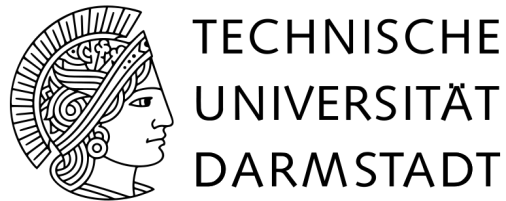


# *Ab initio* calculations as benchmarks for nuclear DFT



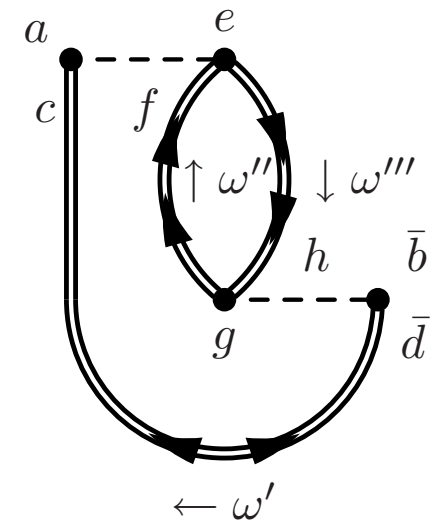
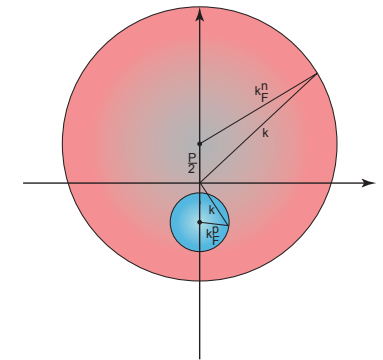
**Vittorio Somà** (TU Darmstadt & EMMI)

## *Nuclear matter:*

- Drischler, Somà, Schwenk, *in preparation*

## *Nuclei:*

- Somà, Duguet, Barbieri, PRC 84 064317 (2011)
- Somà, Barbieri, Duguet, PRC 87 011303(R) (2013)
- Barbieri, Cipollone, Somà, Duguet, Navrátil, arXiv:1211.3315
- Somà, Barbieri, Duguet, arXiv:1304.xxxx
- Somà, Cipollone, Barbieri, Duguet, Navrátil, *in preparation*

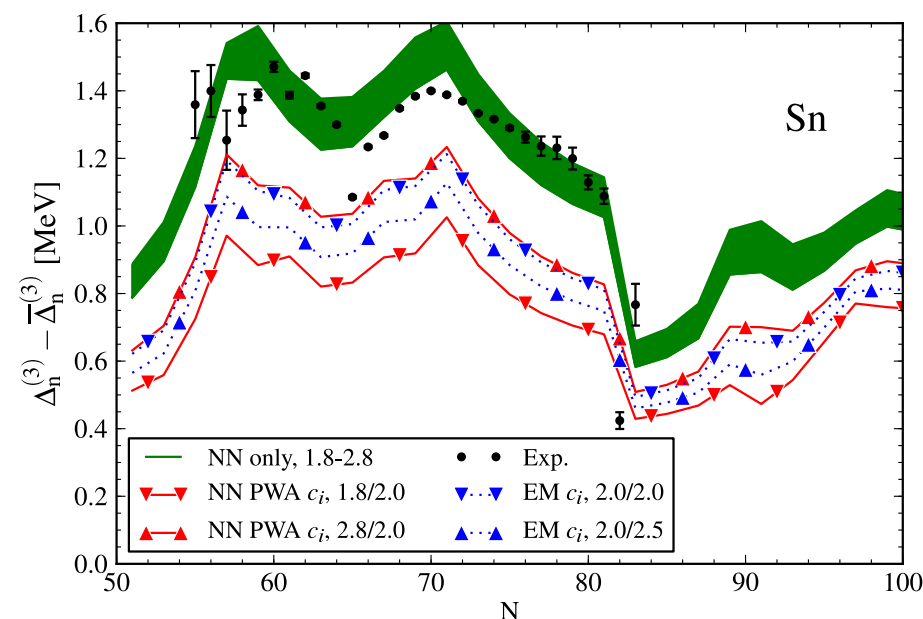


- ✱ Connection with underlying inter-nucleon interactions
  - ⇒ Microscopic perspective on new phenomena
  
- ✱ Predictive power
  - ⇒ Crucial for exotic systems
  - ⇒ Theoretical error estimates possible (and mandatory)
  
- ✱ Towards a more consistent description of structure and reactions
  - ⇒ Crucial for interpretation of new experiments
  
- ✱ Provide benchmarks for other approaches
  - ⇒ Important in the extension of traditional models to the neutron-rich sector

## ✱ Towards *ab initio* EDFs

➡ Density matrix expansion [Gebremariam, Duguet, Bogner 2010]

➡ Low-momentum interactions in the pairing channel [Lesinski *et al.* 2011]



## ✱ Guidance for extensions and improvements of EDFs

➡ UNEDF

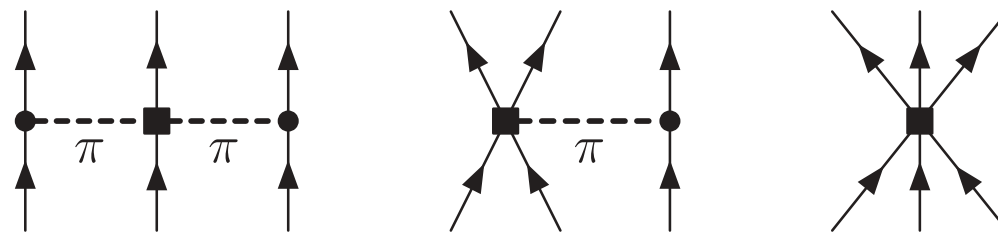
## ✱ *Ab initio* method should share the same features of the EDF

➡ E.g. symmetry breaking (and restoration)

# Advances in the modeling of nuclear forces

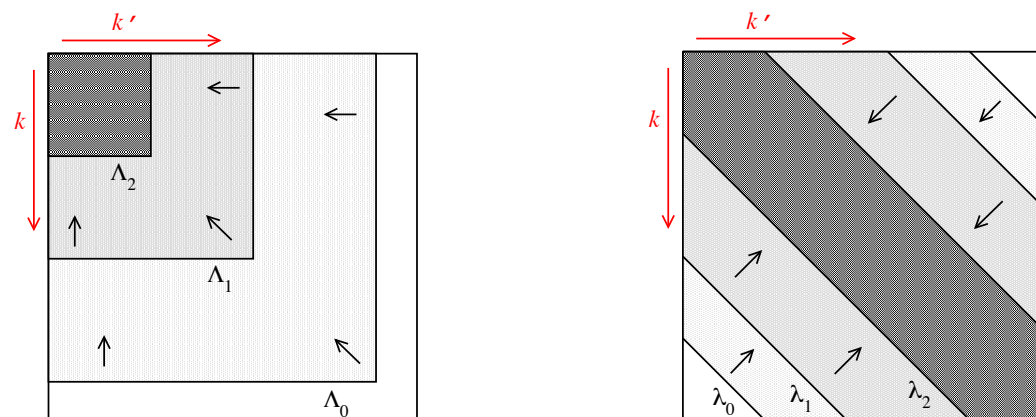
## ✱ Nuclear interactions from chiral EFT

- ⇒ Consistent many-body forces
- ⇒ A way to quantify theoretical errors



## ✱ Renormalization group techniques for NN and 3N forces

- ⇒ Many-body problem more perturbative





✱ Crucial for benchmarking and extrapolations to **neutron-rich sector**

	Asymmetric matter	Medium-mass nuclei
1) Interaction	✓	✗
2) Many-body expansion	✗	✗
3) Model space truncation	✓	✓
4) Numerical algorithms	✓	✓

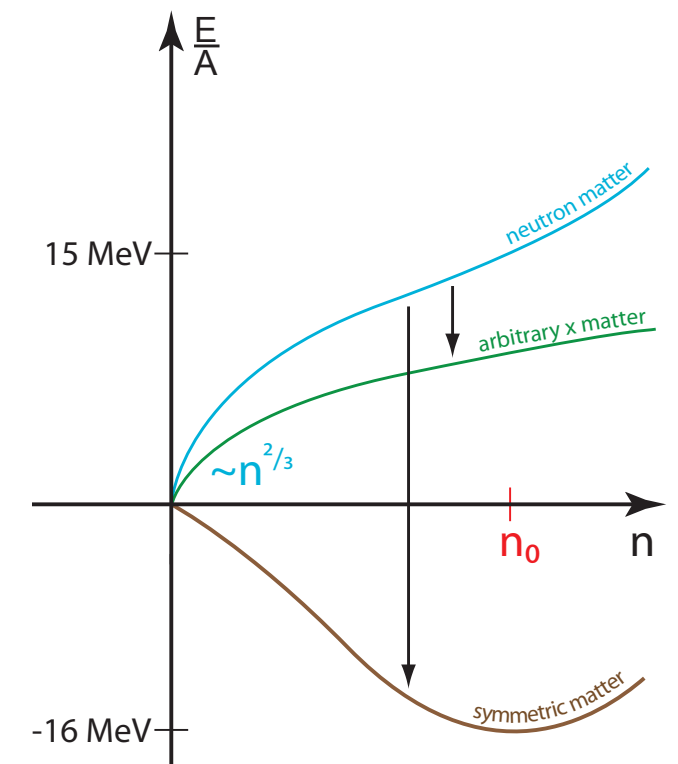
# Asymmetric nuclear matter

- ✱ Applications in astrophysical environments

- ➡ Compact stars
- ➡ Core-collapse supernovae

- ✱ Very few calculations with explicit isospin asymmetry

- ➡ **AFDMC** [Fantoni *et al.* 2008]
  - NN only
- ➡ **SCGF** [Frick *et al.* 2005]
  - NN only, finite temperature
- ➡ **BHF** [Bombaci and Lombardo 1991, Zuo *et al.* 1999, Zuo 2012]



# Asymmetric nuclear matter

## ✱ Applications in astrophysical environments

- ➡ Compact stars
- ➡ Core-collapse supernovae

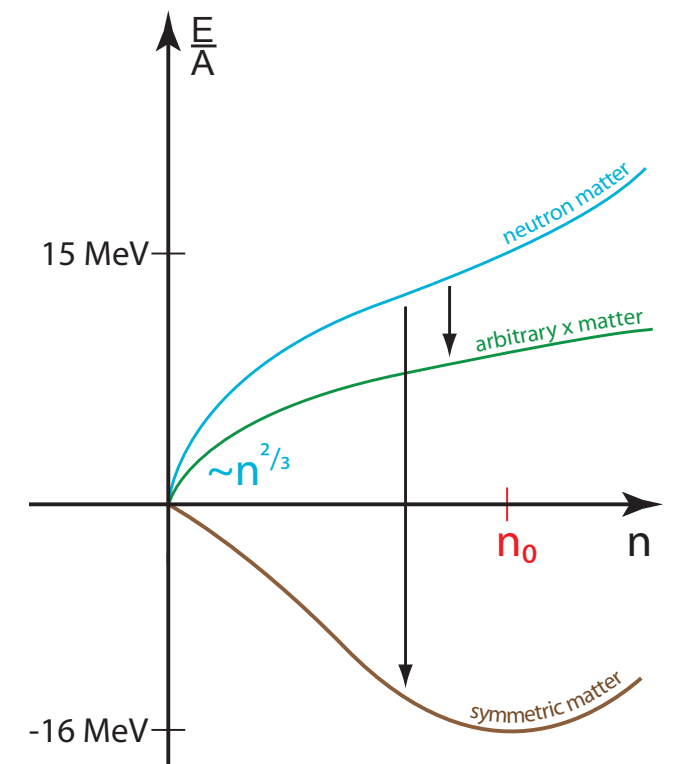
## ✱ Very few calculations with explicit isospin asymmetry

- ➡ **AFDMC** [Fantoni *et al.* 2008]
  - NN only
- ➡ **SCGF** [Frick *et al.* 2005]
  - NN only, finite temperature
- ➡ **BHF** [Bombaci and Lombardo 1991, Zuo *et al.* 1999, Zuo 2012]

## ✱ Parabolic approximation

$$\frac{E(n, \beta)}{A} = \frac{E(n, \beta = 0)}{A} + \beta^2 E_{sym}(n) + \mathcal{O}(\beta^4)$$

where  $\beta \equiv \frac{n_n - n_p}{n_n + n_p}$



## ✱ MBPT(2)

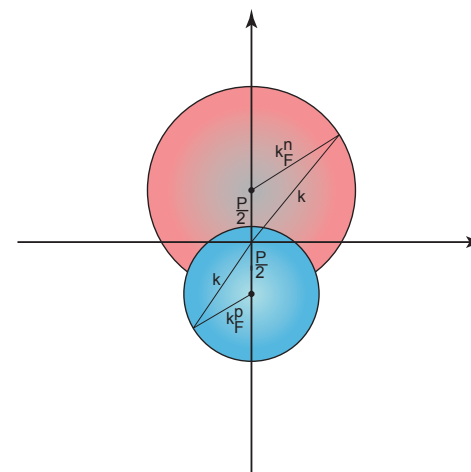
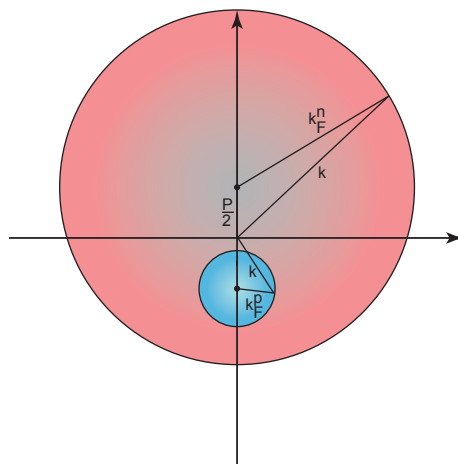
⇒ chiral EFT N<sup>3</sup>LO NN and N<sup>2</sup>LO 3N interactions

⇒ NN evolved to low-momentum (see Gezerlis and Hebeler)

