



The Baryon Antibaryon Symmetry Experiment (BASE)



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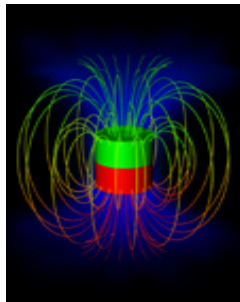


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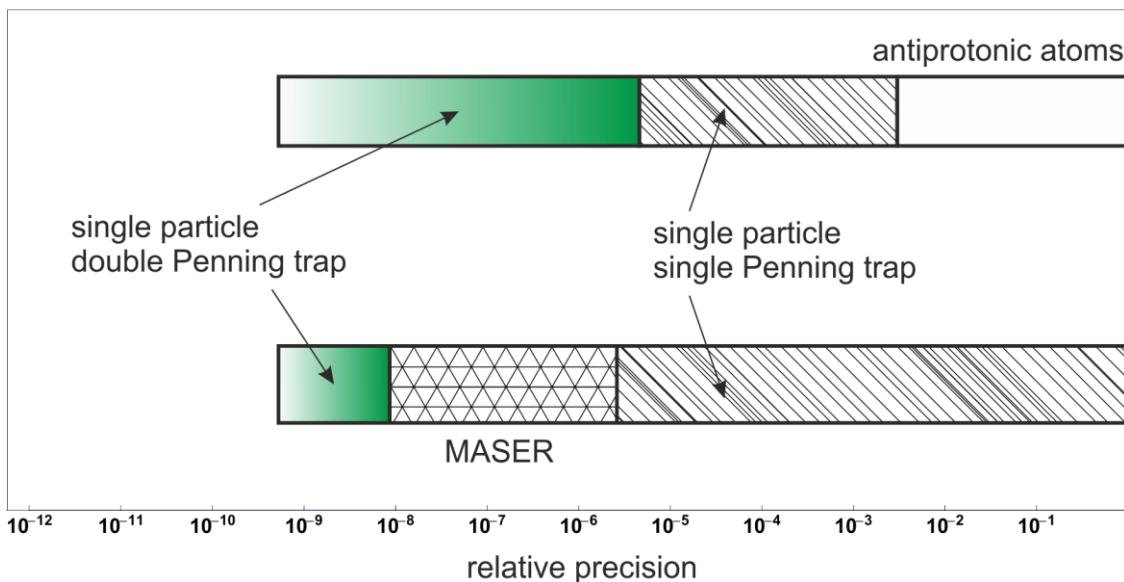


Motivation

High-precision test of the CPT invariance
by comparing the proton and antiproton magnetic moments:



$$\vec{\mu}_{p/\bar{p}} = g_{p/\bar{p}} \frac{q_{p/\bar{p}}}{2m_{p/\bar{p}}} \vec{S}$$



$$g_{\bar{p}} = 5.585690(24)$$

$$g_p = 5.585694713(46)$$



Outline

Motivation of the double-trap method

- Experimental principle of g-factor measurements
- Spin-state detection of a single proton/antiproton
- Double-Penning trap method

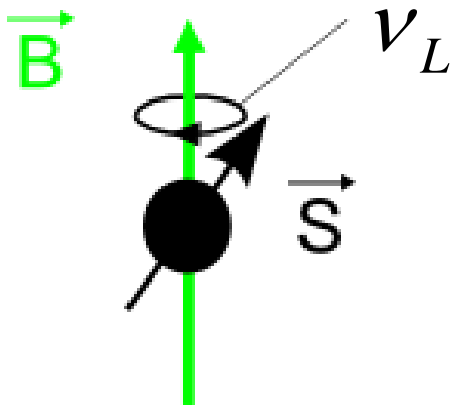
BASE

- Design of the new apparatus
- Status of the implementation in the AD

Experimental principle

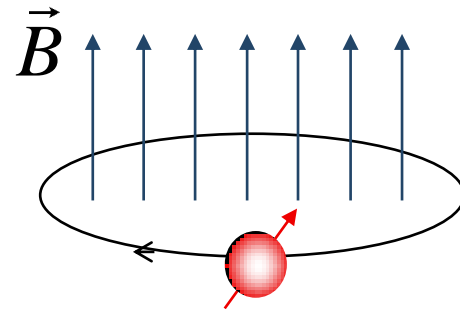
$$\frac{\mu_{p/\bar{p}}}{\mu_N} = \pm \frac{g_{p/\bar{p}}}{2} = \pm \frac{\nu_L}{\nu_C}$$

Spin precession



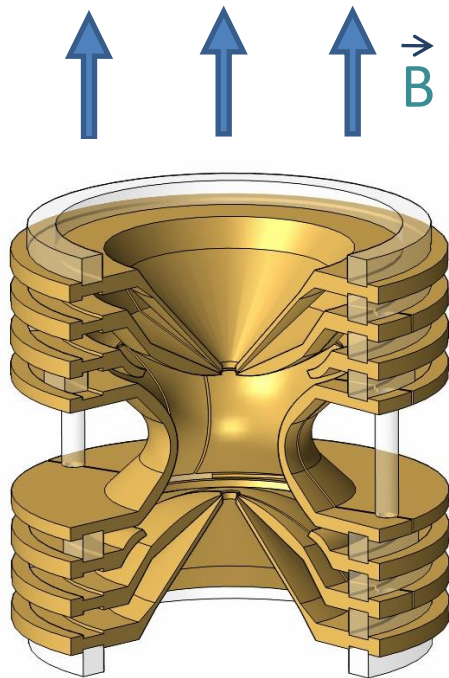
$$\nu_L = \frac{1}{2\pi} \frac{g}{2} \frac{q \cdot B}{m}$$

Cyclotron Motion



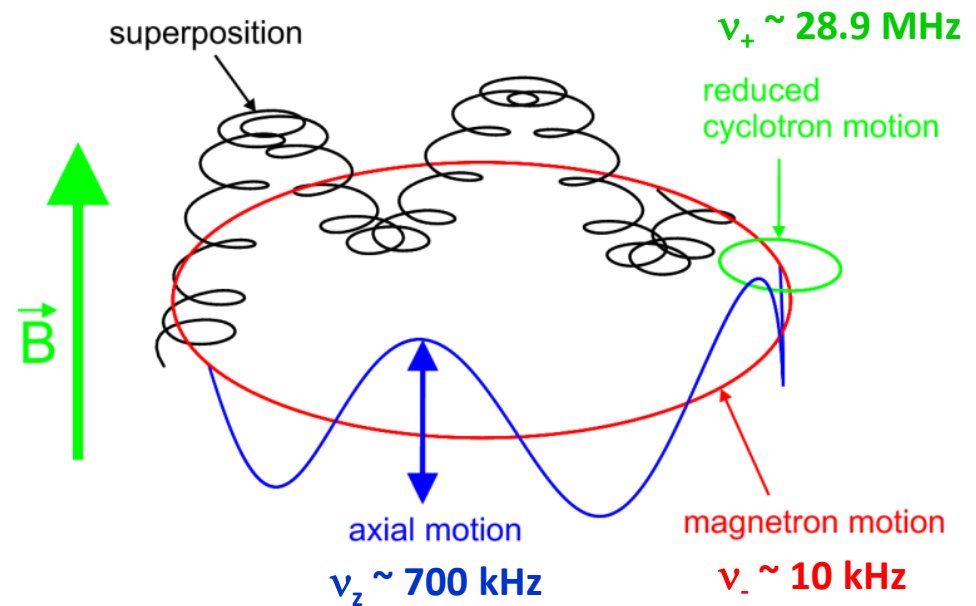
$$\nu_c = \frac{1}{2\pi} \frac{q \cdot B}{m}$$

