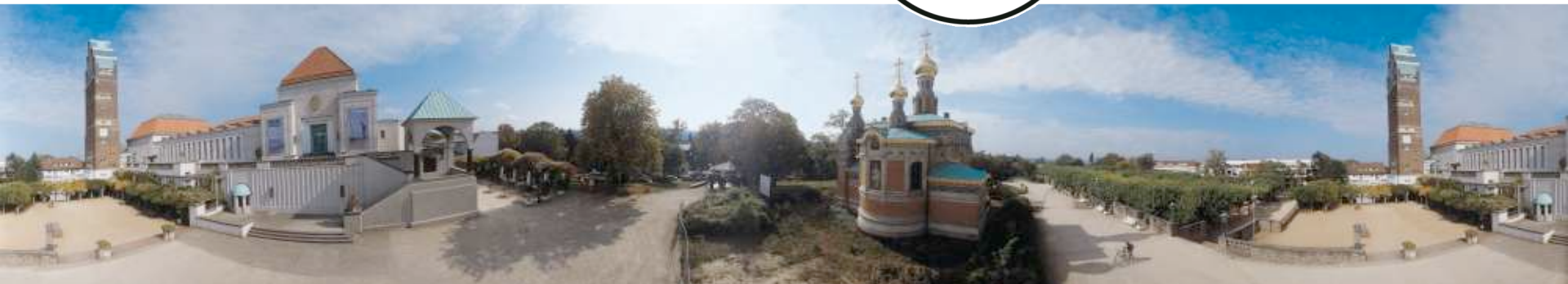


Neutron matter: From cold atoms to astrophysics

Achim Schwenk



TECHNISCHE
UNIVERSITÄT
DARMSTADT



EMMI Workshop “Quark Gluon Plasma meets Cold Atoms III”
Hirschegg, Aug. 27, 2012



DFG



*Minerva
Stiftung*



Bundesministerium
für Bildung
und Forschung

Main points

Advances in nuclear forces and nuclear matter theory

Impact on neutron stars:

provides strong constraints for the equation of state

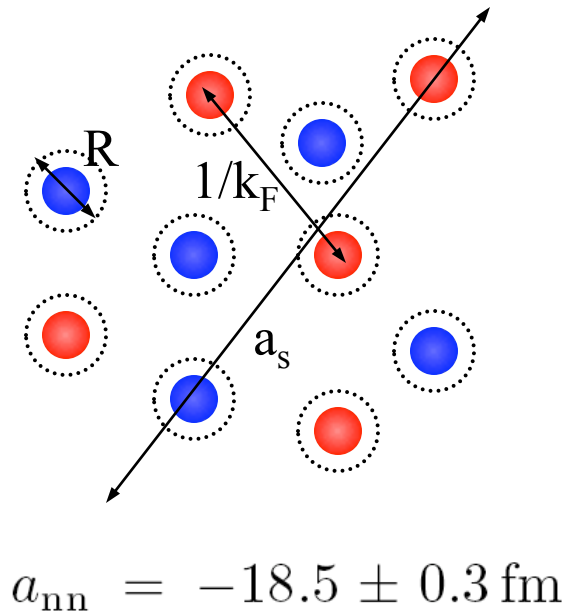
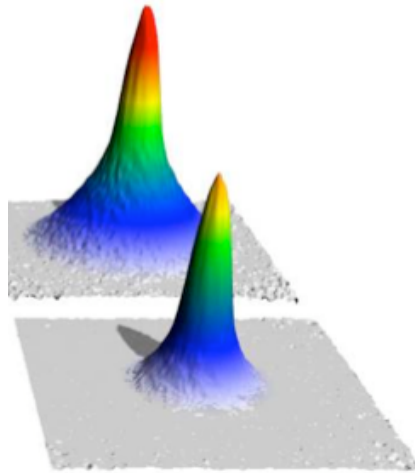
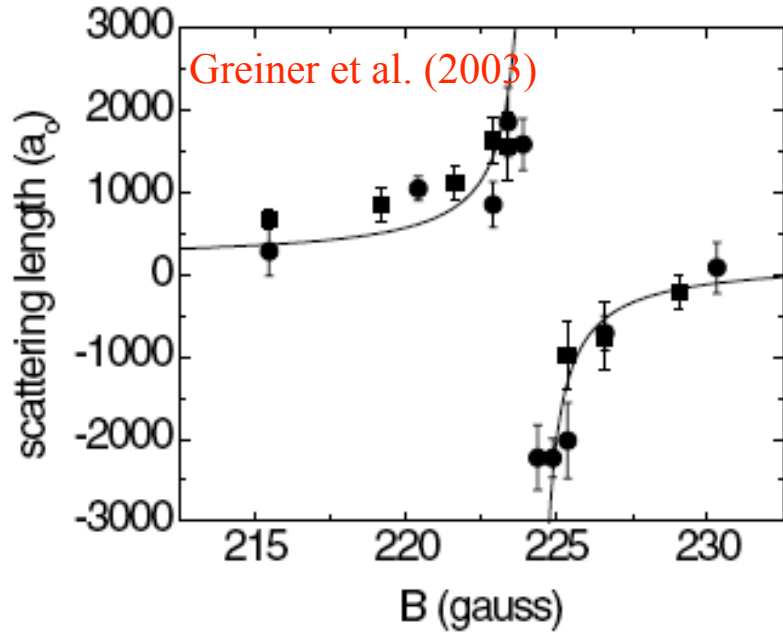
neutron star radius 9.9-13.8 km for $M=1.4 M_{\text{sun}}$ ($\pm 15\%$)

K. Hebeler, J.M. Lattimer, C.J. Pethick, AS, PRL **105**, 161102 (2010) and in prep.
I. Tews, T. Krüger, K. Hebeler, AS, arXiv:1206.0025.

Impact on neutrino-matter interactions

S. Bacca, K. Hally, C.J. Pethick, AS, PRC **80**, 032802(R) (2009) and
S. Bacca, K. Hally, M. Liebendörfer, A. Perego, C.J. Pethick, AS, ApJ (2012).

Large scattering lengths: Universal properties at low densities



$$a_{nn} = -18.5 \pm 0.3 \text{ fm}$$

strong interactions via Feshbach resonances

large for neutrons

dilute Fermi system with large scattering length has **universal properties**

$$0 \leftarrow \frac{1}{a_s} \ll k_F \ll \frac{1}{r_e}, \frac{1}{R}, \dots \rightarrow \infty$$

strongly-interacting
dilute

only Fermi momentum or density sets scale

physics is independent of interaction/system details:

from dilute neutron matter to resonant ${}^6\text{Li}$ or ${}^{40}\text{K}$ atoms in traps

Chiral Effective Field Theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV

	NN	3N	4N
LO $\mathcal{O}\left(\frac{Q^0}{\Lambda^0}\right)$			
NLO $\mathcal{O}\left(\frac{Q^2}{\Lambda^2}\right)$			
N ² LO $\mathcal{O}\left(\frac{Q^3}{\Lambda^3}\right)$			
N ³ LO $\mathcal{O}\left(\frac{Q^4}{\Lambda^4}\right)$			
	+ ...	+ ...	+ ...

limited resolution at low energies,
can expand in powers $(Q/\Lambda_b)^n$

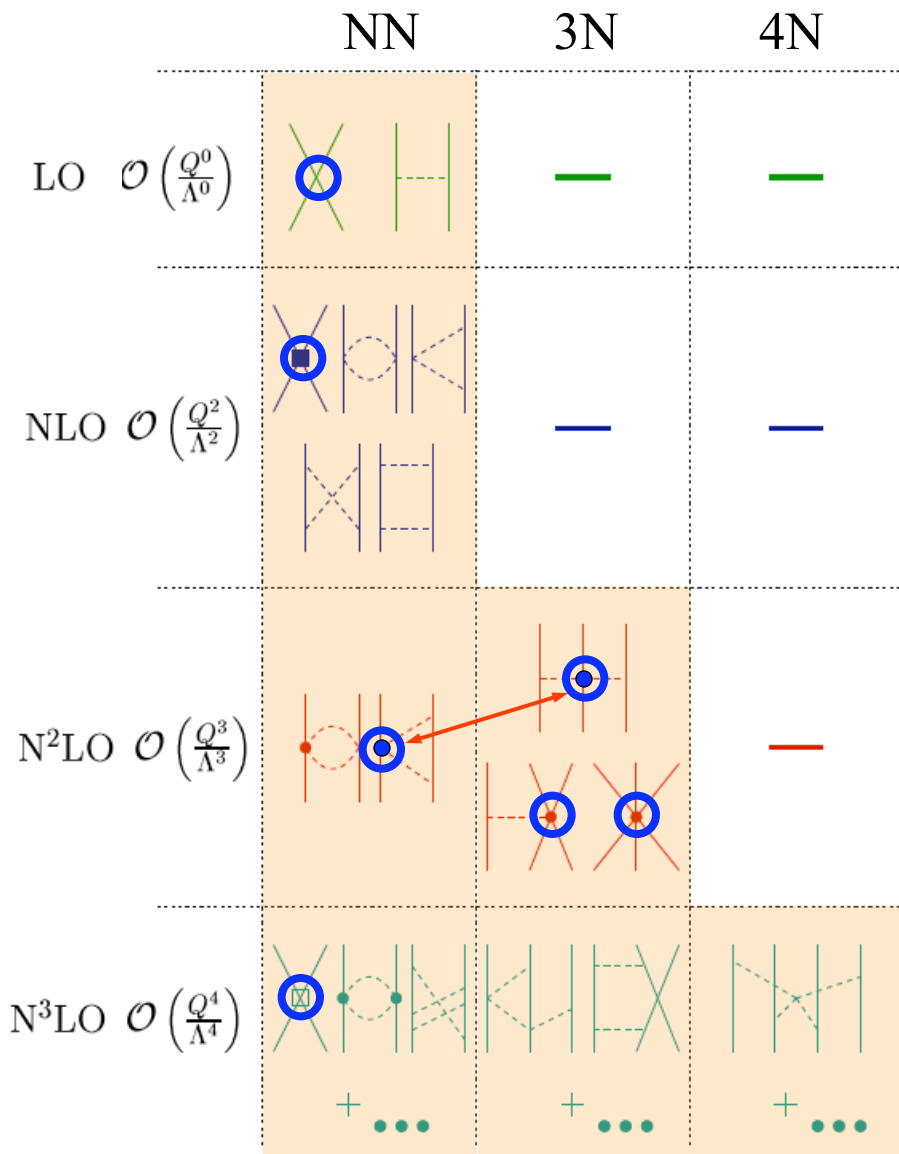
LO, $n=0$ - leading order,
NLO, $n=2$ - next-to-leading order,...

expansion parameter $\sim 1/3$

(compare to multipole expansion
for a charge distribution)

Chiral Effective Field Theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



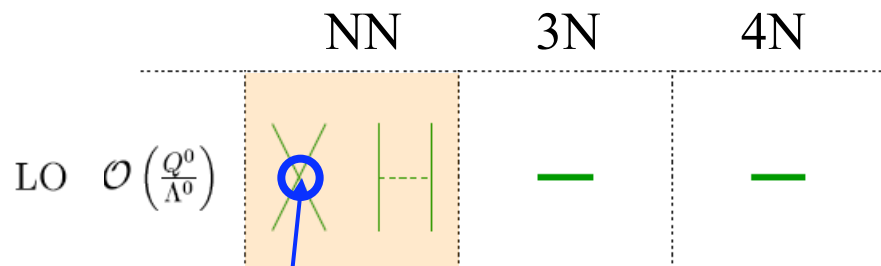
include long-range pion physics

few short-range couplings,
fit to experiment once

systematic: can work to desired
accuracy and obtain **error estimates**

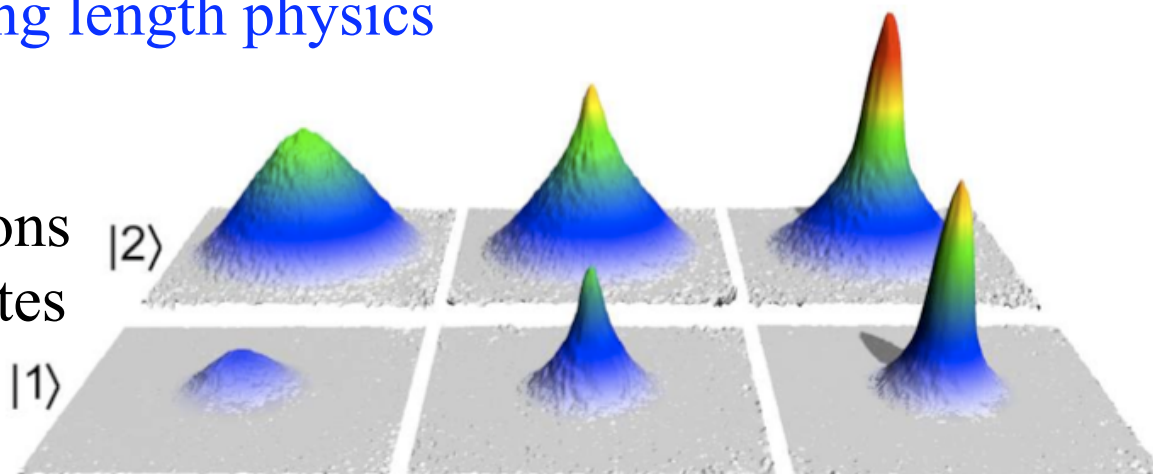
Chiral Effective Field Theory for nuclear forces

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large scattering length physics

${}^6\text{Li}$ fermions
2 spin states

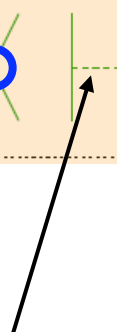
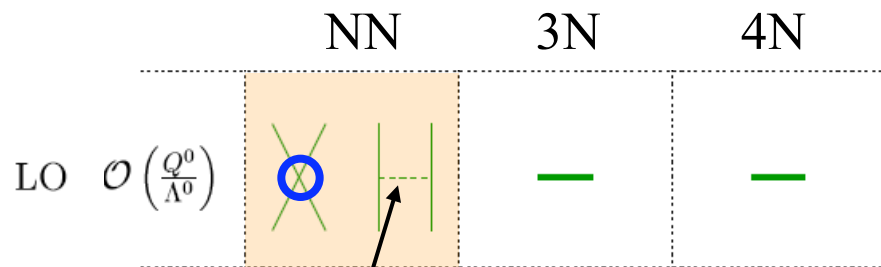


from M. Zwierlein

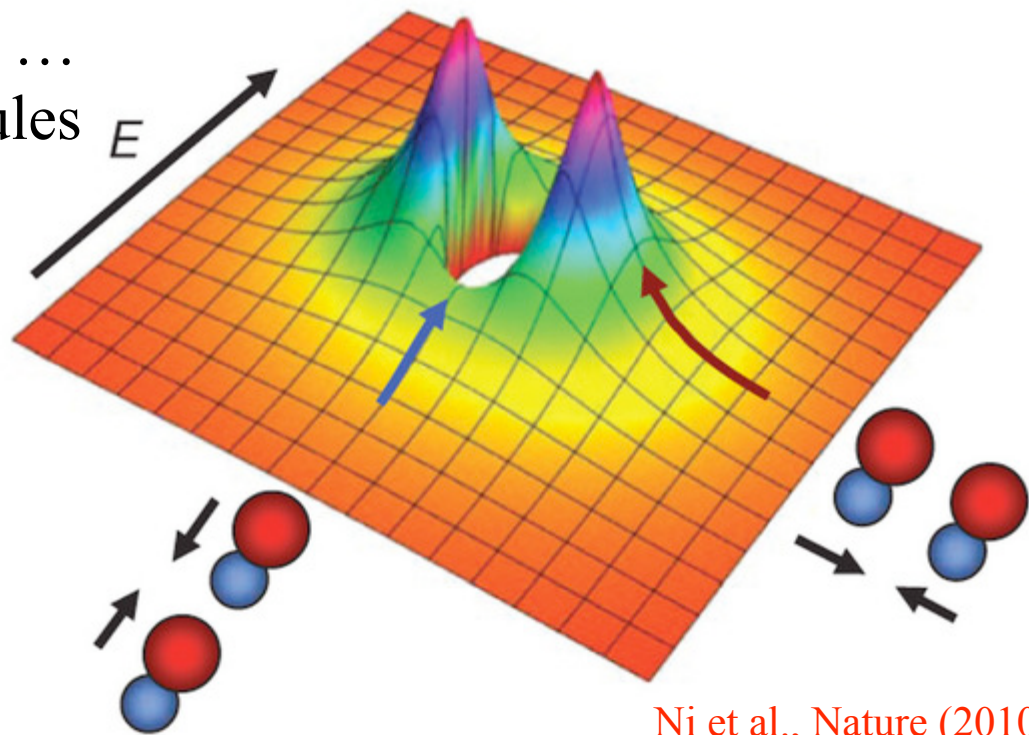
neutrons with same density, temperature and spin polarization
have the same properties!

