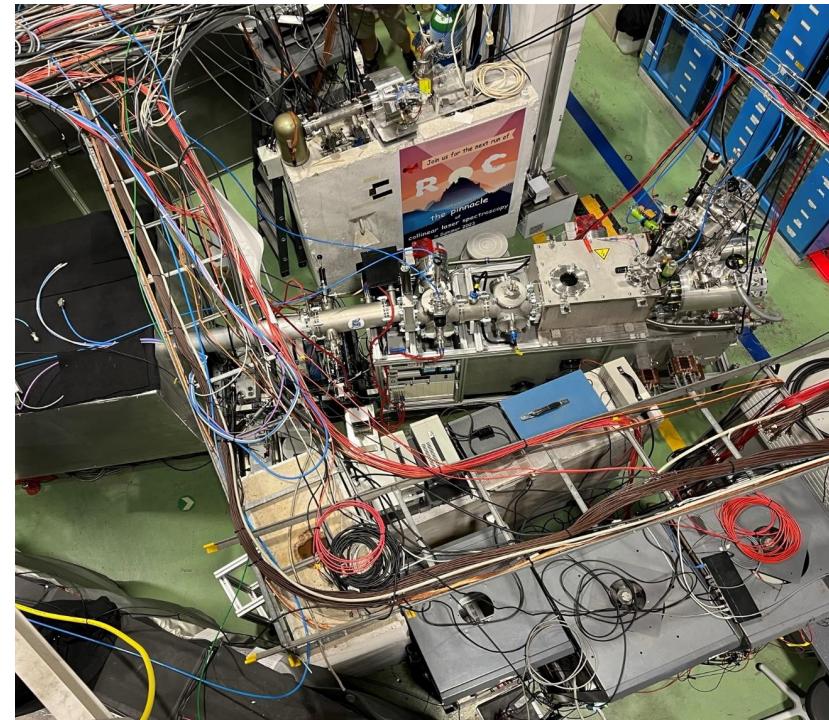
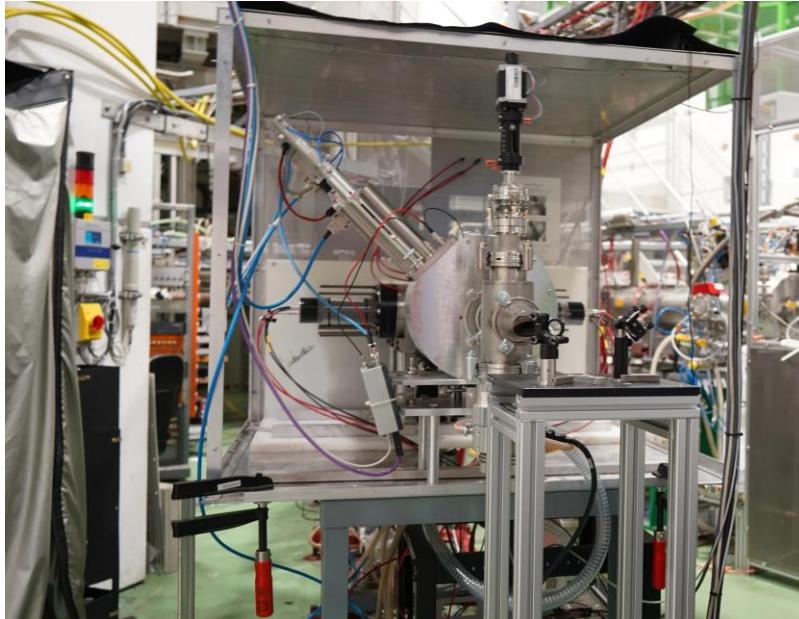


Magic Moments: Exotic Calcium Isotopes in Laser Light



TECHNISCHE
UNIVERSITÄT
DARMSTADT

Wilfried Nörtershäuser
On behalf of the COLLAPS Collaboration



THE UNIVERSITY
of LIVERPOOL

GEFÖRDERT VOM



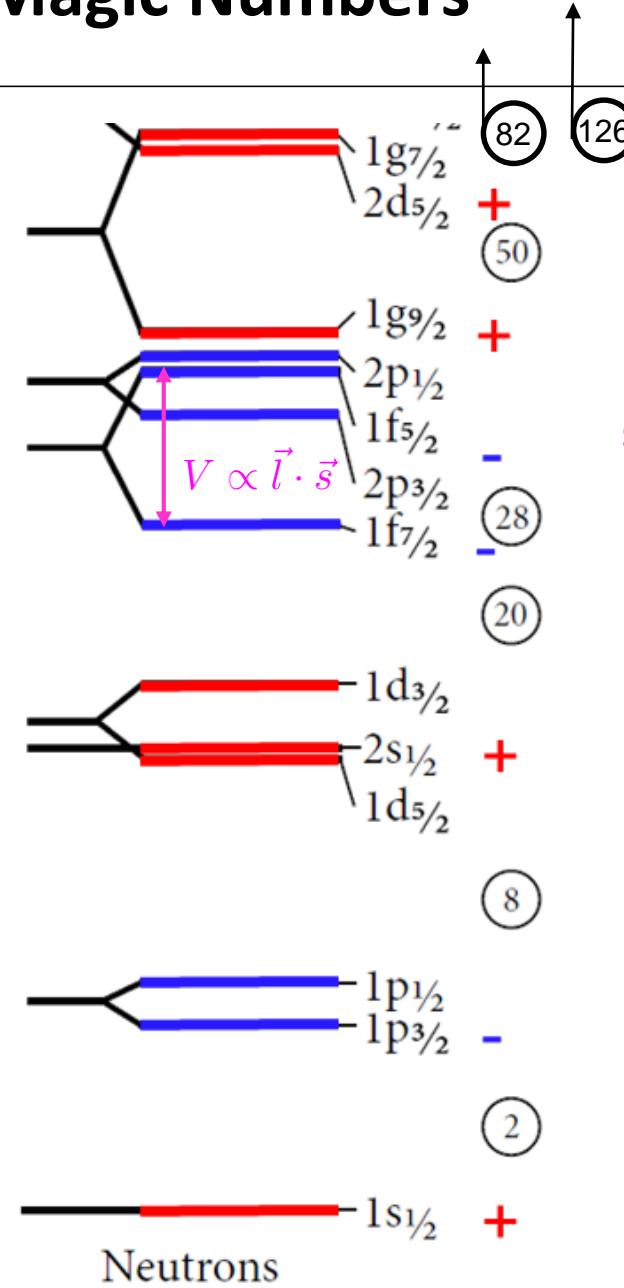
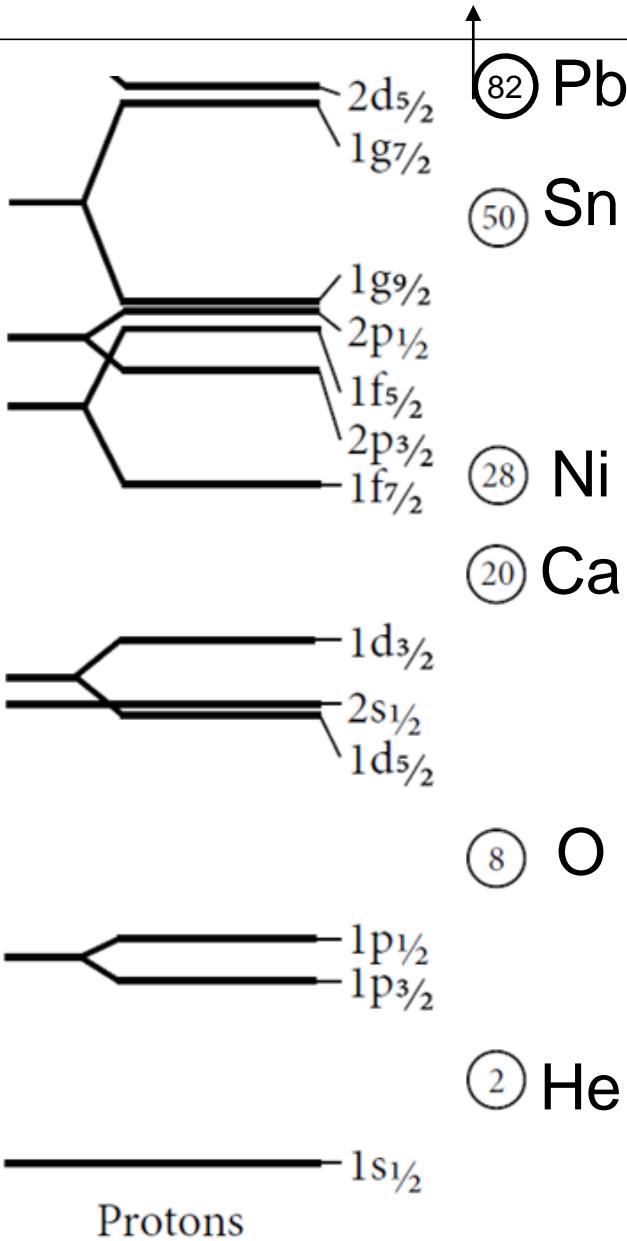
Bundesministerium
für Bildung
und Forschung

Outline



- Basics and Motivation
 - Nuclear Physics: Fingerprints of Magic Numbers
 - Atomic Physics: Nuclear Properties in Optical Spectra
- Nuclear Charge Radii of Calcium Isotopes
 - $^{36-52}\text{Ca}$: New lessons for theory and a tantalizing question
 - Magic moment after a 10-years quest: Laser spectroscopy of $^{53,54}\text{Ca}$
- The N=28 Isotones and the Charge Radius of ^{56}Ni

Magic Numbers

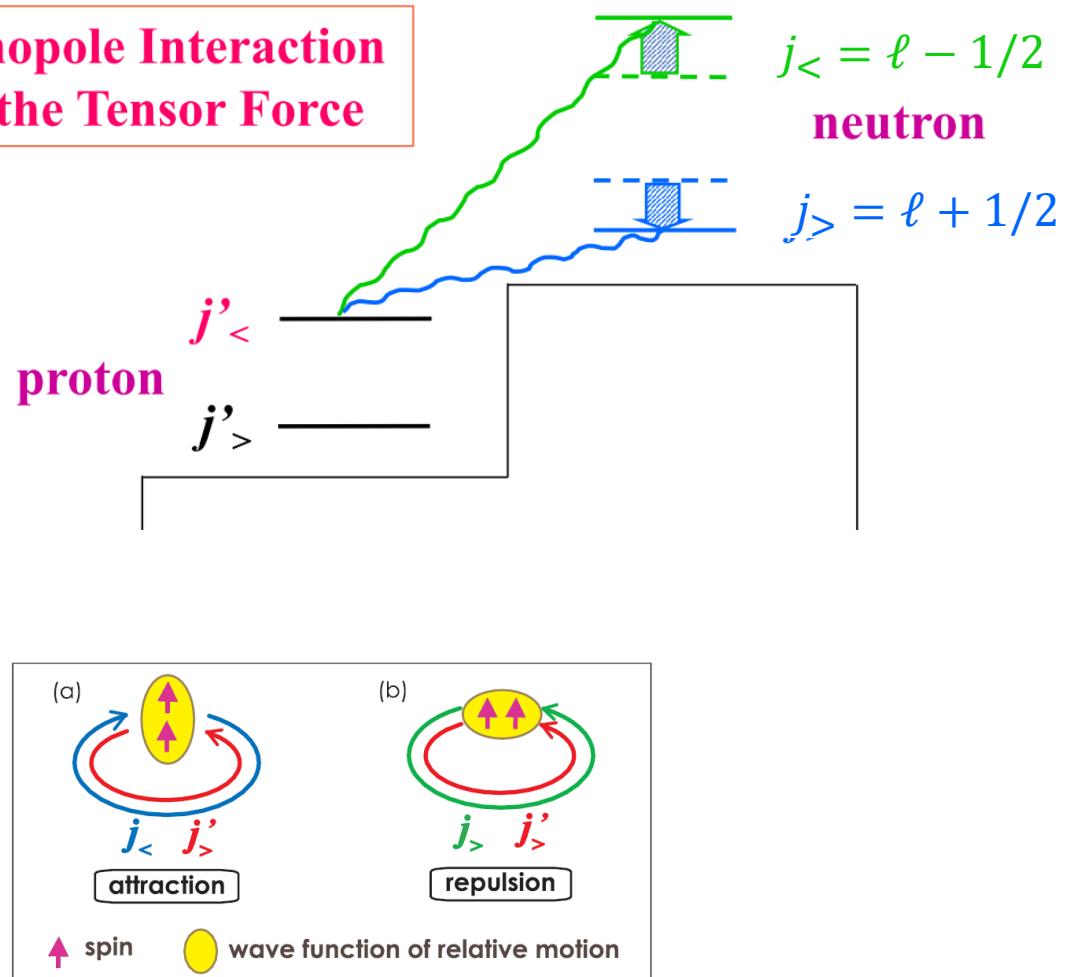

 $^{208}\text{Pb}, ^{132}\text{Sn}$
 $^{100}\text{Sn}, ^{78}\text{Ni}$
 $^{56}\text{Ni}, ^{48}\text{Ca}$
 ^{40}Ca
 ^{16}O
 ^4He

These “traditional”
magic numbers arise
from investigations of
nuclei close to stability.

Shell Evolution in the Calcium Region

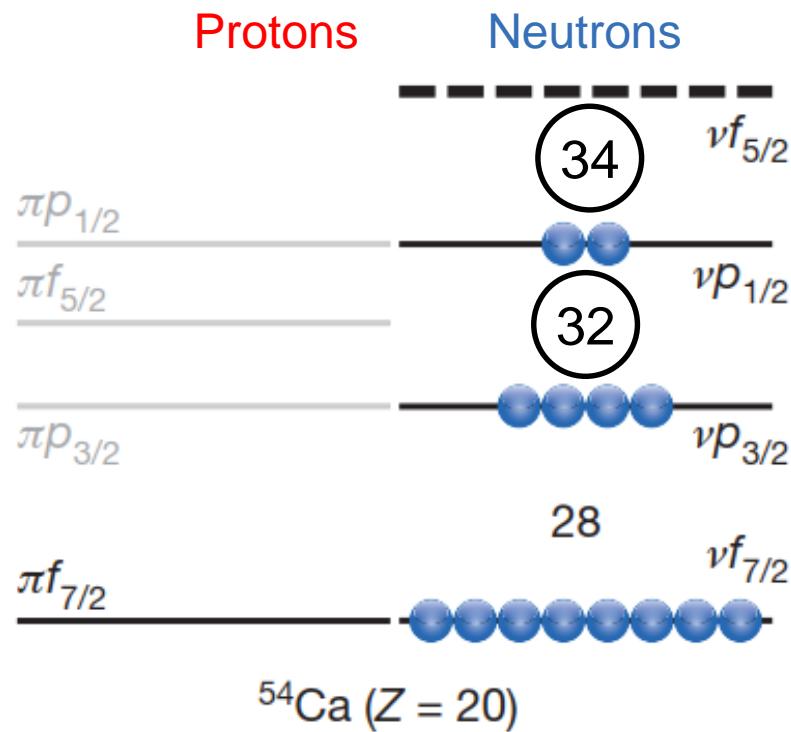


Monopole Interaction of the Tensor Force



Otsuka, PRL 95, 232502 (2005)

How to recognize shell
closures experimentally ?

