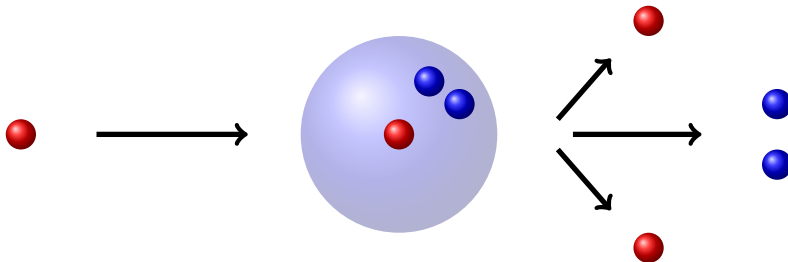
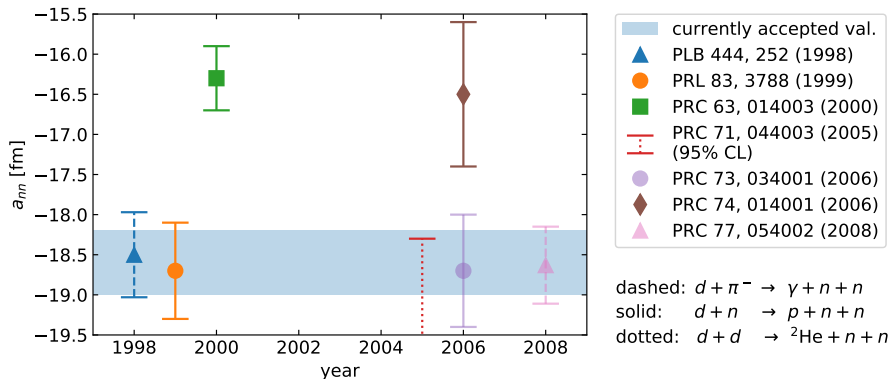


Two neutron distribution in the triton



[T. Kirchner, M. Göbel, HW. Hammer, Phys. Rev. C 111, 044002 (2025)]

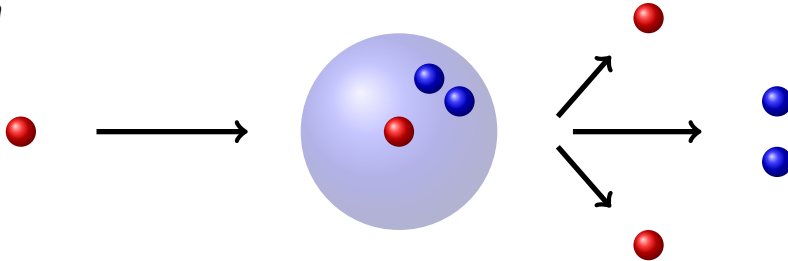
Motivation for determining a_{nn}



[Göbel et al., Phys.Rev.C 104 (2021) 2, 024001]

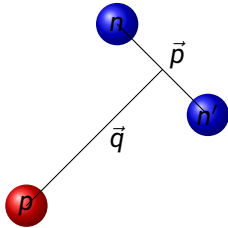
Hard knockout reaction

- goal: obtain nn scattering length a_{nn} with
 ${}^6\text{He}(p, p\alpha)nn$
 $t(p, 2p)nn$



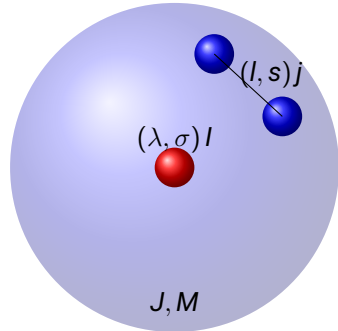
- assumptions:
 - ▣ sudden knockout $\rightarrow$ strong separation of proton and two neutrons in momentum space
 - ▣ nn final-state interactions dominate

Description of the three-body state

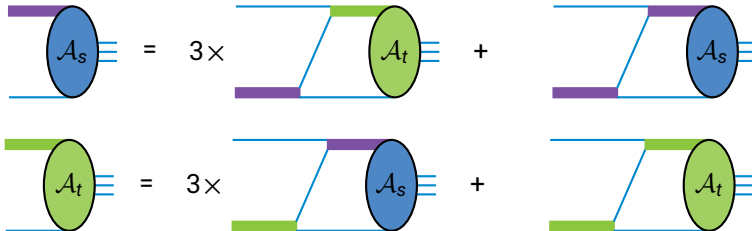


- $^1S_0(nn) + p$ nn singlet
- $^1S_0(np) + n$ np singlet
- $^3S_1(np) + n$ np triplet: deuteron

