

Strange nuclear matter in core-collapse supernovae

I. Sagert

Michigan State University, East Lansing, Michigan, USA

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NSCL

Phase diagram of strongly interacting matter

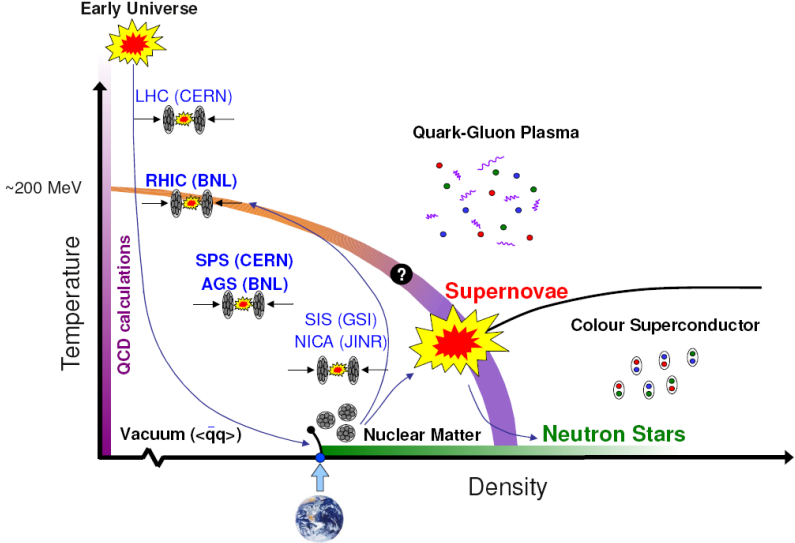
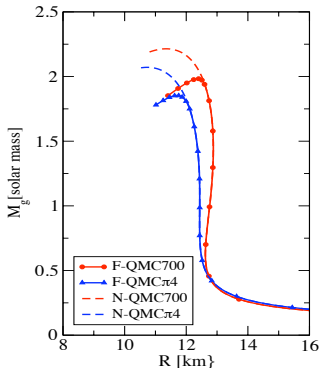
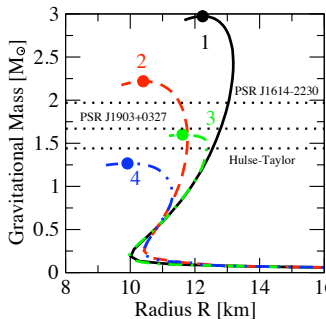


Fig.: Fredrik Sandin

Hyperons in neutron stars

- Large uncertainties in YY interaction, TBF with hyperons
- Microscopic calculations: Inclusion of hyperons softens the EoS (BHF, Schulze & Rijken PRC 84 (2011))
- Low maximum masses of hyperon stars ($< 1.4 M_{\odot}$)

- Quark meson-coupling model (Rikovska-Stone et al., NPA, Vol 792 (2007); Whittenbury et al., arXiv:1204.2614)
- SU(3) non-linear sigma model (Dexheimer and Schramm, PRC, Vol. 18 (2010))
- Relativistic Mean Field (Bednarek and Manka, J. Phys. G, 36 (2009), Weissenborn et al. NPA, Vol. 881 (2012))

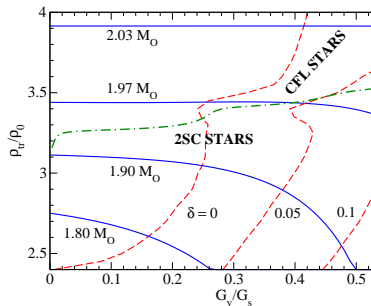
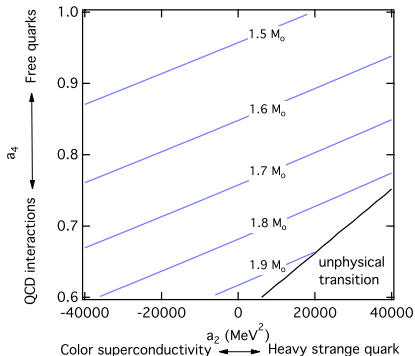


Figures: top: I. Vidana et al., EPL, Vol. 94, 2011, bottom: J. Rikovska Stone et al., NPA, Vol 792, 2007

Quark matter in neutron stars

- Bag models: Ozel et al., ApJL 724 (2010), Weissenborn et al., ApJL 740 (2010)
- Nambu-Jona-Lasinio models: Bonnano & Sedrakhian, Astron. & Astrophys. 539 (2012), Blaschke et al. Phys. Rev. C 80 (2009)
- $O(\alpha_s^2)$ Perturbative QCD: Kurkela et al. Phys. Rev. D 81 (2010)
- Dyson-Schwinger: Klaehn et al. (Phys. Rev C 82 (2010)), Chen et al. (Phys. Rev. D 84, (2011))

Figures: Ozel et al., ApJL 724 (2010), Bonnano & Sedrakhian, Astron. & Astrophys. 539 (2012)



Quark matter in neutron stars - EoS

- Hadronic model: RMF TM1
($M \sim 2.2 M_{\odot}$)

