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# The muon $g - 2$ and the role of hadron physics

**Hartmut Wittig**

Institute for Nuclear Physics, Helmholtz Institute Mainz, and PRISMA<sup>+</sup> Cluster of Excellence,  
Johannes Gutenberg-Universität Mainz

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**DPG Spring Meeting, Justus Liebig University Gießen**

**14 March 2024**



JOHANNES GUTENBERG  
UNIVERSITÄT MAINZ







Braut sich da was zusammen? Im Sommer 2013 wurde der Myonen-Speicherring (auf dem Lastwagen) am Fermilab nahe Chicago angeliefert. Jetzt wurden erste Ergebnisse verkündet.

Foto Fermilab/Reidar Hahn

## Die Macht der Myonen

Ein Teilchen schickt sich an, eine beispiellos erfolgreiche Theorie zu sprengen. Viele Physiker freuen sich wie Bolle. Andere warnen, dazu sei es noch zu früh.

*Von Ulf von Rauchhaupt*

# Neue Zürcher Zeitung

Vor zwei Wochen bekam das Standardmodell der Teilchenphysik nasse Füße. Jetzt steht ihm das Wasser bis zum Hals

Christian Speicher

07.04.2021, 17.00 Uhr



## Frankfurter Allgemeine

HERAUSGEGEBEN VON GERALD BRAUNBERGER, JÜRGEN KAUBE, CARSTEN KNOP, BERTHOLD KOHLER

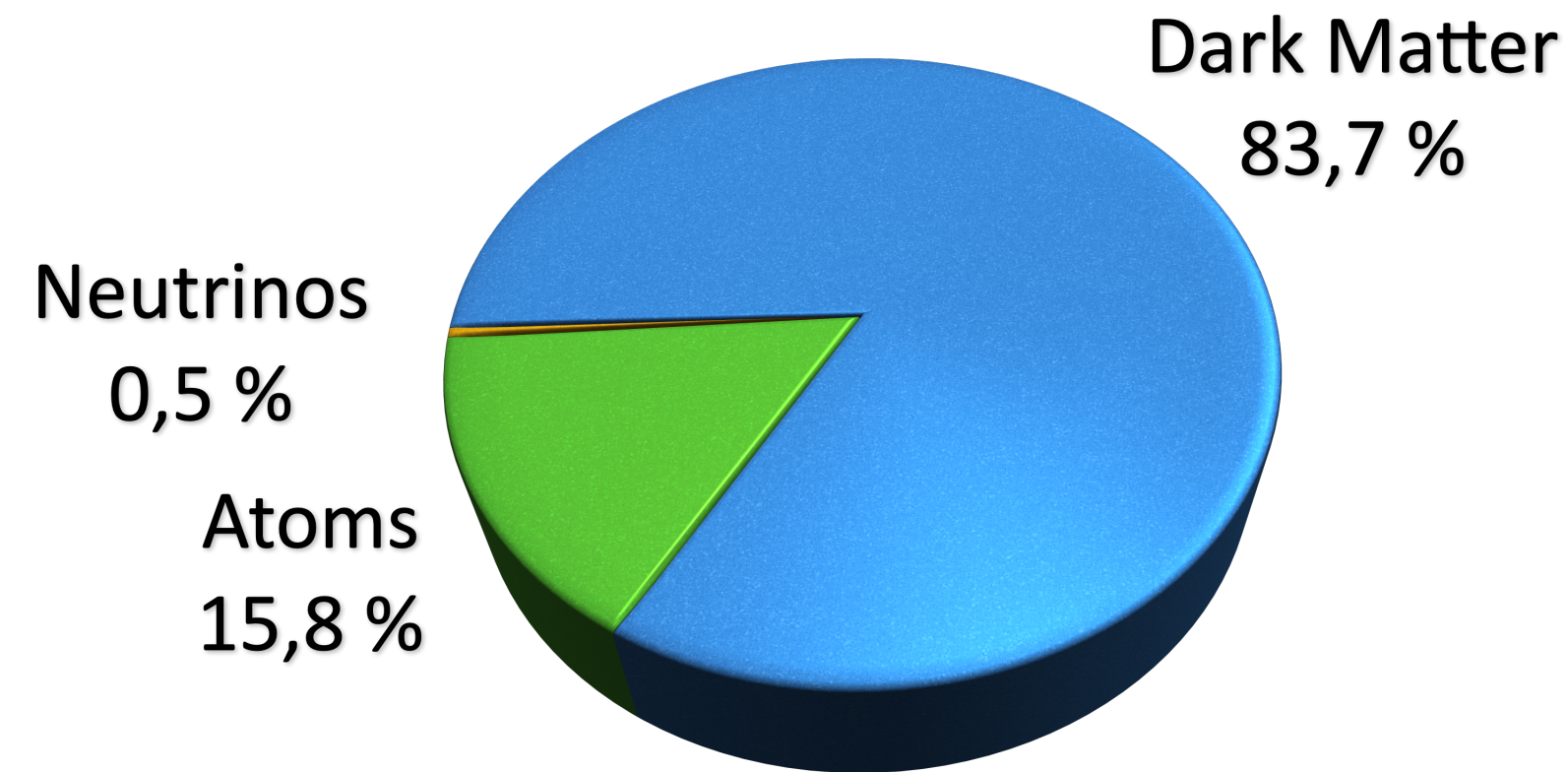
RÄTSELHAFTE MYONEN

## Abschied vom Standardmodell?

VON MANFRED LINDINGER - AKTUALISIERT AM 07.04.2021 - 18:55



# Motivation: The Dark Matter Puzzle

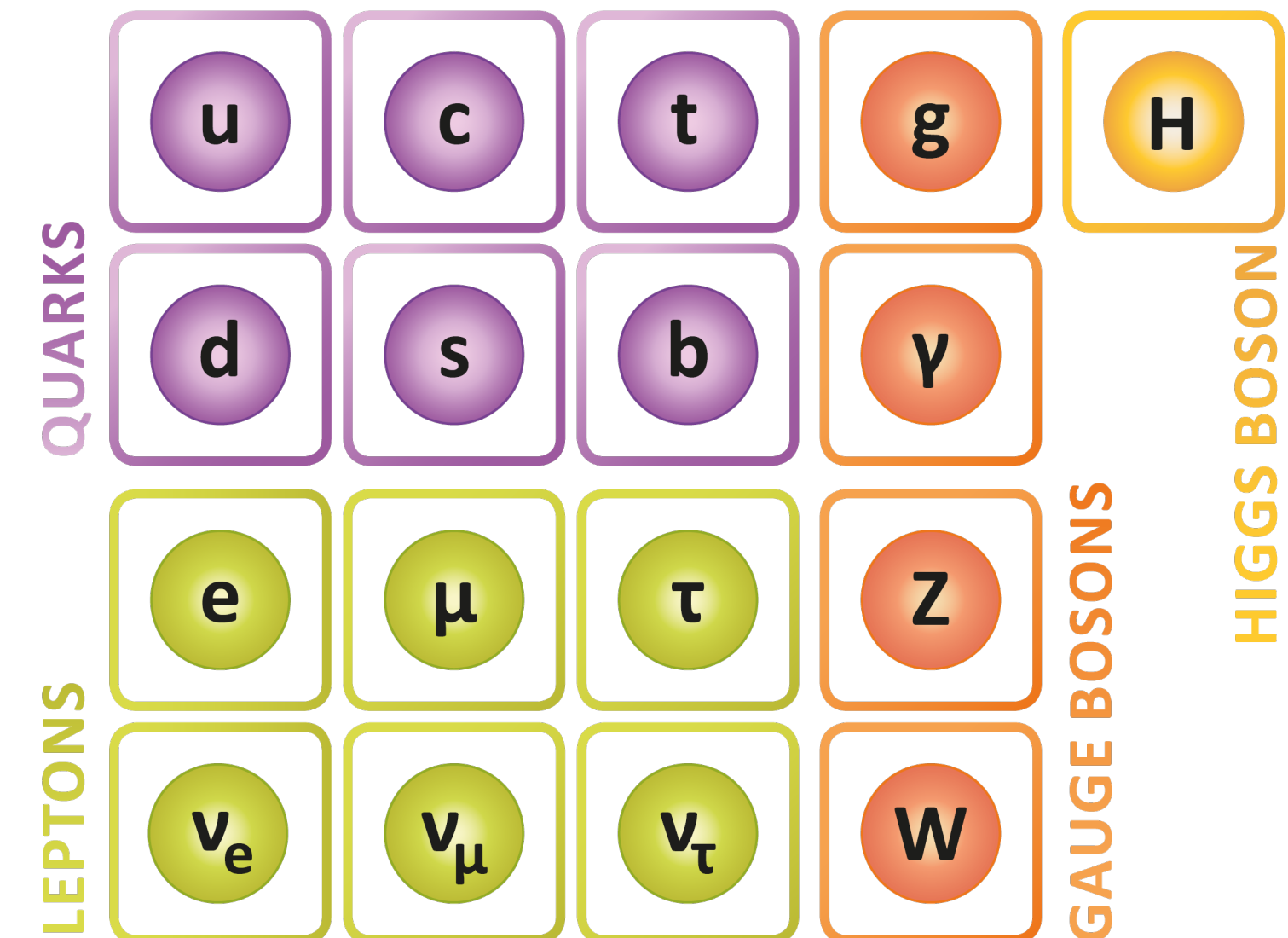


## Astrophysical observations:

- Dark Matter dominates matter density of universe
- No clues on nature of Dark Matter from laboratory experiments

## Standard Model of Particle Physics:

- Quantitative framework for description of known constituents of visible matter
- No Dark Matter candidate among SM particles
- No explanation for baryon asymmetry in universe

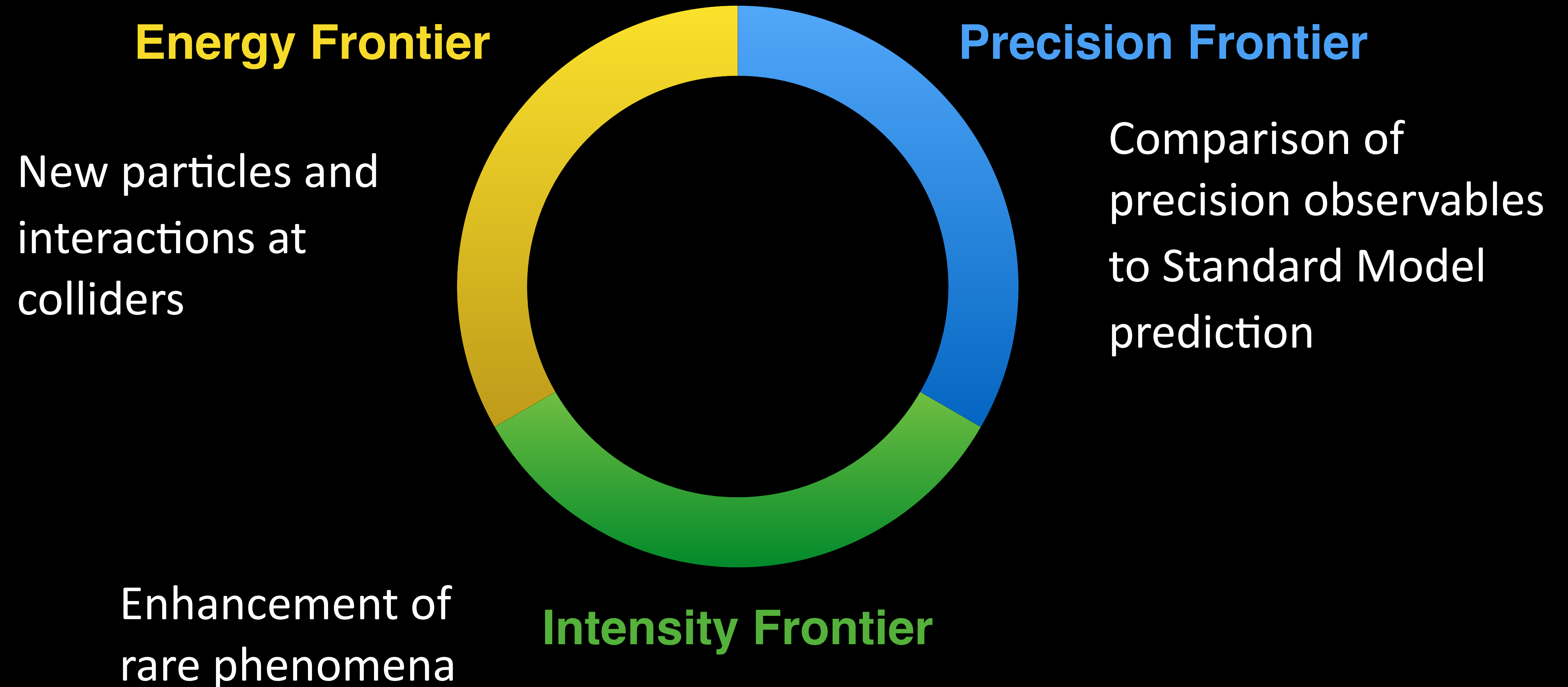


Source: Tanja Labs for PRISMA+

Standard Model does not provide a complete description of Nature



# The Quest for New Physics





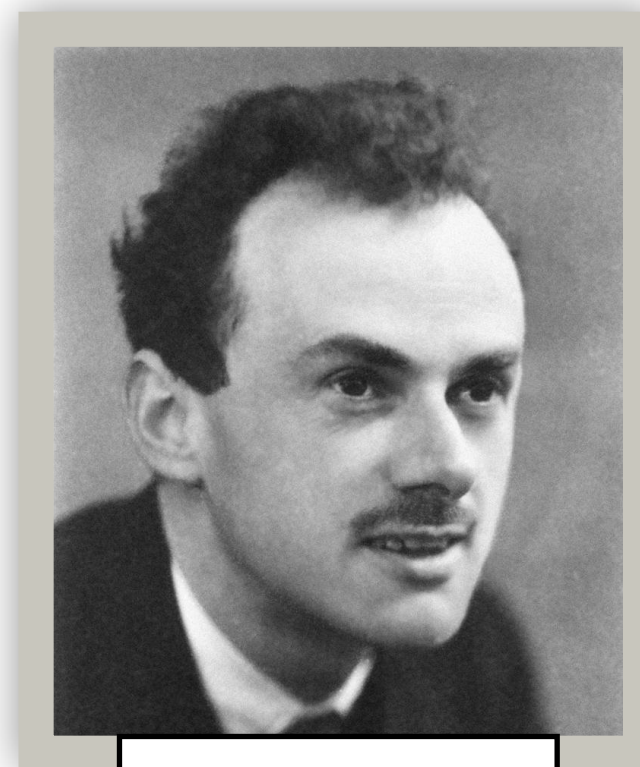
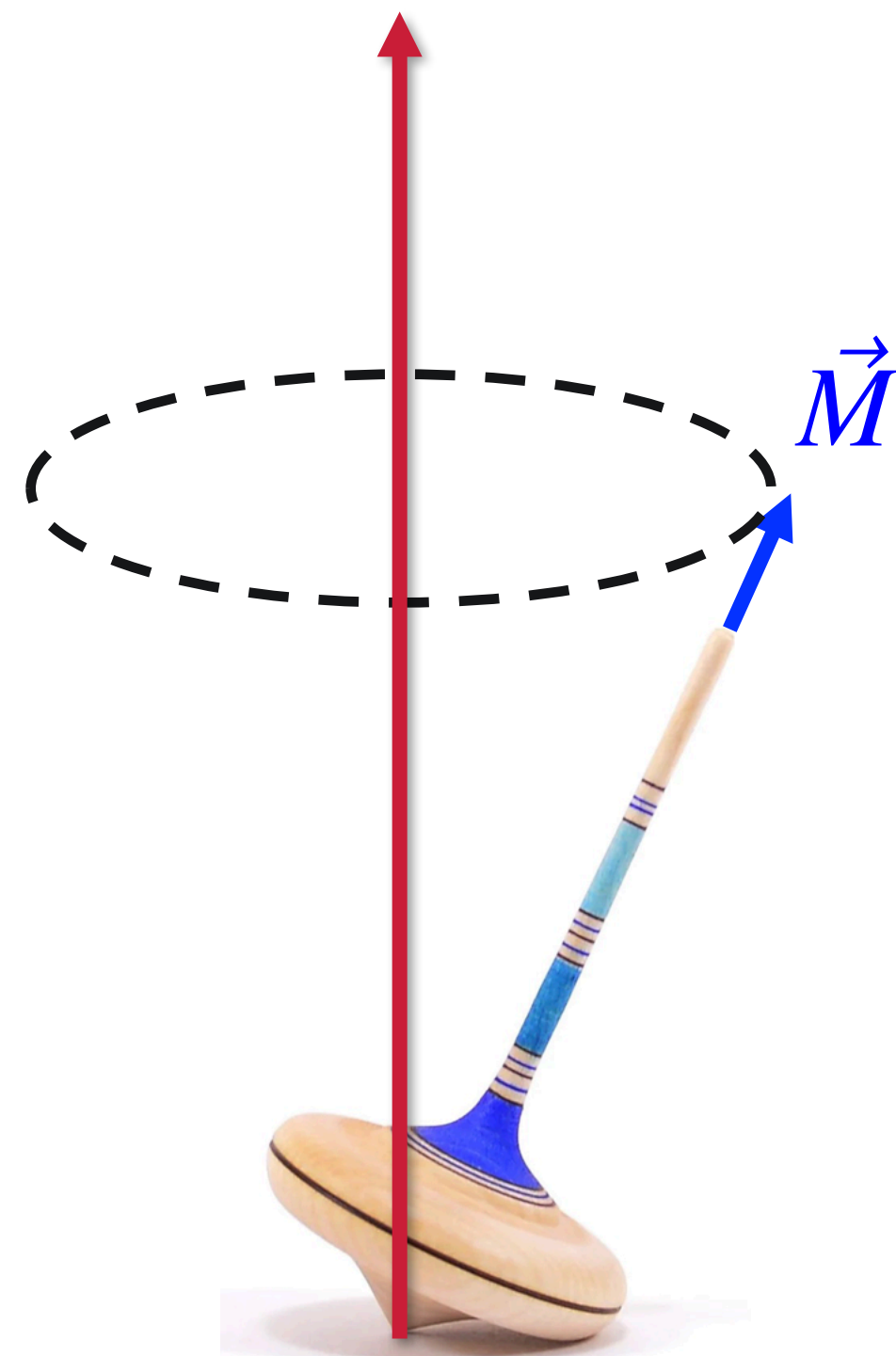
# The anomalous magnetic moment of the muon

Magnetic moment of particle with mass  $m$  and charge  $e$  :

$$\vec{M} = g \frac{e\hbar}{2m} \vec{S}$$

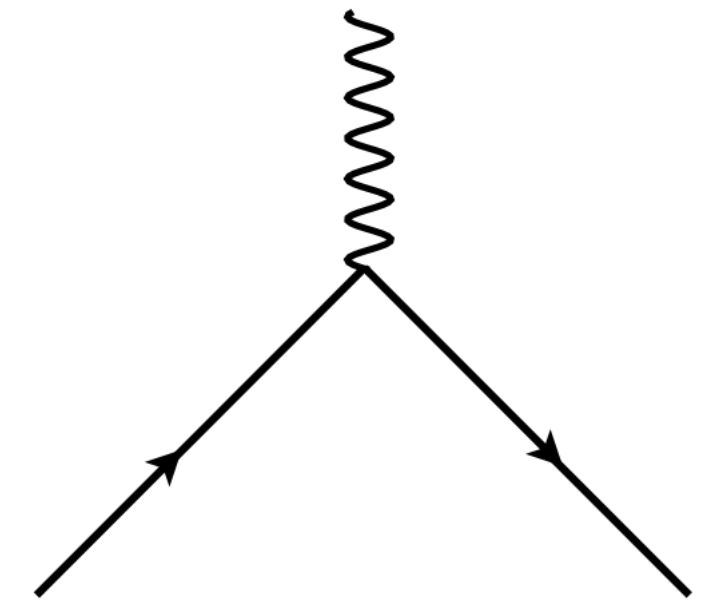
$g$ -factor

Magnetic field  $\vec{B}$



Paul Dirac

$$g = 2$$



Quantum corrections modify Dirac's prediction  $g = 2$

$$g = 2(1 + a), \quad a : \text{anomalous magnetic moment}$$



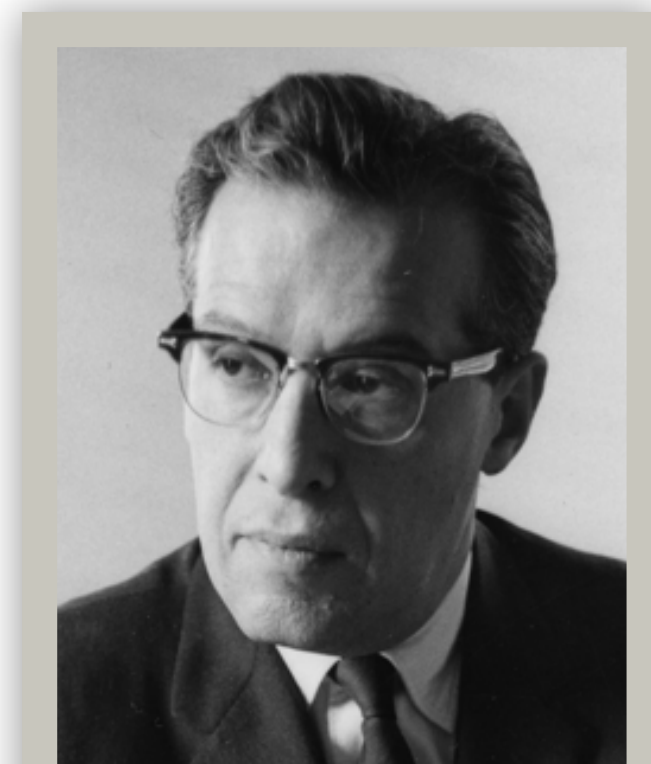
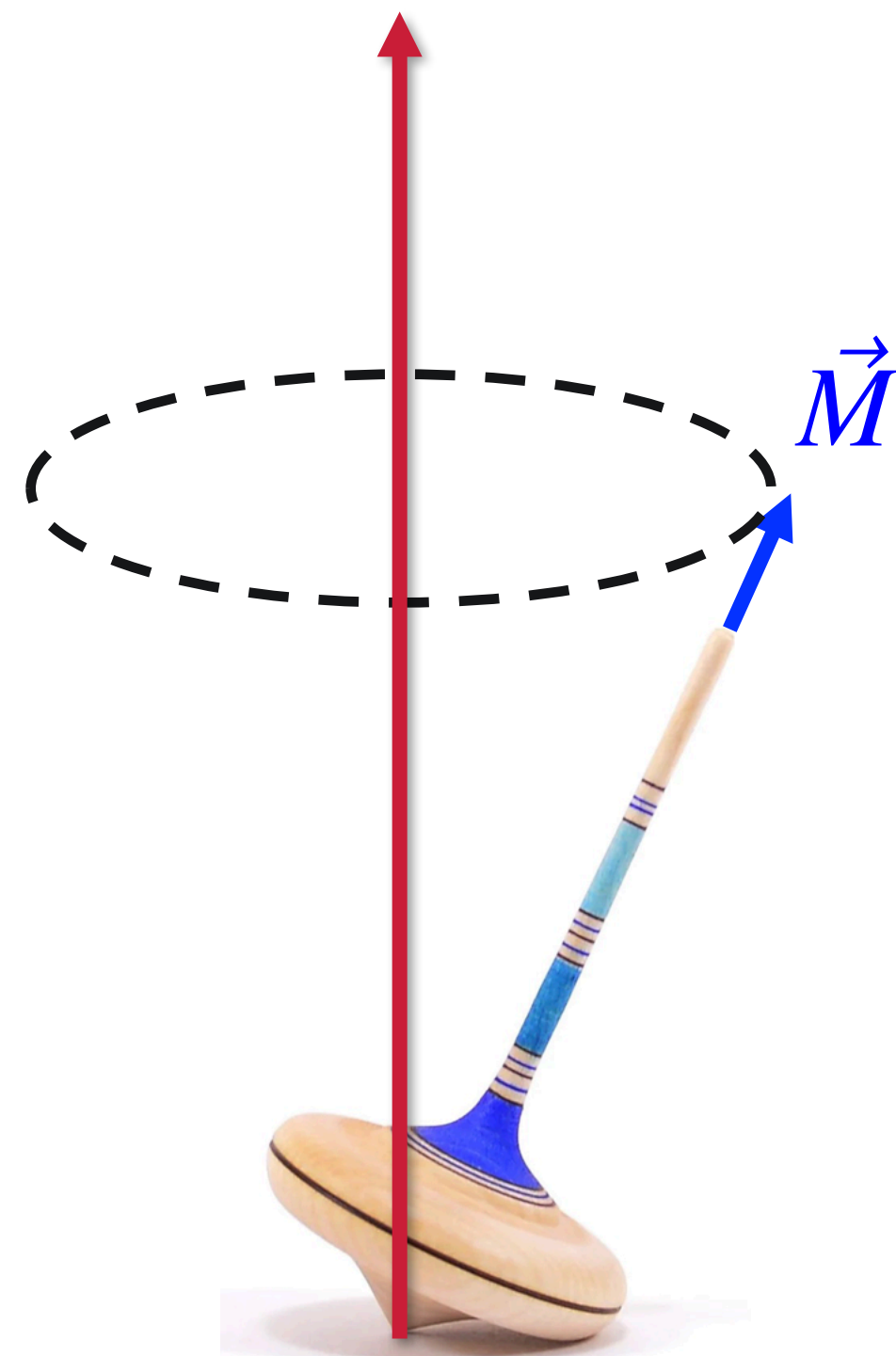
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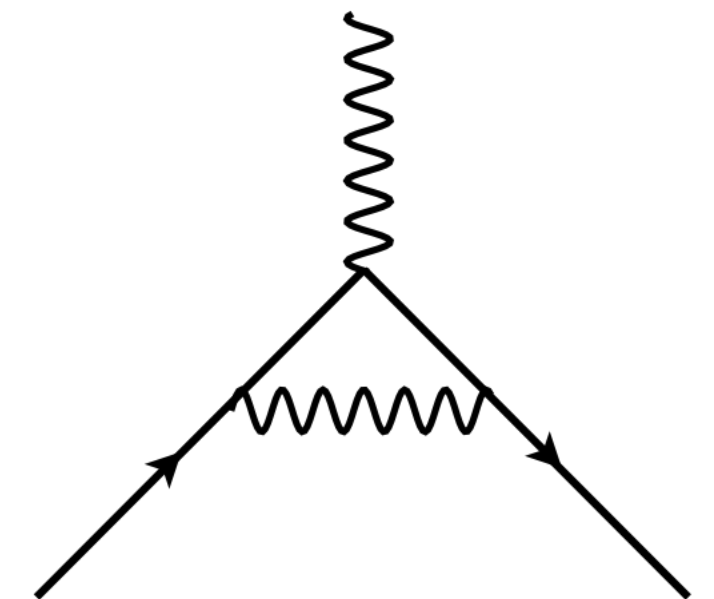
g-factor

