

# The Nature of QCD Phase Transitions

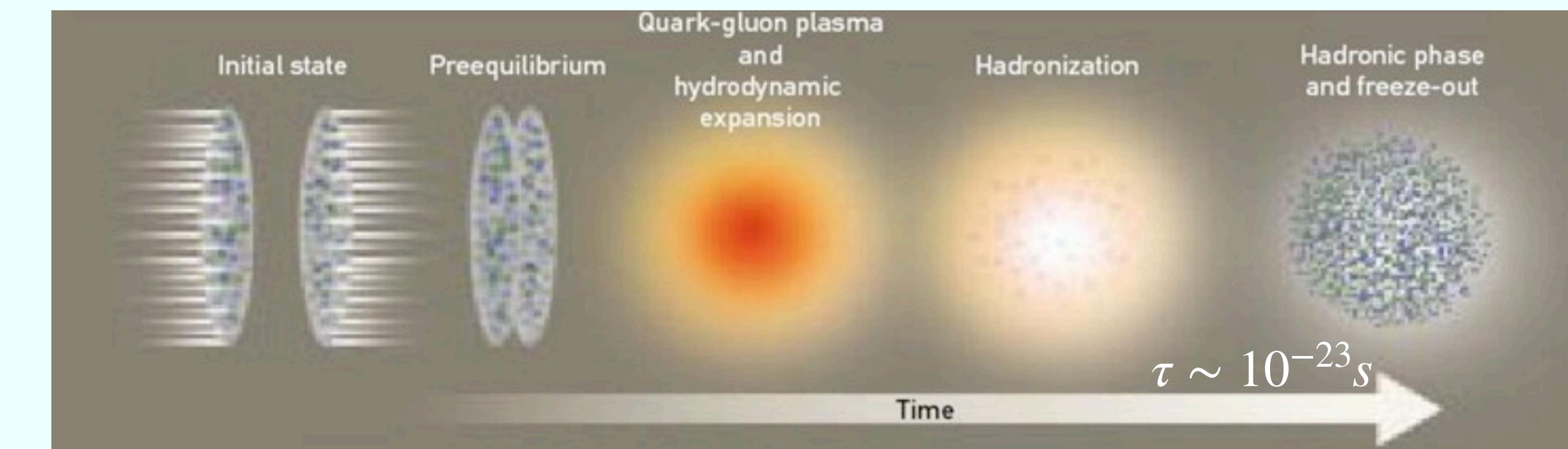
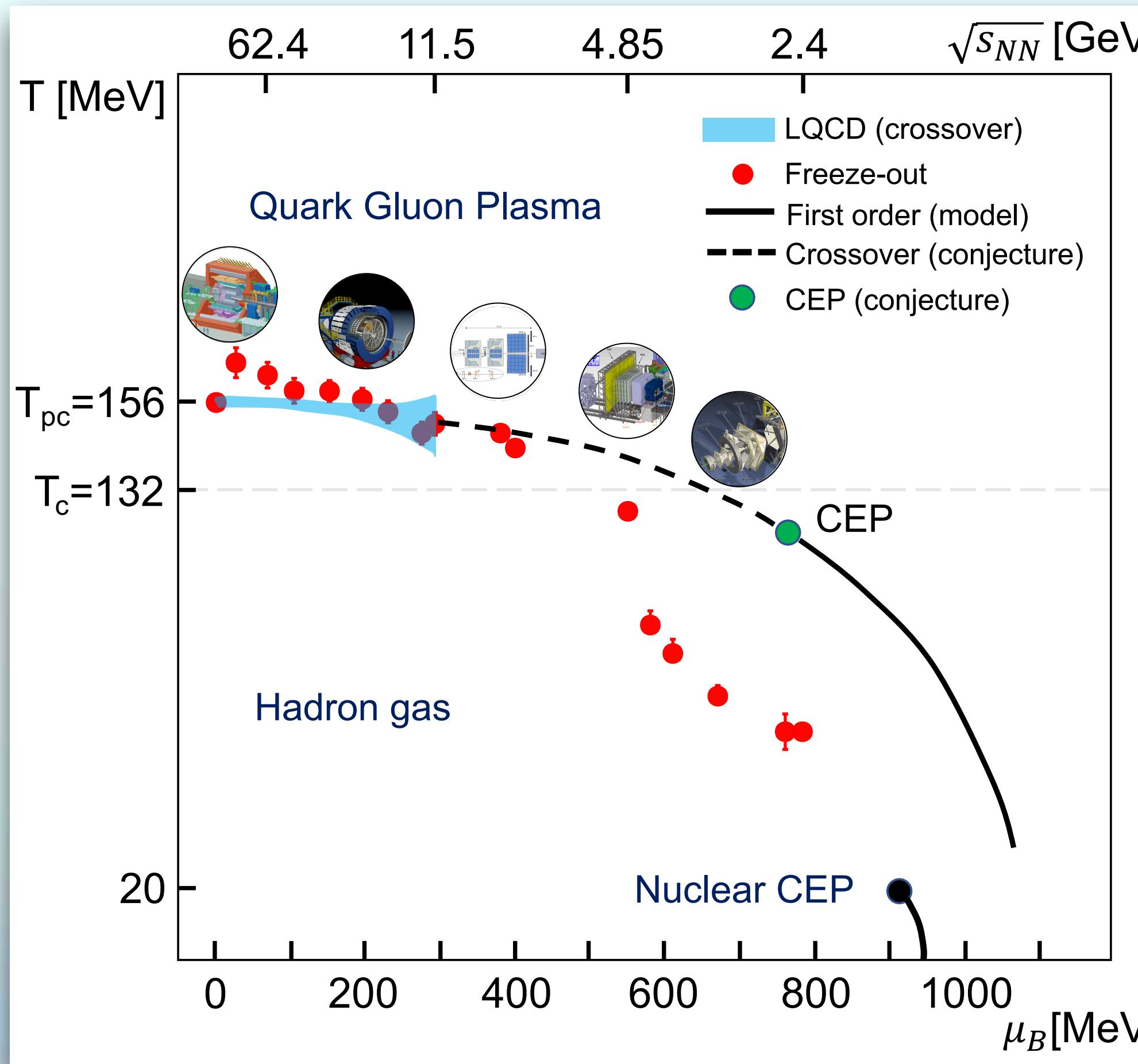
## From cumulants to the Metropolis Algorithm

Anar Rustamov



[a.rustamov@gsi.de](mailto:a.rustamov@gsi.de)  
[a.rustamov@cern.ch](mailto:a.rustamov@cern.ch)

# Phase diagram of QCD matter



$$T_{fo}^{ALICE} = 156.5 \pm 1.5 \pm 3 \text{ MeV(sys)}$$

$$\mu_B \approx 0, V \approx 5280 \text{ fm}^3 \text{ (in one unit of rapidity)} \quad V_{Pb} \approx 1200 \text{ fm}^3$$

$$T_C^{LQCD} = 156.5 \pm 1.5 \text{ MeV}$$

for a thermal system of fixed volume  $V$  and temperature  $T$

$$\frac{\kappa_n(N_B - N_{\bar{B}})}{VT^3} \equiv \frac{1}{VT^3} \frac{\partial^n \ln Z(V, T, \mu_B)}{\partial (\mu_B/T)^n} = \frac{\partial^n P/T^4}{\partial (\mu_B/T)^n} \equiv \hat{\chi}_n^B$$

experiment

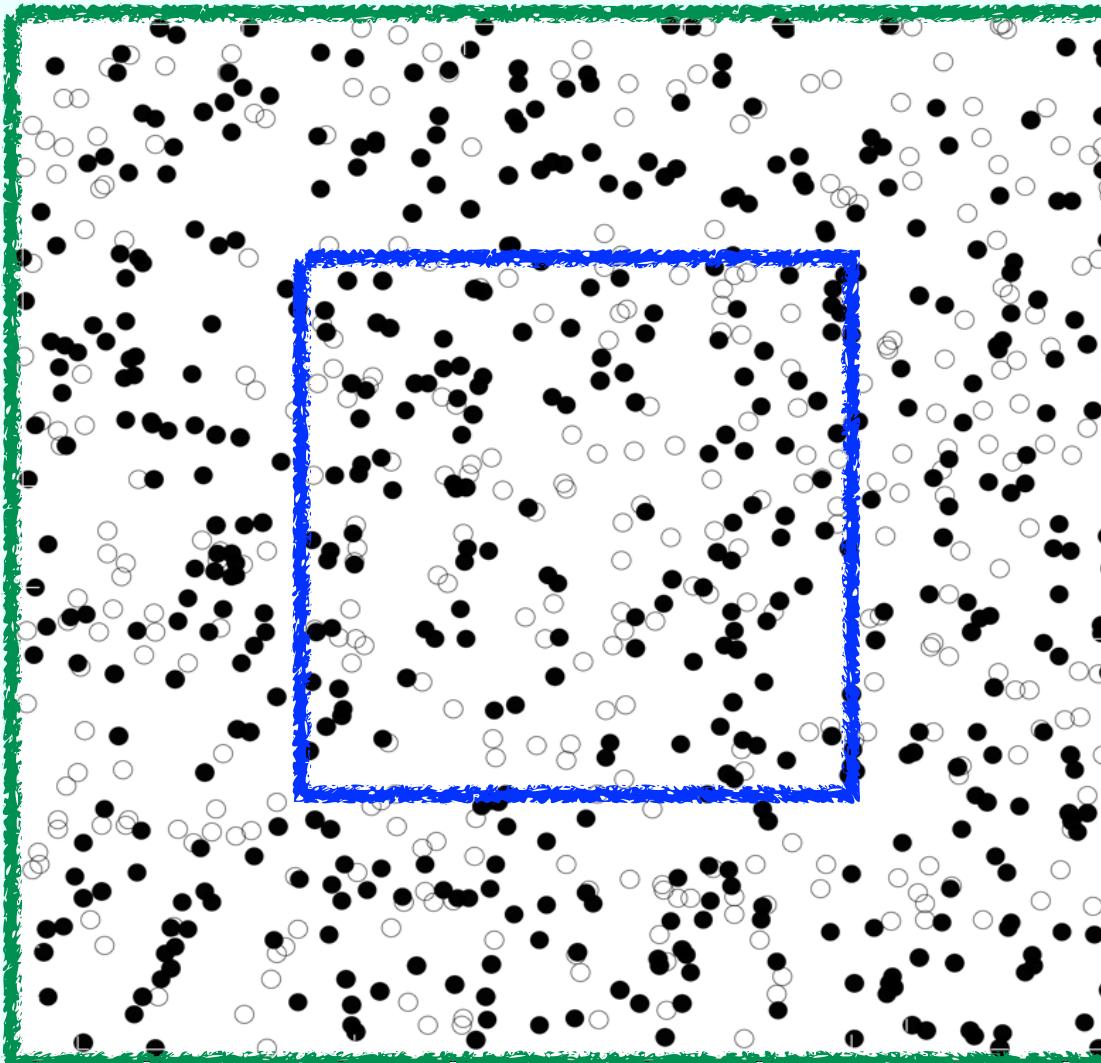
theory

P. Braun-Munzinger, A.R., J. Stachel, 50 years of QCD, EPJ C 83 (2023) 1125

FO: A. Andronic, P. Braun-Munzinger, K. Redlich and J. Stachel, Nature 561, 321–330 (2018)

IQCD: A. Bazavov et al., (HotQCD), PLB 795 (2019) 15-21

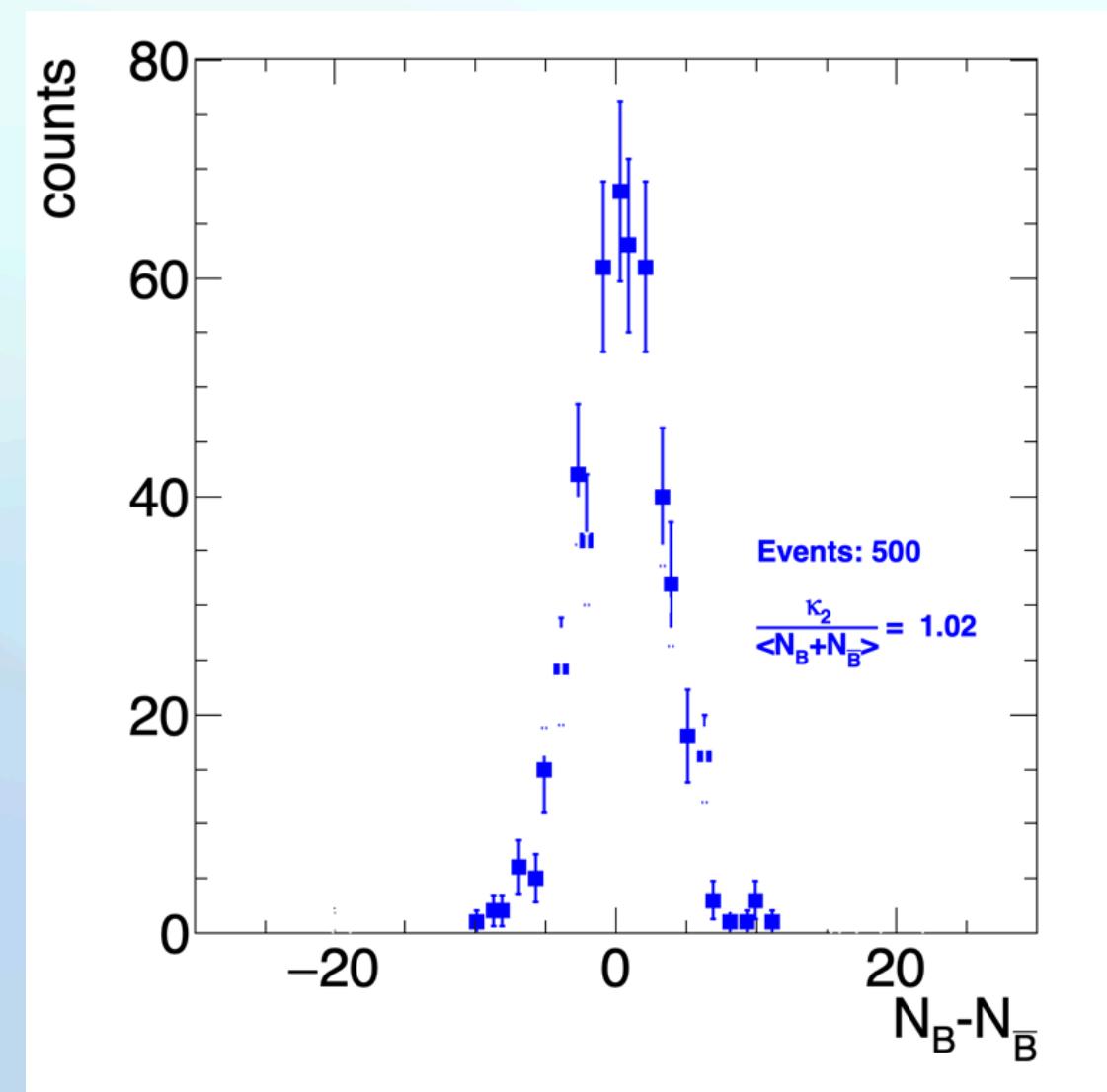
# Probing EoS with event-by-event fluctuations



- $\Delta N = N_B - N_{\bar{B}}$  occurs with probability  $p(\Delta N)$  (measured)
- $r^{th}$  order central moment:  $\mu_r = \sum_{\Delta N} (\Delta N - \langle \Delta N \rangle)^r p(\Delta N)$
- first 4 cumulants:  $\kappa_1 = \langle \Delta N \rangle$ ,  $\kappa_2 = \mu_2$ ,  $\kappa_3 = \mu_3$ ,  $\kappa_4 = \mu_4 - 3\mu_2^2$

$$\frac{\kappa_n(N_B - N_{\bar{B}})}{VT^3} \equiv \frac{1}{VT^3} \frac{\partial^n \ln Z(V, T, \mu_B)}{\partial (\mu_B/T)^n} = \frac{\partial^n P/T^4}{\partial (\mu_B/T)^n} \equiv \hat{\chi}_n^B$$

**experiment**



- “volume” fluctuates
- exact conservations
- measures **net-protons**

**theory**

- volume is fixed
- conservations on average
- predicts for **net-baryons**

## **Experimental challenges**

- Volume fluctuations
- Conservation laws
- First ALICE results

## **Experiment vs. Theory**

- Canonical Thermodynamics
- Comparison to STAR results
- Metropolis algorithm
- Comparison to ALICE results
- Outlook

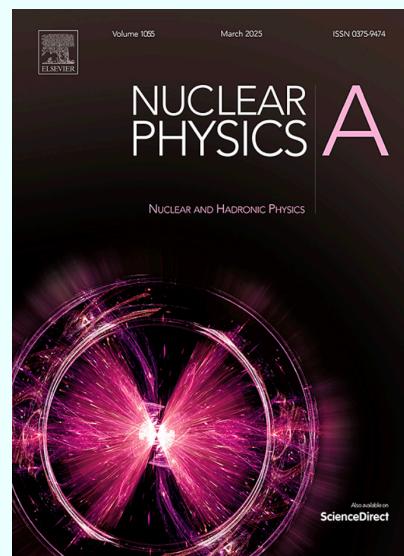
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# Bridging experiment with theory



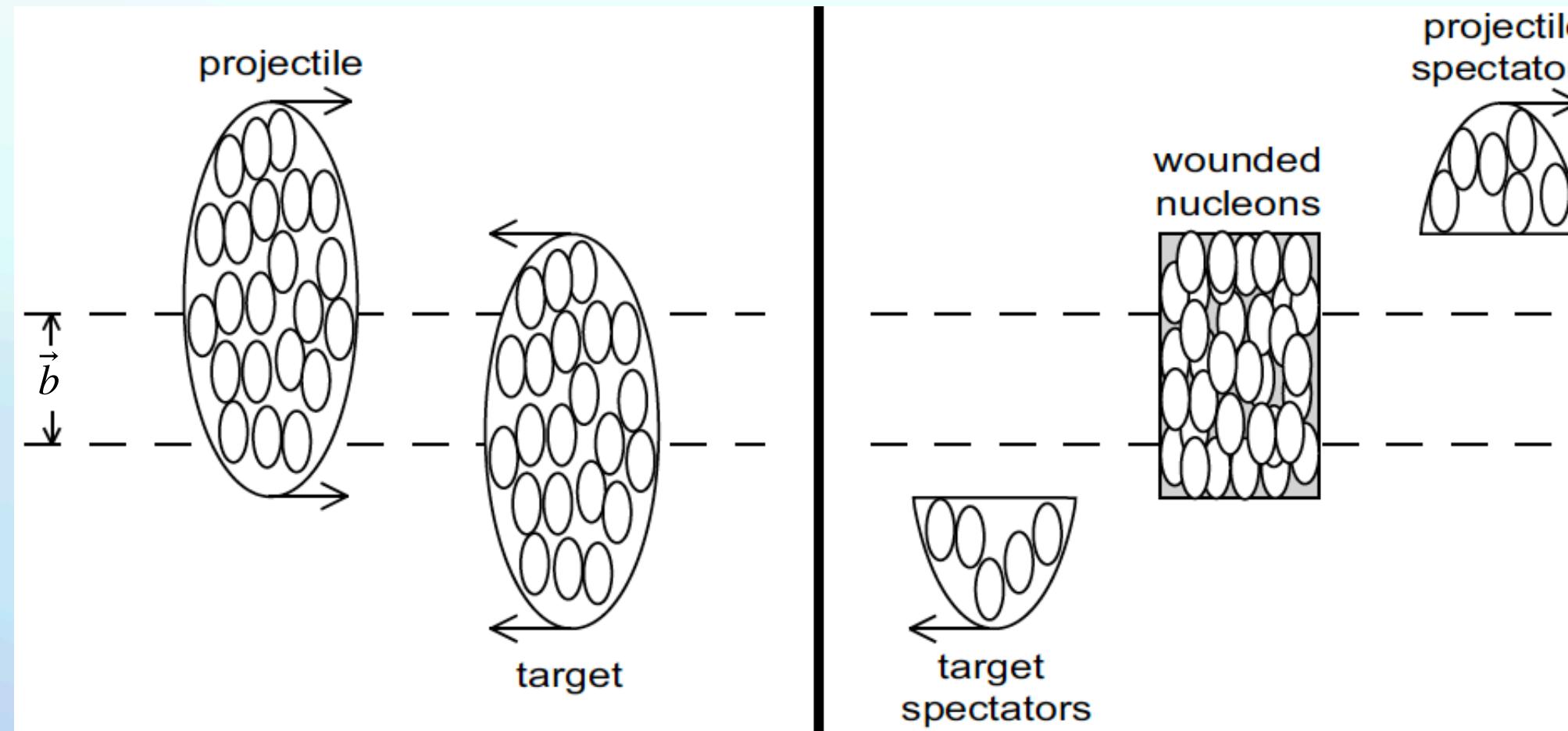
Bridging the gap between event-by-event fluctuation measurements and theory predictions in relativistic nuclear collisions, [P. Braun-Munzinger, A. Rustamov, J. Stachel, Nucl. Phys. A 960 \(2017\) 114-130.](#)

## fixed volume/sources

$$\langle N \rangle \sim V, \quad \langle N \rangle \sim N_W$$

$$\kappa_2(N) \sim V, \quad \kappa_2(N) \sim N_W$$

[V. Skokov, B. Friman, and K. Redlich, Phys.Rev. C88 \(2013\) 034911](#)

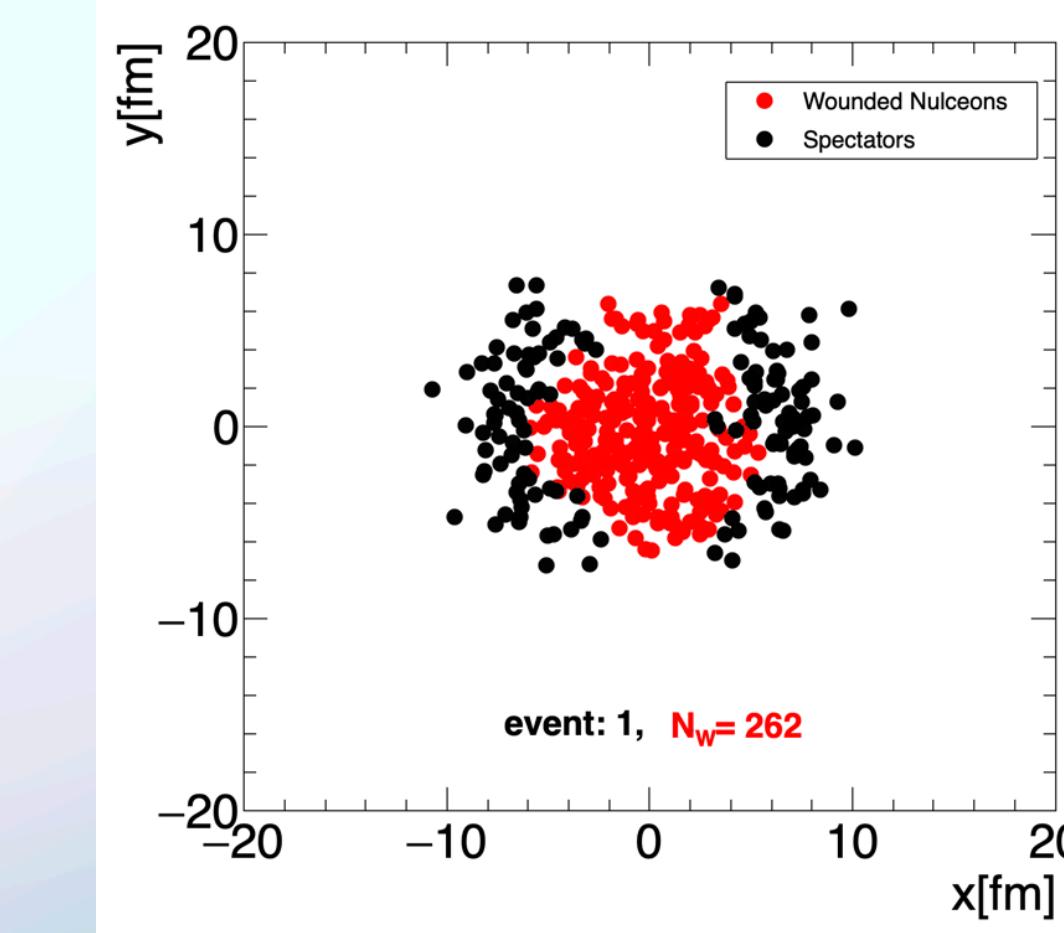


## fluctuating sources

$$\langle N \rangle = \langle n \rangle \langle N_W \rangle$$

$$\kappa_2(N) = \kappa_2(n) \langle N_W \rangle +$$

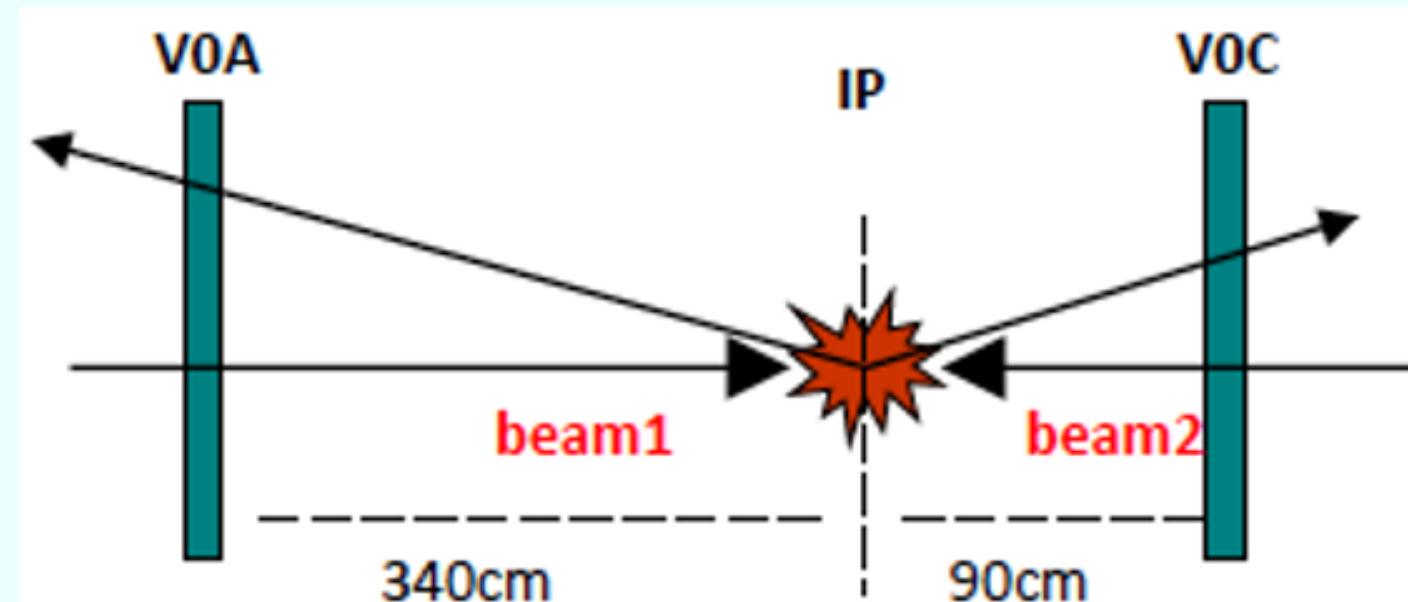
$$\langle N \rangle^2 \frac{\kappa_2(N_W)}{\langle N_W \rangle^2}$$



