

Hypernuclei: Stability and prospects to study the compound nuclear formation and its fragmentation

Kaushal Thakkar and P C Vinodkumar

Department of Physics, Sardar Patel University,
Vallabh Vidyanagar-388120, Gujarat, India

OVERVIEW

- Introduction
- Motivation

Part 1

- BE of orbitally excited Hypernuclei
- Various potential to study Λ -N interaction
- Result

Part 2

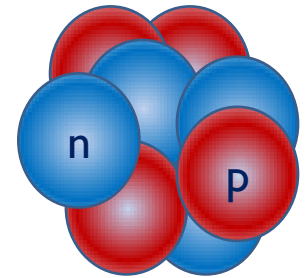
- Alpha Decay of heavy Hypernuclei
- Half life time of heavy Hypernuclei
- Summary

INTRODUCTION



Nucleus

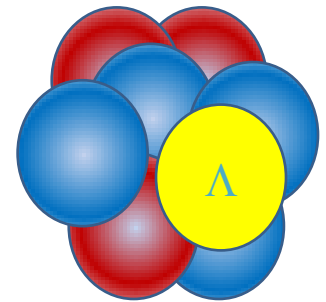
many-body system of protons and neutrons



Hypernucleus

Hyperon in a nucleus

a new degree-of-freedom, “strangeness”.



MOTIVATION

- Properties of hypernucleus are of considerable interest among the intermediate energy regime as it stands at the intersection of nuclear and particle physics.
- The enhanced experimental efforts have thus created renewed interest in the theoretical computations of the various properties of hypernuclei. [1, 2].

[1] E. Hiyama and T. Yamada, Progress in Part. and Nuc. Phy. **63**, 339-395 (2009).

[2] O. Hashimoto and H. Tamura, Progress in Part. and Nuc. Phy. **57**, 564-653 (2006).

BE OF HYPERNUCLEI

- To calculate the binding energy of hypernuclei, we have solved Schrodinger equation numerically.

$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + V(r) \right] \psi(r) = \epsilon \psi(r)$$

$$\mu = \frac{m_{core} m_{\Lambda}}{m_{core} + m_{\Lambda}}$$

VARIOUS POTENTIAL FOR Λ -N INTERACTION

1. The Maeda–Schmid s-wave potential (MS) (sum of two Woods–Saxon functions)

$$V_{N\Lambda}(r) = \frac{V_{rep}}{1 + \exp((r - R_{rep})/a_{rep})} - \frac{V_{att}}{1 + \exp((r - R_{att})/a_{att})}$$

2. The Isle potential (sum of two Gaussians)

$$V_{N\Lambda}(r) = V_R \exp\left[-\left(\frac{r}{b_R}\right)^2\right] - V_A \exp\left[-\left(\frac{r}{b_A}\right)^2\right]$$

3. $U(r) = -D \left[\text{Cosh}^2\left(\frac{r}{R}\right) \right]^{-1}$

4. Woods-saxon Potential $V(r) = -\frac{V_0}{1 + \exp((r - R_{cen})/a_{cen})}$

[1] S. Maeda *et al.* Few-Body Problems in Physics, vol. II, Elsevier, Amsterdam, 1984, 379.

[2] Y. Kurihara, Y. Akaishi, H. Tanaka, Phys. Rev. C 31 (1985) 971.

[3] C G Koutroulos, J. Phys. G. Nucl. Part. Phys. 17 (1991) 1069-1076

[4] C. H. Cai *et al.* Europhys. Lett., 64 (4), pp. 448–453 (2003)

