GAUSSIAN APPROXIMATION WITHIN LINEAR SIGMA MODELS

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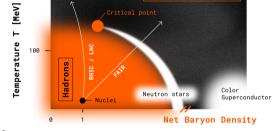


IN COLLABORATION WITH:

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QCD phase diagram Large density: Effective models



Quarks and Gluons

Mean-field approximation

pros:

- Clear and simple
- It works well! (mostly)

How to go beyond?

- Functional methods
- · Add further correction, e.g. in a gaussian approximation

What will be shown is one way to go beyond mean-field

cons:

200

- It can miss something important
- It can generate nonphysical effects

Start with a simple Yukawa model (1 meson + 1 fermion):

$$\mathcal{L}_{Y} = \frac{1}{2} \partial_{\mu} \phi \partial^{\mu} \phi - V_{cl}(\phi) + \bar{\psi}(i\partial - g\phi)\psi, \qquad V_{cl}(\phi) = \frac{1}{2}m^{2}\phi^{2} + \frac{\lambda}{4!}\phi^{4}$$

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To get the thermodynamics:

partition function $Z \Rightarrow$ effective potential Ω from $Z = e^{-i\mathcal{V}_4\Omega}$

Partition function

$$Z = \int \mathcal{D}ar{\psi}\mathcal{D}\psi\mathcal{D}\phi e^{i\int_{\mathsf{X}}\mathcal{L}_{\mathsf{Y}}(\mathsf{X})}$$

Integrate out fermions

$$\int \mathcal{D}\bar{\psi}\mathcal{D}\psi e^{i\int_{X}\bar{\psi}(i\mathcal{S}^{-1})\psi} = \operatorname{Det}(i\mathcal{S}^{-1}) \qquad \Rightarrow \qquad Z = \int \mathcal{D}\phi e^{iS}, \qquad (S = S_m + S_f)$$

$$\phi \to \Phi + \varphi$$

Φ: homogeneous background, φ : fluctuating field

 $Z = \int \mathcal{D}\phi e^{iS(\phi)}$

 $\phi
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 $Z = \int \mathcal{D}\phi e^{iS(\phi)}$

Mean-field approximation: Set $\varphi \rightarrow 0 \Rightarrow Z = e^{iS(\Phi)}$

$$S(\Phi) = \int_{X} (-\frac{1}{2}m^{2}\Phi^{2} - \frac{1}{4!}\Phi^{2}) + S_{f}(\Phi)$$

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 $Z = \int \mathcal{D}\phi e^{iS(\phi)}$

Mean-field approximation: Set $\varphi \rightarrow 0 \Rightarrow Z = e^{iS(\Phi)}$

$$\Omega = \frac{1}{2}m^2\Phi^2 + \frac{1}{4!}\Phi^2 + i\mathrm{tr}\int_k \log(i\mathcal{S}^{-1})$$

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$$\Omega = \frac{1}{2}m^2\Phi^2 + \frac{1}{4!}\Phi^2 + i\operatorname{tr}\int_k \log(i\mathcal{S}^{-1})$$

To go beyond: Expansion in ϕ

$$S(\Phi + \varphi) = S(\Phi) + \varphi \left. \frac{\delta S(\phi)}{\delta \phi} \right|_{\phi = \Phi} + \varphi \left. \frac{\delta S(\phi)}{\delta \phi \delta \phi} \right|_{\phi = \Phi} \varphi + \dots$$

Gaussian integral for mesons

 $Z = \int \mathcal{D}\phi e^{i\mathsf{S}(\phi)}$

Gaussian approximation

$$\int \mathcal{D}\varphi \ e^{\frac{i}{2}\int_{x}\varphi(i\mathcal{G}^{-1})\varphi} = \operatorname{Det}(i\mathcal{G}^{-1})^{-\frac{1}{2}}$$

$$\Omega = \frac{1}{2}m^{2}\Phi^{2} + \frac{1}{4!}\Phi^{2} + i\operatorname{tr}\int_{k}\log(i\mathcal{S}^{-1}) - \frac{i}{2}\operatorname{tr}\int_{k}\log(i\mathcal{G}^{-1})$$

- What is the small parameter?
- How we get closer to the physical case?

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No small parameter We are always "there"

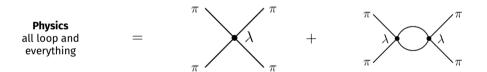
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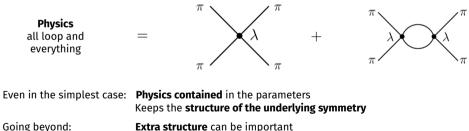
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Extra structure can be important

Simple models with:

- Mesons as main degrees of freedom
- Chiral symmetry as basic structure
- Constituent quarks in a Yukawa type term

