



Electron-beam based Neutron Sources

Roadmap for Future Accelerators- iFAST WP5.2 Workshop
3rd of September, 2024

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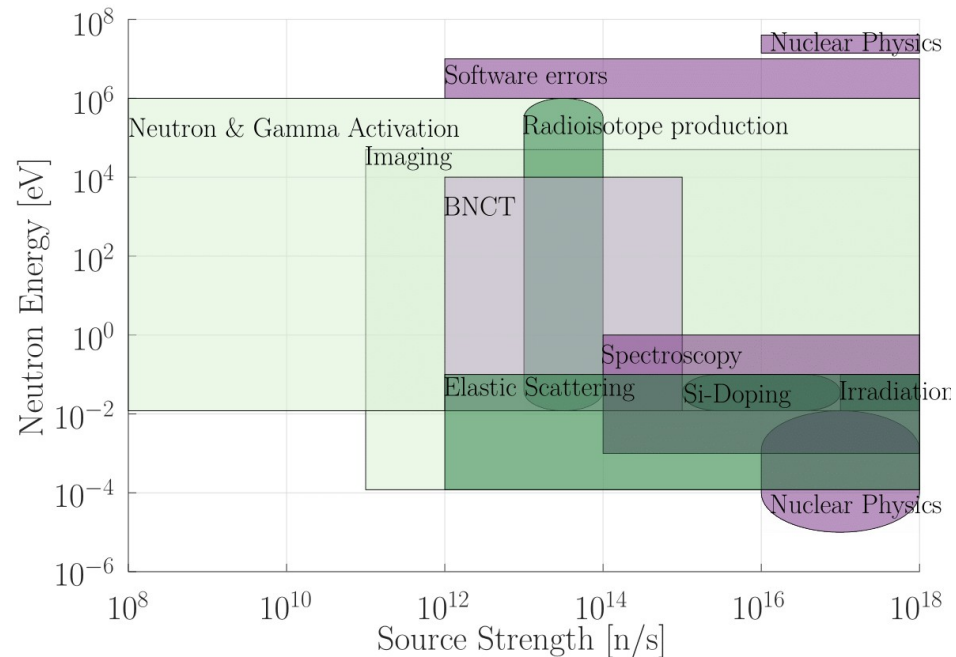
Is an electron linac a suitable driver for neutron production?

To answer this question, I will discuss:

- I. The **necessity** for neutron sources (Introduction)
- II. Neutron production **mechanism** with electrons
- III. Unmoderated neutron **spectrum characterization**
- IV. Comparison with the **state-of-the-art**
- V. Neutron moderation and **brightness/brilliance** discussion
- VI. **VULCAN** neutron source

I. Necessity for neutron sources

- **Uprising demand:** Wide variety of research areas make use of neutrons
 - Not only research: Industrial and medical applications! [1, 2]

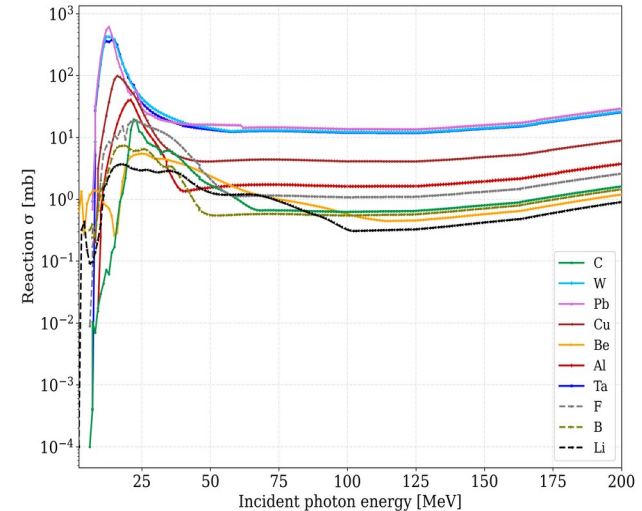


II. Neutron production with accelerators

- Neutron sources migrating from nuclear reactors to **accelerator-based facilities** [3]
- **Hadron**-based machines. **Direct** processes:
 - Spallation
 - Controlled nuclear reaction:
$$p + {}^7\text{Li} \rightarrow n + {}^7\text{Be}$$
- **Electron**-based machines. **Indirect** process:
 - Bremsstrahlung + Photonuclear reaction

↓
High Z

↓
High σ



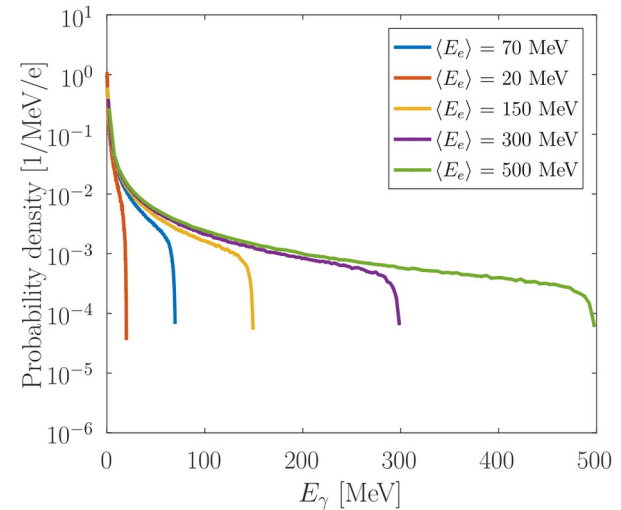
ENDF-v.VIII Photonuclear cross sections.

II. Neutron production with accelerators

- Neutron sources migrating from nuclear reactors to **accelerator-based facilities** [3]
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 - Spallation
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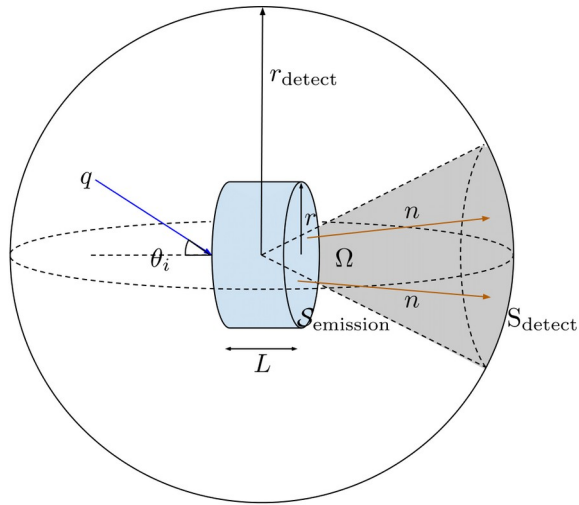


Bremstrahlung spectrum for different electron beams against tungsten.

