

Color superconductivity beyond Mean Field Approximation

Ugo Mire

In collaboration with: Bernd-Jochen Schaefer

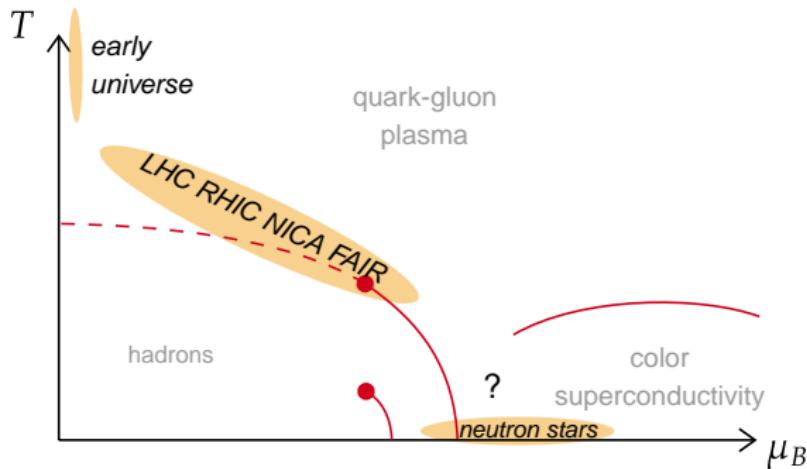
Hirschgägg, September 12, 2023

Outline

1. Color Superconductivity and 2SC Phase
2. Functional Renormalization Group
3. Phase Structure and Astrophysical Applications

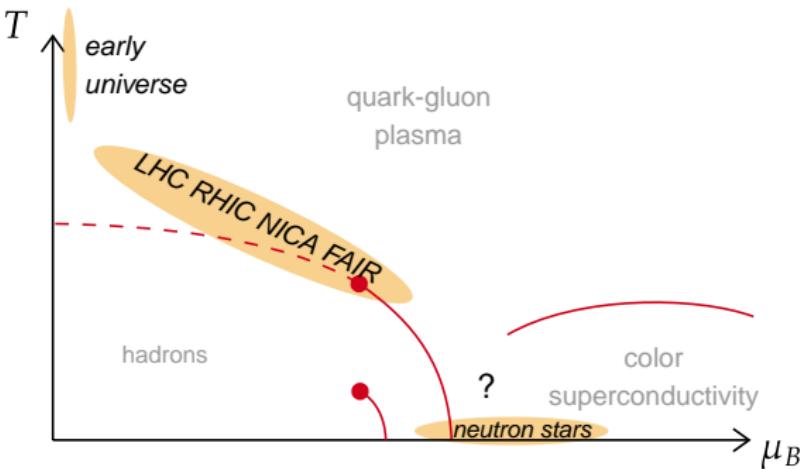
Color Superconductivity and 2SC Phase

Why study color superconductivity?



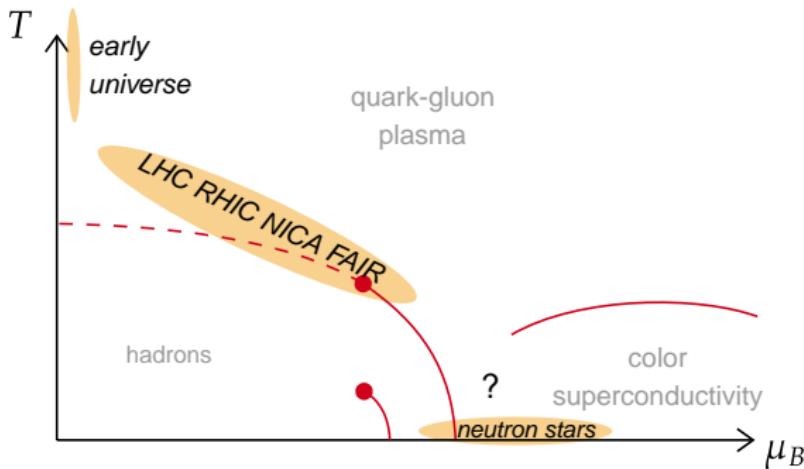
- Relevant for **astrophysical observations**.

Why study color superconductivity?



- Relevant for **astrophysical observations**.
- **Natural in QCD**: attractive gluon exchange \rightarrow Cooper pairing and superconductivity at high densities.

Why study color superconductivity?



- Relevant for **astrophysical observations**.
- **Natural in QCD**: attractive gluon exchange \rightarrow **Cooper pairing** and superconductivity at high densities.

↑
which pairing channel?

2SC Phase

2SC phase: condensates of the shape

$$\Delta_A = \langle q^T C \gamma_5 \tau_2 \lambda_A q \rangle, \quad A = 2, 5, 7$$

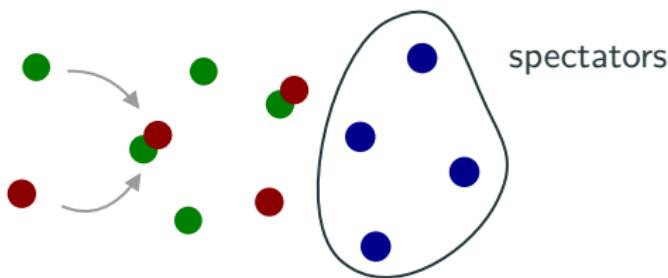
Main characteristics:

some matrix structure in flavor,
color and Dirac space

- Chiral symmetry is not broken.
- Color symmetry is "broken" (Higgs mechanism), but not fully:

$$SU_c(3) \rightarrow SU_c(2) .$$

- Only two color participate in superconductivity.



Functional Renormalization Group

Why use the FRG?

Four-quark interactions in QCD

- Low densities: scalar-pseudoscalar ($\sigma, \vec{\pi}$) channel dominates

$$(\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5\tau_i\psi)^2$$

→ quantitative description of chiral crossover using FRG.

[Fu, Pawłowski & Rennecke; 1909.02991]

Why use the FRG?

Four-quark interactions in QCD

- Low densities: **scalar-pseudoscalar** (σ , $\vec{\pi}$) channel dominates

$$(\bar{\psi}\psi)^2 - (\bar{\psi}\gamma_5\tau_i\psi)^2$$

→ quantitative description of chiral crossover using FRG.

- High densities: **diquark** channel dominates

[Fu, Pawlowski & Rennecke; 1909.02991]

[Braun, Leonhardt & Pospiech; 1909.06298]

$$(q^T C \gamma_5 \tau_2 \lambda_A q) (\bar{q} \gamma_5 \tau_2 \lambda_A C \bar{q}^T)$$

→ a natural step is to include the diquark channel.

Quark-meson-diquark Model

Quark-meson-diquark model action:

$$S_{\text{QMD}} = \int_x \left\{ \bar{q} \left(\not{\partial} - \mu \gamma_0 + g_\phi (\sigma + i \gamma_5 \vec{\pi} \cdot \vec{\tau}) \right) q \right. \\ \left. + \frac{g_\Delta}{2} \left(\Delta_A q^T C \gamma_5 \tau_2 \lambda_A q - \Delta_A^* \bar{q} \gamma_5 \tau_2 \lambda_A C \bar{q}^T \right) \right. \\ \left. + \frac{1}{2} (\partial_\mu \sigma)^2 + \frac{1}{2} (\partial_\mu \vec{\pi})^2 \right. \\ \left. + ((\partial_\nu + \delta_{\nu 0} 2\mu) \Delta_A^*) (\partial_\nu - \delta_{\nu 0} 2\mu) \Delta_A \right. \\ \left. + U(\sigma^2 + \vec{\pi}^2, \Delta_A \Delta_A^*) - c \sigma \right\}$$

Quark-meson-diquark Model

Quark-meson-diquark model action:

quark-meson
scalar-pseudoscalar
channel

$$S_{\text{QMD}} = \int_x \left\{ \bar{q} \left(\not{\partial} - \mu \gamma_0 + g_\phi (\sigma + i \gamma_5 \vec{\pi} \cdot \vec{\tau}) \right) q \right.$$

\swarrow

$$\begin{aligned} &+ \frac{g_\Delta}{2} \left(\Delta_A q^T C \gamma_5 \tau_2 \lambda_A q - \Delta_A^* \bar{q} \gamma_5 \tau_2 \lambda_A C \bar{q}^T \right) \\ &+ \frac{1}{2} (\partial_\mu \sigma)^2 + \frac{1}{2} (\partial_\mu \vec{\pi})^2 \\ &+ ((\partial_\nu + \delta_{\nu 0} 2\mu) \Delta_A^*) (\partial_\nu - \delta_{\nu 0} 2\mu) \Delta_A \\ &\left. + U(\sigma^2 + \vec{\pi}^2, \Delta_A \Delta_A^*) - c\sigma \right\} \end{aligned}$$

Quark-meson-diquark Model

Quark-meson-diquark model action:

quark-meson
scalar-pseudoscalar
channel

$$S_{\text{QMD}} = \int_x \left\{ \bar{q} \left(\not{\partial} - \mu \gamma_0 + g_\phi (\sigma + i \gamma_5 \vec{\pi} \cdot \vec{\tau}) \right) q \right.$$