- Quark model: Bag model

$$p = -\Omega_{QM} = -\left(\sum_i \Omega_i + \frac{3\mu^4}{4\pi^2}(1 - a_4) + B_{eff}\right)$$

$$a_4 = \left(1 - \frac{2\alpha_s}{\pi}\right)$$

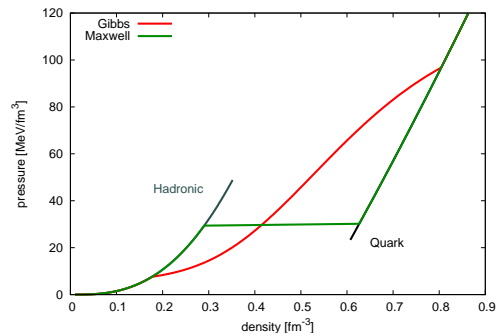
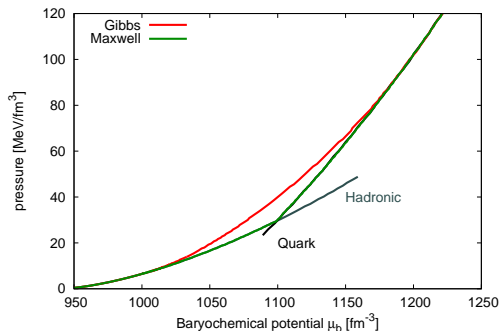
- $T^q = T^h, p^q = p^h, \mu_i^q = \mu_i^h$
 $\mu_s = \mu_d$

- Local charge neutrality: Maxwell transition

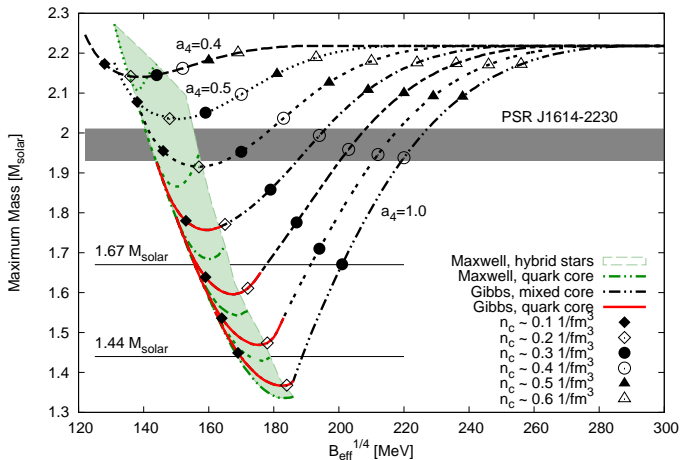
- Global charge neutrality: Gibbs transition

- Mixed phase with quark fraction

$$\chi = \frac{V_Q}{V_Q + V_H}, \quad 0 < \chi < 1$$

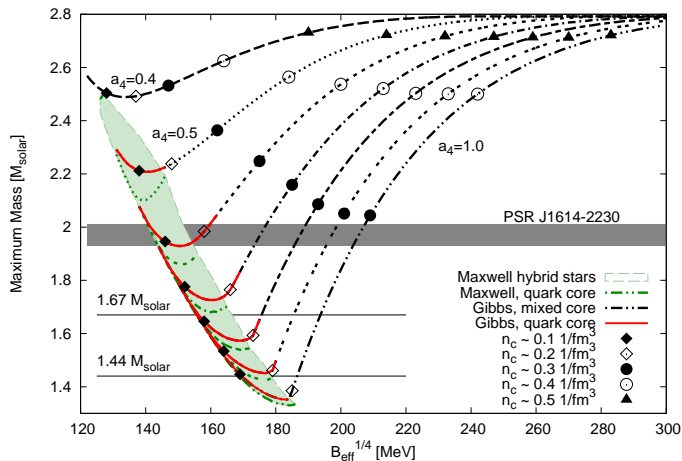


Hybrid star maximum masses with RMF TM1



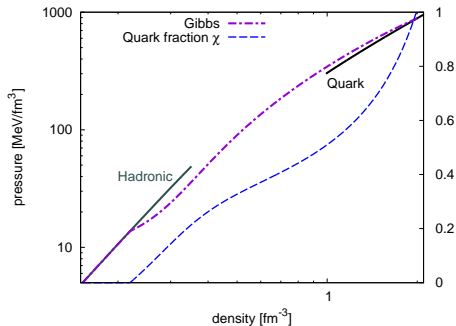
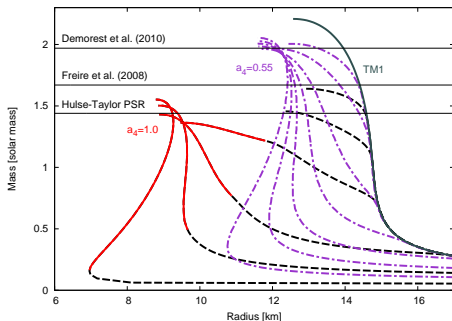
- For hadronic RMF TM1: Massive hybrid stars with low n_{crit} contain only a mixed phase

Hybrid star maximum masses with RMF NL3



- For hadronic RMF NL3: Massive hybrid stars with low n_{crit} can contain a pure quark matter core

Massive hybrid stars



- Stiff quark EoSs can be very similar to stiff nucleonic EoSs (Alford et al., ApJ 629 (2005))
- Not only the stars' masses but also their radii are similar
- How to distinguish hybrid stars from neutron stars?
- Cooling? Viscosity? Supernovae $\rightarrow p = f(n_b, Y_p, T)$!