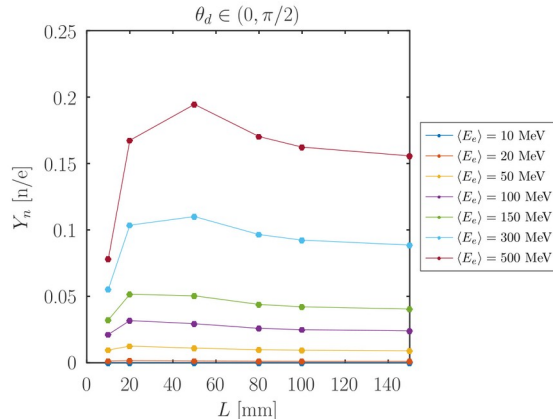
II. Neutron production with electrons

- Single tungsten target where **both** processes occur
- G4beamlines simulations [4]
 - Optimal dimensions: $r = 50\text{mm}$; $L = 80\text{ mm}$

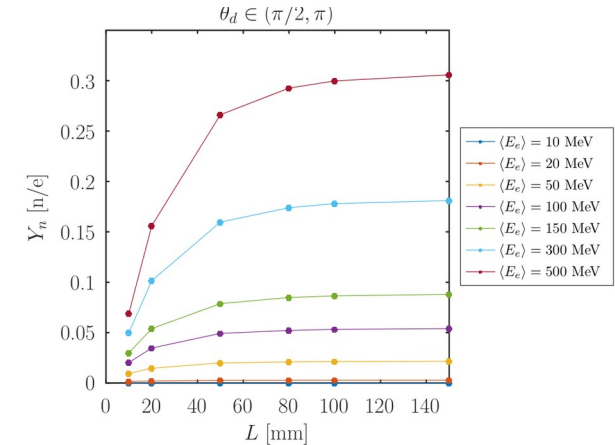
$$Y_n(E_e) \equiv \frac{N_n(E_e)}{N_e}$$



Neutron production setup



Forward neutron yield at $r = 50\text{ mm}$.
 $\Theta_i = 0\text{ deg}$



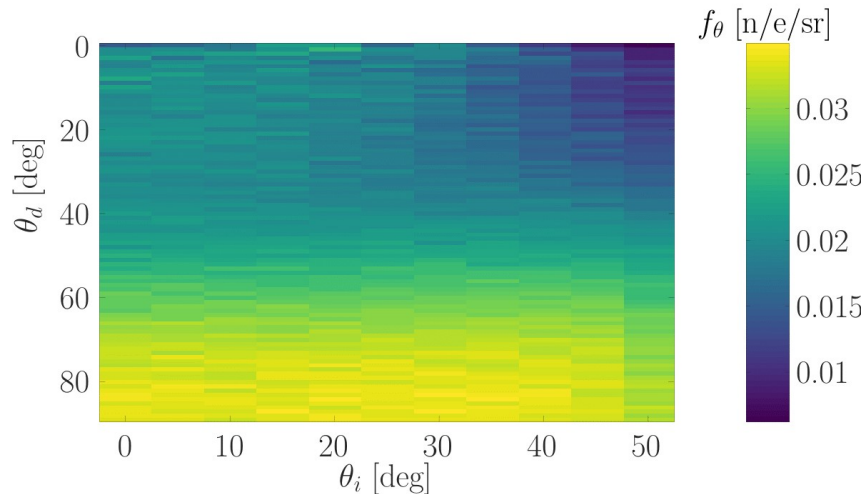
Backward neutron yield at $r = 50\text{ mm}$.
 $\Theta_i = 0\text{ deg}$

III. Unmoderated neutron spectrum

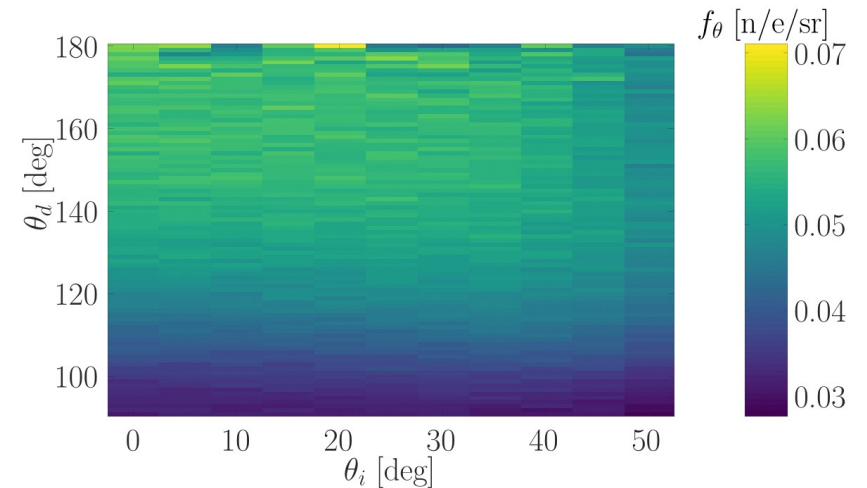
- For the optimized W target, we note:
 - Isotropy in incidence direction θ_i

$$f_{\Omega}(\varphi, \theta_d; E_e) \equiv \frac{d^2 Y_n}{d\Omega}(\varphi, \theta_d)$$

$$f_{\theta}(\theta_d; E_e) \equiv 2\pi \int_0^{2\pi} f_{\Omega}(\theta_d, \varphi) d\varphi$$



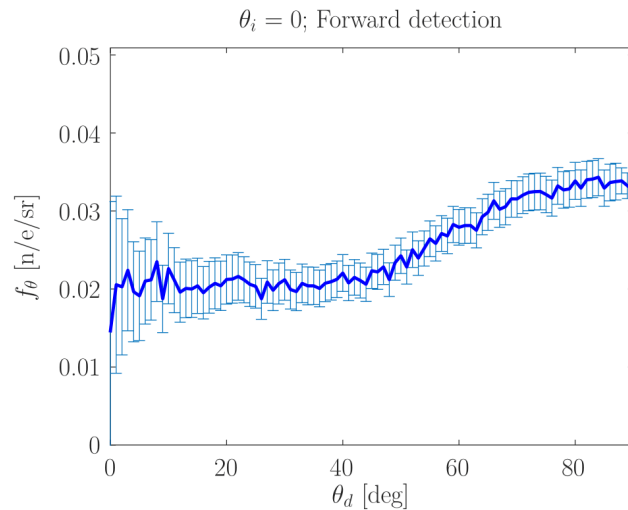
Forward neutron detection dependency with incidence angle and $\langle E_e \rangle = 500 \text{ MeV}$



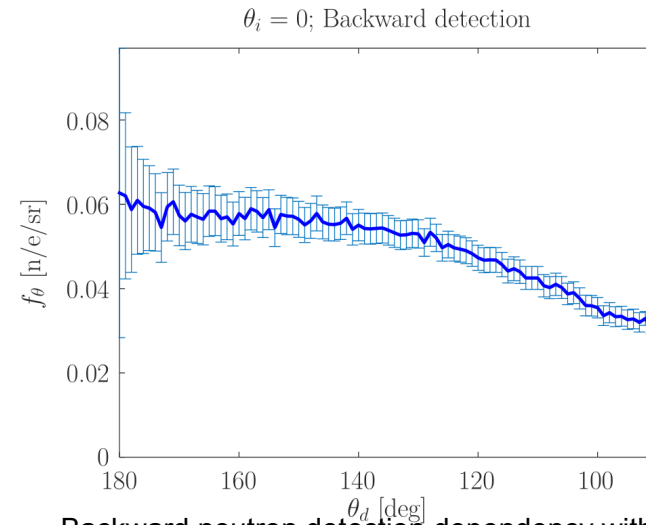
Backward neutron detection dependency with incidence angle and $\langle E_e \rangle = 500 \text{ MeV}$

III. Unmoderated neutron spectrum

- For the optimized W target, we note:
 - Isotropy in incidence direction θ_i
 - Isotropy in polar detecting angle: Up to 40 deg



Forward neutron detection dependency with incidence angle and $\langle E_e \rangle = 500$ MeV and $\theta_i = 0$ deg.