The Penning trap



- Strong homogeneous magnetic field
- Weak electric quadrupole field

Motion in a Penning Trap



$$\nu_c^2 = \nu_+^2 + \nu_-^2 + \nu_z^2$$

Image-current detection

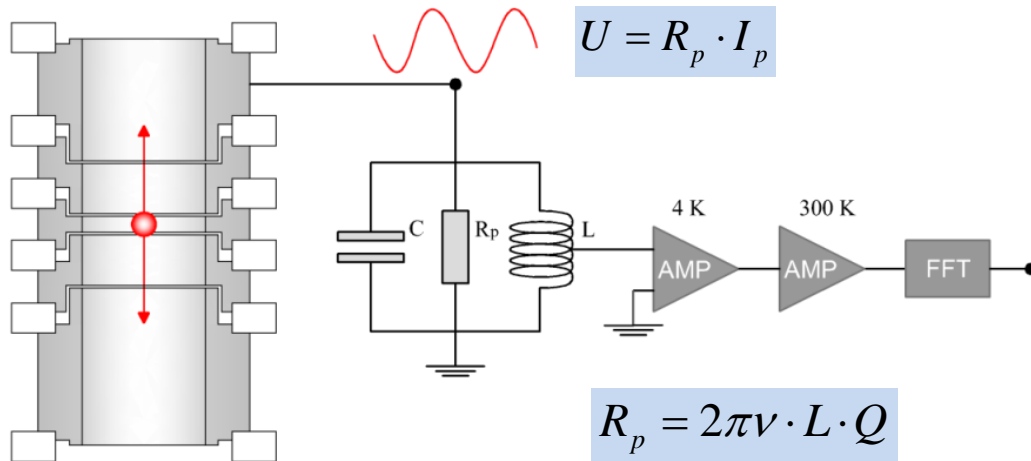
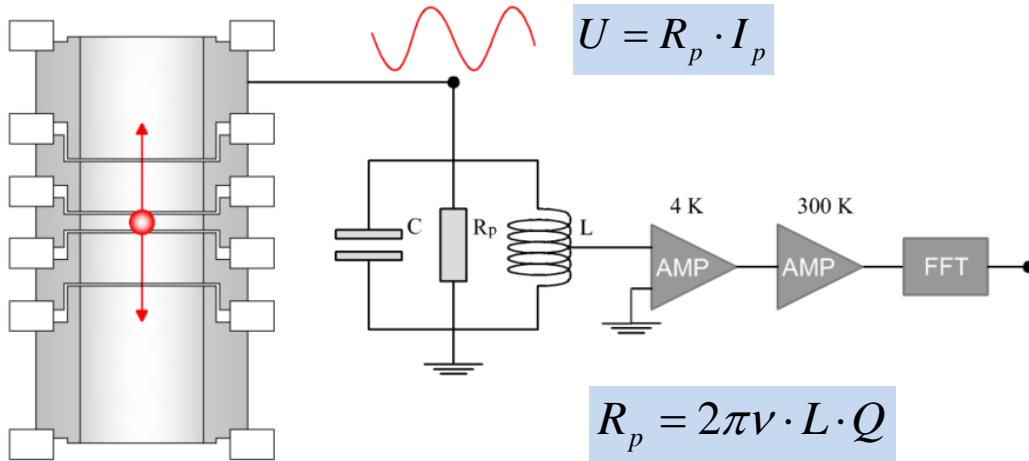


Image-current detection



Detection of the axial frequency

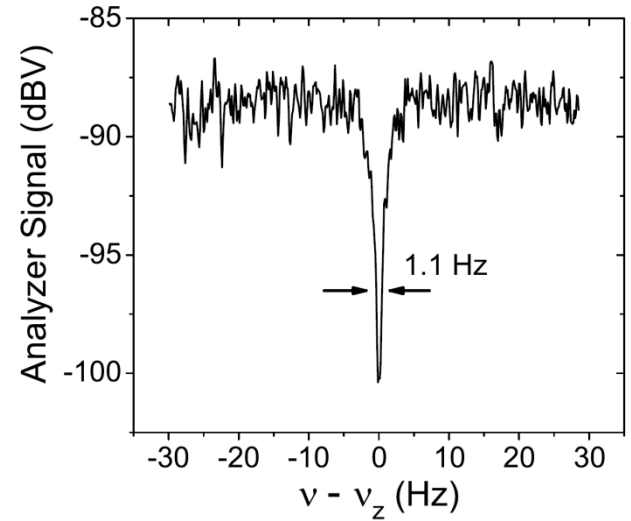
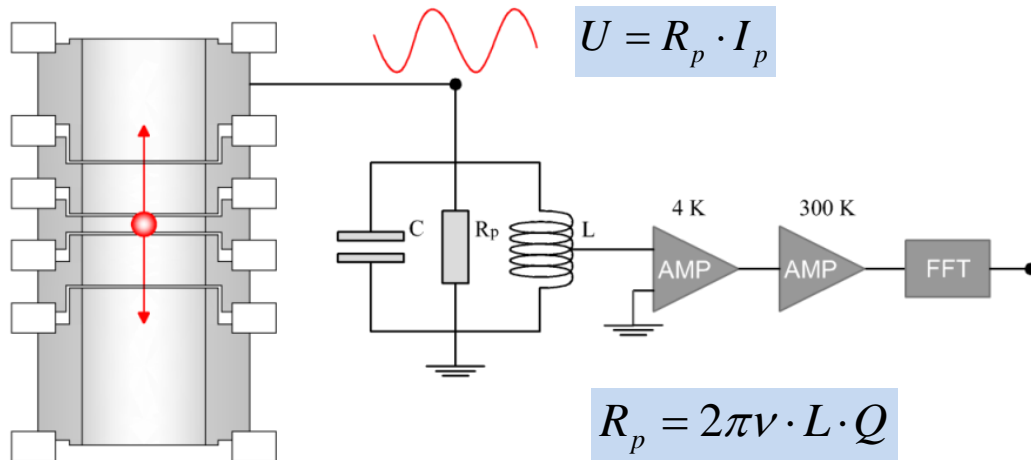
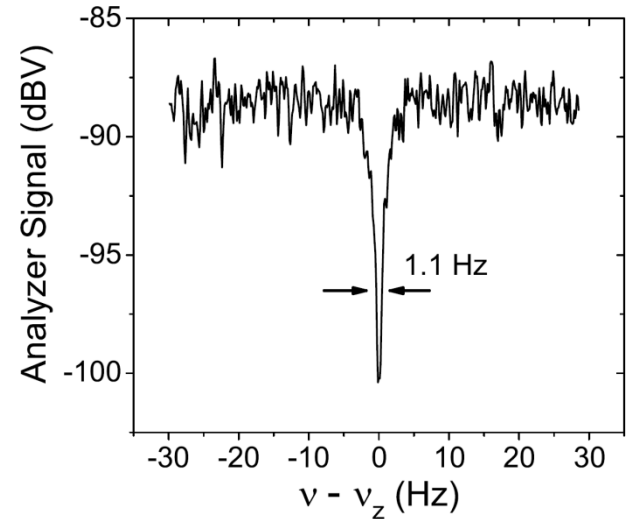


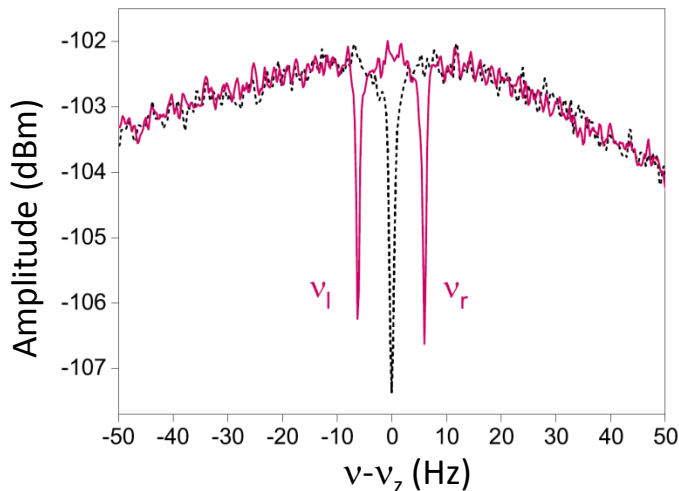
Image-current detection



Detection of the axial frequency



Sideband coupling to detect ν_+ and ν_-



Irradiation of a coupling rf-signal at $(\nu_+ - \nu_z)$ or $(\nu_z + \nu_-)$

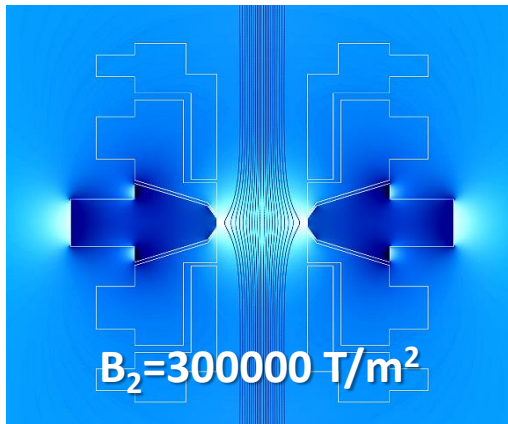
Amplitude modulation of the axial motion

Measurement of the axial and sideband frequencies

Extraction of the free cyclotron frequency

Continuous Stern-Gerlach Effect

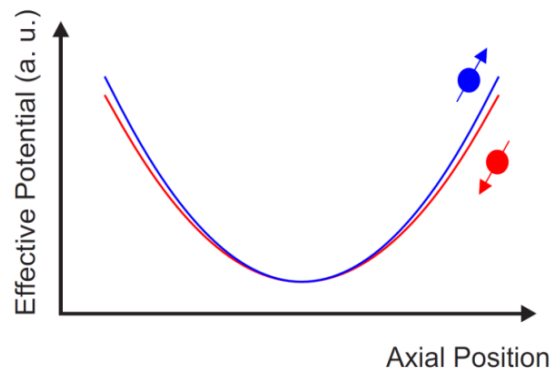
Spin-state detection by coupling of the magnetic moment to the axial motion