Conditions in a core collapse supernova - $15 M_{\odot}$

- Typical conditions after core-bounce:

$$T \sim 10 \text{ MeV}$$

$$Y_p \lesssim 0.3$$

$$n_b \gtrsim n_0$$

- Typical supernova EoSs cover:

$$T : (0 - \geq 100) \text{ MeV}$$

$$Y_p : 0.01 - \geq 0.5$$

$$n_b : (10^5 - \geq 10^{15}) \frac{\text{g}}{\text{cm}^3}$$

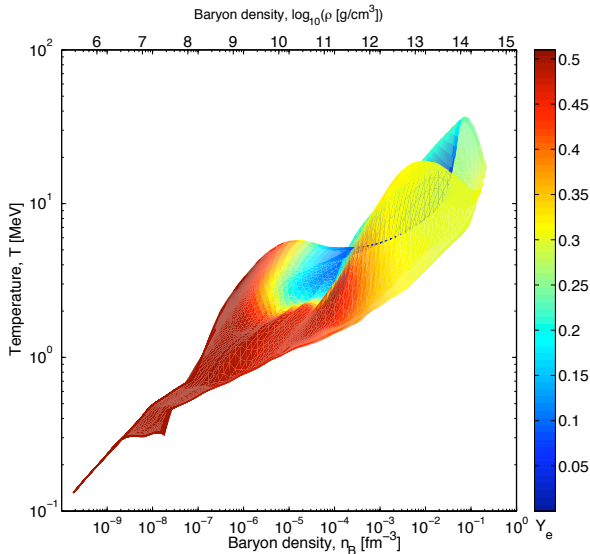


Figure: Fischer et al., ApJS 194, 39 (2011):
Phase space covered in a core collapse
simulation for a $15 M_{\odot}$ progenitor

Conditions in a core collapse supernova - $40 M_{\odot}$

- Typical conditions after

core-bounce:

$$T \sim 10 \text{ MeV}$$

$$Y_p \lesssim 0.3$$

$$n_b \gtrsim n_0$$

- Typical supernova EoSs

cover:

$$T : (0 - \geq 100) \text{ MeV}$$

$$Y_p : 0.01 - \geq 0.5$$

$$n_b : (10^5 - \geq 10^{15}) \frac{\text{g}}{\text{cm}^3}$$

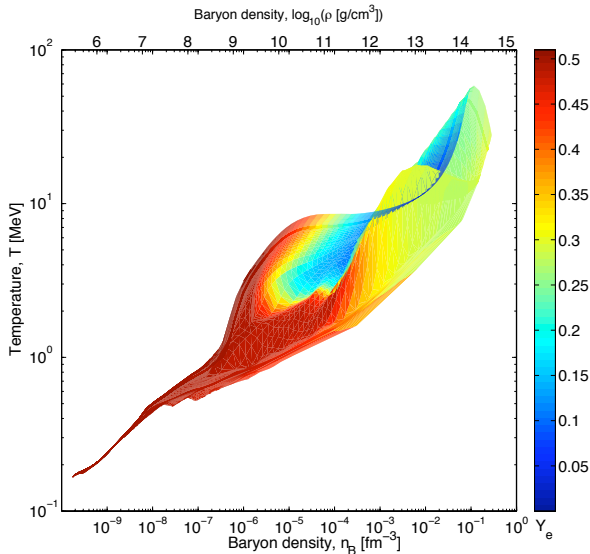


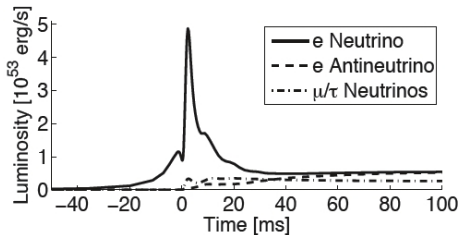
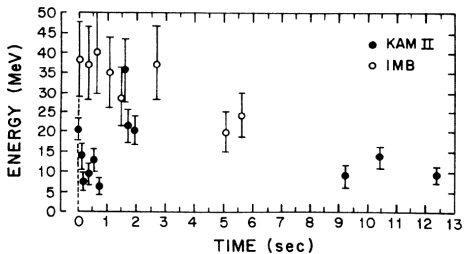
Figure: Fischer et al., ApJS 194, 39 (2011):
Phase space covered in a core collapse
simulation for a $40 M_{\odot}$ progenitor

Supernova observables

- Neutrino signal, gravitational waves, impact on nucleosynthesis, ...

SN1987A:

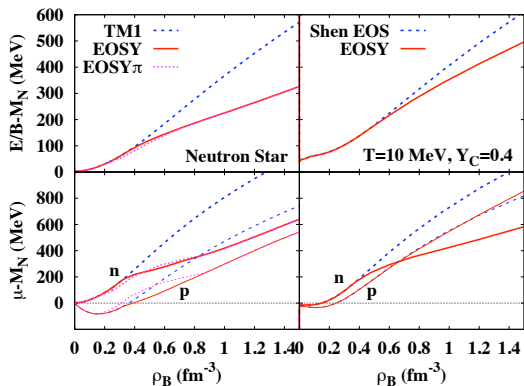
- Supernova explosion of $20M_{\odot}$ progenitor
- Detection of 24 neutrinos during ca. 13 s
- Estimated for emitted energy: $\sim 2 \cdot 10^{53}$ erg
- Next galactic supernova: IceCube and Superkamiokande will observe $\sim 10^3$ neutrinos