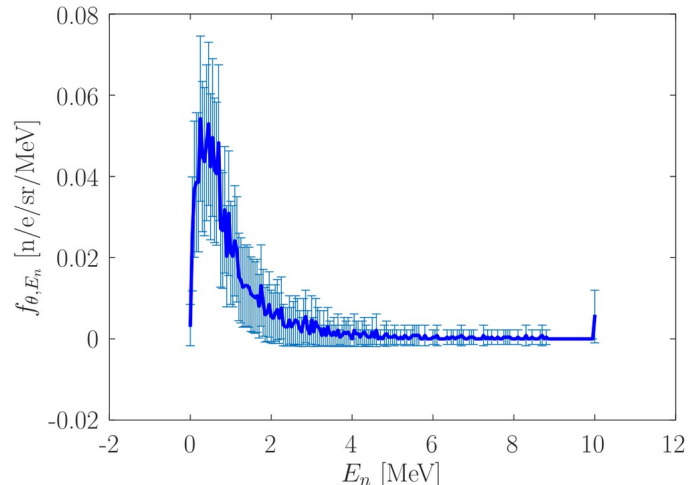


Backward neutron detection dependency with incidence angle and $\langle E_e \rangle = 500$ MeV and $\theta_i = 0$ deg.

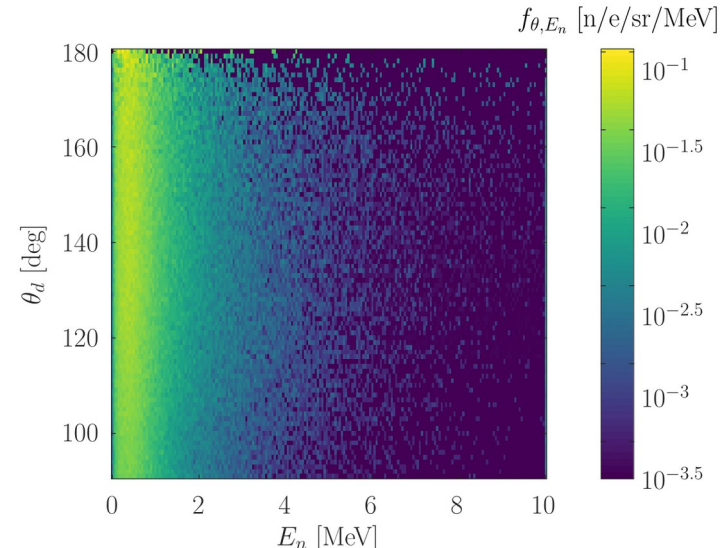
III. Unmoderated neutron spectrum

- For the optimized W target, we note:
 - Isotropy in incidence direction θ_i
 - Isotropy in polar detecting angle: Up to 40 deg
 - Maxwellian neutron emission with $\langle E_n \rangle \sim 1$ MeV

$$f_{\theta, E_n}(\theta_d, E_n; E_e) \equiv \frac{df_{\theta}}{dE_n}$$



Energy distribution for different detecting angles.
 $\langle E_e \rangle = 500$ MeV and $\theta_i = 0$ deg; $\theta_d = 165$ deg.



Energy distribution for different detecting angles.
 $\langle E_e \rangle = 500$ MeV and $\theta_i = 0$ deg.

III. Unmoderated neutron spectrum

- For the optimized W target, we note:
 - Isotropy in incidence direction θ_i
 - Isotropy in polar detecting angle: Up to 40 deg.
 - Maxwellian neutron emission with $\langle E_n \rangle \sim 1$ MeV
 - **Increase of σ_{E_n}** due to high E_n neutrons; **little change in $\langle E_n \rangle$**

$\langle E_e \rangle$ [MeV]	f_θ [10^{-2} n/e/sr]	$\langle E_n \rangle$ [MeV]	σ_{E_n} [MeV]
20	0.056 ± 0.009	0.776	0.625
50	0.41 ± 0.04	0.859	0.741
100	1.02 ± 0.09	0.986	0.973
300	3.4 ± 0.3	1.08	1.15
500	5.8 ± 0.5	1.08	1.17

Energy neutron spectrum details detected at $\theta_d = 165$ deg

IV. High intensity e-linac proposal

- Targeted figure of merit: **Source strength**

$$I_n \equiv I_{e,av} Y_n$$

- Two normal-conducting high-intensity linacs are considered

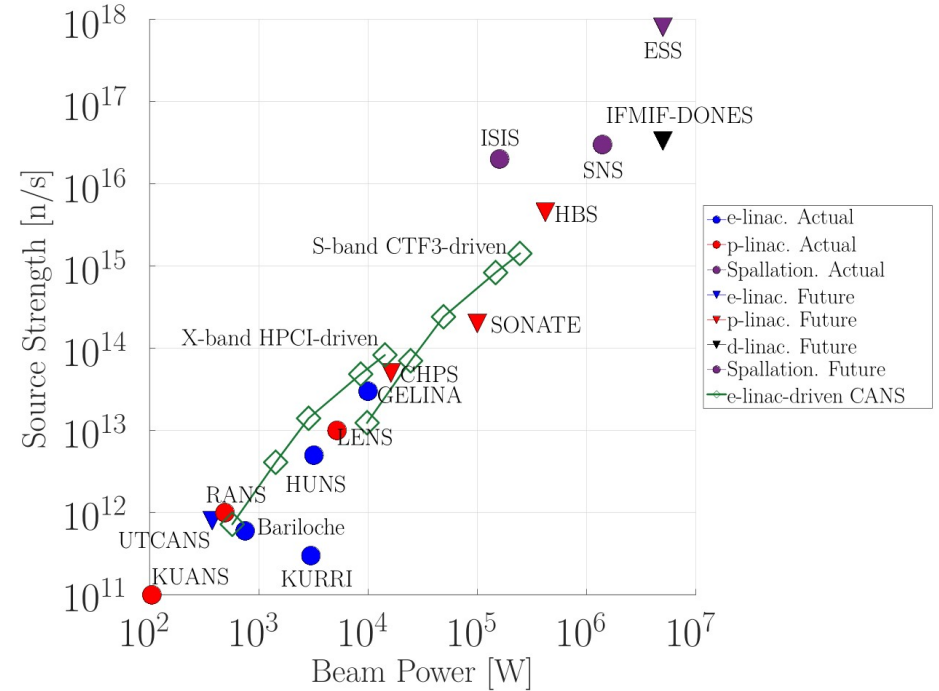
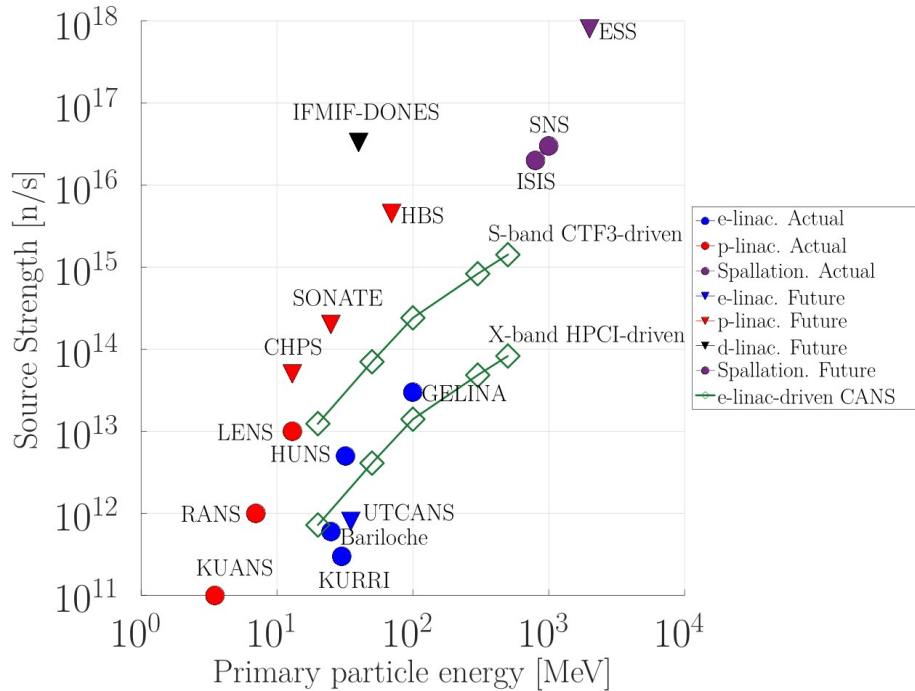
- **HPCI – linac:** S-band Photoinjector + X-band TW structures [5]
- **CTF3 drive-beam linac:** S-band Thermoionic gun + S-band TW structures [6]