⇒ 3N forces at the HF level

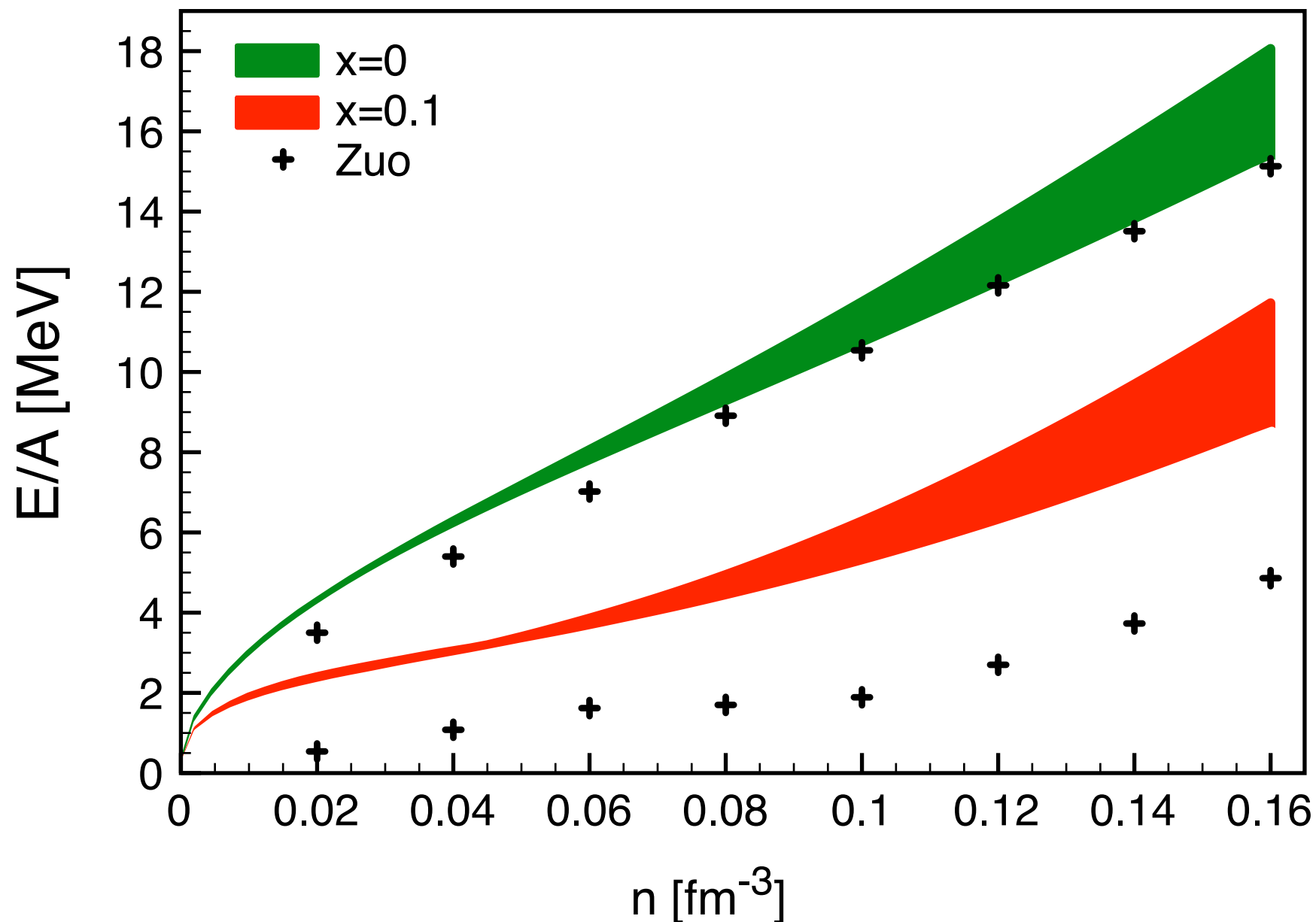
⇒ Small  $x$  limit  $x \equiv \frac{n_p}{n_n + n_p} \longrightarrow V_{3N}|_x \approx V|_{nnn} + 3V|_{nnp}$

## ✱ Non-trivial integration over the two different Fermi spheres



[Drischler 2012]

# Energy of asymmetric matter



✱ NN interaction: chiral  $N^3\text{LO}$  evolved to low momentum [Entem and Machleidt 2003]

✱ Uncertainties: 7 sets of low- $k$  ( $1.8\text{-}2.8\text{ fm}^{-1}$ ) & 3N cutoffs

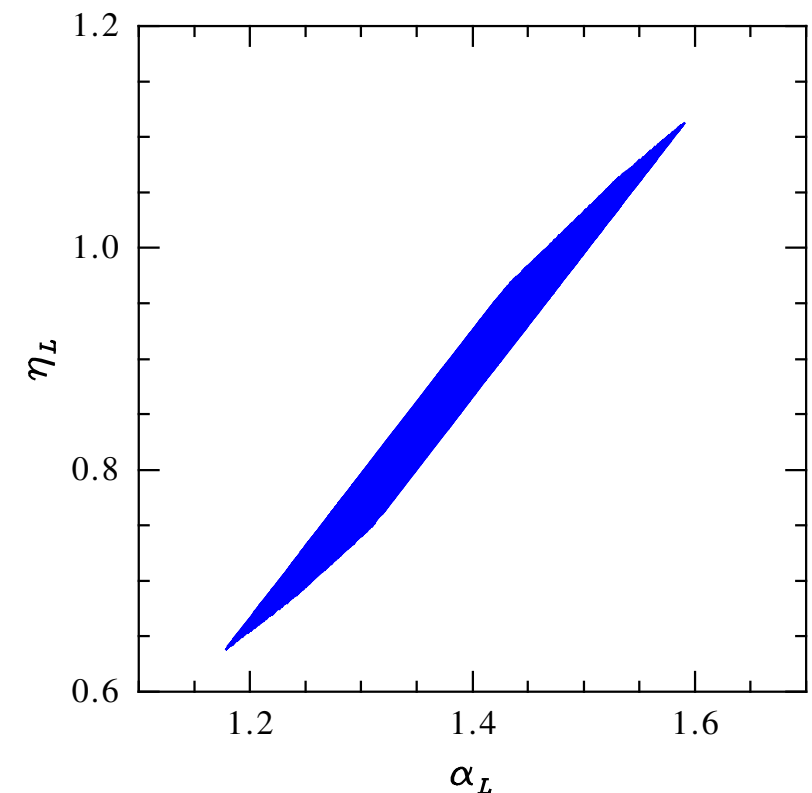
# Quasi-parabolic expansion

✱ Expression for the energy per particle in asymmetric matter

⇒ Kinetic energy + interpolation of interaction energy

$$\frac{E(\bar{n}, x)}{A} = T_0 \left[ \frac{3}{5} \left( x^{\frac{5}{3}} + (1-x)^{\frac{5}{3}} \right) (2\bar{n})^{\frac{2}{3}} \right. \\ \left. - ((2\alpha - 4\alpha_L) x (1-x) + \alpha_L) \bar{n} \right. \\ \left. + ((2\eta - 4\eta_L) x (1-x) + \eta_L) \bar{n}^\gamma \right]$$

[Hebeler, Lattimer, Pethick, Schwenk 2013]



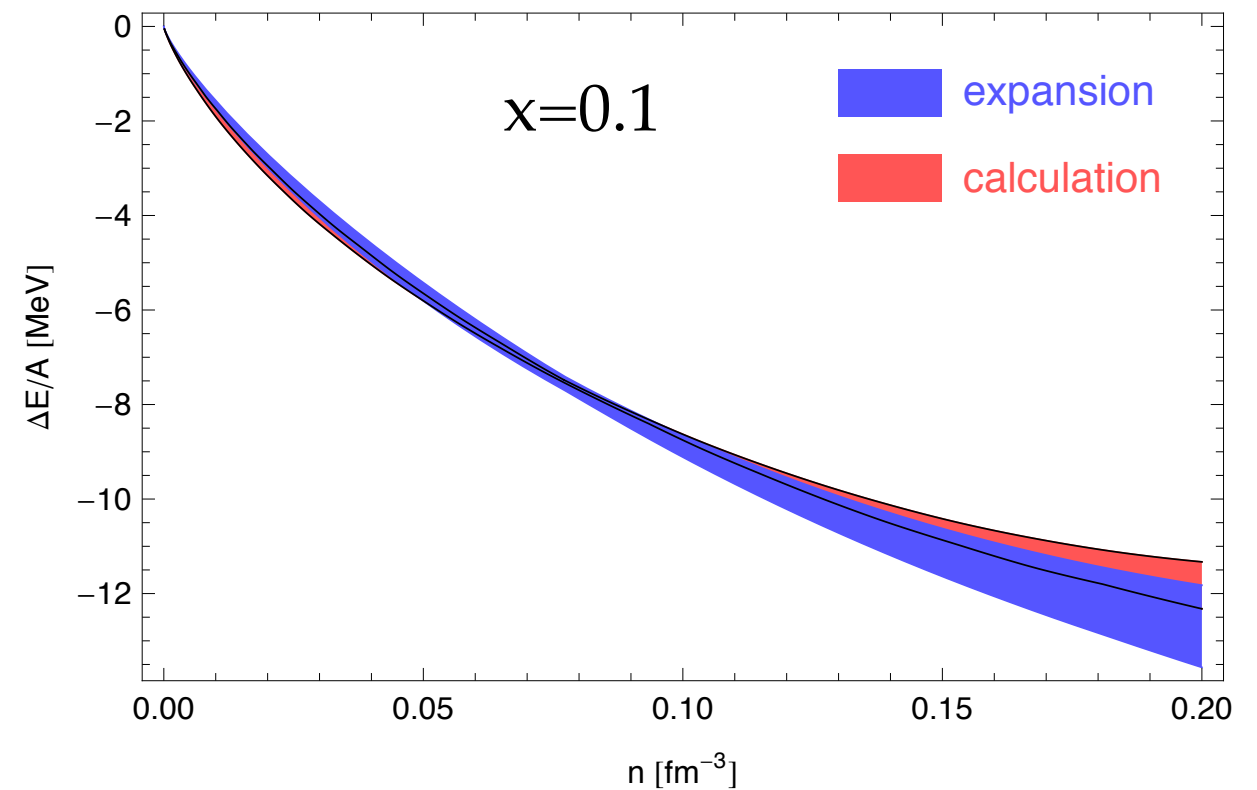
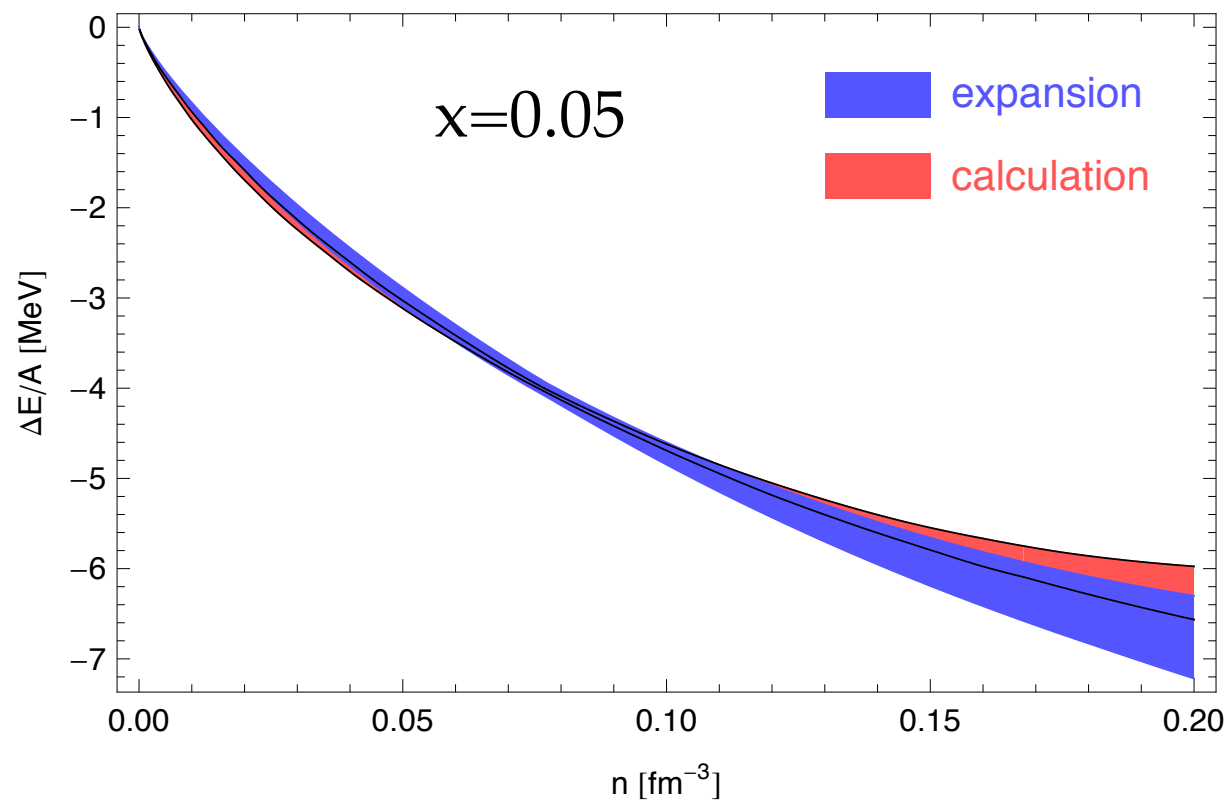
⇒ Parameters determined from

- saturation properties of symmetric matter
- neutron matter energy and pressure calculations

# Validation of quasi-parabolic expansion

✱ Compare energy relative to pure neutron matter

$$\frac{\Delta E(n, x)}{A} = \frac{E(n, x)}{A} - \frac{E(n, x = 0)}{A}$$



⇒ Excellent agreement within uncertainties

⇒ Possible improvements in the expansion



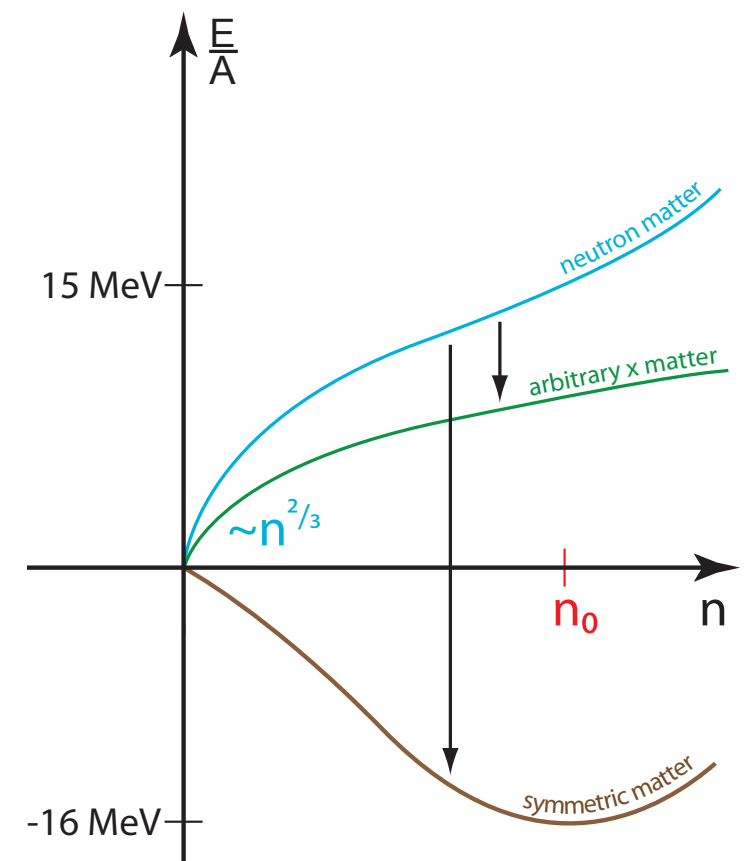
# Conclusions and outlook (part I)

✱ First explicit calculation of asymmetric matter with chiral interactions

- ⇒ Constrain isospin dependence of nuclear EOS
- ⇒ Revisit parabolic approximation for small proton fractions

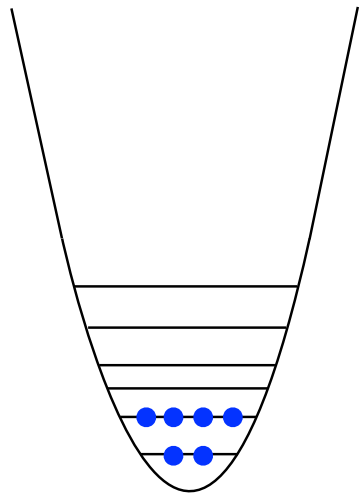
✱ Outlook

- ⇒ Improve many-body expansion
- ⇒ Use consistently evolved NN+3N interactions
- ⇒ Release small proton fraction approximation



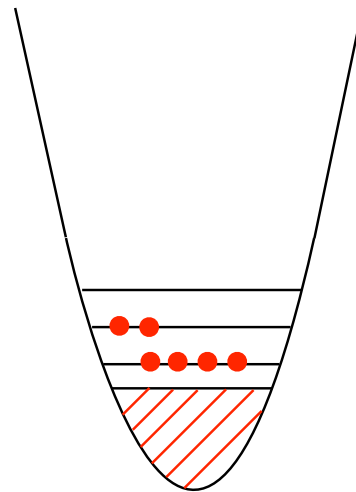
# Open-shell medium-mass nuclei

Light nuclei



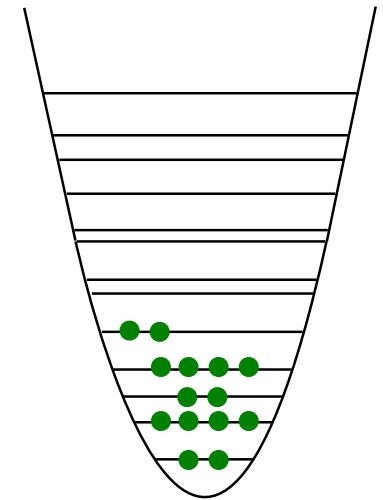
NCSM, GFMC, ....