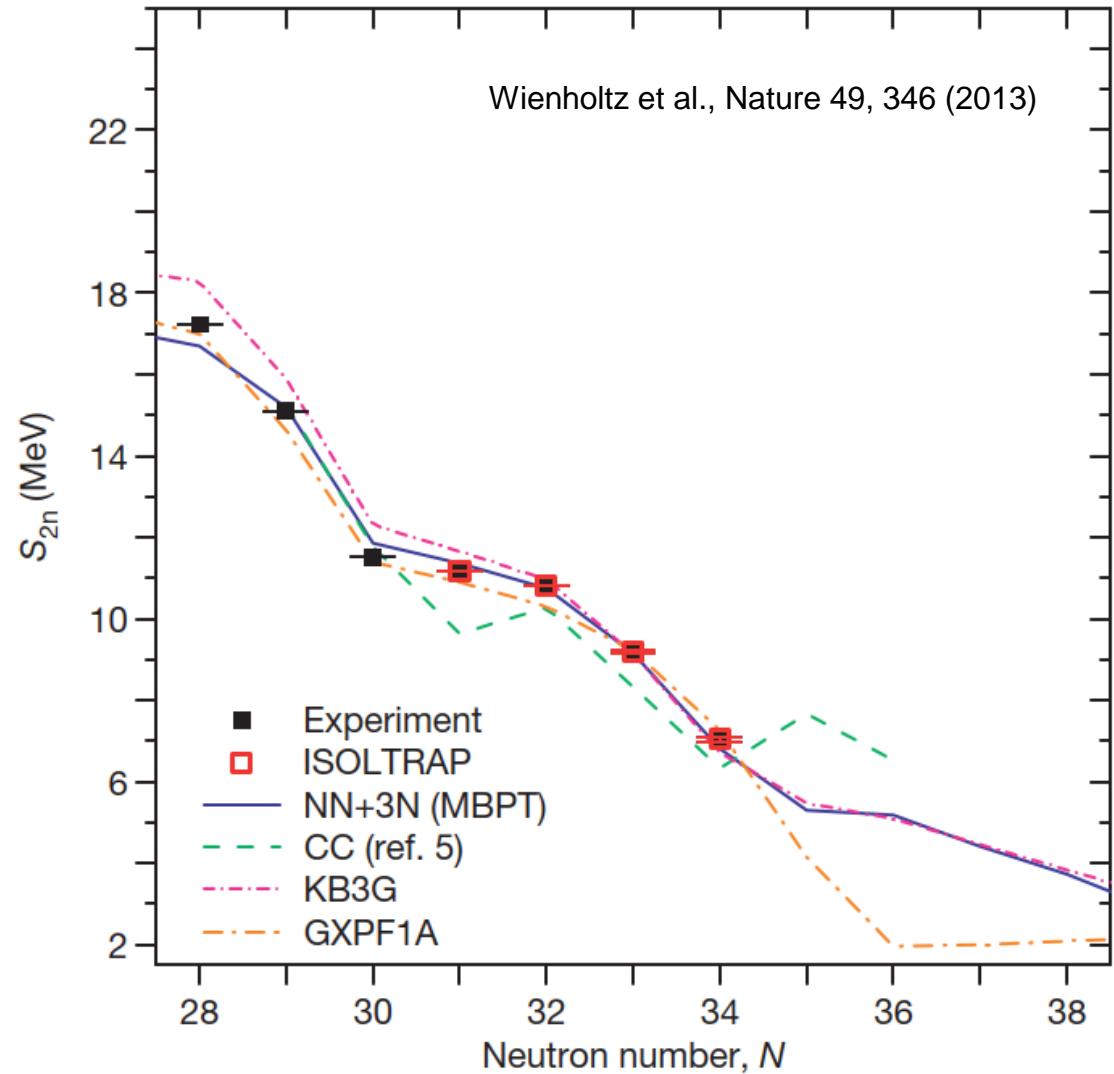


Steppenbeck et al., Nature 502, 207 (2013)

Fingerprints of (Doubly) Magic Nuclei

Decrease in S_{2n} energies

$$S_{2n} = E_B(Z, N) - E_B(Z, N - 2)$$

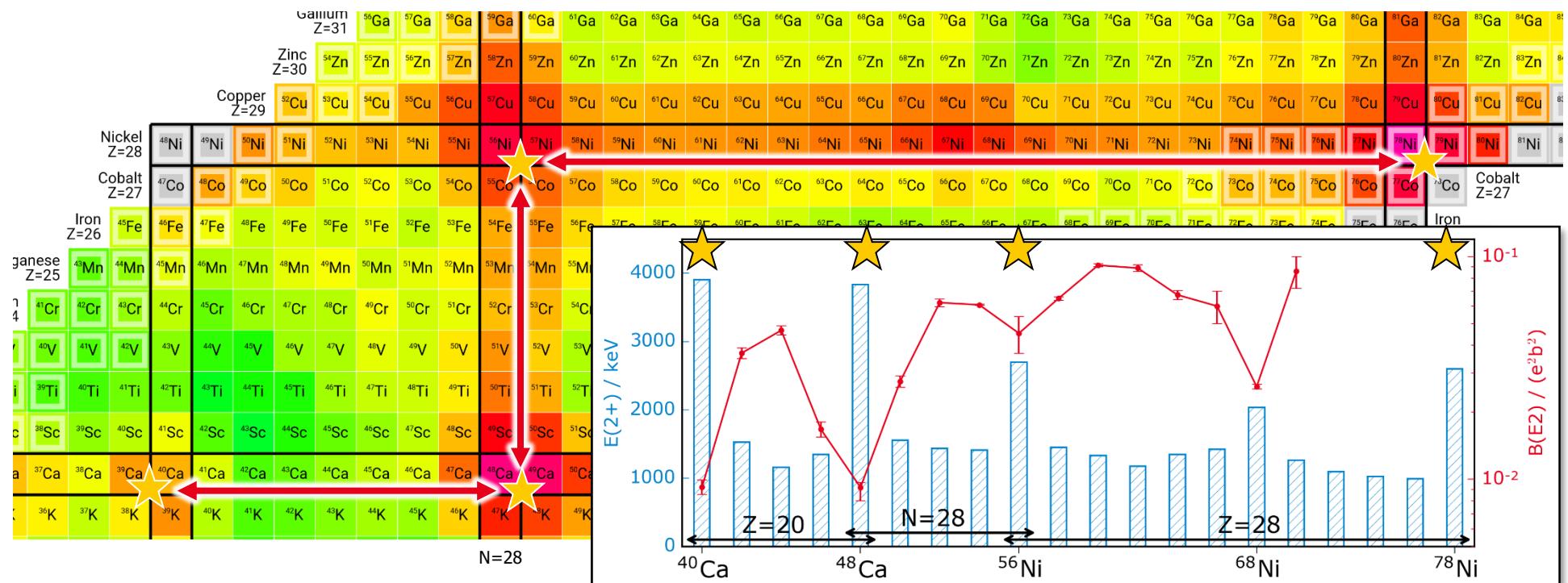
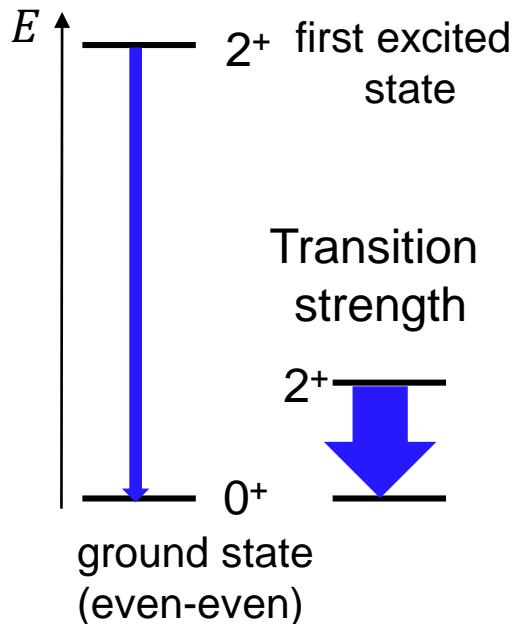


Fingerprints of (Doubly) Magic Nuclei



Decrease in S_{2n} energies

High energy of first excited (2^+) state
Weak transition $2^+ \rightarrow 0^+$



$$B(E2) \propto |\langle I_f | M(E2) | I_i \rangle|^2$$

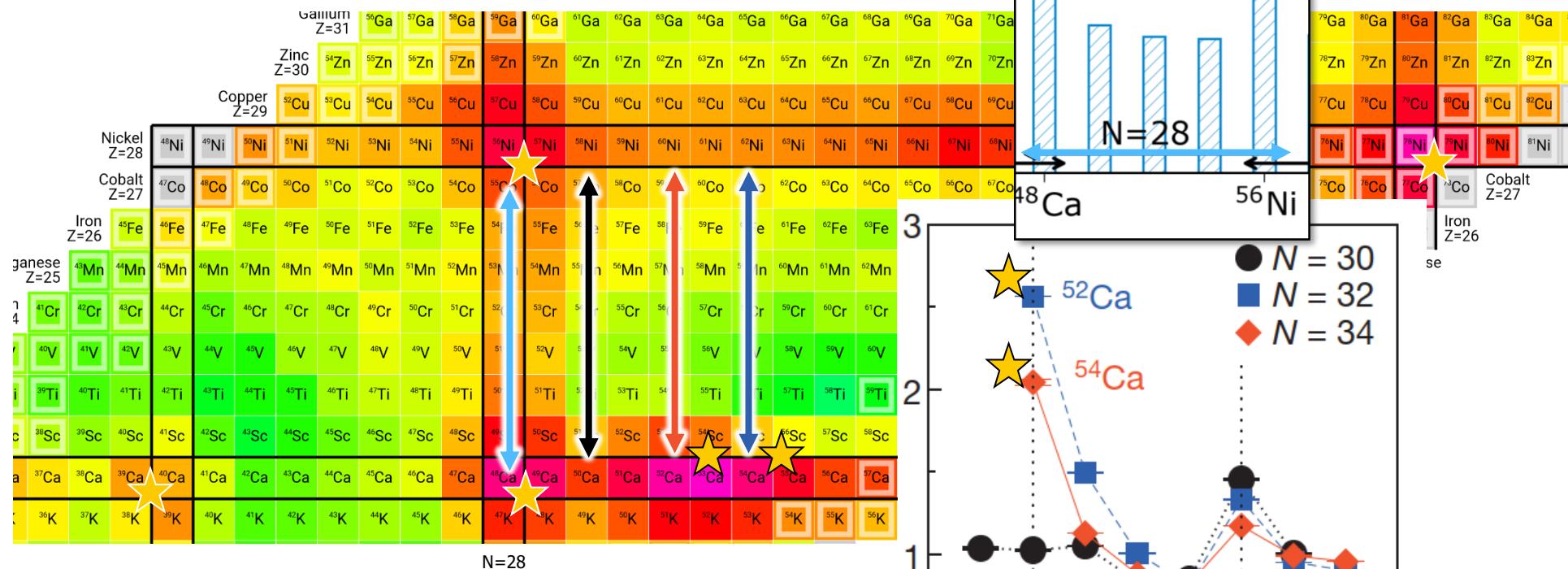
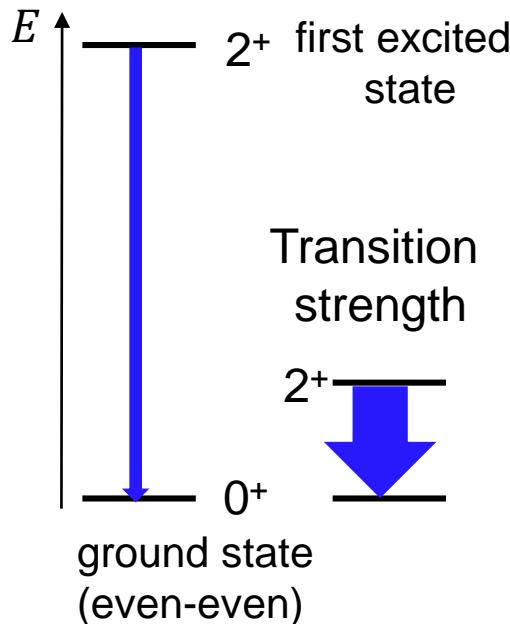
$$\tau[\text{ns}] = \frac{1}{1.22 E_\gamma^5 [\text{MeV}] B(E2) [\text{e}^2 \text{fm}^4]}$$

Fingerprints of (Doubly) Magic Nuclei



Decrease in S_{2n} energies

High energy of first excited (2^+) state
Weak transition $2^+ \rightarrow 0^+$



Fingerprints of (Doubly) Magic Nuclei

Decrease in S_{2n} energies

High energy of first excited (2^+) state
Transition strength $2^+ \rightarrow 0^+$ weak

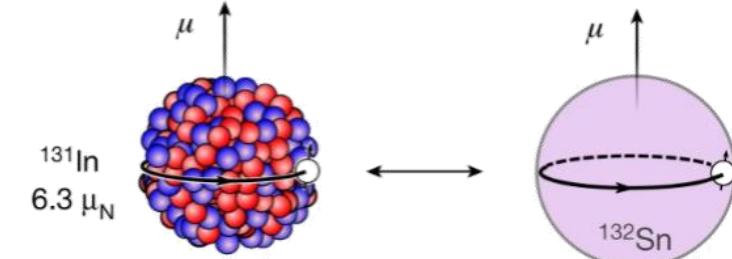
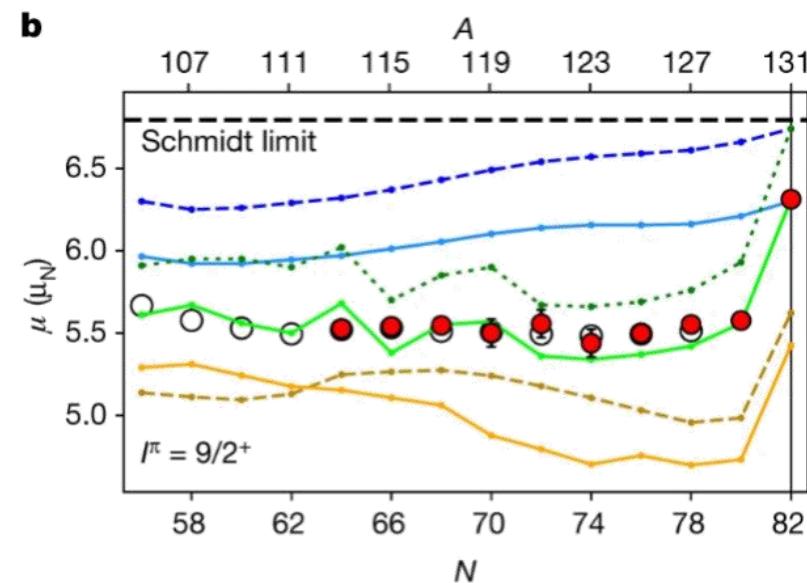
Magnetic moments of doubly magic ± 1

unpaired
nucleon

closed shell

single hole

- Experiment
- Experiments in literature
- · · · VS-IMSRG 1.8/2.0(EM)
- · · · VS-IMSRG N^2LO_{GO}
- · · · DFT HFB without time-odd fields
- · · · DFT HFB with time-odd fields
- · · · DFT HF without time-odd fields
- · · · DFT HF with time-odd fields



Vernon *et al.*, *Nature* **607**, 260 (2022)

Magnetic moments and quadrupole moments of doubly magic ± 1 isotopes are more shell-model like (Schmidt value) than isotopes further away.

