



האוניברסיטה העברית בירושלים
THE HEBREW UNIVERSITY OF JERUSALEM

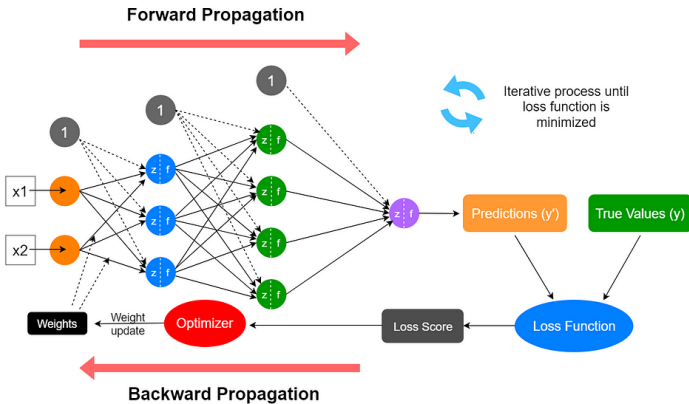


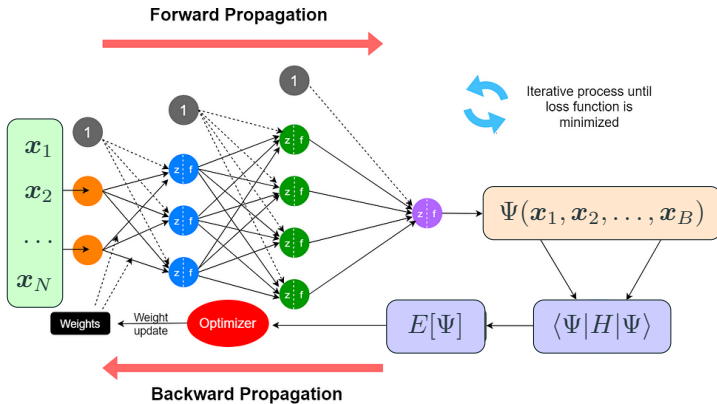
Nuclear Responses with Neural-Network Quantum States

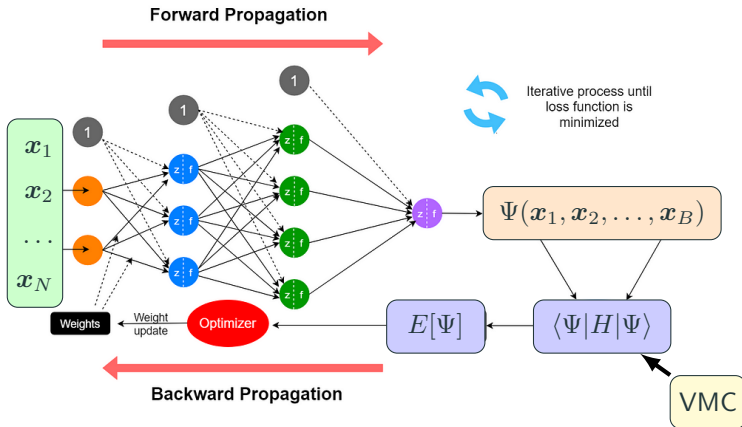
Nir Barnea

EMMI Workshop, January 18-24, 2026

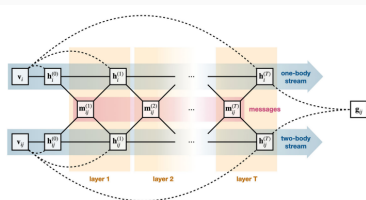
Hirschegg







The MPNN Architecture - Message Passing Neural Network architecture - creates antisymmetrized descriptors of particle location that are already well processed by NNs.