Magnitude	Units	HPCI-linac	CTF3 drive beam linac
f	GHz	12.00	3.00
Q_{bunch}	nC	0.285	2.33
N_{bunches}		1000	2100
$f_{\text{RF-cycle}}$	Hz	100	100
$I_{e,av}$	μA	28.50	489.3

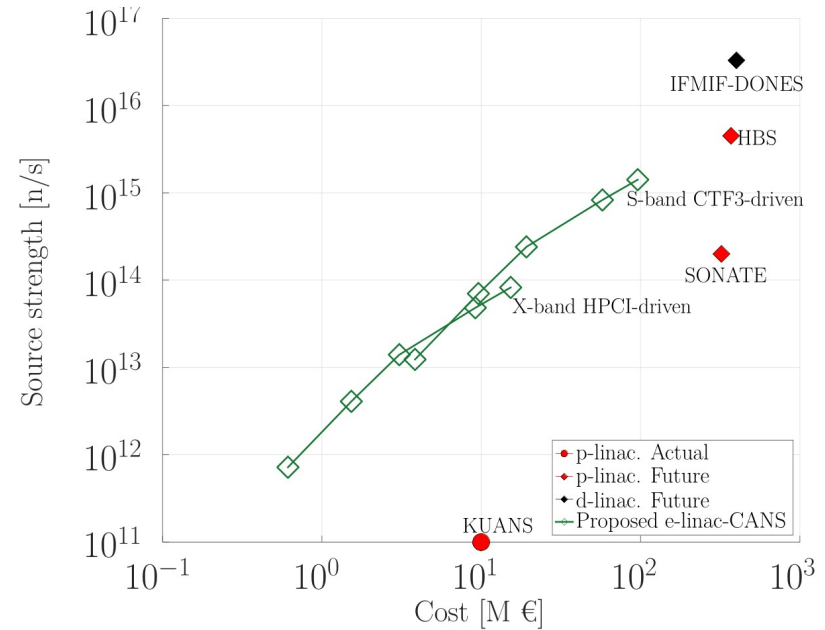
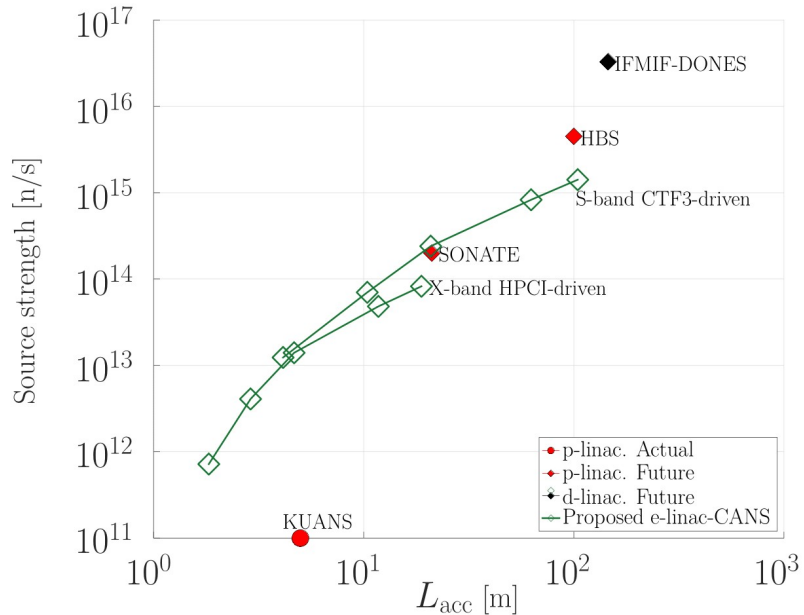
High-intensity compact linac specifications [5, 6]

IV. State-of-the-art comparison

[7]



