

Hybrid Approaches for the Star BES energy range

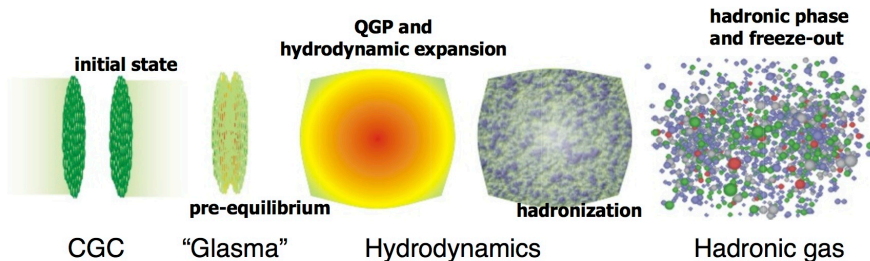
Iurii Karpenko

Faculty of Nuclear Sciences and Physical Engineering, Czech Technical University in Prague



Status quo at high energies (LHC or top RHIC)

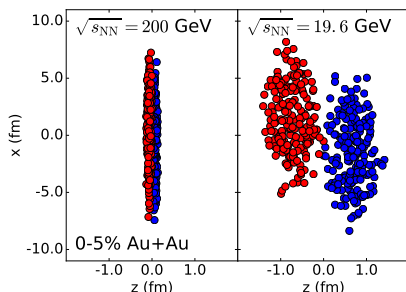
There is a relatively clear time separation between the **initial state** and the **fluid stage**.



When foraging into lower energies using the same tools:

The paradigm of “thin pancakes” gradually loses its applicability.

- There is no boost invariance
- Baryon and electric charge densities are significant
- **Nuclei pass through each other slowly**
(the passage can last as long as subsequent fluid stage)
- There is no clear separation of the initial state and the fluid stage.



picture credit: C. Shen, B. Schenke, Phys. Rev. C 97, 024907 (2018)

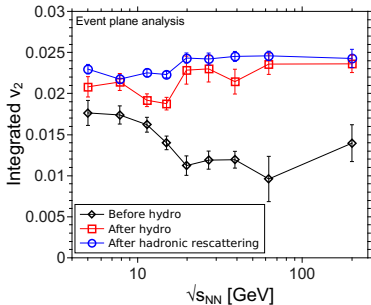
From the last two bullet points:

A lot of evolution is happening before the nuclei have completely passed through each other.

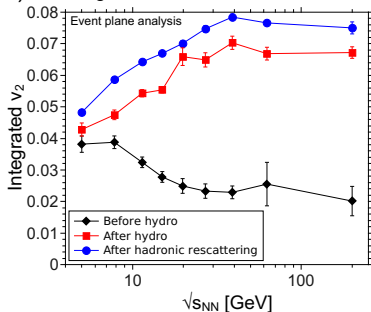
Simulation: UrQMD IS + ideal hydro + UrQMD afterburner

J. Auvinen, H. Petersen, Phys.Rev.C 88 (2013), 064908

a) Charged hadrons, $b = 0 - 3.4$ fm



b) Charged hadrons, $b = 8.2 - 9.4$ fm



Not only the initial state model is important, but the way of switching from initial state to fluid-dynamical picture.

Taxonomy of existing hybrids for RHIC BES and FAIR energies

There is abundance of models for $\sqrt{s_{NN}} = 7 \dots 62$ GeV!

Existing hydro/hybrid models can be classified as follows:

Initial state

- 1 Parametrized initial state + hydro
- 2 Transport + hydro, switch at fixed t or τ
- 3 EPOS
- 4 Transport + hydro, switch dynamically
- 5 Multi-fluid models

Equation of state

- Chiral model EoS
- BM + HRG EoS
- NEOS-B, NEOS-BQS
- BEST EoS

Final conditions

- Particlization at fixed energy density + hadronic cascade (except for multi-fluid models)

All the models feature 3D initial state, 3D hydro and EoS at finite n_B .

Existing hybrids do not contain criticality, and models with criticality aren't hybrids.

EoS in use (1)

All-crossover cases:

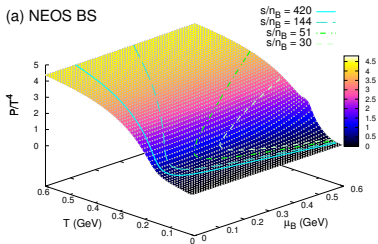
NEOS-B, NEOS-BQ, NEOS-BQS:

- lattice QCD at $\mu_B = 0$
- Taylor expansion in μ_i/T using susceptibilities
- free hadron-resonance gas at low T
- crossover transition at all μ_B

Monnai, Schenke, Shen,

Phys. Rev. C 100, 024907 (2019)

