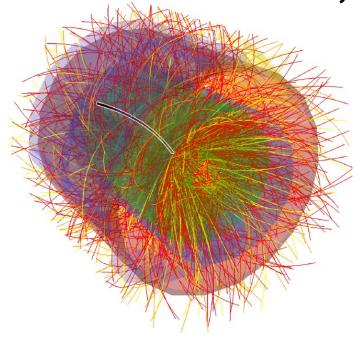
Production of loosely-bound objects in heavy-ion collisions at HADES, RHIC and LHC



April 14, 2025 GSI Physics of High net-baryon Densities - PHD 2025



Benjamin Dönigus

Institut für Kernphysik Goethe Universität Frankfurt



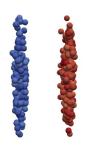


Content

- Introduction
- Nuclei and Exotica
 - (Anti-)nuclei
 - (Anti-)hypernuclei
- Summary & Outlook



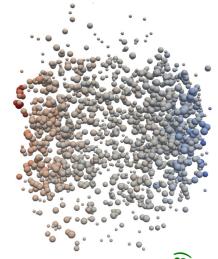
Introduction









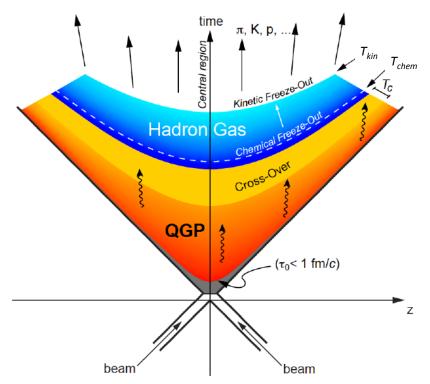


Cartoon of a Ultra-relativistic heavy-ion collision Left to right:

- the two Lorentz contracted nuclei approach,
- collide,
- form a Quark-Gluon Plasma (QGP),
- the QGP expands and hadronizes,
- finally hadrons rescatter and freeze



Introduction

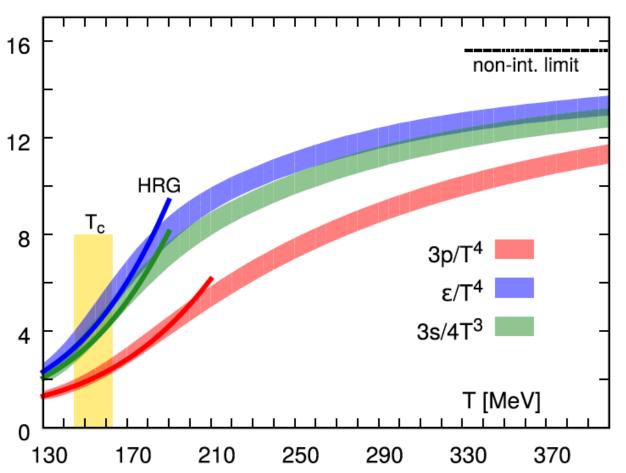


The fireball evolution:

- Starts with a "pre-equilibrium state"
- Forms a Quark-Gluon Plasma phase (if T is larger than T_c)
- At chemical freeze-out, T_{ch}, hadrons stop being produced
- At kinetic freeze-out, T_{fo}, hadrons stop scattering



Lattice QCD results



Lattice QCD tells us where to expect the phase transition

Critical energy density:

 $\epsilon_{\rm C}$ = 0.34 ± 0.16 GeV/fm³

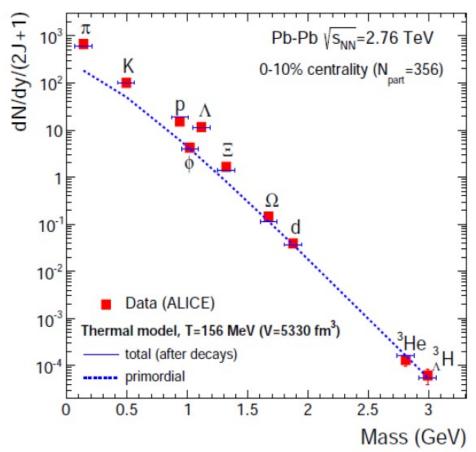
Critical temperature

 $T_C = (157 \pm 2) \text{ MeV}$

A. Bazavov et al. (hotQCD) Phys. Rev. D90 (2014) 094503 & PLB 795 (2019) 15 Similar results from Budapest-Wuppertal group: S. Borsányi et al. JHEP 09 (2010) 073 & PRL 125 (2020) 052001



Temperature of the source



Plot by A. Andronic, GSI-Heidelberg group arXiv:1407.5003 [nucl-ex]

Analogy:

Light source → particle source

 Multiplicity described best with T = 1 900 000 000 000 °C (1,9 trillion degree centigrade)

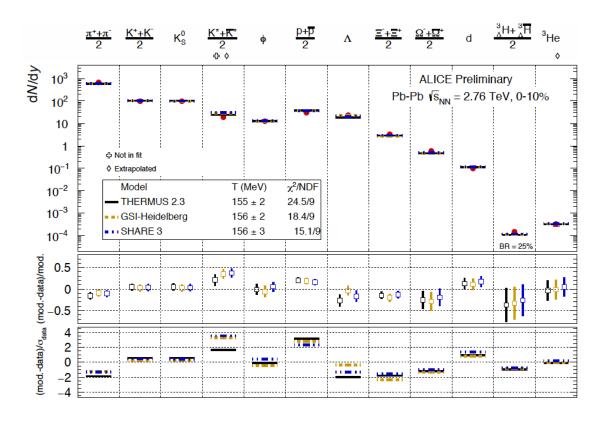
→ 100 000 times hotter than in the interior of the sun!

 $1/40 \text{ eV} = 20 ^{\circ}\text{C}$



Thermal model

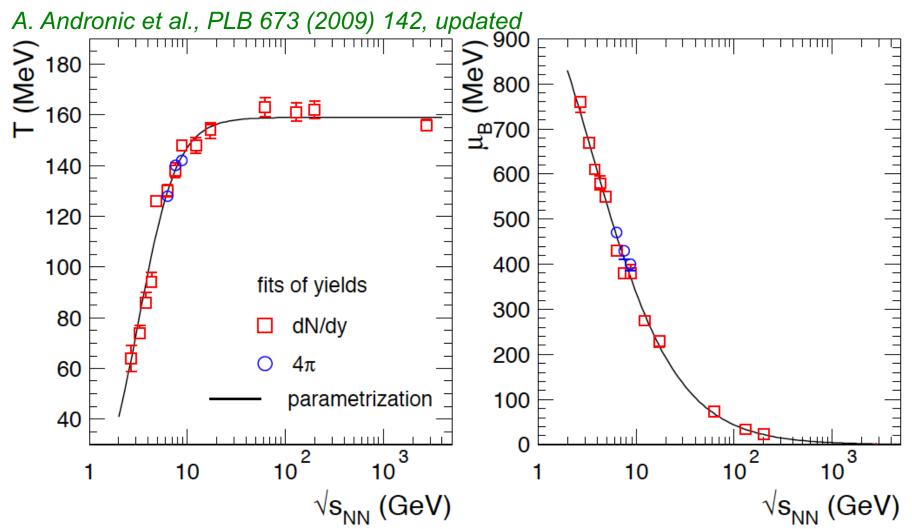
 Statistical (thermal) model with only three parameters able to describe particle yields (grand chanonical ensemble)



- chemical freezeout temperature T_{ch}
- baryo-chemical potential μ_B
- Volume V
- → Using particle yields as input to extract parameters



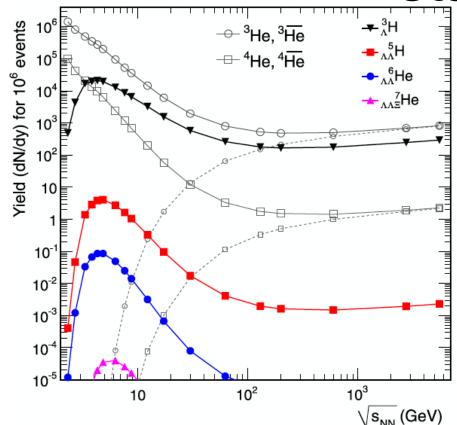
Energy dependence



Thermal model fits show limiting temperature: $T_{lim} = (159 \pm 2) \text{ MeV}$

Predicting yields of bound

states



A. Andronic et al., PLB 697 (2011) 203

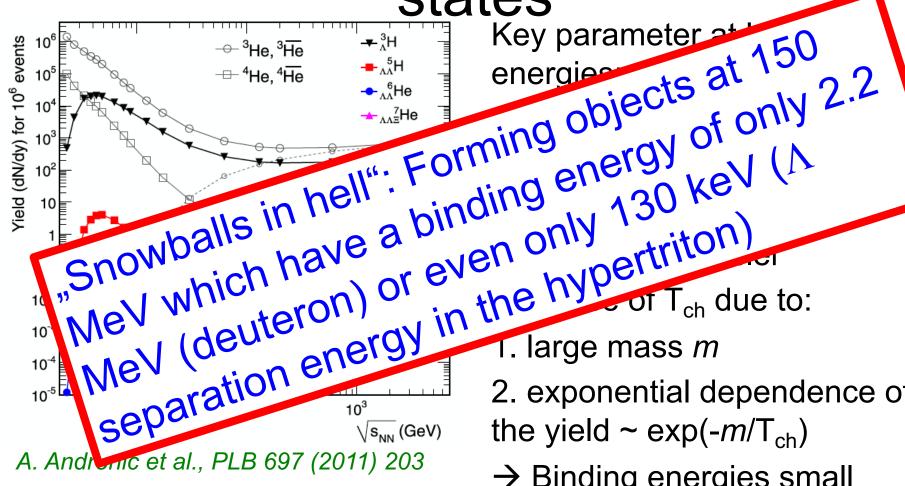
Key parameter at LHC energies:

chemical freeze-out temperature T_{ch}

Strong sensitivity of abundance of nuclei to choice of T_{ch} due to:

- 1. large mass m
- 2. exponential dependence of the yield $\sim \exp(-m/T_{ch})$
- → Binding energies small compared to T_{ch}

Predicting yields of bound states



- 2. exponential dependence of the yield $\sim \exp(-m/T_{ch})$
- → Binding energies small compared to T_{ch}



Snowballs in hell

Recipe for a Hot Universe

 Already mentioned in S. Weinberg's book: The first three minutes quantities from the conditions for thermal equilibrium - for instance, we can make the density of water molecules plus hydrogen ions a little greater or less than 3.3 X 10²² molecules per cubic centimetre by raising or lowering the pressure - so we need to specify them in order to know what is in our glass.

This example also helps us to understand the shifting meaning of what we call 'conserved' quantities. For instance, if our water is at a temperature of millions of degrees, as inside a star, then it is very easy for molecules or ions to dissociate, and for the constituent atoms to lose their electrons. The conserved quantities are then the numbers of electrons and of oxygen and hydrogen nuclei. The density of water molecules plus hydroxyl atoms under these conditions has to be *calculated* from the rules of statistical mechanics rather than specified in advance; of course it turns out to be quite small. (Snowballs are rare in hell.) Actually, nuclear reactions do occur under these condition, so even the numbers of nuclei of each species are not absolutely fixed, but these numbers change so slowly that a star can be regarded as evolving gradually from one equilibrium state to another.