- $|\vec{p}, \vec{q}\rangle_i$, here $|\vec{p}, \vec{q}\rangle_p$
- use absolute momenta: $|p, q; \Omega\rangle_p$



- $|\Psi\rangle = \sum_i |\psi_i\rangle$ with $i \in \left\{ \begin{smallmatrix} (0) \\ nn' \end{smallmatrix}, \begin{smallmatrix} (0) \\ np \end{smallmatrix}, \begin{smallmatrix} (1) \\ np \end{smallmatrix} \right\}$
- Faddeev amplitudes $G_0 t_i |F_i\rangle = |\psi_i\rangle = G_0 V_i |\Psi\rangle$
- coupled equations $|F_i\rangle = \sum_{j \neq i} G_0 t_j |F_j\rangle$ for the interaction channels $\rightarrow t_{nn'}^{(0)} \quad t_{np}^{(0)} \quad t_{np}^{(1)}$



- 3-body force $h(\Lambda)$ fixed to triton binding energy $B = 8.48 \text{ MeV}$

Parameterizing the t-matrix



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effective range expansion

$$\tau_i(p) = \frac{1}{2\pi^2 M} \frac{1}{-1/a_i + r_i p^2/2 - ip}$$

mod. effective range expansion around pole

$$\tau_d(p) = \frac{1}{2\pi^2 M} \frac{1}{-\gamma_d + \rho_d(\gamma_d^2 + p^2)/2 - ip}$$

Parameterizing the t-matrix



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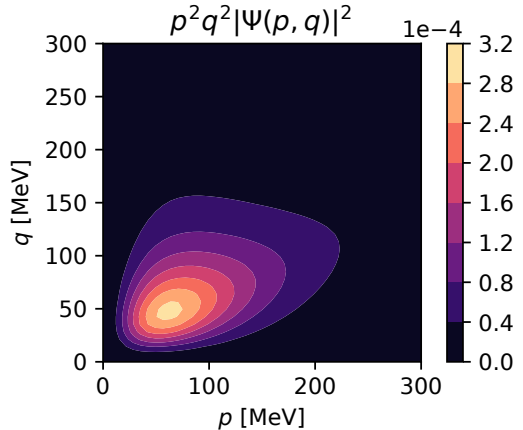
$$\tau_d(p) = \frac{1}{2\pi^2 M} \frac{1}{-\gamma_d + \rho_d(\gamma_d^2 + p^2)/2 - ip}$$

unphysical poles around momenta of $p \approx 150$ MeV:

$$\frac{1}{-1/a_i + r_i p^2/2 - ip} = \underbrace{\frac{1}{-1/a_i - ip}}_{\text{LO t-matrix}} \left(1 - \frac{r_i p^2/2}{-1/a_i - ip} + \dots \right)$$

Obtaining the density distribution

- $|\Omega\rangle_i = |(l, s)j, (\lambda, \sigma)l; J, M\rangle_i$
- wave function in partial wave Ω :
 $\Psi_{p;\Omega}(p, q) = {}_p\langle p, q; \Omega | \Psi \rangle_p$

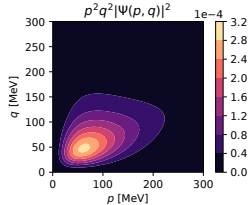


Obtaining the density distribution

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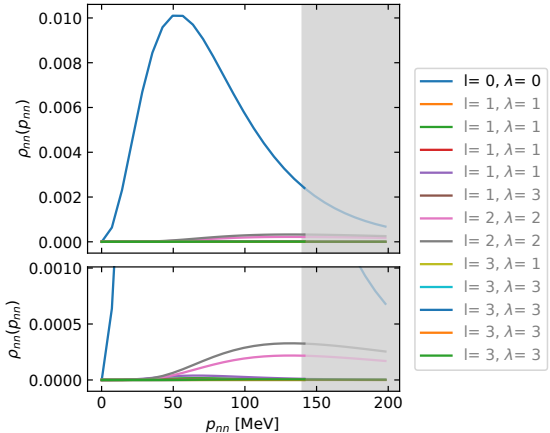
- wave function in partial wave Ω :

