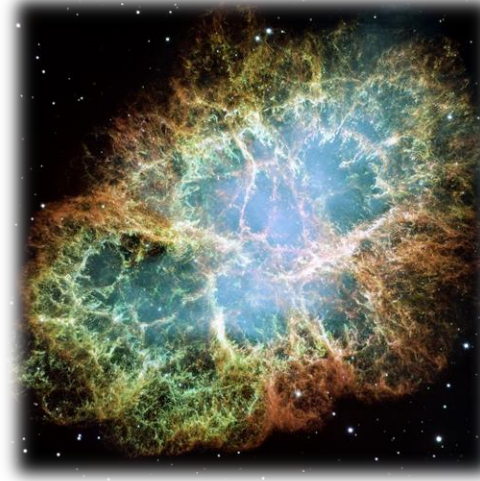


EVOLUTION OF STARS





Crab Pulsar

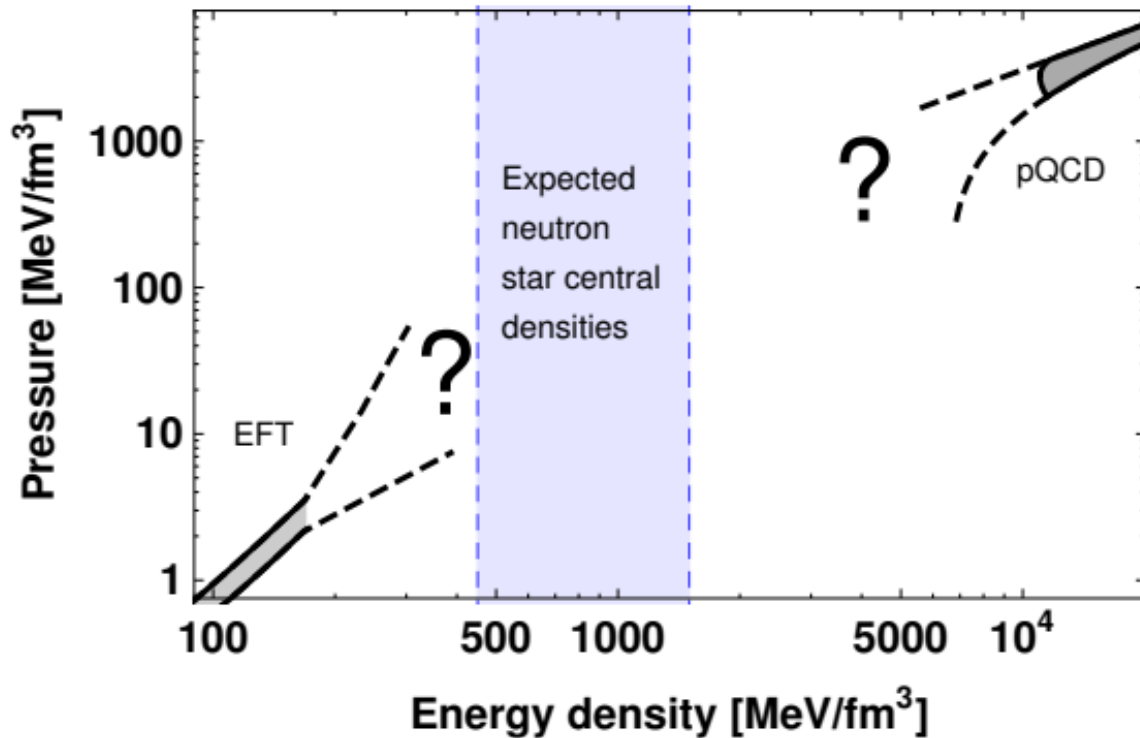


Mass $\sim 1.4 - 2$ solar mass (M_{\odot})

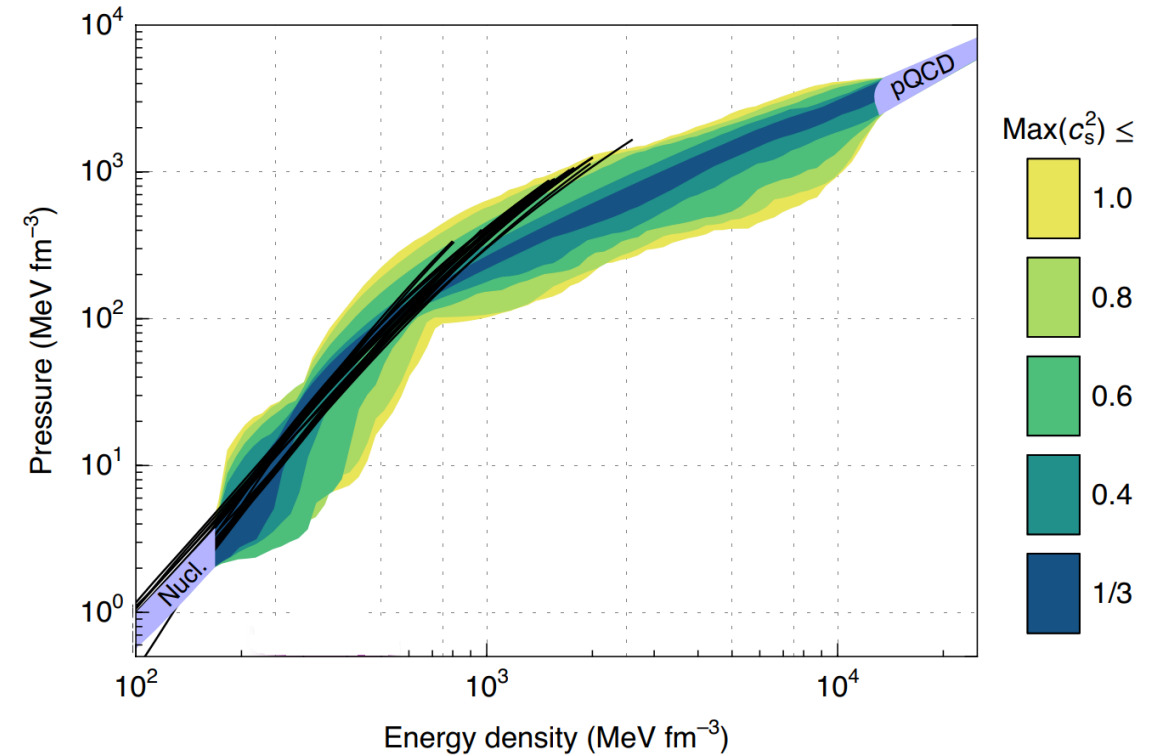
Radius $\sim 10 - 13$ km

**Neutron stars as a natural
laboratory for high density matter**

Equation of state of ultra-high density matter

$$P = P(\varepsilon)$$


Matti Järvinen, *Eur. Phys. J. C* (2022) 82:282



Eemeli Annala, *Nature Phys.* 16 (2020) 9, 907-910

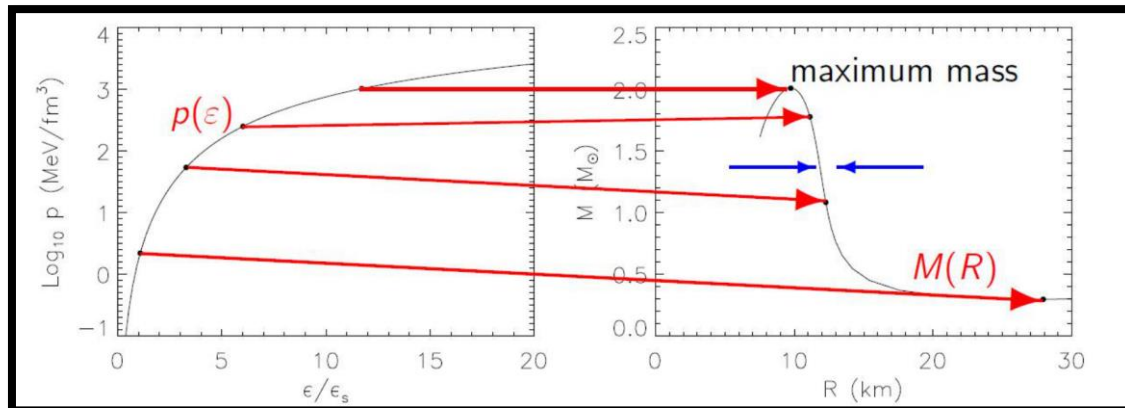
Tolman-Oppenheimer-Volkof (TOV) equations

$$\frac{dP(r)}{dr} = -\frac{GM(r)\varepsilon(r)}{c^2 r^2} \left[1 + \frac{P(r)}{\varepsilon(r)} \right] \left[1 + \frac{4\pi r^3 P(r)}{M(r)c^2} \right] \left[1 - \frac{2GM(r)}{c^2 r} \right]^{-1}$$

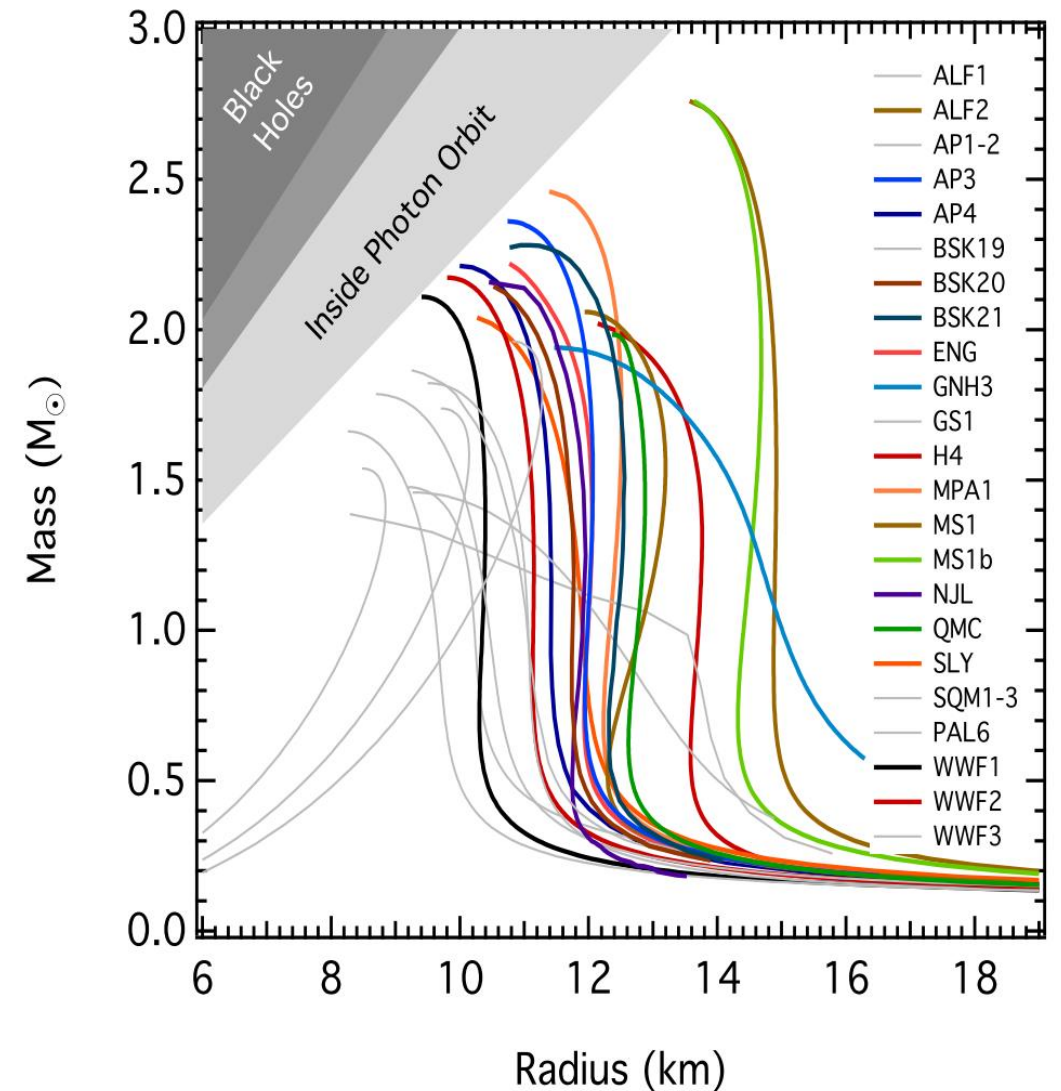
R. C. Tolman, *Phys. Rev.* **55**, 364 (1939).

$$\frac{dM(r)}{dr} = \frac{4\pi r^2 \varepsilon(r)}{c^2}$$

J. R. Oppenheimer and G. M. Volkov, *Phys. Rev.* **55**, 374 (1939).

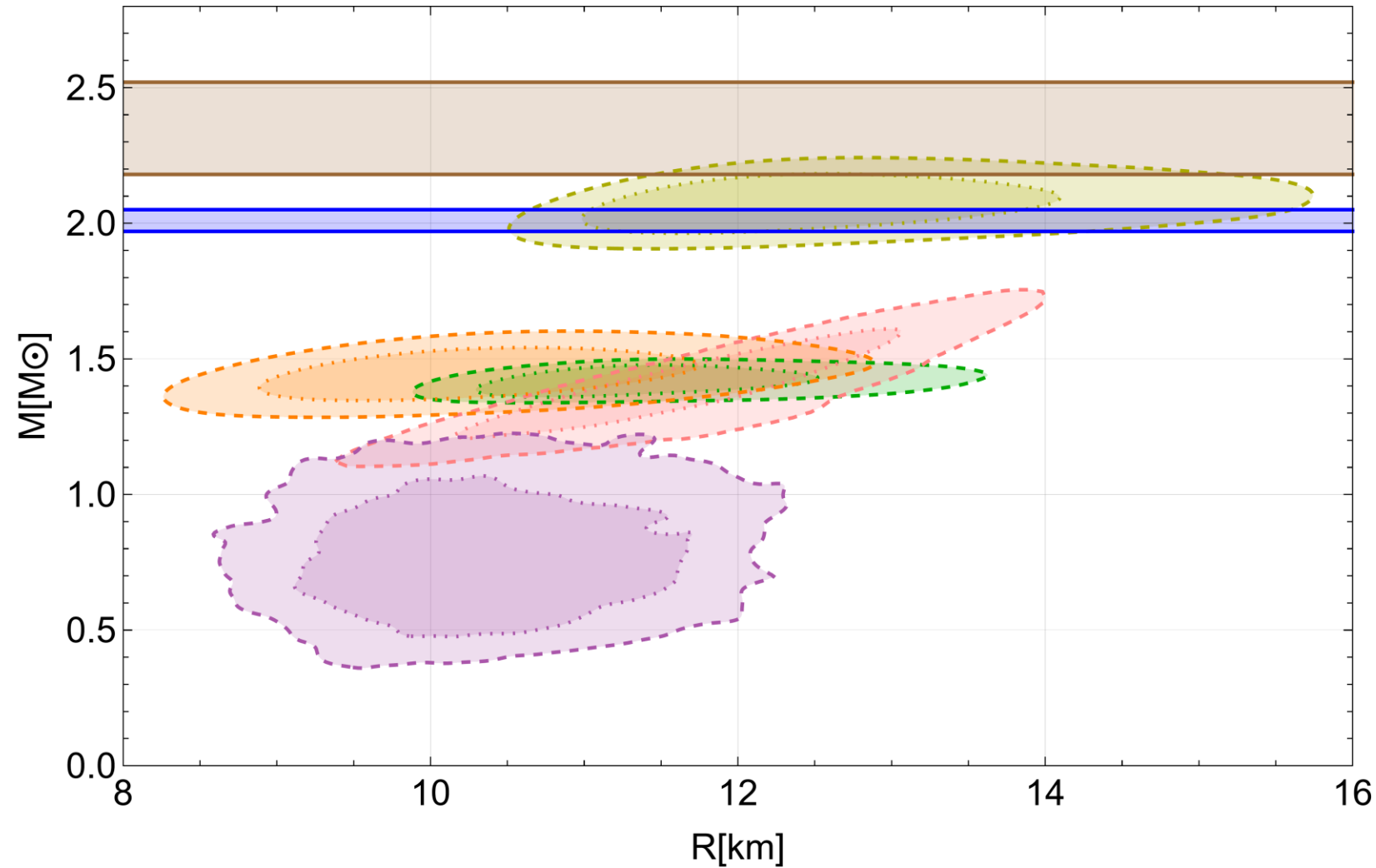


James Latimer, *EPL Web of Conferences* **109**, 07004 (2016)