Magnetic potential:
$$\Phi_M = -\vec{\mu}_{p/\bar{p}} \cdot \vec{B}$$

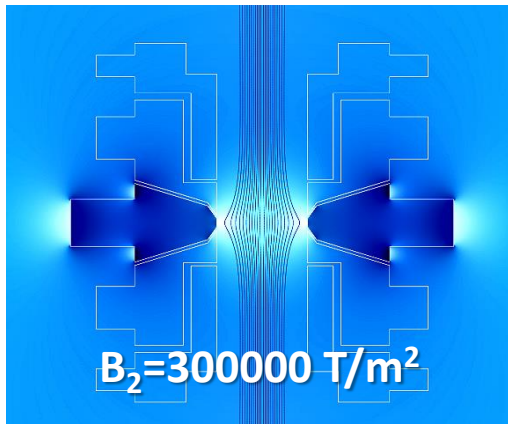
Add a magnetic field perturbation:

$$B_z = B_0 + B_2 \left(z^2 - \frac{\rho^2}{2} \right)$$



Continuous Stern-Gerlach Effect

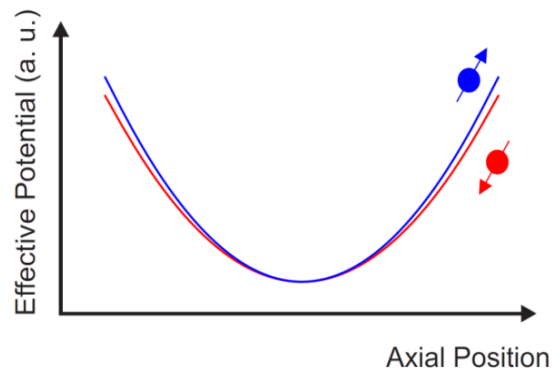
Spin-state detection by coupling of the magnetic moment to the axial motion



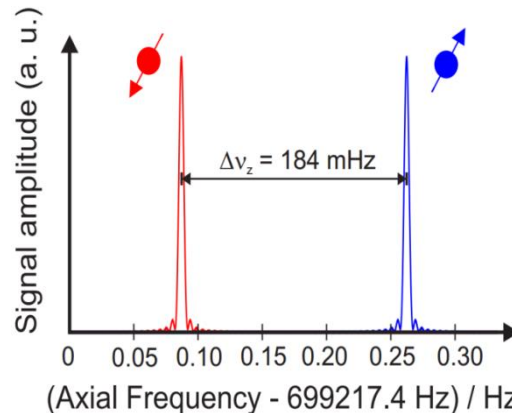
Magnetic potential: $\Phi_M = -\vec{\mu}_{p/\bar{p}} \cdot \vec{B}$

Add a magnetic field perturbation:

$$B_z = B_0 + B_2 \left(z^2 - \frac{\rho^2}{2} \right)$$



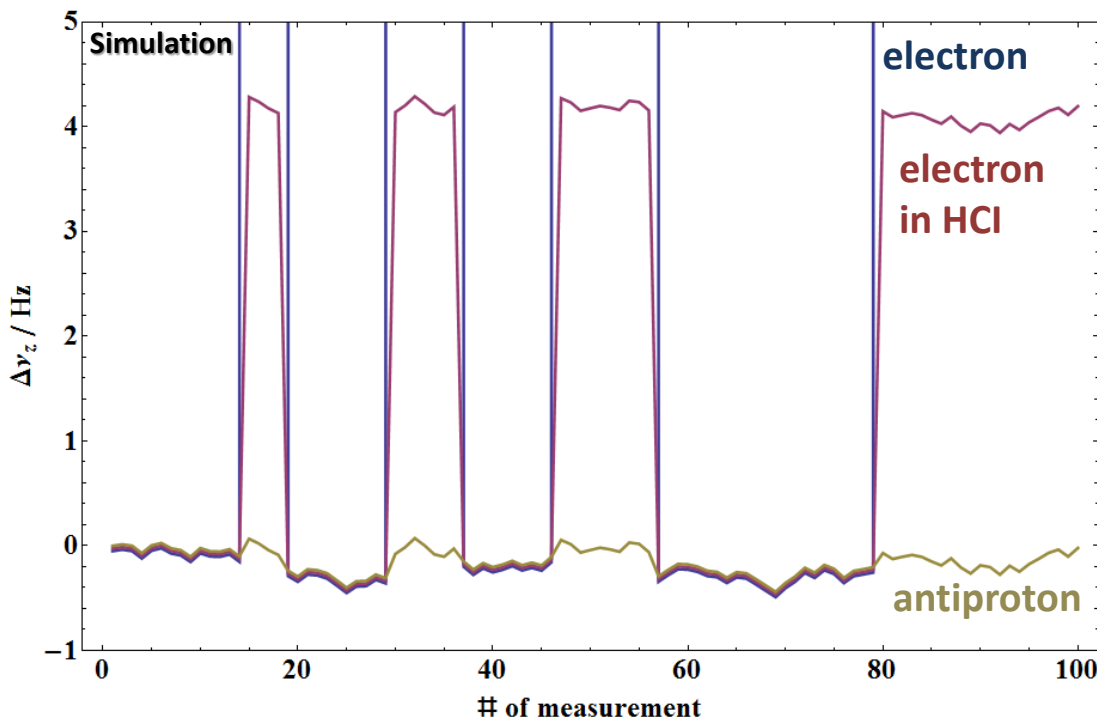
The axial frequency becomes spin-state dependent:



$$\frac{\Delta \nu_z}{\nu_z} \approx \frac{1}{4\pi^2} \frac{\mu}{m} \frac{B_2}{\nu_z^2}$$

Larmor frequency measurement

- A sequence of axial frequency measurements
- Drive the spin-flip with an rf-signal
- Extract the spin-flip probability as function of the drive frequency



Axial frequency difference:

$$\Delta\nu_z \approx \frac{1}{4\pi^2} \frac{\mu}{m} \frac{B_2}{\nu_z}$$

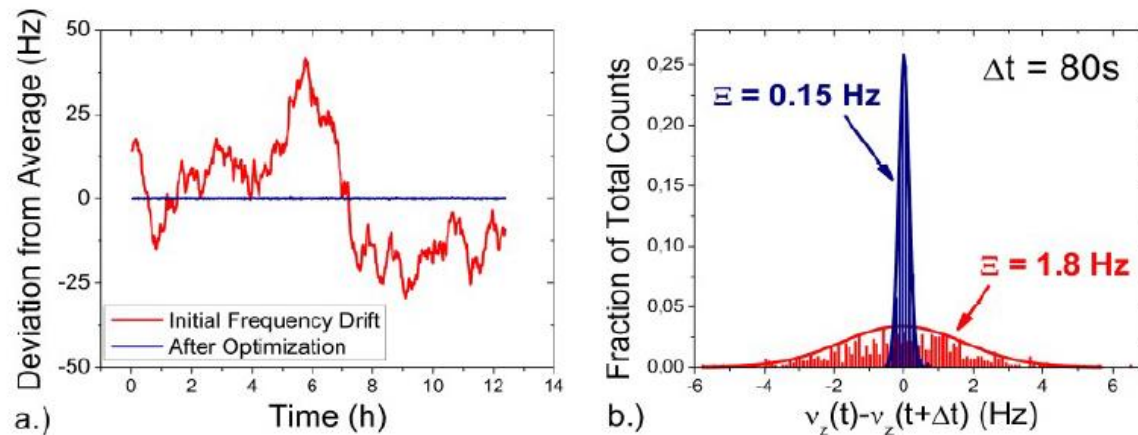
$$B_2 = 300000 \text{ T/m}^2$$

Electron:	228 kHz
Electron in $^{28}\text{Si}^{13+}$:	4.4 Hz
Proton/antiproton:	0.18 Hz

