

Ab-initio studies of few-nucleon reactions of astrophysical interest



UNIVERSITÀ DI PISA

Laura E. Marcucci
University of Pisa
INFN-Pisa



Köln25 - The DPG-Frühjahrstagung
March 10-14, 2025 - University of Köln, Germany

[Microscopic] *ab-initio* studies of few-nucleon reactions of astrophysical interest

- Introduction: when an approach is
 - microscopic
 - *ab-initio*
- Theoretical ingredients:
 - Nuclear interaction
 - Electroweak nuclear currents
- Nuclear reactions of astrophysical interest:
 - reactions of interest for stellar evolution models (SSM)
 - reactions of interest for Big Bang Nucleosynthesis
- ... and beyond:
 - reactions of interest for energy production via nuclear fusion
- Conclusions and outlook

Introduction: microscopic *ab-initio* studies (I)

Nuclear observable X

Cross section of $A_1 + A_2 \rightarrow A + \gamma$

Cross section/polarization obs. for electron scatt. off $^A X$

...

- Phenomenological Cluster methods $\Rightarrow$ sometimes the only way
 - Nucleus = system of clusters depending on the considered situation
 - ${}^6\text{Li} \equiv {}^4\text{He} + {}^2\text{H}$ or ${}^6\text{Li} \equiv {}^4\text{He} + n + p$
 - ${}^7\text{Be} \equiv p + {}^6\text{Li}$ or ${}^3\text{He} + {}^4\text{He}$
 - “Model-dependent” predictions
- Microscopic methods
 - Nucleus = system of A nucleons
 - interacting among themselves $\rightarrow$ **structure**
 - interacting with external electroweak probes $\rightarrow$ **reactions**

Introduction: microscopic ab-initio studies (II)

Microscopic $\rightarrow$ *ab-initio*

Ingredients:

- ① realistic description of nuclear interactions [fixed]
- ② realistic description of electroweak currents [fixed]
- ③ exact (*ab-initio*) method to solve the quantum-mechanical problem
[only controllable approximations are allowed]

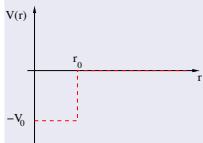
$\Rightarrow$ True predictions for observable X

Ideal case: robust procedure to estimate the theoretical error

Exact means ... ($A = 2$ bound state, $E = -B_d$)

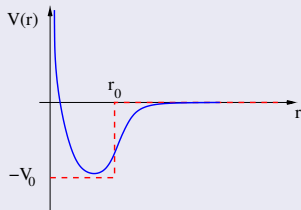
$$\Psi(r) = \frac{u(r)}{r} \quad \longrightarrow \quad u''(r) - \frac{m}{\hbar^2} [V(r) + B_d] u(r) = 0$$

Square well $\rightarrow$ analytical solution



$$\begin{aligned} u(r) &= C_{>} e^{-\alpha r} & r > r_0 & \text{ where } \alpha = \sqrt{m B_d / \hbar^2} \\ u(r) &= C_{<} \sin(kr) & r < r_0 & \text{ where } k = \sqrt{m (V_0 - B_d) / \hbar^2} \end{aligned}$$

Realistic potential $\rightarrow$ numerical solution



Numerical algorithms:

- Numerov method [accurate for $A = 2$]
- **Expansion on a given basis**
- ...

... and the nuclear interaction is much more complicated!

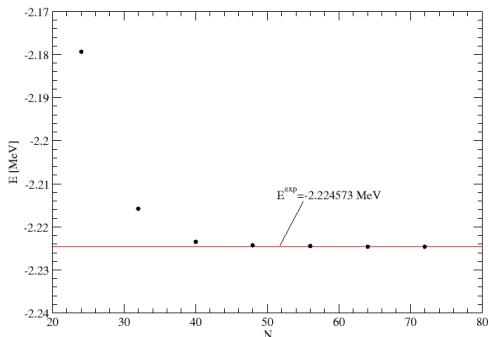
Expansion on a given basis

$$\Psi(\mathbf{r}) = \sum_{\mu} c_{\mu} L_{\mu}(\mathbf{r}) \mathcal{Y}_{\mu}(\hat{\mathbf{r}}) = \sum_{\mu} c_{\mu} \phi_{\mu}$$

- ϕ_{μ} : known functions vs. c_{μ} : unknown coefficients
- Rayleigh-Ritz var. principle $\Rightarrow$ eigenvalue-eigenvector problem $\Rightarrow c_{\mu}$ & E

$$\delta_c \langle \Psi(\mathbf{r}) | H - E | \Psi(\mathbf{r}) \rangle = 0$$

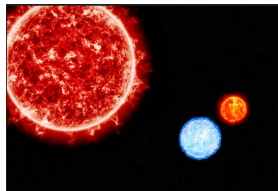
- Phenomenological potential AV18



The nuclear Hamiltonian

$$H = \sum_i \frac{p_i^2}{2m} + \sum_{i < j} \mathbf{V}_{ij} + \sum_{i < j < k} \mathbf{V}_{ijk} + \dots$$

- Non-relativistic approach $\longrightarrow \frac{p_i^2}{2m}$
- Nucleons interact in pairs $\longrightarrow \mathbf{V}_{ij}$
- Nucleons are not point-like: the presence of nucleon k changes the shape of nucleons i, j
 $\longrightarrow \mathbf{V}_{ijk}$ (same as the Earth-Moon-Sun system)
- Empirical evidence of $\mathbf{V}_{ijk}$: $B(A > 2)_{th} < B(A > 2)_{exp}$ with only $\mathbf{V}_{ij}$
- Empirical evidence of no $\mathbf{V}_{ijkl}, \dots$