(a) NEOS BS

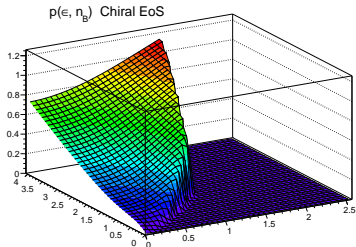


Chiral model EoS

- a single model for hadronic and quark phases
- hadronic SU(3) model
- extension to quark DOF in analogy to PNJL
- crossover transition at all densities

Steinheimer, Schramm, Stocker,

J. Phys. G 38 035001

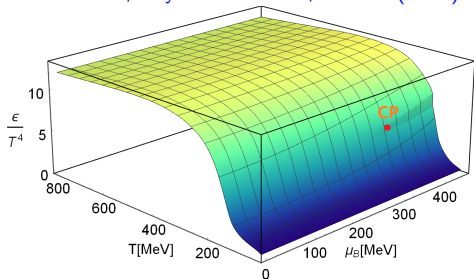


EoS in use (2)

BEST EoS

- lattice QCD at $\mu_B = 0$
- Tylor expansion in μ_i/T using susceptibilities
- critical contribution inspired by a 3D ising model

Parotto et al, [Phys. Rev. C 101, 034901 \(2020\)](#)

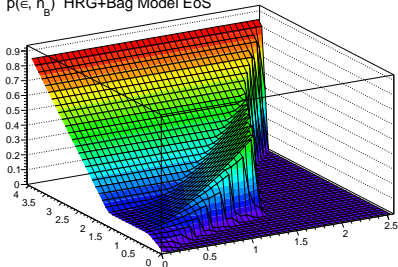


Bag model + HRG EoS a.k.a. EoS Q

- Bag model for QGP phase
- HRG model with repulsion
- Maxwell construction \Rightarrow 1st order PT.

P.F. Kolb, et al, [Phys.Rev. C 62, 054909 \(2000\)](#)

$p(\epsilon, n_B)$ HRG+Bag Model EoS



Hybrids and their initial states

Type 1: Parametrized initial state

Chun Shen, Sahr Alzhrani, [Phys.Rev.C 102 \(2020\) 1, 014909](#)

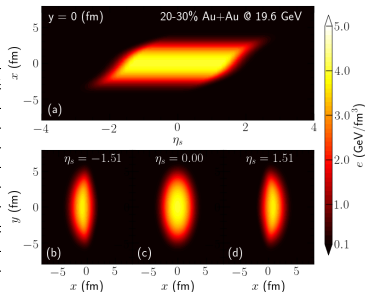
superMC (initial state) + MUSIC (hydro) + iSS (hadron sampling) + UrQMD

$$e(x, y, \eta_s; y_{\text{CM}}) = \mathcal{N}_e(x, y) \times \exp \left[-\frac{(|\eta_s - y_{\text{CM}}| - \eta_0)^2}{2\sigma_\eta^2} \theta(|\eta_s - y_{\text{CM}}| - \eta_0) \right].$$

$$\mathcal{N}_e(x, y) = \frac{M(x, y)}{2 \sinh(\eta_0) + \sqrt{\frac{\pi}{2}} \sigma_\eta e^{\sigma_\eta^2/2} C_\eta} C_\eta = e^{\eta_0} \operatorname{erfc} \left(-\sqrt{\frac{1}{2}} \sigma_\eta \right) + e^{-\eta_0} \operatorname{erfc} \left(\sqrt{\frac{1}{2}} \sigma_\eta \right).$$

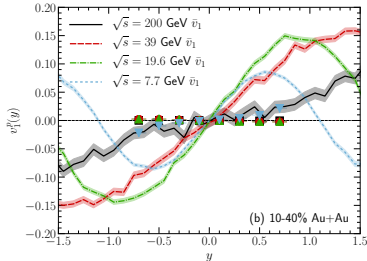
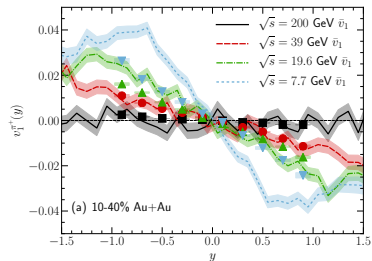
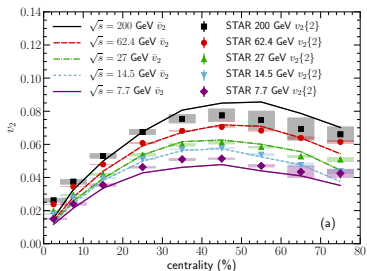
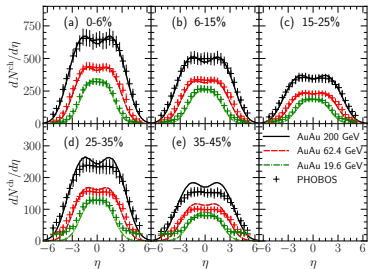
+scaling initial flow, $v_z = z/t$