Chiral Effective Field Theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV

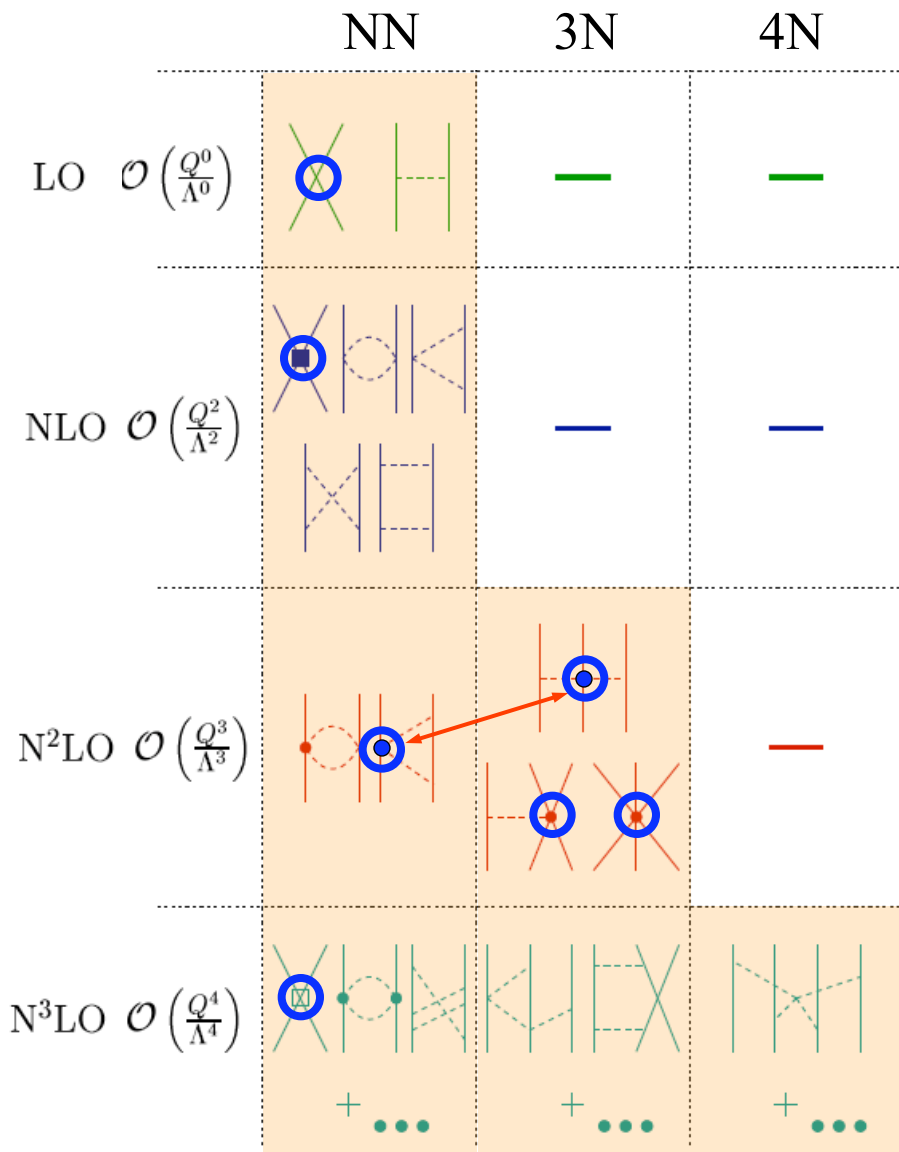


pion tensor/dipole interactions + ...
 → compare to cold polar molecules



Chiral Effective Field Theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



include long-range pion physics

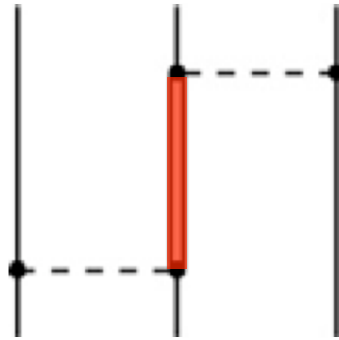
few short-range couplings,
fit to experiment once

systematic: can work to desired
accuracy and obtain **error estimates**

Why are there three-nucleon (3N) forces?

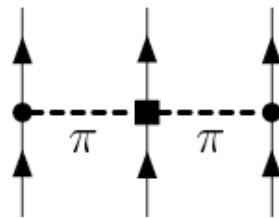
Nucleons are finite-mass composite particles,
can be excited to resonances

dominant contribution from $\Delta(1232 \text{ MeV})$

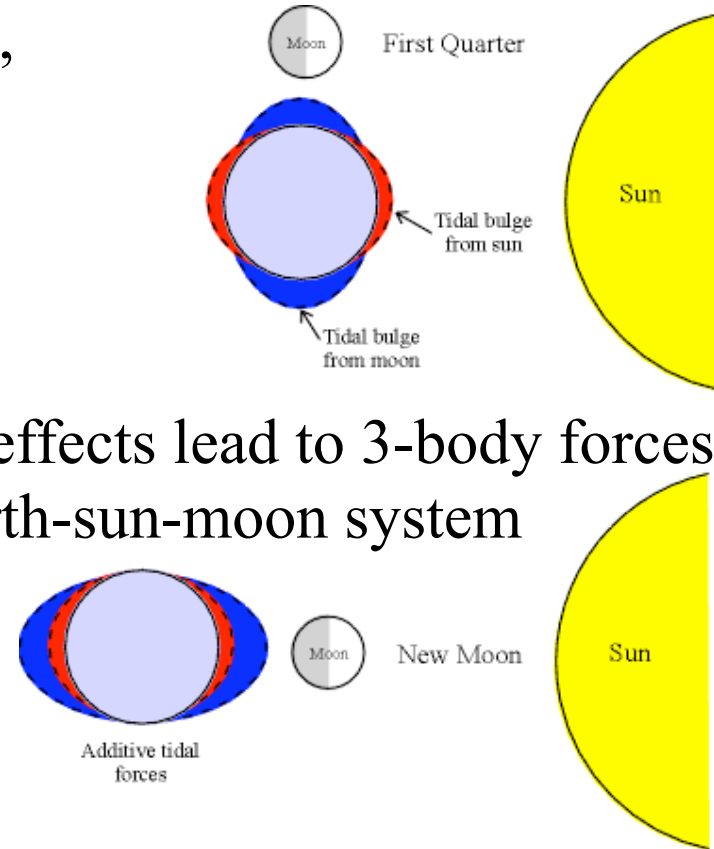


+ many shorter-range parts

in chiral EFT (Delta-less):



tidal effects lead to 3-body forces
in earth-sun-moon system



EFT provides a systematic and powerful approach to organize 3N forces

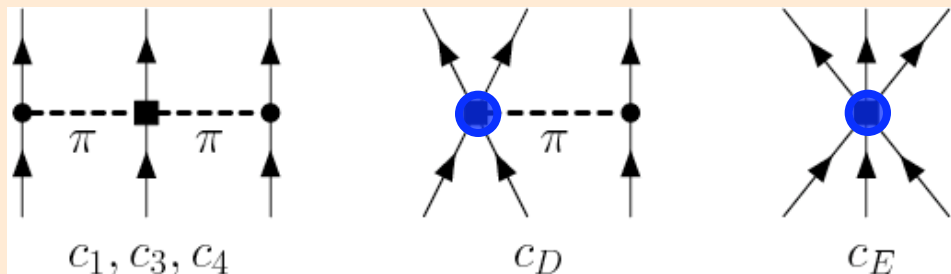
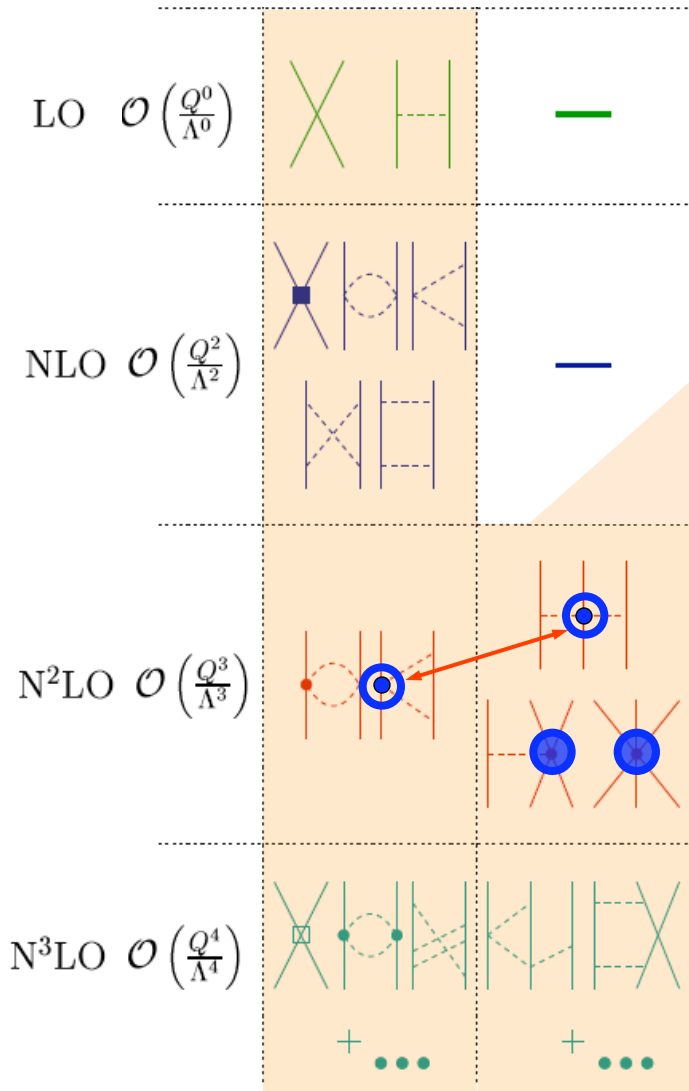
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NN 3N

consistent NN-3N interactions

3N,4N: only 2 new couplings to N³LO



long-range 3N: c_i from π N and NN

$$c_1 = -0.9_{-0.5}^{+0.2}, \quad c_3 = -4.7_{-1.0}^{+1.2}, \quad c_4 = 3.5_{-0.2}^{+0.5}$$

3- and 4-neutron forces are predicted to N³LO ($c_{D,E}$ don't contribute)

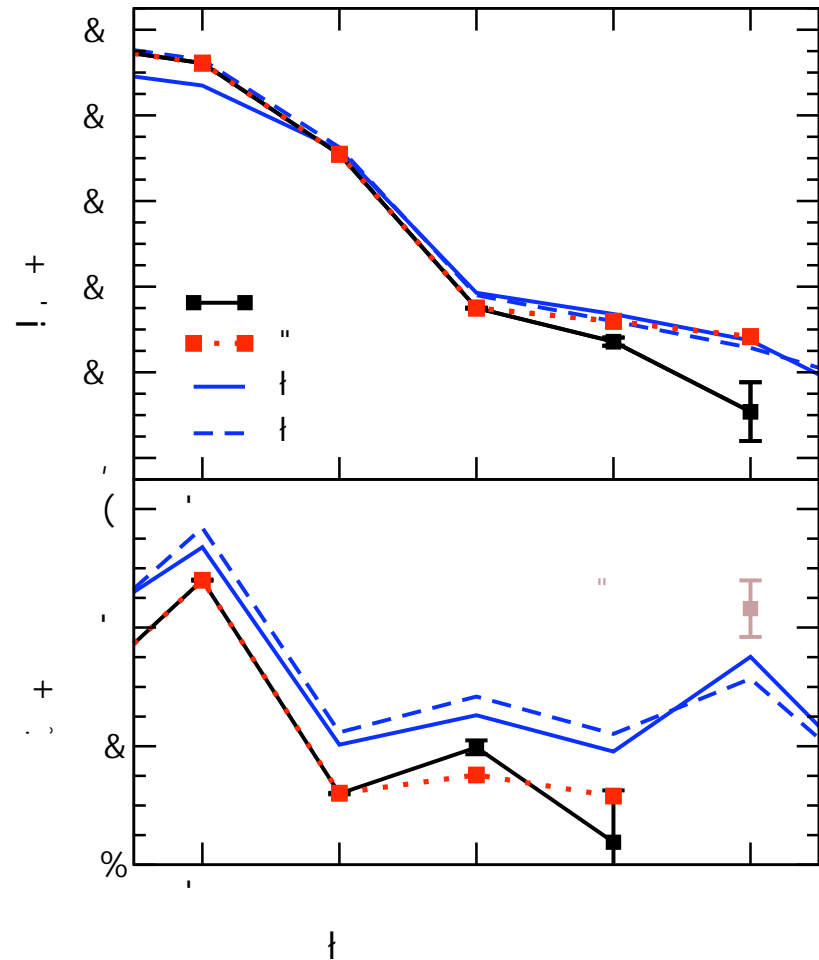
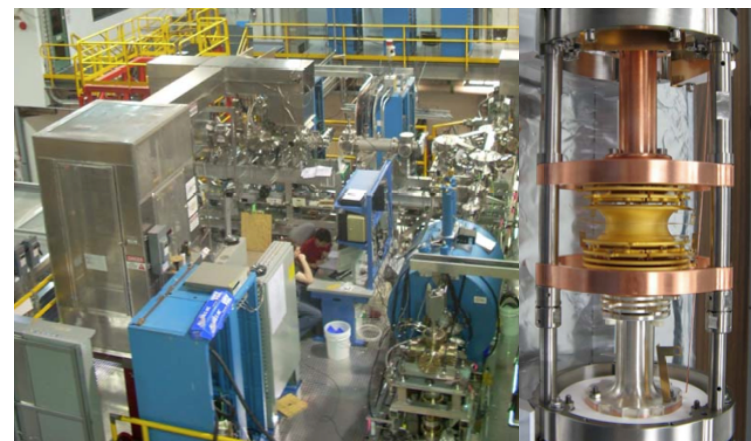
Hebeler, AS (2010)

new $^{51,52}\text{Ca}$ TITAN measurements

^{52}Ca is 1.75 MeV more bound compared to atomic mass evaluation

Gallant et al. (2012)

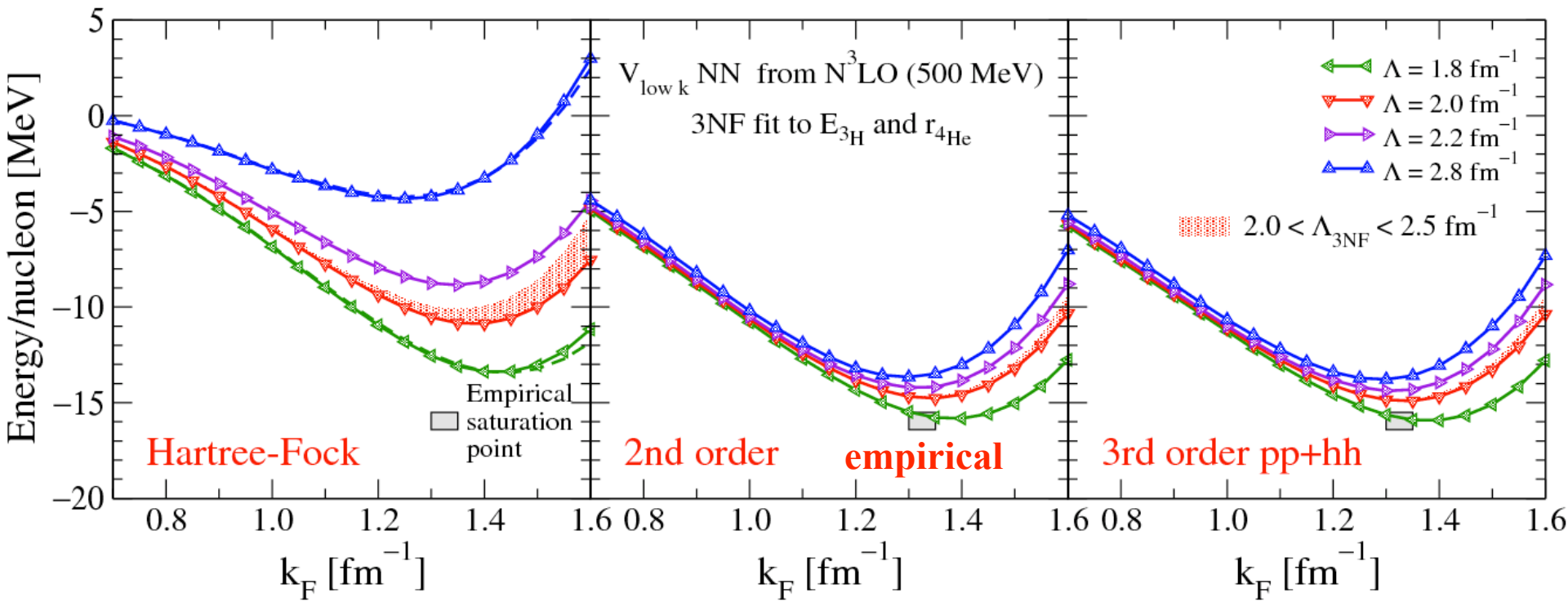
behavior of two-neutron separation energy S_{2n} and odd-even staggering Δ_n agrees with NN+3N predictions