History: phenomenological approach

Theoretical framework until $\simeq 20$ years ago

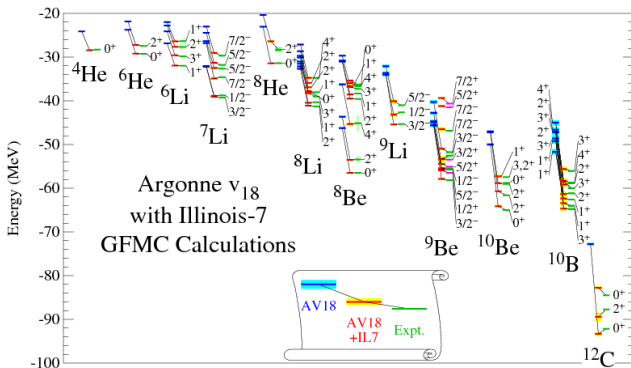
- $V_{ij} \rightarrow \text{AV18}$

$$V_{ij} = \sum_{p=1}^{18} f_p(r_{ij}) O_p(\hat{\mathbf{r}}_{ij})$$

- $O_p(\hat{\mathbf{r}}_{ij})$: **18 operators dictated by theory+data**. They are:

$$\begin{aligned} & [\mathbb{I}, (\boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j), S_{ij}, \mathbf{L} \cdot \mathbf{S}, L^2, L^2(\boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j), (\mathbf{L} \cdot \mathbf{S})^2] \otimes [\mathbb{I}, (\boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j)] \\ & + T_{ij}, T_{ij}(\boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j), T_{ij}S_{ij}, (\tau_{i,z} + \tau_{j,z}) \\ & S_{ij} = 3(\boldsymbol{\sigma}_i \cdot \hat{\mathbf{r}}_{ij})(\boldsymbol{\sigma}_j \cdot \hat{\mathbf{r}}_{ij}) - \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j \quad \& \quad T_{ij} = 3\tau_{i,z}\tau_{j,z} - \boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j \end{aligned}$$

- $f_p(r_{ij})$: **functions with $\simeq 40$ parameters**, fitted to $A = 2$ experimental data $\rightarrow \chi^2/\text{datum} \simeq 1$
- $V_{ijk} \rightarrow \text{UIX or IL7}$: typically **much simpler model**, with 2-3 parameters fitted to $B(A = 3, 4)$



J. Carlson *et al.*, Rev. Mod. Phys. **87**, 1067 (2015)

Very successful, but ...

- no connection with physics of hadrons (QCD)
- difficult to estimate the theoretical uncertainty
- mild connection between potentials and currents

→ **Chiral Effective Field Theory (ChEFT)**

Chiral Effective Field Theory (ChEFT)

Nuclear physics

- low-energy QCD
- nuclear force = residual color force

QCD

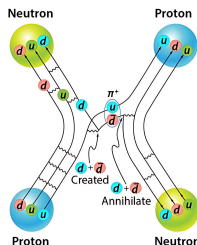
- DOF: quarks, gluons
- non-Abelian theory



non-perturbative at low energies

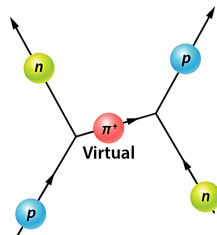
ChEFT

- DOF: nucleons, pions, (Δ -isobar)



QCD

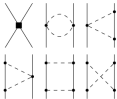
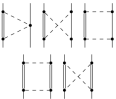
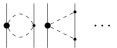
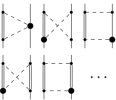
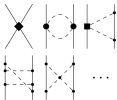
$\gtrsim 1 \text{ GeV}$



ChEFT

$\sim 100 \text{ MeV}$

Two-nucleon interaction

	Δ -less NN Forces	Additional Δ -full graphs
LO $(Q/\Lambda_\chi)^0$		
NLO $(Q/\Lambda_\chi)^2$		
N ² LO $(Q/\Lambda_\chi)^3$		
N ³ LO $(Q/\Lambda_\chi)^4$		

- Lagrangian consistent with low-energy **QCD symmetries**

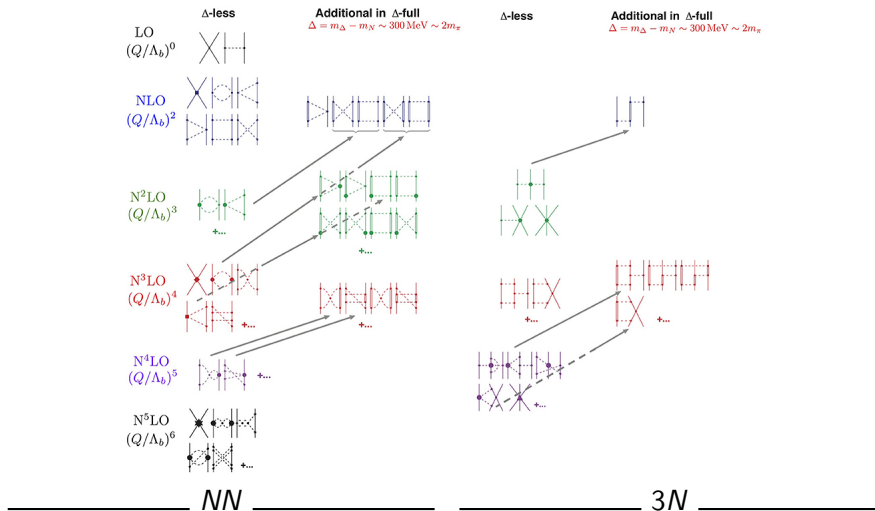
$$\mathcal{L}_{\text{ChEFT}} = \mathcal{L}_{\pi\pi} + \mathcal{L}_{\pi N} + \mathcal{L}_{NN} + (\mathcal{L}_{\pi N\Delta})$$