Medium-mass nuclei



Microscopic SM, ....

Medium-mass nuclei



GF, CC, IM-SRG, ....

⇒ Configuration interaction limited to small valence/model spaces

⇒ Usual expansion schemes fail to account for pairing correlations

⇒ Limited to doubly-closed-shell  $\pm 1$  and  $\pm 2$  nuclei

# Going open-shell: Gorkov's idea

---

- ✱ Bogoliubov algebra + Green's function theory
- ✱ Keep the simplicity of a single-reference
- ✱ Address explicitly the non-perturbative formation of Cooper pairs
  - ⇒ Formulate the expansion scheme around a Bogoliubov vacuum
  - ⇒ Breaking of particle-number conservation (eventually restored)

# Going open-shell: Gorkov's formalism

	<u>Dyson</u>	<u>Gorkov</u>
◦ Reference state	$ \psi_0^A\rangle$	$ \Psi_0\rangle \equiv \sum_A^{\text{even}} c_A  \psi_0^A\rangle$
◦ One-body propagator	$G_{ab} = \begin{array}{c} \parallel \\ \uparrow \\ \parallel \end{array}$	$\mathbf{G}_{ab} = \begin{pmatrix} G_{ab}^{11} & G_{ab}^{12} \\ G_{ab}^{21} & G_{ab}^{22} \end{pmatrix} = \left( \begin{array}{c} \parallel \uparrow \parallel \uparrow \\ \parallel \downarrow \parallel \downarrow \end{array} \right)$
◦ Hamiltonian	$H$	$\Omega = H - \mu A$
◦ G.S. energy	$E_0^A$	$\Omega_0 = \sum_{A'}  c_{A'} ^2 \Omega_0^{A'} \approx E_0^A - \mu A$

## ✱ Set of 4 Green's functions

$$i G_{ab}^{11}(t, t') \equiv \langle \Psi_0 | T \{ a_a(t) a_b^\dagger(t') \} | \Psi_0 \rangle \equiv \begin{array}{c} a \\ \updownarrow \\ b \end{array}$$

$$i G_{ab}^{12}(t, t') \equiv \langle \Psi_0 | T \{ a_a(t) \bar{a}_b(t') \} | \Psi_0 \rangle \equiv \begin{array}{c} a \\ \updownarrow \\ \bar{b} \end{array}$$

$$i G_{ab}^{21}(t, t') \equiv \langle \Psi_0 | T \{ \bar{a}_a^\dagger(t) a_b^\dagger(t') \} | \Psi_0 \rangle \equiv \begin{array}{c} \bar{a} \\ \updownarrow \\ b \end{array}$$

$$i G_{ab}^{22}(t, t') \equiv \langle \Psi_0 | T \{ \bar{a}_a^\dagger(t) \bar{a}_b(t') \} | \Psi_0 \rangle \equiv \begin{array}{c} \bar{a} \\ \updownarrow \\ \bar{b} \end{array}$$



$$\mathbf{G}_{ab}(\omega) = \mathbf{G}_{ab}^{(0)}(\omega) + \sum_{cd} \mathbf{G}_{ac}^{(0)}(\omega) \Sigma_{cd}^*(\omega) \mathbf{G}_{db}(\omega)$$

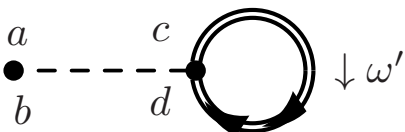
Gorkov equations

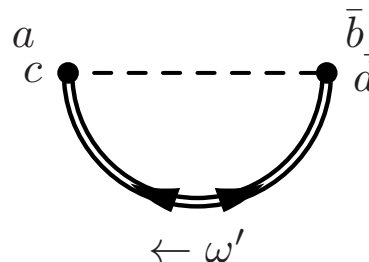
[Gorkov 1958]



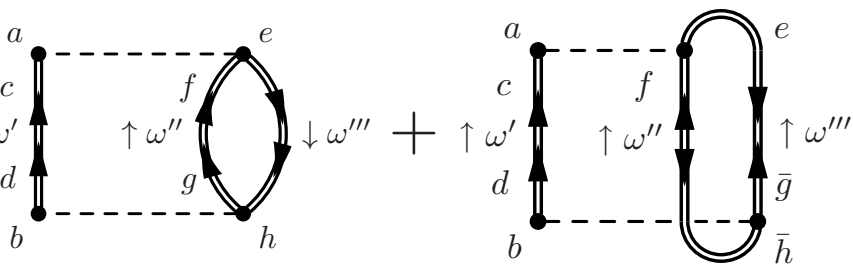
# Self-energy expansion

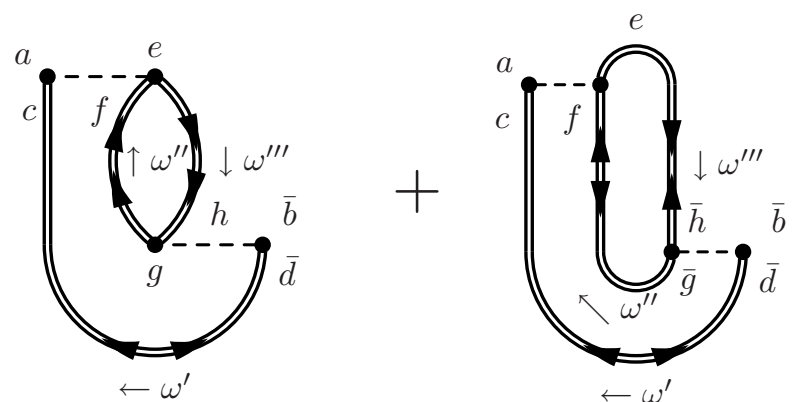
✱ 1<sup>st</sup> order  $\Rightarrow$  energy-independent self-energy

$$\Sigma_{ab}^{11(1)} =$$


$$\Sigma_{ab}^{12(1)} =$$


✱ 2<sup>nd</sup> order  $\Rightarrow$  energy-dependent self-energy

$$\Sigma_{ab}^{11(2)}(\omega) =$$


$$\Sigma_{ab}^{12(2)}(\omega) =$$


✱ Gorkov equations  $\longrightarrow$  energy *dependent* eigenvalue problem

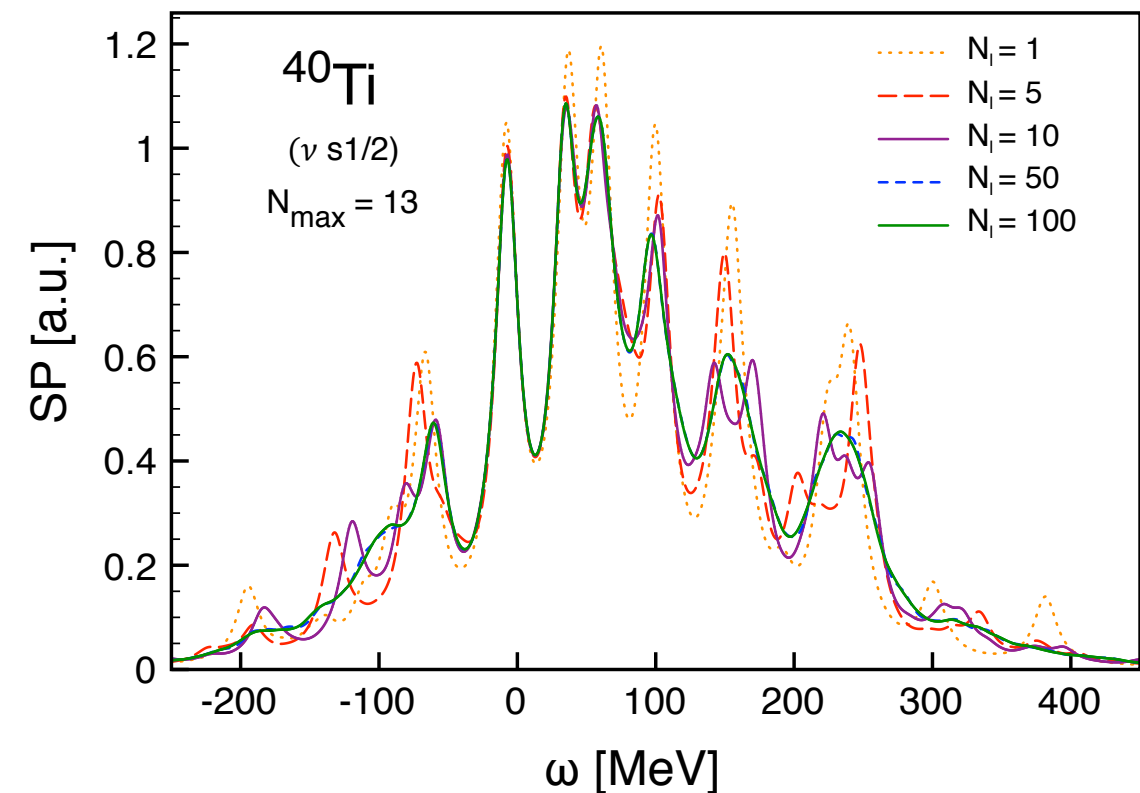
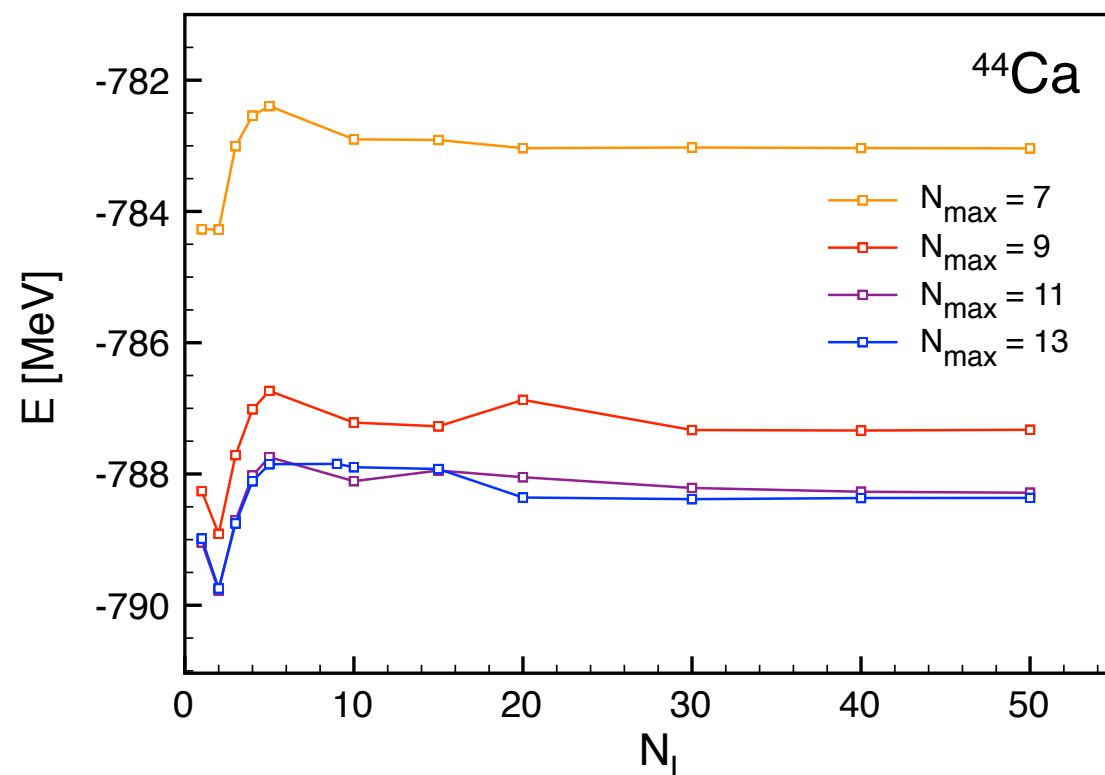
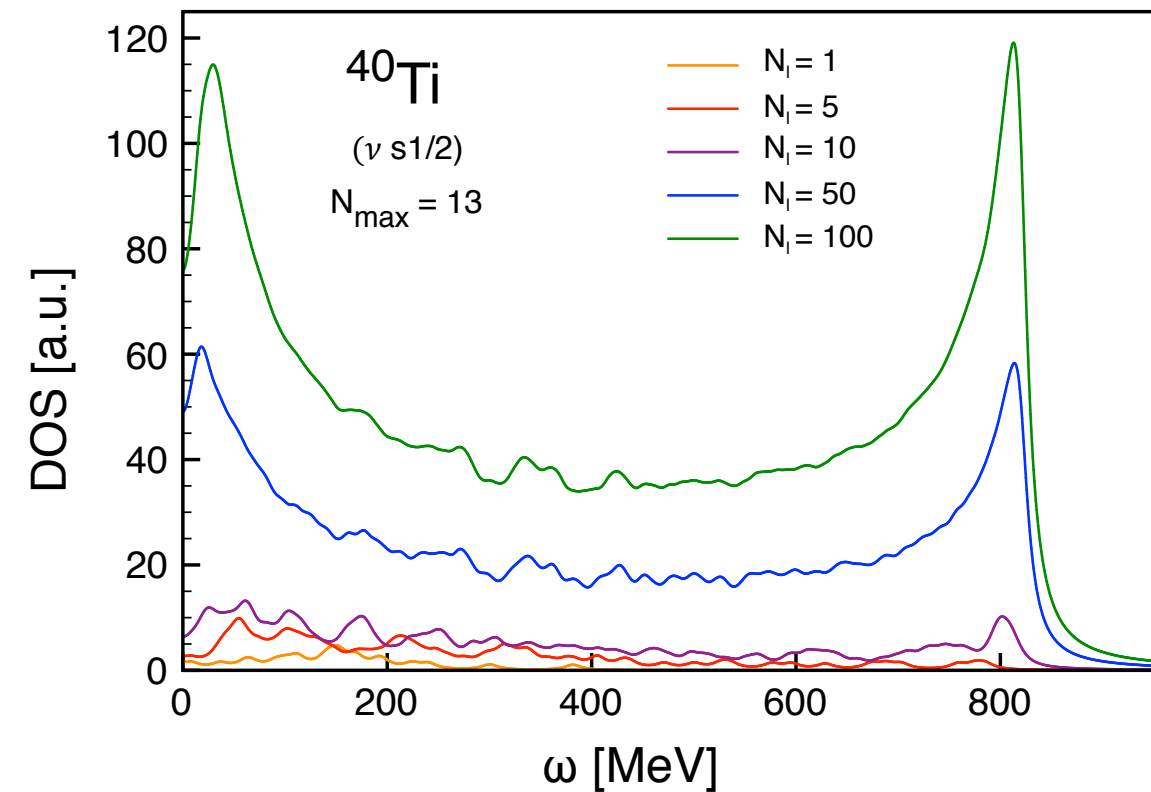
$$\sum_b \begin{pmatrix} t_{ab} - \mu_{ab} + \Sigma_{ab}^{11}(\omega) & \Sigma_{ab}^{12}(\omega) \\ \Sigma_{ab}^{21}(\omega) & -t_{ab} + \mu_{ab} + \Sigma_{ab}^{22}(\omega) \end{pmatrix} \Big|_{\omega_k} \begin{pmatrix} \mathcal{U}_b^k \\ \mathcal{V}_b^k \end{pmatrix} = \omega_k \begin{pmatrix} \mathcal{U}_a^k \\ \mathcal{V}_a^k \end{pmatrix}$$