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Lagrangian:

$$\begin{split} \mathcal{L} &= \frac{1}{2} \mathrm{Tr} \left(\partial_{\mu} \phi^{\dagger} \partial^{\mu} \phi \right) - \frac{1}{2} m^{2} \mathrm{Tr} \left(\phi^{\dagger} \phi \right) - \lambda_{1} \left(\mathrm{Tr} (\phi^{\dagger} \phi) \right)^{2} - \lambda_{2} \mathrm{Tr} \left((\phi^{\dagger} \phi)^{2} \right) \\ &+ \mathrm{Tr} (H\phi) + c_{1} \left(\mathrm{Det} (\phi) + \mathrm{Det} (\phi^{\dagger}) \right) + \bar{\psi} \left(i \partial - g \mathcal{M} \right) \right) \psi \end{split}$$

with $\phi = {\sf S} + i$ P, containing pion, (kaon,) eta, sigma, etc..., while ${\cal M} = {\sf S} + i \gamma_5 {\sf P}$

Simple models with:

- Mesons as main degrees of freedom
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Lagrangian:

$$\mathcal{L} = \frac{1}{2} \operatorname{Tr} \left(\partial_{\mu} \phi^{\dagger} \partial^{\mu} \phi \right) - \frac{1}{2} m^{2} \operatorname{Tr} \left(\phi^{\dagger} \phi \right) - \lambda_{1} \left(\operatorname{Tr}(\phi^{\dagger} \phi) \right)^{2} - \lambda_{2} \operatorname{Tr} \left((\phi^{\dagger} \phi)^{2} \right)$$
$$+ \operatorname{Tr}(H\phi) + c_{1} \left(\operatorname{Det}(\phi) + \operatorname{Det}(\phi^{\dagger}) \right) + \bar{\psi} \left(i\partial - g\mathcal{M} \right) \right) \psi$$

with $\phi = {\sf S} + i$ P, containing pion, (kaon,) eta, sigma, etc..., while ${\cal M} = {\sf S} + i \gamma_5 {\sf P}$

- + Vector and axial-vector mesons (ELSM)
- + Polyakov loop to mimic the confinement
- + Further hadronic fields to describe more dofs

Thermodynamics: Mean-field level effective potential

- Classical potential.
- Fermionic one-loop correction (vacuum part properly renormalized)

$$\Omega_{MF}(T, \mu_q) = V_{cl} + i \mathrm{tr} \int_{K} \log \left(i S_0^{-1}(T, \mu_q) \right)$$

Field equations (FE):

$$\frac{\partial \Omega_{MF}}{\partial \phi_{N}} = \frac{\partial \Omega_{MF}}{\partial \phi_{S}} = \frac{\partial \Omega_{MF}}{\partial \bar{\Phi}} = \frac{\partial \Omega_{MF}}{\partial \Phi} = 0$$

Curvature meson masses:

$$M_{(MF)ab}^{2} = \left. \frac{\partial^{2} \Omega_{MF}}{\partial \varphi_{a} \partial \varphi_{b}} \right|_{\{\varphi_{i}\}=0}$$

Thermodynamics: Gaussian effective potential

- Classical potential.
- Fermionic one-loop correction (vacuum part properly renormalized)
- Meson fluctuations are included (only the pion for now and thermal fluctuations)

$$\Omega_{G}(T,\mu_{q}) = \Omega_{MF}(T,\mu_{q}) - \frac{i}{2} \operatorname{tr} \int_{K} \log \left(i G^{-1}(T,\mu_{q},M_{(MF)}^{2}) \right)$$

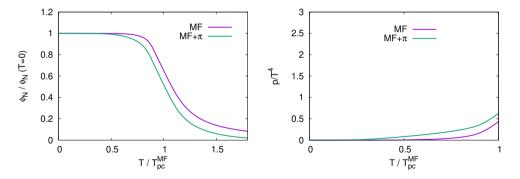
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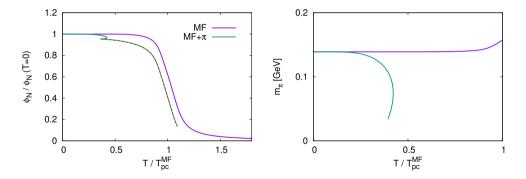
$$\mathsf{M}^{2}_{(\mathsf{MF})ab} = \left. \frac{\partial^{2} \Omega_{\mathsf{MF}}}{\partial \varphi_{a} \partial \varphi_{b}} \right|_{\{\varphi_{i}\}=0}$$

The transition temperature decreases (good news for ELSM) The mesonic pressure corrects the $T < T_{pc}$ behavior



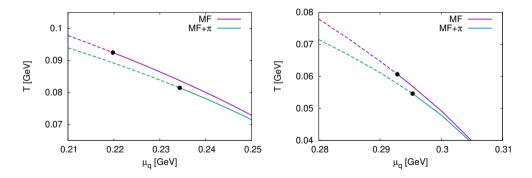
What can go wrong? The pion mass!

This is non-physical



A slight shift in the critical endpoint Note: the meson fluctuations have no $\mu_{\rm B}$ dependence

($\mu_B = 3\mu_q$ in this case)



SUMMARY

- Effective models are simple and useful tools to investigate the phase diagram.
- The usual mean-field approximation works surprisingly well.
- We implemented a beyond mean-field, Gaussian approximation.
- The Gaussian approximation slightly modifies the phase diagram but also qualitatively justifies the mean-field results.
- The problem with the pion mass shows that something might be incomplete.
- How to handle properly the mesonic vacuum fluctuations?
- What are the beyond mean-field effects? E.g. what can be seen for the critical scaling?
- Meson fluctuations in the meson masses?

THANK YOU!