Magnetic field  $\vec{B}$



Julian Schwinger

$$g = 2 \left( 1 + \frac{\alpha}{2\pi} \right)$$



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Electromagnetic, weak and strong interactions contribute to  $a$



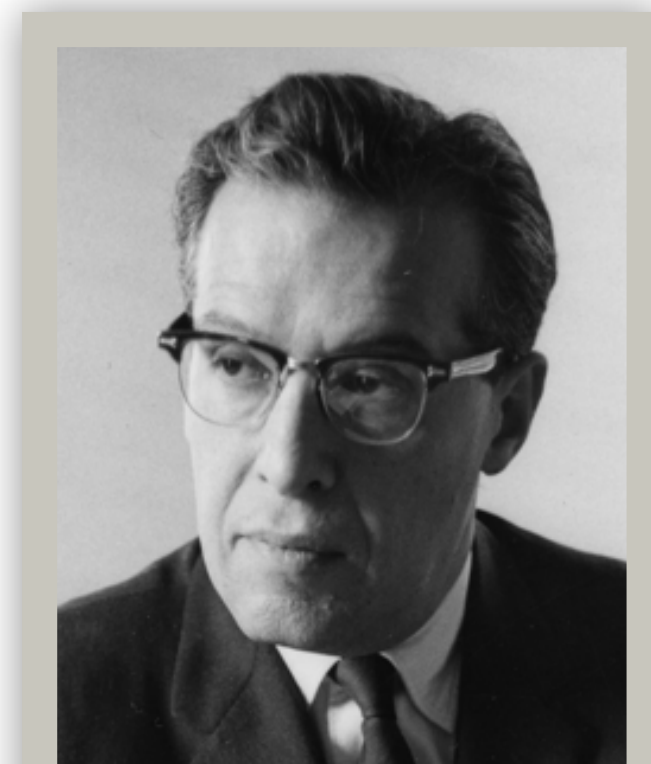
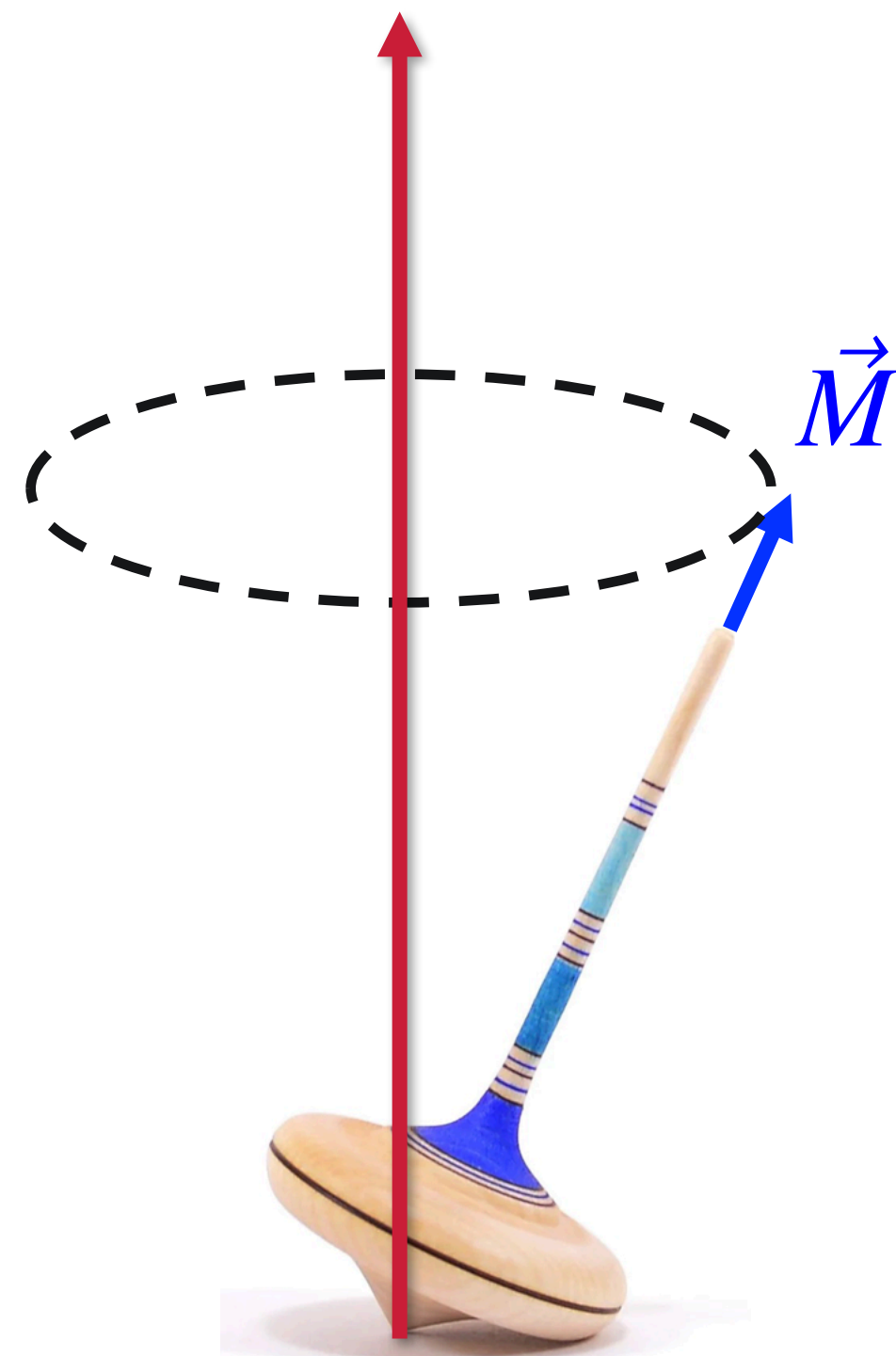
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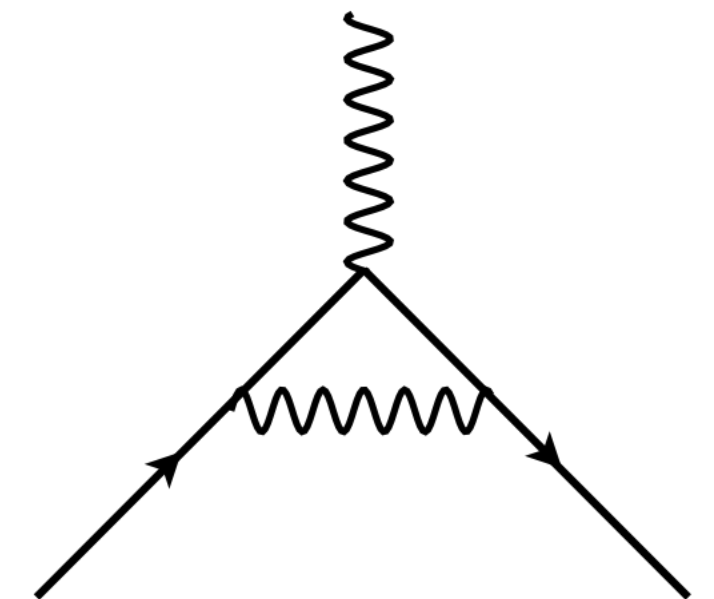
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Beyond leading order: distinct values of  $a_e$ ,  $a_\mu$  and  $a_\tau$



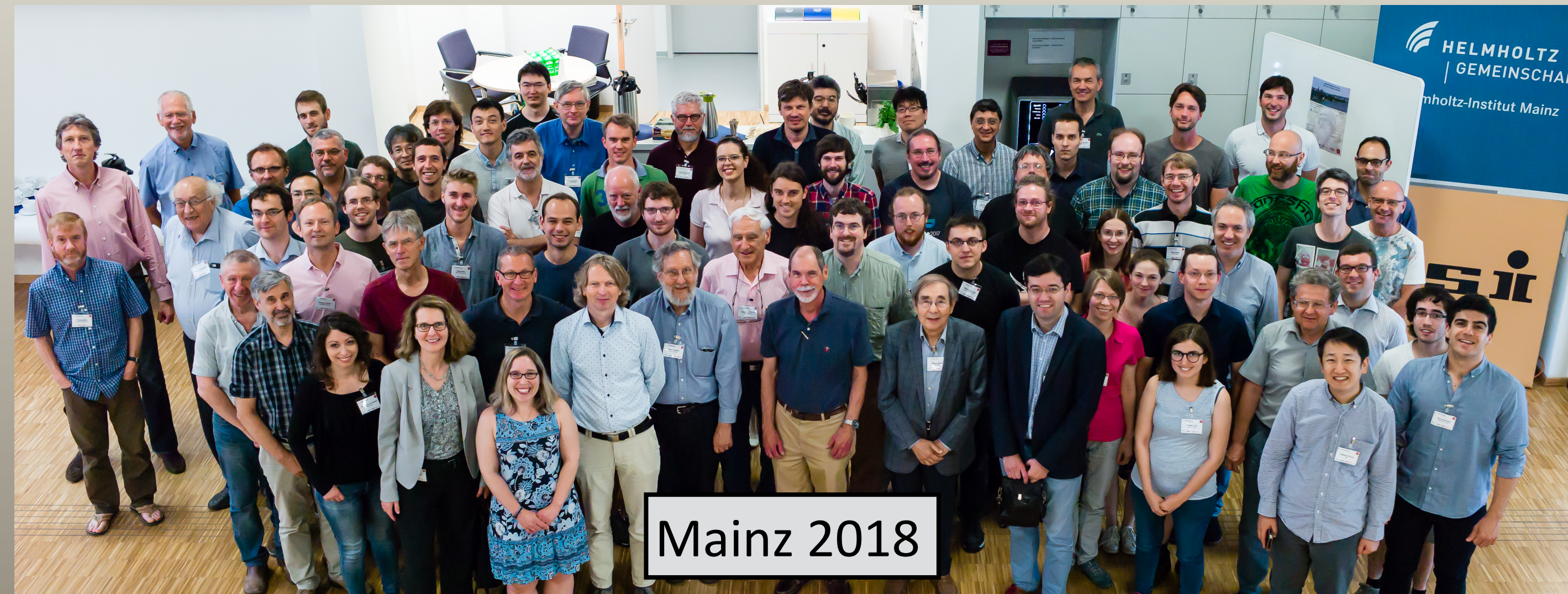
# Muon $g - 2$ Theory Initiative

Founded in 2017

Agree on common SM prediction

Focus on hadronic contributions

Prospects for increased precision





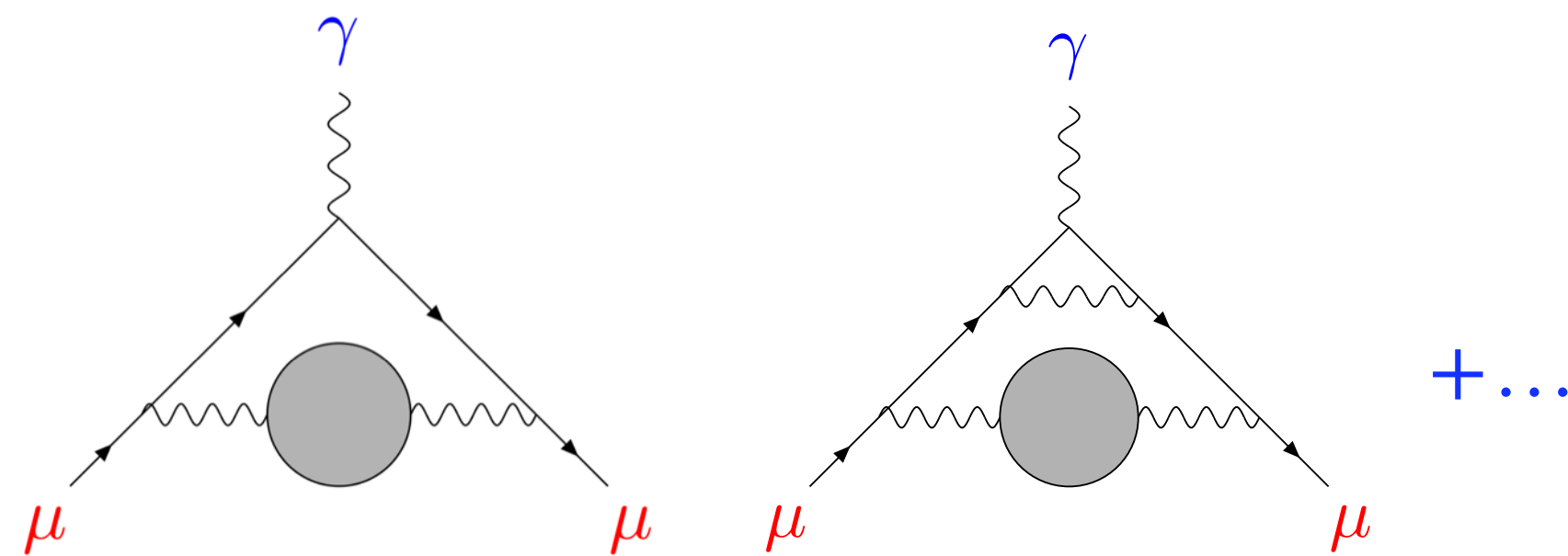
# Standard Model prediction for muon $g - 2$

[2020 White Paper]

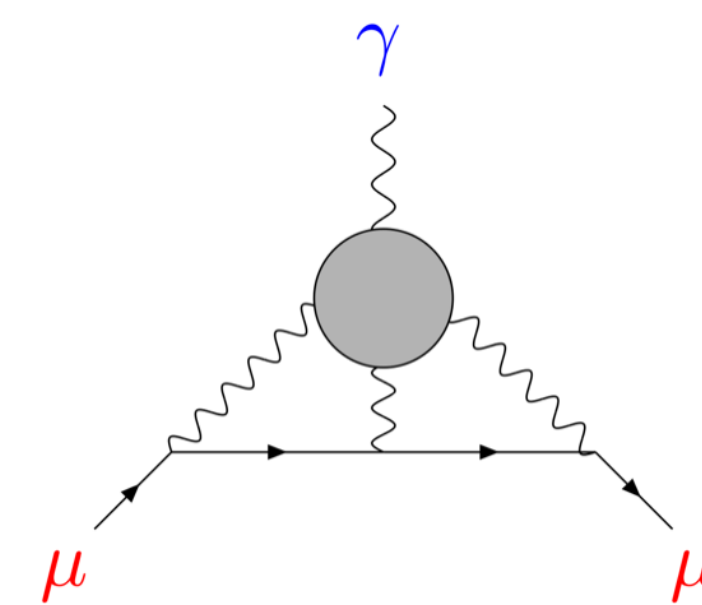
QED:	$116\,584\,718.9(1) \times 10^{-11}$	0.001 ppm	
Weak:	$153.6(1.0) \times 10^{-11}$	0.01 ppm	
Hadronic vacuum polarisation:	$6845(40) \times 10^{-11}$	0.34 ppm	[0.6%]
Hadronic light-by-light scattering:	$92(18) \times 10^{-11}$	0.15 ppm	[20%]

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$$a_{\mu}^{\text{SM}} = a_{\mu}^{\text{QED}} + a_{\mu}^{\text{weak}} + a_{\mu}^{\text{hvp}} + a_{\mu}^{\text{hlbl}} = 116\,591\,810(43) \times 10^{-11} \quad 0.37 \text{ ppm}$$



Hadronic vacuum polarisation (HVP)



Hadronic light-by-light scattering (HLbL)

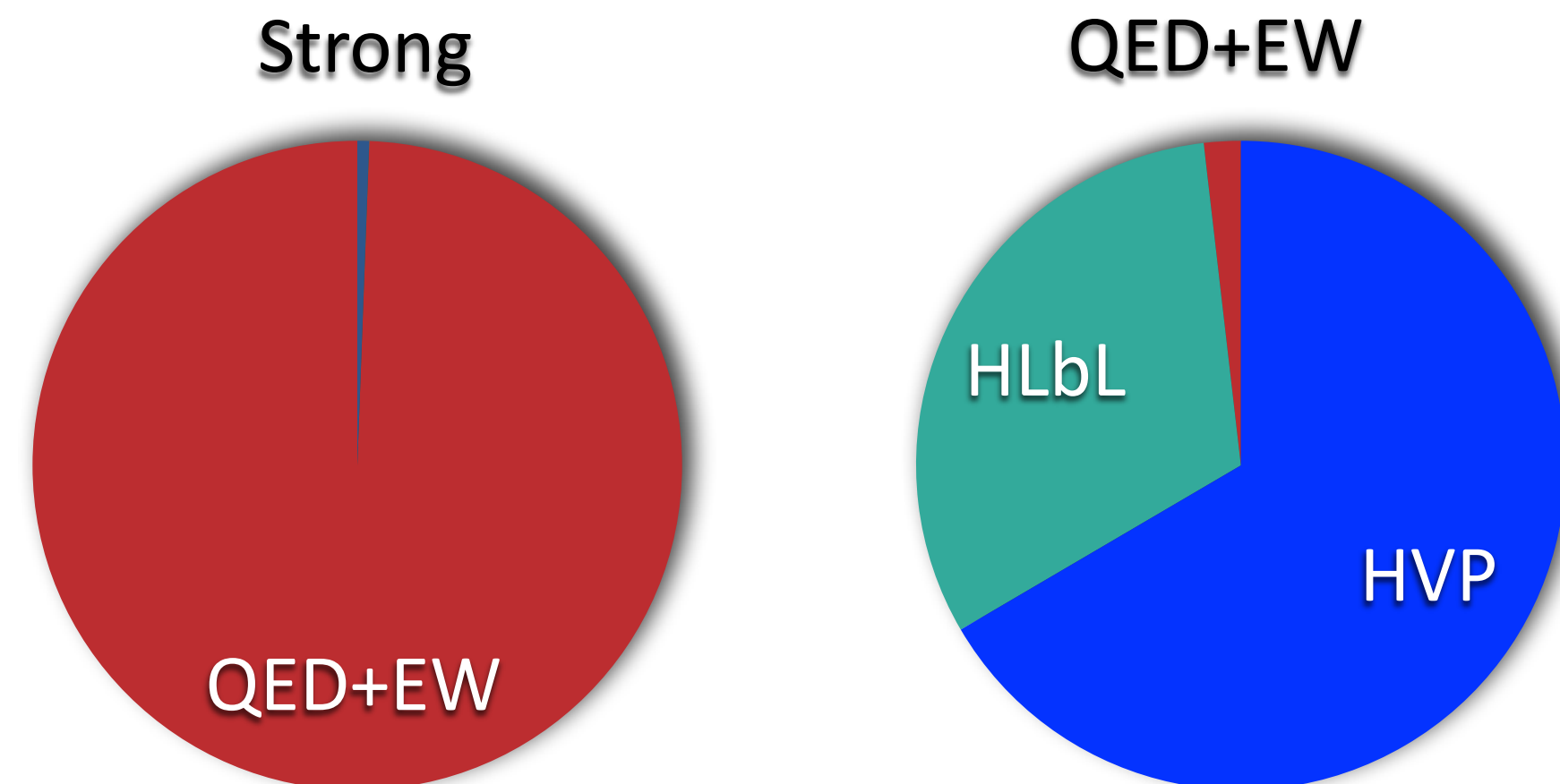
[Aoyama et al., Phys. Rep. 887 (2020) 1, arXiv:2006.04822]

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- QED and electroweak contributions account for 99.994% of the SM prediction for  $a_{\mu}$
- Error is dominated by strong interaction effects





# Standard Model prediction for muon $g - 2$

[2020 White Paper]

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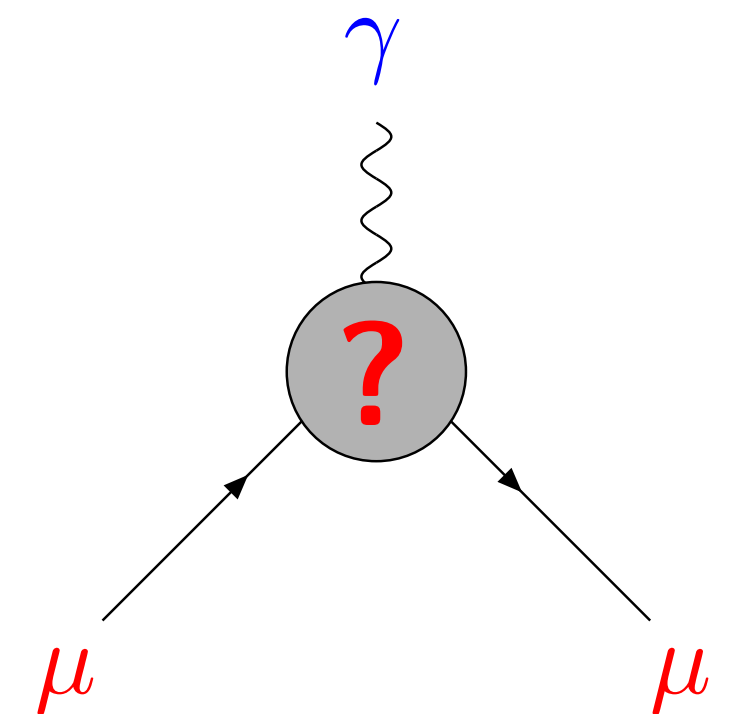

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**Standard Model vs. experiment:**  $a_{\mu}^{\text{exp}} \stackrel{?}{=} a_{\mu}^{\text{QED}} + a_{\mu}^{\text{weak}} + a_{\mu}^{\text{hvp}} + a_{\mu}^{\text{hlbl}} + a_{\mu}^{\text{BSM}}$

**Why the muon?**

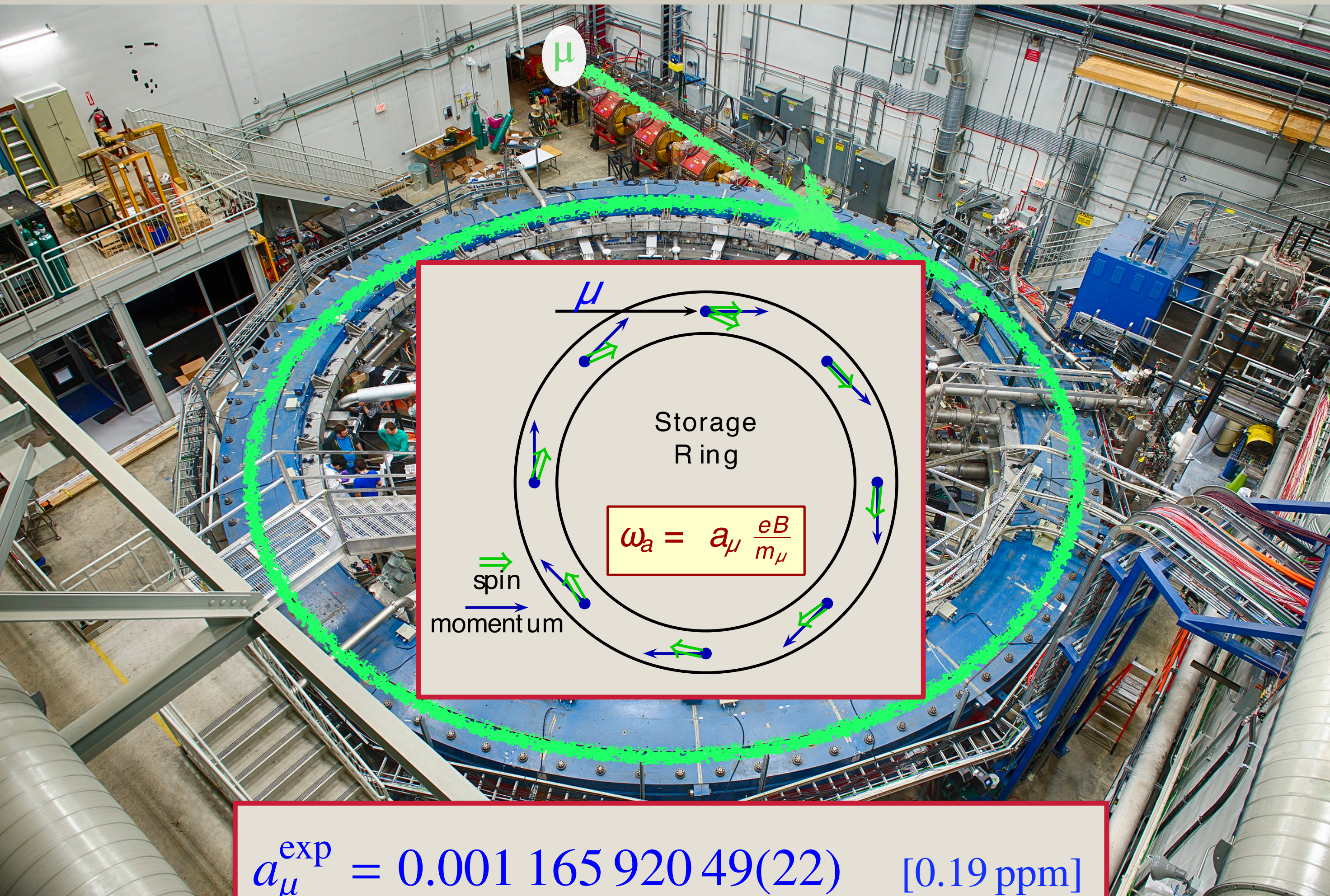
$$a_{\ell}^{\text{BSM}} \propto m_{\ell}^2 / M_{\text{BSM}}^2 \quad \ell = e, \mu, \tau$$

→ sensitivity of  $a_{\mu}$  enhanced by  $(m_{\mu}/m_e)^2 \approx 4.3 \times 10^4$  relative to  $a_e$





