STORAGE RING EXPERIMENTS

at the Interface

of Atomic and Nuclear Physics

C. Brandau^{1,2}, <u>Chr. Kozhuharov^{1,3}</u>, Yu. A. Litvinov^{1,4}

¹GSI; ²EMMI; ³BfA & VBL; ⁴Universität Heidelberg

EMMI Workshop The ²²⁹mTh 'Nuclear Isomer Clock' GSI Darmstadt, Germany, September 25 – 27, 2012

Why Storage Rings?

"New directions in science are launched by new tools much more often than by new concepts.

The effect of a concept-driven revolution is to explain old things in new ways.

The effect of a tool-driven revolution is to discover new things that have to be explained."

Freeman Dyson, Imagined Worlds



The Role of Failure

You can't possibly get a good technology going without an enormous number of failures. It's a universal rule. If you look at bicycles, there were thousands of weird models built and tried before they found the one that really worked. You could never design a bicycle theoretically. Even now, after we've been building them for 100 years, it's very difficult to understand just why a bicycle works – it's even difficult to formulate it as a mathematical problem. But just by trial and error, we¹ found out how to do it, and the error was essential.

Freeman Dyson

(interview by Stewart Brand)

http://www.wired.com/wired/archive/6.02/dyson.html

¹ Caveat - cf. also H.M. Enzensberger, *Hammerstein oder Der Eigensinn*, Suhrkamp, Frankfurt/M, 2008 (footnote by C.K.)



Freeman Dyson 2007 (Wikipedia) Institute for Advanced Studies

Present GSI Accelerators

Heavy Ion Synchrotron,SIS, 2 AGeV for A/q=2 (1 AGeV U)

- Beams of all (chemical) elements and all stable isotopes: from hydrogen to uranium

- Broad range of energies:

from thermal to relativistic energies (2AGeV)

- Secondary beams of unstable (radioactive) nuclei;

ground state, isomers

- Unique beam properties:

well-defined charge states, cooled and stored beams

- Decelerated and cooled species HITRAP (4°K) Fragment Separator

Experimental Storage Ring ESR

- (pions)



ESR (Photograph by A. Zschau)



B. Franzke, NIM B 24/25 (1987) 18

Storage Rings: Electron Cooled Ion Beams



Cooling = narrowing velocity, size and divergence of stored ion beams



'Phase transition' to a linear ion chain



Recording the Schottky-noise





ILIMA: Masses and Halflives



Two-Body Beta Decay



Г н



<u>SCHOTTKY SIGNAL OF THE</u> BOUND-STATE β-DECAY ¹⁸⁷Re⁷⁵⁺ \rightarrow ¹⁸⁷Os⁷⁵⁺



 $T_{1/2}$ (¹⁸⁷Re⁷⁵⁺) = 33(2) y

Orbital Electron Capture

Conventional EC-theory:

W. Bambynek et al., Rev. Mod. Phys. 49, 1977

Gamow-Teller allowed transition $1^+ \square 0^+$

S-electron density at the nucleus:

 $|f_{S}(0)|^{2} \Box 1/n^{3}$

 P_{EC} (neutral atom) $\Box 2 \Box 1/n^3 = 2.4$

 P_{K} (H-like) $\Box 1 \Box 1/1^{3} = 1$

Conclusion: H-Like ion should have 41% longer half-life





3.39 min

140.

1 +

Orbital Electron Capture Decay of Few-Electron Ions



Orbital Electron Capture Decay of Few-Electron Ions



Gamow-Teller transition $1^+ \rightarrow 0^+$







F. Nolden et al., Nucl. Inst. Meth. A659 (2011) 69

10

1. Intertactor

1213

Three Parent He-Like ¹⁴²Pm lons



Photorecombination in cosmic plasmas

Courtesy D.W. Savin, Columbia Astrophysics Lab (CAL), New York, N.Y.



Radiative Recombination, Dielectronic Recombination and Nuclear Excitation by Electron Capture



THE PROPERTIES OF NEUTRAL AND IONIZED ATOMIC OXYGEN AND THEIR INFLUENCE ON THE UPPER ATMOSPHERE

By H. S. W. MASSEY, F.R.S. AND D. R. BATES, University College, London

Reports on Progress in Physics, vol. IX (1942-43)

