

EXPLORING PROTO-NEUTRON STAR MINIMUM MASS LIMITS THROUGH CHIRALLY CONSTRAINED NUCLEAR EOS

Selina Kunkel, Stephan Wüsteb and Jürgen Schaffner-Bielich

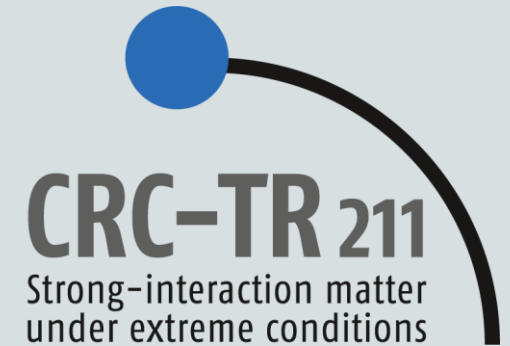
Institut für Theoretische Physik

Goethe Universität Frankfurt am Main

EMMI Workshop GSI Darmstadt 2025



Talk based on *Phys. Rev. C* 111, 035807 (2025)



EVOLUTION OF MASSIVE STARS

$<8M_{\odot} \rightarrow$ White Dwarf



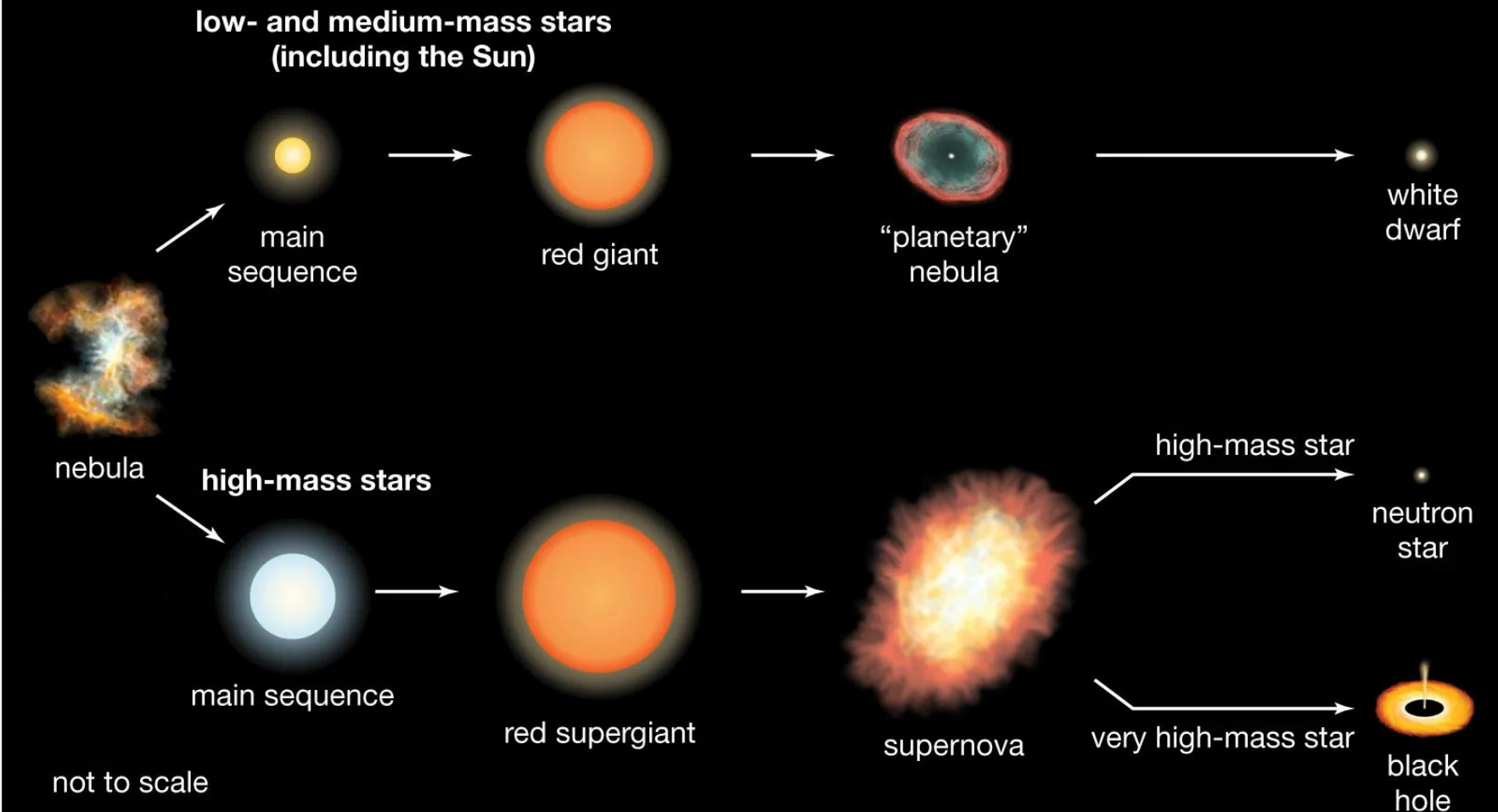
$8M_{\odot} - 25M_{\odot} \rightarrow$ Neutron Star



$> 25M_{\odot} \rightarrow$ Black Hole



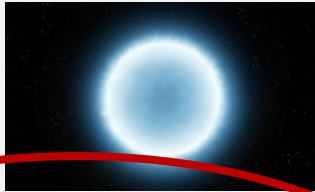
Stellar evolution



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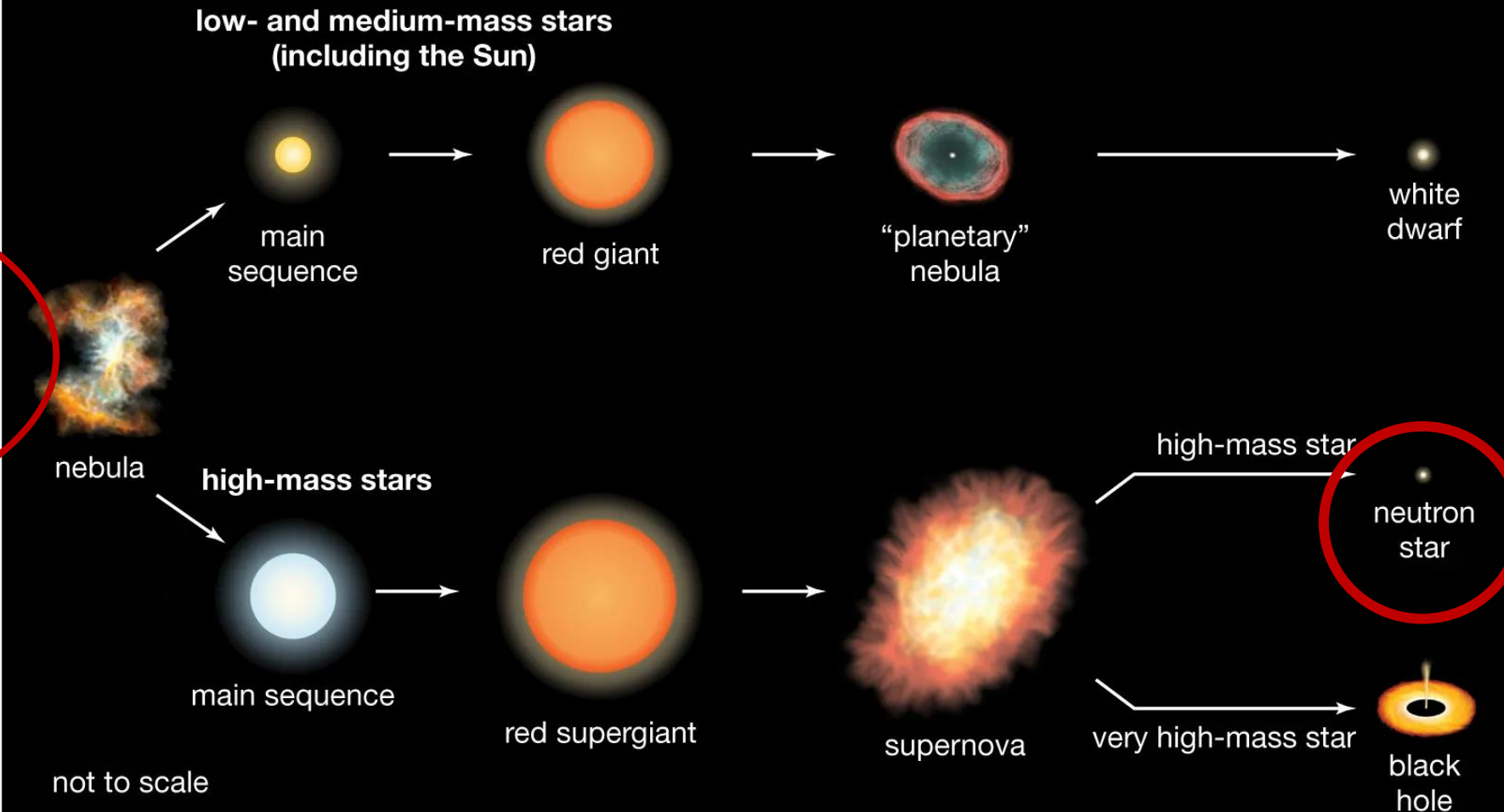
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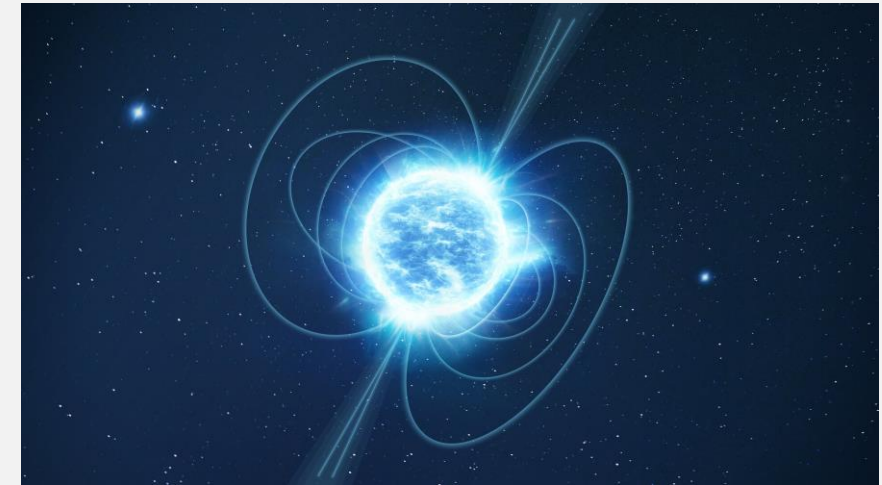


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<https://www.britannica.com/science/star-astronomy/Star-formation-and-evolution>

NEUTRON STARS

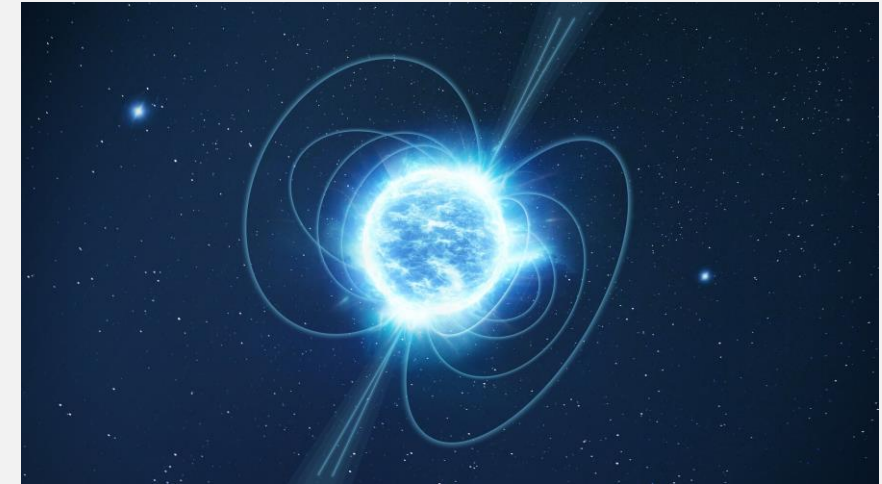
	Neutron Star	Earth
Mass	$1.4M_{\odot} \approx 10^{30}kg$	$10^{24}kg$
Radius	$12km$	$6371km$
Density	$10^{17}kg/m^3$	$10^5kg/m^3$
Surface gravity	$10^{12}m/s^2$	$9.81m/s^2$
Rotation period	$1.4ms - 30s$	$24h$
Magnetic field	$10^4T - 10^{11}T$	$25\mu T - 65\mu T$
Central temperatures	10^9K	10^3K