AGS predictions

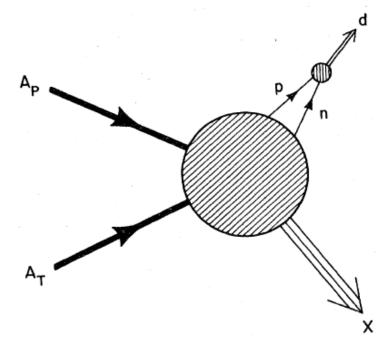
- Already at AGS the predictions from thermal model and coalescence were very close
- Looking at production yields alone might not be enough to disentangle models

P. Braun-Munzinger, J. Stachel, J. Phys. G: Nucl. Part. Phys. 21 (1995) L17

	Thermal model		
Particles	T=0.120 GeV	T = 0.140 GeV	Coalescence model
d	15	19	11.7
t+ ³ He	1.5	3.0	0.8
α	0.02	0.067	810.0
H(I	0.09	0.15	0.07
5 ΛΛΗ	3.5×10^{-5}	2.3×10^{-4}	4×10^{-4}
⁶ ΛΛ He	7.2×10^{-7}	7.6×10^{-6}	1.6×10^{-5}
⁷ Ξ ⁰ ΛΛHe	4.0×10^{-10}	9.6×10^{-9}	4×10^{-8}
10St-8	1.6×10^{-14}	7.3×10^{-13}	
2St-9	1.6×10^{-17}	1.7×10^{-15}	
14St-11	6.2×10^{-21}	1.4×10^{-18}	
16St-13	2.4×10^{-24}	1.2×10^{-21}	
²⁰ St ⁻¹⁶	9.6×10^{-3i}	2.3×10^{-27}	



Coalescence



J. I. Kapusta, PRC 21, 1301 (1980)

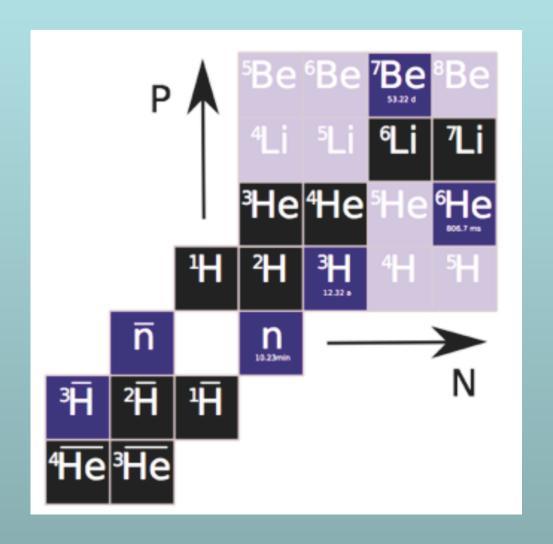
Nuclei are formed by protons and neutrons which are nearby and have similar velocities (after kinetic freezeout)

Produced nuclei

- → can break apart
- → created again by final-state coalescence



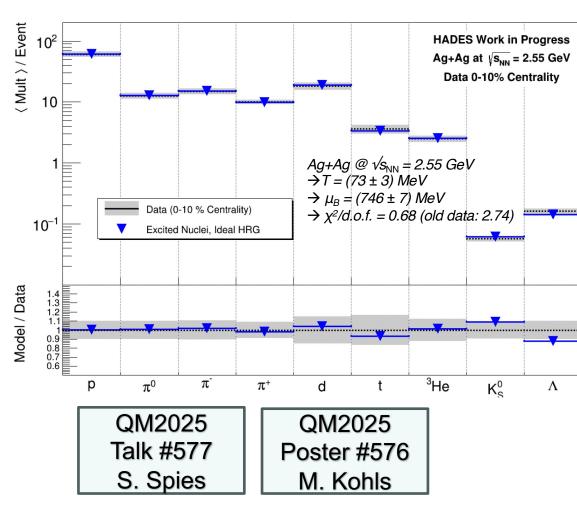
(Anti-)Nuclei





Thermal model: HADES

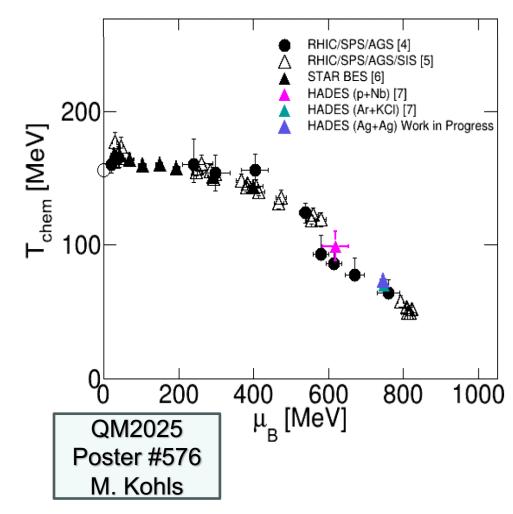
- For HADES
 Ag+Ag data the
 thermal model fit
 describes the
 production of all
 investigated
 yields, including
 nuclei
- Rather good description





Phase diagram

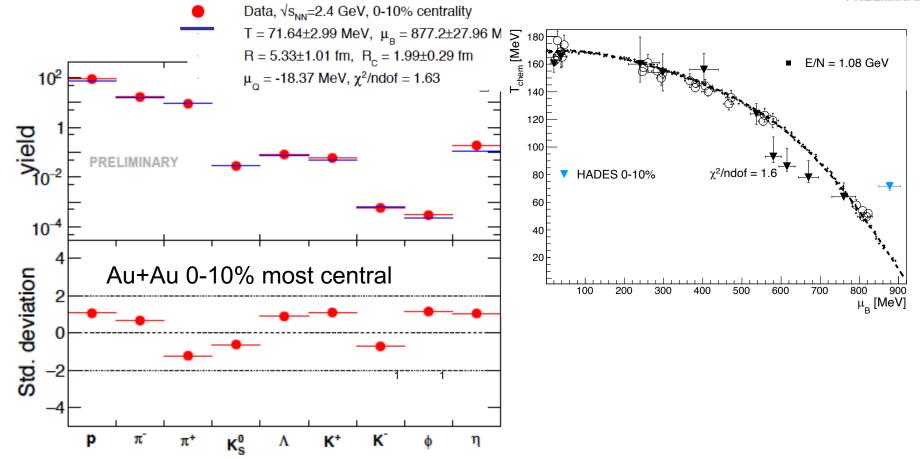
- For HADES Ag+Ag
 data the thermal
 model fit describes
 the production of all
 investigated yields,
 including nuclei
- Rather good description
- Including nuclei leads to nice agreements with the trend of the world data





Phase diagram: HADES

PRELIMINARY

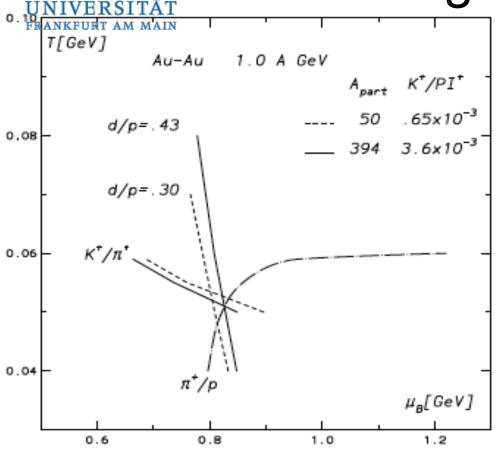


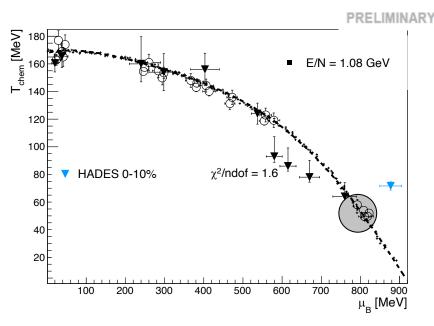
– Freeze-out point stays at higher T and μ_B also for 0-10% most central events

M. Lorenz 2019

GOETHE

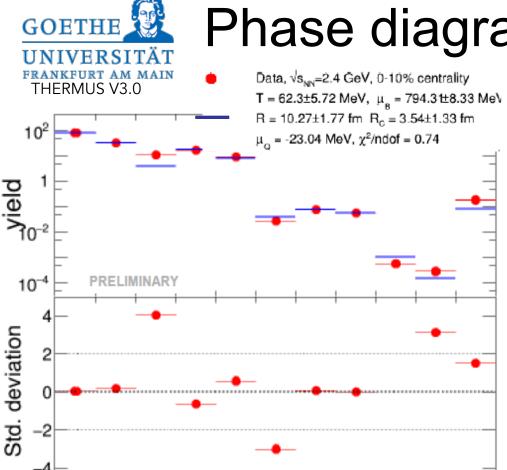
Phase diagram: HADES



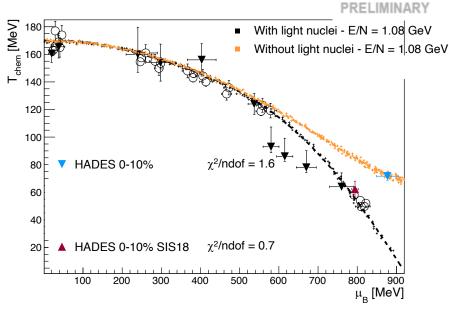


- J. Cleymans, H. Oeschler, K. Redlich, Phys.Rev. C59 (1999))
- Freeze-out points previously estimated based on ratios of p, d, K^+ , π^+
- Light nuclei are not included in Thermus V2.3 Switch to Thermus V3.0 or Therma Fist https://github.com/vlvovch/Thermal-FIST HD 2025 - Benjamin Dönigus

M. Lorenz 2019



Phase diagram: HADES



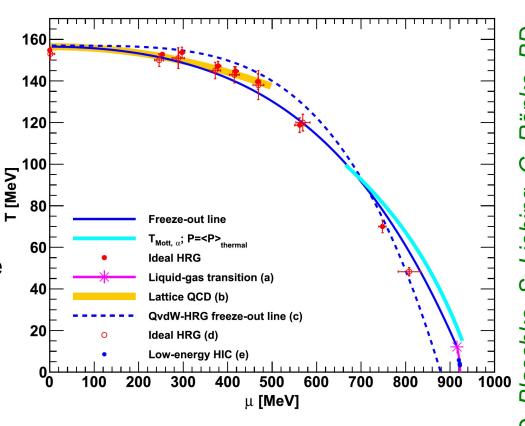
- Fit to HADES data consistent with previous works when same hadron yields are used
- E/N=1.08 GeV with or without light nuclei?
- Light nuclei are important to define chemical freeze-out line at high μ_B .
- J. Cleymans, H. Oeschler, K. Redlich, Phys.Rev. C59 (1999))
- Freeze-out points previously estimated based on ratios of p, d, K^+ , π^+
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M. Lorenz 2019



Phase diagram: Mott α

- If one maps the energy dependence onto a μ_B-T plane one gets experimental access to the QCD phase diagram
- Using parametrizations of the thermal model fits one can get an approximation of the freeze-out curves
- Particular models can even describe their shape, e.g. Mott curves

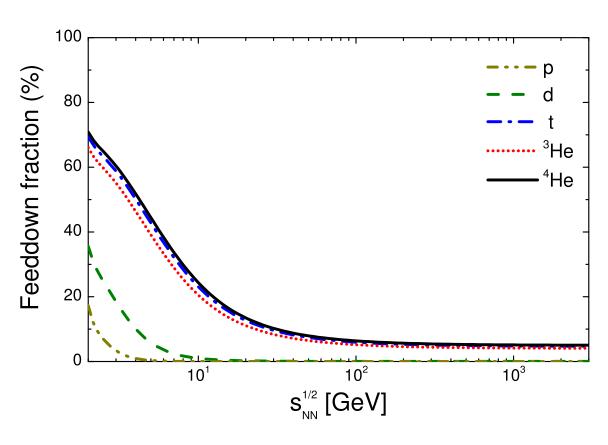


Poster #1096 D. Blaschke et al.



