

# THEORY OF MULTIQUARK STATES

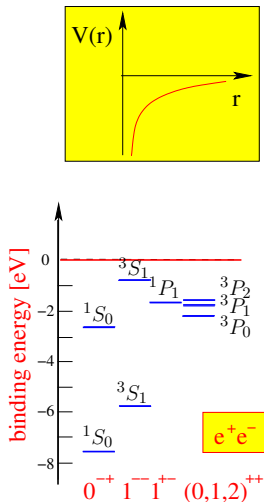
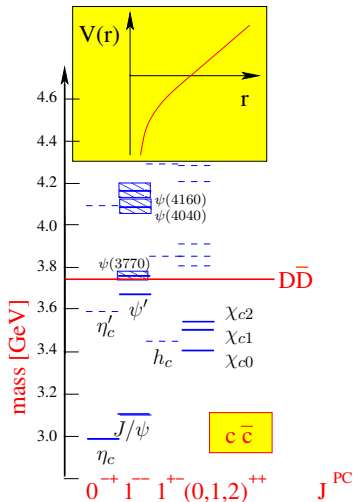
## A Challenge to Our Understanding of the Particle Spectrum

March 14, 2024 | Christoph Hanhart | IKP/IAS Forschungszentrum Jülich



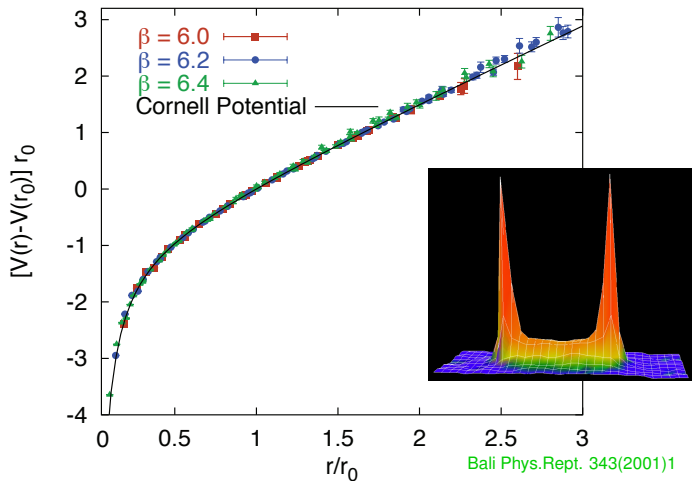
# CHARMONIUM BEFORE 2003

Quark-Model: Eichten et al. PRD 17 (1978)



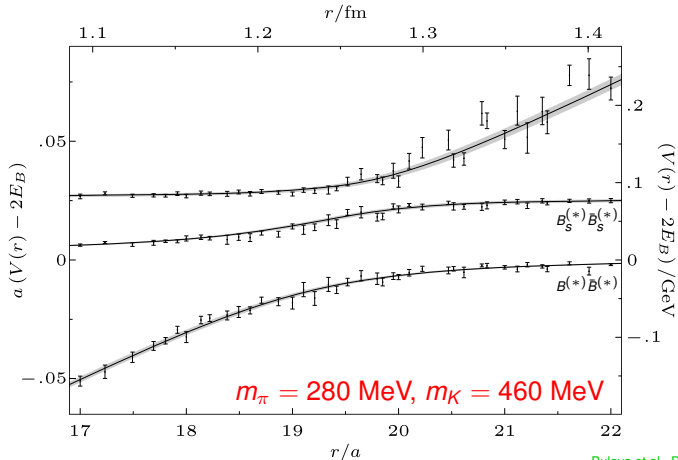
# $\bar{Q}Q$ POTENTIAL FROM LATTICE

Potential of two static color sources



# STRINGBREAKING

Adding light quarks:  $Q\bar{Q} \rightarrow (Q\bar{q})(\bar{Q}q)$  becomes possible

