The puzzles surrounding the muon anomalous magnetic moment

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The Quest for New Physics

No evidence for Beyond-Standard Model particles from collider experiments

Overwhelming evidence for dark sector from astrophysical observations

Standard Model does not provide a complete description of Nature

Energy Frontier

New particles and interactions at colliders

> Enhancement of rare phenomena

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Precision Frontier

Comparison of precision observables to SM predictions







Lepton anomalous magnetic moments as probes for New Physics

Magnetic moment of particle with spin \vec{S} and charge e:



g = 2

Quantum corrections modify Dirac's prediction g = 2

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g = 2(1 + a), a: anomalous magnetic moment



Lepton anomalous magnetic moments as probes for New Physics

Magnetic moment of particle with spin \vec{S} and charge e:



$$g = 2\left(1 + \frac{\alpha}{2\pi}\right)$$

Quantum corrections modify Dirac's prediction g = 2

Electromagnetic, weak and strong interactions contribute to *a*

Beyond leading order: distinct values of a_e , a_μ and a_τ



g = 2(1 + a), a: anomalous magnetic moment



Muon g - 2 Theory Initiative

Founded in 2017 Agree on common SM prediction Focus on hadronic contributions White Paper published in 2020 Update foreseen in early 2025



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Hadronic vacuum polarisation (HVP)

[Aoyama et al., Phys. Rep. 887 (2020) 1, arXiv:2006.04822]

$\log 2$	[2020 White F	'ape
$34718.9(1) \times 10^{-11}$	0.001 ppm	
$153.6(1.0) \times 10^{-11}$	0.01 ppm	
$6845(40) \times 10^{-11}$	0.34 ppm [0	.6%]
$92(18) \times 10^{-11}$	0.15 ppm [2	0%]
$1810(43) \times 10^{-11}$	0.37 ppm	



Hadronic light-by-light scattering (HLbL)





Standard Model prediction for muQED:116 58Weak:116 58Hadronic vacuum polarisation:116 58Hadronic light-by-light scattering:116 59
$$a_{\mu}^{SM} = a_{\mu}^{QED} + a_{\mu}^{weak} + a_{\mu}^{hvp} + a_{\mu}^{hlbl} = 116 59$$

- QED and electroweak contributions account for 99.994% of the SM prediction for a_{μ}
- Error is dominated by strong interaction effe

[2020 Wh	ite Pape
0.001 ppm	
0.01 ppm	
0.34 ppm	[0.6%]
0.15 ppm	[20%]
0.37 ppm	
QED+	-EW
HLbL	HVP
	[2020 Whi 0.001 ppm 0.01 ppm 0.34 ppm 0.15 ppm 0.37 ppm QED+





Standard Model prediction for muon
$$g - 2$$

QED: 116 584 718.9 (1) × 10⁻¹¹
Weak: 153.6(1.0) × 10⁻¹¹
Hadronic vacuum polarisation: 6845(40) × 10⁻¹¹
Hadronic light-by-light scattering: 92(18) × 10⁻¹¹
 $a_{\mu}^{SM} = a_{\mu}^{QED} + a_{\mu}^{weak} + a_{\mu}^{hvp} + a_{\mu}^{hlbl} = 116 591 810(43) × 10^{-11}$
Standard Model vs. experiment: $a_{\mu}^{exp} \stackrel{?}{=} a_{\mu}^{QED} + a_{\mu}^{weak} + a_{\mu}^{hvp} + a_{\mu}^{hlbl}$

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Standard Model vs. experiment: $a_{\mu}^{exp} \stackrel{?}{=} a_{\mu}^{QED} + a_{\mu}^{weak} + a_{\mu}^{exp}$
Why the muon?
 $a_{\ell}^{BSM} \propto m_{\ell}^2/M_{BSM}^2$ $\ell = e, \mu, \tau$
 \rightarrow sensitivity of a_{μ} enhanced by $(m_{\mu}/m_e)^2 \approx 4.3$

ction f	for muon $g-2$	[2020 Wh	ite Paper
	$116584718.9(1) \times 10^{-11}$	0.001 ppm	
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nt: a_{μ}^{ex}	$a_{\mu} \stackrel{?}{=} a_{\mu}^{\text{QED}} + a_{\mu}^{\text{weak}} + a_{\mu}^{\text{hvp}} + a_{\mu}^{\text{hlbl}} $	$+ a_{\mu}^{BSM}$	
$m_\ell^2/M_{\rm BSN}^2$	M $\ell = e, \mu, \tau$	μ	

 $(m_{\mu}/m_e)^2 \approx 4.3 \times 10^4$ relative to a_e















New Physics on the horizon?