Feryal Özel & Paulo Freire
Ann.Rev.Astron.Astrophys. **54** (2016) 401-440

Mass-Radius observational constraints



NICER X-ray Telescope



Keck Optical Telescopes



Green Bank Radio Telescope



PSR J0348+0432, John Antoniadis et al, [Science, 26 April 2013, Vol: 340, Issue: 6131](#)

PSR J0030+0451, S. Vinciguerra et al., [Astrophys. J. 961, 62 \(2024\)](#) HESS J1731-347, Doroshenko et al, [Nature Astron., 6, 1444, 2022](#)

PSR J0740+6620, T. Salmi et al., [Astrophys. J. 974, 294 \(2024\)](#)

PSR J0437-4715, D. Choudhury et al., [Astrophys. J. Lett. 971, L20 \(2024\)](#)

PSR J0614-3329, Lucien Mauviard et al, [arXiv:2506.14883](#)

PSR J0952-0607, W. Romani et al, [ApJL 934 L17, 2022](#)



GW170817: Observation of Gravitational Waves from a Binary Neutron Star Inspiral

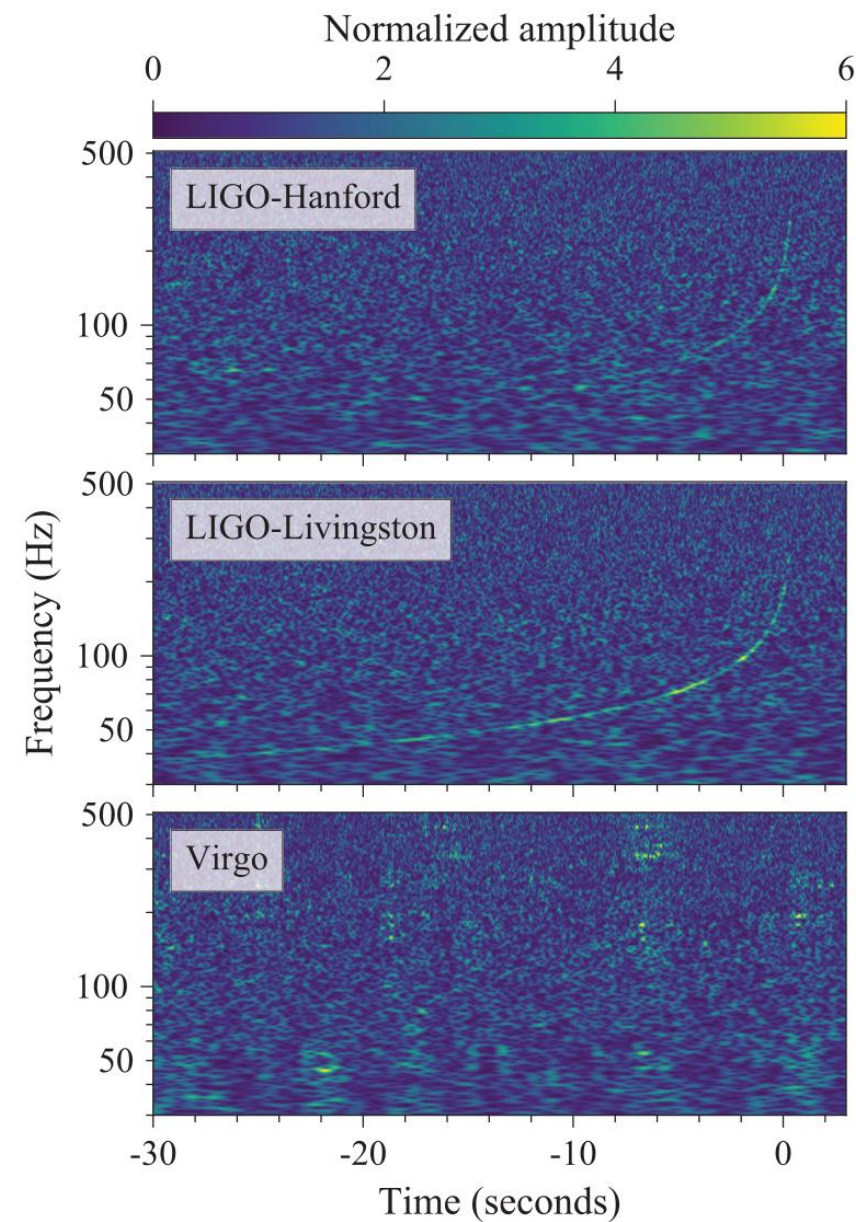
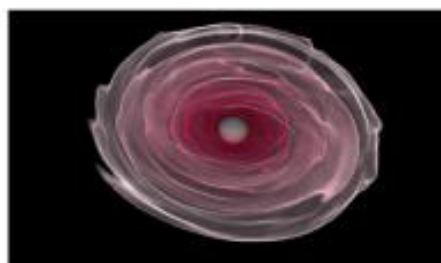
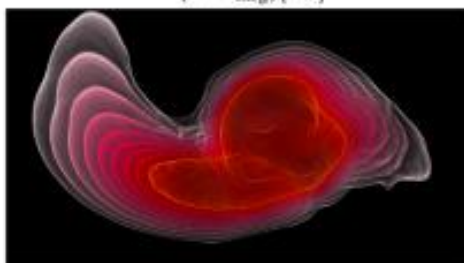
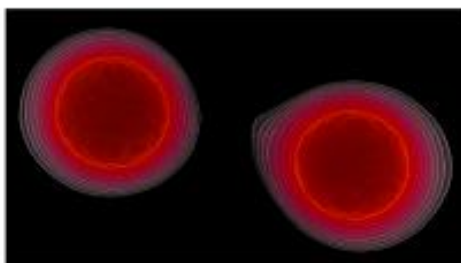
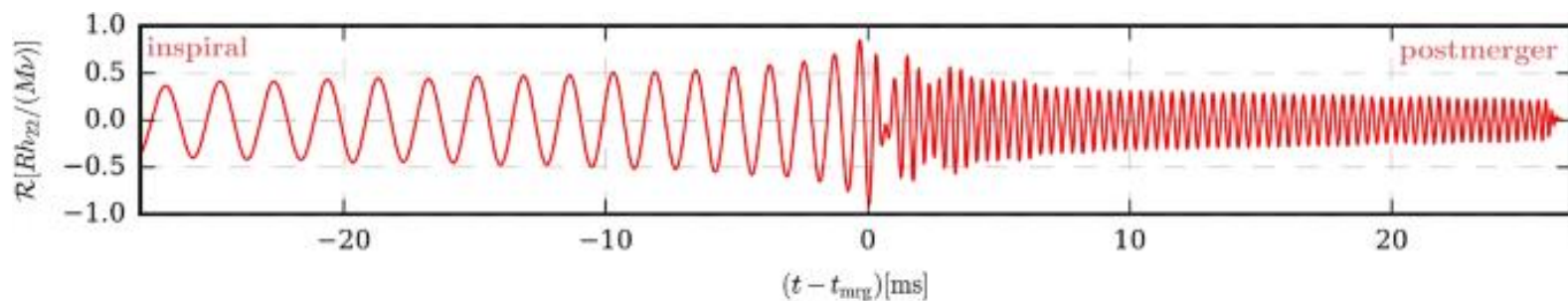
B. P. Abbott *et al.**

(LIGO Scientific Collaboration and Virgo Collaboration)

(Received 26 September 2017; revised manuscript received 2 October 2017; published 16 October 2017)



**LIGO-Virgo
Gravitational-Wave
detectors**



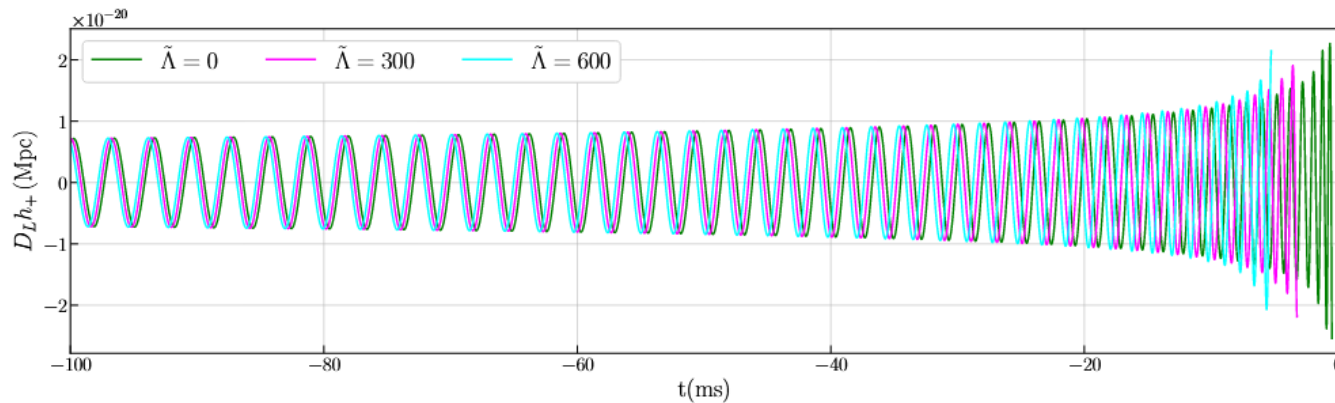
Deformation of the neutron star in the binary systems due to the tidal force

Dimensionless tidal deformability: $\Lambda = \frac{\lambda_t}{M^5} = \frac{2}{3} k_2 \left(\frac{R}{M} \right)^5$

R and M are the mass and radius of star

k_2 : Dimensionless tidal love number

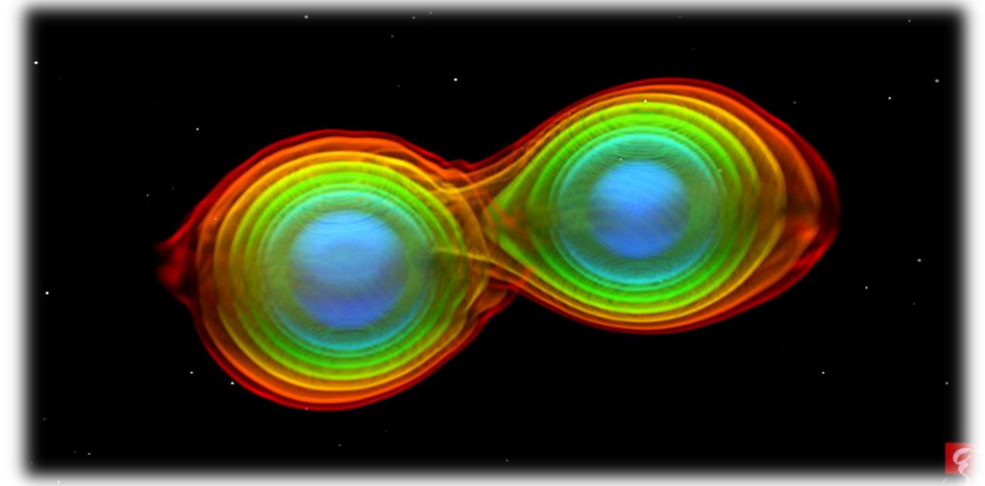
Tanja Hinderer, et al. *Phys.Rev.D*81:123016,2010



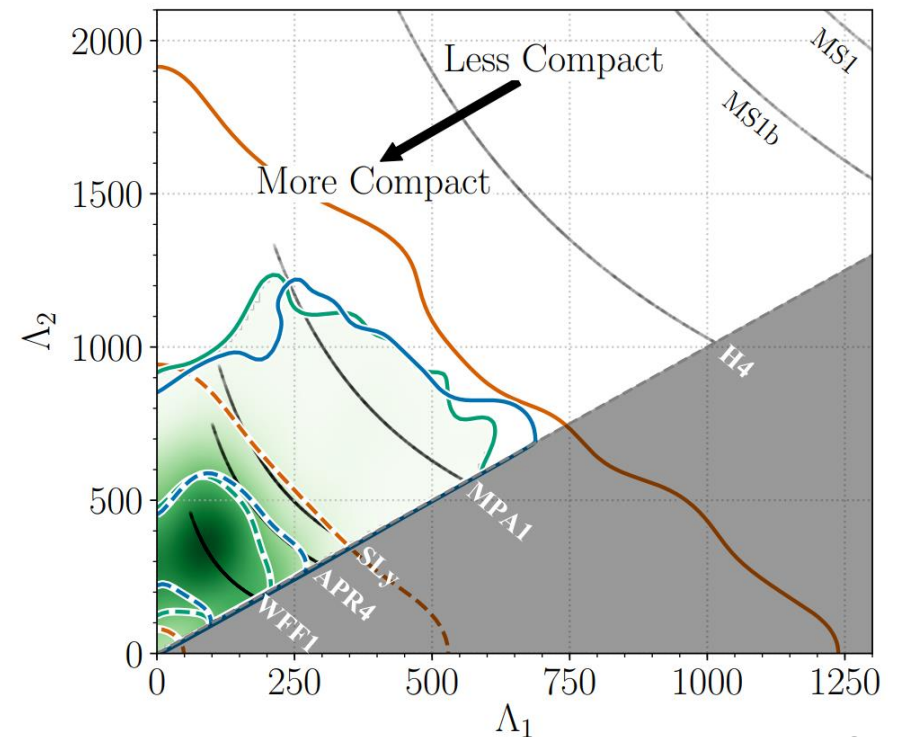
Katerina Chatziioannou, *Gen.Rel.Grav.* 52 (2020) 11, 109

Dimensionless tidal deformability GW170817
 $70 \leq \Lambda \leq 580$ for $M = 1.4 M_\odot$