↖
quark-diquark
interaction

$$\begin{aligned} &+ \frac{g_\Delta}{2} \left(\Delta_A q^T C \gamma_5 \tau_2 \lambda_A q - \Delta_A^* \bar{q} \gamma_5 \tau_2 \lambda_A C \bar{q}^T \right) \\ &+ \frac{1}{2} (\partial_\mu \sigma)^2 + \frac{1}{2} (\partial_\mu \vec{\pi})^2 \\ &+ ((\partial_\nu + \delta_{\nu 0} 2\mu) \Delta_A^*) (\partial_\nu - \delta_{\nu 0} 2\mu) \Delta_A \\ &\left. + U(\sigma^2 + \vec{\pi}^2, \Delta_A \Delta_A^*) - c\sigma \right\} \end{aligned}$$

Quark-meson-diquark Model

Quark-meson-diquark model action:

$$S_{\text{QMD}} = \int_x \left\{ \bar{q} \left(\not{\partial} - \mu \gamma_0 + g_\phi (\sigma + i \gamma_5 \vec{\pi} \cdot \vec{\tau}) \right) q \right.$$

↖
quark-meson
scalar-pseudoscalar
channel

$$+ \frac{g_\Delta}{2} \left(\Delta_A q^T C \gamma_5 \tau_2 \lambda_A q - \Delta_A^* \bar{q} \gamma_5 \tau_2 \lambda_A C \bar{q}^T \right)$$

↖ quark-diquark
interaction

$$+ \frac{1}{2} (\partial_\mu \sigma)^2 + \frac{1}{2} (\partial_\mu \vec{\pi})^2$$
$$+ ((\partial_\nu + \delta_{\nu 0} 2\mu) \Delta_A^*) (\partial_\nu - \delta_{\nu 0} 2\mu) \Delta_A$$

↖ diquark kinetic
term, couples to μ

$$\left. + U(\sigma^2 + \vec{\pi}^2, \Delta_A \Delta_A^*) - c\sigma \right\}$$

Quark-meson-diquark Model

Quark-meson-diquark model action:

$$S_{\text{QMD}} = \int_x \left\{ \bar{q} \left(\not{\partial} - \mu \gamma_0 + g_\phi (\sigma + i \gamma_5 \vec{\pi} \cdot \vec{\tau}) \right) q \right.$$

quark-meson
scalar-pseudoscalar
channel

$$\begin{aligned} &+ \frac{g_\Delta}{2} \left(\Delta_A q^T C \gamma_5 \tau_2 \lambda_A q - \Delta_A^* \bar{q} \gamma_5 \tau_2 \lambda_A C \bar{q}^T \right) \\ &+ \frac{1}{2} (\partial_\mu \sigma)^2 + \frac{1}{2} (\partial_\mu \vec{\pi})^2 \\ &+ ((\partial_\nu + \delta_{\nu 0} 2\mu) \Delta_A^*) (\partial_\nu - \delta_{\nu 0} 2\mu) \Delta_A \\ &\left. + U(\sigma^2 + \vec{\pi}^2, \Delta_A \Delta_A^*) - c\sigma \right\} \end{aligned}$$

quark-diquark
interaction

diquark kinetic
term, couples to μ

meson and diquark interactions,
arbitrary potential constrained by
symmetries

Quark-meson-diquark Model

Quark-meson-diquark model action:

$$S_{\text{QMD}} = \int_x \left\{ \bar{q} \left(\not{\partial} - \mu \gamma_0 + g_\phi (\sigma + i \gamma_5 \vec{\pi} \cdot \vec{\tau}) \right) q \right.$$

↖
quark-meson
scalar-pseudoscalar
channel

$$+ \frac{g_\Delta}{2} \left(\Delta_A q^T C \gamma_5 \tau_2 \lambda_A q - \Delta_A^* \bar{q} \gamma_5 \tau_2 \lambda_A C \bar{q}^T \right)$$

↖ quark-diquark
interaction

$$+ \frac{1}{2} (\partial_\mu \sigma)^2 + \frac{1}{2} (\partial_\mu \vec{\pi})^2$$
$$+ ((\partial_\nu + \delta_{\nu 0} 2\mu) \Delta_A^*) (\partial_\nu - \delta_{\nu 0} 2\mu) \Delta_A$$
$$+ U(\sigma^2 + \vec{\pi}^2, \Delta_A \Delta_A^*) - c\sigma \left. \vphantom{\int_x} \right\}$$

↖ diquark kinetic
term, couples to μ

meson and diquark interactions,
arbitrary potential constrained by
symmetries

explicit chiral symmetry
breaking term

Truncation and Regulator Choice

Some FRG details:

- Local potential approximation (LPA): only scale dependent quantity is the potential $U_k(\sigma, \Delta)$.
 - σ : chiral condensate.
 - Δ : diquark condensate.
 - $U_k \rightarrow \infty$ number of couplings.

Truncation and Regulator Choice

Some FRG details:

- Local potential approximation (LPA): only scale dependent quantity is the potential $U_k(\sigma, \Delta)$.
 - σ : chiral condensate.
 - Δ : diquark condensate.
 - $U_k \rightarrow \infty$ number of couplings.
- 3d Litim regulator

$$R_k = (k^2 - \vec{p}^2) \theta(k^2 - \vec{p}^2) .$$

Expect back bending from studies of the quark-meson model → more on that later.

Truncation and Regulator Choice

Some FRG details:

- Local potential approximation (LPA): only scale dependent quantity is the potential $U_k(\sigma, \Delta)$.
 - σ : chiral condensate.
 - Δ : diquark condensate.
 - $U_k \rightarrow \infty$ number of couplings.
- 3d Litim regulator

$$R_k = (k^2 - \vec{p}^2) \theta(k^2 - \vec{p}^2) .$$

Expect back bending from studies of the quark-meson model → more on that later.

- Initial conditions:

$$U_\Lambda = a_1 \sigma^2 + a_2 \sigma^4 + b_1 \Delta^2 + b_2 \Delta^4$$

Flow Equation

$$\partial_t U_k(\sigma, \Delta) = - \left(q_r, q_g \right) - \left(q_b \right) + \frac{1}{2} \left(\Delta_2, \sigma \right) + \frac{1}{2} \left(\pi \right) + \frac{1}{2} \left(\Delta_5, \Delta_7 \right)$$

The equation shows the time derivative of the flow function U_k as a sum of four terms. The first term is a solid circle with two arrows pointing clockwise, labeled q_r, q_g . The second term is a solid circle with one arrow pointing clockwise, labeled q_b . The third term is a dashed circle with one arrow pointing clockwise, labeled Δ_2, σ . The fourth term is a dashed circle with no arrow, labeled π . The fifth term is another dashed circle with no arrow, labeled Δ_5, Δ_7 . Each term has a coefficient of $\frac{1}{2}$.

Flow Equation

depends on $E_\Delta = \sqrt{(\epsilon_k \pm \mu)^2 + g_\Delta^2 \Delta^2}$

with $\epsilon_k = \sqrt{k^2 + g_\phi^2 \sigma^2}$



$$\partial_t U_k(\sigma, \Delta) = - \left(\text{Diagram } q_r, q_g \right) - \left(\text{Diagram } q_b \right) + \frac{1}{2} \left(\text{Diagram } \Delta_2, \sigma \right) + \frac{1}{2} \left(\text{Diagram } \pi \right) + \frac{1}{2} \left(\text{Diagram } \Delta_5, \Delta_7 \right)$$

The equation shows the time derivative of the flow function U_k as a sum of five terms, each represented by a Feynman diagram. The first term is a solid circle with a clockwise arrow and two external lines labeled q_r (red) and q_g (green). The second term is a solid circle with a clockwise arrow and one external line labeled q_b (blue). The third term is a dashed circle with a clockwise arrow and one external line labeled Δ_2, σ . The fourth term is a dashed circle with a clockwise arrow and one external line labeled π . The fifth term is a dashed circle with a clockwise arrow and two external lines labeled Δ_5, Δ_7 . Each diagram has a tensor symbol (\otimes) at its top vertex.

Flow Equation

depends on $E_\Delta = \sqrt{(\epsilon_k \pm \mu)^2 + g_\Delta^2 \Delta^2}$

with $\epsilon_k = \sqrt{k^2 + g_\phi^2 \sigma^2}$

$$\partial_t U_k(\sigma, \Delta) = - \left(\text{Diagram } q_r, q_g \right) - \left(\text{Diagram } q_b \right) + \frac{1}{2} \left(\text{Diagram } \Delta_2, \sigma \right) + \frac{1}{2} \left(\text{Diagram } \pi \right) + \frac{1}{2} \left(\text{Diagram } \Delta_5, \Delta_7 \right)$$

↓

coupling between one diquark
and the sigma meson

Flow Equation

depends on $E_\Delta = \sqrt{(\epsilon_k \pm \mu)^2 + g_\Delta^2 \Delta^2}$

with $\epsilon_k = \sqrt{k^2 + g_\phi^2 \sigma^2}$

$$\partial_t U_k(\sigma, \Delta) = - \left(\text{Diagram } q_r, q_g \right) - \left(\text{Diagram } q_b \right) + \frac{1}{2} \left(\text{Diagram } \Delta_2, \sigma \right) + \frac{1}{2} \left(\text{Diagram } \pi \right) + \frac{1}{2} \left(\text{Diagram } \Delta_5, \Delta_7 \right)$$