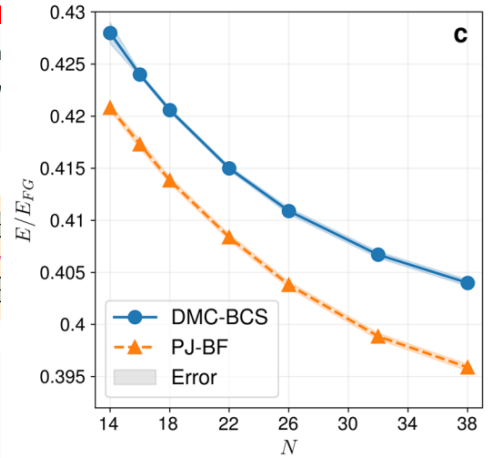
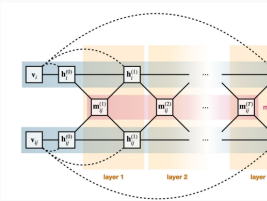
$$\Psi(X) = e^{J(X)} \Phi_{PJ}(X)$$

$$\Phi_{PJ}(X) = \text{pf}(P) \quad ; \quad \text{pf}^2(P) = \text{Det}(P)$$

$$P \equiv \begin{pmatrix} 0 & \phi(x_1, x_2) & \cdots & \phi(x_1, x_N) \\ -\phi(x_2, x_1) & 0 & \cdots & \phi(x_2, x_N) \\ \vdots & \vdots & \ddots & \vdots \\ -\phi(x_N, x_1) & -\phi(x_N, x_2) & \cdots & 0 \end{pmatrix}$$

J. Kim, G. Pescia, B. Fore, N. Nys, G. Carleo, S. Gandolfi, M. Hjorth-Jensen, and A. Lovato, Communications Physics 7, 148 (2024)

The MPNN Arcl architecture - creates states that are already well-suited for the task



location

$\phi^2(P) = \text{Det}(P)$

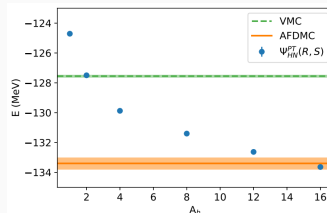
$$\begin{pmatrix} \phi(x_1, x_N) \\ \phi(x_2, x_N) \\ \vdots \\ 0 \end{pmatrix}$$

J. Kim, G. Pescia, B. Fore, N. Nys, G. Carleo, S. Gandolfi, M. Hjorth-Jensen, and A. Lovato, Communications Physics 7, 148 (2024)

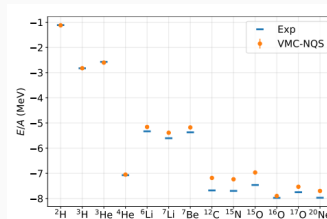
A prominent example of an **antisymmetric wavefunction** is the **Hidden Nucleons architecture**. This extension of the Slater determinant allows much more flexibility.

$$\Psi_{HN}(X) \equiv \begin{vmatrix} \phi_v(X) & \chi_v(X_h) \\ \chi_h(X) & J_h(X_h) \end{vmatrix},$$

Flexibility can be enhanced by incorporating a symmetric Jastrow like terms.