Confronting the SM prediction with the E989 measurement

 $a_u^{\text{exp}} = 116\,592\,049(22) \times 10^{-11}$ [0.19 ppm]

 $a_{\mu}^{\rm SM} = 116\,591\,810(43) \times 10^{-11}$ [0.37 ppm]

 $\Rightarrow a_{\mu}^{\exp} - a_{\mu}^{SM} = (249 \pm 48) \cdot 10^{-11} \quad [5.1\sigma]$

Standard Model prediction:

- White paper estimate based on "data-driven" evaluation of HVP contribution: dispersion integrals and hadronic cross sections
- Lattice QCD result for HVP with comparable precision [Borsányi et al., Nature 593 (2021) 7857]

 $a_{\mu}^{\exp} - a_{\mu}^{SM} \Big|_{BMWc}^{hvp, LO} = (105 \pm 61) \cdot 10^{-11} \quad [1.7 \sigma]$







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Requires independent confirmation







Hadronic light-by-light scattering



 a_{μ}^{hlbl} : Uncontroversial — contributes 0.15 ppm to the total SM uncertainty of 0.37 ppm

 \rightarrow Focus on refinements and further reduction of uncertainty



Hadronic models

Hadronic models, data-driven method and Lattice QCD produce compatible results White paper recommended value: $a_{\mu}^{\text{hlbl}} = (92 \pm 18) \cdot 10^{-11}$ **Recent lattice calculations:**

 $a_{\mu}^{\text{hlbl, LO}} = \begin{cases} (109.6 \pm 14.7) \cdot 10^{-11} & \text{Mainz/CLS} \\ (124.7 \pm 15.2) \cdot 10^{-11} & \text{RBC} \end{cases}$

[Chao et al., EPJC 81 (2021) 651; EPJC 82 (2022) 664; Blum et al., arXiv:2304.04423]





Hadronic vacuum polarisation from dispersion theory

Analyticity, unitarity & optical theorem imply:

$$m = \int \frac{ds}{\pi(s-q^2)} \operatorname{Im} m$$

$$a_{\mu}^{\text{hvp,LO}} = \left(\frac{\alpha m_{\mu}}{3\pi}\right)^2 \int_{m_{\pi^0}^2}^{\infty} ds \, \frac{R_{\text{had}}(s)\hat{K}(s)}{s^2}, \quad R_{\text{h}}$$

Hadronic effects cannot be treated in perturbation theory

- Use experimental data for $R_{had}(s)$ in the low-energy regime ("data-driven approach") Standard Model prediction is subject to experimental uncertainties





Decade-long effort to measure e^+e^- cross sections

 $\sqrt{s} \leq 2 \,\text{GeV}$: sum of exclusive channels $\sqrt{s} > 2 \text{ GeV}$: inclusive channels, narrow resonances, perturbative QCD



 $a_{\mu}^{\text{hvp, LO}} = \left(\frac{\alpha m_{\mu}}{3\pi}\right)^2 \int_{m^2_{-}}^{\infty} ds \, \frac{R_{\text{had}}(s)\tilde{K}(s)}{s^2}$



























 $\sqrt{s} \leq 2 \,\text{GeV}$: sum of exclusive channels



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Decade-long effort to measure e^+e^- cross sections

