

# Recent development on the quasi-static PIC codes QuickPIC and QPAD

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Oct 4th, 2024

ICAP'24



# QuickPIC/QPAD Collaboration

- Weiming An, Weiyu Meng, Rong Tang, Hainan Wang, Yueran Tian, Zhihao Xu
- Qianqian Su, Viktor Decyk, Lance Hildebrand, Yujian Zhao, Thamine Dalichaouch, Warren Mori
- Fei Li





# A Brief History of QuickPIC

## FACET & FACET II



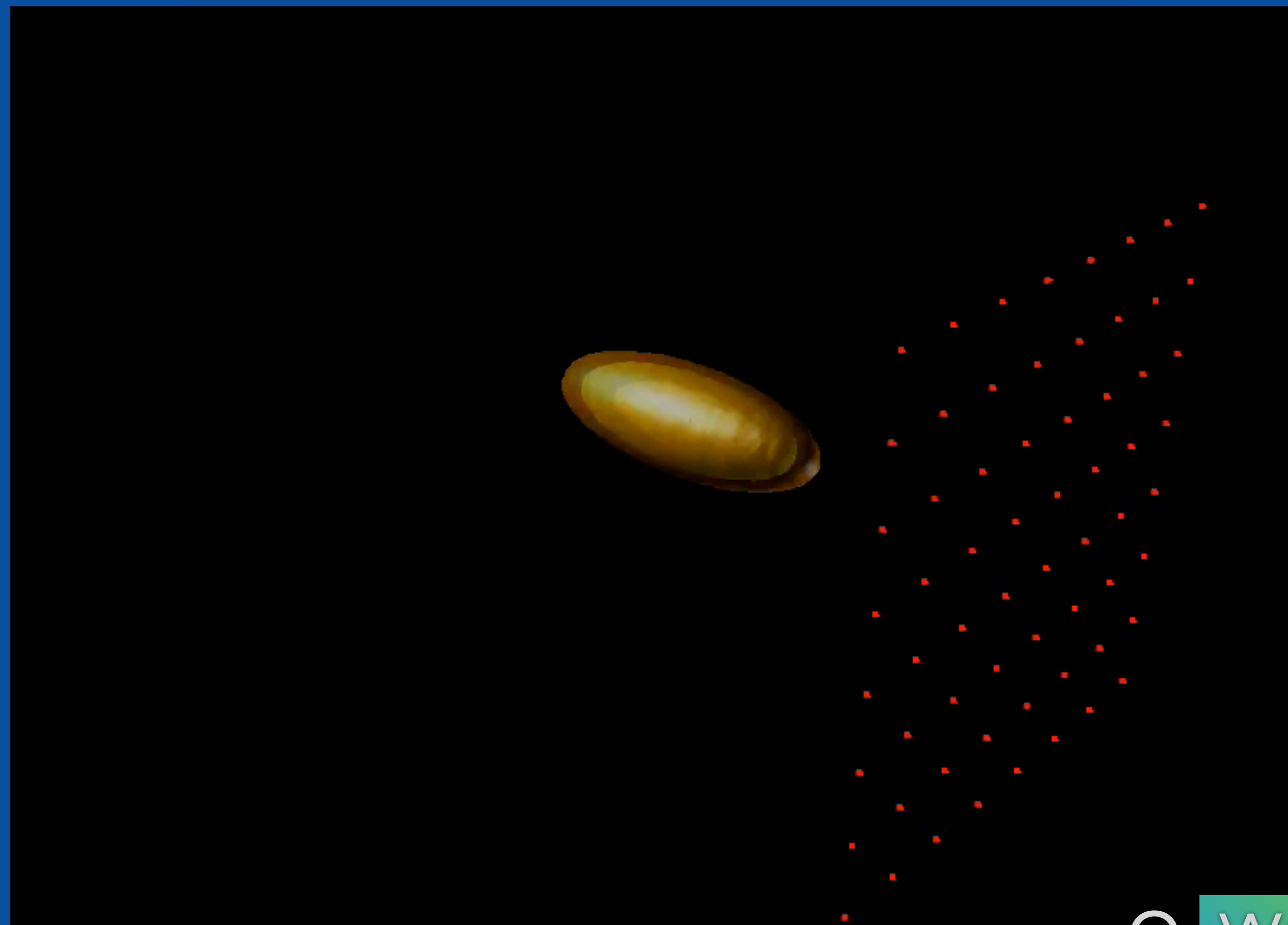
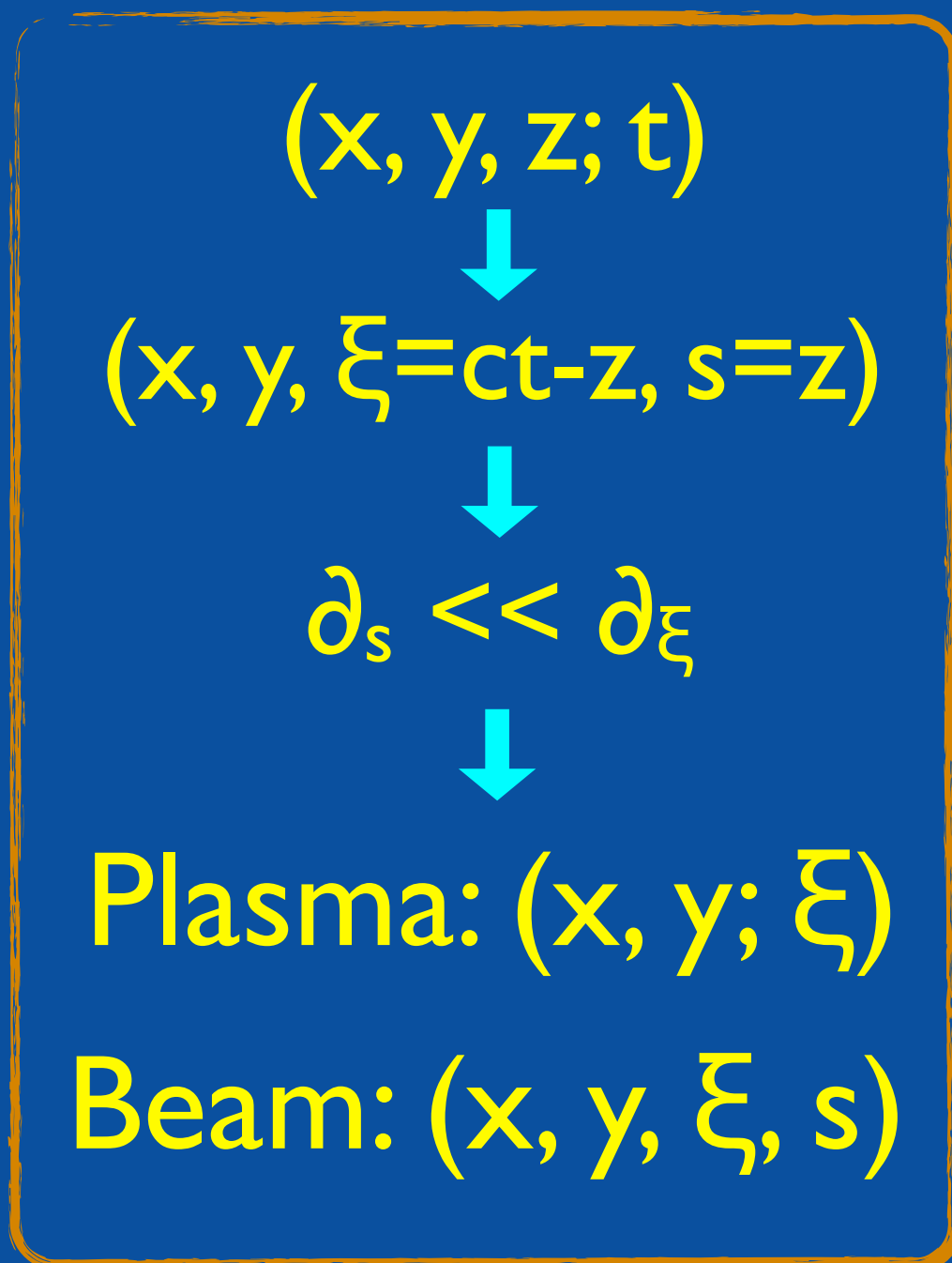
Plasma simulation has greatly impacted on PBA research.



# The Quasi-Static PIC code

Embeds a 2D PIC code inside a 3D PIC code

## Quasi-Static Approximation\*



QuickPIC is a 3D parallel Quasi-Static PIC code, which is developed based on the framework UPIC.

<https://gitee.com/bnu-plasma-astrophysics-sg/quick-pic-open-source>

W.An et. al., JCP 250, 165 (2013)

Q WAKE  
 S LCODE  
 A HIPACE



# QPAD\*: QuickPIC with Azimuthal Fourier Decomposition

- Azimuthal decomposition is a technique to speed up the simulation without much loss of accuracy. Some explicit PIC codes (e.g. Calder, quasi-3D OSIRIS, FBPIC, PLARES-PIC) have employed this algorithm to achieve 100x ~ 1000x speed-up.
- None of current quasi-static PIC codes has this feature.
- QPAD is a newly developed code based on part of the framework of open source QuickPIC.

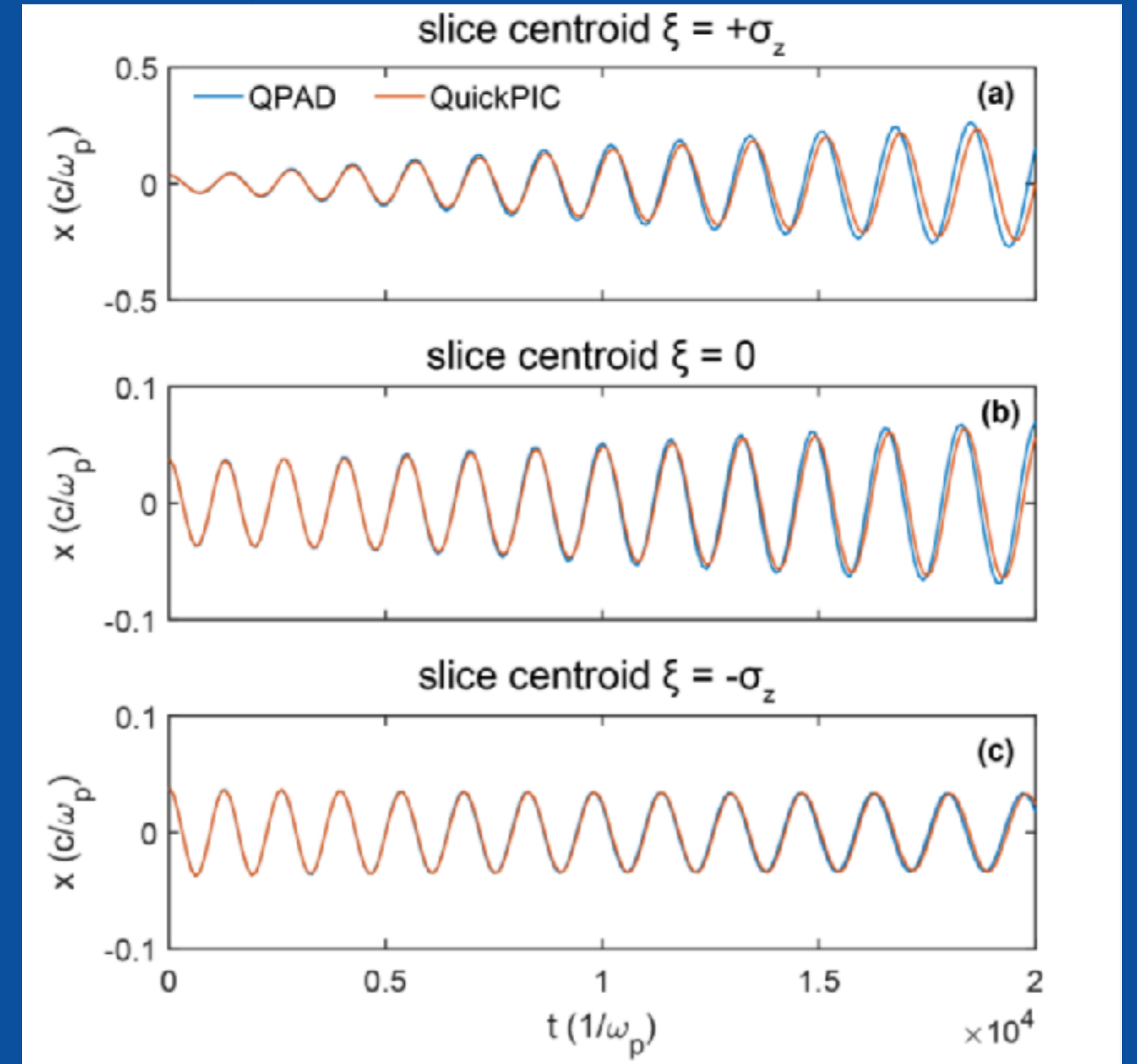
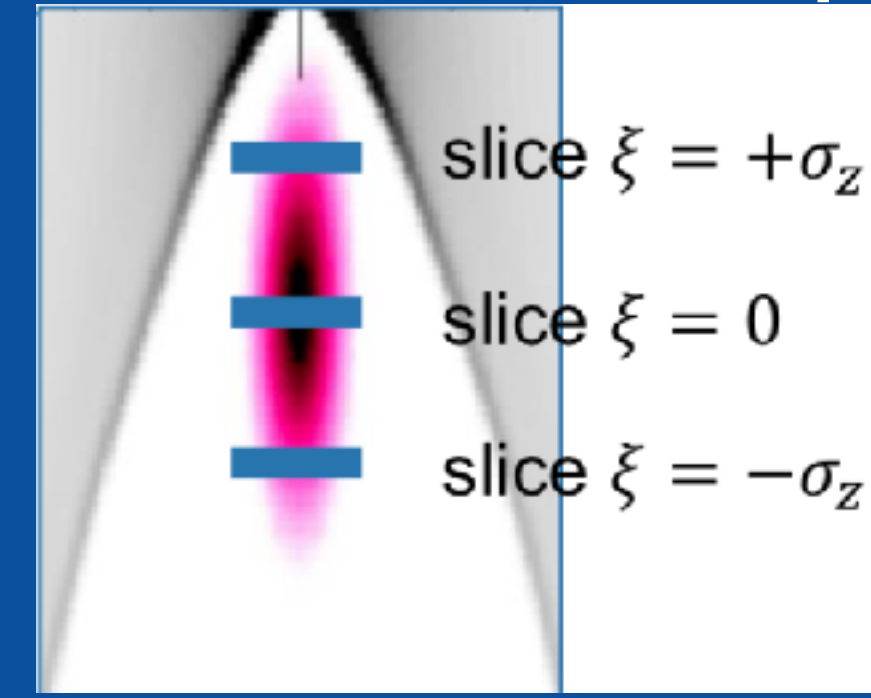
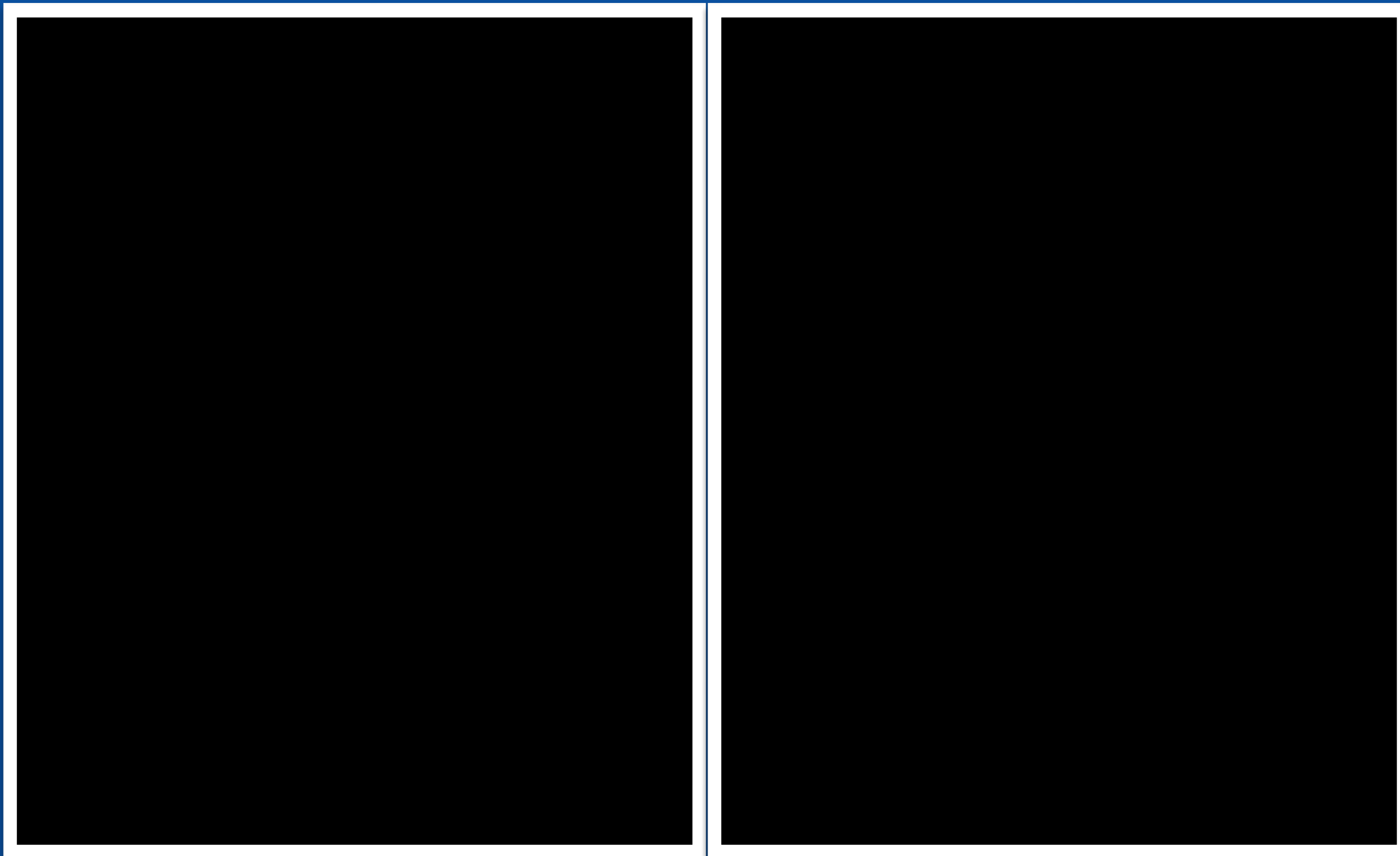
Fei Li, Weiming An\*, Viktor K. Decyk, Xinlu Xu, Mark J. Hogan, Warren B. Mori, et. al., "A quasi-static particle-in-cell algorithm based on an azimuthal Fourier decomposition for highly efficient simulations of plasma-based acceleration: QPAD", *Computer Physics Communications* 261, 107784 (2021).



