

Deeply Virtual Compton Scattering at Jlab/CLAS

Baptiste GUEGAN

PANDA collaboration meeting 09/10/12

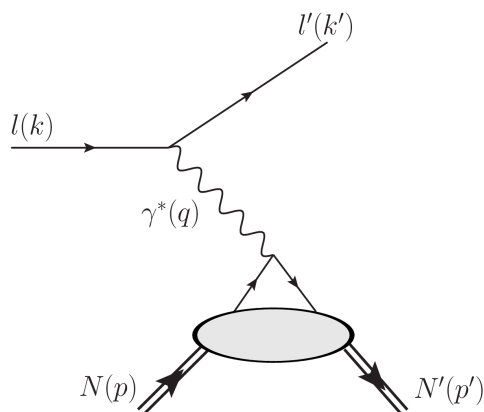
IPN Orsay



Nucleon structure using the electromagnetic probe

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Elastic scattering



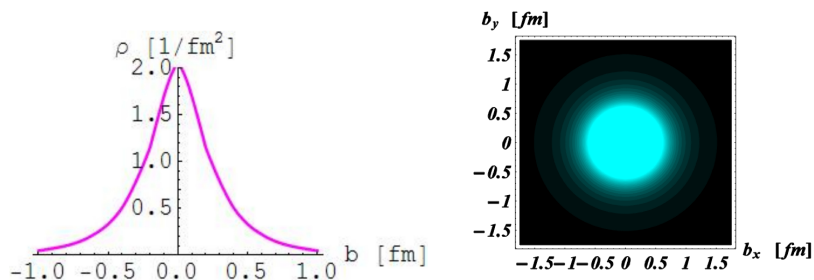
Local non-forward
matrix element:

$$\langle p' | \bar{q}(0) O q(0) | p \rangle$$

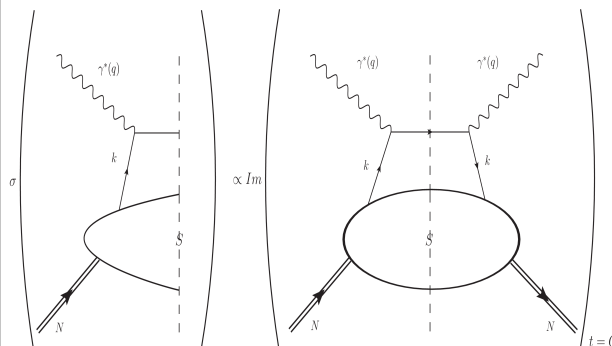
→ Form factors: $F_1(t), F_2(t), G_A(t), G_P(t)$



Interpreted as the transverse localisation
of partons in the nucleon (independently
of their longitudinal momentum)



Deep inelastic scattering



Non-Local forward
matrix element:

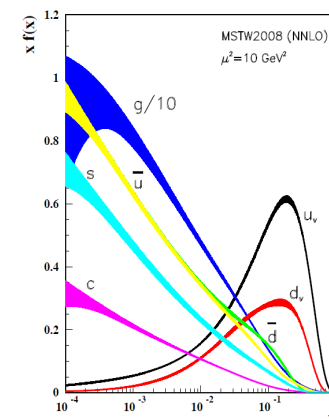
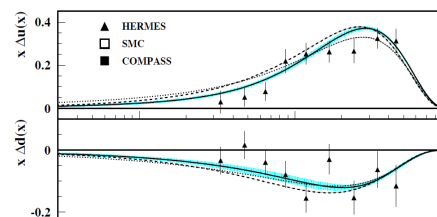
$$\langle p | \bar{q}(0) O q(y) | p \rangle$$

→ Structure functions: $f_2(x), g_1(x)$



Interpreted as the longitudinal momentum
distribution

(no information on the transverse localisation)



Generalized Parton Distributions

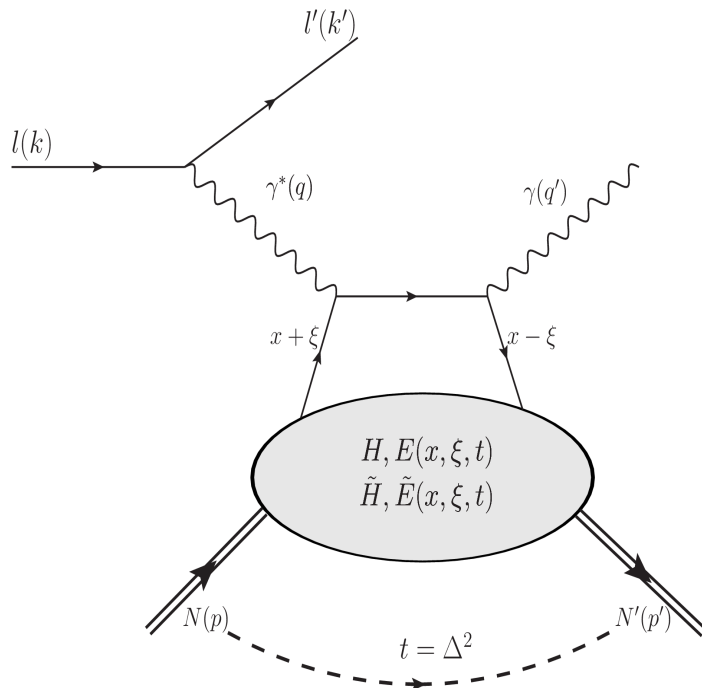
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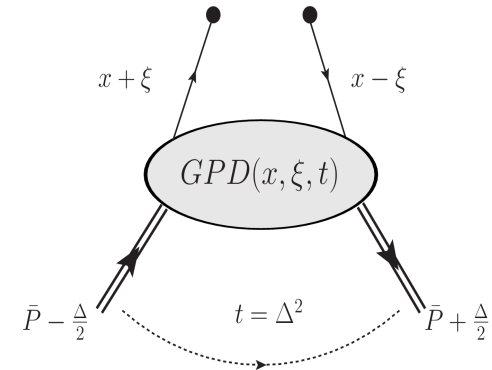
Extension to the non-local, non-forward matrix element:

$$\langle p' | \bar{q}(0) O q(y) | p \rangle$$

Accessed with the Deeply Virtual Compton Scattering :



At large Q^2 , the process can be factorized and described by 4 Generalized Parton Distributions



$(x+\xi)$ and $(x-\xi)$: longitudinal momentum fractions of quarks
 $t = \Delta^2 = (p' - p)^2$: squared momentum transfer

Unpolarized GPDs :

(quark helicity independent)

➡ $H(x, \xi, t)$ Vector

➡ $E(x, \xi, t)$ Tensor

Polarized GPDs:

(quark helicity dependent)

➡ $\tilde{H}(x, \xi, t)$ Axial - Vector

➡ $\tilde{E}(x, \xi, t)$ Pseudoscalar

GPDs properties

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- In the forward limit ($t=\xi=0$), the GPDs reduce to PDFs:

$$H(x,0,0)=q(x) \quad \tilde{H}(x,0,0)=\Delta q(x)$$

No similar equations for E and \tilde{E} . New information !

- The first moment of the GPDs reduce to the FFs:

$$\int_{-1}^1 dx H(x, \xi, t) = F_1(t) \quad \int_{-1}^1 dx E(x, \xi, t) = F_2(t)$$

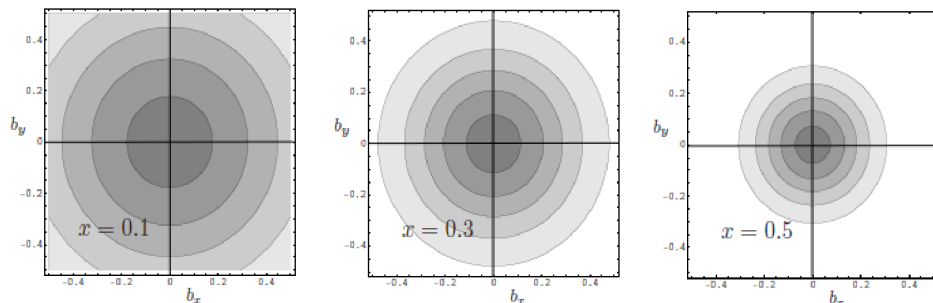
$$\int_{-1}^1 dx \tilde{H}(x, \xi, t) = G_A(t) \quad \int_{-1}^1 dx \tilde{E}(x, \xi, t) = G_P(t)$$