# Fermilab experiment E989



$$a_\mu^{\text{exp}} = 0.001\,165\,920\,49(22) \quad [0.19 \text{ ppm}]$$

[Aguillard et al., Phys Rev Lett 131 (2023) 16, 161802]





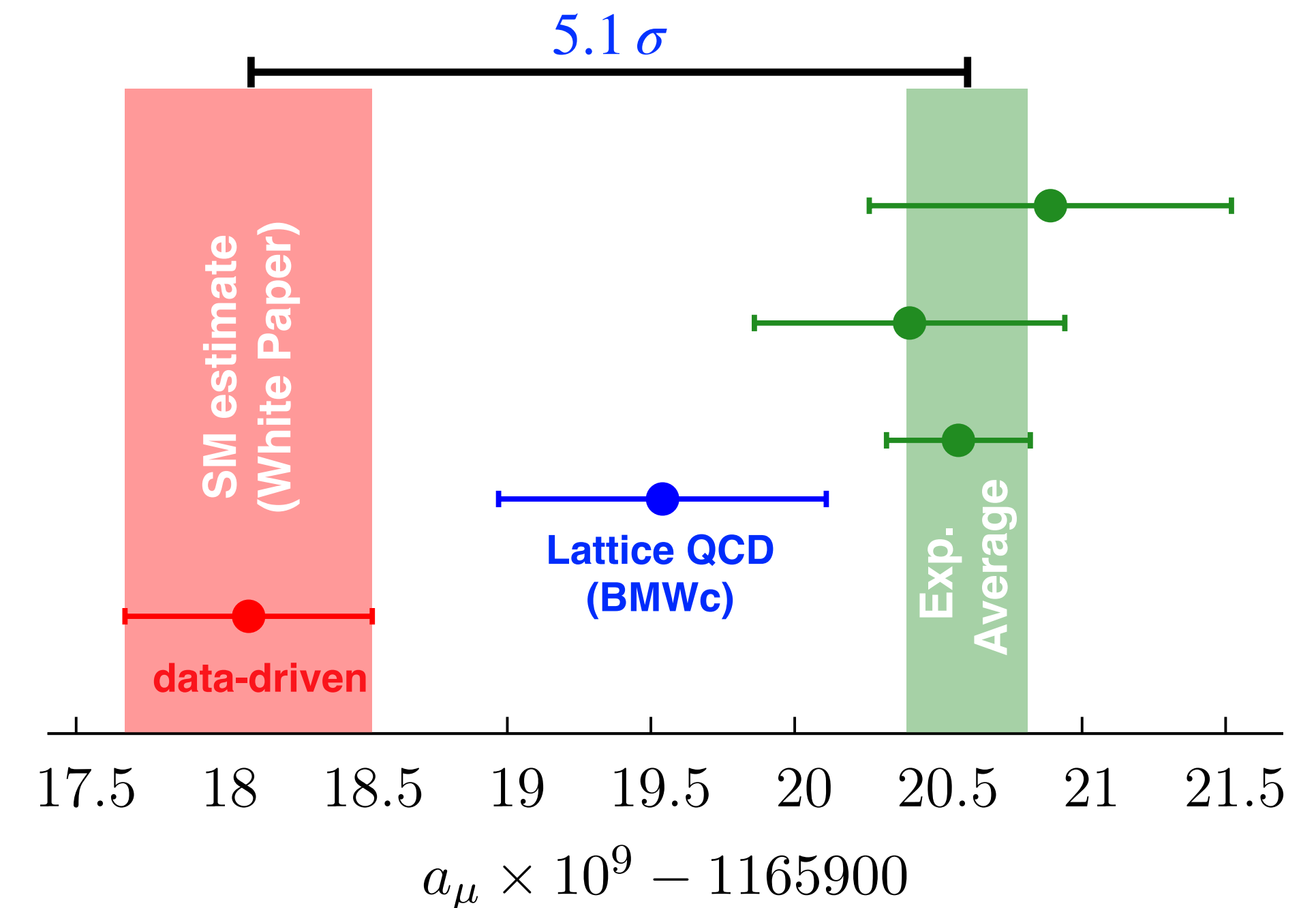
# New Physics on the horizon?

Confronting the SM prediction with the E989 measurement

$$a_{\mu}^{\text{exp}} = 116\,592\,049(22) \times 10^{-11} \quad [0.19 \text{ ppm}]$$

$$a_{\mu}^{\text{SM}} = 116\,591\,810(43) \times 10^{-11} \quad [0.37 \text{ ppm}]$$

$$\Rightarrow a_{\mu}^{\text{exp}} - a_{\mu}^{\text{SM}} = (249 \pm 48) \cdot 10^{-11} \quad [5.1\sigma]$$



Standard Model prediction:

- White paper estimate based on “**data-driven**” evaluation of (leading-order) HVP contribution: dispersion integrals and hadronic cross sections
- Lattice QCD result for HVP with comparable precision *[Borsányi et al., Nature 593 (2021) 7857]*

$$a_{\mu}^{\text{exp}} - a_{\mu}^{\text{SM}} \Big|_{\text{BMWc}}^{\text{hvp, LO}} = (105 \pm 61) \cdot 10^{-11} \quad [1.7\sigma]$$

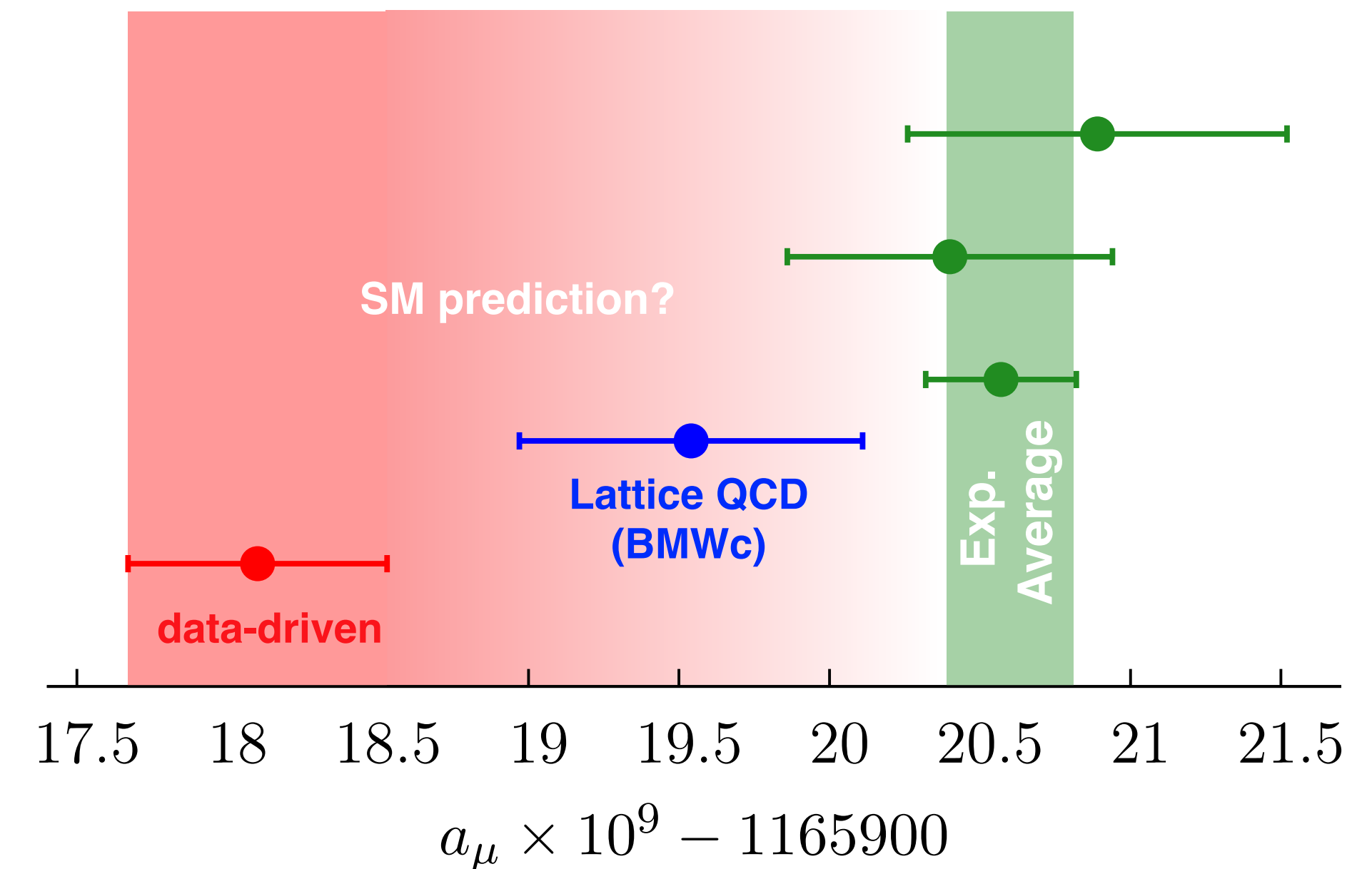
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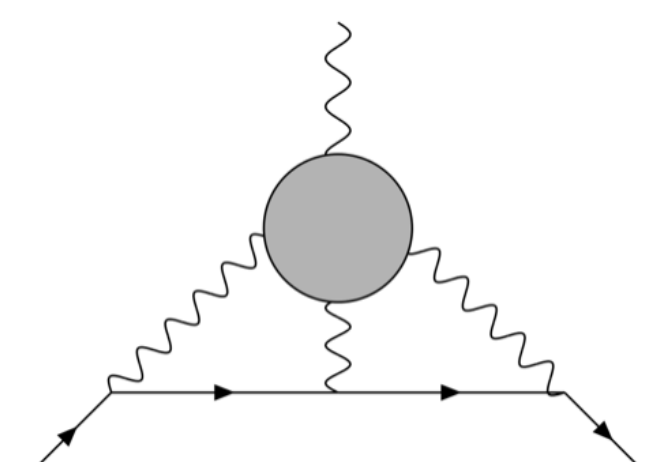
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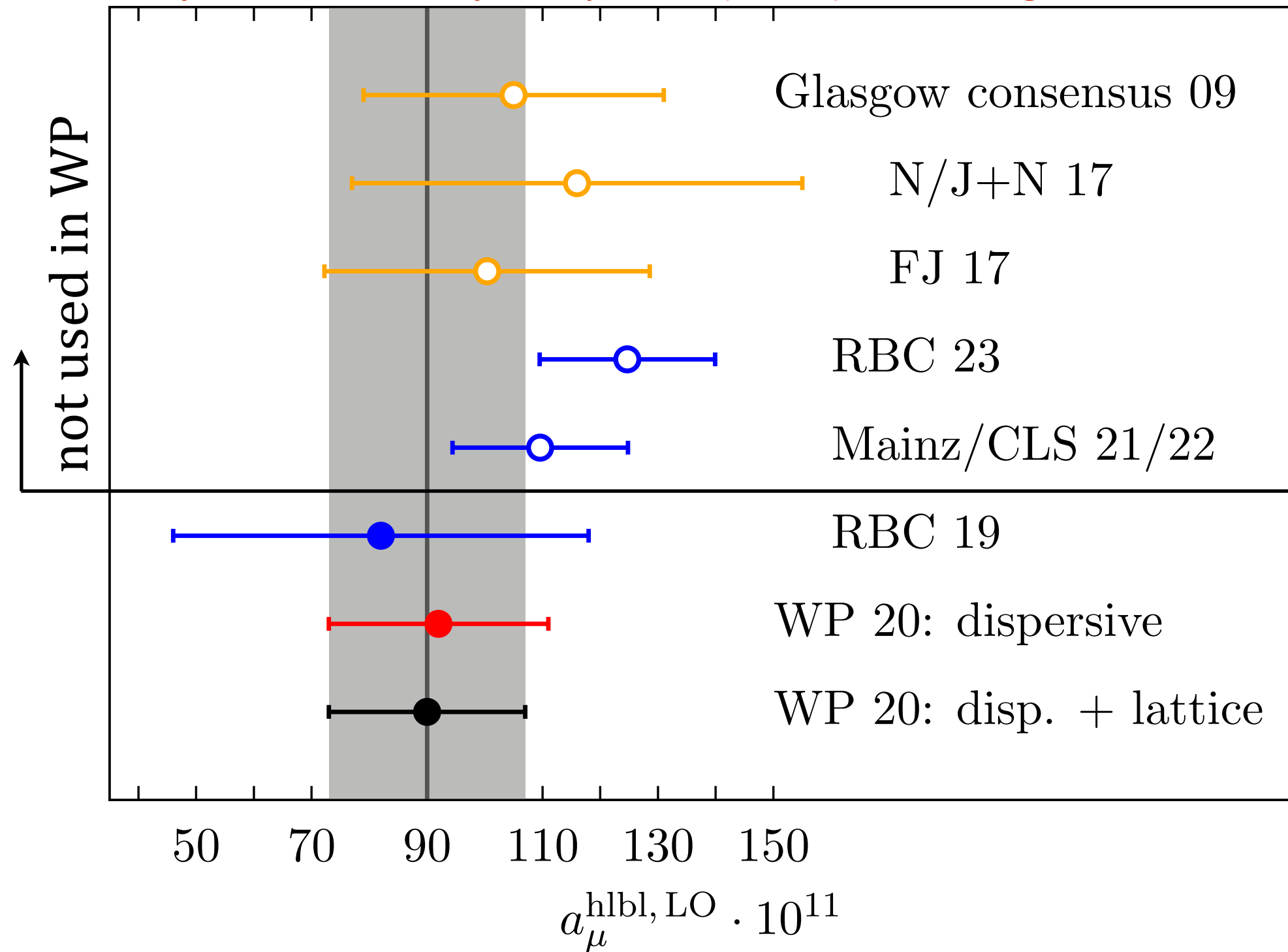
**Requires independent confirmation**



# Hadronic light-by-light scattering



[Aoyama et al., Phys. Rep. 887 (2020) 1; Colangelo et al., arXiv:2203.15810]



Hadronic models  
+ pQCD

Lattice QCD  
(+ QED)

Data-driven

Hadronic models, data-driven method and Lattice QCD produce compatible results

White paper recommended value:

$$a_\mu^{\text{hlbl}} = (92 \pm 18) \cdot 10^{-11}$$

Recent lattice calculations:

$$a_\mu^{\text{hlbl, LO}} = \begin{cases} (109.6 \pm 14.7) \cdot 10^{-11} & \text{Mainz/CLS} \\ (124.7 \pm 15.2) \cdot 10^{-11} & \text{RBC} \end{cases}$$

[Chao et al., EPJC 81 (2021) 651; EPJC 82 (2022) 664; Blum et al., arXiv:2304.04423]

$a_\mu^{\text{hlbl}}$  : **Uncontroversial** — contributes **0.15 ppm** to the total SM uncertainty of **0.37 ppm**

→ Focus on refinements and further reduction of uncertainty

# Hadronic vacuum polarisation from dispersion theory

Analyticity, unitarity & optical theorem imply:

$$\begin{aligned}
 \text{Diagram 1} &= \int \frac{ds}{\pi(s - q^2)} \text{Im} \text{Diagram 2} & 2 \text{Im} \text{Diagram 2} &= \sum_{\text{had}} \int d\Phi \left| \text{Diagram 3} \right|^2 \\
 & & & \propto \sigma(e^+ e^- \rightarrow \text{hadrons})
 \end{aligned}$$

$$a_{\mu}^{\text{hvp, LO}} = \left( \frac{\alpha m_{\mu}}{3\pi} \right)^2 \int_{m_{\pi^0}^2}^{\infty} ds \frac{R_{\text{had}}(s) \hat{K}(s)}{s^2}, \quad R_{\text{had}}(s) = \frac{3s}{4\pi (\alpha(s))^2} \sigma(e^+ e^- \rightarrow \text{hadrons}) \quad \text{“R-ratio”}$$

Hadronic effects cannot be treated in perturbation theory

- Use experimental data for  $R_{\text{had}}(s)$  in the low-energy regime (“data-driven approach”)
- Standard Model prediction is subject to experimental uncertainties

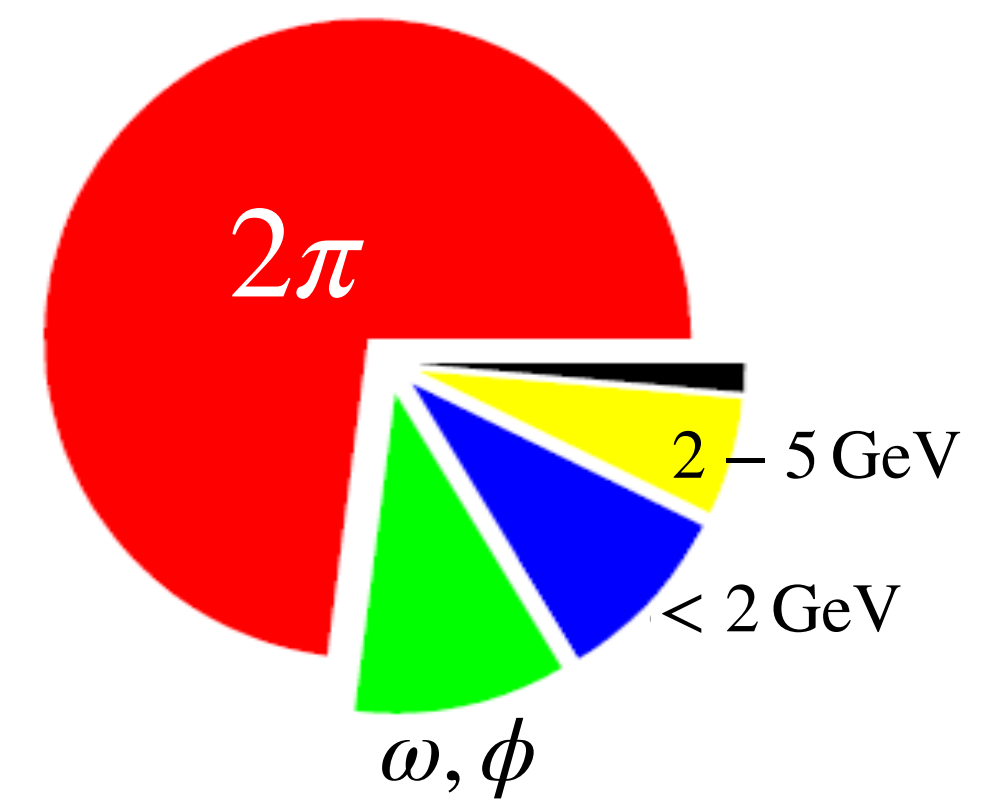
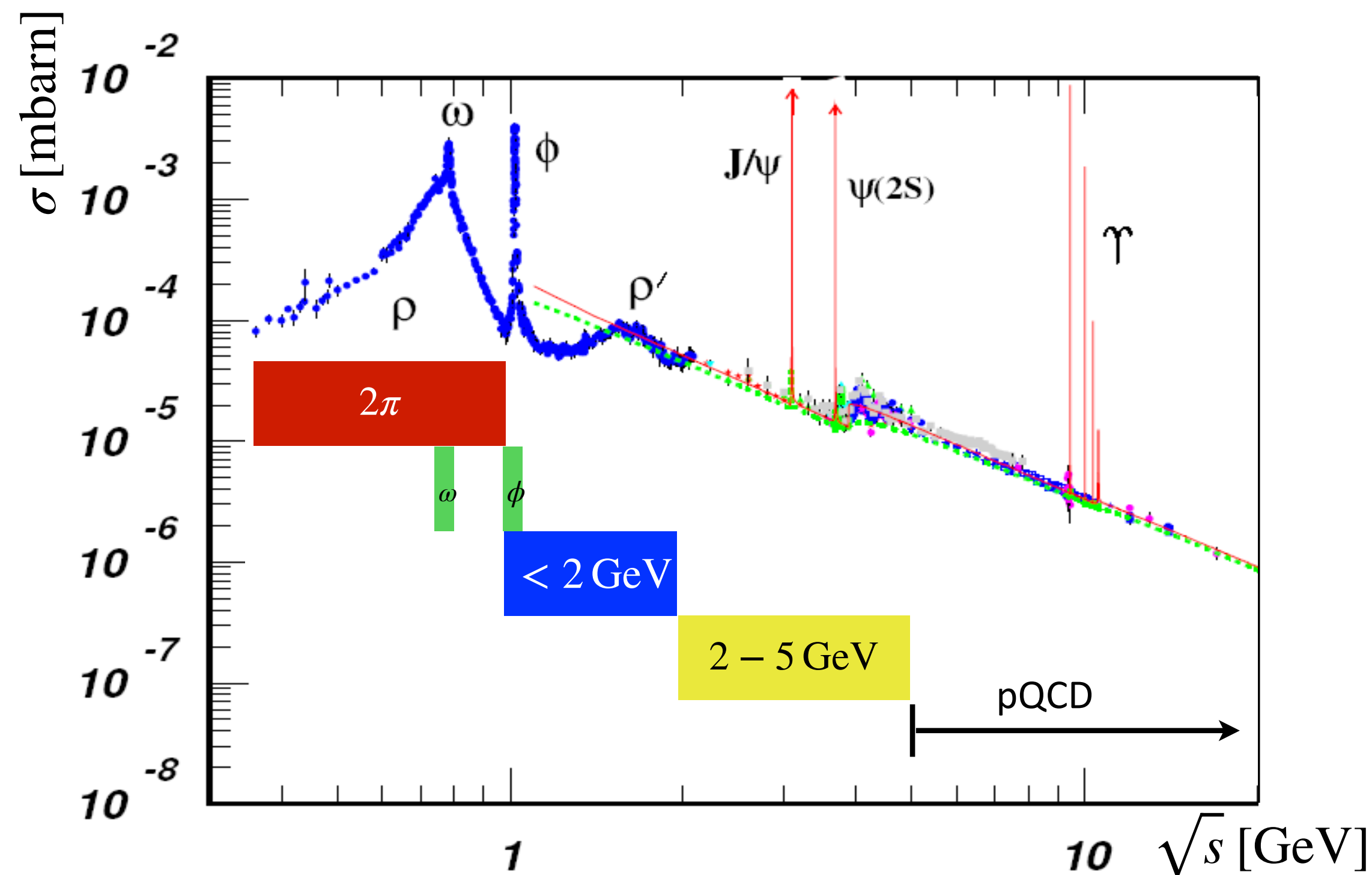
# Hadronic vacuum polarisation: Data-driven approach

Decade-long effort to measure  $e^+e^-$  cross sections

$$a_\mu^{\text{hvp, LO}} = \left(\frac{\alpha m_\mu}{3\pi}\right)^2 \int_{m_{\pi^0}^2}^{\infty} ds \frac{R_{\text{had}}(s) \hat{K}(s)}{s^2}$$

$\sqrt{s} \lesssim 2 \text{ GeV}$ : sum of exclusive channels

$\sqrt{s} > 2 \text{ GeV}$ : inclusive channels, narrow resonances, perturbative QCD