# QPAD\*: QuickPIC with Azimuthal Fourier Decomposition

QuickPIC

QPAD(4 modes)



drive beam:  $k_p \sigma_r = 0.14, k_p \sigma_z = 0.48, \Lambda = 1.8, \gamma = 20000$   
trailing beam:  $k_p \sigma_r = 0.14, k_p \sigma_z = 0.24, \Lambda = 1.1, \gamma = 20000, k_p \Delta_{\text{off}} = 0.0375$



# Laser Module in QuickPIC

\*Weiyu Meng et. al., in preparation.

Laser envelope equation:

$$2 \frac{\partial}{\partial s} \left( -ik_0 + \frac{\partial}{\partial \xi} \right) \hat{\mathbf{a}} - \nabla_{\perp}^2 \hat{\mathbf{a}} = \chi \hat{\mathbf{a}}$$

Finite difference:

$$\left[ -ik_0 - \frac{1}{4} \Delta_s (\nabla_{\perp}^2 + \chi^n) \right] a^{n+\frac{1}{2}} + \frac{\partial a^{n+\frac{1}{2}}}{\partial \xi} = \left[ -ik_0 + \frac{1}{4} \Delta_s (\nabla_{\perp}^2 + \chi^n) \right] a^{n-\frac{1}{2}} + \frac{\partial a^{n-\frac{1}{2}}}{\partial \xi}.$$

$$\left( -\frac{1}{4} \Delta_s \nabla_{\perp}^2 + \frac{3}{2\Delta_{\xi}} - \frac{1}{4} \Delta_s \chi_j \right) a_{R,j} + k_0 a_{I,j} = \frac{1}{2\Delta_{\xi}} (4a_{R,j-1} - a_{R,j-2}) + S_{R,j},$$

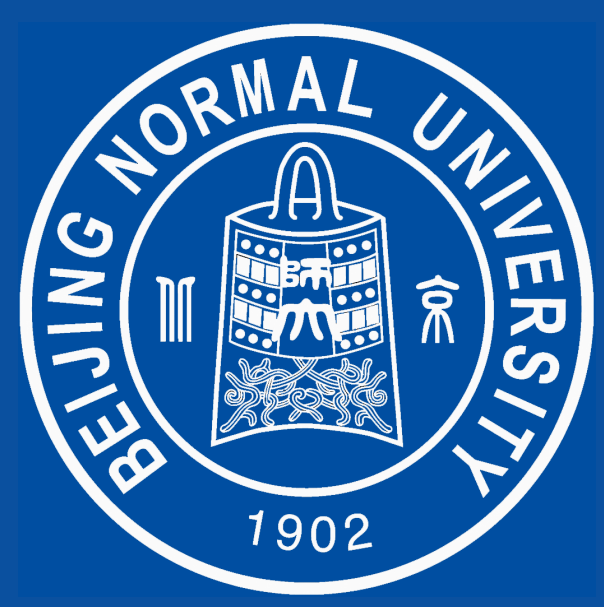
$$\left( -\frac{1}{4} \Delta_s \nabla_{\perp}^2 + \frac{3}{2\Delta_{\xi}} - \frac{1}{4} \Delta_s \chi_j \right) a_{I,j} - k_0 a_{R,j} = \frac{1}{2\Delta_{\xi}} (4a_{I,j-1} - a_{I,j-2}) + S_{I,j}$$

where:

$$S_{R,j} = \frac{1}{4} \Delta_s (\nabla_{\perp}^2 a_{R,j} + \chi_j a_{R,j}) + k_0 a_{I,j} + D_{\xi} a_{R,j}$$

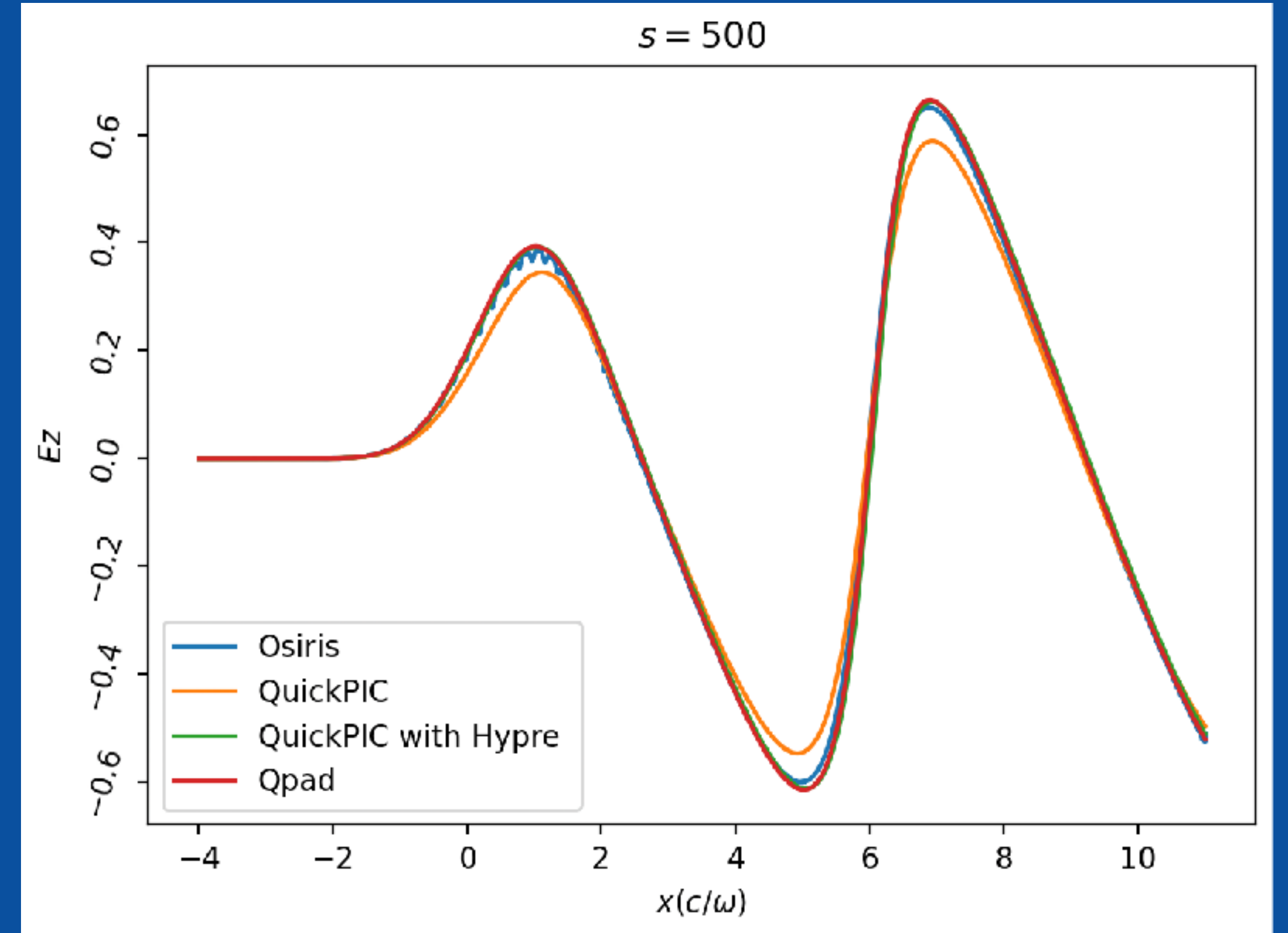
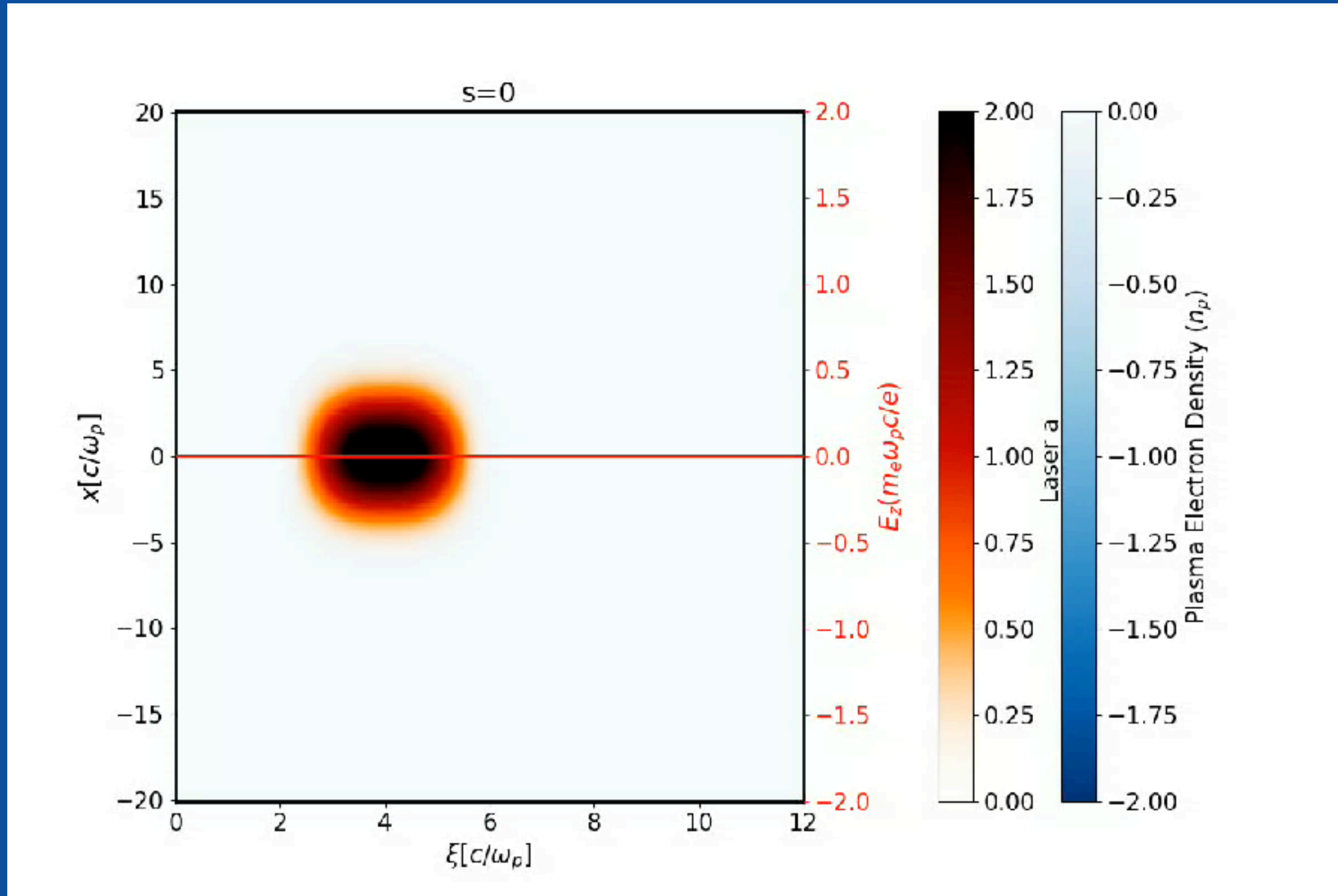
$$S_{I,j} = \frac{1}{4} \Delta_s (\nabla_{\perp}^2 a_{I,j} + \chi_j a_{I,j}) - k_0 a_{R,j} + D_{\xi} a_{I,j},$$

can be solved with Hypre\*



# Laser Module in QuickPIC

\*Weiyu Meng et. al., in preparation.