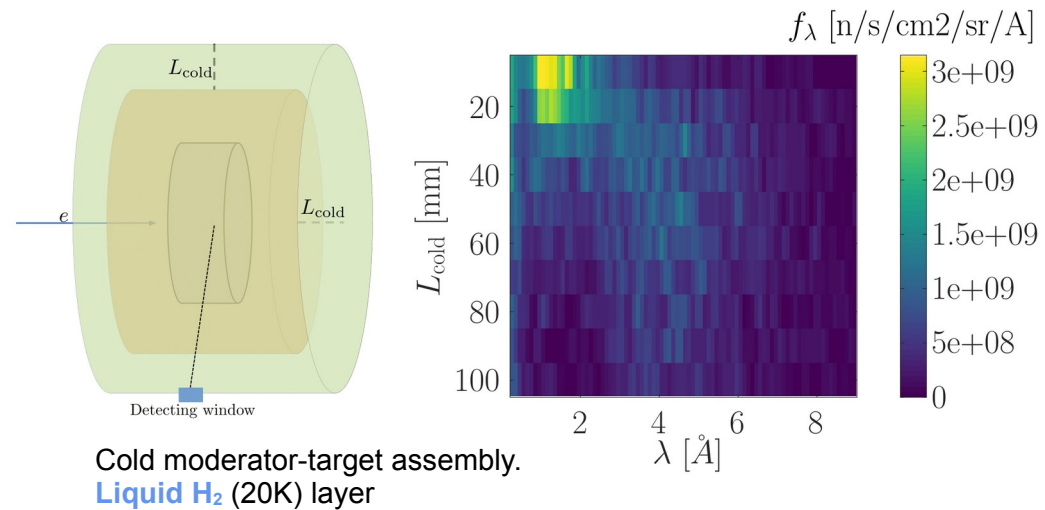
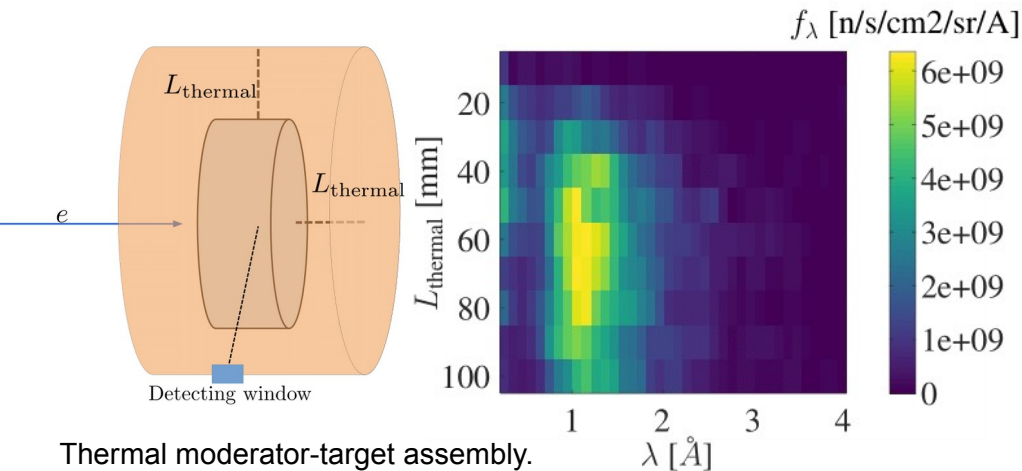
IV. State-of-the-art comparison



V. Thermal and cold neutron moderation

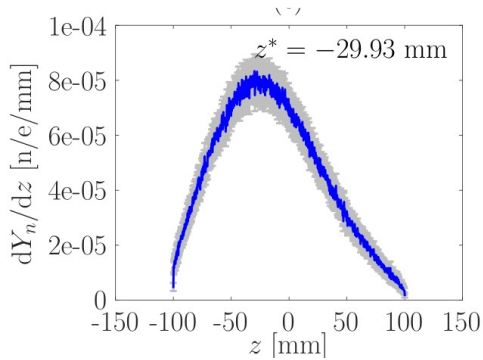
- Material science: **Diffraction** and **imaging** experiments
 - Require moderated neutrons – rich **H compounds**
- Targeted figure of merit: **Average brightness**

$$f_{\lambda}(\theta, \lambda; E_e) \equiv \frac{d^3 I_{n, av}}{dS_{\text{emission}} d\Omega d\lambda}$$

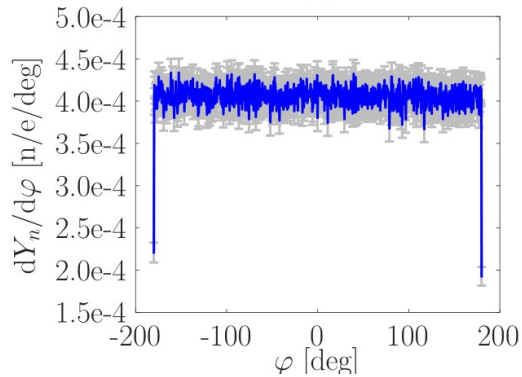
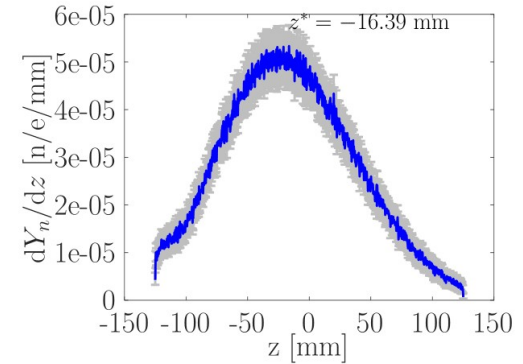


- Optimal dimensions:** $L_{\text{thermal}} = 60 \text{ mm}$; $L_{\text{cold}} = 25 \text{ mm}$.

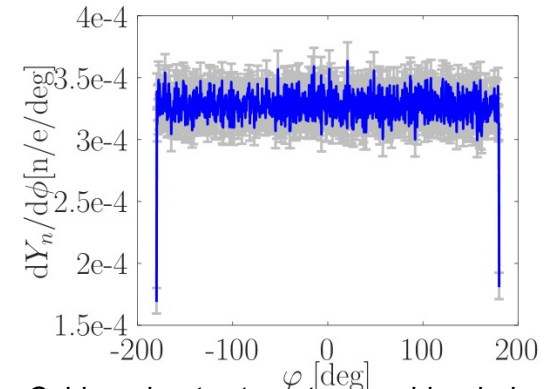
V. Thermal and cold neutron detection



Backward/Lateral maximum emission



Azimuthal isotropy



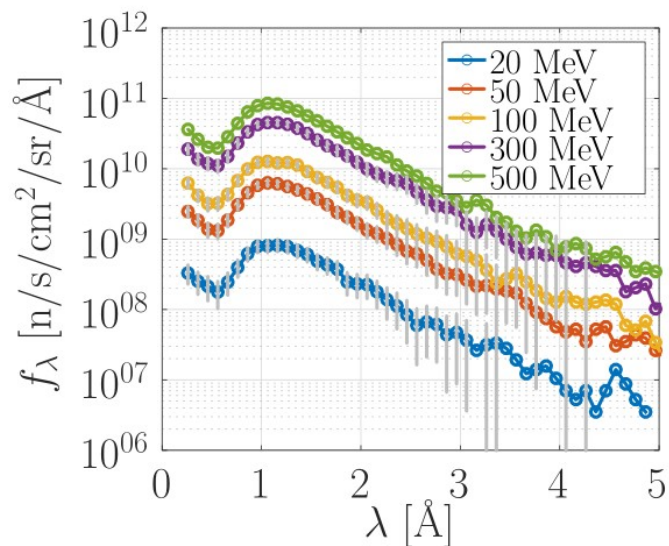
Thermal moderator-target assembly window location (up) and azimuthal dependency of the yield (down)

Cold moderator-target assembly window location (up) and azimuthal dependency of the yield (down)

V. Average Brightness

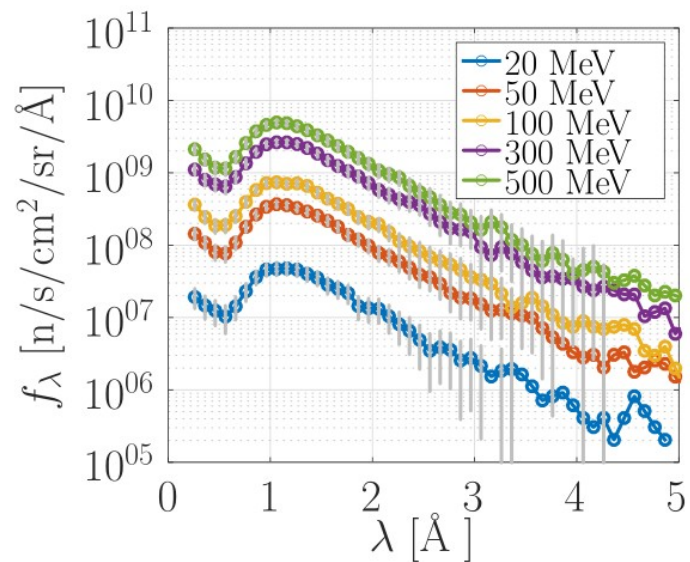
- Proportional to electron intensity (CTF3 > HPCI)

CTF3 d.b - thermal



Thermal moderator-target average brightness for different electron energies for CTF3 drive beam linac.