**Wounded nucleons,  $N_W$ : Nucleons which collided at least once inelastically**

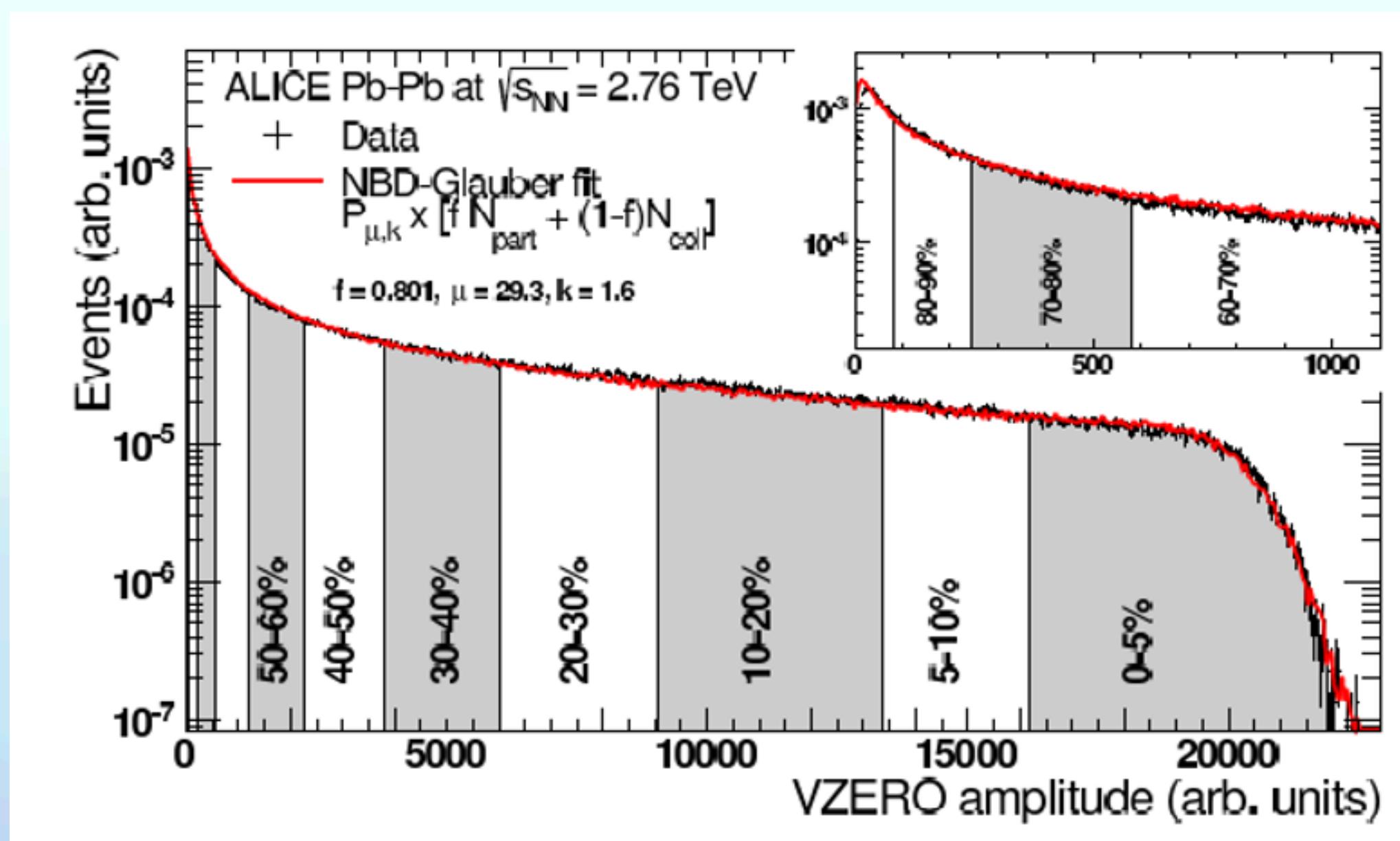
# Contributions from wounded nucleon fluctuations

model of independent sources

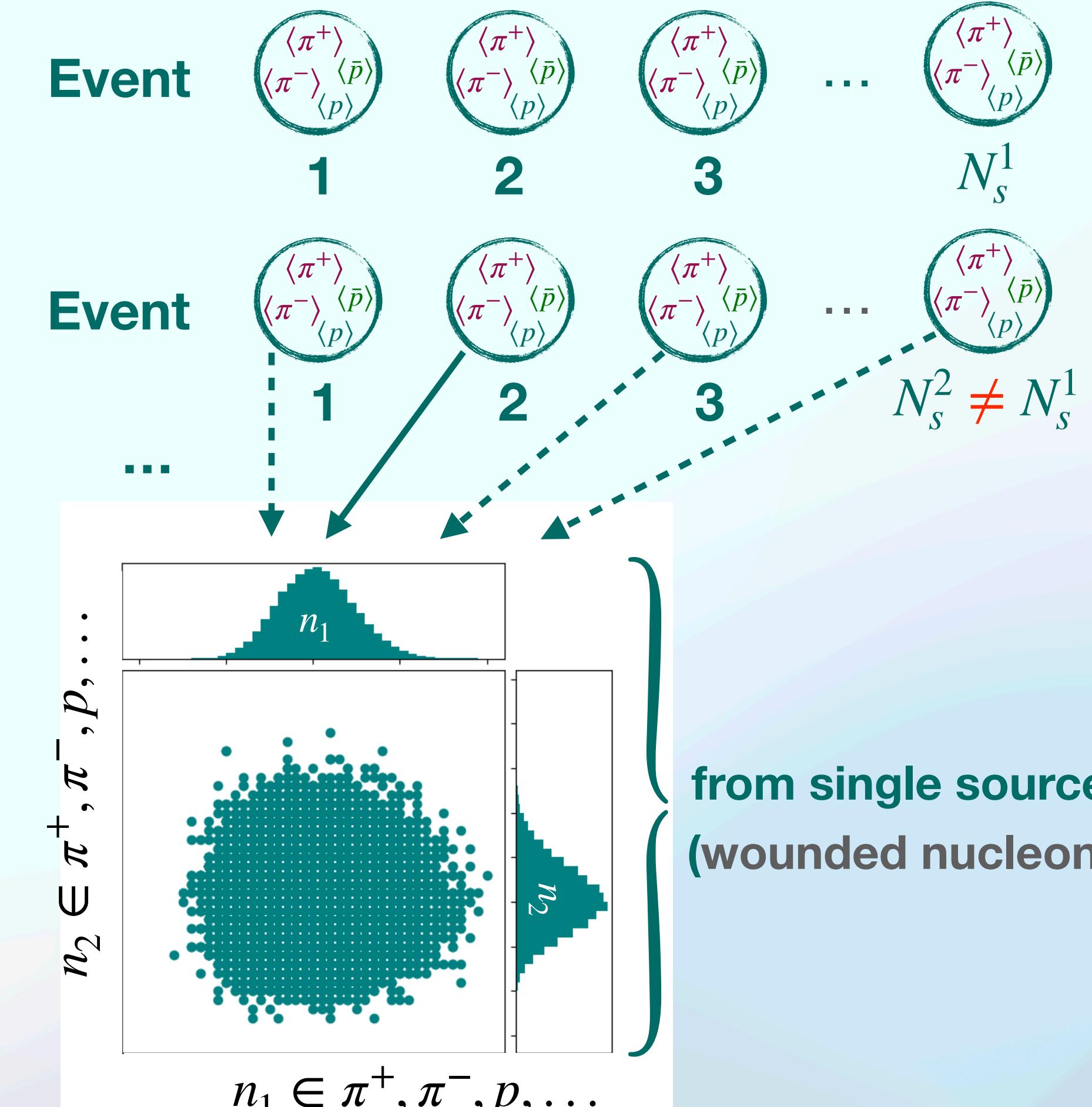


$$2.8 < \eta < 5.1$$

$$-3.7 < \eta < -1.7$$



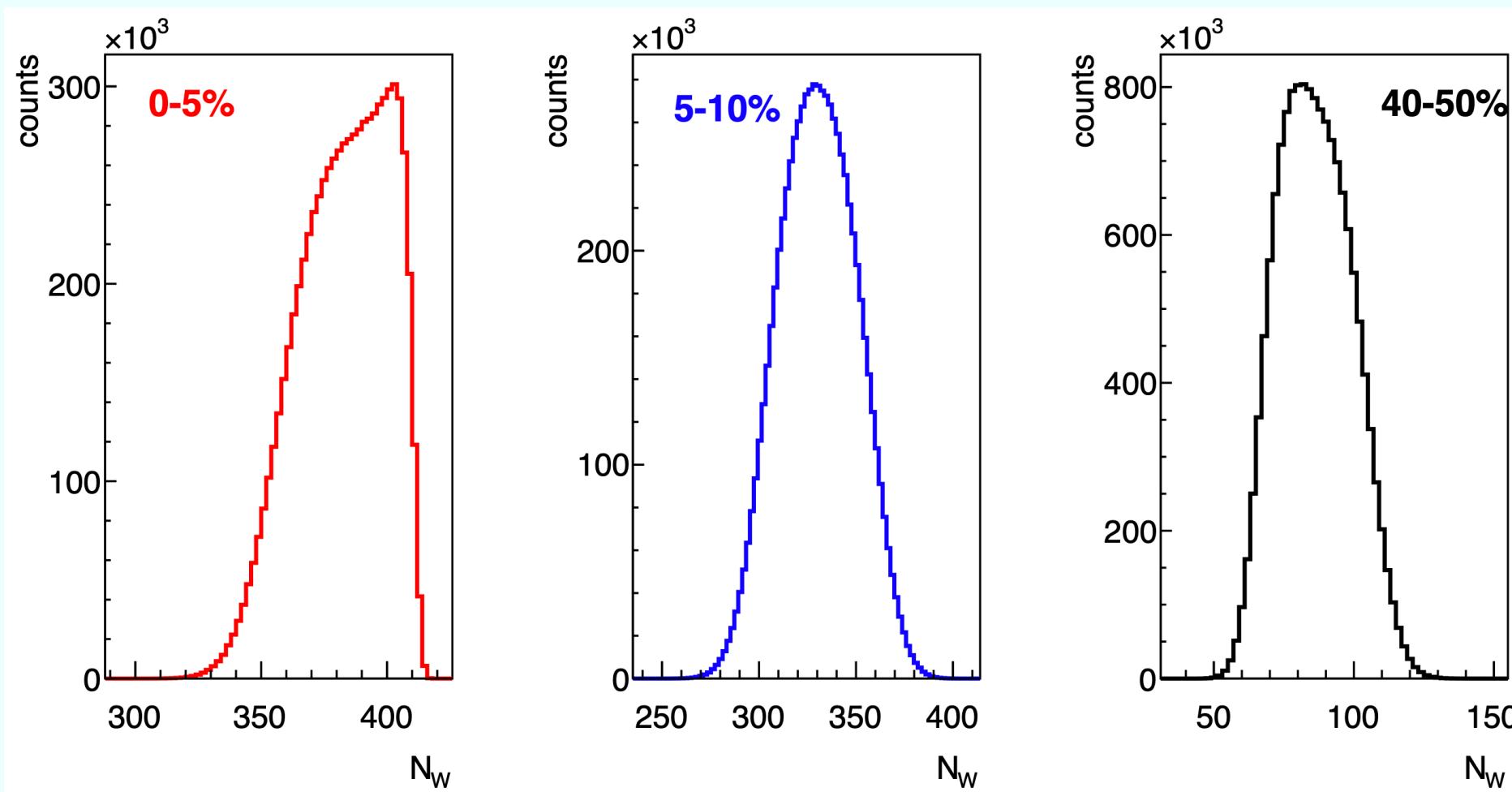
ALICE Phys.Rev. C88 (2013) no.4, 044909



A. R., R. Holzmann, J. Stroth, NPA 1034 (2023) 122641

V. Koch, R. Holzmann, A. R., J. Stroth, Nucl.Phys.A 1050 (2024) 122924

# Volume fluctuations



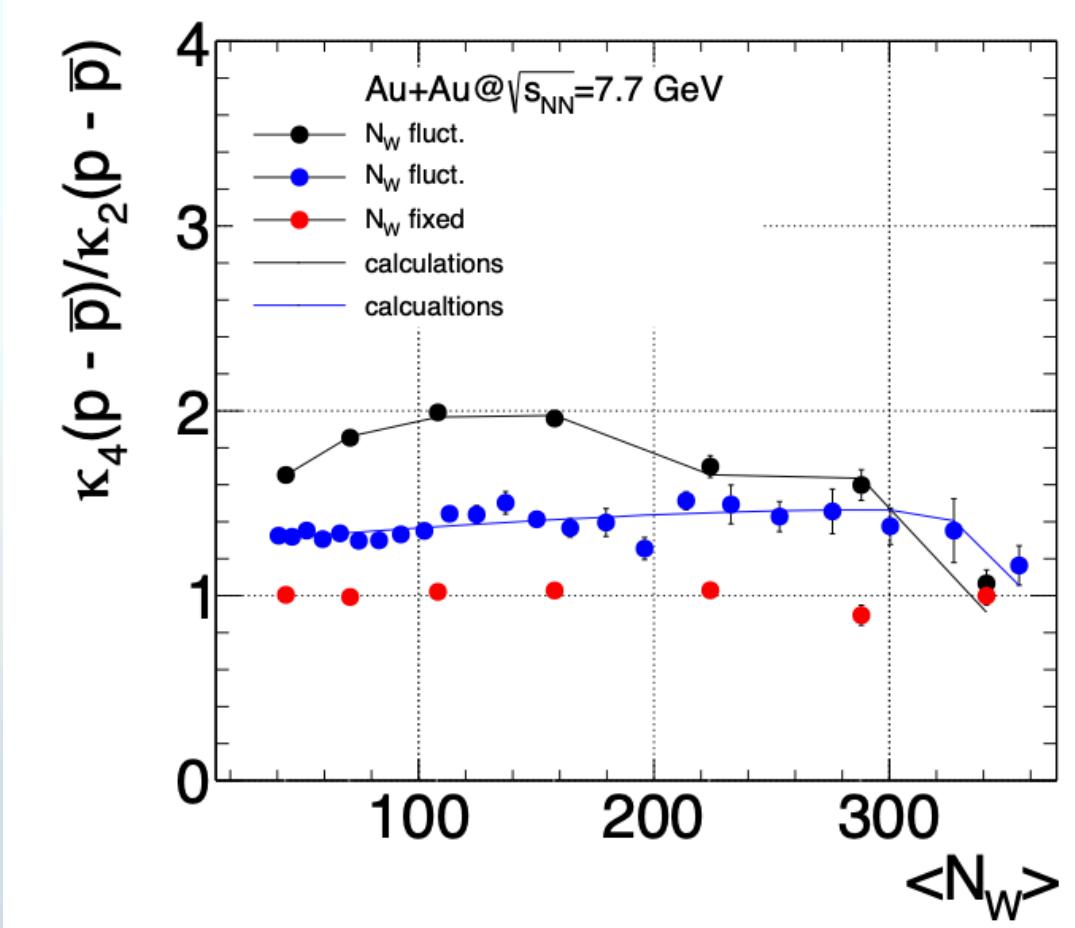
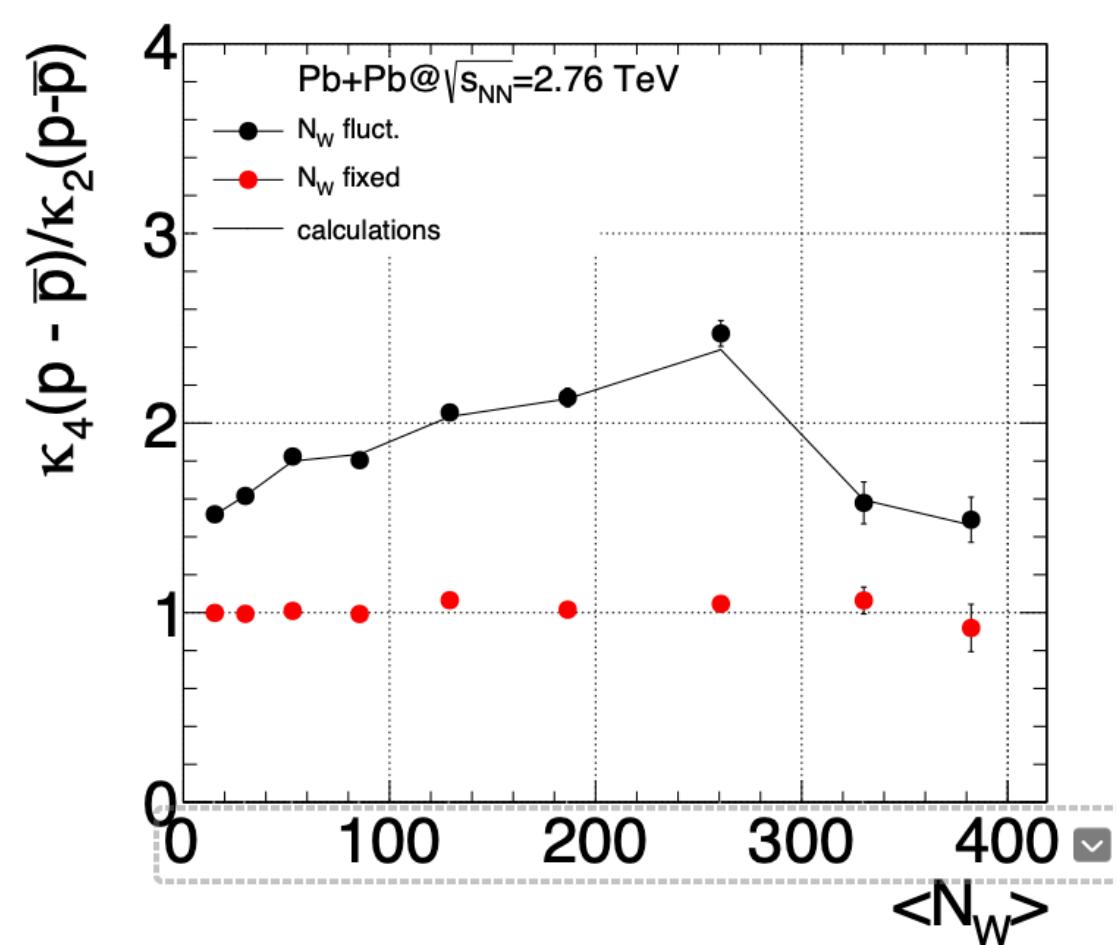
$$\kappa_2(N - \bar{N}) = \kappa_2(n - \bar{n})\langle N_W \rangle + \langle N - \bar{N} \rangle^2 \frac{\kappa_2(N_W)}{\langle N_W \rangle^2}$$

vanishes for ALICE

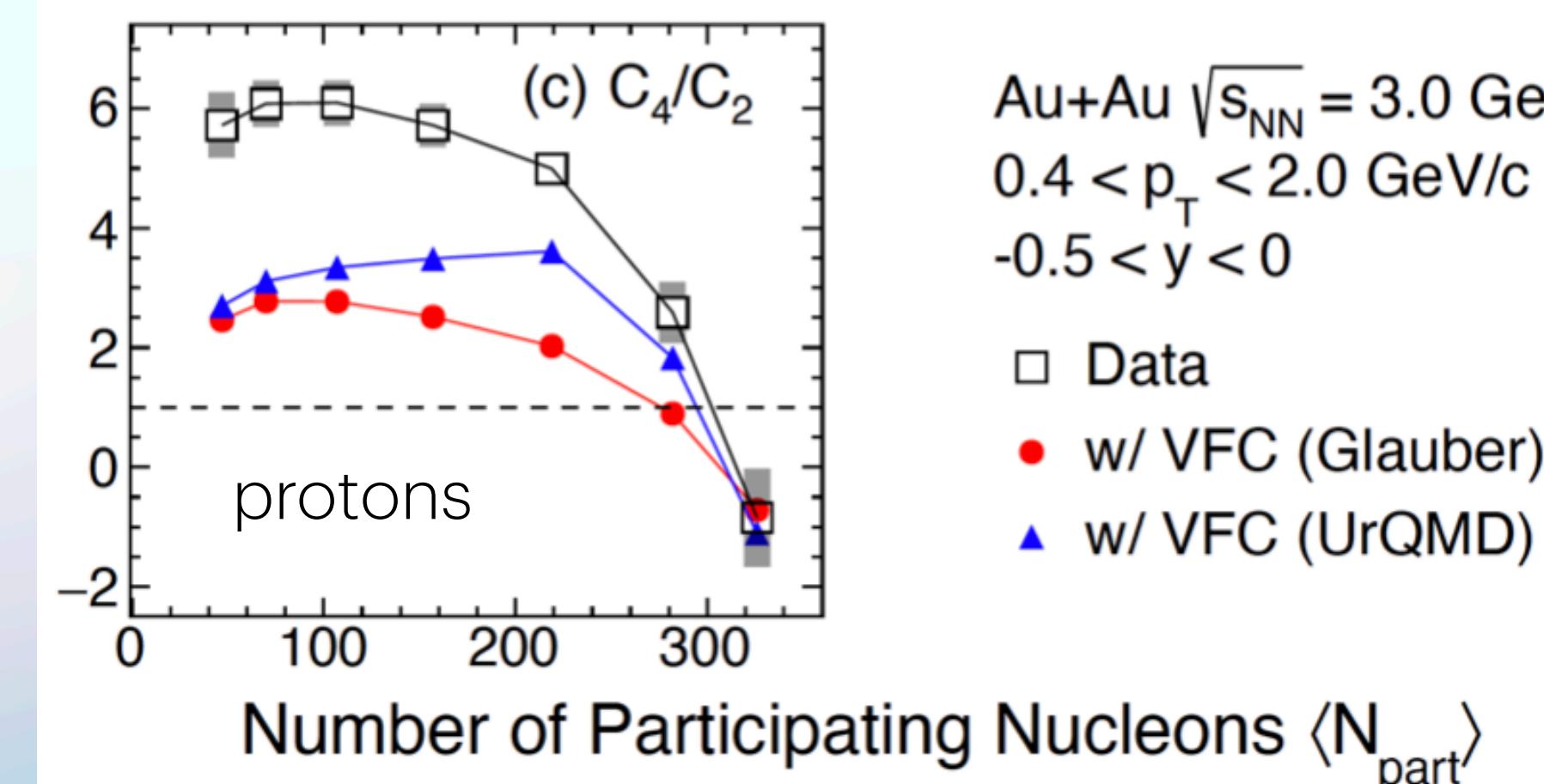
$$\begin{aligned} \kappa_4(\Delta N) &= \langle N_W \rangle \kappa_4(\Delta n) + 4 \langle \Delta n \rangle \kappa_3(\Delta n) \kappa_2(N_W) \\ &+ 3 \kappa_2^2(\Delta n) \kappa_2(N_W) + 6 \langle \Delta n \rangle^2 \kappa_2(\Delta n) \kappa_3(N_W) + \langle \Delta n \rangle^4 \kappa_4(N_W) \end{aligned}$$

may be negative

**Some predictions** P. Braun-Munzinger, A. R., J. Stachel, NPA 960 (2017) 114-130



**STAR experiment** PRL. 128 (2022) 20, 202303



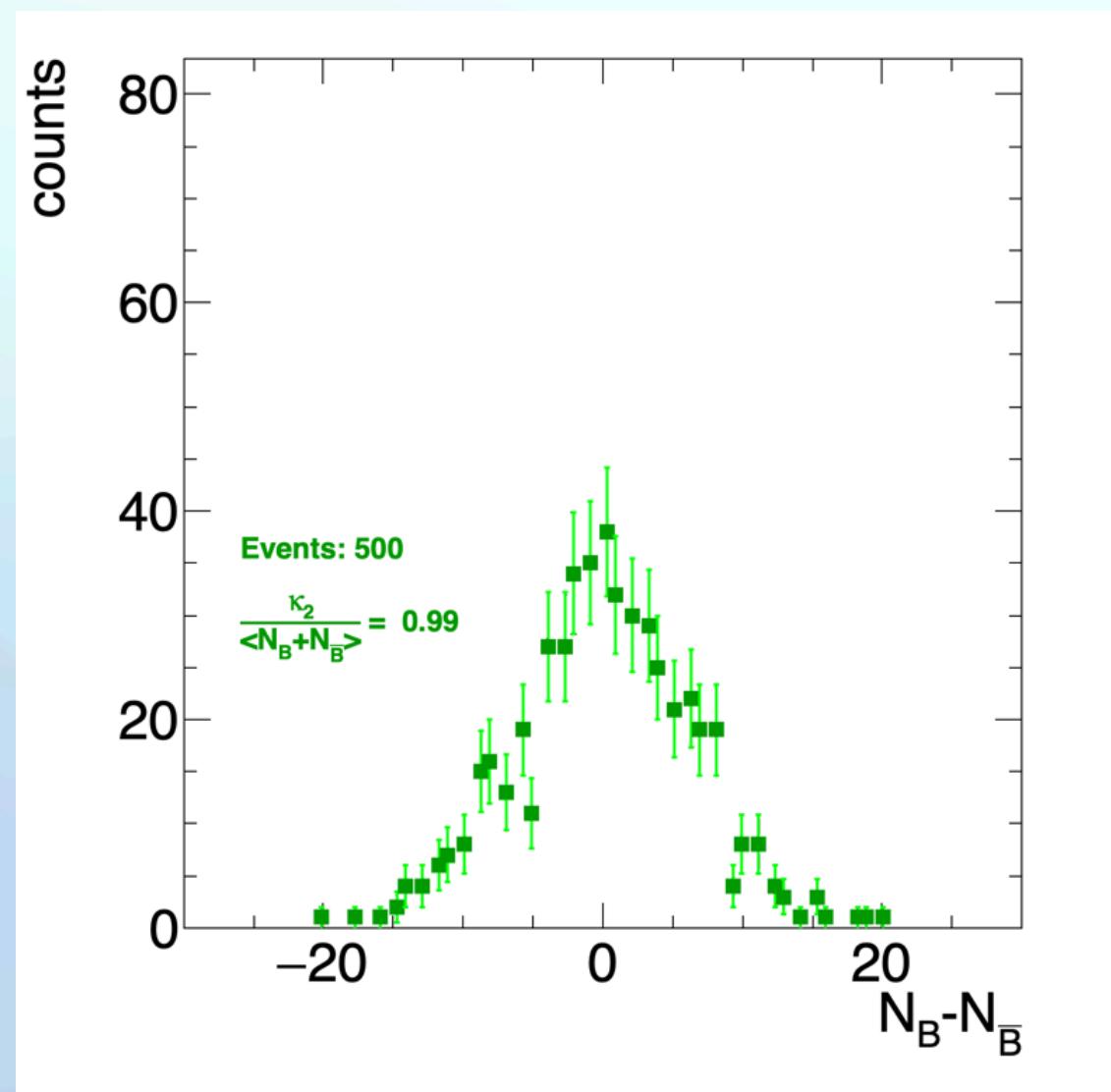
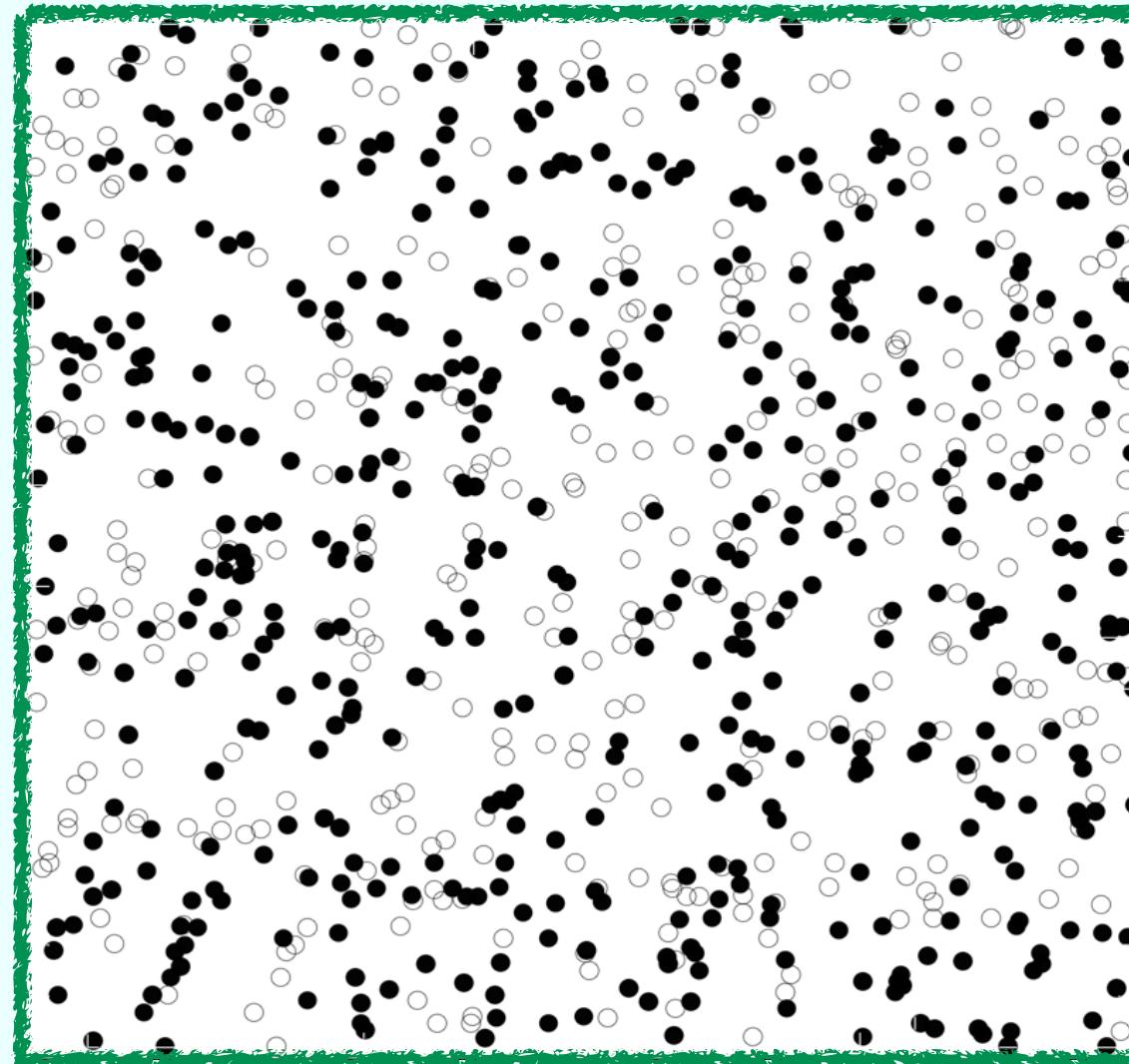
## Experimental challenges