- The second moment of (E+H) when $t \rightarrow 0$: total angular momentum \rightarrow EMC : $\approx 30\%$

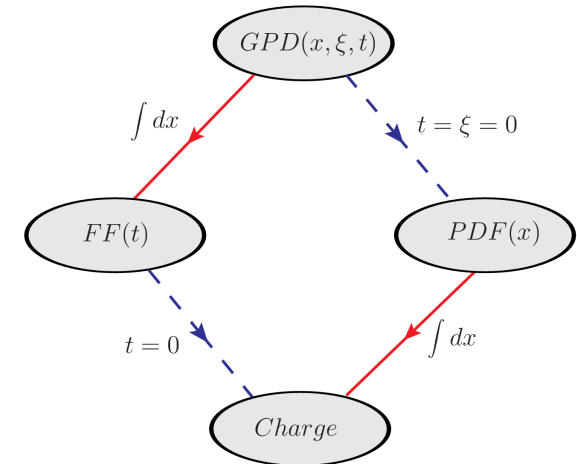
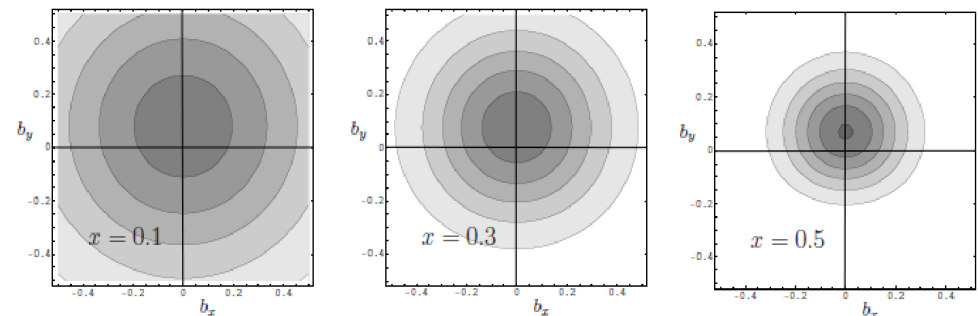
$$\int_{-1}^1 dx x [H^q(x, \xi, 0) + E^q(x, \xi, 0)] = 2 J_{quarks} \quad \frac{1}{2} = (S_{quarks} + L_{quarks}) + J_{gluons}$$

- At $\xi = 0$, a GPD is a x-decomposition of the form factor:

u-quark distribution in a unpolarized proton



u-quark distribution in a polarized proton



Exclusive electroproduction of a photon

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Contribution from both DVCS and Bethe-Heitler (undistinguishable experimentally):

$$\sigma_{(ep \rightarrow ep\gamma)} \propto |T^{DVCS} + T^{BH}|^2$$

$\sigma(ep \rightarrow ep\gamma) =$

DVCS

Bethe-Heitler (BH)

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- \mathcal{T}_{BH} At low t , the nucleon FFs (Dirac, Pauli) are well known so that \mathcal{T}_{BH} is precisely calculable

- $\mathcal{T}^{DVCS} \propto \int_{-1}^1 dx \frac{f(x, \xi, t)}{x \pm \xi \mp i\epsilon} = \mathcal{P} \int_{-1}^1 dx \frac{f(x, \xi, t)}{x \pm \xi} \pm i\pi f(x = \mp \xi, \xi, t)$

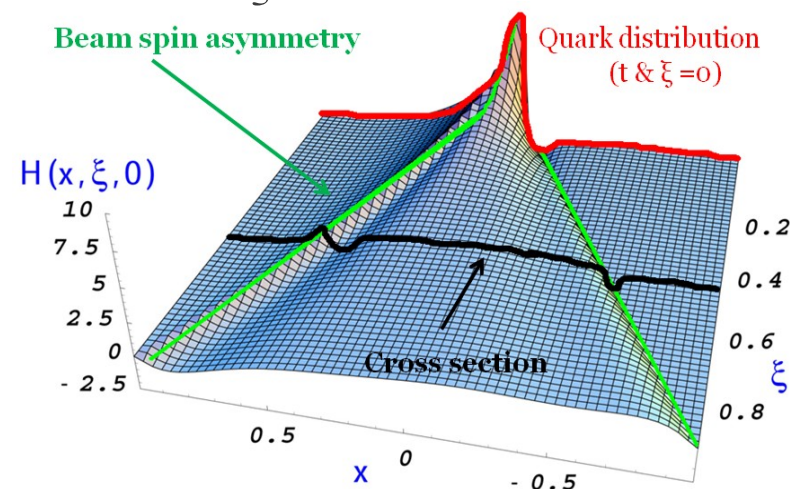
➡ GPDs appear in the **real part** through an integral over x

➡ GPDs appear in the **imaginary part** at the lines $x = \pm \xi$

With a polarized beam and an unpolarized target, one can measure 2 observables:

$$\frac{d^4\sigma}{dtdQ^2 dx_B d\phi} \approx |T^{BH}|^2 + 2T^{BH} \cdot \text{Re}(T^{DVCS}) + |T^{DVCS}|^2$$

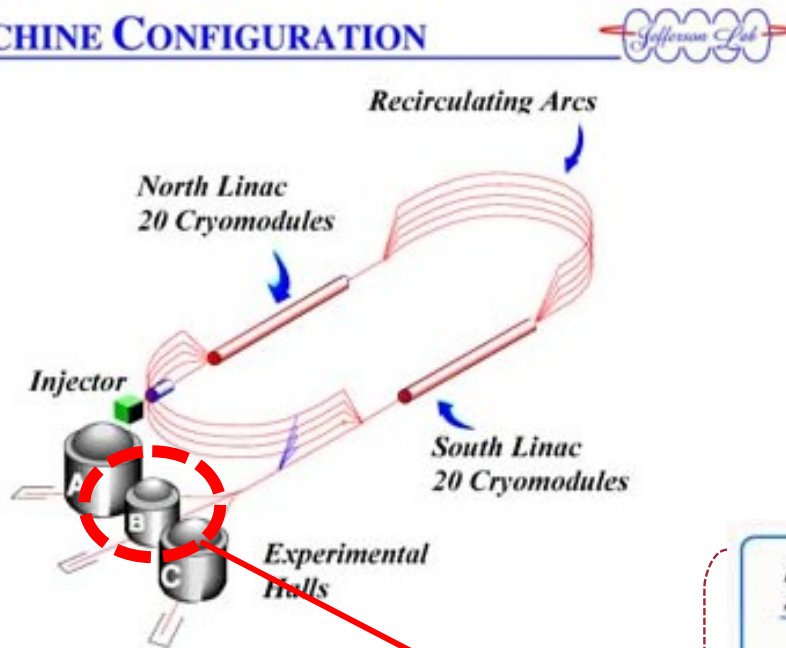
$$\frac{d^4\vec{\sigma} - d^4\overleftarrow{\sigma}}{dtdQ^2 dx_B d\phi} \approx 2T^{BH} \cdot \text{Im}(T^{DVCS}) + \left[|T^{\vec{DVCS}}|^2 - |T^{\overleftarrow{DVCS}}|^2 \right]$$



CLAS (*CEBAF Large Acceptance Spectrometer*) @ JLAB (*Jefferson Laboratory*)

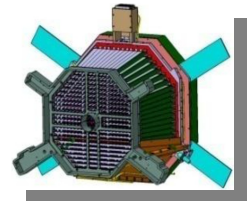
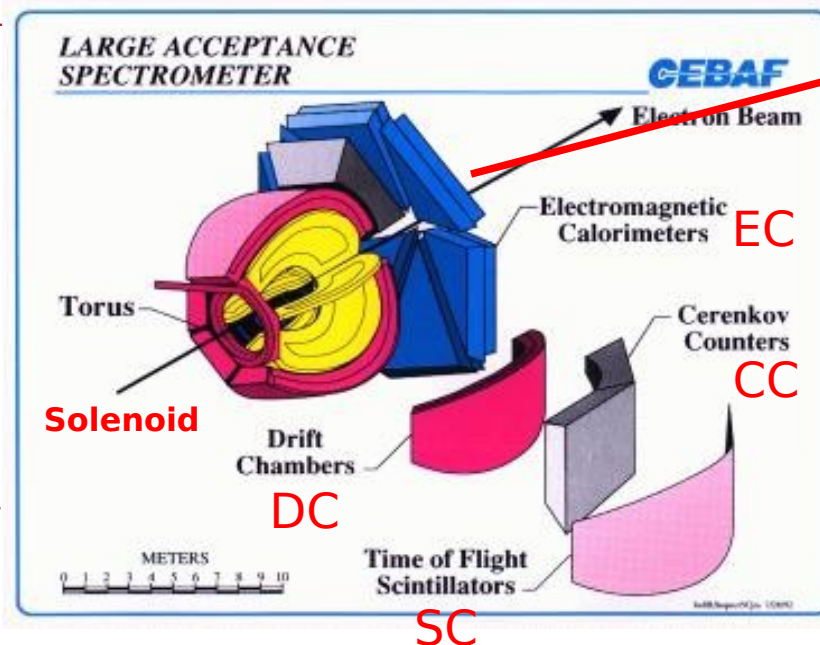
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MACHINE CONFIGURATION