↗

$$\sim \coth \frac{\epsilon_\pi}{2T} \text{ with } \epsilon_\pi = \sqrt{k^2 + 2\partial_{\sigma^2} U_k}$$

coupling between one diquark
and the sigma meson

Flow Equation

depends on $E_\Delta = \sqrt{(\epsilon_k \pm \mu)^2 + g_\Delta^2 \Delta^2}$

with $\epsilon_k = \sqrt{k^2 + g_\phi^2 \sigma^2}$

coupling between one diquark
and the sigma meson



$$\partial_t U_k(\sigma, \Delta) = - \left(q_r, q_g \right) - \left(q_b \right) + \frac{1}{2} \left(\Delta_2, \sigma \right) + \frac{1}{2} \left(\pi \right) + \frac{1}{2} \left(\Delta_5, \Delta_7 \right)$$



$$\sim \coth \frac{\epsilon_\pi}{2T} \text{ with } \epsilon_\pi = \sqrt{k^2 + 2\partial_{\sigma^2} U_k}$$



$$\sim \coth \frac{\epsilon_\Delta - 2\mu}{2T} \text{ with } \epsilon_\Delta = \sqrt{k^2 + 2\partial_{\Delta^2} U_k}$$

Flow Equation

depends on $E_\Delta = \sqrt{(\epsilon_k \pm \mu)^2 + g_\Delta^2 \Delta^2}$

with $\epsilon_k = \sqrt{k^2 + g_\phi^2 \sigma^2}$

$$\partial_t U_k(\sigma, \Delta) = - \left(q_r, q_g \right) - \left(q_b \right) + \frac{1}{2} \left(\Delta_2, \sigma \right) + \frac{1}{2} \left(\pi \right) + \frac{1}{2} \left(\Delta_5, \Delta_7 \right)$$

$$\sim \coth \frac{\epsilon_\pi}{2T} \text{ with } \epsilon_\pi = \sqrt{k^2 + 2\partial_{\sigma^2} U_k}$$

coupling between one diquark
and the sigma meson

$$\sim \coth \frac{\epsilon_\Delta - 2\mu}{2T} \text{ with } \epsilon_\Delta = \sqrt{k^2 + 2\partial_{\Delta^2} U_k}$$

problem with this part

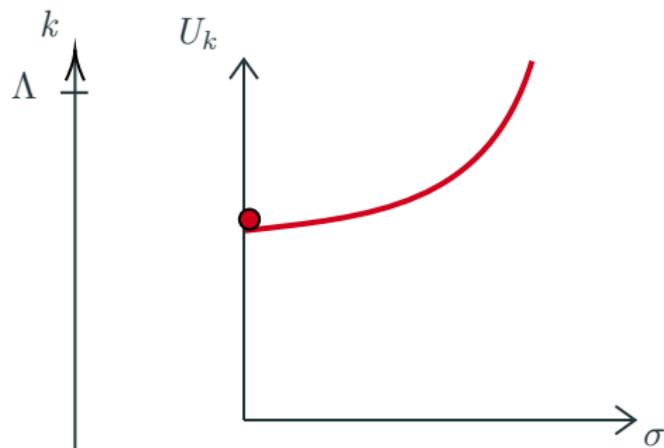
First Trouble: Bosonic Diquark Loops

Pion loops

$$\partial_t U_k = \frac{1}{\epsilon_\pi} \coth \frac{\epsilon_\pi}{2T} + \dots$$

with

$$\epsilon_\pi = \sqrt{k^2 + 2\partial_{\sigma^2} U_k}$$



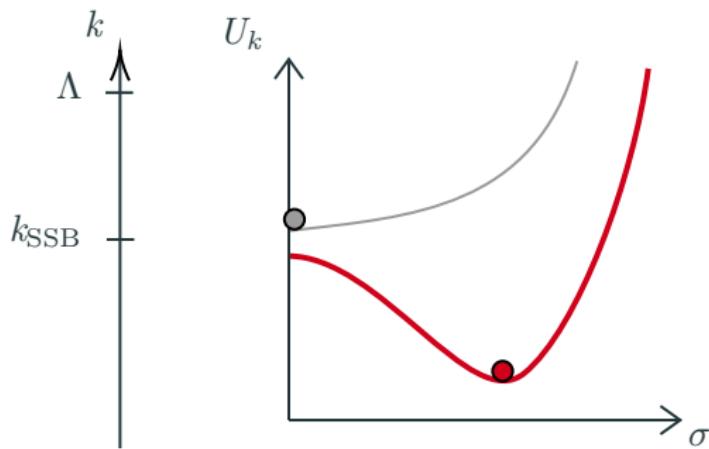
First Trouble: Bosonic Diquark Loops

Pion loops

$$\partial_t U_k = \frac{1}{\epsilon_\pi} \coth \frac{\epsilon_\pi}{2T} + \dots$$

with

$$\epsilon_\pi = \sqrt{k^2 + 2\partial_{\sigma^2} U_k}$$



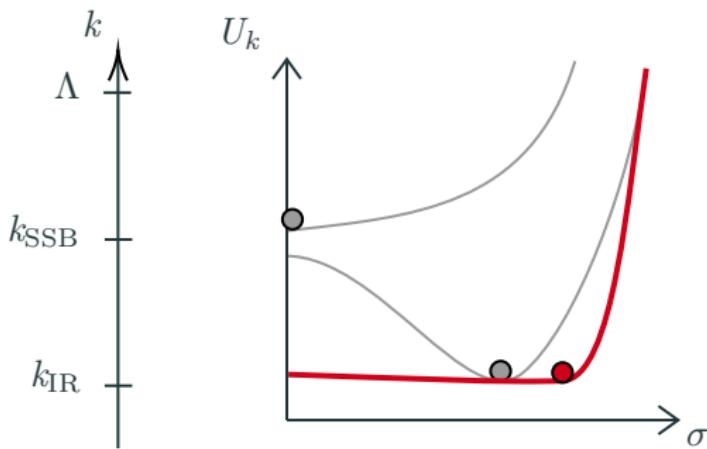
First Trouble: Bosonic Diquark Loops

Pion loops

$$\partial_t U_k = \frac{1}{\epsilon_\pi} \coth \frac{\epsilon_\pi}{2T} + \dots$$

with

$$\epsilon_\pi = \sqrt{k^2 + 2\partial_{\sigma^2} U_k}$$



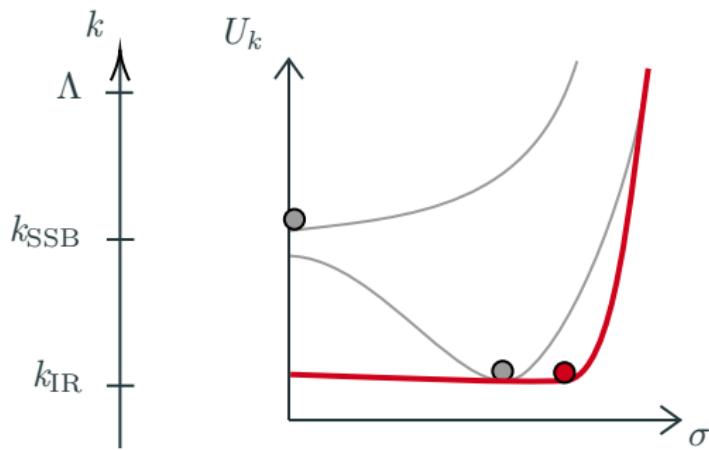
First Trouble: Bosonic Diquark Loops

Pion loops

$$\partial_t U_k = \frac{1}{\epsilon_\pi} \coth \frac{\epsilon_\pi}{2T} + \dots$$

with

$$\epsilon_\pi = \sqrt{k^2 + 2\partial_{\sigma^2} U_k}$$

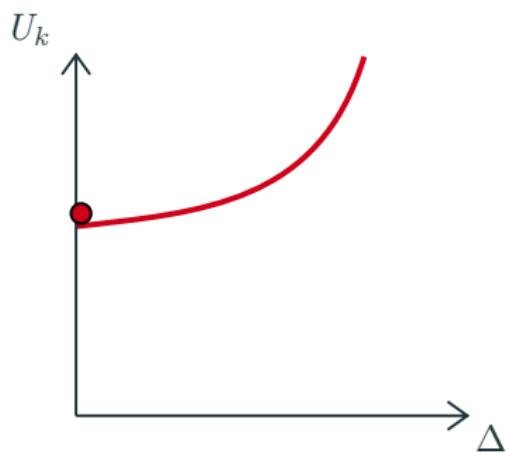


Diquark loops

$$\partial_t U_k = \frac{1}{\epsilon_\Delta} \coth \frac{\epsilon_\Delta - 2\mu}{2T} + \dots$$

with

$$\epsilon_\Delta = \sqrt{k^2 + 2\partial_{\Delta^2} U_k}$$



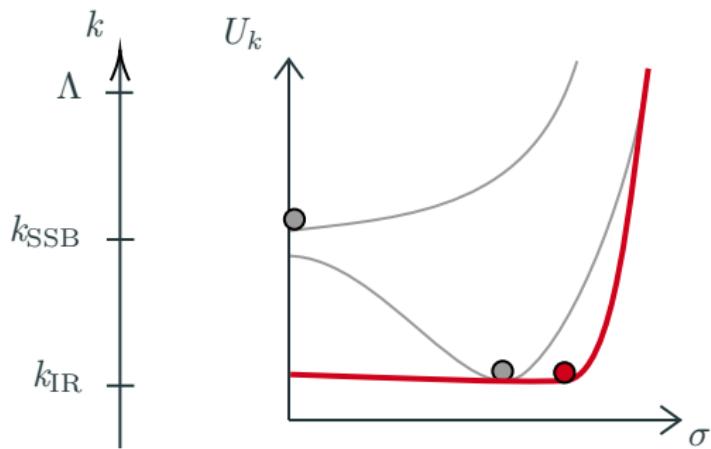
First Trouble: Bosonic Diquark Loops

Pion loops

$$\partial_t U_k = \frac{1}{\epsilon_\pi} \coth \frac{\epsilon_\pi}{2T} + \dots$$