Two-pion channel contributes  $\approx 70\%$



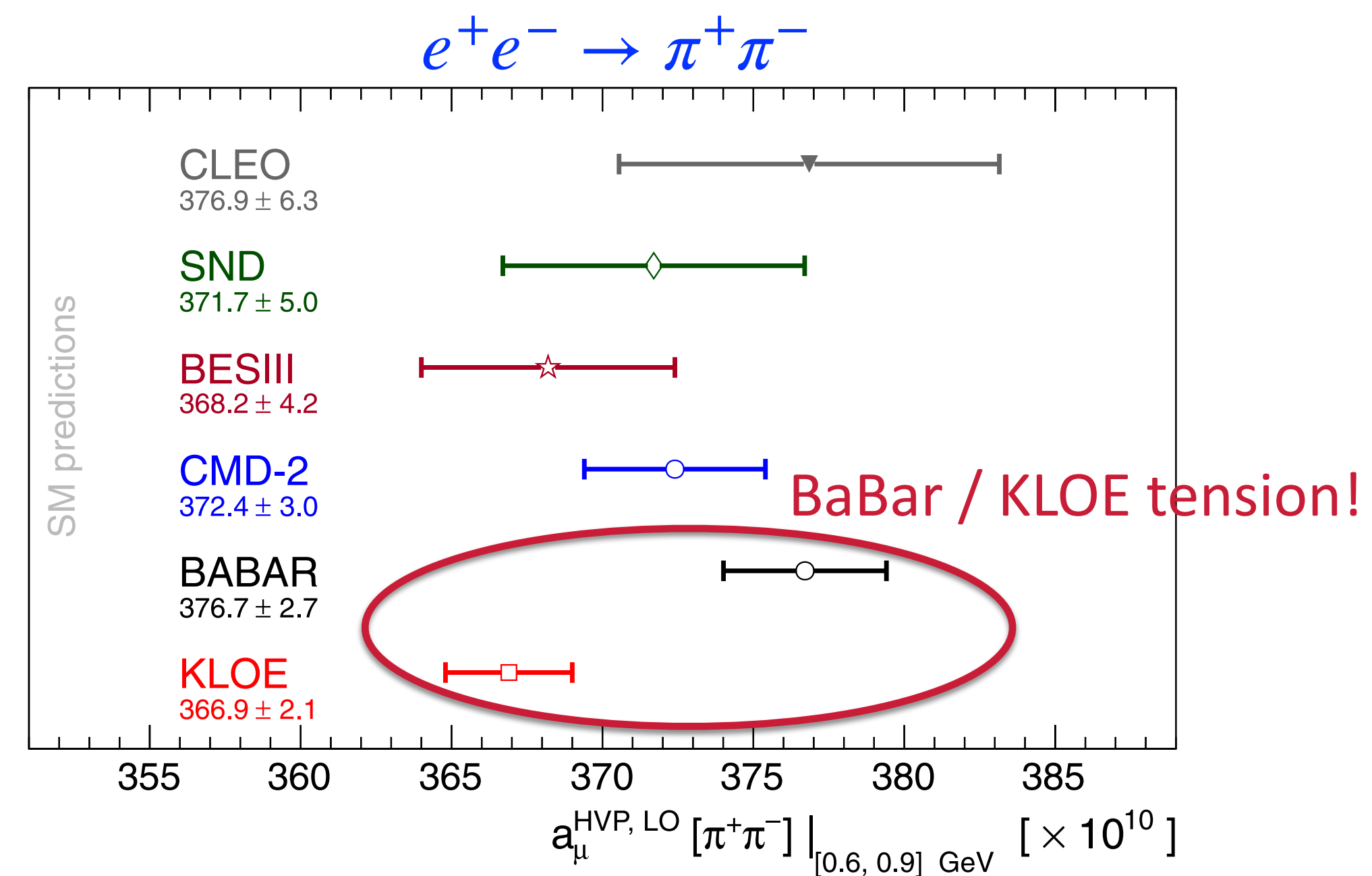
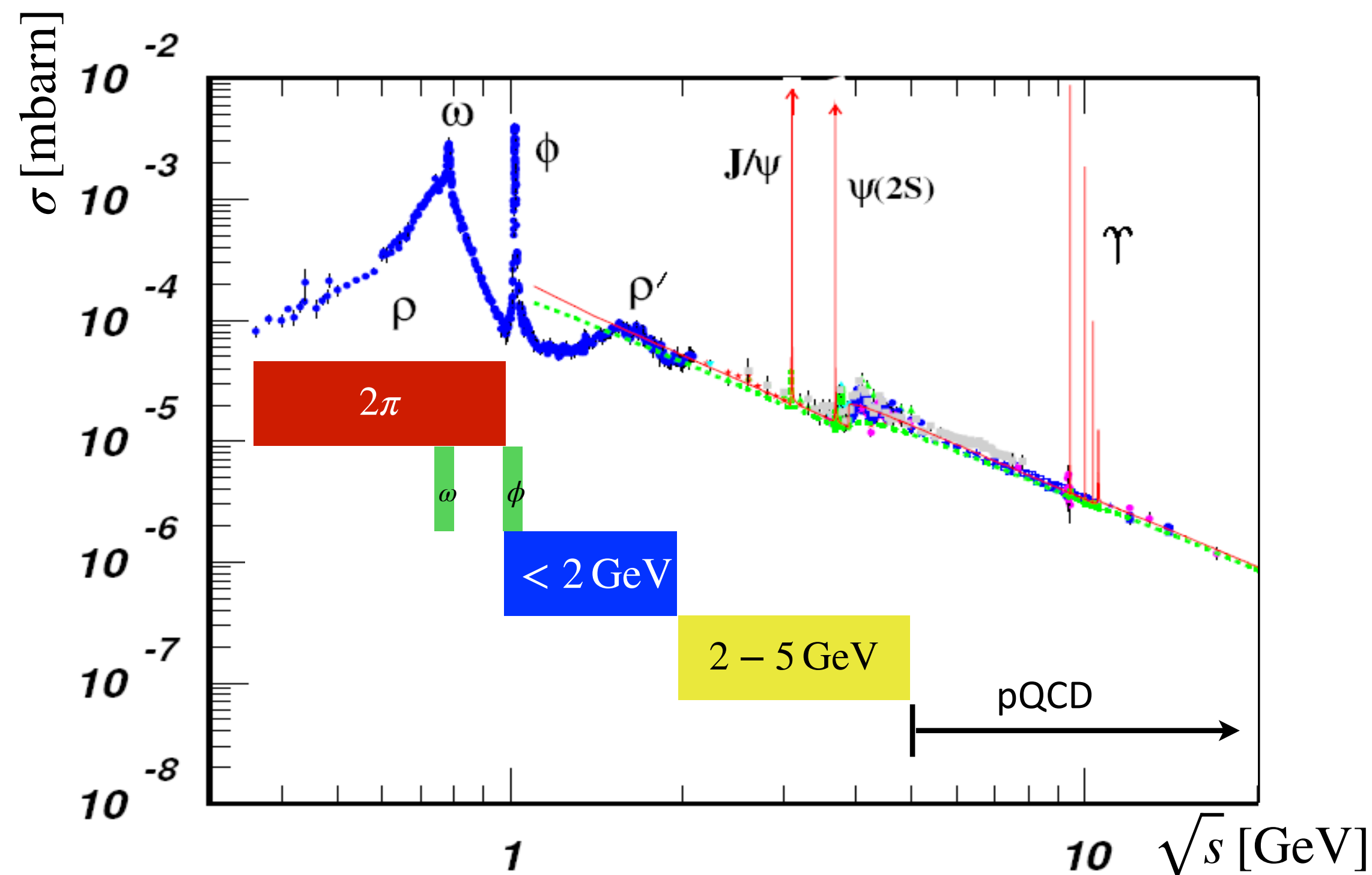
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$\sqrt{s} \lesssim 2 \text{ GeV}$ : sum of exclusive channels

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- White Paper recommended value (2020):

$$\begin{aligned} a_\mu^{\text{hvp, LO}} &= 693.1(2.8)_{\text{exp}}(2.8)_{\text{syst}}(0.7)_{\text{DV+QCD}} \times 10^{-10} \\ &= 693.1(4.0) \times 10^{-10} \quad [0.6\%] \end{aligned}$$

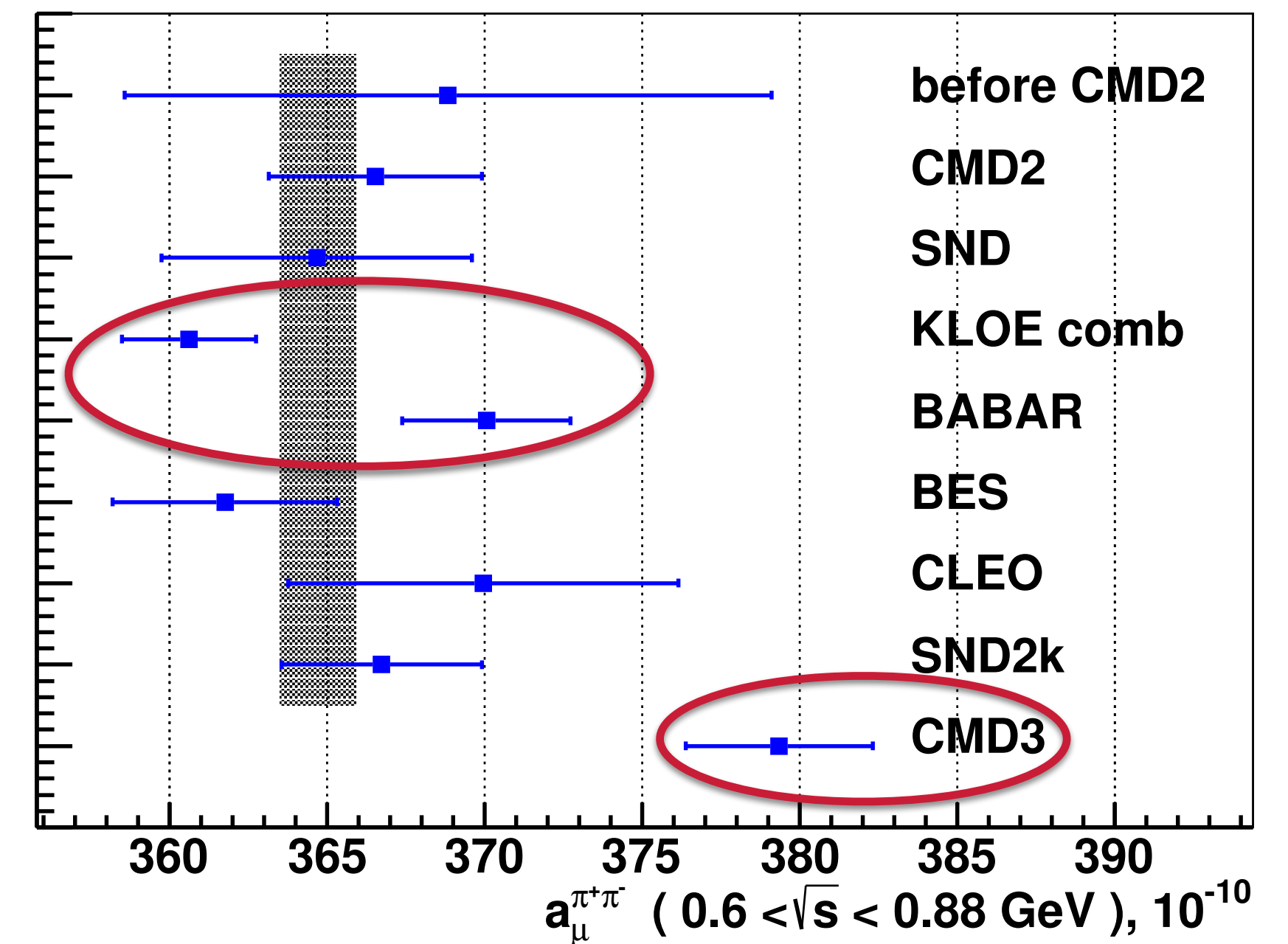
(accounts for tensions in the data and differences between analyses)

- Recent results in the  $\pi^+\pi^-$  channel by CMD-3:

→ further tension among  $e^+e^-$  data

$$a_\mu^{\text{hvp, LO}} = 707.6(3.4)_{\text{exp}}(0.7)_{\text{DV+QCD}} \times 10^{-10}$$

(my own estimate)



[Ignatov et al. (CMD-3 Collab.), arXiv:2302.08834]

# Lattice QCD

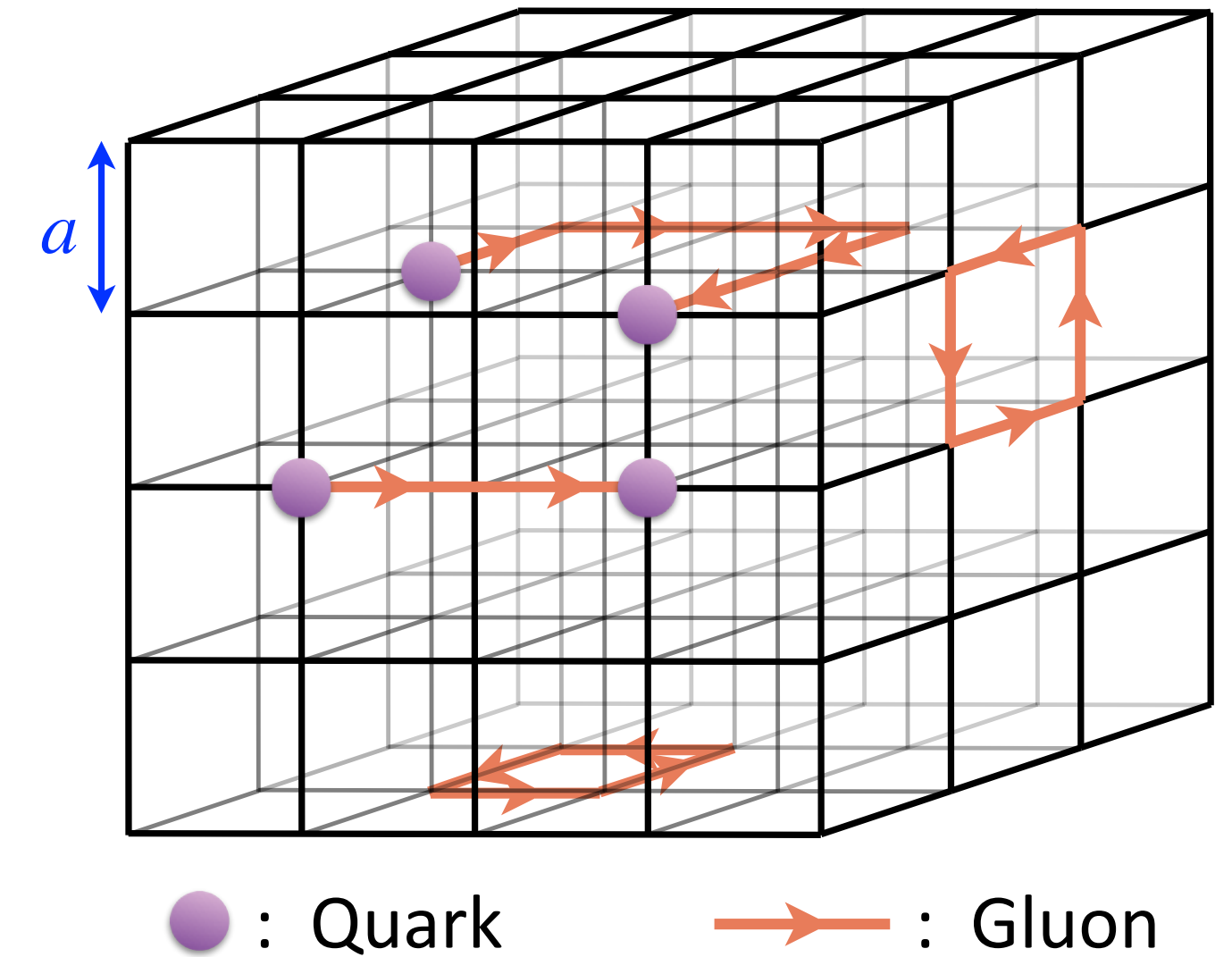
*Non-perturbative treatment of strong interaction via regularised Euclidean path integrals*

Lattice spacing:  $a$ ,  $x_\mu = n_\mu a$ ,  $a^{-1} = \Lambda_{UV}$

Expectation value:  $\langle \Omega \rangle = \frac{1}{Z} \int \prod_{x,\mu} dU_\mu(x) \Omega e^{-S_G^{\text{eff}}[U]}$

Procedure:

- Choose discretisation of QCD action
- Evaluate  $\langle \Omega \rangle$  via **Monte Carlo Integration**:  
generate **ensembles** of gauge configurations via a **Markov chain**
- **Ensemble average**:  $\langle \Omega \rangle \simeq \bar{\Omega}$       Statistical error:  $\sqrt{\bar{\Omega}^2 - \bar{\Omega}^2} \propto 1/N_{\text{cfg}}^{1/2}$
- Extrapolate observables to the **continuum limit**:  $a \rightarrow 0$  and tune quark masses to physical values

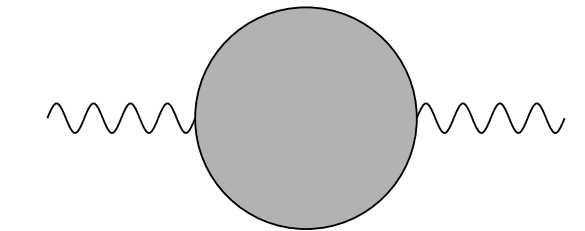




# Hadronic vacuum polarisation from Lattice QCD

Lattice QCD does **NOT** determine the  $R$ -ratio from first principles

Time-momentum representation (TMR): *[Bernecker & Meyer EPJA 47 (2011) 148]*



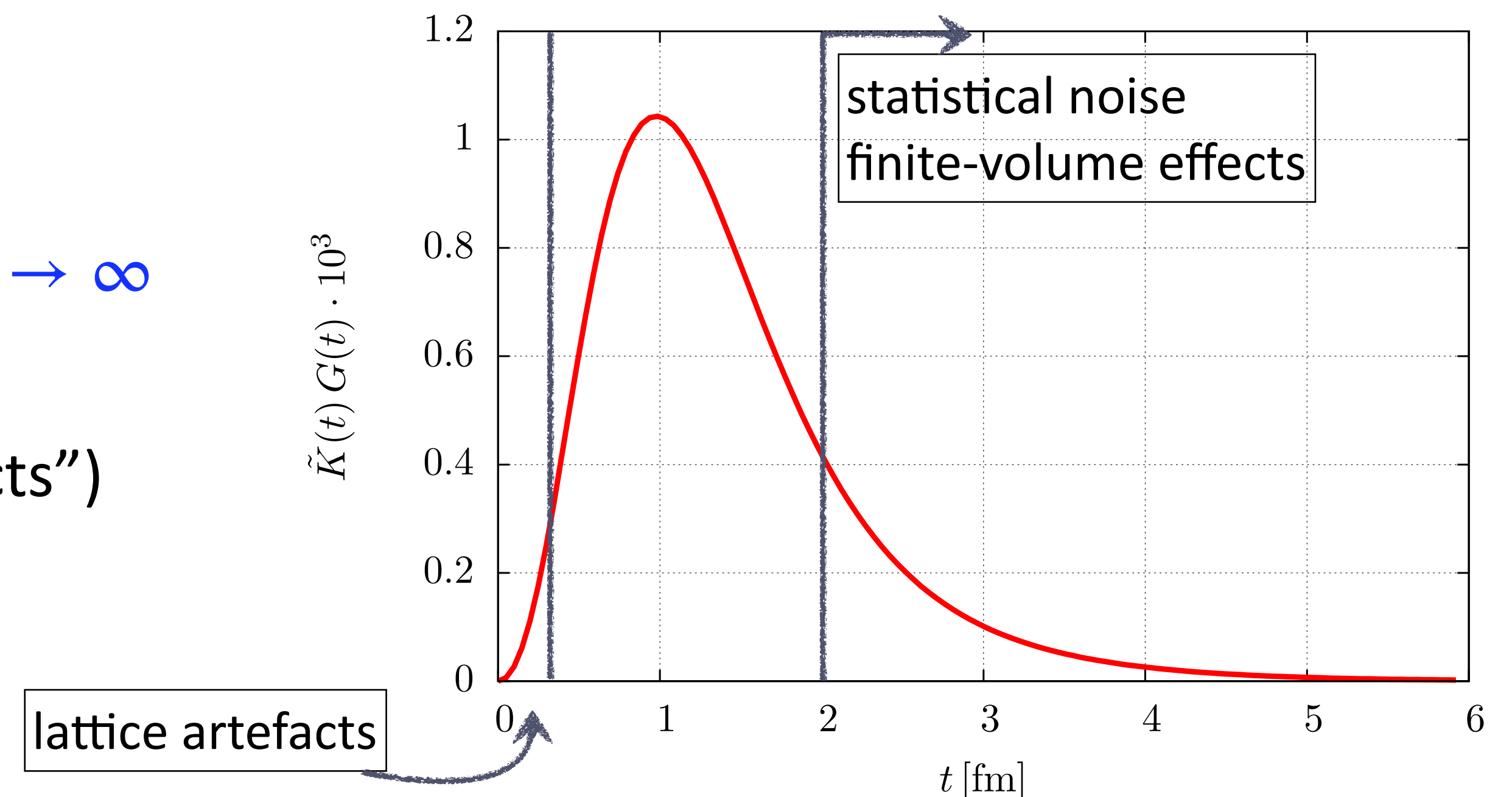
$$a_{\mu}^{\text{hvp}} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^{\infty} dt \tilde{K}(t) G(t), \quad G(t) = -a^3 \sum_{\vec{x}} \langle J_k^{\text{e.m.}}(\vec{x}, t) J_k^{\text{e.m.}}(0) \rangle$$

( $\tilde{K}(t)$ : known analytically)

- No reliance on experimental data, except for simple input quantities → scale setting, calibration
- **Not** sensitive to exclusive hadronic channels

## Challenges

- Exponentially increasing statistical noise as  $t \rightarrow \infty$
- Correct for finite-volume effects
- Control discretisation effects (“lattice artefacts”)
- Include isospin-breaking corrections



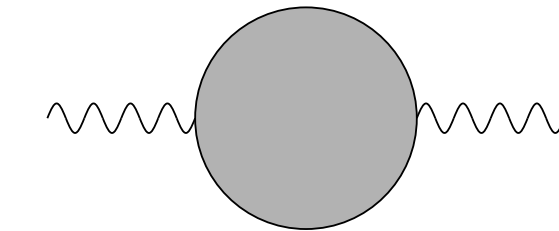


# Hadronic vacuum polarisation from Lattice QCD

Lattice QCD does **NOT** determine the  $R$ -ratio from first principles

Time-momentum representation (TMR): *[Bernecker & Meyer EPJA 47 (2011) 148]*

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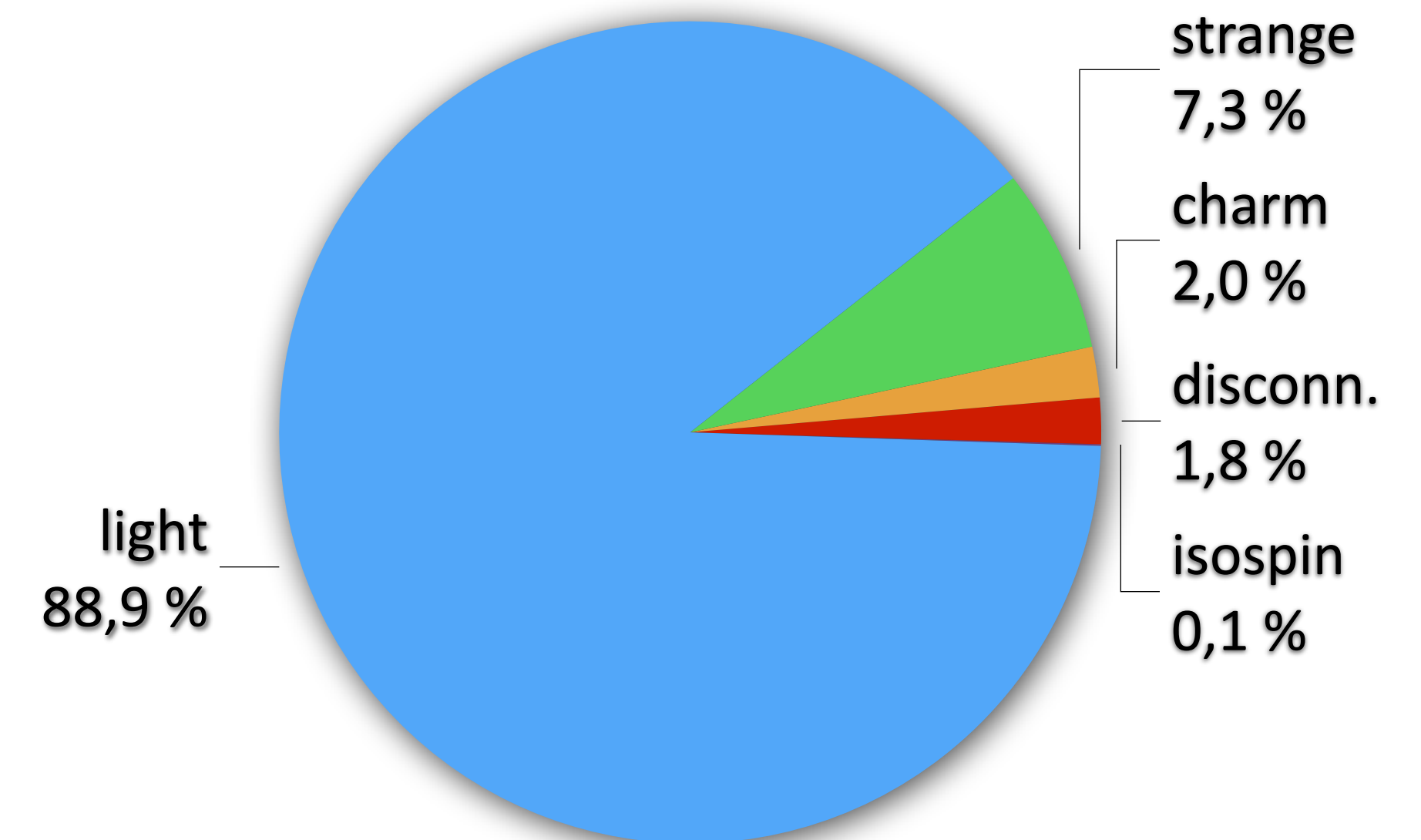
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**Light-quark connected contribution dominates**





# Controlling the long-distance tail of $G(t)$

- Long-distance tail of the light quark contribution to  $G(t)$ : limiting factor for overall statistical precision
- Correlator dominated by isovector two-pion contribution