$\sqrt{s_{\text{NN}}}$ (GeV)	τ_0 (fm/c)	η_0	σ_η	$\eta_{B,0}$	$\sigma_{B,\text{in}}$	$\sigma_{B,\text{out}}$
AuAu & dAu @ 200	1.0	2.5	0.6	3.5	2.0	0.1
AuAu & dAu @ 62.4	1.0	2.25	0.3	2.7	1.9	0.2
AuAu & dAu @ 39	1.3	1.9	0.3	2.2	1.6	0.2
AuAu@27	1.4	1.6	0.3	1.8	1.5	0.2
AuAu & dAu @ 19.6	1.8	1.3	0.3	1.5	1.2	0.2
AuAu@14.5	2.2	1.15	0.3	1.4	1.15	0.2
AuAu@7.7	3.6	0.9	0.2	1.05	1.0	0.1
PbPb@17.3	1.8	1.25	0.3	1.6	1.2	0.2
PbPb@8.77	3.5	0.95	0.2	1.2	1.0	0.1



EoS: NEOS-BQS, [1902.05095](#).

Carve the initial state to fit the data.



Type 2: Transport model for initial state

UrQMD (initial state) + fluidisation at fixed τ_0 + vHLLC (hydro) + UrQMD
IK, Huovinen, Petersen, Bleicher, [Phys.Rev. C91 \(2015\) no.6, 064901](#)

- Initial state from $t = 0$ till $\tau = \tau_0$: UrQMD
- At $\tau = \tau_0$:
Gaussian smearing of energy, momentum and charges at fluidization:

$$\Delta P_{ijk}^\alpha = P^\alpha \cdot C \cdot \exp\left(-(\Delta x_i^2 + \Delta y_j^2)/R_\perp^2 - \Delta \eta_k^2 \gamma_\eta^2 \tau_0^2 / R_\eta^2\right)$$

$$\Delta N_{ijk}^0 = N^0 \cdot C \cdot \exp\left(-(\Delta x_i^2 + \Delta y_j^2)/R_\perp^2 - \Delta \eta_k^2 \gamma_\eta^2 \tau_0^2 / R_\eta^2\right)$$

from each hadron that crosses the $\tau = \tau_0$ surface

- Assume that the resulting energy and momentum corresponds to a flowing fluid:

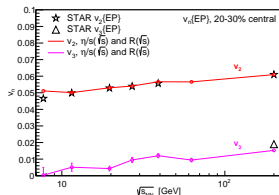
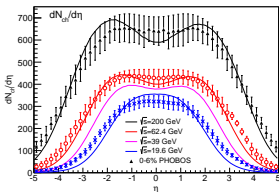
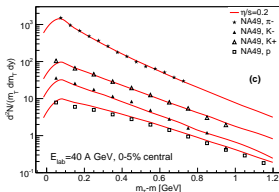
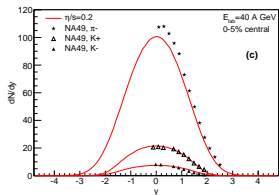
$$E/\Delta V = (\varepsilon + p)(u^0)^2 - p, \quad P^i/\Delta V = (\varepsilon + p)u^0 u^i$$

- Run fluid dynamics.
- ...
- compute the observables

Type 2: Transport model for initial state

UrQMD (initial state) + fluidisation at fixed τ_0 + vHLLC (hydro) + UrQMD
 IK, Huovinen, Petersen, Bleicher, [Phys.Rev. C91 \(2015\) no.6, 064901](#)

Fix parameters of fluidisation procedure (R_{\perp} , R_{η}) and shear viscosity η/s to fit the data.



EoS: chiral model EoS.

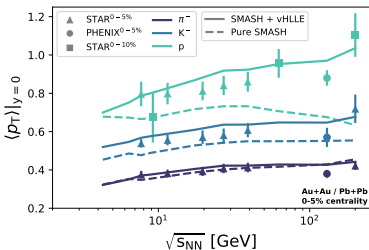
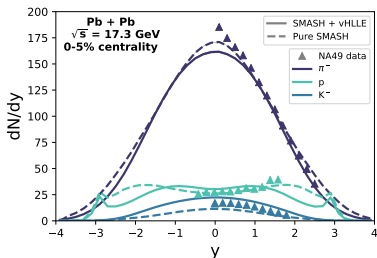
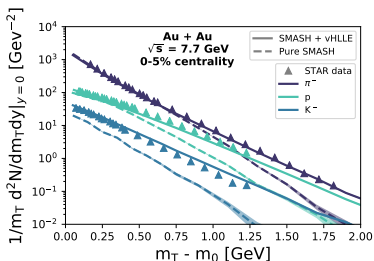
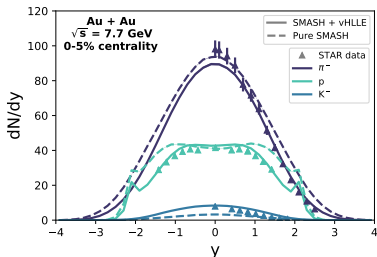
- \Rightarrow decent agreement with a mix of RHIC BES + NA49 + PHOBOS data