- Volume fluctuations
- Conservation laws
- First ALICE results

## Experiment vs. Theory

- Canonical Thermodynamics
- Comparison to STAR results
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# Ideal Gas in GCE



$$\frac{\kappa_n(N_B - \bar{N}_B)}{VT^3} = \frac{1}{VT^3} \frac{\partial^n \ln Z(V, T, \mu_B)}{\partial (\mu_B/T)^n} \equiv \hat{\chi}_n^B$$

## Ideal Gas in Grand Canonical Ensemble

**particles**

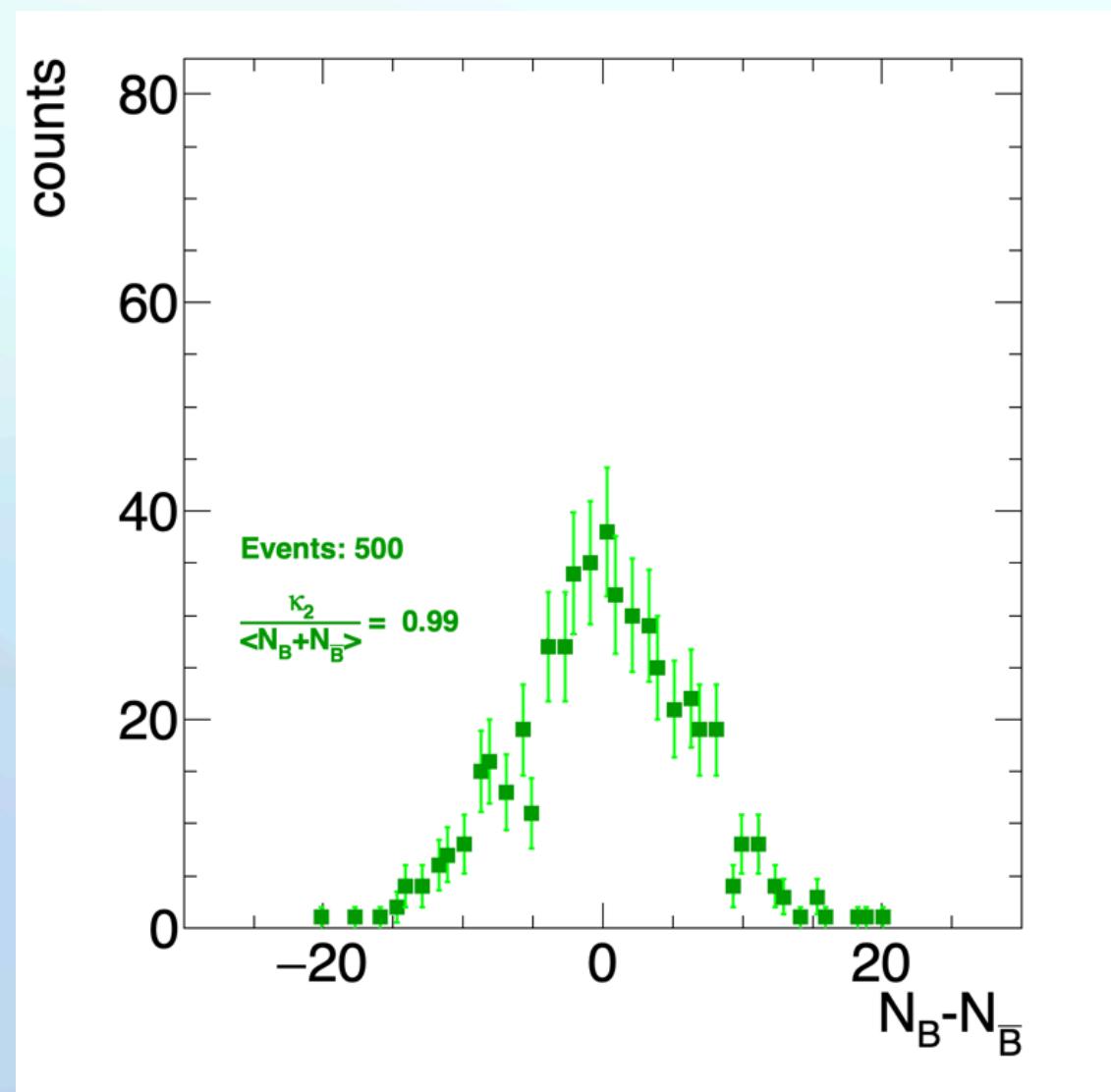
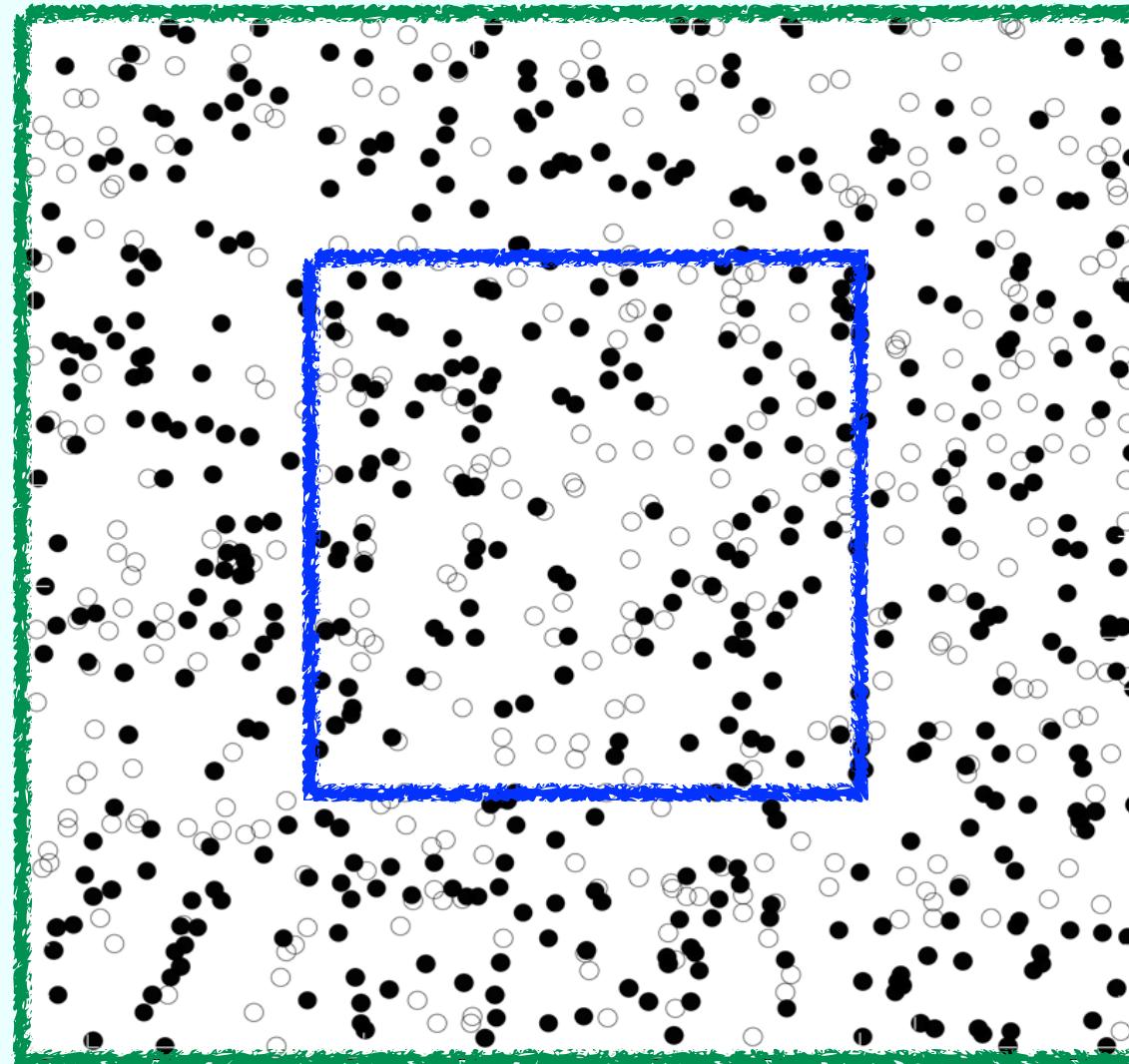
$$\kappa_n(N) = \langle N \rangle \text{ (Poisson distribution)}$$

**net-particles**

$$\kappa_n(N - \bar{N}) = \langle N \rangle + (-1)^n \langle \bar{N} \rangle \text{ (Skellam distribution)}$$

for example:  $\frac{\kappa_2(N - \bar{N})}{\langle N + \bar{N} \rangle} = \frac{\langle N + \bar{N} \rangle}{\langle N + \bar{N} \rangle} = 1$

# Ideal Gas in GCE



$$\frac{\kappa_n(N_B - \bar{N}_B)}{VT^3} = \frac{1}{VT^3} \frac{\partial^n \ln Z(V, T, \mu_B)}{\partial (\mu_B/T)^n} \equiv \hat{\chi}_n^B$$

## Ideal Gas in Grand Canonical Ensemble

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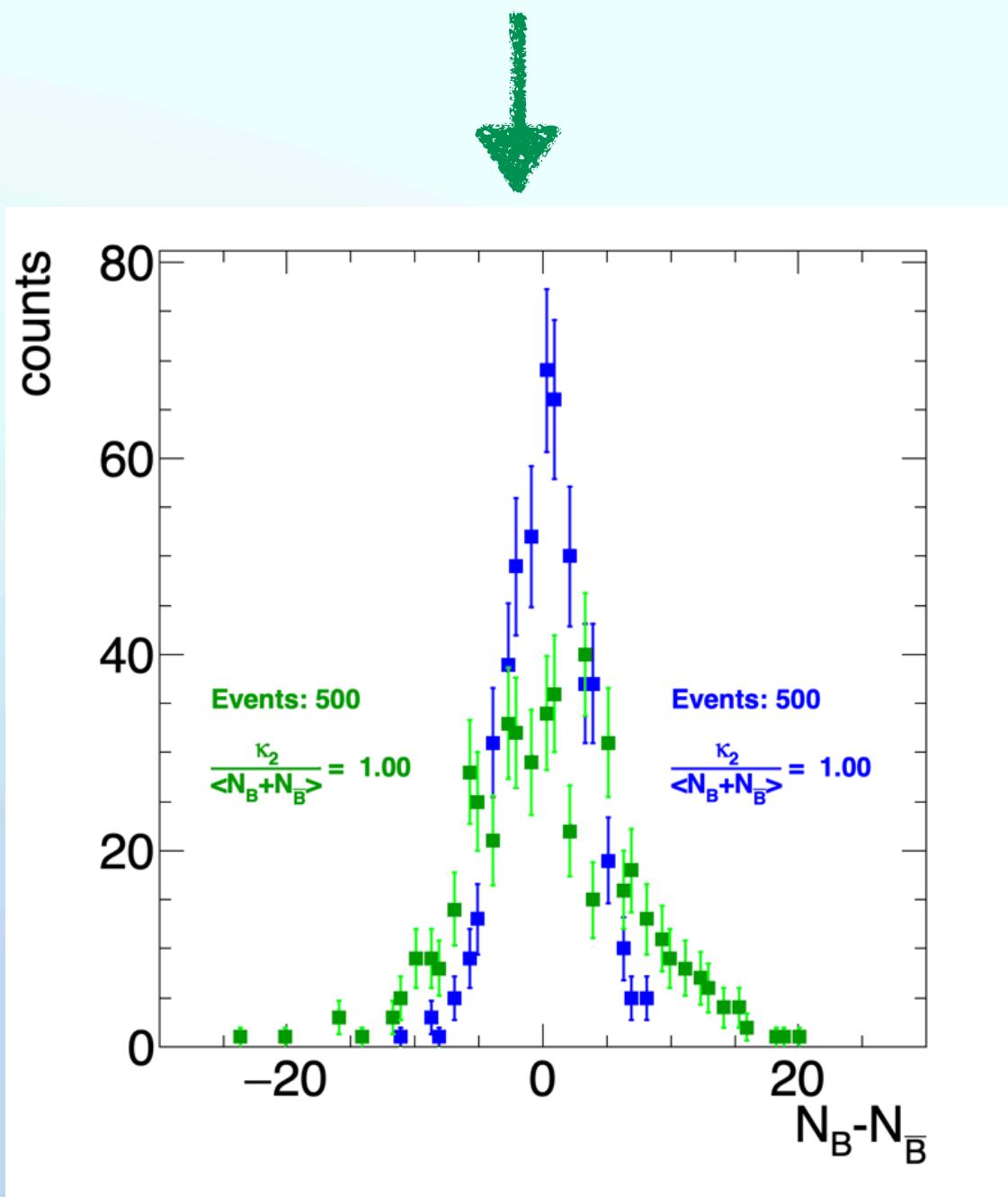
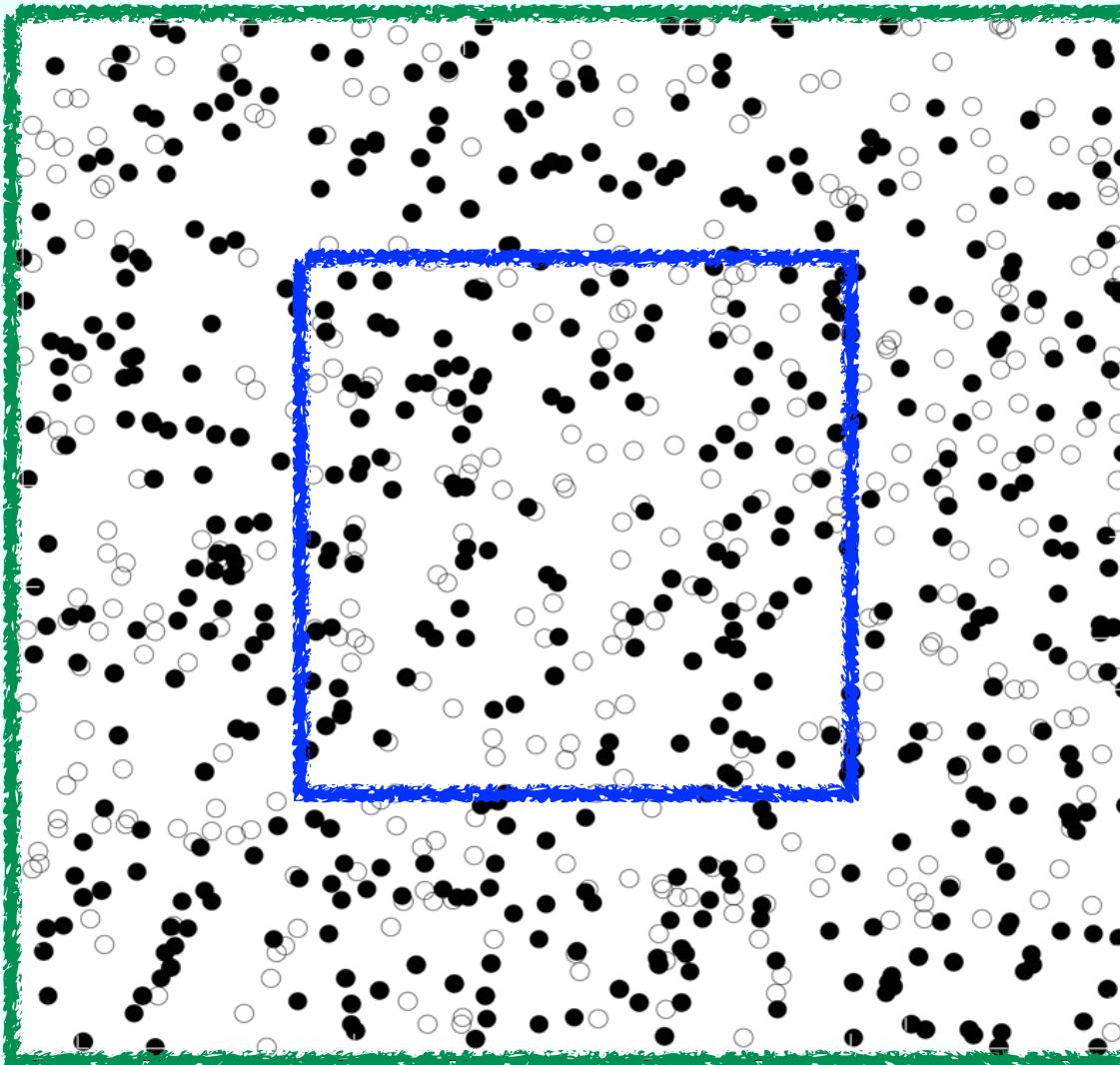
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# Ideal Gas in GCE + conservation laws



$$\frac{\kappa_2(N_B - N_{\bar{B}})}{\langle N_B + N_{\bar{B}} \rangle} = \alpha \quad \frac{\kappa_2(N_B - N_{\bar{B}})}{\langle N_B + N_{\bar{B}} \rangle} + 1 - \alpha$$

$$\alpha = \langle N_B \rangle / \langle N_{\bar{B}} \rangle$$

**Unity in GCE**

P. Braun-Munzinger, A.R., J. Stachel, NPA 960 (2017) 114-130

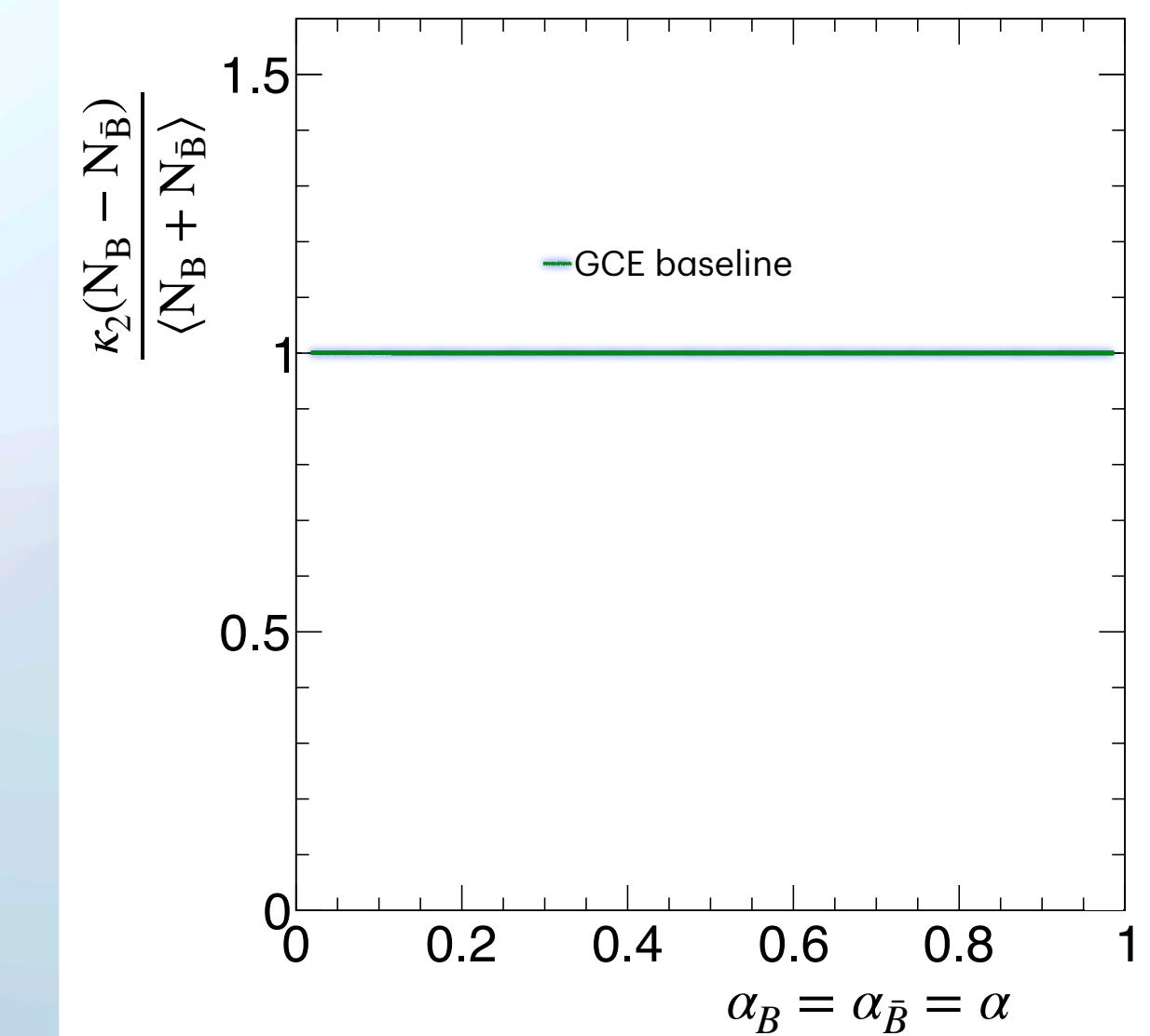
If baryon number is conserved in full phase space

$$\frac{\kappa_2(N_B - N_{\bar{B}})}{\langle N_B + N_{\bar{B}} \rangle} = 1 - \alpha$$

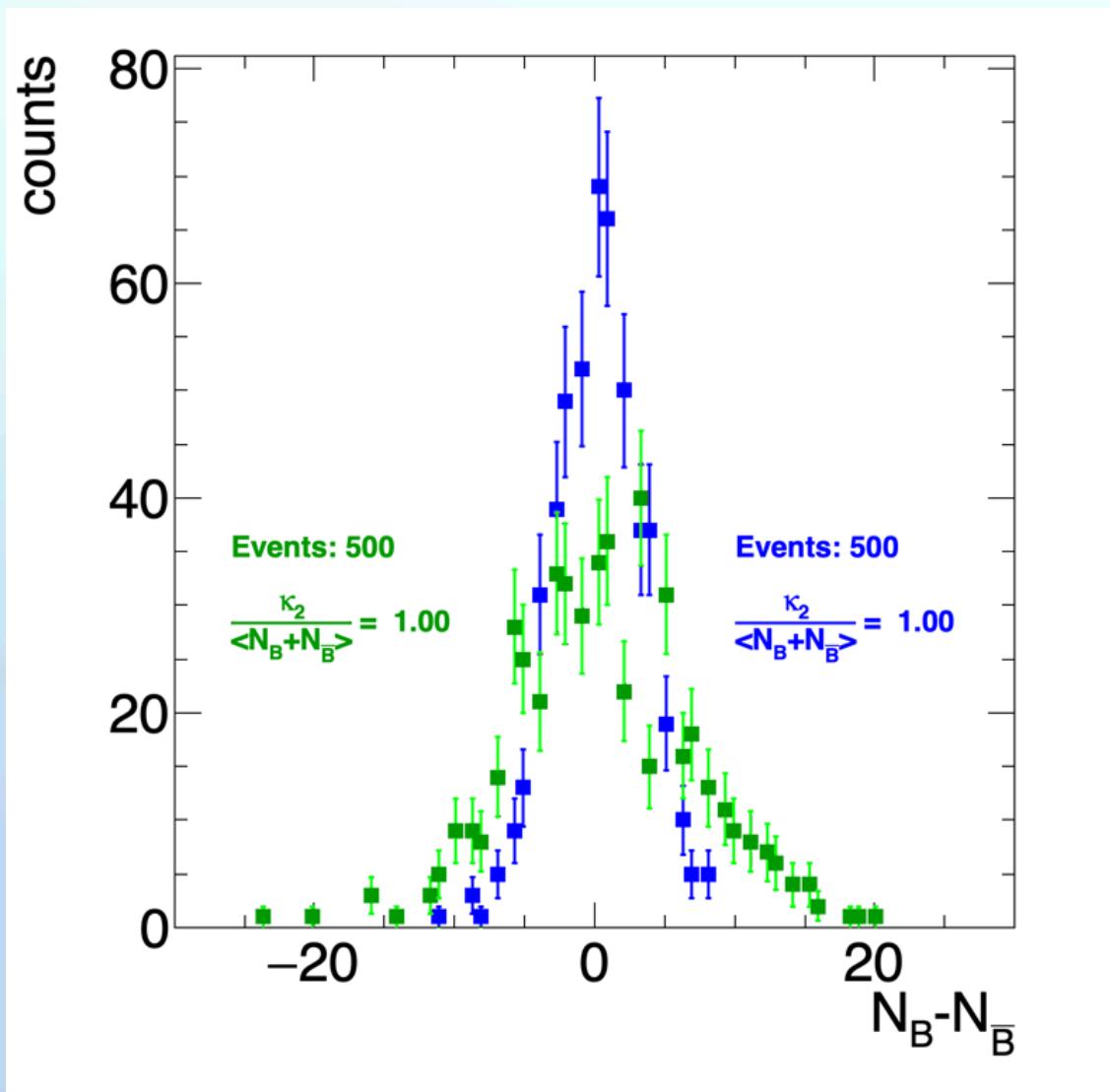
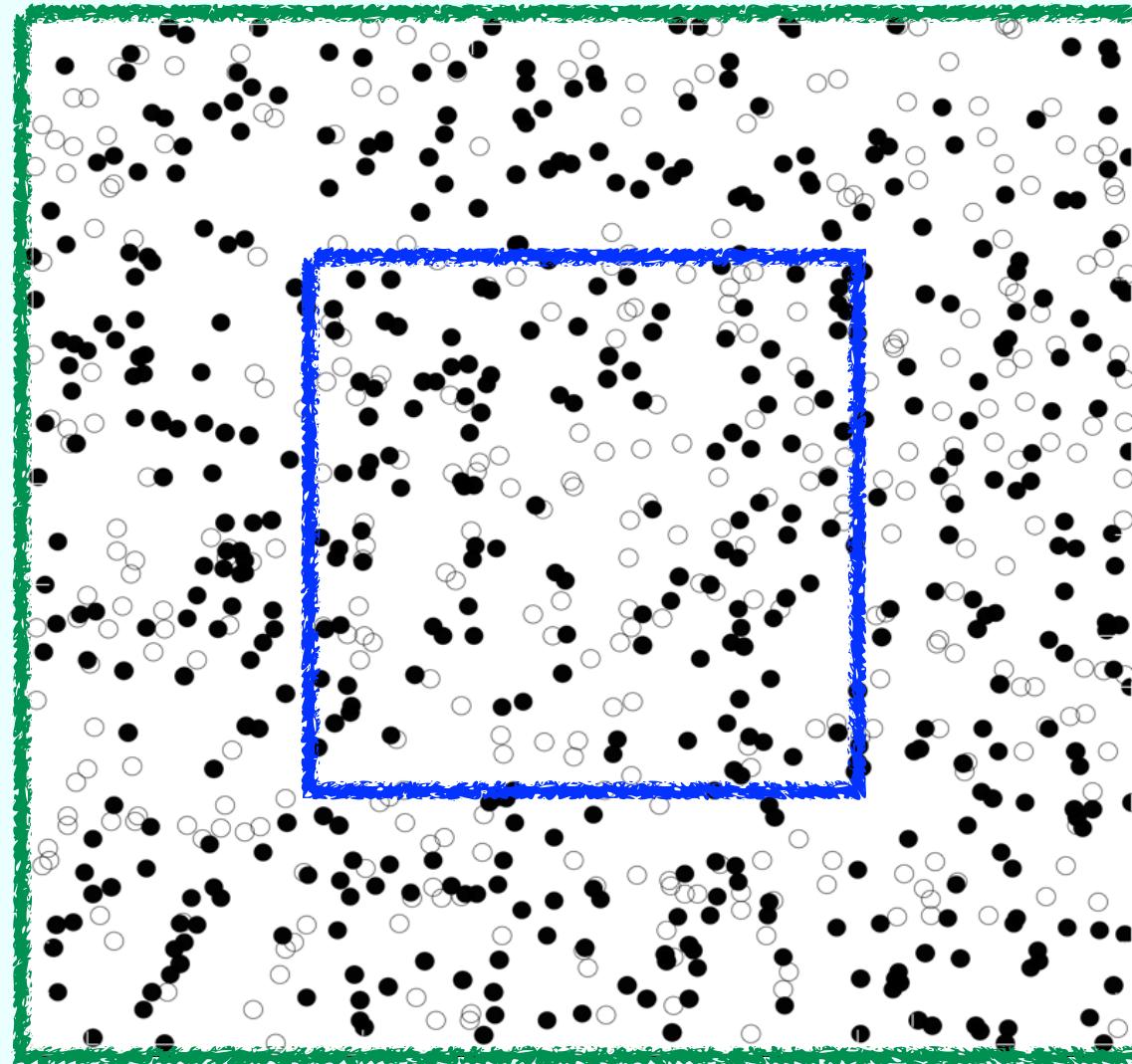
$$\frac{\kappa_2(N_B - N_{\bar{B}})}{\langle N_B + N_{\bar{B}} \rangle} = 0$$

$$\frac{\kappa_2(N_B - N_{\bar{B}})}{\langle N_B + N_{\bar{B}} \rangle} = 1$$

$$\frac{\kappa_2(N_B - N_{\bar{B}})}{\langle N_B + N_{\bar{B}} \rangle} = 1 - \alpha$$