- **low-momentum expansion** in terms of Q/Λ_χ

$$Q \sim m_\pi \quad \Lambda_\chi \sim m_\rho$$

- high-energy DOF $\Rightarrow$ **contact terms**
- LECs $\Rightarrow$ fit to experimental data

Hierarchy of nuclear force



Regularization scheme and cutoff function

$$\mathcal{L}_{\text{ChEFT}} \implies V$$

Divergent terms appear for

- high momentum
- short distances

Regularization of $V \longrightarrow$ **cutoff functions**

- $f(p', p; \Lambda)$ or $f(r', r; R) \implies$ **non-local** $V \rightarrow p$ -space
- $f(p' - p; \Lambda)$ or $f(r; R) \implies$ **local** $V \rightarrow r$ -space [better with Coulomb]
- $f(p' - p; \Lambda)$ for π -exchange and $f(p', p; \Lambda)$ for contact terms $\implies$ **semi-local** $V \rightarrow p$ -space

Regularization scheme and cutoff function

$$\mathcal{L}_{\text{ChEFT}} \implies V$$

Divergent terms appear for

- high momentum
- short distances

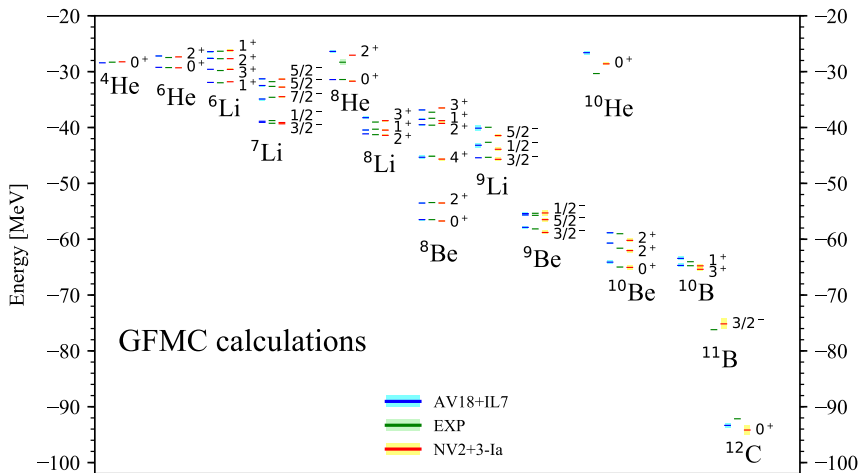
Regularization of $V \longrightarrow$ **cutoff functions**

- $f(p', p; \Lambda)$ or $f(r', r; R) \implies$ **non-local** $V \rightarrow p$ -space
- $f(p' - p; \Lambda)$ or $f(r; R) \implies$ **local** $V \rightarrow r$ -space [better with Coulomb]
- $f(p' - p; \Lambda)$ for π -exchange and $f(p', p; \Lambda)$ for contact terms $\implies$
semi-local $V \rightarrow p$ -space

Name	Space	Order(s)	Locality	Δ	max E_{LAB}	Λ or (R_L, R_S)
EMN450	p	LO – N ⁴ LO	non-local	-less	300 MeV	450 MeV
EMN500	p	LO – N ⁴ LO	non-local	-less	300 MeV	500 MeV
EMN550	p	LO – N ⁴ LO	non-local	-less	300 MeV	550 MeV
NVIa	r	N ³ LO	local	-full	125 MeV	(1.2, 0.8) fm
NVIb	r	N ³ LO	local	-full	125 MeV	(1.0, 0.7) fm
NVIIa	r	N ³ LO	local	-full	200 MeV	(1.2, 0.8) fm
NVIIb	r	N ³ LO	local	-full	200 MeV	(1.0, 0.7) fm

[EMN] D. R. Entem, R. Machleidt, and Y. Nosyk, Phys. Rev. C **96**, 024004 (2017)

[NV] M. Piarulli et al., Phys. Rev. C **91**, 024003 (2015)

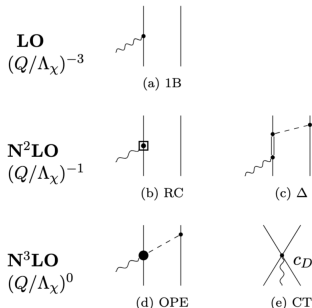


Same degree of success/accuracy as the phenomenological approach!