Type 2: Transport model for initial state

SMASH (initial state) + fluidisation at fixed τ_0 + vHLLC (hydro) + SMASH

A. Schäfer, IK, Xiang-Yu Wu, J. Hammelmann, H. Elfner, *Eur.Phys.J.A* 58 (2022) 11, 230

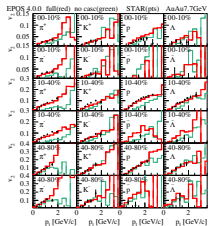
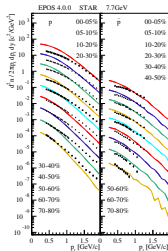
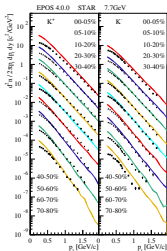
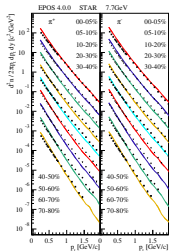
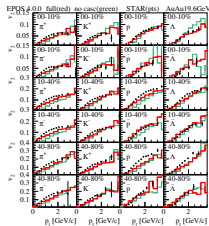
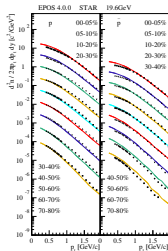
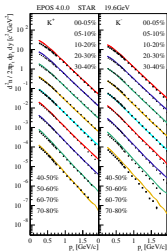
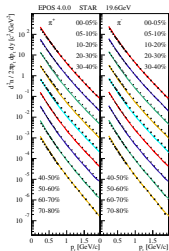
Same set of free parameters.



EPOS 4.0.0

Werner, Jahan, IK, Pierog, Stefaniak, Vintache, [2401.11275](#)

EPOS IS + fluidisation at τ_0 + vHLLC + UrQMD. **core-corona included.** BEST EoS



Type 3: Dynamical fluidisation (strings)

Chun Shen, Bjoern Schenke, [Phys. Rev. C 97, 024907 \(2018\)](#)

Decelerated strings melt into fluid, at different proper times.

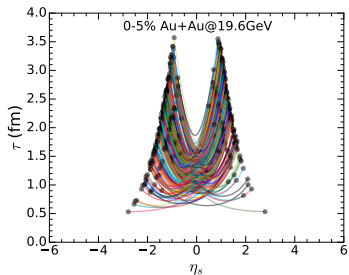
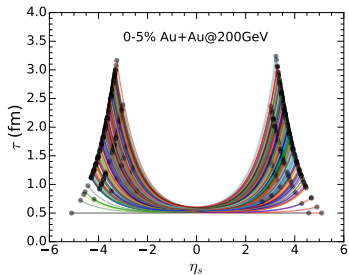
EoS: unclear.

- 1 string per colliding nucleon pair
- The string ends decelerate according to:

$$\frac{dE}{dz} = -\sigma, \quad \frac{dP_z}{dz} = -\sigma$$

- Once string ends come to a halt, its evolution stops.

- The products fluidize.

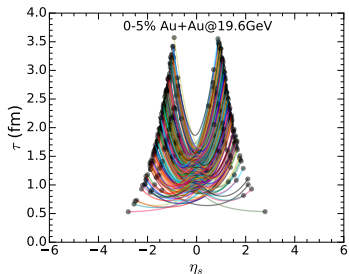
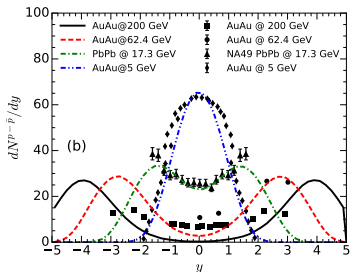
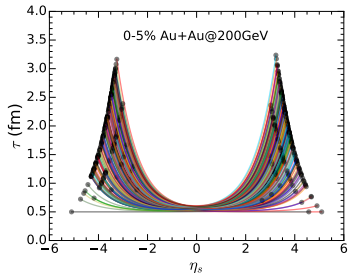
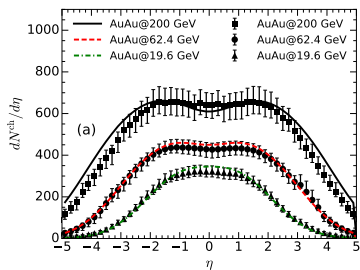


Type 3: Dynamical fluidisation (strings)

Chun Shen, Bjoern Schenke, [Phys. Rev. C 97, 024907 \(2018\)](#)

Decelerated strings melt into fluid, at different proper times.

EoS: unclear.