It's not that simple

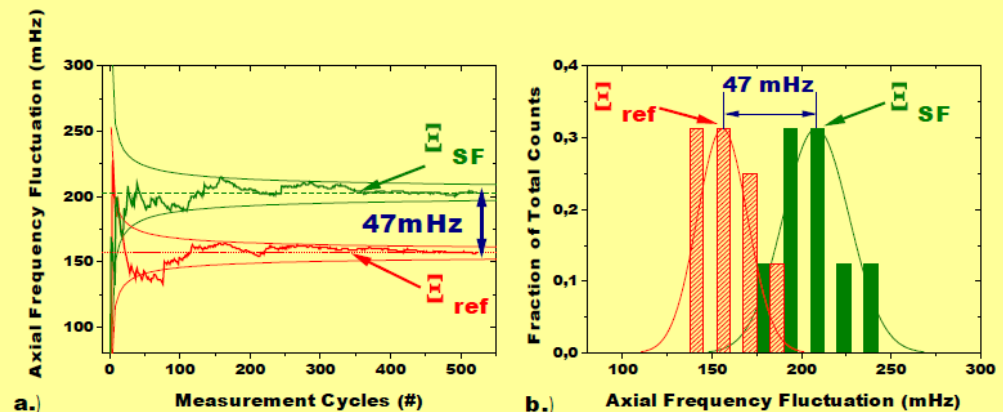
Magnetic bottle coupling: $\Delta \nu_z = \frac{1}{4\pi^2 m \nu_z} \frac{B_2}{B_0} (dE_+ + dE_-) \rightarrow 1\text{Hz}/\mu\text{eV}$

Axial frequency fluctuations:



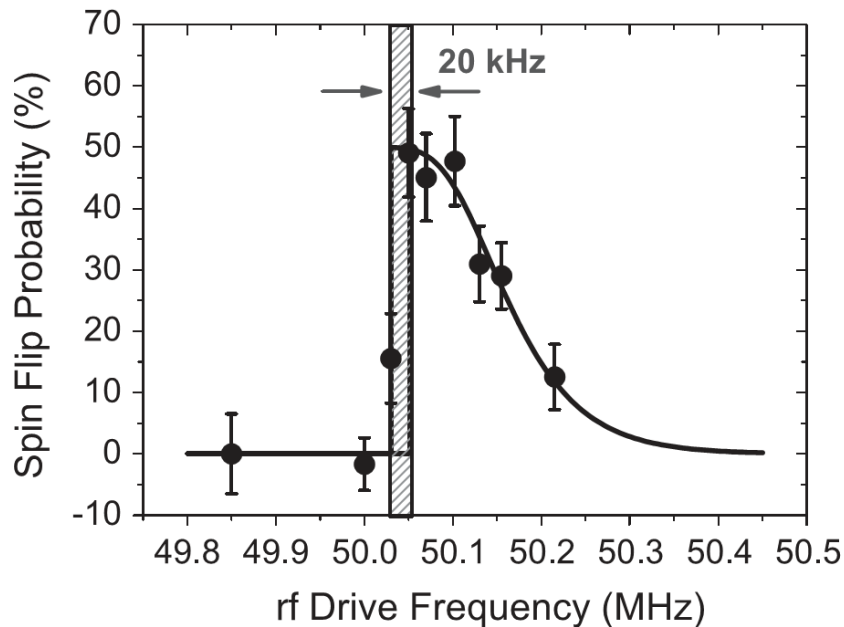
Statistical detection
of spin-flips:

$$\Xi_{SF} = \sqrt{\Xi_{ref}^2 + \Delta \nu_{z,SF}^2 \cdot P_{SF}(\nu_L)}$$



Measurement of the Larmor frequency

Measurement of the Larmor-frequency with the statistical detection method:

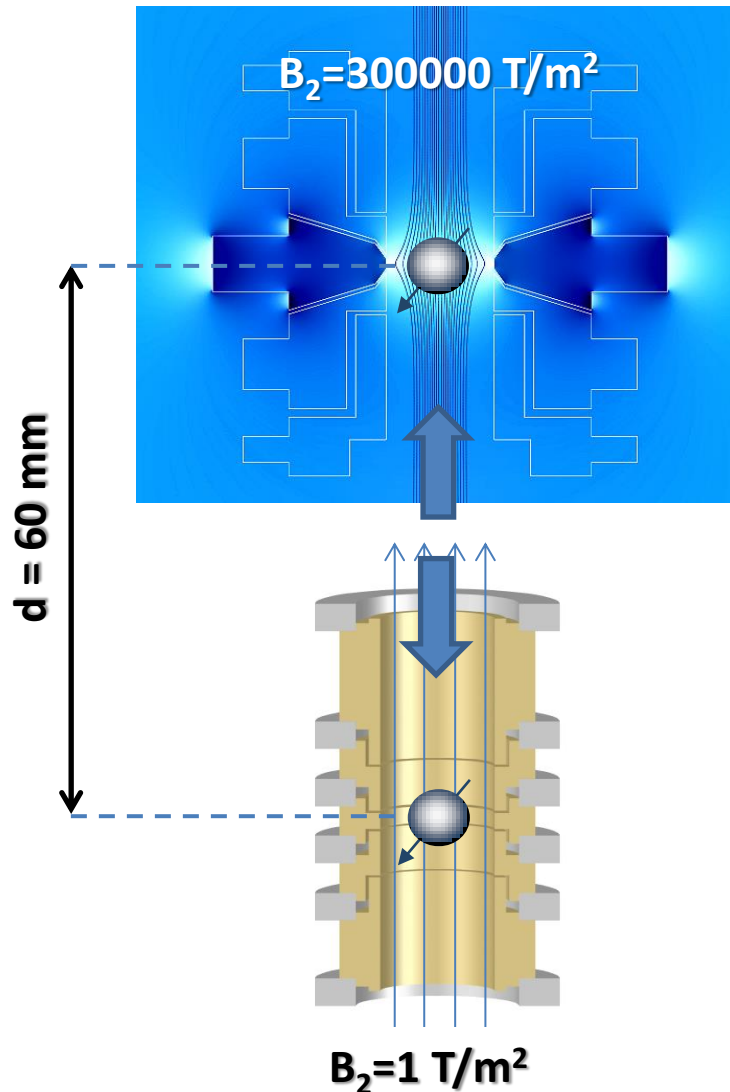


$$\nu_L = \nu_{L,0} \left(1 + \frac{B_2}{B_0} \cdot \langle z^2 \rangle \right)$$

Lineshape is a convolution of the unperturbed Rabi-resonance and Boltzmann-distributed axial energy.

This limits the resolution of the frequency measurement on the 10^{-6} level!

The Double Trap Method



Spatial separation of the frequency measurement and the spin-state detection

- Drive the spin until the spin-state is known

Requires that the spin-state is known!

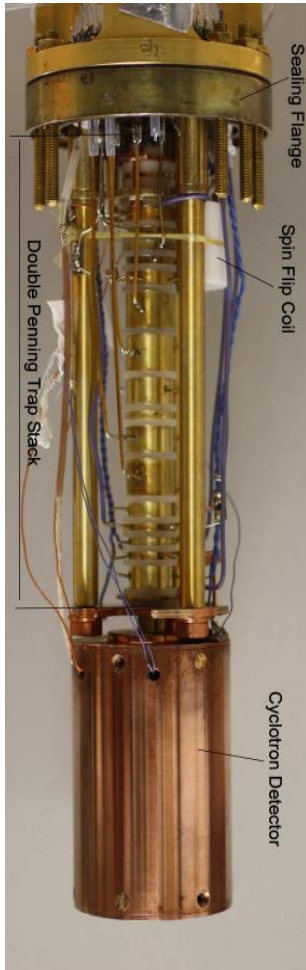
- Transport the particle to the precision trap
- Measure ν_c and drive the spin flip for the Larmor resonance
- Transport the particle to the analysis trap
- Determine the spin state

Simple idea, BUT...

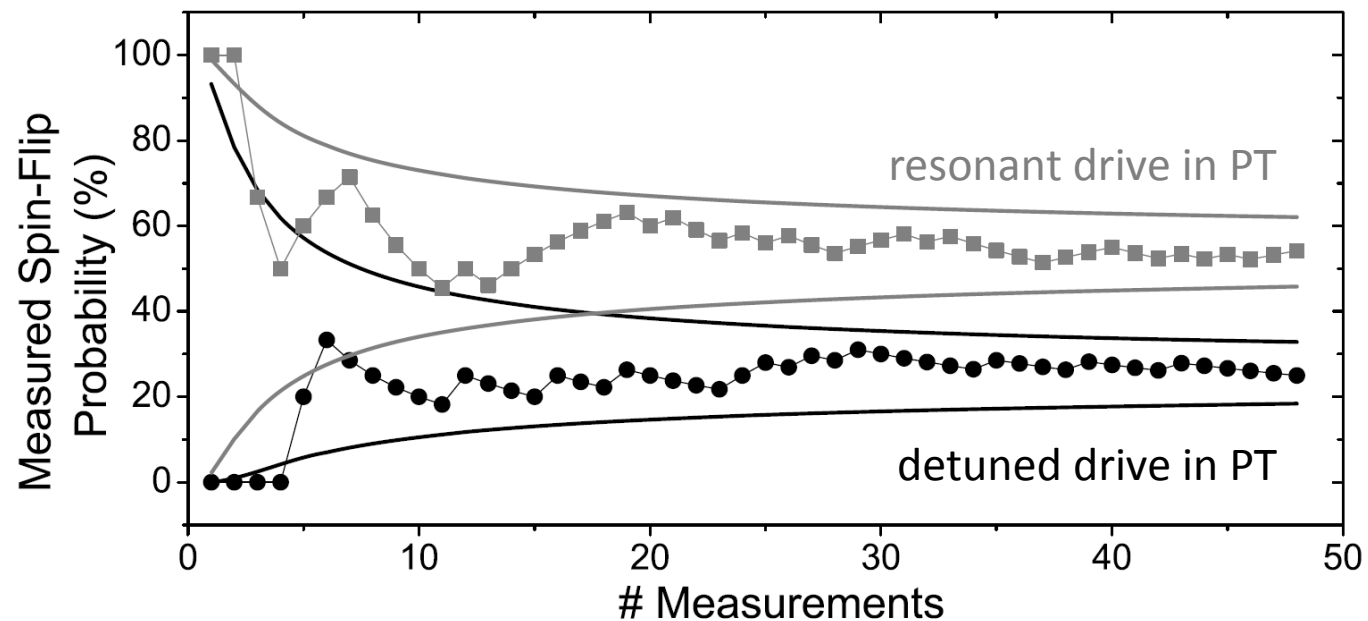
Single spin-flip resolution is required!