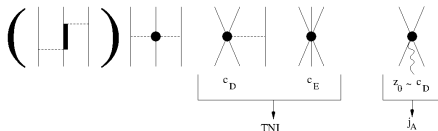
Electroweak current in ChEFT

Adding **electroweak field** as DOF

Axial current

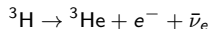


- Similar expansion for the electromagnetic current
- **Interplay interaction-current in ChEFT**



• **LECs fitted to experimental data**

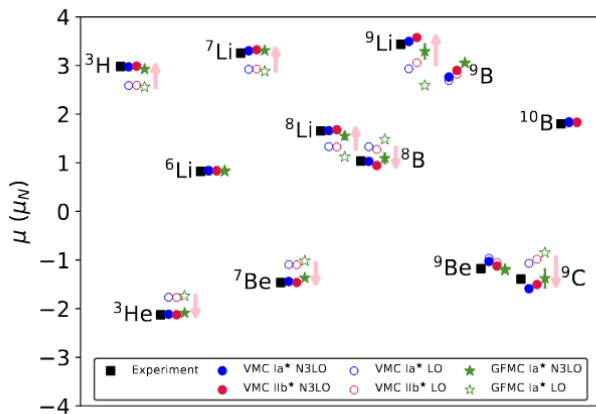
- **5 LECs** in $j_\gamma \rightarrow \mu(A=2,3)$, deuteron $G_M(q^2)$ and $d(e, e')pn$ at threshold
- z_0 -**LEC** in $j_A \rightarrow B(^3\text{H})$ & ^3H β -decay half-life



A. Gnech and R. Schiavilla, Phys. Rev. C **106**, 044001 (2022)

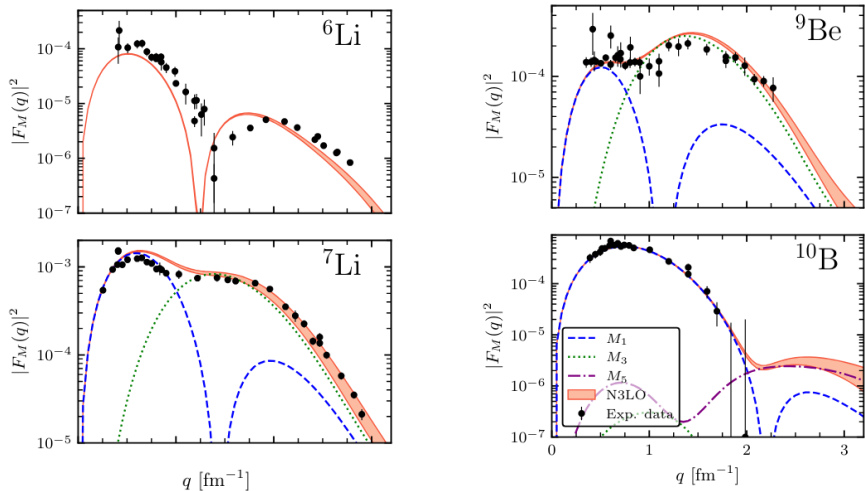
A. Gnech *et al.*, Phys. Rev. C **109**, 035502 (2024)

Magnetic structure of light nuclei: magnetic moments



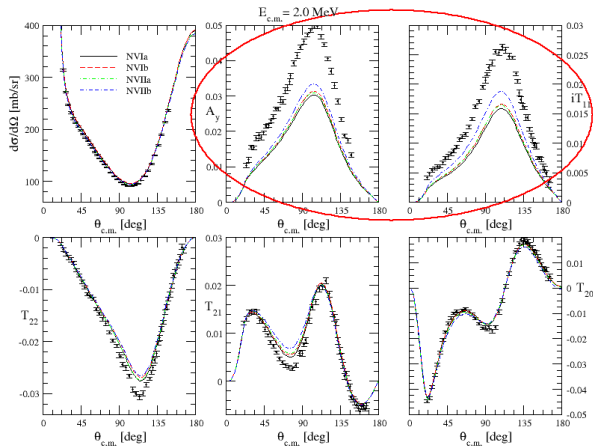
G. Chambers-Wall *et al.*, Phys. Rev. Lett. **133**, 212501 (2024)

Magnetic structure of light nuclei: magnetic form factors



G. Chambers-Wall *et al.*, Phys. Rev. Lett. **133**, 212501 (2024)

Not everything works ... $p + d \rightarrow p + d$ elastic scattering



→ **A_y-puzzle** or **iT₁₁-puzzle**

$$\vec{p} + d \rightarrow \vec{p} + d$$

$$A_y \propto \frac{\sigma_p(\uparrow) - \sigma_p(\downarrow)}{\sigma_p(\uparrow) + \sigma_p(\downarrow)}$$

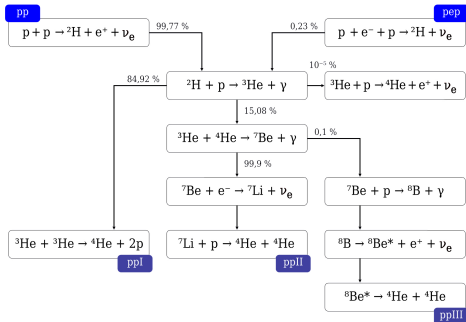
$$\vec{d} + p \rightarrow \vec{d} + p$$

$$iT_{11} \propto \frac{\sigma_d(\uparrow) - \sigma_d(\downarrow)}{\sigma_d(\uparrow) + \sigma_d(\downarrow)}$$

