

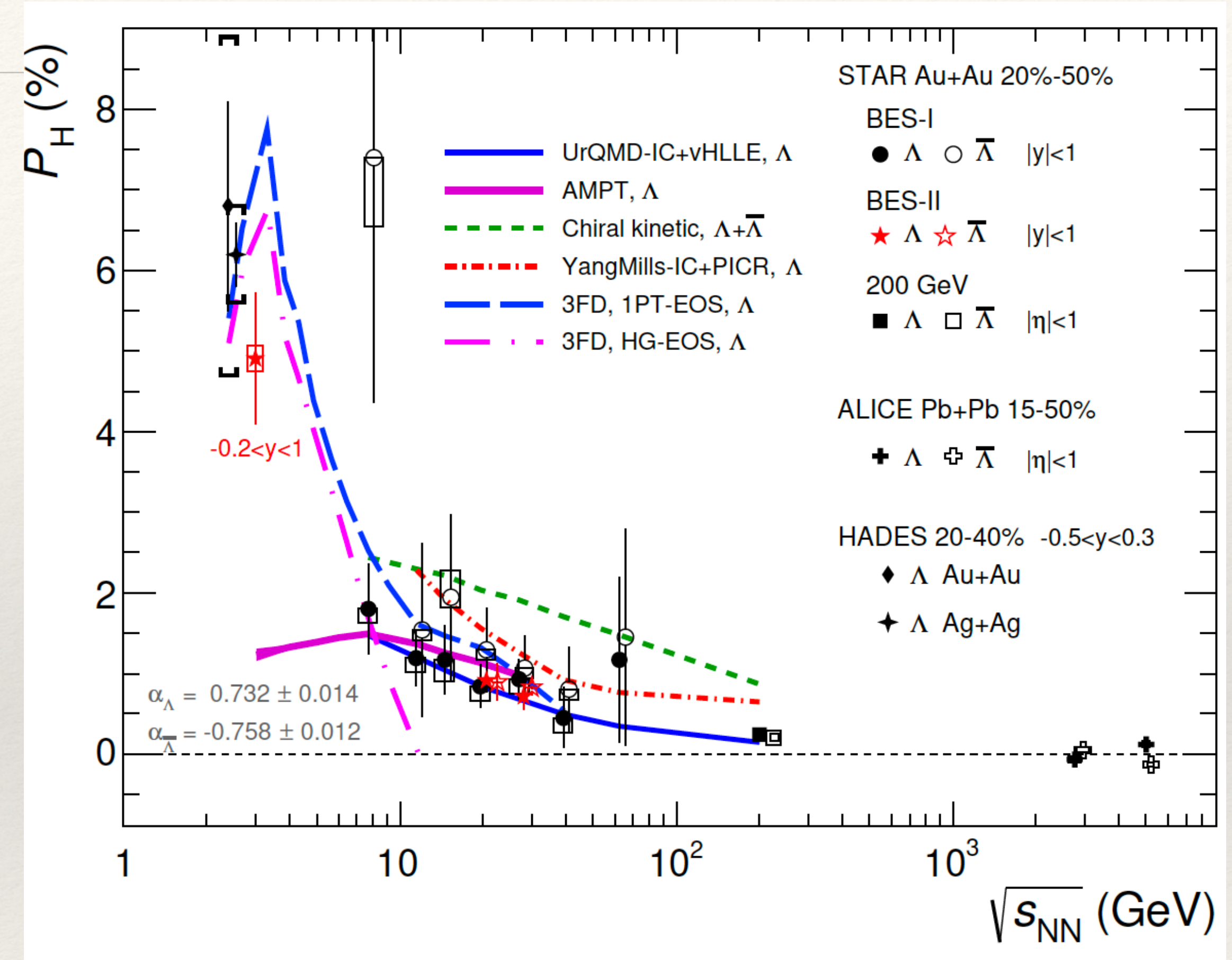
Application range of perfect spin hydrodynamics

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In collaboration with Z. Drogosz, W. Florkowski, S. Bhadury, S. Kar
talk based on: [2506.01537](#) (Phys. Rev. D. 112 2025; [2505.02657](#))

Spin hydrodynamics

- ◆ Relativistic hydro → data interpretation
- ◆ ‘Spin’ measurements → hydrodynamics + spin
- ◆ Various theoretical methods: density operators, entropy principle, grad of hydrodynamic fields, Lagrangian effective theories ...
(see works of V. Ambrus, F. Becattini, W. Florkowski, D. Rischke, R. Ryblewski, E. Speranza, ...)
- ◆ Our formalisms: Kinetic theory & Wigner function
- ◆ One of the issues: applicability range.