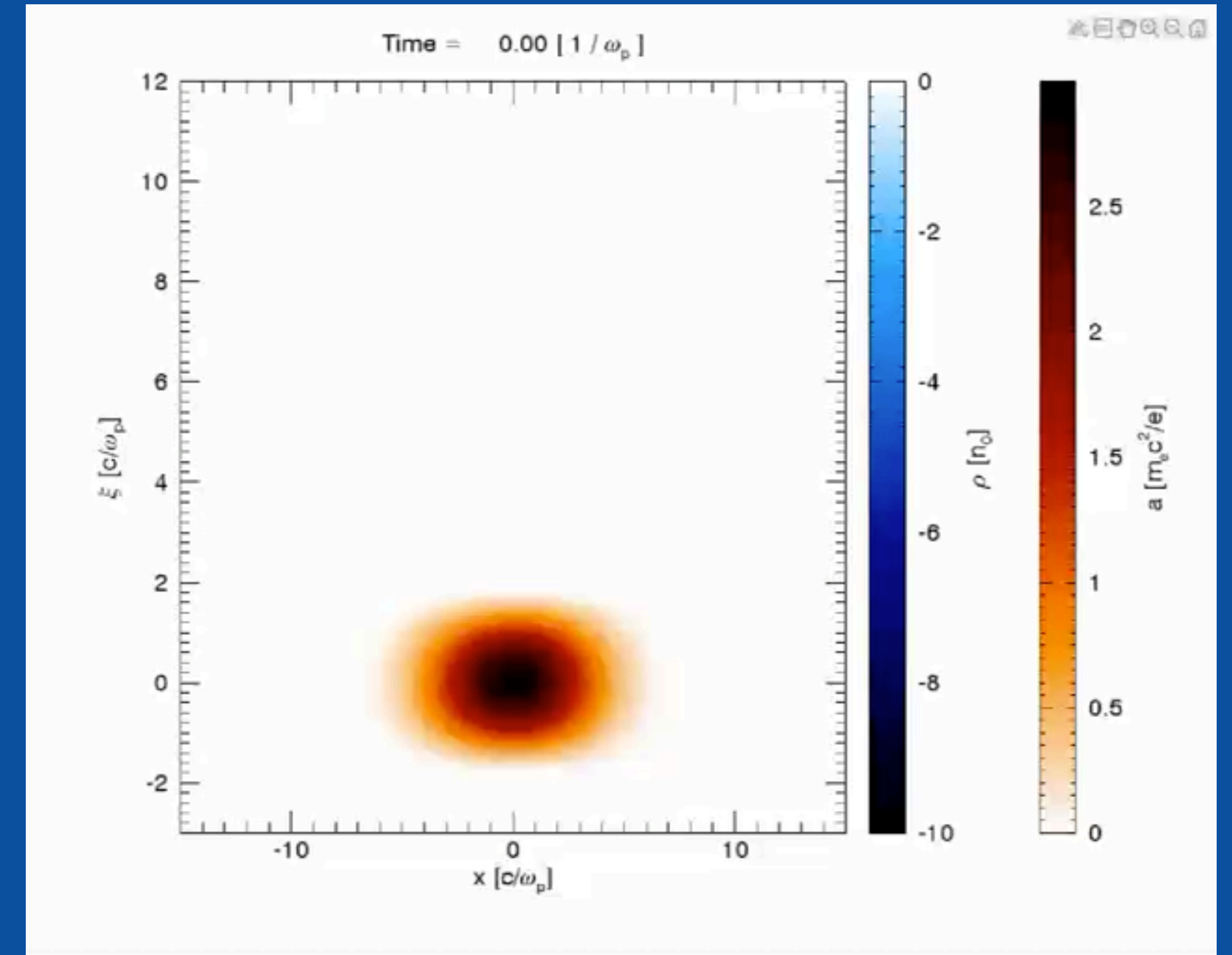
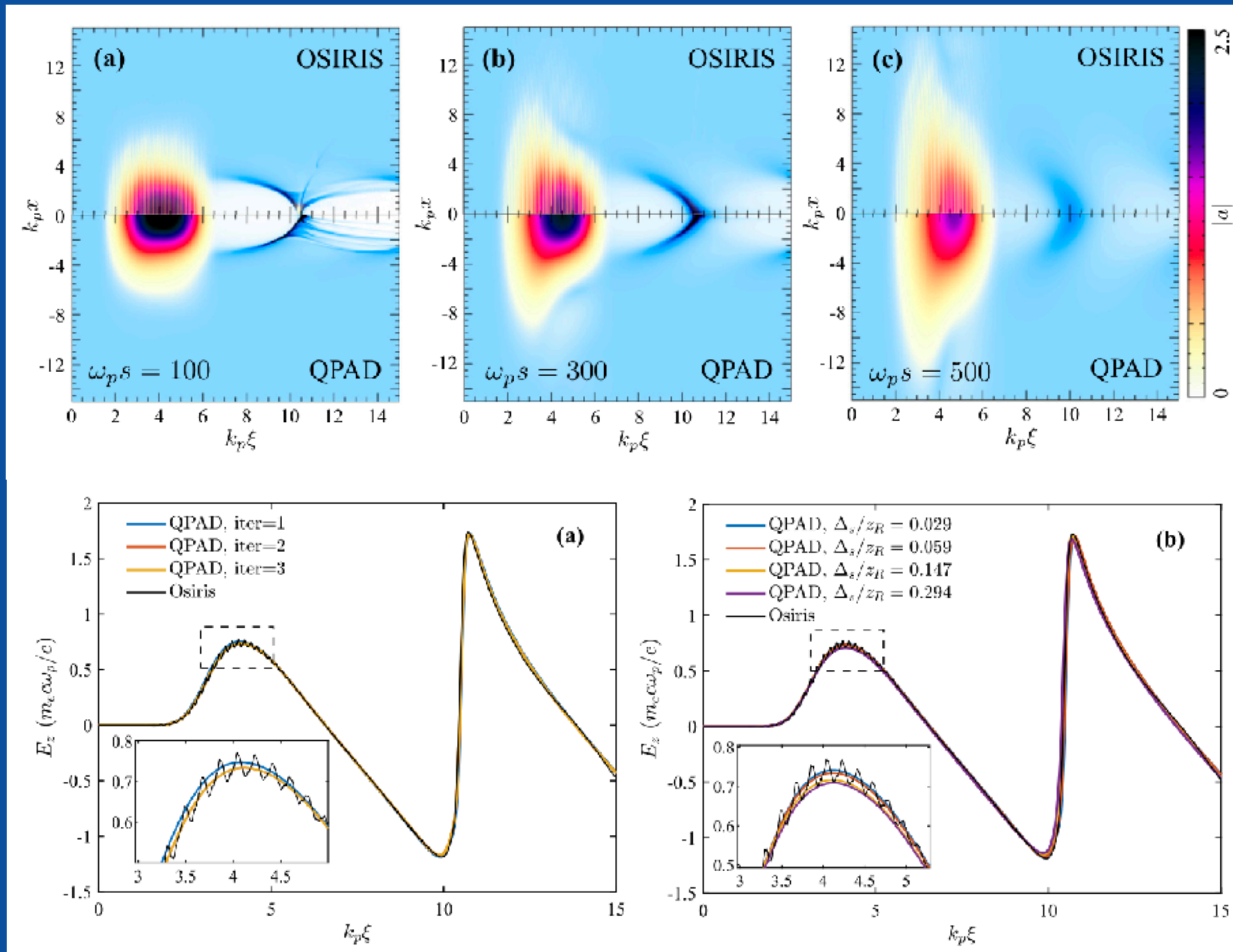
QuickPIC vs Osiris vs Qpad





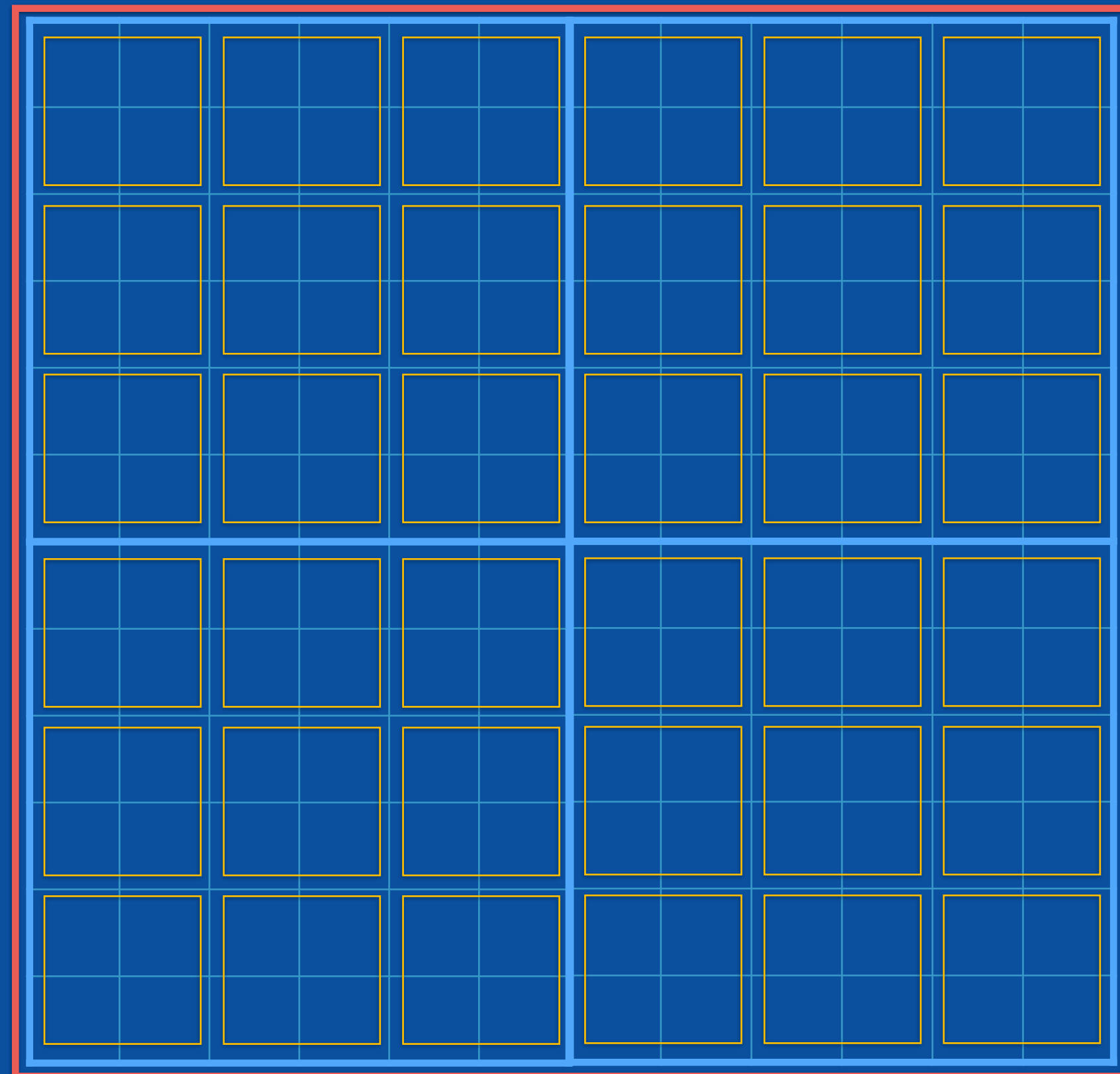
# Laser Module in QPAD

\*Fei Li, Weiming An, Frank S. Tsung, Viktor K. Decyk, Warren B. Mori, JCP 470 111599 (2022).





# QuickPIC on GPU (GPU+MPI)



Simulation Box

MPI + OpenMP

MPI + GPU

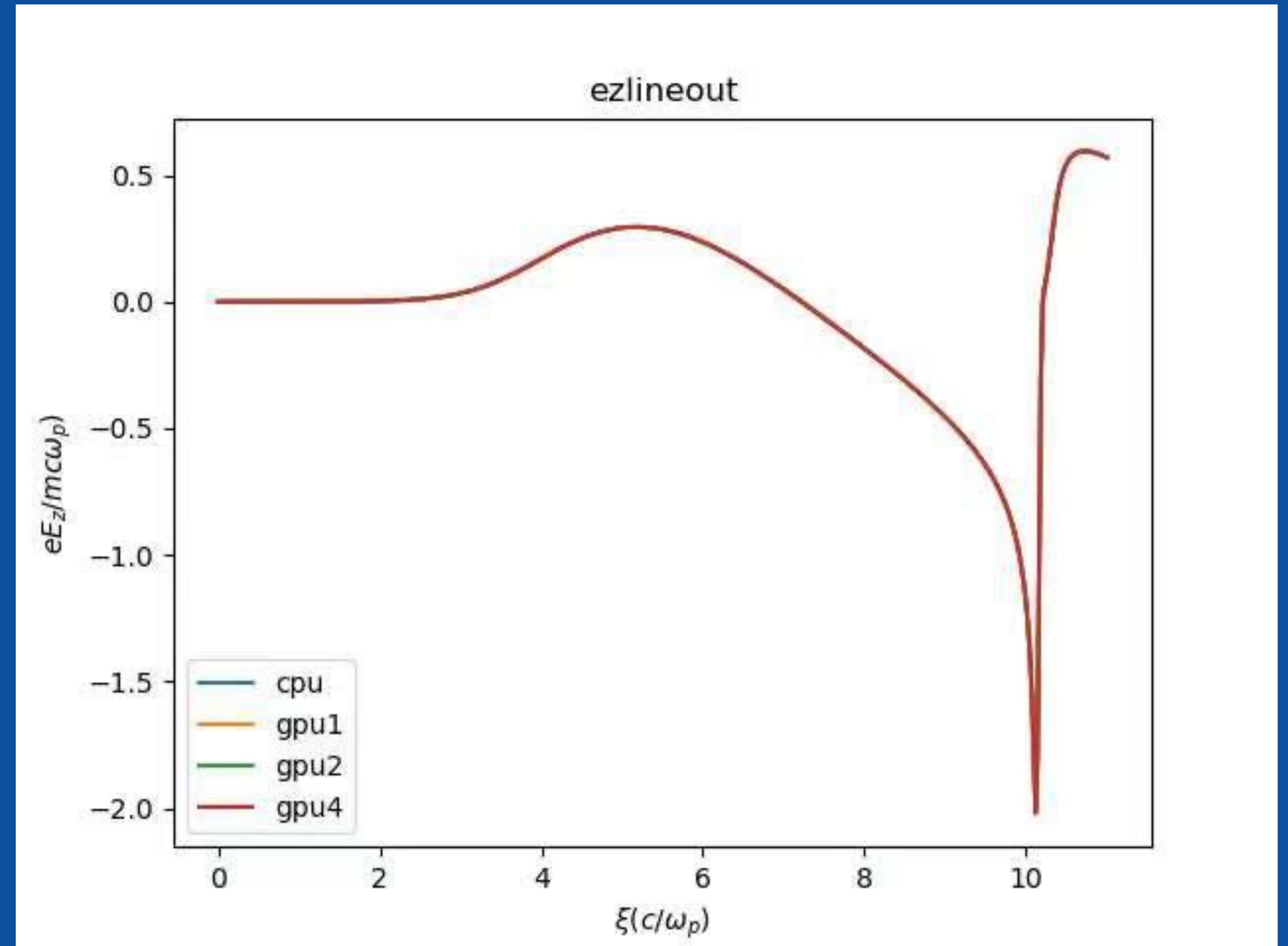
Particle Manager

FFT



# QuickPIC on GPU (GPU+MPI)

- 2D particle and field solver functions completely on GPU
- using CUDA-aware MPI, GPU+MPI parallel simulation





# The explicit solver in QuickPIC/QPAD

## Iterative Solution

$$\nabla_{\perp}^2 \vec{B}_{\perp}^{\text{new}} - \vec{B}_{\perp}^{\text{new}} = \hat{z} \times \left( \frac{\partial}{\partial \xi} \vec{J}_{\perp} + \nabla_{\perp} J_z \right) - \vec{B}_{\perp}^{\text{old}}$$

## Explicit Solution\*

$$\begin{aligned} & \nabla_{\perp}^2 \vec{B}_{\text{plasma},\perp} - \frac{q_e(\rho - J_z)_e}{m_e(1 - \frac{q_e}{m_e}\psi_e)} \vec{B}_{e,\perp} S(\vec{x}_{\perp} - \vec{x}_{i\perp}) - \frac{q_i(\rho - J_z)_i}{m_i(1 - \frac{q_i}{m_i}\psi_i)} \vec{B}_{i,\perp} S(\vec{x}_{\perp} - \vec{x}_{i\perp}) \\ & = \hat{z} \times \left\{ \frac{q}{\text{Volume}} \sum_s \frac{q_s}{m(1 - \frac{q_s}{m_s}\psi_s)} \left( \frac{-\gamma_s \nabla_{\perp} \psi_s}{1 - \frac{q_s}{m_s}\psi_s} - \hat{z} \times \vec{B}_{\text{beam},\perp} + \frac{\vec{p}_{s\perp} \times B_{sz} \hat{z}}{1 - \frac{q_s}{m_s}\psi_s} + \frac{\vec{p}_{s\perp} (E_{sz} + \nabla_{\perp} \psi_s \cdot \frac{\vec{p}_{s\perp}}{1 - \frac{q_s}{m_s}\psi_s})}{1 - \frac{q_s}{m_s}\psi_s} \right) S(\vec{x}_{\perp} - \vec{x}_{s\perp}) \right\} \\ & - \hat{z} \times \left[ \frac{q}{\text{Volume}} \nabla_{\perp} \cdot \sum_s \frac{\vec{p}_{s\perp} \vec{p}_{s\perp}}{\left(1 - \frac{q_s}{m_s}\psi_s\right)^2} S(\vec{x}_{\perp} - \vec{x}_{s\perp}) \right] + \hat{z} \times \nabla_{\perp} J_z \end{aligned}$$