Thermal model

- For the thermal model description of production yields, feed-down is an important ingredient
- All light hadron production yields are populated strongly by resonances
- Seems to not be the case for (hyper-) nuclei at the LHC

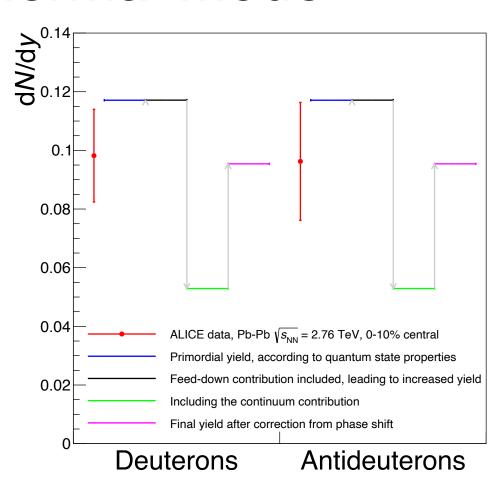


V. Vovchenko, BD, B. Kardan, M. Lorenz, H. Stoecker, Phys.Lett.B 809 (2020) 135746



Thermal model

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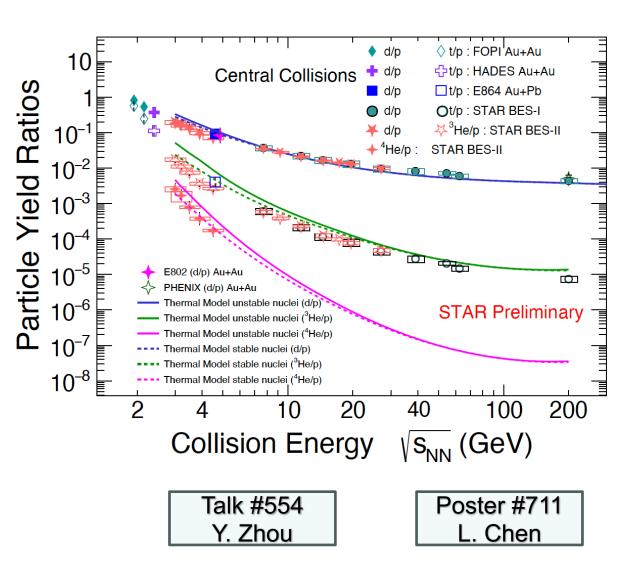


BD, G. Röpke, D. Blaschke, Phys. Rev. C 106 (2022) 044908



UNIVERSITÄT Light nuclei: yield ratios

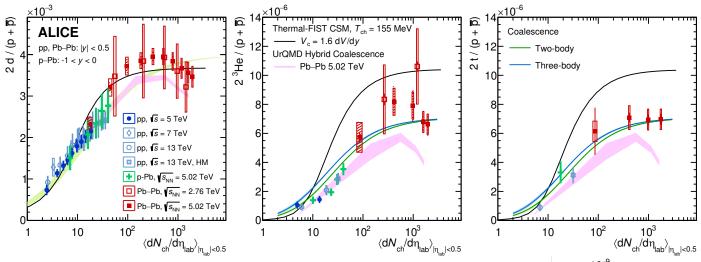
A/p ratio:
 better
 described by
 thermal model
 with only
 stable nuclei



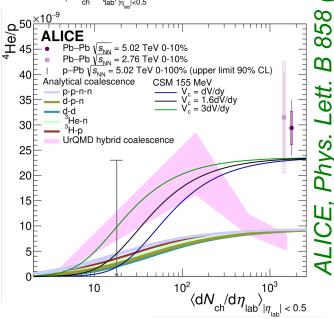


Phys.Rev.C 107 (2023) 064904

Light nuclei: model description



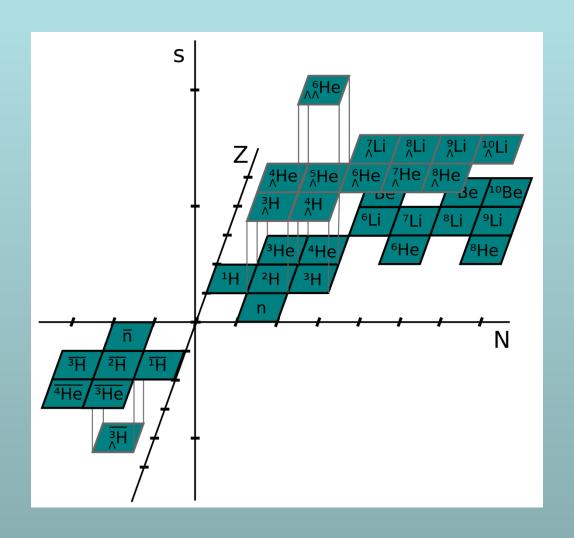
- Model comparisons in ratios of A/p
- d/p possible to describe by coalescence and thermal models
- ³He seems to prefer coalescence, ³H prefers coalescence
- ⁴He in contrary clearly prefers themal production



24

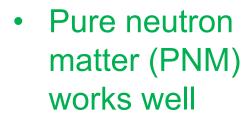


Hypernuclei



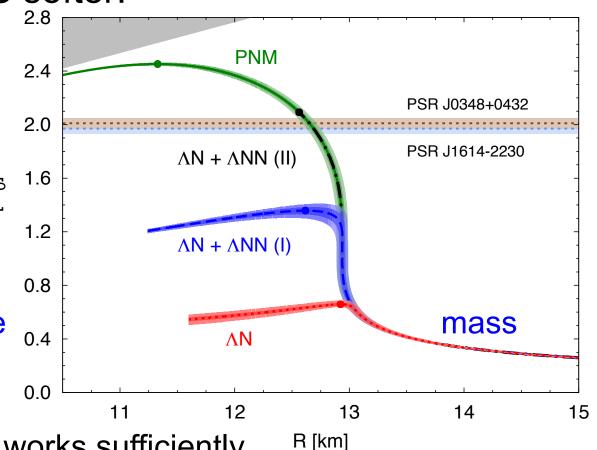
Neutron stars and interactions

 Hyperon puzzle in neutron stars → hyperons make the EOS softer:



- Known ∧N interaction
 → way to soft
- Including ΛNN forces brings the mass slightly up

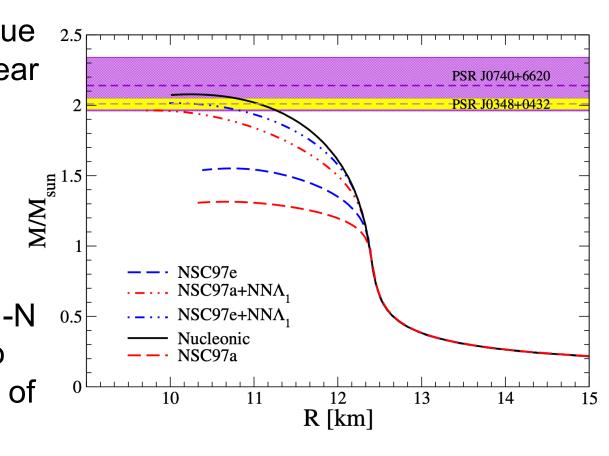
• Only additional 11 12 12 ANN interaction works sufficiently





Hypernuclei

- Hypernuclei are unique probes to study nuclear structure
- Single Λ-hypernuclei are major source of extracting Λ-N interaction
- Correct Λ-N and Λ-N-N interaction needed to understand structure of neutron stars



D. Logoteta et al., Astron. Astrophys. 646 (2021) A55



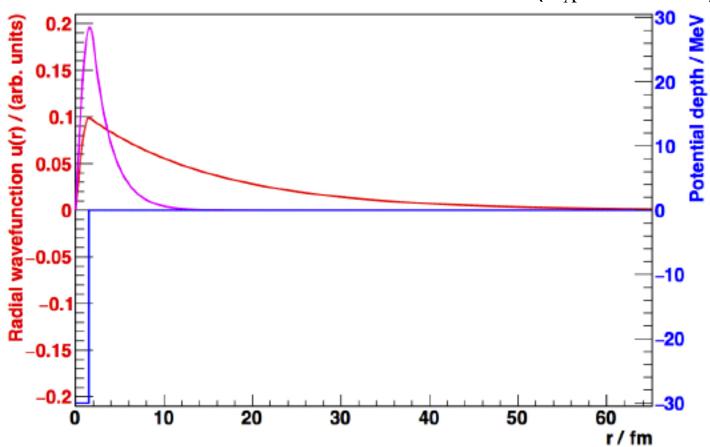
Hypertriton

Bound state of Λ , p, n $m = 2.991 \text{ GeV}/c^2 (B_{\Lambda} = 130 \text{ keV})$



Hypertriton

Bound state of Λ , p, n $m = 2.991 \text{ GeV}/c^2 (B_{\Lambda} = 130 \text{ keV})$

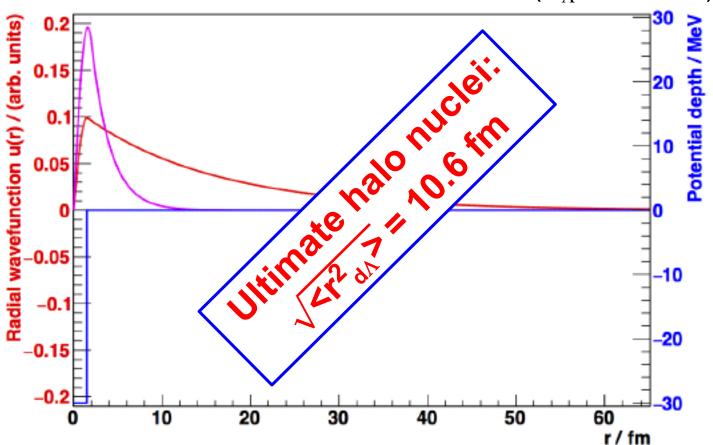


P. Braun-Munzinger, BD, Nucl. Phys. A 987 (2019) 144



Hypertriton

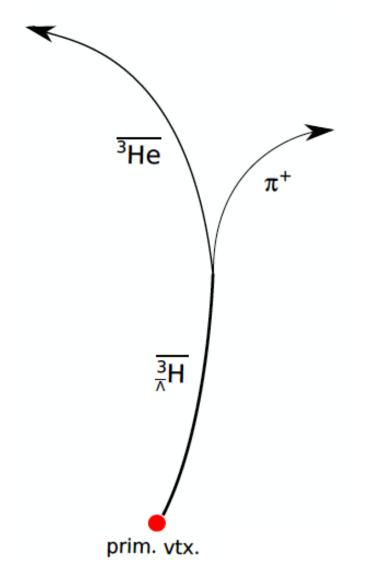
Bound state of Λ , p, n $m = 2.991 \text{ GeV}/c^2 (B_{\Lambda} = 130 \text{ keV})$



P. Braun-Munzinger, BD, Nucl. Phys. A 987 (2019) 144



UNIVERSITÄT Hypertriton Identification



Bound state of Λ , p, n $m = 2.991 \text{ GeV}/c^2 (B_{\Lambda} = 130 \text{ keV})$

→ Radius of about 10.6 fm Decay modes:

$$^{3}_{\Lambda}H \rightarrow^{3}He + \pi^{-}$$
 $^{3}_{\Lambda}H \rightarrow^{3}H + \pi^{0}$
 $^{3}_{\Lambda}H \rightarrow d + p + \pi^{-}$
 $^{3}_{\Lambda}H \rightarrow d + n + \pi^{0}$

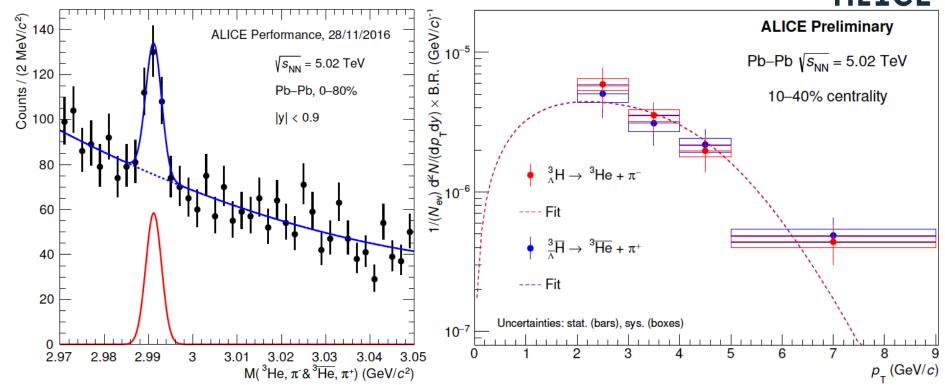
- + anti-particles
- → Anti-Hypertriton first observed by STAR Collaboration:

Science 328,58 (2010)



Hypertriton signal



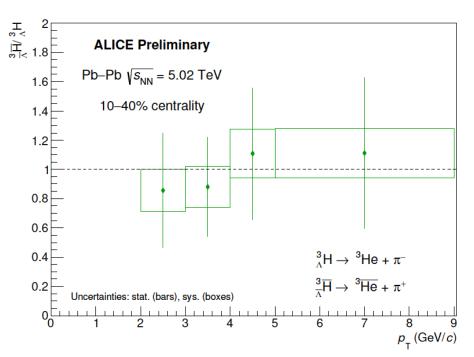


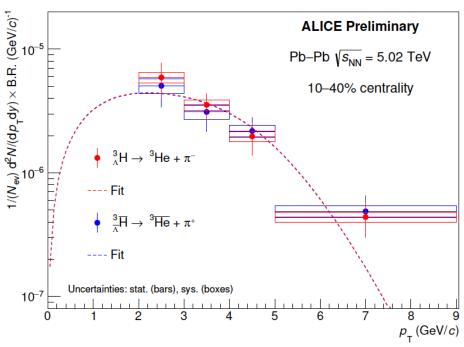
- Clear signal reconstructed by decay products
- Spectra can also be described by Blast-Wave model
 - → Hypertriton flows as all other particles



Hypertriton spectra







• Anti-hypertriton/Hypertriton ratio consistent with unity vs. p_{T}



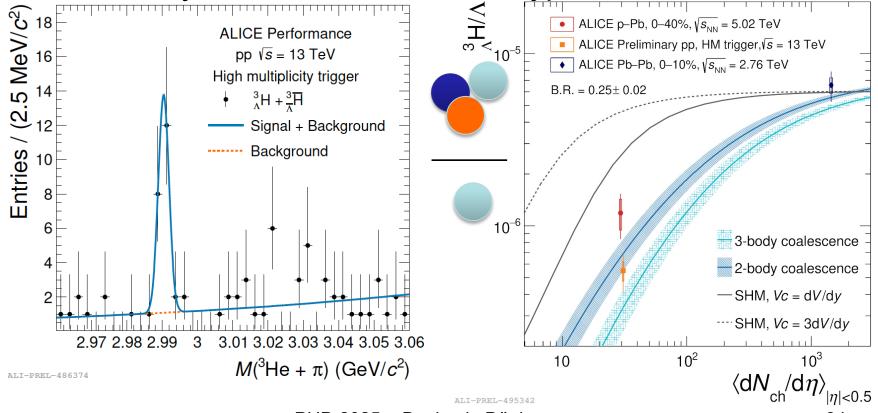
Hypertriton in pp & p-Pb



Hypertriton signal recently also extracted in pp and p-Pb collisions

Stronger separation between models as for other particle

ratios, mainly due to the size of the hypertriton



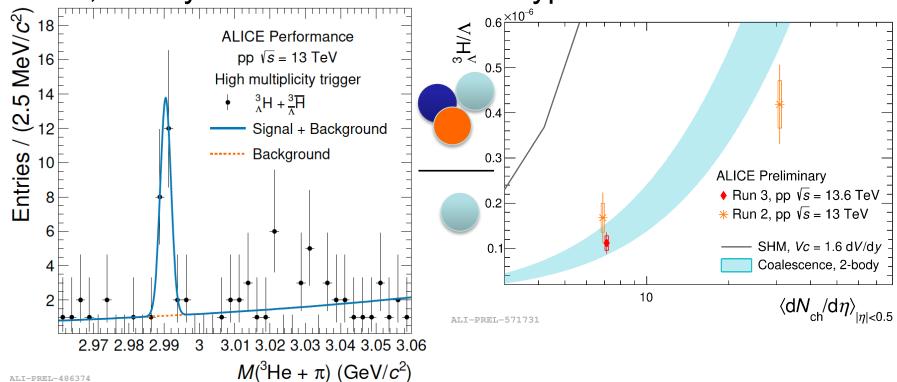


ALI-PREL-486374

Hypertriton in pp & p-Pb



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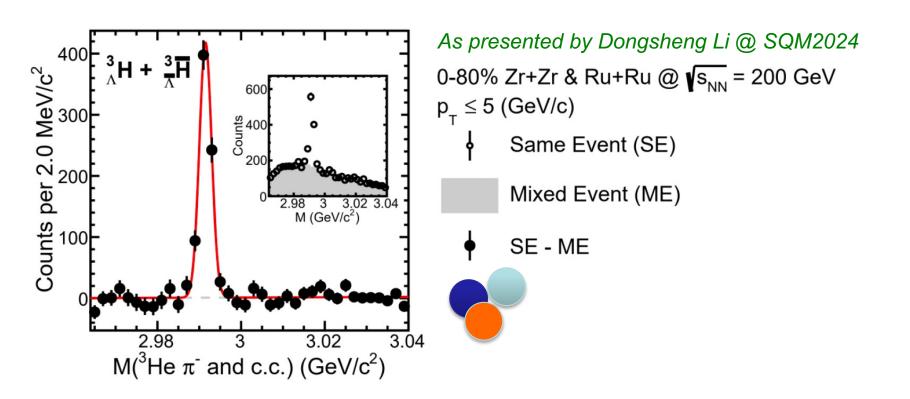




Hypertriton at RHIC



Hypertriton signal recently also extracted in isobar collisions

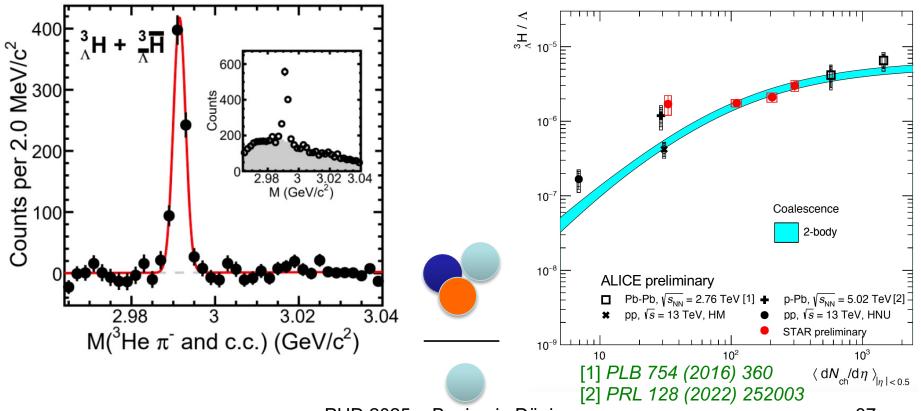




Hypertriton/Λ ratio



- Hypertriton signal recently also extracted in isobar collisions
- Stronger separation between models as for other particle ratios, mainly due to the size of the hypertriton

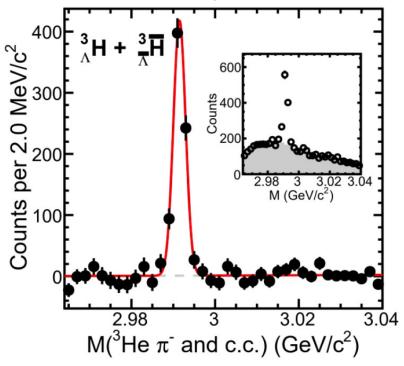


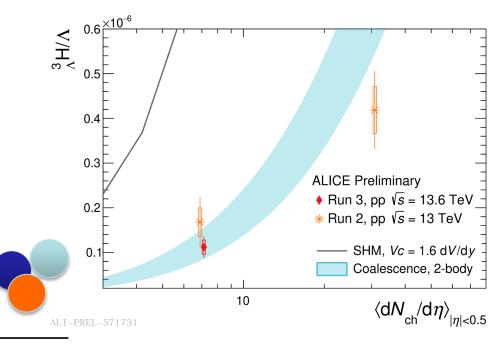


Hypertriton/∧ ratio



- Hypertriton signal recently also extracted in isobar collisions
- Stronger separation between models as for other particle ratios, mainly due to the size of the hypertriton

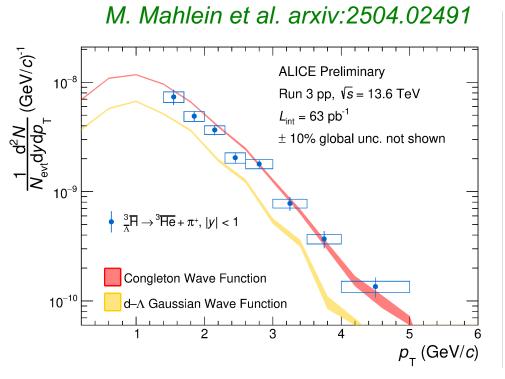






Hypertriton: small systems

- Strong suppression in small systems due to large size of the object compared to the system size
- Coalescence model favoured over canonical statistical model



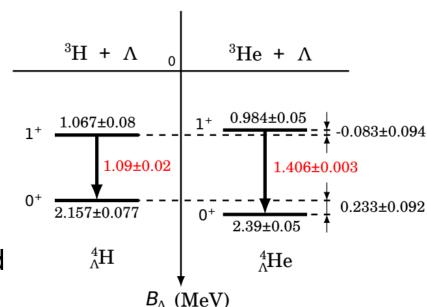




- Large suppression expected for A
 = 4 hypernuclei by the SHM wrt A =
 3
- A = 4 hypernuclei are more bound and each has an excited state

Phys. Rev. Lett. 115, 222501 (2015)

- The yields of these hypernuclei are enhanced with respect to the ground state due to the feed-down from excited states
- Also the yields of the SHM scale with the spin-degeneracy
- Resulting in a total enhancement of a factor 4 for both hypernuclei



M. Schäfer, N. Barnea, A. Gal, Phys.Rev.C 106, L031001 (2022)

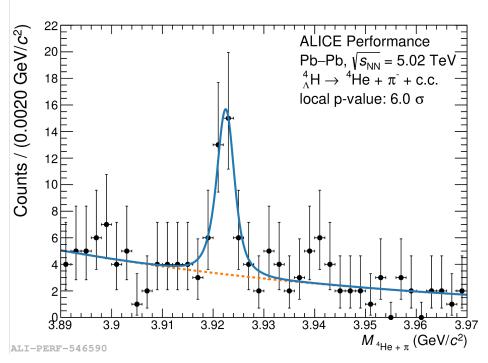




- For the first time, we are able to reconstruct A = 4 (anti)hypernuclei at the LHC and determine their production yield
- (Anti)hyperhydrogen-4 invariant-mass spectrum in Run 2 Pb-Pb collisions at 5.02 TeV
- Examined in the two-body decay:

$$^{4}_{\Lambda}\text{H} \rightarrow ^{4}\text{He} + \pi^{-} + \text{c.c.}$$

Reaching a local p-value of 6σ





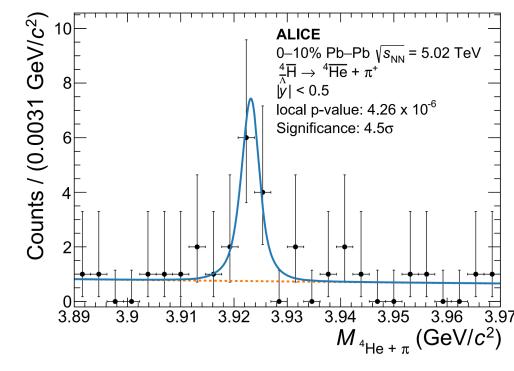


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$$^{4}_{\Lambda}\text{He} \rightarrow ^{3}\text{He} + p + \pi^{-} + \text{c.c.}$$