Outlook: g-factor proton Mainz

Double Penning trap system – Mainz University



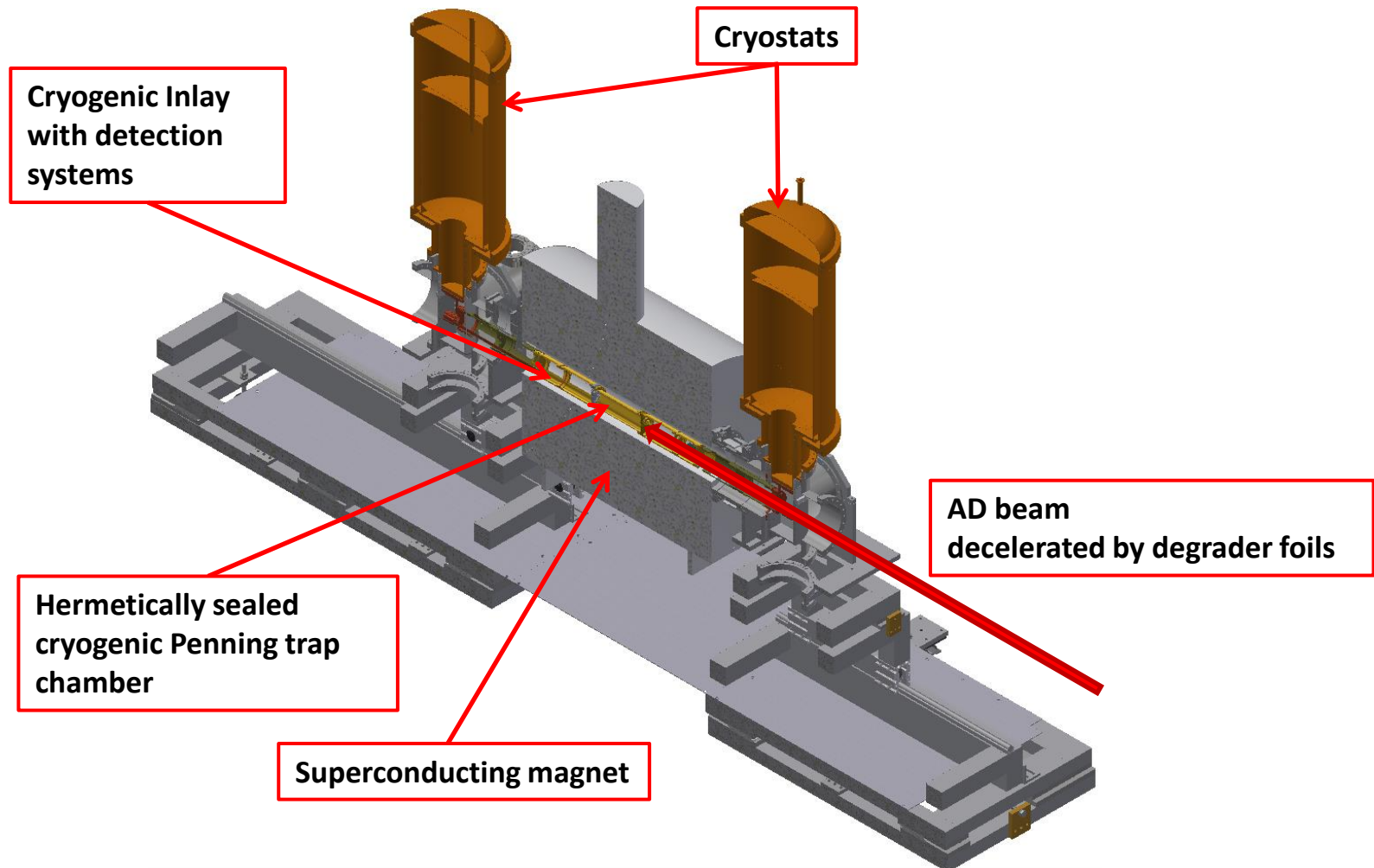
Application of the double-trap scheme with a single proton



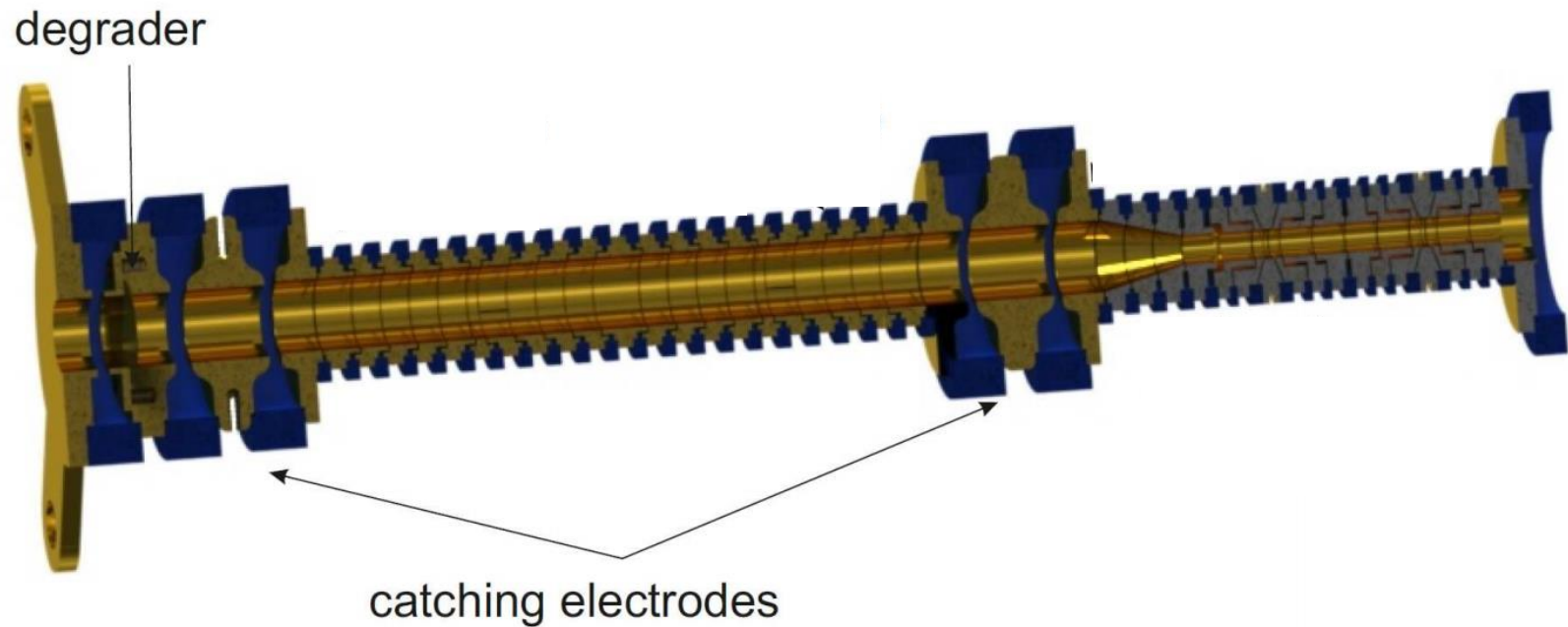
More on this hot topic on **Friday 10.40 h** by **Andreas Mooser!**

The BASE apparatus

Setup of a new experiment in the AD hall to apply the double-trap method to the antiproton