with

$$\epsilon_\pi = \sqrt{k^2 + 2\partial_{\sigma^2} U_k}$$

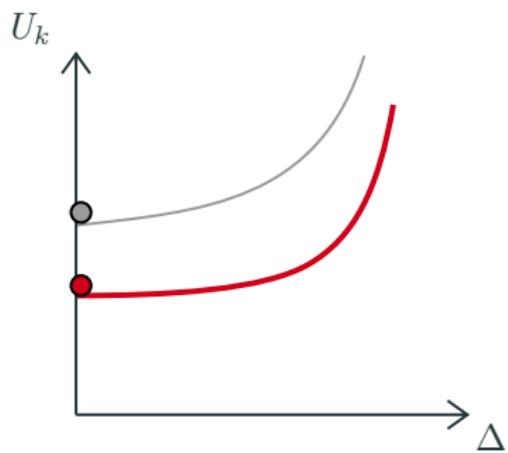


Diquark loops

$$\partial_t U_k = \frac{1}{\epsilon_\Delta} \coth \frac{\epsilon_\Delta - 2\mu}{2T} + \dots$$

with

$$\epsilon_\Delta = \sqrt{k^2 + 2\partial_{\Delta^2} U_k}$$



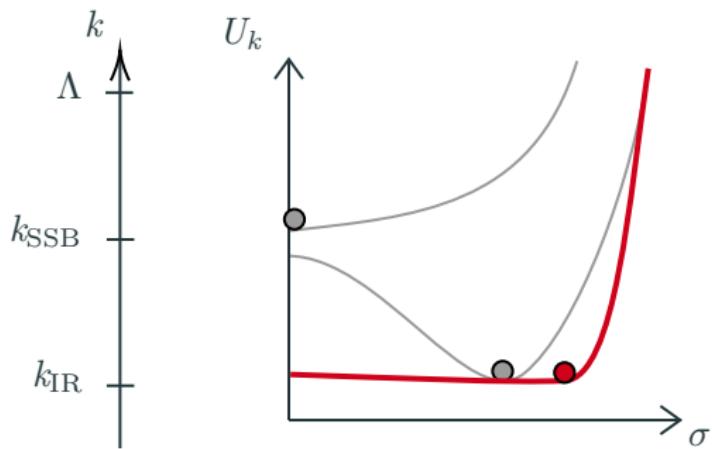
First Trouble: Bosonic Diquark Loops

Pion loops

$$\partial_t U_k = \frac{1}{\epsilon_\pi} \coth \frac{\epsilon_\pi}{2T} + \dots$$

with

$$\epsilon_\pi = \sqrt{k^2 + 2\partial_{\sigma^2} U_k}$$

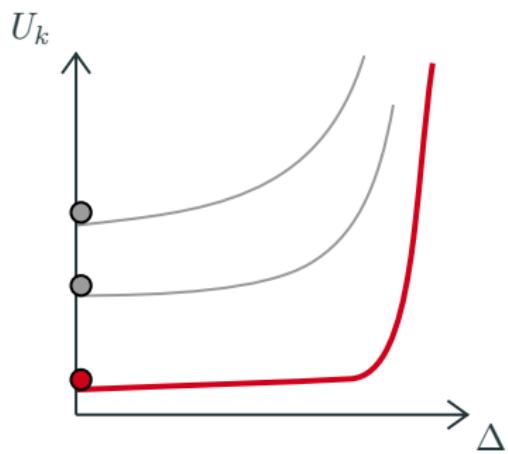


Diquark loops

$$\partial_t U_k = \frac{1}{\epsilon_\Delta} \coth \frac{\epsilon_\Delta - 2\mu}{2T} + \dots$$

with

$$\epsilon_\Delta = \sqrt{k^2 + 2\partial_{\Delta^2} U_k}$$



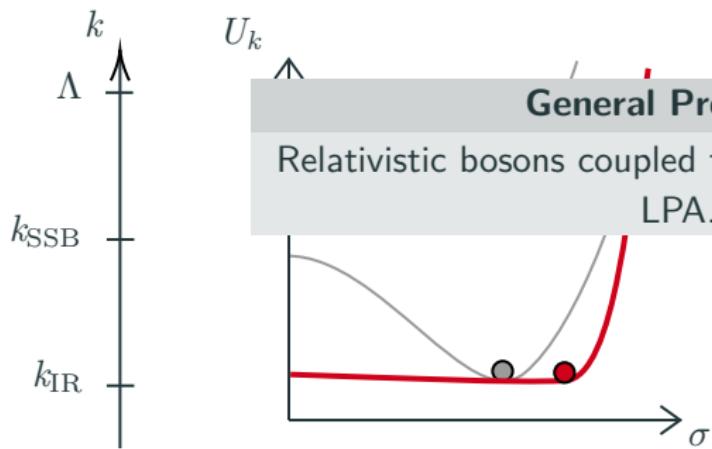
First Trouble: Bosonic Diquark Loops

Pion loops

$$\partial_t U_k = \frac{1}{\epsilon_\pi} \coth \frac{\epsilon_\pi}{2T} + \dots$$

with

$$\epsilon_\pi = \sqrt{k^2 + 2\partial_{\sigma^2} U_k}$$

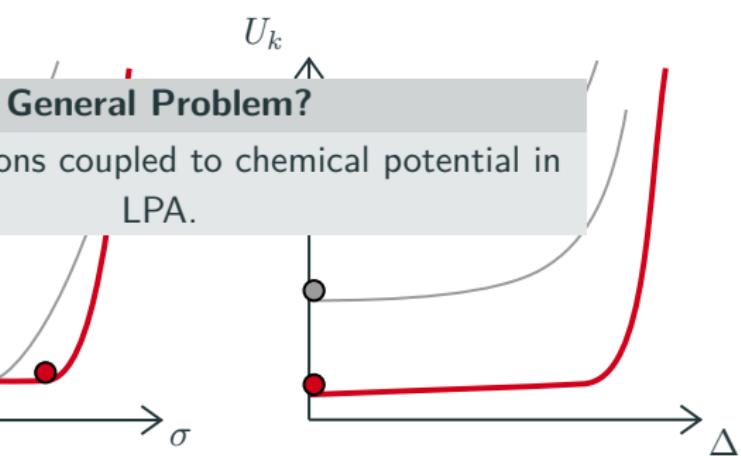


Diquark loops

$$\partial_t U_k = \frac{1}{\epsilon_\Delta} \coth \frac{\epsilon_\Delta - 2\mu}{2T} + \dots$$

with

$$\epsilon_\Delta = \sqrt{k^2 + 2\partial_{\Delta^2} U_k}$$

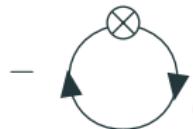


General Problem?

Relativistic bosons coupled to chemical potential in LPA.

Second Trouble: Gapped Quark Loop

Gapped quark loop:



$\textcolor{red}{q}_r, \textcolor{green}{q}_g$

$$\propto \frac{\epsilon_k - \mu}{\sqrt{(\epsilon_k - \mu)^2 + g_\Delta^2 \Delta^2}} \xrightarrow{\partial_{\Delta^2}} \propto \frac{\epsilon_k - \mu}{((\epsilon_k - \mu)^2 + g_\Delta^2 \Delta^2)^{3/2}}$$

diverges at the Fermi-surface



Second Trouble: Gapped Quark Loop

Gapped quark loop:

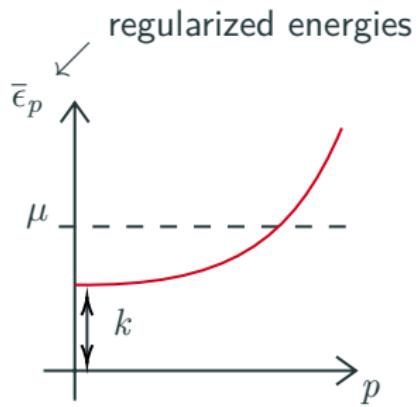
-  $q_r, q_g \propto \frac{\epsilon_k - \mu}{\sqrt{(\epsilon_k - \mu)^2 + g_\Delta^2 \Delta^2}} \xrightarrow{\partial_{\Delta^2}} \propto \frac{\epsilon_k - \mu}{((\epsilon_k - \mu)^2 + g_\Delta^2 \Delta^2)^{3/2}}$

diverges at the Fermi-surface



Possible solution: **flow around the Fermi-surface**

[Braun, Dörnfeld, Schallmo, Töpfel; 2008.05978]



usual momentum

regulator

Second Trouble: Gapped Quark Loop

Gapped quark loop:

-  $q_r, q_g \propto \frac{\epsilon_k - \mu}{\sqrt{(\epsilon_k - \mu)^2 + g_\Delta^2 \Delta^2}} \xrightarrow{\partial_{\Delta^2}} \infty \frac{\epsilon_k - \mu}{((\epsilon_k - \mu)^2 + g_\Delta^2 \Delta^2)^{3/2}}$