Reaching a significance of 4.5σ

ALICE Collaboration, arXiv:2410.17769



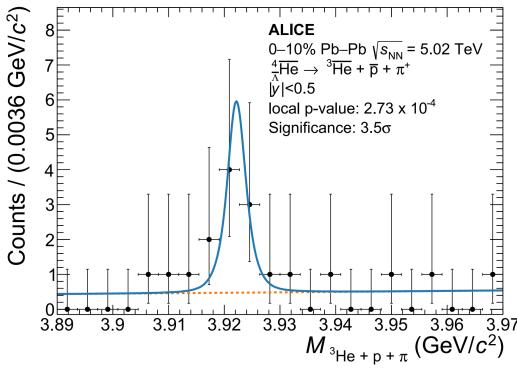




- For the first time, we are able to reconstruct A = 4 (anti)hypernuclei at the LHC and determine their production yield
- First observation
 of the antihyperhelium-4
 in Run 2 Pb-Pb collisions at
 5.02 TeV

 Reaching a significance of 3.5σ

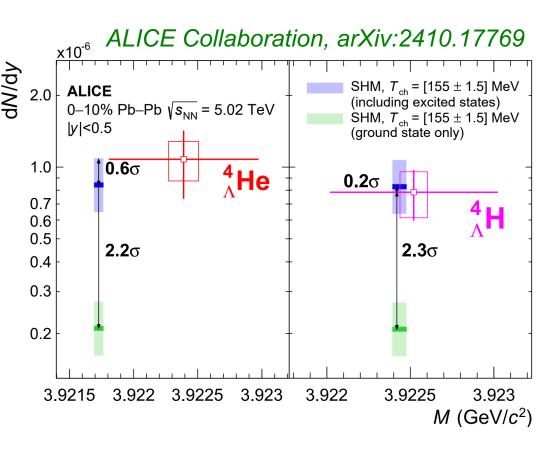








- First measurement of the (anti)hyperhelium-4 production yield
- Testing the dependence of ⇒ the yields of the SHM with ≤ the spin-degeneracy
- Our yields confirm the hypothesis of excited states for both (anti)hypernuclei within 2σ
- currently dominated by statistical uncertainties
- with more data, a high precision measurement will be feasible (like for the Λ hyperon)

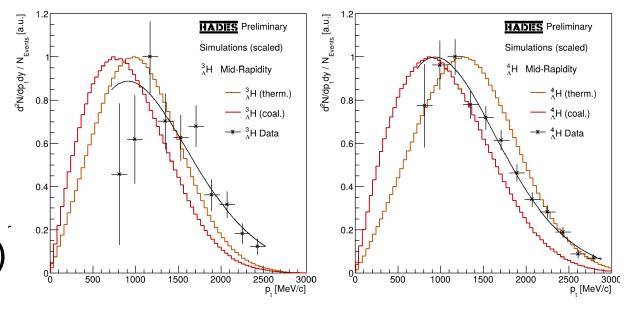




Hypernuclei: HADES

- Lifetimes of investigated hypernuclei consistent with world data
- p_T and y
 distributions
 hint to (thermal)
 production

QM2025 Talk #577 S. Spies

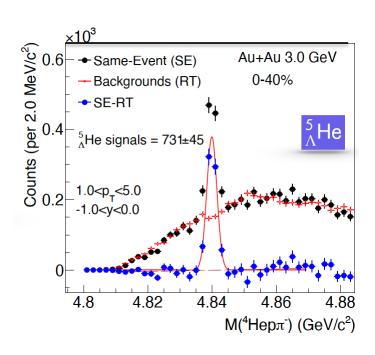




Hypernucleus: ⁵ΛHe



 Highlight: clear signal of hyperhelium-5 in Au+Au at 3.0 GeV



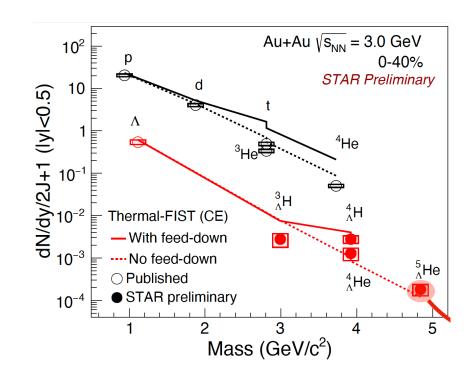
QM2025 Talk #554 Y. Zhou



Hypernucleus: $^{5}_{\Lambda}$ He



- Highlight: clear signal of hyperhelium-5 in Au+Au at 3.0 GeV
- Comparison to thermal model seems to work only up to A=2, even when excited states are considered



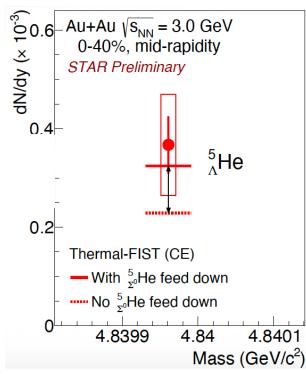
QM2025 Talk #554 Y. Zhou



Hypernucleus: ⁵ΛHe



- Highlight: clear signal of hyperhelium-5 in Au+Au at 3.0 GeV
- Comparison to thermal model seems to work only up to A=2, even when excited states are considered
- Hypothesis of A=5 Σhypernucleus would help explaining the data



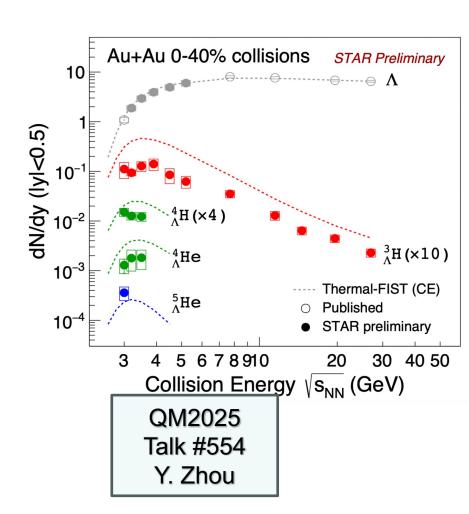
QM2025 Talk #554 Y. Zhou



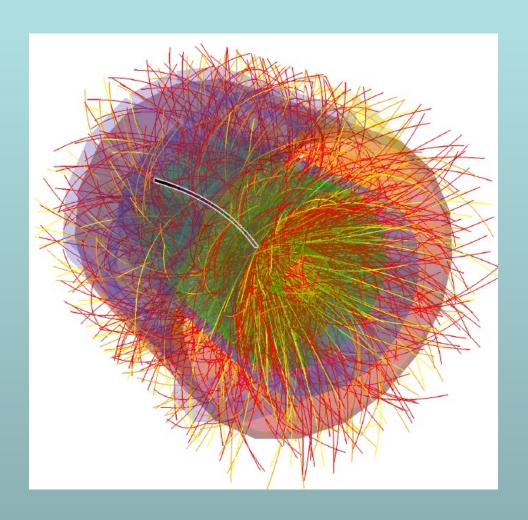
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- Comparison to thermal model seems to work only up to A=2, even when excited states are considered
- Hypothesis of A=5 Σhypernucleus would help explaining the data
- Energy dep. shows slight offset for all hypernuclei

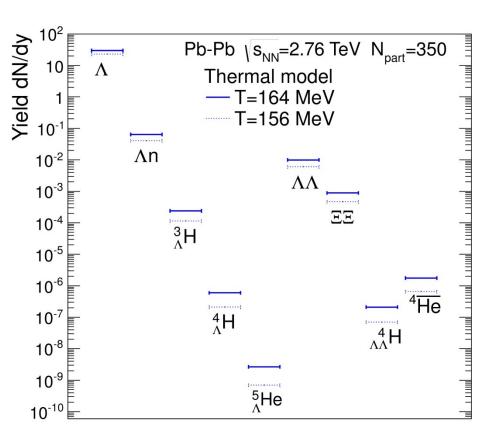


Outlook & Summary





Outlook



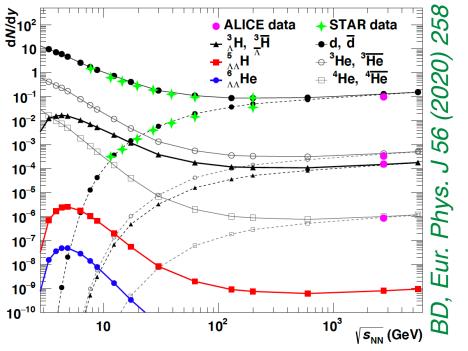
A. Andronic, private communication, model described in A. Andronic et al., PLB 697, 203 (2011) and references therein

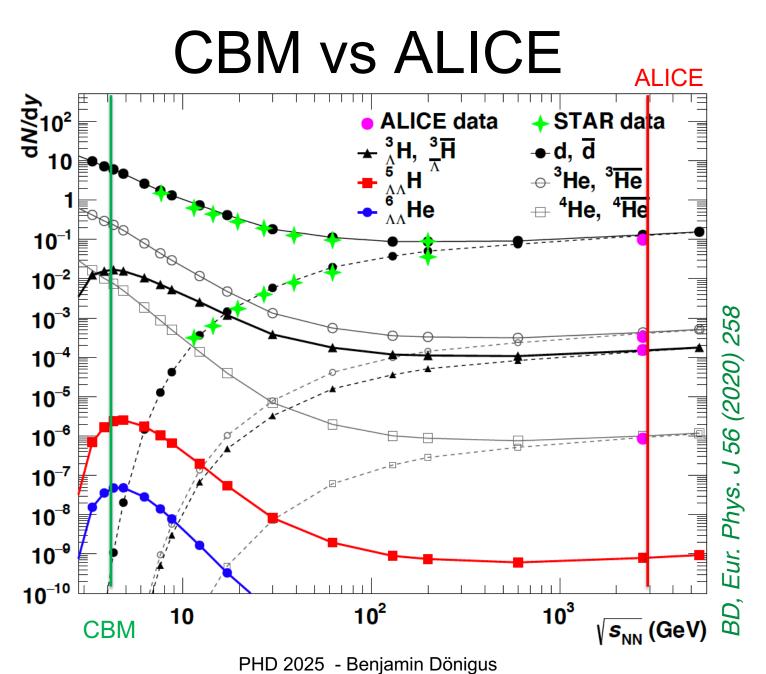
- Explore QCD and QCD inspired model predictions for (unusual) multi-baryon states
- Search for rarely produced anti- and hyper-matter
- Test model predictions, e.g. thermal and coalescence



Conclusion

- ALICE@LHC, HADES@GSI and STAR@RHIC are well suited to study light (anti-)(hyper-) nuclei and perform searches for exotic bound states (A<5)
- Copious production of loosely bound objects measured by ALICE, HADES and STAR as predicted by the thermal model
- Models describe the data rather well
- Ratios vs. multiplicity trend described by both models
 only tension: Alpha vs. ³_^H
- New and more precise data can be expected in the next years





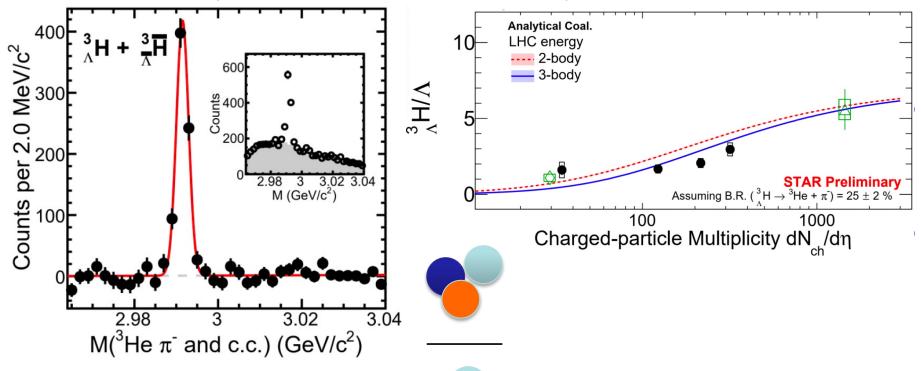
Backup



Hypertriton/Λ ratio



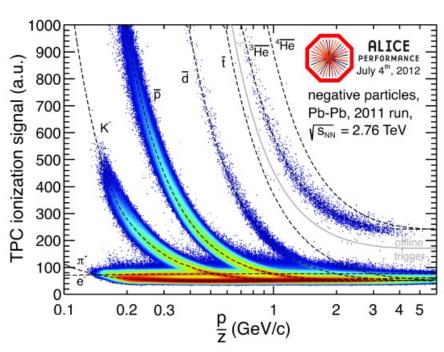
- Hypertriton signal recently also extracted in isobar collisions
- Stronger separation between models as for other particle ratios, mainly due to the size of the hypertriton