# Ideal Gas in GCE + conservation laws



$$\frac{\kappa_2(N_B - N_{\bar{B}})}{\langle N_B + N_{\bar{B}} \rangle} = \alpha \quad \frac{\kappa_2(N_B - N_{\bar{B}})}{\langle N_B + N_{\bar{B}} \rangle} + 1 - \alpha$$

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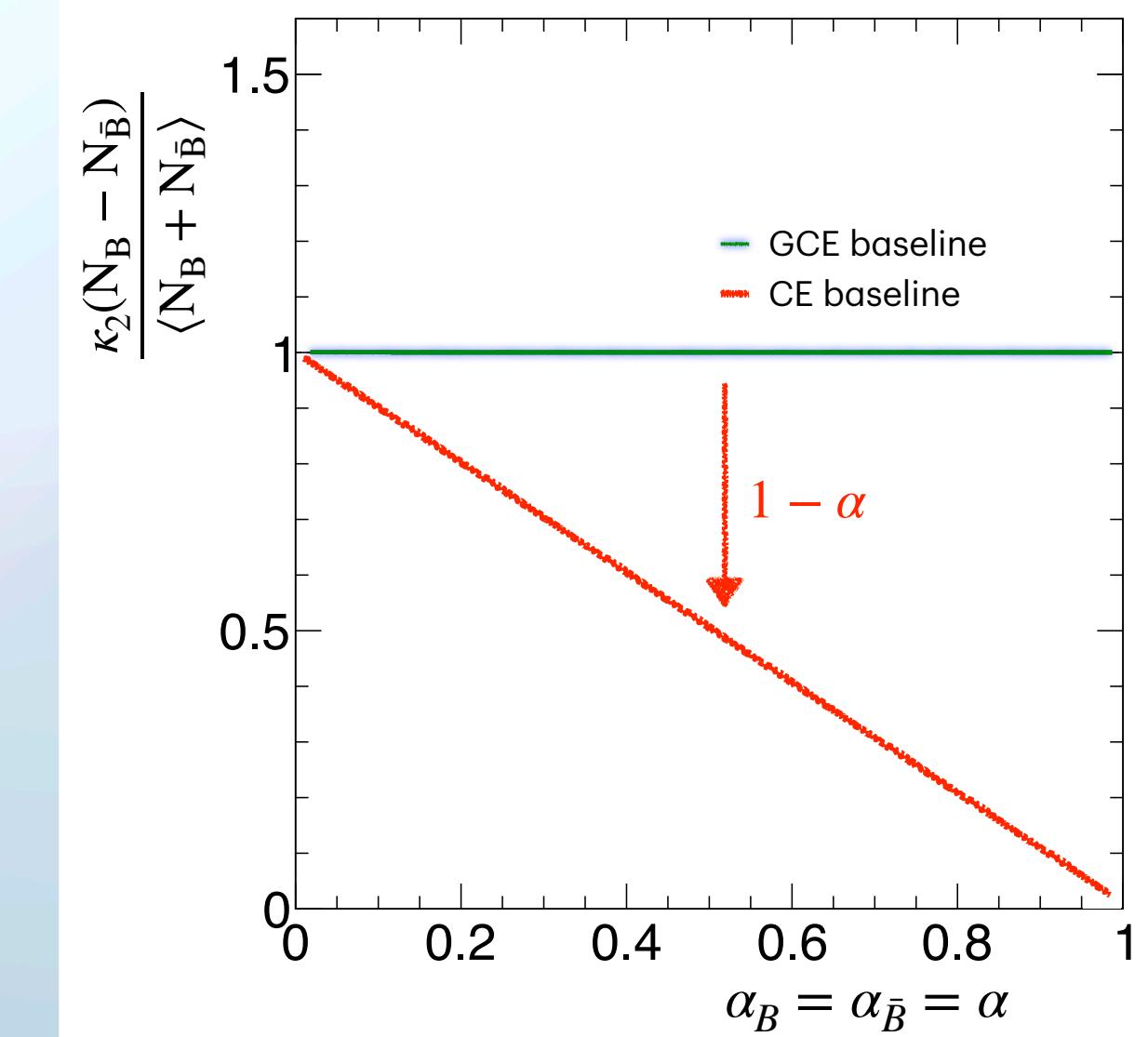
$$\frac{\kappa_2(N_B - N_{\bar{B}})}{\langle N_B + N_{\bar{B}} \rangle} = 0$$

**Grand Canonical  
(conservation on average)**

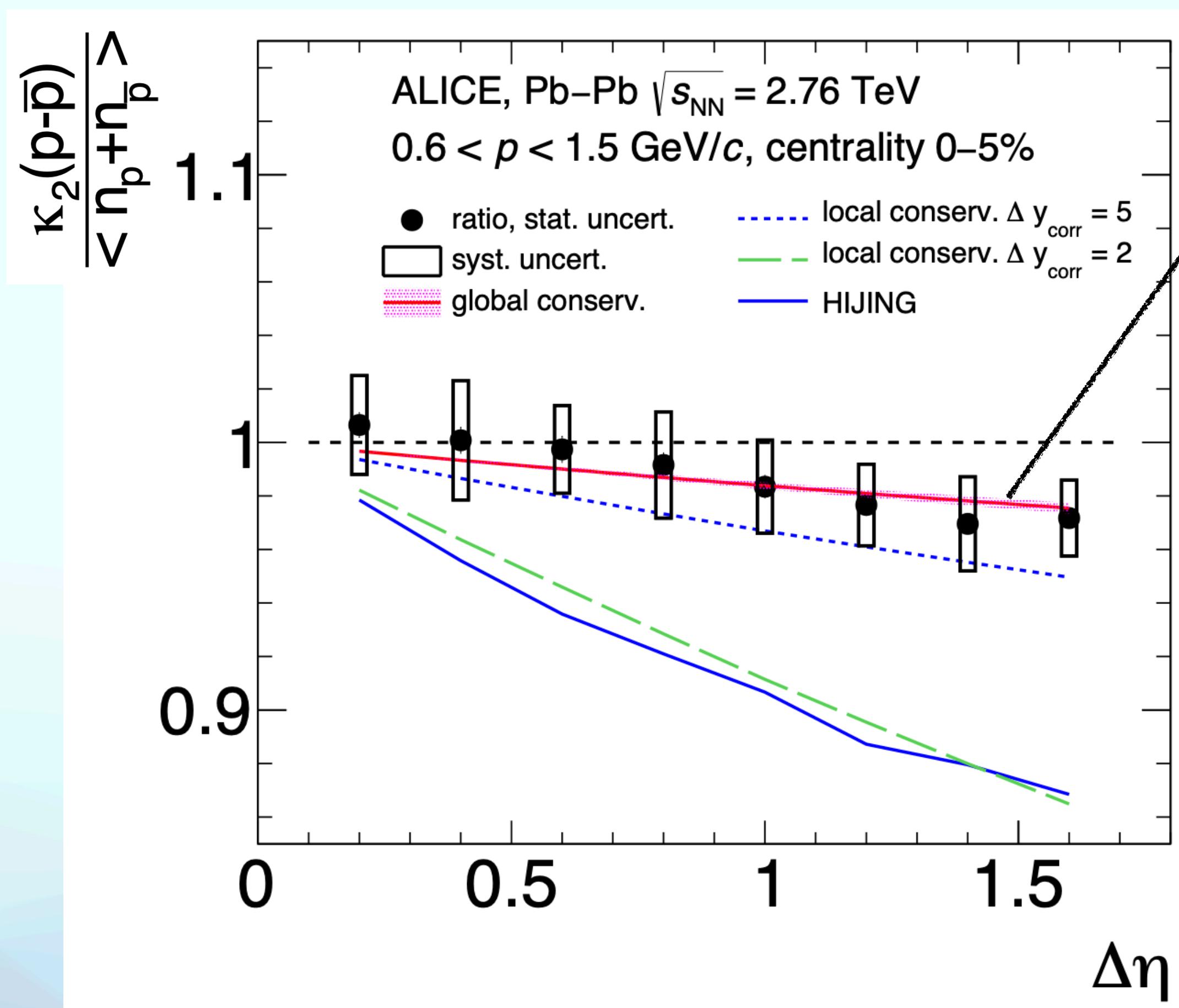
**Exact conservation**

$$\frac{\kappa_2(N_B - N_{\bar{B}})}{\langle N_B + N_{\bar{B}} \rangle} = 1$$

$$\frac{\kappa_2(N_B - N_{\bar{B}})}{\langle N_B + N_{\bar{B}} \rangle} = 1 - \alpha$$

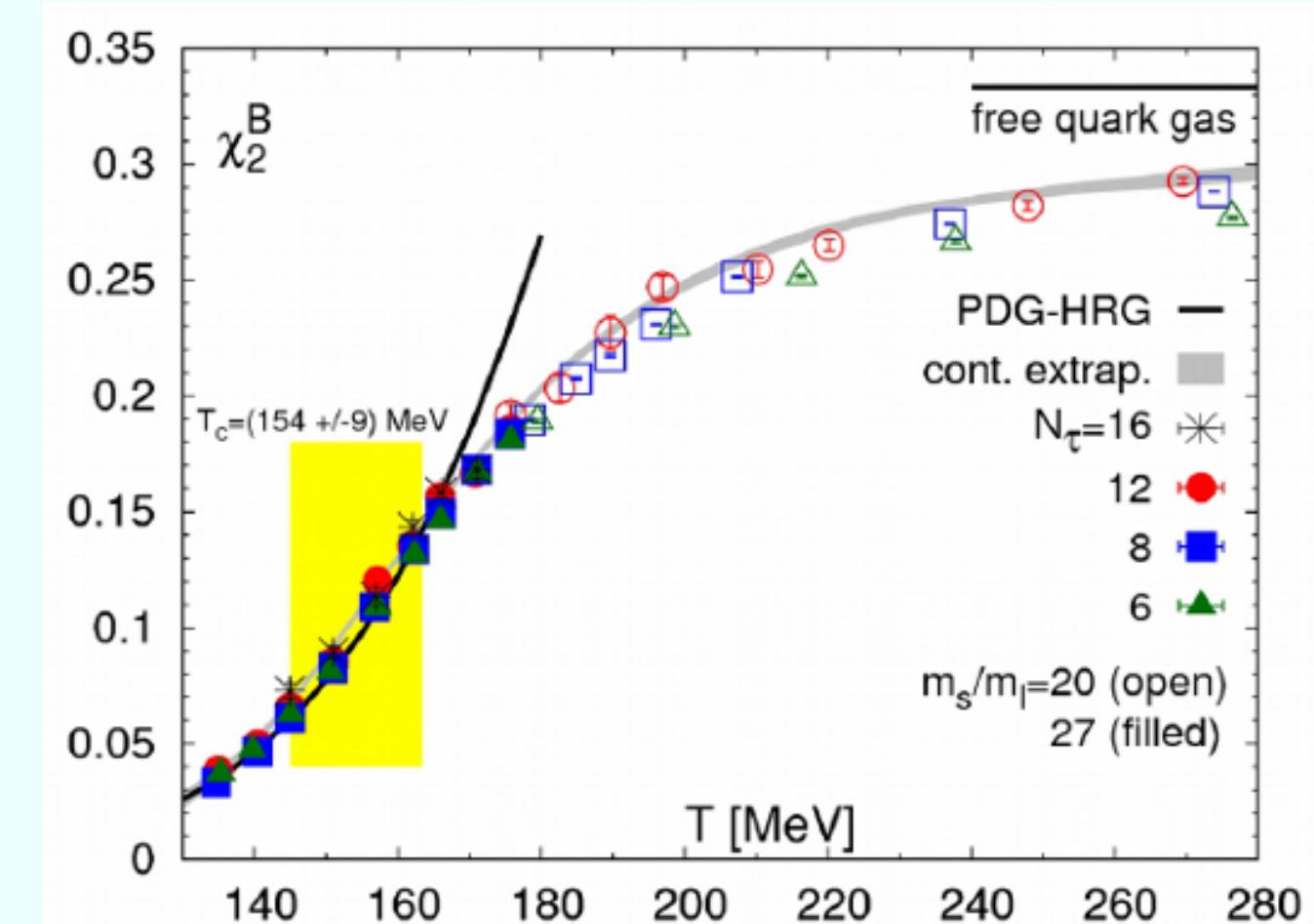


# First Alice results (Identity Method)



$$1 - \alpha$$

$$\alpha_{\Delta\eta} = \frac{\langle n_p \rangle^{acc}}{\langle N_B \rangle^{4\pi}}$$



A. Bazavov et al [HotQCD], PRD 101 (2020) 074502  
A. Bazavov et al., Phys.Rev. D85 (2012) 054503

**first verification of LQCD results**

## Consequences:

- Support for the validity of the HRG model
- Further support for freeze-out at the phase boundary

A. R., Nucl.Phys.A 967 (2017) 453-456 (QM 17)

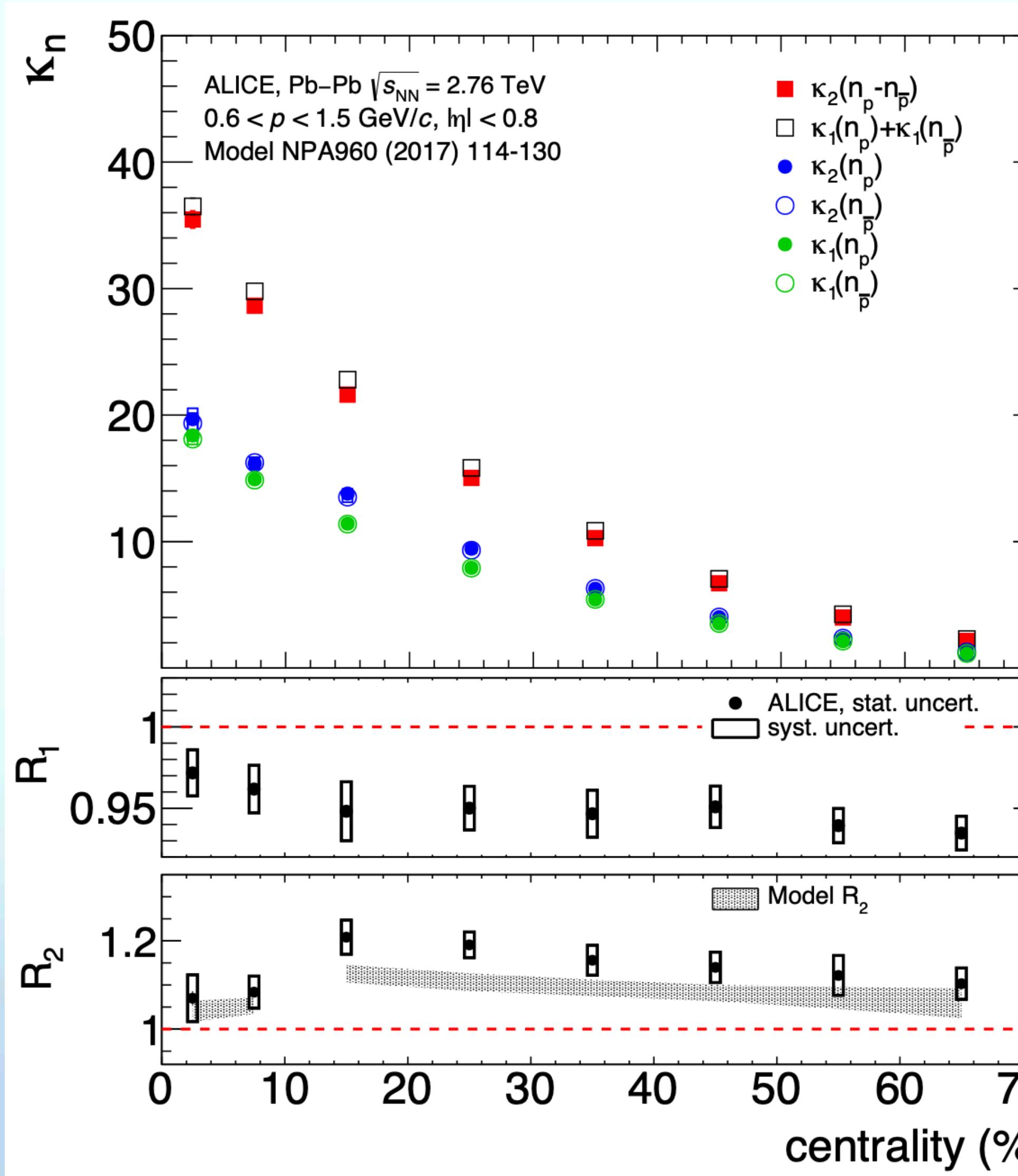
**ALICE:** Phys. Lett. B 807 (2020) 135564,  
Phys. Lett. B (2022) 137545

## Identity Method

A.R., M. I. Gorenstein, PRC 86, 044906 (2012)

M. Arslandok, A.R., NIM A946, 162622 (2019)  
A. R., Phys.Rev.C 110 (2024) 6, 064910

# First ALICE results (Identity Method)



$$\kappa_2(N - \bar{N}) = \kappa_2(n - \bar{n}) \langle N_W \rangle + \langle N - \bar{N} \rangle^2 \frac{\kappa_2(N_W)}{\langle N_W \rangle^2}$$

P. Braun-Munzinger, A. R., J. Stachel, NPA 960 (2017) 114-130

A. R., Nucl.Phys.A 967 (2017) 453-456 (QM 17)

**ALICE:** Phys. Lett. B 807 (2020) 135564,  
Phys. Lett. B (2022) 137545

## Experimental verification:

$$R_1 = \kappa_2(p - \bar{p}) / \langle n_p + n_{\bar{p}} \rangle$$

○ not influenced by volume fluctuations

$$R_2 = \kappa_2(p) / \langle n_p \rangle$$

○ affected by volume fluctuations

## Experimental challenges

- Volume fluctuations
- Conservation laws
- First ALICE results

## Experiment vs. Theory

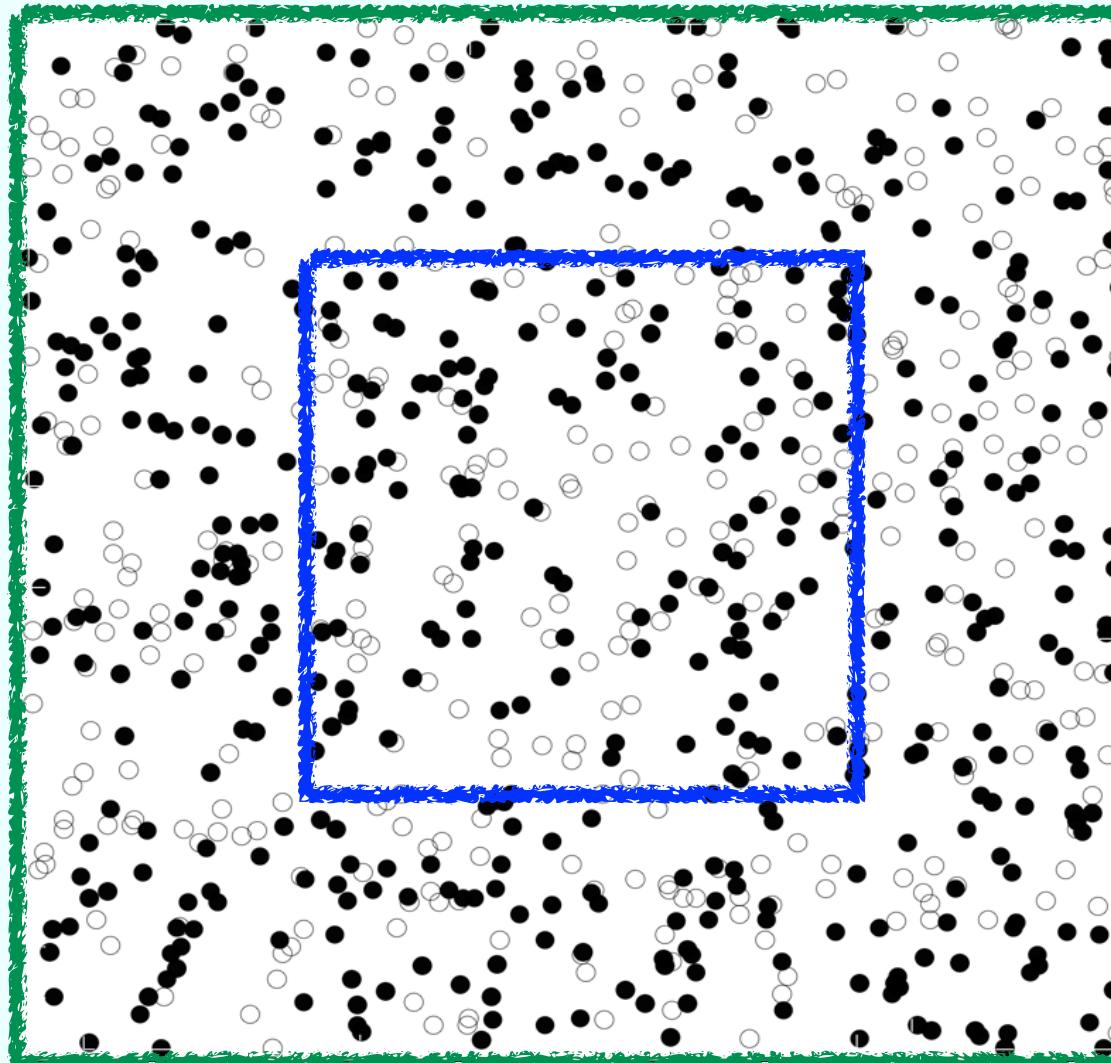
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P. Braun-Munzinger, A.R., J. Stachel, NPA 982 (2019) 307-310 (QM 18)

P. Braun-Munzinger, A.R., J. Stachel, e-Print: 1907.03032 [nucl-th] (2019)

A. R., NPA 1005 (2021) 121858(QM 19)

# Ideal gas EoS plus global baryon number conservation



- exploiting Canonical Ensemble in the full phase space
- no fluctuations in  $4\pi$  (like in experiments)

$$\frac{\kappa_2(B - \bar{B})}{\langle n_B + n_{\bar{B}} \rangle} = 1 - \frac{\alpha_B \langle n_B \rangle + \alpha_{\bar{B}} \langle n_{\bar{B}} \rangle}{\langle n_B + n_{\bar{B}} \rangle} + \left( z^2 - \langle N_B \rangle \langle N_{\bar{B}} \rangle \right) \frac{(\alpha_B - \alpha_{\bar{B}})^2}{\langle n_B + n_{\bar{B}} \rangle}$$

CE suppression

$\langle N_B \rangle, \langle N_{\bar{B}} \rangle$  - in  $4\pi$

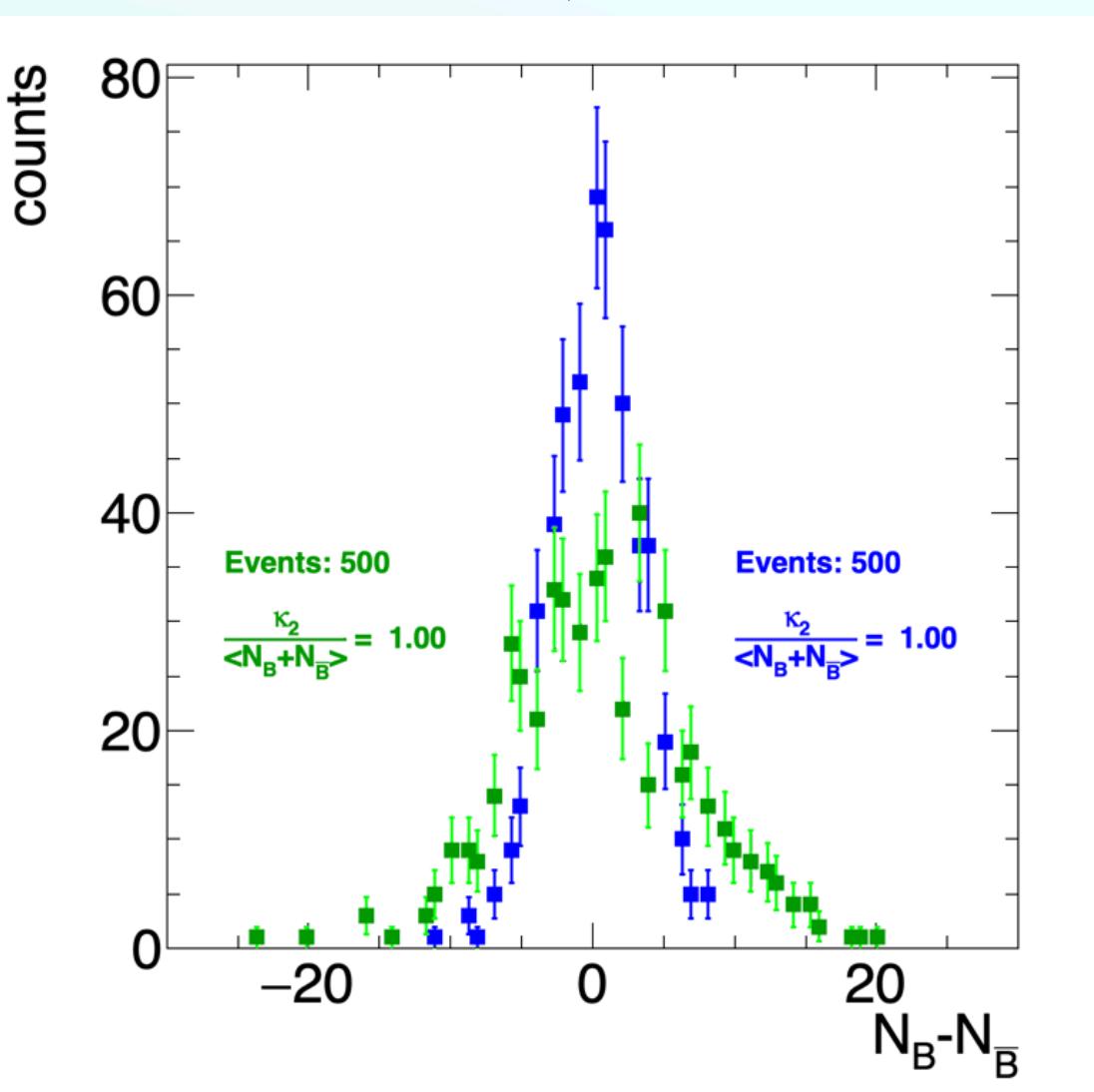
$\langle n_B \rangle, \langle n_{\bar{B}} \rangle$  - inside acceptance

$\alpha_B = \langle n_B \rangle / \langle N_B \rangle$  - acceptance for  $B$

$\alpha_{\bar{B}} = \langle n_{\bar{B}} \rangle / \langle N_{\bar{B}} \rangle$  - acceptance for  $\bar{B}$