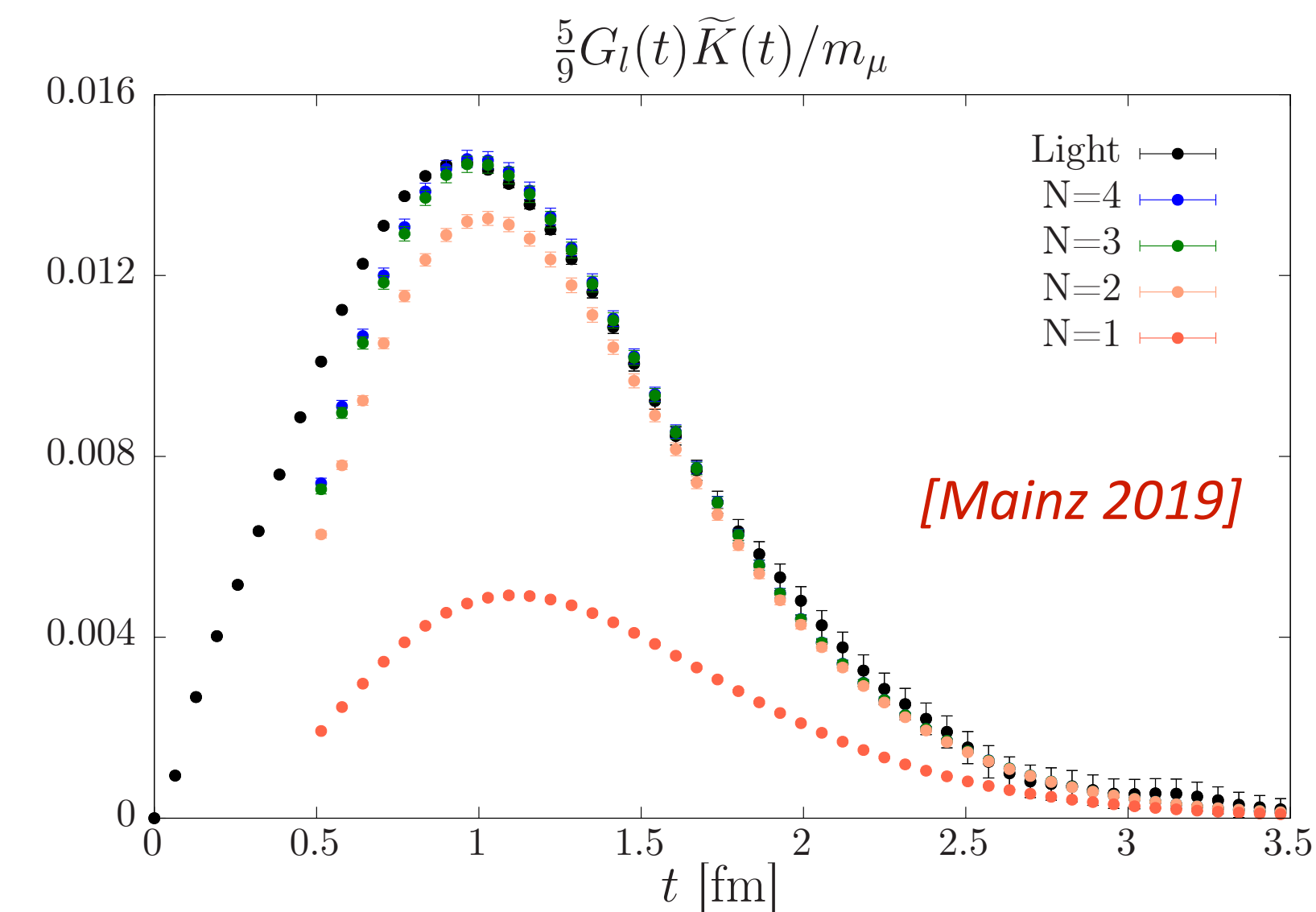
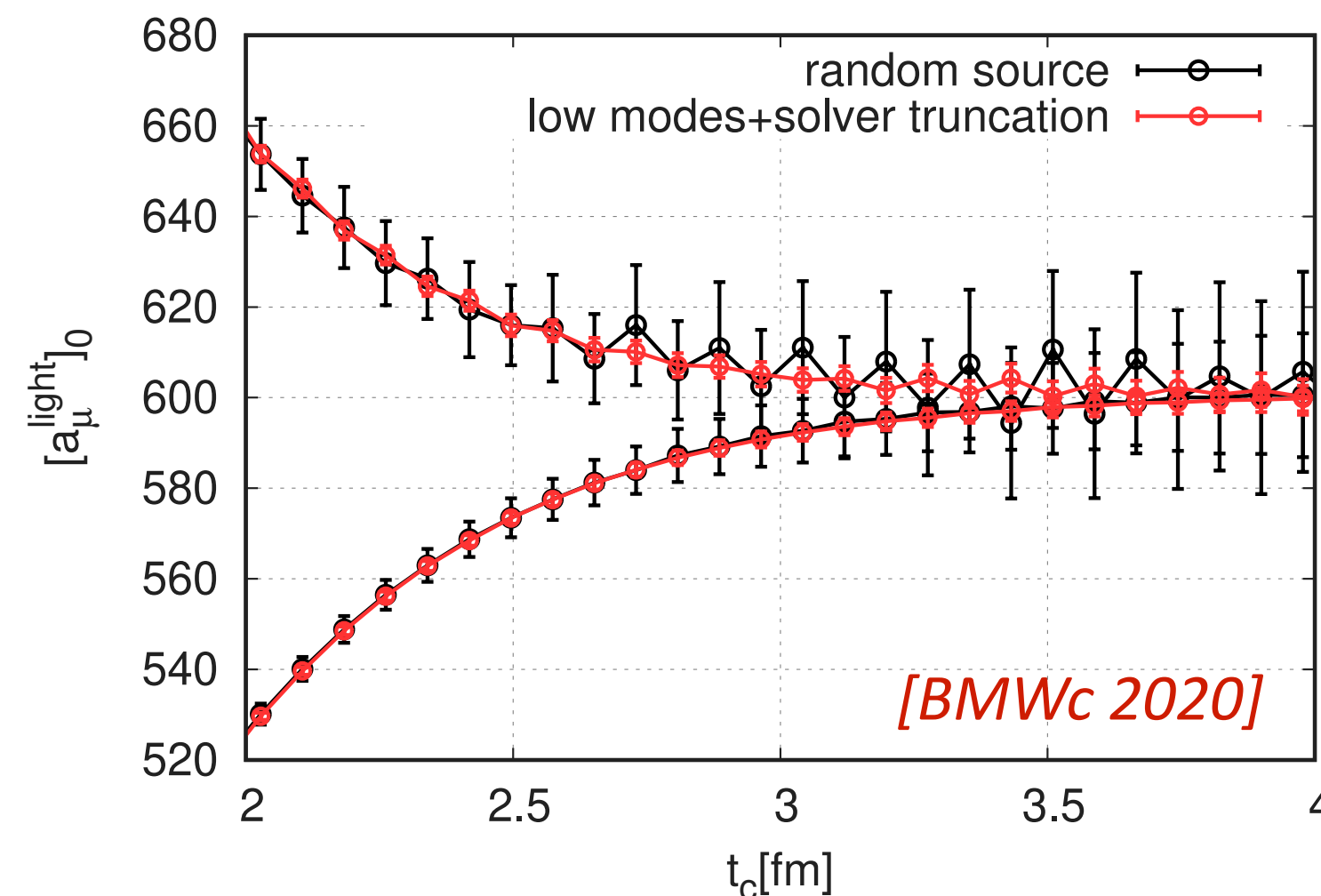
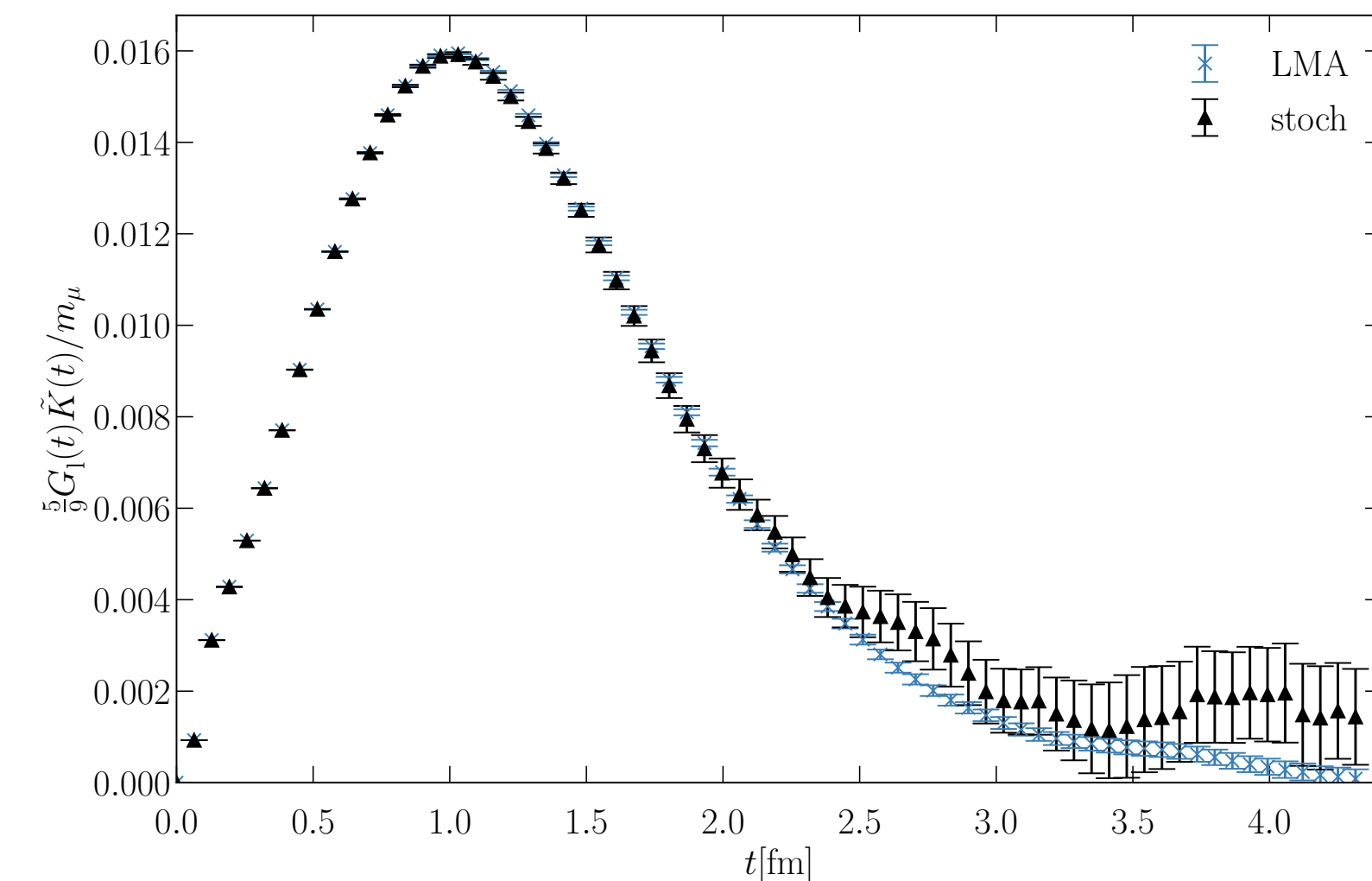
## Strategies:

- Dedicated calculations of the spectrum in isovector channel and/or pion form factor  $F_\pi(\omega)$

- “Bounding method”:

$$0 \leq G(t) \leq G(t_c) \frac{G^{\pi\pi}(t)}{G^{\pi\pi}(t_c)}$$

- Noise-reduction methods: AMA, LMA, truncated solver
- Machine Learning





# Common discretisations of the quark action

Computational cost depends significantly on the chosen discretisation

“Fermion doubling problem”

## Rooted staggered quarks:

- remnant fermion doublers — “tastes”
- correct analytically for taste-induced lattice artefacts
- used by:  
BMW, Fermilab-HPQCD-MILC, ABGP,...

## Wilson quarks:

- no doublers; chiral symmetry broken explicitly
- “exceptional configurations”:  
negative eigenvalues of Wilson-Dirac operator
- used by: Mainz/CLS, ETM, PACS

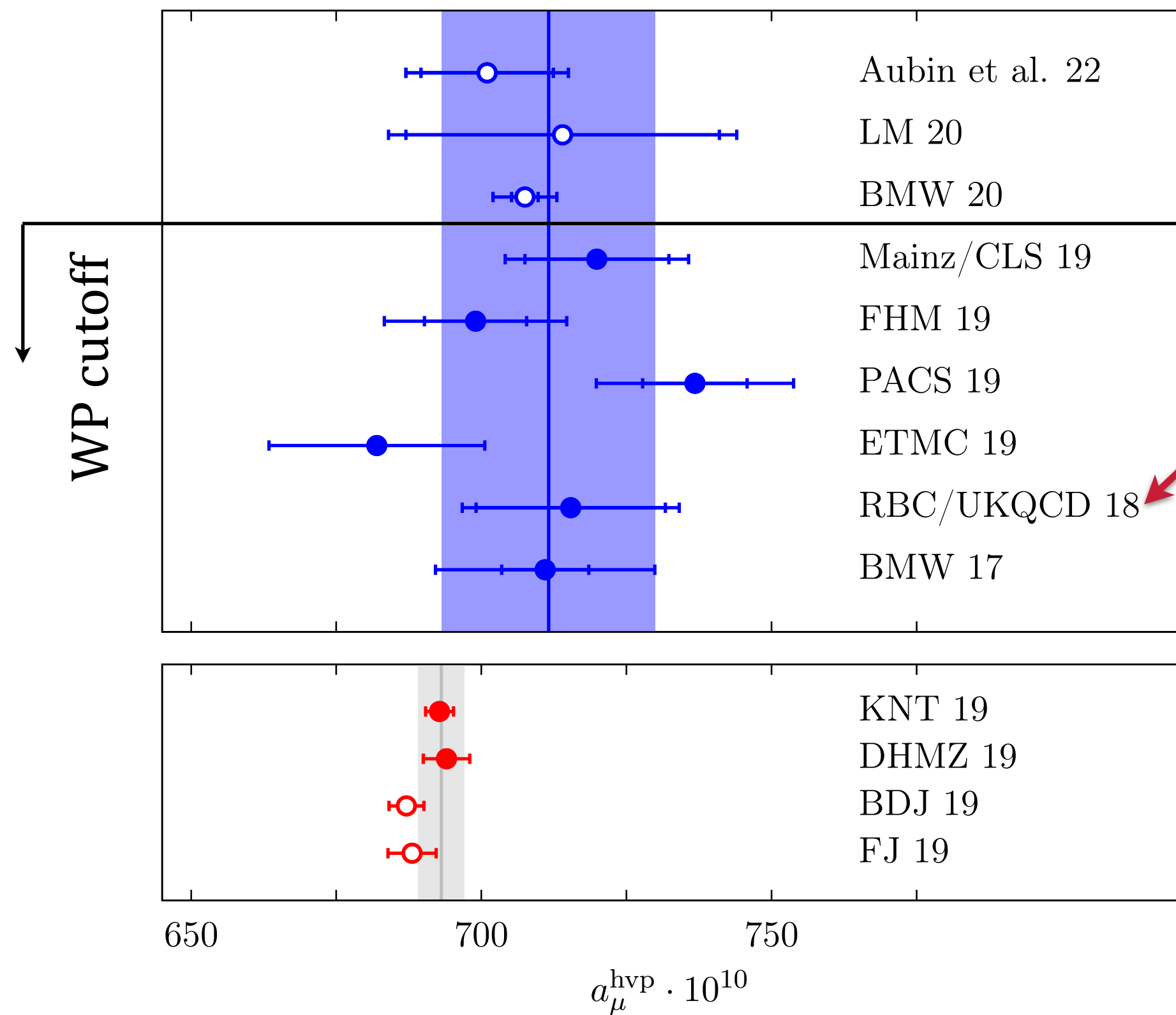
## Domain wall /overlap quarks:

- no doublers; chiral symmetry breaking exponentially small
- live in five dimensions (dwf)
- evaluate sign function of “conventional” action (ovlp)
- used by: RBC/UKQCD,  $\chi$ QCD,...

computational cost



# HVP in Lattice QCD



**RBC/UKQCD** [Blum et al., Phys. Rev. Lett. 121 (2018) 022003]

- Domain wall fermions
- Two ensembles:  $a = 0.114, 0.084$  fm at  $m_\pi^{\text{phys}}$
- Leading isospin-breaking corrections included
- Naive continuum extrapol'n in  $a^2$  including estimated  $a^4$ -term

$$a_\mu^{\text{hvp, LO}} = (715.4 \pm 16.3 \pm 9.2) \cdot 10^{-10} \quad [2.6\%]$$

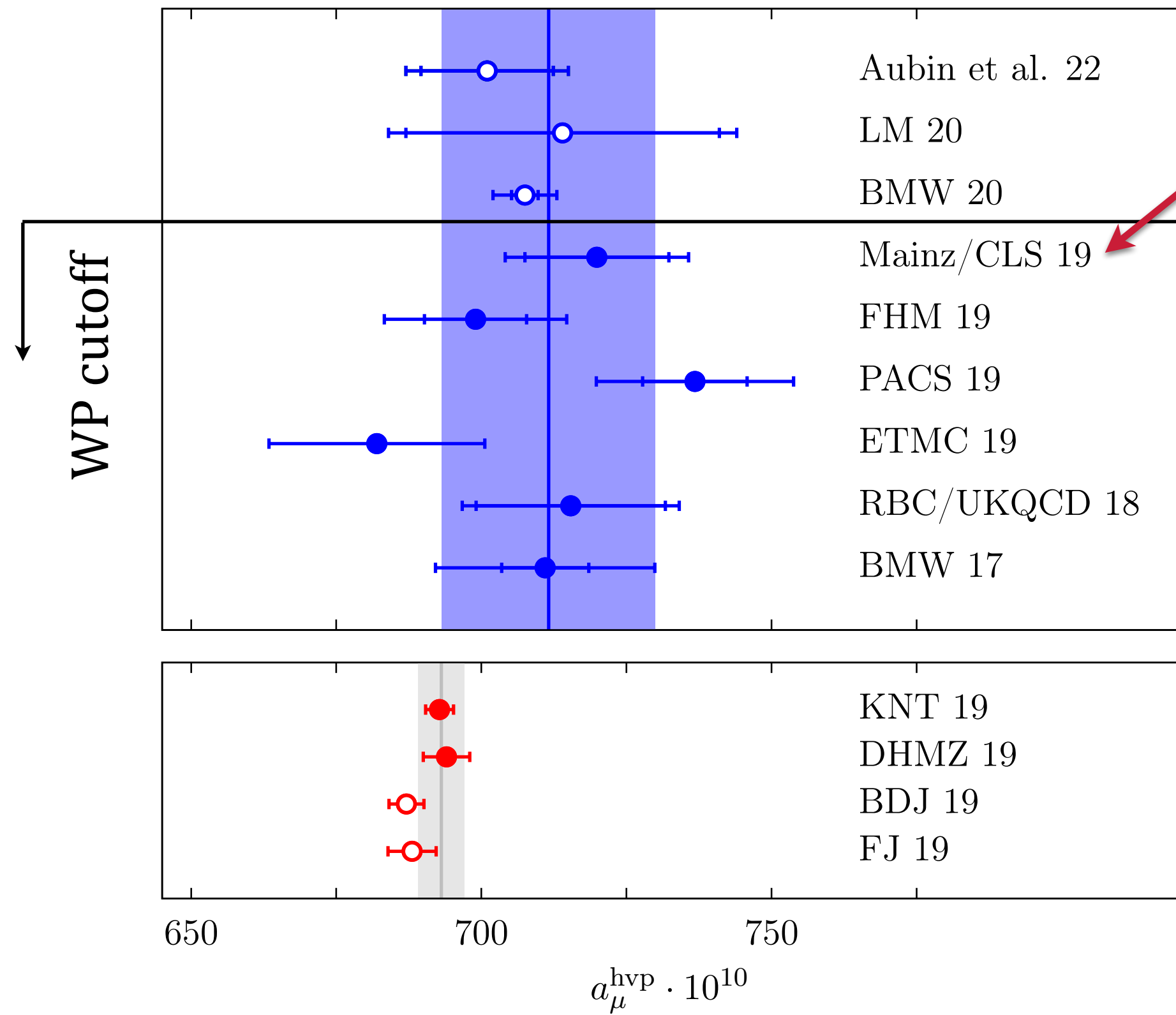
White Paper:

R-ratio:  $a_\mu^{\text{hvp, LO}} = (693.1 \pm 4.0) \cdot 10^{-10} \quad [0.6\%]$

LQCD:  $a_\mu^{\text{hvp, LO}} = (711.6 \pm 18.4) \cdot 10^{-10} \quad [2.6\%]$



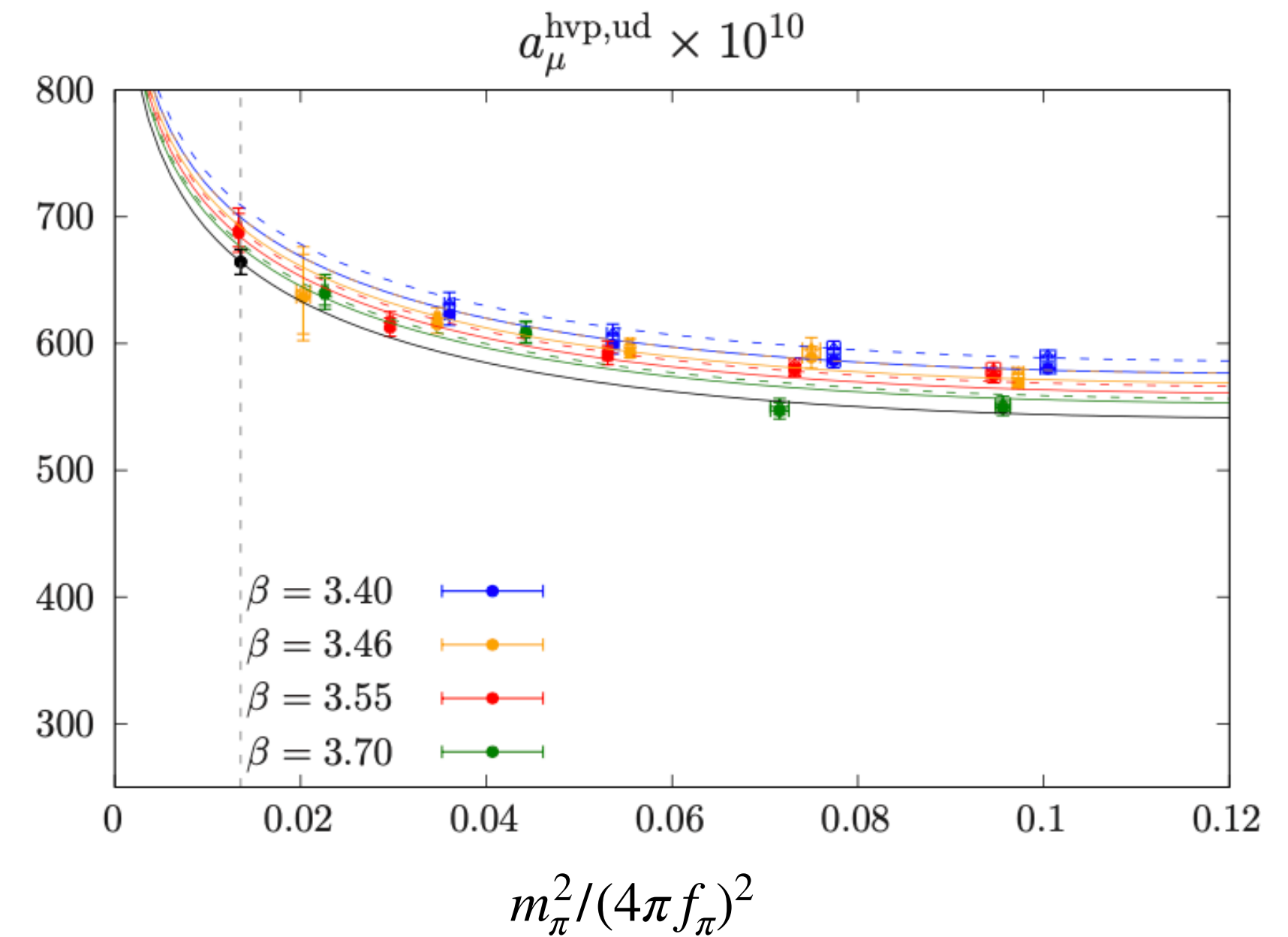
# HVP in Lattice QCD



## Mainz/CLS

[Gérardin et al., Phys. Rev. D 100 (2019) 014510]

- $O(a)$  improved Wilson fermions
- Four lattice spacings:  $a = 0.085 - 0.050$  fm
- Pion masses  $m_\pi = 130 - 420$  MeV
- Isospin-breaking correction by ETMC added to error
- Simultaneous chiral and continuum extrapolation



White Paper:

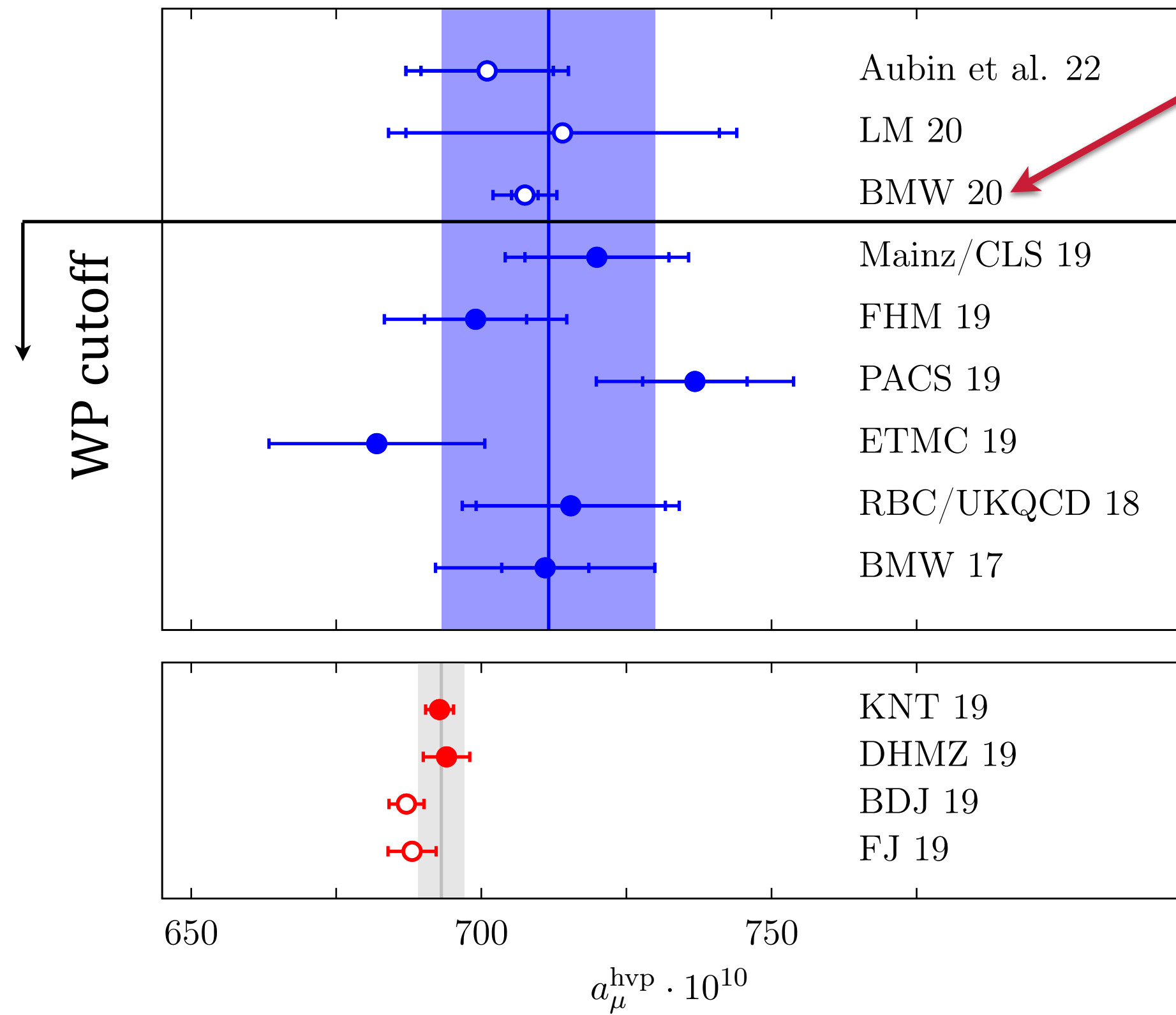
$R$ -ratio:  $a_\mu^{\text{hvp,LO}} = (693.1 \pm 4.0) \cdot 10^{-10}$  [0.6%]

LQCD:  $a_\mu^{\text{hvp,LO}} = (711.6 \pm 18.4) \cdot 10^{-10}$  [2.6%]

$a_\mu^{\text{hvp,LO}} = (720.0 \pm 12.6 \pm 9.9) \cdot 10^{-10}$  [2.2%]