 $\sqrt{s} \lesssim 2 \,\text{GeV}$: sum of exclusive channels

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	• •		
	DHMZ19	KNT19	Difference
$\pi^+\pi^-$	507.85(0.83)(3.23)(0.55)	504.23(1.90)	3.62
$\pi^+\pi^-\pi^0$	46.21(0.40)(1.10)(0.86)	46.63(94)	-0.42
$\pi^+\pi^-\pi^+\pi^-$	13.68(0.03)(0.27)(0.14)	13.99(19)	-0.31
$\pi^+\pi^-\pi^0\pi^0$	18.03(0.06)(0.48)(0.26)	18.15(74)	-0.12
K^+K^-	23.08(0.20)(0.33)(0.21)	23.00(22)	0.08
$K_S K_L$	12.82(0.06)(0.18)(0.15)	13.04(19)	-0.22
$\pi^0\gamma$	4.41(0.06)(0.04)(0.07)	4.58(10)	-0.17
Sum of the above	626.08(0.95)(3.48)(1.47)	623.62(2.27)	2.46
[1.8, 3.7] GeV (without $c\overline{c}$)	33.45(71)	34.45(56)	-1.00
$J/\psi, \psi(2S)$	7.76(12)	7.84(19)	-0.08
$[3.7,\infty)$ GeV	17.15(31)	16.95(19)	0.20
Total $a_{\mu}^{\text{HVP, LO}}$	$694.0(1.0)(3.5)(1.6)(0.1)_{\psi}(0.7)_{\rm DV+QCD}$	692.8(2.4)	1.2

$$a_{\mu}^{\text{hvp, LO}} = \left(\frac{\alpha m_{\mu}}{3\pi}\right)^2 \int_{m_{\pi^0}^2}^{\infty} ds \, \frac{R_{\text{had}}(s)\hat{K}(s)}{s^2}$$

[Aoyama et al., Phys. Rep. 887 (2020) 1, arXiv:2006.04822]





Decade-long effort to measure e^+e^- cross sections

 $\sqrt{s} \leq 2 \,\text{GeV}$: sum of exclusive channels $\sqrt{s} > 2 \text{ GeV}$: inclusive channels, narrow resonances, perturbative QCD

• White Paper recommended value (2020):

 $a_{\mu}^{\text{hvp, LO}} = 693.1(2.8)_{\text{exp}}(2.8)_{\text{syst}}(0.7)_{\text{DV+QCD}} \times 10^{-10}$ $= 693.1(4.0) \times 10^{-10}$ [0.6%]

(accounts for tensions in the data and differences between analyses)

$$a_{\mu}^{\text{hvp, LO}} = \left(\frac{\alpha m_{\mu}}{3\pi}\right)^2 \int_{m_{\pi^0}^2}^{\infty} ds \, \frac{R_{\text{had}}(s)\hat{K}(s)}{s^2}$$









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(accounts for tensions in the data and differences between analyses)

• Recent results in the $\pi^+\pi^-$ channel by CMD-3: \rightarrow further tension among e^+e^- data

 $a_{\mu}^{\text{hvp, LO}} = 707.6(3.4)_{\text{exp}}(0.7)_{\text{DV+QCD}} \times 10^{-10}$

(my own estimate)

 $a_{\mu}^{\text{hvp, LO}} = \left(\frac{\alpha m_{\mu}}{3\pi}\right)^2 \int_{m^2}^{\infty} ds \, \frac{R_{\text{had}}(s)K(s)}{s^2}$



[Ignatov et al. (CMD-3 Collab.), Phys. Rev. D109 (2024) 112002]







Lattice under Chromodynamics (QCD)

 Gauge theory of the strong interaction Non-perturbativaitreatenentoffotreageinteractiongivia regularised Euclidean path i

Lat**tiettisp** AGA: $a, x_{\mu} = n_{\mu}a, a^{-1} = \Lambda_{UV}$ • Ab initio treatment on discretised space-time Expectation value. v_{q} be v_{q} Monte $C_{q} v_{\mu}$ v_{μ} $v_$