SELECTED RESULTS

for **NUCLEAR REACTIONS** of **ASTROPHYSICAL INTEREST**

The $p + p \rightarrow d + e^+ + \nu_e$ reaction (MS thesis of V. Barlucchi)



- first reaction of the pp-chain $\rightarrow$ solar- ν fluxes
- **suppressed** by
 - weak interaction
 - Coulomb barrier

Low energy ($E \sim 6$ keV) $\Rightarrow$ **quantum tunnelling** $\rightarrow \sigma(E) \propto P_G(E) = e^{-\pi\alpha\sqrt{m_p/E}}$

S-factor

- isolate nuclear physics
- varies slowly at low energies

$$\begin{aligned}
 S(E) &= E \sigma(E) P_G^{-1}(E) \\
 &\simeq S(0) + S'(0)E + \frac{1}{2}S''(0)E^2 + \frac{1}{6}S'''(0)E^3
 \end{aligned}$$

Expansion parameters - fixed current order

$S(E)$ calculated in the $E = 3 - 30$ keV range $\longrightarrow$ Taylor expansion to get $S(0)$, $S'(0)$, etc.

j^A fixed at N³LO]

Potential	$S(0)$ [10^{-23} MeV fm ²]	$S'(0)/S(0)$ [MeV ⁻¹]	$S''(0)/S(0)$ [MeV ⁻²]	$S'''(0)/S(0)$ [MeV ⁻³]
EMN450 N ³ LO	4.090	10.59	347.2	-6917
EMN500 N ³ LO	4.091	10.59	346.9	-6911
EMN550 N ³ LO	4.089	10.59	347.0	-6914
NVIa	4.062	10.61	348.9	-7099
NVIb	4.061	10.61	348.7	-7100
NVIIa	4.036	10.61	348.9	-7096
NVIIb	4.030	10.61	348.7	-7095

To be noticed:

- **NV models I** vs. **NV models II**
- **EMN ChEFT non-local** pot. vs. **NV ChEFT local** pot.
- differences in the d wave functions
- $S'(0)/S(0)$, $S''(0)/S(0)$ etc $\rightarrow$ model-independent

Theoretical error

Theoretical uncertainties $\rightarrow$ **truncation chiral expansion** of

- interactions
- currents

$S(E)$ calculated at $N^n\text{LO}$

$$S_n(E) = S_{\text{ref}}(E) \sum_{k=0}^n c_k(E) \left(\frac{Q(E)}{\Lambda_b} \right)^k$$

- $Q(E) = \frac{(m_p E)^4 + m_\pi^8}{(m_p E)^{7/2} + m_\pi^7}$
- $\Lambda_b = 600 \text{ MeV}$

Truncation error $\sigma_n(E) \simeq$ contribution of order $n+1$

$$\sigma_n(E) = \max \left[\left\{ |S_k(E) - S_{k-1}(E)| \left(\frac{Q(E)}{\Lambda_b} \right)^{n-k+1} \right\}_{k=0, \dots, n} \right]$$

J. A. Melendez *et al.*, Phys. Rev. C **100**, 044001 (2019)

Truncation error - interaction

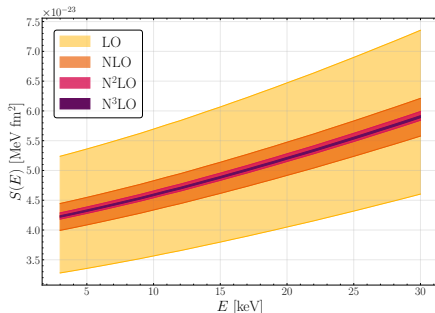
Current fixed at

- $N^2\text{LO}$ for EMN LO, NLO
- $N^3\text{LO}$ for EMN $N^2\text{LO}$, $N^3\text{LO}$ and NV

EMN

Potential	$\sigma_0^{\text{Int}}(0)$	$\sigma_1^{\text{Int}}(0)$	$\sigma_2^{\text{Int}}(0)$	$\sigma_3^{\text{Int}}(0)$
EMN450	0.940	0.216	0.050	0.011
EMN500	0.944	0.218	0.051	0.011
EMN550	0.955	0.220	0.051	0.012

$\sigma_n^{\text{Int}}(0)$ in unit of $10^{-23} \text{ MeV fm}^2$



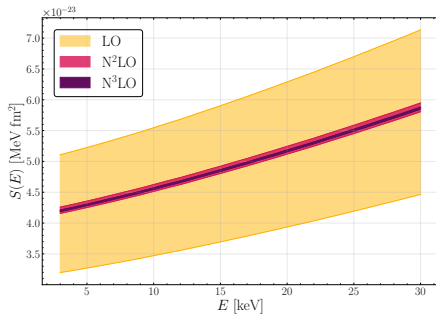
NV

$$\sigma_{\text{NV}}^{\text{Int}}(E) = \frac{S_{\text{NVIIa}}(E) - S_{\text{NVIIb}}(E)}{2}$$

→

$$\sigma_{\text{NV}}^{\text{Int}}(0) = 0.016 \times 10^{-23} \text{ MeV fm}^2$$

Truncation error - current



Interaction fixed at

- $N^3\text{LO}$ in both EMN and NV

Potential	$\sigma_0^{\text{Cur}}(0)$	$\sigma_2^{\text{Cur}}(0)$	$\sigma_3^{\text{Cur}}(0)$
EMN450	0.928	0.049	0.013
EMN500	0.931	0.049	0.011
EMN550	0.933	0.049	0.011
NVIa	0.924	0.049	0.011
NVIb	0.923	0.049	0.011
NVIIa	0.918	0.049	0.011
NVIIb	0.917	0.049	0.011