- 1500 physicists, ~30 countries, operational since end 97
- Two LINACs, 3 experimental halls
- $I_{\text{max}}=20 \text{ nA}$, $E_{\text{max}}=6 \text{ GeV}$, 100% duty cycle
- Luminosity: Hall B= $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$

- Large acceptance detector
- Toroidal magnetic field (2,5 T)



DVCS photons are mostly emitted at forward angles

Identification of the final state particles

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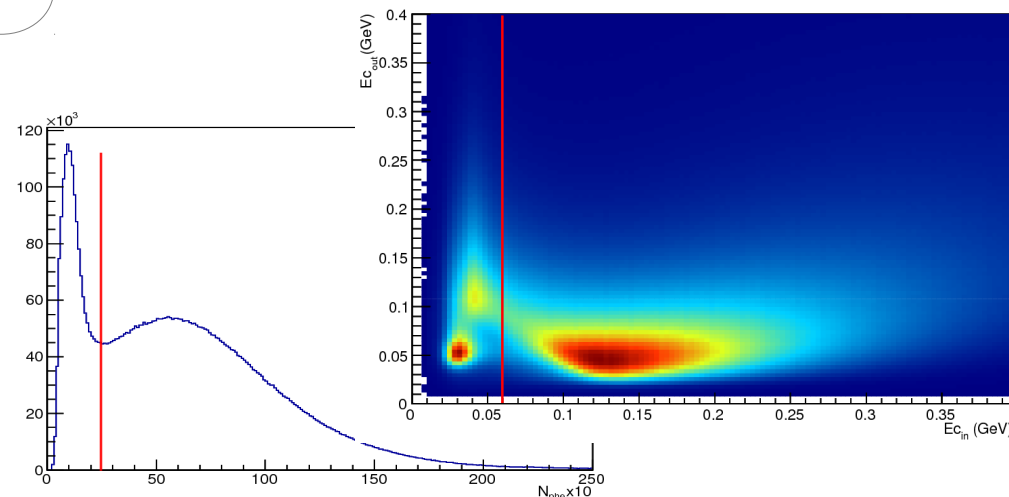
$$ep \rightarrow ep \gamma$$

➔ Electron :

Identification with

DC + CC + EC

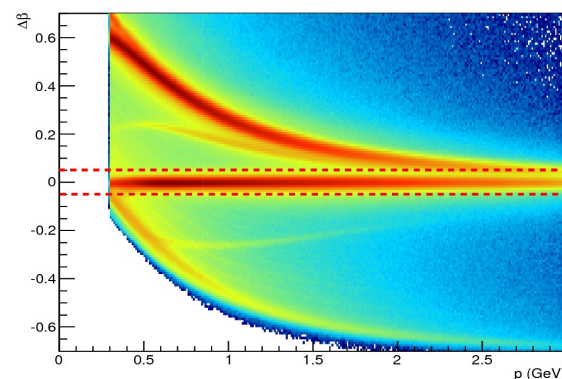
Rejection of the π^- with EC and CC



➔ Proton:

Identified with DC + SC

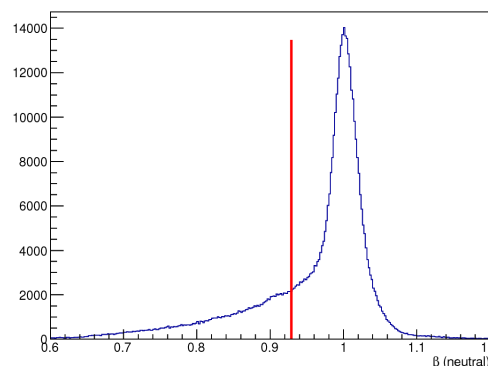
$$\text{Cut on: } \Delta\beta = \beta_{\text{measured}}^{\text{SC}} - \beta_{\text{calculated}}^{\text{DC}}(M_p) = \frac{d}{ct} - \frac{p}{\sqrt{p^2 + M_p^2}}$$



➔ Photon :

Identified with IC and/or EC

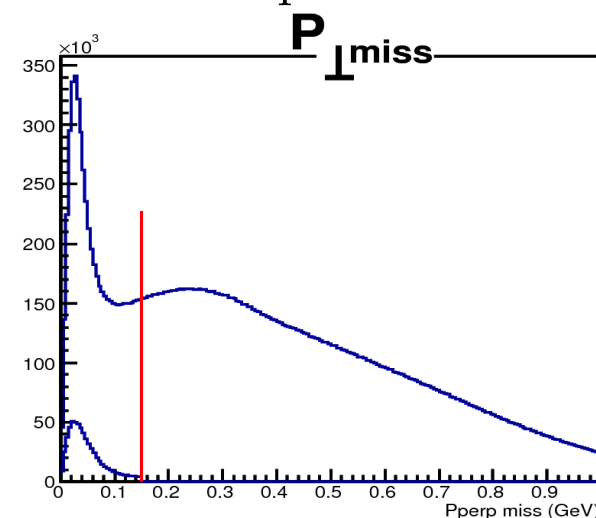
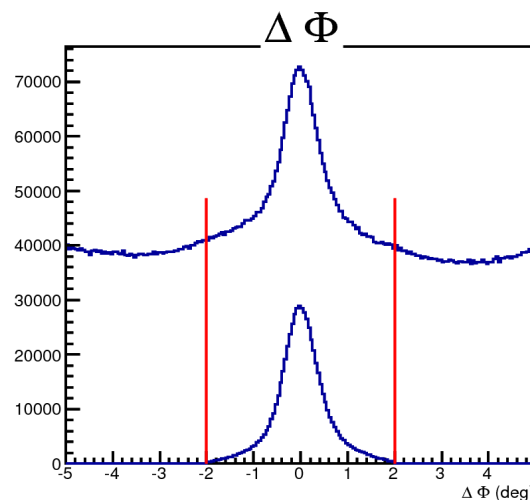
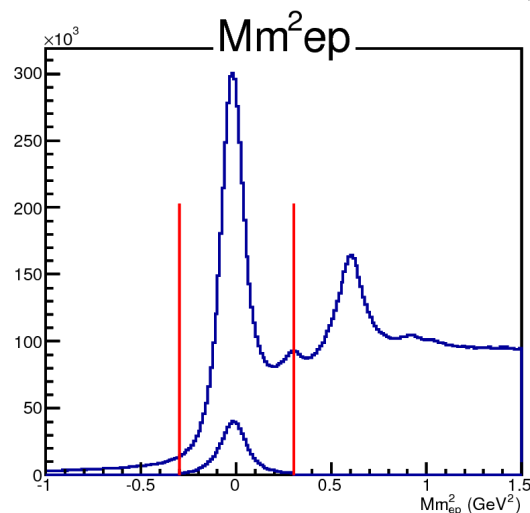
Rejection of the neutrons in EC
with $\beta_{\text{neutral}} > 0.92$



