

Dilepton results in SMASH

Renan Góes-Hirayama

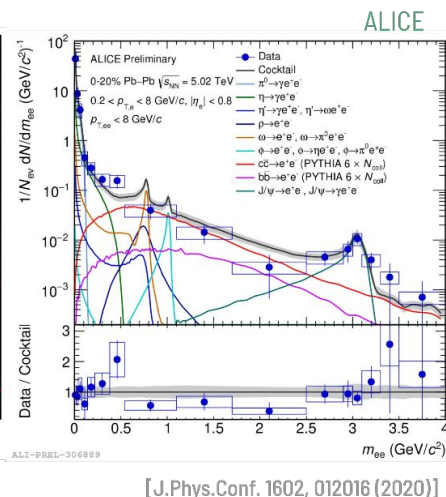
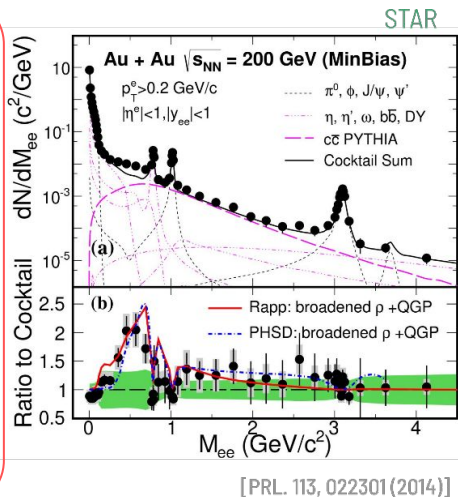
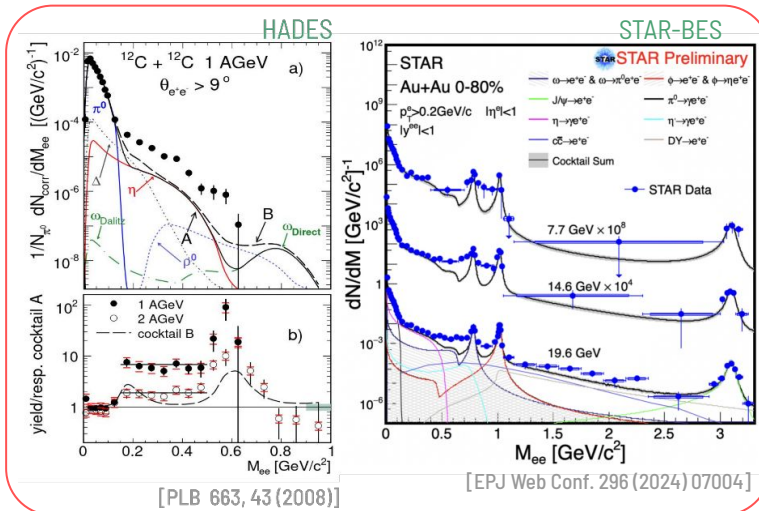
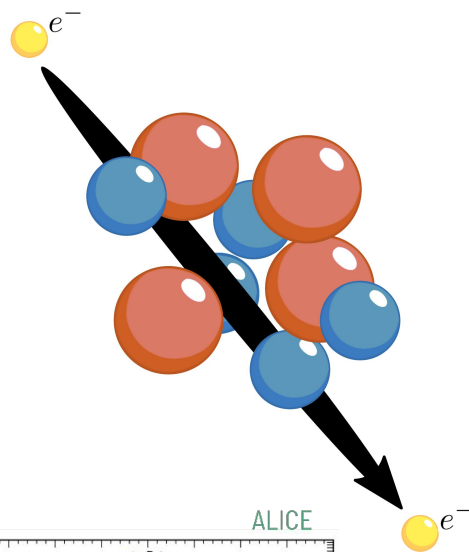
in collaboration with Hannah Elfner

November 11th, 2025

EMMI Workshop: Collective phenomena and the
Equation-of-State of dense baryonic matter

Dileptons as probes

- No strong interaction \Rightarrow leave the medium undisturbed
- Multi-messenger: spectrometer, chromimetry, barometer, thermometer, ...
- Experimental difficulties:
 - rare probes $\text{BR}(h \rightarrow l^+l^-) \sim 10^{-5}$
 - combinatorial background



The SMASH approach



<https://smash-transport.github.io/>

Hadrons

- Hadrons evolved with the relativistic Boltzmann equation (BUU)

$$p^\mu \partial_\mu f_i(x, p) + m_i F^\alpha \partial_\alpha^p f_i(x, p) = C_{\text{coll}}^i$$

- Scatterings determined geometrically from "bottom-up" cross sections

$$\pi d_{\text{trans}}^2(a, b) < \sigma_{\text{tot}}(a, b) = \sum_R \sigma_{ab \rightarrow R} + \sum_{cd} \sigma_{ab \rightarrow cd}$$

- Mass-dependent width for hadronic decays

$$\frac{\text{Prob}(R \text{ decays in } \Delta t)}{\Delta t} = \Gamma_R^{\text{vac}}(m) = \sum_{ab} \Gamma_{R \rightarrow ab}(m)$$

Vacuum properties a priori

[J. Weil et al, PRC 94 (2016) 5, 054905]

Dileptons

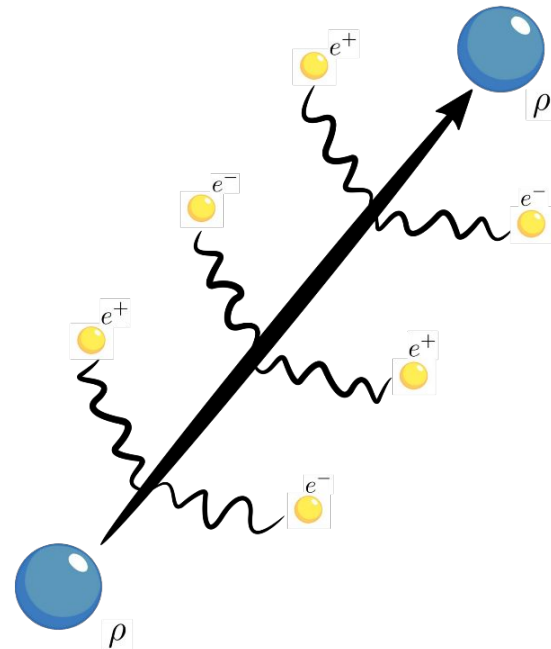
- Electromagnetic coupling is much smaller than strong coupling
- Stable hadrons decay at the end
- Time-integration: perturbative treatment for dilepton emission

$$w_{R \rightarrow l+l-}(\tau) = \int_{t_0}^{\tau-t_0} \frac{dt}{\gamma} \Gamma_{R \rightarrow l+l-}(m_R)$$