$\sigma_n^{\text{Cur}}(0)$ in unit of $10^{-23} \text{ MeV fm}^2$

Final results

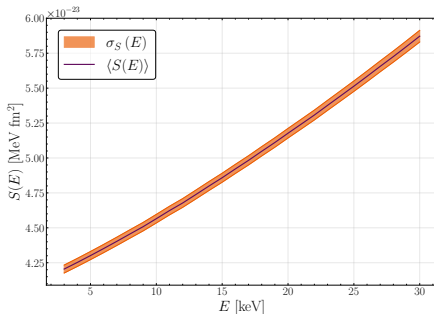
A. Gnech *et al.*, Phys. Rev. C **109**, 035502 (2024): procedure used for $\mu^- + d \rightarrow n + n + \nu_\mu$

$$\langle S(E) \rangle = \sum_i S_i(E) P_i$$

$$P_i = \begin{cases} \frac{1}{6} & \text{for EMN potentials} \\ \frac{1}{8} & \text{for NV potentials} \end{cases}$$

$$\sigma_S^2(E) = \langle S^2(E) \rangle - \langle S(E) \rangle^2 + \sum_i \sigma_i^2(E) P_i$$

$$\sigma_i(E) = \sigma_i^{\text{Int}}(E) + \sigma_i^{\text{Cur}}(E)$$



$$S_{\text{EMN+NV}}(0) = 4.069 (1 \pm 0.007) \times 10^{-23} \text{ MeV fm}^2$$

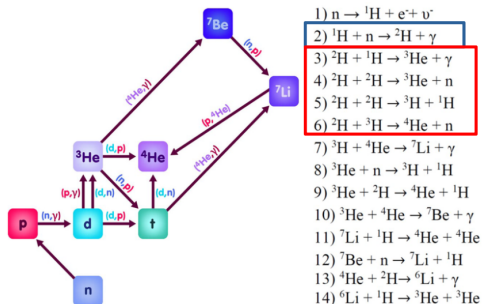
Trunc.

To be compared with: Solar Fusion III

[B. Acharya *et al.*, arXiv: 2405.06470 (2024) to be published in Rev. Mod. Phys.]

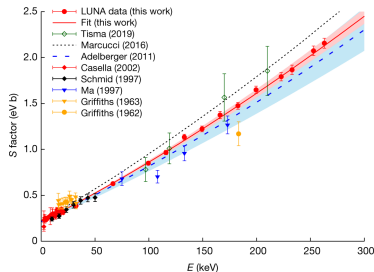
$$S(0) = 4.090 \underset{\text{Trunc.}}{(1 \pm 0.005)} \underset{\text{Syst.}}{\pm 0.010} \times 10^{-23} \text{ MeV fm}^2 \quad ! \text{ no local ChEFT potential}$$

BBN and d abundance: the ${}^2\text{H}(p, \gamma){}^3\text{He}$ reaction



LUNA data

V. Mossa *et al.*, Nature **587**, 210 (2020)



$$10^5 (\text{D}/\text{H})_{\text{BBN}} = 2.52 \pm 0.03 \pm 0.06$$

vs.

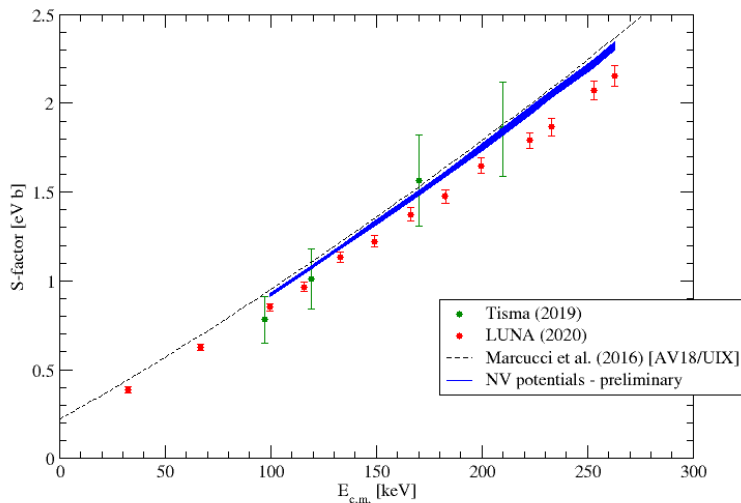
$$10^5 (\text{D}/\text{H})_{\text{exp}} = 2.527 \pm 0.030$$

R.J. Cooke *et al.*, Astrophys. J. **885**, 102 (2018)

----- Phenomenological approach (AV18/UIX) $\rightarrow$ what is the theoretical uncertainty?