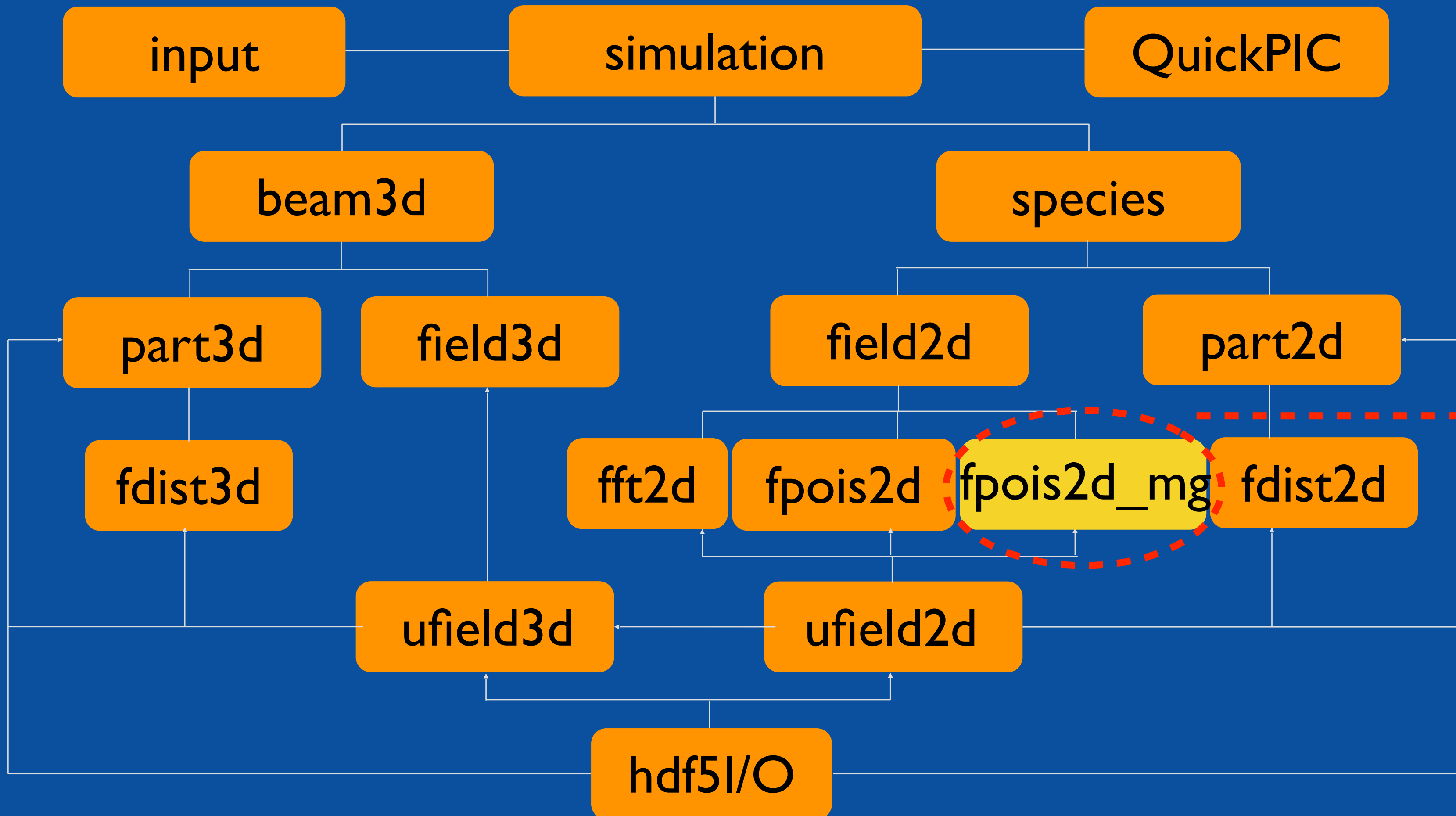
**solved with finite difference method**

	$\xi-1/2\Delta\xi$	$\xi$	$\xi+1/2\Delta\xi$	$\xi+\Delta\xi$	$\xi+3/2\Delta\xi$
Known quantities	$\psi, \vec{J}, \vec{B}, \vec{E}$	$\vec{p}, \gamma$	$\vec{x}_{\perp}$		
Quantities calculated before the iteration			$\rho - J_z$		
			$\psi$		
			$\nabla_{\perp} \psi$		
Quantities predicted or corrected			$E_z, B_z$		
				$\vec{p}_{\perp}$	
			$\vec{p}_{\perp}, \vec{J}, p_z, \frac{\partial \vec{p}_{\perp}}{\partial \xi}, \frac{\partial \vec{J}_{\perp}}{\partial \xi}, \gamma$		
			$\vec{B}_{\perp}$		
Quantities known after the iteration			$E_z, B_z$		
			$\vec{E}_{\perp}$	$\vec{p}, \gamma$	$\vec{x}_{\perp}$

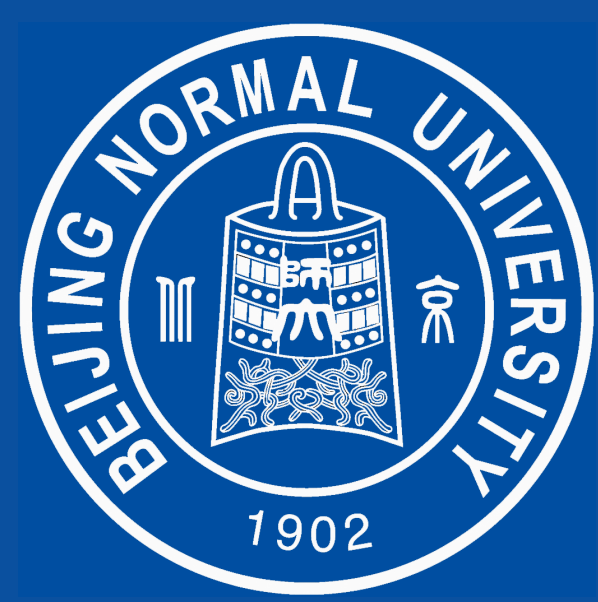
\*Wang, Tianhong, et al., Physical Review Accelerators and Beams 25.10 (2022): 104603.



# The explicit solver in QuickPIC/QPAD

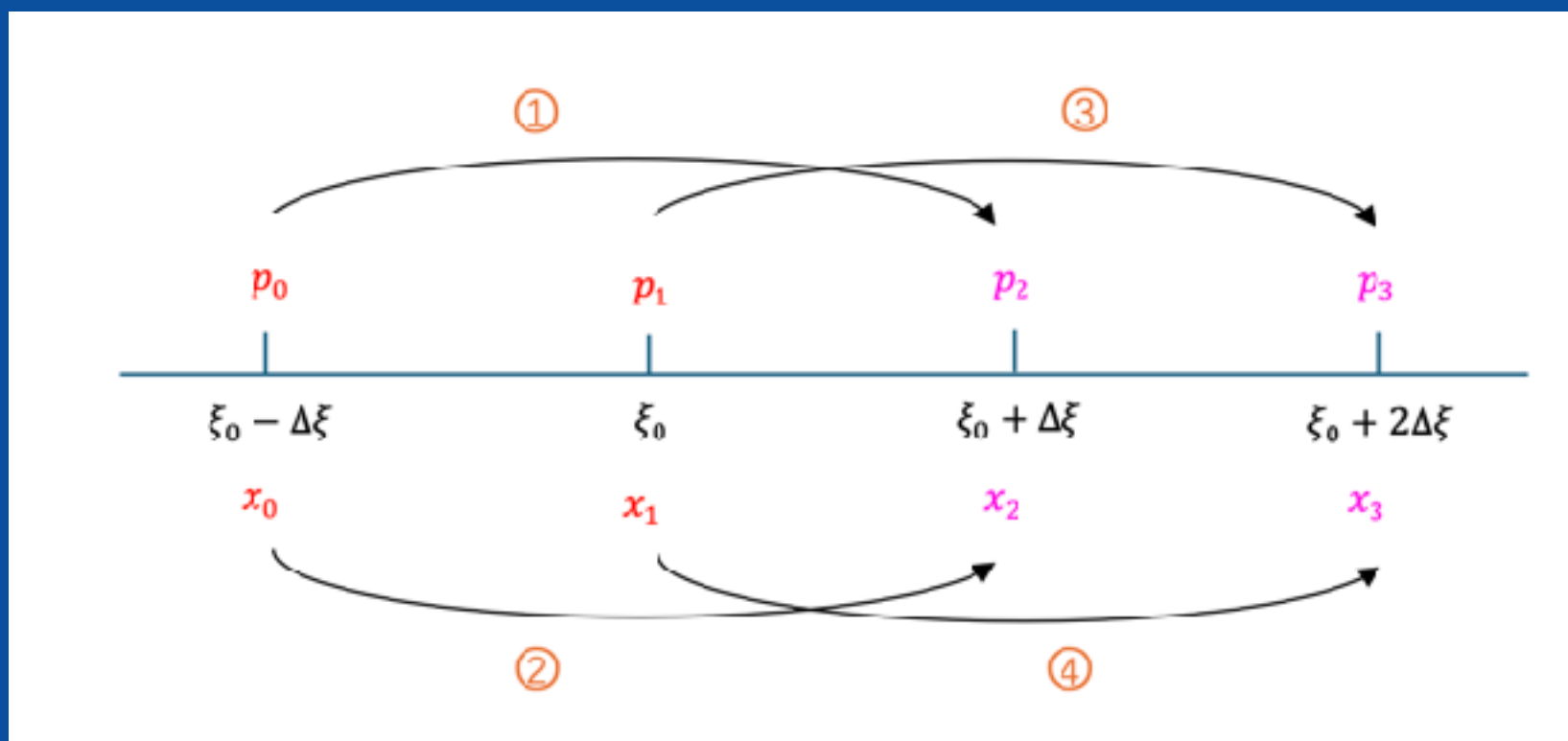


This module is used to solve Poisson equations on a two-dimensional grid. With the HYPRE library, it provides a set of interfaces for setting and solving Poisson or modified Helmholtz equations with different boundary conditions.



# The explicit solver in QuickPIC/QPAD

## algorithm design



$$\vec{x}_0 \longrightarrow \vec{x}_2, \quad \vec{p}_1 \longrightarrow \vec{p}_3$$