Fingerprints of (Doubly) Magic Nuclei

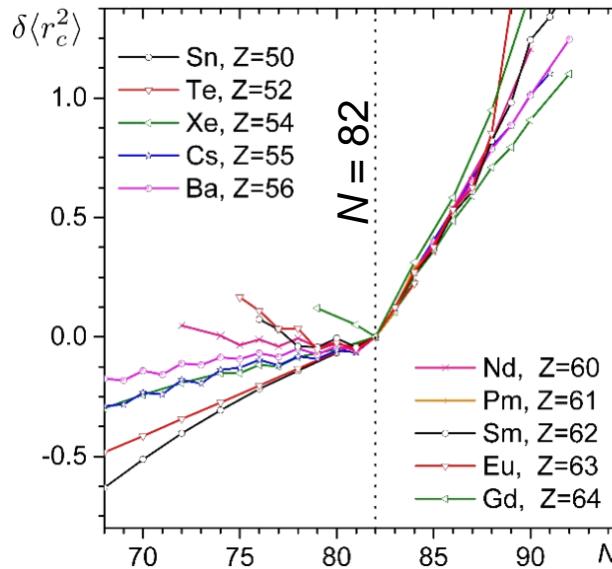
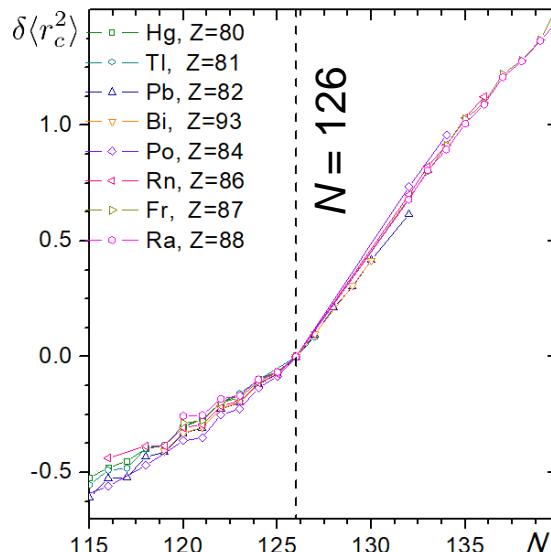


Decrease in S_{2n} energies

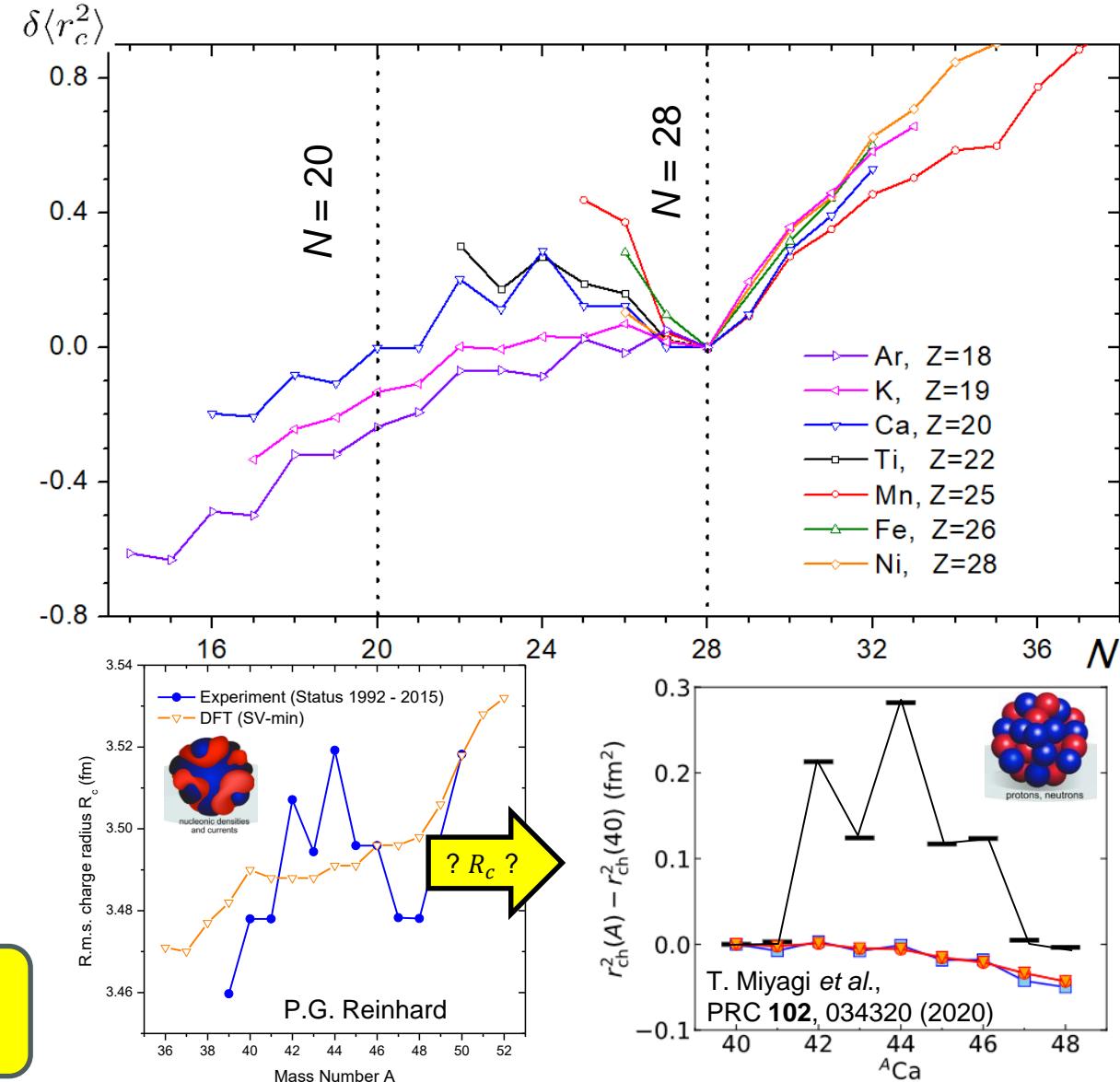
High energy of first excited (2^+) state
Transition strength $2^+ \rightarrow 0^+$ weak

Magnetic moments of doubly magic $\square 1$

Charge radius „kink“



How can we determine
charge radii and moments
of short-lived isotopes?

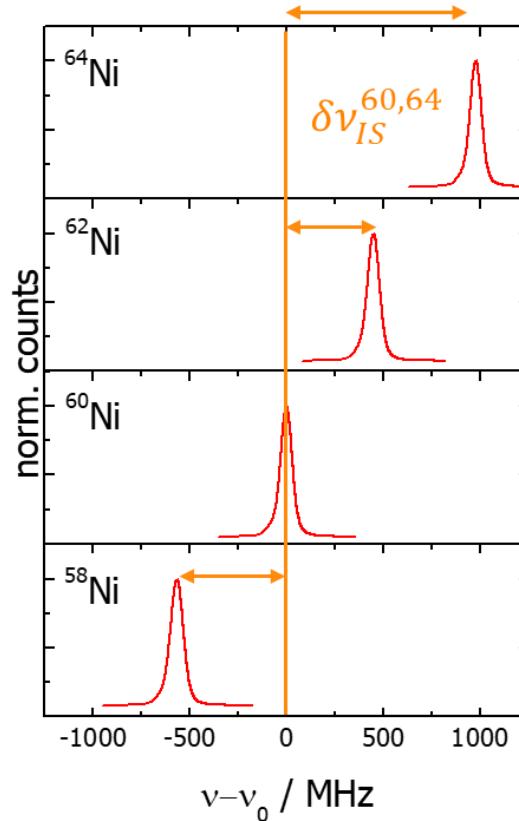


P.G. Reinhard

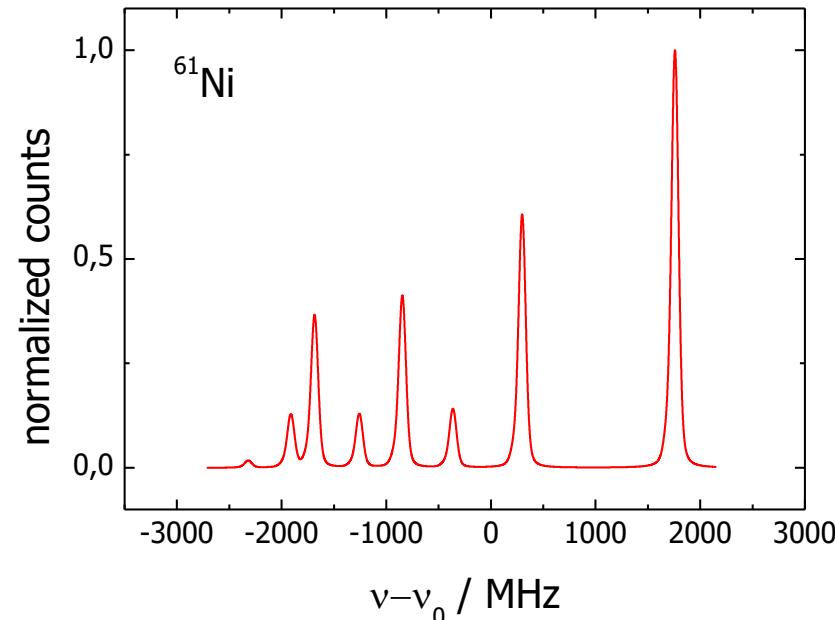
T. Miyagi et al.,
PRC 102, 034320 (2020)

