

# Quark matter in Neutron Stars

Bernd-Jochen Schaefer



September 12<sup>th</sup>, 2023



**EMMI Workshop**  
**Functional Methods in Strongly Correlated  
Systems (FUNSCS2023)**  
Hirschgärtl, Kleinwalsertal, Austria, September 10 - September 15, 2023

# agenda

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- Hybrid and quark star matter based on a nonperturbative equation of state

Konstantin Otto (Giessen U.), Micaela Oertel (LUTH, Meudon), Bernd-Jochen Schaefer (Giessen U.)

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- Nonperturbative quark matter equations of state with vector interactions

Konstantin Otto (Giessen U.), Micaela Oertel (LUTH, Meudon), Bernd-Jochen Schaefer (Giessen U.)

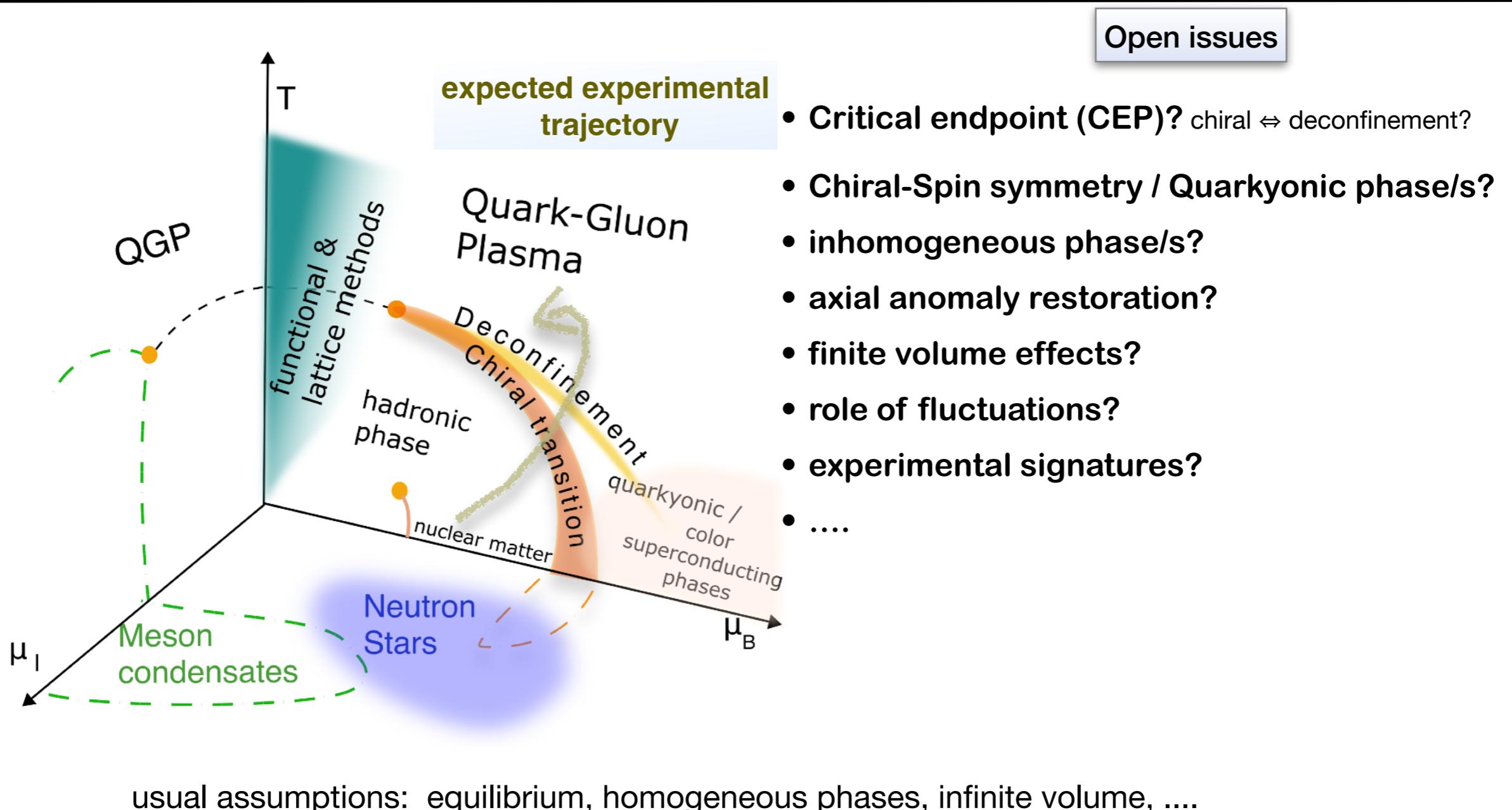
Published in: *Eur.Phys.J.ST* 229 (2020) 22-23, 3629-3649 • e-Print: [2007.07394](#) [hep-ph]

- Regulator scheme dependence of the chiral phase transition at high densities

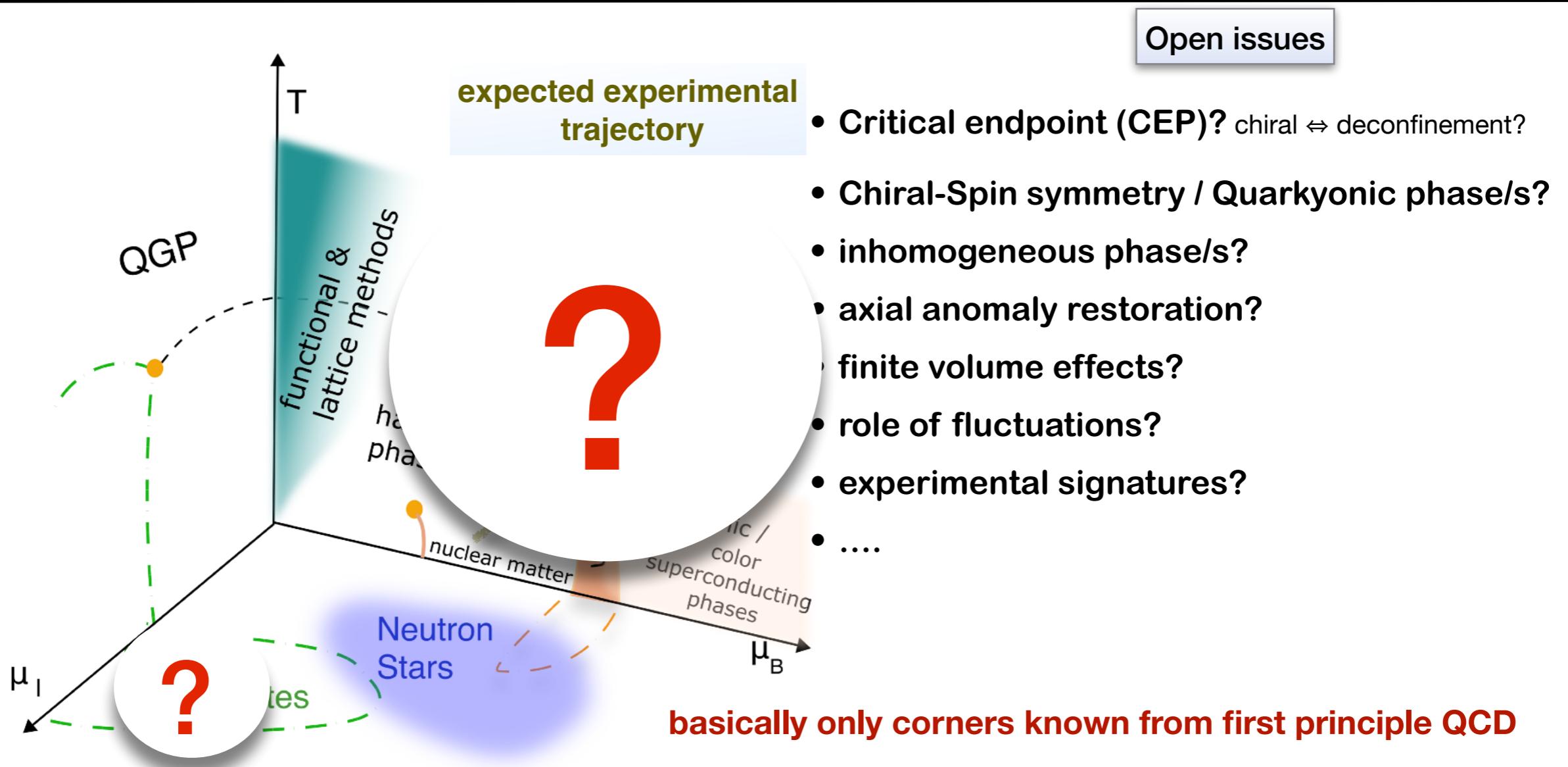
Konstantin Otto (Giessen U.), Christopher Busch (Giessen U.), Bernd-Jochen Schaefer (Giessen U.)

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# conjectured QCD phase structure



# conjectured QCD phase structure



alternative to HIC to probe cold dense QCD matter  $\rightarrow$  massive neutron star (e.g. PSR J0348+0432 )

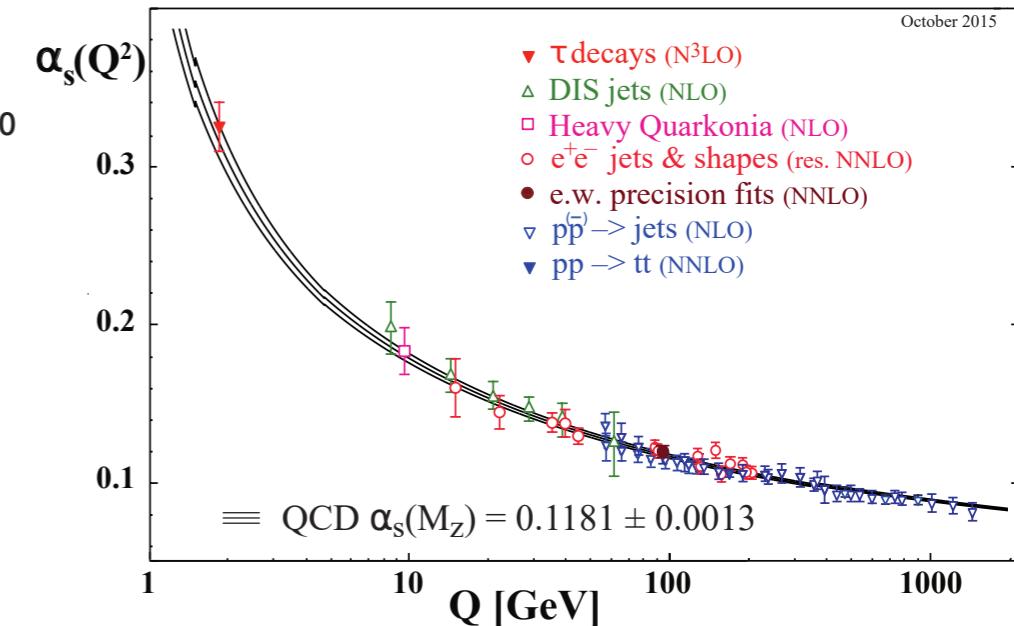
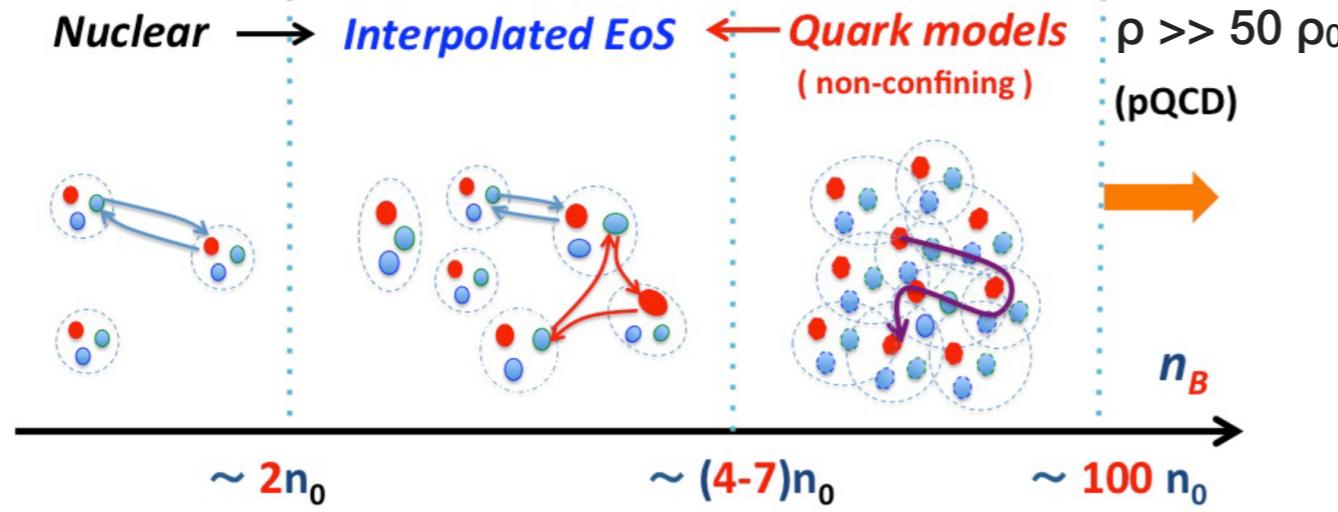
cold dense QCD matter: only effective low-energy realisation of QCD: e.g. (P)QM models

ultimate goal: microscopic description of EoS guided by QCD first principle

# EoS for dense matter

Kojo et al.: three window model of dense matter

[Baym 2018]



Nuclear phase:  
1-2 meson/quark  
exchanges

EoS from  
**nuclear physics**  
 $\rho < 2\rho_0$   $\chi$ EFT

interpolated EoS  
many meson/quark  
exchanges

system gradually changes  
from hadronic to quark matter  
- diquarks, colored quarks virtually ...  
- role of strangeness / hyperons

$2\rho_0 < \rho < 7\rho_0$  Neutron stars

Quark phase:  
quarks no longer  
specific to baryons

**mostly mean-field investigations**  
like NJL-type or phenomenological  
models

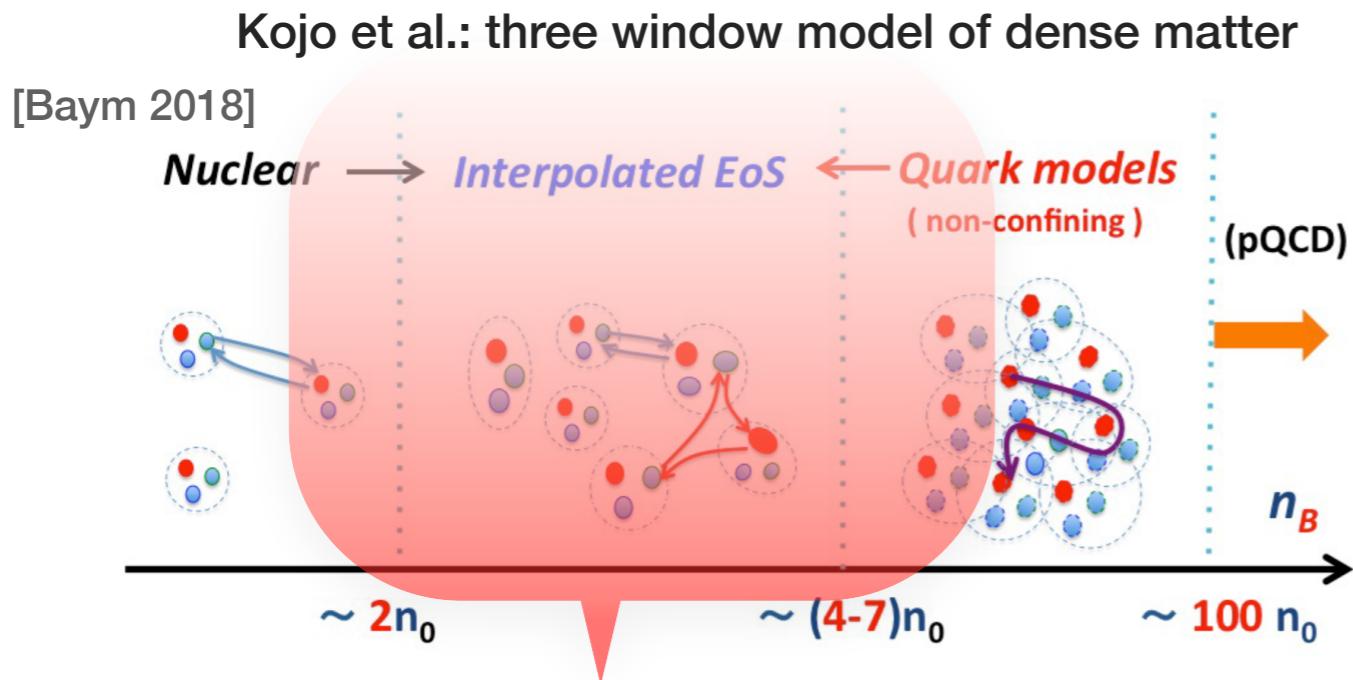
→ upgrade with FRG methods

[Hebeler, Lattimer, Pethick, Schwenk et al. 2010 ]