$$\mathcal{U}_a^{k*} \equiv \langle \Psi_k | \bar{a}_a^\dagger | \Psi_0 \rangle$$

$$\mathcal{V}_a^{k*} \equiv \langle \Psi_k | a_a | \Psi_0 \rangle$$

# Krylov projection

- ⇒ Krylov projection of Gorkov's matrix
- ⇒ Lanczos-like algorithm
- ⇒ Conserves moments of spectral functions
- ⇒ Results become independent of  $N_1$

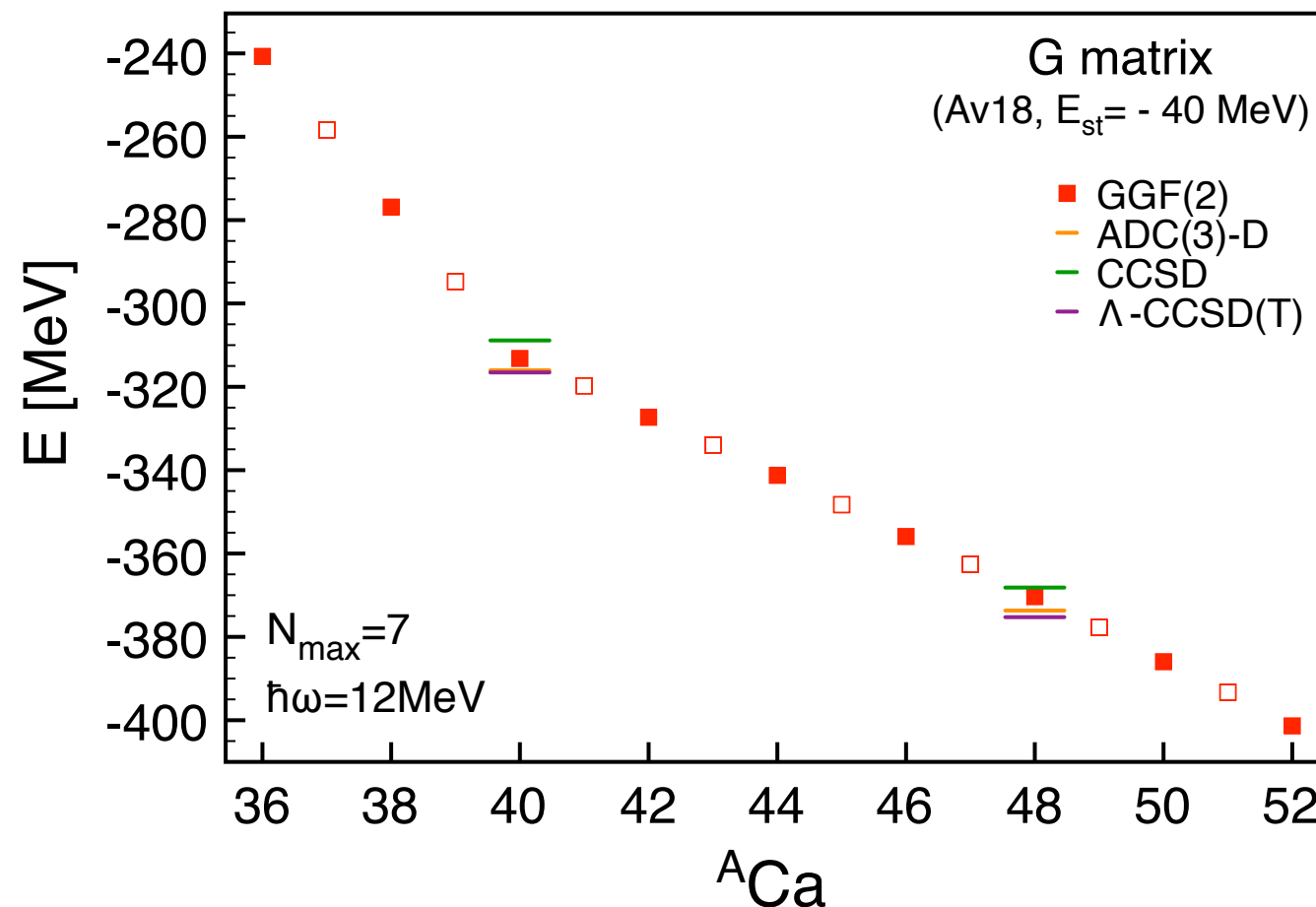




# Benchmark with coupled cluster method

✱ Energy from Galitskii-Koltun sum rule

$$E_0^A = \frac{1}{4\pi i} \int_{C_{\uparrow}} d\omega \operatorname{Tr}_{\mathcal{H}_1} [\mathbf{G}^{11}(\omega) [\mathbf{T} + (\mu + \omega) \mathbf{1}]]$$



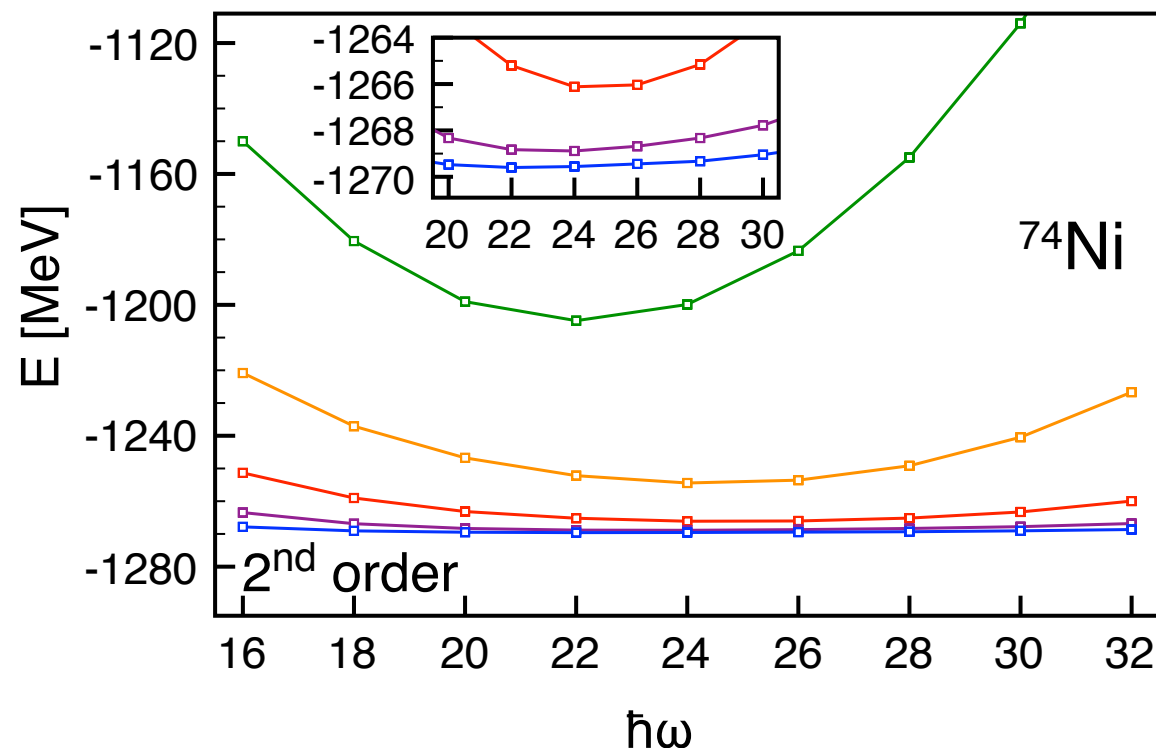
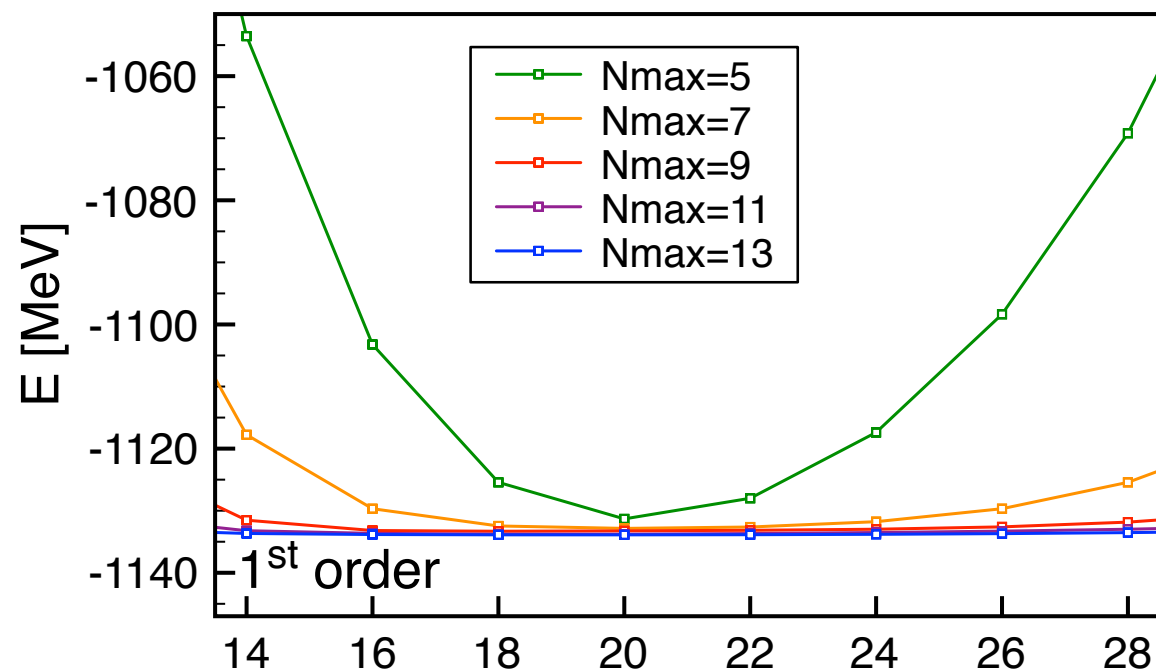
⇒ GGF and CC quantitatively similar

⇒ GGF(3) expected to reach  $\Lambda$ -CCSD(T) accuracy

(CC results courtesy of G. Hagen)

# Towards medium / heavy open-shell

✱  $^{74}\text{Ni}$



⇒ NN interaction:  
chiral  $\text{N}^3\text{LO}$  SRG-evolved to  $2.0 \text{ fm}^{-1}$

[Entem and Machleidt 2003]

⇒ Very good convergence

⇒ From  $N=13$  to  $N=11 \rightarrow 200 \text{ keV}$

$$E(N=13) = -1269.6 \text{ MeV}$$

$$E(N=\infty) = -1269.7(2) \text{ MeV}$$

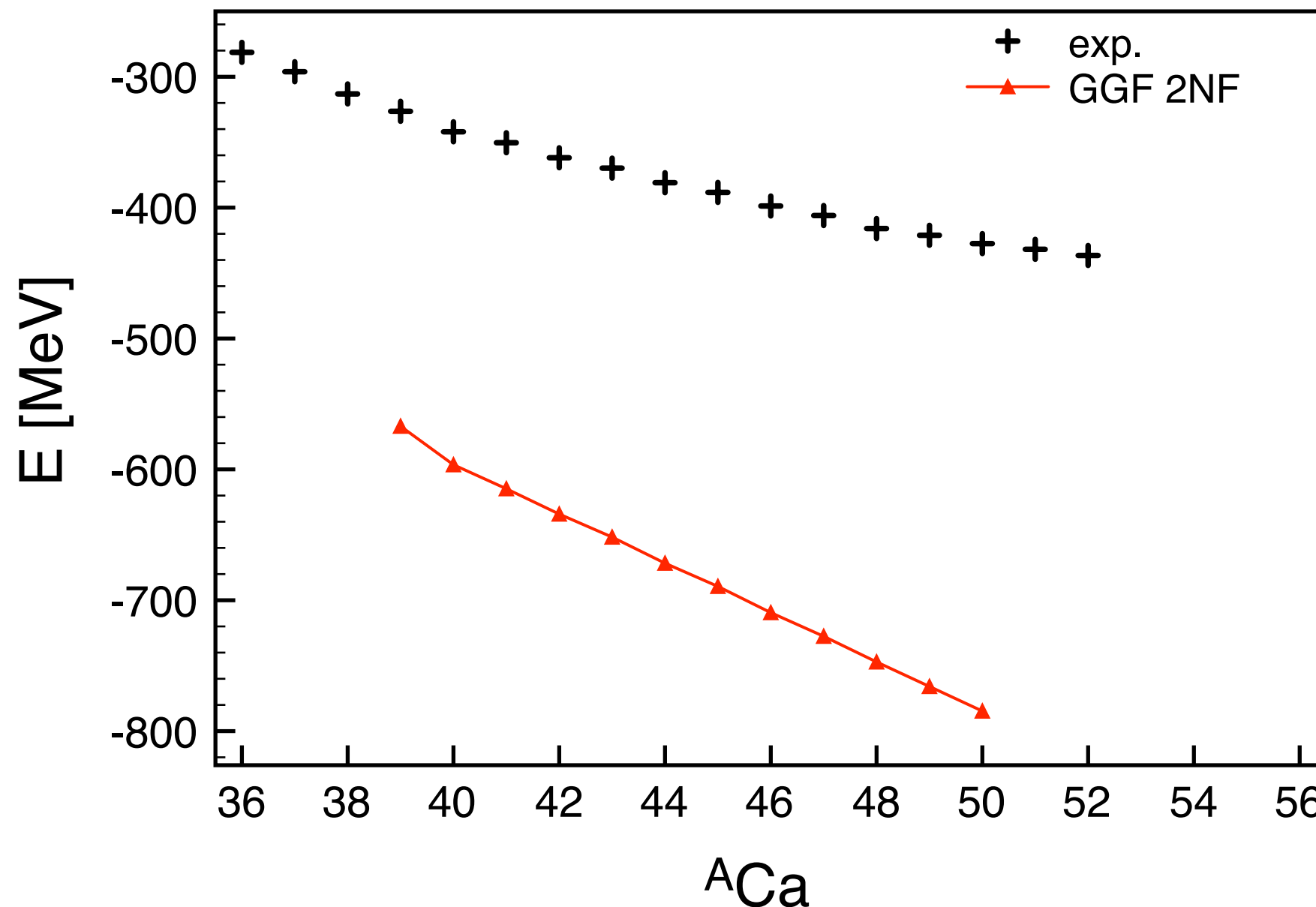
(Extrapolation to infinite model space from  
[Furnstahl, Hagen, Papenbrok 2012]  
and [Coon et al. 2012])