Kinetic theory — classical spin description

Local equilibrium distribution in spin-extended phase space

$$f^{\pm}(x, p, s) = \exp \left(-p_{\mu} \beta^{\mu} \pm \frac{\mu}{T} + \frac{1}{2} \omega_{\mu\nu} s^{\mu\nu} \right), \quad \beta^{\mu} = \frac{u^{\mu}}{T}, \quad u'^{\mu} = (1, \vec{0})$$

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$$\omega_{\mu\nu} = \frac{\Omega_{\mu\nu}}{T} \text{ spin chemical potential}$$

Spin polarization tensor

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$$\omega_{\mu\nu} = \begin{bmatrix} 0 & e^1 & e^2 & e^3 \\ -e^1 & 0 & -b^3 & b^2 \\ -e^2 & b^3 & 0 & -b^1 \\ -e^3 & -b^2 & b^1 & 0 \end{bmatrix}$$

Electric- and magneticlike vectors:

$$\vec{e} = (e^1, e^2, e^3), \quad \vec{b} = (b^1, b^2, b^3)$$

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spin vector

$s \perp p$

$$\omega_{\mu\nu} = \frac{\Omega_{\mu\nu}}{T}$$

spin chemical potential

$$s^{\mu\nu} = \frac{1}{m} \epsilon^{\mu\nu\alpha\beta} p_{\alpha} s_{\beta}$$

$$s_{\beta} = (0, \vec{\mathfrak{s}})$$

Spin polarization tensor

internal angular momentum

[Mathisson, Acta Phys. Polon. 6, 163 (1937)]

$$\mathfrak{s}^2 = 3/4$$

eigenvalue of the Casimir operator in SU(2)

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Physical quantities obtained from

$$\int dP \int dS f^{\pm}(x, p, s) \dots$$

Perfect spin hydro: spin tensor conserved — no spin-orbit coupling, can be added by dissipation

[Florkowski, Hontarenko, PRL 134 '25; Drogoz, Florkowski, Hontarenko PRD 110 '25]

Scalar density in kinetic theory

classical spin

Tensors can be obtained by taking derivatives w.r.t. β^μ and $\omega_{\mu\nu}$

$$n = \int dP \int dS \exp \left(-p_\mu \beta^\mu + \frac{1}{2} \omega_{\mu\nu} s^{\mu\nu} \right) = \int dP \exp \left(-p_\mu \beta^\mu \right) \int dS \exp \left(\frac{1}{2} \omega_{\mu\nu} s^{\mu\nu} \right)$$

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$$\int dS \exp \left(\frac{1}{2} \omega_{\mu\nu} s^{\mu\nu} \right) = \frac{2 \sinh(\mathfrak{s} |\vec{b}_*|)}{\mathfrak{s} |\vec{b}_*|} \quad \text{with} \quad \vec{b}_* = \frac{1}{m} \left(E_p \vec{b} - \vec{p} \times \vec{e} - \frac{\vec{p} \cdot \vec{b}}{E_p + m} \vec{p} \right)$$

Spin measure

* particle rest frame

$$\mathfrak{s}^2 = 3/4$$

eigenvalue of the
Casimir operator in SU(2)

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* particle rest frame

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eigenvalue of the
Casimir operator in SU(2)

$$n = 2 \int dP' \exp \left(-\frac{E_{p'}}{T} \right) \frac{\sinh(\mathfrak{s} |\vec{b}'_*|)}{\mathfrak{s} |\vec{b}'_*|}$$

local rest frame

$$u'^\mu = (1, \vec{0})$$

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classical spin

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divergence possible!