Impact of 3N forces on nuclear matter

chiral 3N forces fit to light nuclei
predict nuclear matter saturation
with theoretical uncertainties

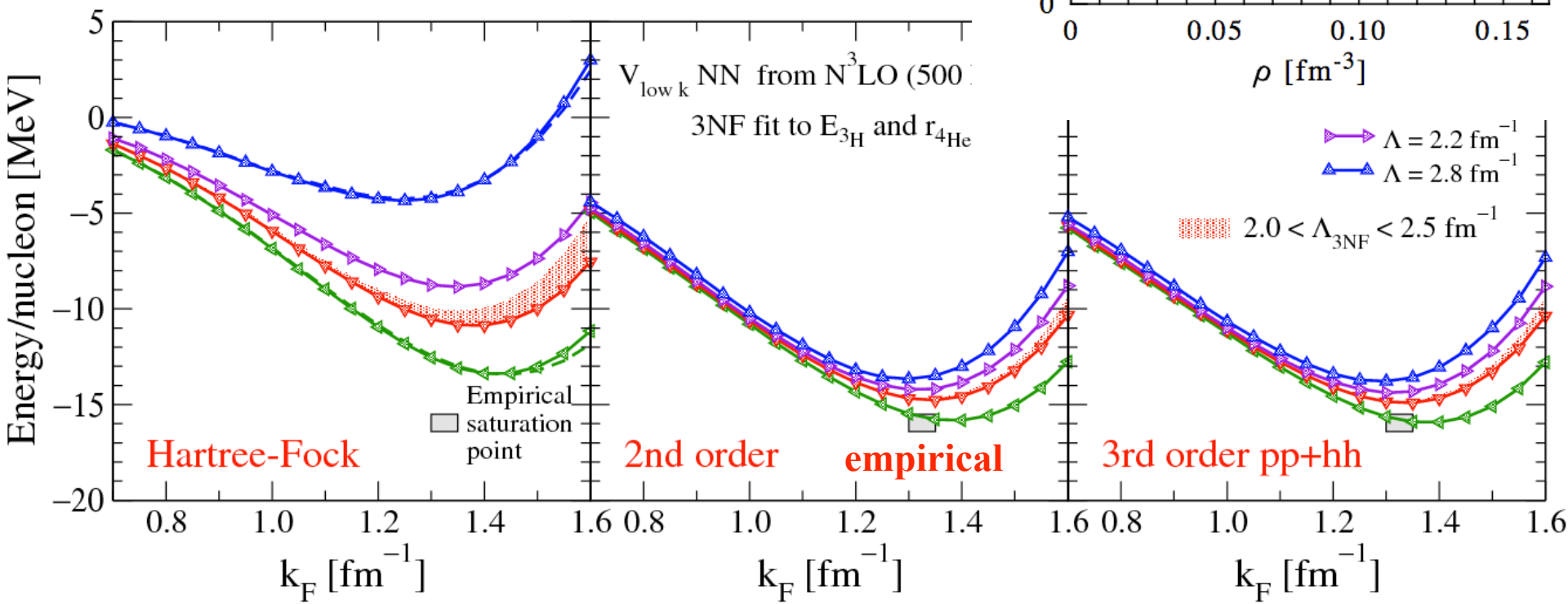
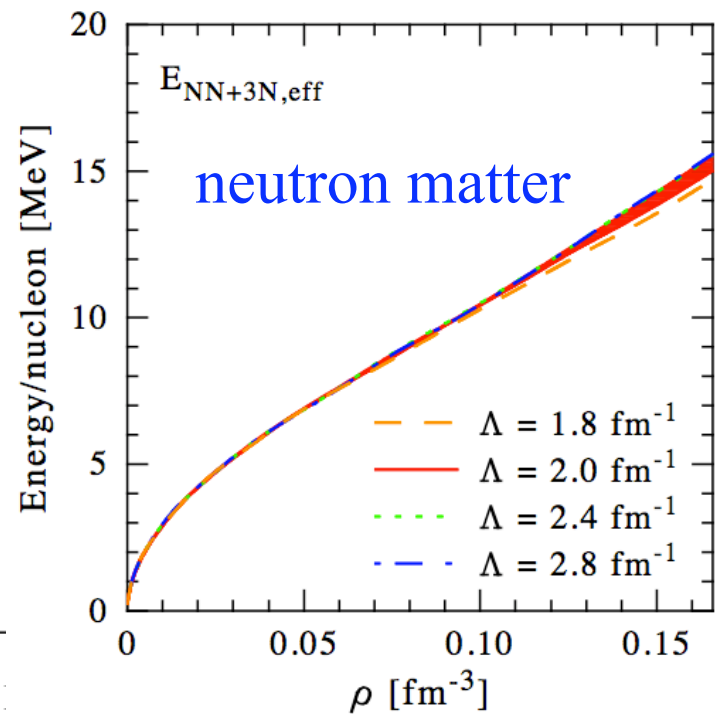
Hebeler et al. (2011), Bogner et al. (2005)



Impact of 3N forces on neutron matter

neutron matter is simpler system,
only long-range parts of 3N forces
contribute (c_1 and c_3)

Hebeler, AS (2010)



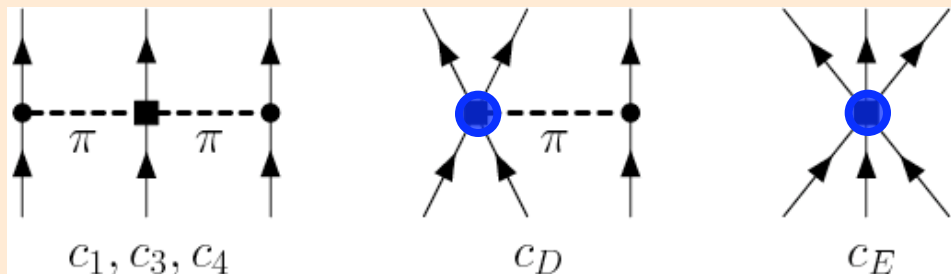
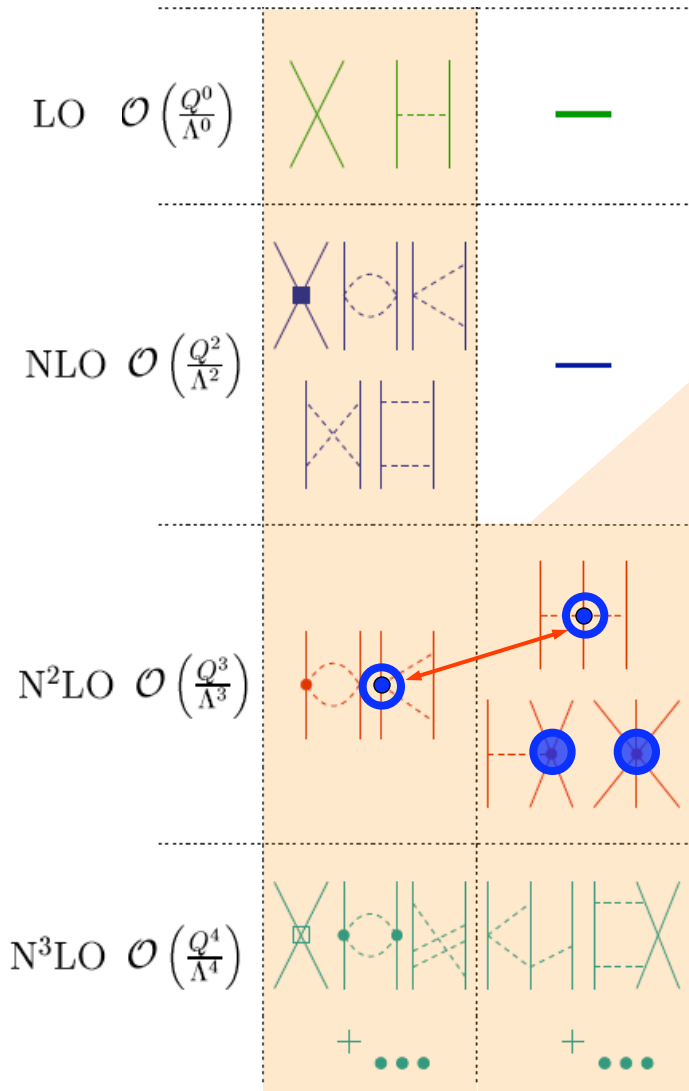
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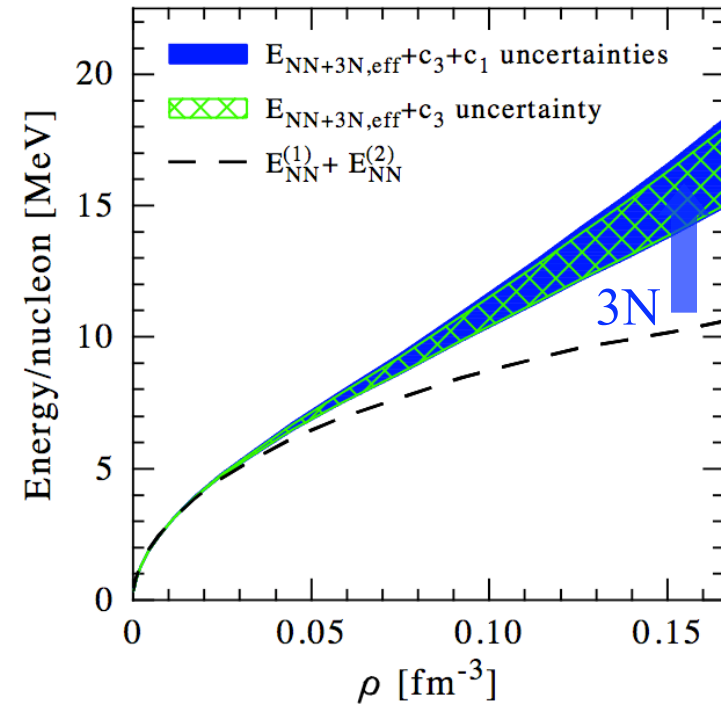
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Hebeler, AS (2010)

Impact of 3N forces on neutron matter

neutron matter uncertainties
dominated by 3N forces (c_3 coupling)

Hebeler, AS (2010)



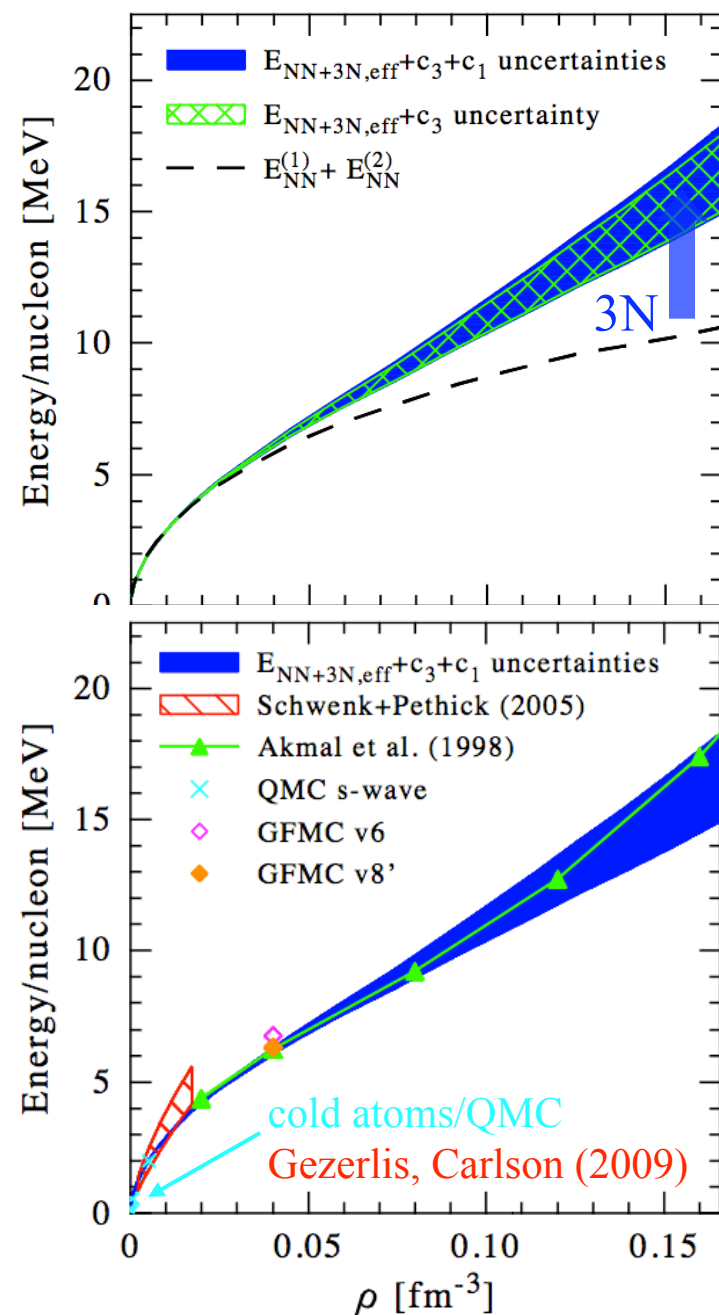
Impact of 3N forces on neutron matter

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Hebeler, AS (2010)

other microscopic calculations within band
(but without uncertainties)

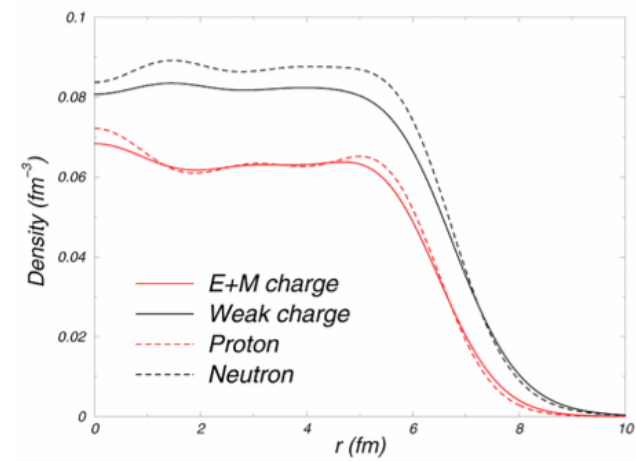


Neutron skin of ^{208}Pb

probes neutron matter energy/pressure,
neutron matter band predicts

neutron skin of ^{208}Pb : 0.17 ± 0.03 fm

Hebeler et al. (2010)