https://www.esa.int/ESA_Multimedia/Images/2024/03/What_is_a_neutron_star

NEUTRON STARS

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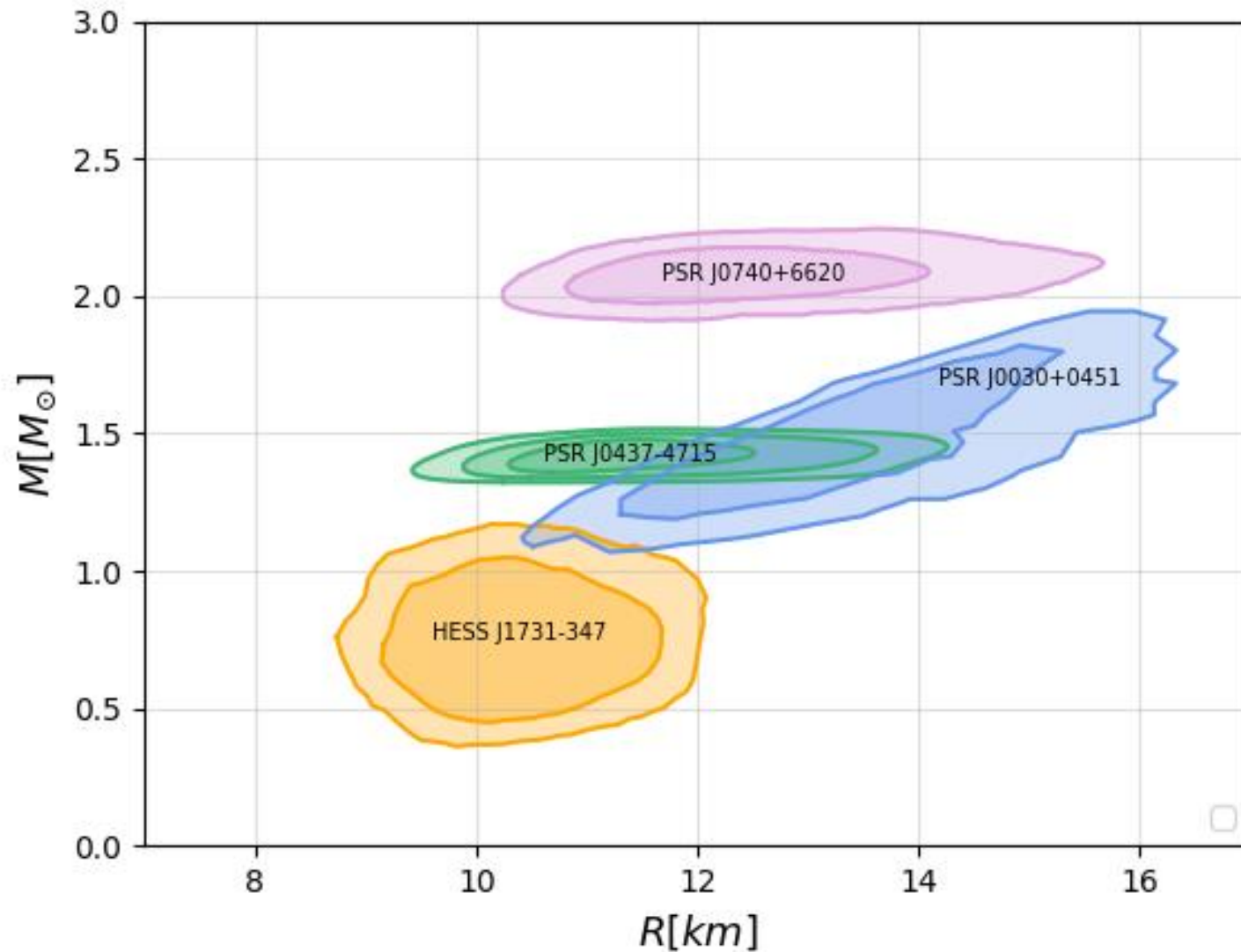


https://www.esa.int/ESA_Multimedia/Images/2024/03/What_is_a_neutron_star

*Neutron Stars are the densest objects in the universe which we can physically describe **and** observe!*

WHY STUDY MINIMAL MASSES OF PROTO-NEUTRON STARS?

Mass-Radius measurements of astrophysical objects

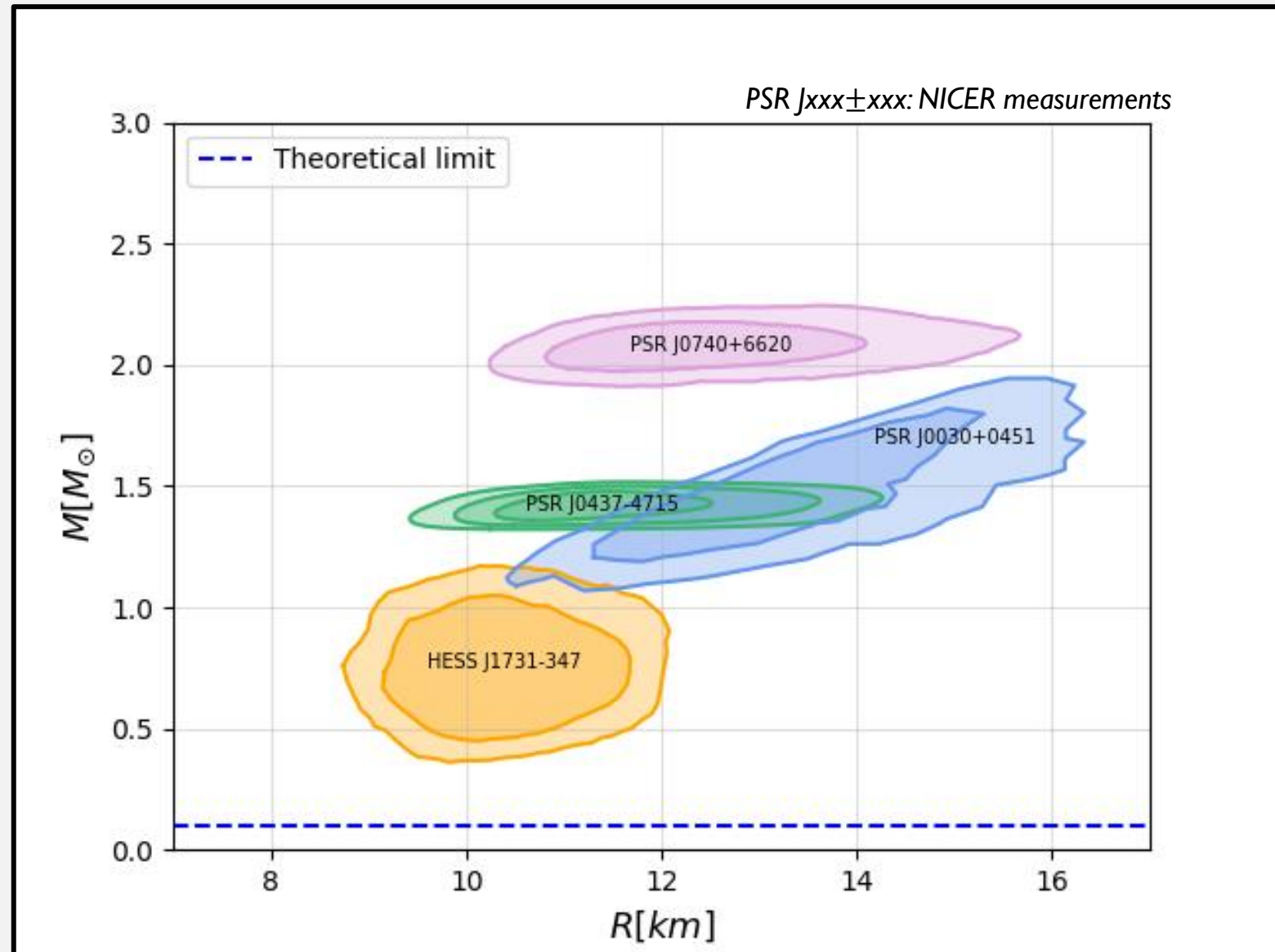


WHY STUDY MINIMAL MASSES OF PROTO-NEUTRON STARS?

Theory:

Lower mass limit of a cold neutron star at

$$M_{\min} \sim 0.1 M_{\odot}$$



WHY STUDY MINIMAL MASSES OF PROTO-NEUTRON STARS?

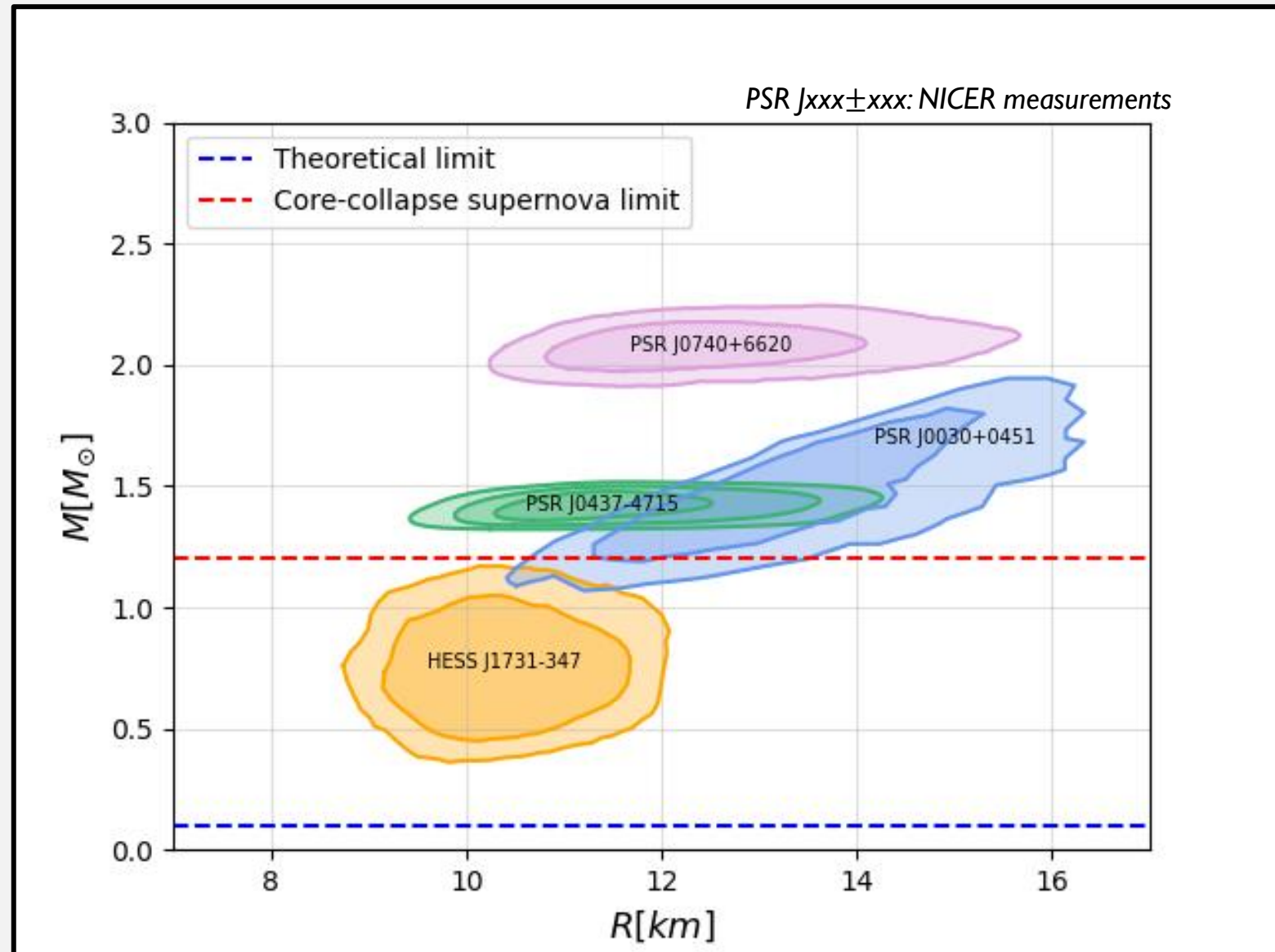
Theory:

Lower mass limit of a cold neutron star at $M_{min} \sim 0.1 M_{\odot}$

Core-collapse supernova simulations:

- Hot PNS (proto-neutron star) with $M_{min} \sim 1.2 M_{\odot}$ is formed
- Cools down to cold PNS

HESS J1731-347 too light!