[Schaffner-Bielich et al. 2008 ]

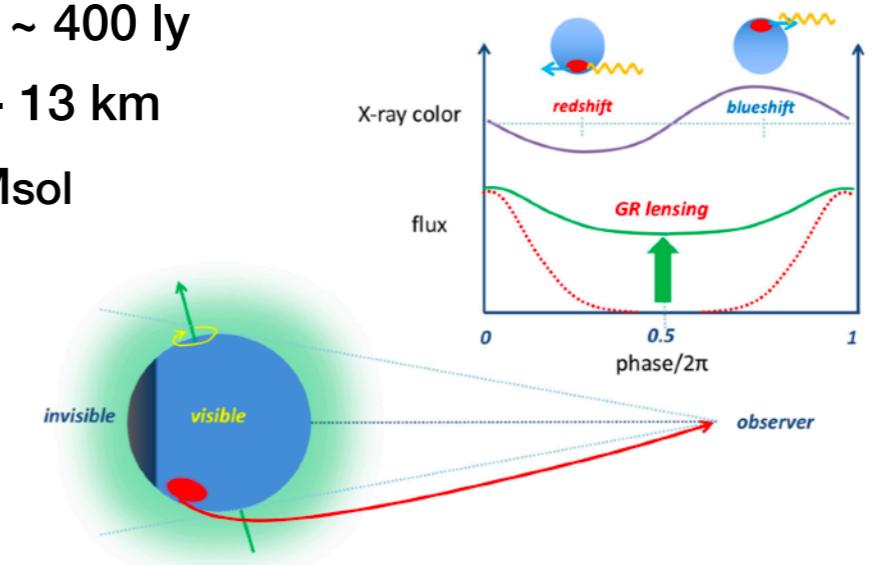
[Blaschke, Fischer, Oertel et al. 2018 ]

# experimental facts



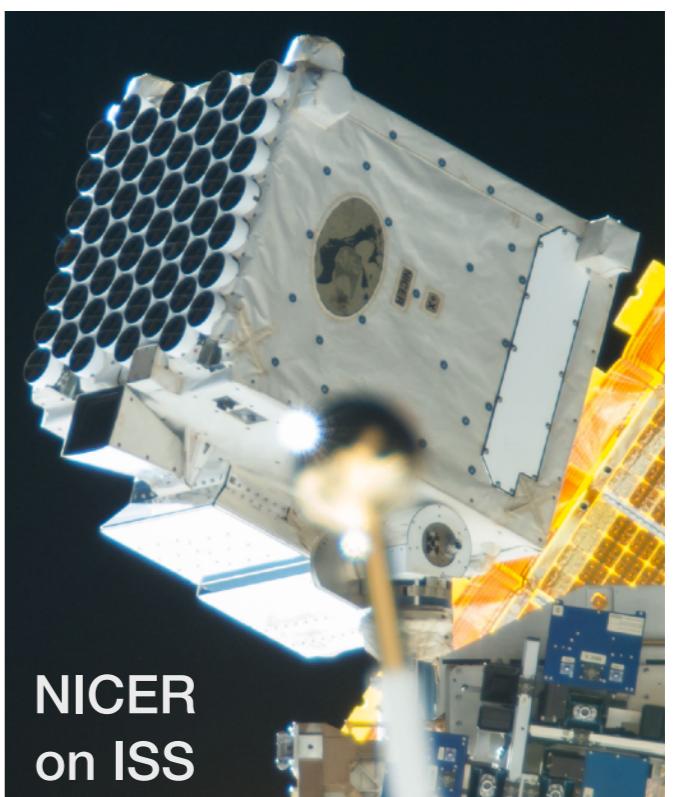
- $M_{\max} = 2.05 M_{\odot}$  &  $R \sim 12.8$  km ;  $\rho_{\text{central}} \sim 5\rho_0$
- $M = 1.4 M_{\odot}$  &  $R = 12.8$  km ;  $\Delta R = 0$
- $M = 2.4 M_{\odot}$  &  $R = 12.8$  km ;  $\sim$  as  $1.4 M_{\odot}$  stars
- PSR J0740+6620 & J0030+0451  
**both  $M (> M_{\odot})$  &  $R$  measured**  
GW170817 constraints tidal deformability  $\rightarrow R$
- **detection phase transition possible:**  
increase post-merger dominant oscillation frequency  
only a few events expected (even w/ 3rd generation detectors)

nearest  $\sim 400$  ly  
 $R \sim 10 - 13$  km  
 $M > 2 M_{\odot}$



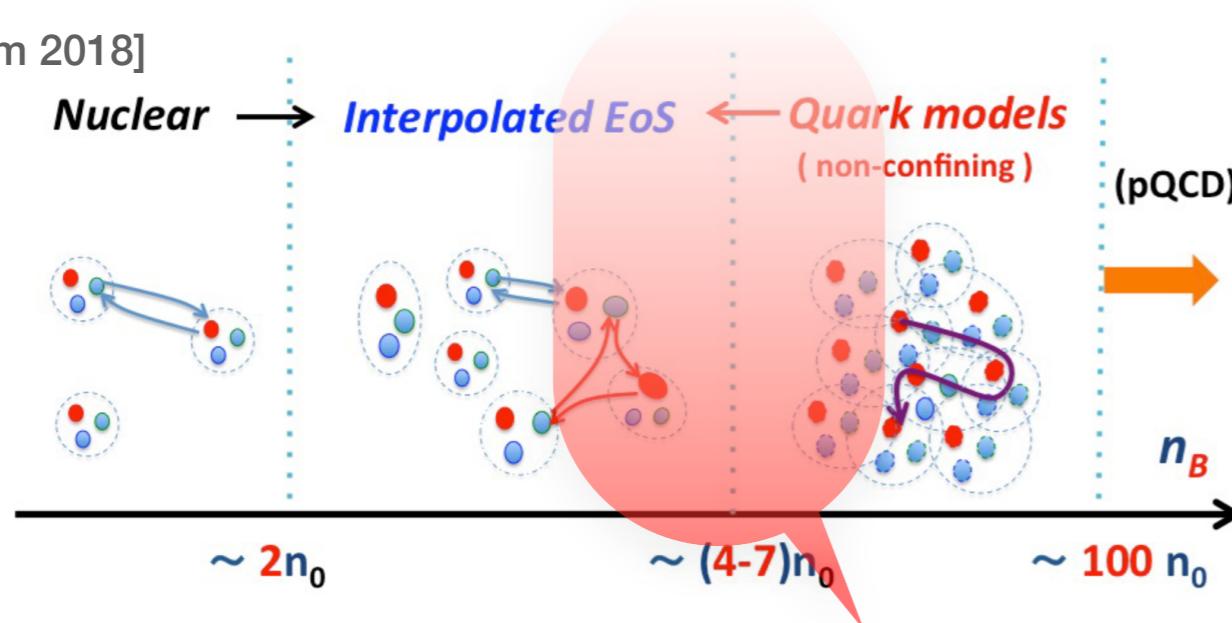
future:  $\sim 2035$

- new 3<sup>rd</sup> generation detectors:
- (American) Cosmic Explorer
- (European) Einstein Telescope



# transition from hadronic to quark matter

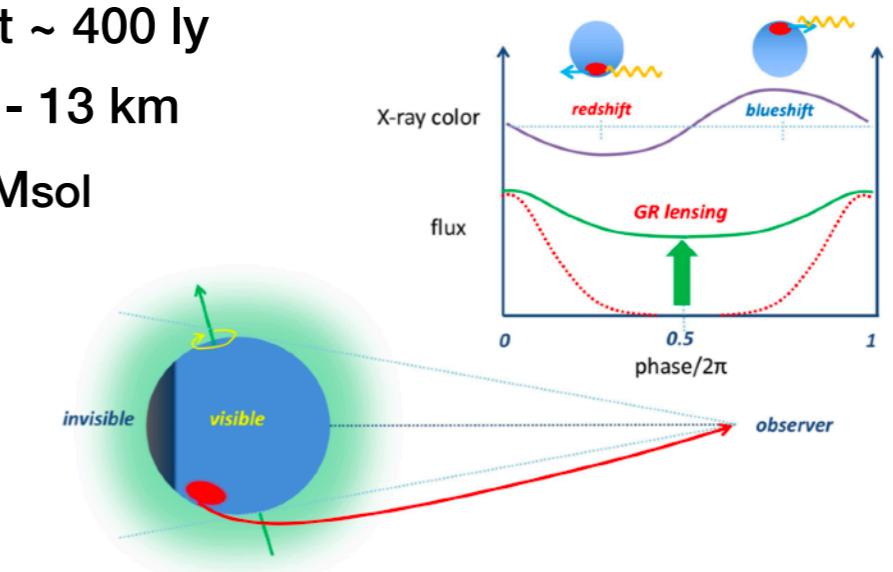
[Baym 2018]



nearest  $\sim 400$  ly

$R \sim 10 - 13$  km

$M > 2$  Msol



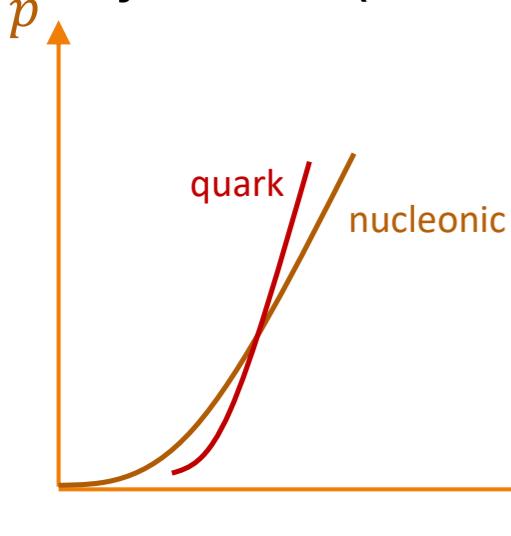
several possibilities (if transition): Maxwell construction or continuous interpolation

first-order transition

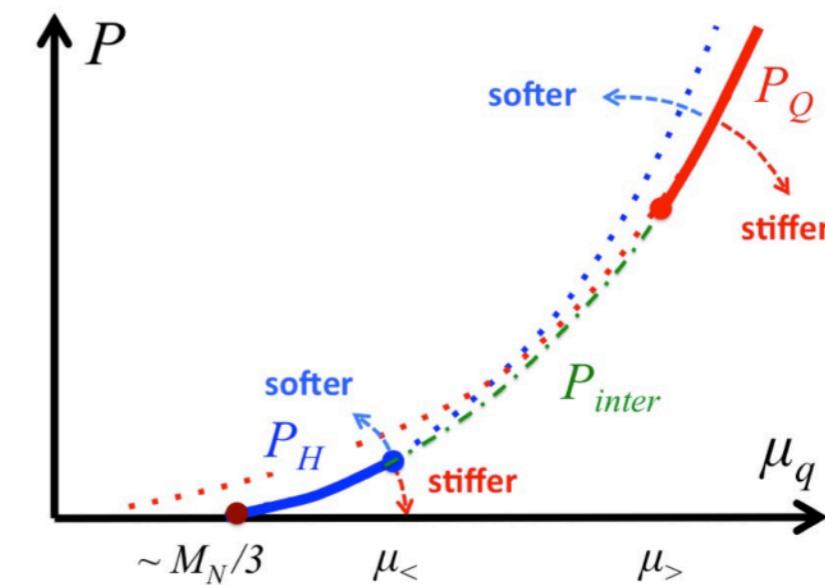


quark-hadron continuity

hybrid star (w/ small quark core)

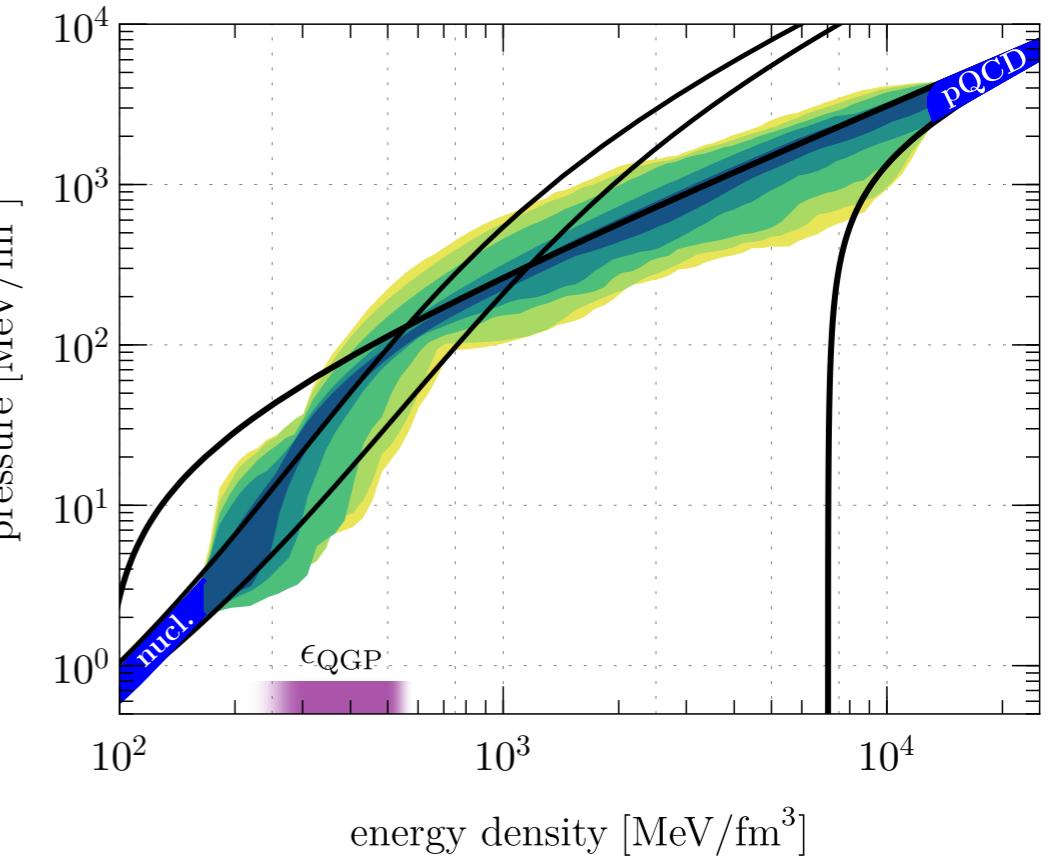
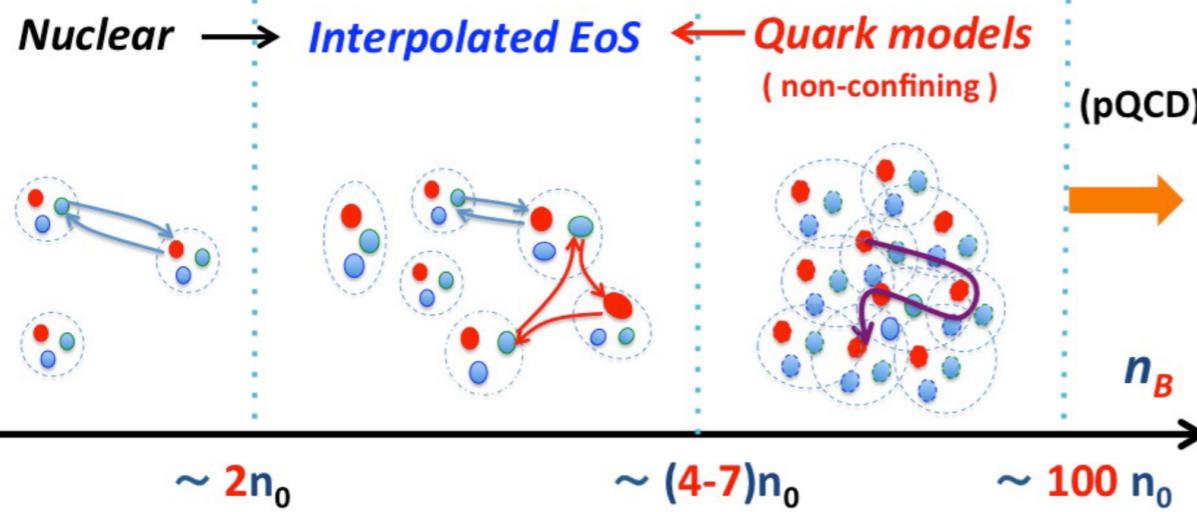