A. Lovato et al., Phys. Rev. R 4, 043178 (2022)

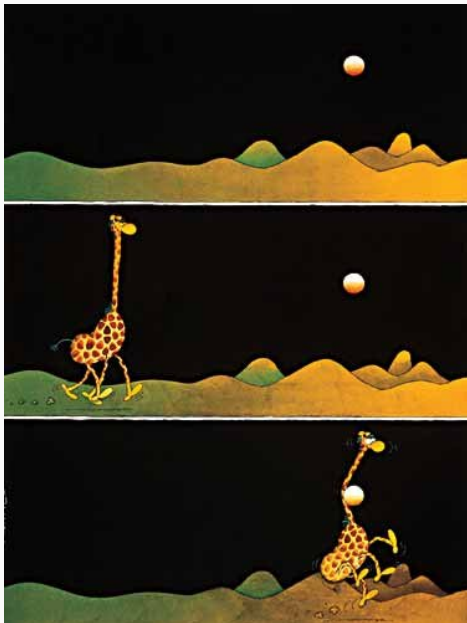


A. Lovato et al. (2024)



- **General-purpose NQS / VMC frameworks**
 - **NetKet** (Python, JAX)
Spin, Bose/Fermi; RBM, CNN, scalable VMC
 - **NeuralQuantumStates.jl** (Julia)
Early Julia-based NQS; pedagogical and research-oriented
- **Quantum chemistry-focused NQS**
 - **PauliNet** (Python, PyTorch)
Fermi w.f. with explicit physical priors (Slater, Jastrow)
 - **FermiNet** (Python, JAX)
Electronic structure
 - **DeepErwin** (Python, TensorFlow/JAX)
Deep-learning-based variational ansatz for molecules
- **Hybrid quantum / differentiable programming libraries**
 - **PennyLane** (Python)
NQS-like variational w.f. and differentiable quantum models
 - **TorchQuantum** (Python, PyTorch)
Neural-network-driven quantum state representations

What about
reactions?



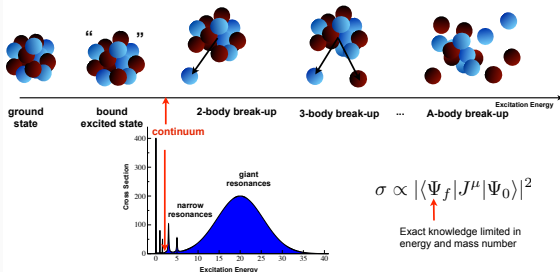


The EW interaction Hamiltonian: $\hat{H}_{EW} = \int dx \hat{A}_\mu(x) \hat{J}^\mu(x)$

Example - the photo-absorption **cross-section**

$$\sigma(\omega) = 4\pi^2 \alpha \omega \sum_J [\mathcal{R}_{E_J}(\omega) + \mathcal{R}_{M_J}(\omega)]$$

$$\mathcal{R}_\Theta(\omega) = \sum_f |\langle \Psi_f | \hat{\Theta} | \Psi_0 \rangle|^2 \delta(E_f - E_0 - \omega)$$



The Lorentz Integral Transform (LIT)



- Response in continuum

$$\mathcal{R}(\omega) = \sum_f |\langle \Psi_f | \hat{\Theta} | \Psi_0 \rangle|^2 \delta(E_f - E_0 - \omega)$$

- Lorentz integral transform (LIT) method

$$\mathcal{L}(\sigma, \Gamma) = \int d\omega \frac{\mathcal{R}(\omega)}{(\sigma - \omega)^2 + \Gamma^2} = \langle \tilde{\Psi} | \tilde{\Psi} \rangle$$

V. Efros, W. Leidemann, G.
Orlandini, PLB 238, 130
(1994)

- The LIT equation

$$(H - E_0 - \sigma + i\Gamma) |\tilde{\Psi}\rangle = \hat{\Theta} |\Psi_0\rangle$$

The LIT:

- The LIT equation is a Schrödinger equation with a **source**
- $|\tilde{\Psi}\rangle$ has bound-state asymptotic behavior
- $\mathcal{R}(\omega)$ is obtained inverting the LIT

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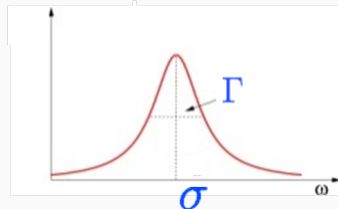
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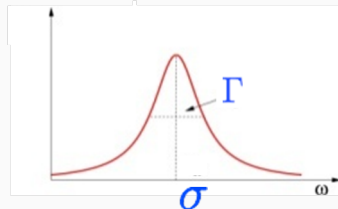
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The LIT:

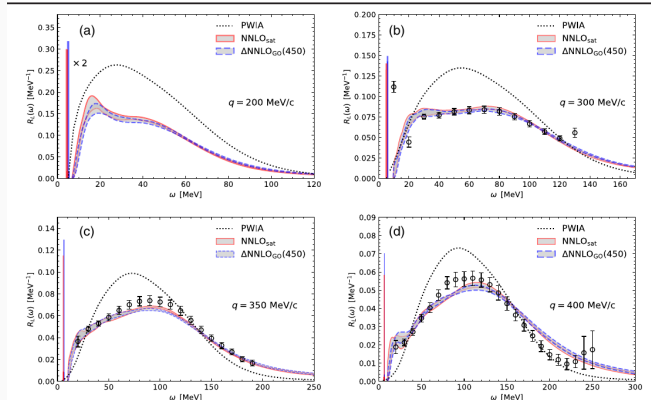
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The CC-LIT method

Longitudinal response of ^{40}Ca

PHYSICAL REVIEW LETTERS **127**, 072501 (2021)



J. E. Sobczyk, B. Acharya, S. Bacca, and G. Hagen



- The LIT variational principle
- The convergence lemma
- The H_{LIT} ground state gap
- The fidelity and the accuracy of the LIT



The LIT equation

$$(\hat{H} - z)|\tilde{\Psi}\rangle = |R\rangle \quad ; \quad |R\rangle \equiv \hat{\Theta}|\Psi_0\rangle,$$

Here $z = E_0 + \sigma - i\Gamma$.

Lemma:

The solution $\tilde{\Psi}$ is unique, and minimizes the functional

$$\mathcal{I}[\tilde{\Psi}] = |(\hat{H} - z)|\tilde{\Psi}\rangle - |R\rangle|^2$$

Proof:

If $|\tilde{\Psi}_0\rangle$ is the **exact** LIT WF then $\mathcal{I}[\tilde{\Psi}_0] = 0$.

Otherwise, for $|\tilde{\Psi}\rangle = |\tilde{\Psi}_0 + \delta\tilde{\Psi}\rangle$ one has

$$\begin{aligned} \mathcal{I}[\tilde{\Psi}] &= |(\hat{H} - z)|\tilde{\Psi}_0 + \delta\tilde{\Psi}\rangle - |R\rangle|^2 \\ &= |(\hat{H} - z)|\delta\tilde{\Psi}\rangle|^2 \geq \Gamma^2 \langle \delta\tilde{\Psi} | \delta\tilde{\Psi} \rangle \geq 0. \end{aligned}$$



Lemma: The relative accuracy of the LIT function $\mathcal{L}$ is given by

$$\left| \frac{\Delta \mathcal{L}}{\mathcal{L}} \right| = \mathcal{O} \left(\frac{\sqrt{\mathcal{I}[\tilde{\Psi}]}}{\Gamma \sqrt{\mathcal{L}}} \right) \quad ; \quad \Delta \mathcal{L} = \langle \tilde{\Psi} | \tilde{\Psi} \rangle - \langle \tilde{\Psi}_0 | \tilde{\Psi}_0 \rangle$$

In the limit $\Gamma \rightarrow 0$, $\sigma \approx \omega$, and $\mathcal{L} \rightarrow (\pi/\Gamma)\mathcal{R}$,

Hence, we may further suggest that

$$\left| \frac{\Delta \mathcal{R}}{\mathcal{R}} \right| = \mathcal{O} \left(\sqrt{\frac{\mathcal{I}[\tilde{\Psi}]}{\Gamma \mathcal{R}}} \right)$$

Implication:

For a desired relative accuracy in $\mathcal{R}$, we need to calculate $\mathcal{I}$ with a precision proportional to Γ .



Given the **LIT** equation

$$(\hat{H} - z)|\tilde{\Psi}\rangle = |R\rangle \quad ; \quad |R\rangle \equiv \hat{\Theta}|\Psi_0\rangle,$$

we introduce the LIT Hamiltonian:

$$H_{LIT} \equiv (\hat{H} - z^*)(1 - \hat{P}_R)(\hat{H} - z) \quad ; \quad \hat{P}_R = \frac{|R\rangle\langle R|}{\langle R|R\rangle}$$

Observations:

- H_{LIT} is Hermitian
- The LIT w.f. $|\tilde{\Psi}\rangle$ minimizes H_{LIT}
- $|\tilde{\Psi}\rangle$ is the zero energy eigenstates of H_{LIT}
- The spectrum of H_{LIT} has a gap above the ground state.