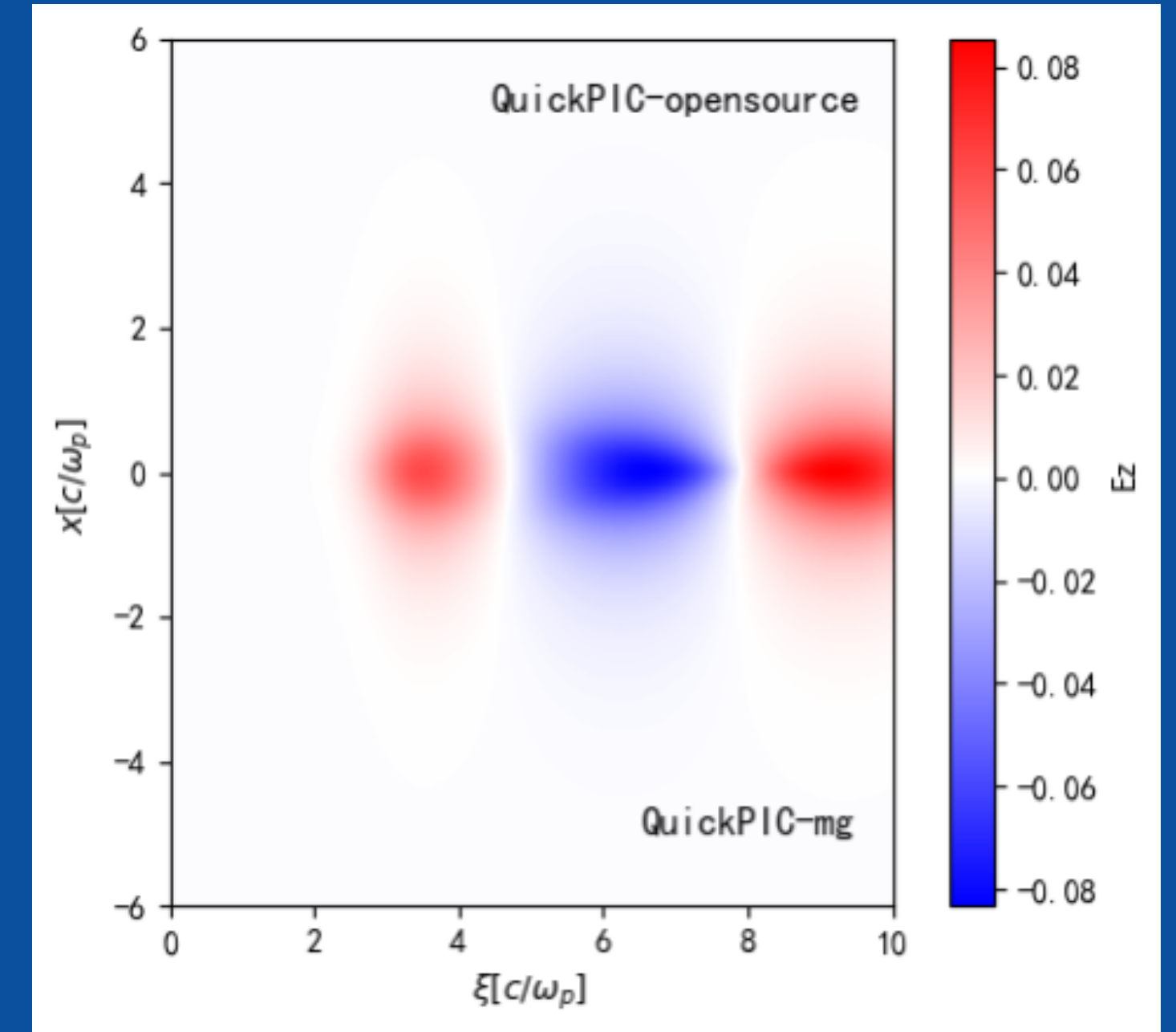
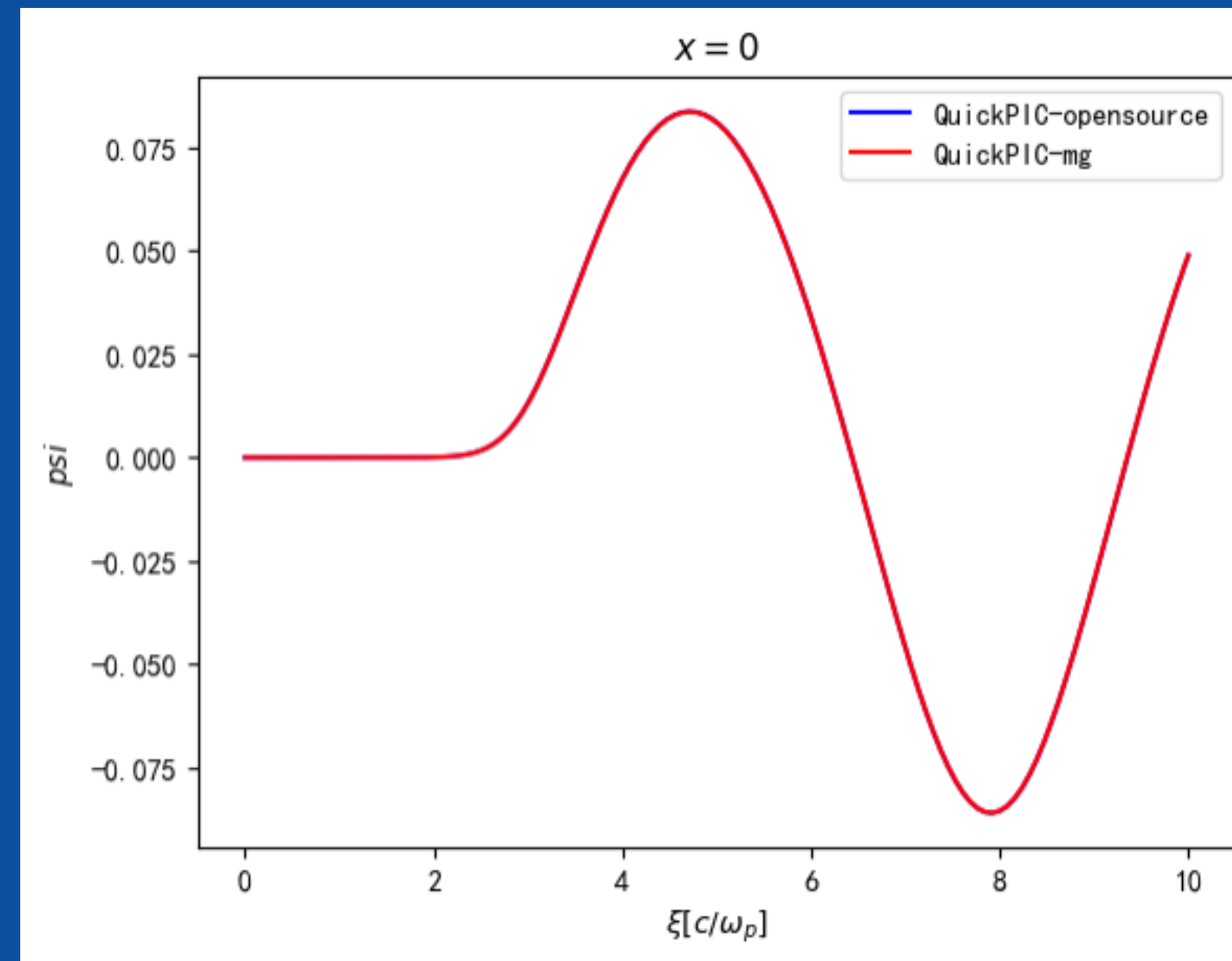
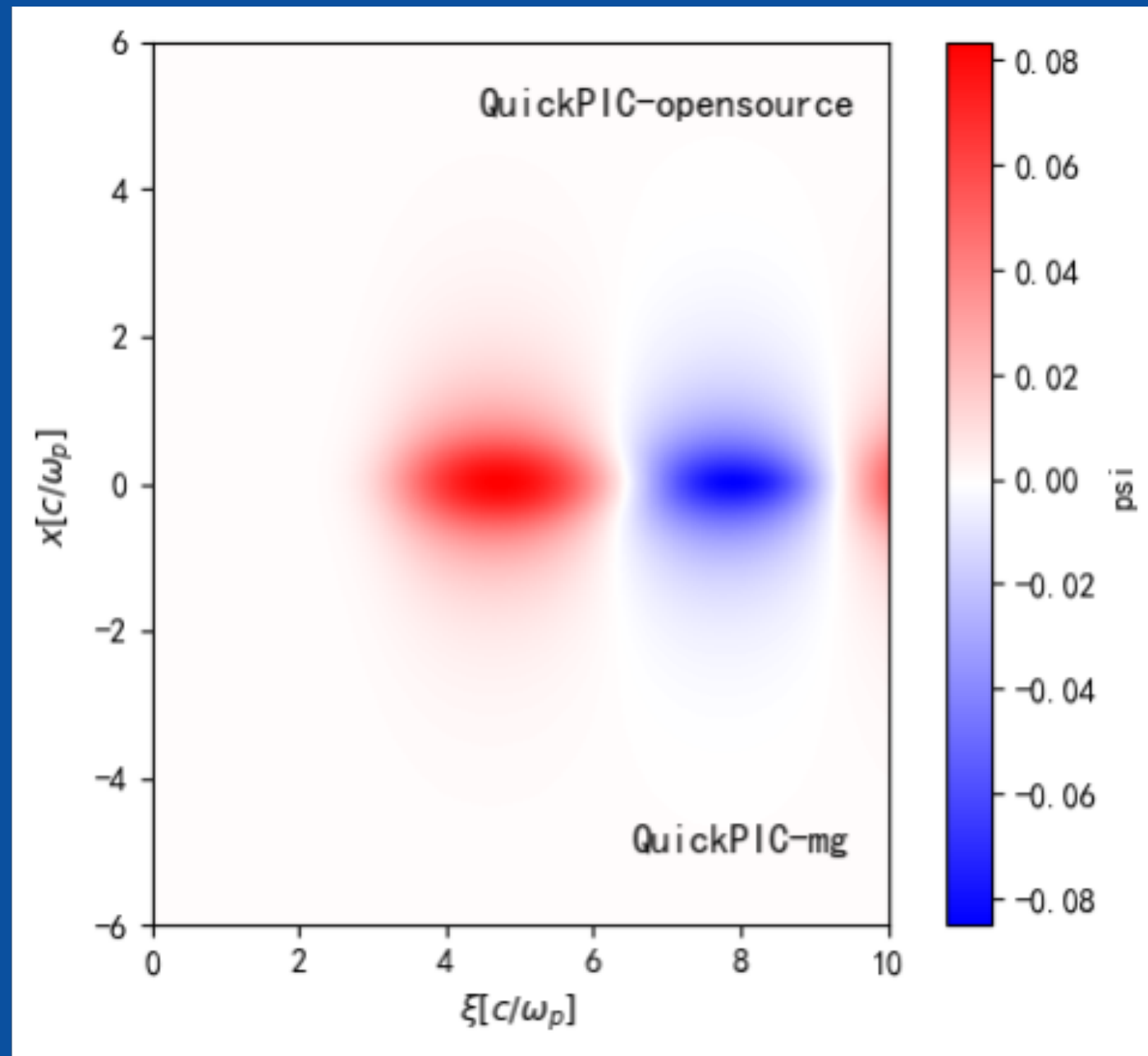
$$\vec{x}_1 \longrightarrow \vec{x}_3, \quad \vec{p}_2 \longrightarrow \vec{p}_4$$

.....

	$\xi - \Delta\xi$	$\xi$	$\xi + \Delta\xi$	$\xi + 2\Delta\xi$
Known quantities	$\vec{p}, \vec{x}$	$\vec{p}, \vec{x}$		
		$\rho - J_z$		
		$\psi$		
		$\nabla_{\perp}\psi$		
		$\vec{J}, B_z, E_z$		
		$\vec{B}_{\perp}, \vec{E}_{\perp}$		
Quantities calculated			$\vec{p}, \vec{x}$	
			$\rho - J_z$	
			$\psi$	
			$\nabla_{\perp}\psi$	
			$\vec{J}, B_z, E_z$	
			$\vec{B}_{\perp}, \vec{E}_{\perp}$	
				$\vec{p}, \vec{x}$



# The explicit solver in QuickPIC/QPAD





# The explicit solver in QuickPIC/QPAD

- In QPAD, we need to perform a Fourier expansion of the field quantities.

$$\frac{\frac{q}{m}(\rho - J_z)}{1 - \frac{q}{m}\psi} B_{\perp} = \frac{\sum_{n_1=-\infty}^{+\infty} a_{n_1} e^{-in_1\phi}}{1 + \sum_{n_2=-\infty}^{+\infty} b_{n_2} e^{-in_2\phi}} B_{\perp} = \sum_{n_3=-\infty}^{+\infty} u_{n_3} e^{-in_3\phi} \sum_{n_4=-\infty}^{+\infty} B_{\perp, n_4} e^{-in_4\phi} = \sum_{n_5=-\infty}^{+\infty} X_{\perp, n_5} e^{-in_5\phi}$$

- The transverse magnetic field equation for each mode :

$$\left( \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} - \frac{(k \pm 1)^2}{r^2} \right) B_{\pm, k} + X_k = S_{\pm, k}, \quad k = [-m, m], \quad X_k = \sum_{n=k-m} B_n U_{k-n}$$

We use Hypre[1,2] to solve these linear equations.

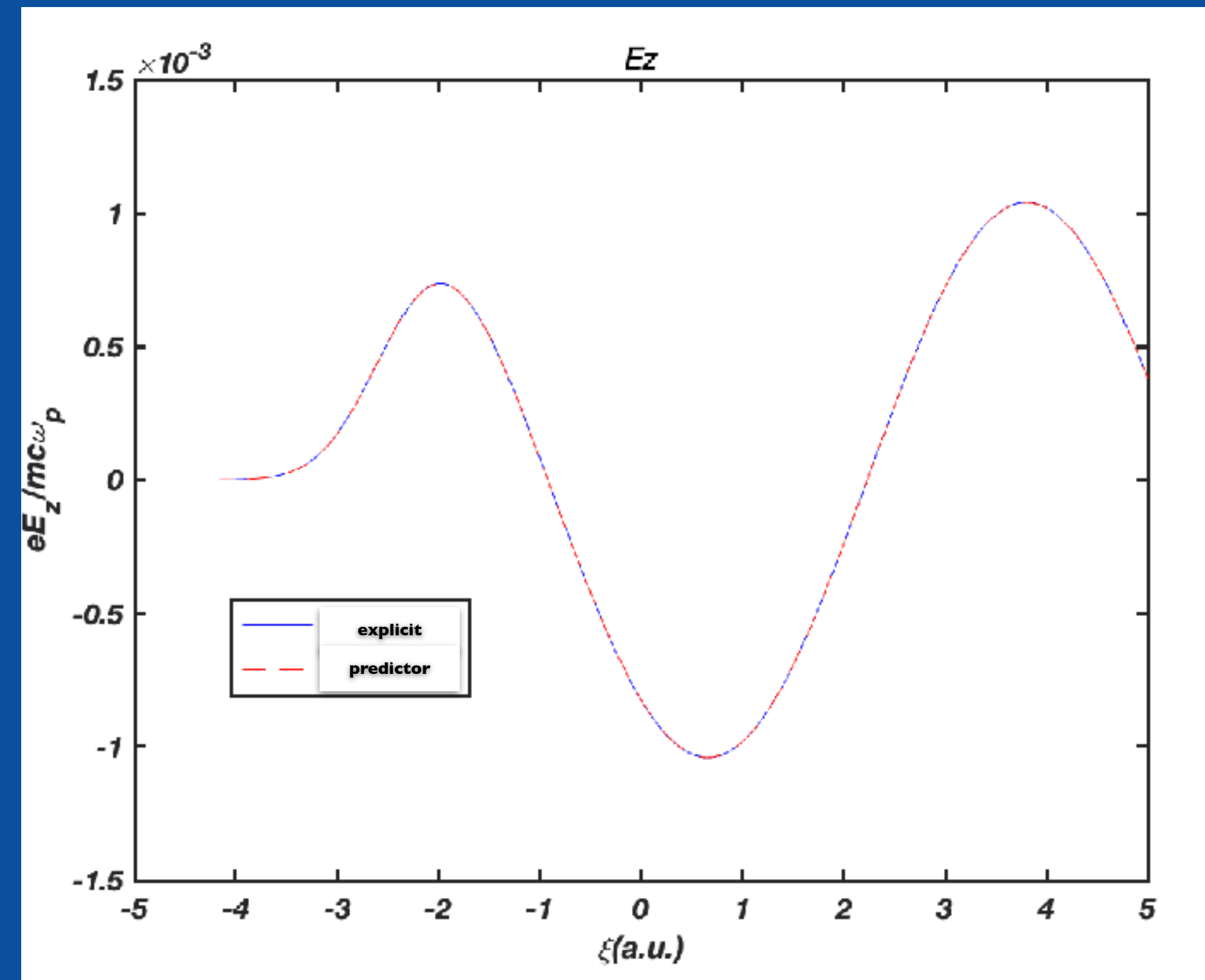
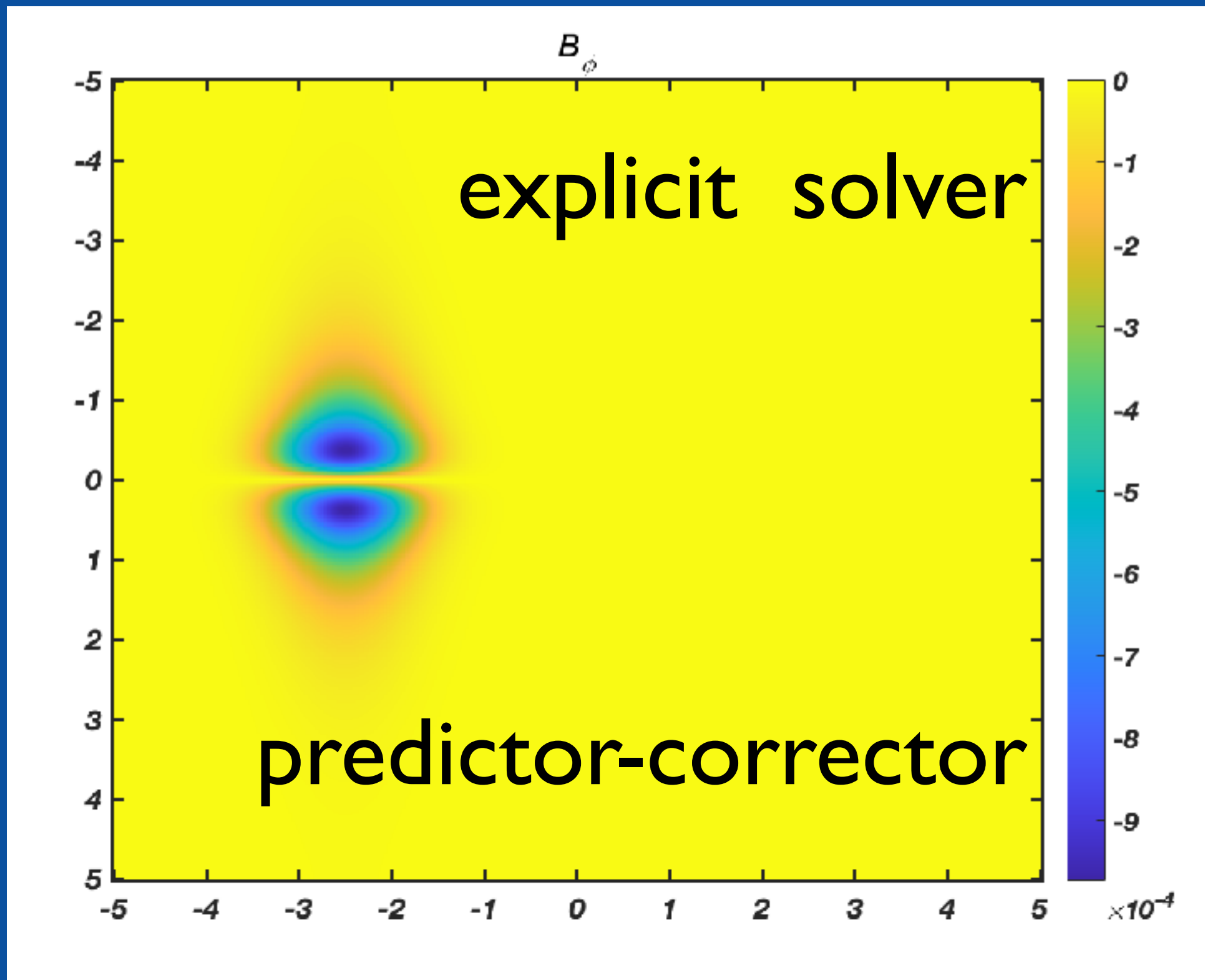
[1] <https://github.com/hypre-space/hypre/tags>

[2] Falgout R D, Jones J E, Yang U M. Springer Berlin Heidelberg, 2006: 267-294.