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Lower mass limit of a cold neutron star at $M_{min} \sim 0.1 M_{\odot}$

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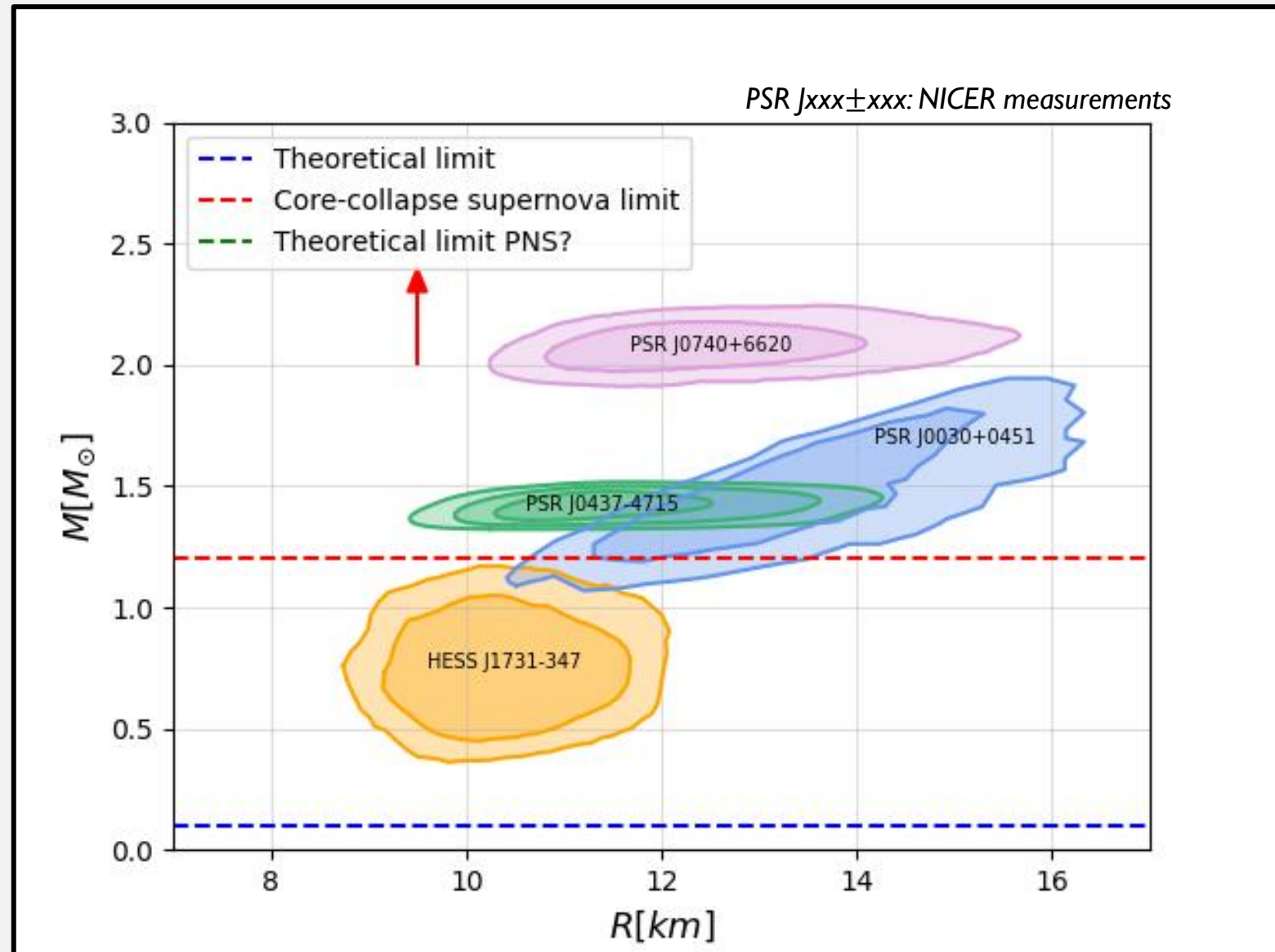
- Hot PNS (proto-neutron star) with $M_{min} \sim 1.2 M_{\odot}$ is formed
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HESS J1731-347 too light!

Theory for proto-neutron star:

- Mass constraint of hot PNS gives constraint on mass of cold PNS!

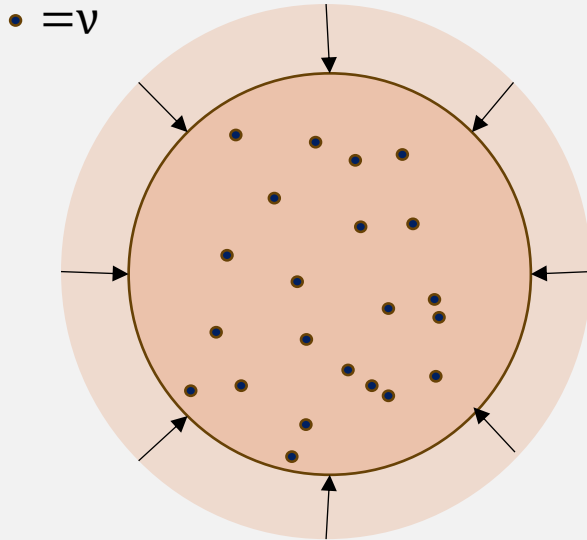
Is HESS J1731-347 still a potential neutron star candidate?



PROTO-NEUTRON STAR EVOLUTION

Collapse

$t = 0\text{s}$

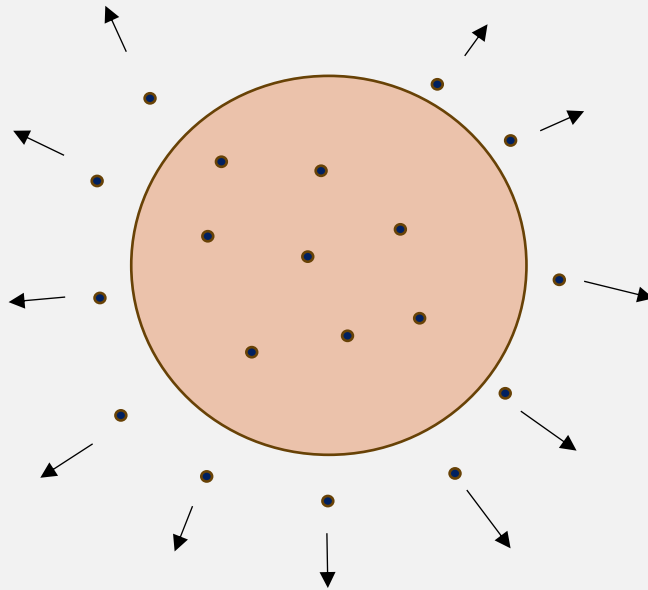


$$s = 1$$

$$Y_L = 0.4$$

Deleptonization

$t = 0.5\text{-}1\text{s}$

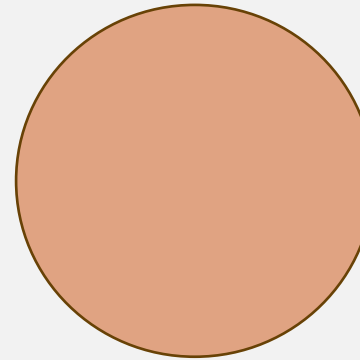


$$s = 1 - 3$$

$$Y_L = 0 - 0.4$$

Cooling

$t = 15\text{s}$

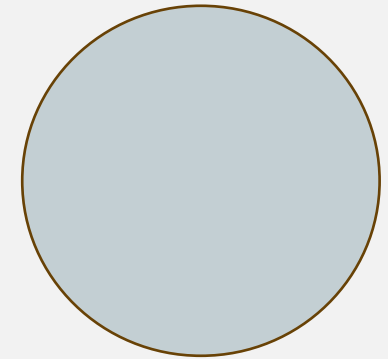


$$s = 2$$

$$Y_V = 0$$

Cold Neutron Star

$t = 10^6\text{yr}$



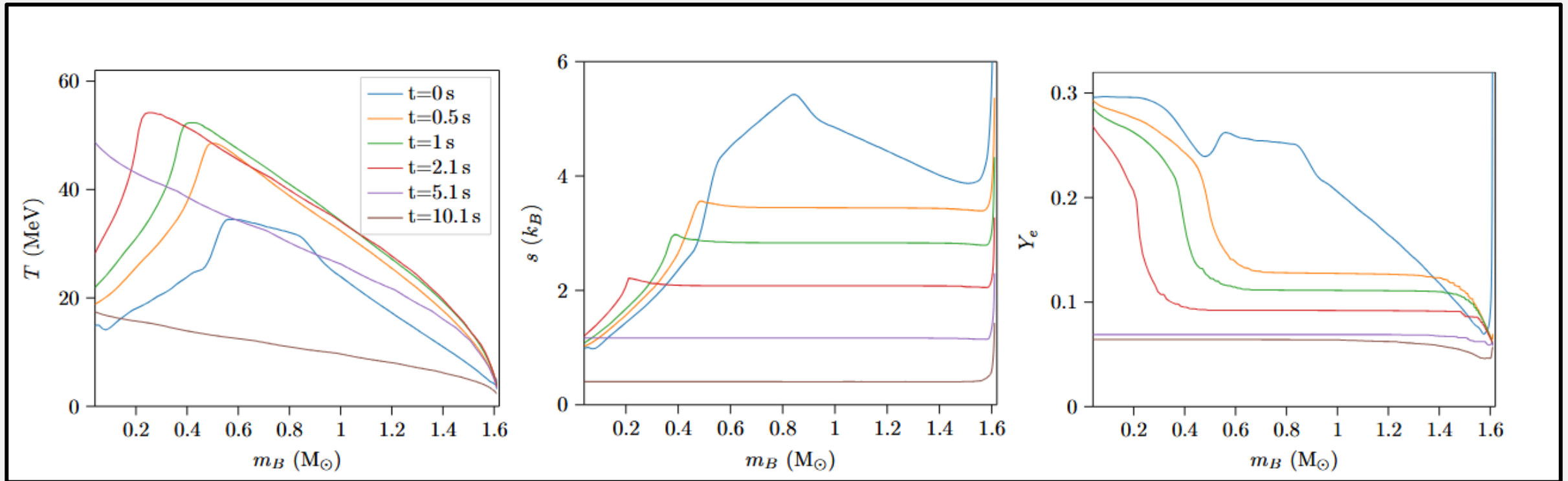
$$s = 0$$

$$Y_V = 0$$

Entropy per baryon: $s = \frac{S}{A}$

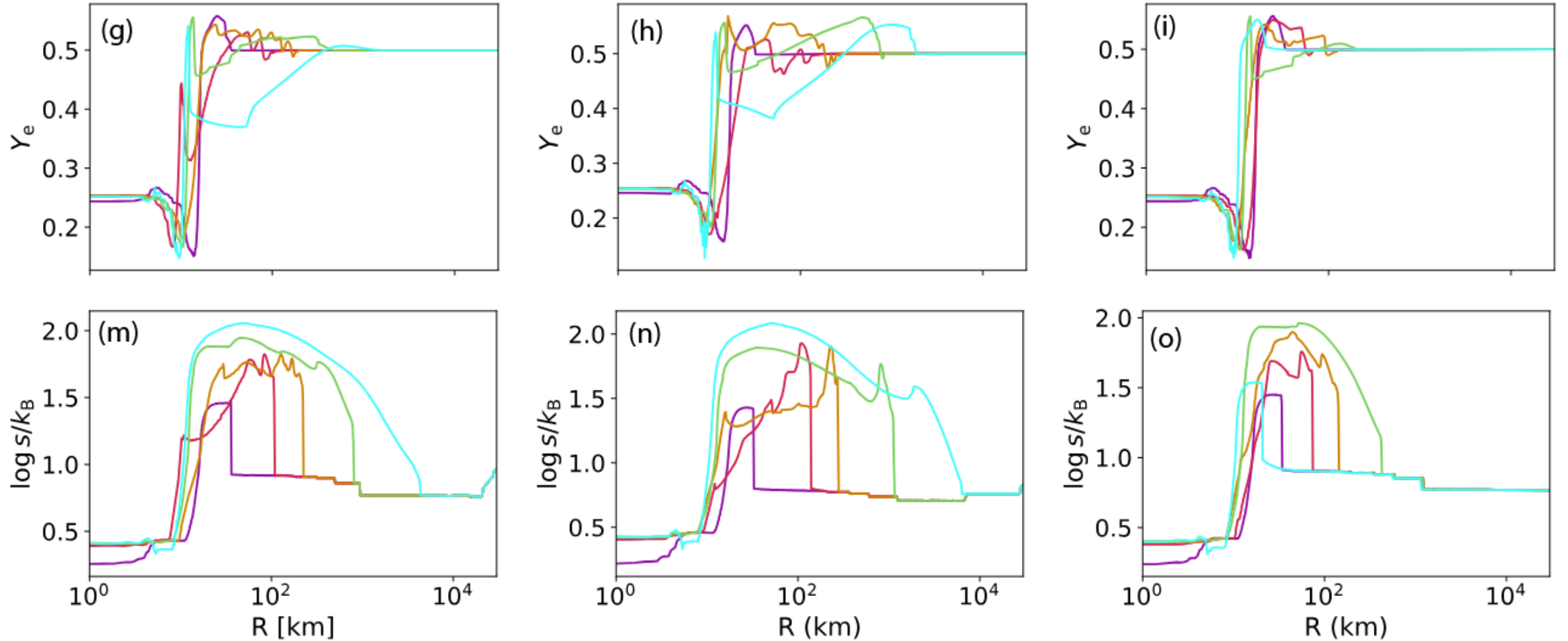
Lepton fraction: $Y_L = \frac{(n_e - n_{e^+}) + (n_{\nu_e} - n_{\bar{\nu}_e})}{(n_n - n_{\bar{n}}) + (n_p - n_{\bar{p}})}$