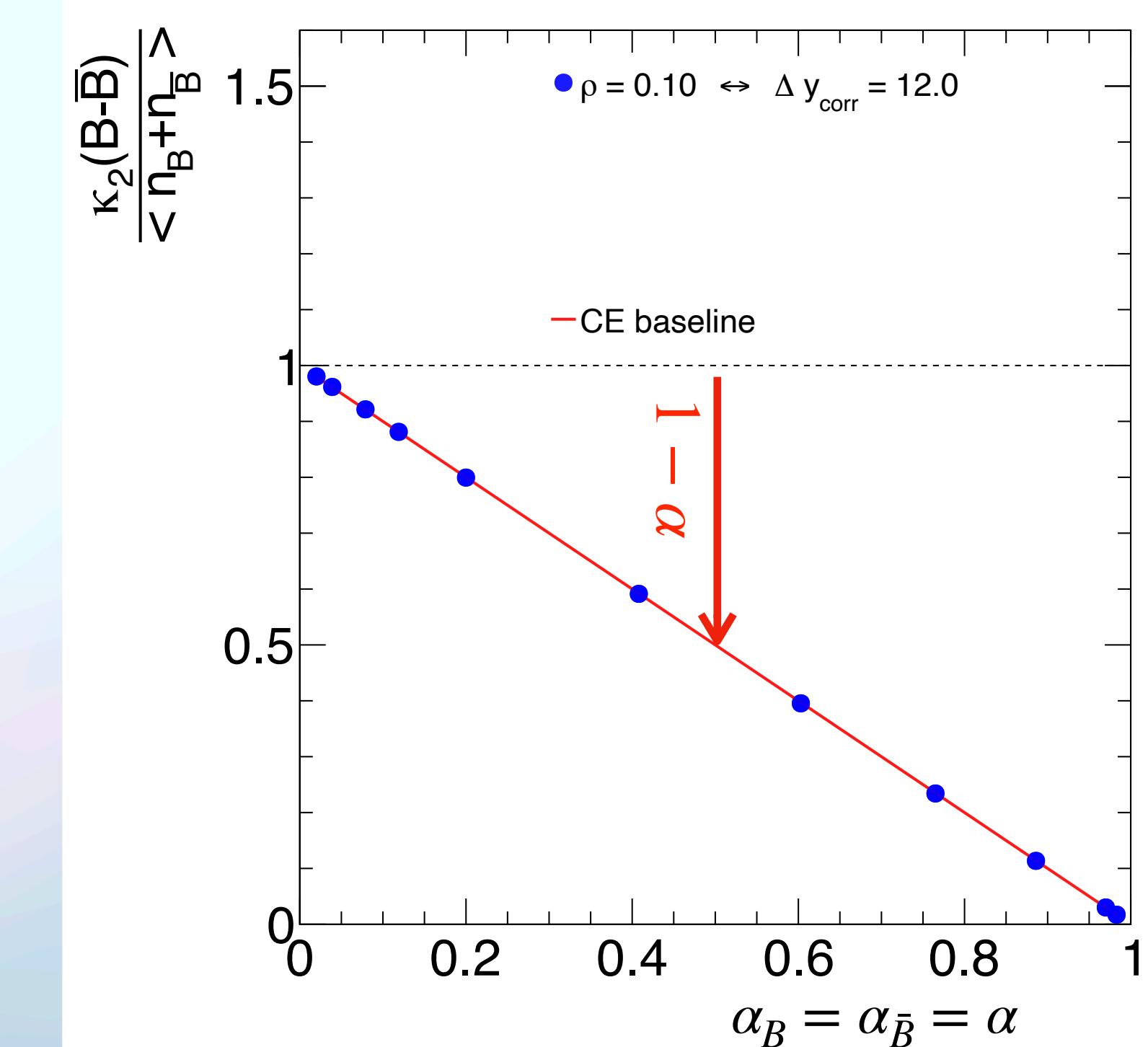
$z$  - single baryon partition function

in general:  $\alpha_B \neq \alpha_{\bar{B}}$

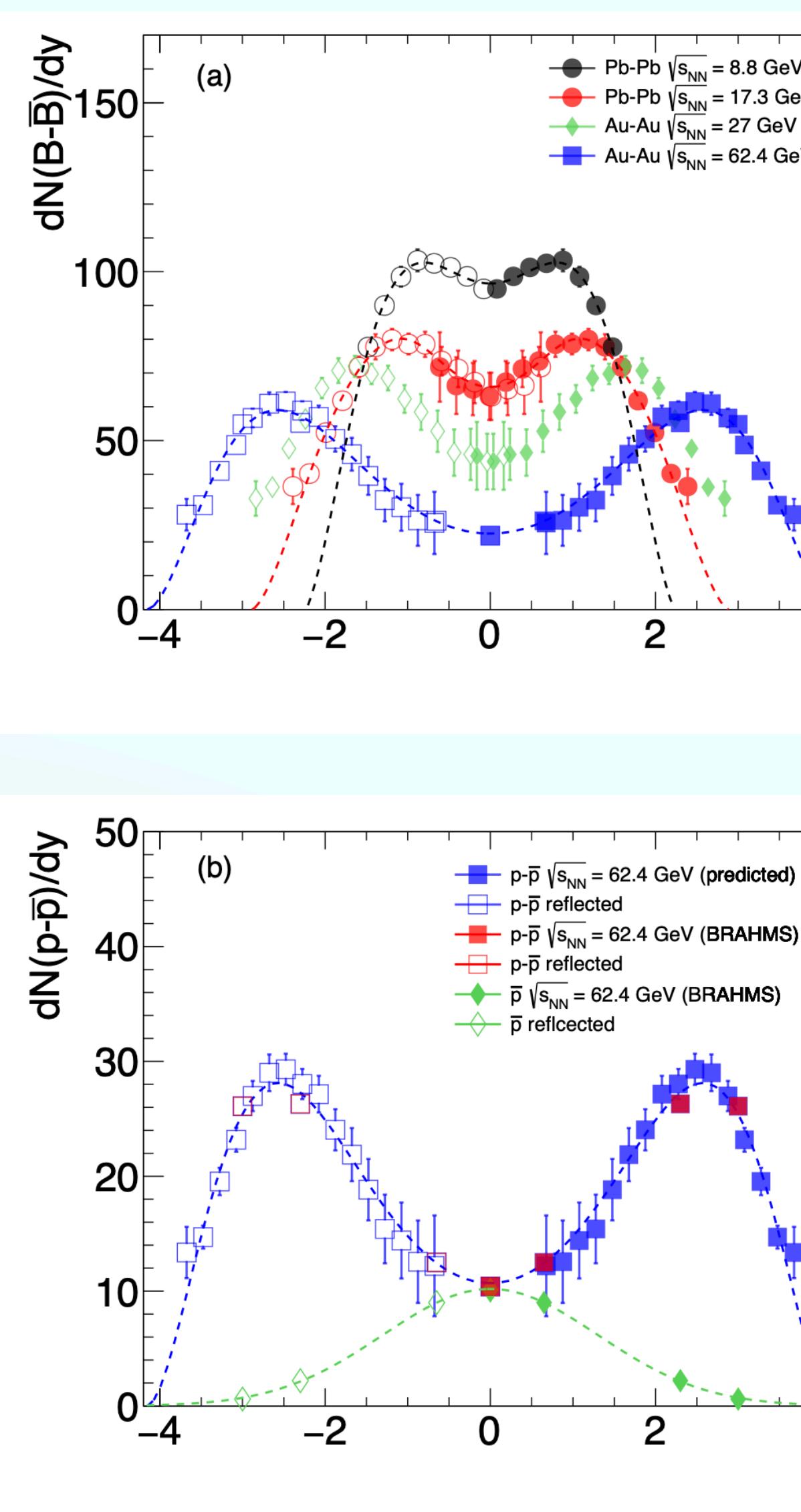


P. Braun-Munzinger, B. Friman, K. Redlich, A.R., J. Stachel , NPA 1008  
(2021) 122141

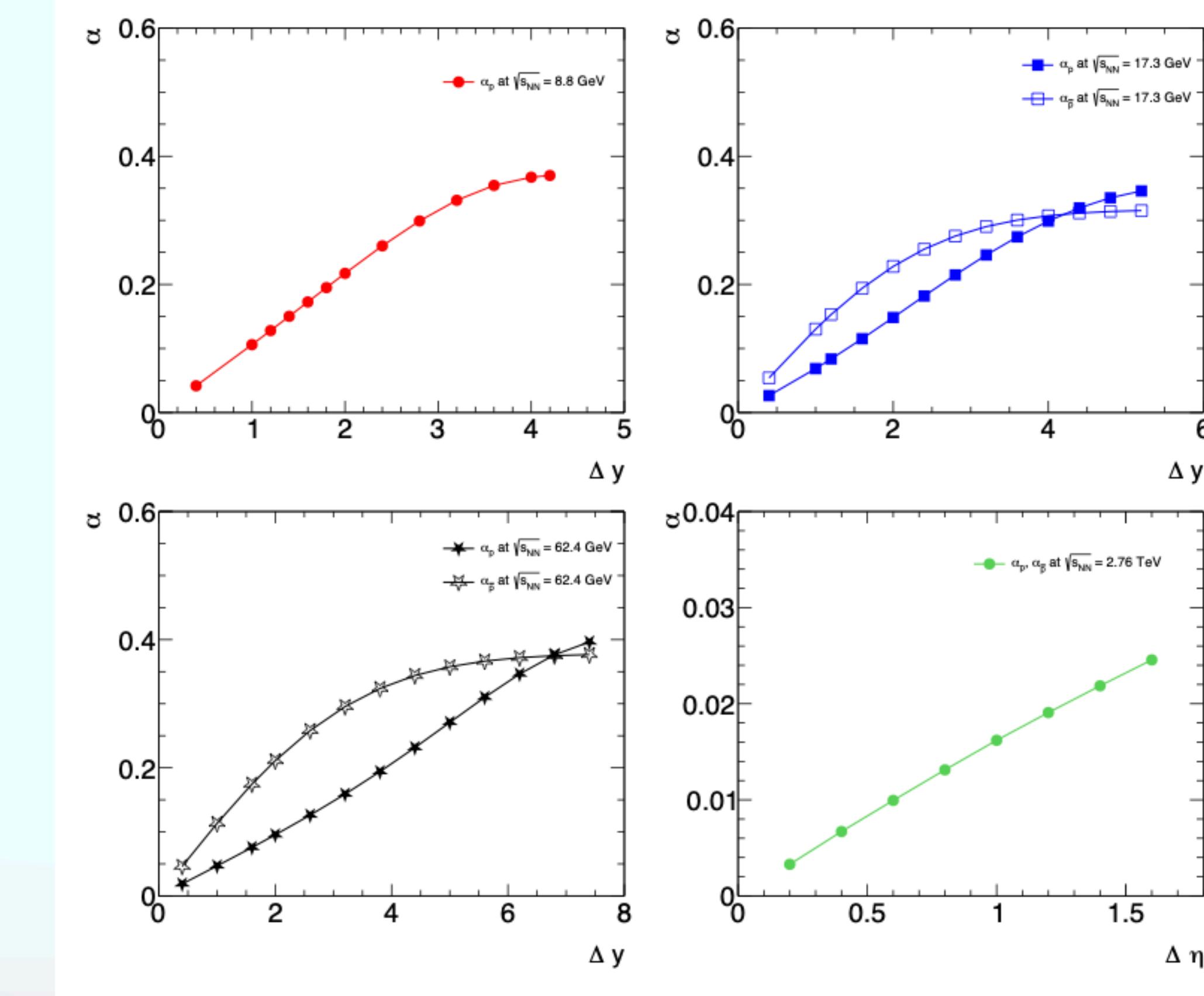
A. Bzdak, V. Koch, V. Skokov, Phys.Rev.C 87 (2013) 1, 014901



# Experimental acceptance



P. Braun-Munzinger, B. Friman, K. Redlich, A.R., J. Stachel , NPA 1008 (2021) 122141

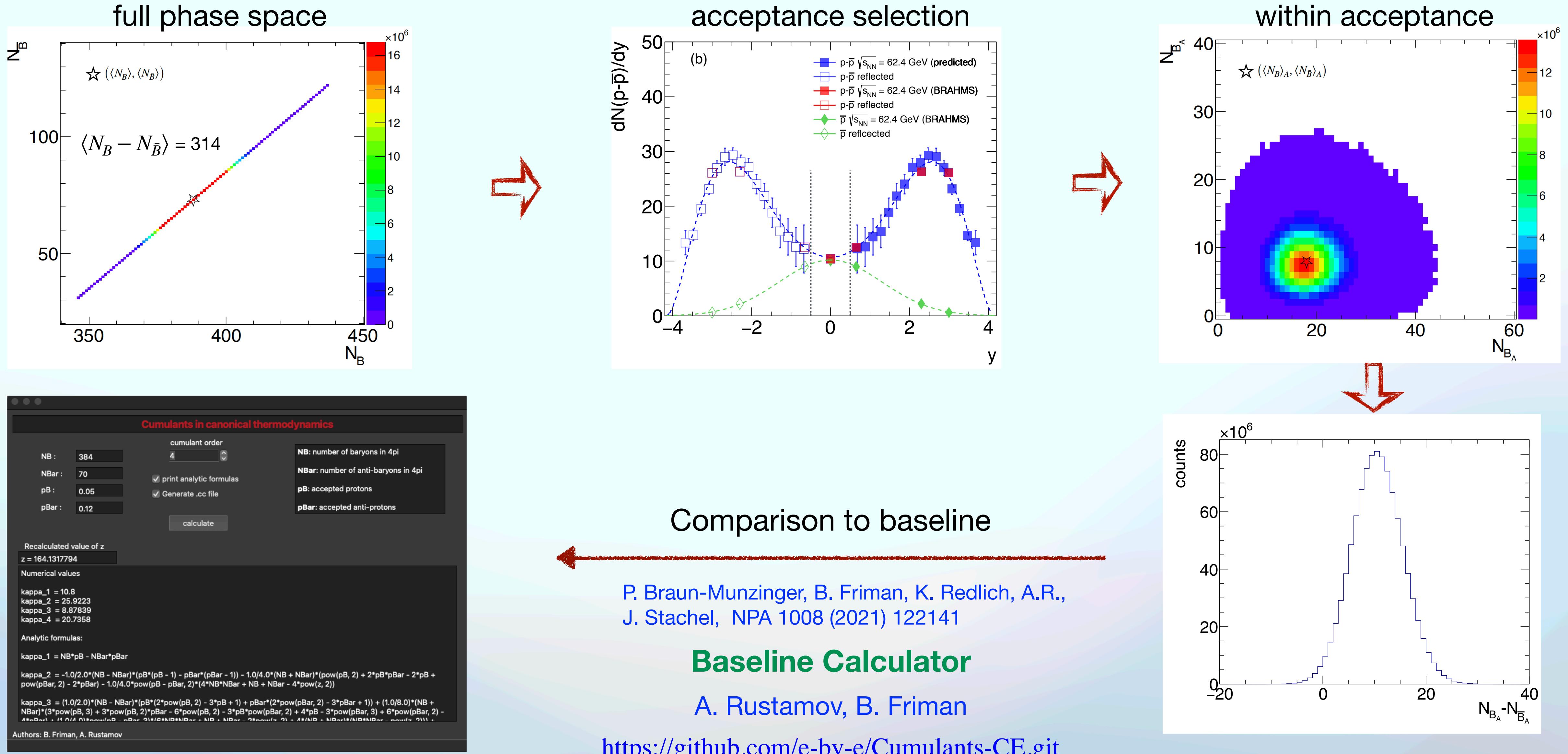


$$\alpha_p = \frac{\gamma_p \int_{y_{min}}^{y_{max}} \left[ \frac{dn_p}{dy} \right] dy}{\langle N_B \rangle}$$

$$\alpha_{\bar{p}} = \frac{\gamma_{\bar{p}} \int_{y_{min}}^{y_{max}} \left[ \frac{dn_{\bar{p}}}{dy} \right] dy}{\langle N_{\bar{B}} \rangle}$$

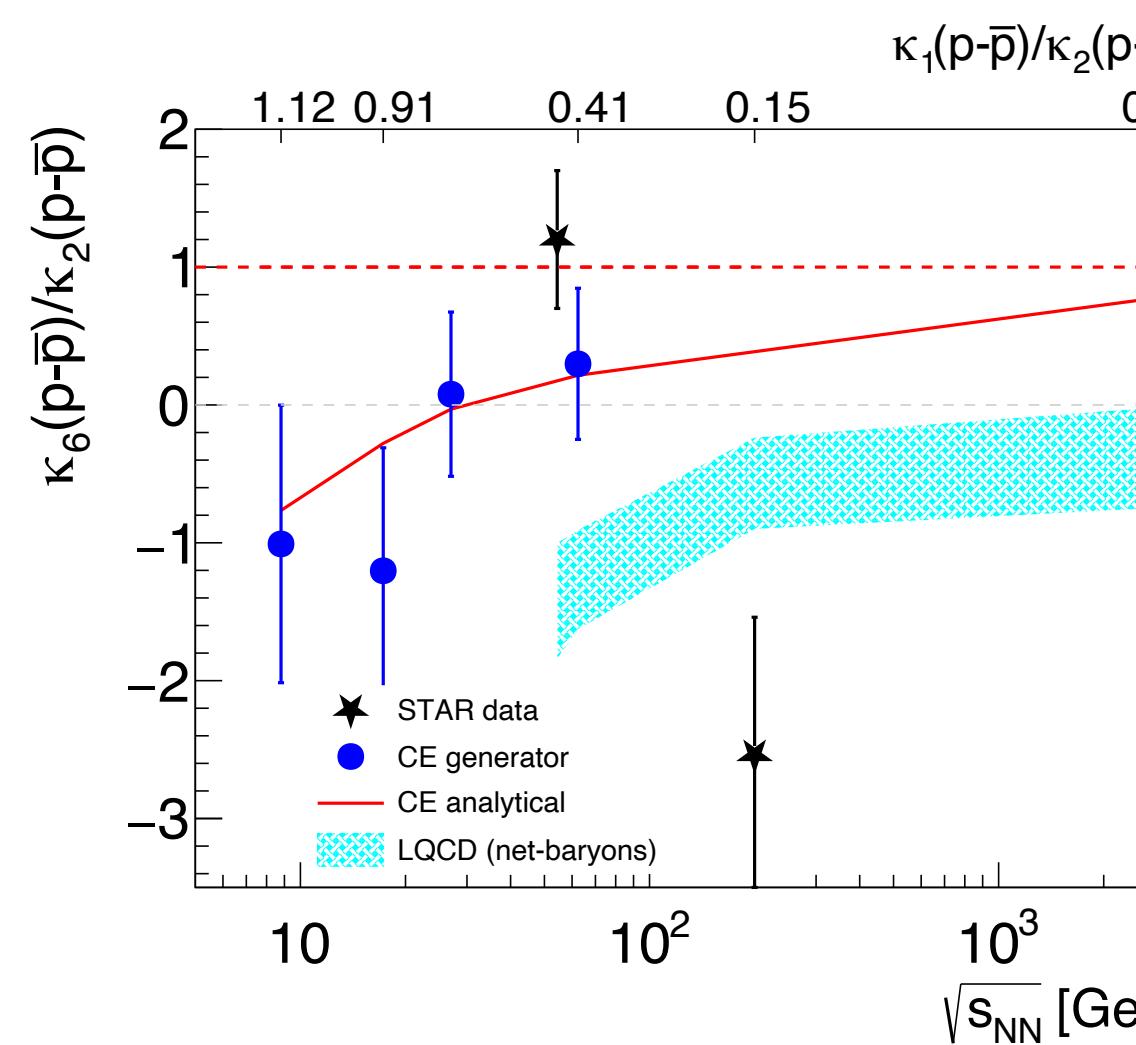
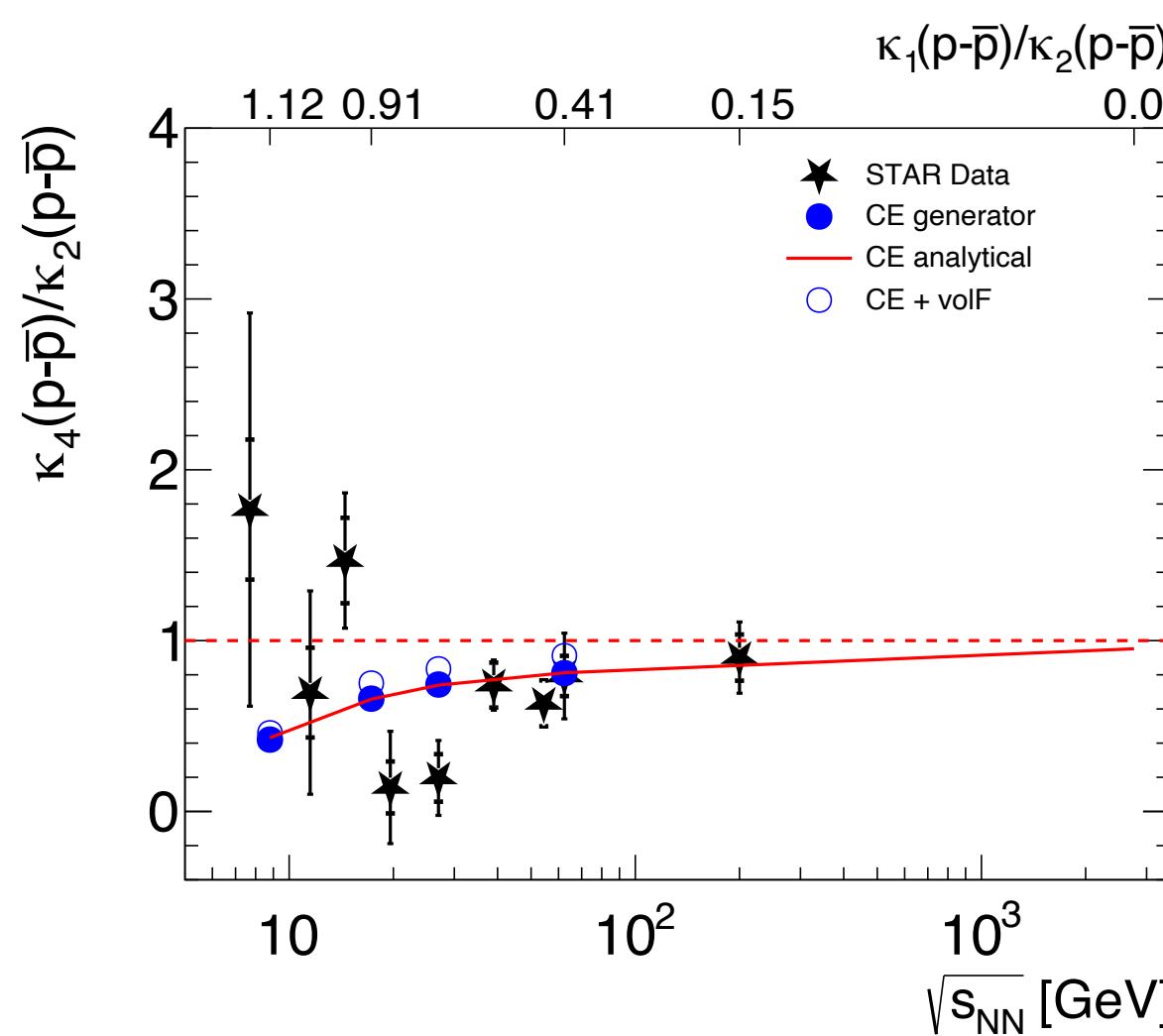
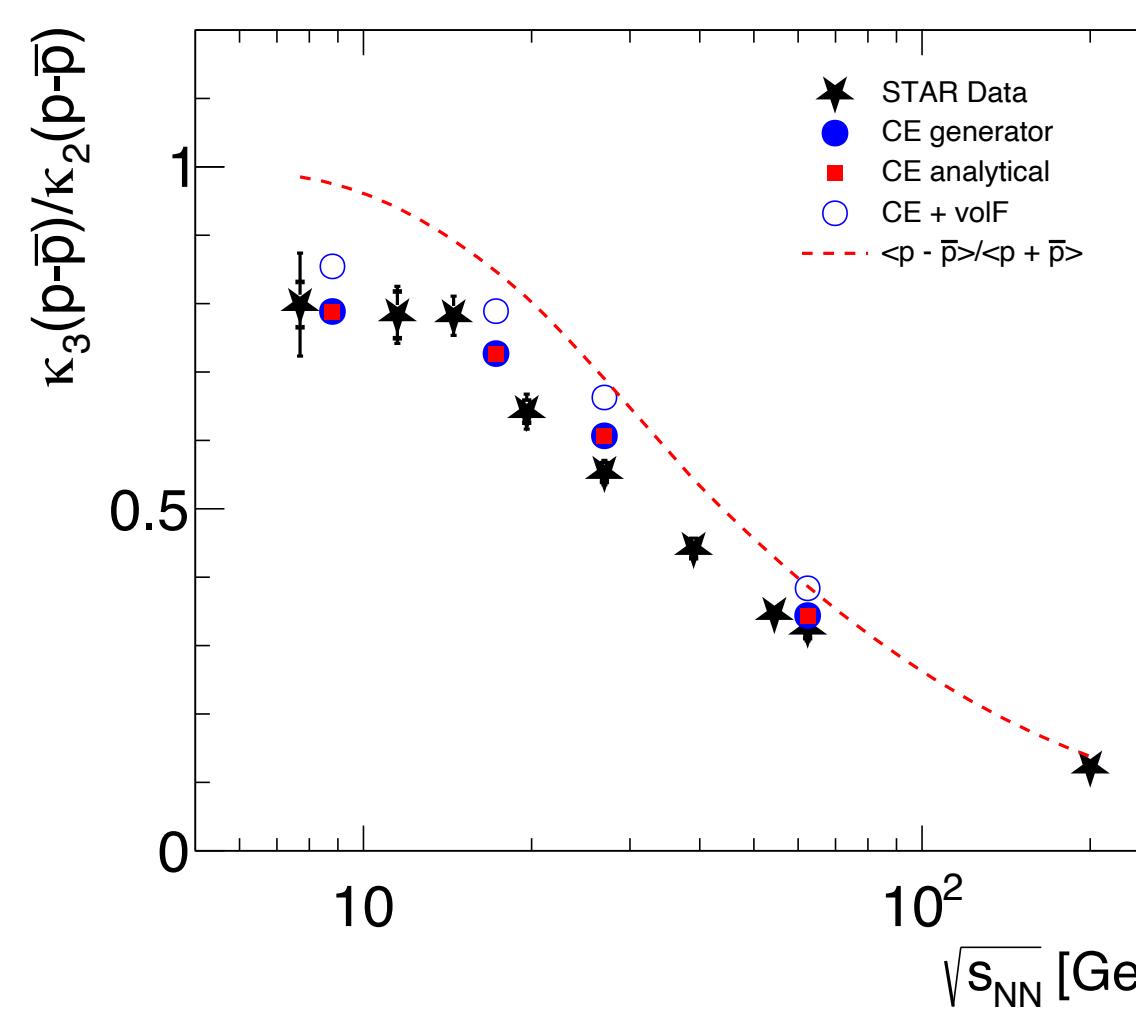
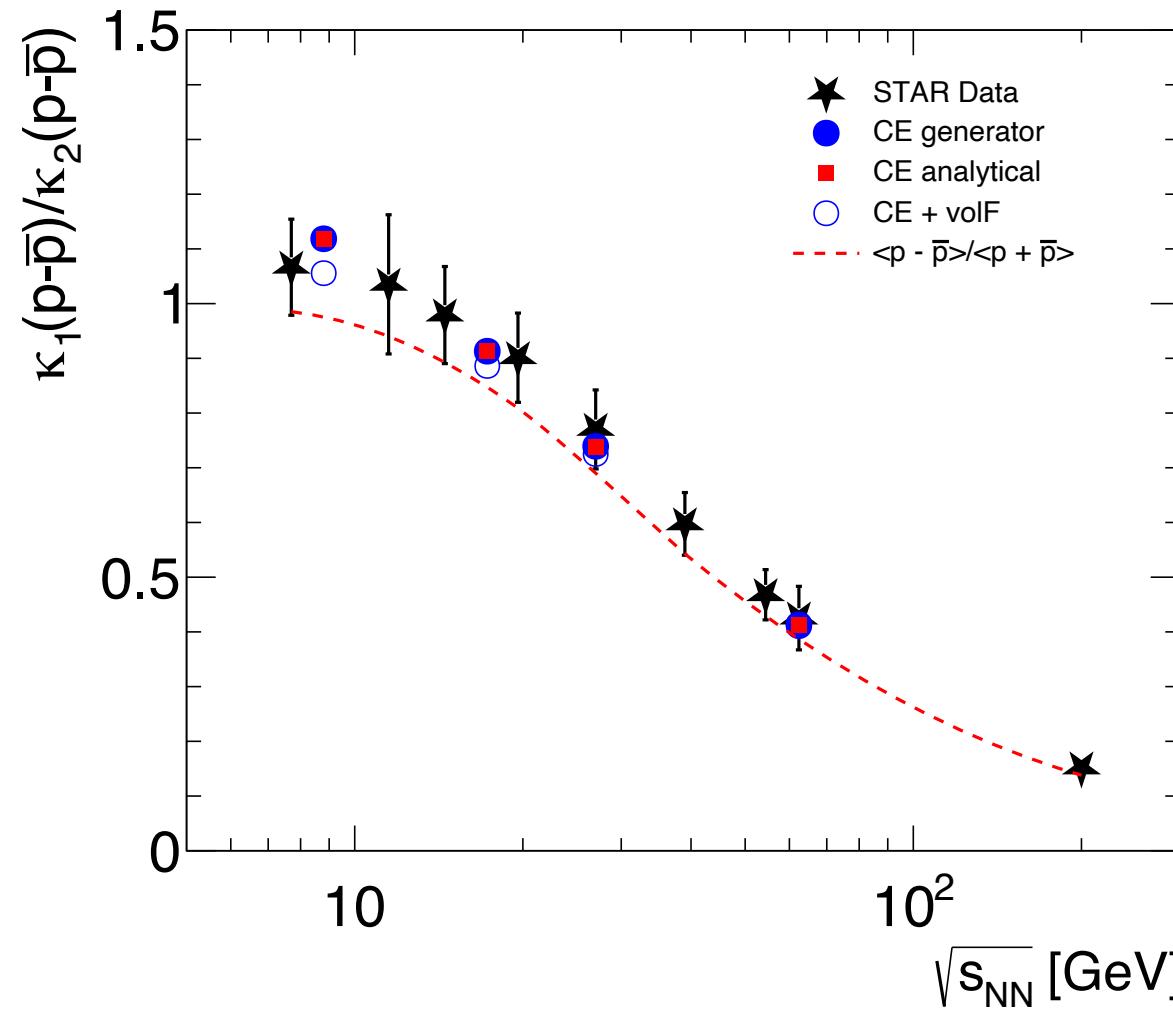
extraction of acceptances for **protons (solid symbols)**  
and **antiprotons (open symbols)**