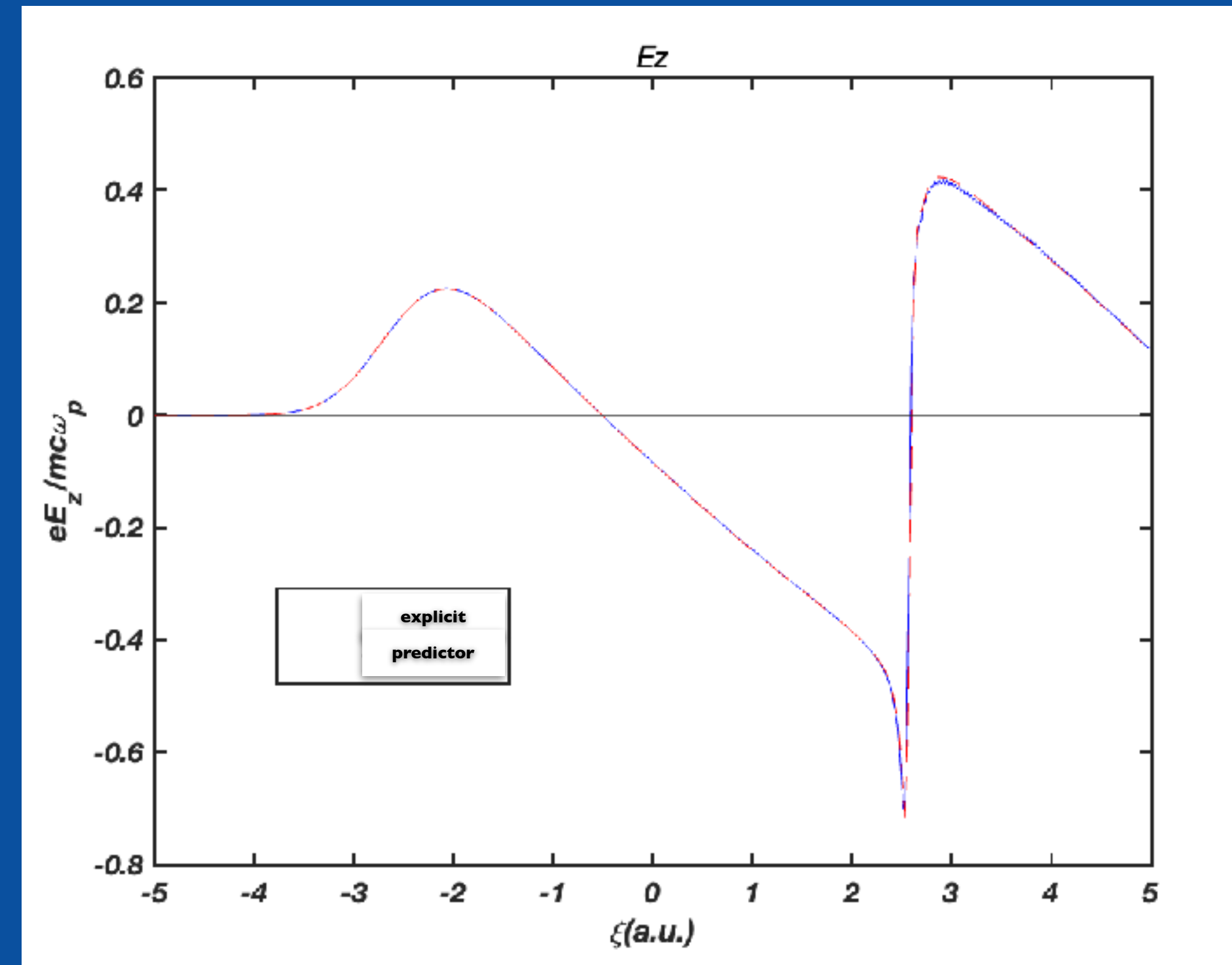
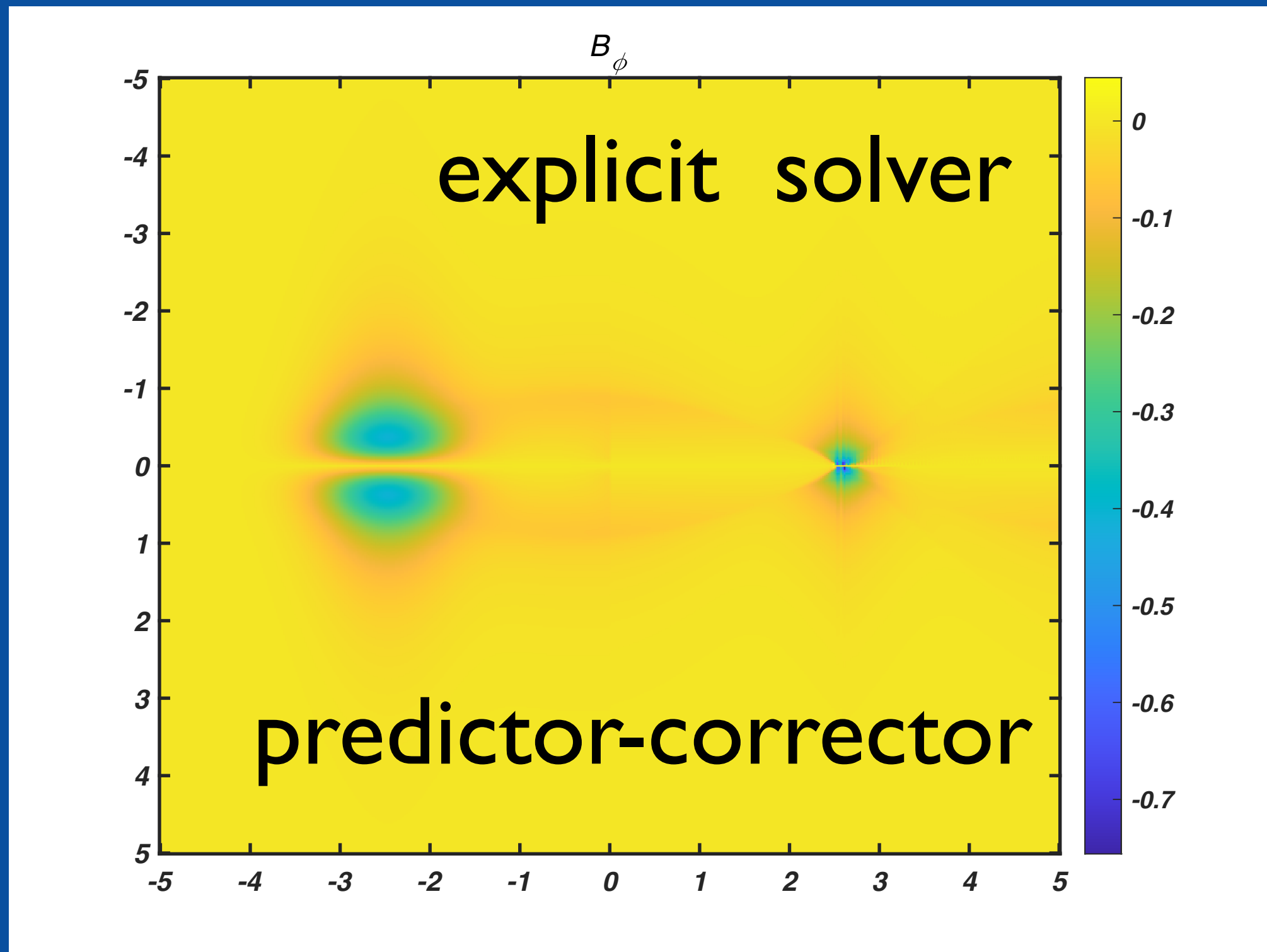


# The explicit solver in QuickPIC/QPAD





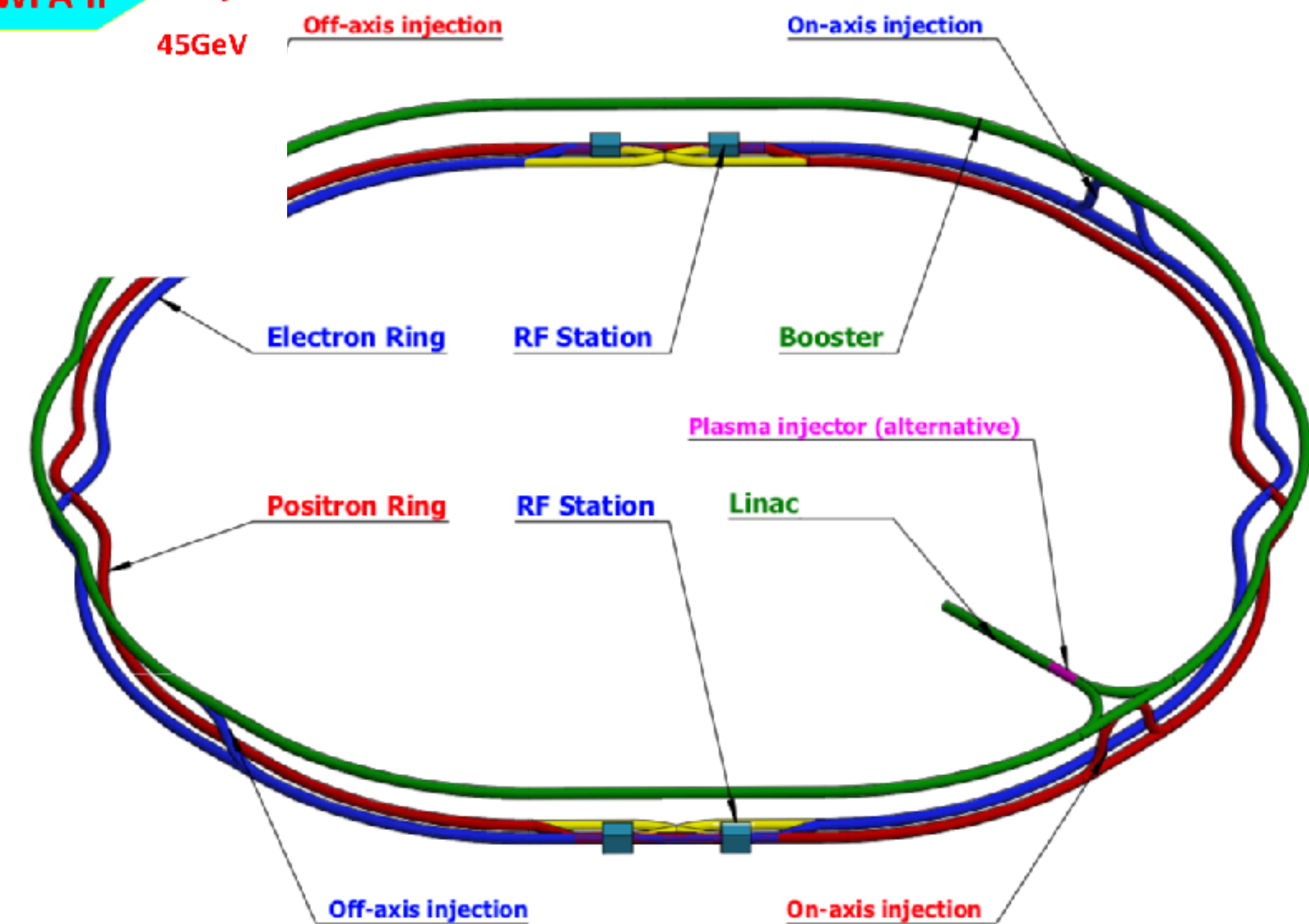
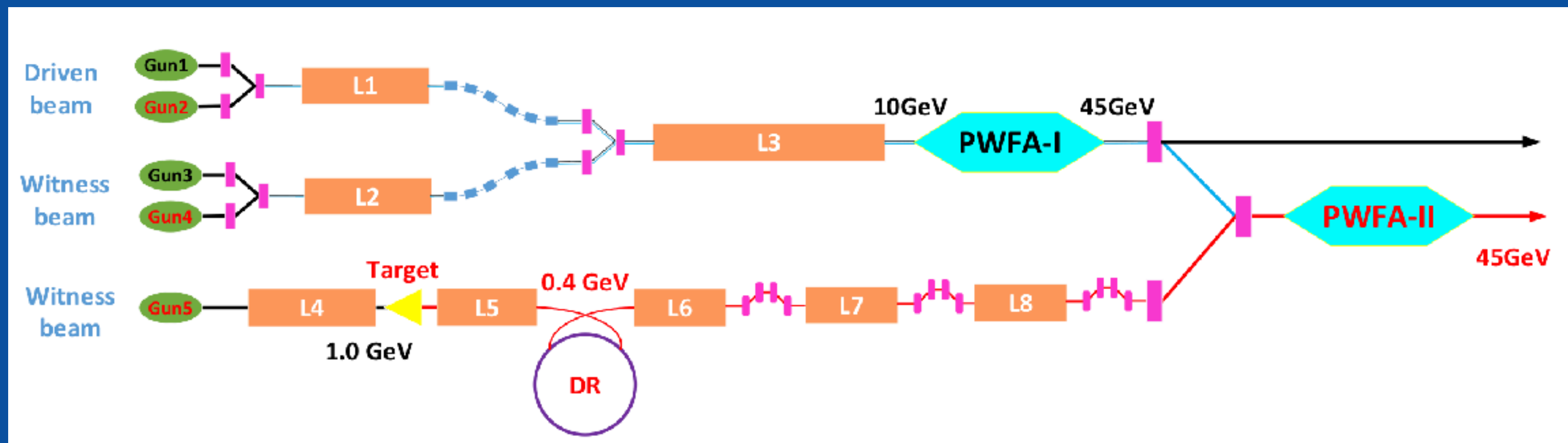
# The explicit solver in QuickPIC/QPAD