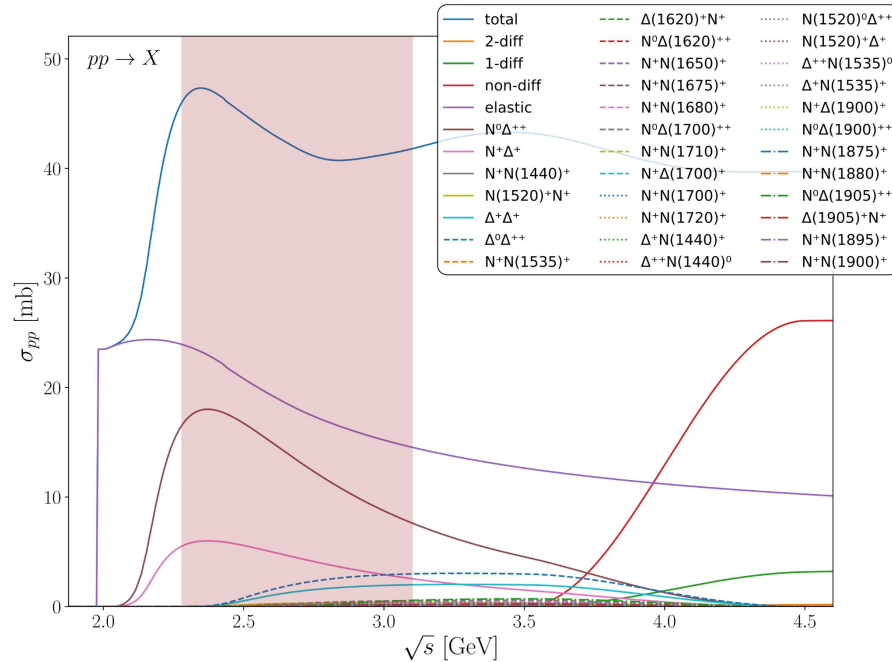
[U. Heinz and K. S. Lee, Nucl. Phys. A 544, 503 (1992)]

[G.-Q. Li and C. M. Ko, Nucl. Phys. A 582, 731 (1995)]

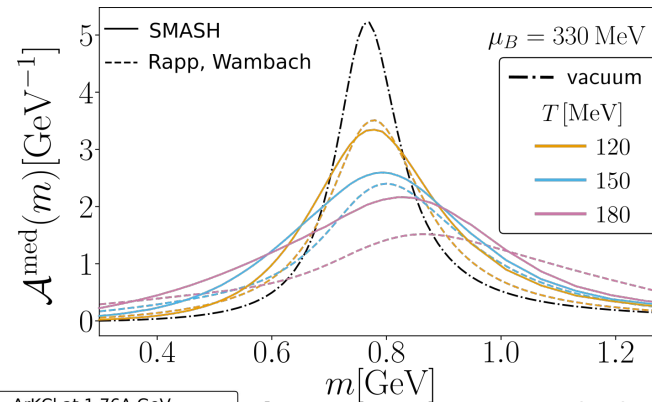
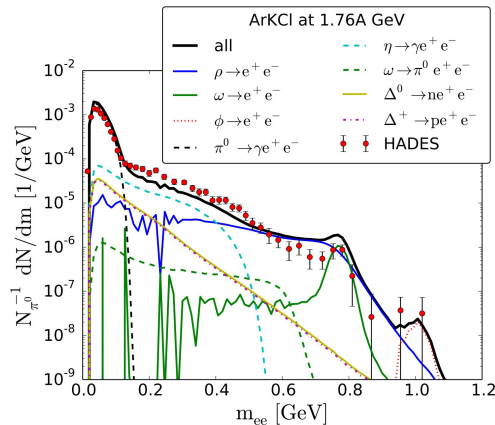
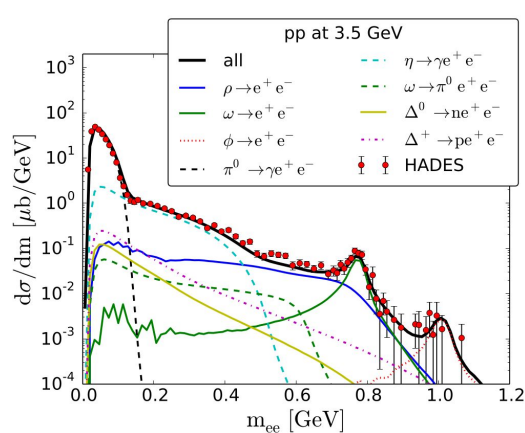
- At every time step the particle radiates a lepton pair, carrying the shining weight



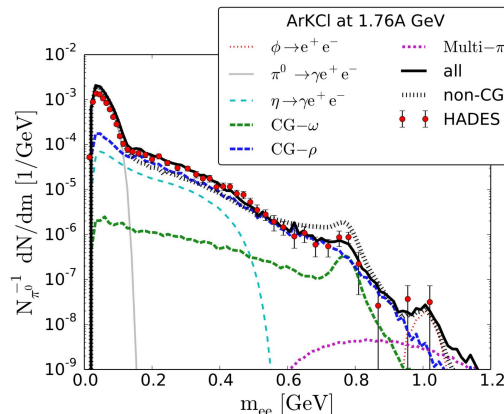
The SMASH approach – NN cross sections



The SMASH approach – dilepton production



- Data well described in pp collisions, but not in ArKCl
- Missing in-medium modifications
- Alternative: coarse graining



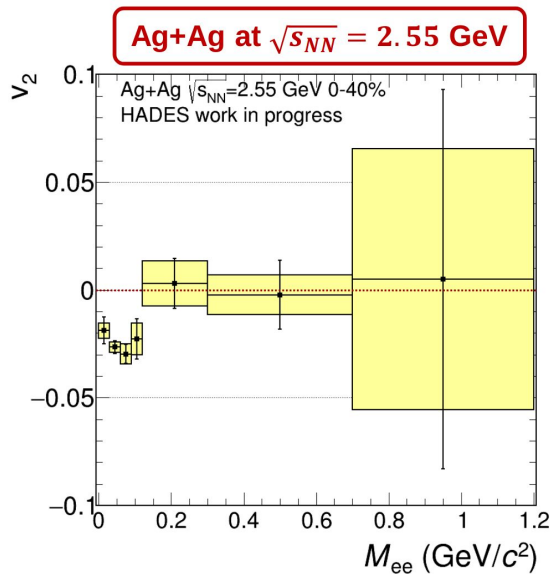
[R. Rapp, J. Wambach, Eur.Phys.J.A 6 (1999) 415-420]
 [H. van Hees, R. Rapp, Nucl.Phys.A 806 (2008) 339-387]
 [RH, J. Staudenmaier, H. Elfner, PRC 107 (2022) 2, 025208]

Medium modifications are not fully reproduced by collisional broadening from vacuum interactions