Quark and hyperon matter in core collapse supernovae

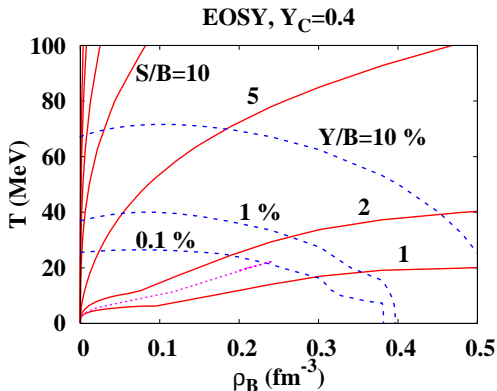
- Migdal, Chernoustan, Mishustin, Phys. Lett. B 83 (1979)
- Takahara and Sato, Astrophys. and Space Science 119 (1986)
- Drago and Tambini, Journal of Phys. G 25 (1999)
- Gentile et al., Astrophys. Journal 414 (1993)
- Pons et al., ApJ 513 (1999), Pons et al., Phys.Rev.Lett. 86 (2001)
- Nakazato et al., Phys. Rev. D 77 (2008)
- Sumiyoshi et al., ApJL, 690 (2009)
- I.S. Fischer et al., Phys. Rev. Lett. 102 (2009)
- ...
- Ishizuka et al., Journal of Phys. G 35 (2008)
- Shen et al., Astrophys. Journal Suppl. 197 (2011)
- Oertel, Fantina, and Novak, Phys. Rev. C 85 (2012)

Hyperons in core-collapse of light progenitor stars



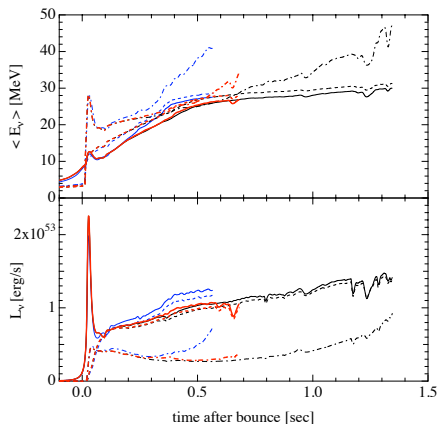
- Shen et al. supernova EoS (RMF TM1, $M_{max} \sim 2.1M_{\odot}$) extended to hyperons and thermal pions ($M_{max} \sim 1.6M_{\odot}$)
- $(U_{\Lambda}^{(N)}, U_{\Sigma}^{(N)}, U_{\Xi}^{(N)}) = (-30\text{MeV}, +30\text{MeV}, -15\text{MeV})$

Hyperons in core-collapse of light progenitor stars



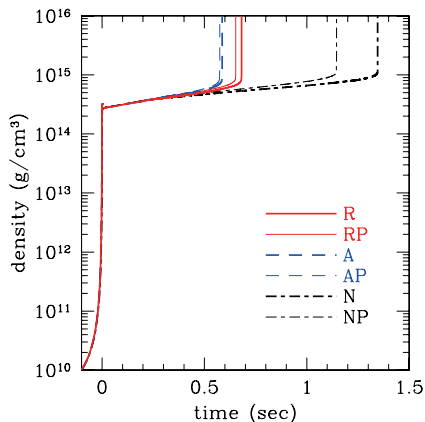
- Simulation an adiabatic collapse of an iron core from $15 M_{\odot}$ progenitor
- No neutrino transport
- Small hyperon fraction $\sim 0.1\%$ has no effect on the supernova dynamics

Hyperons in core-collapse SNe of massive progenitor stars



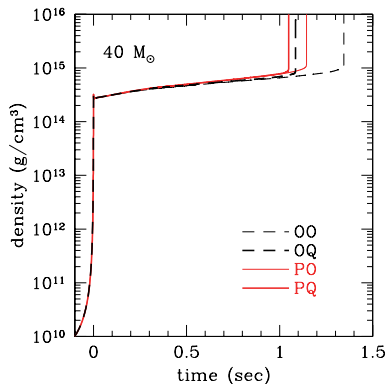
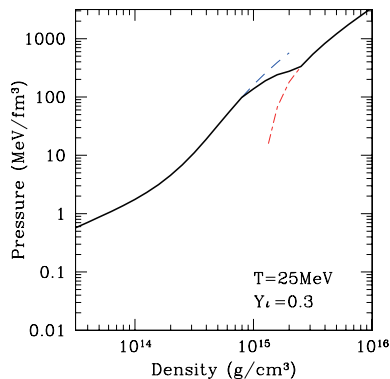
- Core-collapse supernova simulation of $40M_{\odot}$ progenitor with **hyperon EoS** from Ishizuka et al., JPG. 35 (2008)
- Comparison to H. Shen EoS and **Lattimer-Swesty EoS** ($K_0 = 180$ MeV)
- Hyperons appear around 500ms after core bounce and accelerate black hole production