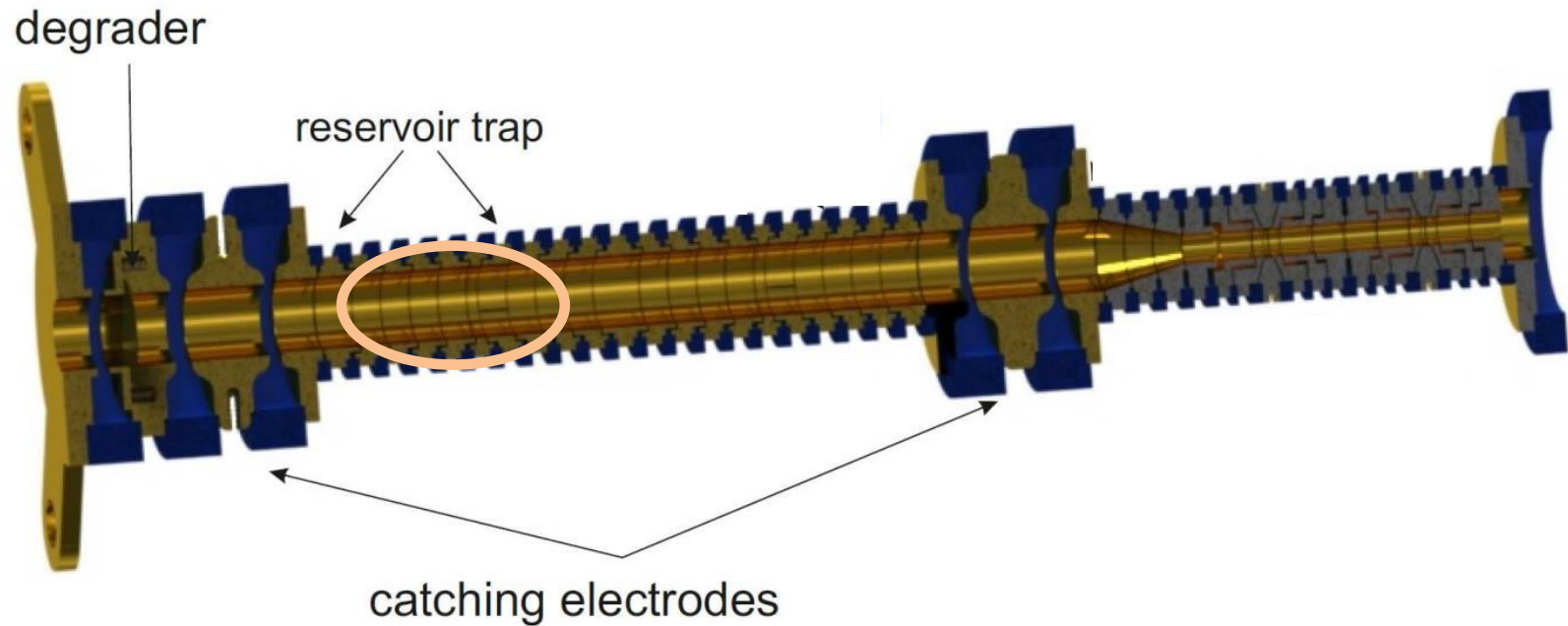


The BASE Penning trap stack



The BASE Penning trap stack

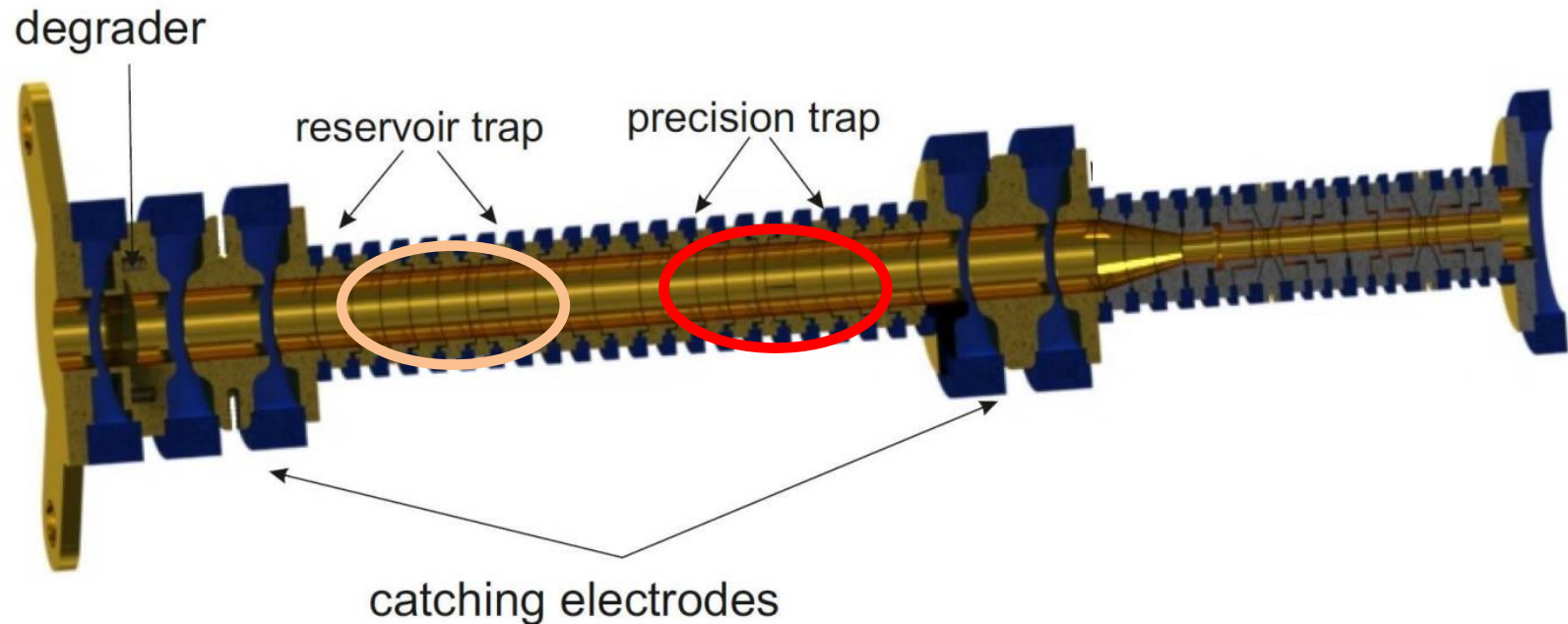
Storage of an Antiproton cloud



The BASE Penning trap stack

Cyclotron frequency measurement
Application of the spin-flip drive

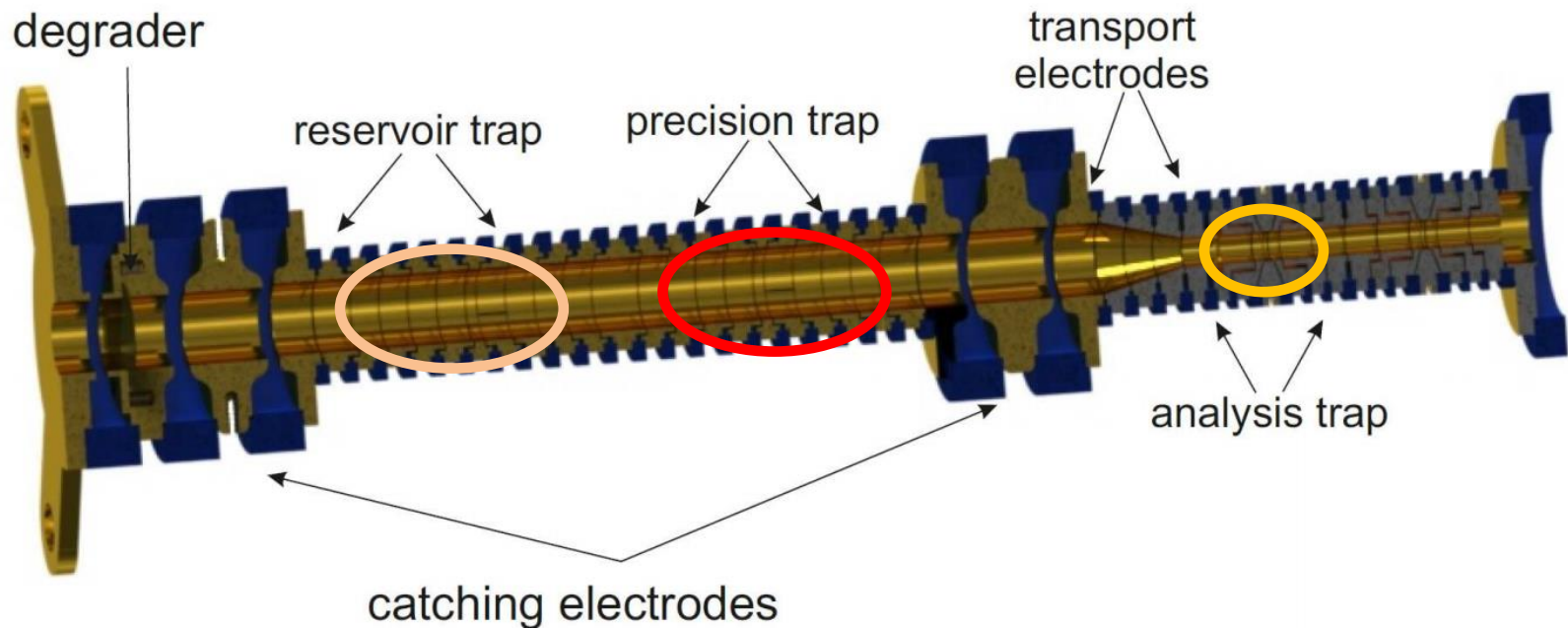
Storage of an Antiproton cloud



The BASE Penning trap stack

Cyclotron frequency measurement
Application of the spin-flip drive

Storage of an Antiproton cloud

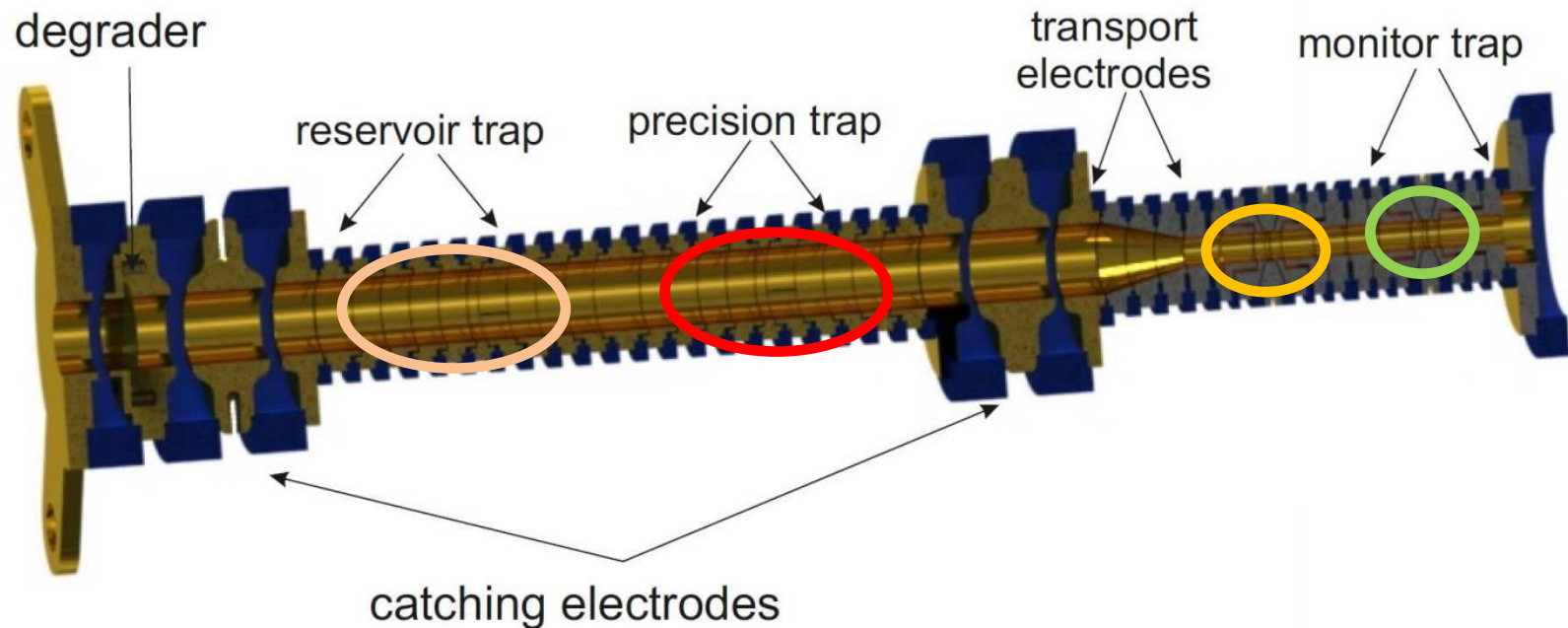