The Lorentz Integral Transform is obtained from the solution of the inhomogeneous equation:

$$\underbrace{(\hat{H} - z)}_{|L\rangle} |\tilde{\Psi}\rangle = \underbrace{\hat{O}}_{|R\rangle} |\Psi_0\rangle$$

We solve the LIT equation maximizing the **FIDELITY**

$$\mathcal{F} = \frac{\langle L|R\rangle\langle R|L\rangle}{\langle L|L\rangle\langle R|R\rangle}$$

Upper bounds on the LIT uncertainty:

$$\Delta\mathcal{L}(\omega_0, \Gamma) \leq \mathcal{D} \frac{\mathcal{N}^{-1}||R\rangle|}{\Gamma} \sqrt{\frac{1 - \mathcal{F}}{\mathcal{F}}},$$

where

$$\mathcal{D} = \min \left(\left| (1 - P_R) |\tilde{\Psi}\rangle \right|, \left| (1 - P_R) \frac{H}{\sqrt{\sigma^2 + \Gamma^2}} |\tilde{\Psi}\rangle \right| \right)$$

NNLIT

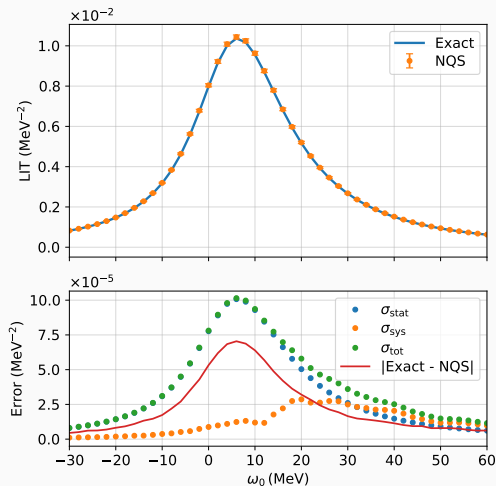


Key Idea:

A variational Monte Carlo framework to compute dynamical nuclear response functions

- Neural-network quantum states (NQS)
- Lorentz Integral Transform (LIT)
- Both ground state Ψ_0 and LIT state $\tilde{\Psi}$ represented by NQS
- Optimized via the neural net

- Avoids explicit treatment of continuum states
- Compact representation
- Efficient scaling with particle number
- Rigorous upper bounds for systematic errors derived