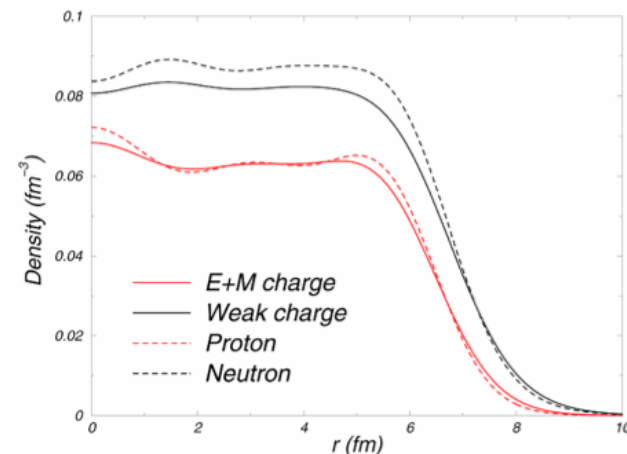


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in excellent agreement with extraction from complete E1 response

$0.156 + 0.025 - 0.021$ fm

PRL 107, 062502 (2011)

PHYSICAL REVIEW LETTERS

week ending
5 AUGUST 2011

Complete Electric Dipole Response and the Neutron Skin in ^{208}Pb

A benchmark experiment on ^{208}Pb shows that polarized proton inelastic scattering at very forward angles including 0° is a powerful tool for high-resolution studies of electric dipole (E1) and spin magnetic dipole (M1) modes in nuclei over a broad excitation energy range to test up-to-date nuclear models. The extracted E1 polarizability leads to a neutron skin thickness $r_{\text{skin}} = 0.156^{+0.025}_{-0.021}$ fm in ^{208}Pb derived within

PREX: neutron skin from parity-violating electron-scattering at JLAB

electron exchanges Z-boson, couples preferentially to neutrons

goal II: ± 0.06 fm

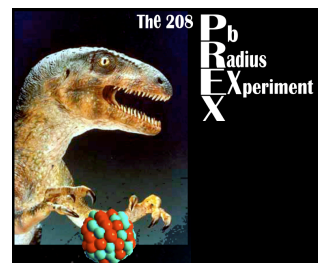
PRL 108, 112502 (2012)

PHYSICAL REVIEW LETTERS

week ending
16 MARCH 2012

Measurement of the Neutron Radius of ^{208}Pb through Parity Violation in Electron Scattering

We report the first measurement of the parity-violating asymmetry A_{PV} in the elastic scattering of polarized electrons from ^{208}Pb . A_{PV} is sensitive to the radius of the neutron distribution (R_n). The result $A_{\text{PV}} = 0.656 \pm 0.060(\text{stat}) \pm 0.014(\text{syst})$ ppm corresponds to a difference between the radii of the neutron and proton distributions $R_n - R_p = 0.33^{+0.16}_{-0.18}$ fm and provides the first electroweak observation of the neutron skin which is expected in a heavy, neutron-rich nucleus.



Symmetry energy and pressure of neutron matter

neutron matter band predicts
symmetry energy S_v and
its density dependence L

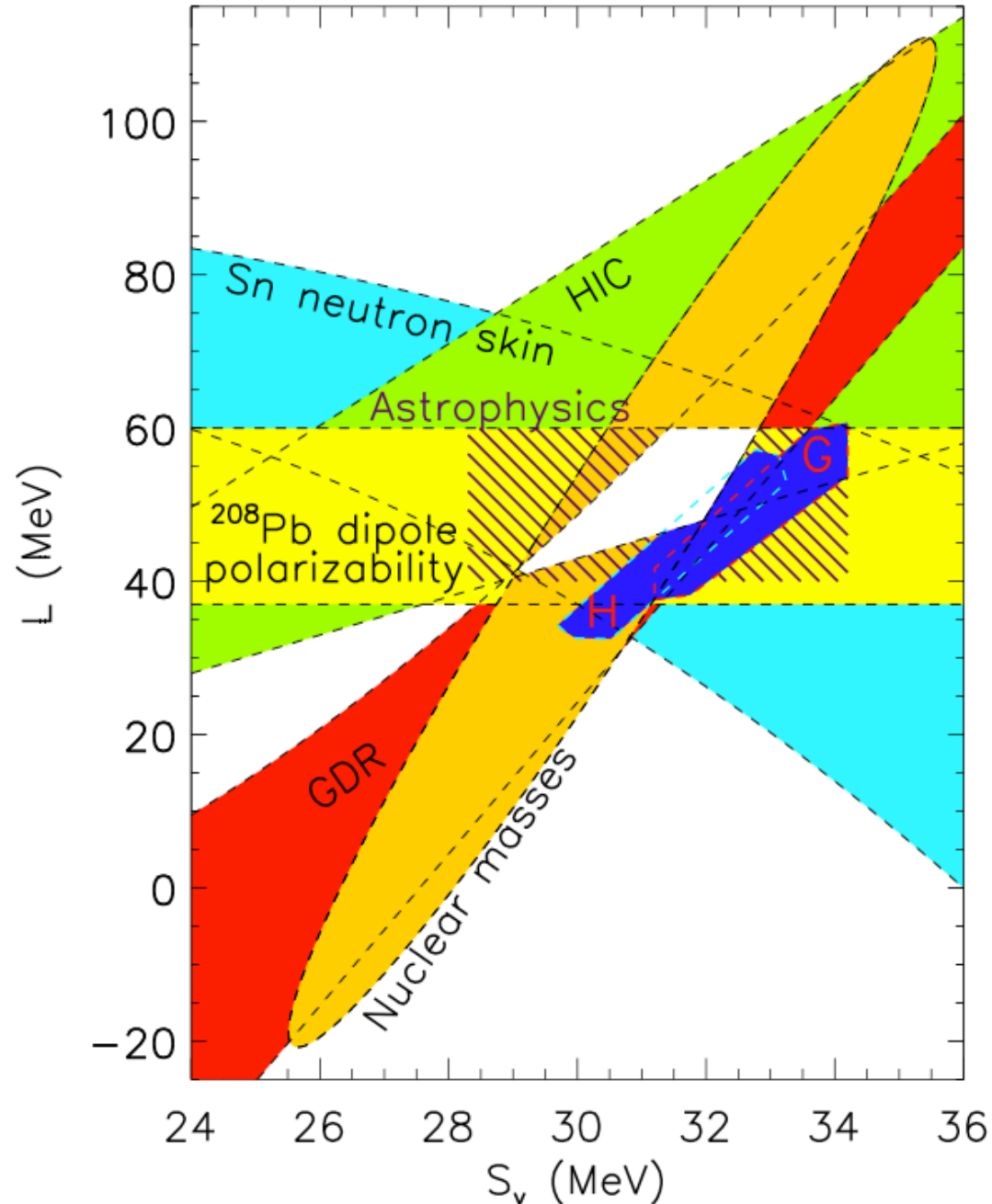
comparison to experimental
and observational constraints
Lattimer, Lim (2012)

neutron matter constraints

H: Hebeler et al. (2010) and in prep.

G: Gandolfi et al. (2011)

predicts correlation
but not range of S_v and L



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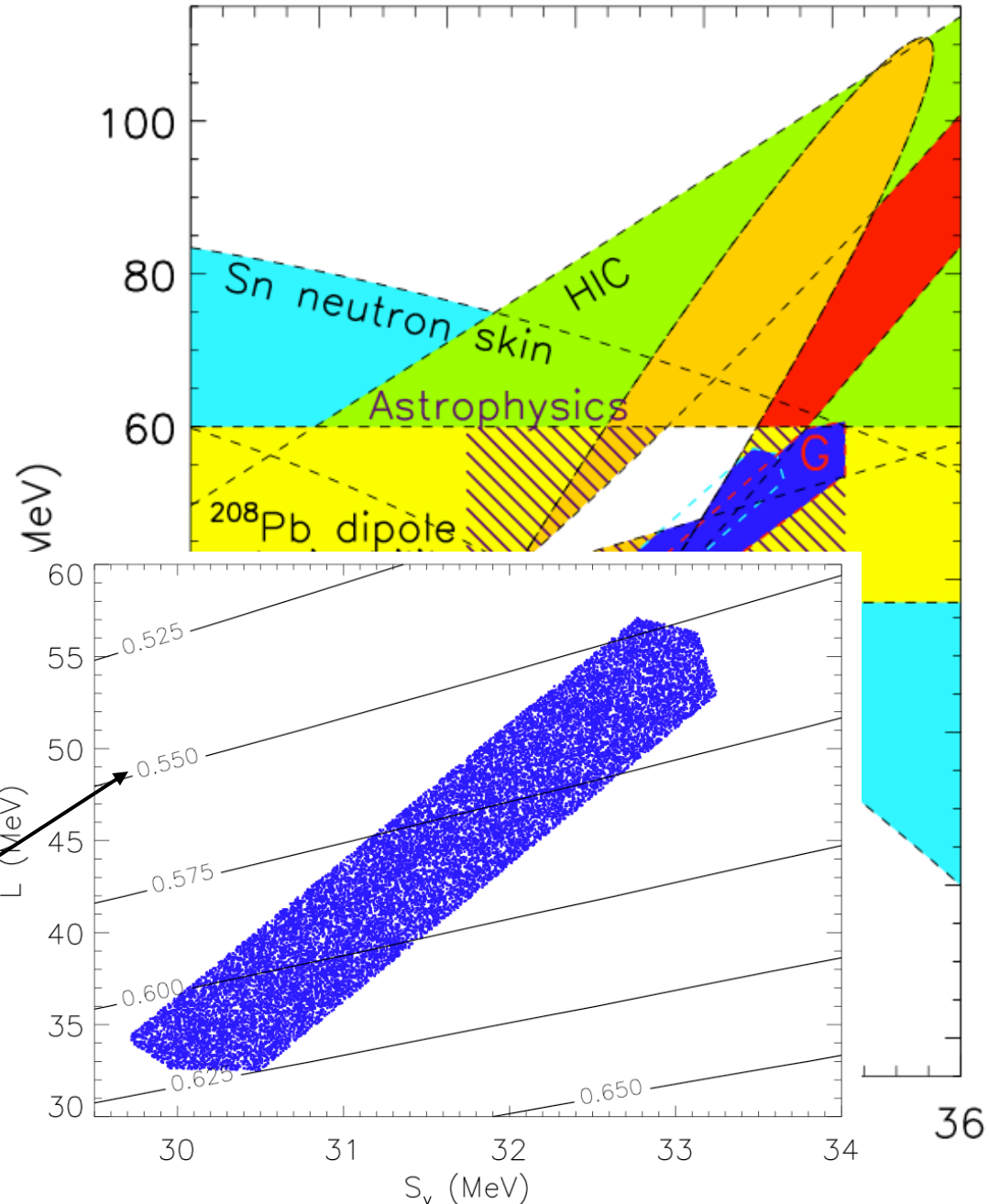
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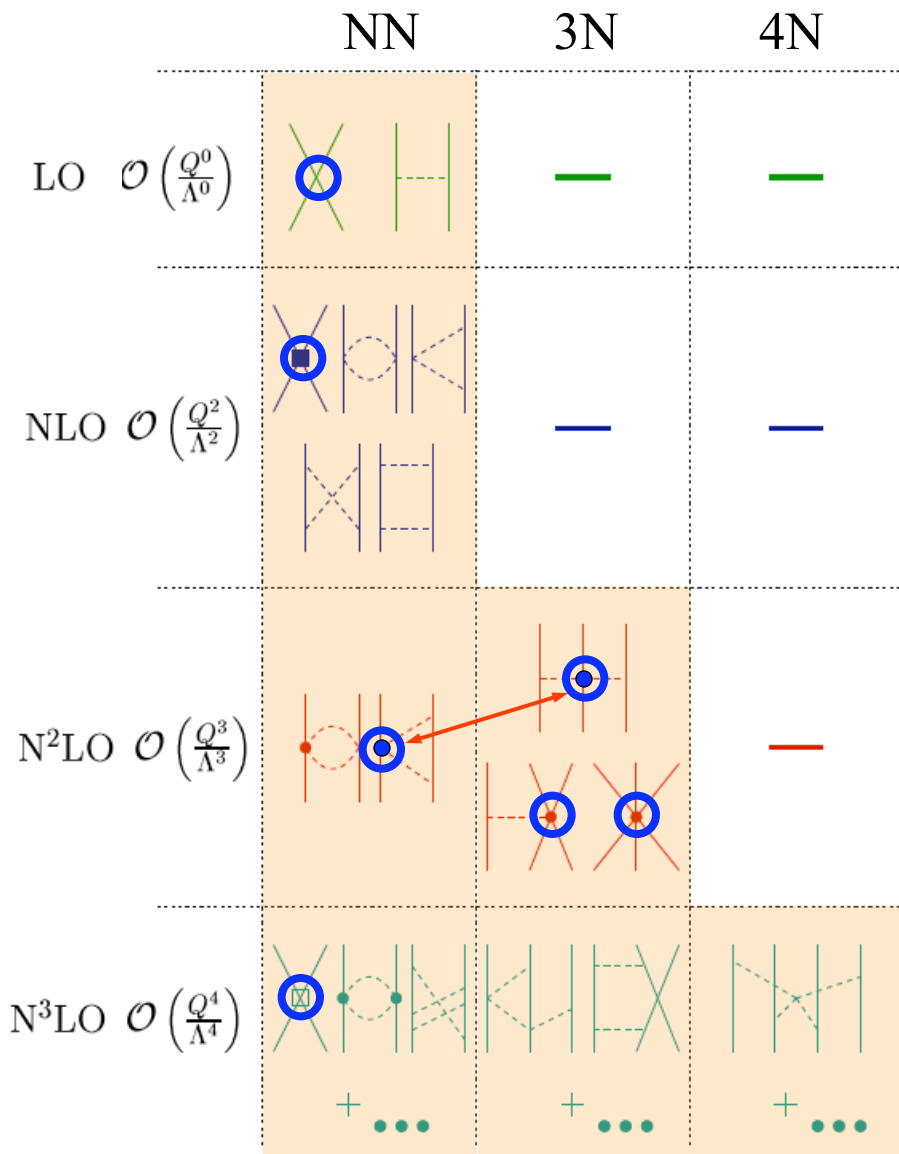
predicts correlation
but not range of S_v and L