→ quark EoS  
softer than  
nucleonic EoS  
@ higher densities



# conflicting constraints on EoS

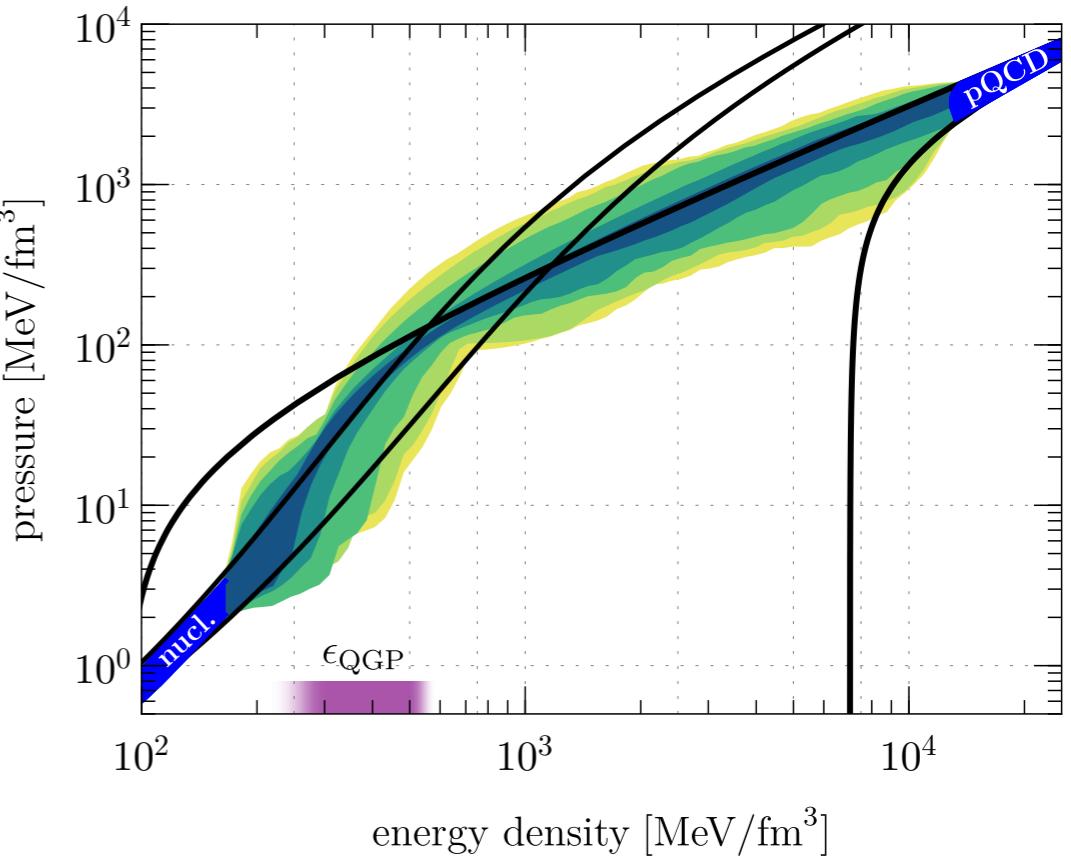
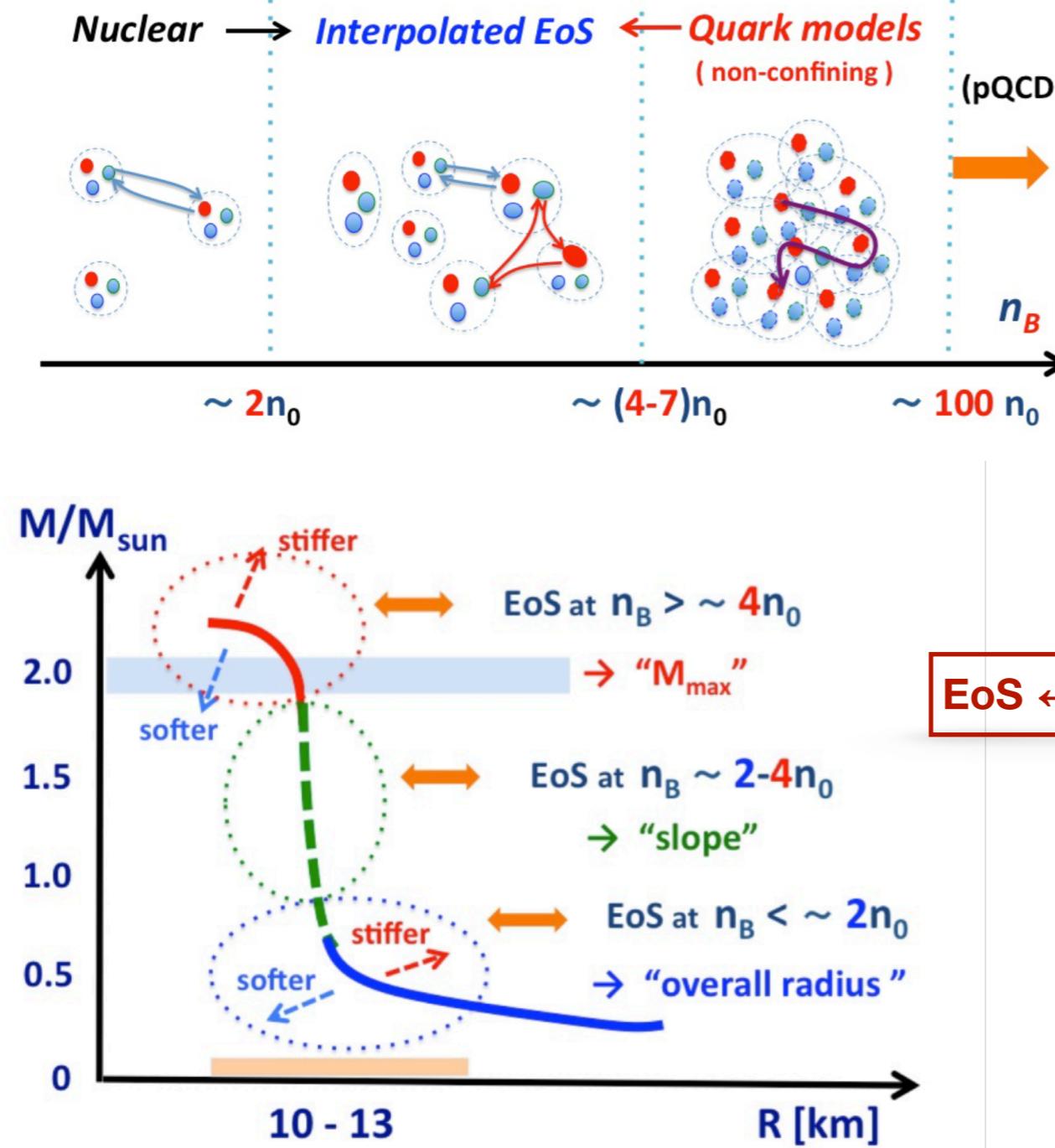
[Baym 2018]



EoS  $\leftrightarrow$  TOV equation  $\leftrightarrow$  M-R relation (observables)

# conflicting constraints on EoS

[Baym 2018]



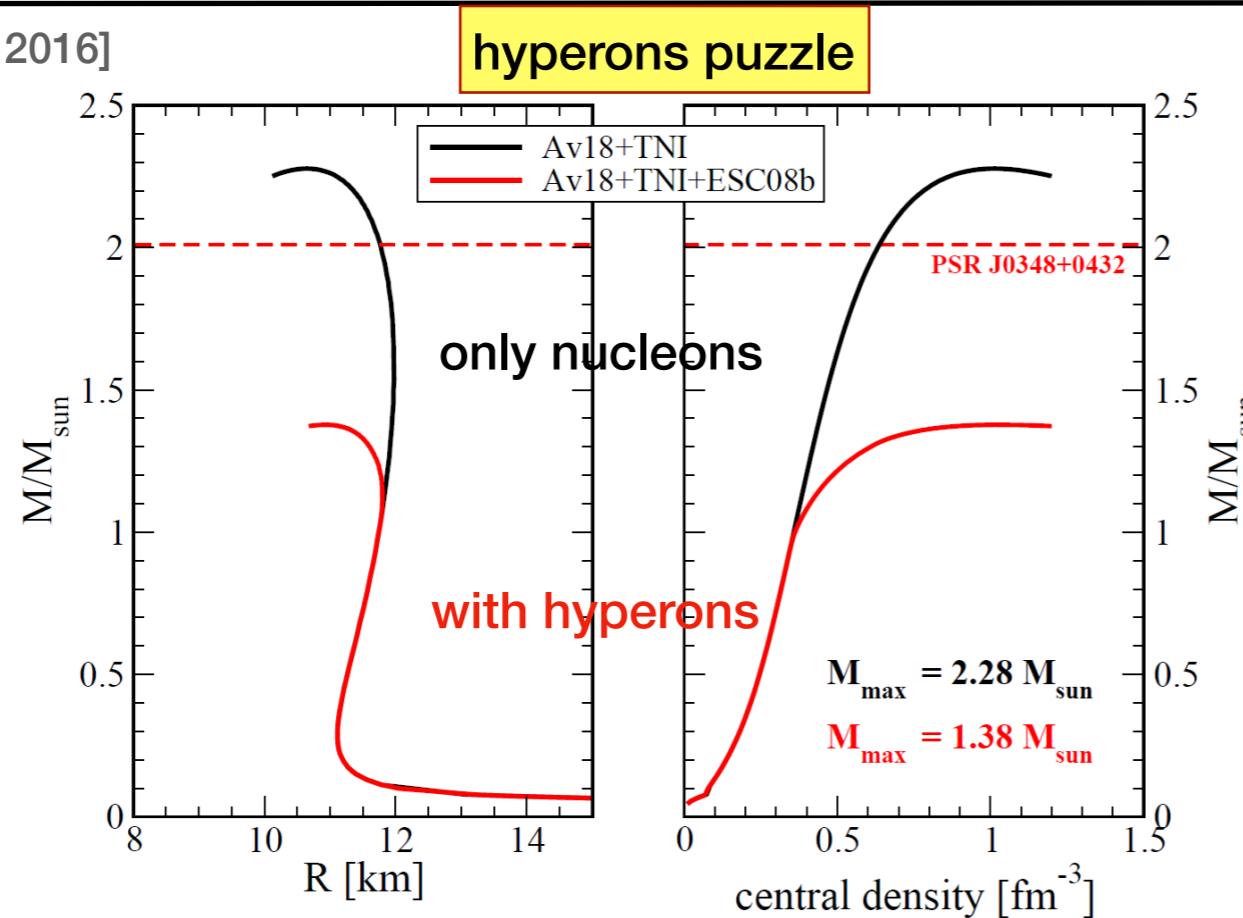
**EoS ↔ TOV equation ↔ M-R relation (observables)**

three constraints on the EoS:

1. stiff enough (@high density) →  $2M_\odot$
2. soft enough (@low density) → Radius
3. speed of sound < 1

# unsolved puzzles / open issues

[Bombaci 2016]



Further constraints:

causality

charge neutrality:  $n_p = n_e + n_\mu$

$\beta$ -equilibrium:  $\mu_n = \mu_p + \mu_e$

simplification:

→ electrons and muons as  
free Fermi gas in EoS

General problems (physical theory input required):

→ hyperon puzzle

onset of strangeness in hadronic phase or quark phase

→ soften EoS

[Djapo, BJS, Wambach 2010]

→ masquerade problem

many EoS look similar → similar M-R relation  
increasing #dof soften EoS,  
repulsive interactions stiffen EoS

[Alvarez-Castillo, Blaschke 2014]

# Functional Renormalization Group

## ■ Wetterich Equation (average effective action)

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \partial_t R_k \left( \frac{1}{\Gamma_k^{(2)} + R_k} \right)$$

$k \partial_k \Gamma_k[\phi] \sim \frac{1}{2}$

[Wetterich 1993]

$$t = \ln(k/\Lambda)$$

$$\Gamma_k^{(2)} = \frac{\delta^2 \Gamma_k}{\delta \phi \delta \phi}$$

$R_k$  regulators

shape function conditions:

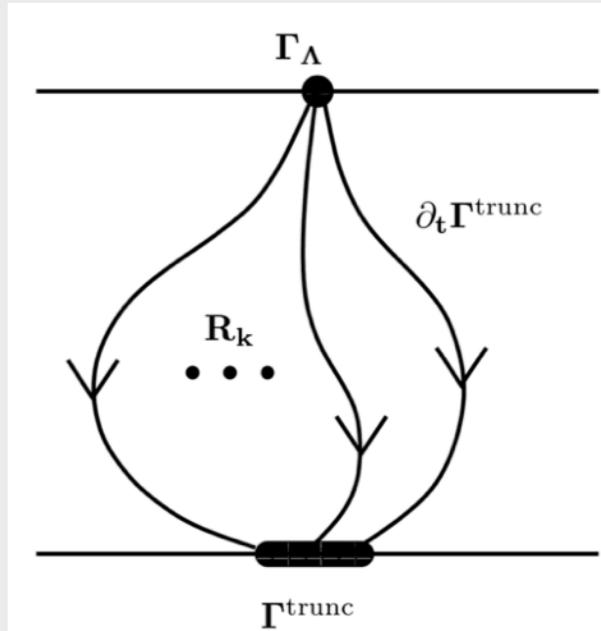
- $R_k(p^2) = p^2 r(p^2/k^2)$
- $\lim_{p^2/k^2 \rightarrow \infty} R_k(p^2) = 0$
- $\lim_{p^2/k^2 \rightarrow 0} R_k(p^2) > 0 (= k^2)$
- $\lim_{k \rightarrow \infty} R_k(p^2) \rightarrow \infty$