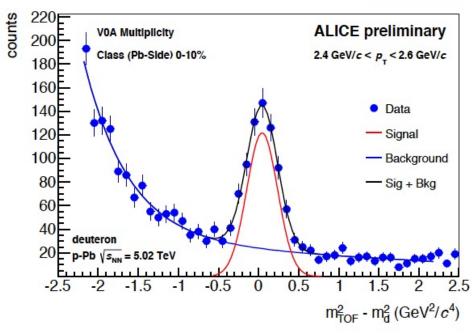




Particle Identification







Low momenta:

Nuclei are identified using the d*E*/d*x* measurement in the Time Projection Chamber (TPC)

Higher momenta:

Velocity measurement with the Time-of-Flight (TOF) detector is used to calculate the m^2 distribution



Anti-Alpha

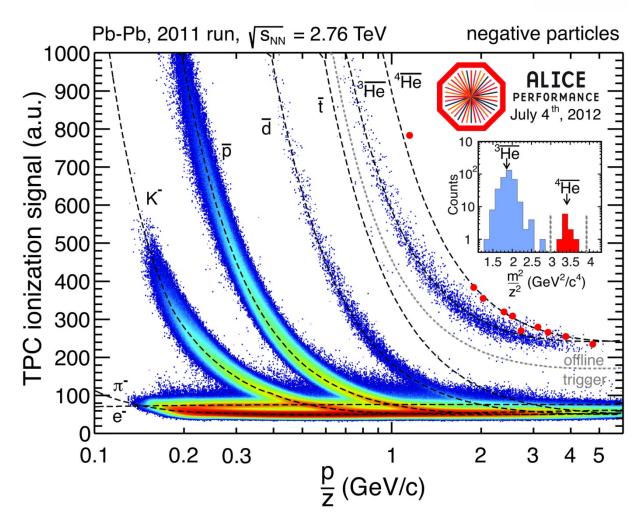




For the full statistics of 2011 ALICE identified 10 Anti-Alphas using TPC and TOF

STAR observed the Anti-Alpha in 2010:

Nature 473, 353 (2011)

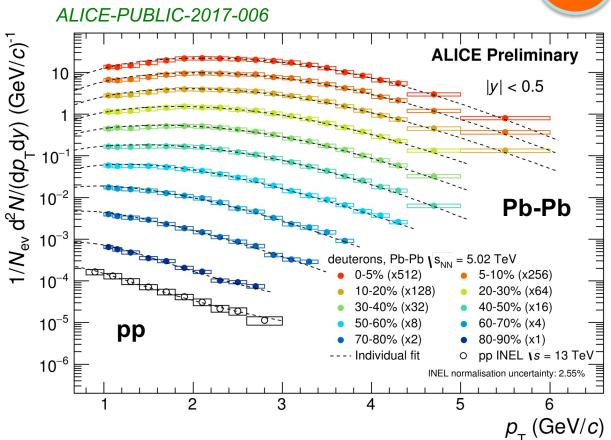




Deuterons





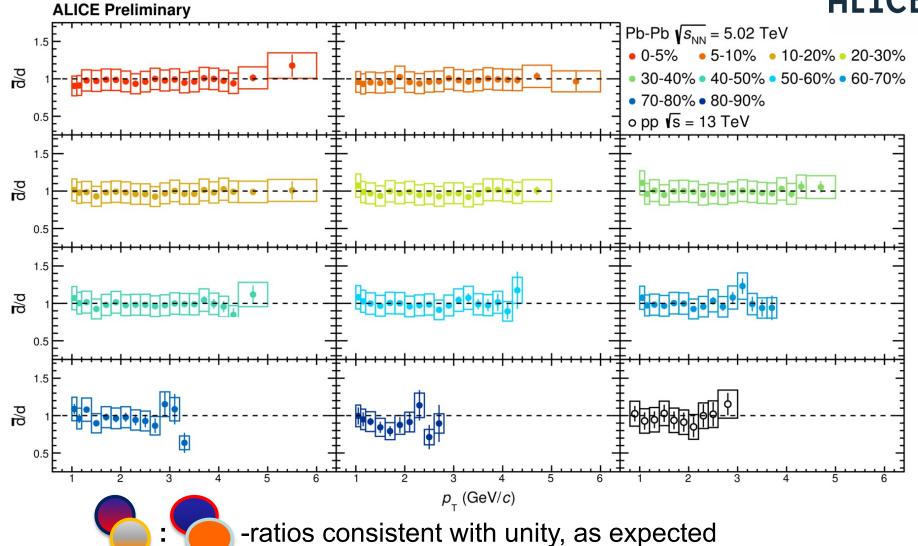


- p_T spectra getting harder for more central collisions (from pp to Pb-Pb) \rightarrow showing clear radial flow
- Blast-Wave fits describe the data in Pb-Pb very well
- No hint for radial flow in pp



(Anti-)Deuteron ratio



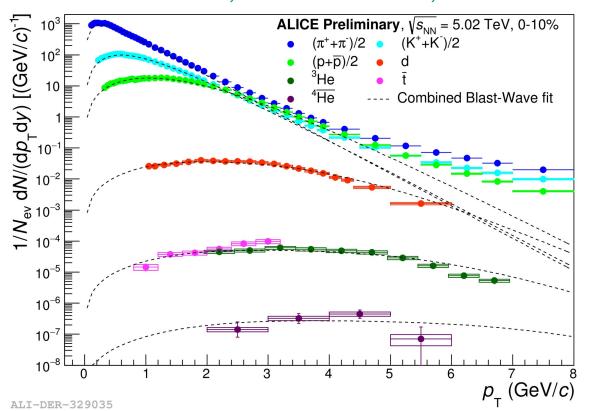




Combined Blast-Wave fit



ALICE Collaboration, arXiv:1910.07678, see also arXiv:2311.11758



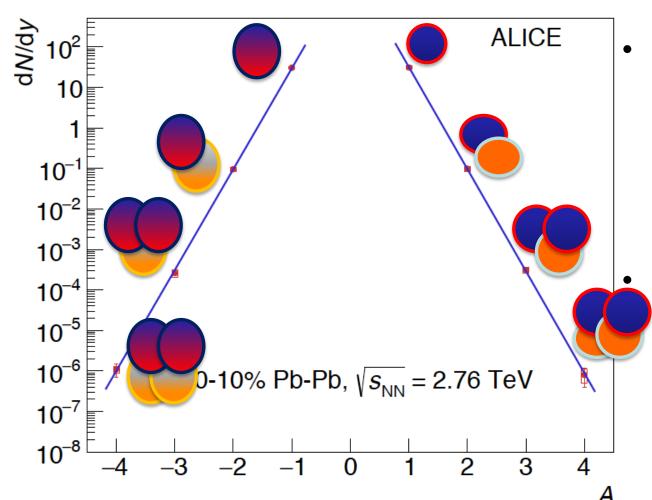
• Simultaneous Blast-Wave fit of π^+ , K⁺, p, d, t, 3 He and 4 He spectra for central Pb-Pb collisions leads to values for $\langle \beta \rangle$ and $T_{\rm kin}$ close to those obtained when only π ,K,p are used

All particles are described rather well with this simultaneous fit



Mass dependence





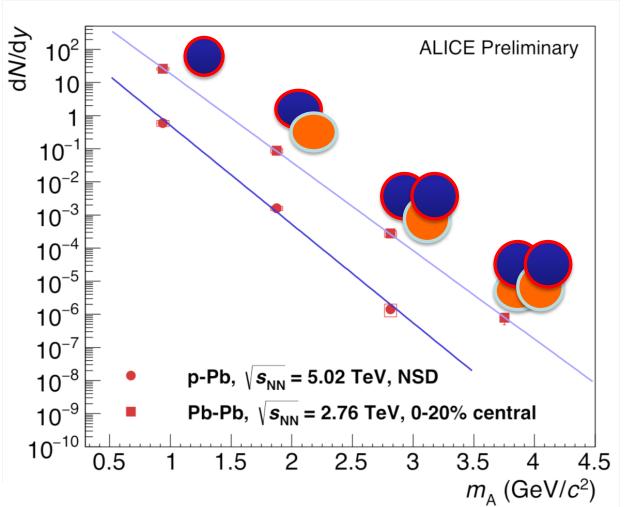
Production of (anti-) nuclei is follwing an exponential, and decreases with mass as expected from thermal model In Pb-Pb the "penalty factor" for each additional baryon ~300 (for particles and antiparticles)

ALICE Collaboration, arXiv:1710.07531, NPA 971, 1 (2018)



Mass dependence



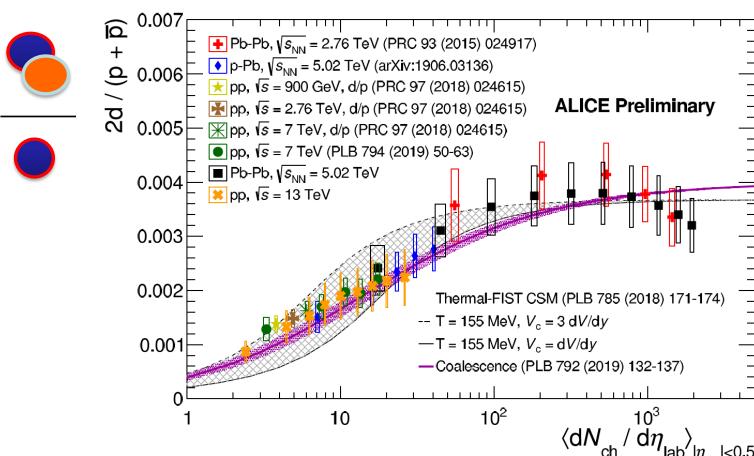


- Production of (anti-)
 nuclei is follwing an
 exponential, and
 decreases with
 mass as expected
 from thermal model
- In Pb-Pb the "penalty factor" for each additional baryon ~300, in p-Pb ~600 and in pp ~1000



d/p vs. multiplicity





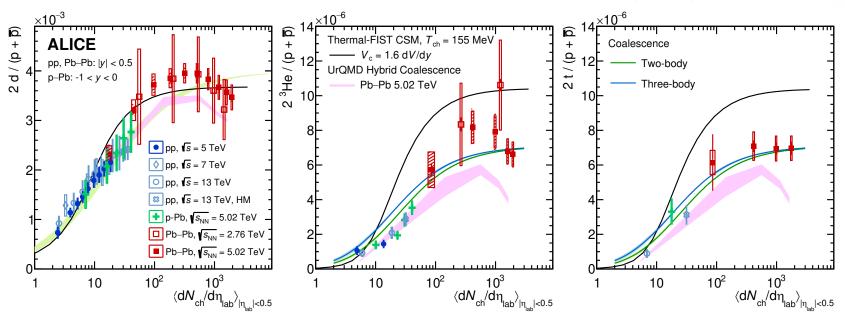
d/p ratio rather well described by coalescence and (canonical) thermal model



Ratios vs. multiplicity



ALICE Collaboration, arXiv:2211.14015, Phys.Rev.C 107 (2023) 064904



- d/p ratio rather well described by coalescence and (canonical) thermal model
- Some tension for 3 He/p and 3 H/p over p_{T}