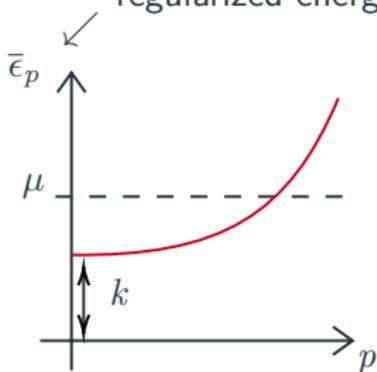
diverges at the Fermi-surface



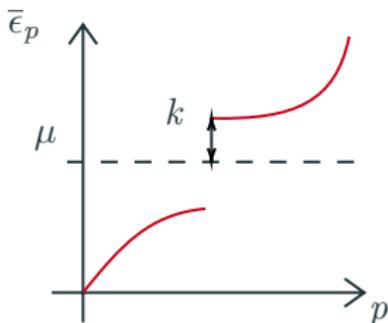
Possible solution: **flow around the Fermi-surface**

[Braun, Dörnfeld, Schallmo, Töpfel; 2008.05978]

regularized energies



usual momentum
regulator



Fermi-surface
regulator

Second Trouble: Gapped Quark Loop

Gapped quark loop:

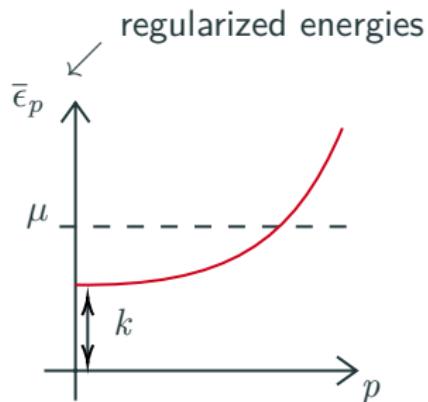
-  $q_r, q_g \propto \frac{\epsilon_k - \mu}{\sqrt{(\epsilon_k - \mu)^2 + g_\Delta^2 \Delta^2}} \xrightarrow{\partial_{\Delta^2}} \infty \frac{\epsilon_k - \mu}{((\epsilon_k - \mu)^2 + g_\Delta^2 \Delta^2)^{3/2}}$

diverges at the Fermi-surface

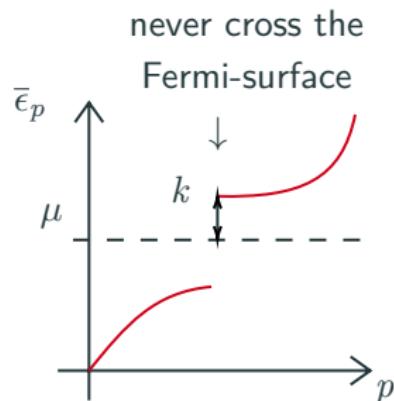


Possible solution: **flow around the Fermi-surface**

[Braun, Dörnfeld, Schallmo, Töpfel; 2008.05978]



usual momentum
regulator

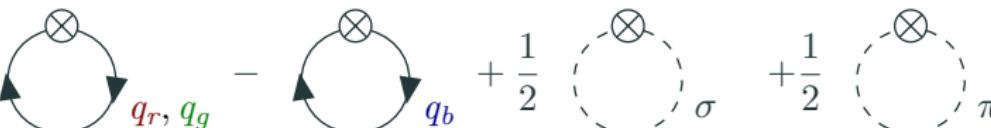


Fermi-surface
regulator

Phase Structure and Astrophysical Applications

Quark-meson-diquark Model without Diquark Loops

Diquarks only at the mean field level: none of the previously mentioned problems are present.

$$\partial_t U_k = - \left(q_r, q_g \right) - \left(q_b \right) + \frac{1}{2} \left(\sigma \right) + \frac{1}{2} \left(\pi \right)$$


Quark-meson-diquark Model without Diquark Loops

Diquarks only at the mean field level: none of the previously mentioned problems are present.

$$\partial_t U_k = - \underset{\uparrow}{\text{---}} \circlearrowleft \otimes \circlearrowright \text{---} q_r, q_g - \circlearrowleft \otimes \circlearrowright \text{---} q_b + \frac{1}{2} \circlearrowleft \otimes \circlearrowright \text{---} \sigma + \frac{1}{2} \circlearrowleft \otimes \circlearrowright \text{---} \pi$$

diquarks still present here

$$E_{\Delta}^{\pm} = \sqrt{(\epsilon_k \pm \mu)^2 + g_{\Delta}^2 \Delta^2}$$

Quark-meson-diquark Model without Diquark Loops

Diquarks only at the mean field level: none of the previously mentioned problems are present.

$$\partial_t U_k = - \left(q_r, q_g \right) - \left(q_b \right) + \frac{1}{2} \left(\sigma \right) + \frac{1}{2} \left(\pi \right)$$

Diagram illustrating the model components:

- Left term: A solid circle with a clockwise arrow and a crossed circle symbol inside. Below it is the text q_r, q_g .
- Middle term: A solid circle with a clockwise arrow and a crossed circle symbol inside. Below it is the text q_b .
- Right term: Two dashed circles with clockwise arrows and crossed circle symbols inside. The first is labeled σ and the second is labeled π .

Annotations:

- A vertical arrow points upwards from the q_r, q_g term to the q_b term, with the text "diquarks still present here".
- A horizontal arrow points from the σ term to the π term, with the text "still include fluctuations from pion and sigma mesons".

$$E_{\Delta}^{\pm} = \sqrt{(\epsilon_k \pm \mu)^2 + g_{\Delta}^2 \Delta^2}$$

Quark-meson-diquark Model without Diquark Loops

Diquarks only at the mean field level: none of the previously mentioned problems are present.

$$\partial_t U_k = - \left(q_r, q_g \right) - \left(q_b \right) + \frac{1}{2} \left(\sigma \right) + \frac{1}{2} \left(\pi \right)$$

Diagram illustrating the model terms:

- q_r, q_g : A solid circle with a clockwise arrow and a crossed circle symbol (\otimes) at the top.
- q_b : A solid circle with a clockwise arrow.
- σ : A dashed circle with a clockwise arrow and a crossed circle symbol (\otimes) at the top.
- π : A dashed circle with a clockwise arrow.

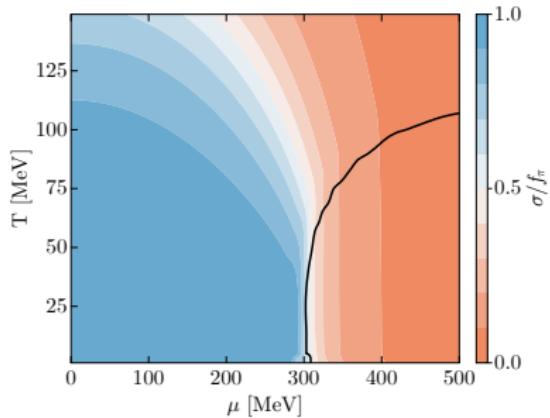
Annotations:

- "diquarks still present here" is written below the first two terms.
- "still include fluctuations from pion and sigma mesons" is written below the last two terms.

$$E_{\Delta}^{\pm} = \sqrt{(\epsilon_k \pm \mu)^2 + g_{\Delta}^2 \Delta^2}$$

- Look at phase structure.
- First astrophysical applications.

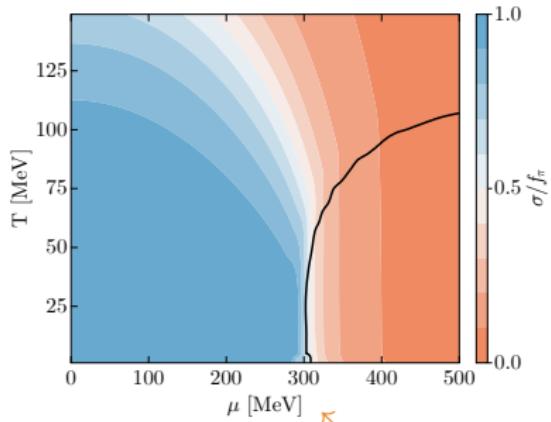
Phase Diagram



Phase diagram.

- Expected phase structure.

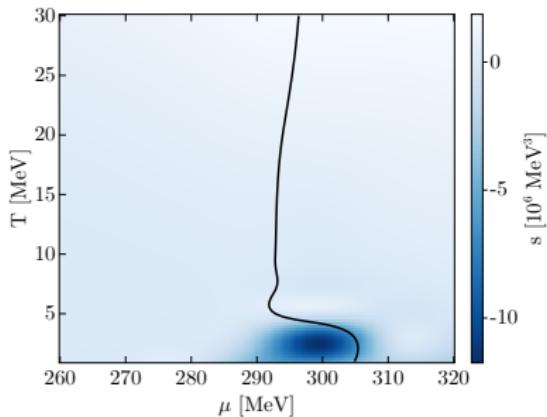
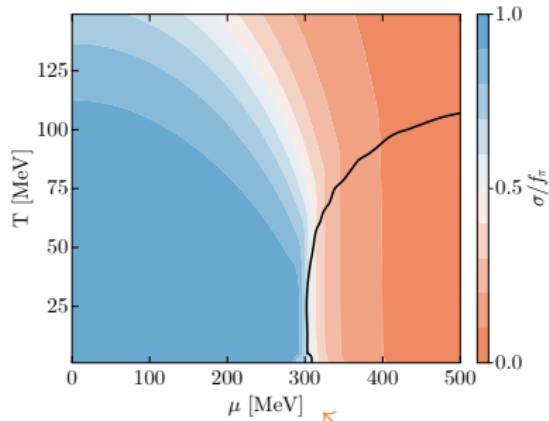
Phase Diagram



Phase diagram What's going on here?

- Expected phase structure.

Phase Diagram



Phase diagram What's going on here? Entropy density s .