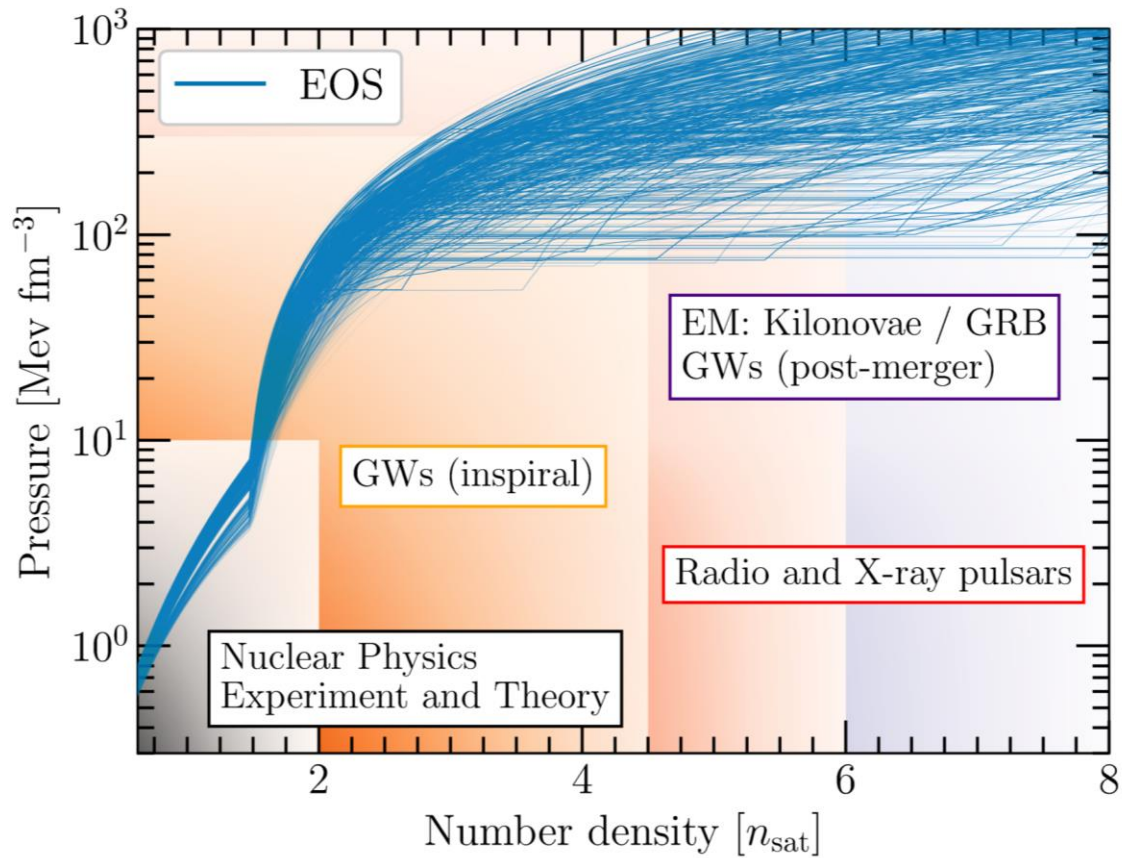
LIGO Scientific and Virgo Collaborations, *Phys.Rev.Lett.* 121 (2018) 16



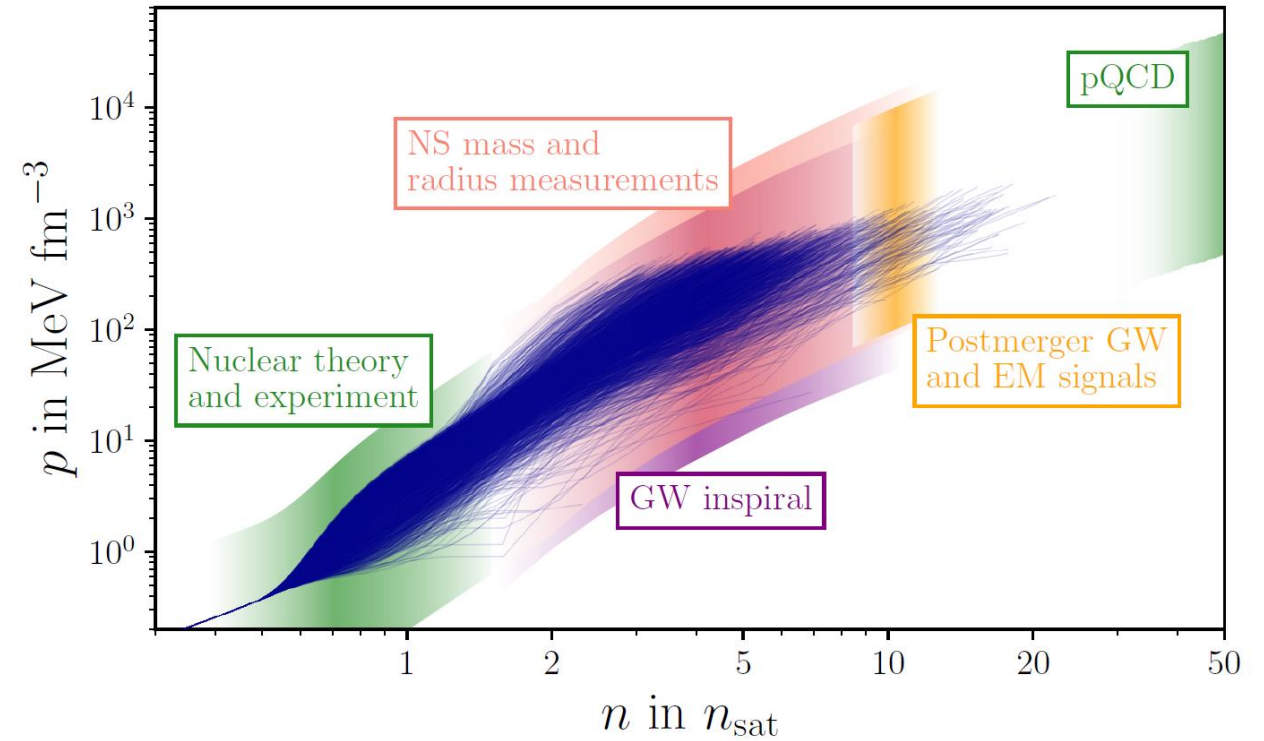
Tim Dietrich et al. Max Planck Institute



Constraining the EOS of high-density matter by multi-messenger observations of neutron stars

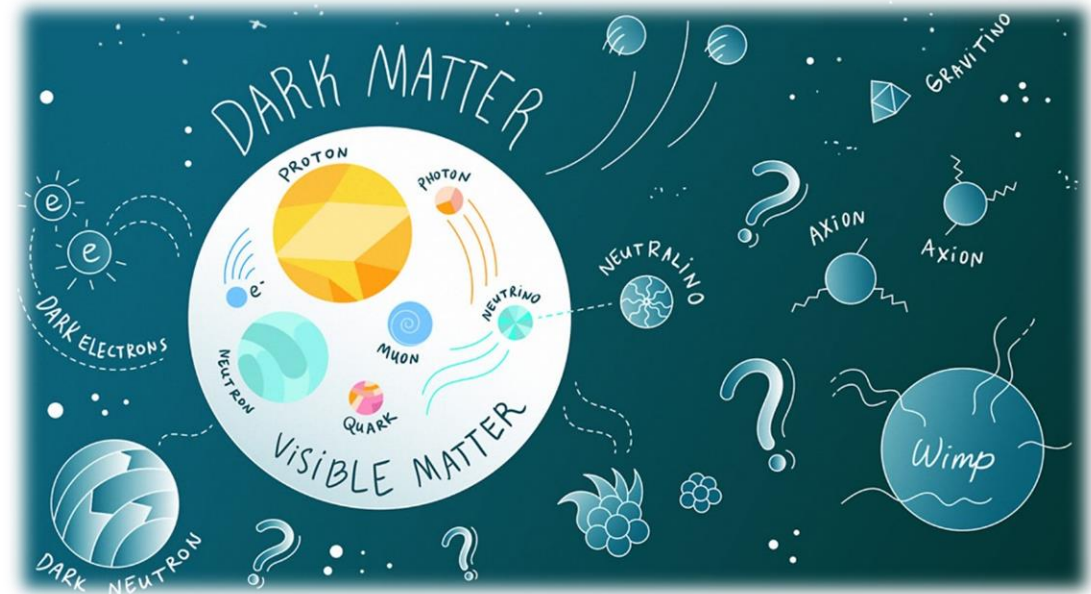
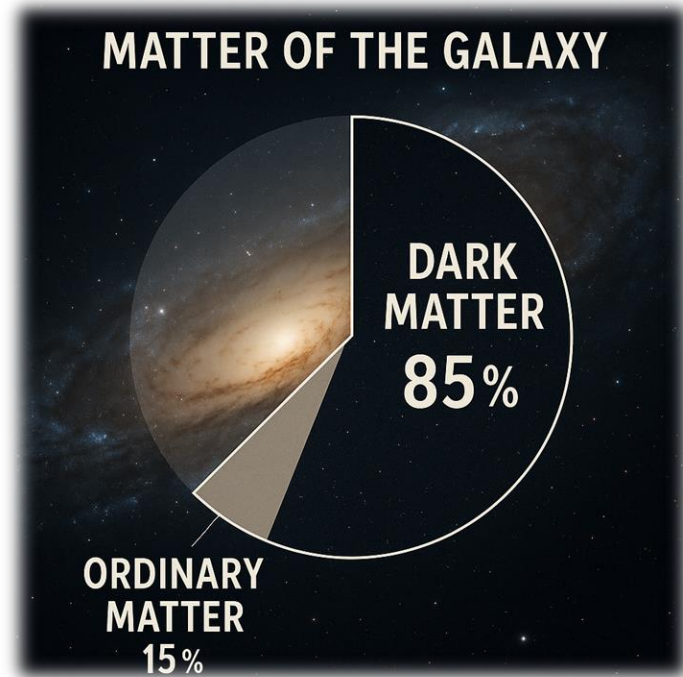


Peter T. H. Pang, et al. , *Nature Commun.* 14 (2023) 1, 8352



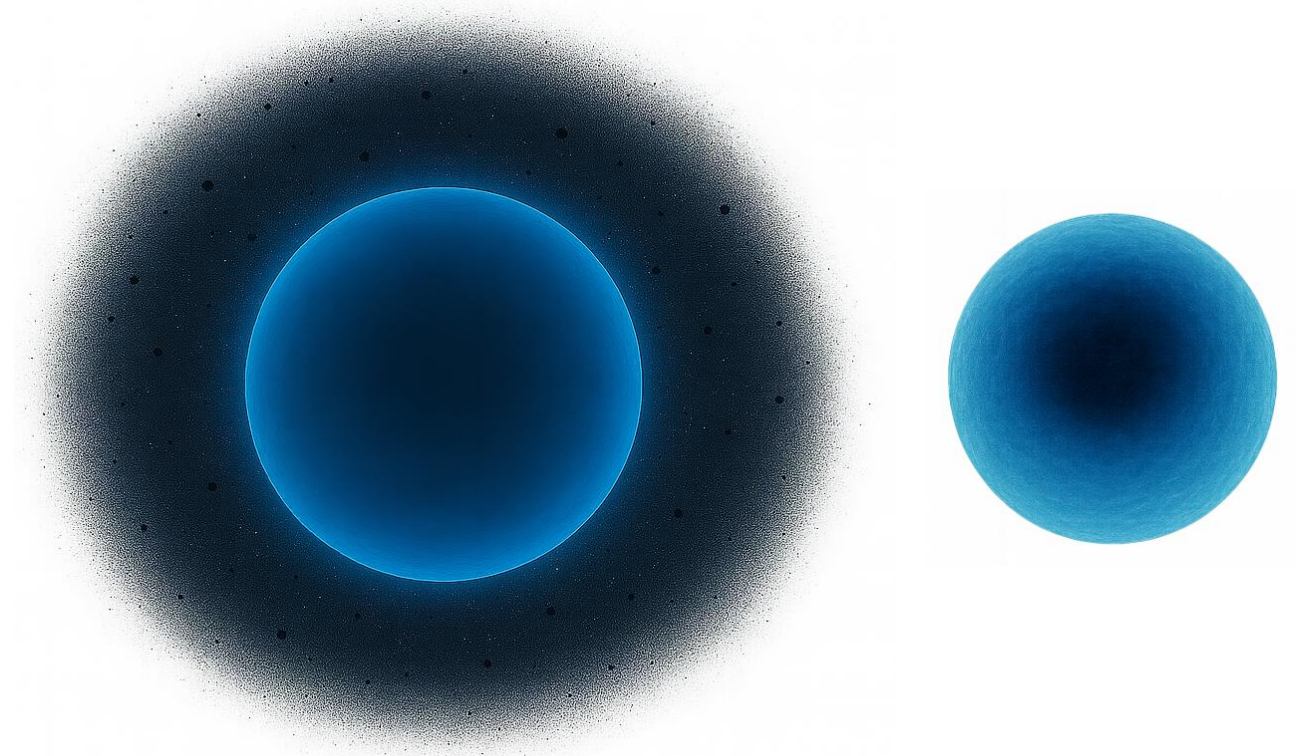
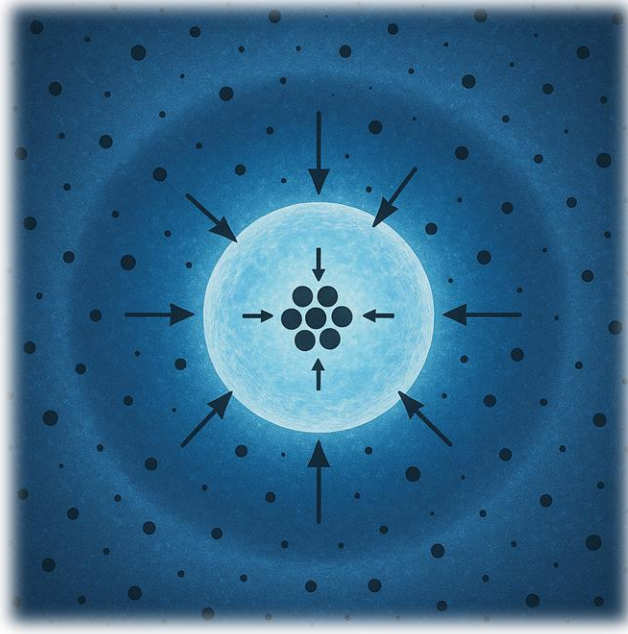
Hauke Koehn et al. *Phys.Rev.X* 15 (2025) 2, 021014

Dark matter constitutes about 85% of the matter in the universe, yet it remains a challenging puzzle in physics.



Main goal is to constrain the parameter space of the dark matter models and looking for signature

Neutron stars, due to their high densities and strong gravity, are promising astrophysical laboratories for probing the properties of dark matter.