Challenges for Lattice QCD calculations

Procedure: Noise problem: exponential growth of statistical fluctuations

- Bias from unsuppressed excited-state contributions
- Choose discretisation of OCD action Extrapolation to continuum limit: $a \rightarrow 0$
- Evaluate $\langle \Omega \rangle$ via Monte Carlo Integration:

generate ensembles of gauge configurations via a iviarkov chain

- Statistical error: $\sqrt{\Omega^2 \overline{\Omega}^2} \propto 1/N_{cfo}^{1/2}$ • Ensemble average: $\langle \Omega \rangle \simeq \Omega$



• Extrapolate observables to the continuum limit: $a \rightarrow 0$ and tune quark masses to physical values







Hadronic vacuum polarisation from Lattice QCD

Lattice QCD does **NOT** determine the *R*-ratio from first principles Time-momentum representation (TMR): [Bernecker & Meyer EPJA 47 (2011) 148]

$$a_{\mu}^{\text{hvp}} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^\infty dt \,\tilde{K}(t) \,G(t), \quad G(t) = -a^3 \sum_{\vec{x}} \left\langle J_k^{\text{e.m.}}(\vec{x}, t) J_k^{\text{e.m.}}(0) \right\rangle$$

- Not sensitive to exclusive hadronic channels



 $(\tilde{K}(t):$ known analytically)

• No reliance on experimental data, except for simple input quantities \rightarrow scale setting, calibration





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Challenges

- Exponentially increasing statistical noise as $t \rightarrow \infty$
- Correct for finite-volume effects
- Control discretisation effects ("lattice artefacts")
- Include isospin-breaking corrections



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Challenges

- Exponentially increasing statistical noise as $t \rightarrow \infty$
- Correct for finite-volume effects
- Control discretisation effects ("lattice artefacts")
- Include isospin-breaking corrections

Light-quark connected contribution dominates

• No reliance on experimental data, except for simple input quantities \rightarrow scale setting, calibration



 $\sim \sim \sim$

 \sim



Common discretisations of the quark action

Computational cost depends significantly on the chosen discretisation

"Fermion doubling problem"



Domain wall /overlap quarks:

- no doublers; chiral symmetry breaking exponentially small
- live in five dimensions (dwf)
- evaluate sign function of "conventional" action (ovlp)
- used by: RBC/UKQCD, χ QCD,...

Wilson quarks:

- no doublers; chiral symmetry broken explicitly
- "exceptional configurations":
 - negative eigenvalues of Wilson-Dirac operator
 - used by: Mainz/CLS, ETM, PACS







HVP in Lattice QCD



White Paper:

 $a_{\mu}^{\text{hvp, LO}} = (693.1 \pm 4.0) \cdot 10^{-10}$ *R*-ratio: [0.6%] LQCD: $a_u^{\text{hvp, LO}} = (711.6 \pm 18.4) \cdot 10^{-10}$ [2.6%] RBC/UKQCD [Blum et al., Phys. Rev. Lett. 121 (2018) 022003]

- Domain wall fermions
- Two ensembles: a = 0.114, 0.084 fm at m_{π}^{phys}
- Leading isospin-breaking corrections included
- Naive continuum extrapol'n in a^2 including estimated a^4 -term

 $a_{\mu}^{\text{hvp,LO}} = (715.4 \pm 16.3 \pm 9.2) \cdot 10^{-10}$ [2.6%]

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- O(*a*) improved Wilson fermions
- Four lattice spacings: $a = 0.085 0.050 \,\mathrm{fm}$
- Pion masses $m_{\pi} = 130 420 \,\mathrm{MeV}$
- Isospin-breaking correction by ETMC added to error
- Simultaneous chiral and continuum extrapolation



 $a_{\mu}^{\text{hvp,LO}} = (720.0 \pm 12.6 \pm 9.9) \cdot 10^{-10}$ [2.2%]



HVP in Lattice QCD



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BMWC [Borsányi et al., Nature 593 (2021) 7857]

- Rooted staggered fermions
- Six lattice spacings: a = 0.132 0.064 fm
- Physical pion mass throughout
- Correct for taste-breaking before continuum extrapol'n
- Final result selected from distribution of different fits