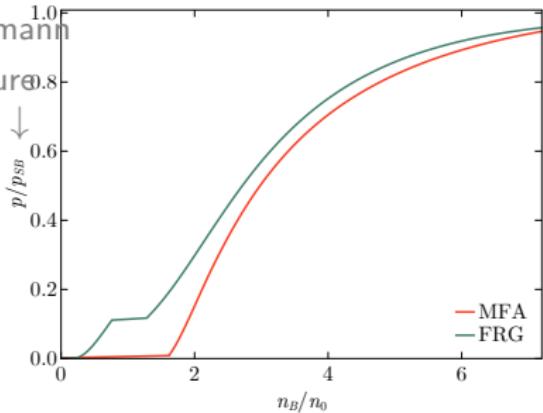
- Expected phase structure.
- Litim regulator in LPA: expect negative entropy **but** diquarks reduce the size of the region.

Equation of State

Stefan-

Boltzmann

Pressure

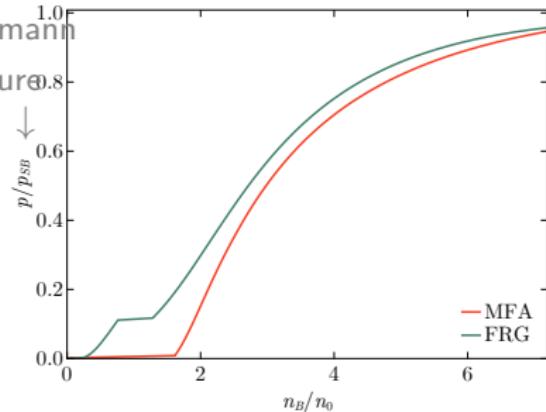


Equation of state $p(n)$ at $T = 1$

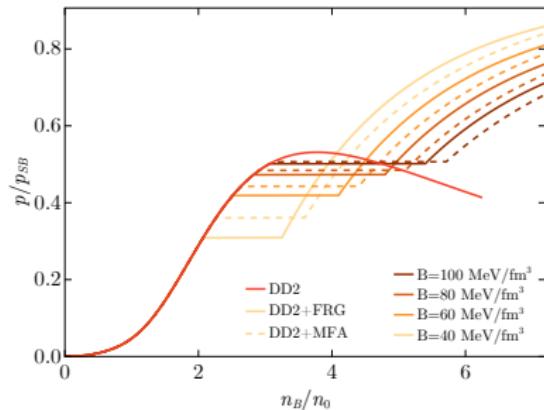
Equation of State

Stefan-

Boltzmann



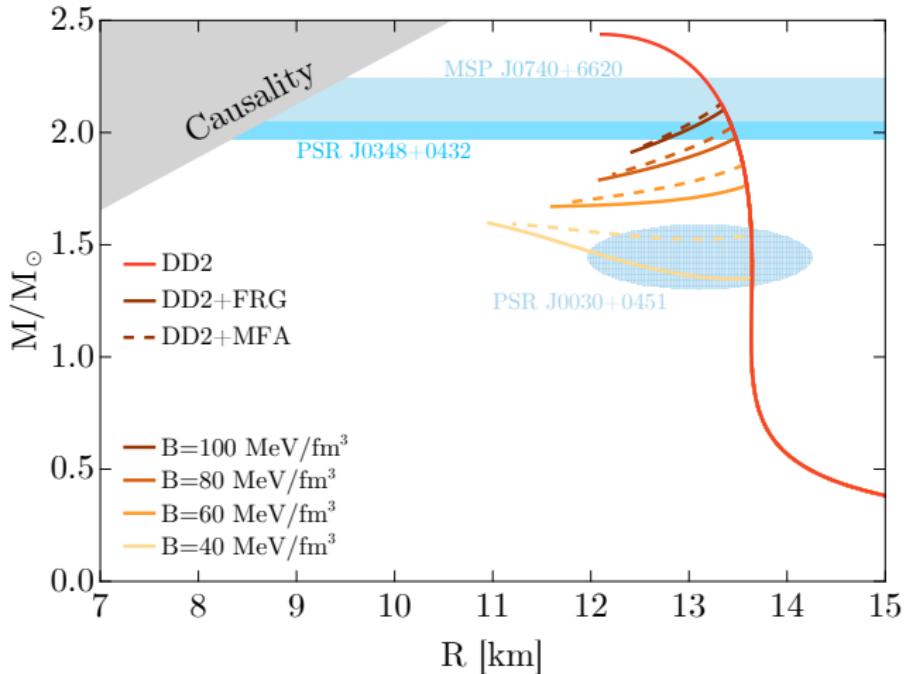
Equation of state $p(n)$ at $T = 1$



Maxwell construct with hadronic eos DD2

- Maxwell construct for astrophysical applications.
- No crossing with hadronic equation of state (DD2) → introduce **bag constant** B .

Mass-radius Relationship



- Superconducting core → mostly unstable with current diquark parameters.

Charge Neutral Matter

More realistic EoS: impose **neutrality conditions**

Charge neutrality (and β -equilibrium)

$$\frac{2}{3}n_u - \frac{1}{3}n_d - n_e = 0 \quad (\text{and } u \leftrightarrow d + e^+ + \nu_e)$$

Need more down quark than up quark \rightarrow introduce different chemical potential: μ_{up} and μ_{down} .

Charge Neutral Matter

More realistic EoS: impose **neutrality conditions**

Charge neutrality (and β -equilibrium)

$$\frac{2}{3}n_u - \frac{1}{3}n_d - n_e = 0 \quad (\text{and } u \leftrightarrow d + e^+ + \nu_e)$$

Need more down quark than up quark \rightarrow introduce different chemical potential: μ_{up} and μ_{down} .

- Stress on **2SC pairing** controlled by $\mu_Q = \mu_{\text{up}} - \mu_{\text{down}}$.

Charge Neutral Matter

More realistic EoS: impose **neutrality conditions**

Charge neutrality (and β -equilibrium)

$$\frac{2}{3}n_u - \frac{1}{3}n_d - n_e = 0 \quad (\text{and } u \leftrightarrow d + e^+ + \nu_e)$$

Need more down quark than up quark \rightarrow introduce different chemical potential: μ_{up} and μ_{down} .

- Stress on 2SC pairing controlled by $\mu_Q = \mu_{\text{up}} - \mu_{\text{down}}$.
- Also color neutrality $n_r = n_b = n_g \rightarrow$ not a problem for 2SC.

Charge Neutral Matter

More realistic EoS: impose **neutrality conditions**

Charge neutrality (and β -equilibrium)

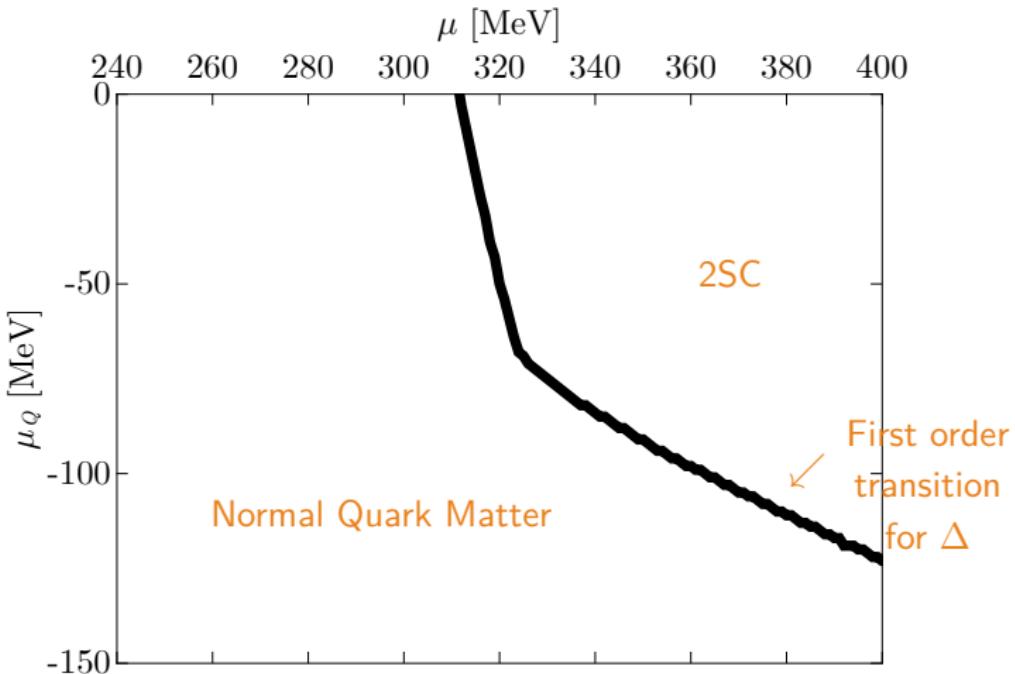
$$\frac{2}{3}n_u - \frac{1}{3}n_d - n_e = 0 \quad (\text{and} \quad u \leftrightarrow d + e^+ + \nu_e)$$

Need more down quark than up quark \rightarrow introduce different chemical potential: μ_{up} and μ_{down} .