# Calcium isotopic chain

✱ NN only

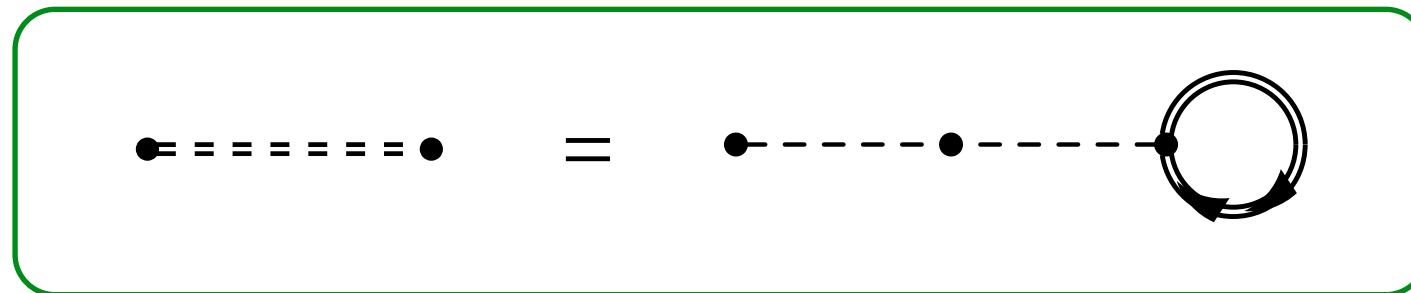
⇒ Systematic along isotopic/isotonic becomes available

⇒ Overbinding (increasing with  $A$ ): need for three-body forces



## ✱ Inclusion of 3NF as effective 2NF

⇒ Average over the 3<sup>rd</sup> nucleon in each nucleus



⇒ Additional term in the Galitskii-Koltun sum rule

$$E_0^A = \frac{1}{4\pi i} \int_{C_\uparrow} d\omega \operatorname{Tr}_{\mathcal{H}_1} [\mathbf{G}^{11}(\omega) [\mathbf{T} + (\mu + \omega) \mathbf{1}]] - \frac{1}{2} \langle \Psi_0 | W | \Psi_0 \rangle$$

## ✱ 3N interaction: chiral N<sup>2</sup>LO (400 MeV) SRG-evolved to 2.0 fm<sup>-1</sup> [Navrátil 2007]

⇒ Fit to **three-** and **four-body** systems only

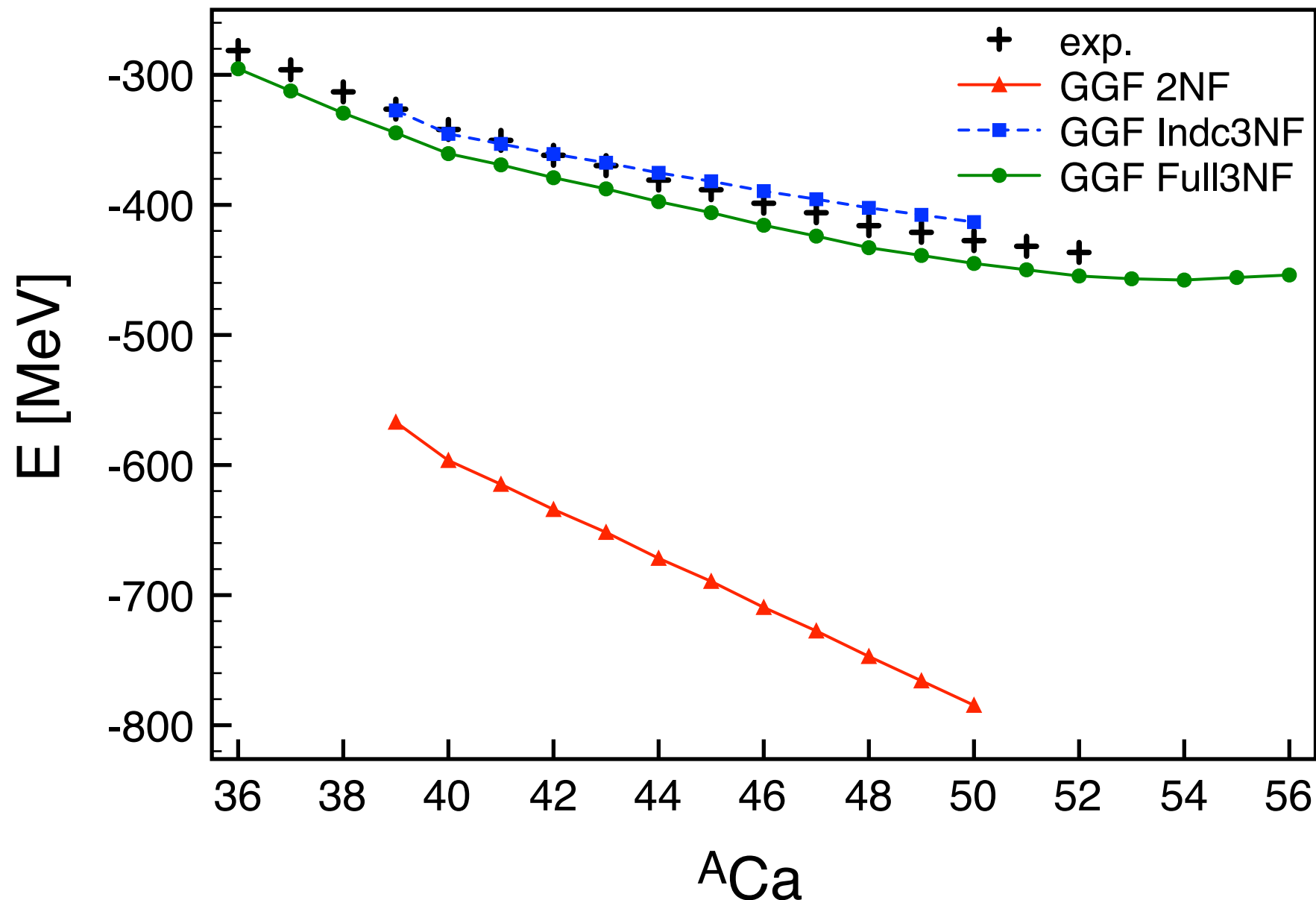
⇒ Modified cutoff to reduce induced 4N contributions [Roth *et al.* 2012]

# Calcium isotopic chain

✱ First *ab initio* calculation of the whole Ca chain with NN + 3N forces

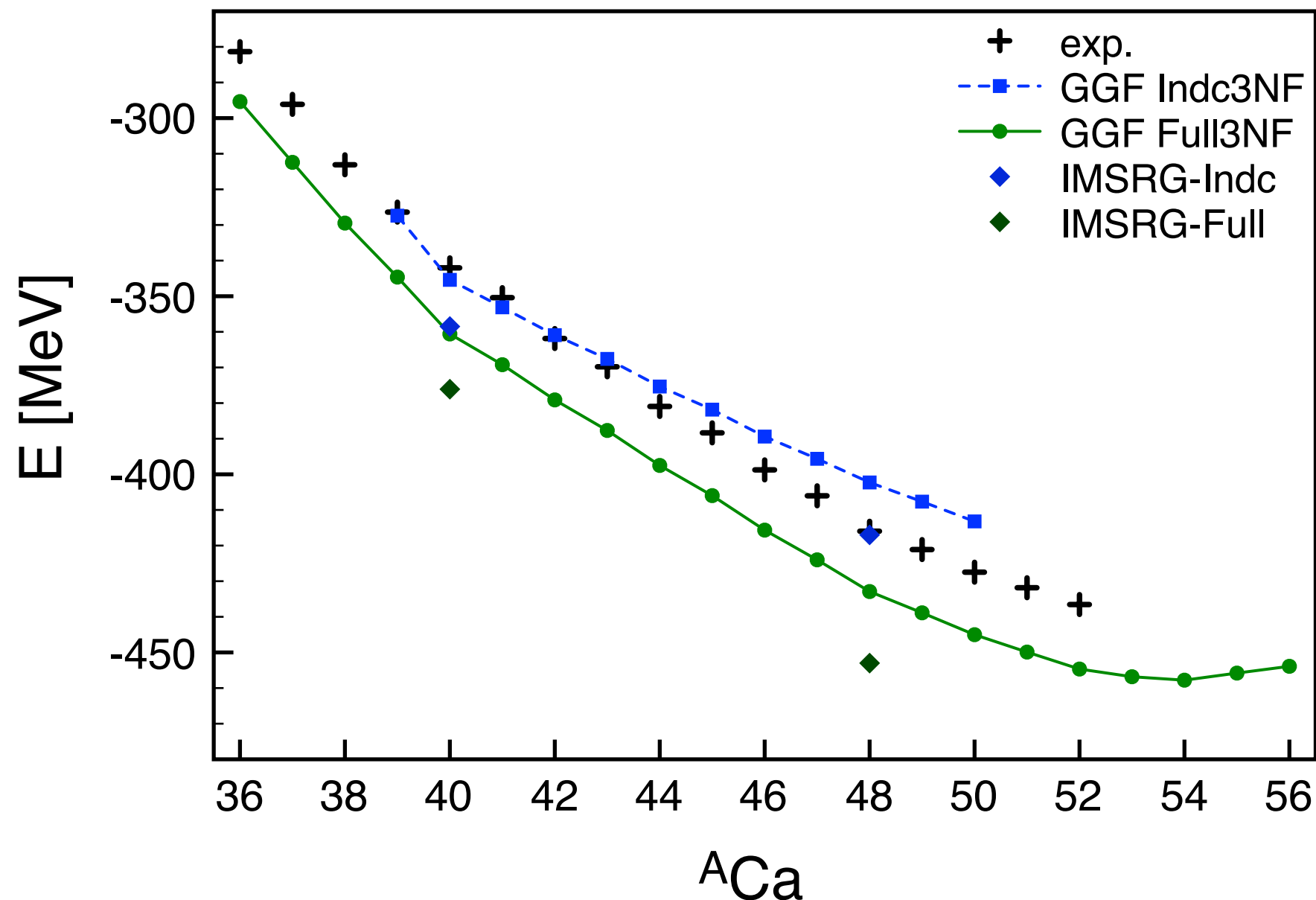
⇒ 3NF bring energies close to experiment

⇒ Induced 3NF and full 3NF investigated



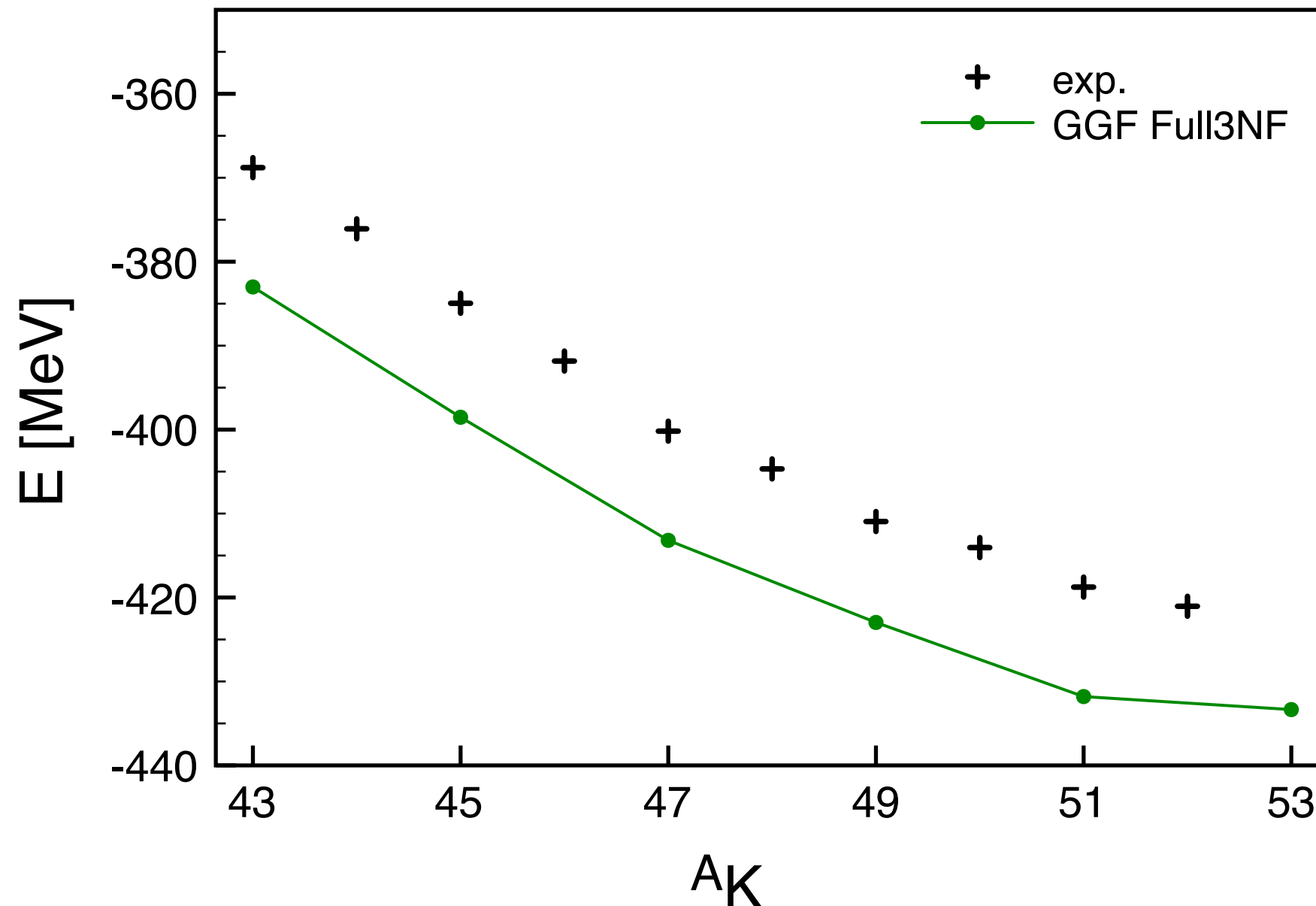
# Calcium isotopic chain

- ⇒ Original 3NF correct the energy curvature
- ⇒ Good agreement with IM-SRG (quantitative when 3<sup>rd</sup> order included)