## ■ Ansatz effective action    Quark-Meson truncation in LPA (LO derivative expansion)

$$\Gamma_k = \int d^4x \bar{q} [i\gamma_\mu \partial^\mu - g(\sigma + i\vec{\tau}\vec{\pi}\gamma_5)] q + \frac{1}{2}(\partial_\mu \sigma)^2 + \frac{1}{2}(\partial_\mu \vec{\pi})^2 + V_k(\phi^2)$$

$$V_{k=\Lambda}(\phi^2) = \frac{\lambda}{4}(\sigma^2 + \vec{\pi}^2 - v^2)^2 - c\sigma$$

arbitrary potential



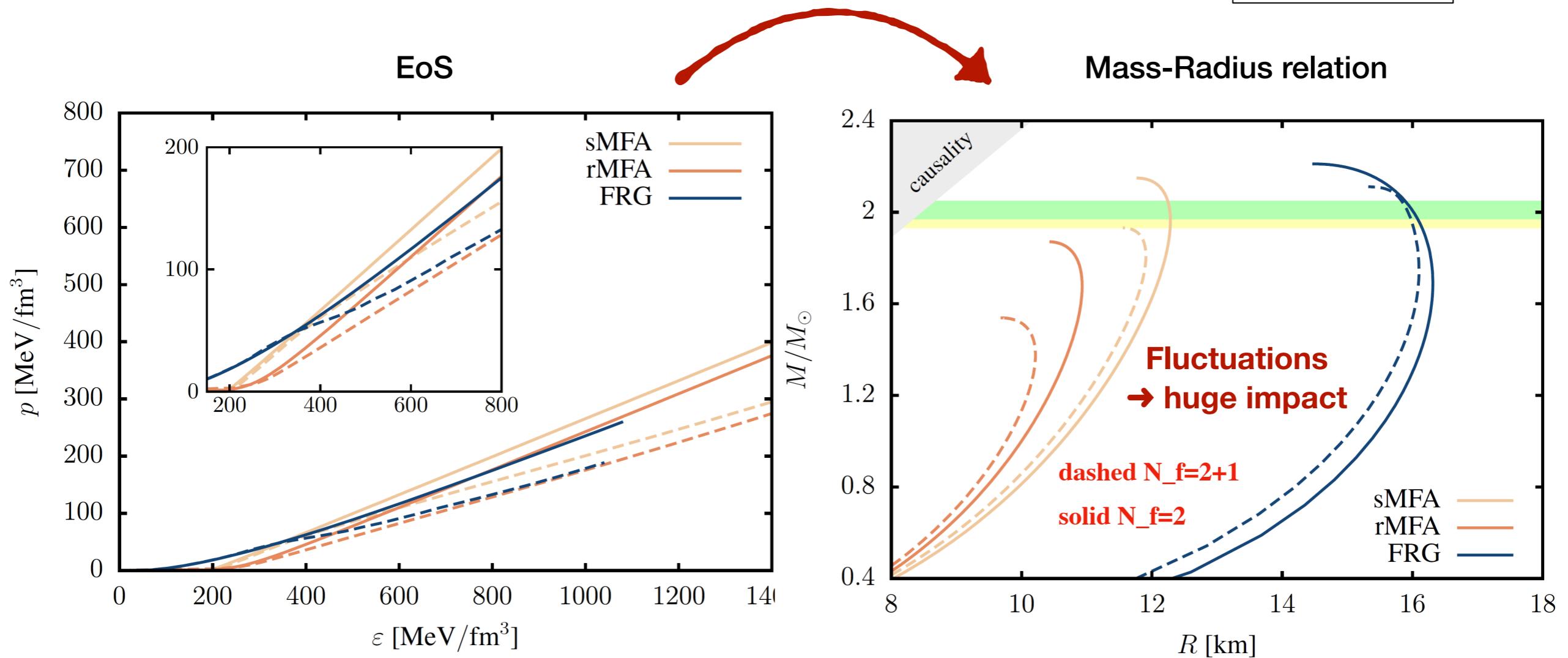
# Impact of fluctuations on EoS

[Otto, Oertel, BJS 2020]

Impose  $\beta$ -equilibrium and charge neutrality conditions

Tolman-Oppenheimer-Volkoff  
equations

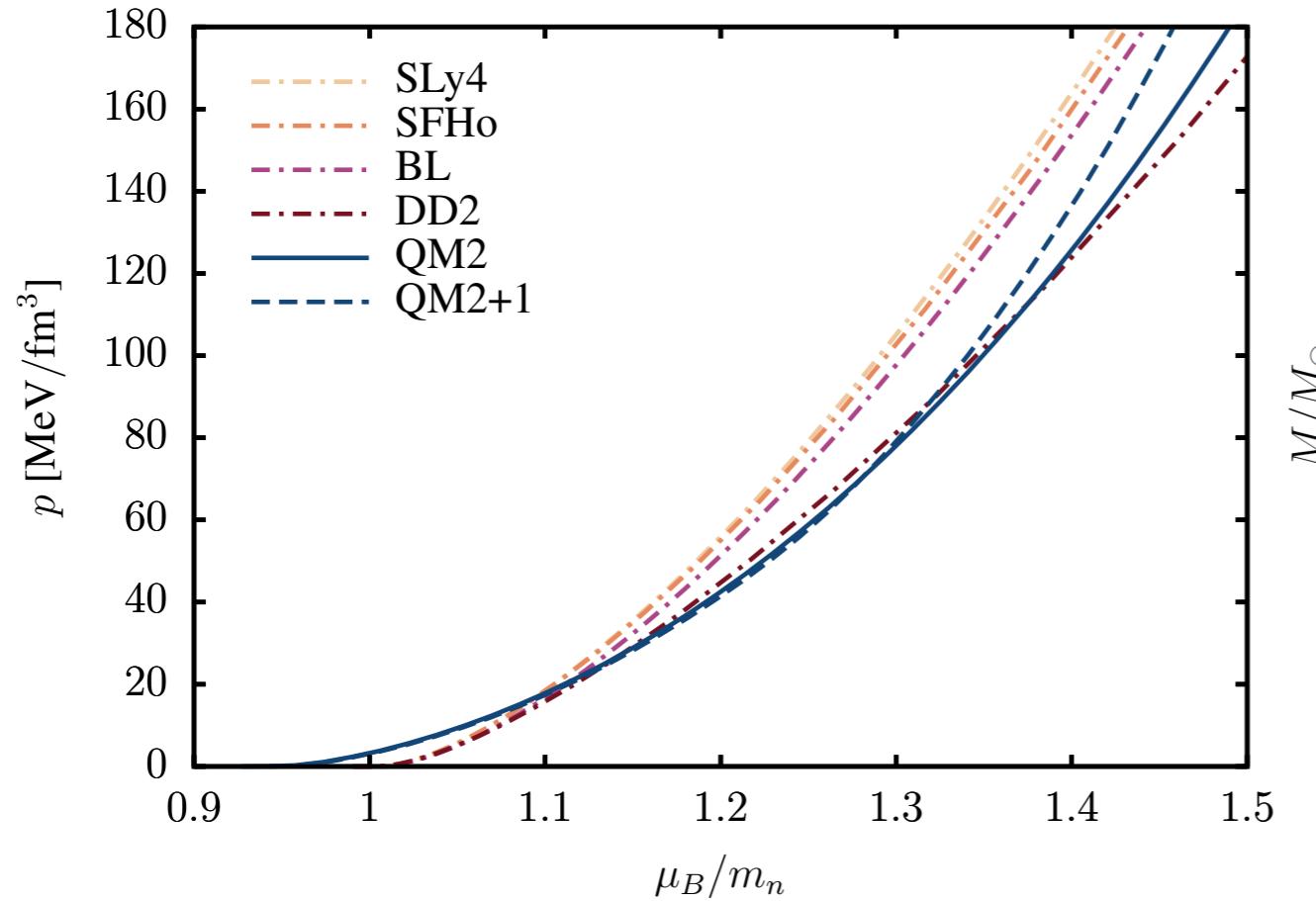
$$\begin{aligned}\mu_u &= \mu_q - \frac{2}{3} \mu_e \\ \mu_d &= \mu_q + \frac{1}{3} \mu_e \\ \mu_s &= \mu_q + \frac{1}{3} \mu_e\end{aligned}$$



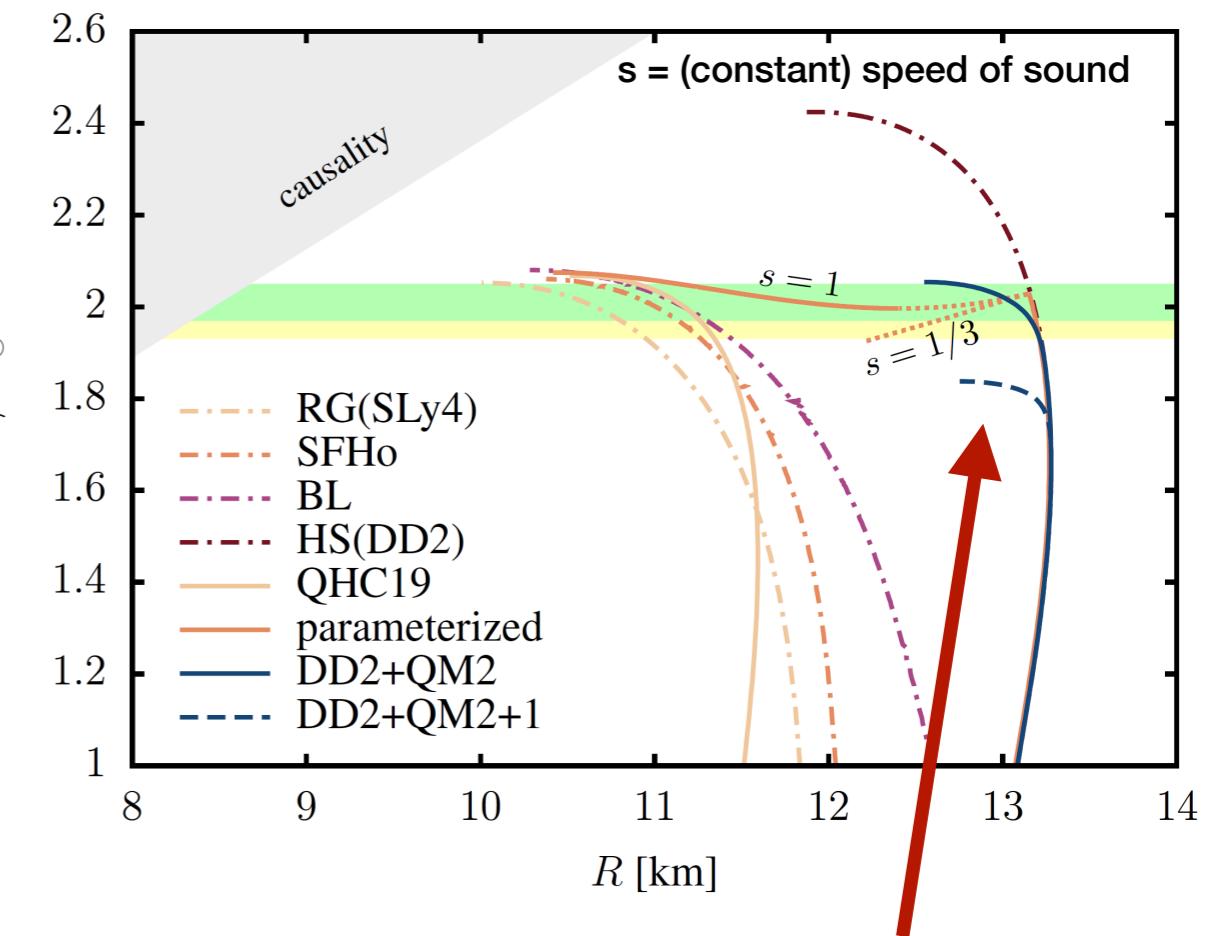
# Hybrid star construction possible? - yes

[Otto, Oertel, BJS 2020]

combine nuclear EoS (DD2) with FRG QM truncation



→ continuous nuclear-hybrid branch



2  $M_\odot$  limit violated for  $N_f = 2+1$

can a repulsive vector interaction remedy this behavior?

# vector mesons & the FRG EoS

[Otto, Oertel, BJS 2020]

[Rennecke 2015]  
[Pereira, Stiele, Costa 2020]

- Yukawa type interaction of temporal component and mean-field potential

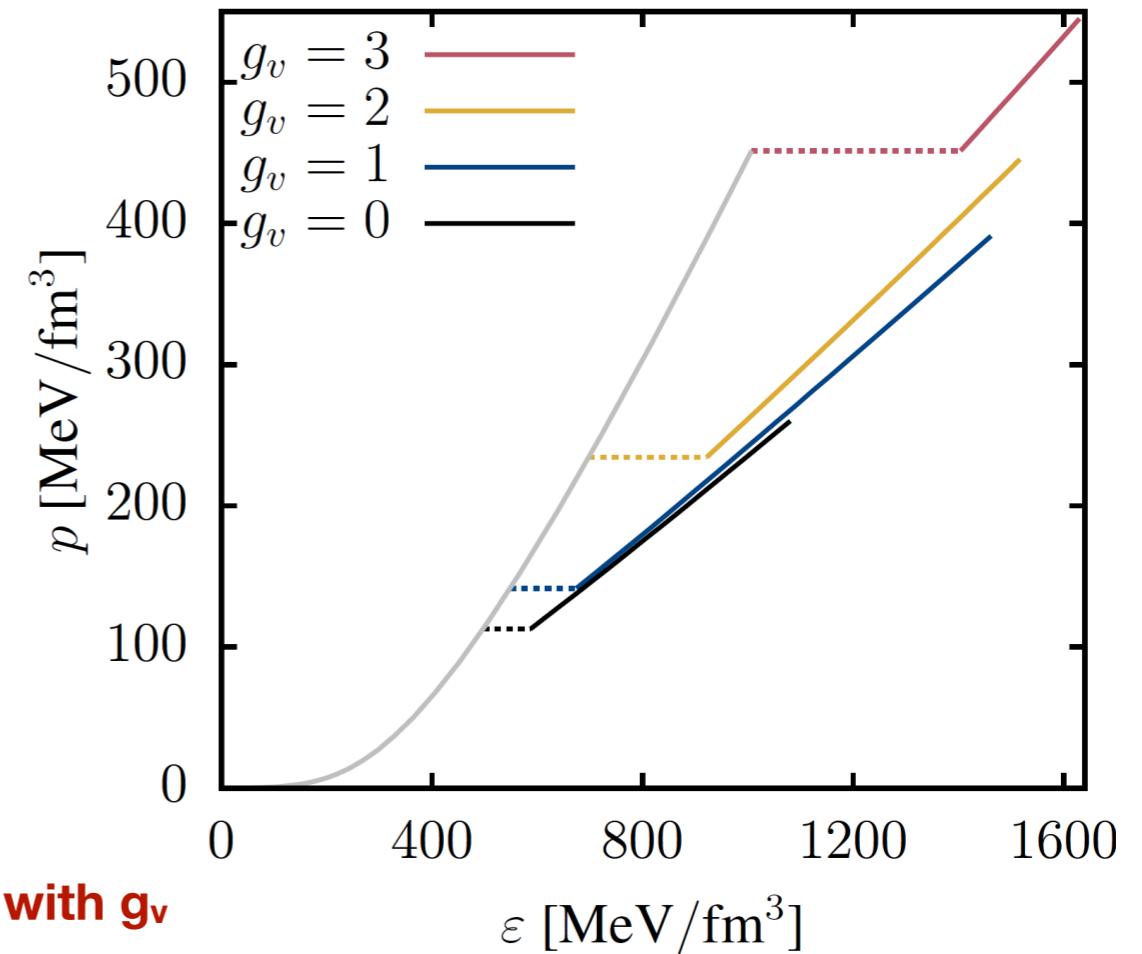
$$\Gamma_{\text{vec}} = \int_x \left[ \frac{g_v}{2} \bar{q} \gamma_0 \text{diag}_f(\omega, \omega, \sqrt{2}\phi) q - \frac{1}{2} (m_\omega^2 \omega^2 + m_\phi^2 \phi^2) \right]$$