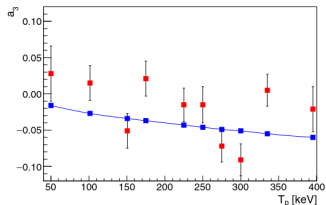
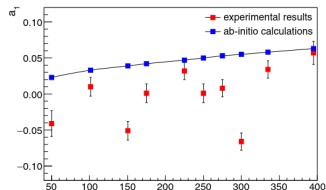
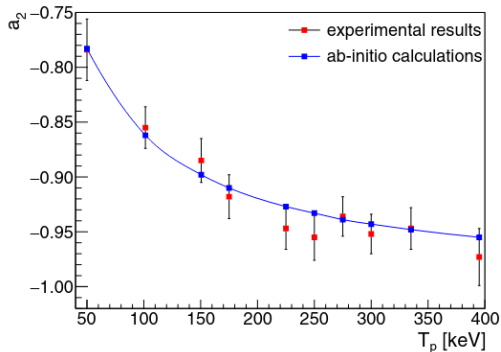
$\Rightarrow$ ChEFT [work in progress]

The $^2\text{H}(p, \gamma)^3\text{He}$ reaction - ChEFT - PRELIMINARY



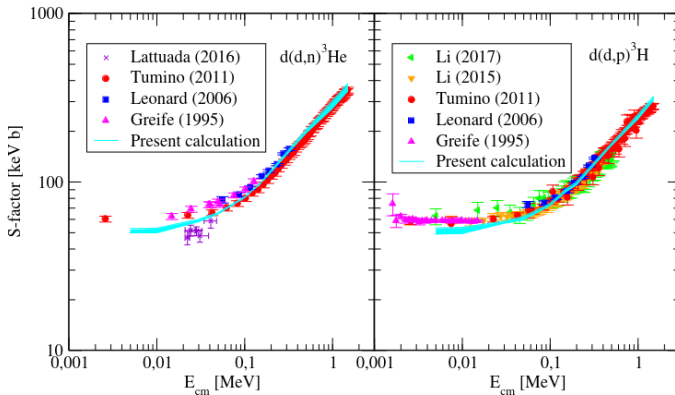
The $^2\text{H}(p, \gamma)^3\text{He}$ reaction - angular distribution

$$\frac{d\sigma}{d\Omega} = A_0 \left[1 + \sum_{\ell=1}^3 a_{\ell} P_{\ell}(\cos \theta_{c.m.}) \right]$$



K. Stöckel *et al.*, Phys. Rev. C **110**, L032801 (2024)

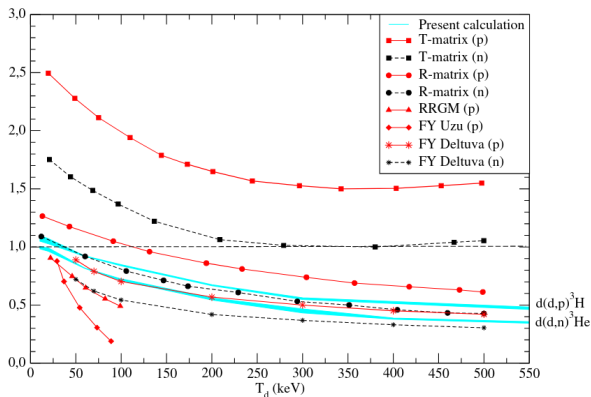
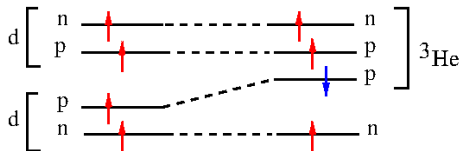
BBN and d abundance: $d(d, p)^3\text{H}$ and $d(d, n)^3\text{He}$



M. Viviani *et al.*, Phys. Rev. Lett. **130**, 122501 (2023)

Not only nuclear astrophysics: the “quintic” suppression factor

$\vec{d}(\vec{d}, n)^3\text{He}$ & $\vec{d}(\vec{d}, p)^3\text{H}$ suppressed in S -wave



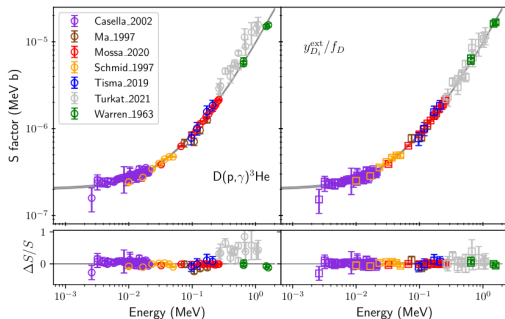
⇒ “neutron lean” reactors

What have we learned:

- Lot of reactions at reach ($A = 2, 3, 4 \dots$)
- In general “good” *ab-initio* predictions
- Predictions accompanied by theoretical error estimate
- Some “puzzles” still there
 - ① $p + d \rightarrow {}^3\text{He} + \gamma$
 - ② A_y -puzzle in $p + d$ elastic scattering

Puzzles: 1) $p + d \rightarrow {}^3\text{He} + \gamma$

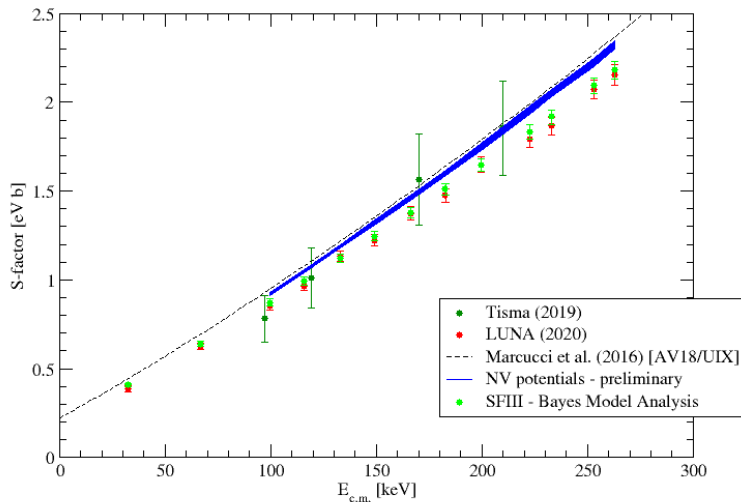
- calculation to be performed with the non-local ChEFT potentials (EMN)
- calculation to be extended to higher energy values