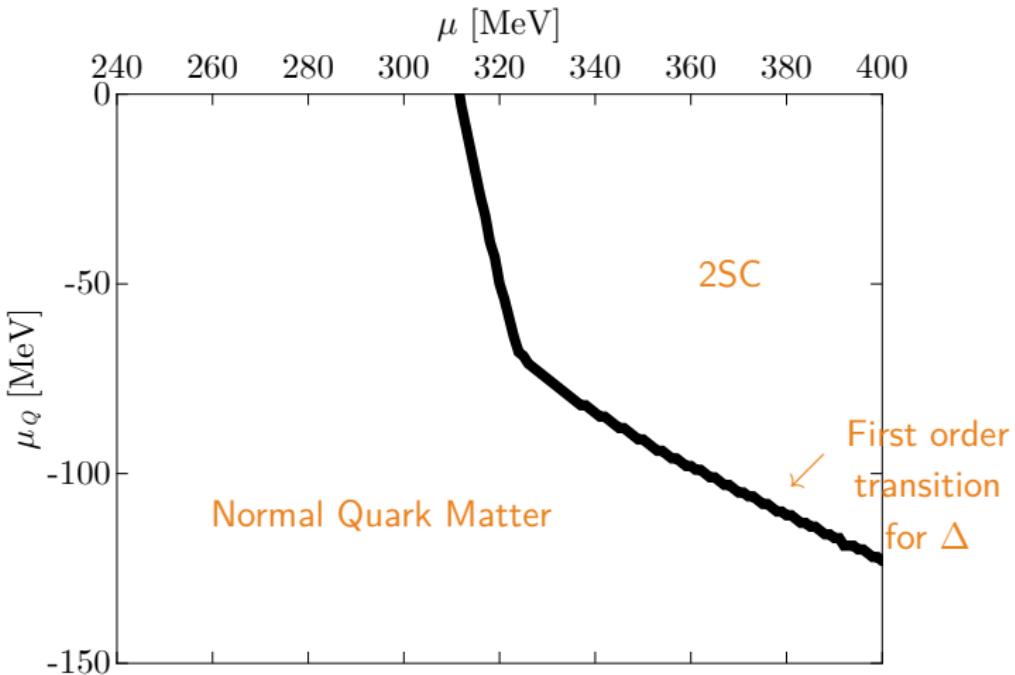
- Stress on 2SC pairing controlled by $\mu_Q = \mu_{\text{up}} - \mu_{\text{down}}$.
- Also color neutrality $n_r = n_b = n_g \rightarrow$ not a problem for 2SC.

Next: preliminary results in meanfield at $T = 1$ MeV.

Phase structure in the (μ, μ_Q) plane



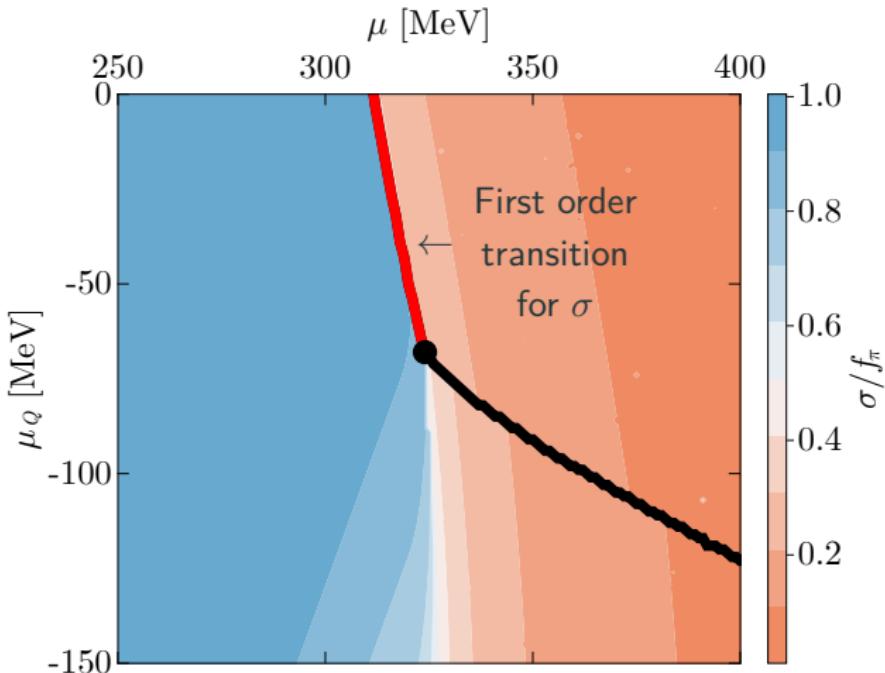
Phase structure in the (μ, μ_Q) plane



Question

What happen to the chiral condensate?

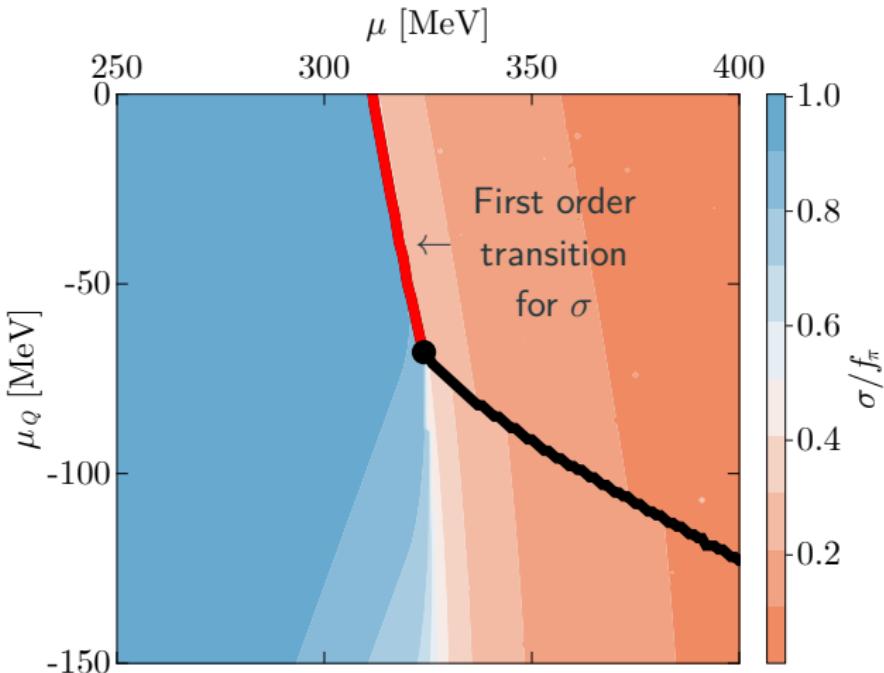
Phase structure in the (μ, μ_Q) plane



Question

What happen to the chiral condensate?

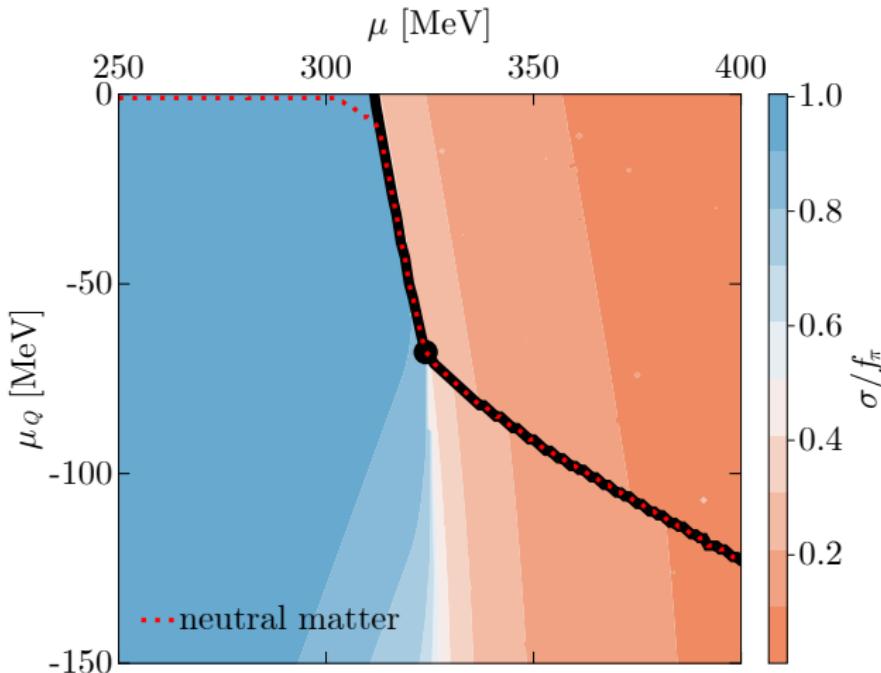
Phase structure in the (μ, μ_Q) plane



Question

Where is neutral matter in this phase diagram?

