## INTERNATIONAL SCHOOL OF NUCLEAR PHYSICS, 46th COURSE, Erice

# Role of isospin asymmetry in the onset of quark matter in neutron stars

#### Pavlo Panasiuk<sup>1</sup>,

Oleksii Ivanytskyi<sup>2</sup>, Violetta Sagun<sup>3</sup>, David Blaschke<sup>2,4,5</sup>, Tim Dietrich<sup>6,7</sup>

<sup>1</sup>Centre for Physics of the University of Coimbra, <sup>2</sup>University of Wroclaw, <sup>3</sup>University of Southampton, <sup>4</sup>Helmholtz-Zentrum Dresden-Rossendorf, <sup>5</sup>CASUS, <sup>6</sup>University of Potsdam, <sup>7</sup>Max Planck Institute

17 Sep 2025





Introduction

2 Equation of state

Introduction

2 Equation of state

### The problem

- $\Rightarrow$  Terrestrial experiments on nuclear matter provide knowledge at density near  $n_{\rm sat}=0.15{\rm fm}^{-3}$ 
  - → Rough shape of energy density, pressure etc is known
  - $\rightarrow$  Certain intuition around  $n_{sat}$  is formulated
  - ightarrow In particular, we suspect that quarks only emerge at much higher density
- ⇒ It is tempting to apply the same intuition to neutron stars (unearthly conditions)
  - → But the intuition might be severely misleading

#### Aim of the talk:

- to demonstrate that the <u>quark onset density</u> within neutron stars may strongly deviate from earthly expectations
- ② to establish a relation between symmetric and asymmetric onset densities, supported by constraints

### Symmetric and asymmetric matter

The difference between neutron star matter and earthly nuclear matter lies in their typical isospin asymmetry  $I = \frac{n_I}{n_B}$ .

- $\Rightarrow$  Nuclei : symmetric matter  $n_n \approx n_p \Rightarrow I \approx 0$
- $\Rightarrow$  NSs : highly asymmetric matter  $n_n \gg n_p \Rightarrow I \gg 0$  (usually  $I \approx 0.8$ )

Isospin asymmetry induces repulsion between n, p

⇒ quark onset is easier to reach?

Introduction

2 Equation of state

### Equation of state in use

To describe onset properties we need hybrid EoS.

Hadronic EoS: DDTCY [Courtesy of S. Typel]

Relativistic mean-field description with nucleons and hyperons, with scalar attraction and vector repulsion employed. Marginally fits into flow constraint [Danielewicz et. al 2002], but does not describe  $2M_{\odot}$  NSs.

Quark EoS: 3F nonlocal NJL [O. Ivanytskyi, 2025]

Nambu-Jona-Lasinio model with nonlocal current-current interactions in scalar, vector and diquark channels, with latter controlled by corresponding dimensionless couplings  $(\eta_V, \eta_D)$ .

- $\Rightarrow$   $(\eta_V, \eta_D)$  priorly serve as free parameters
- ⇒ *uds* particle content

For neutron star matter, charge neutrality and  $\beta$ -equilibrium is imposed.

### Matching scheme 1

Maxwell construction (bridge between hadronic (HP) and quark phases (QP)) ensures mechanical and baryon chemical equilibrium on the phase boundary

$$P_H(\mu_{HP}^{ons}) = P_Q(\mu_{QP}^{ons})$$
  
 $\mu_{HP}^{ons} = \mu_{QP}^{ons}$ 

This type of matching ensures nonzero density jump  $n_{QP}^{ons}-n_{HP}^{ons}$ , which corresponds to slope difference in  $P/\mu$  plane.

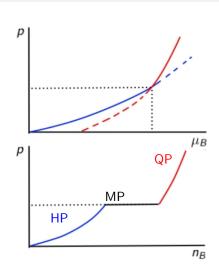


Figure: Sample Maxwell construction in  $P/\mu$  and P/n planes

### Matching scheme 2

#### Procedure:

⇒ EoS matching for a uniform grid of couplings

$$(\eta_V, \eta_D) \in [0.00, 1.00] \times [0.12, 0.5];$$

- ⇒ Seek Maxwell crossings;
- ⇒ Only keep models which have Maxwell crossings in symmetric and neutron star EoS simultaneously.