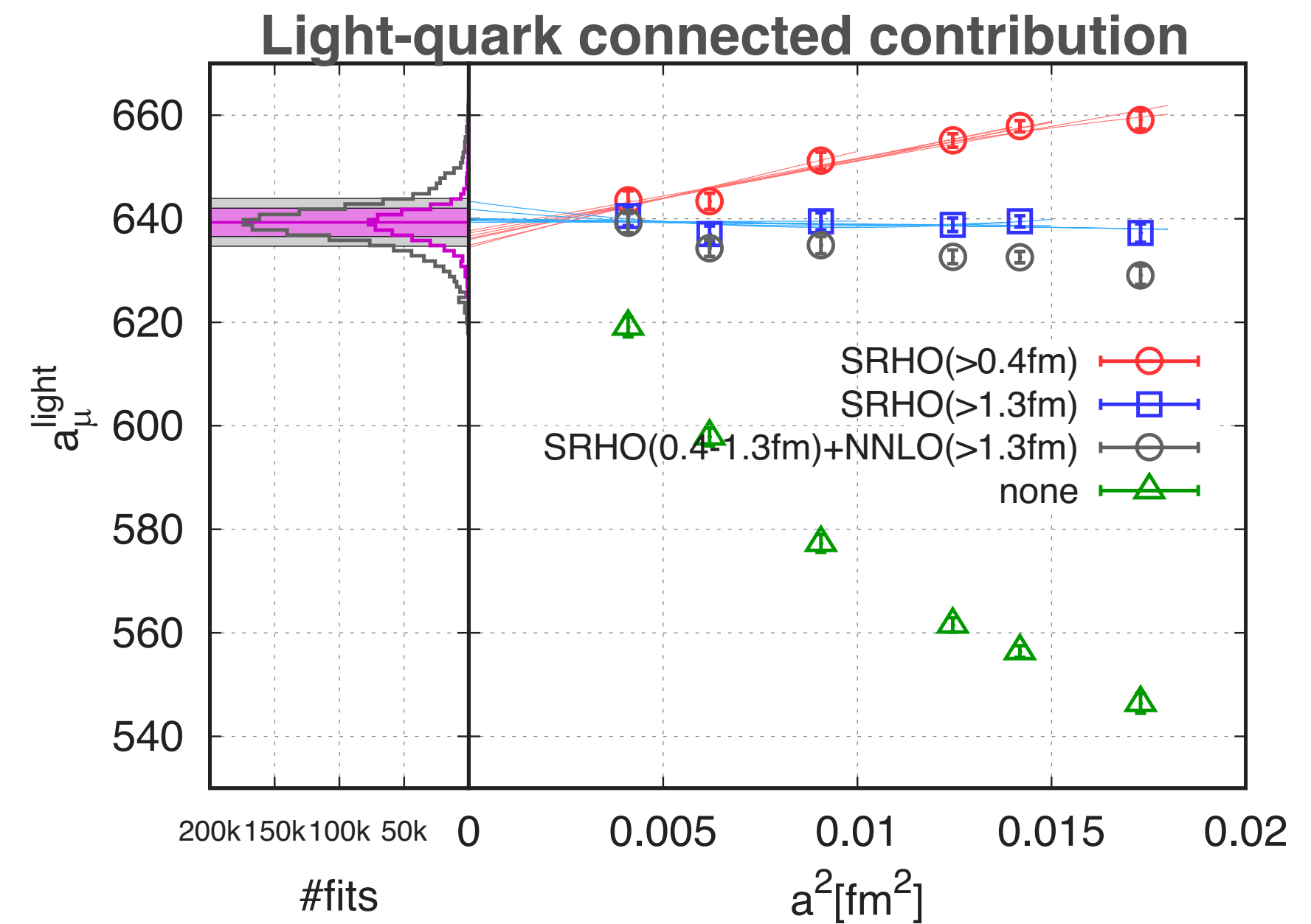
# HVP in Lattice QCD



## BMWc

[Borsányi et al., Nature 593 (2021) 7857]

- Rooted staggered fermions
- Six lattice spacings:  $a = 0.132 - 0.064$  fm
- Physical pion mass throughout
- Correct for taste-breaking before continuum extrapol'n
- Final result selected from distribution of different fits



$$a_{\mu}^{\text{hvp, LO}} = (707.5 \pm 2.3 \pm 5.0) \cdot 10^{-10} \quad [0.8\%]$$

White Paper:

$$R\text{-ratio: } a_{\mu}^{\text{hvp, LO}} = (693.1 \pm 4.0) \cdot 10^{-10} \quad [0.6\%]$$

$$\text{LQCD: } a_{\mu}^{\text{hvp, LO}} = (711.6 \pm 18.4) \cdot 10^{-10} \quad [2.6\%]$$

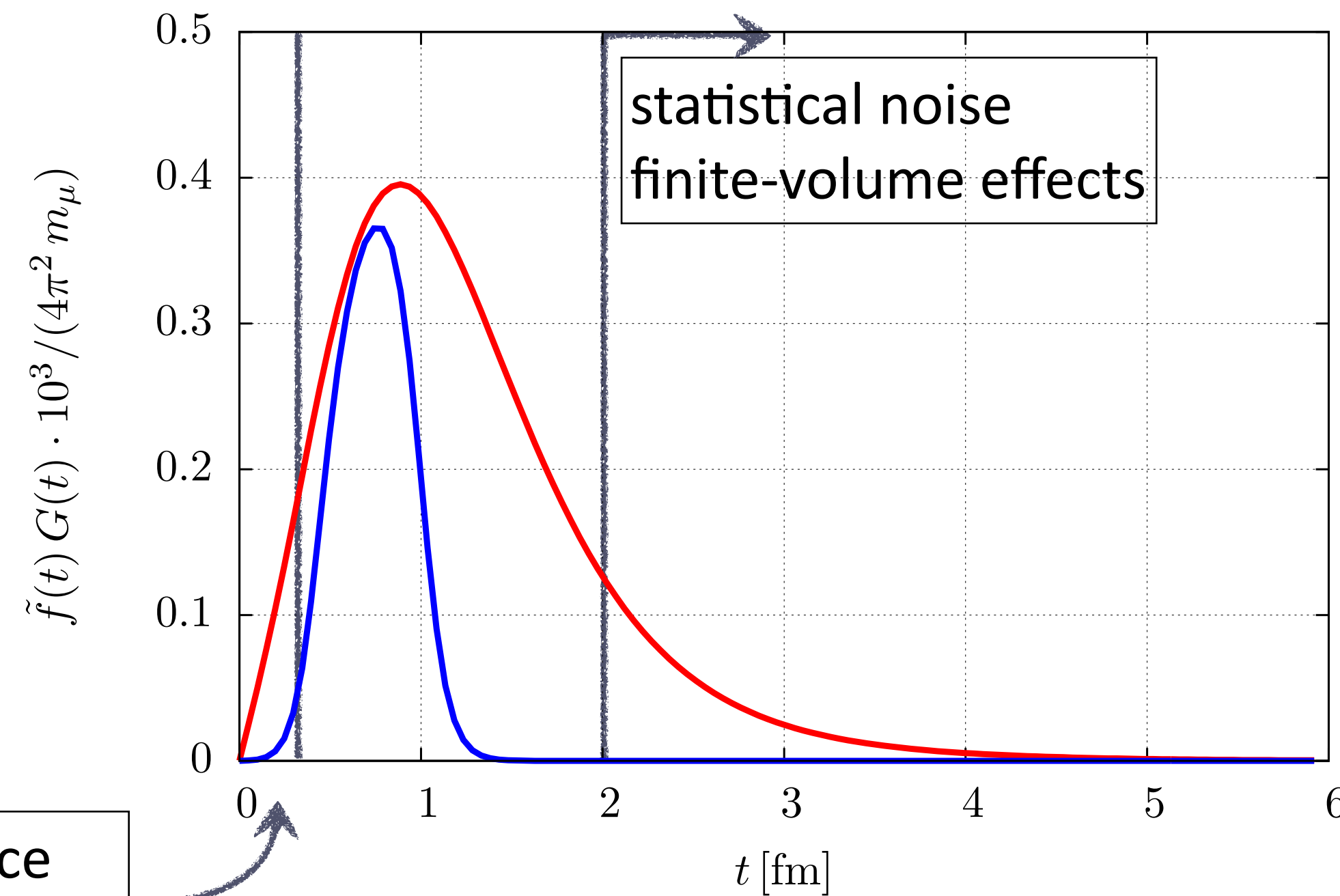
# Window observables

[Blum et al., Phys. Rev. Lett. 121 (2018) 022003]

**Idea:** restrict integration to “unproblematic” regions

→ reduce statistical fluctuations and systematic effects

$$a_{\mu}^{\text{hvp, win}} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^{\infty} dt \tilde{K}(t) G(t) W(t; t_0, t_1)$$



Intermediate-distance window:

$$W^{\text{ID}}(t; t_0, t_1) = \Theta(t, t_0, \Delta) - \Theta(t, t_1, \Delta)$$

$$\Theta(t, t', \Delta) = \frac{1}{2} [1 + \tanh(t - t')/\Delta]$$

$$t_0 = 0.4 \text{ fm}, t_1 = 1.0 \text{ fm}, \Delta = 0.15 \text{ fm}$$

- Finite-volume correction reduced to 0.25%
- Uncertainty dominated by statistics

→ Benchmark quantity for sub-contribution of HVP

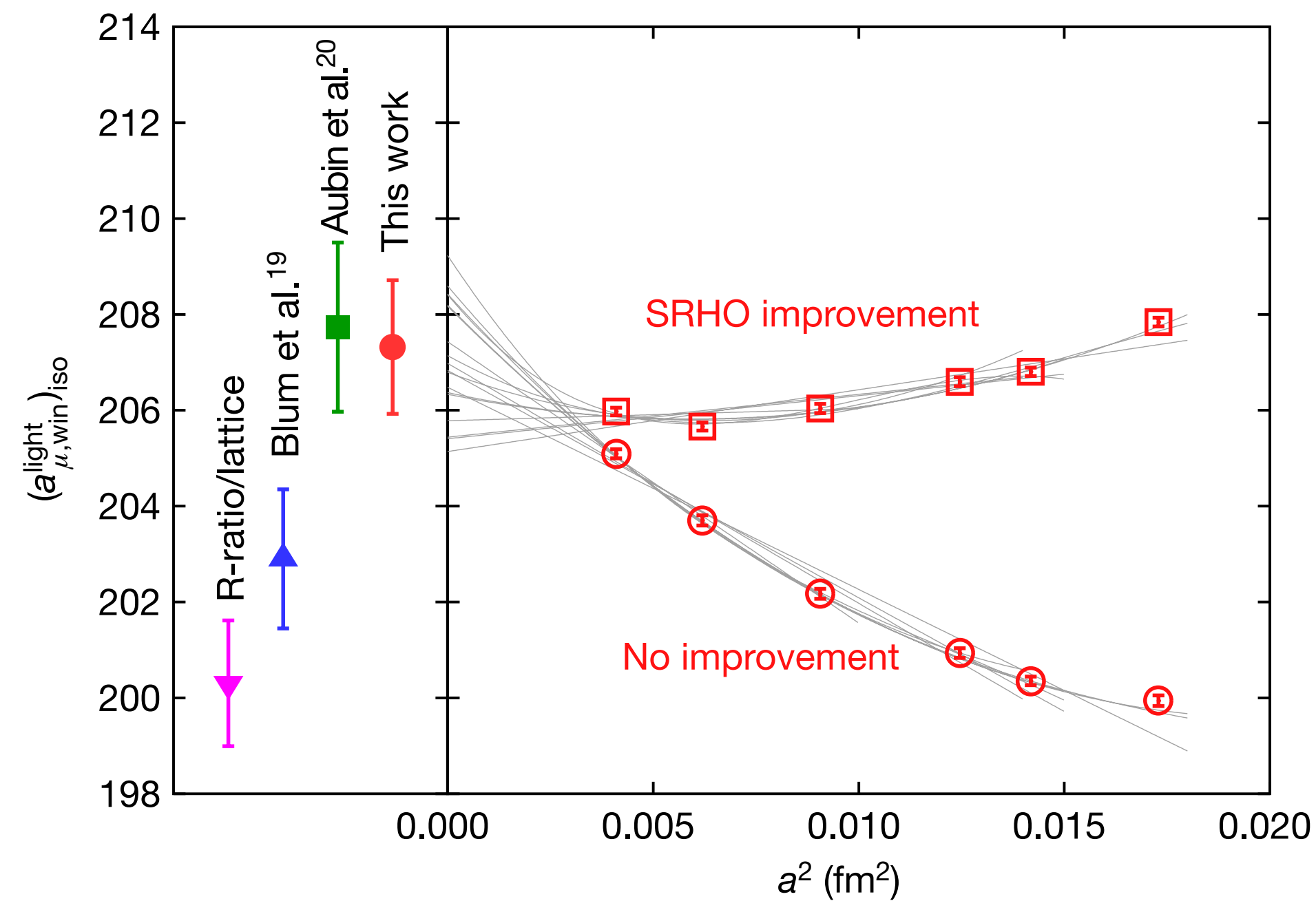
Data-driven approach:  $a_{\mu}^{\text{win}} = (229.4 \pm 1.4) \cdot 10^{-10}$  [Colangelo et al., Phys Lett B833 (2022) 137313]

(Excluding the 2023 CMD-3 result for  $e^+e^- \rightarrow \pi^+\pi^-$ )



# Intermediate window observable in Lattice QCD

BMWc: Rooted staggered quarks



$$a_{\mu}^{\text{win,ud}} = (207.3 \pm 0.4 \pm 1.3) \cdot 10^{-10}$$

[Borsányi et al., Nature 593 (2021) 7857]

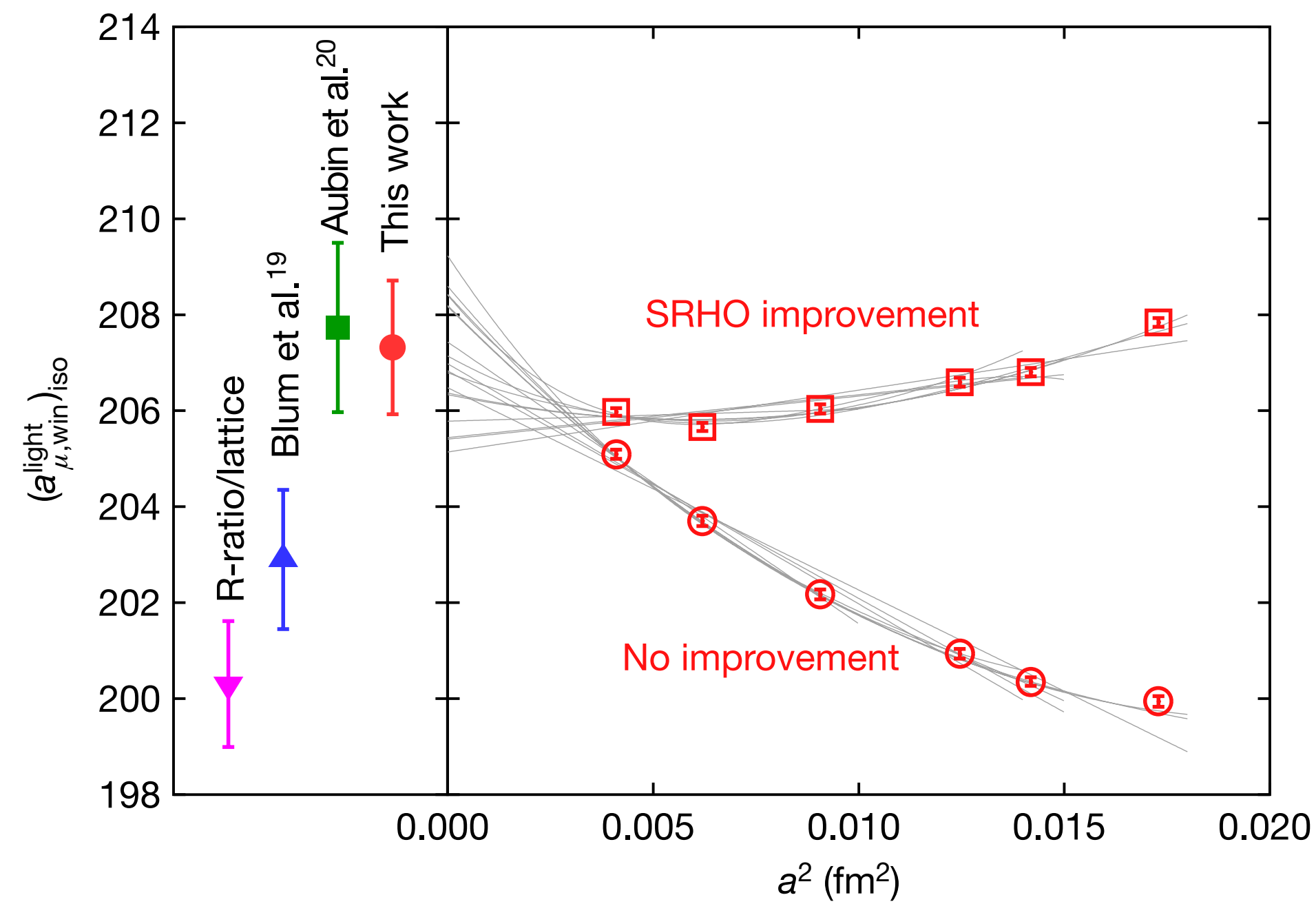
Mainz/CLS:  $O(a)$  improved Wilson quarks

- Extension to six lattice spacings:  
 $a = 0.099 - 0.035 \text{ fm}$
- Pion masses  $m_{\pi} = 130 - 420 \text{ MeV}$
- Two discretisations of the vector current:  
local and conserved
- Simultaneous chiral and continuum extrapolation
- Isospin-breaking correction included

[Cè et al., Phys Rev D106 (2022) 114502]

# Intermediate window observable in Lattice QCD

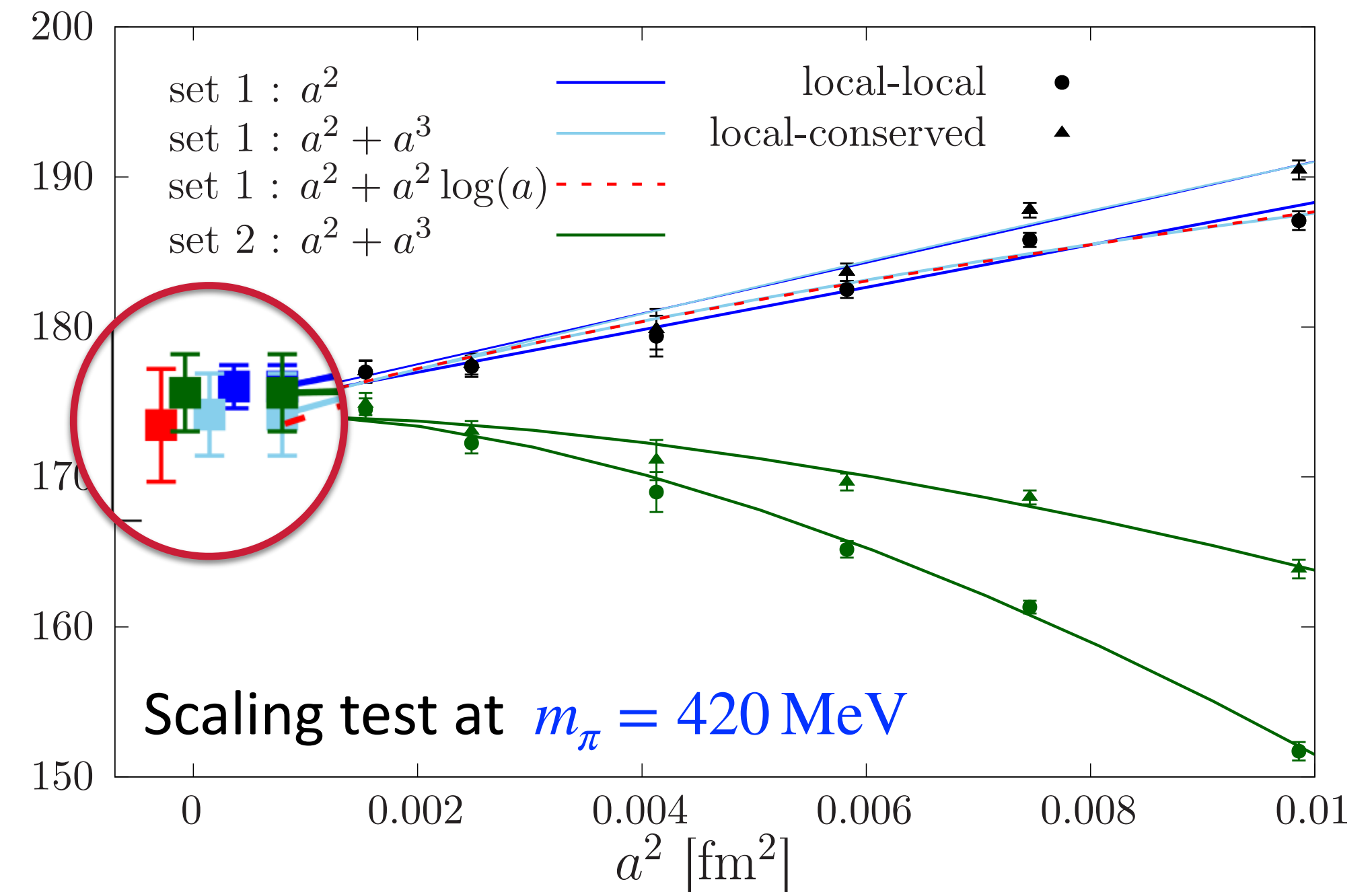
## BMWc: Rooted staggered quarks



$$a_{\mu}^{\text{win,ud}} = (207.3 \pm 0.4 \pm 1.3) \cdot 10^{-10}$$

[Borsányi et al., Nature 593 (2021) 7857]

## Mainz/CLS: $O(a)$ improved Wilson quarks

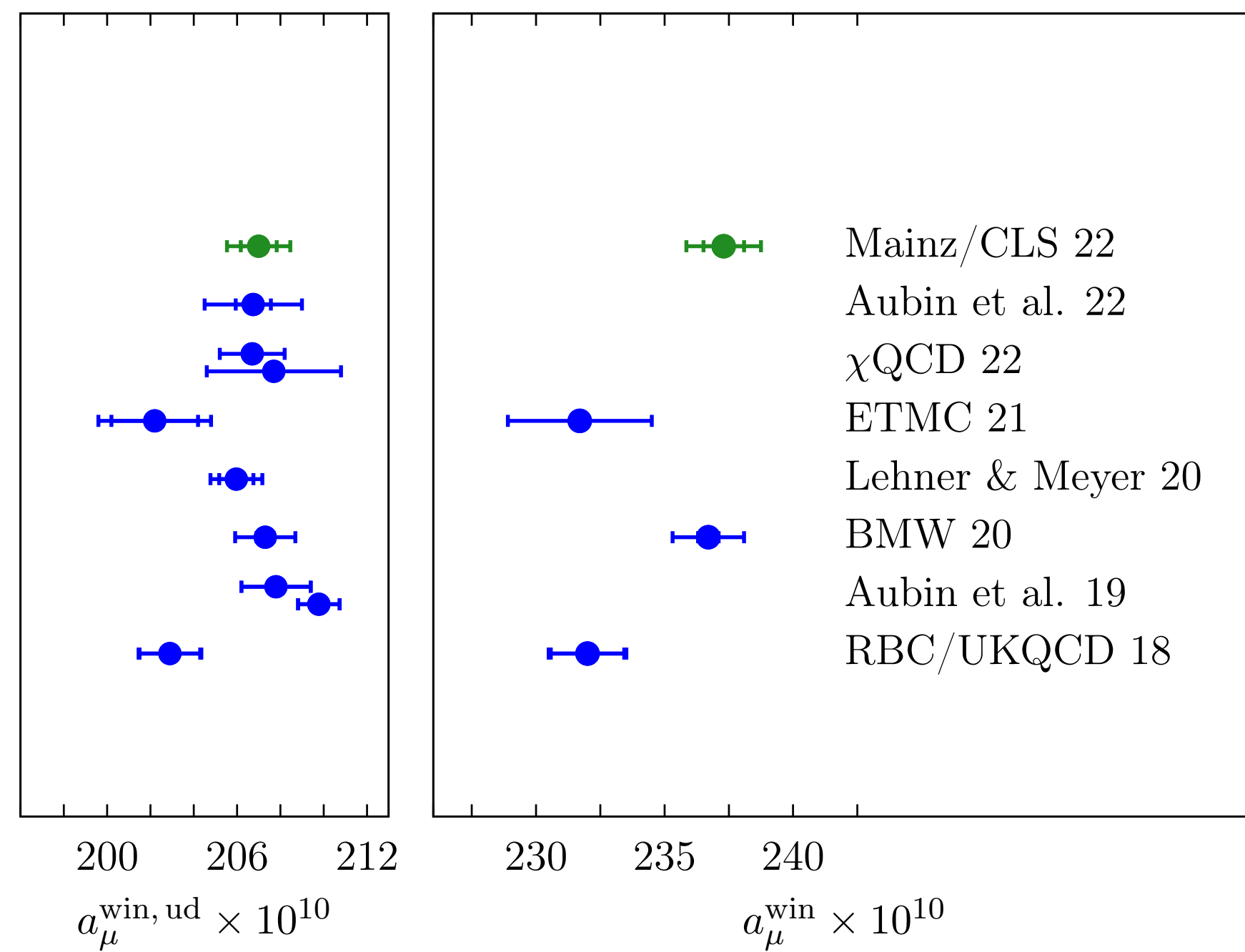


$$a_{\mu}^{\text{win,ud}} = (207.0 \pm 0.8 \pm 1.2) \cdot 10^{-10}$$

[Cè et al., Phys Rev D106 (2022) 114502]