	$E = 262.9 \text{ keV}$	222.8 keV	166.1 keV	99.5 keV
LUNA ¹	2.156(58)	1.791(45)	1.375(36)	0.850(22)
Bayes Model Analysis ²	2.181(45)(3)	1.834(38)(2)	1.376(28)(1)	0.873(18)(4)
NV-ChEFT	2.32(3)	1.96(3)	1.46(2)	0.92(2)

¹ V. Mossa *et al.*, Nature **587**, 210 (2020)

² A. Walker-Loud, <https://github.com/nrp-g/leaner>, Solar Fusion III, arXiv: 2405.06470 (2024)



Puzzles: 2) the A_y -puzzle in $p + d$ elastic scattering

N3LO $2N$ contact interactions

E. Filandri *et al.*, Few-Body Syst. **65**, 57 (2024)

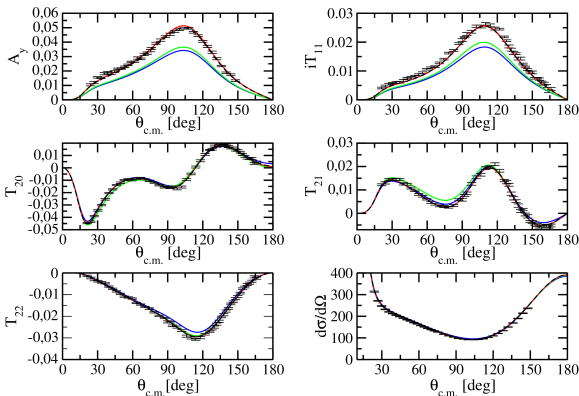
$$V_{2N}^{\text{N3LO}} = \sum_{i=1}^{15} D_i \mathcal{O}_i(\mathbf{k}, \mathbf{Q}) + iD_{16} \mathbf{k} \cdot \mathbf{Q} \mathbf{Q} \times \mathbf{P} \cdot (\boldsymbol{\sigma}_1 - \boldsymbol{\sigma}_2) + D_{17} \mathbf{k} \cdot \mathbf{Q} (\mathbf{k} \times \mathbf{P}) \cdot (\boldsymbol{\sigma}_1 \times \boldsymbol{\sigma}_2)$$

$$\mathbf{k} = \mathbf{p}' - \mathbf{p}$$

$$\mathbf{Q} = \frac{\mathbf{p}' + \mathbf{p}}{2}$$

$$\mathbf{P} = \mathbf{p}_1 + \mathbf{p}_2$$

The $\mathbf{P} = 0$ in $2N$,
but $\mathbf{P} \neq 0$ in $3N$
 $\Rightarrow$ Unitary-transformation
 $\rightarrow V_{3N}$



Conclusions and outlook

Conclusions

Ab-initio approaches are a powerful tool for

- providing **insights** on the nature of **nuclear interactions** and nuclear systems more in general;
- providing **predictions** together with a **robust estimate of theoretical uncertainties**.

Outlook: many ideas

- Move to $A > 3, 4$
- Move to **hypernuclear systems**
 - ChEFT for hypernuclear interactions and currents
 - *ab-initio* approaches for bound and scattering hypernuclear states

Collaborators

- A. Kievsky and M. Viviani (INFN-Pisa)
- **A. Grassi** (Post-doc, PRIN2022, Pisa Univ.)
- **V. Barlucchi** (MS, Pisa Univ. → PhD Mainz Univ. + Pisa Univ.)
- **E. Proietti** (PhD, Pisa Univ.)
- **M. Sagina** (MS → PhD, Pisa Univ.)
- **E. Filandri** (Post-doc, ECT*-Trento)
- L. Girlanda (Salento Univ.)
- A. Gnech (ODU and JLab-USA)
- M. Piarulli (WUSTL)



Finanziato
dall'Unione europea
NextGenerationEU



Ministero
dell'Università
e della Ricerca



Italiadomani
PIANO NAZIONALE
DI RIPARAZIONE E RESILIENZA



UNIVERSITÀ DI PISA

Collaborators

- A. Kievsky and M. Viviani (INFN-Pisa)
- **A. Grassi** (Post-doc, PRIN2022, Pisa Univ.)
- **V. Barlucchi** (MS, Pisa Univ. → PhD Mainz Univ. + Pisa Univ.)
- **E. Proietti** (PhD, Pisa Univ.)
- **M. Sagina** (MS → PhD, Pisa Univ.)
- **E. Filandri** (Post-doc, ECT*-Trento)
- L. Girlanda (Salento Univ.)
- A. Gnech (ODU and JLab-USA)
- M. Piarulli (WUSTL)

Thank you for your attention!



Finanziato
dall'Unione europea
NextGenerationEU



Ministero
dell'Università
e della Ricerca



Italiadomani
PIANO NAZIONALE
DI RIPARAZIONE E RESILIENZA



UNIVERSITÀ DI PISA