THE PHENOMENOLOGICAL POTENTIAL

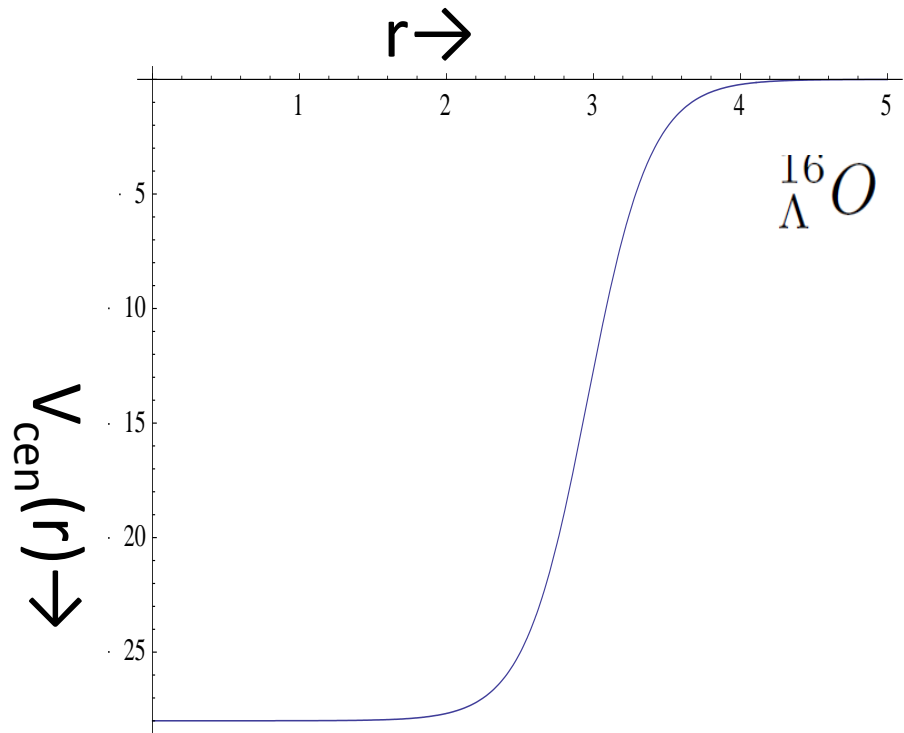
$$V_{cen}(r - R_{cen}) = \frac{-V_0}{1 + e^{\frac{r - R_{cen}}{a_{cen}}}}$$
$$V_0 = U_0 + U_1 \frac{A - 2Z}{A}$$
$$R_{cen/so} = r_0 A^{\frac{1}{3}}$$

$$r_0 = 1.2 \text{ fm}; a_{so} = a_{cen} = 0.215 \text{ fm}$$

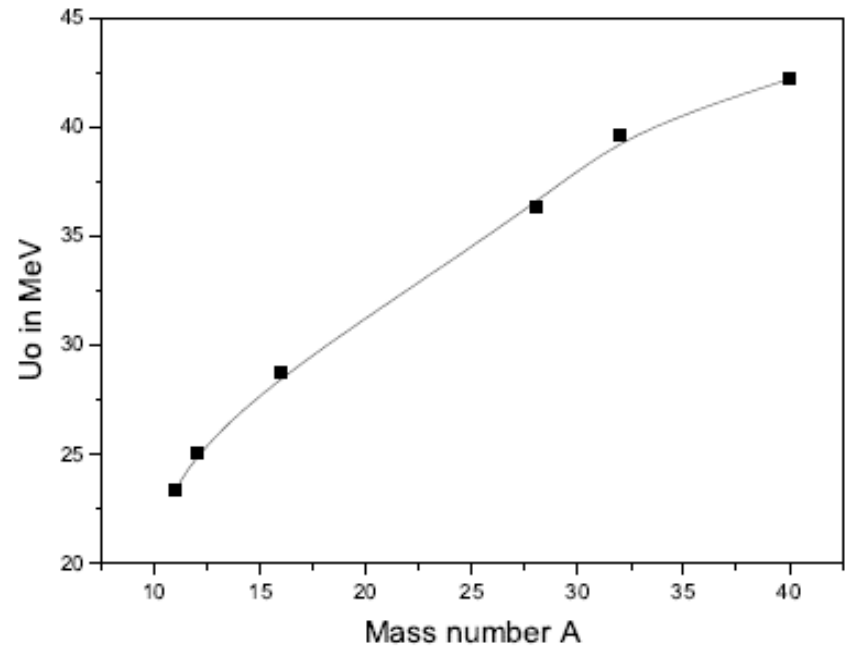
$$U_1 = -0.075 \text{ MeV}; U_{so} = 3.75 \text{ MeV}$$

$$V_{so}(r) = -\frac{2 U_{so}}{r a_{so}} [j(j + 1) - l(l + 1) - s(s + 1)] * \frac{\text{Exp}\left[\frac{r - R_{so}}{a_{so}}\right]}{\left[1 + \text{Exp}\left[\frac{r - R_{so}}{a_{so}}\right]\right]^2}$$

C. H. Cai et al. Europhys. Lett., 64 (4), pp. 448–453 (2003)



Nature of Woods-Saxon Potential



Variation of Potential depth Vs. A

RESULT

TABLE I: Binding Energy ($-B_\Lambda$) of S and P-states Λ Hypernuclei in MeV.