Type 3: Dynamical fluidisation

JAM + dynamical fluidisation + hydro + JAM

Y. Akamatsu, M. Asakawa, T. Hirano, M. Kitazawa, K. Morita, K. Murase, Y. Nara, C. Nonaka, A. Ohnishi, [Phys. Rev. C 98, 024909 \(2018\)](#)

JAM IS: HIJING string excitation + PYTHIA6 fragmentation + rescatterings of produced hadrons.

Hadrons are converted to fluid if the local energy density $e > e_f = 0.5 \text{ GeV}/\text{fm}^3$.

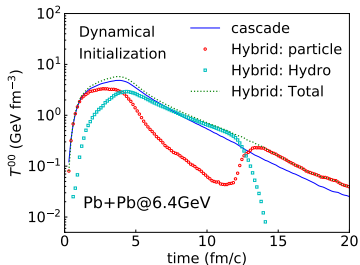
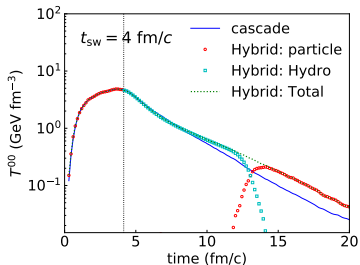
$$\partial_\mu T_f^{\mu\nu} = J^\nu, \quad \partial_\mu N_f^\mu = \rho$$

$$J^\mu(r) = \frac{1}{\Delta t} \sum_i p_i^\mu G(r - r_i(t))$$

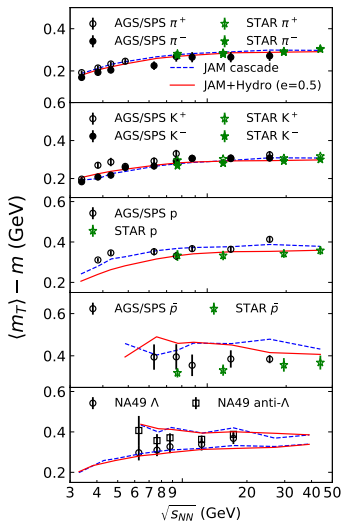
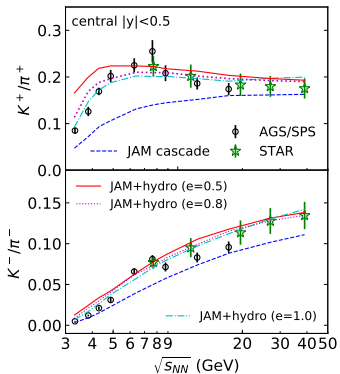
$$\rho(r) = \frac{1}{\Delta t} \sum_i B_i G(r - r_i(t))$$

$G(r)$ is a Gaussian smearing profile.

EoS: EoS Q, a.k.a. BM+HRG EoS



JAM + dynamical fluidisation + hydro + JAM

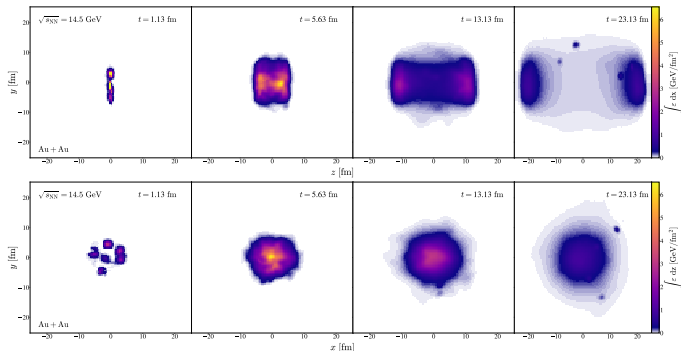
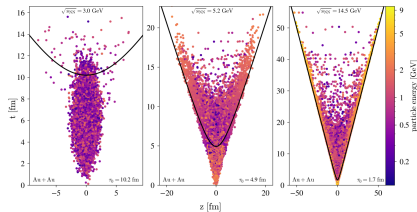


Type 3: Dynamical fluidisation (SMASH cascade)

SMASH + dynamical fluidisation + hydro + SMASH

WIP: Renan Hirayama, Zuzana Paulíniová, IK,
Hannah Elfner

Same idea but with SMASH cascade.



Does transport prove the fluid dynamics is applicable?

Not really.

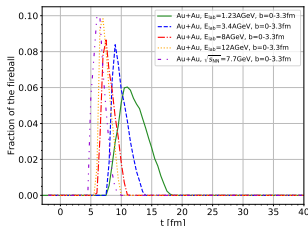
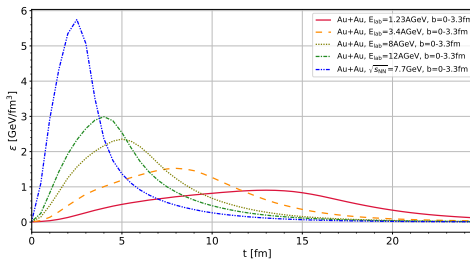
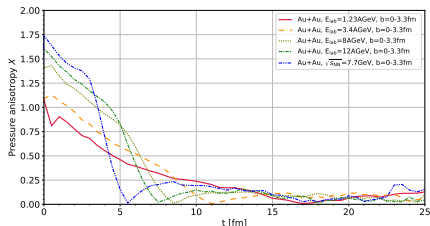
Gabriele Inghirami, Hannah Elfner, *Eur. Phys. J. C* 82, 796 (2022)