Criterion of convergence

$$n = 2 \int dP' \exp \left(-\frac{E_{p'}}{T} \right) \frac{\sinh(\mathfrak{z} | \vec{b}'_* |)}{\mathfrak{z} | \vec{b}'_* | }$$

* particle rest frame

' local (fluid's) rest frame

$$\vec{b}' = \gamma (\vec{b} - \vec{v} \times \vec{e}) - \frac{\gamma^2}{\gamma + 1} \vec{v} (\vec{v} \cdot \vec{b})$$

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$$\tau \equiv -\frac{E_{p'}}{T} + \mathfrak{z} |\vec{b}'_*|, \quad \max(\tau) < 0$$

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For large $|\vec{p}'|$:

$$\tau = -\frac{|\vec{p}'|}{T} + \mathfrak{z} \frac{|\vec{p}'|}{m} |\vec{b}' - \hat{p}' \times \vec{e}' - (\hat{p}' \cdot \vec{b}) \vec{p}|$$

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$$\mathfrak{s}^2 = 3/4$$

$$\mathfrak{s} \sqrt{\vec{b}'^2 + \vec{e}'^2 + 2 |\vec{e}' \times \vec{b}'|} < \frac{m}{T}$$

$\vec{b}, \vec{e}$ — components of $\omega^{\mu\nu}$
spin polarization tensor

Criterion of convergence

$$n = 2 \int dP' \exp \left(-\frac{E_{p'}}{T} \right) \frac{\sinh(\xi |\vec{b}'_*|)}{\xi |\vec{b}'_*|}$$

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$$\xi^2 = 3/4$$

$$\xi \sqrt{\vec{b}'^2 + \vec{e}'^2 + 2 |\vec{e}' \times \vec{b}'|} < \frac{m}{T}$$

$\vec{b}, \vec{e}$ — components of $\omega^{\mu\nu}$
spin polarization tensor

see also the result of Abboud, Gavassino, Singh, Speranza, arXiv:2506.19786

Local equilibrium Wigner function

NEW

$$W^{\pm}(x, k) = \pm \frac{1}{2} \int dP \delta^{(4)}(k \mp p) (p_{\mu} \gamma^{\mu} \pm m) X^{\pm}(x, p)$$

$$X^{\pm}(x, p) = \exp \left[\pm \xi(x) - \beta_{\mu}(x) p^{\mu} + \gamma_5 a_{\mu} \gamma^{\mu} \right], \quad a_{\mu}(x, p) = -\frac{1}{4m} \epsilon^{\mu\nu\alpha\beta} \omega_{\alpha\beta}(x) p^{\nu}$$

Differs from widely used version by Becattini, Chandra, Del Zanna, Grossi, *Annals Phys.* 338 '13

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♣ Accurate definition of the mean polarization vector.

◆ Spin density matrix consistent with Landau-Lifshitz.

✱ Conserved currents agree with kinetic theory at any order in ω up to numerical factor. [Drogosz, arXiv:2509.0601]

★ Perfect spin hydro is non-linearly causal and its equations are symmetric hyperbolic.

[S.Bhadury, Z.Drogosz, W.Florkowski, S.Kar, V.M. [2505.02657](#)]

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Thermodynamics obtained from $\int dP \operatorname{tr}_4 [\cdots X^{\pm}(x, p) \cdots]$
traces over the spinor space

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Applicability range?