Nuclei	BE for S-State in MeV			BE for P-State in MeV		
	Present	Other	Expt. [5]	Present	Other	Expt. [5]
${}^{11}_\Lambda B$	10.029	10.163 [6] 10.28 [7]	10.20			
${}^{12}_\Lambda C$	10.8076	10.928 [6] 10.97 [7]	10.80			
${}^{16}_\Lambda O$	12.5031	13.243 [6] 13.15 [7]	12.50	5.812	2.544 [6] 1.92 [7]	2.50
${}^{28}_\Lambda Si$	16.0025	16.930 [6] 16.95 [7]	16.00	7.490	8.099 [6] 6.039 [7] 5.83 [8]	7.00
${}^{32}_\Lambda S$	17.5251	17.665 [6] 17.77 [7]	17.50	8.220	9.324 [6] 7.047 [7] 7.28 [8]	8.10
${}^{40}_\Lambda Ca$	18.7022	18.785 [6] 19.09 [7]	18.70	8.780	11.248 [6] 8.718 [7] 9.56 [8]	11.00

6) C. H. Cai, L. Li, Y. H. Tan and P. Z. Ning, *Europhys. Lett.*, **64** (4): 448-453 (2003).

7) C G Koutroulos, *J. Phys. G. Nucl. Part. Phys.* **17**: 1069-1076, (1991).

ALPHA DECAY OF HYPERNUCLEI

we extend our understanding of the decay theory of the normal α -emitter nuclei to study the decay properties of heavy hypernuclei such as

1. Q -value for the decay,
2. Tunnelling probabilities,
3. Half life time and
4. Decay constant of heavy Λ -hypernuclei

ALPHA DECAY OF HYPERNUCLEI

Alpha (hyper-alpha) decay of a hypernucleus can be expressed as



From the Q value it is found that spontaneous decay of hyper- α from the hypernuclei are not energetically possible.

Q-VALUE FOR THE α DECAY

To calculate the Q-value for α decay of heavy hypernuclei we should know the value of BE and SE of hypernuclei.

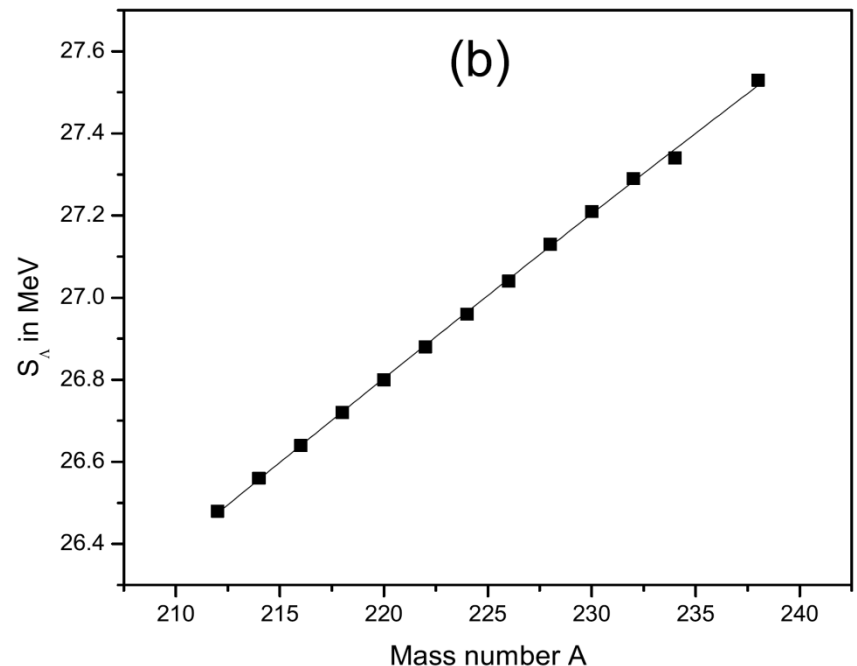
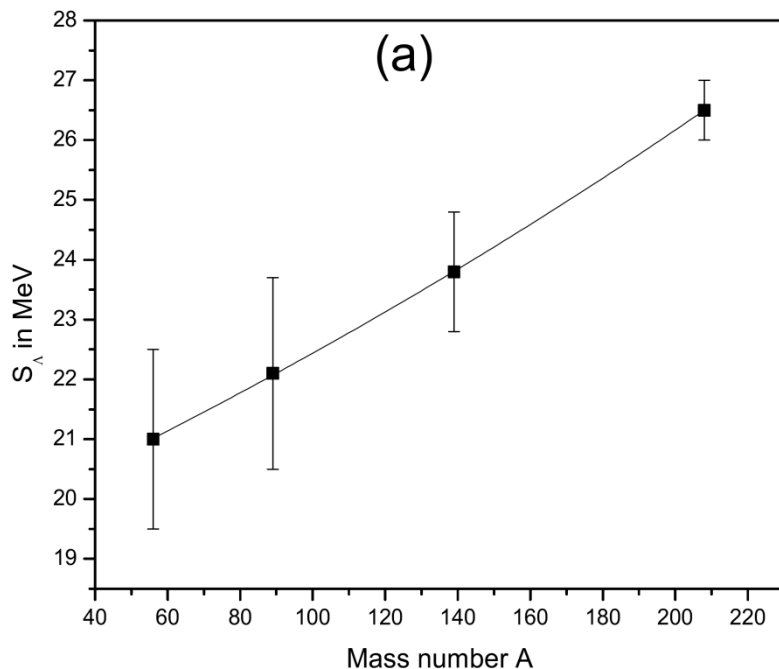
$$\text{BE} \left(\begin{smallmatrix} A_Y \\ Z_Y \end{smallmatrix} X \right) = \text{BE} \left(\begin{smallmatrix} A \\ Z \end{smallmatrix} X \right) + S_Y \left(\begin{smallmatrix} A_Y \\ Z_Y \end{smallmatrix} X \right)$$