PROTO-NEUTRON STAR SIMULATIONS



Pascal et al., *Mon. Not. R. Astron. Soc.* 511, 1, 356–370 (2022)

PROTO-NEUTRON STAR SIMULATIONS

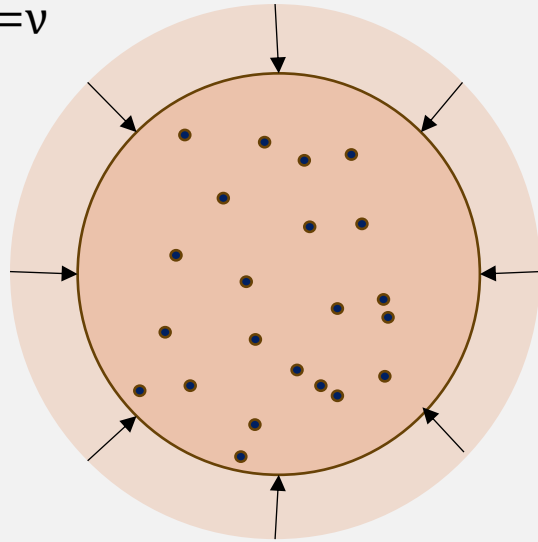


PROTO-NEUTRON STAR EVOLUTION

Collapse

$t = 0\text{s}$

• = ν

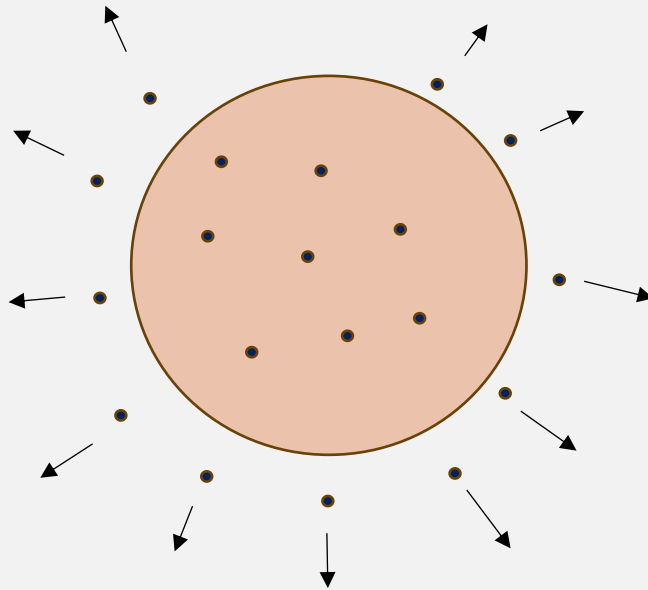


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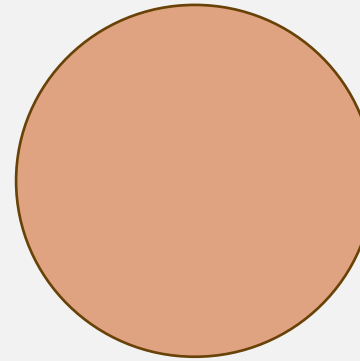


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Cooling

$t = 15\text{s}$

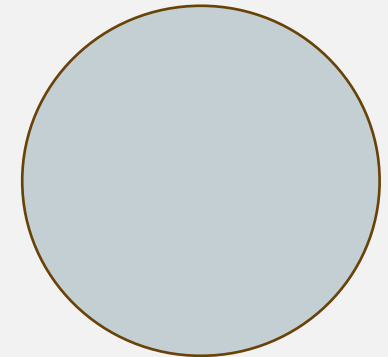


$$s = 2$$

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Cold Neutron Star

$t = 10^6\text{yr}$



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Entropy per baryon: $s = \frac{S}{A}$

Lepton fraction: $Y_L = \frac{(n_e - n_{e^+}) + (n_{\nu_e} - n_{\bar{\nu}_e})}{(n_n - n_{\bar{n}}) + (n_p - n_{\bar{p}})}$

WHY STUDY IT WITH CHIRAL EFT?

Low masses of neutron stars → Low central densities → Framework which describes matter at low densities

Chiral EFT

Description of nucleon
interaction at low
energies

Relevant degrees of
freedom are nucleons
and pions

Based on chiral
symmetry and chiral
symmetry breaking

Pion as goldstone
boson

Power counting
scheme

Expansion in $\left(\frac{p}{\Lambda_\chi}\right)^v$

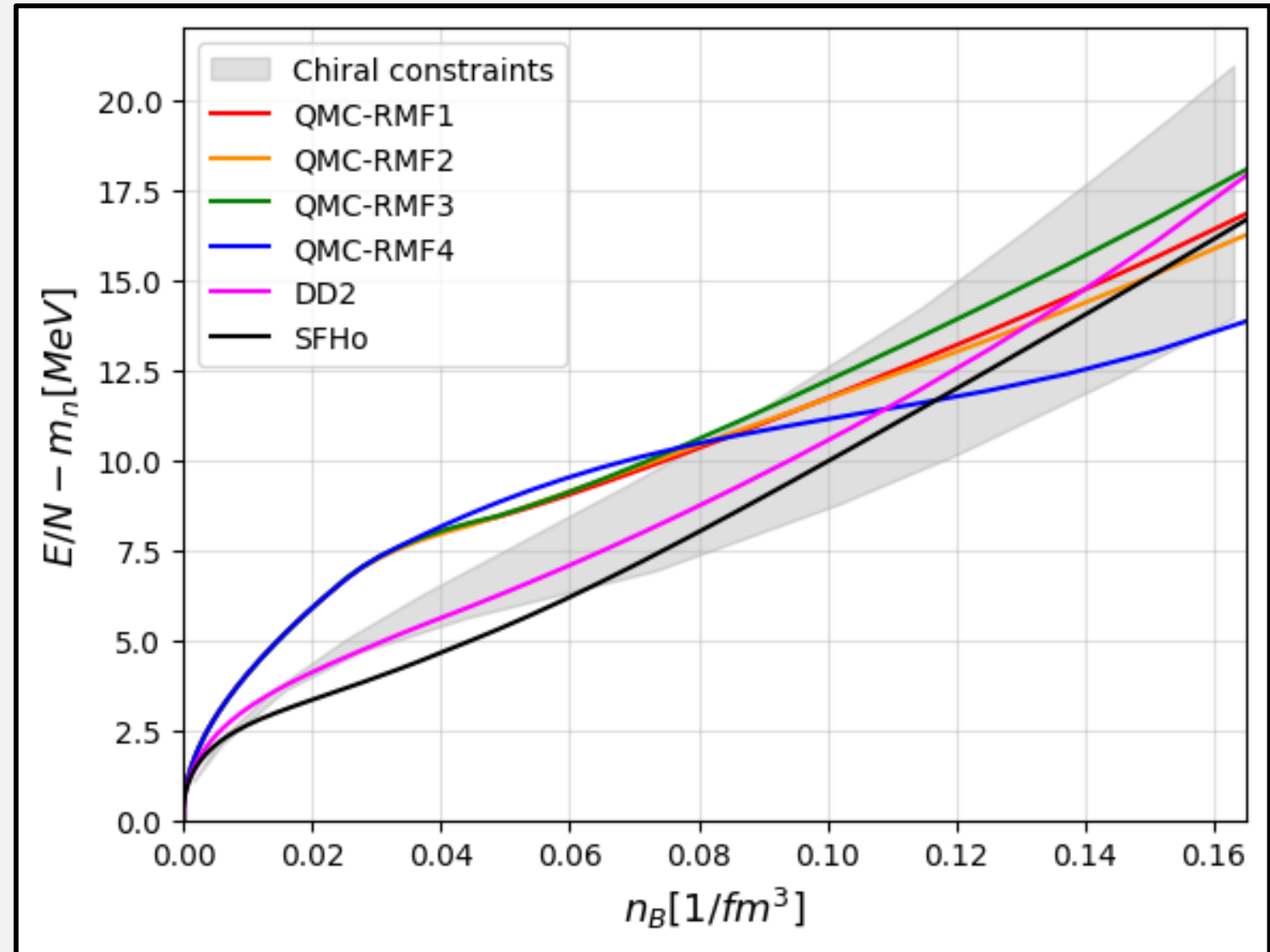
EQUATIONS OF STATE

QMC-RMF (Alford et al., *Phys. Scripta* 98, 125302 (2023))

DD2 (Typel et al., *Phys. Rev. C* 81, 015803 (2010))

SFH0 (Steiner, Hempel & Fischer, *Astrophys. J.* 774, 17 (2013))

- Fulfill chiral EFT constraints for the binding energy at $T = 0$
- Relativistic Mean Field Models



Chiral constraints taken from: Krüger et al., *Phys. Rev. C* 88, 025802 (2013)

TOLMAN-OPPENHEIMER-VOLKOFF EQUATIONS

Structure of a static, spherically symmetric star in General Relativity *

- Mass conservation

$$\overset{\text{mass}}{\frac{dm}{dr}} = 4\pi r^2 \overset{\text{energy density}}{\epsilon}$$

$\nwarrow$
radius

- Hydrostatic equilibrium

$$\overset{\text{pressure}}{\frac{dP}{dr}} = -\frac{Gm}{r^2} \epsilon \left(1 + \frac{P}{\epsilon c^2} \right) \left(1 + \frac{4\pi r^3 P}{mc^2} \right) \left(1 - \frac{2Gm}{rc^2} \right)^{-1}$$

*R. C. Tolman, *Phys. Rev.* **55**, 364 (1939);
J. R. Oppenheimer & G. M. Volkoff, *Phys. Rev.* **55**, 374 (1939)

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Input: Pressure and Energy Density

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Input: Pressure and Energy Density



Output: Mass and Radius

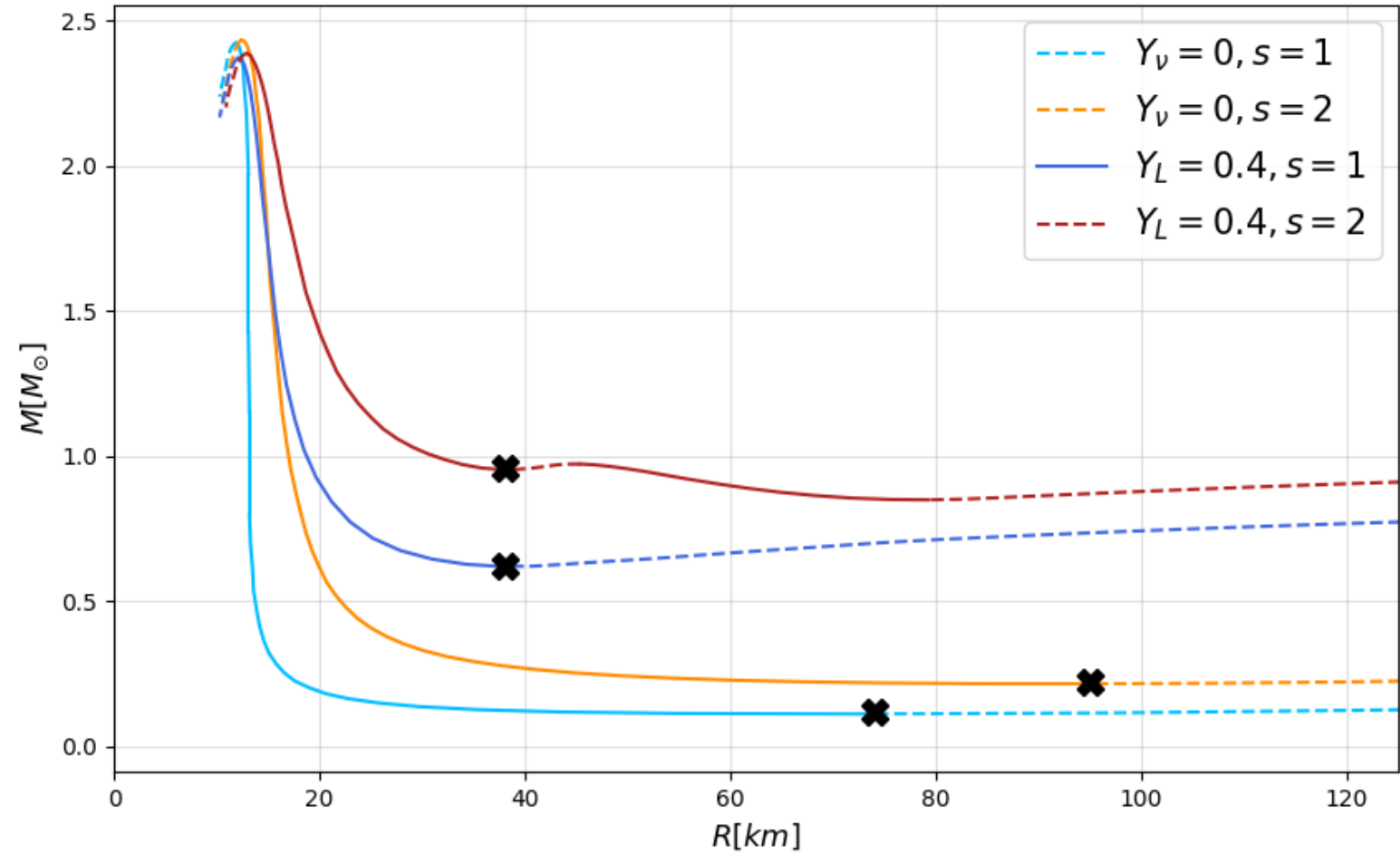
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RESULTS – MASS-RADIUS CURVES