Exclusivity cuts and background

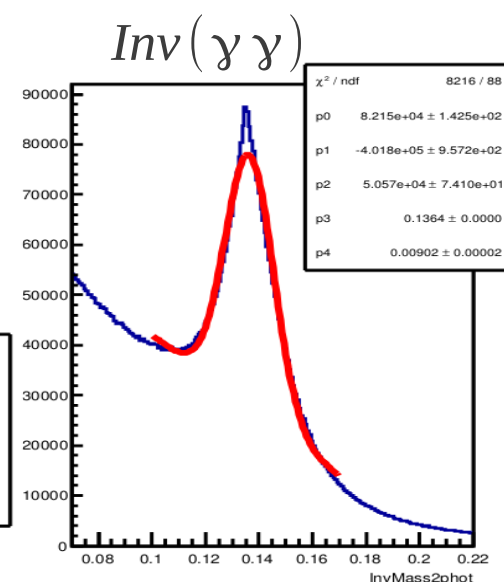
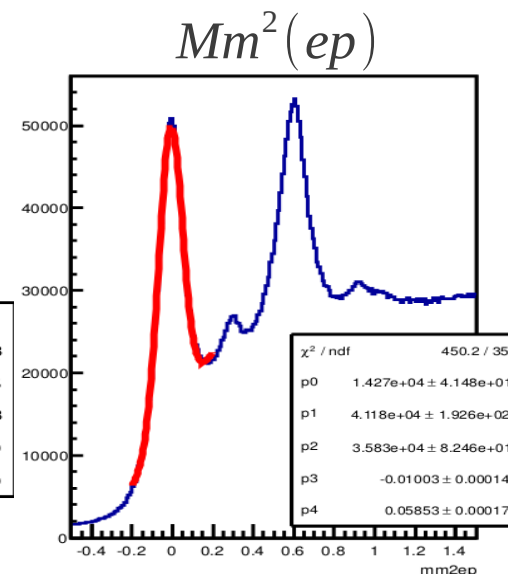
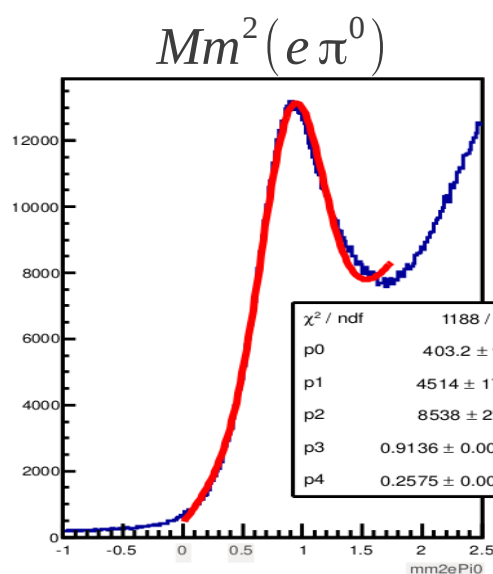
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- Additional cuts are applied to ensure the exclusivity of the DVCS process:



- Background :
 $ep \rightarrow ep \pi^0 \rightarrow ep \gamma(\gamma)$

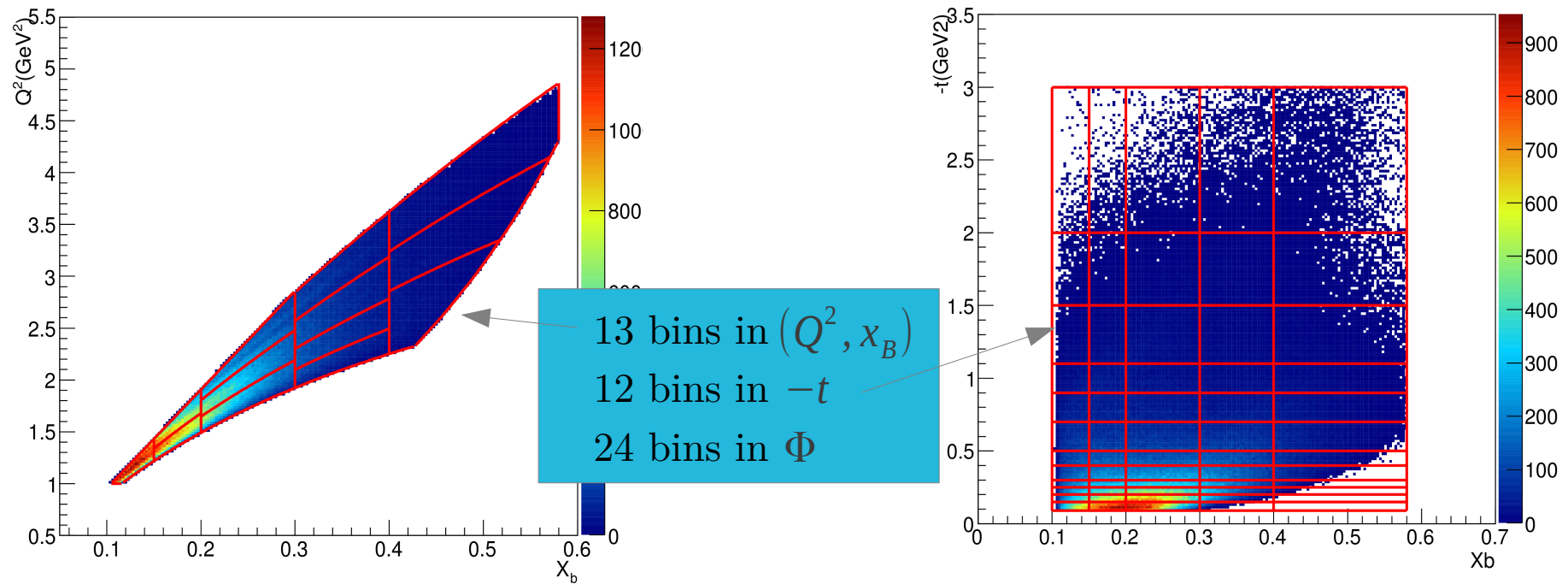
Between 10 to 20 %



Kinematic coverage of the e1-DVCS data and binning

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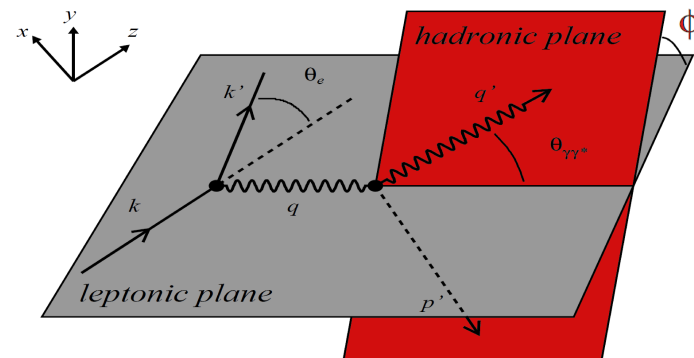
$$Q^2 > 1, 0.1 < x_B < 0.58, 21 < \theta_e < 45, p_e > 0.8, W > 2$$



The kinematics of the DVCS reaction is defined by 4 independent variables :