Window observables

 $a_{\mu}^{\text{hvp, win}} = \left(\frac{\alpha}{\pi}\right)^2 \int_{0}^{\infty} dt \,\tilde{K}(t) \,G(t) \,W(t; t_0, t_1)$ **Idea:** restrict integration to "unproblematic" regions → reduce statistical fluctuations and systematic effects 0.5Intermediate-distance window: statistical noise 0.4finite-volume effects $W^{\text{ID}}(t; t_0, t_1) = \Theta(t, t_0, \Delta) - \Theta(t, t_1, \Delta)$ 0.3 $\Theta(t, t', \Delta) = \frac{1}{2} \left[1 + \tanh(t - t') / \Delta \right]$ 0.2 $t_0 = 0.4 \,\text{fm}, t_1 = 1.0 \,\text{fm}, \Delta = 0.15 \,\text{fm}$



Data-driven approach: $a_{\mu}^{\text{win}} = (229.4 \pm 1.4) \cdot 10^{-10}$ [Colangelo et al., Phys Lett B833 (2022) 137313]

[Blum et al., Phys. Rev. Lett. 121 (2018) 022003]

- Finite-volume correction reduced to 0.25%
- Uncertainty dominated by statistics

Benchmark quantity for sub-contribution of HVP \rightarrow

(Excluding the 2023 CMD-3 result for $e^+e^- \rightarrow \pi^+\pi^-$)





Intermediate window observable in Lattice QCD

BMWc: Rooted staggered quarks



 $a_{\mu}^{\text{win,ud}} = (207.3 \pm 0.4 \pm 1.3^{\text{(jight)}_{\text{ise}}} \cdot 10^{-10}$

[Borsányi et al., Nature 593 (2021) 7857]

Mainz/CLS: O(a) improved Wilson quarks



 $a_{\mu}^{\text{win,ud}} = (207.0 \pm 0.8 \pm 1.2) \cdot 10^{-10}$

[Cè et al., Phys Rev D106 (2022) 114502]



Window observable: Lattice QCD vs. *R*-ratio



• Tension of 3.8σ in the window observable evaluated from e^+e^- data* and four lattice calculations $= (6.8 \pm 1.8) \cdot 10^{-10}$ [3.8 σ]

$$a_{\mu}^{\mathrm{win}}\Big|_{\langle \mathrm{lat} \rangle} - a_{\mu}^{\mathrm{win}}\Big|_{e^+e^-} =$$

• Subtract *R*-ratio result $a_{\mu}^{\text{win}}|_{e^+e^-}$ from WP estimate and replace by lattice average $a_{\mu}^{\text{win}}|_{\langle \text{lat} \rangle}$:

 $a_{\mu}^{\exp} - a_{\mu}^{SM} \Big|_{e^+e^- \to \langle lat \rangle}^{\text{win}} = (18.1 \pm 4.8) \cdot 10^{-10} \quad [3.8 \,\sigma]$

*excluding the CMD-3 result

- Dominant light-quark contribution confirmed for wide range of discretisation with sub-percent precision
- Significant tension with results based on the *R*-ratio*
 - $a_{\mu}^{\rm win} = (229.4 \pm 1.4) \cdot 10^{-10}$ *R*-ratio estimate: $a_{\mu}^{\rm win} = (236.16 \pm 1.09) \cdot 10^{-10}$ Lattice average:
 - (RBC/UKQCD 23, ETMC 22, Mainz/CLS 22, BMW 20) [HW, arXiv:2306.04165]



What can we learn from a_{μ}^{win} ?

Primary observable in lattice calculations: vector correlator G(t)

$$G(t) \equiv -\frac{a^3}{3} \sum_{k} \sum_{\vec{x}} \left\langle j_k^{\text{em}}(\vec{x}, t) j_k^{\text{em}}(0) \right\rangle = \frac{1}{12\pi^2} \int_{m_{\pi^0}^2}^{\infty} d(\sqrt{s}) R(s)^{\text{lat}} s \, e^{-\sqrt{s}t}$$
$$a_{\mu}^{\text{win}} \Big|_{\text{lat}} > \left. a_{\mu}^{\text{win}} \right|_{e^+e^-} \text{ implies that } R(s)^{\text{lat}} > R(s)^{e^+e^-} \text{ in some interval of } \sqrt{s}$$