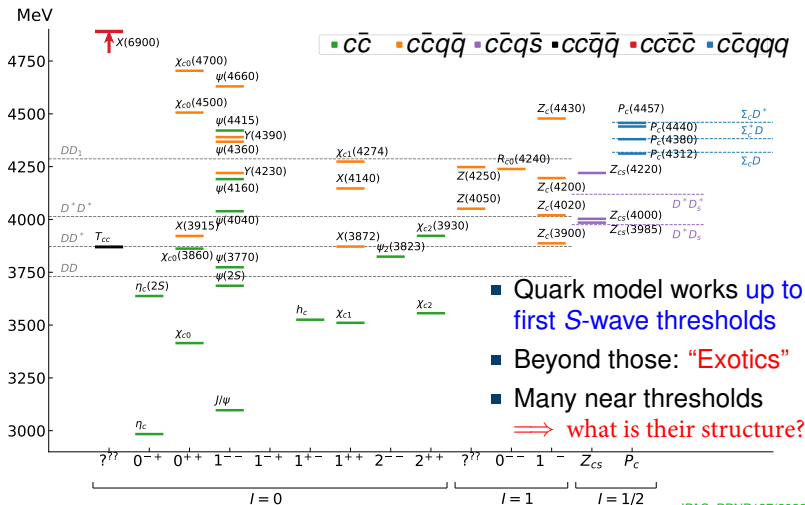


Bulava et al., PLB793(2019)493

⇒ string breaking for energies above first open flavor threshold



# ... PROVIDES NEW PHENOMENA



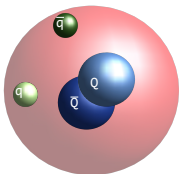
- Quark model works up to first  $S$ -wave thresholds
- Beyond those: “Exotics”
- Many near thresholds  $\implies$  what is their structure?

JPAC, PPNP127(2022)103981

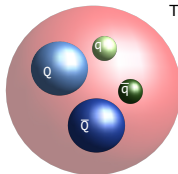


# THE NAME OF THE GAME:

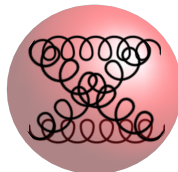
HADRO-  
QUARKONIUM



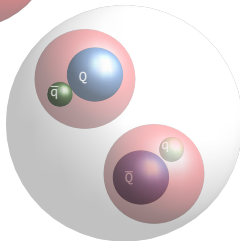
TETRAQUARK



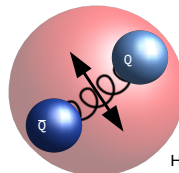
GLUEBALL



HADRONIC  
MOLECULE



HYBRID



Picture by Soeren Lange

How can one disentangle the different multi-quark structures?



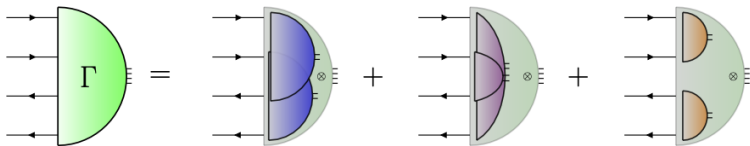
# DISCLAIMER

In what follows I focus on individual scenarios, in the spirit of

What if the exotics were pure .... (put your favourite approach)

So far little work on the possible mixing of the structures.

Promising approach: Dyson-Schwinger equations for multiquark states



J. Hoffer, G. Eichmann and C. S. Fischer, [arXiv:2402.12830 [hep-ph]].

So far exploratory studies for spectra, no uncertainty estimates yet ....