Solid: Stable

Dashed: Unstable

✱ = minimum mass configuration



RESULTS – MASS-RADIUS CURVES

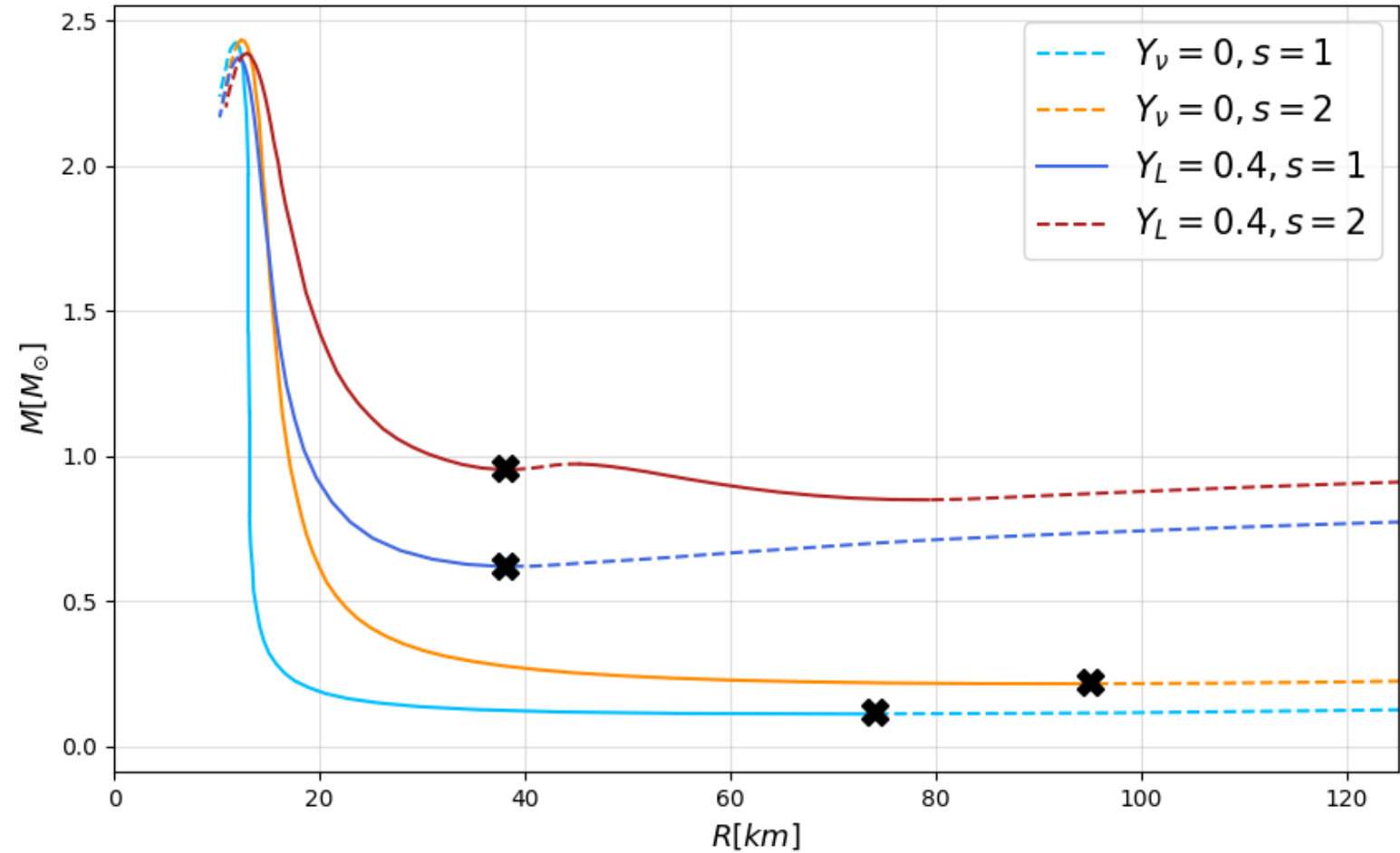
$$s \uparrow \rightarrow M_{\min} \uparrow$$

Additional thermal pressure!

Solid: Stable

Dashed: Unstable

✱ = *minimum mass configuration*



RESULTS – MASS-RADIUS CURVES

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Additional thermal pressure!

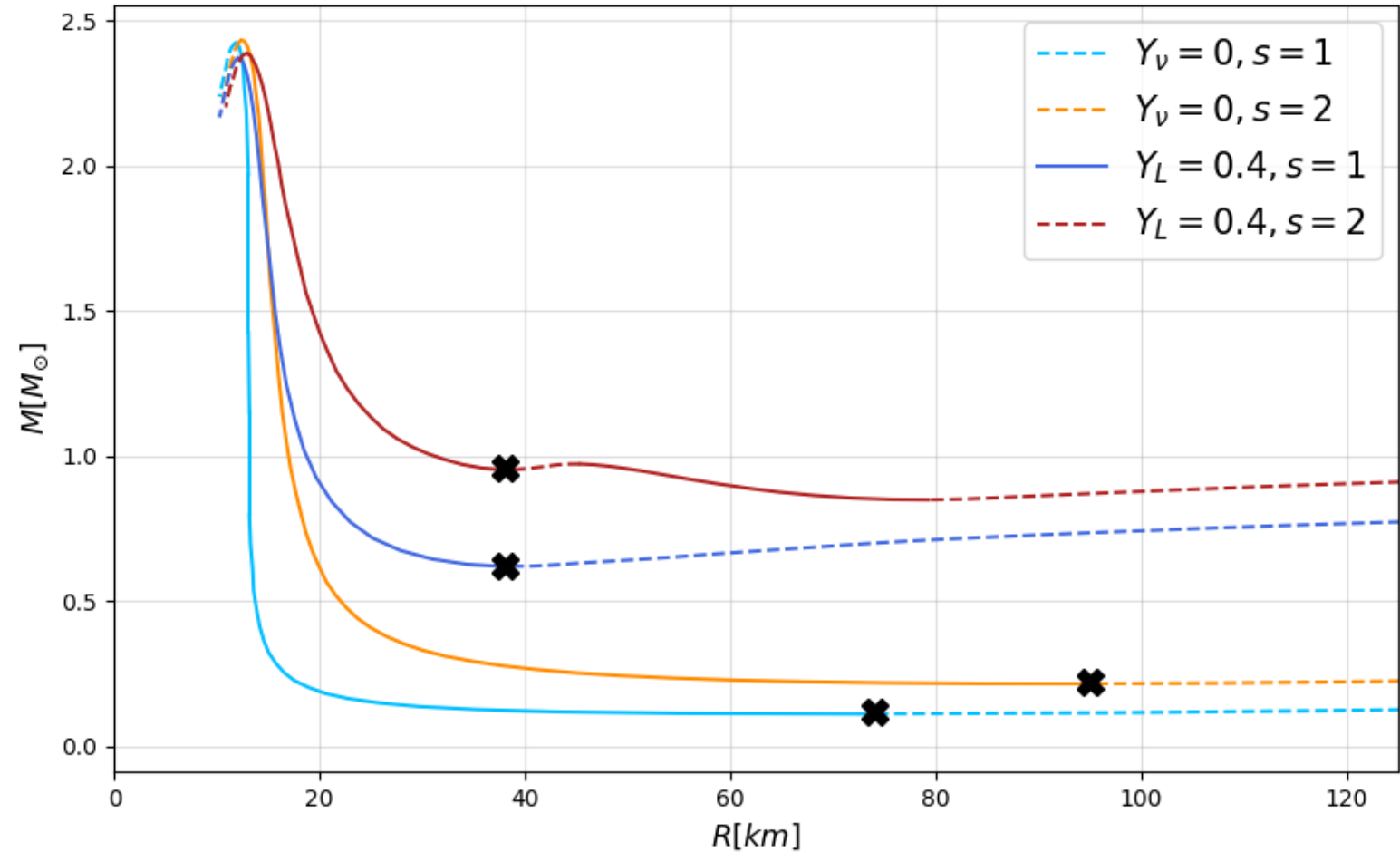
$$Y_L \uparrow \rightarrow M_{\min} \uparrow$$

Additional lepton pressure!

Solid: Stable

Dashed: Unstable

$\times$ = *minimum mass configuration*



RESULTS – MASS-RADIUS CURVES

$$s \uparrow \rightarrow M_{\min} \uparrow$$

Additional thermal pressure!

$$Y_L \uparrow \rightarrow M_{\min} \uparrow$$

Additional lepton pressure!

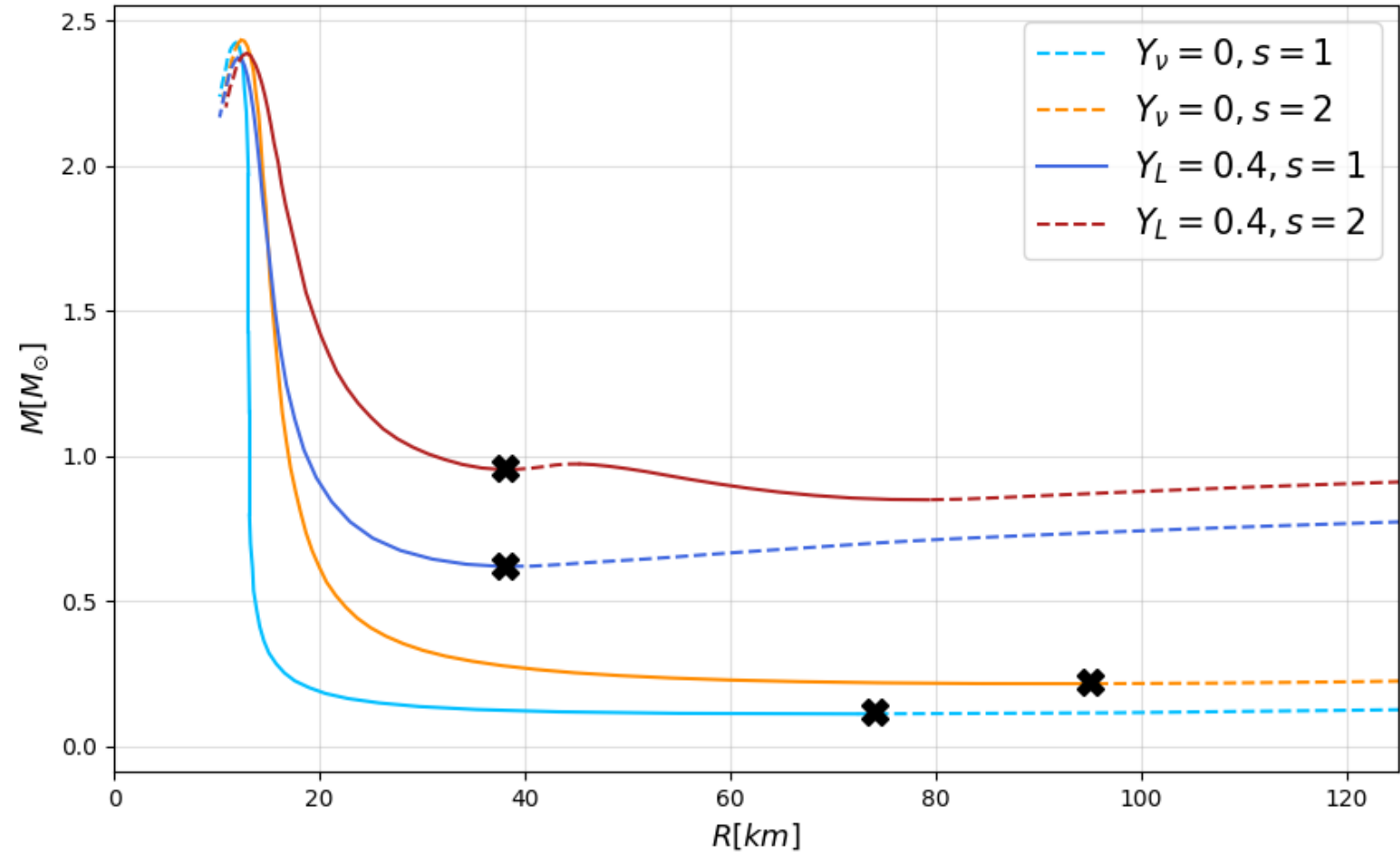
$$Y_L \uparrow \rightarrow M_{\max} \downarrow$$

Less neutron pressure!