- effectively shifts the chemical potentials:

$$\tilde{\mu}_u = \mu_q - \frac{2}{3} \mu_e - \frac{g_v}{2} \omega$$

$$\tilde{\mu}_d = \mu_q + \frac{1}{3} \mu_e - \frac{g_v}{2} \omega$$

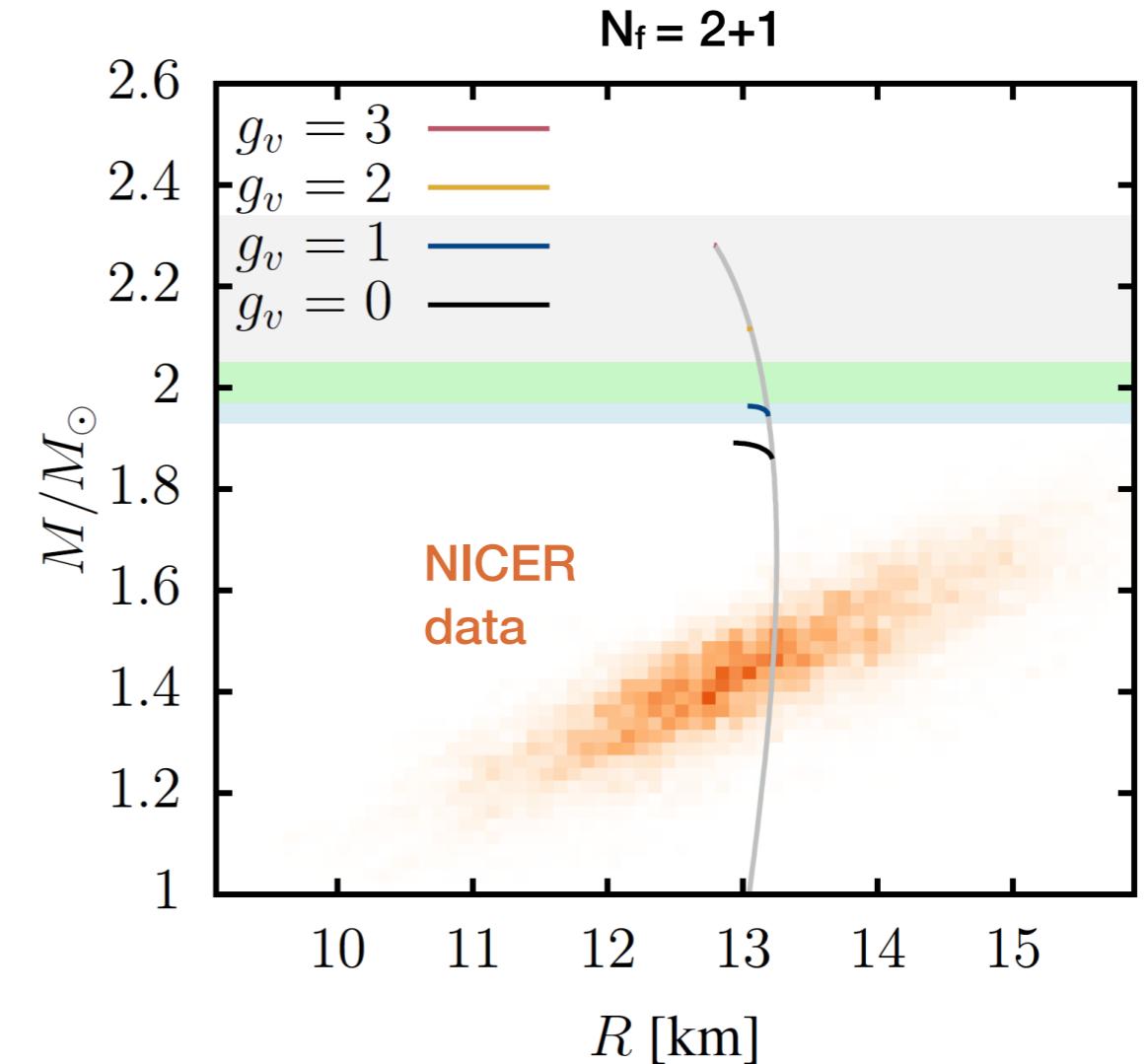
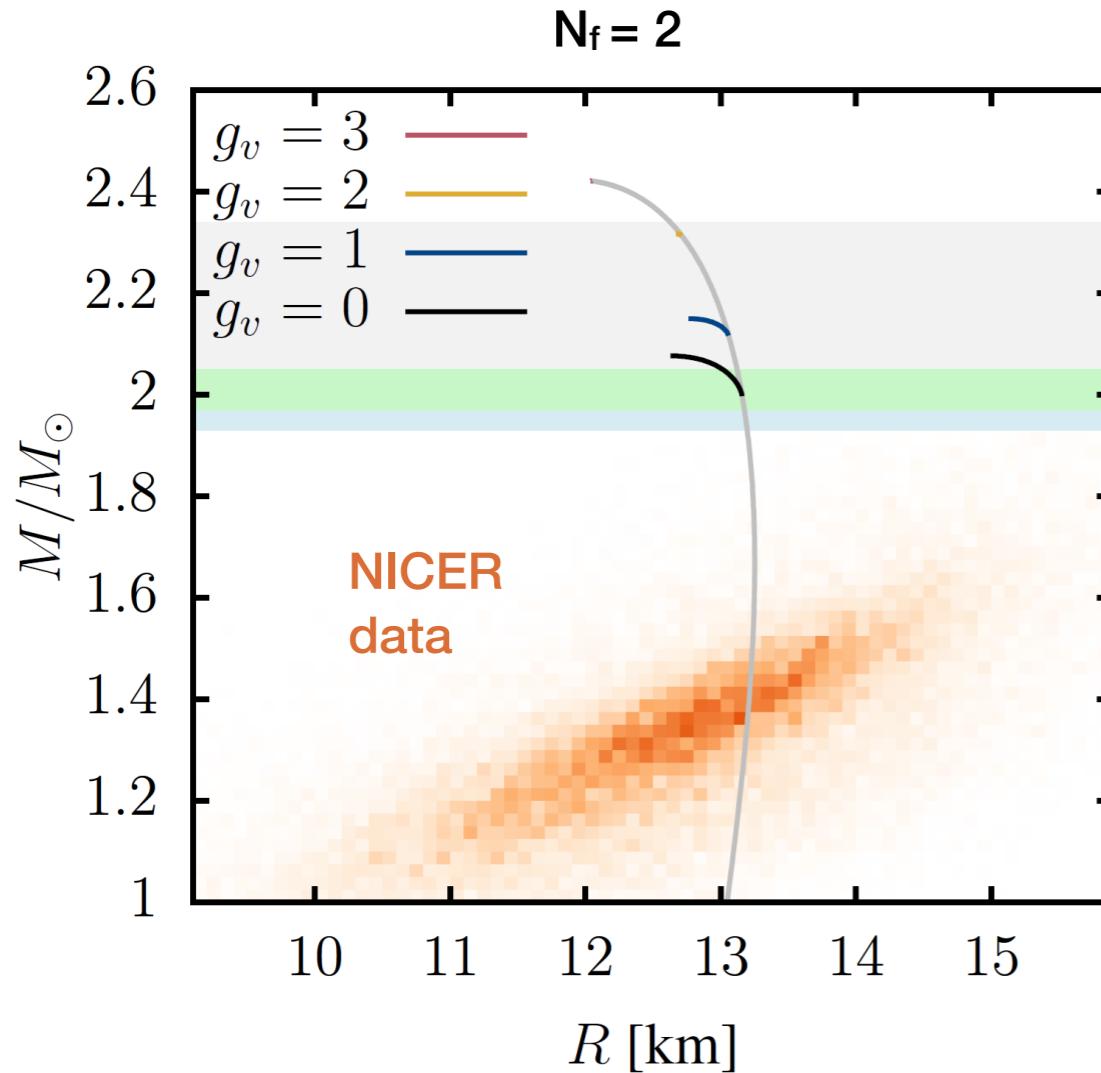
$$\tilde{\mu}_s = \mu_q + \frac{1}{3} \mu_e - \frac{g_v}{\sqrt{2}} \phi$$



→ energy gap and transition pressure increases with  $g_v$

# Mass-Radius relations

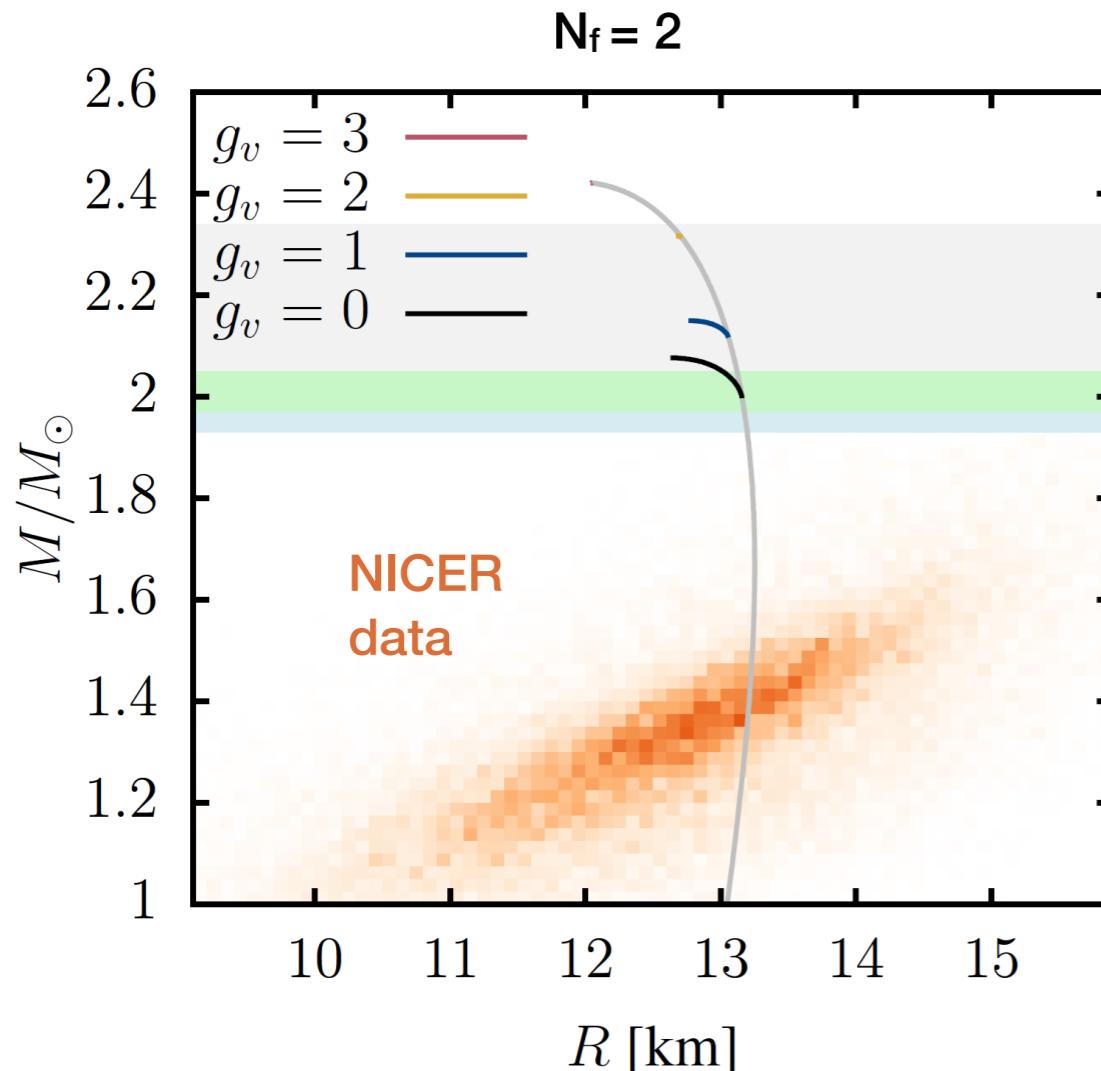
[Otto, Oertel, BJS 2020]



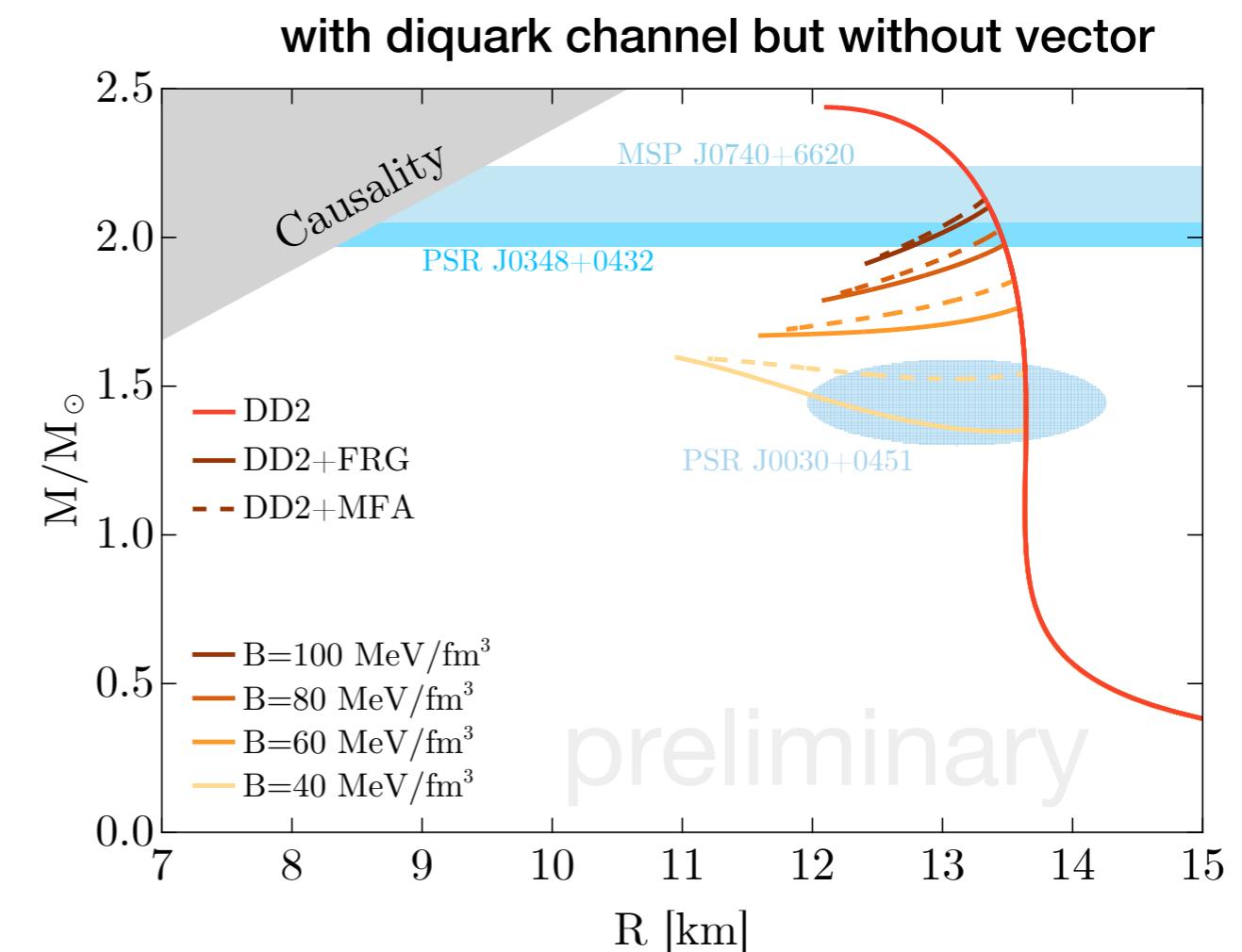
- including strange quarks: **finite vector coupling is needed to achieve  $2M_\odot$  limit**
- at the same time: **larger vector coupling lead to smaller quark cores!**

# Mass-Radius relations

[Otto, Oertel, BJS 2020]



[Mire, BJS ... next talk]