# QuickPIC/QPAD @ CEPC

## Design CEPC Plasma Injector



IHEP

THU

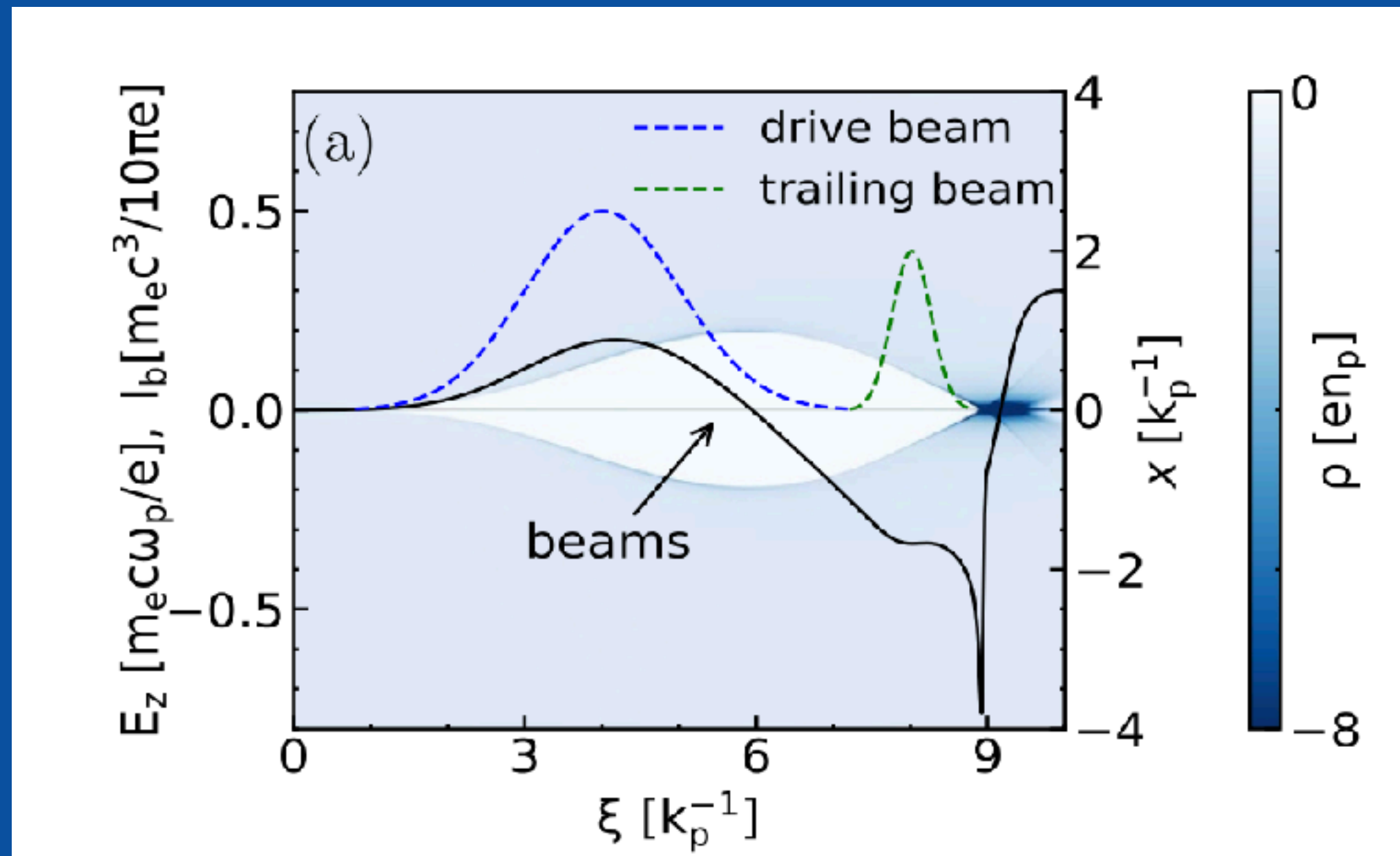
BNU

Fig. 1. CEPC layout. The blue one corresponds to the electron ring and the red one corresponds to the positron ring. There are two interaction regions, two RF regions and four injection regions in the collider.<sup>1</sup>



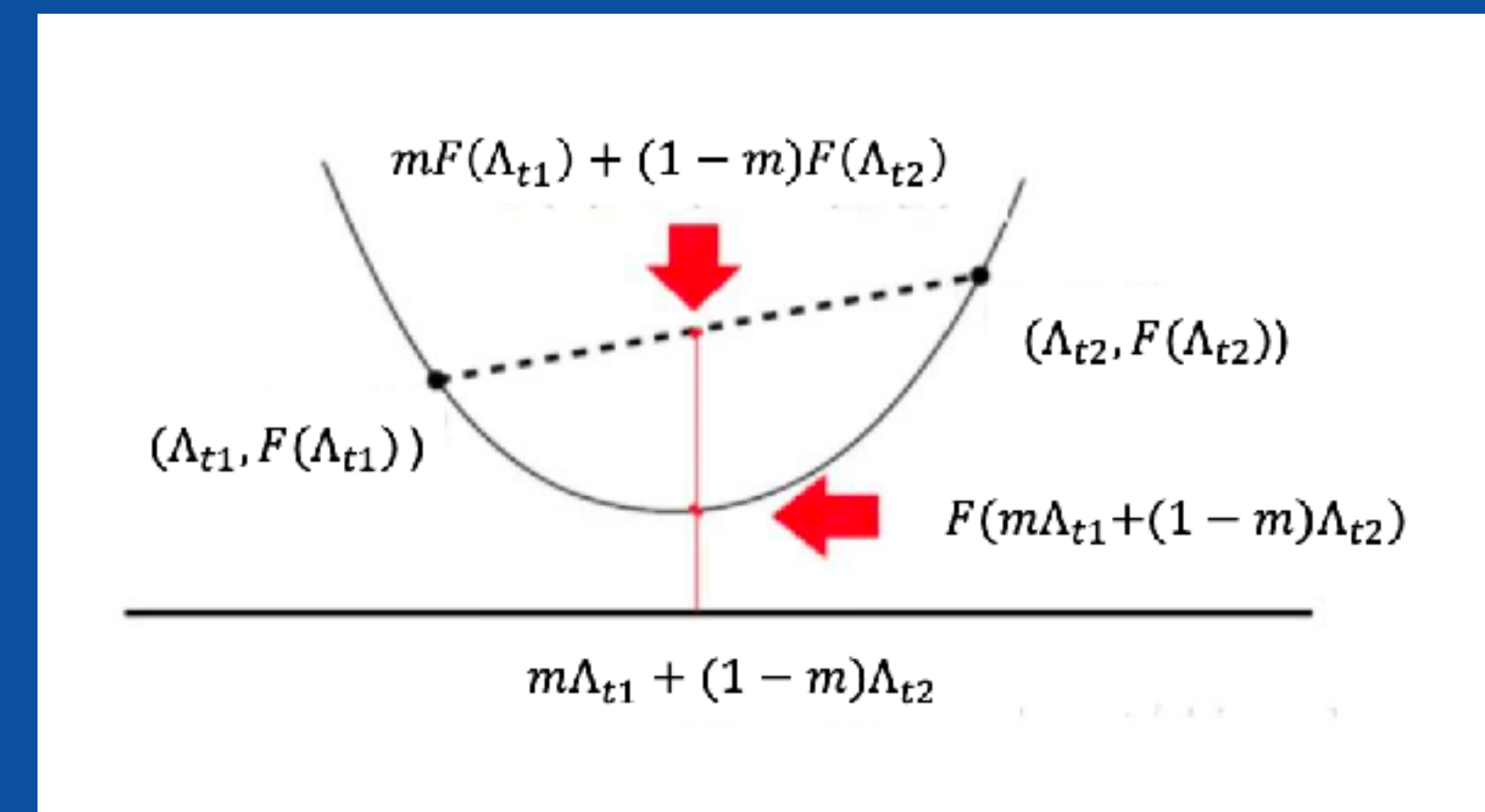
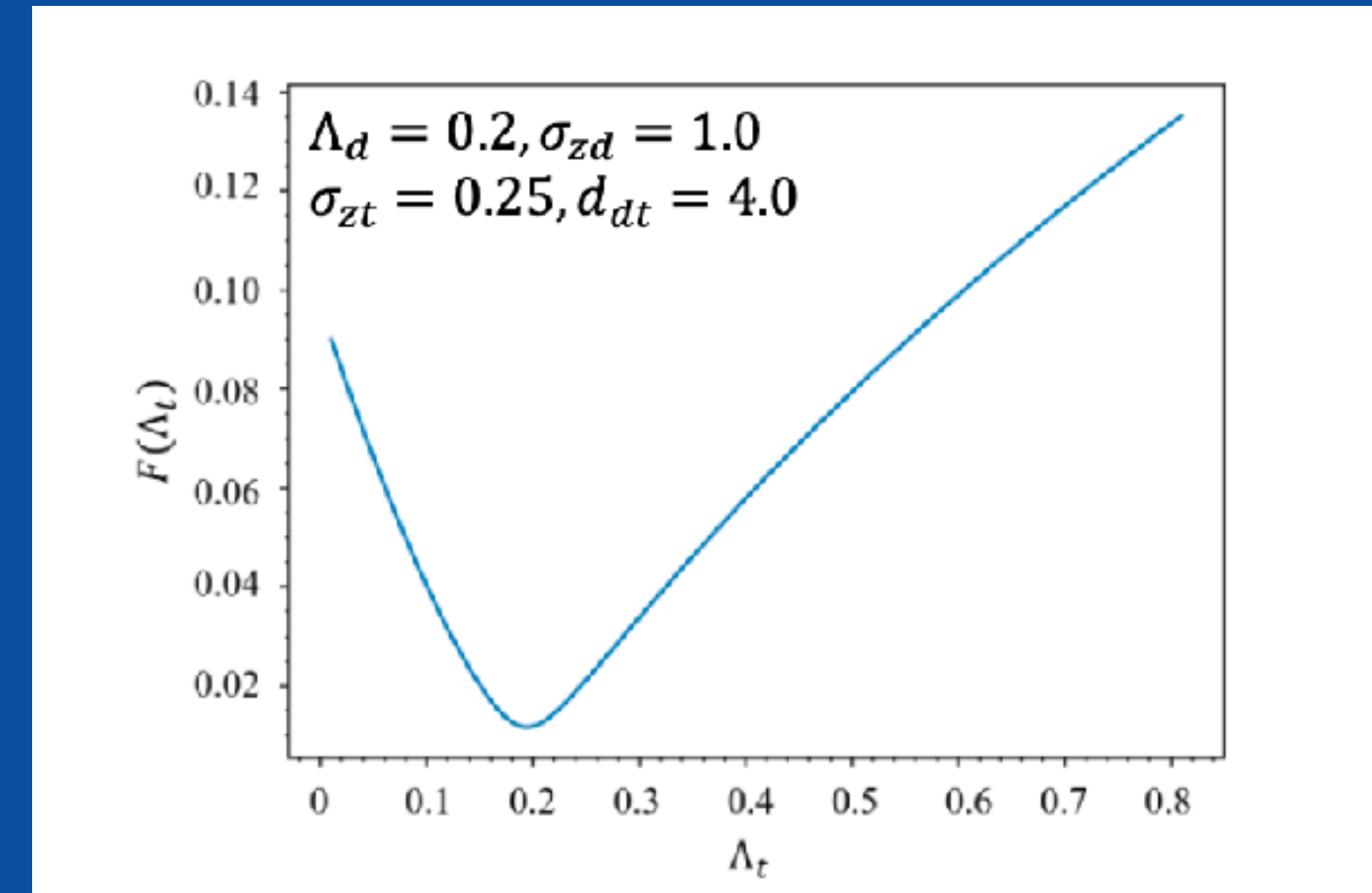
# Find the Optimal Beam Loading

\*Xiaoning Wang, et. al., Plasma Phys. Control. Fusion 64, 065007 (2022).

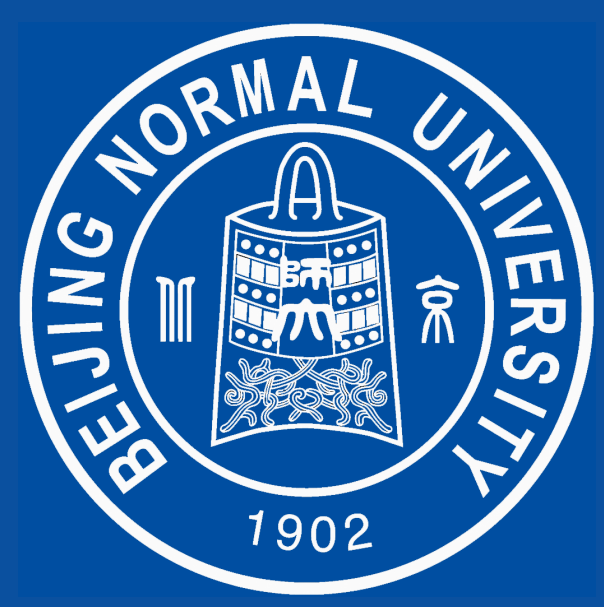


The Object Function

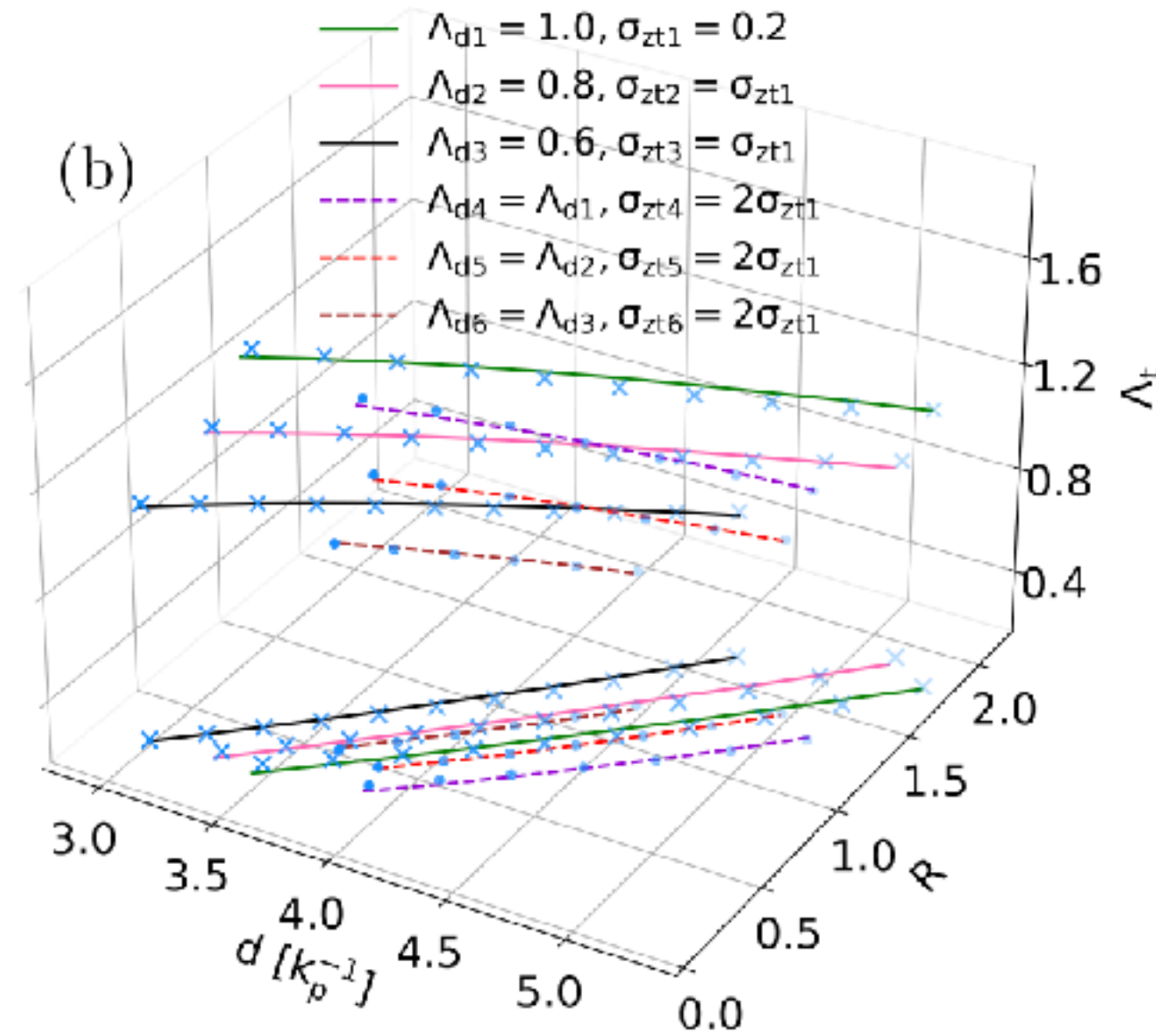
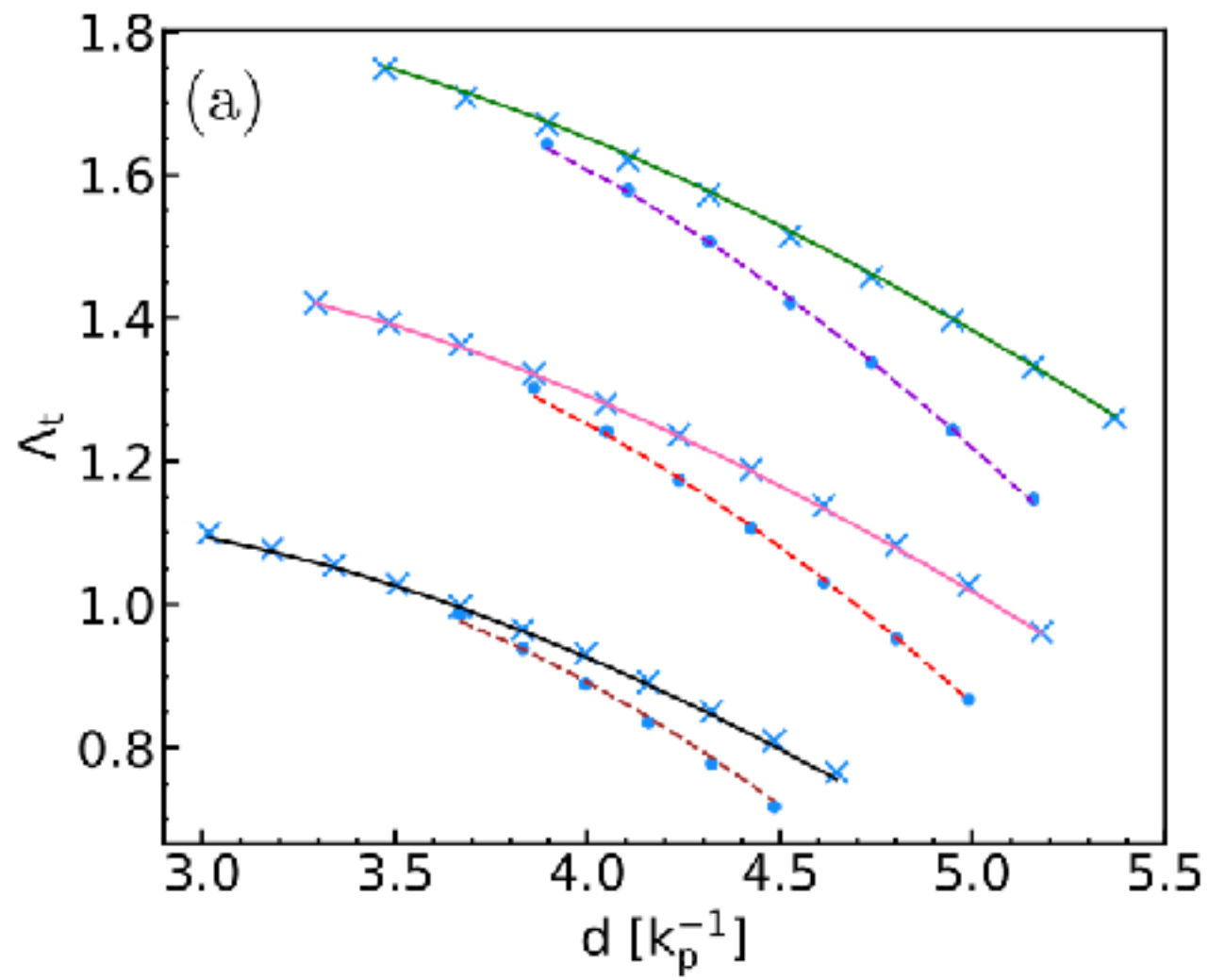
$$F(\Lambda_t) = \sqrt{\frac{\int_{\xi_s}^{\xi_e} (E_z(\xi))^2 \lambda_{bt}(\xi) d\xi}{\int_{\xi_s}^{\xi_e} \lambda_{bt}(\xi) d\xi} - \left( \frac{\int_{\xi_s}^{\xi_e} E_z(\xi) \lambda_{bt}(\xi) d\xi}{\int_{\xi_s}^{\xi_e} \lambda_{bt}(\xi) d\xi} \right)^2}$$