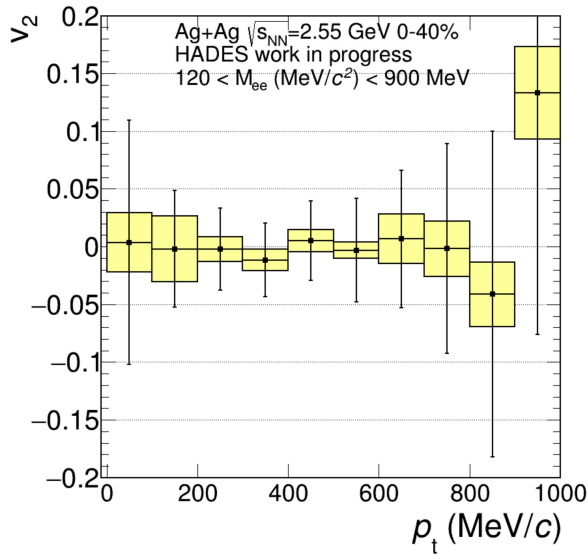
Elliptic flow and finding the Δ

Based on Phys.Rev.C 110 (2024) 6, 064903

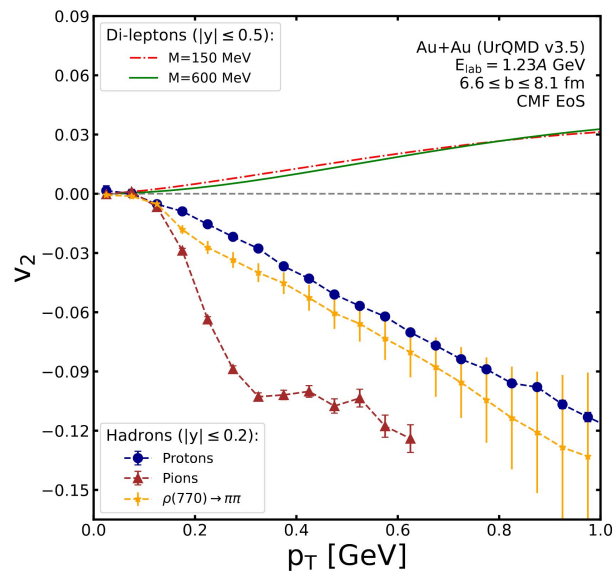
Tension?



[Niklas Schild, 15 minutes ago]



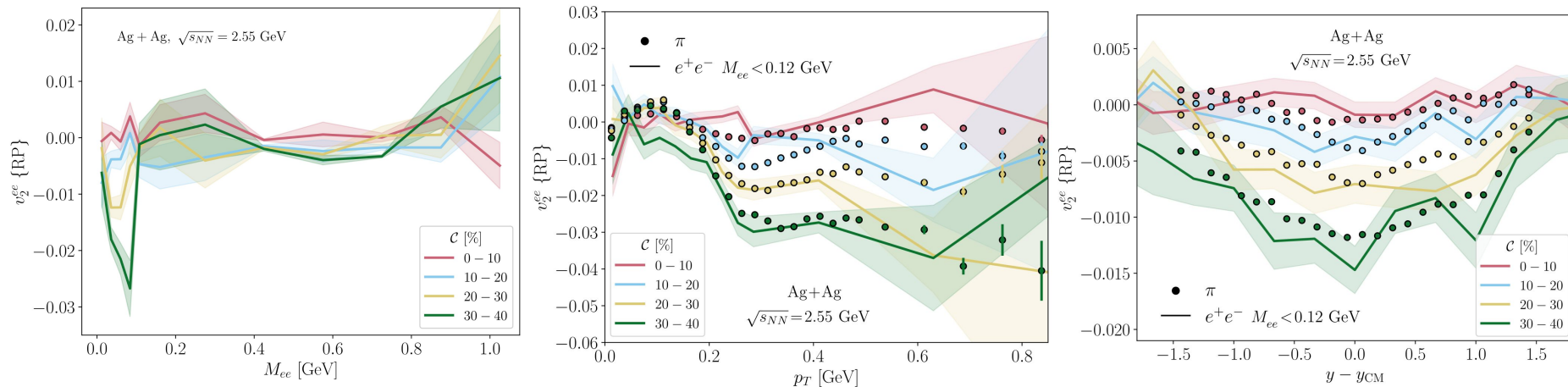
[Niklas Schild, PoS HardProbes2023 (2024) 072]



[Reichert et al. Phys.Lett.B 841 (2023) 137947]

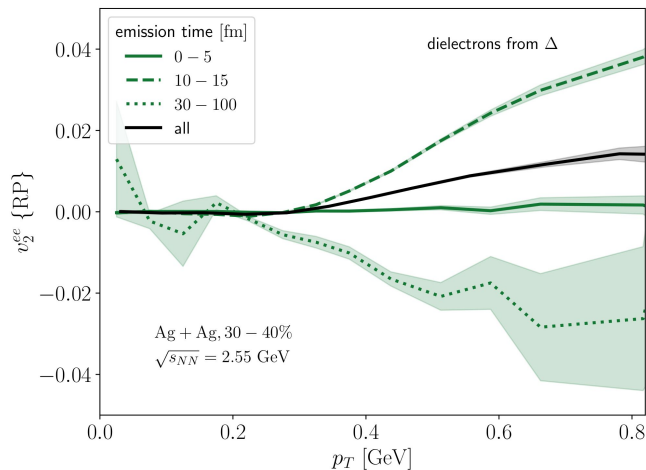
Reaction plane results

$$v_n\{\text{RP}\} = \langle\langle e^{in\phi}\rangle\rangle$$

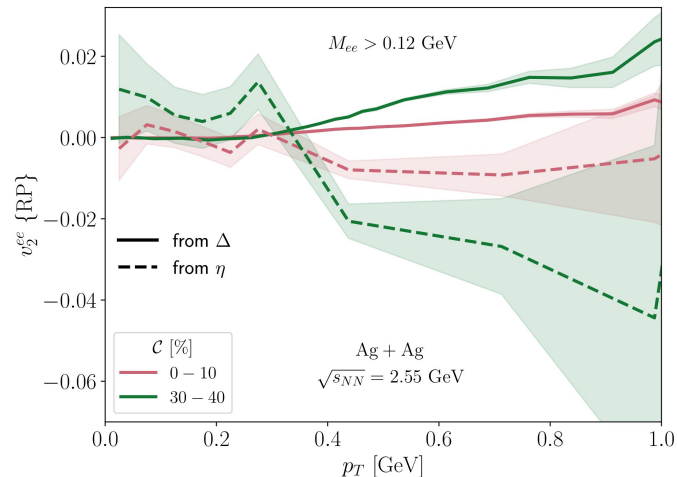


- Centrality dependence follows expected trend (less central, more flow)
- Self-consistency: low-mass SMASH dileptons follow own charged pion flow
- Overall consistency to HADES
- “Zero” flow for $M_{ee} > m_{\pi^0}$ ← Does this mean that resonances do not develop anisotropic flow?