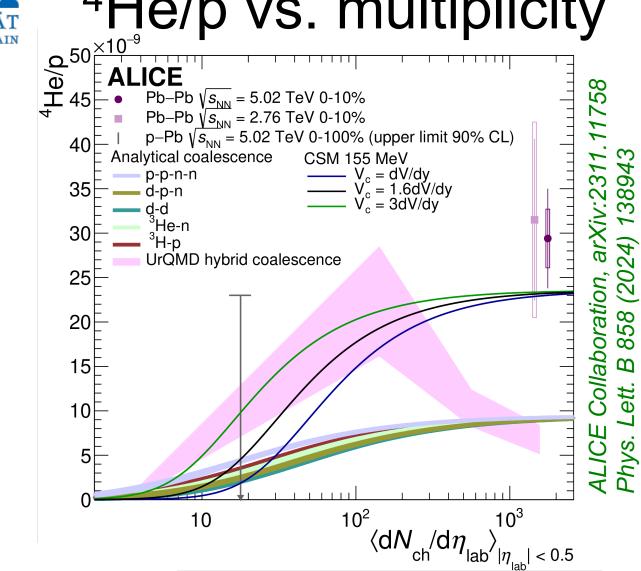






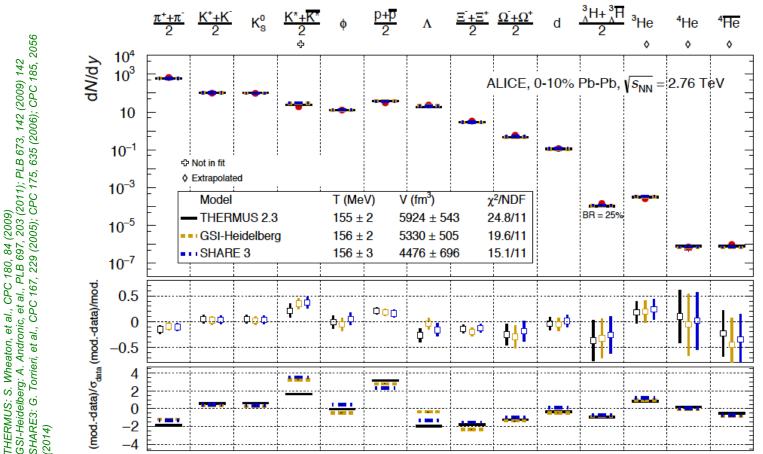
⁴He/p vs. multiplicity





⁴He/p ratio significantly better described by the thermal model

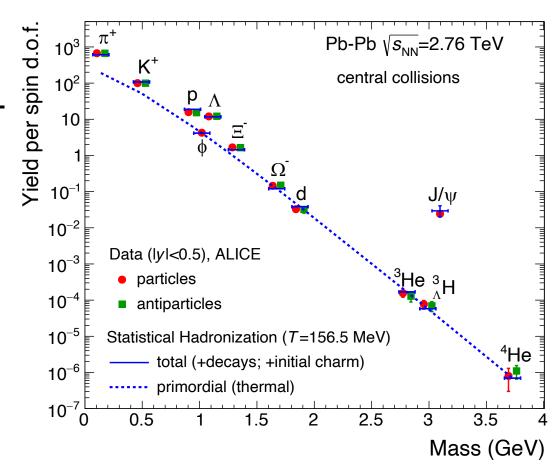




 Different model implementations describe the production probability, including light nuclei and hyper-nuclei, rather well at a temperture of about T_{ch} =156 MeV



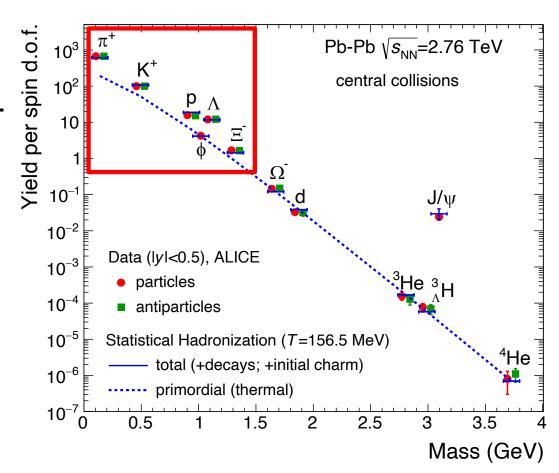
- For the thermal model description of production yields, feeddown is an important ingredient
- All light hadron production yields are populated strongly by resonances
- Seems to not be the case for (hyper-)nuclei



A. Andronic et al., Phys.Lett.B 797 (2019) 134836



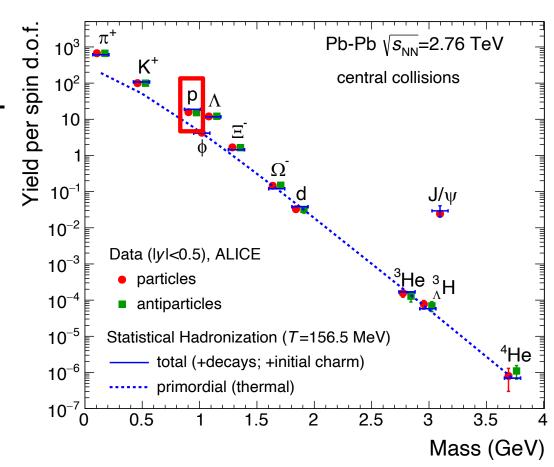
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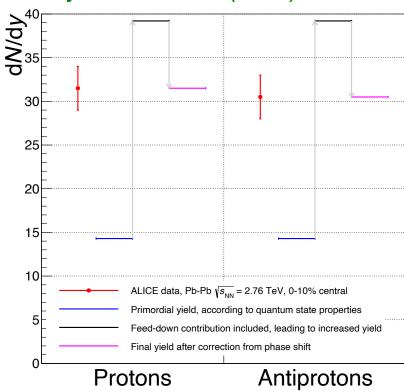


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BD, G. Röpke, D. Blaschke, Phys. Rev. C 106 (2022) 044908

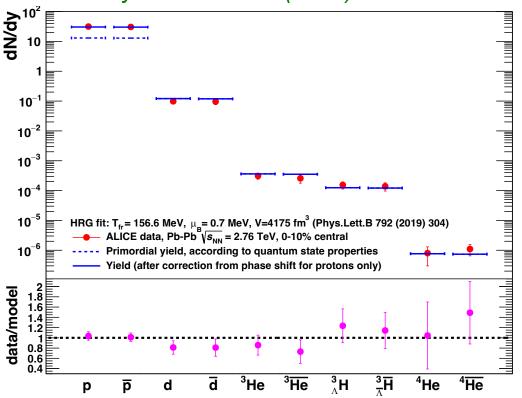


A. Andronic et al., Phys.Lett.B 797 (2019) 134836; Nature 561 (2018) 7723, 321; Phys.Lett.B 697 (2011) 203; Phys.Lett.B 792 (2019) 304



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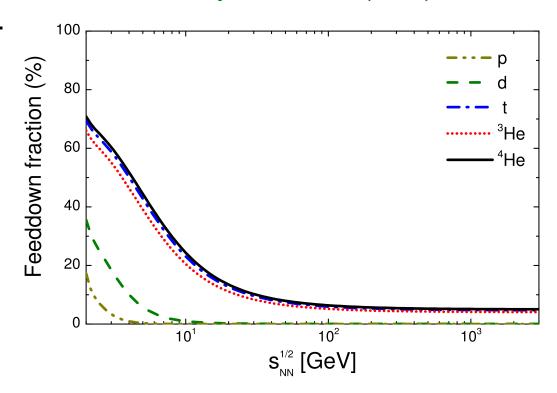


A. Andronic et al., Phys.Lett.B 797 (2019) 134836; Nature 561 (2018) 7723, 321; Phys.Lett.B 697 (2011) 203; Phys.Lett.B 792 (2019) 304



- For the thermal model description of production yields, feeddown is an important ingredient
- All light hadron production yields are populated strongly by resonances
- Seems to not be the case for (hyper-)nuclei at LHC

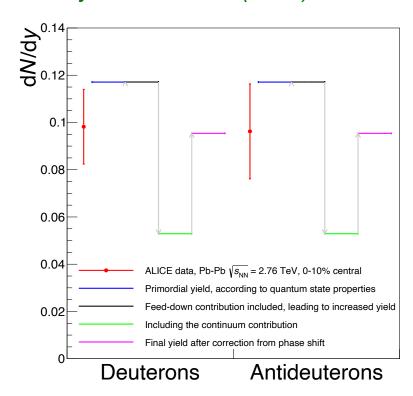
V. Vovchenko, BD, B. Kardan, M. Lorenz, H. Stoecker, Phys.Lett.B 809 (2020) 135746