Accretion of dark matter during different stages of stars' and neutron stars' life

Dark matter production in the neutron star matter or supernova explosions

Dark matter halo
around neutron star

Dark matter core
inside neutron star

D.R.K, S. Shakeri, V. Sagun, O. Ivanytskyi,

Bosonic dark matter in neutron stars and its effect on gravitational wave signal

Phys. Rev. D 105, 023001 (2022) [159 citations]

D.R.K, S. Shakeri, V. Sagun, O. Ivanytskyi,

Tidal deformability as a probe of dark matter in neutron stars

World Scientific pp. 3713-3731 (2023) [21 citations]

S. Shakeri, D.R.K

Bosonic Dark Matter in Light of the NICER Precise Mass-Radius Measurements

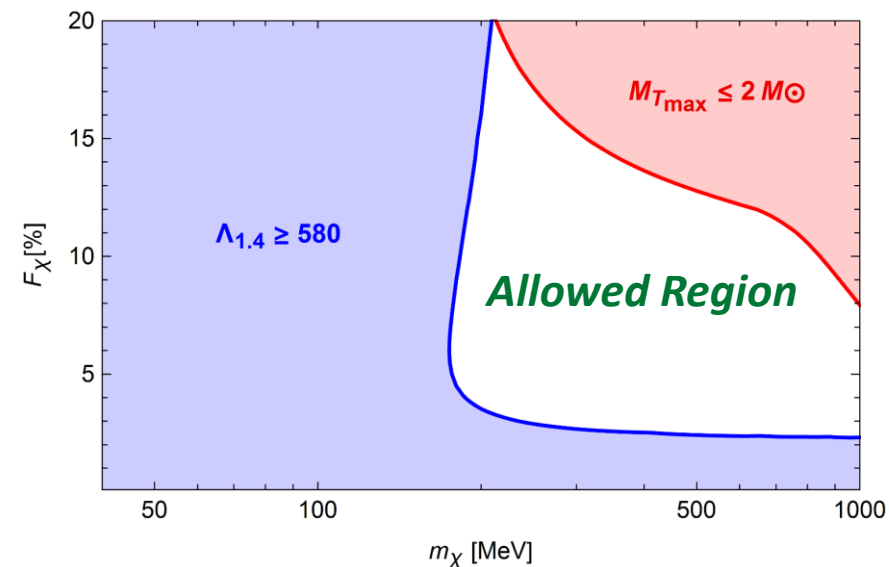
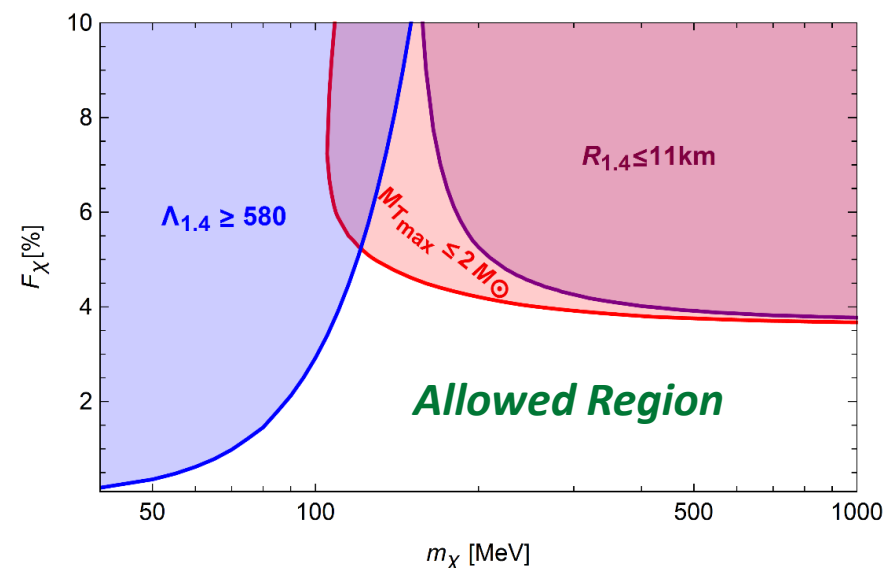
Phys. Rev. D 109, 043029 (2024) [77 citations]

D.R.K, M. Shahrbafe, S. Shakeri, S. Typel

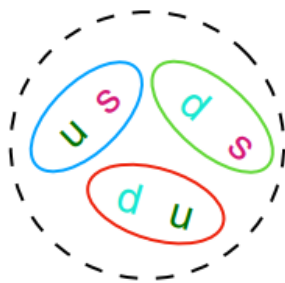
Exploring the distribution and impact of bosonic dark matter in neutron stars

Particles 7 1, 201-213 (2024) [25 citations]

Accretion of dark matter inside neutron stars



**Sexaquark a bosonic
dark matter candidate**



Glennys R. Farrar

arXiv:2201.01334, review paper

[Home](#) > [International Journal of Theoretical Physics](#) > [Article](#)

A Stable H-Dibaryon: Dark Matter, Candidate Within QCD?

Published: June 2003

Volume 42, pages 1211–1218, (2003) [Cite this article](#)

Sexaquark (Mass: ~ 2 GeV)

(S=uuddss): bosonic, spin-color-flavor singlet, deeply bound

Q=0, B = 2, s= -2

If $m_s \leq (m_\Lambda + m_p + m_e) = 2054$ MeV, it will decay with a lifetime more than the age of the universe.

OPEN ACCESS

Search for a Stable Six-Quark State at BABAR

[J. P. Lees](#)¹, [V. Poireau](#)¹, [V. Tisserand](#)¹, [E. Grauges](#)², [A. Palano](#)³, [G. Eigen](#)⁴, [D. N. Brown](#)⁵, [Yu. G. Kolomensky](#)⁵,
[M. Fritsch](#)⁶ *et al.* (BABAR Collaboration)

Show more ▾

Phys. Rev. Lett. **122**, 072002 – Published 22 February, 2019

OPEN ACCESS

Coalescence production of sexaquark with three diquarks in high-energy nuclear collisions

[Zhi-Lei She](#) ¹, [An-Ke Lei](#)^{2,3}, [Dai-Mei Zhou](#)^{2,*}, [Larissa V. Bravina](#) ^{3,†}, [Evgeny E. Zabrodin](#) ^{3,4}, [Sonia Kabana](#) ^{5,‡}, and [Vipul Bairathi](#) ⁵

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Phys. Rev. D **112**, 034002 – Published 4 August, 2025

[Home](#) > [The European Physical Journal C](#) > [Article](#)


Searching for a dark matter particle with anti-protonic atoms

Letter | [Open access](#) | Published: 18 December 2023

Volume 83, article number 1149, (2023) [Cite this article](#)

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[Michael Doser](#) , [Glennys Farrar](#) & [Georgy Kornakov](#)

Sexaquark cannot be a dark matter candidate


Phys. Lett. B 857 (2024) 138973

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Letter

Hypernuclear constraints on the existence and lifetime of a deeply bound H dibaryon


Avraham Gal

Racah Institute of Physics, The Hebrew University, Jerusalem 9190401, Israel



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Cosmology and terrestrial signals of sexaquark dark matter

[Marianne Moore](#) ^{*} and [Tracy R. Slatyer](#)[†]

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Phys. Rev. D **110**, 023515 – Published 11 July, 2024

Sexaquark dilemma in neutron stars and its solution by quark deconfinement

[M. Shahrbafe](#) ^{1,*}, [D. Blaschke](#) ^{1,2,3,†}, [S. Typel](#) ^{4,5,‡}, [G. R. Farrar](#) ^{6,§}, and [D. E. Alvarez-Castillo](#) ^{2,7,||}

Show more ▾

Phys. Rev. D **105**, 103005 – Published 4 May, 2022

For the first time Sexaquark dark matter (?) or exotic particle was investigated via its production in high density matter of neutron stars

Hadronic EOS:

- A generalized relativistic density functional (GRDF) approach
- DD2Y-T (includes hyperons) *Mon. Not. Roy. Astron. Soc., 502, 3476 2021*
- Exotic matter component: sexaquark (S) included in hadronic sector
- $m_S^* = m_S(1 + x_S \frac{n_b}{n_0})$
- Parameters: $m_S \in [1885 - 2054] \text{ MeV}$, $x_S = 0.03$
- x_S is the slope parameter and represents the coupling strength

Quark matter EOS :

nlNJL model

- The most likely parameters were selected within a Bayesian analysis (BA)
- Mapped to the constant speed of sound (CSS) parameterization
- M. Shahrbafe, et al. 2023, Phys. Rev. D, 107, 054011*
- Two sets of Parameters (η_D, η_V, c_S^2) have been used
(0.70, 0.10, 0.44) & (0.70, 0.11, 0.44)

**The model RIC-DD2Y-T- S_{m_S}
describes a neutron star that includes
Nucleons, Hyperons,
Sesquark **dark matter (?)** or exotic particle
and Quark matter**

Phase transition:

- Smooth **crossover phase transition** from hadronic matter to deconfined quark matter
- Replacement Interpolation Construction (**RIC**)

*M. Shahrbafe, D. Blaschke, S. Khanmohamadi,
J. Phys. G: Nucl. Part. Phys. 47 (2020) 115201*

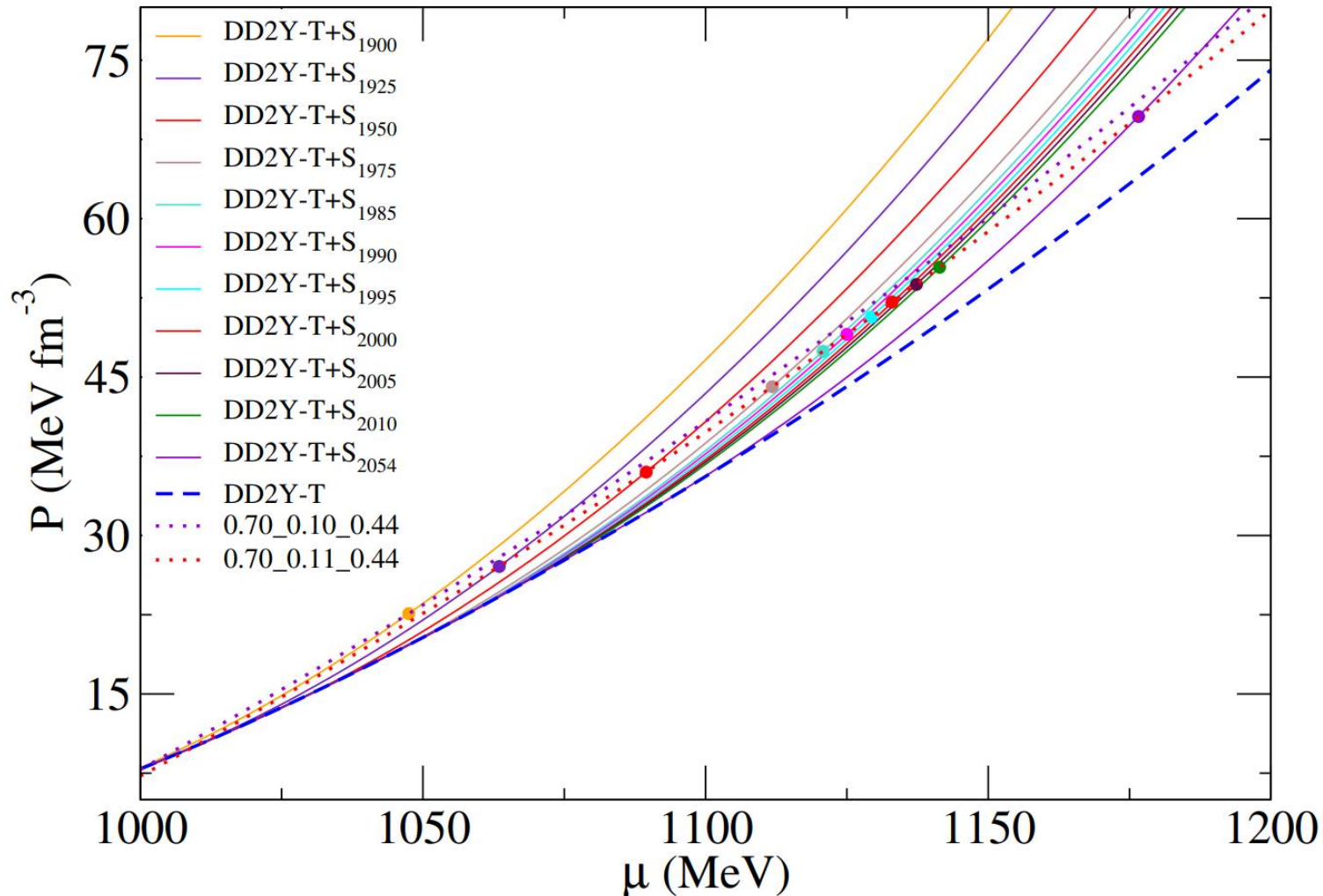
A. Ayriyan et al. Phys. Rev. C 97, 045802

Observational Probes of the Neutron Star Equation of State with Hyperons, Bosonic Dark Matter, and Quark Matter

Mahboubeh Shahrbafe^{1,2,3*}, Davood Rafiei Karkevandi^{1**}, Alexander Ayriyan^{1,4***}, and Stefan Typel^{5,6}