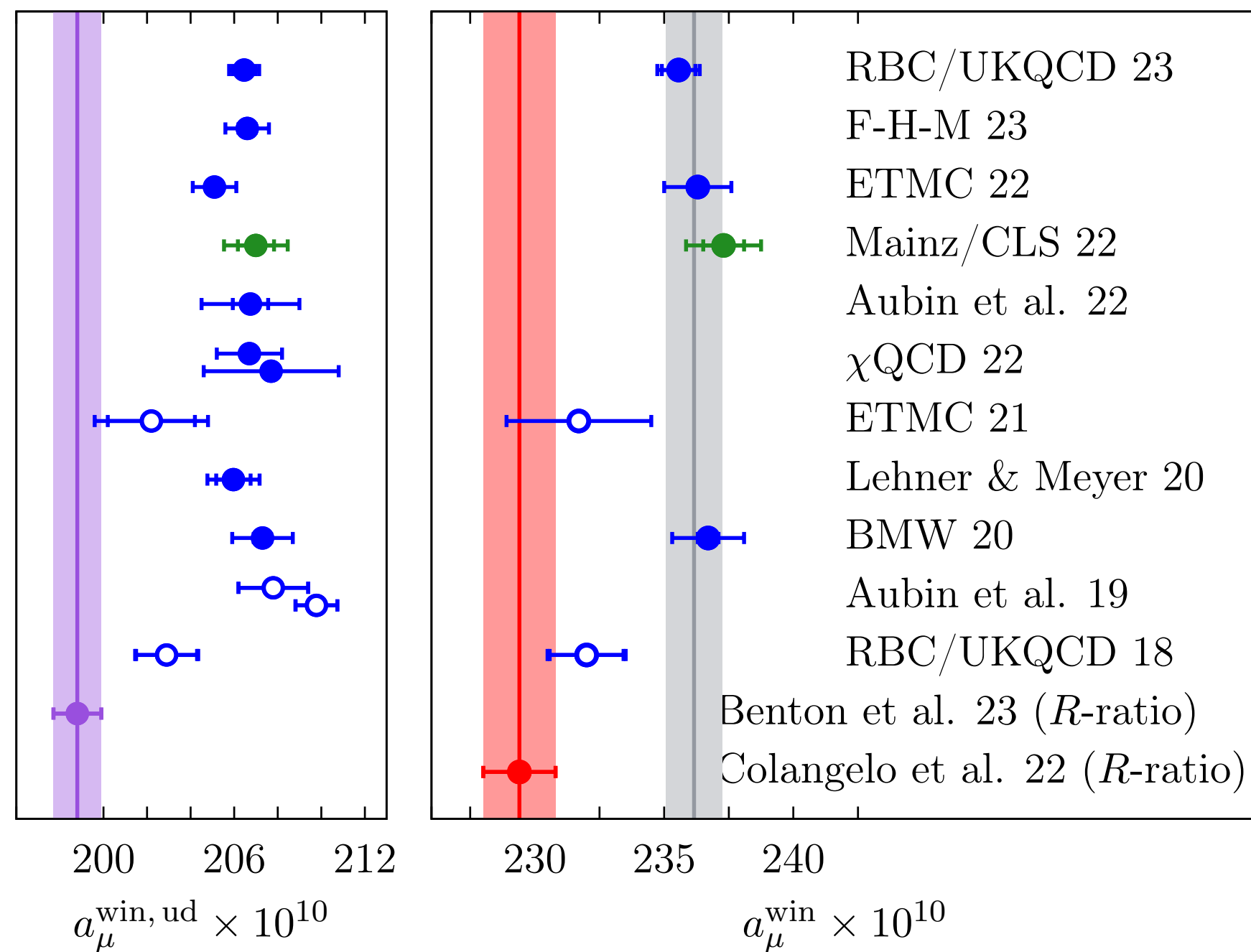


# Window observable: Lattice QCD vs. $R$ -ratio



Left: dominant light-quark contribution to  $a_\mu^{\text{win}}$   
Right: including sub-leading contributions

# Window observable: Lattice QCD vs. $R$ -ratio



- Dominant light-quark contribution confirmed for wide range of discretisation with sub-percent precision
- Significant tension with results based on the  $R$ -ratio\*

**$R$ -ratio estimate:**  $a_\mu^{\text{win}} = (229.4 \pm 1.4) \cdot 10^{-10}$

**Lattice average:**  $a_\mu^{\text{win}} = (236.16 \pm 1.09) \cdot 10^{-10}$

(RBC/UKQCD 23, ETMC 22, Mainz/CLS 22, BMW 20)

[HW, arXiv:2306.04165]

- Tension of  $3.8\sigma$  in the window observable evaluated from  $e^+e^-$  data\* and four lattice calculations

$$a_\mu^{\text{win}}|_{\langle \text{lat} \rangle} - a_\mu^{\text{win}}|_{e^+e^-} = (6.8 \pm 1.8) \cdot 10^{-10} \quad [3.8 \sigma]$$

- Subtract  $R$ -ratio result  $a_\mu^{\text{win}}|_{e^+e^-}$  from WP estimate and replace by lattice average  $a_\mu^{\text{win}}|_{\langle \text{lat} \rangle}$ :

$$a_\mu^{\text{exp}} - a_\mu^{\text{SM}}|_{e^+e^- \rightarrow \langle \text{lat} \rangle}^{\text{win}} = (18.1 \pm 4.8) \cdot 10^{-10} \quad [3.8 \sigma]$$

\*excluding the CMD-3 result

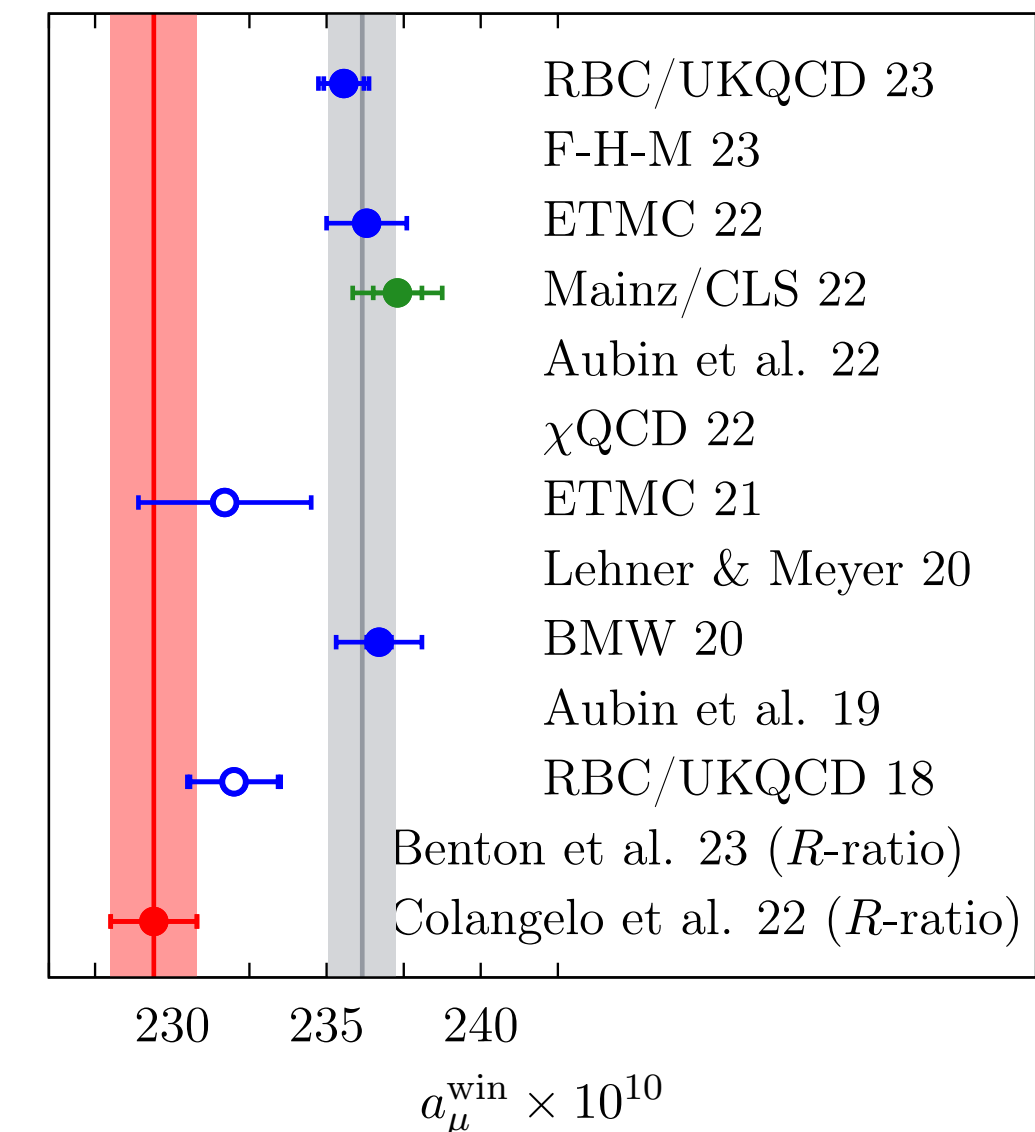


# What can we learn from $a_\mu^{\text{win}}$ ?

Primary observable in lattice calculations: vector correlator  $G(t)$

$$G(t) \equiv -\frac{a^3}{3} \sum_k \sum_{\vec{x}} \langle j_k^{\text{em}}(\vec{x}, t) j_k^{\text{em}}(0) \rangle = \frac{1}{12\pi^2} \int_{m_{\pi^0}^2}^{\infty} d(\sqrt{s}) R(s)^{\text{lat}} s e^{-\sqrt{s}t}$$

$a_\mu^{\text{win}}|_{\text{lat}} > a_\mu^{\text{win}}|_{e^+e^-}$  implies that  $R(s)^{\text{lat}} > R(s)^{e^+e^-}$  in some interval of  $\sqrt{s}$



Energy interval  $600 \leq \sqrt{s} \leq 900 \text{ MeV}$  contributes the same fraction to  $a_\mu^{\text{hvp}}$  and  $a_\mu^{\text{win}}$

$\sqrt{s}$ interval	$a_\mu^{\text{hvp}}$	$(a_\mu^{\text{hvp}})^{\text{SD}}$	$(a_\mu^{\text{hvp}})^{\text{ID}}$	$(a_\mu^{\text{hvp}})^{\text{LD}}$	$\bar{\Pi}(1 \text{ GeV}^2)$
Below 0.6 GeV	15.5	1.5	5.5	23.5	8.2
0.6 to 0.9 GeV	58.3	23.1	54.9	65.4	52.6
Above 0.9 GeV	26.2	75.4	39.6	11.1	39.2
Total	100.0	100.0	100.0	100.0	100.0

[Cè et al., Phys Rev D106 (2022) 114502]

# What can we learn from $a_\mu^{\text{win}}$ ?

- Phenomenological model for  $R$ -ratio predicts *[Mainz/CLS, Cè et al., Phys Rev D 106 (2022) 114502]*

$$\sqrt{s} = 600 - 900 \text{ MeV: } \frac{R(s)^{\text{lat}}}{R(s)^{e^+e^-}} = 1 + \epsilon \Rightarrow \frac{(a_\mu^{\text{hvp}})^{\text{lat}}}{(a_\mu^{\text{hvp}})^{e^+e^-}} \approx \frac{(a_\mu^{\text{win}})^{\text{lat}}}{(a_\mu^{\text{win}})^{e^+e^-}} = 1 + 0.6\epsilon$$

- Lattice average vs.  $R$ -ratio:  $(a_\mu^{\text{win}})^{\text{lat}} / (a_\mu^{\text{win}})^{e^+e^-} = 1.030(8)$   
 $\Rightarrow R(s)^{\text{lat}}$  is enhanced by 5% relative to  $R(s)^{e^+e^-}$  for  $\sqrt{s} = 600 - 900 \text{ MeV}$
- If confirmed, it would imply that BMW's estimate might be too low....

## Similar conclusions

- Dispersive treatment of pion form factor *[Colangelo, Hoferichter, Stoffer, PLB 814 (2021) 136073]*
- “Energy-smeared”  $R$ -ratio from lattice data *[ETMC, Alexandrou et al., PRL 130 (2023) 241901]*



# More windows....

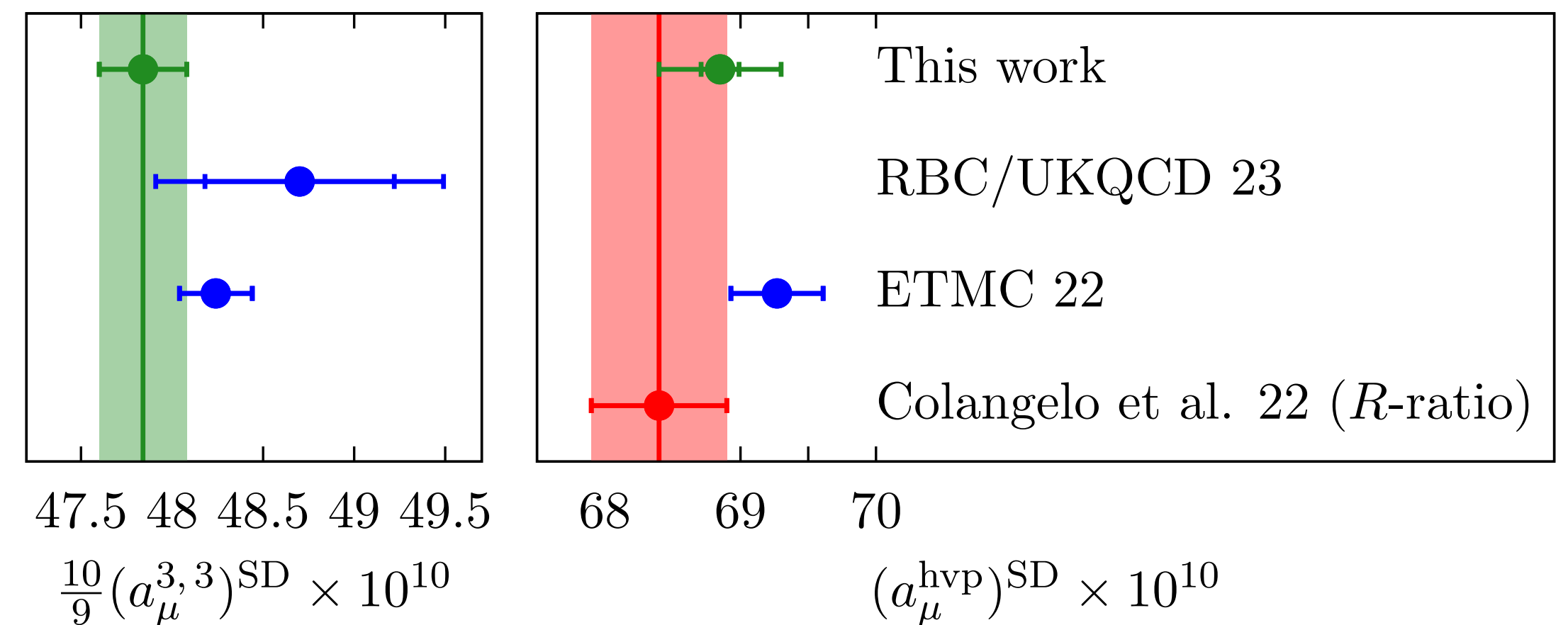
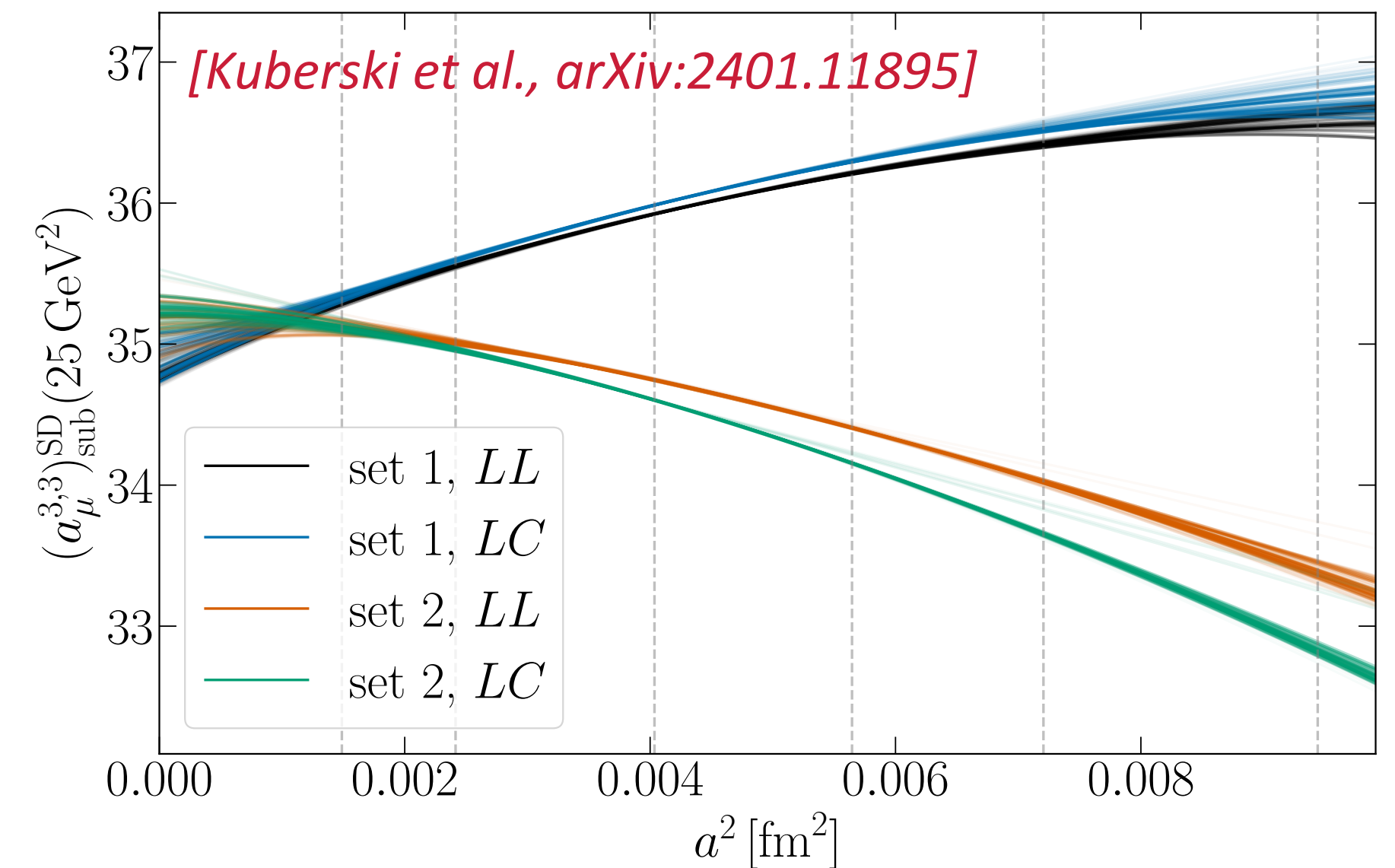
## Short-distance window:

- Finite-volume correction negligible
- Uncertainty dominated by control over lattice artefacts

$$(a_\mu^{\text{win}})^{\text{SD}} = (68.85 \pm 0.15 \pm 0.42) \times 10^{-10}$$

## Hadronic model:

- 5% enhancement of  $R(s)^{\text{lat}}$  for  $0.6 \text{ GeV} \leq \sqrt{s} \leq 0.9 \text{ GeV}$  increases  $(a_\mu^{\text{win}})^{\text{SD}}$  by  $+1 \times 10^{-10}$
- Expectation confirmed by lattice calculations



# Hadronic running of electromagnetic coupling

Electromagnetic coupling is energy-dependent:

$$\alpha^{-1} = 137.035\,999\dots \quad \alpha(q^2) = \frac{\alpha}{1 - \Delta\alpha(q^2)} \quad \alpha^{-1}(M_Z^2) = 127.951 \pm 0.009$$

Correlation between  $a_\mu^{\text{hvp}}$  and the hadronic running of  $\Delta\alpha_{\text{had}}$ :

$$\Delta\alpha_{\text{had}}(q^2) = -\frac{\alpha q^2}{3\pi} \oint_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{s(s - q^2)}, \quad a_\mu^{\text{hvp}} = \left(\frac{\alpha m_\mu}{3\pi}\right)^2 \int_{m_{\pi^0}^2}^{\infty} ds \frac{R(s) \hat{K}(s)}{s^2}$$

Euclidean momenta

$\Delta\alpha_{\text{had}}(-Q^2)$  accessible in lattice QCD via the same correlator  $G(t)$  with a different kernel function:

$$\Delta\alpha_{\text{had}}(-Q^2) = \frac{\alpha}{\pi} \frac{1}{Q^2} \int_0^\infty dt G(t) \left[ Q^2 t^2 - 4 \sin^2 \left( \frac{1}{2} Q^2 t^2 \right) \right]$$

Hadronic running at  $Z$ -pole:  $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) \rightarrow$  key quantity in global electroweak fit



# Evaluation of $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$ and comparison with EW precision data

Adler function approach, aka. “Euclidean split technique”

$$\begin{aligned}\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) &= \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) \quad \leftarrow \text{lattice QCD} \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2)] \quad \leftarrow \text{perturbative Adler function} \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-M_Z^2)] \quad \leftarrow \text{pQCD}\end{aligned}$$

# Evaluation of $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$ and comparison with EW precision data

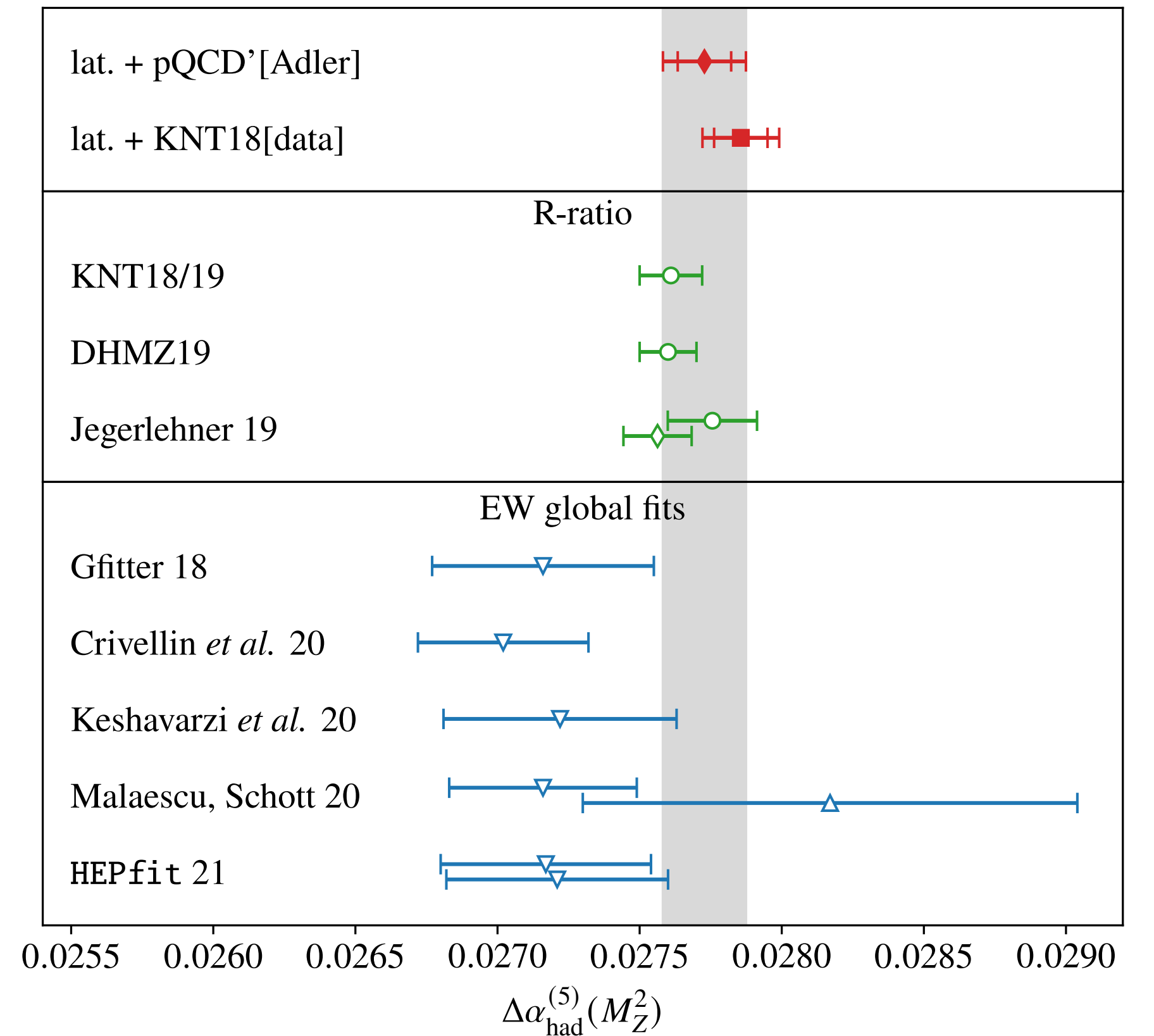
Adler function approach, aka. “Euclidean split technique”

$$\begin{aligned} \Delta\alpha_{\text{had}}^{(5)}(M_Z^2) &= \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2)] \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-M_Z^2)] \end{aligned}$$

$$\Rightarrow \Delta\alpha_{\text{had}}^{(5)}(M_Z^2) = 0.027\,73(9)_{\text{lat}}(2)_{\text{btm}}(12)_{\text{pQCD}}$$

[Mainz/CLS, Cè *et al.*, *JHEP* 08 (2022) 220, *arXiv:2203.08676*]

- No inconsistency with global electroweak fit!