# The strategy



# Results from STAR vs. canonical baseline

**the first quantitative and most precise canonical baselines**



remarkable agreement between canonical baseline and STAR BESI data

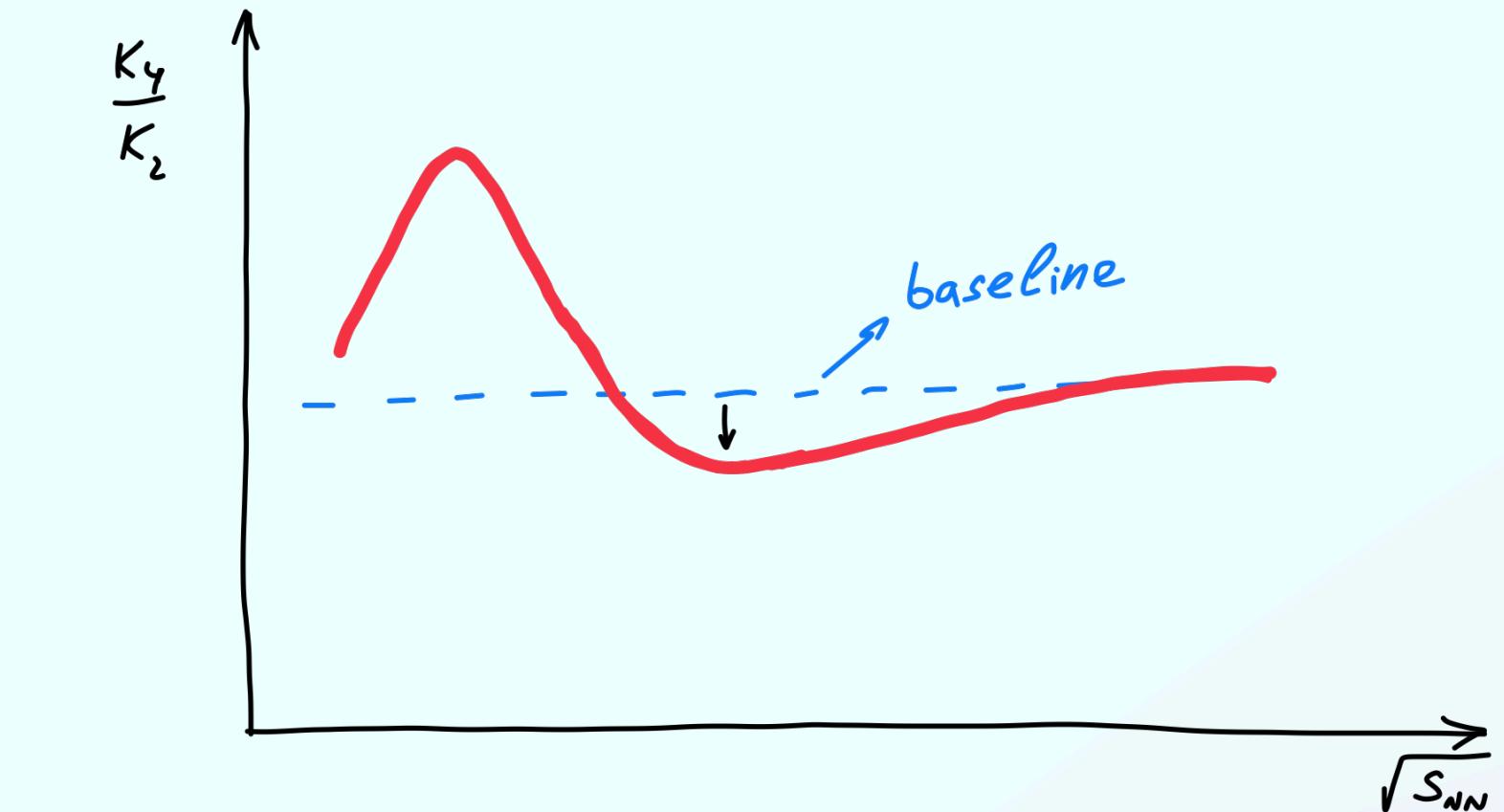
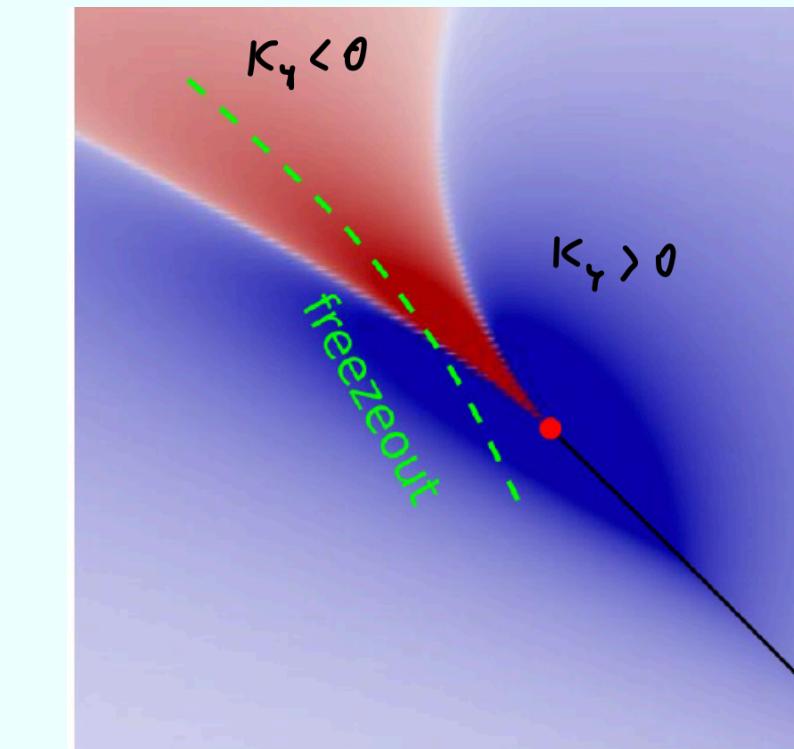
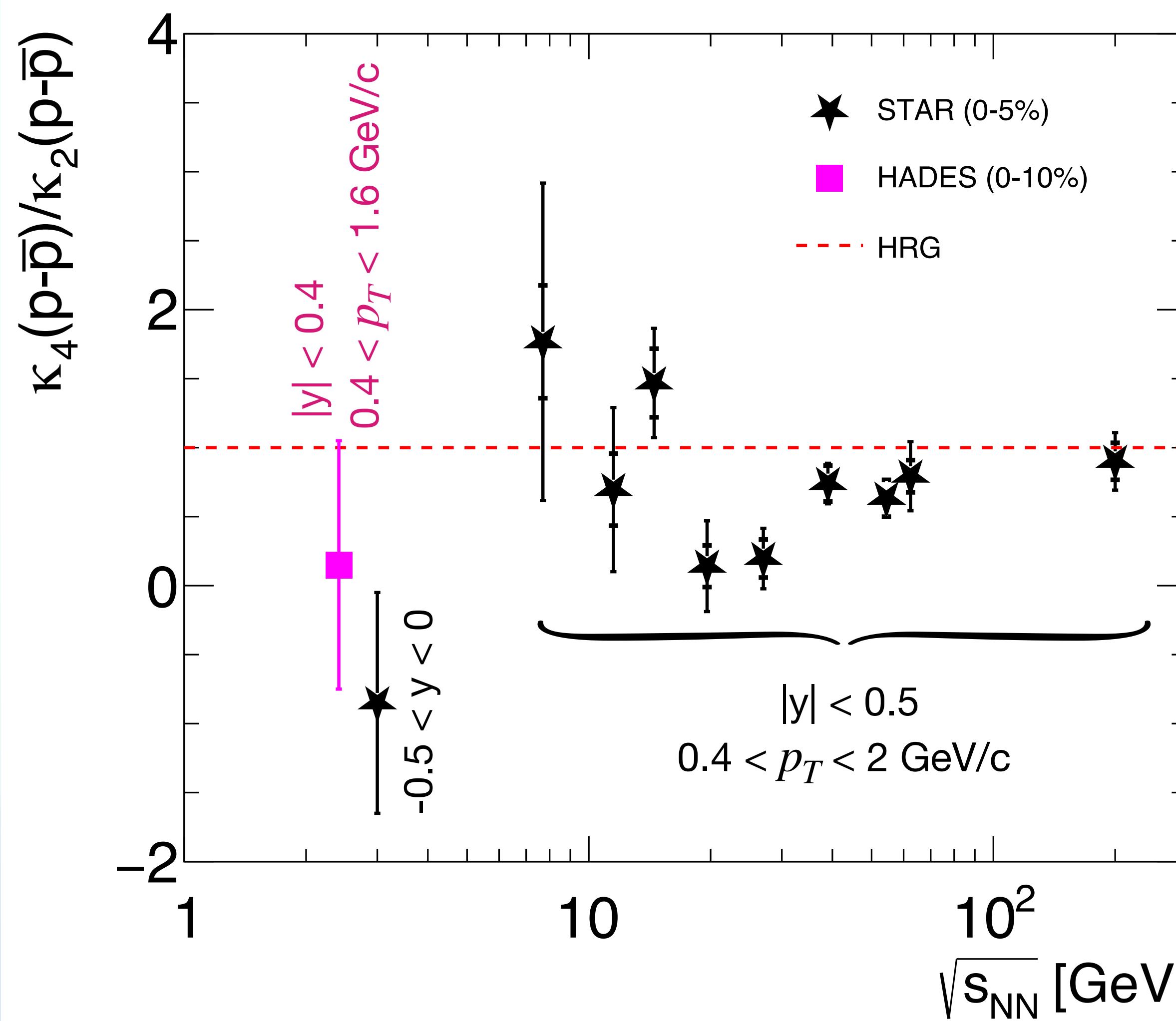
- significant reduction of canonical baseline for  $\kappa_6/\kappa_2$  going from positive values at LHC to negative values at lower energies

- STAR DATA for  $\kappa_6/\kappa_2$  is not consistent with the LQCD predictions

**STAR:** PRL 126 (2021) 9, 092301, PRL 130 (2023) 8, 082301  
P. Braun-Munzinger, B. Friman, K. Redlich, A.R., J. Stachel, NPA 1008  
(2021) 122141

# Energy excitation function of $\kappa_4/\kappa_2$ in central Au-Au collisions

HADES: Phys.Rev.C 102 (2020) 2, 024914  
 STAR: Phys.Rev.Lett. 126 (2021) 9, 092301



a dip in the excitation function is generic

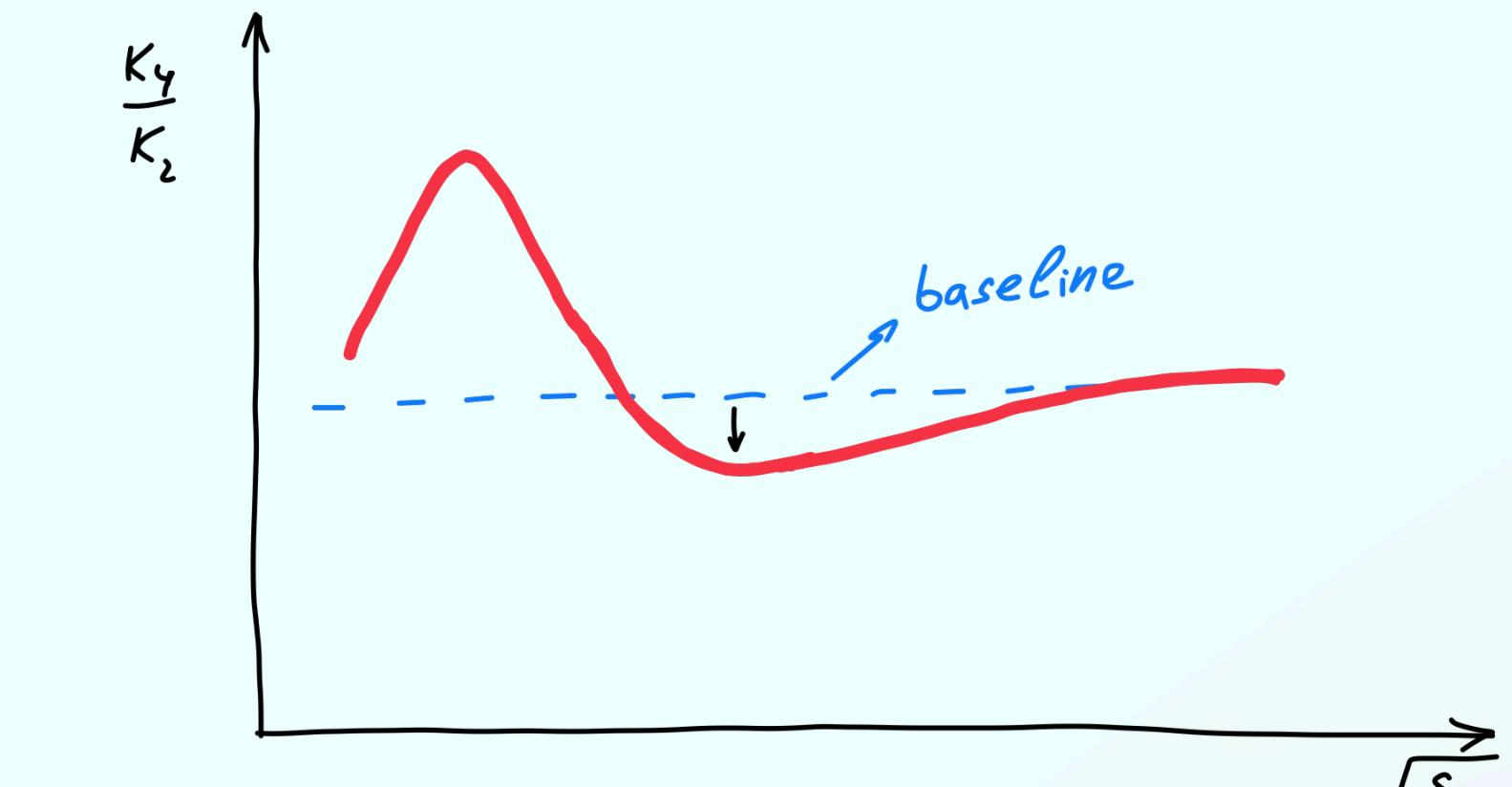
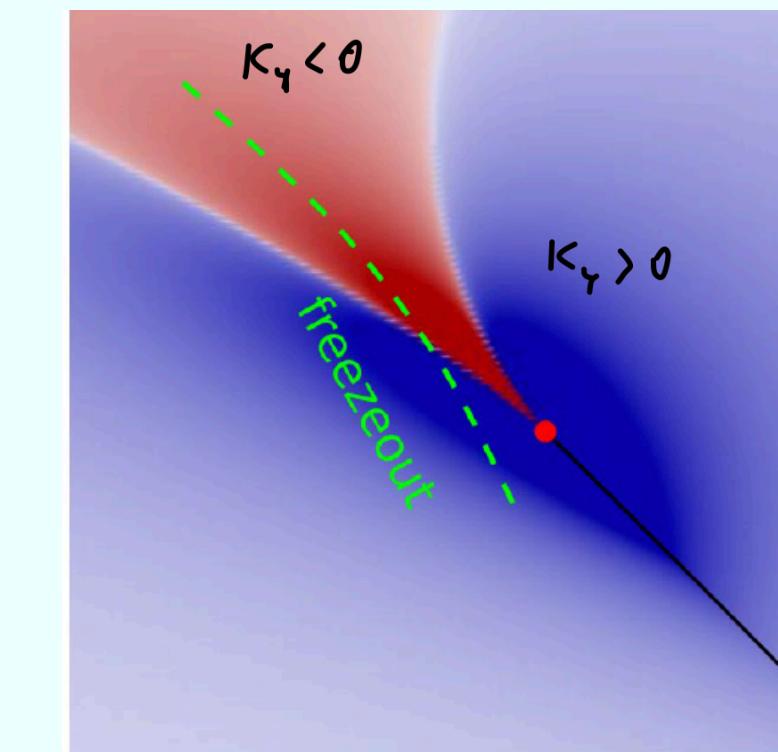
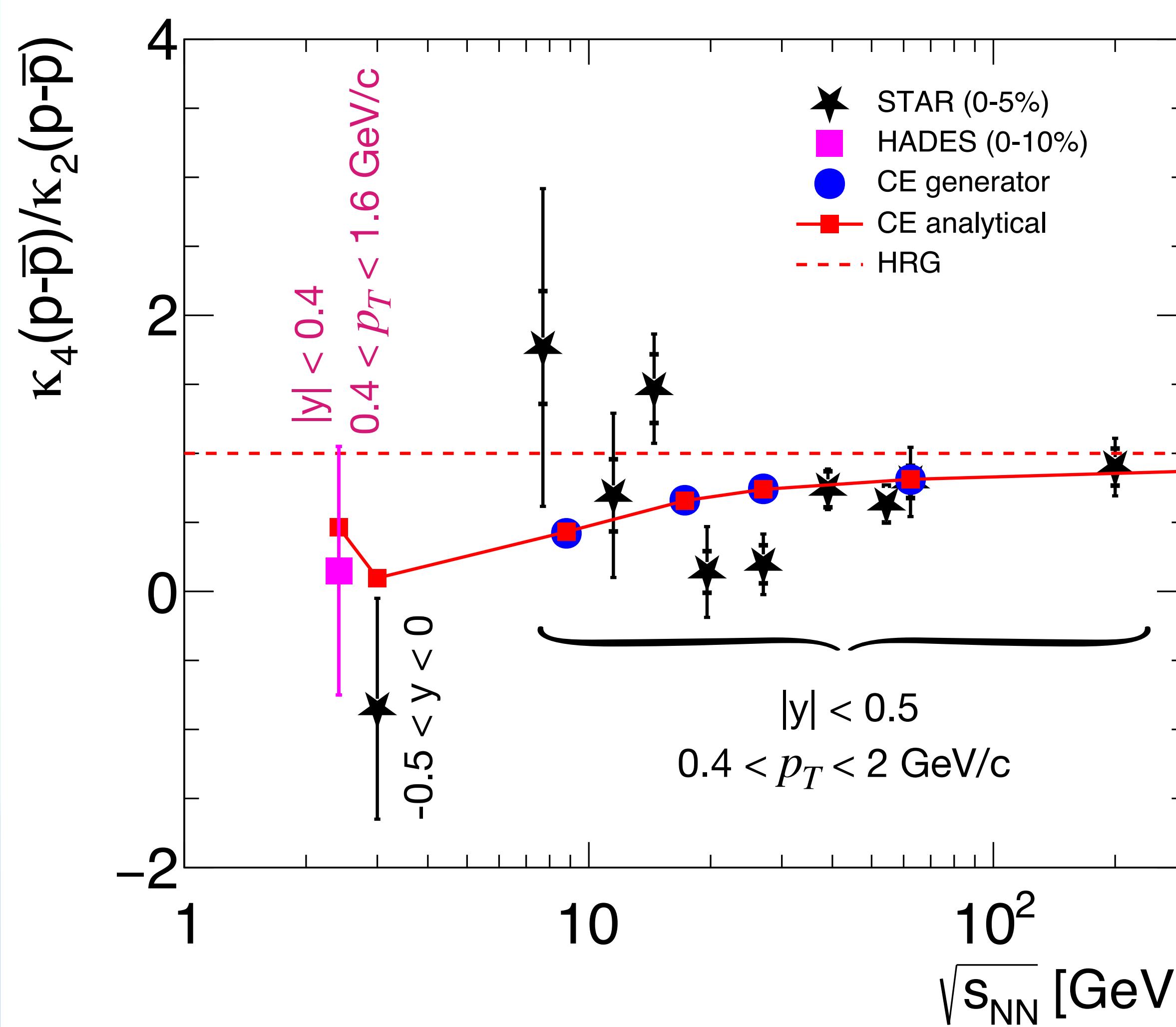
M. Stephanov, PRL102.032301(2009), PRL107.052301(2011)  
 M.Cheng et al, PRD79.074505(2009)

STAR: Phys.Rev.Lett. 126 (2021) 9, 092301

“non-monotonic behavior with a significance of  $3.1\sigma$  relative to GCE expectation”

# Energy excitation function of $\kappa_4/\kappa_2$ in central Au-Au collisions

HADES: Phys.Rev.C 102 (2020) 2, 024914  
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 M.Cheng et al, PRD79.074505(2009)

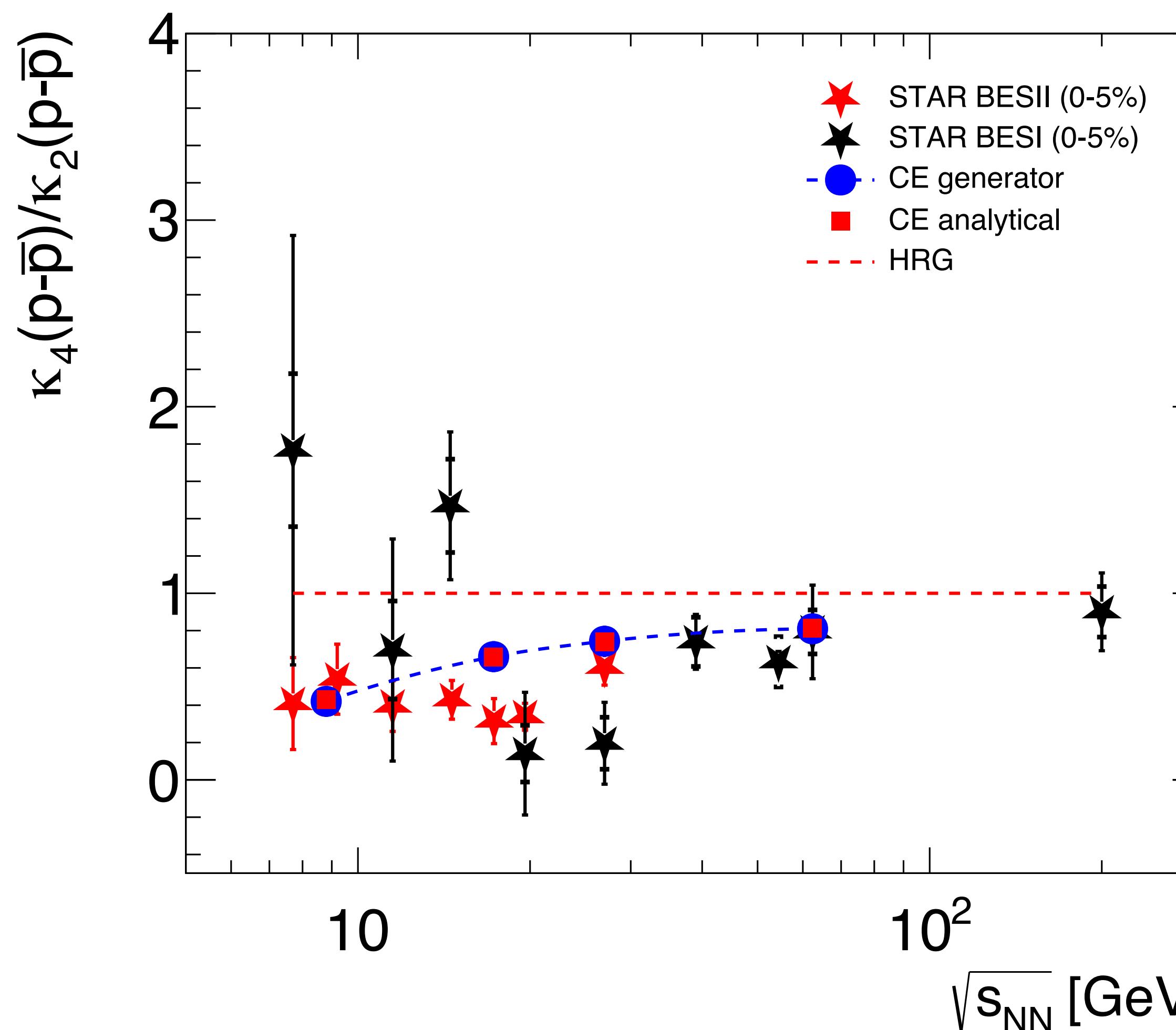
STAR: Phys.Rev.Lett. 126 (2021) 9, 092301

“non-monotonic behavior with a significance of  $3.1\sigma$   
 relative to GCE expectation”

**CE Baseline:** P. Braun-Munzinger, B. Friman, K. Redlich, A.R., J. Stachel, NPA 1008  
 (2021) 122141

no statistically significant difference between the data and  
 the canonical baseline (KS test:  $1.2\sigma$ ,  $\chi^2$  test:  $1.5\sigma$ )

# STAR BES I vs. BES II DATA, $\kappa_4/\kappa_2$

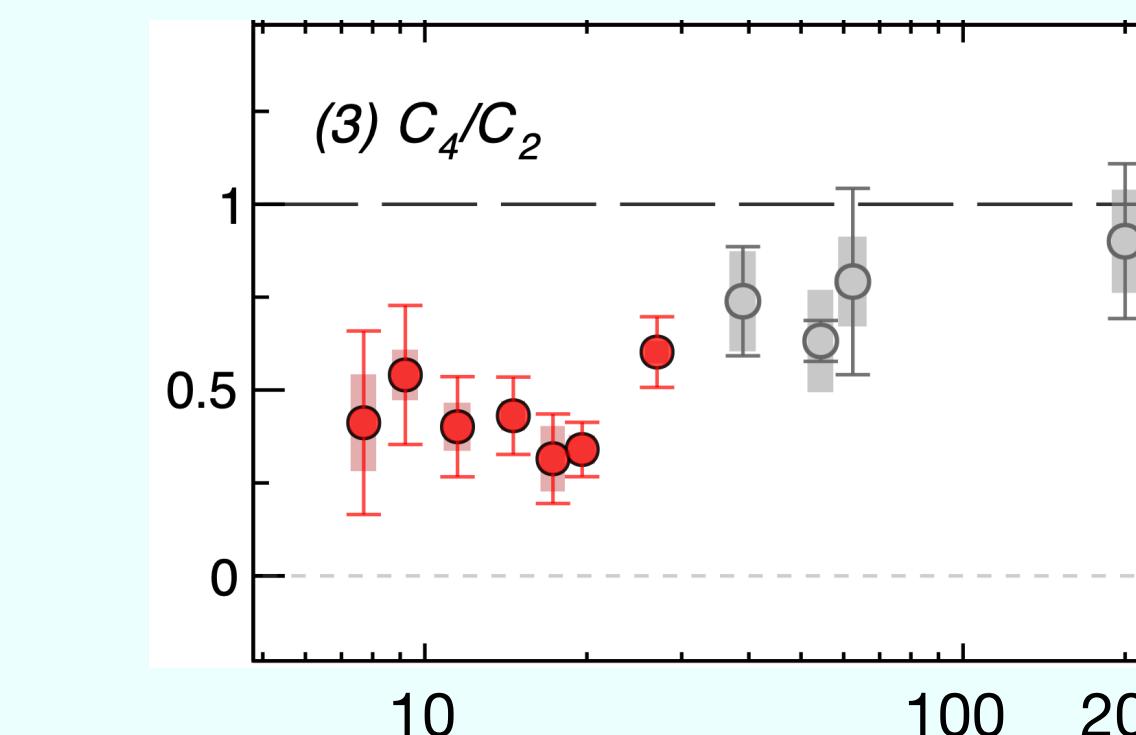


NEW STAR data points are digitized from the pdf plot!