- including strange quarks: **finite vector coupling is needed to achieve  $2M_\odot$  limit**
- at the same time: **larger vector coupling lead to smaller quark cores!**

so far so good ... BUT

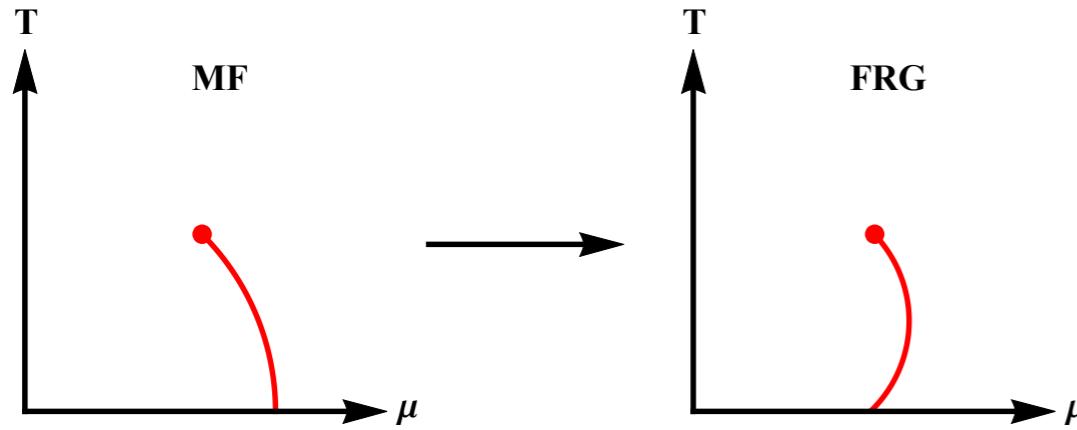
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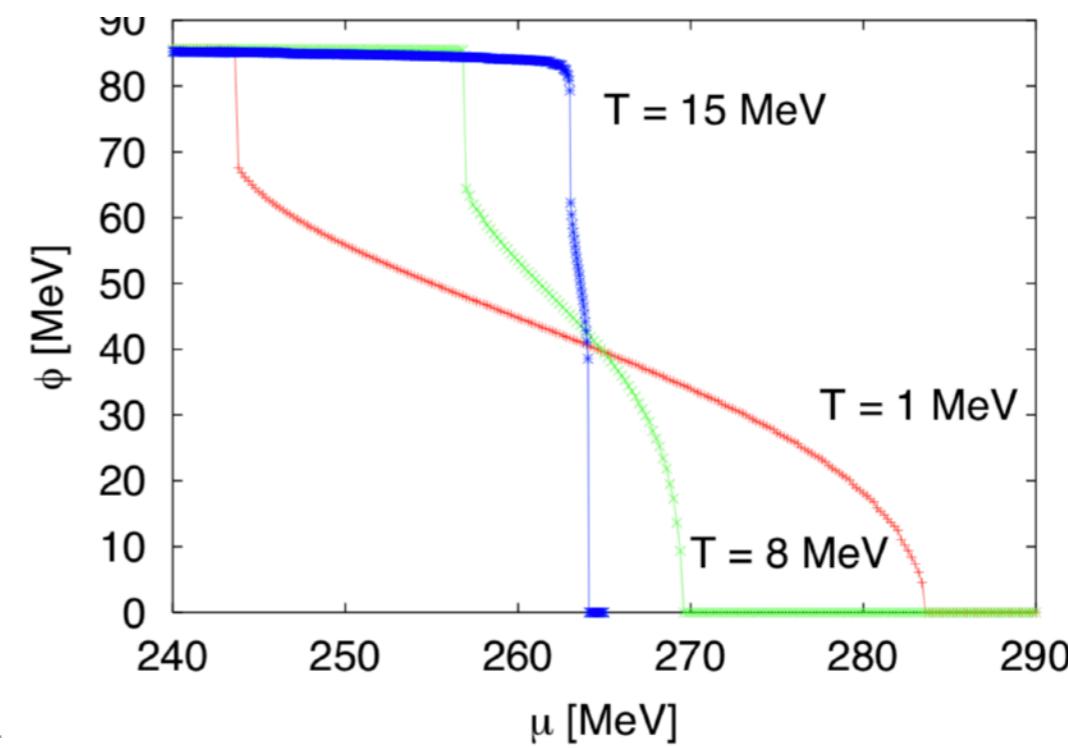
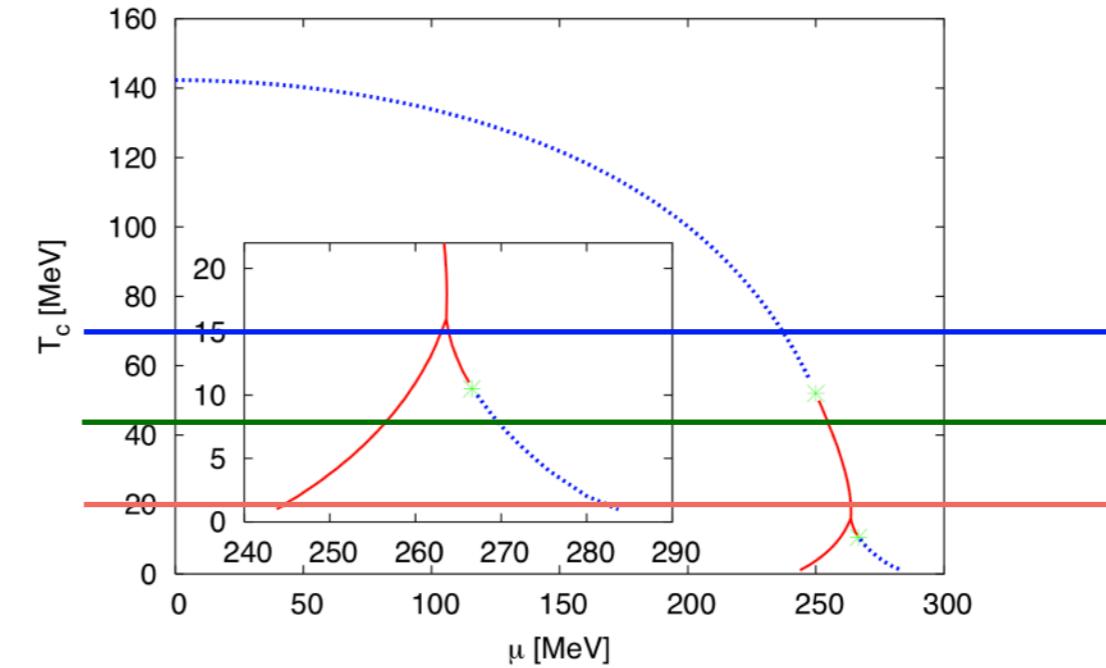
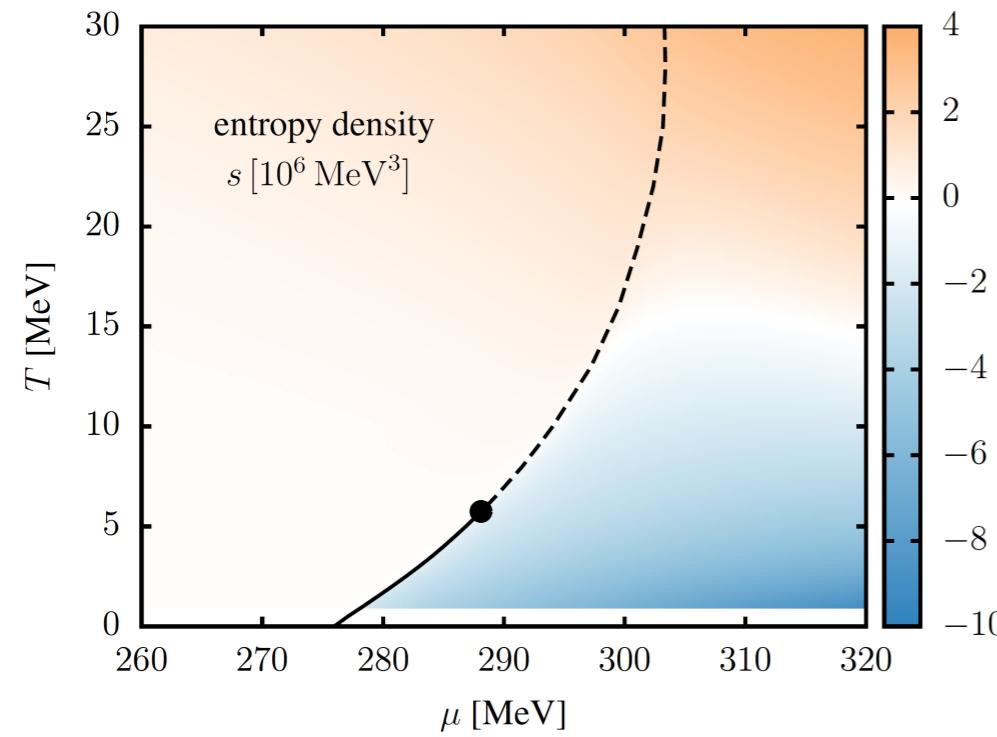
# back-bending / negative entropy density

[R-A Tripolt, BJS, L von Smekal, J Wambach 2018]

[BJS, Wambach 2005]

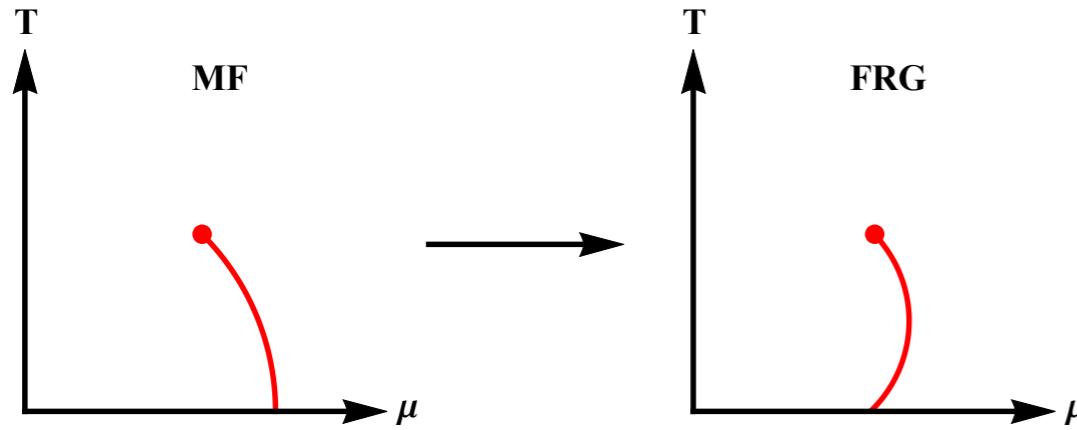


- Phase diagram quark-meson model
- Entropy density:  $s/T^3$

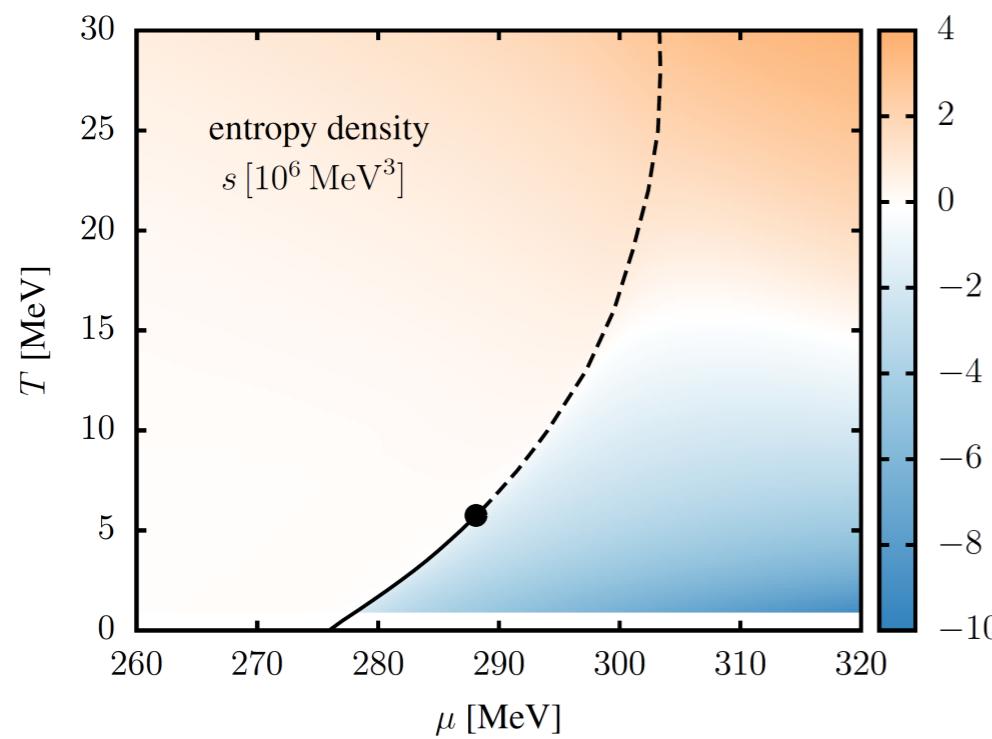


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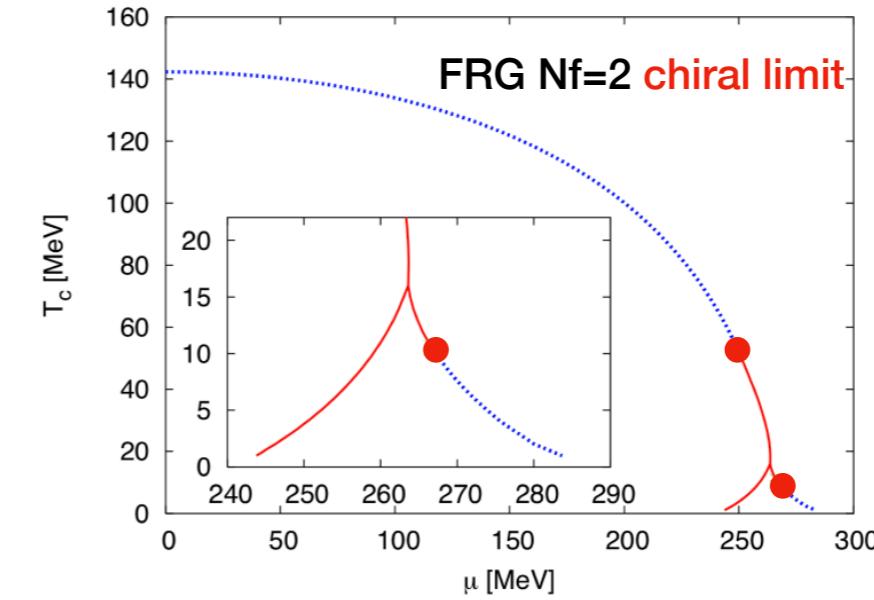
[R-A Tripolt, BJS, L von Smekal, J Wambach 2018]



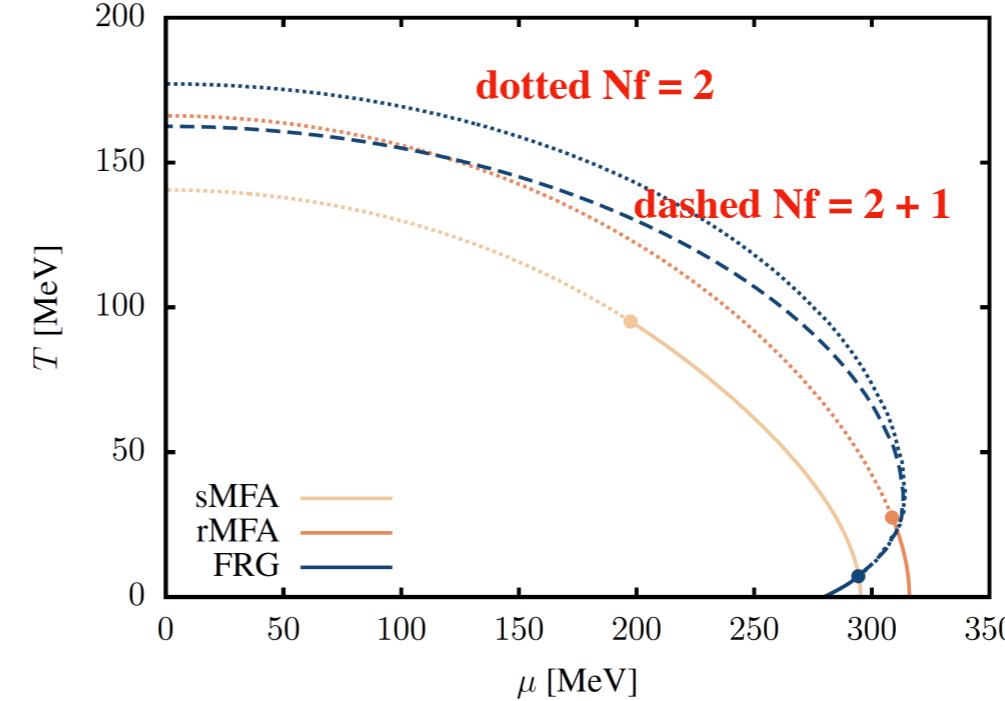
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[BJS, Wambach 2005]