Reaction plane results



No!



Cancellation effects:

Zero flow is only on average

- Detected dileptons emitted throughout evolution, and flow of the medium changes direction
- Typical radiation of each source happens at different times

short-lived Δ vs. “final state” resonance η

Tagging the dileptons

$$\mathbf{q}_n^X(\Omega) = \frac{1}{N_X(\Omega)} \sum_{k \in \Omega}^{N_X} e^{in\phi_k} \quad \text{Event flow vector}$$

Measure

Tagged scalar product method

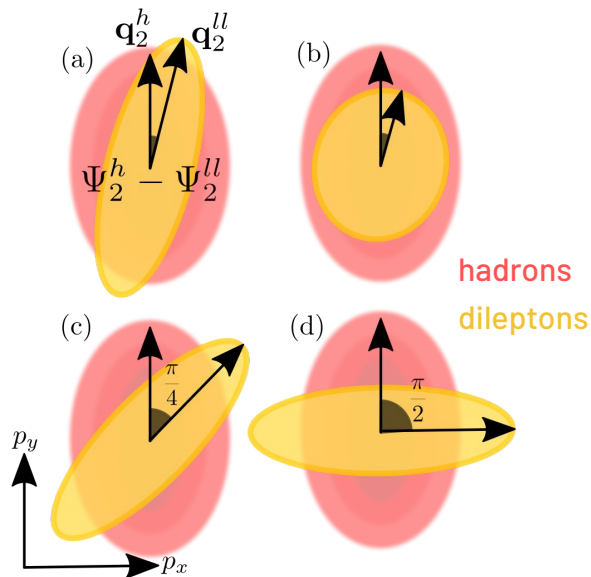
Measure differential flow of dileptons from global flow of a hadron species

$$v_{ll}\{\text{SP}|h\}(\Omega) \equiv \frac{\text{Cov}[\mathbf{q}_{ll}(\Omega), \mathbf{q}_h]}{\sqrt{\langle |\mathbf{q}_h|^2 \rangle}} \quad \text{Cov}(\mathbf{A}, \mathbf{B}) = \langle \mathbf{A} \cdot \mathbf{B}^\dagger \rangle - \langle \mathbf{A} \rangle \cdot \langle \mathbf{B} \rangle$$

Usually very small

$$v_{ll}\{\text{SP}|h\}(\Omega) = \frac{\langle q_{ll}(\Omega) q_h \cos n(\Psi_{ll} - \Psi_h) \rangle}{\sqrt{\langle |\mathbf{q}_h|^2 \rangle}} - \underbrace{\frac{v_{ll}\{\text{RP}\}(\Omega) v_h\{\text{RP}\}}{\sqrt{\langle |\mathbf{q}_h|^2 \rangle}}}_{\text{Usually very small}}$$

"Free" measurement!



$$v_2(a) > v_2(b) > 0$$

$$v_2(c) = 0 \quad v_2(d) < 0$$

Usual geometric interpretation lost

Proposed for LHC energies:

[Jean-François Paquet et al. PRC 93, 044906 (2016)]

[Gojko Vujanovic et al. PRC 101, 044904 (2020)]

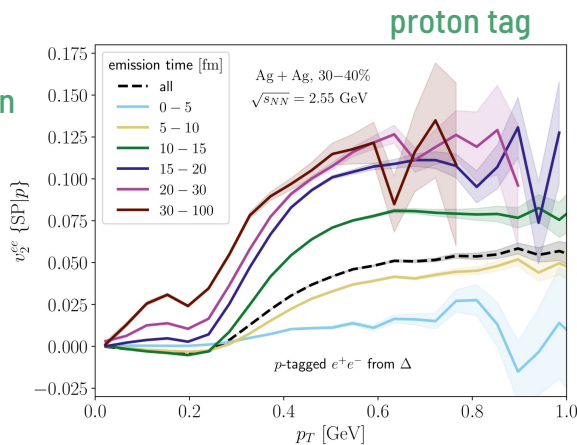
Time evolution

$$\Delta \rightarrow N e^+ e^- \text{ and } \Delta \rightarrow N \pi$$

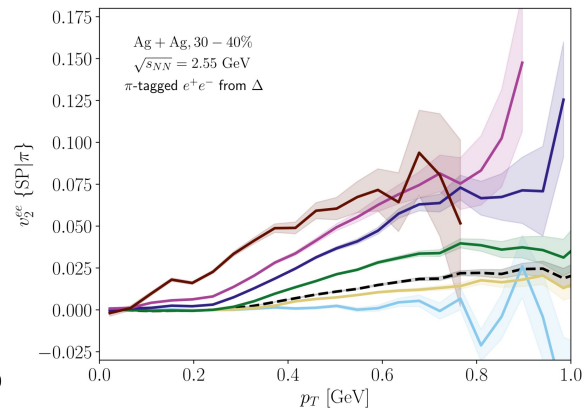
- Δ contribution starts at 0 and increases as medium expands
- Expansion: Δ collides elastically with the medium, loses momentum
- Typical emission at ~ 10 fm
- Vector mesons not linked to the protons but decay directly into pions

Tagging selects the relevant sources

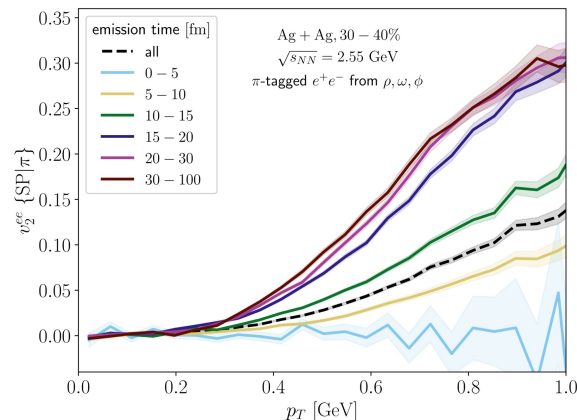
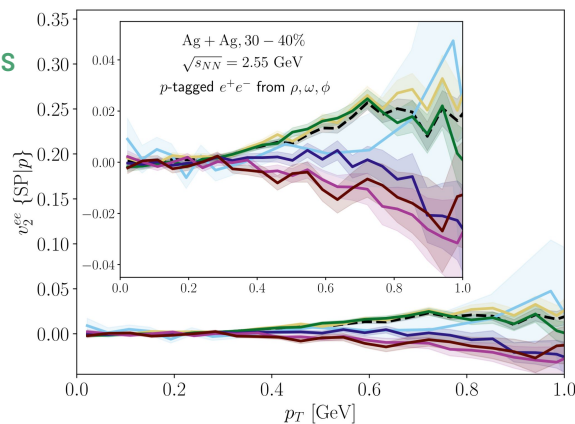
Δ baryon