Hyperons in core-collapse SNe of massive progenitor stars



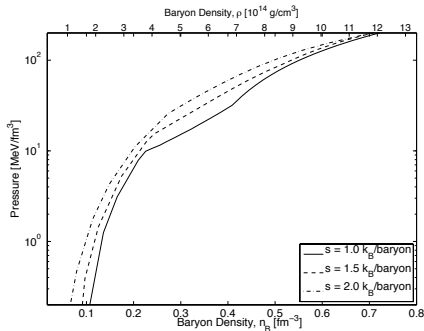
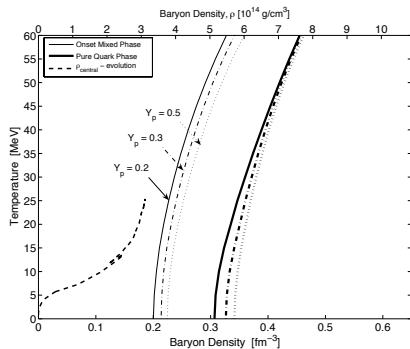
- Variation of stiffness in hyperon EoS via Σ hyperon potential
- $A = U_{\Sigma}^{(N)} = -30\text{MeV}$; $R = U_{\Sigma}^{(N)} = +30\text{MeV}$
- Softening due to appearance of hyperons ($M_{\text{max}} \sim 1.6M_{\odot}$) \rightarrow earlier black hole formation

Quark matter in core-collapse SNe of massive progenitor stars



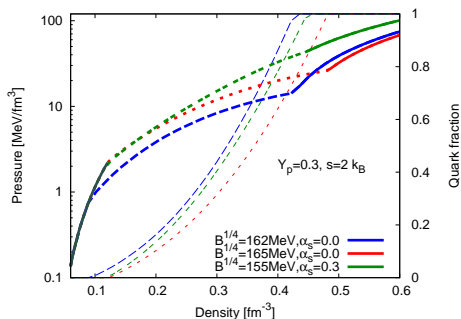
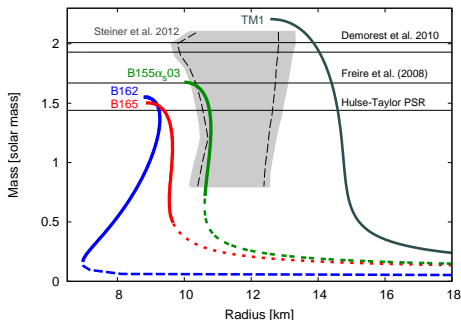
- Shen et al. EoS with strange quark matter and thermal pions
- Quark bag model: $B^{1/4} \sim 209$ MeV ($M_{max} \sim 1.8 M_{\odot}$)
- Core-collapse supernova simulation of 40M_⊙ progenitor
- Black hole formation ~ 0.07 ms after onset of phase transition

Quark matter in core-collapse SNe of light progenitors - Stiff EoS



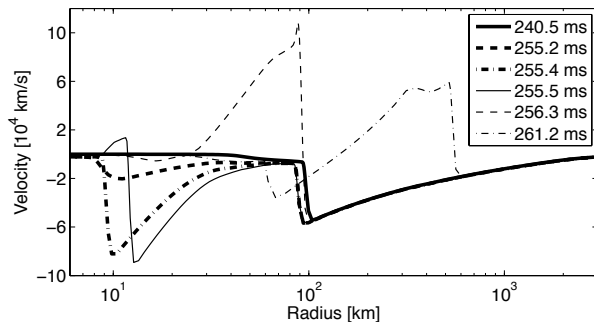
- Shen EoS & phase transition to quark matter with PNJL model ($M_{\text{max}} \sim 2 M_{\odot}$)
- Deconfinement to up and down quark matter, strangeness appears later
- 1D Supernova simulation of a $15 M_{\odot}$ progenitor
- GR hydrodynamics and Boltzmann neutrino transport in 1D (Liebendoerfer et al. 2004)
- Due to high critical density no phase transition during post-bounce accretion phase

Quark matter in core-collapse SNe of light progenitors - Soft EoS



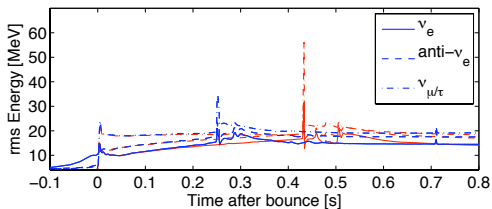
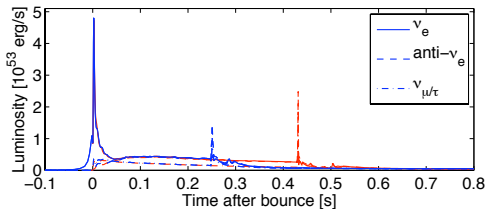
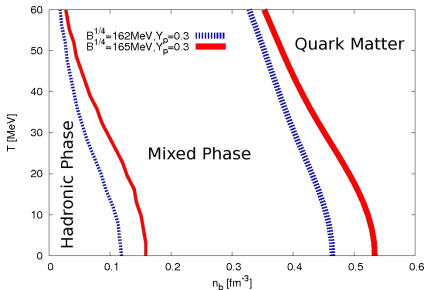
- Shen et al. EoS with quark bag model $\Omega_{QM} = \sum_{i=u,d,s,e} \Omega_i + \alpha_s \frac{\mu^4}{4\pi^3} + B$
- Progenitors: $10.8 M_{\odot}$, $13 M_{\odot}$, and $15 M_{\odot}$
- 1D GR hydrodynamics and Boltzmann neutrino transport (Liebendoerfer et al. 2004)
- $B^{1/4} = 162\text{MeV}$, 165MeV and $B^{1/4} = 155\text{MeV}$ & $\alpha_s = 0.3$

Second Shock Wave from Phase Transition



- Mixed phase is present after core-bounce
- Collapse of the proto neutron star due to transition to pure quark matter 200ms - 400ms after core bounce
- Formation of second shock wave which leads to the explosion of the star