Solid: Stable

Dashed: Unstable

$\times$ = *minimum mass configuration*



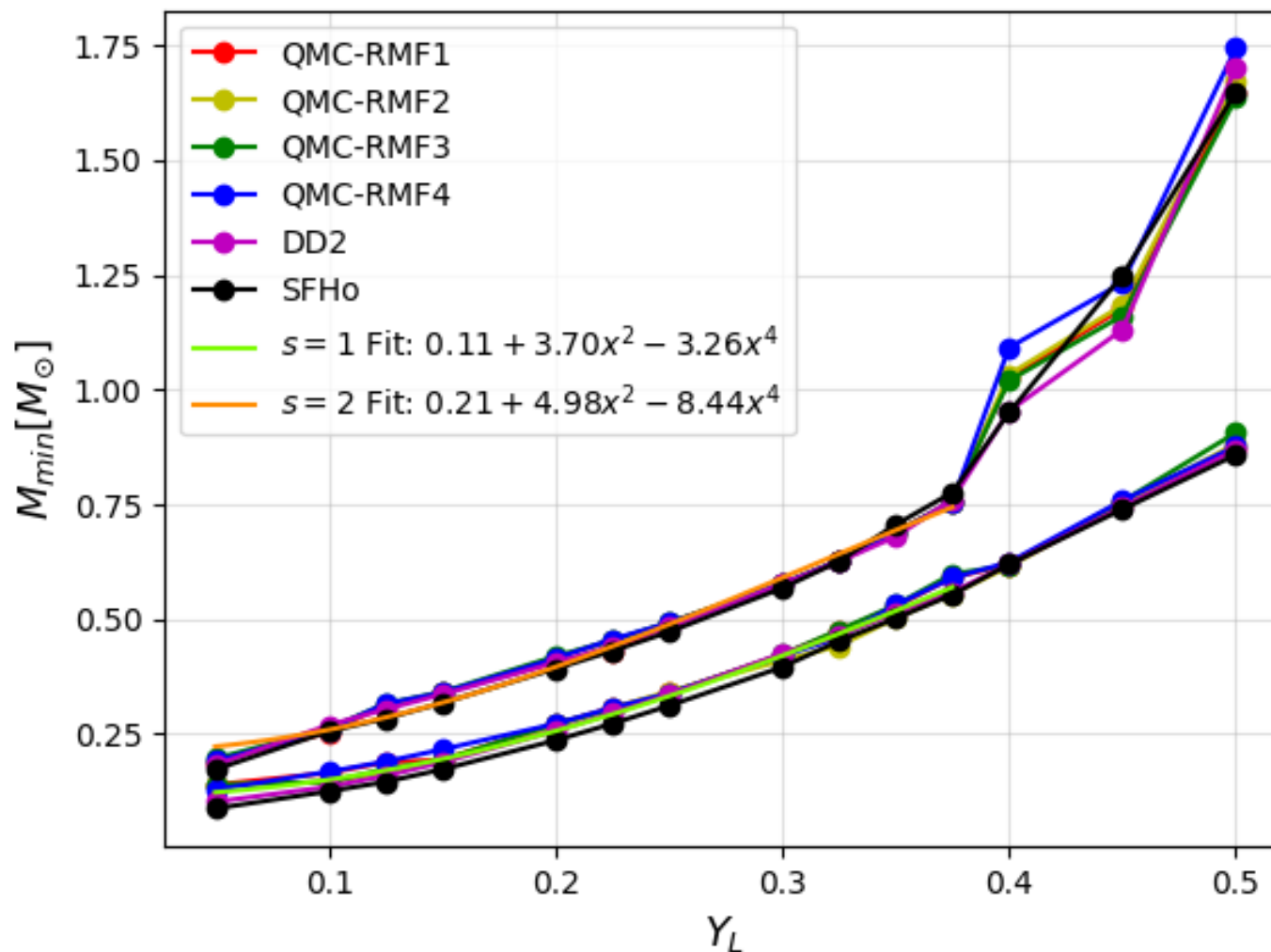
MINIMAL MASSES, RADII AND CENTRAL DENSITIES

$Y_L = 0.4, s = 1$				$Y_v = 0, s = 2$			
EOS	$M_{\min} (M_{\odot})$	R (km)	n_B (fm ⁻³)	EOS	$M_{\min} (M_{\odot})$	R (km)	n_B (fm ⁻³)
QMC-RMF1	0.617	37.9	0.149	QMC-RMF1	0.238	101	0.00821
QMC-RMF2	0.615	36.6	0.163	QMC-RMF2	0.238	101	0.00823
QMC-RMF3	0.618	38.0	0.145	QMC-RMF3	0.238	101	0.00821
QMC-RMF4	0.624	36.1	0.164	QMC-RMF4	0.238	101	0.00821
DD2	0.620	38.3	0.134	DD2	0.216	94.9	0.0121
SFHo	0.613	37.0	0.153	SFHo	0.195	85.0	0.0161

- Captured neutrinos before the explosion result in: **Bigger minimum mass, smaller radius, higher densities**
→ **more compact**
- Loss of neutrinos after the explosion result in: **Smaller minimum mass, bigger radius, lower densities**
→ **less compact**

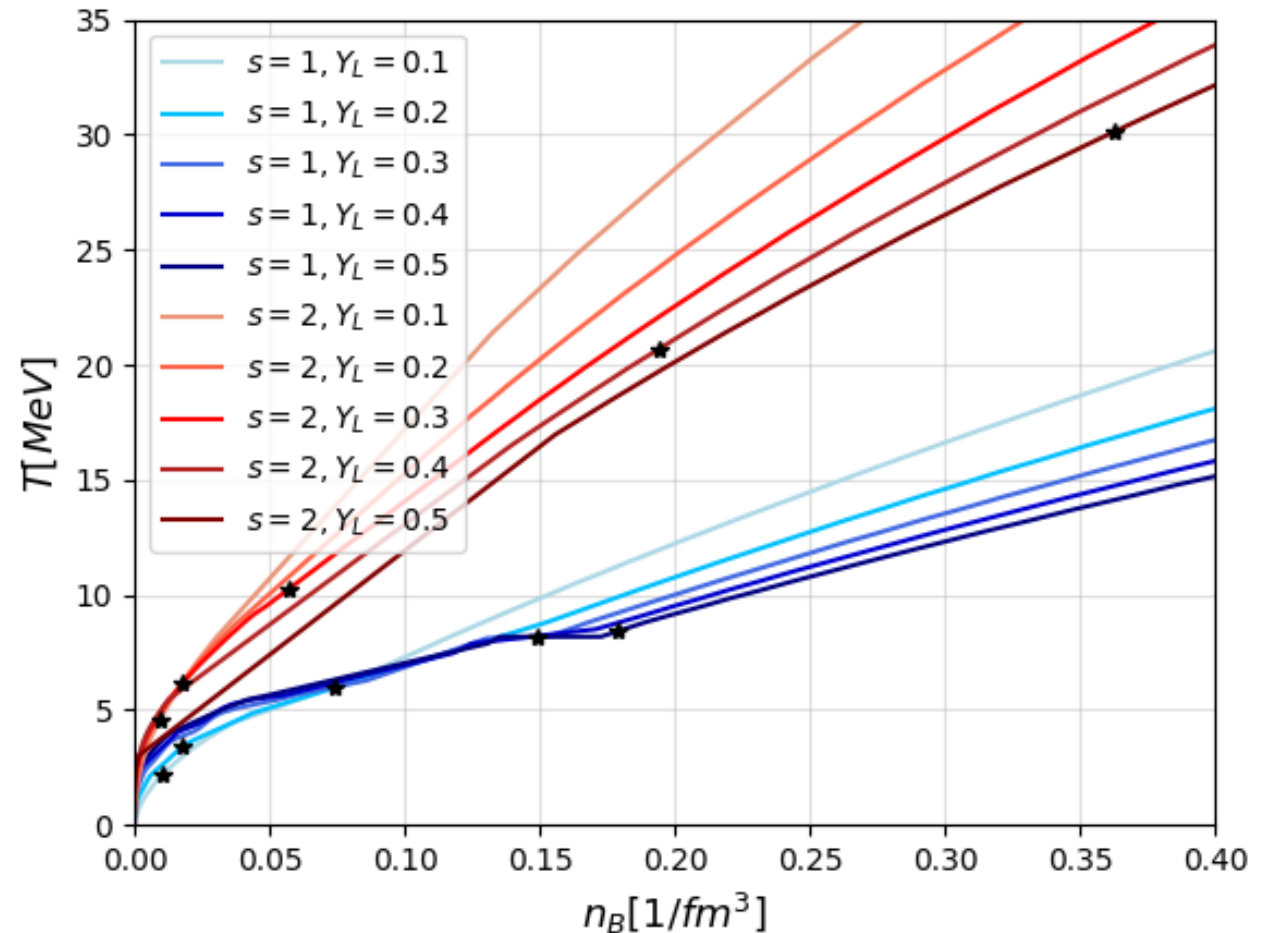
RESULTS – MINIMAL MASSES FOR DIFFERENT LEPTON FRACTIONS

- Increasing minimal mass with increasing lepton fraction
- Higher minimal masses for higher entropy
- Non-linear dependency
- EoS show same behaviour → Universal relation for EoS fulfilling chiral EFT
- Common fit for all EoS up to $Y_L \leq 0.375$