From $E_{\text{lab}} = 1.23$ GeV up to $\sqrt{s} = 7.7$ GeV

$$T^{\mu\nu} = \sum_i \frac{p_i^\mu p_i^\nu}{p_i^0} K(\mathbf{r} - \mathbf{r}_i, \mathbf{p}_i)$$

$$X \equiv \frac{|\langle T_L^{11} \rangle - \langle T_L^{22} \rangle| + |\langle T_L^{22} \rangle - \langle T_L^{33} \rangle| + |\langle T_L^{33} \rangle - \langle T_L^{11} \rangle|}{\langle T_L^{11} \rangle + \langle T_L^{22} \rangle + \langle T_L^{33} \rangle},$$

$$Y \equiv \frac{3(|\langle T_L^{12} \rangle| + |\langle T_L^{23} \rangle| + |\langle T_L^{13} \rangle|)}{\langle T_L^{11} \rangle + \langle T_L^{22} \rangle + \langle T_L^{33} \rangle}$$

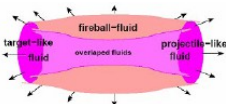


For most of the system, when simulated with transport, $T^{\mu\nu}$ does not look like a fluid.

Type 4: Everything is Fluid



3-Fluid Dynamics



Produced particles
populate mid-rapidity
⇒ fireball fluid

Baryon Stopping

JINR,
24.08.10

Model

Rapidity
Density

Fit

Reduced
curvature

Trajectories

Crossover

Summary

Target-like fluid:

$$\partial_\mu J_t^\mu = 0$$

Leading particles carry bar. charge

$$\partial_\mu T_t^{\mu\nu} = -F_{tp}^\nu + F_{ft}^\nu$$

exchange/emission

Projectile-like fluid:

$$\partial_\mu J_p^\mu = 0,$$

$$\partial_\mu T_p^{\mu\nu} = -F_{pt}^\nu + F_{ip}^\nu$$

Fireball fluid:

$$J_f^\mu = 0,$$

Baryon-free fluid

$$\partial_\mu T_f^{\mu\nu} = F_{pt}^\nu + F_{ip}^\nu - F_{fp}^\nu - F_{ft}^\nu$$

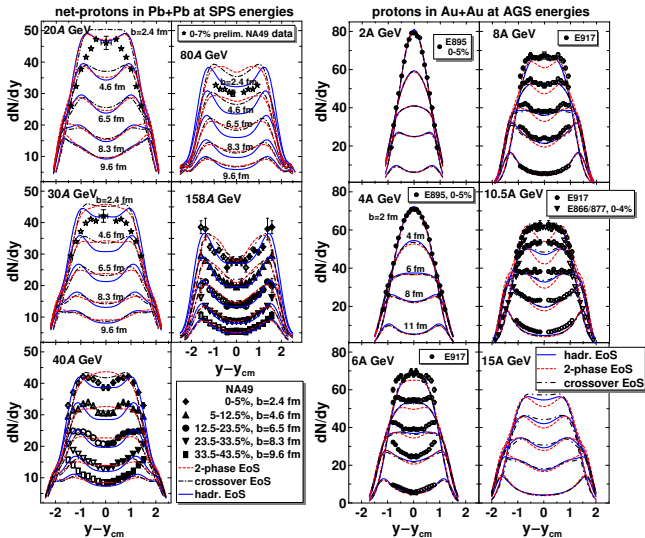
Source term Exchange

The **source term** is delayed due to a formation time $\tau \sim 1 \text{ fm}/c$

Total energy-momentum conservation:

$$\partial_\mu (T_p^{\mu\nu} + T_t^{\mu\nu} + T_f^{\mu\nu}) = 0$$

<http://theory.gsi.de/~ivanov/mfd/>



What also works:

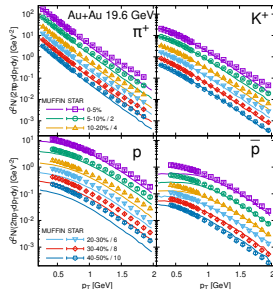
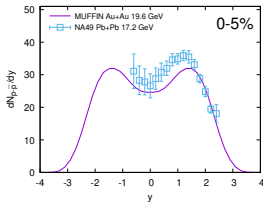
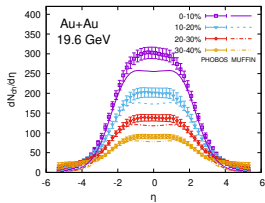
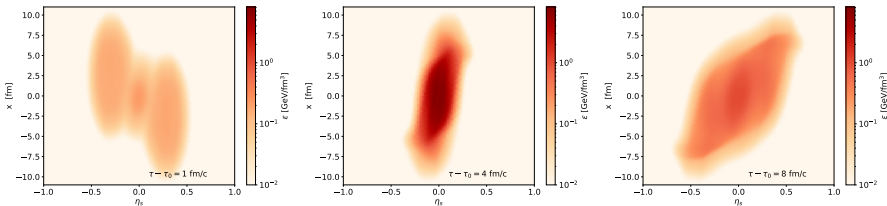
- pion dN/dy
- kaon dN/dp_T
Yu.B. Ivanov, *Phys. Rev. C* 87, 064905
- proton, pion dN/dp_T
Yu.B. Ivanov, *Phys. Rev. C* 89, 024903 (2014)
- elliptic flow
Ivanov, Soldatov, *Phys. Rev. C* 91, 024914 (2015)

Relatively low sensitivity to the PT type in the EoS; rather worse agreement for purely hadronic EoS.

Type 4: Improved 3-fluid dynamics: MUFFIN

3-fluid version of vHLLE + hadron sampling + SMASH

WIP but published version corresponds to Ivanov's friction: Cimerman, IK, Tomasik, Huovinen, *Phys.Rev.C* 107 (2023) 4, 044902

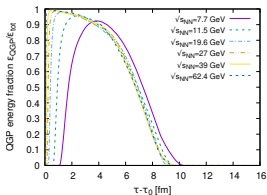


Acceptable data reproduction with chiral model EoS (crossover)

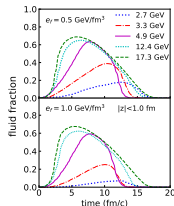
Generic conclusions from the modelling

- Hybrid models generally fit dN/dy , dN/dp_T and v_2 **together**
- high-density medium is formed down to 7.7 GeV. According to EoS it has to be in the QGP phase.

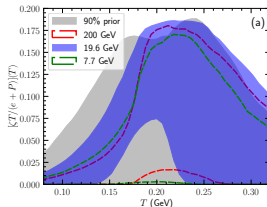
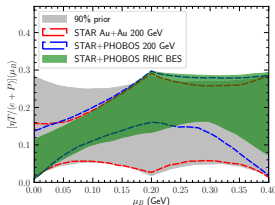
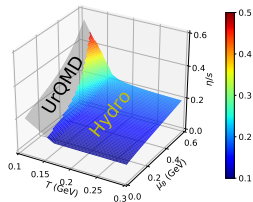
MUFFIN (MFH)



JAM+hydro



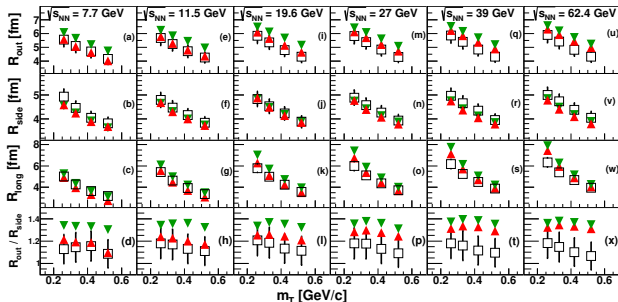
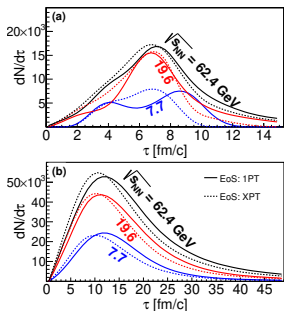
- η/s decreases with T (Geom.IS [2003.05852](#); String deceleration [2310.10787](#)).



Some models (MFH-based) do without viscosity but overestimate v_2 .

EoS sensitivity: HBT?

UrQMD (initial state) + fluidisation at fixed τ_0 + vHLLC (hydro) + UrQMD
Batyuk, IK, Lednický, Malinina, Mikhaylov, Rogachevsky, Wielanek, [Phys. Rev. C 96, 024911 \(2017\)](#)



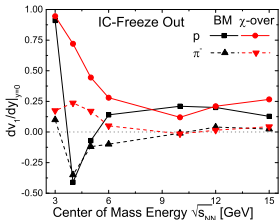
1PT = 1st order PT, XPT = crossover; crossover EoS is red, 1PT EoS is green

There is weak EoS sensitivity, crossover EoS is preferred.

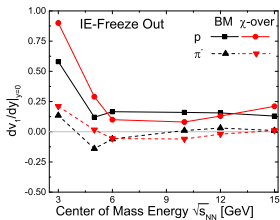
EoS sensitivity: directed flow?