The properties of neutral and ionized atomic oxygen

67

be got rid of in some way. If the pressure is so low that no third body is likely to be sufficiently close to receive the energy, only two modes of energy disposal are possible. Either it is directly radiated or else is communicated to a second electron in the O^+ ion, raising it to an excited orbital. In the latter case, the state in which the incident and second electron are both in excited orbitals is not stable and can either revert to the initial condition by re-emission of an electron or, by emission of radiation, undergo a transition to a stable state of the neutral atom, leading thus to recombination. This process is the inverse of that known as auto-ionization, and its contribution cannot be ignored. Since it depends on interaction between two electrons we refer to it henceforward as dielectronic recombination. We now consider separately the contributions from the two modes of energy disposal. 20 years later - Astrophys. J. 139, 776 (1964)

DIELECTRONIC RECOMBINATION AND THE TEMPERATURE OF THE SOLAR CORONA*

The process of dielectronic recombination in plasmas appears to be of much greater possible importance at high temperatures than has been previously realized (Massey and Bates 1942; Pery-Thorne and Garton 1960; Bates and Dalgarno 1962; Seaton 1962).

TABLE 1

DIELECTRONIC RECOMBINATION COEFFICIENTS FOR $Fe^{+15} + e$

Τ°Κ	a_d (tot) cm ² sec ⁻¹
106	1 1×10-10
2×10^{6}	$6 4 \times 10^{-11}$
3×10 ⁶	4.4×10^{-11}
4×10^{6}	$3 3 \times 10^{-11}$

The simplified formulation given above may easily be generalized. The full discussion will be given elsewhere together with the results of further and more precise calculations now in progress. The case of the Mg, Ca, and Fe ions should be of particular astrophysical interest.

I am indebted to the Institute for Advanced Study, Princeton, New Jersey, and to the Physics Department of New York University for their hospitality.

Alan Burgess¹

January 8, 1964 INSTITUTE FOR ADVANCED STUDY PRINCETON, NEW JERSEY

¹ Permanent address: Department of Physics, University College London.

40 years later – J.B..A. Mitchell et al, PRL. 50, 335 (1983), D.S. Belic et al PRL 50 339 (1983), P.F. Dittner et al. PRL 51 31 (1983) → first experimental observations

49 Years later (20 years ago): First GSI-ESR PRL

VOLUME 69, NUMBER 19

PHYSICAL REVIEW LETTERS

9 NOVEMBER 1992

Dielectronic and Radiative Recombination of Lithiumlike Gold

W. Spies, ^(a) A. Müller, ^(a) J. Linkemann, ^(a) A. Frank, and M. Wagner ^(b) Institut für Kernphysik, Universität Glessen, D-6300 Glessen, Germany

C. Kozhuharov, B. Franzke, K. Beckert, F. Bosch, H. Eickhoff, M. Jung, O. Klepper, W. König, P. H. Mokler, R. Moshammer, F. Nolden, U. Schaaf, P. Spädtke, and M. Steck Gesellschaft für Schwerionenforschung (GSI), D-6100 Darmstadt, Germany

> P. Zimmerer, N. Grün, and W. Scheid Institut für Theoretische Physik. Universität Giessen, D-6300 Giessen, Germany

> > M. S. Pindzola and N. R. Badnell

Department of Physics, Auburn University, Auburn, Alabama 36849 (Received 15 June 1992)

We report the first measurements on radiative and dielectronic recombination (RR and DR) of very highly charged ions at energies $E_{e.m.}$ as low as 0 to 50 eV. Novel techniques were employed at the heavy-ion storage ring ESR to obtain absolute recombination rates of Au^{3b+}(1s²2s). The increase of the RR rate for $E_{e.m.} \rightarrow 0$ could be recorded and single DR resonances in Au^{3b+}(1s²2p_{3/2}6l_j) resolved. The experimental data allow sensitive judgement of calculations based on the Bethe-Salpeter approach to RR combined with results for DR from perturbative-relativistic, semirelativistic, and fully relativistic theories. Significant discrepancies are found only in the energy region of the $2p_{3/2}6l_{3/2}$ resonances.



Cooled ion beams in well-defined ionic states

DR-Measurement

- Merged electron and ion beams
- Drift-tube defined variation of the relative ion-electron velocities
- Recombination
- Separation by the ESR dipol magnet

$$\alpha(E_{CM}) = \left\langle v_{rel} \sigma(E_{CM}) \right\rangle \\ = \frac{1}{1 - \beta_e \beta_i} \frac{R}{N_{Ion} n_e L_{U}}$$

Single particle detection (4π)



Dielectronic Recombination of Li-like Gold



Li-like Xe⁵¹⁺



Li-like Xe⁵¹⁺ (preliminary)

Continuum electron are captures into series of Rydberg state up to the series limit for both types of excitation of the 2s electron:

 $\begin{array}{c} 2s \rightarrow 2p_{1/2} \\ 2s \rightarrow 2p_{3/2} \end{array}$