Standard Model can accommodate a larger value for  $a_\mu$  without contradicting electroweak precision data



## Summary and outlook

Hadron physics holds the key in the quest for new physics via  $(g - 2)_\mu$

No straightforward interpretation of the Fermilab E989 experiment

Discrepant determinations of the HVP contribution:

- Tensions between lattice QCD and  $e^+e^-$  hadronic cross sections\*
- Tension in  $\pi^+\pi^-$  channel between BaBar vs. KLOE and CMD-3 vs. all other results

Analyses / re-analyses of  $e^+e^-$  data in progress: BaBar, BESIII, CMD-3, KLOE

Lattice QCD to produce more results for HVP contribution with sub-percent precision

Experimental measurement of the HVP contribution by MUonE experiment

Fermilab E989 to analyse data from Runs 4–6

Update of White Paper expected by  $\approx$  Dec 2024

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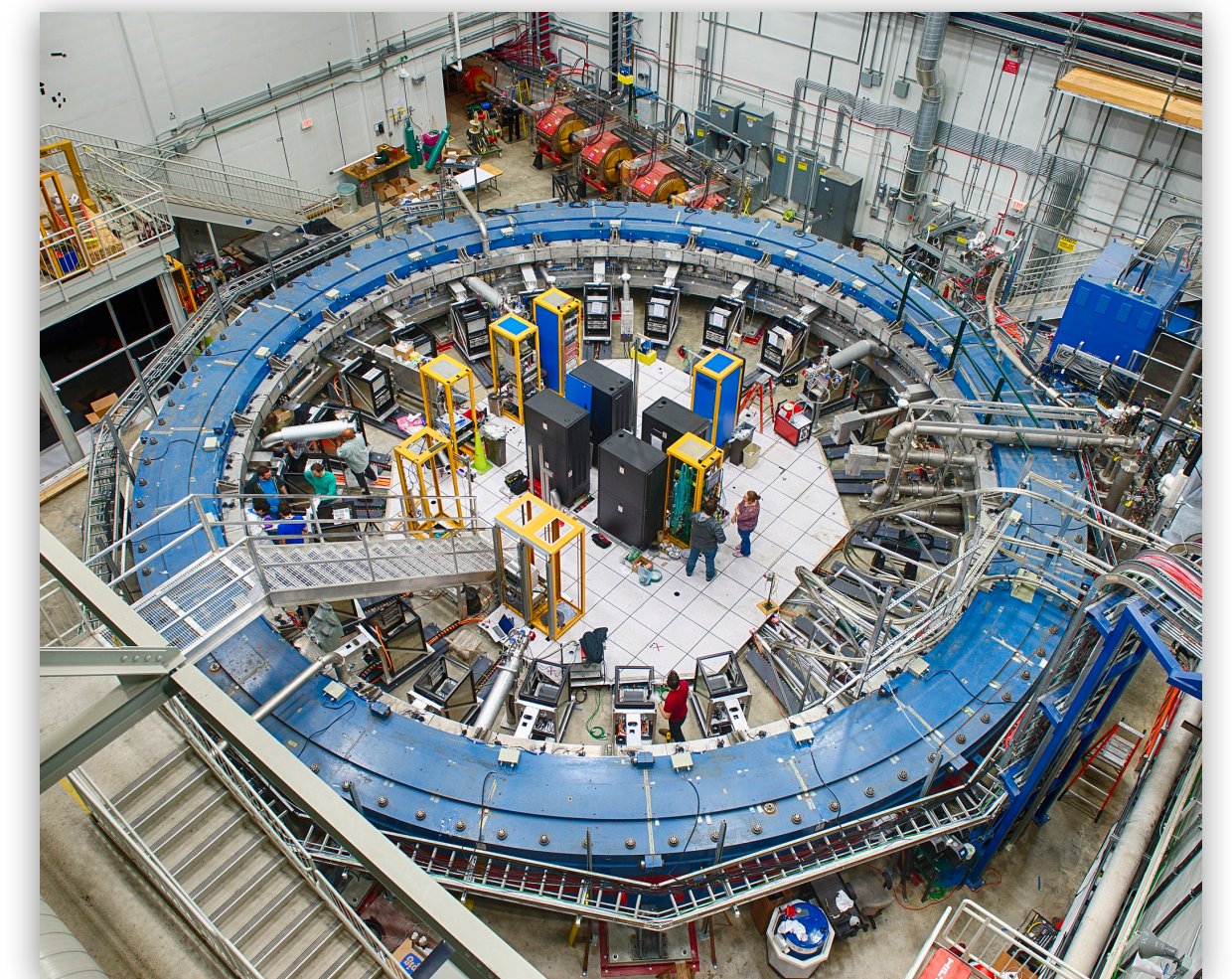
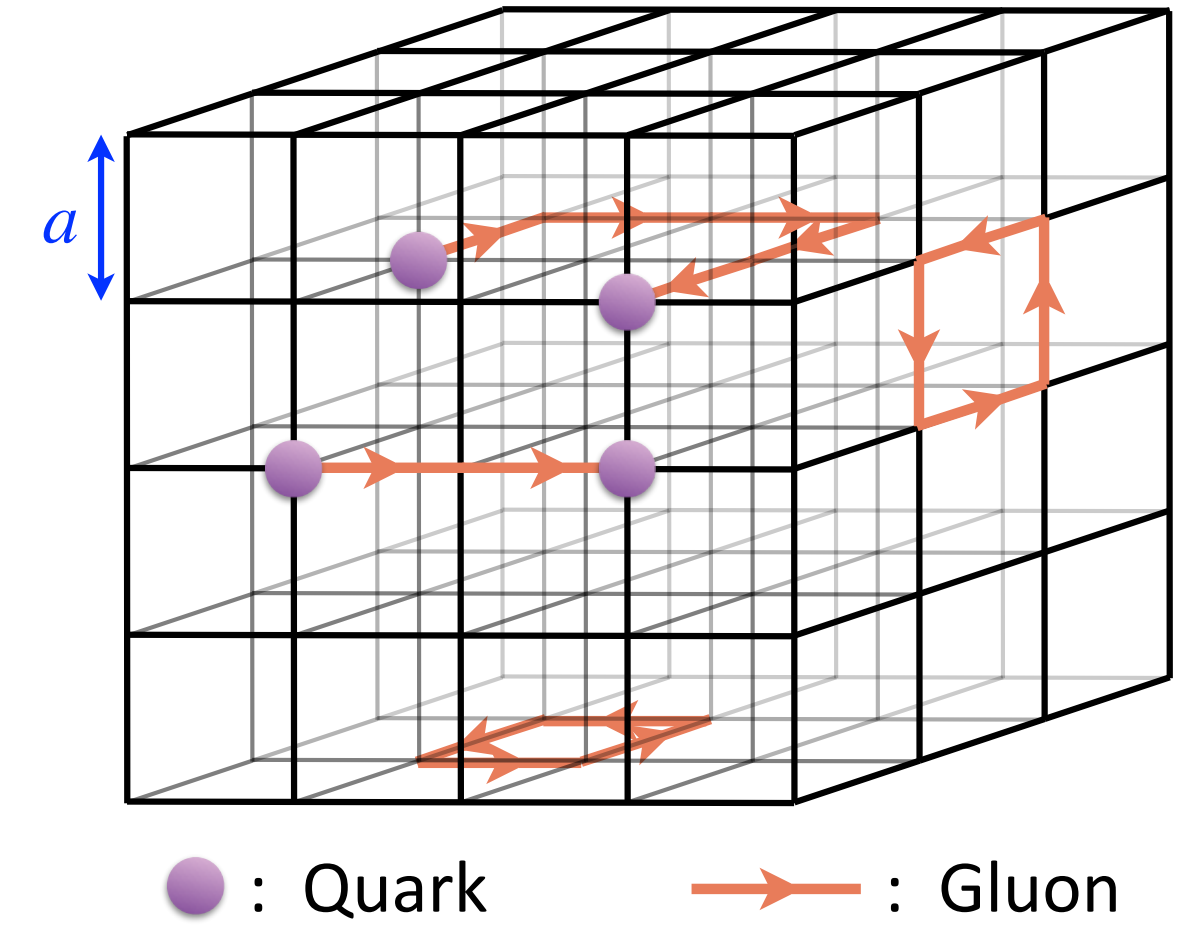
\*pre-2023





It's **not** magnitude that matters  
but significance!

**Thank you!**





# Backup

# QED contributions to $a_\mu$

QED contribution has been worked out to in perturbation theory to 5-loop order:

PRL **109**, 111808 (2012)

PHYSICAL REVIEW LETTERS

week ending  
14 SEPTEMBER 2012

## Complete Tenth-Order QED Contribution to the Muon $g - 2$

Tatsumi Aoyama,<sup>1,2</sup> Masashi Hayakawa,<sup>3,2</sup> Toichiro Kinoshita,<sup>4,2</sup> and Makiko Nio<sup>2</sup>

<sup>1</sup>*Kobayashi-Maskawa Institute for the Origin of Particles and the Universe (KMI), Nagoya University, Nagoya, 464-8602, Japan*

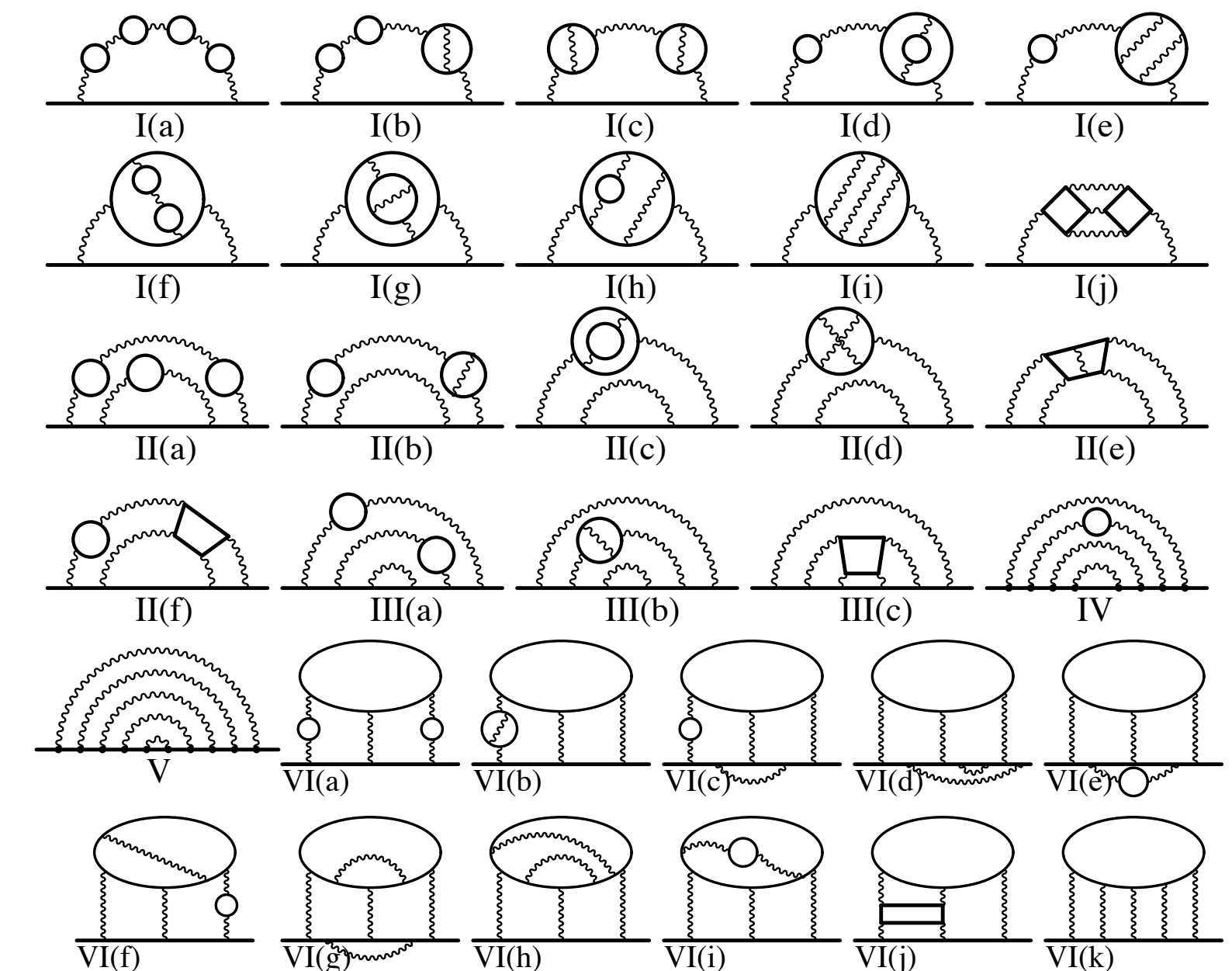
<sup>2</sup>*Nishina Center, RIKEN, Wako, Japan 351-0198*

<sup>3</sup>*Department of Physics, Nagoya University, Nagoya, Japan 464-8602*

<sup>4</sup>*Laboratory for Elementary Particle Physics, Cornell University, Ithaca, New York 14853, USA*

(Received 24 May 2012; published 13 September 2012)

SM	116	591	810	100	%	#diagrams
QED(tot)	116	584	718.931	99,9939	%	
2	116	140	973.321	99,6133	%	1
4		413	217.626	0,3544	%	9
6		30	141.902	0,0259	%	72
8			381.004	0,0003	%	891
10			5.078	$4 \cdot 10^{-6}$	%	12672

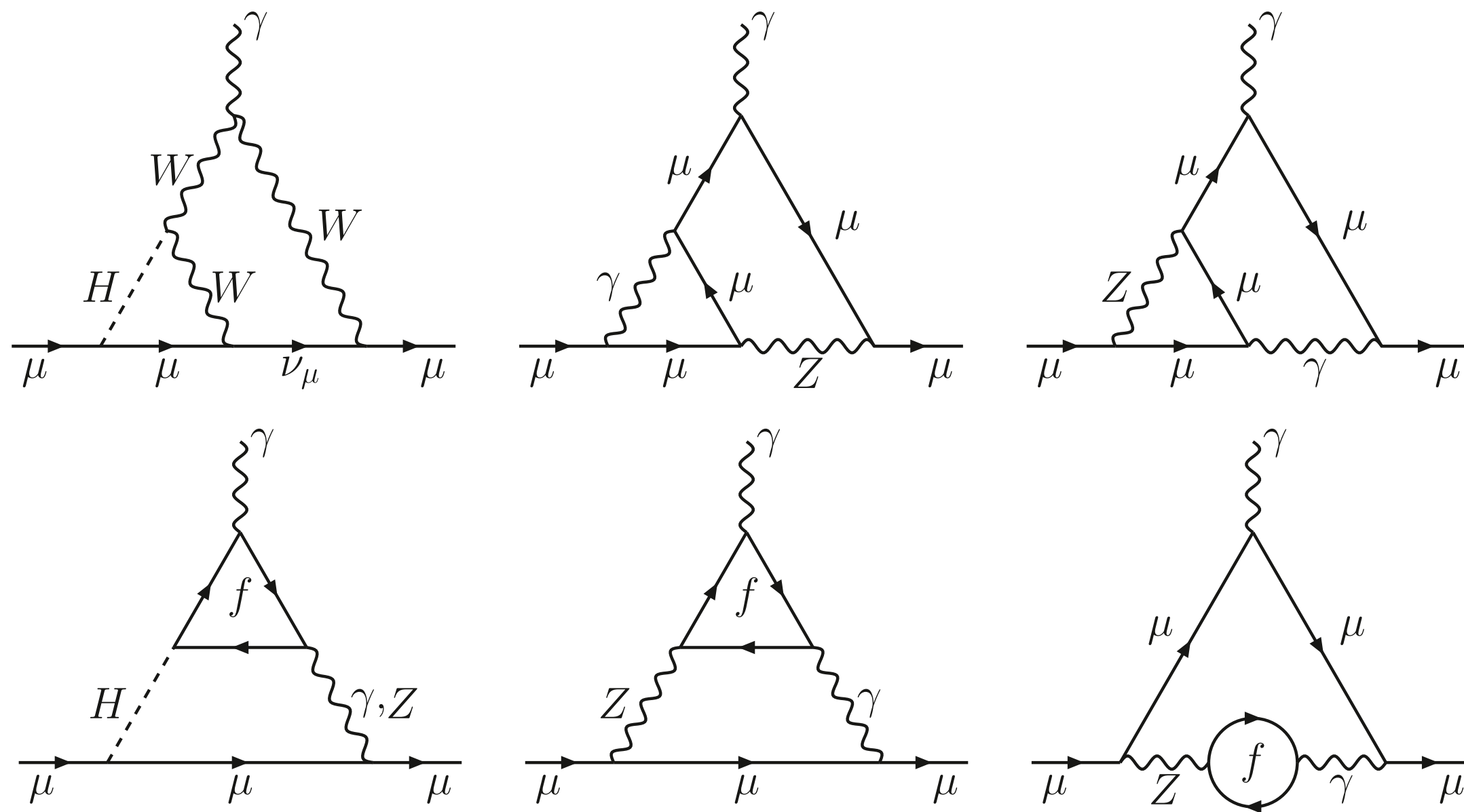




# Electroweak contributions to $a_\mu$

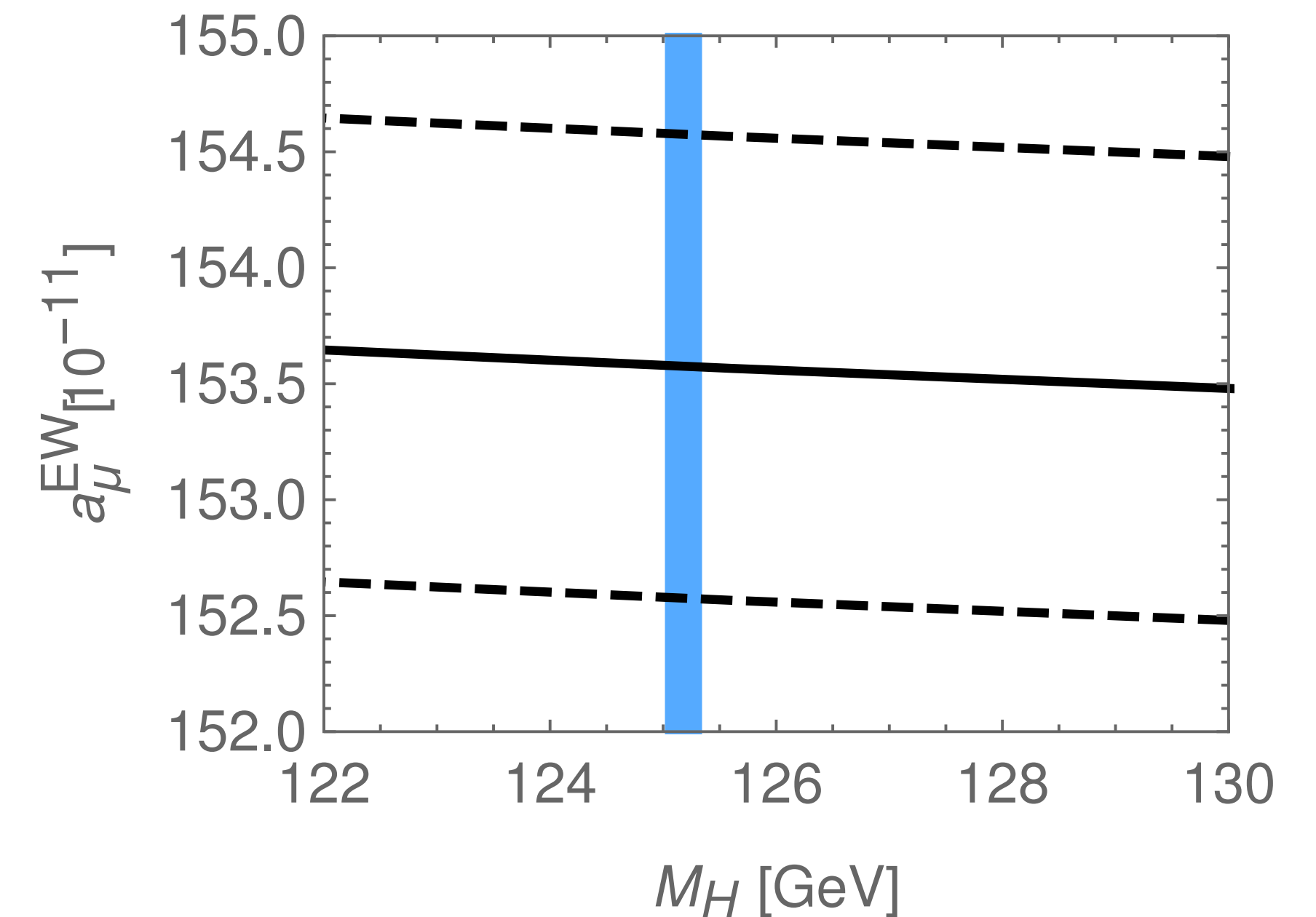
Weak contributions known to leading three-loop order

Sample two-loop diagrams:



$$a_\mu^{\text{weak}} = (153.6 \pm 1.0) \times 10^{-11}$$

Dependence on the Higgs mass



[Gnendiger et al., arXiv:1306.5546]

[Aoyama et al., Phys. Rep. 887 (2020) 1, arXiv:2006.04822]

# Hadronic cross section data and their analysis

Cross section data collected either via energy scan (VEPP-2000, ...) or using the ISR technique (BaBar, CLEO, KLOE, BESIII,...)

Determine the ISR luminosity:

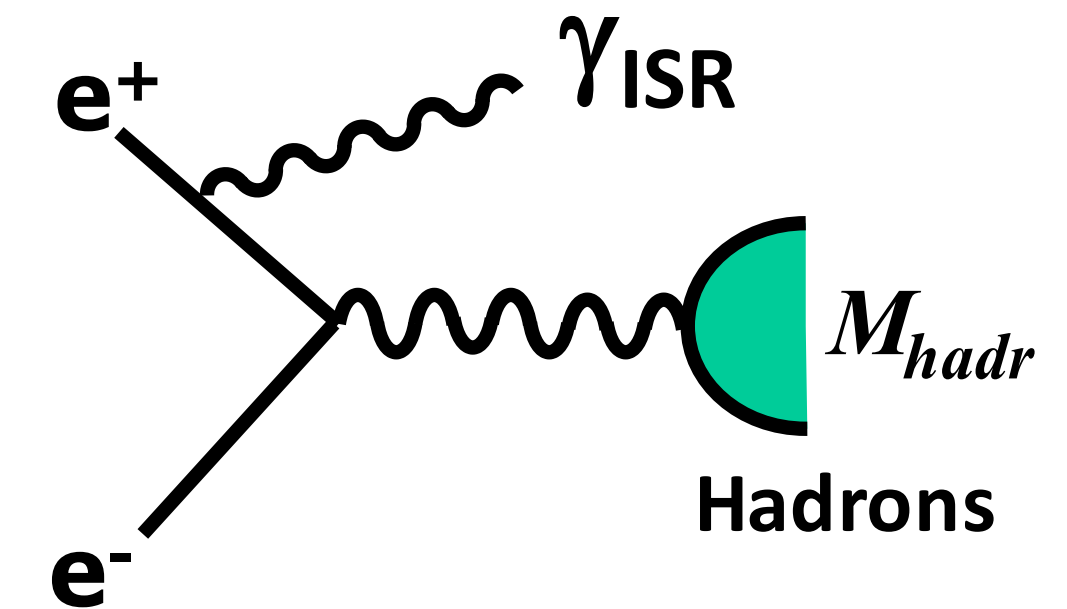
$$\frac{d\sigma_{\text{ISR}}(\sqrt{s'})}{d\sqrt{s'}} = \frac{2\sqrt{s'}}{s} W(s, E_\gamma, \theta_\gamma) \sigma(\sqrt{s'}) \quad W : \text{ radiator function}$$

Monte Carlo event generators for  $e^+e^- \rightarrow \text{hadrons}(\gamma)$ , e.g. BABAYAGA, PHOKHARA

Several variants of ISR analyses:

- Tagged vs. untagged ISR photon
- Measured radiator function vs. event generator

Other issues: particle ID and  $\pi$ - $\mu$  separation





# Hadronic vacuum polarisation: Data-driven approach

Decade-long effort to measure  $e^+e^-$  cross sections

$$a_\mu^{\text{hvp, LO}} = \left(\frac{\alpha m_\mu}{3\pi}\right)^2 \int_{m_{\pi^0}^2}^{\infty} ds \frac{R_{\text{had}}(s) \hat{K}(s)}{s^2}$$

$\sqrt{s} \lesssim 2 \text{ GeV}$ : sum of exclusive channels

$\sqrt{s} > 2 \text{ GeV}$ : inclusive channels, narrow resonances, perturbative QCD

	DHMZ19	KNT19	Difference
$\pi^+\pi^-$	507.85(0.83)(3.23)(0.55)	504.23(1.90)	3.62
$\pi^+\pi^-\pi^0$	46.21(0.40)(1.10)(0.86)	46.63(94)	-0.42
$\pi^+\pi^-\pi^+\pi^-$	13.68(0.03)(0.27)(0.14)	13.99(19)	-0.31
$\pi^+\pi^-\pi^0\pi^0$	18.03(0.06)(0.48)(0.26)	18.15(74)	-0.12
$K^+K^-$	23.08(0.20)(0.33)(0.21)	23.00(22)	0.08
$K_S K_L$	12.82(0.06)(0.18)(0.15)	13.04(19)	-0.22
$\pi^0\gamma$	4.41(0.06)(0.04)(0.07)	4.58(10)	-0.17
Sum of the above	626.08(0.95)(3.48)(1.47)	623.62(2.27)	2.46
[1.8, 3.7] GeV (without $c\bar{c}$ )	33.45(71)	34.45(56)	-1.00
$J/\psi, \psi(2S)$	7.76(12)	7.84(19)	-0.08
[3.7, $\infty$ ) GeV	17.15(31)	16.95(19)	0.20
Total $a_\mu^{\text{HVP, LO}}$	694.0(1.0)(3.5)(1.6)(0.1) $_{\psi}$ (0.7) $_{\text{DV+QCD}}$	692.8(2.4)	1.2

[Aoyama et al., Phys. Rep. 887 (2020) 1, arXiv:2006.04822]

# White Paper estimate of the HVP contribution

## Merging procedure

- Determine  $a_\mu^{\text{hvp, LO}}$  as the sum of simple averages of individual hadronic channels (DHMZ, KNT, CHKS below 1 GeV), including correlations in the data
- Experimental and theoretical uncertainties:  
use maximum error estimate in each channel (except  $\pi^+\pi^-$ ) of either DHMZ or KNT
- Extra systematic uncertainty of  $\Delta \equiv \frac{1}{2} |\text{DHMZ} - \text{KNT}|$  in each channel, except  $\pi^+\pi^-$ ;  
For  $\pi^+\pi^-$  : use maximum of  $\Delta$  and tension between BaBar and KLOE
- Add uncertainty associated with duality violations and perturbative QCD

$$a_\mu^{\text{hvp, LO}} = 693.1(2.8)_{\text{exp}}(2.8)_{\text{syst}}(0.7)_{\text{DV+QCD}} \times 10^{-10} = 693.1(4.0) \times 10^{-10} \quad [0.6\%]$$