# Production of loosely-bound objects in heavy-ion collisions at RHIC and LHC



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- Introduction
- Nuclei and Exotica
  - (Anti-)nuclei
  - (Anti-)hypernuclei
- Summary & Outlook







### Zoo of hadrons

#### **Baryons**



Proton (p)  $\rightarrow$  uud Neutron (n)  $\rightarrow$  udd Lambda ( $\Lambda$ )  $\rightarrow$  uds



#### <u>Mesons</u>

 $\pi$ -Meson  $\rightarrow d\bar{u}$ K-Meson  $\rightarrow u\bar{s}$ 

- Hadrons are consisting of quarks, anti-quarks und gluons
- Strangeness as new quark flavour not part of every-day matter, but is created for instance in high-energy particle collisions
- Theoretical description of hadrons through quantum chromo dynamics (QCD)



### Collisions

- Nuclei are accelerated to high energies, i.e. speeds close to the speed of light, and are then collided
- This leads to the creation of (new) particles that can be detected in the experiments surrounding the collision point



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# Experiments: STAR

STAR

GOETHE

UNIVERSITÄT FRANKFURT AM MAIN



**Experiments: ALICE** 

GOETHE

UNIVERSITÄT FRANKFURT AM MAIN







Pb

Pb

# Interlude: Centrality

Central Pb-Pb collision: High multiplicity = large  $dN/d\eta$ High number of tracks (more than 2000 tracks in the detector)

Peripheral Pb-Pb collision: Low multiplicity = small  $dN/d\eta$ Low number of tracks (less than 100 tracks in the detector)



#### Introduction



Time  $\rightarrow$ 

Cartoon of a Ultra-relativistic heavy-ion collision

Left to right:

- the two Lorentz contracted nuclei approach,
- collide,
- form a Quark-Gluon Plasma (QGP),
- the QGP expands and hadronizes,
- finally hadrons rescatter and freeze

Plot by S. Bass, Duke University; http://www.phy.duke.edu/research/NPTheory/QGP/transport/evo.jpg





beam

#### The fireball evolution:

- Starts with a "pre-equilibrium state"
- Forms a Quark-Gluon Plasma phase (if T is larger than  $T_c$ )

beam

- At chemical freeze-out, T<sub>ch</sub>, hadrons stop being produced
- At kinetic freeze-out, T<sub>fo</sub>, hadrons stop scattering



### Lattice QCD results



A. Bazavov et al. (hotQCD) Phys. Rev. D90 (2014) 094503 & PLB 795 (2019) 15 Similar results from Budapest-Wuppertal group: S. Borsányi et al. JHEP 09 (2010) 073 & PRL 125 (2020) 052001



#### Temperature of the source



Analogy:

Light source  $\rightarrow$  particle source

 Multiplicity described best with T = 1 900 000 000 000 °C (1,9 trillion degree centigrade)

 $\rightarrow$  100 000 times hotter than in the interior of the sun!

1/40 eV = 20 °C

Plot by A. Andronic, GSI-Heidelberg group arXiv:1407.5003 [nucl-ex]



# Thermal model

• Statistical (thermal) model with only three parameters able to describe particle yields (grand chanonical ensemble)



- chemical freezeout temperature T<sub>ch</sub>
- baryo-chemical
   potential μ<sub>B</sub>
- Volume V
- → Using particle yields as input to extract parameters



A. Andronic et al., PLB 673 (2009) 142, updated



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# Predicting yields of bound

![](_page_17_Figure_1.jpeg)

Key parameter at LHC energies:

chemical freeze-out temperature T<sub>ch</sub>

Strong sensitivity of

abundance of nuclei

to choice of T<sub>ch</sub> due to:

1. large mass m

2. exponential dependence of the yield ~  $exp(-m/T_{ch})$ 

→ Binding energies small compared to  $T_{ch}$ 

![](_page_18_Picture_0.jpeg)

#### (Anti-)Nuclei

![](_page_18_Figure_2.jpeg)

![](_page_19_Picture_0.jpeg)

#### Coalescence

![](_page_19_Figure_2.jpeg)

J. I. Kapusta, PRC 21, 1301 (1980)

Nuclei are formed by protons and neutrons which are nearby and have similar velocities (after kinetic freezeout)