Spin-state detection

The BASE Penning trap stack

Cyclotron frequency measurement
Application of the spin-flip drive

Storage of an Antiproton cloud

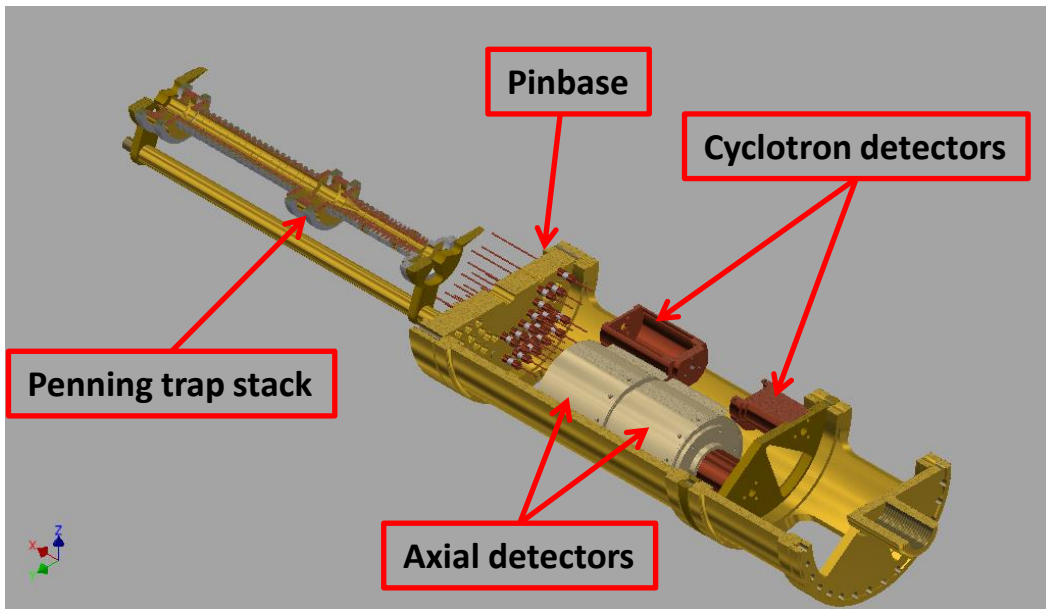


Spin-state detection

Magnetic field monitoring

The BASE detection systems

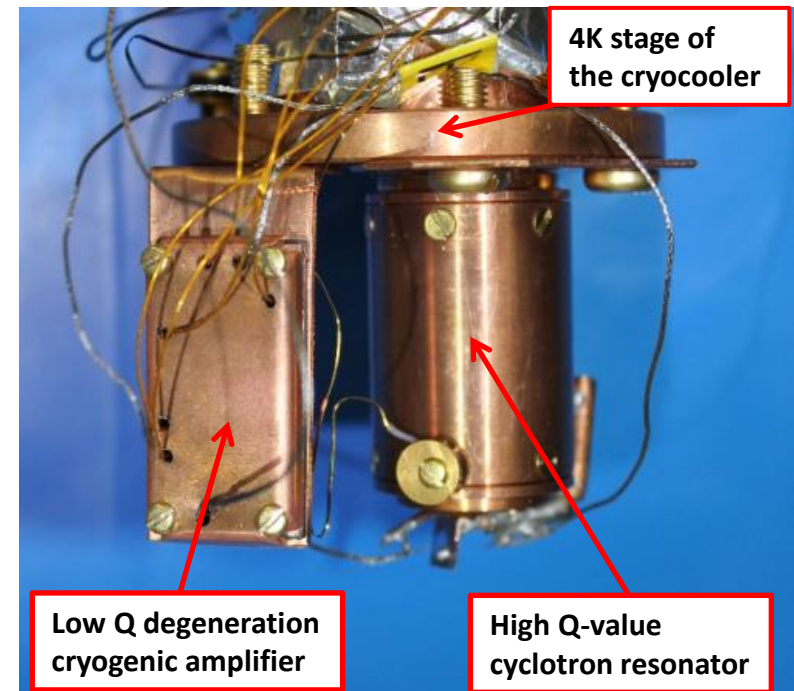
Currently: development of the detection systems



- Four axial detection systems
NbTi resonators with NbTi coil
- Two cyclotron detection systems
Copper resonator with copper coil