Nuclear Observables in the Optical Hyperfine Structure

Isotope Shift



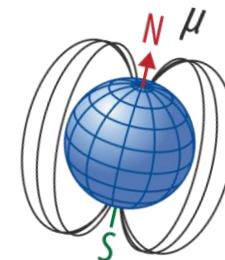
Hyperfine Splitting



Spin I



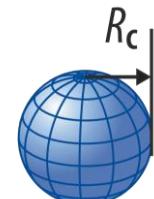
Magnetic Dipole Moment μ



Electric Quadrupole Moment Q_s

Nuclear Size

$$\delta \langle r_c^2 \rangle^{AA'}$$



$$Q_s = 0$$



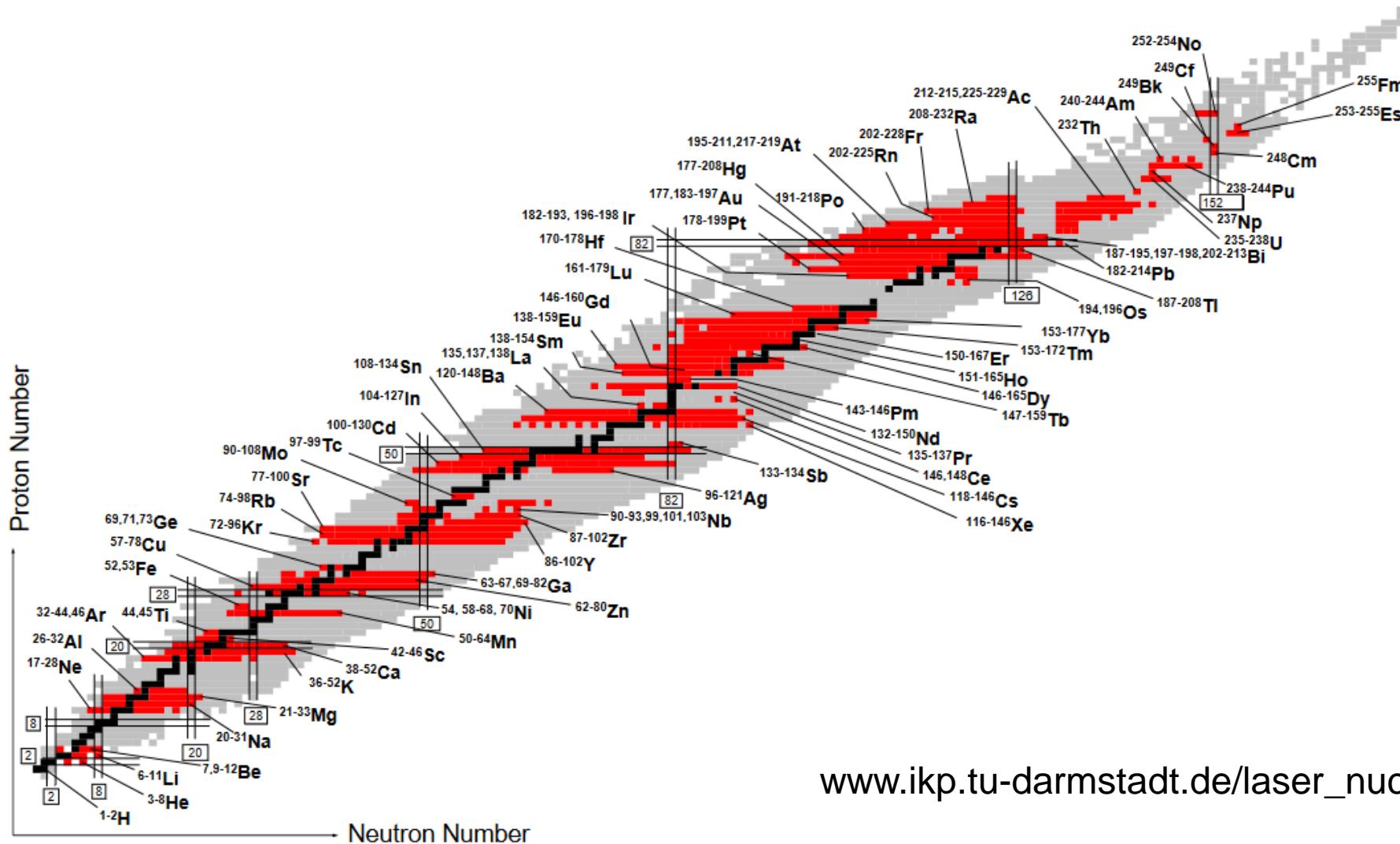
$$Q_s < 0$$



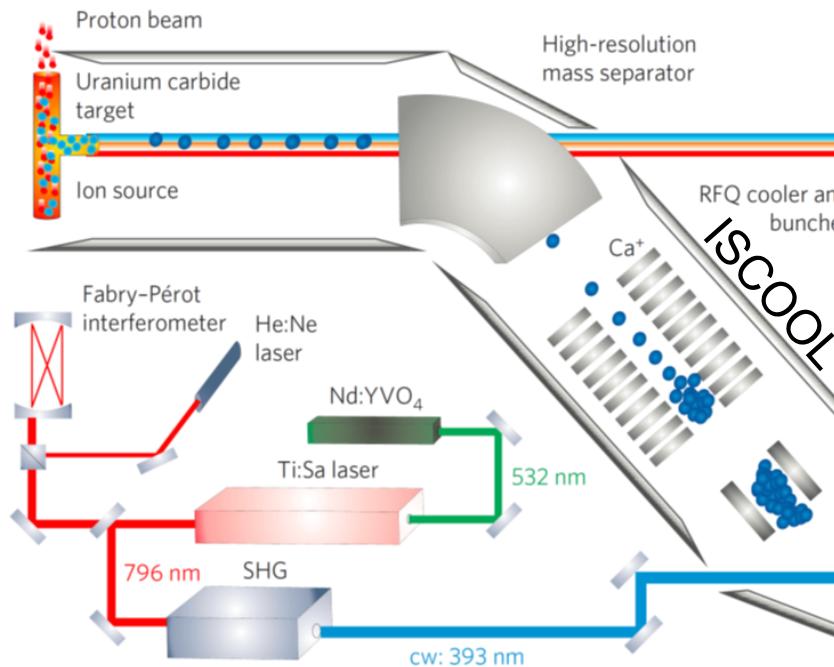
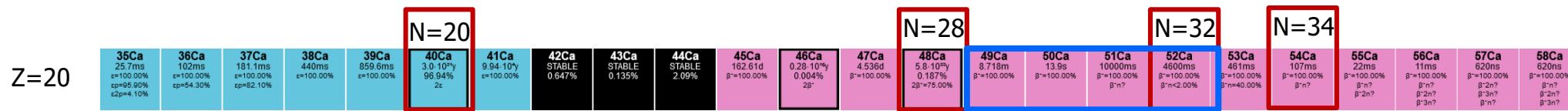
$$Q_s > 0$$



The Laser Nuclear Chart



Measurement of $^{49-52}\text{Ca}$ using Collinear Laser Spectroscopy and Bunched-Beam Detection



ISCOOL (typical)

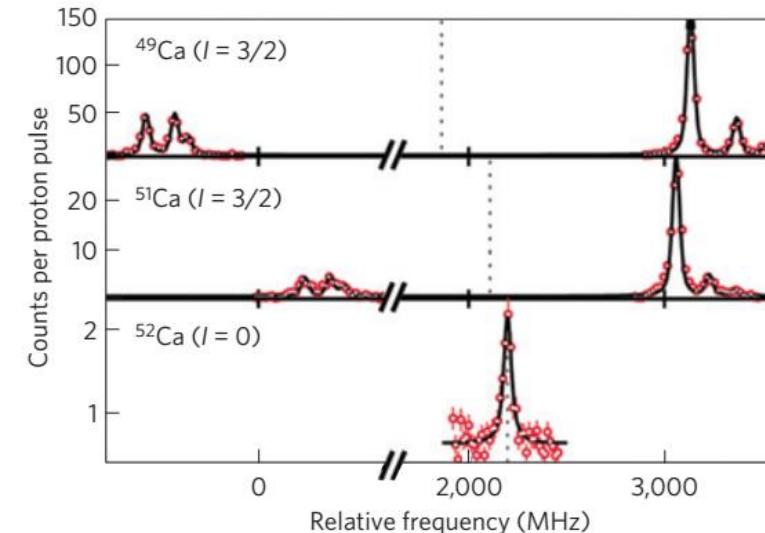
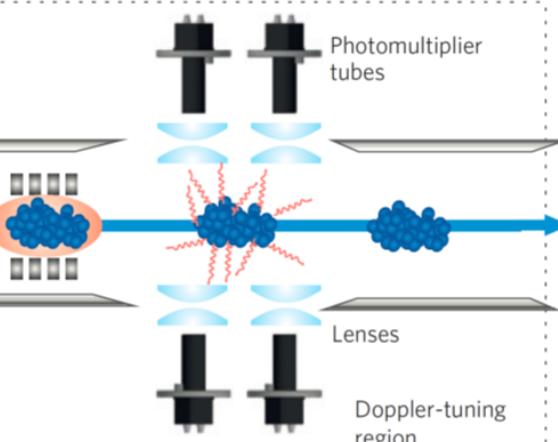
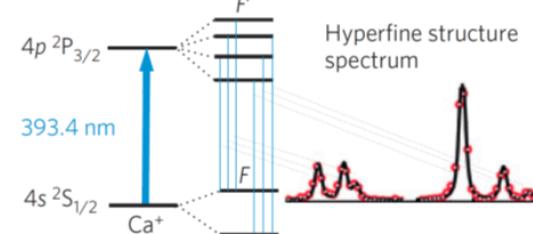
$$t_{\text{acc}} = 50 \text{ ms}$$

$$t_{\text{bunch}} = 5 \mu\text{s}$$

$$\frac{t_{\text{bunch}}}{t_{\text{acc}}} = \frac{0.005 \text{ ms}}{50 \text{ ms}} = 10^{-4}$$

background reduction

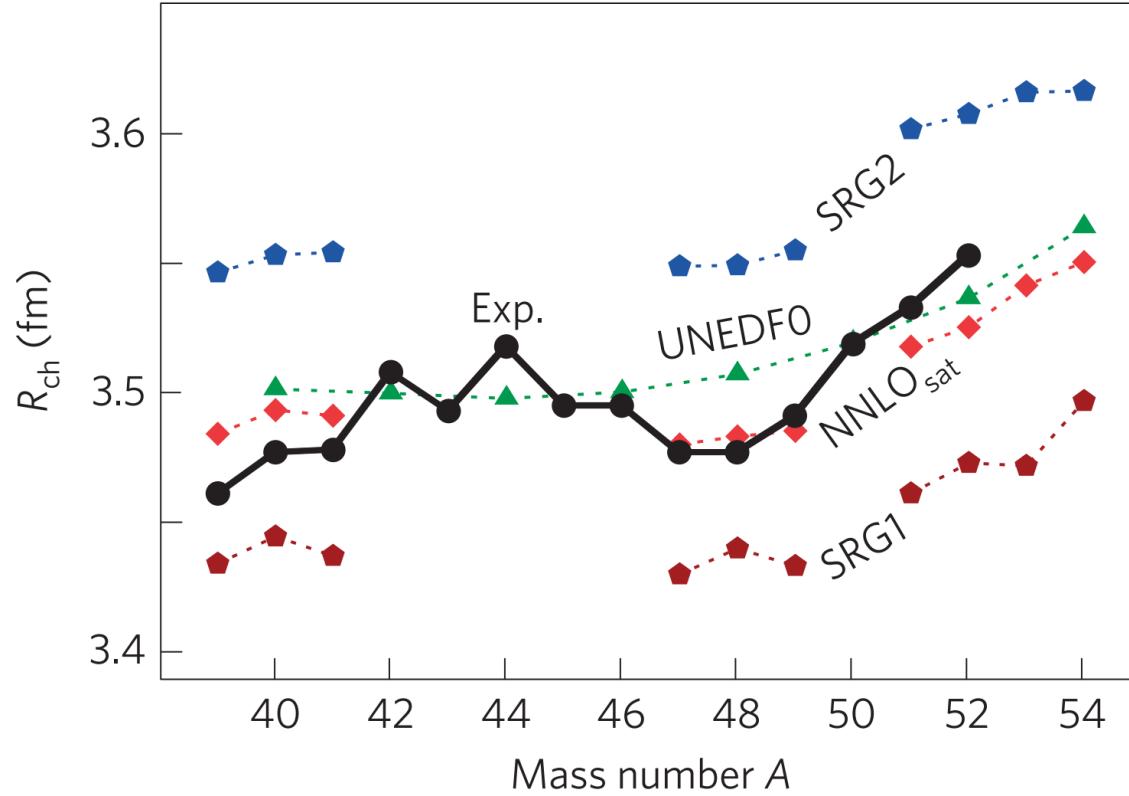
Level scheme for resonance excitation



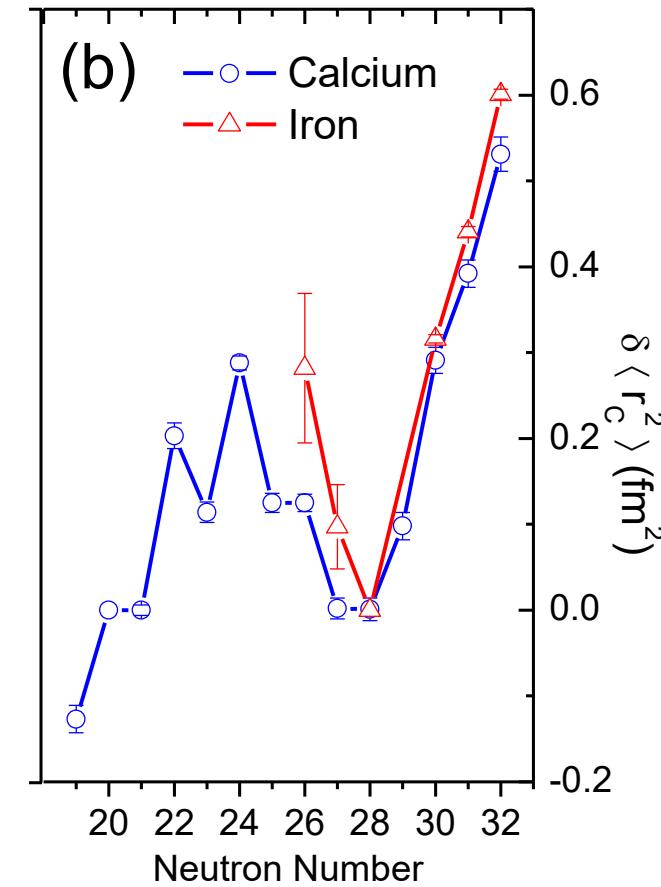
R.F. Garcia Ruiz et al.,
Nature Physics **12**, 594 (2016)

Sensitivity improved
100 x

Unexpected Strong Rise up to ^{52}Ca

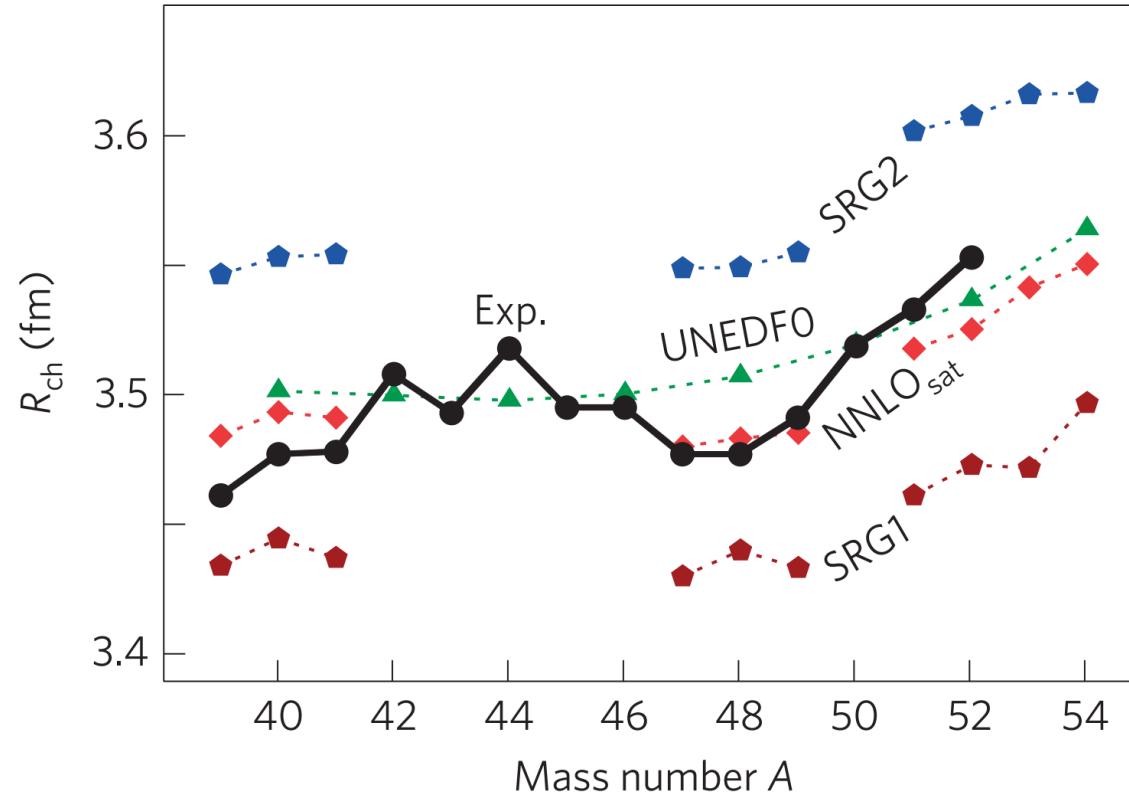


EFT-based interaction NNLO_{sat} does a particularly good job for absolute charge radii.