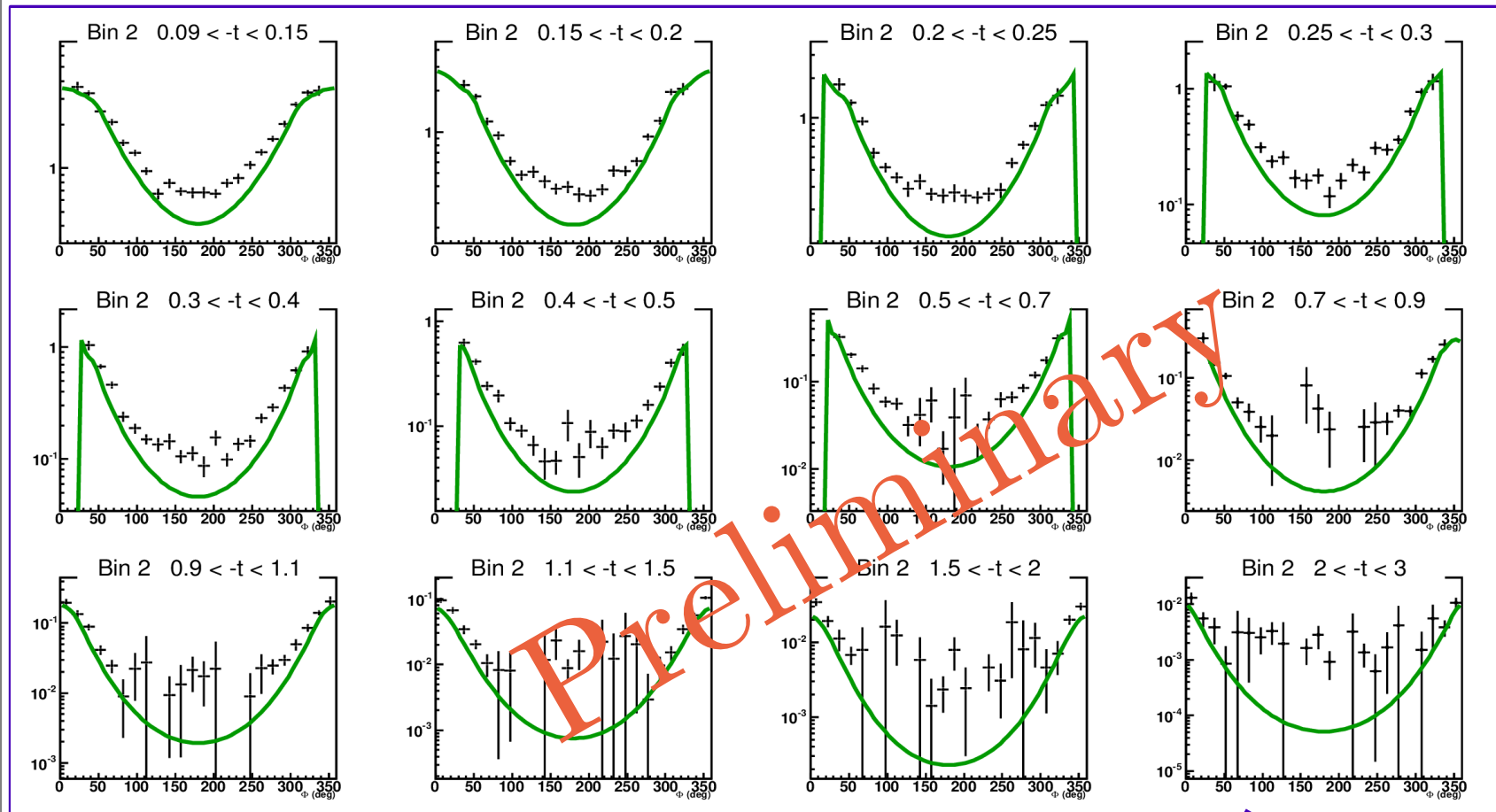
$$Q^2, x_B, -t, \Phi$$

→ 4-dimensional binning



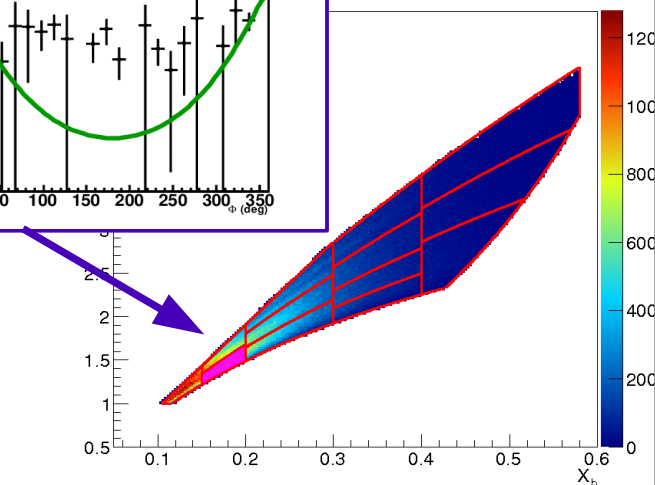
Unpolarized cross sections

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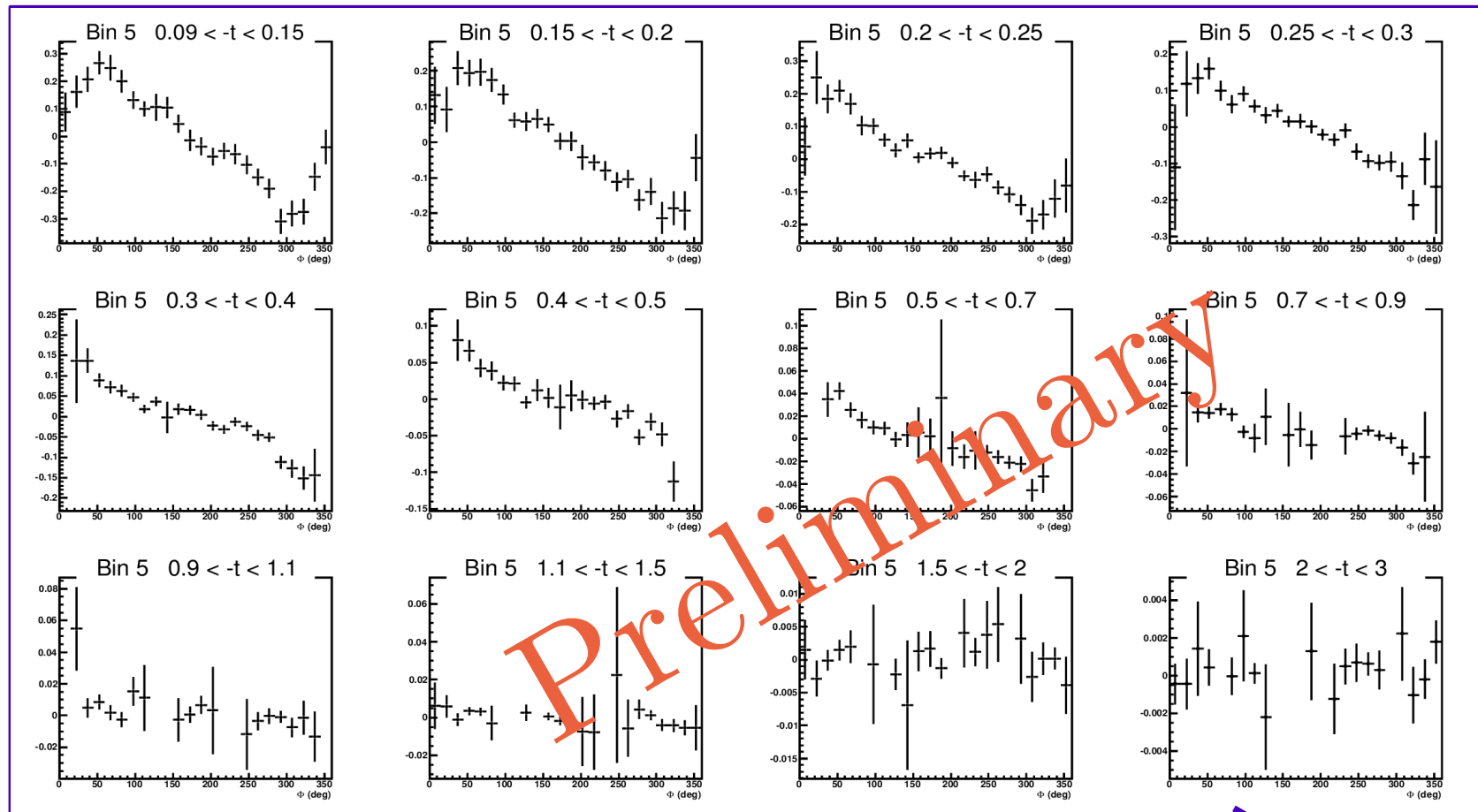
● $\frac{d^4 \sigma_{ep \rightarrow ep \gamma}}{dQ^2 dx_B dt d\Phi} \text{ (nb/GeV}^4\text{)}$

— Bethe-Heitler integrated over the bin

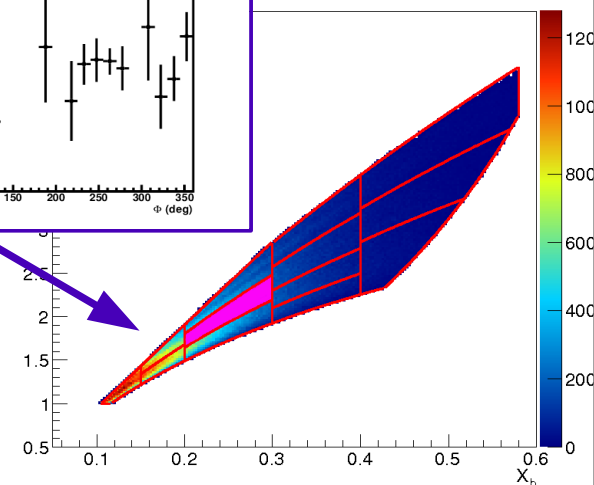


Difference of polarized cross sections

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$$\bullet \frac{1}{2} \left(\frac{d^4 \sigma_{ep \rightarrow ep \gamma}^{\rightarrow}}{dQ^2 dx_B dt d\Phi} - \frac{d^4 \sigma_{ep \rightarrow ep \gamma}^{\leftarrow}}{dQ^2 dx_B dt d\Phi} \right) (nb/GeV^4)$$



