

Phase Diagram of Two-Color QCD in a Dyson-Schwinger Approach



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Quark Gluon Plasma meets Cold Atoms - Episode III
Hirschegg, Austria

Pascal Büscher, Michael Buballa, Jochen Wambach
Institut für Kernphysik, Technische Universität Darmstadt

Features of Two-Color QCD (QC_2D)

Formalism

Results

Summary & Outlook

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(Selected) Features of Two-Color QCD(QC₂D)



→ cf. talk by Lorenz v. Smekal (16:30 today)

Dirac operator obeys antiunitary symmetry

- ▶ no sign problem in fermion determinant at $\mu \neq 0$
- ▶ results from lattice calculations available

Pauli-Gürsey symmetry (at $\mu = m_q = 0$)

- ▶ extended flavor symmetry ($SU(N_f) \times SU(N_f) \times U(1)_B \rightarrow SU(2N_f)$)
- ▶ pions and diquarks degenerate

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Dyson-Schwinger Equation for Quarks (qDSE)



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$$S^{-1}(p) = Z_2 S_0^{-1}(p) + Z_{1F} g^2 \int \frac{d^4 q}{(2\pi)^4} \gamma_\mu \frac{\tau^a}{2} S(q) \Gamma_\nu^b(p, q) D_{\mu\nu}^{ab}(k = p - q)$$

Dyson-Schwinger Equation for Quarks (qDSE) ...and its Truncation



$$\text{---}\bullet\text{---}^{-1} = \text{---}\bullet\text{---}^{-1} + \text{---}\bullet\text{---}^{-1} \text{---}\bullet\text{---}$$

$$S^{-1}(p) = Z_2 S_0^{-1}(p) + Z_{1F} g^2 \int \frac{d^4 q}{(2\pi)^4} \gamma_\mu \frac{\tau^a}{2} S(q) \Gamma_\nu^b(p, q) D_{\mu\nu}^{ab}(k = p - q)$$



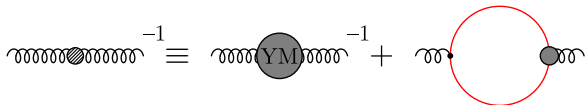


$$\text{---}\bullet\text{---}^{-1} = \text{---}\bullet\text{---}^{-1} + \text{---}\bullet\text{---}^{-1} \text{---}\bullet\text{---}$$

●: ansatz for the vertex


$$\Gamma_\mu^a(k) = \tilde{Z}_3 \tau^a \gamma_\mu \left(\frac{d_1}{d_2 + k^2} + \frac{k^2}{\Lambda^2 + k^2} \left(\frac{\beta_0 \alpha(\nu) \text{Ln} \left[\frac{k^2}{\Lambda^2} + 1 \right]}{4\pi} \right)^{2\delta} \right)$$

Dyson-Schwinger Equation for Gluons (gDSE) ...and its Truncation



$$\text{Gluon propagator with blob}^{-1} \equiv \text{Gluon propagator with YM blob}^{-1} + \text{Gluon propagator with quark loop}$$

$$D_{\mu\nu}(k) = \frac{Z_T(k)}{k^2 + G(k)} P_{\mu\nu}^T(k) + \frac{Z_L(k)}{k^2 + F(k)} P_{\mu\nu}^L(k)$$

 : quark loop in Hard-Dense/Hard-Thermal Loop Approximation

$$G(k_4, |\vec{k}|) = m_{g,T}^2(T, \mu, k) \dots, \quad F(k_4, |\vec{k}|) = 2m_{g,L}^2(T, \mu, k) \dots$$

with

$$m_{g,T/L}^2(T, \mu, k) \sim N_f \left(\frac{\pi}{3} T^2 + \frac{1}{\pi} \mu^2 \right) Z_{T/L}(k) \Gamma(k)$$