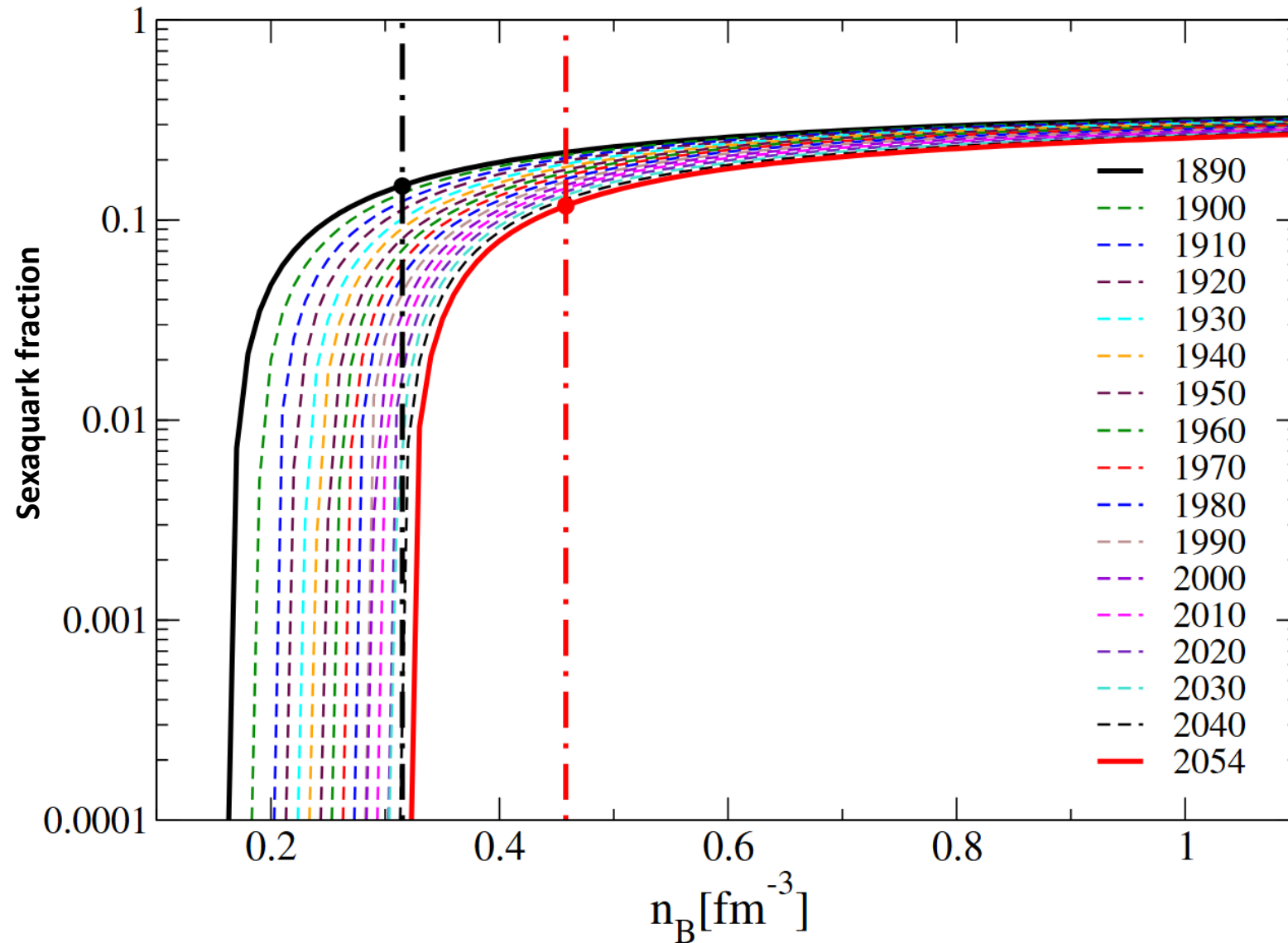
arXiv:2402.18686v2

Accepted for publication in
Astronomy & Astrophysics



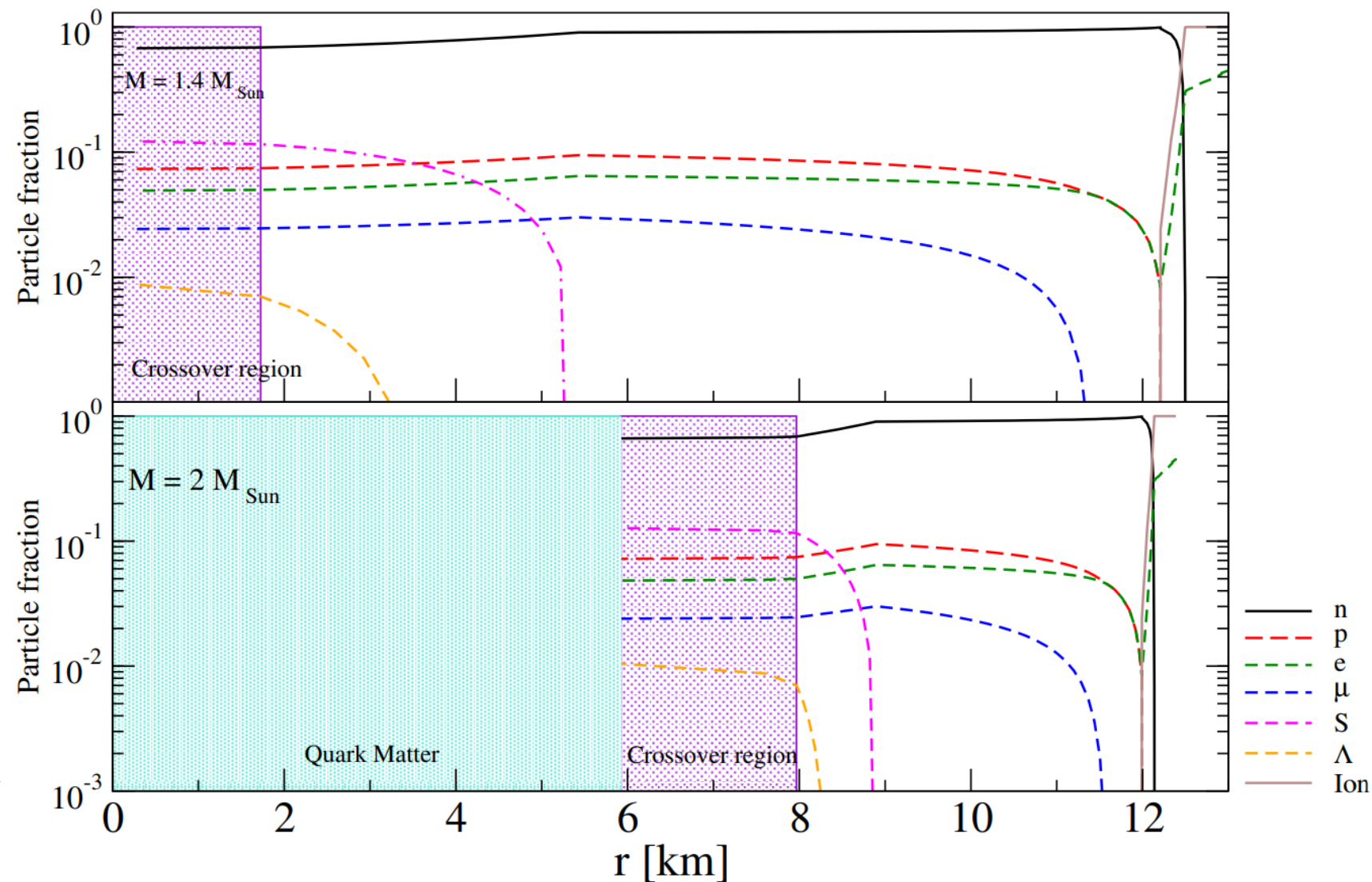
**Different masses of
Sexaquark particles
were considered for its
the lowest value of
coupling strength with
the medium.**

The maximum Sexaquark fraction allowed in our models varies from about 12% for the highest S-particle mass to 15% for the lightest one.

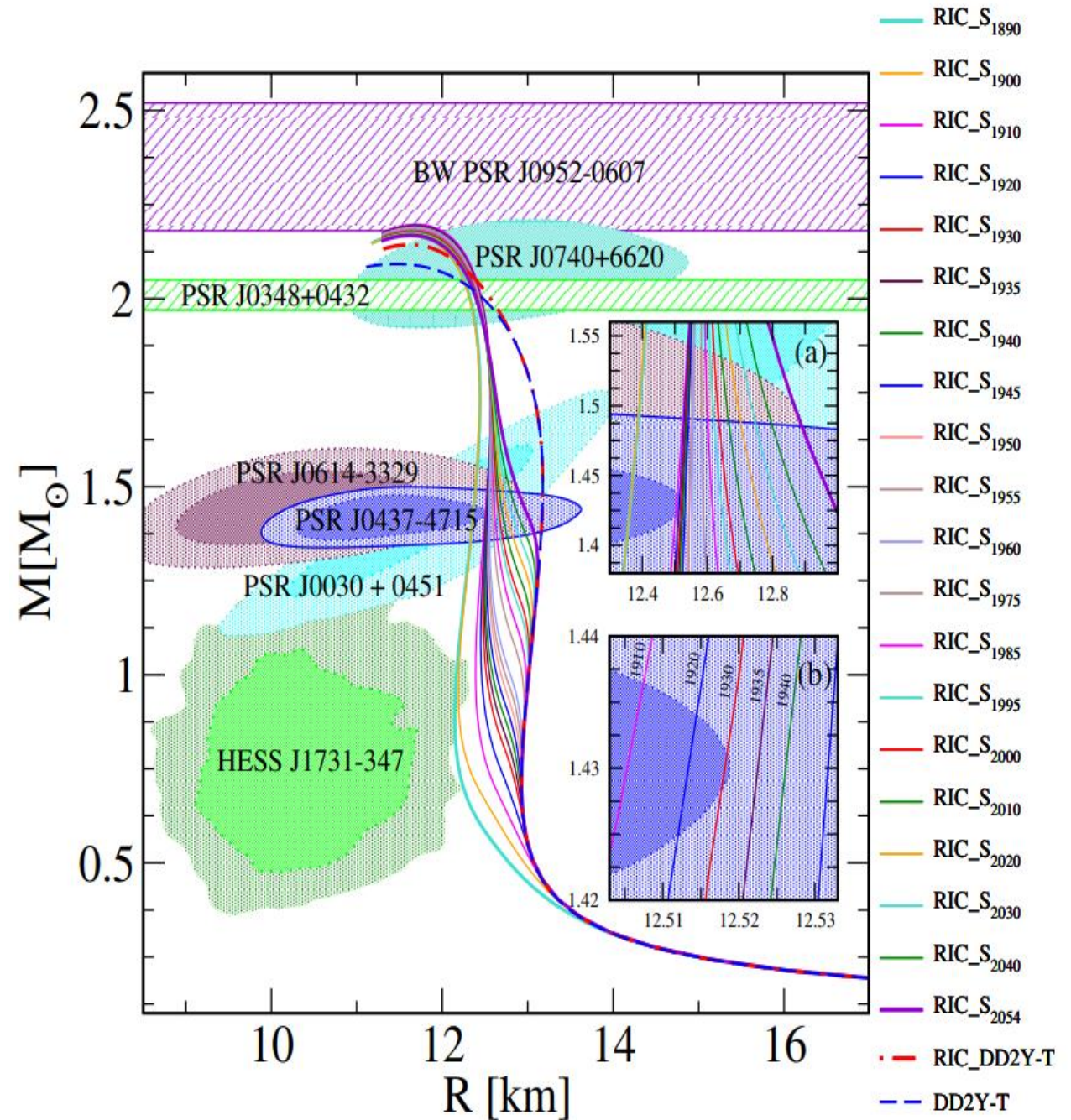


RIC-DD2Y-T+S₂₀₅₀, x = 0.03

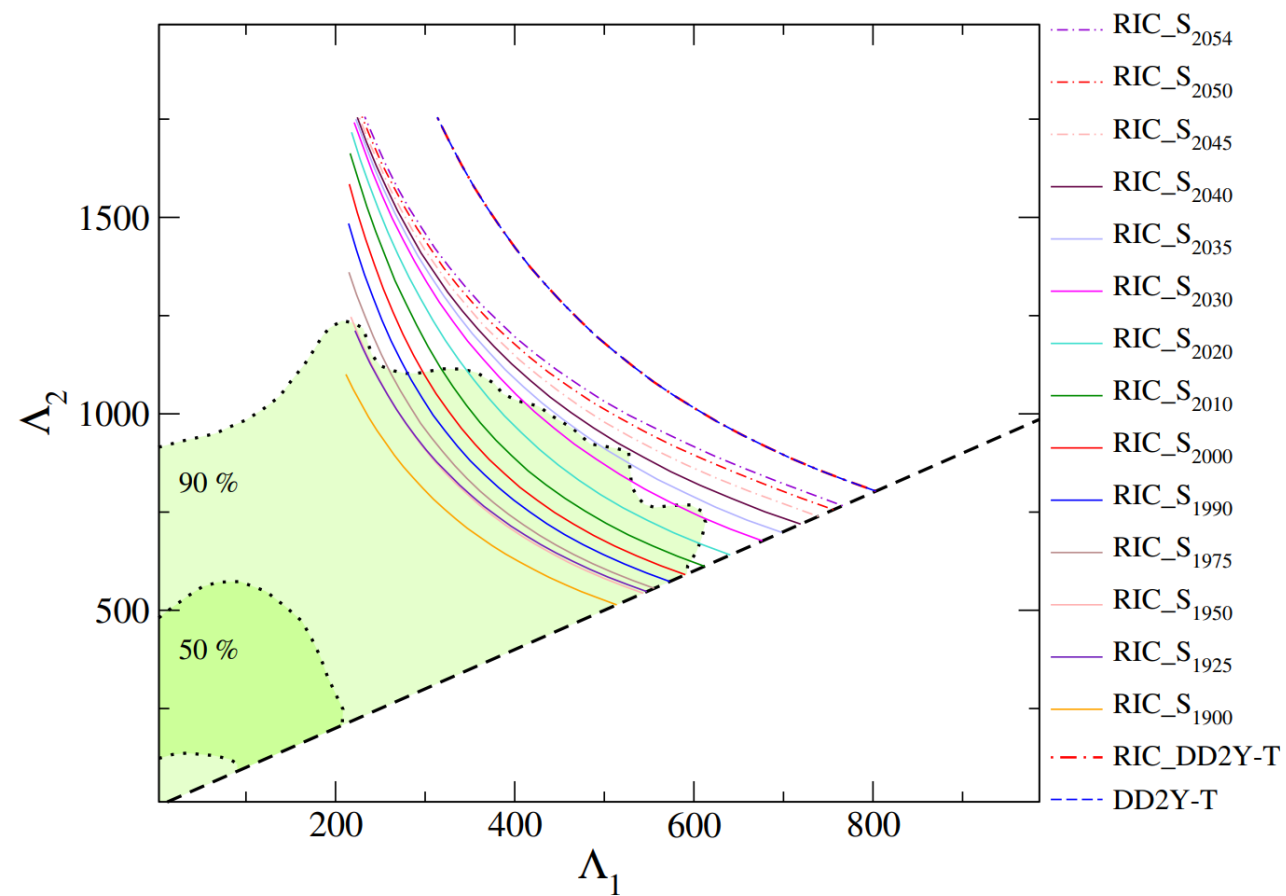
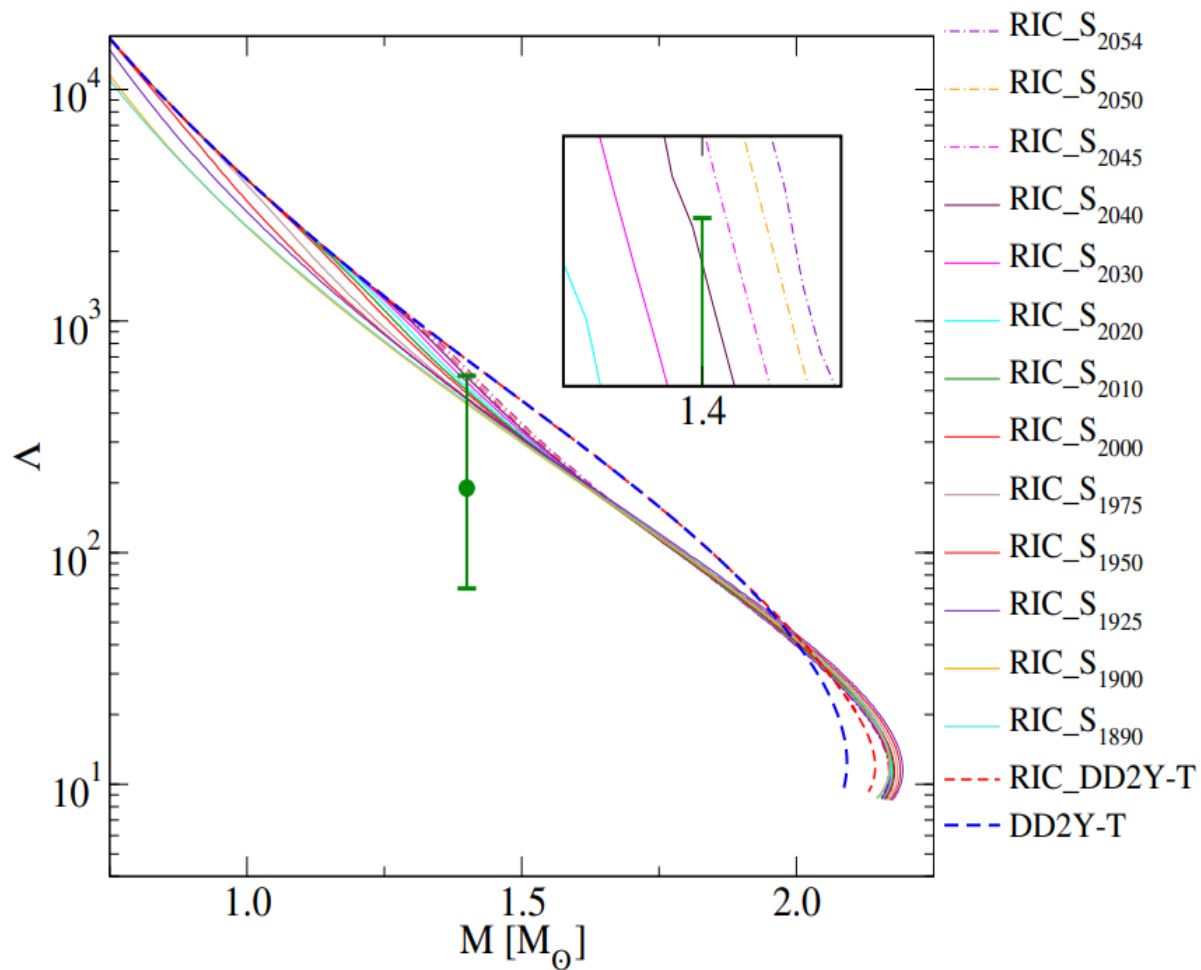
Particle fraction inside
the neutron star
for $1.4M_{\odot}$ (up)
and $2M_{\odot}$ (down)



Mass-Radius
observational constraints
for neutron stars are
considered to probe the
model and different
masses of Sexaquarks

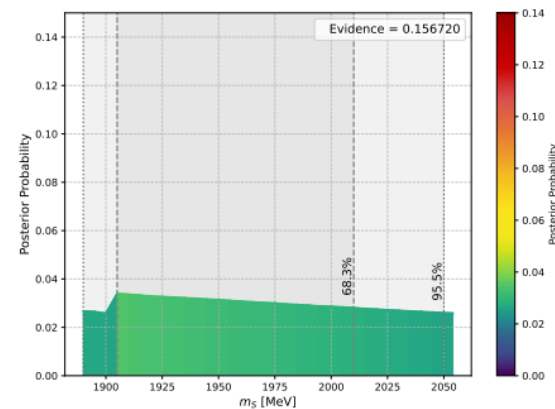


Tidal deformability constraints from gravitational waves of neutron stars merger

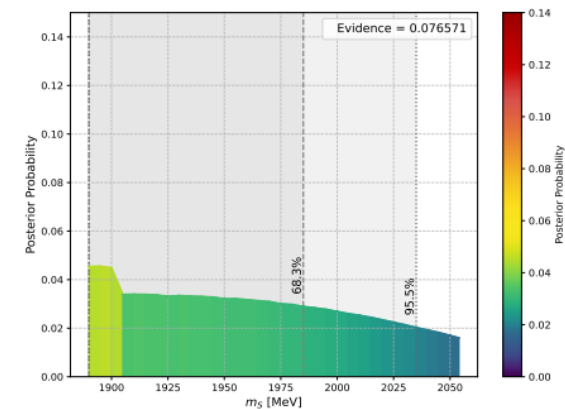


Bayesian analysis based on all available observational constraints for neutron stars