lower limit on crust-core
transition density



Chiral Effective Field Theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV

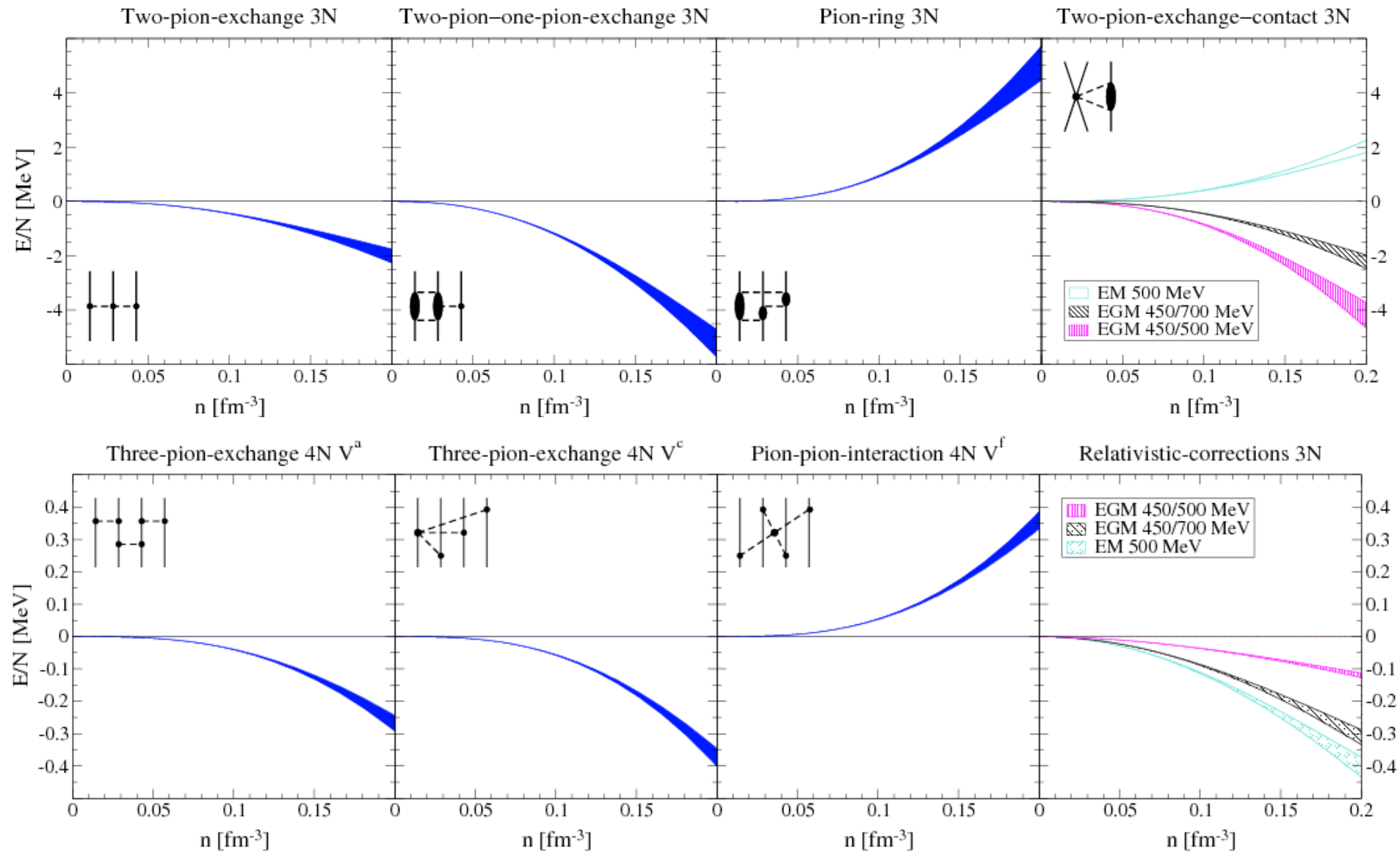


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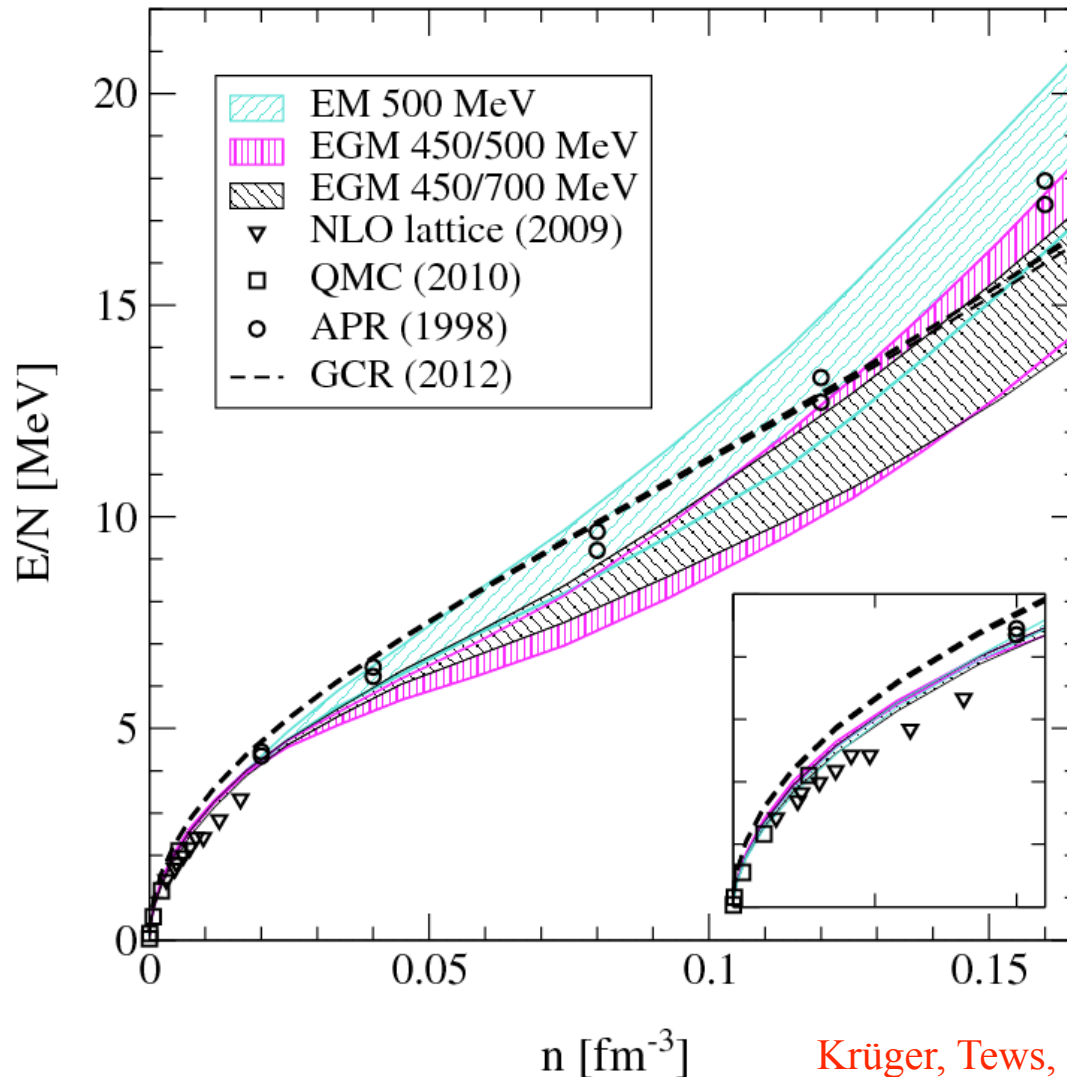
Complete N^3LO calculation of neutron matter



Complete N^3 LO calculation of neutron matter

first complete N^3 LO result

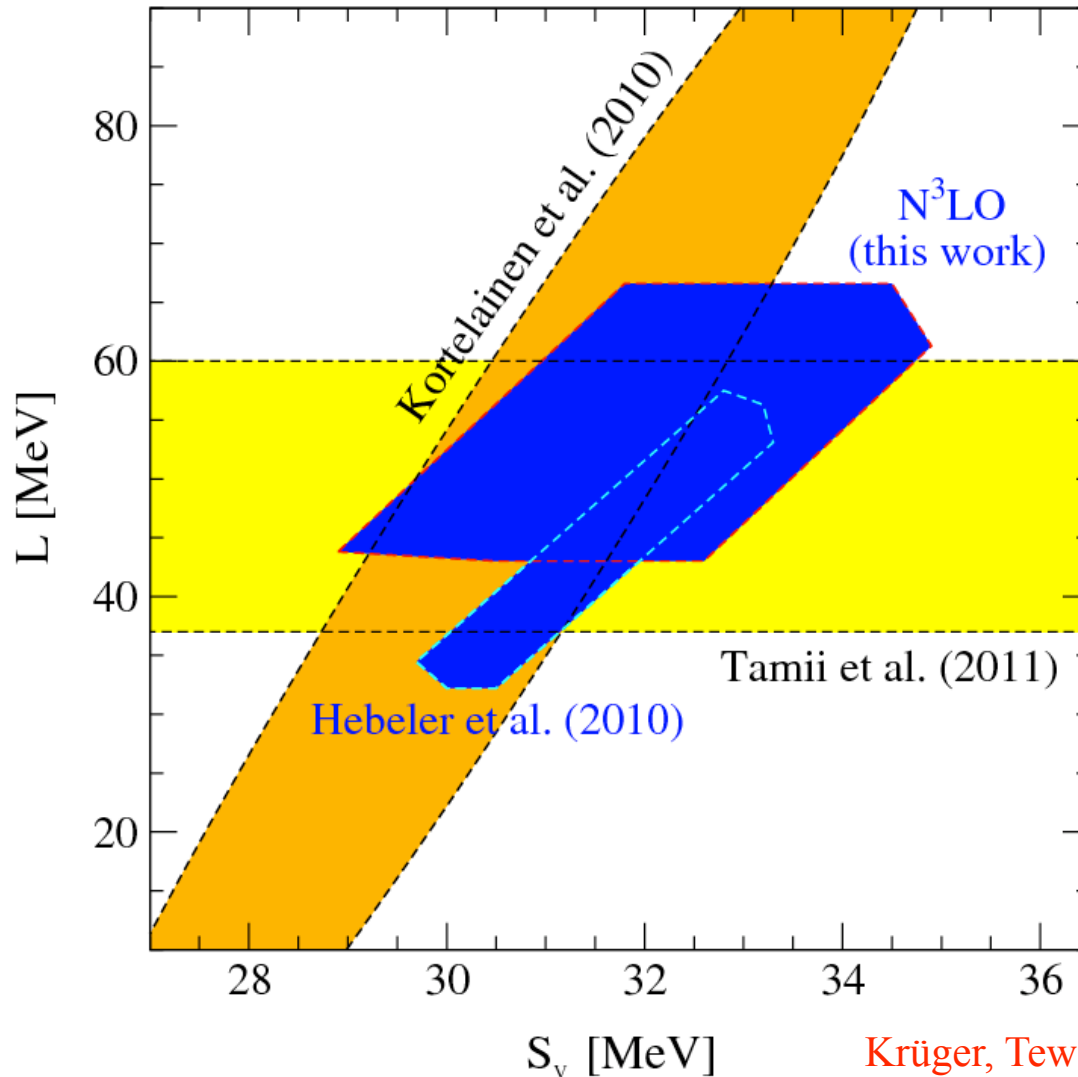
no RG evolution necessary, includes uncertainties from bare NN, 3N, 4N



Complete N^3LO calculation of neutron matter

first complete N^3LO result

no RG evolution necessary, includes uncertainties from bare NN, 3N, 4N



N^3LO correlation
broader because more
density dependences

Discovery of the heaviest neutron star

A two-solar-mass neutron star measured using Shapiro delay

P. B. Demorest¹, T. Pennucci², S. M. Ransom¹, M. S. E. Roberts³ & J. W. T. Hessels^{4,5}

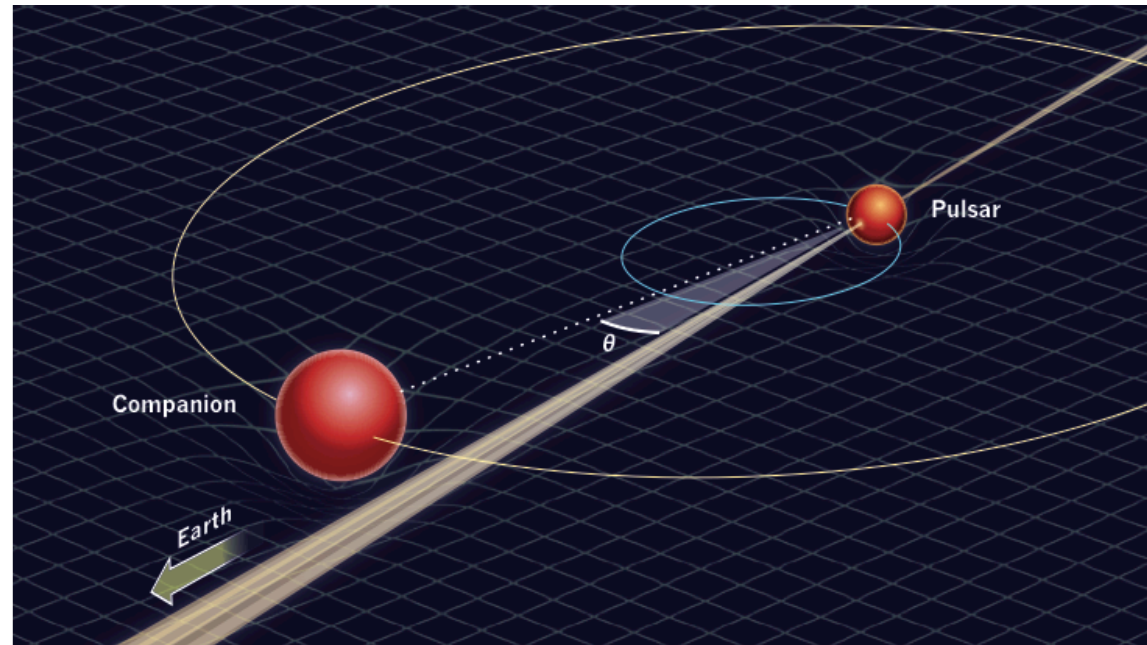
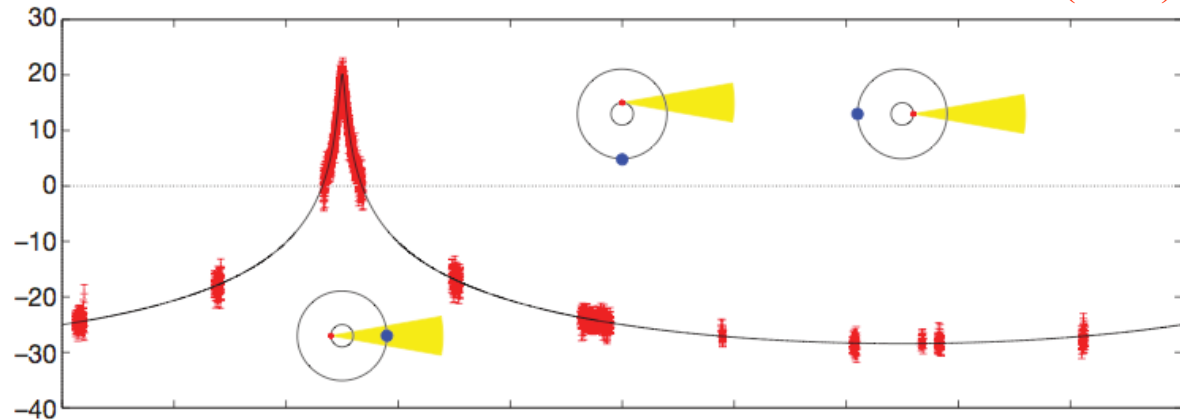
Nature (2010)

direct measurement of
neutron star mass from
increase in signal travel
time near companion

J1614-2230

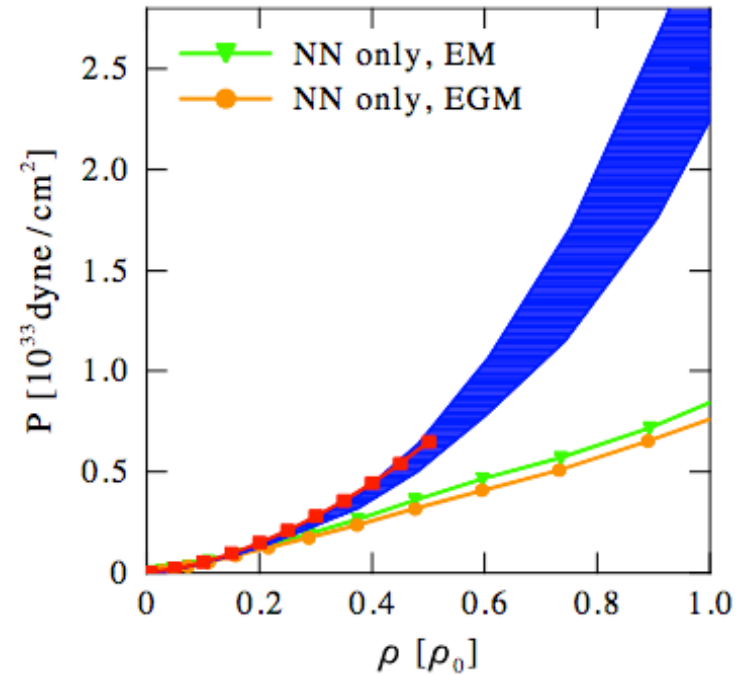
most edge-on binary
pulsar known (89.17°)
+ massive white dwarf
companion ($0.5 M_{\text{sun}}$)

heaviest neutron star
with $1.97 \pm 0.04 M_{\text{sun}}$



Impact on neutron stars Hebeler et al. (2010) and in prep.