J. Steinheimer, J. Auvinen, H. Petersen, M. Bleicher, H. Stöcker, Phys. Rev. C 89 (2014) 054913, arXiv:1402.7236

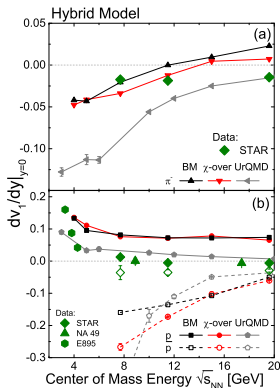
- 1-fluid model with iso-time freezeout:
sign change of dv_1/dy with 1st-order PT EoS



- 1-fluid model with iso-T freezeout:
NO sign change of dv_1/dy with 1st-order PT EoS



- Full hybrid model: no sign change of dv_1/dy , weak EoS dependence and no agreement with the data



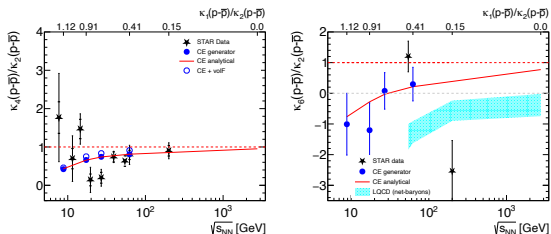
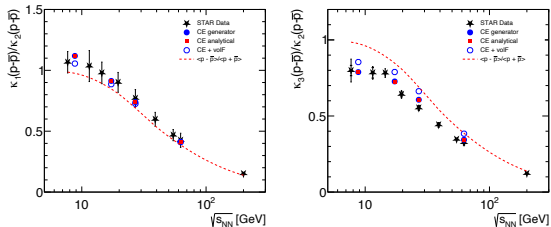
Proton Number Cumulants

Non-critical baseline for proton cumulants is non-trivial

...as a function of collision energy

Braun-Munzinger, Friman, Redlich, Rustamov, Stachel, *Nucl. Phys. A* 1008, 122141 (2021)

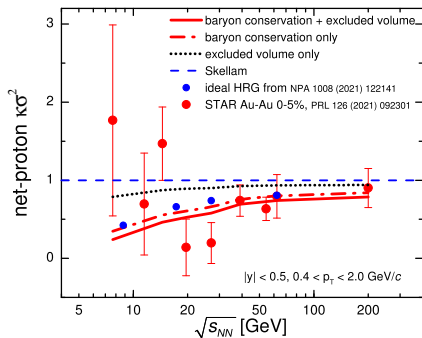
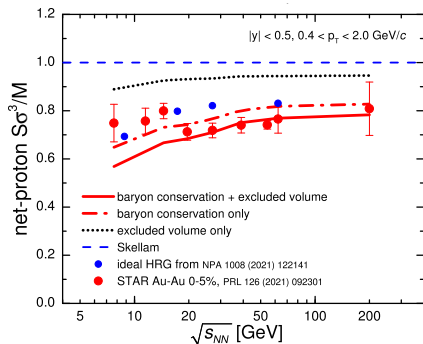
Effect of global baryon number conservation + volume fluctuations



Non-critical baseline from 3D hydro

Vovchenko, Koch, Shen, *Phys. Rev. C* 105, 014904 (2022)

Parametrized initial state (Type 1) + MUSIC
Global baryon conservation using SAM-2.0 method

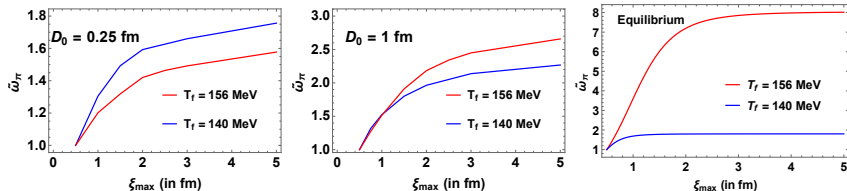


Progress on critical fluctuations

Maneesha Pradeep, Jamie Karthein, Krishna Rajagopal, Misha Stephanov, Yi Yin,
[Phys.Rev.D 106 \(2022\) 3, 036017](#)

Hydro+ : deterministic hydrodynamics + non-hydrodynamic modes ϕ
No concrete result just yet.

$$\tilde{\omega}_\pi = \frac{\omega_\pi}{\omega_\pi^{\text{nc}}}, \text{ where } \omega_A = \frac{\langle \delta N_A^2 \rangle}{\langle N_A \rangle}$$



where D is a diffusion coefficient which inversely scales with the correlation length ξ .

Conclusions

- There is abundance of hybrid models for the RHIC BES energy range, with different assumptions about initial state dynamics (or absence thereof).
- Most of the models describe most of the basic observables.
- EoS sensitivity is generally not strong; it is not clear yet what observables would be very sensitive to the type of PT in the EoS.
- Probably a future Bayesian analysis will tell us whether some EoS are excluded by the coherent data reproduction
- Baryon/proton cumulants: we know a non-critical baseline.
- There are no hybrids with critical behaviour included yet.
- We need more modelling and **more consistent experimental data**.