Dyson-Schwinger Equation for Gluons (gDSE) ...and its Truncation

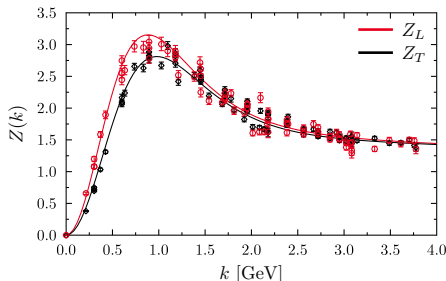
$$\text{gluon}^{-1} \equiv \text{gluon}^{-1} \text{YM} \text{gluon}^{-1} + \text{gluon} \text{loop} \text{gluon}$$

$$D_{\mu\nu}(k) = \frac{Z_T(k)}{k^2 + G(k)} P_{\mu\nu}^T(k) + \frac{Z_L(k)}{k^2 + F(k)} P_{\mu\nu}^L(k)$$



use input from Lattice
(C. S. Fischer et al., EPJ
C68 (2010))

right: $T = 109 \text{ MeV}$



Dyson-Schwinger Equation for Quarks (qDSE) ...in the Nambu-Gorkov Formalism

$$\text{---} \bullet \text{---} = \begin{pmatrix} S^+ & T^- \\ T^+ & S^- \end{pmatrix}$$

$$\begin{array}{c} \text{---} \bullet \text{---} \\ \text{---} \bullet \text{---} \end{array} = \begin{pmatrix} \Sigma^+ & \Phi^- \\ \Phi^+ & \Sigma^- \end{pmatrix}$$

$$\text{---} \bullet \text{---} = D_{\mu\nu}(k)$$

$$\bullet = \begin{pmatrix} \tau^a & \\ & -\tau^{aT} \end{pmatrix} \gamma_\mu \Gamma(p, q)$$

τ : SU(2) generators in color space

$$\begin{aligned} \Sigma^+(p) &\propto \int d^4 q \quad \tau^a \gamma^\mu S^+(q) \tau^a \gamma^\nu & \Gamma(p, q) D_{\mu\nu}(p - q) \\ \Phi^+(p) &\propto \int d^4 q \quad -\tau^{aT} \gamma^\mu T^+(q) \tau^a \gamma^\nu & \Gamma(p, q) D_{\mu\nu}(p - q) \end{aligned}$$

Pauli-Gürsey Symmetry ...in the Formalism

Consider color-flavor product:

$$\begin{aligned}\Sigma^+(p) &\propto \int d^4q \quad \tau^a \gamma^\mu S^+(q) \tau^a \gamma^\nu & \Gamma_\mu(p, q) D_{\mu\nu}(p - q) \\ \Phi^+(p) &\propto \int d^4q \quad -\tau^{aT} \gamma^\mu T^+(q) \tau^a \gamma^\nu & \Gamma_\mu(p, q) D_{\mu\nu}(p - q)\end{aligned}$$

color-flavor components of the propagators:

$$\begin{aligned}S^+(p) &= \mathbf{1} & s^+(p) \\ T^+(p) &= \tau^2 \sigma^2 & t^+(p)\end{aligned}$$

σ : SU(2) generators in flavor space

$$\begin{aligned}\Sigma^+(p) &\propto \int d^4q \quad 3 \gamma^\mu S^+(q) \gamma^\nu & \Gamma(p, q) D_{\mu\nu}(p - q) \\ \Phi^+(p) &\propto \int d^4q \quad 3 \gamma^\mu T^+(q) \gamma^\nu & \Gamma(p, q) D_{\mu\nu}(p - q)\end{aligned}$$

\Rightarrow same coupling for $\bar{q}q$ and qq !