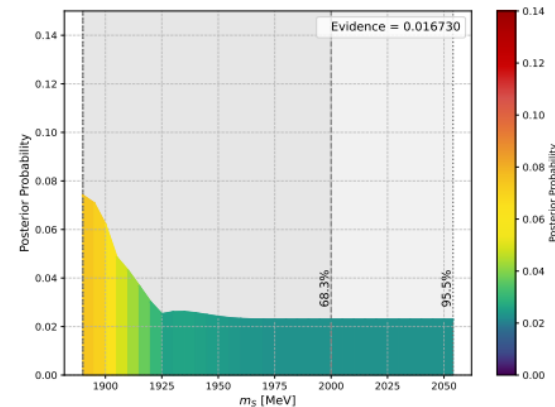
Our results disfavors Sexaquark
particle masses
higher than 1980 MeV
strongly favors $m_s \lesssim 1935$ MeV
for the lowest of coupling strength



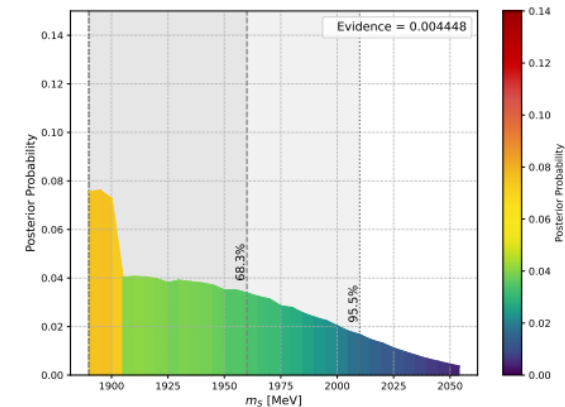
(a) (VI)



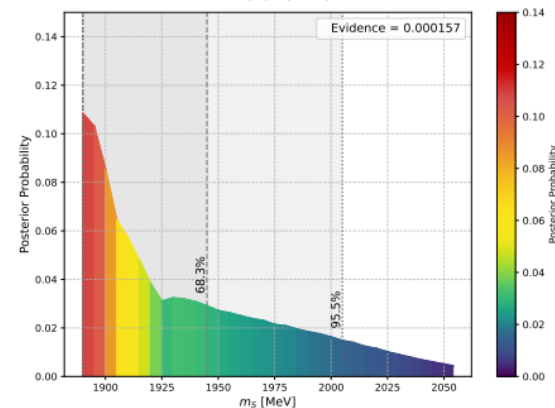
(a) (VIII)



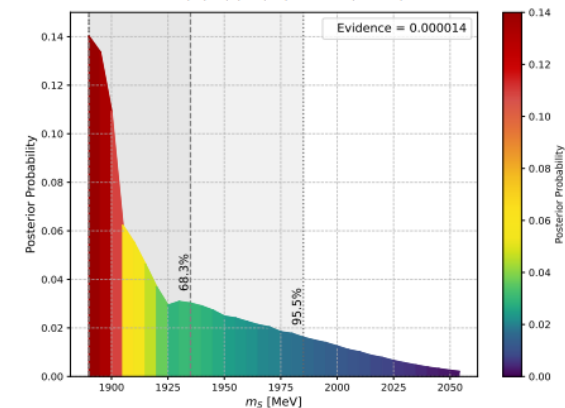
(b) (VII)



(b) (I)-(V) and (VIII)



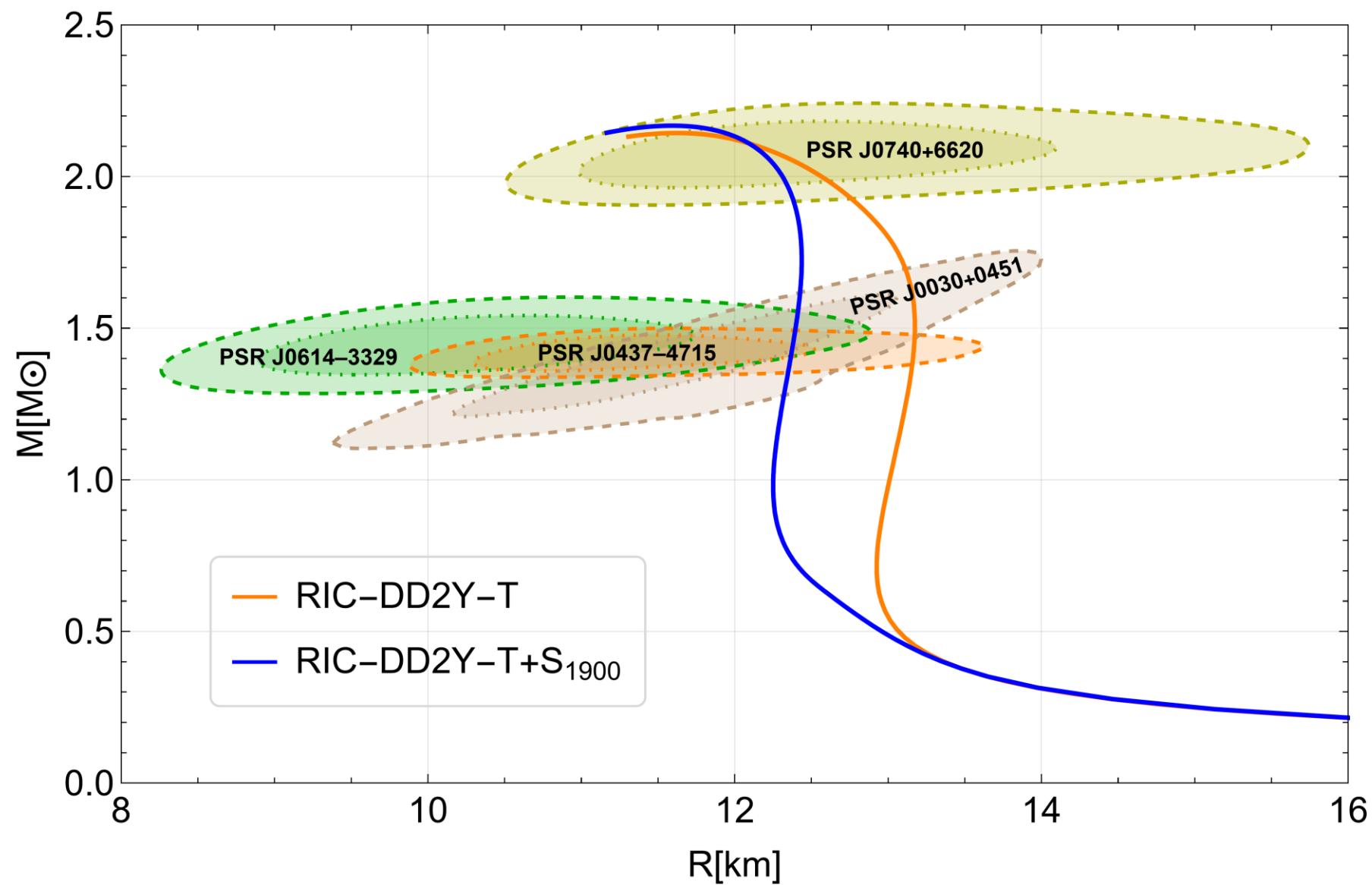
(c) (I)-(VII)



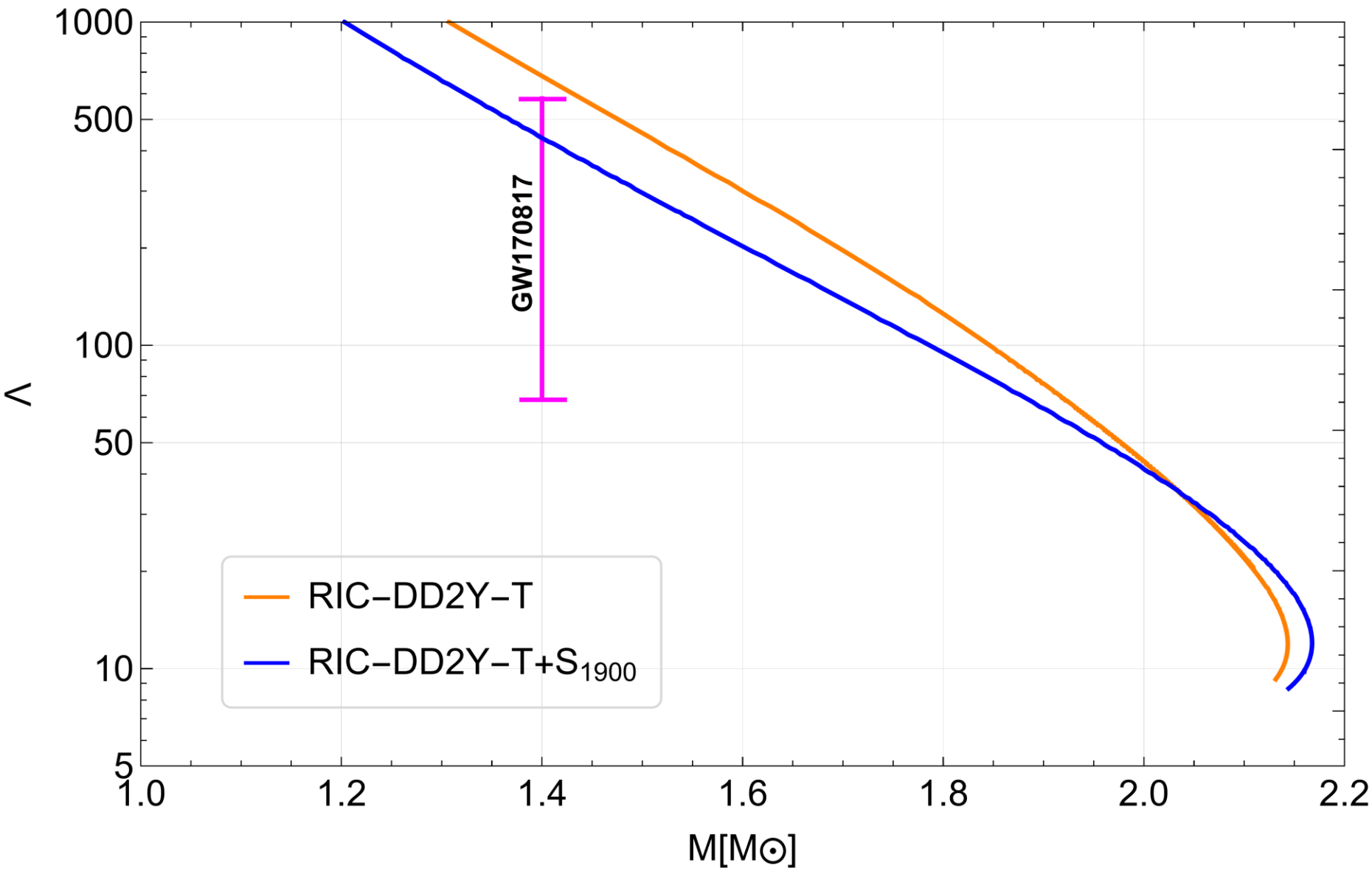
(c) (I)-(VIII)

Orange curve : without Sexaquark

Blue curve : with Sexaquark, $m_s = 1900 \text{ MeV}$ and $x_s = 0.03$

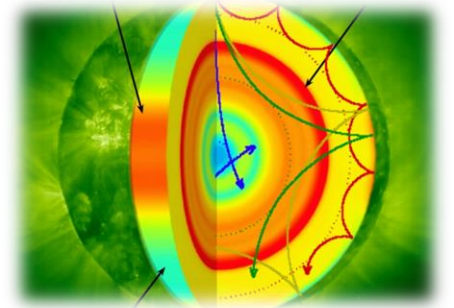
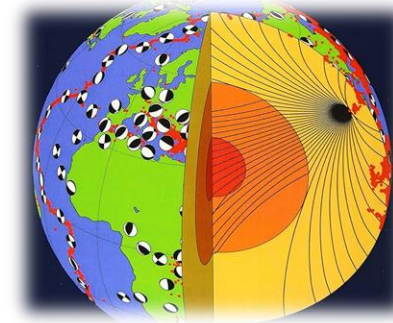


Presence of sexaquark softens the equation of the state and make it consistent with $\Lambda_{1.4}$ limit

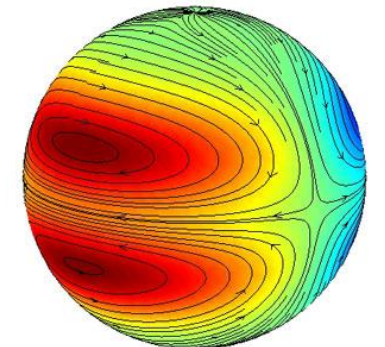
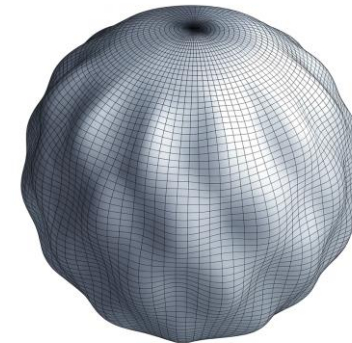


Radial and non-radial oscillations in neutron stars

Seismology and Helioseismology for earth and sun



Asteroseismology for neutron stars
the interior structure and
the equation of state can be probed



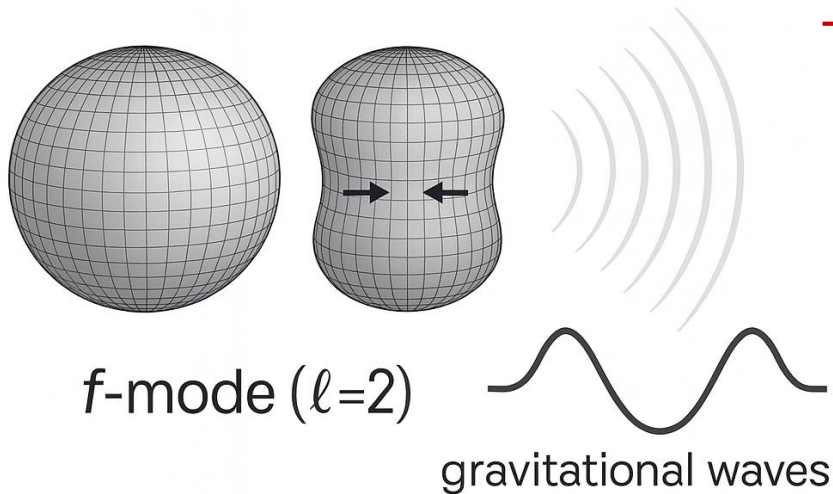
Of particular interest are the non-radial modes, they distort the star's spherical shape and thus serve as a sources of gravitational waves.

THE ASTROPHYSICAL JOURNAL, Vol. 149, September 1967

NON-RADIAL PULSATION OF GENERAL-RELATIVISTIC STELLAR
MODELS. I. ANALYTIC ANALYSIS FOR $l \geq 2^*$

KIP S. THORNE[†]
California Institute of Technology, Pasadena

AND
ALFONSO CAMPOLATTARO
University of California, Irvine
Received February 24, 1967



The fundamental mode (f-mode) oscillation in a neutron star

The frequency is between 1.3 and 2.8 kHz
could be detectable by current and future gravitational wave detectors

The f-mode oscillations are expected to be
damped over time

mainly through the emission of gravitational waves
The damping time is predicted to be relatively short,
around 0.1–0.5 seconds.