Extraction of Compton Form Factors (CFFs)

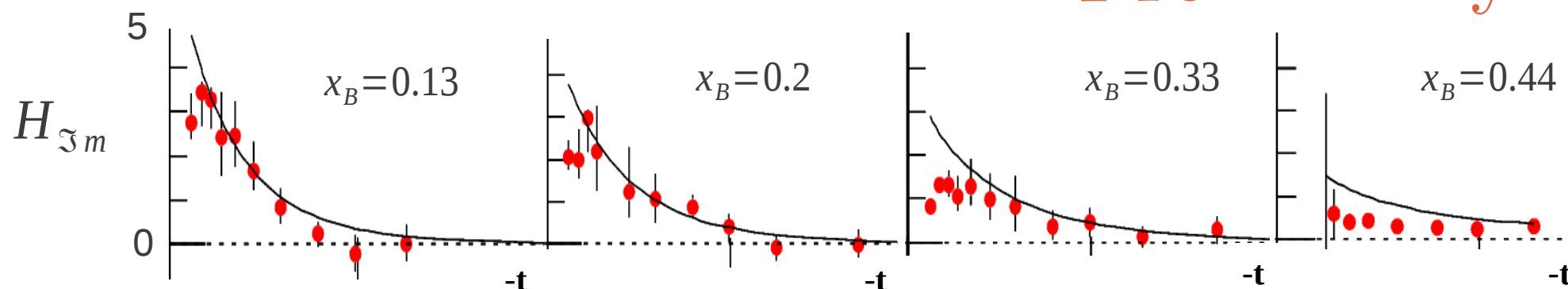
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Model-independent fitting procedure:

At fixed (Q^2, x_B, t) , extraction of the CFFs from the DVCS observables

CFFs are combinations of GPDs

Preliminary



● Results of model-independent fit — VGG model prediction

$H_{\Im m}$: the t-slope reflects the size of the probed object:

Large t-slope \iff extended object

Small t-slope \iff localized object

➡ The sea quarks (low x) spread to the periphery of the nucleon

➡ The valence quarks (large x) remain in the center

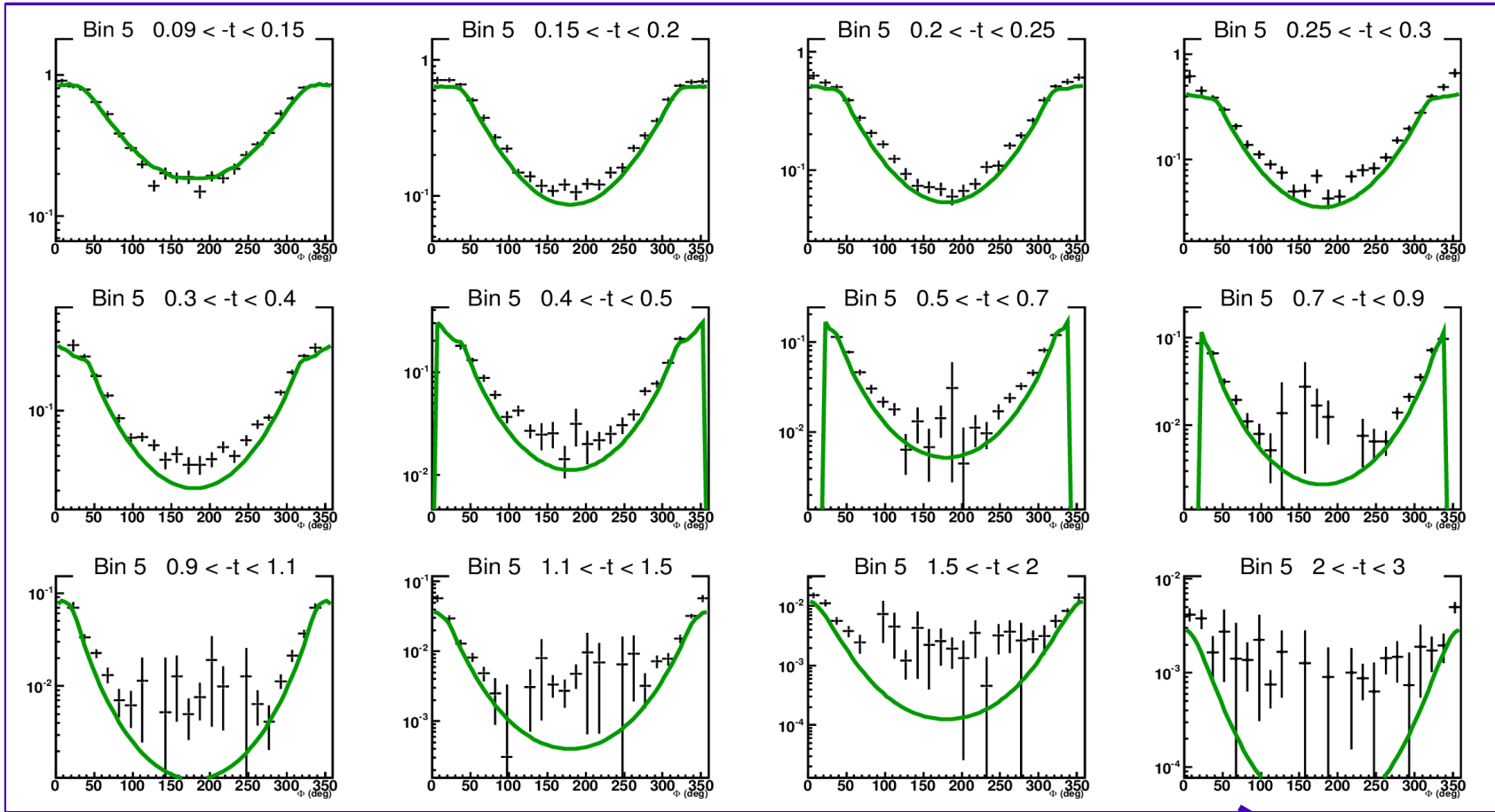
Conclusion

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- Extraction of DVCS unpolarized and polarized cross sections in the largest kinematic domain ever explored in the valence region
- Extraction of Compton Form Factors by fitting simultaneously the unpolarized and polarized is in progress
- Dedicated GPD program at Jlab 12GeV :
 - ➡ Target spin asymmetry A_{UT} , A_{LT}
 - ➡ Target spin asymmetry A_{LU}
 - ➡ DVMP: pseudoscalar/vector mesons