First and Second Neutrino Bursts



- Second shock wave passes neutrinospheres \rightarrow second neutrino burst dominated by antineutrinos
- For $B^{1/4} = 165 \text{ MeV}$ second neutrino burst is ~ 200 ms later than for $B^{1/4} = 162 \text{ MeV}$

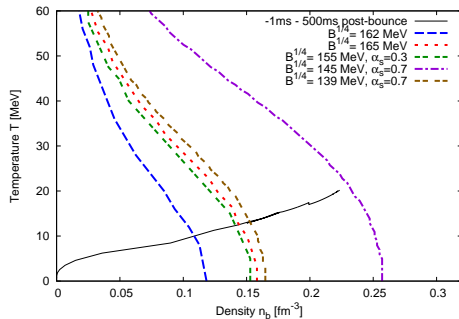
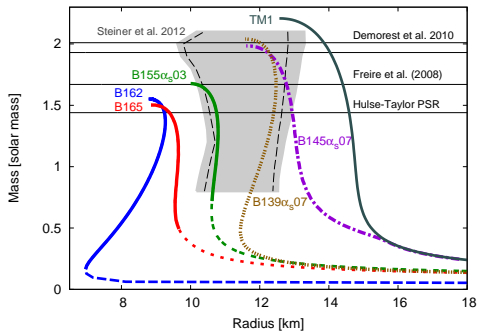
Fig: T.Fischer, Neutrino luminosities and rms neutrino energies, at 500km for $10 M_{\odot}$ progenitor

Parameter study (Fischer et al., ApJS 194, 39 (2011))

Prog. M_{\odot}	$B^{1/4}$ MeV	t_{pb} ms	ρ_c 10^{14}g/cm^3	T_c MeV	M_{pns} M_{\odot}	E_{expl} 10^{51}erg	M_{max} M_{\odot}
10.8	162	240	6.61	13.14	1.431	0.373	1.55
10.8	165	428	6.46	14.82	1.479	1.194	1.50
13	162	235	6.49	13.32	1.465	0.232	1.55
13	165	362	7.23	16.38	1.496	0.635	1.50
15	162	209	7.52	17.15	1.608	0.420	1.55
15	165	276	7.59	16.25	1.641	u	1.50
15	155, $\alpha_s = 0.3$	326	5.51	17.67	1.674	0.458	1.67

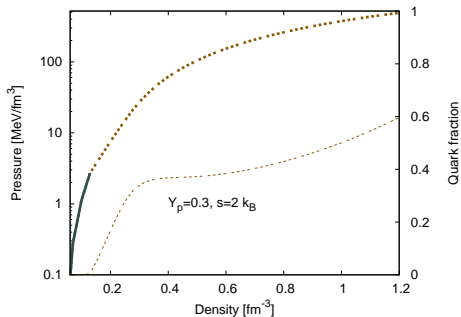
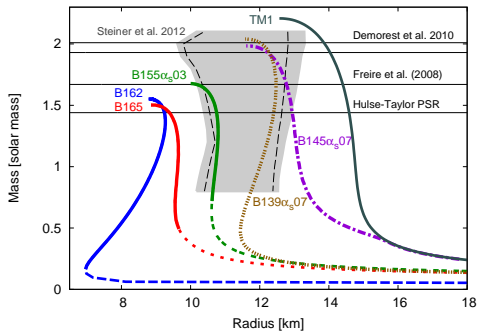
- Higher critical density:
 - More massive proto neutron star with deeper gravitational potential
 - Stronger second shock and larger explosion energies
 - Second neutrino burst later with larger peak luminosities
- More massive progenitor: earlier onset of phase transition and more massive proto neutron star

Quark matter in core-collapse SN of light progenitors - stiff EoS



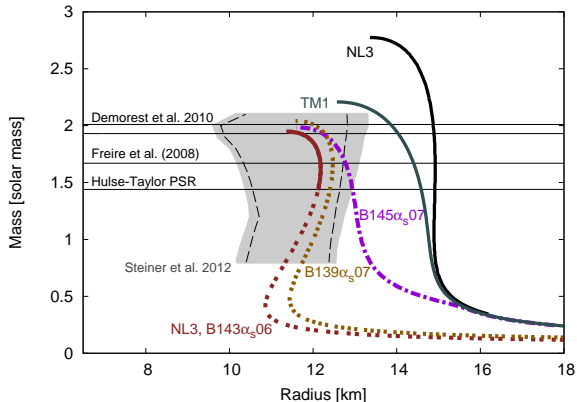
- For $B^{1/4} = 145$ MeV & Collapse of $15 M_{\odot}$ and $30 M_{\odot}$ progenitors
- Phase transition too late ~ 1 s after bounce \rightarrow No second collapse
- For $B^{1/4} = 139$ MeV: Earlier phase transition, but no second collapse

Quark matter in core-collapse SN of light progenitors - stiff EoS



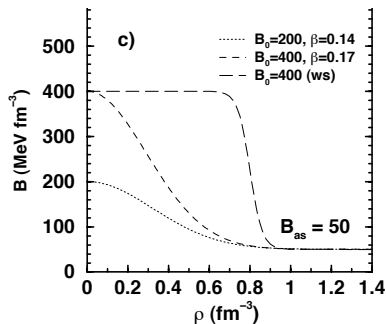
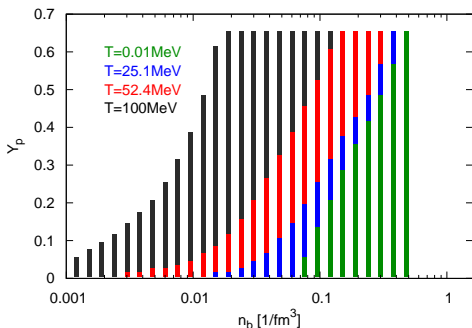
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Second collapse & stiff hybrid EoS (?)