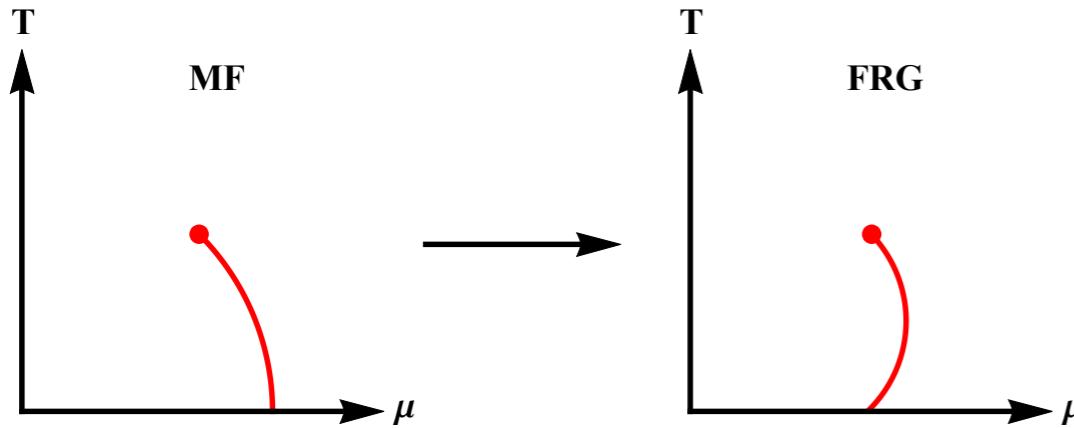


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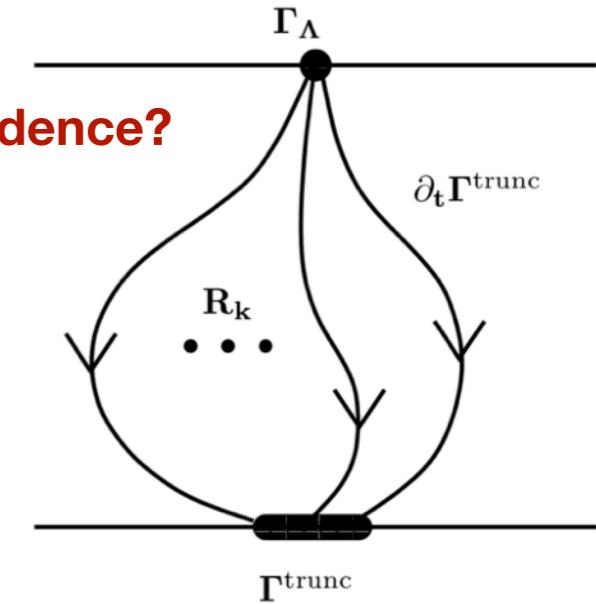
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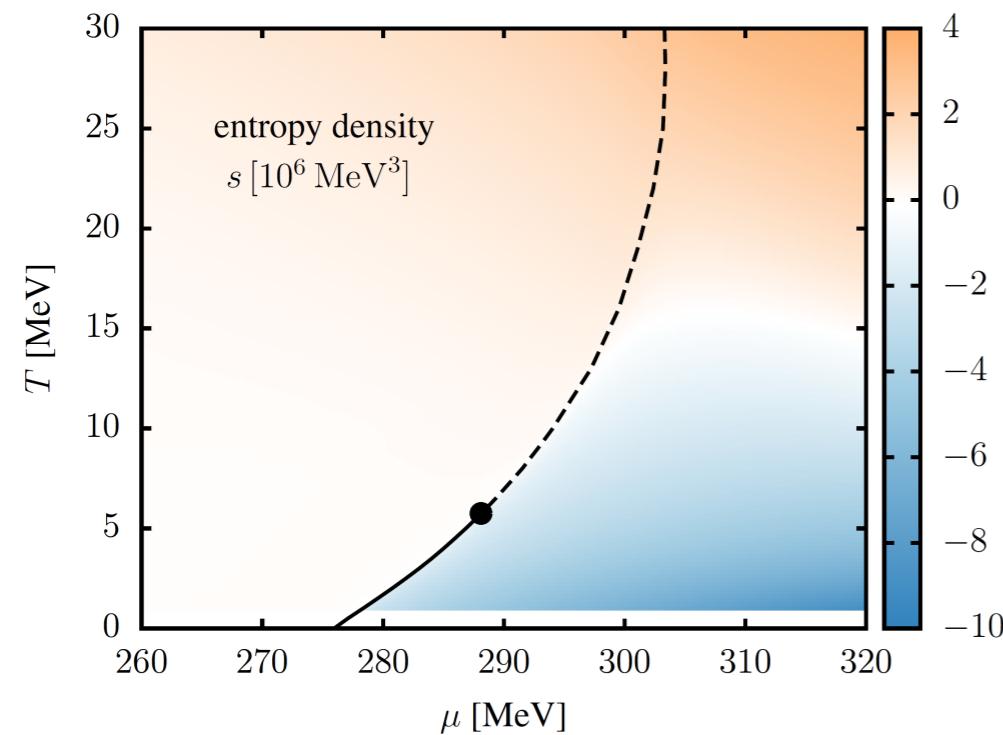


→ Regulator scheme dependence?

less pronounced when more channels are included  
e.g. pairing channel  
s. next talk by Ugo Mire



- Phase diagram quark-meson model
- Entropy density:  $s/T^3$



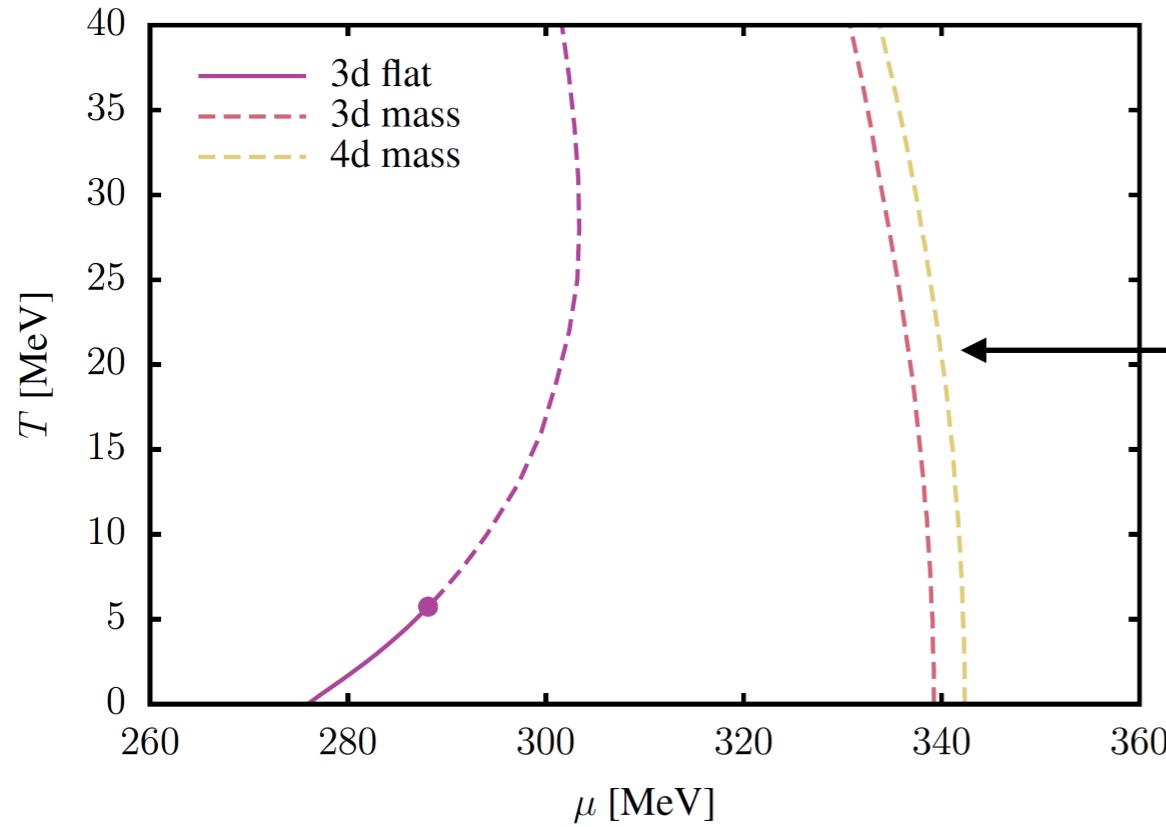
- fermionic regulator

$$R_k^F(p, \mu) = R_k^F(\tilde{p}, 0) \quad \tilde{p} = \begin{pmatrix} p_0 + i\mu \\ \vec{p} \end{pmatrix}$$

shift required to preserve Silver Blaze property ( $T=0$ )  
(necessary but not sufficient)

- relative shift in the cutoff-scales between bosonic & fermionic regulator  
→ is needed beyond LPA
- e.g. field-dependent Yukawa-coupling  
(multi quark-antiquark-meson scatterings)

# Chiral transition at low temperature

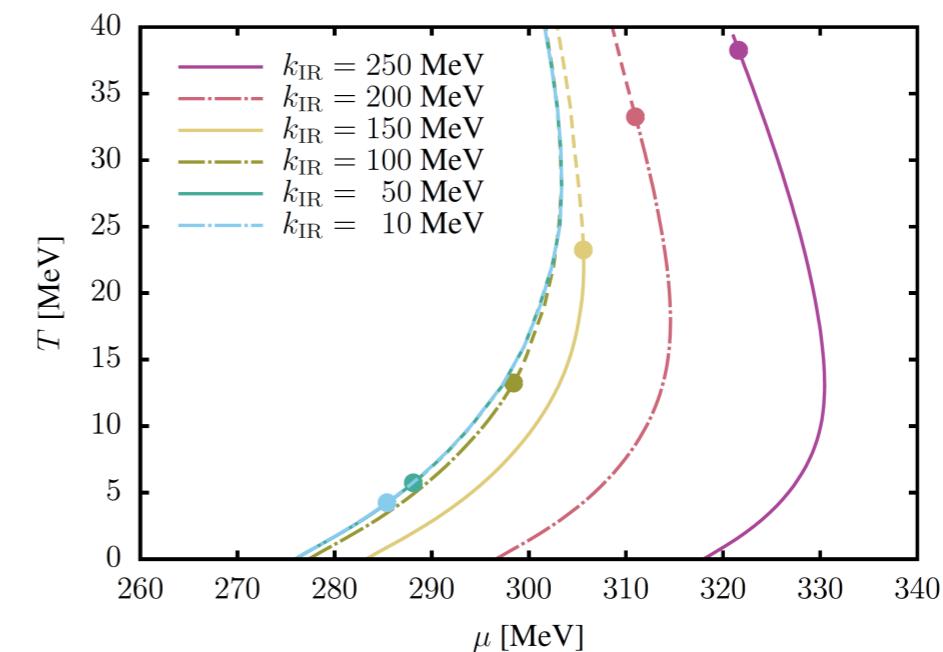


no back-bending with  
Callan-Symanzik mass-like regulator  
small differences between 3d and 4d regulators

purely crossover  $\rightarrow$  pseudocritical  $\mu_c$   
larger than  $m_q$

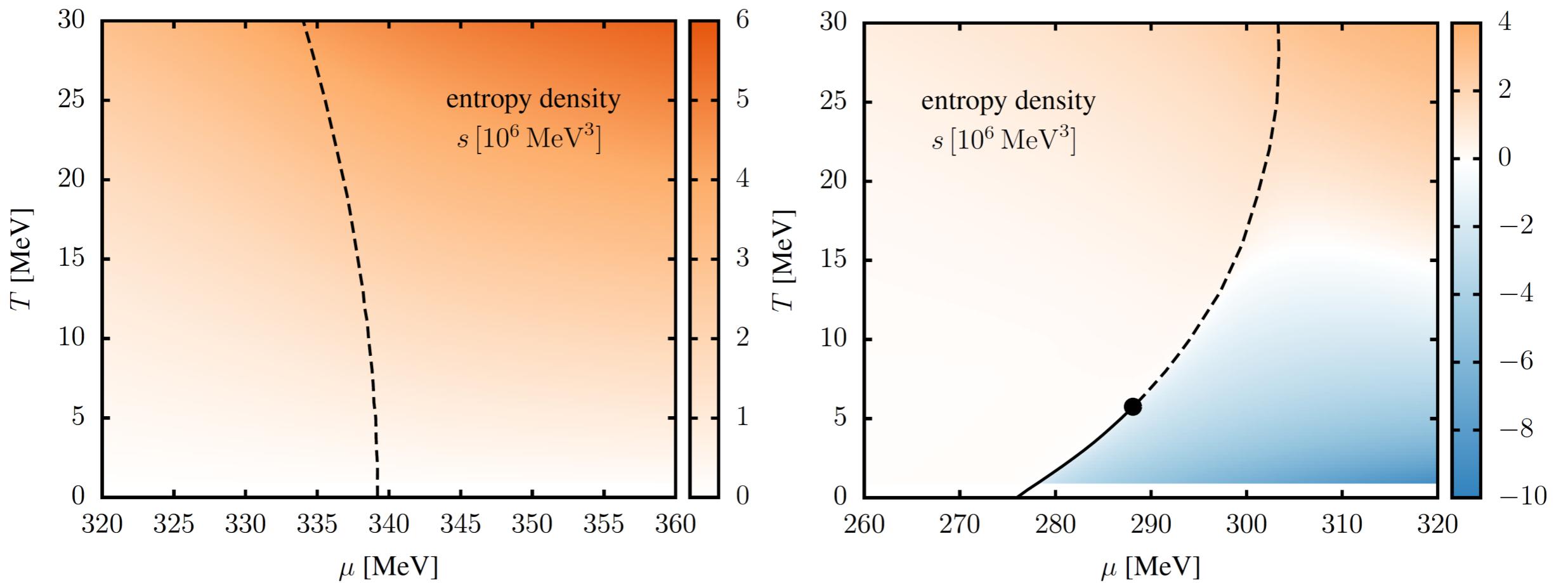
could finite IR cutoff play a role?  $\rightarrow$  No

transition line shifts and CEP moves down  
but back-bending over large  $k_{IR}$  range



# Chiral transition at low temperature

→ no negative entropy density anymore for Callan-Symanzik mass-like regulator



# EoS from QCD

- QCD procedure: start @ $O(100 \text{ GeV})$  (deep high-energy perturbative region)

[Braun et al. 2012++]

$\Lambda \sim O(10 \text{ GeV})$       quarks, gluons

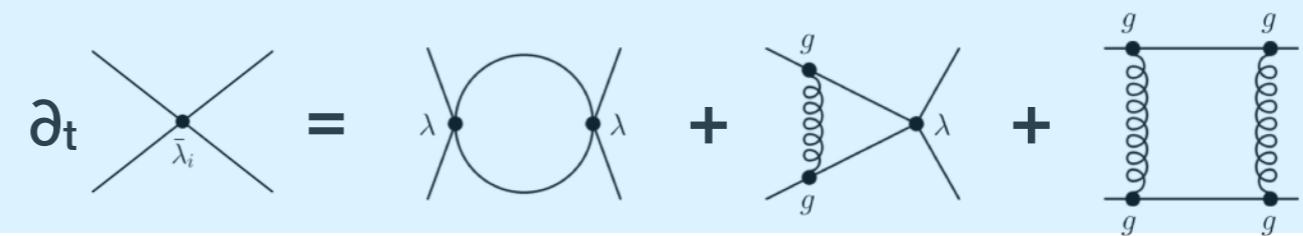
↓  
symmetric regime  
quarks, gluons → mesons

$k_\Phi$   
↓  
symmetry-broken regime  
quarks, diquarks, mesons etc

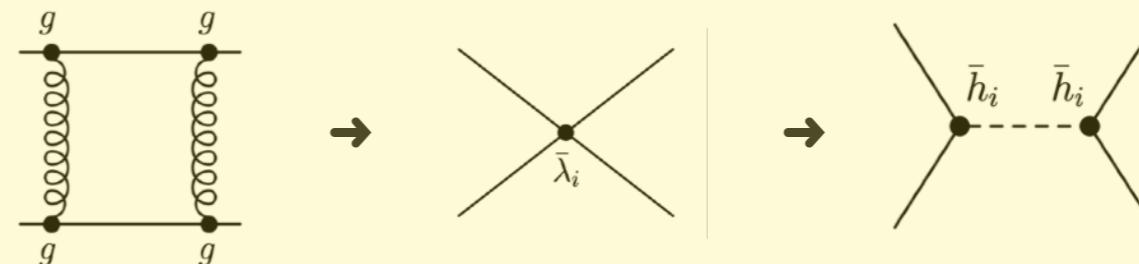
$k_{\text{IR}}$

$$S = \int d^4x \left\{ \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \bar{\psi} (i\partial^\mu + \bar{g} A^\mu + i\gamma_0 \mu) \psi \right\}$$

- quark-gluon vertex → many quark self-interaction channels
- dynamical hadronisation:  
4-quark correlators → bound states /resonances