# HEAVY TETRAQUARKS

- Straightforward extension of the quark model

M. Gell-Mann, PL8(1964)214

- Mesons as diquark–anti-diquark systems

Jaffe, PRD15(1977)267, Maiani et al., PRD71(2005)014028

- Building blocks: Spin 0 and 1 heavy-light diquarks

- Separated by potential well

Selem and Wilczek, hep-ph/0602128; Maiani et al., PLB778(2018)247

- Spin-spin interaction dominant within diquarks

Maiani et al. PRD89(2014)114010

- and tensor force,  $S_{12}$ , needs to be considered

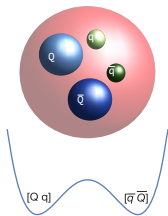
Ali et al. EPJC78(2018)29

$$M = 2M_Q + \frac{B_Q}{2} L^2 + 2a_Y L \cdot S + \frac{b_Y}{4} S_{12} + 2\kappa_{cq} (S_q \cdot S_c + c.c.)$$

alternative approaches, e.g., Cui et al., HEPNP31(2007)7; Stancu, JPG37(2010) 075017; Bhavsar et al., NPA1000(2020)121856

→ Already many ground states

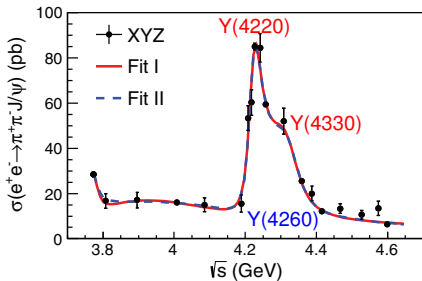
→ Each level has isovector and isoscalar state (cf.  $\rho$  and  $\omega$ )



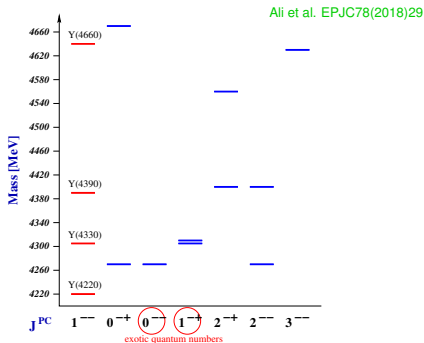


# RESULTS FOR NEGATIVE PARITY STATES

- four  $1^{--}$  ground states:  $P$ -wave &  $[0, 0]_0, [1, 0]_1 + [0, 1]_1, [1, 1]_0, [1, 1]_2$
- BESIII claims 2 in  $J/\psi\pi\pi$



BESIII, PRL118(2017)092001



- Without tensor force very light  $3^{--}$
- Many more states predicted than observed!

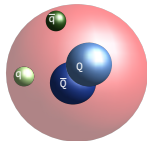
Cleven et al., PRD 92(2015)014005

Maybe since **di-quark picture too restrictive/constraining?**

Richard et al., PRD95(2017)054019



# HADROCHARMONIUM



- Extra states are viewed as **compact  $\bar{Q}Q$**  surrounded by light quarks

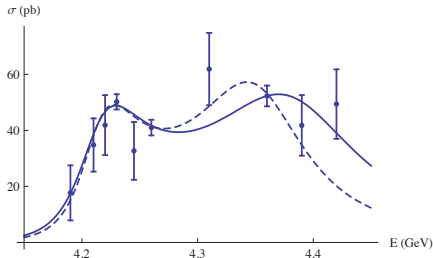
M. B. Voloshin, PPNP61(2008)455

- Provides natural explanation why, e.g.,  $Y(4260)$  is seen in  $J/\psi\pi\pi$  final state but not in  $\bar{D}D$

- Heavy quark spin symmetry demands that **spin of the core is conserved** in decay to charmonia

- Explaining  $e^+e^- \rightarrow h_c\pi\pi$  needs **mixing** between states with  $s_{\bar{c}c} = 0$  and  $s_{\bar{c}c} = 1$  leading to  $Y(4260)$  and  $Y(4360)$

Li & Voloshin MPLA29(2014)1450060



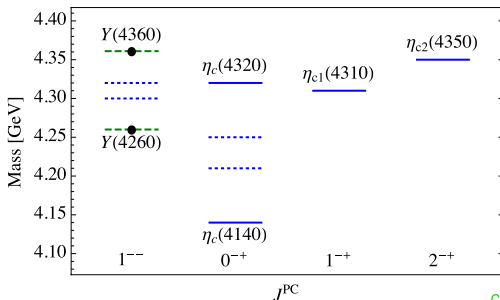
# HADROCHARMONIUM: NEW STATES

The above mentioned mixing suggests for the unmixed states:

$$\Psi_3 \sim (1^{--})_{c\bar{c}} \otimes (0^{++})_{q\bar{q}} \quad \Psi_1 \sim (1^{+-})_{c\bar{c}} \otimes (0^{-+})_{q\bar{q}},$$

where the heavy cores are  $\psi'$  and  $h_c$ .