**A. Pandav, CPOD 2024**

Note: We prefer to plot  $C_1/C_2$

Notation:  $C_i \rightarrow \kappa_i$



**The NEW data show significantly reduced uncertainties**

CE Baseline: P. Braun-Munzinger, B. Friman, K. Redlich, A.R., J. Stachel, NPA 1008 (2021) 122141

P. Braun-Munzinger, A. R., N. Xu, Annual Review of Nuclear and Particle Science (under preparation)

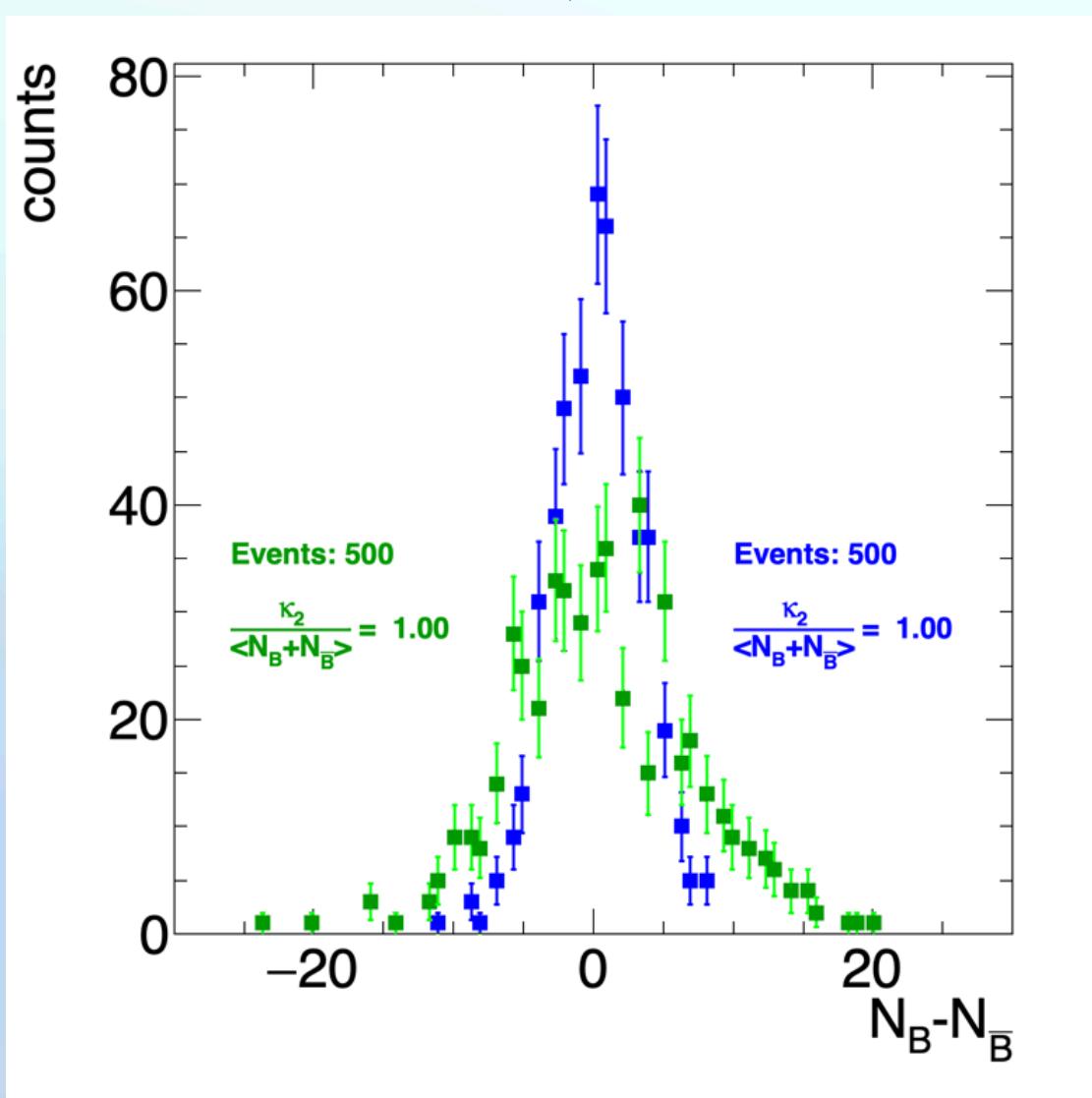
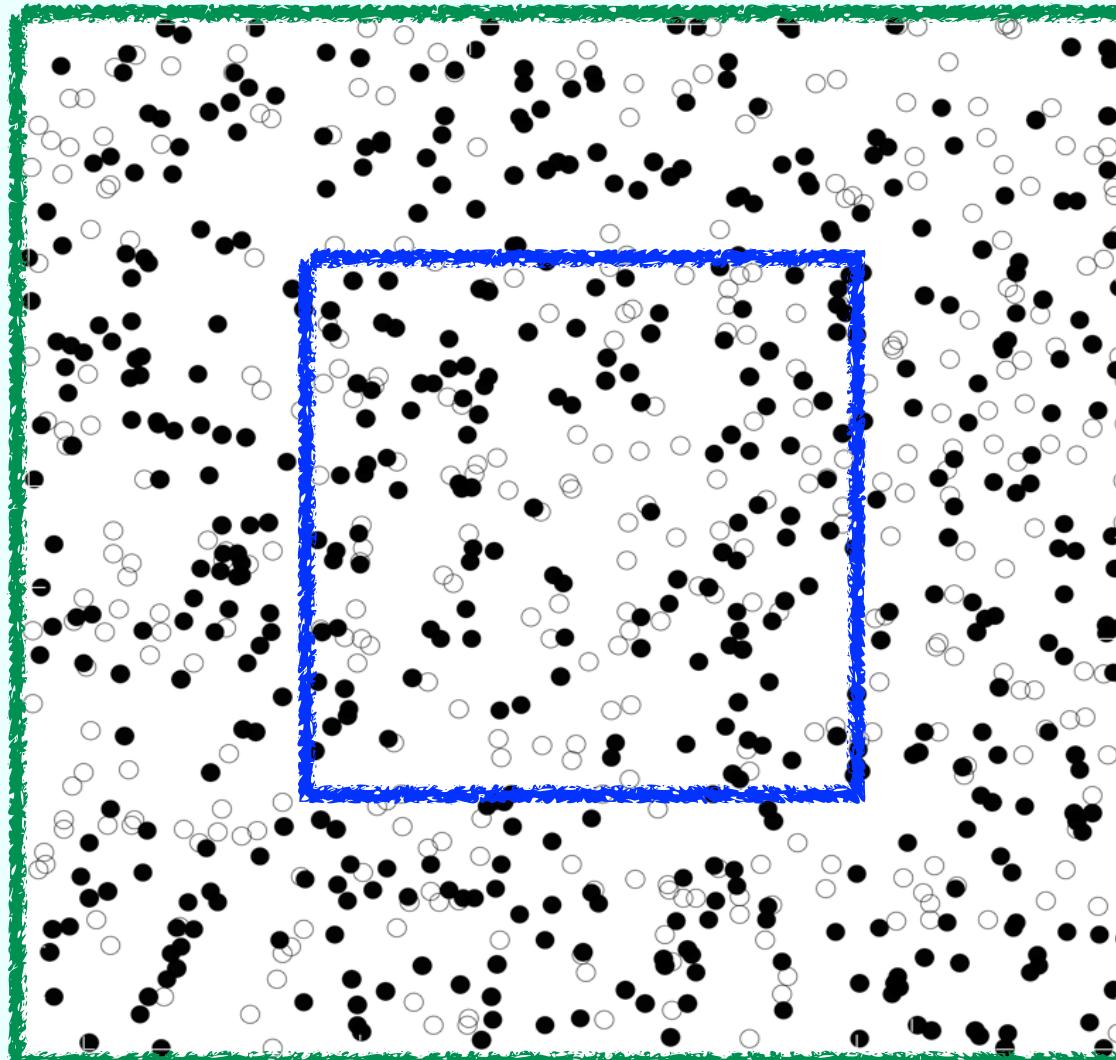
## Experimental challenges

- Volume fluctuations
- Conservation laws
- First ALICE results

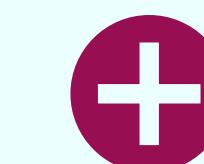
## Experiment vs. Theory

- Canonical Thermodynamics
- Comparison to STAR results
- Metropolis algorithm
- Comparison to ALICE results
- Outlook

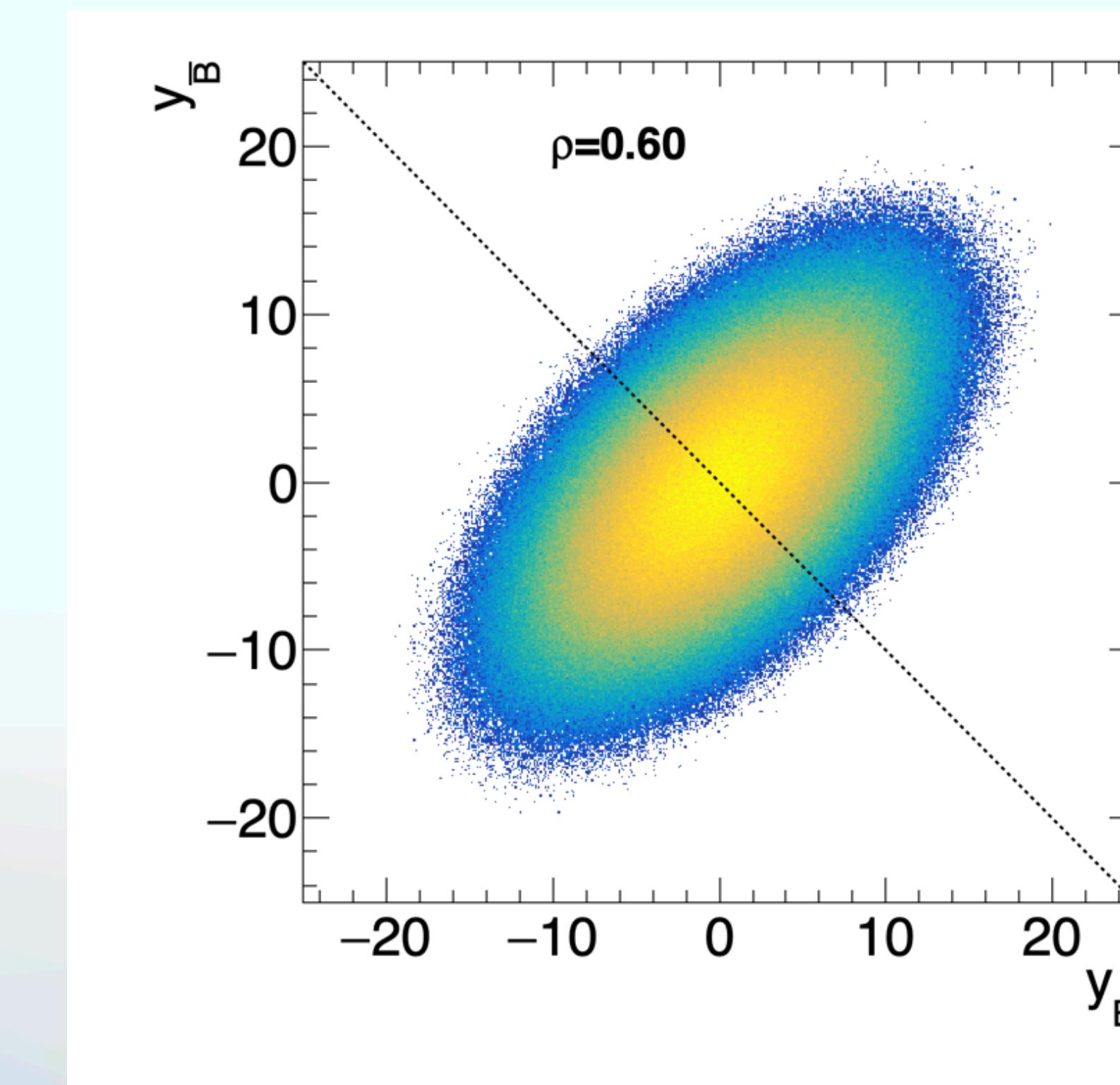
# Ideal gas EoS plus local conservation laws



- exploiting Canonical Ensemble in the full phasespac
- no fluctuations in  $4\pi$  (like in experiments)



**correlations in rapidity space (local conservations)**



# Metropolis algorithm (Simulated annealing)

start with uncorrelated  $\{y_B\}, \{y_{\bar{B}}\}$

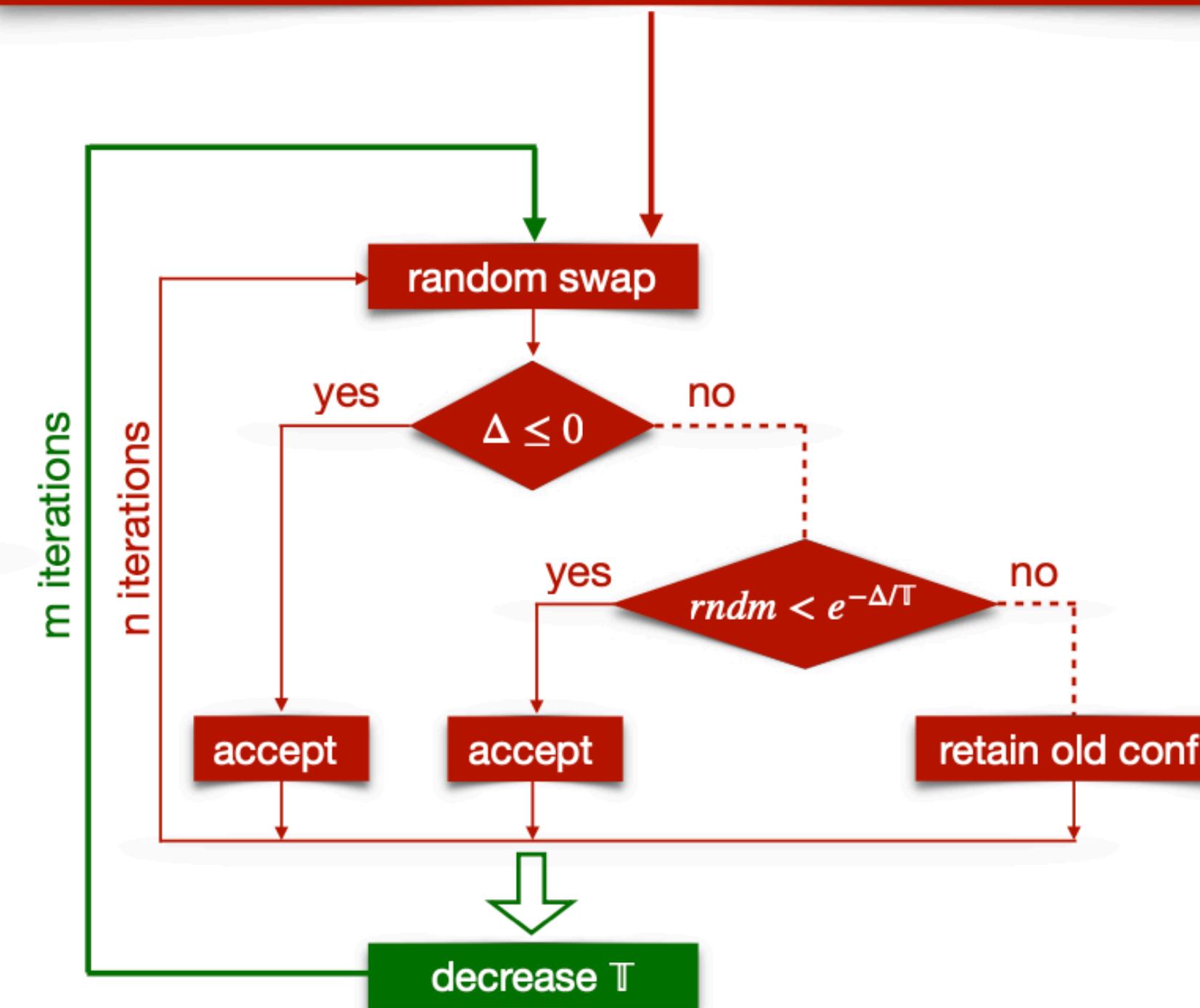
$$P_0(y^{\bar{B}}) = \begin{pmatrix} y_1^B & y_1^{\bar{B}} \\ y_2^B & y_2^{\bar{B}} \\ y_3^B & y_3^{\bar{B}} \\ y_4^B & y_4^{\bar{B}} \\ \vdots & \vdots \end{pmatrix}$$

initial  
 $\rho_0 = 0$

$$P_1(y^{\bar{B}}) = \begin{pmatrix} y_1^B & y_1^{\bar{B}} \\ y_2^B & y_2^{\bar{B}} \\ y_3^B & y_3^{\bar{B}} \\ y_4^B & y_4^{\bar{B}} \\ \vdots & \vdots \end{pmatrix}$$

$3 \leftrightarrow 1$   
 $\rho_1 \neq 0$

iteratively swap  $\{y_{\bar{B}}\}$ , start with the high value of temperature  $T$



$$\rho_n = \frac{\text{cov}[y_B, P_n(y_{\bar{B}})]}{\sigma_{y_B} \sigma_{y_{\bar{B}}}}$$

$$\Delta = |\rho_n - \rho| - |\rho_{n-1} - \rho|$$

$\rho$ : desired corr. coefficient

works for arbitrary rapidity distributions

A. R., P. Braun-Munzinger, J. Stachel, QM 2022

P. Braun-Munzinger, K. Redlich, A. R., J. Stachel, JHEP 08 (2024) 113

# Details of implementation

$$Z_B(V, T) = \sum_{N_B=0}^{\infty} \sum_{N_{\bar{B}}=0}^{\infty} \frac{(\lambda_B z_B)^{N_B}}{N_B!} \frac{(\lambda_{\bar{B}} z_{\bar{B}})^{N_{\bar{B}}}}{N_{\bar{B}}!} \delta(N_B - N_{\bar{B}} - B) = \left( \frac{\lambda_B z_B}{\lambda_{\bar{B}} z_{\bar{B}}} \right)^{\frac{B}{2}} I_B(2z \sqrt{\lambda_B \lambda_{\bar{B}}})$$

$B$  net baryon number, conserved in each event

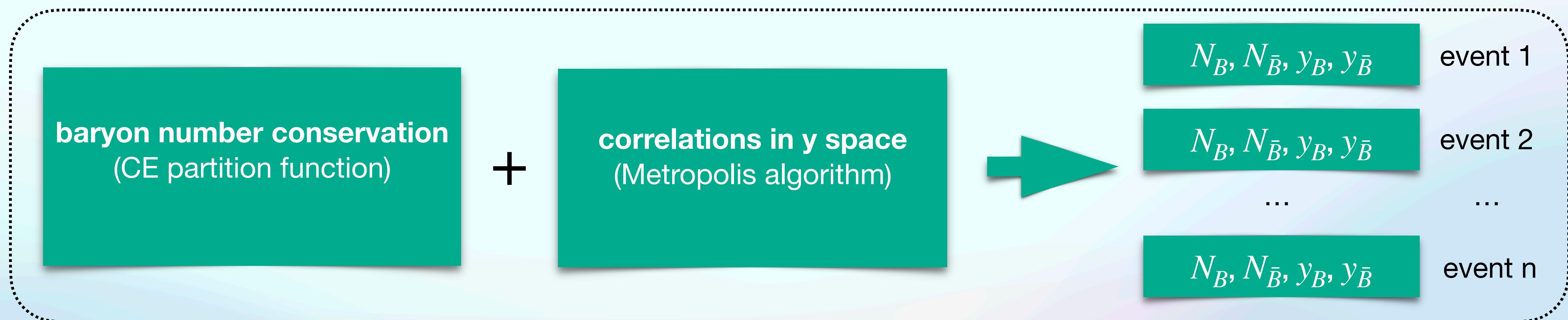
A. R., P. Braun-Munzinger, J. Stachel, QM 2022

$I_B$  modified Bessel function of the first kind

P. Braun-Munzinger, K. Redlich, A. R., J. Stachel, JHEP 08 (2024) 113

$z_B, z_{\bar{B}}$  single particle partition functions for baryons, anti baryons

$\lambda_B, \lambda_{\bar{B}}$  auxiliary parameters for calculating cumulants of baryons, anti baryons



## Input from experiments

- baryon rapidity distributions

- measured (canonical)  $\langle N_B \rangle, \langle N_{\bar{B}} \rangle$

$z = \sqrt{z_B z_{\bar{B}}}$  is calculated by solving

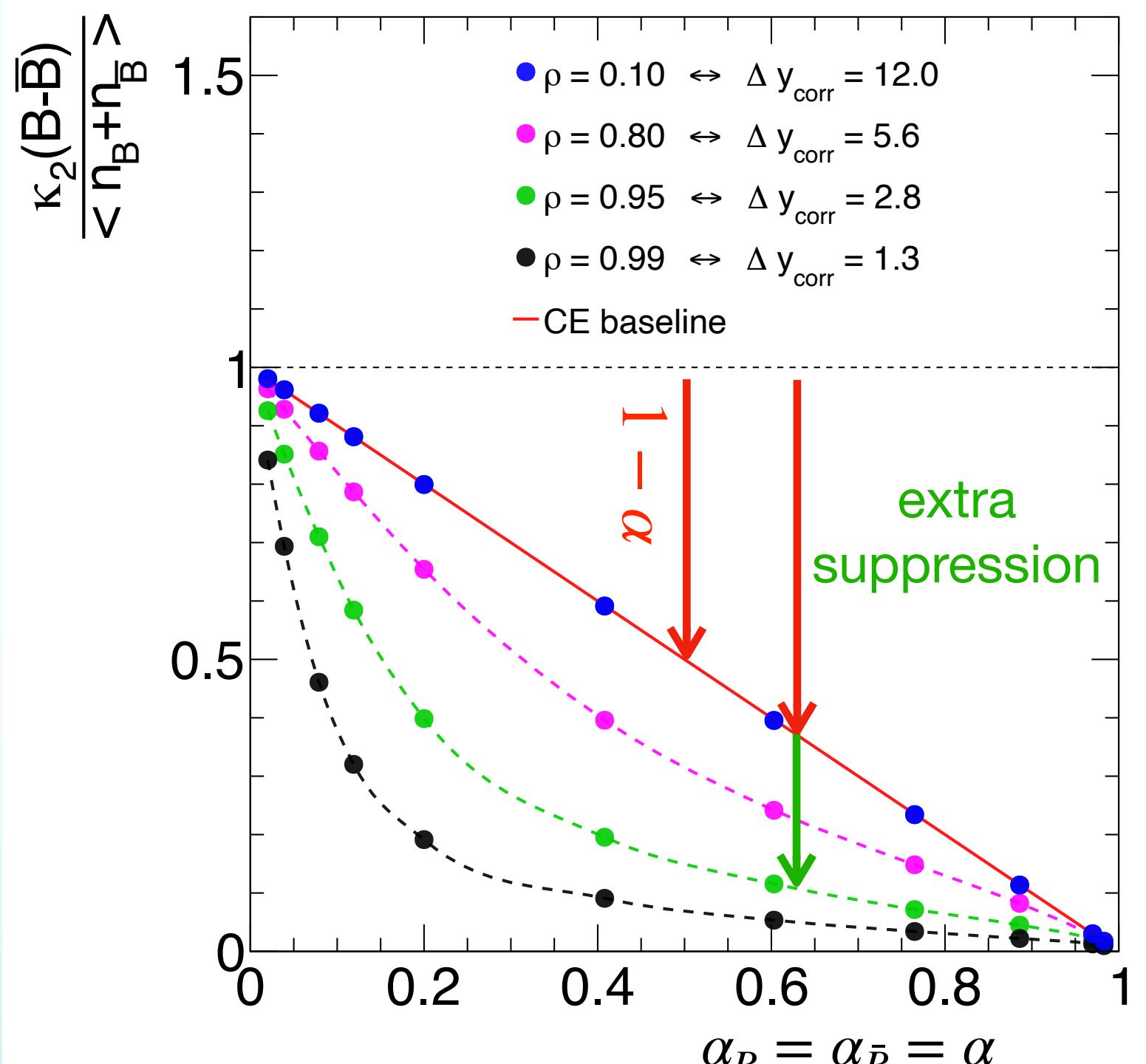
$$\langle N_B \rangle = \lambda_B \frac{\partial \ln Z_B}{\partial \lambda_B} \Bigg|_{\lambda_B, \lambda_{\bar{B}} = 1} = \frac{I_{B-1}(2z)}{I_B(2z)}$$

# ALICE Results (Identity Method)

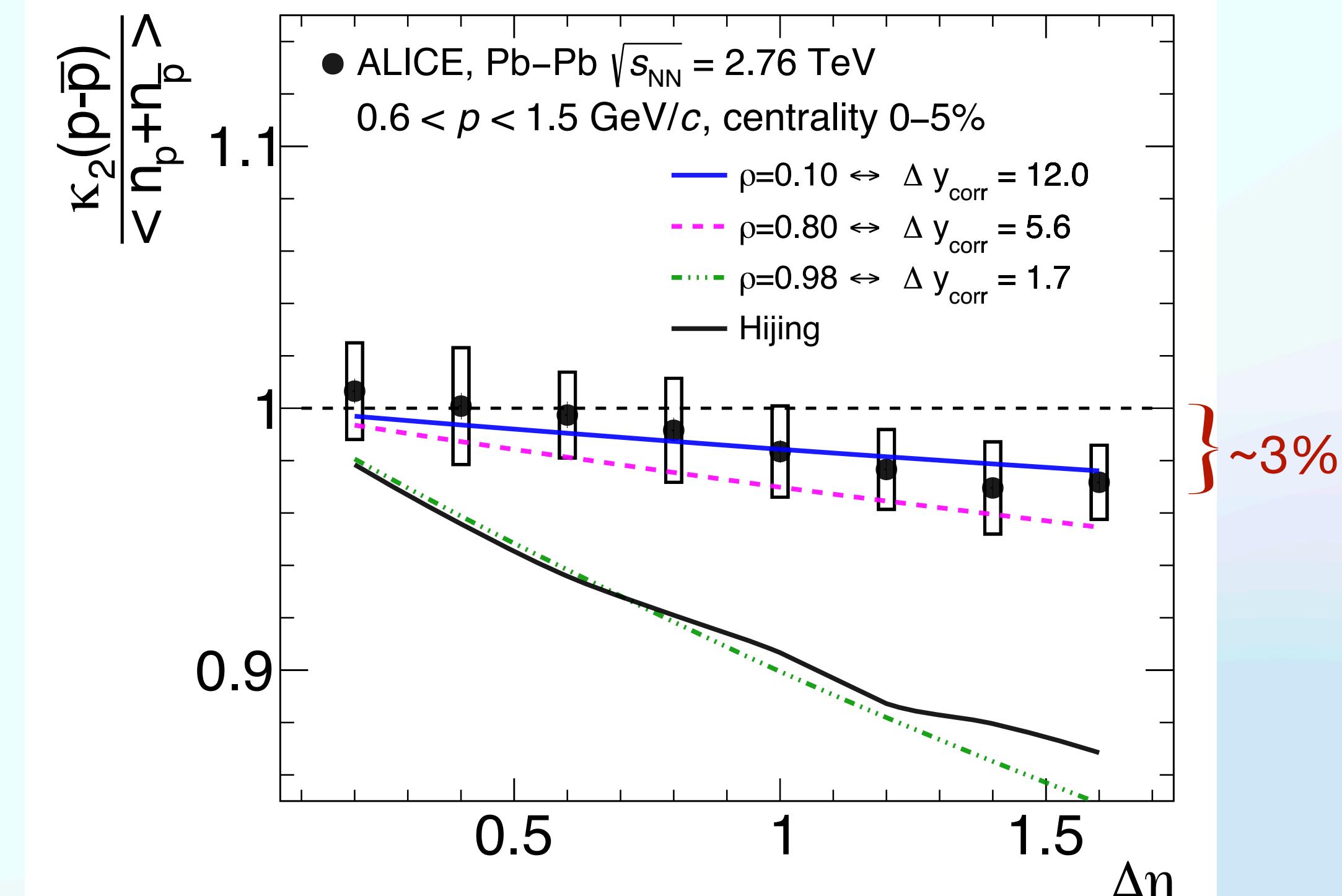
P. Braun-Munzinger, B. Friman, K. Redlich, A.R., J. Stachel, NPA 1008 (2021) 122141