The first term corresponding to the binding energy of core nucleus (excluding hyperon) and the second term indicates separation energy of hyperon in the hypernucleus.

SE OF HPERNUCLEI

- using the experimentally known Λ -separation energies of various hypernuclei we have obtained an empirically fitted formula for the lambda separation energy for $A \geq 56$.

$$S_{\Lambda} = 19.27 e^{0.0015 A}$$



TUNNELLING PROBABILITY

$$T = \text{Exp} \left[-\frac{2}{h} \int_b^R [2 m_\alpha V(r) - E_\alpha]^{0.5} \right] dr$$

$$V(r) = \frac{2 Z_d e^2}{4 \pi \epsilon r} \qquad E_\alpha = \frac{m_d Q}{m_\alpha + m_d}$$

- Z_d - atomic number of (hyper) nucleus.
- m_d - mass of the daughter (hyper) nucleus.
- m_α - mass of the α particle.
- E_α - KE of the α particle.

HALF LIFE TIME

$$T_{1/2} = \frac{\ln 2}{\nu T P_0}$$

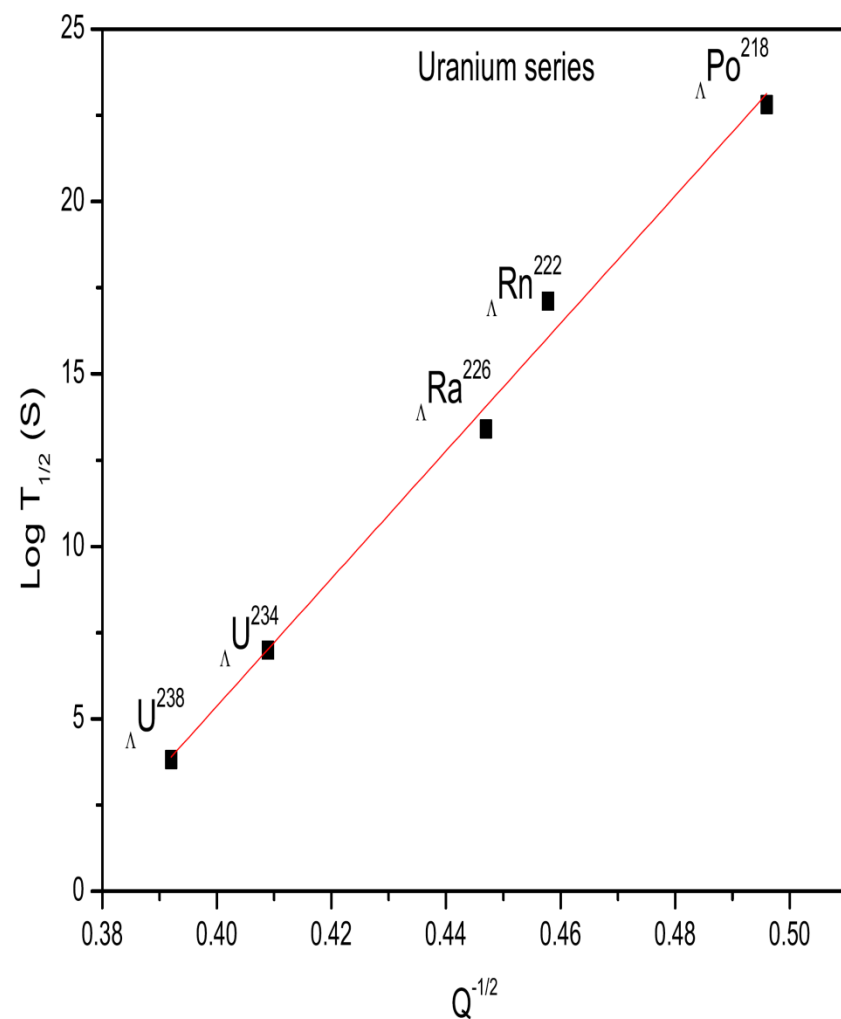
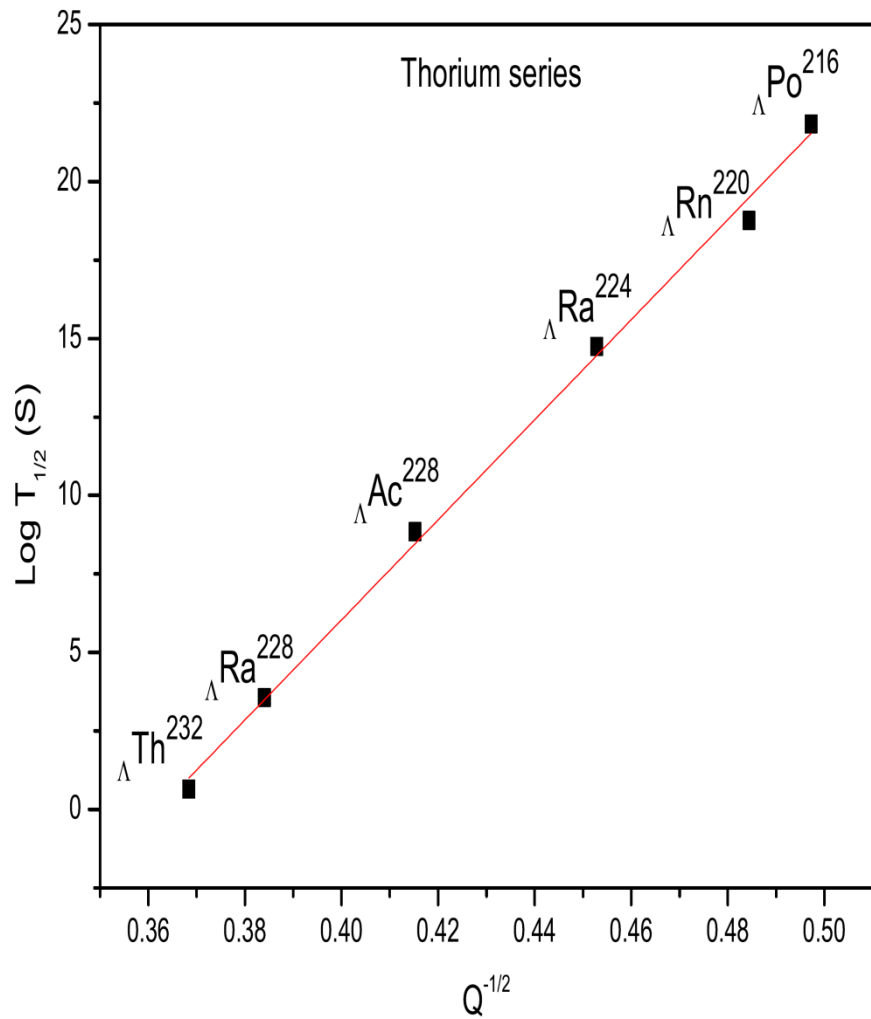
- ν - Frequency of α -collision at the hypernuclear surface.
- T - tunnelling probability.
- P_0 - is the α cluster preformation probability.