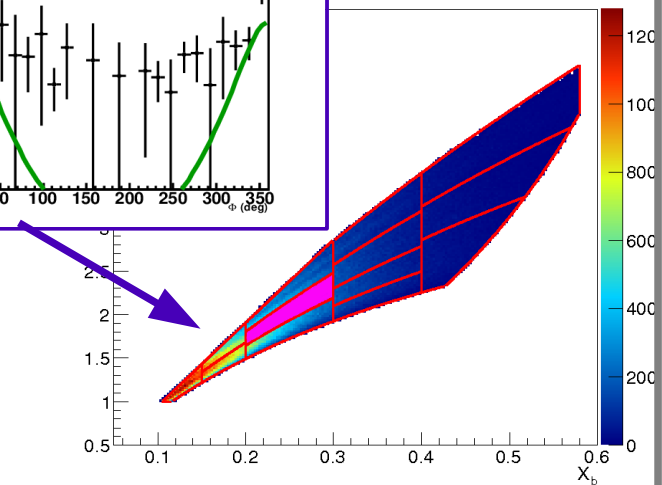
Unpolarized cross sections

15



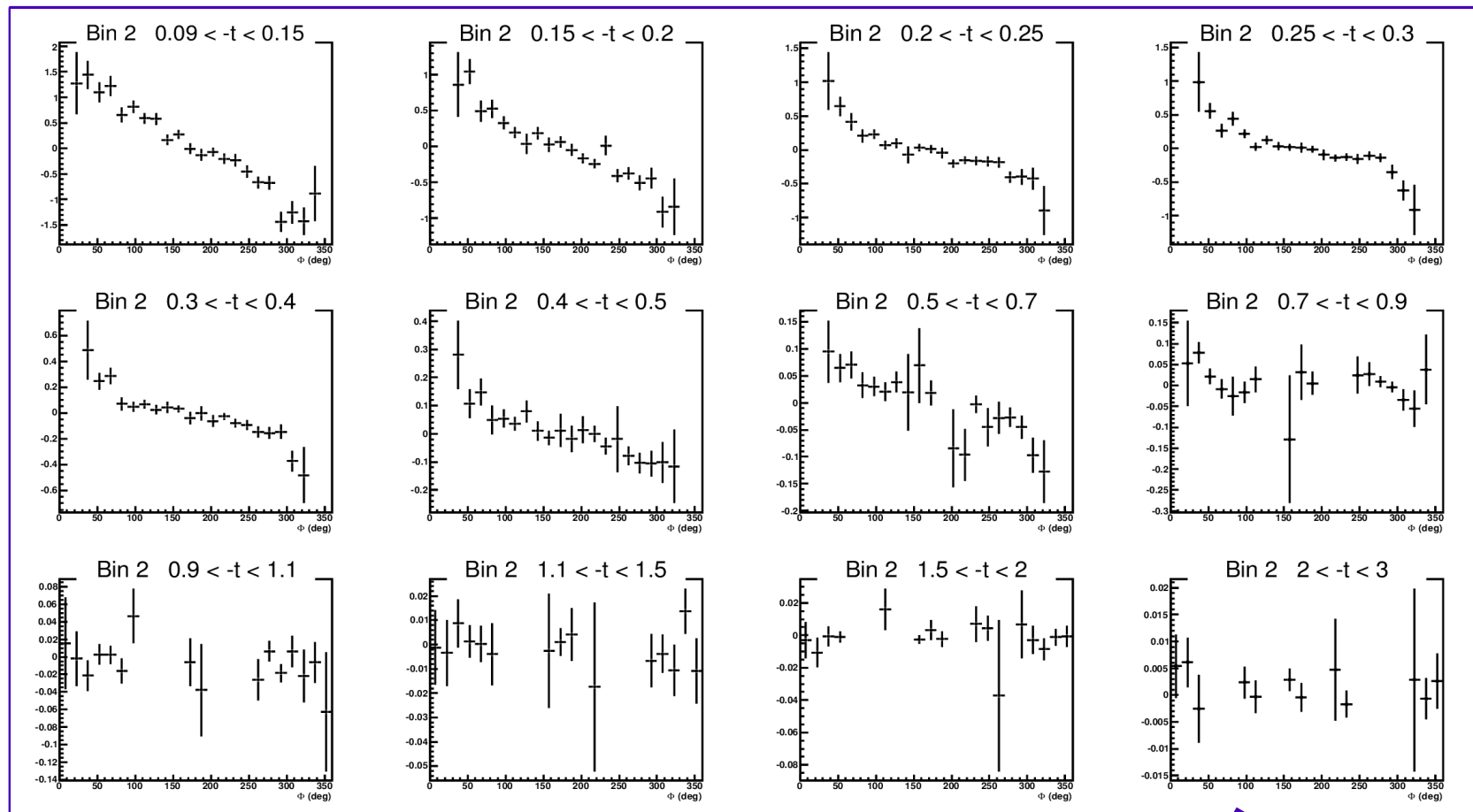
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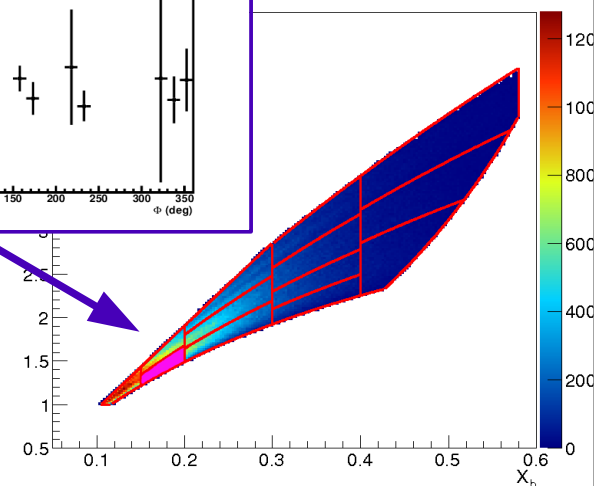


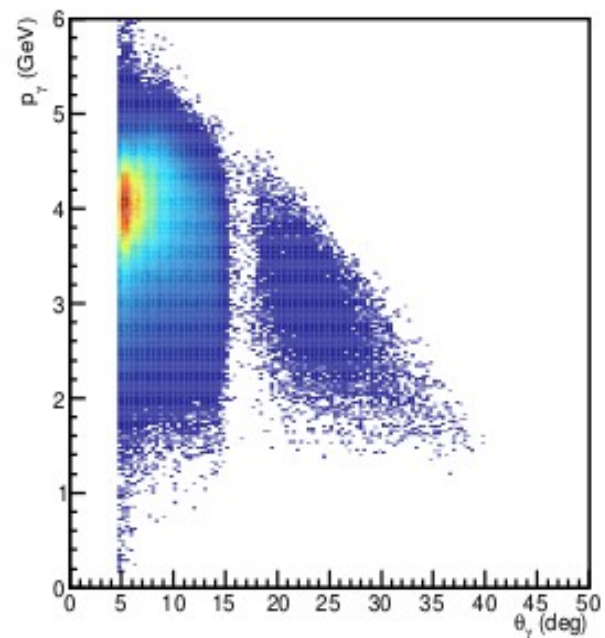
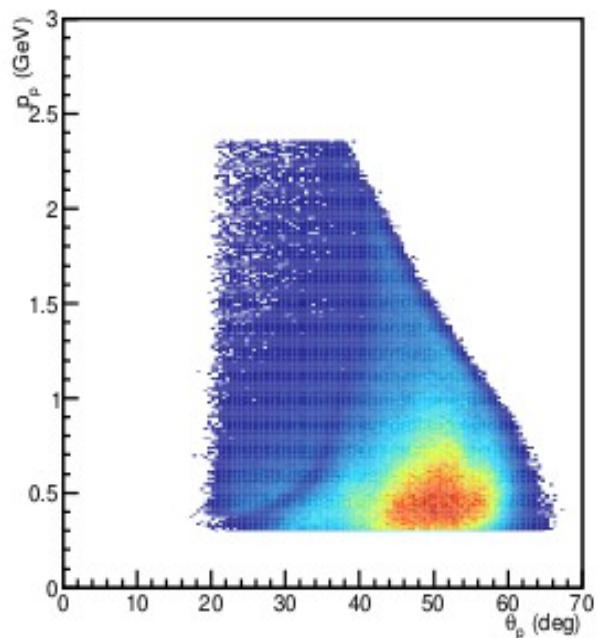
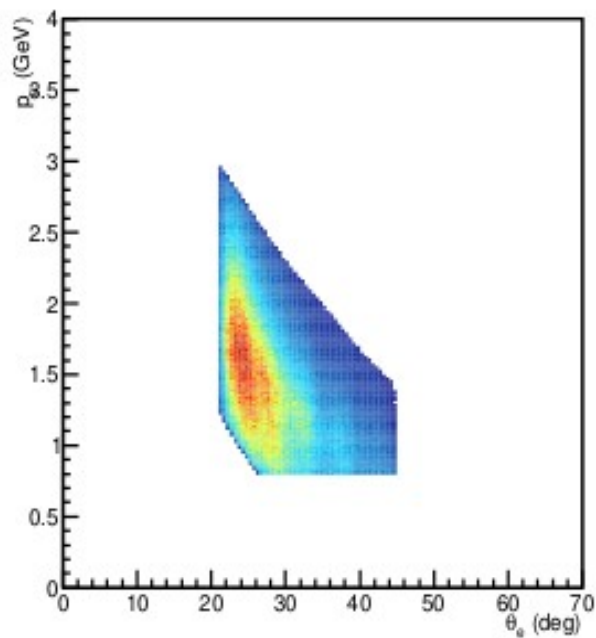
Difference of polarized cross sections