→ get spin partners via  $\psi' \rightarrow \eta'_c$  and  $h_c \rightarrow \{\chi_{c0}, \chi_{c1}, \chi_{c2}\}$



Cleven et al., PRD 92(2015)014005

Special feature: very light  $0^{-+}$  state that should not decay to  $D^* \bar{D}$



# HADRONIC MOLECULES

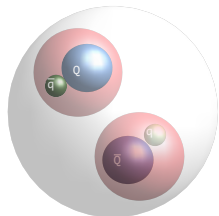
review article: Guo et al., Rev. Mod. Phys. 90(2018)015004

- are few-hadron states, **bound by the strong force**
- **do exist**: light nuclei.  
e.g. **deuteron as  $pn$**  & **hypertriton as  $\Lambda d$**  bound state
- are located typically **close to relevant continuum threshold**;  
e.g., for  $E_B = m_1 + m_2 - M$  ( $\gamma = \sqrt{2\mu E_B}$ ;  $\mu = m_1 m_2 / (m_1 + m_2)$ )
  - $E_B^{\text{deuteron}} = 2.22 \text{ MeV}$  ( $\gamma = 40 \text{ MeV}$ )
  - $E_B^{\text{hypertriton}} = (0.13 \pm 0.05) \text{ MeV}$  (to  $\Lambda d$ ) ( $\gamma = 26 \text{ MeV}$ )
- **can be identified in observables (Weinberg compositeness)**:

$$\frac{g_{\text{eff}}^2}{4\pi} = \frac{4M^2\gamma}{\mu}(1 - \lambda^2) \rightarrow a = -2 \left( \frac{1 - \lambda^2}{2 - \lambda^2} \right) \frac{1}{\gamma}; \quad r = - \left( \frac{\lambda^2}{1 - \lambda^2} \right) \frac{1}{\gamma}$$

where  $(1 - \lambda^2)$ =probability to find molecular component in bound state wave function

Are there mesonic molecules?



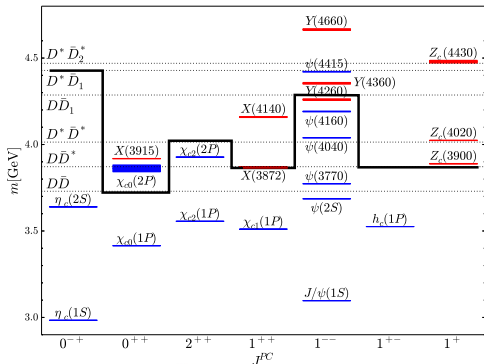
# GENERAL CONSIDERATIONS

Constituents must be narrow. Heavy candidates ( $M, \Gamma$  in MeV)

$D(0^-, M = 1865, \Gamma \simeq 0)$ ;  $D^*(1^-, M = 2007, \Gamma \simeq 0.1)$

$D_1(1^+, M = 2420, \Gamma \simeq 30)$ ;  $D_2^*(2^+, M = 2460, \Gamma \simeq 50)$

$D_0(2400)$  and  $D_1(2430)$  with  $\Gamma = 300$  MeV too broad ...