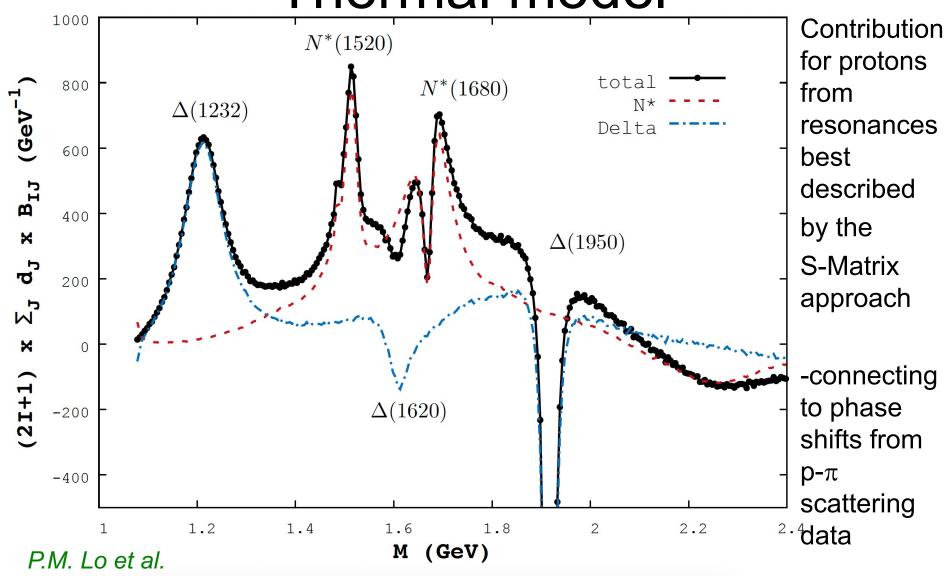
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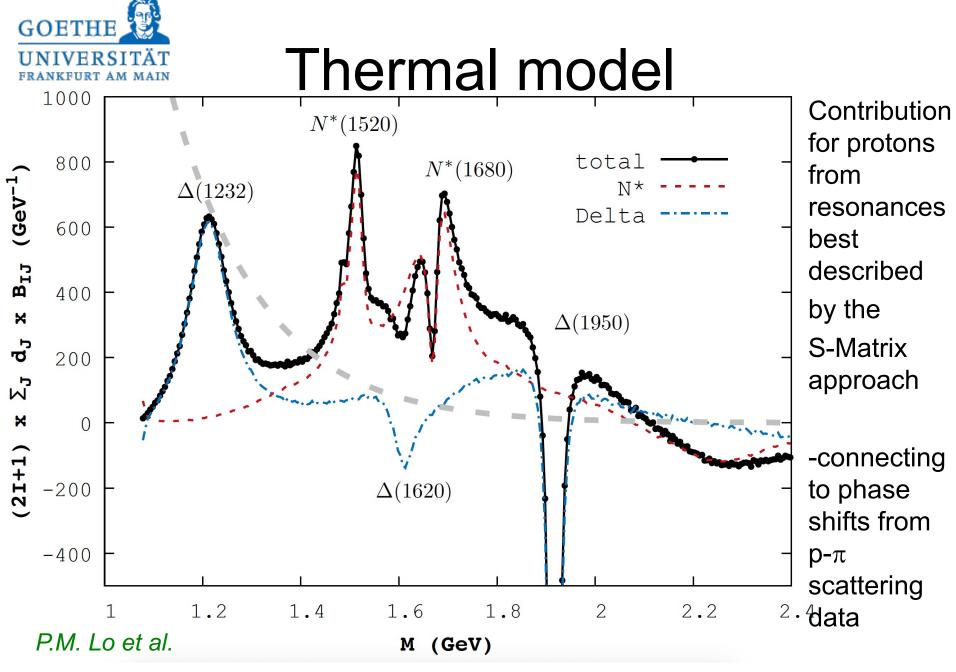
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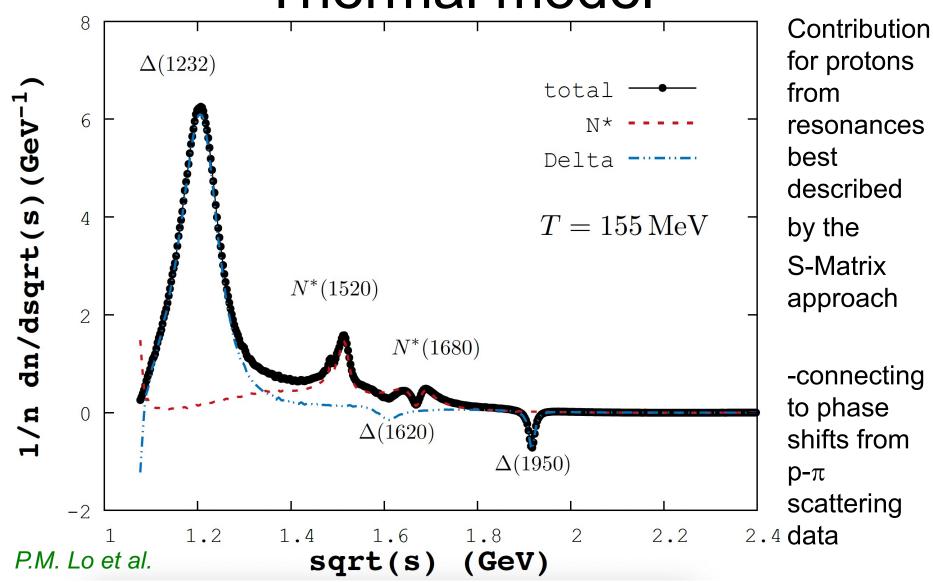
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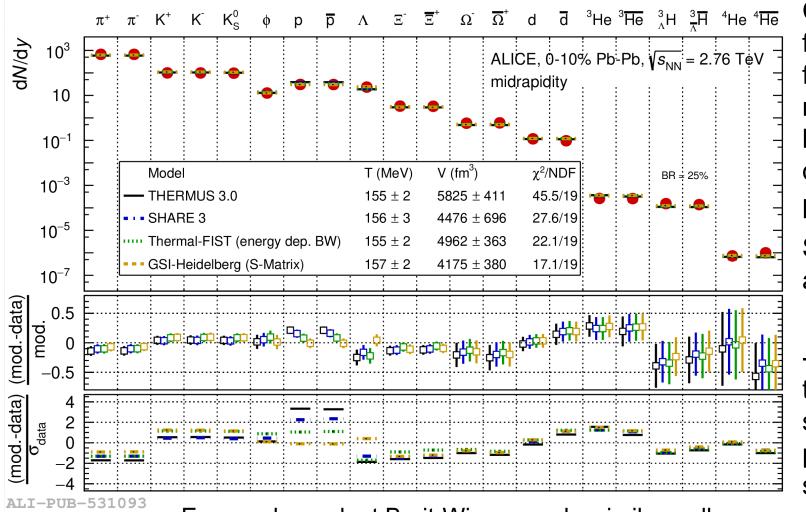












Contribution for protons from resonances best described by the S-Matrix approach

-connecting to phase shifts from p-π scattering data

Energy-dependent Breit-Wigner works similar well

ALICE Collaboration, arXiv:2211.04384 [nucl-ex]

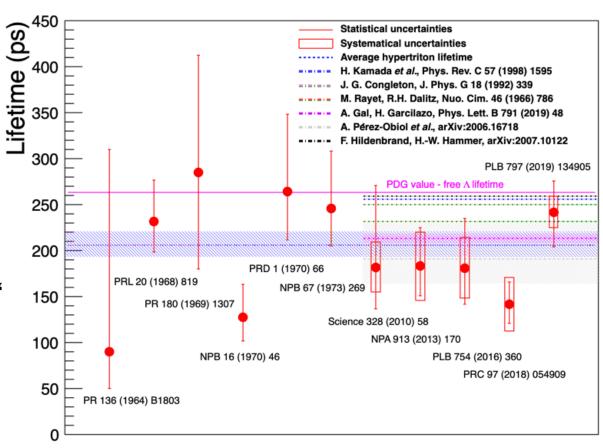


Hypertriton "Puzzle"

• Recently measured lifetimes are significantly below the lifetime of the free Λ \rightarrow new ALICE results agree with the

world average of all known measurements and with the free Λ lifetime

 Most recent calculations include "final-state" interaction and agree well with the data



BD, Eur. Phys. J 56 (2020) 258

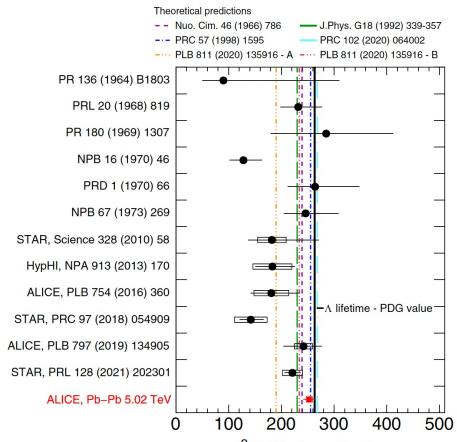


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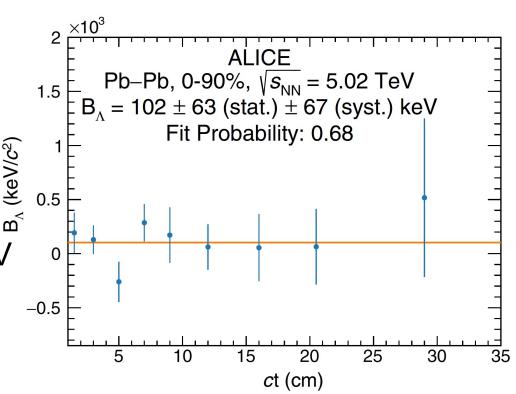
ALICE Collaboration, Phys. Rev. Lett. 131 (2023) 102302 ³ H lifetime (ps)



Binding Energy



- Current studies show a better constraint and smaller statistical uncertainties
- The value obtained by this fit is $B_{\Lambda} = (102 \pm 63 \pm 67) \text{ keV}$
- Is compatible with the theoretical predictions



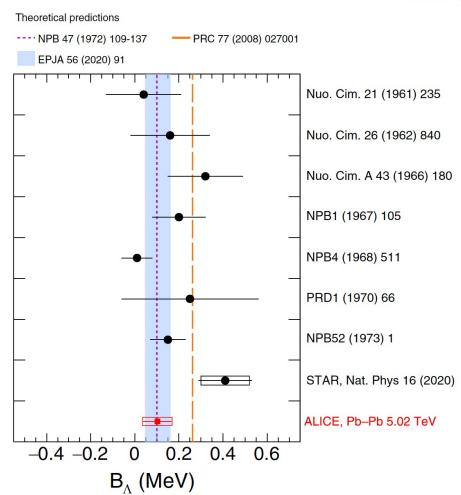
ALICE Collaboration, Phys. Rev. Lett. 131 (2023) 102302



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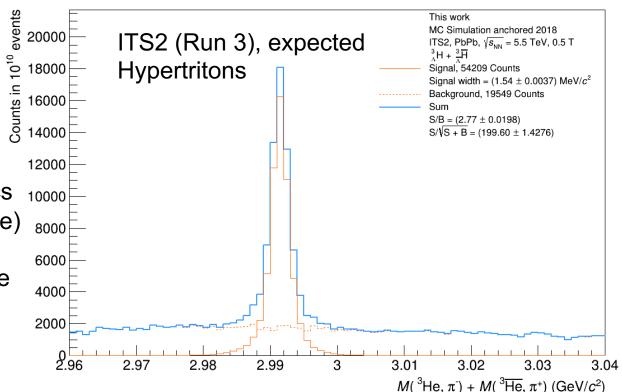
ALICE Collaboration, Phys. Rev. Lett. 131 (2023) 102302



Expectations



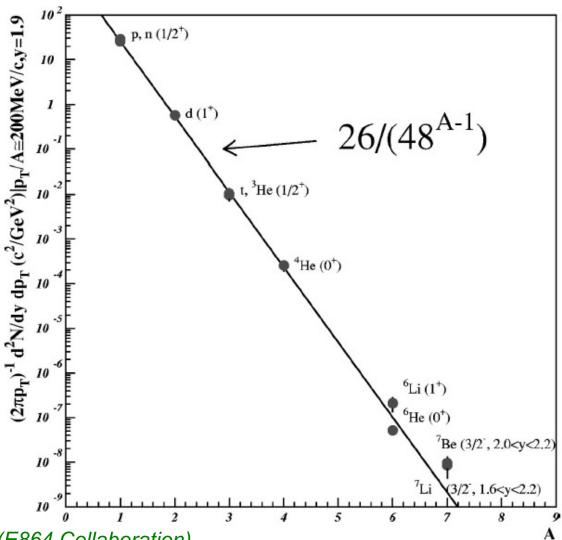
- Run 2 of the LHC has started in 2015 and for Pb-Pb collisions ~ factor 10 increase expected in statistics
- Run 3 & Run 4 of LHC will deliver much more statistics (50 kHz Pb-Pb collision rate)
- Upgraded ALICE detector will be able to cope with the high luminosity
- TPC Upgrade: GEMs for continous readout



- ITS Upgrade: less material budget and more precise tracking for the identification of hyper-nuclei
- Physics which is now done for A = 2 and A = 3 (hyper-)nuclei
 will be done for A = 4



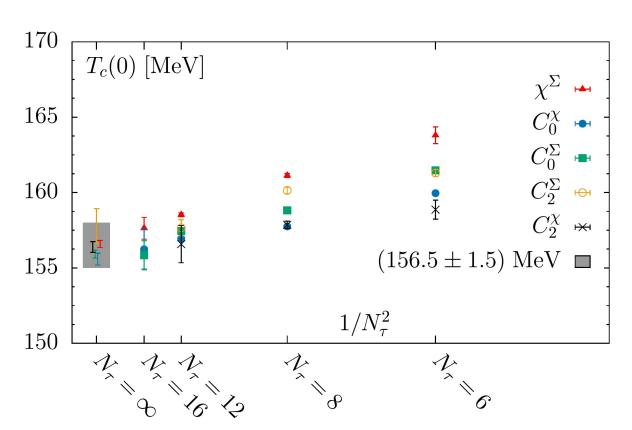
E864 nuclei result



T.A. Armstrong et al. (E864 Collaboration), Phys. Rev. C 61 (2000) 064908



Lattice QCD results



Lattice QCD tells us where to expect the phase transition

Critical energy density: $\epsilon_C = 0.34 \pm 0.16 \text{ GeV/fm}^3$

Critical temperature $T_C = (156.5 \pm 1.5) \text{ MeV}$

A. Bazavov et al. (hotQCD) PLB 795 (2019) 15 Similar results: S. Borsányi et al. (Budapest-Wuppertal group) PRL 125 (2020) 052001



Hypertriton/Λ ratio



- Hypertriton signal recently also extracted in isobar collisions
- Stronger separation between models as for other particle ratios, mainly due to the size of the hypertriton