#### Limitations:

- ⇒ Only 1st order PT is considered;
- ⇒ Cases, where crossings only exist in symmetric or neutron star EoS, are discarded.

This restriction defines the  $(\eta_V, \eta_D)$  sample space.

Introduction

2 Equation of state

### Coupling space

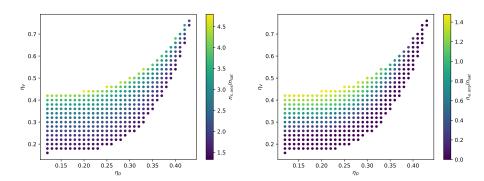


Figure:  $(\eta_V, \eta_D)$  space vs quark onset density in symmetric matter.

Figure:  $(\eta_V, \eta_D)$  space vs quark onset density in neutron star matter.

- ⇒ Asymmetric onset densities appear much lower (factor of 3?)
- $\Rightarrow$  Symmetric onset densities reach 4.5 $n_{sat}$   $\rightarrow$  available experiments not violated

### Onset space

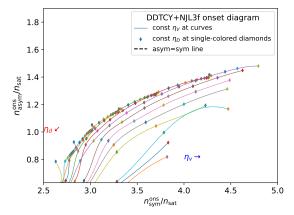
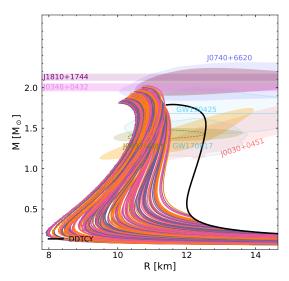


Figure:  $(\eta_V, \eta_D)$  mapping to  $(n_{asym}^{ons}, n_{sym}^{ons})$  variables.

- $\Rightarrow n_{asym} n_{sym}$  relation is evidently strongly capped
  - → Support for "early" deconfinement in neutron stars.
- $\Rightarrow$  Below  $1n_{sat}$  is reachable in neutron stars, while maintaining  $\sim 4n_{sat}$  in symmetric matter
  - n<sub>sat</sub> does not carry the same intuition in neutron stars

### Astrophysical sector



We run the EoS sample space through TOV equations to compare against astro constraints. Evidently, elimination with Bayesian analysis is possible. We apply

- ⇒ NICER constraints, J0740+6620 in particular;
- ⇒ GW170817 tidal deformability.

At the same time, flow constraint is marginally maintained for symmetric EoS.

Figure: TOV solutions for DDTCY-NJL3f.

### Credible couplings

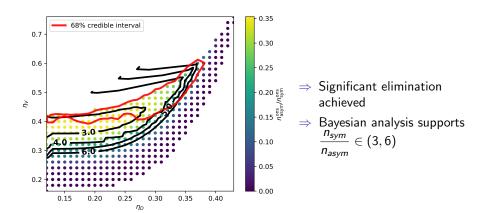


Figure: Coupling diagram with  $1\sigma$  credible region (credits to A. Ayriyan). Levels indicate onset ratio.

#### Conclusions

- ⇒ The quark onset density of electrically neutral matter at beta equilibrium may exhibit significantly lower than in the symmetric case.
- ⇒ Bayesian analysis suggests relation between symmetric and asymmetric onset densities starting from 3, which elucidates the distinction between heavy ion and neutron star regimes.
- $\Rightarrow$  Notably, the <u>analysis does not disfavor</u> asymmetric onset densities <u>below 1</u>  $n_{sat}$ .

### Backup : Flow constraint

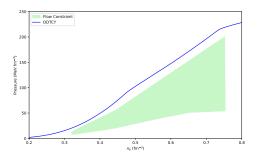


Figure: Symmetric DDTCY vs flow constraint band.

### Backup: Onset constraint

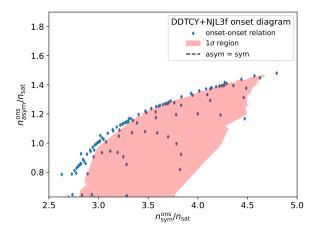


Figure: Onset space with 68% CL from BA.