# Potassium isotopic chain

- ✱ Exploit the odd-even formalism: application to K
  - ➡ Trend and agreement similar to calcium
  - ➡ Future: consistent description of medium-mass driplines

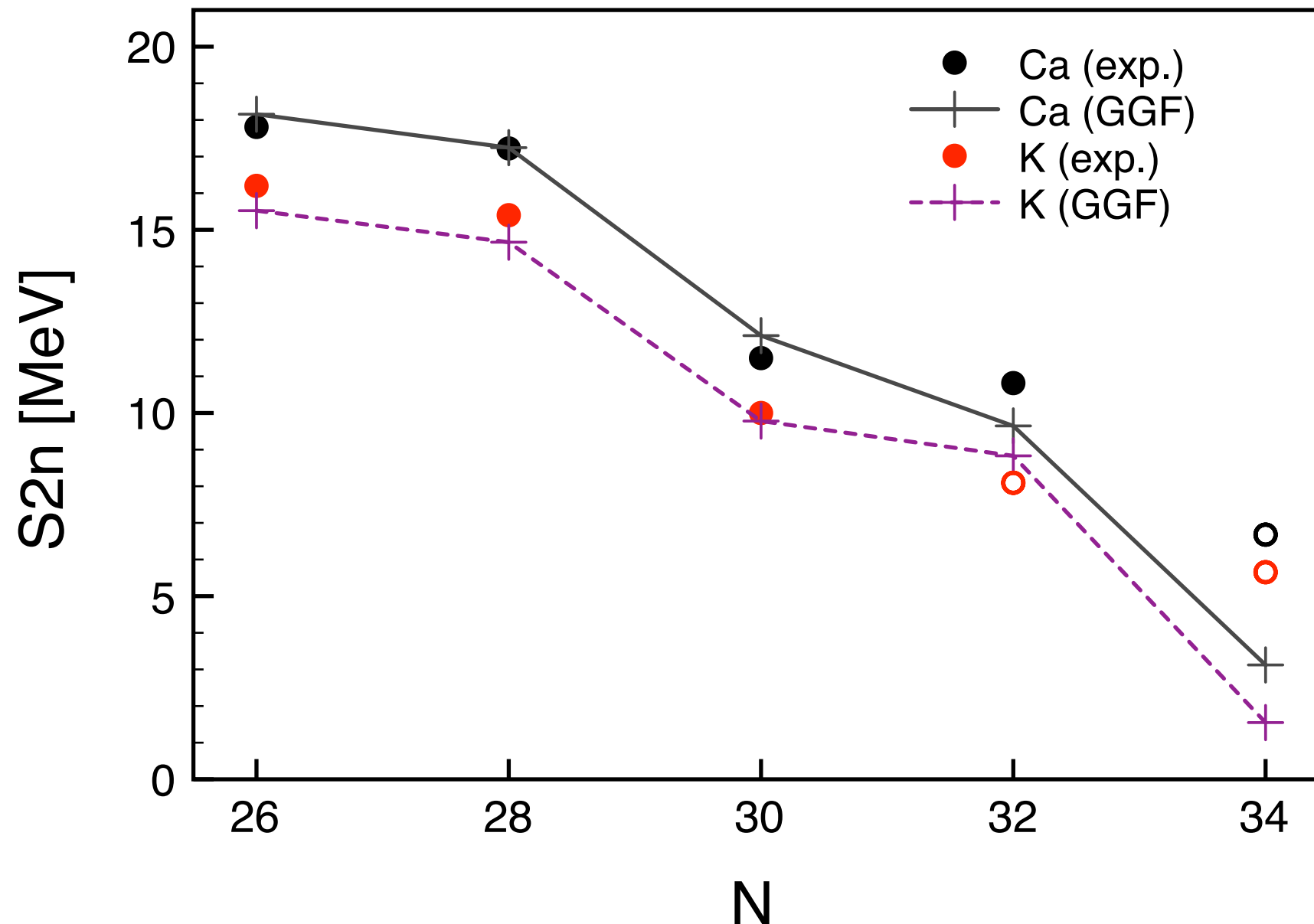


# Two-neutron separation energies

✱ Neutron-rich extremes of the nuclear chart

➡ Good agreement with measured  $S_{2n}$

➡ Towards a quantitative *ab initio* description of the medium-mass region



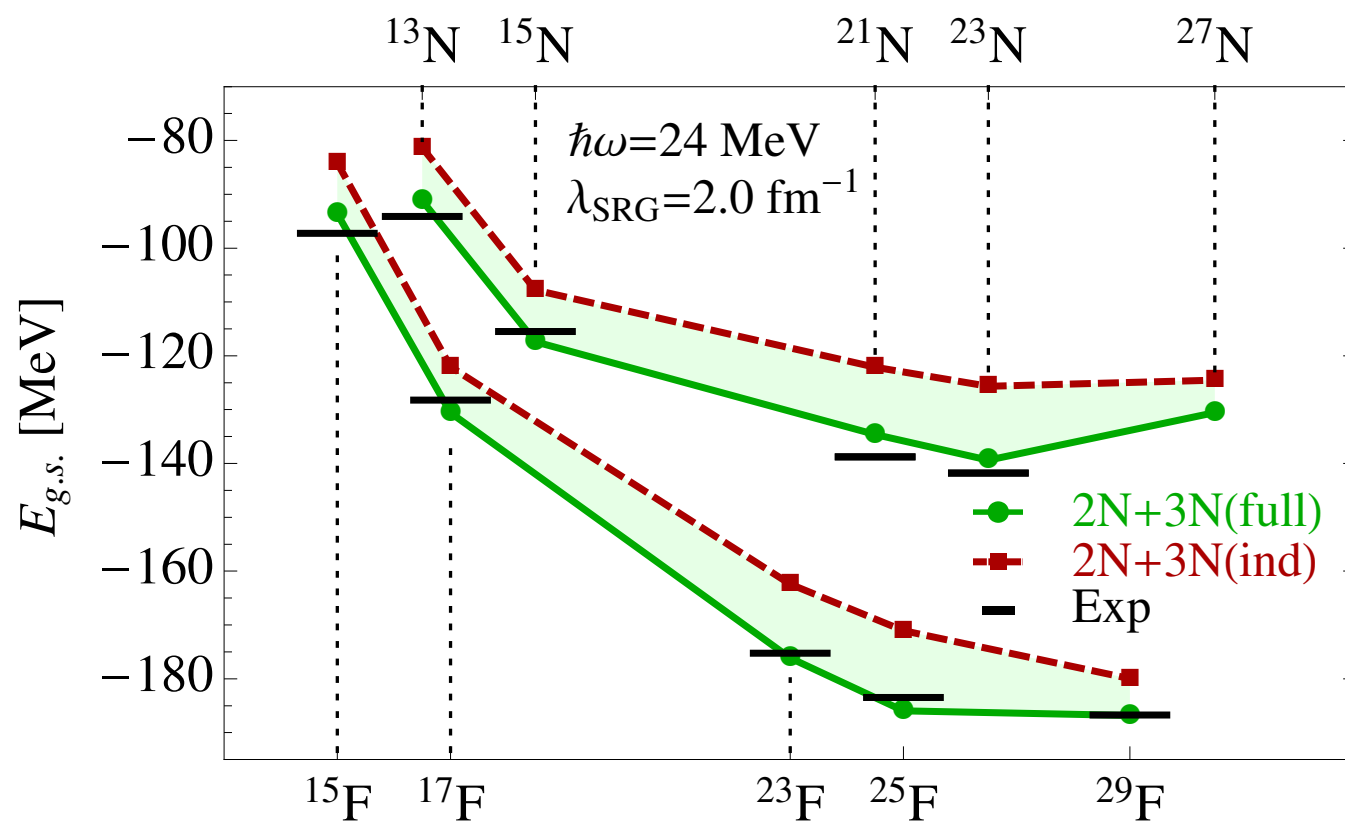


# Benchmarks and chiral EFT interactions

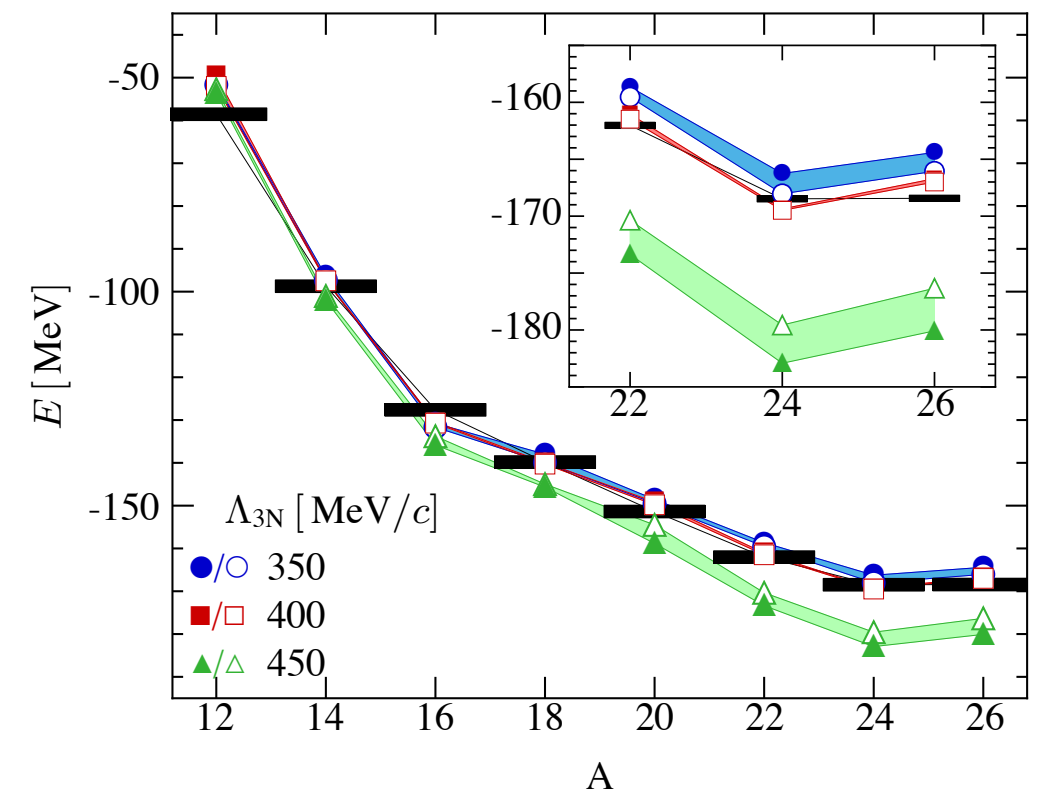
✱ *Ab initio* calculations as a test for chiral EFT interactions

✱ Different approaches agree in O and Ca chains

➡ Current chiral NN+3N forces overbind medium/heavy-mass nuclei



[Cipollone, Barbieri, Navrátil, 2013]



[Hergert *et al.*, 2013]

# Spectrum and spectroscopic factors

## ✱ Separation energy spectrum

$$G_{ab}^{11}(\omega) = \sum_k \left\{ \frac{\mathcal{U}_a^k \mathcal{U}_b^{k*}}{\omega - \omega_k + i\eta} + \frac{\bar{\mathcal{V}}_a^{k*} \bar{\mathcal{V}}_b^k}{\omega + \omega_k - i\eta} \right\}$$

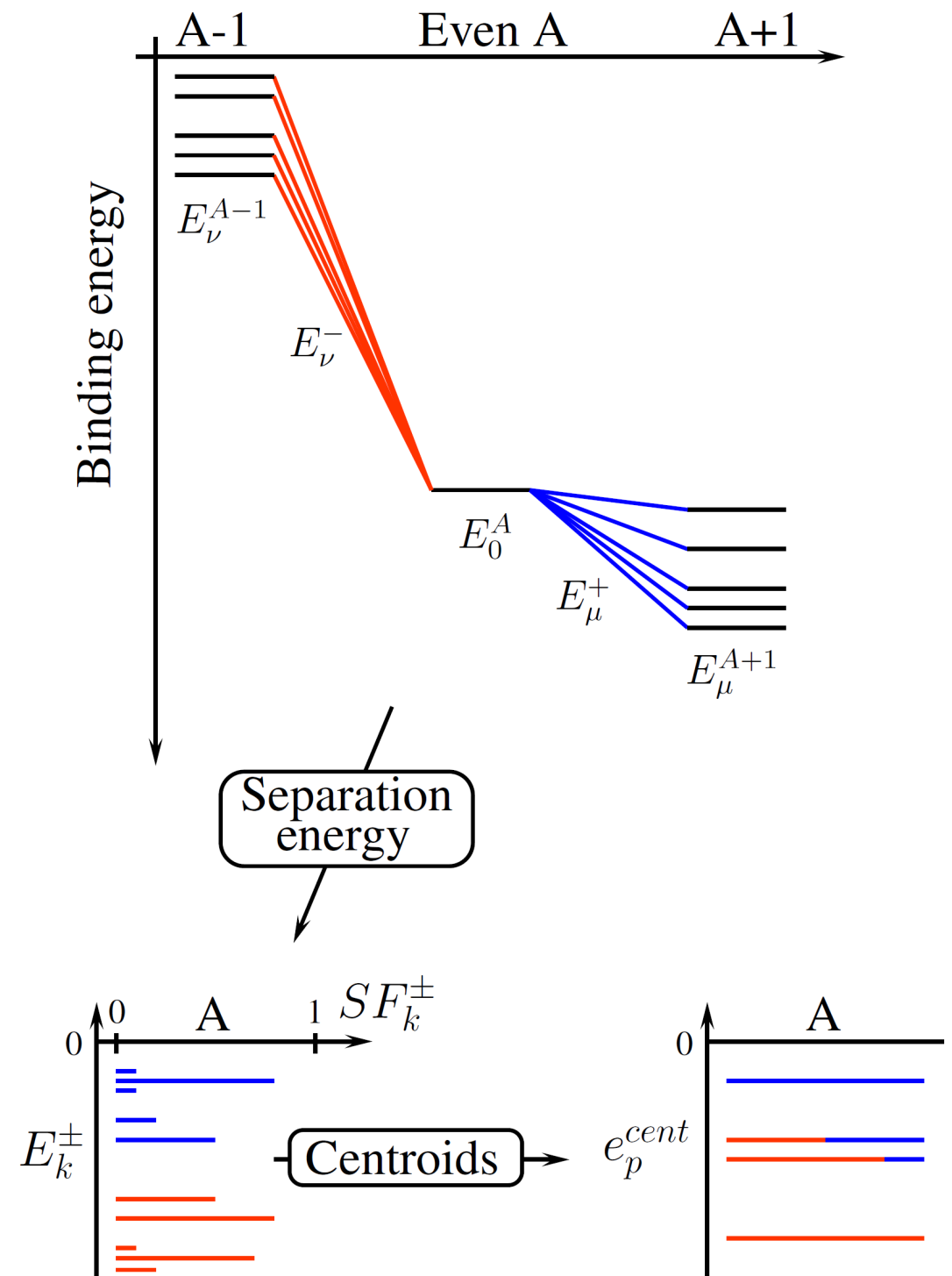
Lehmann representation

where 
$$\begin{cases} \mathcal{U}_a^{k*} \equiv \langle \Psi_k | a_a^\dagger | \Psi_0 \rangle \\ \mathcal{V}_a^{k*} \equiv \langle \Psi_k | \bar{a}_a | \Psi_0 \rangle \end{cases}$$

and 
$$\begin{cases} E_k^{+(A)} \equiv E_k^{A+1} - E_0^A \equiv \mu + \omega_k \\ E_k^{-(A)} \equiv E_0^A - E_k^{A-1} \equiv \mu - \omega_k \end{cases}$$