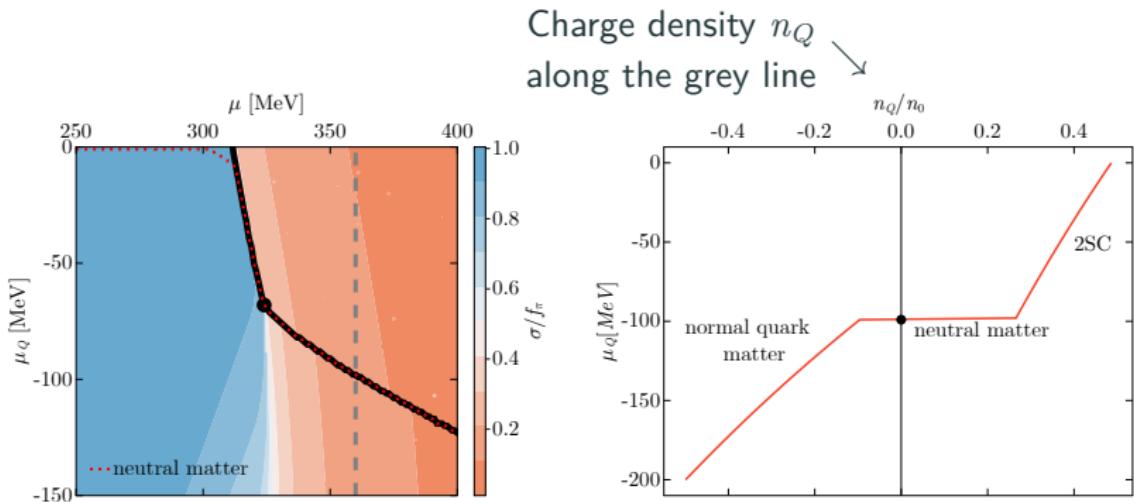
Phase structure in the (μ, μ_Q) plane



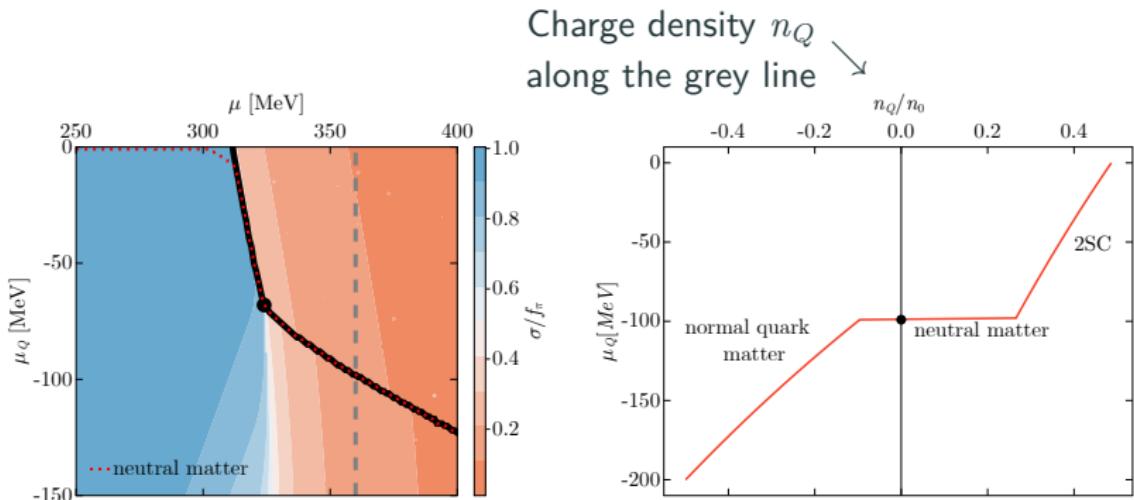
Question

Where is neutral matter in this phase diagram?

Neutral Matter and First Order Transition



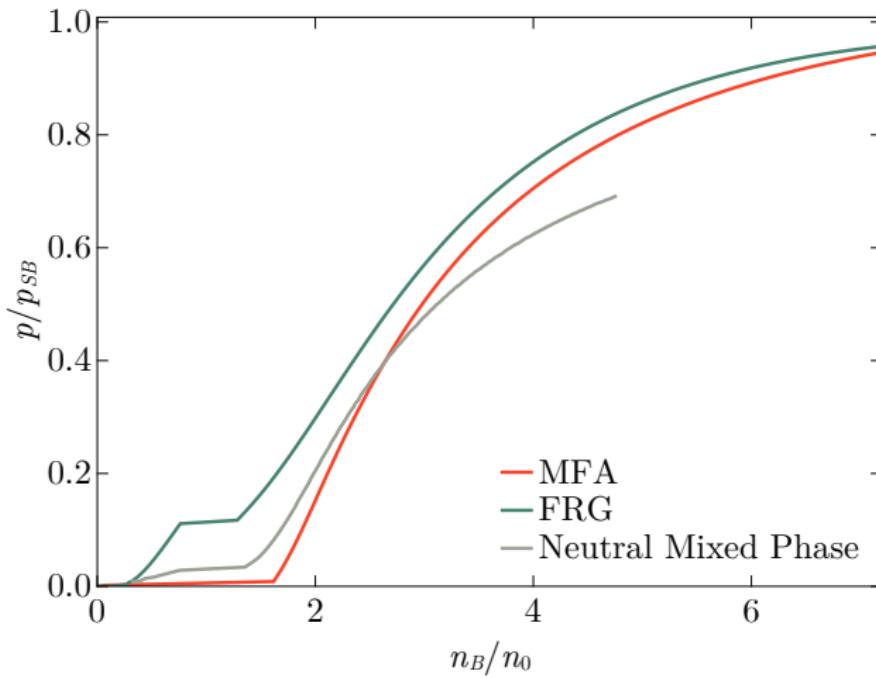
Neutral Matter and First Order Transition



- Density of neutral matter: inside a **first order region**.
- Suggest the presence of a **mixed phase**: normal quark matter + 2SC phase.

Mixed Equation of State

Neutral equation of state: right amount of normal quark matter and 2SC phase to reach neutrality.



Summary and Outlook

- Quark-meson-diquark model: model chiral transition and 2SC color superconducting phase.
- FRG resolution faces two problems:
 - Diquarks couple to μ : cannot flow from symmetry restored phase to symmetry broken phase.
 - Divergence at the Fermi-surface: possible resolution with Fermi-surface regulator.
- Negative entropy after the chiral transition: better with diquarks at mean-field level.
- Neutrality condition suggest the presence of a mixed phase: normal quark matter and 2SC phase.

Backup Slides

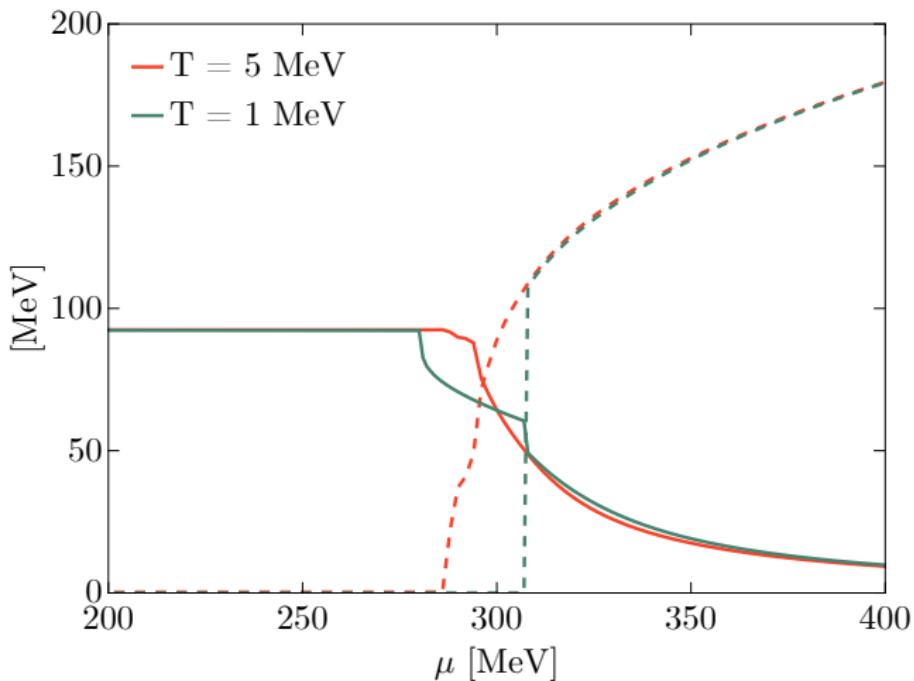
Full Flow Equation

$$\begin{aligned}
\partial_t U_k = & \frac{k^5}{12\pi^2} \left\{ \frac{3}{\epsilon_\pi} \coth \frac{\epsilon_\pi}{2T} + \frac{2}{\epsilon_{\Delta,0}} \left[\coth \frac{\epsilon_{\Delta,0} - 2\mu}{2T} + \coth \frac{\epsilon_{\Delta,0} + 2\mu}{2T} \right] \right. \\
& + \sum_{i=1}^3 \frac{\alpha_2 z_i^4 - \alpha_1 z_i^2 + \alpha_0}{(z_i^2 - z_{i+1}^2)(z_i^2 - z_{i+2}^2)} \frac{1}{z_i} \coth \frac{z_i}{2T} \Big\} \\
& - \frac{k^5}{3\pi^2} \left\{ \frac{2}{\epsilon_k} \left[\frac{E_k^+}{E_\Delta^+} \tanh \frac{E_\Delta^+}{2T} + \frac{E_k^-}{E_\Delta^-} \tanh \frac{E_\Delta^-}{2T} \right] \right. \\
& \left. \left. + \frac{1}{\epsilon_k} \left[\tanh \frac{\epsilon_k^+}{2T} + \tanh \frac{\epsilon_k^-}{2T} \right] \right\}
\end{aligned}$$

with

$$\begin{aligned}
\epsilon_k^\pm &= \sqrt{k^2 + g_\phi^2 \rho^2} \pm \mu = \epsilon_k \pm \mu & E_\pi &= \sqrt{k^2 + 2U_{k,\rho}} \\
E_\Delta^\pm &= \sqrt{(\epsilon_k \pm \mu)^2 + g_\Delta^2 d^2} & \epsilon_{\Delta,0} &= \sqrt{k^2 + 2U_{k,d}}
\end{aligned}$$

Chiral and diquark condensates



Mixed Phase