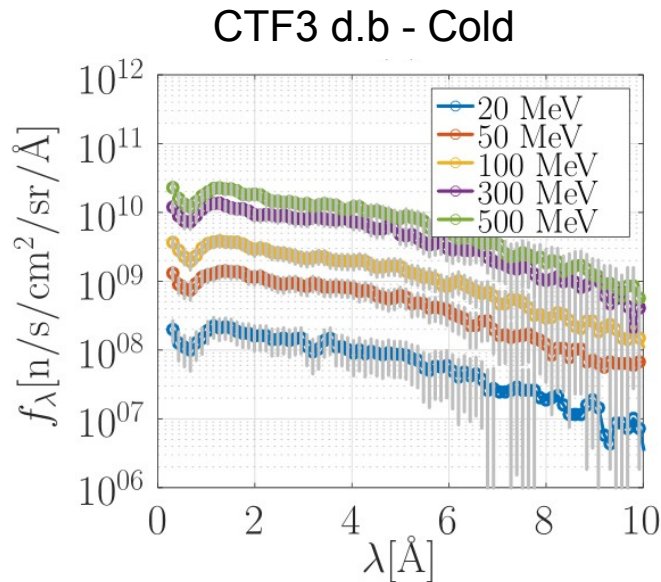
HPCI - thermal



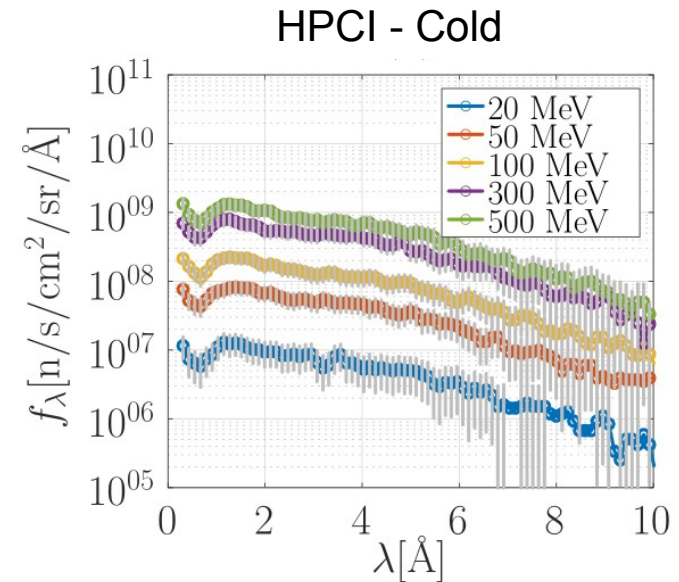
Thermal moderator-target average brightness for different electron energies for HPCI linac.

V. Average Brightness

- Proportional to electron intensity (CTF3 > HPCI)



Cold moderator-target average brightness for different electron energies for CTF3 drive beam linac.



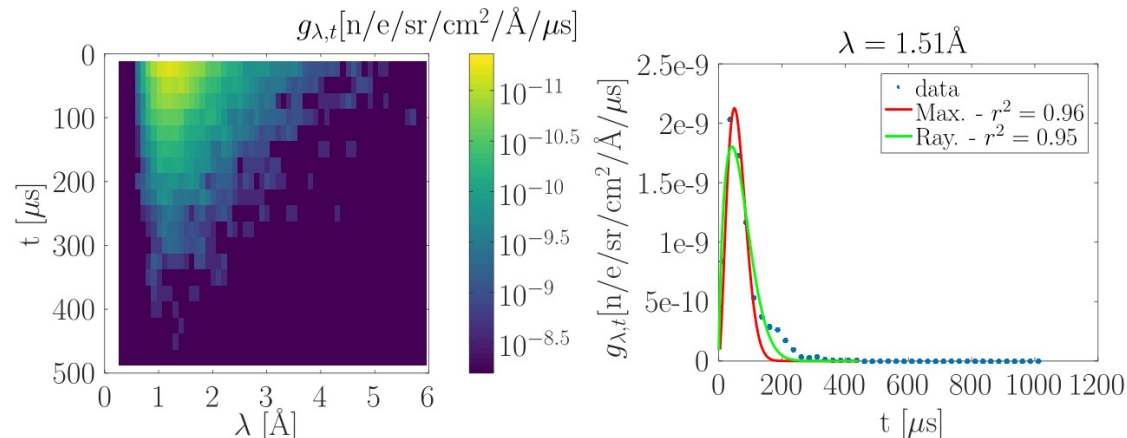
Cold moderator-target average brightness for different electron energies for HPCI linac.

V. Peak brightness

- Time-resolution of the brightness spectrum
 - Convolves electron pulse with neutron response

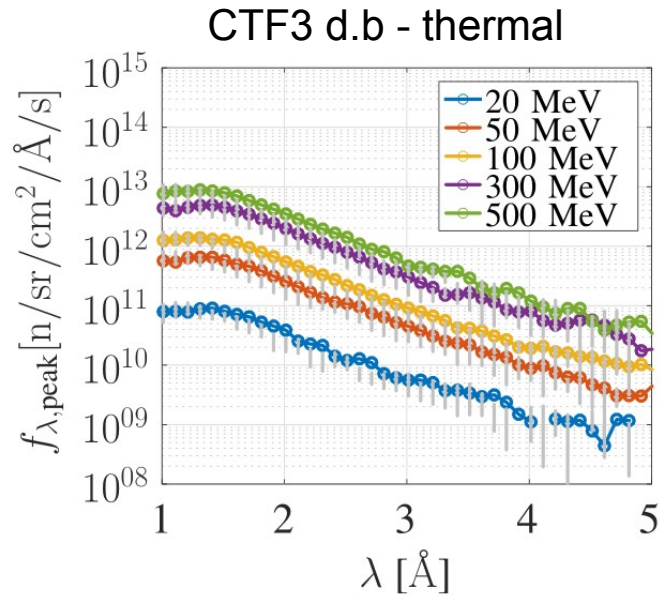
$$f_{\lambda, \text{peak}} \equiv \max_{t>0} \left(\frac{dq_e}{dt} \otimes g_{\lambda,t} \right) \quad g_{\lambda,t} \equiv \frac{{}^4Y_n}{dS_{\text{emission}} d\Omega d\lambda dt}$$

- Cold and thermal **neutron** responses extend **several μs** . GHz **electron** pulses extend **few ns**