Produced nuclei

- → can break apart
- → created again by final-state coalescence

![](_page_20_Picture_0.jpeg)

![](_page_20_Figure_1.jpeg)

#### Low momenta:

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Nuclei are identified using the d*E*/dx measurement in the Time Projection Chamber (TPC)

![](_page_20_Figure_4.jpeg)

ALICE

#### Higher momenta:

Velocity measurement with the Time-of-Flight (TOF) detector is used to calculate the  $m^2$ distribution

![](_page_21_Picture_0.jpeg)

Anti-Alpha

![](_page_21_Picture_2.jpeg)

![](_page_21_Picture_3.jpeg)

For the full statistics of 2011 ALICE identified 10 Anti-Alphas using TPC and TOF

STAR observed the Anti-Alpha in 2010: *Nature 473, 353 (2011)* 

![](_page_21_Figure_6.jpeg)

![](_page_22_Figure_0.jpeg)

- *p*<sub>T</sub> spectra getting harder for more central collisions (from pp to Pb-Pb) → showing clear radial flow
- Blast-Wave fits describe the data in Pb-Pb very well
- No hint for radial flow in pp

![](_page_23_Picture_0.jpeg)

#### (Anti-)Deuteron ratio

![](_page_23_Picture_2.jpeg)

![](_page_23_Figure_3.jpeg)

![](_page_24_Picture_0.jpeg)

#### **Combined Blast-Wave fit**

![](_page_24_Picture_2.jpeg)

#### ALICE Collaboration, arXiv:1910.07678, see also arXiv:2311.11758

![](_page_24_Figure_4.jpeg)

- Simultaneous Blast-Wave fit of  $\pi^+$ , K<sup>+</sup>, p, d, t, <sup>3</sup>He and <sup>4</sup>He spectra for central Pb-Pb collisions leads to values for  $\langle \beta \rangle$  and  $T_{kin}$  close to those obtained when only  $\pi$ ,K,p are used
- All particles are described rather well with this simultaneous fit

# Mass dependence

![](_page_25_Picture_1.jpeg)

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ALICE Production of (anti-) nuclei is follwing an exponential, and decreases with mass as expected from thermal model In Pb-Pb the "penalty factor" for each additional baryon ~300 (for particles and antiparticles)

ALICE Collaboration, arXiv:1710.07531, NPA 971, 1 (2018)

# Mass dependence

![](_page_26_Picture_1.jpeg)

![](_page_26_Figure_2.jpeg)

ALICE

- Production of (anti-) nuclei is follwing an exponential, and decreases with mass as expected from thermal model
- In Pb-Pb the "penalty factor" for each additional baryon ~300, in p-Pb ~600 and in pp ~1000

![](_page_27_Figure_0.jpeg)

ALICE

d/p ratio rather well described by coalescence and (canonical) thermal model

# Ratios vs. multiplicity

![](_page_28_Picture_1.jpeg)

ALICE Collaboration, arXiv:2211.14015, Phys.Rev.C 107 (2023) 064904

![](_page_28_Figure_3.jpeg)

- d/p ratio rather well described by coalescence and (canonical) thermal model
- Some tension for <sup>3</sup>He/p and <sup>3</sup>H/p over  $p_{T}$

![](_page_29_Figure_0.jpeg)

<sup>4</sup>He/p ratio significantly better described by the thermal model

![](_page_30_Picture_0.jpeg)

![](_page_30_Figure_1.jpeg)

![](_page_30_Figure_2.jpeg)

Different model implementations describe the production probability, including light nuclei and hyper-nuclei, rather well at a temperture of about  $T_{ch}$  =156 MeV

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ALICE Collaboration, arXiv:1710.07531 NPA 971, 1 (2018) NPA

![](_page_31_Picture_0.jpeg)

#### Hypernuclei

![](_page_31_Figure_2.jpeg)

15

PSR J0348+0432

PSR J1614-2230

mass

14

### Neutron stars and interactions

**PNM** 

 $\Lambda N + \Lambda NN$  (II)