Chiral Condensate

$$\langle \bar{q}q \rangle \propto \int d^4q \text{Tr} S^+(q)$$

- ▶ integral UV-divergent: quadratic $\sim m$, logarithmic $\sim \langle \bar{q}q \rangle$
- ▶ regularize with condensate at high temperature $\langle \bar{q}q \rangle' \equiv \langle \bar{q}q \rangle - \langle \bar{q}q \rangle|_{\mu=0}^{T=600 \text{ MeV}}$

Diquark Condensate

$$\langle qq \rangle \equiv \langle q^T C \tau^2 \sigma^2 \gamma_5 q \rangle \propto \int d^4q \text{Tr} [\tau^2 \sigma^2 \gamma_5 T^+(q)]$$

- ▶ integral converges

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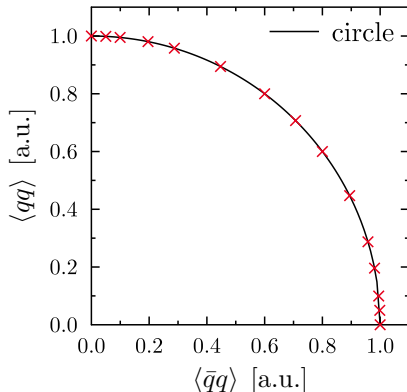
Pauli-Gürsey Symmetry ...in Numerical Calculations

$$\mu = m = 0$$

$$\langle \bar{q}q \rangle \propto \int d^4q \text{Tr} S^+(q)$$

$$\langle qq \rangle \equiv \langle q^T C \tau^2 \sigma^2 \gamma_5 q \rangle$$

$$\propto \int d^4q \text{Tr} [\tau^2 \sigma^2 \gamma_5 T^+(q)]$$



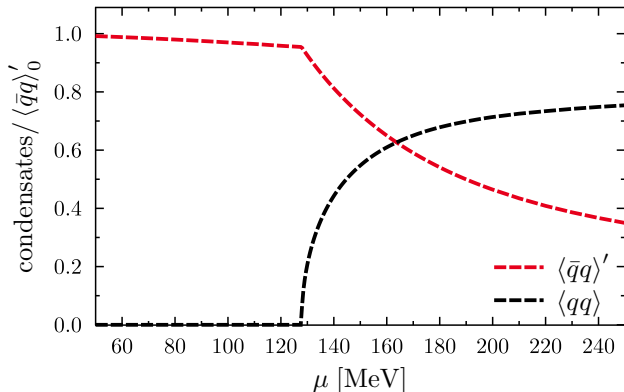
Silver Blaze Property (at finite quark mass)

- ▶ for $T = 0$, thermodynamic observables independent of μ unless μ exceeds lowest excitation threshold
- ▶ for QC₂D: observables unchanged for $\mu < m_\pi/2$

Silver Blaze Property (at finite quark mass)

for $T = 0$ and $\mu < m_\pi/2$

- ▶ observables should be independent of μ



▶ $m \neq 0$

▶ $T = 20$ MeV

▶ $\langle \bar{q}q \rangle'_0 \equiv \langle \bar{q}q \rangle - \langle \bar{q}q \rangle_{\mu=0}^{T=600 \text{ MeV}}$

▶ $\langle \bar{q}q \rangle'_0 \equiv \langle \bar{q}q \rangle'_{\mu=0}^{T=20 \text{ MeV}}$

$$\text{Diagram 1}^{-1} \equiv \text{Diagram 2}^{-1} + \text{Diagram 3}$$

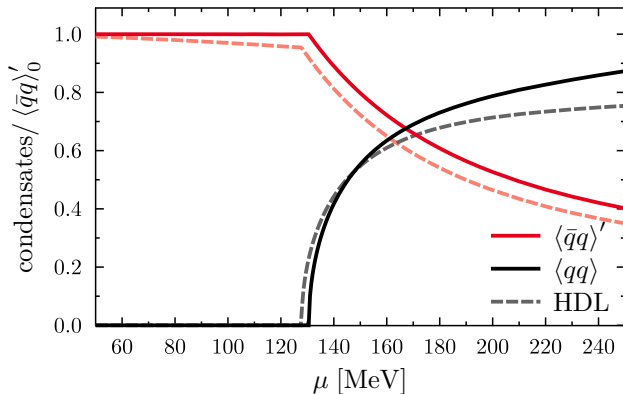
The diagram shows the decomposition of a self-energy correction to a gluon line. The first term is a gluon line with a shaded circular loop. The second term is a gluon line with a grey circular loop labeled 'YM'. The third term is a gluon line with a red circular loop and a grey circular loop.