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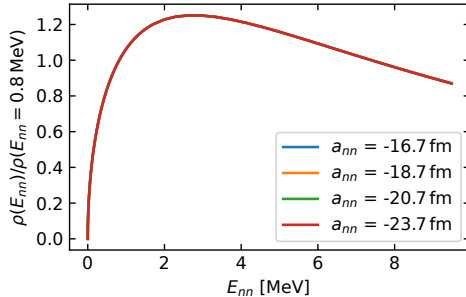
- density distribution

$$\rho_{nn}(p) = \sum_{\Omega} \int dq p^2 q^2 |\Psi_{p;\Omega}(p, q)|^2$$



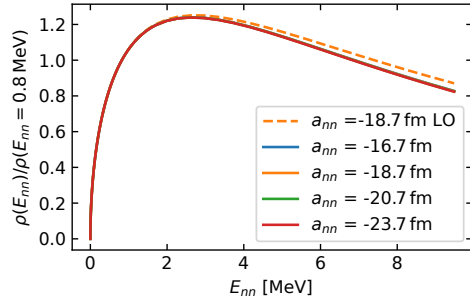
Results for the groundstate

LO groundstate



$a_{np} [\text{fm}]$	$\gamma_d [\text{MeV}]$
-23.736	45.7021

NLO groundstate



$r_{nn} [\text{fm}]$	$r_{np} [\text{fm}]$	$\rho_d [\text{fm}]$
2.84	2.704	1.752

nn-final-state interactions included via the t-matrix

$$\begin{aligned}\psi_p^{(wFSI)}(p, q) = & \psi_p(p, q) + \frac{2}{\pi} \underbrace{\frac{1}{a_{nn}^{-1} - \frac{r_{nn}}{2} p^2 + ip}}_{\text{nn t-matrix}} \\ & \times \left[\int_0^\Lambda dp' \frac{p'^2 \psi_p(p', q) - p^2 \psi_p(p, q)}{p^2 - p'^2} - \left(\frac{i\pi}{2} - \frac{1}{2} \ln \left(\frac{\Lambda + p}{\Lambda - p} \right) \right) p \psi_p(p, q) \right]\end{aligned}$$

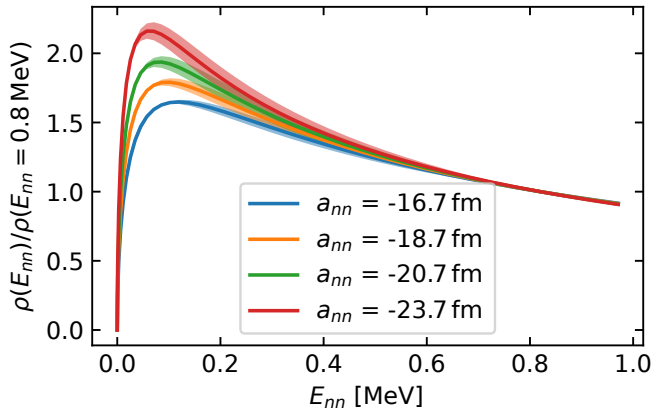
FSI results with uncertainties



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- NLO groundstate curve with nn FSI applied
- conservative error band given by the difference to the LO curve
- all curves normed at $E_{nn} = 0.8 \text{ MeV}$

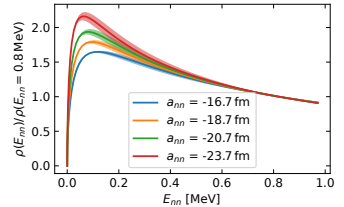
possible to distinguish a_{nn}
within 2 fm



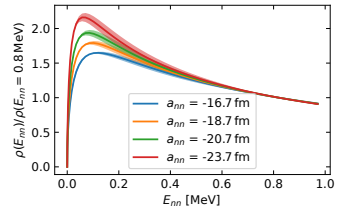
[T. Kirchner, M. Göbel, HW. Hammer,
Phys. Rev. C 111, 044002 (2025)]

Summary & Outlook

- calculate NLO groundstate curve for the triton
- include nn-FSI
- distinguish peaks for different a_{nn} within the uncertainties

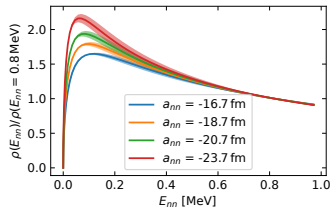


- calculate NLO groundstate curve for the triton
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- Going to NNLO: include another three-body force (e.g. determine from radius)

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- Going to NNLO: include another three-body force (e.g. determine from radius)

Thank you for your attention!