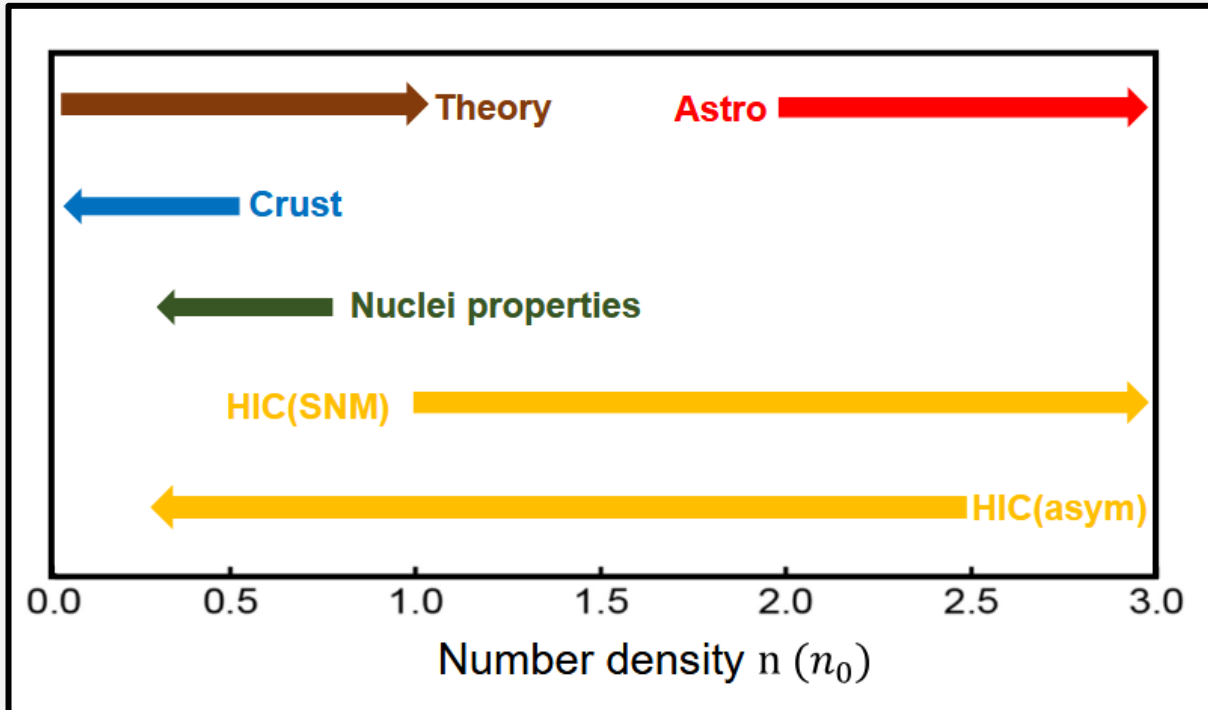


RESULTS - ISENTROPES

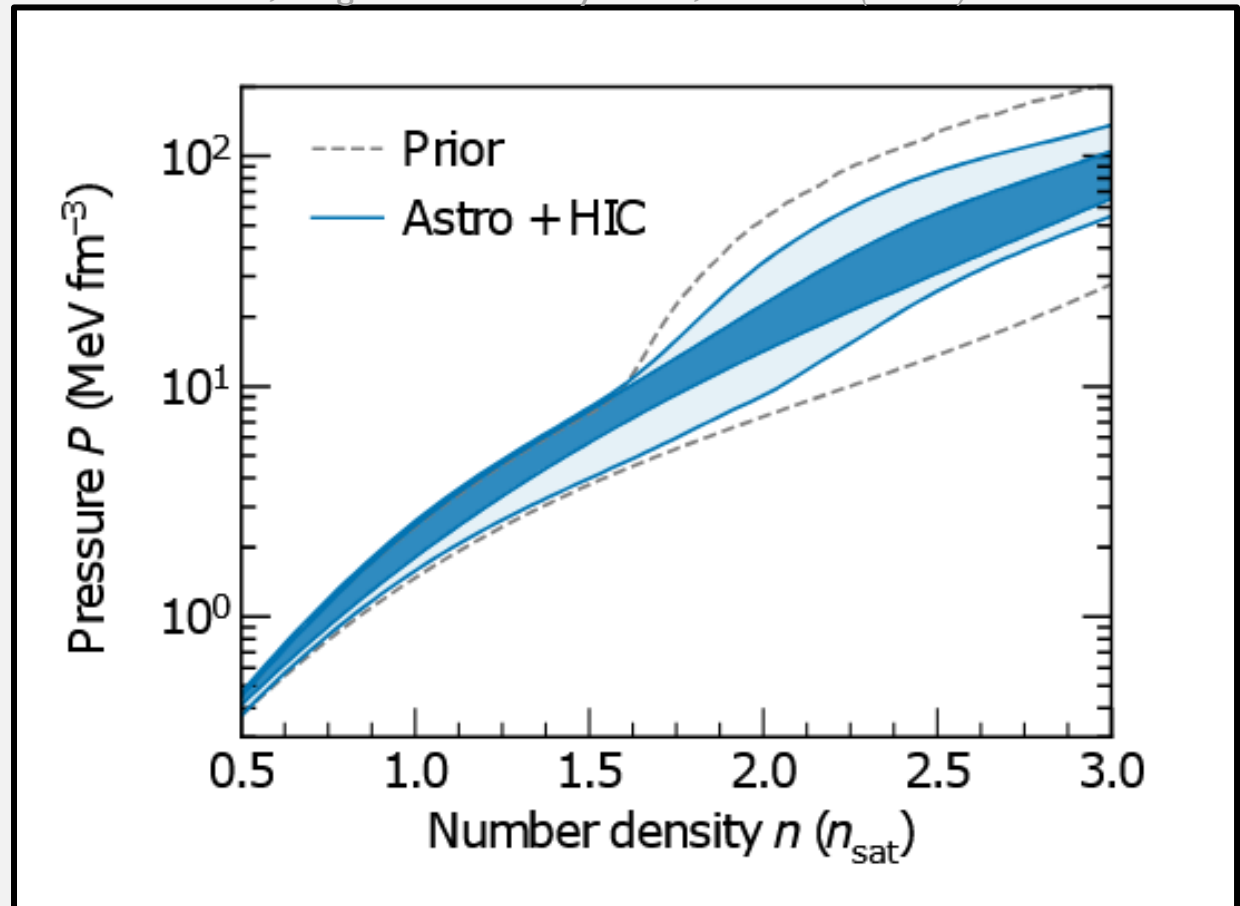
- High Y_L and high $s \rightarrow$ high n_B
- Increase in temperature and density for $(Y_L = 0.4, s = 2)$ and $(Y_L = 0.5, s = 2)$ is huge
- At those densities we are not in chiral band anymore
 $\rightarrow$ noticable differences in M_{min}



CONSTRAINTS FROM HIC AND ASTRO



Heavy-Ion Collisions and astrophysical constraints both describe matter at $2.0 - 3.0n_0$



Combination of Heavy-Ion Collisions with astrophysical constraints narrow down pressure region

- PNS: $T \approx 10 - 50 \text{ MeV}$ & $Y_p \approx 0.1 - 0.4$
- HIC create similar conditions after collision



Heavy-Ion Collisions probe regions in temperature and proton fraction that are found in Proto-Neutron Star matter and can therefore help to constrain PNS properties!

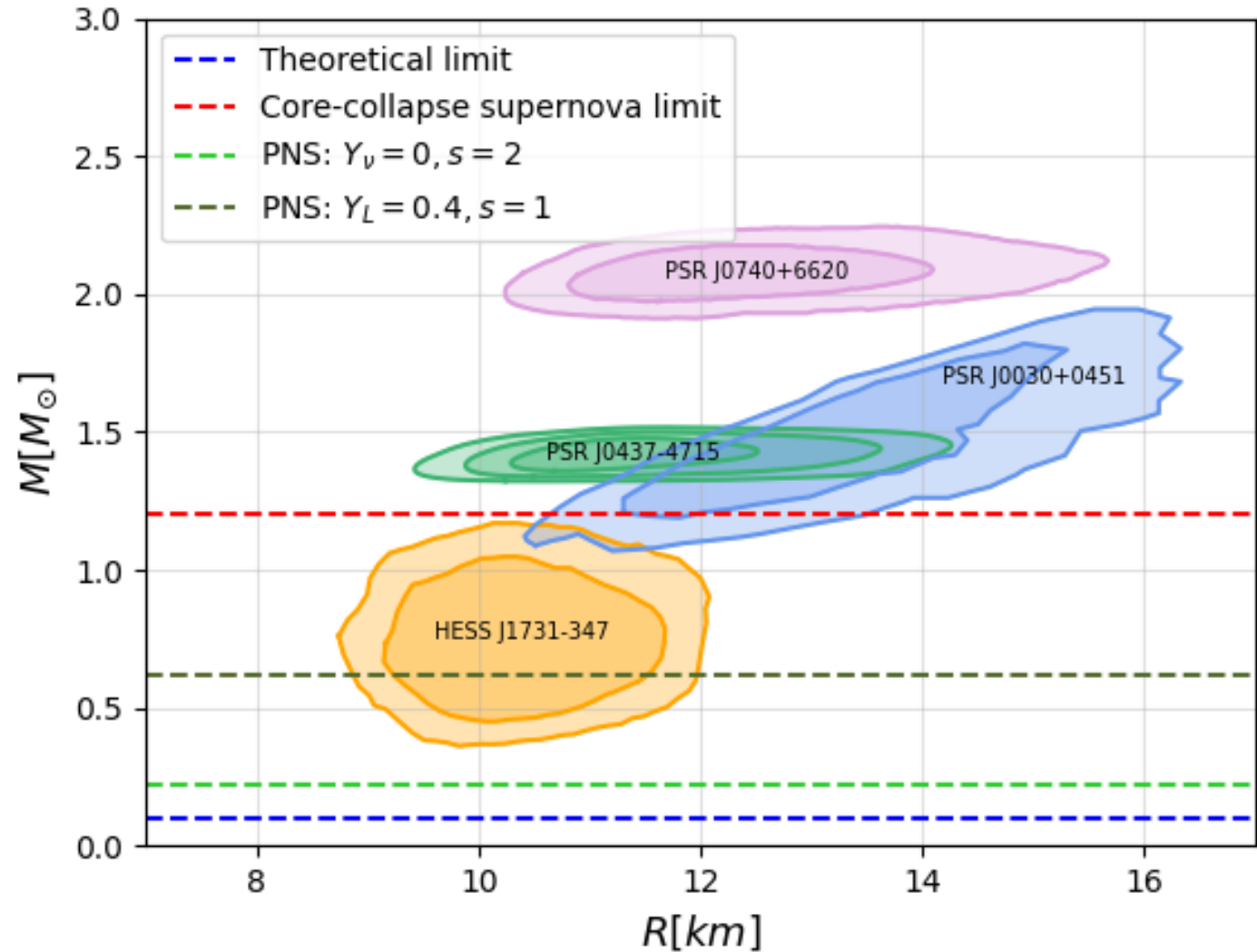
SUMMARY

- Minimal mass of proto-neutron stars with chirally constrained equations of state:

- $t = 0: Y_L = 0.4 \text{ \& } s = 1$
 $\rightarrow M_{min} = 0.62M_{\odot}$

- $t = 15s: Y_v = 0 \text{ \& } s = 2$
 $\rightarrow M_{min} = 0.22M_{\odot}$

- Chiral EFT gives very good estimate on M_{min} of PNS!
- Information from HIC could help constrain the equation of state for M_{max} of PNS!



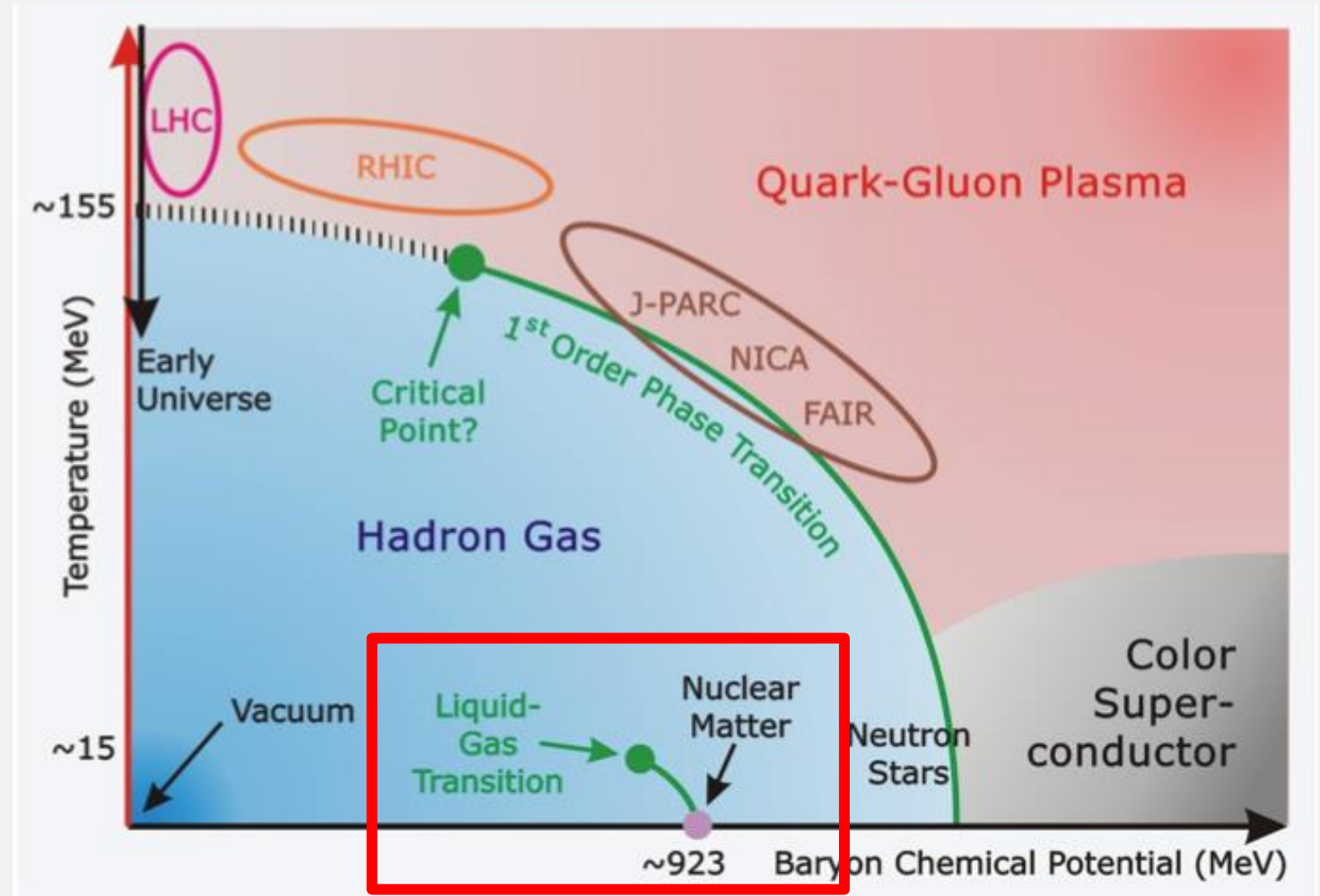
BACKUP

CHIRAL HIRARCHY

	2N Force	3N Force	4N Force	5N Force
LO $(Q/\Lambda_\chi)^0$				
NLO $(Q/\Lambda_\chi)^2$				
NNLO $(Q/\Lambda_\chi)^3$				
N³LO $(Q/\Lambda_\chi)^4$				
N⁴LO $(Q/\Lambda_\chi)^5$				
N⁵LO $(Q/\Lambda_\chi)^6$				