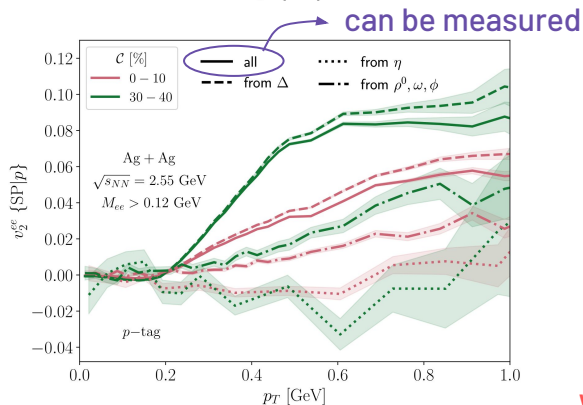
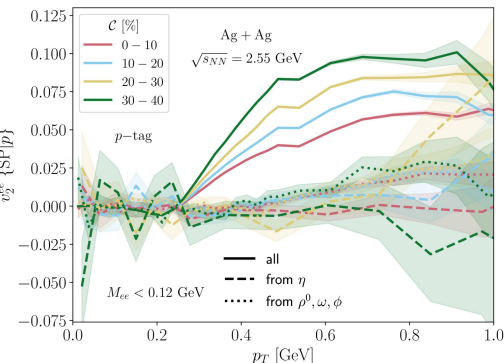
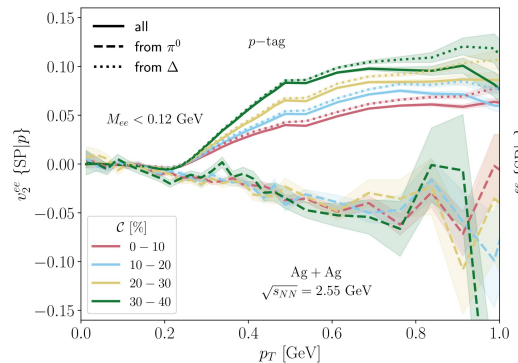
pion tag



vector mesons

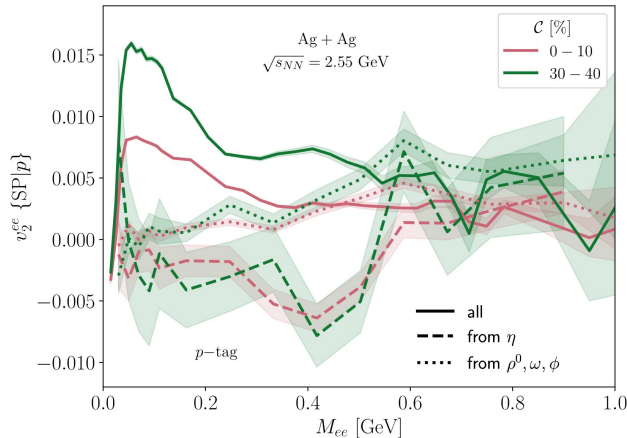
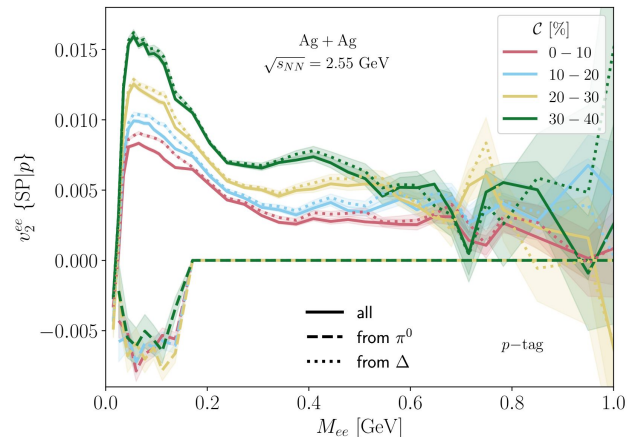


Momentum and mass dependence



Even though Δ is not the strongest dilepton production channel, it's more strongly correlated to the proton flow

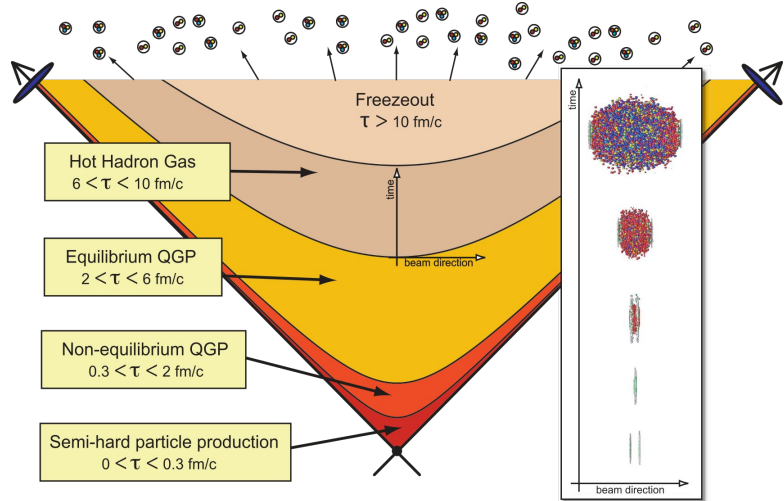
We can use the p-tag to track the Δ !



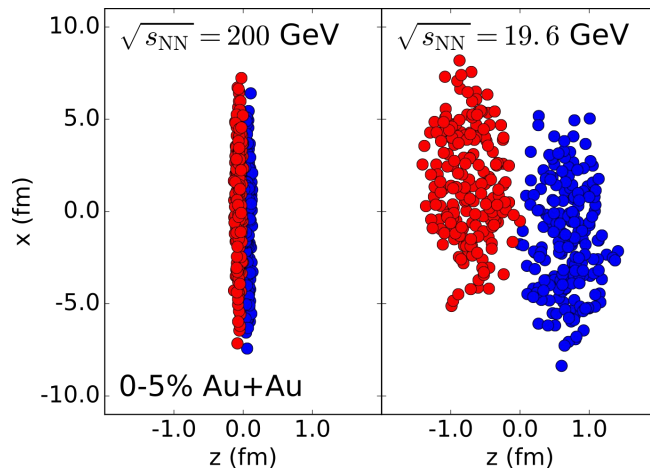
Dilepton production in a hybrid model

Based on 2507.19389 and 2510.02862

The standard HIC pictures



[M. Strickland, A Phys. Pol. B 45, 2355 (2014)]



[C. Shen & B. Schenke, Phys.Rev.C 97, 024907 (2018)]