## ✱ Spectroscopic factors

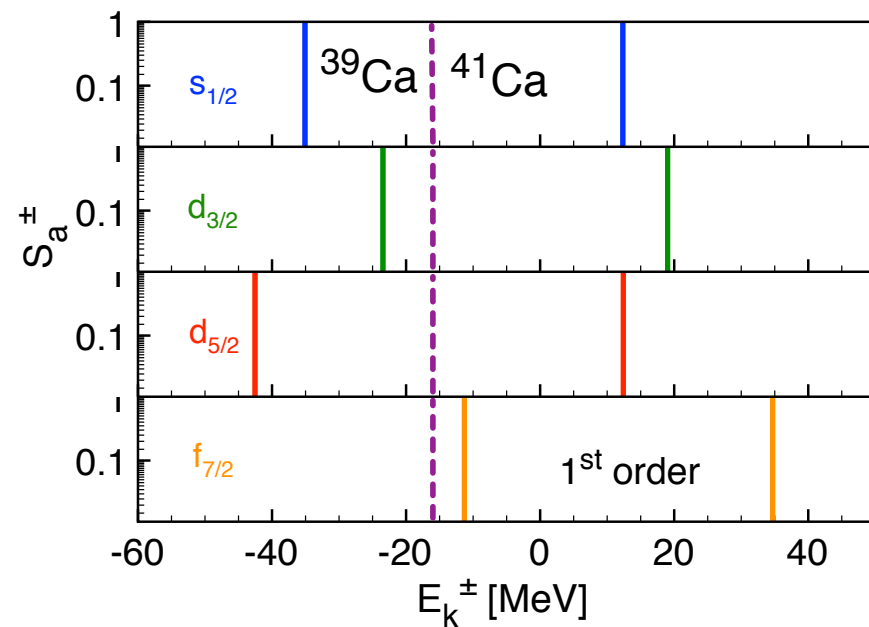
$$\begin{aligned} SF_k^+ &\equiv \sum_{a \in \mathcal{H}_1} |\langle \psi_k | a_a^\dagger | \psi_0 \rangle|^2 = \sum_{a \in \mathcal{H}_1} |\mathcal{U}_a^k|^2 \\ SF_k^- &\equiv \sum_{a \in \mathcal{H}_1} |\langle \psi_k | a_a | \psi_0 \rangle|^2 = \sum_{a \in \mathcal{H}_1} |\mathcal{V}_a^k|^2 \end{aligned}$$



[figure from Sadoudi]

# Spectral strength distribution

Dyson 1<sup>st</sup> order (HF)

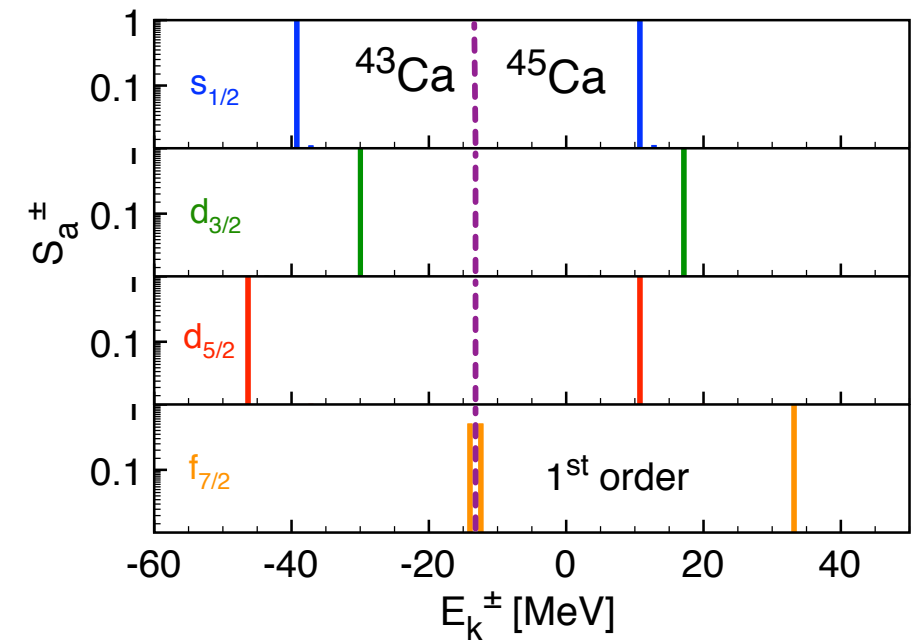


Fragmentation

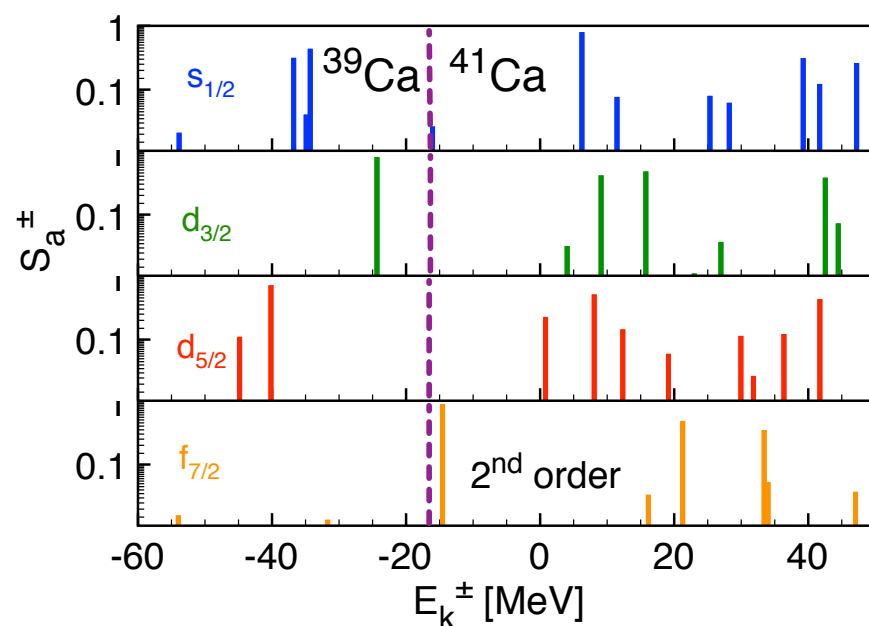
Static pairing



Gorkov 1<sup>st</sup> order (HFB)