$$\text{Red Loop} \sim N_f \left(\frac{\pi}{3} T^2 + \frac{1}{\pi} \mu^2 \right)$$

⇒ explicit dependence on μ violates Silver Blaze property

Silver Blaze Property (at finite quark mass)

without HDL:



▶ $m \neq 0$

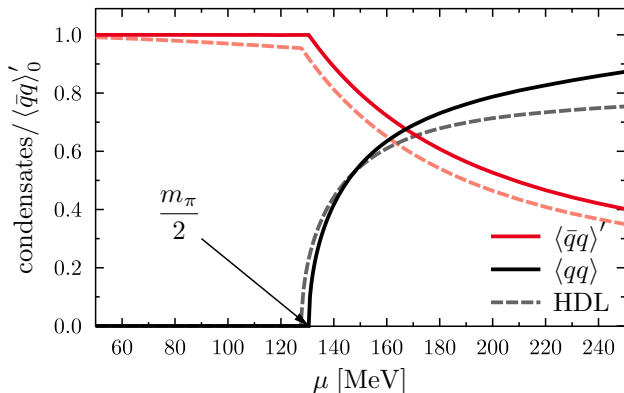
▶ $T = 20 \text{ MeV}$

▶ $\langle \bar{q}q \rangle' \equiv$
 $\langle \bar{q}q \rangle - \langle \bar{q}q \rangle|_{\mu=0}^{T=600 \text{ MeV}}$

▶ $\langle \bar{q}q \rangle'_0 \equiv \langle \bar{q}q \rangle'|_{\mu=0}^{T=20 \text{ MeV}}$

Silver Blaze Property (at finite quark mass)

without HDL:



▶ $m \neq 0$

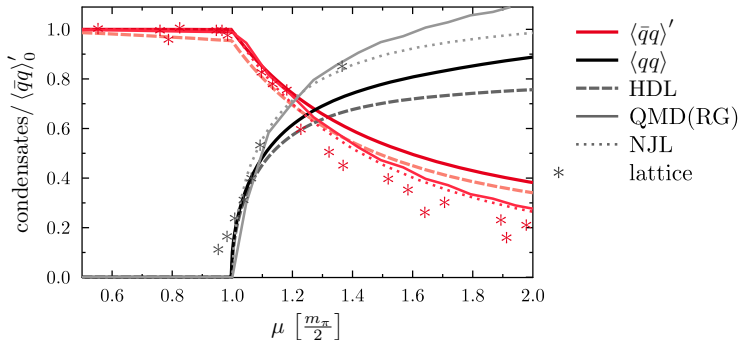
▶ $T = 20 \text{ MeV}$

▶ $\langle \bar{q}q \rangle'_0 \equiv \langle \bar{q}q \rangle - \langle \bar{q}q \rangle|_{\mu=0}^{T=600 \text{ MeV}}$

▶ $\langle \bar{q}q \rangle'_0 \equiv \langle \bar{q}q \rangle'|_{\mu=0}^{T=20 \text{ MeV}}$