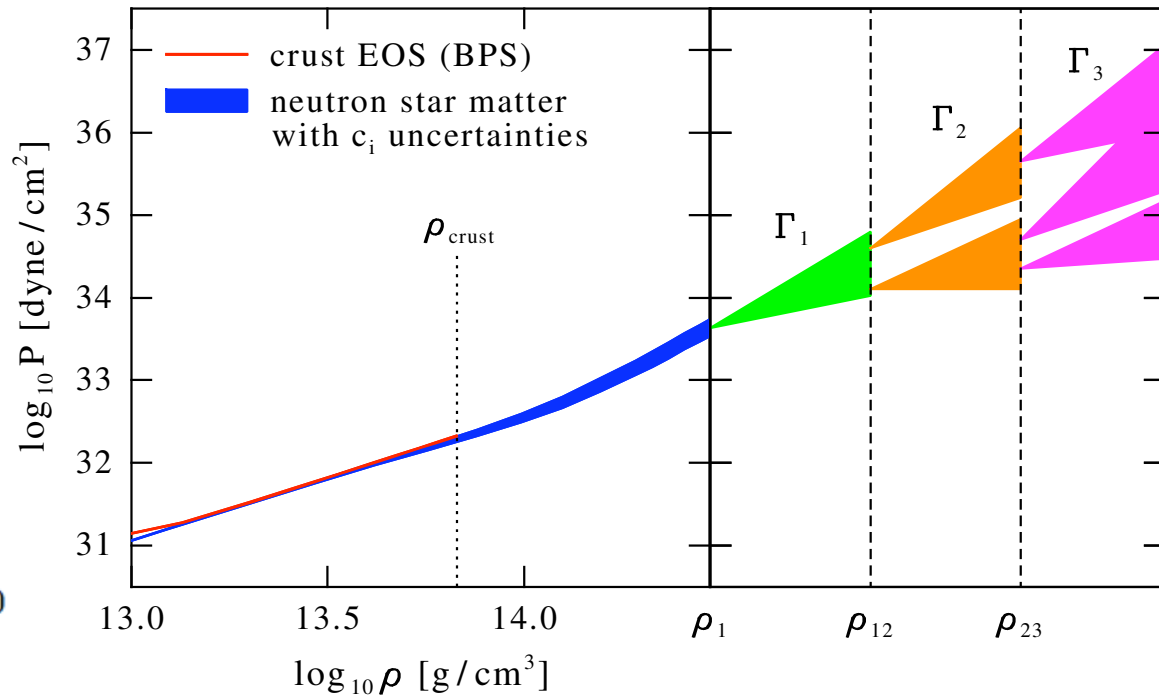
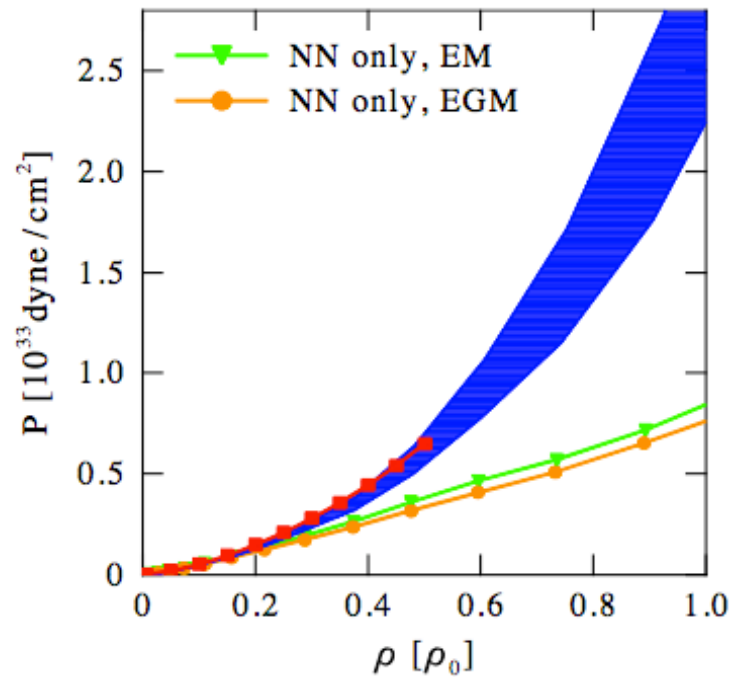
Equation of state/pressure for **neutron-star matter** (includes small $Y_{e,p}$)



pressure below nuclear densities agrees with standard crust equation of state only after 3N forces are included

Impact on neutron stars Hebeler et al. (2010) and in prep.

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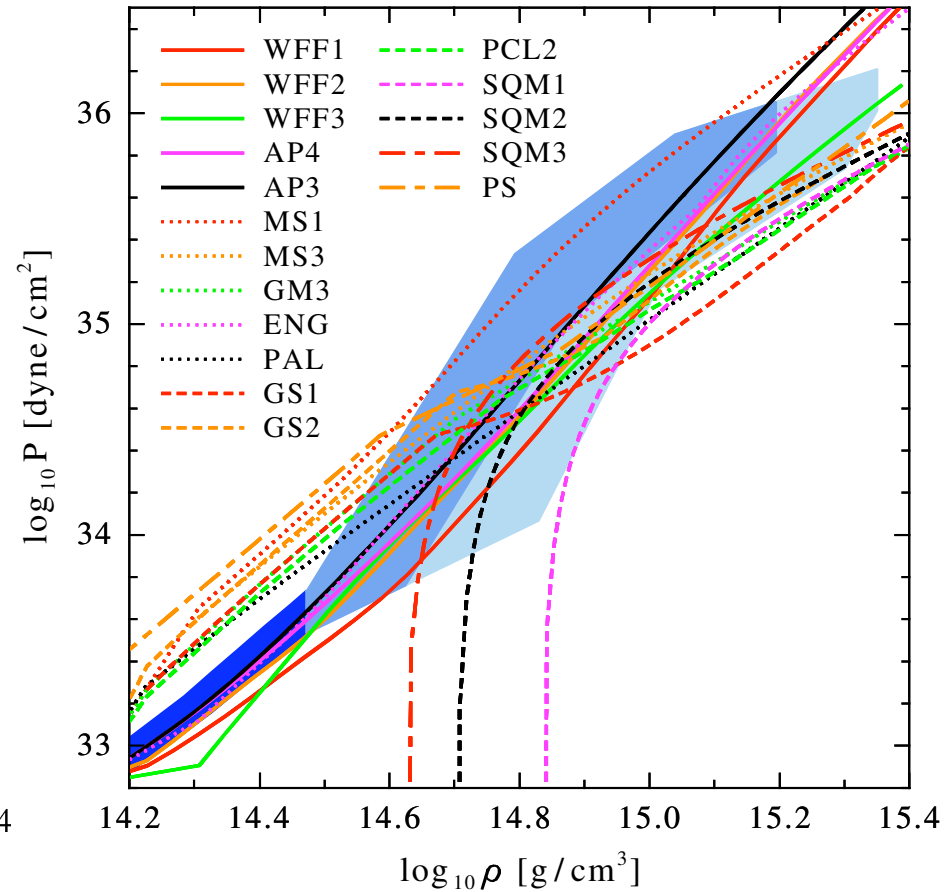
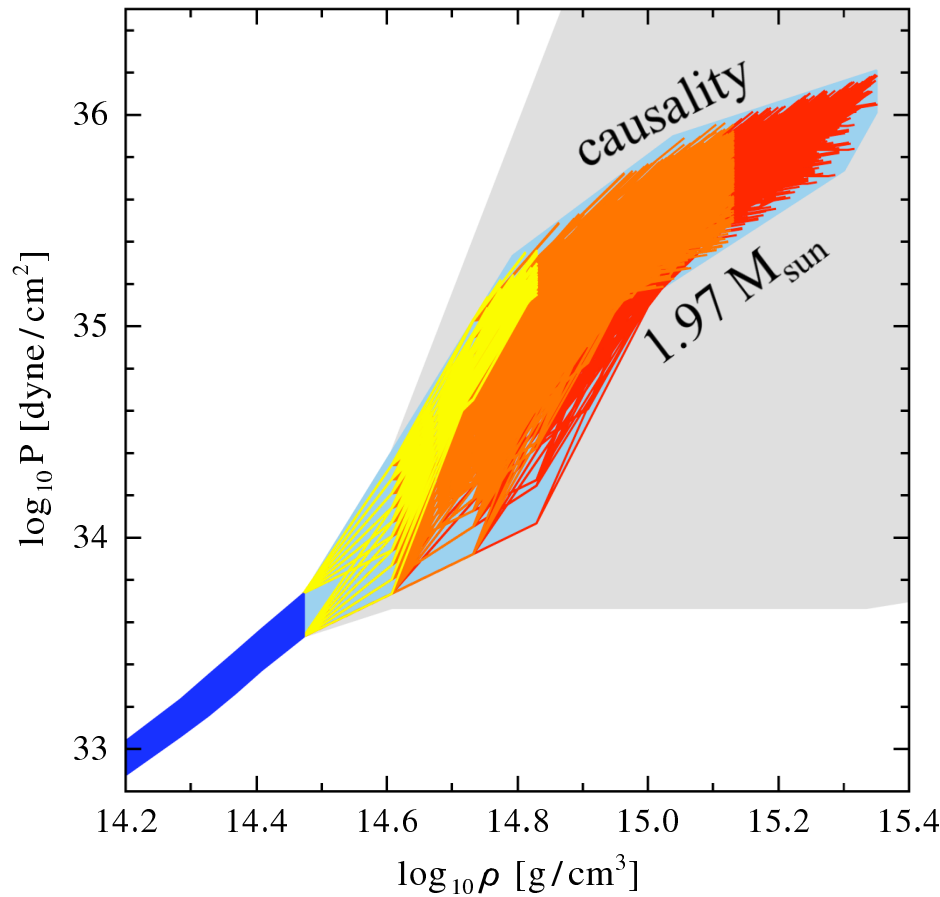


pressure below nuclear densities agrees with standard crust equation of state only after 3N forces are included

extend uncertainty band to higher densities using piecewise polytropes
allow for soft regions

Pressure of neutron star matter

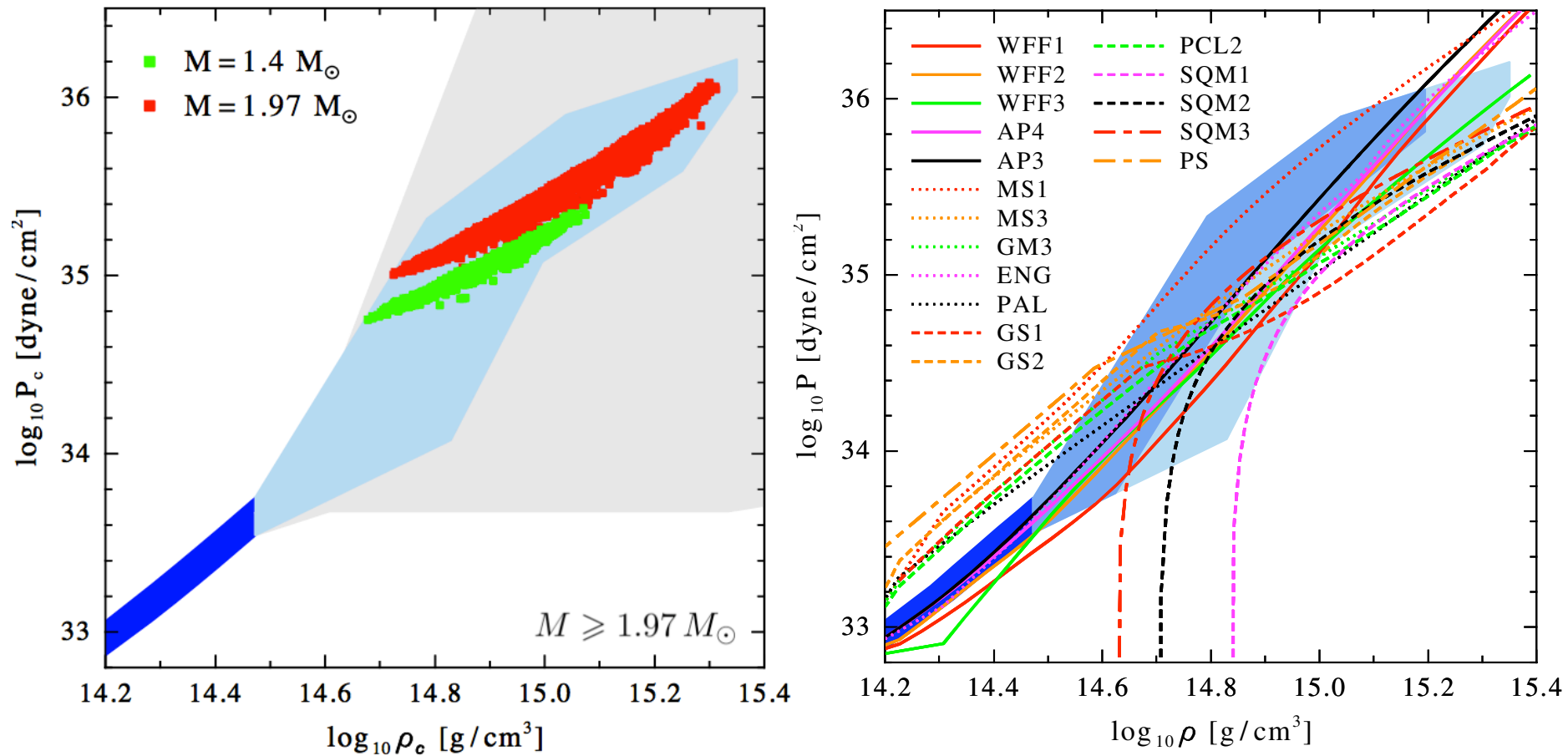
constrain polytropes by causality and require to support $1.97 M_{\text{sun}}$ star



low-density pressure sets scale, chiral EFT interactions provide strong constraints, ruling out many model equations of state

Pressure of neutron star matter

constrain polytropes by causality and require to support $1.97 M_{\text{sun}}$ star

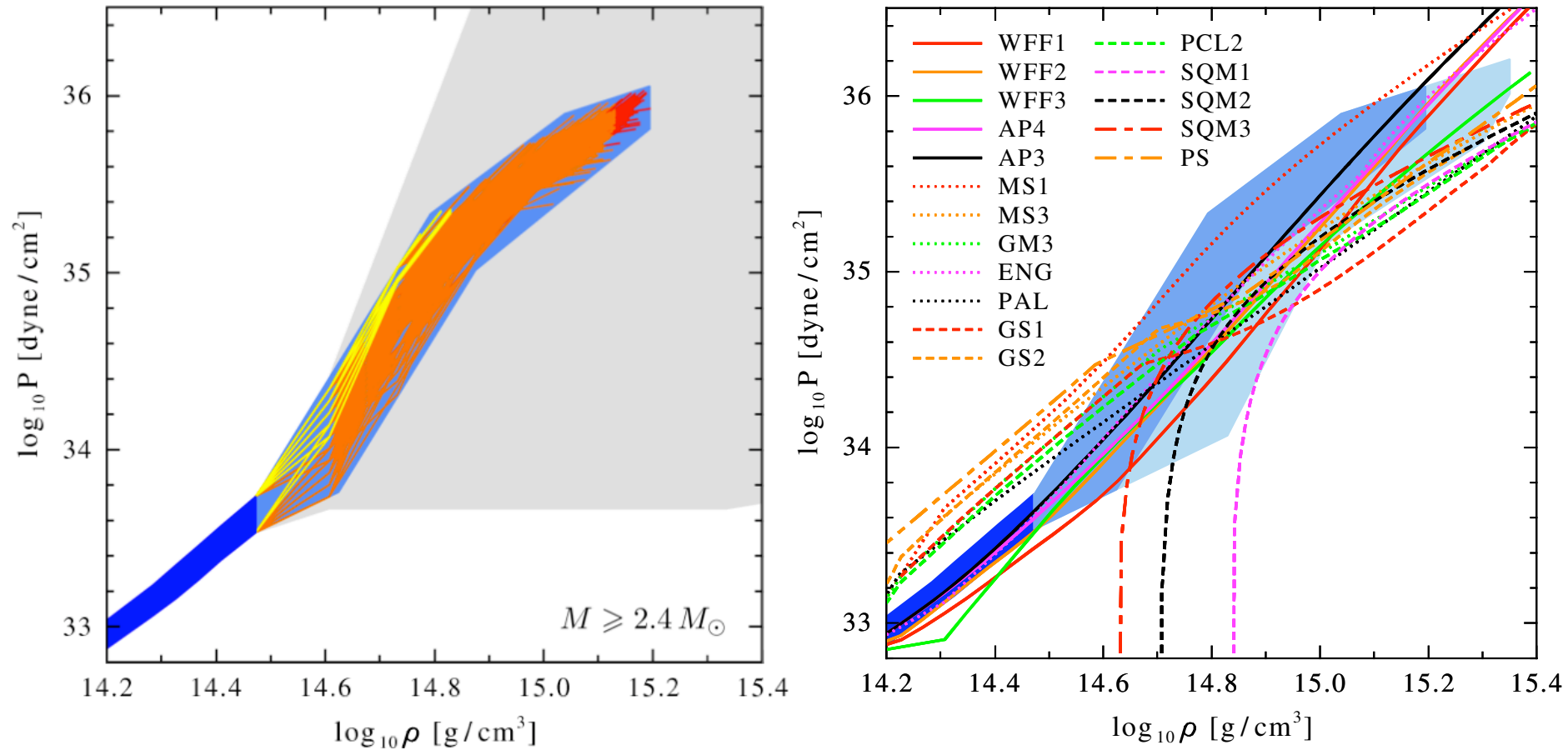


low-density pressure sets scale, chiral EFT interactions provide strong constraints, ruling out many model equations of state

central densities for $1.4 M_{\text{sun}}$ star: $1.7\text{-}4.4 \rho_0$

Pressure of neutron star matter

constrain polytropes by causality and require to support $1.97 M_{\text{sun}}$ star