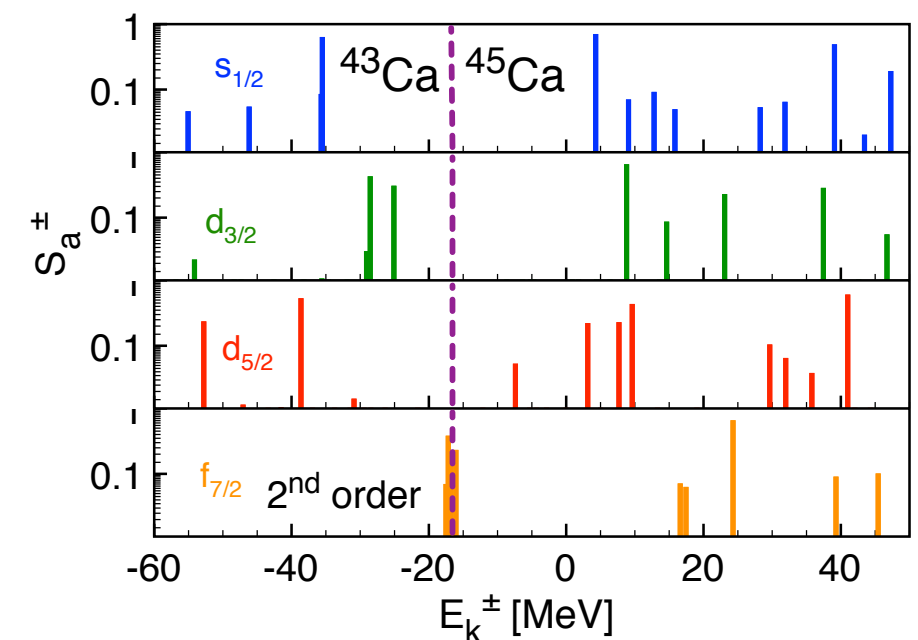
Dyson 2<sup>nd</sup> order



Dynamical  
fluctuations

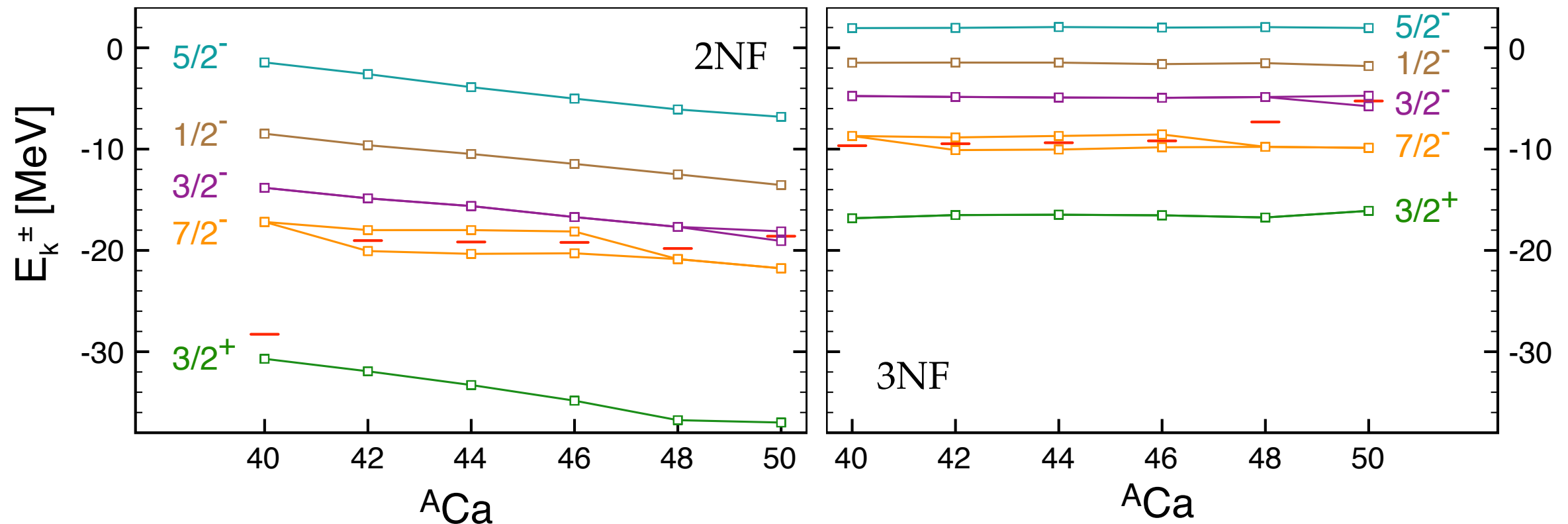


Gorkov 2<sup>nd</sup> order



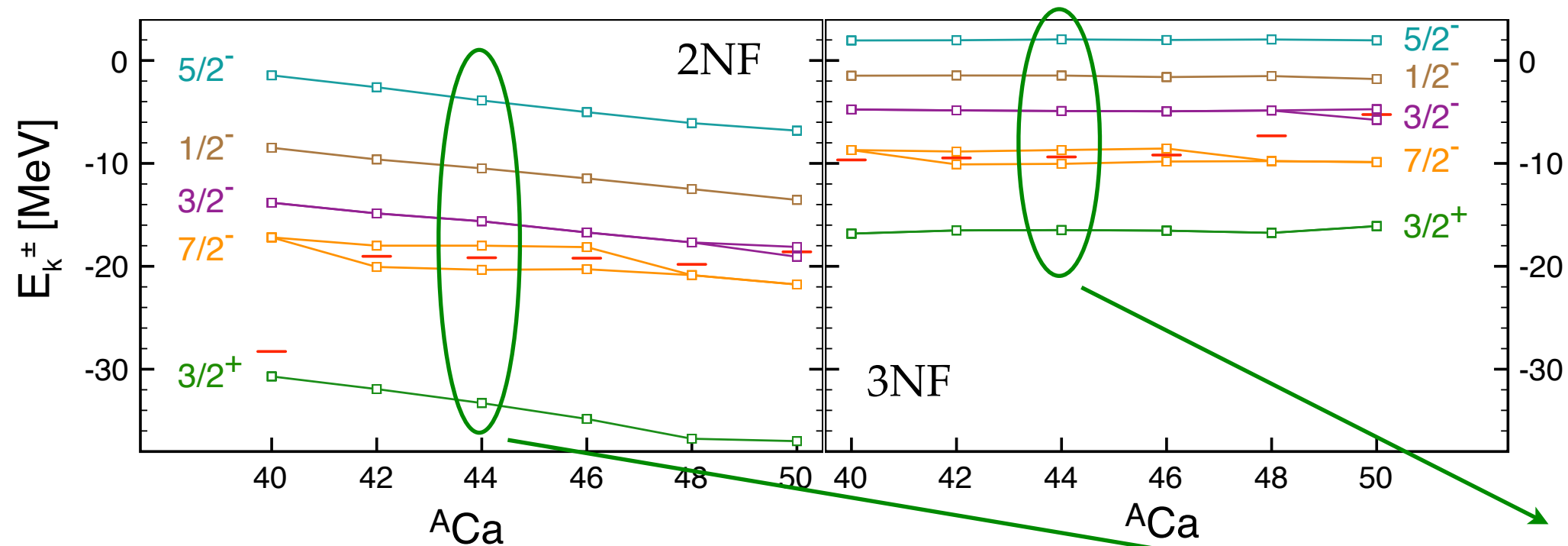
# Shell structure evolution

## ✱ One-neutron separation energies



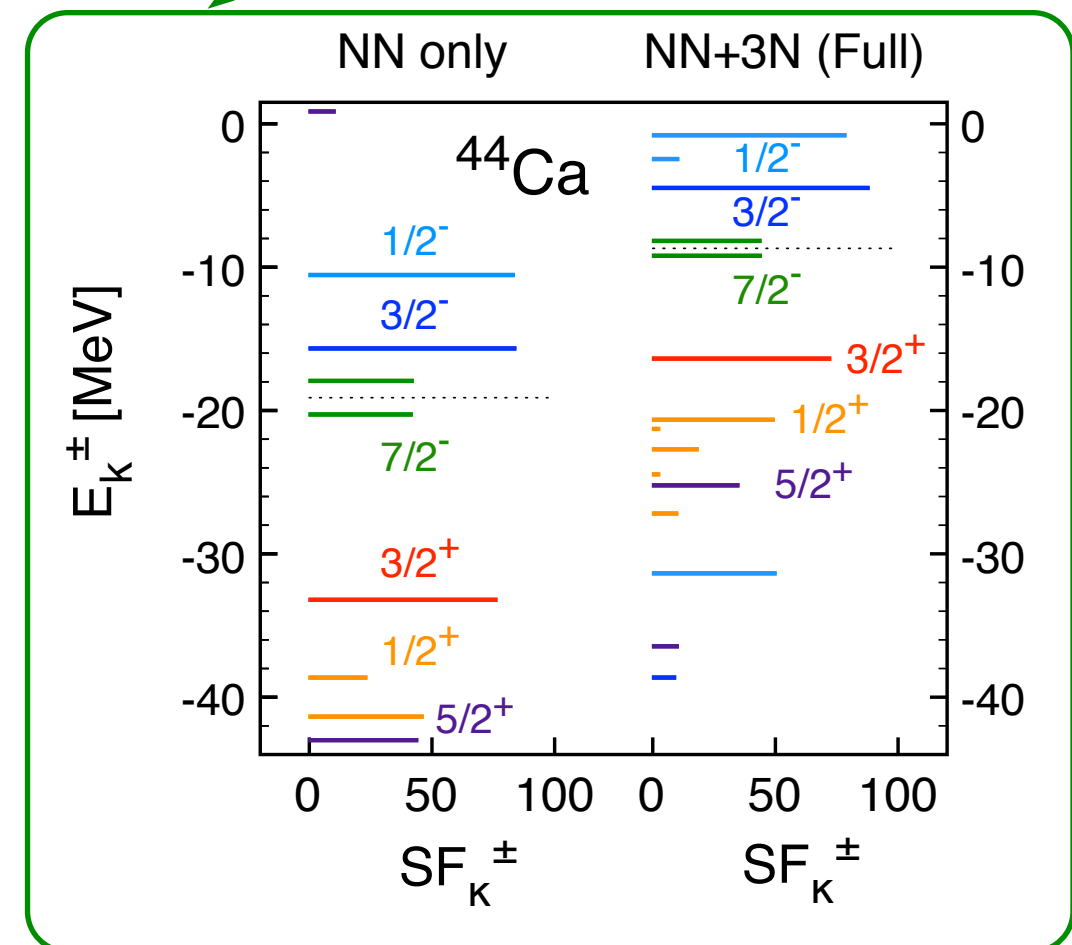
- ⇒ Static and dynamic pairing correlations
- ⇒ 3NF significantly compress spectrum

# Shell structure evolution



✱ NN vs. NN+3NF in  $^{44}Ca$

- ⇒ Similar degree of fragmentation
- ⇒ SF of main peaks not much affected

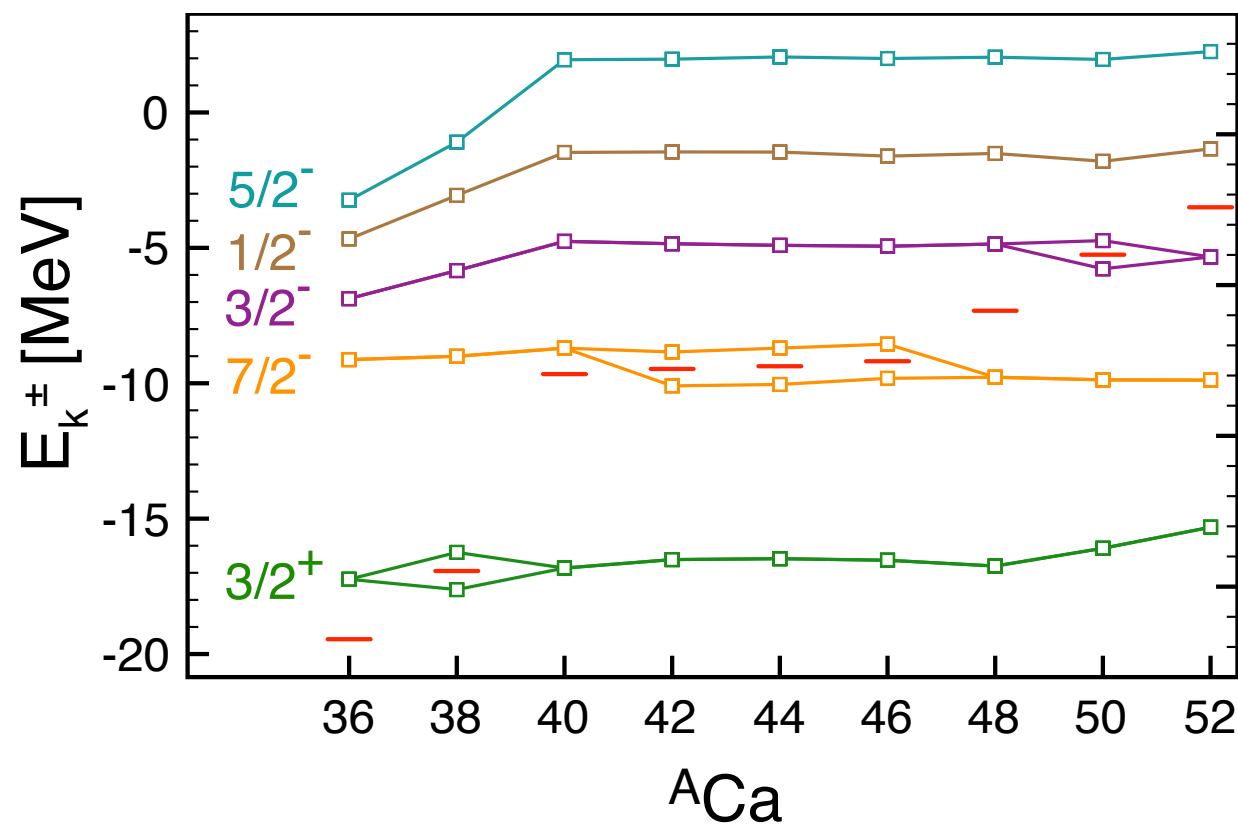


✱ ESPE collect fragmentation of “single-particle” strengths from both  $N \pm 1$

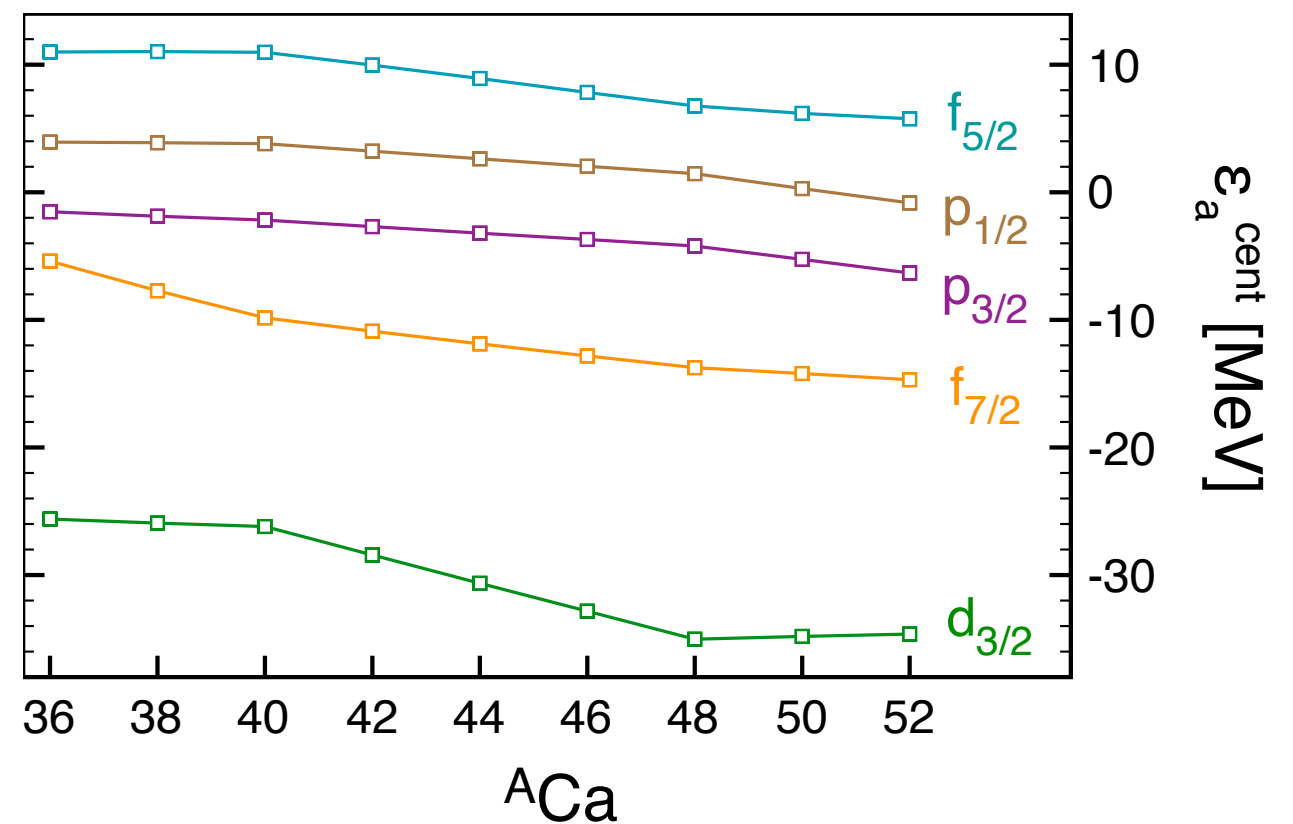
$$\epsilon_a^{cent} \equiv h_{ab}^{cent} \delta_{ab} = t_{aa} + \sum_{cd} \bar{V}_{acad}^{NN} \rho_{dc}^{[1]} + \sum_{cdef} \bar{V}_{acdaef}^{NNN} \rho_{efcd}^{[2]} \equiv \sum_k \mathcal{S}_k^{+a} E_k^+ + \sum_k \mathcal{S}_k^{-a} E_k^-$$

[Baranger 1970, Duguet and Hagen 2011]

Separation energies



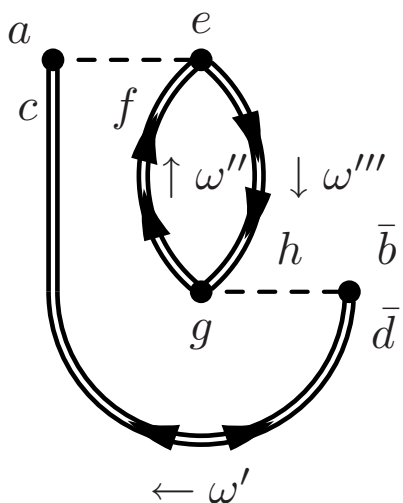
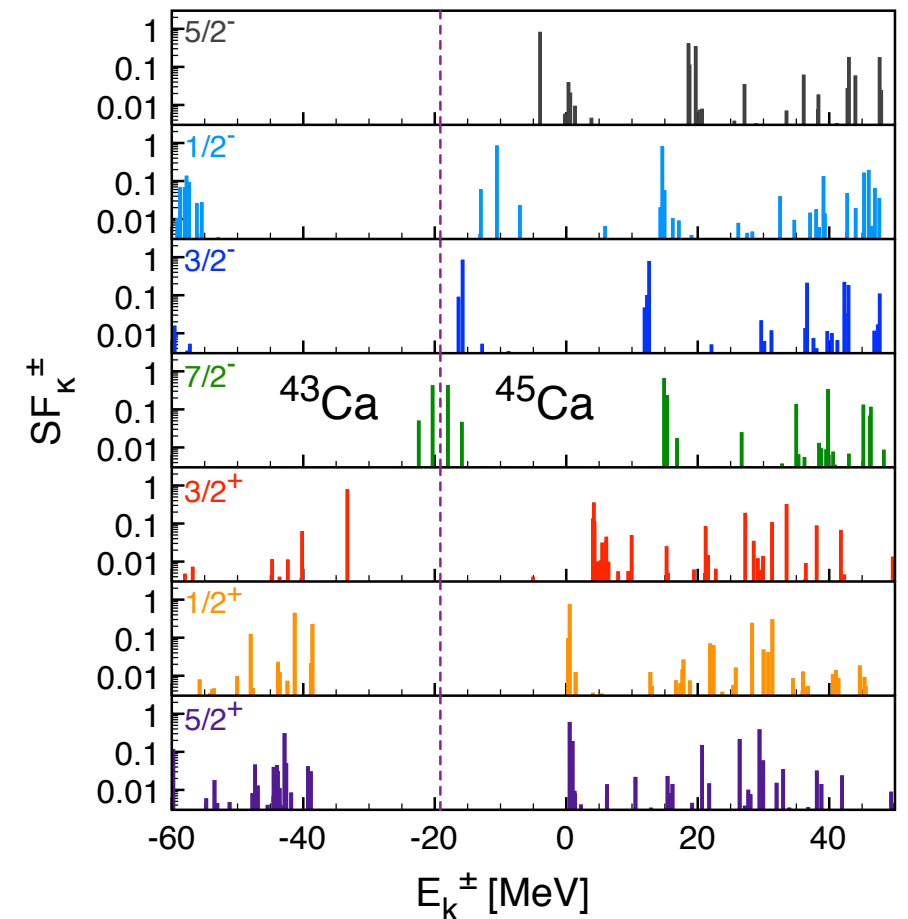
Centroids



# Conclusions and outlook (part II)

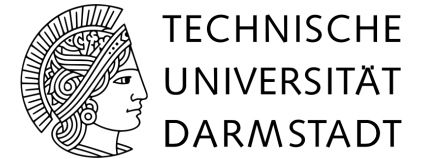
## ✱ Gorkov-Green's functions:

- ⇒ Manageable route to (near) degenerate systems
- ⇒ First *ab initio* description of medium-mass chains
- ⇒ 2NF + 3NF: towards predictive calculations
- ⇒ Energies: quantitative agreement
- ⇒ Spectra: study of shell structure evolution



- ✱ Improvement of the self-energy expansion
- ✱ Proper coupling to the continuum
- ✱ Formulation of **particle-number restored** Gorkov theory
- ✱ Towards consistent description of structure and reactions

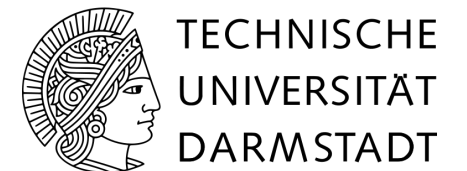
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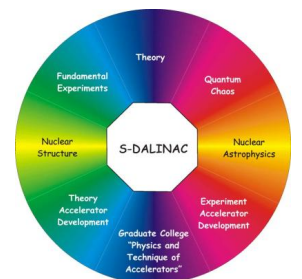


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