Silver Blaze Property

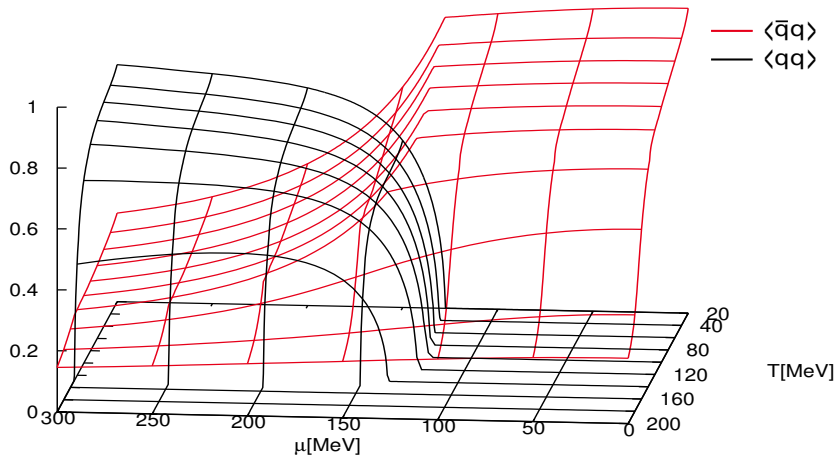
...comparison with other approaches



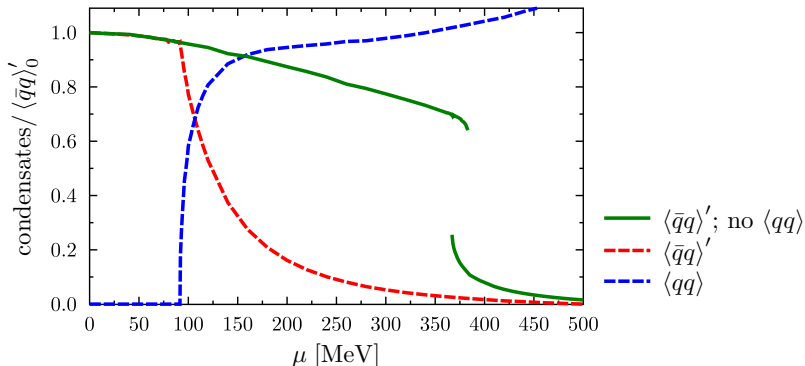
non-DSE plots from N. Strodthoff et al., 1112.5401v2[hep-ph]

lattice data from S. Hands et al., EPJ C17 (2000); EPJ C22 (2001)

Phase Diagram including T Direction



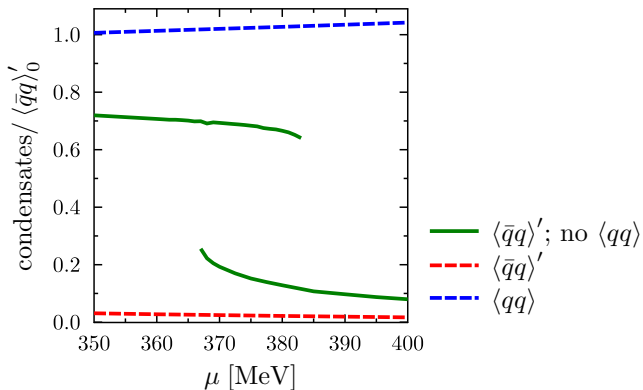
if diquarks are suppressed



half the current quark mass as before

$$\langle \bar{q}q \rangle' \equiv \langle \bar{q}q \rangle - \langle \bar{q}q \rangle \Big|_{\mu=500 \text{ MeV}}^{T=20 \text{ MeV}}$$

if diquarks are suppressed



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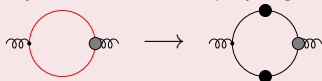
Summary & Outlook

Introduced a Dyson-Schwinger approach to calculate QC_2D which shows

- ▶ Pauli-Gürsey symmetry
- ▶ Silver Blaze property

Outlook

- ▶ improve truncation (in progress):



- ▶ consider other quantities, e.g. the dual condensate
- ▶ calculate m_π from Bethe-Salpeter equations
⇒ cross-check with m_π from Silver Blaze property



Thank you!

Also thanks to Daniel Müller and Nils Strodthoff for fruitful discussions.