Energy interval $600 \le \sqrt{s} \le 900 \,\text{MeV}$ contributes the same fraction to a_{μ}^{hvp} and a_{μ}^{win}

\sqrt{s} interval	$a_{\mu}^{ m hvp}$	$(a_{\mu}^{\mathrm{hvp}})^{\mathrm{SD}}$	$(a_{\mu}^{\mathrm{hvp}})^{\mathrm{ID}}$	$(a_{\mu}^{\mathrm{hvp}})^{\mathrm{LD}}$	$\bar{\Pi}(1 \text{ GeV}^2)$
Below 0.6 GeV	15.5	1.5	5.5	23.5	8.2
0.6 to 0.9 GeV	58.3	23.1	54.9	65.4	52.6
Above 0.9 GeV	26.2	75.4	39.6	11.1	39.2
Total	100.0	100.0	100.0	100.0	100.0



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[Cè et al., Phys Rev D106 (2022) 114502]

What can we learn from a_{μ}^{win} ?

• Phenomenological model for *R*-ratio predicts

 $\sqrt{s} = 600 - 900 \text{ MeV}$: $\frac{R(s)^{\text{lat}}}{R(s)^{e^+e^-}} = 1 + 1$

- Lattice average vs. *R*-ratio: $(a_{\mu}^{\text{win}})^{\text{lat}}/(a_{\mu}^{\text{win}})^{e^+e^-} = 1.030(8)$ \Rightarrow $R(s)^{\text{lat}}$ is enhanced by 5% relative to $R(s)^{e^+e^-}$ for $\sqrt{s} = 600 - 900 \text{ MeV}$
- If confirmed, it would imply that BMW's estimate might be too low....

Similar conclusions

- Dispersive treatment of pion form factor
- "Energy-smeared" *R*-ratio from lattice data

[Mainz/CLS, Cè et al., Phys Rev D 106 (2022) 114502]

$$\epsilon \implies \frac{(a_{\mu}^{\text{hvp}})^{\text{lat}}}{(a_{\mu}^{\text{hvp}})^{e^+e^-}} \approx \frac{(a_{\mu}^{\text{win}})^{\text{lat}}}{(a_{\mu}^{\text{win}})^{e^+e^-}} = 1 + 0.6\epsilon$$

[Colangelo, Hoferichter, Stoffer, PLB 814 (2021) 136073] [ETMC, Alexandrou et al., PRL 130 (2023) 241901]



Update by BMW Collaboration

• More statistics, added lattice spacing (a = 0.048 fm)



 $a_{\mu}^{\text{hvp,LO}} = (707.5 \pm 2.3 \pm 5.0) \cdot 10^{-10} \quad [0.8\%] \quad \Rightarrow \quad (714.5 \pm 2.2 \pm 2.5) \cdot 10^{-10}$

[Boccaletti et al., 2407.10913]

[0.5%]



Hadronic running of electromagnetic coupling

Electromagnetic coupling is energy-dependent:

 $\alpha^{-1} = 137.035999...$ $\alpha(q^2) = \frac{\alpha}{1 - \Delta \alpha(q^2)}$ $\alpha^{-1}(M_Z^2) = 127.951 \pm 0.009$

Correlation between $a_{\mu}^{\rm hvp}$ and the hadronic running of $\Delta \alpha_{\rm had}$:

$$\Delta \alpha_{\text{had}}(q^2) = -\frac{\alpha q^2}{3\pi} \mathcal{J}_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{s(s-q^2)},$$

Euclidean momenta

 $\Delta \alpha_{had}(-Q^2)$ accessible in lattice QCD via the same correlator G(t) with a different kernel function: ∞

$$\Delta \alpha_{\rm had}(-Q^2) = \frac{\alpha}{\pi} \frac{1}{Q^2} \int_0^\infty$$