- **Upper panel:** LIT of the ^2H photo-absorption dipole response

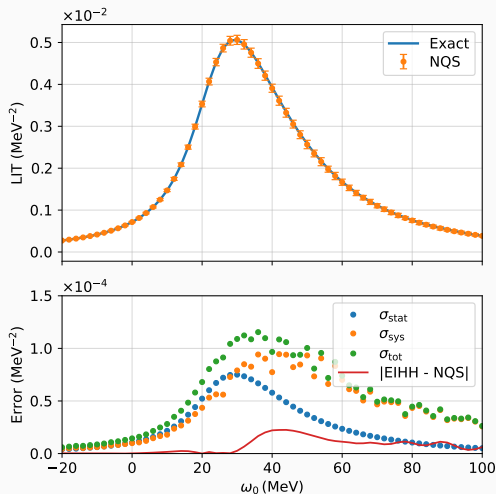
- $\Gamma = 10$ MeV

- **Lower panel:** NQS errors

- Statistical - MC sampling

- Systematic - LIT error-bound lemma

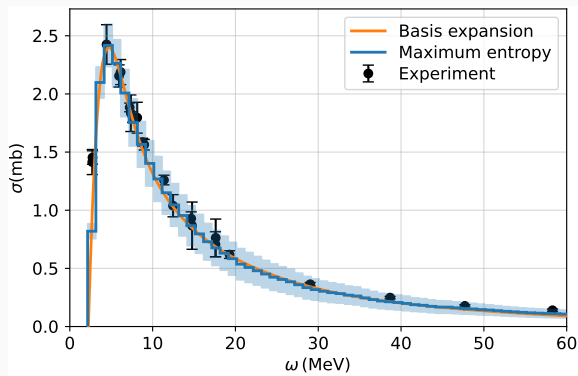
- Total - sum of the two



- **Upper panel:** LIT of the ^4He photo-absorption dipole response
- $\Gamma = 10$ MeV
- **Lower panel:** NQS errors
- Statistical - MC sampling
- Systematic - LIT error-bound lemma
- Total - sum of the two



- Pionless EFT inspired potential [A]
- Excellent accuracy at the LIT level
- Inversion - Basis expansion + regulation, Maximum entropy
- Error bars - sys+stat+inver

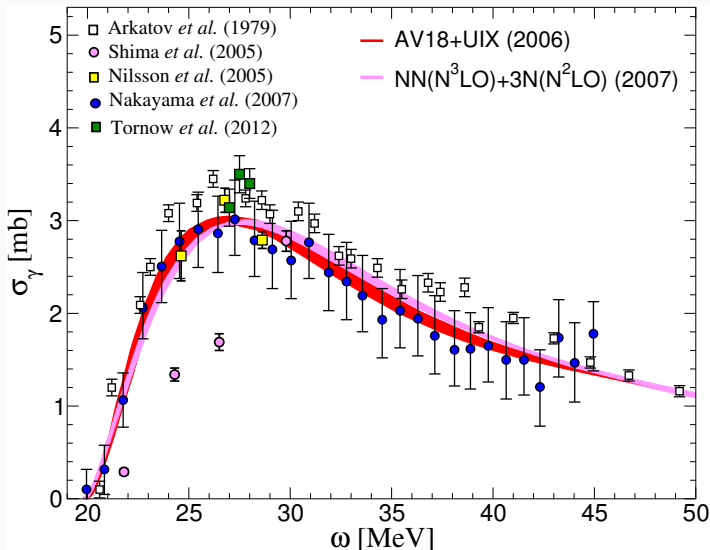


[A] R. Schiavilla, et. al., Phys. Rev. C 103, 054003 (2021).

^4He photoabsorption cross-section $\sigma_\gamma(\omega)$

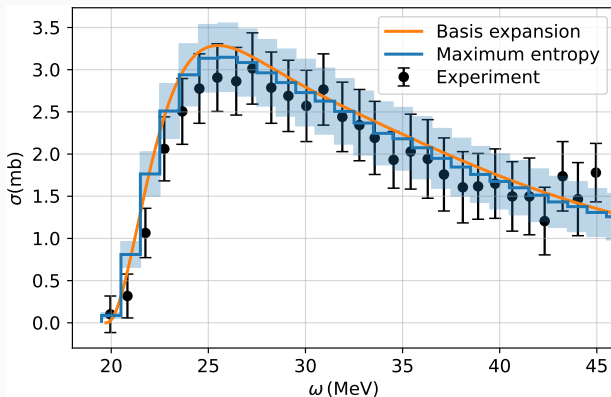


$$\sigma_\gamma(\omega) = 4\pi^2\alpha\omega\mathcal{R}_{D1}(\omega)$$





- Pionless EFT inspired potential [A]
- Comparison with highly accurate EIHH calculations
- LIT reproduced at percent-level accuracy
- Photoabsorption cross-section:
 - Good agreement with the experimental data



Summary



- **NQS+LIT** provides an accurate and scalable framework for nuclear responses
- Robust propagation of statistical and systematic uncertainties
- **Successful** description of ^2H and ^4He photoabsorption data
- Outlook:
 - Extension to nuclei with $A \leq 20$
 - Improved interactions (local χEFT)
 - Applications beyond nuclear physics (atoms, molecules)



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Thank you !