- Choose stiffer hadronic EoS (NL3) \rightarrow Stronger phase transition and pure quark matter
- Caution 1: NL3 parameter set is too stiff with regard to heavy-ion data
- Caution 2: Mass-Radius relation for relevant density region look very similar

Second collapse & stiff hybrid EoS (?)



- Quark EoS that is more flexible at higher densities and can stiffen
- Density dependent Bag constant (Burgio et al. Phys.Lett. B526 (2002))

$$B(\rho) = B_{as} + (B_0 - B_{as}) \exp \left[-\beta \left(\frac{\rho}{\rho_0} \right)^2 \right]$$
- Soft mixed phase and stiff high density quark EoS

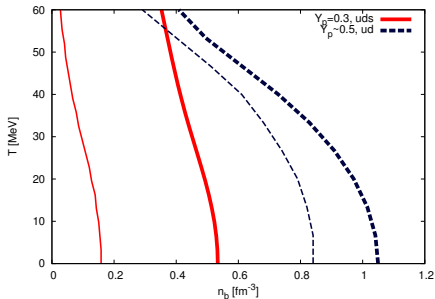
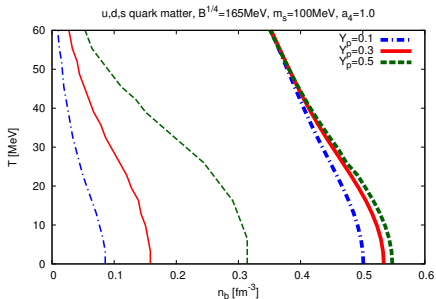
Summary & Conclusion

- Variety of studies for quark matter and hyperons in neutron stars and core-collapse supernovae
- Need more input on high density nuclear matter and hyperon interactions from heavy ion experiments
- Softening of the quark or hyperon EoS impacts the duration and spectrum of the supernova neutrino signal
- **But:** EoSs have to be stiff at high density! Only a few studies use EoSs which reproduce a neutron star mass of $M_{NS} = 1.97 \pm 0.04 M_{\odot}$
- Will we be able to distinguish between effects from (stiff) hyperonic, quark, or hadronic EoSs ?
- All supernova EoSs containing hyperons or quarks assume weak equilibrium with respect to strangeness production \leftrightarrow Is this a valid assumption? If not - how do we evolve strangeness in astrophysical simulations ?

With

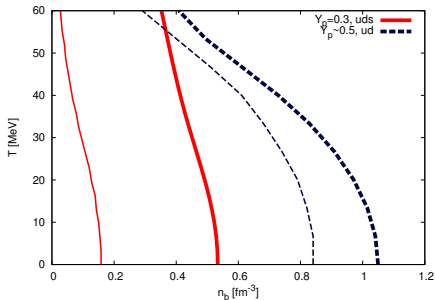
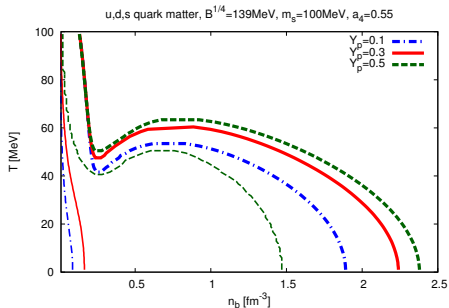
- T. Fischer
- M. Hempel
- M. Liebendoerfer
- A. Mezzacappa
- G. Pagliara
- J. Schaffner-Bielich
- F.-K. Thielemann

Low critical density - Neutron stars vs. heavy ion collision



	Neutron stars & supernovae	Heavy ion collisions
Proton Fraction Y_p	≤ 0.3	~ 0.5
Strangeness production	weak interaction: $d \leftrightarrow s$	$s\bar{s}$
Phase transition to	uds quark matter	ud quark matter

Low critical density - Neutron stars vs. heavy ion collision



	Neutron stars & supernovae	Heavy ion collisions
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Strangeness production	weak interaction: $d \leftrightarrow s$	$s\bar{s}$
Phase transition to	uds quark matter	ud quark matter

Core-collapse and second shock

- Mixed phase of quarks and hadrons appears at core bounce in the center

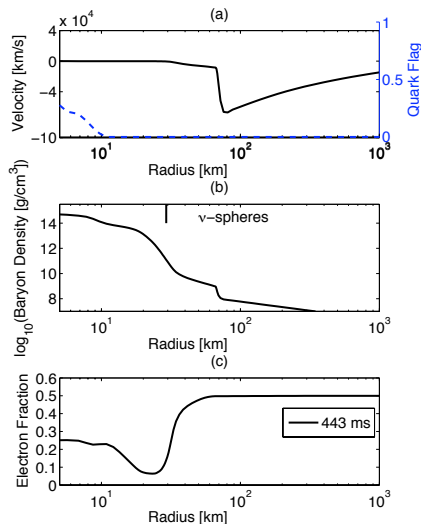
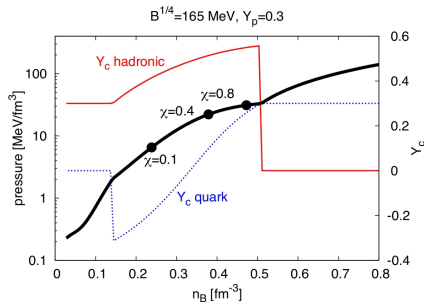