Gravitational wave asteroseismology

By calculating the universal relation between the f-mode frequency (and damping time) and the star's fundamental properties (mass, radius, tidal, EOS).



The screenshot shows the website for 'Monthly Notices of the Royal Astronomical Society'. The header is purple with the title in white. Below the header is a navigation bar with links: Issues, Advance Access, More content, Submit, Alerts, About, and a dropdown for 'Monthly Notices of the'. The main content area has a purple sidebar on the left with the journal's logo and 'Volume 299, Issue 4, October 1998'. The main text area is white and features the article title 'Towards gravitational wave asteroseismology' with a 'FREE' badge. Below the title is the authors' names 'Nils Andersson, Kostas D. Kokkotas' and the publication details: 'Monthly Notices of the Royal Astronomical Society, Volume 299, Issue 4, October 1998, Pages 1059–1068, <https://doi.org/10.1046/j.1365-8711.1998.01840.x>'. At the bottom of the article information is 'Published: 01 October 1998' and 'Article history'. Below this is a row of icons for PDF, Split View, Cite, Permissions, and Share.

Monthly Notices
of the Royal Astronomical Society

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JOURNAL ARTICLE

Towards gravitational wave asteroseismology FREE

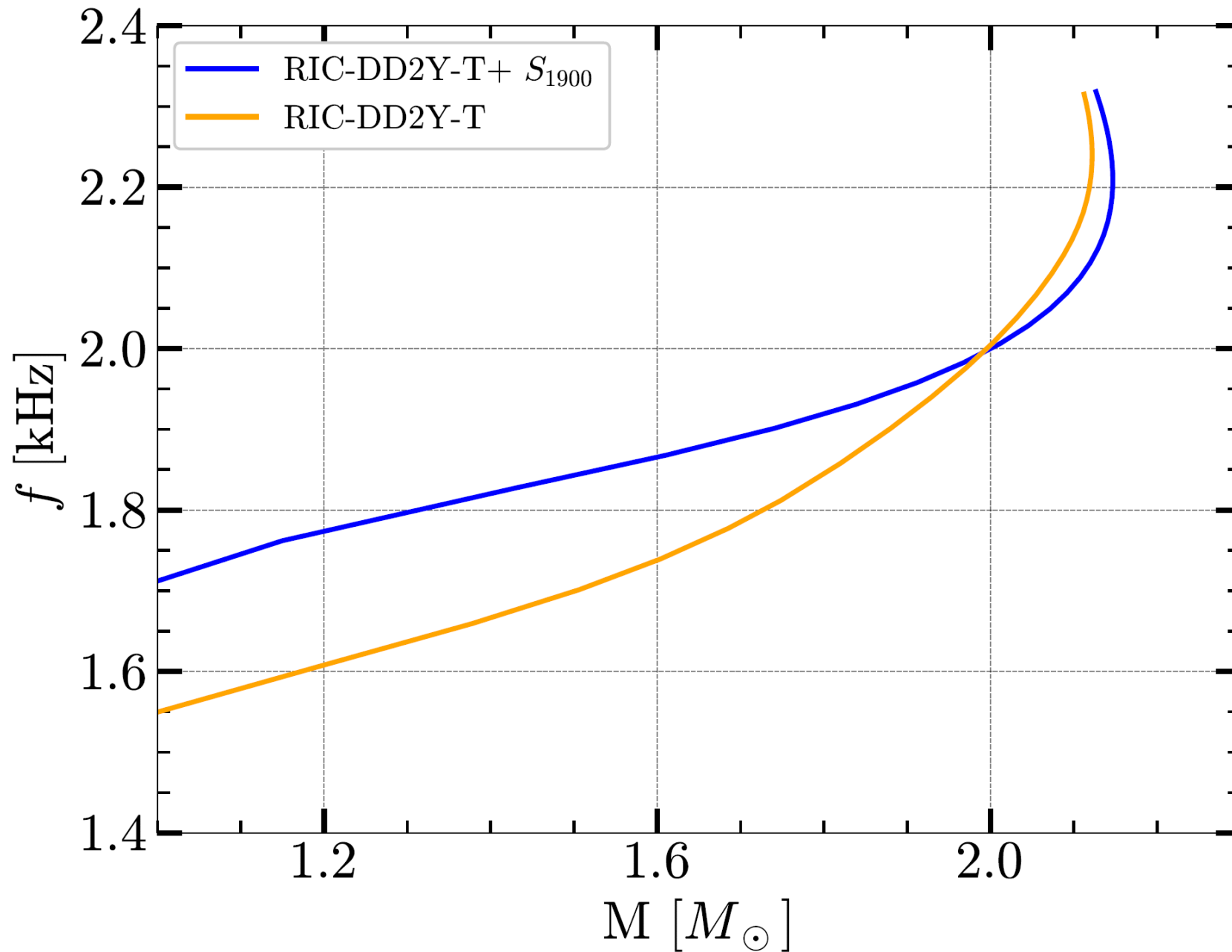
Nils Andersson, Kostas D. Kokkotas

Monthly Notices of the Royal Astronomical Society, Volume 299, Issue 4, October 1998,
Pages 1059–1068, <https://doi.org/10.1046/j.1365-8711.1998.01840.x>

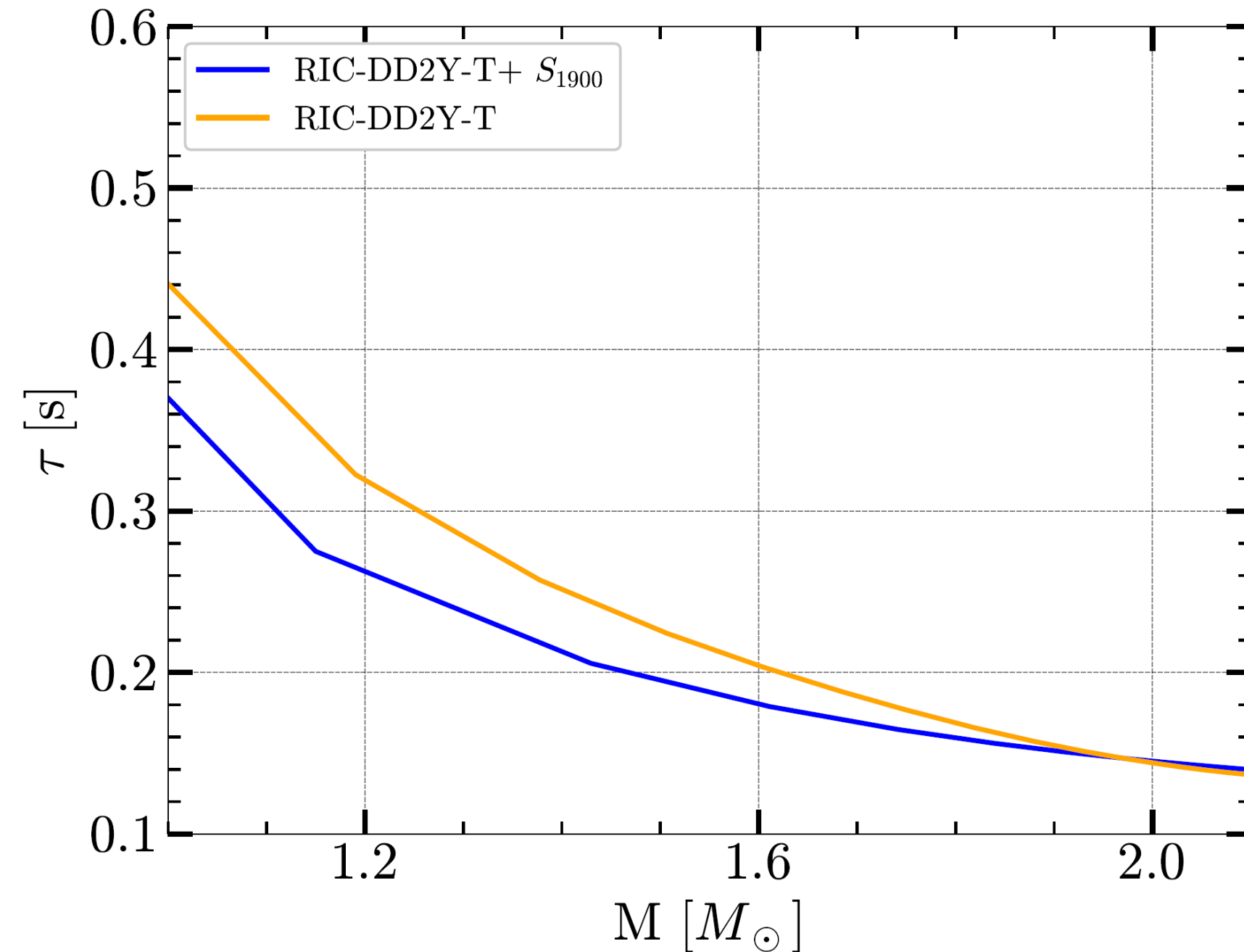
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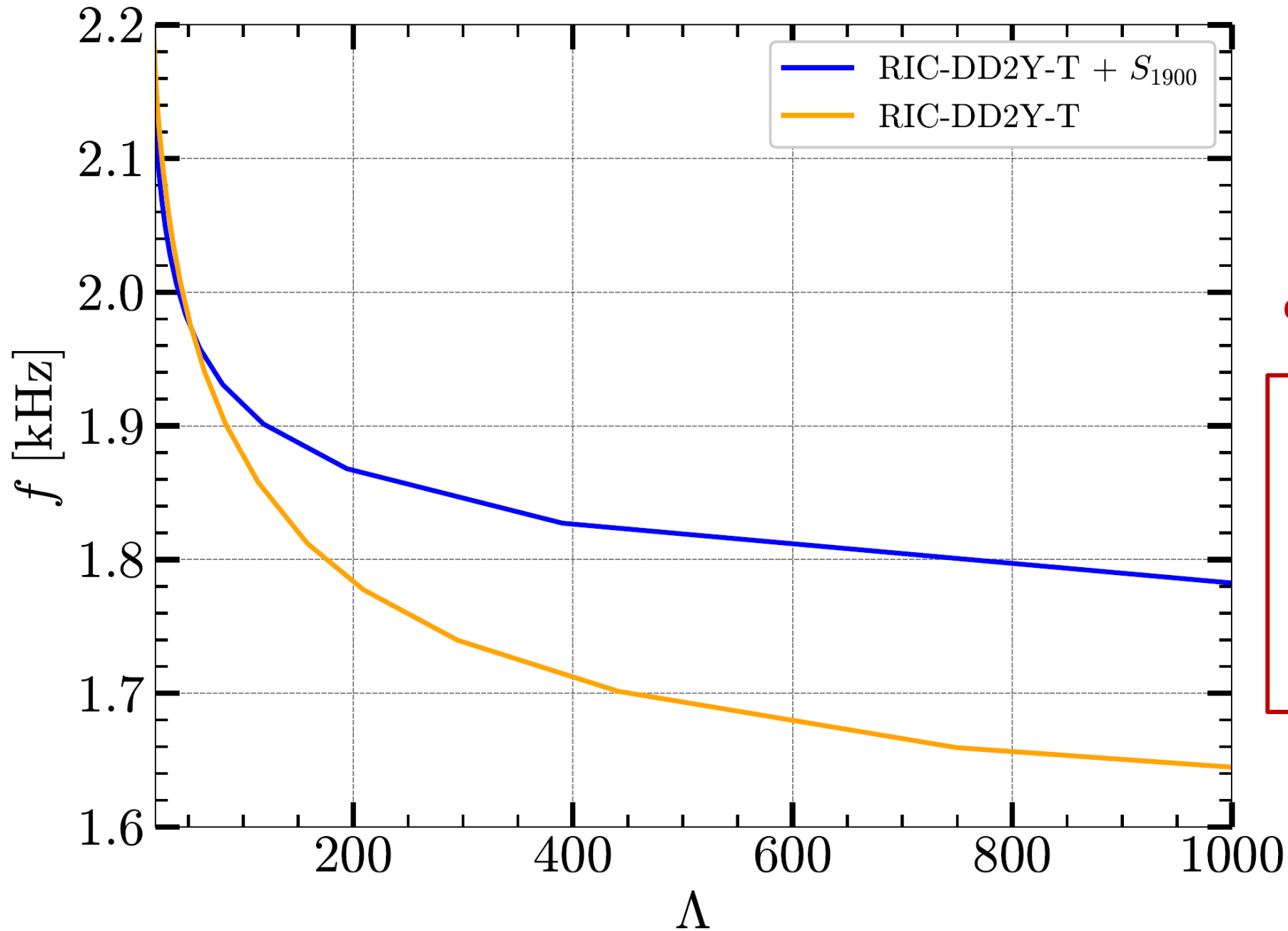
Andersson and Kokkotas
opening up the field of
gravitational wave
asteroseismology in which
they propose the first
**asteroseismological universal
relations for oscillation
modes of neutron stars.**



The presence of **Sexaquark** directly contributes to the **softening of the EOS** and **higher f-mode frequencies**.

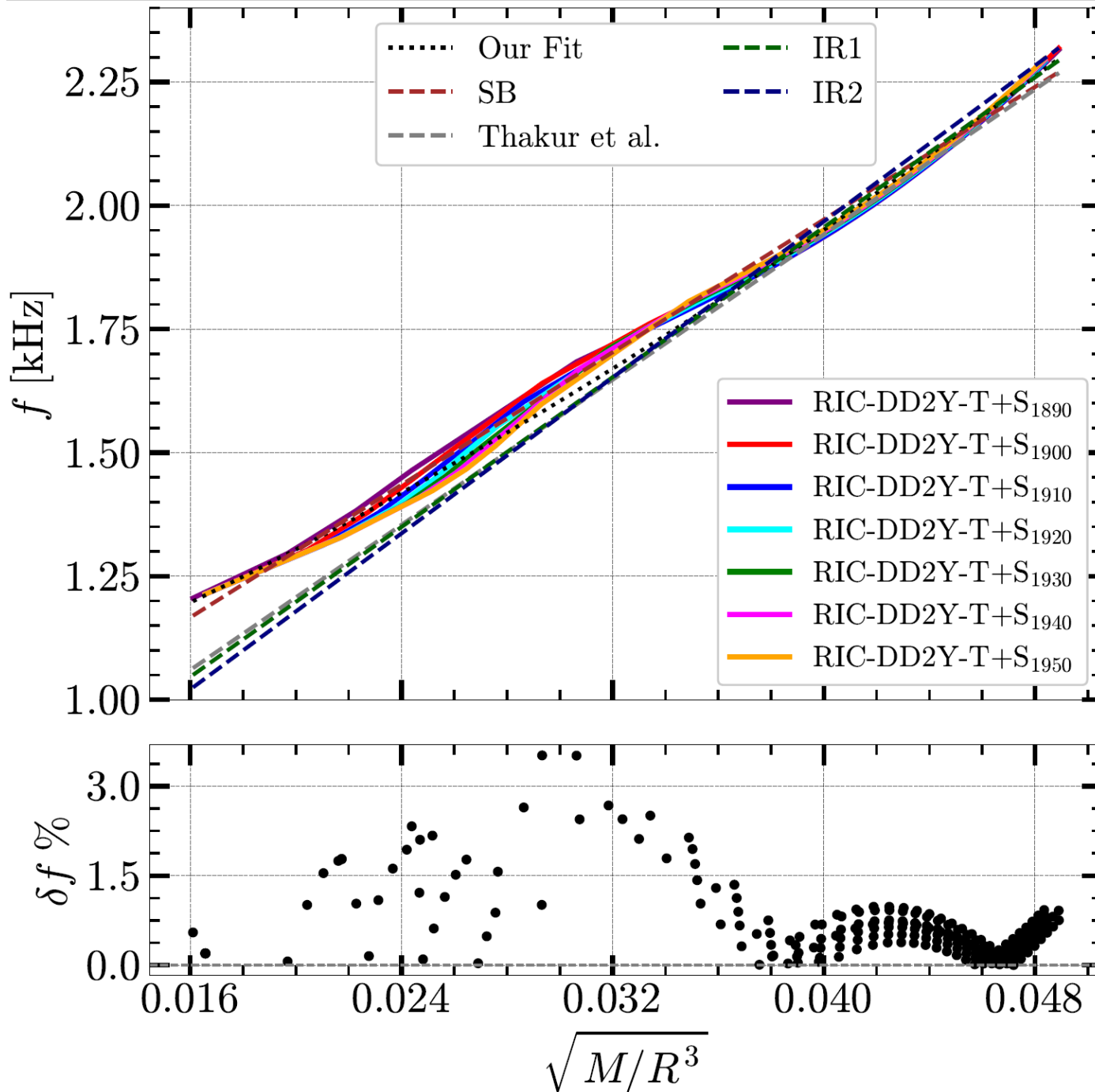


The increased compactness caused by Sexaquark **enhances metric perturbations**, and enables more efficient emission of **gravitational waves**, so that **shortens** the damping time



The f-mode has close connection with the tidal deformability of neutron stars

Independent or simultaneous measurements of the **f-mode frequency and tidal deformability** can therefore provide valuable constraints on the EOS and the underlying composition of NSs.