 $\Lambda N + \Lambda NN (I)$ 

ΛN

- Hyperon puzzle in neutron stars → hyperons make the EOS softer:
- Pure neutron matter (PNM) works well
- Known ∧N interaction → way to soft
- Including ANN 0.8 forces brings the mass slightly up
- Only additional 0.0 11 12 13
   ΛNN interaction works sufficiently <sup>R [km]</sup>

2.4

2.0

1.6

1.2

![](_page_32_Figure_7.jpeg)

![](_page_33_Picture_0.jpeg)

## Hypernuclei

- Hypernuclei are unique probes to study nuclear structure
- Single Λ-hypernuclei are major source of extracting Λ-N interaction
- Correct Λ-N and Λ-N-N interaction needed to understand structure of neutron stars

![](_page_33_Figure_5.jpeg)

D. Logoteta et al., Astron. Astrophys. 646 (2021) A55

![](_page_34_Picture_0.jpeg)

#### Hypertriton

Bound state of  $\Lambda$ , p, n m = 2.991 GeV/c<sup>2</sup> (B<sub> $\Lambda$ </sub> =130 keV)

![](_page_35_Picture_0.jpeg)

### Hypertriton

Bound state of  $\Lambda$ , p, n m = 2.991 GeV/ $c^2$  (B<sub> $\Lambda$ </sub> =130 keV)

![](_page_35_Figure_3.jpeg)

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![](_page_36_Picture_0.jpeg)

### Hypertriton

Bound state of  $\Lambda$ , p, n m = 2.991 GeV/c<sup>2</sup> (B<sub> $\Lambda$ </sub> =130 keV)

![](_page_36_Figure_3.jpeg)

![](_page_36_Figure_4.jpeg)

# **GOETHE WIVERSITÄT** Hypertriton Identification

![](_page_37_Figure_1.jpeg)

Bound state of  $\Lambda$ , p, n  $m = 2.991 \text{ GeV}/c^2 (B_{\Lambda} = 130 \text{ keV})$   $\rightarrow$  Radius of about 10.6 fm Decay modes:

$${}^{3}_{\Lambda} H \rightarrow^{3} H e + \pi^{-}$$

$${}^{3}_{\Lambda} H \rightarrow^{3} H + \pi^{0}$$

$${}^{3}_{\Lambda} H \rightarrow d + p + \pi^{-}$$

$${}^{3}_{\Lambda} H \rightarrow d + n + \pi^{0}$$

+ anti-particles

→ Anti-Hypertriton first observed by STAR Collaboration:

Science 328,58 (2010)

![](_page_38_Picture_0.jpeg)

![](_page_38_Picture_1.jpeg)

![](_page_38_Figure_2.jpeg)

- Clear signal reconstructed by decay products
- Spectra can also be described by Blast-Wave model
   → Hypertriton flows as all other particles

![](_page_39_Picture_0.jpeg)

### Hypertriton spectra

![](_page_39_Picture_2.jpeg)

![](_page_39_Figure_3.jpeg)

• Anti-hypertriton/Hypertriton ratio consistent with unity vs.  $p_{T}$ 

![](_page_40_Picture_0.jpeg)

![](_page_40_Picture_1.jpeg)

- Hypertriton signal recently also extracted in pp and p-Pb collisions
- Stronger separation between models as for other particle ratios, mainly due to the size of the hypertriton

Hypertriton in pp & p-Pb

![](_page_40_Figure_4.jpeg)

![](_page_41_Picture_0.jpeg)

![](_page_41_Picture_1.jpeg)

- Hypertriton signal recently also extracted in pp and p-Pb collisions
- Stronger separation between models as for other particle ratios, mainly due to the size of the hypertriton

![](_page_41_Figure_4.jpeg)

Hypertriton in pp & p-Pb

![](_page_42_Picture_0.jpeg)

# Hypertriton at RHIC

![](_page_42_Picture_2.jpeg)

 Hypertriton signal recently also extracted in isobar collisions

![](_page_42_Figure_4.jpeg)

![](_page_43_Picture_0.jpeg)

# Hypertriton/Λ ratio

![](_page_43_Picture_2.jpeg)

- Hypertriton signal recently also extracted in isobar collisions
- Stronger separation between models as for other particle ratios, mainly due to the size of the hypertriton