RESULT

Reaction	Q	TP	$T_{\frac{1}{2}}$ Sec	$\lambda \text{ Sec}^{-1}$	$\frac{T_{1/2}(\Lambda)}{T_{1/2}(\Lambda=0)}$
${}_{\Lambda}^{232}\text{Th} \rightarrow {}_2^4\text{He} + {}_{\Lambda}^{228}\text{Ra}$	4.045	3.04025×10^{-42}	6.87124×10^{21}	1.00855×10^{-22}	2.17
${}_{\Lambda}^{228}\text{Ra} \rightarrow {}_2^4\text{He} + {}_{\Lambda}^{224}\text{Rn}$	4.262	3.55441×10^{-39}	5.69377×10^{18}	1.21712×10^{-19}	0.042
${}_{\Lambda}^{228}\text{Ac} \rightarrow {}_2^4\text{He} + {}_{\Lambda}^{224}\text{Fr}$	4.877	3.34728×10^{-35}	5.65206×10^{14}	1.2261×10^{-15}	0.406
${}_{\Lambda}^{224}\text{Ra} \rightarrow {}_2^4\text{He} + {}_{\Lambda}^{220}\text{Rn}$	5.800	2.44318×10^{-29}	7.06046×10^8	9.81523×10^{-10}	0.867
${}_{\Lambda}^{220}\text{Rn} \rightarrow {}_2^4\text{He} + {}_{\Lambda}^{216}\text{Po}$	6.782	4.29783×10^{-24}	3.69×10^3	1.87792×10^{-4}	0.028
${}_{\Lambda}^{216}\text{Po} \rightarrow {}_2^4\text{He} + {}_{\Lambda}^{212}\text{Pb}$	7.368	3.37336×10^{-21}	4.48419	0.154543	0.0233

Reaction	Q	TP	$T_{\frac{1}{2}}$ Sec	$\lambda \text{ Sec}^{-1}$	$\frac{T_{1/2}(\Lambda)}{T_{1/2}(\Lambda=0)}$
${}_{\Lambda}^{238}\text{U} \rightarrow {}_2^4\text{He} + {}_{\Lambda}^{234}\text{Th}$	4.060	3.28743×10^{-43}	6.38966×10^{22}	1.08456×10^{-23}	59.523
${}_{\Lambda}^{234}\text{U} \rightarrow {}_2^4\text{He} + {}_{\Lambda}^{230}\text{Th}$	4.770	1.47727×10^{-37}	1.30582×10^{17}	5.30701×10^{-18}	4.28
${}_{\Lambda}^{226}\text{Ra} \rightarrow {}_2^4\text{He} + {}_{\Lambda}^{222}\text{Rn}$	5.000	7.23074×10^{-34}	2.57677×10^{13}	2.68941×10^{-14}	0.13875
${}_{\Lambda}^{222}\text{Rn} \rightarrow {}_2^4\text{He} + {}_{\Lambda}^{218}\text{Po}$	5.980	1.71563×10^{-27}	9.87359×10^6	7.01872×10^{-8}	0.0108
${}_{\Lambda}^{218}\text{Po} \rightarrow {}_2^4\text{He} + {}_{\Lambda}^{214}\text{Pb}$	6.500	2.43351×10^{-24}	6.637×10^3	1.044×10^{-4}	2.71×10^{-3}

RESULT



SUMMARY

We find that the hypernuclei are stable against the hyper-alpha decay because the Λ -binding energy of hyper-alpha is much smaller than those with heavy hypernuclei, that reduces the formation of hyper-alpha cluster within the heavy hypernuclei.

It is found that the presence of a hyperon, Λ inside the normal α -emitter could change their life times considerably.

For instance, the half life times of ${}^{232}_{\Lambda}Th$, ${}^{234}_{\Lambda}U$ and ${}^{238}_{\Lambda}U$ are found to be 2, 4 and 60 times larger than their normal counterpart of ${}^{232}Th$, ${}^{234}U$ and ${}^{238}U$ respectively.

WORK PLAN

We found that the heavy hypernucleus are stable against the α cluster decay. In future we will plan to study the decay of heavy cluster like ${}^{12}_Y C$, ${}^{16}_Y O$, ${}^{28}_Y Si$ etc,. from heavy hypernucleus.

We will try to study the fusion of light hypernuclei and will study the change in Coulomb energy.

Thank You