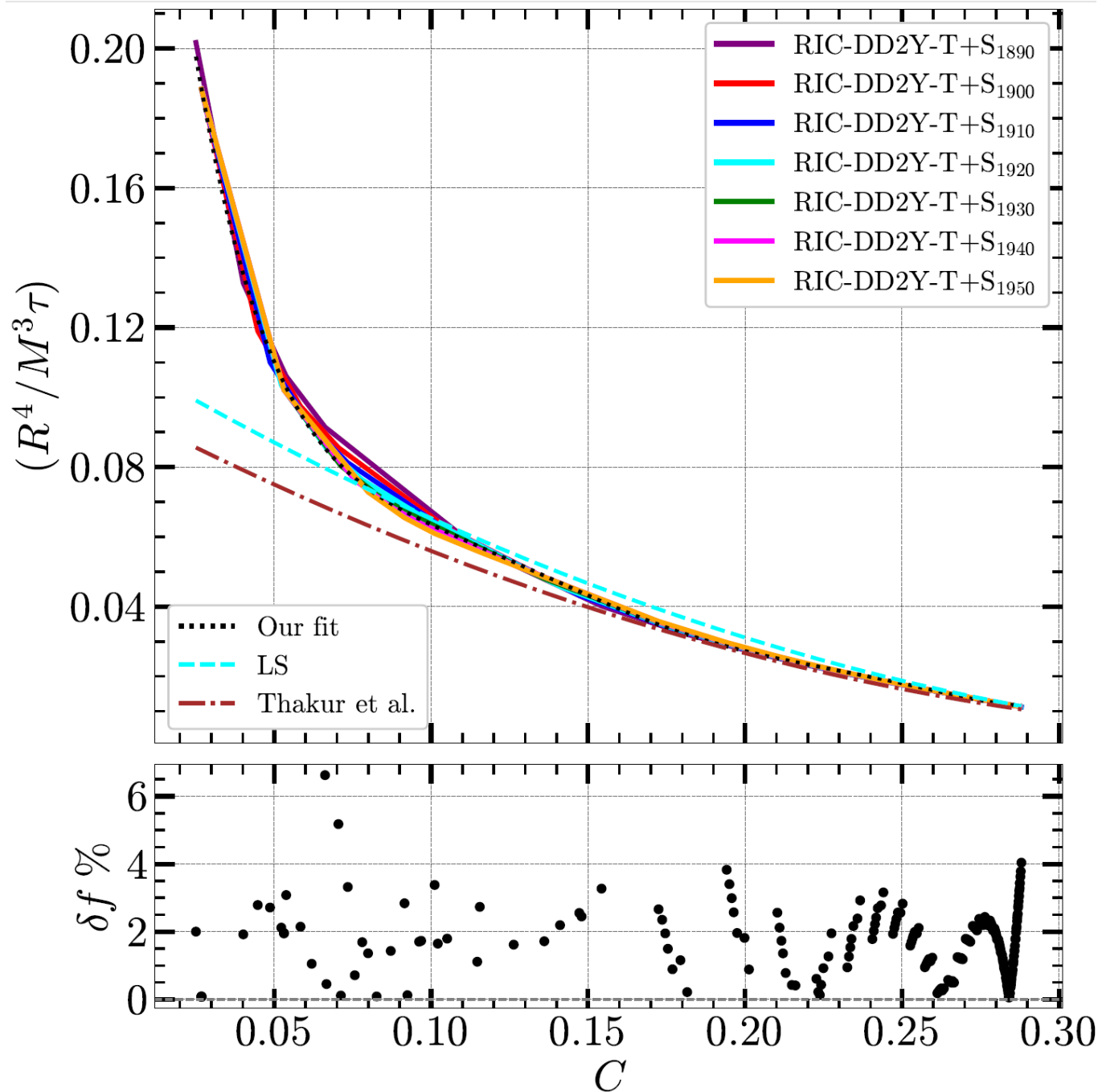


The **f-mode frequency** scales approximately with the **square root of the neutron star's average density**.

$$\frac{f}{\text{kHz}} = a + b \left(\sqrt{\frac{M}{R^3}} \right) + c \left(\sqrt{\frac{M}{R^3}} \right)^2$$

where $a = 0.840846$, $b = 18.569058$ and $c = 229.339724$

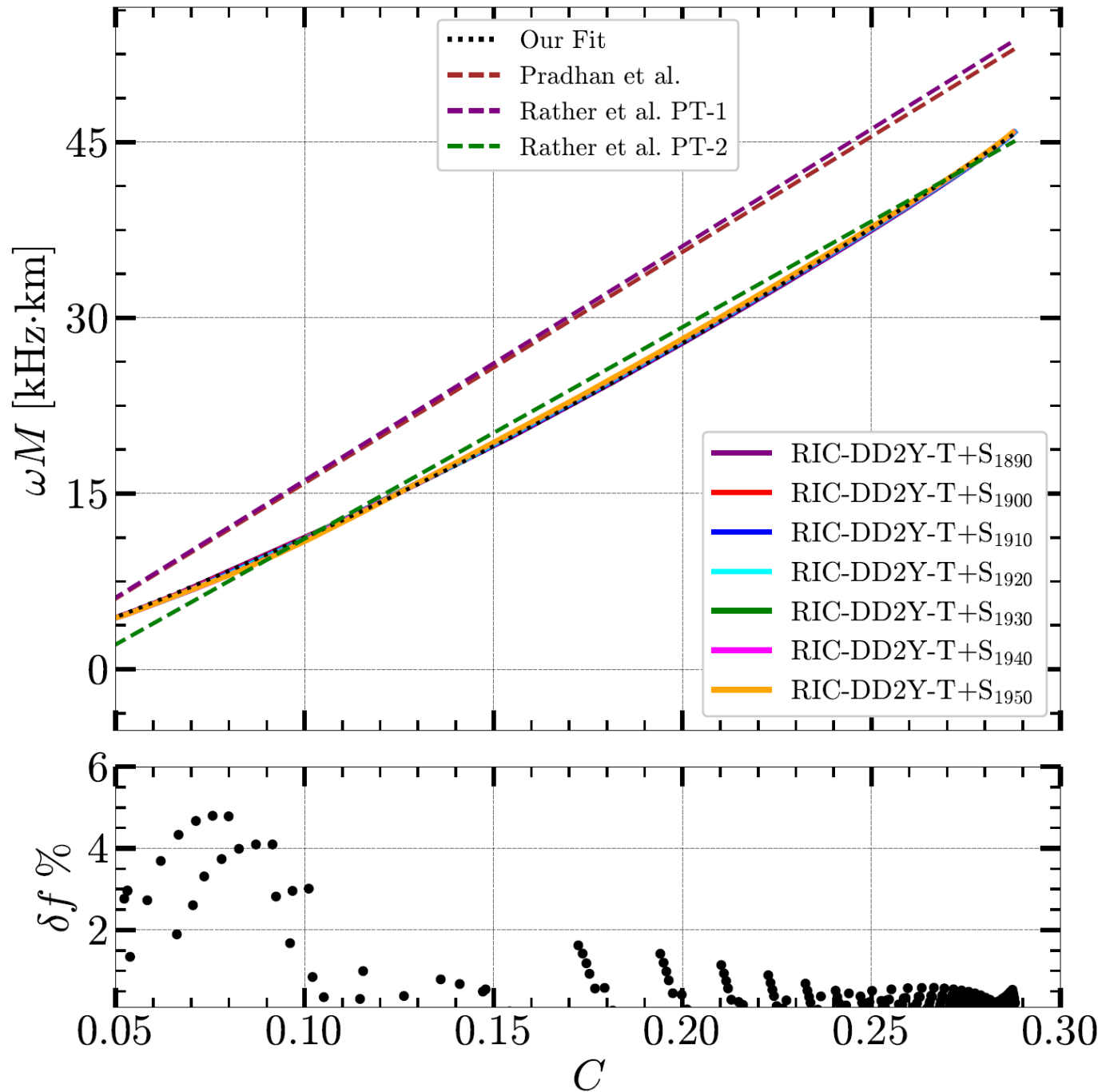
M. Shahrbaaf, P. Thakur, [D.R.K.](#), [arXiv:2510.08115](#)



**The quasi-universal relation
between the
scaled damping-time
variable and the
compactness**

$$\frac{R^4}{M^3 \tau} = a + bC + cC^2 + dC^3 + eC^4 + fC^5 + gC^6$$

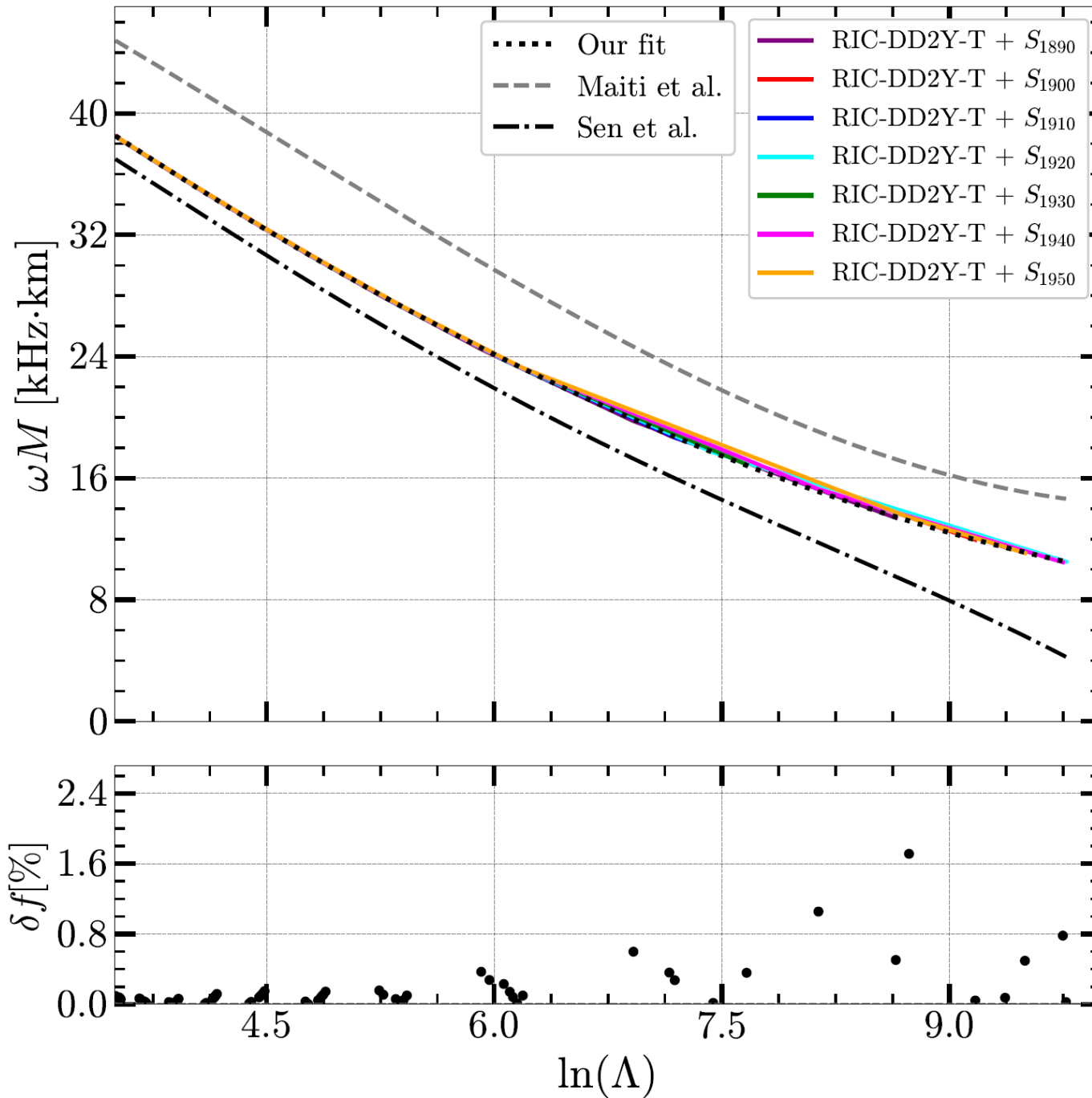
where, $a = 0.4095642$, $b = -12.26236$, $c = 186.0756$, $d = -1505.978$, $e = 6603.648$, $f = -14837.04$ and $g = 13389.98$



The mass scaled angular frequency of f-mode, ωM (with $\omega = 2\pi f$), as a function of the stellar compactness

$$\omega M = a C^2 + b C + c$$

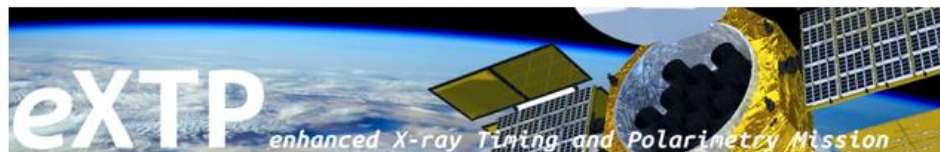
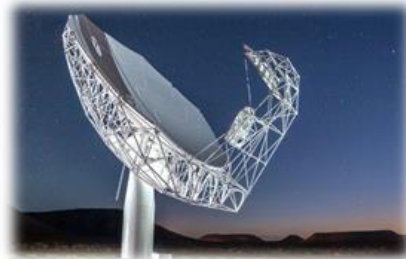
where $a = 199.734977$, $b = 106.082697$, $c = -1.365759$ have units of $\text{kHz} \cdot \text{km}$



The variation of the **scaled frequency**, ωM , as a function of the dimensionless **tidal deformability** parameter (Λ)

$$\omega M = a + b (\ln \Lambda) + c (\ln \Lambda)^2 + d (\ln \Lambda)^3$$

where $a = 67.426668$, $b = -9.797683$, $c = 0.470981$, and $d = -0.006726$



**Golden Age
of neutron stars research**

**Investigating the
properties of ultra
high density matter**

**We may shed
light on the nature
of dark matter**



Tomb of Cyrus the Great in Pasargad Iran

Photo by Sepehr Moelini

**What we know is a drop
What we do not know is an ocean**