Hadronic running at Z-pole: $\Delta \alpha_{had}^{(5)}(M_Z^2) \rightarrow key quantity in global electroweak fit$

$$a_{\mu}^{\text{hvp}} = \left(\frac{\alpha m_{\mu}}{3\pi}\right)^2 \int_{m_{\pi^0}^2}^{\infty} ds \, \frac{R(s) \, \hat{K}(s)}{s^2}$$

$$dt G(t) \left[Q^2 t^2 - 4 \sin^2 \left(\frac{1}{2} Q^2 t^2 \right) \right]$$

Evaluation of $\Delta \alpha_{had}^{(5)}(M_Z^2)$ and comparison with EW precision data

Adler function approach, aka. "Euclidean split technique"

 $\Delta \alpha_{\rm had}^{(5)}(M_Z^2) = \Delta \alpha_{\rm had}^{(5)}(-Q_0^2) \quad \leftarrow \text{ lattice QCD}$

$$+[\Delta \alpha_{\rm had}^{(5)}(-M_Z^2) - \Delta \alpha_{\rm had}^{(5)}(-Q_0^2)]$$

 $+ [\Delta \alpha_{\text{had}}^{(5)}(M_Z^2) - \Delta \alpha_{\text{had}}^{(5)}(-M_Z^2)] \quad \leftarrow \text{ pQCD}$

- $\binom{2}{0} \leftarrow \text{perturbative Adler function}$



Evaluation of $\Delta \alpha_{had}^{(5)}(M_Z^2)$ and comparison with EW precision data

Adler function approach, aka. "Euclidean split t

$$\Delta \alpha_{\rm had}^{(5)}(M_Z^2) = \Delta \alpha_{\rm had}^{(5)}(-Q_0^2)$$

$$+[\Delta \alpha_{\rm had}^{(5)}(-M_Z^2) - \Delta \alpha_{\rm had}^{(5)}(-Q_0^2)]$$

+
$$[\Delta \alpha_{\text{had}}^{(5)}(M_Z^2) - \Delta \alpha_{\text{had}}^{(5)}(-M_Z^2)$$

 $\Rightarrow \quad \Delta \alpha_{\text{had}}^{(5)}(M_Z^2) = 0.02773(9)_{\text{lat}}(2)_{\text{btm}}(12)_{\text{pQCE}}$

[Mainz/CLS, Cè et al., JHEP 08 (2022) 220, arXiv:2203.

No inconsistency with global electroweak fit

Standard Model can accommodate a larger value for a_{μ} without contradicting electroweak precision data

technique"	lat. + pQCD'[Adler]
	lat. + KNT18[data]
	R-ratio
	KNT18/19
)]	DHMZ19 HOH
	Jegerlehner 19
	EW global fits
	Gfitter 18
	Crivellin <i>et al.</i> 20
)	Keshavarzi <i>et al.</i> 20
	Malaescu, Schott 20
.08676]	HEPfit 21
0	0255 0.0260 0.0265 0.0270 0.0275 0.0280 0.0285 0.029
t!	$\Delta lpha_{ m had}^{(5)}(M_Z^2)$





Summary and outlook

No straightforward theoretical interpretation of the Fermilab E989 experiment Discrepant determinations of the HVP contribution:

- Tensions between lattice QCD and e^+e^- hadronic cross sections*
- Tension in $\pi^+\pi^-$ channel between BaBar vs. KLOE and CMD-3 vs. all other results

Analyses / re-analyses of e^+e^- data in progress: BaBar, BESIII, CMD-3, KLOE Experimental measurement of the HVP contribution by MUonE experiment Fermilab E989 prepares to release result including data from Runs 4–6 Update of White Paper expected by \approx Dec 2024, including new lattice results(?)

- Lattice QCD to produce more results for HVP contribution with sub-percent precision





^{*}pre-2023



September 9-13, 2024 @ KEK, Tsukuba, Japan

https://conference-indico.kek.jp/event/257

Gilberto Colangelo (University of Bern) Michel Davier (University of Paris-Saclay and CNRS, Orsay), co-chair Aida X. El-Khadra (University of Illinois), chair Martin Hoferichter (University of Bern) Christoph Lehner (University of Regensburg), co-chair Lee Roberts (Boston University)

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