High energy initial conditions: nuclear passing time $\tau_0 = \frac{2R/\gamma}{v_z} = 2R \left(\frac{s_{NN}}{4m_N^2} - 1 \right)^{-1/2}$
 ($20 \leq \sqrt{s_{NN}}/\text{GeV} \leq 200$)

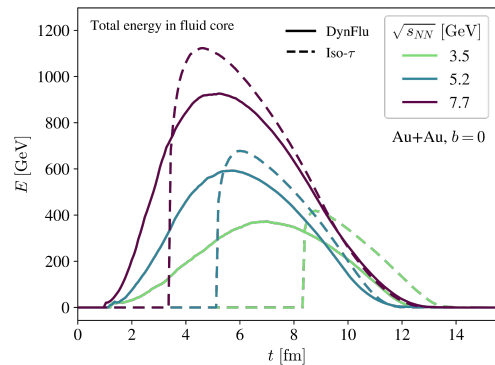
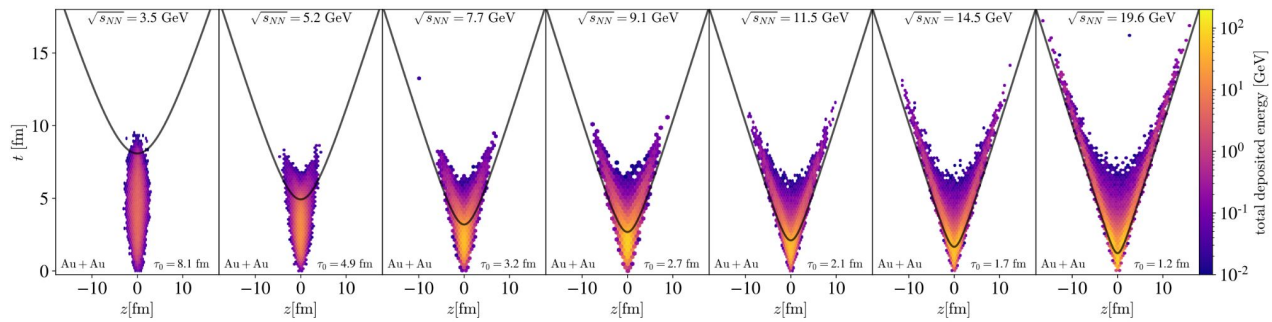
[I. Karpenko, P. Huovinen, & H. Elfner, Phys.Rev.C 91 (2015) 6, 064901]

Low/intermediate beam energies: nuclei are not pancakes!

($3 \leq \sqrt{s_{NN}}/\text{GeV} \leq 20$)

How do we start a hydro evolution in this case?

Dynamic fluidization



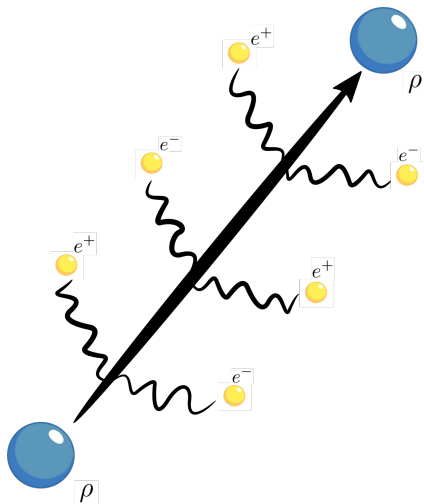
- SMASH Hadrons in regions of *enough* energy density become part of the “core” IC
- Given as sources for the fluid evolution in vHLL Hydro
- After particlization, *core+corona* are merged for rescattering Sampler Afterburner