BFGS Method



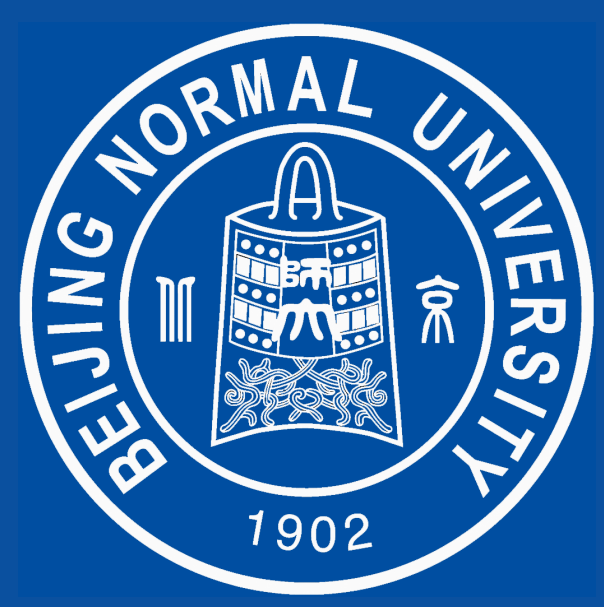
# Polynomial Fitting for $\Lambda_t$



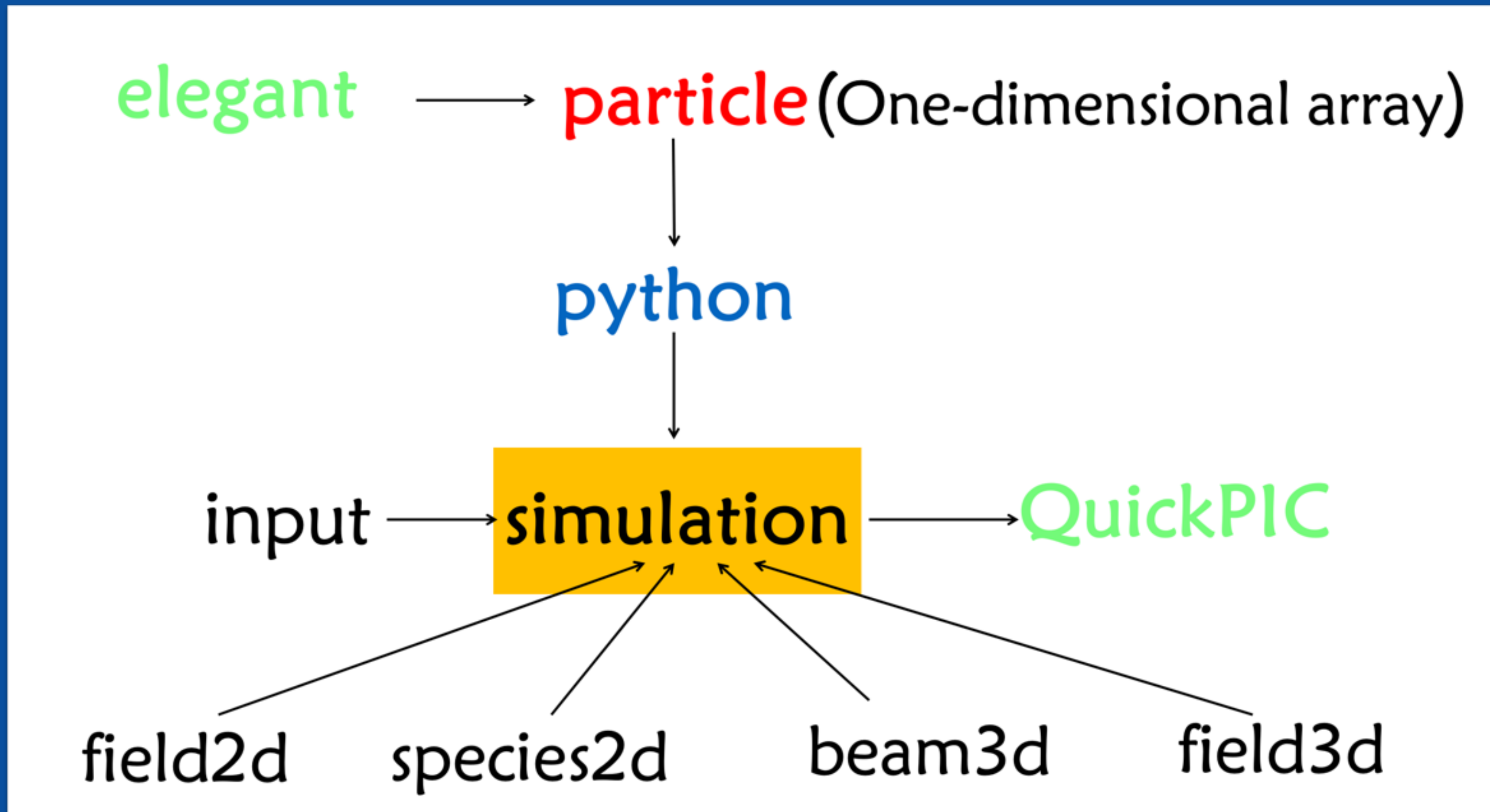
$$\Lambda_t = f(\Lambda_d, \sigma_{zd}, \sigma_{zt}, d_{dt}): r^2=0.999 \text{ at degree}=3$$

$$\begin{aligned} \Lambda_t = & h_0 + h_1\Lambda_d + h_2\sigma_{zd} + h_3\sigma_{zt} + h_4d + h_5\Lambda_d^2 + h_6\Lambda_d\sigma_{zd} \\ & + h_7\Lambda_d\sigma_{zt} + h_8\Lambda_d d + h_9\sigma_{zd}^2 + h_{10}\sigma_{zd}\sigma_{zt} + h_{11}\sigma_{zd}d \\ & + h_{12}\sigma_{zt}^2 + h_{13}\sigma_{zt}d + h_{14}d^2 + h_{15}\Lambda_d^3 + h_{16}\Lambda_d^2\sigma_{zd} \\ & + h_{17}\Lambda_d^2\sigma_{zt} + h_{18}\Lambda_d^2d + h_{19}\Lambda_d\sigma_{zd}^2 + h_{20}\Lambda_d\sigma_{zd}\sigma_{zt} \\ & + h_{21}\Lambda_d\sigma_{zd}d + h_{22}\Lambda_d\sigma_{zt}^2 + h_{23}\Lambda_d\sigma_{zt}d + h_{24}\Lambda_d d^2 \\ & + h_{25}\sigma_{zd}^3 + h_{26}\sigma_{zd}^2\sigma_{zt} + h_{27}\sigma_{zd}^2d + h_{28}\sigma_{zd}\sigma_{zt}^2 + h_{29}\sigma_{zd}\sigma_{zt}d \\ & + h_{30}\sigma_{zd}d^2 + h_{31}\sigma_{zt}^3 + h_{32}\sigma_{zt}^2d + h_{33}\sigma_{zt}d^2 + h_{34}d^3, \quad (2) \end{aligned}$$

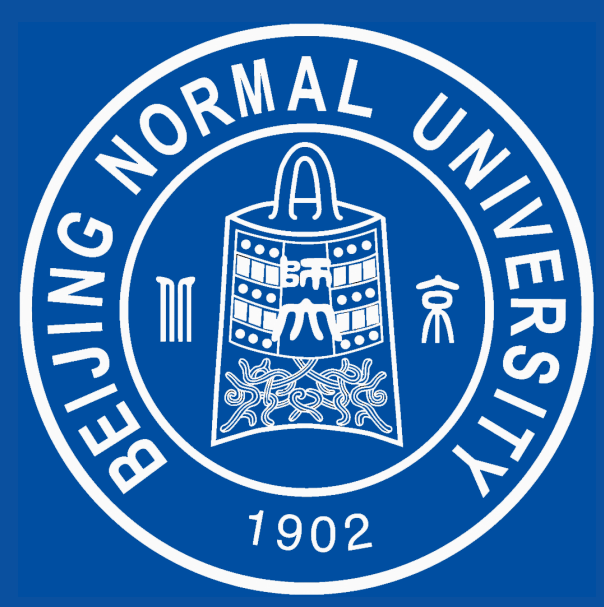
$$R = f(\Lambda_d, \sigma_{zd}, \sigma_{zt}, d_{dt}, \Lambda_t): r^2=0.99 \text{ at degree}=2$$



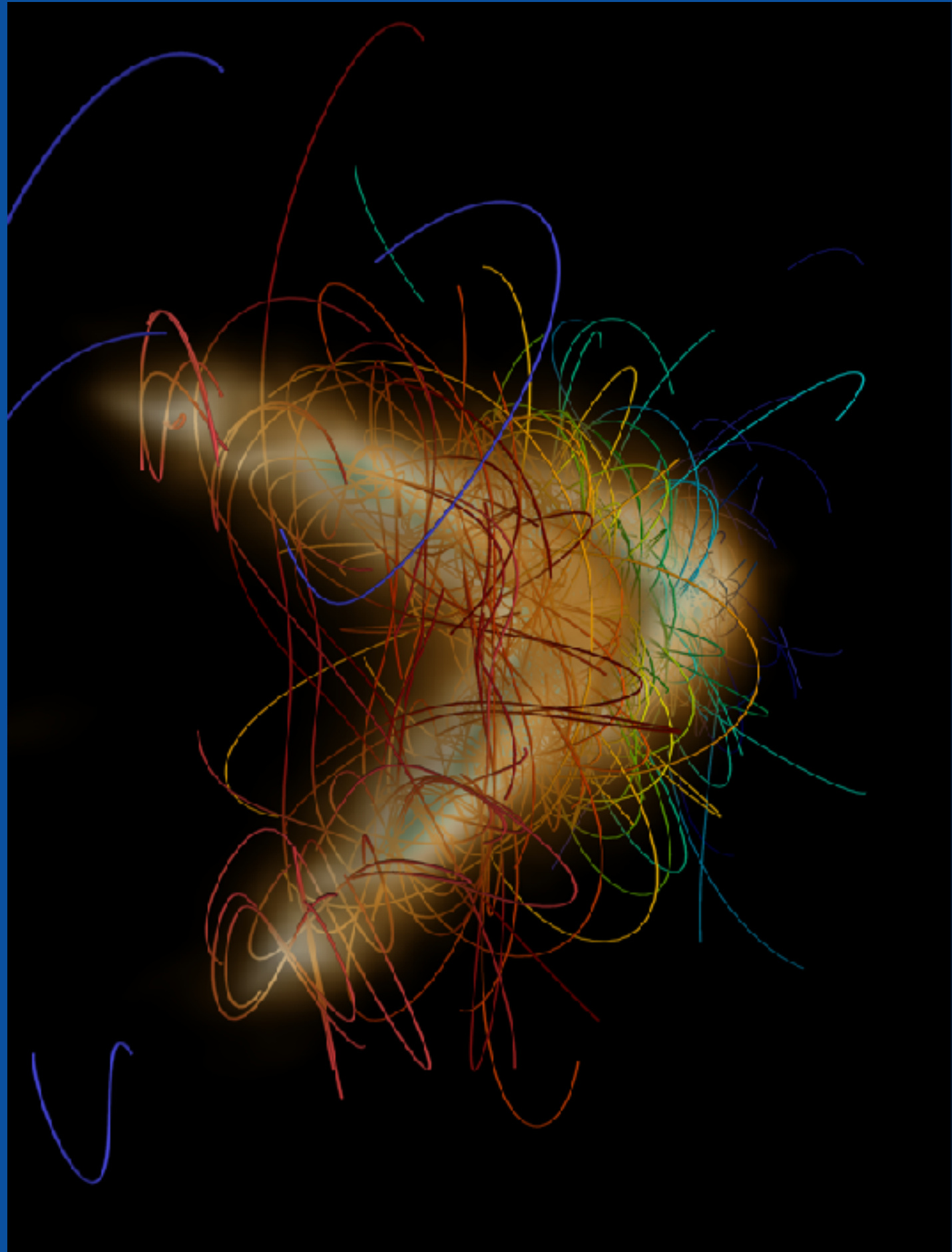
# Plasma-based Accelerator Optimization Framework



CTYPES



# Thanks!



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Oct 4th, 2024