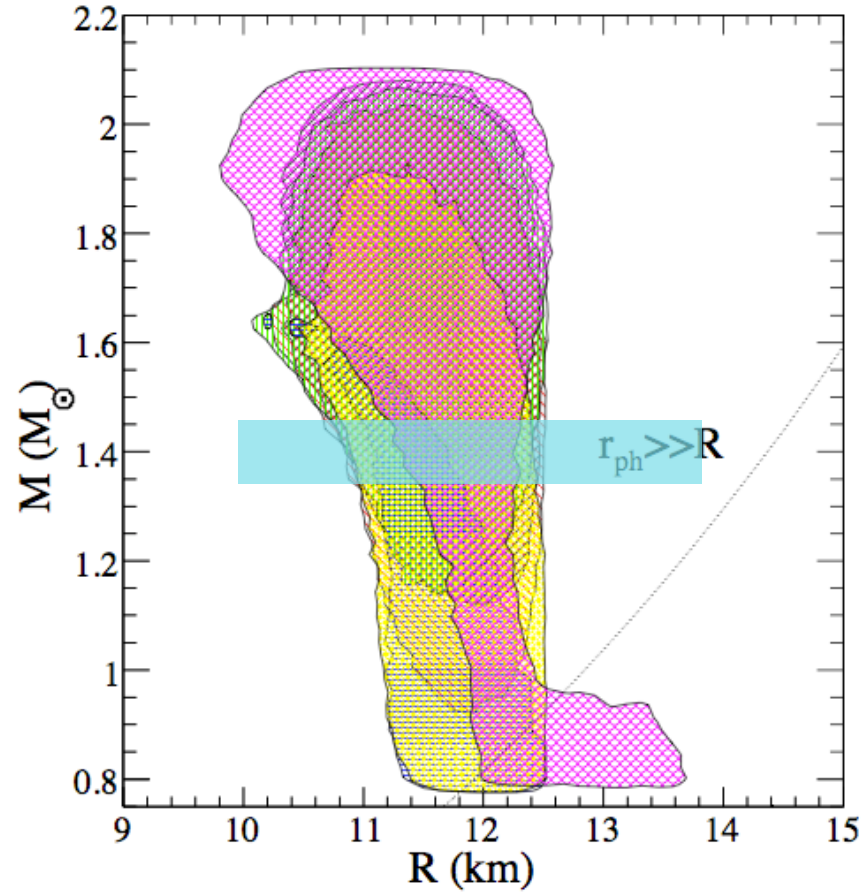
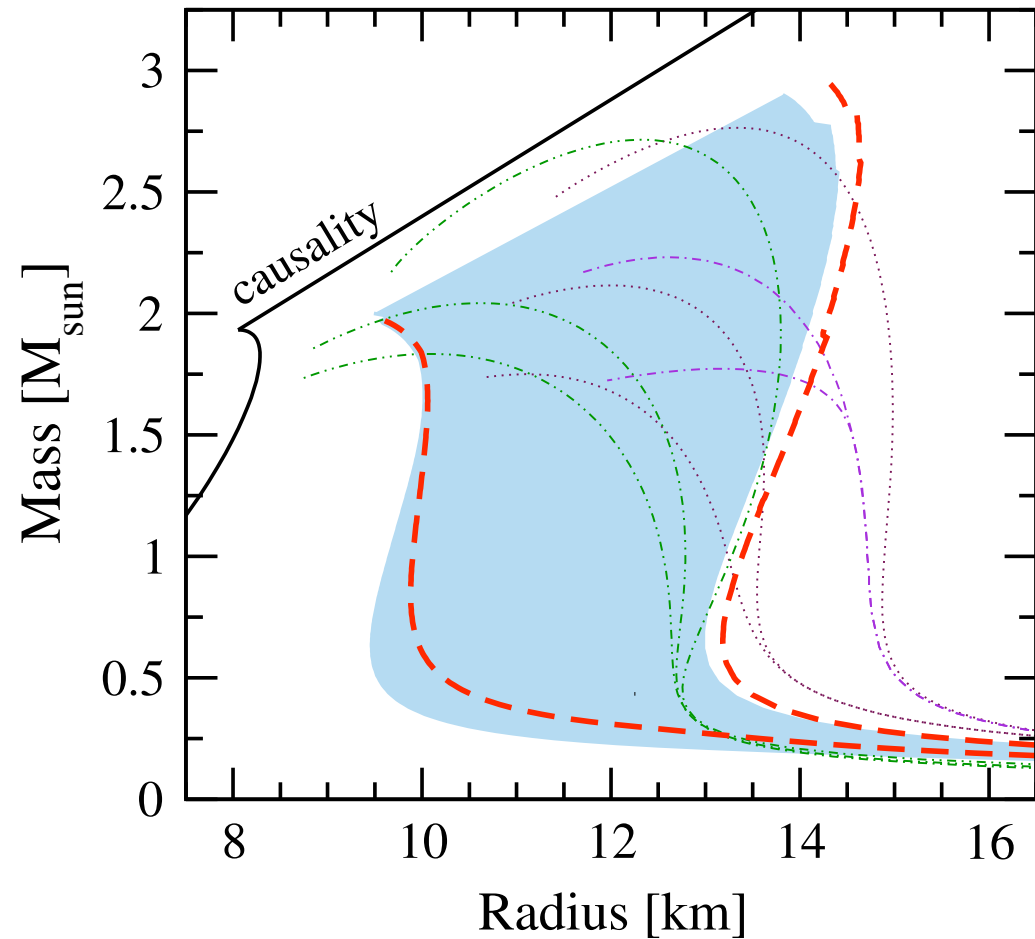


low-density pressure sets scale, chiral EFT interactions provide strong constraints, ruling out many model equations of state

darker blue band for $2.4 M_{\text{sun}}$ star

Neutron star radius constraints

uncertainty from many-body forces and general extrapolation



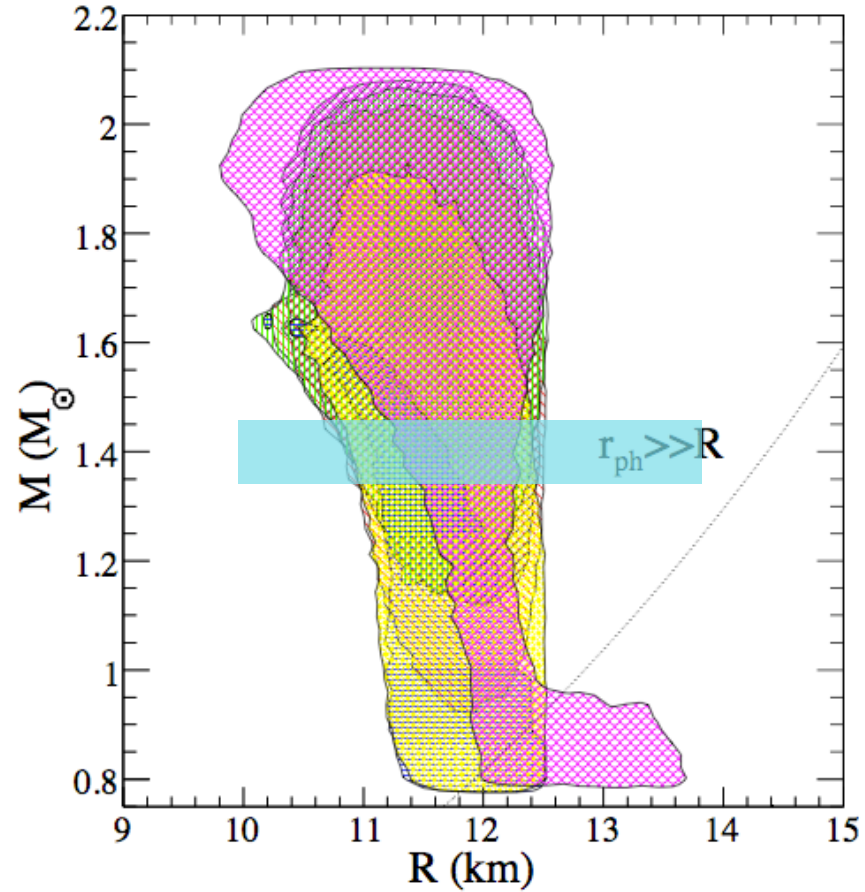
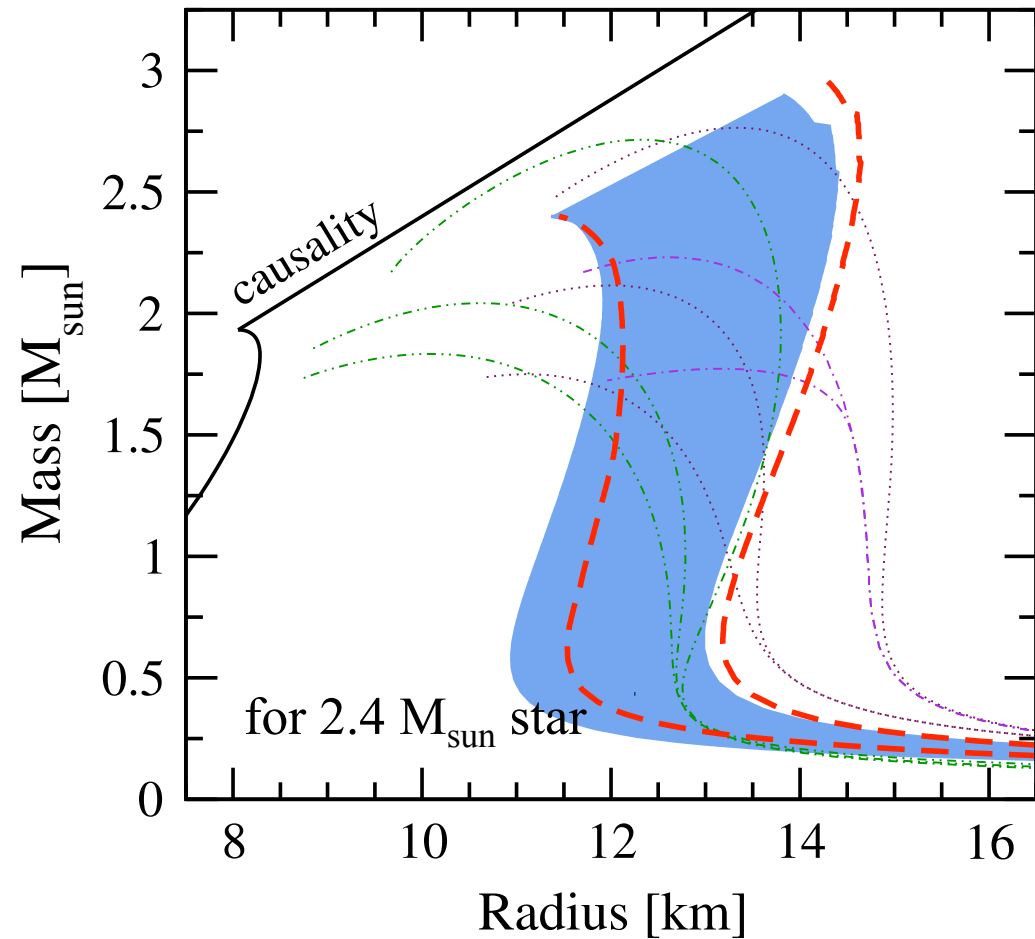
constrains neutron star radius: 9.9-13.8 km for $M=1.4 M_{\text{sun}}$ ($\pm 15\%$!)

consistent with extraction from X-ray burst sources [Steiner et al. \(2010\)](#)

provides important constraints for EOS for core-collapse supernovae

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Neutron-star merger and gravitational waves

explore sensitivity to neutron-rich matter in neutron-star merger and gw signal

Bauswein, Janka (2012) and A. Bauswein et al., arXiv:1204.1888.

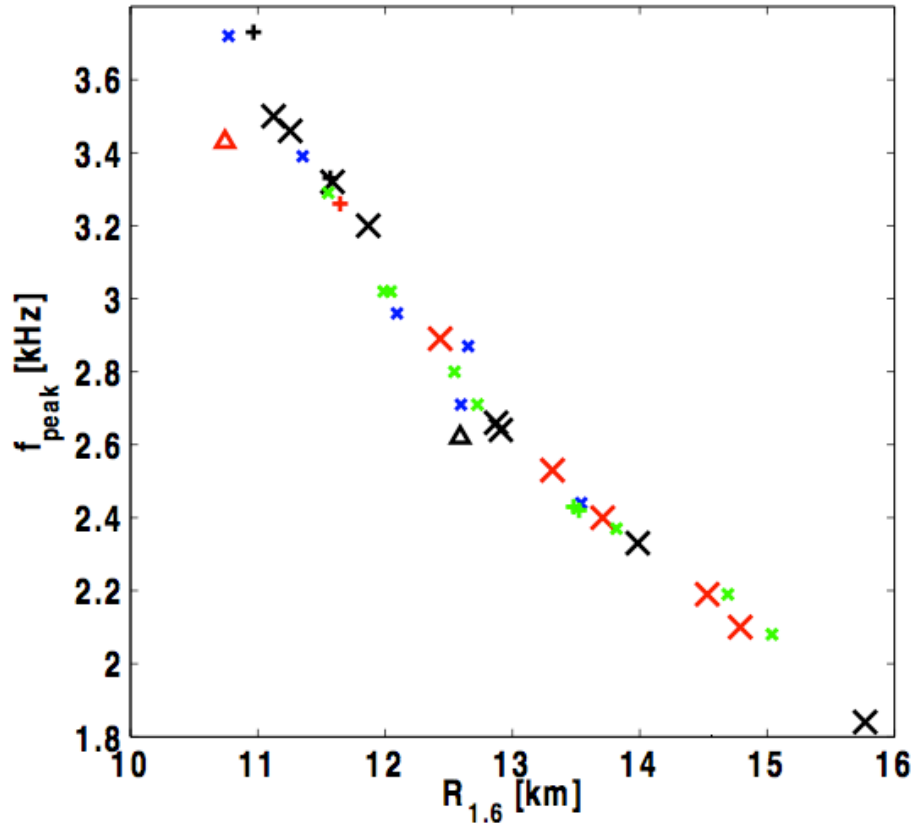


FIG. 10: Peak frequency of the postmerger GW emission versus the radius of a nonrotating NS with $1.6 M_{\odot}$ for different EoSs. Symbols have the same meaning as in Fig. 8.