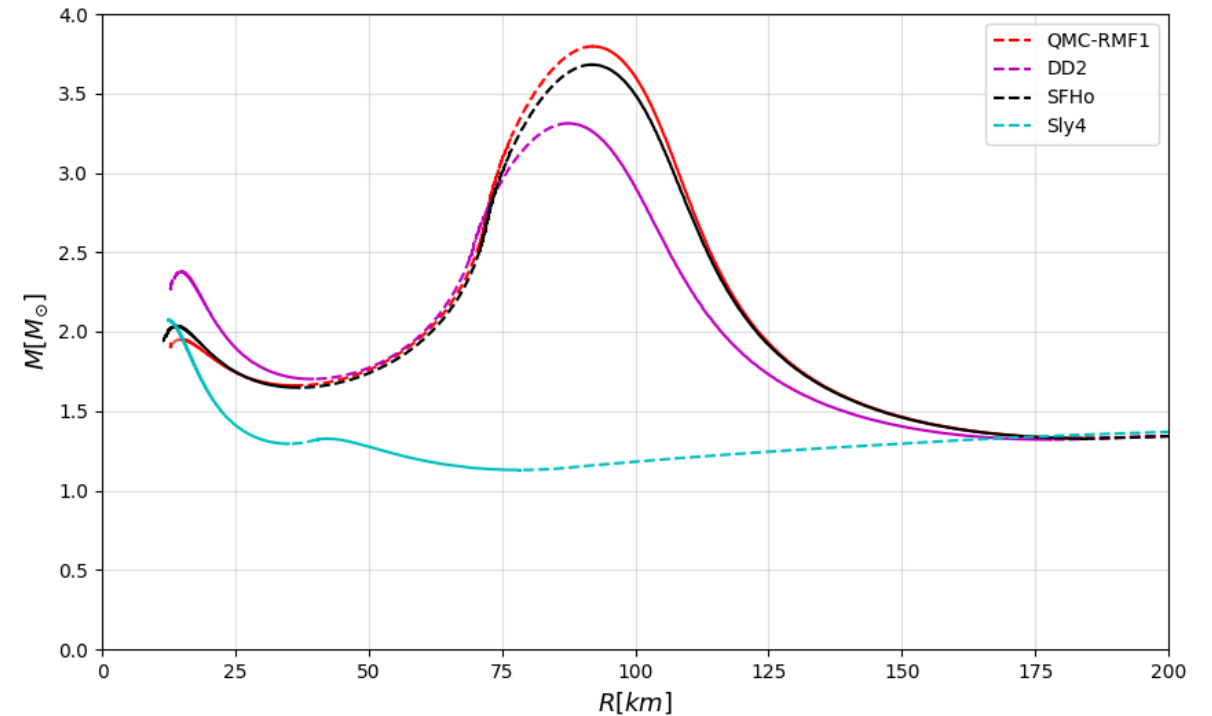
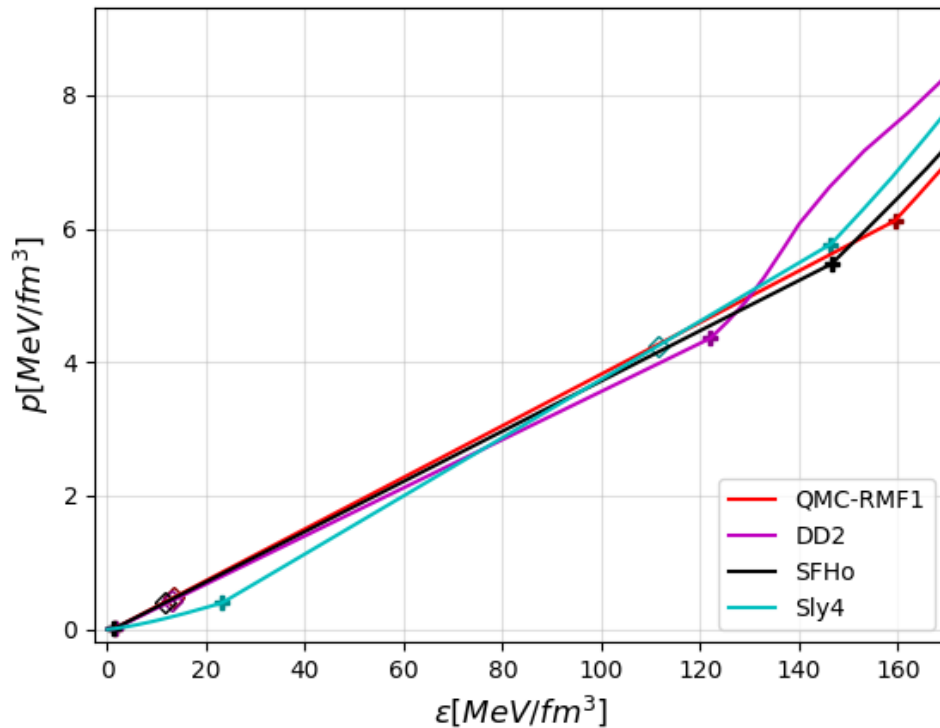
BRIEF EXCURSION: LIQUID-GAS PHASE TRANSITION

- First order phase transition between
 - A liquid of nucleons at higher densities
 - A gas of nuclei at lower densities
- Two globally conserved quantities in PNS: Y_L and $N_B \rightarrow$ Gibbs construction
- Mixture of both phases at the crust-core transition in proto-neutron stars
- What we do: Linear interpolation
- What is physically right? Gibbs construction



IMPACT OF LIQUID GAS PHASE TRANSITION FOR HIGH LEPTON FRACTION

$$Y_L = 0.5$$



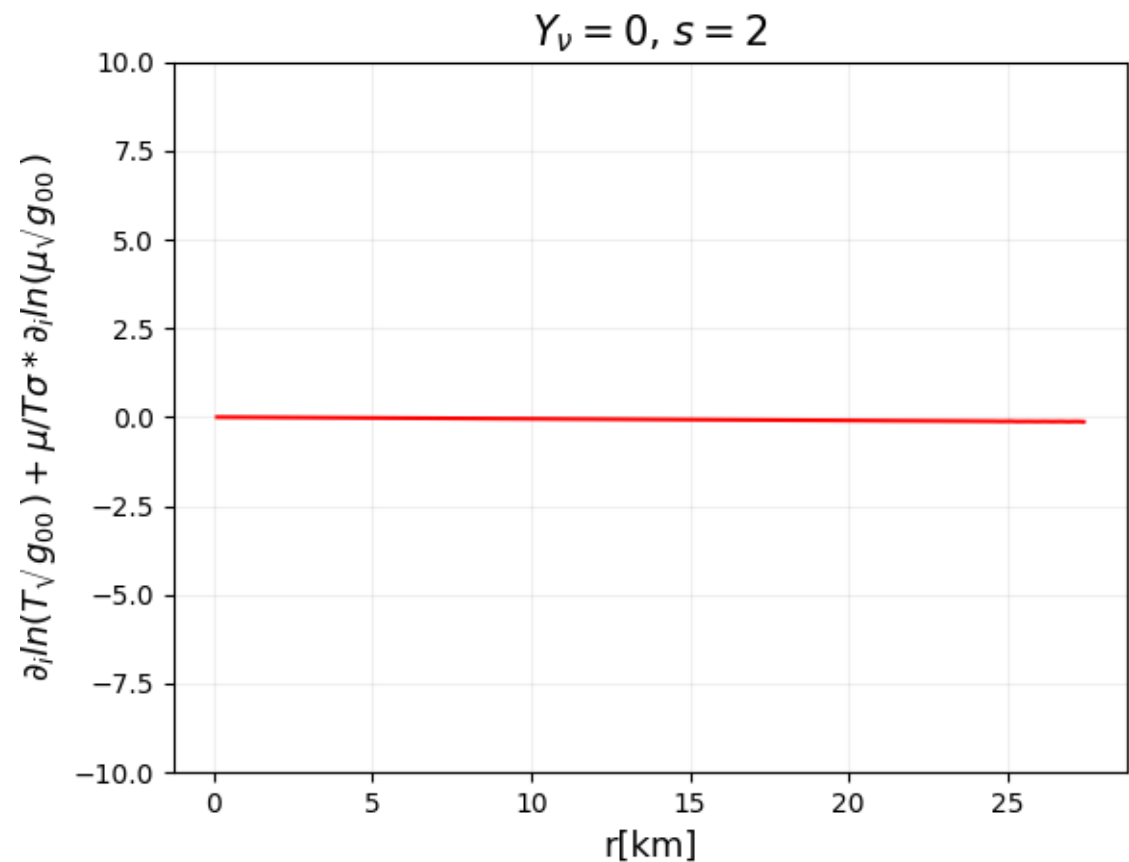
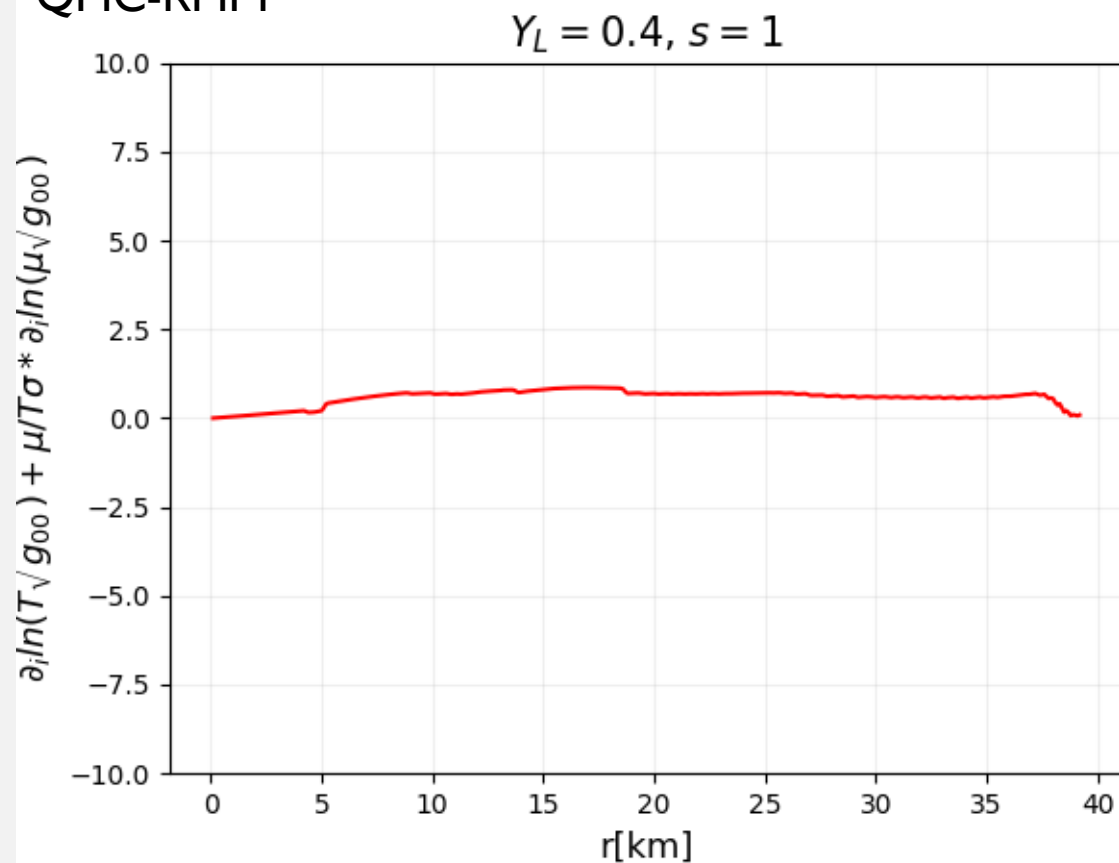
- High $Y_L \rightarrow$ Huge gap between crust and core, which needs to be constructed

- Leads to a huge second maximum in the mass-radius curve

TOLMAN-EHRENFEST-KLEIN LAW

$$\partial_i \ln(T \sqrt{g_{00}}) + \frac{\mu}{T \sigma} \partial_i \ln(\mu \sqrt{g_{00}}) = 0^*$$

QMC-RMFI



*R. C. Tolman & P. Ehrenfest, *Phys. Rev.* **36**, 1791 (1930);
O. Klein, *Rev. Mod. Phys.* **21**, 531 (1949)

OUTLOOK

- Proto-neutron stars with color-superconducting matter
- Which color-superconducting phases can we reach in proto-neutron stars?
- Which maximal masses do we get for the same conditions (same lepton fraction and entropy per baryon)?