Fig: T.Fischer, $B^{1/4} = 165 \text{ MeV}$, $10M_{\odot}$ progenitor

Core-collapse and second shock

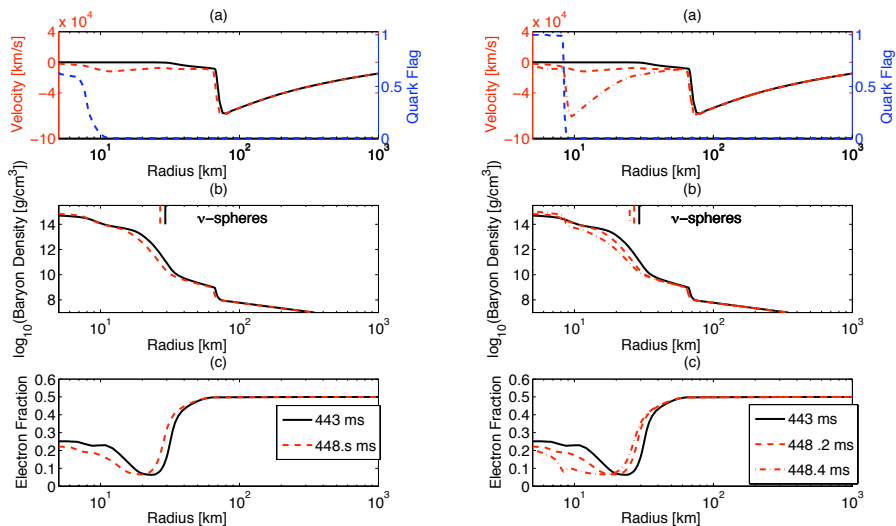


Fig: T.Fischer, $B^{1/4} = 165$ MeV, $10M_{\odot}$ progenitor

Core-collapse and second shock

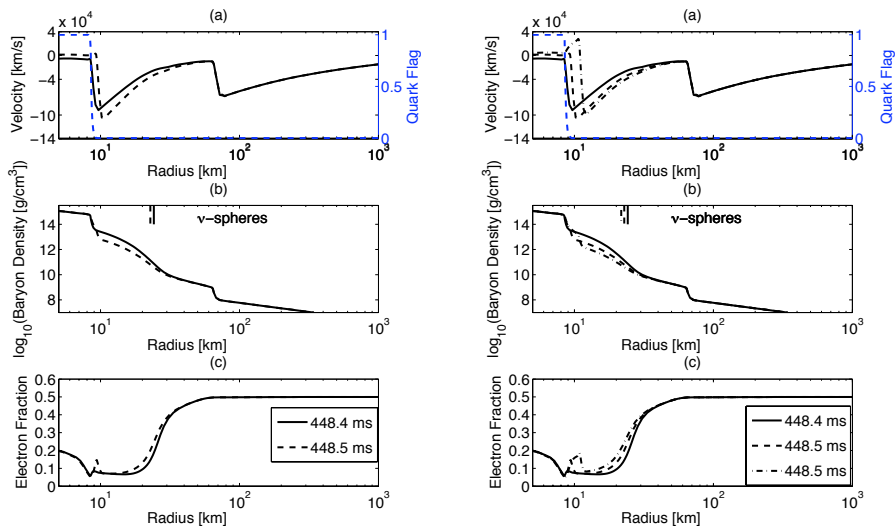


Fig: T.Fischer, $B^{1/4} = 165$ MeV, $10M_{\odot}$ progenitor

Core-collapse and second shock

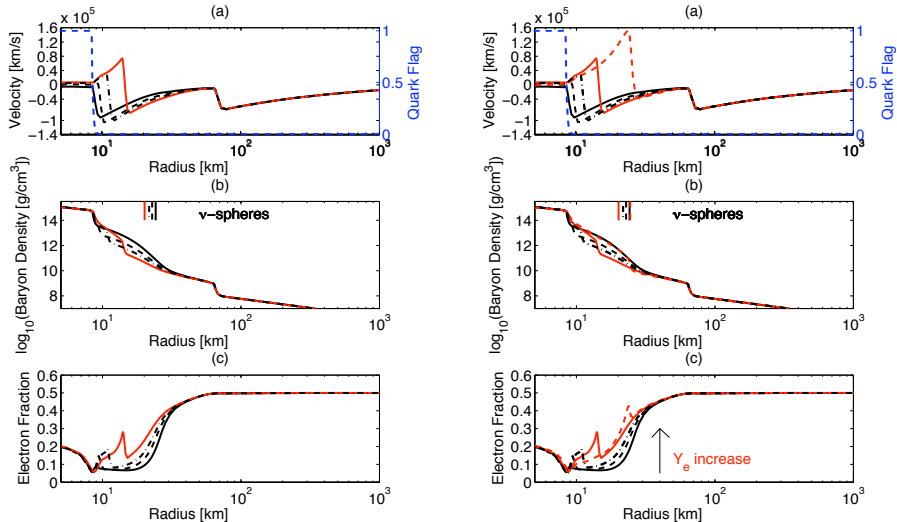


Fig: T.Fischer, $B^{1/4} = 165$ MeV, $10M_{\odot}$ progenitor