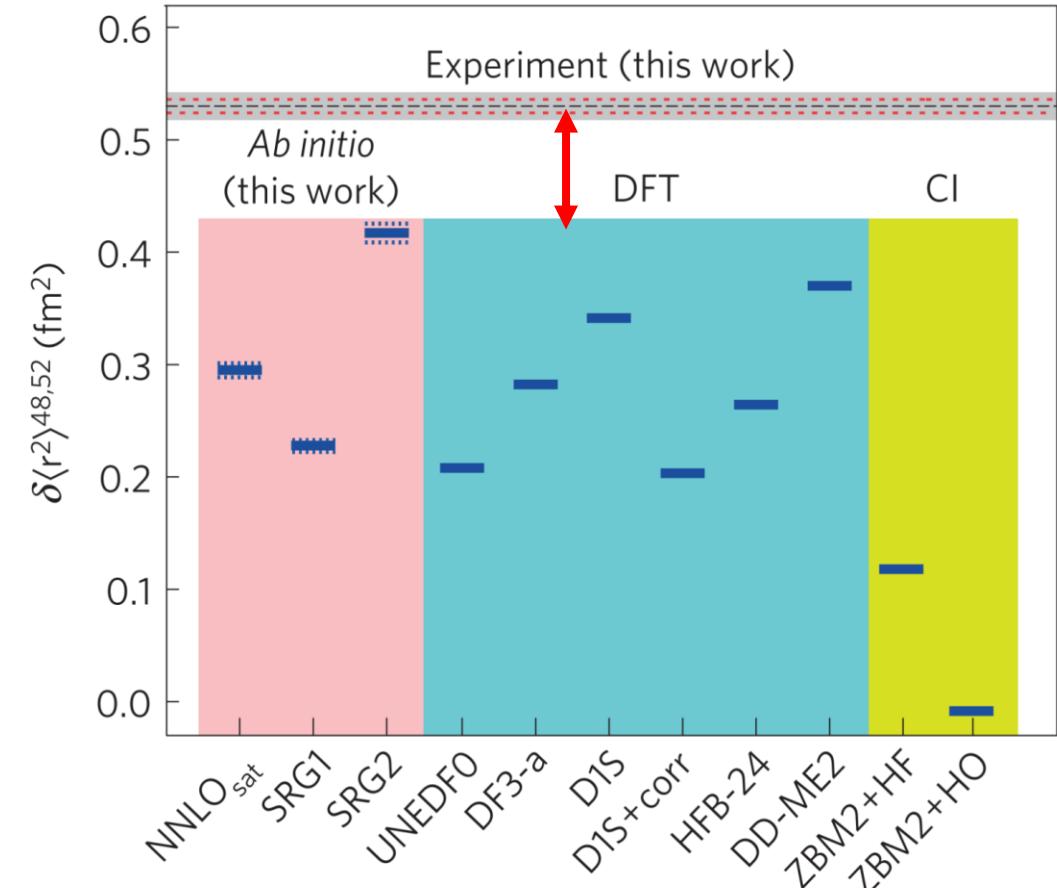


No significant difference to charge radii with more protons in $f_{7/2}$ shell
→ no indication for a shell-closure at $N=32$

Unexpected Strong Rise up to ^{52}Ca



EFT-based interaction NNLO_{sat} does a particularly good job for absolute charge radii.



But all theories underestimate rise of charge radii.

The Fayans Functional

$$\mathcal{E}_{\text{Fy}} = \mathcal{E}_{\text{Fy}}^{\text{v}}(\rho) + \mathcal{E}_{\text{Fy}}^{\text{s}}(\rho) + \mathcal{E}_{\text{Fy}}^{\text{s}}(\rho, \mathbf{J}) + \mathcal{E}_{\text{Fy}}^{\text{pair}}(\rho, \breve{\rho})$$

$$\mathcal{E}_{\text{Fy},q}^{\text{pair}} = \frac{4\varepsilon_F}{3\rho_{\text{sat}}} \check{\rho}_q^2 \left[f_{\text{ex},+}^\xi + h_{1+}^\xi x_{\text{pair}}^\gamma + h_\nabla^\xi r_s^2 (\nabla x_{\text{pair}})^2 \right]$$

↑ pairing density gradient term in pairing density

Fayans et al., Nucl. Phys. A 676, 49 (2000)

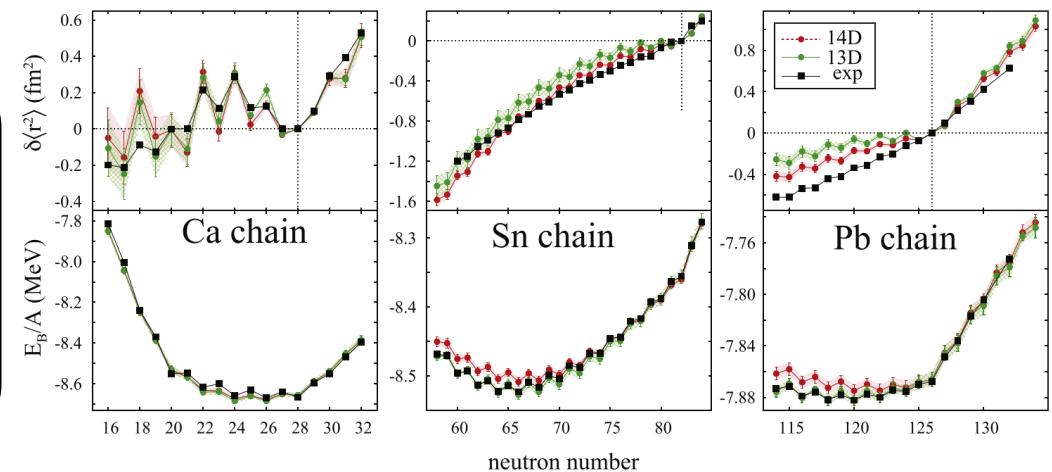
$$\rho_0 = \rho_n + \rho_p, \rho_1 = \rho_n - \rho_p$$

isoscalar isovector

$$x_t = \frac{\rho_t}{\rho_{\text{sat}}}, \quad x_{\text{pair}} = \frac{\rho_0}{\rho_{\text{pair}}}$$

normalized densities

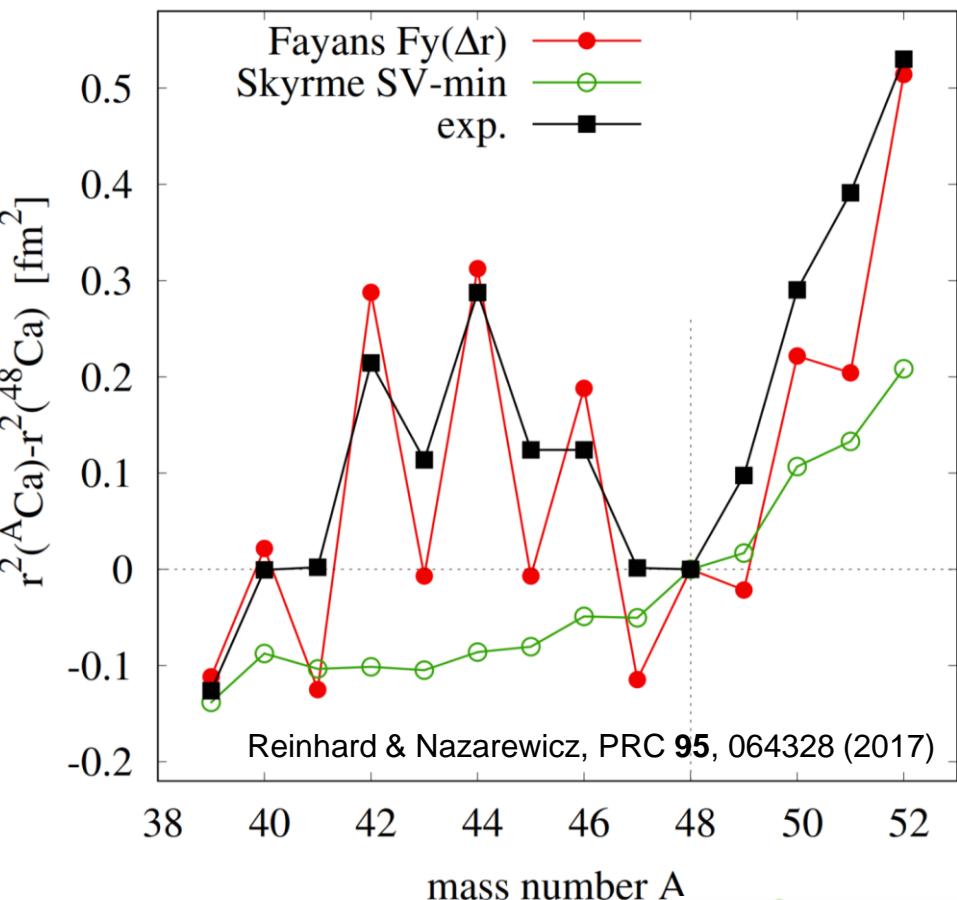
Can this
optimized
functional be
used in other
parts of the
nuclear chart ?



Reinhard et al., J. Phys. G 51, 105101 (2024)

Used Ca radii for optimization:

$$\delta \langle r_c^2 \rangle^{40,48}, \quad \delta \langle r_c^2 \rangle^{44,48}, \quad \delta \langle r_c^2 \rangle^{52,48}$$



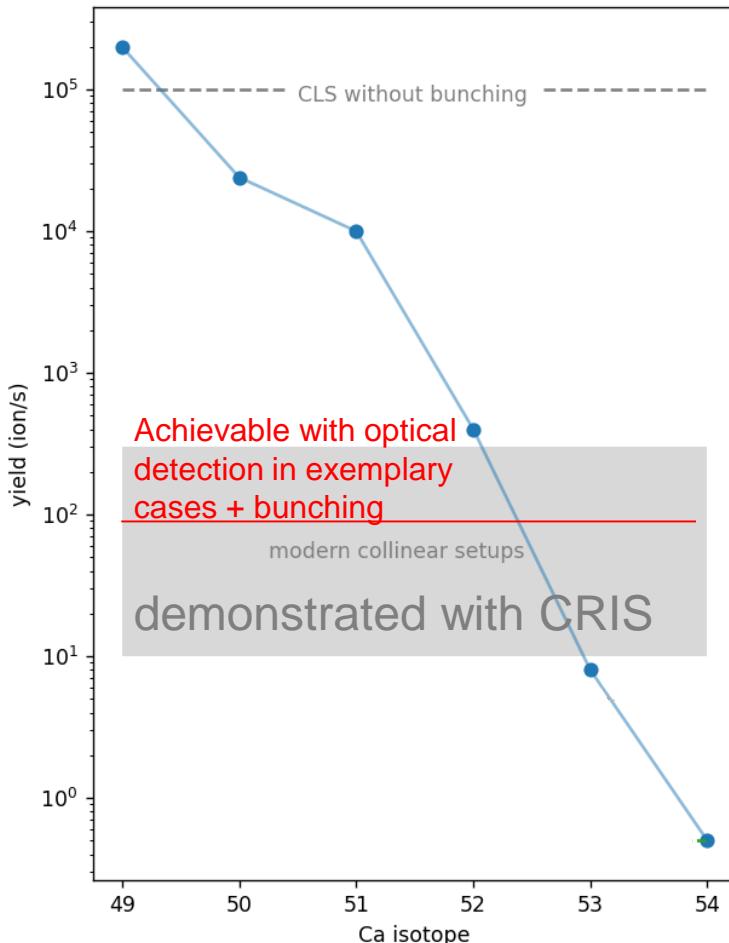
Reinhard & Nazarewicz, PBC 95, 064328 (2017)

Towards ^{54}Ca



Expected production rates: 10^{-10} to 10^{-1} ions/s

Ca yields and sensitivity of similar setups



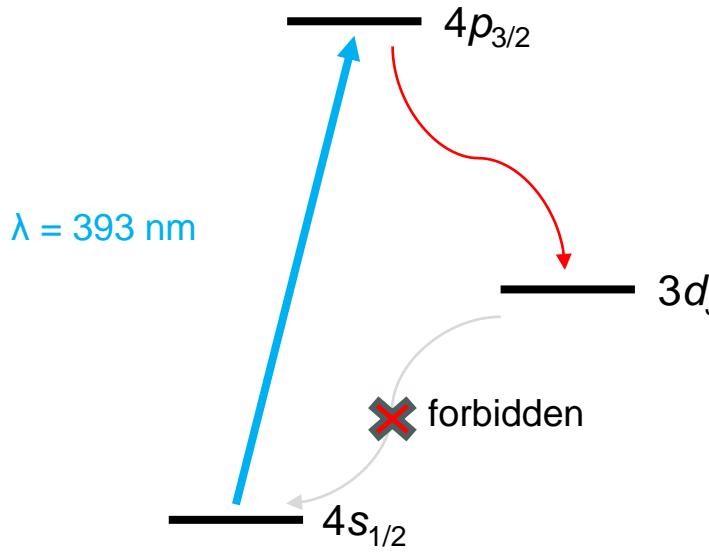
Radiation detection after
Optical pumping and state-selective
Charge exchange
Acronym by P. Lievens *et al.*



The ROC Detection Principle

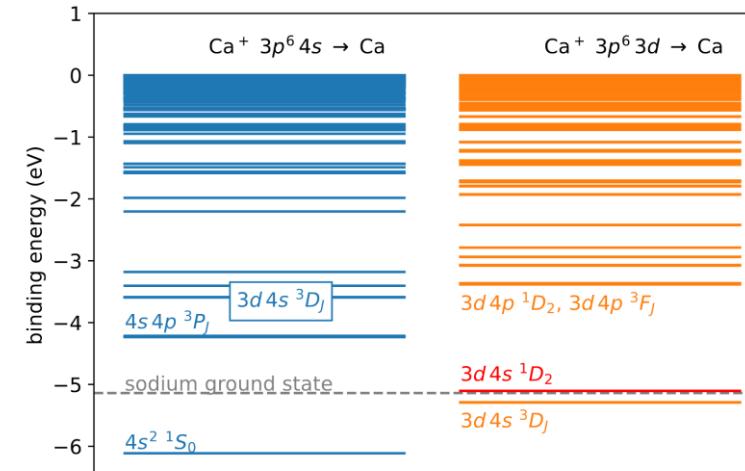


Relevant Ca^+ level scheme



Optical Pumping

Reaction: $\text{Ca}^+ + \text{Na} \rightarrow \text{Ca} + \text{Na}^+ + \Delta E$
 ΔE energy depends on initial ionic state!