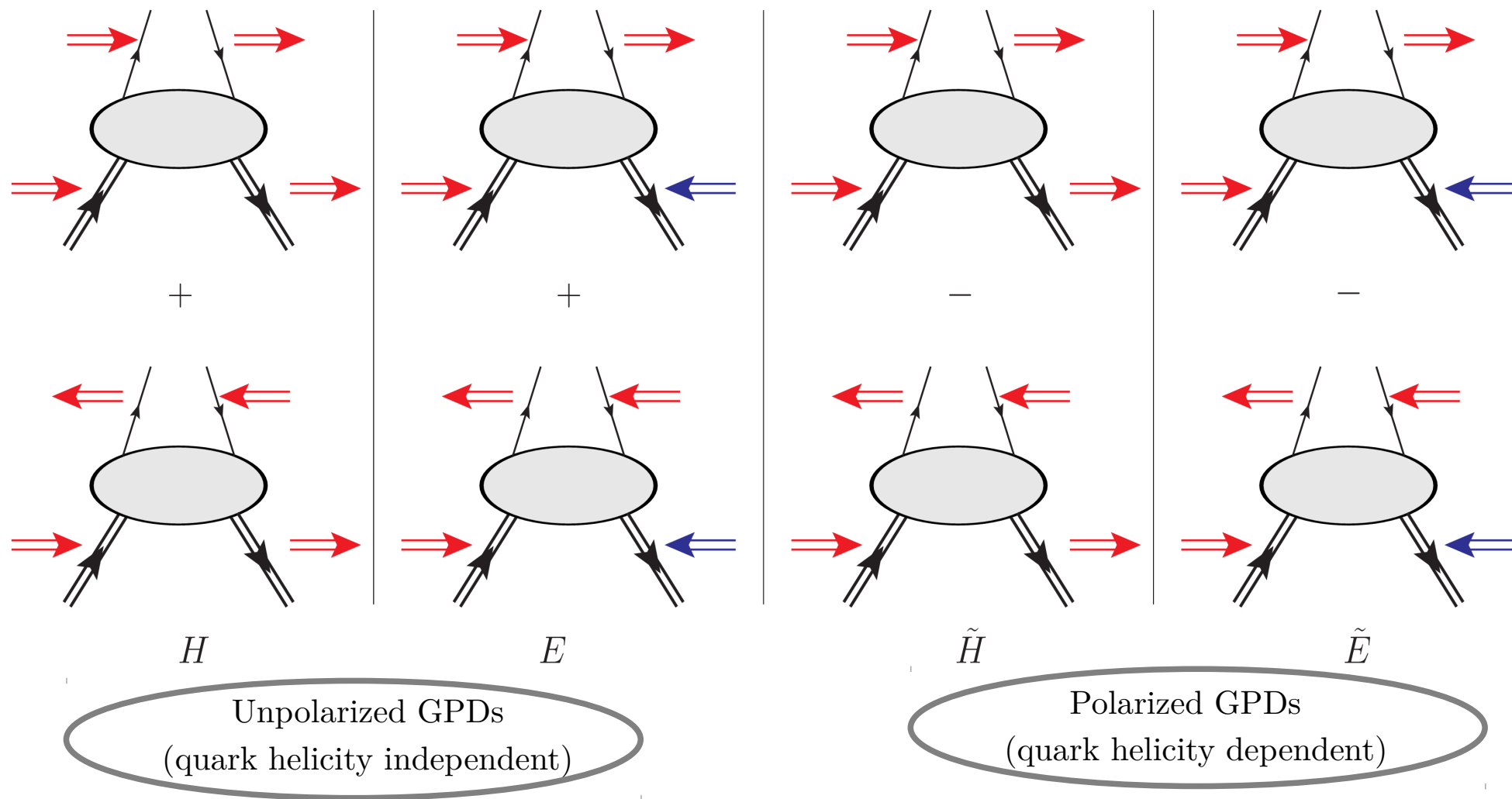
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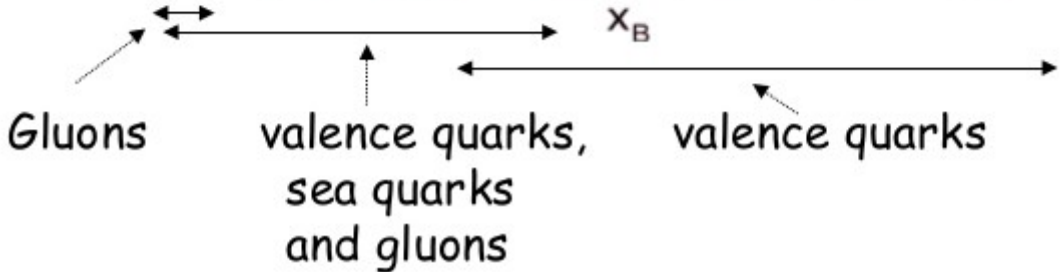
$$\bullet \frac{1}{2} \left(\frac{d^4 \sigma_{ep \rightarrow ep \gamma}^{\rightarrow}}{dQ^2 dx_B dt d\Phi} - \frac{d^4 \sigma_{ep \rightarrow ep \gamma}^{\leftarrow}}{dQ^2 dx_B dt d\Phi} \right) (nb/GeV^4)$$







Four different possible spin-helicity transition of the nucleon-quark system



Extraction of Compton Form Factors (CFFs)

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$$\mathcal{T}^{DVCS} \propto \int_{-1}^1 dx \frac{f(x, \xi, t)}{x \pm \xi \mp i\epsilon} = \mathcal{P} \int_{-1}^1 dx \frac{f(x, \xi, t)}{x \pm \xi} \pm i\pi f(x = \mp \xi, \xi, t)$$

Real part	Imaginary part
$\mathcal{H}_{Re}(\xi, t) = \mathcal{P} \int_0^1 dx [H(x, \xi, t) - H(-x, \xi, t)] C^+(x, \xi)$	$\mathcal{H}_{Im}(\xi, t) = H(\xi, \xi, t) - H(-\xi, \xi, t)$
$\mathcal{E}_{Re}(\xi, t) = \mathcal{P} \int_0^1 dx [E(x, \xi, t) - E(-x, \xi, t)] C^+(x, \xi)$	$\mathcal{E}_{Im}(\xi, t) = E(\xi, \xi, t) - E(-\xi, \xi, t)$
$\tilde{\mathcal{H}}_{Re}(\xi, t) = \mathcal{P} \int_0^1 dx [\tilde{H}(x, \xi, t) + \tilde{H}(-x, \xi, t)] C^-(x, \xi)$	$\tilde{\mathcal{H}}_{Im}(\xi, t) = \tilde{H}(\xi, \xi, t) + \tilde{H}(-\xi, \xi, t)$
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$$\text{With: } C^\pm(x, \xi) = \frac{1}{x - \xi} \pm \frac{1}{x + \xi}$$

Extraction of Compton Form Factors (CFFs)

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	Sensitivity	Experiment
$\sigma_{unp} = \sigma^{\rightarrow} + \sigma^{\leftarrow}$	$\propto \mathcal{H}_{Re}$	H1 ((2001), (2005), (2008)) ZEUS ((2003) , (2009)) Hall-A (2006) Hall-B (E1-DVCS experiment: data under analysis)
$\sigma_{pol} = \sigma^{\rightarrow} - \sigma^{\leftarrow}$	$\propto \mathcal{H}_{Im}$	Hall-A (2006) Hall-B (E1-DVCS experiment: data under analysis)
\mathcal{A}_C	$\propto \mathcal{H}_{Re}$	HERMES ((2007) , (2008) , (2009))
\mathcal{A}_{LU}	$\propto \mathcal{H}_{Im}$	HERMES ((2001) , (2009)) Hall-B ((2001) , (2008))
\mathcal{A}_{UL}	$\propto \mathcal{H}_{Im}, \tilde{\mathcal{H}}_{Im}$	Hall-B (2006) HERMES (2010) Hall-B (Eg1-DVCS experiment: data under analysis)
\mathcal{A}_{LL}	$\propto \mathcal{H}_{Re}, \tilde{\mathcal{H}}_{Re}$	HERMES (2010) Hall-B (Eg1-DVCS experiment: data under analysis)
\mathcal{A}_{UT}	$\propto \mathcal{E}_{Im}$	HERMES (2008) Hall-B proposal
\mathcal{A}_{LT}	$\propto \mathcal{H}_{Re}, \mathcal{E}_{Re}$	HERMES (2011) Hall-B proposal

Model-independent fit, at fixed (Q^2, x_B, t) of DVCS observables
with MINUIT + MINOS

8 unknowns (the CFFs), non-linear problem, strong correlations

Bounding the domain of variation of the CFFs (5xVGG)

Extraction of Compton Form Factors (CFFs)

22

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Bounding the domain of variation of the CFFs (5xVGG)