NEEC

first mentioned by Goldanskii & Namiot Phys. Lett. 62B (1976)

Natural line widths $10^{-8} - 10^{-5}$ eV Resonance strength 1 beV

A. Pálffy, Z. Harman, W. Scheidt, PRA 73(2006)012715

Isotope	$E_{\rm c}({\rm keV})$	Туре	$Y_{\rm n}(1/{\rm s})$	$\Gamma_{d}(eV)$	S(b eV)
¹⁷⁴ Yb	4.89	E2	1.79 · 10 ⁸	4.85 · 10 ⁻⁸	0.09
¹⁷³ Yb	7.07	M1	7.32 · 10 ⁹	4.80 · 10 ⁻⁶	1.26
¹⁸⁵ ₇₅ Re	42.19	M1	$2.62 \cdot 10^{10}$	$2.36 \cdot 10^{-5}$	1.34
¹⁸⁷ ₇₅ Re	51.08	M1	$2.50 \cdot 10^{10}$	2.47 · 10 ⁻⁵	1.16

convolution with realistic electron energy distribution

$$\frac{1}{s\sqrt{2\pi}}\exp\left(-\frac{(E-E_0)^2}{2s^2}\right)$$

Main Idea, Feasibility, Proposal



Available online at www.sciencedirect.com



Physics Letters B 661 (2008) 330-334

PHYSICS LETTERS B

www.elsevier.com/locate/physletb

Nuclear excitation by electron capture followed by fast x-ray emission

Adriana Pálffy^{a,*}, Zoltán Harman^a, Christophor Kozhuharov^b, Carsten Brandau^b, Christoph H. Keitel^a, Werner Scheid^c, Thomas Stöhlker^{b,d}

^a Max-Planck-Institut für Kernphysik, Saupfercheckweg 1, 69117 Heidelberg, Germany
 ^b Gesellschaft für Schwerionenforschung (GSI), Planckstr. 1, 64291 Darmstadt, Germany
 ^c Institut für Theoretische Physik, Justus-Liebig-Universität Giessen, Heinrich-Buff-Ring 16, 35392 Giessen, Germany
 ^d Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Philosophenweg 12, 69120 Heidelberg, Germany

Received 23 October 2007; received in revised form 5 February 2008; accepted 5 February 2008

Available online 20 February 2008

Editor: V. Metag

(Challenge - only a few events per minute)

What would happen...

... if capture does not occur into the ground state!



NEECX into L shell of ²³⁸U

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Basic Idea for a NEECx experiment

<u>Illustrative Example</u>: ²³⁸U

- The first excited state in ²³⁸U decays predominantly via L-IC (44.9 keV)
- Bare uranium ions bombard cold electron target with the appropriate energy → → an electron is captured into the L shell and the nucleus is excited.
- K-vacancies in uranium decay in 10⁻¹⁷ sec, the L-electron goes to the K-shell, i.e. this partial decay width dominates the total one.
- The competing process of radiative recombination, RR, is very fast.
- The excited nucleus leaves the RR-background zone
- The excited nucleus can decay only via gamma emission i.e. slower by the conversion coefficient (270)
- The gammas can be detected in a background free zone in coincidence with ions which have captured one electron
- The colder the electrons the better.
- (The experimental proposal is currently being prepared.)

NEEC with and without x-ray

		He-, Be-, C-like ions	bare ions
AZX	nlj	S _{NEEC} (b eV)	$S_{NEECX}(beV)$
	$2s_{1/2}$	$8.8 imes10^{-3}$	$5.37 imes 10^{-2}$
$^{238}_{92}$ U	$2p_{1/2}$	$9.2 imes10^{-3}$	1.58
-	$2p_{3/2}$	$2.7 imes10^{-3}$	1.16
	$2s_{1/2}$	$5.9 imes10^{-3}$	$2.04 imes 10^{-2}$
²³² Th	$2p_{1/2}$	$7.7 imes10^{-3}$	$6.53 imes10^{-1}$
	$2p_{3/2}$	$2.6 imes10^{-3}$	$5.51 imes 10^{-1}$

IC channel is suppressed \rightarrow gain of two orders of magnitude

Prolongation of nuclear state lifetime via NEECX

	neutral	bare/H-like		Li-, B-, N-like
AZX	$ au_0$ (ns)	$ au_\gamma$ (ns)	nlj	$ au_{\mathit{nl}_j}$ (ns)
			2 <i>s</i> _{1/2}	36.2
238 ₉₂ U	0.292	185	$2p_{1/2}$	1.54
			$2p_{3/2}$	0.67
			$2s_{1/2}$	50.1
²³² Th	0.497	150	$2p_{1/2}$	2.35
			$2p_{3/2}$	1.04