Dilepton emission

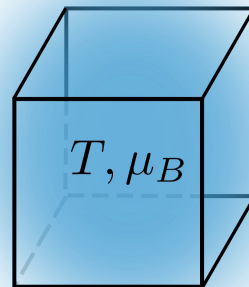


Shining method: radiate weighted dileptons every time step



vHLE

Thermal rates: generate dilepton for each source in each hydro cell



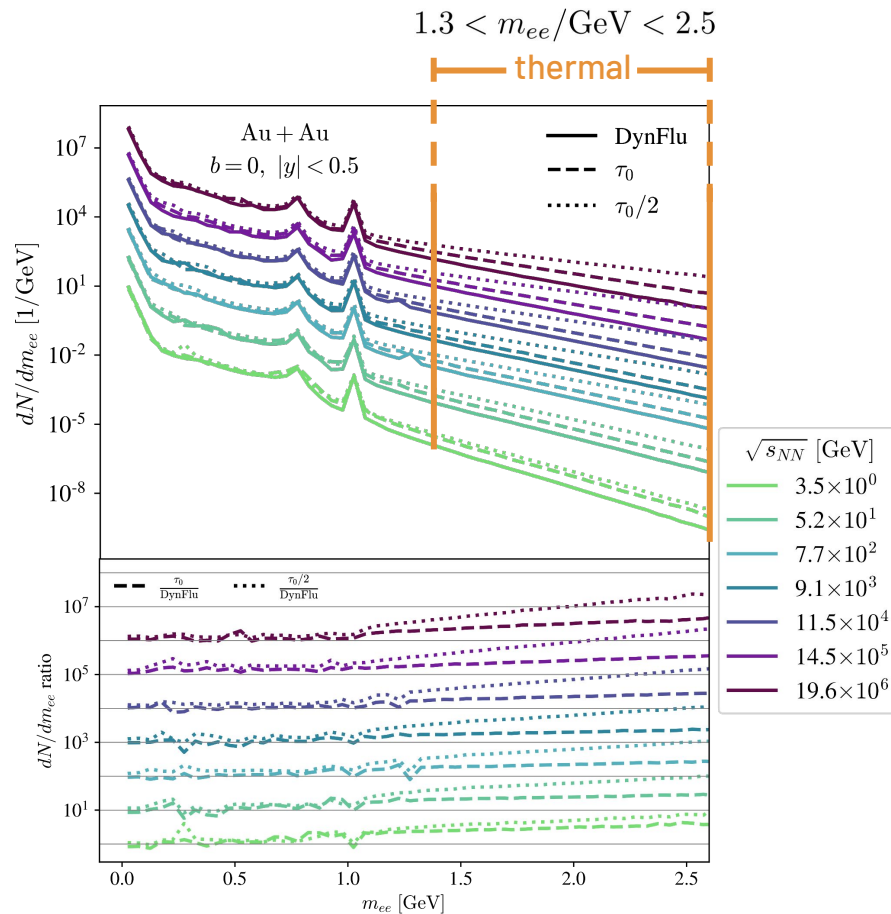
Invariant mass spectra

→ LMR: little sensitivity to IC type

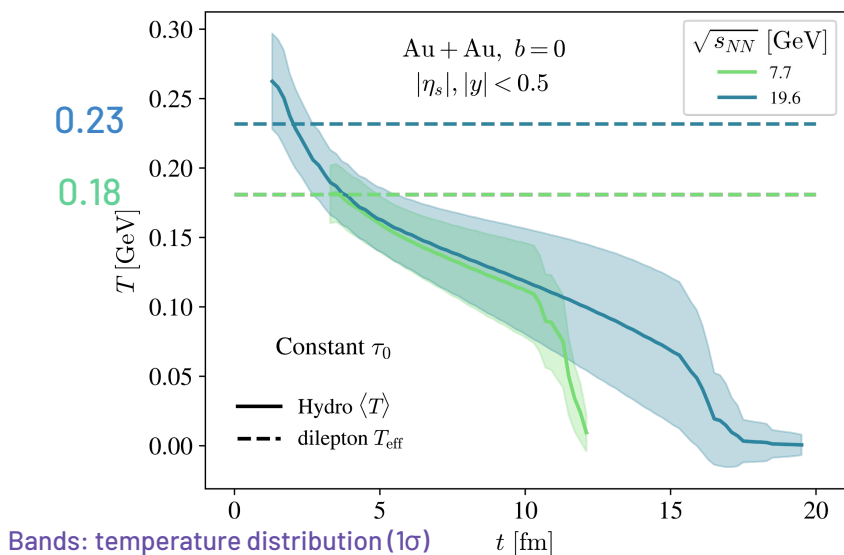
→ IMR: both yield and angle are affected

$$\frac{dN}{dm} \propto m^{3/2} e^{-m/T}$$

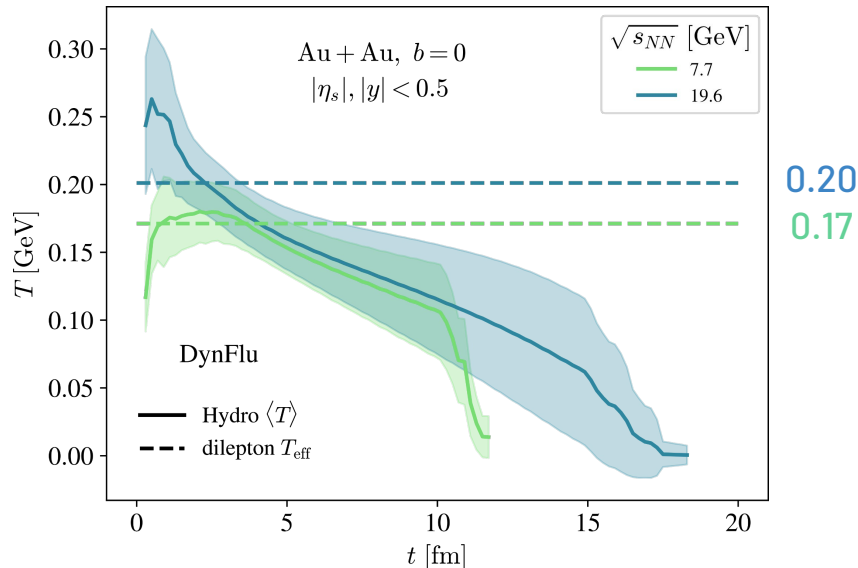
Hotter fluid → harder dileptons



Dileptons as thermometers



Bands: temperature distribution (1σ)

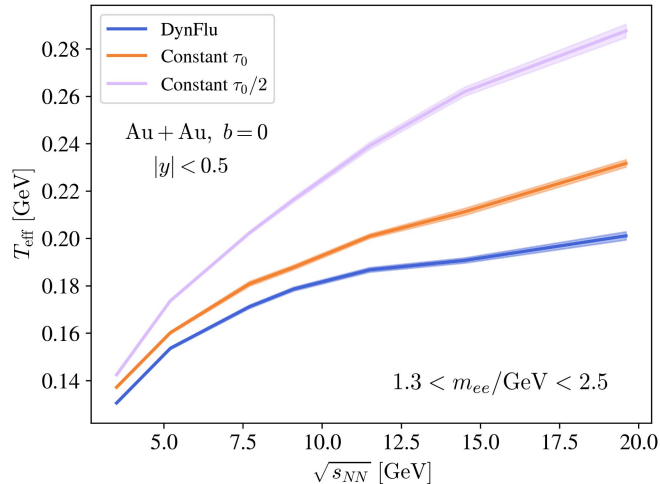
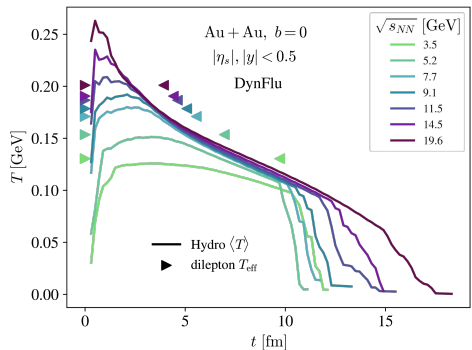
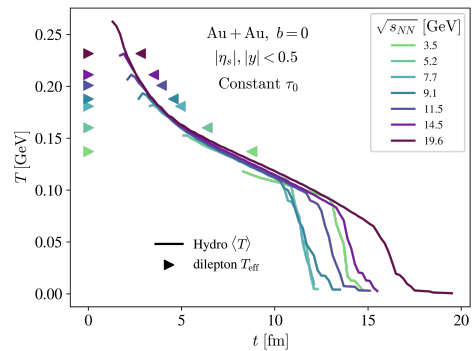


Hydro curves are similar, but smaller effective temperatures



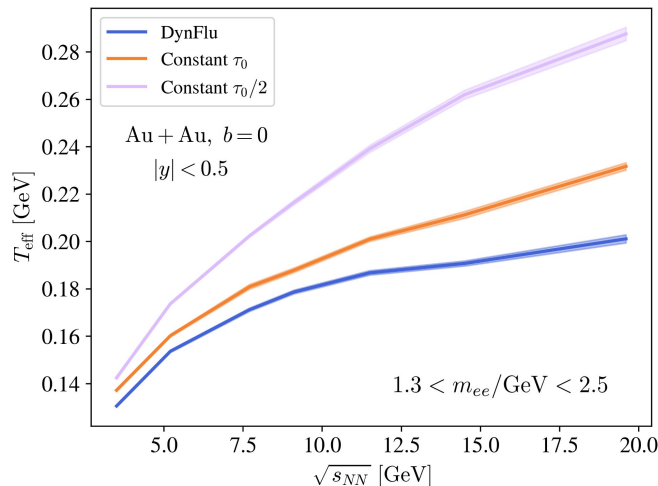
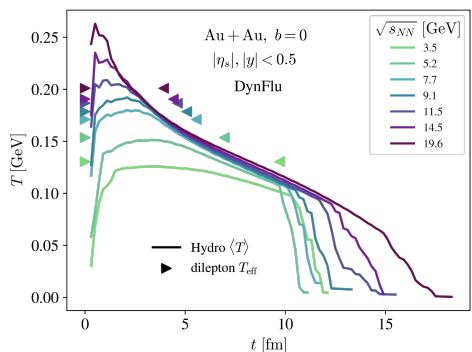
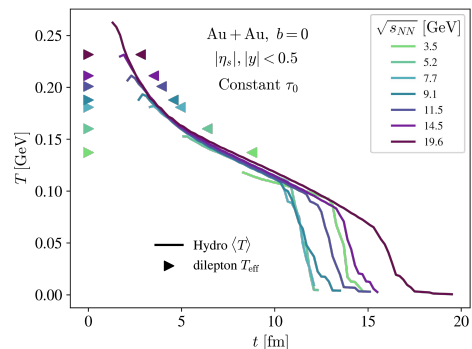
fewer hot early hydro cells