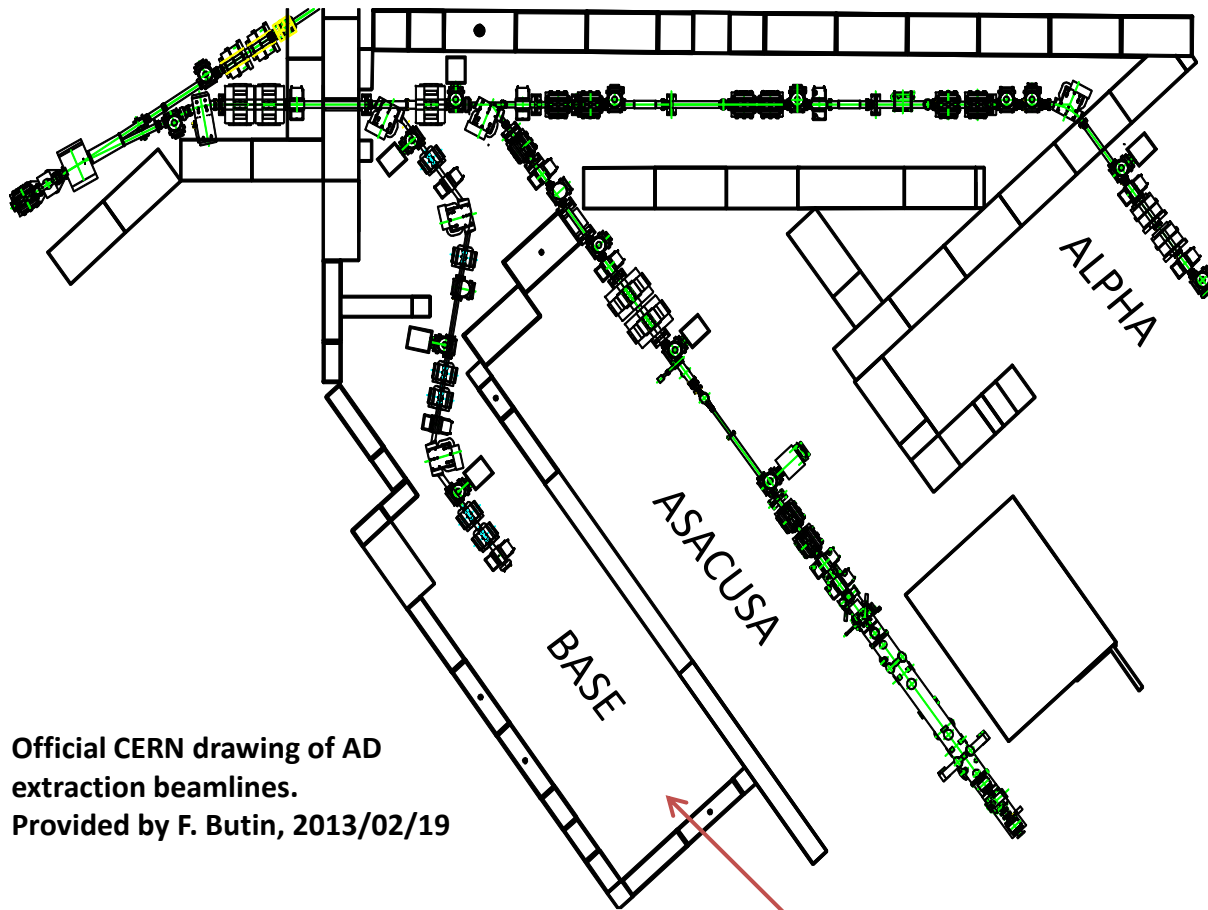
Cyclotron detectors
commissioned and
tested:

Q – 1500 - 3500



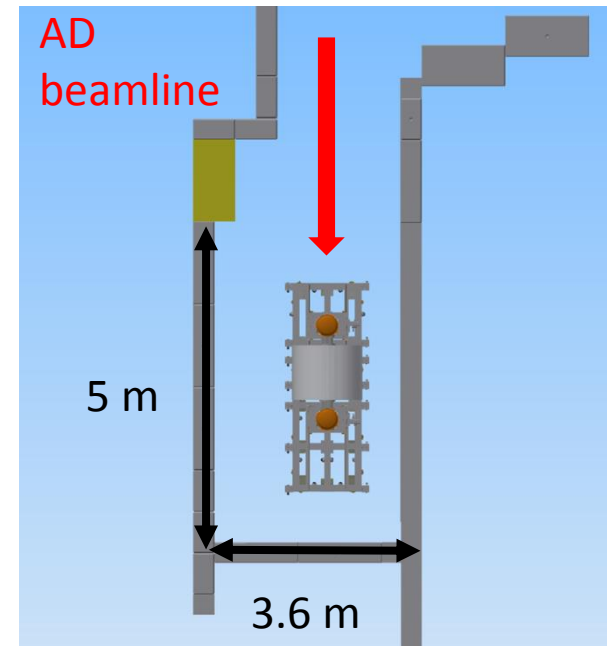
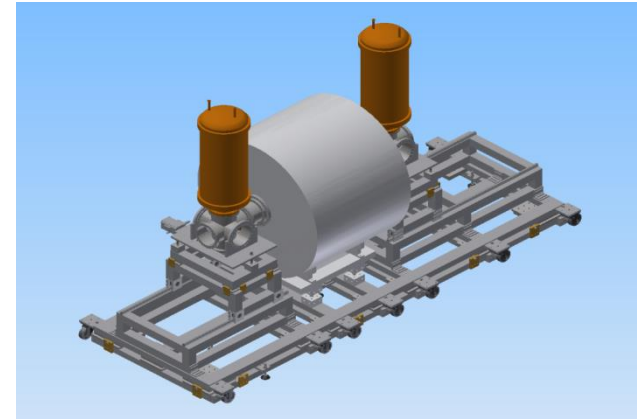
Implementation of BASE in the AD

**BASE approved by the CERN research board
Installation of BASE in the AD during LS1**



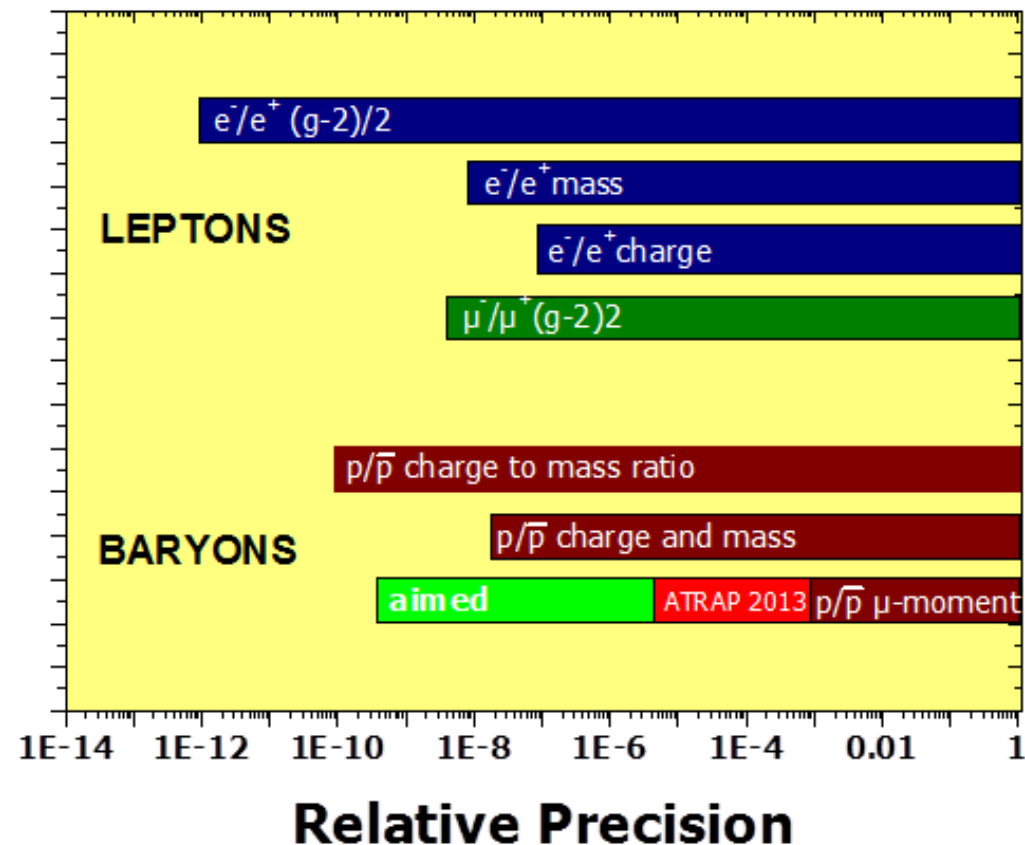
Official CERN drawing of AD
extraction beamlines.
Provided by F. Butin, 2013/02/19

**New experimental zone
proposed by AD**



Conclusions

- BASE aims for a measurement of the antiproton magnetic moment with ppb precision
- Statistical spin-state detection resulted in a ppm measurement of the proton magnetic moment
- Single spin-flip resolution was achieved
- The double-trap method was applied for the first time to a proton
- Installation of BASE in the AD is in progress



Funding



Thank you for your attention!