■ Explains mass gap

$$M_{Y(4260)} - M_{X(3872)} = 388 \text{ MeV} \simeq$$

$$M_{D_1(2420)} - M_{D^*} = 410 \text{ MeV}$$

■ Predicts, e.g.,

$$M(0^-) - M(1^-) \simeq M_{D^*} - M_D \simeq +100 \text{ MeV}$$

microscopic calc.: T. Ji et al. PRL129(2022)102002

For tetraquark:

$$M(0^-) - M(1^-) \simeq + 50 \text{ MeV}$$

Ali et al. EPJC78(2018)29

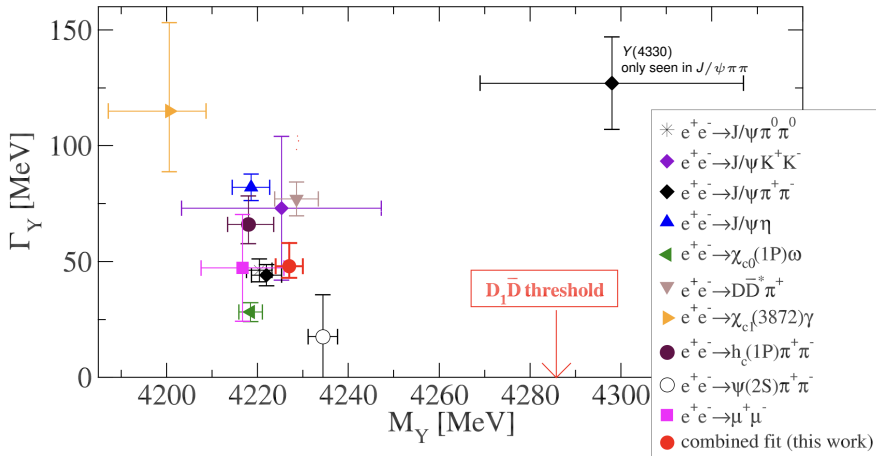
For hadrocharmonium:

$$M(0^-) - M(1^-) \simeq -100 \text{ MeV}$$

Cleven et al., PRD 92 (2015) 014005



# EXAMPLE: $Y(4230)$ AS $D_1\bar{D}$ MOLECULE

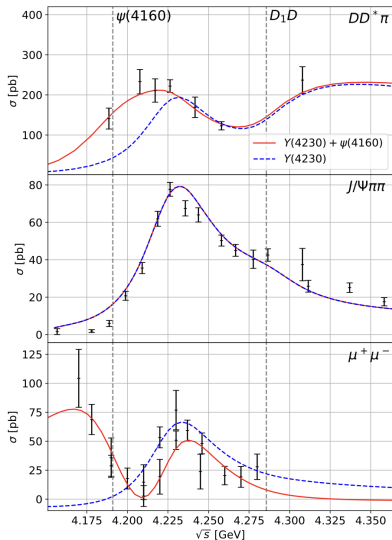


- Inclusion of  $D_1\bar{D}$  intermediate states ( $g_{YD_1D}$  large for molecule)
- Inclusion of charmonium  $\psi(4160)$  ( $M_{\psi(4160)} = 4191$  MeV)

L. von Detten, V. Baru, CH, Q. Wang, D. Winney, Q. Zhao; arxiv: 2402.03057



# IMPACT OF $\psi(4160)$



Well established  $\bar{c}c$  state

Parameters from RPP2023:

2023 update of R. L. Workman *et al.* [PDG], PTEP2022 (2022)083C01

$$m_{\psi(4160)} = (4191 \pm 5) \text{ MeV}$$

$$\Gamma_{\psi(4160)} = (70 \pm 10) \text{ MeV}$$

Experimental extractions:

$$D^0 D^{*-} \pi^+ : \Gamma_{\gamma} = (77 \pm 6.3 \pm 6.8) \text{ MeV}$$

BESIII, PRL130(2023) 121901

$$J/\psi \pi^+ \pi^- : \Gamma_{\gamma} = (41.8 \pm 2.9 \pm 2.7) \text{ MeV}$$

BESIII, PRD106(2022)072001

in both cases  $\psi(4160)$  omitted

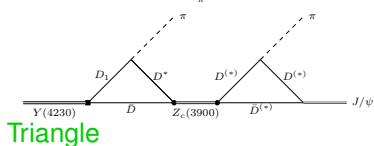
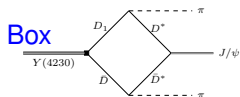
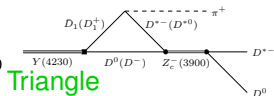
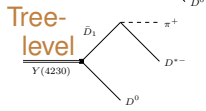
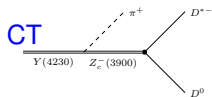
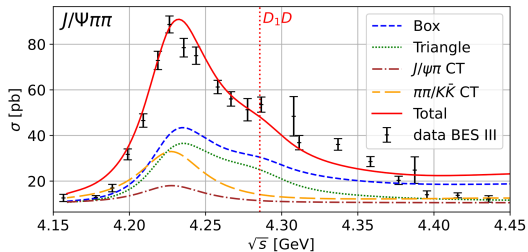
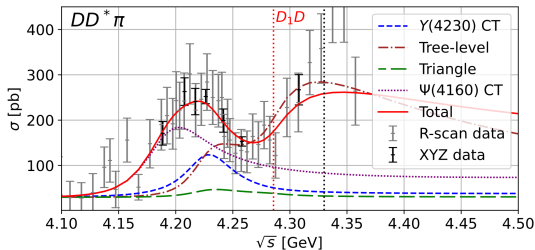
$$\mu^+ \mu^- : \Gamma_{\gamma} = (47.2 \pm 22.8 \pm 10.5) \text{ MeV}$$

BESIII, PRD102(2020)112009

with  $\psi(4160)$  included



# ROLE OF $D_1\bar{D}$ CUT





# SUMMARY AND PERSPECTIVES

- These are exciting times in (heavy meson) spectroscopy
- The recent and future data have the potential to allow us to identify the prominent components in XYZ states

## to-do for experiment

- **Continue** great performance! Especially needed:
  - data for **different quantum numbers** and
  - data for **line shapes**

## to-do for theory

- Provide more predictions for the **different scenarios**
- Go beyond most simple approaches  
e.g. study **interplay of regular quarkonia with exotics**
  - potentially significant mixing

Kalashnikova et al., PRD80(2009)074004; Takizawa et al., PTEP(2013)093D01; Ortega et al., JPG 40(2013)065107;  
Coito et al., EPJC73(2013)2351; Cincioglu et al., EPJC76(2016)576

- negligible mixing

van Beveren et al., PLB641(2006)265; Hammer et al., EPJA52(2016)330, C.H. et al., PRD106(2022)114003

Thanks a lot for your attention



## Backup Slides



# QQ $\bar{q}\bar{q}$ TETRAQUARKS

Recently growing number of claims for those tetraquarks, e.g.

- from QCD sum rules
- from lattice QCD
- from phenomenology

Du et al., PRD87(2013)014003

Francis et al. PRL118(2017)142001

Ader et al., PRD 25(1982)2370

E.g. from the last work

Karliner and Rosner, PRL119(2017)202001; Eichten and Quigg, PRL119(2017)202002

$$m(QQ\bar{q}\bar{q}) - m(QQq) \simeq m(\bar{Q}\bar{q}\bar{q}) - m(\bar{Q}q)$$

exploiting heavy quark-diquark symmetry:

expansion in  $r_{QQ}/r_q \sim \Lambda_{\text{QCD}}/(M_Q v)$

Savage and Wise, PLB248(1990)177

One finds: The heavier QQ, the more deeply bound QQ $\bar{q}\bar{q}$

cc: From mass of  $\Sigma_{cc}$ : At most very shallow

State found ( $T_{cc}$ ) — consistent with molecule

LHCb, Nature Commun. 13(2022)3351

bb: No bb-baryon mass available; all groups agree that  
there should be a tetraquark structure



# OTHER CHANNELS

L. von Detten, V. Baru, CH, Q. Wang, D. Winney, Q. Zhao; arxiv: 2402.03057