Dileptons as thermometers



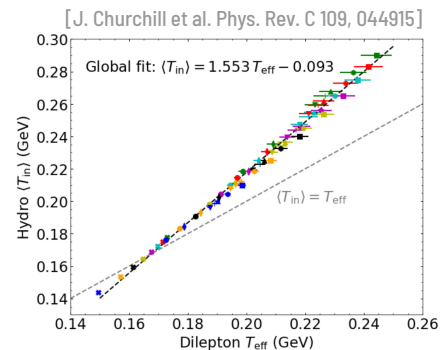
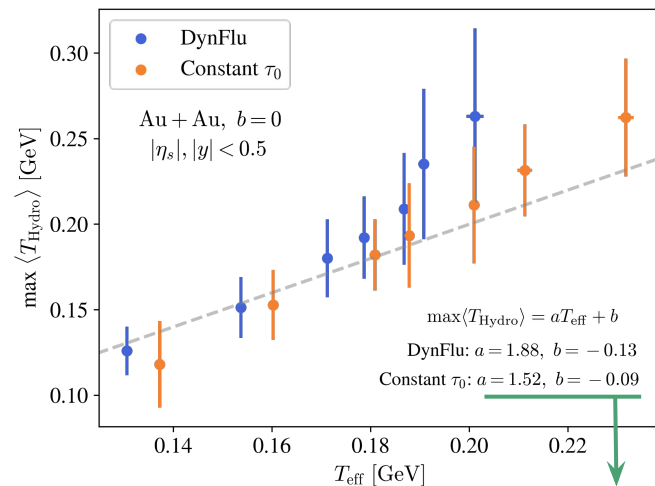
- IMR dileptons track early stages of hydro
- Low energies: extracted T_{eff} above $\langle T_{\text{Hydro}} \rangle$
- Dynamic fluid is colder

Dileptons as thermometers



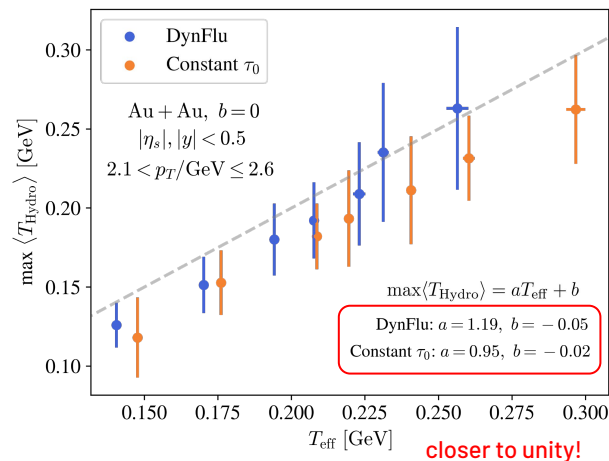
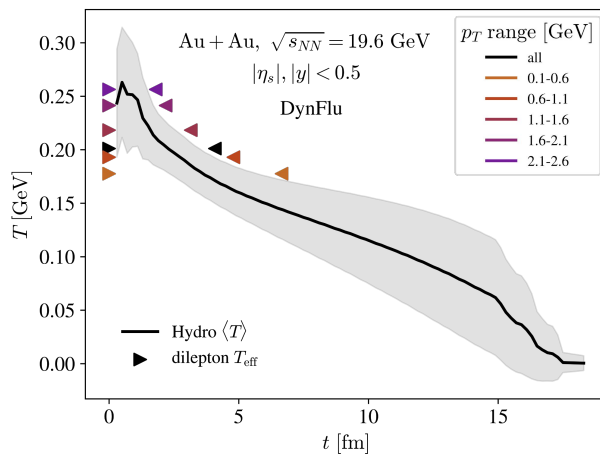
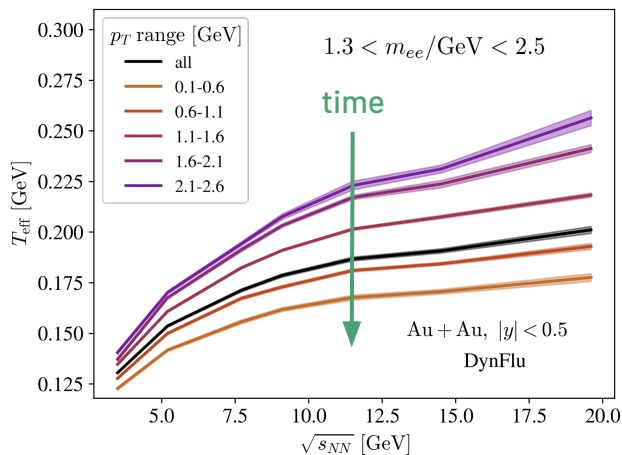
- IMR dileptons track early stages of hydro
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Different interpretations of measurement



Dileptons as thermoclocks

Early fluid → hotter cells → higher p_T dileptons



Different p_T cuts probe different times!

Summary

Dilepton flow at HADES

- Cancellation between opposing sources may lead to “zero” flow
- Delta contribution can be tracked by correlating dilepton with protons

Production in hybrid

- Dilepton effective temperature highly dependent on initialisation
- Going to high p_T

Backup

Dynamic fluidization

Philosophy: *dense systems thermalize*

1. Energy momentum tensor for particles: $T^{\mu\nu}(\mathbf{x}) = \sum_i \frac{p_i^\mu p_i^\nu}{p_i^0} K(\mathbf{x} - \mathbf{x}_i)$ Truncated gaussian smearing
2. Fluidization time of string fragmentation products: $f_t = 0.25 \text{ fm}$ (rest frame)
3. Hadron j becomes **core** if fluidization time has passed and $T^{00}(\mathbf{x}_j) > m_j K(0) + \epsilon_{\text{threshold}}$
4. Core and corona particles only interact elastically → 0.5 GeV/fm^3