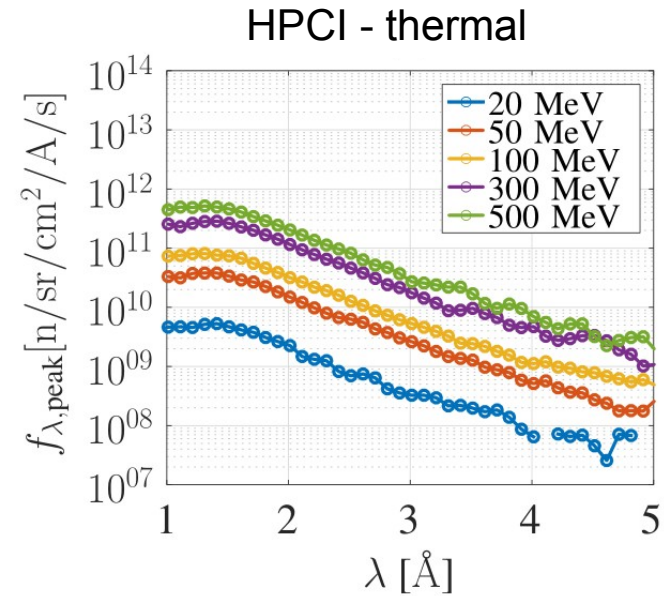


V. Peak brightness

- For the case of electron trains (few ns), the peak brightness is just the normalization of g_λ to the total train charge.



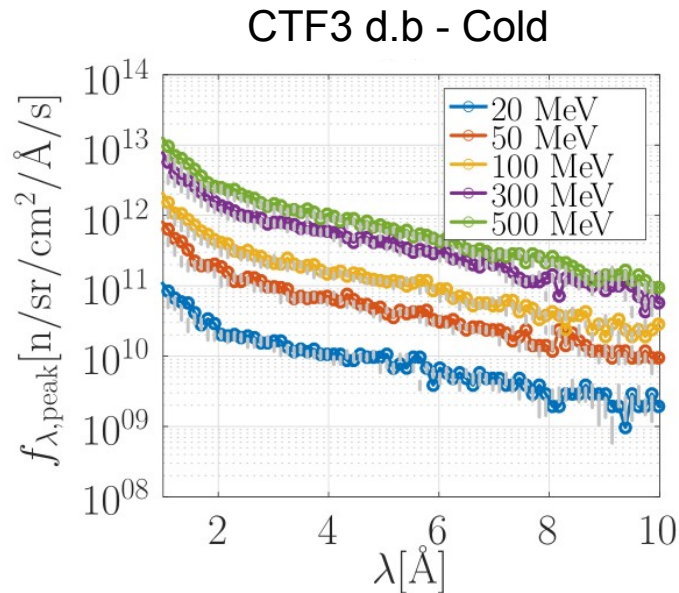
Thermal moderator-target peak brightness for different electron energies for CTF3 drive beam linac.



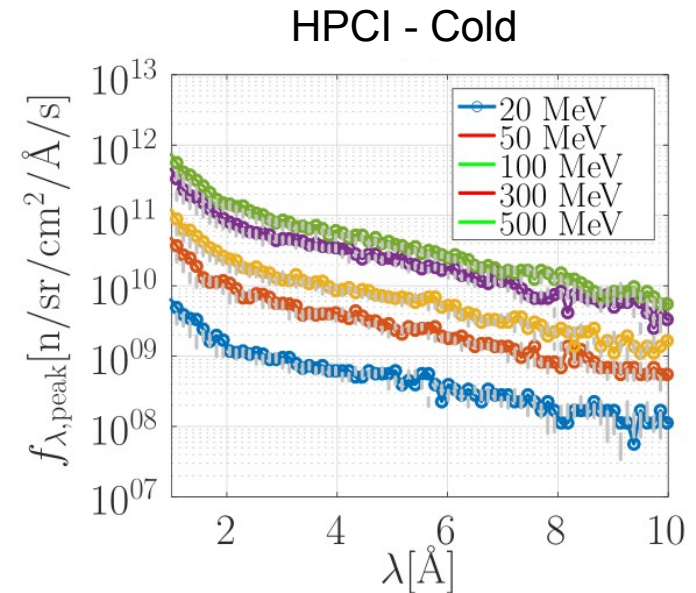
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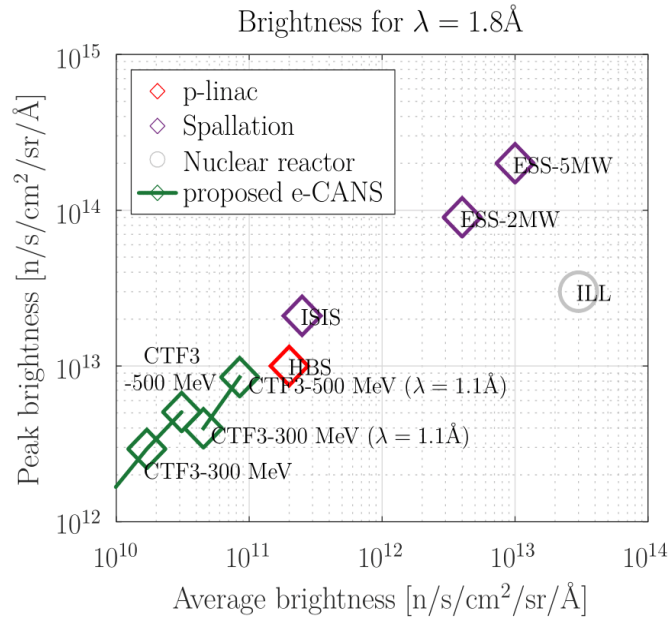
Cold moderator-target peak brightness for different electron energies for CTF3 drive beam linac.



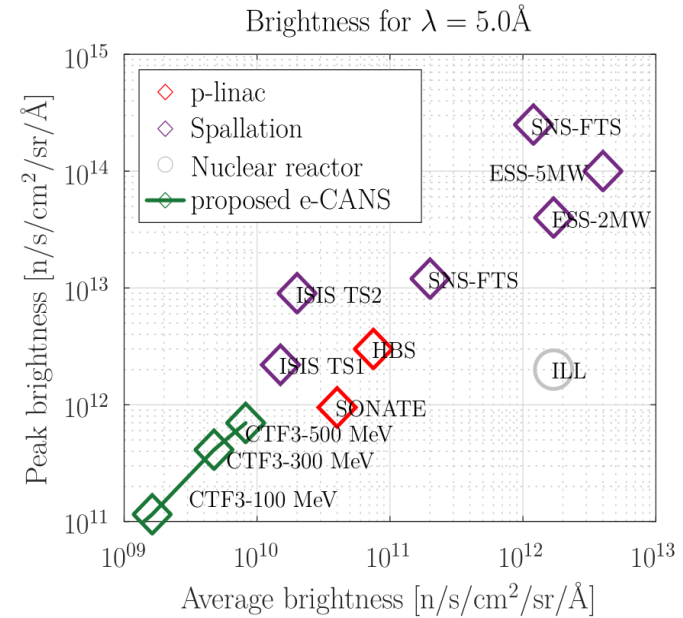
Cold moderator-target peak brightness for different electron energies for HPCI linac.

V. Brightness State-of-the-art comparison

- [7] [8] [9]



Thermal neutron brightness state-of-the-art.



Cold neutron brightness state-of-the-art.