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$$\omega_{\alpha\beta} = \begin{bmatrix} 0 & e^1 & e^2 & e^3 \\ -e^1 & 0 & -b^3 & b^2 \\ -e^2 & b^3 & 0 & -b^1 \\ -e^3 & -b^2 & b^1 & 0 \end{bmatrix}$$

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$$-\omega_{\max} \leq e^i, b^i \leq \omega_{\max}$$

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Applicability range?

$$\omega_{\max} \sqrt{\frac{1+v}{1-v}} < \frac{2}{2+\sqrt{2}} \frac{m}{T}$$

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Applicability range?

longitudinal relativistic
Doppler effect

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Local equilibrium Wigner function

Criteria of convergence

$$\frac{1}{2} \sqrt{\vec{b}'^2 + \vec{e}'^2 + 2|\vec{e}' \times \vec{b}'|} < \frac{m}{T}$$

$$\vec{b}' = \gamma(\vec{b} - \vec{v} \times \vec{e}) - \frac{\gamma^2}{\gamma + 1} \vec{v}(\vec{v} \cdot \vec{b})$$

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More restrictive: $(\vec{b}, \vec{v}, \vec{e}) \perp$ to each other,
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 $|\vec{b}| = |\vec{e}|, (\vec{v} \times \vec{e}) \parallel \vec{b}$

but easier to apply: need $\omega_{\max}, v$ only

vs. $\vec{b} = (b_1, b_2, b_3), \vec{e} = (e_1, e_2, e_3), \vec{v} = (v_1, v_2, v_3)$

Variety of the application ranges

Variables	Kinetic theory – classical spin	Wigner function – quantum spin
$\omega_{\max}, v$	$\mathfrak{S} \omega_{\max} \left(2 + \sqrt{2}\right) \sqrt{\frac{1+v}{1-v}} < \frac{m}{T}$	$\frac{1}{2} \omega_{\max} \left(2 + \sqrt{2}\right) \sqrt{\frac{1+v}{1-v}} < \frac{m}{T}$
$\vec{e}, \vec{b}, v$	$\mathfrak{S} \sqrt{\vec{b}^2 + \vec{e}^2 + 2 \vec{e} \times \vec{b} } \sqrt{\frac{1+v}{1-v}} < \frac{m}{T}$	$\frac{1}{2} \sqrt{\vec{b}^2 + \vec{e}^2 + 2 \vec{e} \times \vec{b} } \sqrt{\frac{1+v}{1-v}} < \frac{m}{T}$
$\vec{e}, \vec{b}, \vec{v}$	$\mathfrak{S} \sqrt{\vec{b}'^2 + \vec{e}'^2 + 2 \vec{e}' \times \vec{b}' } < \frac{m}{T}$	$\frac{1}{2} \sqrt{\vec{b}'^2 + \vec{e}'^2 + 2 \vec{e}' \times \vec{b}' } < \frac{m}{T}$

in local rest frame

The bound depends on the „level” of given variables

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$\vec{e}, \vec{b}, \vec{v}$	$\mathfrak{s} \sqrt{\vec{b}'^2 + \vec{e}'^2 + 2 \vec{e}' \times \vec{b}' } < \frac{m}{T}$	$\frac{1}{2} \sqrt{\vec{b}'^2 + \vec{e}'^2 + 2 \vec{e}' \times \vec{b}' } < \frac{m}{T}$

in local rest frame

The bound depends on the „level” of given variables

$$\mathfrak{s}^2 = 3/4$$

Practical applicability

Perfect spin hydro with conserved spin tensor $\omega_{\mu\nu} \rightarrow$ last stages of HIC / HIC at moderate energies.

$$\frac{1}{2} \omega_{\max} \left(2 + \sqrt{2} \right) \sqrt{\frac{1+v}{1-v}} < \frac{m}{T}$$

$$\omega_{\max} < 0.2, \quad m = 0.4 - 1.1 \text{ GeV}, \quad T = 0.15 \text{ GeV}, \quad v = 0.5$$

$$m/T \approx 3-10, \quad \sqrt{(1+v)/(1-v)} \approx 1.7$$

Values of ω cannot be arbitrary — like $\mu < m$ for bosons

Abboud, Gavassino, Singh, Speranza, arXiv:2506.19786

Summary

- ❖ Proper formulation of the local equilibrium Wigner function for particles with spin $1/2$ — **long-standing problem solved** — principal tool in spin hydrodynamics.
- ❖ Criteria constraining the application of perfect spin hydro have been defined.
- ❖ Our perfect spin hydro is non-linearly causal and its equations are symmetric hyperbolic.