![](_page_43_Figure_5.jpeg)

![](_page_44_Picture_0.jpeg)

# A = 4 hypernuclei

- Large suppression expected for A
   = 4 hypernuclei by the SHM wrt A =
   3
- A = 4 hypernuclei are more bound and each has an excited state
   Phys. Rev. Lett. 115, 222501 (2015)
- The yields of these hypernuclei are enhanced with respect to the ground state due to the feed-down from excited states
- Also the yields of the SHM scale with the **spin-degeneracy**
- Resulting in a total enhancement of a factor 4 for both hypernuclei BD, EPJ Web Conf. 276 (2023) 04002

![](_page_44_Picture_8.jpeg)

M. Schäfer, N. Barnea, A. Gal, Phys.Rev.C 106, L031001 (2022)

![](_page_44_Picture_10.jpeg)

![](_page_45_Picture_0.jpeg)

![](_page_45_Picture_1.jpeg)

- For the first time, we are able to reconstruct A = 4 (anti)hypernuclei at the LHC and determine their production yield
- (Anti)hyperhydrogen-4 invariant-mass spectrum in Run 2
   Pb-Pb collisions at 5.02 TeV
- Examined in the two-body decay:

$$^{4}_{\Lambda}H \rightarrow {}^{4}He + \pi^{-} + c.c.$$

Reaching a local p-value of 6σ

![](_page_45_Figure_7.jpeg)

![](_page_45_Picture_8.jpeg)

![](_page_46_Picture_0.jpeg)

![](_page_46_Picture_1.jpeg)

- For the first time, we are able to reconstruct A = 4 (anti)hypernuclei at the LHC and determine their production yield
- (Anti)hyperhelium-4 invariant-mass spectrum in Run 2
   Pb-Pb collisions at 5.02 TeV
- Examined in the three-body decay:

 $^{4}_{\Lambda}\text{He} \rightarrow {}^{3}\text{He} + p + \pi^{-} + \text{c.c.}$ 

Reaching a local p-value of 4.4σ

![](_page_46_Figure_7.jpeg)

![](_page_46_Picture_9.jpeg)

![](_page_47_Picture_0.jpeg)

![](_page_47_Picture_1.jpeg)

- For the first time, we are able to reconstruct A = 4 (anti)hypernuclei at the LHC and determine their production yield
- First observation of the antihyperhelium-4 in Run 2 Pb-Pb collisions at 5.02 TeV

![](_page_47_Figure_4.jpeg)

![](_page_47_Figure_5.jpeg)

![](_page_47_Picture_6.jpeg)

![](_page_48_Picture_0.jpeg)

![](_page_48_Picture_1.jpeg)

- First measurement of the (anti)hyperhelium-4 production yield
- Testing the dependence of the yields of the SHM with the spin-degeneracy
- Our yields confirm the hypothesis of excited states for both (anti)hypernuclei within 2σ
- currently dominated by statistical uncertainties
- with more data, a high precision measurement will be feasible (like for the Λ hyperon)

![](_page_48_Figure_7.jpeg)

![](_page_48_Picture_9.jpeg)

#### **Outlook & Summary**

![](_page_49_Picture_1.jpeg)

![](_page_50_Picture_0.jpeg)

### Outlook

![](_page_50_Picture_2.jpeg)

![](_page_50_Figure_3.jpeg)

A. Andronic, private communication, model described in A. Andronic et al., PLB 697, 203 (2011) and references therein

- Explore QCD and QCD inspired model predictions for (unusual) multi-baryon states
- Search for rarely produced anti- and hyper-matter
- Test model predictions, e.g. thermal and coalescence

![](_page_51_Picture_0.jpeg)

### Conclusion

![](_page_51_Picture_2.jpeg)

- ALICE@LHC and STAR@RHIC are well suited to study light (anti-)(hyper-) nuclei and perform searches for exotic bound states (A<5)</li>
- Copious production of loosely bound objects measured by ALICE and STAR as predicted by the thermal model
- Models describe the data rather well
- Ratios vs. multiplicity trend described by both models
   only tension: Alpha vs. <sup>3</sup><sub>A</sub>H
- New and more precise data can be expected in the next years

![](_page_51_Figure_8.jpeg)