VI. VULCAN

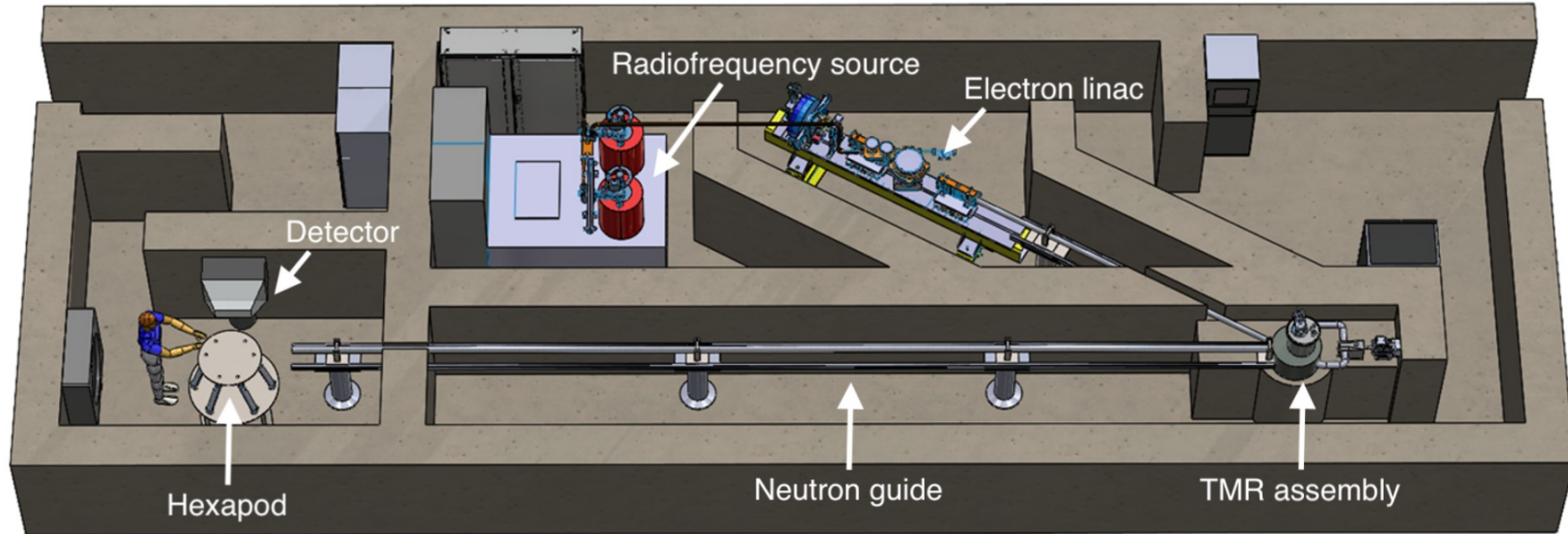
- **Commercial off-the-shelf** CANS (compact accelerator-based neutron source) [10]
- **VULCAN** = **V**ersatile **UL**tra **C**ompact **A**ccelerator-based **N**eutron source
- Collaboration between DAES SA and CERN → Industrial implementation

- Targeted applications:
 - In-situ **analysis** of **battery** and fuel cell **electrodes**
 - **Measurements** of internal **stresses** of **metallic** and **ceramic** components

VI. VULCAN

Parameter	Unit	Value
Energy	MeV	35
Energy spread	MeV	<5
Transverse size	mm	< 5
Transverse Jitter	mm	< 2
Train length	μs	<1
Train frequency	Hz	100-200

Electron beam requirements.



VI. VULCAN

REQUIREMENTS

- Beam power: > 1 kW
 - Average beam current > 29 μA
- Peak beam Current > 290 mA
- Length: < 10 m
- Cost: < 5 M€

ACCELERATOR DESIGN CHOICES

- **Thermoionic** gun
 - High average intensities
- High gradient RF cavities (**3-12 GHz**)
 - TW: Compatible with pulse compressor
- RF power source: **klystron**
 - Peak power in 5-50 MW

Conclusions

- Neutrons are produced from electron beams by **Bremsstrahlung** + **Photonuclear excitation**
- Neutron production is a trade off between beam power, cost and length. Electron-linac-based neutron sources serve as **affordable** and middle-flux options
 - Eg: VULCAN – **compact**, suited for **industrial** purposes
- Electron linacs are suitable for **multi-purpose facilities** since the unmoderated energy spectrum does not vary strongly with the initial electron energy
 - Different intensities can be achieved while keeping the same moderating scheme can be adopted for different values of $\langle E_e \rangle$.
- High-energy and high-intensity electron linacs (like CTF3 drive beam linac at 300, 500 MeV) can provide **bright neutron beams comparable** to proton-linac-based and spallation sources.

Further work

- Experimental verification of simulations- that is currently in progress at CLEAR.
- Specific target-moderator design to meet the requirements of a particular application
 - Further engineering aspects to be considered
- VULCAN: Beam dynamics and EM simulations ongoing
 - CDR in writing phase.

References

- 1) Y. Kiyanagi, “Neutron applications developing at compact accelerator-driven neutron sources,” *AAPPS Bulletin*, vol. 31, pp. 1–19, 2021.
- 2) *Compact Accelerator Based Neutron Sources*, ser. TECDOC Series. Vienna: INTERNATIONAL ATOMIC ENERGY AGENCY, 2021, no. 1981. [Online]. Available: <https://www.iaea.org/publications/14948/compact-accelerator-based-neutron-sources>
- 3) J. M. Carpenter, “The development of compact neutron sources,” *Nature Reviews Physics*, vol. 1, no. 3, pp. 177–179, 2019.
- 4) T. Roberts, “G4beamline user’s guide,” *Muons, Inc*, pp. 3468–3470, 2013.
- 5) A. Latina, V. Musat, R. Corsini, L. A. Dyks, E. Granados, A. Grudiev, S. Stapnes, P. Wang, W. Wuensch, E. Cormier, and G. Santarelli, “A compact inverse compton scattering source based on x-band technology and cavity-enhanced high average power ultrafast lasers,” in 67th *ICFA Adv. Beam Dyn. Workshop Future Light Sources Conference Proceedings*, 2023, pp. 257–260
- 6) G. Geschonke and A. Ghigo, “Ctf3 design report,” Tech. Rep., 2002.

References

- 7) J. Olivares Herrador, L.Wroe, A. Latina, et al. “Neutron production using compact linear electron accelerators”. JACoW **IPAC2024** (2024), MOPR93 doi:10.18429/JACoW-IPAC2024-MOPR93
- 8) T. Brückel, T. Gutberlet, J. Baggemann, S. Böhm, P. Doege, J. Fenske, M. Feygenson, A. Glavic, O. Holderer, S. Jaksch et al., *Conceptual Design Report-Jülich High Brilliance Neutron Source (HBS)*. Forschungszentrum Jülich GmbH, Zentralbibliothek, Verlag 2020.
- 9) F. Ott, A. Menelle, and C. Alba-Simionesco, “The sonate project, a french cans for materials sciences research,” in *EPJ Web of Conferences*, vol. 231. EDP Sciences, 2020, p. 0100.
- 10) L.M Wroe, A.Latina, J. Olivares Herrador et. al. “Accelerator design choices for a Compact Electron-Driven, pulsed neutron source”, in JACoW **LINAC2024** (2024).

Thanks for your attention



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