- different time scales of NEEC and RR → spatial separation of emitted photons
- RR photons $\simeq 10^{-14} 10^{-16}$

•
$$\gamma$$
 photons $\simeq 10^{-9} - 10^{-11}$

 $au \simeq$ 10 - 100 ns \rightarrow 2 - 20 m

Stochastic Cooling



Stochastic cooling is in particular efficient for hot ion beams Cooling time τ scales as N_{ion} / bandwith

C

HOLTZ



Stochastic cooling (self-correction of trajectory)



The Nobel Prize in Physics 1984

"for their decisive contributions to the large project, which led to the discovery of the field particles W and Z, communicators of weak interaction"

http://www.nobel.se/physics/laureates/1984/



Carlo Rubbia 1/2 of the prize Italy CERN Geneva, Switzerland b. 1934



Simon van der Meer 1/2 of the prize the Netherlands

CERN Geneva, Switzerland b. 1925 The Nobel Prize in Physics 1984 Press Release Presentation Speech

Carlo Rubbia

Autobiography Nobel Lecture Banquet Speech

Simon van der Meer Autobiography Nobel Lecture

1983

1985 🕑

(GO)

The 1984 Prize in: <u>Physics</u> Chemistry Physiology or Medicine Literature Peace Economic Sciences

Find a Laureate:

Name

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DR of H-like U⁹¹⁺

D. Bernhardt et al. Phys. Rev. <u>A83</u> 020701(R) (2011)

Breit Interaction in dielectronic recombination of hydrogenlike uranium



Energy Resolution at ESR



HITRAP so far built / designed



CRYRING@ESR



Study Group

Norbert Angert Angela Bräuning-Demian Hakan Danared Wolfgang Enders Mats Engström Bernhard Franzke Anders Källberg Oliver Kester Michael Lestinsky Yuri Litvinov Markus Steck Thomas Stöhlker



2nd Realistic Scenario for a NEEC Experiment



Basic Idea



Why Li-Like Heavy-Ions?

2s_{1/2} -2p_{1/2} Energy Splitting for Li-like Ions



Radial density distribution

of low-lying electron orbitals



Low-Energy PR Spectrum of Li-like Neodymium (¹⁵⁰Nd⁵⁷⁺)



Li-like ion $(\Delta n = 0)$ 2 series of resonances: $E (2s_{1/2} \rightarrow 2p_{1/2}) = 139.2 \text{ eV}$ \Rightarrow n_{min} = 18 and to $E(2s_{1/2} \rightarrow 2p_{3/2}) = 729.1 \text{ eV}$ different overlap of 2s and 2p wavefunctions with the

2p wavefunctions with the nucleus
=> shift of a whole pattern of resonances.

Li-like ¹⁴²Nd⁵⁷⁺ vs. ¹⁵⁰Nd⁵⁷⁺



C. Brandau, et al., PRL 100 (2008) 073201

Li-like ¹⁴²Nd⁵⁷⁺ vs. ¹⁵⁰Nd⁵⁷⁺



^ANd⁵⁷⁺ DR-IS and Change in Mean Square Radius



C. Brandau, et al., PRL **100** (2008) 073201.Y.S. Kozhedub, et al., PRA **77** (2008) 032501.

Z. Harman, et al., in preparation.

Production of Li-like (!) Exotic Ions

complexities :

- + energy loss and straggling in thick target
- + distribution of charge states
- + cooling times ~ 1-5 min (for hot fragments far off β_{Cool} ?)
- => beam loss due to recombination in cooler (~95 % after 5 min)

Preparation of Li-like Exotic Beams in the ESR or: the Storage Ring as an Isotope Separator



Isotope Shift and Hyperfine Effects in the Dielectronic Recombination of In-Flight Synthesized ^AU⁸⁹⁺ (A=236, 237, 238)



Nov 2010 Test Run: Brevium - ^{234m}Pa



Kazimierz Fajans

Brevium, 1913



Otto Hahn and Lise Meitner Protoactinium, 1917 (1918)



Latin. Man's natural language. Spoils your style. Useful for reading the inscriptions on public fountains. Beware of quotations in Latin: they always conceal something improper. Gustave Flaubert

DR of 0⁻- Isomers in ²³⁴Pa⁸⁸⁺



preliminary

Conclusions

- A storage ring for high-Z ions in well-defined charge states, equipped with electron cooling capabilities for well-defined ion velocities and furnished with some nice instrumentation is a very good tool for nuclear and atomic physics studies.
- GPAC-accepted ^{229m}Th-Lol and proposal.
- Beam Time!?
- PhD students?