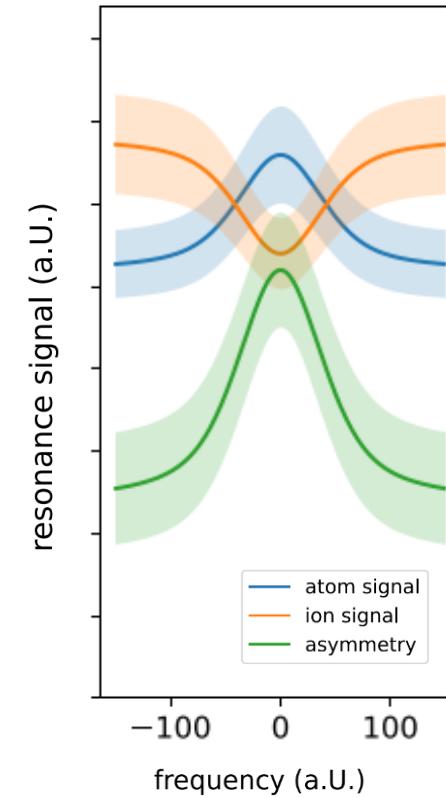


Charge exchange from $\text{Ca}^+ d$ -state is “quasi-resonant” with the sodium ground state ($\Delta E \approx 0$)

→ Larger cross-section, higher CE probability

State-Selective Charge Exchange

$$\varepsilon_{\text{cec,gs}} = 30\%$$

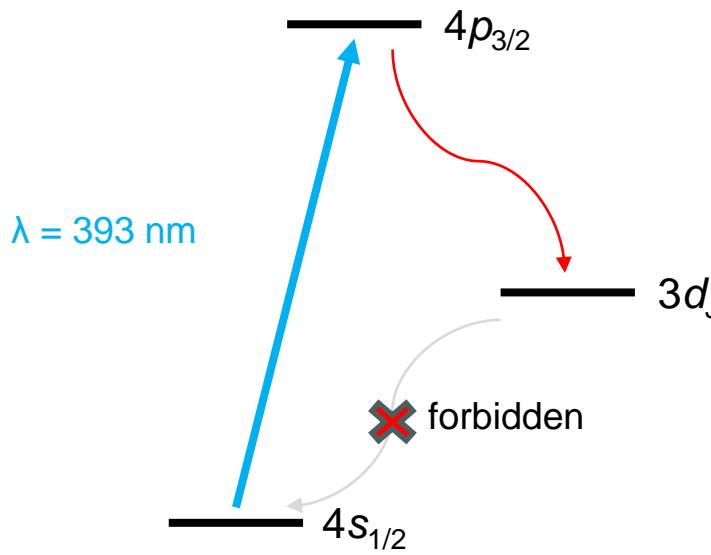


$$s_{\text{asym}} = \frac{s_{\text{atom}} - s_{\text{ion}}}{s_{\text{atom}} + s_{\text{ion}}}$$

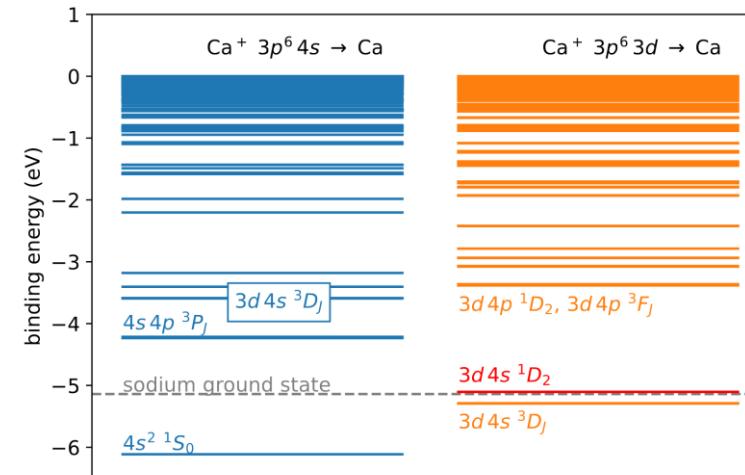
ROC (Radiation detection after Optical pumping and state selective Charge exchange)

The ROC Detection Principle

Relevant Ca^+ level scheme



Reaction: $\text{Ca}^+ + \text{Na} \rightarrow \text{Ca} + \text{Na}^+ + \Delta E$
 ΔE energy depends on initial ionic state!



Charge exchange from $\text{Ca}^+ d$ -state is “quasi-resonant” with the sodium ground state ($\Delta E \approx 0$)
 → Larger cross-section, higher CE probability

Optical Pumping

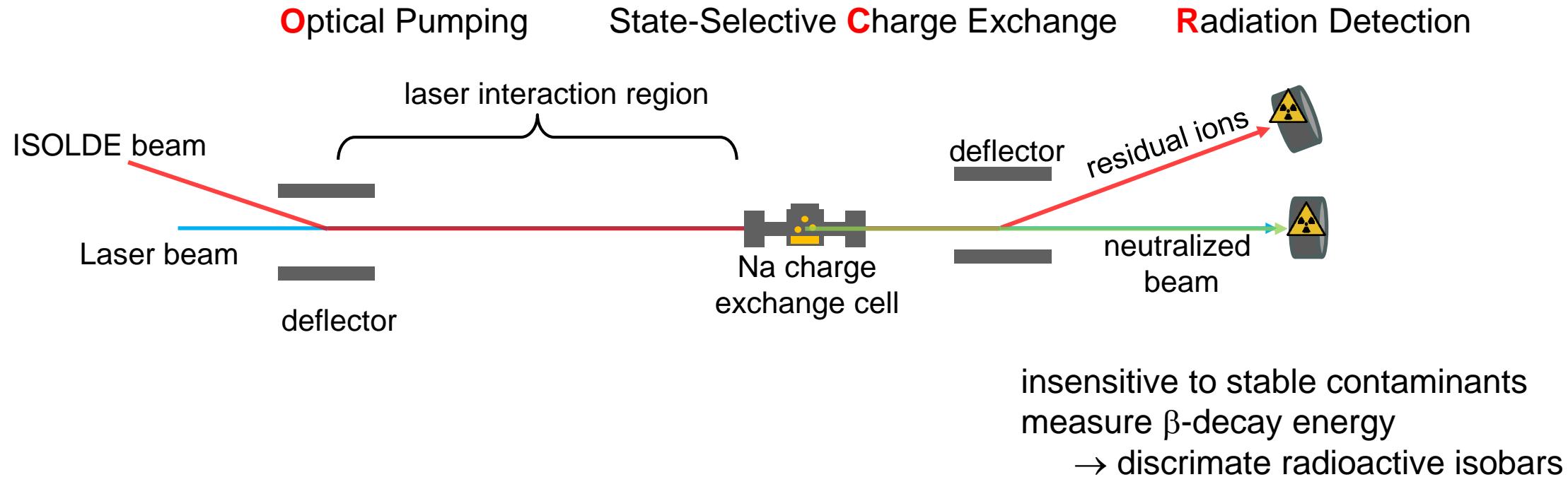
State-Selective Charge Exchange

ROC (Radiation detection after Optical pumping and state selective Charge exchange)

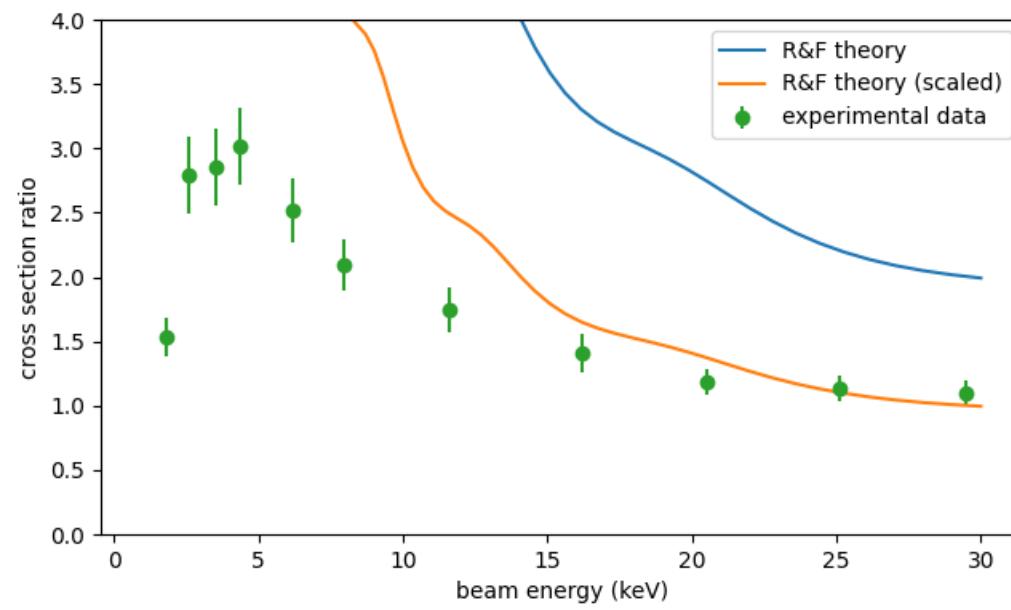
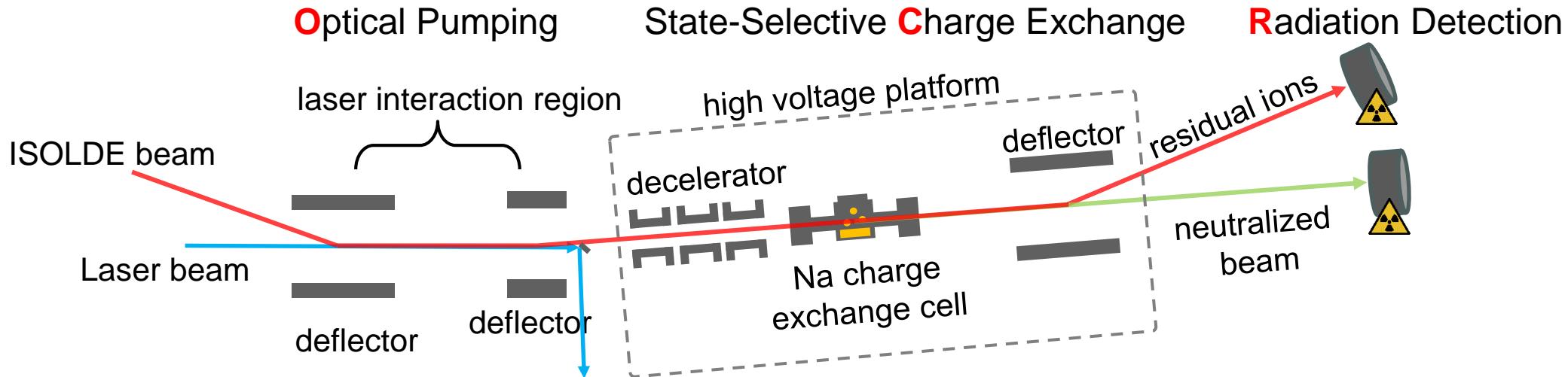


Radiation Detection

ROC – Experimental Setup



Charge exchange as a state detector



Engineering Challenge

Maximum difference in charge exchange cross section at **4 keV**
→ beam deceleration
→ CEC on HV platform
→ short HV section, neutralized beam can't be refocussed

Real setup



Installed at the **COLLAPS** setup
at ISOLDE/CERN:

Two experimental campaigns:

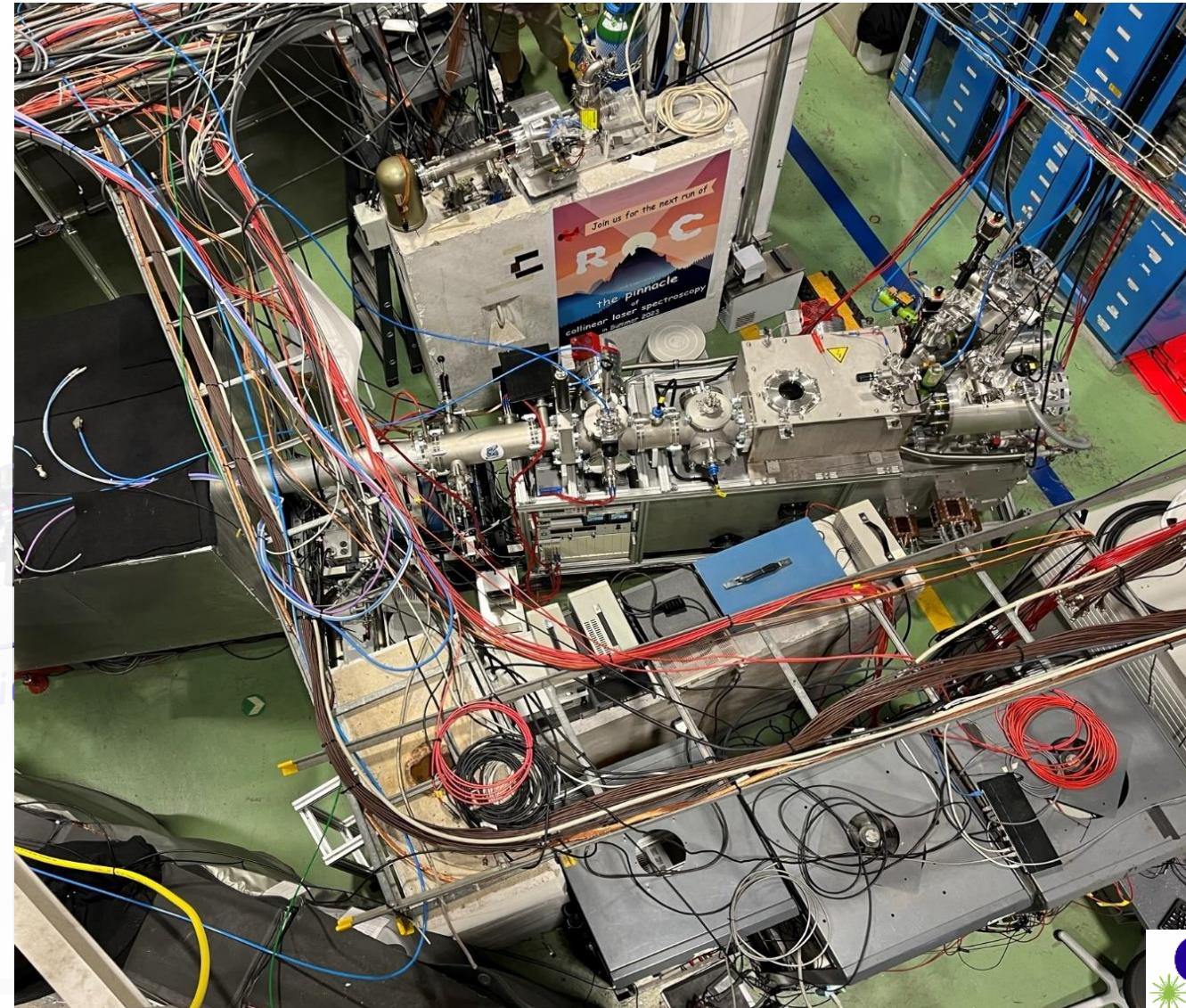
2023 → ^{53}Ca

2024 → ^{54}Ca

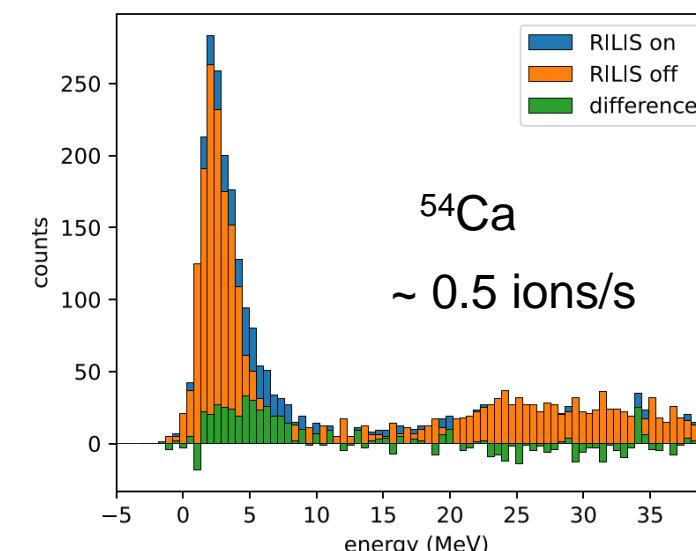
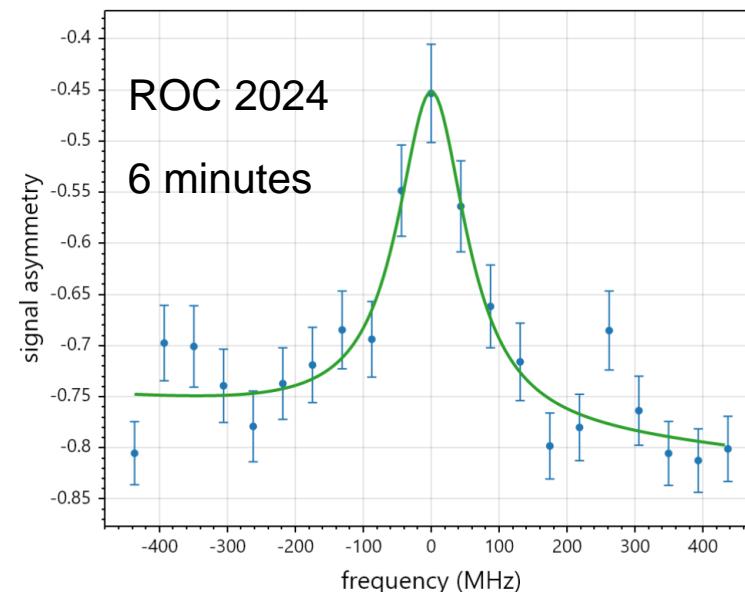
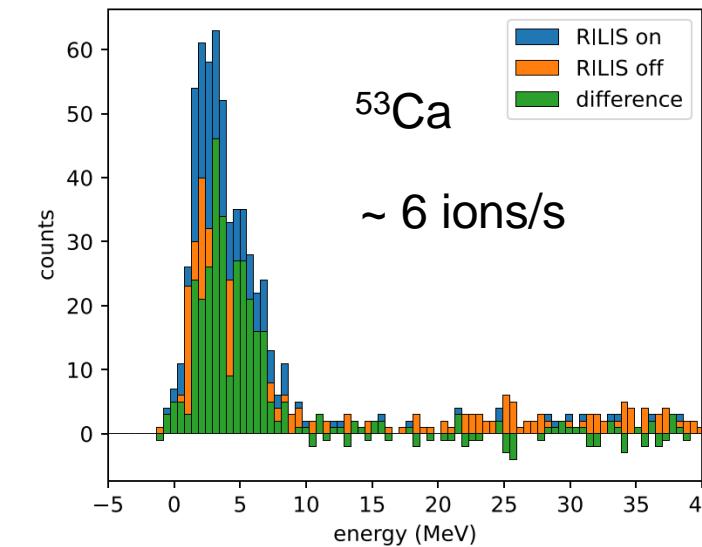
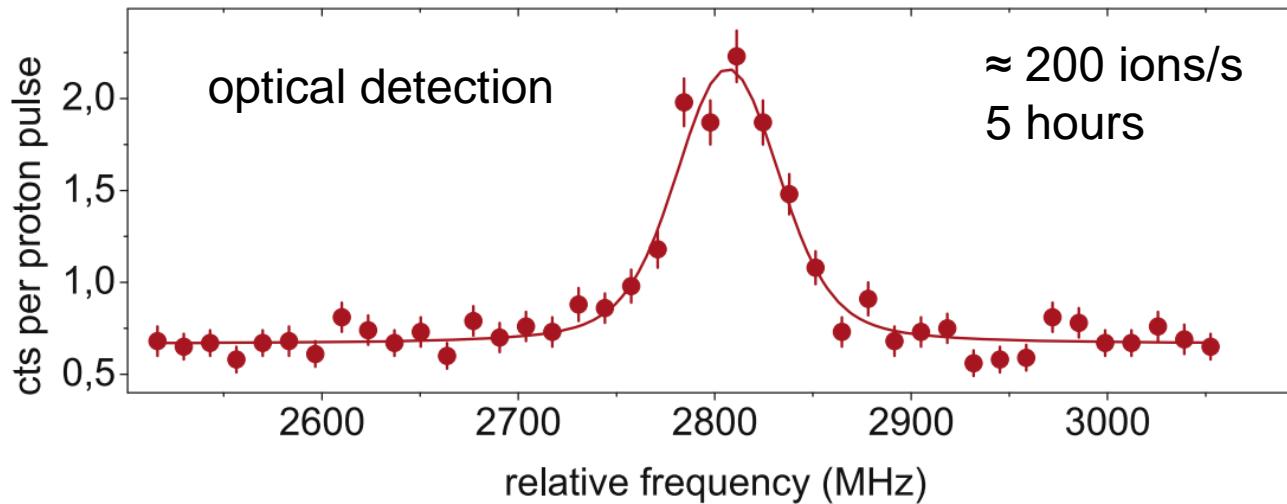
COLLAPS

laser/ion overlap

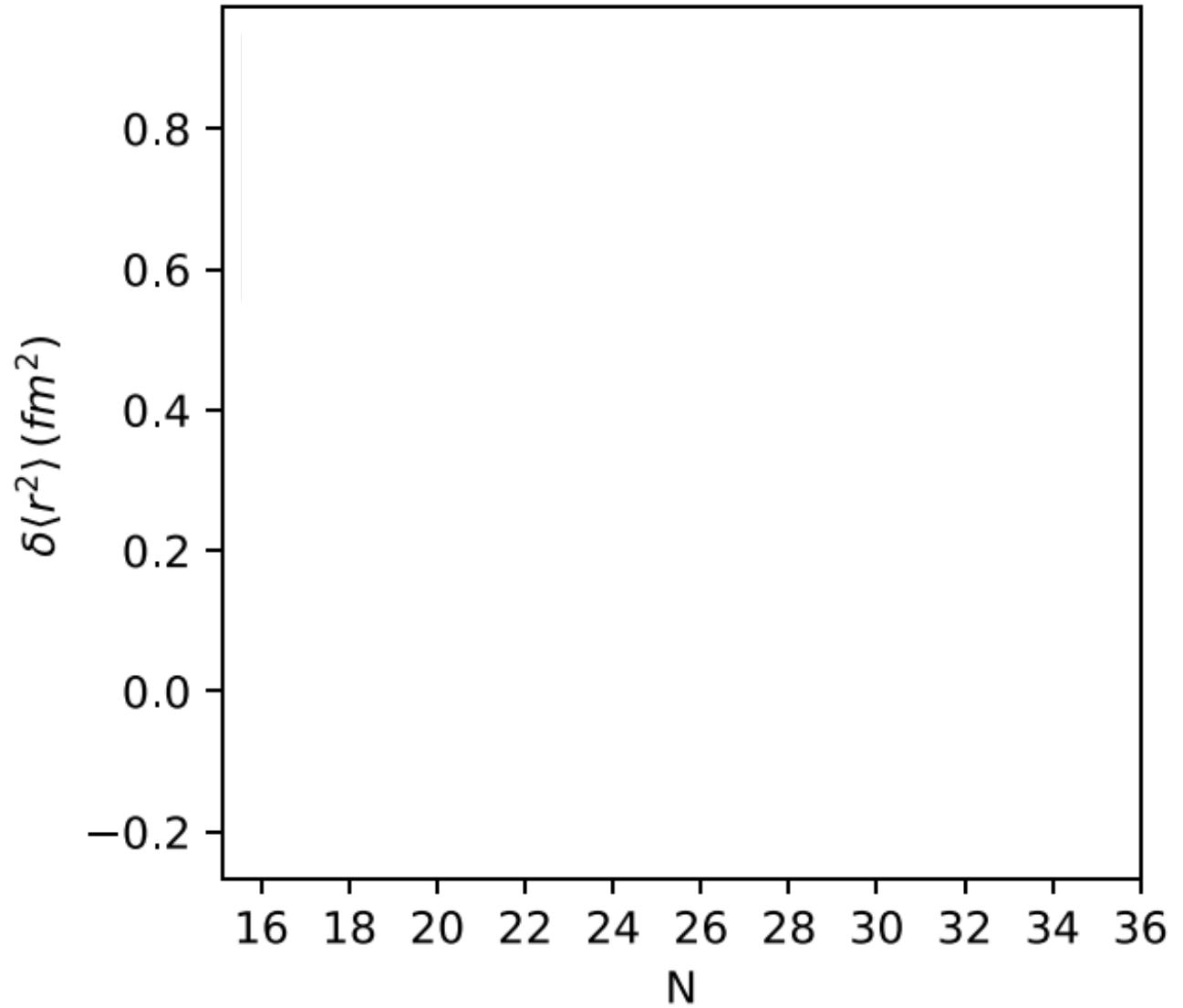
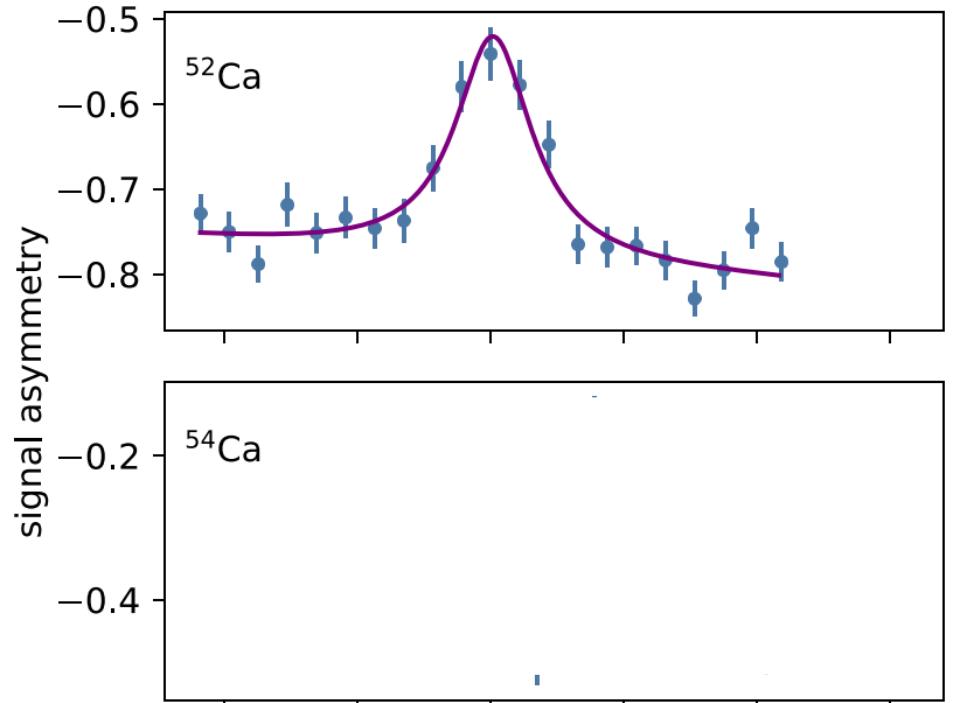
optical
focus/steering



Results: ^{52}Ca as a Sensitivity Test

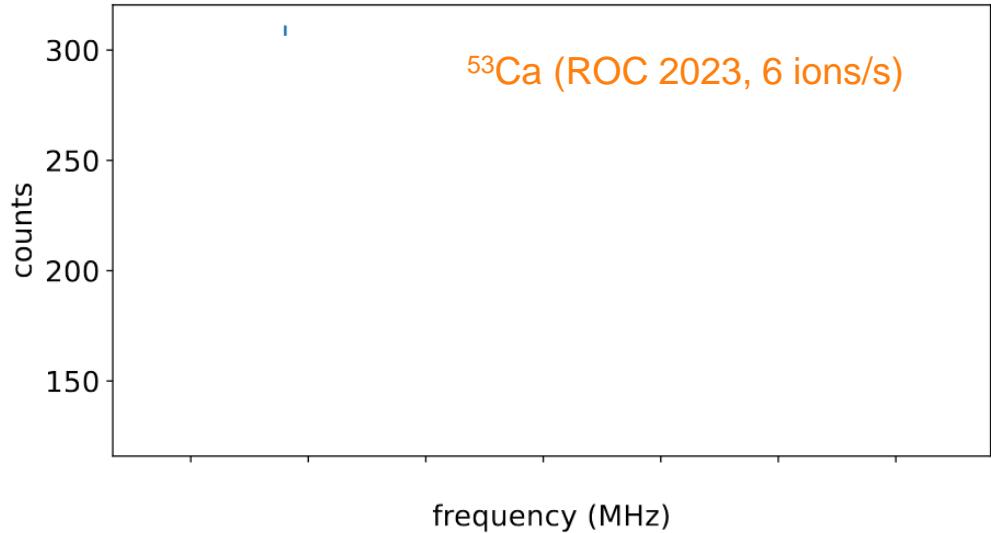


Isotope Shift and Charge Radii of $^{53,54}\text{Ca}$



$$\delta \langle r^2 \rangle^{A,40} = \frac{1}{F} \cdot \left(\delta \nu^{A,40} - K \cdot \frac{m_A - m_{40}}{m_A \cdot m_{40}} \right)$$