P. Braun-Munzinger, K. Redlich, A.R., J. Stachel, JHEP 08 (2024) 113

**ALICE:** Phys. Lett. B 807 (2020) 135564, Phys. Lett. B (2022) 137545

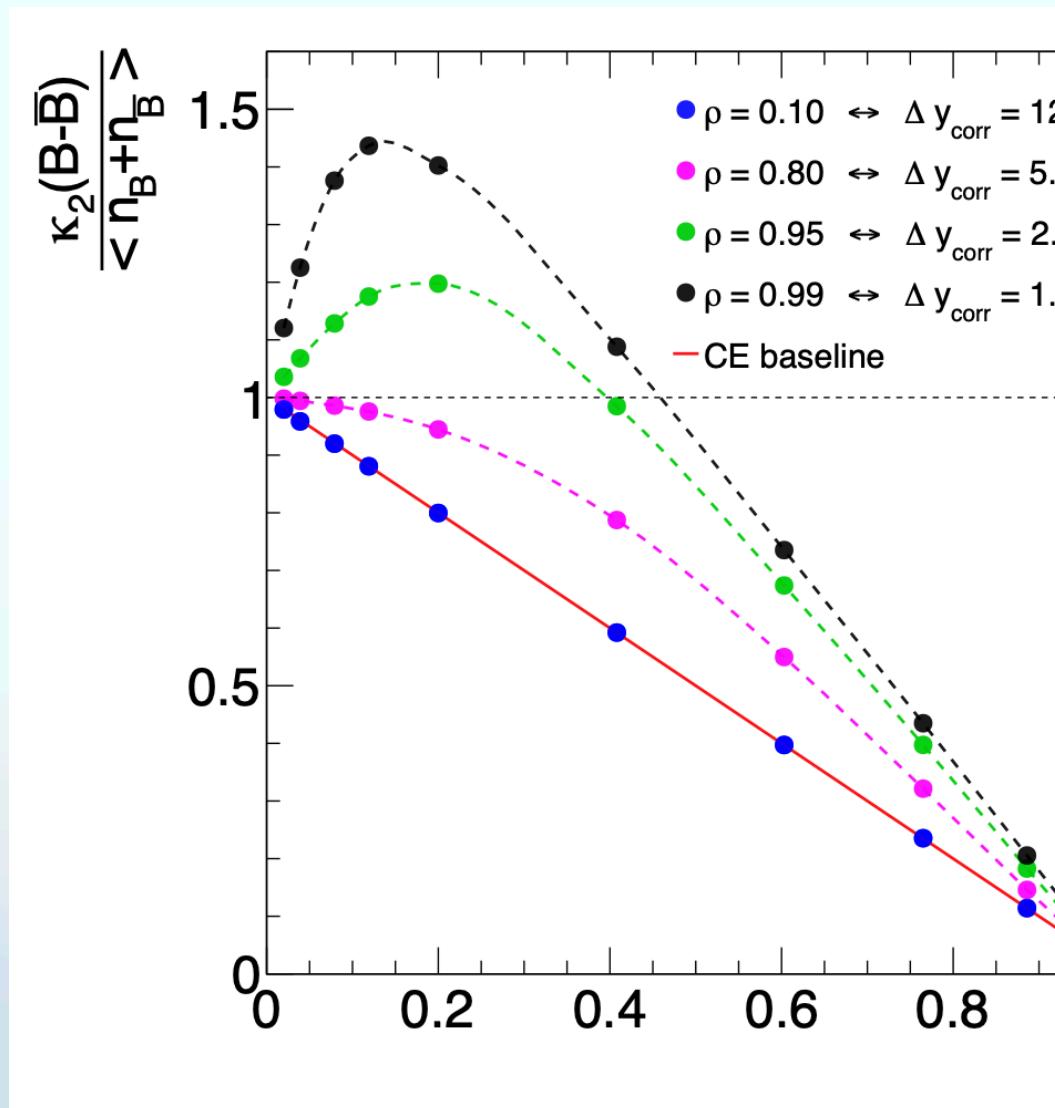
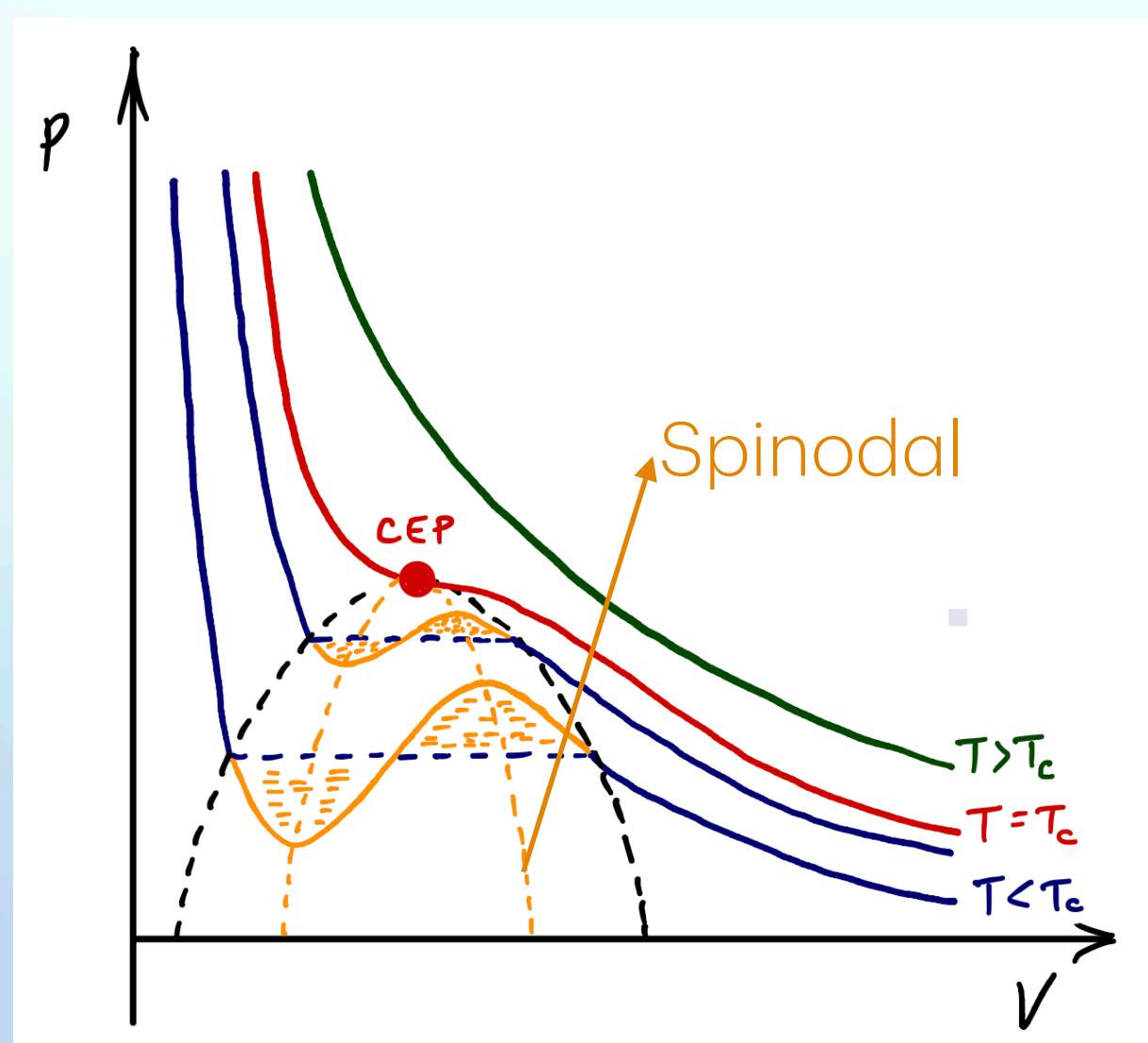
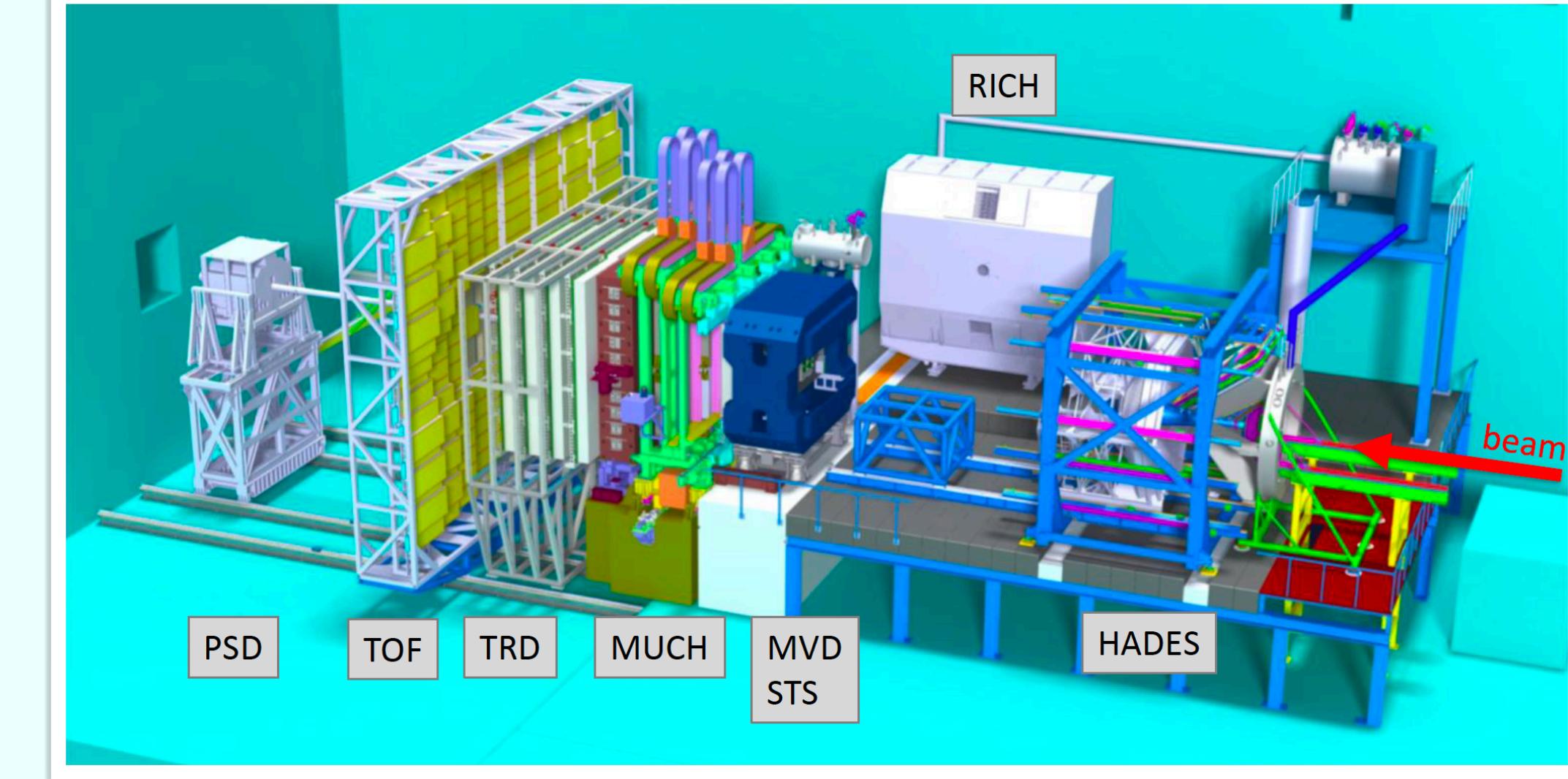
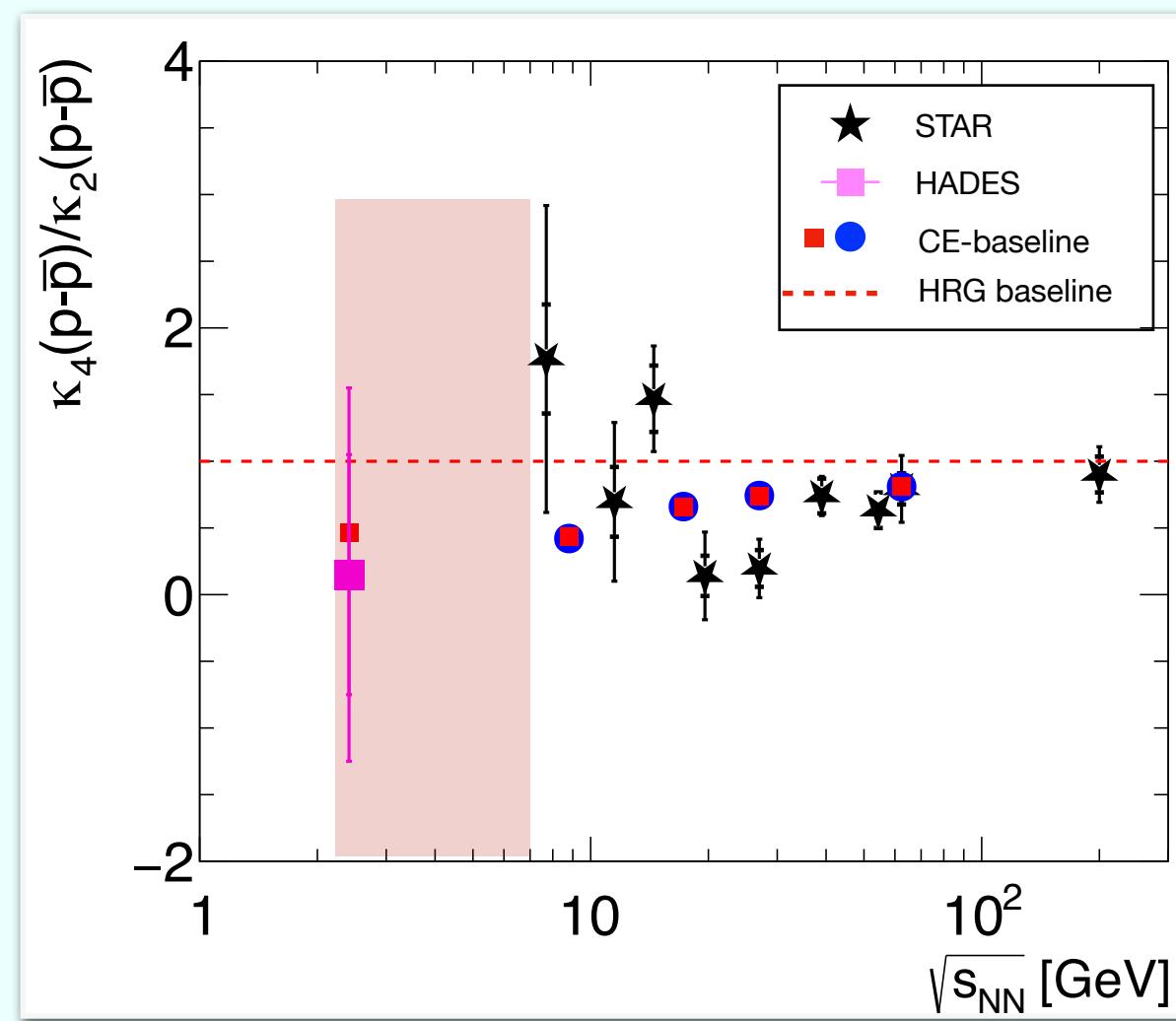


essential to constrain  
baryon production  
mechanism



- Alice data: best description with  $\rho = 0.1$  ( $\Delta y_{corr} = 12$ )  $\leftrightarrow$  **Global baryon number conservation**
- Agreement with LQCD predictions
- Calls into question baryon production mechanism in Hijing (Lund String Fragmentation)

# Near future, CBM

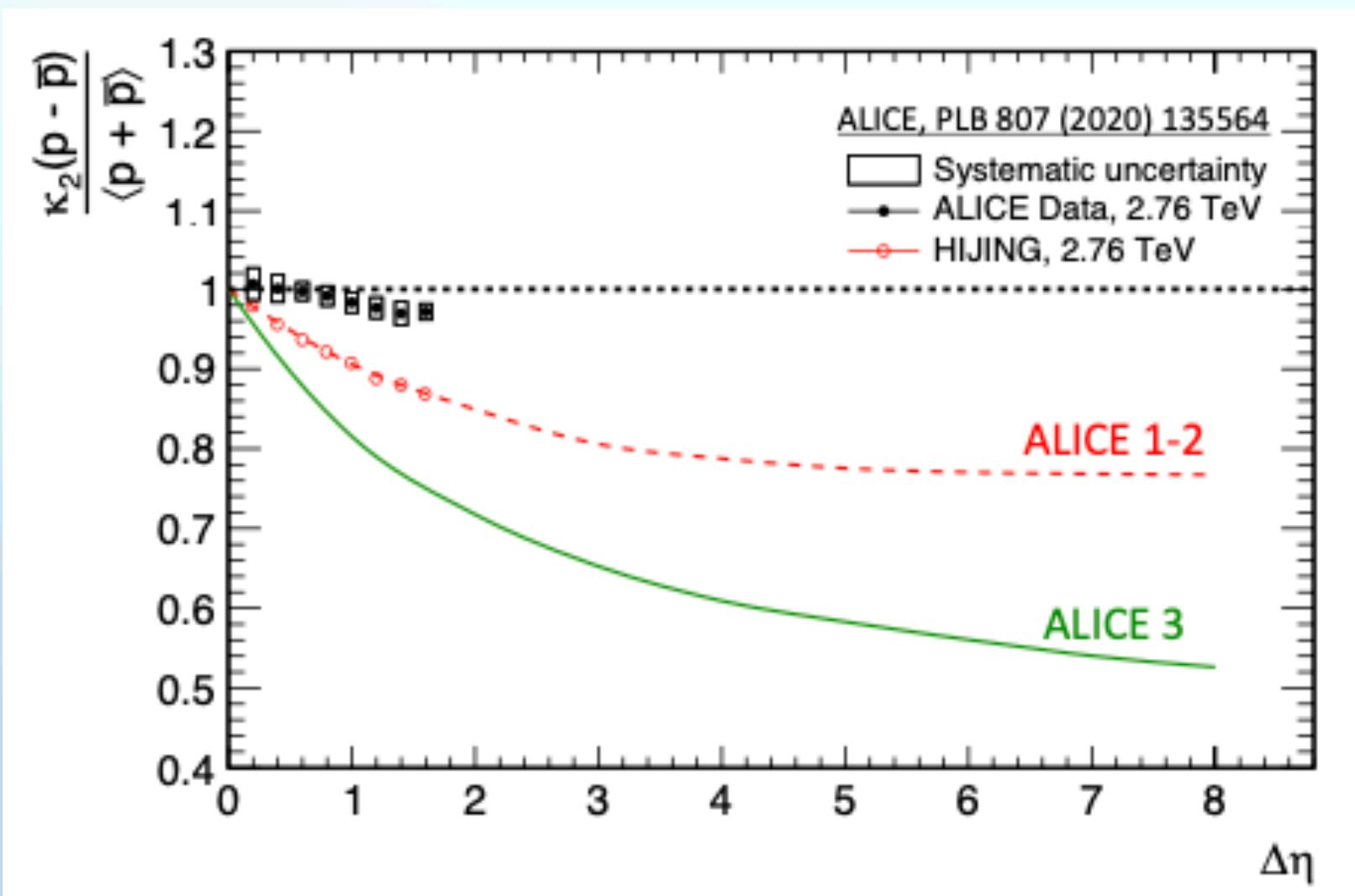
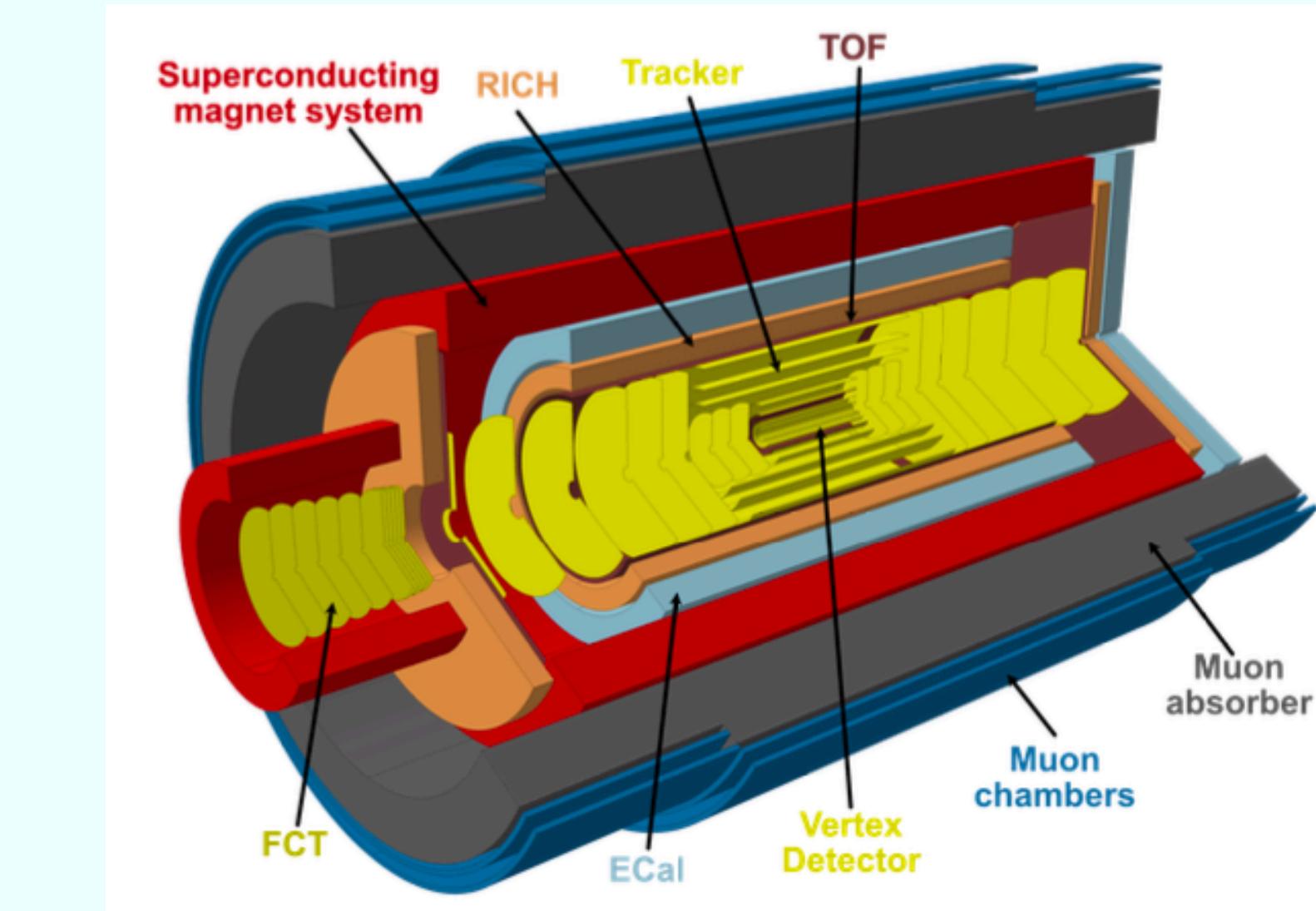
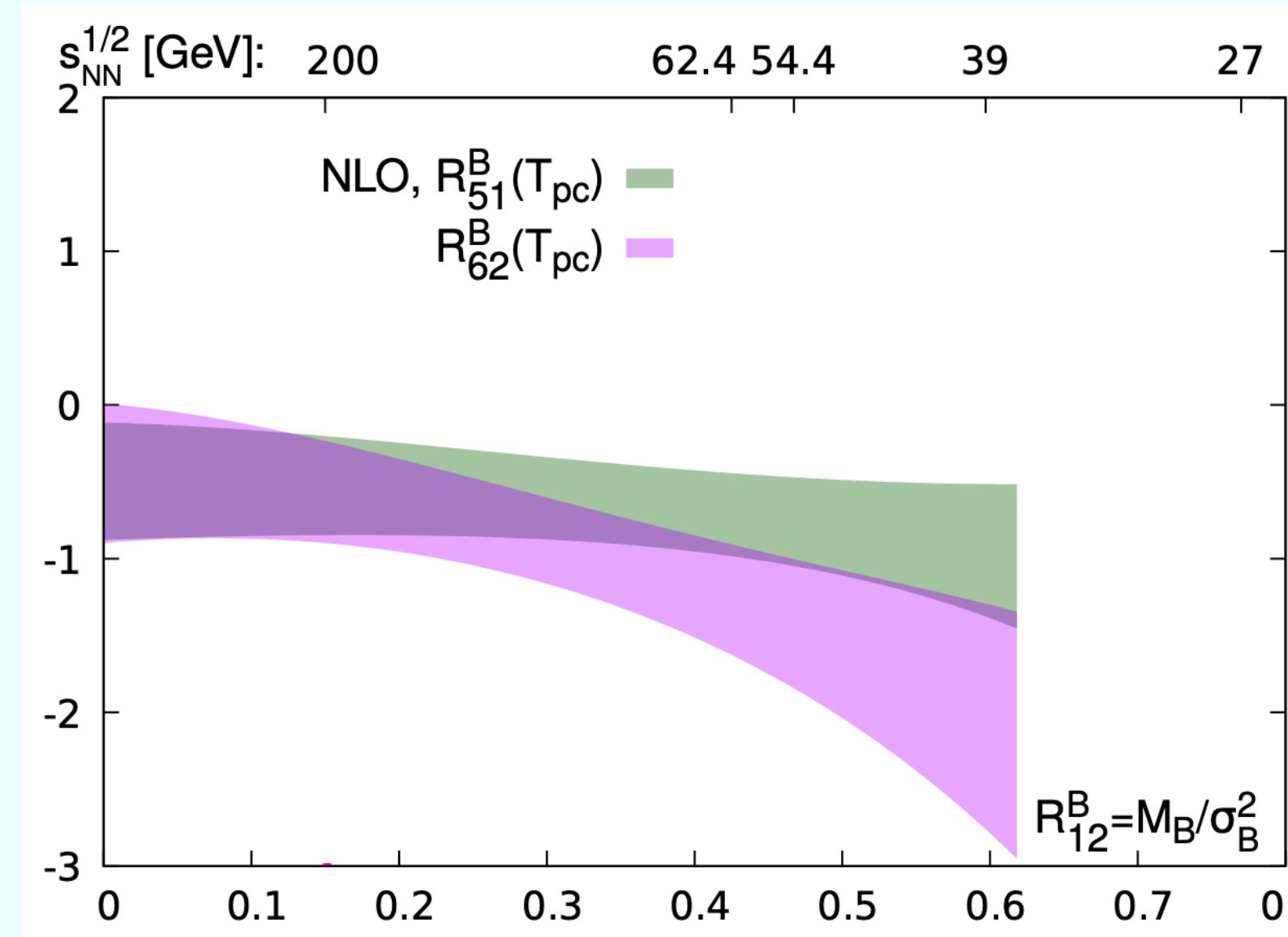


- Systematic measurements of fluctuations stemming from critical point
- Measuring fluctuations induced by spinodal decomposition
- Search for cluster formation

P. Braun-Munzinger, K. Redlich, A. R., J. Stachel, JHEP 08 (2024) 113  
 C. Sasaki, B. Friman, K. Redlich, Phys.Rev.D 77 (2008) 034024

# Near future, ALICE3

e-Print: 2211.02491 [physics.ins-det]



## Acceptance coverages

- **ALICE 1-2:**  $0.6 < p < 1.5 \text{ GeV}/c$ ,  $|\eta| < 0.8$
- **ALICE 3:**  $0.3 < p < 10 \text{ GeV}/c$ ,  $|\eta| < 4$

**Opens new avenues, such as study of charm fluctuations**



## Happy Birthday, Dear Johanna!

May your journey, both in physics and in life, be filled with breakthrough moments, smooth trajectories, and just the right amount of fluctuations to keep things interesting.