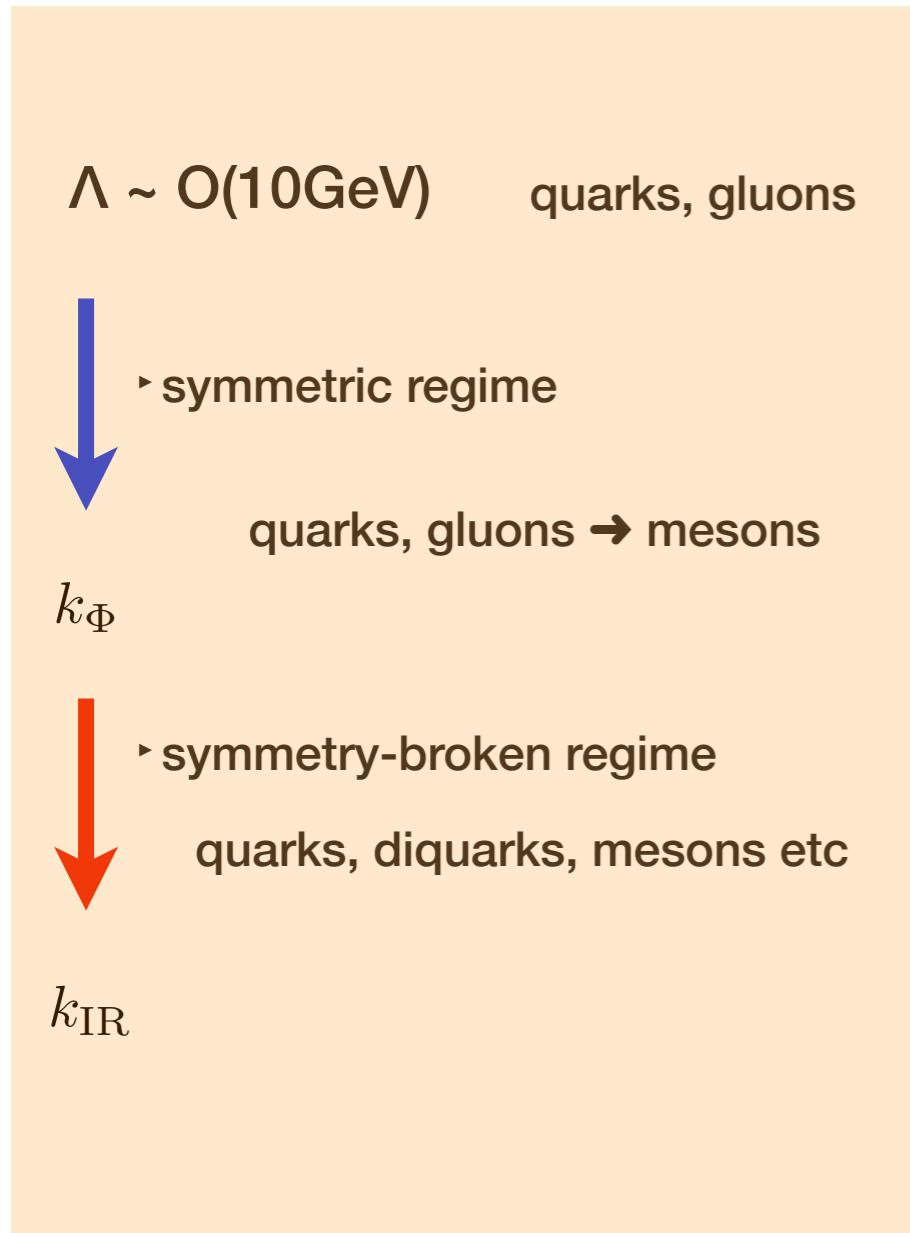
- general picture:



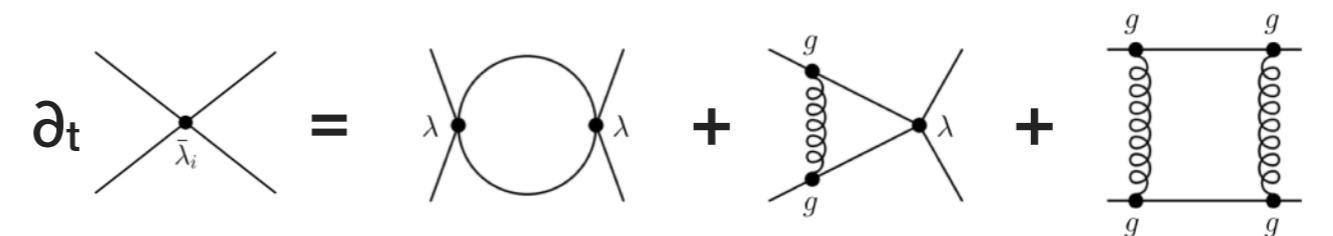
# EoS from QCD

- QCD procedure: start @ $\mathcal{O}(100 \text{ GeV})$  (deep high-energy perturbative region)

[Braun et al. 2012++]



$$S = \int d^4x \left\{ \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \bar{\psi} (i\partial^\mu + \bar{g} A^\mu + i\gamma_0 \mu) \psi \right\}$$



- symmetry breaking → condensates

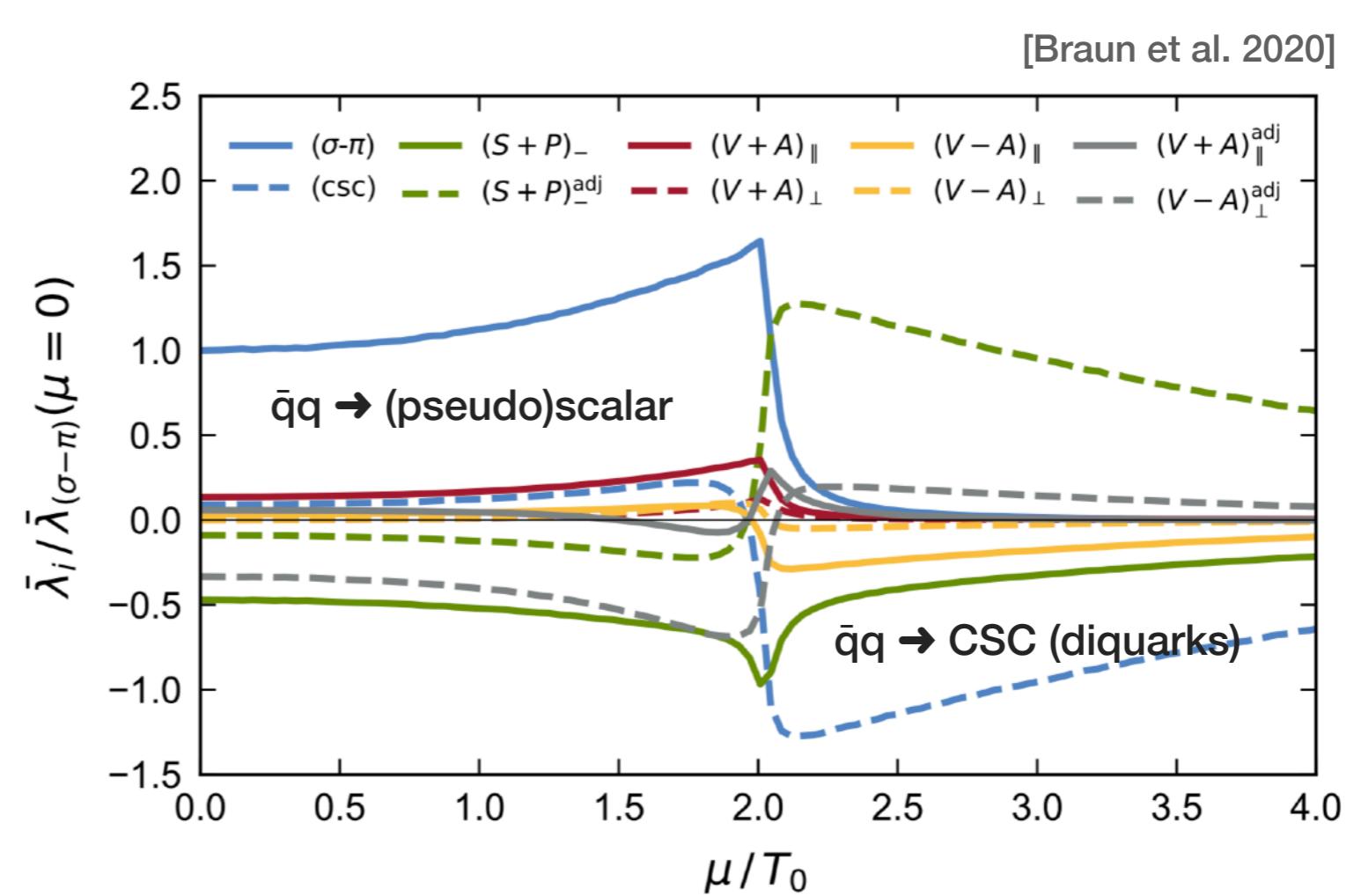
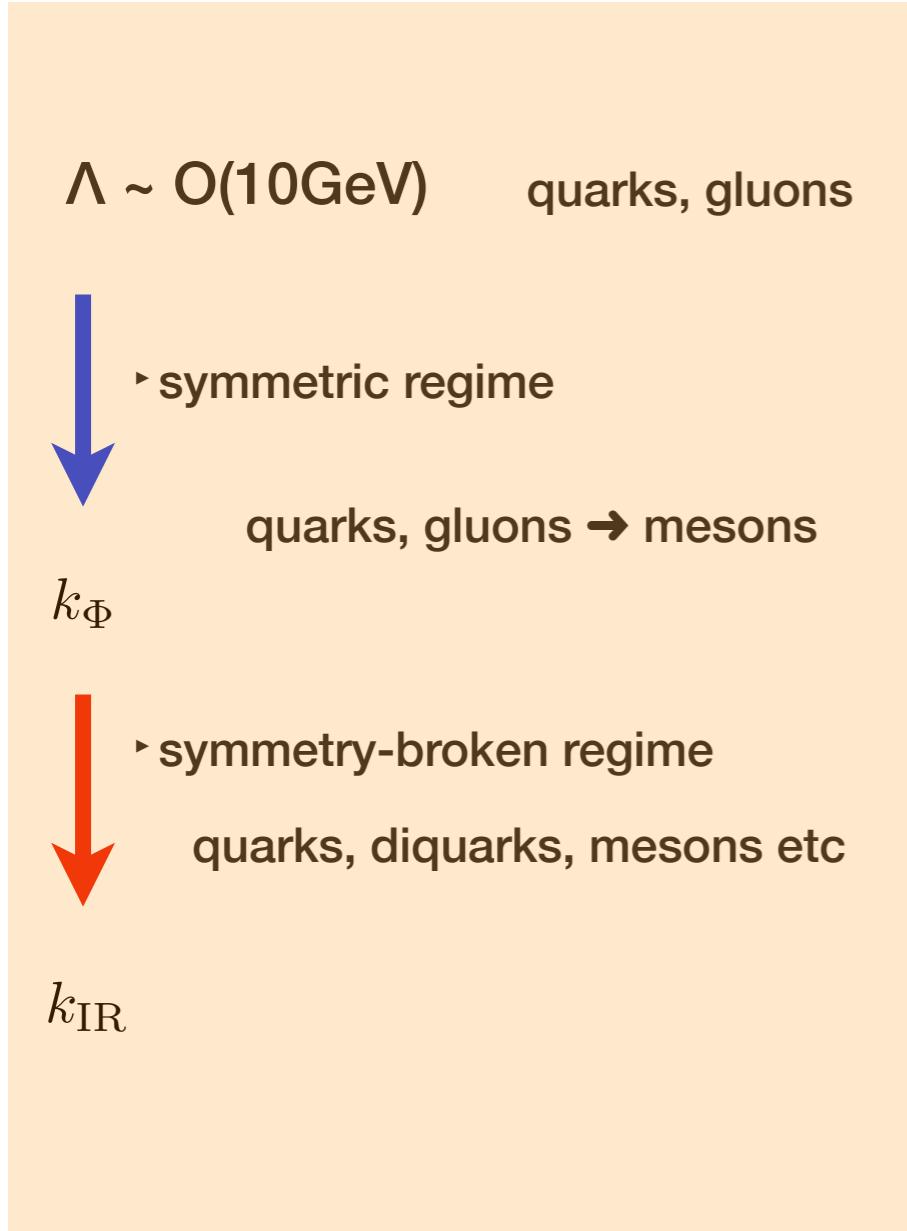
→ onset Landau-pole-type behavior  
 $\lambda \sim 1/m$

→ Ginzburg-Landau effective potential

Quark-meson-diquark truncation

# EoS from QCD

- QCD procedure: start @ $O(100 \text{ GeV})$  (deep high-energy perturbative region) [Braun et al. 2012++]



# Summary

Three home messages:

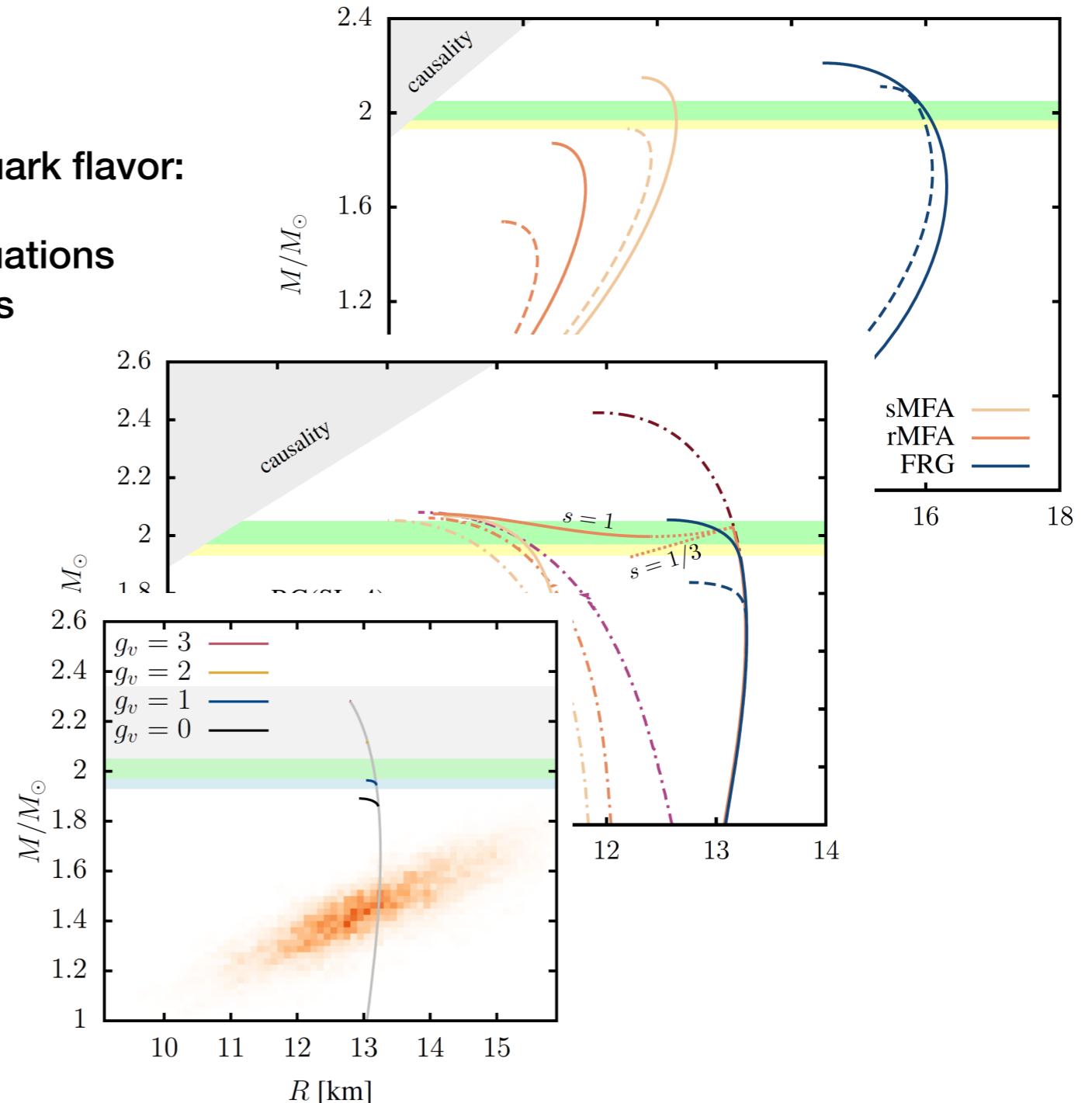
1. EoS with the FRG for two and three quark flavor:

→ significant impact of fluctuations  
on M-R relation for NSs

2. hybrid stars are possible

Non-zero vector coupling needed  
→ to reach  $2 M_{\odot}$  with strangeness

similar findings with pairing channels  
see next talk by Ugo Mire



# Summary

3. in LPA no back-bending / negative entropy density  
for  
CS mass-like regulators

CS type regulators closer to  
poles compared to flat regulator  
→ (vacuum) flows numerically harder