The Magnetic Moment of ^{53}Ca



$$\mu(^{53}\text{Ca}) = \mu(^{43}\text{Ca}) \cdot \frac{A(^{53}\text{Ca}) \cdot I(^{53}\text{Ca})}{A(^{43}\text{Ca}) \cdot I(^{43}\text{Ca})}$$

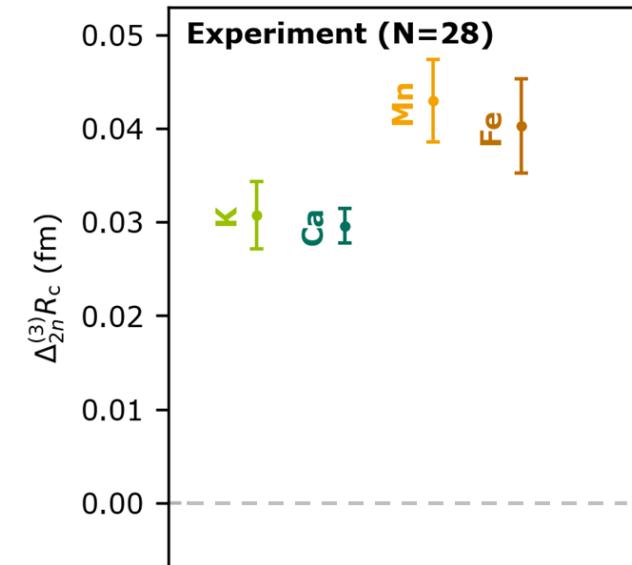
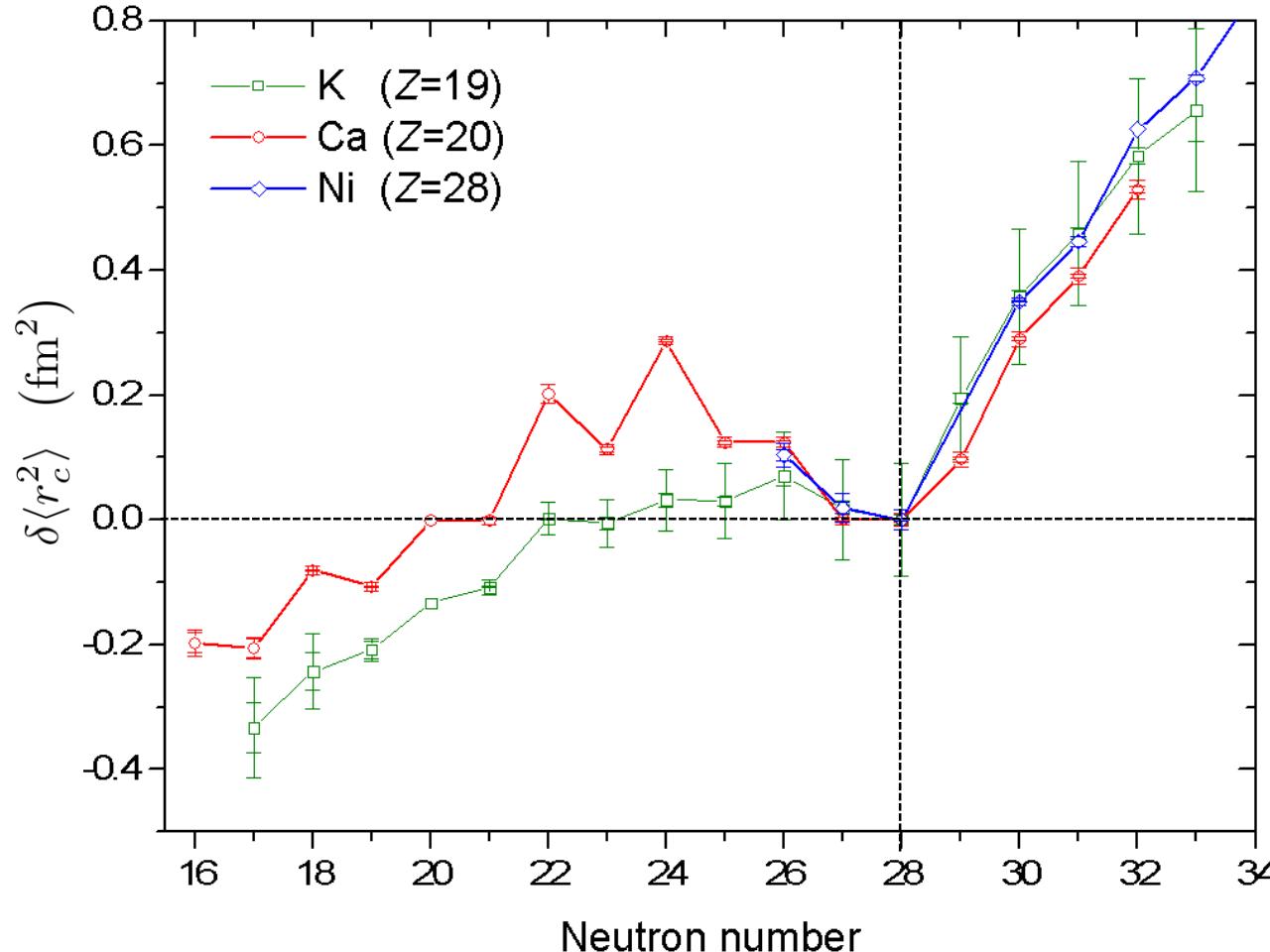
Comparison to Nickel



Three-point two-neutron difference:

$$\Delta_{2n}^{(3)} R_c(N) \equiv \frac{1}{2} [(R_c(N+2) - R_c(N)) - (R_c(N) - R_c(N-2))]$$

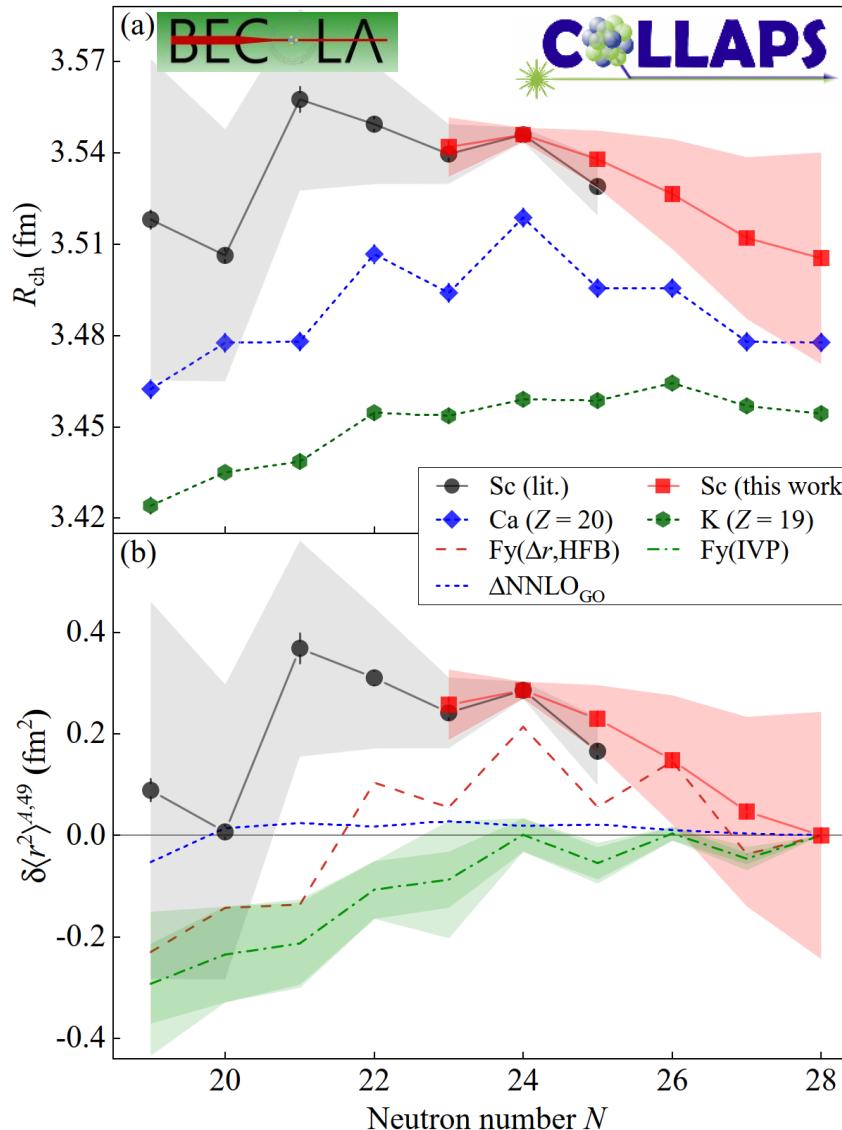
$$\Delta_{2n}^{(3)} R_c(N) \equiv \frac{1}{2} [R_c(N+2) - 2R_c(N) + R_c(N-2)]$$



F. Sommer et al., Phys. Rev. Lett. 129 132501 (2022)

Even though N=28 is only a „weak“ shell closure in Ni, the „strength“ of the kink is identical to Ca!

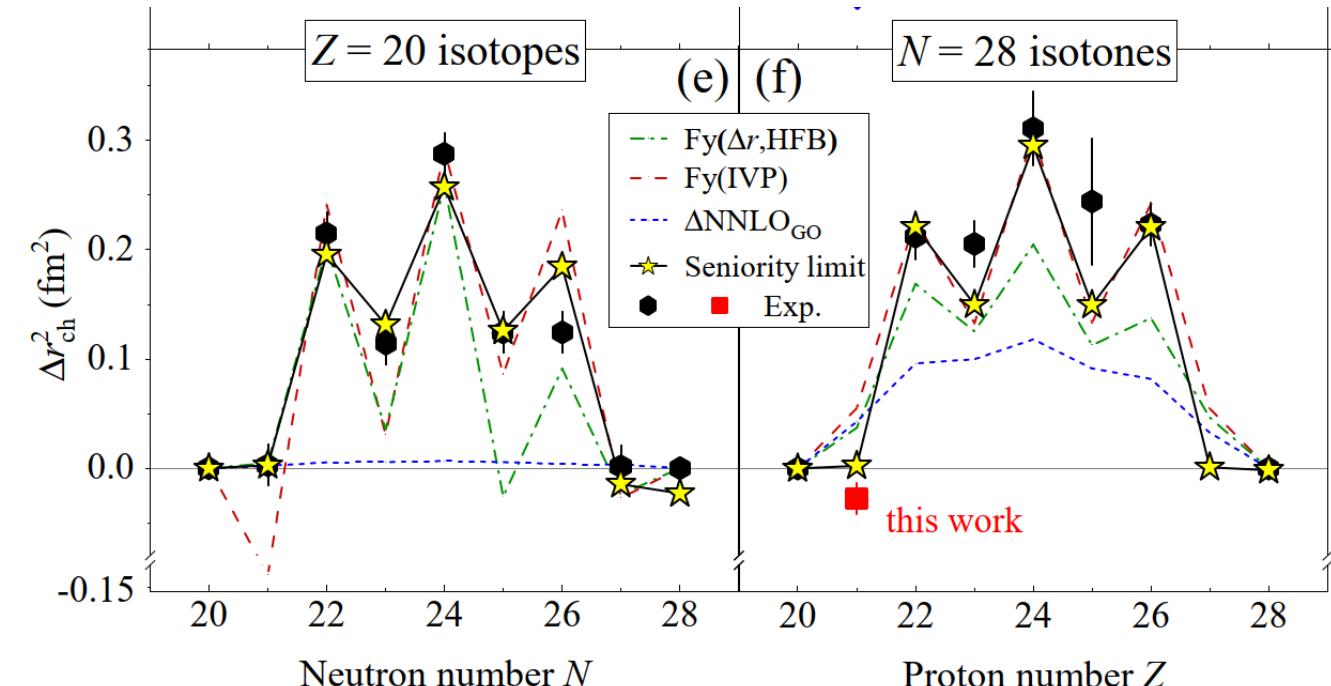
Scandium (Z=21) Results and the N=28 Isotones



In Sc ($Z=21$) a kink appears at $N = 20$.

König et al., PRL 131, 102501 (2023)

A similar pattern as in Ca appears:
 $R_c(^{49}\text{Sc}) \approx R_c(^{41}\text{Sc})$



Removing a linear trend in the charge radii along the $N=28$ isotones, gives a similar trend from $Z=20$ to $Z=28$ as in the Ca isotopes from $N=20$ to $N=28$

Summary



Laser spectroscopy is a „universal tool“ to study charge radii and nuclear moments

Charge radii changes are very sensitive to nuclear structure details and *can* herald indications for shell closures

The interplay between measurements and theory has lead to an improved understanding of nuclear structure and a continuous development of theoretical tools

We have proven the ROC technique being sensitive at the 1 atom/s level and determined the moments and charge radii of $^{53,54}\text{Ca}$

The microscopic origin of the kink and the odd-even staggering is still not fully understood.

Mi, 14:00

HK 25: Focus Session II: Accurate Nuclear Charge Radii of Light Elements

K. König: First laser spectroscopic measurements of charge radii along the carbon isotope chain

F. Wauters: Precision radii of light elements using Metallic Magnetic Calorimeters

R. Roth: Precision Radii from the No-Core Shell Model via Neural Networks

THANK YOU



TECHNISCHE
UNIVERSITÄT
DARMSTADT



KU LEUVEN



THE UNIVERSITY
of LIVERPOOL

GEFÖRDERT VOM



Bundesministerium
für Bildung
und Forschung

Local Team at ISOLDE:
Liss Vasquez Rodriguez
Tim Lellinger
Peter Plattner
Edward Matthews