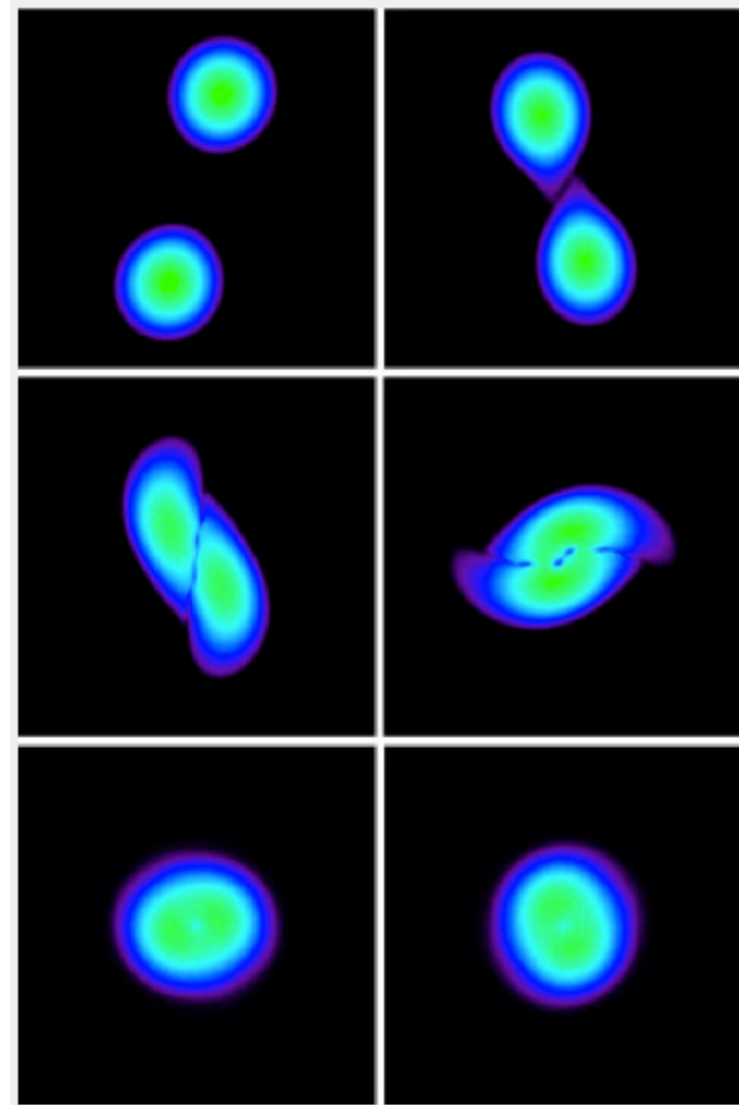
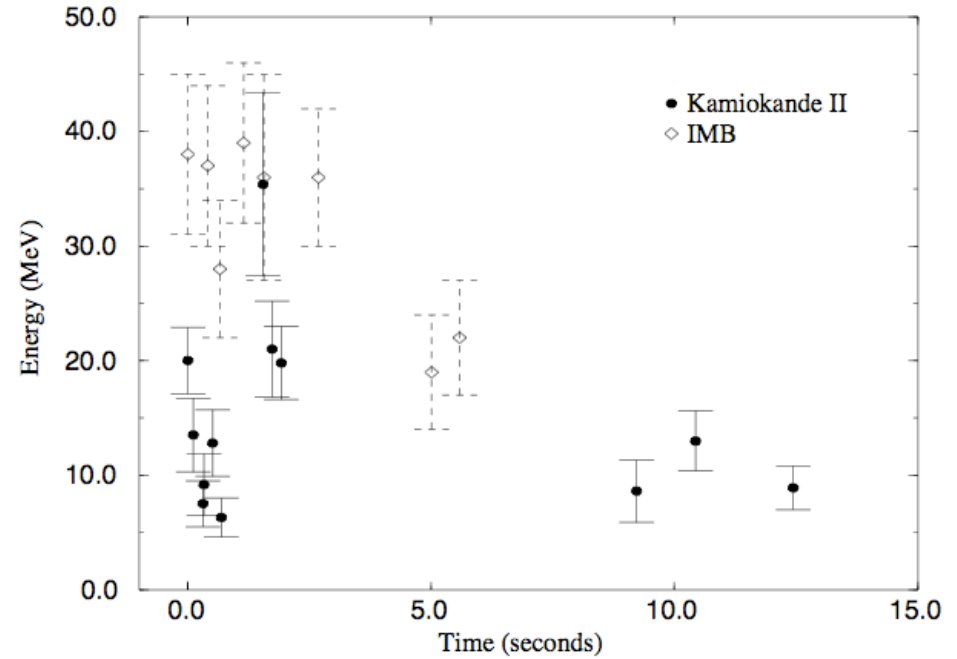


Fig. 1: Various snapshots of the collision of two neutron stars initially revolving around each other. The sequence simulated by the computer covers only 0.03 seconds. The two stars orbit each other counterclockwise (top left) and quickly come closer (top right). Finally they collide (centre left), merge (centre right), and form a dense, superheavy neutron star (bottom). Strong vibrations of the collision remnant are noticeable as deformations in east-west direction and in north-south direction (bottom panels). (Simulation: Andreas Bauswein and H.-Thomas Janka/MPA)

Neutrinos under extreme conditions



anti-electron neutrinos from SN1987a

Relevant conditions in core-collapse supernovae

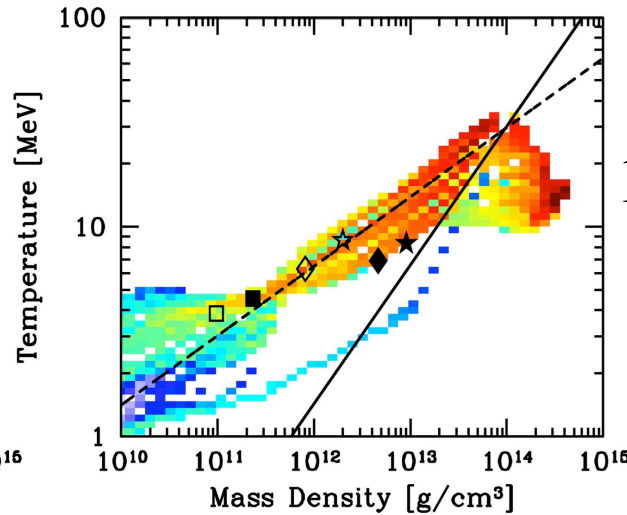
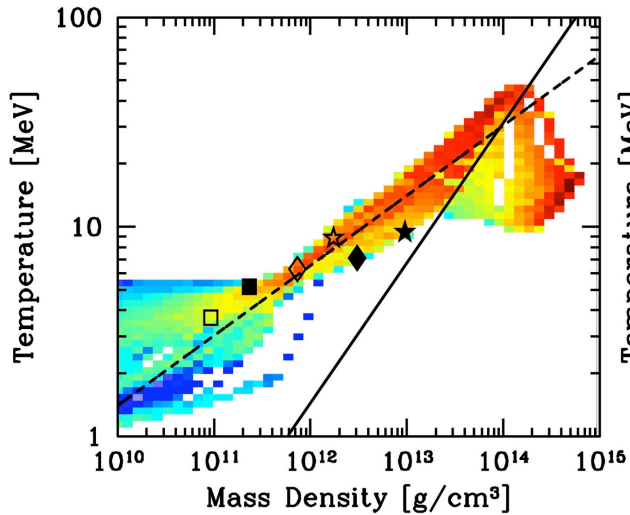
15 M_⊙ progenitor

S. Bacca et al. (2012)

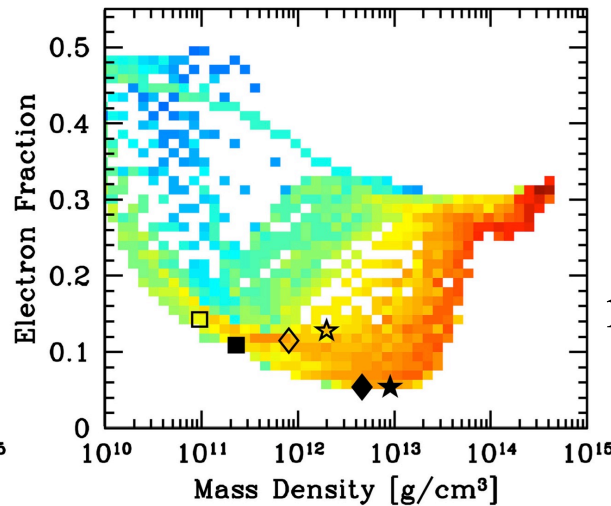
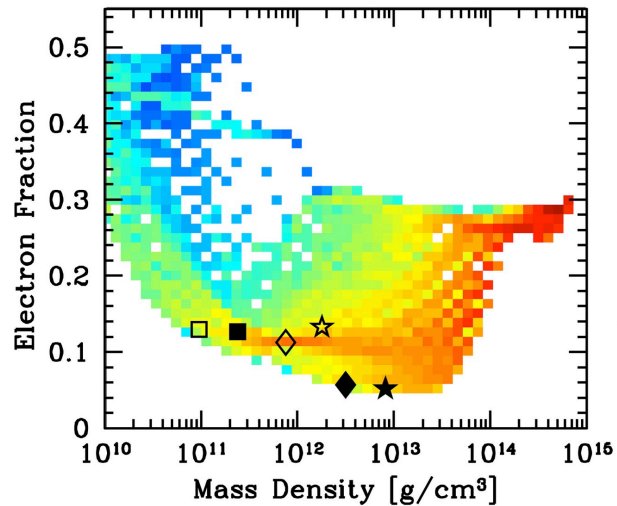
simulations by M. Liebendörfer et al.

Lattimer–Swesty EOS

Shen EOS



partially degenerate



neutron-rich

crucial densities below nuclear matter density $\sim 10^{13}$ - 10^{14} g/cm³
(high densities: neutrinos trap; low densities: few interactions)

Neutrino rates from chiral effective field theory

processes involving two nucleons play a special role [Friman,...](#) [Suzuki, Raffelt,...](#)

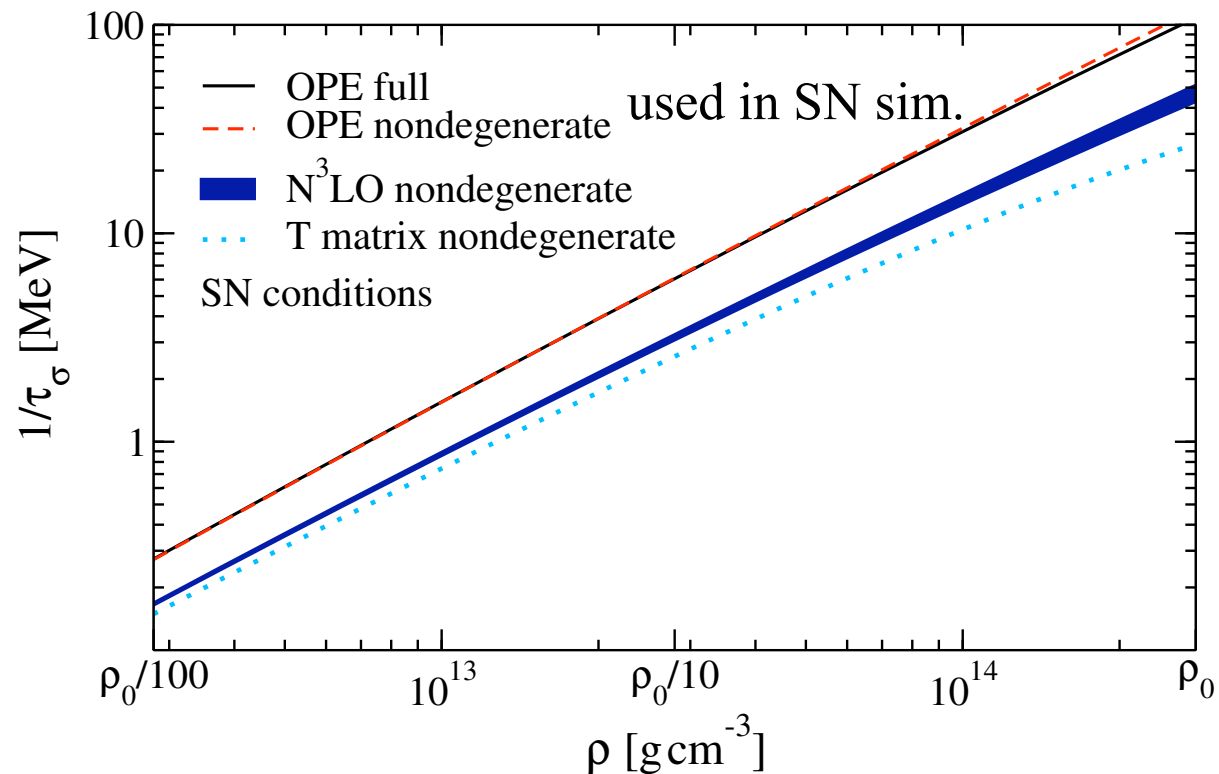
$NN \leftrightarrow NN\nu\bar{\nu}$ key for muon and tau neutrino production in supernovae
(and neutron stars crust and core cooling)

determined by spin relaxation time = rate of change of nucleon spin
through collisions

first neutrino rates
based on chiral EFT
[Bacca et al. \(2009,2012\)](#)

shorter-range interactions
reduce rates for neutrons

towards chiral EFT rates in supernova simulations



Energy transfer in neutrino scattering from nucleons

mean-square neutrino energy transfer in $\nu nn \leftrightarrow \nu nn$

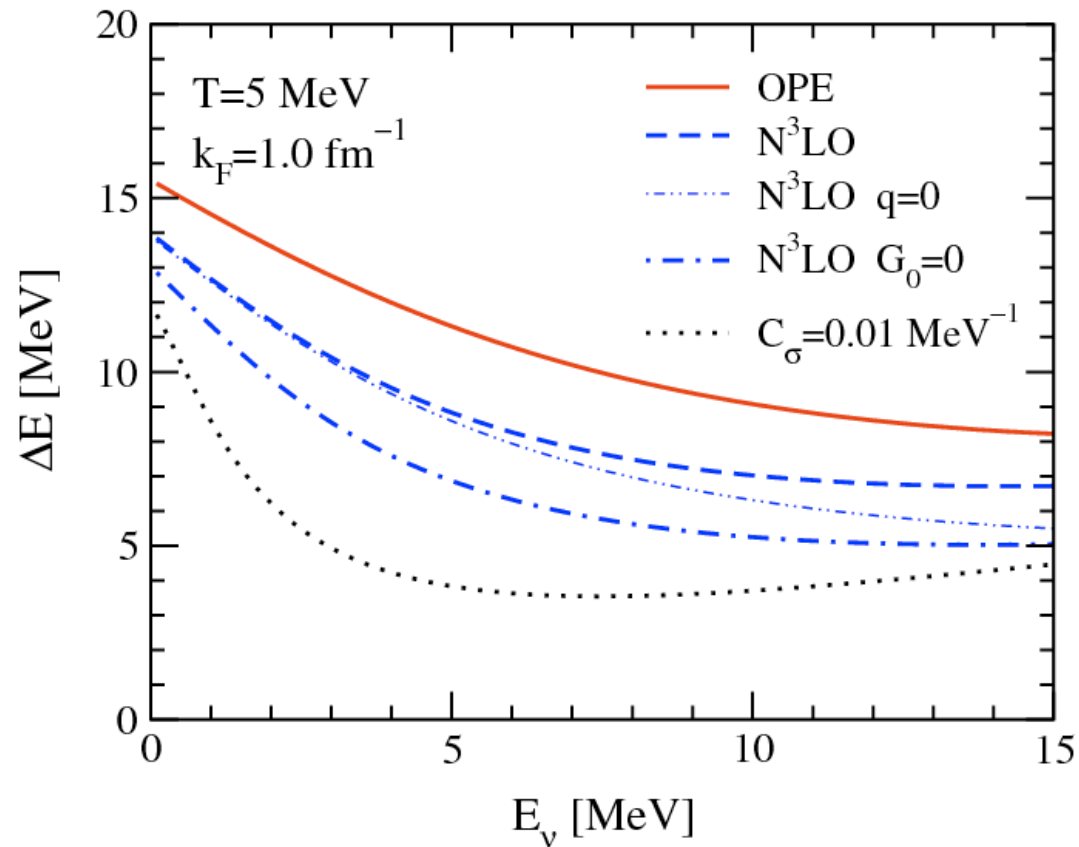
$$(\Delta E)^2 = \frac{\int d\mathbf{p}'_\nu (E_\nu - E'_\nu)^2 \Gamma(E_\nu - E'_\nu, p_\nu - p'_\nu)}{\int d\mathbf{p}'_\nu \Gamma(E_\nu - E'_\nu, p_\nu - p'_\nu)}$$

leads to heating,

NN analogue of inelastic
excitations of nuclei
(but post-collapse)

energy transfer significant,
dominates over recoil effects

not included in simulations



Main points and summary

Chiral EFT interactions provide strong constraints for EOS,
3N forces are a frontier for neutron-rich nuclei/matter

exciting intersections with cold atoms at low densities

dominant uncertainty of neutron (star) matter below nuclear densities
also key to explain neutron-rich nuclei

neutron star radius 9.9-13.8 km for $M=1.4 M_{\text{sun}}$ ($\pm 15\%$)

towards chiral EFT rates in supernova simulations,
can constrain structure factors with cold atoms at low densities